

Title: Quarks and Leptons as Nambu-Goldstone Fermions Under $E_7/SO(10)$

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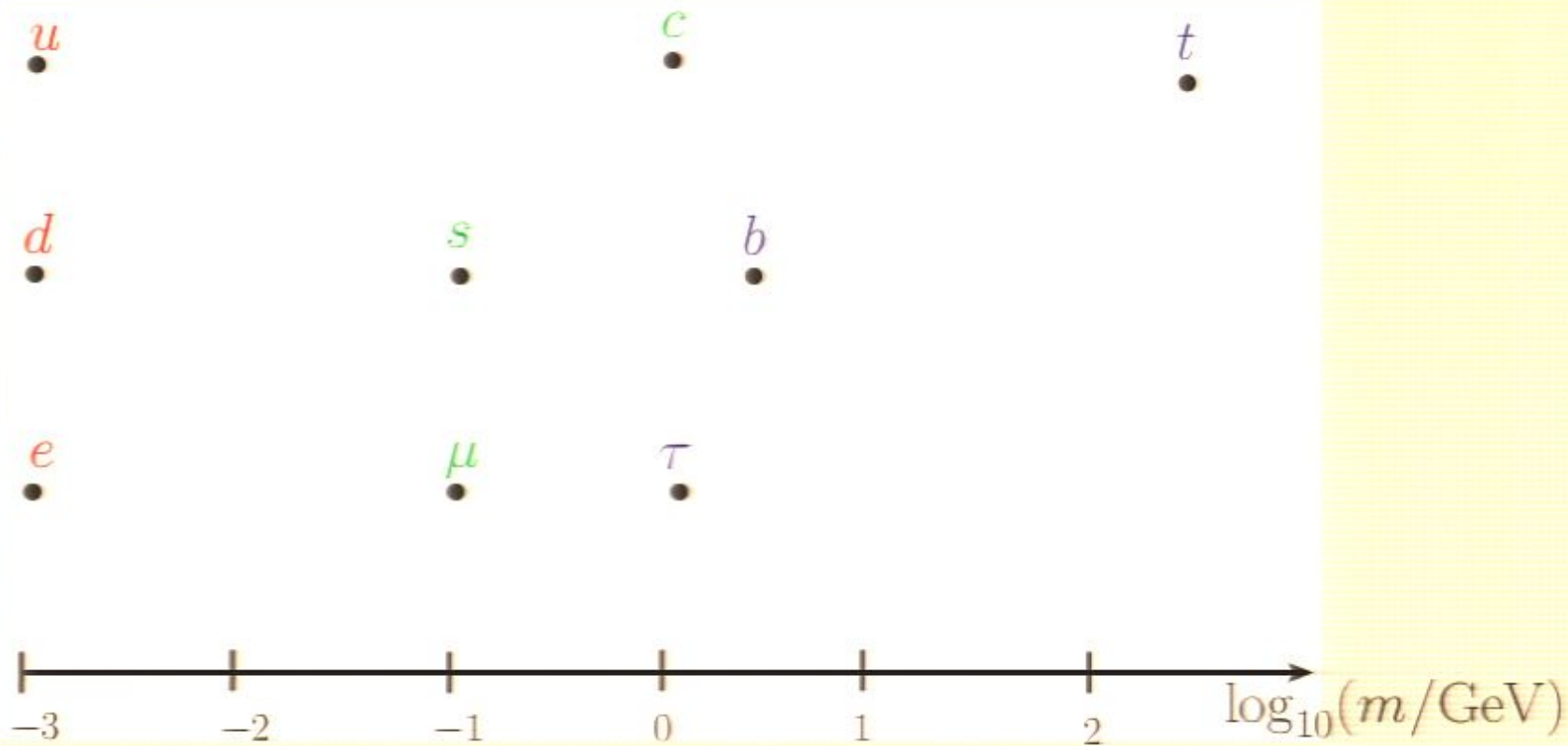
URL: <http://pirsa.org/11080069>

Abstract: The hierarchy of the Yukawa couplings is an outstanding problem of the standard model. We present a class of models in which the first and second generation fermions are SUSY partners of pseudo-Nambu-Goldstone bosons that parameterize an $E_7/SO(10)$ Kahler manifold, explaining the small values of these fermion masses relative to those of the third generation. We consider experimental constraints on this scenario, and find that the simplest model with universal gaugino masses is already ruled out by the LHC. However, models with non-universal gaugino masses will likely be excluded only by direct dark matter searches.

Hierarchy of Yukawa couplings?

