

Title: Overview of the Challenges

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URL: <http://pirsa.org/11070044>

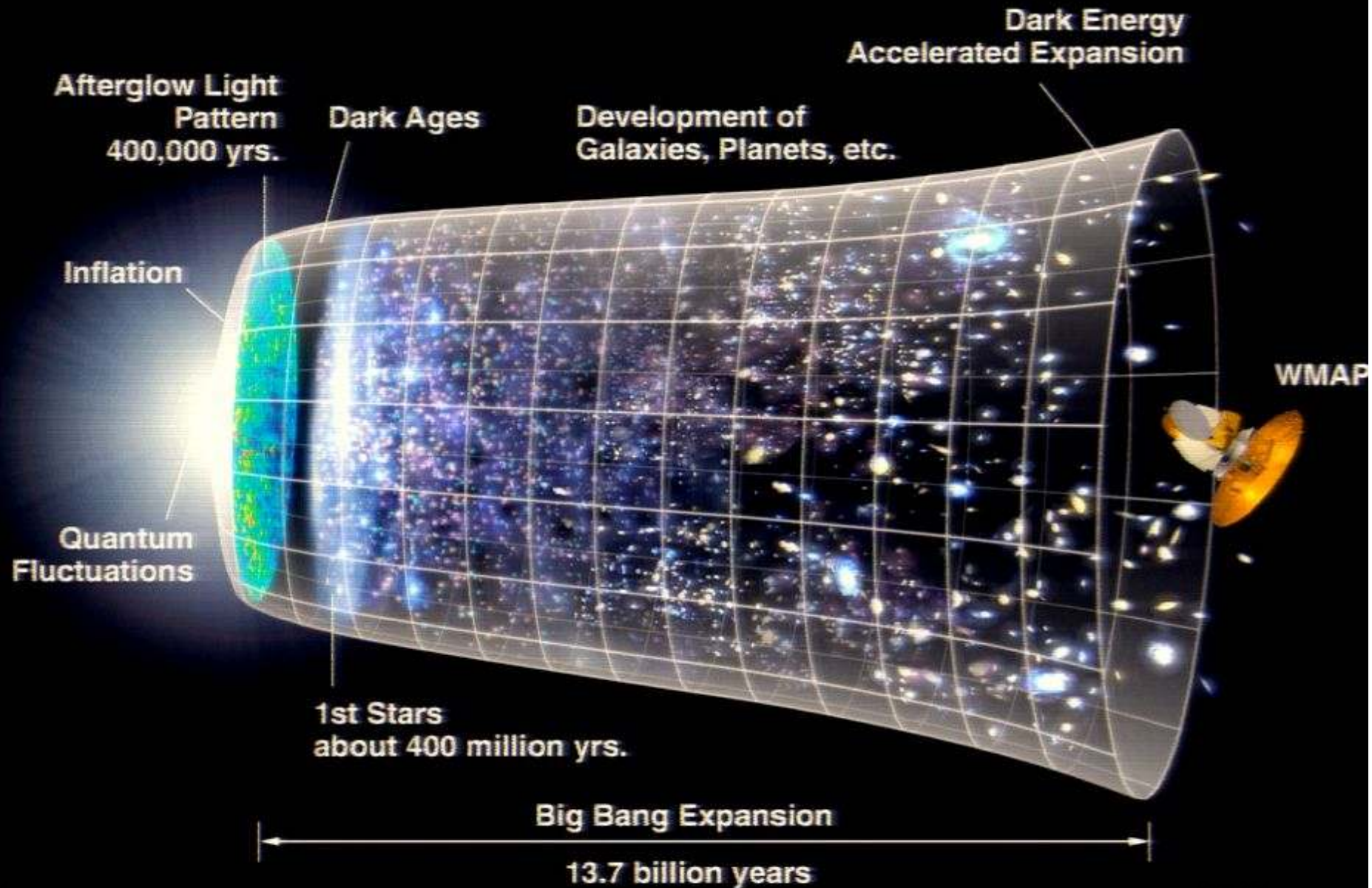
Abstract:

Challenges for Early Universe Cosmology

Neil Turok, Perimeter Institute

- * Concordance Model
- * Anthropics
- * Measures

Concordance Model



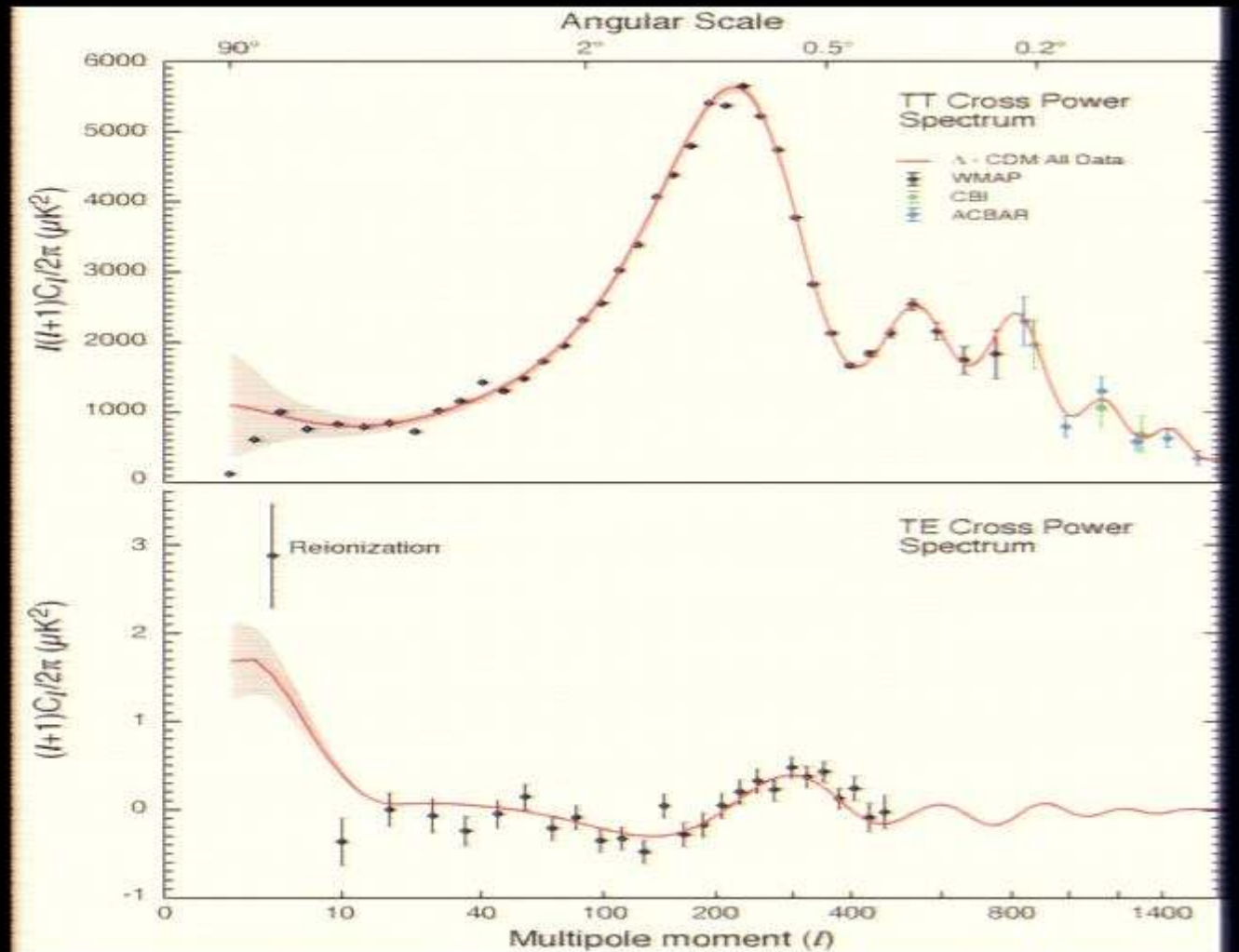
Success!

fluctuation level:
temperature

Zeldovich, Peebles+Yu 70's
Bond+Efstathiou 80's

polarization

Coulson, Crittenden, NT 94



100

10

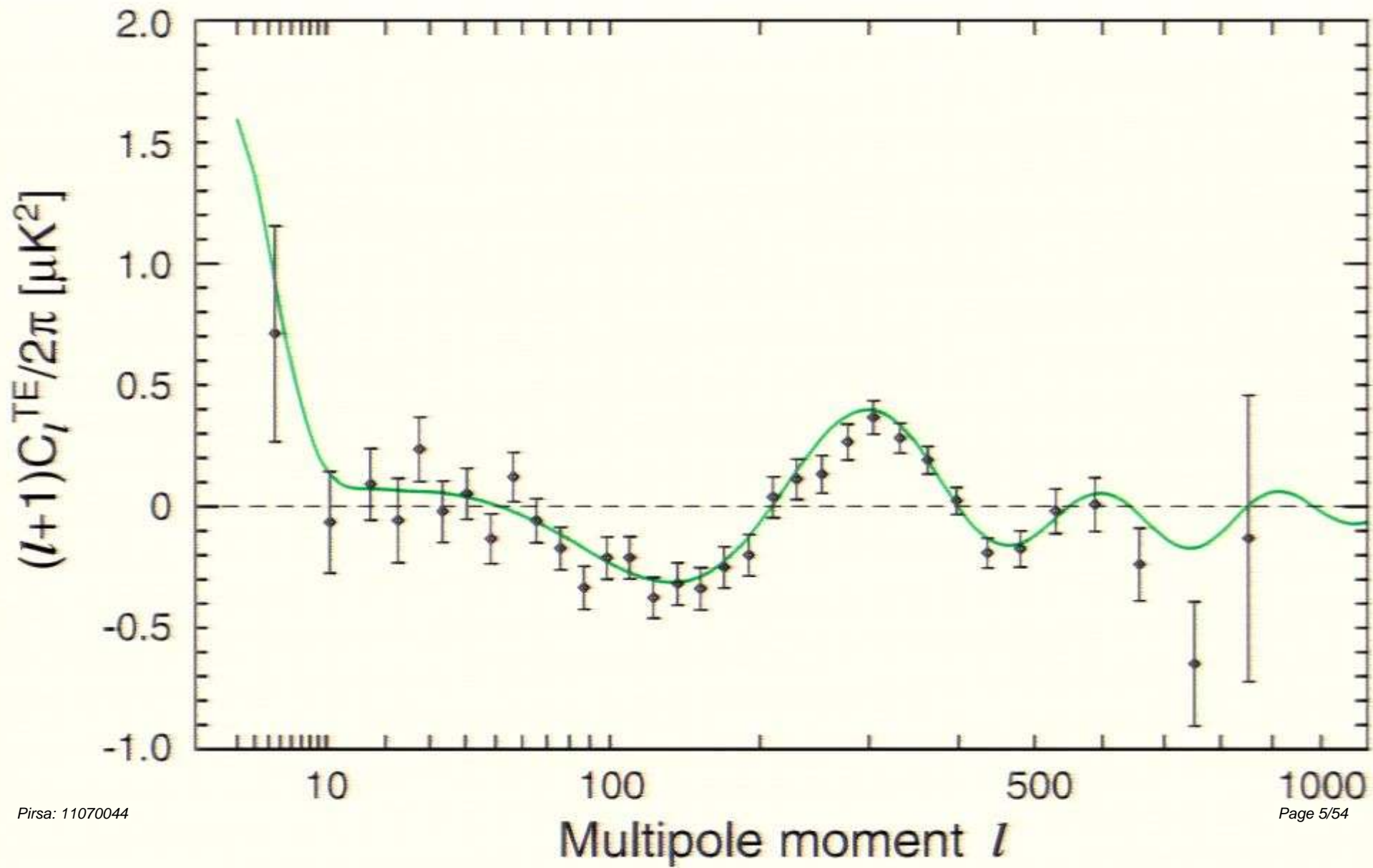
1

.1

($l = 2\pi/\theta$)

