

Title: Cosmic Strings

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Abstract:

# Cosmic Strings

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August 9, 2011

# Outline

Introduction and Motivation

String Dynamics

Cusp Formations

Extra Dimensions

A New Approach

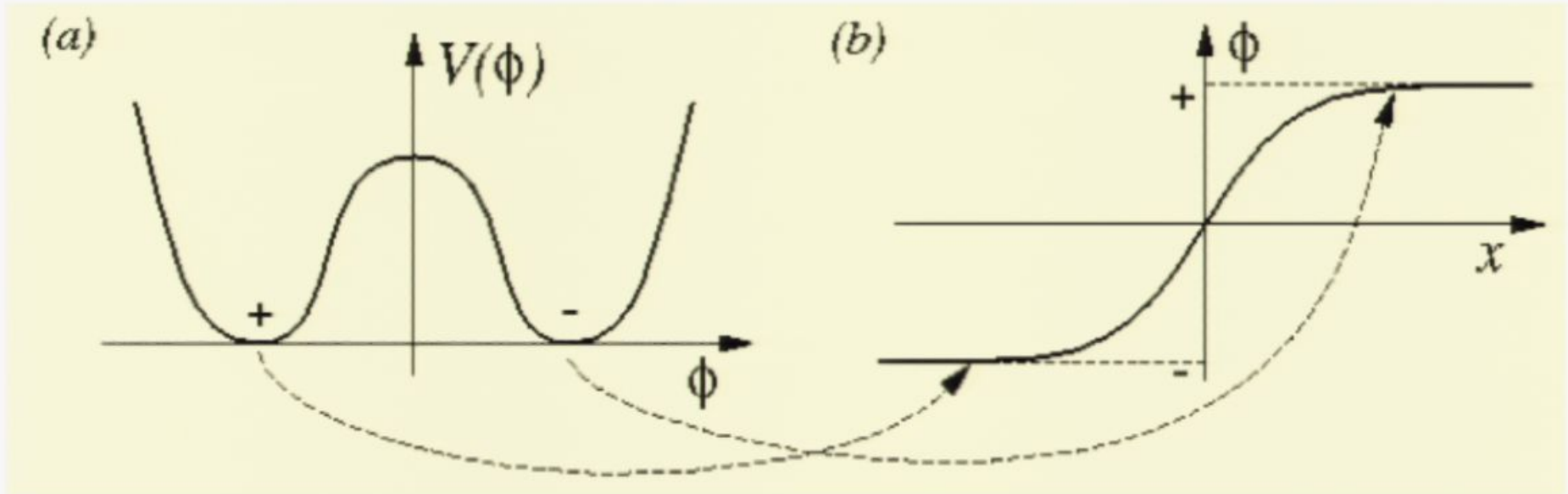
Conclusions

# Phase Transitions and Topological Defects

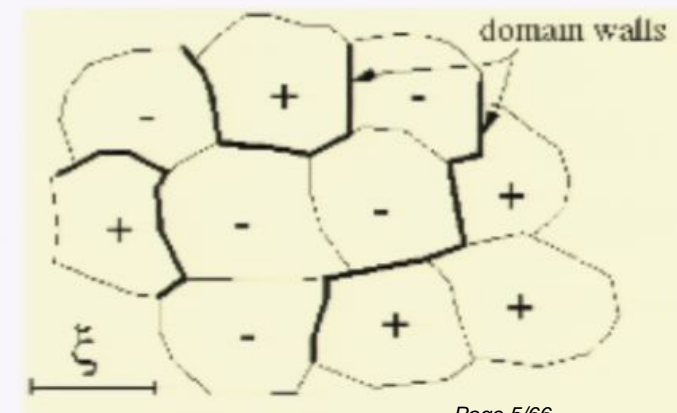
- ▶ Phase transitions can have a wide variety of important implications, including the formation of topological defects - cosmic strings, domain walls, and monopoles.
- ▶ The type of defect formed is determined by the symmetry properties of the matter and the nature of the phase transition.



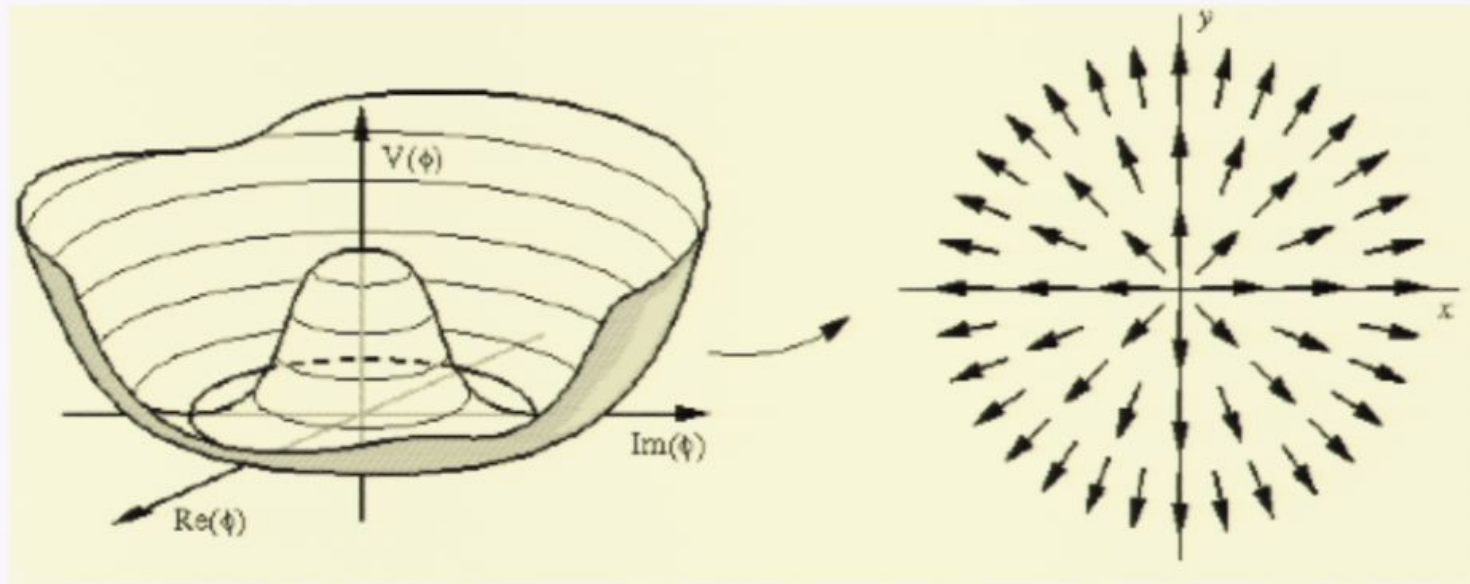
# Domain Walls - Discrete Symmetry



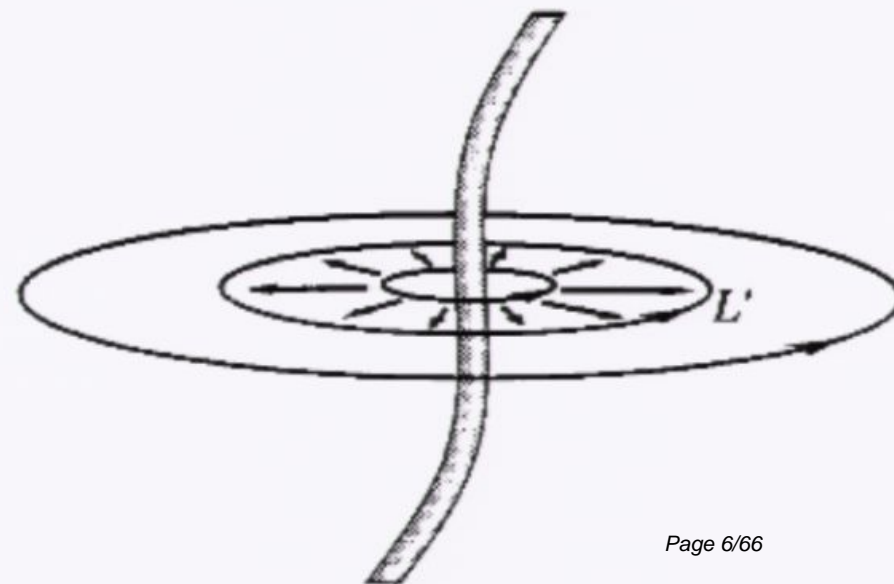
Topological defects are the “boundaries” between these regions with different choices of minima.



# Cosmic Strings - $U(1)$ Symmetry



The "Mexican hat" potential has a "hole" in the vacuum manifold.



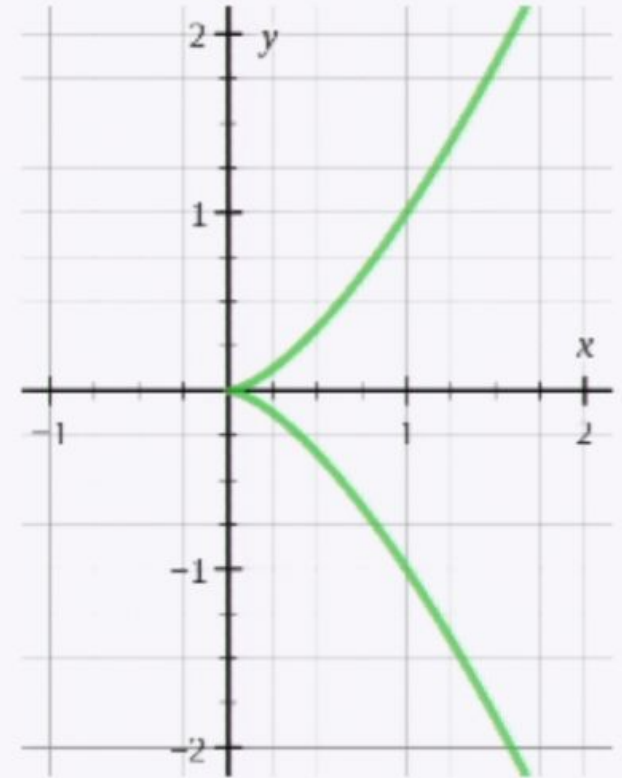
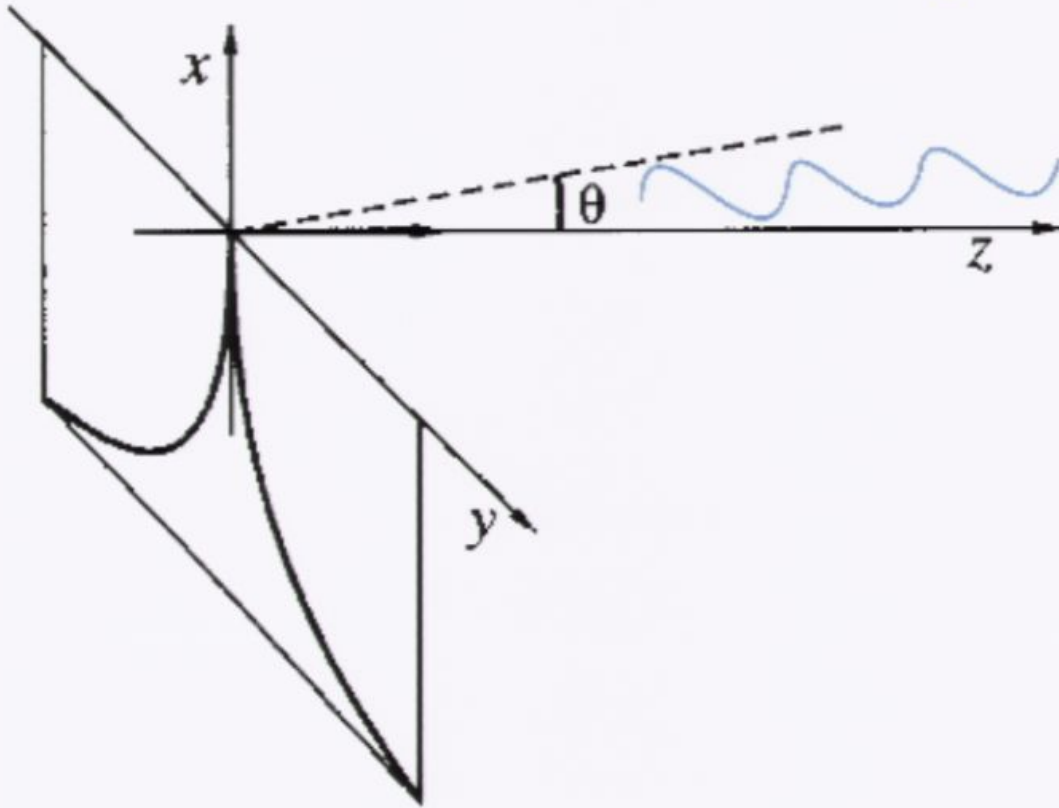
## Why do we care?

- ▶ Link to the physics of the very early universe
- ▶ Cosmic strings could be the “seeds” that led to the formation of the large-scale structures
- ▶ String evolution corresponds to minimal surfaces



# How do we look for a cosmic string?

## Gravitational wave bursts from cusps



Typical cusp:

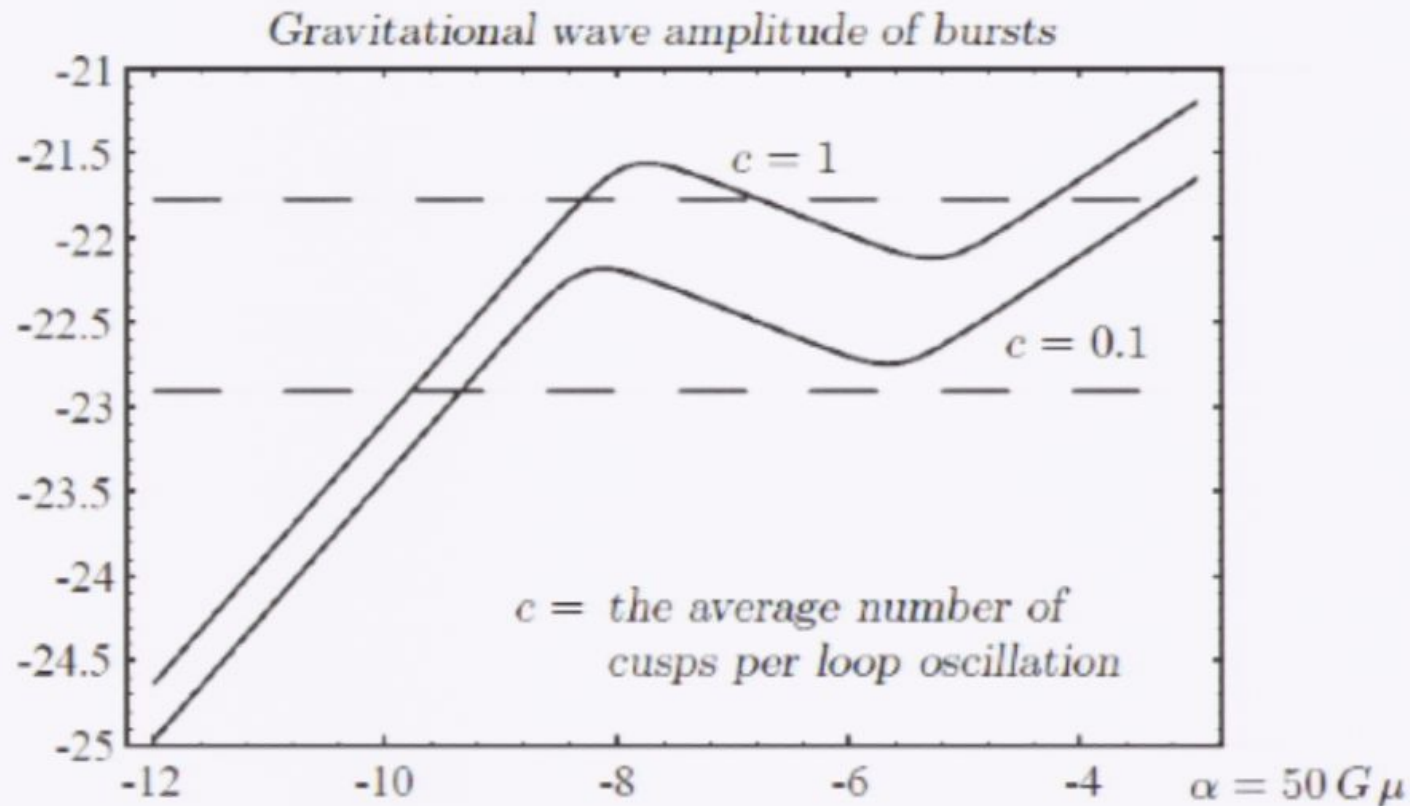
$$y^2 = x^3$$

$$\mathbf{r}(s) = (s^2, s^3)$$



# Have we seen a cosmic string?

- ▶ Nope
- ▶ but LIGO is working hard



## Question:

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How do we (better) model cusp formation in cosmic strings?

