Title: Effects of Disorder on Superfluids

Date: Aug 09, 2011 03:30 PM

URL: http://pirsa.org/11080073

Abstract:

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Effects of Disorder on Superfluids

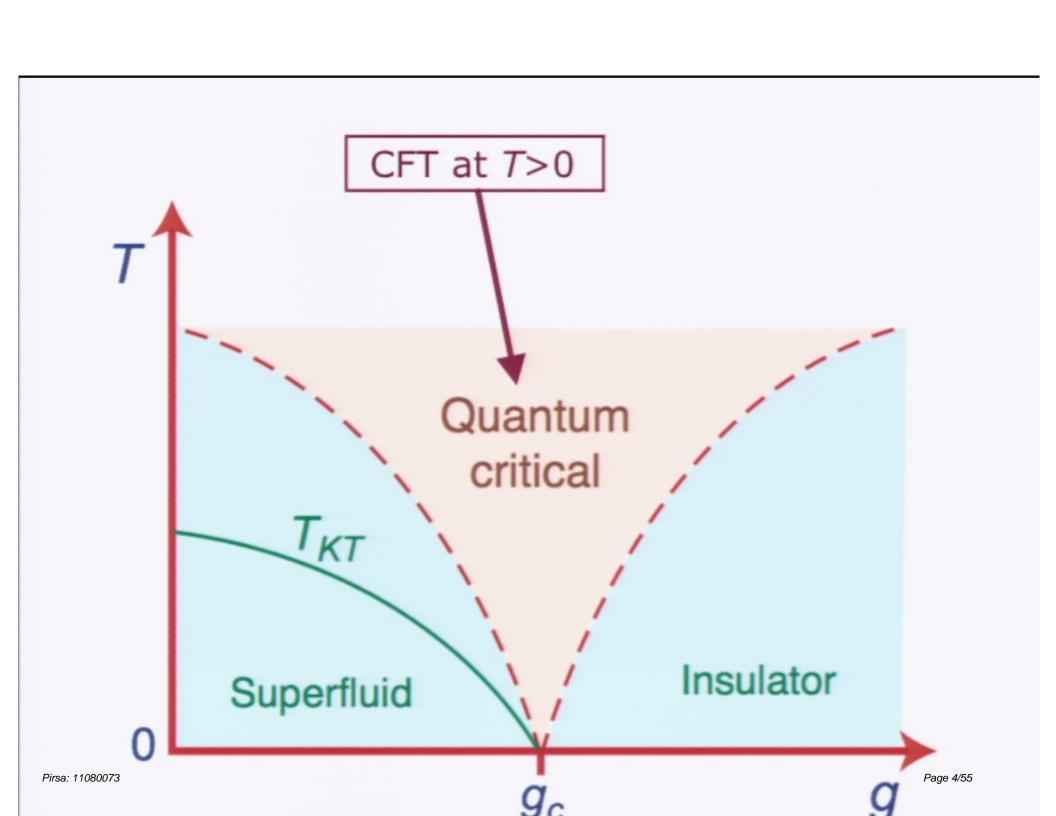
Nilesh Tripuraneni

Superconductivity/ Superfluidity

Landau/Ginzburg $\iff F = \alpha (T - T_c) |\varphi|^2 + \beta |\varphi|^4 + ...$

Broken U(1) symmetry

B.C.S. Theory ⇒ fermions ⇒ Cooper pair ⇒ condense into boson superfluid



AdS/CFT Correspondence

N=4 SYM in 3+1 \Leftrightarrow IIB on $AdS_5 \times S^5$

The precise statement:

$$\left\langle e^{\int_{d^4x\phi_0(\vec{x})O(\vec{x})}\right\rangle_{CFT} = Z_{String} \left[\phi(\vec{x}, z=0) \equiv \phi_0(\vec{x})\right]$$

Generalized to "Gauge/Gravity" Duality

The Dictionary

rongly Coupled Field Theory \Leftrightarrow Weak Gravity

$$_{\text{String}} \left[\phi(\vec{x}, z = 0) = \phi_0(\vec{x}) \right] \approx e^{-S_{classical}} \left[B.C. \rightarrow \phi_0(\vec{x}) \right]$$

Fields in AdS

⇔ Local Operators in CFT

Spin ⇔ Spin

Mass

⇔ Scaling Dimension

Black Holes

⇔ Thermodynamics

Gauged Symmetries

Global Symmetries

Gravity Dual to Superconductor/ Superfluid

Einstein-Maxwell-charged scalar

$$= \int \sqrt{-g} d^4 x \left(R + \frac{6}{L^2} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \left| \nabla_{\mu} \psi - i q A_{\mu} \psi \right|^2 - m^2 |\psi|^2 \right)$$

$$m^2_{\it eff} = m^2 + q^2 g^{tt} A_t^2 \Leftrightarrow {\it Lower temperature} \Rightarrow g^{tt} {\it closer to double horizon} \Rightarrow {\it more instability}$$