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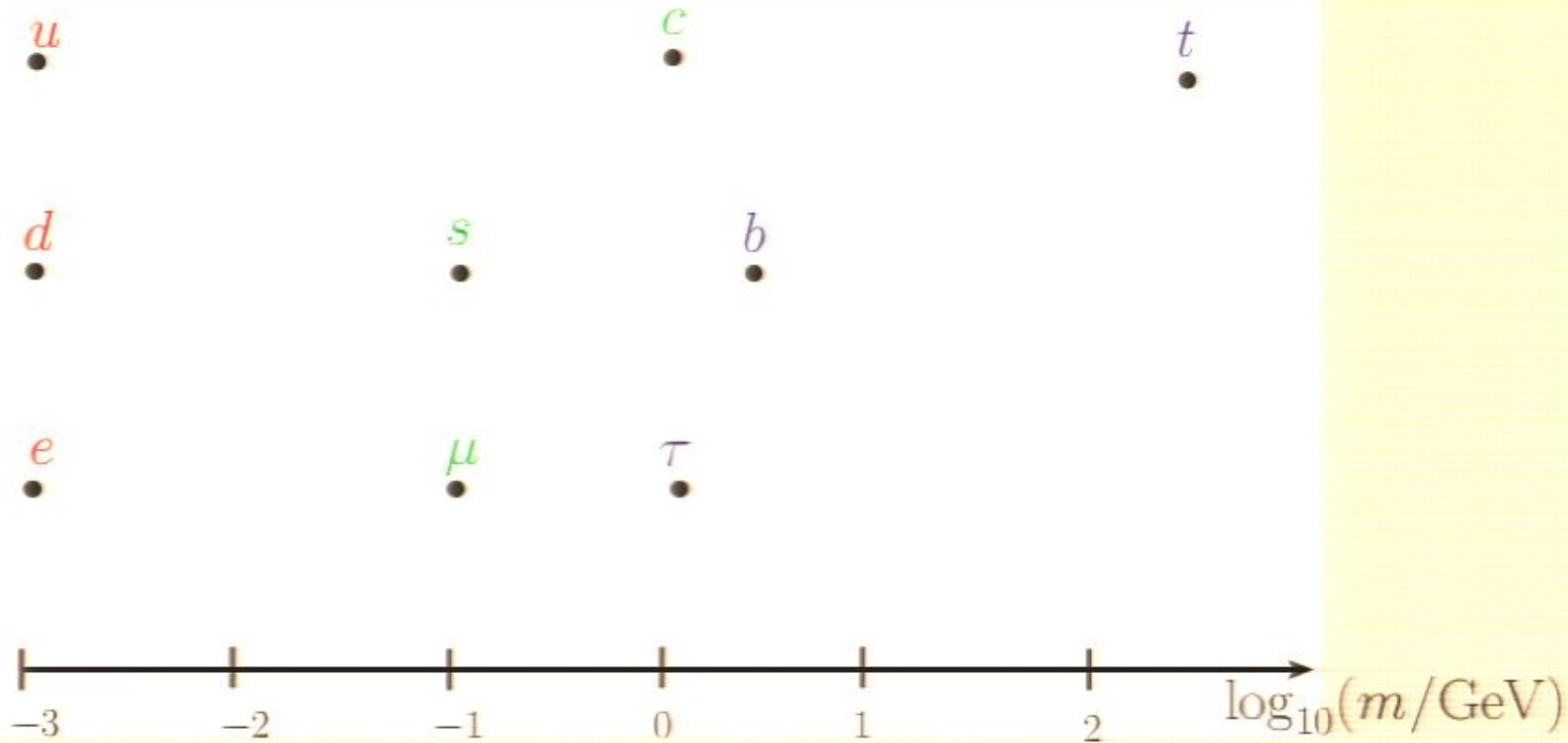
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Ad-hoc \Rightarrow dynamical?

Field content of standard model

Gauged adjoint and light chiral fundamentals

- Gauge some group $H \subset G$ in real representation
- Where do light chiral fields in fundamental of H come from?

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Hint from pions

- Take $SU(2)/U(1)$ toy model
- Broken X and $Y \Rightarrow Z = X + iY, Z^* = X - iY$
- $SU(2) \sim S^3 \implies S^3/U(1) \sim \mathbb{C}P^1$

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- Broken X and $Y \Rightarrow Z = X + iY, Z^* = X - iY$
- $SU(2) \sim S^3 \implies S^3/U(1) \sim \mathbb{C}P^1$
- Gauge $U(1)$, toy pions in fundamental complex representation
- Need fermionic analogue

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Other problems ...

Problems of the Standard Model

- Higgs hierarchy problem
- Dark matter
- Gauge unification
- ⋮

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Problems of supersymmetry

- μ problem

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Problems of supersymmetry

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- Flavor/CP problem $\Rightarrow m_{0(3)} \gg m_{0(1,2)} ?$
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Natural way to accomplish this? Follow hint ...

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How to construct one?

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Under global symmetry breaking $G \rightarrow H$,

$$\mathcal{L}_{NGB}^{NL} = - (\partial^\mu \pi_i) g^{ij} (\pi/f_\pi) (\partial_\mu \pi_j)$$

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Need:

$$2n_{NGF} = n_{NGB} + n_{QNGF}$$

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“Quasi”-NGBs n_{QNGB} minimal if G/H is Kähler!

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Incorporate in supergravity ...

Broken d.o.f.'s parameterize Kähler manifold

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Supersymmetric generalization of nonlinear σ model:

$$\mathcal{L}_{SUSY}^{NL} = K(\Phi, \Phi^\dagger) \Big|_{\theta\theta\bar{\theta}\bar{\theta}}$$

Then for scalars:

$$\mathcal{L}_{SUSY}^{NL} = -\partial^\mu \phi^{*i} \left[\frac{\partial^2 K(\phi^*, \phi)}{\partial \phi^{*i} \partial \phi^j} \right] \partial_\mu \phi^j$$

such that

$$\frac{\partial g_{ij}}{\partial \phi^{*k}} = \frac{\partial g_{kj}}{\partial \phi^{*i}}, \quad \frac{\partial g_{ij}}{\partial \phi^k} = \frac{\partial g_{ik}}{\partial \phi^j}$$

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Technical points

Invariance of K

- Kähler transformation:

$$K(\Phi, \Phi^\dagger) \rightarrow K(\Phi, \Phi^\dagger) + F(\Phi) + F^\dagger(\Phi^\dagger)$$

- Under supergravity,

$$L = \left[\Sigma \Sigma^\dagger e^{-K(\Phi, \Phi^\dagger)} \right]_D + \left[\Sigma^3 W(\Phi) \right]_F$$

- For finite W , $\Sigma \rightarrow \Sigma e^F$ not sufficient!

