

Title: Lie Groups & Lie Algebras - Lecture 2

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URL: <http://pirsa.org/11080122>

Abstract:

Commutator $\leftarrow [\hat{t}_a, \hat{t}_b] = f_{ab}^c \hat{t}_c$

$\hat{t}_a \hat{t}_b - \hat{t}_b \hat{t}_a$ $\hat{t}_a * \hat{t}_b = f_{ab}^c \hat{t}_c$

$\hat{t}_a + (\hat{t}_b + \hat{t}_c) + \hat{t}_a + (\hat{t}_b + \hat{t}_c)$

a, b, c

• $\hat{t}_a * \hat{t}_b = -\hat{t}_b * \hat{t}_a$

• $\hat{t}_a * (\hat{t}_b * \hat{t}_c) - (\hat{t}_a * \hat{t}_b) * \hat{t}_c = \hat{t}_b * (\hat{t}_a * \hat{t}_c)$

$$b) = \hat{f}_{ab}^c \hat{t}_c$$

$$= \hat{f}_{ab}^c \hat{t}_c$$

$$\hat{t}_b * \hat{t}_a$$

$$c) -(\hat{t}_a + \hat{t}_b) * \hat{t}_c = \hat{t}_b * (\hat{t}_a * \hat{t}_c)$$

$$\hat{t}_a + (\hat{t}_b + \hat{t}_c) + \hat{t}_c * (\hat{t}_a * \hat{t}_b) + \hat{t}_b * (\hat{t}_c * \hat{t}_a) = 0$$

a, b, c

$$\hat{t}_a + (\hat{t}_b + \hat{t}_c) + t_c * (t_a + t_b) + t_b * (t_c + t_a) = 0$$

