

Title: Conformal Cyclic Cosmology: Some Striking New Observational Support

Date: Jul 14, 2011 12:10 PM

URL: <http://pirsa.org/11070010>

Abstract:

CRAZY?

CRAZY?

YES!

YES!

But it provides new insights into several outstanding problems of cosmology:

- second law of thermodynamics
- "dark energy"
- "dark matter"
- black-hole "information paradox"
- structure of the microwave background
- inflation? x

- No eternal inflation
- No ordinary inflation
- No Boltzmann brains
(B-brains)
- No D-branes
- No p-branes
- No strings
- No extra dimensions
- No instantons
- No AdS-CFT
- No quantum gravity
- No symm. breaking
- No bubble
- No tunnelling
- No anthropics
- No landscape

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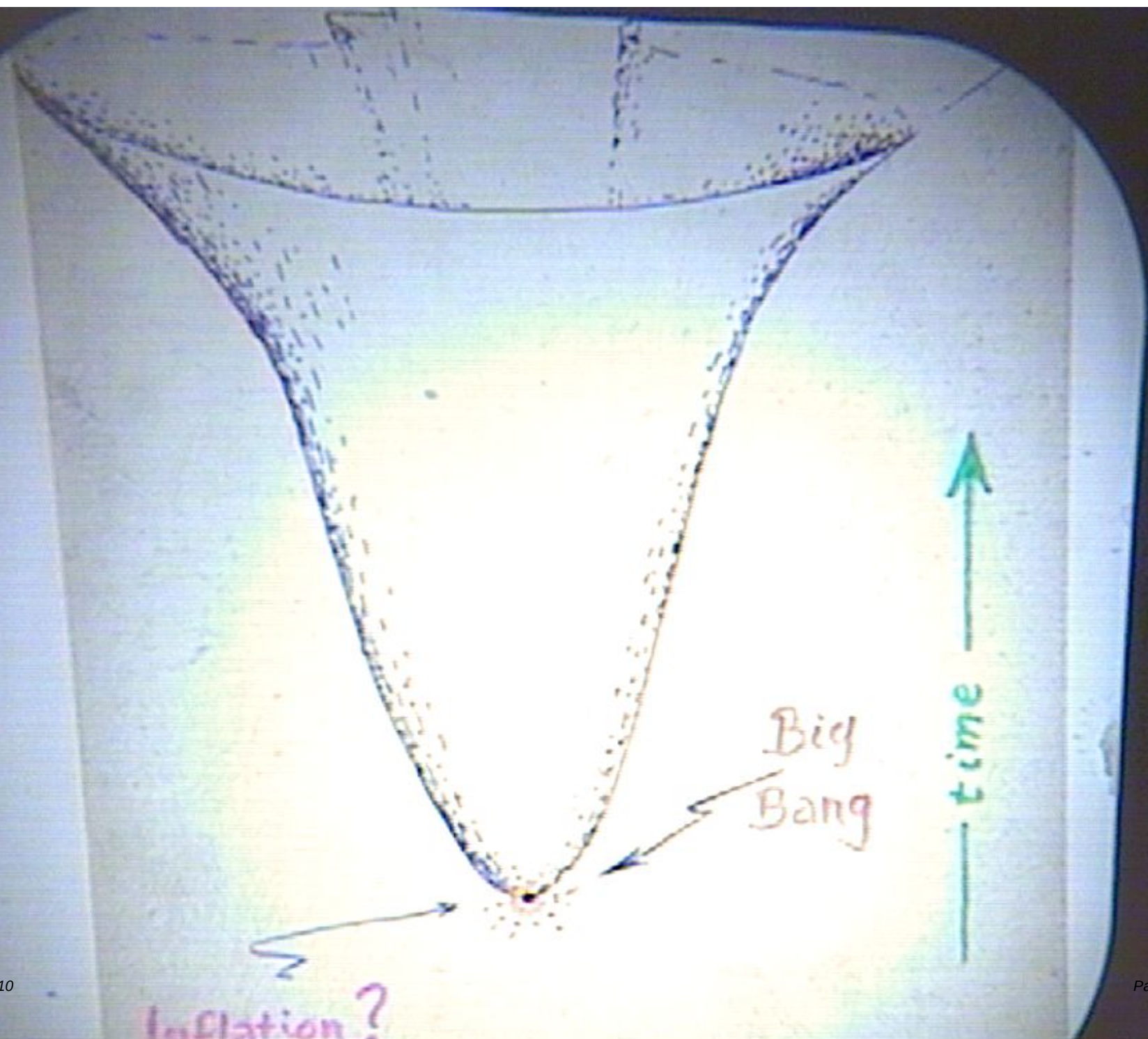
- Only standard classical Einstein $\Lambda(>0)$ -general relativity
— with a slight extra twist!

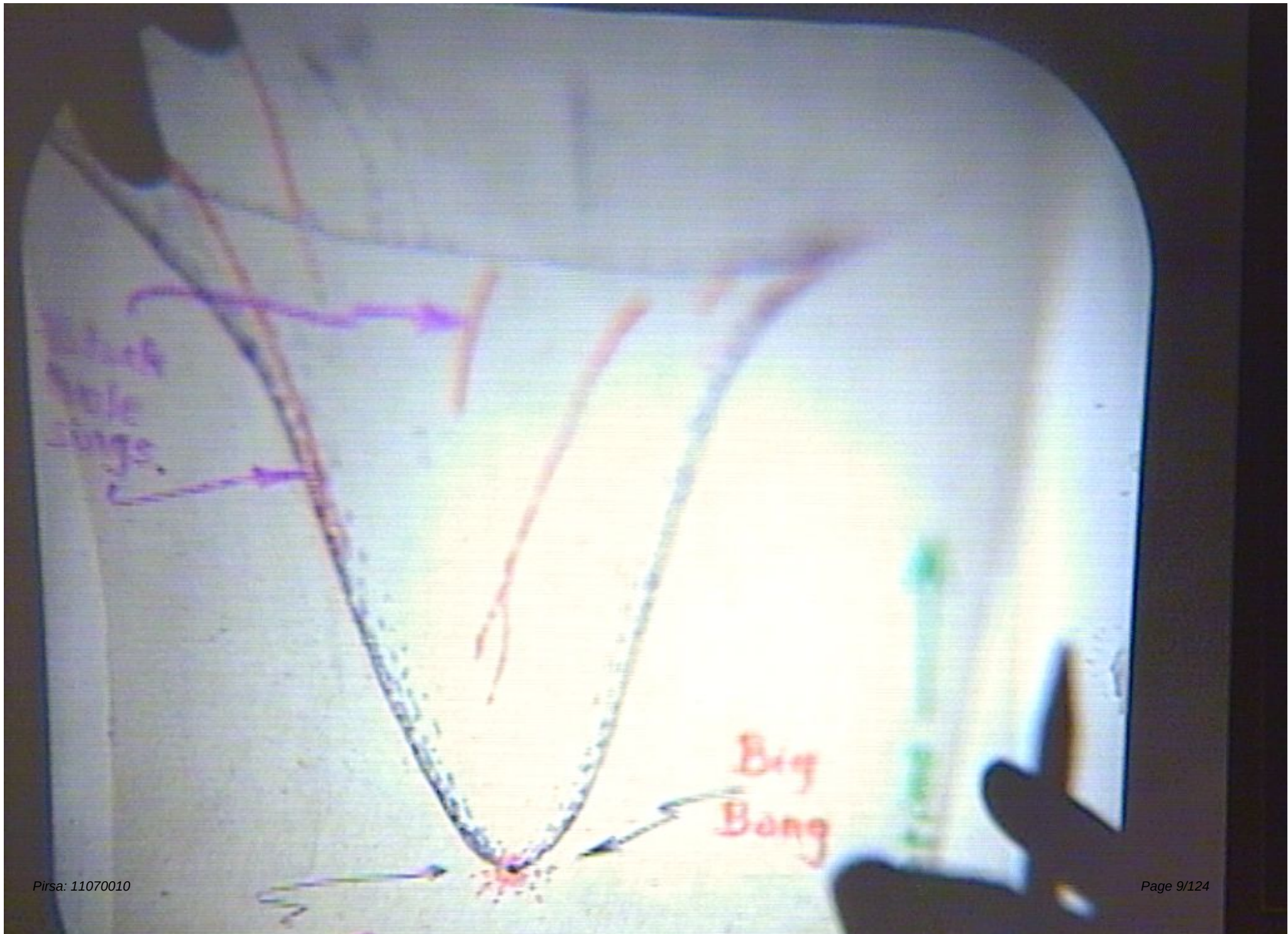
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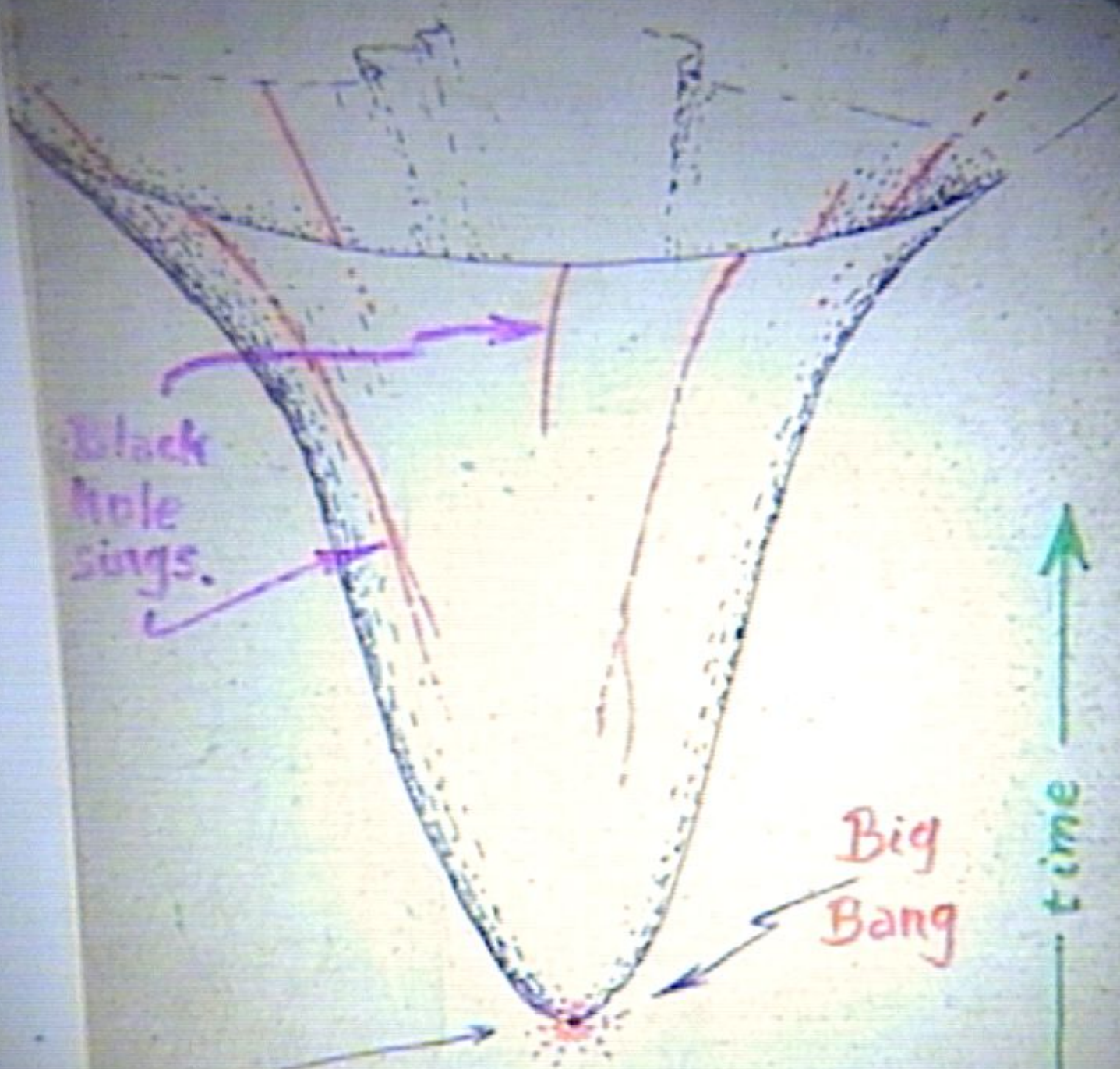




Black
hole
sings.

Big
Bang

time



Fundamental raised by the 2nd law of thermodynamics & the nature of the Big Bang:

Why did the BB not involve any white holes (time-reversed black holes) ?

Not answered by inflation or by any (current) theory of quantum

Fundamental problem
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Why did the BB
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Not answered by
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quantum

initial
collapsing
material

time $\rightarrow$ up to 10^{10}

singularity

Hawking
radiation



Hawking
evaporation
of a black
hole



Hawking
evaporation
of a black
hole

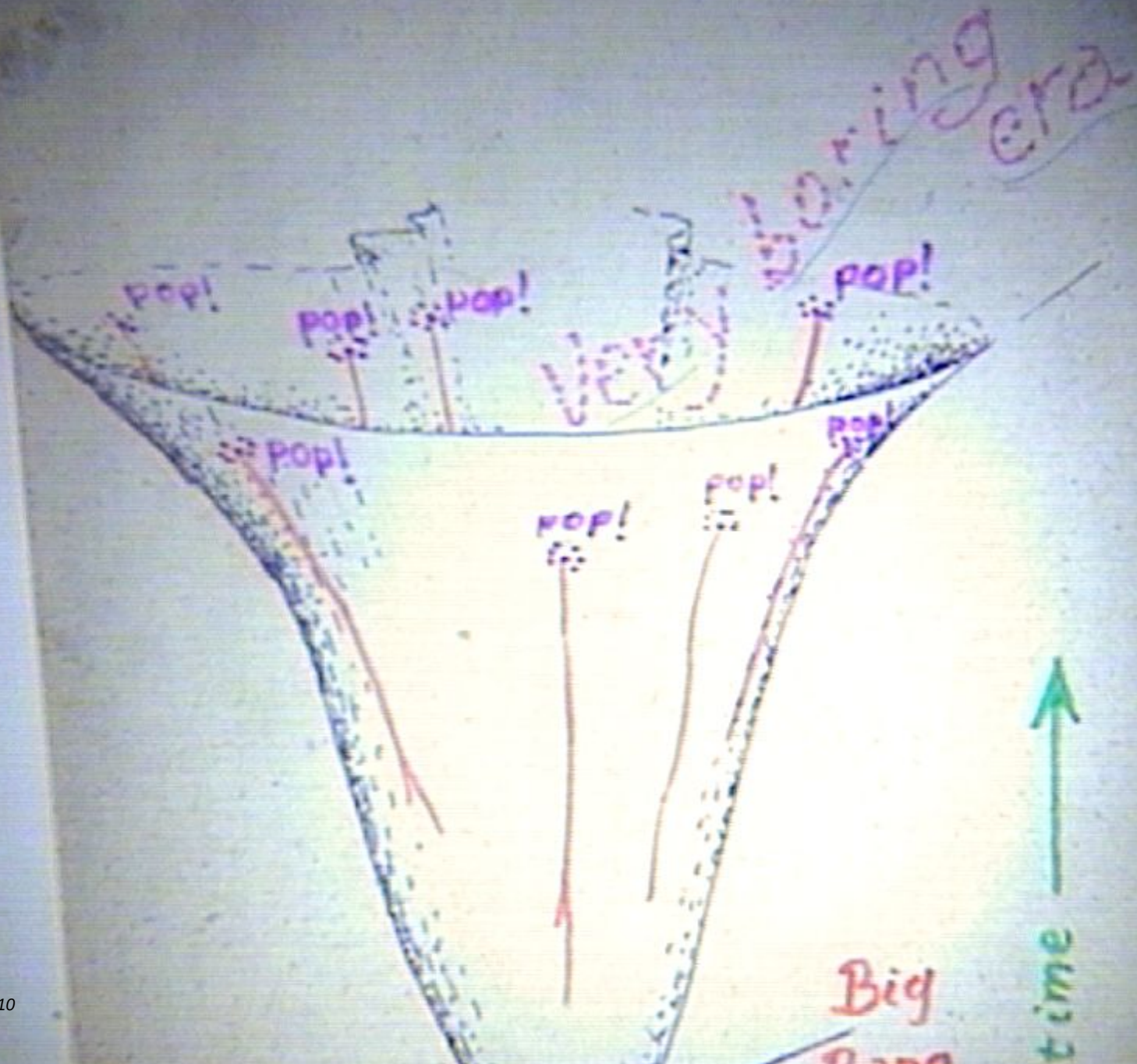


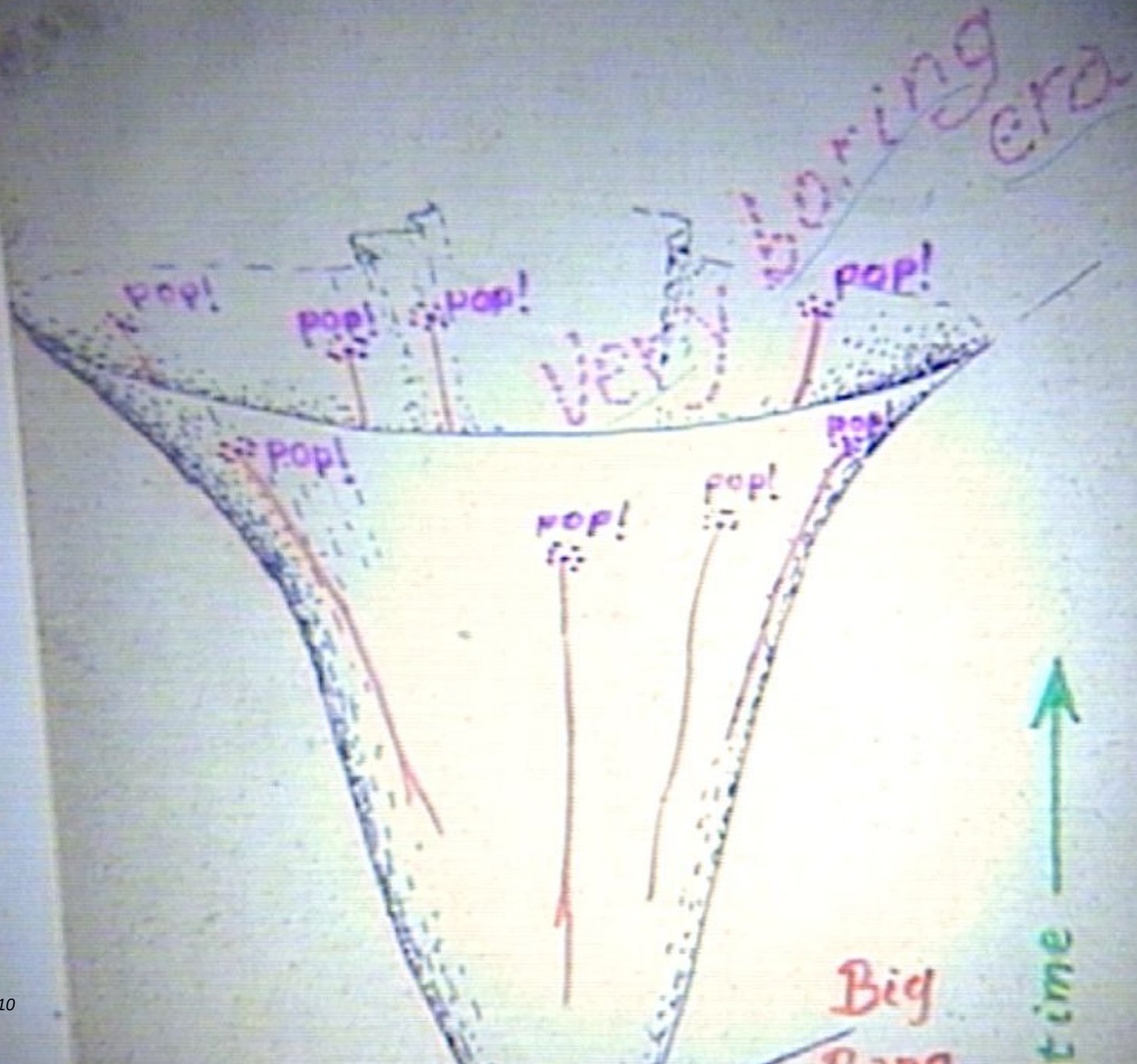
initial
collapsing

time $\rightarrow$ up to 10^{100} yrs.

singularity

Hawking
radiation





Planck: $E = h\nu$

Einstein: $E = mc^2$

$$\therefore \nu = m \times \left(\frac{c^2}{h} \right)$$

particle of
rest-mass m

frequency ν

metric

g_{ab}

10 components
per point

conformal
structure
cone itself

9 components
per point

Planck: $E = h\nu$

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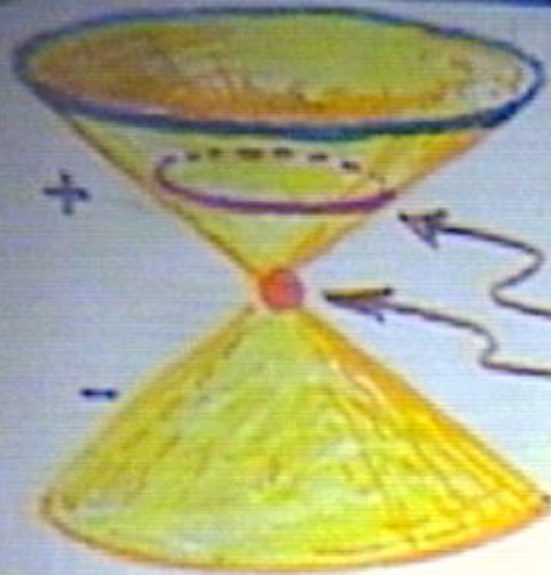
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Space-time picture



