Title: The Hidden Phase Space of our Universe

Date: Jun 22, 2011 09:00 AM

URL: http://pirsa.org/11060065

Abstract: By combining insights from black holes and string theory we argue for the existence of a hidden phase space associated with an underlying fast dynamical system, which is largely invisible from a macroscopic point of view. The dynamical system is influenced by slow macroscopic observables, such as positions of objects. This leads to a collection of reaction forces, whose leading order Born Oppenheimer force is determined by the general principle that the phase space volume of the underlying system is preserved. We propose that this adiabatic force is responsible for inertia and gravity. This fact allows us to calculate the hidden phase space volume from the known laws of inertia and gravity. We find that in a cosmological setting the appearance of dark energy is naturally explained by the finite temperature of the underlying system. The adiabatic approximation that leads to the usual laws of inertia and gravity breaks down in the neighborhood of horizons. In this regime the reaction force degenerates into an entropic force, and the laws of inertia and gravity receive corrections due to thermal effects. A simple estimate of these effects leads to the conclusion that they coincide with observed phenomena attributed to dark matter.

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Introduction:

Gravity is a macroscopic force that dominates in IR.

Yet it knows about microscopic information (UV).

But why? Is it a UV/IR conspiracy?

Or is it due to a principle?

Basic idea:

To arrive at our paradigm in terms of space time, matter, and forces we ignore (non-local) UV DOF.

This requires a separation of time scales: breaks down near horizons and in cosmology.

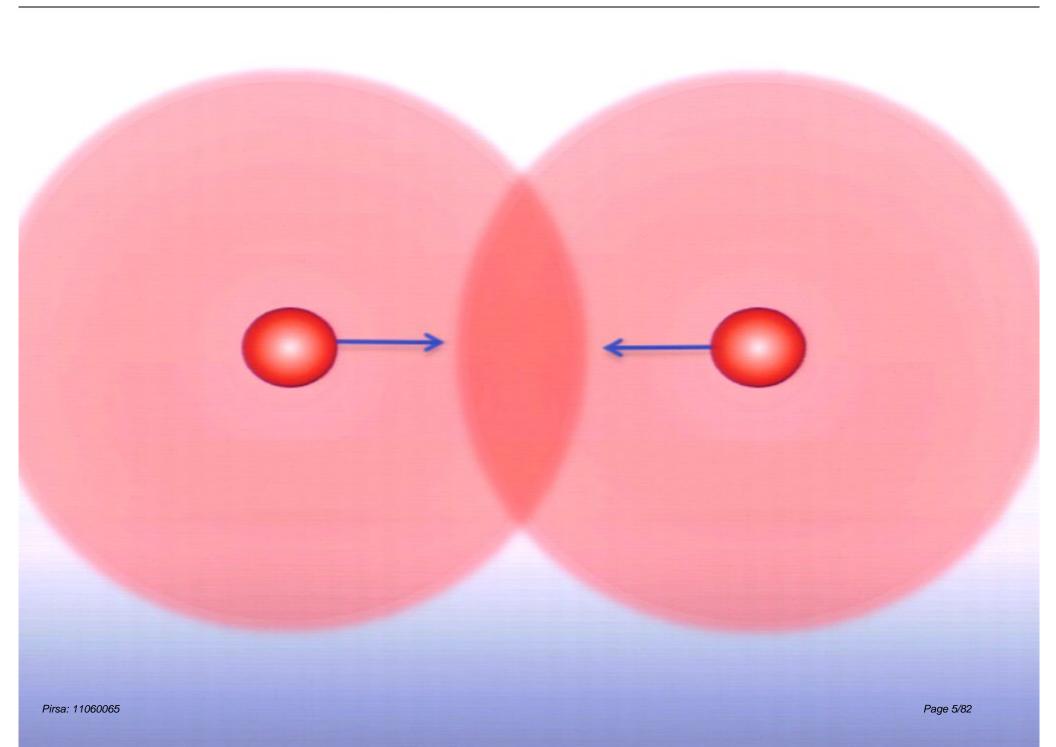
The distinction between matter and space time will disappear and forces should be seen as emergent.

Gravity, and the other forces, are reaction forces due to the fast microscopic dynamics of the

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Anderlying dynamical system





Outline:

- Motivation: BH's + String-/matrix theory.
- Adiabatic reaction forces.
- Hidden phase space: proposal.
- Inertia/gravity as adiabatic reaction force.
- Speculations about dark energy/dark matter.

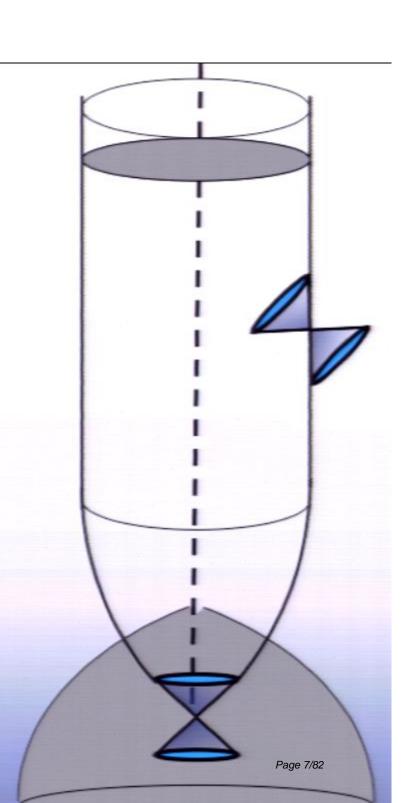
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GRAVITATIONAL COLLAPSE:

What happens to the phase space occupied by the fermions?

NEUTRON STAR

Degenerate Fermions



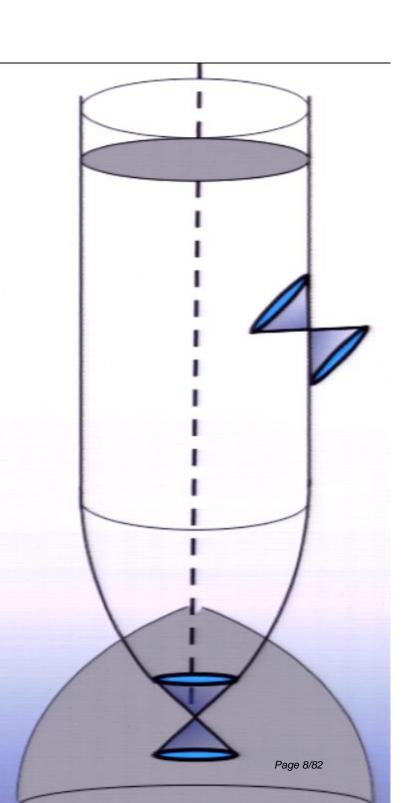
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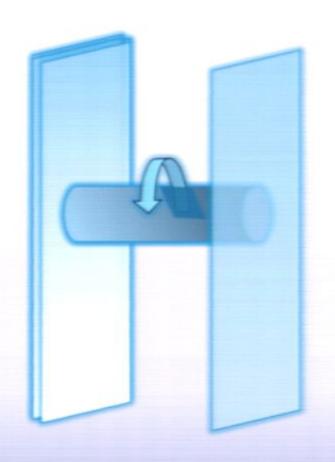
String theory/ AdS-CFT:

$$\Psi(x) \leftrightarrow \operatorname{tr}(X^n)$$

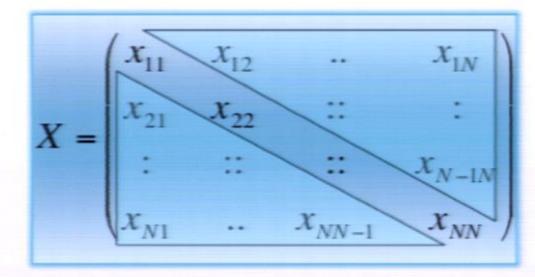
Bulk fields correspond to (short) traces of boundary fields.



Gravity results from integrating out DOF, associated with the emergence of space-time

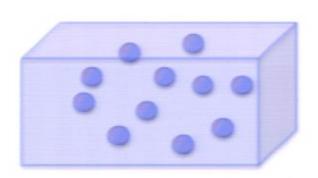


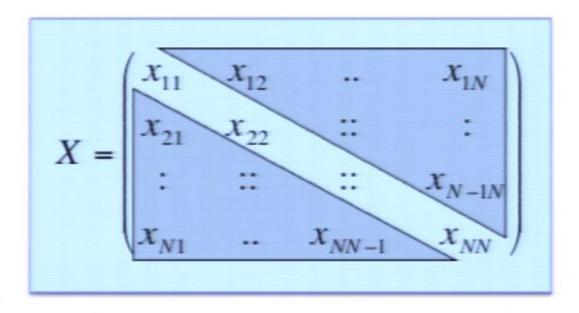
Coordinates are matrices



$$\operatorname{tr}\left(\dot{X}_{I}^{2}+\left[X_{I},X_{J}\right]^{2}+\Psi^{*}X_{I}\Gamma^{I}\Psi\right)$$

Eigenvalues: describe position of matter Off-diagonal DOF induce forces.