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# Nambu-Goto Action

- ▶ We can parametrize a cosmic string as

$$X^\mu = X^\mu(\tau, \sigma),$$

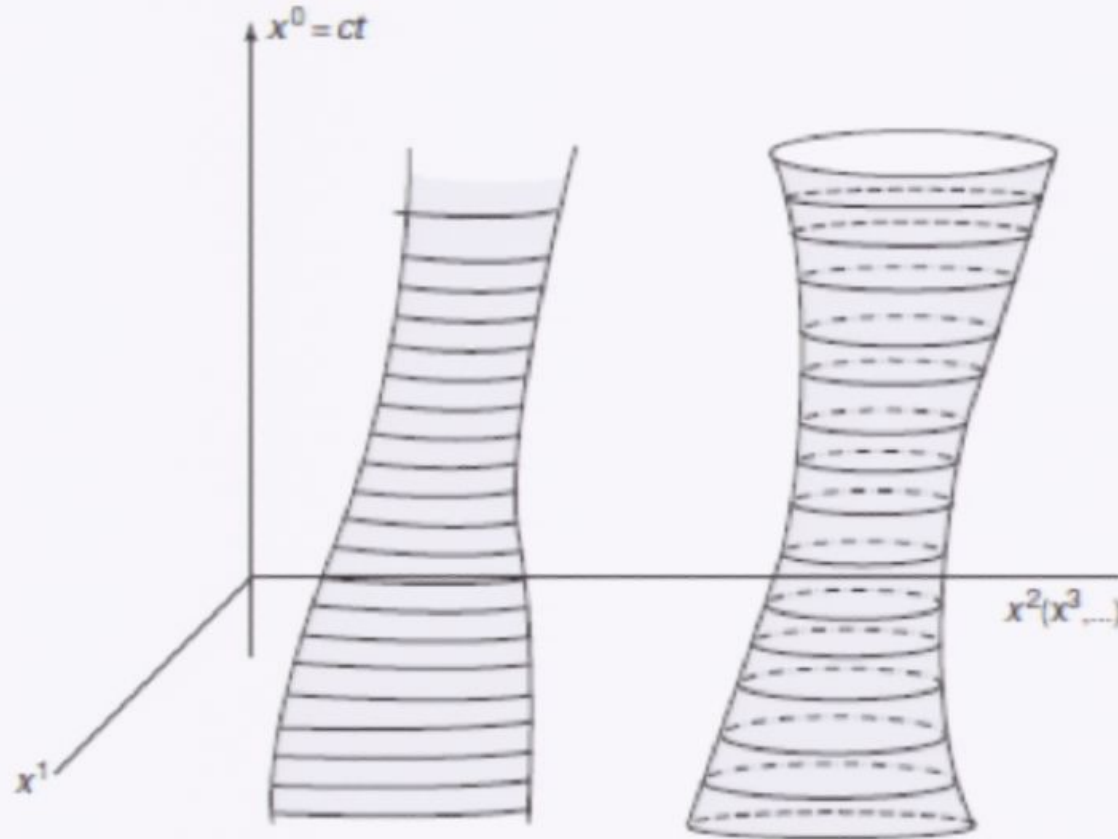
where  $\mu = 0, 1, 2, 3$ .

- ▶ The Nambu-Goto action is given by

$$\mathcal{S} = -\mu \int d\tau \int d\sigma \sqrt{-\gamma}$$

where  $\gamma_{\alpha\beta} = \eta_{\mu\nu} \frac{\partial X^\mu}{\partial \xi^\alpha} \frac{\partial X^\nu}{\partial \xi^\beta}$  is the pull-back metric.

# String Evolution as Minimal Surface



- ▶ String moving in spacetime traces out a minimal surface
- ▶ Action is proportional to area, and hence *reparametrization invariant*.

# Equation of Motion

## ► First Variation

$$\delta S = -\mu \int_{\tau_f}^{\tau_i} d\tau \int_0^{\sigma_0} d\sigma \delta X^\mu \left( \frac{\partial}{\partial \tau} \left( \frac{\partial \mathcal{L}}{\partial \dot{X}^\mu} \right) + \frac{\partial}{\partial \sigma} \left( \frac{\partial \mathcal{L}}{\partial X'^\mu} \right) \right)$$

## ► Equation of Motion

$$\frac{\partial}{\partial \tau} \left( \frac{(\dot{X} \cdot X')X' - (X')^2 \dot{X}}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}} \right) + \frac{\partial}{\partial \sigma} \left( \frac{(\dot{X} \cdot X')\dot{X} - (\dot{X})^2 X'}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}} \right) = 0$$

# Temporal Gauge

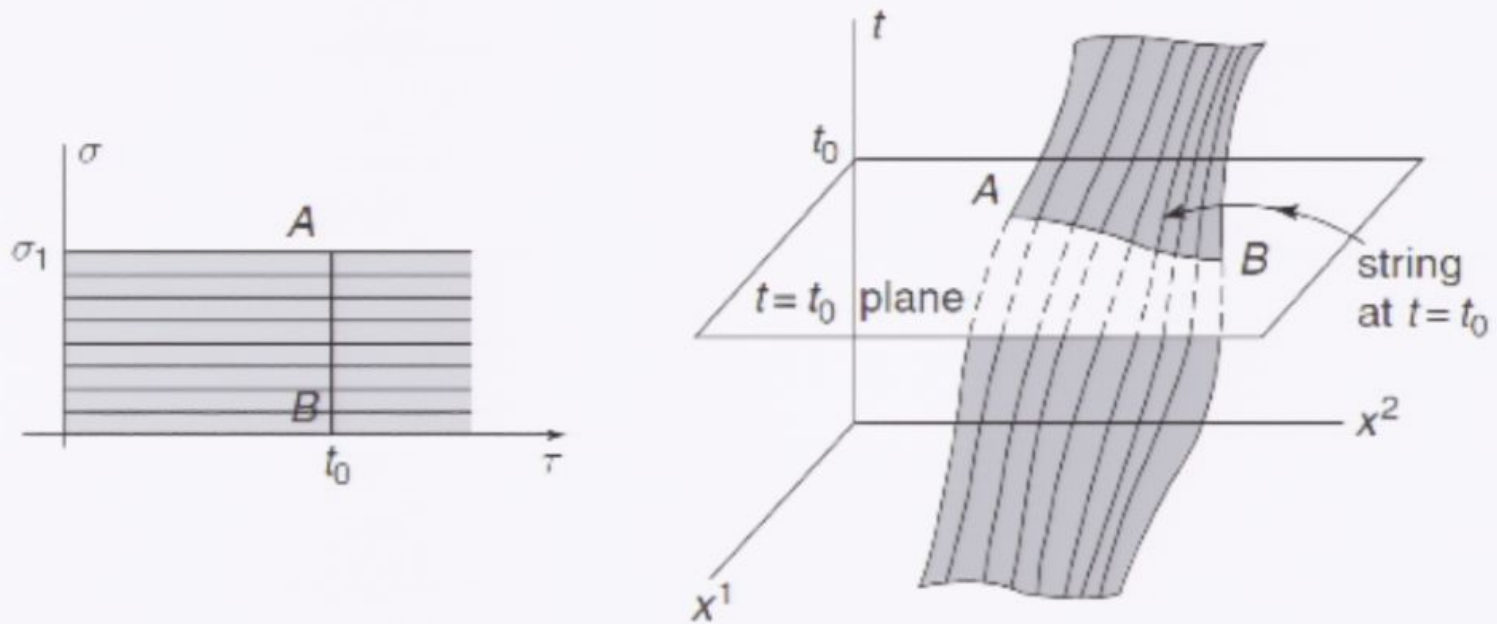


Figure:  $X^\mu(\tau, \sigma) = (\tau, \mathbf{r}(\tau, \sigma))$

# Orthogonal Gauge

- ▶ How do we get

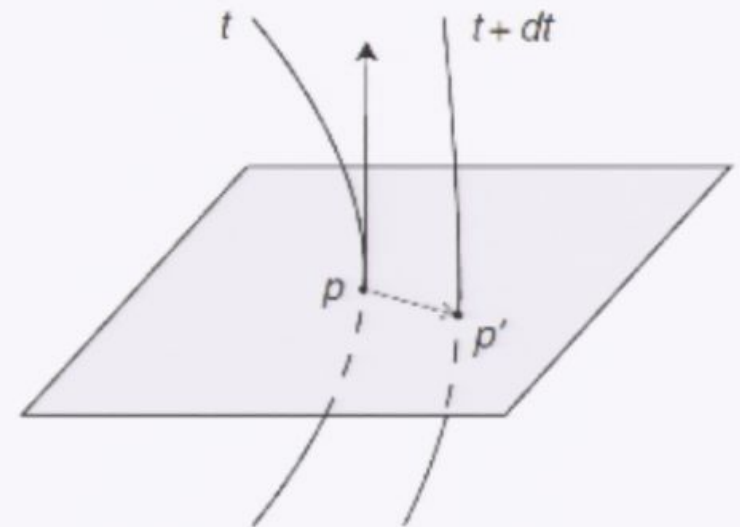
$$\dot{\mathbf{r}} \cdot \mathbf{r}' = 0?$$

- ▶ Reparametrize

$$(\tau, \sigma) \mapsto \mathbf{r}(\tau, f(\tau, \sigma)),$$

where  $f$  solves the linear transport equation

$$f_t = -\frac{\mathbf{r}_\tau \cdot \mathbf{r}_\sigma}{|\mathbf{r}_\sigma|^2} f_\sigma$$



Transverse velocity

$$\mathbf{v}_\perp = \dot{\mathbf{r}} - (\dot{\mathbf{r}} \cdot \mathbf{r}')\mathbf{r}'$$



# Wave Equation and Constraints

- ▶ Wave Equation

$$\frac{\partial^2 \mathbf{r}}{\partial \tau^2} - \frac{\partial^2 \mathbf{r}}{\partial \sigma^2} = 0$$

- ▶ Constraint can be written as

$$\left( \frac{\partial \mathbf{r}}{\partial \tau} \right)^2 + \left( \frac{\partial \mathbf{r}}{\partial \sigma} \right)^2 = 1$$

- ▶ Periodic boundary conditions

## Representation of Solutions

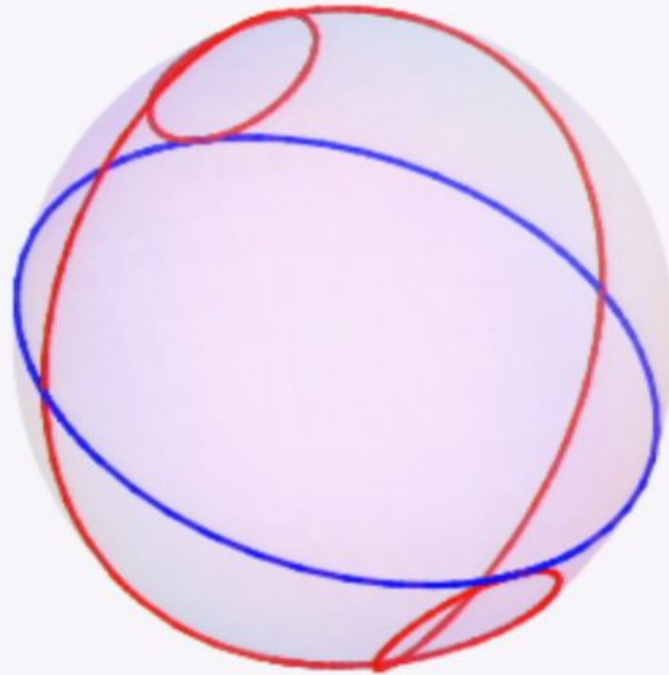
Using the d'Alembert formula, we have

$$\mathbf{r}(\tau, \sigma) = \frac{1}{2} \left( \mathbf{a}(\tau + \sigma) + \mathbf{b}(\tau - \sigma) \right)$$

where

$$|\mathbf{a}'|^2 = |\mathbf{b}'|^2 = 1$$

# Kibble-Turok Sphere



Left- and Right-movers on the Kibble-Turok sphere.

## Kibble-Turok Strings (1982)

$$\mathbf{r}_{KT}(s, t) = \frac{1}{2} \begin{pmatrix} (1 - \alpha) \sin u + \frac{1}{3} \alpha \sin 3u \\ -(1 - \alpha) \cos u - \frac{1}{3} \alpha \cos 3u \\ -\sqrt{\alpha(1 - \alpha)} \cos u \end{pmatrix} + \frac{1}{2} \begin{pmatrix} \sin v \\ -\cos v \\ 0 \end{pmatrix},$$

where  $u = t + s$ ,  $v = t - s$ , and  $\alpha$  is an arbitrary parameter between 0 and 1.

Previous Next Zoom Move Text Select Annotate Sidebar

### Nambu-Goto Strings (1982)

$$r_{KT}(s, t) = \frac{1}{2} \begin{pmatrix} (1 - \alpha) \sin u + \frac{1}{3} \alpha \sin 3u \\ -(1 - \alpha) \cos u - \frac{1}{3} \alpha \cos 3u \\ -\sqrt{\alpha(1 - \alpha)} \cos u \end{pmatrix} + \frac{1}{2} \begin{pmatrix} \sin v \\ -\cos v \\ 0 \end{pmatrix},$$

where  $u = t + s$ ,  $v = t - s$ , and  $\alpha$  is an arbitrary parameter between 0 and 1.

- ▼ talk
  - Introduction and Motiv...
  - String Dynamics
  - Cusp Formations
  - Extra Dimensions
  - A New Approach
  - Conclusions

### Nambu-Goto Strings (1982)



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## ▼ Kibble-Turok Strings

T.W.B Kibble and Neil Turok. Self-intersection of Cosmic Strings. Phys. Lett., B116:141, 1982.

Manipulate[

```

a[u_] = {(1 - α) Cos[u] + α Sin[3 u] / 3, -(1 - α) Cos[u] - α Cos[3 u] / 3,
  -Sqrt[α (1 - α)] Cos[u]};
b[v_] = {Sin[v], -Cos[v], 0};
ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange → 1],
{{α, 0.5}, 0, 1}, {t, 0, 2 π}]

```

Slide 2 of 5

## ▼ Turok Strings

Neil Turok. Grand Unified Strings and Galaxy Formation. Nucl. Phys., B242:520, 1984.

Manipulate[

```

a[u_] = {(1 - α) Cos[u] + α Sin[3 u] / 3, -(1 - α) Cos[u] - α Cos[3 u] / 3,
  -Sqrt[α (1 - α)] Cos[u]};

```



```
ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange -> 1],
{{alpha, 0.5}, 0, 1}, {t, 0, 2 pi}]
```

Slide 2 of 5

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Neil Turok. Grand Unified Strings and Galaxy Formation. Nucl. Phys., B242:520, 1984.

```
Manipulate[
  a[u_] = {(1 - alpha) Cos[u] + alpha Sin[3 u] / 3, -(1 - alpha) Cos[u] - alpha Cos[3 u] / 3,
    -Sqrt[alpha (1 - alpha)] Cos[u]};
  b[v_] = {Sin[phi] Cos[v], -Cos[phi] Cos[v], -Sin[phi] Cos[v]};
  ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange -> 1],
  {{alpha, 0.5}, 0, 1}, {{phi, pi / 2}, 0, pi}, {t, 0, 2 pi}]
```

Slide 3 of 5

## ▼ Chen, DiCarlo, and Hotes (CDH) Strings

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```
ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange -> 1],
{{alpha, 0.5}, 0, 1}, {t, 0, 2 pi}]
```

Slide 2 of 5

## ▼ Turok Strings

Neil Turok. Grand Unified Strings and Galaxy Formation. Nucl. Phys., B242:520, 1984.

```
In[1]:= Manipulate[
  a[u_] = {(1 - alpha) Cos[u] + alpha Sin[3 u] / 3, -(1 - alpha) Cos[u] - alpha Cos[3 u] / 3,
    -Sqrt[alpha (1 - alpha)] Cos[u]};
  b[v_] = {Sin[v], -Cos[phi] Cos[v], -Sin[phi] Cos[v]};
  ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange -> 1],
  {{alpha, 0.5}, 0, 1}, {{phi, pi / 2}, 0, pi}, {t, 0, 2 pi}]
```

Slide 3 of 5

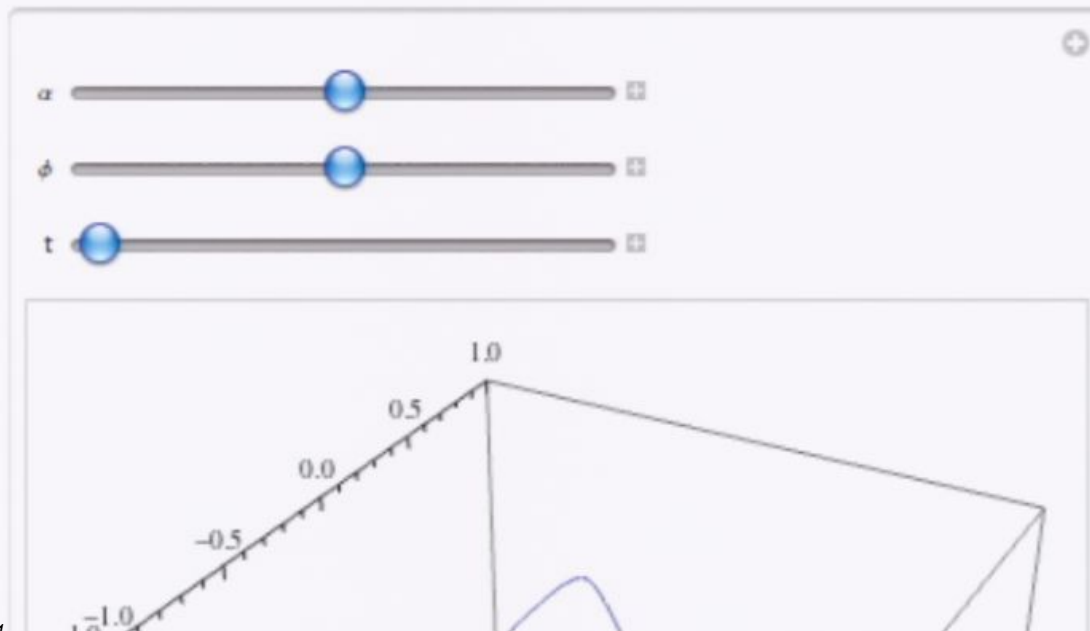
## ▼ Chen, DiCarlo, and Hotes (CDH) Strings

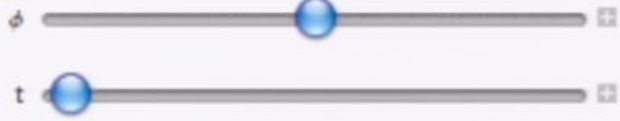
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## ▼ Turok Strings

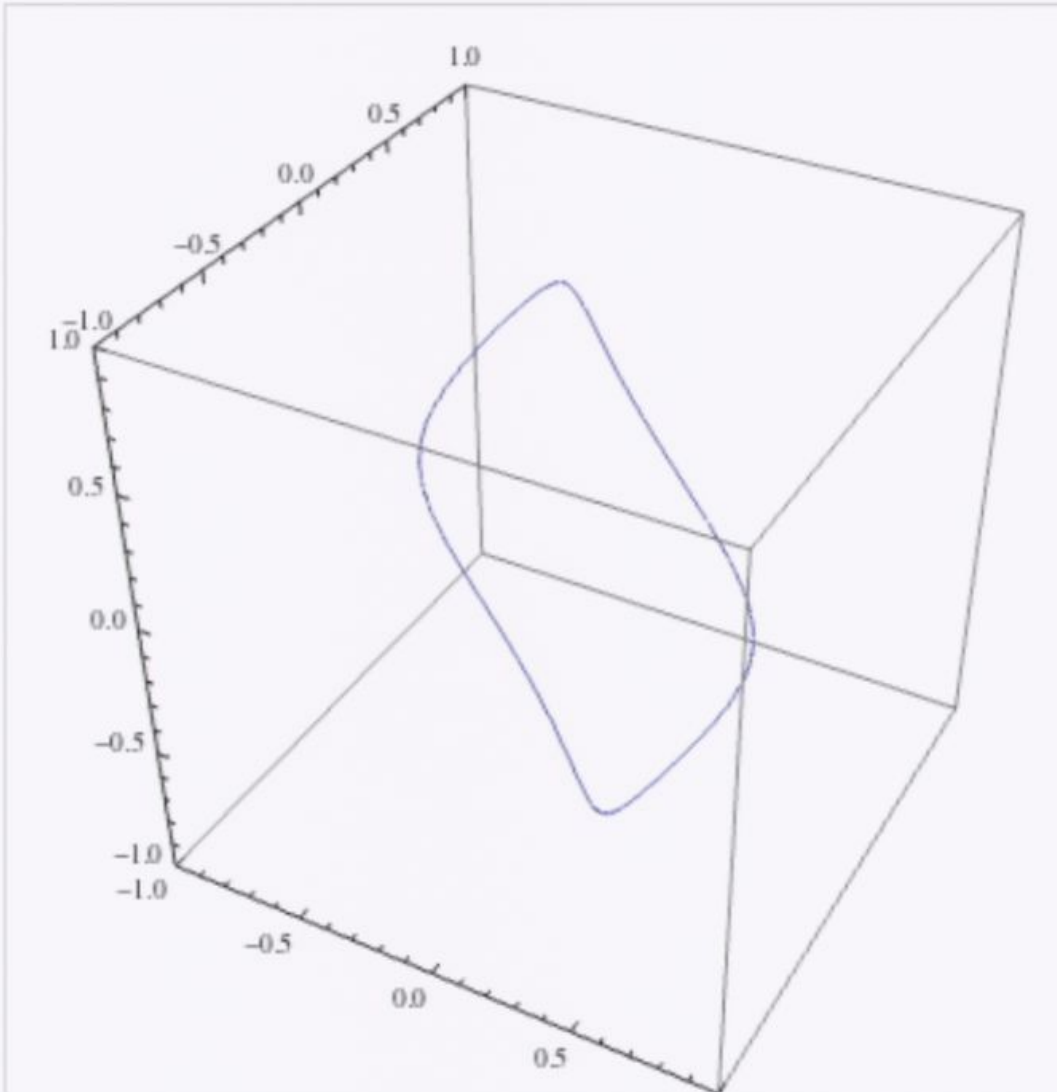
Neil Turok. Grand Unified Strings and Galaxy Formation. Nucl. Phys., B242:520, 1984.

