

Title: Statistical Mechanics - Lecture 1

Date: Oct 03, 2011 10:30 AM

URL: <http://pirsa.org/11100025>

Abstract:

## Where do we come from?

Undergraduate Institution:

Level of Instruction:

Different in US, Britain, Europe, Canada

Stat Mech Courses:

Thermo Course:



Concepts which specifically belong to statistical physics:  
P.W.Anderson “More IS Different”  
Physics statement, Math statement

Not in few particle quantum mechanics or in Classical Mechanics  
Physics of Large systems is different from physics of small ones  
19C early 20C getting few degrees of freedom straight  
later looking for effects of infinities  $N \rightarrow \infty$ ,  $t \rightarrow \infty$

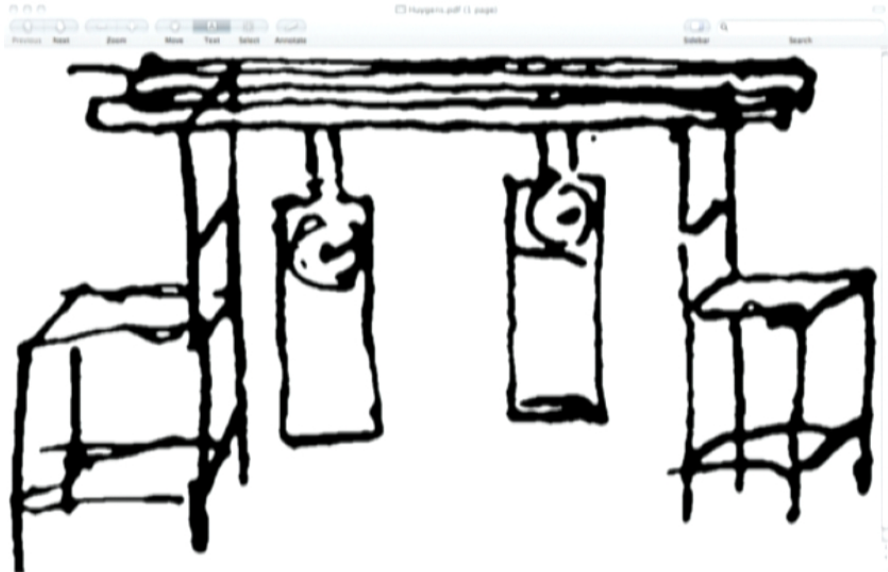
- Temperature
- Butterflies-Lorenz

## Concepts which specifically belong to statistical physics: P.W.Anderson “More IS Different” Physics statement, Math statement

Not in few particle quantum mechanics or in Classical Mechanics  
Physics of Large systems is different from physics of small ones  
19C early 20C getting few degrees of freedom straight  
later looking for effects of infinities  $N \rightarrow \infty$ ,  $t \rightarrow \infty$

- Temperature
- Butterflies-Lorenz

Example of Moire is Different:  
Christian Huygens clock synchronization



$$\mathcal{L} \quad \ddot{\theta}_1 = \omega_1^2 \sin(\theta_1) + k(\theta_1 - \theta_2)$$

$$\ddot{\theta}_2 = \omega_2^2 \sin(\theta_2) - k(\theta_1 - \theta_2)$$



$\mathcal{L}$

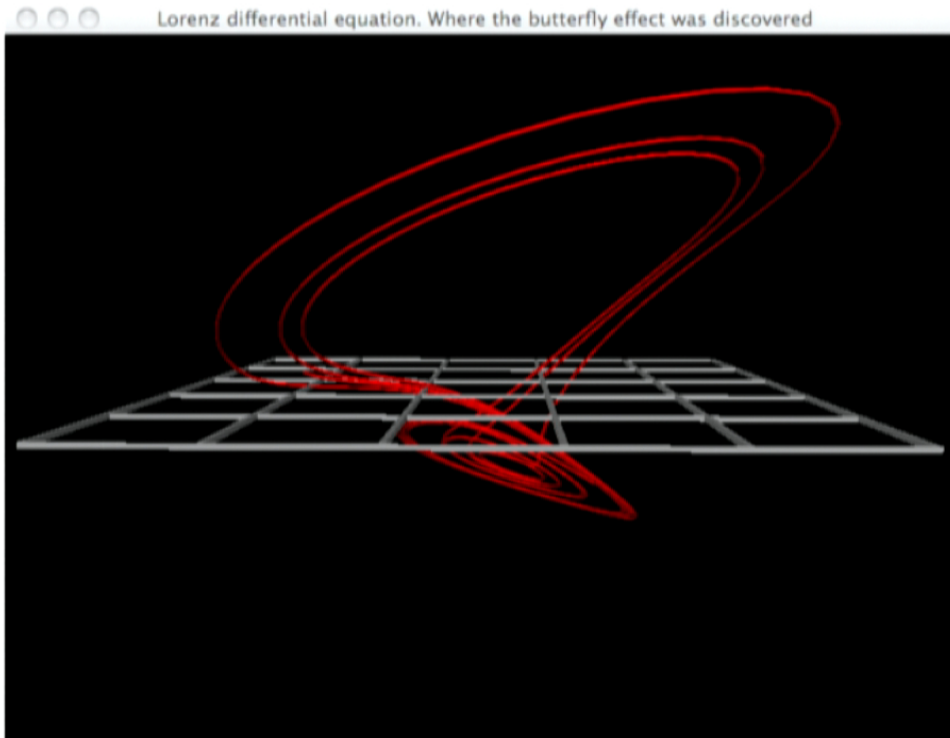
$$\ddot{\theta}_1 = \omega_1^2 \sin \theta_1 + k(\theta_1 - \theta_2)$$

$$\ddot{\theta}_2 = \omega_2^2 \sin \theta_2 - k(\theta_1 - \theta_2)$$

$$\omega_1^2 \neq \omega_2^2$$

## Butterfly effect--go for app

Lorenz, Edward N. (March 1963). "Deterministic Nonperiodic Flow". *Journal of the Atmospheric Sciences* 20 (2): 130–141. One of the most important scientific papers in the 1960s. Prior to this people believed that the behavior of physical systems was essentially predictable.



An orbit:  
after Ed. Lorenz

## How did we learn about the butterfly effect Here I express things as an algorithm

Ed Lorenz put together a very simple model of the weather: Three variables  $r=(r.x, r.y, r.z)$  a simple differential equation:

$$\text{drdt} = \text{vector} \left( \begin{array}{l} -8.0/3*r.x + r.y*r.z, \\ -10*r.y + 10*r.z, \\ -r.y*r.x + 28*r.y - r.z \end{array} \right);$$

$$r = r + \text{drdt}*dt$$

$$t=t+dt$$

Loop through this tons of times and you get

initial values:

