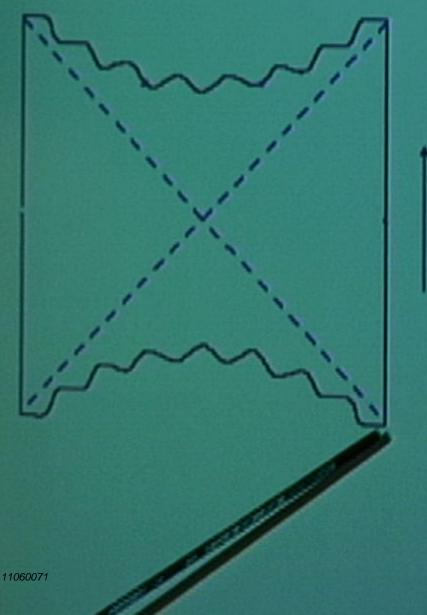
Title: Inside the horizon with holographic Wilsonian RG

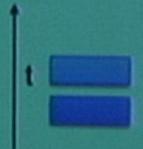
Date: Jun 21, 2011 11:00 AM

URL: http://pirsa.org/11060071

Abstract:

Pirsa: 11060071 Page 1/396

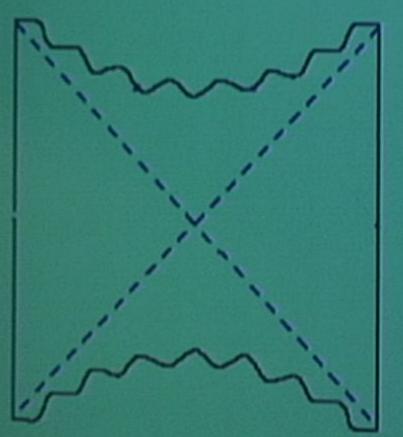


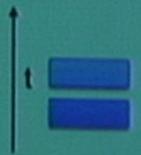


Certain Super Yang-Mills at finite temperature on S3.

Example: N=4 SYM with gauge group SU(N)

Pirsa: 11060071

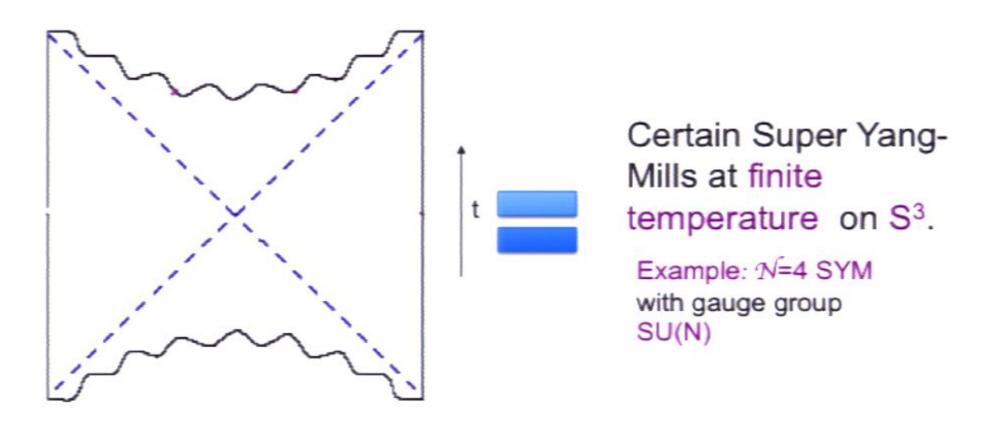




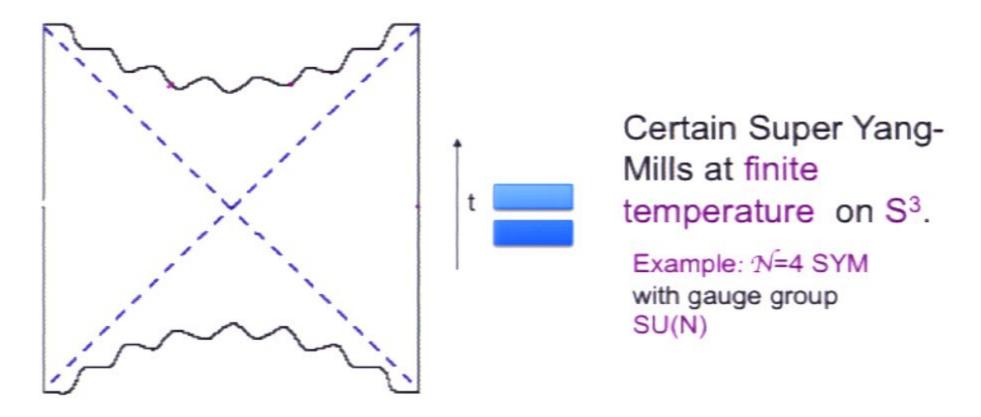
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Example: N=4 SYM with gauge group SU(N)

Pirsa: 11060071 Page 3/39

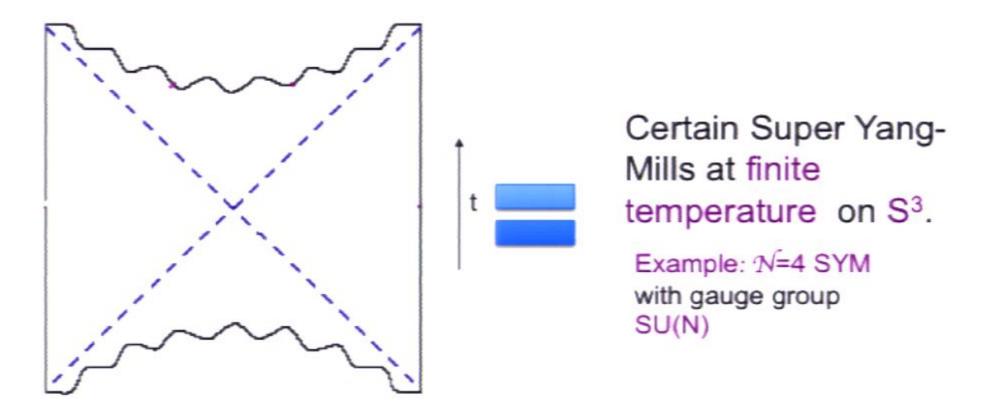


Pirsa: 11060071 Page 4/396



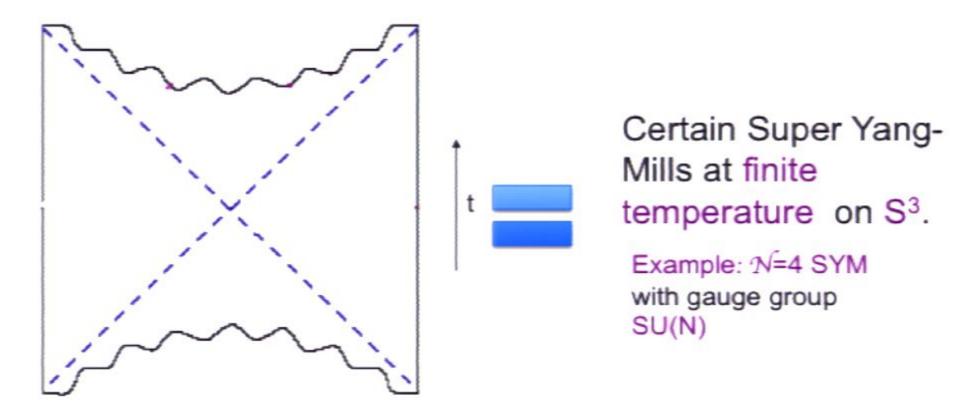
Challenges: What happens to the horizon and singularity in quantum gravity?

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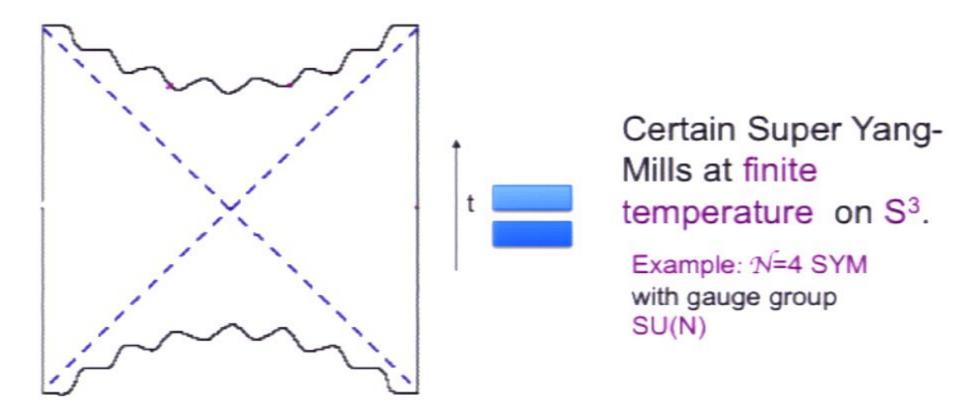
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Pirsa: 11060071



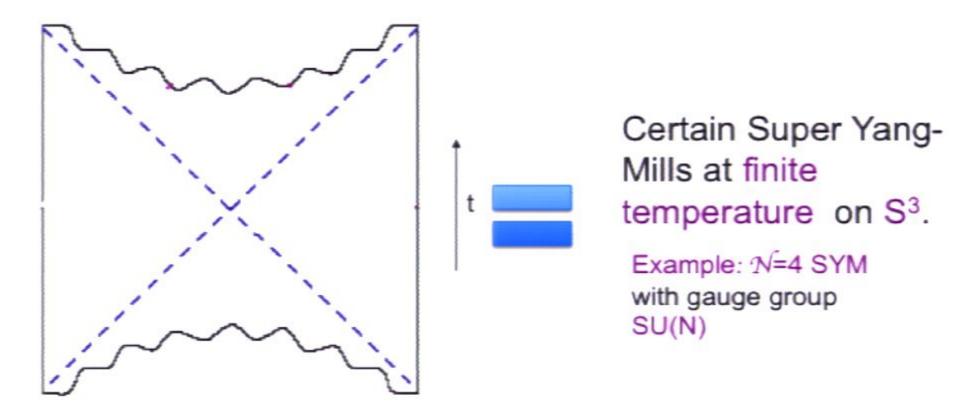
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DWisat 100007 understand the region inside the horizon using AdS/@730T?



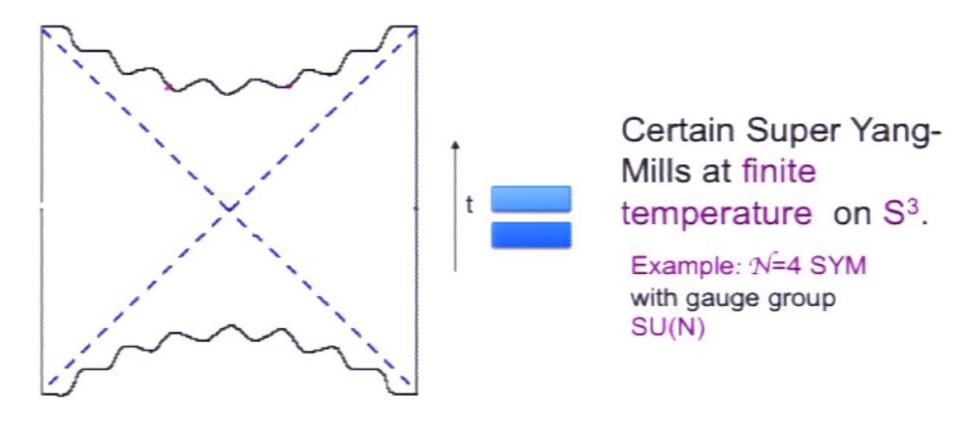
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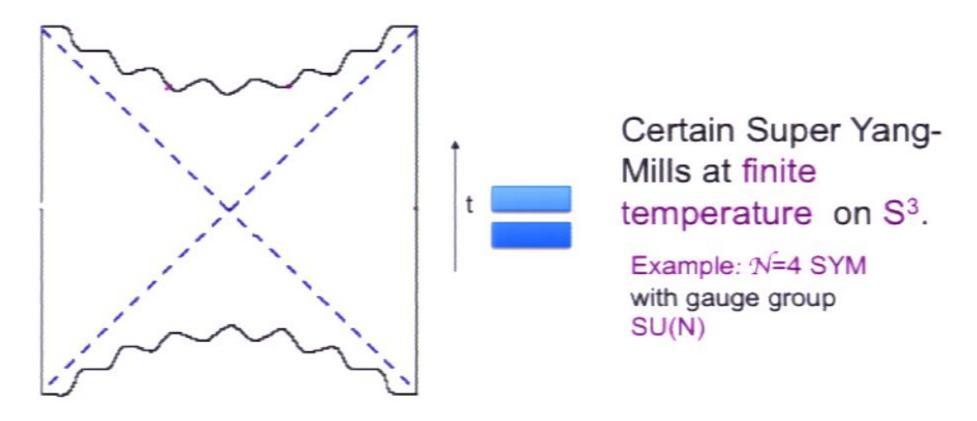
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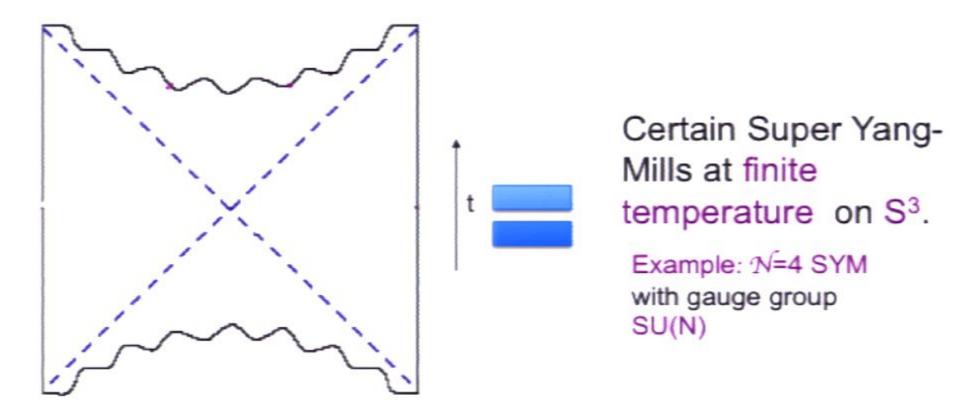
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Plan

Holographic Wilsonian RG

Faulkner, HL, Rangamani, arxiv: 1010:4036

Heemskerk and Polchinski: 1010:1264

Pirsa: 11060071 Page 13/396

Plan

Holographic Wilsonian RG

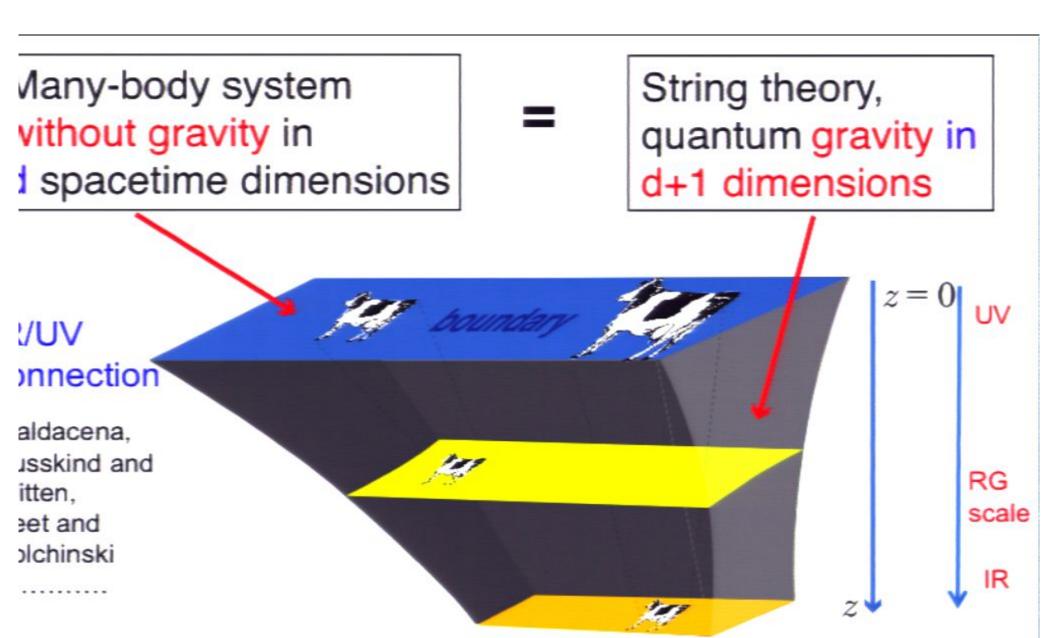
Faulkner, HL, Rangamani, arxiv: 1010:4036

Heemskerk and Polchinski: 1010:1264

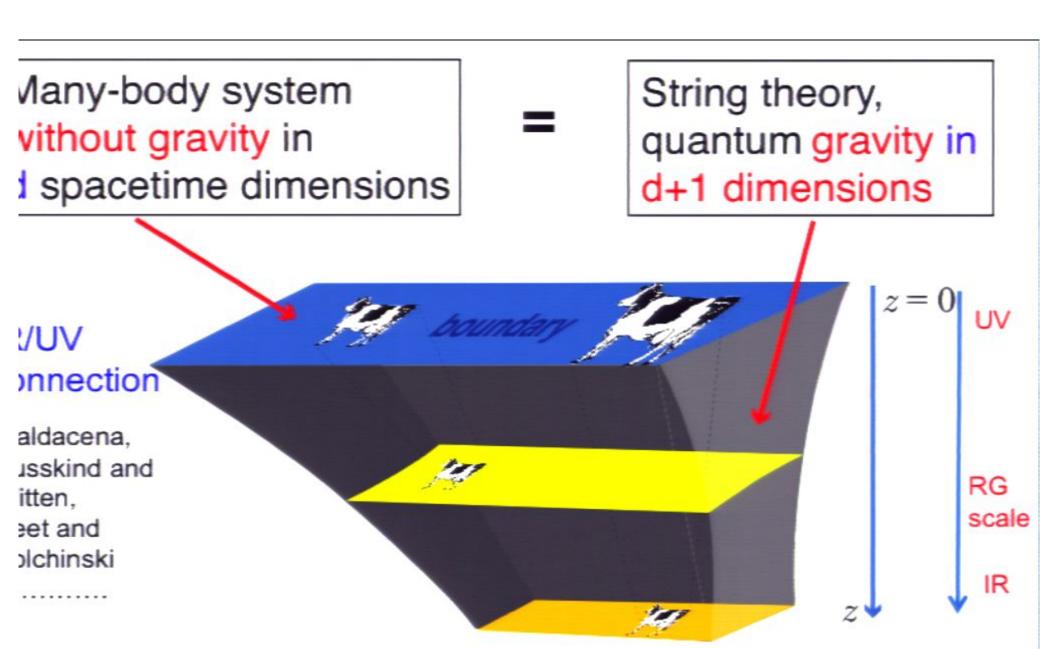
 Some speculations on how to describe physics beyond the black hole horizon using AdS/CFT

Festuccia, HL, hep-th/0506202, hep-th/0611098

Pirsa: 11060071 Page 14/396



Pirsa: 11060071 Page 15/396



The extra bulk radial direction: geometrization of the Pirsa: 100071 ormalization group flow of the boundary system!

