Title: Almost quantum theory

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Abstract: Modal quantum theory (MQT) is a discrete model that is similar in structure to ordinary quantum theory, but based on a finite field instead of complex amplitudes. Its interpretation involves only the "modal" concepts of possibility and impossibility rather than quantitative probabilities. Despite its very simple structure, MQT nevertheless includes many of the key features of actual quantum physics, including entanglement and nonclassical computation. In this talk we describe MQT and explore how modal and probabilistic theories are related. Under what circumstances can we assign probabilities to a given modal structure?

Almost Quantum Theory

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Denison University

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PIRSA: 10100050 (MDW)

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Axioms

0) Systems exist.

1) Associated with each is a complex vector space $\mathcal{H}$.

2) Measurements correspond to orthonormal bases $|e_i\rangle$ on $\mathcal{H}$.

3) States correspond to density operators ρ on $\mathcal{H}$.

4) Systems combine by tensor producting their vector spaces, $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$.

5) When no measurement is performed, states evolve by unitary maps U.

Chris Fuchs, "*The Oyster and the Quantum*"

- Where does the elaborate mathematical structure of quantum theory "come from"?
- How would quantum theory change if we modified the axioms?
- What is the role of complex numbers in quantum theory?
- Can we develop a "toy model" of quantum theory that is much simpler but has many of the same general features?

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interference
entanglement
non-classical computation
....

An imaginary world

A world without probability

Probability: Events have probabilities

$$0 \leq p(x) \leq 1$$

normalization: $\sum_x p(x) = 1$

In $N \gg 1$ trials, event x occurs N_x times.

With high probability, $p(x) \approx \frac{N_x}{N}$

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Possibility: Some events are possible

$$P = \{x, x', \ldots\}$$

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In N trials, the set of events that occur is $R \subseteq P$

Modal logic explores the ideas of possibility and necessity (propositions p, $\Diamond p$, $\Box p = \sim\Diamond(\sim p)$, etc.).

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Naive connection:

$x \in P \Leftrightarrow p(x) \neq 0$

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Actual Quantum Theory (AQT)

<u>States</u>

Hilbert space H over field C

Pure state is vector $|\psi\rangle$

$$\langle\psi|\psi\rangle = 1$$

<u>Measurement</u>

Orthonormal basis $\{|k\rangle\}$ for H

$$|\psi\rangle = \sum_k \psi_k |k\rangle$$

Probability: $P(k) = |\psi_k|^2$

<u>Time evolution</u>

$$|\psi\rangle \longrightarrow U|\psi\rangle \quad (U \text{ unitary})$$

<u>Composite systems</u>

$$H^{12} = H^1 \otimes H^2$$

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Modal Quantum Theory (MQT)

States

Vector space V over field F

Pure state is vector $|\alpha)$

$$|\alpha) \neq 0$$

Measurement

Basis $\{|k)\}$ for V

$$|\alpha) = \sum_k \alpha_k |k)$$

Possibility: $\alpha_k \neq 0$

Time evolution

$$|\alpha) \longrightarrow T|\alpha) \quad (T \text{ invertible})$$

Composite systems

$$V^{12} = V^1 \otimes V^2$$

"Mobits"

Simplest possible situation: $F = Z_2$ and $\dim V = 2$

Only three possible states $|0)\,, |1)\,, |\sigma) = |0) + |1)$

Three basis sets (X, Y, Z)

$$\boxed{\begin{array}{rcl} |0_x) & = & |1) \\ |1_x) & = & |\sigma) \end{array}} \quad \boxed{\begin{array}{rcl} |0_y) & = & |\sigma) \\ |1_y) & = & |0) \end{array}} \quad \boxed{\begin{array}{rcl} |0_z) & = & |0) \\ |1_z) & = & |1) \end{array}}$$

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Measurement basis includes $|0)$
Is this result possible?

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Y basis: $|\sigma) = |0_y)$

1_y not possible

Z basis: $|\sigma) = |0_z) + |1_z)$

0_z possible

The dual view

A better way: Think about the dual space V^*.

$$\underset{V \text{ basis}}{\{|a)\}} \leftrightarrow \underset{V^* \text{ basis}}{\{(a|\}}$$

NB: Correspondence $|a) \leftrightarrow (a|$ depends on the *entire* basis.

