

Title: Solving AdS/CFT Y-System

Date: Aug 10, 2011 10:00 AM

URL: <http://pirsa.org/11080029>

Abstract:

Solving AdS/CFT Y-system

Sébastien Leurent
LPT-ENS (Paris)

[arXiv:1108.soon] N. Gromov, V. Kazakov, SL & D. Volin


[arXiv:1007.1770] V. Kazakov & SL

[arXiv:1010.2720] N. Gromov, V. Kazakov, SL & Z. Tsuboi

[arXiv:1010.4022] V. Kazakov, SL & Z. Tsuboi

Outline



- 1 Motivation : Understanding Y-system better
- 2 Solving Hirota \Leftrightarrow Q-functions
- 3 Understanding symmetries from strong coupling 

Solving
AdS/CFT
Y-system

Leurent

tivation

unctions

nmetries

Thermodynamic Bethe Ansatz



short operators
infinite time periodicity
Path integral $Z \sim e^{-RE_0(L)}$



Long operators
finite time-periodicity
 \Rightarrow finite temperature

“free Energy” : $f(L) = E_0(L)$

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + (\text{Source Terms})$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}



- Extra assumption : Excited states obey the same equations.

Each state correspond to a different solution of Y-system.
 characterized by its zeroes and poles

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + \text{Source Terms}$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}

- Extra assumption : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + \text{Source Terms}$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}

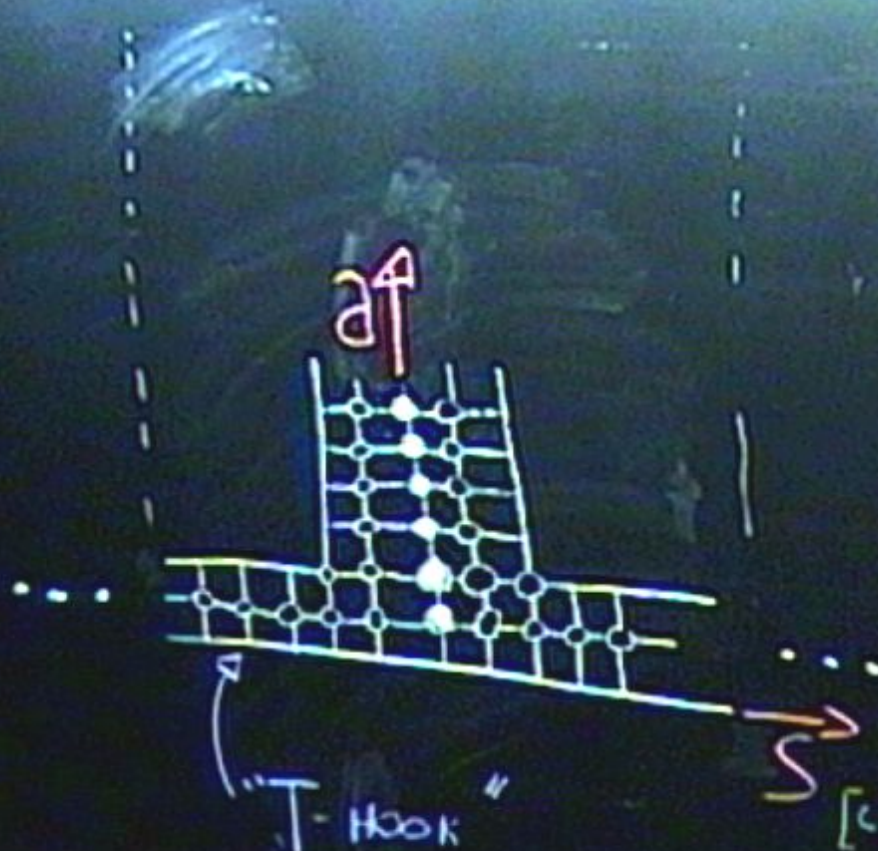


- Extra assumption : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles.

SPECTRAL PROBLEM

[CANNY 03]
[AF 03]
[AF 03]



[CAMBIA FIORANTI TATRO 03]

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + (\text{Source Terms})$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}



- Extra assumption : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles

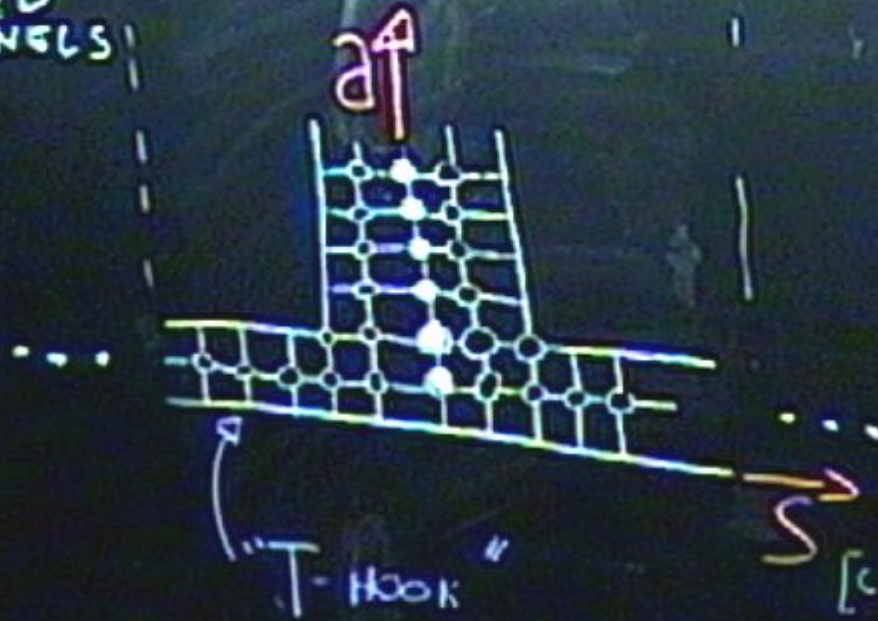
SPECTRAL PROBLEM

* TBA EQUATIONS

COMPLICATED KERNELS

∞ set of NLIES

[CAF 03]
[CAF 03]
[CAF 03]



[CANTALIA FIORAVANTI TATEO 03]

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) \\ + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + \langle \text{Source Terms} \rangle$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}

- **Extra assumption** : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) \\ + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + \langle \text{Source Terms} \rangle$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$

- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}

- **Extra assumption** : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles



Y-system and Hirota equation

Y-system Equation

The TBA integral equation imply the 'local' relation

$$Y_{a,s}^+ Y_{a,s}^- = \frac{1+Y_{a,s+1}}{1+(Y_{a+1,s})^{-1}} \frac{1+Y_{a,s-1}}{1+(Y_{a-1,s})^{-1}}$$

[Gromov Kazakov Vieira 09]

where $Y_{a,s}^\pm = Y_{a,s}(u \pm \frac{i}{2})$

- change of variable $Y_{a,s} = \frac{T_{a+1,s} T_{a,s-1}}{T_{a-1,s} T_{a,s}}$

Hirota equation

$$T_{a,s}^- T_{a,s}^+ = T_{a+1,s} T_{a-1,s} + T_{a,s-1} T_{a,s+1}$$



Gauge freedom

Y-functions and Hirota equation are invariant under gauge

transformations $T_{a,s} \rightarrow \frac{g_1^{a-s}}{g_2^{a-s}} \frac{g_3^{a-s-1}}{g_4^{a-s-1}} T_{a,s}$

$f(x) \equiv f(u = x/2)$

Y-system and Hirota equation

Y-system Equation

The TBA integral equation imply the 'local' relation

$$Y_{a,s}^+ Y_{a,s}^- = \frac{1+Y_{a,s+1}}{1+(Y_{a+1,s})^{-1}} \frac{1+Y_{a,s-1}}{1+(Y_{a-1,s})^{-1}}$$