Pirsa: 11060071 Page 17/396

One should organize a system by energy scales.

Pirsa: 11060071 Page 18/396

- One should organize a system by energy scales.
- Higher energy degrees of freedom should be properly integrated out, their physical effects encoded in the resulting low energy effective action.

Pirsa: 11060071 Page 19/396

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Pirsa: 11060071 Page 20/396

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Pirsa: 11060071

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$$I_{\mathrm{UV}}[\Phi,\Lambda]$$
 Flow is reversible if keeping head act I_{UV}

Page 22/396

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Page 23/396

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hersa@wact luv

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$$\begin{split} Z &= \int_{\Lambda_0} D\Phi \, e^{-I_0[\Phi,\Lambda_0]} \\ &= \int_{\Lambda} D\Phi \, e^{-I_0[\Phi,\Lambda]-I_{\rm UV}[\Phi,\Lambda]} \end{split} \qquad I_{\rm UV}[\Phi,\Lambda] \end{split}$$
 Flow is reversible if keeping

Page 24/396

Question:

What is the counterpart of the Wilsonian effective action in the dual gravity description?

Pirsa: 11060071 Page 25/396

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While there have been much work on holographic RG, the Wilsonian point of view has not been properly developed.

Will comment on relation with earlier work at the end.

Pirsa: 11060071 Page 26/396

Question:

What is the counterpart of the Wilsonian effective action in the dual gravity description?

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Will comment on relation with earlier work at the end.

Will work in the large N limit, i.e. classical bulk gravity.

Pirsa: 11060071 Page 27/396

Pirsa: 11060071 Page 28/396

$$ds^2 = -g_{tt}dt^2 + g_{zz}dz^2 + g_{ii}d\vec{x}^2$$

z = 0: boundary, asymptotic AdS, metric only depend on z.

Pirsa: 11060071 Page 29/396

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Pirsa: 11060071 Page 30/396

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$$E_{\text{boundary}} = \sqrt{g_{tt}(z)} E_{\text{prop}}$$

$$E_{\rm prop} \sim \frac{1}{R}$$

Typical physical process at z:
$$E_{
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Pirsa: 11060071

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 34/396

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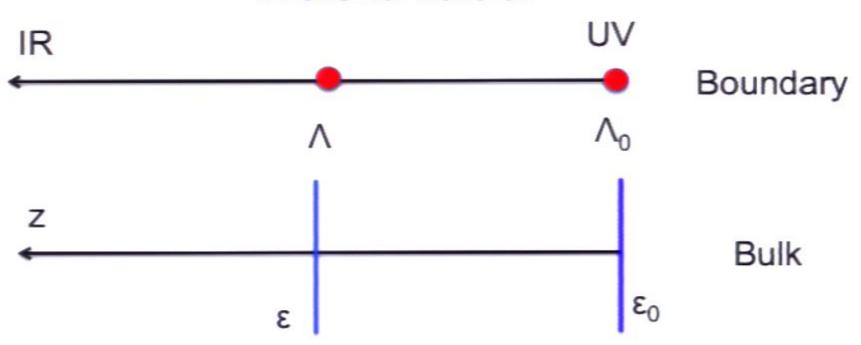
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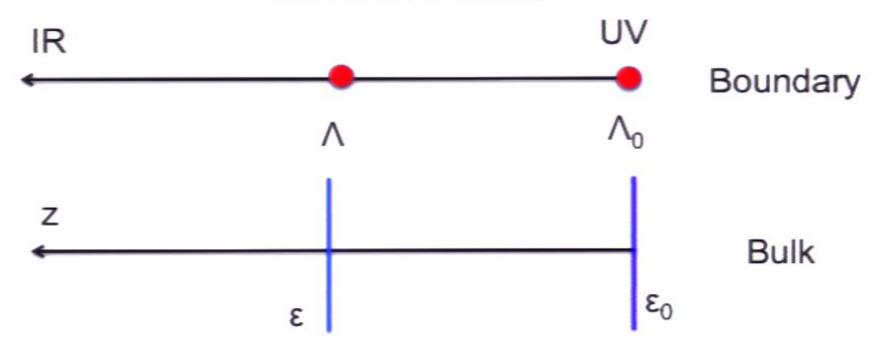
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Pirsa: 11060071 Page 38/396

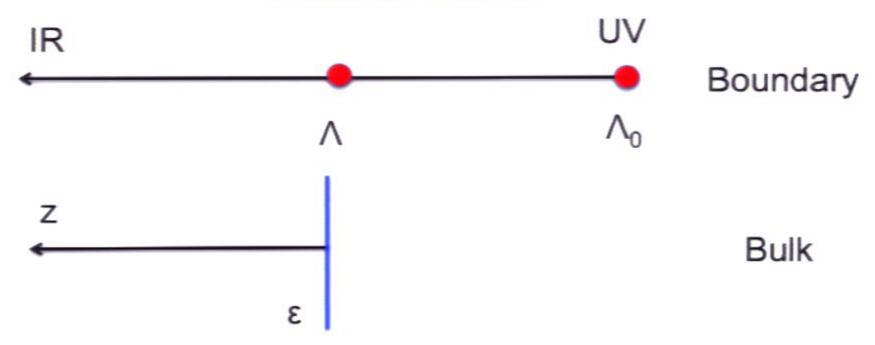


Pirsa: 11060071 Page 39/396



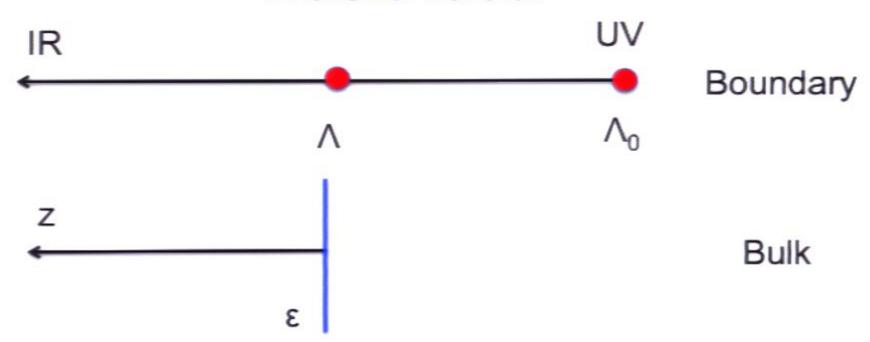
Integrate out bulk d.o.f. in the region $\varepsilon > z > \varepsilon_0$

Pirsa: 11060071 Page 40/396



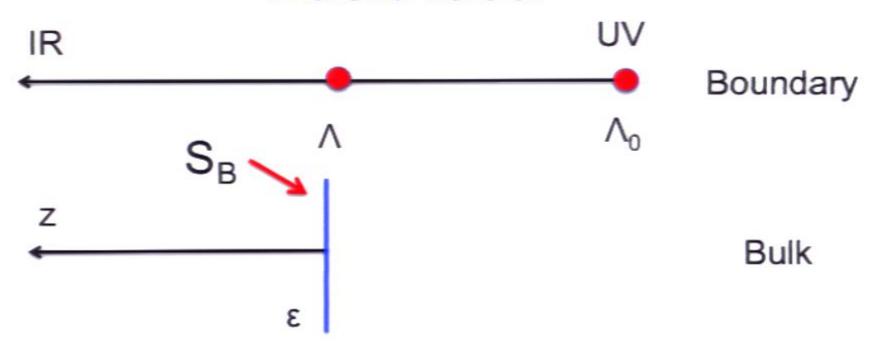
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Pirsa: 11060071 Page 41/396



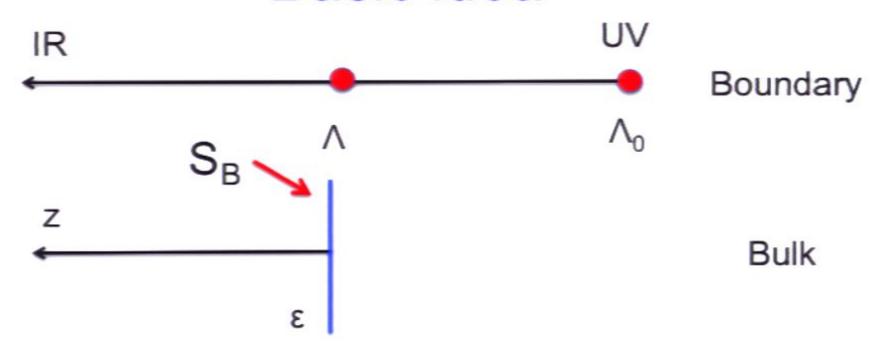
Integrate out bulk d.o.f. in the region $\varepsilon > z > \varepsilon_0$

$$\begin{aligned} I_{\text{bulk}} &= \int_{z \ge \epsilon_0} D\phi \, e^{-S_0[\phi]} \\ &= \int_{z \ge \epsilon} D\phi \, e^{-S_0[\phi] - S_B[\epsilon, \phi(z = \epsilon, x^{\mu})]} \\ &= \int_{z > \epsilon} D\phi \, e^{-S_0[\phi] - S_B[\epsilon, \phi(z = \epsilon, x^{\mu})]} \end{aligned}$$



Integrate out bulk d.o.f. in the region $\varepsilon > z > \varepsilon_0$

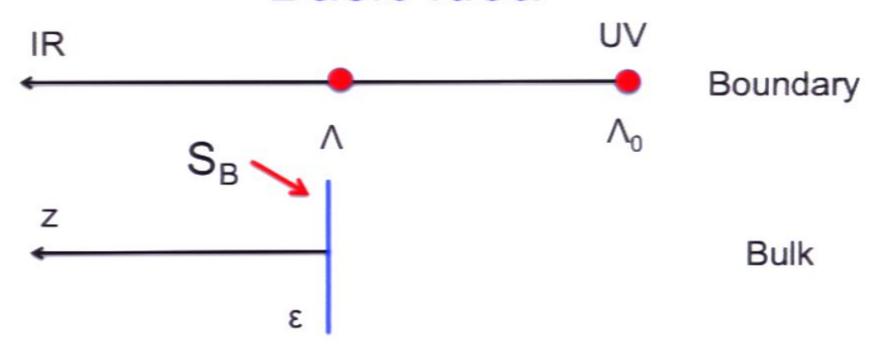
$$\begin{aligned} T_{\text{bulk}} &= \int_{z \ge \epsilon_0} D\phi \, e^{-S_0[\phi]} \\ &= \int_{z \ge \epsilon} D\phi \, e^{-S_0[\phi] - S_B[\epsilon, \phi(z = \epsilon, x^{\mu})]} \\ &= \int_{z > \epsilon} D\phi \, e^{-S_0[\phi] - S_B[\epsilon, \phi(z = \epsilon, x^{\mu})]} \end{aligned}$$



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S_B: boundary action for the remaining part of the bulk spacetime. Page 44/396



Integrate out bulk d.o.f. in the region $\varepsilon > z > \varepsilon_0$

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S_B: boundary action for the remaining part of the bulk spacetime. Page 45/396

Pirsa: 11060071 Page 46/396

Legendre transform

 $I_{\mathrm{UV}}[\Phi,\Lambda]$



$$S_B[\phi,\epsilon]$$

Pirsa: 11060071 Page 47/396

Legendre transform

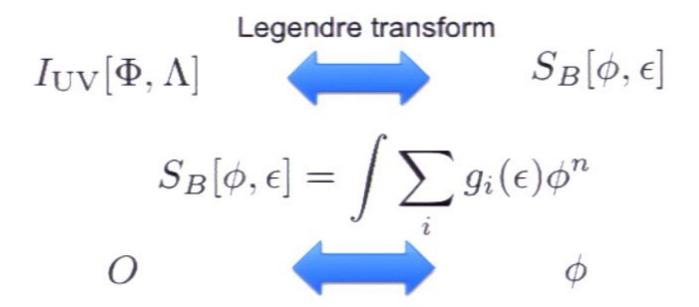
$$I_{\mathrm{UV}}[\Phi, \Lambda]$$



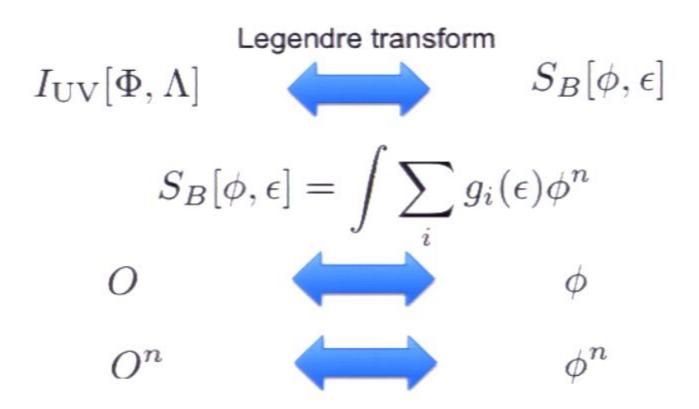
$$S_B[\phi,\epsilon]$$

$$S_B[\phi, \epsilon] = \int \sum_i g_i(\epsilon) \phi^n$$

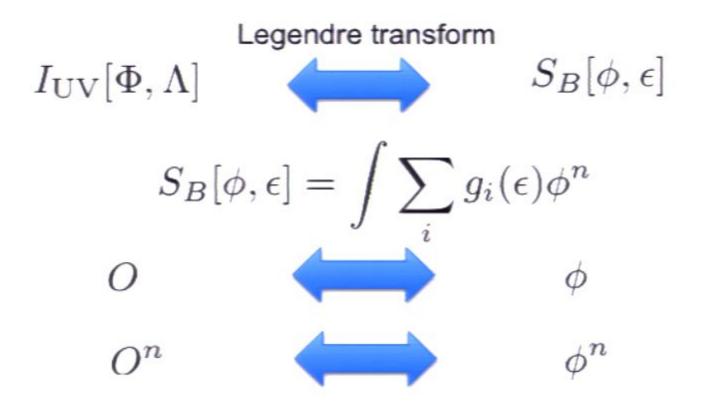
Pirsa: 11060071 Page 48/396



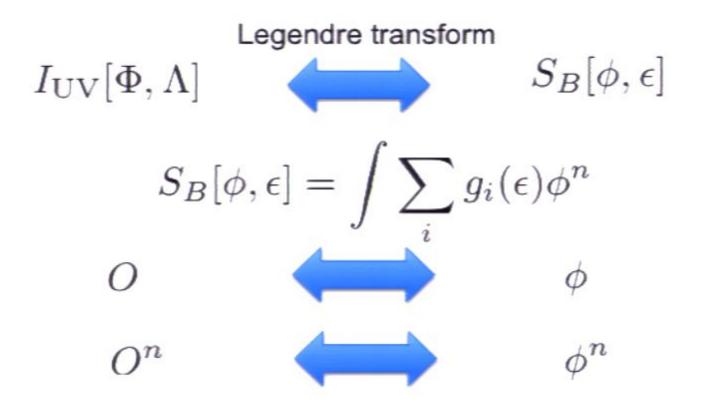
Pirsa: 11060071 Page 49/396



Pirsa: 11060071 Page 50/396



Generically multiple-trace operators will be induced along the flow.



Generically multiple-trace operators will be induced along the flow.

Consistent with field theory expectations.

Miao Li (2000)

Pirsa: 11060071 Page 53/396

Physics should not depend on where we choose z=ε suface

Pirsa: 11060071 Page 54/396

Physics should not depend on where we choose z=ε suface



Flow equation for $S_B[\epsilon]$

Pirsa: 11060071 Page 55/396

Physics should not depend on where we choose z=ε suface



Flow equation for $S_B[\varepsilon]$

Semi-classical limit: Hamilton-Jacobi equation

$$\partial_{\epsilon} S_B[\phi, \epsilon] = -\int_{z=\epsilon} d^d x \, H\left(\phi, \Pi = \frac{\delta S_B}{\delta \phi}\right)$$

H: bulk Hamiltonian corresponding to z-foliation.