Angle on Sky (Degrees)

WMAP 5 year TE



good evidence for ...

nearly flat FRW universe:

$$\Omega_{\Lambda} : \Omega_{CDM} : \Omega_B : \Omega_v : \Omega_{\gamma} \sim 0.7 : 0.25 : 0.05 : 0.003 : 0.0003$$

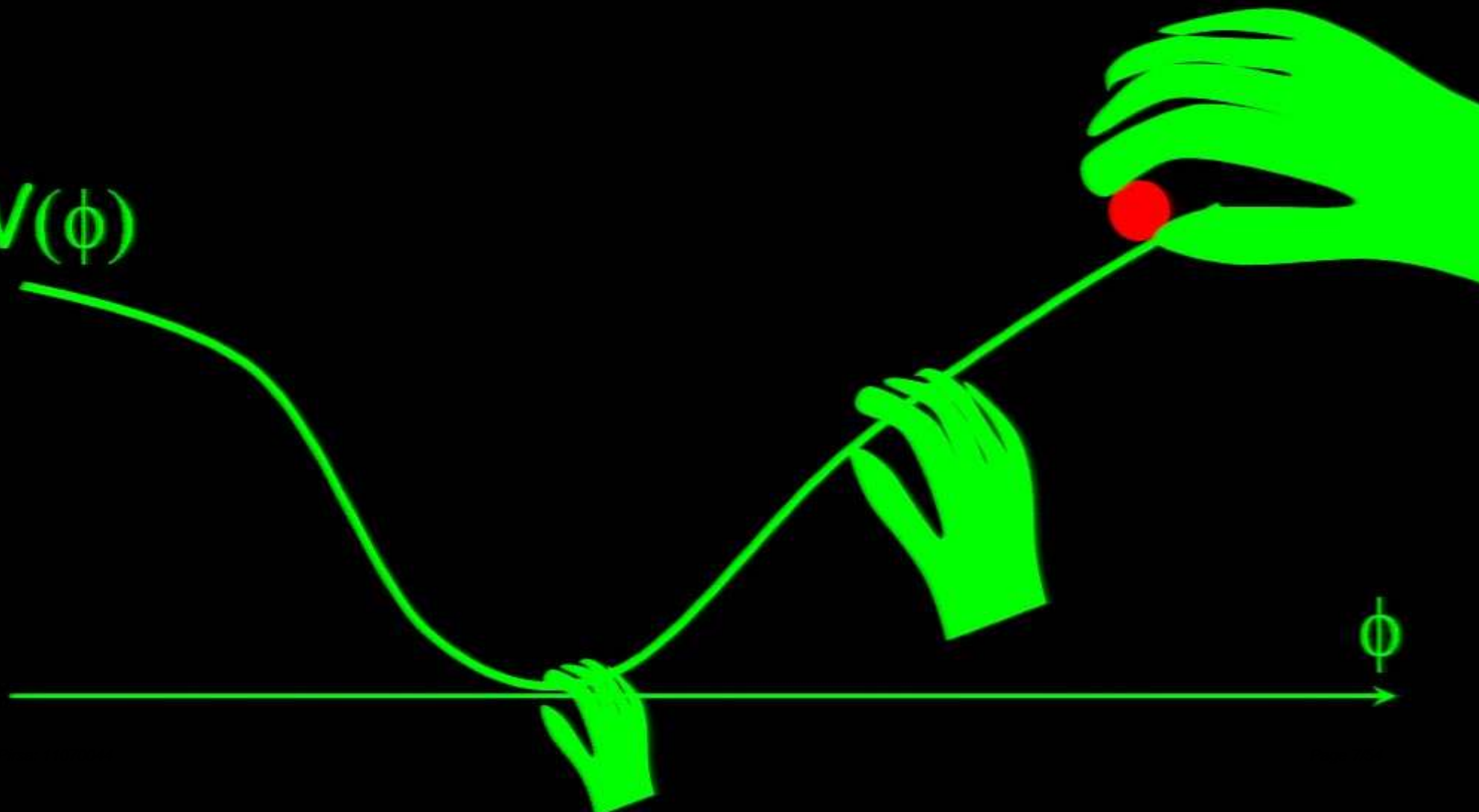
primordial perturbations

- * linear
- * growing mode
- * nearly scale-invariant
- * nearly 'adiabatic'
- * nearly Gaussian

universe is: geometrically simple
compositionally complex

inflation

$V(\phi)$



Challenges for Inflation

- * initial conditions
- * fine-tuned potentials
- * $\Lambda \sim 10^{-120}$; $\Lambda_{\text{I}} \sim 10^{-15}$
- * eternal inflation
 - "anything that can happen will happen: and it will happen an infinite number of times" A. Guth, 2002
- * landscape
 - measure problem
- * reliance on anthropic arguments

Anthropics: the universe is the way it is because of a (gigantic) selection effect

why are we then not "close to extinction"?

why is the universe **so** geometrically simple?

- * very nearly flat
- * very nearly spatially homogeneous
- * the fluctuations are scale-invariant on scales much larger than galaxy scales

for anthropic arguments to become convincing, we need well-defined

- * measures
- * projections

The Measure Problem

What is the natural measure on the space of cosmological solutions?

What is the likelihood of a universe like ours, in a given physical model? eg inflation, cyclic,

Two key ingredients

I: Penrose critique of inflation - Hamiltonian evolution almost never turns a generic state into a highly unusual state. Canonical measure is invariant.

II: Counting of states in gravity can only be done in an asymptotic region where global properties of spacetime become sharp

In a specific setup, we shall construct a canonical measure and show a universe like ours is extremely unlikely in slow-roll inflationary models

- standard slow-roll inflation

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{1}{3} \rho_\phi - \frac{k}{a^2}; \quad \rho_\phi = \frac{1}{2} \dot{\phi}^2 + V(\phi)$$

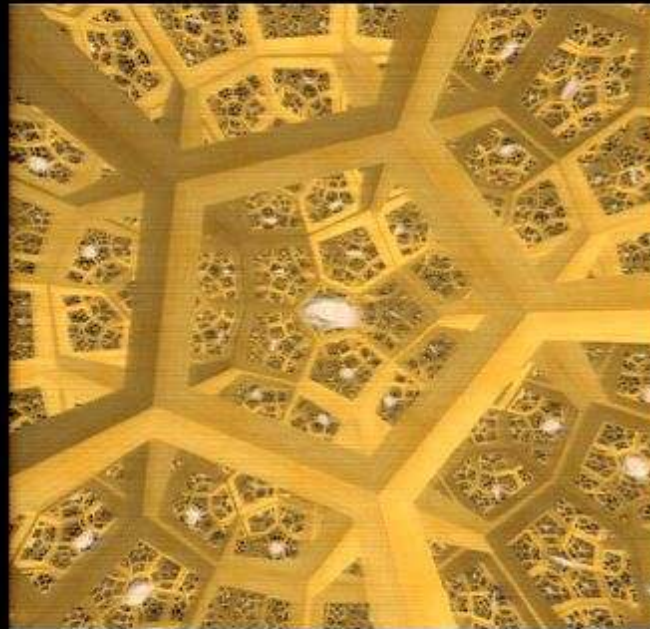
$$\ddot{\phi} + 3 \frac{\dot{a}}{a} \dot{\phi} = -V_{,\phi}$$

$$\Rightarrow \dot{H} = -\frac{1}{2} \dot{\phi}^2 + \frac{k}{a^2}$$

- Hamiltonian and time reversal invariant

(focus on FRW spacetimes and assume $V(\phi)$ is monotonic away from its min)

Assume $k=-1$ (so a and H are monotonic),
zero Λ , and compactify the spatial slices



- * a mathematical device to keep everything finite: the results do not depend on the compactification volume
- * (but in fact has been advocated as a very natural setup for chaotic inflation)