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Solution criteria

- G/H must be non-compact [Bagger and Witten, Phys. Lett. B **118**, 103 (1982)]
- H must have no $U(1)$ factors [Kugo and Yanagida, Prog. Theor. Phys. **124**, 555 (2010)]

Technical points (cont'd)

- General form of Kähler potential:

$$K = v^2 F(\det[\Psi_\alpha^\dagger \Psi^\alpha])$$

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- v is scale of G breaking, $F(x)$ freedom from broken $U(1)$'s

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- For a simple example manifold similar to Pati-Salam, see:

SKM, Nojiri, Sudano, Yanagida, JHEP **01**, 131 (2011)

Supergravity and SUSY breaking

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- Couple to supergravity:

$$\Delta\mathcal{L} = \frac{3}{4} \int d^2\theta d^2\bar{\theta} \mathcal{E} (\bar{D}\bar{D} - 8R) \exp\left(\frac{1}{3}K(\Psi, \Psi^\dagger)\right)$$

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- Introduce hidden SUSY breaking sector:

$$K = K(\Psi, \Psi^\dagger) + Z^\dagger Z + \dots$$

- NG bosons remain massless, but QNGB d.o.f.'s from broken $U(1)$'s inside Ψ get mass \Rightarrow “novino” superfields

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Choosing $G = E_7$

Formal motivations for E_7

- Appears in $N = 8$ supergravity
- Nonlinear realization gives the “simplest field theory”
[Arkani-Hamed, Cachazo, Kaplan arXiv:0808.1446]

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GUT-like representations

- $SU(5) \supset (5 + 5^*)_H \oplus 10 \oplus 5^*$
- $SO(10) \supset 10_H \oplus 16 \rightarrow (5 + 5^*)_H \oplus 10 \oplus 5^* \oplus 1$

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 - $E_7/[SU(5) \times U(1)^3] \rightarrow (10 \oplus 5^*) \times 3 + 5_H$
 - $E_7/[SO(10) \times U(1)^2] \rightarrow (10 \oplus 5^*) \times 2 + (5 + 5^*)_H$

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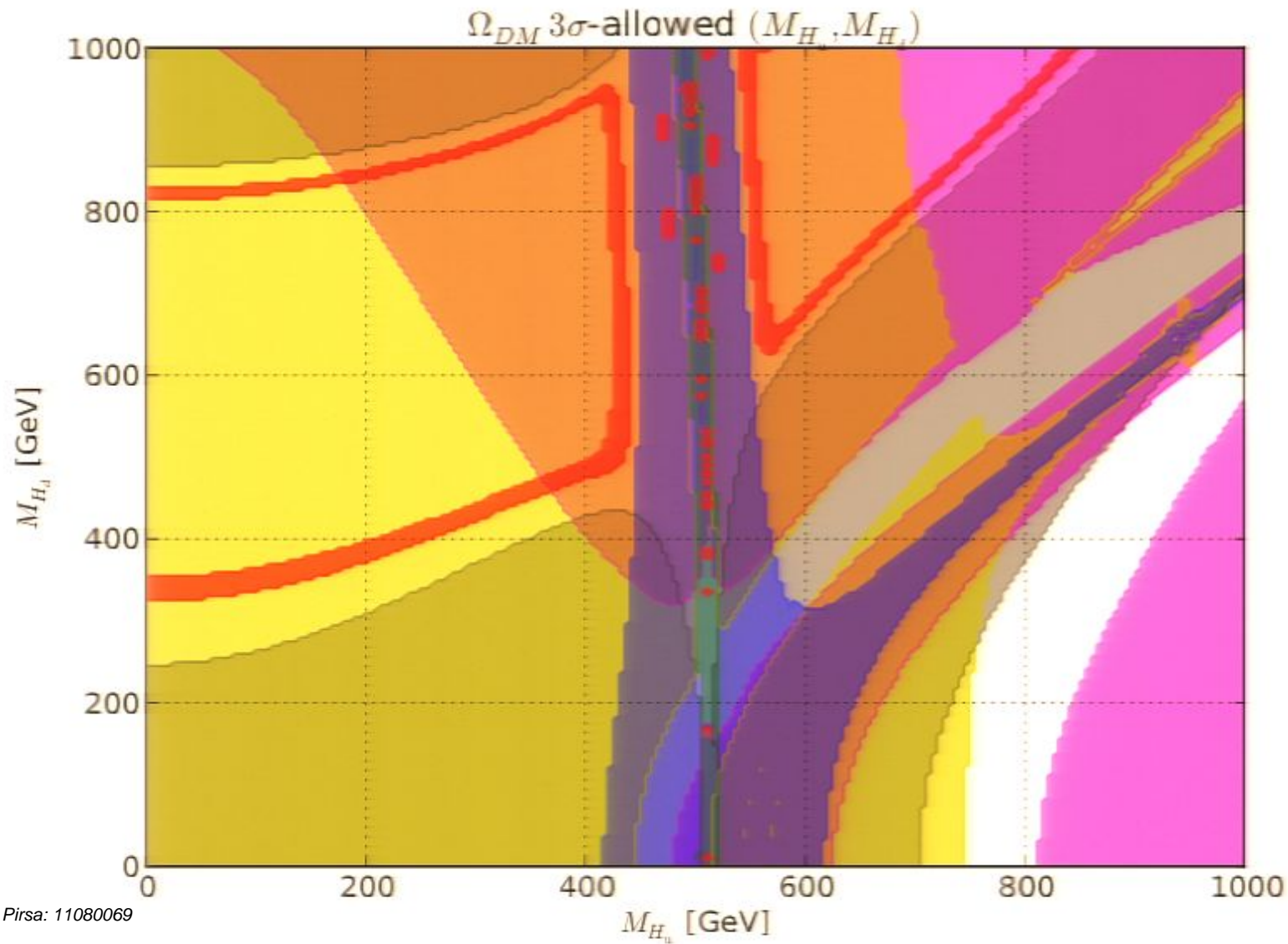
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- E_6 too small, E_8 gives mirror families

Light 5*, Light 10 (3rd generation)