$$dy=.01$$

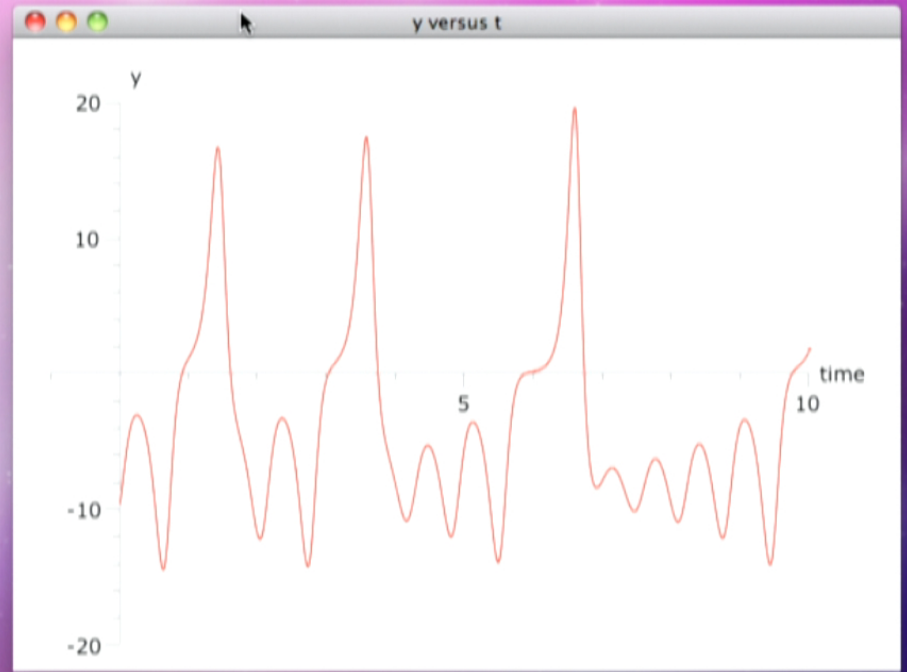
$$\text{maxt}=20$$

$$dt = 0.01$$

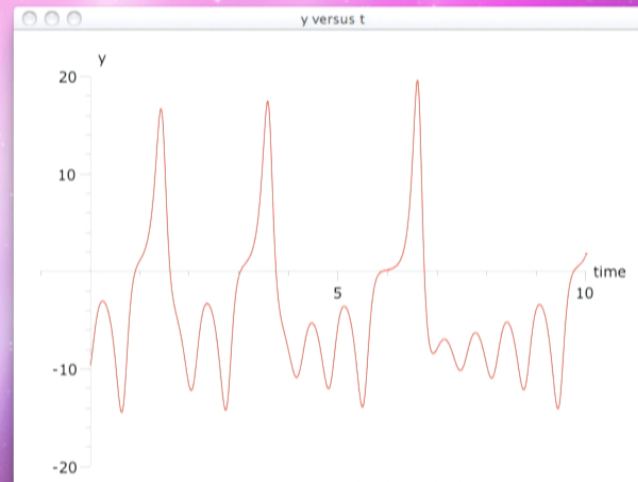
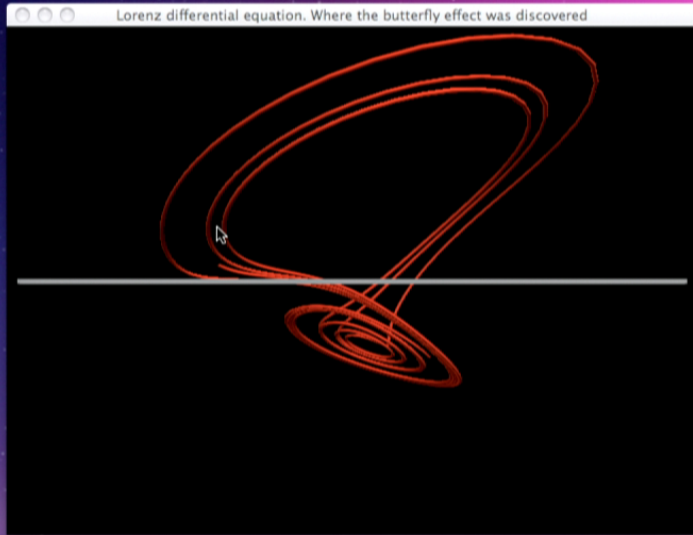
$$r = \text{vector}(35, -10, -7) \text{ \#red}$$

$$r1=\text{vector}(35, -10+dy, -7) \text{ \#green}$$

and this was called the butterfly effect





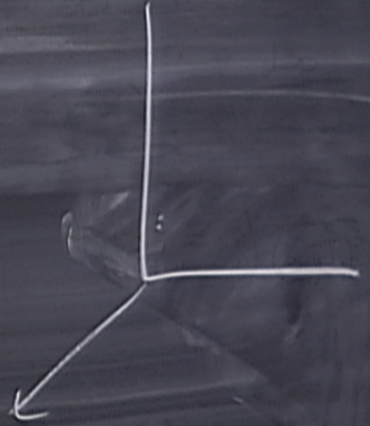




dice

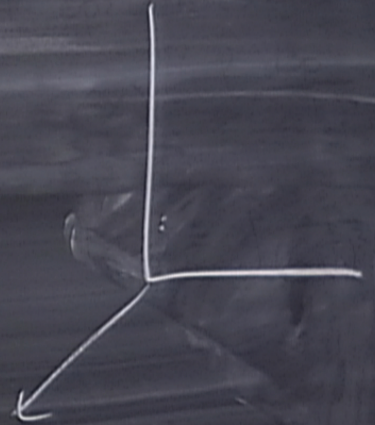
$\alpha$

$r_\alpha$   
relative probs.