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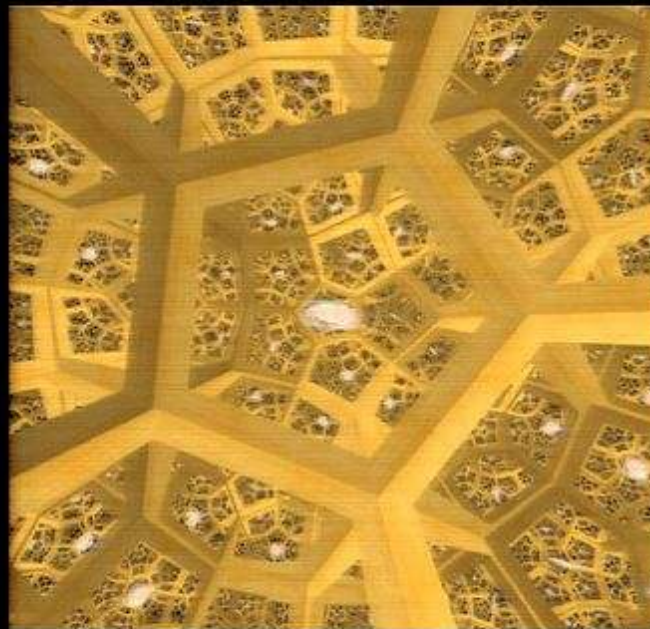
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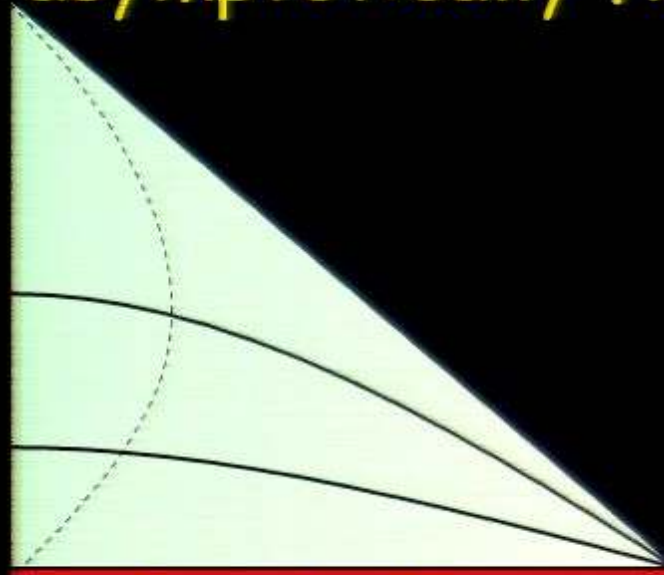
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2-parameter family of solutions with an initial singularity:

asymptotically flat



singularity

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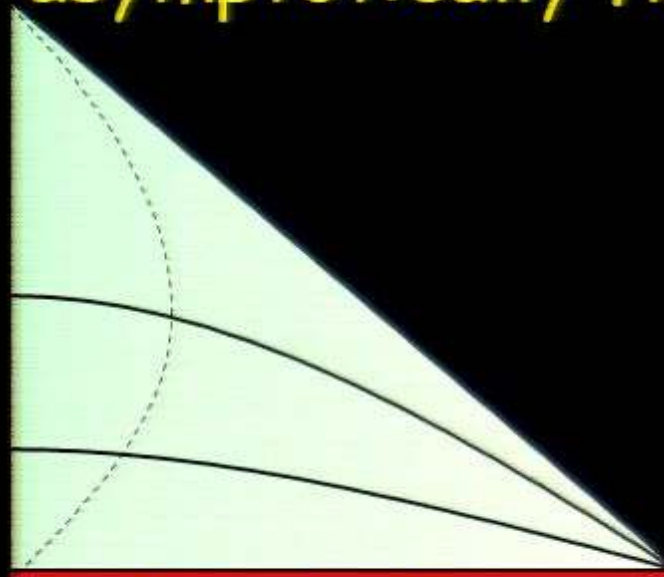
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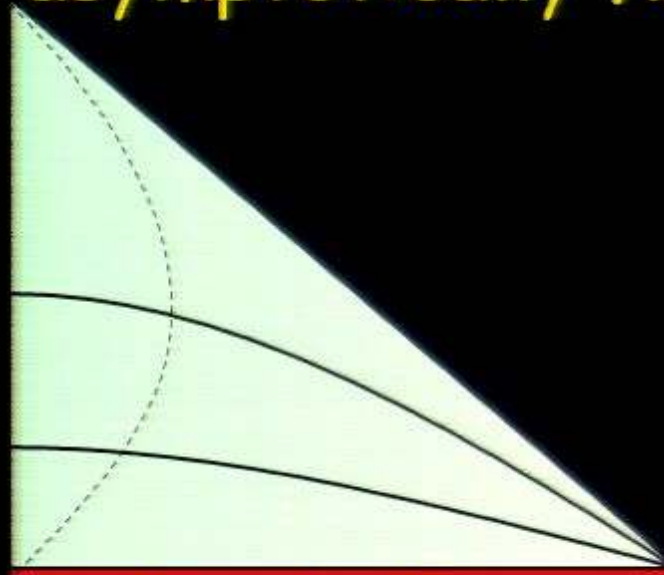
singularity

Criteria for a Measure

- (i) Positive, normalisable
- (ii) Independent of slicing or coordinates on either space-time or field space
- (iii) Independent of ad hoc external structures eg cutoffs, comoving observers, "volume" factors .
- (iv) Natural extension of canonical quantum measure for fluctuations to the background (why use for one but not the other???)

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Canonical measure on space of solutions

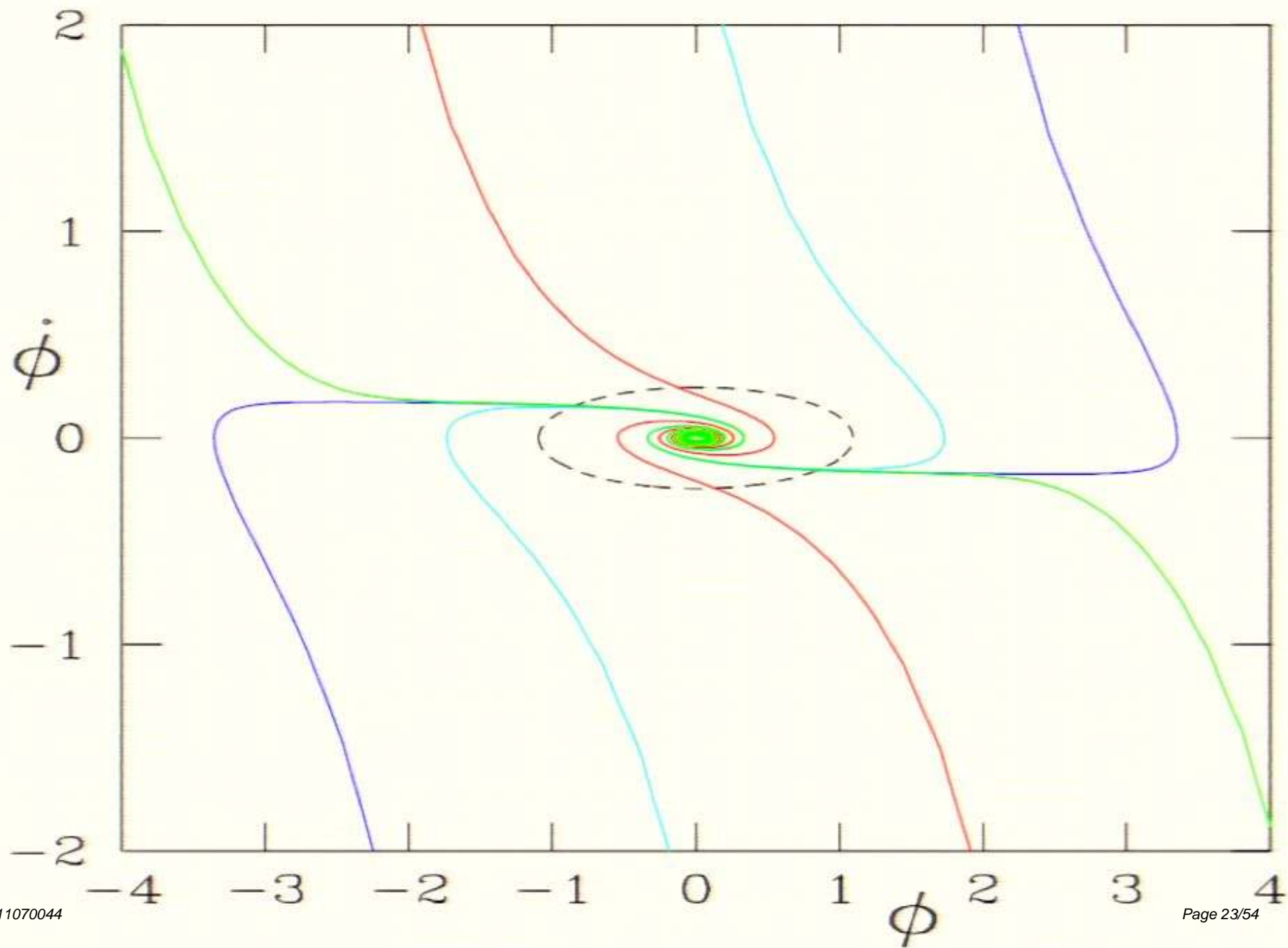
$$\omega_c = dp_a \wedge da + dp_\phi \wedge d\phi$$

$$\int_{\Sigma} \omega_c \Big|_{\mathbb{H}=0}$$

Liouville
Gibbons, Hawking, Stewart
Hawking, Page
Hollands Wald
Kofman, Linde, Mukhanov
Gibbons, NT
Carroll, Tan

with Σ pierced once by every trajectory
e.g. $a=\text{const}$ or $H=\text{const}$