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$$S_B[\phi,\epsilon] = \int \sum_i g_i(\epsilon)\phi^n$$
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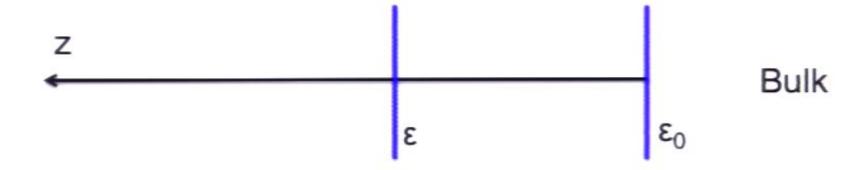
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Pirsa: 11060071 Page 100/396

Standard AdS/CFT procedure: Solve classical equations vith boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).

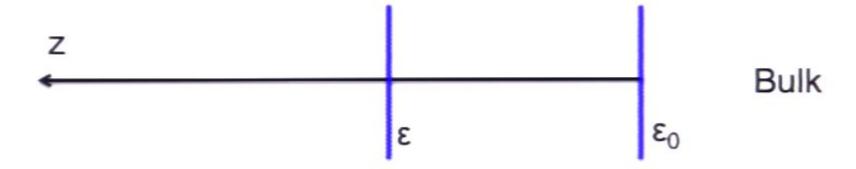
Pirsa: 11060071 Page 101/396

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Pirsa: 11060071 Page 102/396

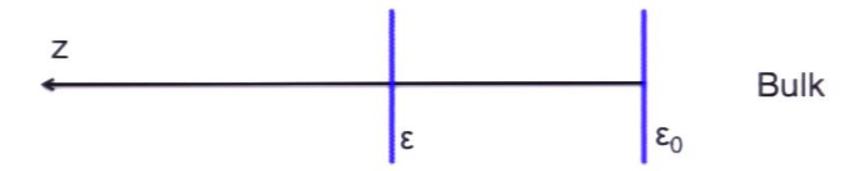
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Pirsa: 11060071 Page 103/396

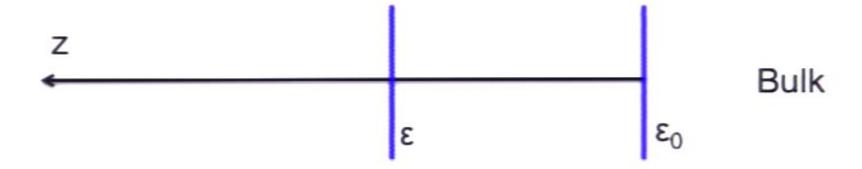
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Pirsa: 11060071 Page 104/396

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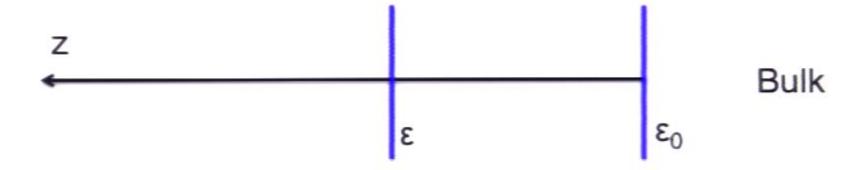
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Shift the cutoff surface



flow of boundary conditions

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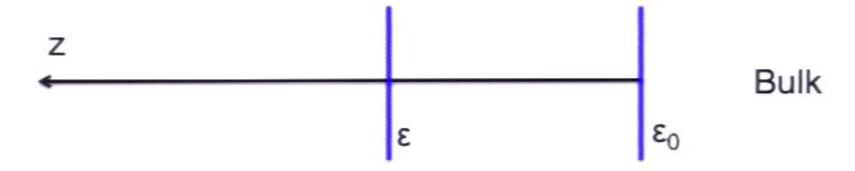
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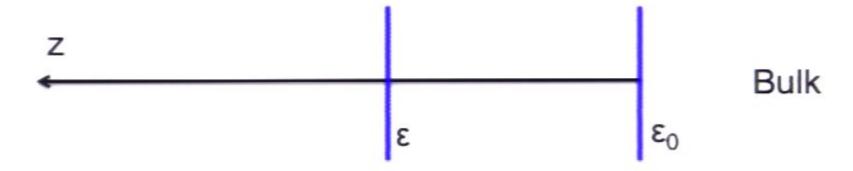
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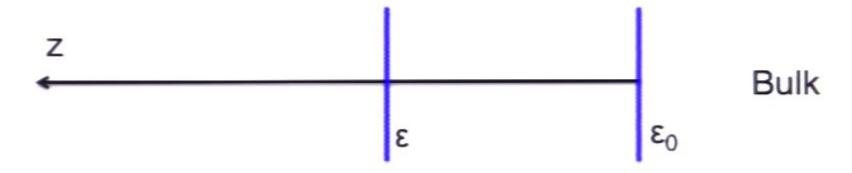


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Page 108/396

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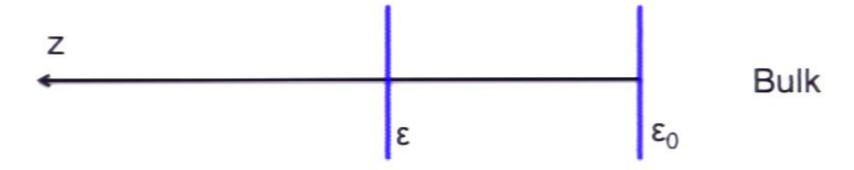
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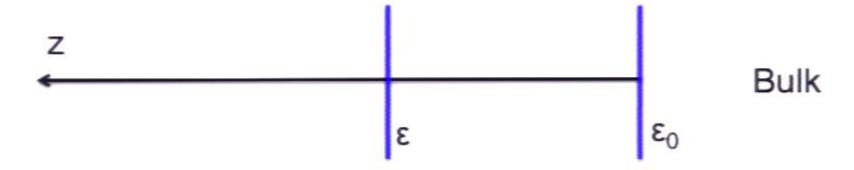


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Page 110/396

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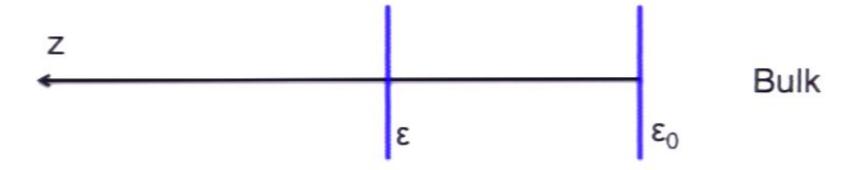


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Page 111/396

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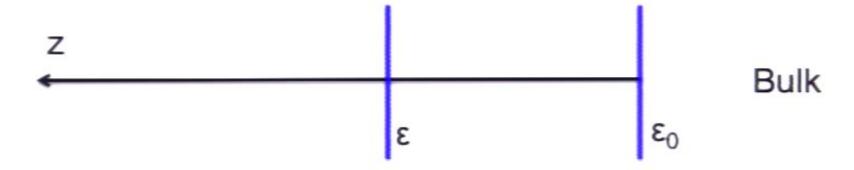


flow of boundary conditions



Page 112/396

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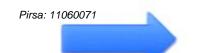


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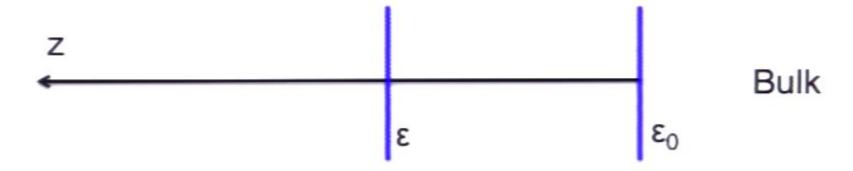


flow of boundary conditions



flow of S_R

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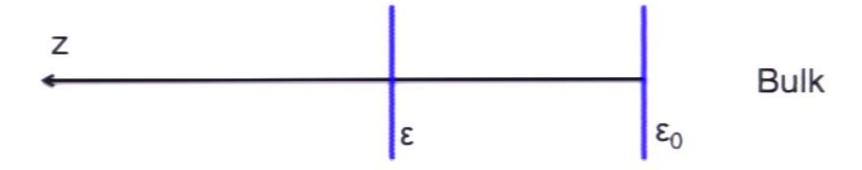


flow of boundary conditions



Page 114/396

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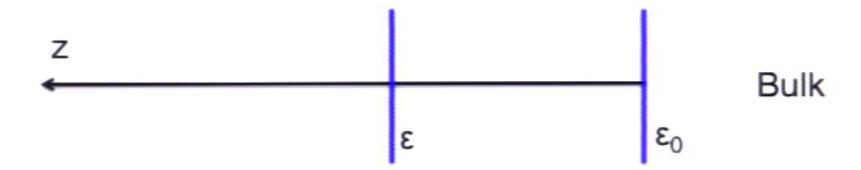
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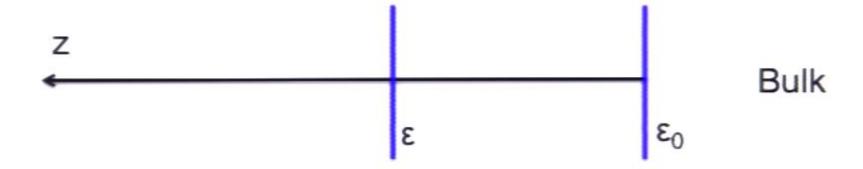


flow of boundary conditions



Page 116/396

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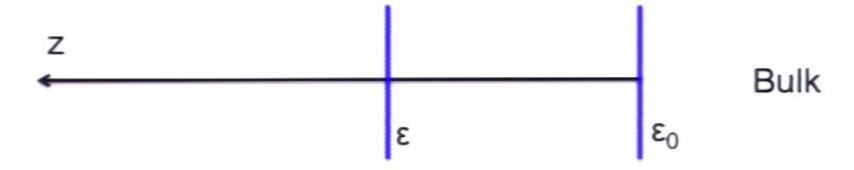
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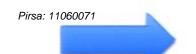


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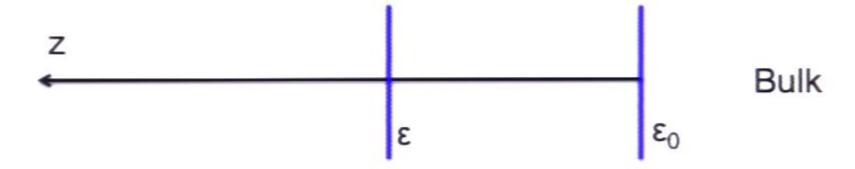


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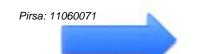


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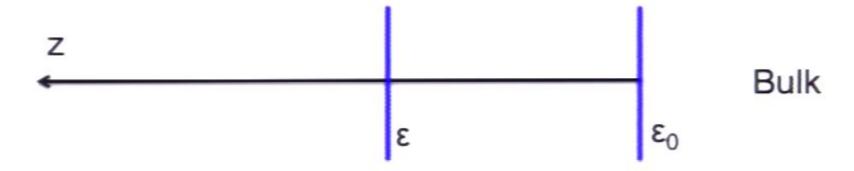


flow of boundary conditions



flow of S_P

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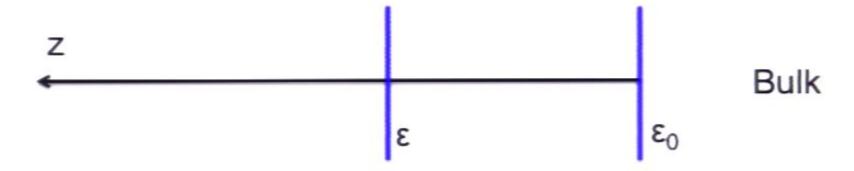


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Page 120/396

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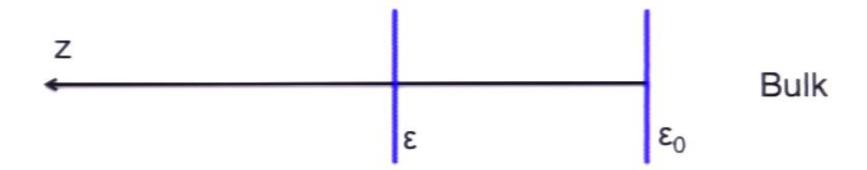


flow of boundary conditions



Page 121/396

Standard AdS/CFT procedure: Solve classical equations with boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).



s we shift the boundary to $z = \varepsilon$, we should not change physics, e. keep the same classical solution, which requires different bundary conditions at $z = \varepsilon$.

Shift the cutoff surface

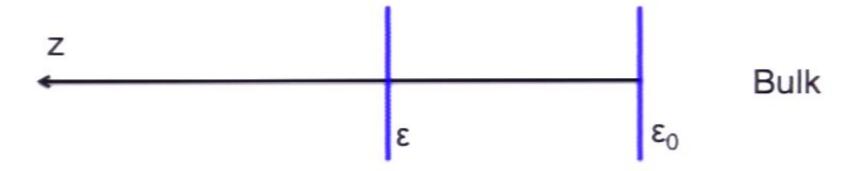


flow of boundary conditions



Page 122/396

Standard AdS/CFT procedure: Solve classical equations vith boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).

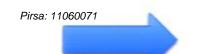


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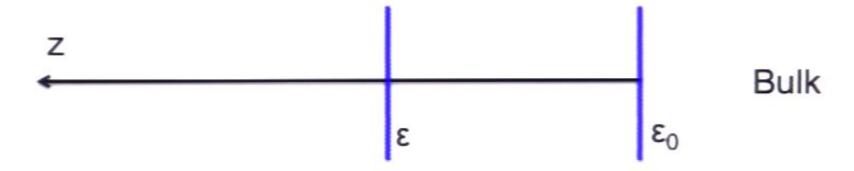
Shift the cutoff surface



flow of boundary conditions



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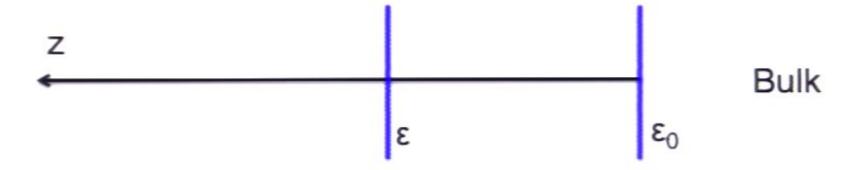


flow of boundary conditions



Page 124/396

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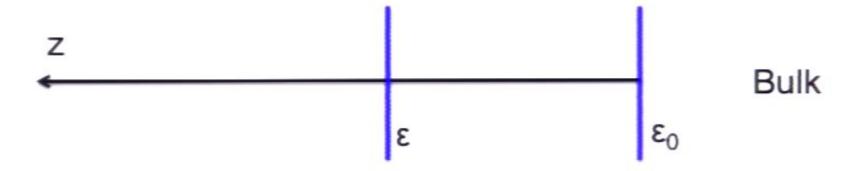
Shift the cutoff surface



flow of boundary conditions



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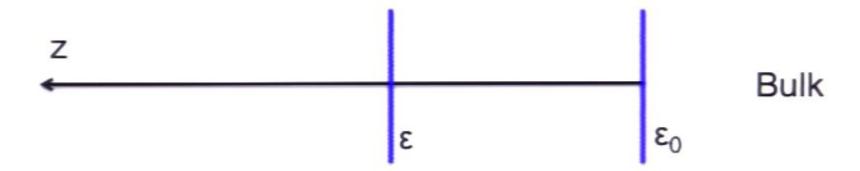


flow of boundary conditions



Page 126/396

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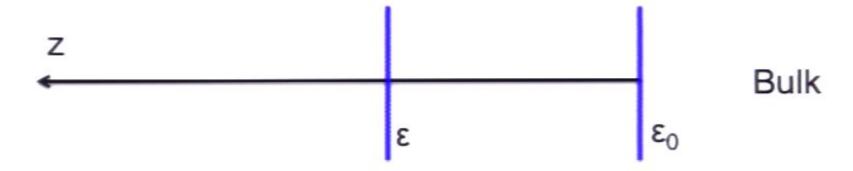


flow of boundary conditions



Page 127/396

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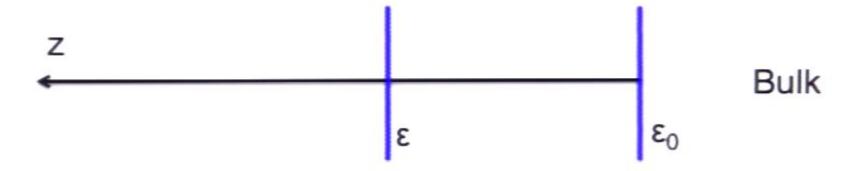


flow of boundary conditions



Page 128/396

Standard AdS/CFT procedure: Solve classical equations vith boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).