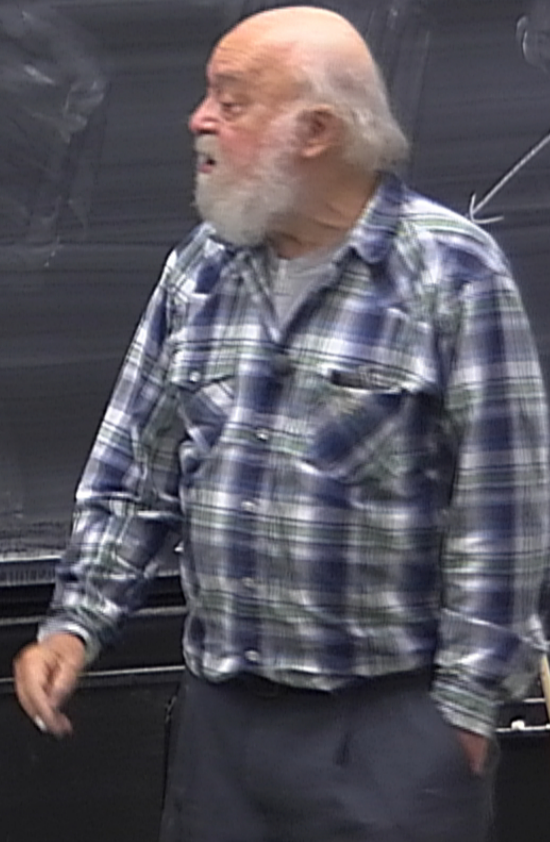


$\alpha$  dice  
 $r_\alpha = 1$   
relative probs.  
 $P_\alpha$

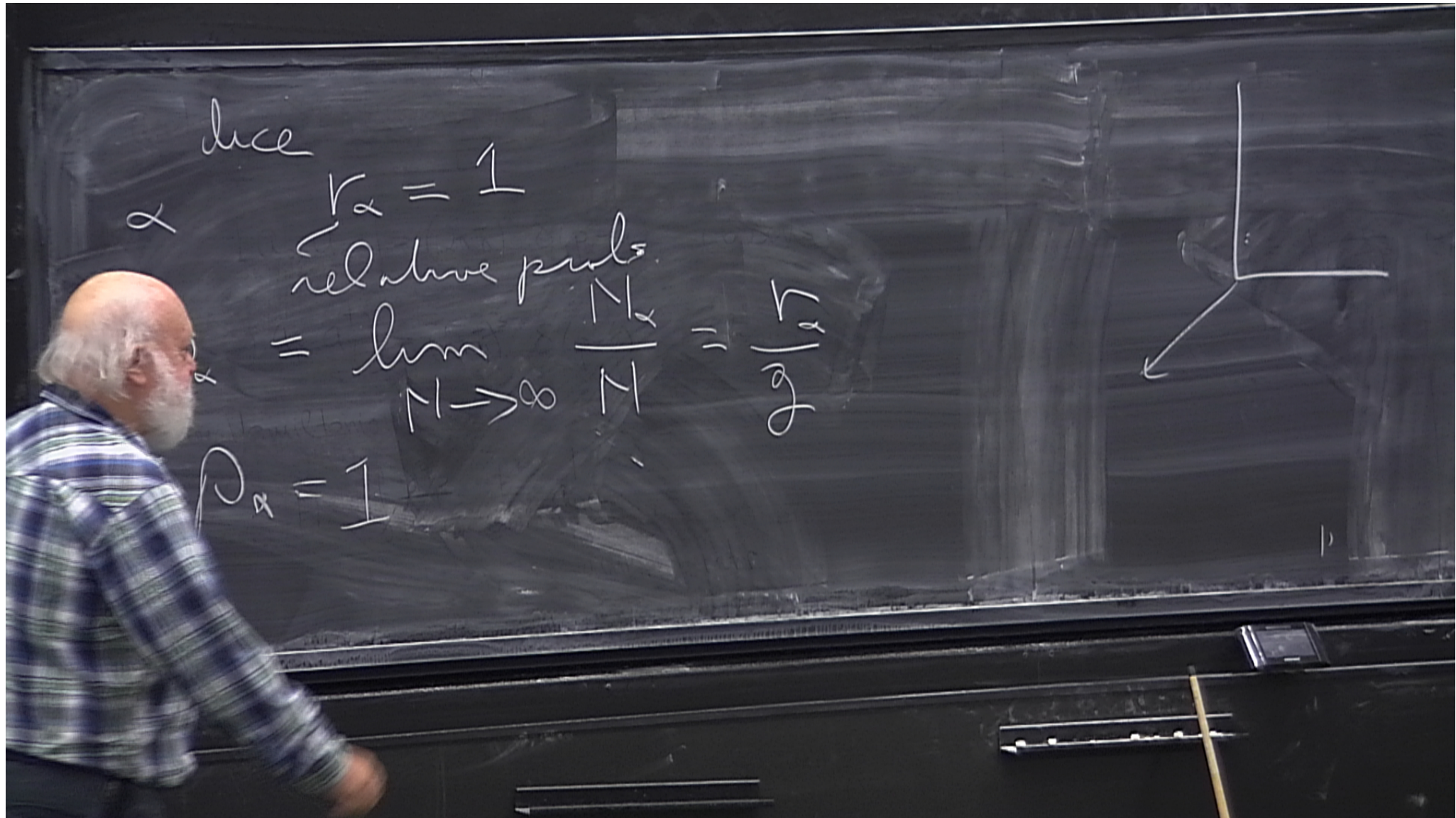




dice  
 $\alpha$   $r_\alpha = 1$   
relative probs.  
$$P_\alpha = \lim_{N \rightarrow \infty} \frac{N_\alpha}{N}$$







dice

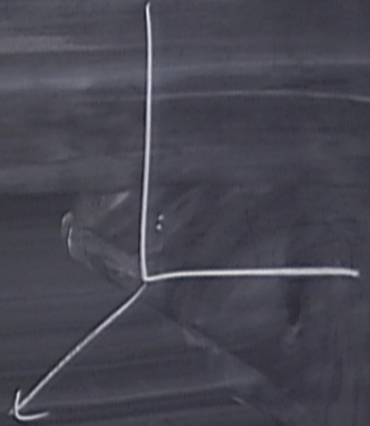
$\alpha$

$$r_\alpha = 1$$

relative probs

$$= \lim_{N \rightarrow \infty} \frac{N_\alpha}{N} = \frac{r_\alpha}{g_\alpha}$$

$$\rho_\alpha = 1$$





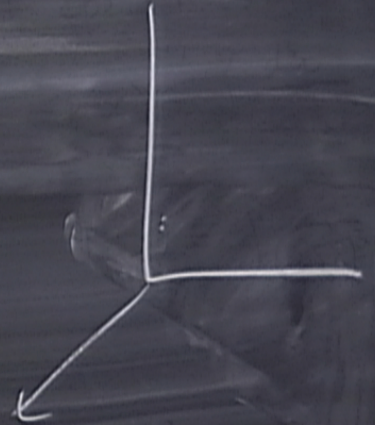
$\alpha$  dice

$$r_\alpha = 1$$

relative probs

$$p_\alpha = \lim_{N \rightarrow \infty} \frac{N_\alpha}{N} = \frac{r_\alpha}{g}$$

$$\sum_\alpha p_\alpha = 1 \quad p_\alpha \geq 0$$





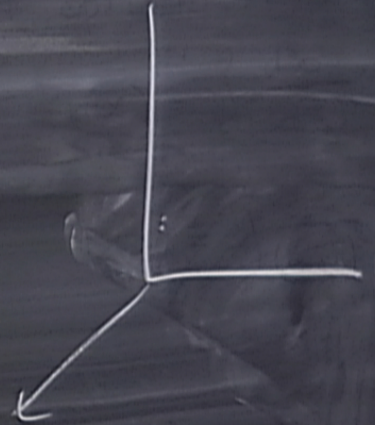
dice

$$r_\alpha = 1$$

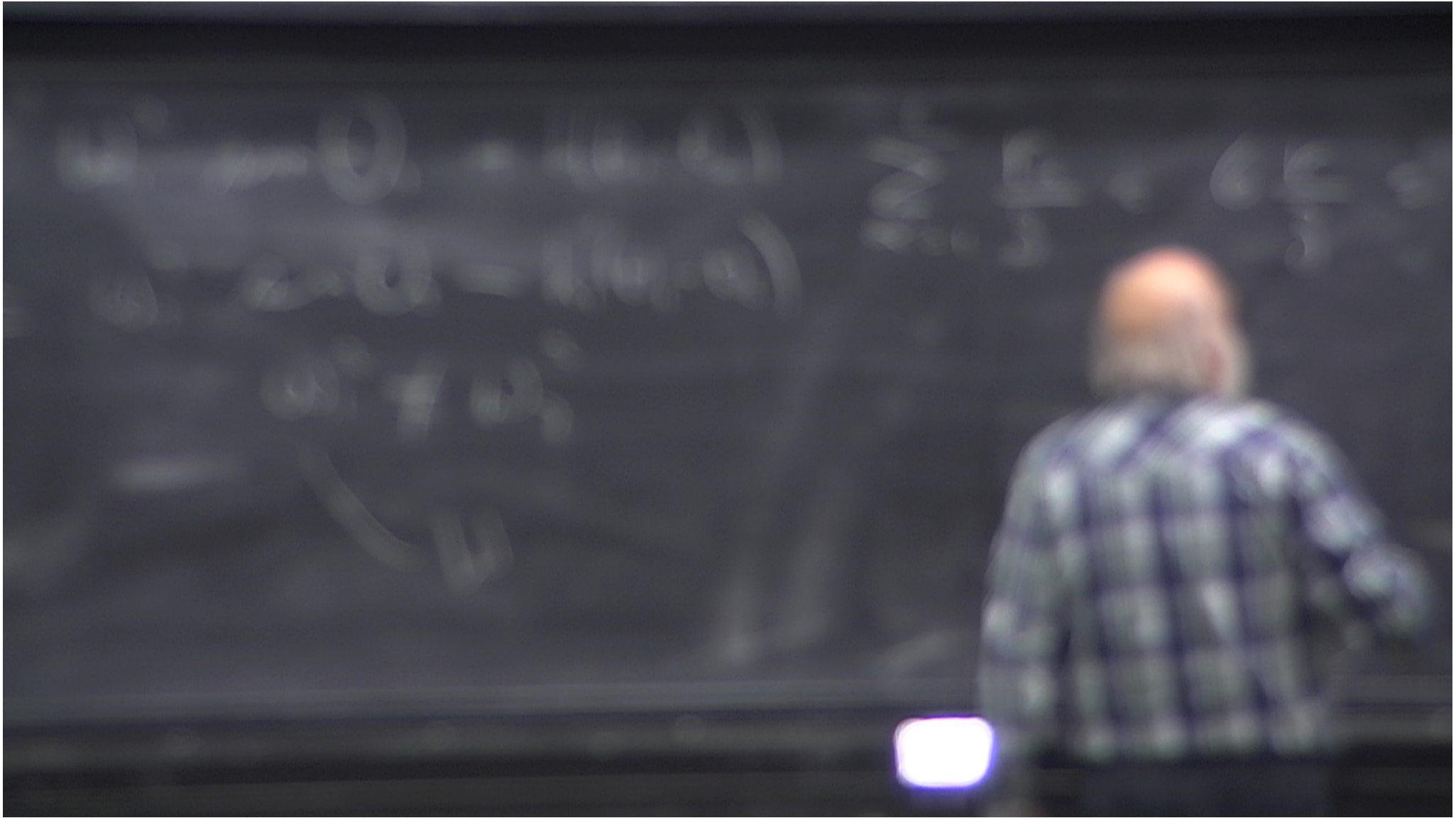
relative probs

$$p_\alpha = \lim_{N \rightarrow \infty} \frac{N_\alpha}{N} = \frac{r_\alpha}{g}$$

