

Title: A toy model of quantum gravity in cosmological spacetimes

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URL: <http://pirsa.org/11060058>

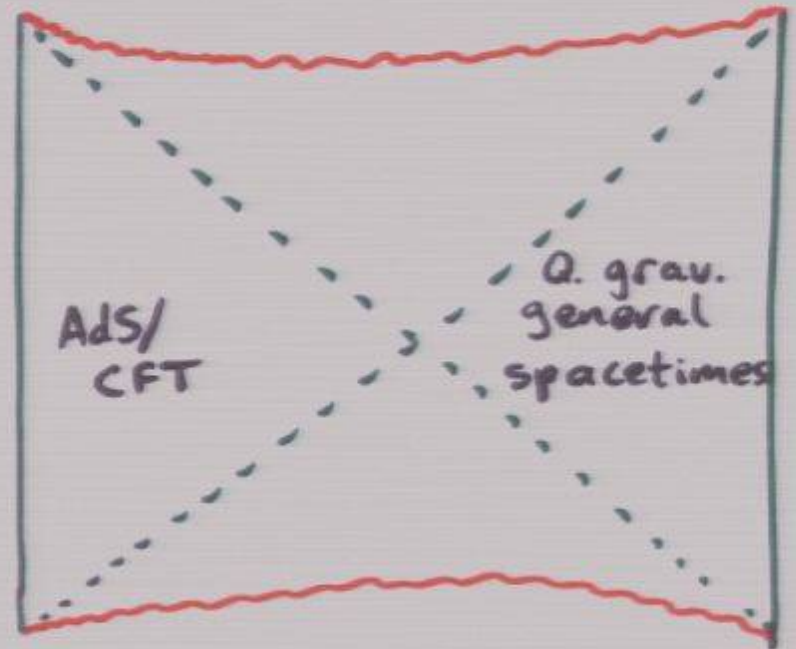
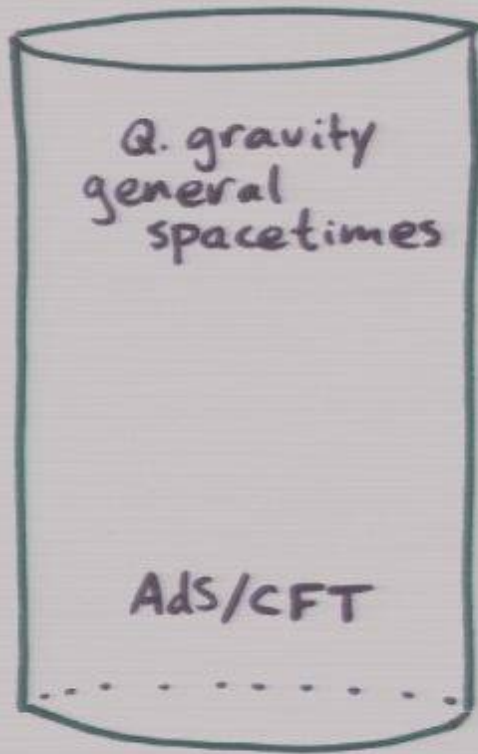
Abstract: In this talk, I attempt to gain insight into the description of quantum gravity on cosmological spacetimes by considering the physics of families of accelerating observers in spacetimes which admit non-perturbative descriptions vis AdS/CFT.

BIG QUESTION: What does a nonperturbative description of quantum gravity for general spacetimes look like?

WE KNOW: Nonperturbative description of quantum gravity for asymptotically AdS, etc...

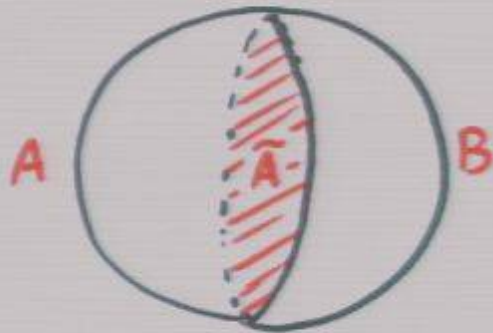
Does this help?