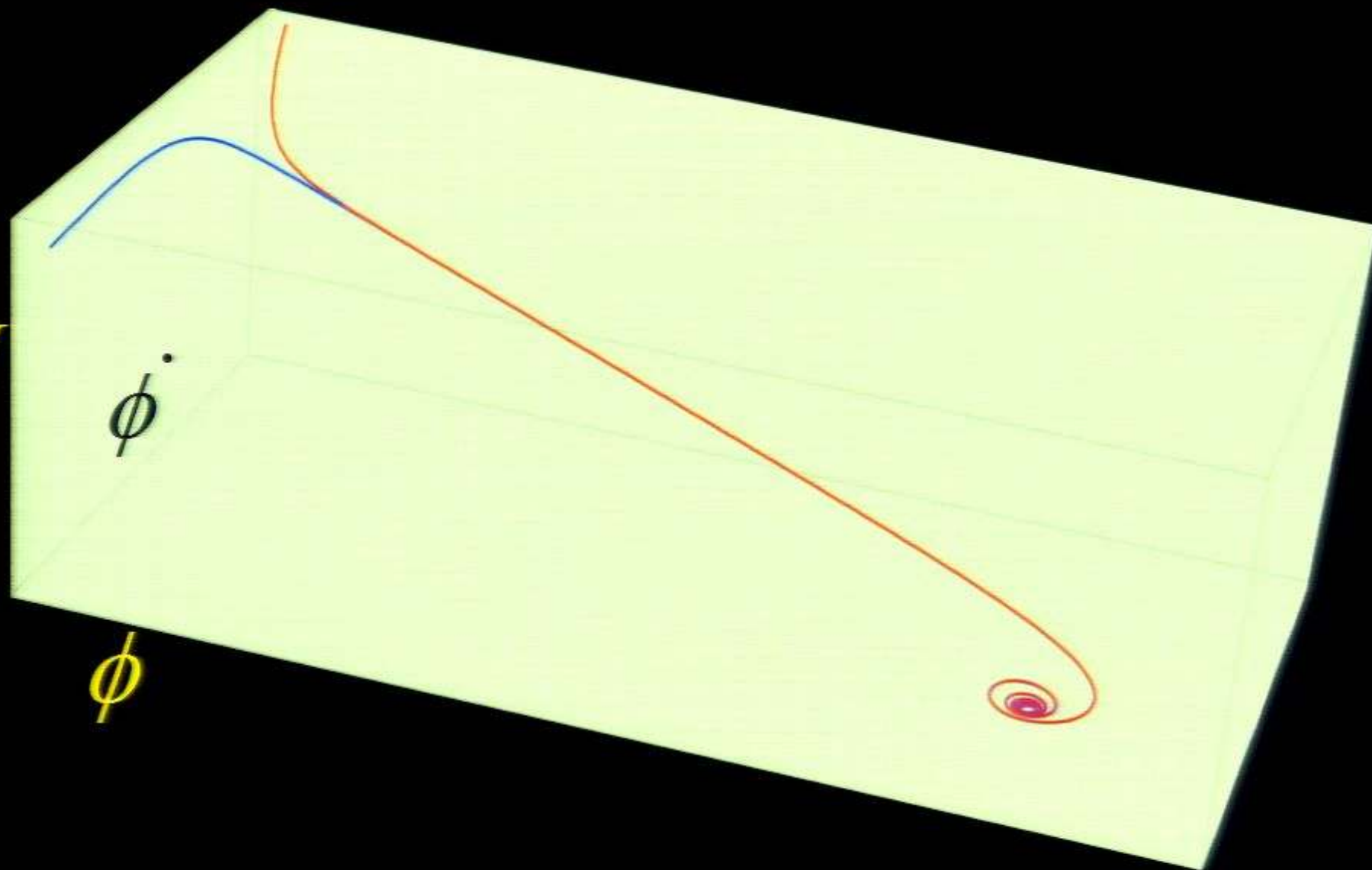
Satisfies all of conditions (i)-(iv)
except normalisability, because Σ is not
compact (because \mathbb{H} isn't positive)

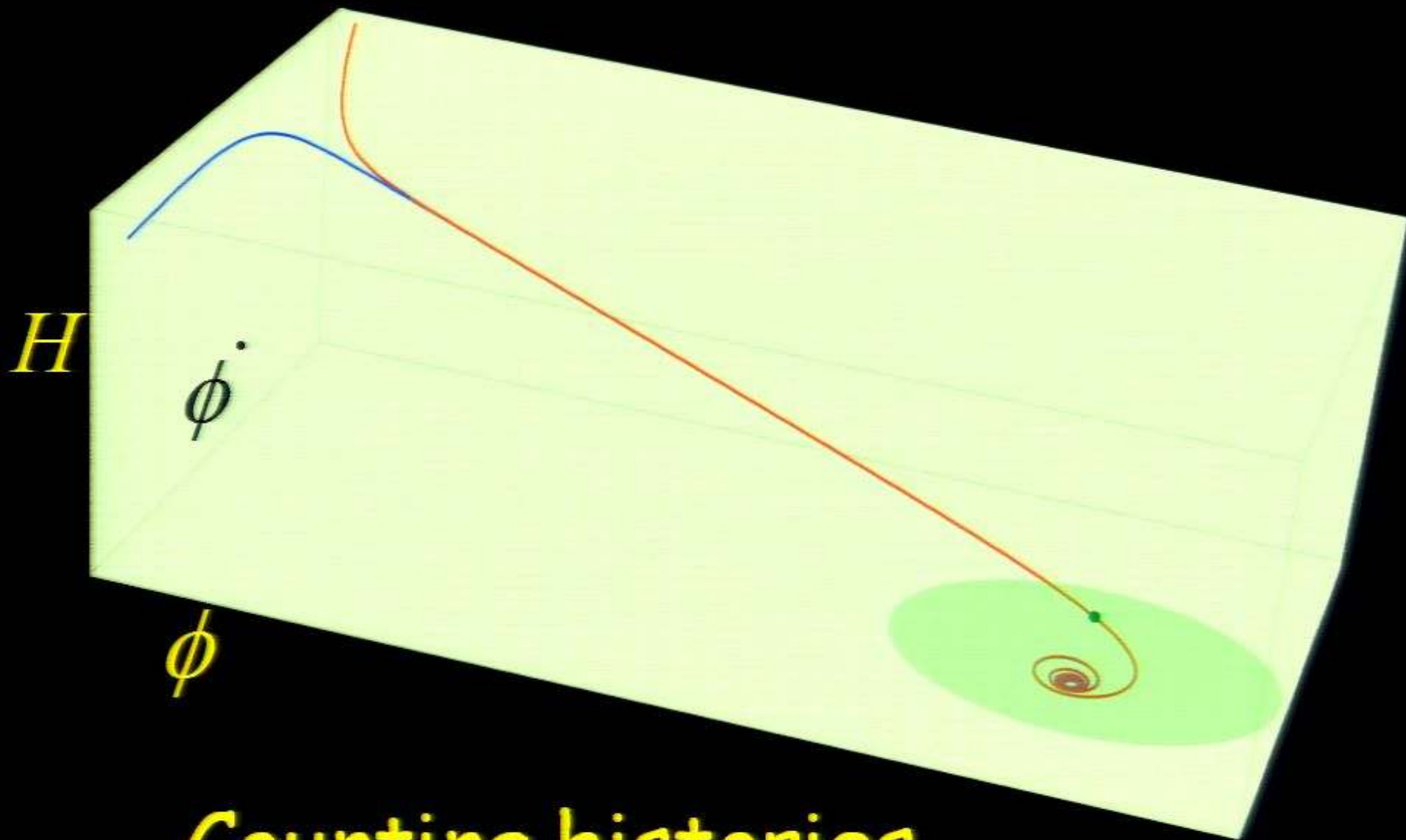


H

ϕ

ϕ





Counting histories

Flat space canonical (Gibbs) ensemble.

Cannot just integrate over Liouville: instead, we maximise entropy $S = -\sum p_i \ln p_i$

subject to $E = \sum p_i E_i$ (assuming E_i b.b.)

Note: in information theory approach, max ent principle is very general, can even be applied to non-equilibrium situations (see e.g. papers of E.T. Jaynes)

But in GR, $\mathbb{H}=0$ on all physical states, so cannot constrain its expectation value

What do we do?

In $k=-1$, zero Λ cosmologies, matter density is diluted away at large a \rightarrow gravity becomes negligible, expansion becomes adiabatic

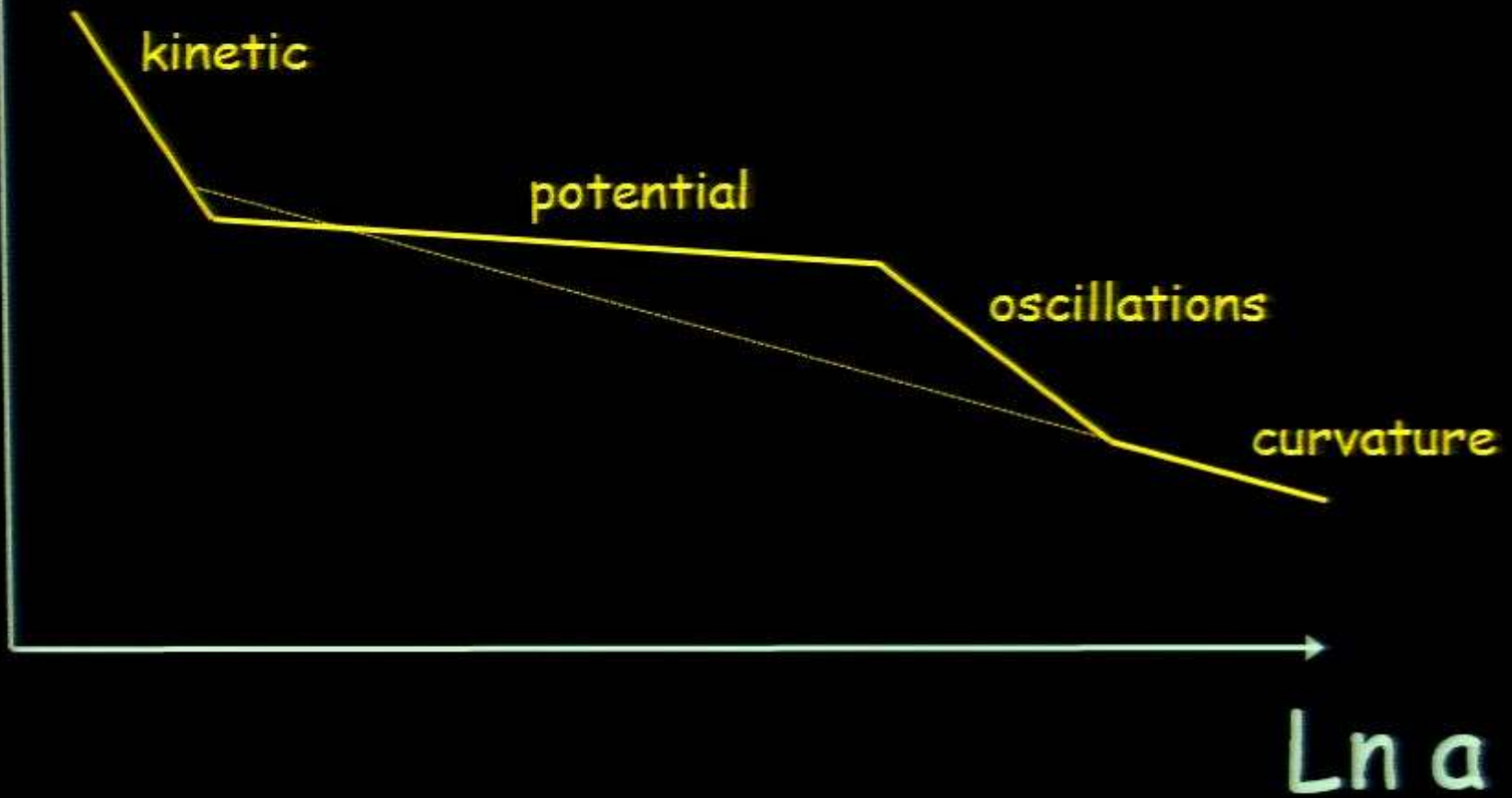
\rightarrow entropy reduces to that of the matter (inc grav waves), and is an adiabatic invariant