The off diagonal modes carry a <u>positive</u> gravitational energy

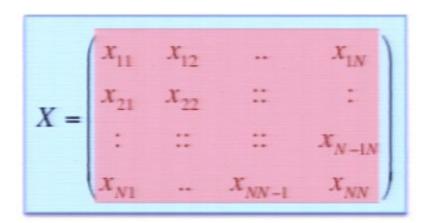
$$E_g = \frac{1}{8\pi G} \int \left| \nabla \Phi \right|^2$$

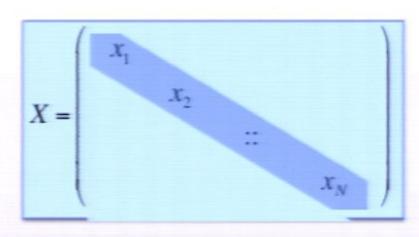
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The phase space occupied by the fermions goes into the "off diagonal" modes.

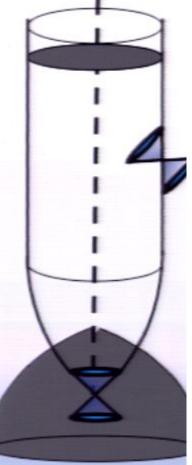
Eigenvalues and off diagonal modes equilibrate and together form a thermal state describing the black hole.

Space-time in its usual seasonsceases to exist.

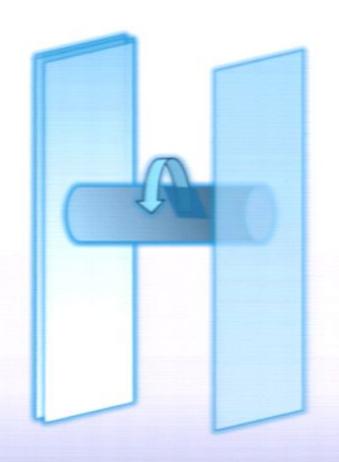




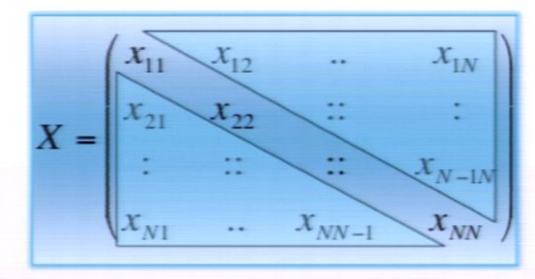
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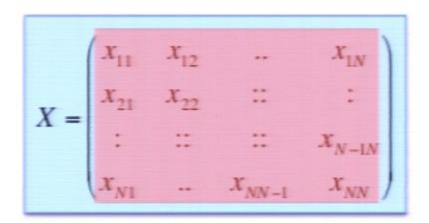
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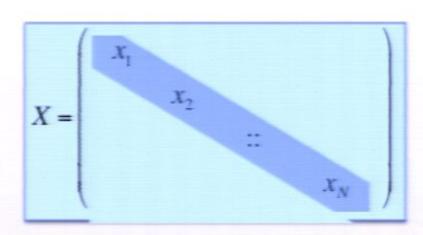
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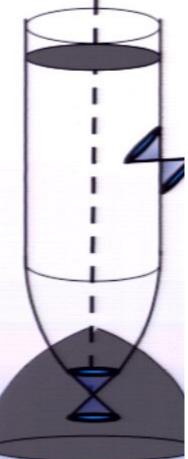
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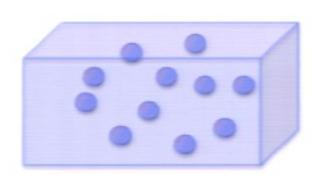
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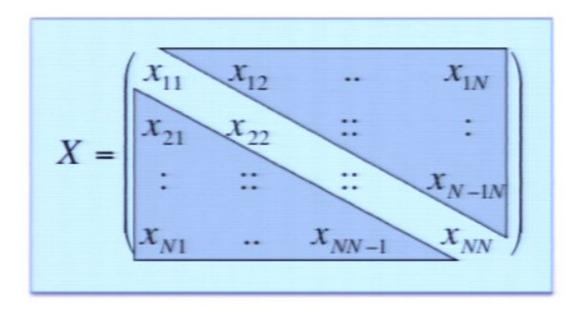




Degenerate Fermions



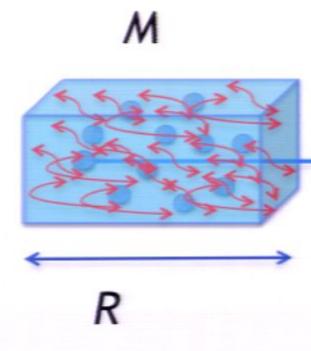




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lack Hole Horizon

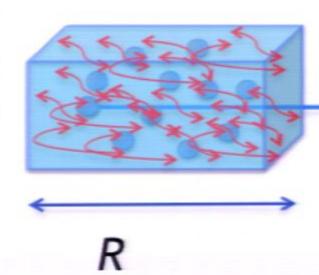


Box with a gas of particles.

What is the maximal entropy inside the box?

lack Hole Horizon

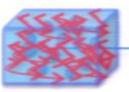
M



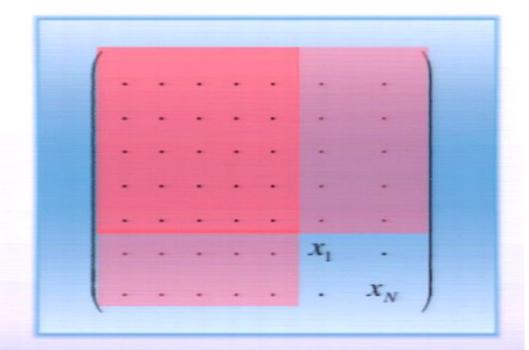
$$\Delta W = MgR = T\Delta S$$

$$T = \frac{\hbar g}{2\pi} \implies \Delta S = 2\pi \frac{McR}{R}$$

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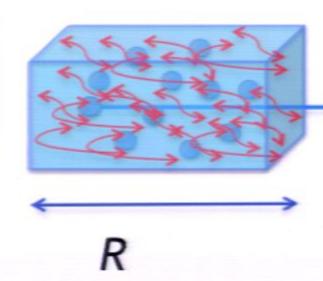


$$N = \frac{McR}{\hbar}$$



lack Hole Horizon

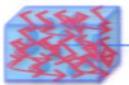
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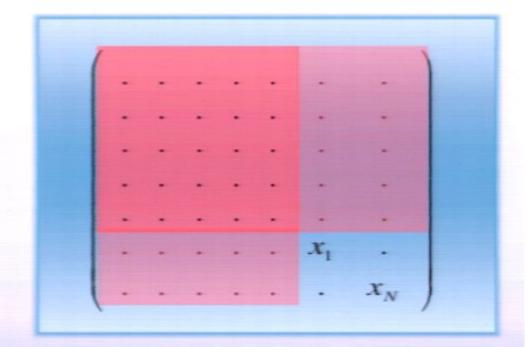
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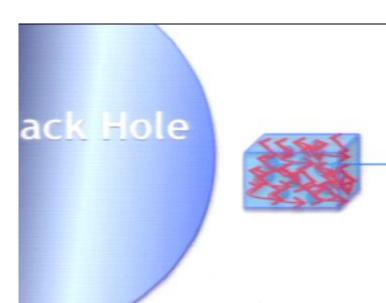
$$T = \frac{\hbar g}{2\pi} \implies \Delta S = 2\pi \frac{McR}{R}$$

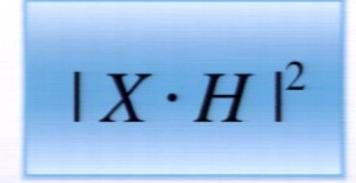
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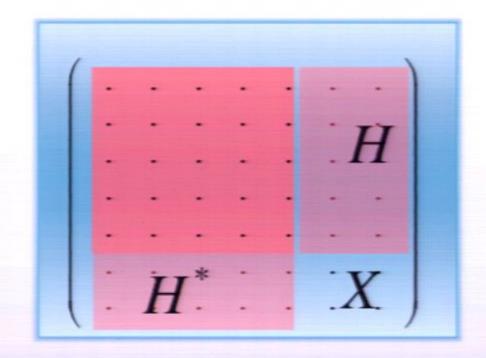