2 possibilities:



(conjectural) Lessons from AdS/CFT

I) Spacetime emerges from entangling degrees of freedom in the nonperturbative description



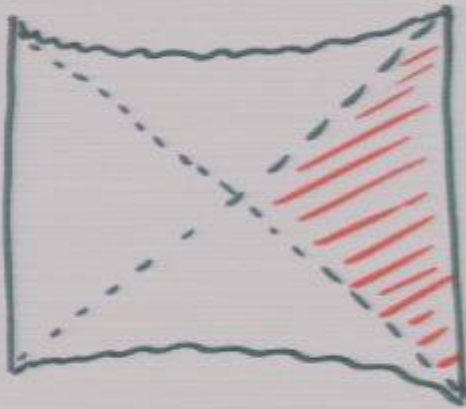
$$\sum e^{-\beta E_i} \langle E_i | \otimes \langle E_i | = \langle \text{X} \rangle$$

$$\frac{\text{area}(\tilde{A})}{4G_N} = \text{entanglement entropy}$$

Ryu, Takayanagi

← Measures of entanglement give direct information about geometry

II) certain density matrices describe (causal) patches of spacetime see 0907.2939

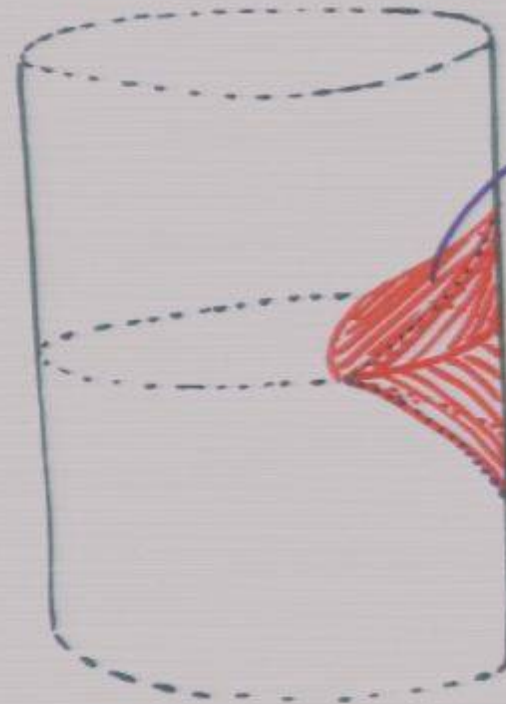


full spacetime:

$$\sum e^{-\frac{\beta E_i}{2}} |E_i\rangle \otimes |E_i\rangle = |\psi\rangle$$

one side:

$$\text{tr}_B |\psi\rangle \langle \psi| = \rho_{\text{Th}}^A$$



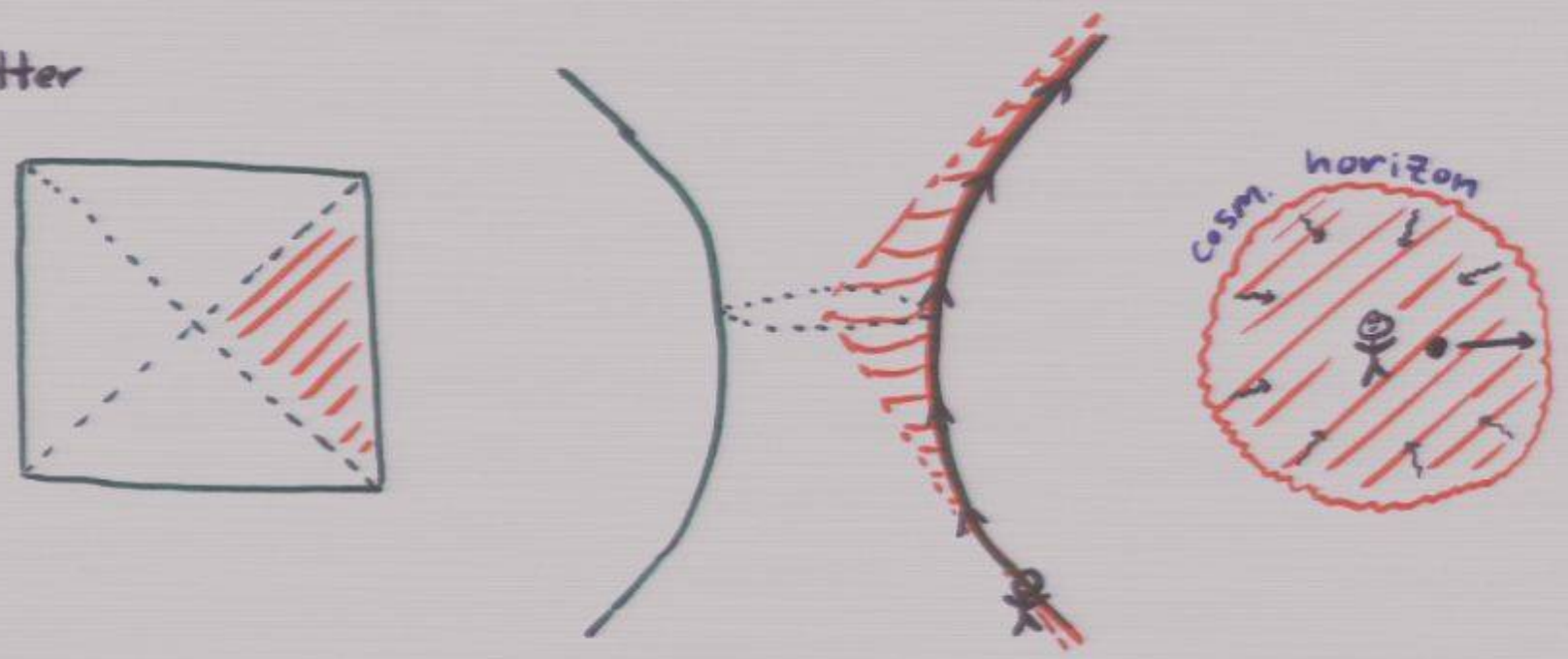
dual to CFT
on $H^d \times \mathbb{R}$
in mixed state
 ρ_{Th}

Von Neumann entropy of ρ :
area of a horizon or
minimal surface

COSMOLOGICAL SPACETIMES

- No causal patch has complete information about spacetime

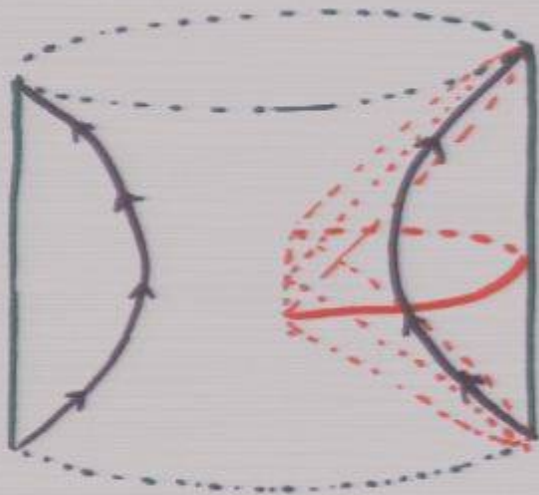
e.g. de Sitter



Plausible that data accessible to single observer encoded in density matrix for some system

AN AdS ANALOGY

- consider family of accelerating observers in AdS



- Each one: causal patch -
"hyperbolic black hole"