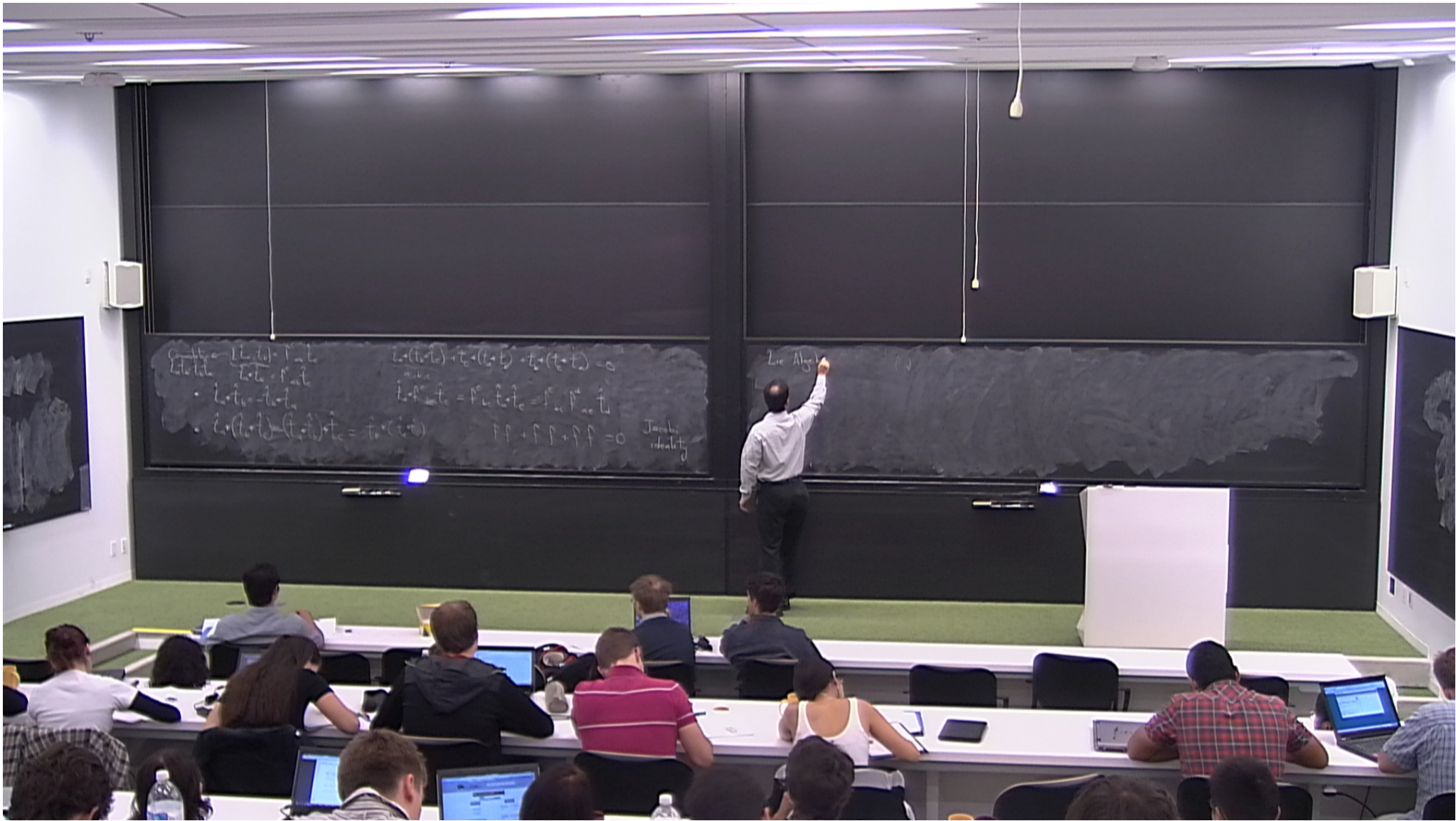
a, b, c

$$\hat{t}_a + f_{bc}^e \hat{t}_e = f_{bc}^e \hat{t}_a + \hat{t}_e = f_{bc}^e f_{ae}^d \hat{t}_d$$

$$t_c = \hat{t}_b * (\hat{t}_a + \hat{t}_c)$$

$$f f + f f + f f = 0$$

Jacobi
identity.



Poincaré Group.

- Lorentz Transformations
Translations.

$$\bullet (\Lambda, a) \longrightarrow \mathbb{U}(\Lambda, a)$$

- Infinitesimal transformations: ϵ^μ

$$x'^\mu = \Lambda^\mu_\nu x^\nu + a^\mu$$

$$|\psi\rangle \rightarrow \mathbb{U}(\Lambda, a)|\psi\rangle$$

$$\bar{K}^\mu = \delta^\mu_\nu + \omega^\mu_\nu$$

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$$\bar{K}^\mu = \int \mathcal{H}_\nu + \omega^\mu_\nu$$

Incaré Group.

transformations

→ $\mathbb{U}(1)$

1) transformations:

$$X'^{\mu} = \Lambda^{\mu}_{\nu} X^{\nu} + a^{\mu}$$

$$|\psi\rangle \rightarrow \mathbb{U}(\Lambda, a) |\psi\rangle$$

$$\in \mathbb{U}(1)$$

$$\bar{X}^{\mu} = \delta^{\mu}_{\nu} + \omega^{\mu}_{\nu}$$

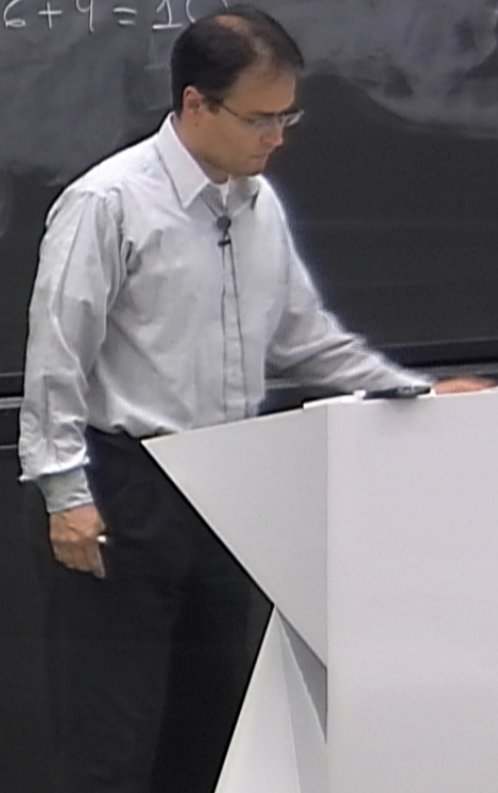
$$\eta_{\mu\nu} \Lambda^{\mu}_{\rho} \Lambda^{\nu}_{\sigma} = \eta_{\rho\sigma}$$

$$\eta_{\rho\sigma} (\delta^{\mu}_{\rho} + \omega^{\mu}_{\rho}) (\delta^{\nu}_{\sigma} + \omega^{\nu}_{\sigma}) = \eta_{\rho\sigma}$$

$$\eta_{\rho\sigma} + \omega_{\nu\rho}$$

$$\Lambda = \text{diag}(-1, 1, 1, 1)$$

$\omega_{\mu\nu} = -\omega_{\nu\mu}$ \times of independent parameters $\frac{4 \times 3}{2} = 6$ dimension of the Lorentz group = 6.
Translation \rightarrow 4 parameters $\Rightarrow \dim(\text{Poincaré}) = 6 + 4 = 10$



$$\eta = \text{diag}(-1, 1, 1, 1)$$

- $\omega_{\mu\nu} = -\omega_{\nu\mu}$ \times of independent parameters. $\frac{4 \times 3}{2} = 6$ dimension of the Lorentz group =
- Translation \rightarrow 4 parameters $\Rightarrow \dim(\text{Poincaré}) = 6 + 4 = 10$
 - \rightarrow Abelian $\sqcup(\mathbb{1}, a)$
- Lorentz $\sqcup(\Lambda, 0)$