Space picture



Special
Relativity

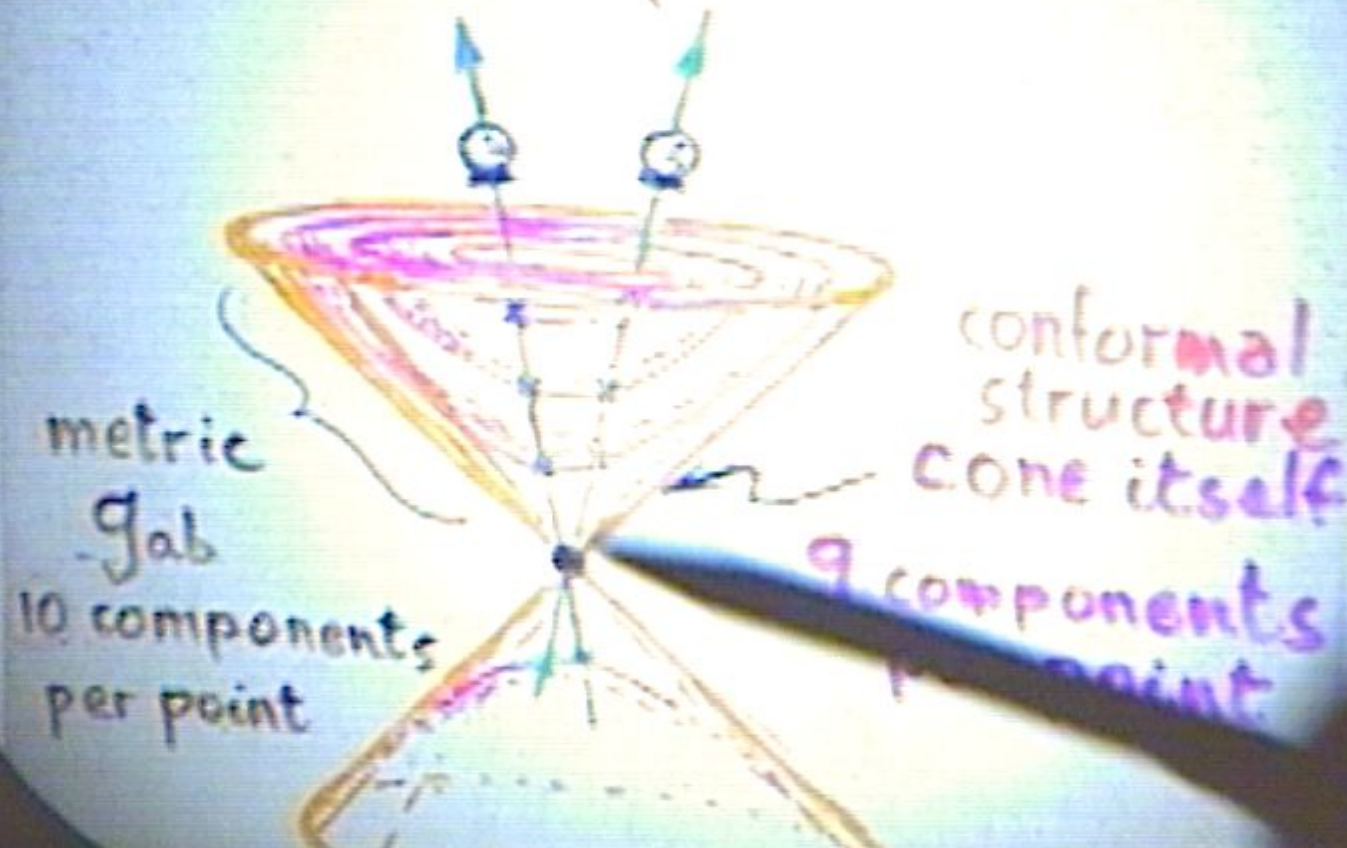


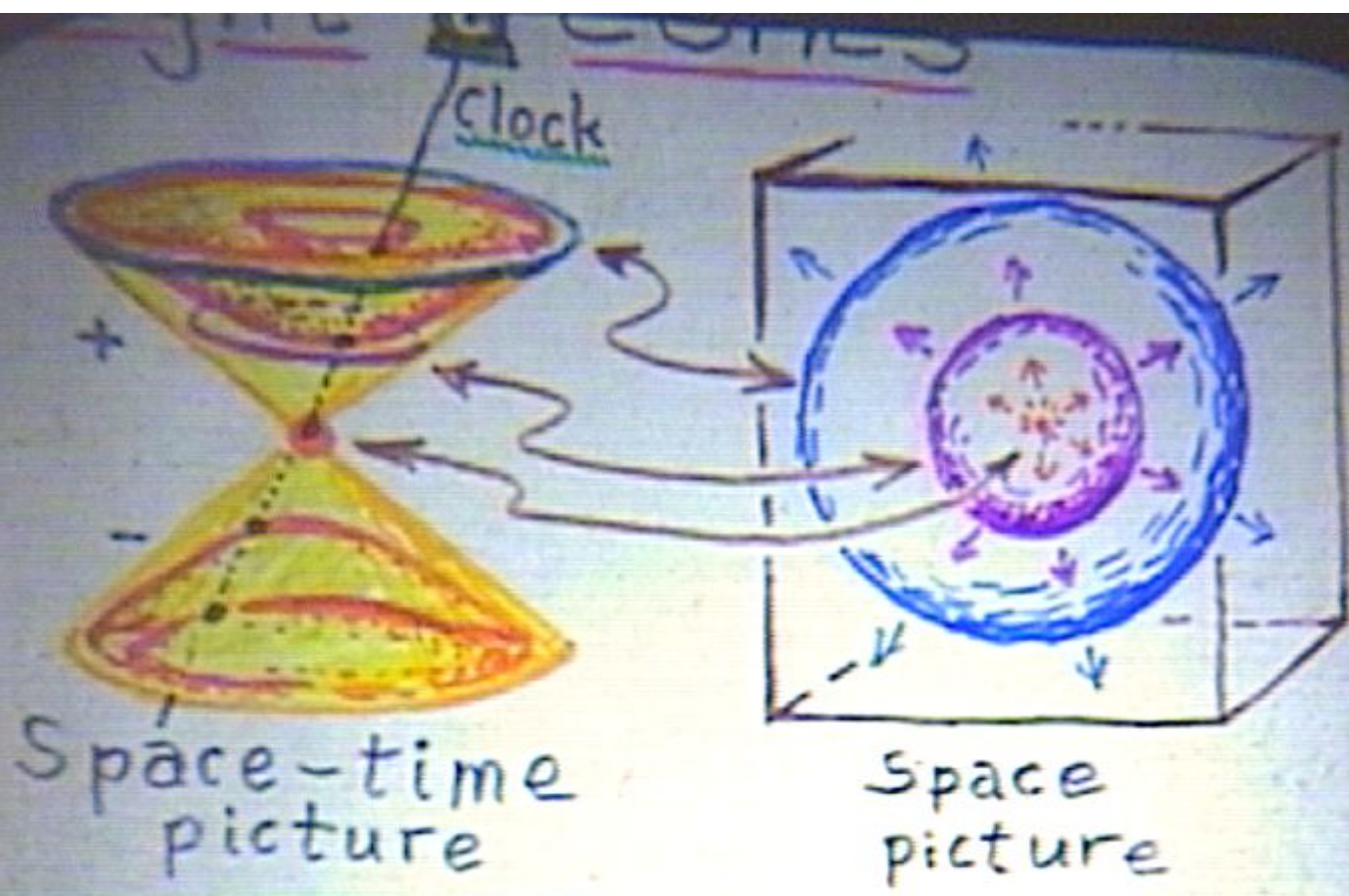
General
Relativity

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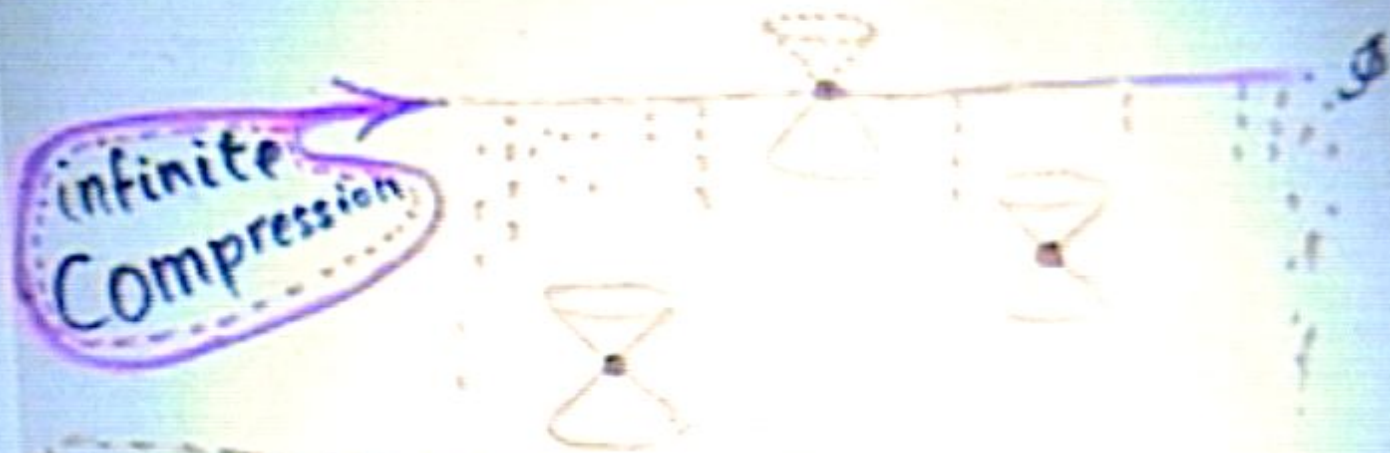
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The extremely remote future

Much matter collapses to black holes. Eventually expanding universe cools to lower than black-holes' Hawking temperature. Then, the hole gradually evaporates away, very slowly until it disappears.

$\sim 10^{66}$ yrs for $M_{\odot}$, $\sim 10^{90}$ yrs. for galactic mass

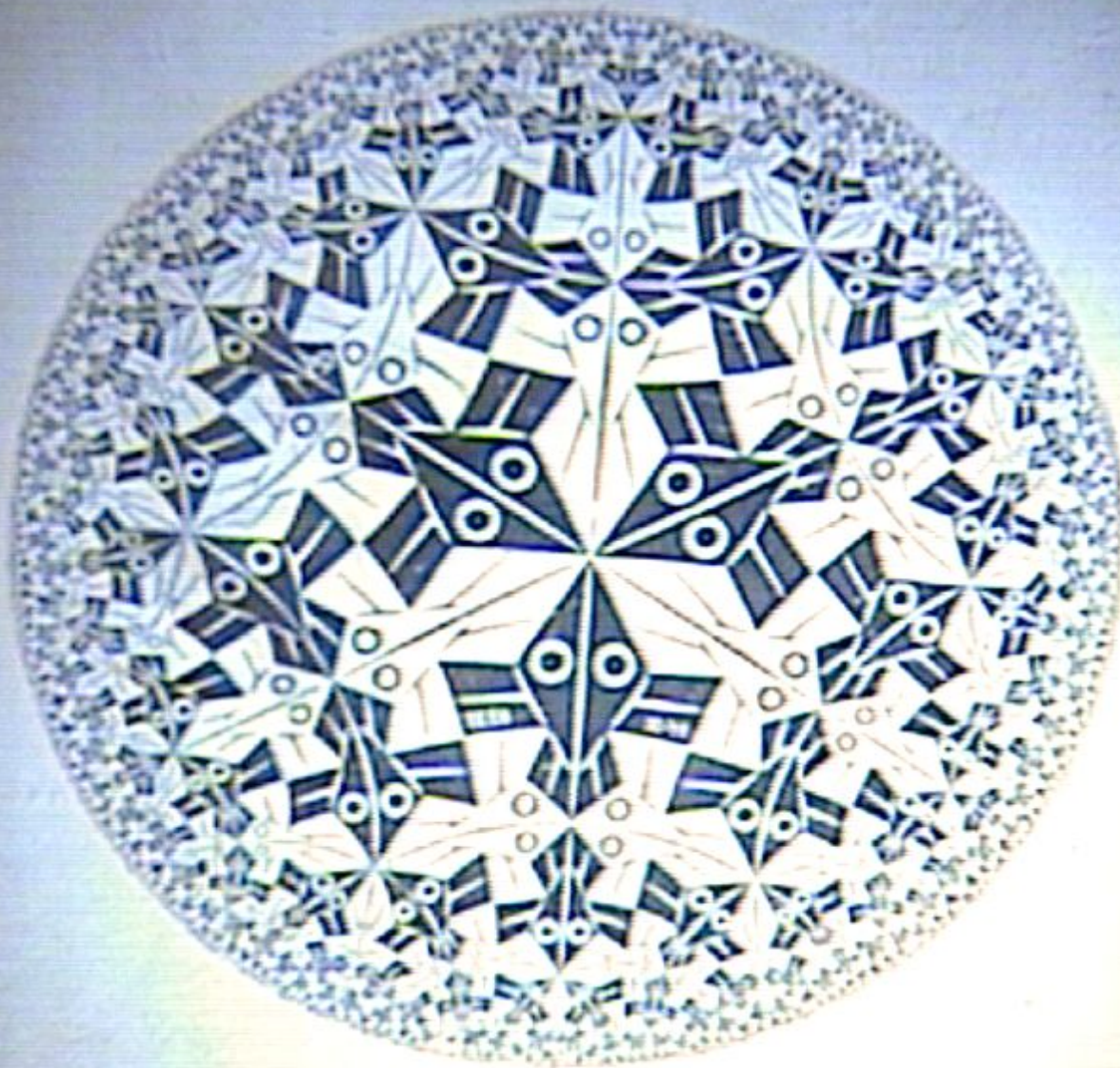
POP! it disappears



The extremely remote future

Much matter collapses to black holes. Eventually expanding universe cools to lower than black holes' Hawking temperature. Then, the hole gradually evaporates away, very slowly until it disappears. POP!

$\sim 10^{66}$ yrs for $M_{\odot}$, $\sim 10^{90}$ yrs. for galactic mass



COMP



The extremely remote future

Much matter collapses to black holes. Eventually expanding universe cools to lower than black-holes' Hawking temperature. Then, the hole gradually evaporates away, very slowly $\sim 10^{64}$ yrs. for Mo, $\sim 10^{90}$ yrs. for galactic ~~pop~~ ^{pop} it disappears. With only massless ingredients left, the universe loses track of time. No way to build a clock!

Eternity is no time at all, for a photon!
Conformal geometry
Conformal $\infty(\mathcal{F})$ is spacelike if $\Lambda > 0$. Mathematical

regain electrons?
mass evaporation?

Extremely remote future

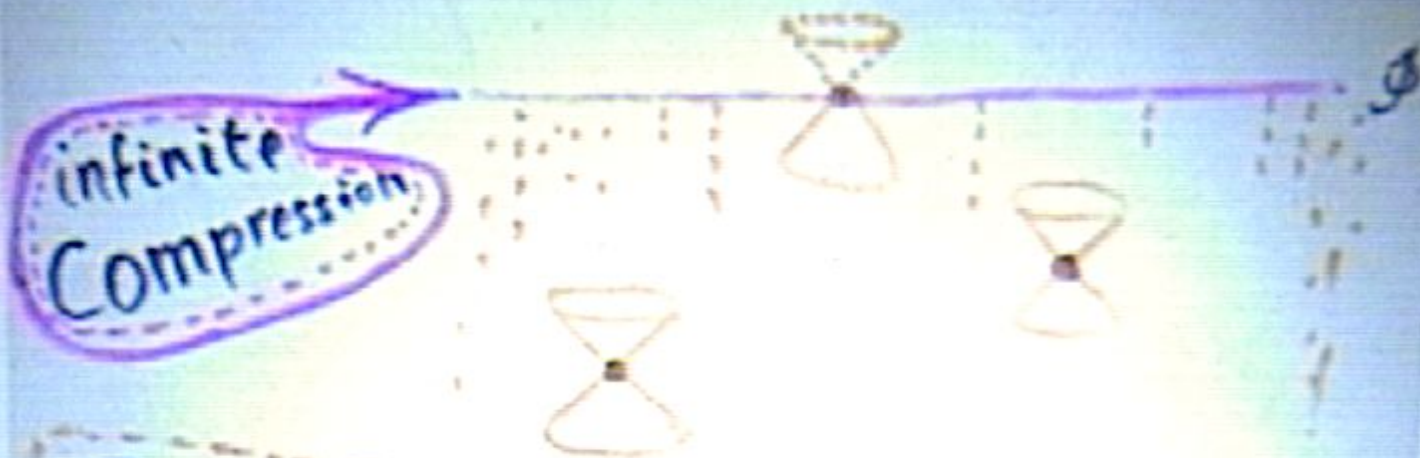
Much matter collapses to black holes. Eventually expanding universe cools to lower than black holes' Hawking temperature. Then, the hole gradually evaporates away. Very slowly until ^{300,000 yrs for $M_{\odot}$} $\sim 10^{90}$ yrs. for galactic ~~it~~ ^{it disappears}. With only massless ingredients left, the universe loses track of time. No way to build a clock!

Eternity is no time at all, for a photon!

Conformal geometry

conformal $\infty(\mathcal{I})$ is spacelike ∞ ^{time} $\Lambda > 0$. Mathematical trick!

rogue electrons?
mass evaporation?



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Compression



The extremely remote future

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 $\sim 10^{66}$ yrs for Mo, $\sim 10^{90}$ yrs. for galactic
until ~~it~~ ^{it disappears} it disappears. With only massless ingredients left, the universe loses track of time. No way to build a clock!

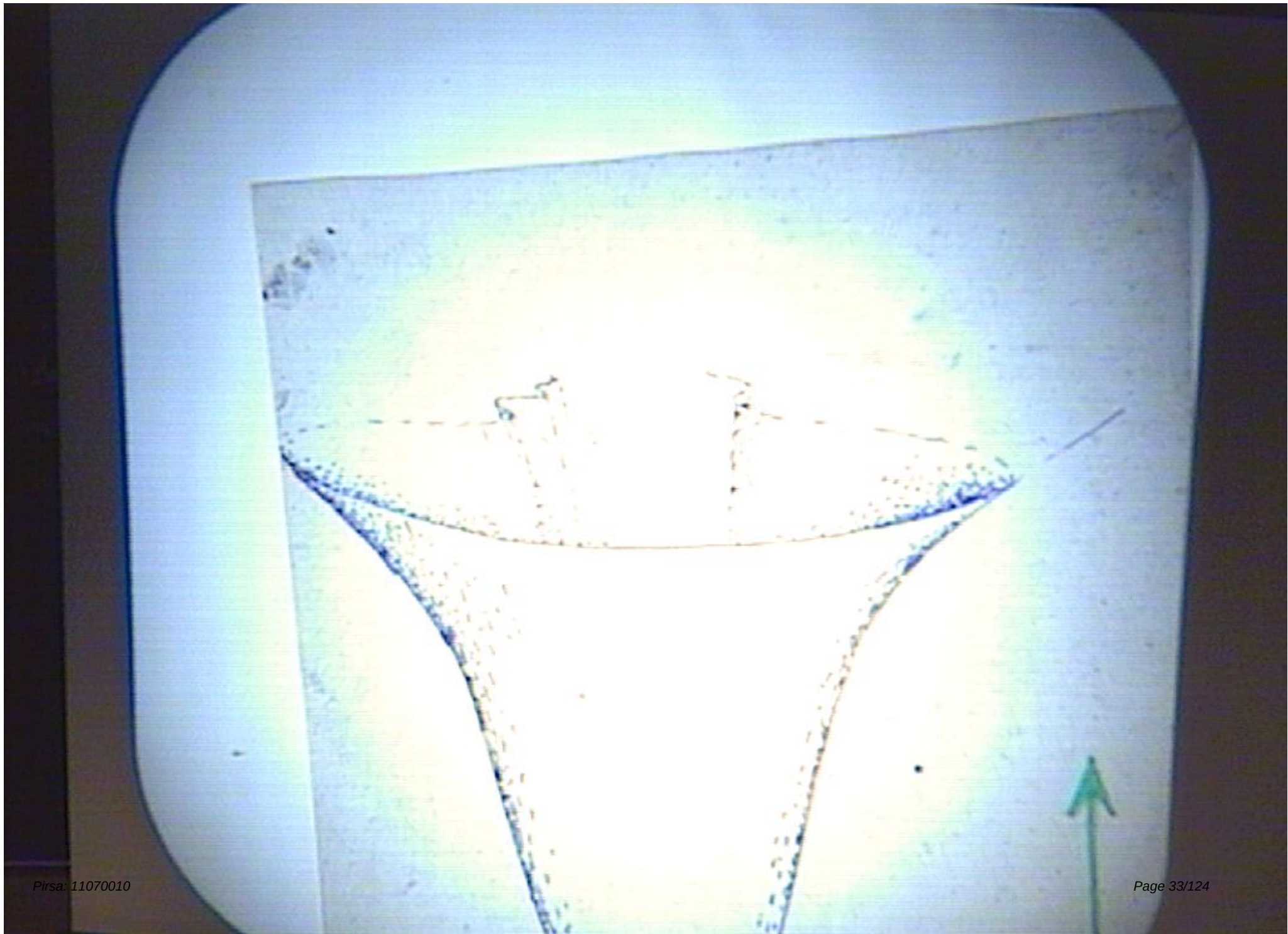
Eternity is no time at all, for a photon!

Conformal geometry

Conformal $\infty(\mathcal{F})$ is spacelike

0 time

How electrons?
How evaporation?



Weyl Curvature Hypothesis

The Big Bang must have been subject to a HUGE constraint — to a region of phase space no larger than 1 part in $10^{10^{124}}$

The specialness in the Big Bang appears to be only in Gravitation i.e. the Weyl curvature appears to have been $= 0$ at the Big Bang whereas it diverges wildly in black holes. But an awkward condition to state mathematically. Tod's proposal: space-time extendible conformally (i.e. as light-cone structure) to BEFORE Big Bang

(infinite expansion)



— to a region of phase space
no larger than 1 part in 10^{124}

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Just a mathematical
trick! ? Temperatures
become so large at the Big Bang
rest masses become

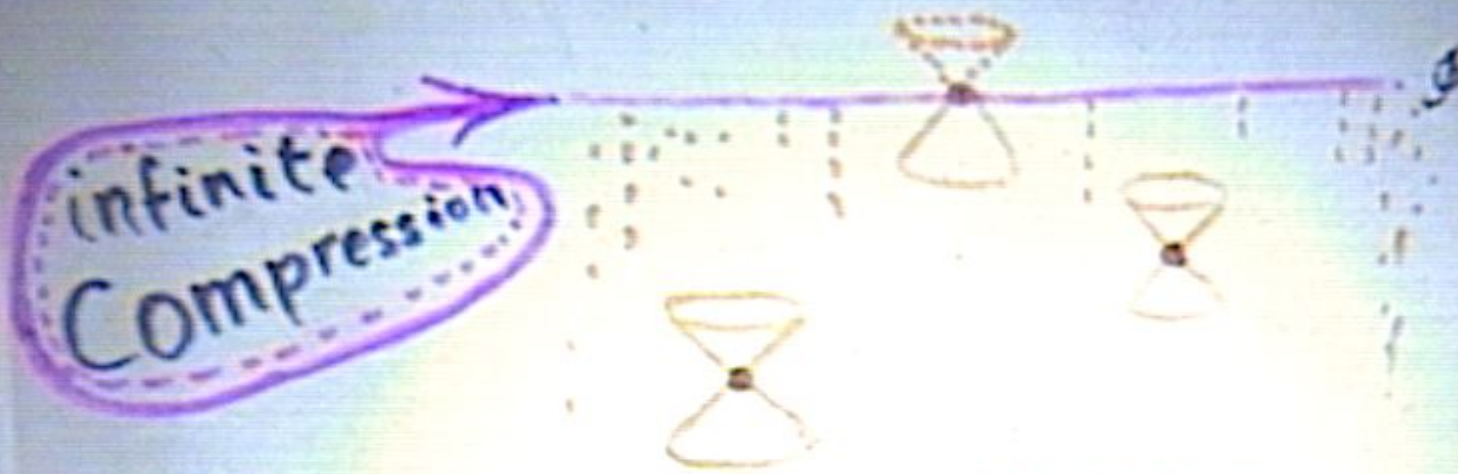
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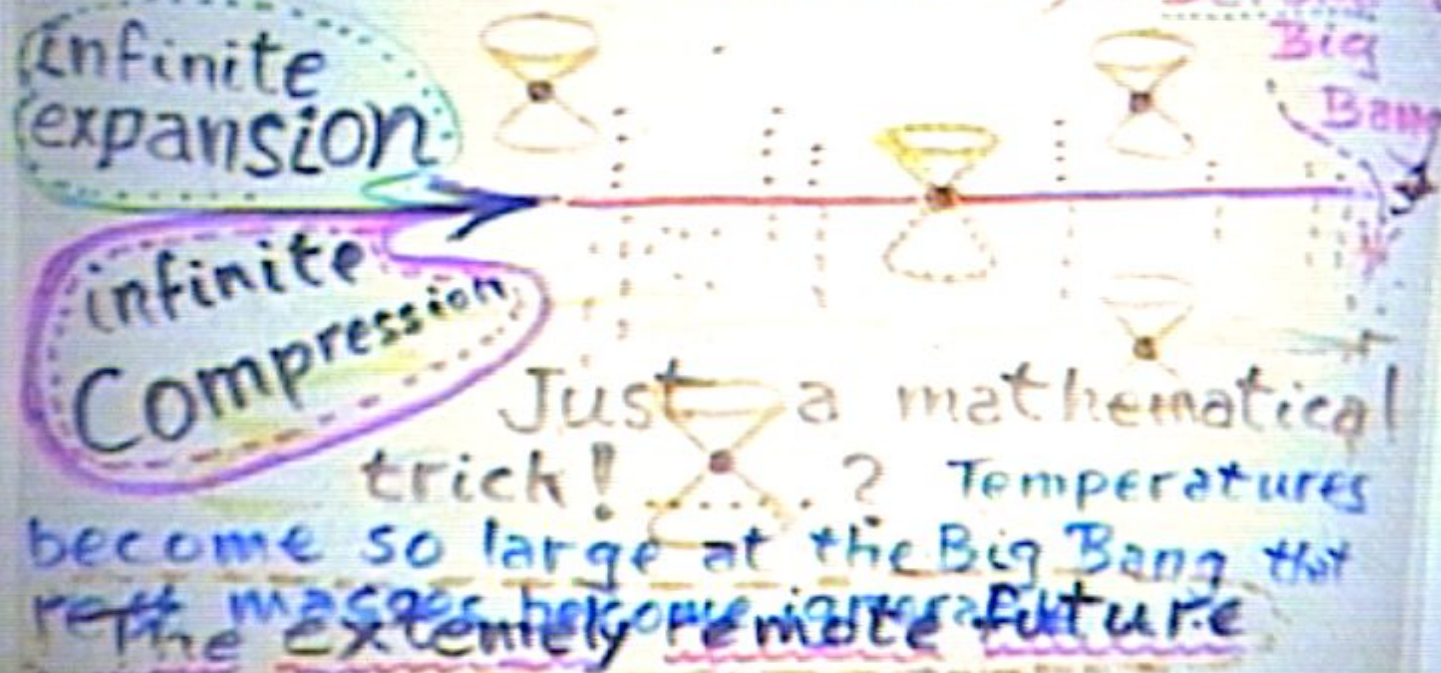
no larger than 1 part in 10

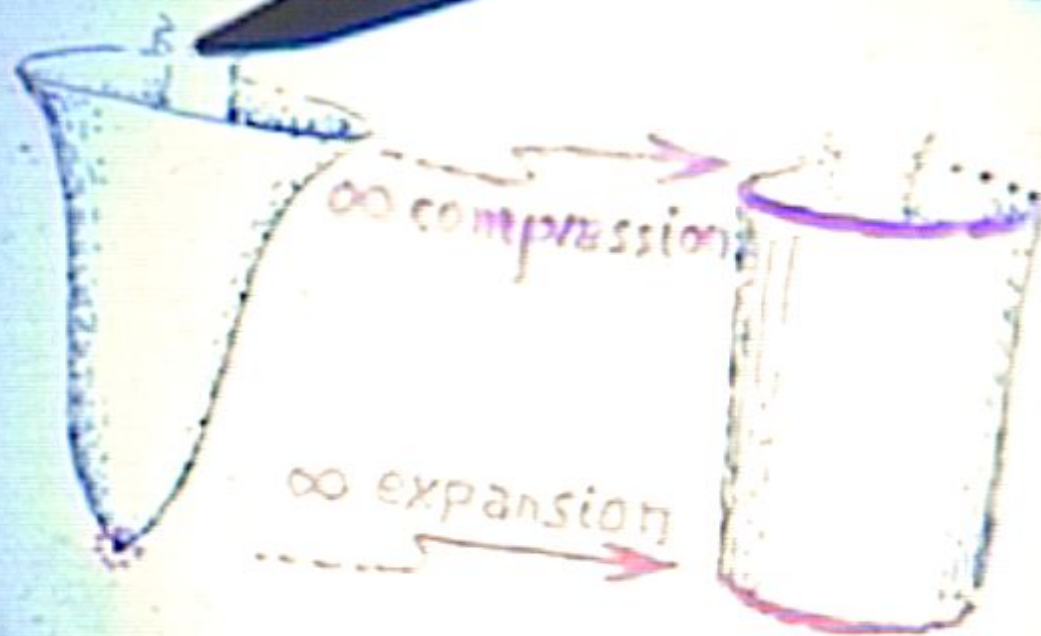
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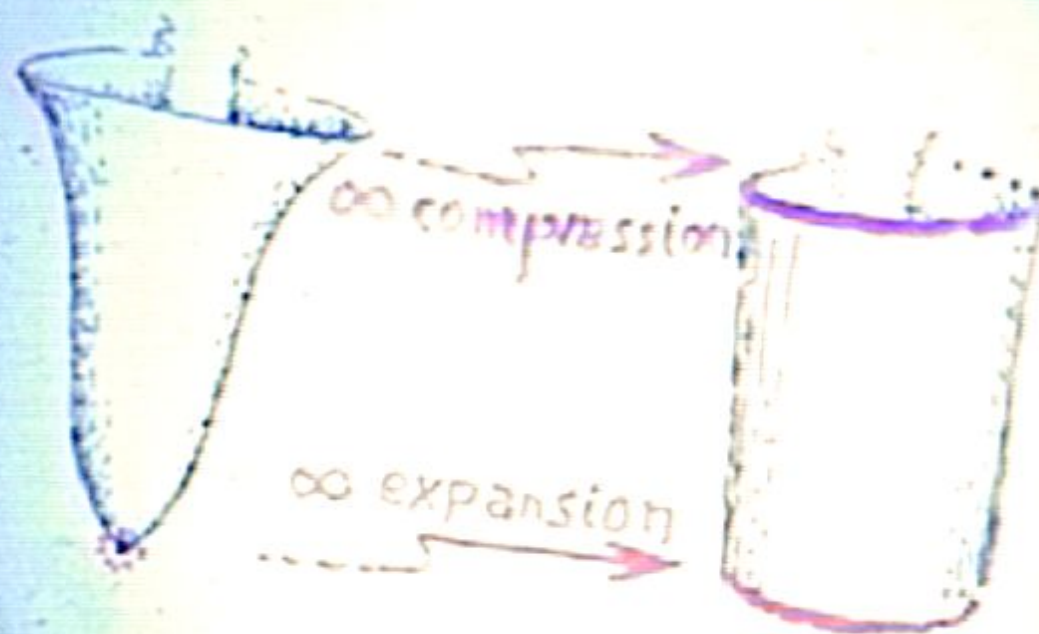