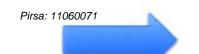


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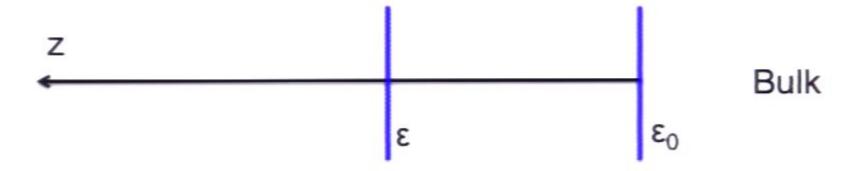


flow of boundary conditions



flow of S_P

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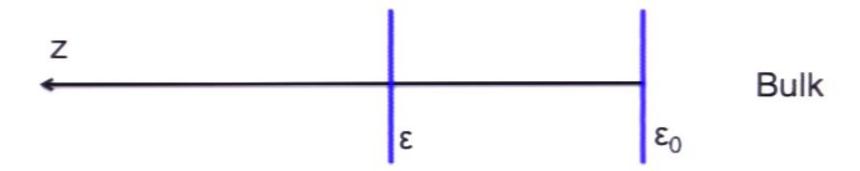


flow of boundary conditions



Page 130/396

Standard AdS/CFT procedure: Solve classical equations with boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).

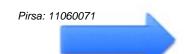


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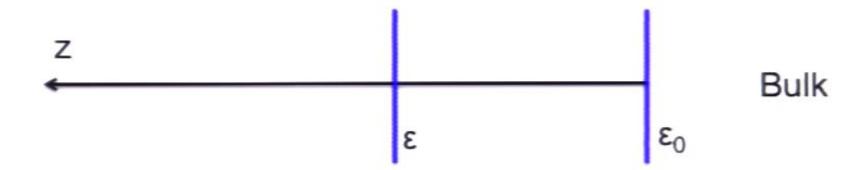


flow of boundary conditions



Page 131/396

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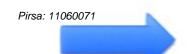


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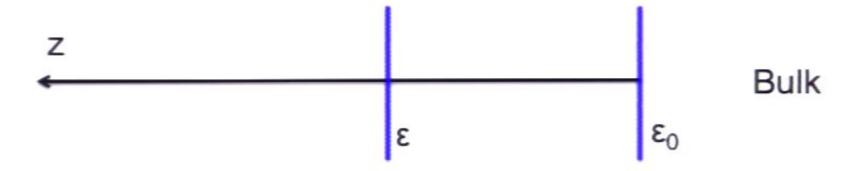


flow of boundary conditions



Page 132/396

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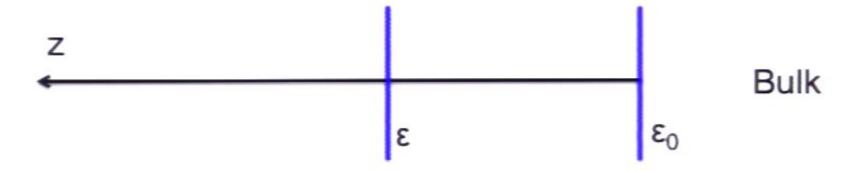


flow of boundary conditions



Page 133/396

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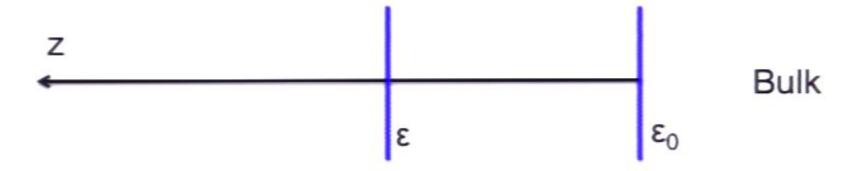


flow of boundary conditions



Page 134/396

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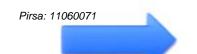


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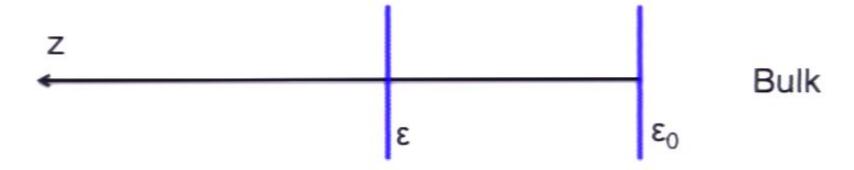


flow of boundary conditions



Page 135/396

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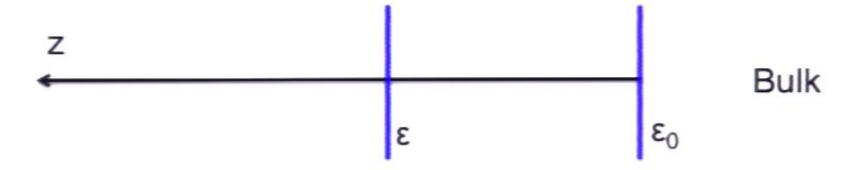


flow of boundary conditions



Page 136/396

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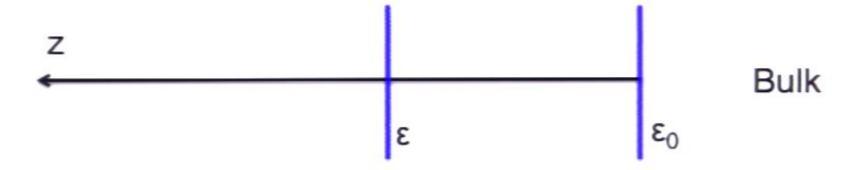
flow of boundary conditions

Page 137/396



flow of S_p

Standard AdS/CFT procedure: Solve classical equations with boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).



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Shift the cutoff surface

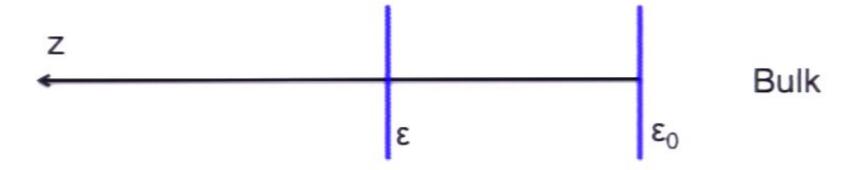


flow of boundary conditions



Page 138/396

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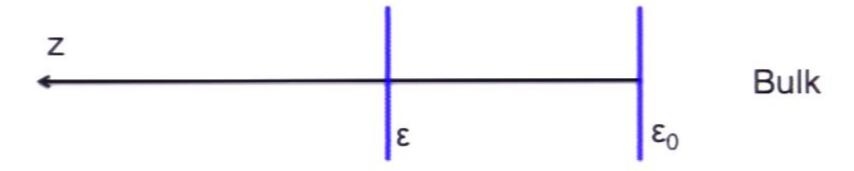


flow of boundary conditions



flow of S_P

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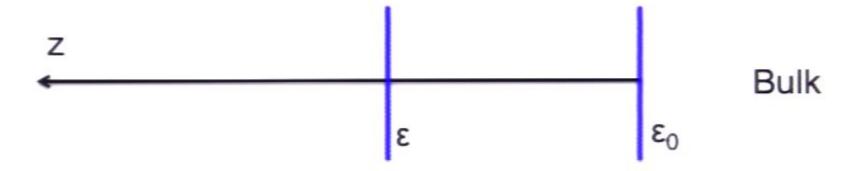


flow of boundary conditions



Page 140/396

Standard AdS/CFT procedure: Solve classical equations vith boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).

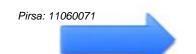


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Shift the cutoff surface

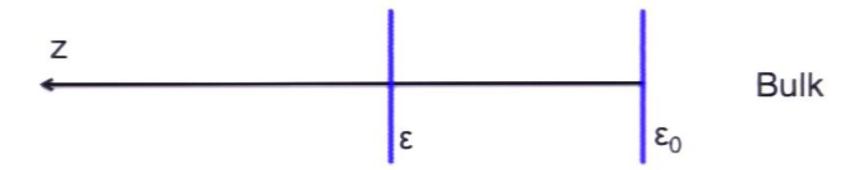


flow of boundary conditions



flow of S_R

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Shift the cutoff surface

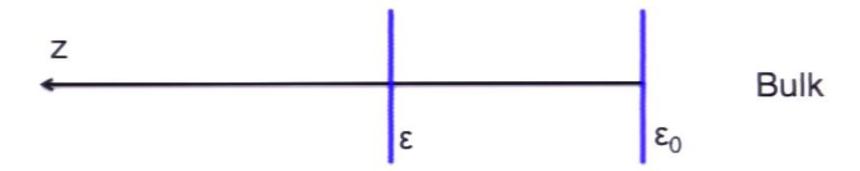


flow of boundary conditions



Page 142/396

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Shift the cutoff surface

flow of S_R

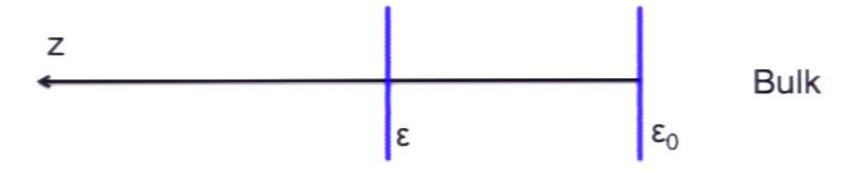


flow of boundary conditions



Page 143/396

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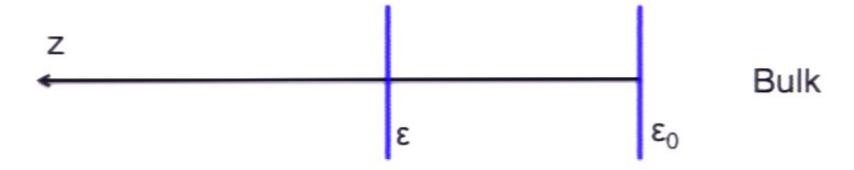
Shift the cutoff surface



flow of boundary conditions



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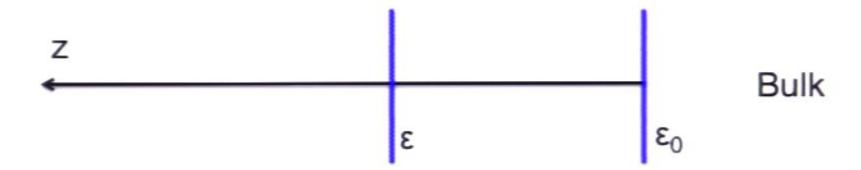


flow of boundary conditions



Page 145/396

Standard AdS/CFT procedure: Solve classical equations with boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).



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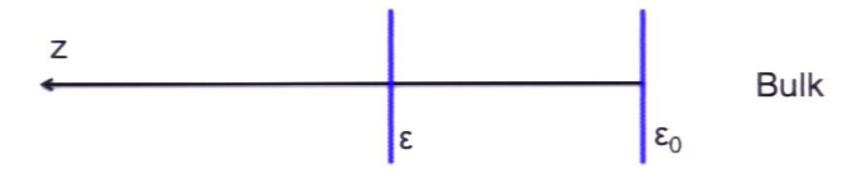


flow of boundary conditions



Page 146/396

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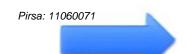


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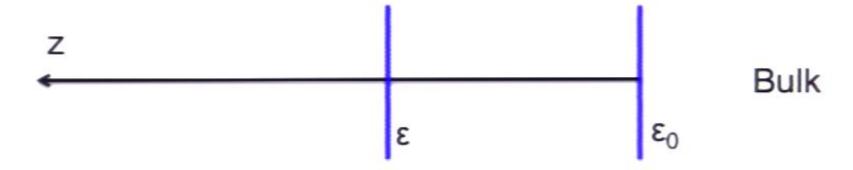


flow of boundary conditions



Page 147/396

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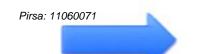


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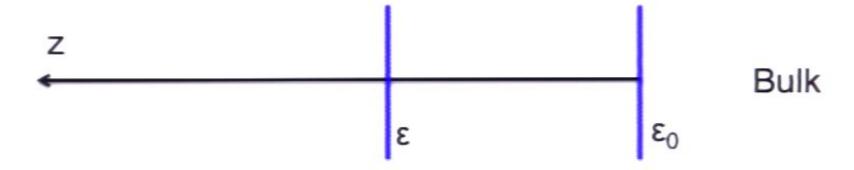


flow of boundary conditions



flow of S_R

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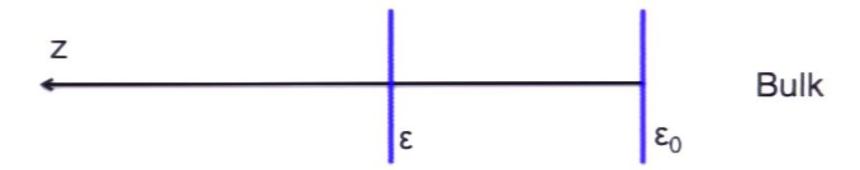


flow of boundary conditions



Page 149/396

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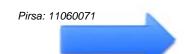


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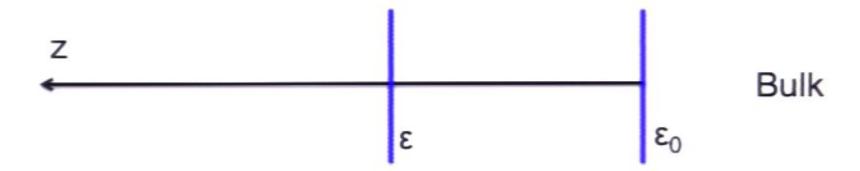


flow of boundary conditions



Page 150/396

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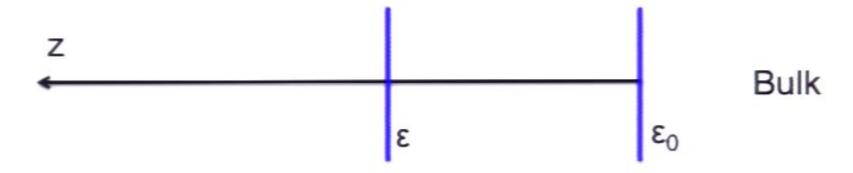


flow of boundary conditions



Page 151/396

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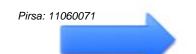


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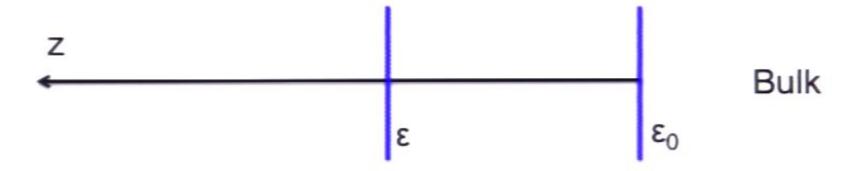


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Page 152/396

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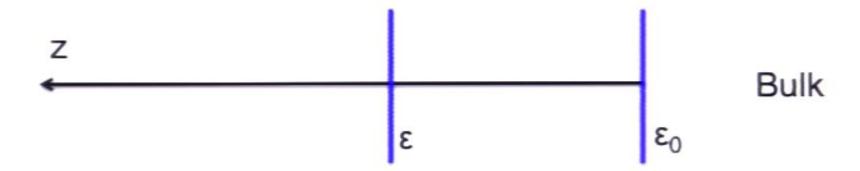


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Page 153/396

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Shift the cutoff surface



flow of boundary conditions



flow of S_R

Pirsa: 11060071 Page 155/396

Physics should not depend on where we choose z=ε suface



Flow equation for S_B [ε]

Semi-classical limit: Hamilton-Jacobi equation

$$\partial_{\epsilon} S_B[\phi, \epsilon] = -\int_{z=\epsilon} d^d x \, H\left(\phi, \Pi = \frac{\delta S_B}{\delta \phi}\right)$$

H: bulk Hamiltonian corresponding to z-foliation.

$$S_B[\phi,\epsilon] = \int \sum_i g_i(\epsilon)\phi^n$$
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Flow equation for S_B [ε]

Semi-classical limit: Hamilton-Jacobi equation

$$\partial_{\epsilon} S_B[\phi, \epsilon] = -\int_{z=\epsilon} d^d x \, H\left(\phi, \Pi = \frac{\delta S_B}{\delta \phi}\right)$$

H: bulk Hamiltonian corresponding to z-foliation.

$$S_B[\phi,\epsilon] = \int \sum_i g_i(\epsilon)\phi^n$$
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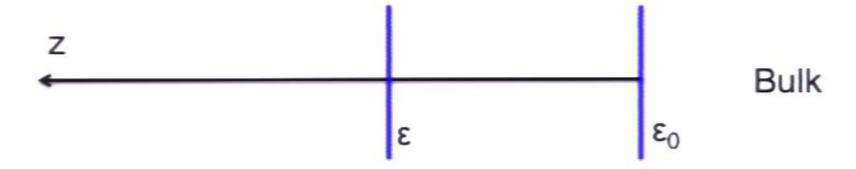
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Pirsa: 11060071 Page 172/396

Standard AdS/CFT procedure: Solve classical equations vith boundary conditions (e.g. Dirichlet) at $z = \varepsilon_0$ (and regularity condition in the interior).