$$\sum_\alpha p_\alpha = 1 \quad p_\alpha \geq 0$$









$$\omega_1^2 \sin \theta_1 + k(\theta_1 - \theta_2)$$

$$\omega_2^2 \sin \theta_2 - k'(\theta_1 - \theta_2)$$

$$\omega_1^2 \neq \omega_2^2$$

$\omega$

$$\sum_{\alpha=1}^6 \frac{r_{\alpha}}{3} = 6 \frac{r}{3} =$$

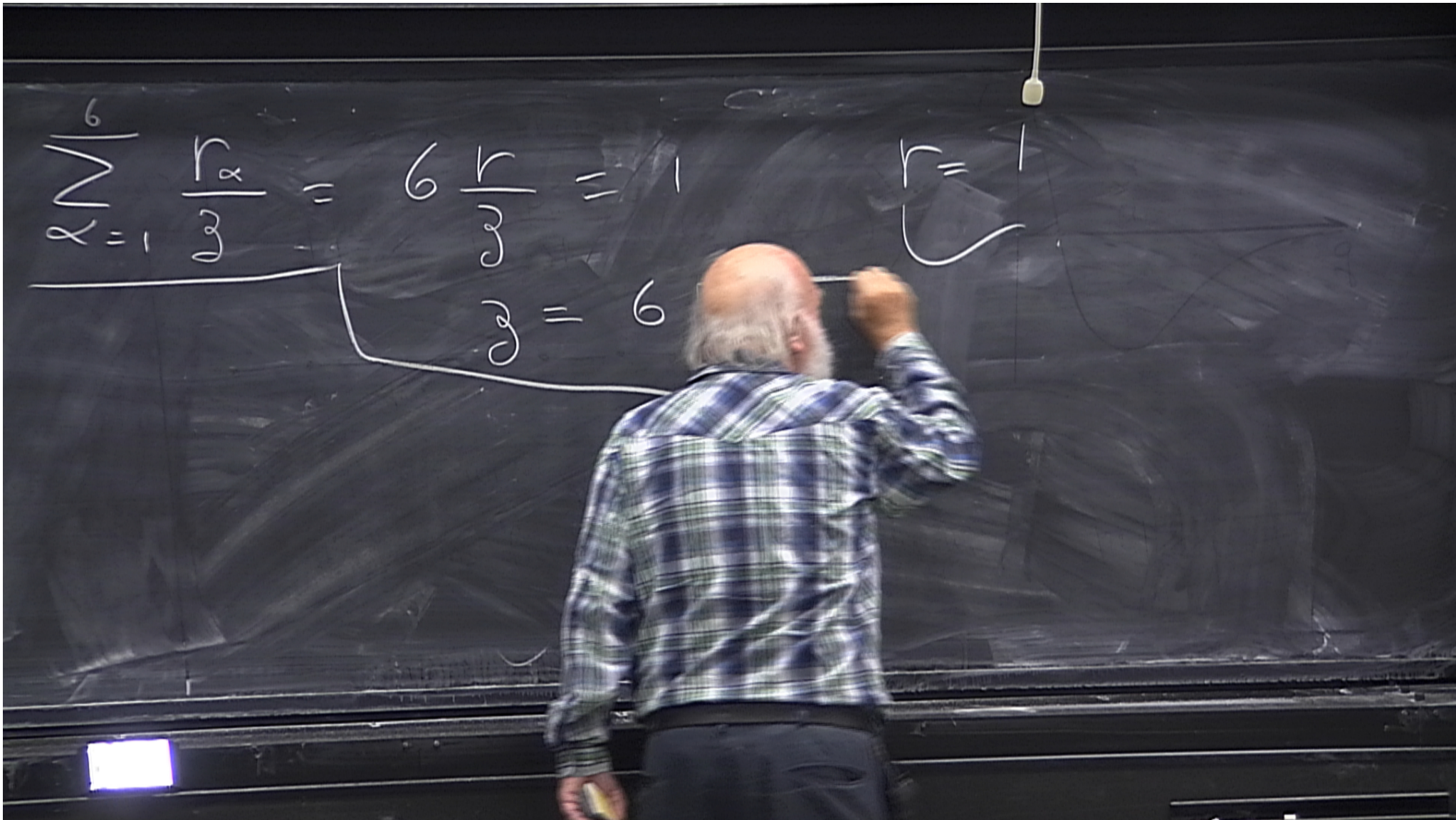


$$\sum_{k=1}^6 \frac{r}{3} = 6 \frac{r}{3} = 1$$

$$r = \frac{1}{6}$$









$$\sum_{\alpha=1}^6 \frac{r}{3} = 6 \frac{r}{3} = 1$$

$$r = 1$$

$$3 = 6$$

M P 8



$$\theta_2) \quad \sum_{\alpha=1}^6 \frac{r_\alpha}{3} = 6 \frac{r}{3} = 1 \quad r=1$$

$$-\theta_2) \quad \frac{\quad}{3=6}$$

C.M

$$d\gamma = \prod_j (dp_j) (dq_j) \quad j=1, 2, \dots, 7$$



$$\sum_{\alpha=1}^6 \frac{r_{\alpha}}{j} = 6 \frac{r}{j} = 1 \quad r=1$$

$$j = 6$$

Gibbs Boltzmann  
Maxwell

$$j = 1, 2, \dots, 7$$

$C$   
 $dX$   
 $g$   
 $(p_j)(dq_j)$



$$\frac{r_\alpha}{3} = 6 \frac{r}{3} = 1 \quad r=1$$

$$\Omega = \prod_{j=1}^M (dp_j)^2 (dq_j)^2$$

$$\Omega = 6 \quad j=1, 2, \dots, N$$

Gibbs Boltzmann  
Maxwell

a priori you have  
an equal chance of  
being anywhere in  $\Omega$



$$Q \rho_Q \langle Q \rangle = \frac{\sum_{\alpha} \langle \alpha | Q | \alpha \rangle}{\sum_{\alpha} \langle \alpha | \alpha \rangle} = 1$$

$\rho_Q$



$$Q \text{ in } \langle Q \rangle = \frac{\sum_{\alpha} \langle \alpha | Q | \alpha \rangle}{\sum_{\alpha} \langle \alpha | \alpha \rangle} = 1$$

$$\rho_{\alpha} = \frac{1}{2M} \frac{1}{M}$$



$$r_x = e^{-\beta E_x} \quad \rho_x = \frac{r_x}{Z} \quad Z = \sum_x e^{-\beta E_x}$$
$$\beta = \frac{1}{k_B T}$$



$$r_{\alpha} = d\gamma e^{-\beta E(\gamma)}$$
$$Z = \int d\gamma e^{-\beta E(\gamma)}$$