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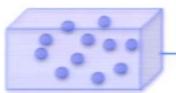


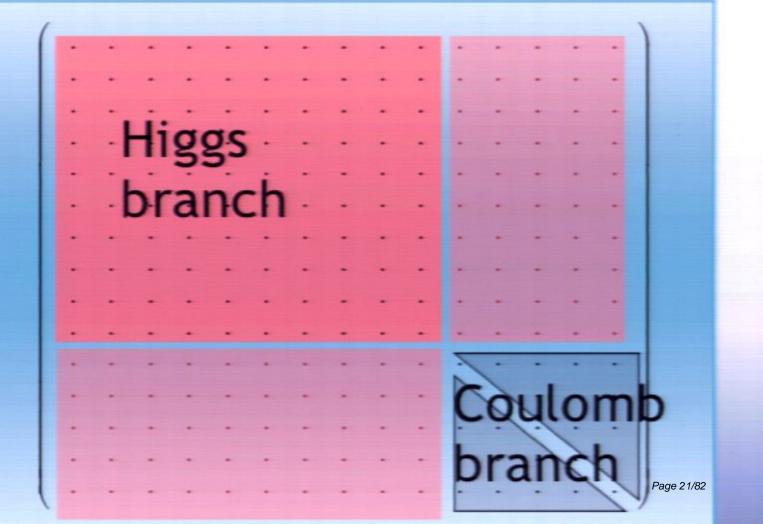




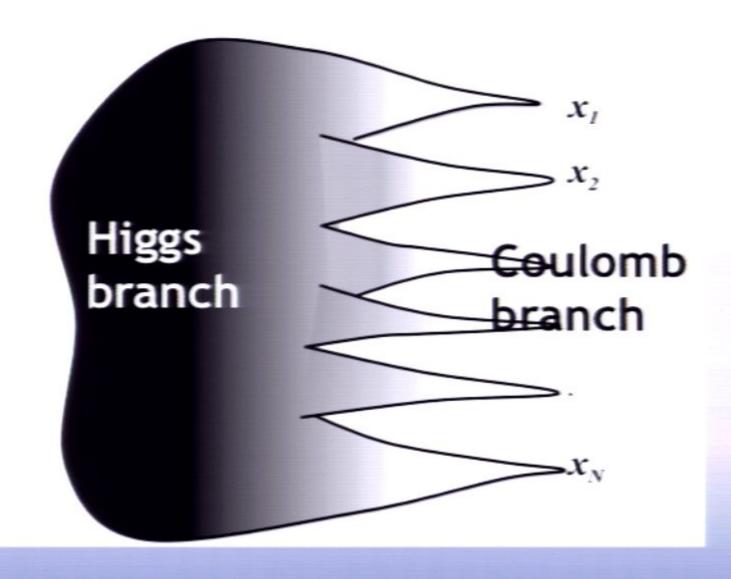
Matrices gets "Higgsed"
Eigenvalues disappear as flat directions
Pisc 100061 lomb branch goes over into Higgs branch

ack Hole





PHASE SPACE



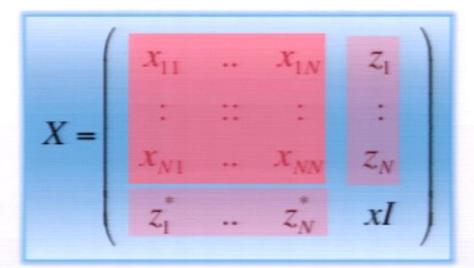
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ack Hole Horizon

$$X = \begin{pmatrix} x_{11} & \dots & x_{1N} & z_1 \\ \vdots & \vdots & \vdots & \vdots \\ x_{N1} & \dots & x_{NN} & z_N \\ z_1^* & \dots & z_N^* & xI \end{pmatrix}$$

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$$F = mg$$

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$$F = T \frac{\partial S}{\partial x}$$

$$T = \frac{\hbar g}{2\pi} \implies \frac{\partial S}{\partial x} = 2\pi \frac{mc}{\hbar}$$

Adiabatic reaction forces:

When a fast dynamical system is driven by a slow system the fast reacts back on the slow and creates a reaction force.

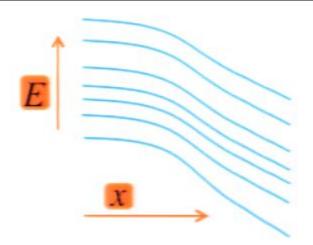
When the time scales are widely separated the force is determined by the principle that the phase space volume is preserved.

In quantum mechanics this is a consequence of the Born-Oppenheimer approximation.

Adiabatic Reaction Force

Harmonic oscillator with slowly varying frequency

$$H(p,q;x) = \frac{1}{2} \left(p^2 + \omega^2(x) q^2 \right)$$



Bohr-Sommerfeld action integral satisfies the basic identiy

$$J = \frac{1}{2\pi} \oint p dq$$

$$dE = \omega \ dJ - Fdx$$

Adiabatic reaction force:

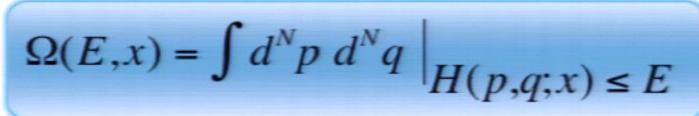
$$F = -\left(\frac{\partial E}{\partial x}\right)_{J} = \omega \left(\frac{\partial J}{\partial x}\right)_{E}$$
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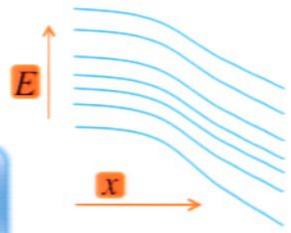
$$\frac{1}{\omega} = \left(\frac{\partial J}{\partial E}\right)_{x}$$