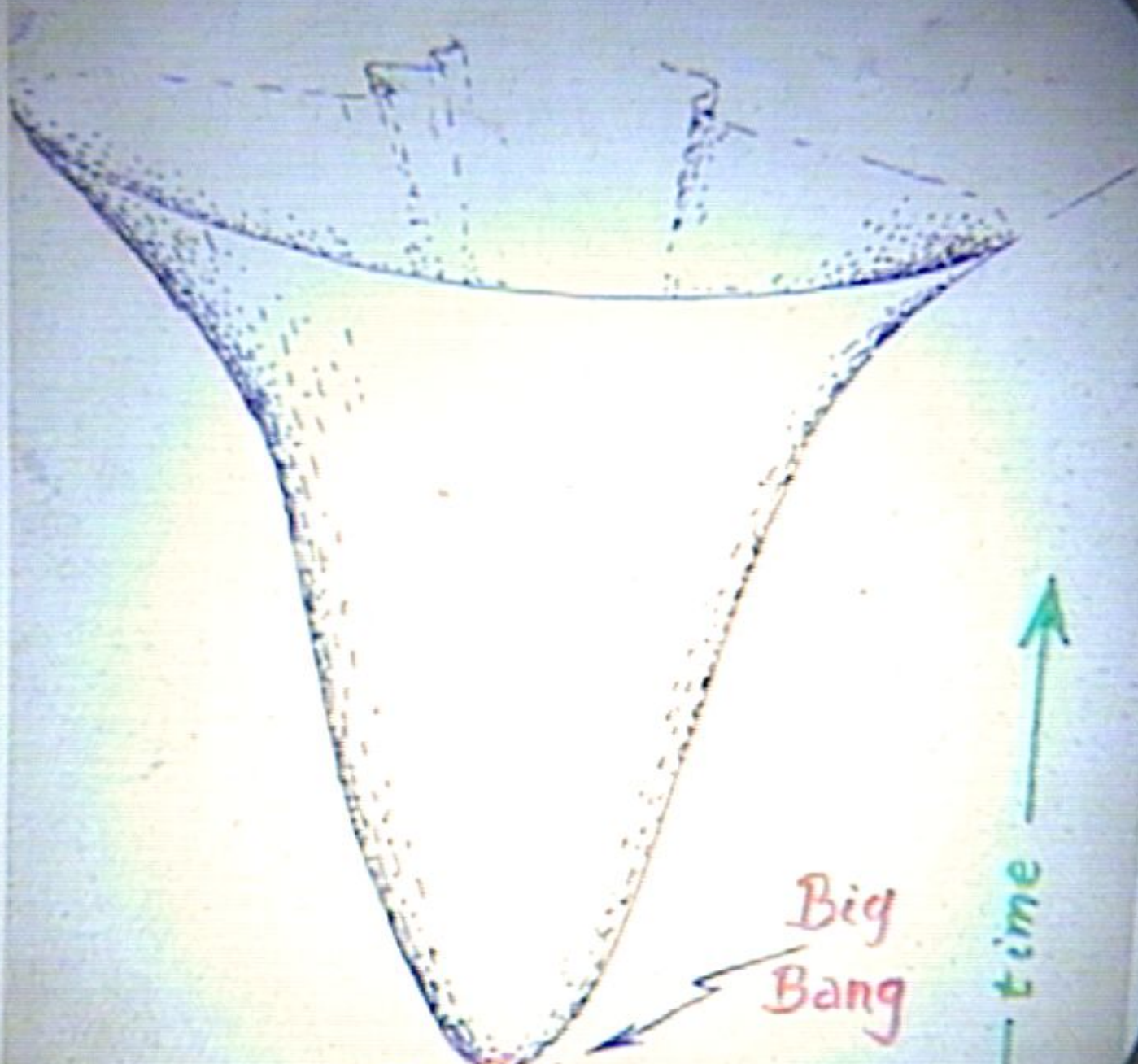
Just a mathematical trick!? Temperatures become so large at the Big Bang that rest masses become ignorable

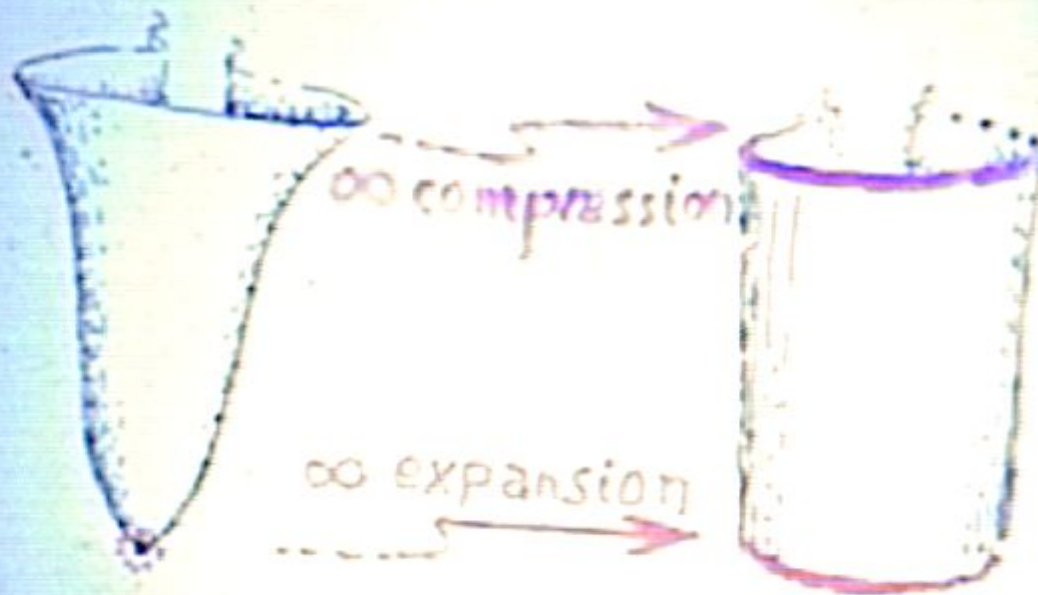
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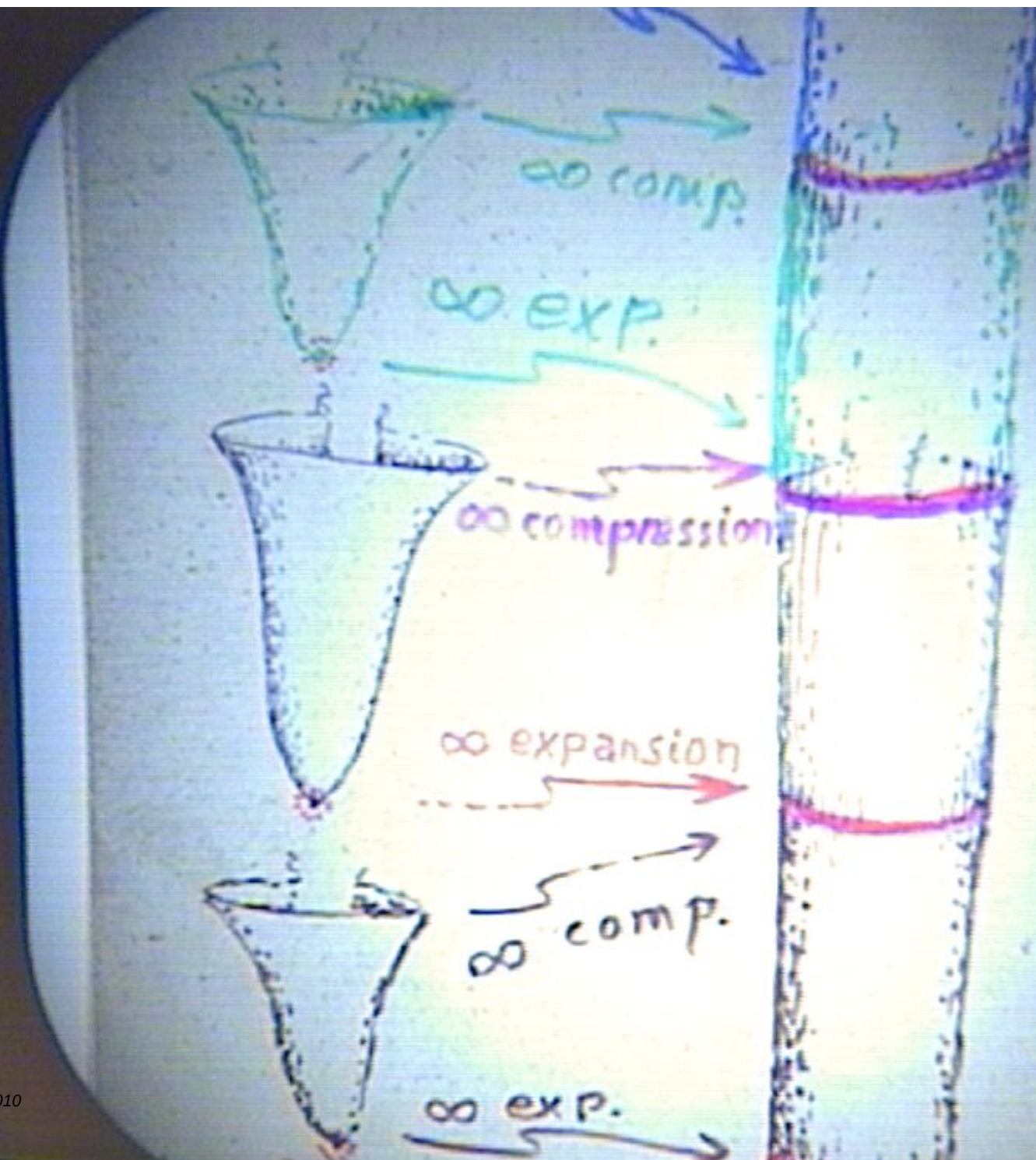


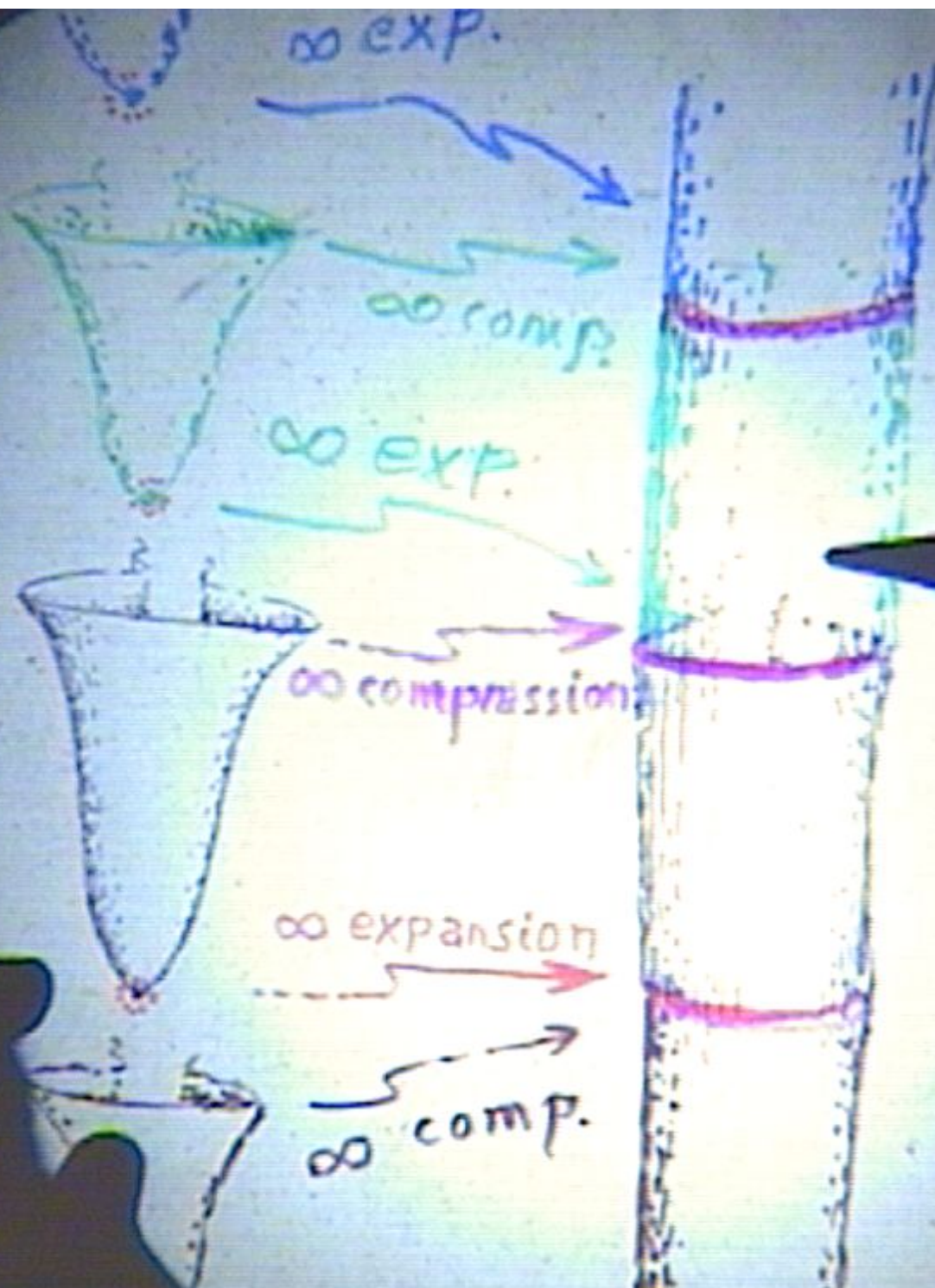


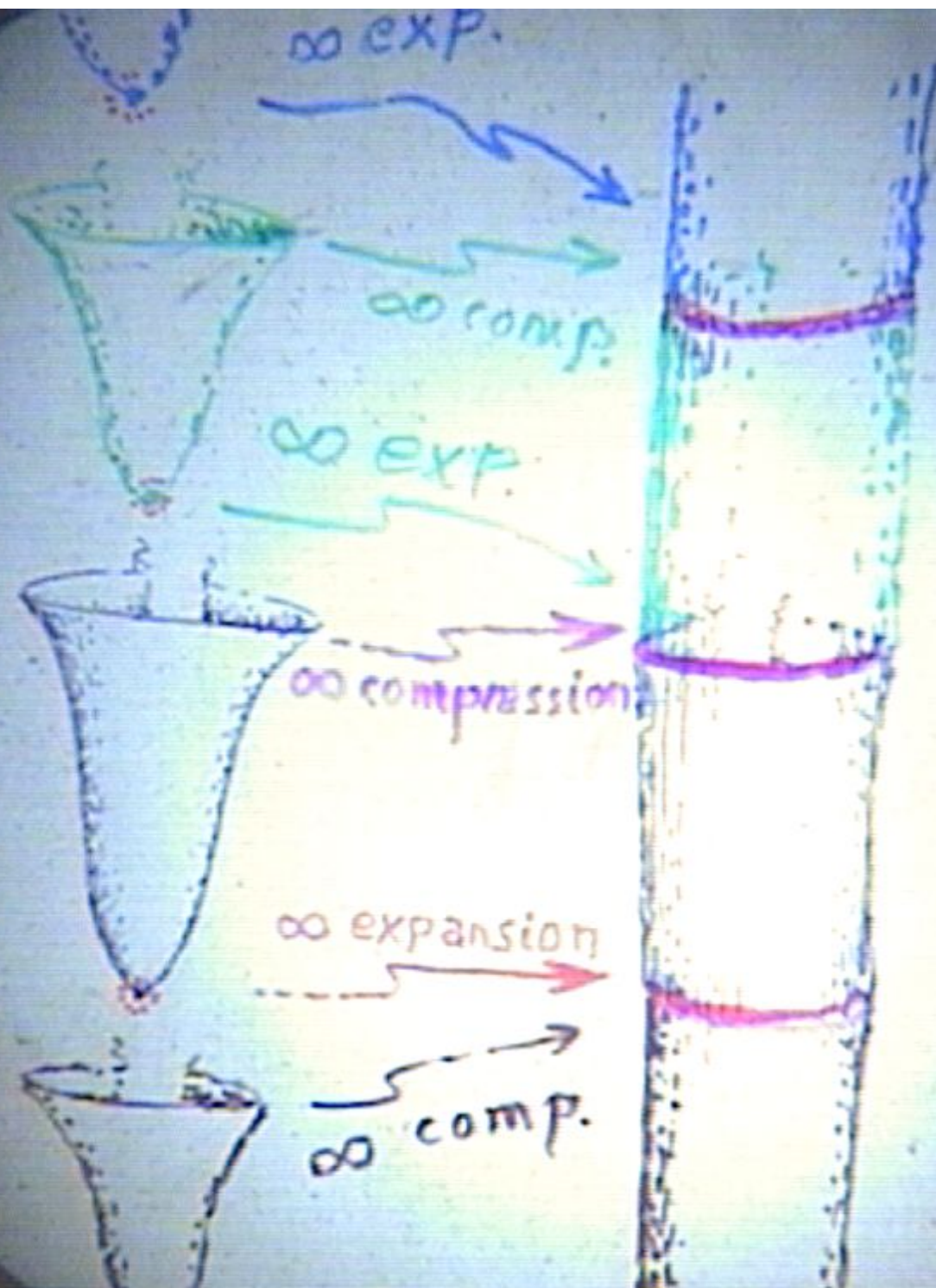












Proposal for the origin of:

- the extremely low entropy in gravitational degrees of freedom
- the detailed departures from uniformity

Conformal Cyclic Cosmology CCC

Claim: • & • both come from the evolution of purely CLASSICAL (non-quantum) p.d.e.s governing the Big Bang!

(all my friends working in QG)

But, you will say: surely, near the Big Bang, space-time radii of curvature will get down to Planck length, so quantum foam,

- in gravitational degrees of freedom
- the detailed departures from uniformity

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But, you will say: surely, near the Big Bang, space-time radii of curvature will get down to $\sim$ planck length, so quantum foam, space-time chaos, discrete space-time, non-commutative geometry, ... or ???
it can't have been like this, or

the detailed departures from uniformity

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I say: it can't have been like this, or the extremely low-entropy conformally invariant state must have arisen.

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I say: it can't have been like this, or
the extremely low-entropy conformally
smooth Big Bang would not have arisen.
Weyl curvature hypothesis tells us it's
radii of Weyl curvature that would need
get down to $\sim$ planck length for foam, chaos

the evolution of purely classical
(non-quantum) p.d.e.s governing
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All my friends
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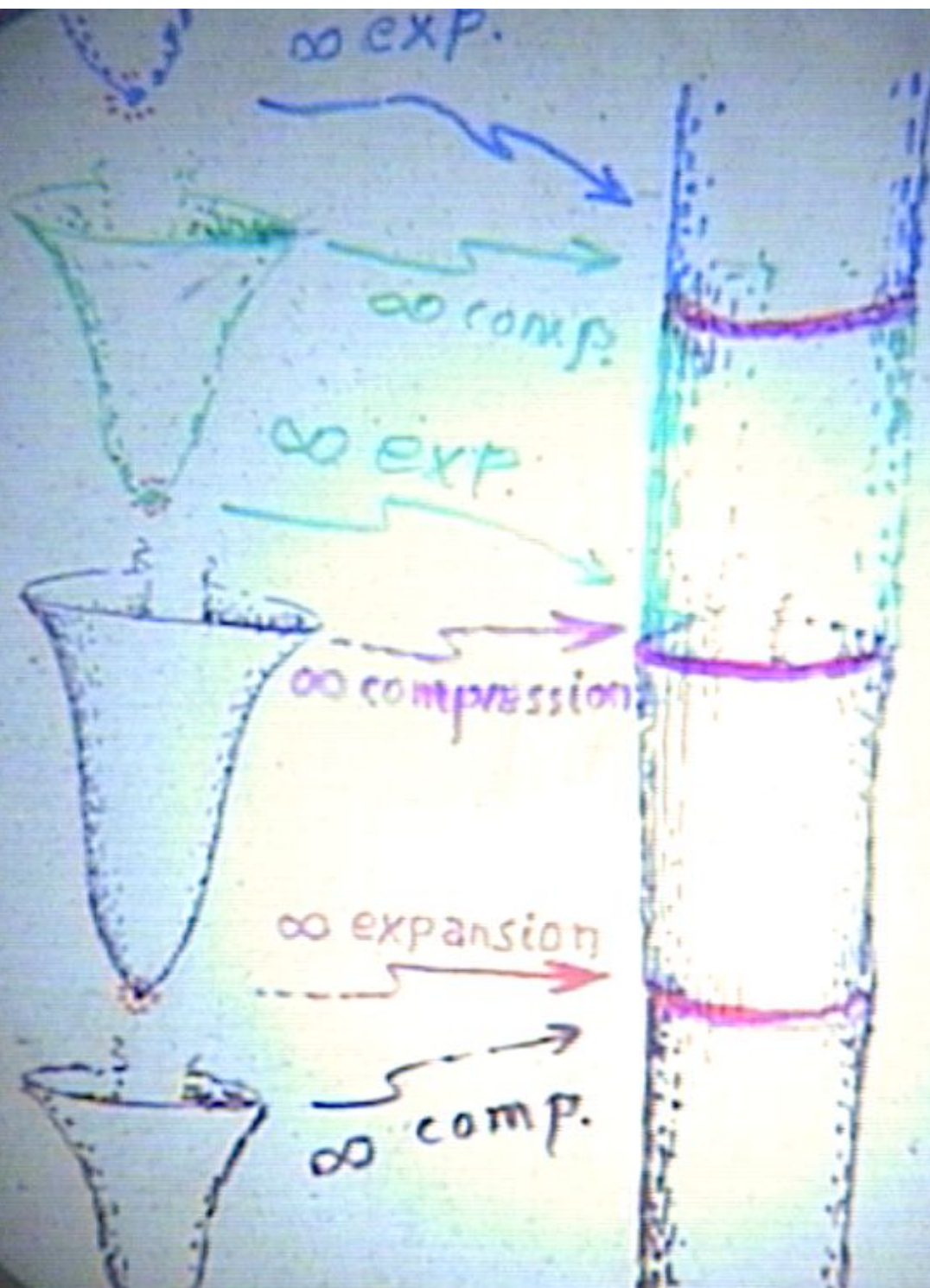
- the extremely low energy
- in gravitational degrees of freedom
- the detailed departures from uniformity

