

Title: Overview of the Challenges

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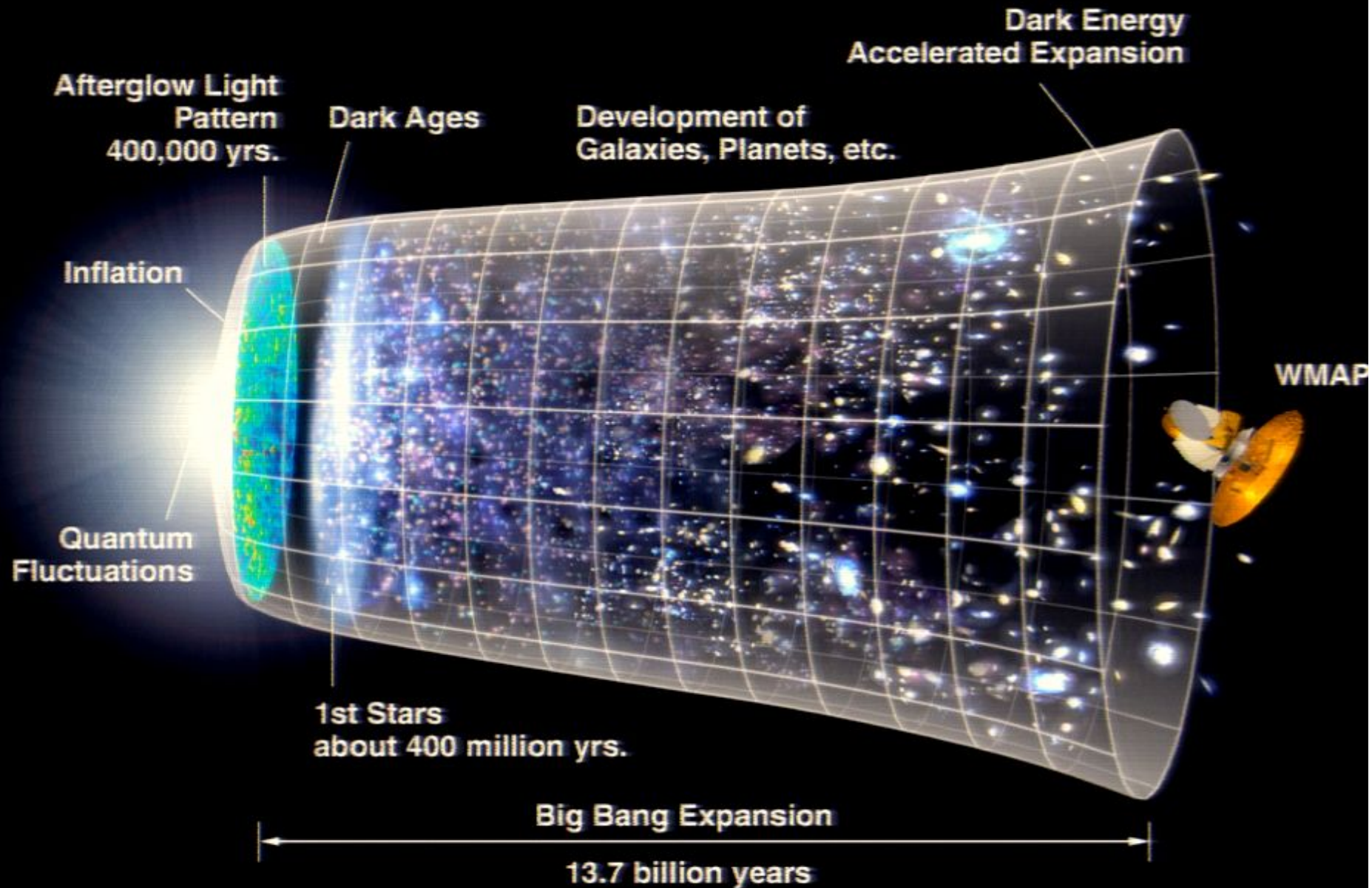
Abstract:

# Challenges for Early Universe Cosmology

Neil Turok, Perimeter Institute

- \* Concordance Model
- \* Anthropics
- \* Measures

# Concordance Model



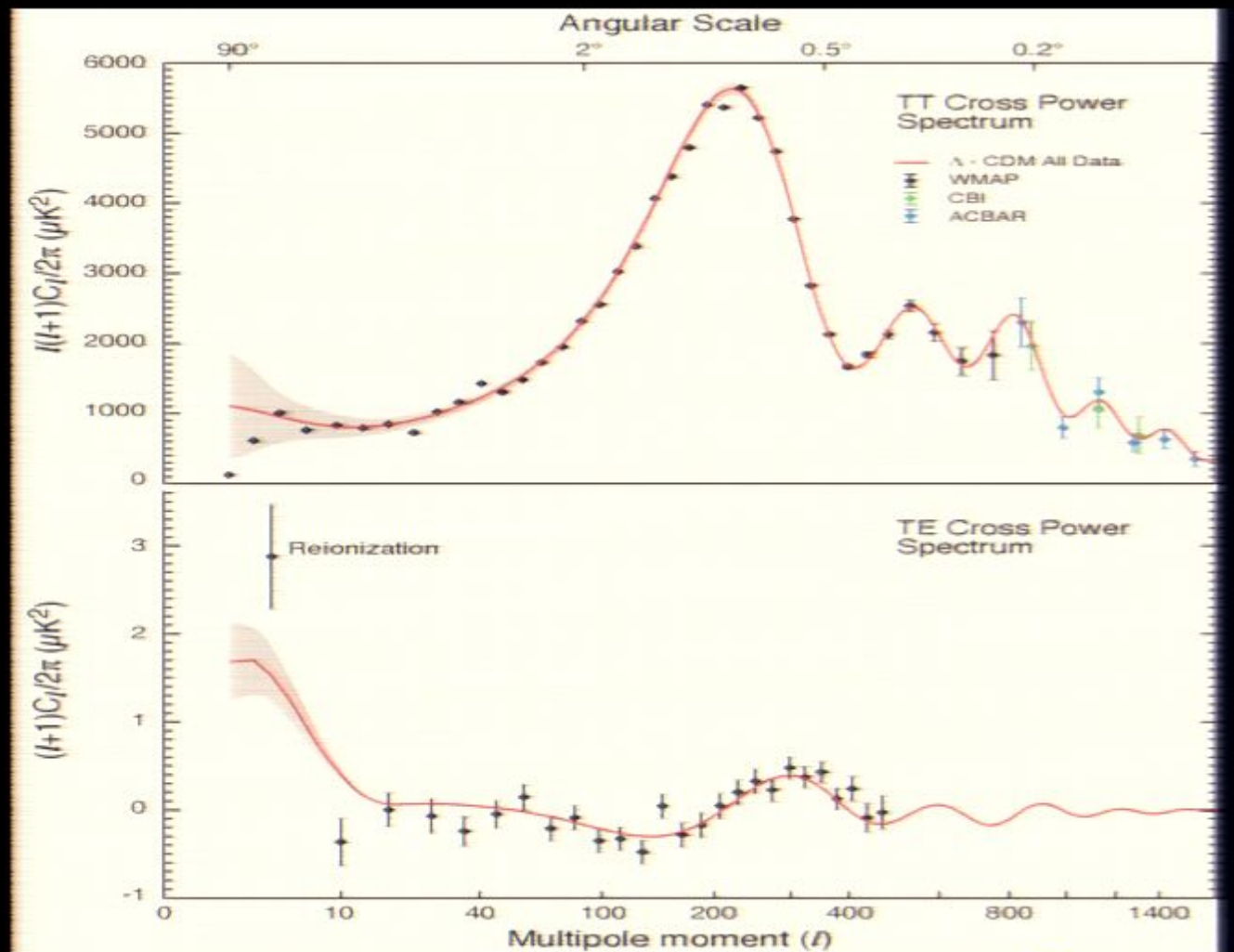
# Success!