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Adiabatic Reaction Force

For a (chaotic/ergodic) system with many DOF





the phase space volume is an adiabatic invariant. Formally it defines an entropy

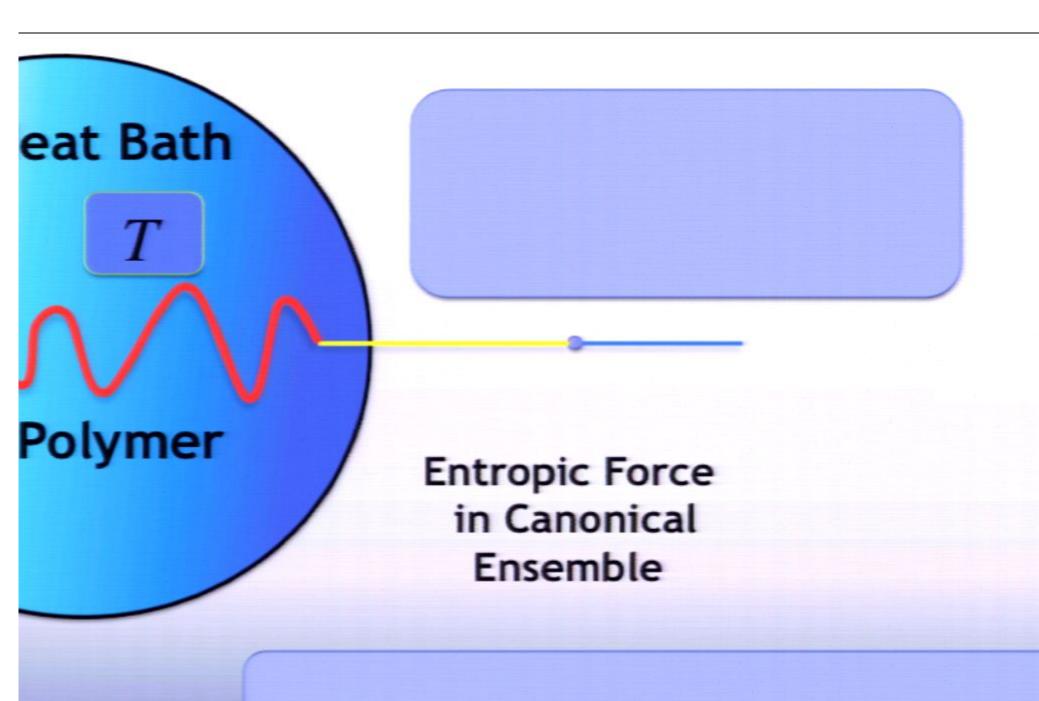
$$S = k_B \log \Omega$$

$$dE = T dS - F dx$$

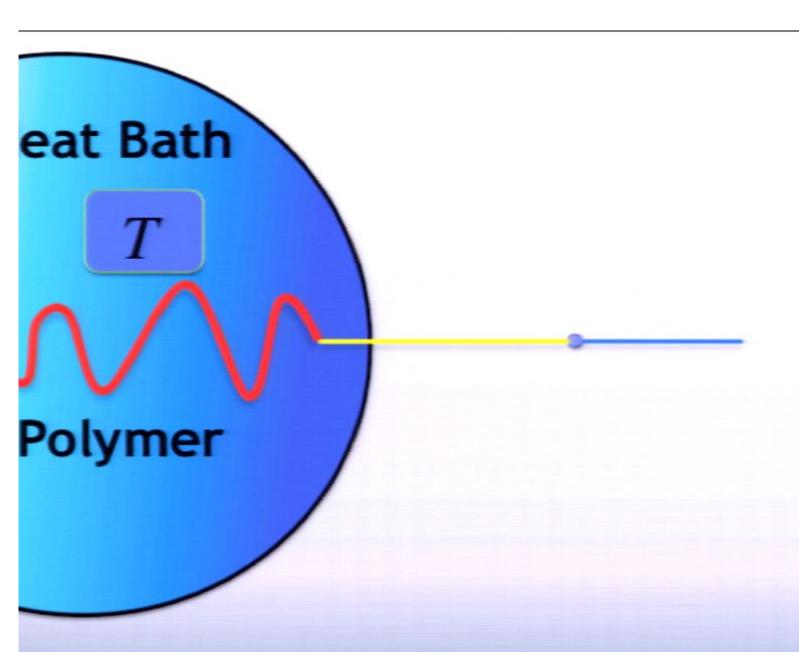
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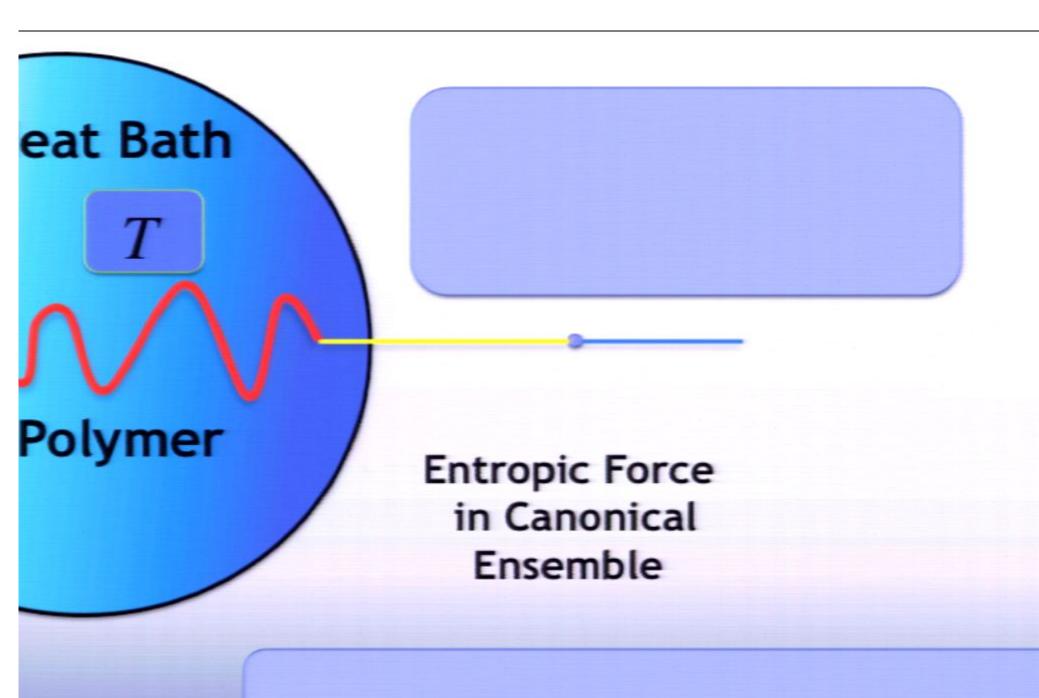
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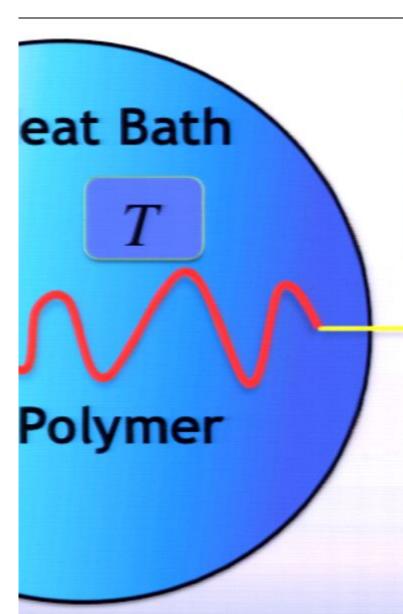
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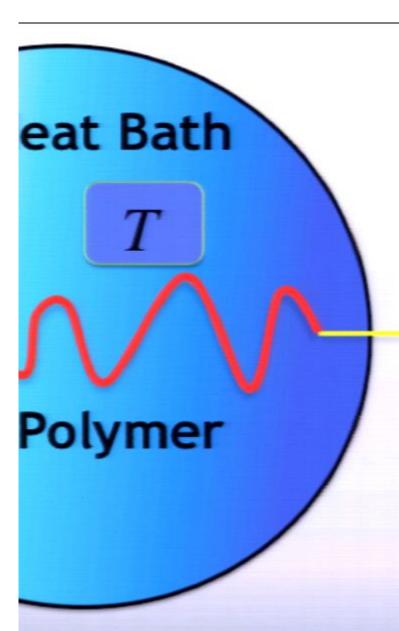
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$$1/T = \partial_E S(E, x)$$
$$F/T = \partial_x S(E, x)$$

in Canonical Ensemble



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Saddle point equations

in Canonical Ensemble

$$Z(T,F) = \int dE dx \ \Omega(E,x) \ e^{-(E+Fx)/K}$$

Consider a fast system with phase space volume

$$\Omega(E,x)$$

but we want to integrate out its DOF and describe it effectively by a small number of degrees of freedom.

How would we do it?

We have to make sure that the reaction forces remain the same

Replace all degrees of freedom by just one

$$\int \Omega(E,x) \ e^{-E/T} = \int [dp \, dq] e^{I[p,q;x]/\hbar}$$

$$I[p,q;] = \oint (pdq - H(p,q;x)dt)$$

Action/Entropy correspondence

$$S=2\pirac{J}{\hbar}$$

$$T = \frac{\hbar\omega}{2\pi}$$

Action/Entropy correspondence in Gravity

$$S = 2\pi \frac{J}{\hbar}$$

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Gibbons-Hawking-Wald formalism

$$I = 2\pi(J - E/\omega) = S - E/T$$

$$I = \frac{2\pi}{\omega} \left[\int_{hor} Q - \int_{\infty} Q \right]$$

Gravity is an adiabatic reaction force

Due to a fast dynamical system underlying our universe.

Inertia (gravity) is the leading order reaction force.

The separation of time scales breaks down at horizons.

Other forces arise from higher order corrections

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Gravity is an adiabatic reaction force

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Gravity degenerates into an (adiabatic) entropic force.

In this regime the fast system serves as a heat bath

This leads to thermal fluctuations in inertial frame

Dark energy and dark matter are associated with these thermal effects.