```
▼In[1]:= Manipulate[  
  a[u_] = {(1 -  $\alpha$ ) Cos[u] +  $\alpha$  Sin[3 u] / 3, -(1 -  $\alpha$ ) Cos[u] -  $\alpha$  Cos[3 u] / 3,  
    -Sqrt[ $\alpha$  (1 -  $\alpha$ )] Cos[u]};  
  b[v_] = {Sin[v], -Cos[ $\phi$ ] Cos[v], -Sin[ $\phi$ ] Cos[v]};  
  ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange -> 1],  
  {{ $\alpha$ , 0.5}, 0, 1}, {{ $\phi$ ,  $\pi/2$ }, 0,  $\pi$ }, {t, 0, 2  $\pi$ }]
```



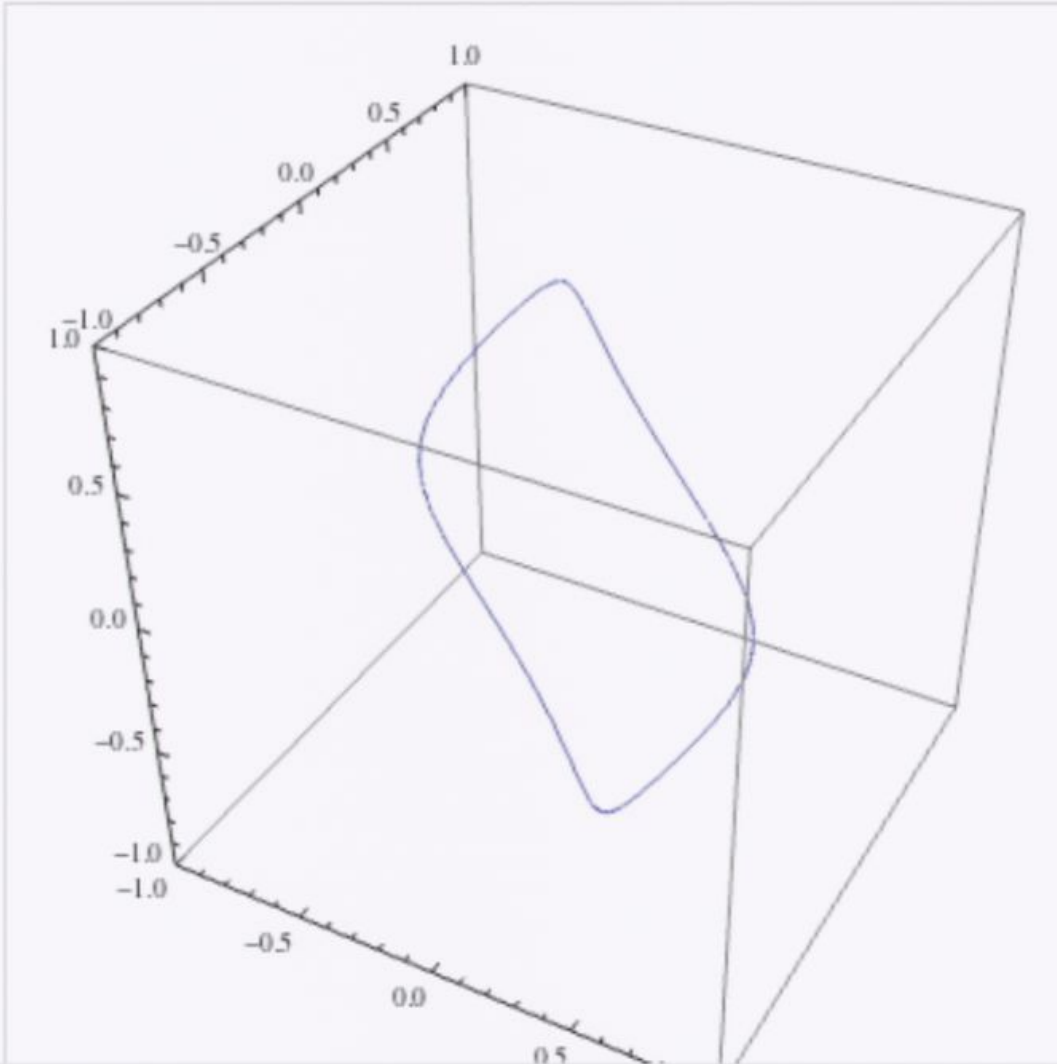


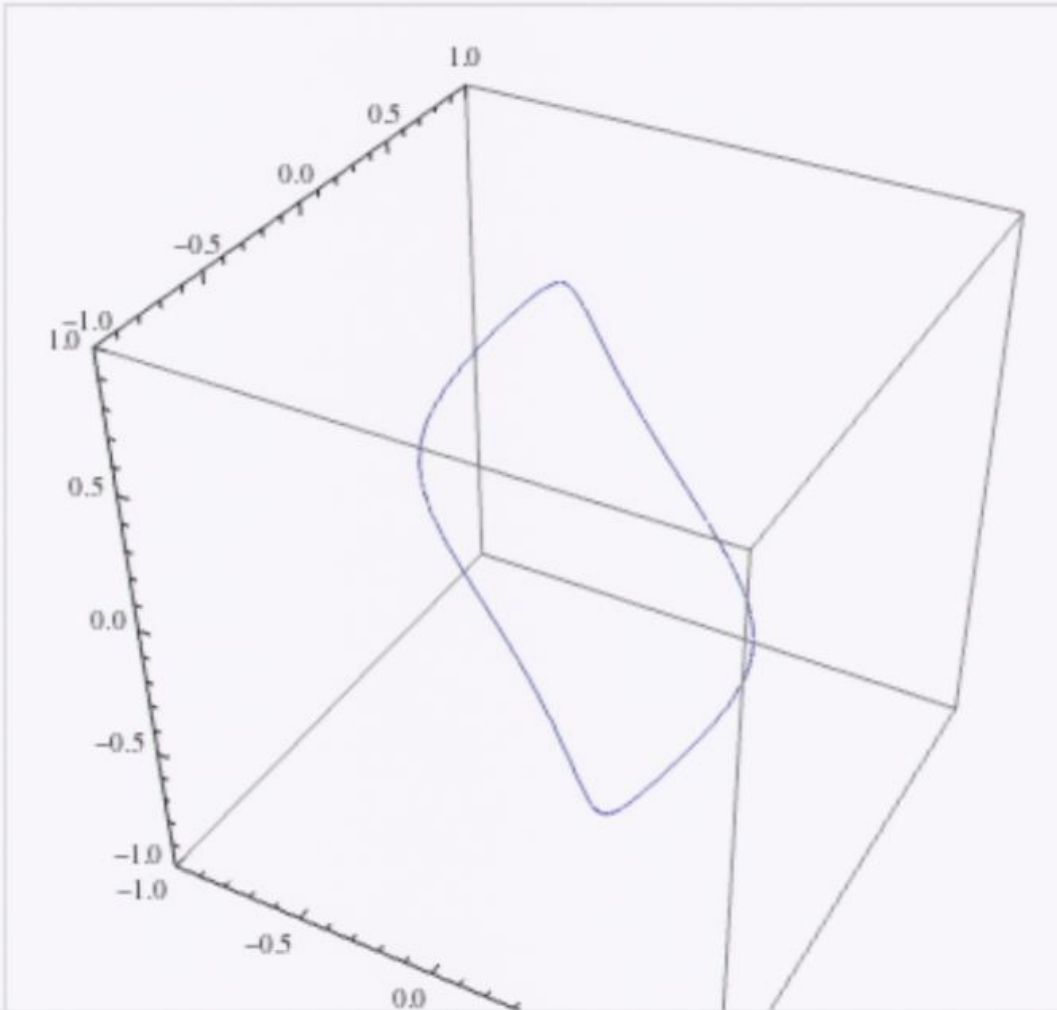
Out[1]=





Out[1]=

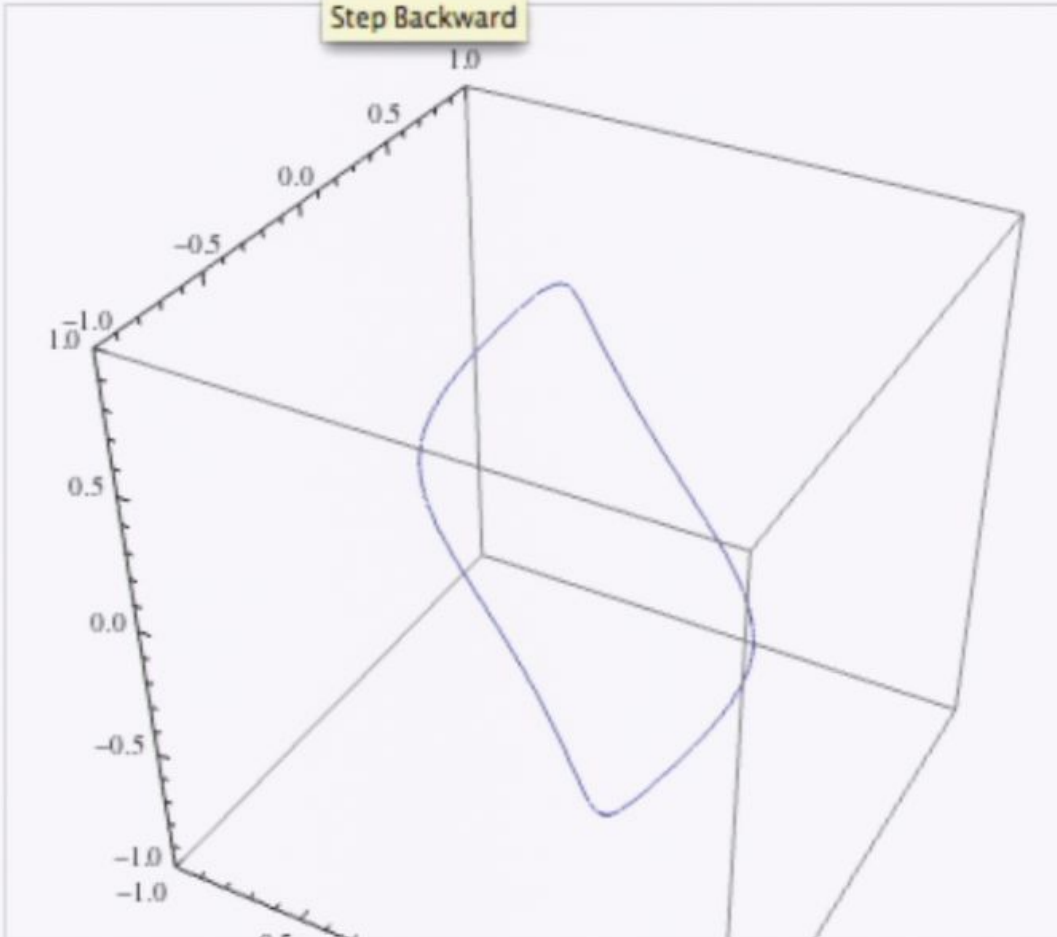








Step Backward



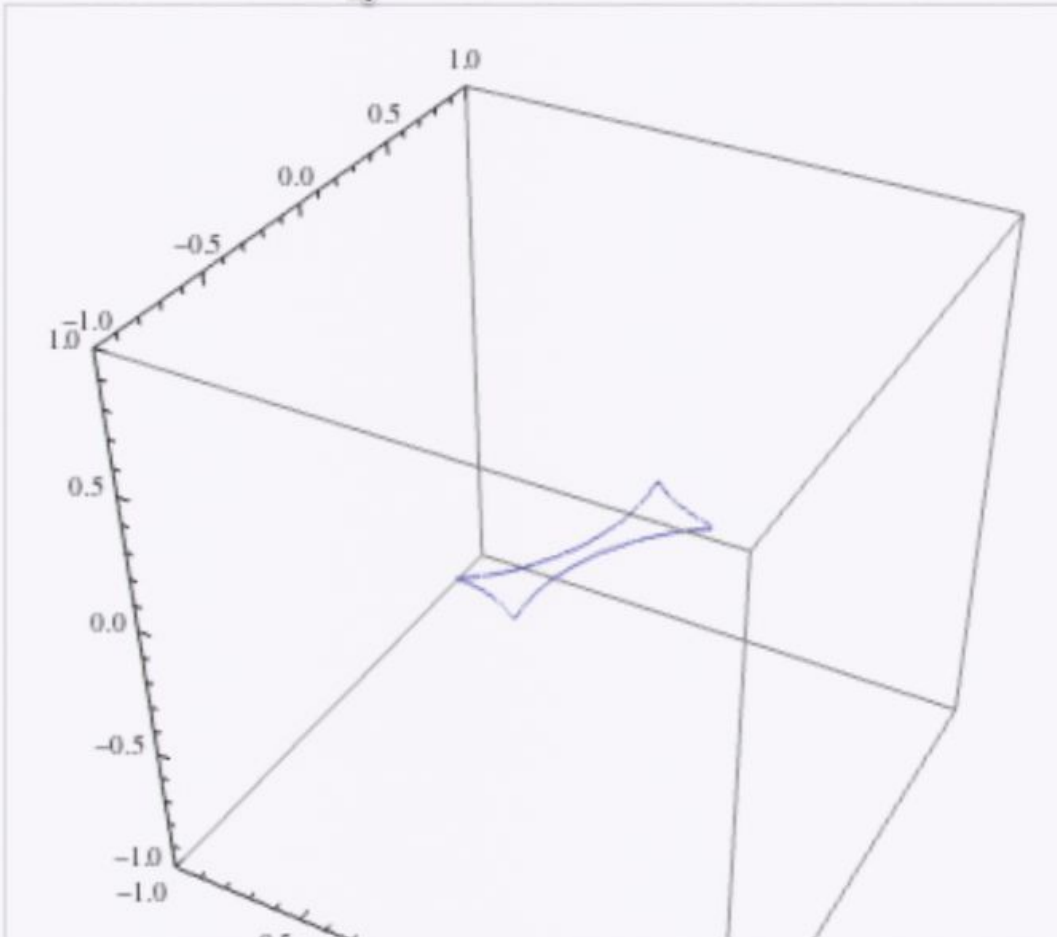
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1.87431



Out[1]=



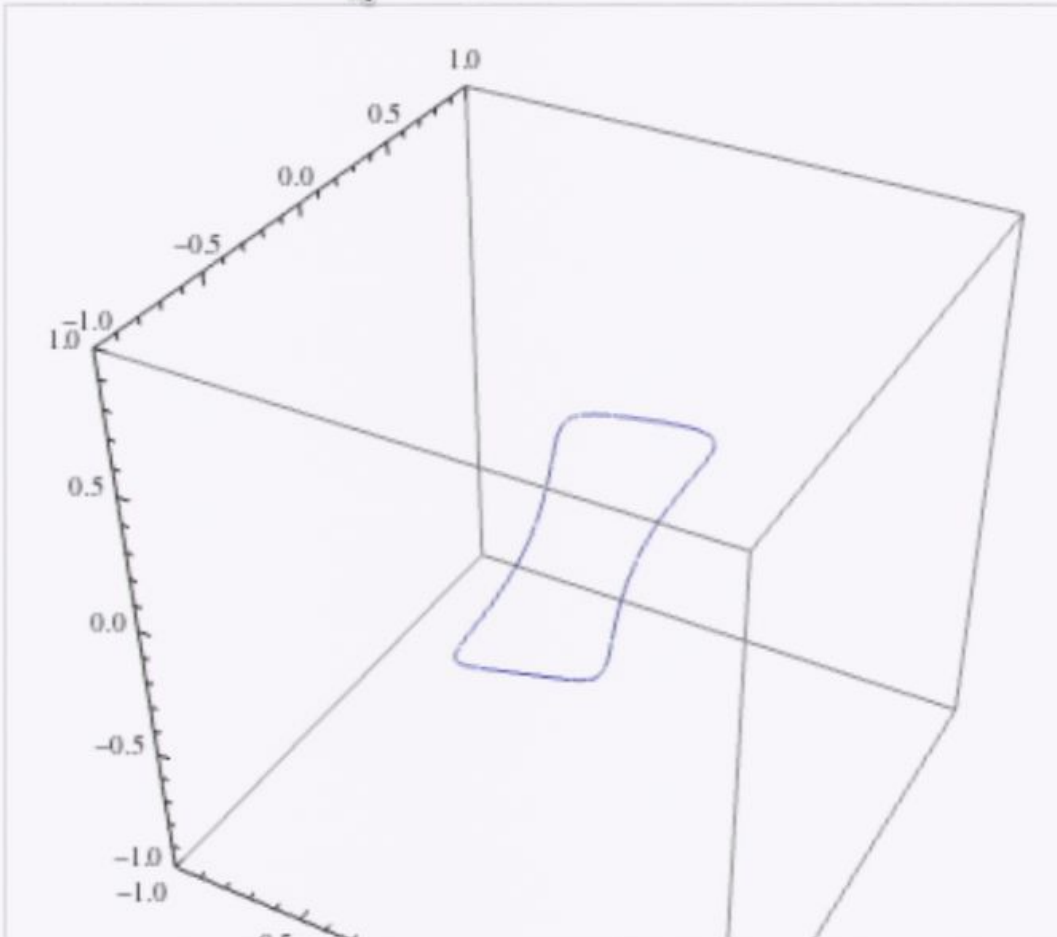




5.35758



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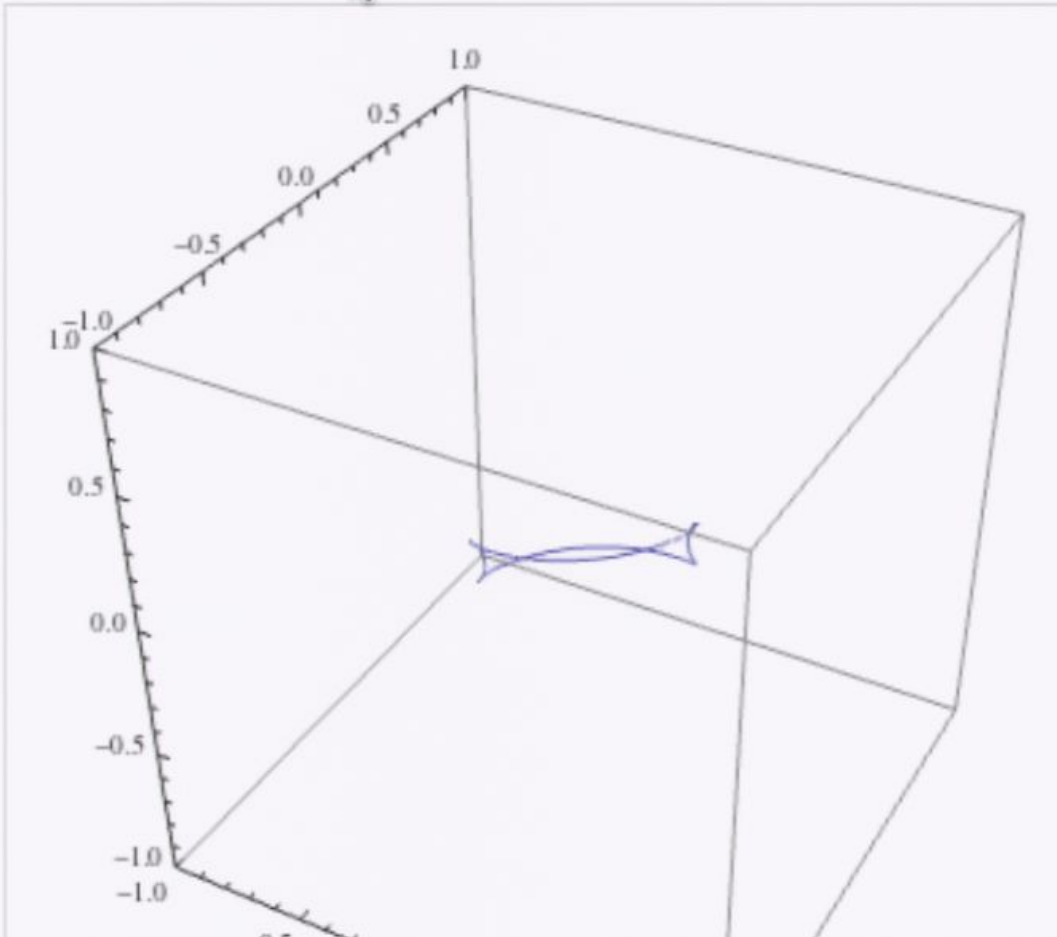




1.69839



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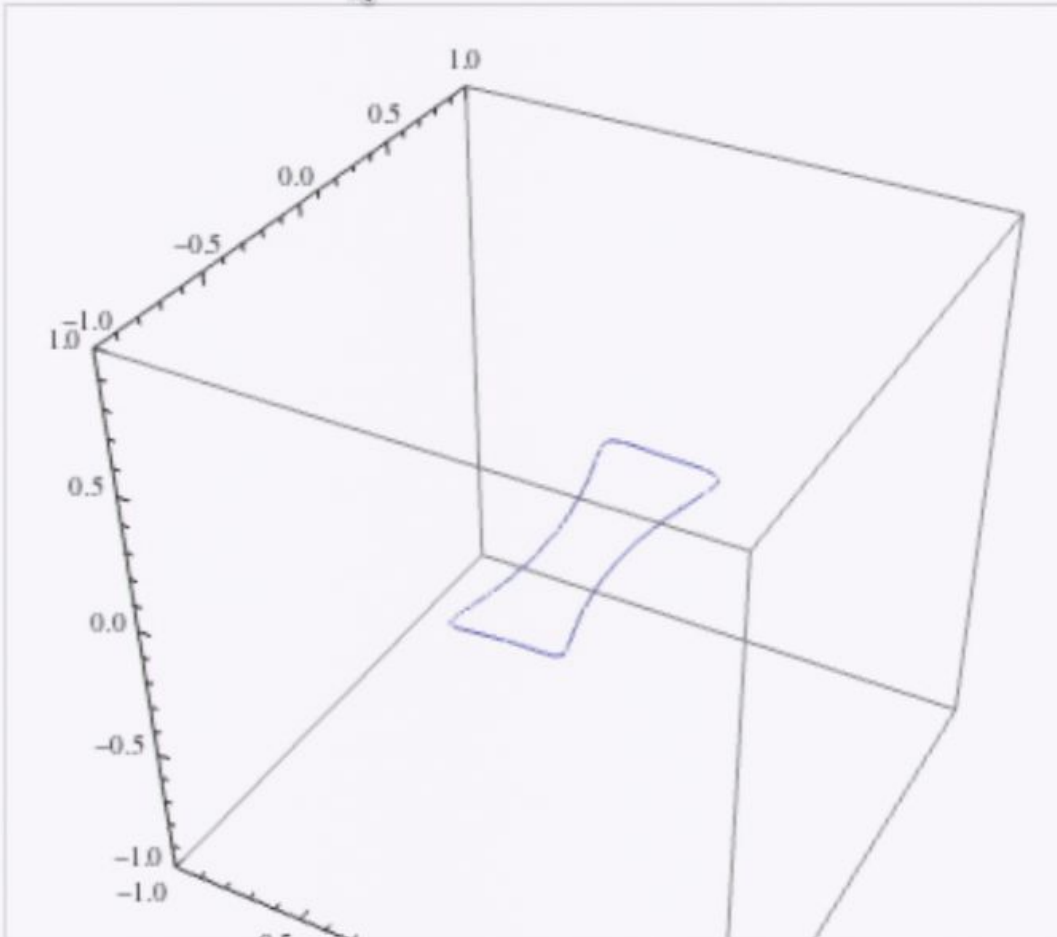