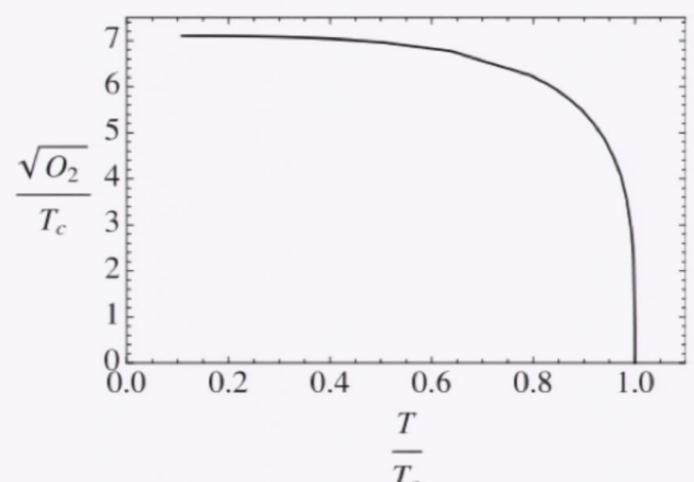
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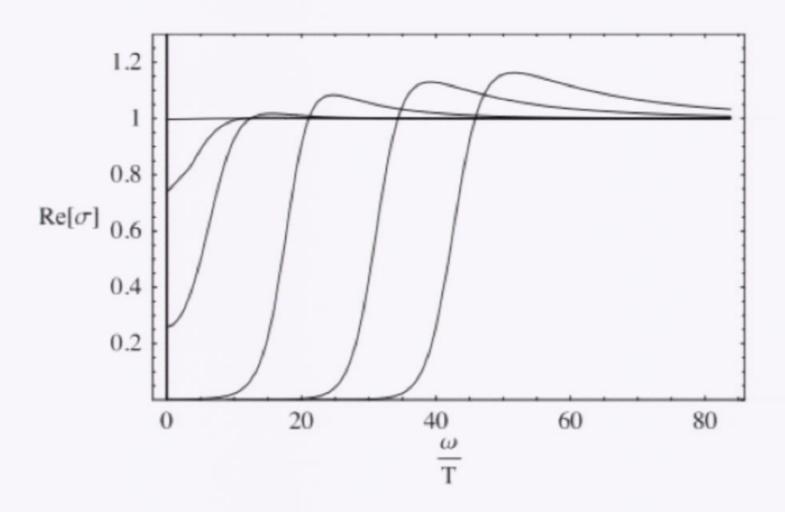
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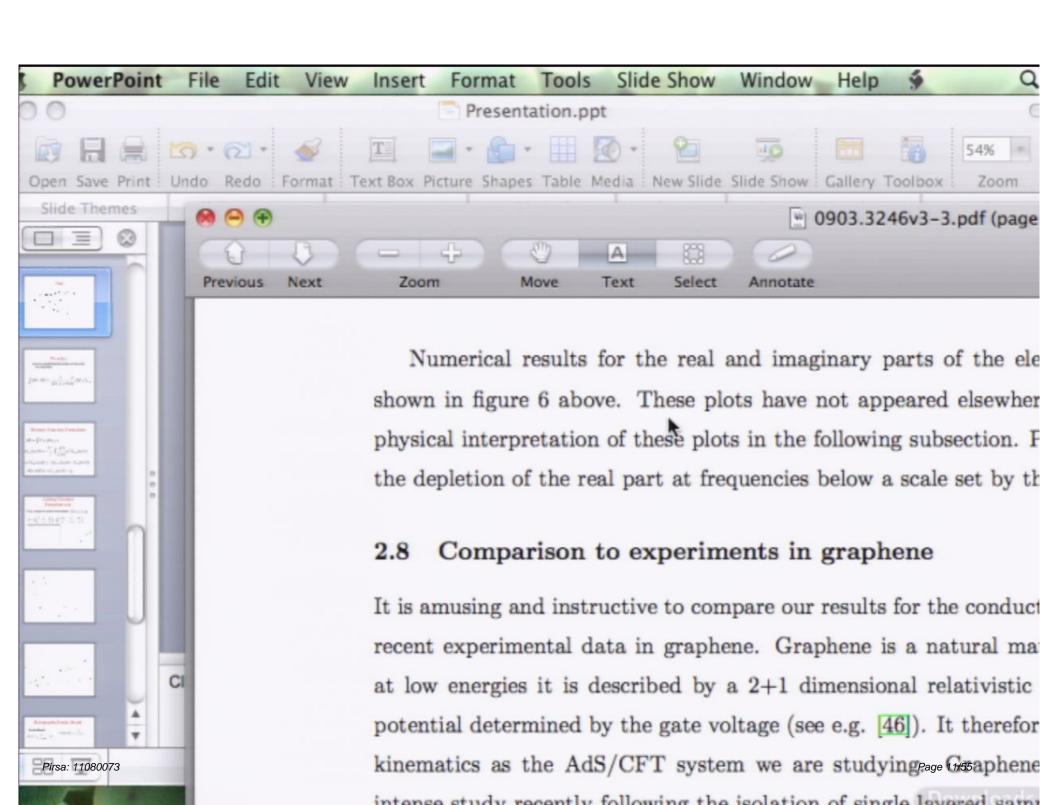
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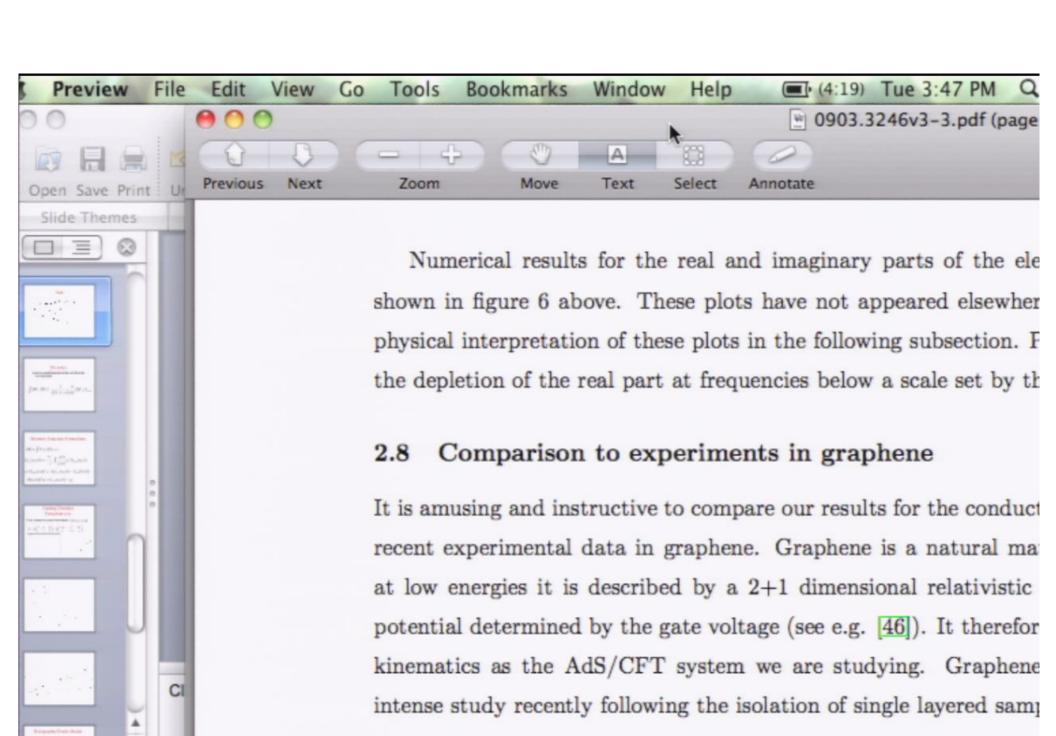
Condensation



Gap







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L=1 and furthermore to scale the horizon to $r_{+}=1$. However, one then needs to undo this scaling to recover physical units.

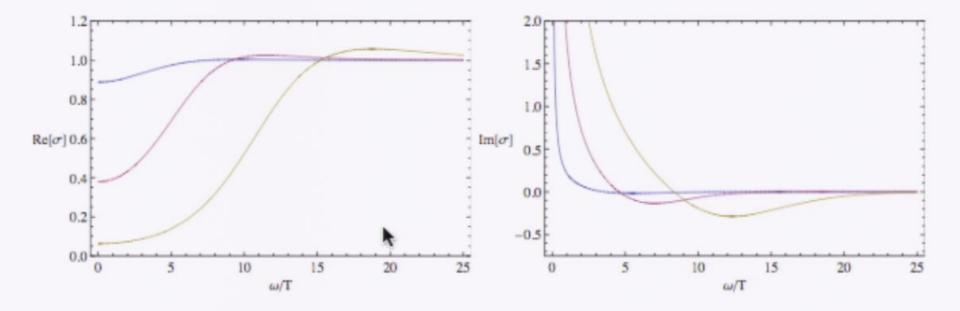
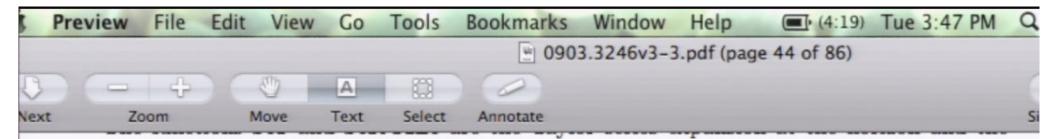


Figure 6: The real (left) and imaginary (right) parts of the electrical conductivity computed via AdS/CFT as described in the text. The conductivity is shown as a function of frequency. Different curves correspond to different values of the chemical potential at fixed temperature.