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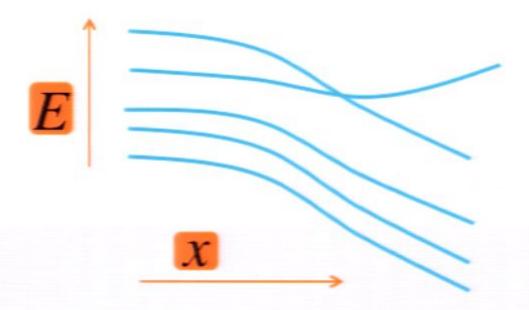
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Berry Phase and Crossing Eigenvalues



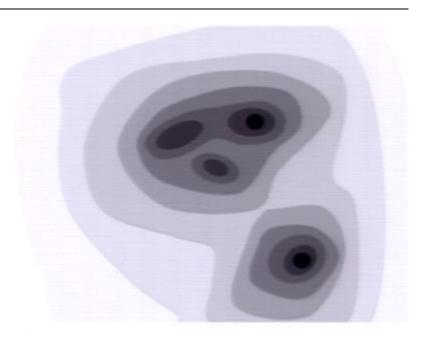
$$H = \begin{pmatrix} z & x + iy \\ x - iy & -z \end{pmatrix} = \vec{x} \cdot \vec{\sigma}$$

$$\vec{B} = \frac{\hat{x}}{4\pi |\vec{x}|^2}$$

Dirac monopool

At the locus of coinciding eigenvalues one can construct a non-abelian Berry connection.

$$A_{ij} = \left\langle \psi_i \middle| d\psi_j \right\rangle$$



Question:

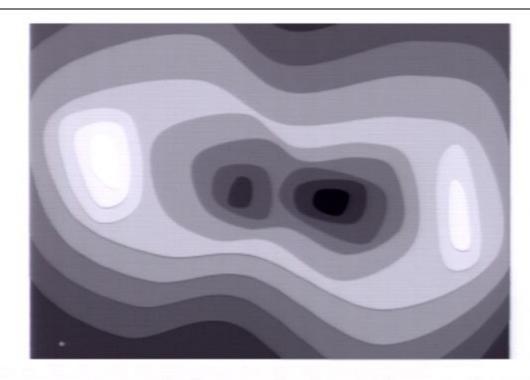
What is the total phase space volume associated with a region of space contained within a particular surface given the amount of energy inside of it?

Assumption:

We consider a quasi static situation and a region that is surrounded by an equipotential surface.

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What is the total phase space volume associated with a region of space contained within a particular surface given the amount of energy inside of it?

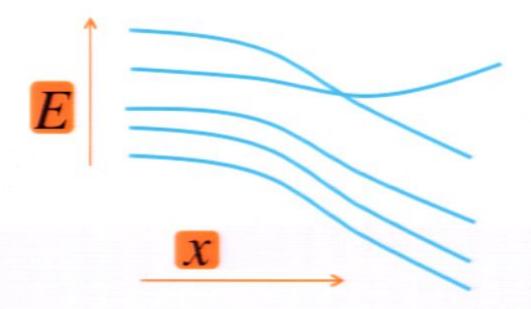
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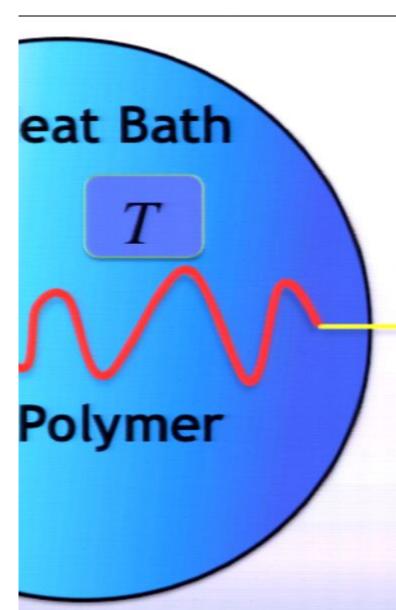
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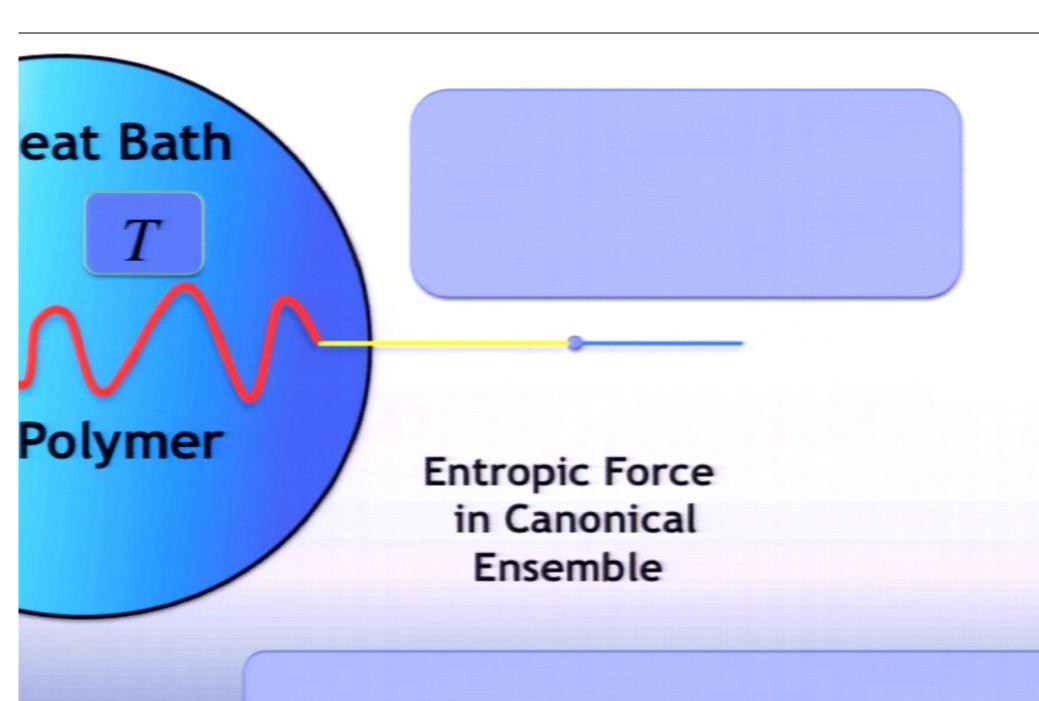
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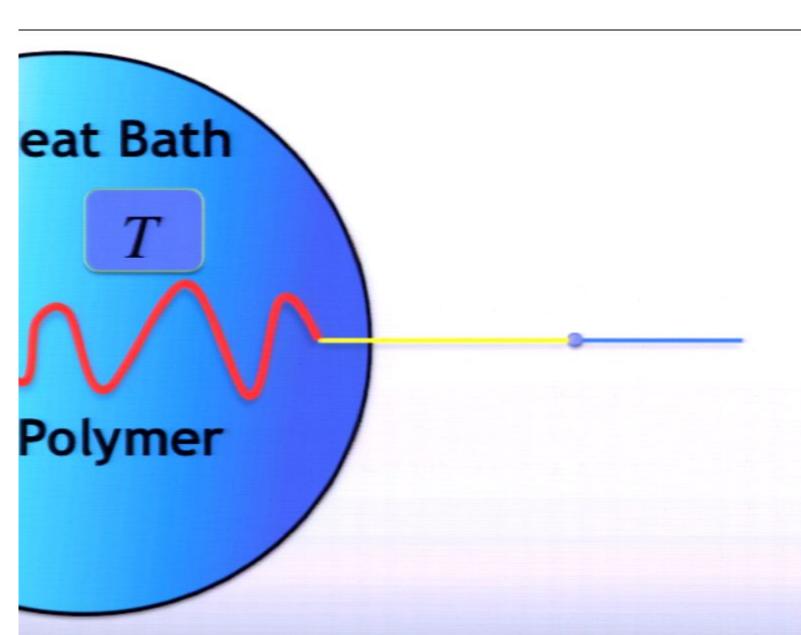
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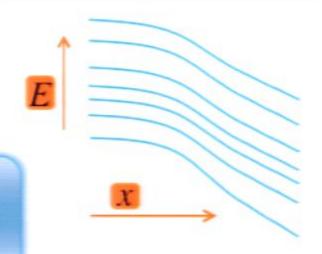




Adiabatic Reaction Force

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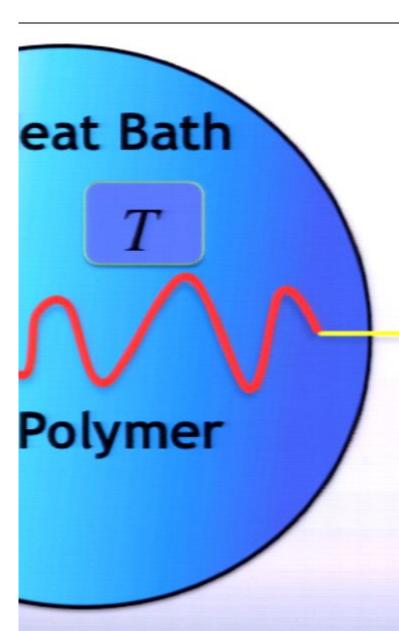
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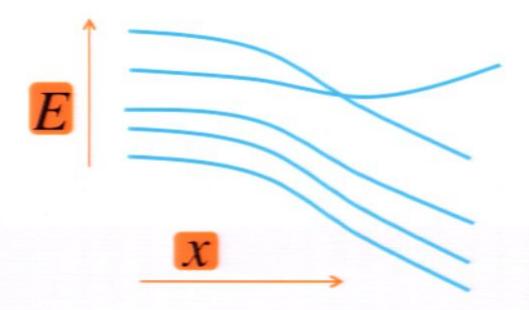
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Berry Phase and Crossing Eigenvalues



