

Title: Electromagnetic self-force as a cosmic censor

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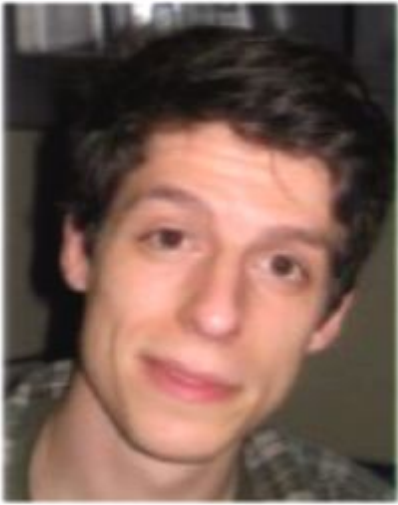
URL: <http://pirsa.org/11050063>

Abstract: Hubeny identified a scenario in which a charged particle falling toward a near-extreme Reissner-Nordstrom black hole can penetrate the black hole and drive it beyond the extremal limit, thereby giving rise to an apparent violation of cosmic censorship. A version of this scenario, relevant to a Kerr black hole and involving a particle with orbital and/or spin angular momentum, was recently examined by Jacobson and Sotiriou (following up on earlier work by Hod); here also the black hole is driven beyond the extremal limit. The Hubeny analysis was inconclusive, however, because in her scenario the particle crosses the horizon with a near-vanishing acceleration; the test-body acceleration is of the same order of magnitude as the acceleration produced by the particle's own electromagnetic self-force, which was not fully incorporated in the analysis. In this talk we report on our computation of the electromagnetic self-force acting on a charged particle falling radially toward a Reissner-Nordstrom black hole, and we reveal whether the self-force acts as a cosmic censor by preventing the particle from reaching the event horizon.



# Electromagnetic self-force as a cosmic censor

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University of Guelph



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# Outline

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- ▶ Cosmic censorship
- ▶ Electromagnetic self-force
- ▶ Self-force in curved spacetime
- ▶ Self-force in Reissner-Nordström spacetime
- ▶ Conclusion

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**Progress report: no punch line!**

# Violations of cosmic censorship

## • Critical collapse

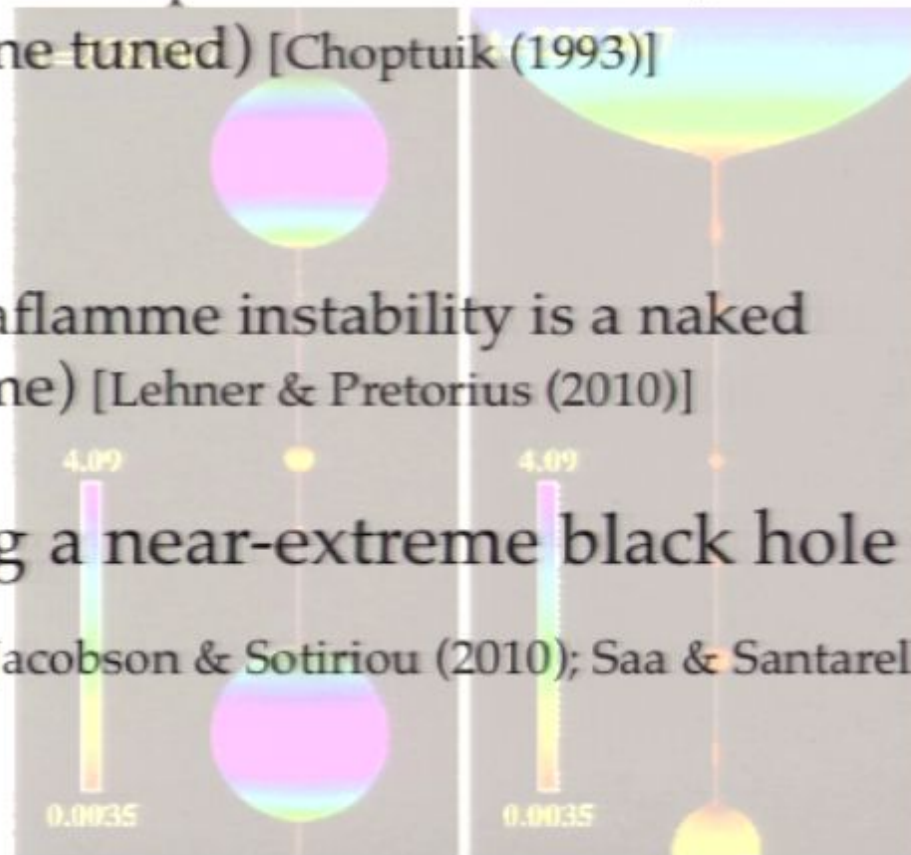
Choptuik's critical solution represents a massless, naked singularity (infinitely fine-tuned) [Choptuik (1993)]