The gap becomes deeper at larger chemical potential. We have set g=1 in (115). Page 13/58



derivative thereof, respectively. In performing numerics it is generally convenient to set L=1 and furthermore to scale the horizon to $r_+=1$. However, one then needs to undo this scaling to recover physical units.

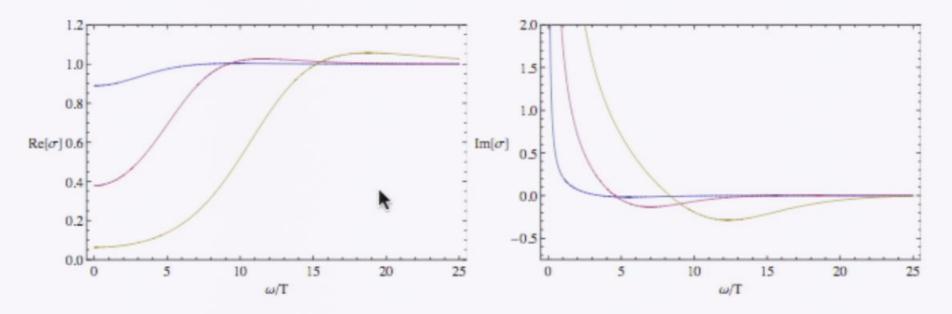


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Different curves correspond to different values of the chemical potential at fixed temperature.



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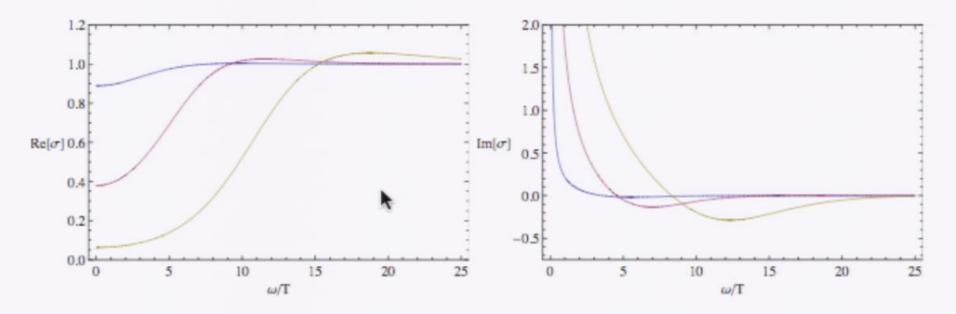
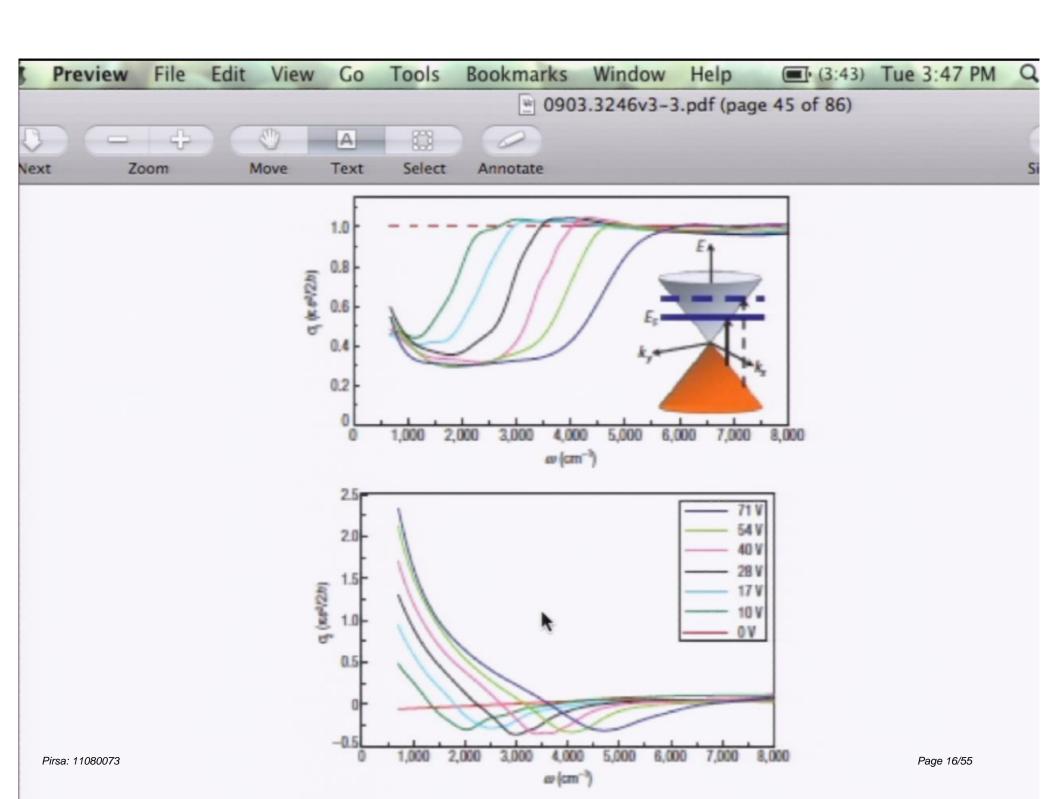


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Fields

$$\psi = \frac{O_1}{r} + \frac{O_2}{r^2} + \cdots \quad \phi = \mu - \frac{\rho}{r} + \cdots$$

$$g = r^2 + a + \frac{b}{r} \dots \qquad e^{-\chi} g = r^2 - \frac{c}{2r} \dots$$

$$e^{\chi} \rightarrow e^{\chi} a^2, t \rightarrow at, \phi \rightarrow \phi/a$$

Pirsa: 11080073 $(\psi_{\scriptscriptstyle H}, E_{\scriptscriptstyle H}) \mapsto (\mu, \rho, O_{\scriptscriptstyle 1}, O_{\scriptscriptstyle 2}, M)$ Page 17/55