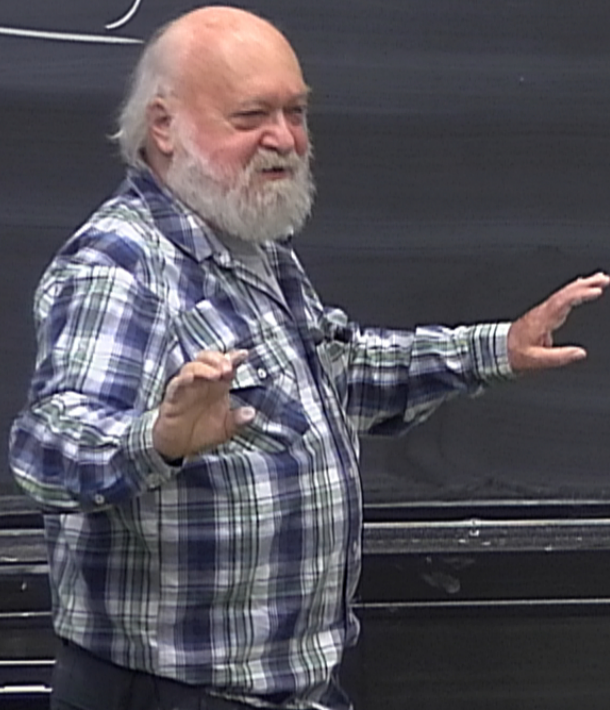
$$\langle Q \rangle = \frac{1}{Z} \int d\gamma Q e^{-\beta E(\gamma)}$$





$$r_{\alpha} = d\gamma e^{-\beta E(\gamma)}$$
$$Z = \int d\gamma e^{-\beta E(\gamma)}$$

$$\langle Q \rangle = \frac{1}{Z} \int d\gamma e^{-\beta E(\gamma)} Q(\gamma)$$





$$\langle Q \rangle = \frac{1}{Z} \int d\gamma e^{-\beta E(\gamma)} Q(\gamma)$$

$$E(\gamma) = \mathcal{H}(\gamma) = \sum_{j=1}^N \frac{p_j^2}{2M} + \sum_{j < k} V(r_j - r_k)$$



$$\sigma_z = \pm 1$$

$$\mathcal{H} = -\mu_B \sigma_z B_z$$

$$-\beta \mathcal{H} = h \sigma_z$$

$$h = \frac{\beta \mu}{k T}$$



$$\sigma_z = \pm 1 \quad \mathcal{H} = -\mu_B \sigma_z B_z$$

$$h = \frac{B_z \mu_B}{k}$$

$$r_{\pm 1} = \frac{1}{e^{-\beta \mathcal{H}(\sigma_z = \pm 1)}} = e^{+\beta \mathcal{H}}$$



$$\sigma_z = \pm 1 \quad \mathcal{H} = -\mu_B \sigma_z B_z$$

$$h = \frac{\mu_B B_z}{k T}$$

$$r_{\pm} = \frac{1}{2} e^{-\beta \mathcal{H}(\sigma_z = \pm 1)}$$
$$r = e^{+h \sigma}$$

$$- \beta \mathcal{H} = h \sigma_z$$
$$z = e^h + e^{-h} = 2 \cosh h$$



$$\langle \sigma \rangle = \sum_{\sigma=\pm 1} \sigma \frac{e^{h\sigma}}{Z} = \frac{e^h - e^{-h}}{Z}$$
$$= \tanh(h)$$



$$\langle \sigma \rangle = \sum_{\sigma=\pm 1} \sigma \frac{e^{h\sigma}}{Z} = \frac{e^h - e^{-h}}{Z}$$

$$= \tanh(h)$$

$$\begin{array}{ccc} \sigma_1 & \sigma_2 & \sigma_3 \\ \uparrow & \uparrow & \uparrow \\ 1 & 2 & 3 \end{array}$$

$$\tau_2 \tau_k = l \in \mathbb{Z}$$

$$\tau_1 \tau_2 = 1 \tau_3$$

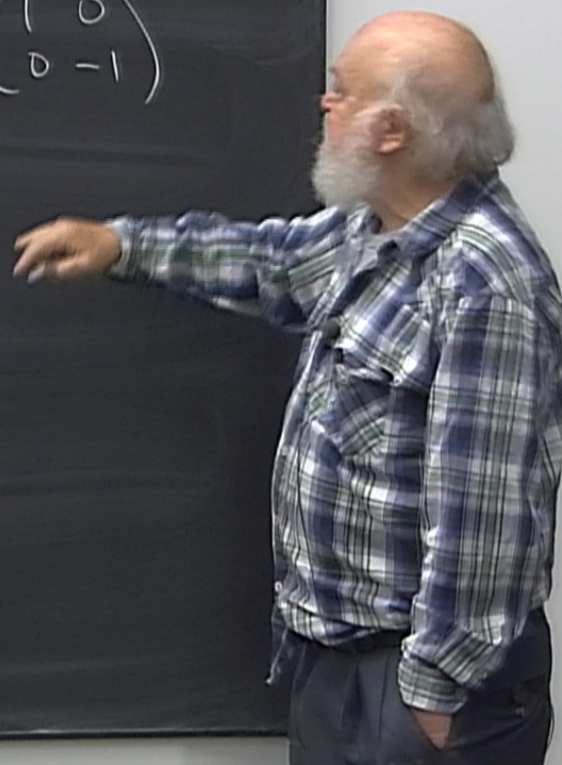


$$\tau_2^2 = \underline{1}$$

$$\underline{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$





$$\tau_2^2 = \mathbb{1}$$

$$\mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\tau_3 \tau_1 = \tau_2$$

$$= \tau_2$$

$$\tau_2 = \begin{pmatrix} 0 & -\lambda \\ \lambda & 0 \end{pmatrix}$$

$$\sigma_2 = \pm i$$

$$r_2 = 1e^{-\beta H_0}$$

$$r = e^{+\beta H_0}$$

$$p_{\sigma_1} = e^{-\beta H_0}$$



$$\tau_2^2 = \mathbb{1} \quad \mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\tau_3 \tau_1 = i \tau_2$$

$$\tau_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\tau_j \tau_k = i \epsilon_{jkl} \tau_l + \delta_{jk} \mathbb{1}$$



$$\tau_2^2 = \mathbb{1}$$

$$\mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

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$$\tau_3 \tau_1 = \tau_2$$

$$\tau_2 = \begin{pmatrix} 0 & -\lambda \\ \lambda & 0 \end{pmatrix}$$

$$\tau_j \tau_k = \lambda \epsilon_{jkl} \tau_l + \delta_{jk} \mathbb{1}$$

$$\epsilon_{123} = 1$$



$$\tau_2^2 = \mathbb{1} \quad \mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Pauli  
Spin matrices

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\tau_3 \tau_1 = i \tau_2$$

$$\tau_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\tau_j \tau_k = i \epsilon_{jkl} \tau_l + \delta_{jk} \mathbb{1}$$

$$\epsilon_{123} = 1$$



$$\tau_j \tau_k + \tau_k \tau_j = 0$$

$$j \neq k$$

$$\tau_j^2 = \mathbb{1}$$

$$\mathbb{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Pauli  
spin matrices

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\tau_3 \tau_1 = i \tau_2$$

$$\tau_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\tau_j \tau_k = i \epsilon_{jkl} \tau_l + \delta_{jk} \mathbb{1}$$

$$\epsilon_{123} = 1$$



$$\tau_j \tau_k + \tau_k \tau_j = 0$$

for  $j \neq k$

$$-\beta H = h_1 \tau_1 + h_2 \tau_2$$

$$\tau_j^2 = \mathbb{1}$$

Pauli spin matrices

$$\tau_3 \tau_1 = i \tau_2$$

$$\tau_k = \frac{1}{i} \epsilon_{jkl} \tau_l$$

$$\epsilon_{123}$$



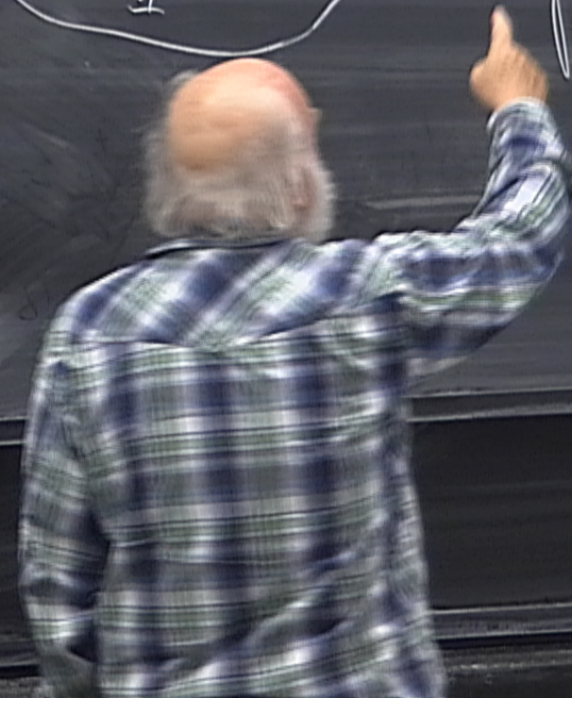
$$Q = \sum_{\alpha} \langle \alpha | Q | \alpha \rangle$$