[Gromov Kazakov Vieira 09]

where $Y_{a,s}^\pm = Y_{a,s}(u \pm \frac{i}{2})$

• change of variable $Y_{a,s} = \frac{T_{a+1,s} T_{a,s-1}}{T_{a-1,s} T_{a,s}}$

Hirota equation

$$T_{a,s}^- T_{a,s}^+ = T_{a+1,s} T_{a-1,s} - T_{a,s-1} T_{a,s+1}$$

Gauge freedom

Y-functions and Hirota equation are invariant under gauge

transformations $T_{a,s} \rightarrow g_1^{[a-s]} g_2^{[a-s]} g_3^{[-a+s]} g_4^{[-a+s]} T_{a,s}$

$f(\lambda) \equiv f(u \pm \frac{i}{2})$

SPECTRAL PROBLEM

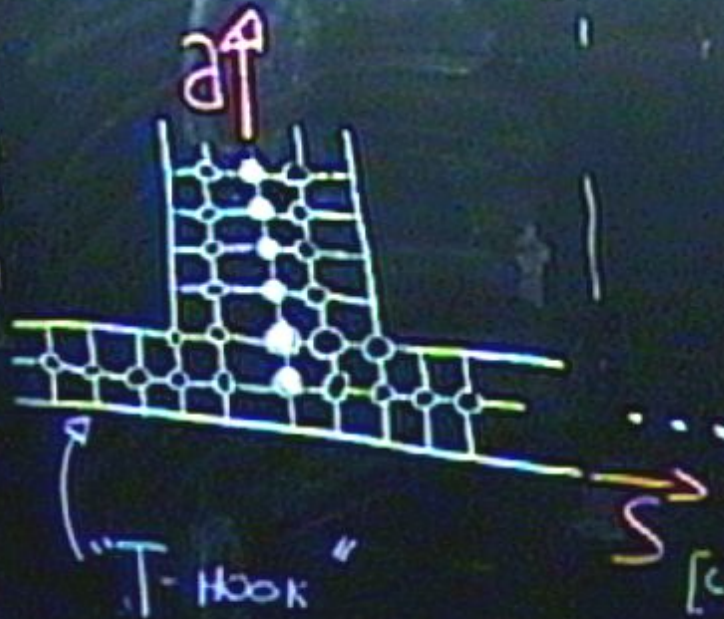
* TBA EQUATIONS

COMPLICATED KERNELS

∞ set of NLIES

[4444 03]
[444 03]
[44 03]

$$y^+ y^- = \frac{(1+y)(1+y)}{(1+\frac{1}{y})(1+\frac{1}{y})}$$



[CAMGLIA FIORAVANTI TATEO 03]

Y-system and Hirota equation

Y-system Equation

The TBA integral equation imply the 'local' relation

$$Y_{a,s}^+ Y_{a,s}^- = \frac{1+Y_{a,s+1}}{1+(Y_{a+1,s})^{-1}} \frac{1+Y_{a,s-1}}{1+(Y_{a-1,s})^{-1}}$$

[Gromov Kazakov Vieira 09]

where $Y_{a,s}^\pm = Y_{a,s}(u \pm \frac{i}{2})$

- change of variable $Y_{a,s} = \frac{T_{a+1,s} T_{a,s-1}}{T_{a-1,s} T_{a,s}}$

Hirota equation

$$T_{a,s}^- T_{a,s}^+ = T_{a+1,s} T_{a-1,s} - T_{a,s+1} T_{a,s-1}$$

Gauge freedom

Y-functions and Hirota equation are invariant under gauge

transformations $T_{a,s} \rightarrow g_1^{[a-s]} g_2^{[a-s]} g_3^{[-a-s]} g_4^{[-a-s]} T_{a,s}$

$f(\pm k) \equiv f(u \pm \frac{i}{2})$

Y-system and Hirota equation

Y-system Equation

The TBA integral equation imply the 'local' relation

$$Y_{a,s}^+ Y_{a,s}^- = \frac{1+Y_{a,s+1}}{1+(Y_{a+1,s})^{-1}} \frac{1+Y_{a,s-1}}{1+(Y_{a-1,s})^{-1}}$$

[Gromov Kazakov Vieira 09]

where $Y_{a,s}^\pm = Y_{a,s}(u \pm \frac{i}{2})$

- change of variable $Y_{a,s} = \frac{T_{a+1,s} T_{a,s-1}}{T_{a,s} T_{a+1,s-1}}$

Hirota equation

$$T_{a,s}^- T_{a,s}^+ = T_{a+1,s} T_{a-1,s} - T_{a,s+1} T_{a,s-1}$$

Gauge freedom

Y-functions and Hirota equation are invariant under gauge

transformations $T_{a,s} \rightarrow g_1^{|a-s|} g_2^{|a-s|} g_3^{|-a-s|} g_4^{|-a-s|} T_{a,s}$

$f(\pm \frac{i}{2}) \equiv f(u \pm \frac{i}{2})$

Y-system and Hirota equation

Y-system Equation

The TBA integral equation imply the 'local' relation

$$Y_{a,s}^+ Y_{a,s}^- = \frac{1+Y_{a,s+1}}{1+(Y_{a+1,s})^{-1}} \frac{1+Y_{a,s-1}}{1+(Y_{a-1,s})^{-1}}$$

[Gromov Kazakov Vieira 09]

where $Y_{a,s}^\pm = Y_{a,s}(u \pm \frac{i}{2})$

- change of variable $Y_{a,s} = \frac{T_{a,s+1} T_{a,s-1}}{T_{a+1,s} T_{a-1,s}}$

Hirota equation

$$T_{a,s}^+ T_{a,s}^- = T_{a+1,s} T_{a-1,s} + T_{a,s+1} T_{a,s-1}$$

Gauge freedom

Y-functions and Hirota equation are invariant under gauge

transformations $T_{a,s} \rightarrow g_1^{[a-s]} g_2^{[a-s]} g_3^{[-a-s]} g_4^{[-a-s]} T_{a,s}$

$f(x) \equiv f(u \pm \frac{i}{2})$

SPECTRAL PROBLEM

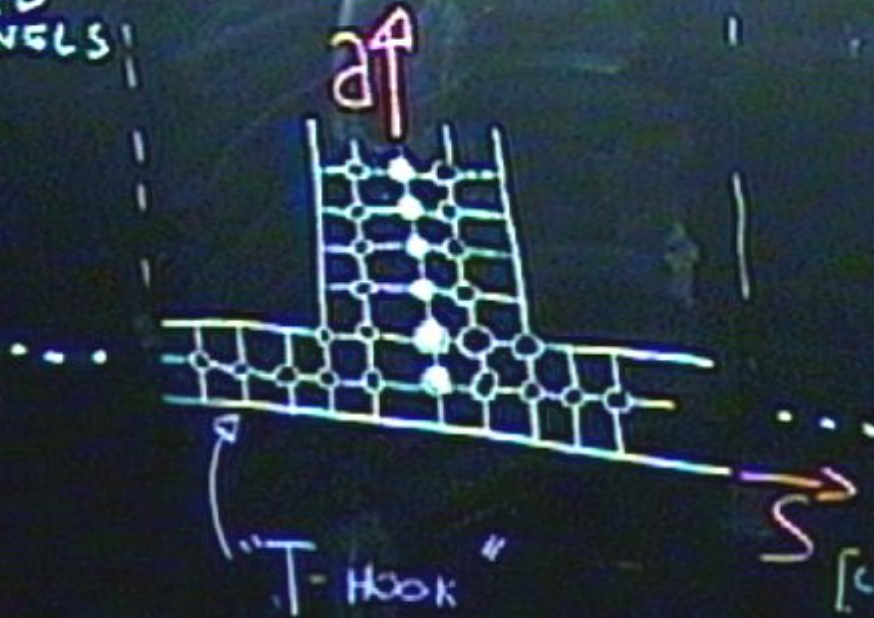
* TBA EQUATIONS

$$\Rightarrow y^+ y^- = \frac{(1+y^+)(1+y^-)}{\left(\frac{1+y^+}{\gamma}\right)\left(\frac{1+y^-}{\gamma}\right)} \Leftrightarrow \text{HAROTA}$$