5.09432

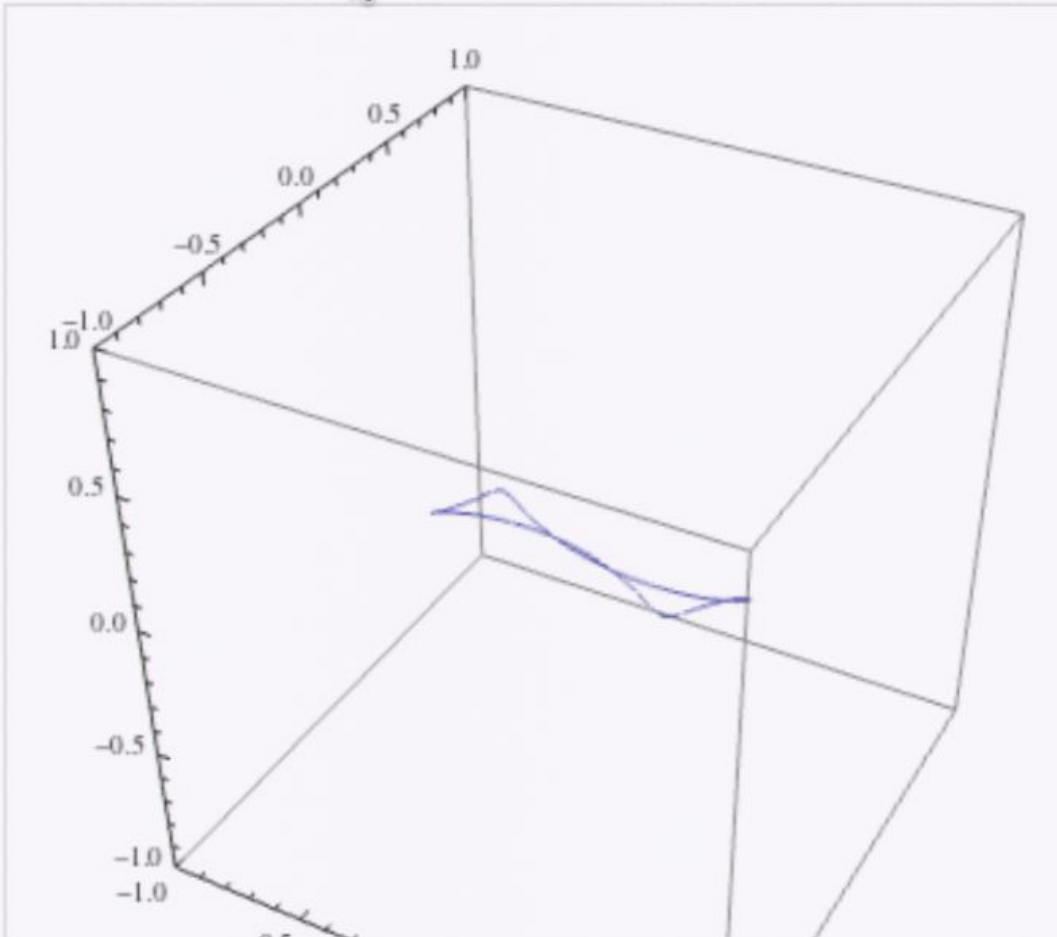


Out[1]=





1.33416



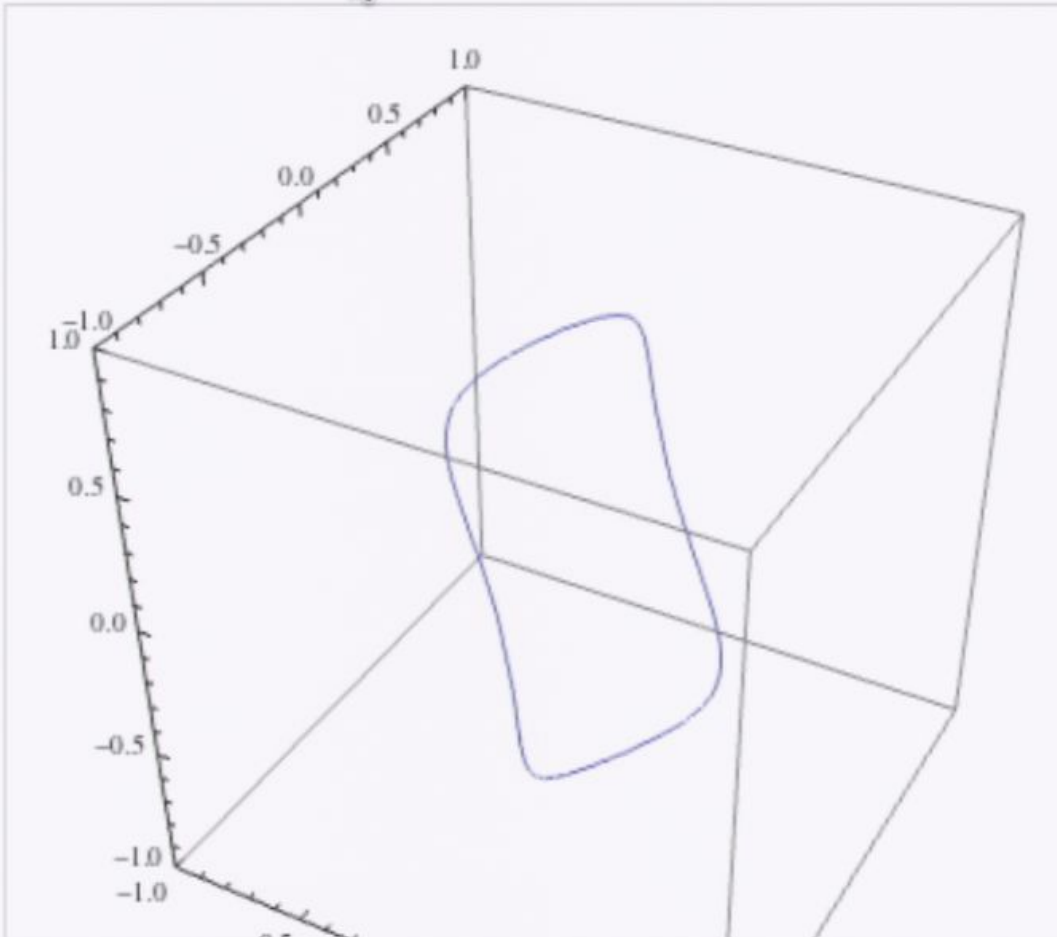
Out[1]=



2.75755

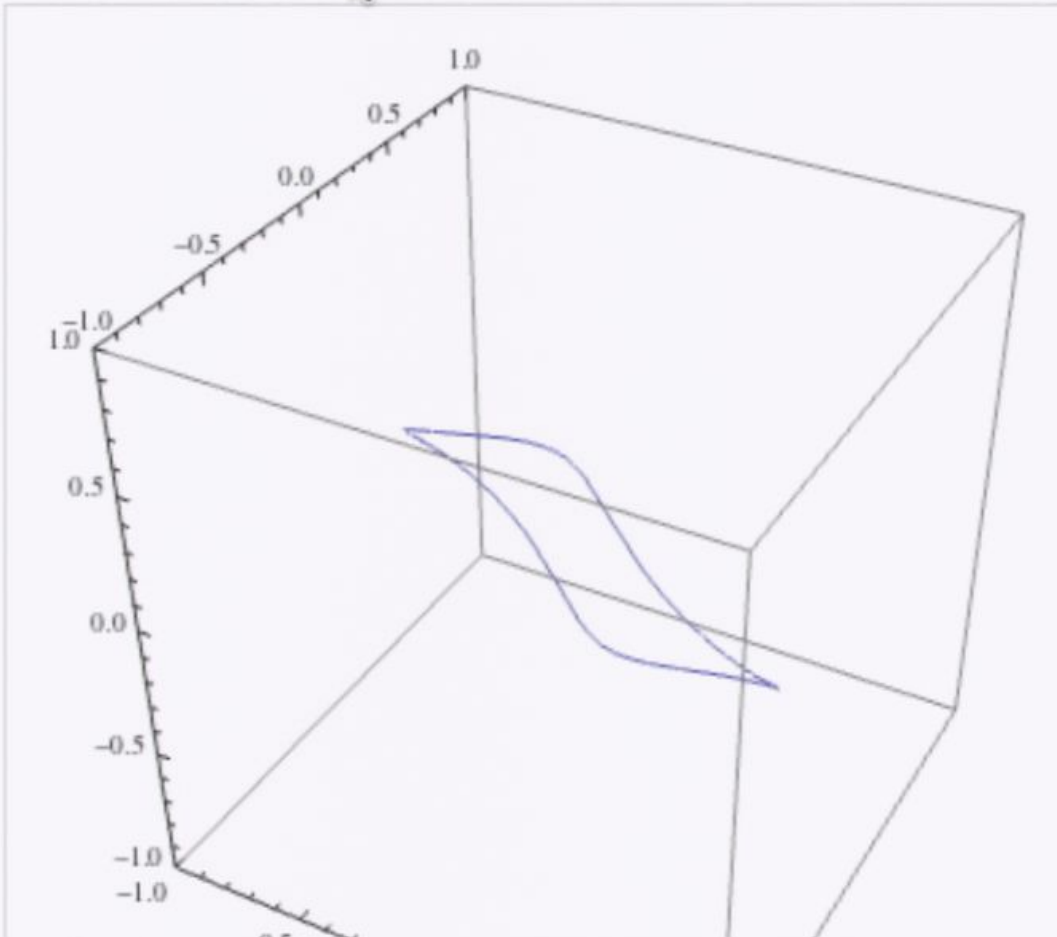


Out[1]=





4.31803 - + >> << →

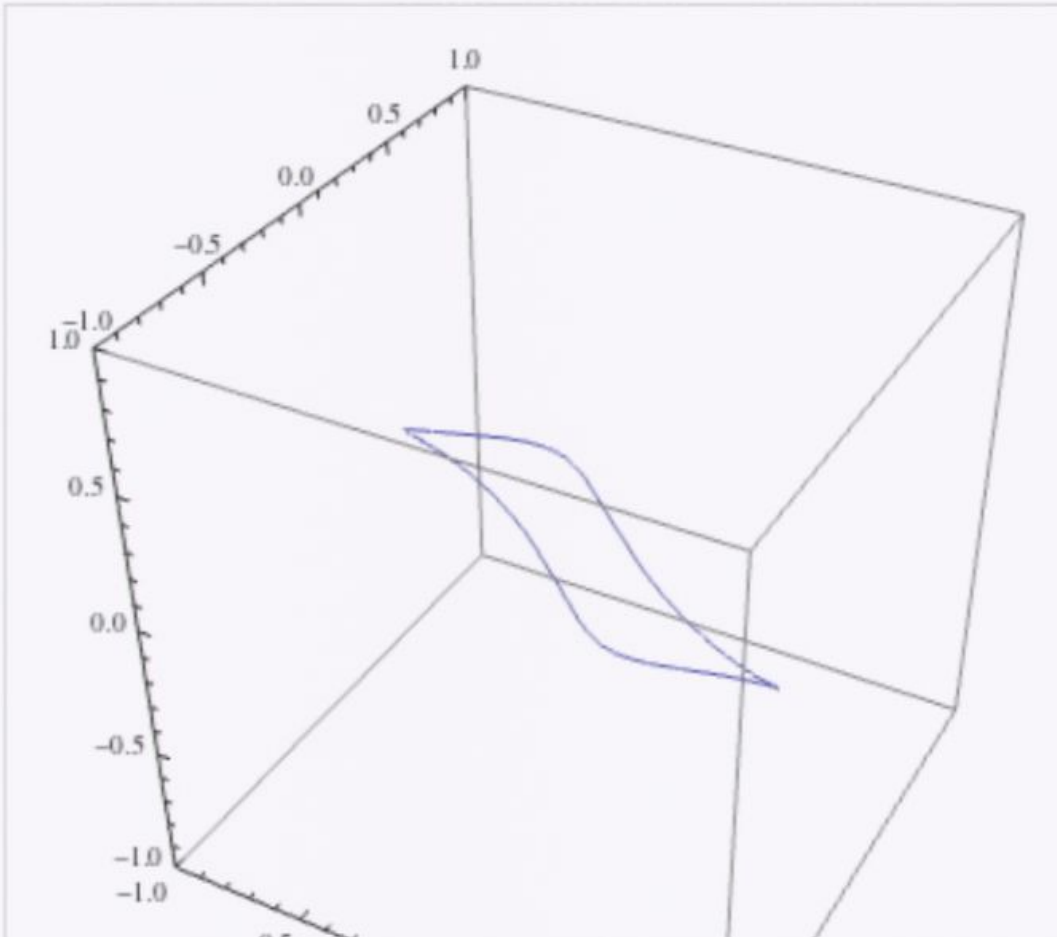




4.31803



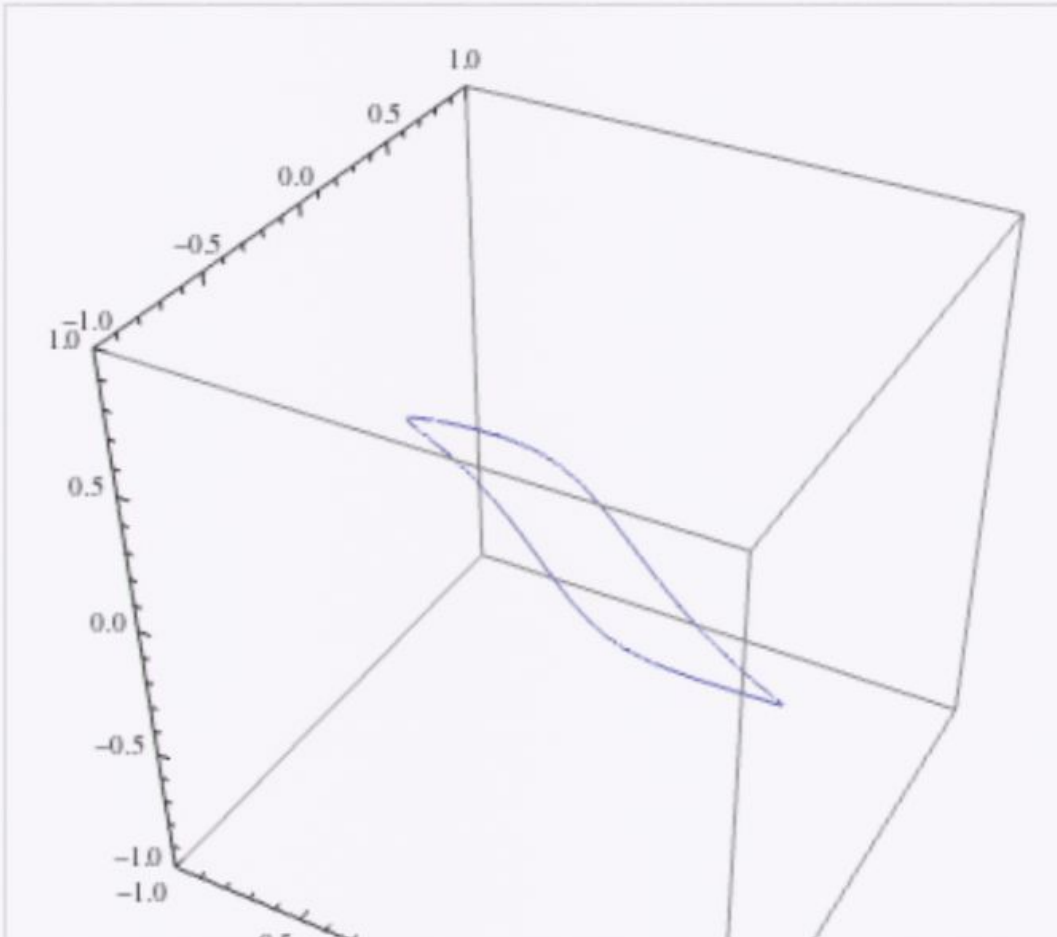
Out[1]=







Out[1]=



## ▼ Chen, DiCarlo, and Hotes (CDH) Strings

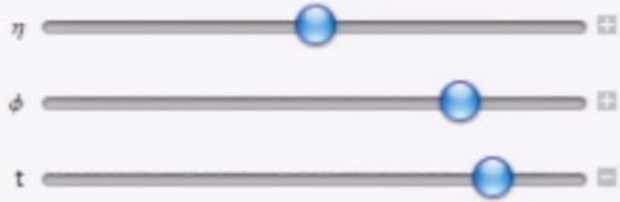
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```
Manipulate[CC = Cos[θ] - Cos[η];
S = Sin[θ] Sin[η];
a[u_] := {CC Cos[θ], S, -CC Sin[θ]} Sin[3 u] / 6 - {-S Cos[θ], CC, S Sin[θ]} Cos[3 u] / 6 +
{2 - CC Cos[θ], -S, CC Sin[θ]} Sin[u] / 2 -
{3 S Cos[θ], Cos[θ] + Cos[η], -Sin[η] (1 - 3 Cos[θ]^2)} Cos[u] / 2;
b[v_] := {1, 0, 0} Sin[v] - {0, Cos[φ], Sin[φ]} Cos[v];
plot1 = ParametricPlot3D[(a[t + s] + b[t - s]) / 2, {s, 0, 2 Pi}, PlotRange → 1],
{{θ, 2 π / 5}, 0, π / 2}, {{η, π / 2}, 0, π}, {{φ, 4 π / 5}, 0, π}, {{t, 0}, 0, 2 π}]
```