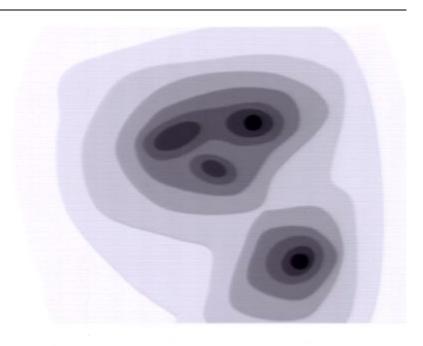
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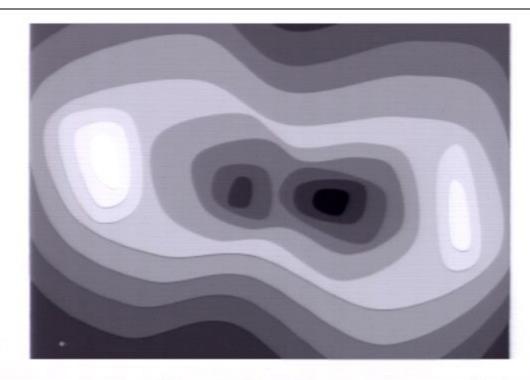
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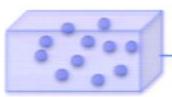
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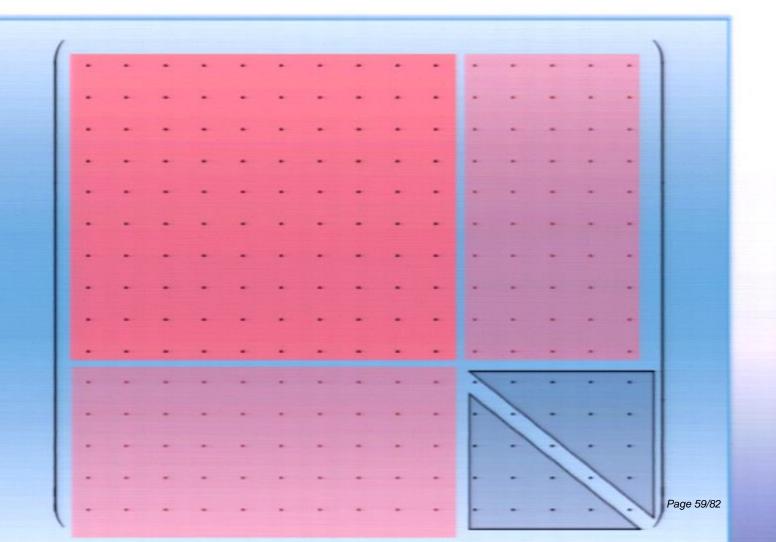
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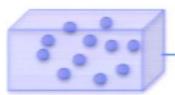
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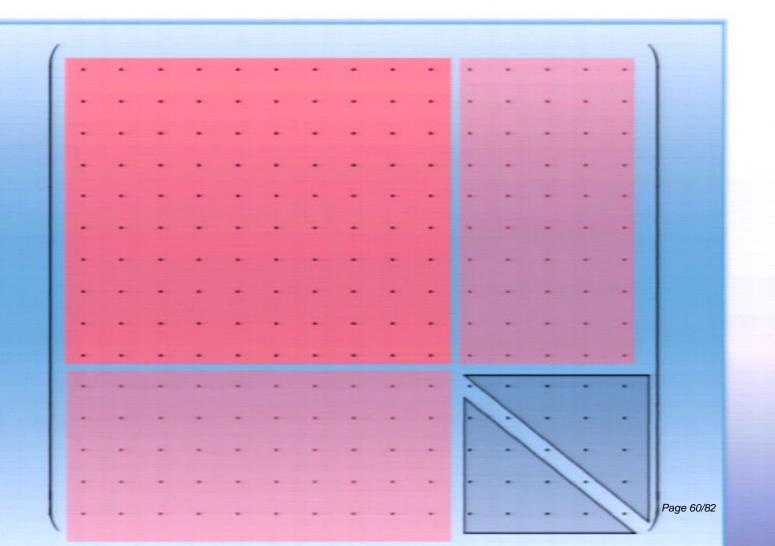
ack Hole



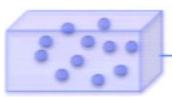


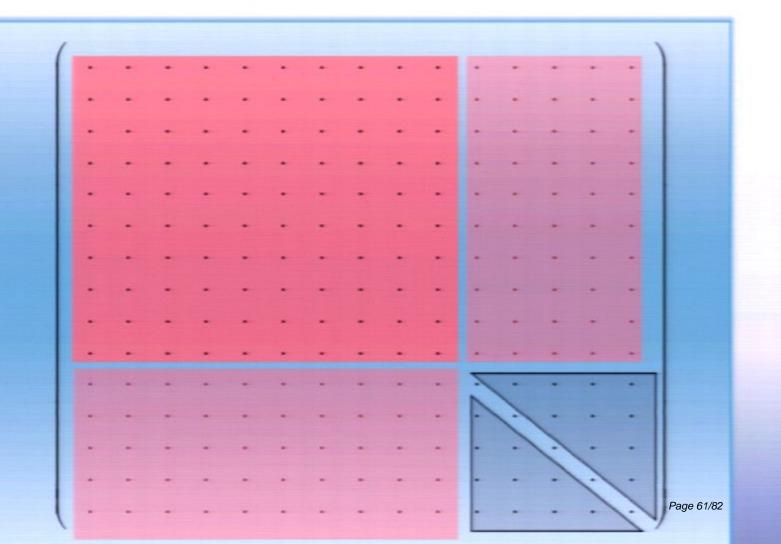
ack Hole





ack Hole





Proposal for phase space volume

$$MR \rightarrow 2L_0$$

$$\frac{A}{4G} \rightarrow 2\pi \frac{c}{12}$$

(Brown Hennaux)

Take Cardy formula

$$S = 2\pi \sqrt{\frac{c}{6}L_0}$$

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Proposal

$$\frac{24L_0}{c} = \frac{16\pi GMR}{A} = 2\Phi$$

For equipotential surfaces (= constant redshift)

$$S = \frac{A}{4G}v$$

$$2\Phi = v^2$$

Proposal for phase space volume

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Proposal for phase space volume

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$$\frac{A}{4G} \to 2\pi \frac{c}{12}$$

(Brown Hennaux)

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Proposal

$$\frac{24L_0}{c} = \frac{16\pi GMR}{A} = 2\Phi$$

For equipotential surfaces (= constant redshift)

$$S = \frac{A}{4G}v$$

$$2\Phi = v^2$$

Cosmological lorizon

De Sitter Space

$$T = \frac{H_0}{2\pi}$$

$$S = \frac{A}{4G}v$$

$$v = H_0 R$$

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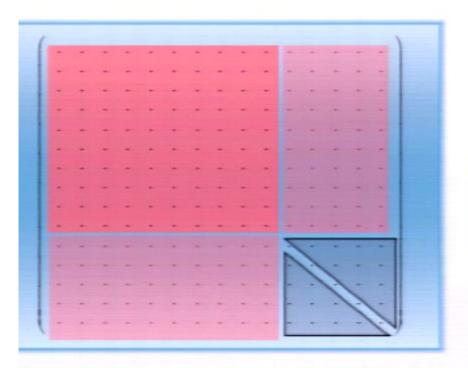
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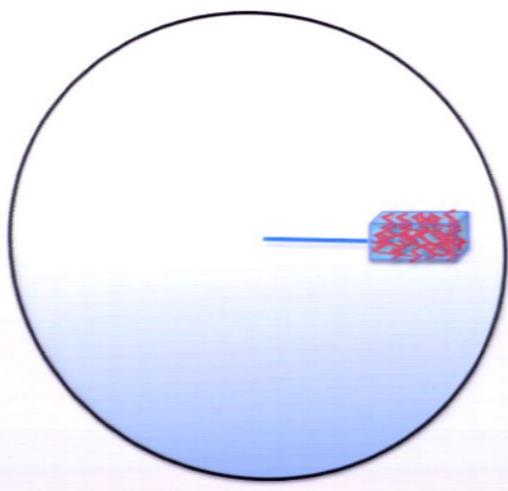
$$X = \begin{pmatrix} x_{11} & \dots & x_{1N} & z_1 \\ \vdots & \vdots & \vdots & \vdots \\ x_{N1} & \dots & x_{NN} & z_N \\ \hline z_1^* & \dots & z_N^* & xI \end{pmatrix}$$

$$F = T \frac{\partial S}{\partial x}$$

$$T = \frac{\hbar H_0}{2\pi} \implies \frac{\partial S}{\partial x} = 2\pi \frac{mv}{\hbar}$$
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$$v = H_{0r}$$