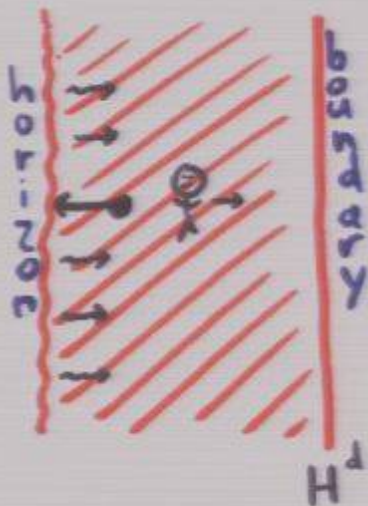
↓ dual to

thermal state of
CFT on $H^d \times \text{time}$

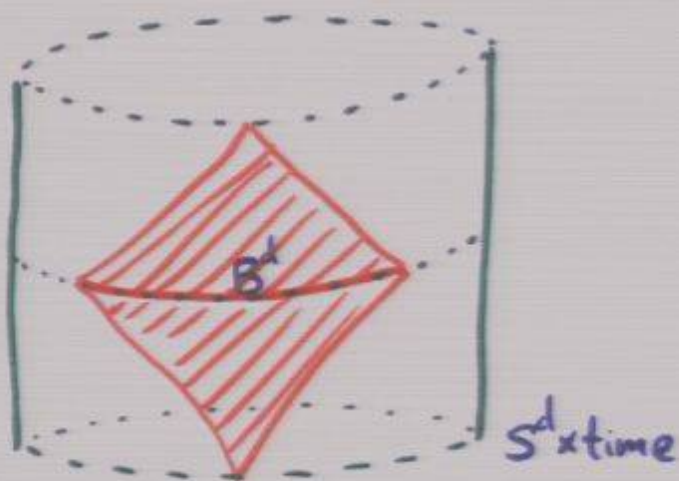
e.g. Myers, Sinha

each patch described by
density matrix ρ_{Th} for
CFT on H^d

$S(\rho) \sim$ horizon area



More general: arbitrary state for global spacetime



- still have description via family of patches
- "causal development" of B^d conformal to $H^d \times \text{time}$

e.g. Casini, Huerta, Myers

\exists map

$$|\psi\rangle_{S^d} \xrightarrow{\text{trace over } B^d} \rho_{B^d} \xrightarrow{\text{unitary}} \rho_{H^d}$$

Patchwise description of spacetime via family of density matrices w. constraints

$$\begin{array}{l} \text{entanglement} \\ \text{entropy} \end{array} S(\rho_{B^d}) = S(\rho_{H^d}) \begin{array}{l} \text{thermal} \\ \text{entropy} \end{array}$$

Density matrices & entanglement in QFT

Recent work: focused on density matrix for spatial region



$$\text{Locality} \Rightarrow \mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_{\bar{A}}$$

define ρ_A

calculate $S(\rho_A)$ (divergent)

$$I(A, B) = S(A) + S(B) - S(A, B) \quad \text{finite}$$

MUTUAL INFORMATION

CONNECTION TO WILSONIAN PICTURE

Usually: description of low-energy physics via
Wilsonian effective action $S_W(\phi_{|\vec{p}|<\mu})$

$|\vec{p}|<\mu$ d.o.f. in mixed state

$$\langle \hat{\phi} | \rho_{|\vec{p}|<\mu} | \tilde{\Phi} \rangle = \frac{1}{Z} \int_{\phi(\tau \rightarrow -\infty) = \hat{\phi}}^{\phi(\tau \rightarrow \infty) = \tilde{\Phi}} [d\phi_{|\vec{p}|<\mu}] e^{-S_W^\mu}$$

$$\int [d\phi] \Theta_\phi e^{-S_W} = \text{Tr}(\rho^\mu \Theta_\phi)$$

NEW OBSERVABLES IN QFT

Given density matrix for subset of momenta, can define observables based on $\text{spec.}(\rho)$

examples: $S(|\vec{p}| < \mu)$

$I(\vec{p}, \vec{q})$

$\text{spectrum}(\rho(\vec{p}))$

Try to calculate these in perturbation theory for weakly coupled field theories

QUESTION: Given a QFT, is there a precise notion of a density matrix describing the IR physics?

work with V. Balasubramanian, M. McDermott

WARNING: the rest of the talk has nothing to do with cosmology or AdS/CFT.