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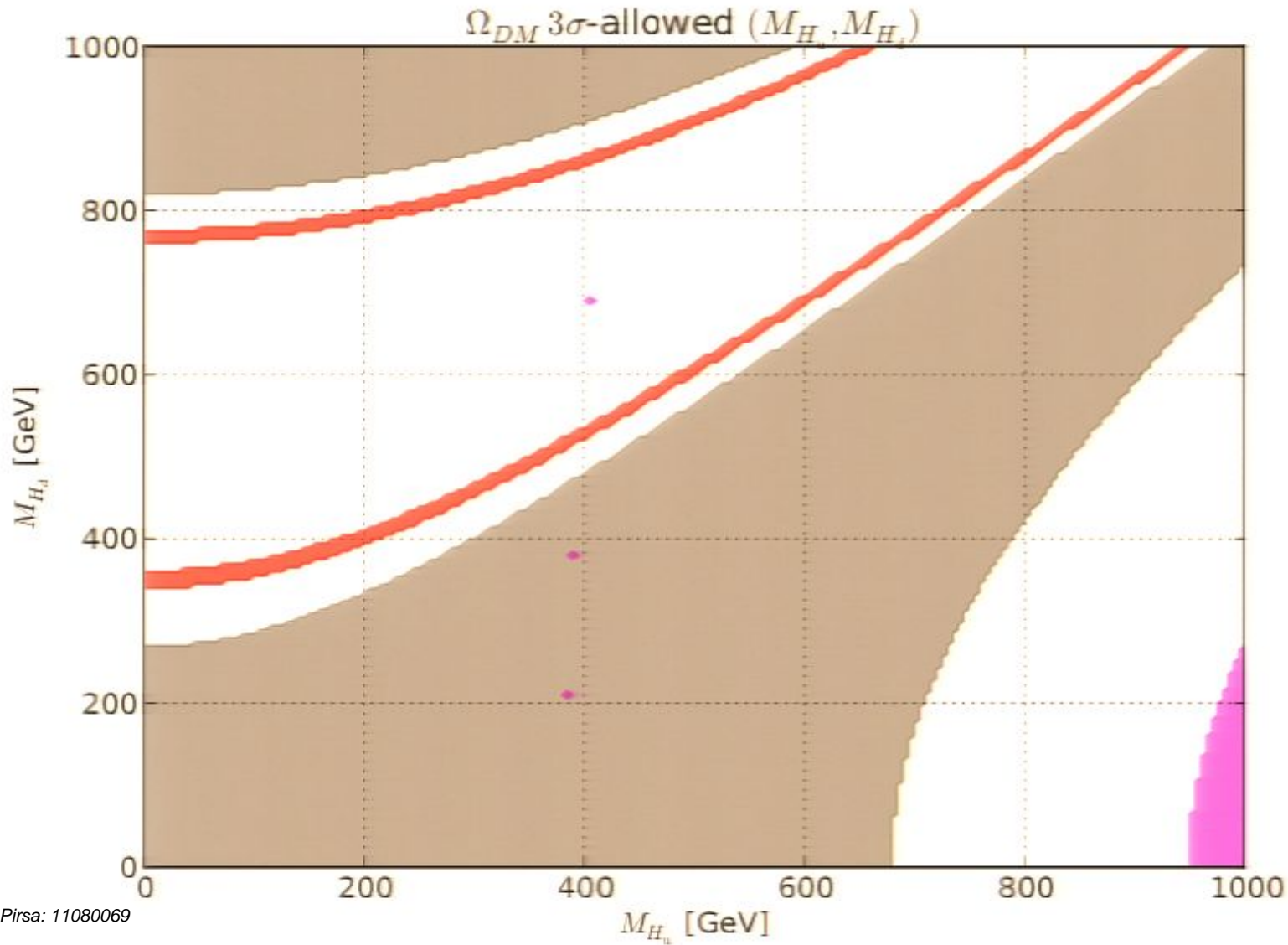
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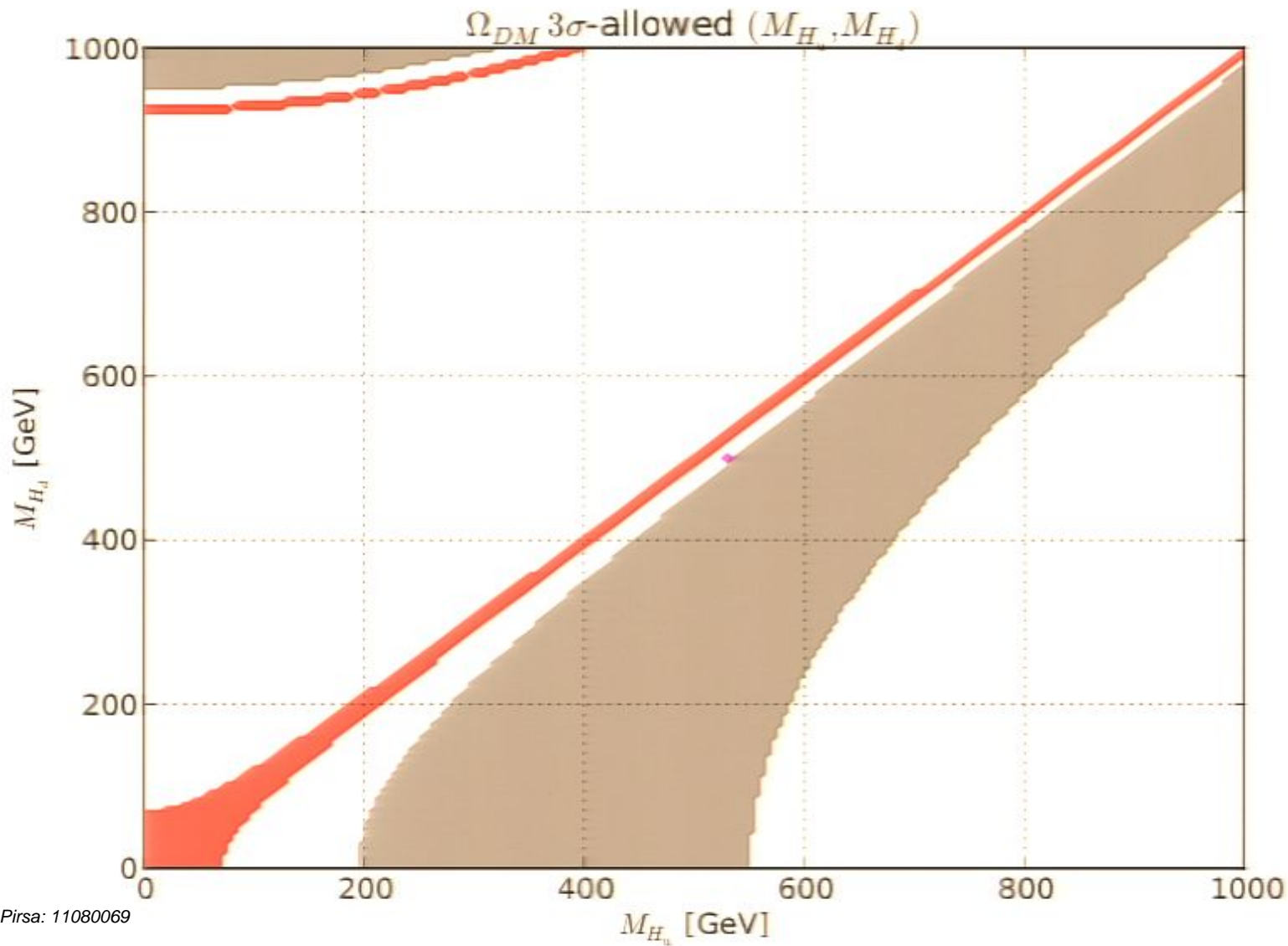
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heavy 5^* , Heavy 10 ($M_{input} = M_P$)