Unitary Operator

$$U(\mathbb{1} + \omega, e) = \mathbb{1} + i\omega_{\mu\nu} J^{\mu\nu} - i\epsilon_{\mu} P^{\mu}$$

$J^{\mu\nu}, P^{\mu}$ Hermitian. \rightarrow convenient.

$-i \epsilon_n P^0$
→ convenient

Physical Meaning.

$$U(\mathbb{1}, a) U(\mathbb{1}, b) = U(\mathbb{1}, a+b)$$
$$b = a \quad U(\mathbb{1}, 2a)$$

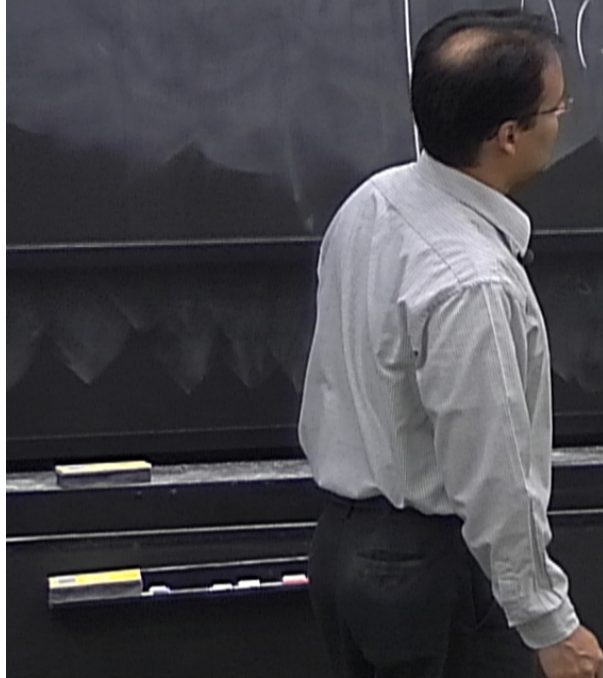
$$U(\mathbb{1}, a) = \lim_{N \rightarrow \infty} U(\mathbb{1}, \frac{a}{N})^N = \lim_{N \rightarrow \infty} \left(\mathbb{1} - i \frac{a}{N} P^0 \right)^N = e^{-i a P^0}$$

$$a^\mu = (t, 0, 0, 0)$$

$$-itP^0$$

$$e$$

$$P^0 = H \text{ Hamiltonian.}$$



$$\lim_{N \rightarrow \infty} \left(\frac{a}{N} \right)$$

$$\vec{P} = (P^1, P^2, P^3) \rightarrow \text{Momentum Operators}$$

$$\vec{J} = (J^{23}, J^{31}, J^{12}) \rightarrow \text{Angular momentum operators}$$

$$U(\mathbb{1}, a) = \lim_{N \rightarrow \infty} U\left(\mathbb{1}, \frac{a}{N}\right) = \lim_{N \rightarrow \infty} \left(\mathbb{1} - i \frac{a_\mu P^\mu}{N}\right) = e$$