Disorder

precise weighting/interaction of disorder not important

$$\int DV \ P[V] \ \frac{1}{Z[V,J=0]} \frac{\delta}{\delta J} Z[V,J]_{J=0}$$

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$$\delta H = \int V(y)O(t,y)$$

$$G_{FF}(\omega,0) = -\frac{\bar{V}^2}{2} \int \frac{d^2k}{(2\pi)^2} k^2 G_{OO}(\omega,k)$$

$$G^{2}G_{PP}(\omega,0) = -(G_{FF}(\omega,0) - G_{FF}(0,0))$$

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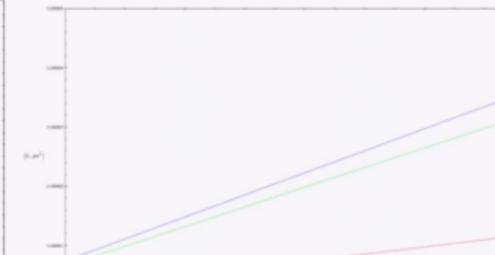
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Adding Disorder Perturbatively

Free, massive scalar field matter $O(t,y) = \phi$

$$\phi'' - \phi' \left(\frac{2}{r} + \frac{g'}{g} - \frac{\chi'}{2} \right) + \phi \left(\frac{e^{\chi} \omega^2}{g^2} - \frac{k^2}{r^2 g} - \frac{m^2}{g} \right)$$





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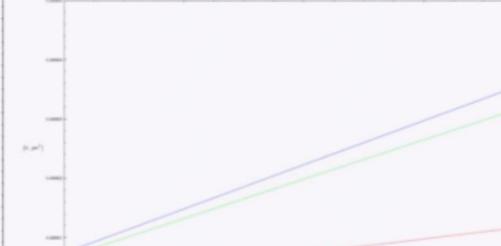
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Gravity Dual to Superconductor/ Superfluid

Einstein-Maxwell-charged scalar

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Schwarzschild-AdS Black Brane

$$ds^{2} = -g(r)e^{-\chi(r)}dt^{2} + \frac{dr^{2}}{g(r)} + r^{2}(dx^{2} + dy^{2})$$

$$\psi = \psi(r)$$
 $A_{\mu} = (\phi(r), 0, 0, 0)$

Equations of Motion

$$y'' + \left(\frac{g'}{g} - \frac{\chi'}{2} + \frac{2}{r}\right)\psi' + \frac{q^2\phi^2 e^{\chi}}{g^2}\psi - \frac{m^2}{g}\psi = 0$$

$$y'' + \left(\frac{\chi'}{2} + \frac{2}{r}\right)\phi' - \frac{2q^2\psi^2}{g}\phi = 0$$

$$\chi' + r\psi'^2 + \frac{rq^2\phi^2\psi^2 e^{\chi}}{g^2} = 0$$

$$y' + \left(\frac{1}{r} - \frac{\chi'}{2}\right)g + \frac{r\phi'^2 e^{\chi}}{4} - 3r + \frac{rm^2\psi^2}{2} = 0$$

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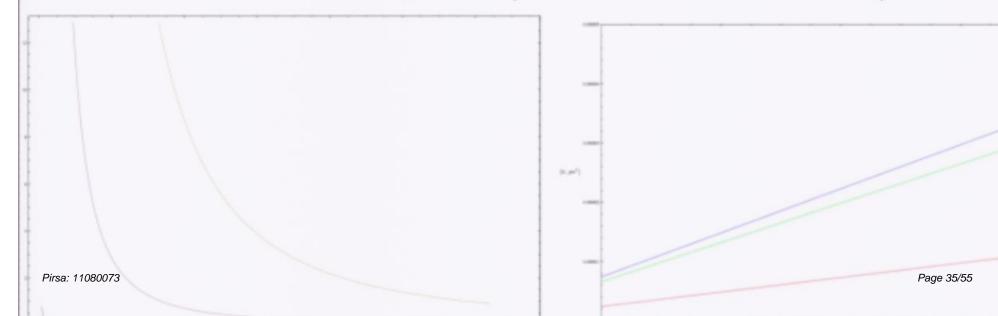
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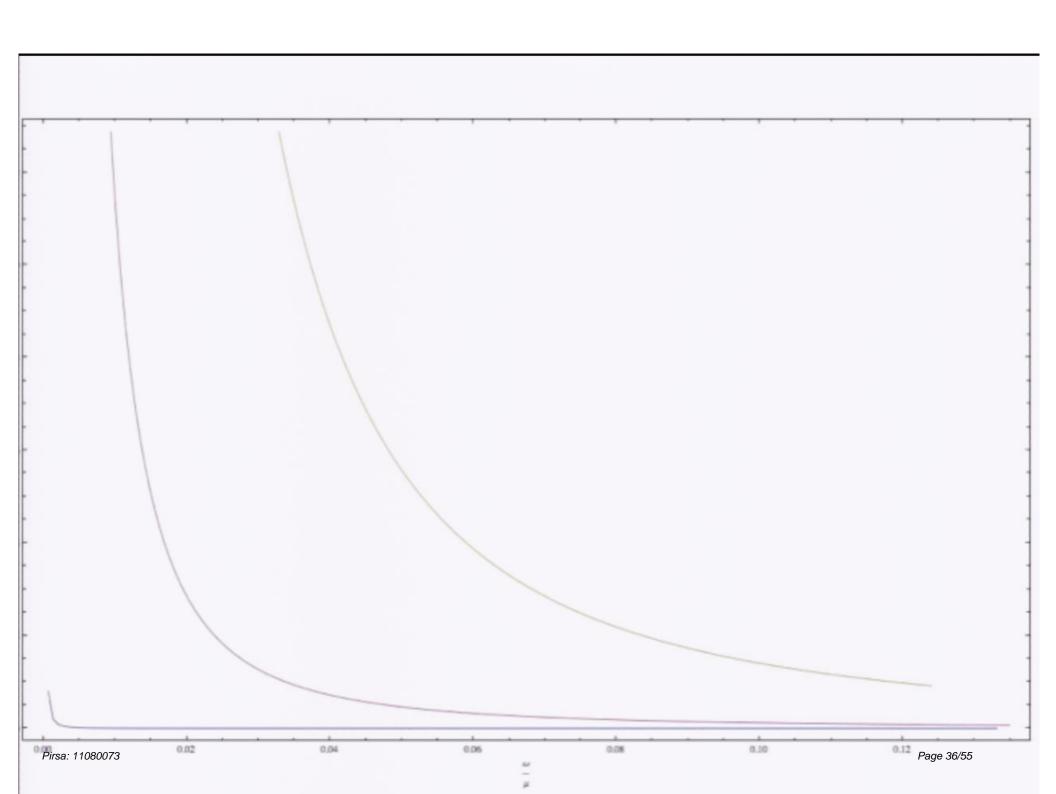
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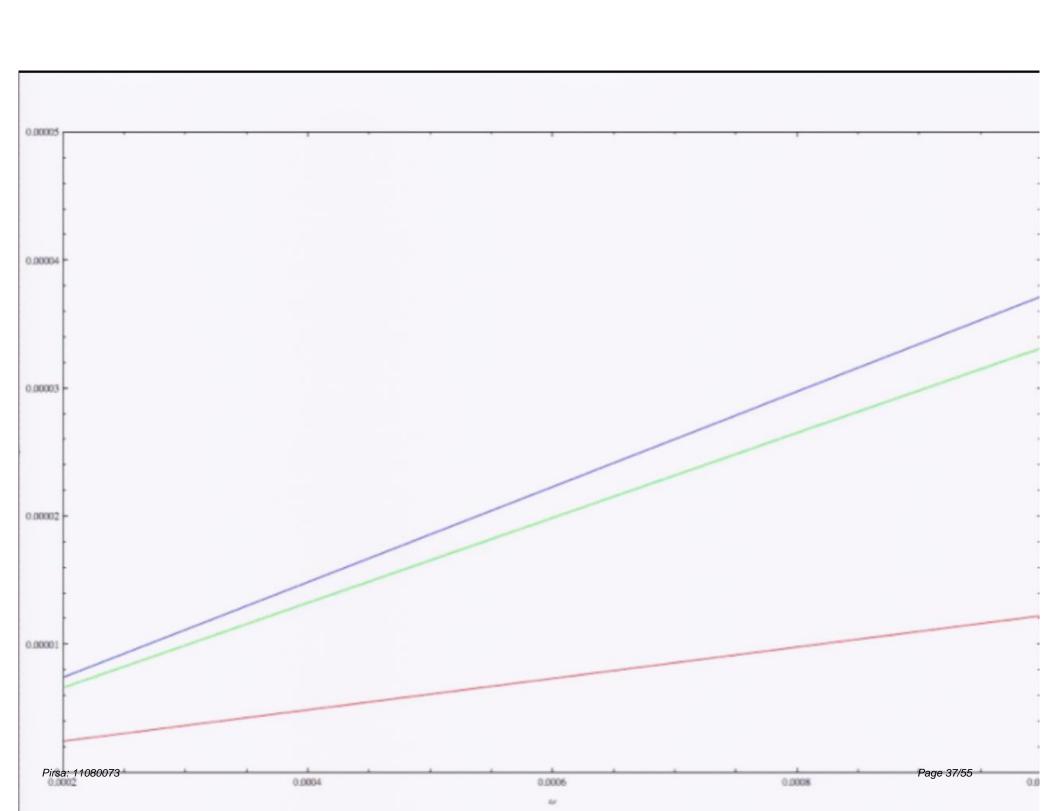
Adding Disorder Perturbatively