## • Black string

Endpoint of Gregory-Laflamme instability is a naked singularity (5D spacetime) [Lehner & Pretorius (2010)]

## • Overcharging / overspinning a near-extreme black hole

[Hubeny (1999); Hod (2002); Jacobson & Sotiriou (2010); Saa & Santarelli (2011)]





# Violations of cosmic censorship

- Critical collapse

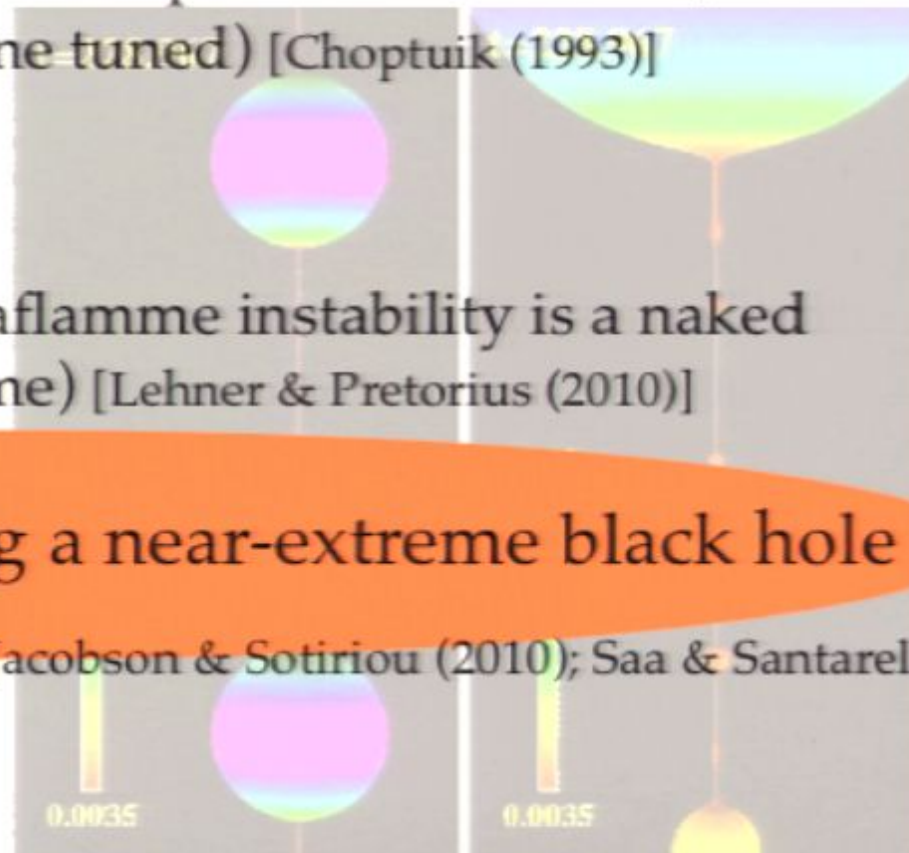
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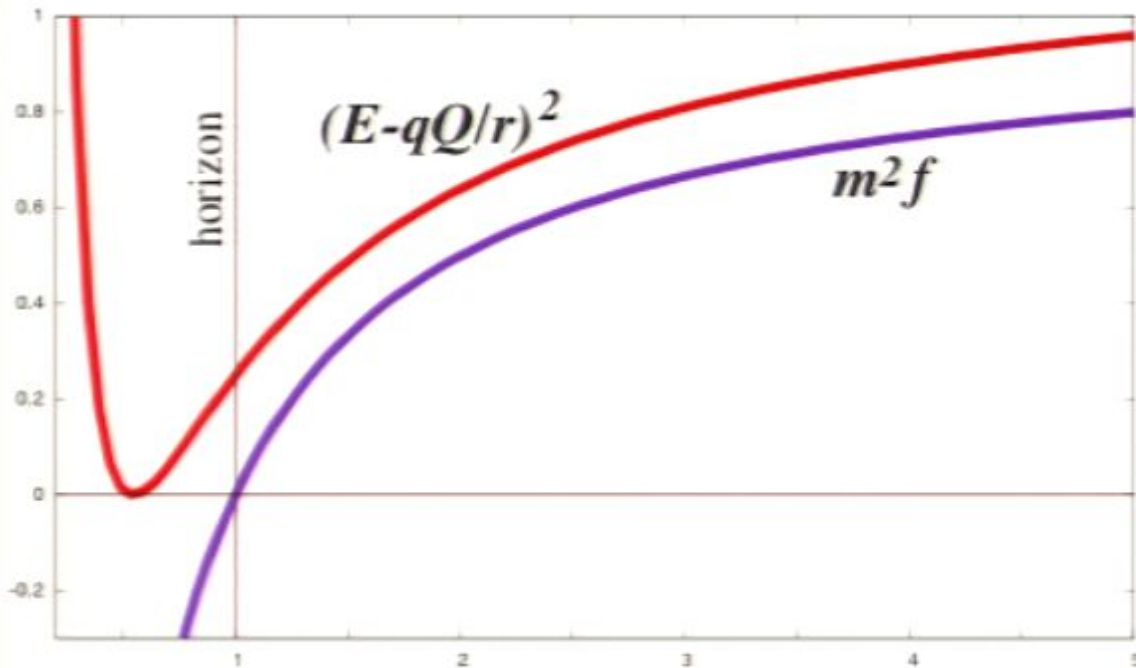
- Overcharging / overspinning a near-extreme black hole