COMPLICATED KERNELS

∞ set of NLIES

[LNV 03]
[LFF 03]
[LAF 03]



[CINGLIA FIORAVANTI TATEO 03]

SPECTRAL PROBLEM

* TBA EQUATIONS

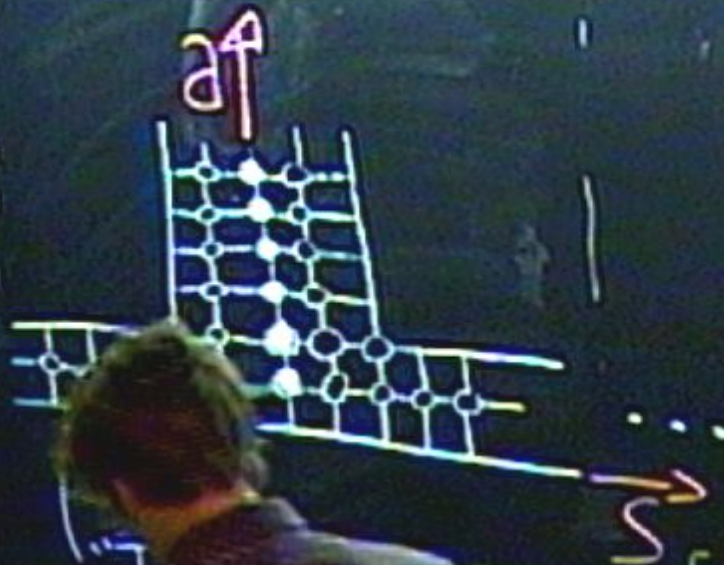
$$y^+ y^- = \frac{(1+y)(1+y)}{(1+\frac{1}{y})(1+\frac{1}{y})}$$

↔ HARDY/GAUGE

COMPLICATED KERNELS

∞ set of NLIES

[EAV 03]
[EAF 03]
[EAF 03]



[CANGIOLA FIORAVANTI TATEO 03]



SPECTRAL PROBLEM

* TBA EQUATIONS

$$y^+ y^- = \frac{(1+y)(1+y)}{(1+\frac{1}{y})(1+\frac{1}{y})}$$

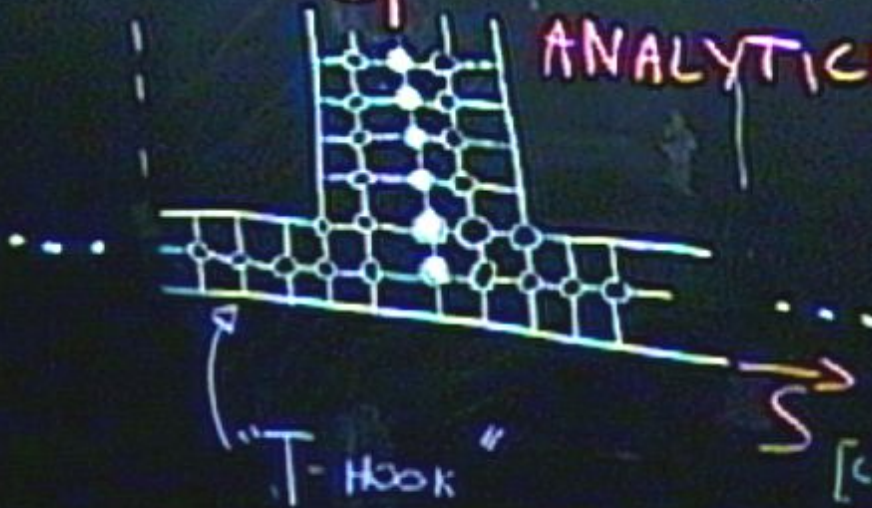
HAROTA/GAUGE

COMPLICATED KERNELS

∞ set of NLIES

EXTRA INFORMATION

ANALYTICITY



[LAP 03]
[LAP 03]
[LAP 03]

[CIRCUITI FIDELI TATEO 03]

SPECTRAL PROBLEM

* TBA EQUATIONS

COMPLICATED KERNELS

∞ set of NLIES

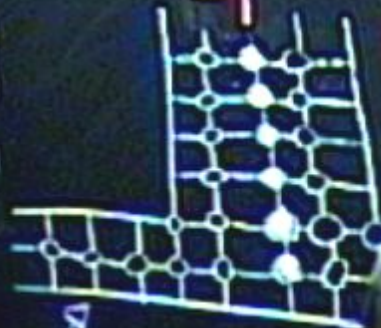
[EUV 03]
[EFT 03]
[AF 03]

$$y^+ y^- = \frac{(1+y|)(1+y)}{(1+\frac{1}{y})(1+\frac{1}{y})}$$

HAROTA/GAUGE

ANALYTICITY

EXTRA INFORMATION



ANALYTICITY

FIORAVANTI TATEO 03)

SPECTRAL PROBLEM

FINITE SET OF EQUATIONS

* TBA EQUATIONS

$$y^+ y^- = \frac{(1+y|/|y)}{(1+y|/|y)}$$

HAROTA/GAUGE

COMPLICATED KERNELS

∞ set of NLIES

ANALYTICITY

↑ EXTRA INFORMATION

ANALYTICITY



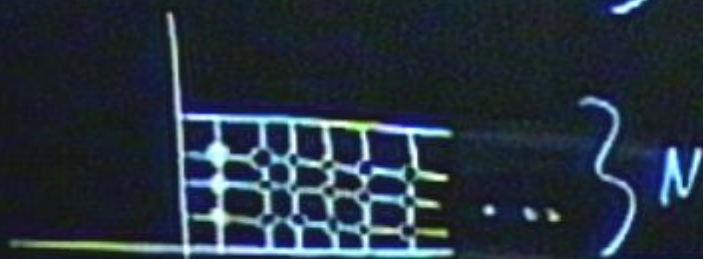
[0000 03]
[0000 03]
[0000 03]

[CANTALIA FIORAVANTI TATEO 03]

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$PSU(2, 2 | 4)$$



$$SU(N) \leftarrow N$$

TBA equations

↪ Equations of the form

$$Y_{a,s}(u) = \sum_{a',s'} K_{a,s}^{(a',s')} \star \log(1 + Y_{a',s'}(u)^{\pm 1}) + \delta_{s,0} L \log \frac{x^{[-a]}}{x^{[+a]}} + \langle \text{Source Terms} \rangle$$

[Gromov Kazakov Kozak Vieira 09]

[Bombardelli Fioravanti Tateo 09] [Autyunov Frolov 09]

$$x^{[\pm a]} = x(u \pm a \frac{i}{2}) = \frac{1}{2} \frac{u \pm a \frac{i}{2}}{g} + \frac{i}{2} \sqrt{4 - \left(\frac{u \pm a \frac{i}{2}}{g} \right)^2}$$


- $Y_{a,s}(u)$ is a function of $a, s \in \mathbb{Z}$ and u in \mathbb{R}



- **Extra assumption** : Excited states obey the same equations.