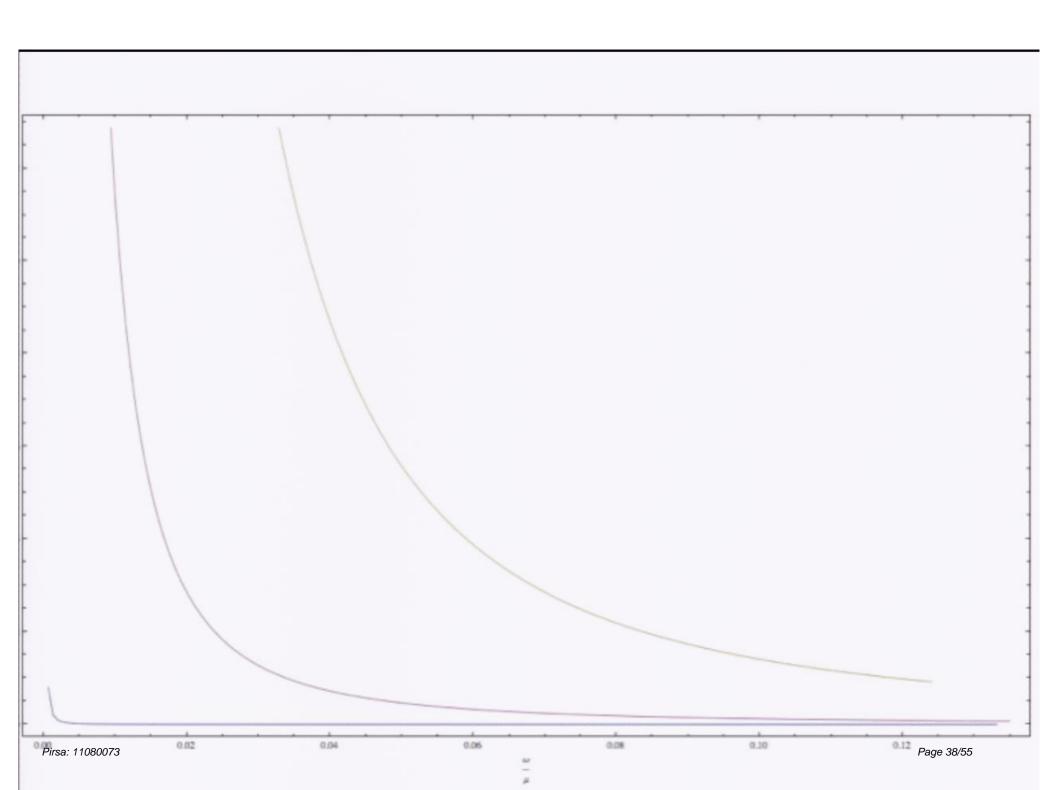
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Holography/Drude Model