Conformal Cyclic Cosmology CCC

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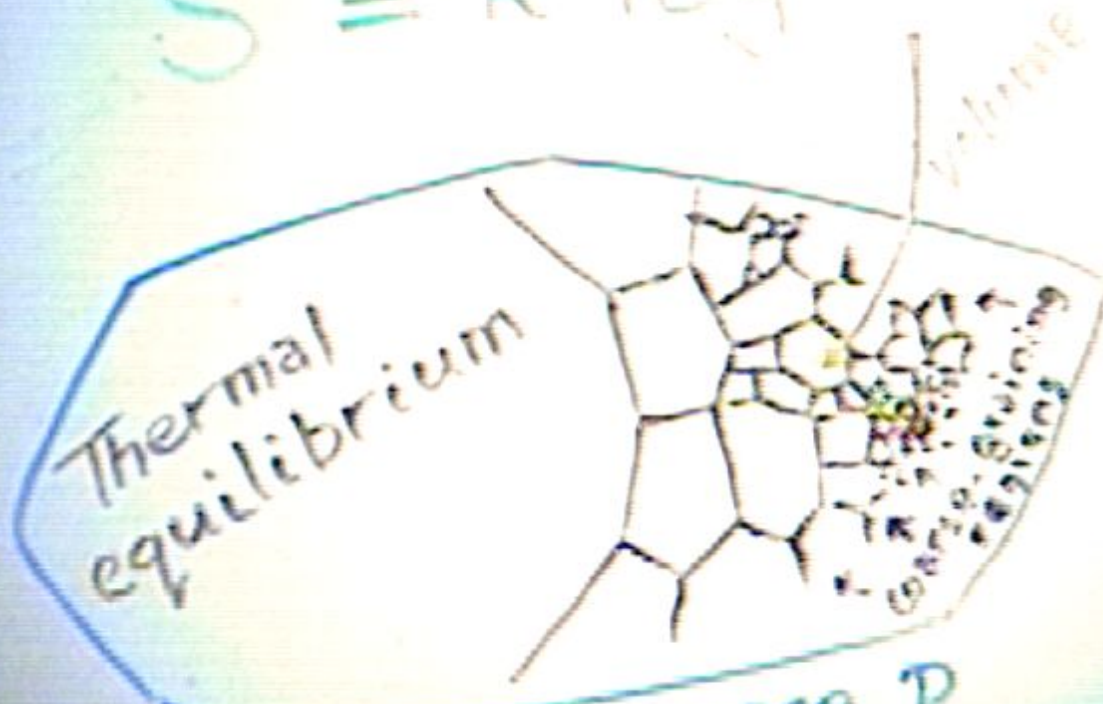
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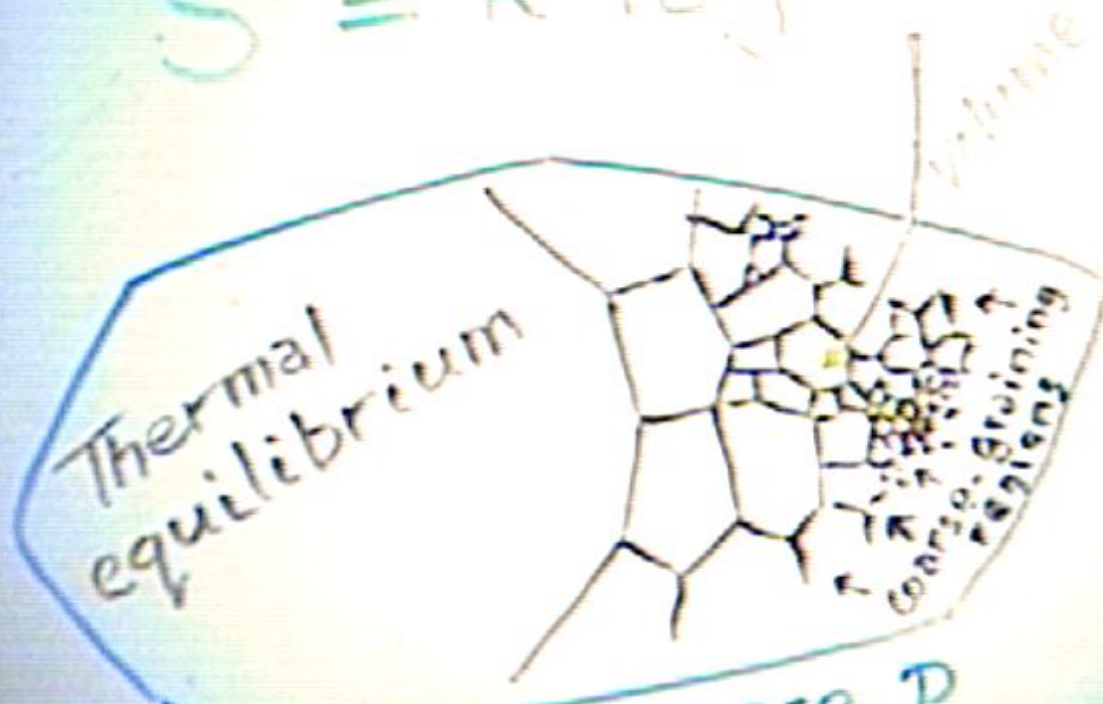
Boltzmann's definition of ENTROPY

$$S = k \log V$$

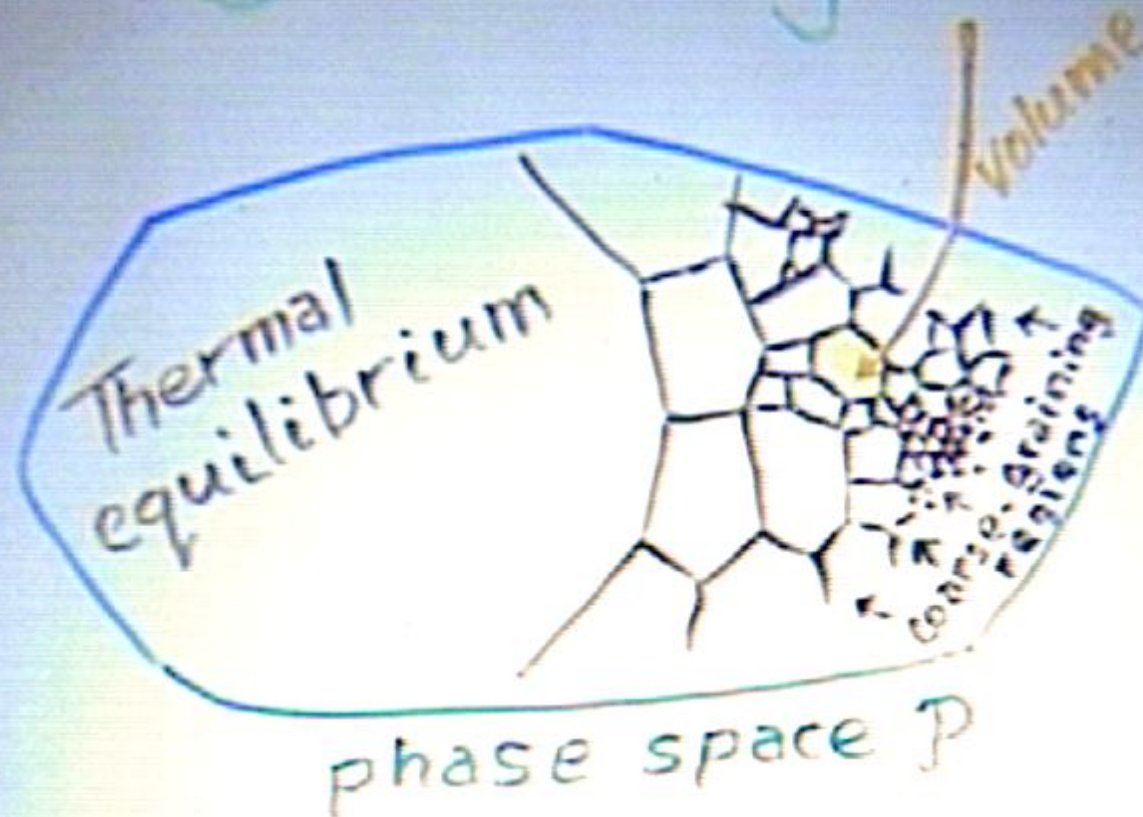


Boltzmann's definition of ENTROPY

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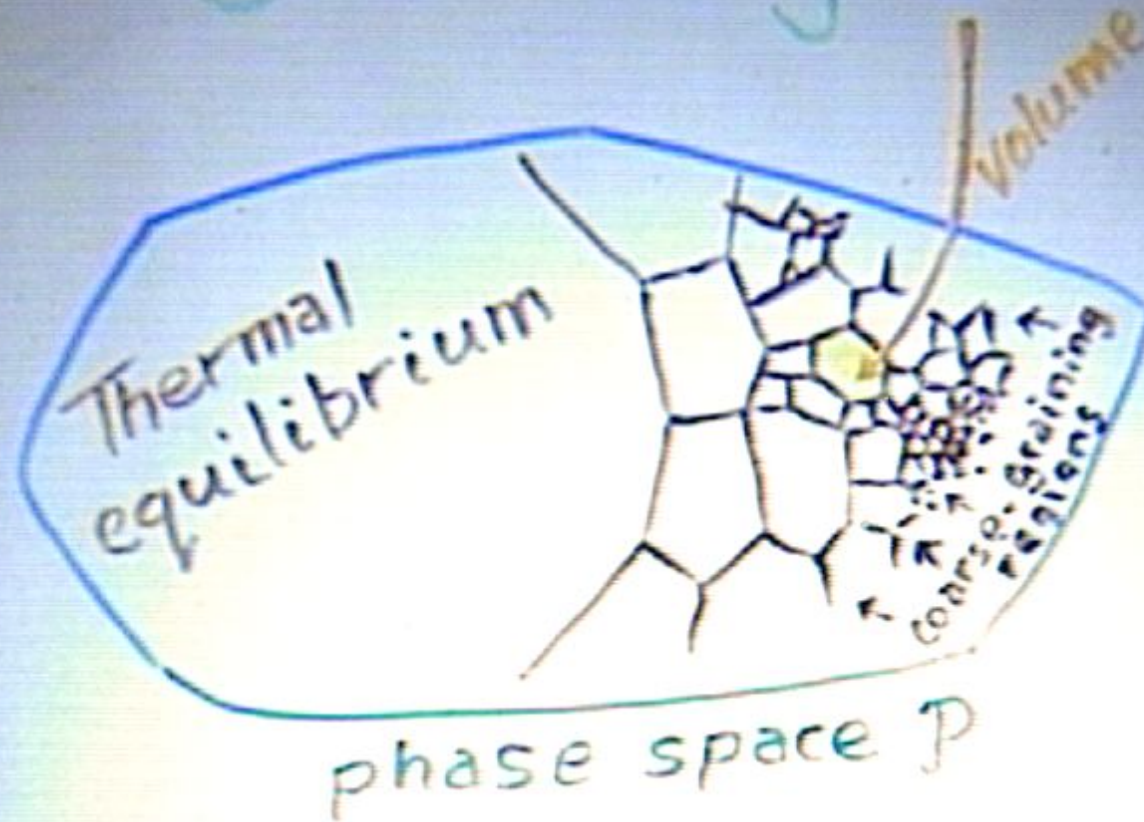
$$S = k \log V$$



Note: $\log(VW) = \log V + \log W$

so entropy is additive
in independent systems

$$S = k \log V$$



Note: $\log(VW) = \log V + \log W$

so entropy is additive
for independent systems

The second law of thermodynamics: how can this be squared with a cyclic universe?

Information (i.e. phase-space volume) gets lost at black holes' singularities

agrees with the young Hawking; disagrees with the old Hawking!

Phase-space volume shrinks enormously with each hole's disappearance

can this be squared with
a cyclic universe?

Information (i.e. phase-
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Phase-space volume
shrinks enormously
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zero of entropy gets
re-set! In effect:

a cyclic universe?

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Phase-space volume shrinks enormously with each hole's disappearance

zero of entropy gets re-set! In effect:

the second law is never violated despite cyclic,

Information (i.e. phase space volume) gets lost at black holes' singularities

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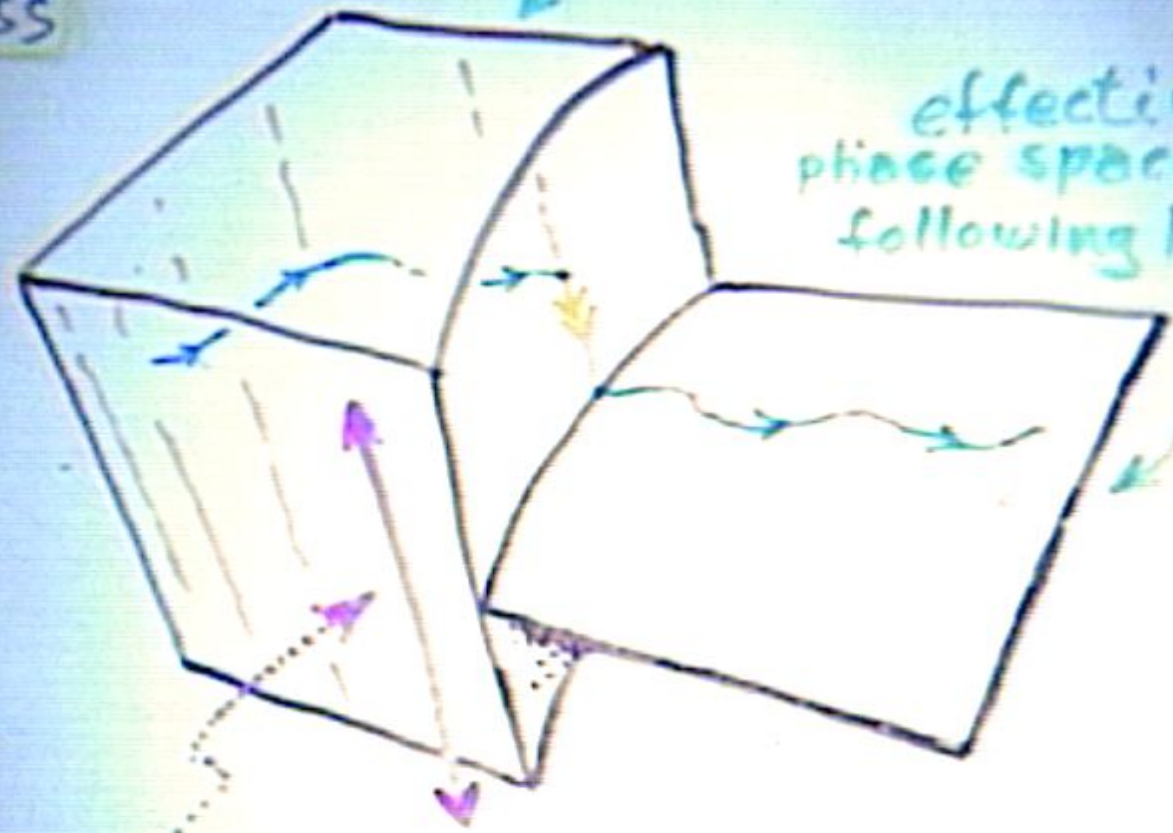
zero of entropy gets re-set! In effect:

the second law is never violated, despite cyclic nature of CCC ("transcended").

phase space
with info.
loss

phase space
prior to info.
loss

effective
phase space
following loss



degrees of freedom lost
in the black hole

pop!

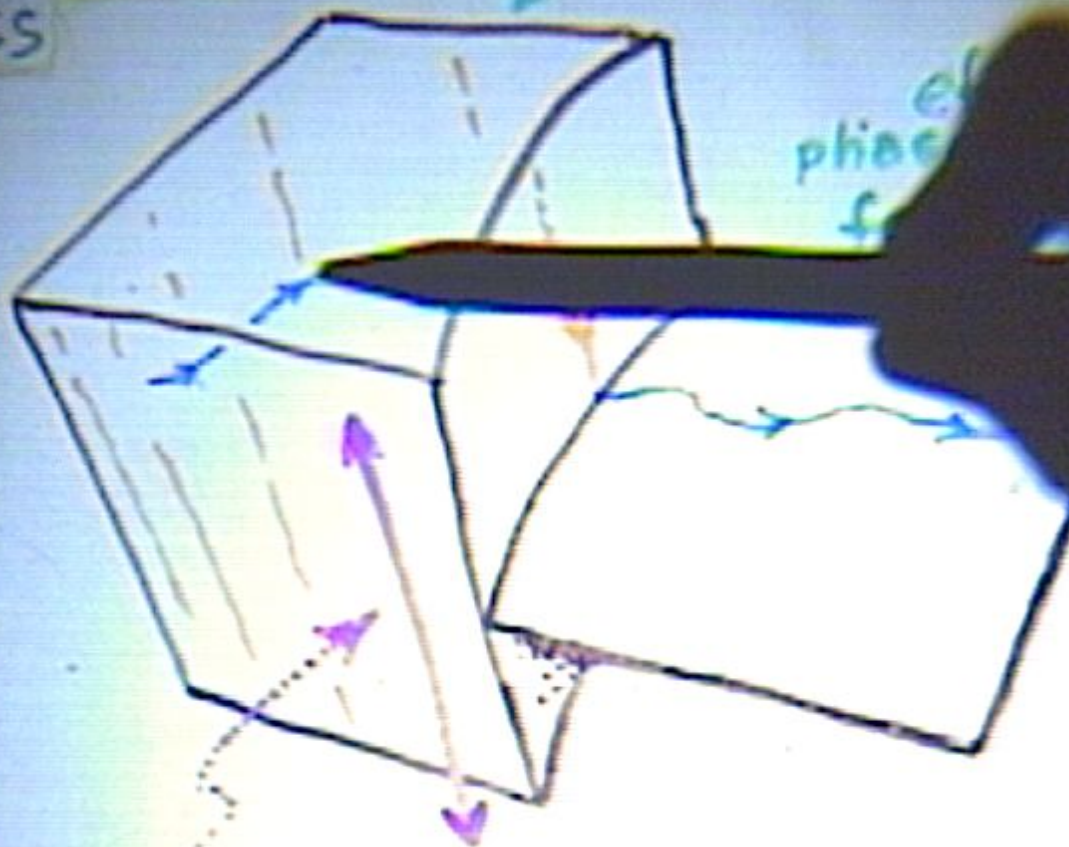
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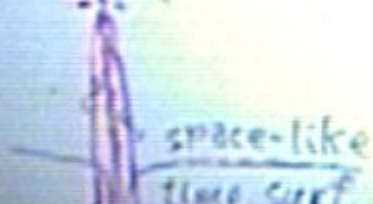
degrees of freedom lost
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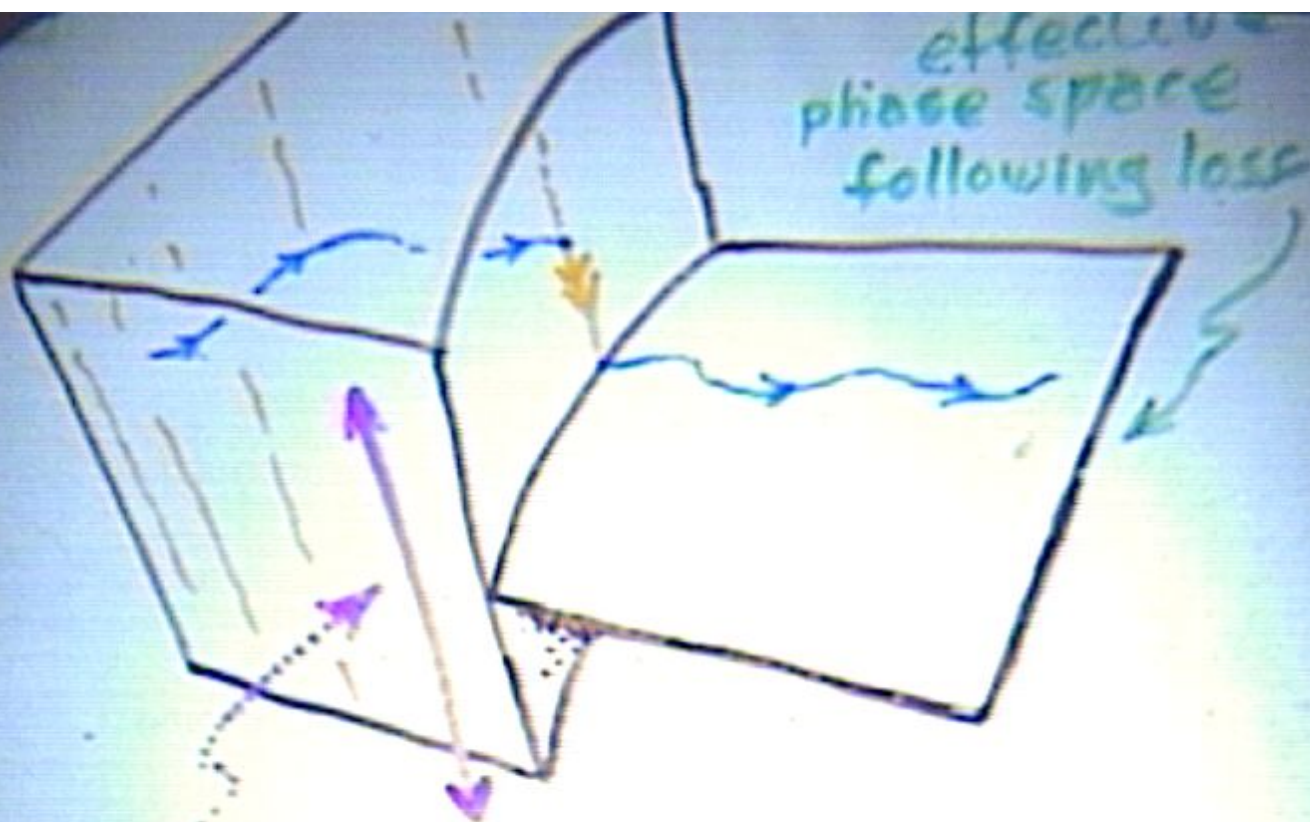
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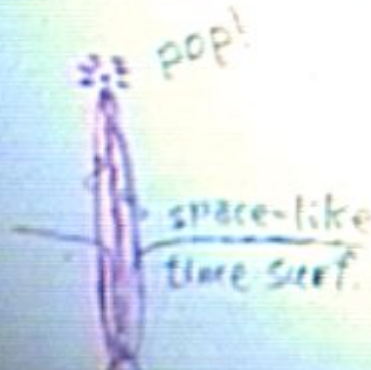
degrees of freedom lost
in the black hole.

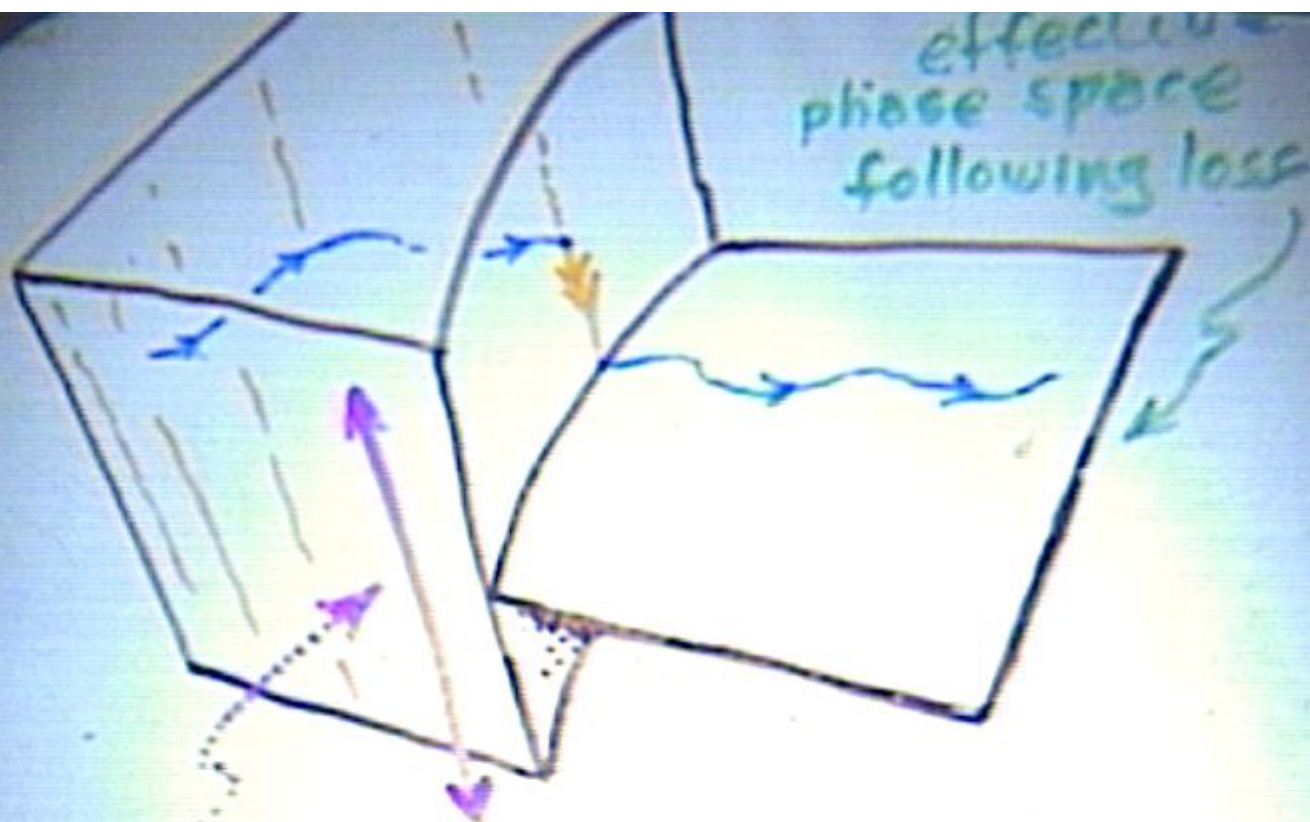
pop!





degrees of freedom lost in the black hole



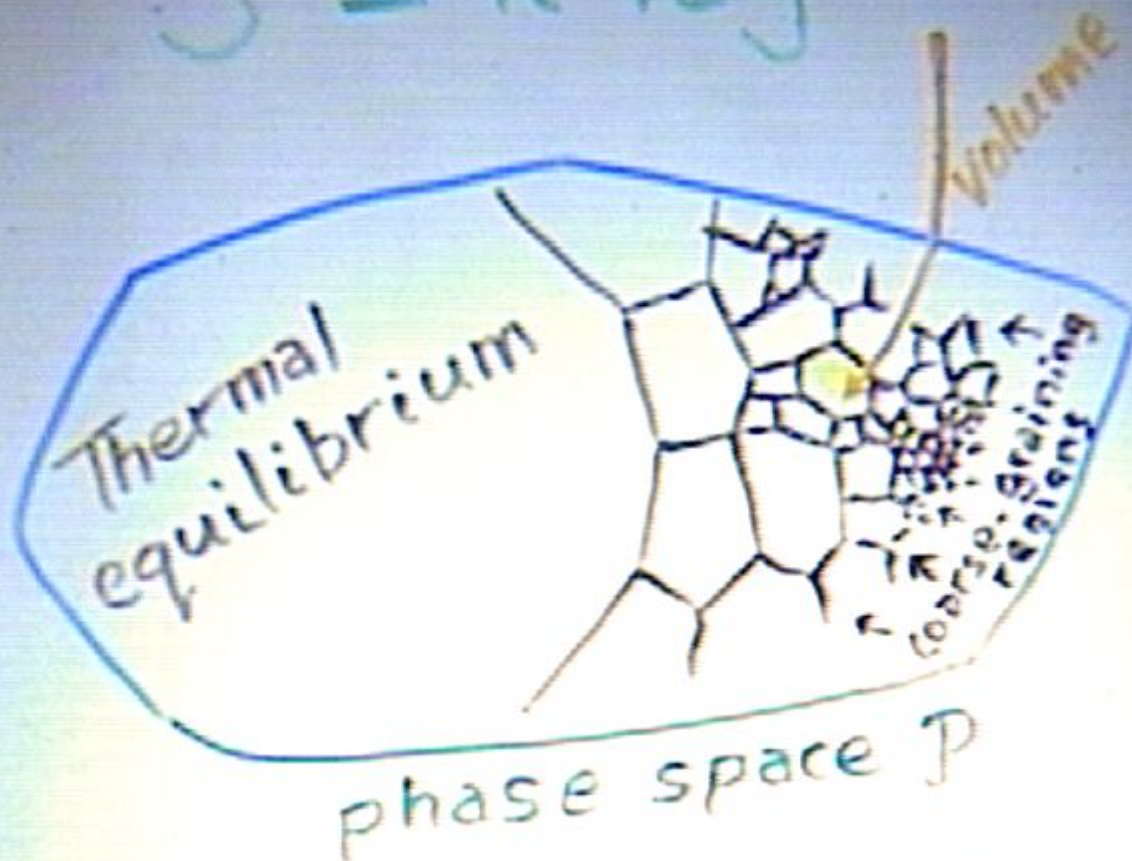


effective
phase space
following loss

degrees of freedom lost
in the black hole



$$S = k \log V$$



Note: $\log(VW) = \log V + \log W$
 so entropy is additive
 independent systems

Top stretched
big bang of
next æon

(Conformal
infinity ($\Lambda > 0$)
of previous æon



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At $\mathcal{X}$, $\Omega \rightarrow \infty$