$$\rho_{\alpha} = \frac{1}{\sum_{\alpha} e^{-\beta E_{\alpha}}}$$

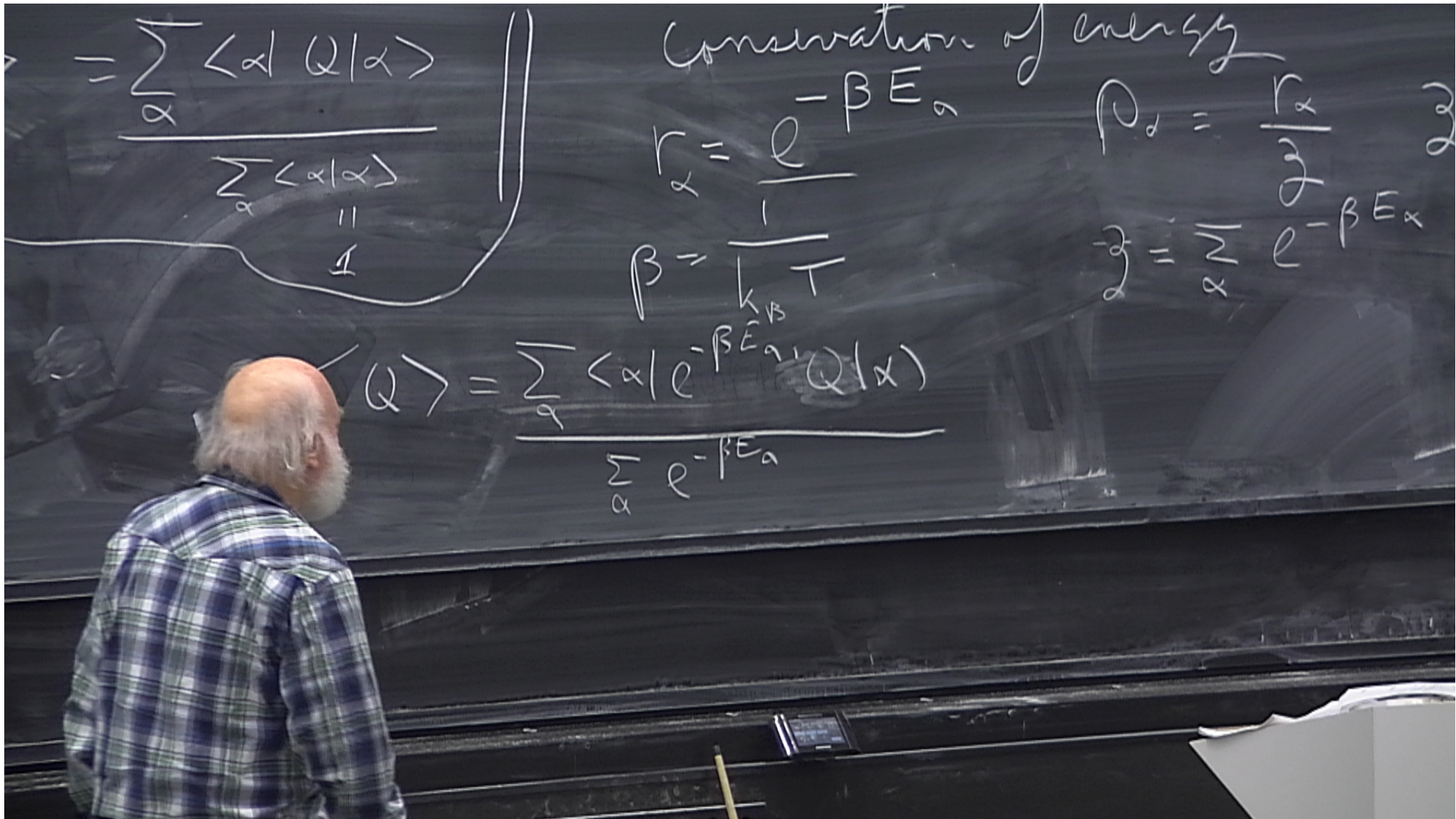
Conservation  
 $-\beta E_{\alpha}$

$$r_{\alpha} = \frac{e^{-\beta E_{\alpha}}}{\sum_{\alpha} e^{-\beta E_{\alpha}}}$$

$$\beta = \frac{1}{k_B T}$$







$$\frac{\sum_{\alpha} \langle \alpha | Q | \alpha \rangle}{\sum_{\alpha} \langle \alpha | \alpha \rangle} = 1$$

Conservation of energy

$$r_{\alpha} = \frac{e^{-\beta E_{\alpha}}}{Z}$$

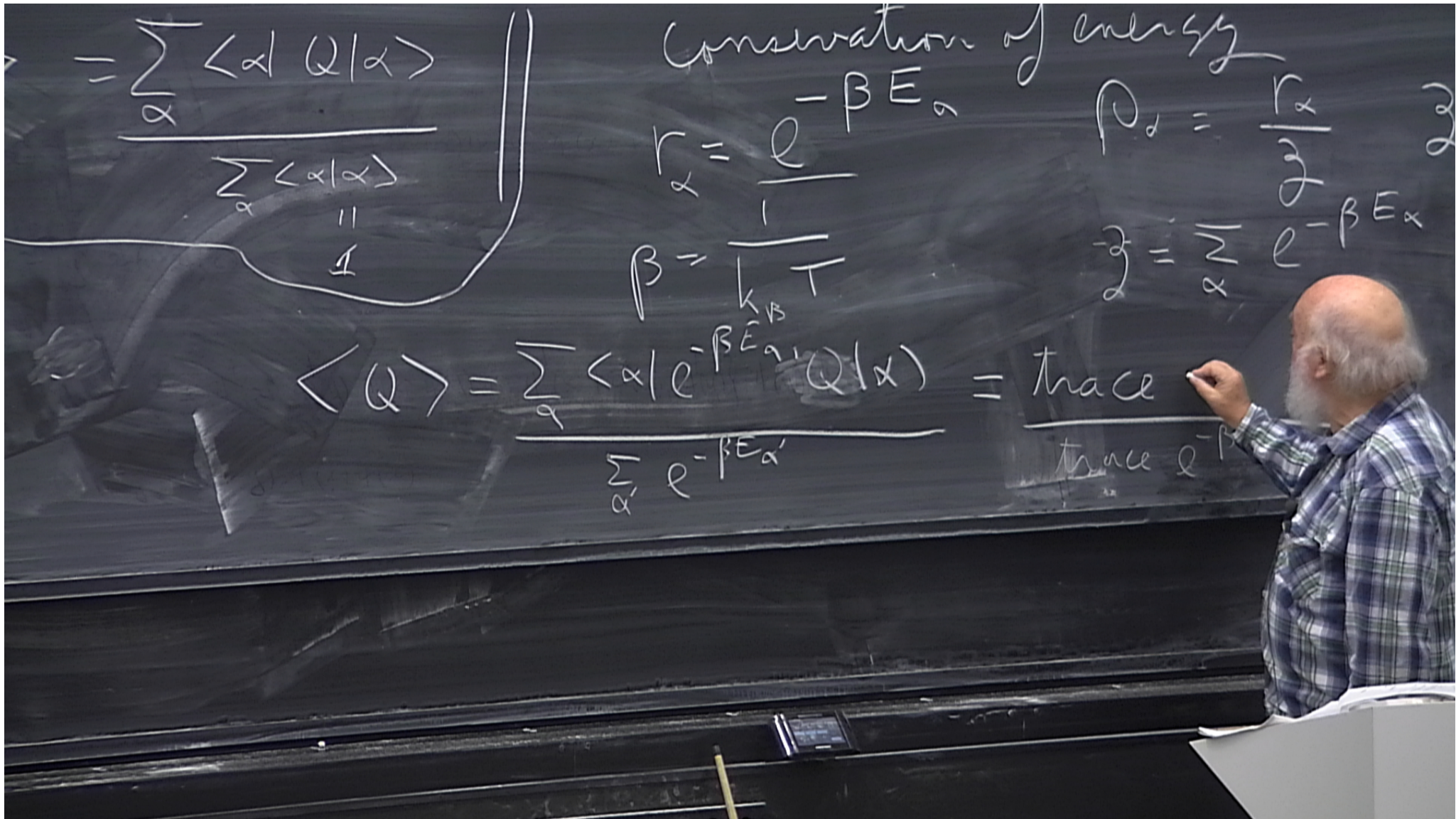
$$\rho_{\alpha} = \frac{r_{\alpha}}{z}$$