Each state correspond to a different solution of Y-system, characterized by its zeroes and poles

Outline

- 1 Motivation : Understanding Y-system better
- 2 Solving Hirota \rightsquigarrow Q-functions
- 3 Understanding symmetries from strong coupling 

Q-functions solve Hirota equation

Solving
AdS/CFT
Y-system
Leurent
ivation
unctions
nmetries

Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \} 3 \\ \} 4 - 3 \end{array} \right.$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

Q-functions solve Hirota equation

Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \} 3 \\ \} 4 - 3 \end{array} \right.$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

Q-functions solve Hirota equation

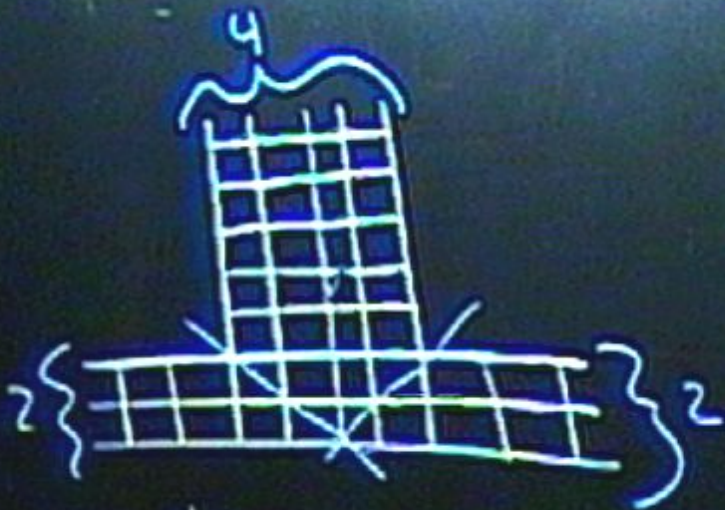
Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \left. \vphantom{\begin{array}{c} q_1^{[+s+2]} \\ q_1^{[+s]} \\ q_1^{[+s-2]} \end{array}} \right\} 3 \\ \left. \vphantom{\begin{array}{c} \tilde{q}_1^{[-s]} \\ \tilde{q}_2^{[-s]} \\ \tilde{q}_3^{[-s]} \\ \tilde{q}_4^{[-s]} \end{array}} \right\} 4 - 3 \end{array} \right.$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$PSU(2, 2 | 4)$$



$$SU(N) \leftarrow N$$

Q-functions solve Hirota equation

Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} 3 \\ \\ 4-3 \end{array}$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

Q-functions solve Hirota equation

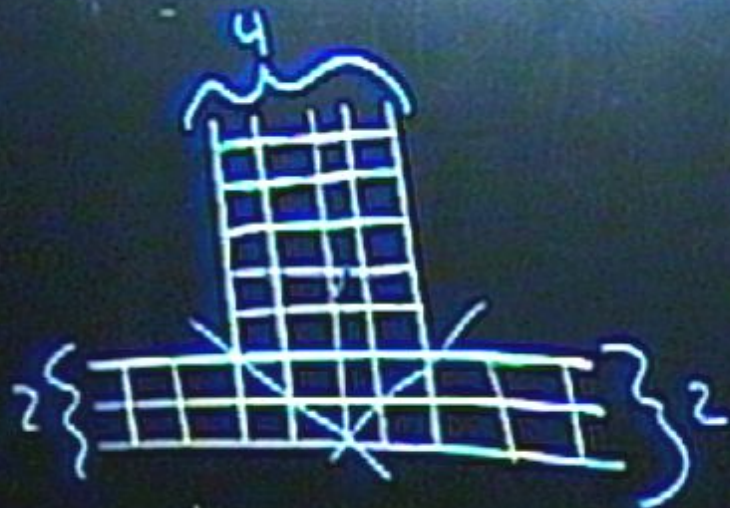
Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} 3 \\ \\ 4-3 \end{array}$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS