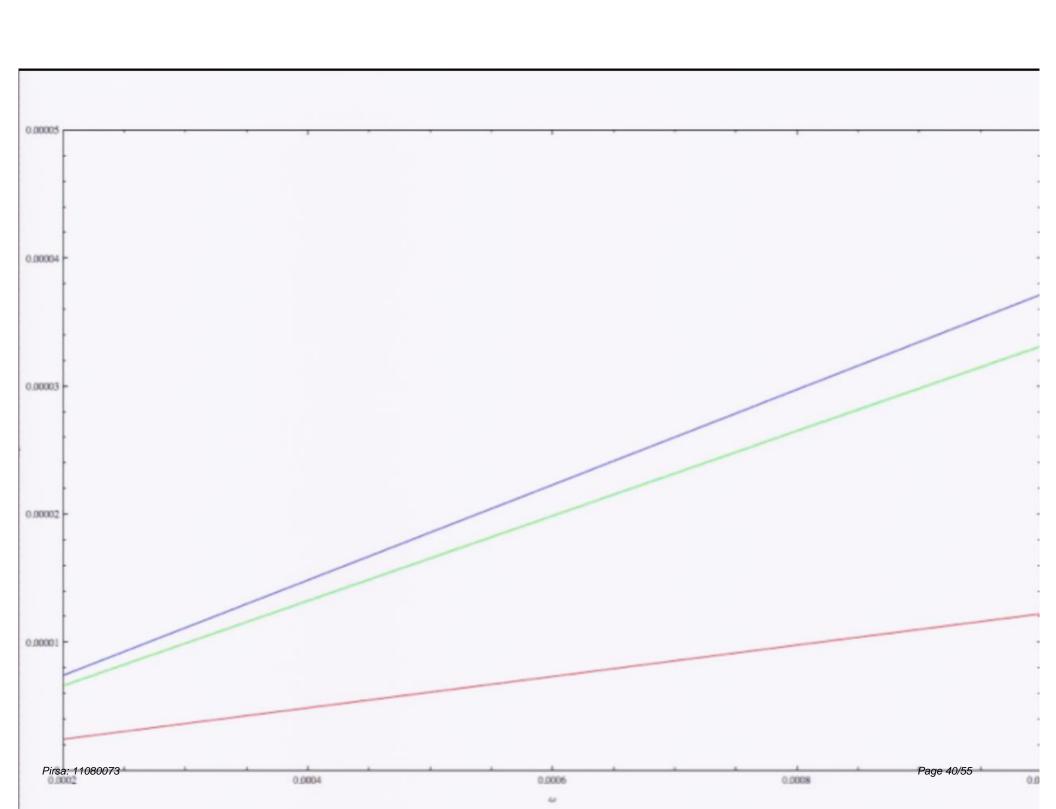
Drude Model

$$\kappa(\omega) = \frac{\tau}{1 - i\omega\tau} \implies$$

$$\operatorname{Re}[\kappa(\omega)] \propto \frac{k\tau}{1+\omega^2\tau^2}$$

Holography+Perturbative Disorder

$$\kappa(\omega) = \frac{1}{i\omega T} + O\left(\frac{1}{\omega^2}\right)$$
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Holography/Drude Model

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Conclusions/Further Thoughts

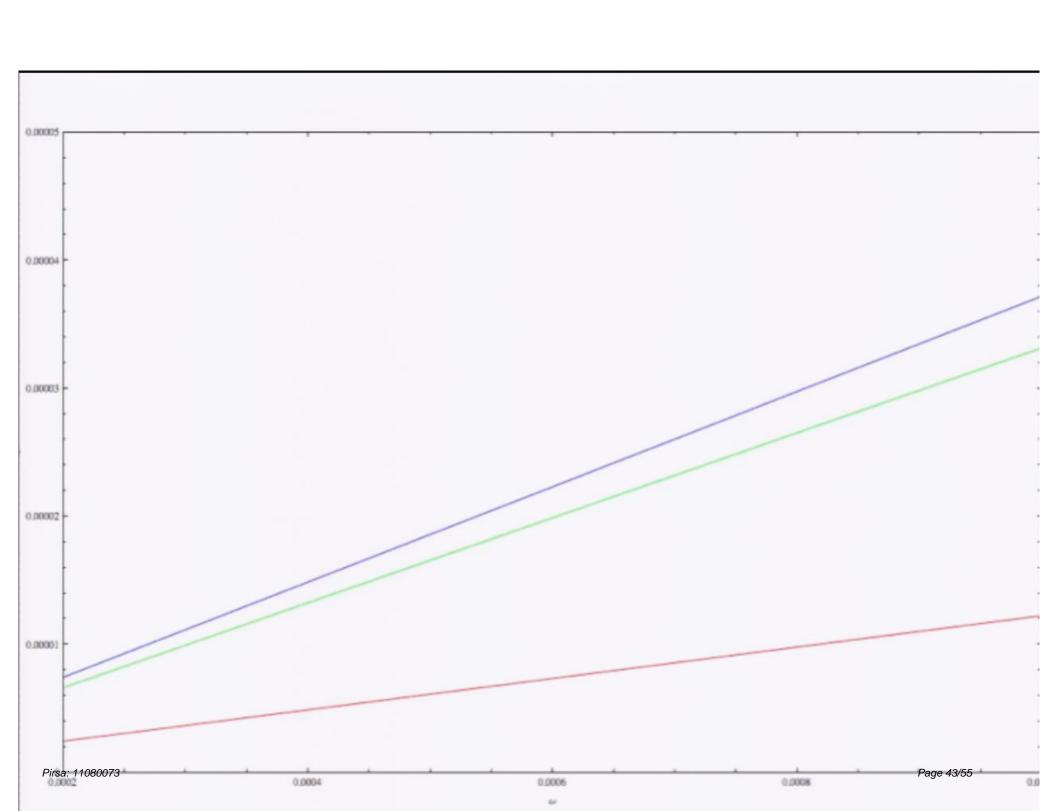
Holography+Perturbative Disorder – agrees with crude Drude model

Disorder Weak in some dimensionless units

Extract temperature dependence from holography +Perturbative Disorder

Consider direct coupling of disorder to order parameter

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Conclusions/Further Thoughts