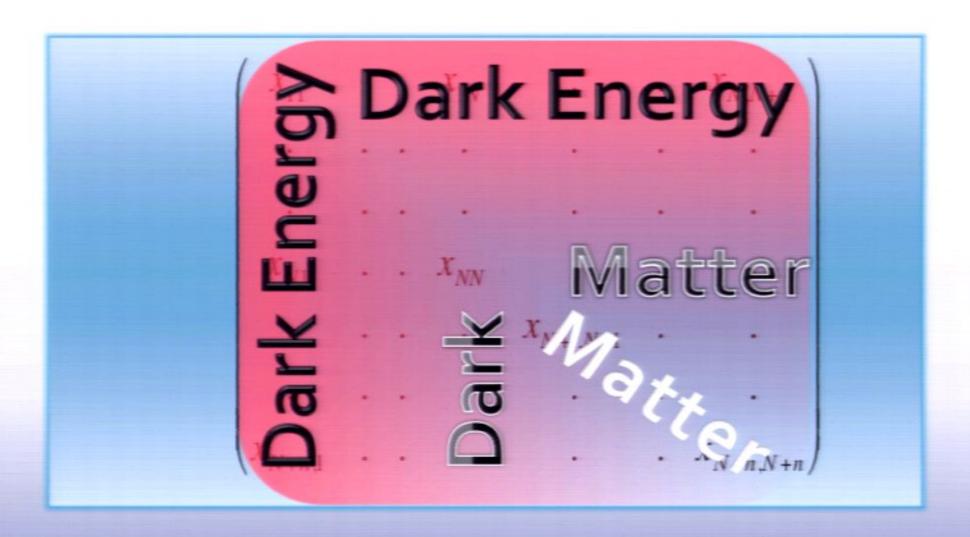




$$N = \frac{McR}{\hbar}$$

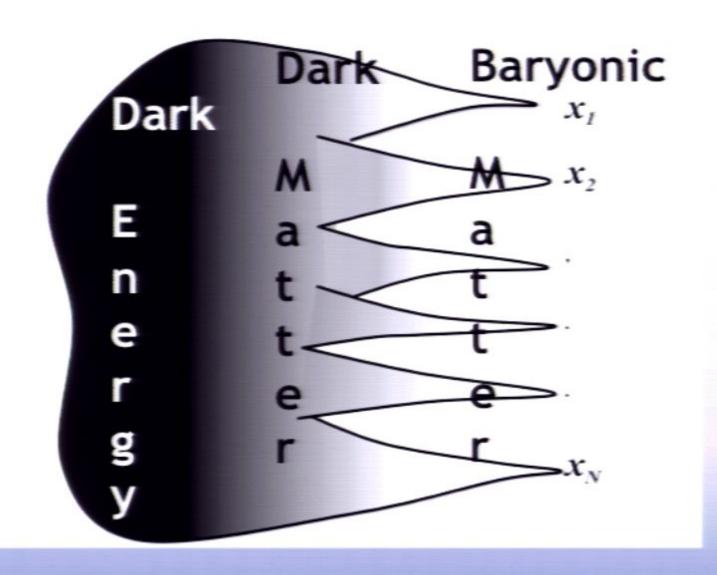
 $T = \frac{\hbar H_0}{2\pi}$

At horizons space and time dissappear,



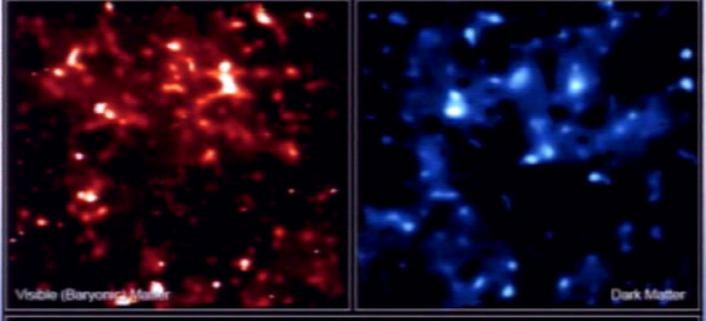
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PHASE SPACE



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Distribution of Visible and Dark Matter • Cosmic Evolution Survey

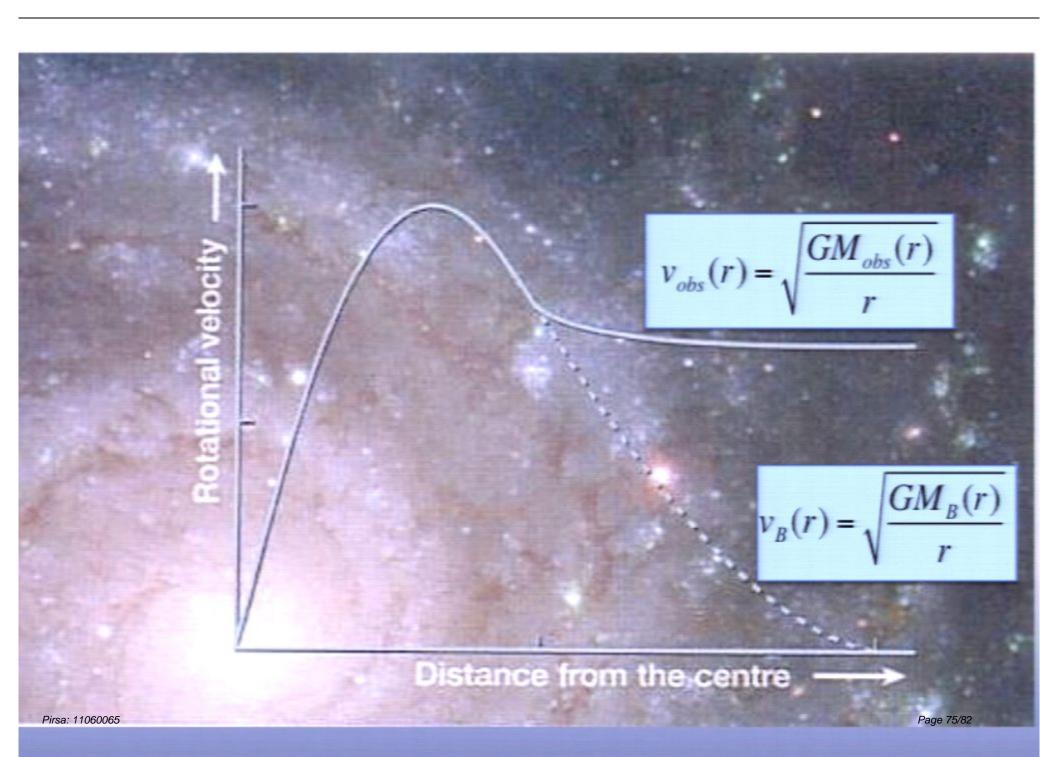
Hubble Space Telescope • Advanced Camera for Surveys

Galaxy rotation curves

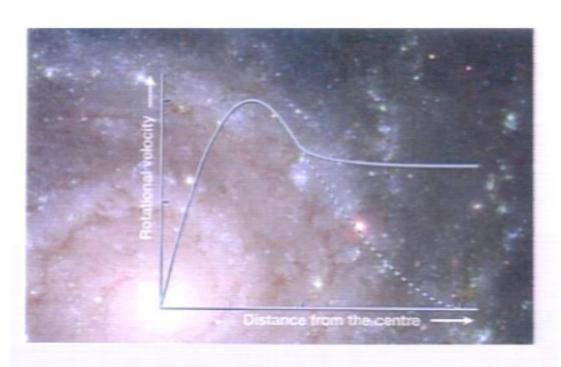


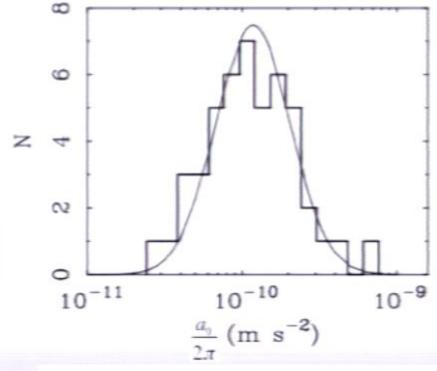
$$\frac{v(r)^2}{r} = \frac{GM(r)}{r^2}$$