fluctuation level:  
temperature

Zeldovich, Peebles+Yu 70's  
Bond+Efstathiou 80's

polarization

Coulson, Crittenden, NT 94

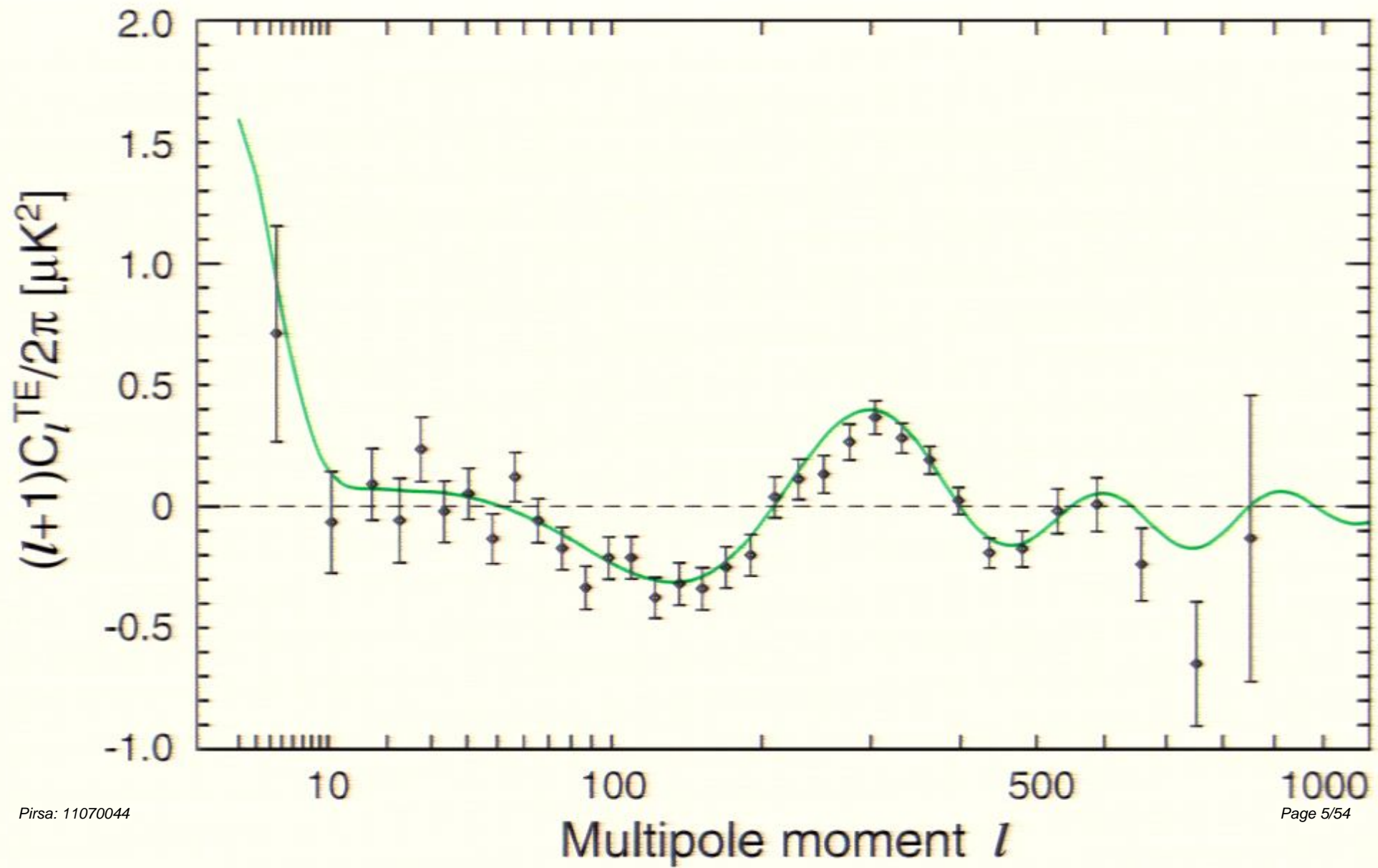


100 10 1 .1 ( $l = 2\pi/\theta$ )

Angle on Sky (Degrees)



# WMAP 5 year TE



# good evidence for ...

nearly flat FRW universe:

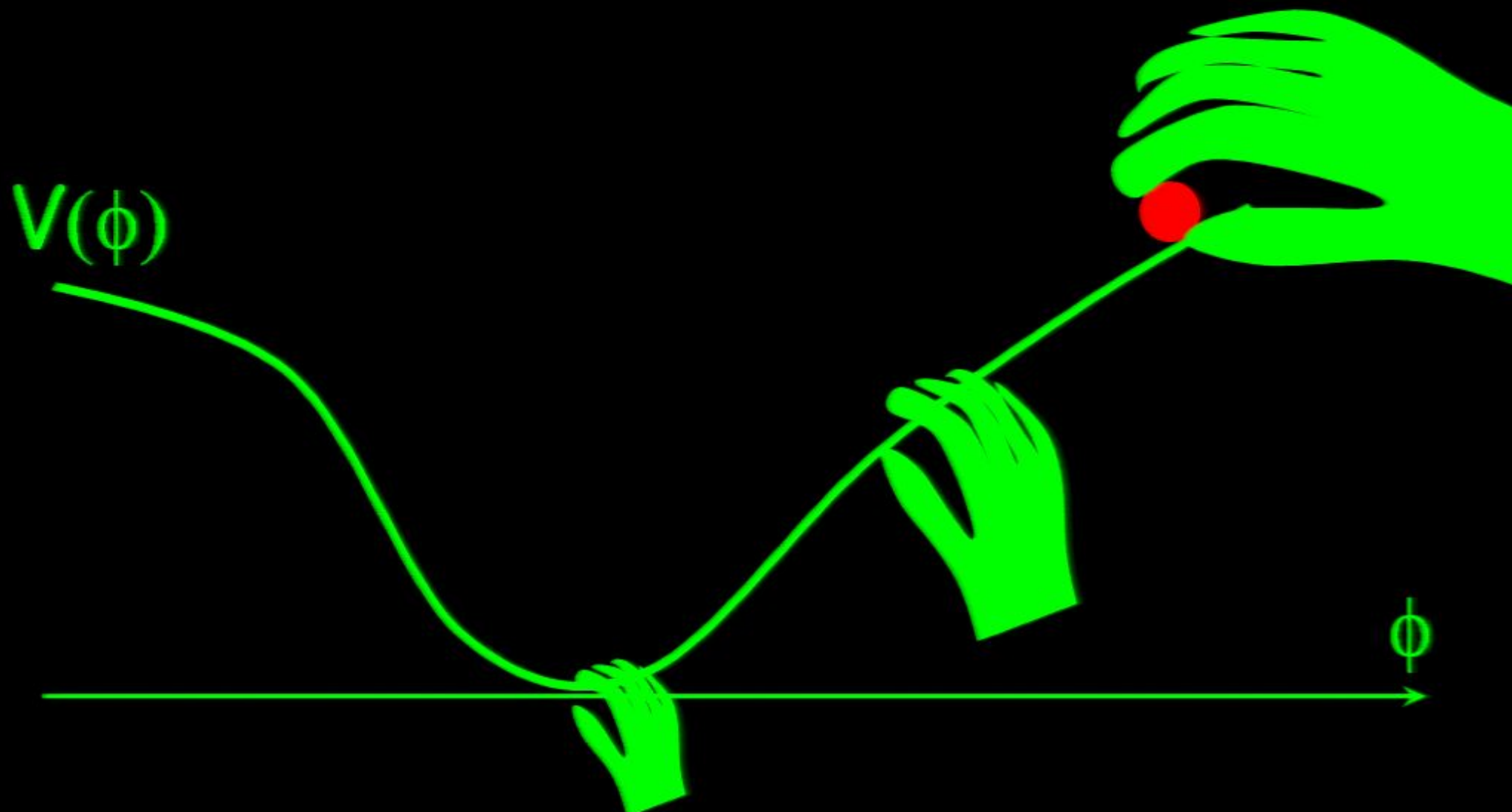
$$\Omega_{\Lambda} : \Omega_{\text{CDM}} : \Omega_{\text{B}} : \Omega_{\text{v}} : \Omega_{\gamma} \sim 0.7 : 0.25 : 0.05 : 0.003 : 0.0003$$

primordial perturbations

- \* linear
- \* growing mode
- \* nearly scale-invariant
- \* nearly 'adiabatic'
- \* nearly Gaussian

universe is: geometrically simple  
compositionally complex

inflation



# Challenges for Inflation

- \* initial conditions
- \* fine-tuned potentials
- \*  $\Lambda \sim 10^{-120}$  ;  $\Lambda_I \sim 10^{-15}$
- \* eternal inflation
  - "anything that can happen will happen:  
and it will happen an infinite number  
of times" A. Guth, 2002
- \* landscape
  - measure problem
- \* reliance on anthropic arguments



Anthropics: the universe is the way it is because of a (gigantic) selection effect

why are we then not "close to extinction"?

why is the universe **so** geometrically simple?

- \* very nearly flat
- \* very nearly spatially homogeneous
- \* the fluctuations are scale-invariant on scales much larger than galaxy scales

for anthropic arguments to become convincing, we need well-defined

- \* measures
- \* projections

# The Measure Problem

What is the natural measure on the space of cosmological solutions?

What is the likelihood of a universe like ours, in a given physical model? eg inflation, cyclic,



# Two key ingredients

- I: Penrose critique of inflation - Hamiltonian evolution almost never turns a generic state into a highly unusual state. Canonical measure is invariant.
- II: Counting of states in gravity can only be done in an asymptotic region where global properties of spacetime become sharp

In a specific setup, we shall construct a canonical measure and show a universe like ours is extremely unlikely in slow-roll inflationary models

- standard slow-roll inflation

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{1}{3}\rho_\phi - \frac{k}{a^2}; \quad \rho_\phi = \frac{1}{2}\dot{\phi}^2 + V(\phi)$$

$$\ddot{\phi} + 3\frac{\dot{a}}{a}\dot{\phi} = -V_{,\phi}$$

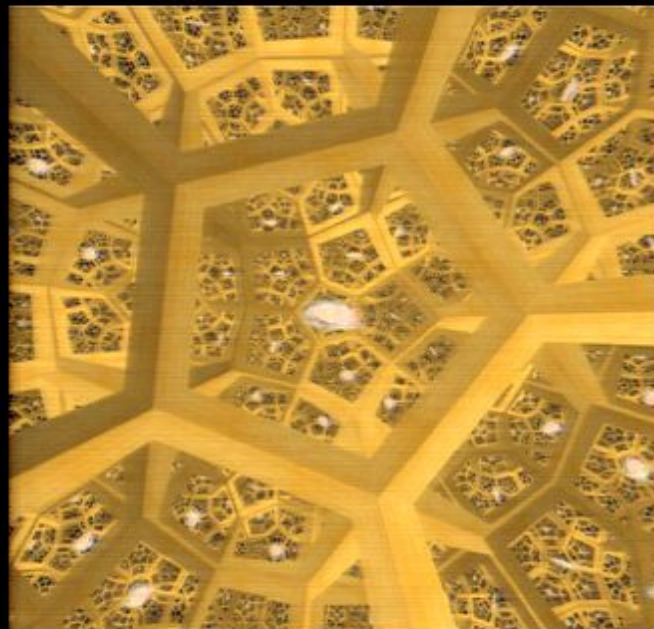
$$\Rightarrow \dot{H} = -\frac{1}{2}\dot{\phi}^2 + \frac{k}{a^2}$$

- Hamiltonian and time reversal invariant

(focus on FRW spacetimes and assume  $V(\phi)$  is monotonic away from its min)



Assume  $k=-1$  (so  $a$  and  $H$  are monotonic),  
zero  $\Lambda$ , and compactify the spatial slices



- \* a mathematical device to keep everything finite: the results do not depend on the compactification volume
- \* (but in fact has been advocated as a very natural setup for chaotic inflation)