$\omega \rightarrow 0$ (ω smooth across $\mathcal{X}$)

In $\hat{\mathcal{X}}$, we assume

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$\check{\nu}$
or
 Λ

Tod stretched
big bang of
next aeon

Conformal
infinity ($\Lambda > 0$)
of previous aeon



$\check{g}_{ab}$

$\hat{g}_{ab}$

g_{ab}

$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At $\mathcal{H}$, $\Omega \rightarrow \infty$

$\omega \rightarrow 0$

In $\check{\mathcal{E}}$, we assume

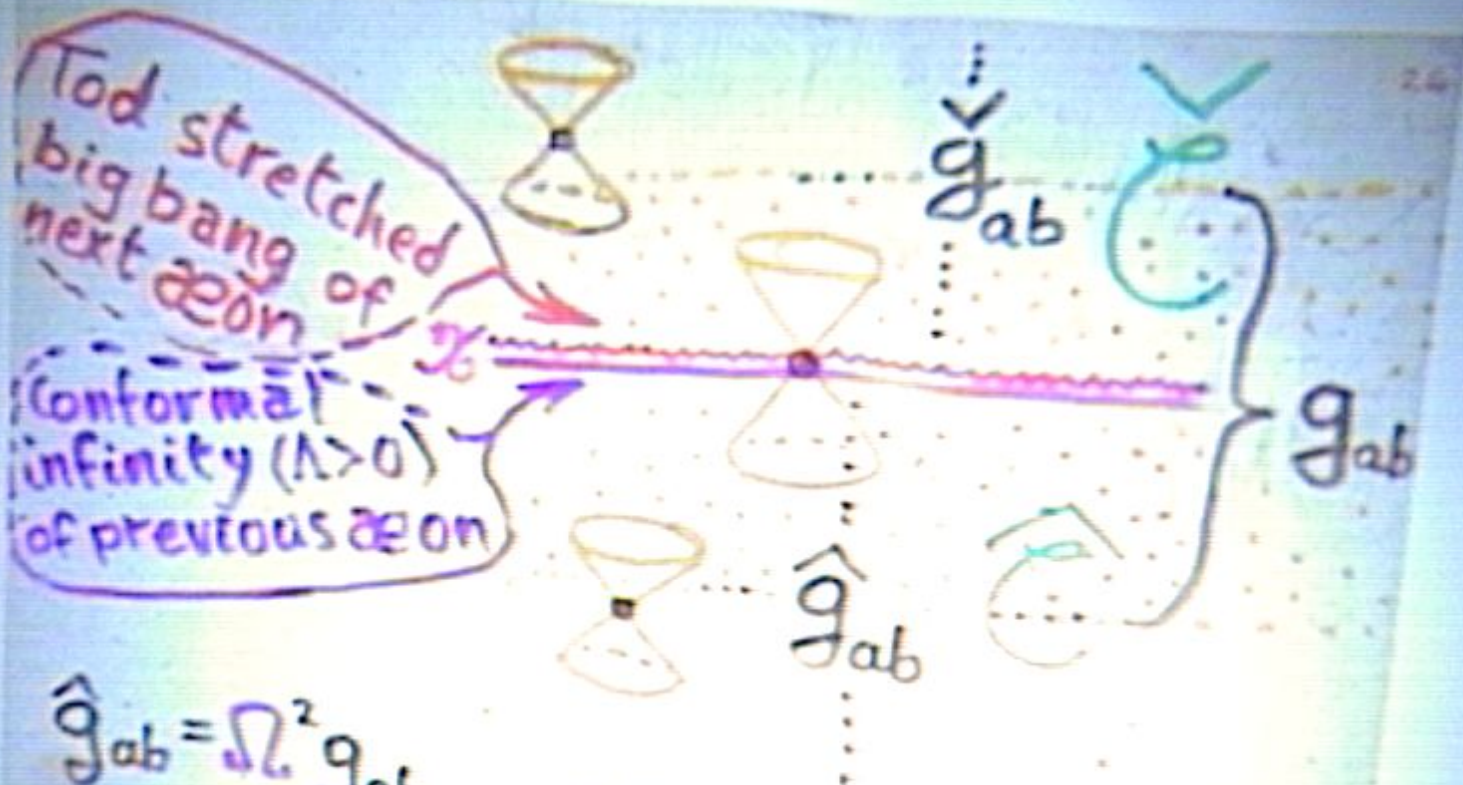
Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab}$$

$$= 8\pi G T_{ab} + \Lambda g_{ab}$$

$\check{\omega}$
or
 Λ

(ω smooth across $\mathcal{H}$)



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At x , $\Omega \rightarrow \infty$

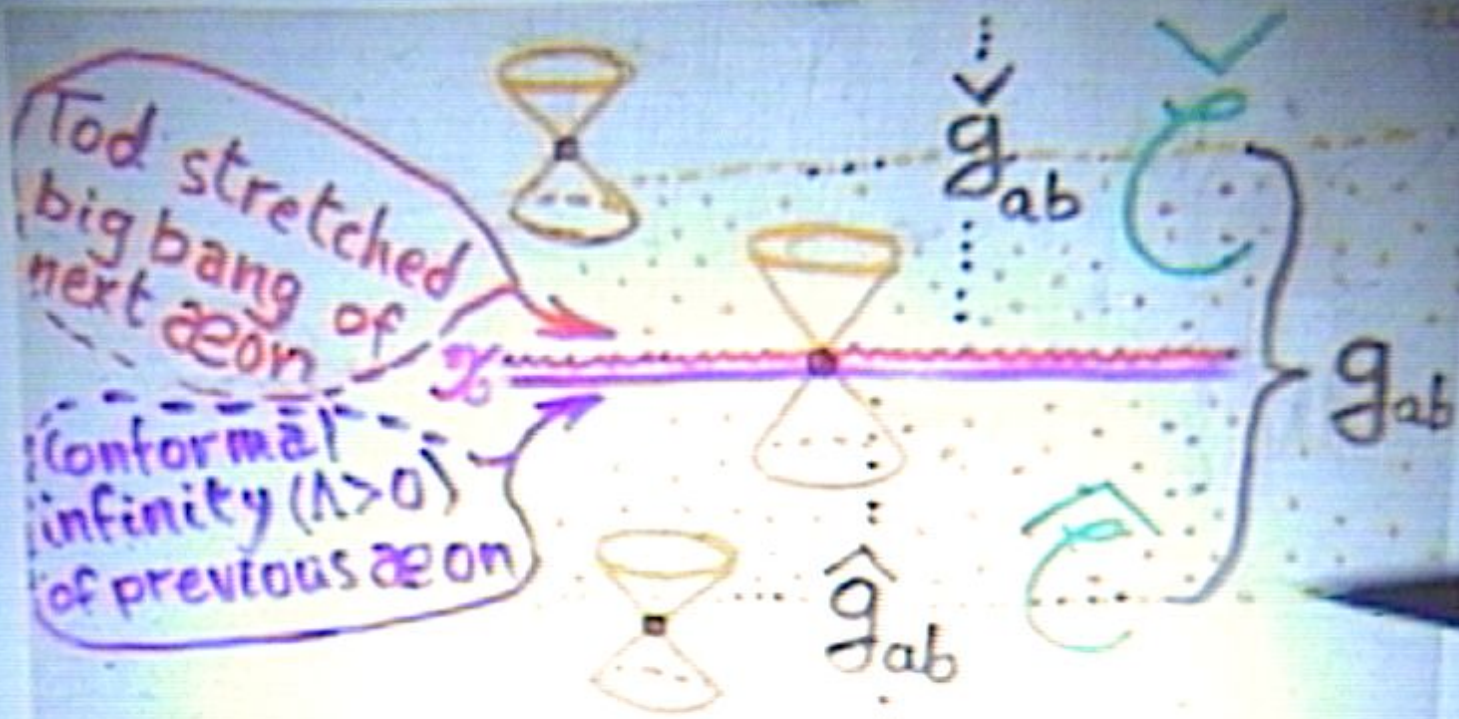
$\omega \rightarrow 0$ (ω smooth across x)

In $\hat{x}$, we assume

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$\check{\nu}$
or
 Λ



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At $\mathcal{X}$, $\Omega \rightarrow \infty$

$\omega \rightarrow 0$ (ω smooth across $\mathcal{X}$)

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$\check{\nu}$
 or
 Λ

In $\hat{\mathcal{E}}$, we assume that all rest-mass has died away (so that $\hat{g}_{ab}$ is smooth across $\mathcal{X}$)

infinity ($\Lambda > 0$)
of previous region



$\hat{g}_{ab}$



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At $\mathcal{H}$, $\Omega \rightarrow \infty$

$\omega \rightarrow 0$ (ω smooth across $\mathcal{H}$)

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$\checkmark$
or
 Λ

In $\hat{\mathcal{E}}$, we assume that all rest-mass has died away (CCC hypothesis) so that $\hat{T}_a^a = 0$, i.e. $\hat{R} = 4\Lambda$.

In $\hat{\mathcal{E}}$, $\check{\mathcal{E}}$, we find that the conf. invar. fields in $\hat{T}_{ab}$ scale, so that

$$\hat{T}_{ab} = \Omega^{-2} T_{ab}$$

$$\check{T}_{ab} = \omega^{-2} T_{ab}$$

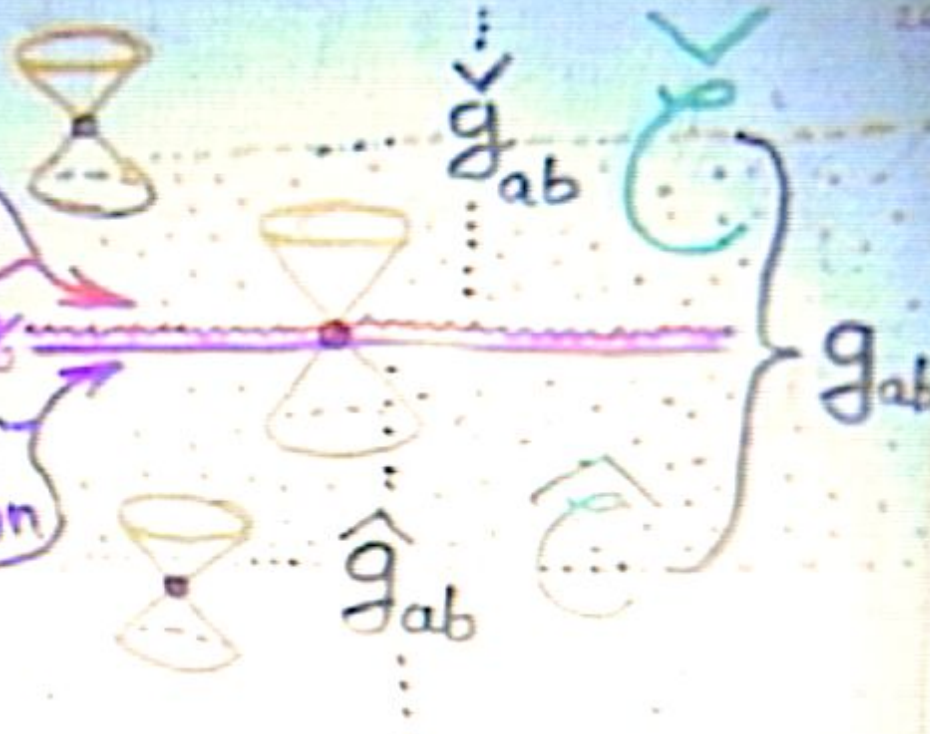
Preserves

$$\nabla^a T_{ab} = 0$$

$\checkmark$
or
 Λ

Top stretched
big bang of
next aeon

Conformal
infinity ($\Lambda > 0$)
of previous aeon



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

$$\check{g}_{ab} = \omega^2 g_{ab}$$

At $\mathcal{X}$, $\Omega \rightarrow \infty$

$\omega \rightarrow 0$ (ω smooth across $\mathcal{X}$)

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$\check{\Lambda}$
or
 Λ

But we find that, in $\check{\mathcal{E}}$, $\check{T}_{ab}$ cannot be the only contribution to the energy tensor $\check{U}_{ab}$ (as I shall call it); Einstein eqns. are:

$$\hat{\mathcal{G}}: \hat{E}_{ab} = 8\pi G \hat{T}_{ab} + \Lambda \hat{g}_{ab}$$

$$\check{\mathcal{G}}: \check{E}_{ab} = 8\pi G \check{U}_{ab} + \Lambda \check{g}_{ab}$$

Reciprocal hypothesis:

$$\omega = -\Omega^{-1}$$

we find that $\check{U}_{ab}$ necessarily acquires a trace in $\check{\mathcal{E}}$: $\check{U}_a^a = \mu$, so whereas $\hat{R} = 4\Lambda$, we have $\check{R} = 4\Lambda + 8\pi G\mu$ where does all this come from. Consider $\check{\mathcal{P}}$ first. How

$$E_{ab} = 8\pi G T_{ab} + \Lambda g_{ab}$$

$$\check{E}_{ab} = 8\pi G \check{U}_{ab} + \Lambda \check{g}_{ab}$$

Reciprocal hypothesis:

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we find that $\check{U}_{ab}$ necessarily acquires a trace in $\check{E}$: $\check{U}_a^a = \mu$, so whereas $\hat{R} = 4\Lambda$, we have $\check{R} = 4\Lambda + 8\pi G\mu$

Where does all this come from? Consider $\check{E}$ first. How do we write the Ein. eqns. in the $\check{g}_{ab}$ metric? If we want $\check{g}$ to have the same "cosmol. const." Λ , i.e. $\check{R} = 4\Lambda$ (as well as $\hat{R} = 4\Lambda$), then we find that $\omega = \Omega$ sats. the

"~~Ein~~-eqn": $(\square + \frac{R}{2})\omega = \frac{2}{3}\Lambda\omega^3$

$$\omega = -\Omega$$

we find that $\check{U}_{ab}$ necessarily acquires a trace in $\check{\epsilon}$: $\check{U}_a^a = \mu$, so whereas $\hat{R} = 4\Lambda$, we have $\check{R} = 4\Lambda + 8\pi G\mu$

Where does all this come from? Consider $\check{\epsilon}$ first. How do we write the Ein. eqns. in the g_{ab} metric? If we want g to have the same "cosmol. const." Λ , i.e. $R = 4\Lambda$ (as well as $\hat{R} = 4\Lambda$), then we find that $\omega = \Omega$ sats. the " ω -eqn.":

$$\left(\square + \frac{R}{6}\right)\omega = \frac{2}{3}\Lambda\omega^3$$

we find that $\check{U}_{ab}$ necessarily
 acquires a trace in $\check{\mathcal{E}}$: $\check{U}_a^a = \mu$, so
 whereas $\hat{R} = 4\Lambda$, we have $\check{R} = 4\Lambda + 8\pi G\mu$

Where does all this come
 from? Consider $\check{\mathcal{E}}$ first. How
 do we write the Ein. eqns. in
 the geo metric? If we want g to
 have the same "cosmol. const." Λ ,
 i.e. $R = 4\Lambda$ (as well as $\hat{R} = 4\Lambda$), then
 we find that $\check{\omega} = \Omega$ satis. the
 " $\check{\omega}$ -eqn.":

$$\left(\square + \frac{R}{6}\right) \check{\omega} = \frac{2}{3} \Lambda \check{\omega}^3$$

This is a conf. invar. self-coupled scalar eqn., since if

$$\tilde{\phi} = \tilde{\Omega}^{-1} \phi,$$

then

$$\left(\tilde{\square} + \frac{\tilde{R}}{6}\right) \tilde{\phi} = \tilde{\Omega}^{-3} \left(\square + \frac{R}{6}\right) \phi$$

The energy ("new improved") tensor for ϕ is: Callan Coleman
Jackiw
Newman-RP

$$T_{ab}[\phi] = \frac{1}{4\pi G} \phi^2 \left\{ \phi \nabla_{(a} \nabla_{b)} \phi^{-1} - \frac{1}{2} R_{(ab)} \right\}$$

Specific choice of const. made

$\langle \cdot \cdot \rangle =$ trace-removed symmetric part

Einstein eqns. in $\tilde{g}$, in terms of g_{ab} are:

$$T_{ab} = T_{ab}[\tilde{\Omega}]$$

Call $\tilde{\Omega}$ the "phantom field"

Ein. eqns. in g . metric take the form
"energy tensor of matter = energy tensor of phantom field"

This is a conf. invariant. self-coupled scalar eqn., since if

$$\tilde{\phi} = \tilde{\Omega}^{-1} \phi,$$

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Call $\tilde{\Omega}$ the "phantom field"

Einstein eqns. in g -metric take the form:

"energy tensor of matter = energy tensor of phantom field."

The reciprocal hypothesis

then $\tilde{\phi} = \tilde{\Omega}^{-1} \phi$,
 $(\tilde{\square} + \frac{\tilde{R}}{6}) \tilde{\phi} = \tilde{\Omega}^{-3} (\square + \frac{R}{6}) \phi$

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 Jakiw
 Newman-RP

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 tensor of phantom field."

The reciprocal hypothesis
 allows us to re-interpret the
 phantom field

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then

$$(\tilde{\square} + \frac{\tilde{R}}{6}) \tilde{\phi} = \tilde{\Omega}^{-3} (\square + \frac{R}{6}) \phi$$

The energy ("new improved") tensor for ϕ is:

Callan Coleman
Jackiw
Newman-RP

$$T_{ab}[\phi] = \frac{1}{4\pi G} \phi^2 \{ \nabla_a \nabla_b \phi^{-1} - \frac{1}{2} R_{ab} \}$$

specific choice
of const. made

$\langle \dots \rangle$: trace-removed
symmetric part

Einstein eqns. in $\tilde{g}$, in terms of g_{ab} are:

$$T_{ab} = T_{ab}[\Omega]$$

Call Ω the "phantom field"

Ein. eqns. in g . metric take the forms:

"energy tensor of matter = energy tensor of phantom field."

The reciprocal hypothesis allows us to re-interpret phantom

then $(\tilde{\square} + \frac{\tilde{R}}{6}) \tilde{\Phi} = \tilde{\Omega}^{-3} (\square + \frac{R}{6}) \Phi$

The energy ("new improved") tensor for Φ is: Callan Coleman
Jackiw
Newman-RP

$$T_{ab}[\Phi] = \frac{1}{4\pi G} \Phi^2 \left\{ \Phi \nabla_{(a} \nabla_{b)} \Phi^{-1} - \frac{1}{2} R_{(ab)} \right\}$$

Specific choice of const. made $\langle \dots \rangle =$ trace-removed symmetric part

Einstein eqns. in $\tilde{g}$, in terms of g_{ab} are:

$$T_{ab} = T_{ab}[\tilde{\Omega}]$$

Call $\tilde{\Omega}$ the "phantom field"

Ein. eqns. in $g_{..}$ metric take the form:
"energy tensor of matter = energy tensor of phantom field."

The reciprocal hypothesis allows us to re-interpret the phantom field as a real field in $\tilde{g}$ which satisfies

$$\hat{g}: \hat{E}_{ab} = 8\pi G \hat{T}_{ab} + \Lambda \hat{g}_{ab}$$

$$\check{g}: \check{E}_{ab} = 8\pi G \check{U}_{ab} + \Lambda \check{g}_{ab}$$

Reciprocal hypothesis:

$$\omega = -\Omega^{-1}$$

we find that $\check{U}_{ab}$ necessarily acquires a trace in $\check{E}$: $\check{U}_a^a = \mu$, so whereas $\hat{R} = 4\Lambda$, we have $\check{R} = 4\Lambda + 8\pi G\mu$

Where does all this come from? Consider $\check{E}$ first. How do we write the Ein. eqns. in the $\check{g}$ metric? If we want $\check{g}$ to have the same "cosmol. const." Λ , i.e. $\check{R} = 4\Lambda$ (as well as $\hat{R} = 4\Lambda$), then we find that $\check{\alpha} = \Omega$ sats. the

$$\hat{\mathcal{E}}: \hat{E}_{ab} = 8\pi G \hat{T}_{ab} + \Lambda \hat{g}_{ab}$$

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However, we now find that the eqn. is the wrong way around to preserve the trace " $\check{U}_a^a = 4\check{\Lambda}$ " and instead we find that there is a rest-mass contribution to the total energy tensor in $\check{C}$

$$\check{U}_{ab} = \check{T}_{ab} + \check{V}_{ab} + \check{W}_{ab}$$

Total $\nearrow$ $\check{U}_{ab}$ = $\check{T}_{ab}$ (massless fields from $\check{C}$) + $\check{V}_{ab}$ (energy tensor of phantom field, because real) + $\check{W}_{ab}$ (remaining piece, which has the trace μ)

We find:

$$\mu = \frac{(\Omega^2 - 1)}{2\pi G} \{3\Pi^a \Pi_a - \Lambda\}$$

The 1-form

$$\Pi = \frac{d\Omega_i}{\Omega^2 - 1} = \frac{d\omega}{1}$$

seems important

to preserve the trace $U_a = 4\pi$
 and instead we find that there
 is a rest-mass contribution to the
 total energy tensor in $\check{\mathcal{E}}$

$$\check{U}_{ab} = \check{T}_{ab} + \check{V}_{ab} + \check{W}_{ab}$$

Total $\nearrow$ massless fields from $\check{\mathcal{E}}$ $\nearrow$ energy tensor of phantom field become real $\nearrow$ remaining piece, which has the trace μ $\nwarrow$

We find:

$$\mu = \frac{(\Omega^2 - 1)}{2\pi G} \{ 3\Pi^a \Pi_a - \Lambda \}$$

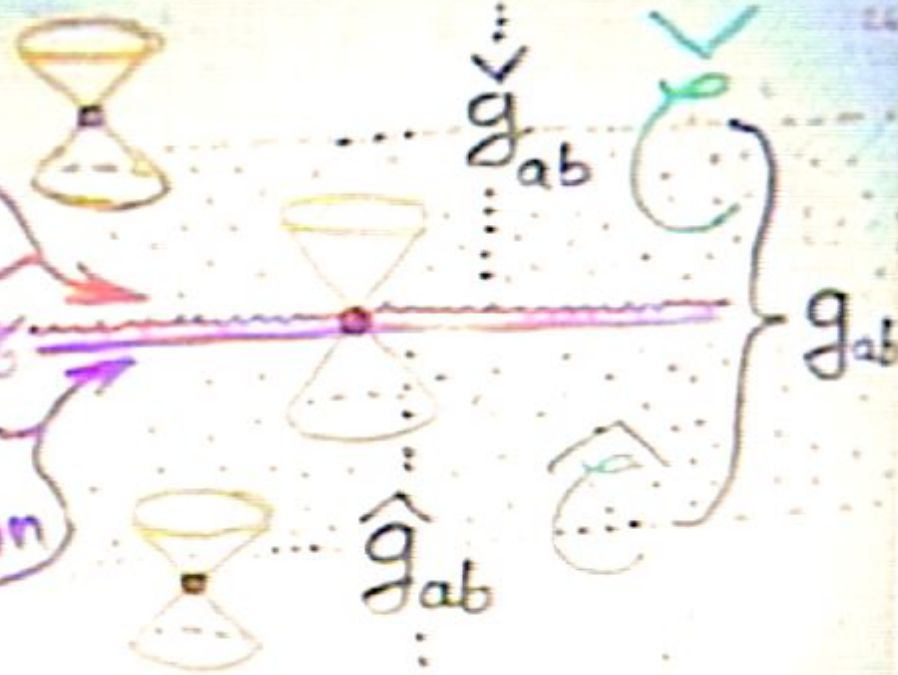
The 1-form

$$\Pi = \frac{d\Omega}{\Omega^2 - 1} = \frac{d\omega}{1 - \omega^2}$$

seems important. It remains smooth
 across $\mathcal{H}$ and is intimately related
 to the reciprocal property

$$\Pi = d\alpha$$

Tod stretched
 big bang of
 next æon
 Conformal
 infinity ($\Lambda > 0$)
 of previous æon



$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab}$$

V
 or
 Λ

Tod stretched
 big bang of
 next æon
 Conformal
 infinity ($\Lambda > 0$)
 of previous æon



g_{ab}

$\hat{g}_{ab}$

g_{ab}

$$\hat{g}_{ab} = \Omega^2 g_{ab}$$

Einstein eqns.

$$E_{ab} = \frac{1}{2} R g_{ab} - R_{ab}$$

V
or
 Λ

total energy tensor in $\mathcal{C}$

$$\check{U}_{ab} = \check{T}_{ab} + \check{V}_{ab} + \check{W}_{ab}$$

Total $\nearrow$ $\check{T}_{ab}$ massless fields from $\mathcal{C}$ $\nearrow$ energy tensor of photon fld. become real $\nearrow$ remaining piece, which has the trace μ

We find:

$$\mu = \frac{(\Omega^2 - 1)}{2\pi G} \{3\Pi^a \Pi_a - \Lambda\}$$

The 1-form

$$\Pi = \frac{d\Omega}{\Omega^2 - 1} = \frac{d\omega}{1 - \omega^2}$$

seems important. It remains smooth across $\mathcal{H}$ and is intimately related to the reciprocal proposal. Note:

$$\Pi = d\tau, \quad -\coth \tau = \Omega, \quad \tanh \tau = \omega$$

τ is a useful "time" parameter $(\tau < 0)$ $(\tau \geq 0)$

rest-mass contribution
total energy tensor in $\mathcal{C}$

$$\check{U}_{ab} = \check{T}_{ab} + \check{V}_{ab} + \check{W}_{ab}$$

Total $\nearrow$ massless fields from $\mathcal{C}$ $\nearrow$ energy tensor of photon fld. become real $\nearrow$ remaining piece, which has the trace μ

We find:

$$\mu = \frac{(\Omega^2 - 1)}{2\pi G} \{3\Pi^a \Pi_a - \Lambda\}$$

The 1-form

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$$\Pi = d\tau, \quad -\coth\tau = \Omega, \quad \tanh\tau = \omega$$

τ is a useful "time" parameter $(\tau < 0)$ $(\tau > 0)$

factor ω (or, equivalently

This fixing would appear
be satisfactorily
if we can force ω
2 scalar conditions
(since it sat's a 2nd order eq.
My guess for best bet is:

$$3 \Pi^a \Pi_a - \Lambda = O(\omega^3)$$

Since this delays the growing
of rest mass for as long as
possible, so it is most in
keeping with the CCC ph

factor ω (or, equivalently $\Omega = -\omega$)

This fixing would appear to be satisfactorily achieved if we can force ω to satisfy 2 scalar conditions over $\mathcal{X}$ (since it satis a 2nd-order eqn).

