Title: Complexity in Fundamental Physics

Date: Jun 23, 2011 05:10 PM

URL: http://pirsa.org/11060070

Abstract:

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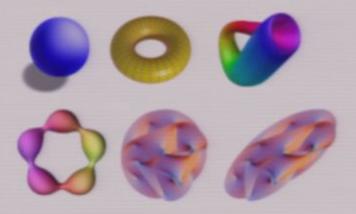
- Motivation
- 2 Inspiration
 - Spin glass models
 - State space structure of the SK model
 - Overlap order parameter
- 3 Application to string theory and cosmology
 - Generalized overlap order parameters
 - Examples of glassy systems in string theory + results
 - Cosmology

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Page 2/248

a: 11060070 Page 3/248

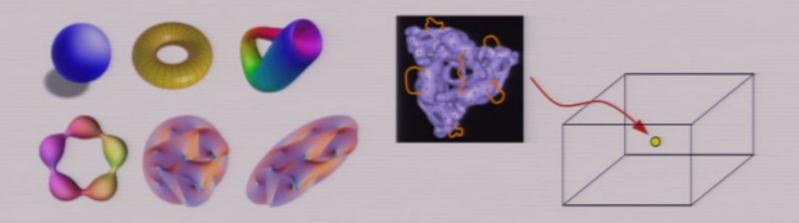
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Quantum gravity (string theory) seems to imply reversal of usual complexity - fundamentalness relation:

Landscape of compactifications

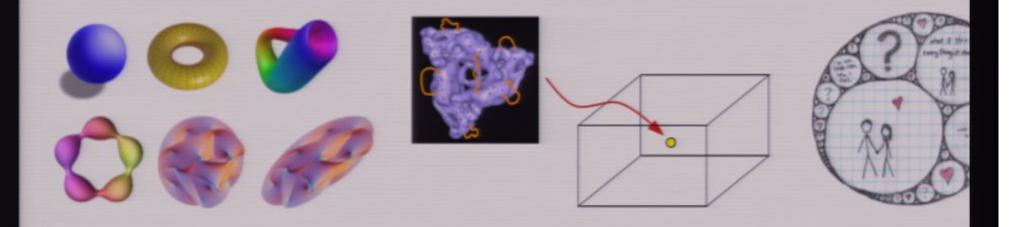
a: 11060070 Page 5/24



Quantum gravity (string theory) seems to imply reversal of usual complexity - fundamentalness relation:

- Landscape of compactifications
- Black hole microstates

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Quantum gravity (string theory) seems to imply reversal of usual complexity - fundamentalness relation:

- Landscape of compactifications
- Black hole microstates
- Eternal inflation

a: 11060070 Page 7/248

a: 11060070 Page 8/248

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- Seems unavoidable to understand positive cc in string theory:
 - All string constructions are of high complexity [KKLT,Silverstein,...]
 - Compactification data not a superselection sector [Coleman-de Luccia, Bousso-Polchinski,...]

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- Interesting mathematical structures in complex systems in general:
 - spin glasses: spontaneous hierarchical organization, overlap order parameters, replica symmetry breaking, ...
 - universality: relaxation times, ...
 - · computational complexity, P vs NP, ...

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- O Holographic understanding of (nonlocal) glassy systems
 - large N = thermodynamic limit
 - microscopic understanding of glasses considered to be one of most important unsolved problems in CM physics.

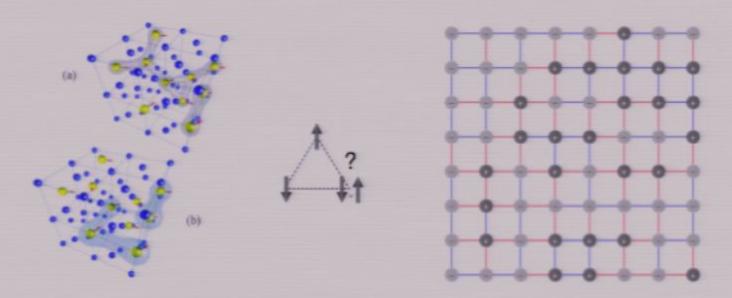
a: 11060070 Page 11/248

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- Approach to understand state space geometry quantum cosmology

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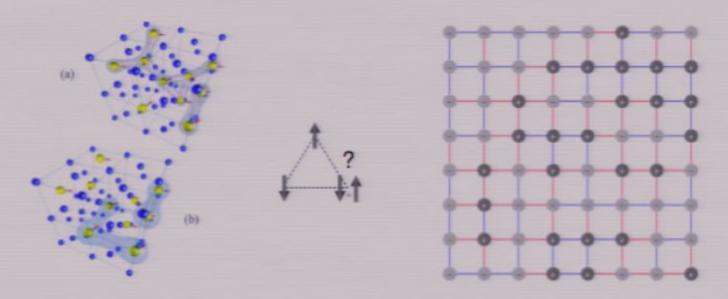
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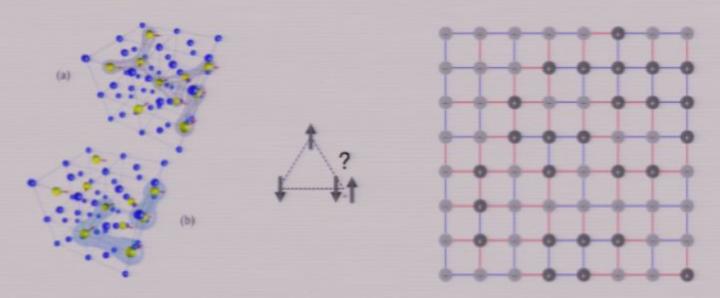
Real world example: Copper-Manganese alloy.

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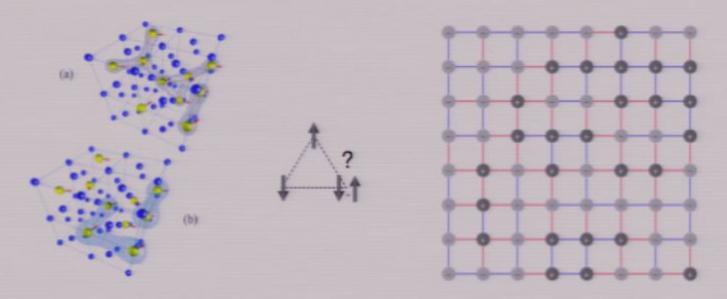
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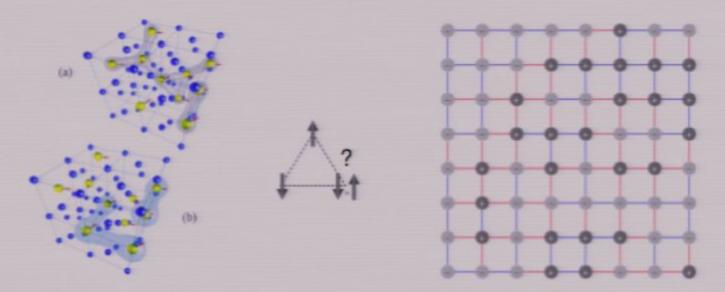
$$H = \sum_{ij} J_{ij} s_i s_j$$
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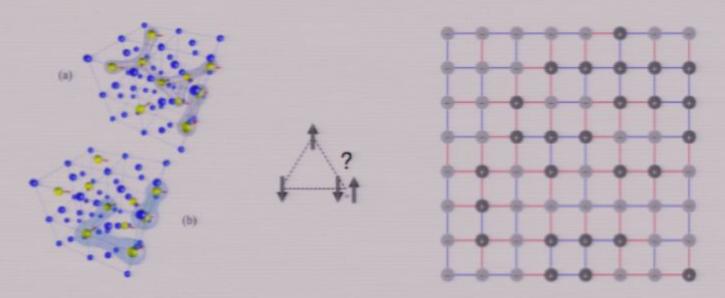
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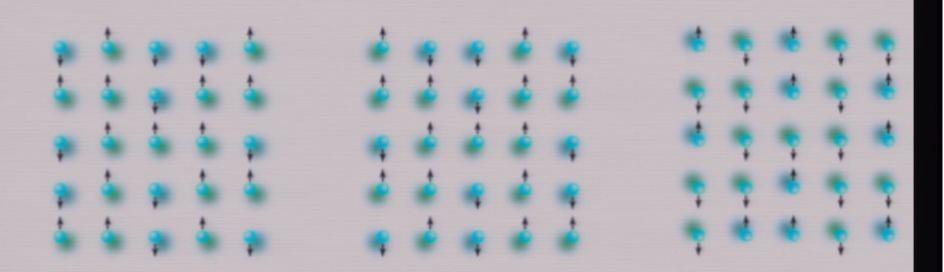


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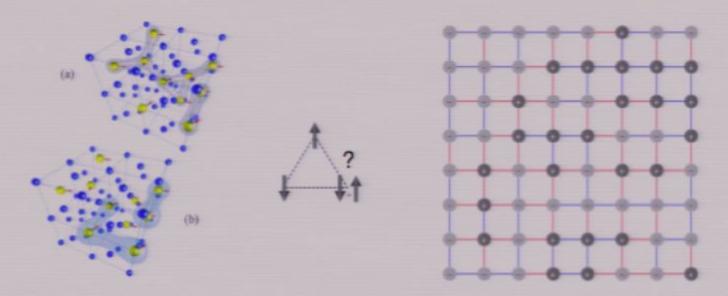
State space structure of the SK model



- Below critical temperature, spins freeze = spin glass phase.
- Different possible freezing patterns possible

 equilibrium states.
- Local magnetization in state α : $m_{i\alpha} \equiv \langle s_i \rangle_{\alpha}$ (depends on T).
- State overlap:

$$q_{\alpha\beta}\equiv \frac{1}{N}\sum_{i}m_{i\alpha}m_{i\beta}$$
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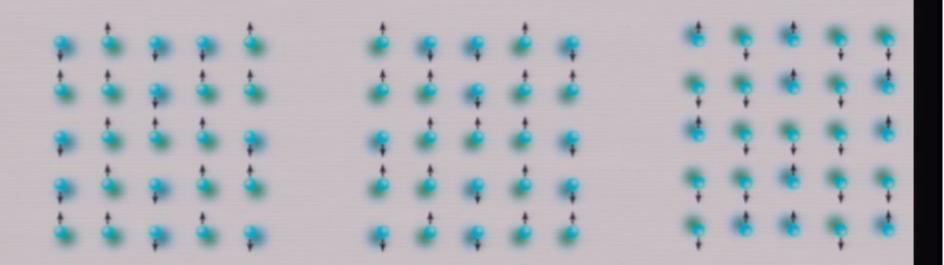


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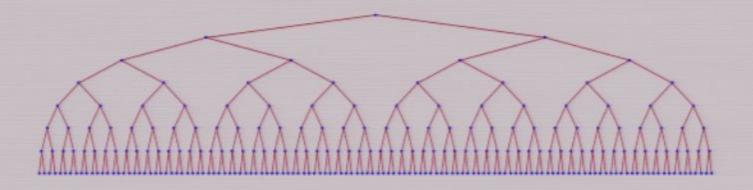
Results

• SK model exactly solvable [Parisi 79]

a: 11060070 Page 23/248

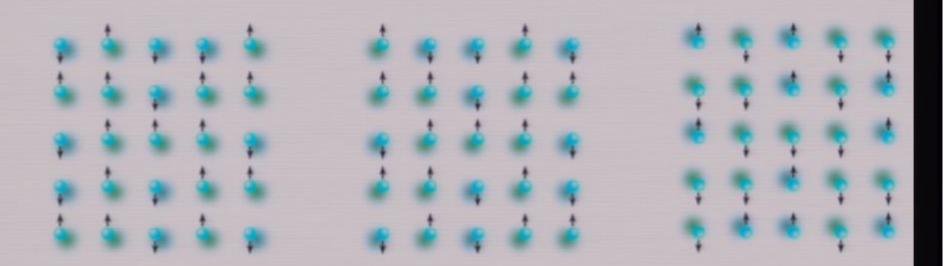
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- Key result: state space is ultrametric (distance $d_{\alpha\beta} \equiv q_{max} q_{\alpha\beta}$)
 - $\Leftrightarrow \forall \alpha, \beta, \gamma : d_{\alpha\beta} \leq \max\{d_{\alpha\gamma}, d_{\beta\gamma}\}.$
 - All triangles isosceles, with unequal side shortest (i.e. largest overlap
 - States organized as leaves of tree.



Analogous to evolution tree: distance = time to common ancestor, equivalently DNA overlap.

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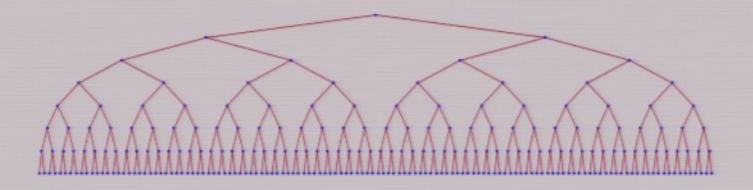
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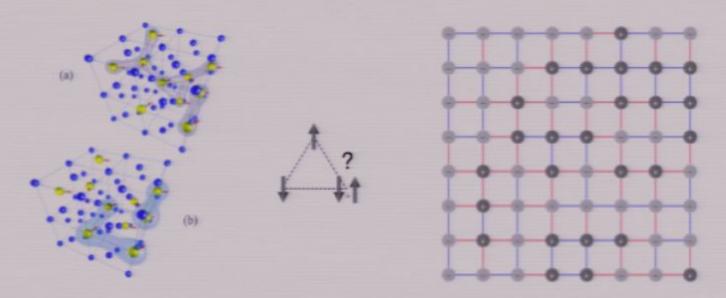
a: 11060070 Page 26/248

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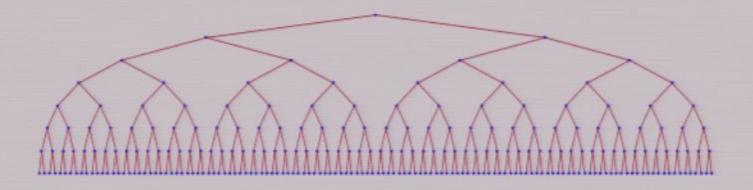
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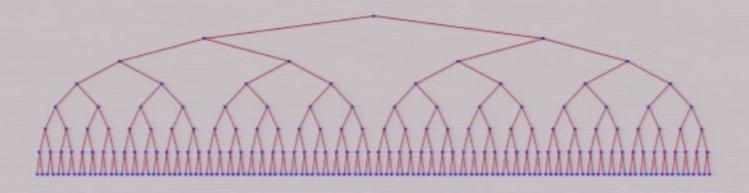
a: 11060070 Page 30/248

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Page 31/248

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Arises purely statically in SK model.

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a: 11060070 Page 33/24

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11060070 Page 34/248

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- E.g. Ising model at low T: $p_G(s) \equiv \frac{1}{Z}e^{-\beta H(\sigma)} = \frac{1}{2}p_+(s) + \frac{1}{2}p_-(s)$ where $p_{\pm}(s) = \lim_{\epsilon \to \mp 0} \lim_{N \to \infty} e^{-\epsilon \sum_i s_i} p_G(s)$.

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Page 35/248

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- Spin glass: $p_G(s) = \sum_{\alpha} w_{\alpha} p_{\alpha}(s) \rightsquigarrow p_{\alpha}(s) = ?$. Also: M = 0 for all T.

Page 30/246

• Definition equilibrium state ("pure state"): p_{α} such that cluster decomposition holds: $\lim_{N\to\infty}\frac{1}{N^2}\sum_{ij}|\langle\mathcal{O}_i\mathcal{O}_j\rangle_{\alpha}-\langle\mathcal{O}_i\rangle_{\alpha}\langle\mathcal{O}_j\rangle_{\alpha}|=0$

a: 11060070 Page 37/24

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$$P(q) = \sum_{s^{(1)}, s^{(2)}} p_G(s^{(1)}) p_G(s^{(2)}) \delta(q - \frac{1}{N} \sum_i s_i^{(1)} s_i^{(2)})$$

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Page 42/246

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Page 43/248

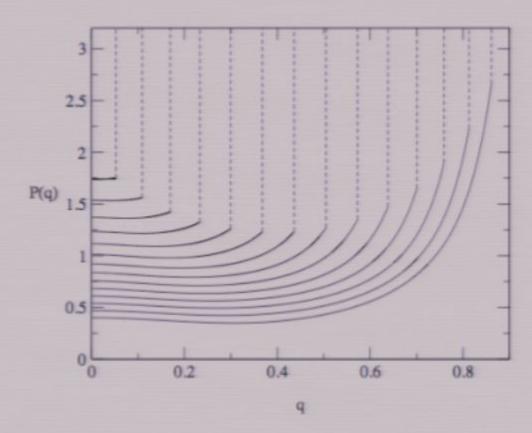
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- Average over disorder $\overline{P(q)} \equiv \int dJ_{ij} \, p(J_{ij}) \, P(q)$ can be computed analytically for SK model.
- a: 11060070 Multiple pure states $\Leftrightarrow P(q)$ nontrivial.

Result for SK model



P(q) for $T/T_c = 0.95, 0.9, ..., 0.3$ [Crisanti-Rizzo '02]

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Triangle distribution:

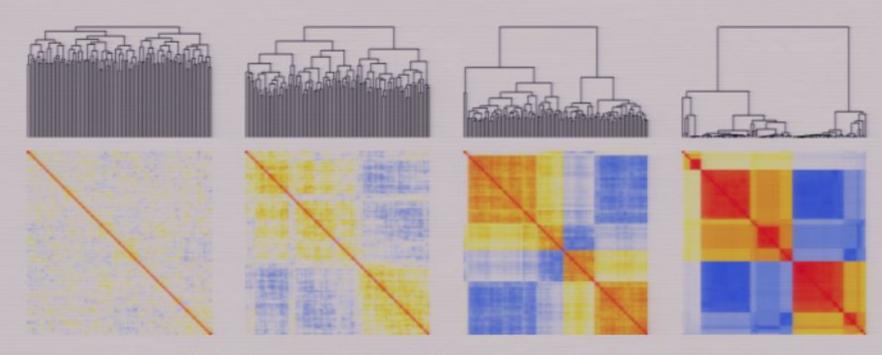
$$P(q_1, q_2, q_3) \equiv \sum_{\alpha\beta\gamma} w_{\alpha}w_{\beta}w_{\gamma}\,\delta(q_1 - q_{\beta\gamma})\delta(q_2 - q_{\gamma\alpha})\,\delta(q_3 - q_{\alpha\beta})$$

• Result for SK model, $q_i \ge 0$:

$$\overline{P(q_1, q_2, q_3)} = \frac{1}{2} \int_0^{q_1} dq \, \overline{P(q)} \, \overline{P(q_1)} \, \delta(q_1 - q_2) \, \delta(q_2 - q_3)$$

$$+ \frac{1}{2} \overline{P(q_1)} \, \overline{P(q_2)} \, \Theta(q_1 - q_2) \, \delta(q_2 - q_3) + \text{perm.}$$

Monte Carlo simulations



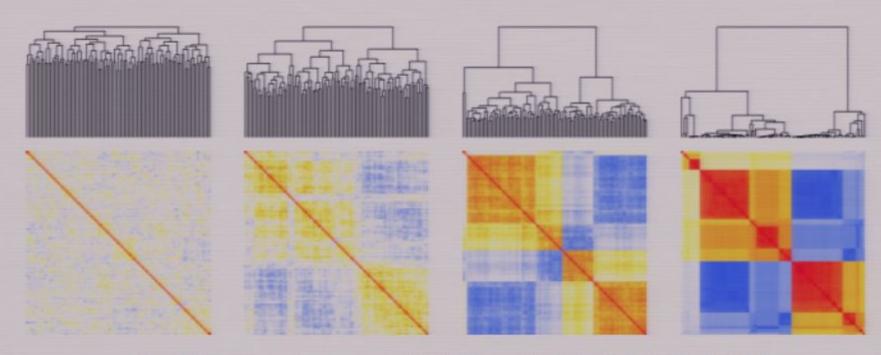
N=800 simulation at $T/T_c = 1.2, 0.86, 0.55, 0.12$

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Generalization P(q) from classical Ising spins to general, nonlinear spac and quantum systems?

Several possibilities

a: 11060070 Page 51/24

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Several possibilities

- whatever comes up naturally when integrating out quenched disord
- overlap operators that violate cluster decomposition
- reduced density matrices

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General definition:

- ϕ_i , $i=1,\cdots,N$ subset of d.o.f. of system.
- $\rho = \sum w_{\alpha} \rho_{\alpha}$, ρ_{α} s.t. cluster decomposition holds.
- ρ_{α}^{i} be reduced density matrix for ϕ_{i} d.o.f.
- Overlap $q_{lphaeta}\equiv rac{1}{N}\sum_{i}\operatorname{Tr}
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Replica formula: $P(q) = \langle \delta(q - \frac{1}{N} \sum_i \chi_i) \rangle_{n=2}$, where χ_i is the operator

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String glasses

(Partial) results for:

Page 57/248 a: 11060070

(Partial) results for:

- Bousso-Polchinski with quenched random metric on flux lattice
 - analytic: replica symmetric solution
 - numeric: exchange monte carlo

a: 11060070 Page 58/248

(Partial) results for:

- Bousso-Polchinski with quenched random metric on flux lattice
 - analytic: replica symmetric solution
 - numeric: exchange monte carlo
- Quiver models (for wrapped D-branes) with quenched random superpotentials

1 age 39/240

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String glasses

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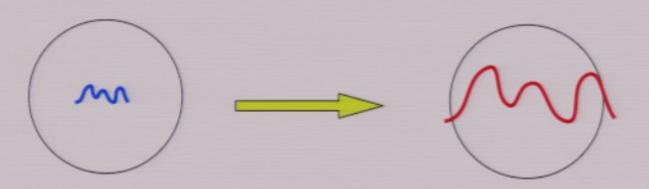
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Cosmology



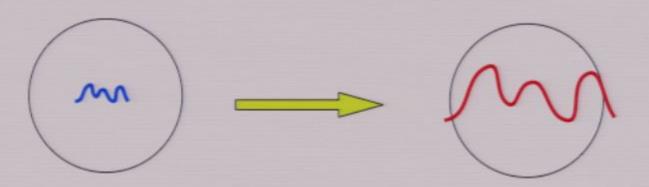
- Quantum fluctuations get exponentially stretched to super-Hubble scales and freeze, becoming effectively classical.
- This dynamically generates analog of spin glass "pure states", in which scalars, metric,... have definite values on large scales and cluster decomposition holds. (Not the case for HH/BD/Euclidean vacuum.)
- Overlap distribution → state space analysis without explicit description of clustering states.
- Useful quantity to compare quantum wave function of universe and "stochastic" measure approaches.

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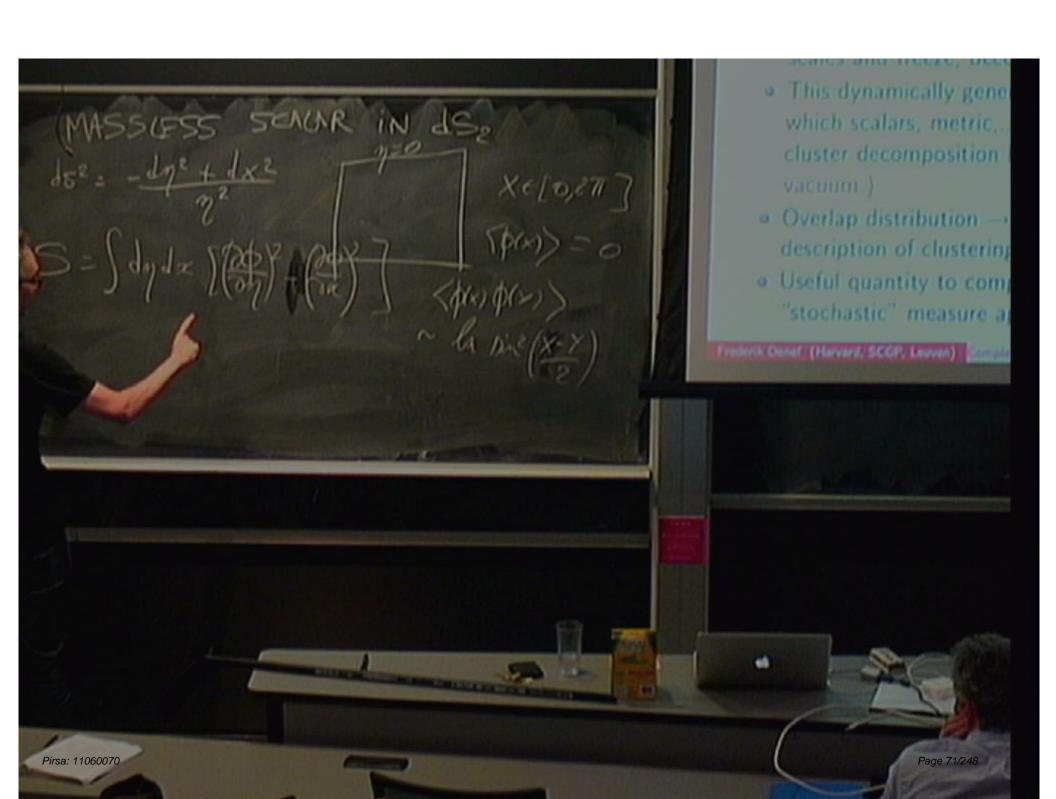
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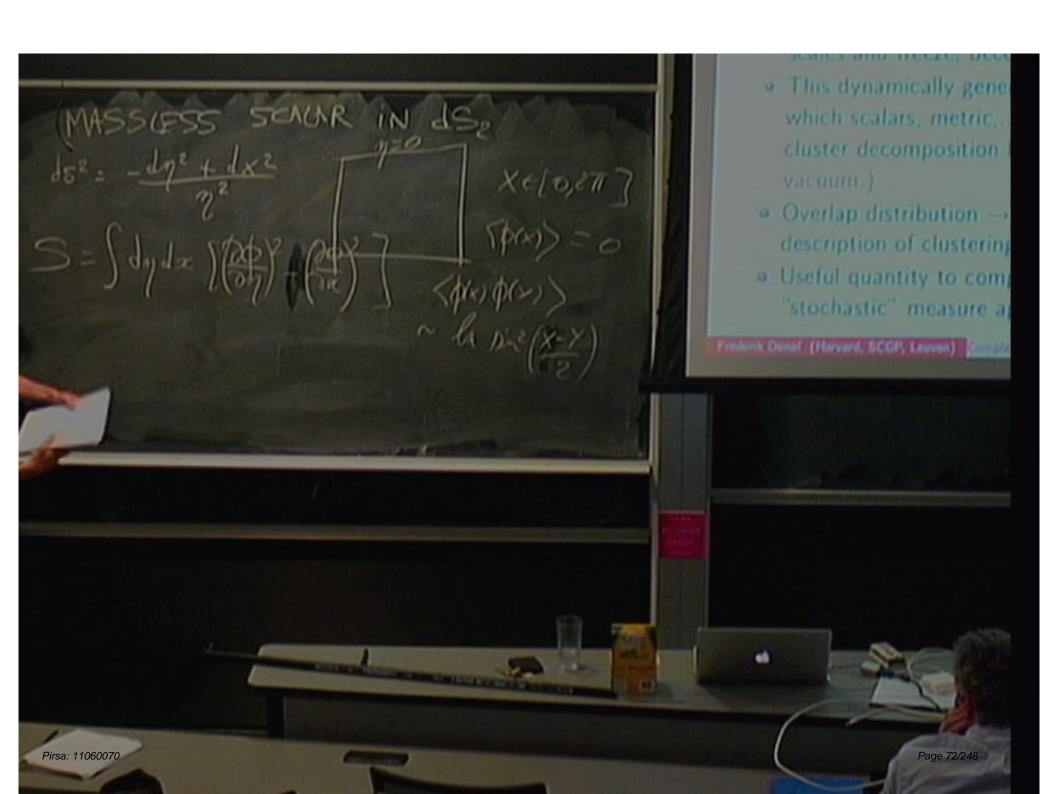
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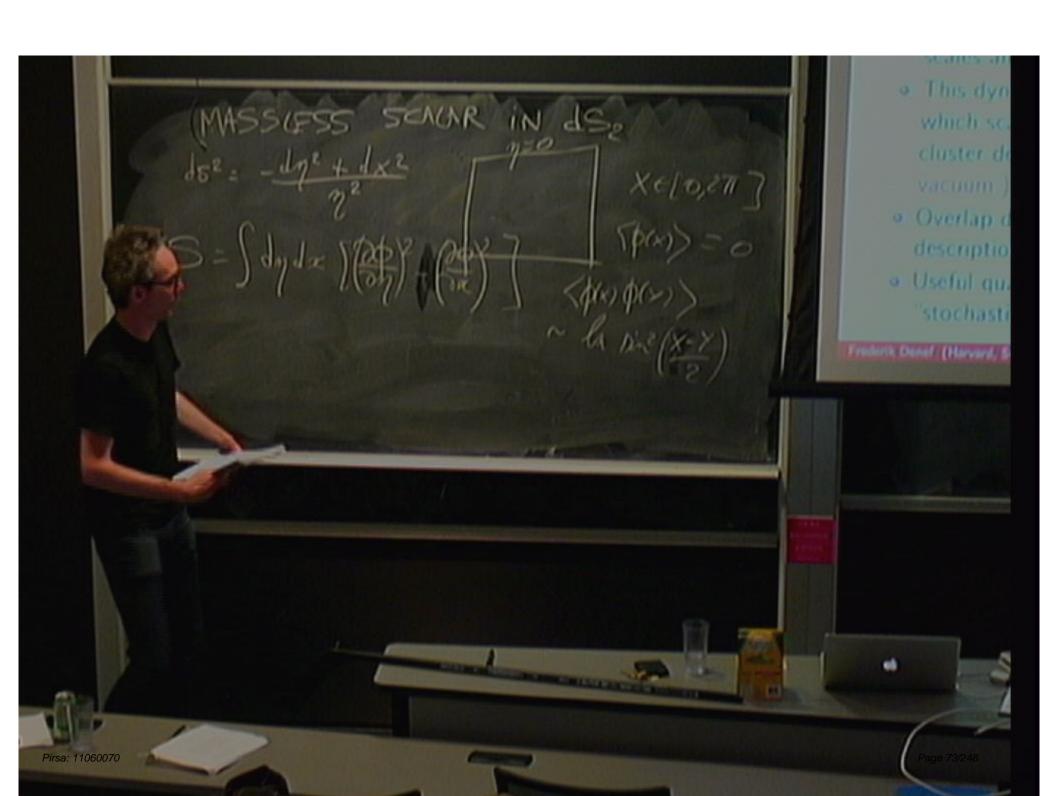
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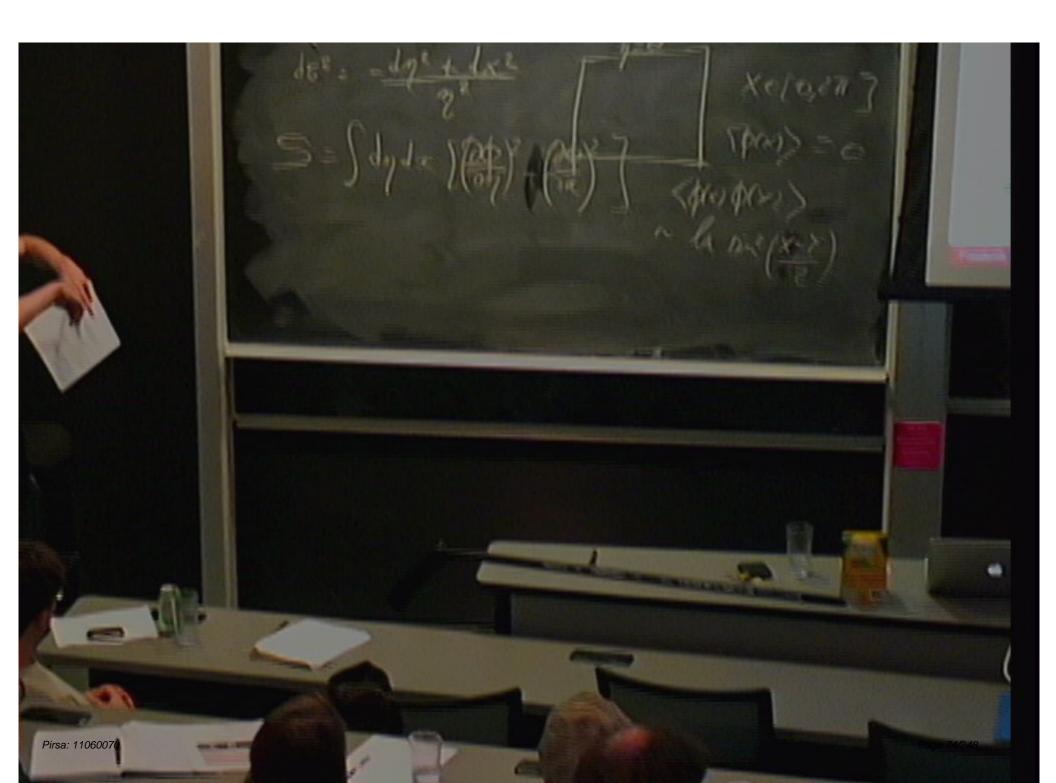


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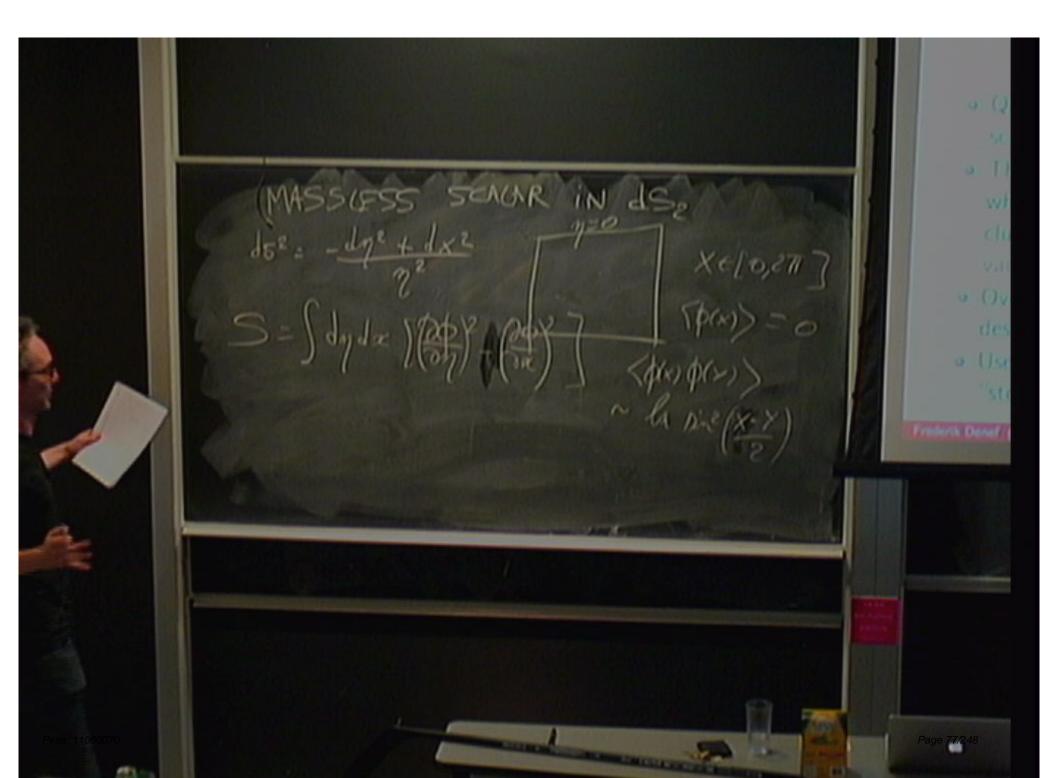






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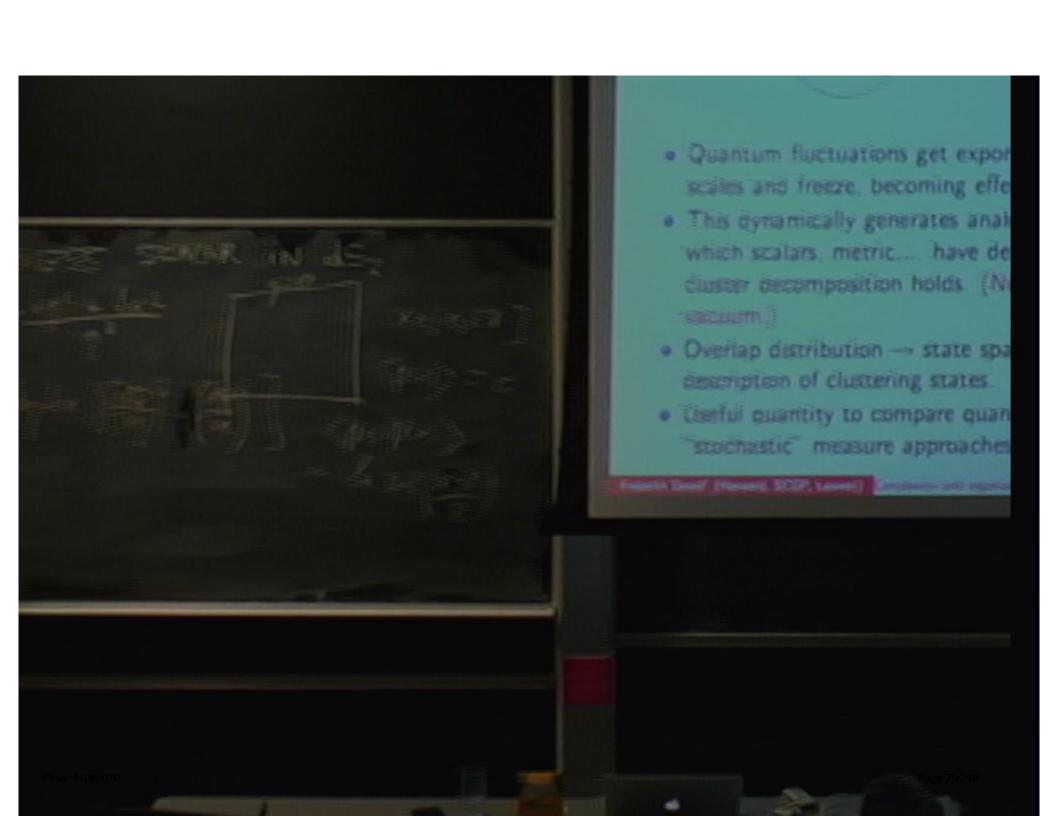


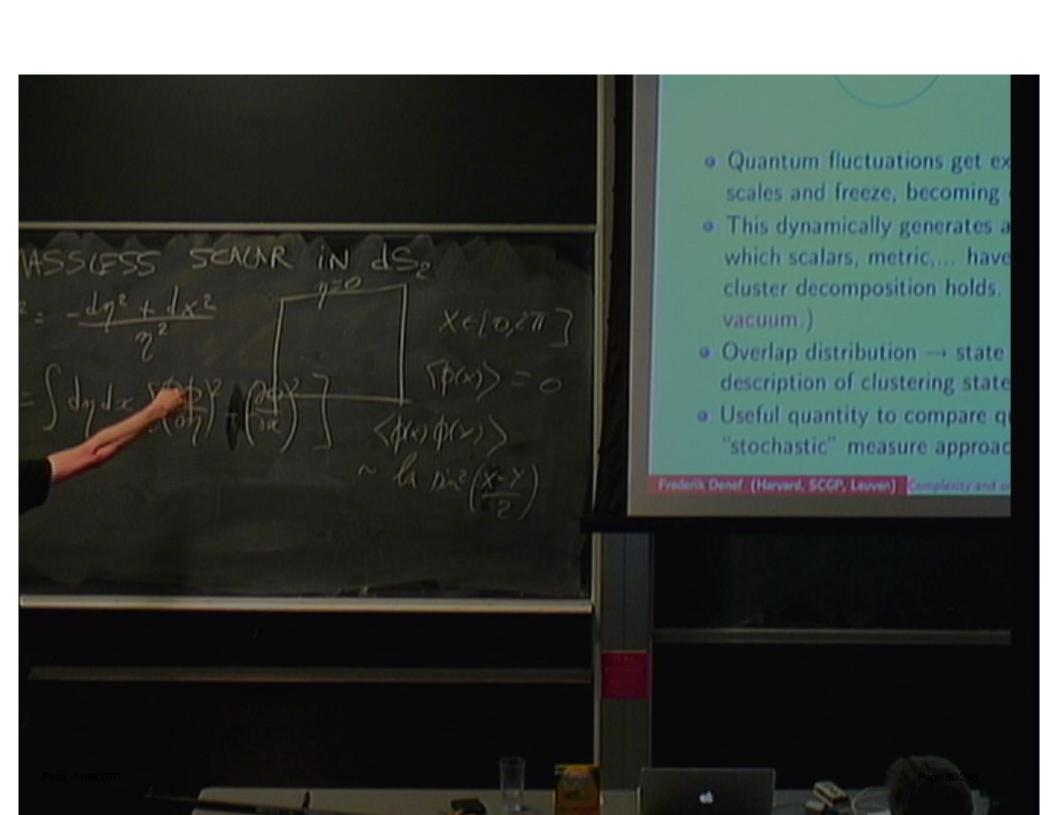
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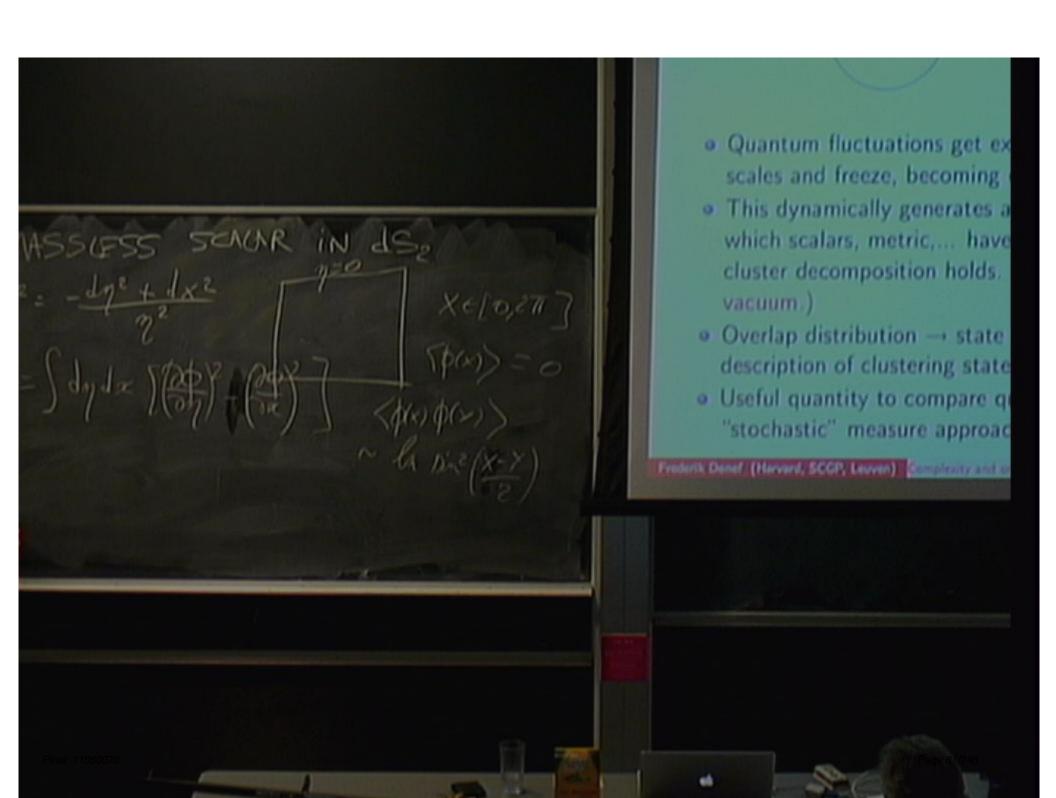
Frederik Desef (Harvard, SCGP, Leaven) Complexity and organization in many tesses.

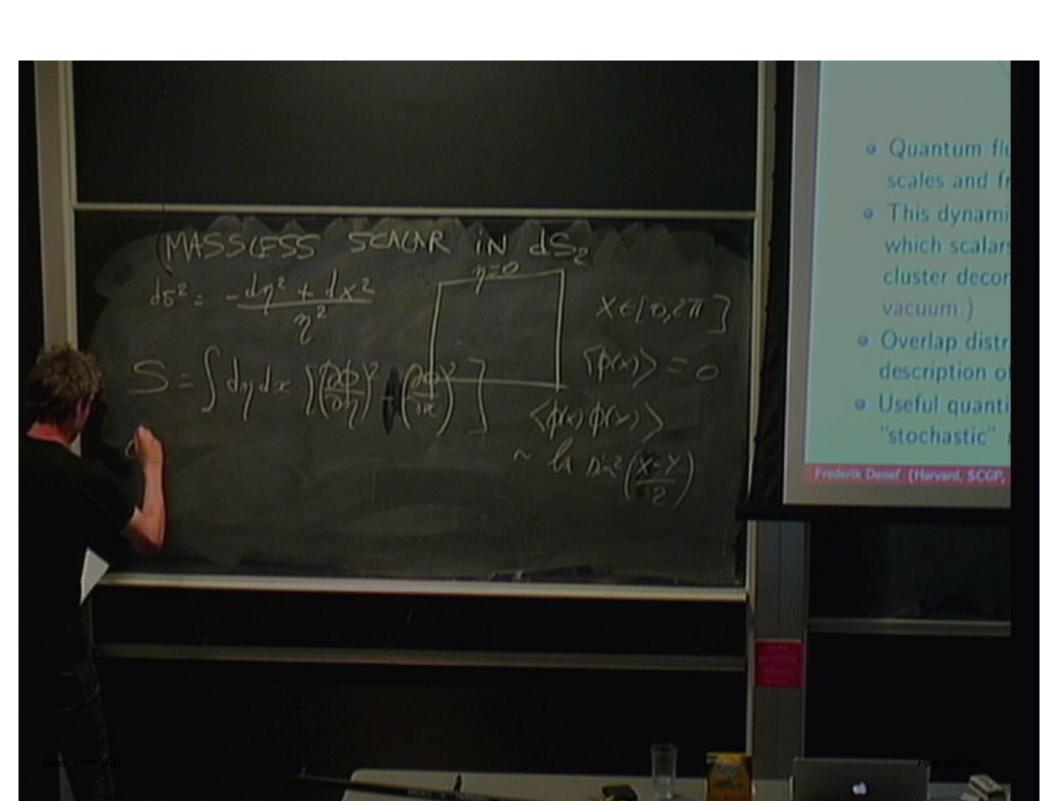
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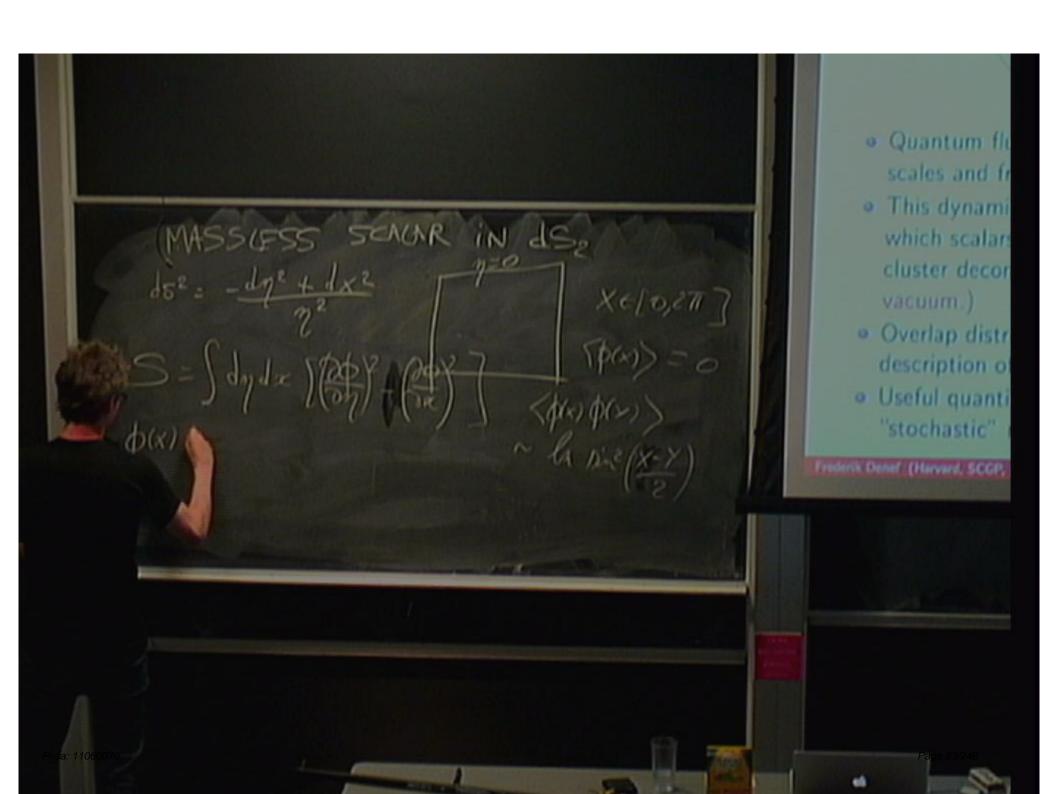
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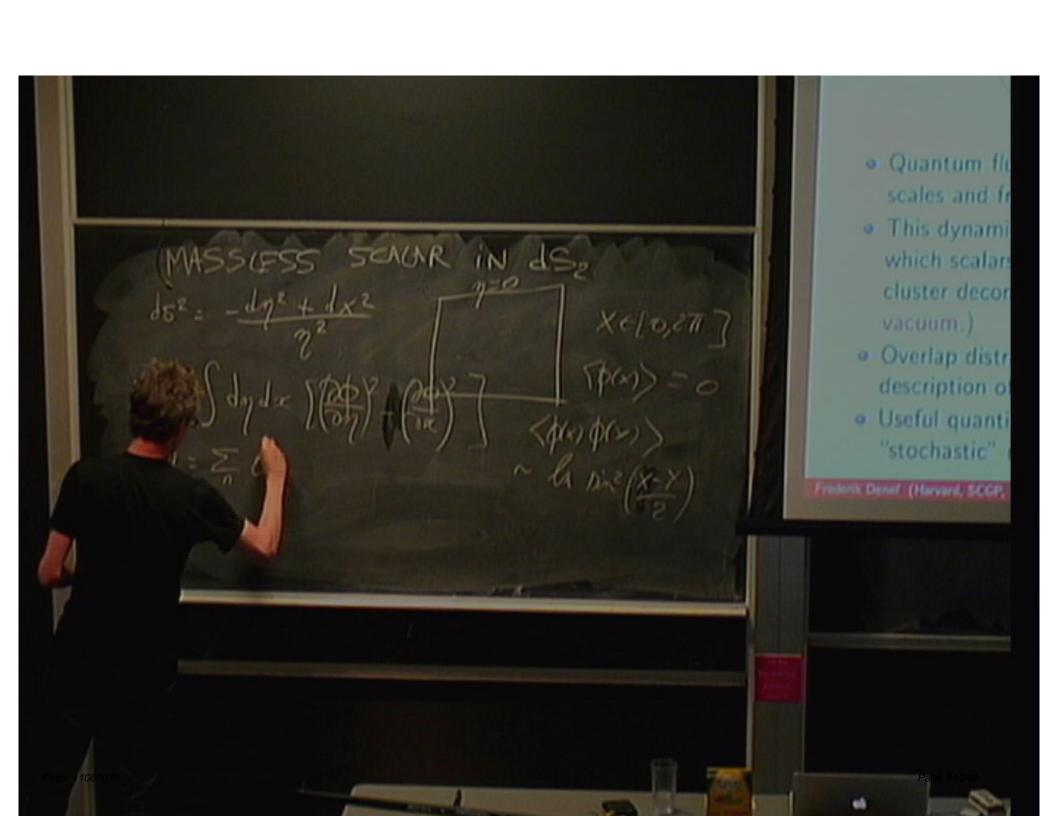