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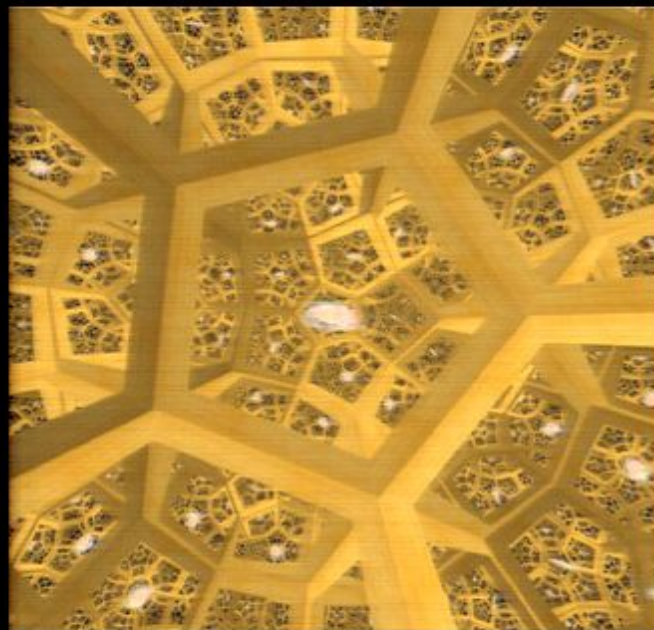
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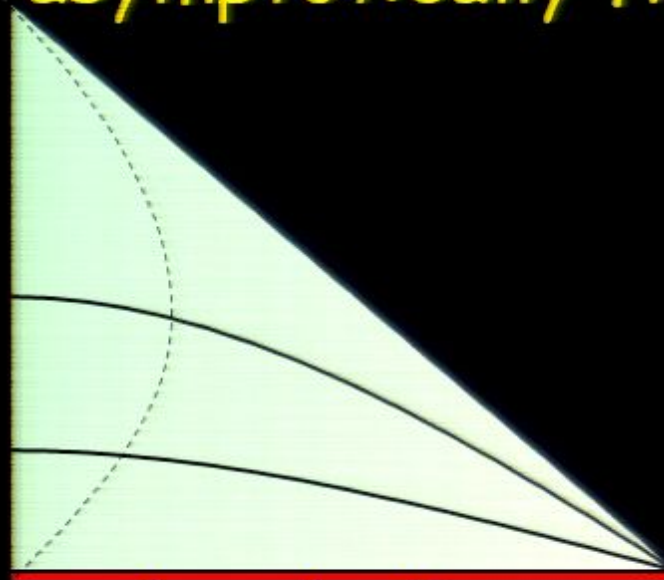
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# 2-parameter family of solutions with an initial singularity:

asymptotically flat



singularity



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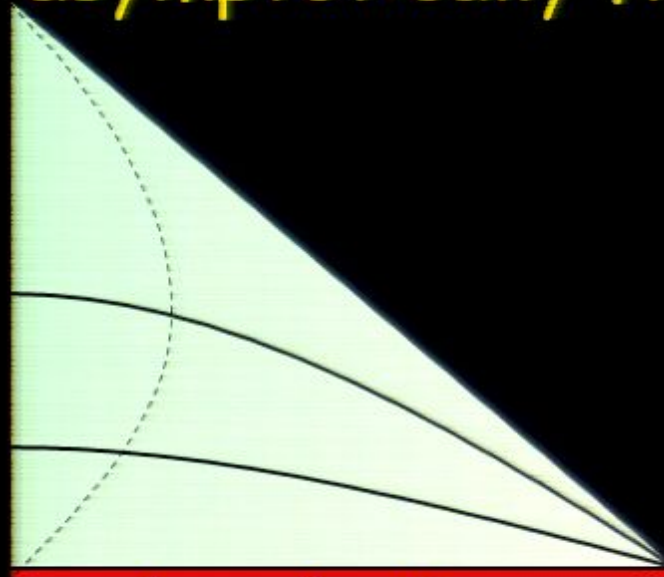
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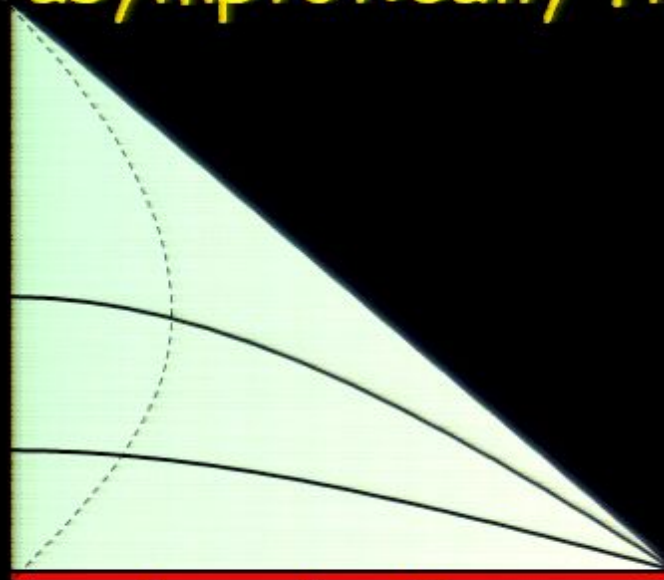
singularity

# Criteria for a Measure

- (i) Positive, normalisable
- (ii) Independent of slicing or coordinates on either space-time or field space
- (iii) Independent of ad hoc external structures eg cutoffs, comoving observers, "volume" factors .
- (iv) Natural extension of canonical quantum measure for fluctuations to the background (why use for one but not the other???)

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## Canonical measure on space of solutions

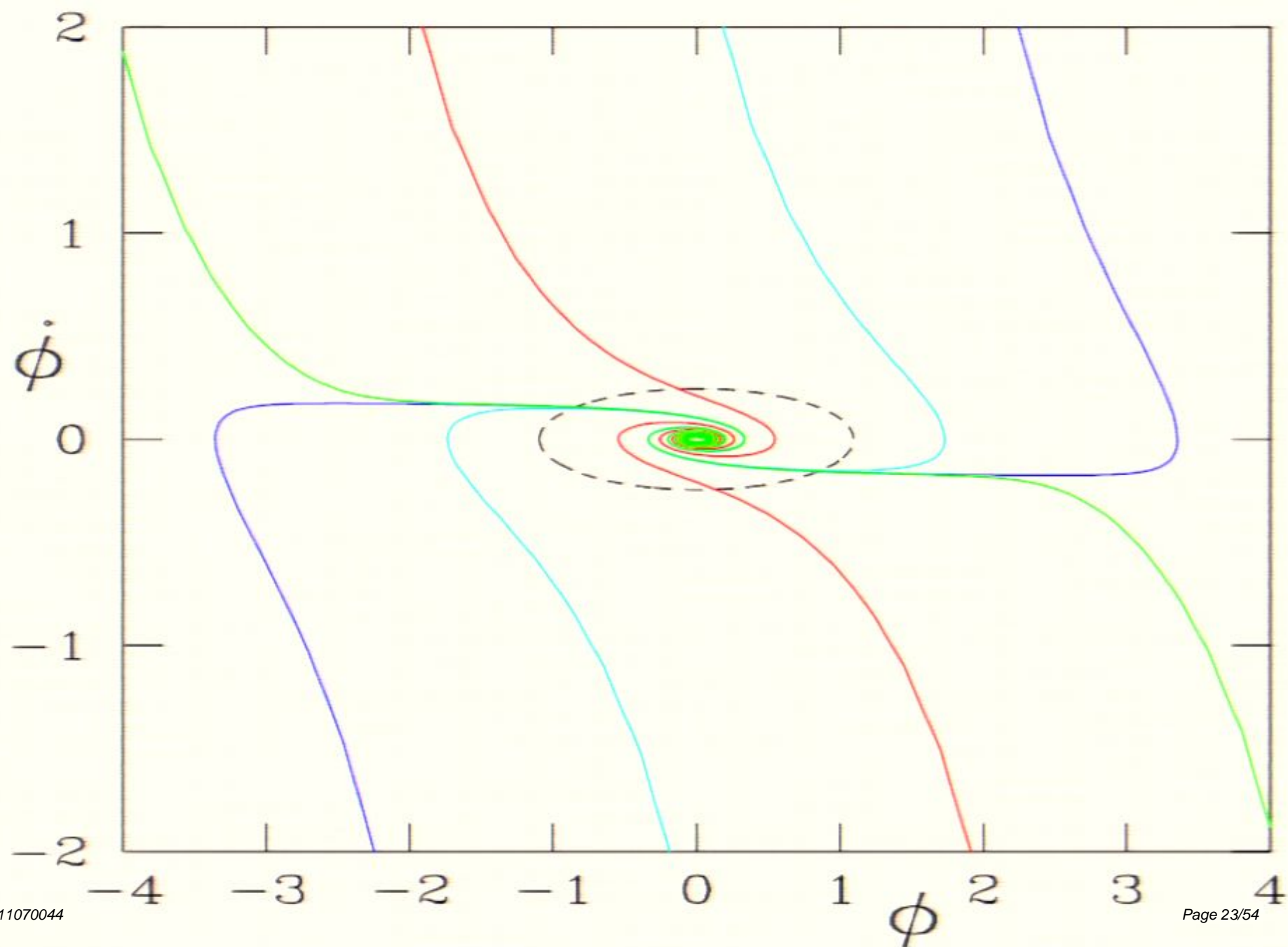
$$\omega_c = dp_a \wedge da + dp_\phi \wedge d\phi$$