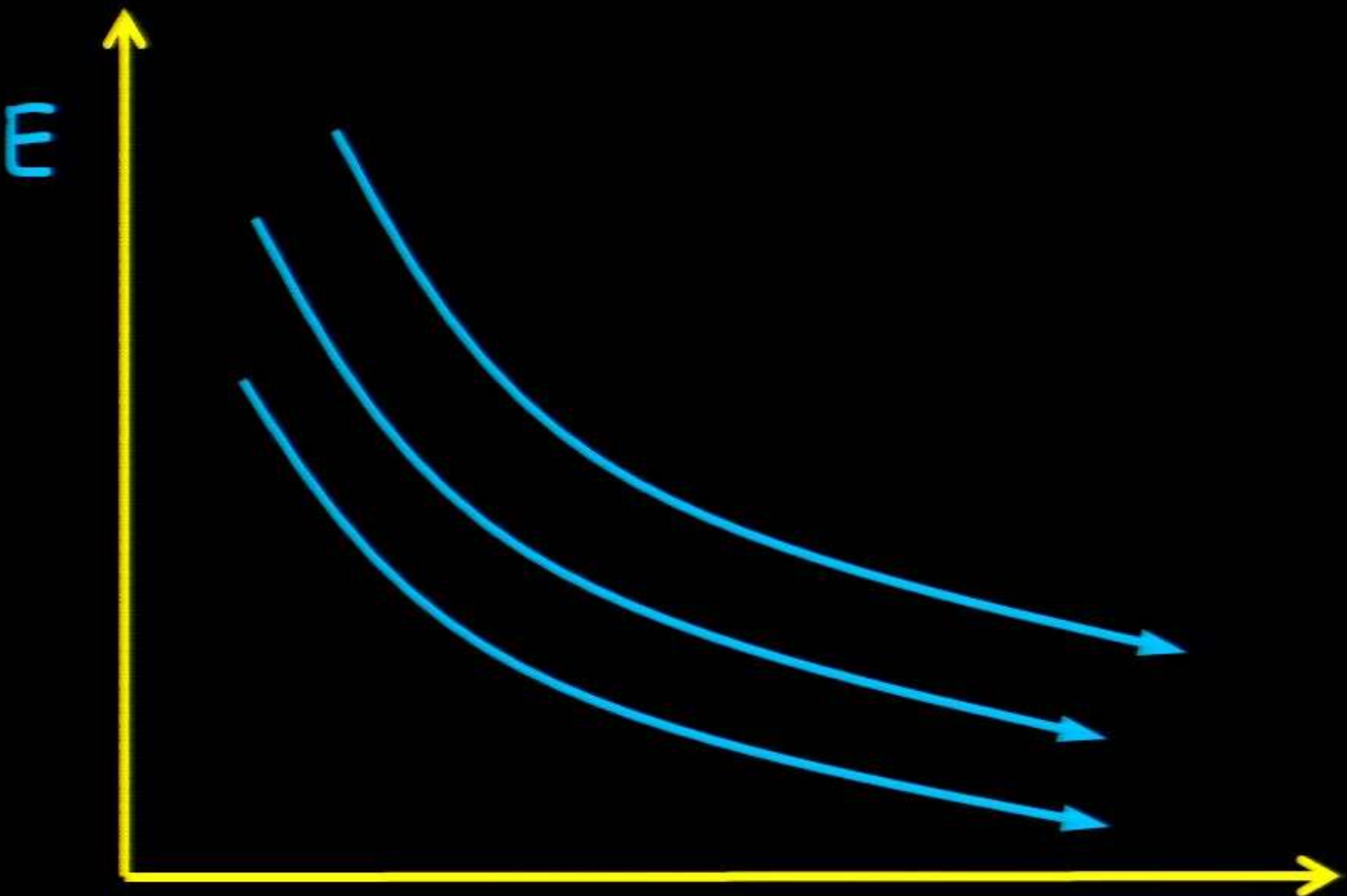
Every trajectory ends up on an **adiabat** curve
 $S_m(E_m, a) = \text{const}$

Natural to label an ensemble of spacetimes by the **asymptotic entropy** $S = S_m$

generic open FRW cosmology

$\text{Ln } \rho$





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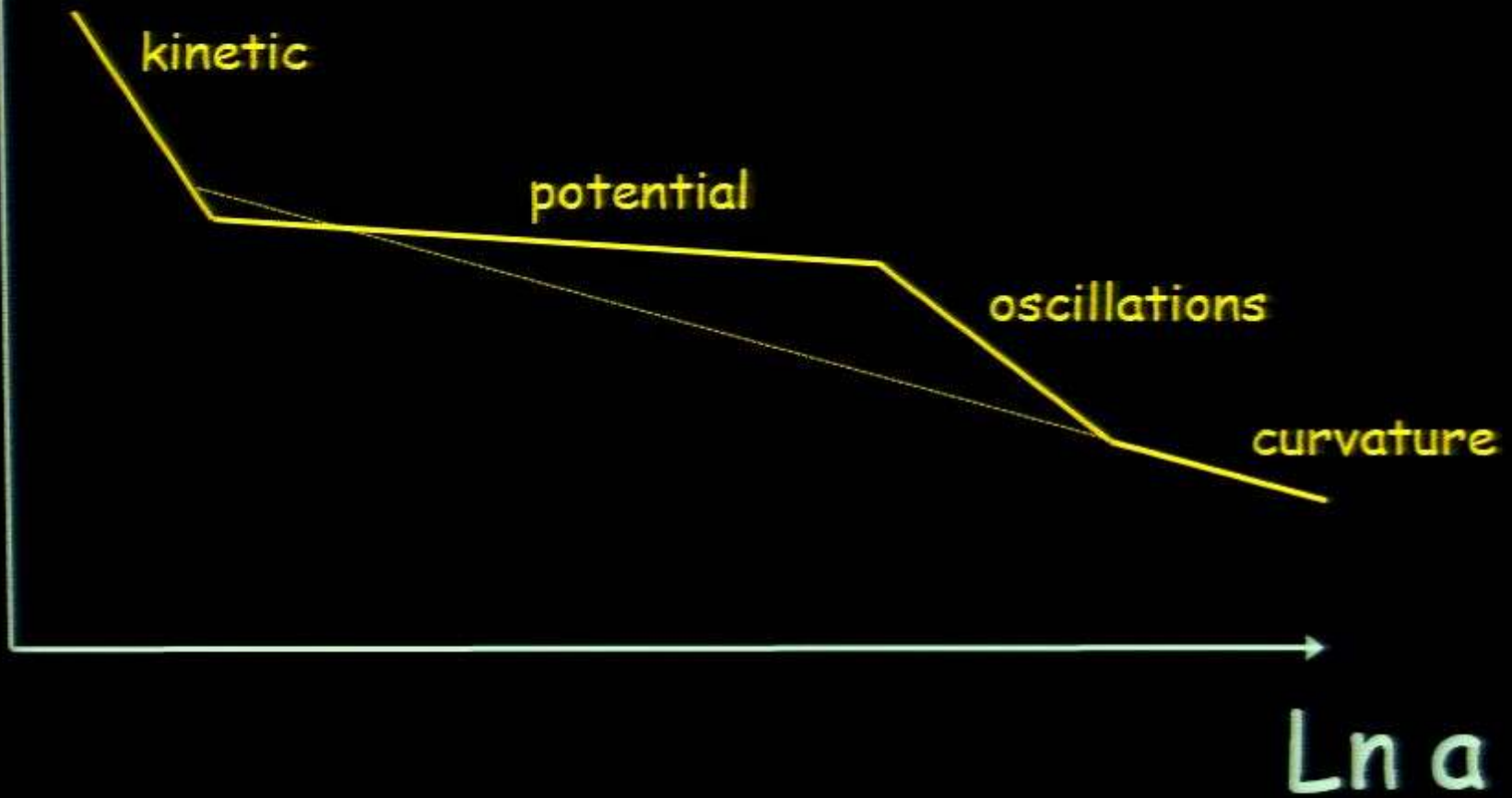
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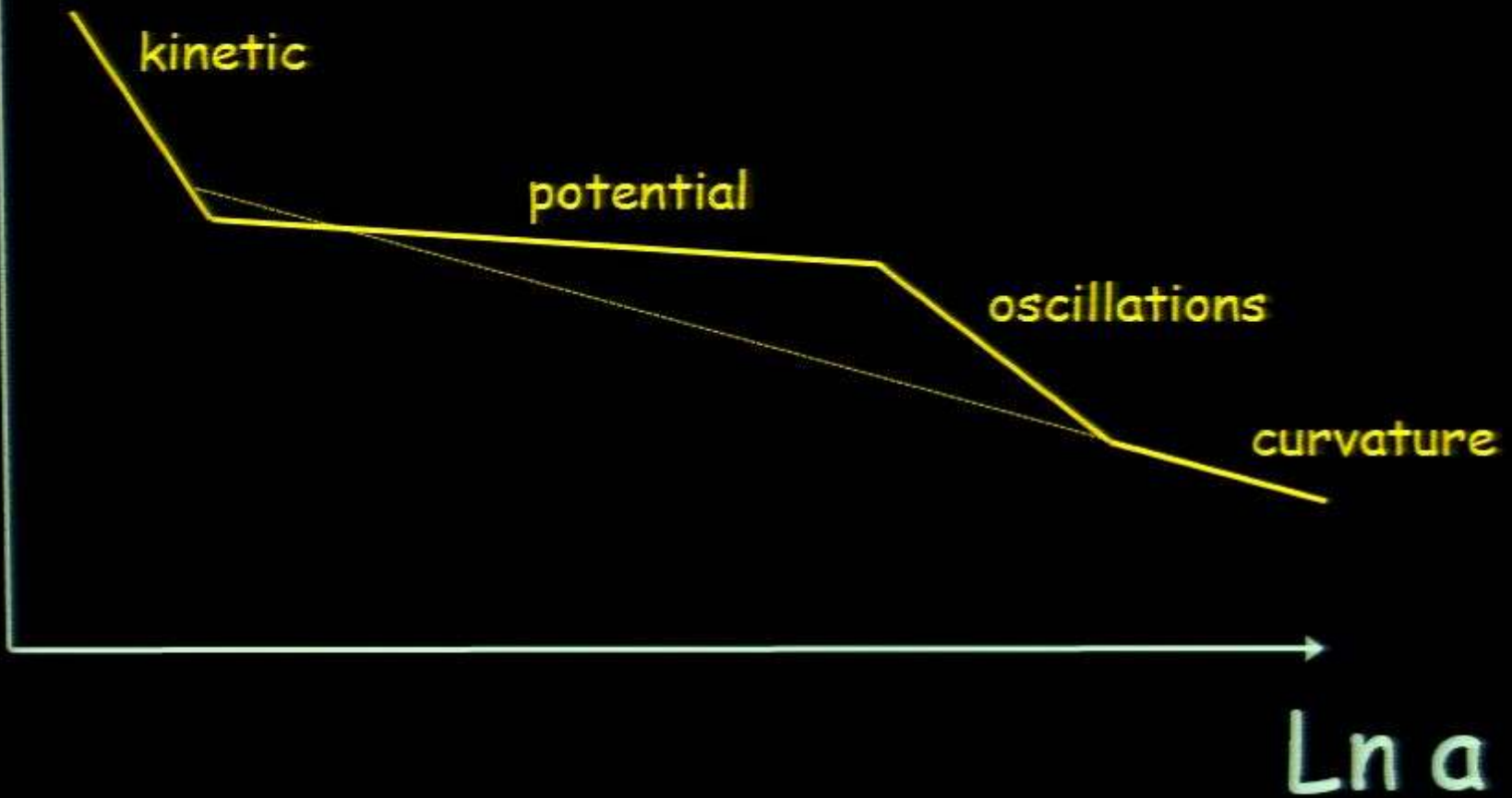
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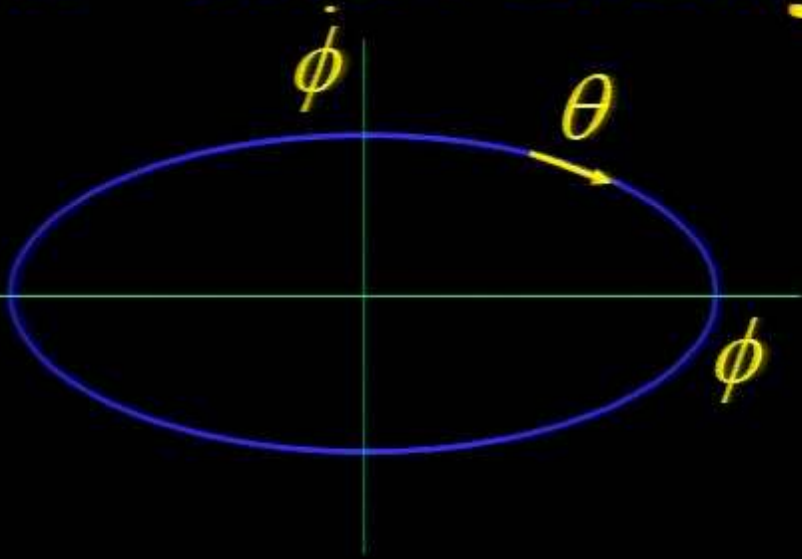
generic open FRW cosmology