$$\beta = \frac{1}{k_B T}$$

$$z = \sum_{\alpha} e^{-\beta E_{\alpha}}$$

$$\langle Q \rangle = \frac{\sum_{\alpha} \langle \alpha | e^{-\beta E_{\alpha}} Q | \alpha \rangle}{\sum_{\alpha} e^{-\beta E_{\alpha}}}$$





$$\frac{\sum_{\alpha} \langle \alpha | Q | \alpha \rangle}{\sum_{\alpha} \langle \alpha | \alpha \rangle} = 1$$

Conservation of energy

$$r_{\alpha} = \frac{e^{-\beta E_{\alpha}}}{Z}$$

$$\rho_{\alpha} = \frac{r_{\alpha}}{Z}$$

$$\beta = \frac{1}{k_B T}$$

$$Z = \sum_{\alpha} e^{-\beta E_{\alpha}}$$

$$\langle Q \rangle = \frac{\sum_{\alpha} \langle \alpha | e^{-\beta E_{\alpha}} Q | \alpha \rangle}{\sum_{\alpha} e^{-\beta E_{\alpha}}} = \frac{\text{trace } Q e^{-\beta H}}{\text{trace } e^{-\beta H}}$$



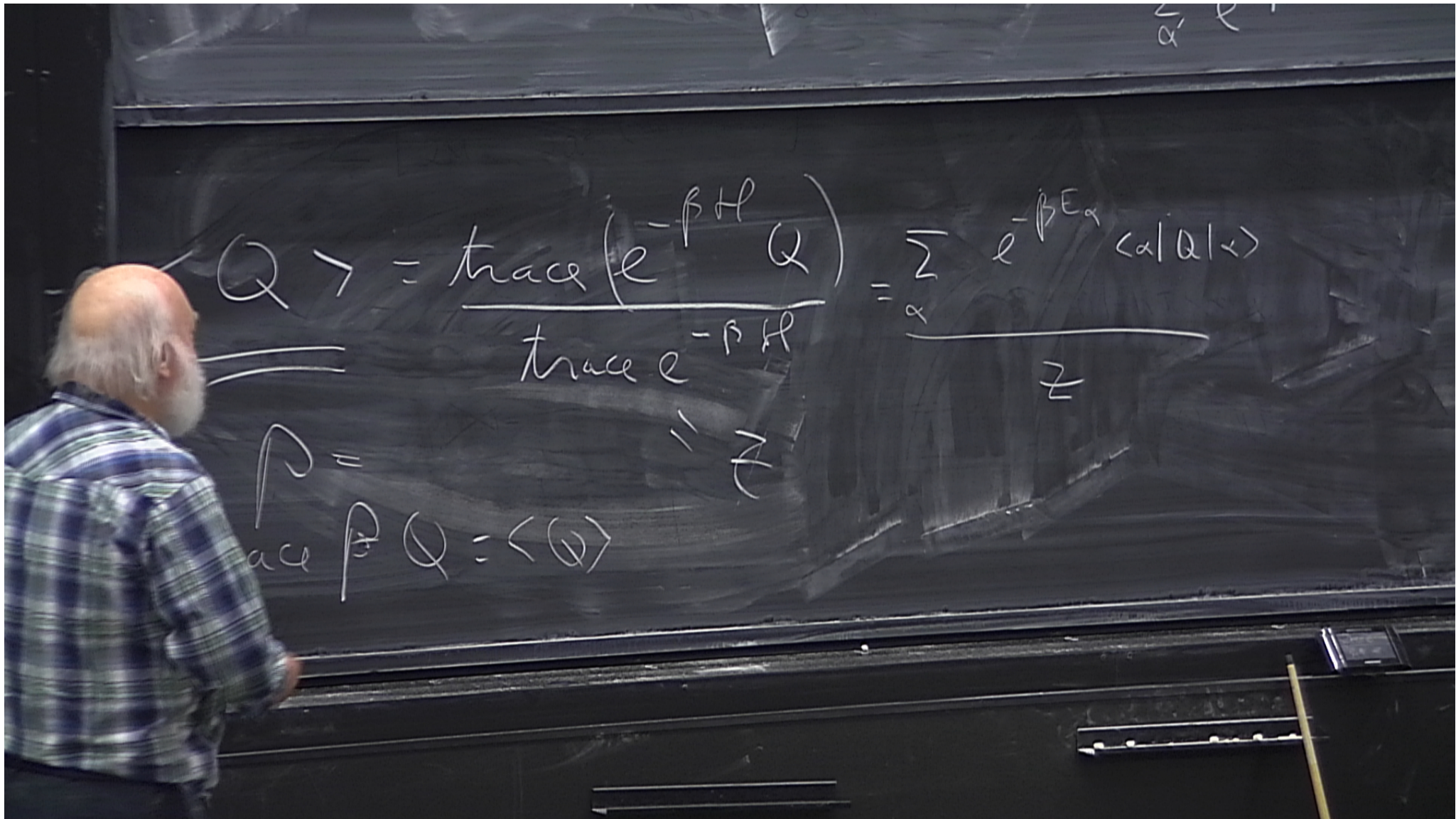
$$\langle Q \rangle = \frac{\text{trace} (e^{-\beta H} Q)}{\text{trace} e^{-\beta H}}$$



$$\langle Q \rangle = \frac{\text{trace} (e^{-\beta H} Q)}{\text{trace} e^{-\beta H}} = \frac{\sum_{\alpha} e^{-\beta E_{\alpha}} \langle \alpha | Q | \alpha \rangle}{Z}$$