$$\int_{\Sigma} \omega_c \Big|_{\mathbb{H}=0}$$

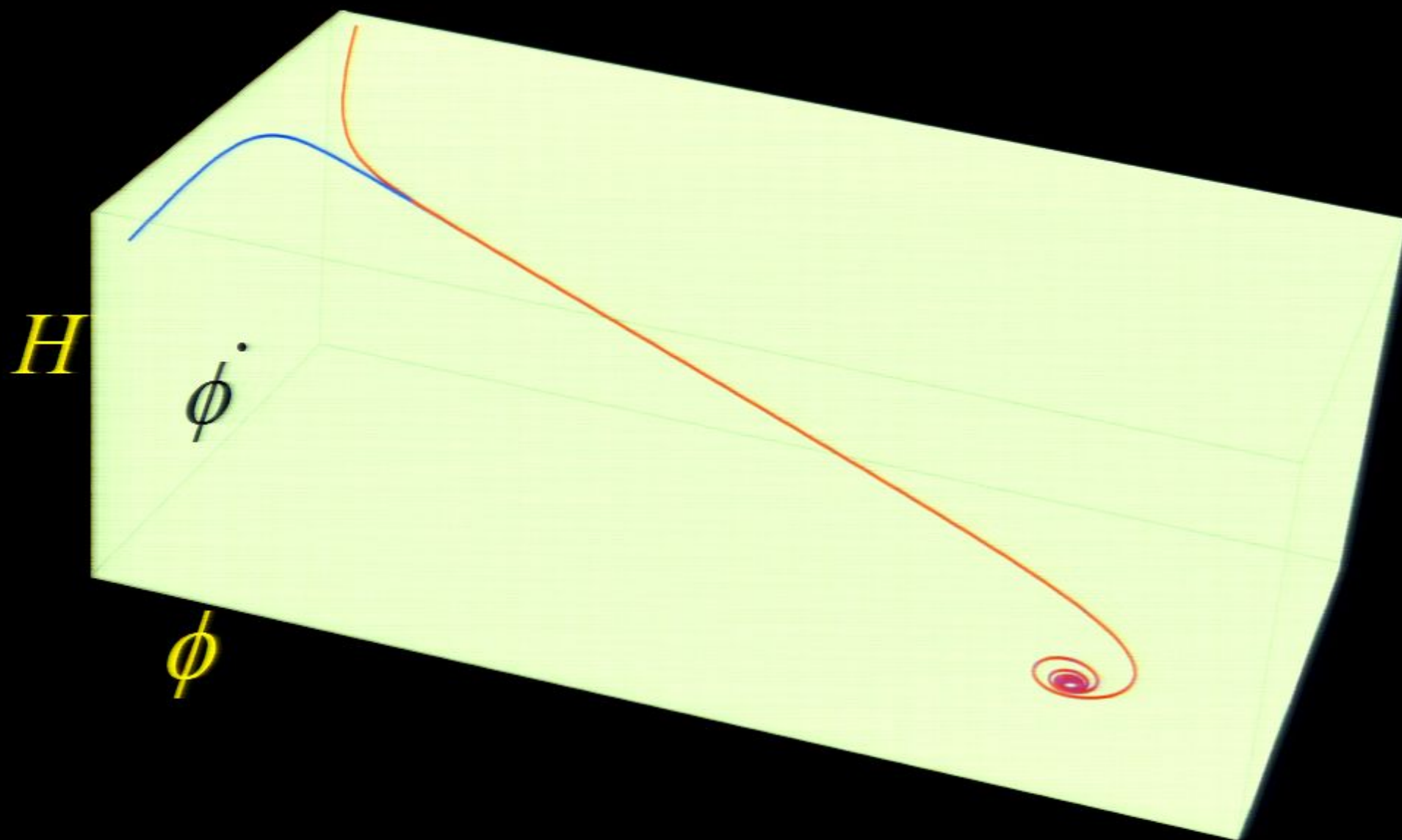
Liouville  
Gibbons, Hawking, Stewart  
Hawking, Page  
Hollands Wald  
Kofman, Linde, Mukhanov  
Gibbons, NT  
Carroll, Tan

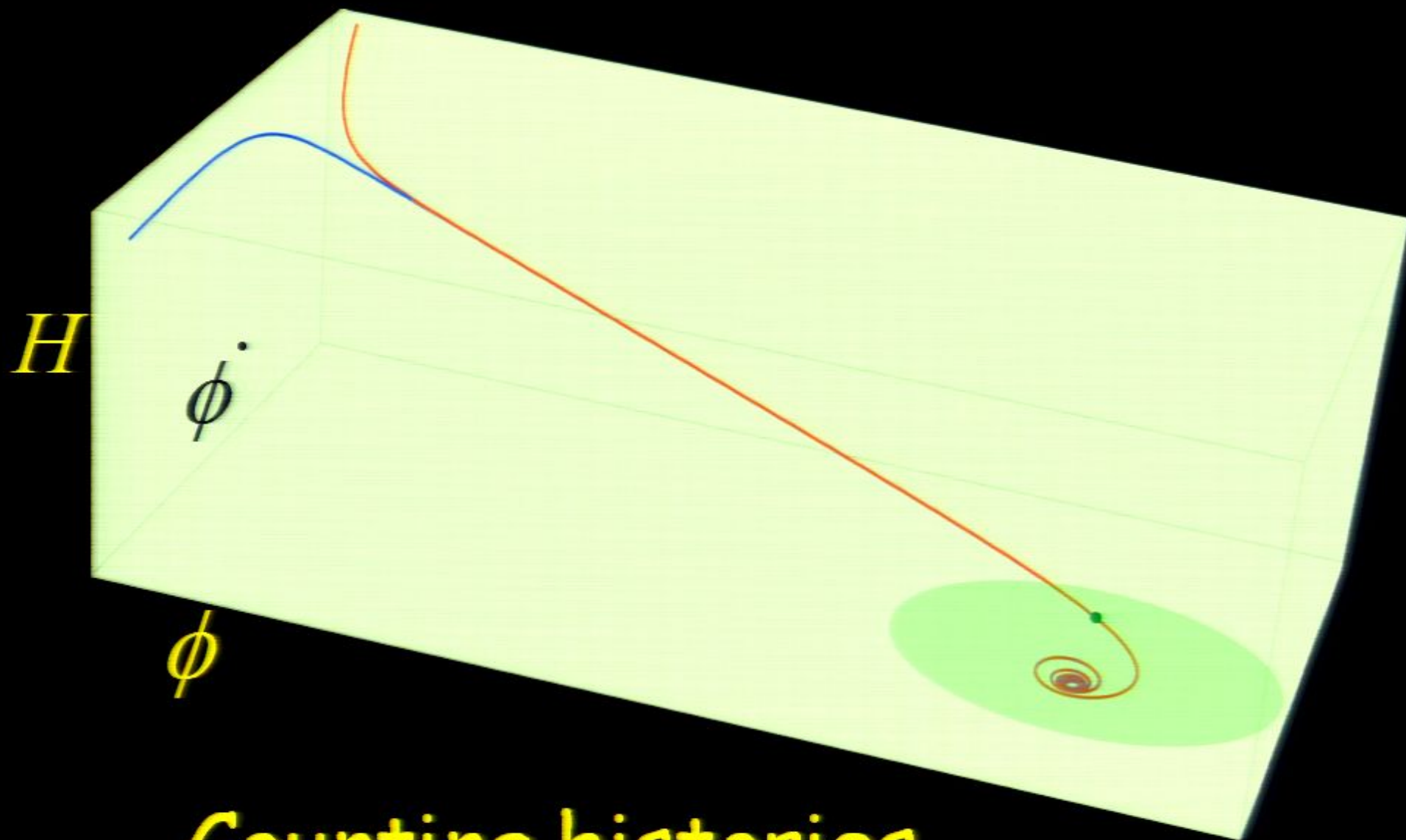
with  $\Sigma$  pierced once by every trajectory  
e.g.  $a=\text{const}$  or  $H=\text{const}$

Satisfies all of conditions (i)-(iv)  
**except** normalisability, because  $\Sigma$  is not  
compact (because  $\mathbb{H}$  isn't positive)









Counting histories

Flat space canonical (Gibbs) ensemble.

Cannot just integrate over Liouville: instead, we maximise entropy  $S = -\sum p_i \ln p_i$

subject to  $E = \sum p_i E_i$  (assuming  $E_i$  b.b.)

Note: in information theory approach, max ent principle is very general, can even be applied to non-equilibrium situations (see e.g. papers of E.T. Jaynes)

But in GR,  $\mathbb{H}=0$  on all physical states, so cannot constrain its expectation value

What do we do?



In  $k=-1$ , zero  $\Lambda$  cosmologies, matter density is diluted away at large  $a \rightarrow$  gravity becomes negligible, expansion becomes adiabatic

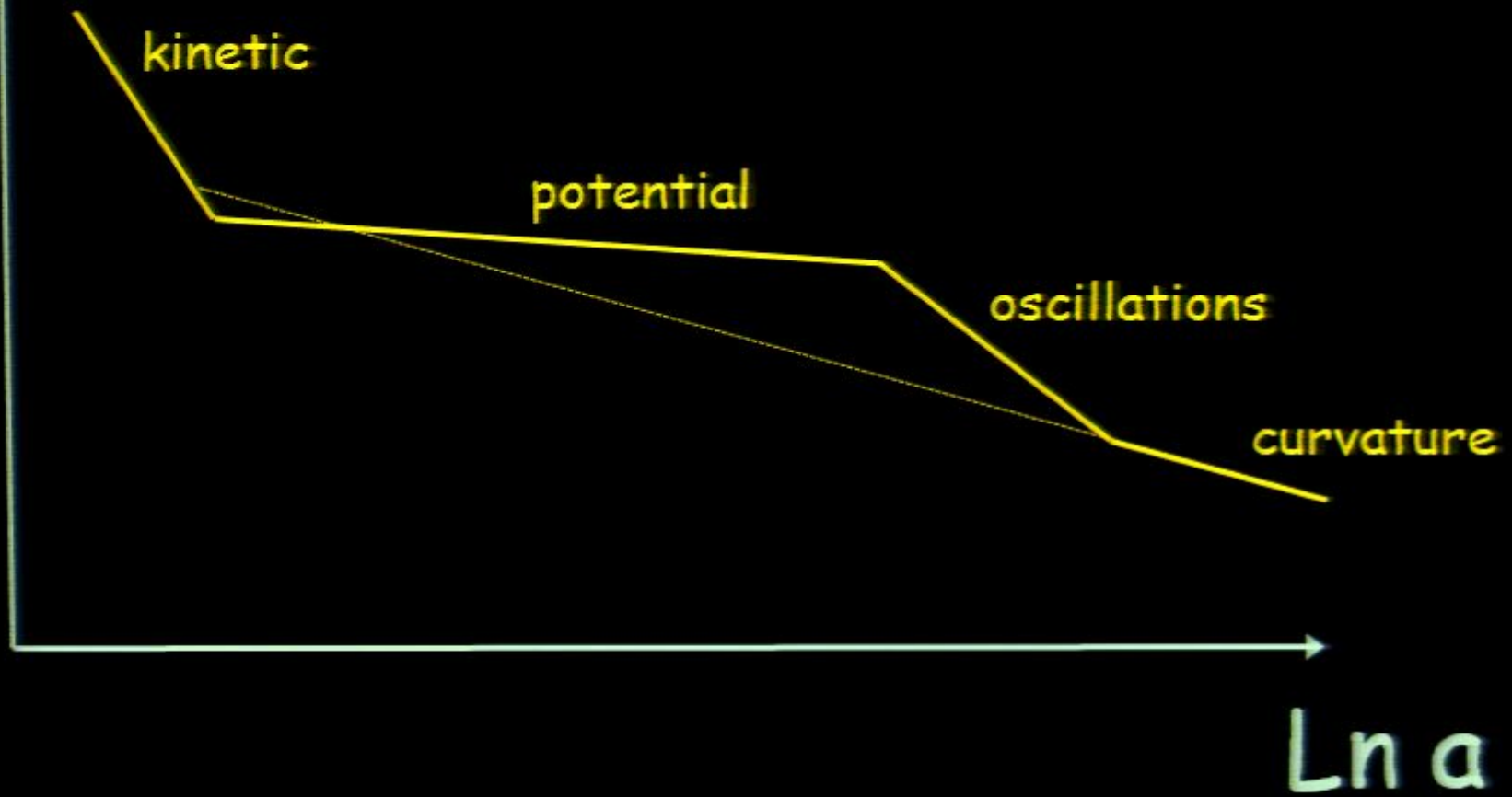
$\rightarrow$  entropy reduces to that of the matter (inc grav waves), and is an adiabatic invariant