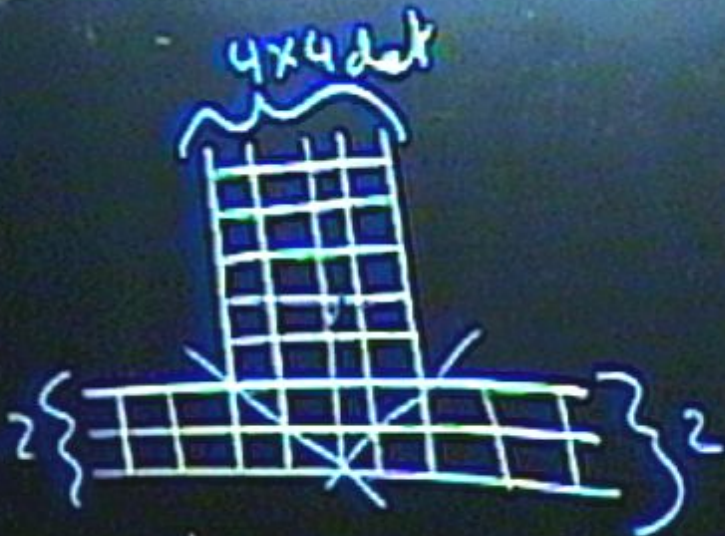


$$PSU(2, 2 | 4)$$

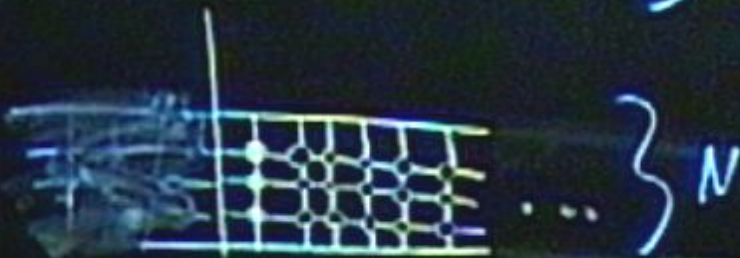


$$SU(N) \quad \leftarrow N$$

SOLVING HIROTA \longrightarrow FINITE NUMBER OF FUNCTIONS

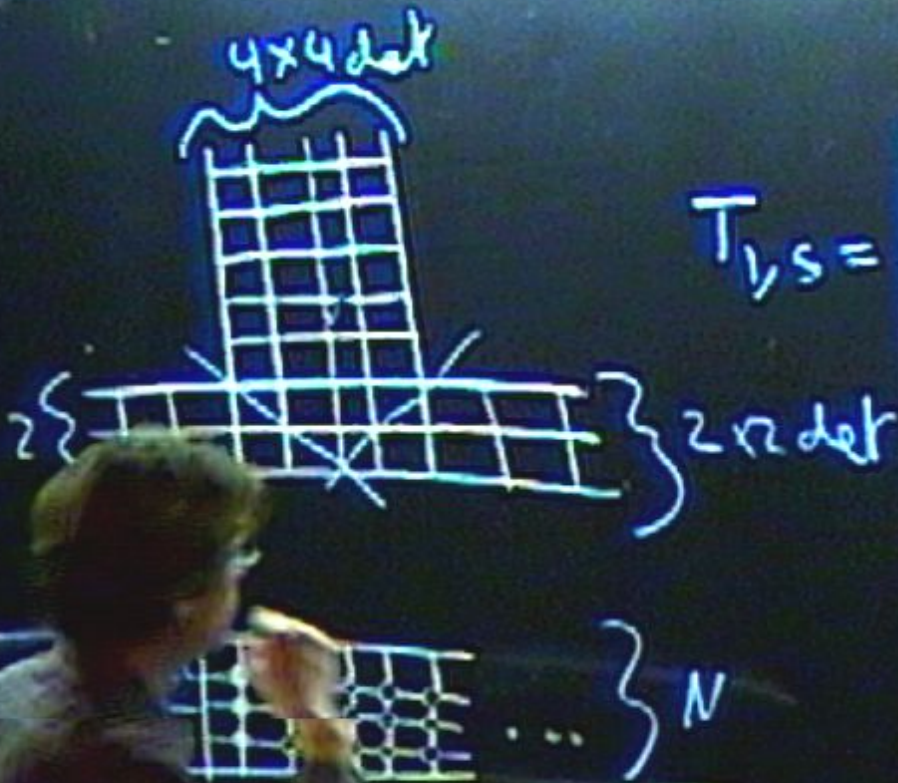


$$\text{PSU}(2, 2 | 4)$$



$$\text{SU}(N) \quad \leftarrow N$$

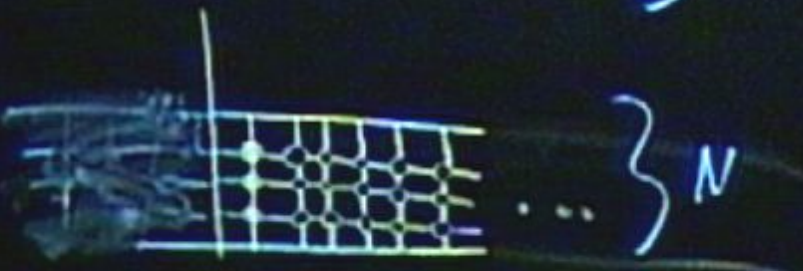
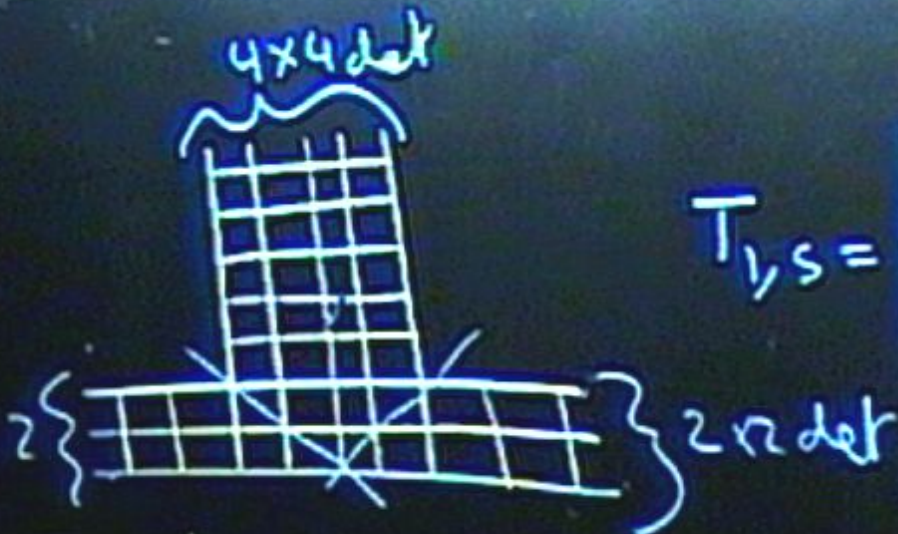
SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$T_{1,s} = \begin{vmatrix} q_1 [s] & q_2 [+s] \\ \bar{q}_1 [s] & -\bar{q}_2 [s] \end{vmatrix}$$

$SU(N) \quad \leftarrow N$

SOLVING HIROTA \longrightarrow FINITE NUMBER OF FUNCTIONS

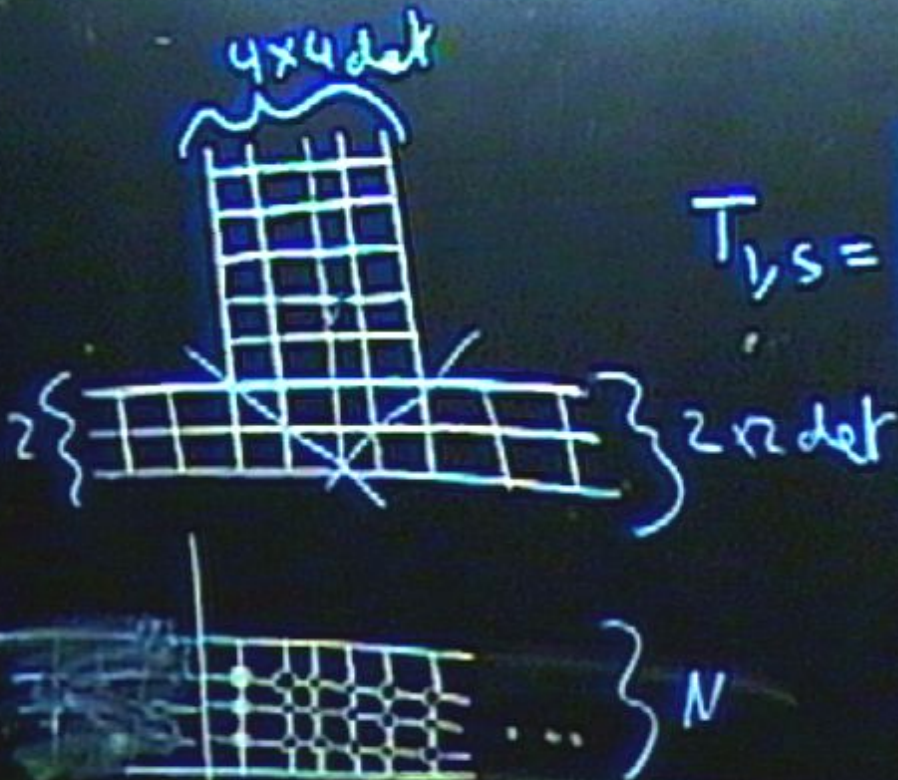


$$T_{1,s} = \begin{vmatrix} q_1 [s] & q_2 [+s] \\ q_1 [s] & -q_2 [s] \end{vmatrix}$$

[GROM, KAZAKOV, SL, TSUBO]

SU(N) \leftarrow N

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$T_{1,5} = \begin{vmatrix} q_1 [1,5] & q_2 [+5] \\ \bar{q}_1 [5,1] & -\bar{q}_2 [5,5] \end{vmatrix}$$

[GROM, KAZAKOV, SL, TSUBO]

$$SU(N) \leftarrow N$$

Q-functions solve Hirota equation

Hirota equation is solved by determinants of Q-functions :
eg. for $SU(4)$,

$$T_{3,s} = \left| \begin{array}{cccc} q_1^{[+s+2]} & q_2^{[+s+2]} & q_3^{[+s+2]} & q_4^{[+s+2]} \\ q_1^{[+s]} & q_2^{[+s]} & q_3^{[+s]} & q_4^{[+s]} \\ q_1^{[+s-2]} & q_2^{[+s-2]} & q_3^{[+s-2]} & q_4^{[+s-2]} \\ \tilde{q}_1^{[-s]} & \tilde{q}_2^{[-s]} & \tilde{q}_3^{[-s]} & \tilde{q}_4^{[-s]} \end{array} \right| \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} 3 \\ \\ 4-3 \end{array}$$

- $q_i^{[+k]} = q_i(u + k \frac{i}{2})$

[Baxter 72], [Pasquier Gaudin 92], [Bazhanov Lykhanov Zamolodchikov 96],
[Derkachov, 99], [Bytsko Teschner 06],
[Bazhanov Frassek Lukowski Meneghelli Staudacher 10]
[Kazakov S.L. Tsuboi 10]

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints:
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size)

SPECTRAL PROBLEM

FINITE SET OF EQUATIONS

* TBA EQUATIONS

$$y^+ y^- = \frac{(1+y|)(1+y)}{(1+y^-|)(1+y^-)}$$

HAROTA GAUGE

ANALYTICITY

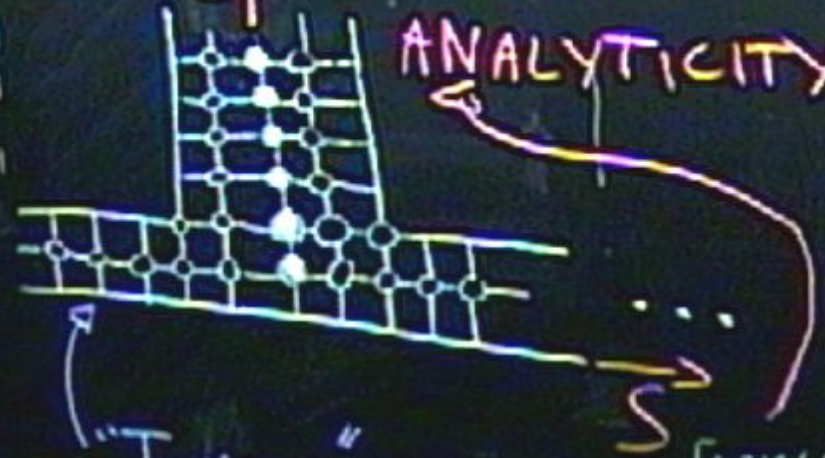
COMPLICATED KERNELS

EXTRA INFORMATION

∞ set of NLIES

ANALYTICITY

- [AF 03]
- [AF 03]
- [AF 03]