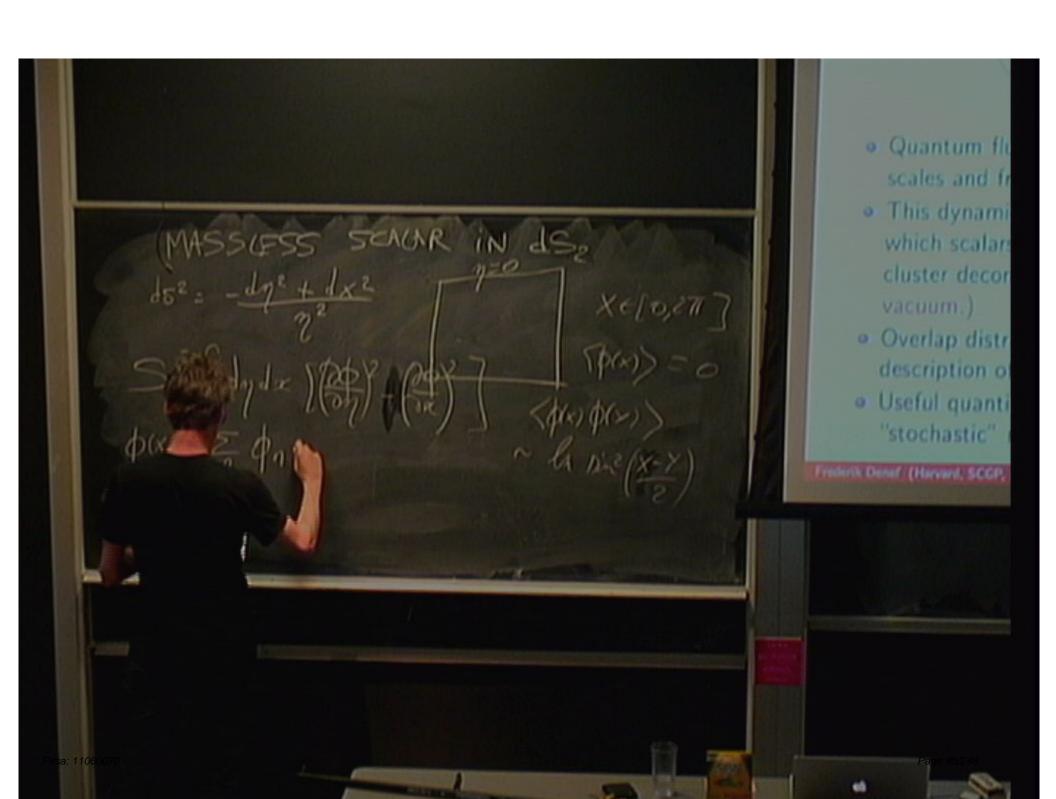


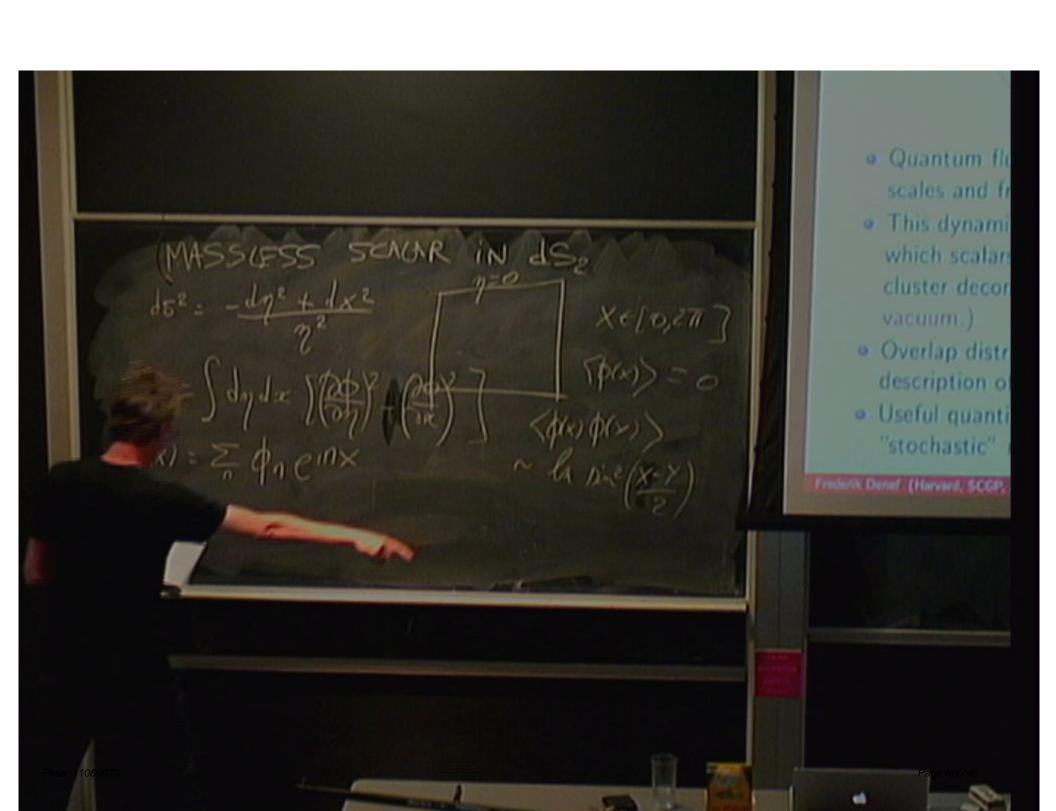


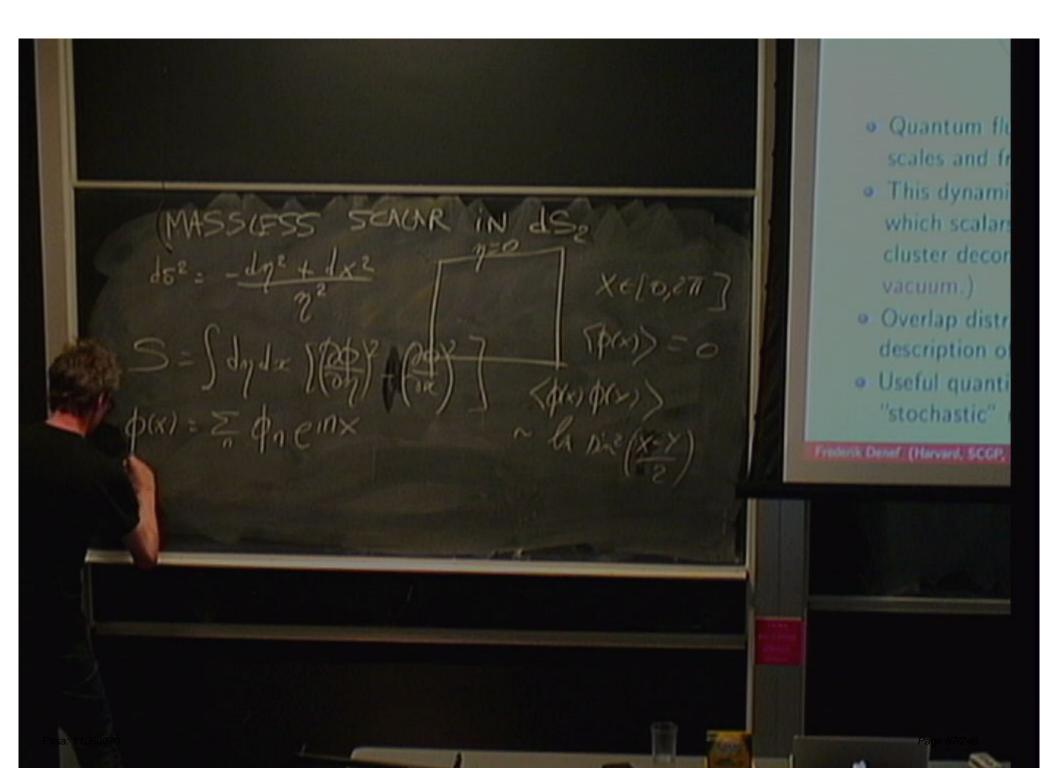


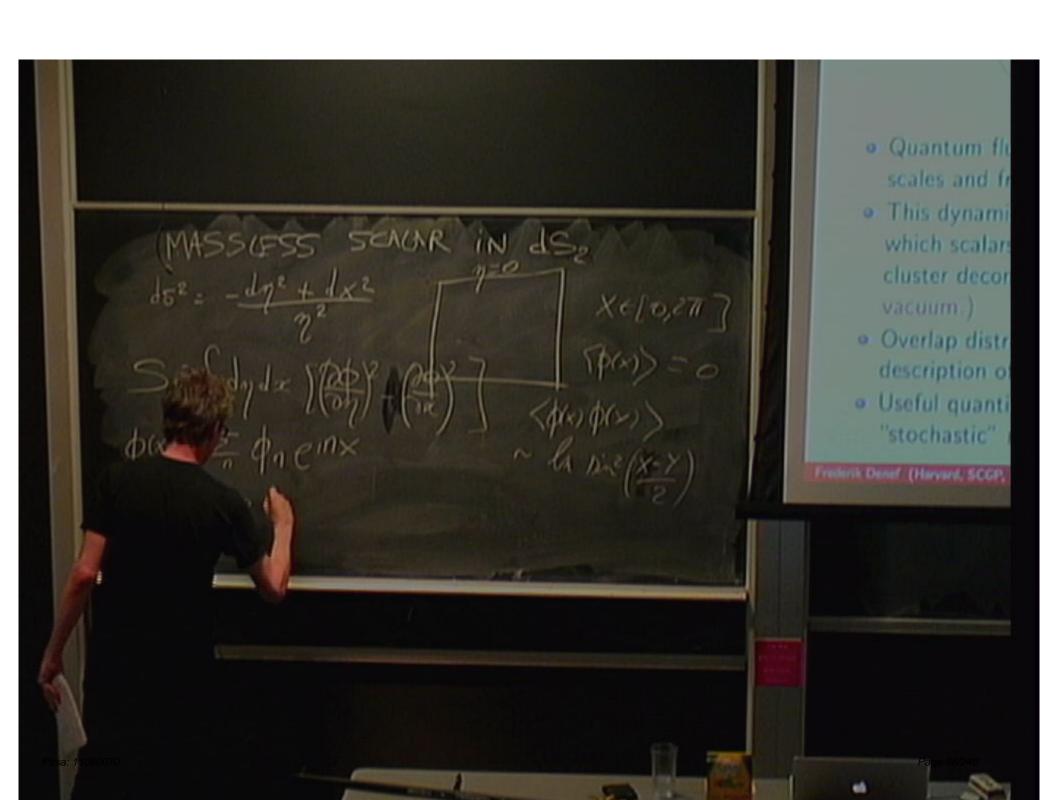


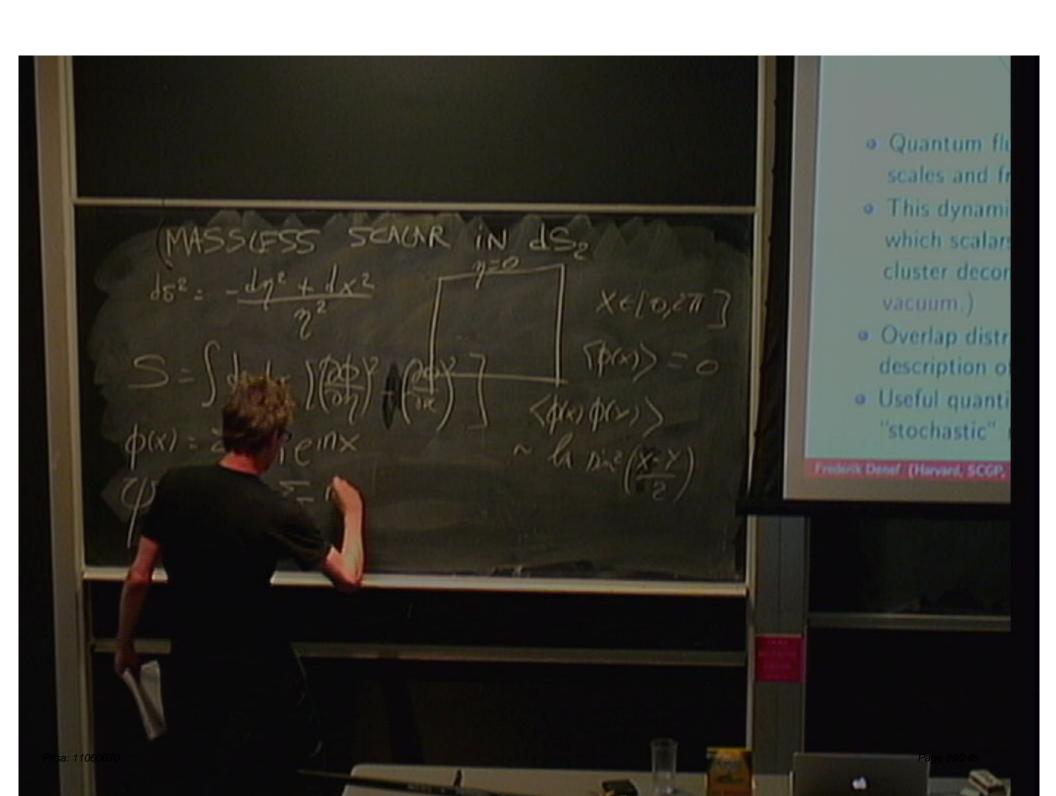


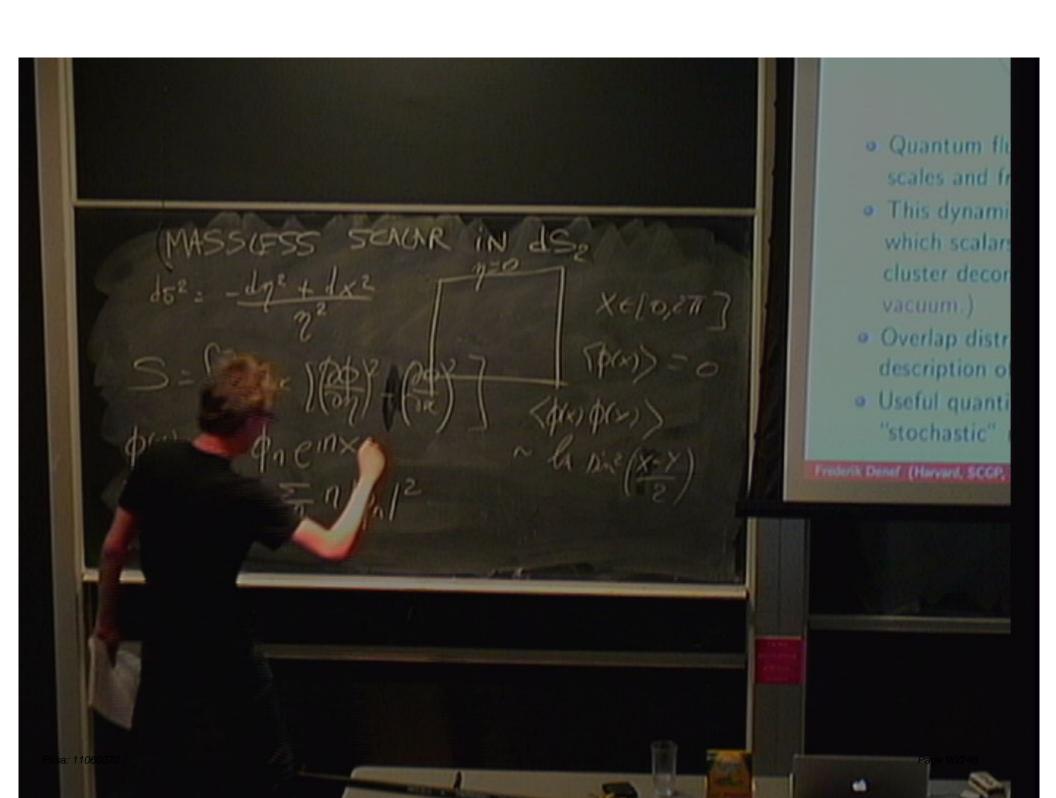


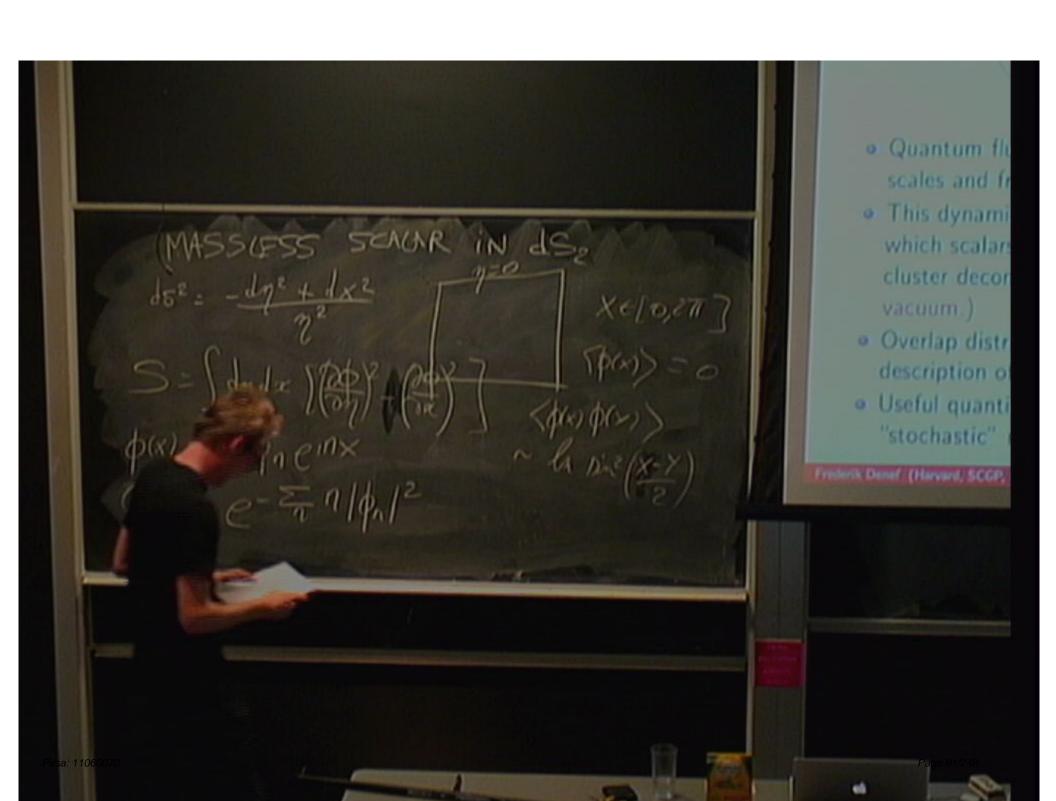


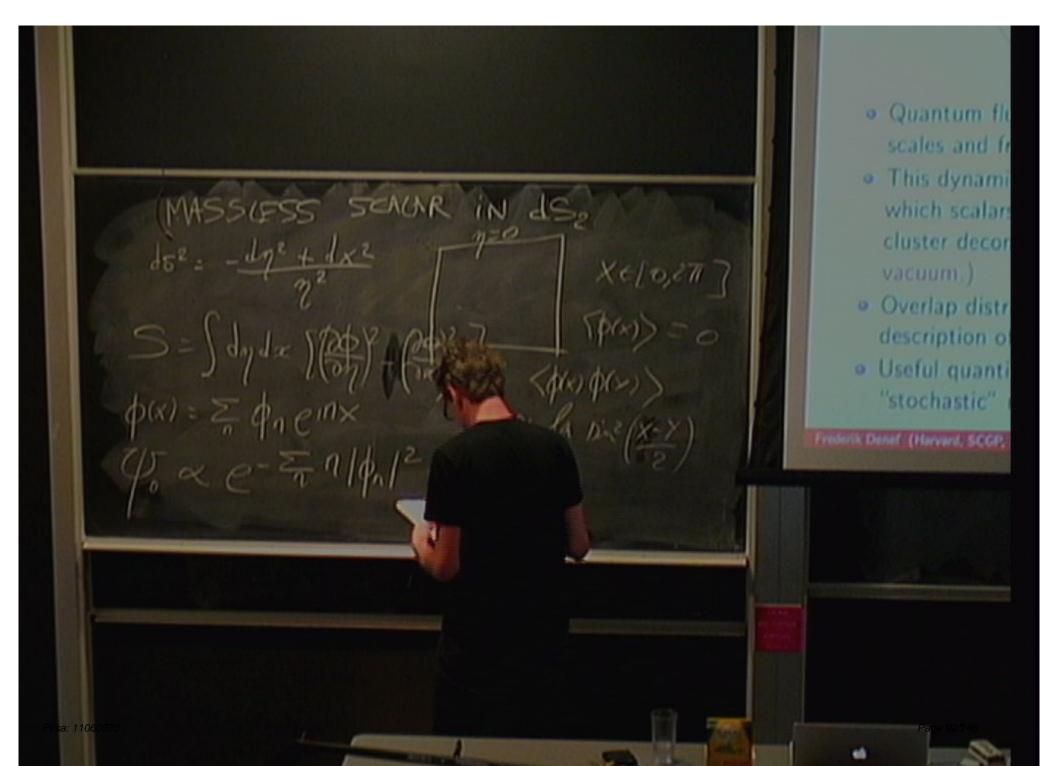


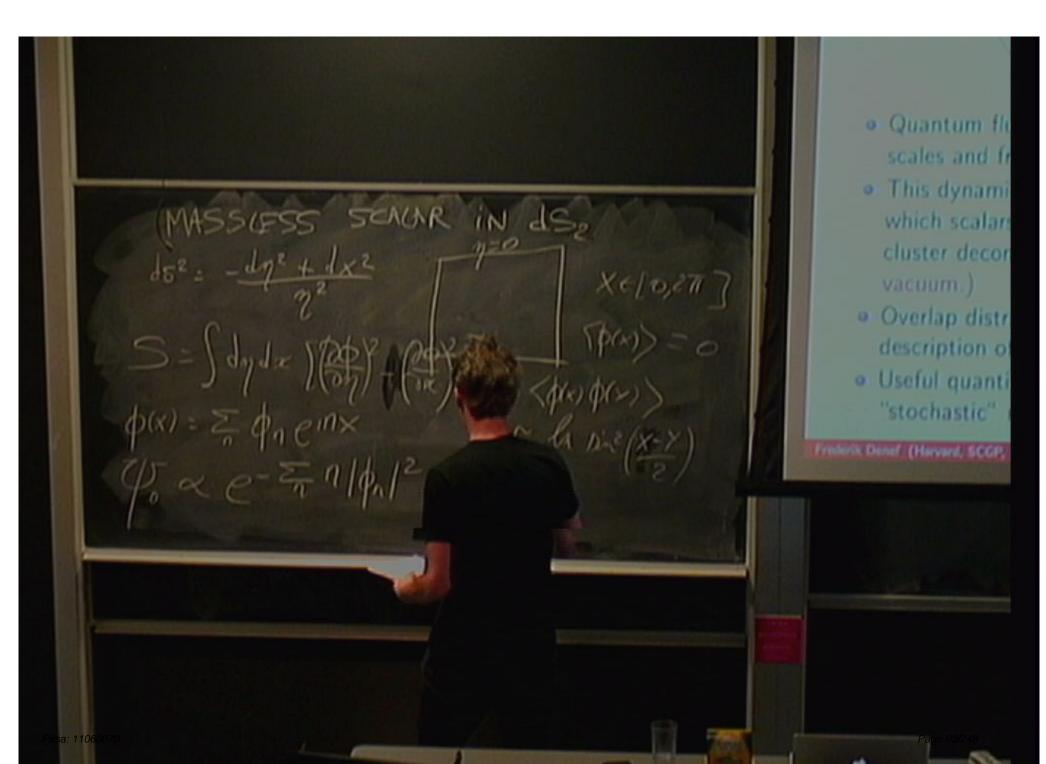


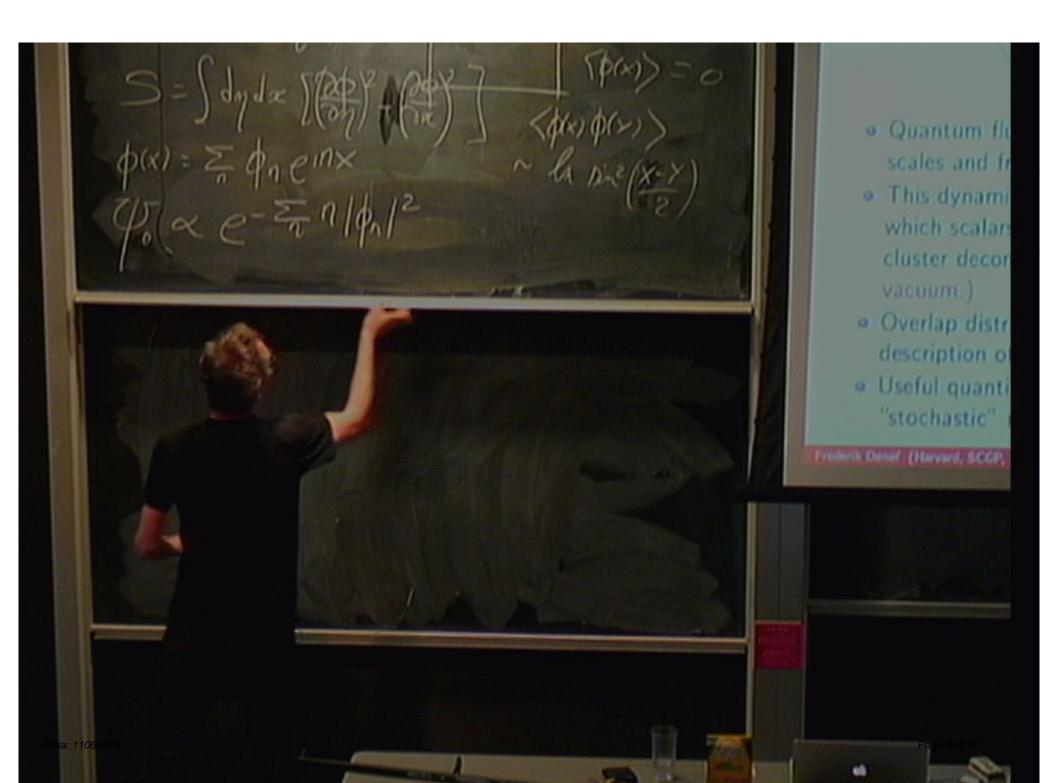


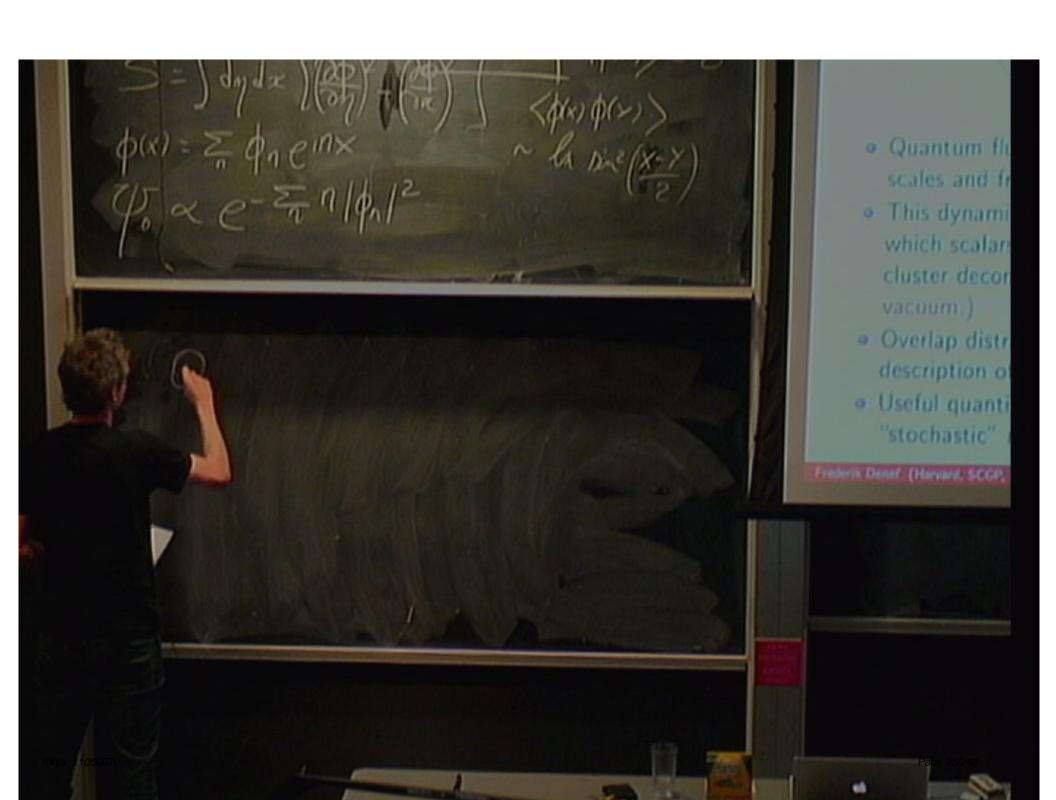


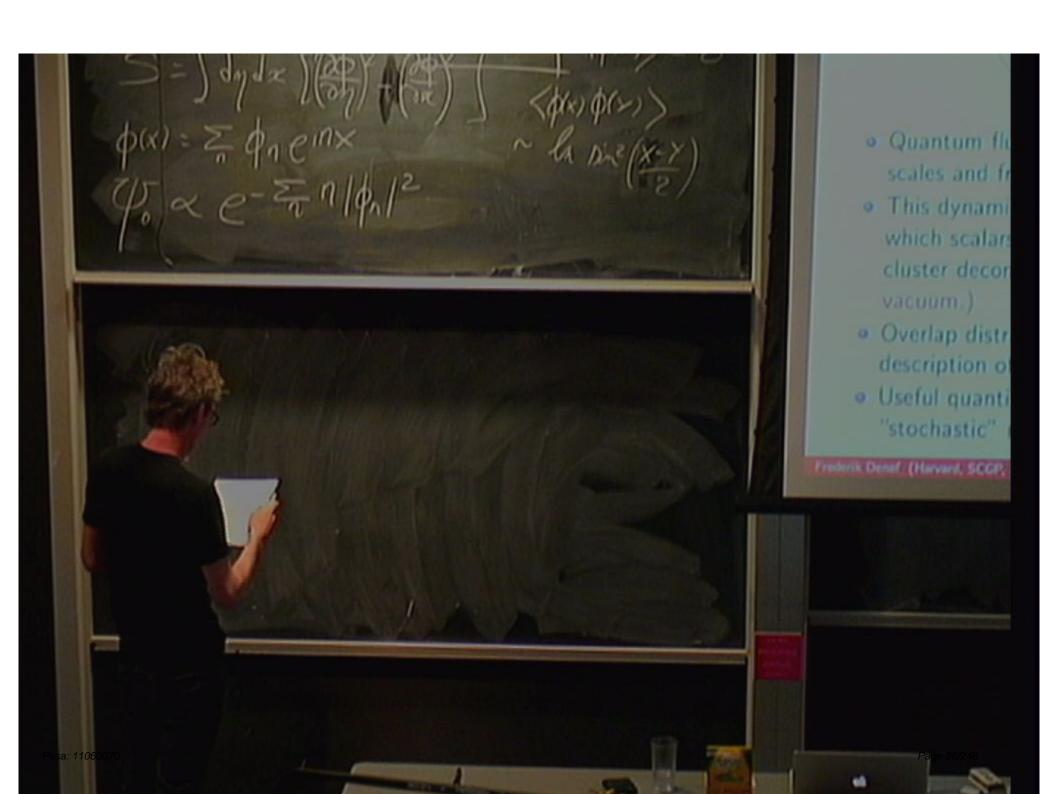


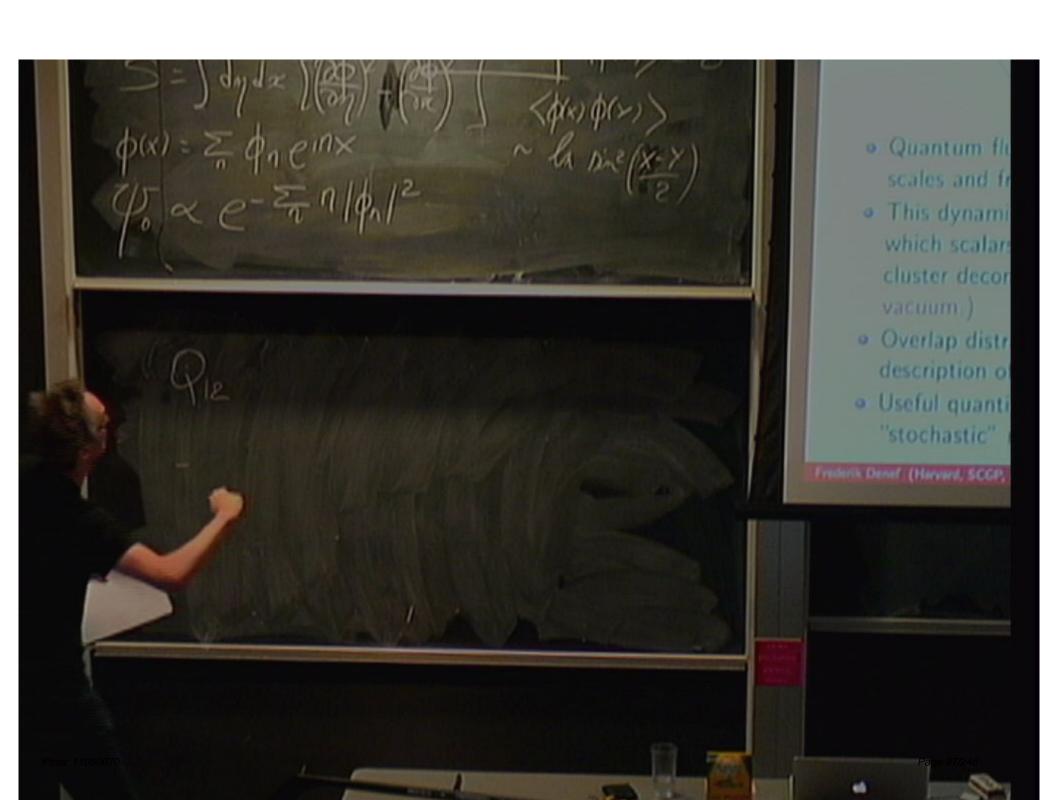


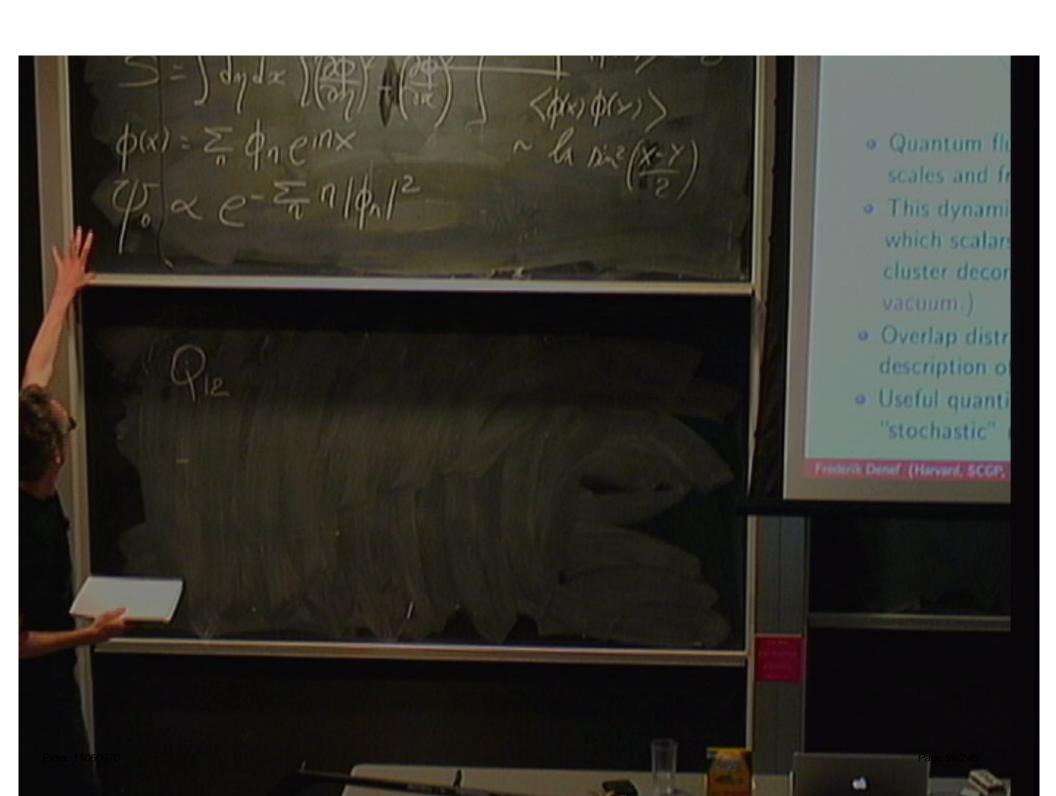


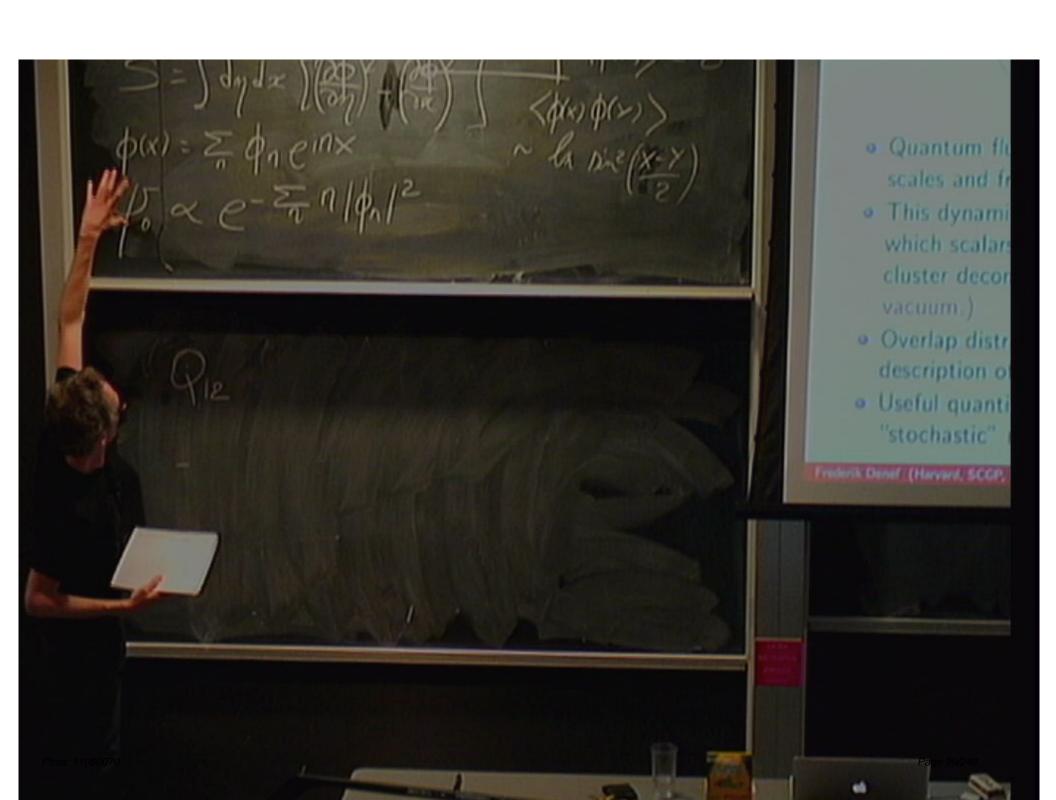


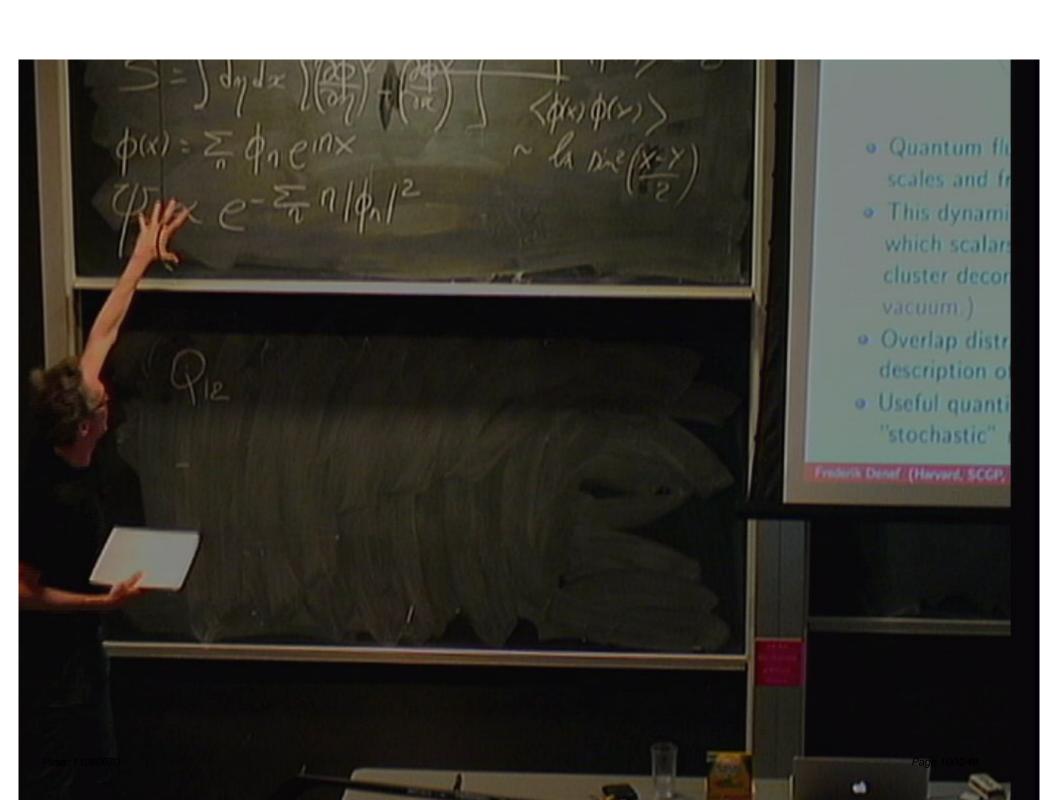


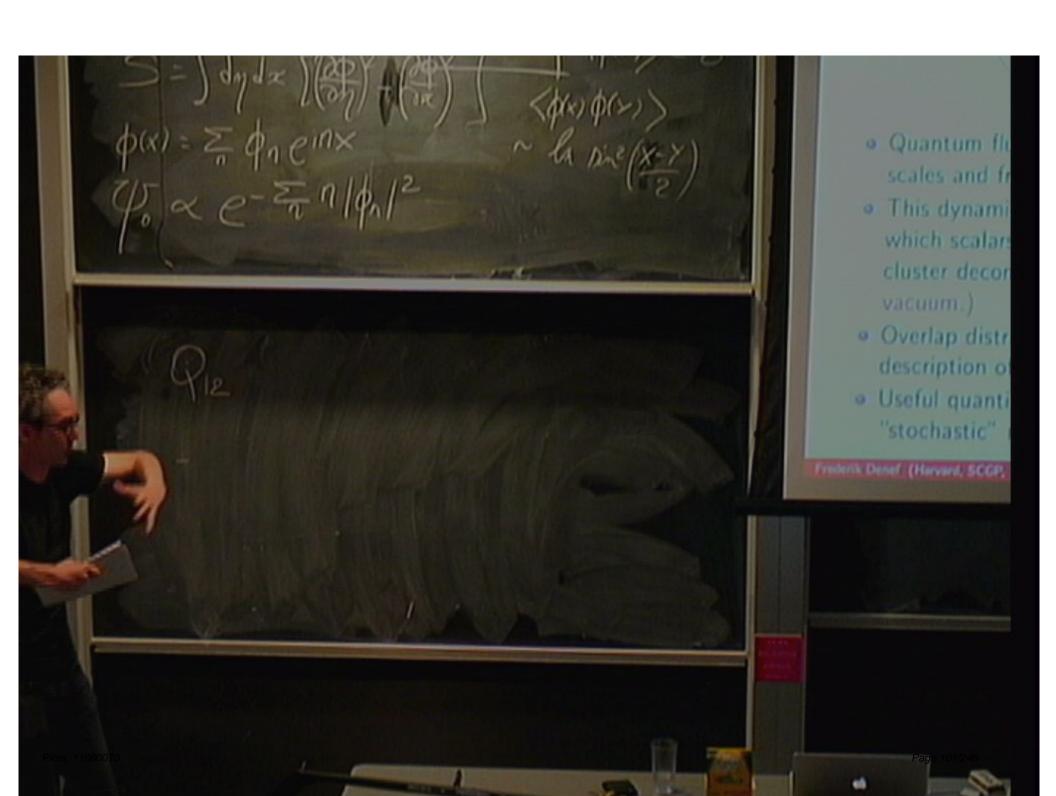


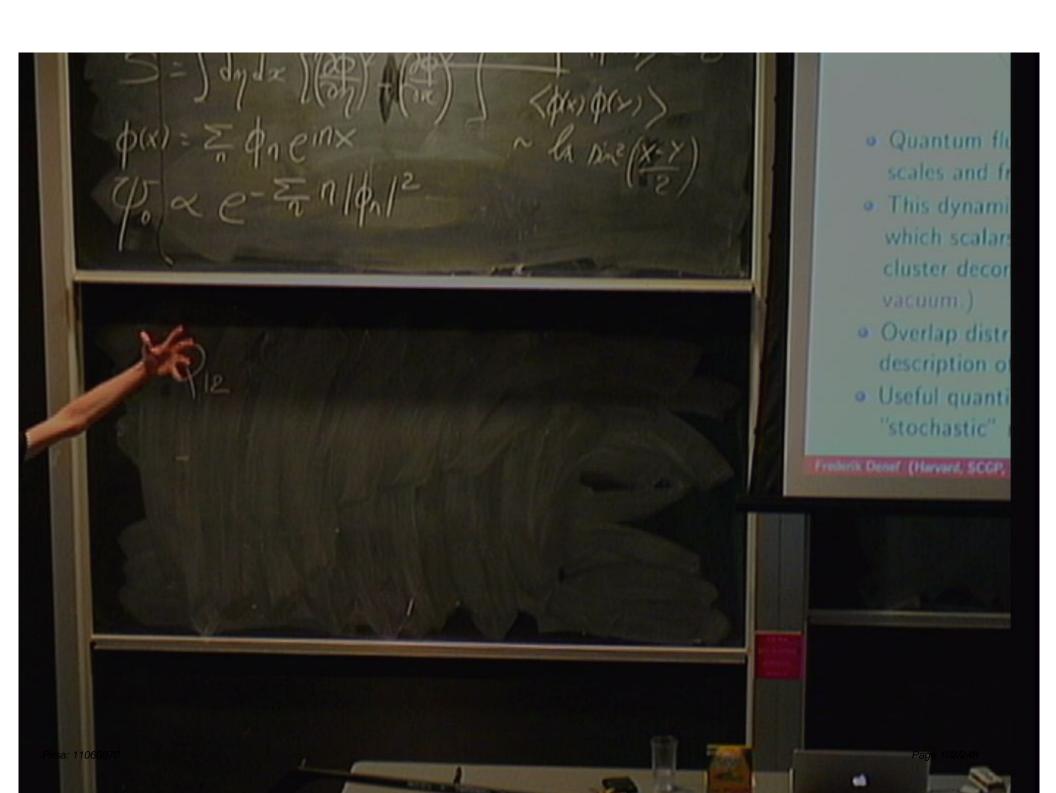


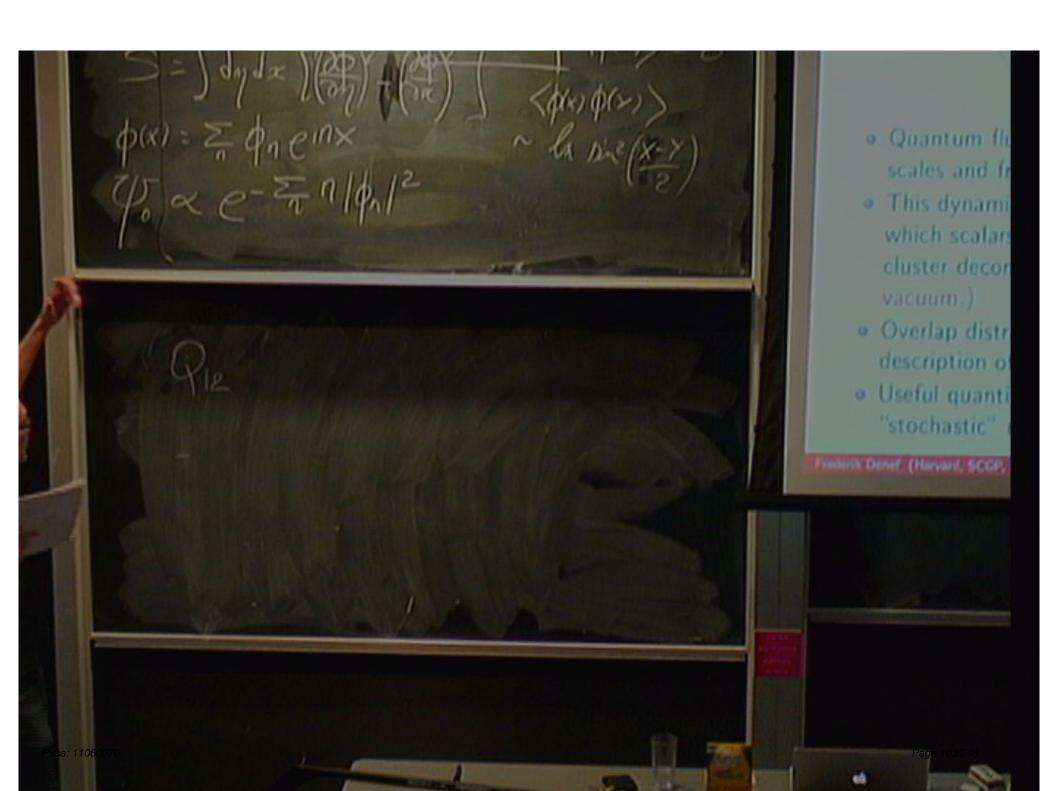


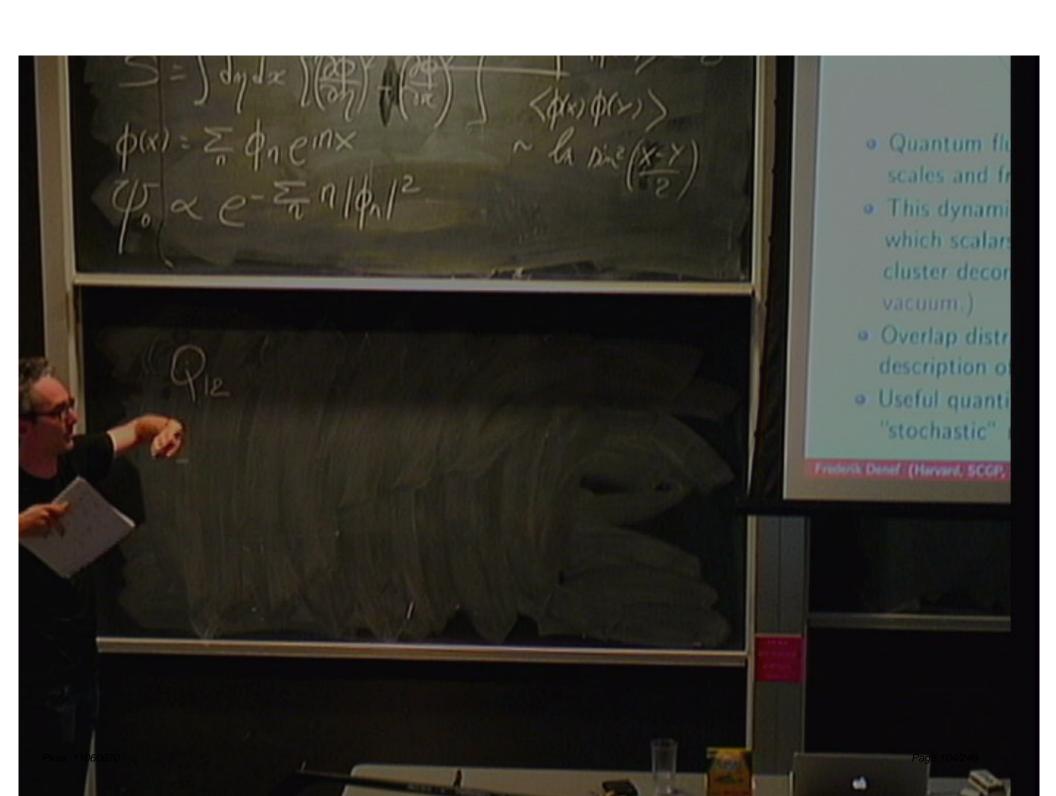


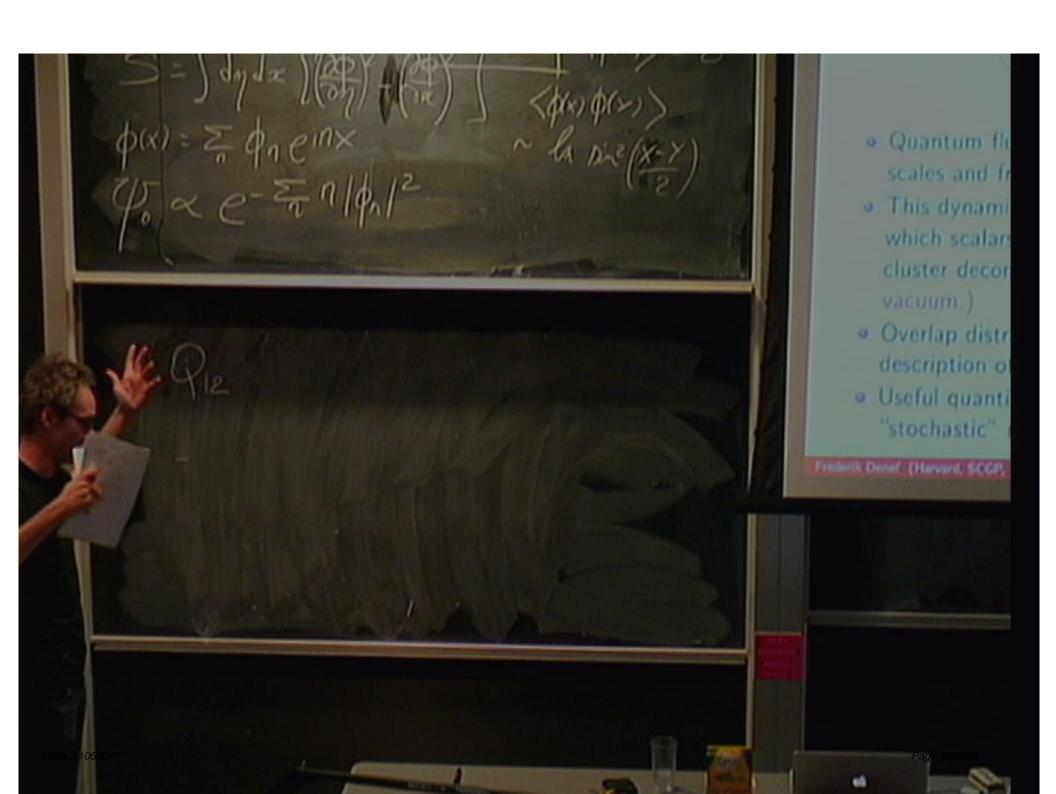


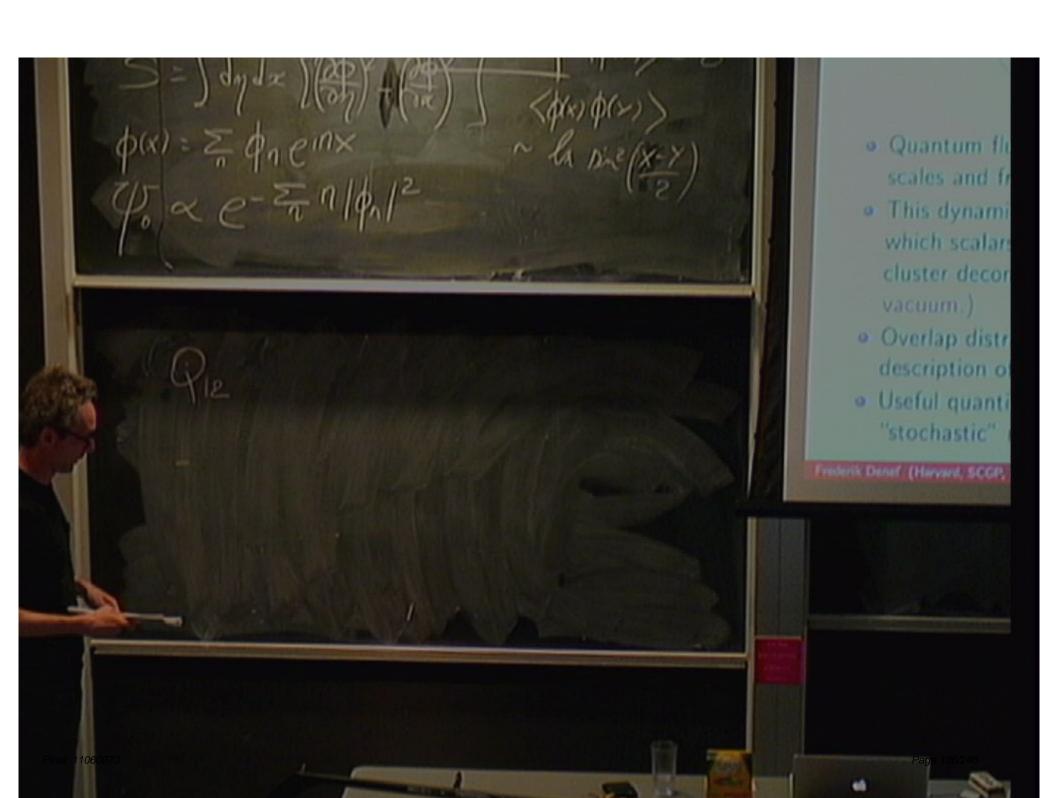