[Hubeny (1999); Hod (2002); Jacobson & Sotiriou (2010); Saa & Santarelli (2011)]



# Hubeny scenario

A charged particle falls toward a near-extreme, Reissner-Nordström black hole...



$$\begin{aligned}
 M &= 1 \\
 Q &= 1 - 2\epsilon^2 \\
 q &= a\epsilon \\
 E &= a\epsilon - 2b\epsilon^2 \\
 m &= c\epsilon
 \end{aligned}
 \qquad
 \begin{aligned}
 a &> 1 \\
 1 &< b < a \\
 c &< \sqrt{a^2 - b^2}
 \end{aligned}$$

Final state is overcharged:

$$Q + q > M + E$$

$$(mr\dot{r})^2 = \left(E - \frac{qQ}{r}\right)^2 - m^2 \left(1 - \frac{2M}{r} + \frac{Q^2}{r^2}\right)$$



# Hubeny scenario

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This analysis ignores the **electromagnetic self-force** acting on the falling particle; the argument is inconclusive.

Incorporating the self-force is important. (This was only partially attempted by Hubeny, based on incomplete information.)

For the Hubeny scenario,

$$m \frac{d^2 r}{d\tau^2} \Big|_{\text{horizon}} \sim \epsilon^2$$

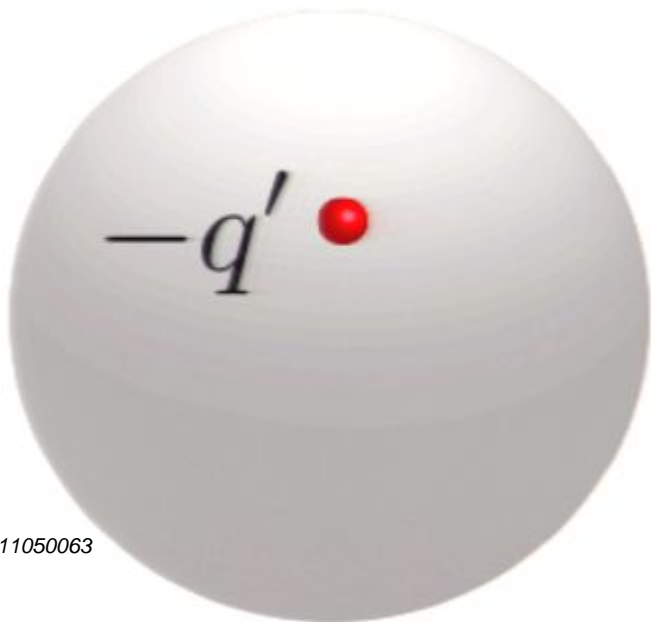
This is the same scaling as the self-force.