"T-Hook"

[CIRCUIT FLOWING TARD 03]

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)

$$T + T = \int_{\Sigma} T_{ab} dx^a dx^b + TT$$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)

$$T^+ T^- = \int_{\mathcal{C}_+} T_{ab} dx^a dx^b + T T$$

$T(u+i\epsilon)$

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints:
 - $\det = 1$
 - invariance under a Σ_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size)

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)

$x \rightarrow T_a$



SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
"CHOOSING THE GAUGE"

* $T_a \approx T_2, a$

* ANALYTICITY STRIPS

$+i\epsilon/2$

R

$-i\epsilon/2$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
"CHOOSING THE GAUGE"

* $T_a \approx T_{2,a}$

* ANALYTICITY STRIPS

==

==

$+i\epsilon/2$

$T_{1,s}$

R

==

==

==

$-i\epsilon/2$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
"CHOOSING THE GAUGE"

* $T_{\alpha\beta} = T_{\beta\alpha}$

* ANALYTICITY STRIPS



$+i\epsilon/2$

$T_{1,1}$

R



$-i\epsilon/2$

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_0 \approx T_2, a$



* ANALYTICITY STRIPS



$+i\gamma_2$

* \mathbb{Z}_2 Symmetry

$T_{1,5}$



R



$-i\gamma_2$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{2,2} = T_{2,-2}$$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{L,3} = T_{L,-3}$$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{\Omega,3} = T_{\Omega,-3}$$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{0,0} = T_{2,0}$



* ANALYTICITY STRIPS



$+i\epsilon/2$

* Z_4 Symmetry

$T_{1,1}$



R

$\chi_{1,1} \sim \lambda_i^s$



$-i\epsilon/2$

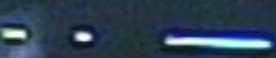
SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{a, \alpha} \approx T_{2, a}$

* ANALYTICITY STRIPS

* \mathbb{Z}_2 Symmetry

$\chi_{1, s} \sim \lambda_i^s = \sum \lambda_i^{-s} \sim \hat{\chi}_{1, -s}$



$+i\gamma/2$

\mathbb{R}

$-i\gamma/2$



\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{\Omega} = T_{\Omega^{-1}}$$

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{L,S} = \tilde{T}_{L,S}$$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{1,3} = \hat{T}_{1,3}$$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{1,s} = \hat{T}_{1,-s}$$

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

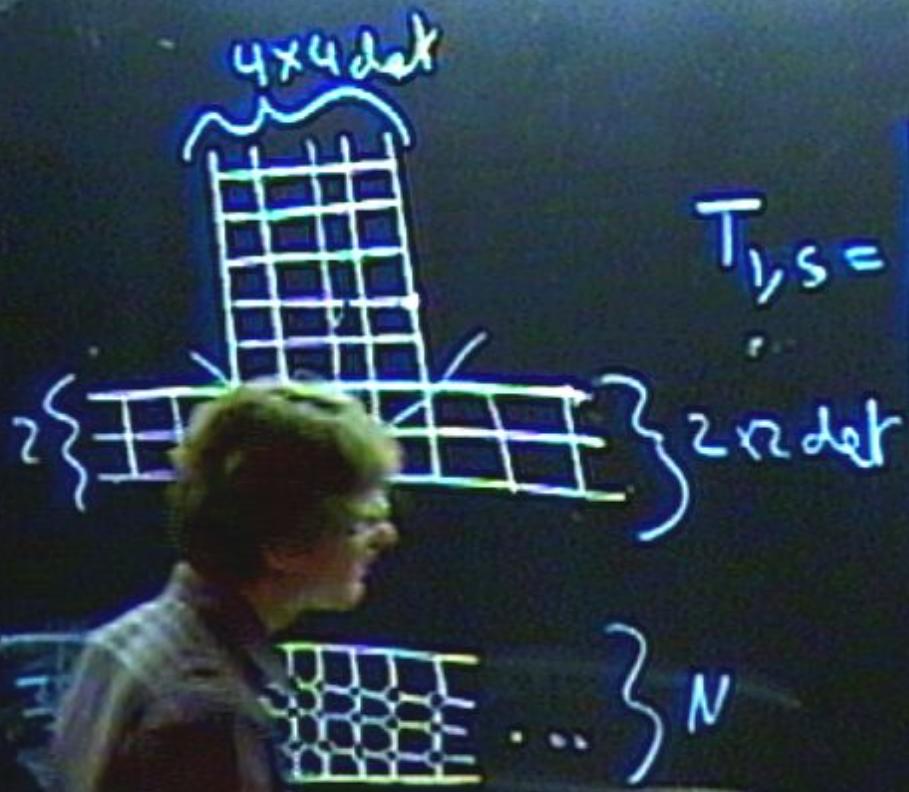
$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{1,s} = \hat{T}_{1,-s}$$

SOLVING HIROTA — FINITE NUMBER OF FUNCTIONS



$$T_{1,s} = \begin{vmatrix} q_1 & q_2 \\ q_1 [E^s] & -q_2 [E^s] \end{vmatrix}$$

[GROMOV, KAZAKOV, SL, TSUBOKAWA]

SU(N) ← N

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{2,2} = T_{2,0}$



* ANALYTICITY STRIPS

* Z_4 Symmetry

$T_{1,s}$



$+i\gamma_2$

R

$T_{1,s} = -\overline{T_{1,-s}}$



$-i\gamma_2$

$\sum x_{1,s} z^s = \text{sdet} \left(\frac{1}{1-s} \right)$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{0,2} = T_{2,0}$

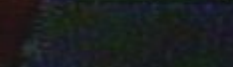
* ANALYTICITY STRIPS

* Z_4 Symmetry

$T_{1,5} = -T_{1,-5}$

$T_{1,5}$

$\sum x_{1,5} z^5 = \text{residue} \left(\frac{1}{1-z^5} \right)$



$+i\pi/2$

R

$-i\pi/2$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{0,2} = T_{2,0}$

* ANALYTICITY STRIPS

* Z_4 Symmetry



$+i\gamma_2$

R

$-i\gamma_2$

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,

$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -j\bar{u} + \frac{1}{2T} \sqrt{2g} \frac{d\bar{u}}{u-\bar{u}}$
- generalizable to all T-functions 

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$T_{1,s} = \begin{vmatrix} 1 & q_2 [s] \\ 1 & -q_2 [s] \end{vmatrix}$$

$\rightarrow [s]$
 $\rightarrow [s]$
 $\rightarrow [s]$

[GROMOV, KAZAKOV, SL, TSUBOI]

$SU(N) \leftarrow N$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{2,2} = T_{2,0}$



* ANALYTICITY STRIPS

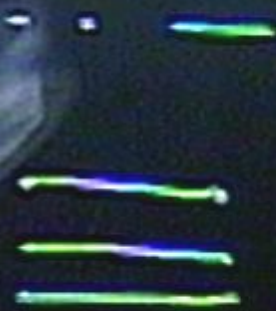
$+i\epsilon_2$

* \mathbb{Z}_4 Symmetry

$T_{1,1}$

R

$$\begin{array}{ccc} \uparrow & & \uparrow \\ T_{1,1} & = & -T_{1,-1} \\ \uparrow & & \uparrow \\ T_{1,0} & = & -T_{1,0} \end{array}$$