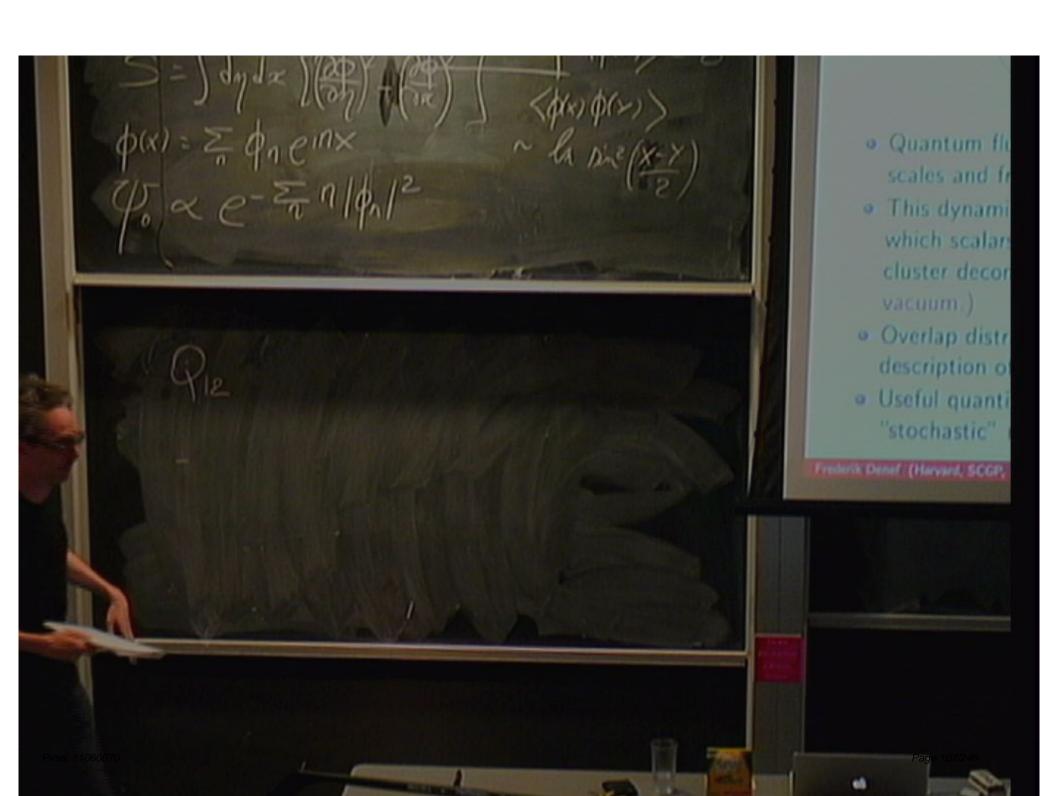


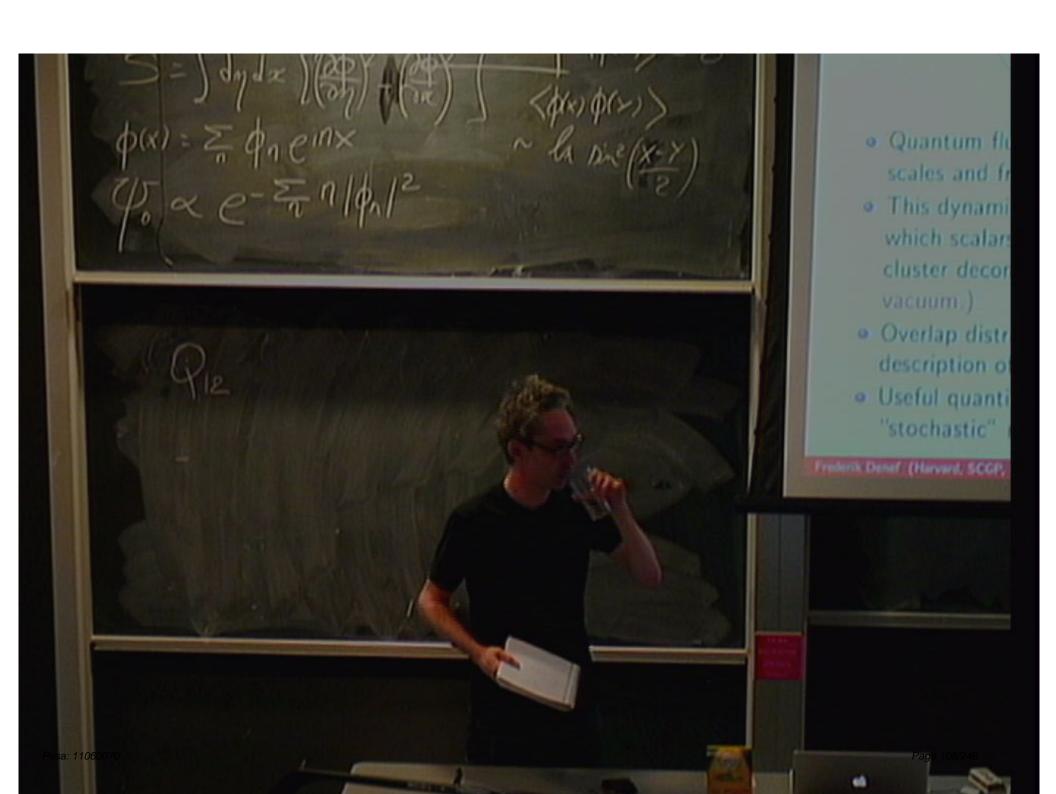


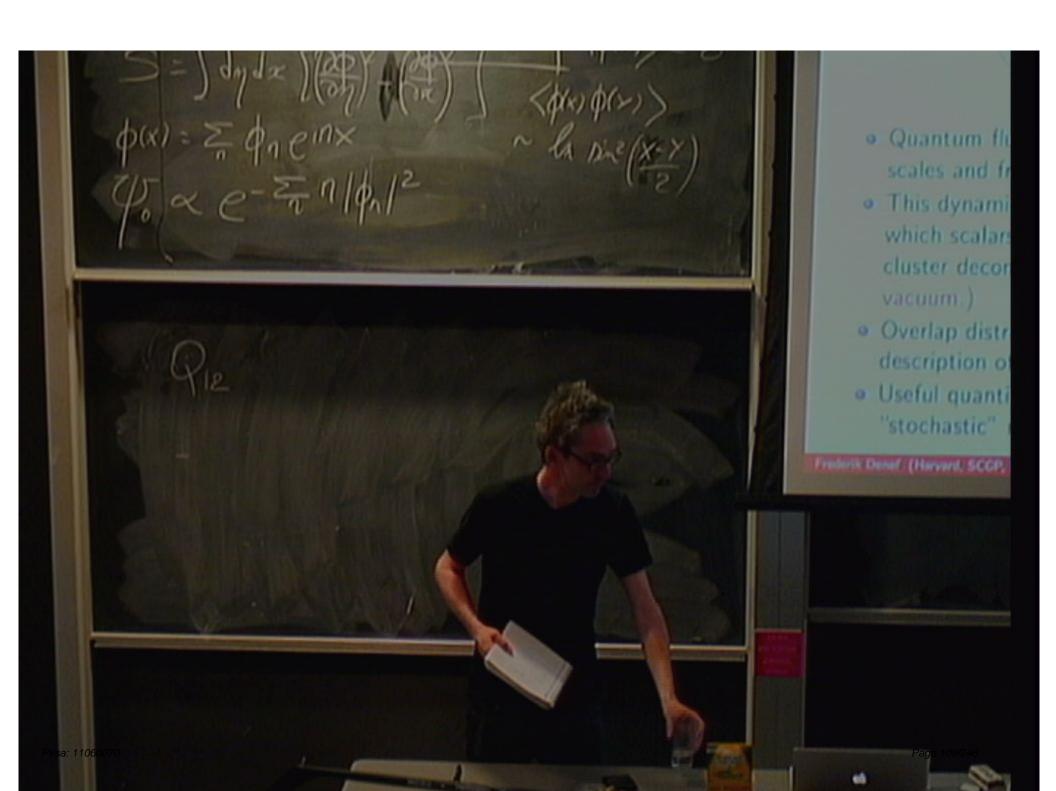


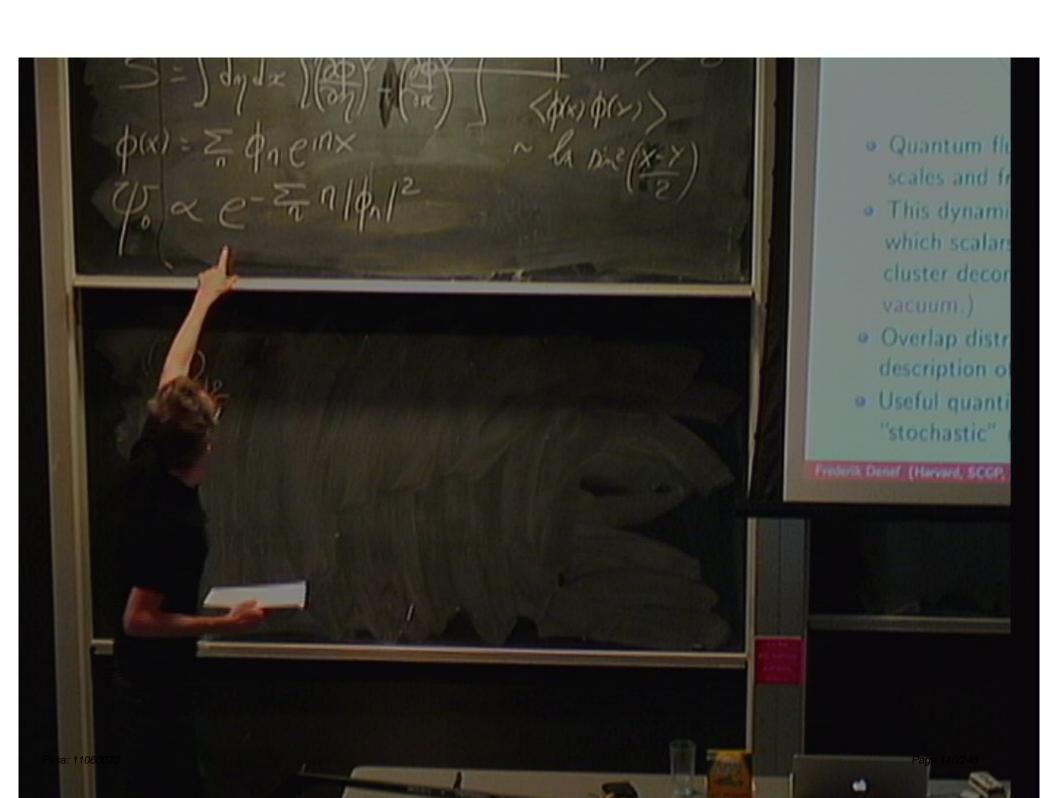


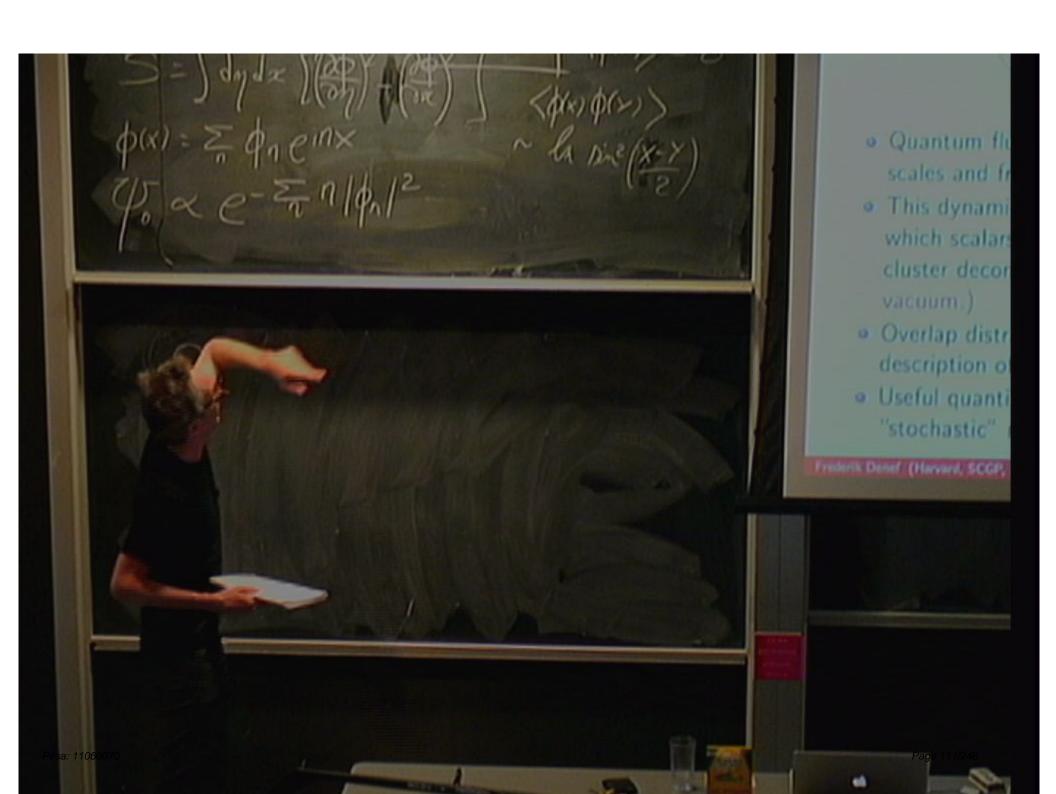


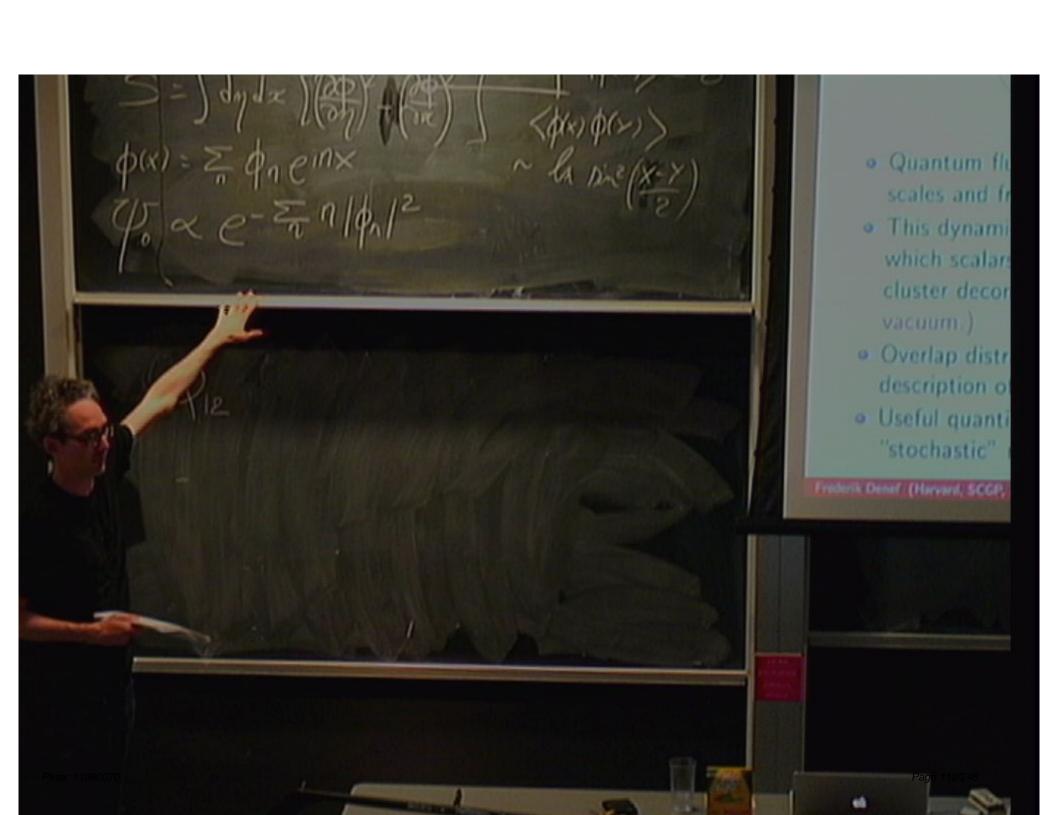


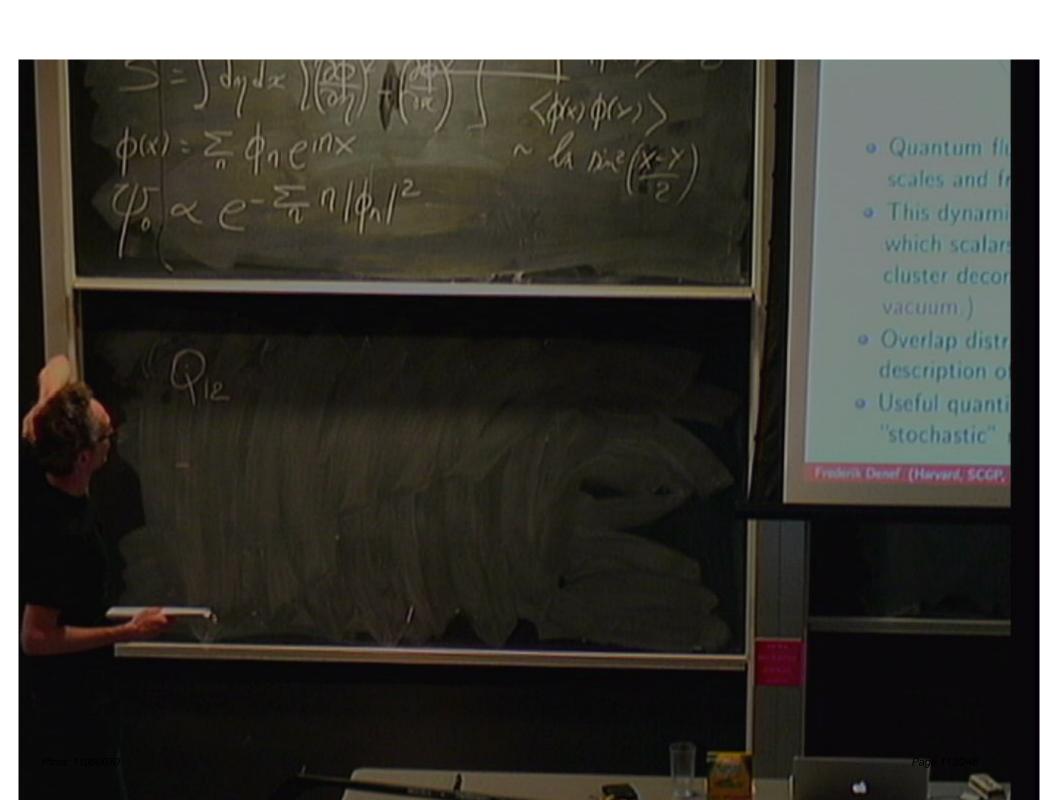


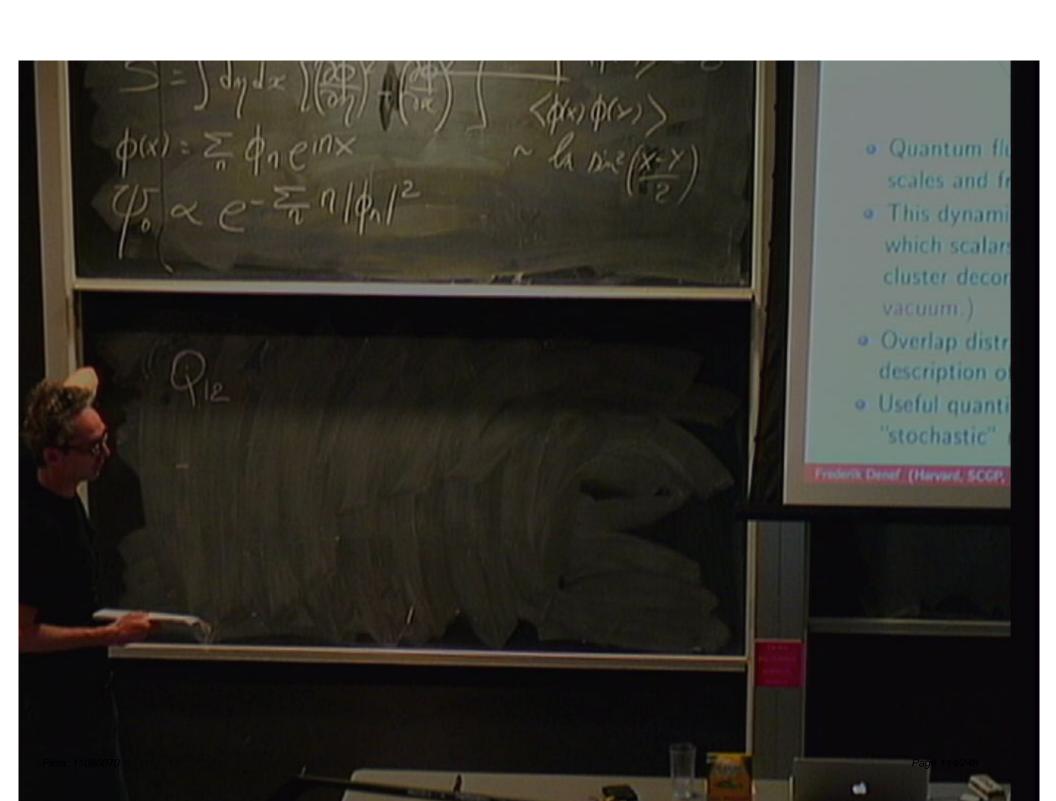


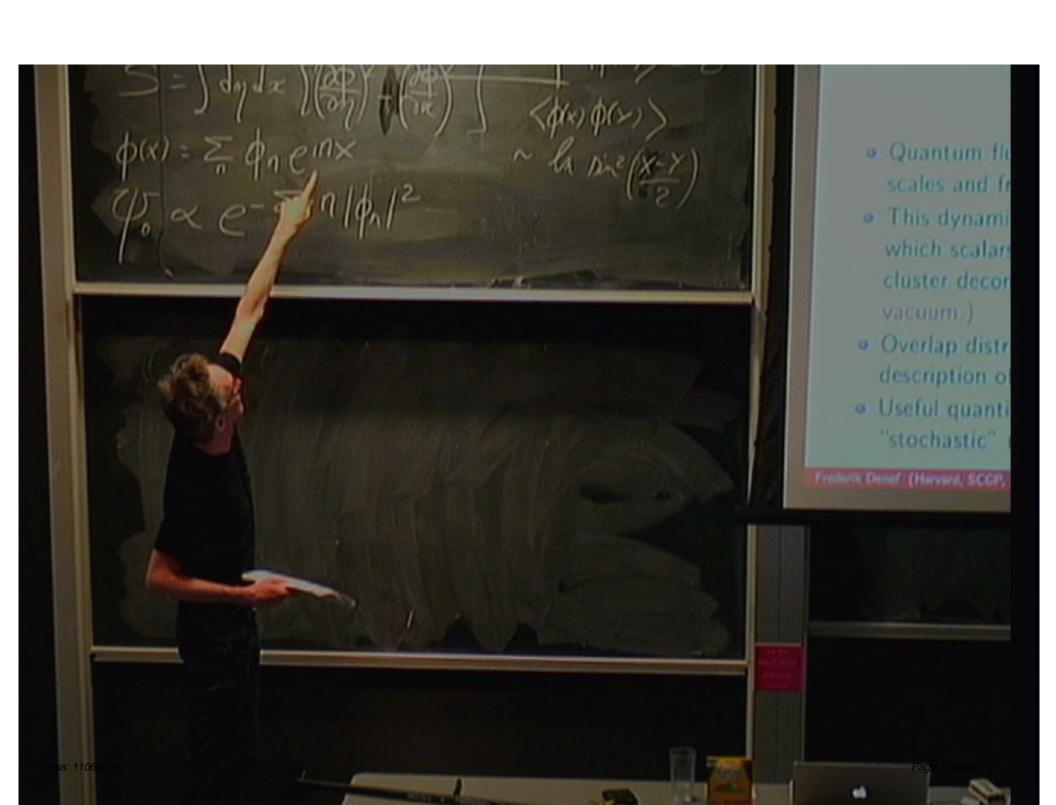


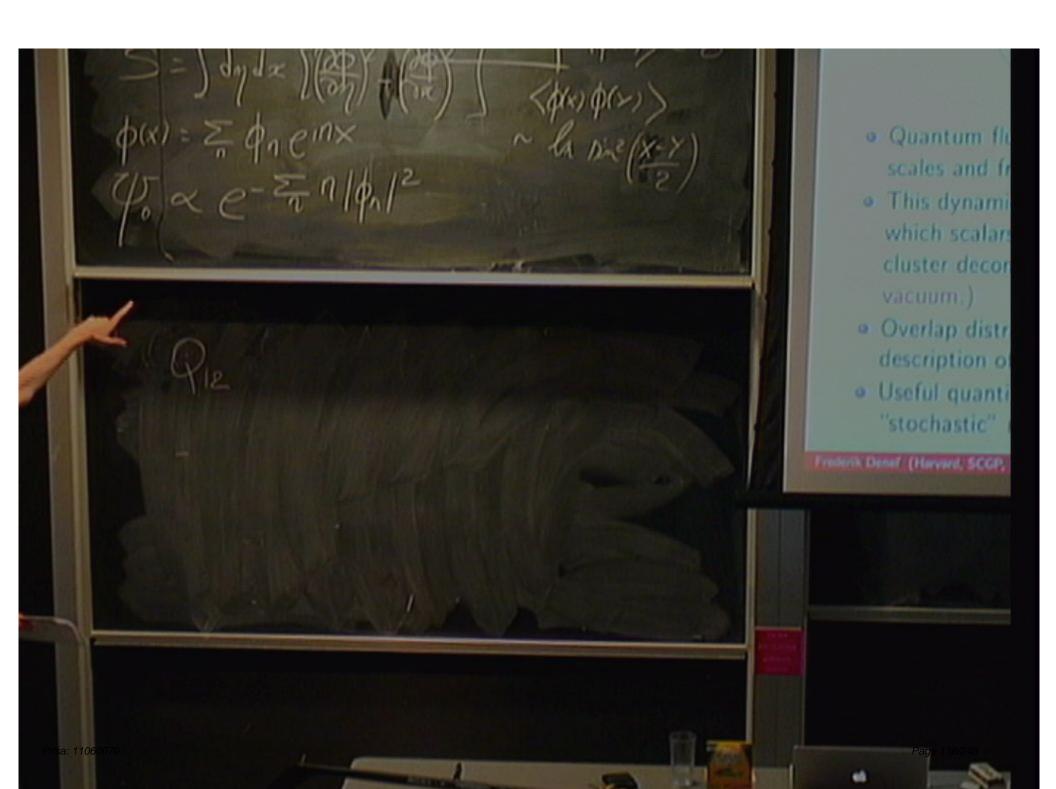


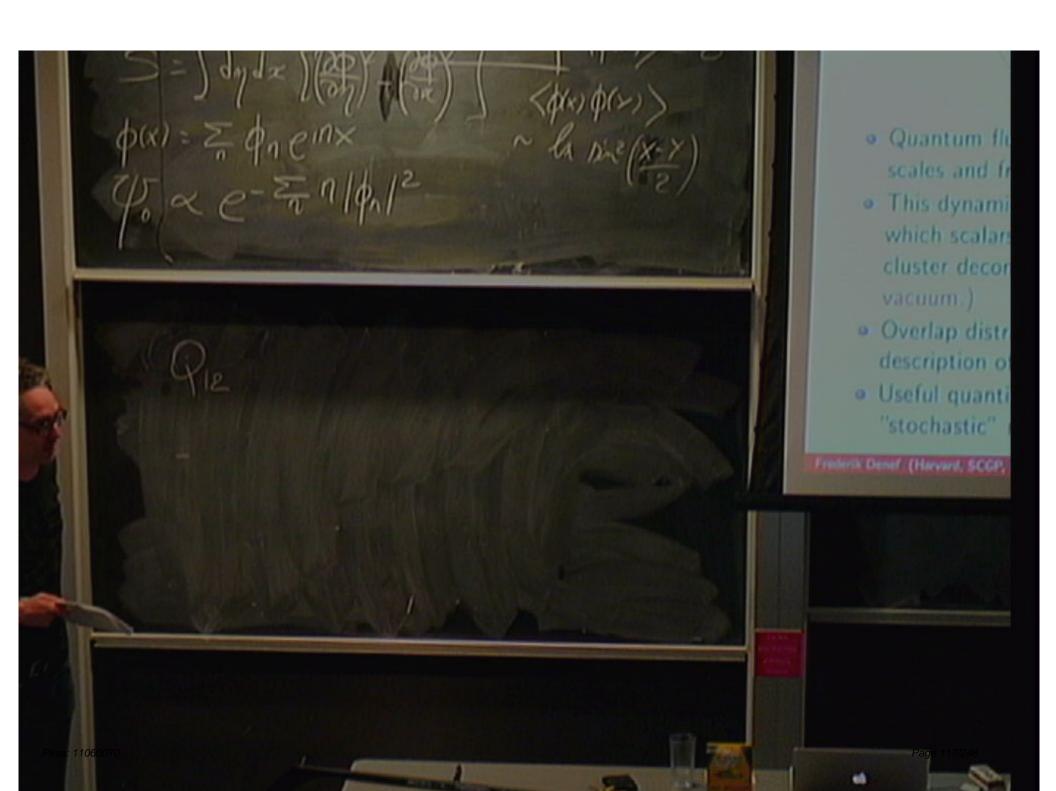


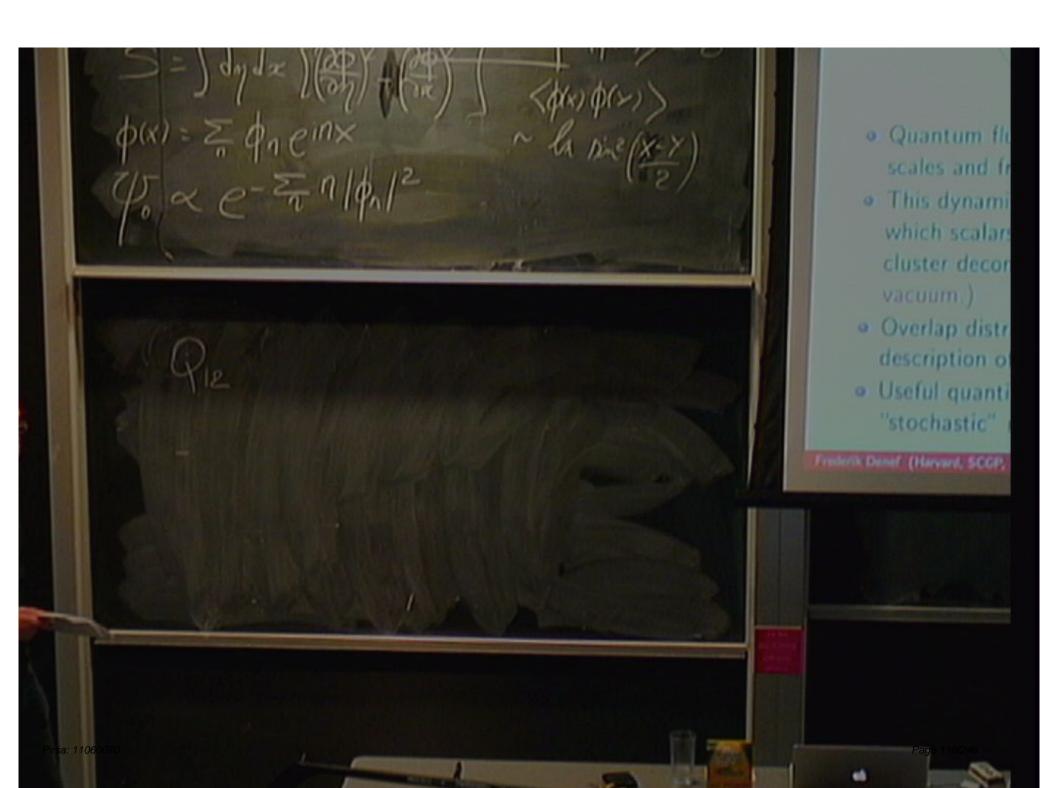


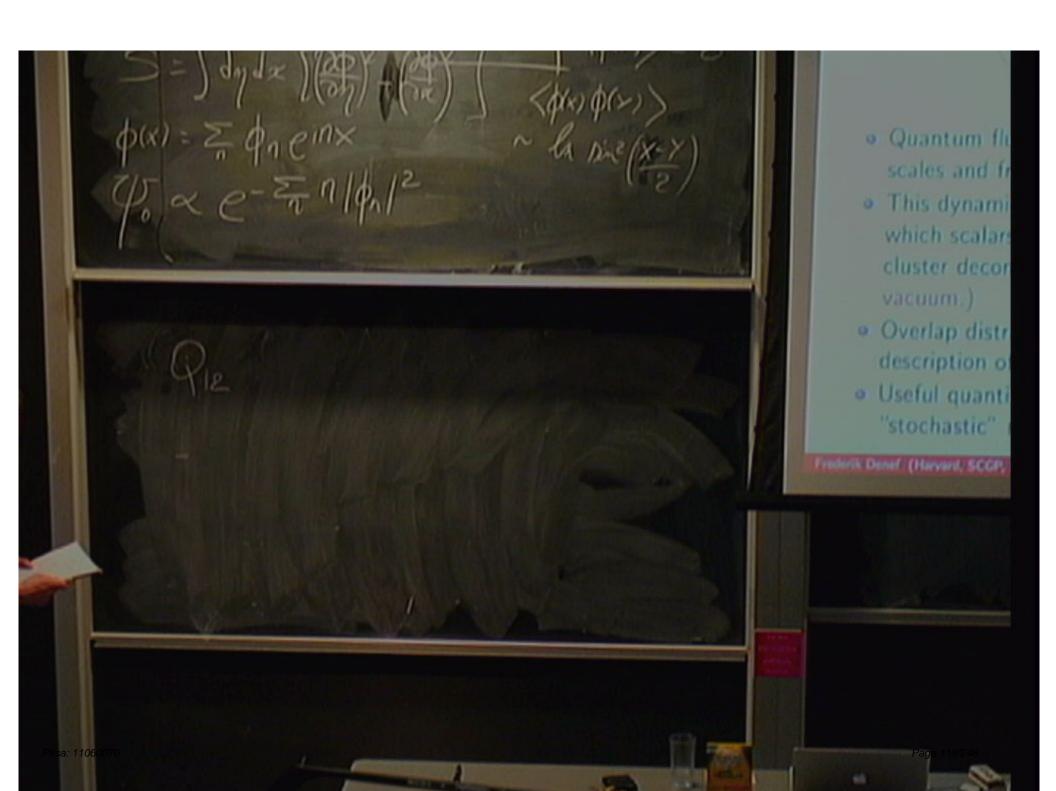


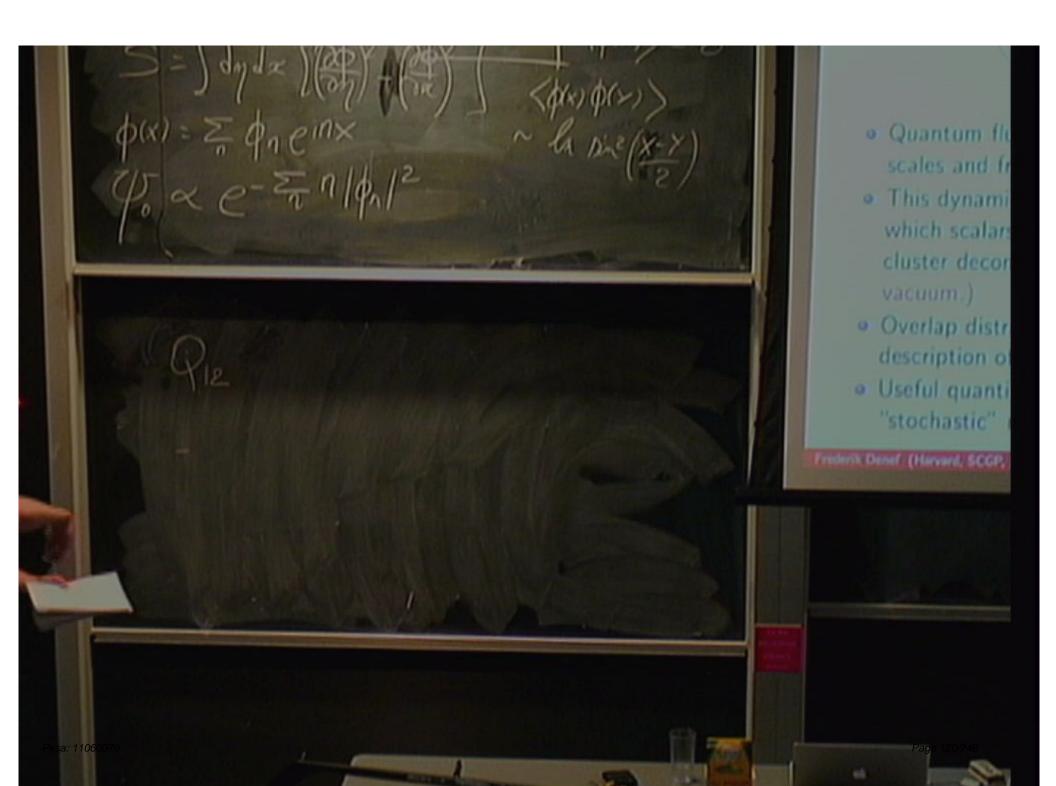


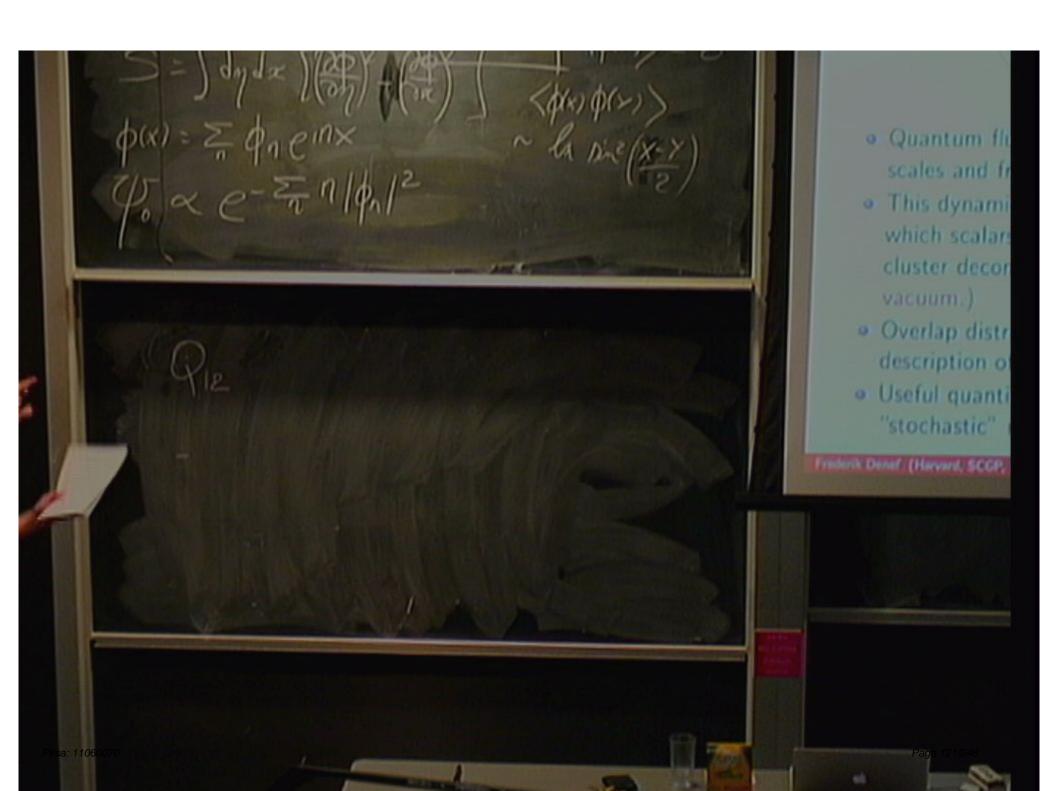


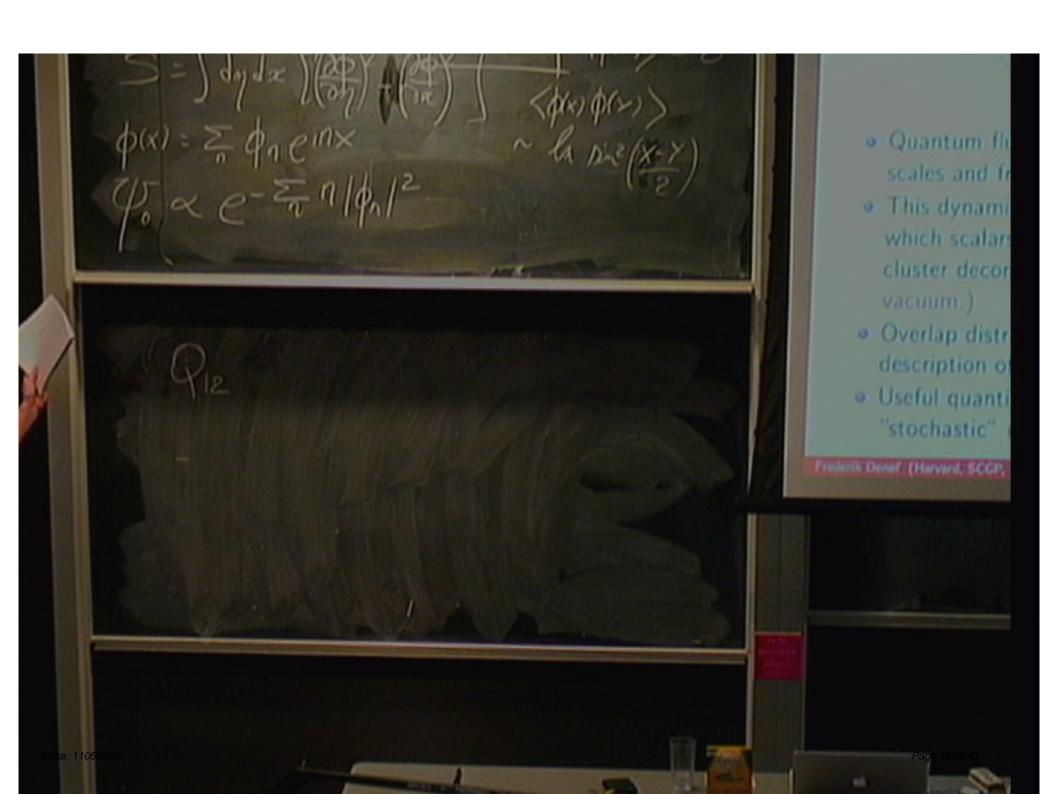


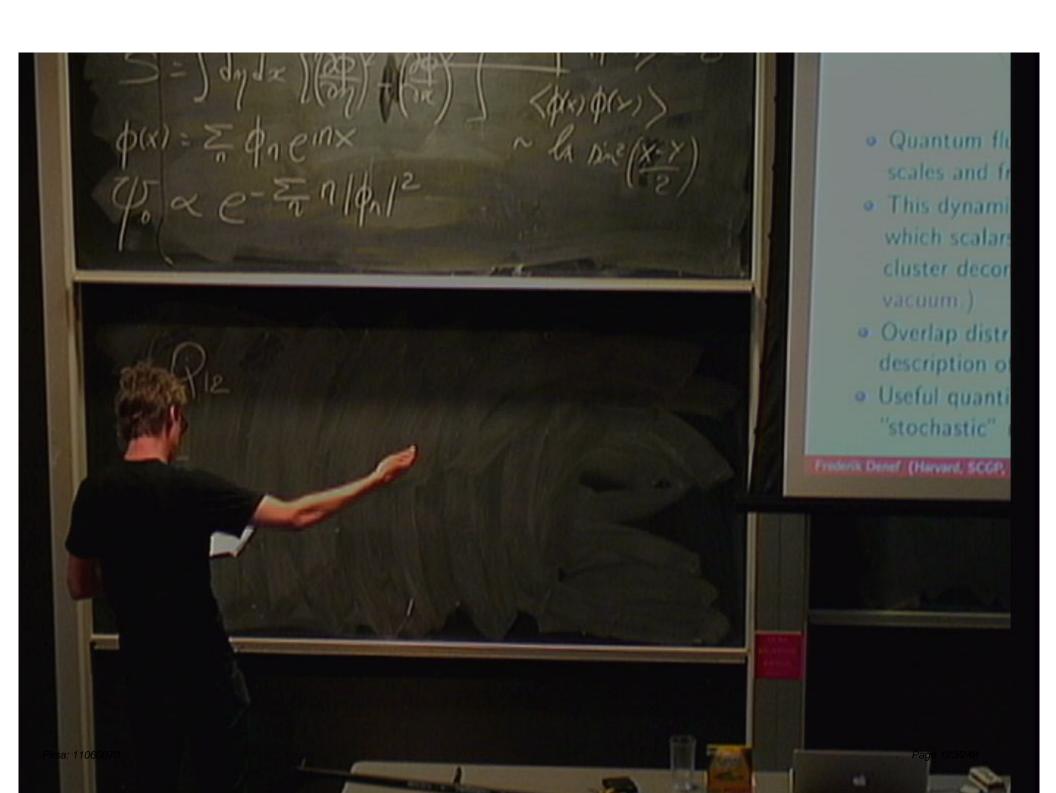


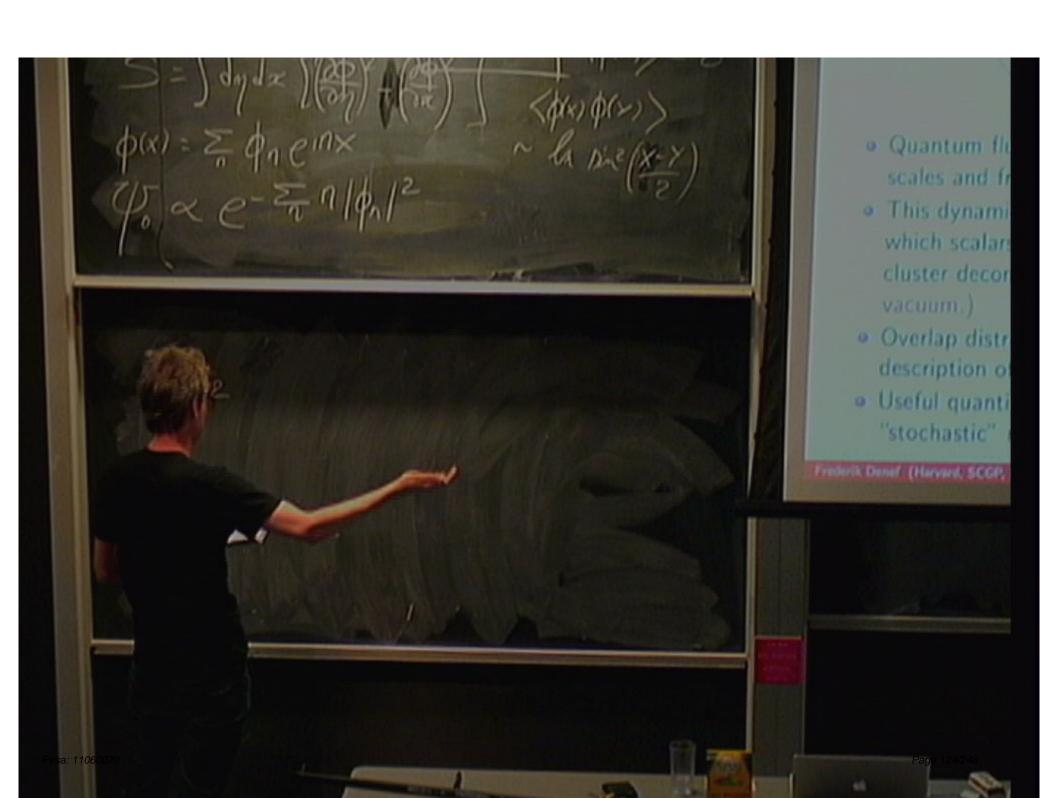


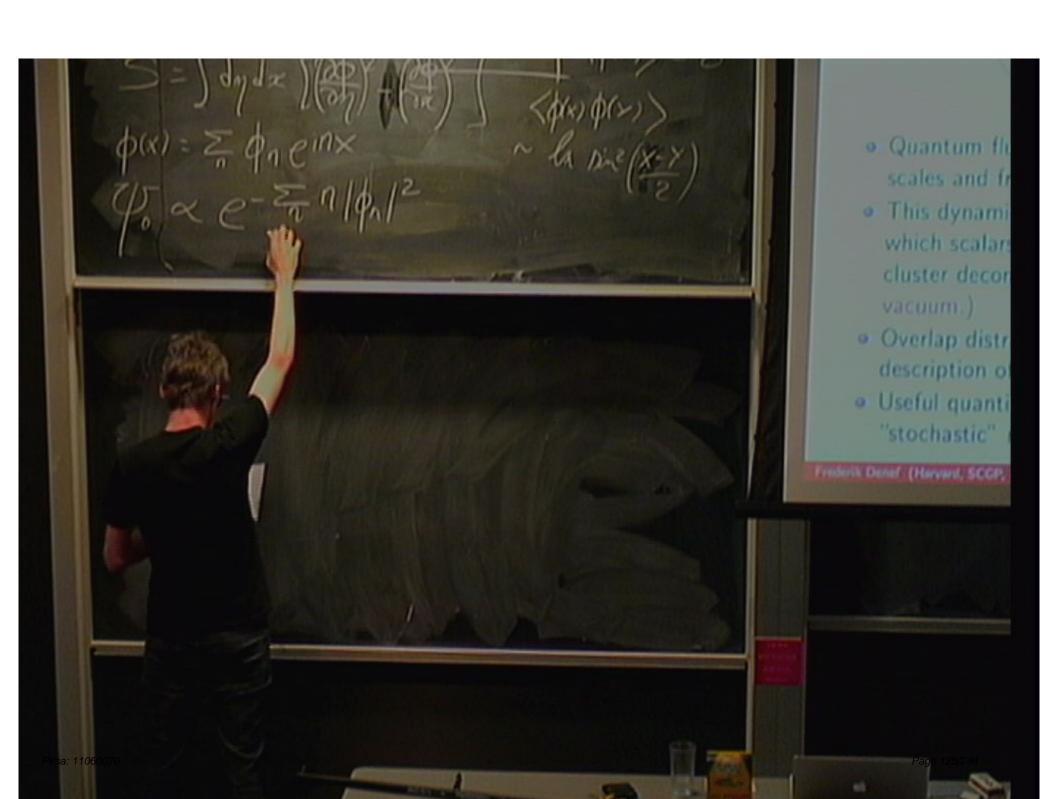


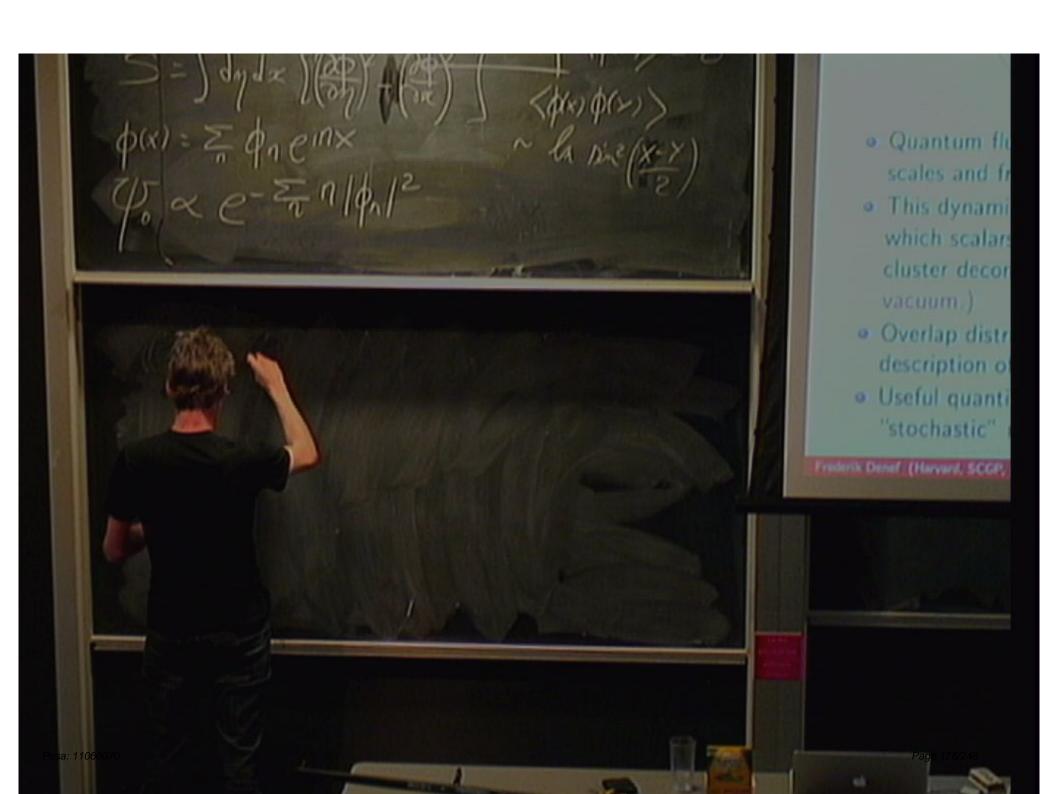


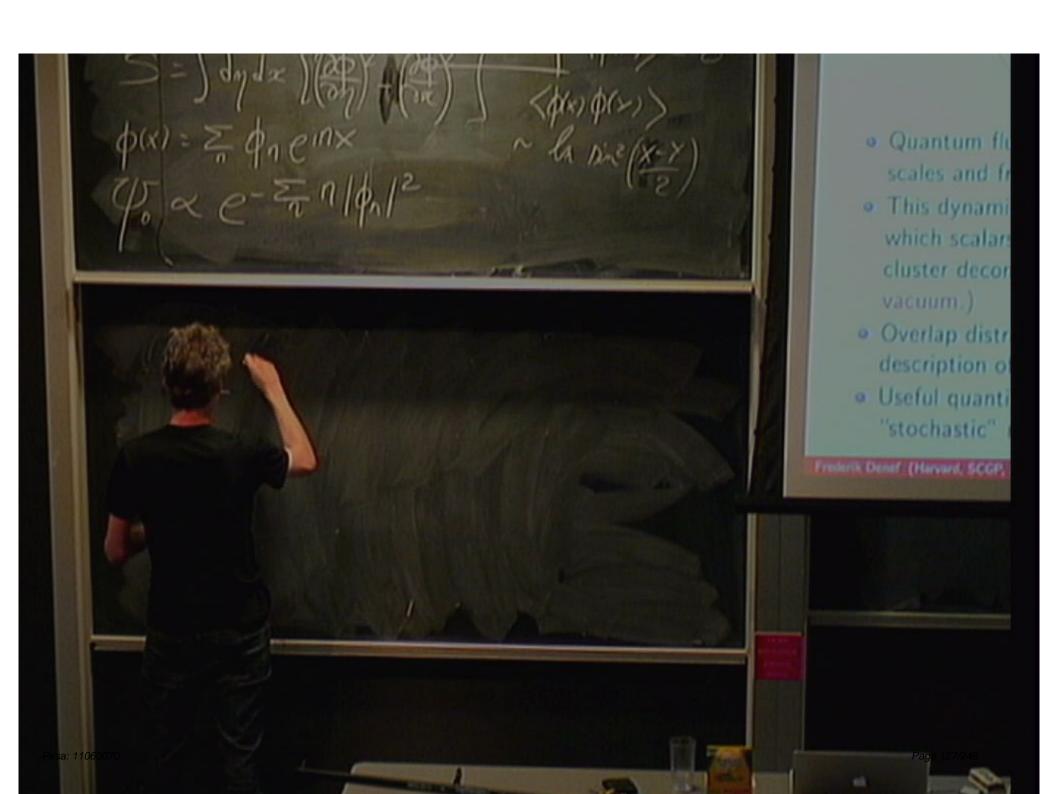