# Electromagnetic self-force

There is nothing intuitive about the self-force.

Consider a static charge outside a Schwarzschild hole...

•  $q$



Grounded conductor:

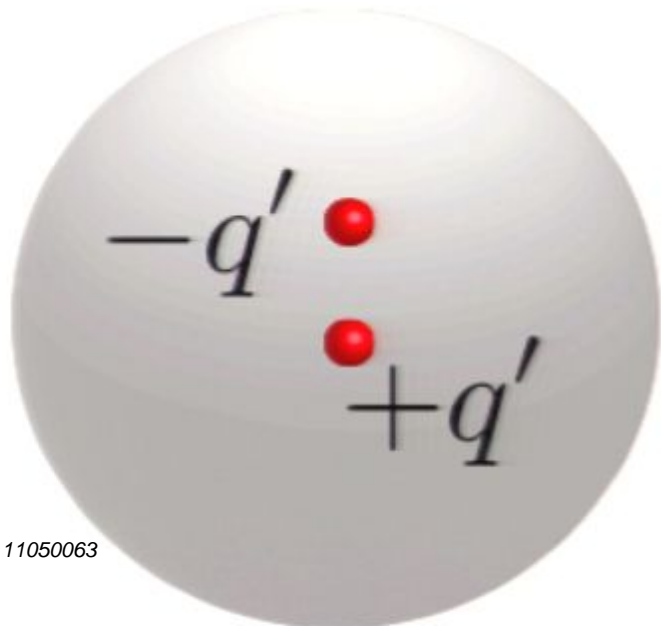
$$f \sim -\frac{q^2 M}{r^3}$$

# Electromagnetic self-force

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Uncharged conductor:

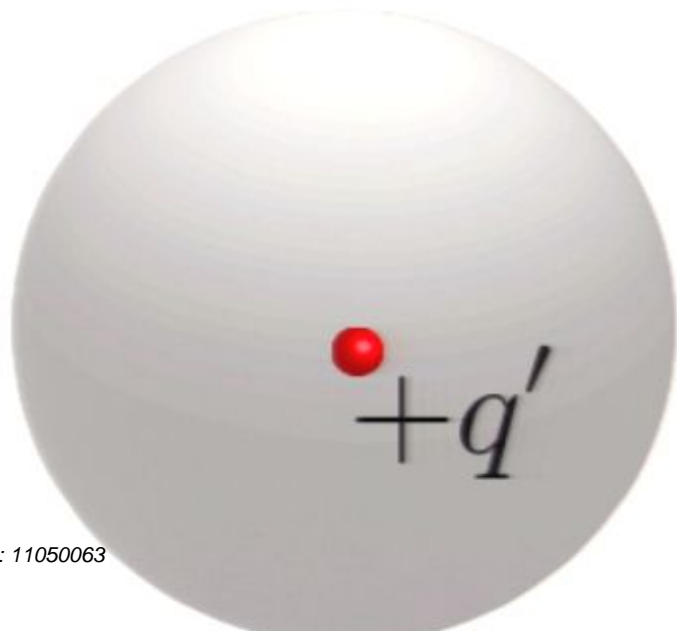
$$f \sim -\frac{q^2 M^3}{r^5}$$

# Electromagnetic self-force

There is nothing intuitive about the self-force.

Consider a static charge outside a Schwarzschild hole...

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Black hole:

[Smith & Will (1980); Zelnikov & Frolov (1982)]

$$f \sim + \frac{q^2 M}{r^3}$$

Against all expectations,  
the self-force is **repulsive**.

# Electromagnetic self-force

There is nothing intuitive about the self-force.

Consider a nonrelativistic point charge falling radially toward another charge...

$$\mathbf{F}_{\text{ext}} = -\frac{qQ}{r^2} \hat{\mathbf{r}}$$

$$\begin{aligned} \mathbf{f}_{\text{self}} &= \frac{2}{3} \frac{q^2}{m} \frac{d\mathbf{F}_{\text{ext}}}{dt} \\ &= \frac{4}{3} \frac{q^3 Q}{m r^3} \mathbf{v} \end{aligned}$$

The self-force **pushes** on the falling charge.