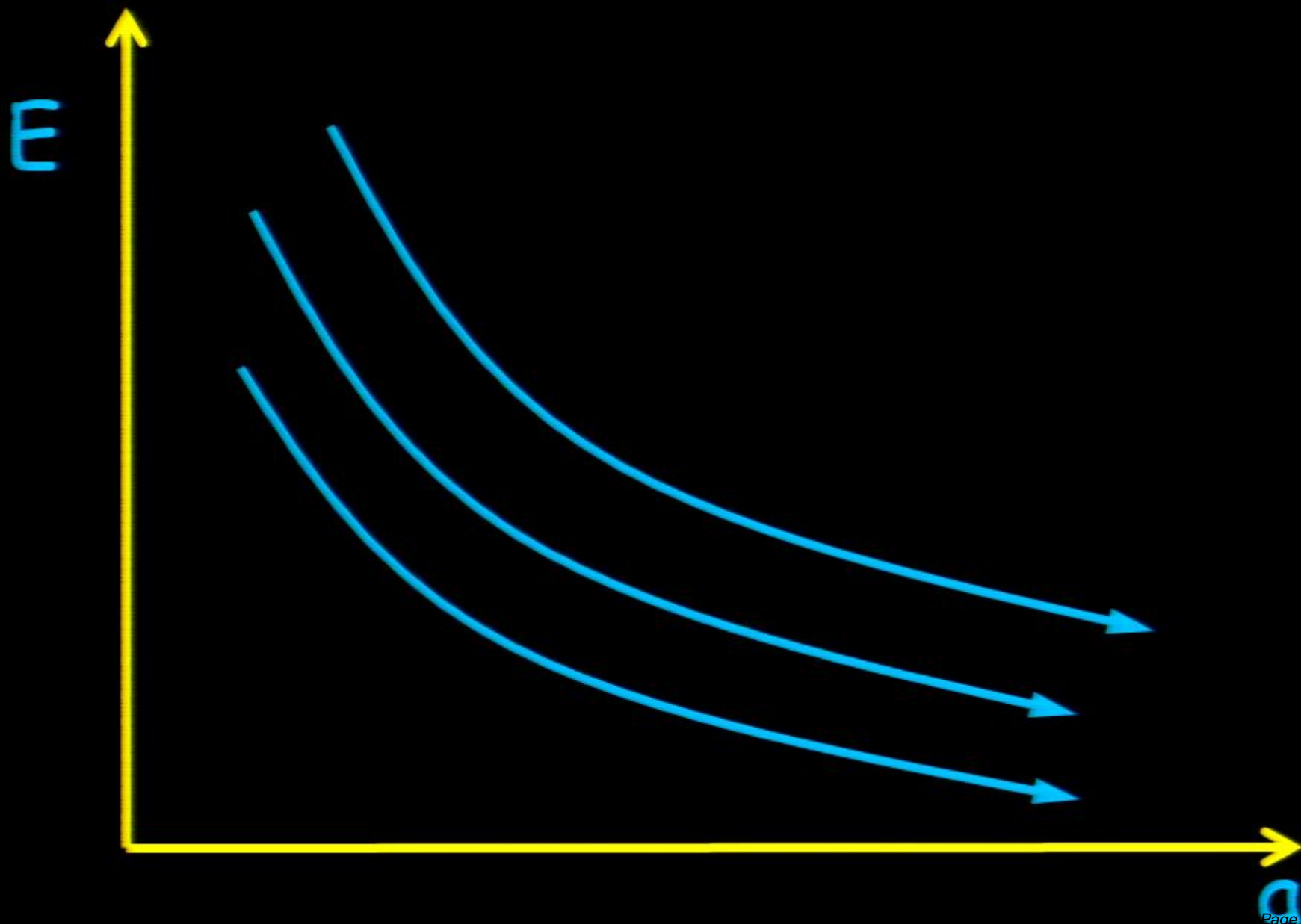
Every trajectory ends up on an **adiabat** curve  
 $S_m(E_m, a) = \text{const}$

Natural to label an ensemble of spacetimes by the **asymptotic entropy**  $S = S_m$

# generic open FRW cosmology

$\text{Ln } \rho$







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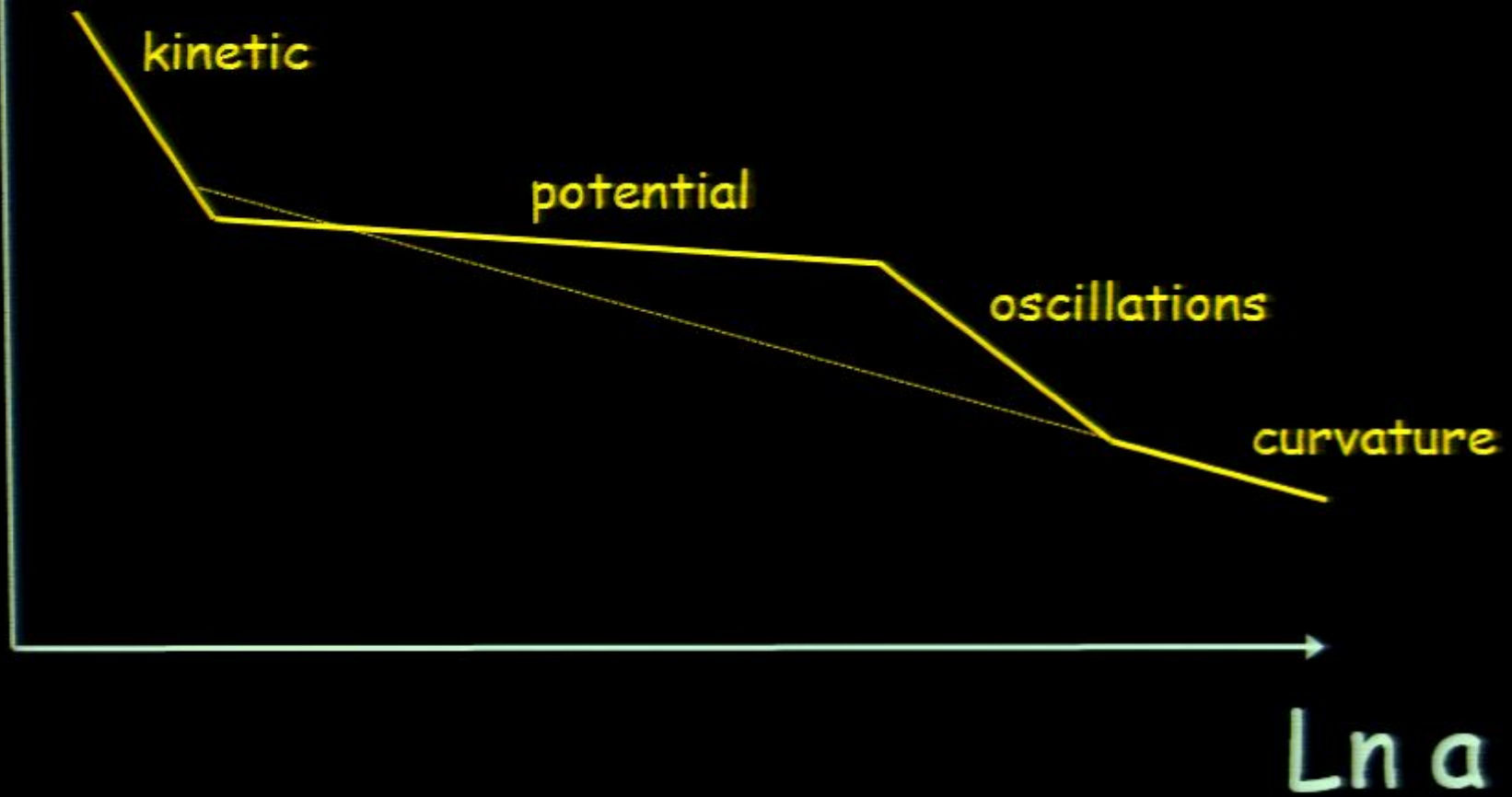
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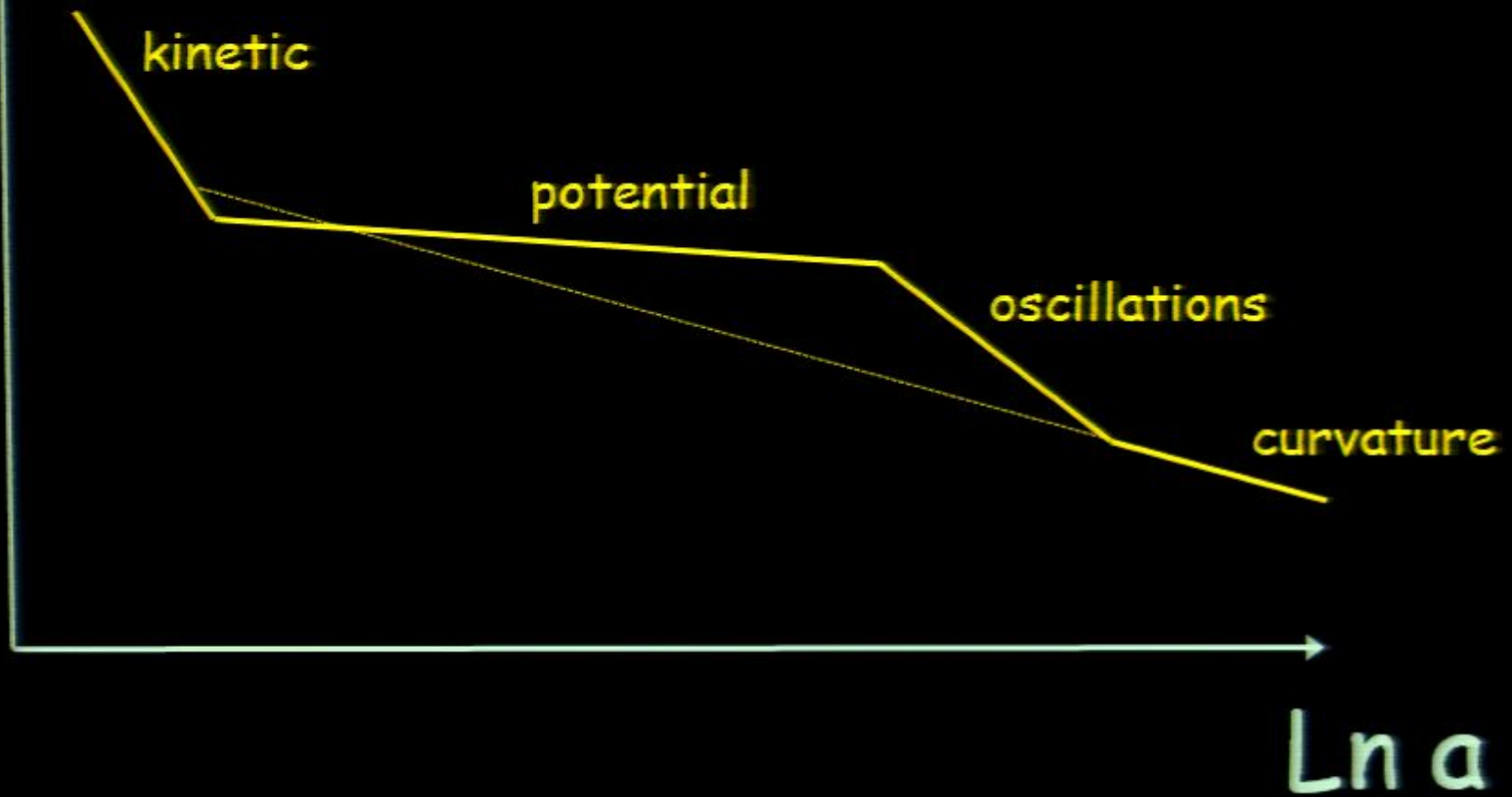
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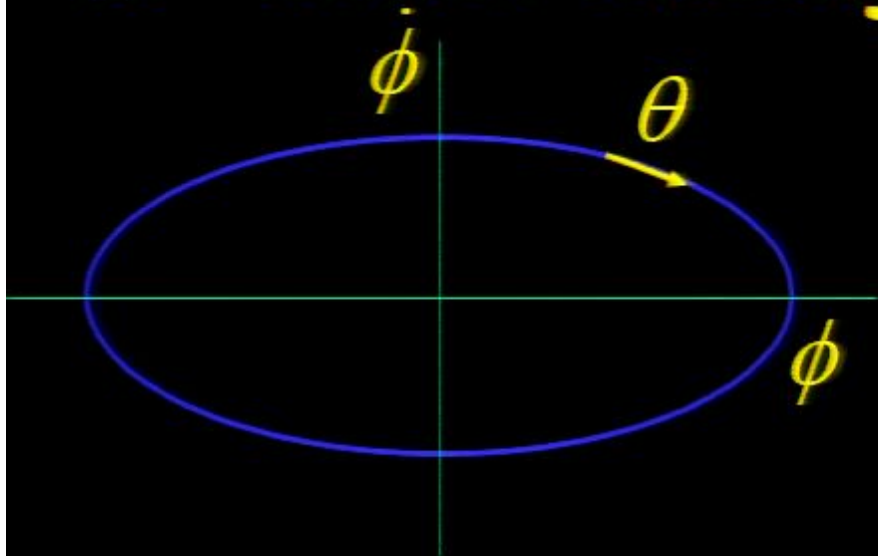
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$\text{Ln } \rho$