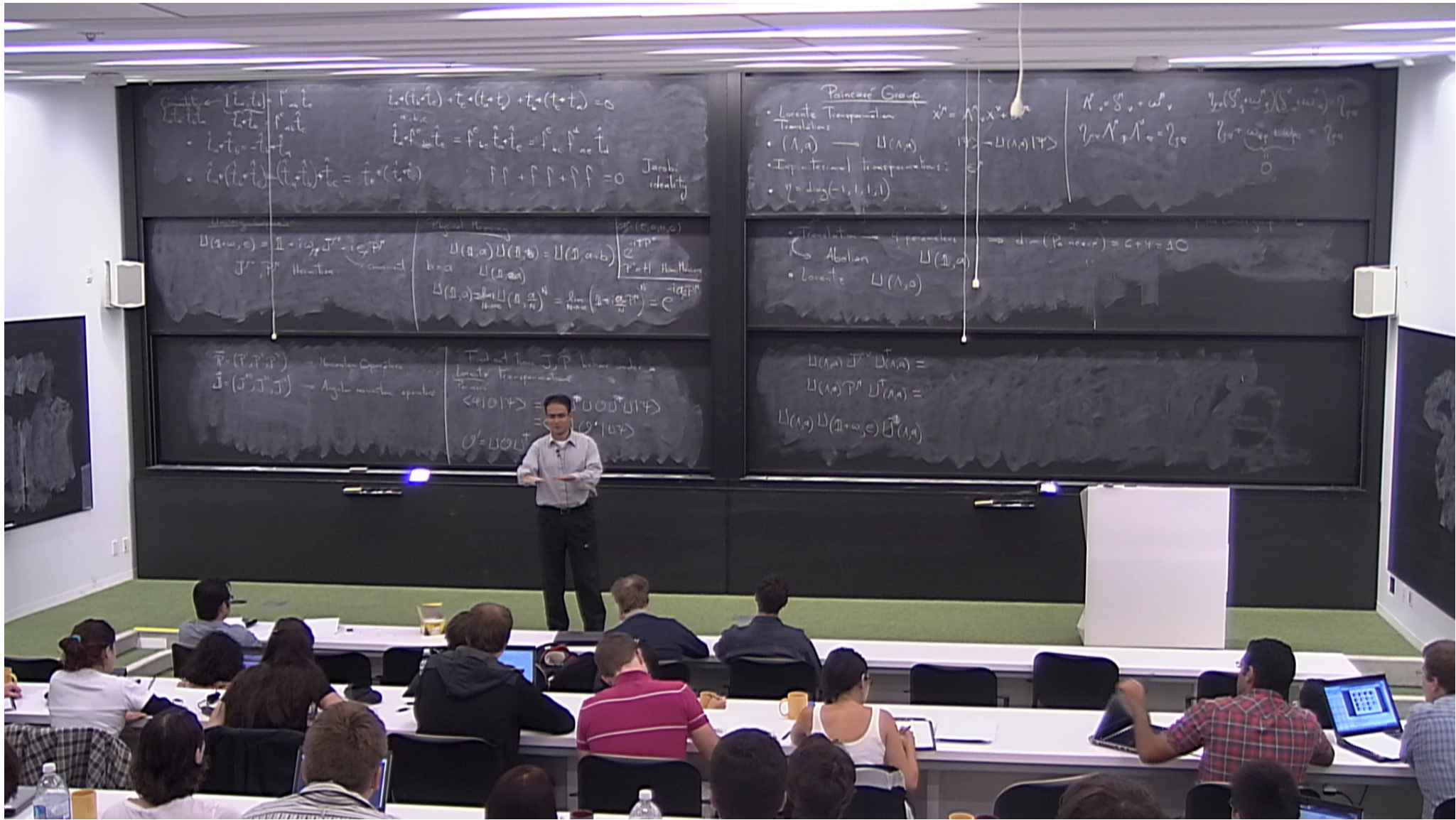
→ Momentum Operators

→ Angular momentum operators

Find out how J, P behave under a
Lorentz transformation
Poincaré

$$\begin{aligned} \langle \varphi | \mathcal{O} | \psi \rangle &= \langle \varphi | U^\dagger U \mathcal{O} U^\dagger U | \psi \rangle \\ &= \langle U\varphi | \mathcal{O}' | U\psi \rangle \end{aligned}$$

$$\mathcal{O}' = U \mathcal{O} U^\dagger$$



$$\mathbb{U}(\lambda, a) \mathbb{J}^{\omega} \mathbb{U}^{\dagger}(\lambda, a) =$$

$$\mathbb{U}(\lambda, a) \mathbb{P}^{\lambda} \mathbb{U}^{\dagger}(\lambda, a) =$$

$$\mathbb{U}(\lambda, a) \mathbb{U}(\mathbb{1} + \omega, \epsilon) \mathbb{U}^{-1}(\lambda, a) = \mathbb{U}(\circ, \circ)$$

$$U(\Lambda, a) \left(\mathbb{I} + i\omega_{\mu\nu} J^{\mu\nu} - i\epsilon_{\mu} P^{\mu} \right) U(\Lambda, a)^{-1} = U(\Lambda, a)$$

$$\mathbb{U}(\lambda, a) \mathbb{J}^{\mu\nu} \mathbb{U}^\dagger(\lambda, a) =$$

$$\mathbb{U}(\lambda, a) \mathcal{P}^\mu \mathbb{U}^\dagger(\lambda, a) =$$

$$\mathbb{U}(\lambda, a) \mathbb{U}(\mathbb{1} + \omega, e) \mathbb{U}^{-1}(\lambda, a) = \mathbb{U}(\lambda, a)$$

$$U(\Lambda, a) \left(\mathbb{I} + i\omega_{\mu\nu} J^{\mu\nu} - i\epsilon_{\mu} P^{\mu} \right) U^{-1}(\Lambda, a) = U \left(\mathbb{I} + \Lambda\omega\Lambda^{-1}, \Lambda\epsilon - \Lambda\omega\Lambda^{-1}a \right)$$