$-i\epsilon_2$

$T_{1,0} = 0$

$q_1 = -q_2$

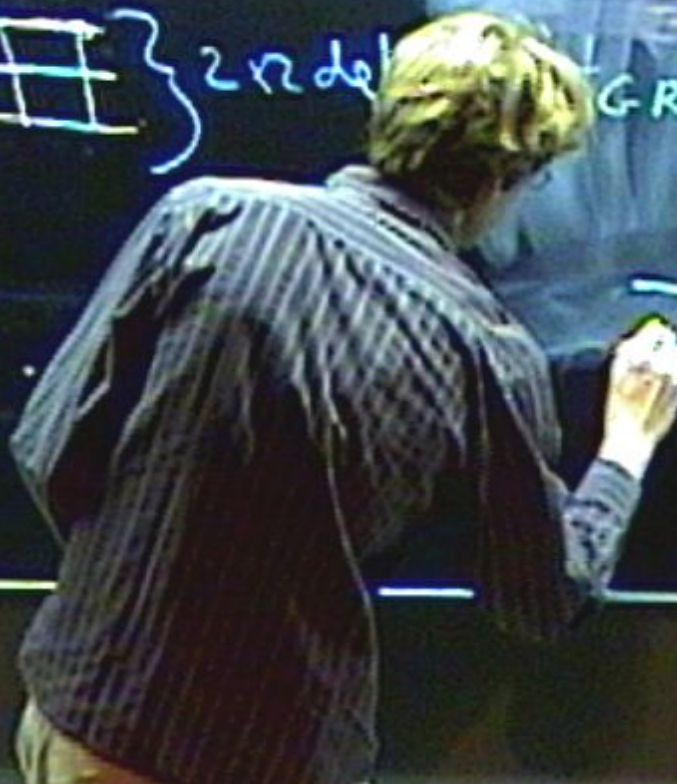
SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



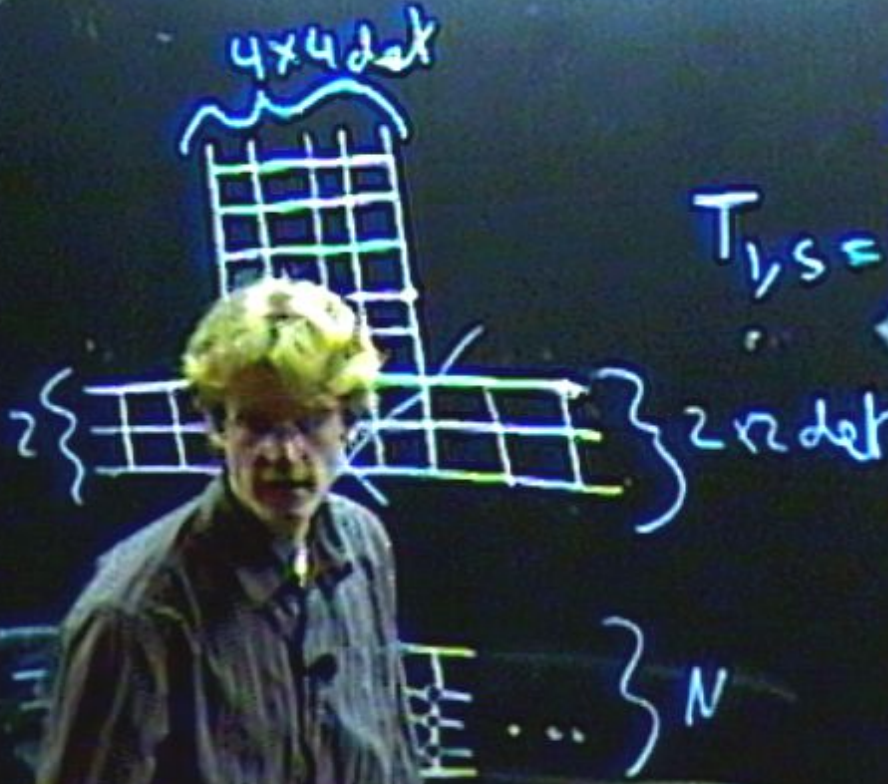
$$T_{1,5} = \begin{vmatrix} 1 & q_2 [+s] \\ 1 & -q_2 [s] \end{vmatrix}$$

$\rightarrow [s]$
 $\rightarrow [s]$
 $\rightarrow [s]$

GROMOV, KAZAKOV, SL, TSUBO



SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



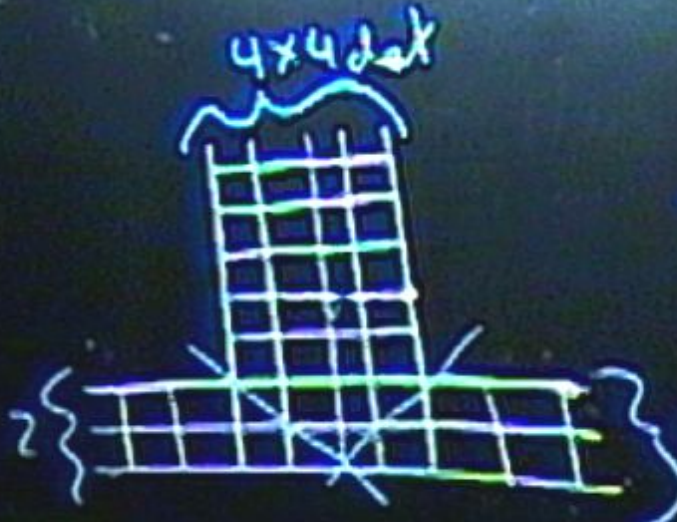
$$T_{1,5} =$$

$$\begin{vmatrix} 1 & q_2 [+s] \\ 1 & -q_2 [-s] \end{vmatrix}$$

[GROMOV, KAZAKOV, SL, TSUBO]

$$q_2 \equiv \dots \equiv R$$

SOLVING HIROTA → FINITE NUMBER OF FUNCTIONS



$$T_{1,5} = \begin{vmatrix} 1 & q_2 \\ -q_2 & 1 \end{vmatrix}$$

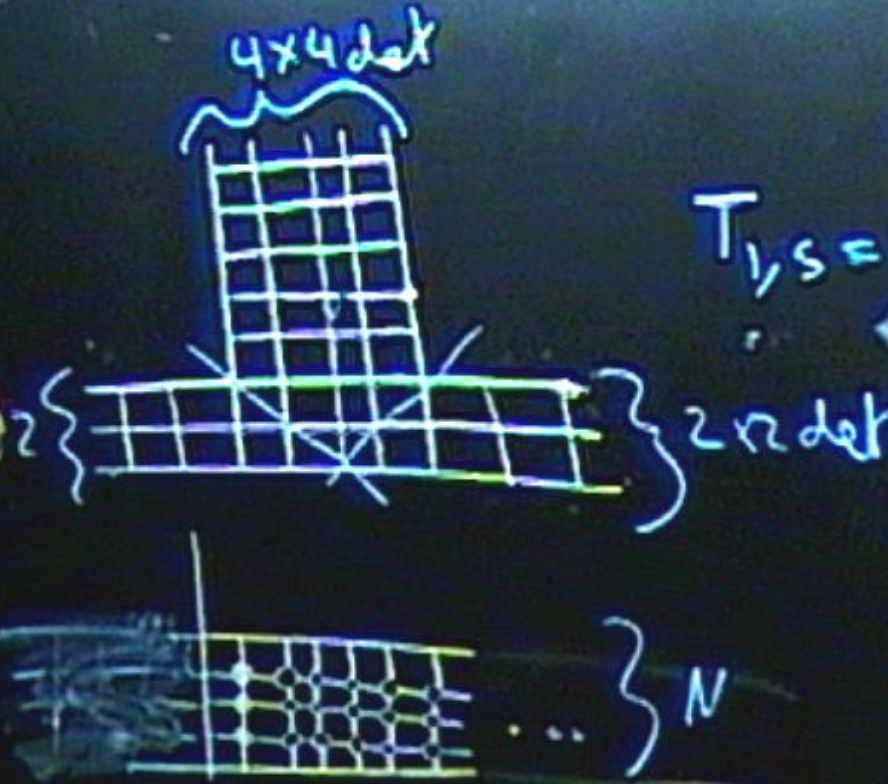


[G-ROM, KAZAKOV, SL, TSUBA]



$$q_2 \begin{vmatrix} \dots & R & q_2 \\ \dots & \dots & \dots \end{vmatrix}$$

SOLVING HIROTA → FINITE NUMBER OF FUNCTIONS



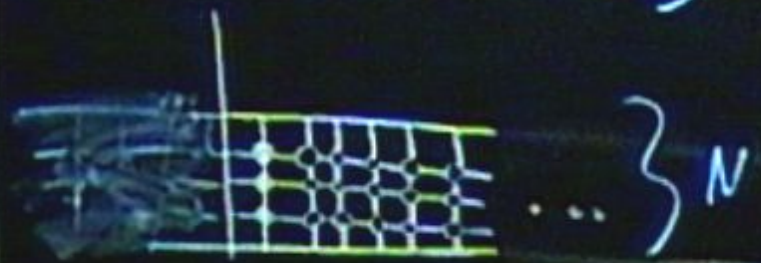
$$T_{1,5} =$$

$$\begin{vmatrix} 1 & q_2 [+s] \\ 1 & -q_2 [+s] \end{vmatrix}$$

[GROMOV, KAZAKOV, SL, TSUBO]

$$q_2 \begin{vmatrix} \text{---} \\ \text{---} \\ \text{---} \end{vmatrix} \dots R \begin{vmatrix} \text{---} \\ \text{---} \\ \text{---} \end{vmatrix}$$