My guess for best bet is:

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This fixing would appear to be satisfactorily achieved if we can force ω to satisfy 2 scalar conditions over $\mathcal{X}$ (since it sat's a 2nd-order eqn).

My guess for best bet is:

$$3\Pi^a\Pi_a - \Lambda = 0(\omega^3)$$

Since this delays the growing of rest mass for as long as possible, so it is most in keeping with the CCC philosophy

... the trace $\bar{U}_a^a = 4\Lambda$
 instead we find that there
 is a rest-mass contribution to the
 total energy tensor in $\tilde{\mathcal{E}}$

$$\tilde{U}_{ab} = \tilde{T}_{ab} + \tilde{V}_{ab} + \tilde{W}_{ab}$$

Total $\nearrow$ massless fields from $\tilde{\mathcal{E}}$ $\nearrow$ energy tensor of phantom fld. become real $\nearrow$ remaining piece, which has the trace μ $\nwarrow$

We find:

$$\mu = \frac{(\Omega^2 - 1)}{2\pi G} \{3\Pi^a \Pi_a - \Lambda\}$$

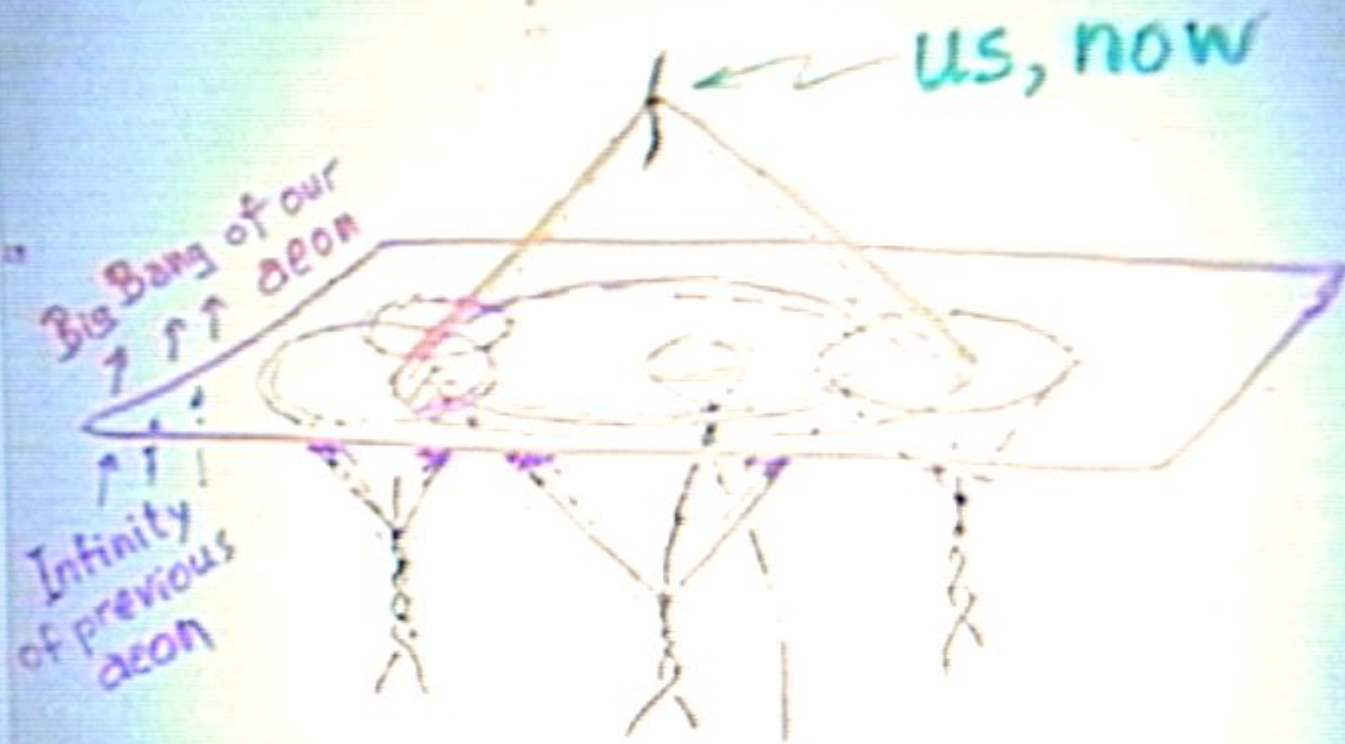
The 1-form

$$\Pi = \frac{d\Omega}{\Omega^2 - 1} = \frac{d\omega}{1 - \omega^2}$$

seems important. It remains smooth

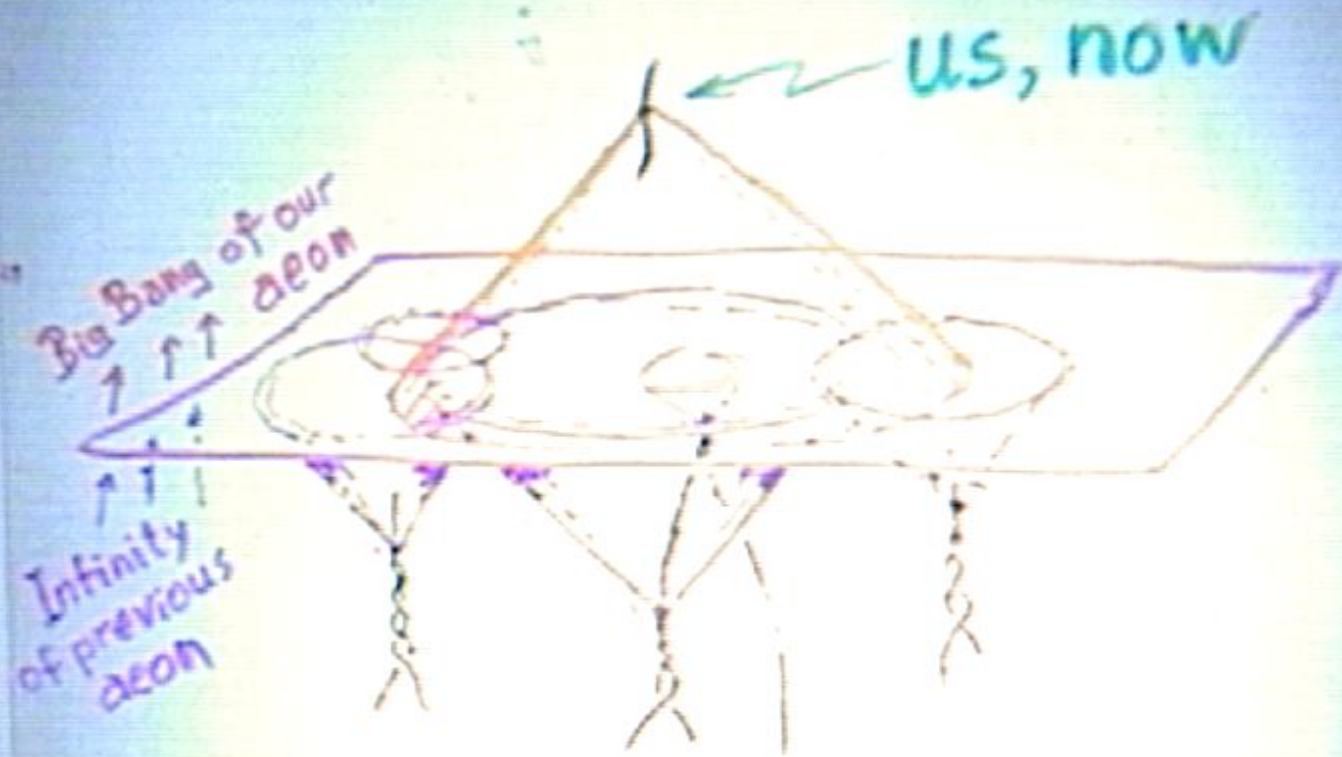
across $\mathcal{H}$ and is intimately related

Observational consequence concerning temperature/density variations in Cosmic Microwave Background



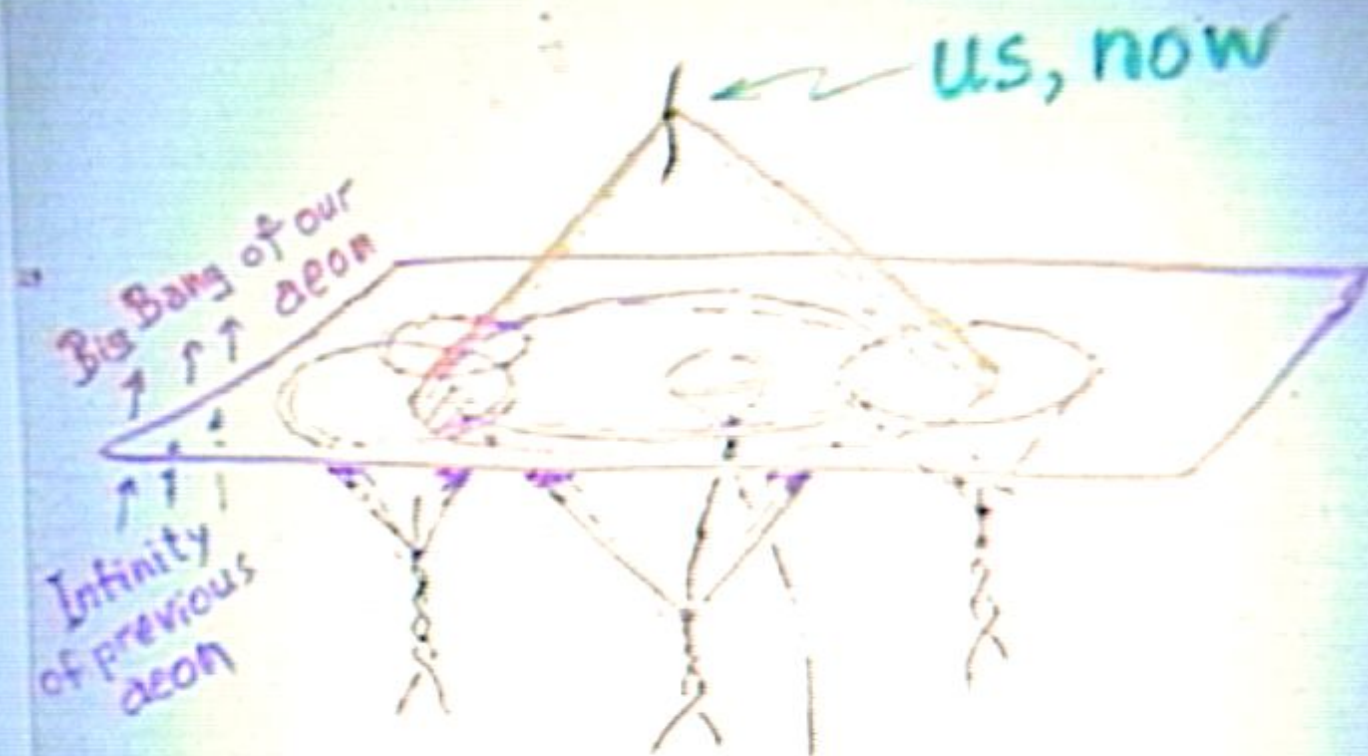
-Think of ripples on a pond, caused by raindrops which have recently stopped falling.
Pattern of ripples

Observational consequences
concerning temperature/density
variations in Cosmic Microwave Background

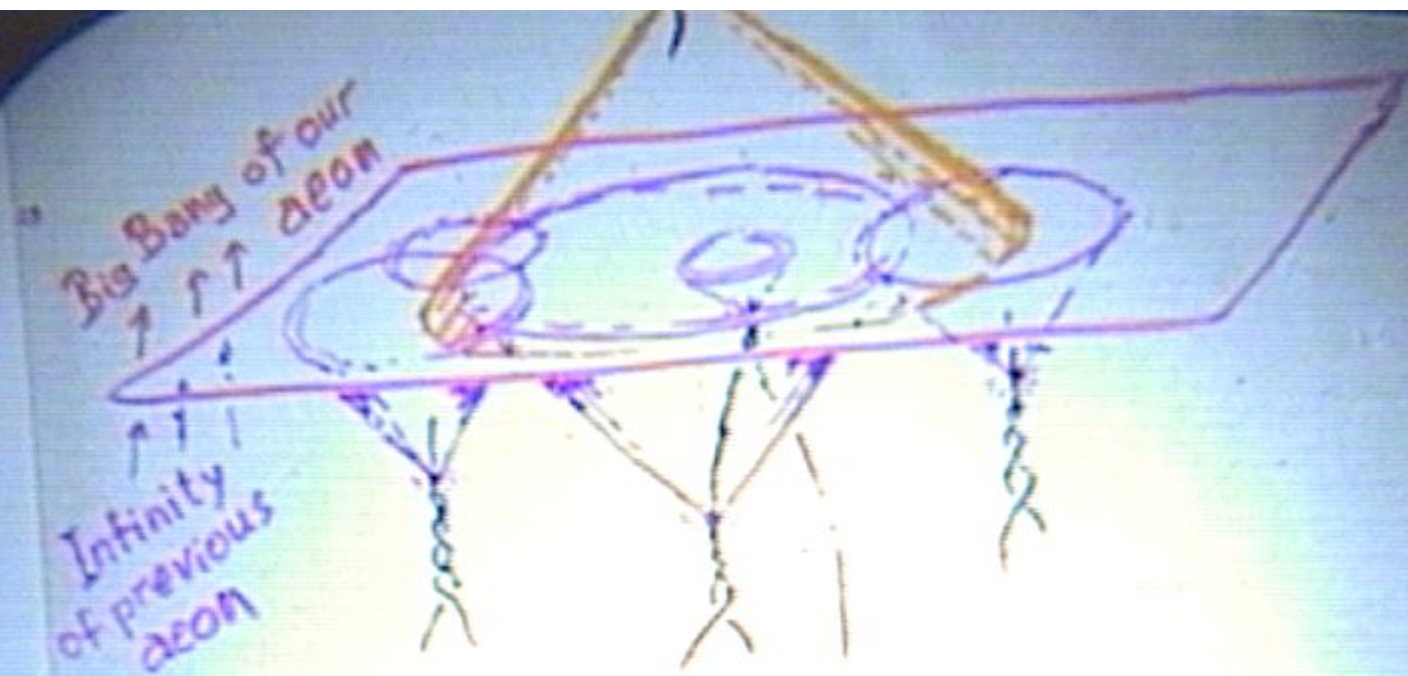


Think of ripples on a pond, caused
by raindrops which have recently
stopped falling.
Pattern of ripples
looks random

Observational consequences concerning temperature/density variations in Cosmic Microwave Background

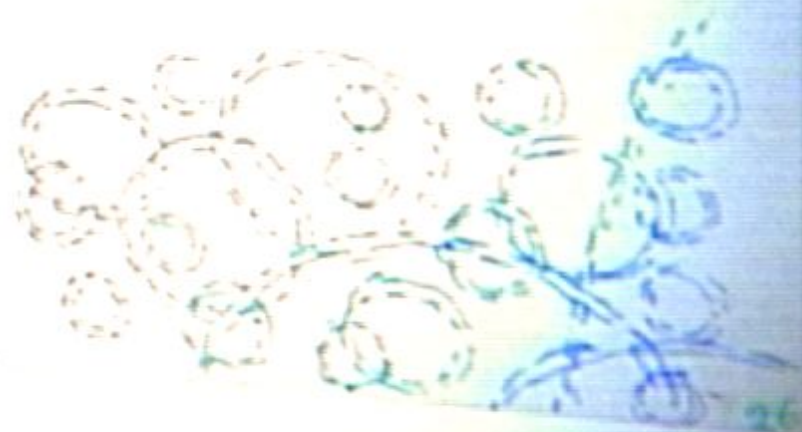


Think of ripples on a pond, caused by raindrops which have recently stopped falling.
Pattern of ripples

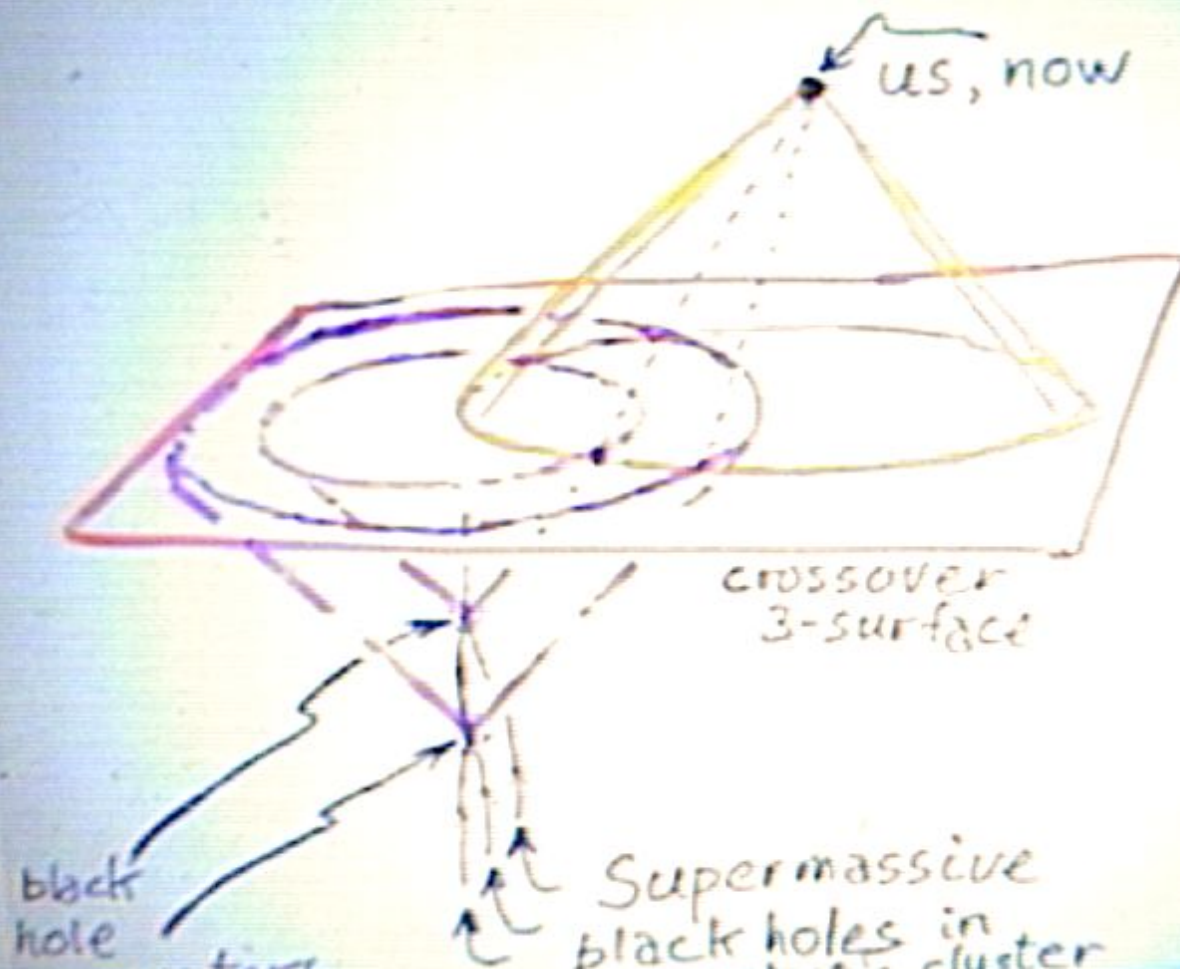


Think of ripples on a pond, caused by raindrops which have recently stopped falling.

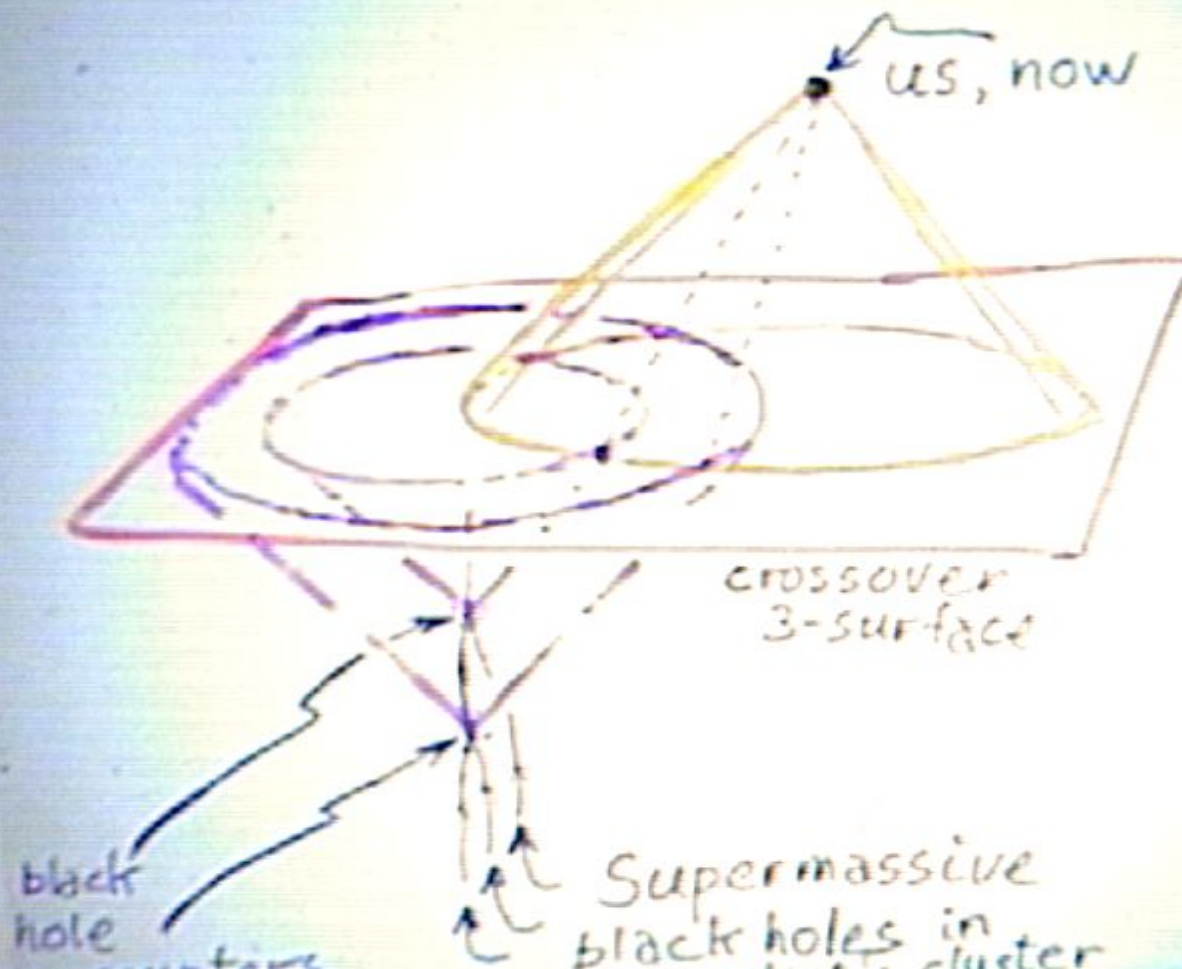
Pattern of ripples looks random at first, but can be analysed into circles by statistical analysis.



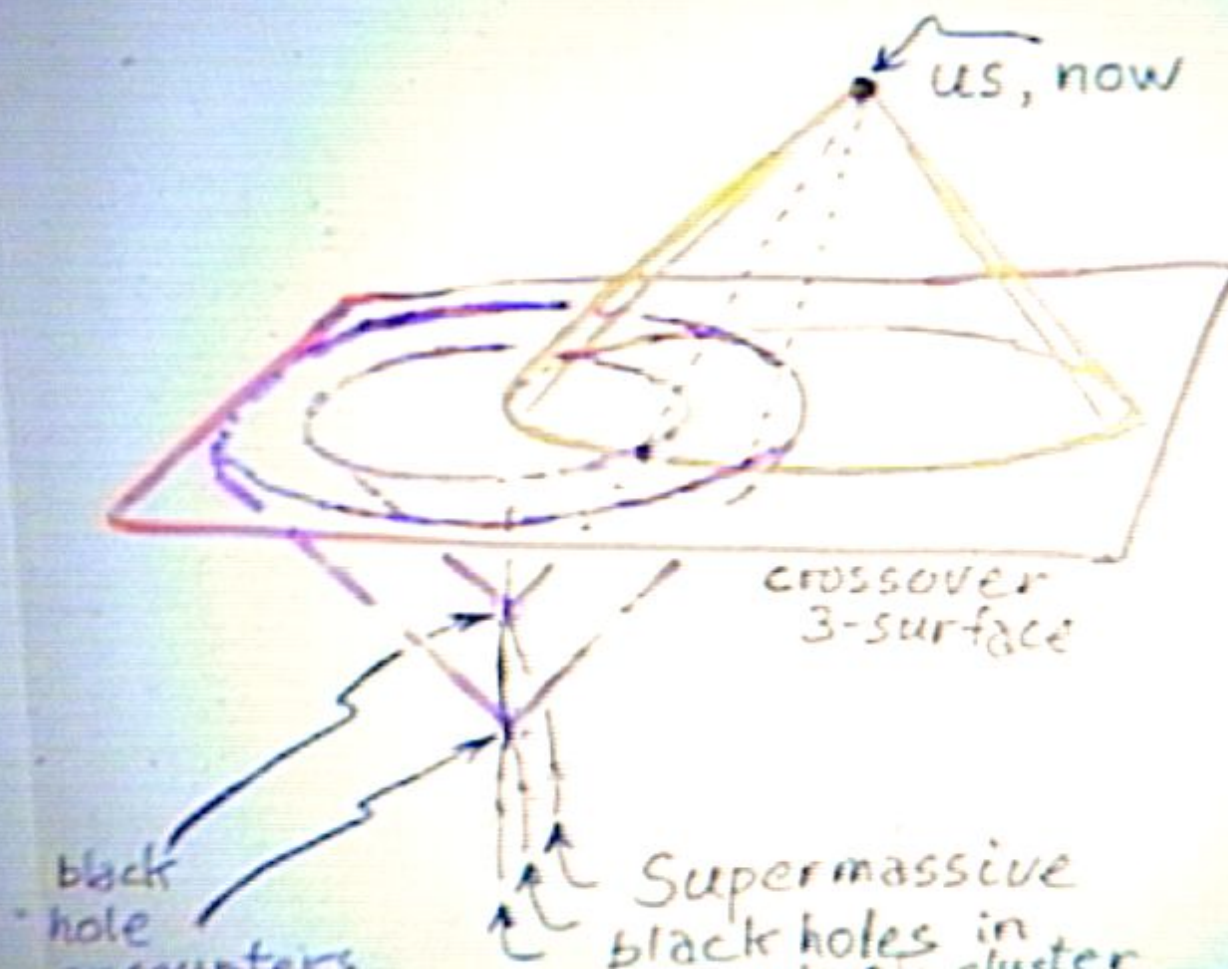
Can we "see" through
into the aeon prior to ours?