Ex:
$$U(\Lambda, a) J^{\mu\nu} U^\dagger(\Lambda, a) = \Lambda_\rho^\mu \Lambda_\sigma^\nu (J^{\rho\sigma} - a^\rho P^\sigma + a^\sigma P^\rho)$$

$$U(\Lambda, a) P^\mu U^\dagger(\Lambda, a) = \Lambda_\rho^\mu P^\rho$$

Find the Poincaré Algebra

$$\Lambda = \mathbb{1} + \tilde{\omega} \quad a^\mu = \tilde{\epsilon}^\mu$$

Poincaré Algebra

$$i [J^{\mu\nu}, J^{\rho\sigma}] = \eta^{\nu\rho} J^{\mu\sigma} - \eta^{\mu\rho} J^{\nu\sigma} - \eta^{\sigma\mu} J^{\rho\nu} + \eta^{\sigma\nu} J^{\rho\mu}$$

$$i [P^{\mu}, J^{\rho\sigma}] = \eta^{\mu\rho} P^{\sigma} - \eta^{\mu\sigma} P^{\rho}$$

$$[P^{\mu}, P^{\nu}] = 0$$

Ex:
$$\Lambda(\Lambda, a) J^{\mu\nu} \Lambda^{\dagger}(\Lambda, a) = \Lambda_{\rho}^{\mu} \Lambda_{\sigma}^{\nu} (J^{\rho\sigma} - a^{\rho} P^{\sigma} + a^{\sigma} P^{\rho})$$

$$\Lambda(\Lambda, a) P^{\mu} \Lambda^{\dagger}(\Lambda, a) = \Lambda_{\rho}^{\mu} P^{\rho}$$

Find the Poincaré Algebra

$$\Lambda = \mathbb{1} + \tilde{\omega} \quad a^{\mu} = \tilde{\epsilon}^{\mu}$$

Poincaré Algebra

$$i [J^{\mu\nu}, J^{\rho\sigma}] = \eta^{\nu\rho} J^{\mu\sigma} - \eta^{\mu\rho} J^{\nu\sigma} - \eta^{\sigma\mu} J^{\rho\nu} + \eta^{\sigma\nu} J^{\rho\mu}$$

$$i [P^{\mu}, J^{\rho\sigma}] = \eta^{\mu\rho} P^{\sigma} - \eta^{\mu\sigma} P^{\rho}$$

$$[P^{\mu}, P^{\nu}] = 0$$

$$L^2, P, J=0$$

Label 1-Particle states: P^M
entz:

$$P^M |p, \sigma\rangle = p^M |p, \sigma\rangle$$

Any other label needed

$$[P^\mu, P^\nu] = 0$$

Label 1-Particle states: P^μ

• Lorentz:

$$U(\Lambda) |P, \sigma\rangle$$

$$U(\Lambda^{-1}, 0) P^\mu U(\Lambda^{-1}, 0)^{-1} = \Lambda^\mu_\nu P^\nu \quad \leftarrow \text{Any other } \Lambda$$

$$P^\mu (U(\Lambda) |P, \sigma\rangle) = (P^\mu U(\Lambda)) |P, \sigma\rangle$$

$$U(\Lambda) \left(U^{-1}(\Lambda) P^\mu U(\Lambda) = \Lambda^\mu_\nu P^\nu \right) P^\mu U(\Lambda)$$

$$= U(\Lambda) \Lambda^\mu_\nu P^\nu |P, \sigma\rangle = U(\Lambda) (\Lambda^\mu_\nu P^\nu) |P, \sigma\rangle = \Lambda^\mu_\nu P^\nu (U(\Lambda) |P, \sigma\rangle)$$

$$[P^\mu, P^\nu] = 0$$

Label 1-Particle states: P^μ

• Lorentz: $U(\Lambda) |P, \sigma\rangle$

$$P^\mu (U(\Lambda) |P, \sigma\rangle) = (P^\mu U(\Lambda)) |P, \sigma\rangle$$

$$= U(\Lambda) \Lambda^\mu_\nu P^\nu |P, \sigma\rangle = U(\Lambda) (\Lambda^\mu_\nu P^\nu) |P, \sigma\rangle = \Lambda^\mu_\nu P^\nu (U(\Lambda) |P, \sigma\rangle)$$

$U(\Lambda) (U^{-1}(\Lambda) P^\mu U(\Lambda) = \Lambda^\mu_\nu P^\nu)$

$U(\Lambda^{-1}, 0) P^\mu U^{-1}(\Lambda^{-1}, 0) = \Lambda^\mu_\nu P^\nu$

$P^\mu |P, \sigma\rangle = p^\mu |P, \sigma\rangle$

Any other...