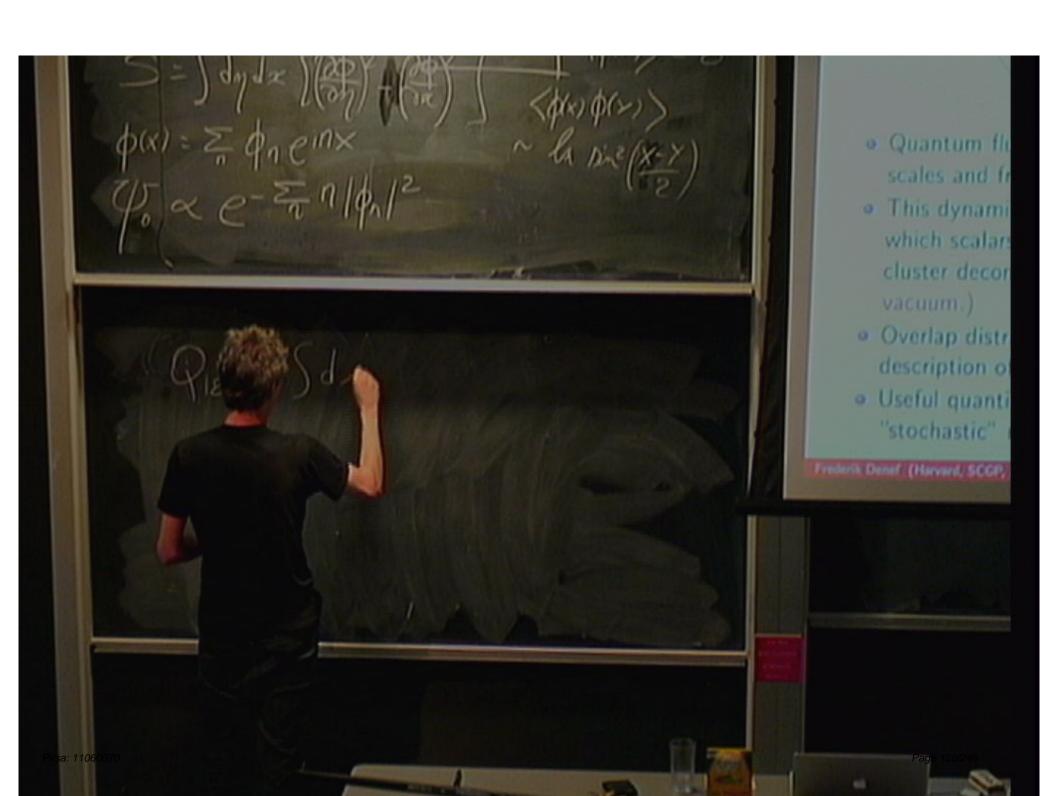


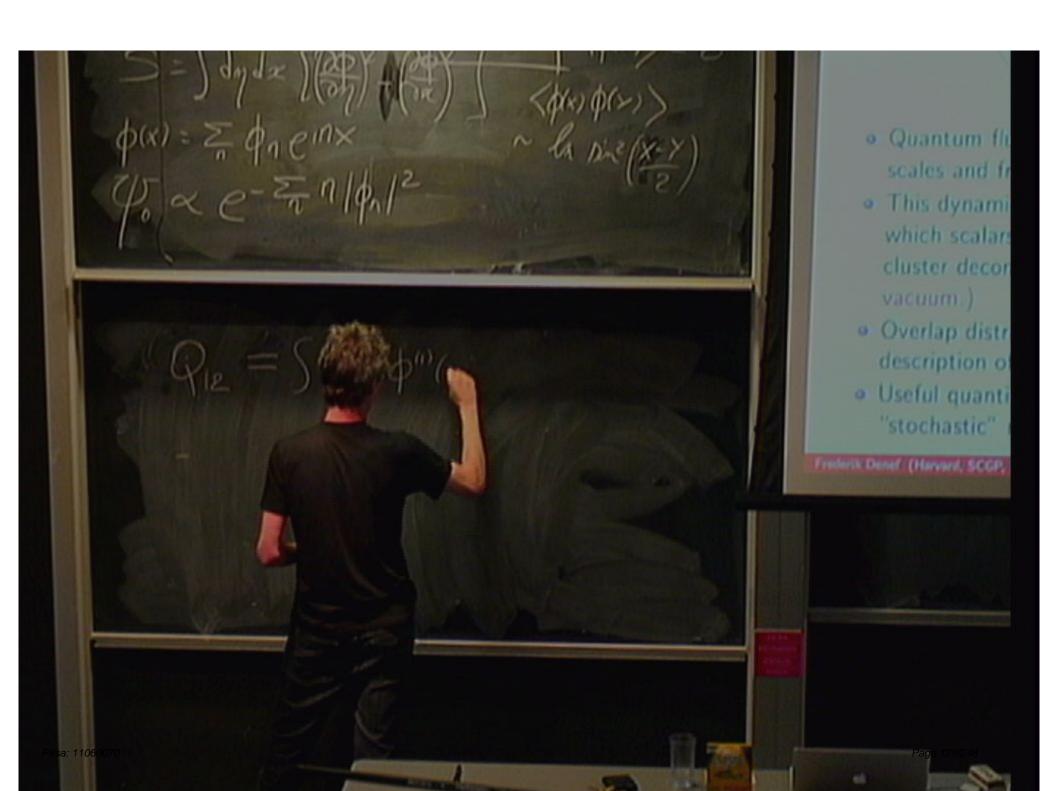


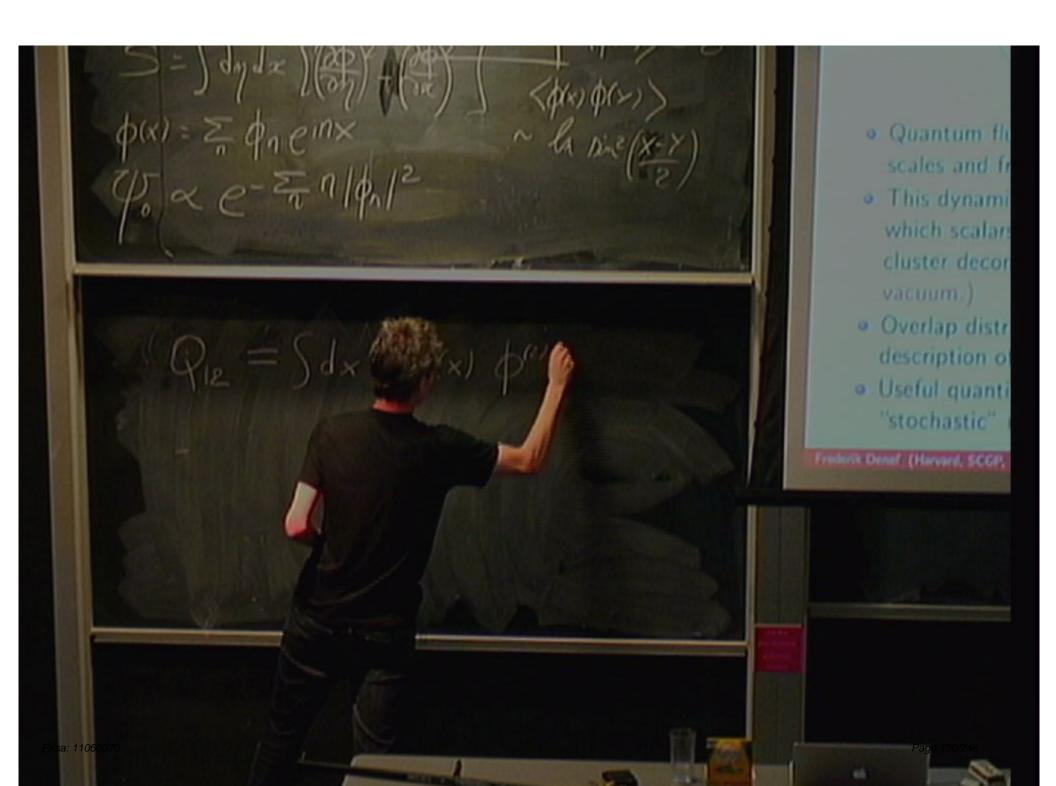


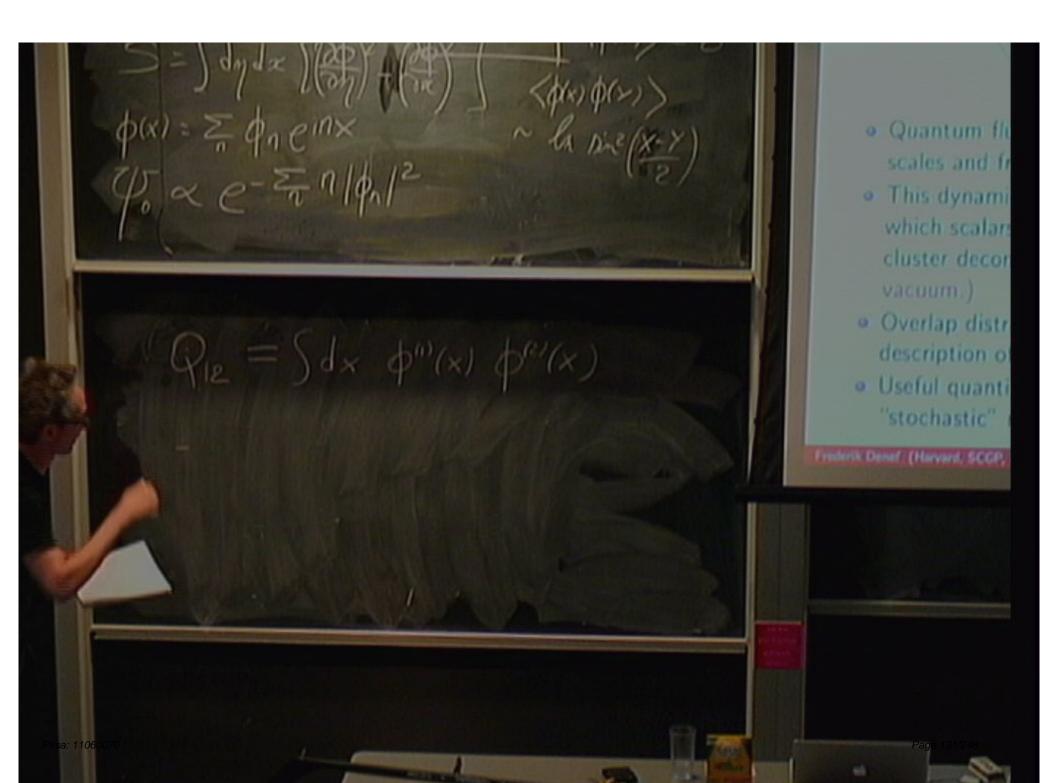


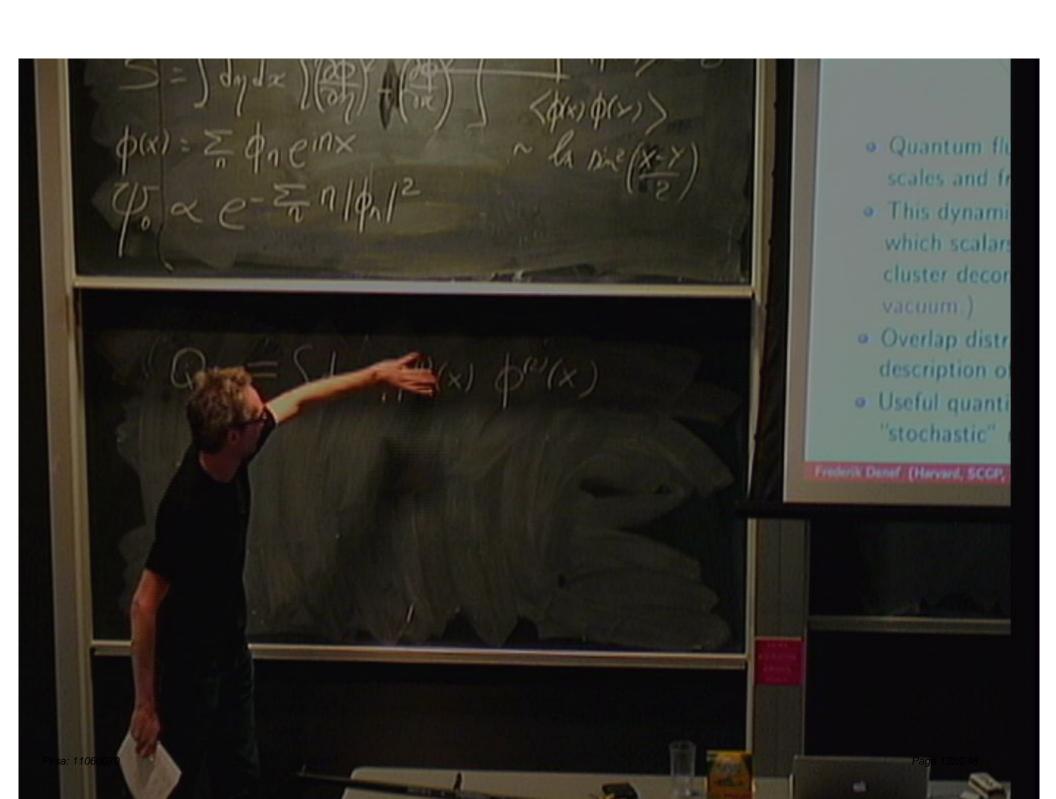


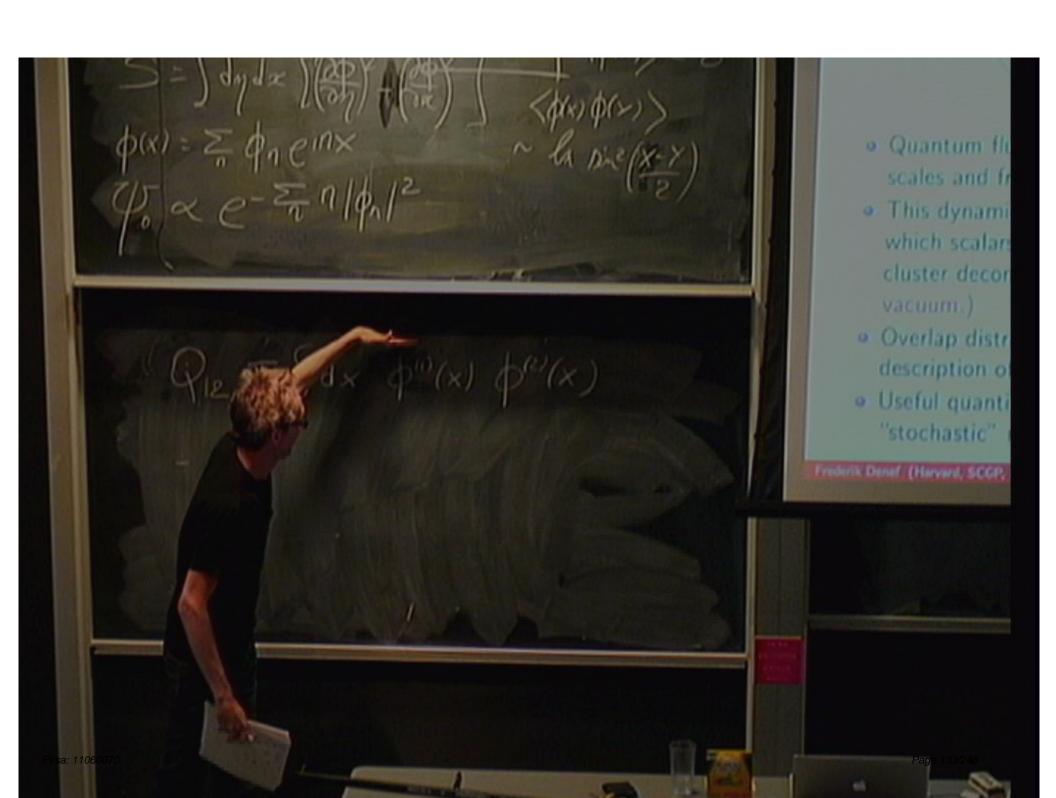


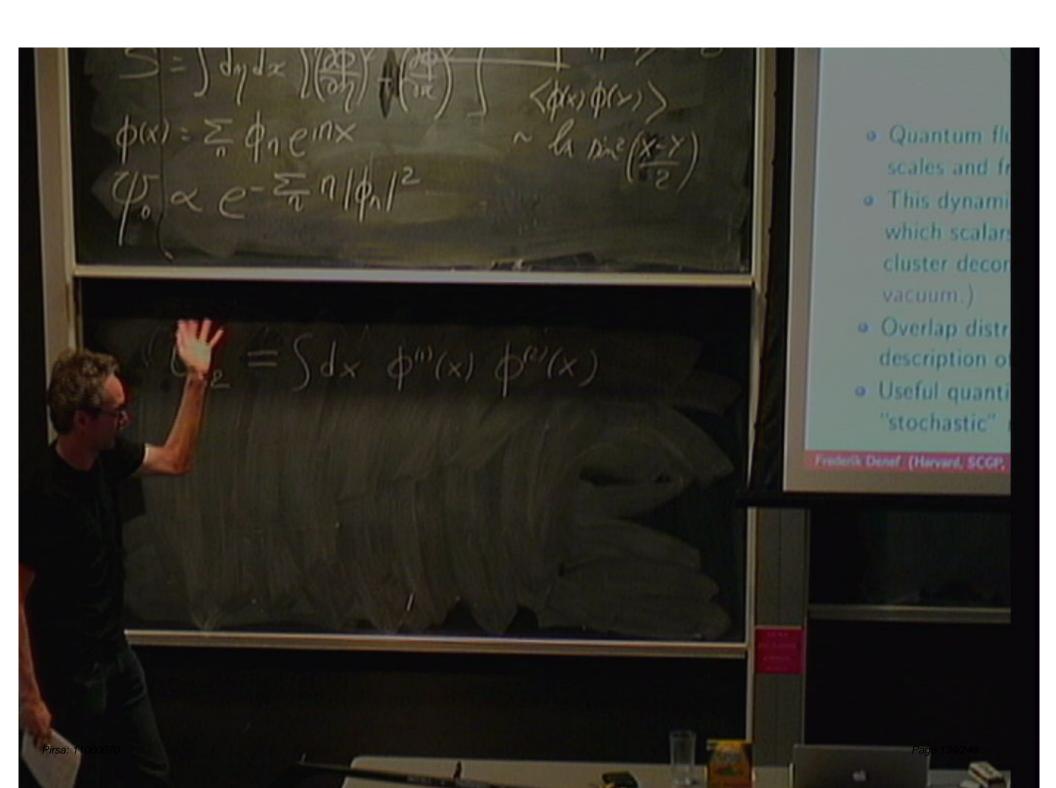


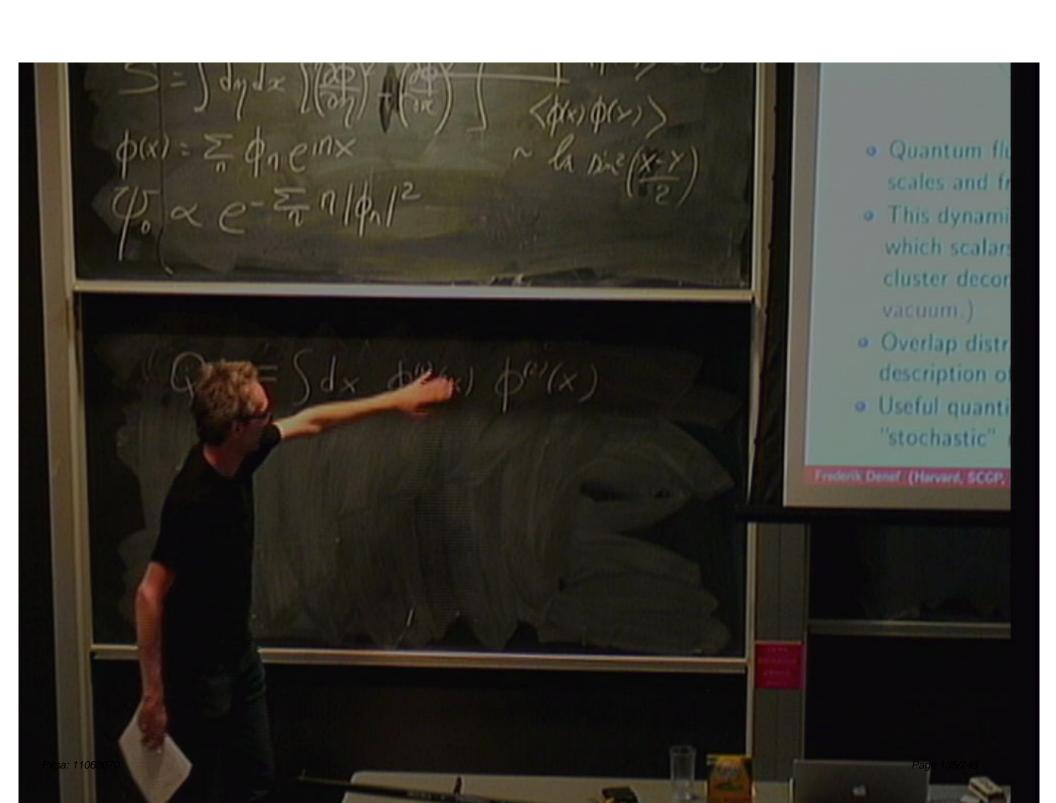


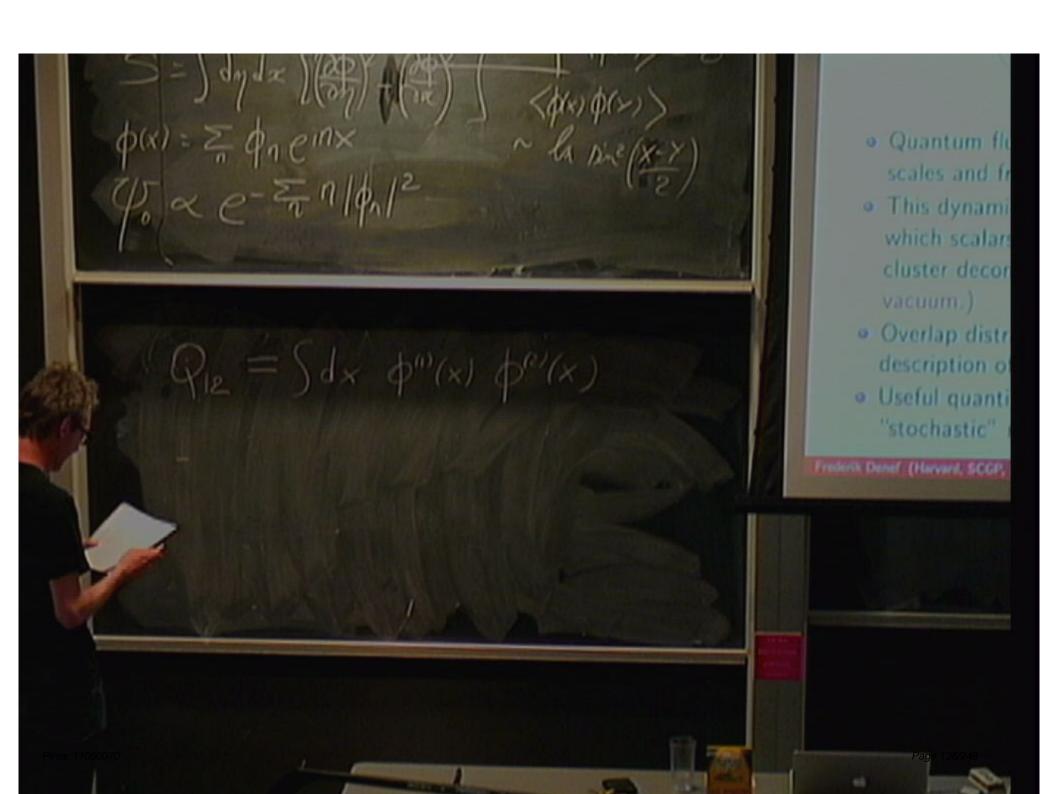


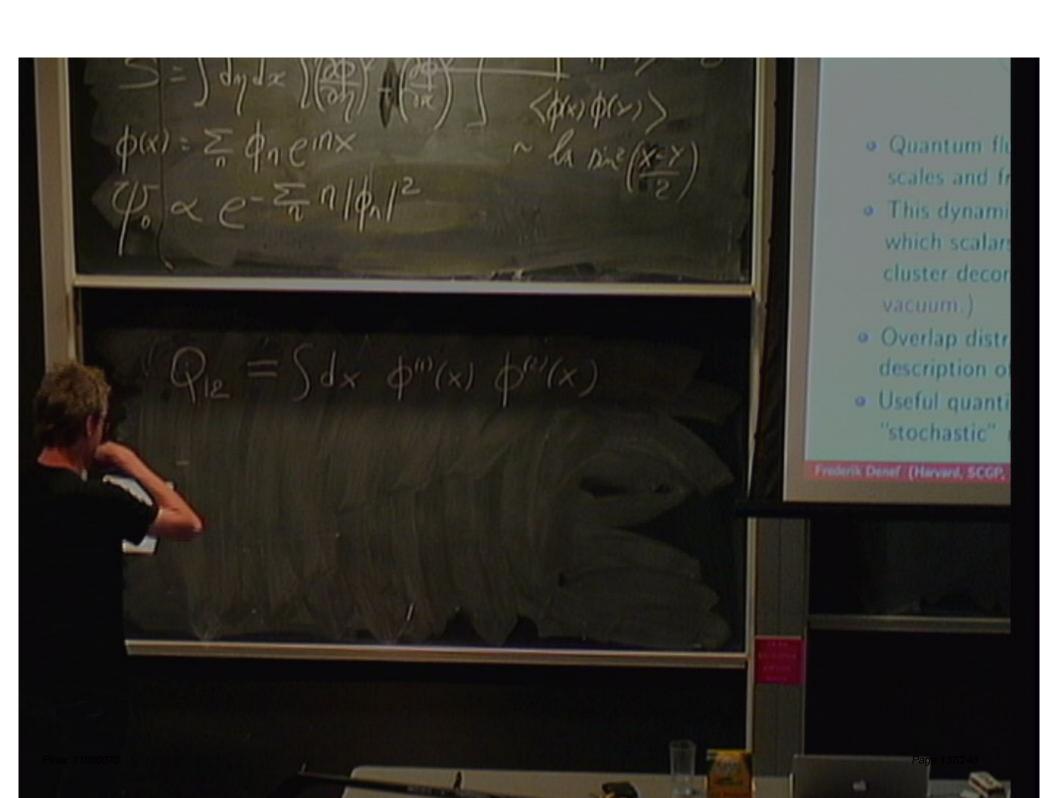


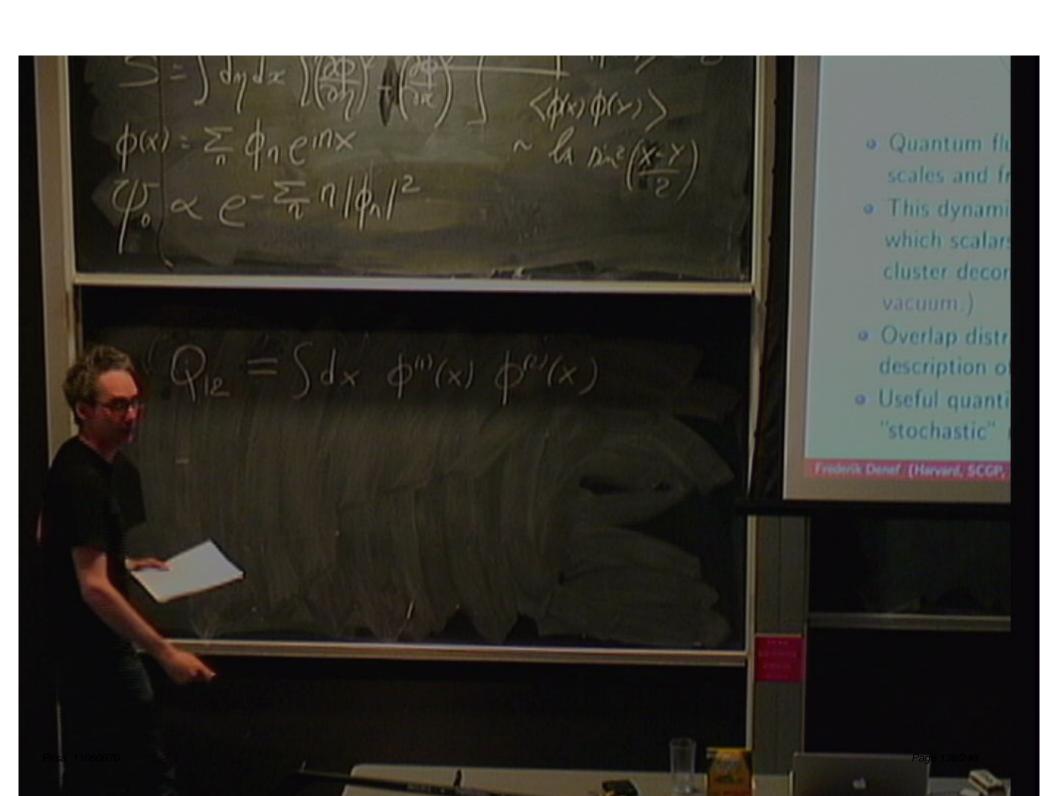


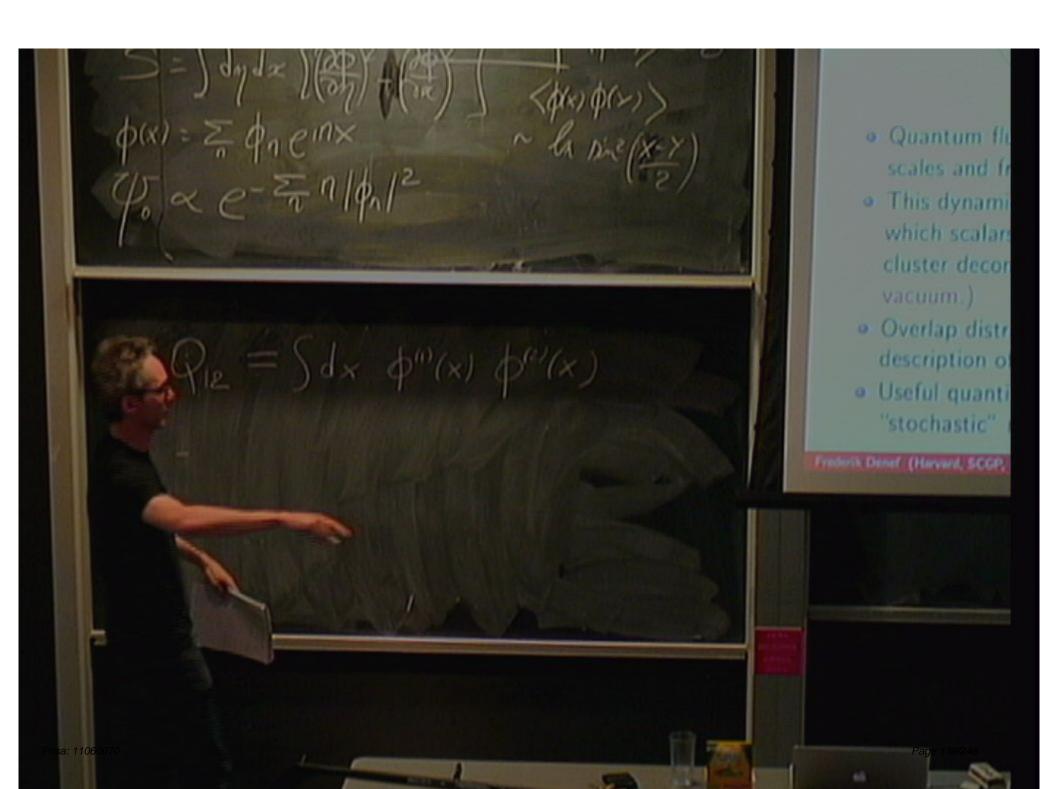


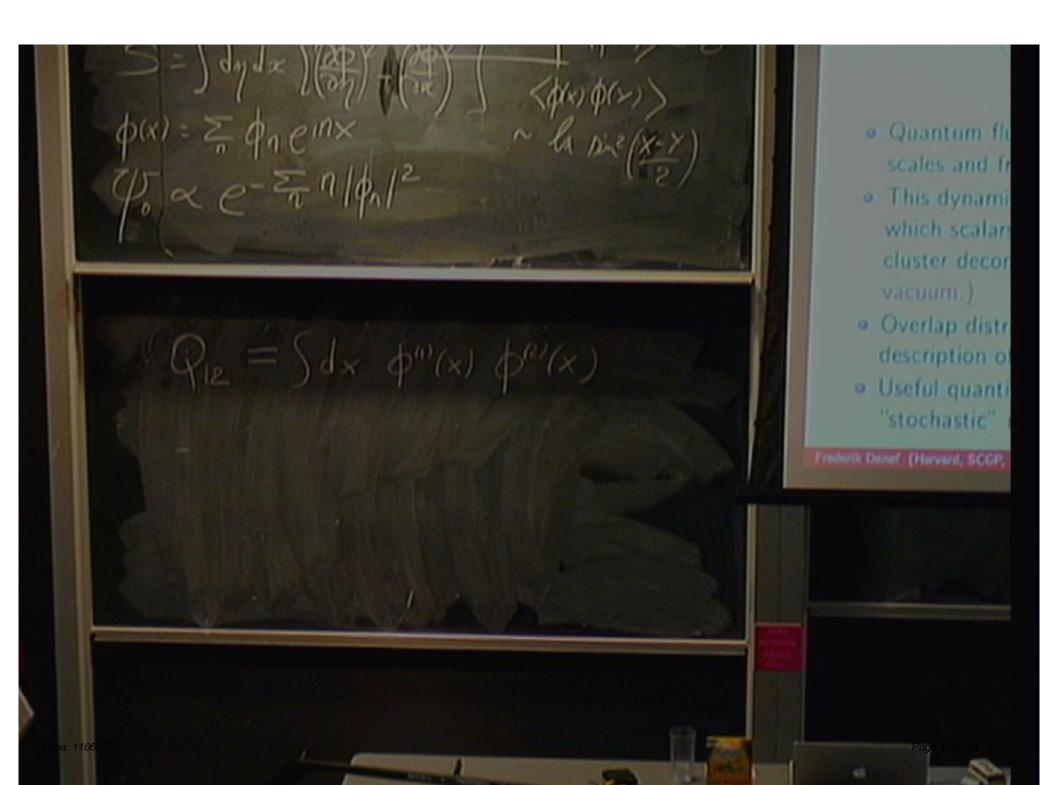


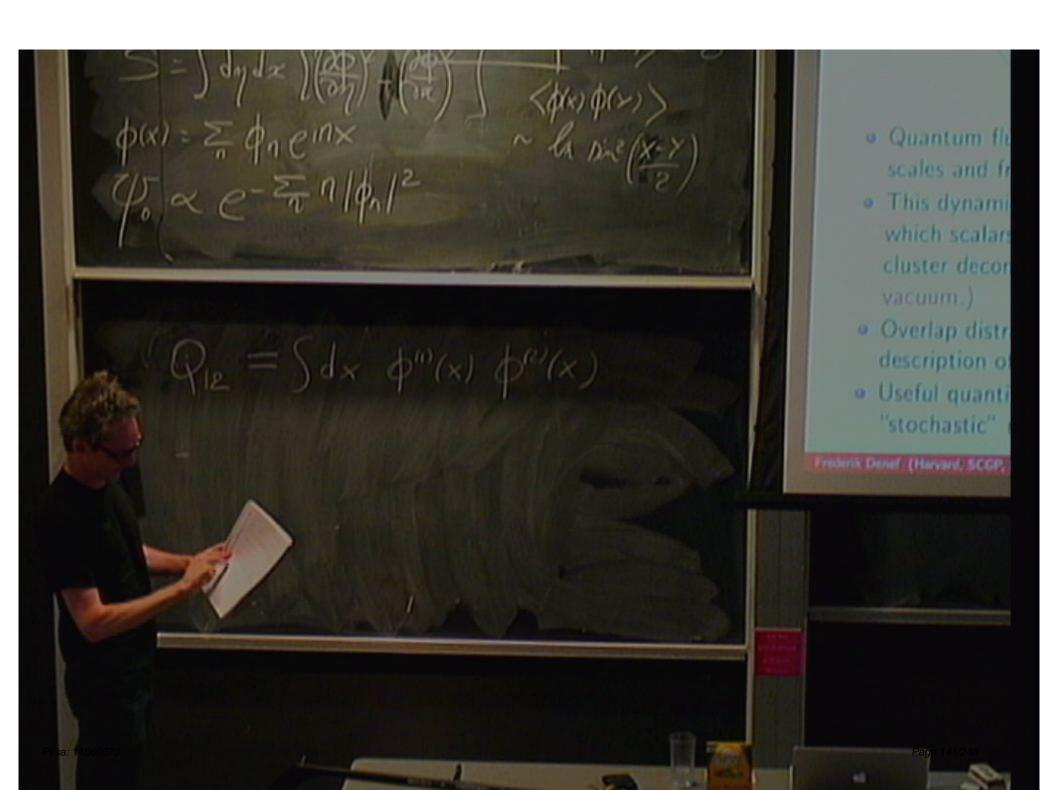


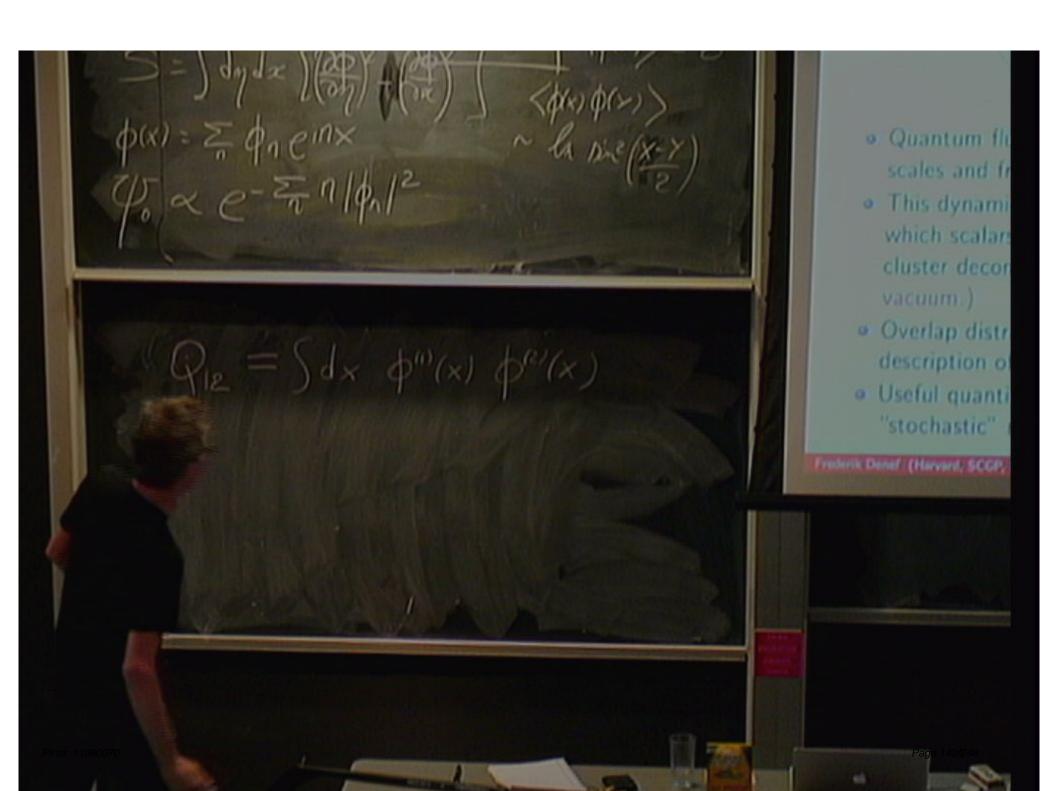


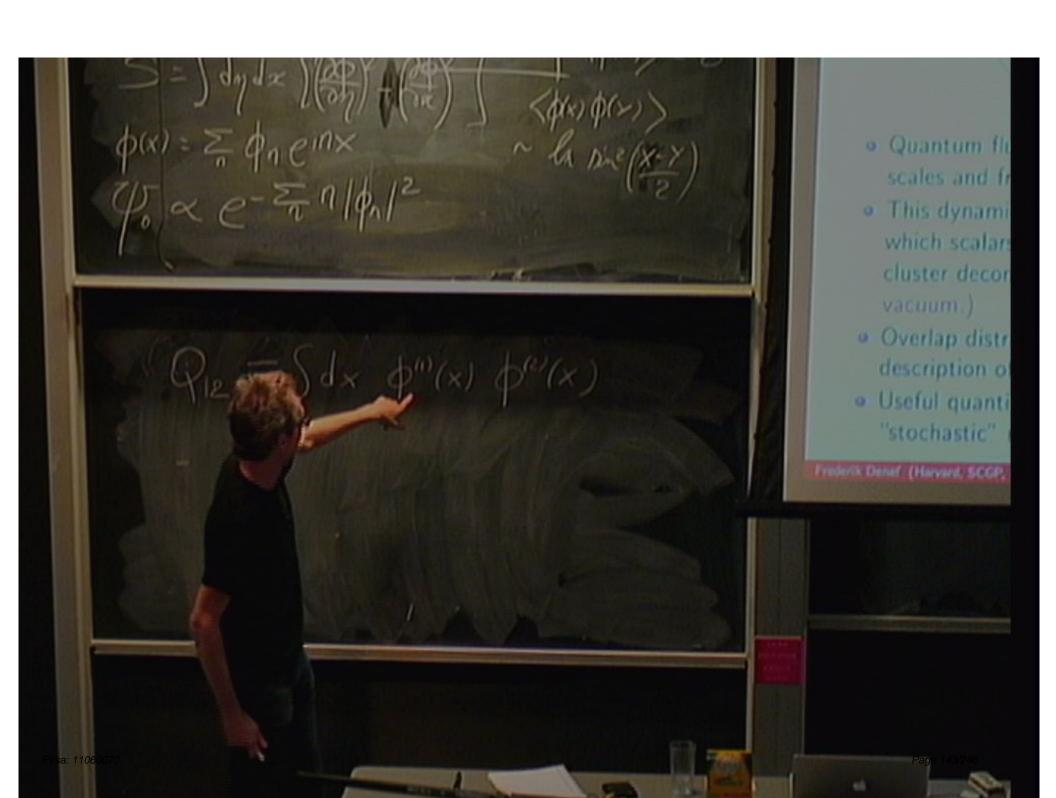


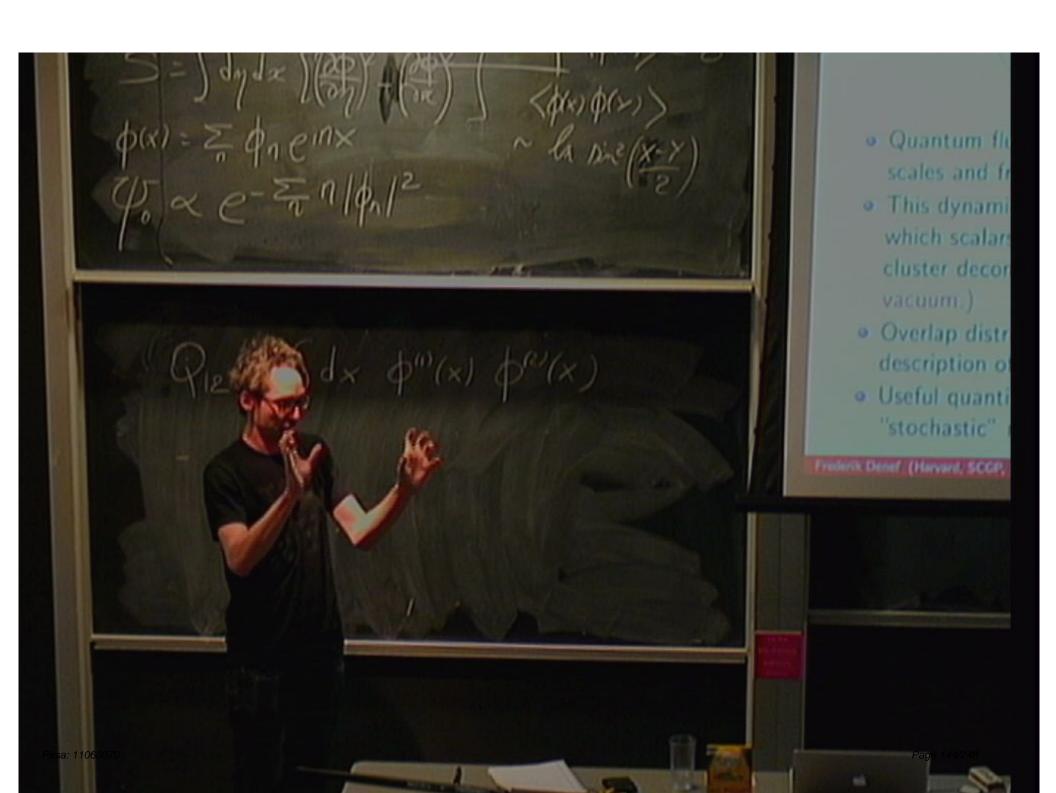


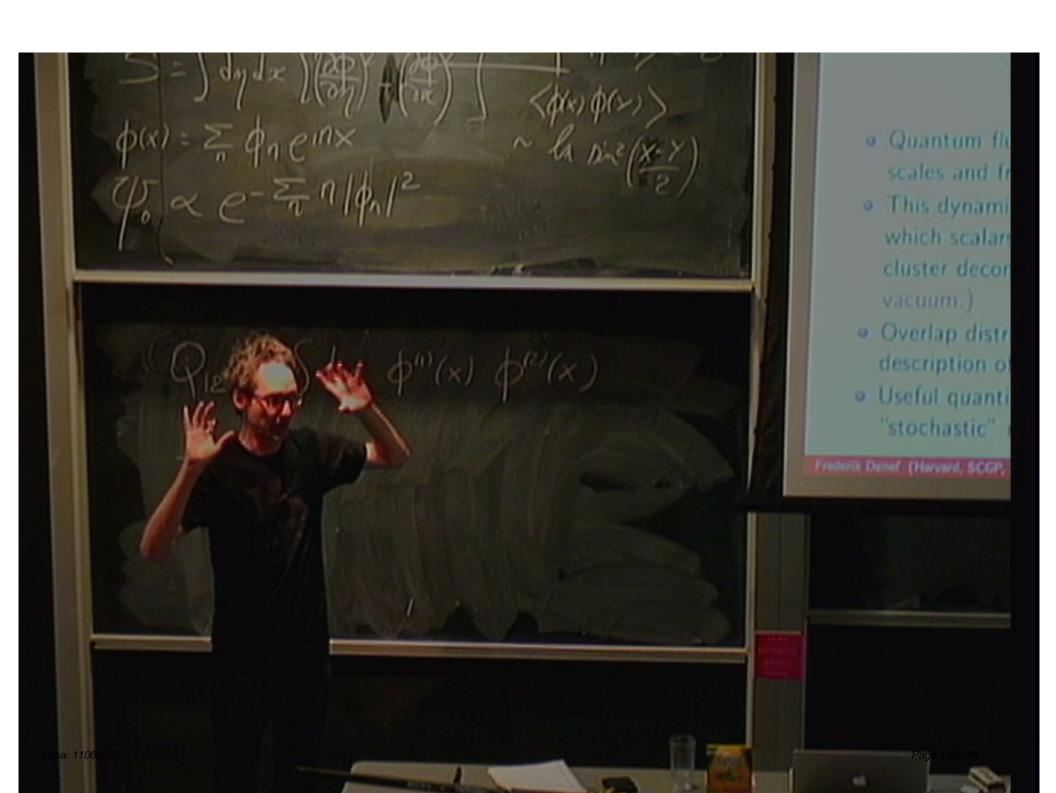


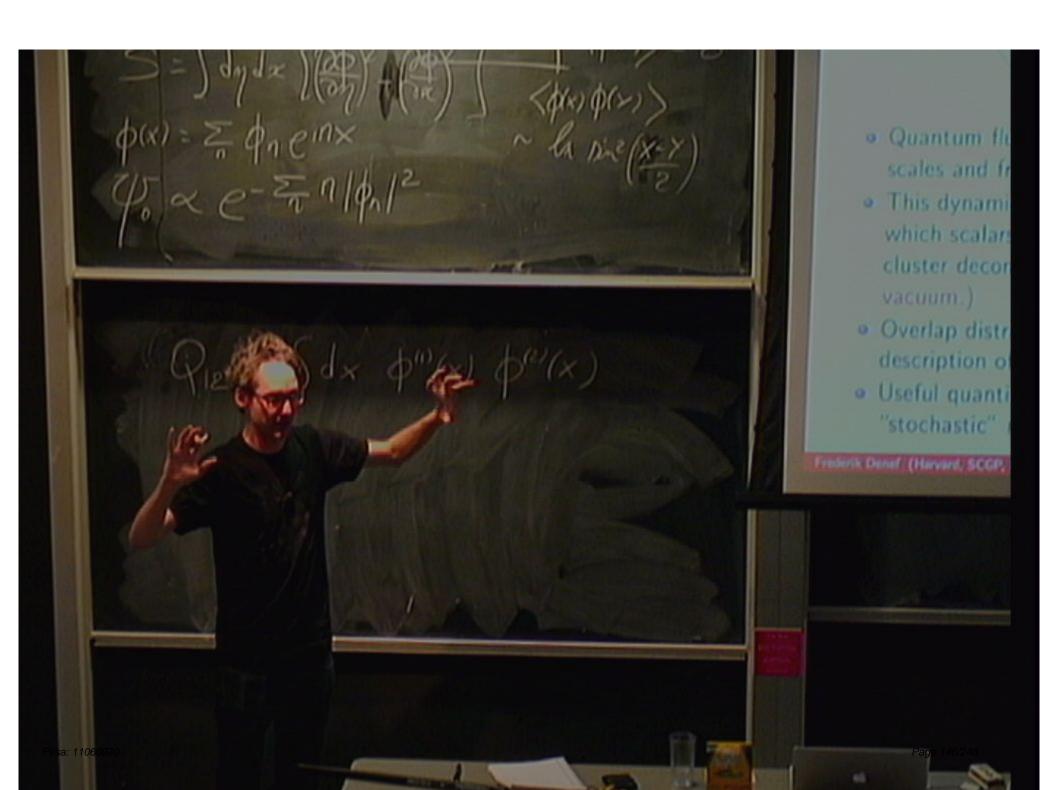


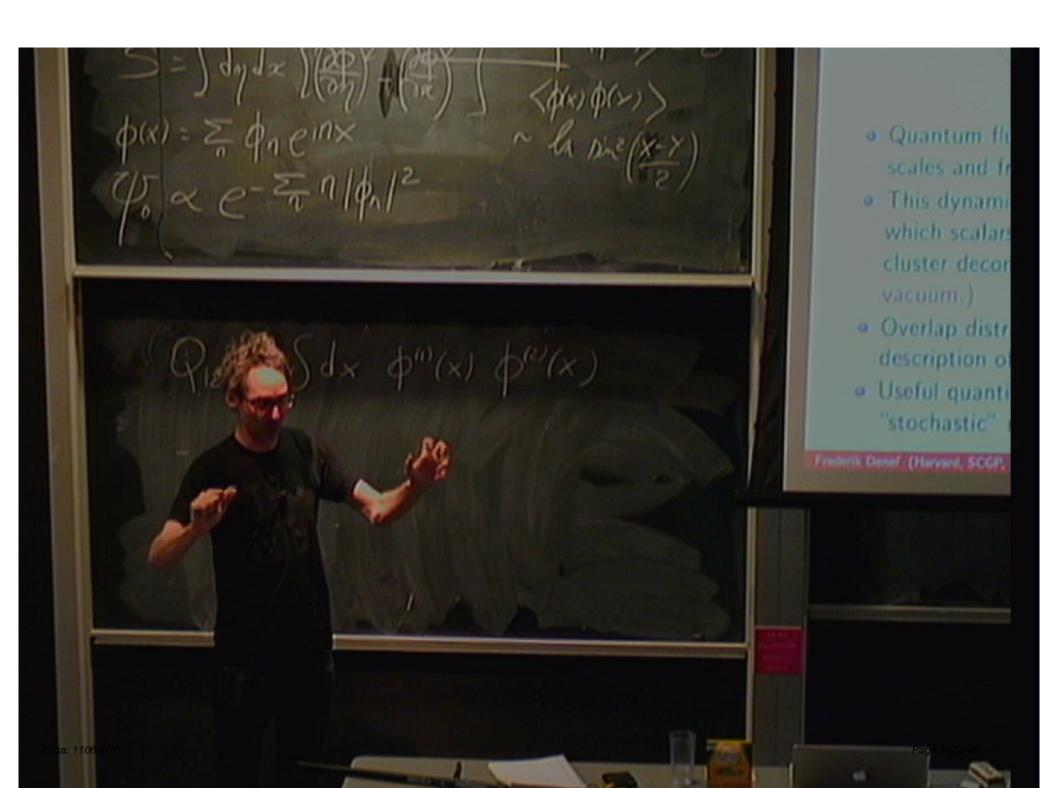


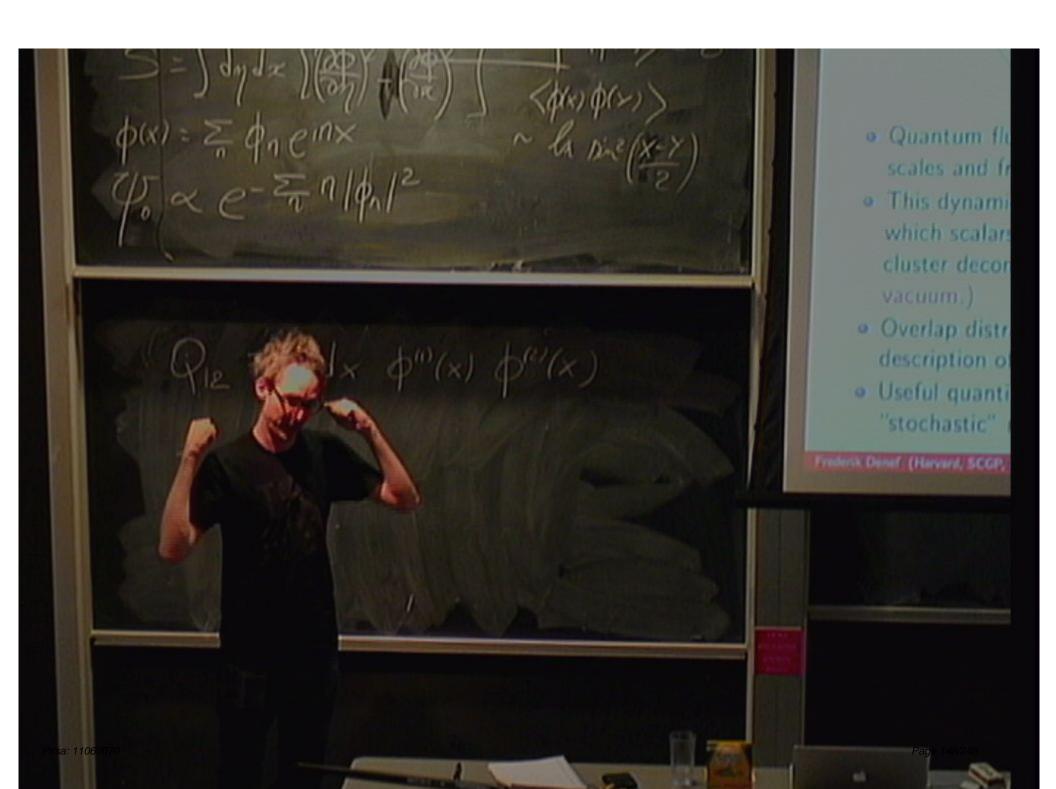


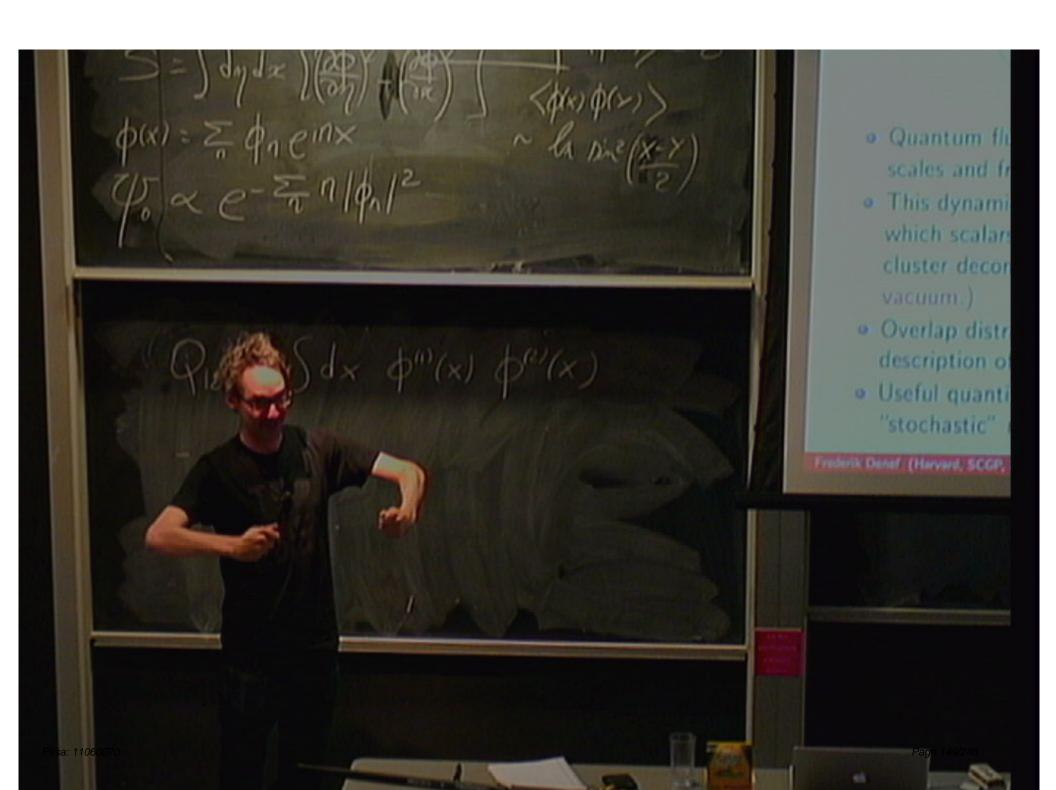