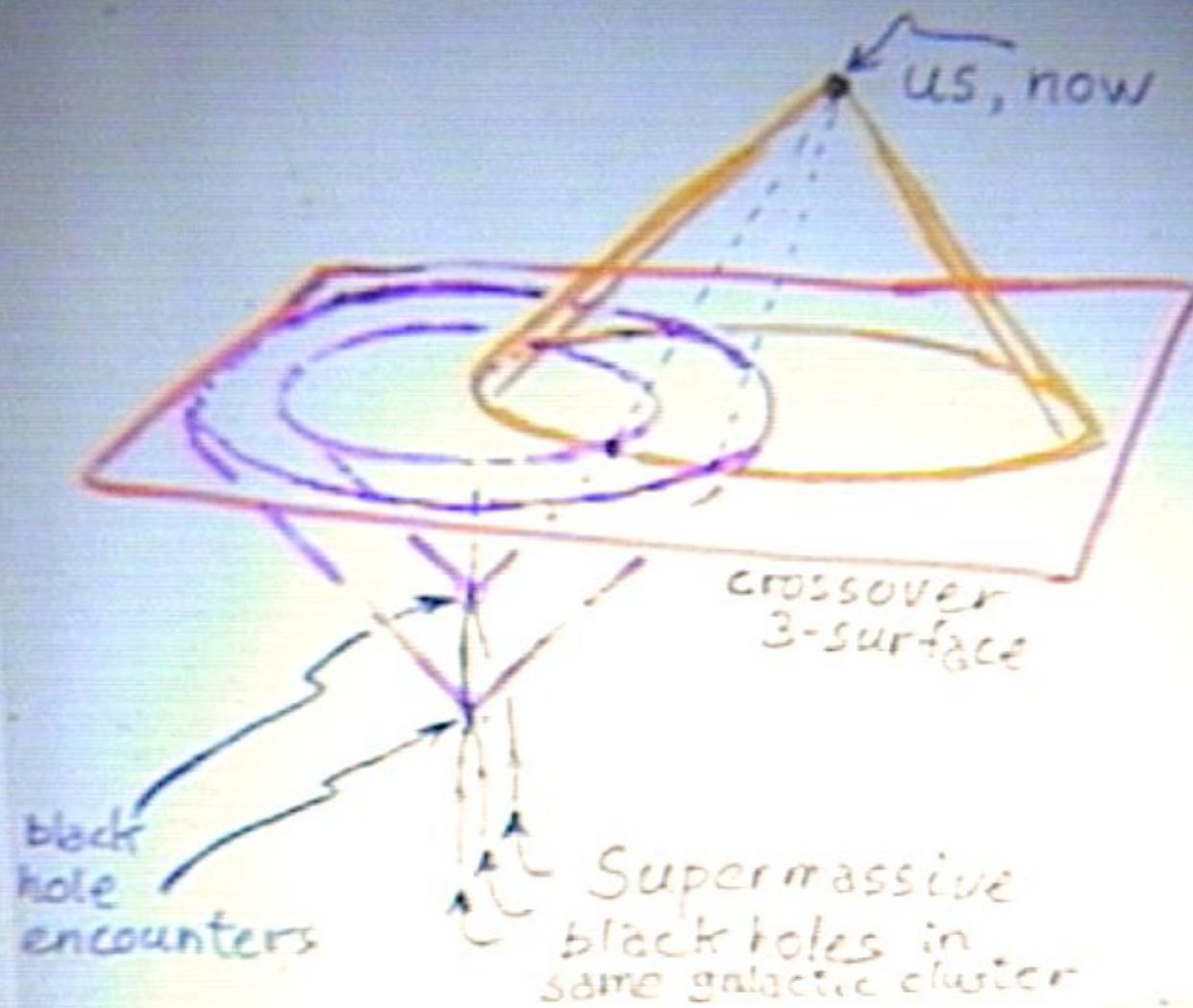


Can we "see" through
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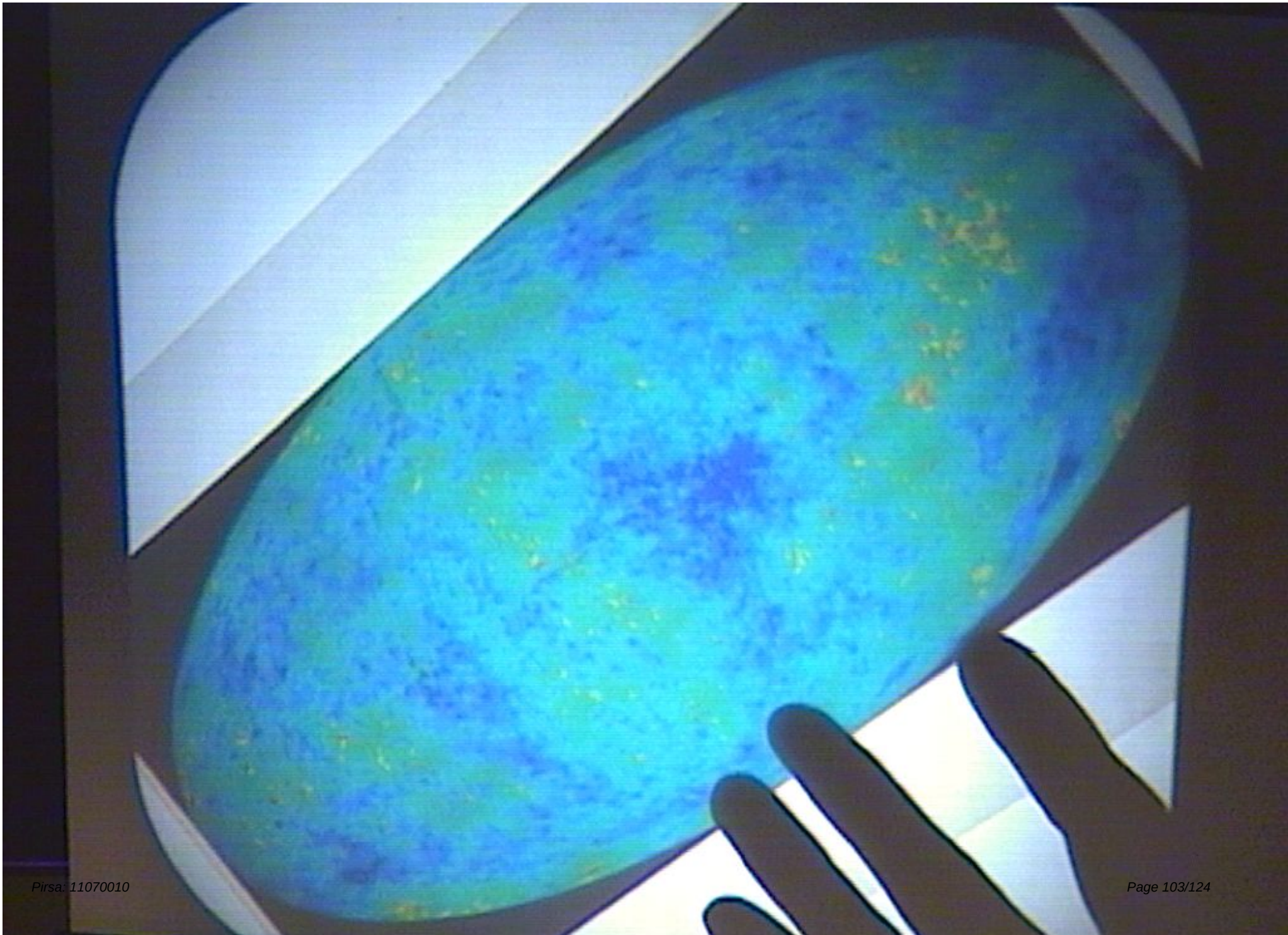
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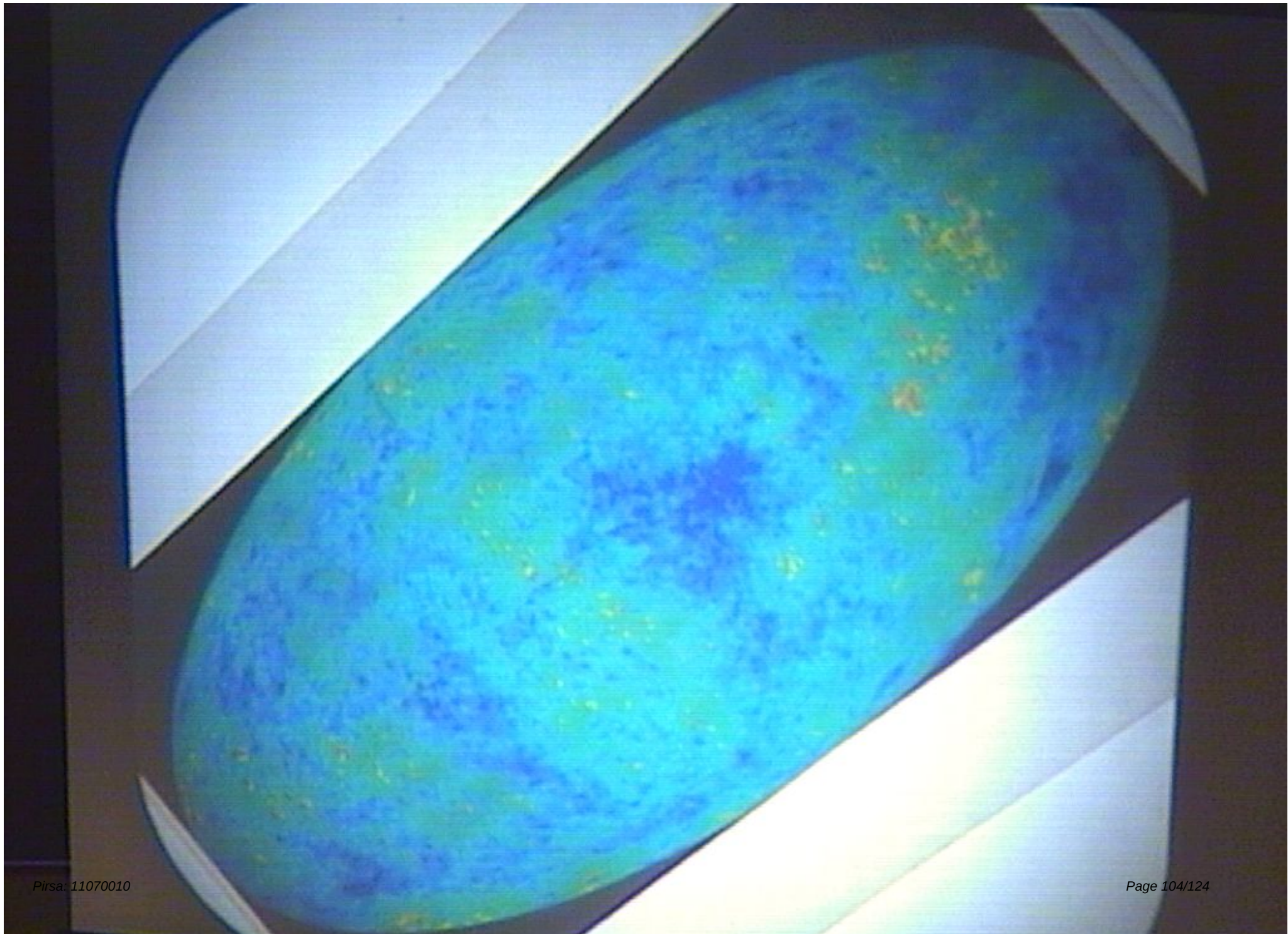




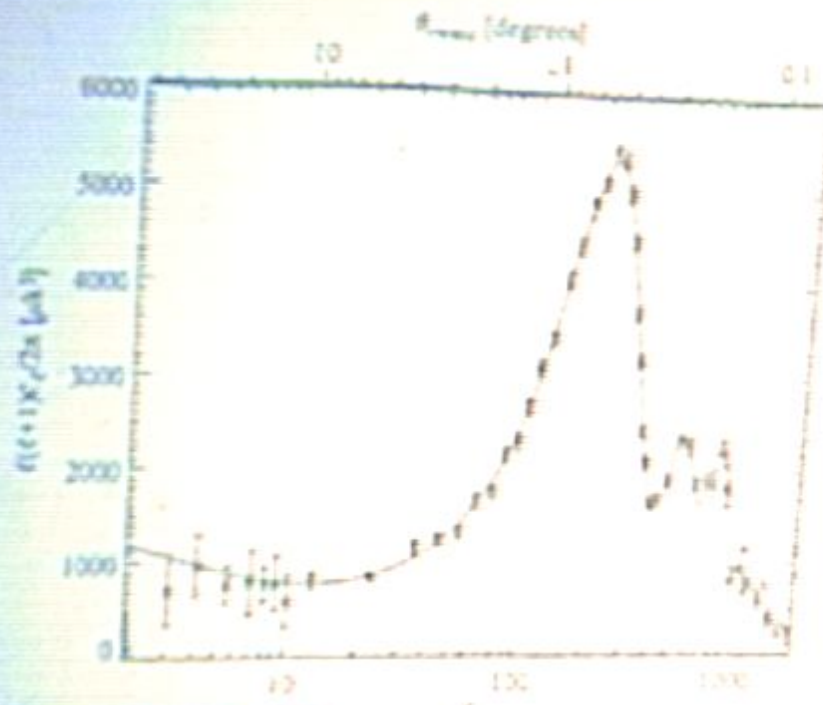
noteworthy
• temperature

• low variance





CMB power spectrum



ℓ -values, in harmonic (ℓ, m)-
CMB temperatures

83143

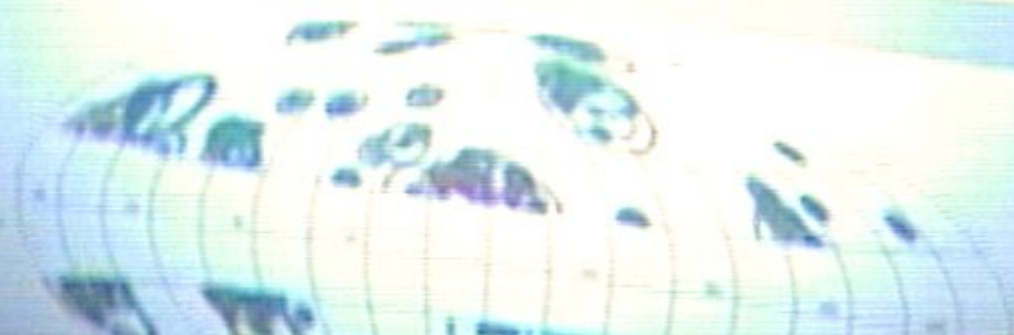
WMAPI (LINE 15, OF 3-10, n 352)

WMAPI (LINE 15, OF 3-10, n 352)

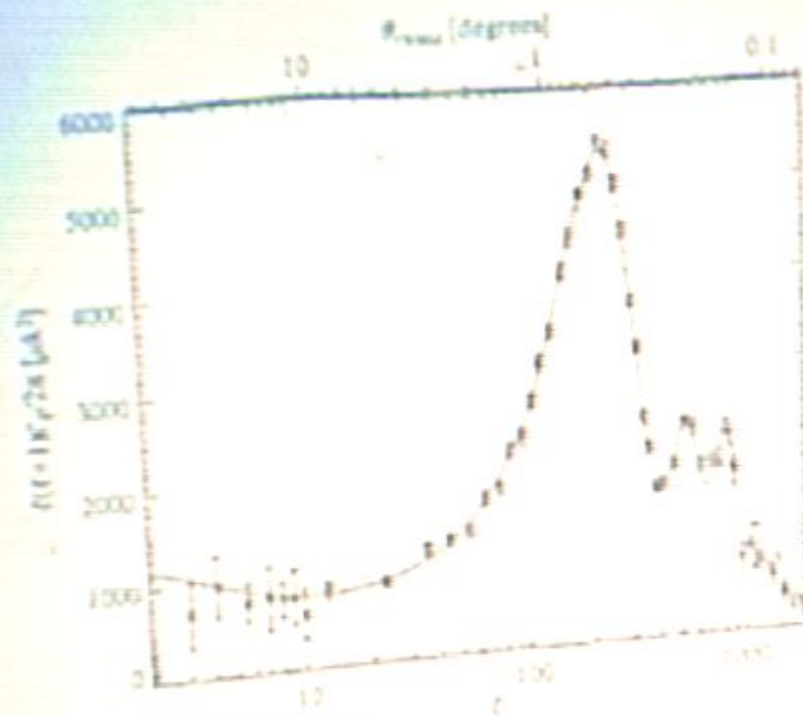


WMAPI (LINE 15, OF 3-10, n 352)

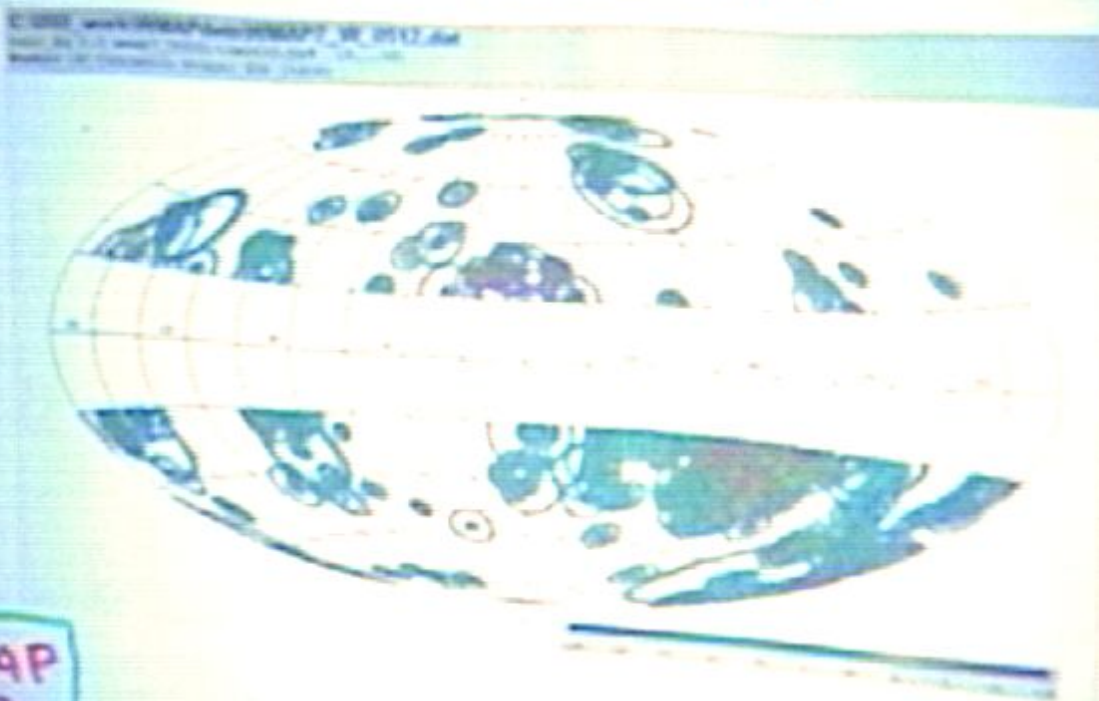
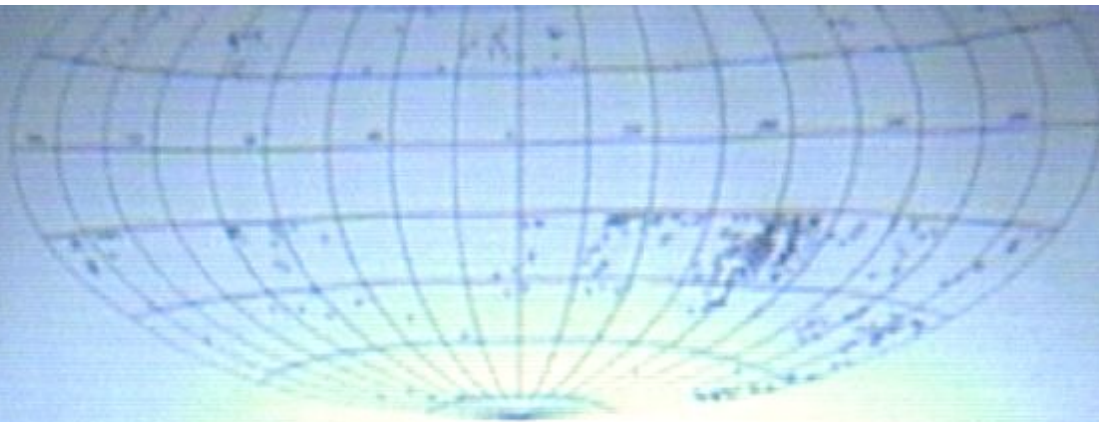
WMAPI (LINE 15, OF 3-10, n 352)



CMB power spectrum



l -values, in harmonic (l, m)-
basis of CMB temperatures

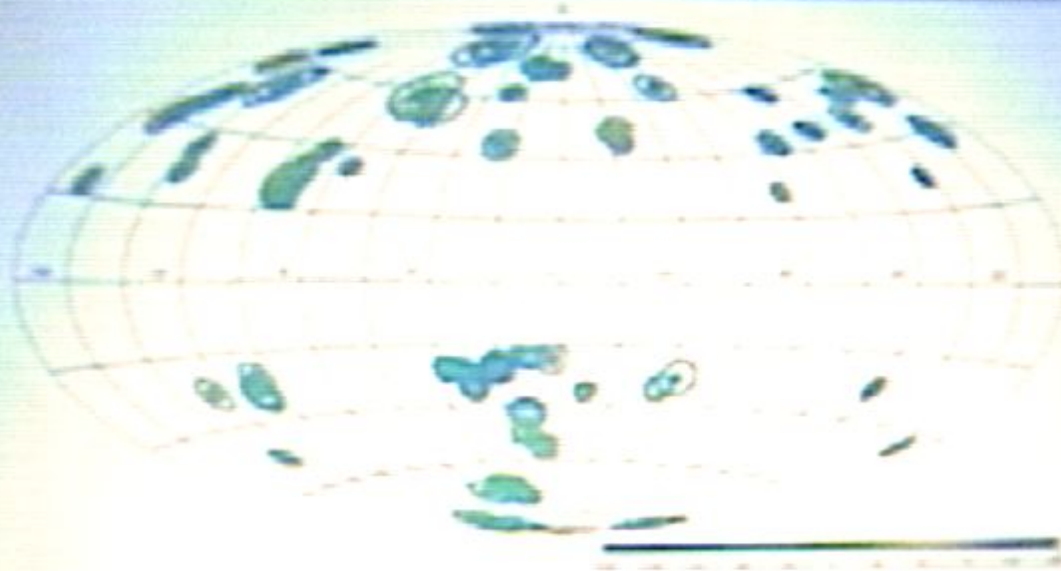


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352

SIM1-01 [Limit 15, cr 3..10, n 79]

THESE WERE THE FIRST SIMULATIONS OF THE UNIVERSE
AND THE FIRST TO SHOW THE UNIVERSE AS A WHOLE
AND NOT JUST A PART OF IT



Conventional
Simulation, but
no overlapped noise

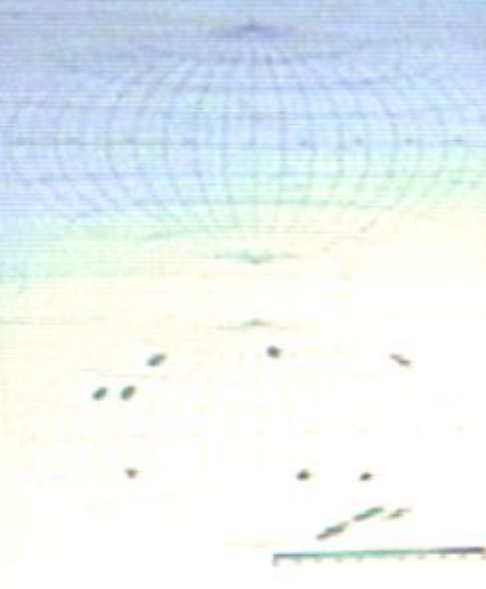
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79
centres

figs



47 centres



19 centres



12 centres

Figure 5. The corresponding steps to those of Figure 4, but where 5 simulations of 1 Mbit/s are used representing WNUF's spectrum with constant in value with overlapped WNUF's area. The differences are striking, notably for every three iterations was 47, 19, 12 respectively, the absence of significant suboptimization and all large circles.