Conclusion:
$$U(\Lambda) |p, \sigma\rangle = \sum_{\sigma'} C_{\sigma, \sigma'}(\Lambda, p) |p, \sigma'\rangle$$

Trick: Note that $p^\mu p_\mu = p^2 = -M^2$ First label.
 $p^2 < 0$ \rightarrow Mass
 $\text{sign}(p^0)$ is Lorentz invariant.

$$L^{\mu\nu} p_{\nu} = 0$$

Let's consider a reference vector K^{μ}

$$K^{\mu} K_{\mu} = -M^2$$

$$p^{\mu} = L^{\mu}{}_{\nu}(p) K^{\nu}$$

→ 4x4 matrix
that leaves η invariant.

$$p^\mu p_\mu = 0$$

Consider a reference vector \underline{k}^μ

$$p'^\mu = L^\mu{}_\nu(p) k^\nu$$

\rightarrow 4×4 matrix
that leaves η invariant.

$$k^2 = -M^2$$

Def: $|p, \sigma\rangle \equiv U(L(p)) |k, \sigma\rangle$

Apply

Trick 2:
$$U(\Lambda)|p, \sigma\rangle = U(\Lambda L(p))|k, \sigma\rangle = U(L^{-1}(\Lambda p)) U(L(p)) U(\Lambda L(p))|k, \sigma\rangle$$

Let's consider all Lorentz transformations that keep k invariant.

$$= U(L(\Lambda p)) U(L^{-1}(\Lambda p) \Lambda L(p))|k, \sigma\rangle$$

$$\begin{aligned} L(p)k &= p & L(\Lambda p)k &= \Lambda p \\ \Lambda L(p)k &= \Lambda p & L^{-1}(\Lambda p) \Lambda L(p)k &= k \end{aligned}$$

Trick 2:
$$U(\Lambda)|p, \sigma\rangle = U(\Lambda L(p))|k, \sigma\rangle = \underbrace{U(L^{-1}(p))}^{-1} U(L(p)) U(\Lambda L(p))|k, \sigma\rangle$$

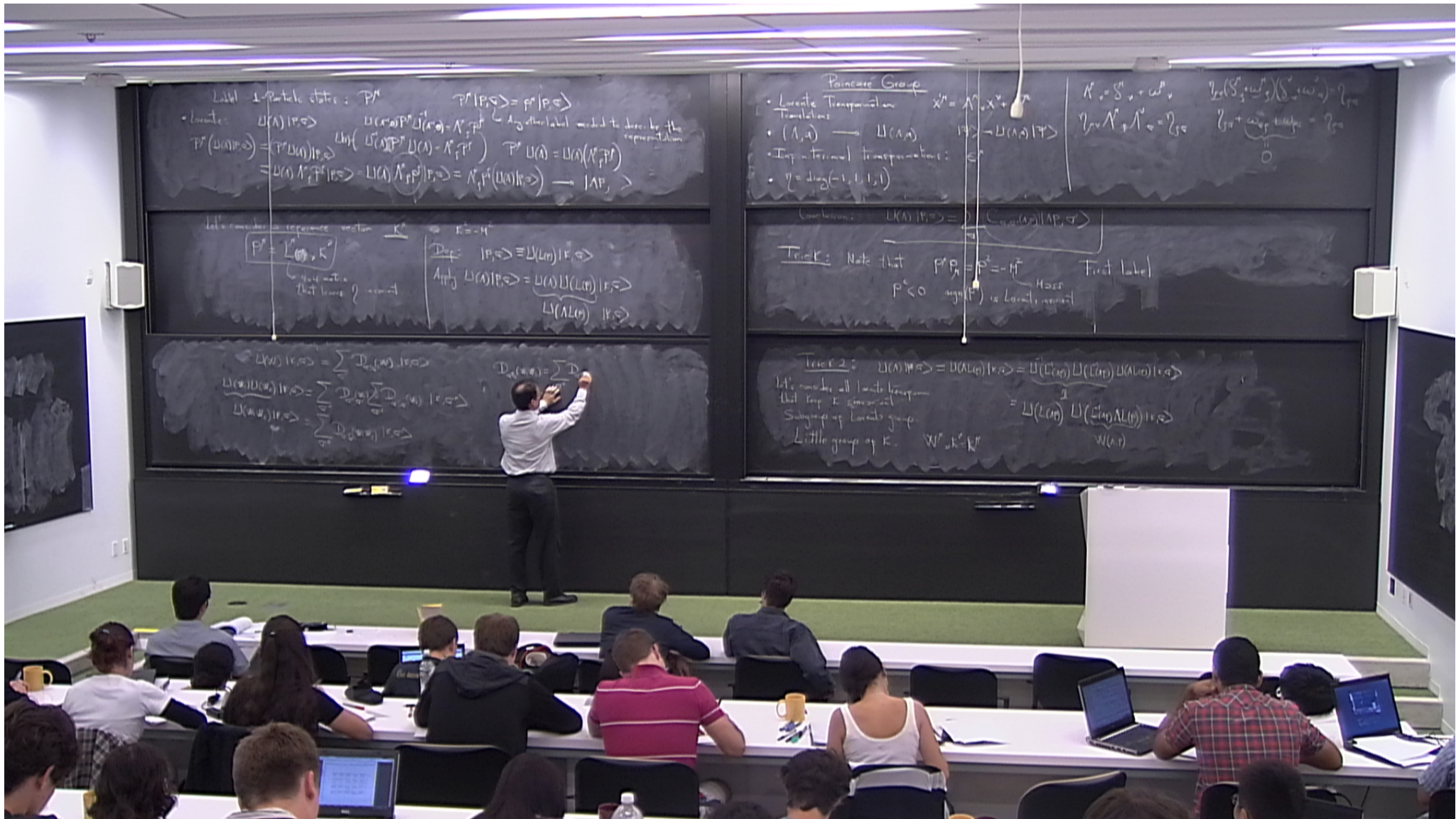
Let's consider all Lorentz transformations that keep k invariant.