## ▼ Rigidly Rotating Strings

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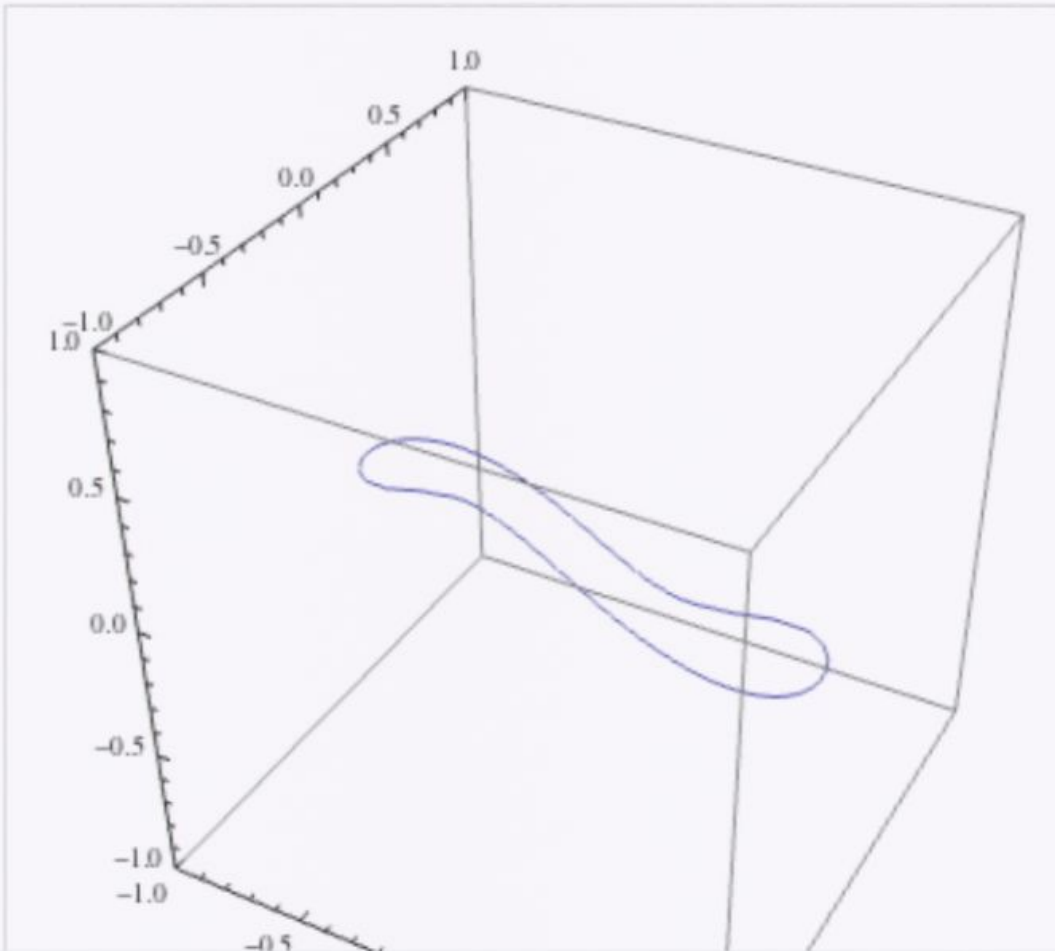
```
Manipulate[
AA = (1 - v^2 + x^2)^2 - 4 x^2:
```



5.42746

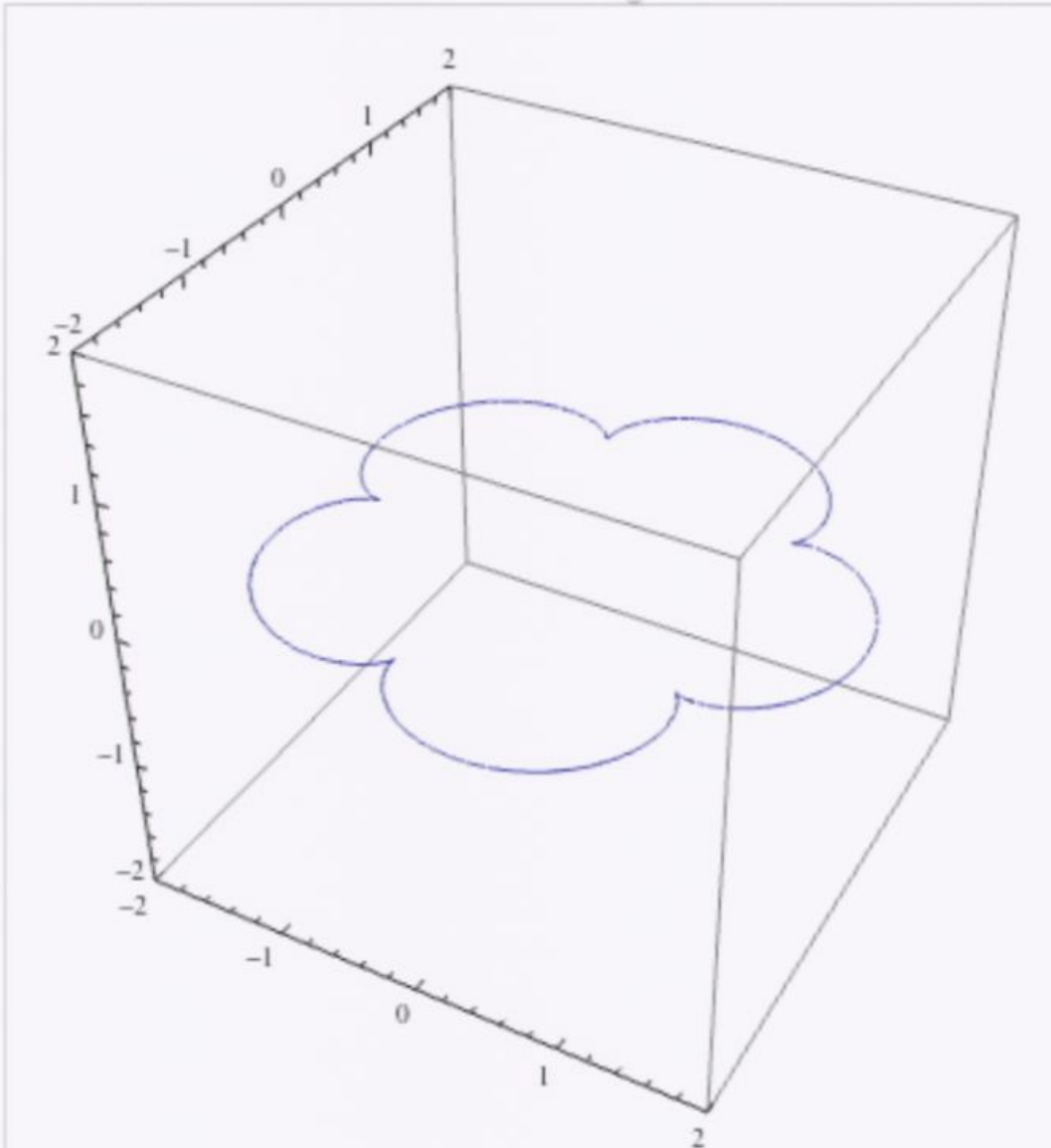


Out[2]=





Out[3]=



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**Cusp Formations**

Extra Dimensions

A New Approach

Conclusions

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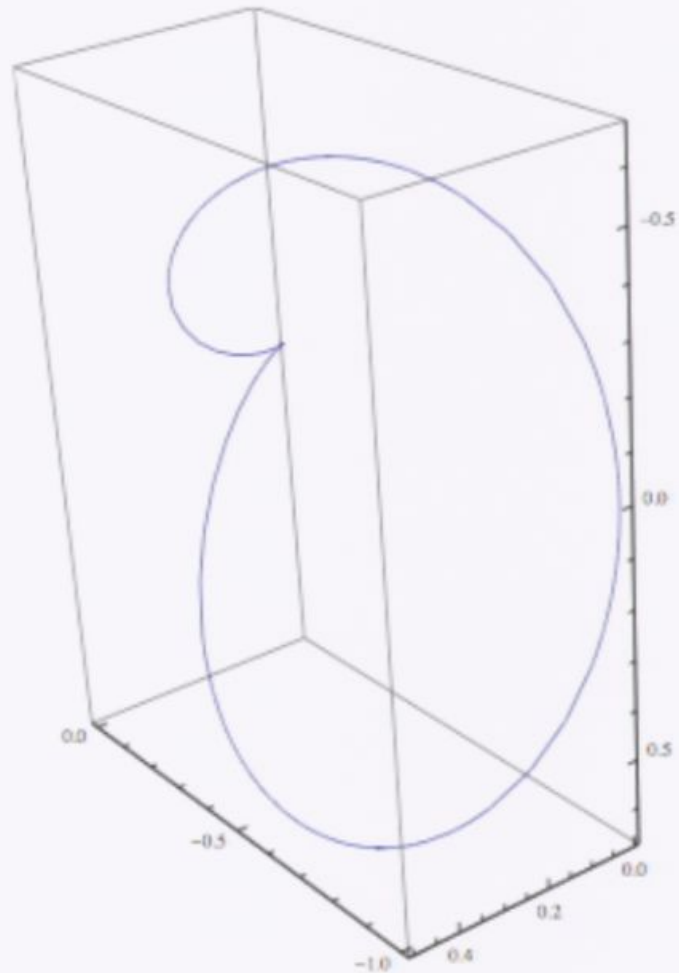
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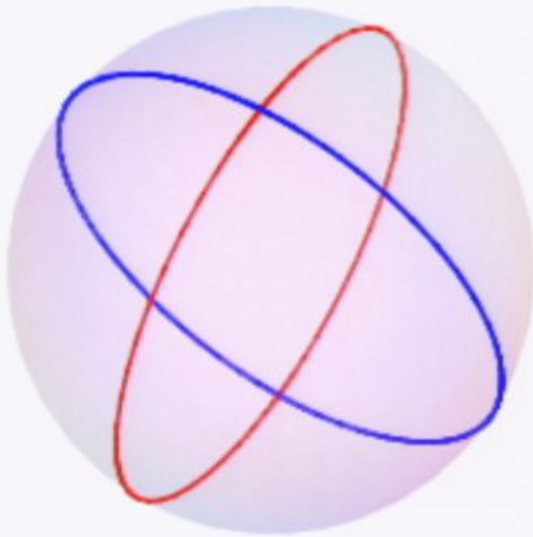
# What is a Cusp?



▶ Points at which the string reaches the speed of light.