$a = \text{const}$  slicing

$$\omega_c = a^3 d\dot{\phi} d\phi$$



$$\rho = \frac{1}{2}(\dot{\phi}^2 + m^2 \phi^2)$$

$$\sim C/a^3, \quad a \rightarrow \infty,$$

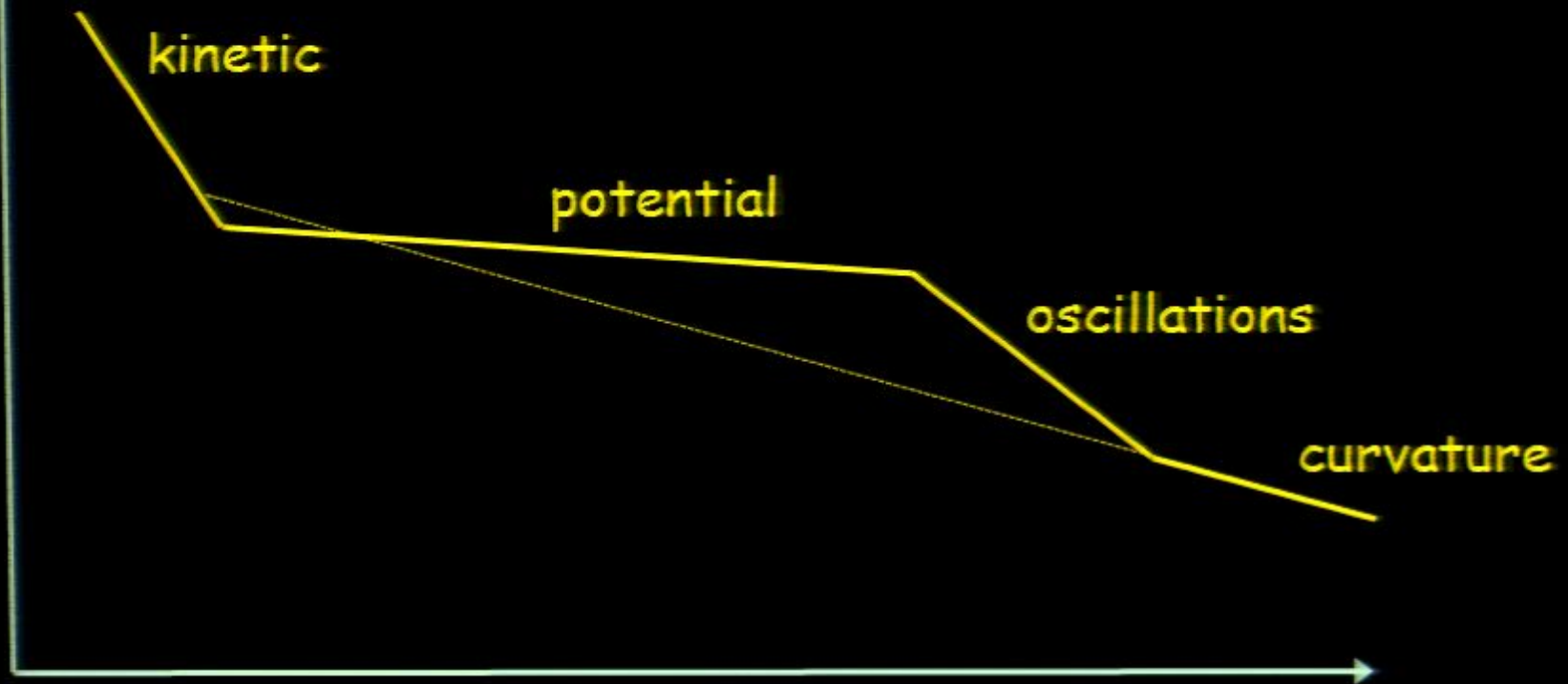
$\rho a^3$  adiabatically conserved

$$ds^2 = \frac{da^2}{1 + 8\pi G \rho a^2 / 3} + a^2 dH_3^2 \approx dM_4^2 - \frac{8\pi G}{3} \frac{C}{a} da^2$$

( $C$  is analogous to the AdM mass)

# generic open FRW cosmology

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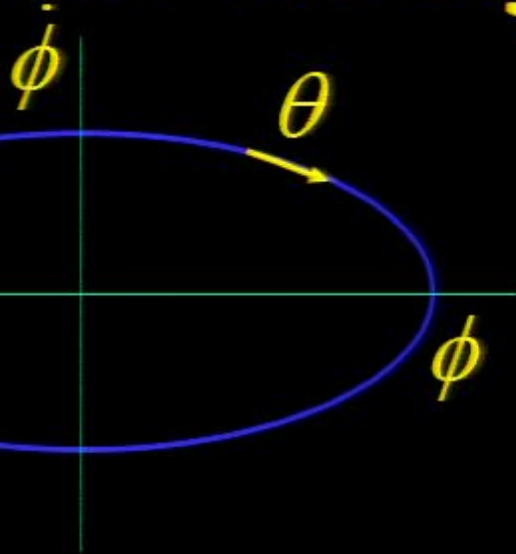


$\text{Ln } a$



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## Including backreaction

$$\begin{aligned}\phi &\sim \frac{(2C)^{\frac{1}{2}}}{ma^{\frac{3}{2}}} \left( \cos \theta + \frac{C \cos^3 \theta}{24a} \right) + \dots, \\ \rho_\phi &\sim \frac{C}{a^3} \left( 1 + \frac{3 \sin 2\theta}{2ma} + \frac{C(2 \cos 2\theta + \sin^2 2\theta)}{24a} \right) + \dots\end{aligned}$$

$$\text{where } \theta = m(a - \frac{5C}{24} \ln(ma)) + \theta_0$$

in large a limit, effect of matter  
on background spacetime (i.e. gravity)  
becomes negligible

we just have flat spacetime, and an  
adiabatically expanding box filled with  
matter