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$SU(5)$ of $E_7/SO(10)$

$$SO(10) \rightarrow SU(5)$$

- Global symmetry must be gauged anyway \Rightarrow gauge $SU(5) \subset SO(10)$
- *Explicitly* break $U(1)^2$ for coupling to supergravity, gives two gravitinos $N_1, N_2 \sim M_P$

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- Explicitly break $U(1)^2$ for coupling to supergravity, gives two gravitinos $N_1, N_2 \sim M_P$
- Superpotential:

$$W = W_Y + W_S + W_H$$

where

$$W_Y = Y_u \cdot 10 \cdot 10 \cdot 5_H + Y_d \cdot 10 \cdot 5^* \cdot 5_H^*$$

$$W_S = M_\nu \cdot 1 \cdot 1 + M_N \cdot N \cdot N$$

$$W_H = \mu \cdot 5_H \cdot 5_H^*$$

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GUT superpotential

Additions to superpotential

$$\Delta W = W_\Sigma + W_{H'}$$

where

$$W_\Sigma = M_\Sigma \text{Tr} \Sigma^2 + \lambda \text{Tr} \Sigma^3$$

$$W_{H'} = \lambda_1 \cdot 5_H \cdot \Sigma \cdot 5_H^*$$

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Outstanding issues

- Set right boundary conditions for Ω_{DM}
- Evade LHC bounds?

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Making $M_{GUT} \approx M_P$

$$\Sigma \sim 24 \supset (8, 1) \oplus (1, 3) \oplus (1, 1) \oplus (3, 2) \oplus (3^*, 2)$$

- Since gauging anyway, gauge un-Higgsed part of 24

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- R -charge of $\text{Tr } \Sigma^3 \neq 2 \implies \lambda \ll 1$ natural $\implies m_{3,8} \ll M_{GUT}$

Choose: $\sim 10^{12}$ GeV

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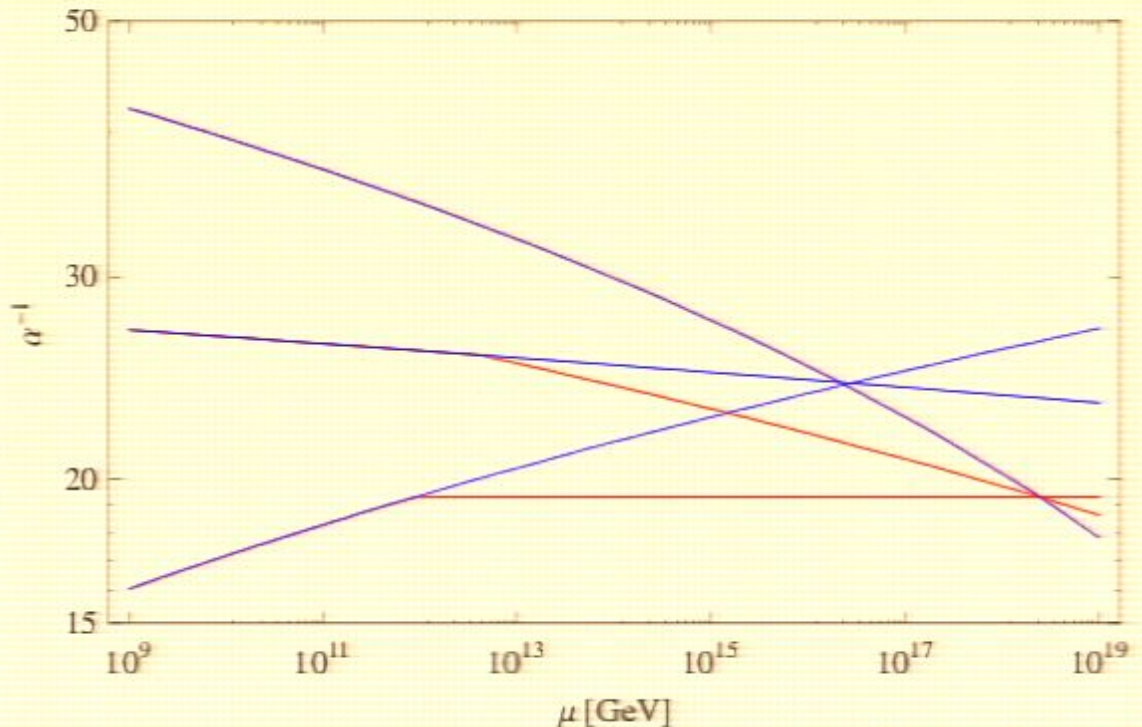
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Solving the double-triplet splitting problem