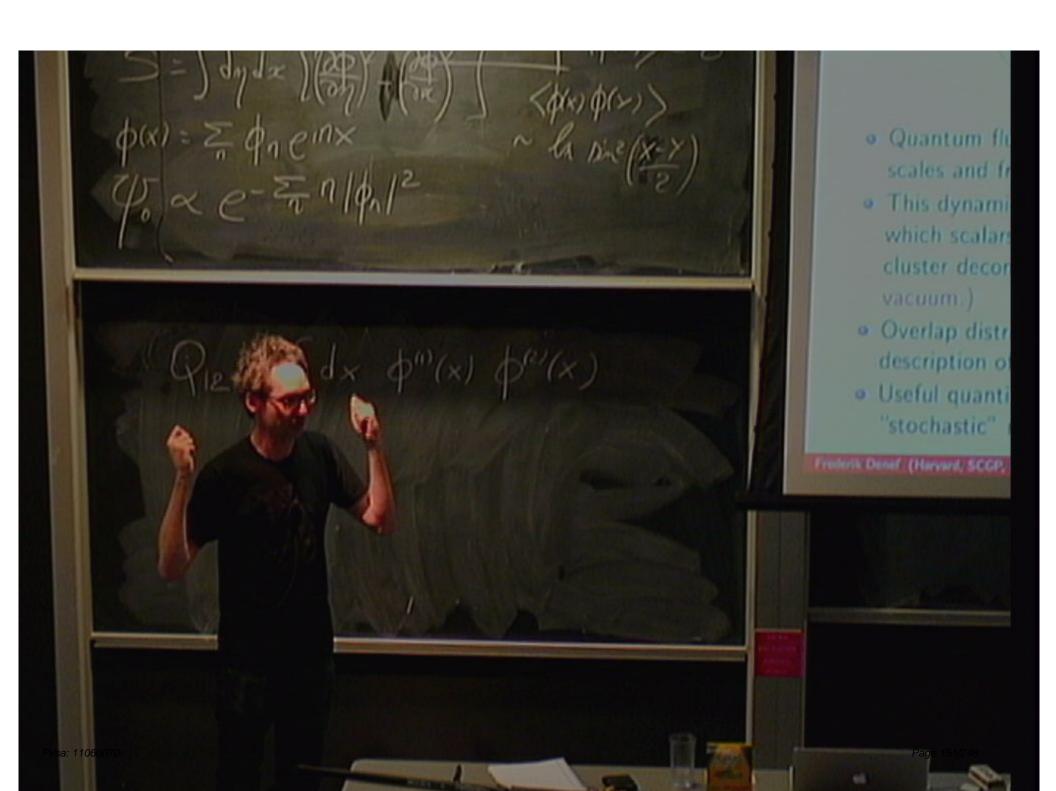


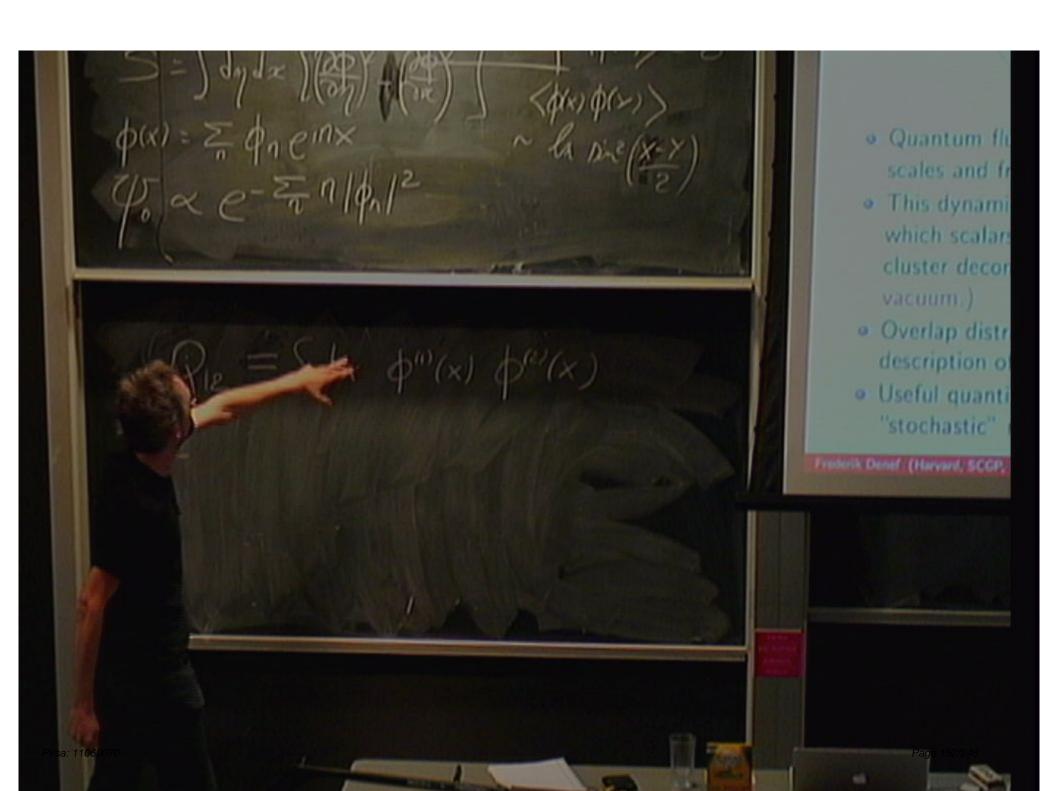


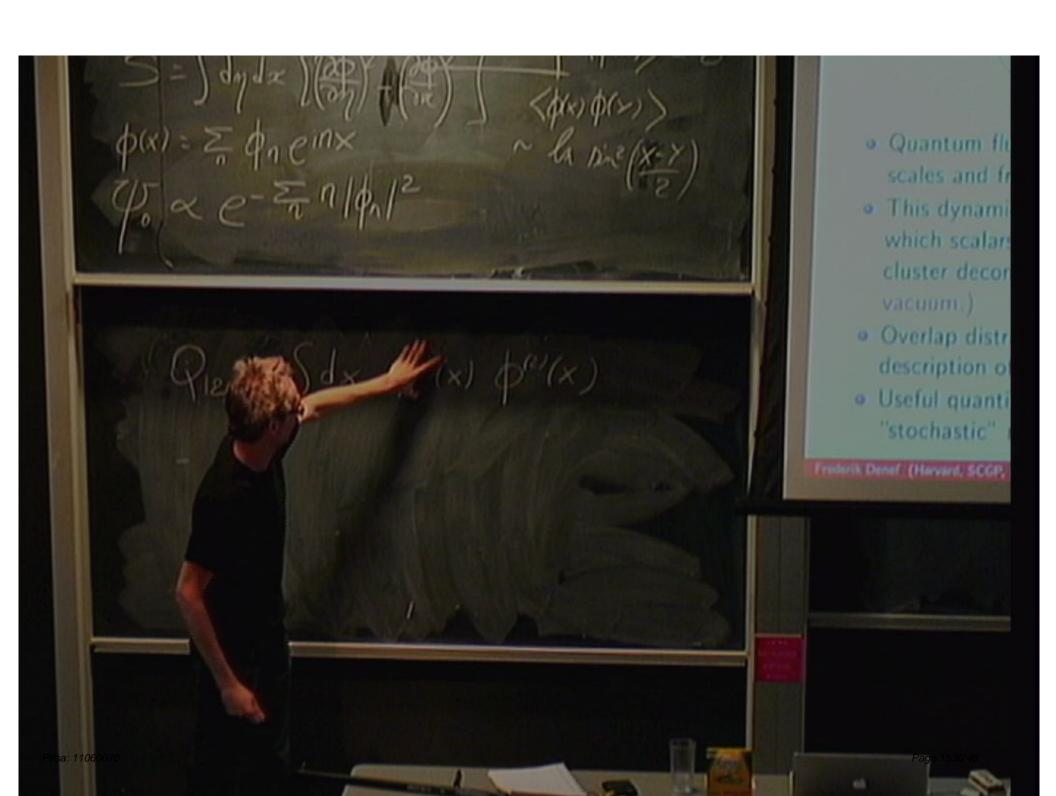


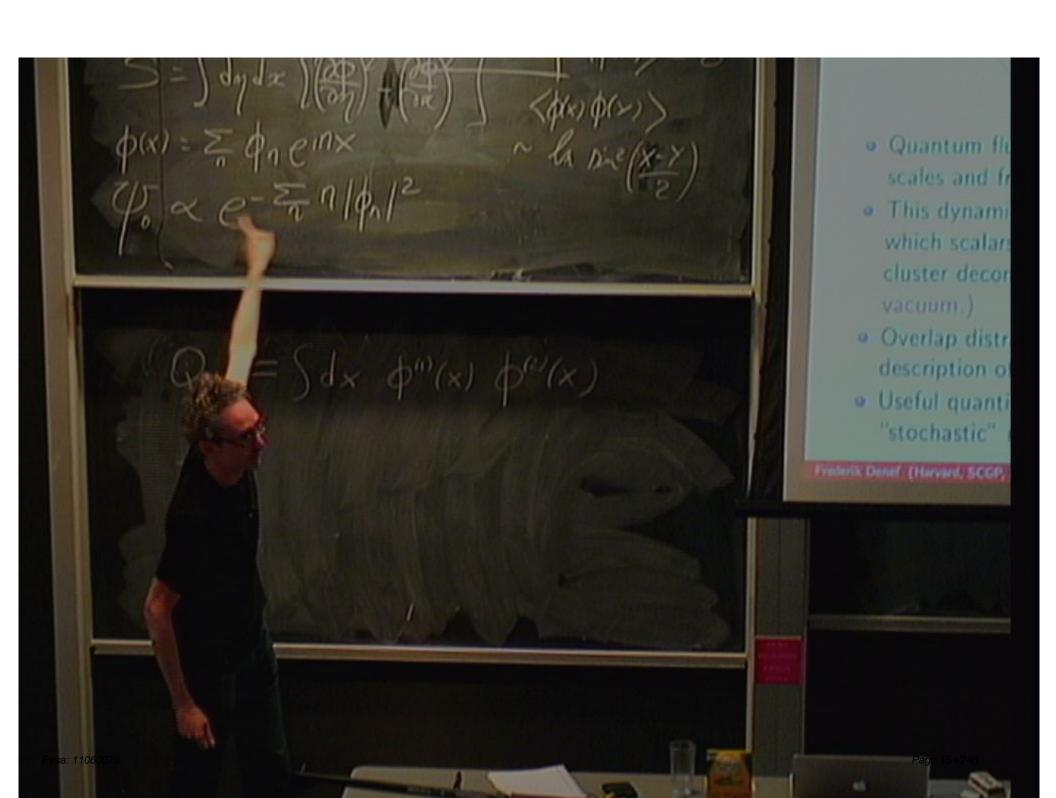


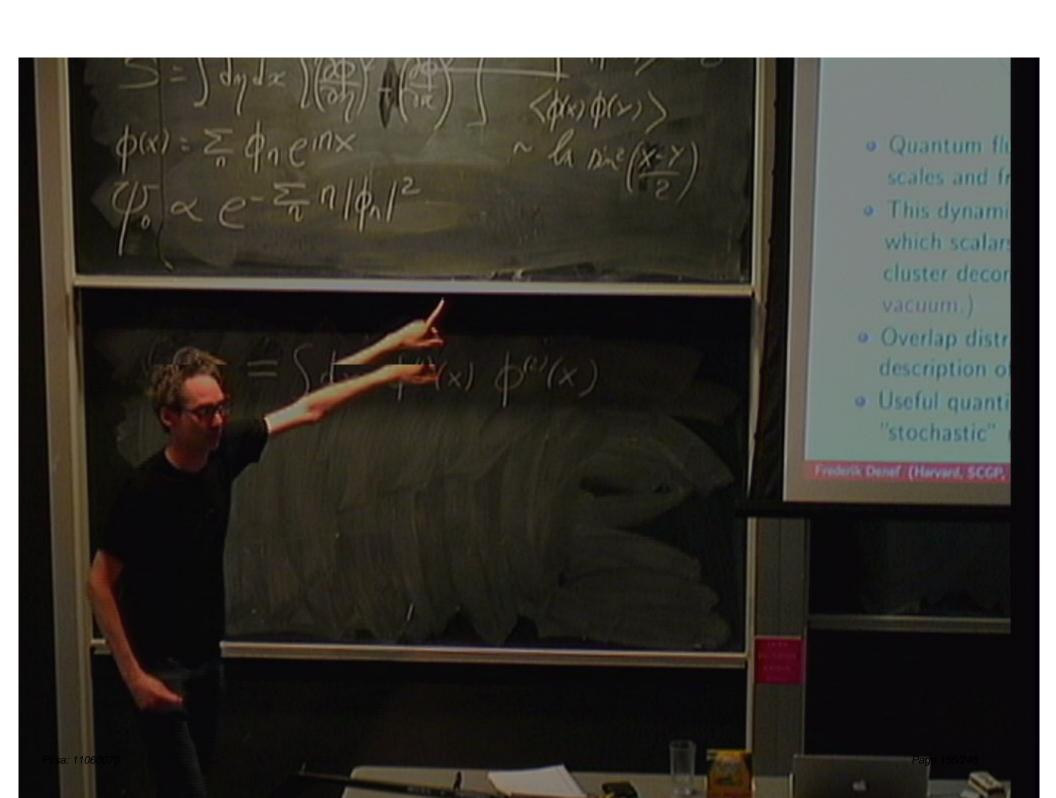


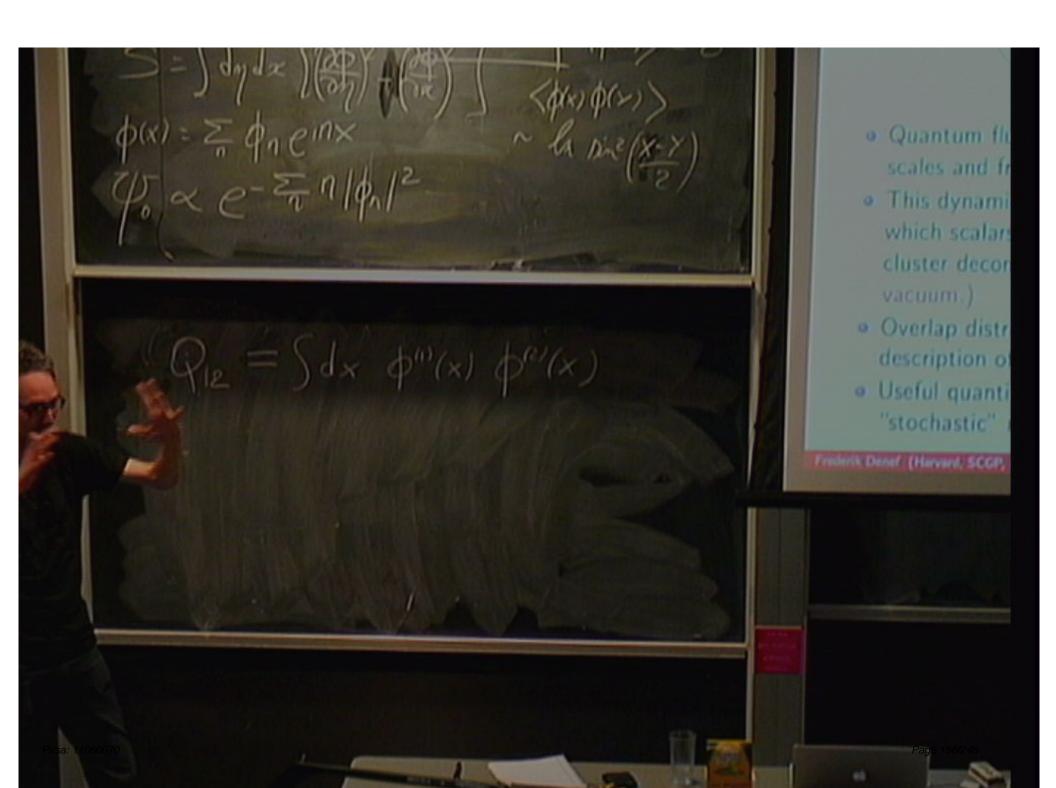


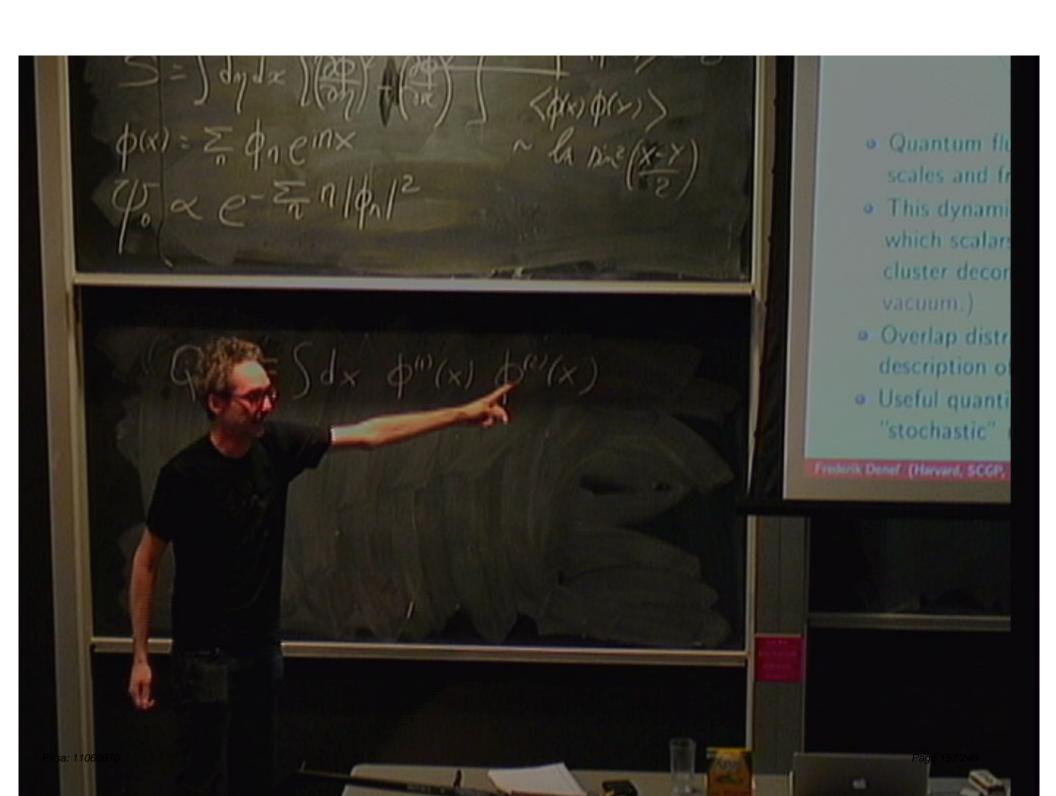


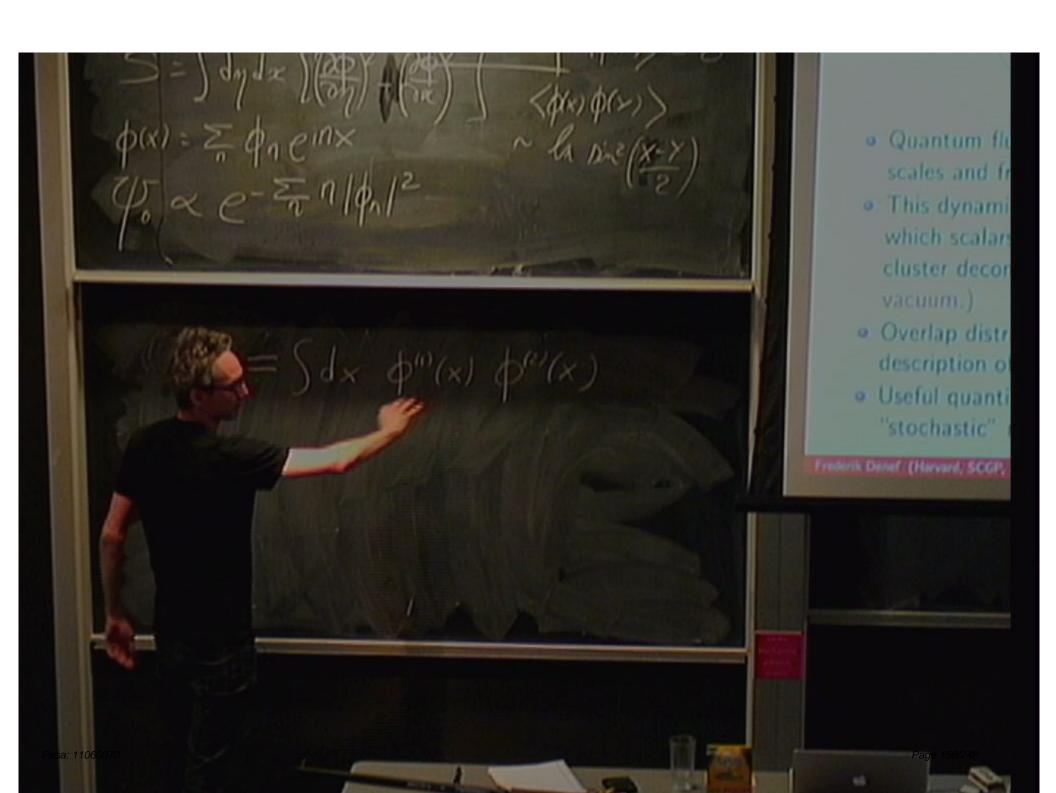


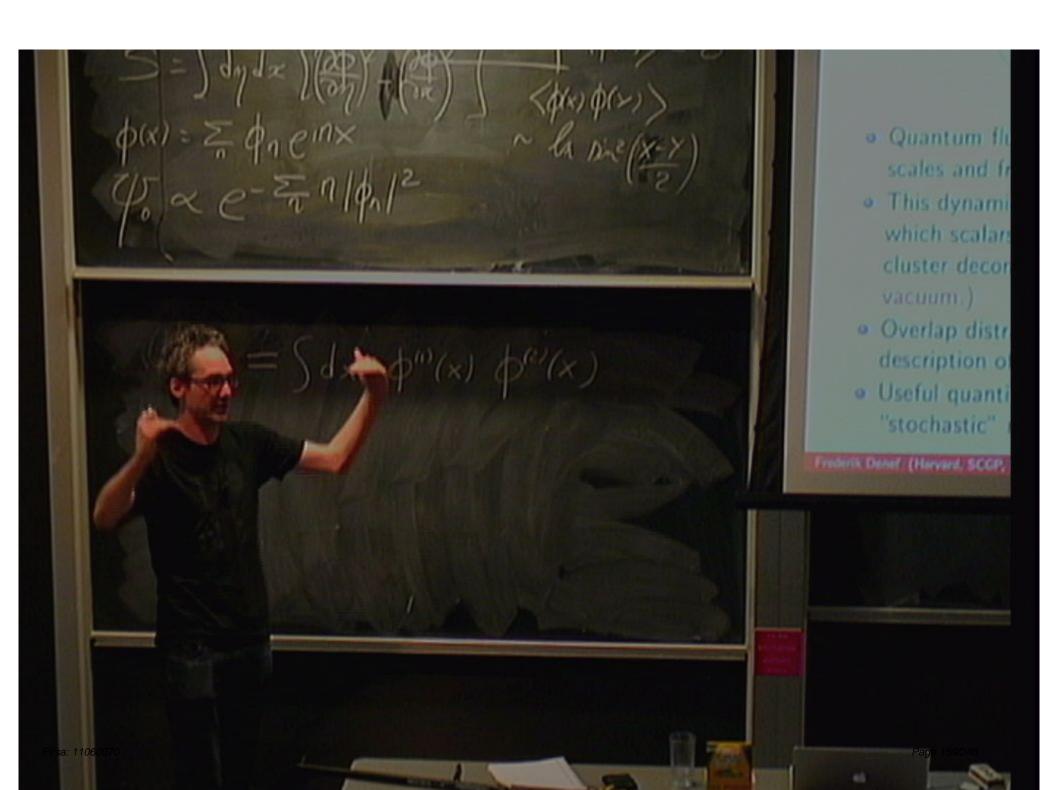


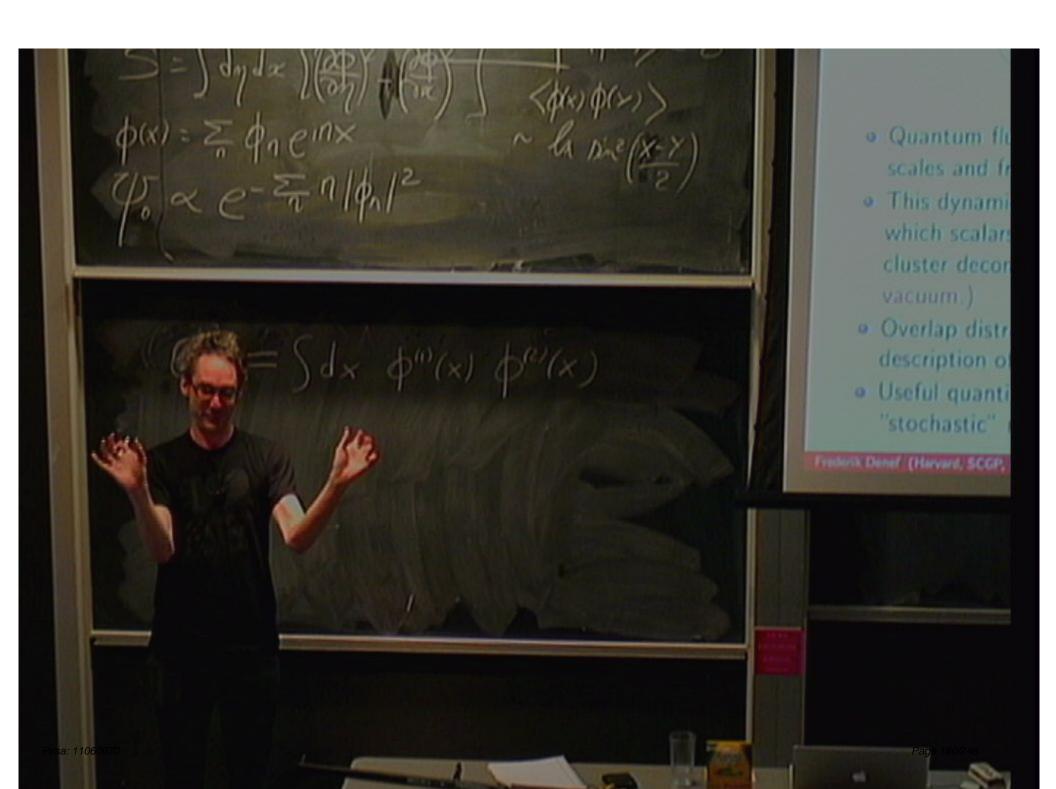


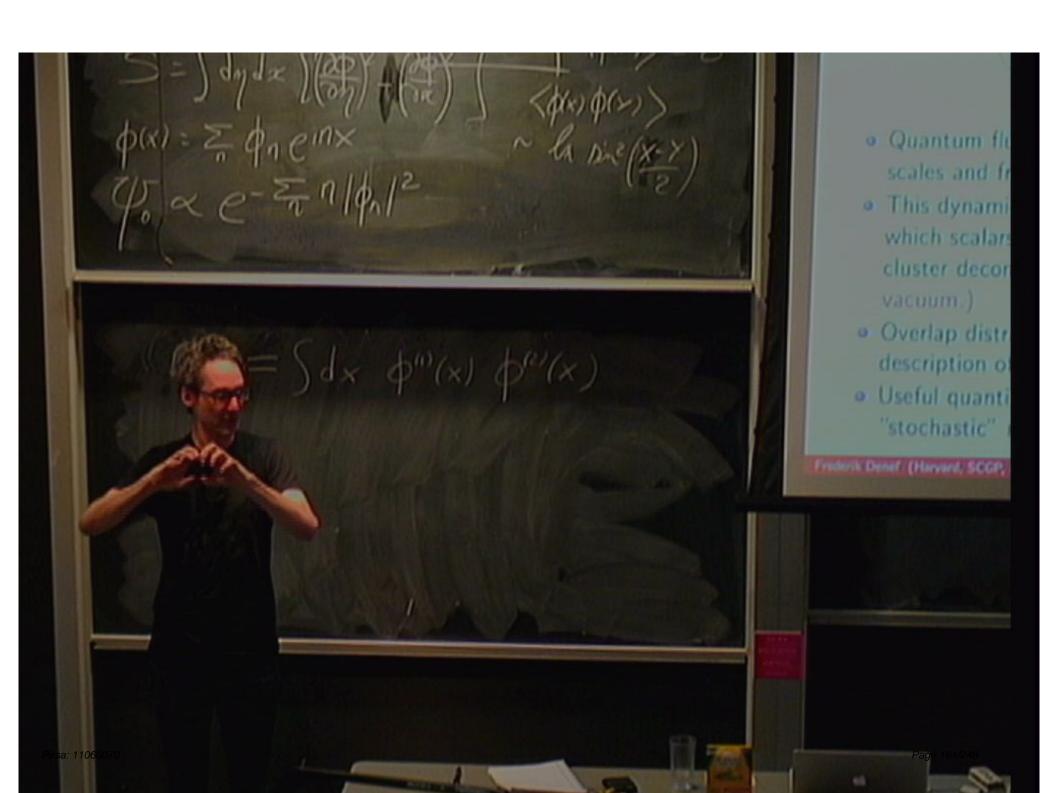


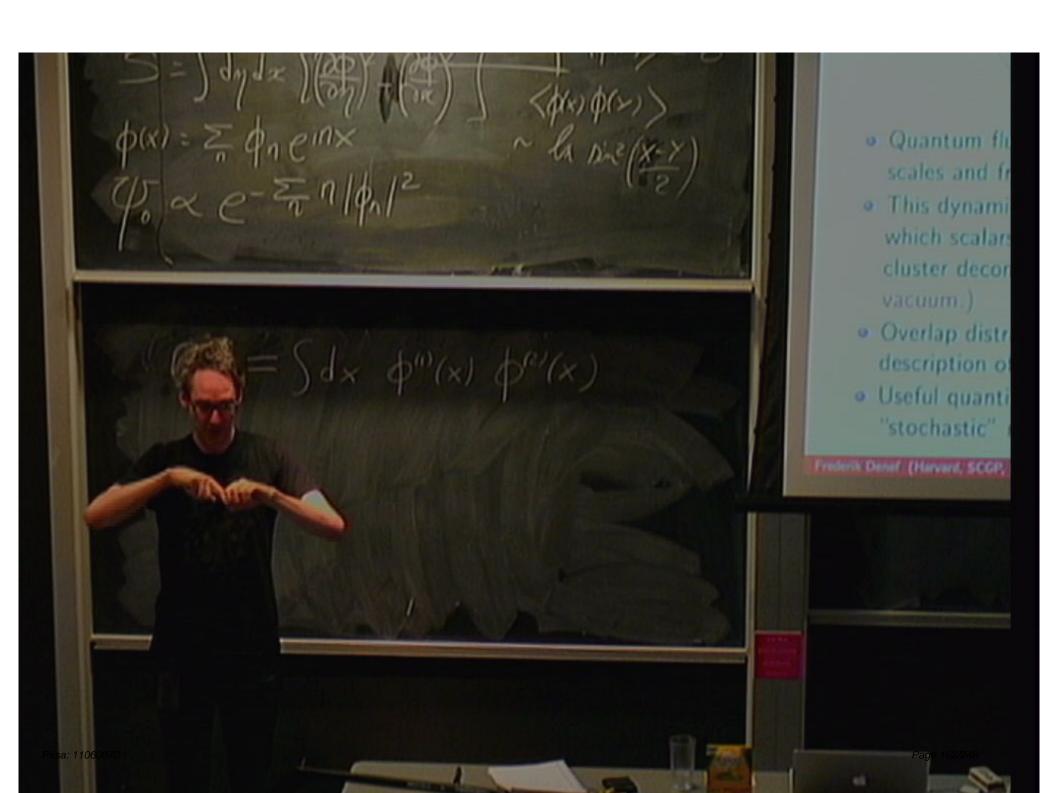




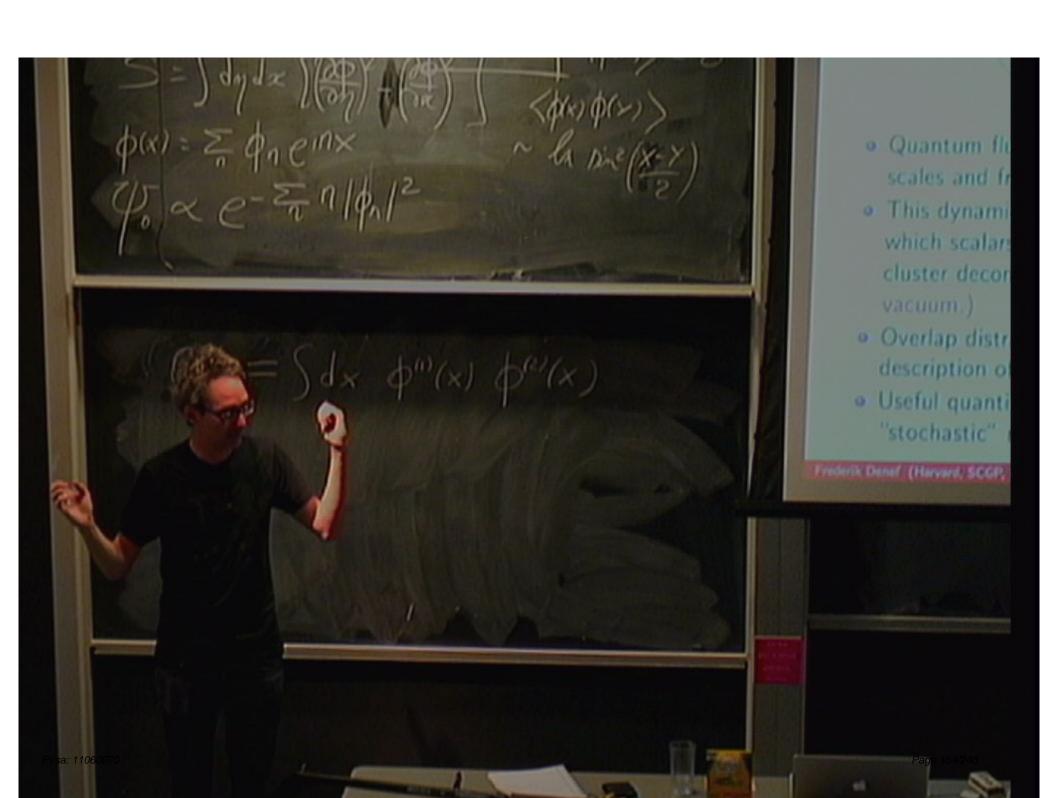


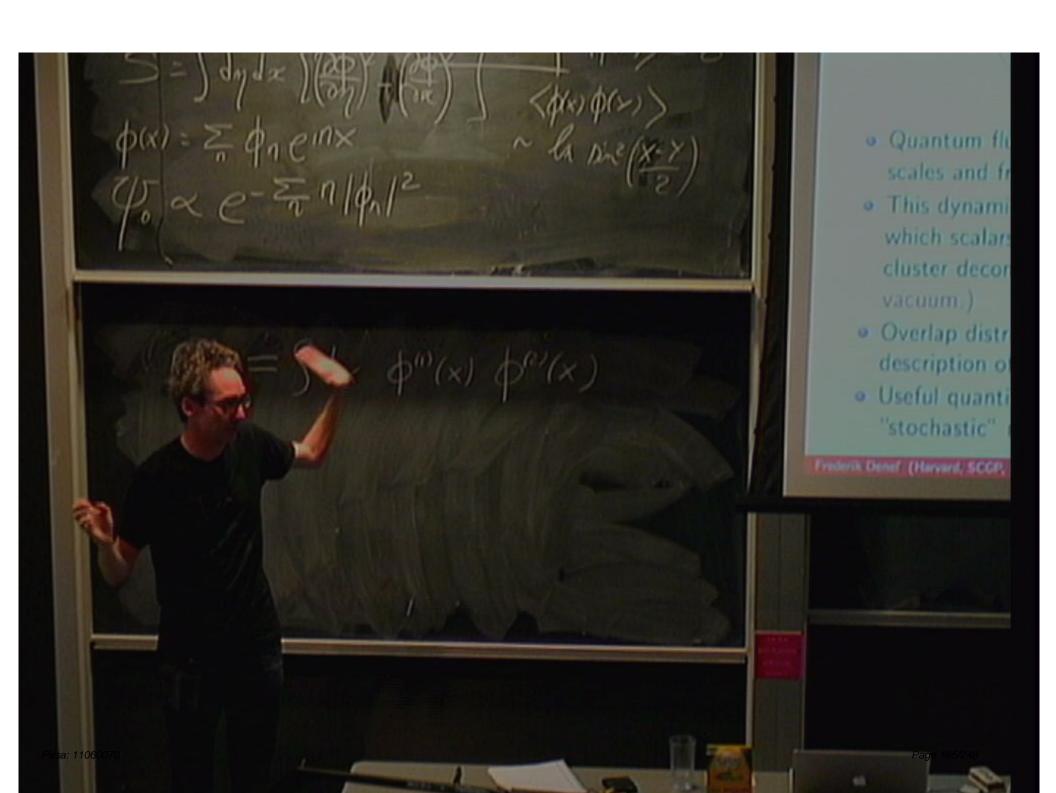


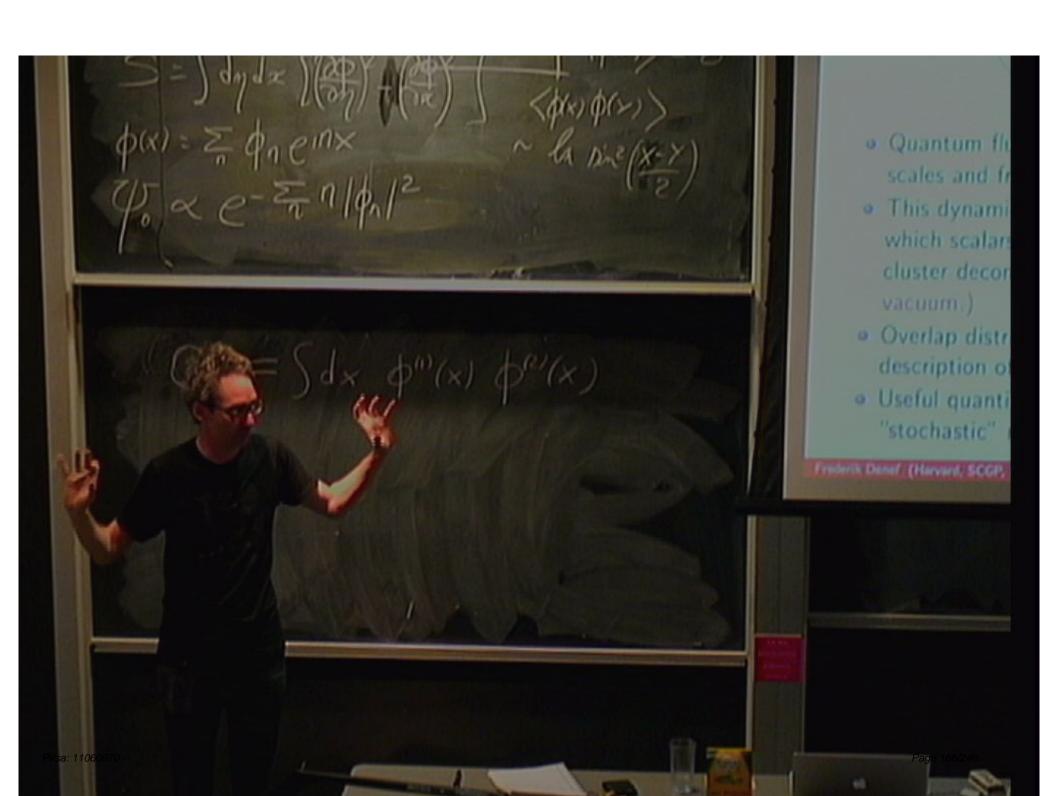


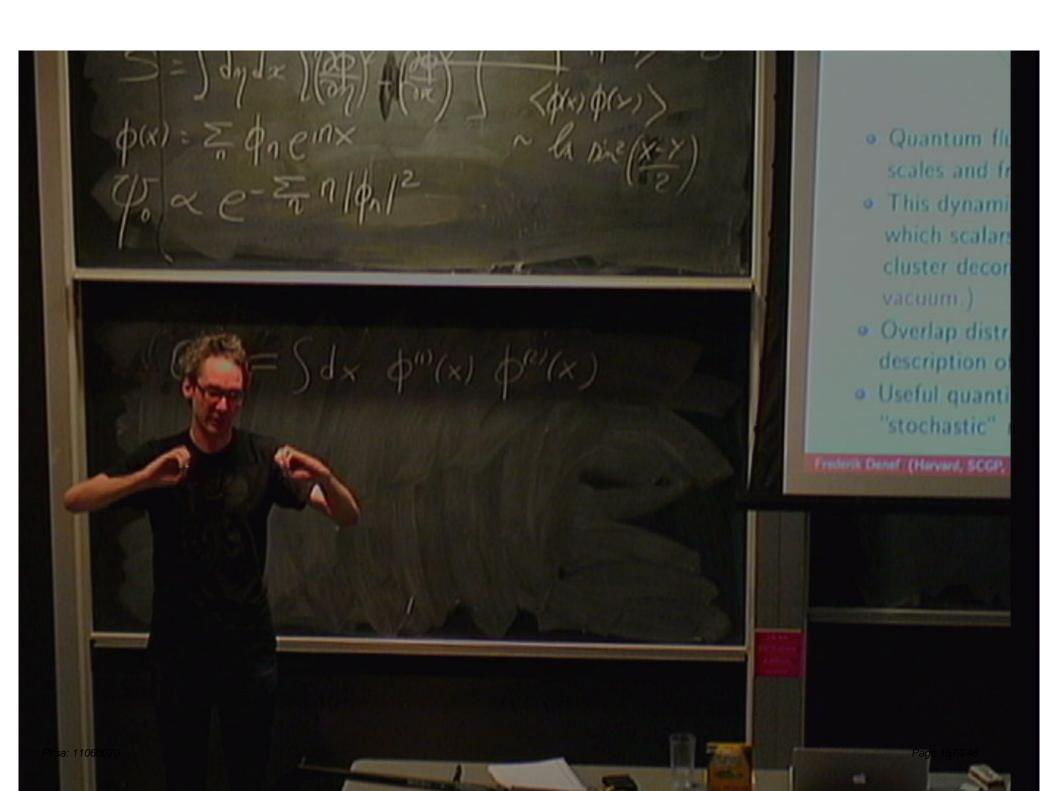


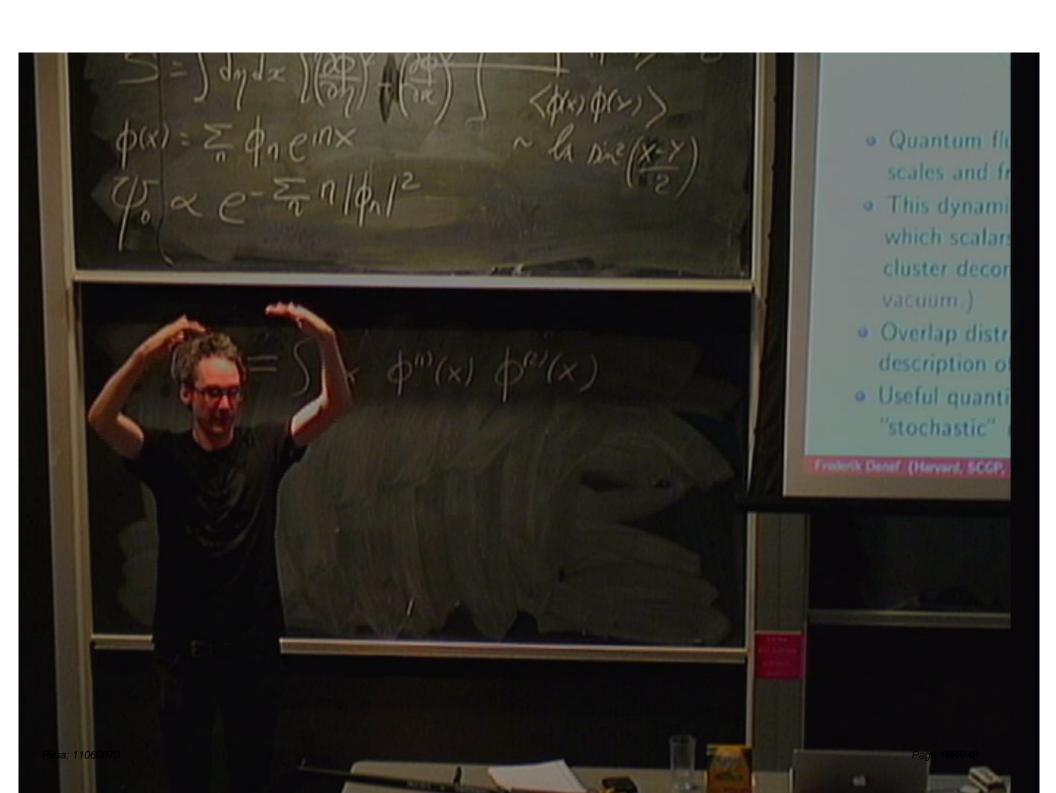


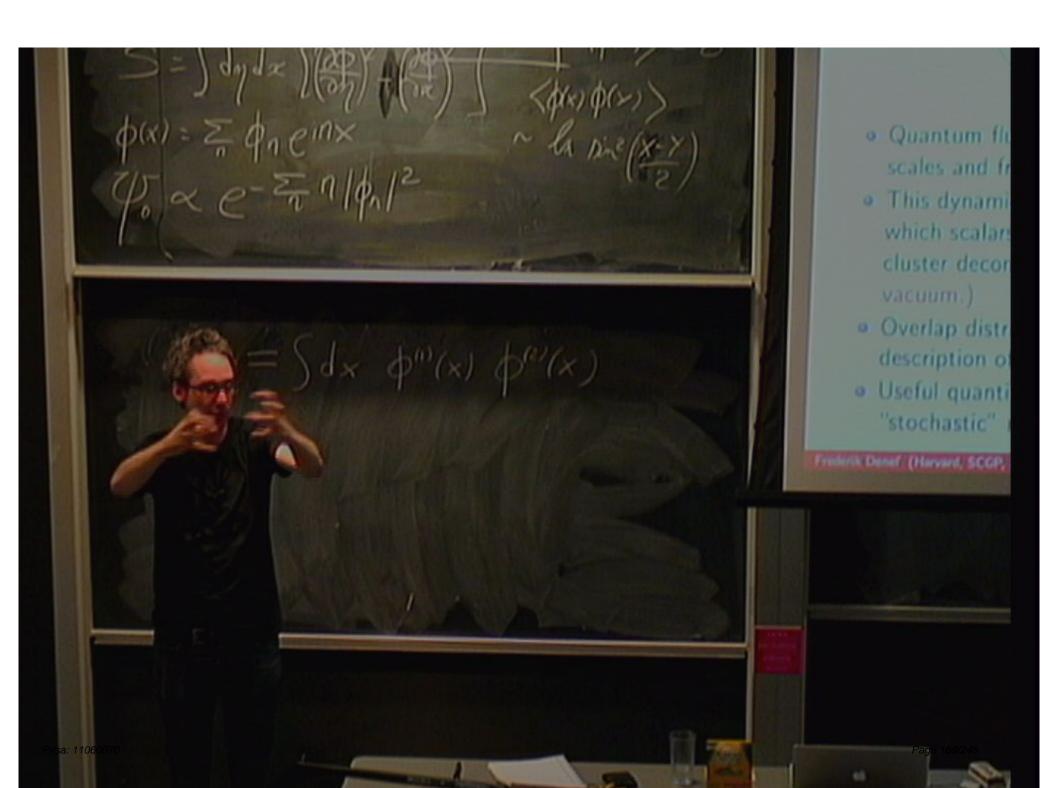


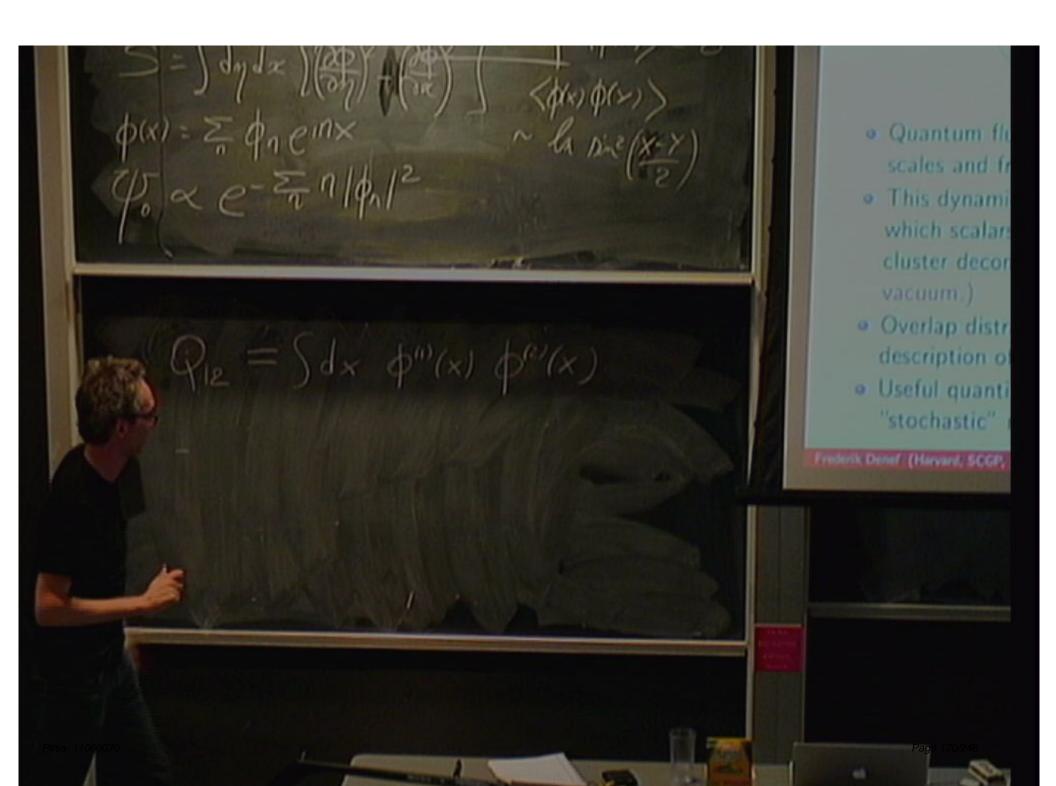


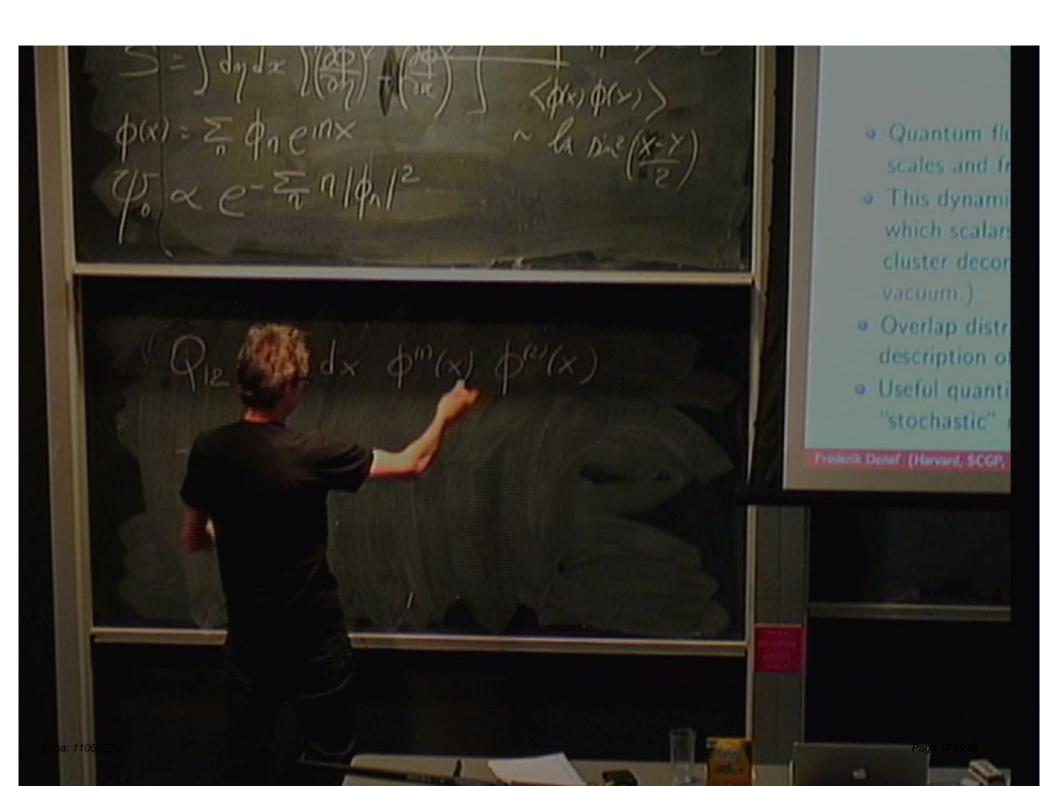


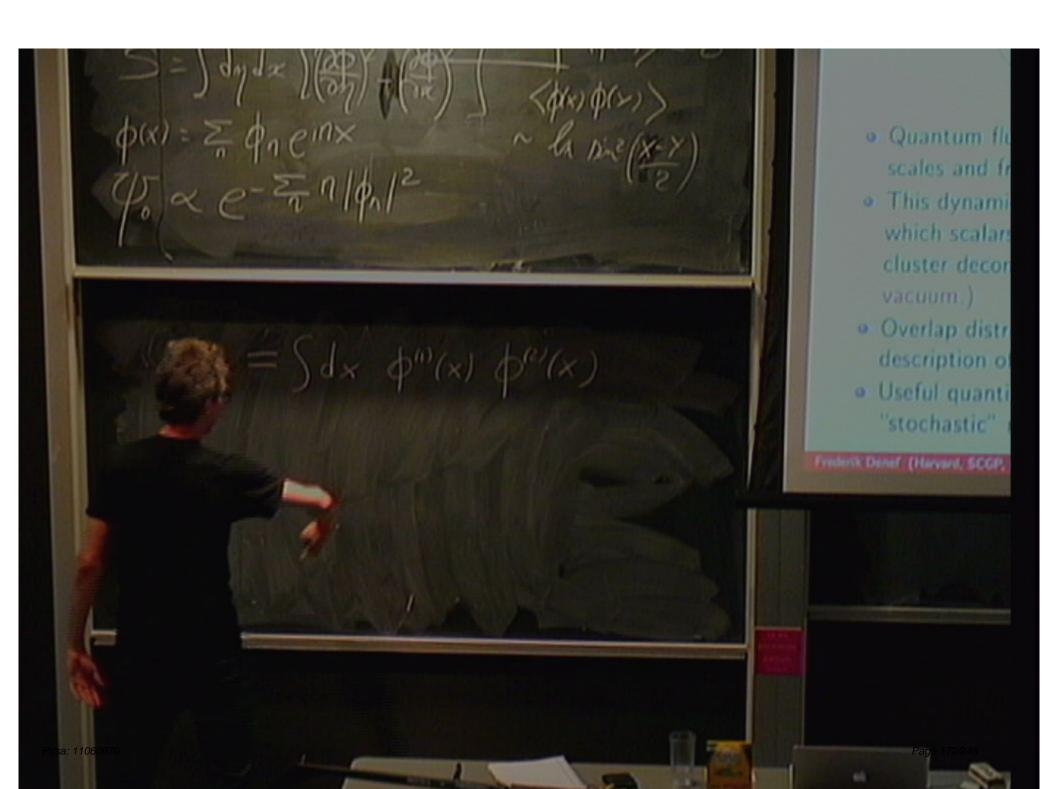