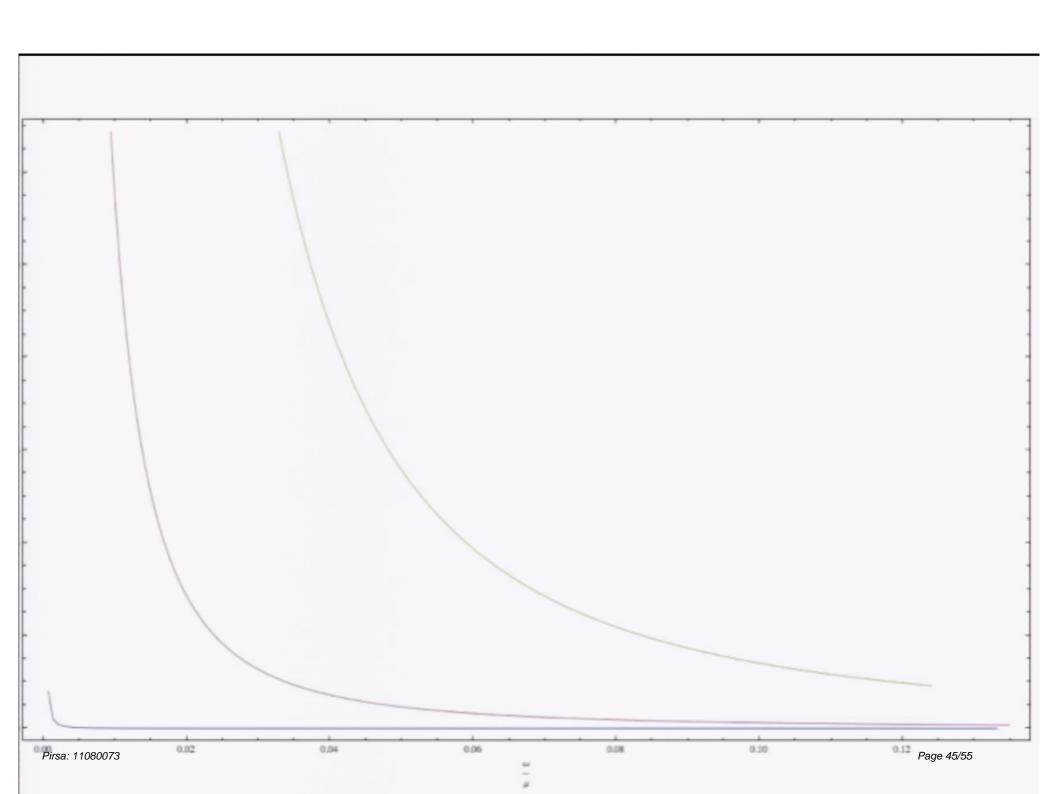
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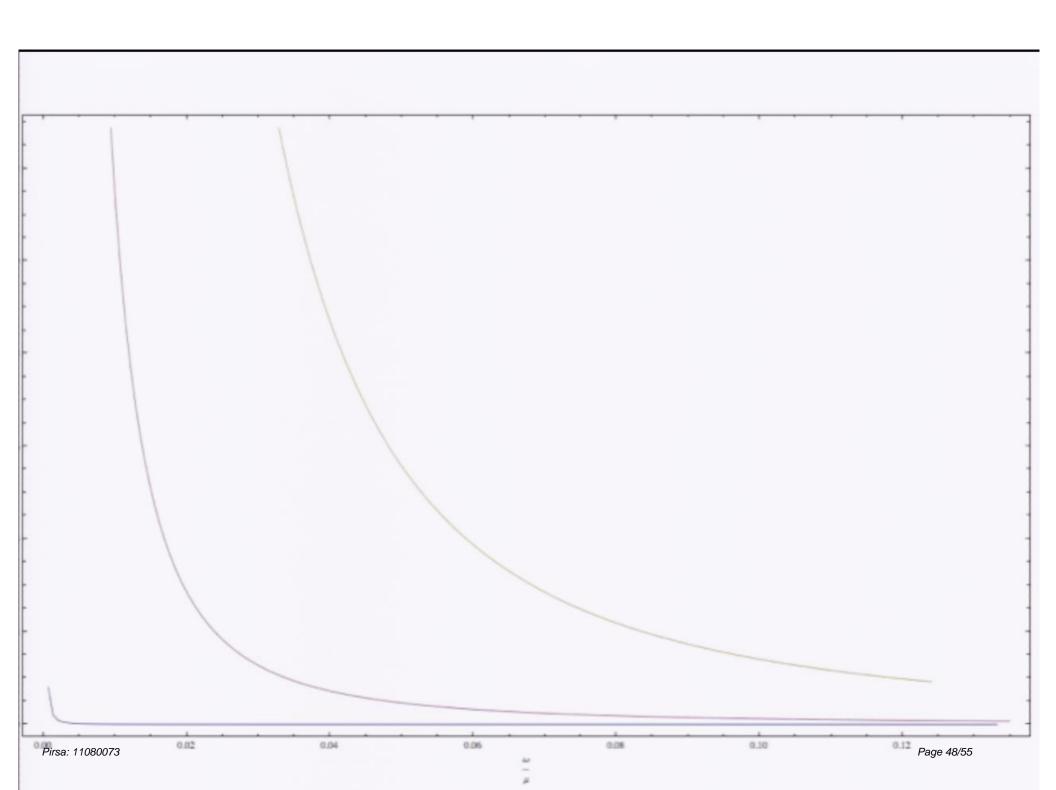
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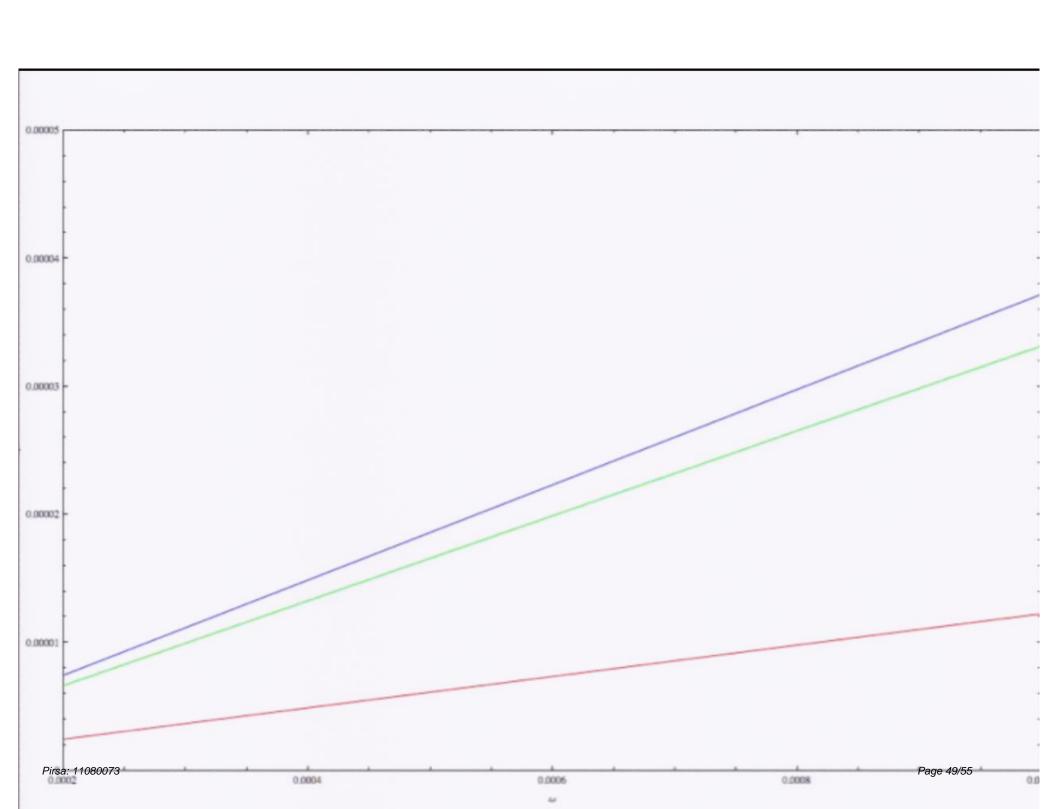
Thanks

Advisors: Janet Hung and Yanwen Shang

Perimeter Institute

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Holography+Perturbative Disorder