▶ Extrinsic curvature diverges

# Cusps as Intersections



- ▶ At intersection,

$$\frac{\partial \mathbf{r}}{\partial t}(t_0, \sigma_0) = \mathbf{a}'(u_0) = \mathbf{b}'(v_0)$$

$$\frac{\partial \mathbf{r}}{\partial \sigma}(t_0, \sigma_0) = \left( \mathbf{a}'(u_0) - \mathbf{b}'(v_0) \right) / 2 = 0$$

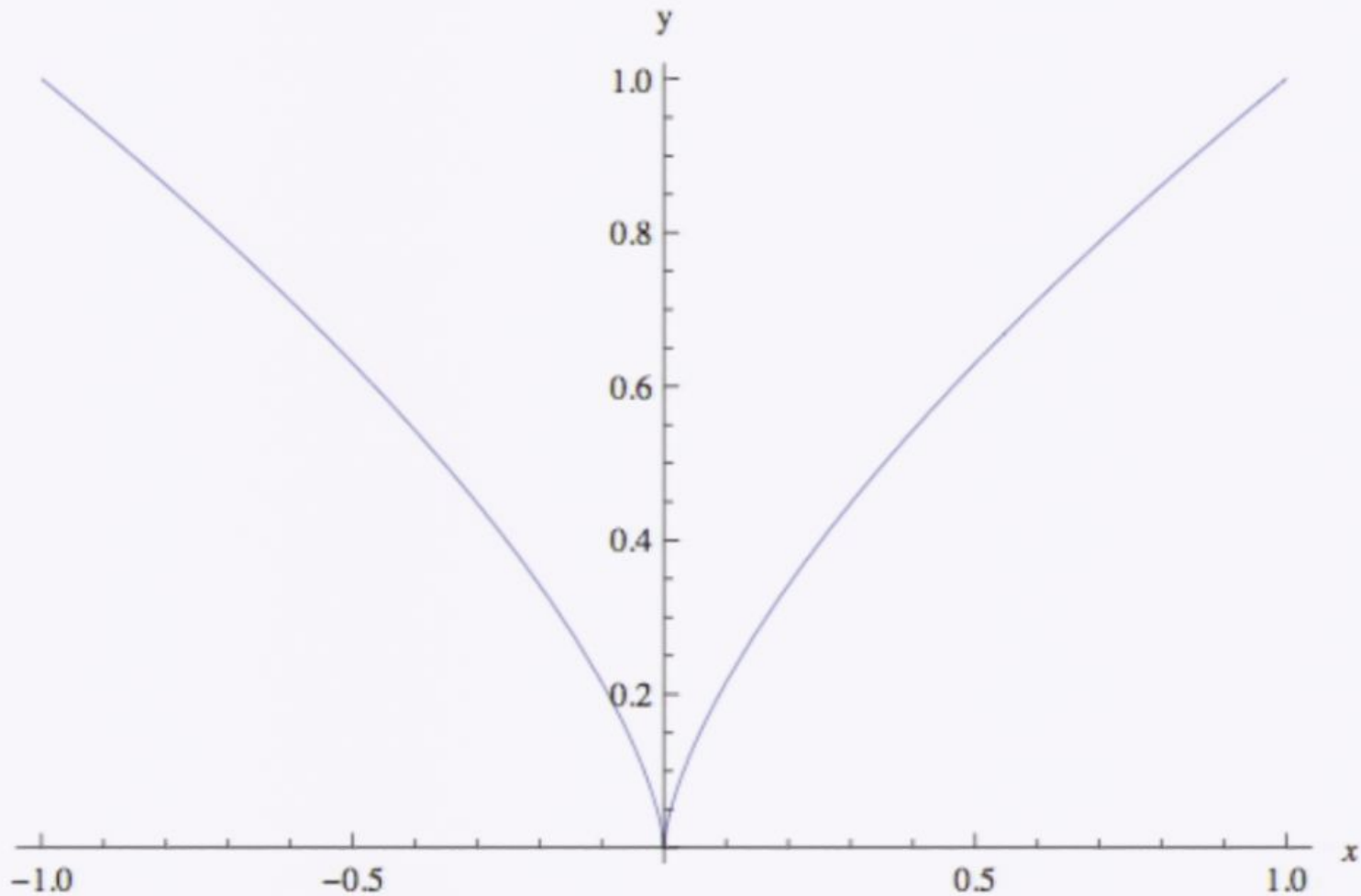
- ▶ Local Expansion:

$$\begin{aligned} \mathbf{r}(t_0, \sigma) - \mathbf{r}_0 &= \frac{(\sigma - \sigma_0)^2}{2} \frac{\partial^2 \mathbf{r}}{\partial \sigma^2} \\ &\quad + \frac{(\sigma - \sigma_0)^3}{6} \frac{\partial^3 \mathbf{r}}{\partial \sigma^3} + \dots \end{aligned}$$

$$\mathbf{a}'(u_0) = \mathbf{b}'(v_0)$$

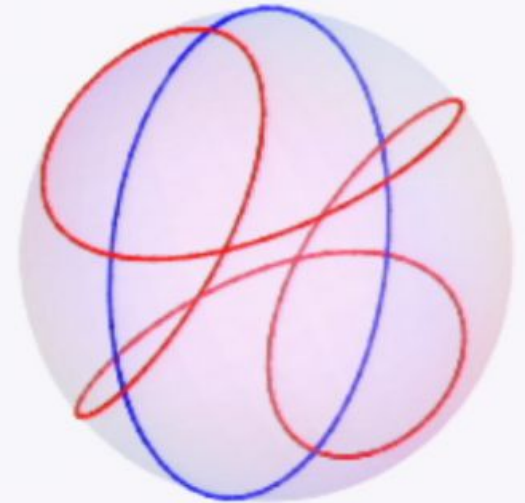
# Prototype of all Cusps

$$\mathbf{r}(s) = (s^3, s^2) \implies y = x^{2/3}$$



## How many cusps?

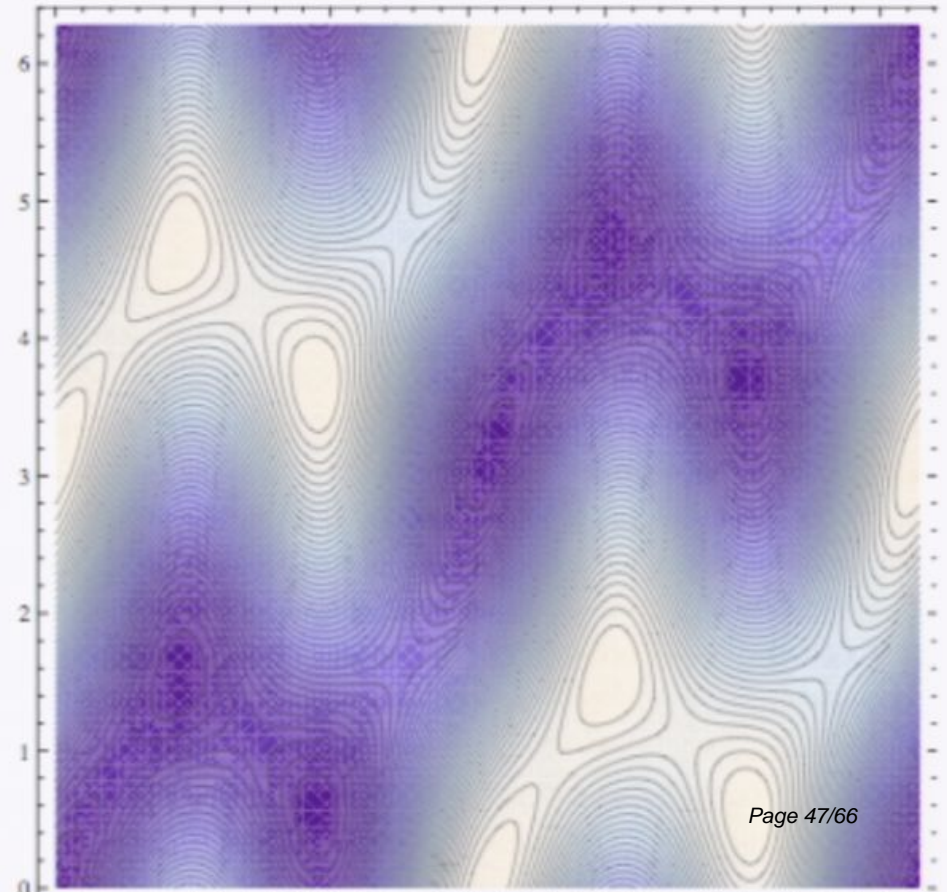
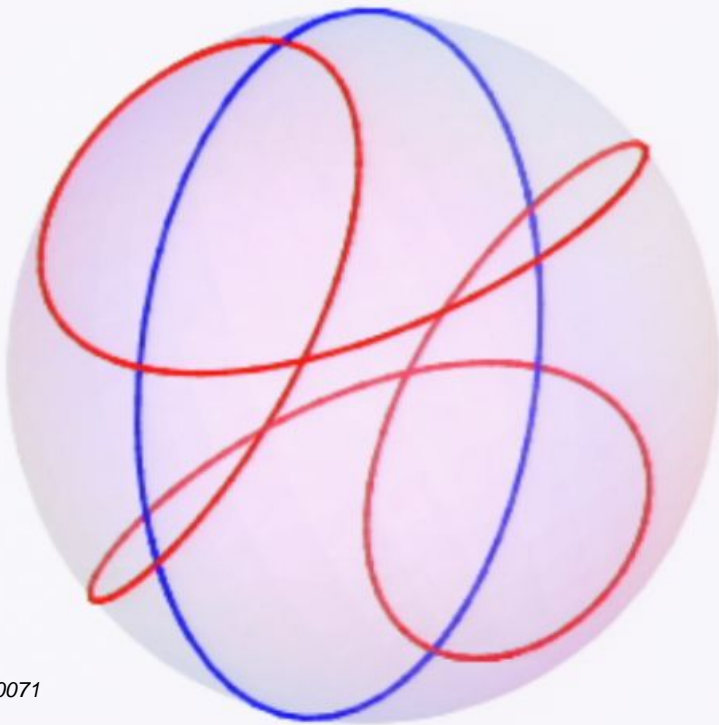
- ▶ It's important to model the number of cusps per oscillation
- ▶ String evolution is fairly complicated, we have to resort to numerical methods
- ▶ Classical solutions already exhibit a variety of behavior



# Cusp Counting as Minimization

Want to minimize

$$f(u, v) = |\mathbf{a}'(u) - \mathbf{b}'(v)|^2$$





# Minimization Techniques

- ▶ **Gradient Descent:**

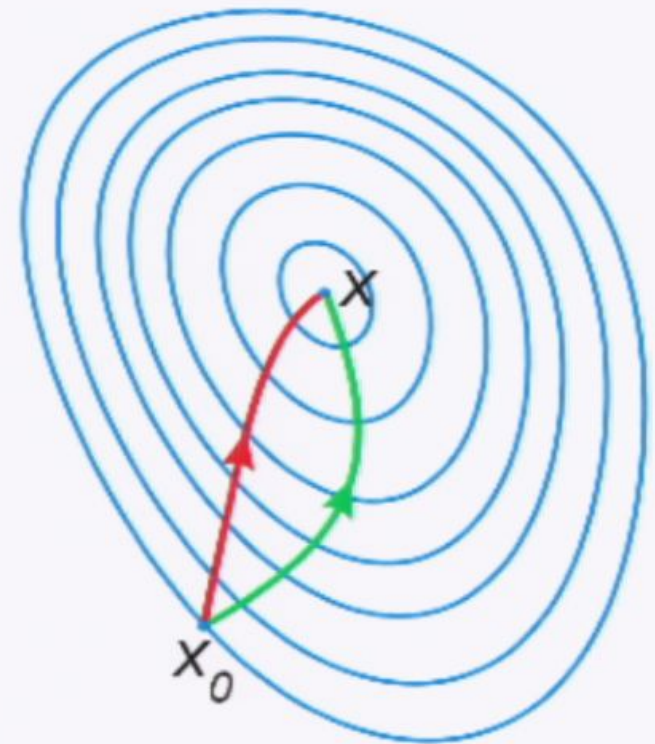
$$x_{n+1} = x_n - \gamma_n \nabla F(x_n)$$

for some small enough  $\gamma_n$

- ▶ **Newton Method:**

$$x_{n+1} = x_n - \frac{f'(x_n)}{f''(x_n)}$$

- ▶ If we start close enough to a minimum, the sequence will converge.



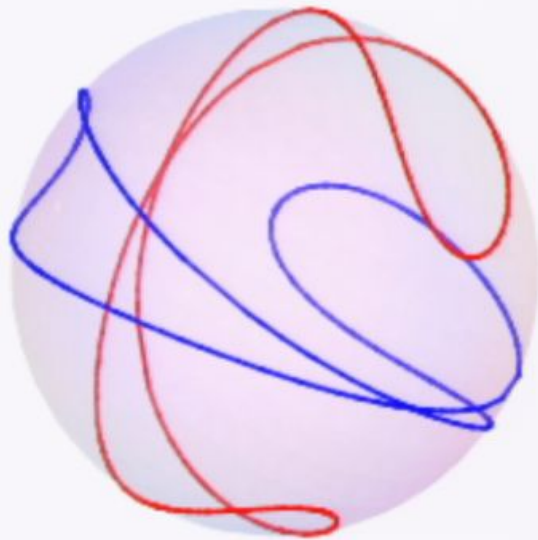


## Our Strategy

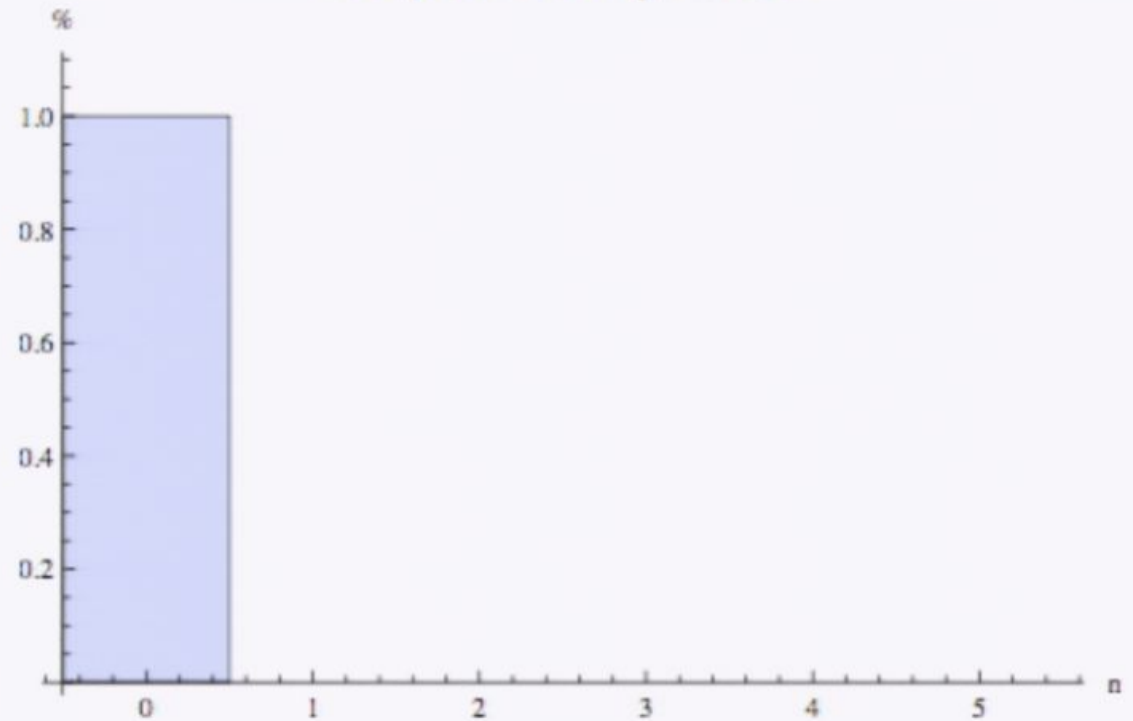
1. Simulate a number of string solutions from a particular family
2. For each string,
  - 2.1 Throw a random points on the  $(u, v)$  space
  - 2.2 Run our minimization procedure, record (in increasing order) any new local minimum we find
  - 2.3 Keep trying until we don't get any more new minima
  - 2.4 Count the number of minima below a cutoff
3. Average over all string solutions

# Numerical Results - checking the simple cases

## Cuspless Strings:

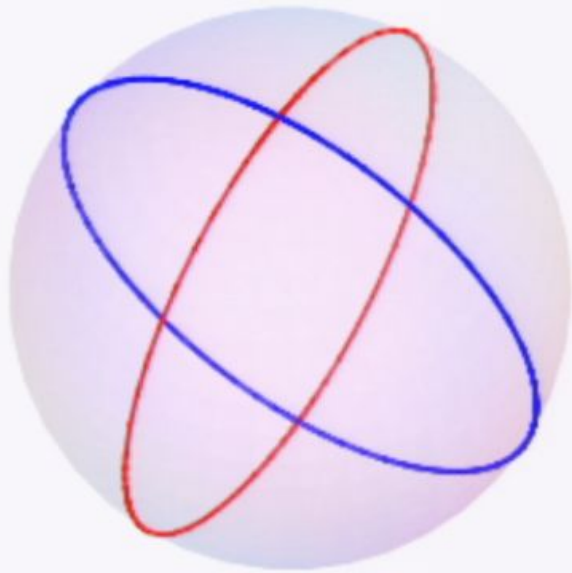


## Cuspless Strings in 3D

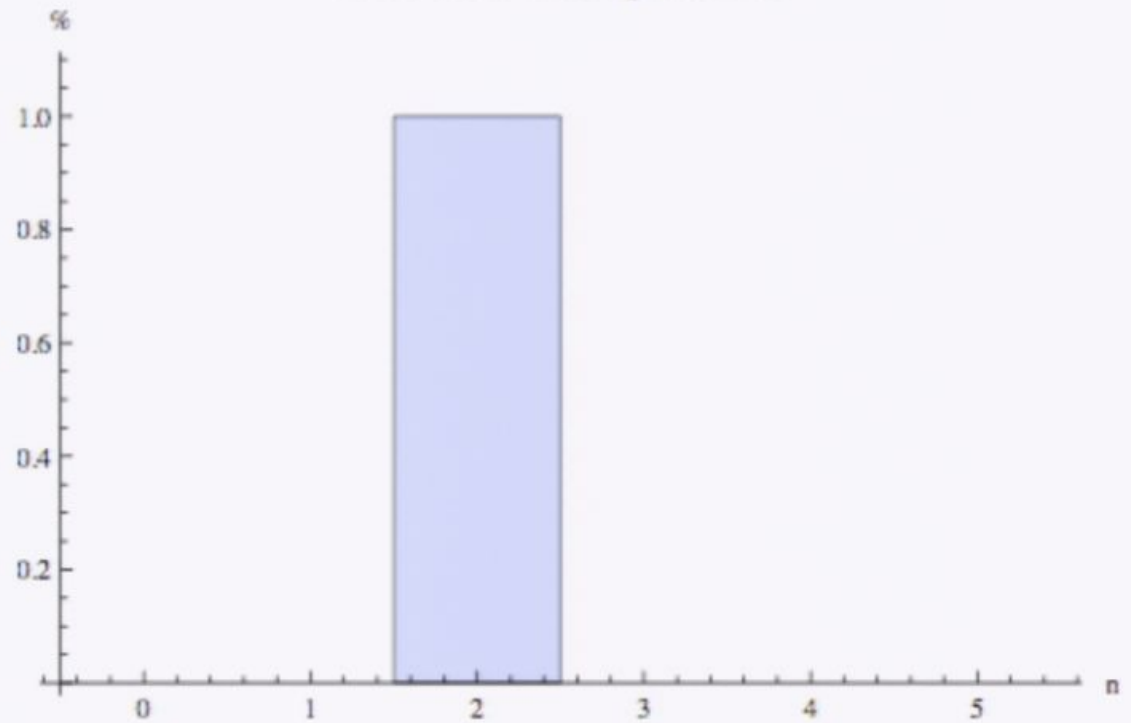


# Numerical Results - checking the simple cases

Simplest Siemens Strings:

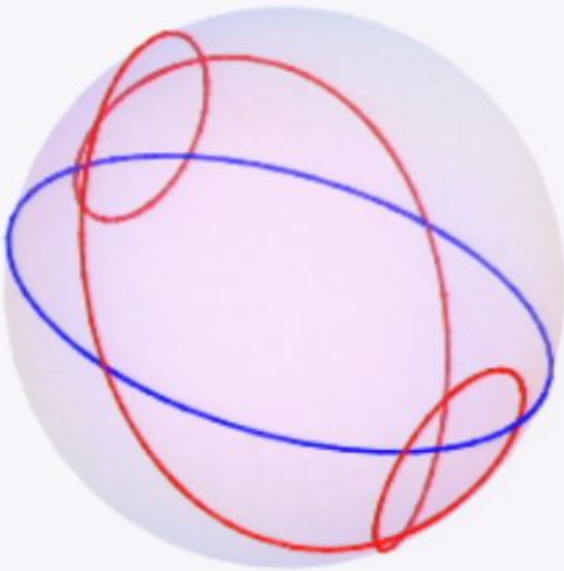


Siemens Strings in 3D

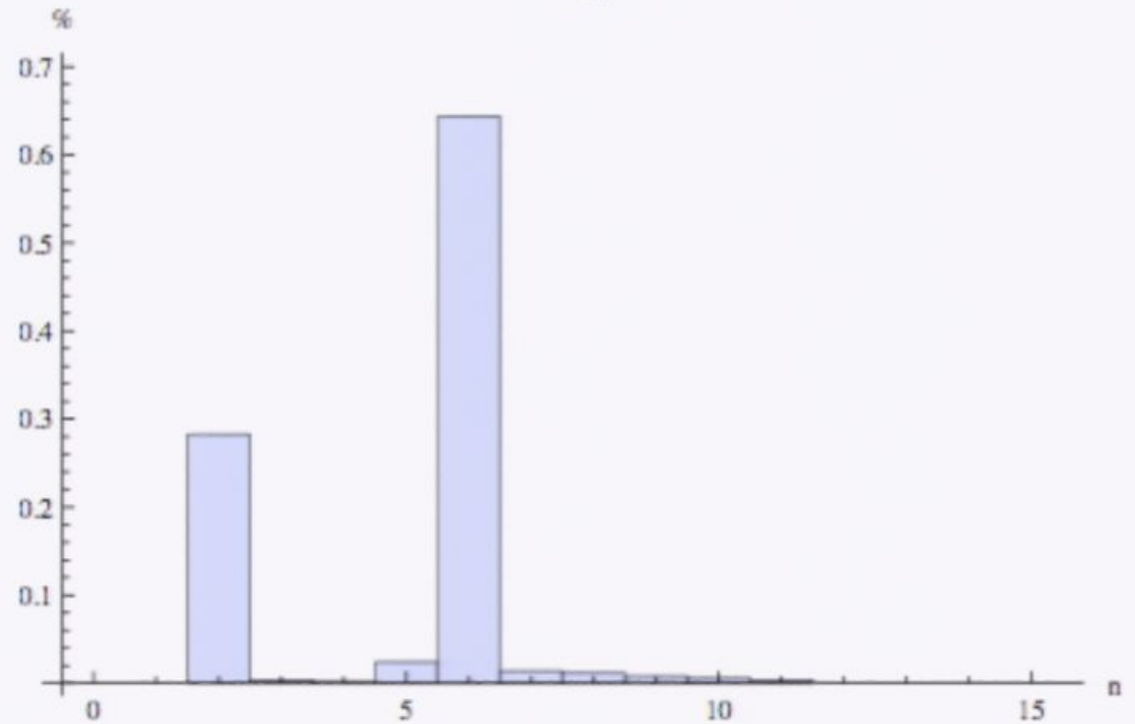


# Numerical Results

CDH Strings:



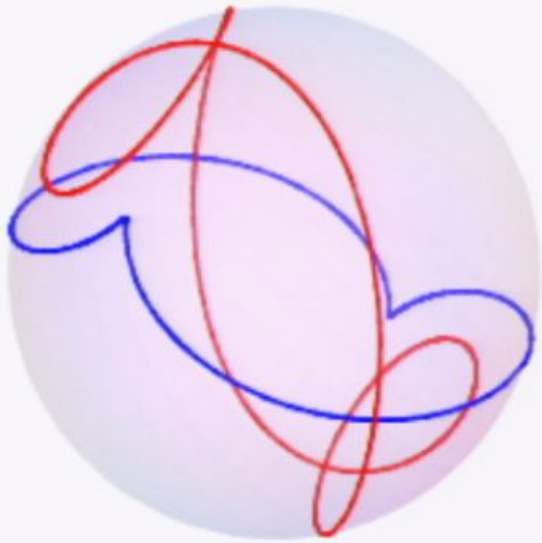
CDH Strings in 3D



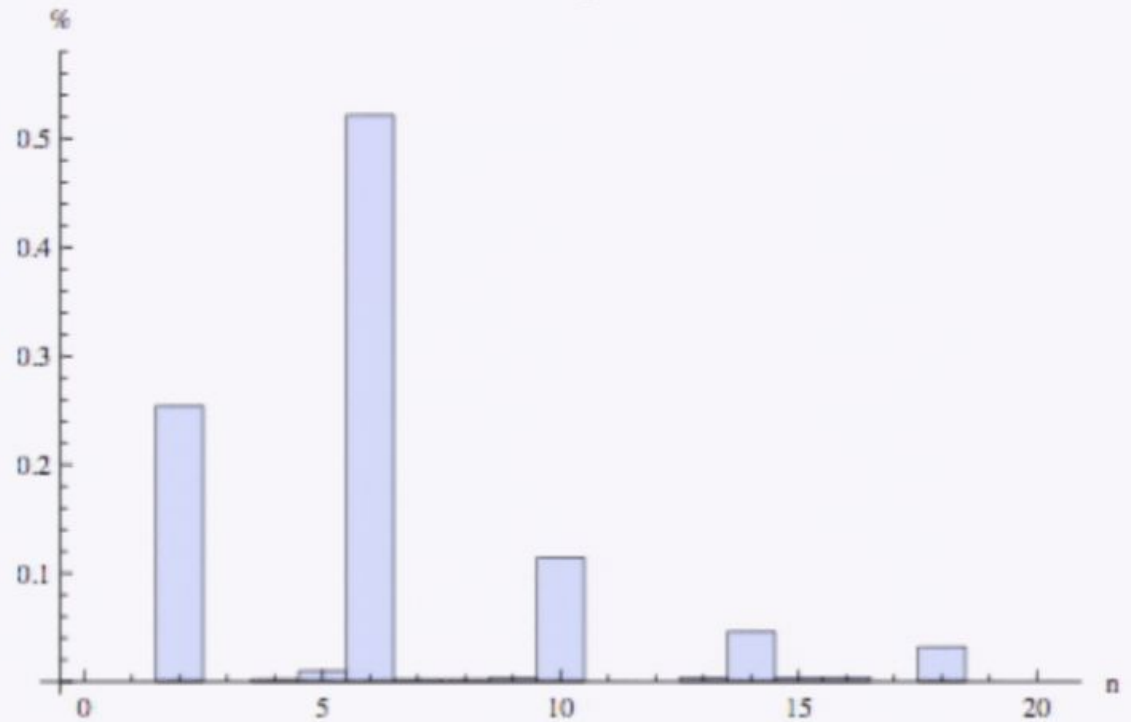
Average = 4.96

# Numerical Results

DES Strings:



DES Strings in 3D



Average = 6.3

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Conclusions



## No More Cusps

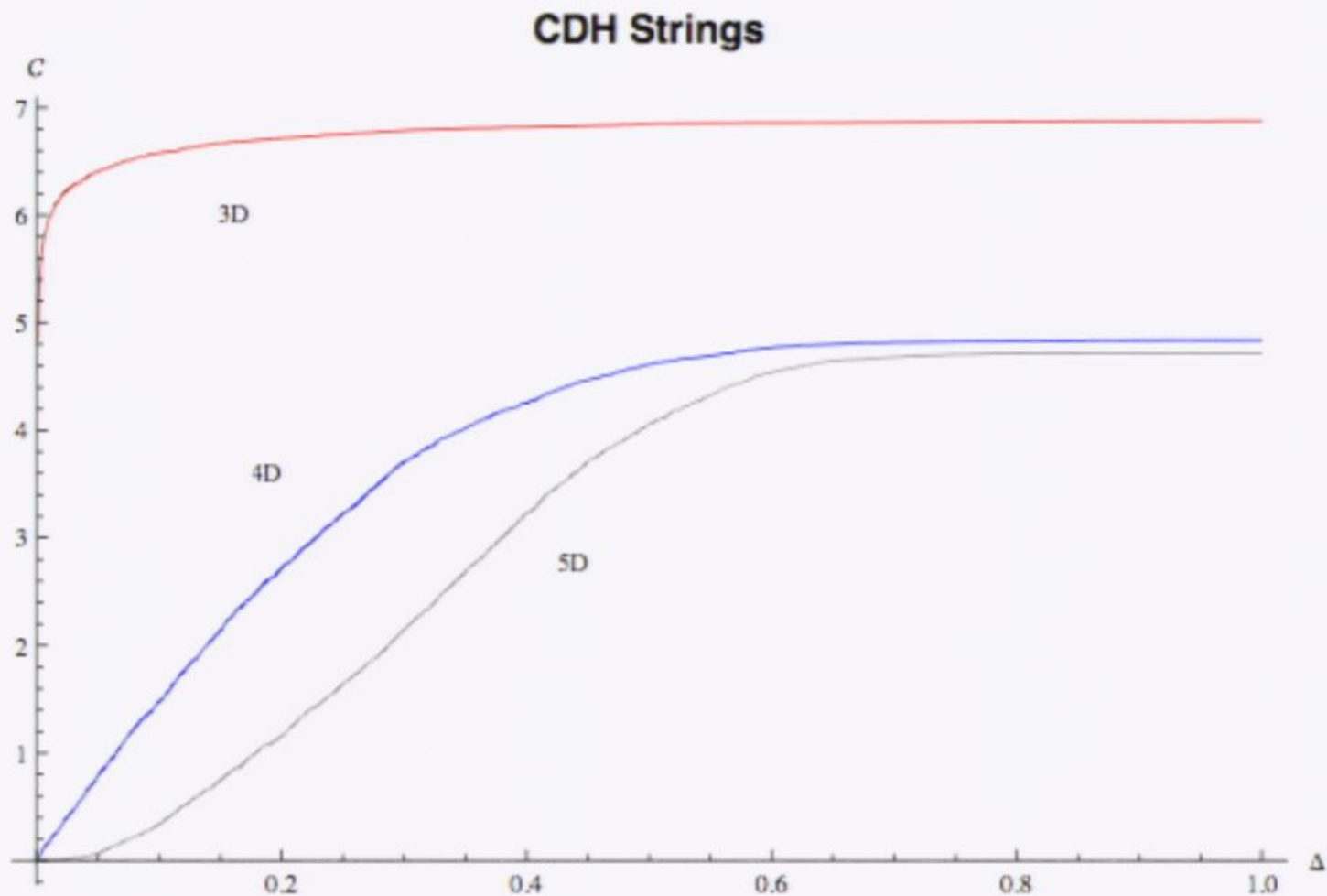
- ▶ Two smooth curves in  $S^3$  don't intersect. (Need to solve three equations but only have two degrees of freedom.)
- ▶ In the presence of extra dimensions, the probability of exact cusp is strictly zero.
- ▶ Near Cusp Events

$$\Delta = \frac{1}{2} \min |\mathbf{a}'(u) - \mathbf{b}'(v)|$$



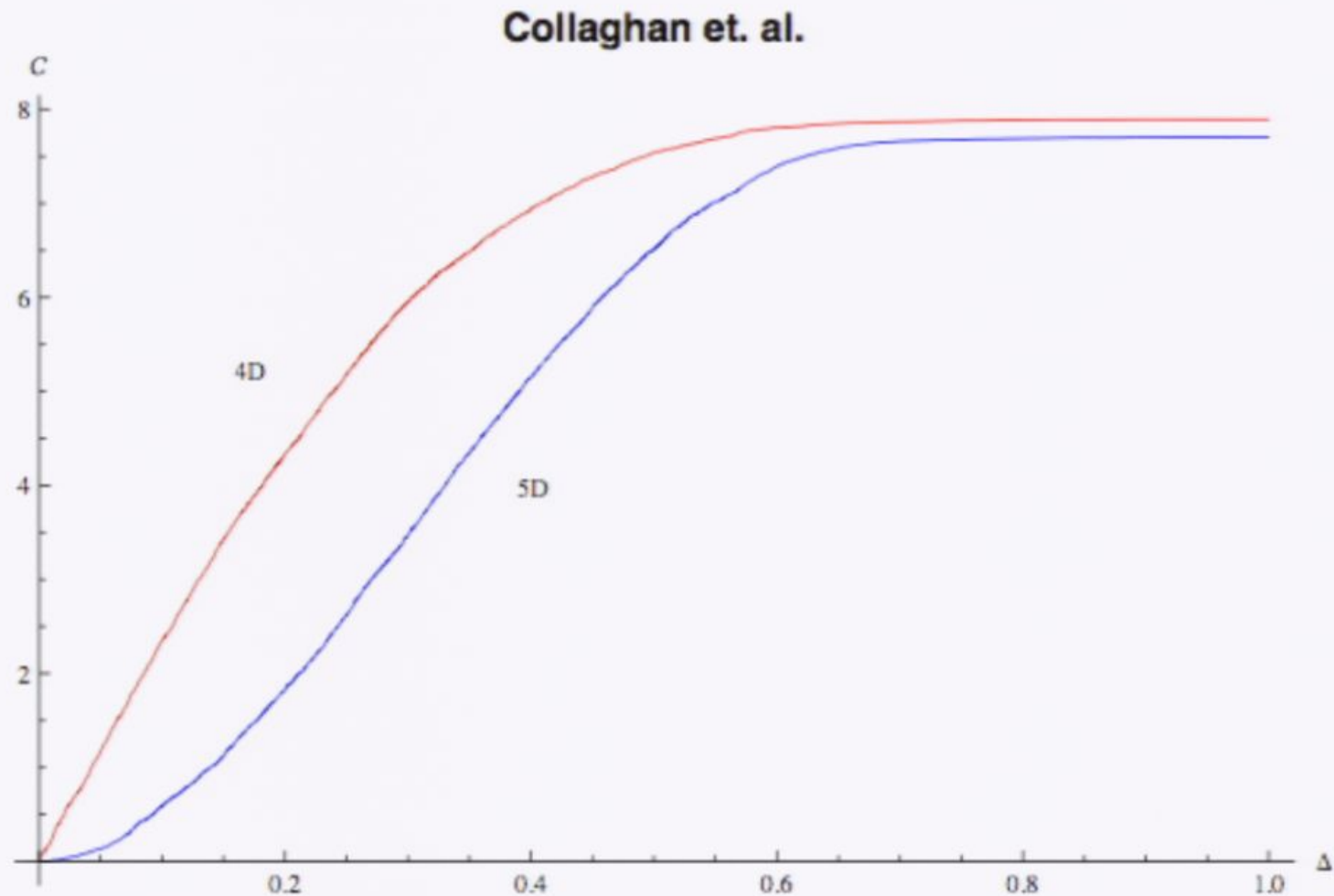
# Results - number of dimensions

## CDH Strings in 3, 4, and 5 spatial dimensions

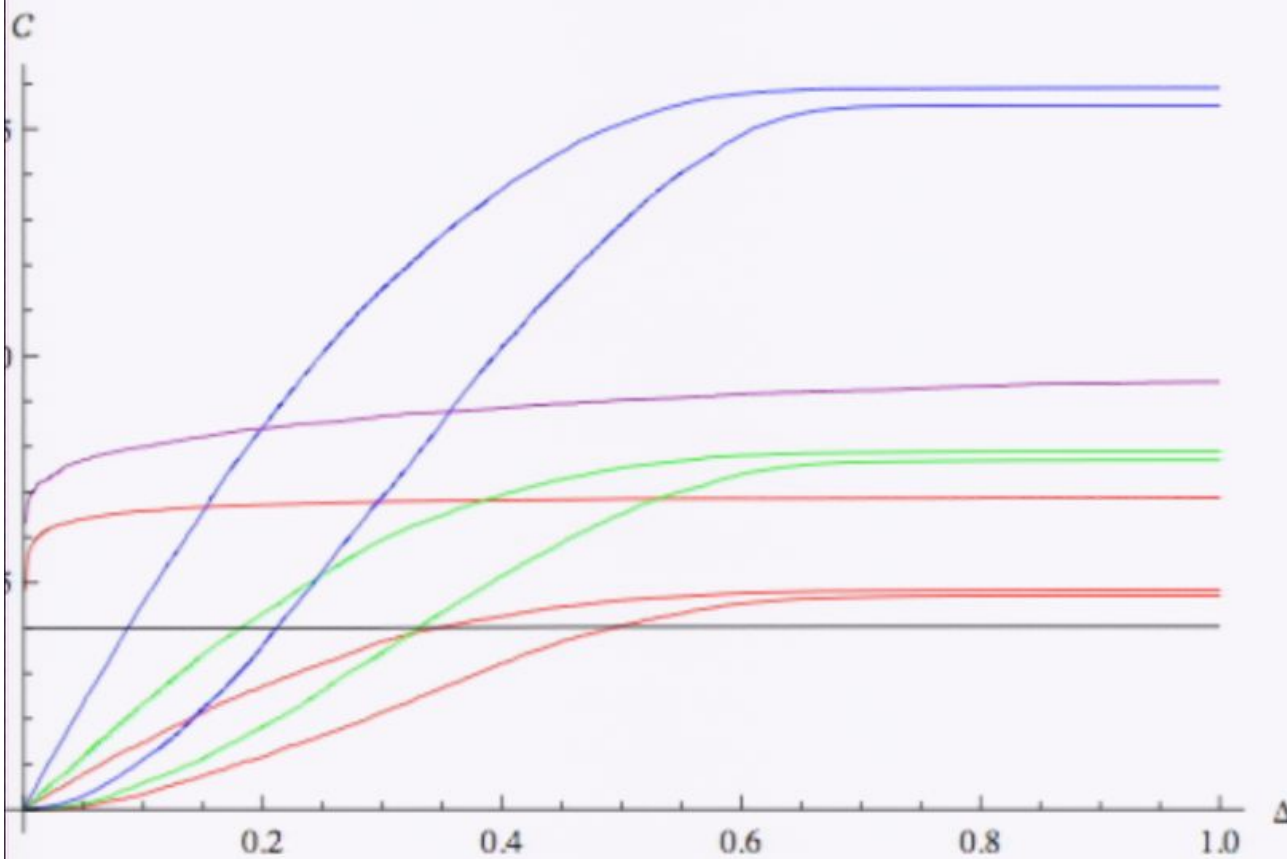


# Results - number of dimensions

Collaghan *et al.* proposed the 1-5/1 Strings



# Results - overview



- ▶ CDH Strings
- ▶ 1-10/1 Strings
- ▶ DES Strings
- ▶ Siemens Strings
- ▶ Collaghan Strings