# Electromagnetic self-force

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$$\mathbf{F}_{\text{ext}} = -\frac{qQ}{r^2} \hat{\mathbf{r}} \qquad \mathbf{f}_{\text{self}} = \frac{2}{3} \frac{q^2}{m} \frac{d\mathbf{F}_{\text{ext}}}{dt}$$
$$= \frac{4}{3} \frac{q^3 Q}{m r^3} \mathbf{v}$$

The self-force **pushes** on the falling charge.

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$$\mathbf{f}_{\text{self}} \cdot \mathbf{v} = \frac{2}{3} q^2 \dot{\mathbf{a}} \cdot \mathbf{v} = -\frac{2}{3} q^2 |\mathbf{a}|^2 + \frac{2}{3} q^2 \frac{d}{dt} (\mathbf{a} \cdot \mathbf{v})$$

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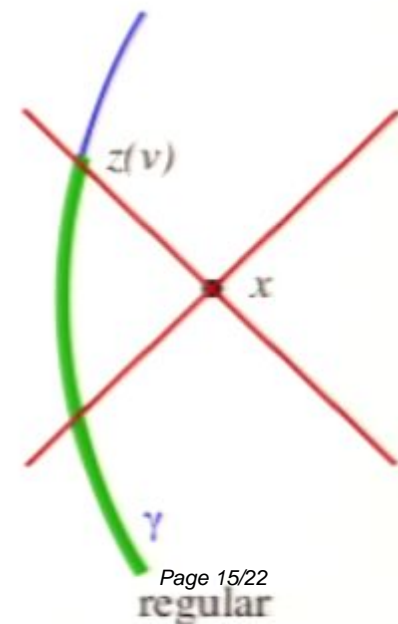
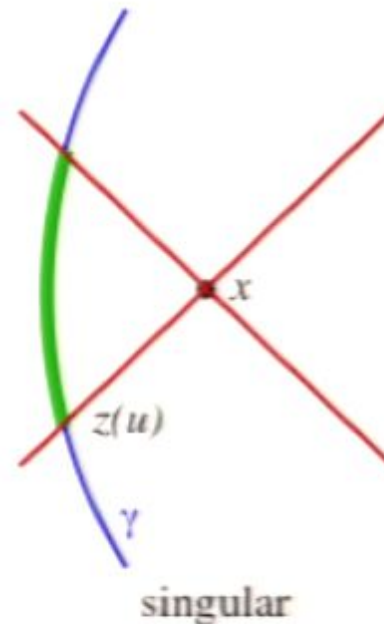
# Self-force in curved spacetime

DeWitt & Brehme (1960):

$$ma^\mu = f_{\text{ext}}^\mu + q^2 (\delta^\mu_\nu + u^\mu u_\nu) \left( \frac{2}{3m} \frac{Df_{\text{ext}}^\nu}{d\tau} + \frac{1}{3} R^\nu_\lambda u^\lambda \right) + 2q^2 u_\nu \int_{-\infty}^{\tau^-} \nabla^{[\mu} G_{\text{ret} \lambda'}^{\nu]}(z(\tau), z(\tau')) u^{\lambda'} d\tau'$$

Detweiler & Whiting (2003):

$$\begin{aligned} F_{\text{ret}}^{\alpha\beta} &= F_S^{\alpha\beta} + F_R^{\alpha\beta} \\ \nabla_\beta F_S^{\alpha\beta} &= 4\pi j^\alpha \\ \nabla_\beta F_R^{\alpha\beta} &= 0 \\ ma^\alpha &= q F_R^{\alpha\beta} u_\beta \end{aligned}$$



# Regularization

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The **singular field** is known locally as an expansion in powers of the distance to the particle.

It can easily be subtracted from the **retarded field**, which is obtained numerically.

The result is a **regular field** that is entirely responsible for the self-force.