Subgroup of Lorentz group.

Little group of k . $W^\mu{}_\nu k^\nu = k^\mu$

$$= U(L(p)) \underbrace{U(L^{-1}(p)\Lambda L(p))}^{\mathbb{1}} U(L(p))|k, \sigma\rangle$$





Let $|k\rangle$ be a particle state: $P^\mu |k\rangle = k^\mu |k\rangle$
 • Lorentz: $U(\Lambda) |k\rangle = U(\Lambda) P^\mu U^\dagger(\Lambda) U(\Lambda) |k\rangle = U(\Lambda) P^\mu |k\rangle = U(\Lambda) k^\mu |k\rangle = (U(\Lambda) k)^\mu |k\rangle$
 $P^\mu U(\Lambda) |k\rangle = U(\Lambda) P^\mu |k\rangle = U(\Lambda) k^\mu |k\rangle = (U(\Lambda) k)^\mu |k\rangle \rightarrow |U(\Lambda)k\rangle$

Let's consider a reference vector $k^\mu = (k^0, \mathbf{k})$
 $P^\mu = L^\mu_{\nu} k^\nu$
 Def: $|k\rangle \equiv U(L) |k\rangle$
 $U(\Lambda) |k\rangle = U(\Lambda) U(L) |k\rangle = U(\Lambda L) |k\rangle$

$U(\Lambda) |k\rangle = \sum_i D_{ij}(\Lambda) |k_i\rangle$
 $U(\Lambda) U(L) |k\rangle = \sum_i D_{ij}(\Lambda) \sum_j D_{jk}(L) |k\rangle$
 $U(\Lambda L) |k\rangle = \sum_i D_{ij}(\Lambda L) |k\rangle$

Poincaré Group
 • Lorentz Transformation $x^\mu \rightarrow \Lambda^\mu_{\nu} x^\nu + a^\mu$
 • Translation $(\Lambda, a) \rightarrow U(\Lambda, a) |k\rangle = U(\Lambda) U(a) |k\rangle$
 • Infinitesimal transformations: ϵ^μ
 • $\eta = \text{diag}(-1, 1, 1, 1)$

Conclusion: $U(\Lambda) |k\rangle = U(\Lambda) U(L) |k\rangle$
 Trick: Note that $P^\mu P_\mu = -M^2$ (Mass)
 $P^\mu < 0$ implies $|k\rangle$ is Lorentz invariant