Missing partner mechanism

- 50 of $SU(5)$ contains color (3, 1), not weak (1, 2)
- Explicit triplet Higgs $5_H + 5_H^*$ to $50 + 50^*$ coupling

$$W_T = \lambda_T \cdot 5_H \cdot \langle 75 \rangle \cdot 50 + \text{c.c.}$$

- Huge threshold corrections?

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Product group unification

- $SU(5)_G \rightarrow SU(5)_G \times SU(3)_H \times U(1)_H$
- Diagonal

$$SU(3)_c \subset SU(3)_G \times SU(3)_H$$

$$U(1)_Y \subset U(1)_G \times U(1)_H$$

Product group unification (1/2)

Breaking

$$\bar{Q}(m + \lambda\Sigma)Q + \frac{1}{2}m_\Sigma\text{Tr}(\Sigma^2) + hHQ\bar{q} + c.c.$$

where

$$Q \sim (5^*, 3, 1), \quad q \sim (1, 3, 1)$$

and

$$\langle Q \rangle \sim vI_3, \quad \langle \Sigma \rangle \sim \text{diag}(3, 3, -2, -2, -2)$$

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Gaugino mass splitting

Modified soft breaking

$$\begin{aligned} \Delta\mathcal{L}_{SSB} = & - \frac{1}{2}M_G\lambda_G\lambda_G - \frac{1}{2}M_{H3}\lambda_{H3}\lambda_{H3} \\ & - \frac{1}{2}M_{H1}\lambda_{H1}\lambda_{H1} + c.c. \end{aligned}$$

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Product gauge unification (2/2)

Gaugino mass splitting

- Re-weighted SM gaugino soft masses:

$$M_3 \implies g_3^2 \left(\frac{M_{H3}}{g_{H3}^2} + \frac{M_G}{g_G^2} \right)$$

$$M_1 \implies g_1^2 \left(\frac{M_{H1}}{15g_{H1}^2} + \frac{M_G}{g_G^2} \right)$$

- Changes in low-energy phenomenology:

$$M_3 = M_{1/2} \implies M_3 \sim 2 \times M_{1/2}$$

$$M_1 \sim M_{1/2}$$

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MSSM analysis

Soft masses

- $m_{3/2} = m_{0(3)} = 1 \text{ TeV}$
- $m_{0(1,2)} = 0$
- $A_0 = 0, m_{H_u} = m_{H_d} = 0$
- $M_{input} = M_P$

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Uncertainties in $\tan \beta$

- b term depends on details of Kähler potential
- Spectrum unaffected, except for small shift in \tilde{t}, \tilde{b} masses
- Direct detection cross section varies from $4 \times 10^{-47} \text{ cm}^2$ ($\tan \beta = 10$) to $4 \times 10^{-46} \text{ cm}^2$ ($\tan \beta = 50$)

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Spectra

Choose SPS1a for comparison since slepton/light gaugino spectra similar:

Field	SPS1a	$M_3 = M_{1/2}$	$M_3 = 500 \text{ GeV}$
$\tilde{\chi}_1^0$	100	100	100
$\tilde{\chi}_2^0, \tilde{\chi}_1^\pm$	200	200	200
$\tilde{l}_{R(1,2)}$	140	140	120
$\tilde{l}_{L(1,2)}$	200	200	200
H_0, A_0, H^\pm	400	1000	1000
$\tilde{\chi}_2^\pm, \tilde{\chi}_{3,4}^0$	380	1000	1000
$\tilde{\tau}, \tilde{\nu}_\tau$	200	1000	1000
\tilde{b}_2, \tilde{t}_2	550	1100	1450
\tilde{g}	600	830	1300
$\tilde{q}_{1,2}$	550	720	1150