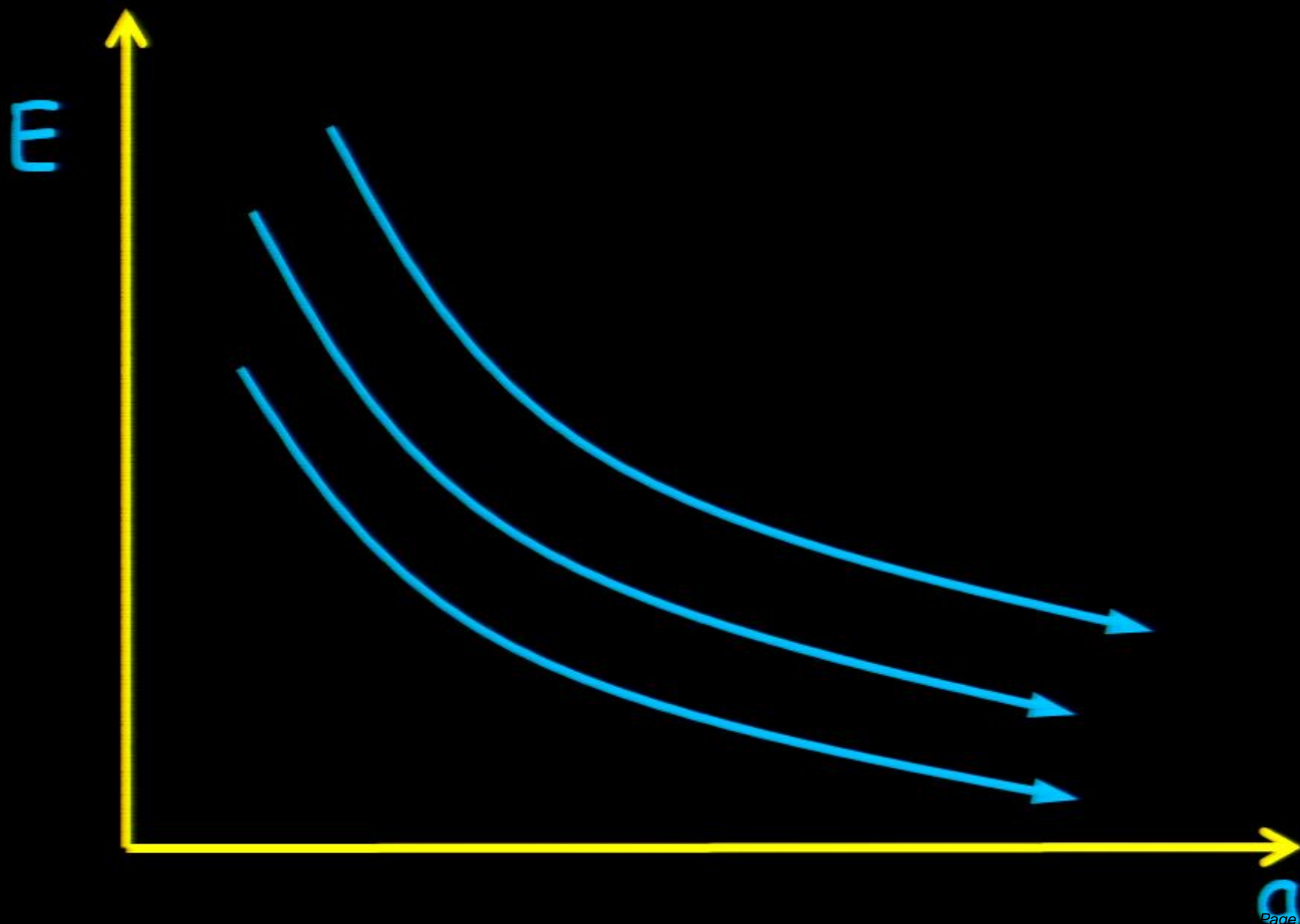
**statistical ensemble: minisuperspace**

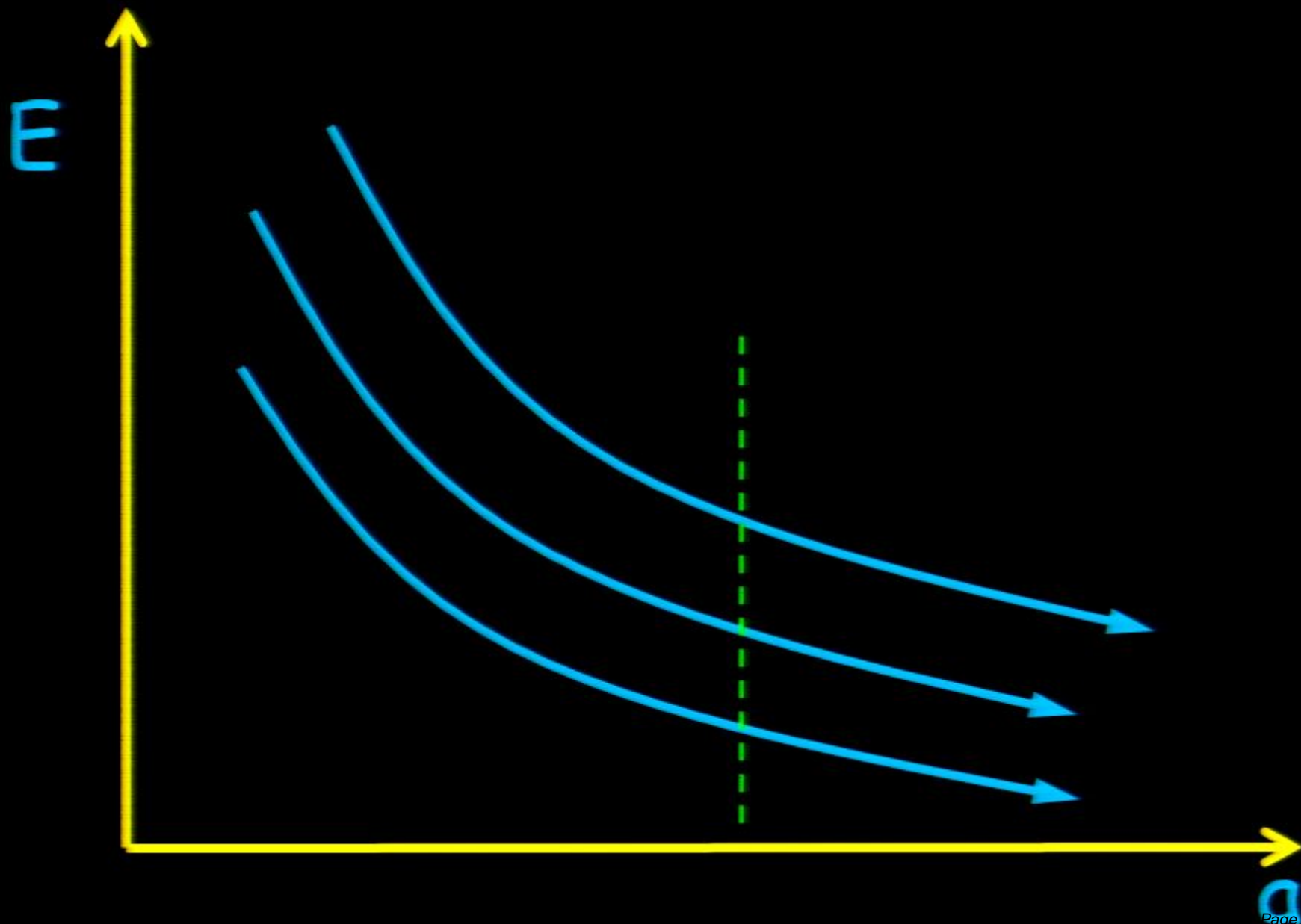
$$\mathbb{H}_m(p_\phi, \phi) = \frac{1}{2} \left( \frac{p_\phi^2}{Ua^3} + Ua^3 V(\phi) \right)$$

$$\langle \mathbb{H}_m \rangle = \frac{\int dp_\phi d\phi e^{-\beta \mathbb{H}_m} \mathbb{H}_m}{\int dp_\phi d\phi e^{-\beta \mathbb{H}_m}} = E(a, \beta)$$