$\text{Ln } \rho$



$a = \text{const}$ slicing

$$\omega_c = a^3 d\dot{\phi} d\phi$$



$$\rho = \frac{1}{2}(\dot{\phi}^2 + m^2 \phi^2)$$

$$\sim C/a^3, \quad a \rightarrow \infty,$$

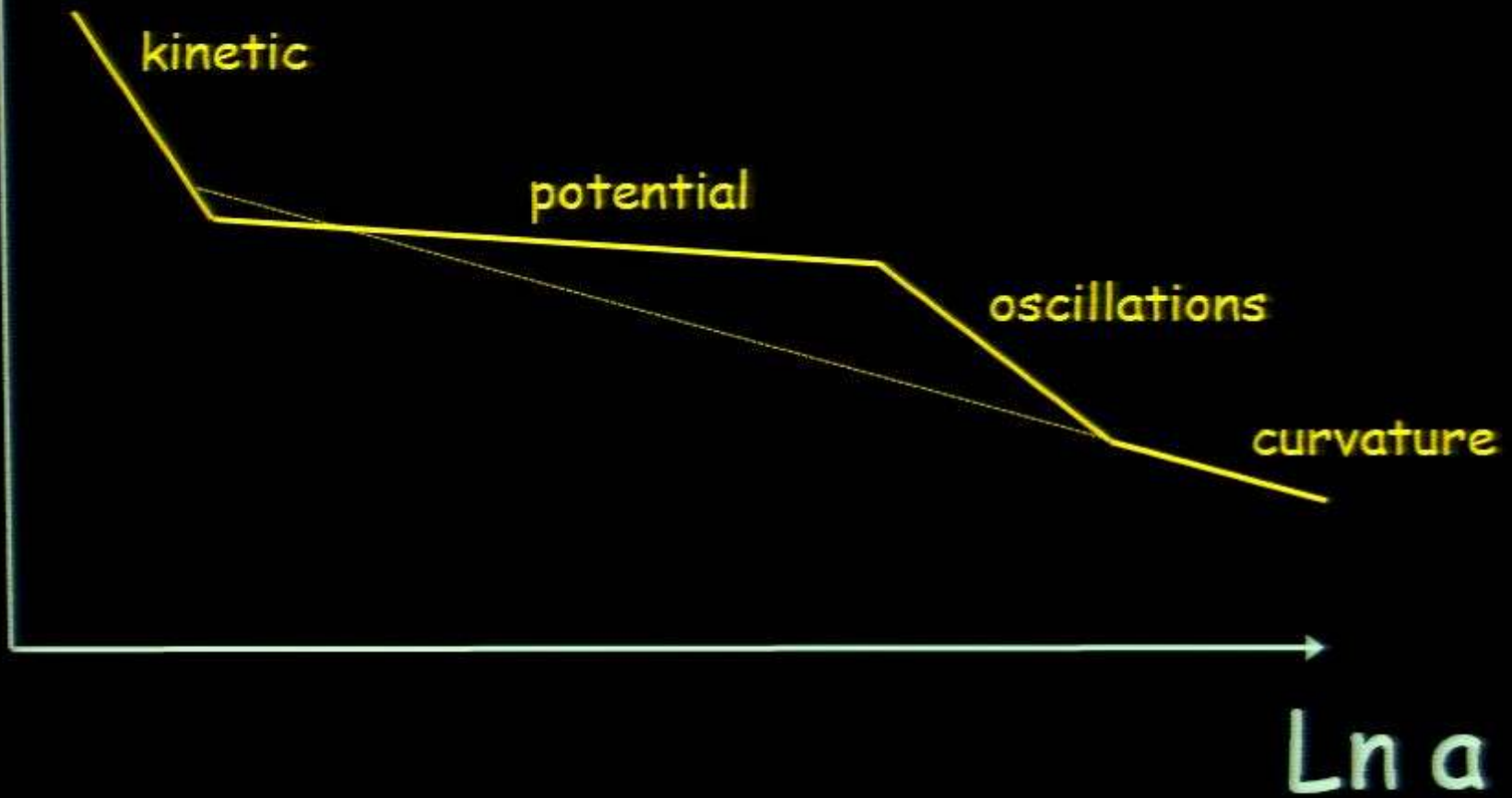
ρa^3 adiabatically conserved

$$ds^2 = \frac{d\alpha^2}{1 + 8\pi G \rho \alpha^2 / 3} + \alpha^2 dH_3^2 \approx dM_4^2 - \frac{8\pi G C}{3} \frac{d\alpha^2}{a}$$

(C is analogous to the AdM mass)

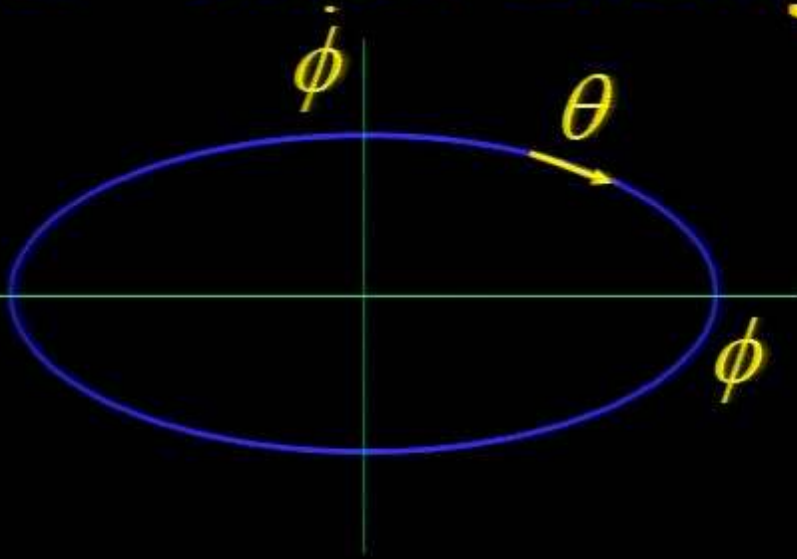
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(C is analogous to the AdM mass)

Including backreaction

$$\phi \sim \frac{(2C)^{\frac{1}{2}}}{ma^{\frac{3}{2}}} \left(\cos \theta + \frac{C \cos^3 \theta}{24a} \right) + \dots,$$
$$\rho_\phi \sim \frac{C}{a^3} \left(1 + \frac{3 \sin 2\theta}{2ma} + \frac{C(2 \cos 2\theta + \sin^2 2\theta)}{24a} \right) + \dots$$

$$\text{where } \theta = m \left(a - \frac{5C}{24} \ln(ma) \right) + \theta_0$$

in large a limit, effect of matter
on background spacetime (i.e. gravity)
becomes negligible

we just have flat spacetime, and an
adiabatically expanding box filled with
matter