The dual view

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V basis

V^* basis

$$(a\,|\phi\,) = \phi_a$$

in $|\phi) = \sum_a \phi_a\,|a)$

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NB: Correspondence $|a) \leftrightarrow (a|$ depends on the *entire* basis.

- A measurement is associated with a basis for V^*. (This always corresponds to a basis for V itself.)
- Each measurement result a is represented by a dual vector $(\,a\,|$ -- the "effect functional".
- Possibility rule: a is possible iff $(a\,|\phi\,) \neq 0$

  This depends only on the state and the effect functional!
- Vectors $|\phi)$ and $c\,|\phi)$ are equivalent ("same state").

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V basis V^* basis

$$(a\,|\phi\,) = \phi_a \quad \text{in } |\phi) = \sum_a \phi_a\,|a)$$

This all looks similar to AQT.

- In AQT, the result α is possible provided $\langle a\,|\psi\,\rangle \neq 0$.
- **However**, the Hilbert space inner product in AQT fixes a natural correspondence $|a\rangle \leftrightarrow \langle a|$.

NB: C ... he *entire* basis.

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Entanglement

A mixed state is a subspace of V.

Two mobits in Z_2-MQT:

$V \otimes V$ contains 16 vectors (15 states), including

- 9 product states $|\alpha, \beta) = |\alpha) \otimes |\beta)$
- 6 entangled states -- e.g., $|R) = |0,0) + |1,1)$

For larger $|F|$ and dim V, entangled states greatly outnumber product states. Most states are entangled.

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Advertisement: For details on this and other results, see Mike Westmoreland's poster tomorrow!

Mixed states, etc.

Mixture = collection of possible states: $M = \{|a\rangle, |b\rangle, \ldots\}$

Mixtures M and M' are equivalent iff they span the same subspace. A mixed state $\langle M \rangle$ is a subspace of V.

Given an entangled state $\left|\Phi^{(12)}\right\rangle = \sum_a \left|a^{(1)}\right\rangle \otimes \left|\phi_a^{(2)}\right\rangle$

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A generalized effect is a subspace $E \subseteq V^*$

E is possible for M iff there exist $(e| \in E$ and $|m) \in M$ such that $(e|m) \neq 0$.

A generalized measurement is a set of effects $\{E_k\}$ that spans V^*

$$\left\langle \bigcup_k E_k \right\rangle = V^*$$

Bugs and features

What MQT **does not** have

- Probabilities, expectations
- (F finite) Continuous sets of states and observables, or continuous time evolution
- Inner product, outer product, orthogonality
- Convexity
- Hermitian conjugation (†)
- Density operators
- Effect operators
- CP maps
- Unextendible product bases

What MQT **does** have

- "Classical" versus "quantum" theories
- Superposition, interference
- Linear dynamics
- Complementary measurements
- Entanglement
- No local hidden variables
- KS theorem, "free will" theorem
- Superdense coding, teleportation, "steering" of mixtures, etc.
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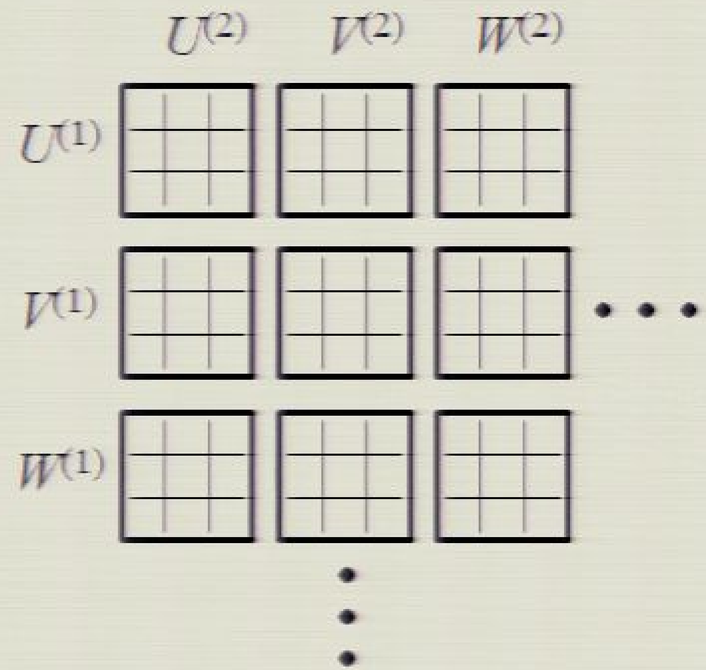
Nothing on the right logically depends on anything on the left.