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Baryonic Tully-Fisher relation

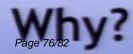


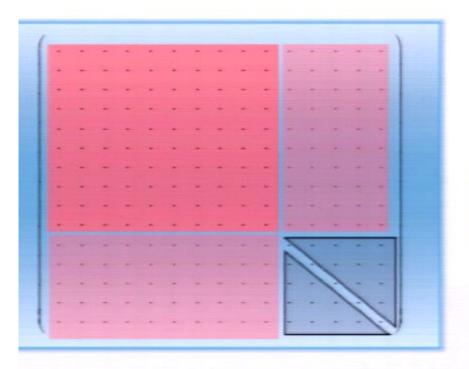


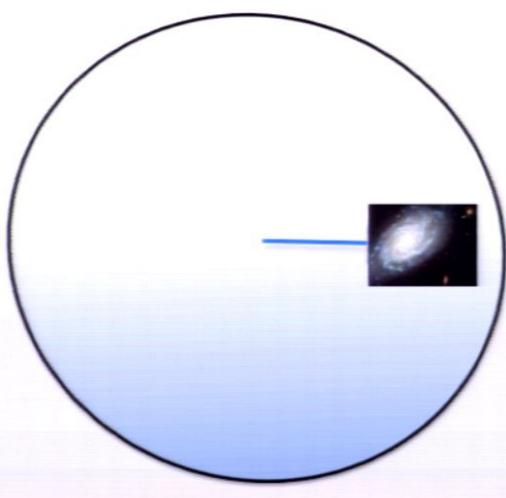
$$V_{obs}^4 = GM_B \frac{a_0}{2\pi}$$

$$\frac{a_0}{2\pi}$$
 = 1.24 ± 0.14 · 10⁻¹⁰ m/s

$$a_0 \approx cH_0$$







$$N = \frac{McR}{\hbar}$$

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 $T = \frac{\hbar H_0}{2\pi}$

Fluctuations in the gravitational field

$$E_{grav} = \frac{1}{8\pi G} \int |\nabla \Phi|^2$$

$$\langle \nabla \Phi \nabla \Phi \rangle = \int [d\Phi] e^{-\frac{1}{8\pi GkT} \int |\nabla \Phi|^2} \nabla \Phi \nabla \Phi$$

$$\left\langle \nabla \Phi \nabla \Phi \right\rangle = \frac{GkT}{d^3}$$

d = short distance cut off

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Fluctuations in the gravitational field

$$\left\langle \nabla \Phi \nabla \Phi \right\rangle = \frac{GkT}{d^3}$$

d = short distance cut off

$$N = \frac{R^3}{d^3}$$

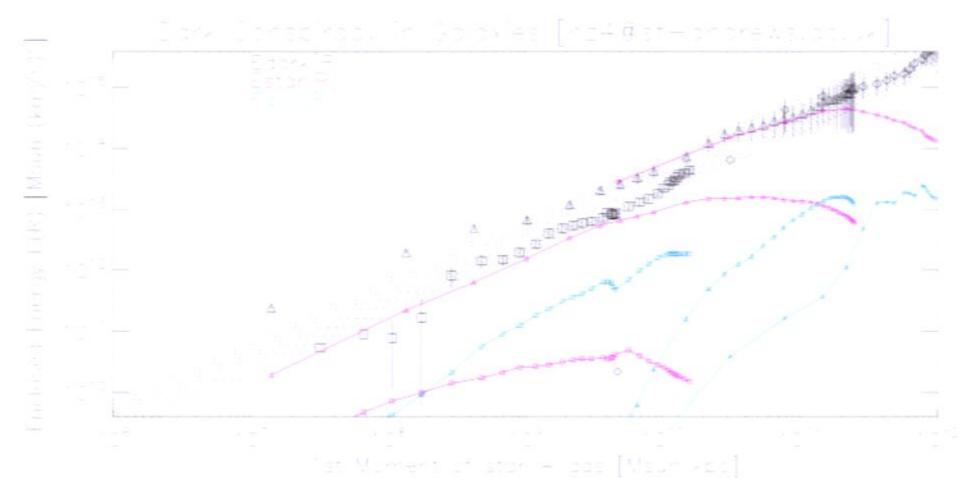
N= maximum number of degrees

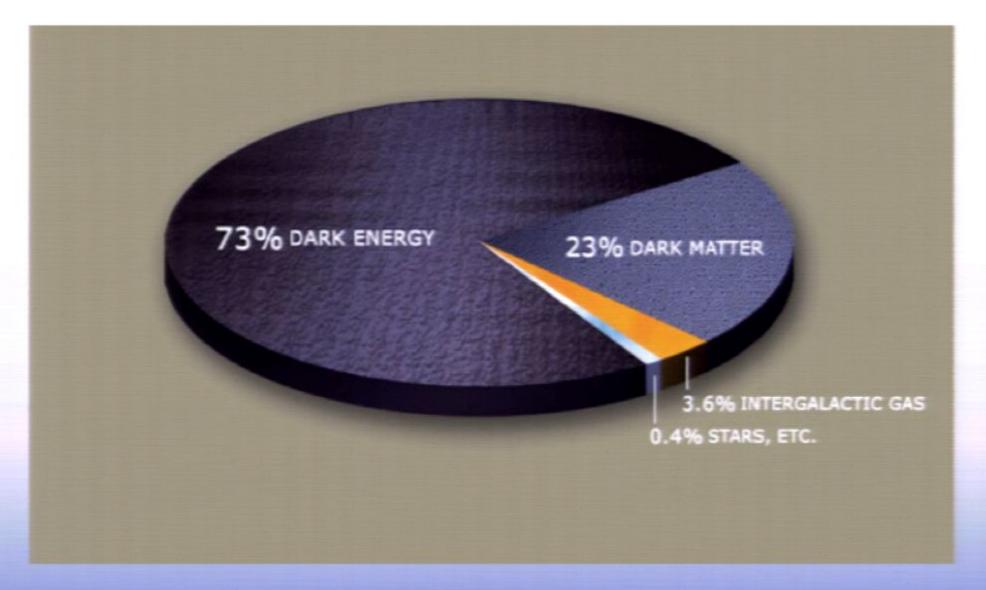
Take N = Bekenstein bound

$$N=rac{McR}{\hbar}$$

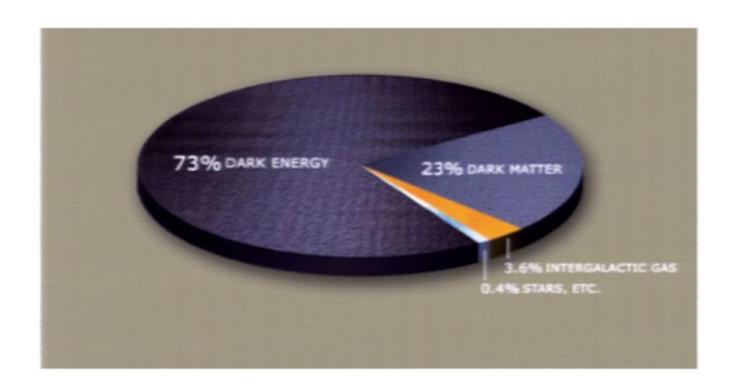
$$\left\langle \nabla \Phi \nabla \Phi \right\rangle = \frac{GMkT}{\hbar R^2}$$

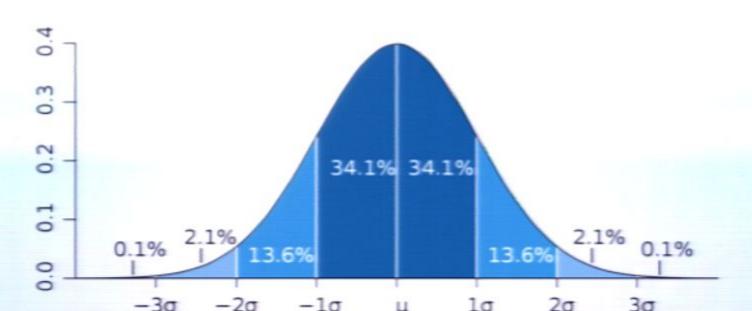
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