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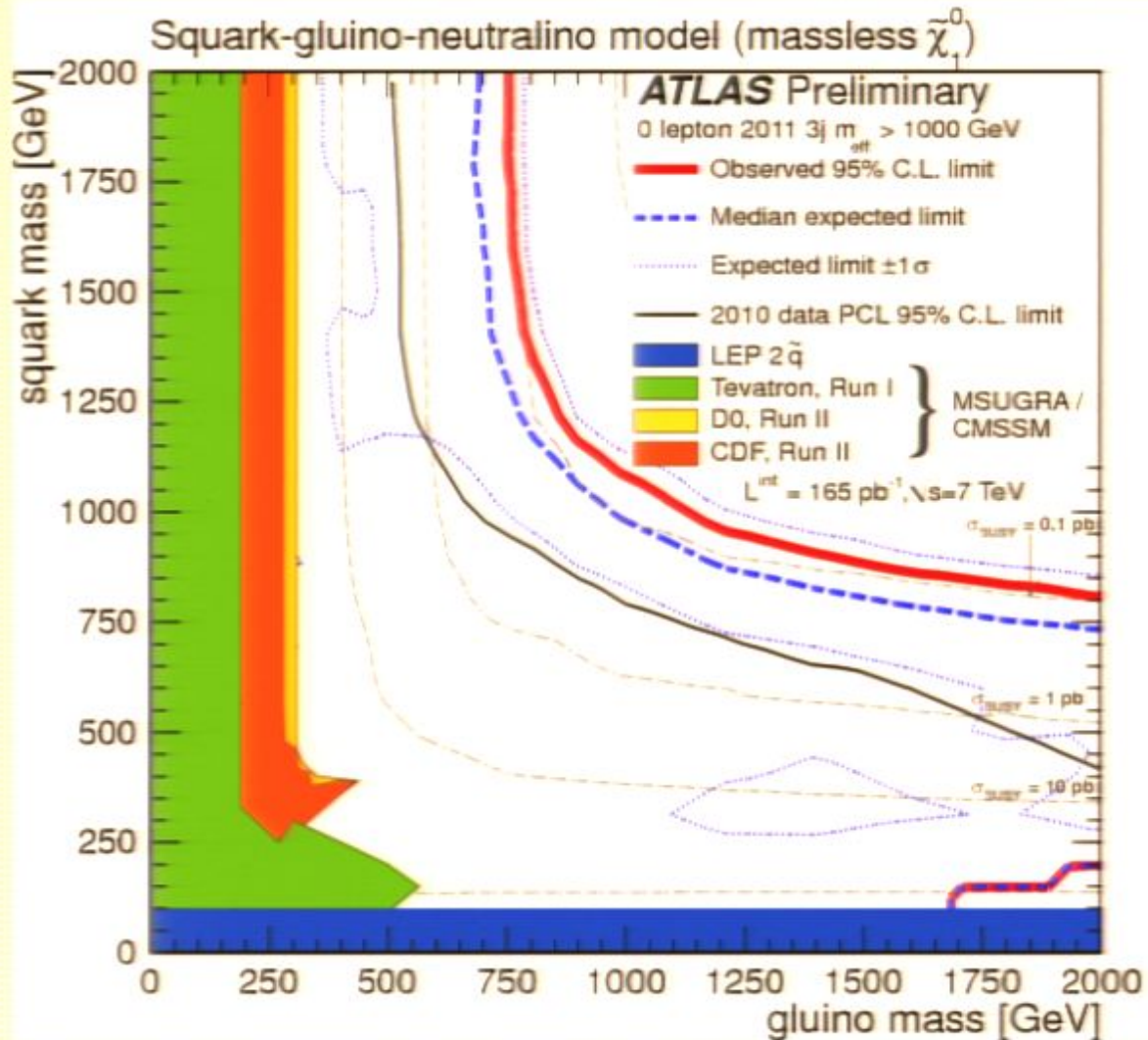
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Newest model-independent constraints



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Constraint on model points

Jets + missing E_T

≥ 3 -jet signal region only 10 ± 2 SM events at 165pb^{-1} :

	SPS1a	$M_3 = M_{1/2}$	$M_3 = 500$ GeV
Prod. σ	4.5 pb	0.6 pb	16 fb
Efficiency	12%	28%	29%
$\mathcal{L}(2\sigma)$	7pb^{-1}	15pb^{-1}	10fb^{-1}

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Constraint on model points

Jets + missing E_T

≥ 3 -jet signal region only 10 ± 2 SM events at 165pb^{-1} :

	SPS1a	$M_3 = M_{1/2}$	$M_3 = 500$ GeV
Prod. σ	4.5 pb	0.6 pb	16 fb
Efficiency	12%	28%	29%
$\mathcal{L}(2\sigma)$	7pb^{-1}	15pb^{-1}	10fb^{-1}

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Multiplicities

Mistag rates: τ -jet (8%), b -jet (5%)

	SPS1a	$M_3 = M_{1/2}$	$M_3 = 500$ GeV
b -jets	18%	5%	5%
τ -jets	8%	9%	9%
$n_l \geq 2$	8%	14%	14%

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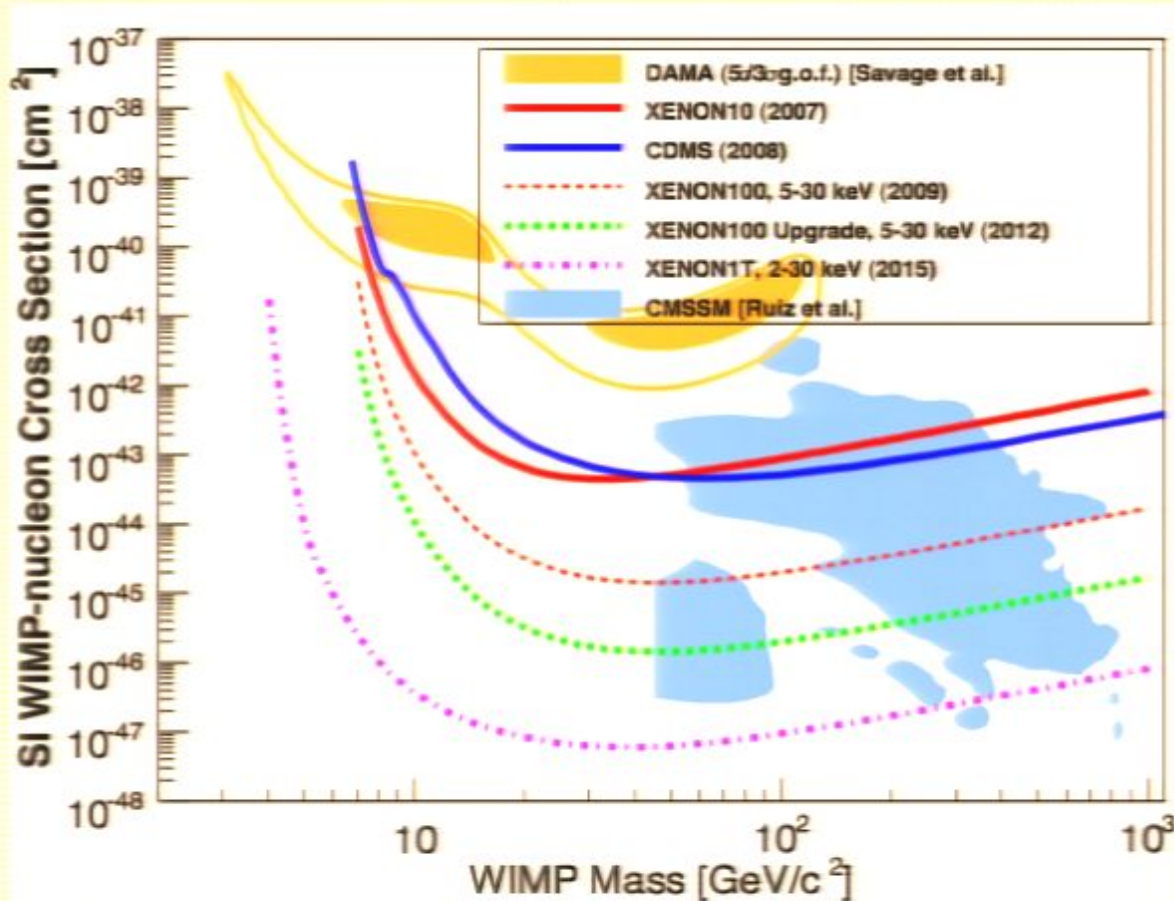
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Dark matter phenomenology