# Self-force in Reissner-Nordström

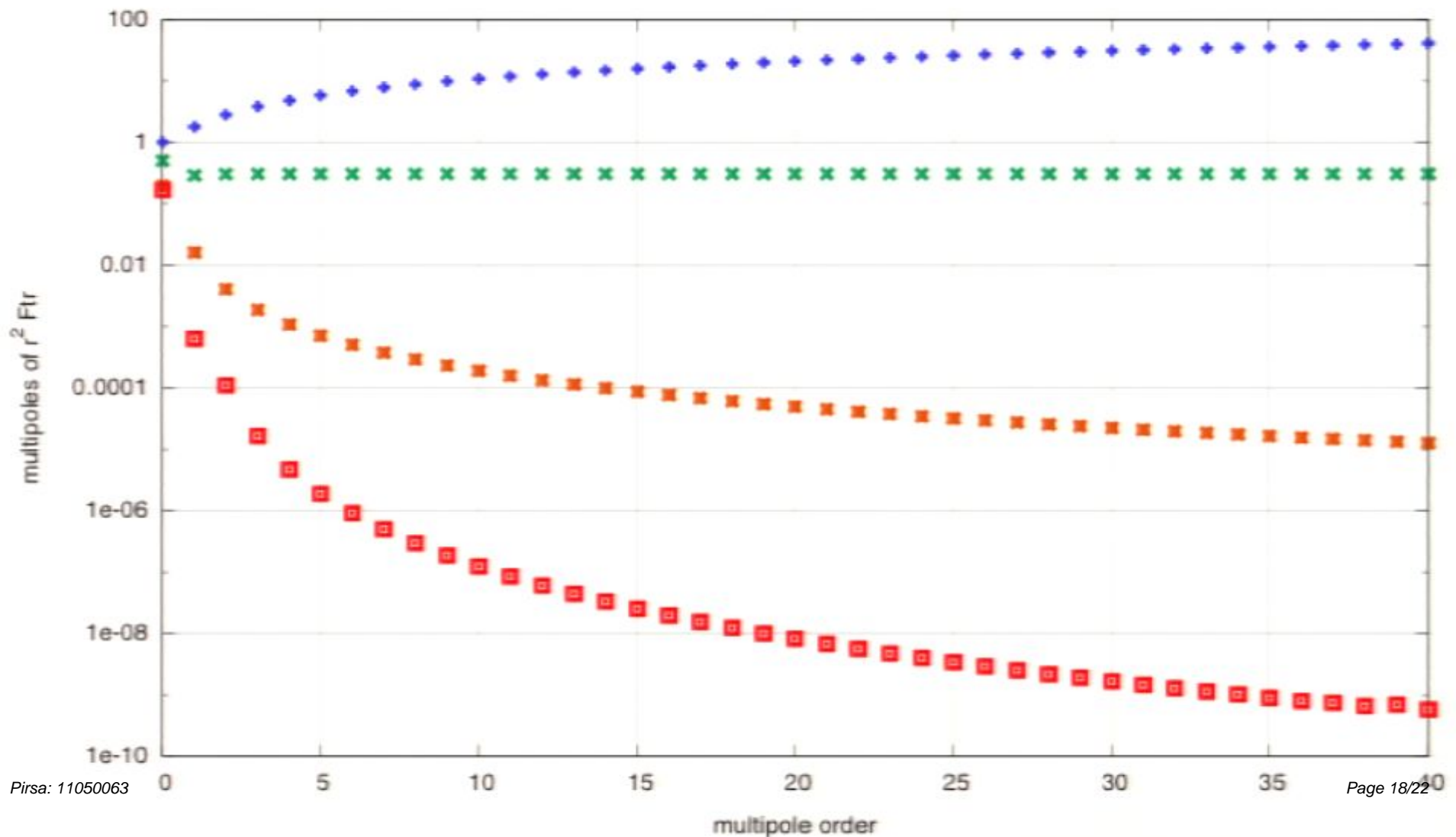
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- ▶ We compute the Maxwell field for a point particle falling radially toward a Reissner-Nordström black hole.
- ▶ The field is decomposed into spherical harmonics, and the reduced equations are integrated with a time-domain, finite-difference code that accounts exactly for the delta-function source.
- ▶ Each spherical-harmonic mode contains a contribution to a **singular field** which exerts no force, and a contribution to a **smooth field** which is entirely responsible for the self-force; the singular contribution is removed (mode-sum regularization).
- ▶ The self-force is computed by summing over all computed modes up to a maximum value of  $l$ ; the large- $l$  tail can be estimated analytically.