General models for two systems

General probabilistic models

- Subsystems (1) and (2)
- Measurements U, V, etc. on each subsystem
- Each joint measurement yields a probability distribution over joint results

$$p\left(u, v | U^{(1)}, V^{(2)}\right)$$



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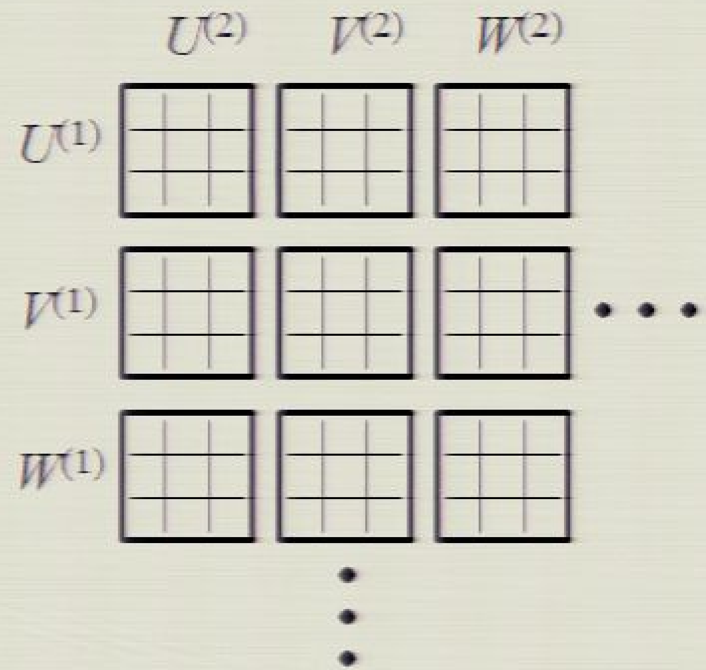

No-signaling principle: The probability of a result on one subsystem does not depend on the choice of measurement on the other subsystem.

$$\sum_{v} p\left(u, v | U^{(1)}, V^{(2)}\right) = p\left(u | U^{(1)}\right)$$

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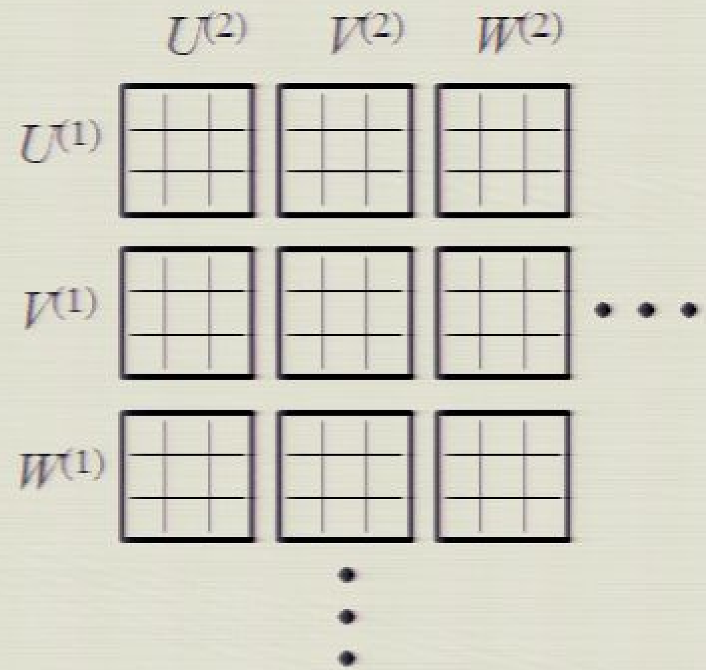