Detection

- $\sigma_{\chi N}$ excludable by XENON1T
- χ_1^0 mostly bino, $\langle \sigma v \rangle_{\chi\chi}$ is p -wave suppressed



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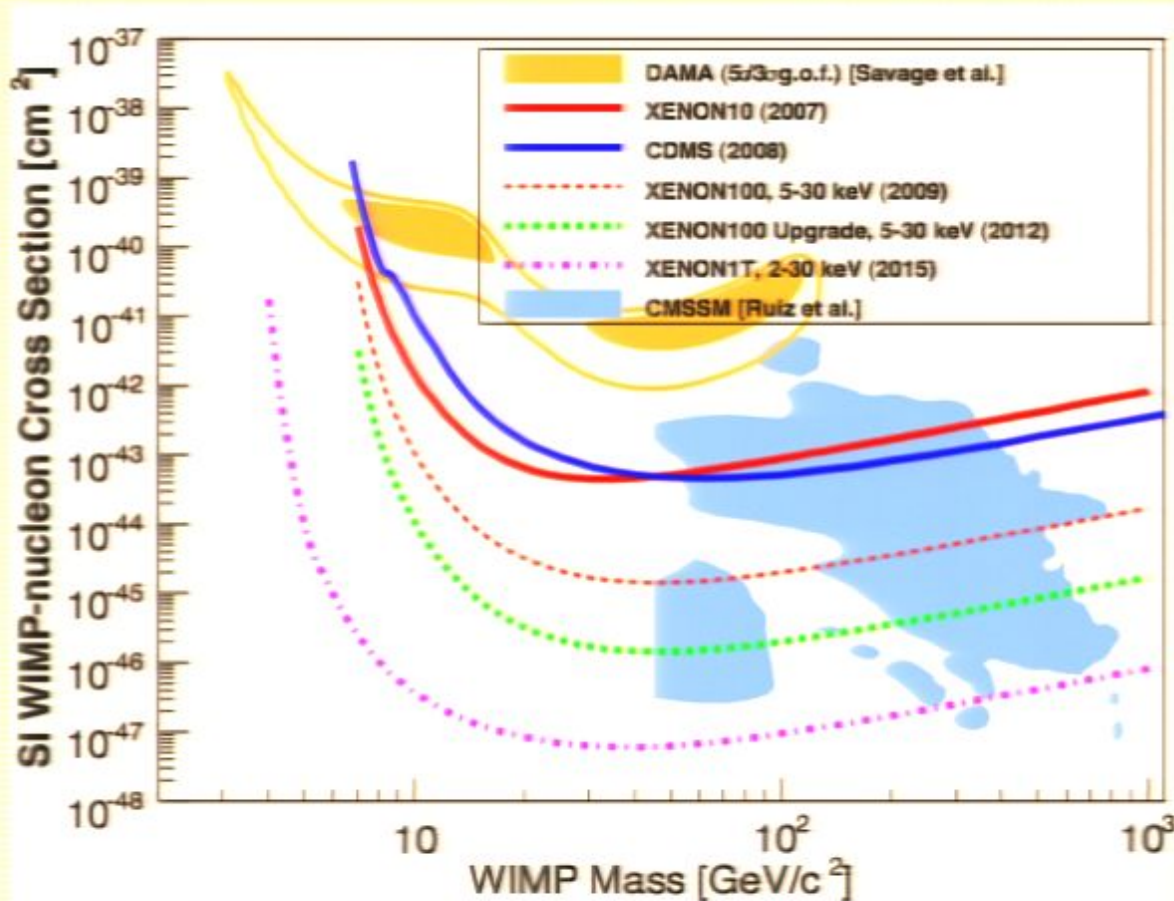
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CP anomalies/flavor constraints

$B_s - \bar{B}_s \rightarrow \mu^+ \mu^+ (\mu^- \mu^-)$ anomaly

- 3.2σ deviation from SM prediction
- Oscillations $M(B_q) - M(\bar{B}_q)$ constrained to be small
- But $b \rightarrow s\gamma$, EDMs close to SM prediction

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Analysis for NG hypothesis

- Endo, Shirai, Yanagida [arXiv:1009.3366]
- Large CP angle $\arg(U_{23})$, $|U_{23}| \sim \lambda^2$ can explain dimuon anomaly
- Heavy 3rd generation suppresses oscillations, $b \rightarrow s\gamma$

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- Heavy 3rd generation suppresses oscillations, $b \rightarrow s\gamma$
- EDMs suppressed by heavy 3rd generation plus hermiticity
- Predicted lepton flavor violations in reach of future experiments

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Created consistent $E_7/SO(10)$ coset model with distinctive collider signatures

Evading LHC constraints weakens flavor predictions

Can be ruled out direct detection in 2–5 years

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