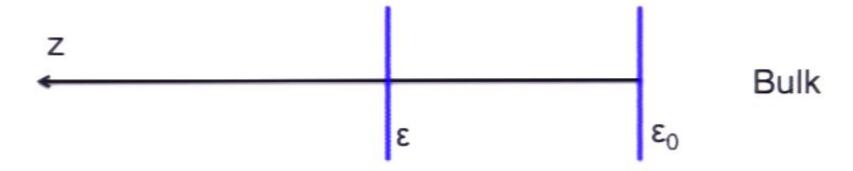
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Shift the cutoff surface



flow of boundary conditions

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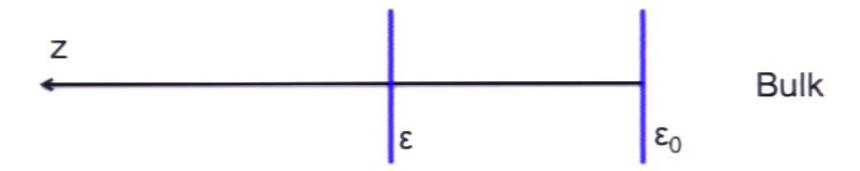


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Page 174/396

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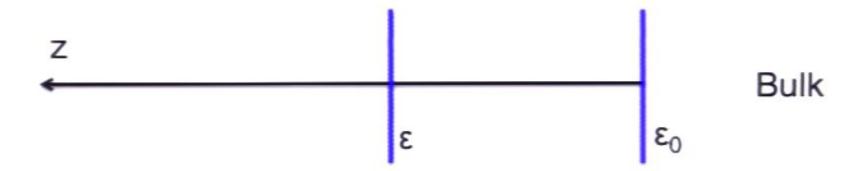


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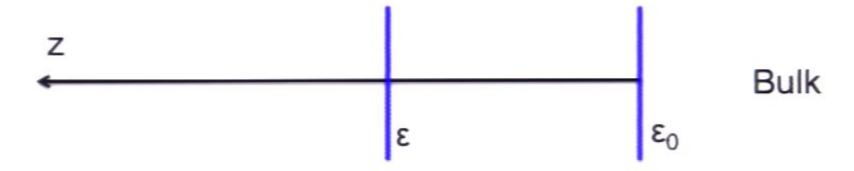


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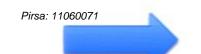


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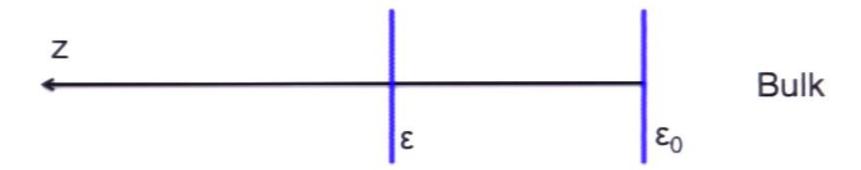


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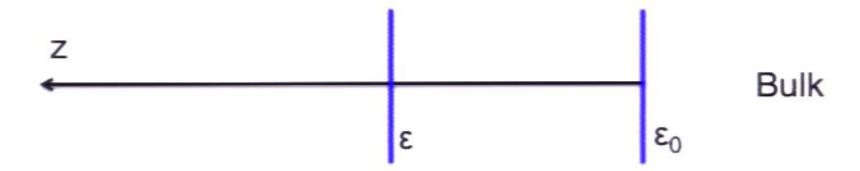


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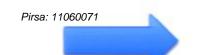


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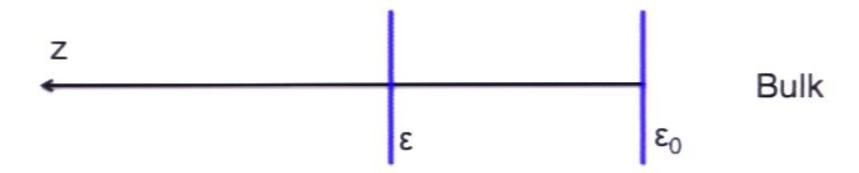


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Pirsa: 11060071 Page 181/396

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Pirsa: 11060071 Page 226/396

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simplify calculations, organize physics better

Pirsa: 11060071 Page 227/396

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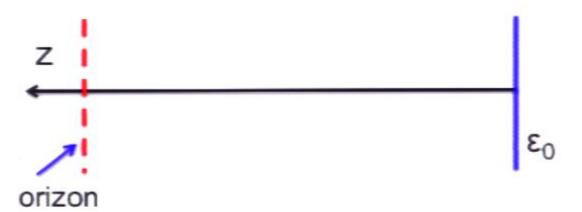
Finite temperature:

Pirsa: 11060071 Page 228/396

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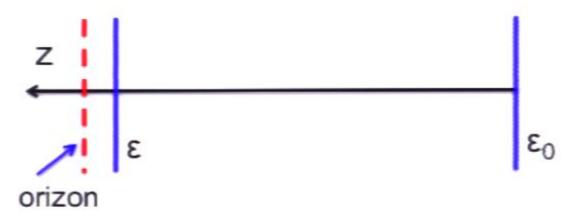


Pirsa: 11060071 Page 229/396

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Hydrodynamics: enough to solve the equations in the near-horizon region with S_B membrane paradigm ...

Pirsa: 11060071 Page 231/396

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Pirsa: 11060071 Page 232/396

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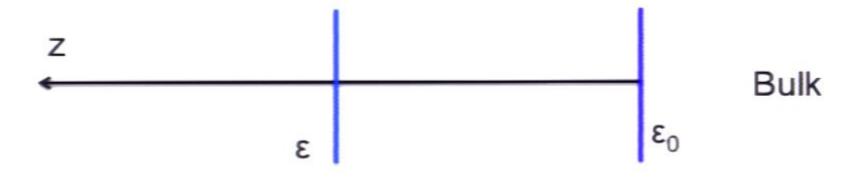
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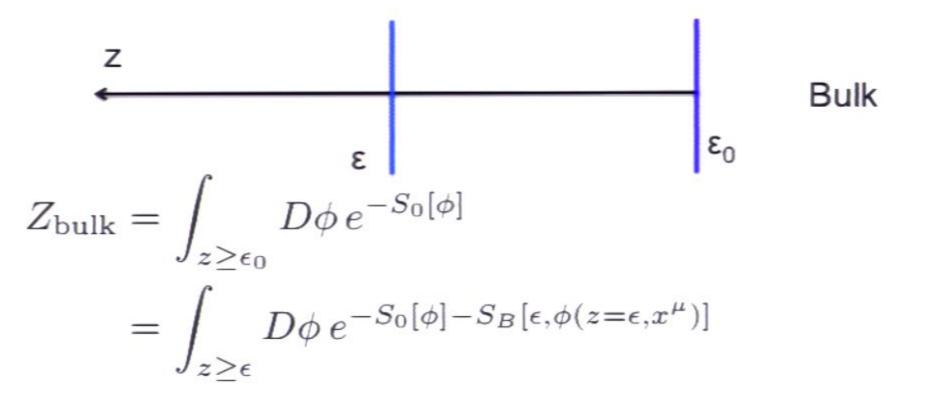
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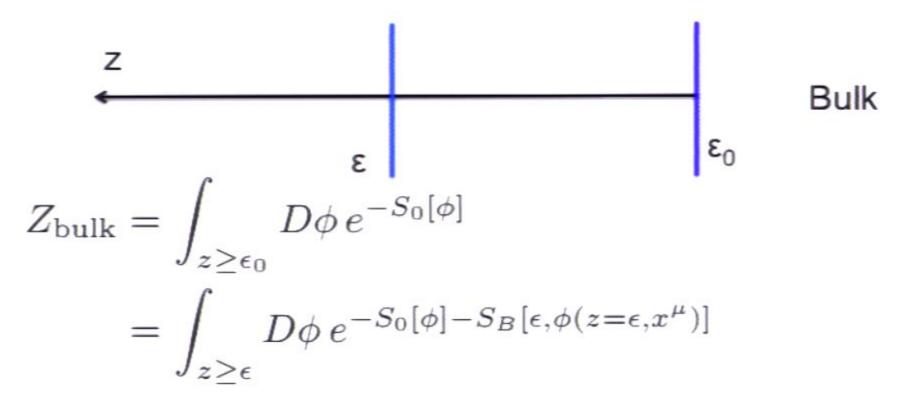
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Pirsa: 11060071 Page 235/396

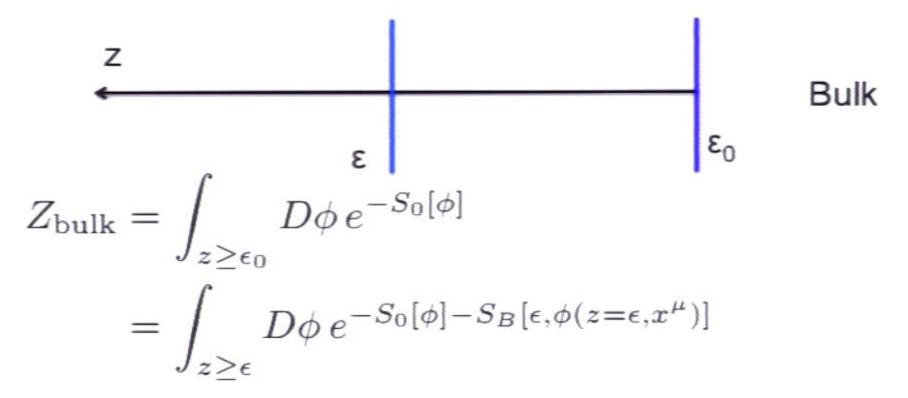


Pirsa: 11060071



de Boer, Verlinde and Verlinde (2000): Flow of $\log Z_{\mathrm{bulk}}[\epsilon_0]$

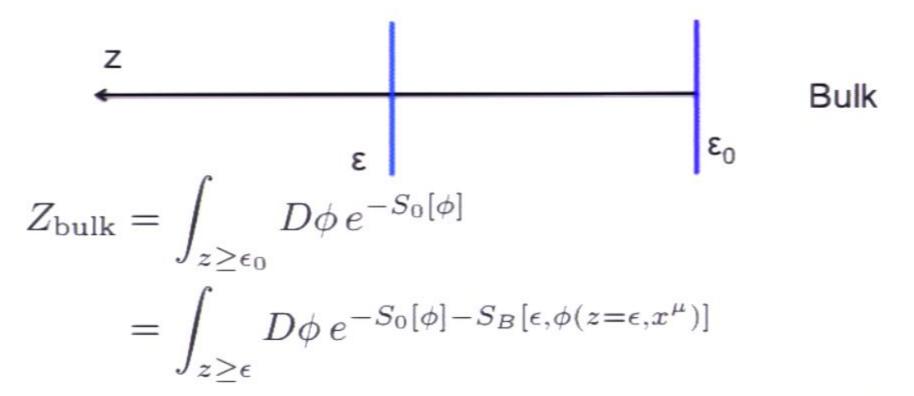
Pirsa: 11060071 Page 237/396



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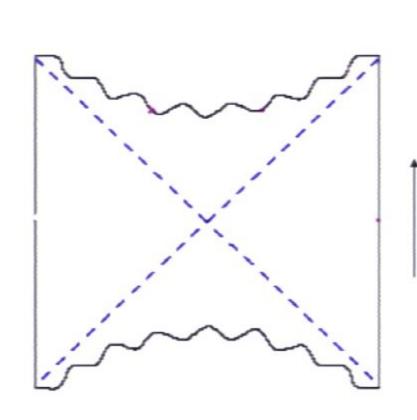


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It depends on IR data, cannot be Wilsonian

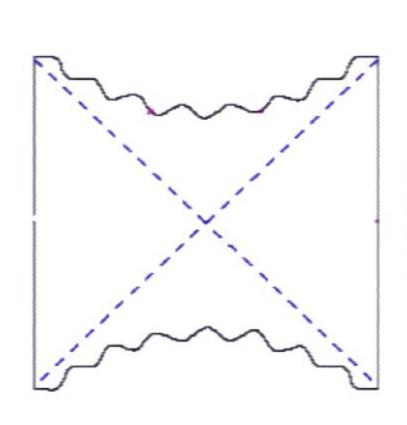
Both $\log Z_{\mathrm{bulk}}[\epsilon_0]$ and S_B satisfy Hamilton-Jacobi equation, but contains different physics

Pirsa: 11060071 Page 240/396



$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

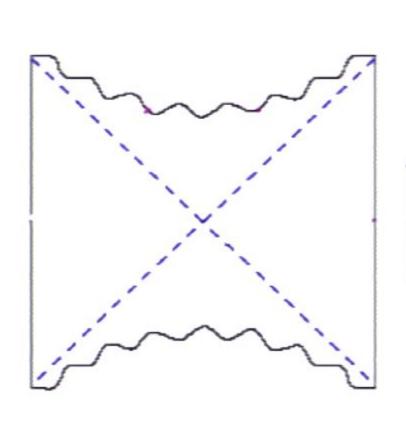
Pirsa: 11060071 Page 241/396



$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

$$f = r^2 + 1 - \frac{\mu}{r^2}$$
 Boundary:
$$f(r \to \infty) = +\infty$$

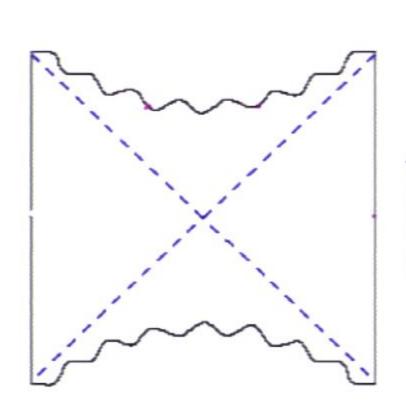
Pirsa: 11060071 Page 242/396



$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

Boundary: $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

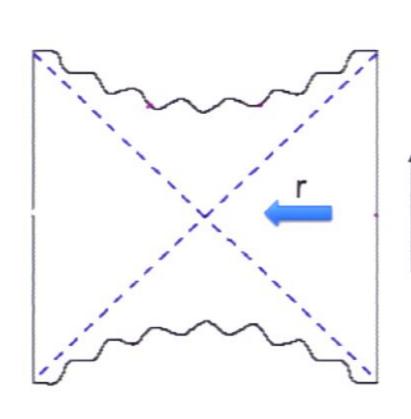


$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

Boundary: $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

Singularity: $f(r \to 0) = -\infty$

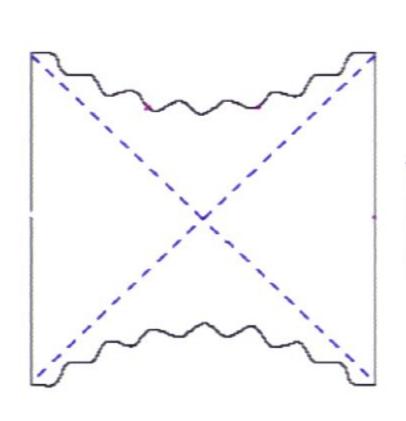


$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

Boundary: $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

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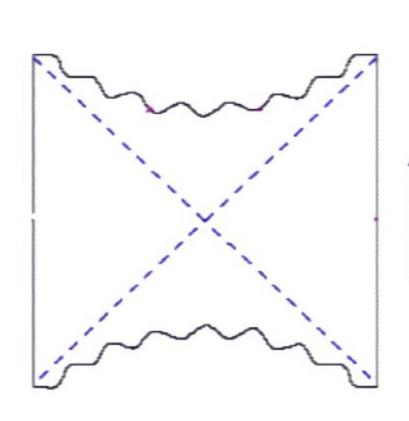


$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

Boundary: $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

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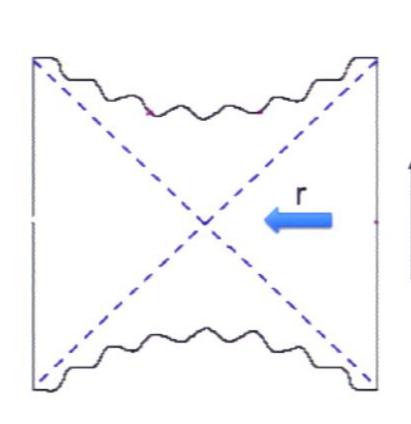


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$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

Boundary: $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

Pirsa: 11060071



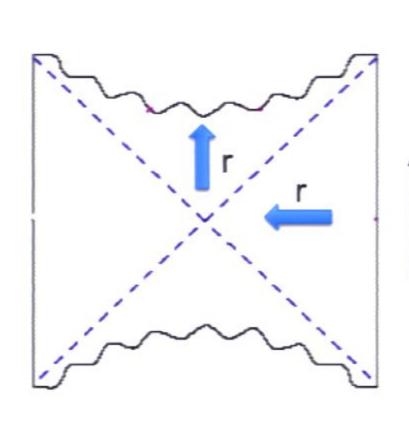
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Horizon: $f(r_0) = 0$

Singularity: $f(r \to 0) = -\infty$

Inside the horizon:



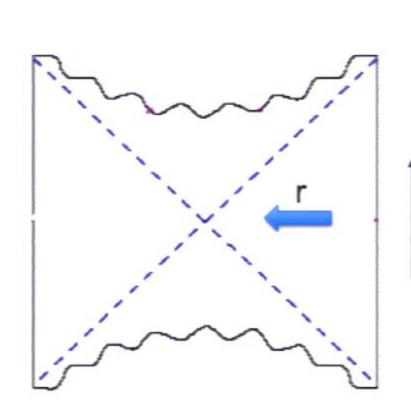
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Horizon: $f(r_0) = 0$

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Inside the horizon:



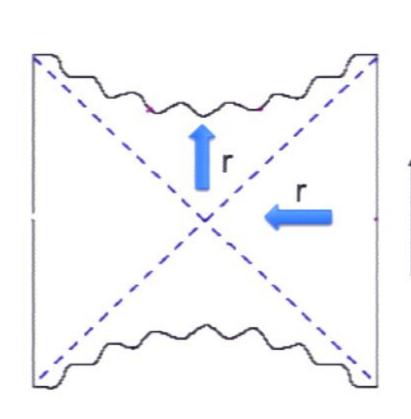
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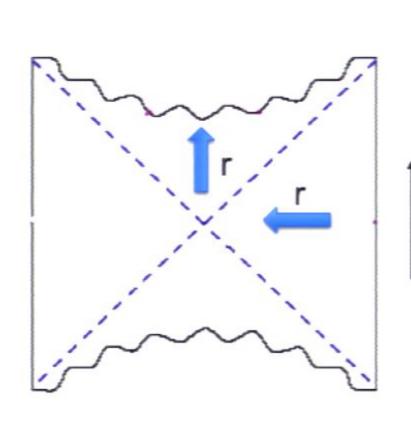
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Inside the horizon: time dynamically generated



$$ds^{2} = -fdt^{2} + \frac{1}{f}dr^{2} + r^{2}d\Omega^{2}$$
$$f = r^{2} + 1 - \frac{\mu}{r^{2}}$$