Average of number of centres
for 10 independent simulations
→ 15 centres

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47 centres



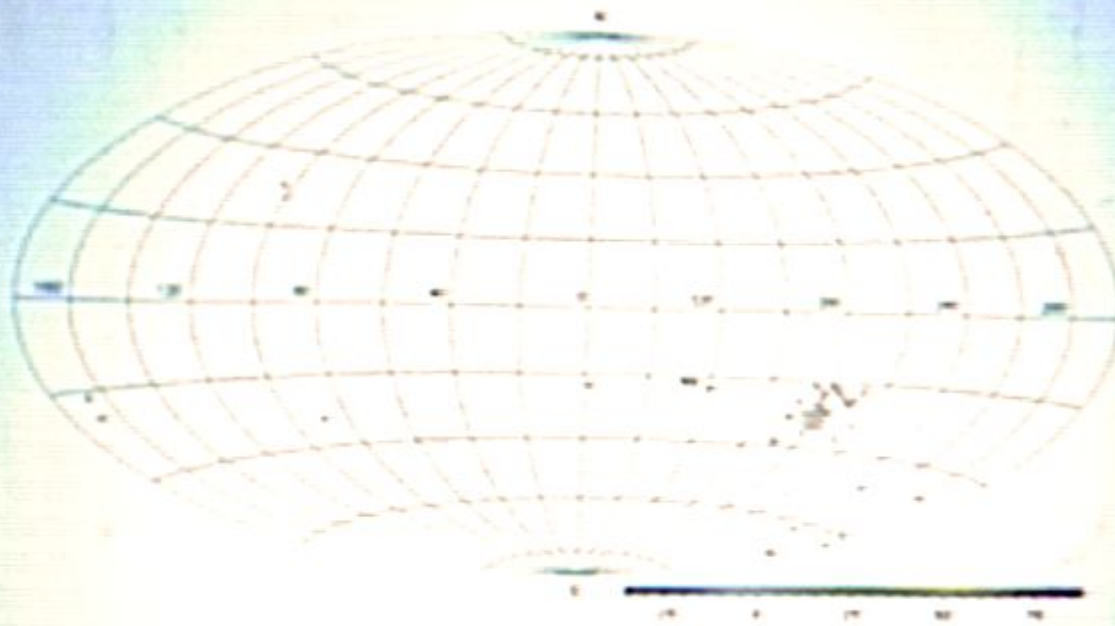
19 centres



12 centres

Average of number of centres
for 10 independent simulations
→ 15 centres

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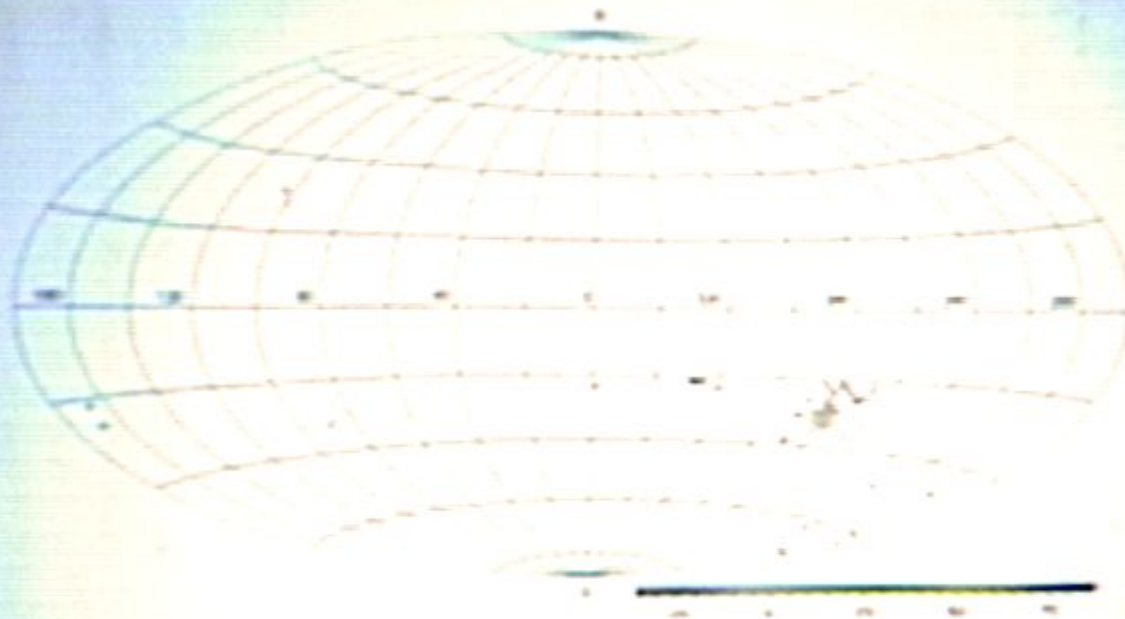


WMAP
data

56 centers of $N \geq 4$ circles

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56
centres

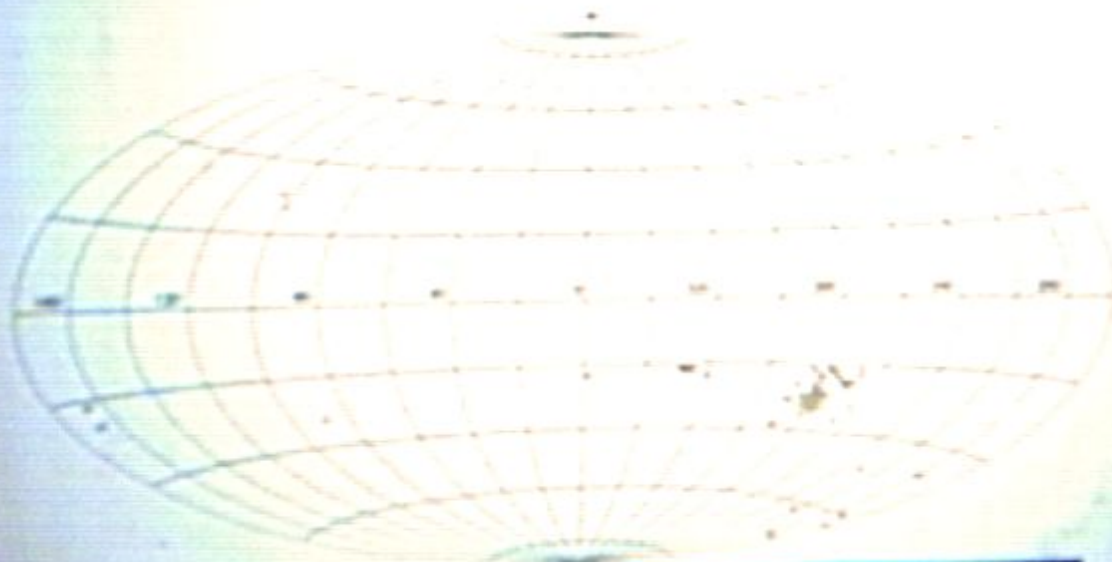
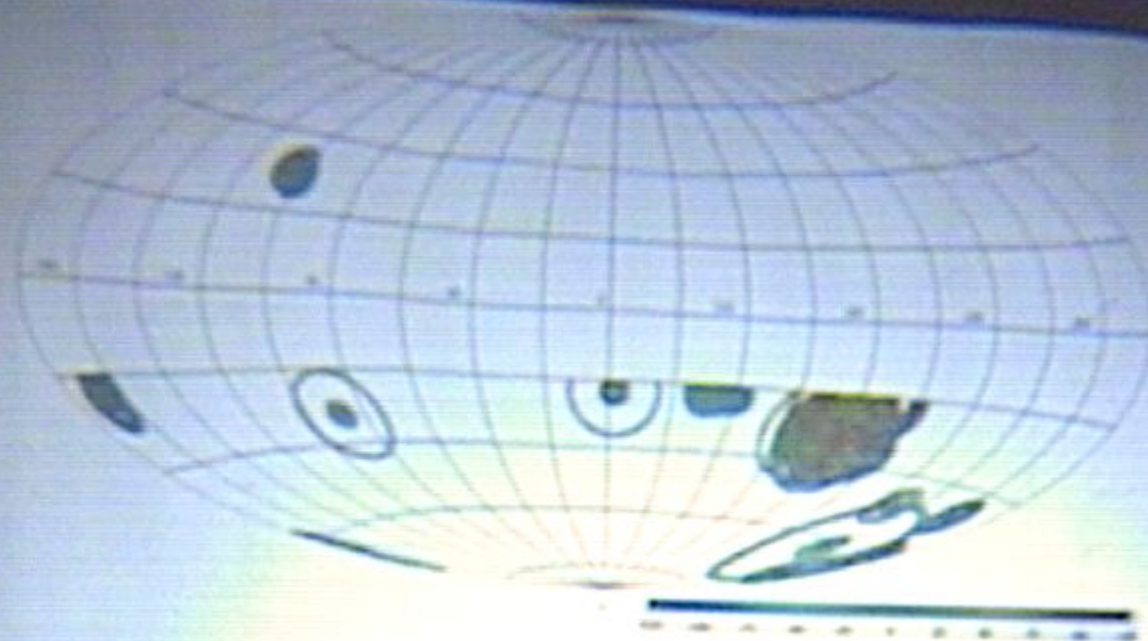


WMAP
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56 centers of $N \geq 4$ circles

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56
centres



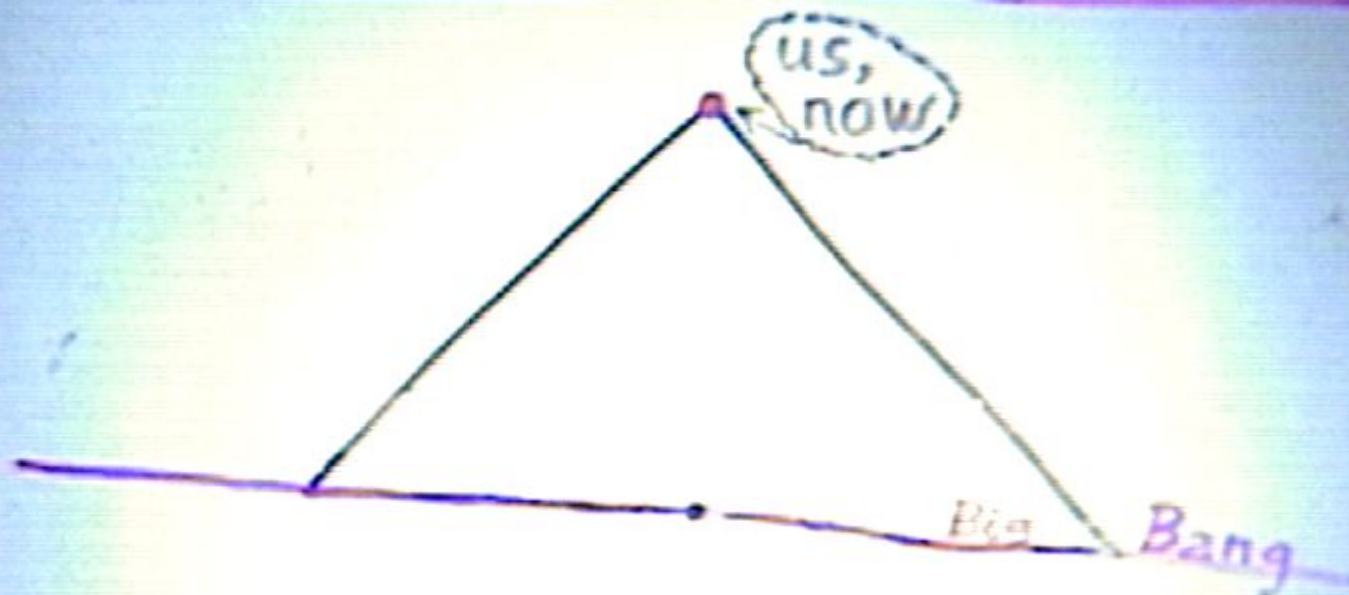
Previous Aeon to Our Own

future infinity →

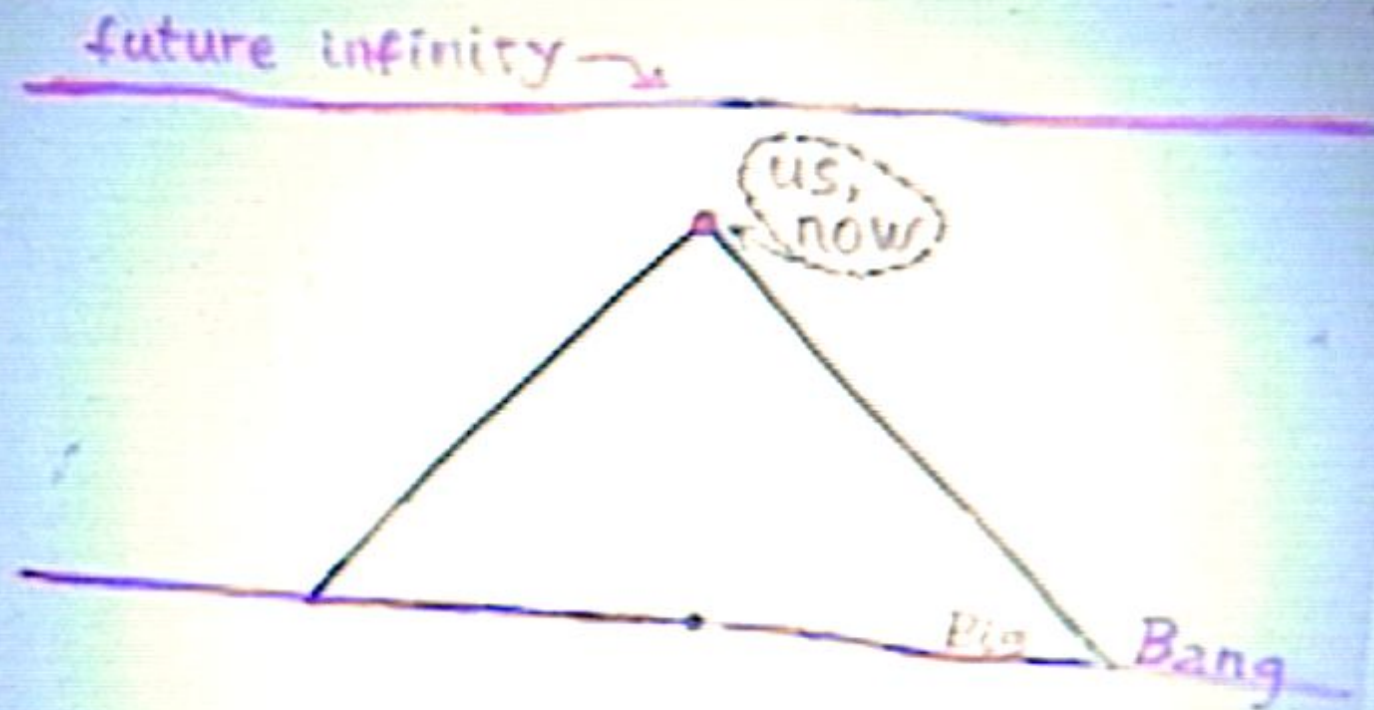


...THE ENCOUNTERS IN THE
Previous Aeon to Our Own

future infinity →



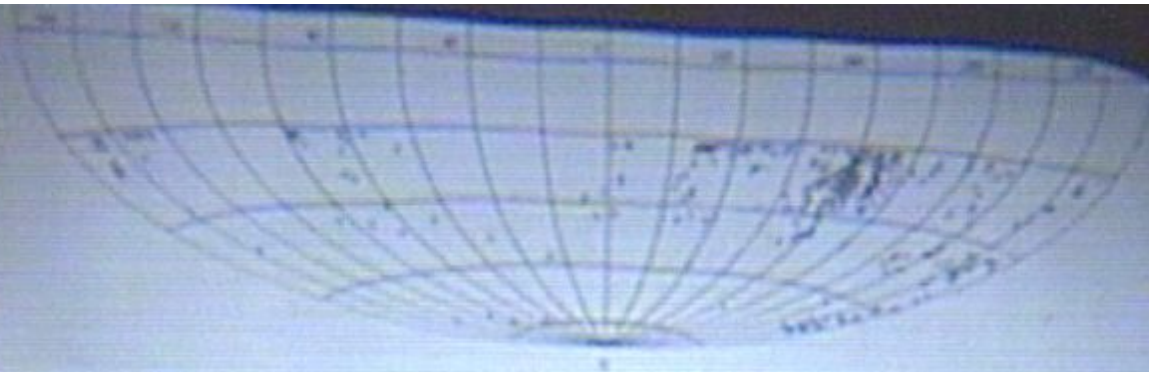
Black-Hole Encounters in the Previous Aeon to Our Own





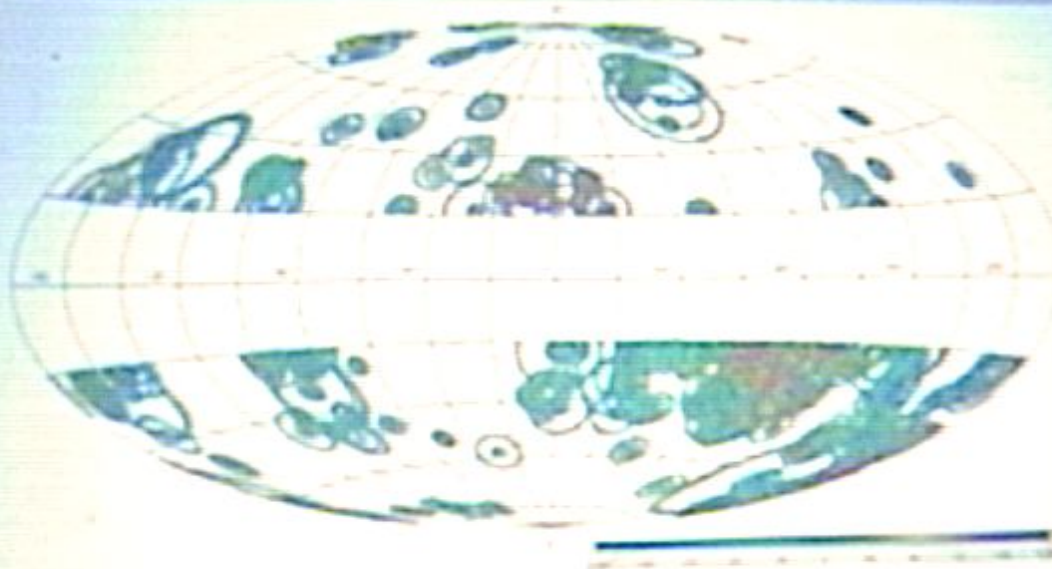
Maximum ring ° radius
~ 20°

Tod
Nelson & Wilson
-Euler



WMAP7 [Limit 15, cr 3..10, n 352]

C:\000\www\WMAP\data\WMAP7_W_0512.dat
 WMAP7 radio galaxy centers (cr 3..10, n 352)
 WMAP7 radio galaxy centers (cr 3..10, n 352)



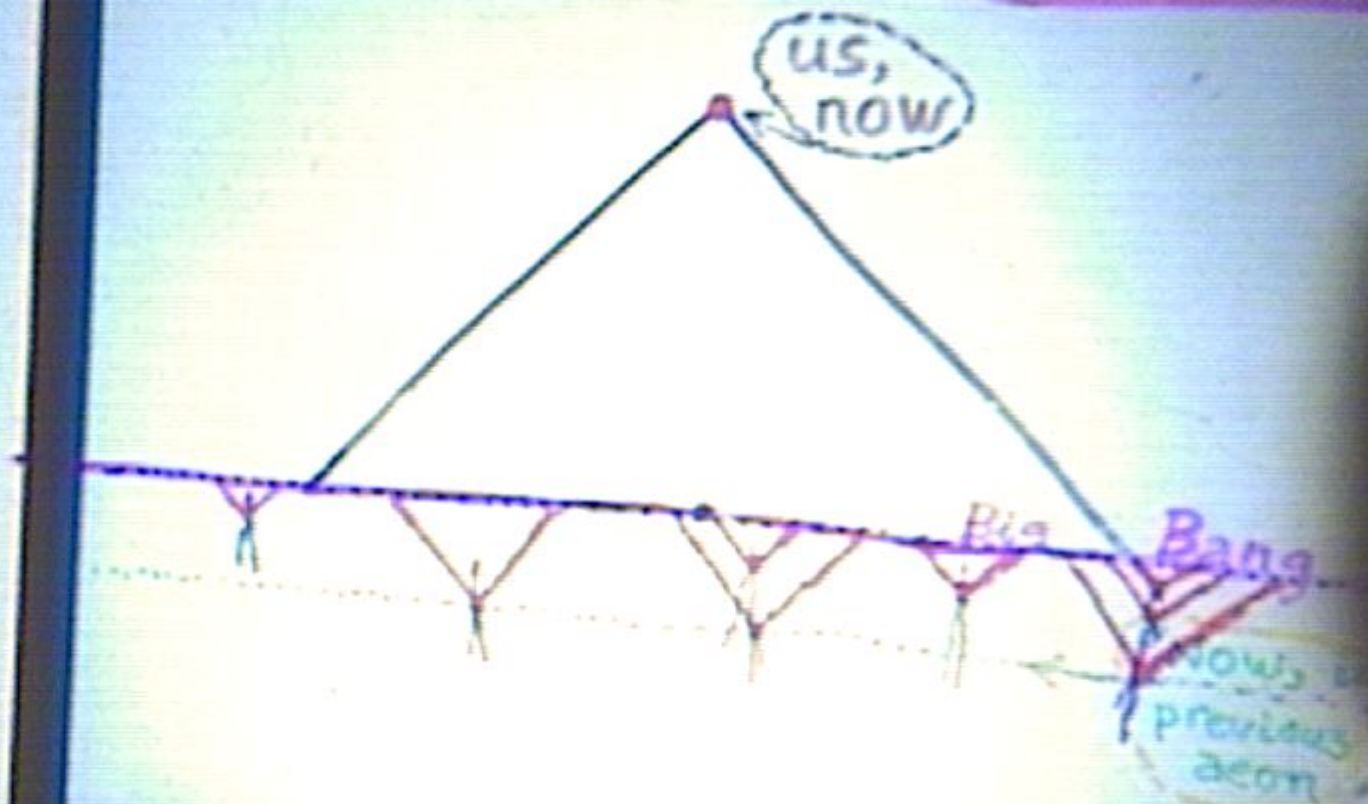
WMAP
data

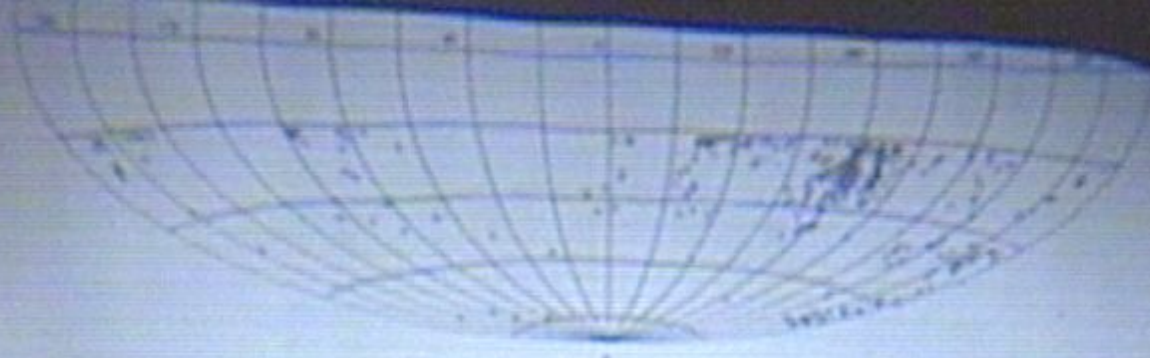
Vahe Gurzadyan

352
centres

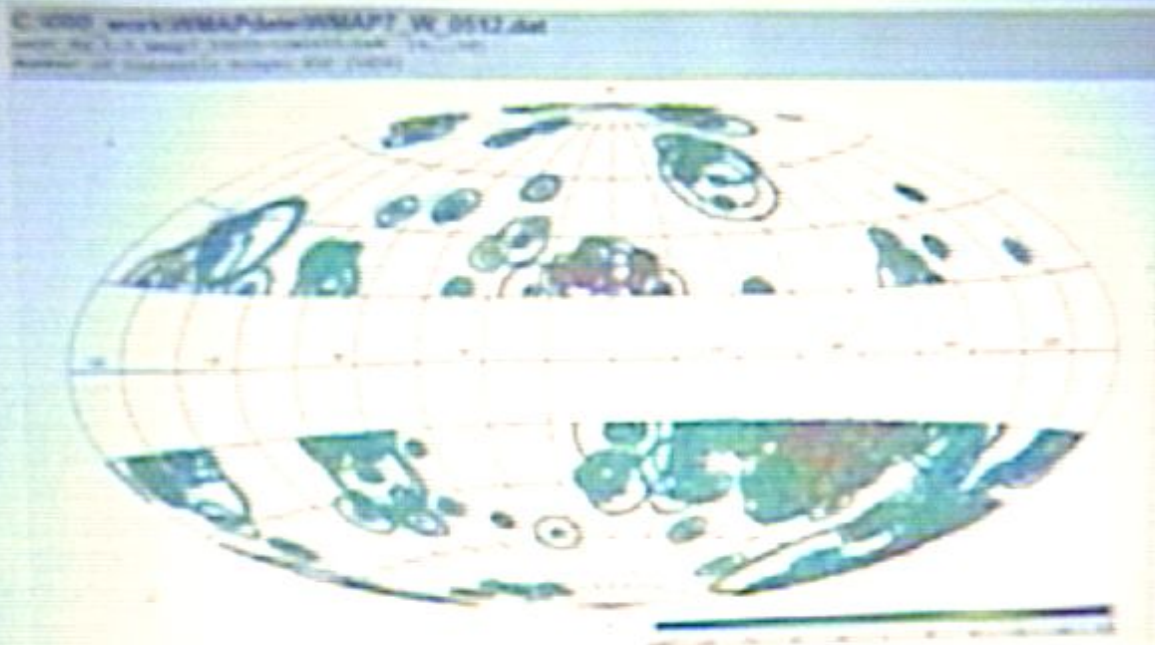
Black-Hole Encounters in the
previous Aeon to Our Own

Future infinity $\rightarrow$





WMAP7 [Limit 15, cr 3..10, n 352]



WMAP
data

Vahe Gurzadyan

352
centres