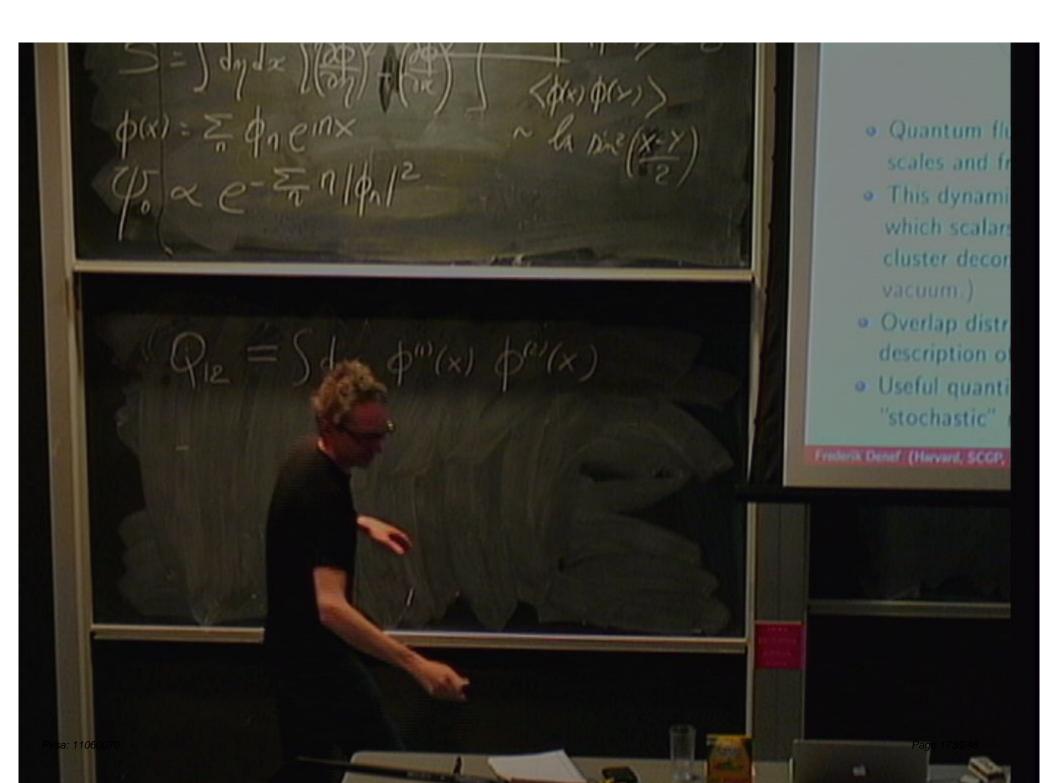


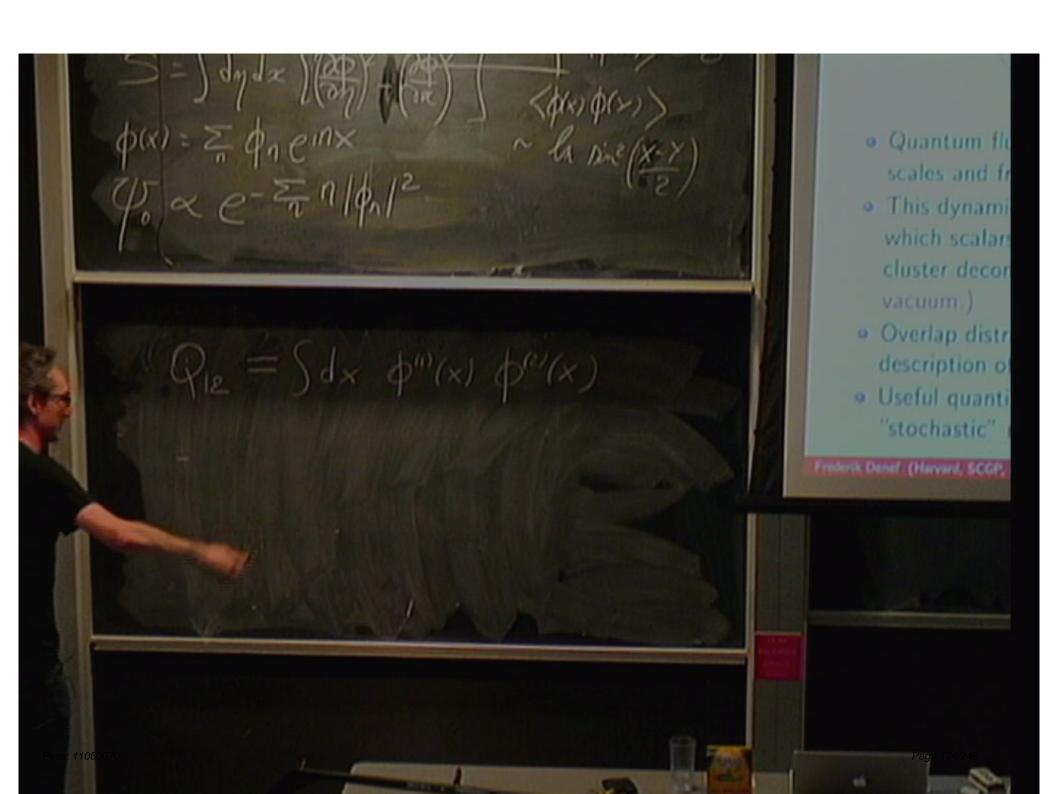


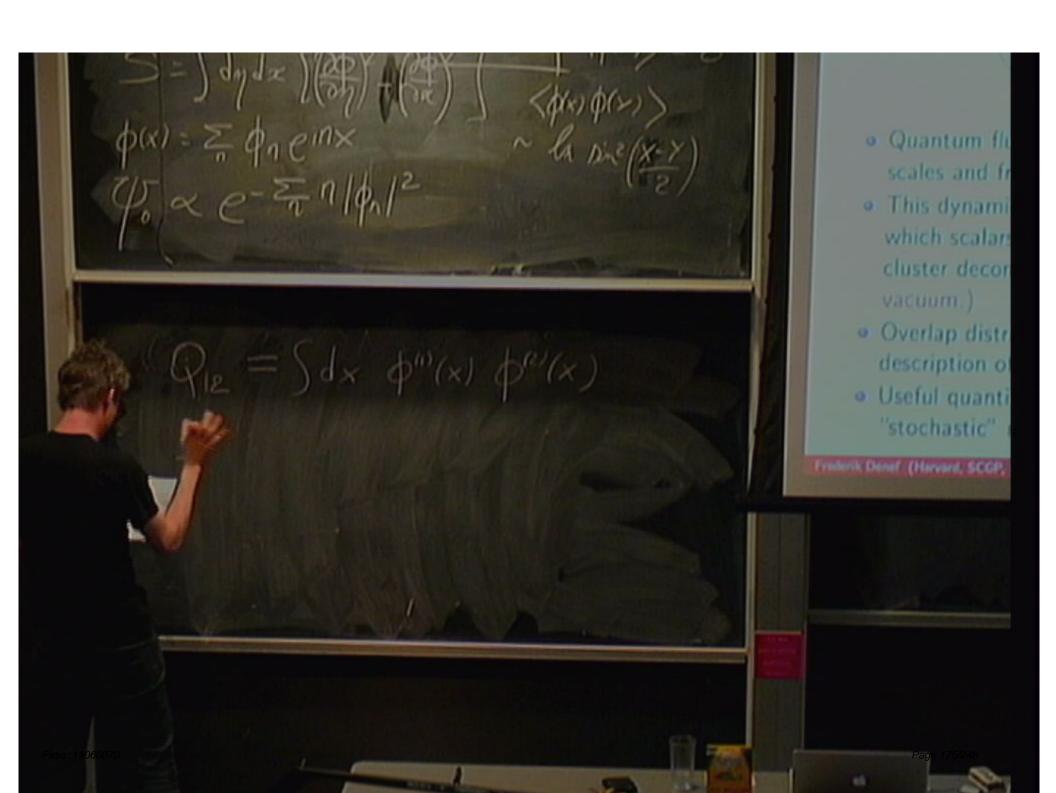


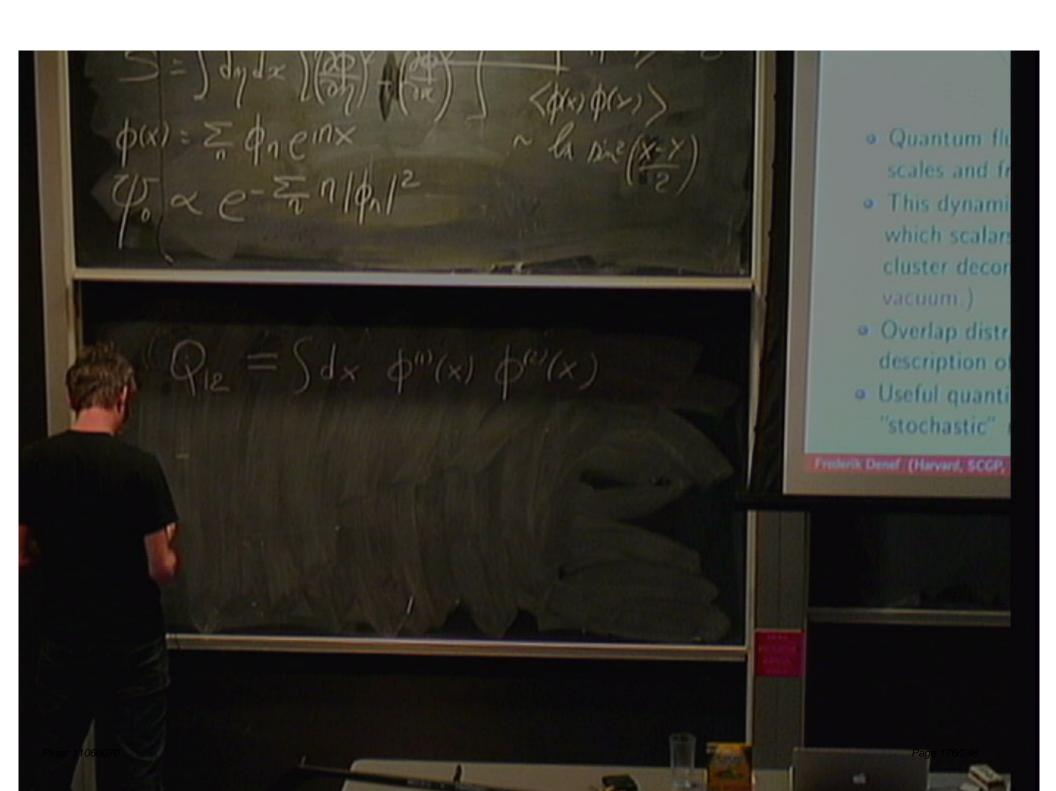


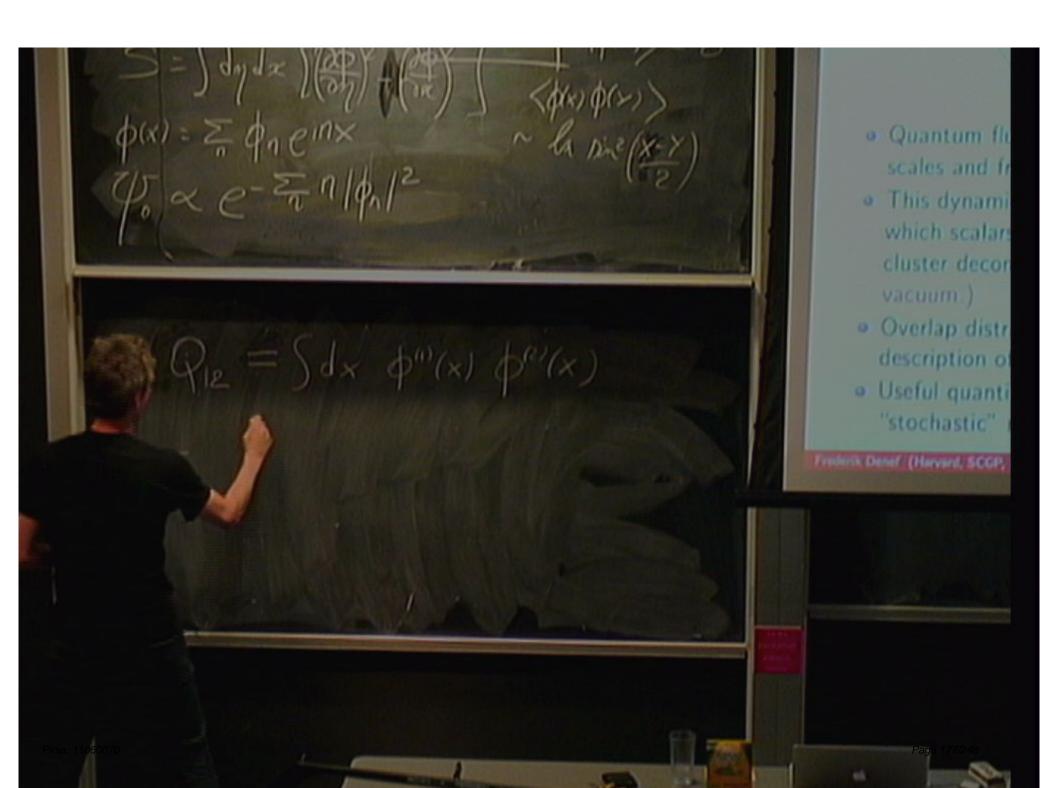


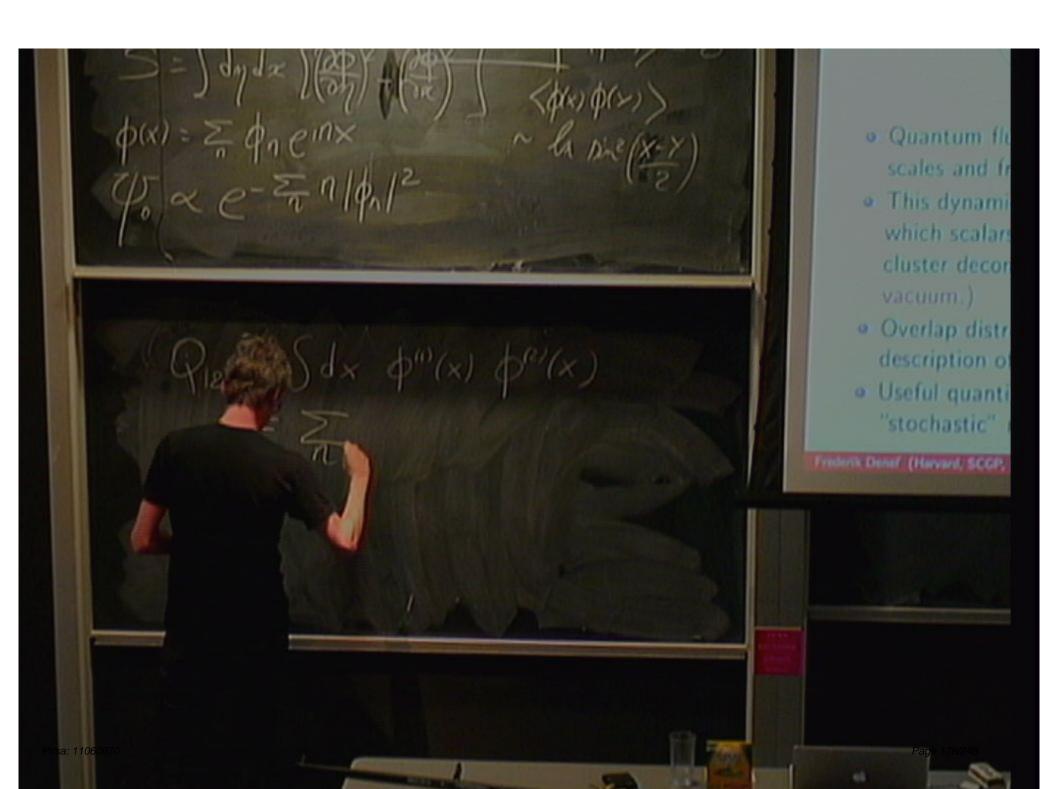


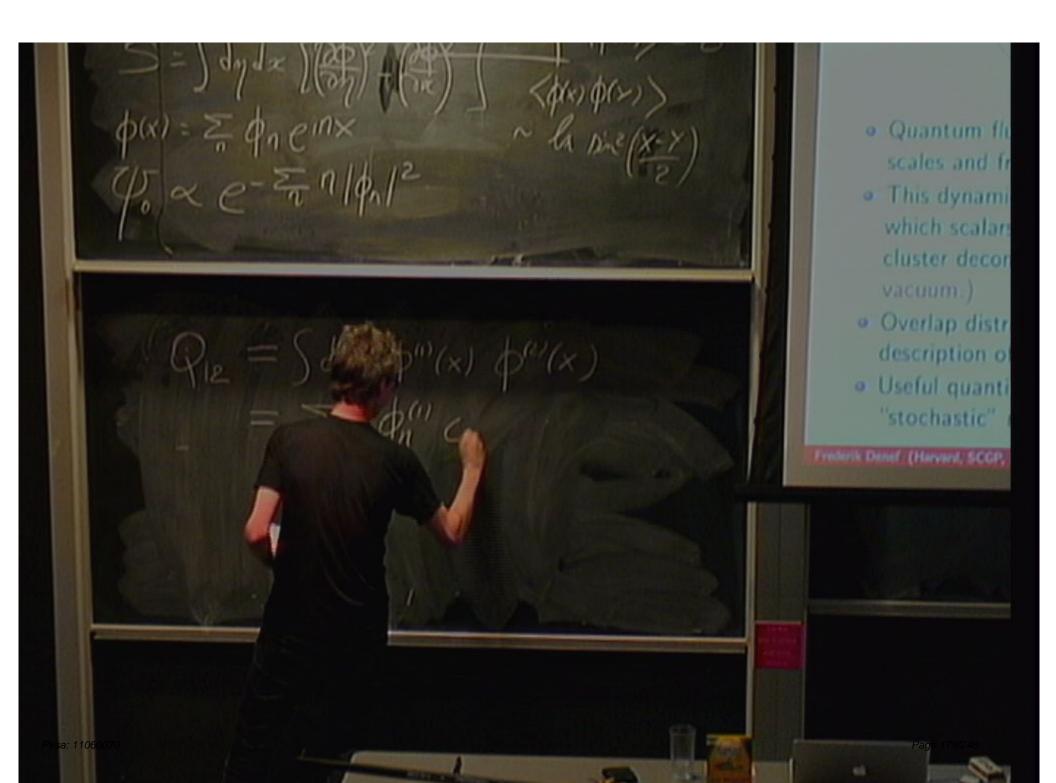


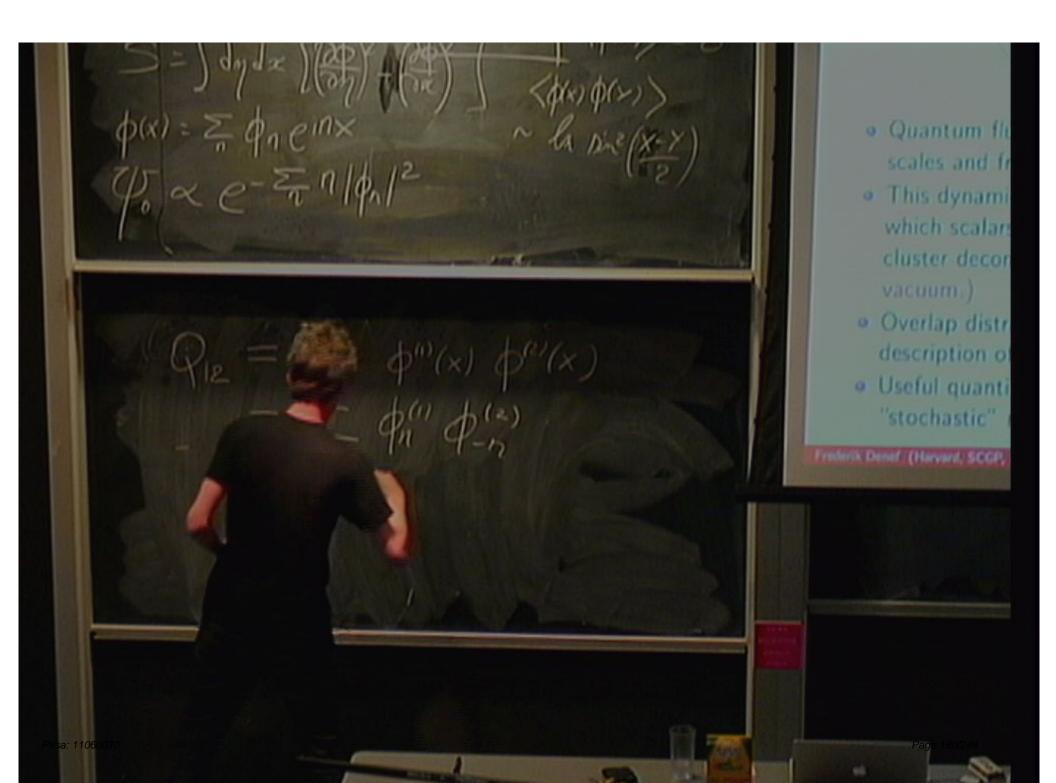


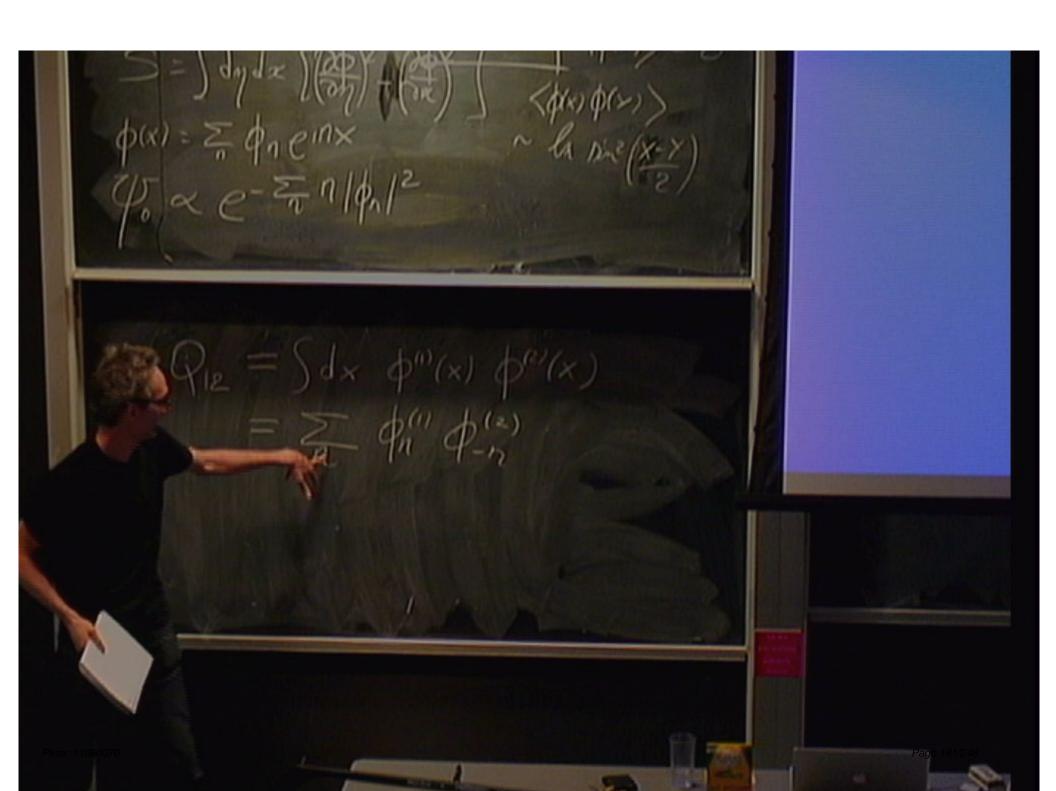


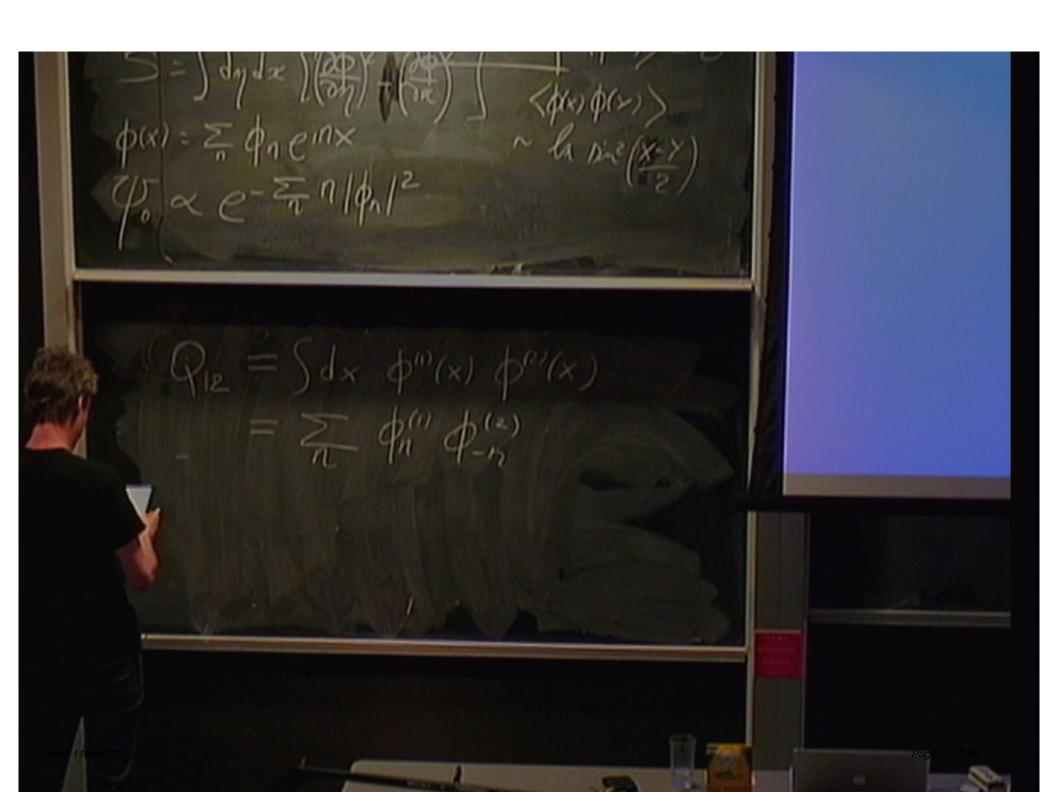


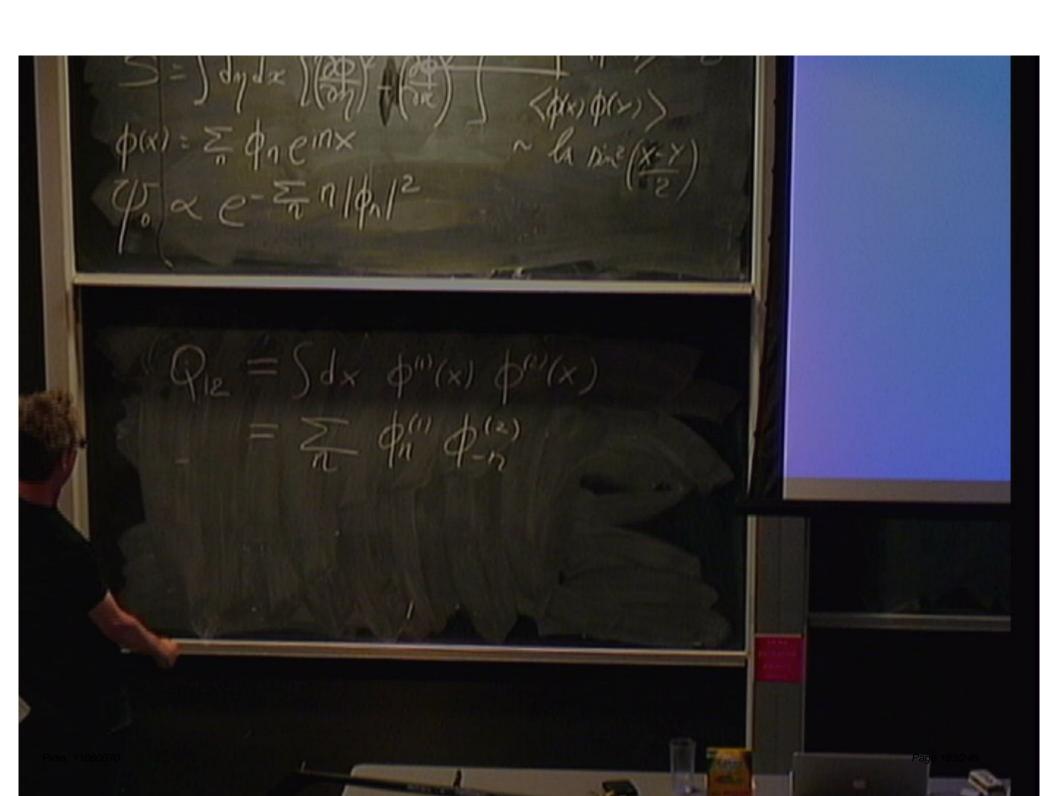


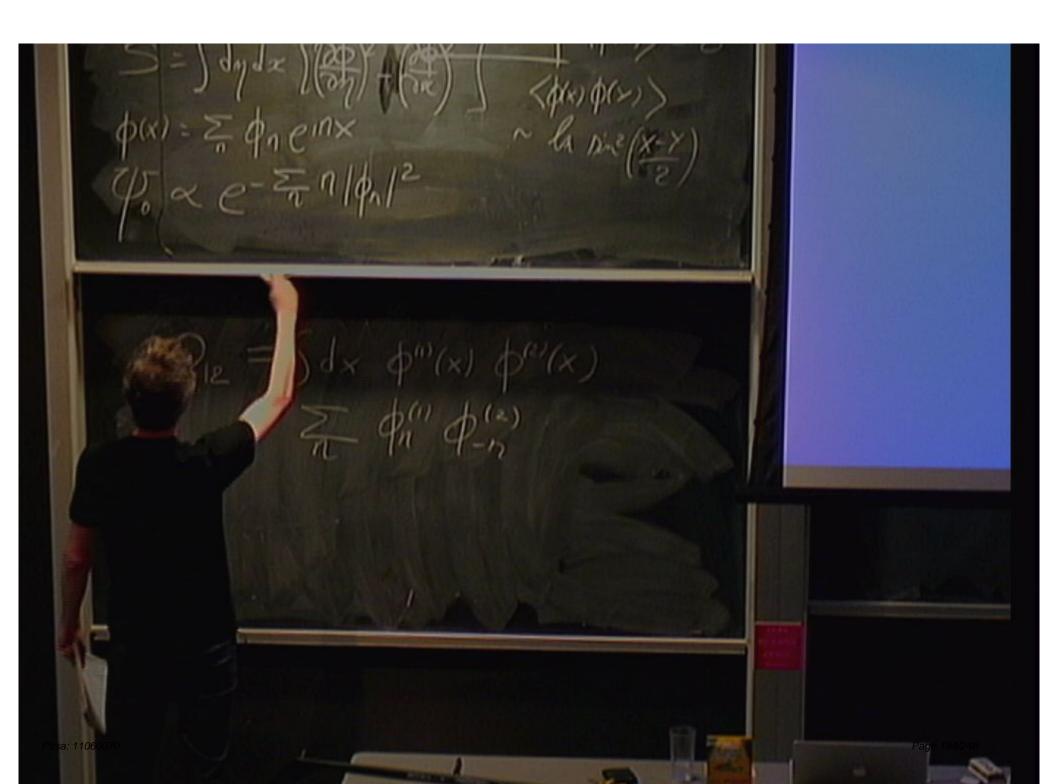


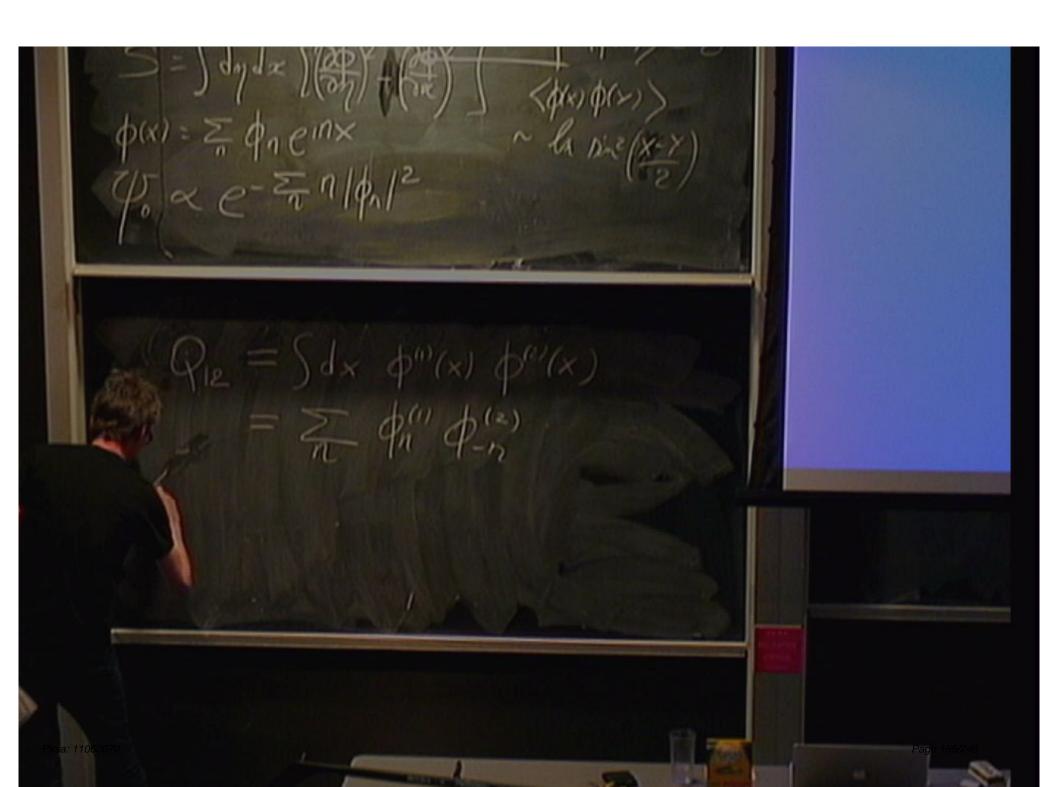


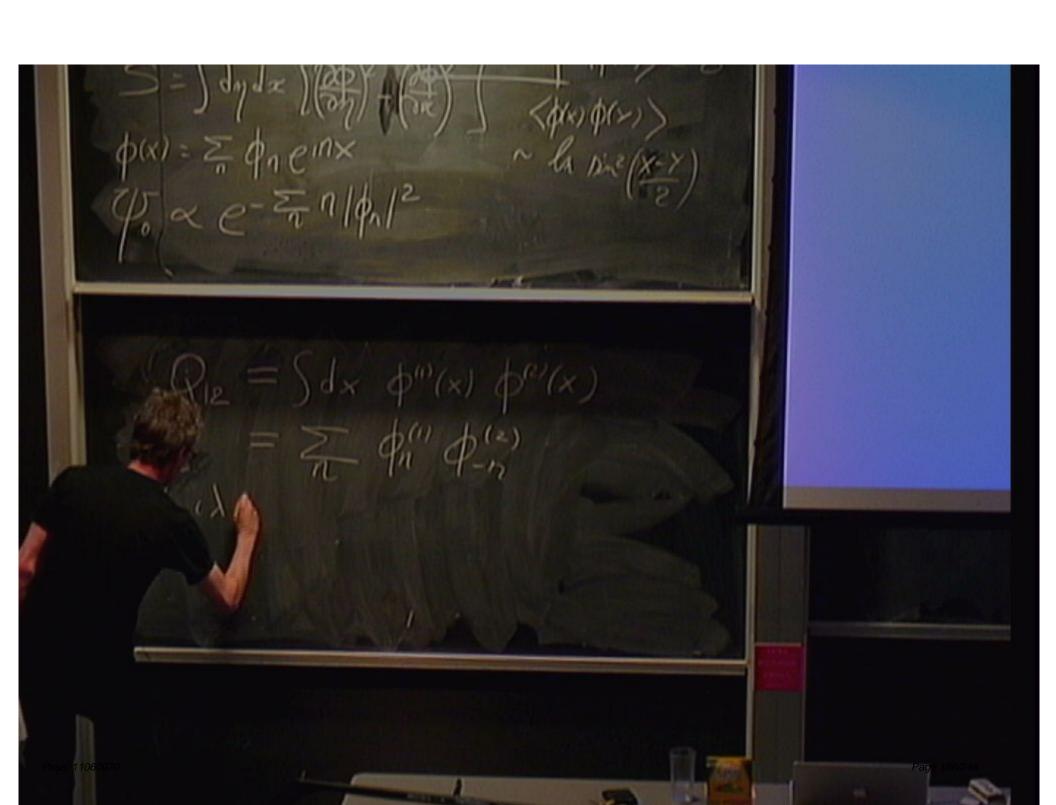










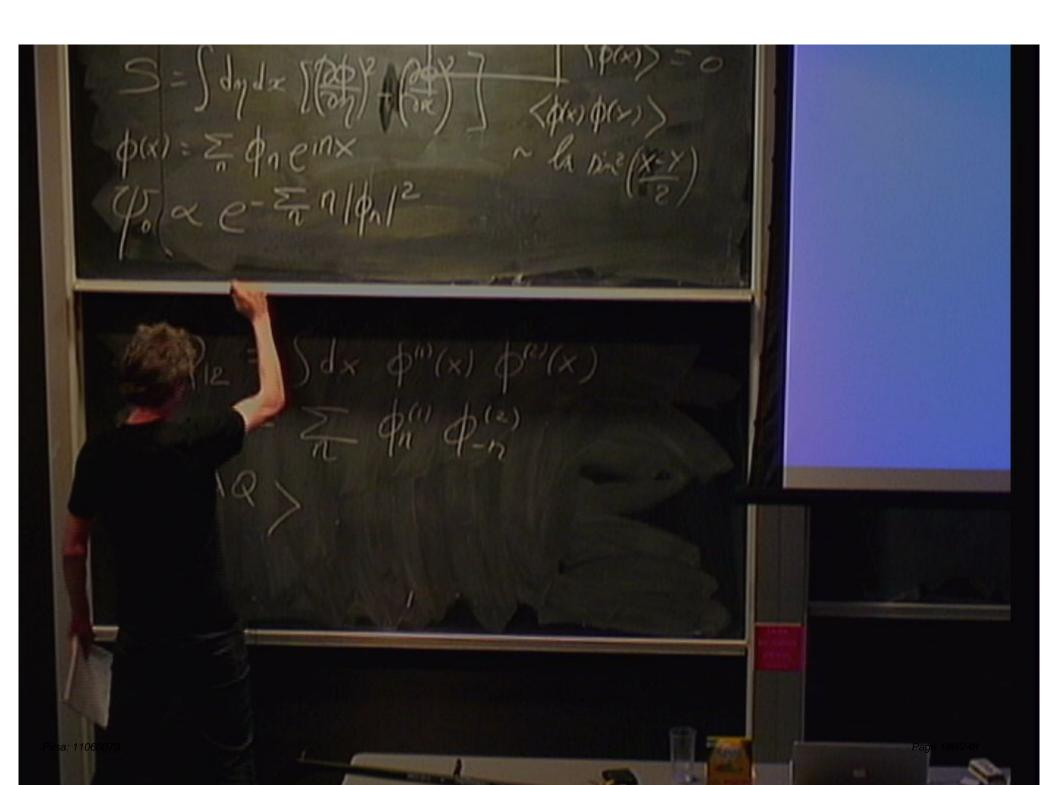


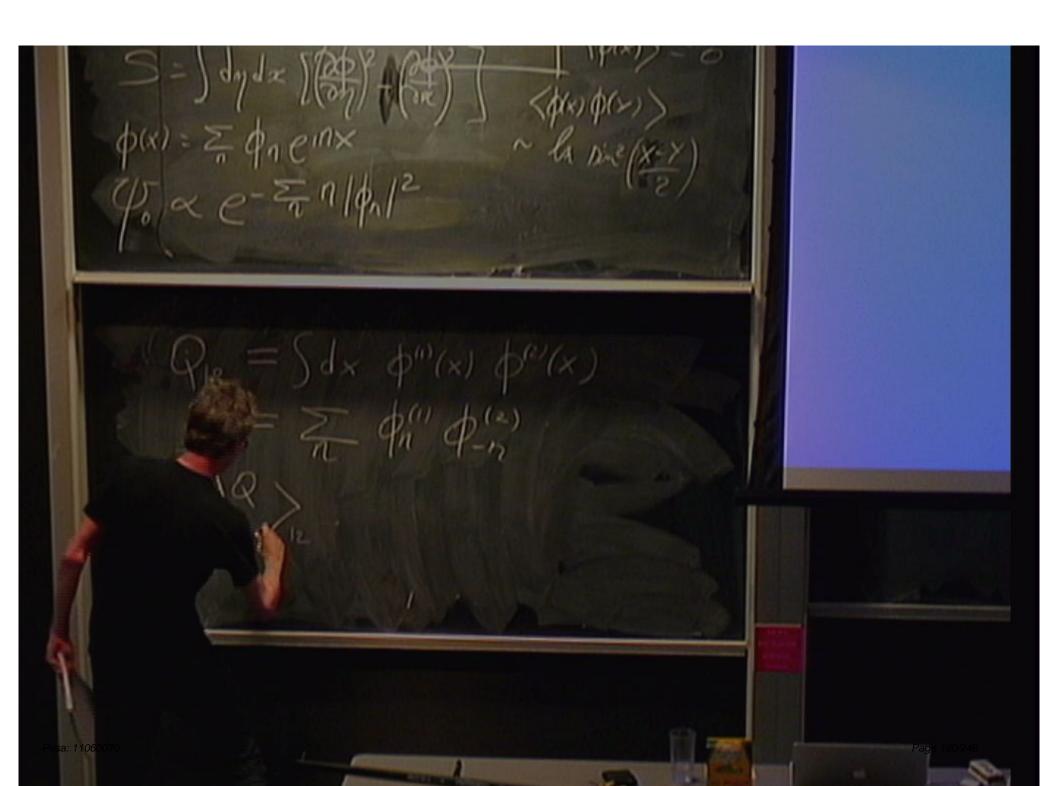
~ la 12= (X-X) φ(x)= = φnemx

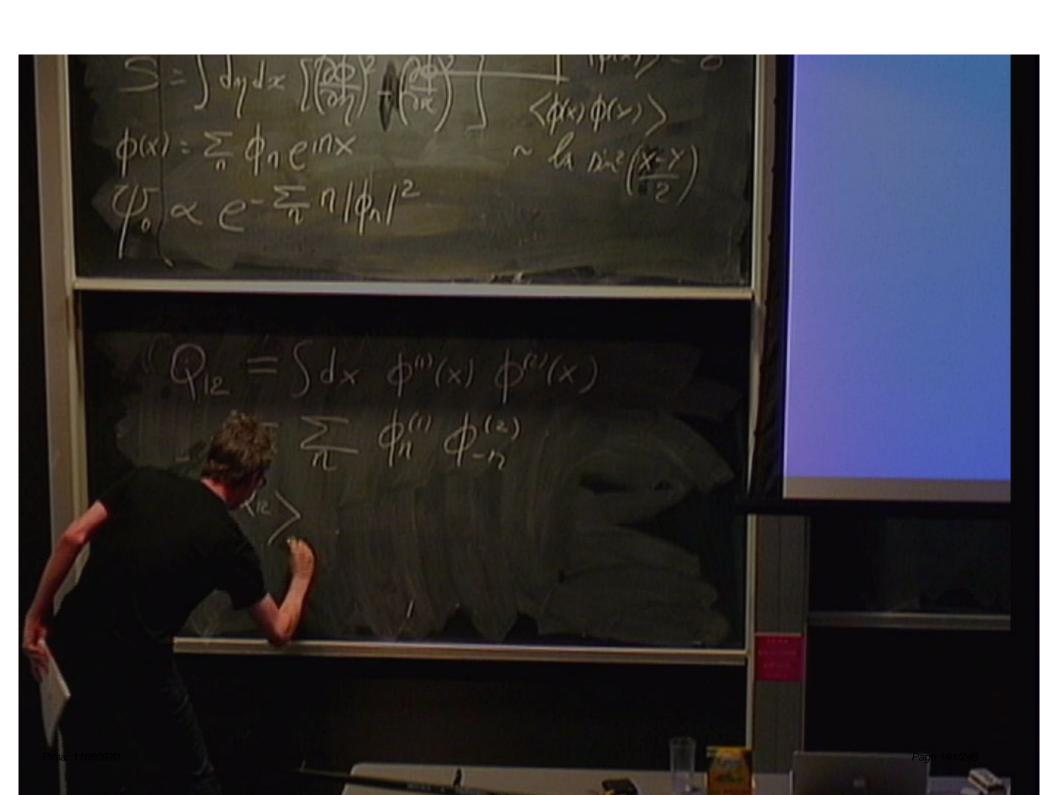
(γ) ~ = = η 1 φη 12 = = p(1) p(2)

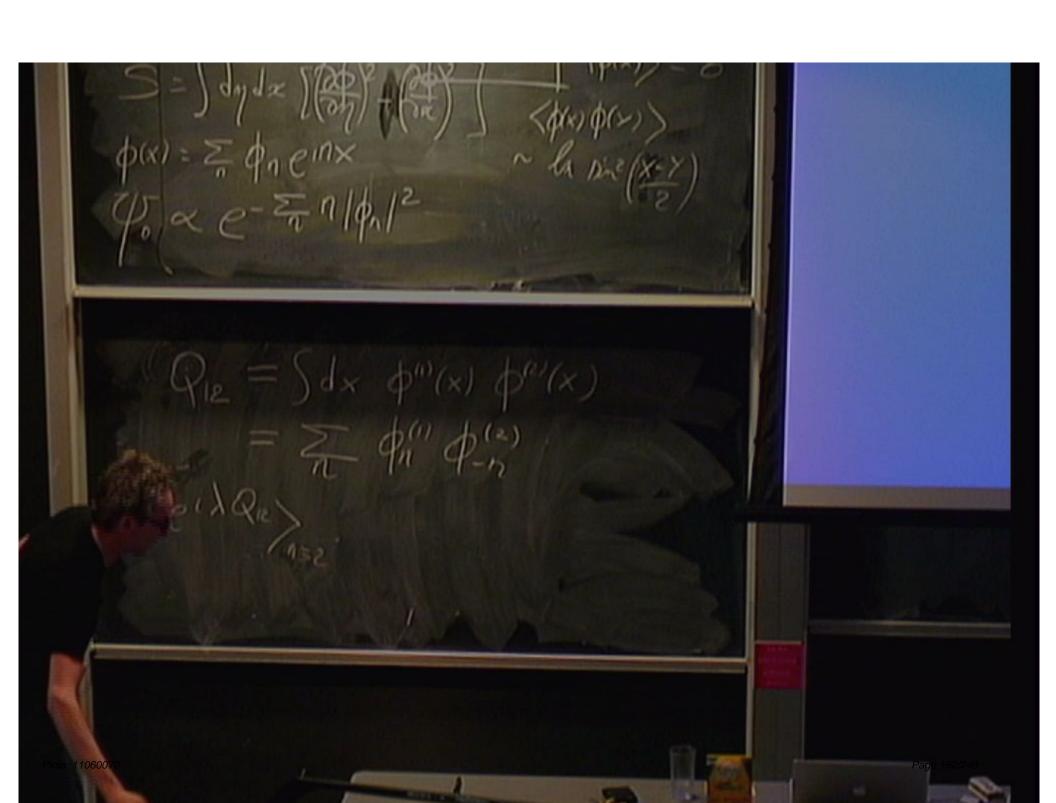
~ pro(x-x) φ(x)= = φnemx

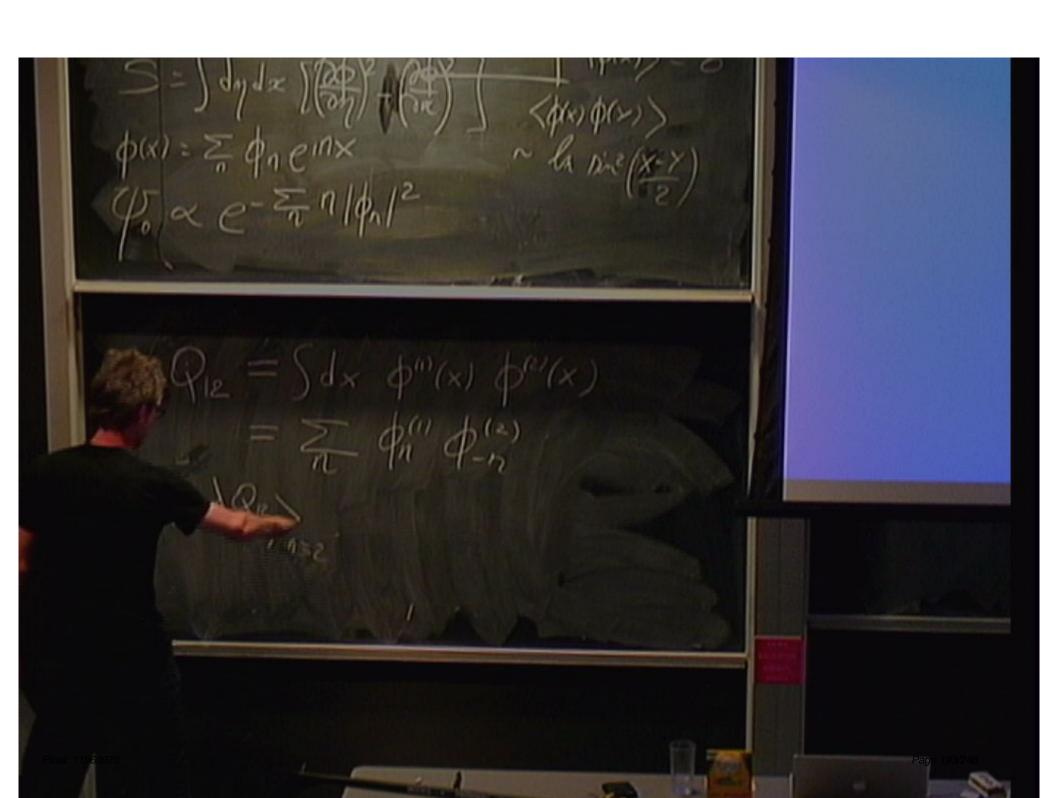
(γ) ~ = = η |φη|² = Z p(1) p(2)