 $f(r \to \infty) = +\infty$

Horizon: $f(r_0) = 0$

Singularity: $f(r \to 0) = -\infty$

Inside the horizon: time dynamically generated

How to describe the region inside the horizon using AdS/CFT?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Pirsa: 11060071 Page 253/396

$$ds^2 = -f dt^2 + \frac{1}{f} dr^2 + r^2 d\Omega^2$$

 Dutside the horizon: $f>0$

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

Pirsa: 11060071 Page 254/396

$$ds^2 = -f dt^2 + \frac{1}{f} dr^2 + r^2 d\Omega^2$$

 Dutside the horizon: $f>0$

$$E_{\rm boundary} = \sqrt{f}E_{\rm prop}$$



boundary: UV

horizon: IR

Pirsa: 11060071 Page 255/396

$$ds^2 = -f dt^2 + \frac{1}{f} dr^2 + r^2 d\Omega^2$$

 Dutside the horizon: $f>0$

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Pirsa: 11060071

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

boundary: UV

horizon: IR

 $E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

Eprop: now a spatial momentum

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

boundary: UV

horizon: IR

 $E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

 $\mathcal{E}_{\text{prop}}$: now a spatial momentum

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$



Singularity:

 $E_{\rm boundary} \to \pm i\infty$

 Z_{prop} : now a spatial momentum

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



horizon: IR

Bulk evolution in r



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Inside the horizon: f < 0

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Singularity:

$$E_{\rm boundary} \to \pm i\infty$$

Horizon: $E_{\text{boundary}} \to 0$

Eprop: now a spatial momentum

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

Singularity:

 $E_{\rm boundary} \to \pm i\infty$

Pprop: now a spatial momentum

Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 22/39?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$



horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$



 $E_{\rm boundary} \to \pm i\infty$

Horizon: $E_{\text{boundary}} \to 0$

 $\mathcal{I}_{\mathrm{prop}}$: now a spatial momentum

Bulk evolution in r (time evolution)



flow in imaginary energy 202/39?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Singularity:

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$



 $E_{\rm boundary} \to \pm i\infty$

 $\mathcal{I}_{\mathrm{prop}}$: now a spatial momentum

Horizon: $E_{\text{boundary}} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 20/39?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

Singularity:

 $E_{\rm boundary} \to \pm i\infty$

Eprop: now a spatial momentum

Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 202/39??

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

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 $E_{\rm boundary} \to \pm i\infty$

Eprop: now a spatial momentum

Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 28/3??

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

Singularity:

 $E_{\rm boundary} \to \pm i\infty$

omentum Horizon: $E_{\text{boundary}} \to 0$

 $\mathbb{Z}_{\mathrm{prop}}$: now a spatial momentum

Bulk evolution in r (time evolution)



flow in imaginary energy 28/39?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$



 $E_{\rm boundary} \to \pm i\infty$

Eprop: now a spatial momentum

Horizon: $E_{\text{boundary}} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 27/39??

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$



 $E_{\rm boundary} \to \pm i\infty$

Eprop: now a spatial momentum Horizo

Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 27/39?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

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boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

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Singularity:

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Horizon: $E_{\text{boundary}} \to 0$

Eprop: now a spatial momentum

Bulk evolution in r time evolution)



flow in imaginary energy 272/396?

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

$$E_{\text{boundary}} = i\sqrt{-f}E_{\text{prop}}$$

Singularity:

 $E_{\rm boundary} \to \pm i\infty$

Eprop: now a spatial momentum Horizo

Horizon: $E_{\mathrm{boundary}} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy, 273/39?

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boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

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 $E_{\rm boundary} \to \pm i\infty$

 $\mathcal{I}_{\mathrm{prop}}$: now a spatial momentum

Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 274/39??

Pirsa: 11060071 Page 275/396

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

Outside the horizon: f > 0

$$E_{\text{boundary}} = \sqrt{f}E_{\text{prop}}$$

boundary: UV

horizon: IR

Singularity:

Bulk evolution in r



flow in energy scale

Inside the horizon: f < 0

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 $E_{\rm boundary} \to \pm i\infty$

 $\mathcal{I}_{\mathrm{prop}}$: now a spatial momentum

Horizon: $E_{\mathrm{boundary}} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy, 272/39??

$$ds^2 = -fdt^2 + \frac{1}{f}dr^2 + r^2d\Omega^2$$

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boundary: UV

horizon: IR

Bulk evolution in r



flow in energy scale

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 $E_{\rm boundary} \to \pm i\infty$

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Horizon: $E_{\rm boundary} \to 0$

Bulk evolution in r (time evolution)



flow in imaginary energy 27/39??

Pirsa: 11060071 Page 278/396

Consider a generic operator O, and

$$G_{+}(t, \vec{x}) = \langle O(t, \vec{x})O(0)\rangle_{\beta}$$

Pirsa: 11060071 Page 279/396

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black hole geometry



Pirsa: 11060071 Page 280/396

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black hole geometry



• continuous spectrum: $\omega \in (-\infty, +\infty)$ (due to presence of horizon)

Pirsa: 11060071 Page 281/396

Consider a generic operator O, and

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• continuous spectrum: $\omega \in (-\infty, +\infty)$ (due to presence of horizon)

Simple poles in the complex frequency plane.

Pirsa: 11060071 Page 282/396

Consider a generic operator O, and

$$G_{+}(t, \vec{x}) = \langle O(t, \vec{x})O(0)\rangle_{\beta}$$



- continuous spectrum: $\omega \in (-\infty, +\infty)$ (due to presence of horizon)
- Simple poles in the complex frequency plane.
- For a given (ω, \vec{k}) , $G_{+}(\omega, \vec{k})$ is uniquely associated with a complex bulk spacelike geodesic.

Pirsa: 11060071 Page 283/396

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Pirsa: 11060071 Page 284/396

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Pirsa: 11060071 Page 285/396

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Pirsa: 11060071 Page 286/396

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Pirsa: 11060071 Page 287/396

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Pirsa: 11060071 Page 288/396

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Pirsa: 11060071 Page 289/396

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black hole geometry $G_+(\omega, \vec{k})$



$$G_{+}(\omega, \vec{k})$$

- continuous spectrum: $\omega \in (-\infty, +\infty)$ (due to presence of horizon)
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$$(\omega,k) \rightarrow r_c(\omega,k)$$

$$G_{+}(t, \vec{x}) = \langle O(t, \vec{x})O(0)\rangle_{\beta}$$



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- Simple poles in the complex frequency plane.
- For a given (ω, \vec{k}) , $G_+(\omega, \vec{k})$ is uniquely associated with a complex bulk spacelike geodesic.
- Each geodesic is characterized by a Pirsa: 11060071 INC POINT

$$(\omega,k) \rightarrow r_c(\omega,k)$$

Consider a generic operator O, and

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black hole geometry $G_+(\omega, \vec{k})$



$$G_+(\omega, \vec{k})$$

- continuous spectrum: $\omega \in (-\infty, +\infty)$ (due to presence of horizon)
- Simple poles in the complex frequency plane.
- For a given (ω, \vec{k}) , $G_+(\omega, \vec{k})$ is uniquely associated with a complex bulk spacelike geodesic.
- Each geodesic is characterized by a Pirsa: 11060071 INC DOINT

$$(\omega,k) \rightarrow r_c(\omega,k)$$

$$G_{+}(t, \vec{x}) = \langle O(t, \vec{x})O(0)\rangle_{\beta}$$



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$$G_{+}(t, \vec{x}) = \langle O(t, \vec{x})O(0)\rangle_{\beta}$$

black hole geometry $G_+(\omega, \vec{k})$



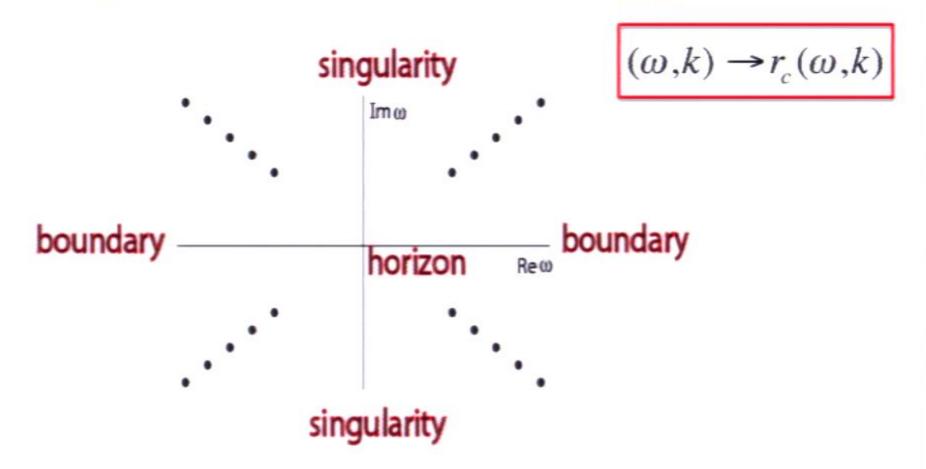
$$G_+(\omega, \vec{k})$$

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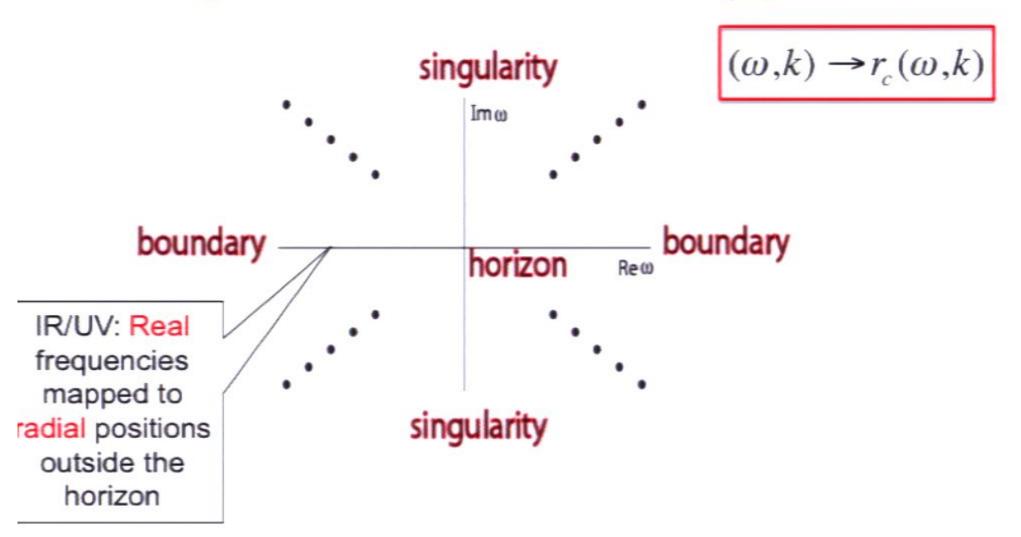
$$(\omega,k) \rightarrow r_c(\omega,k)$$

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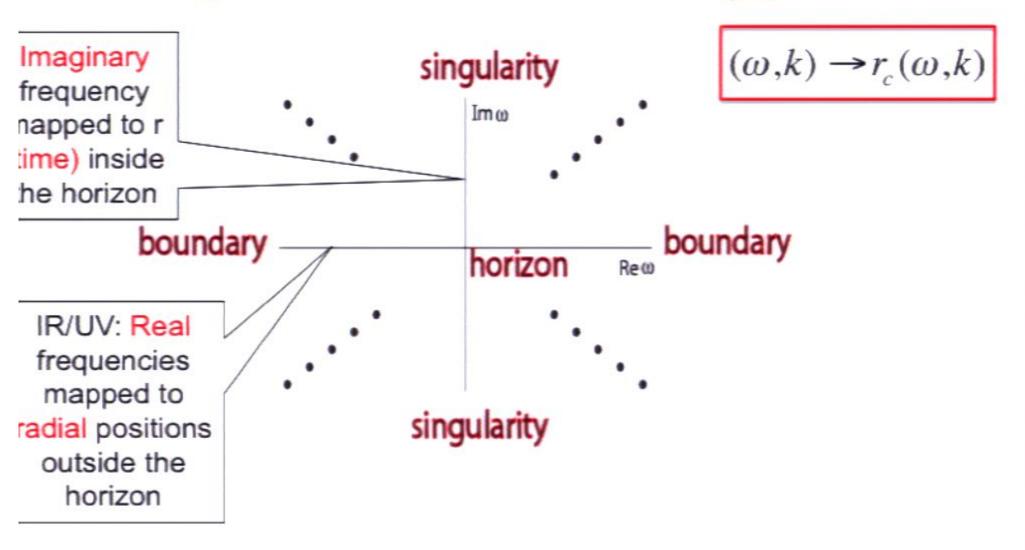
Pirsa: 11060071 Page 305/396