SOLVING HIROTA \rightarrow FINITE NUMBER OF FUNCTIONS



$$T_{1,5} = \begin{vmatrix} 1 & q_2 [+s] \\ 1 & -q_2 [+s] \end{vmatrix}$$

[GROMOV, KAZAKOV, SL, TSUBO]

$$q_2 \begin{vmatrix} \text{---} \\ \text{---} \\ \text{---} \end{vmatrix} \dots R q_2 = \begin{vmatrix} \text{---} \\ \text{---} \\ \text{---} \end{vmatrix}$$

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{\alpha\beta}$ & $T_{2,a}$

* ANALYTICITY STRIPS

* \mathbb{Z}_2 Symmetry

$$T_{1,5} = -T_{1,-5}$$

$$\uparrow \quad \quad \quad \uparrow$$

$$T_{1,0} = -T_{1,0}$$

$$T_{1,0} = 0$$

$T_{1,5}$

$$g_2 = -\bar{g}_2$$

+i/2


R

-i/2

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,


$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,

$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 

\mathbb{Z}_4 symmetry

\mathbb{Z}_4 symmetry of the classical limit

$$\Omega = \hat{C}^{-1}(\Omega^{-1})^T \hat{C}$$

or

$$\{\lambda_i\} = \{1/\lambda_i\}$$

[Bena Polchinski Roiban]

Quantum case

$$T_{L,3} = \hat{T}_{L,3}$$

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

Classical limit

In the classical limit, $g \rightarrow \infty$, and $T_{a,s} \rightarrow T_{a,s}(u/g)$.
 \Rightarrow shifts by $\pm \frac{i}{2}$ in Hirota equation can be neglected.
 $\Rightarrow T_{a,s}(u) = \chi_{a,s}(\Omega(u))$ where $\Omega \in U(2, 2|4)$.

[Gromov Kazakov Tsuboi]

- Actually, the $PSU(2, 2|4)$ symmetry imposes more constraints :
 - $\det = 1$
 - invariance under a \mathbb{Z}_4 transformation

That gives extra symmetries of the characters (generalizing to symmetries of T-functions at finite size).

SYMMETRIES (MOTIVATED BY CLASSICAL LIMIT)
 "CHOOSING THE GAUGE"

* $T_{\alpha\beta} = T_{\beta\alpha}$

* ANALYTICITY STRIPS

* \mathbb{Z}_2 Symmetry

$T_{1,1}$

$$T_{1,1} = -T_{1,-1}$$

$$\uparrow \quad \quad \quad \uparrow$$

$$T_{1,0} = -T_{1,0}$$

$$T_{00} + T_{1,0} = 0$$

* $T_{00} = T_{00}$

$$g_{-} = -g_{+}$$

$+i\eta$


R

$-i\eta$

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,

$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $O(\log T_{\text{sc}}) \sim \frac{1}{\epsilon} \frac{E}{\epsilon}$
 - Exact Bethe equations arise as absence of poles of T-functions
- to be continued
 - currently restricted to symmetric sl_2 "sector" states
 - { numeric efficiency, ... } are to be studied
 - { best FINLIE formulation, ... }
 - application to other Y-systems?
 - BFKL
 - strong coupling construction of T ($\mathcal{T} = \text{trace} \Omega_{\pm}$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system

- analytic properties
- new symmetries
- Finite set of NLIEs
 - $\log T_{\text{reg}} \sim \frac{2\pi}{\beta} \frac{E}{\hbar v}$
 - Exact Bethe equations arise as absence of poles of T-functions

- to be continued

- currently restricted to symmetric \mathfrak{sl}_2 "sector" states
- { numeric efficiency, fast NLIE formulation } are to be studied
- application to other Y-systems?
- BFKL
- strong coupling construction of T (? $T = \text{trace} \Omega_{\text{reg}}$)
- weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$
 - Exact Bethe equations arise as absence of poles of T-functions

- to be continued
 - currently restricted to symmetric sl_2 "sector" states
 - numeric efficiency
 - best FNLIE formulation are to be studied
 - application to other Y-systems?
 - BFKL
 - strong coupling constructions of T (? $T \equiv \text{Tr} \mu^{\Omega}$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$
 - Exact Bethe equations arise as absence of poles of T-functions
- to be continued
 - currently restricted to symmetric sl_2 "sector" states
 - [numeric efficiency]
 - [best FNLIE formulation] are to be studied
 - application to other Y-systems?
 - BFKL
 - strong coupling construction of T (? $T = \text{tr}(\rho \Omega)$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$
 - Exact Bethe equations arise as absence of poles of T-functions
- to be continued
 - currently restricted to symmetric sl_2 “sector” states
 - $\left\{ \begin{array}{l} \text{numeric efficiency} \\ \text{best FiNLIE formulation} \end{array} \right.$ are to be studied
 - application to other Y-systems ?
 - BFKL
 - strong coupling construction of T (? $T = \langle \text{trace} \Omega \rangle$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$
 - Exact Bethe equations arise as absence of poles of T-functions
- to be continued
 - currently restricted to symmetric sl_2 “sector” states
 - $\left\{ \begin{array}{l} \text{numeric efficiency} \\ \text{best FiNLIE formulation} \end{array} \right.$ are to be studied
 - application to other Y-systems ?
 - BFKL
 - strong coupling construction of T (? $T = \langle \text{trace} \Omega \rangle$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$

Really

Thank you !

- TO BE CONTINUED
 - currently restricted to symmetric sl_2 “sector” states
 - $\left\{ \begin{array}{l} \text{numeric efficiency} \\ \text{best FiNLIE formulation} \end{array} \right.$ are to be studied
 - application to other Y-systems ?
 - BFKL
 - strong coupling construction of T (? $T = \langle \text{trace} \Omega \rangle$)
 - weak coupling interpretation of T

Conclusion

- A better understanding of Y-system
 - analytic properties
 - new symmetries
 - Finite set of NLIEs
 - $\partial \log T_{0,0} \xrightarrow{u \rightarrow \infty} \frac{2E}{u}$

Really


Thank you !

- TO BE CONTINUED
 - currently restricted to symmetric sl_2 “sector” states
 - $\left\{ \begin{array}{l} \text{numeric efficiency} \\ \text{best FiNLIE formulation} \end{array} \right.$ are to be studied
 - application to other Y-systems ?
 - BFKL
 - strong coupling construction of T (? $T = \langle \text{trace} \Omega \rangle$)
 - weak coupling interpretation of T

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,


$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,

$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 

Example of the right Strip

For instance, in a real gauge where $q_1 = 1$,

$$\hat{T}_{1,0} = -\bar{q}_2(u) - q_2(u)$$

- Hence $q_2(u) = -iu + \frac{1}{2l\pi} \int_{-2g}^{2g} \frac{\rho(v)}{v-u}$
- generalizable to all T-functions 