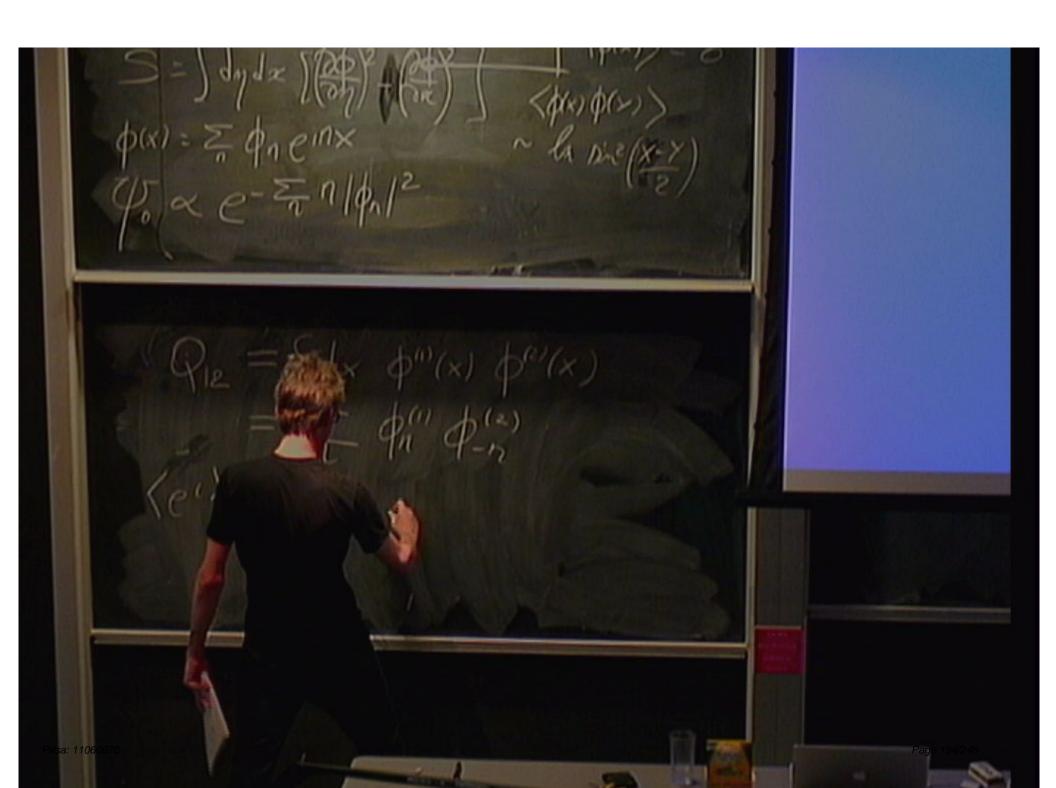


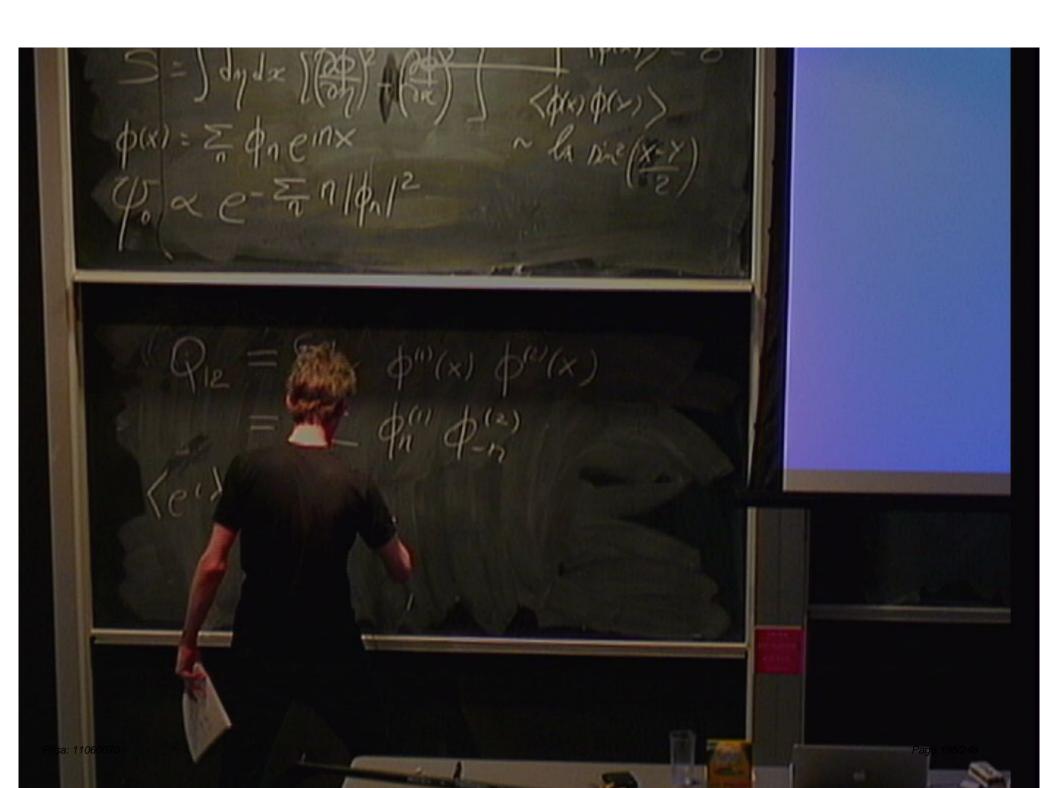


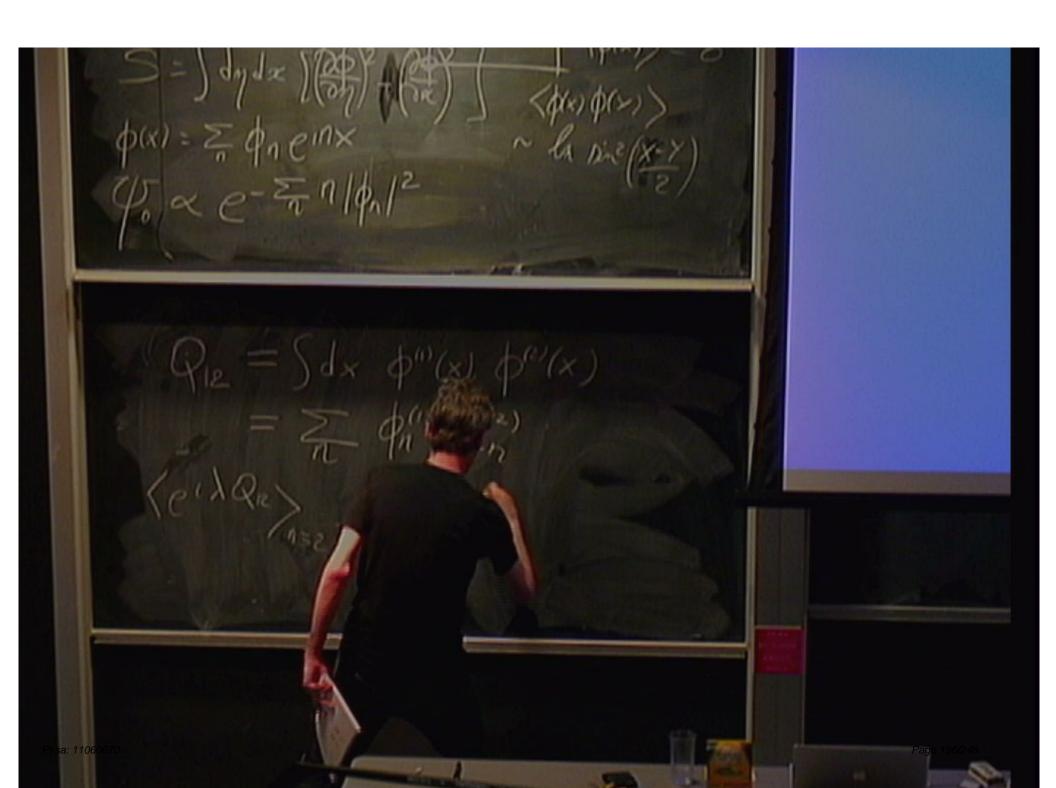


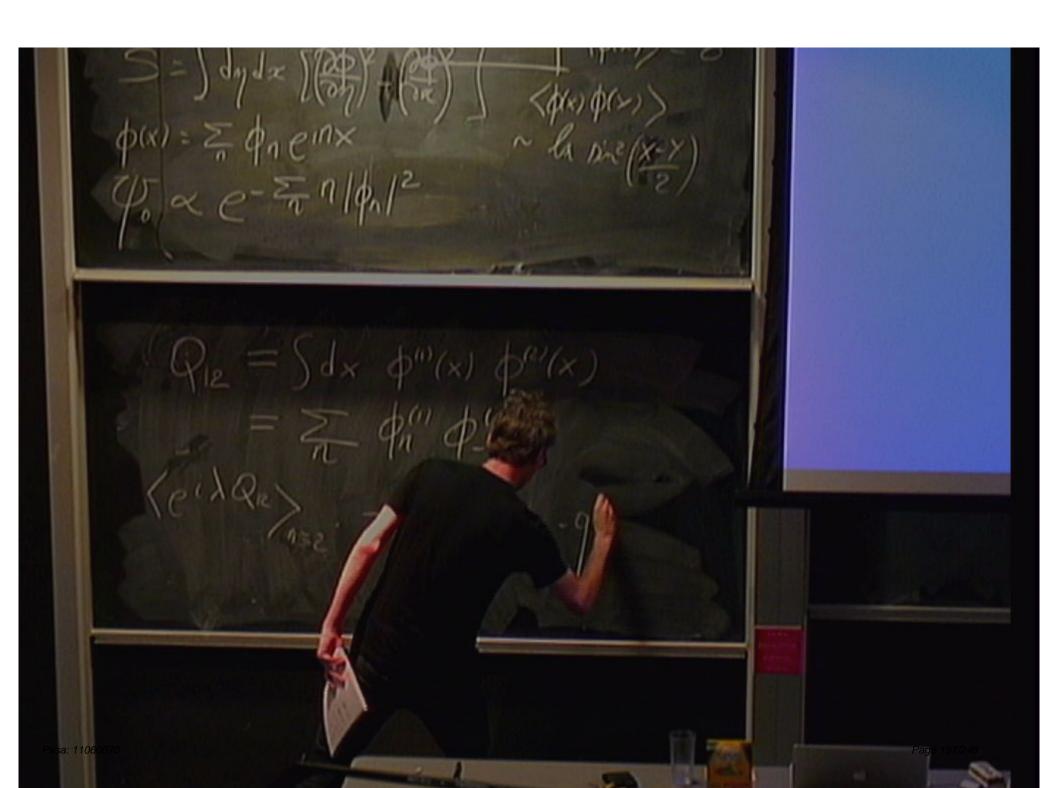


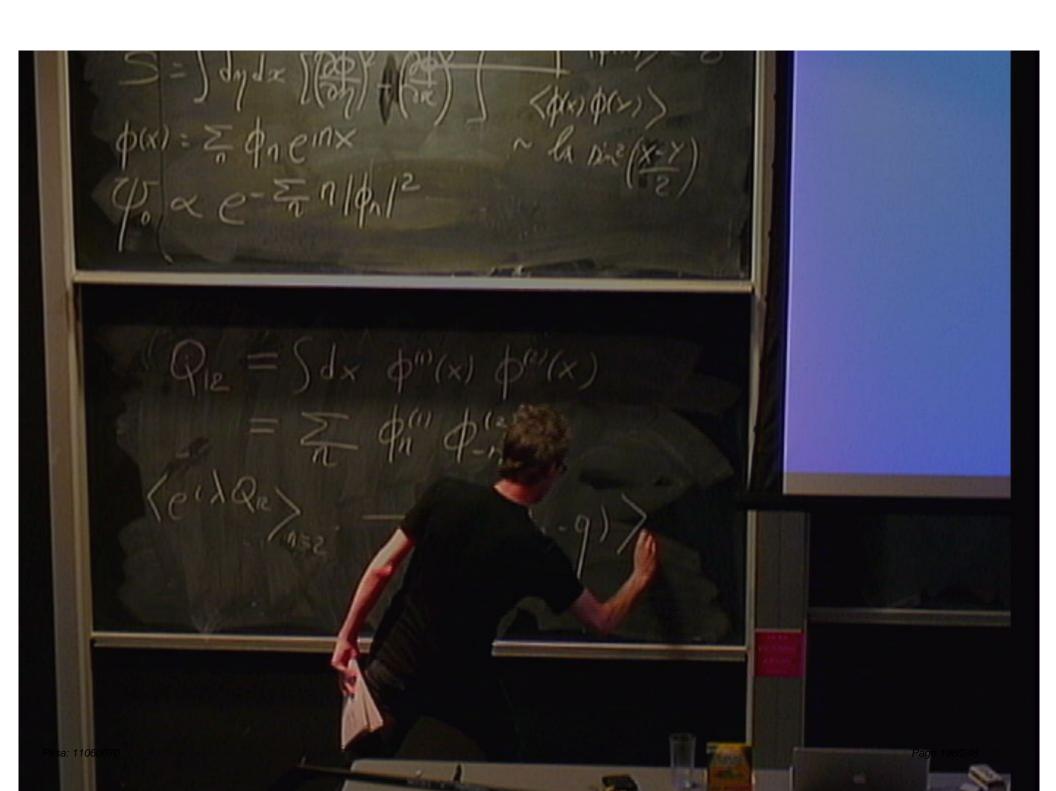






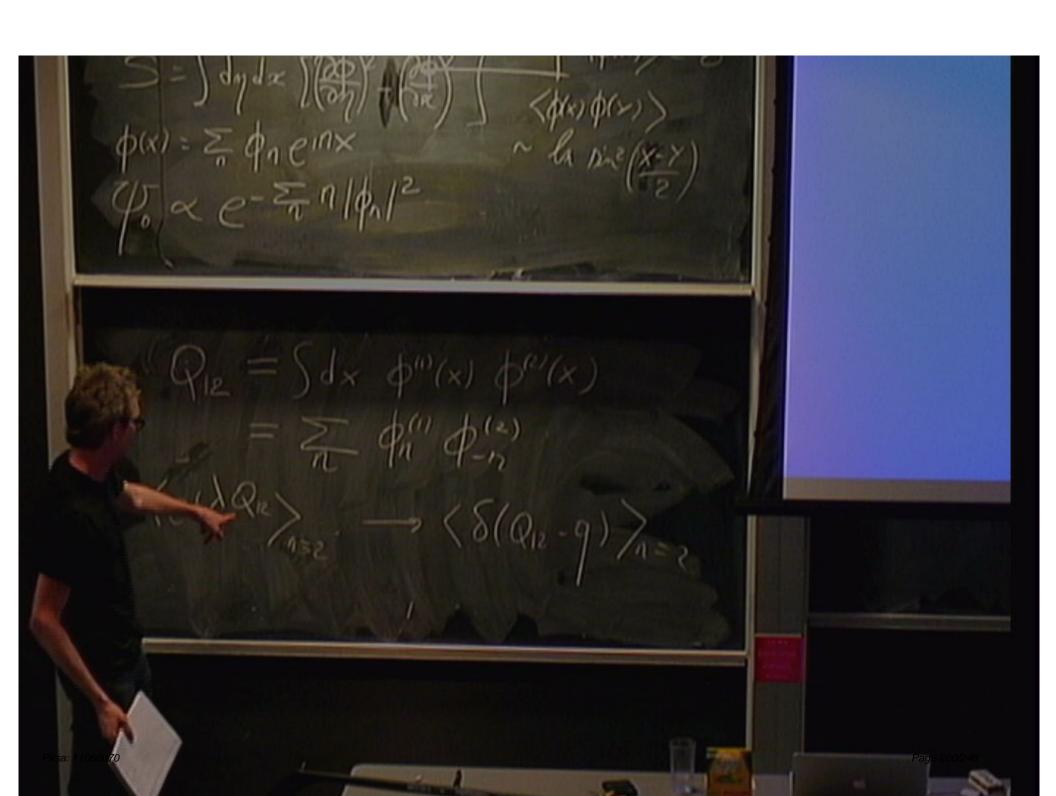


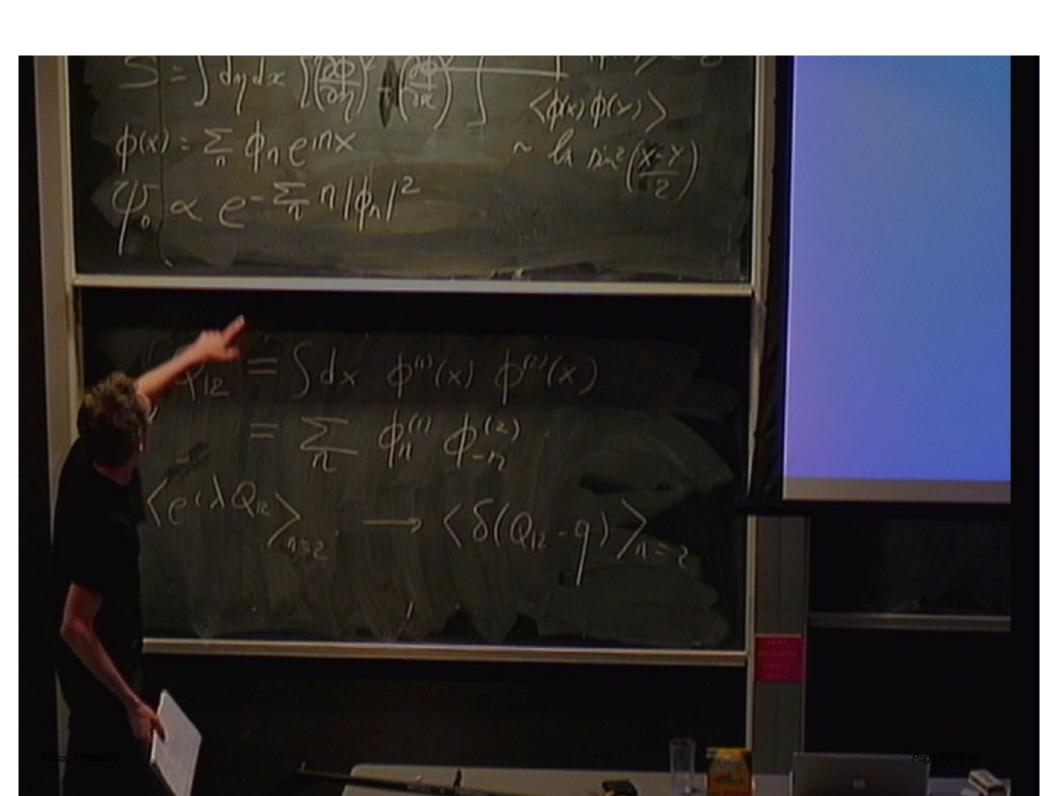


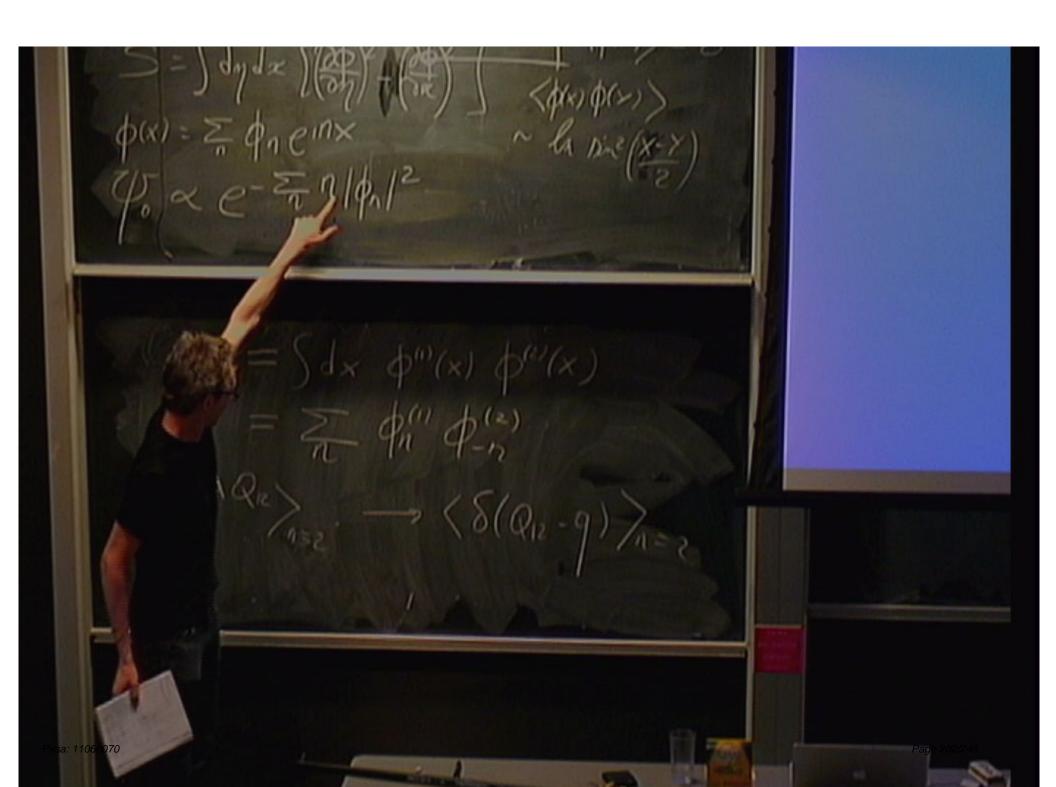


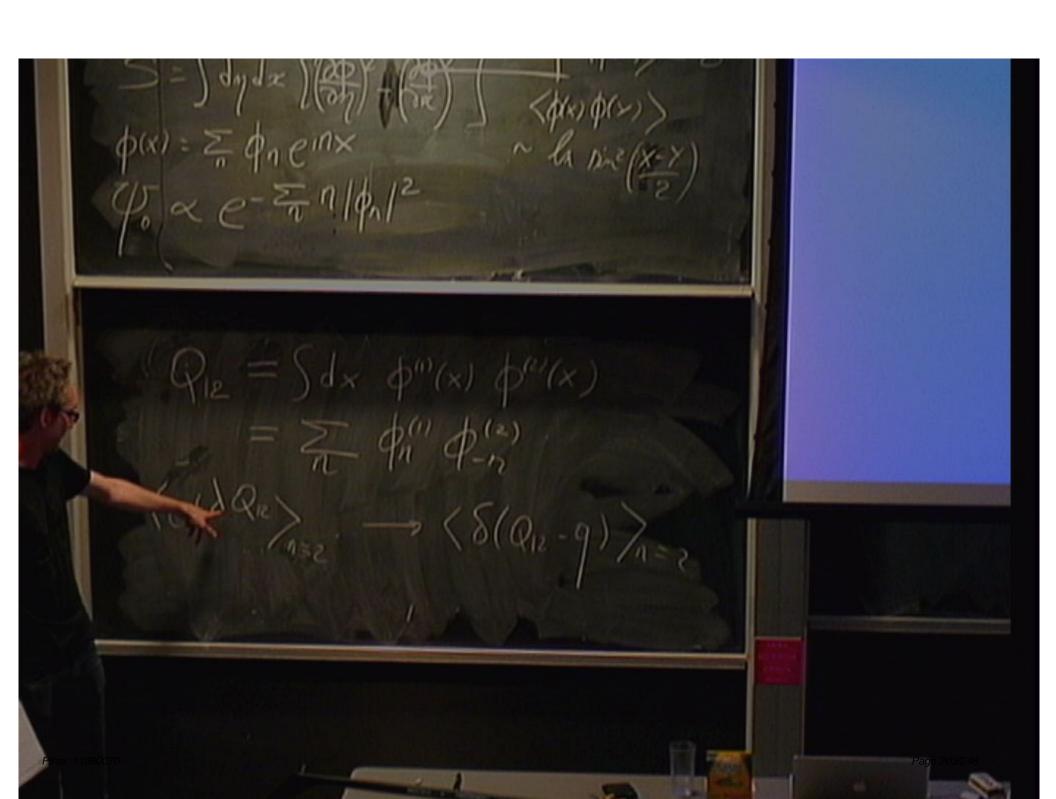
< f(x) \$(x-x) φ(x)= = φnemx

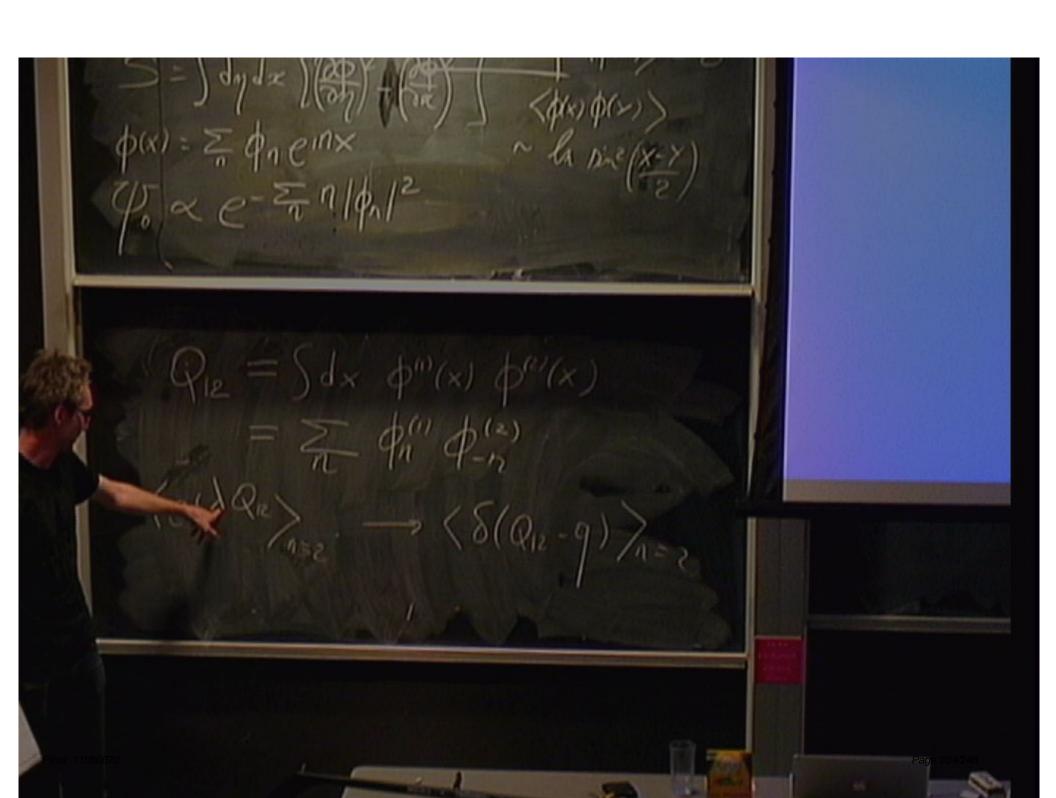
(γ) ~ e- = n | φη | 2 $= \sum_{n} \phi_n^{(i)} \phi_n^{(2)}$ (eil Que) -> (8(Qu-9))/1=2

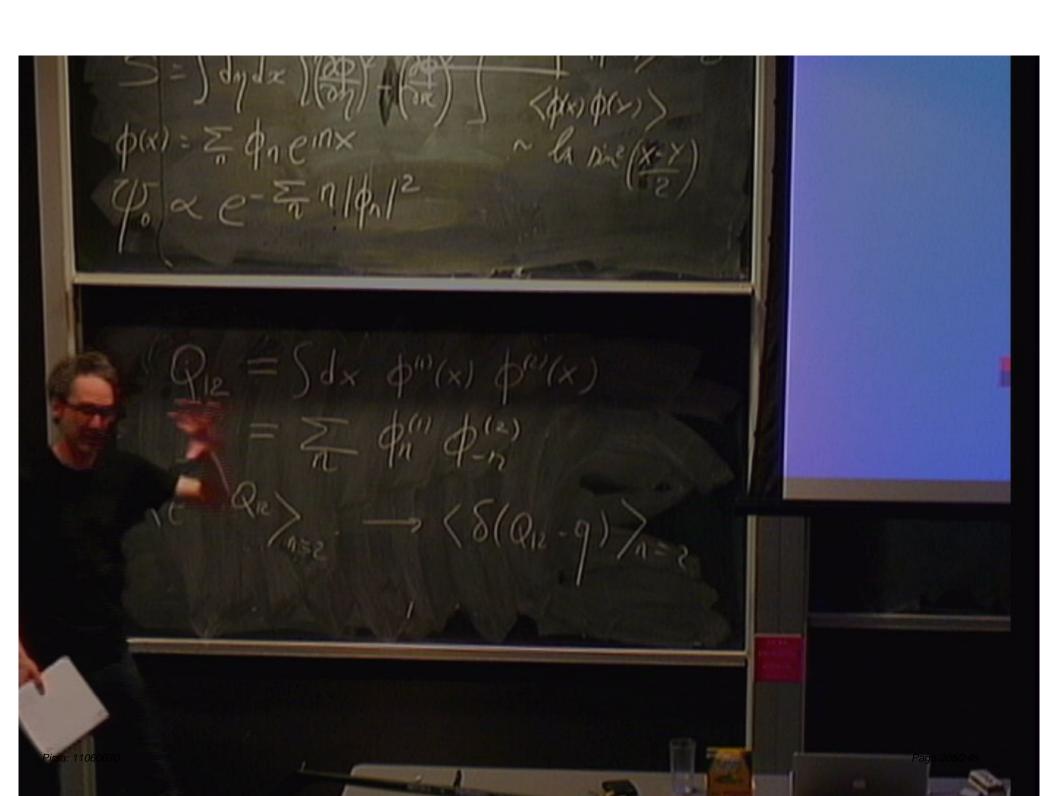




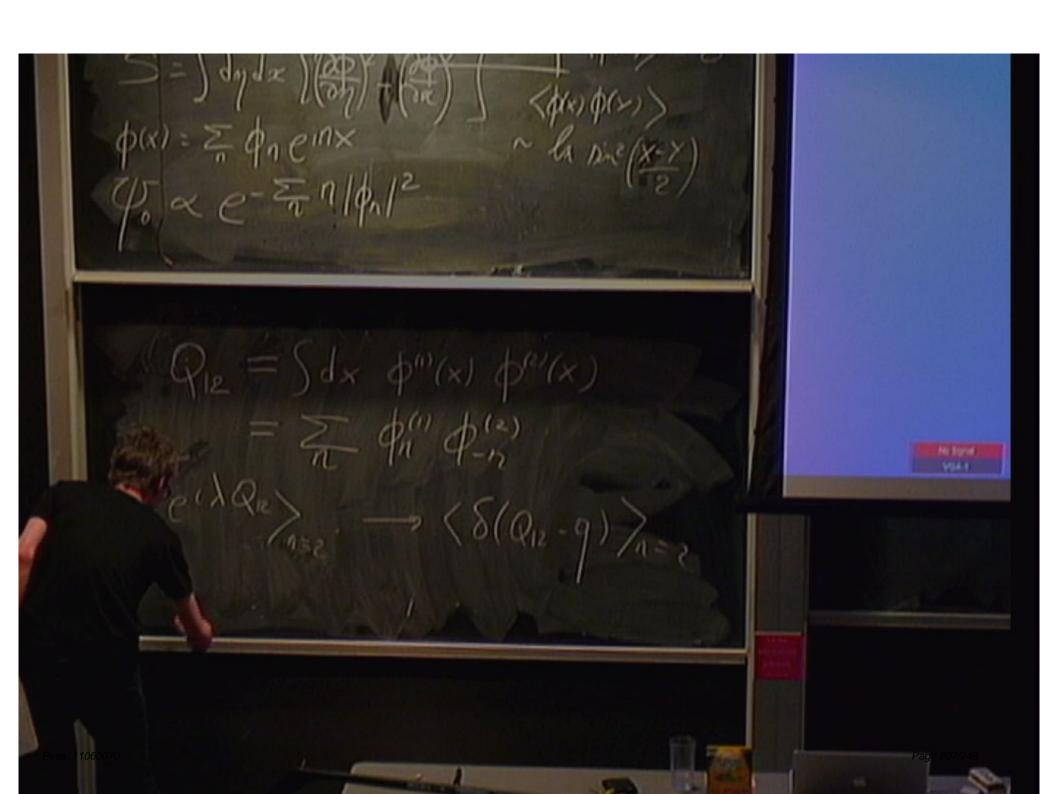


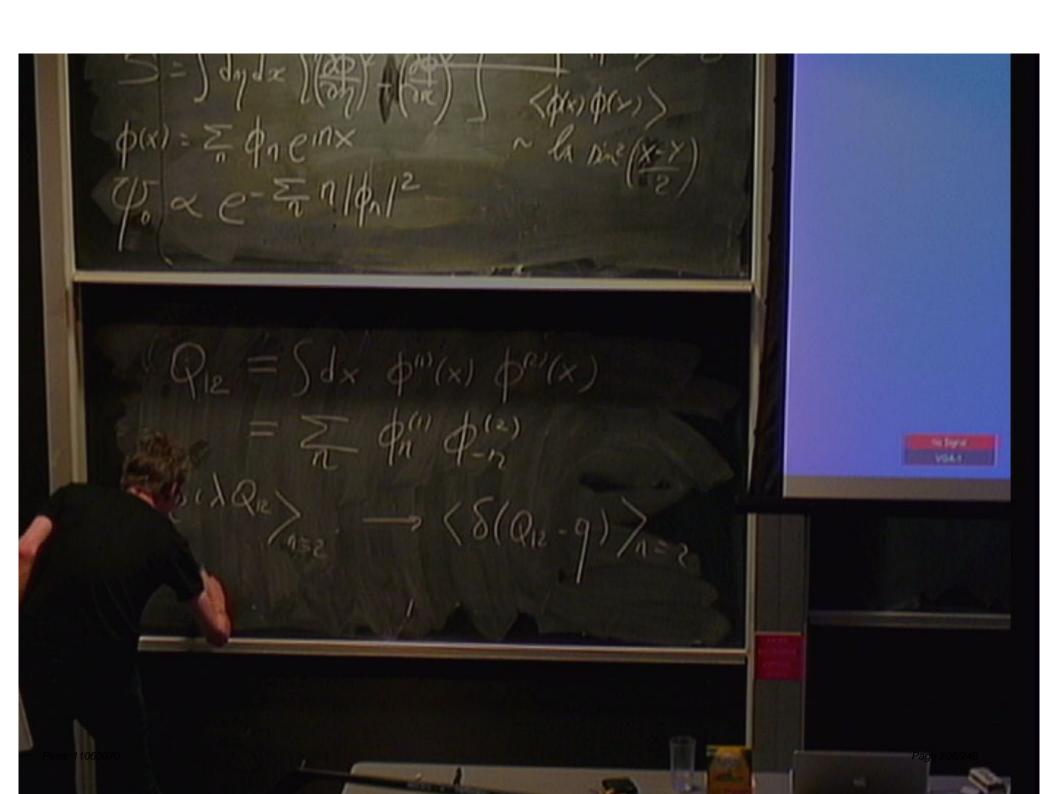


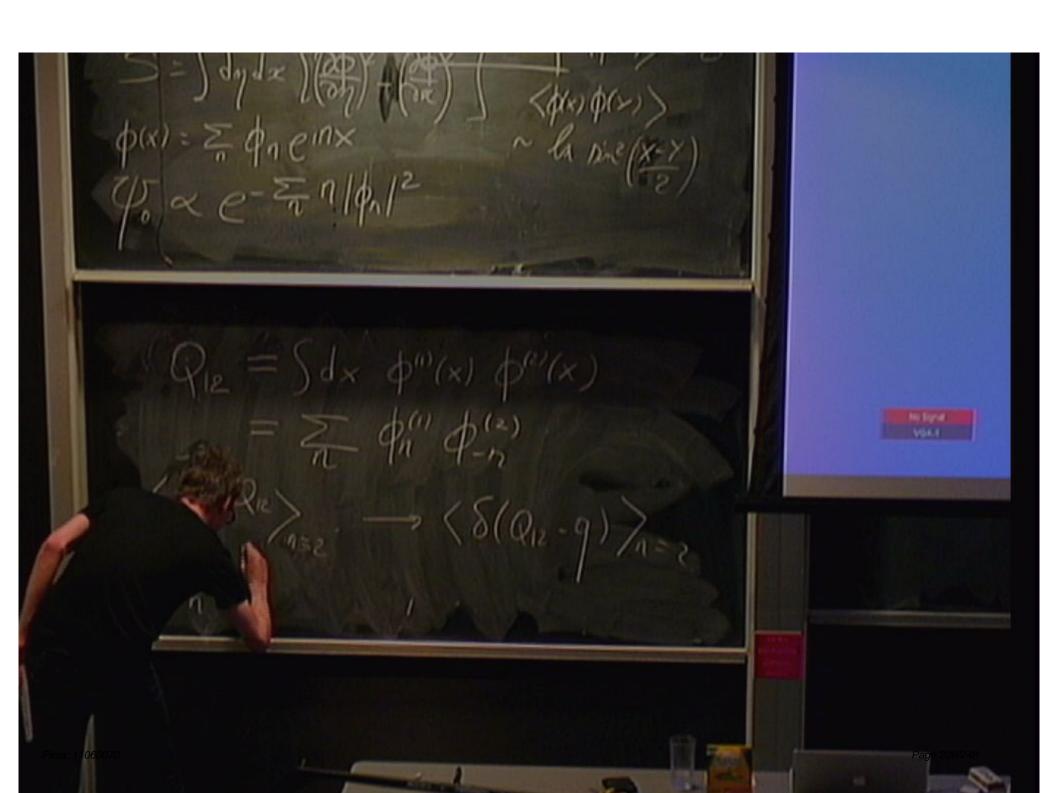


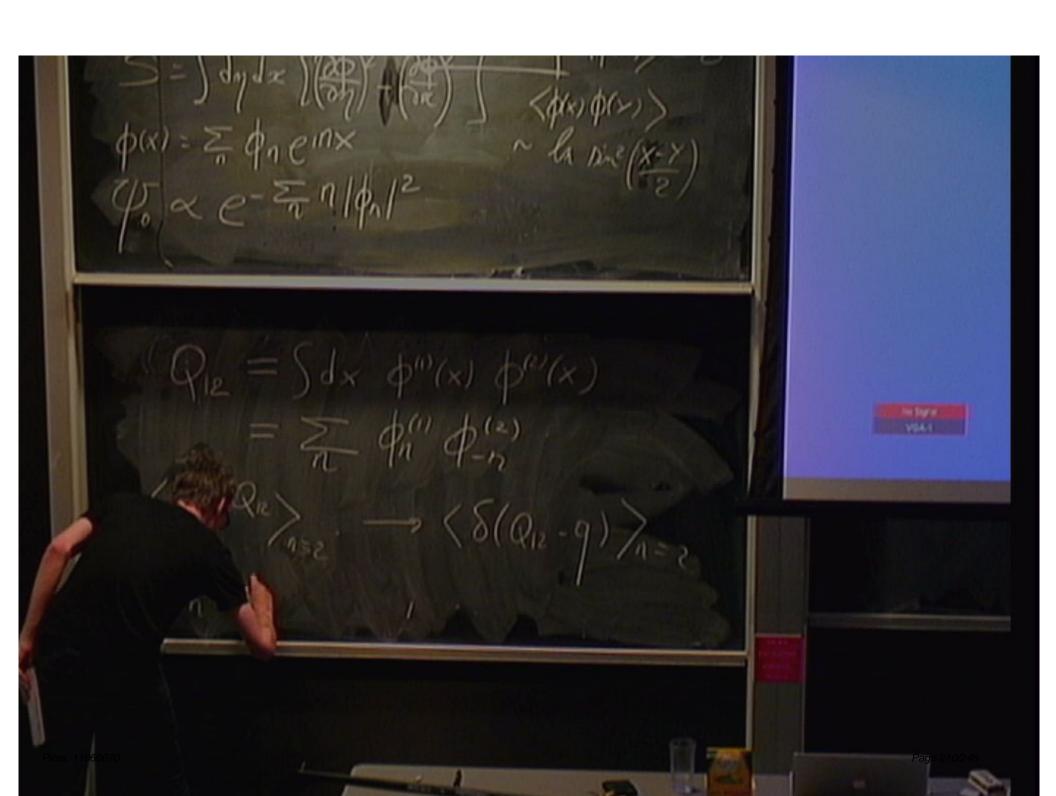


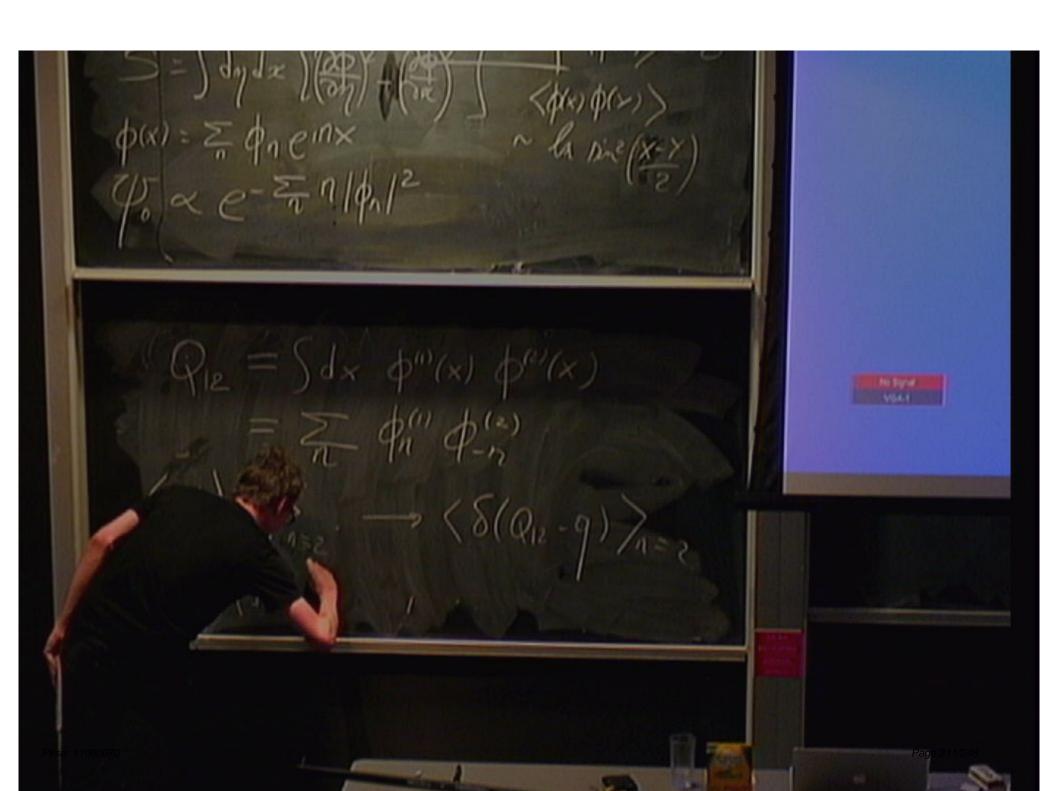
< pre (x-x) p(x) = = pnemx 40 × e- = 1/9/2 $= \sum_{n} \phi_n^{(n)} \phi_n^{(2)}$ (e1) Que -> (8(Qn-9)) /1= 2