$Z$



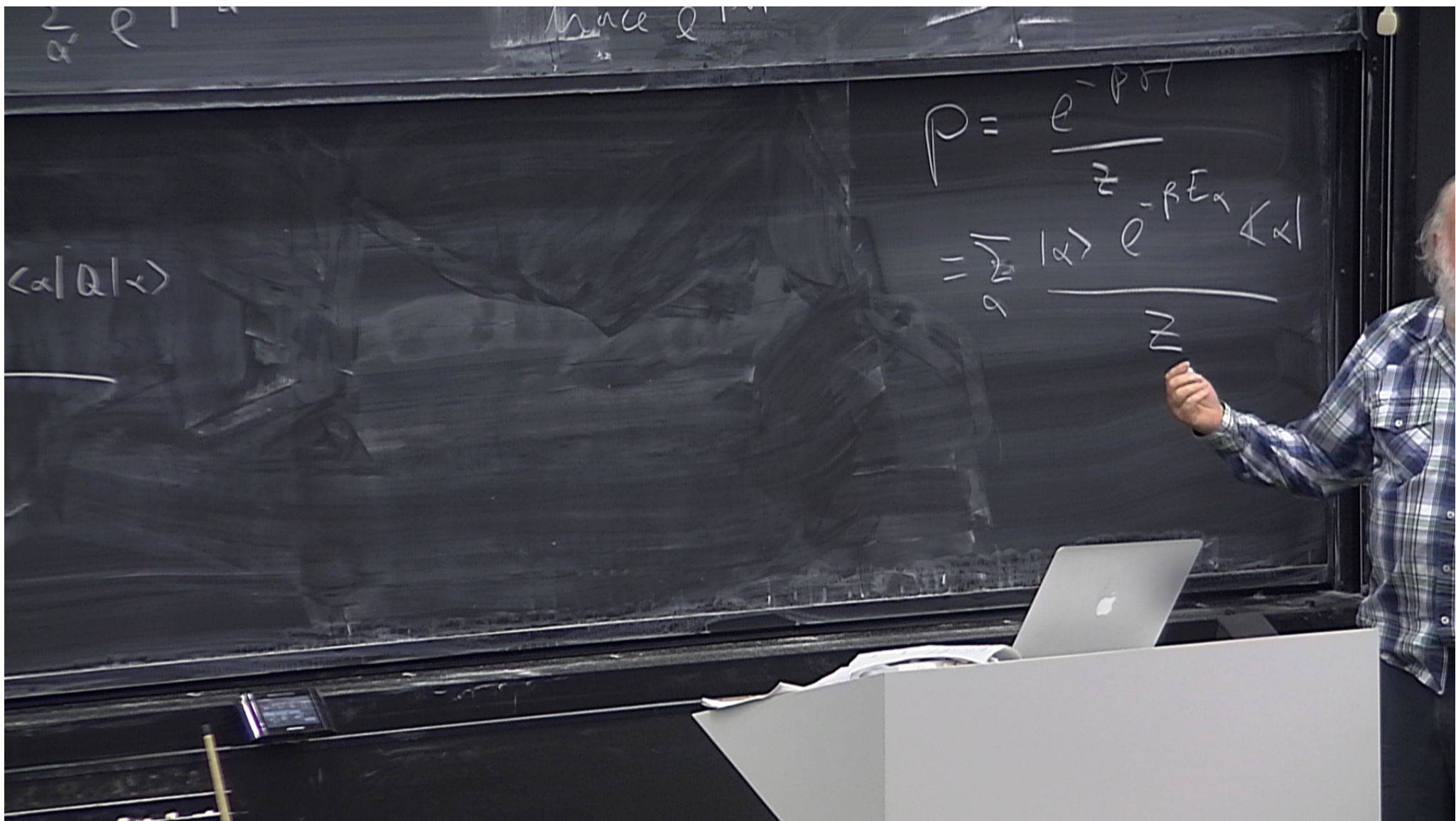


$$\langle Q \rangle = \frac{\text{trace}(e^{-\beta H} Q)}{\text{trace} e^{-\beta H}} = \frac{\sum_{\alpha} e^{-\beta E_{\alpha}} \langle \alpha | Q | \alpha \rangle}{Z}$$

$$\rho = \frac{e^{-\beta H}}{Z}$$

$$\text{trace } \rho Q = \langle Q \rangle$$







$$Z = \sum_{\text{poss.}} r_i$$

$$\tau_j \tau_k + \tau_k \tau_j = 0$$

for  $j \neq k$

$$-\beta H = (h_1 \tau_1 + h_2 \tau_2)$$

$$Z =$$

$$\tau_2^2 =$$

Pauli  
Spin

$\tau$



$$Z = \sum_{\text{poss.}} r_i$$

$$\tau_j \tau_k + \tau_k \tau_j = 0$$

for  $l \neq k$

$$-\beta H = (h_1 \tau_1 + h_2 \tau_2)$$
$$(-\beta H)^2$$

$$\tau_l^2 =$$

Pauli  
spin mat

$$\tau_3$$

$$= \lambda$$



$$Z = \sum_{\text{poss.}} r_i$$

$$\tau_j \tau_k + \tau_k \tau_j = 0$$

for  $j \neq k$

$$\begin{aligned} -\beta H &= (h_1 \tau_1 + h_2 \tau_2) \\ (-\beta H)^2 &= h_1^2 \tau_1^2 + h_2^2 \tau_2^2 \\ &\quad + h_1 h_2 (\tau_1 \tau_2 + \tau_2 \tau_1) \\ &= (h_1^2 + h_2^2) \mathbb{I} \\ &\quad + 0 \end{aligned}$$

$$\tau_j^2 = \mathbb{I}$$

Pauli spin mat

$$\tau_j \tau_k = -\tau_k \tau_j$$



$$Z = \sum_{\text{poss.}} r_i$$

$$(-\beta \mathcal{H}) = \begin{pmatrix} +h \\ -h \end{pmatrix}$$

$$Z = \text{trace } e^{-\beta \mathcal{H}}$$

$$= e^h + e^{-h} = 2 \cosh(\sqrt{h_1^2 + h_2^2})$$

$$\tau_j \tau_k + \tau_k \tau_j = 0$$

for  $j \neq k$

$$-\beta \mathcal{H} = (h_1 \tau_1 + h_2 \tau_2)$$

$$(-\beta \mathcal{H})^2 = h_1^2 \tau_1^2 + h_2^2 \tau_2^2$$

$$+ 2 h_1 h_2 (\tau_1 \tau_2 + \tau_2 \tau_1)$$

$$= (h_1^2 + h_2^2) + 0$$



$$Q = h_1 T_1 + h_2 T_2$$

$$\text{trace } Q = h_1 \text{trace } T_1 + h_2 \text{trace } T_2$$

$$Z = \sum_{\text{pos.}} r_i$$

$$\begin{pmatrix} -\beta & +h \\ & -\beta \end{pmatrix}$$

$$Z = \text{trace} \dots = e^{-\beta \sqrt{h_1^2 + h_2^2}}$$



$$x_{j+1} = r x_j - \lambda x_j^2$$

$$x \rightarrow r x / \lambda$$

$$x_{j+1} = r x_j (1 - x_j)$$

$$0 < r < 1$$



$$x_2^2 \quad 0 < r < 1 \quad x_2 \rightarrow 0 \text{ as } j \rightarrow \infty \quad x_2 \in [0, 1] \\ (0, 1)$$



$0$  as  $j \rightarrow \infty$

$$x_j \in [0, 1] \\ (0, 1)$$

$$x_j = \frac{\text{ratio of} \\ \text{pop}}{\text{max pop}}$$



$$x_{j+1} = r x_j - \Delta x_j^2$$

$\rightarrow r x / \Delta$

$$0 < r < 1$$

$$1 < r < 2$$

$$2 < r < 3.3$$

$$x_j \rightarrow 0 \text{ as } j \rightarrow \infty$$

$$x_j \rightarrow \text{const}$$

$$x_{j+1} = r x_j (1 - x_j)$$

$$x^* = r x^* (1 - x^*) \quad x^* =$$



$$x_{j+1} = r x_j - \Delta x_j^2$$

$\rightarrow r x / \Delta$

$$x_{j+1} = r x_j (1 - x_j)$$

$$x^* = r x^* (1 - x^*) \quad x^* = \frac{1}{r}$$

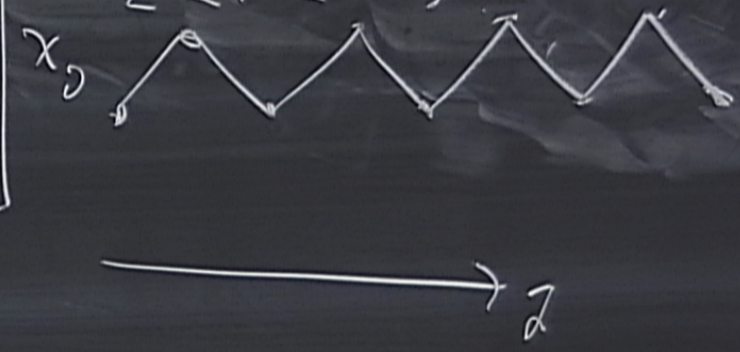
$0 < r < 1$

$1 < r < 2$

$2 < r < 3.3$

$x_j \rightarrow 0$  as

$x_j \rightarrow \text{const}$





$$x_{j+1} = r x_j - \Delta x_j^2$$

$\rightarrow r x / \Delta$

$$x_{j+1} = r x_j (1 - x_j)$$

$$x^* = r x^* (1 - x^*) \quad x^* = \frac{1}{r}$$

$0 < r < 1$

$1 < r < 2$

$2 < r < 3.3$

$x_j \rightarrow 0$  as

$x_j \rightarrow \text{const}$