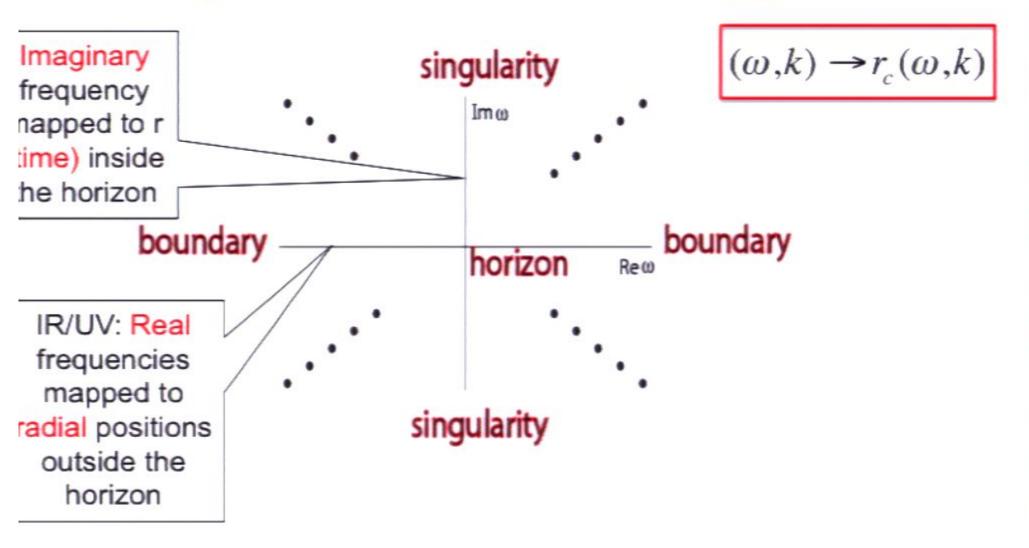
Pirsa: 11060071 Page 306/396



Pirsa: 11060071 Page 307/396



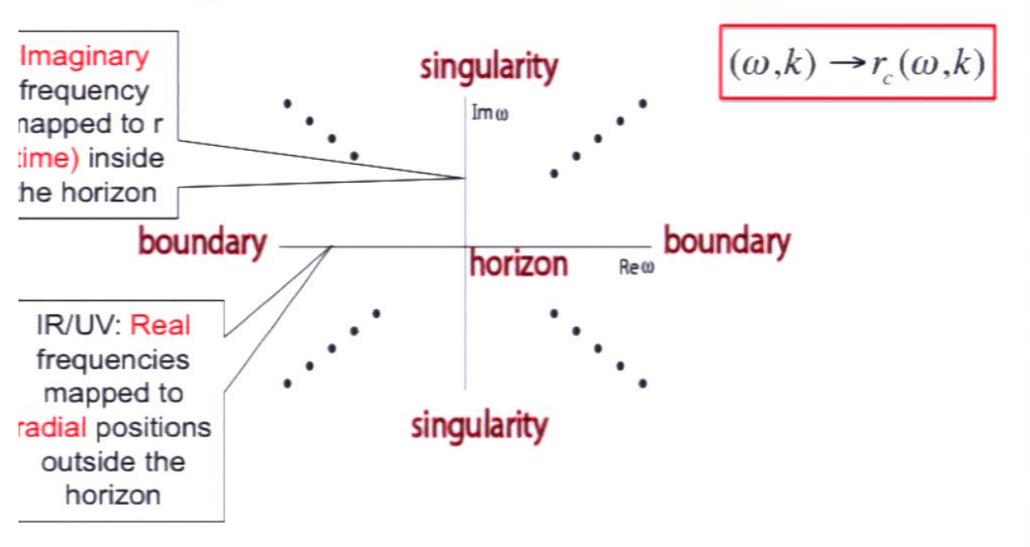
Pirsa: 11060071 Page 308/396



Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

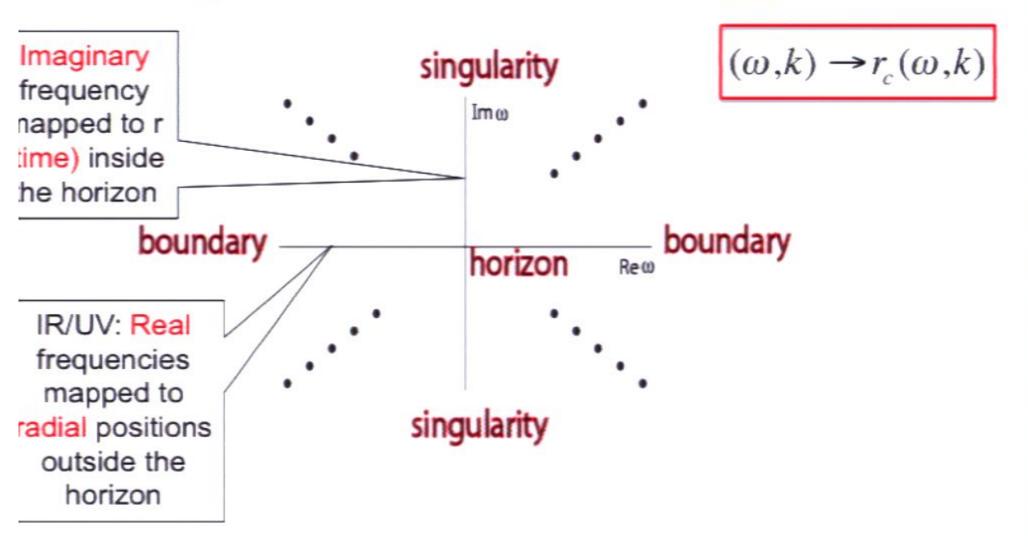
Page 309/396



Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

Page 310/396



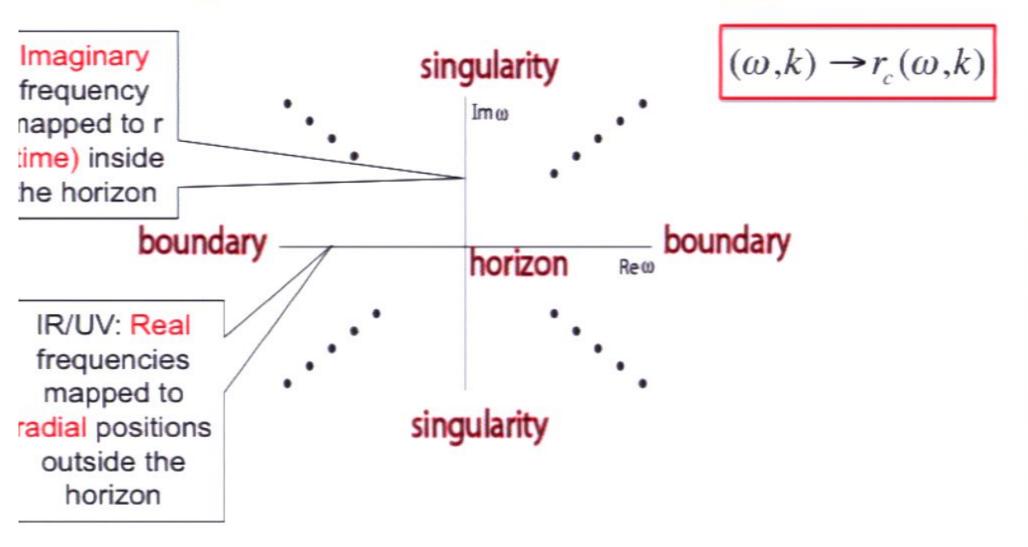
Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

Page 311/396

Pirsa: 11060071 Page 312/396

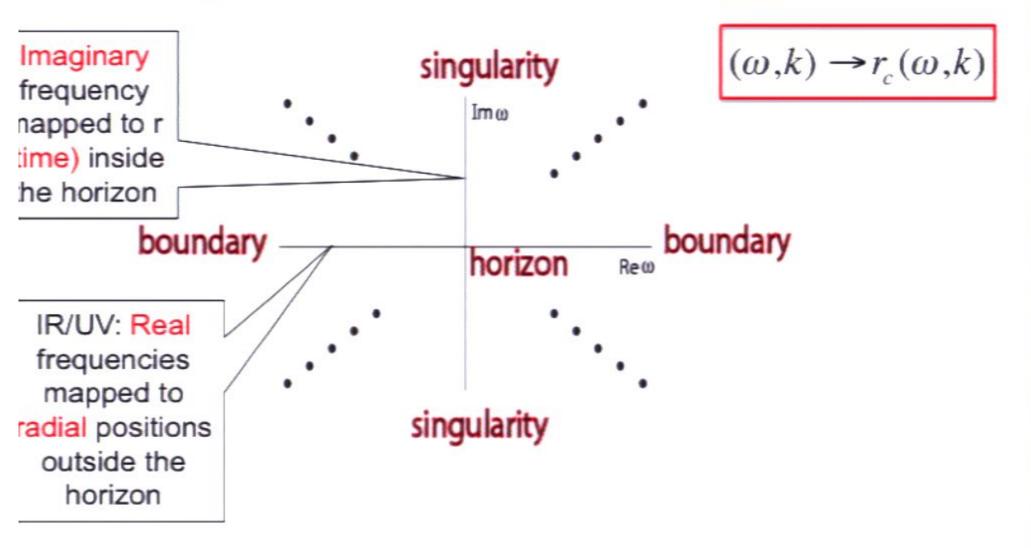
Pirsa: 11060071 Page 313/396



Consistent with picture obtained from the heuristic

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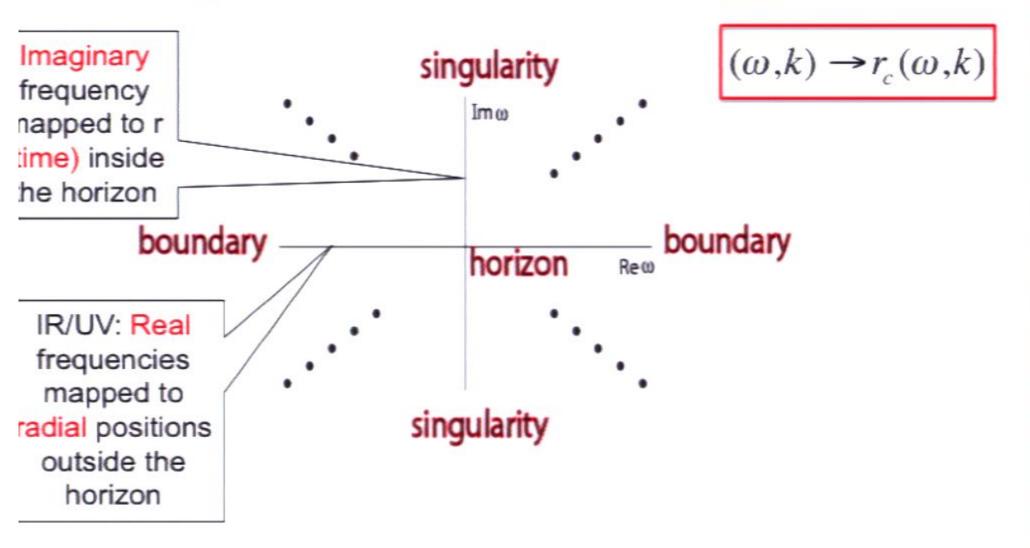
Page 314/396



Consistent with picture obtained from the heuristic

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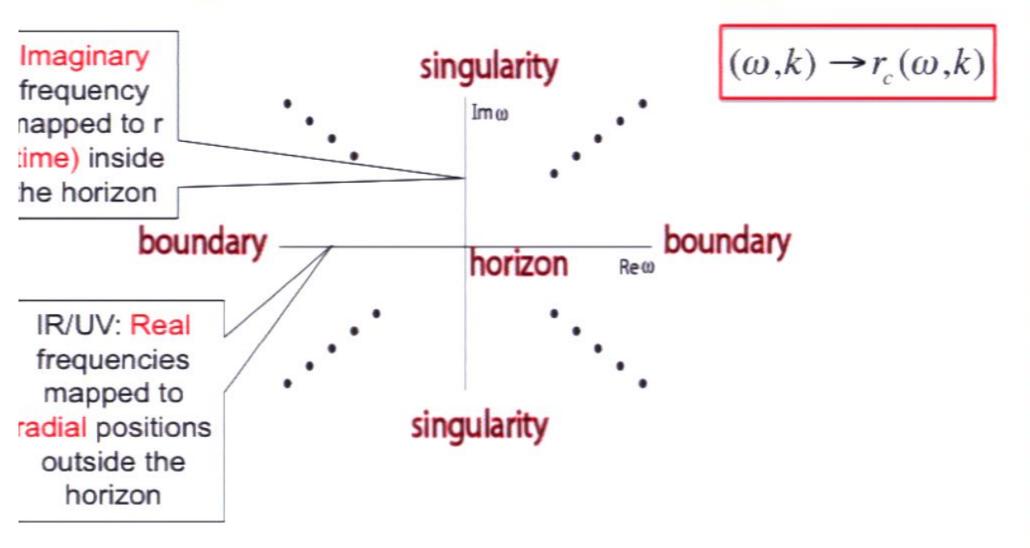
Page 315/396



Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

Page 316/396



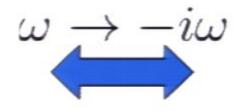
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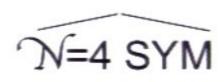
Pirsa: 11060071 "redshift" analysis.

Page 317/396

Pirsa: 11060071 Page 318/396

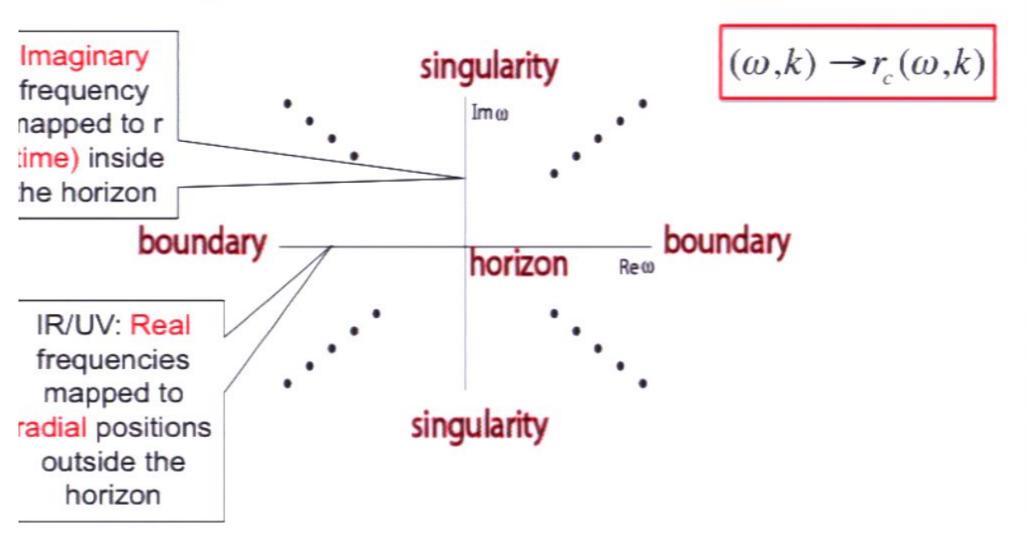






Pirsa: 11060071 Page 319/396

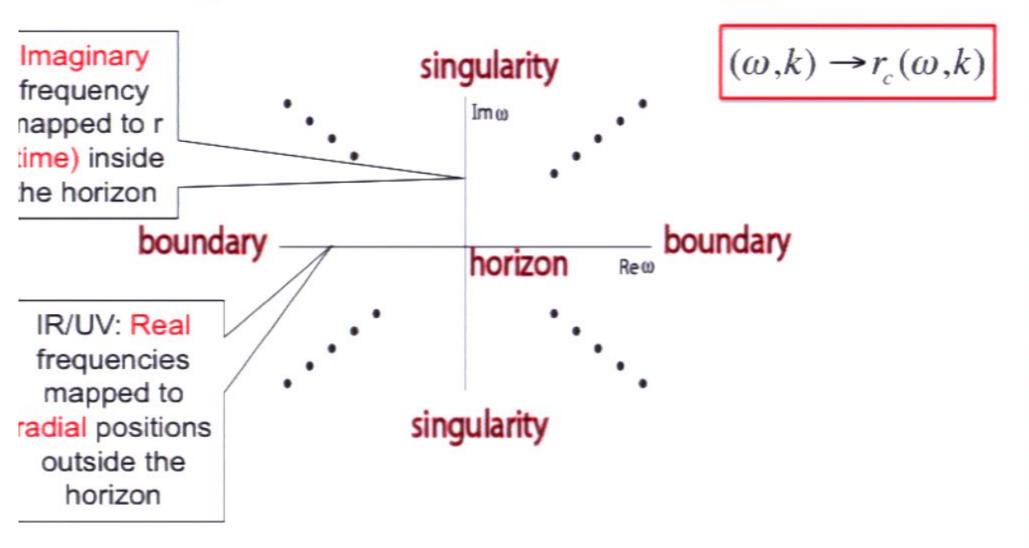
Pirsa: 11060071 Page 320/396



Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

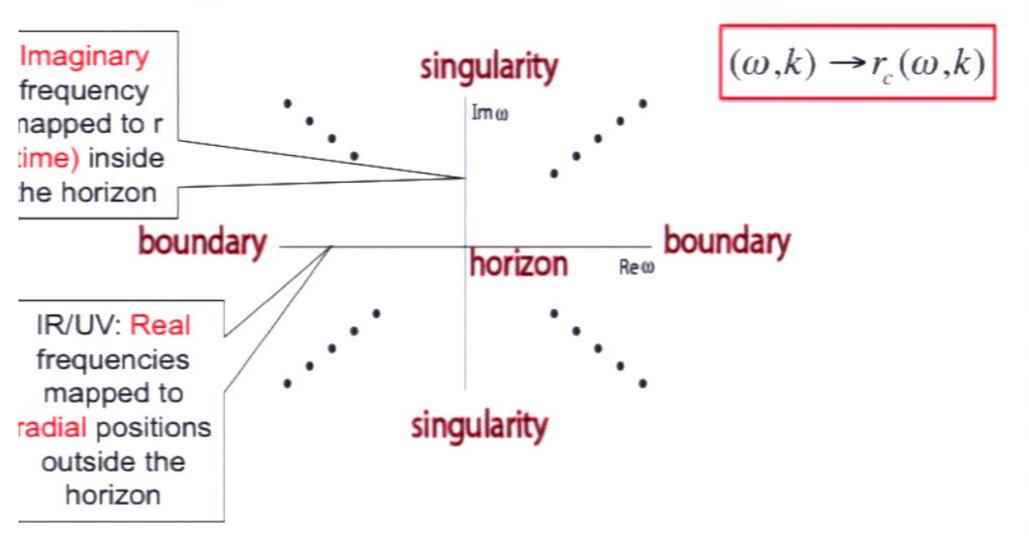
Page 321/396



Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

Page 322/396



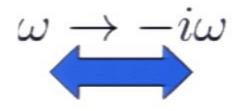
Consistent with picture obtained from the heuristic

Pirsa: 11060071 "redshift" analysis.