- ▶ Mode-sum regularization acts as a robust check on the numerics.

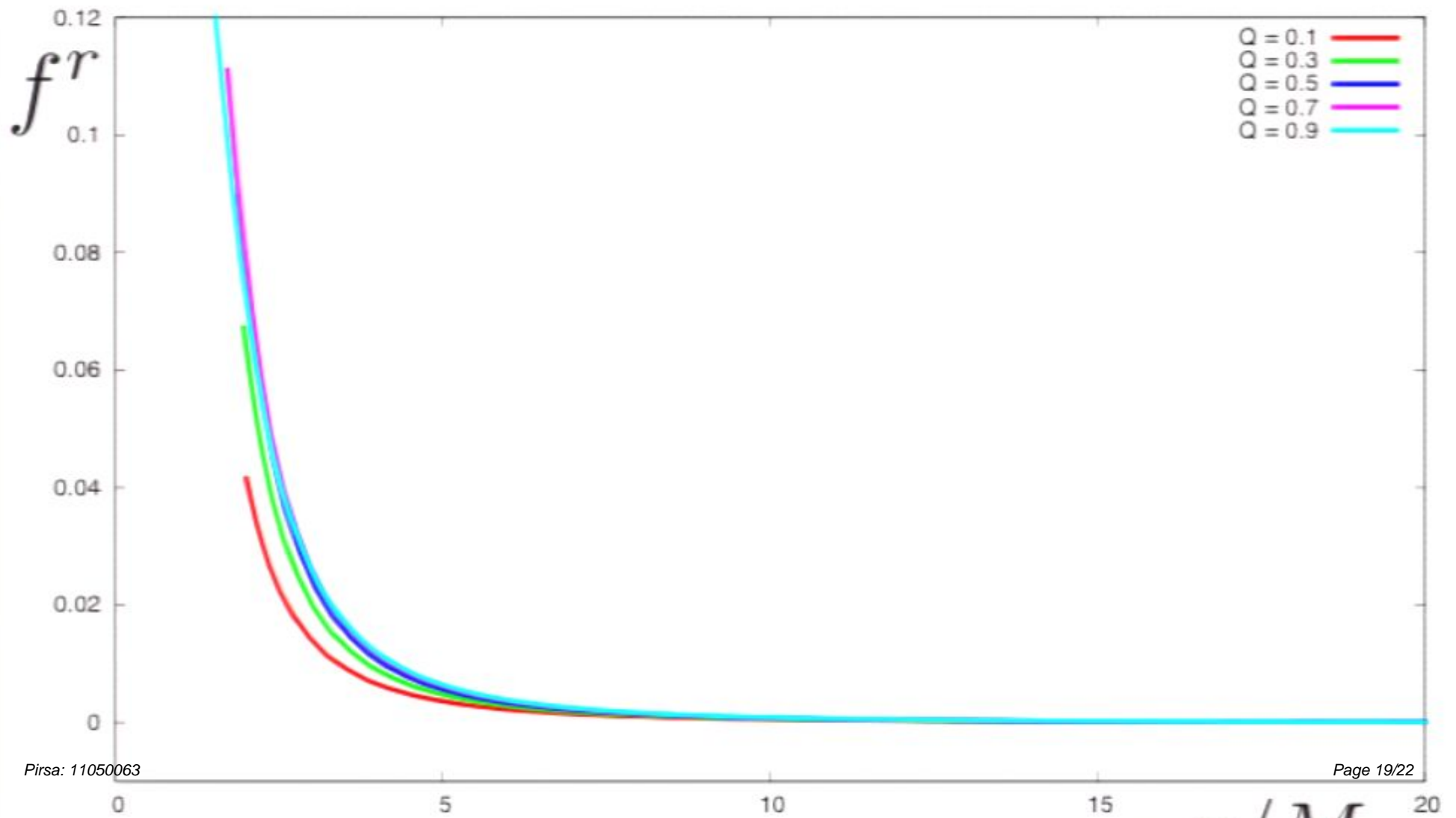


# Mode-sum regularization

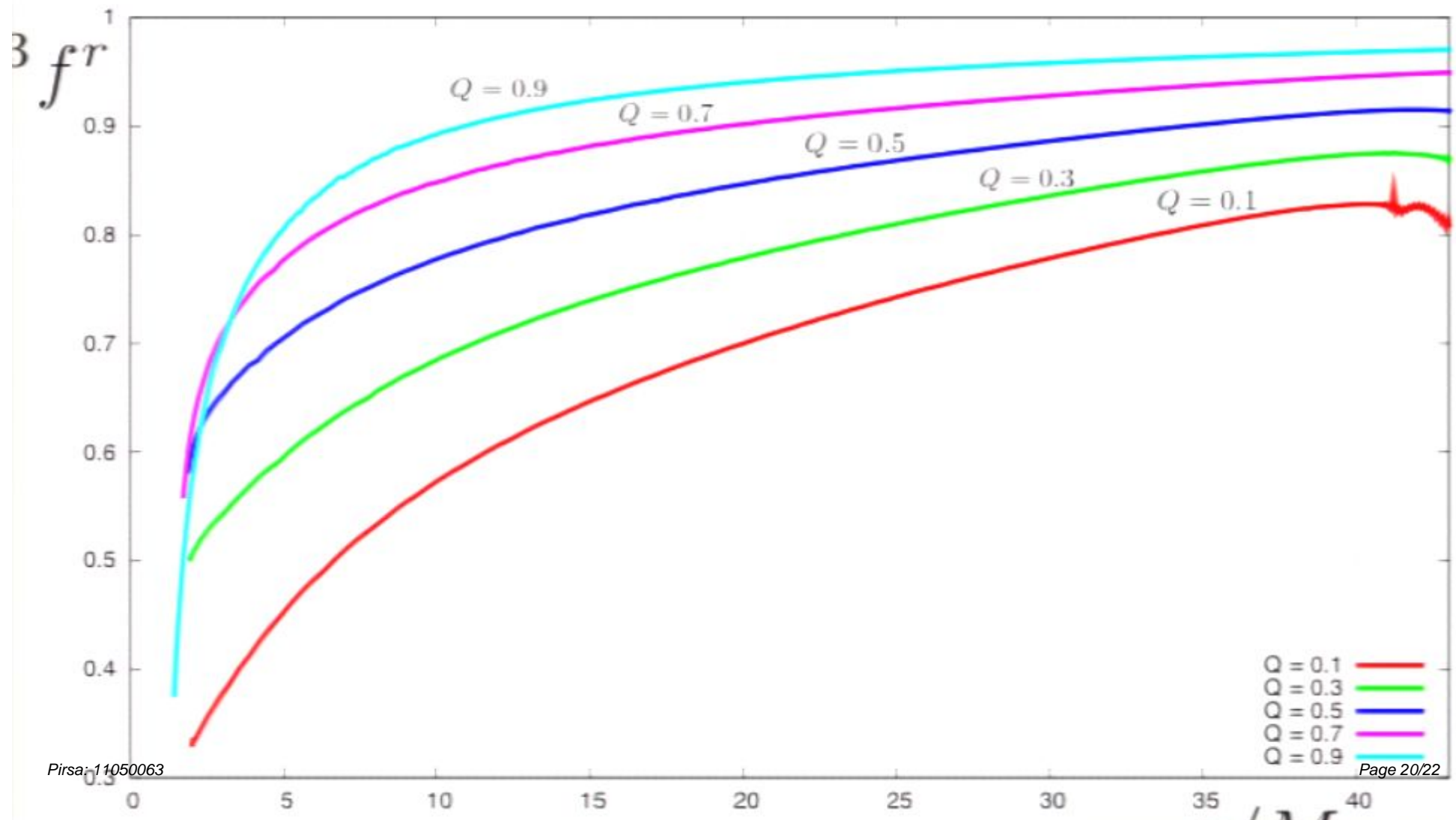




# Results



# Results



# Conclusion

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- ▶ For a point charge falling radially toward a Reissner-Nordström black hole (**starting from rest** --- not a Hubeny trajectory), the electromagnetic self-force is **repulsive and increases with the black-hole charge**.
- ▶ This gives us reason to expect that the self-force may act as a **cosmic censor** to invalidate the Hubeny overcharging scenario (but we are cautious: self-force is not intuitive).
- ▶ Work is currently underway to calculate the self-force for Hubeny trajectories.

# Results