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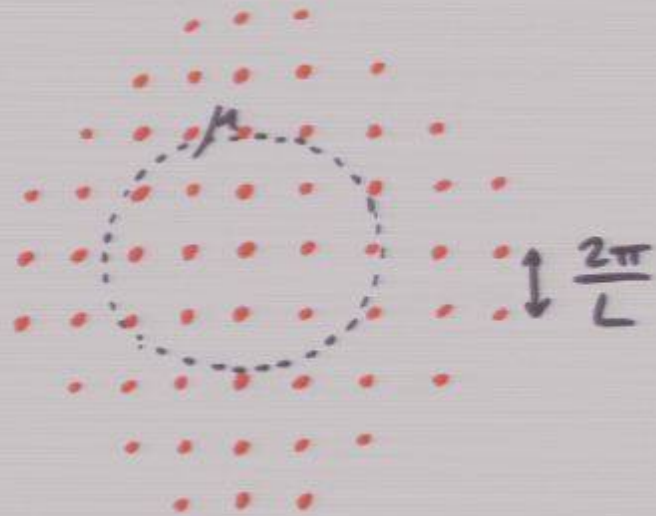
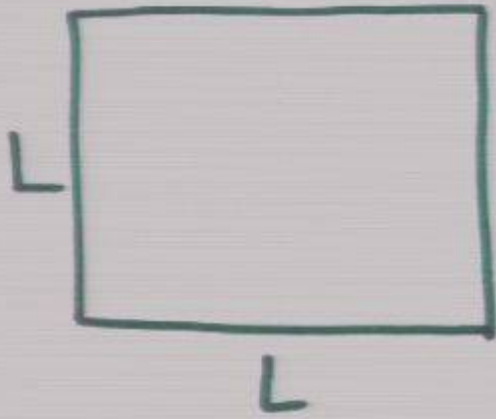
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Start with finite volume:



$$\mathcal{H} = \bigotimes_{p_i} \mathcal{H}_{p_i}$$

Basis $| \{n_i\} \rangle = (a_{p_i}^\dagger)^{n_i} | 0 \rangle$

$\lambda = 0$: vacuum $\bigotimes_{p_i} | 0 \rangle$ No entanglement

$$\lambda \text{ small: } \mathcal{H}_{\text{OFT}} = \mathcal{H}_{|p| < \mu} \otimes \mathcal{H}_{|p| > \mu}$$

$$H = H_{|p| < \mu} + H_{|p| > \mu} + \lambda H_{\text{int}}$$

$$|\Omega\rangle = |0\rangle \otimes |0\rangle + \lambda |n_i\rangle \otimes |N_i\rangle e_{nN} + \dots$$

Compute $|\Omega\rangle$, $\rho_{|p| < \mu}$, $S(\mu)$

in PERTURBATION THEORY

Leading order:

$$S = \sum_{\substack{n \neq 0 \\ N \neq 0}} \frac{|\langle n_i, N_i | H_{int} | 0, 0 \rangle|^2}{E_{n,N}^2} \{-\lambda^2 \log \lambda^2\} + \mathcal{O}(\lambda^2)$$

states with at least one particle $|p| < \mu$ + at least one w. $|p| > \mu$

Infinite volume:

$$S = -V \lambda^2 \log \lambda^2 \int_0^\infty d\tau \int d^d x \langle \mathbb{H}_I(-i\tau, x) P(\mu) \mathbb{H}_I(0, 0) \rangle + \mathcal{O}(\lambda^2)$$

Projects to states with