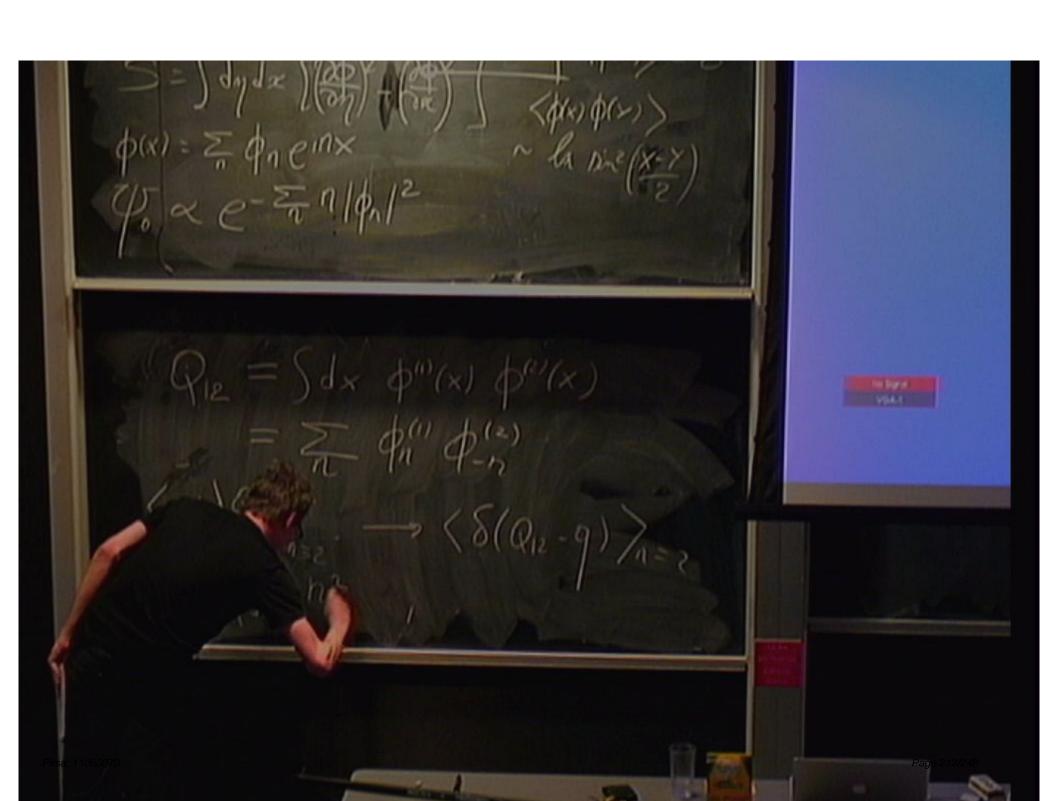


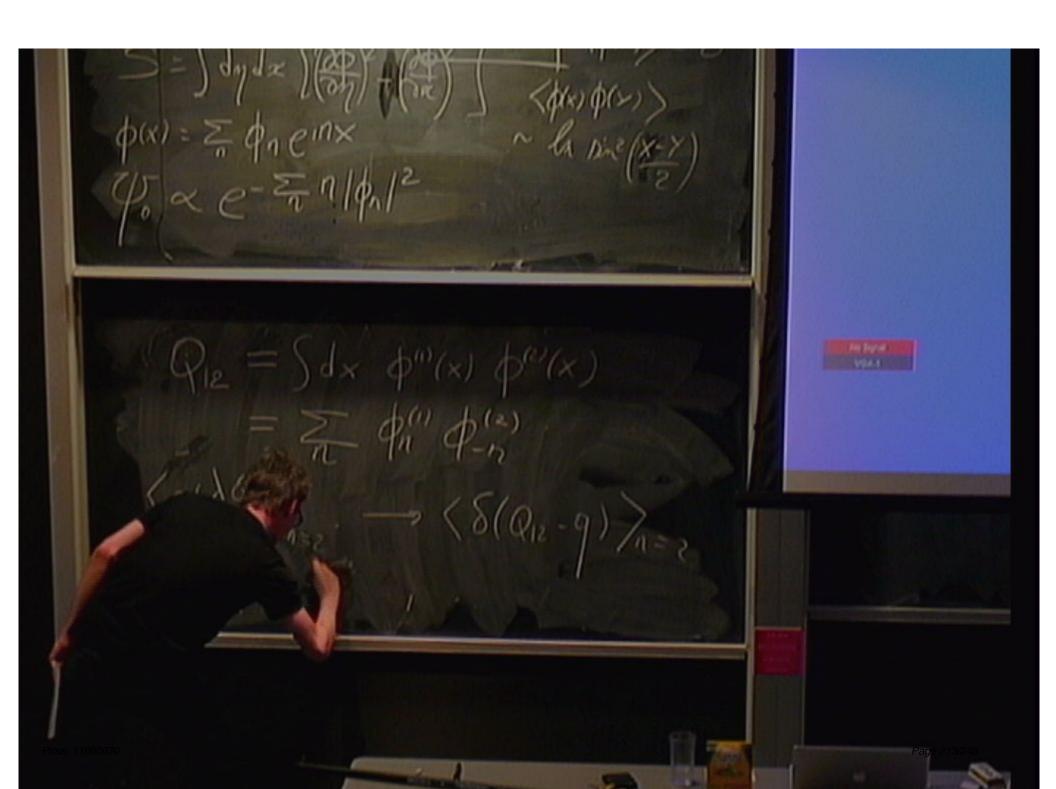


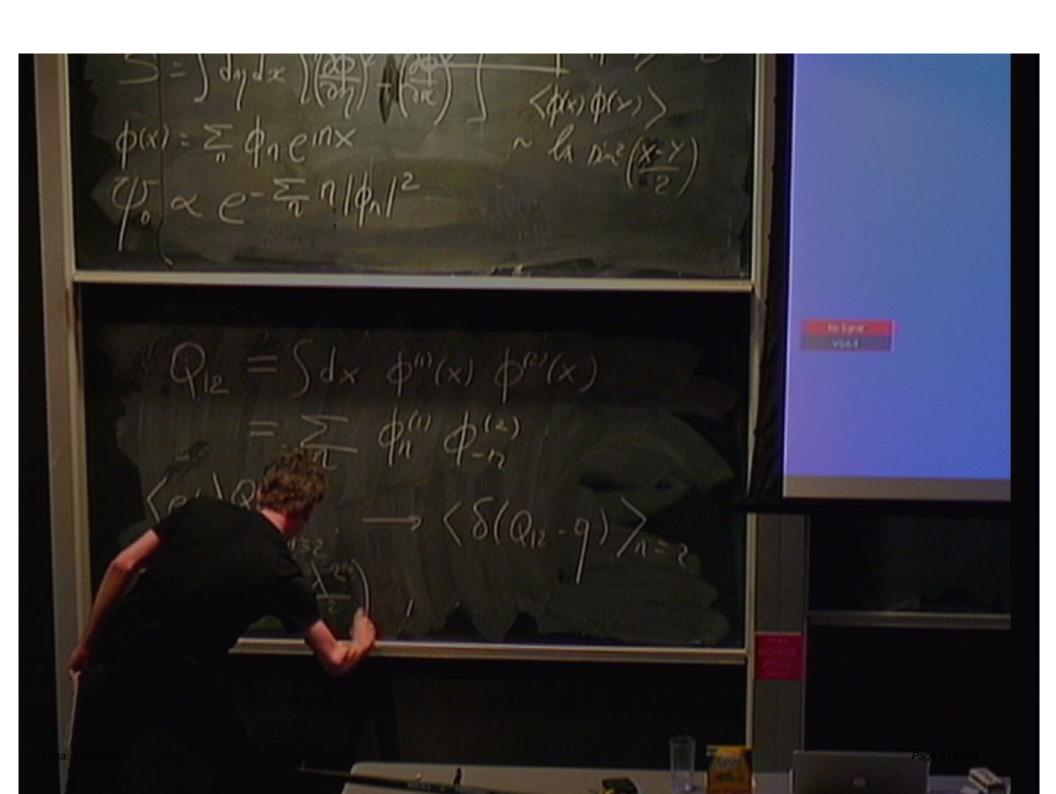


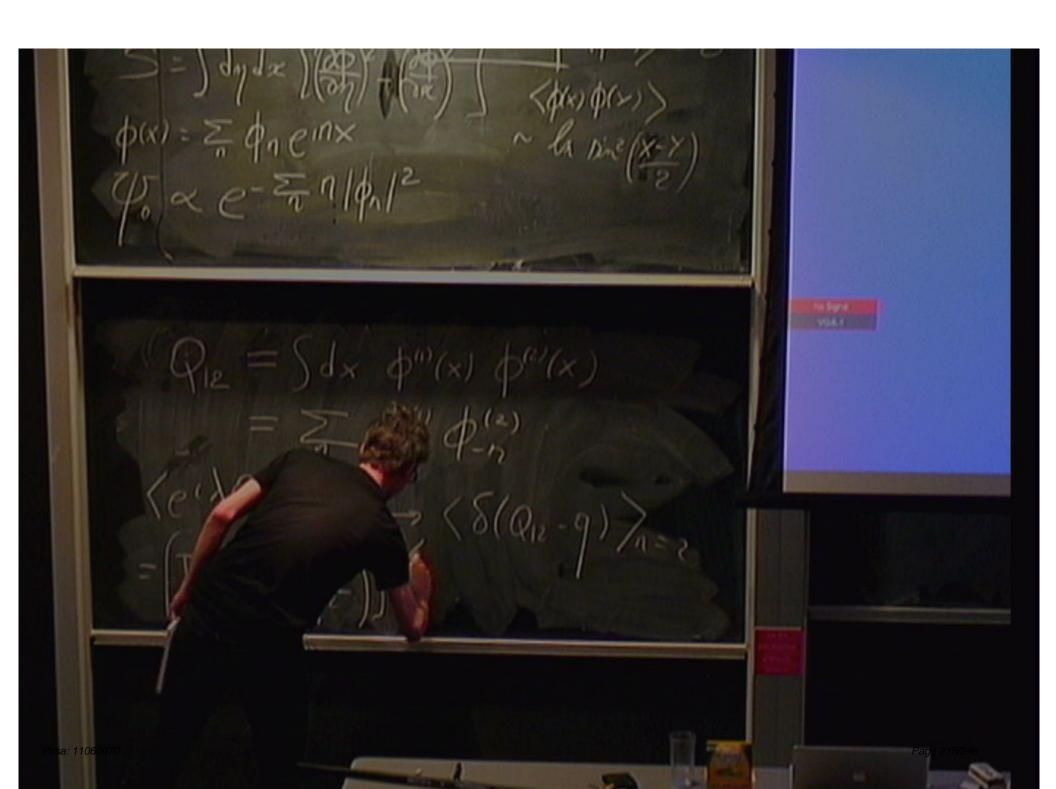










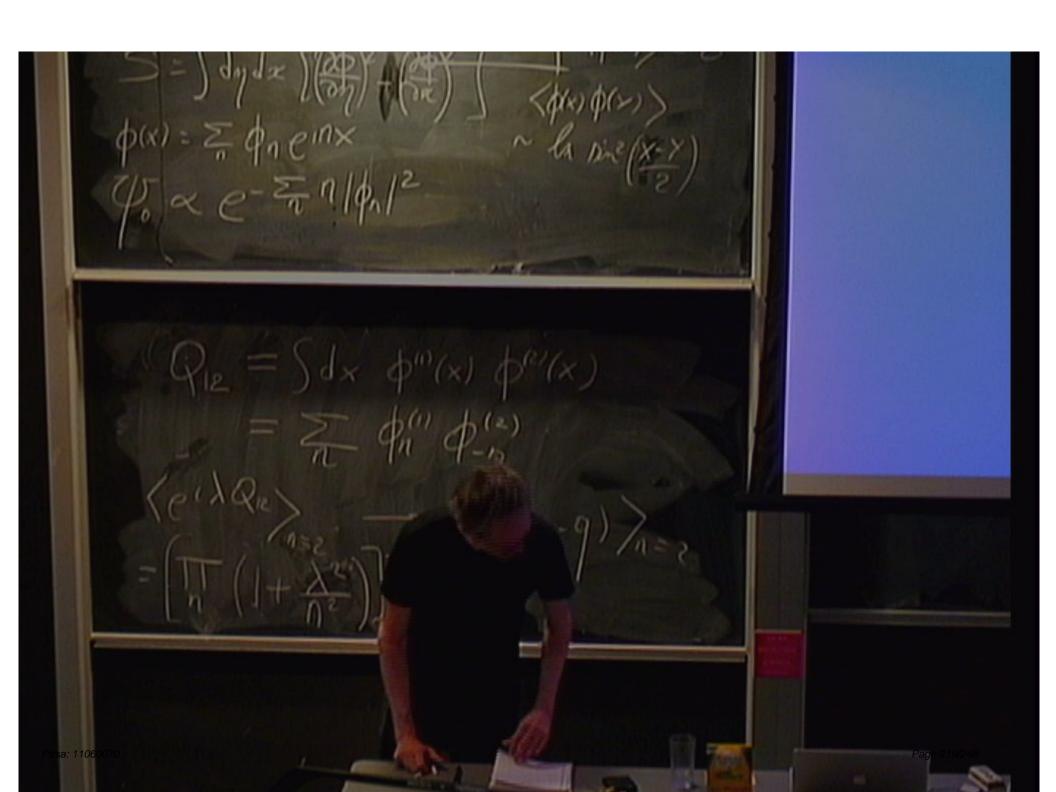


~ la 120 (x-x) φ(x) = = φnemx

(γ) ~ e = = n | φn | 2 = q'' φ''

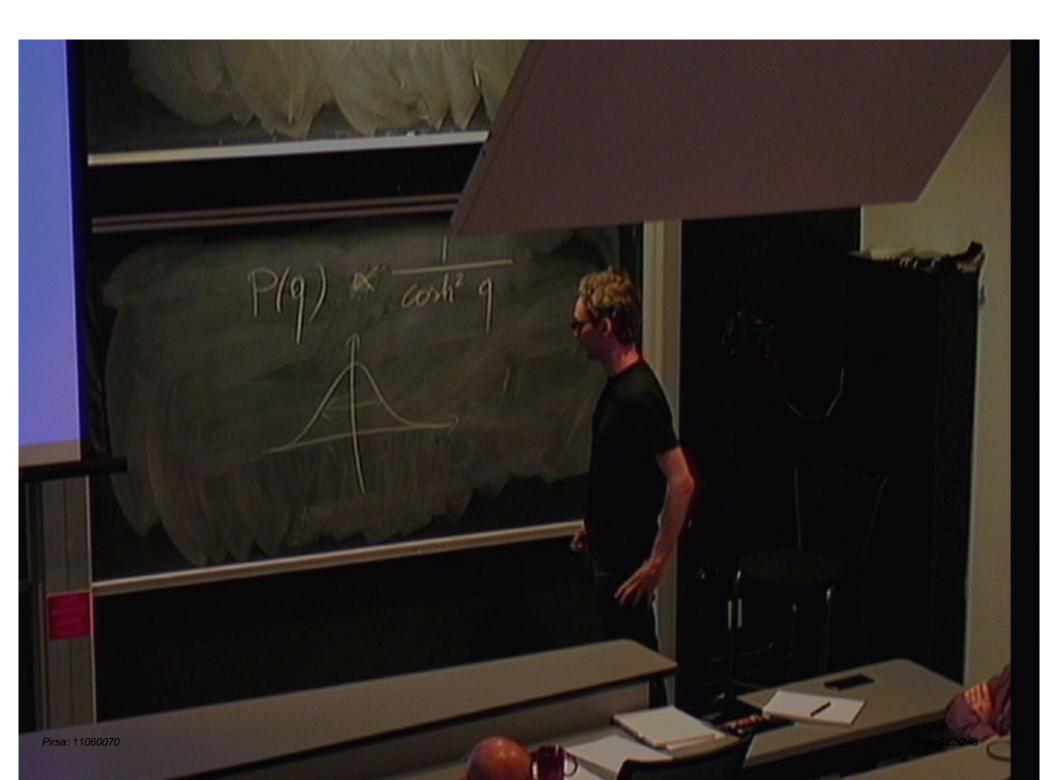
~ la 12 (x=x) p(x) = = pnemx 45 ~ e- = 1/9/2 = = p(1) p(2) = (8(Qn-9))=

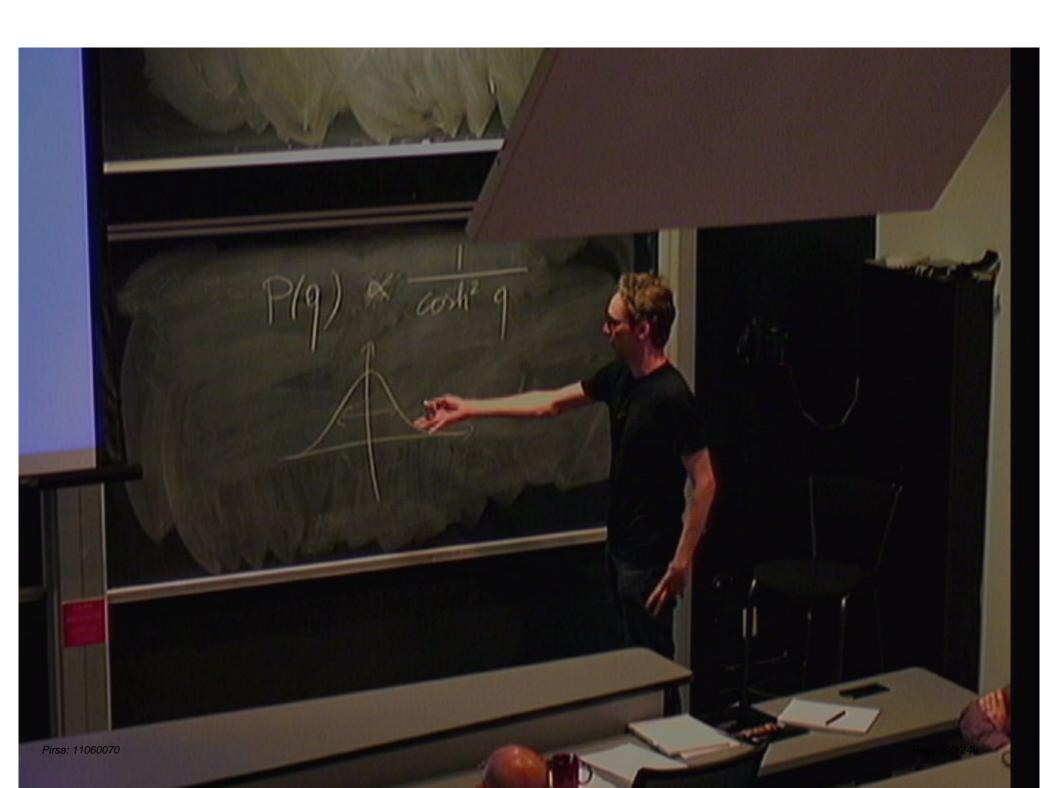
~ la 12 (x= x p(x) = Z pnemx 40 × e- = 1/9/2 = qn φ(2) -3 (8(Qn-9)) /n=3

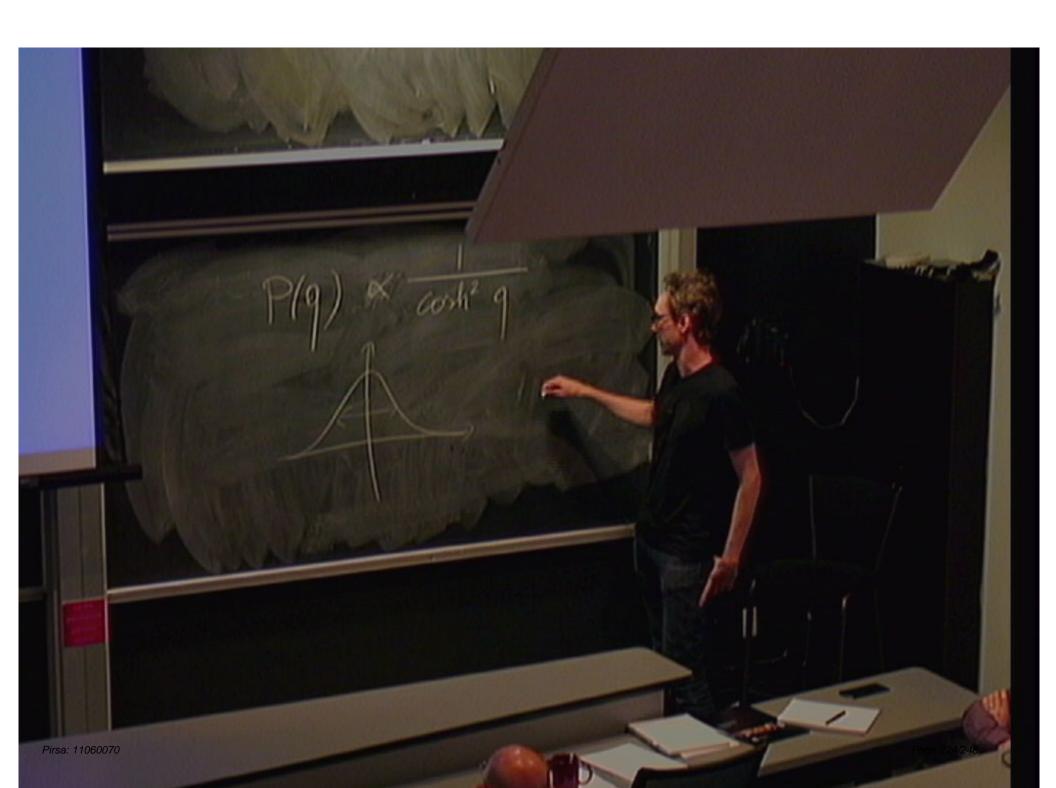


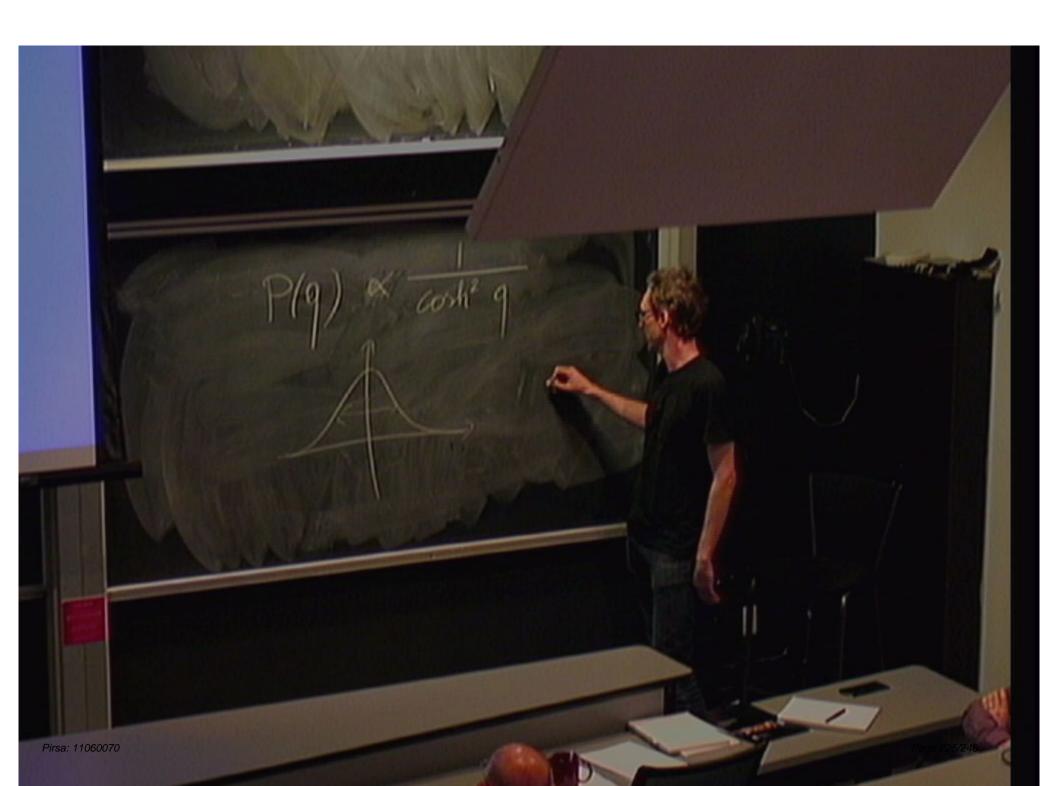
~ la 120 (x- γ) p(x) = Z pnemx 40 × e- = 1/1/2 = qn (2) - (8(Qn-9)) /1=

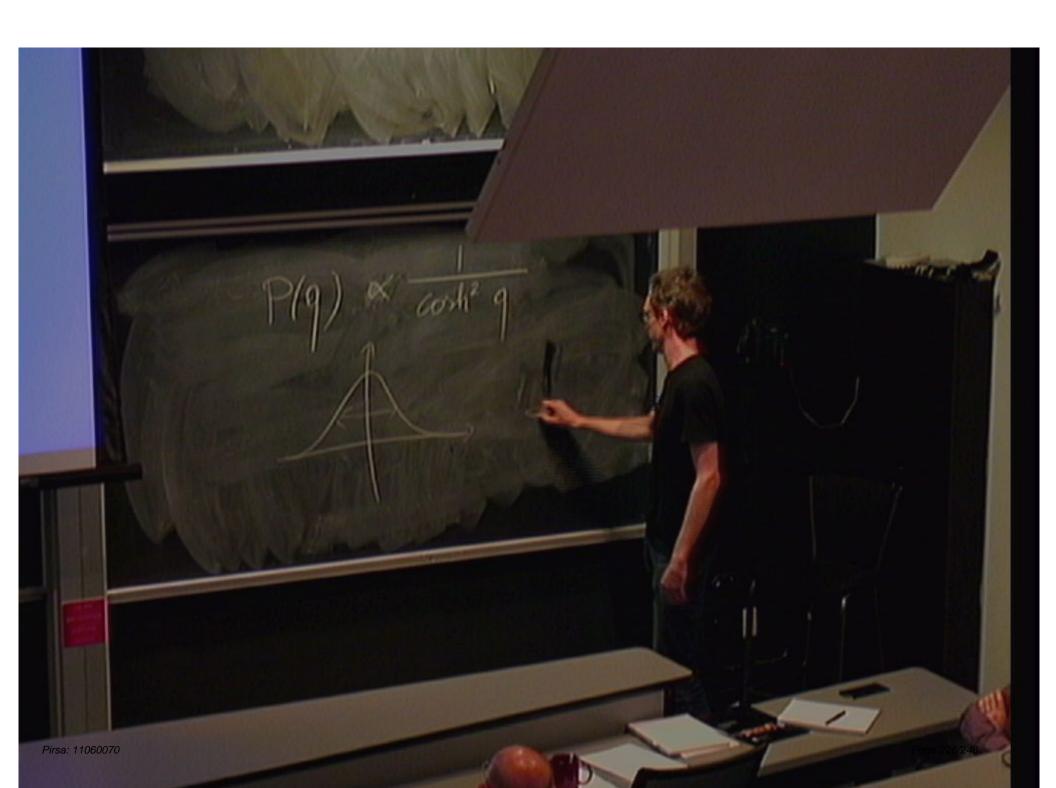
~ la 12 (x- γ) p(x) = Z qnemx 40 ~ e- = 1/9/2 = = Pn (2) - (8(Qn-9)) =

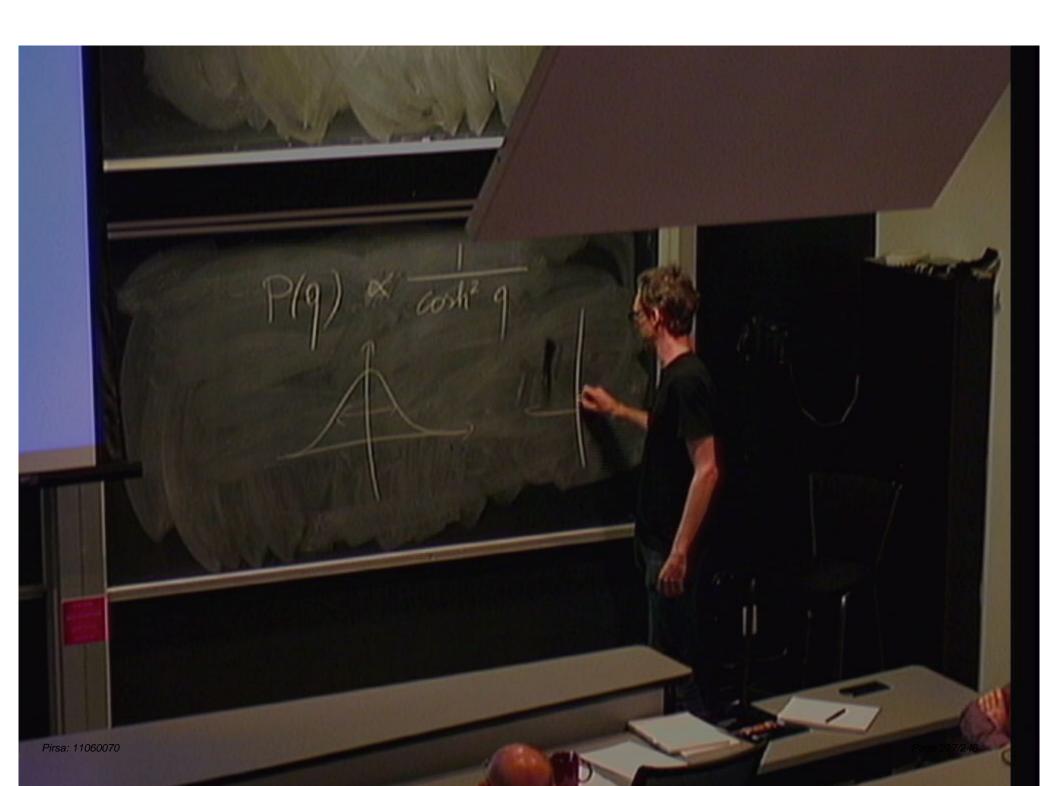


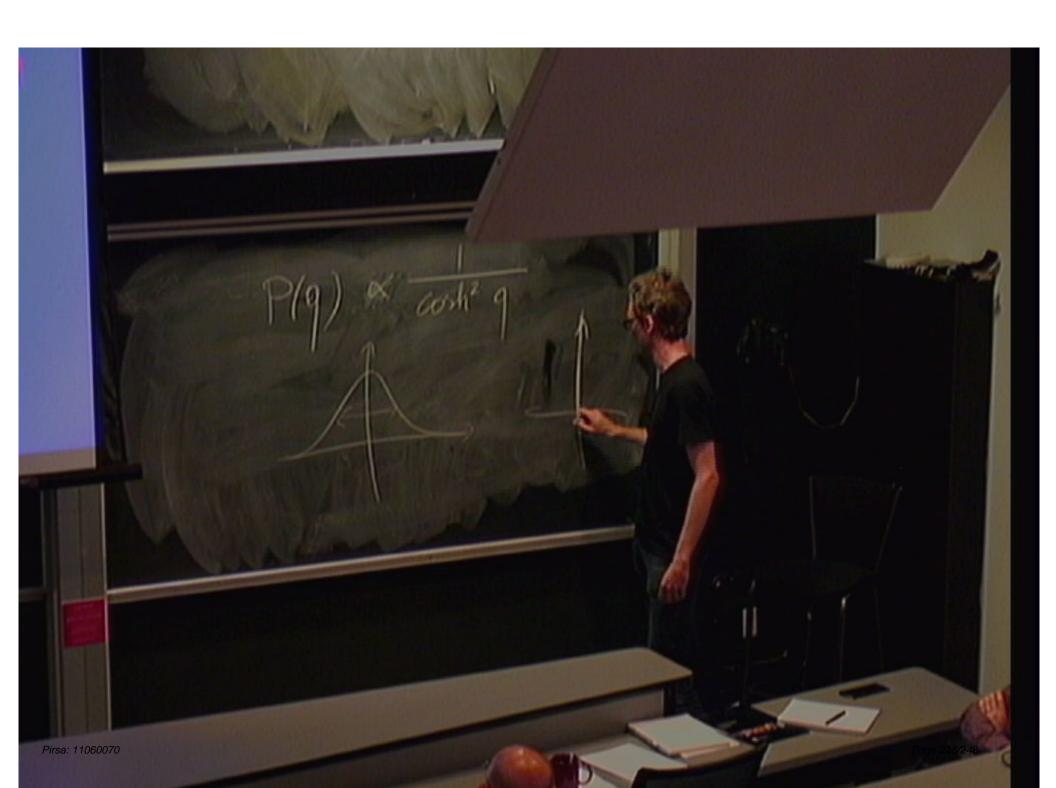


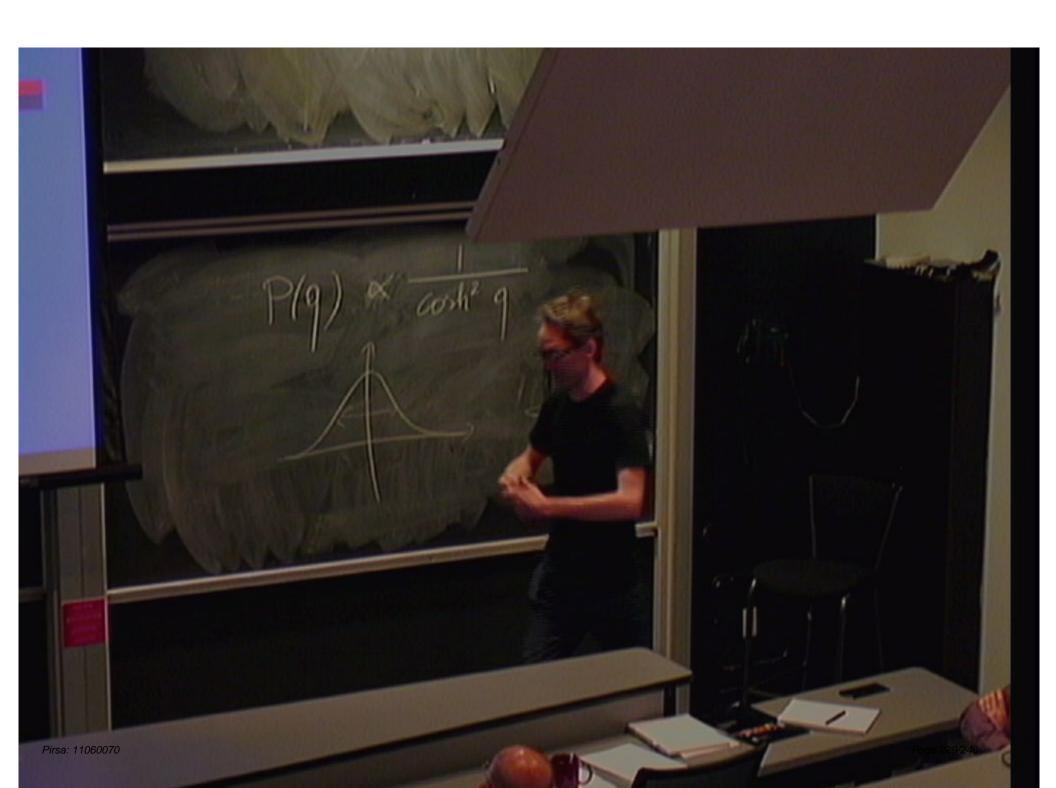


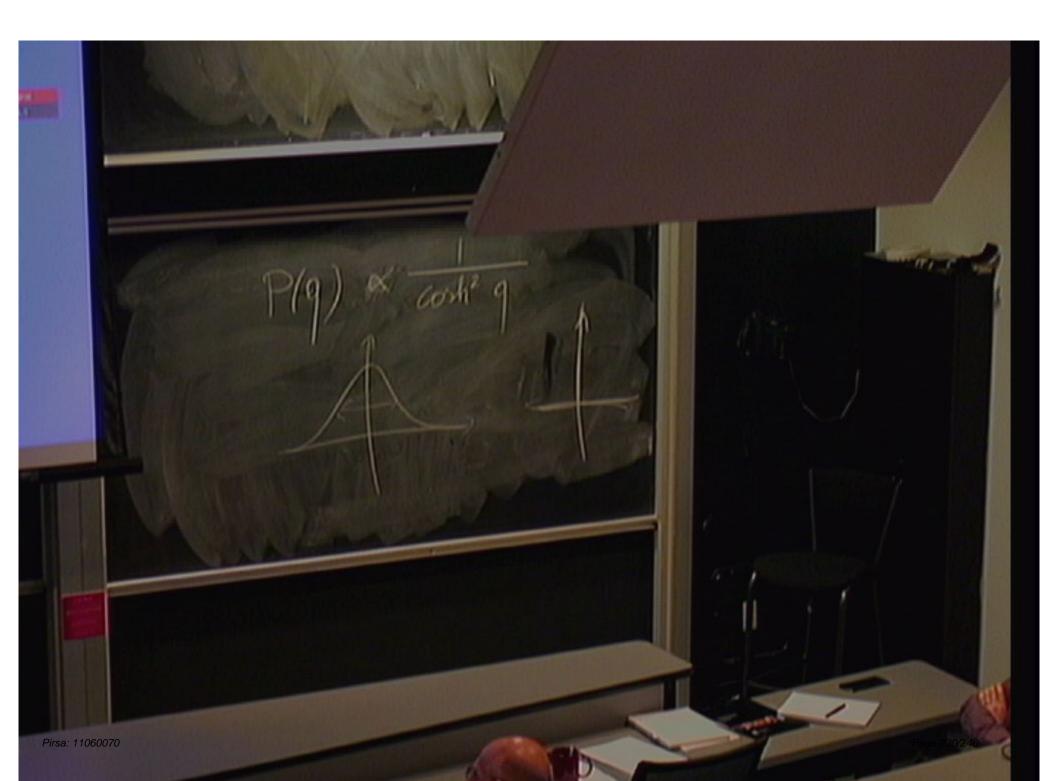


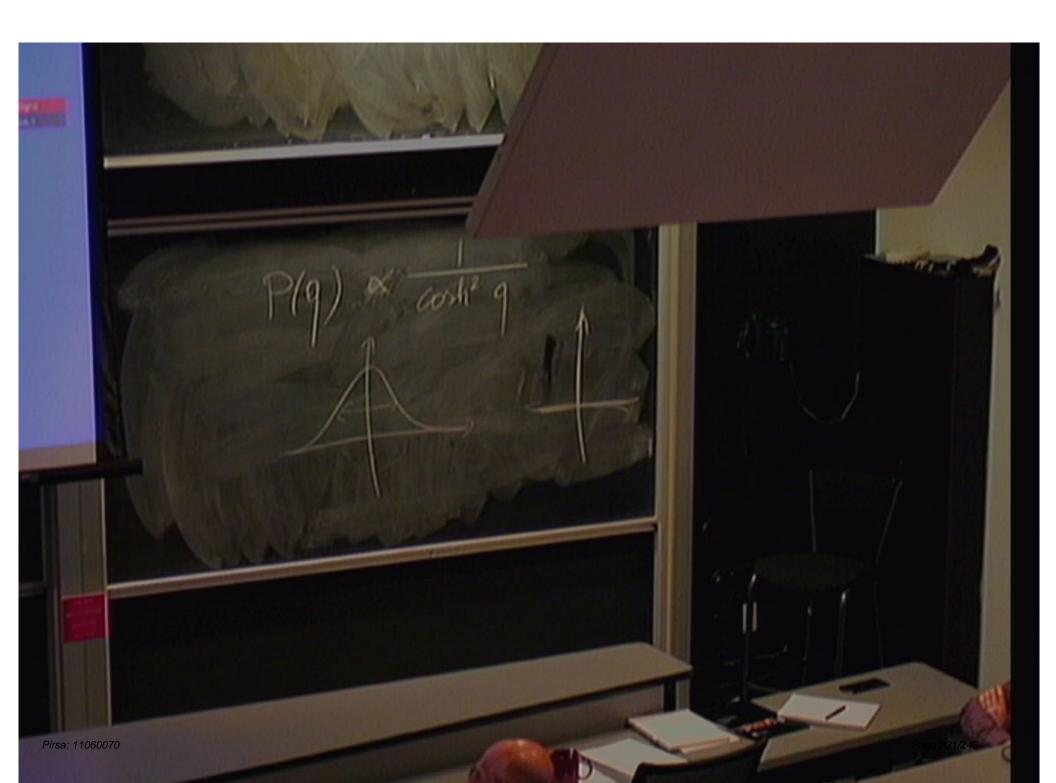


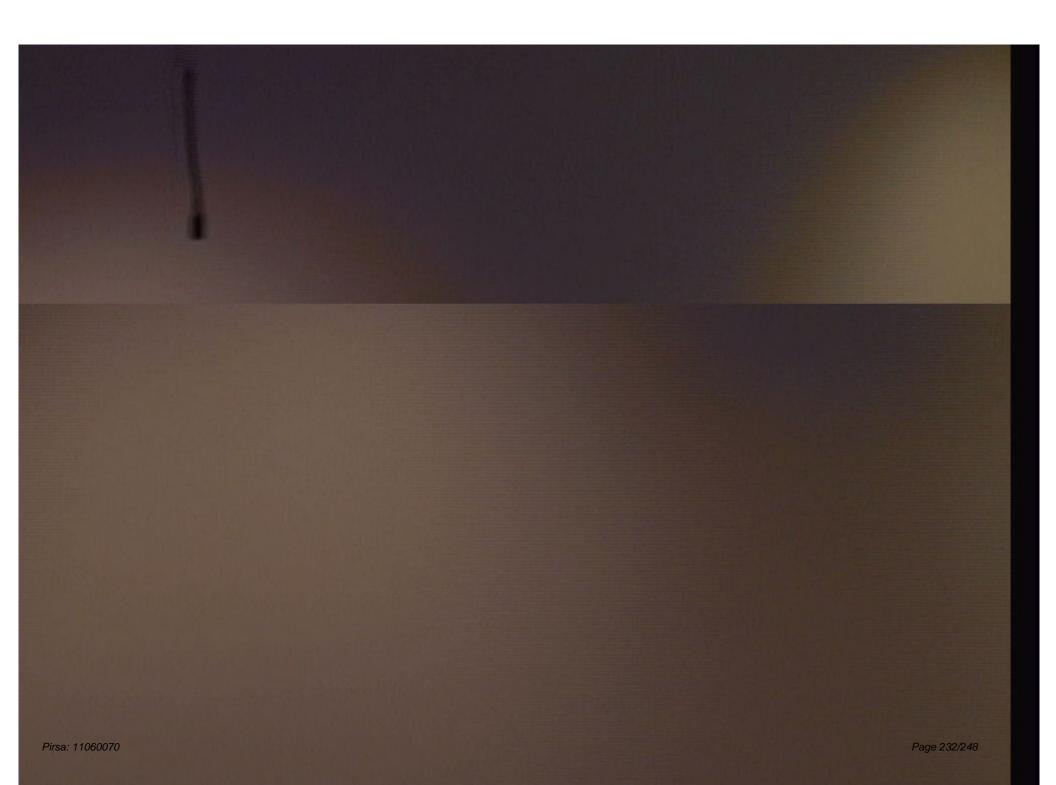


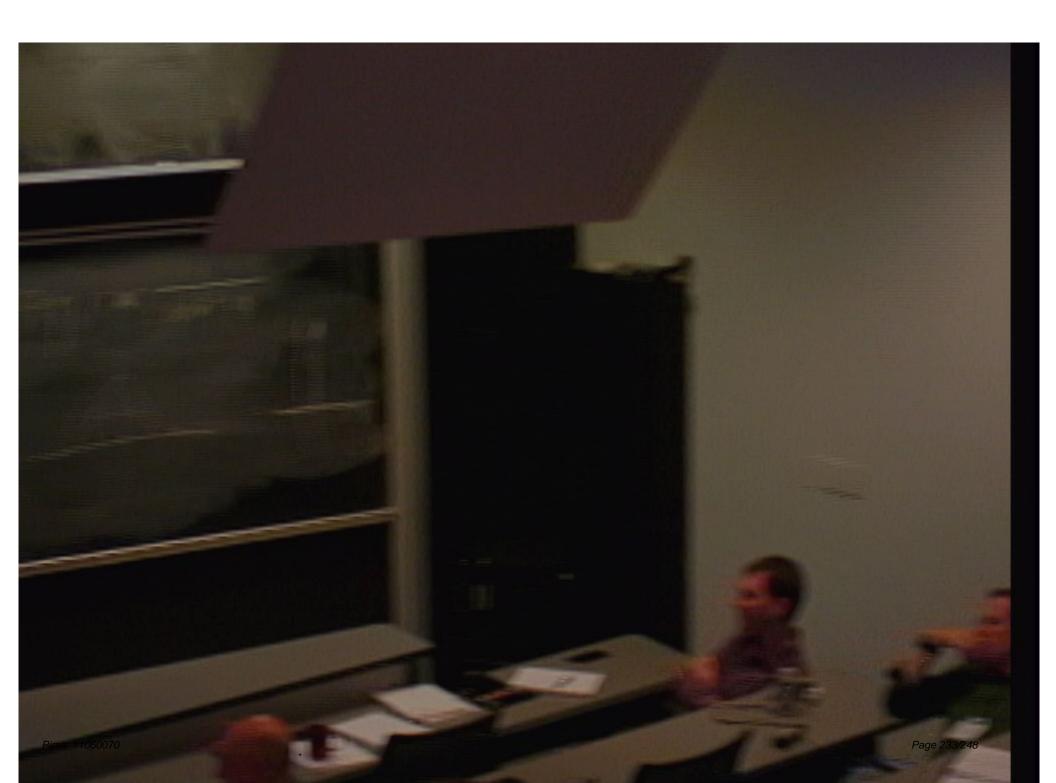












~ la Die (x-x) φ(x)= = φnemx

(γ) ~ e- = n/4/2