**entropy**  $S = S_m = \ln\left(\frac{Ua^3 \rho_\phi}{m}\right) = \text{adiabatic invariant}$

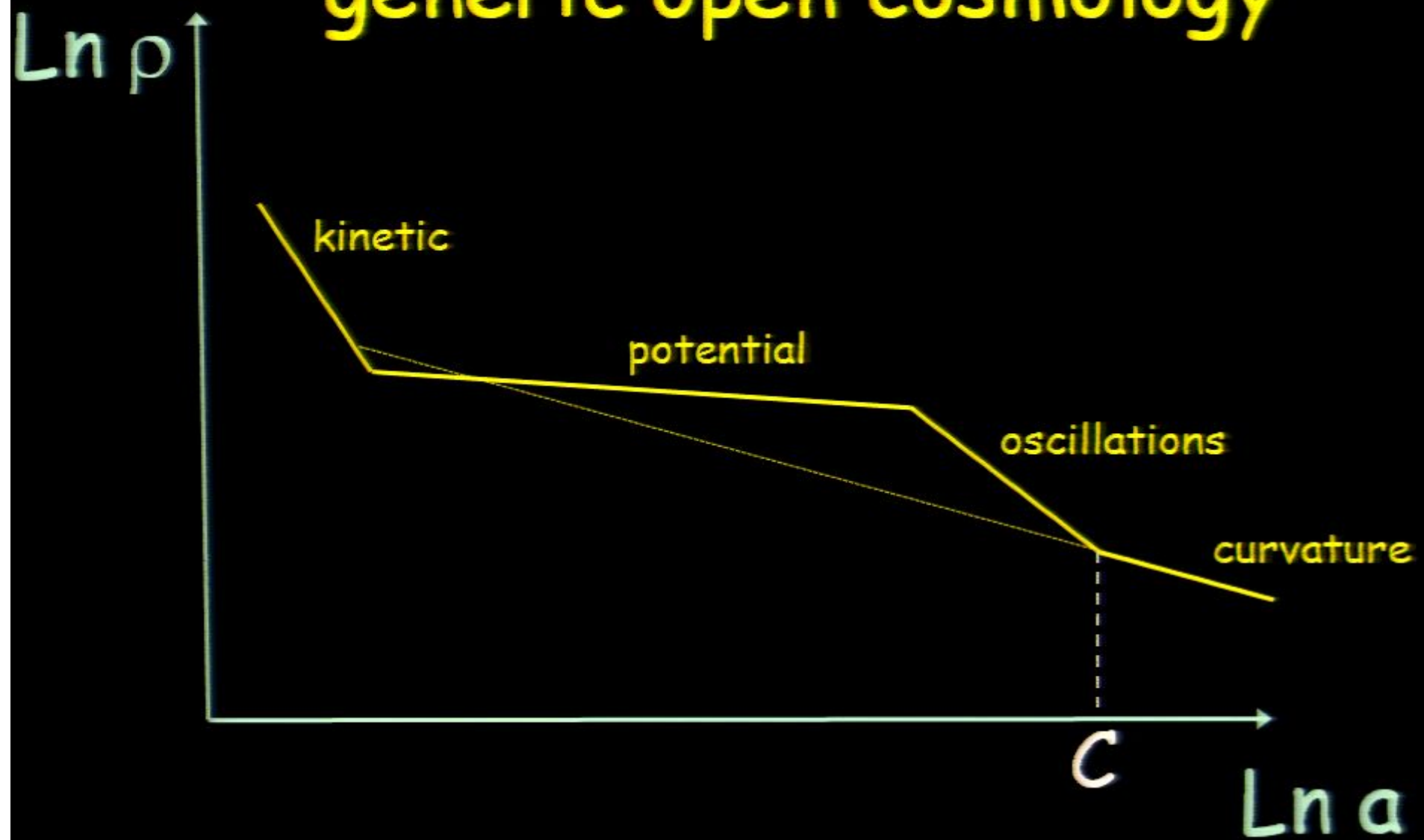
**constant entropy=fixed  $C$**



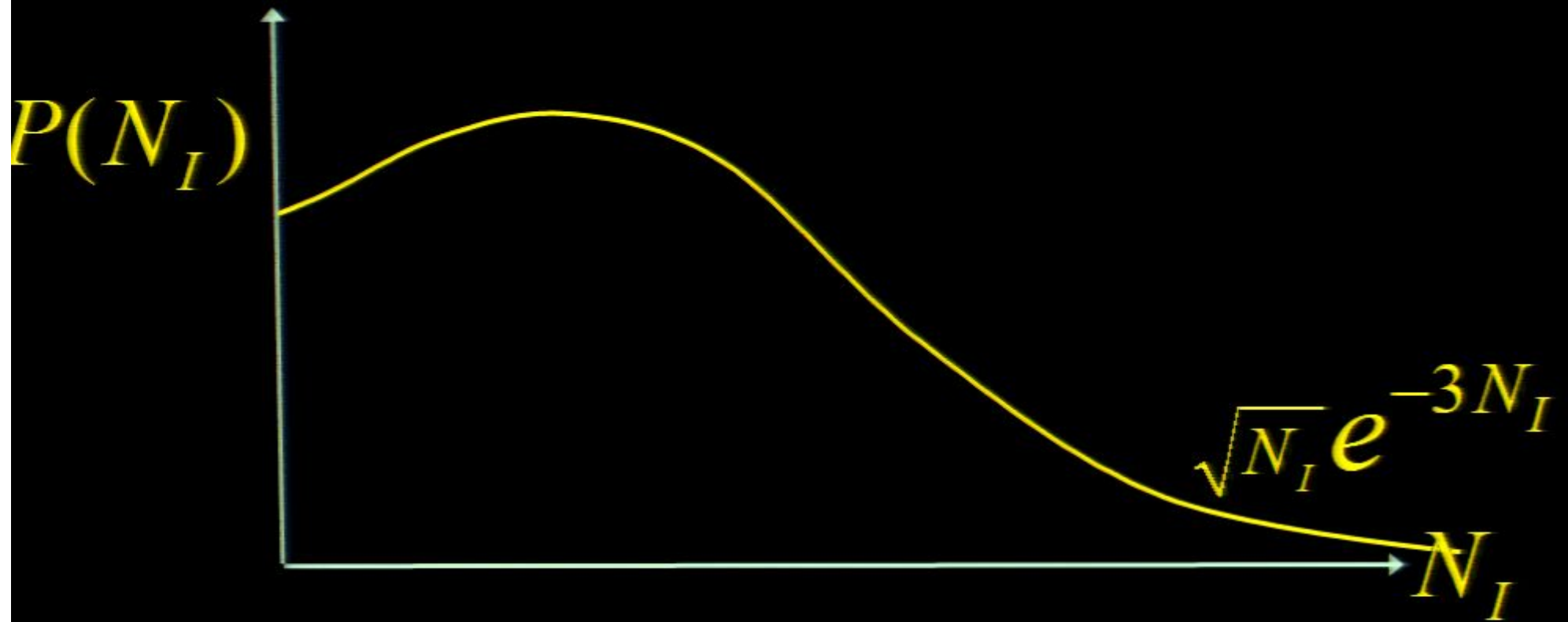




# generic open cosmology



# Canonical measure for inflation



Finite  $C$  result is always lower than  $C = \infty$  result at large  $N_I$

\* Note: "attractor" becomes "repeller" because statistical ensemble defined in asymptotic region where gravity becomes unimportant: the **future**

\* N slow-roll inflaton fields (N-flation) makes problem **worse**

$$P_>(N_I) \propto \prod_i \delta\theta_i \sim e^{-3N_I N}$$

\* this analysis makes precise a problem identified by Penrose long ago (Annals NYAS, 1989)

\* with this canonical measure, slow-roll/'chaotic' inflation **cannot** be considered a viable explanation for the observed state of the cosmos

\* can be extended to landscape with inclusion of complex solutions describing tunneling: false vacua don't help

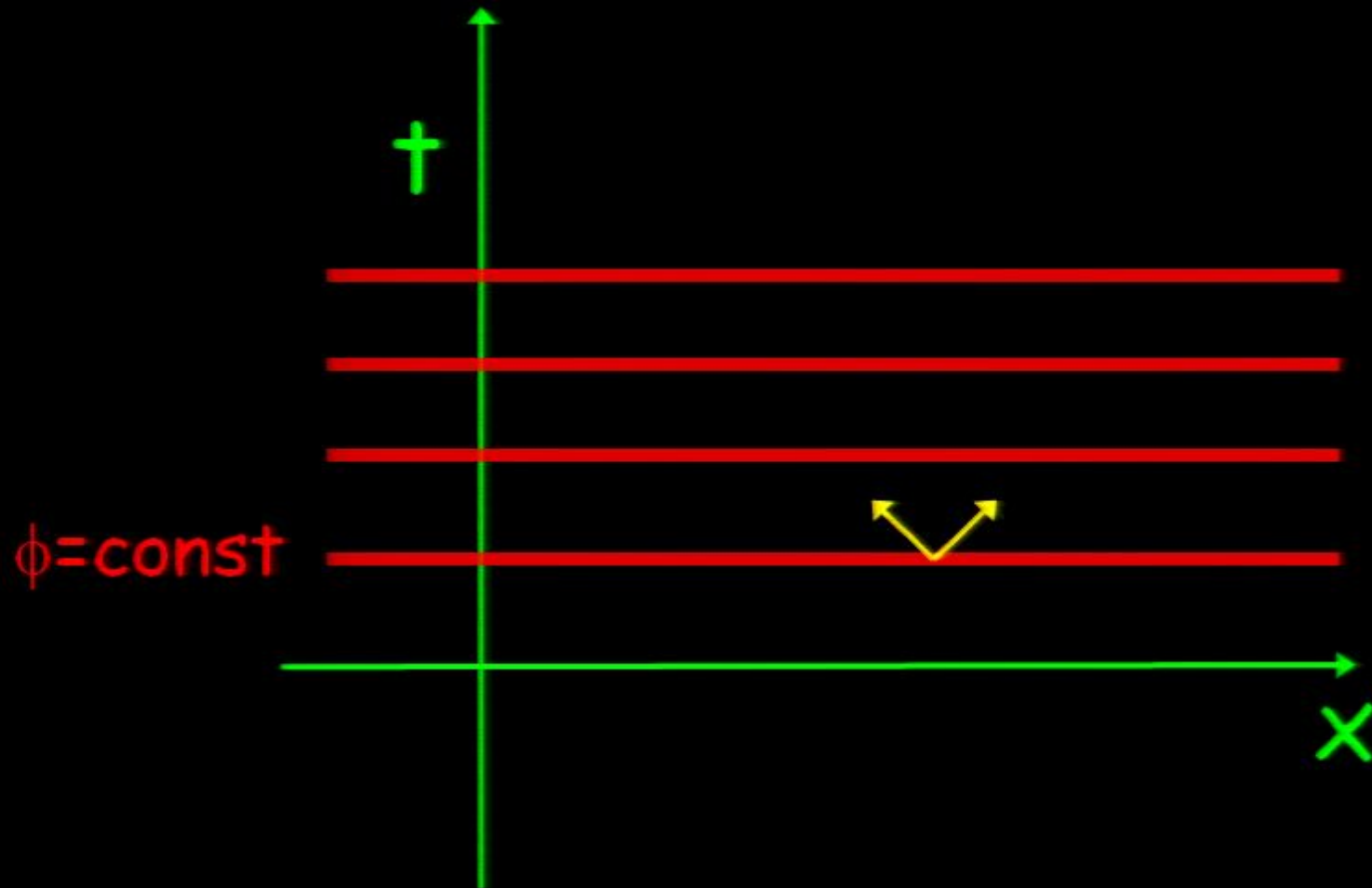


- \* Could add further restrictions  
e.g. insist  $\phi$  bounded  $\rightarrow$  1-parameter family  $\phi_0(S_{\text{final}})$
- \* Large entropy  $\rightarrow$  many e-folds, but large final entropy (or flatness) is then an input, not a prediction
- \* Could restrict to non-singular spacetimes (eg Page) and just reject all  $k=-1$  solutions
- \*  $k=+1$  "bounce" solutions possible, have finite measure but collapsing phase very unstable
- \* How do we justify rejecting cosmic singularities but allowing black holes? What about white holes?

# What could be wrong?

- \* canonical measure?
- \* neglect of: entropy production?

# $\phi$ decays: information stored in radiation



But the relative proportions of phase space corresponding to  $N_I$  or more efolds are preserved (by unitarity or Liouville)



## What could be wrong?

- \* canonical measure?
- \* neglect of: entropy production?
  - : inhomogeneities?
  - : quantum fluctuations?
- \* constraint on final entropy?  
maybe, but what alternative?
- \* global structure?
- \* inflation?

compare "cyclic/ekpyrotic" theory, where gravity is unimportant in the **past**, and according to the corresponding canonical measure, **every** trajectory undergoes near-maximal ekpyrosis

(w/P. Steinhardt )

## Scale Invariant perturbations

- \* near-massless scalar in de Sitter
- \* exponentially flat potentials in collapsing phase (cyclic)
- \* scale-invariant duals via holography

## Summary

concordance cosmology has plenty of challenges

- singularity
- tuning
- reliance on anthropics
- measure: a good one exists!



No Signal

VGA-1

No Signal

VGA-1