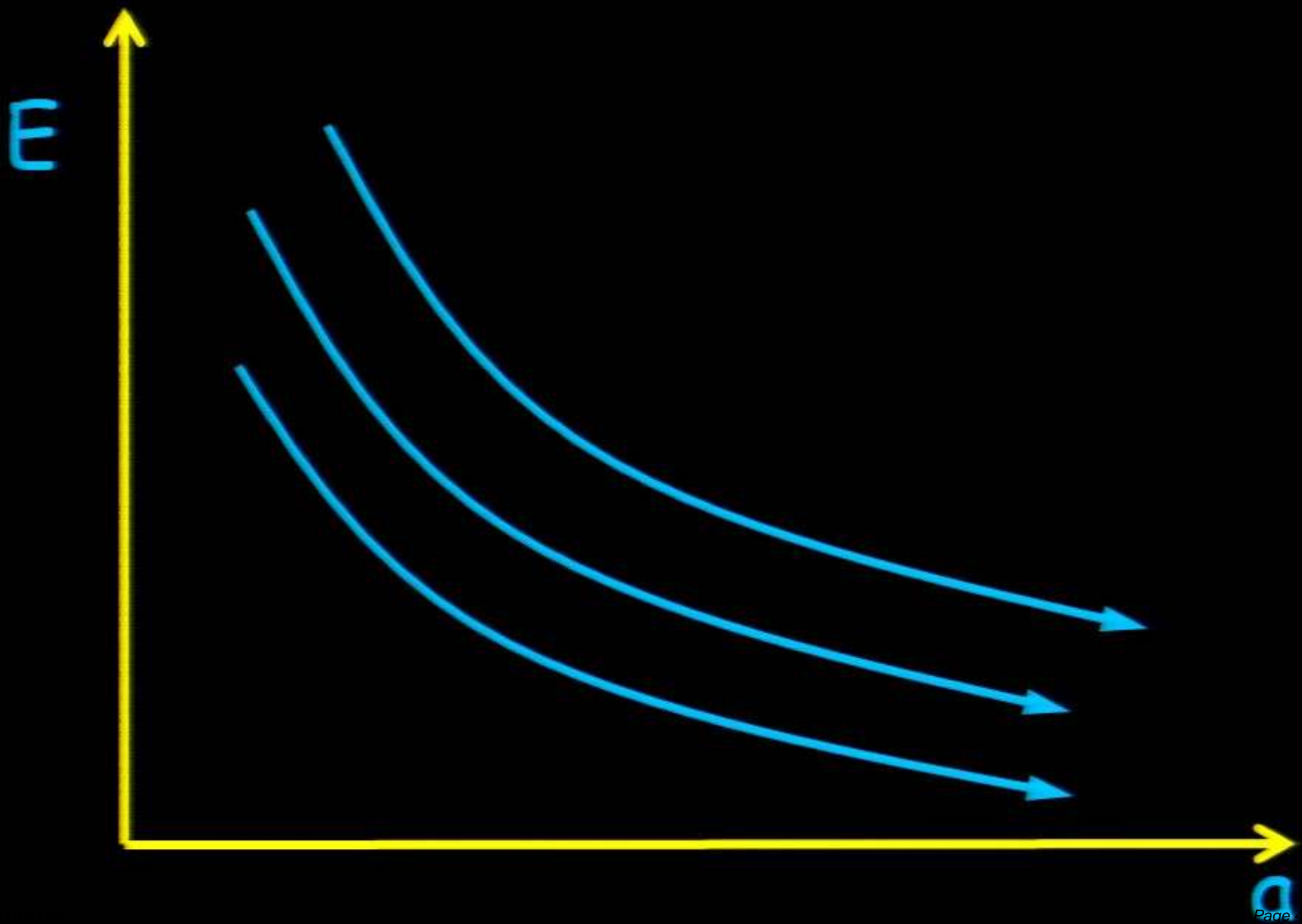
statistical ensemble: minisuperspace

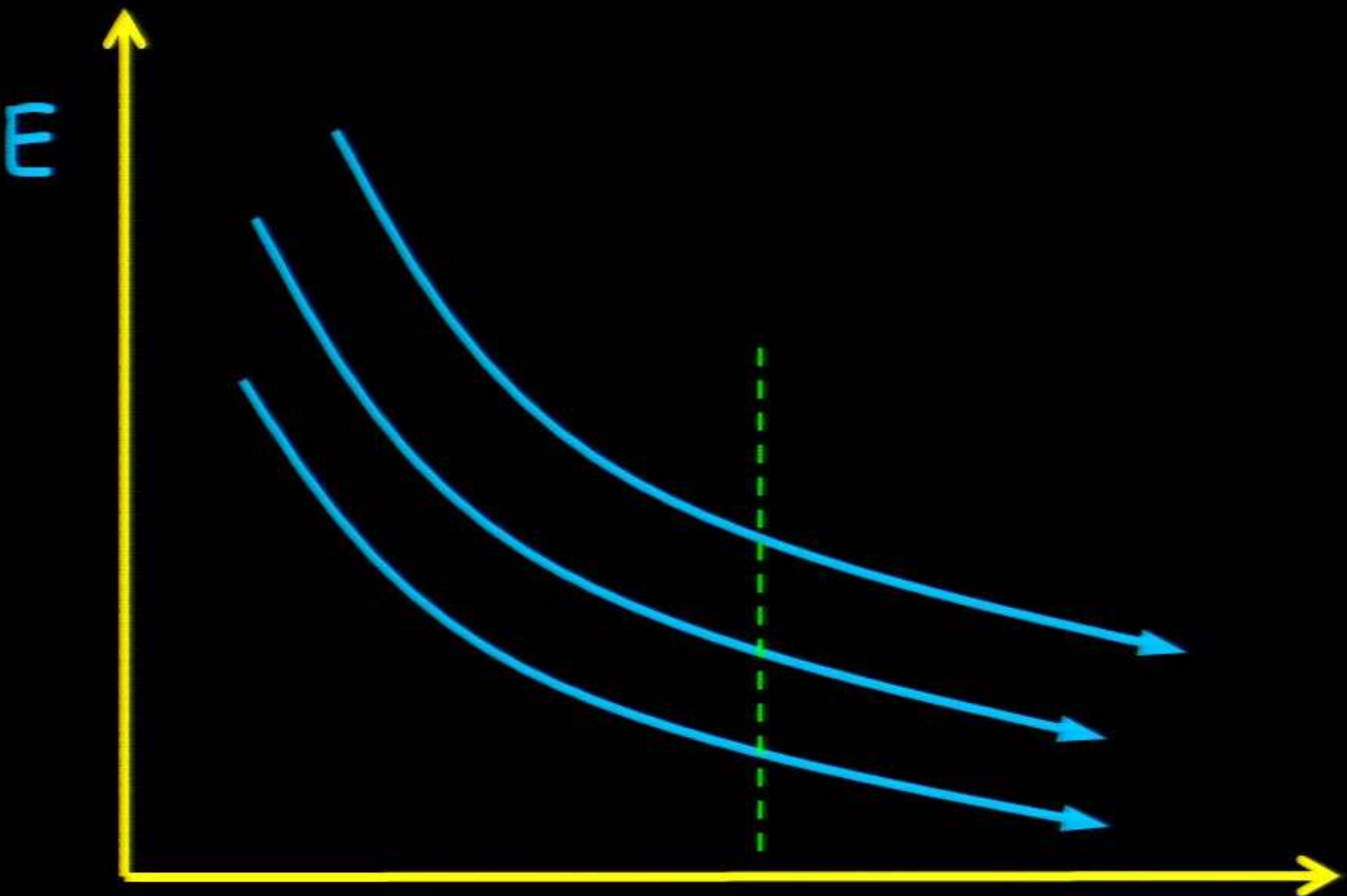
$$H_m(p_\phi, \phi) = \frac{1}{2} \left(\frac{p_\phi^2}{Ua^3} + Ua^3 V(\phi) \right)$$

$$\langle H_m \rangle = \frac{\int dp_\phi d\phi e^{-\beta H_m} H_m}{\int dp_\phi d\phi e^{-\beta H_m}} = E(a, \beta)$$

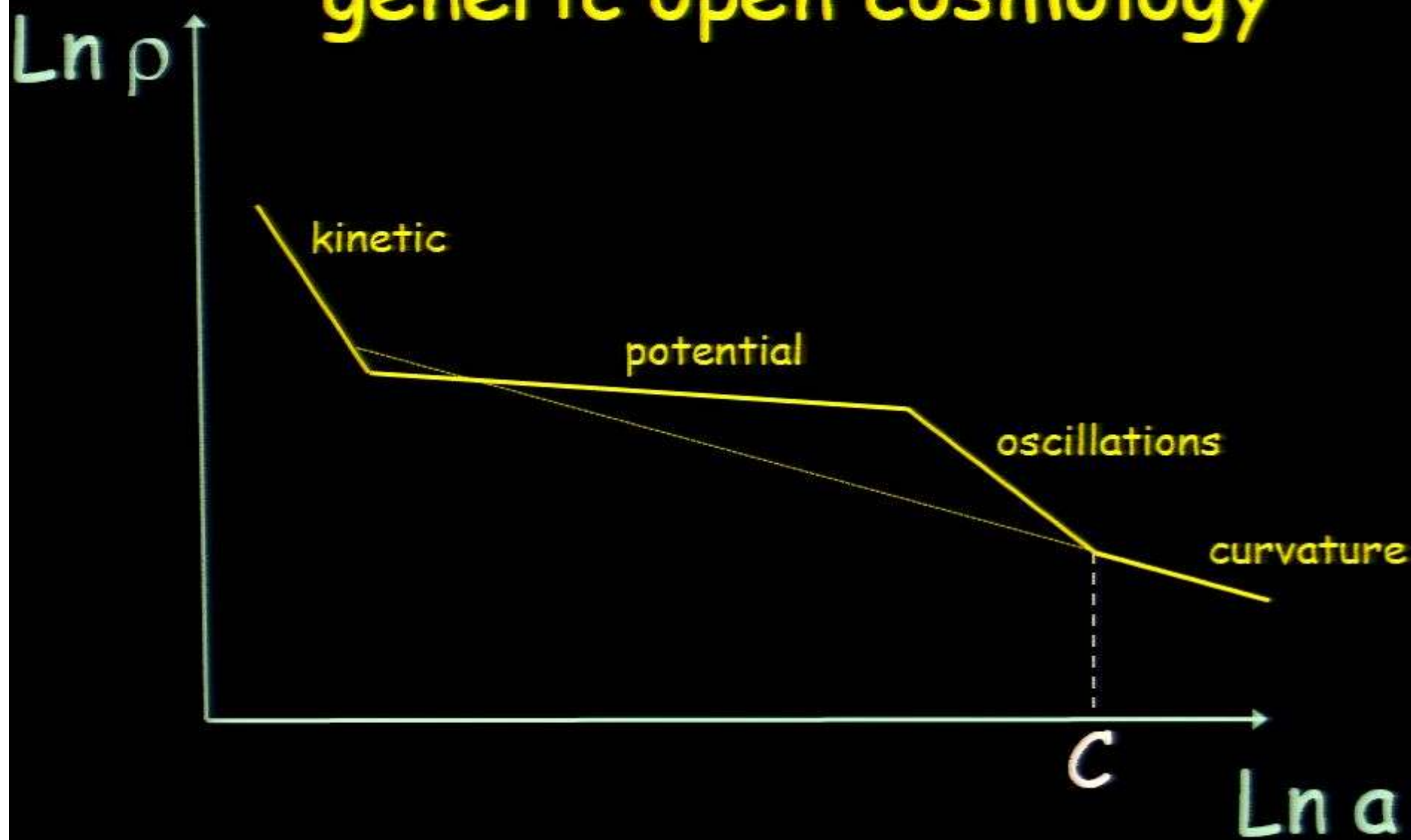
entropy $S = S_m = \ln\left(\frac{Ua^3 \rho_\phi}{m}\right) = \text{adiabatic invariant}$