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Satisfied by AQT

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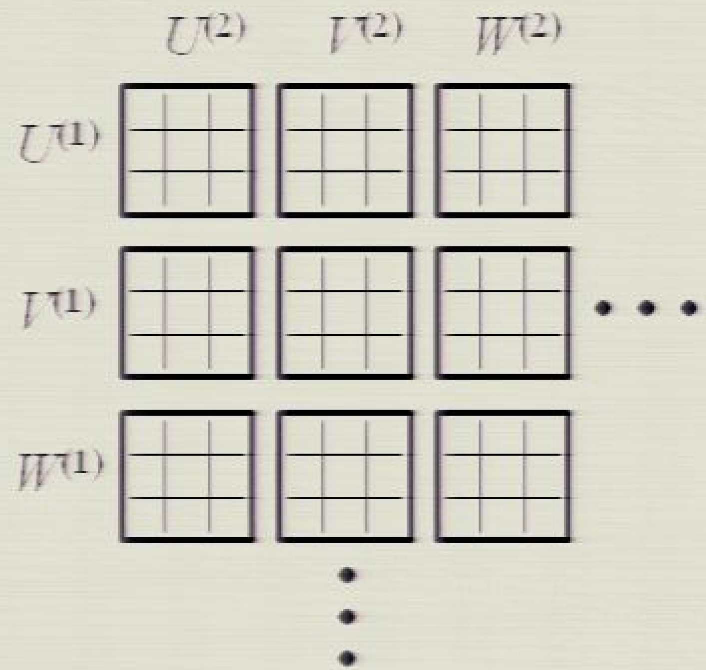

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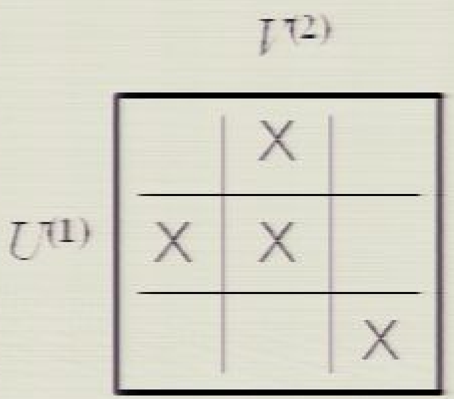

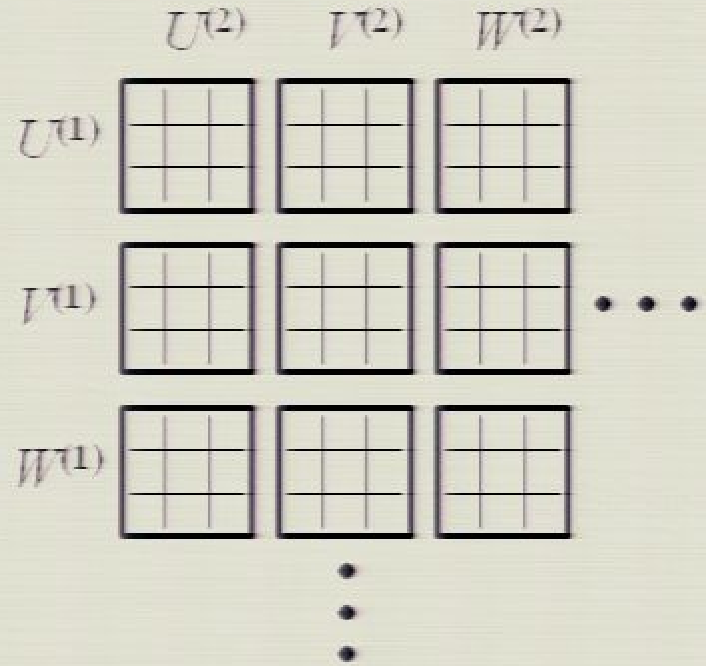

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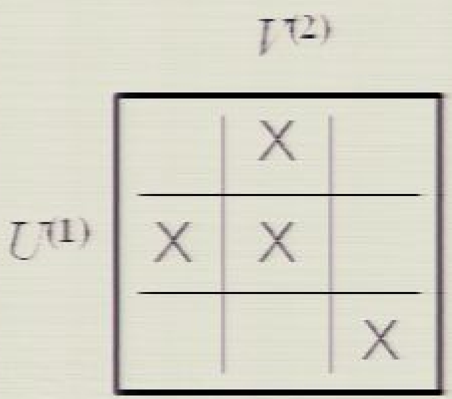