Trick 2: $U(\Lambda) |k\rangle = U(\Lambda L) |k\rangle = U(L) U(\Lambda) |k\rangle$
 Let's consider all Lorentz transformations that keep k invariant
 Subgroup of Lorentz group
 Little group of k : $W^\mu, k^\mu = W^\mu$

$$\underbrace{U(\Lambda L(\rho))}_{K, \sigma} \rightarrow$$

$$\sum_{\sigma'} D_{\sigma', \sigma}(W) |K, \sigma\rangle$$

$$D_{\sigma'', \sigma'}(W_1) \sum_{\sigma} D_{\sigma, \sigma}(W_2) |K, \sigma''\rangle$$

$$D_{\sigma'', \sigma}(W_1 W_2) |K, \sigma''\rangle$$

$$D_{\sigma'', \sigma}(W_1, W_2) = \sum_{\sigma'} D_{\sigma'', \sigma'}(W_1) D_{\sigma', \sigma}(W_2)$$

$$D(W_1, W_2) = D(W_1) D(W_2)$$

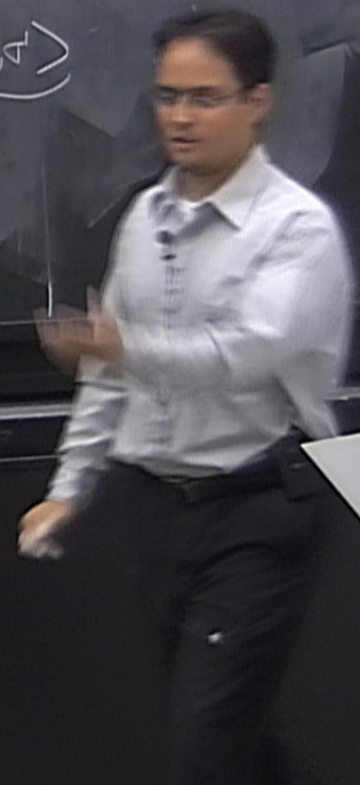
D are \wedge represent of the little group's
Matrix.

Conclusion:

$$U(\Lambda) |P, \sigma\rangle = \sum_{\sigma'} C_{\sigma, \sigma'}(\Lambda, P) | \Lambda P, \sigma' \rangle$$

$$W(\Lambda, P) = L^{-1}(\Lambda P) \Lambda L(P)$$

$$U(\Lambda) |P, \sigma\rangle = U(L(\Lambda P)) \underbrace{U(W(\Lambda, P))}_{\sum_{\sigma'} D_{\sigma' \sigma}(W(\Lambda, P))} |K, \sigma\rangle = \sum_{\sigma'} D_{\sigma' \sigma}(W(\Lambda, P)) \underbrace{U(\Lambda P)}_{\sum_{\sigma'} D_{\sigma' \sigma}(W(\Lambda, P))} |K, \sigma'\rangle$$



Conclusion: $U(\lambda) |P, \sigma\rangle = \sum_{\sigma'} C_{\sigma, \sigma'}(\lambda, P) |P, \sigma'\rangle$ $C_{\sigma, \sigma'} = D_{\sigma', \sigma}(W)$

$W(P, \lambda) = L^{-1}(\lambda P) \Lambda L(P)$

$|P, \sigma\rangle = U(L(\lambda P)) U(W(\lambda, P)) |K, \sigma\rangle = \sum_{\sigma'} D_{\sigma', \sigma}(W(\lambda, P)) \underbrace{U(\lambda P) |K, \sigma\rangle}_{|P, \sigma'\rangle}$

$\sum_{\sigma'} D_{\sigma', \sigma}(W(\lambda, P)) |K, \sigma'\rangle$



$$\sum_{\sigma'} D_{\sigma' \sigma}(W(\Lambda, P)) |K, \sigma'\rangle$$

$$|K, \sigma\rangle$$

Trick 2:

consider all Lorentz transform
 keep K invariant.
 subgroup of Lorentz group.

little group of K .

$$W^\mu{}_\nu K^\nu = K^\mu$$

Case 1: Massive Particles

$$p^2 = -M^2 \quad M > 0$$

$$K = (M, 0, 0, 0)$$

Little group
 $SO(3)$

Case 2: Massless Particles

$$U(W(\lambda, \sigma)) |k, \sigma\rangle \quad |k, \sigma\rangle$$

Case 1: Massive Particles

$$p^2 = -M^2 \quad M > 0$$

$$K = (M, 0, 0, 0)$$

Little group
SO(3)

$$\dim = 2\lambda + 1$$

$$\lambda = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$$

$$(M, \lambda) \rightarrow S^m$$

Form

p.

$$W^\mu \nu k^\nu = k^\mu$$

Case 2: Massless Particles

$$p^2 = 0 \quad K = (1, 0, 0, 1)$$

Little group: Rigid motions in \mathbb{R}^2

$$\sigma = \text{Helicity} \quad \sigma = 0, \pm \frac{1}{2}, \pm 1, \dots$$

$$J_3 |k, \sigma\rangle = \sigma |k, \sigma\rangle$$