# Outline

Introduction and Motivation

String Dynamics

Cusp Formations

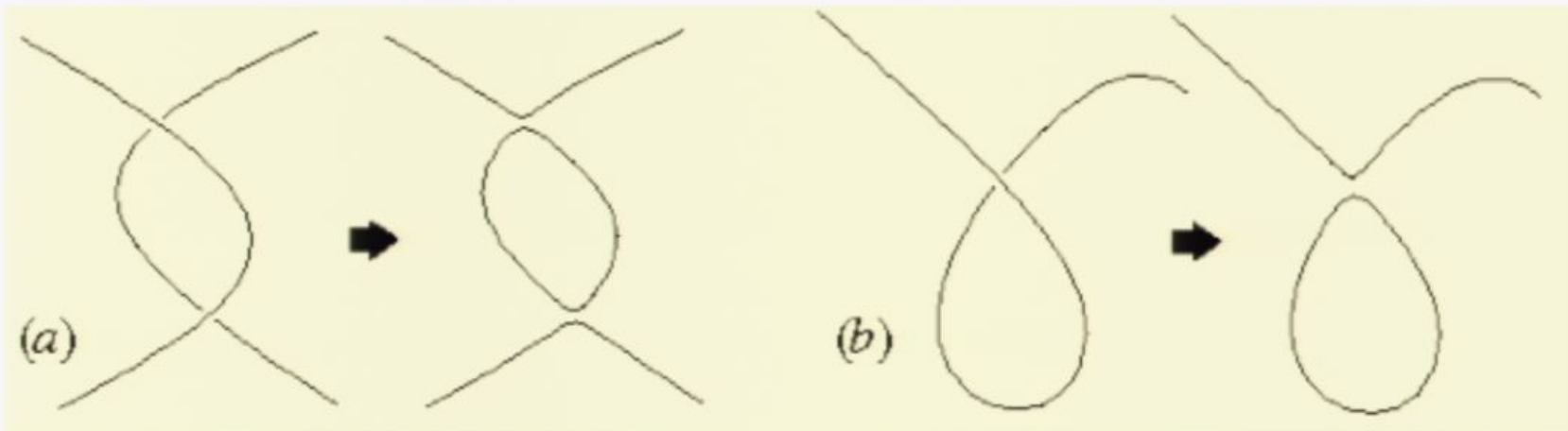
Extra Dimensions

**A New Approach**

Conclusions

# String Intercommutation

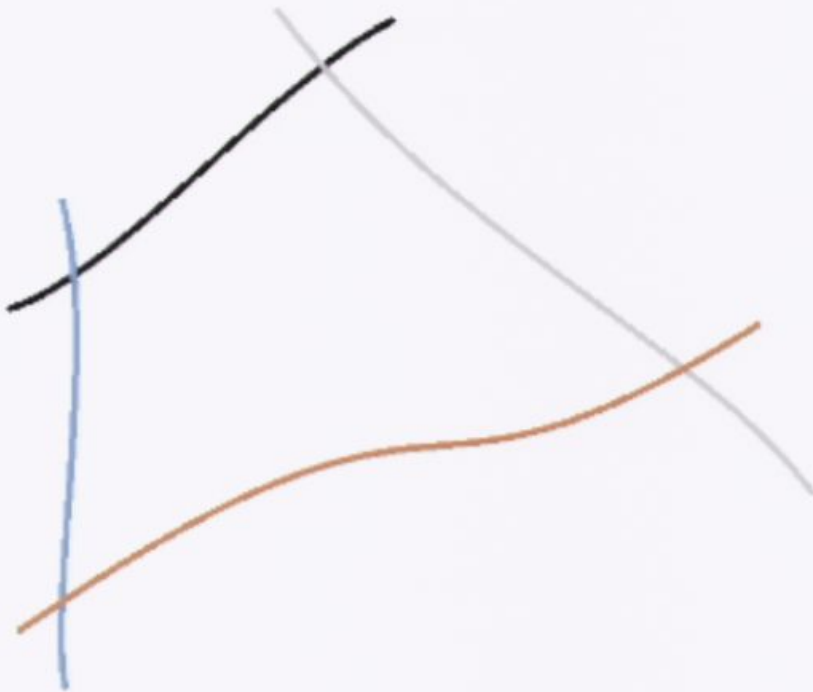
- ▶ Whenever two long strings cross each other, they exchange ends, or “intercommute”.
- ▶ A long string can intercommute with itself, in which case a loop will be produced



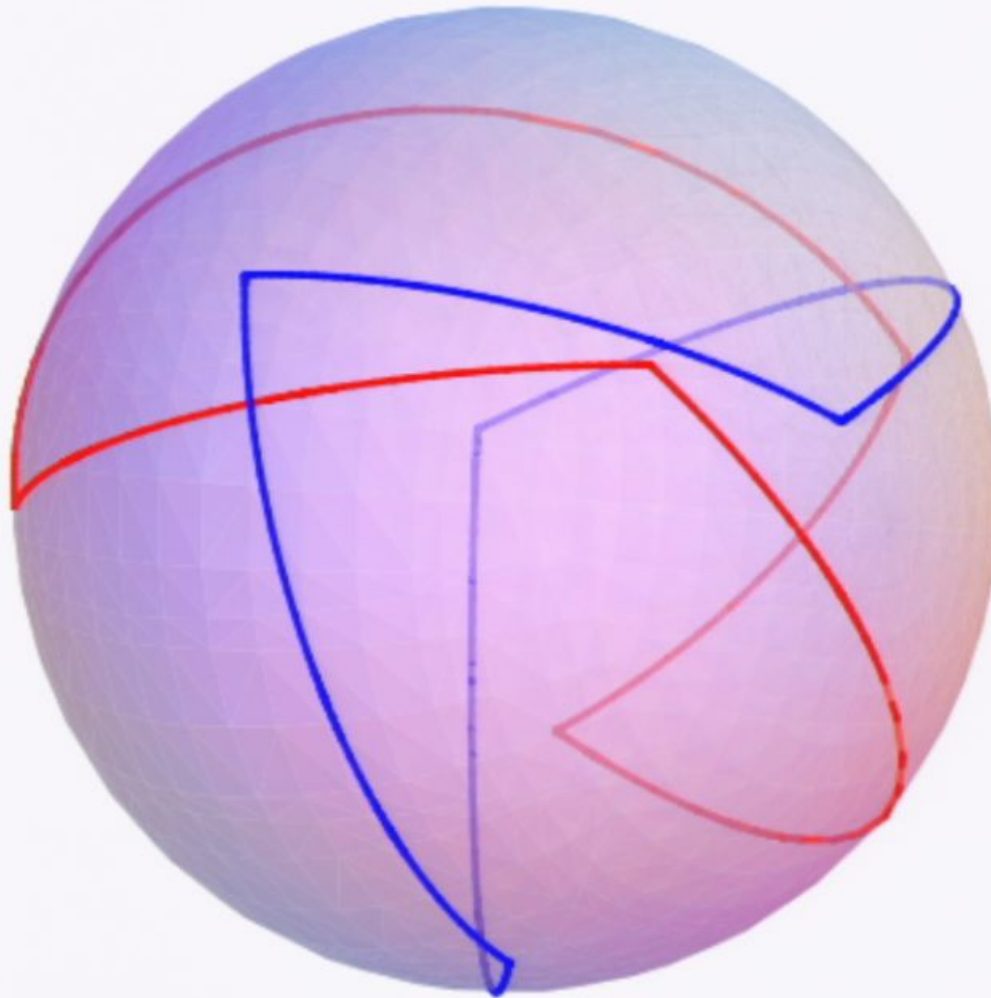


# Network of Strings

From numerical simulation, we model loops as piecewise linear.

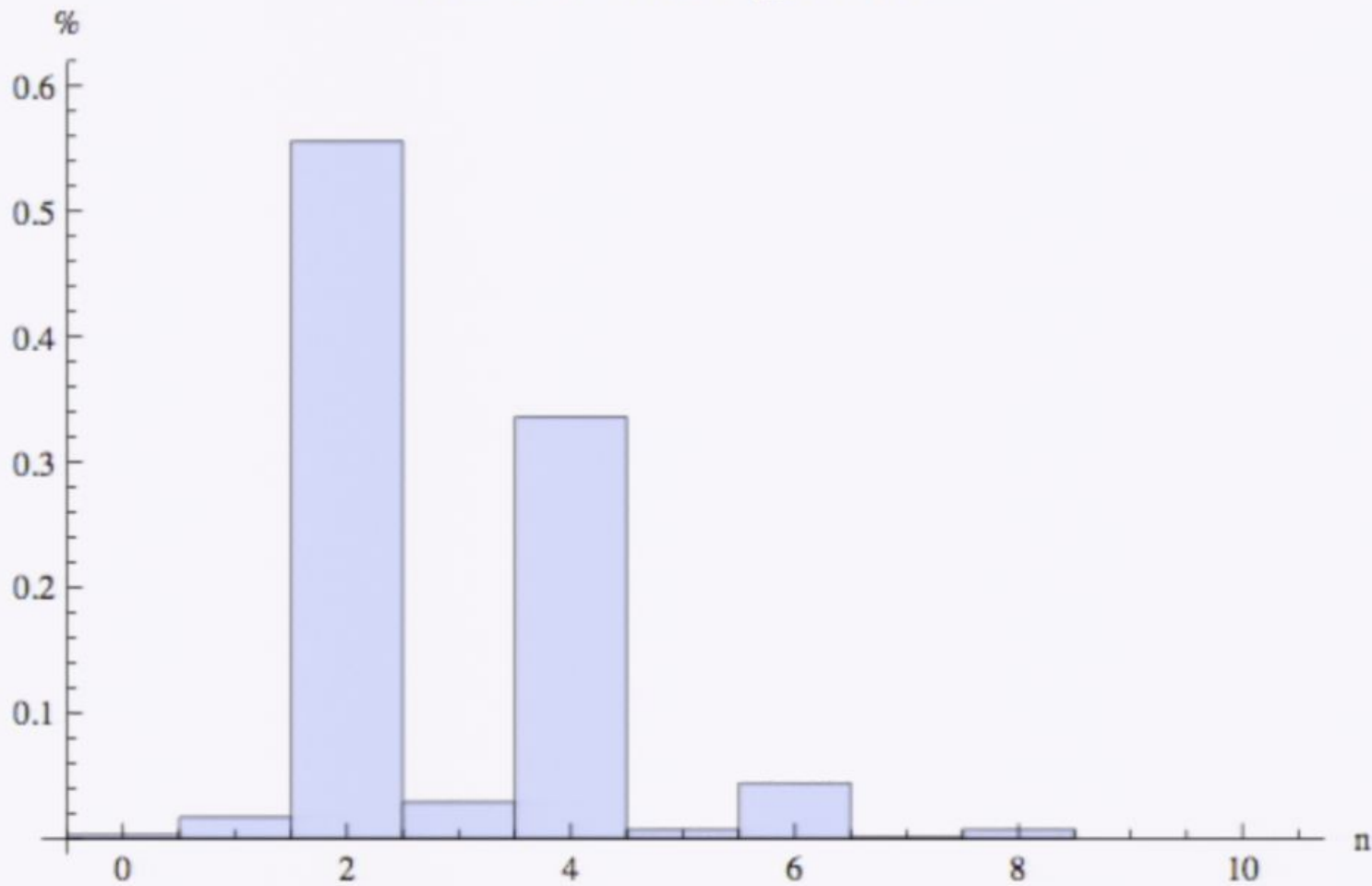


# Random Curves on Kibble-Turok Sphere



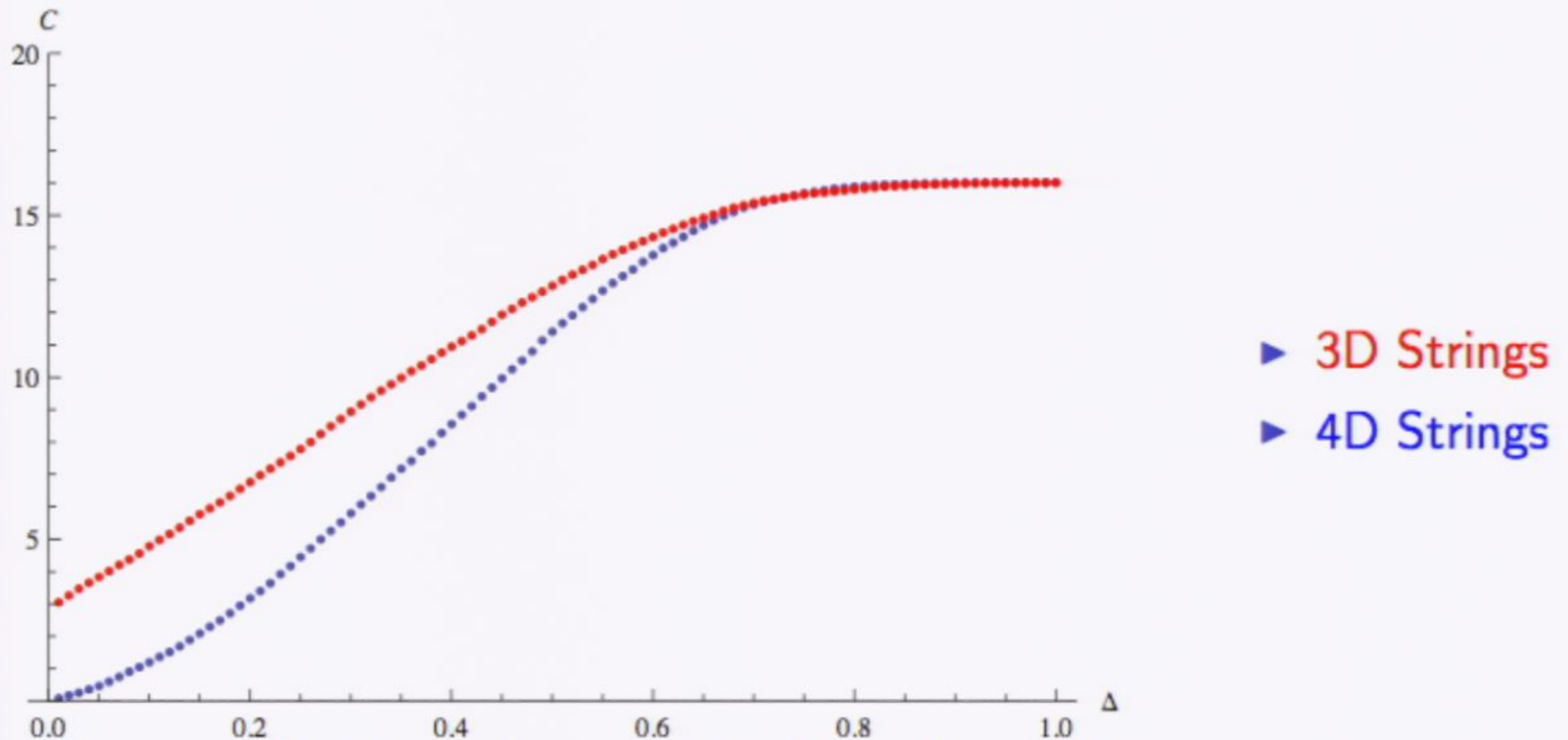
# Numerical Results - 3D

## Random Strings in 3D



Cusps are generic in 3D

# Numerical Results - Extra Dimensions



Same qualitative features, extra dimensions kill cusps.

# Outline

Introduction and Motivation

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**Conclusions**

## Conclusions

- ▶ In 3D, LIGO may have been too conservative with their search.
- ▶ But with extra dimensions, it may be much harder to detect a cosmic string than previously thought.

### Other parts of the story:

- ▶ Consider the evolution of cosmic strings in an expanding universe.
- ▶ Add a Randall-Sundrum type potential in the extra dimensions.