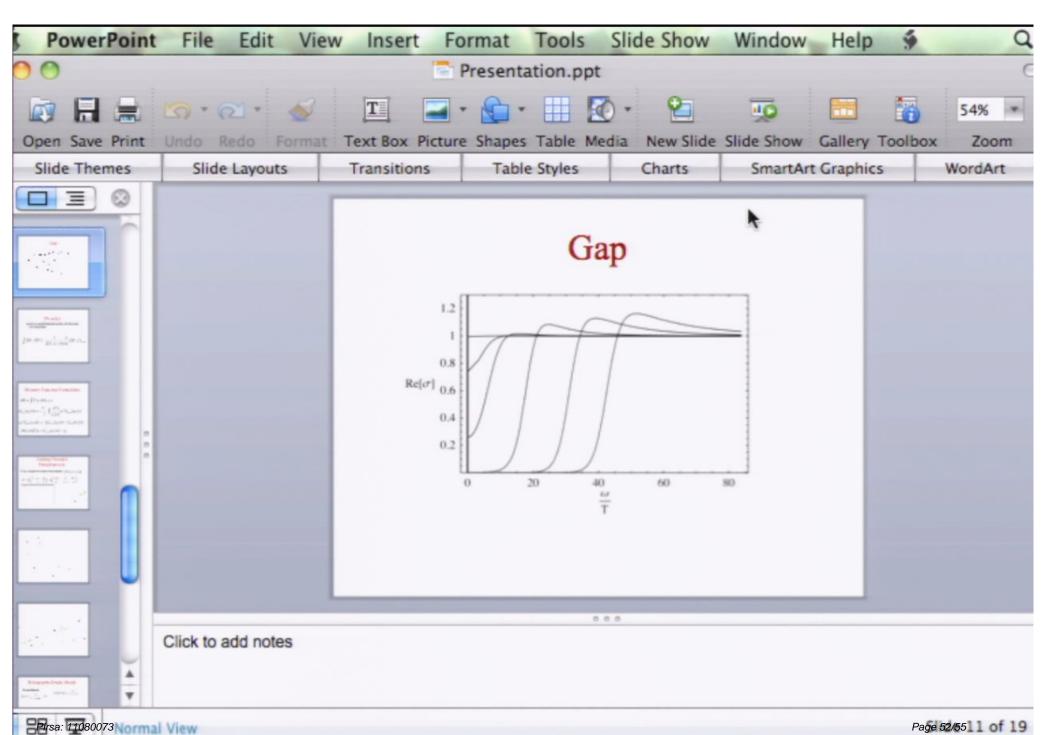
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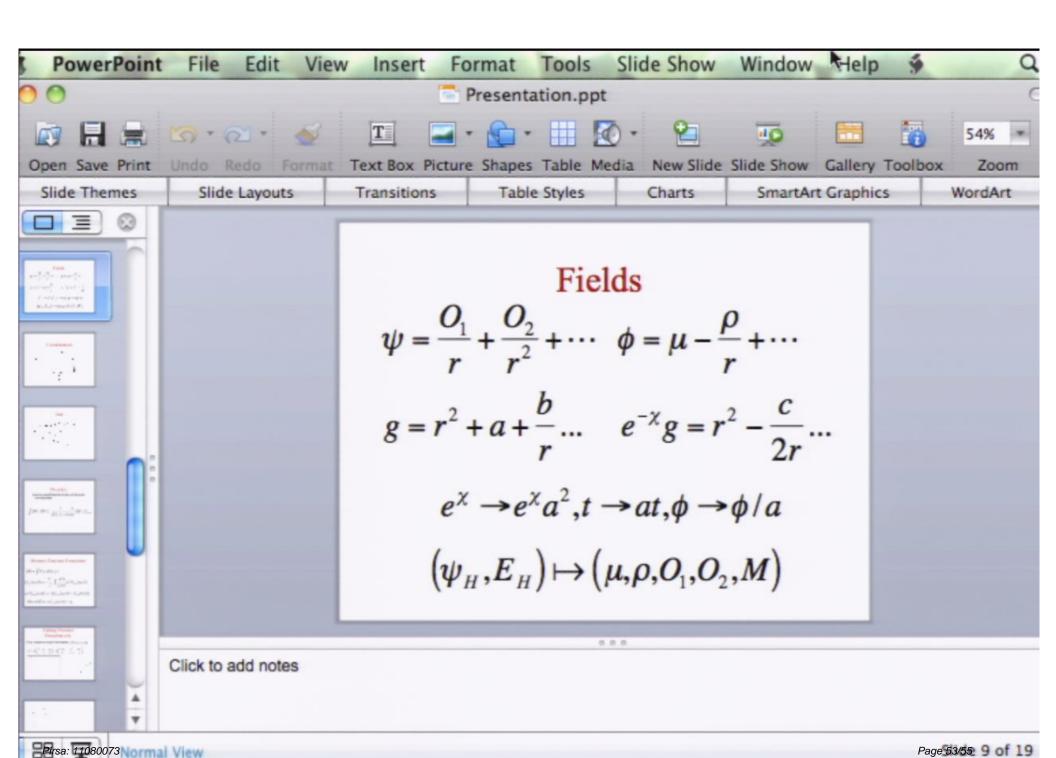
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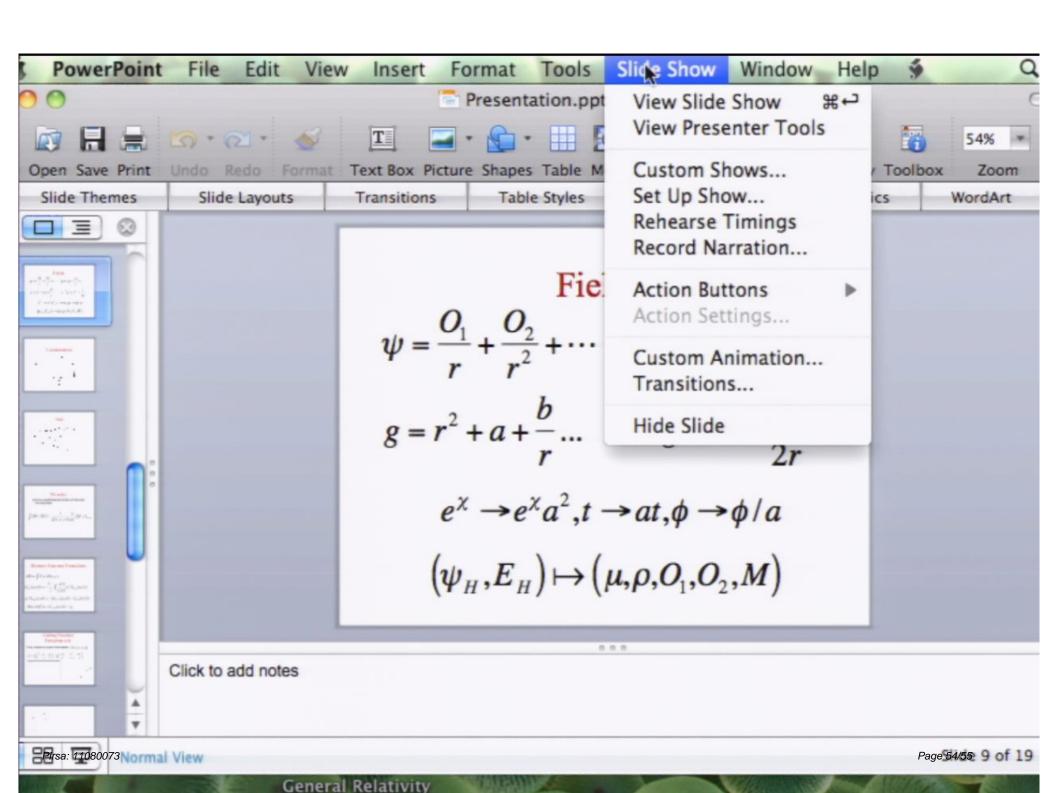
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