Page 323/396

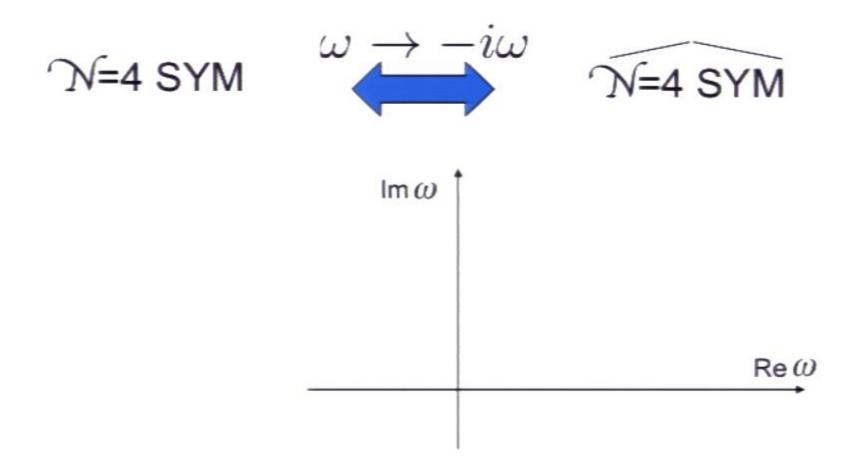
Pirsa: 11060071 Page 324/396

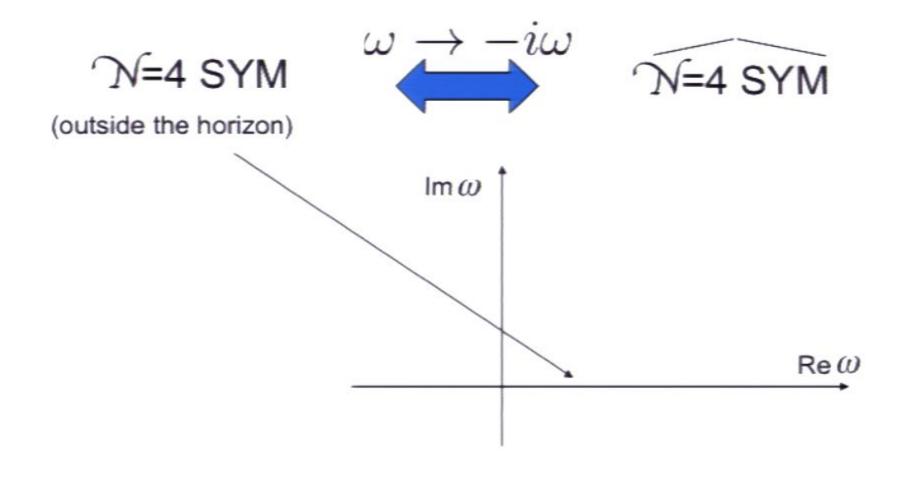




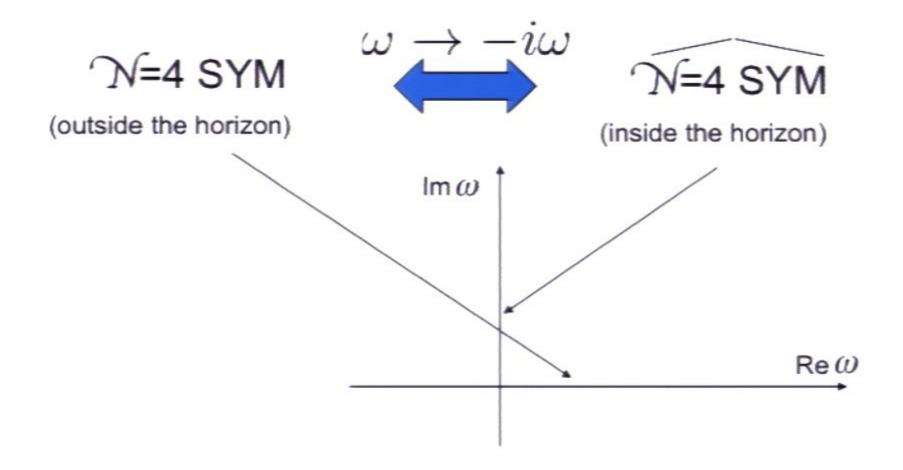


Pirsa: 11060071 Page 325/396

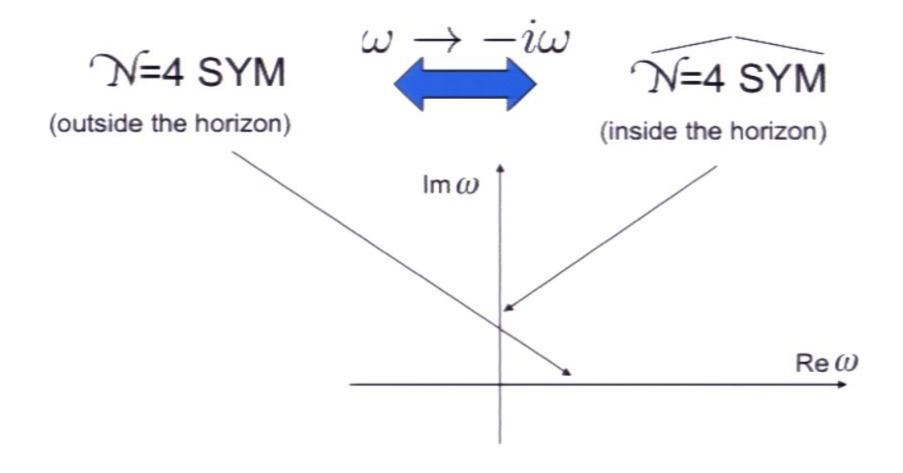




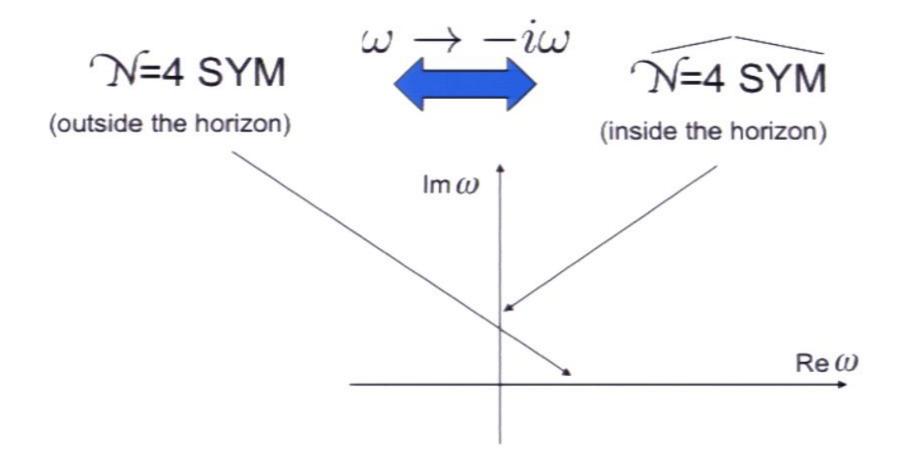
Pirsa: 11060071 Page 327/396



Pirsa: 11060071 Page 328/396



Different from the standard Euclidean analytic continuation!



Different from the standard Euclidean analytic continuation!

GR fnW 13 GE(WE) Rew ×

60071

Page 331/306

Gt=GR-GA GA XX X Rew

GA XX X Rew

Gt=GR-GA (XA X Rew

Gt=GR-GA GA xX X Rew

Gt=GR-GA GA X Rew

Gt=GR-GA GA xX X Rew

Pirsa

Page 337/396

Gt=GR-GA GA x X Rew

Rew

Gt=GR-GA GA XX X Rew

Gt=GR-GA GA

Gt=GR-GA GA X Rew

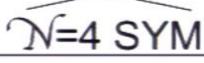
UV of N=4 SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 345/396

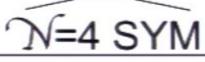
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 346/396

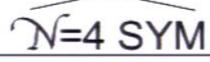
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 347/396

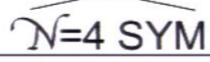
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 348/396

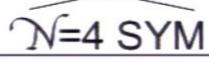
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 349/396

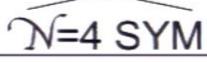
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 350/396

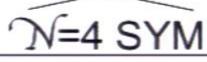
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 351/396

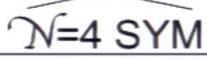
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 352/396

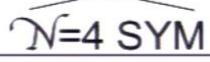
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 353/396

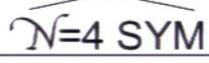
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Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 354/396

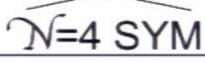
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 355/396

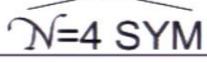
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 356/396

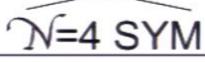
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 357/396

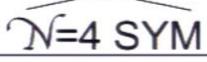
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 358/396

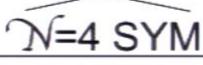
UV of $\mathcal{N}=4$ SYM at finite T



Low energies



Analytic continuation to





UV of N=4 SYM

Boundary of AdS black hole



Horizon (from outside)



Horizon (from inside)



singularity

Page 359/396

The story we just described is at the large N limit.

What happens at finite N?

Pirsa: 11060071 Page 360/396

Correlation functions at finite N

 $\mathcal{N}=4$ SYM theory on S³ at finite N:

Pirsa: 11060071 Page 361/396

Correlation functions at finite N

 $\mathcal{N}=4$ SYM theory on S³ at finite N:

Discrete energy spectrum.

Pirsa: 11060071 Page 362/396

Correlation functions at finite N

 $\mathcal{N}=4$ SYM theory on S³ at finite N:

- Discrete energy spectrum.
- Wightman functions have simple analytic structure:

$$G_{+}(\omega, l) = 2\pi \sum_{m,n} e^{-\beta E_m} \rho_{mn} \delta(\omega - E_n + E_m)$$

A discrete sum of delta functions

Pirsa: 11060071 Page 363/396

Pirsa: 11060071 Page 364/396



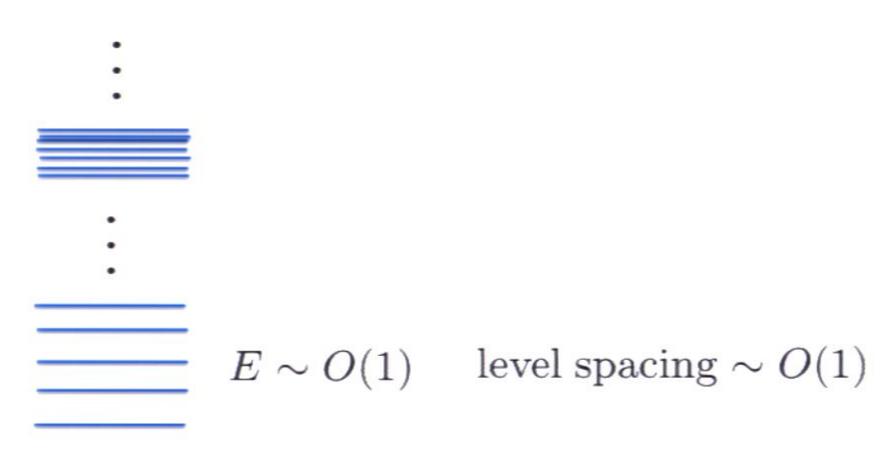
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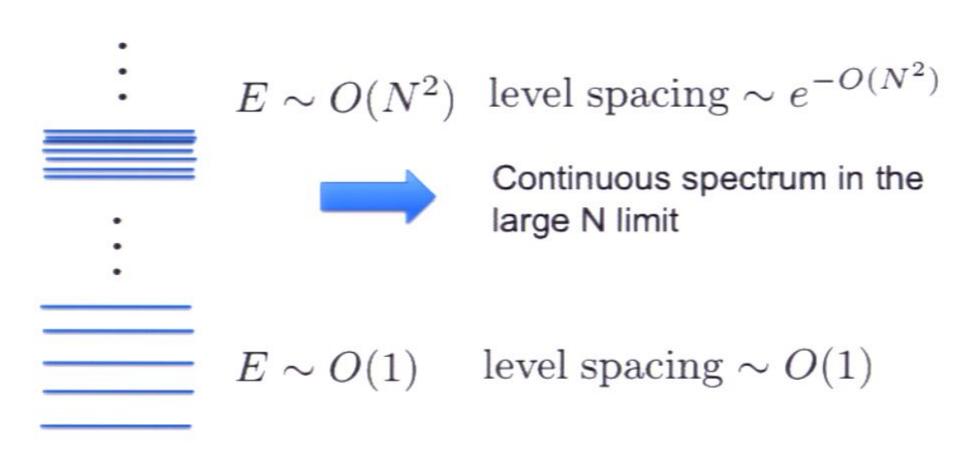
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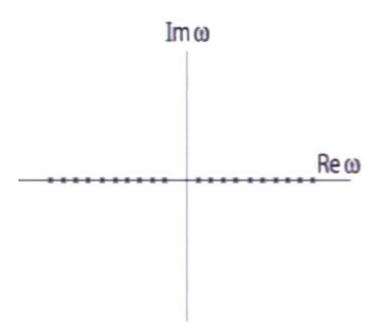


$$\vdots \qquad E \sim O(N^2) \quad \text{level spacing} \sim e^{-O(N^2)}$$

$$\vdots \qquad \qquad E \sim O(1) \quad \text{level spacing} \sim O(1)$$

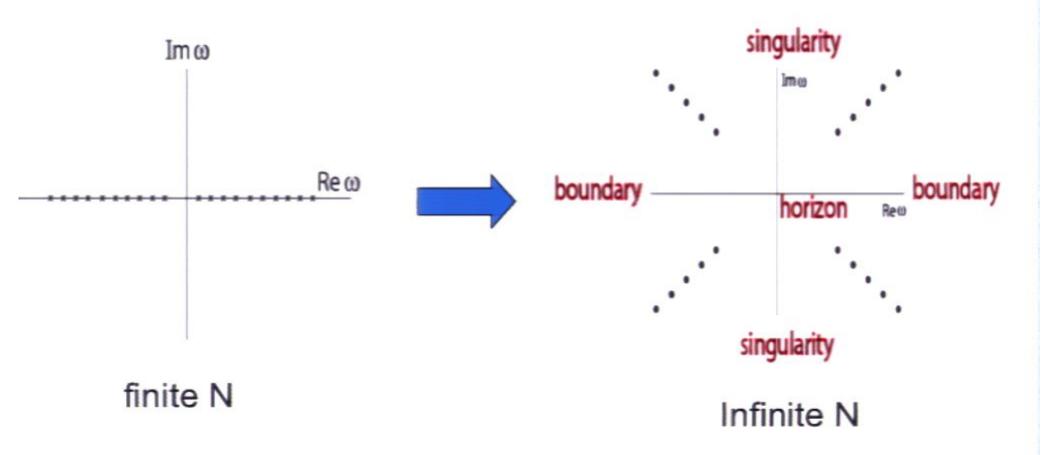


Pirsa: 11060071 Page 369/396

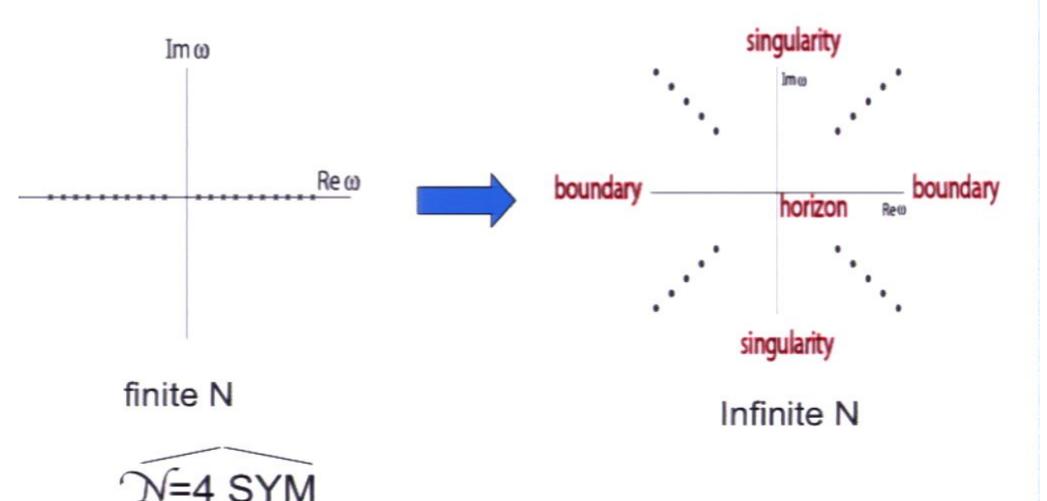


finite N

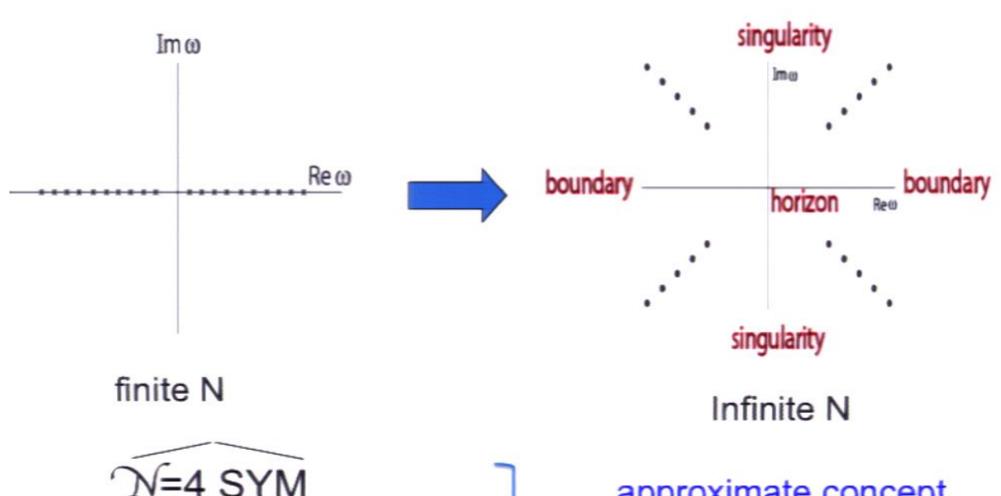
Pirsa: 11060071 Page 370/396



Pirsa: 11060071 Page 371/396



spacetime inside the horizon Pirsa: 11060071 ding singularities:



spacetime inside the horizon

approximate concept appearing only in the large N limit Page 373/396

Pirsa: 11060071 Page 374/396

Pirsa: 11060071 Page 375/396

Likely enough to understand them from some simple matrix quantum mechanics, like

Pirsa: 11060071 Page 376/396

Likely enough to understand them from some simple matrix quantum mechanics, like

$$S = \frac{N}{2} \operatorname{tr} \int dt \left[(D_t M_1)^2 + (D_t M_2)^2 - \omega_0^2 (M_1^2 + M_2^2) + \lambda M_1 M_2 M_1 M_2 \right]$$

Pirsa: 11060071 Page 377/396

Thank You

Pirsa: 11060071 Page 378/396

Thank You

Pirsa: 11060071 Page 379/396

Likely enough to understand them from some simple matrix quantum mechanics, like

$$S = \frac{N}{2} \operatorname{tr} \int dt \left[(D_t M_1)^2 + (D_t M_2)^2 - \omega_0^2 (M_1^2 + M_2^2) + \lambda M_1 M_2 M_1 M_2 \right]$$

Pirsa: 11060071 Page 380/396

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Pirsa: 11060071 Page 381/396

No Signal

VGA-1

Pirsa: 11060071 Page 382/39

Pirsa: 11060071 Page 383/39

No Signal

VGA-1

Pirsa: 11060071 Page 384/39

Pirsa: 11060071 Page 385/39

Page 386/30

Pirsa: 11060071 Page 387/39

Pirsa: 11060071 Page 388/39

Pirsa: 11060071 Page 389/39

Pirsa: 11060071 Page 390/39

No Signal

VGA-1

Pirsa: 11060071 Page 391/39

Pirsa: 11060071 Page 392/39

Pirsa: 11060071 Page 393/39