constant entropy = fixed C

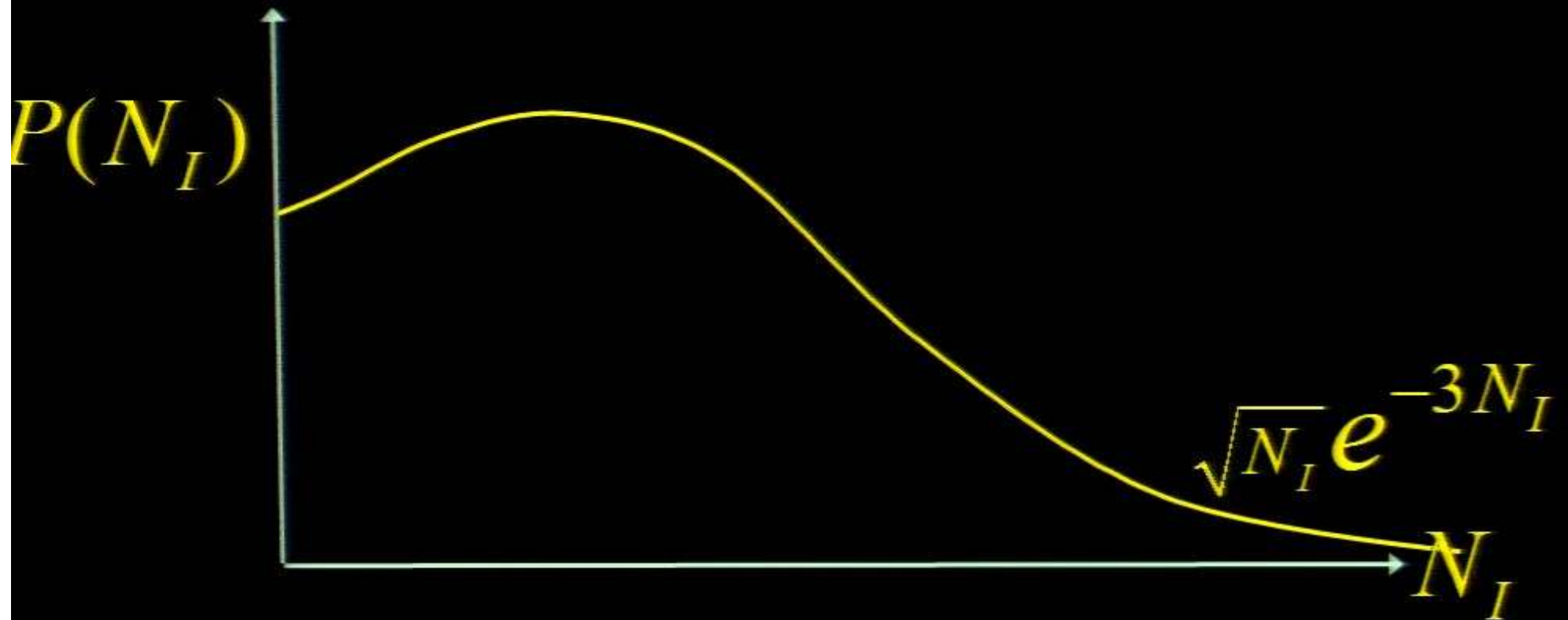




generic open cosmology



Canonical measure for inflation



Finite C result is always lower than $C = \infty$ result at large N_I

* Note: "attractor" becomes "repeller" because statistical ensemble defined in asymptotic region where gravity becomes unimportant: the **future**

* N slow-roll inflaton fields (N-flation) makes problem **worse**

$$P_{>}(N_I) \propto \prod_i \delta\theta_i \sim e^{-3N_I N}$$

* this analysis makes precise a problem identified by Penrose long ago (Annals NYAS, 1989)

* with this canonical measure, slow-roll/'chaotic' inflation **cannot** be considered a viable explanation for the observed state of the cosmos

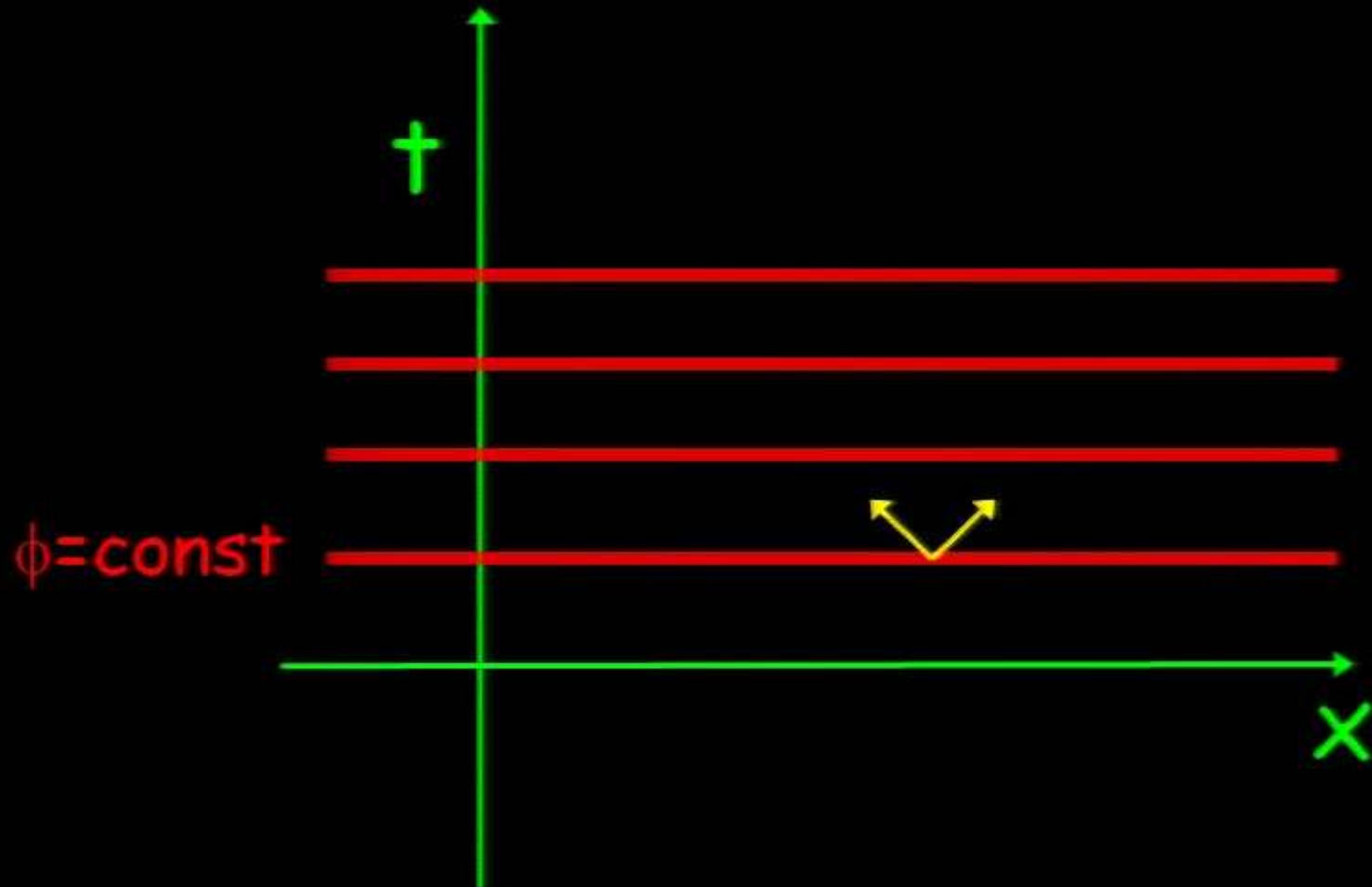
* can be extended to landscape with inclusion of complex solutions describing tunneling: false vacua don't help

- * Could add further restrictions
e.g. insist ϕ bounded \rightarrow 1-parameter family $\phi_0(S_{\text{final}})$
- * Large entropy \rightarrow many e-folds, but large final entropy (or flatness) is then an input, not a prediction
- * Could restrict to non-singular spacetimes (eg Page) and just reject all $k=-1$ solutions
- * $k=+1$ "bounce" solutions possible, have finite measure but collapsing phase very unstable
- * How do we justify rejecting cosmic singularities but allowing black holes? What about white holes?

What could be wrong?

- * canonical measure?
- * neglect of: entropy production?

ϕ decays: information stored in radiation



But the relative proportions of phase space corresponding to N_{I} or more e-folds are preserved (by unitarity or Liouville)

What could be wrong?

- * canonical measure?
- * neglect of: entropy production?
 - : inhomogeneities?
 - : quantum fluctuations?
- * constraint on final entropy?
maybe, but what alternative?
- * global structure?
- * inflation?

compare "cyclic/ekpyrotic" theory, where gravity is unimportant in the **past**, and according to the corresponding canonical measure, **every** trajectory undergoes near-maximal ekpyrosis

(w/P. Steinhardt)

Scale Invariant perturbations

- * near-massless scalar in de Sitter
- * exponentially flat potentials in collapsing phase (cyclic)
- * scale-invariant duals via holography

Summary

concordance cosmology has plenty of challenges

- singularity
- tuning
- reliance on anthropics
- measure: a good one exists!

No Signal

VGA-1

No Signal

VGA-1