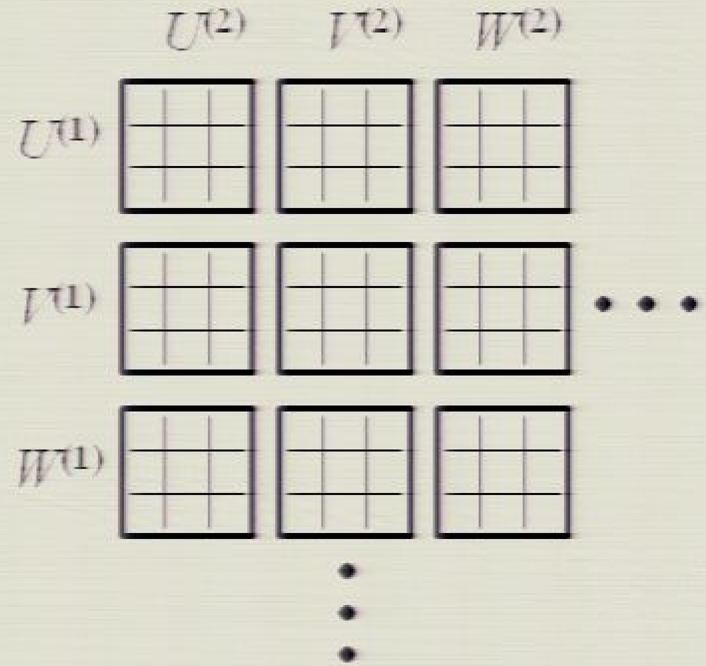

No-signaling principle: The overall possibility of a result on one subsystem does not depend on the choice of measurement on the other subsystem.

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Probabilistic resolution

Can we always replace the X's with probabilities that satisfy the no-signaling principle?

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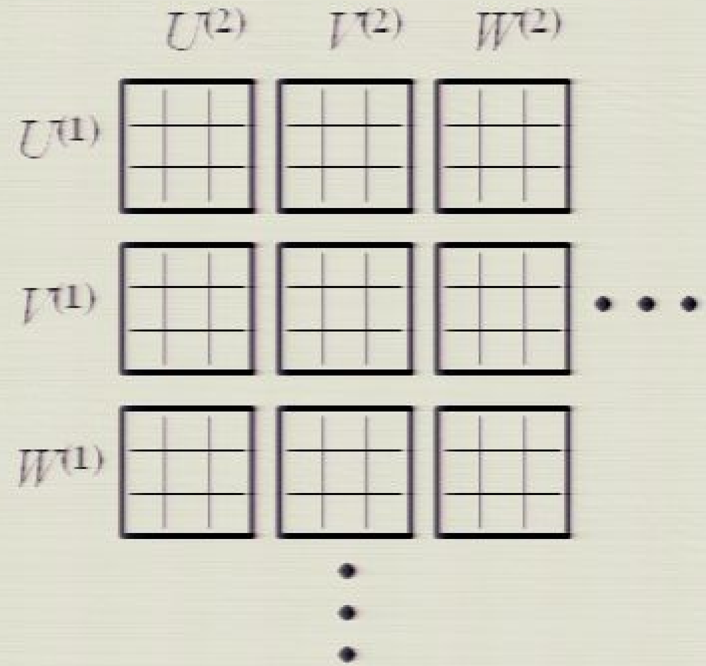

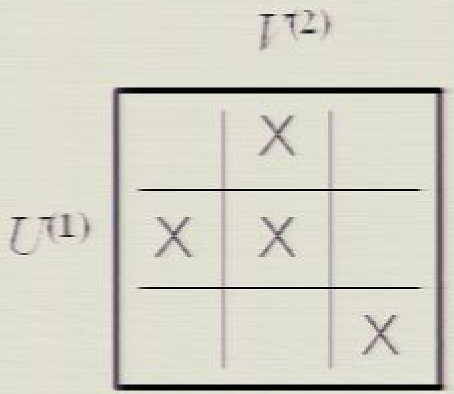

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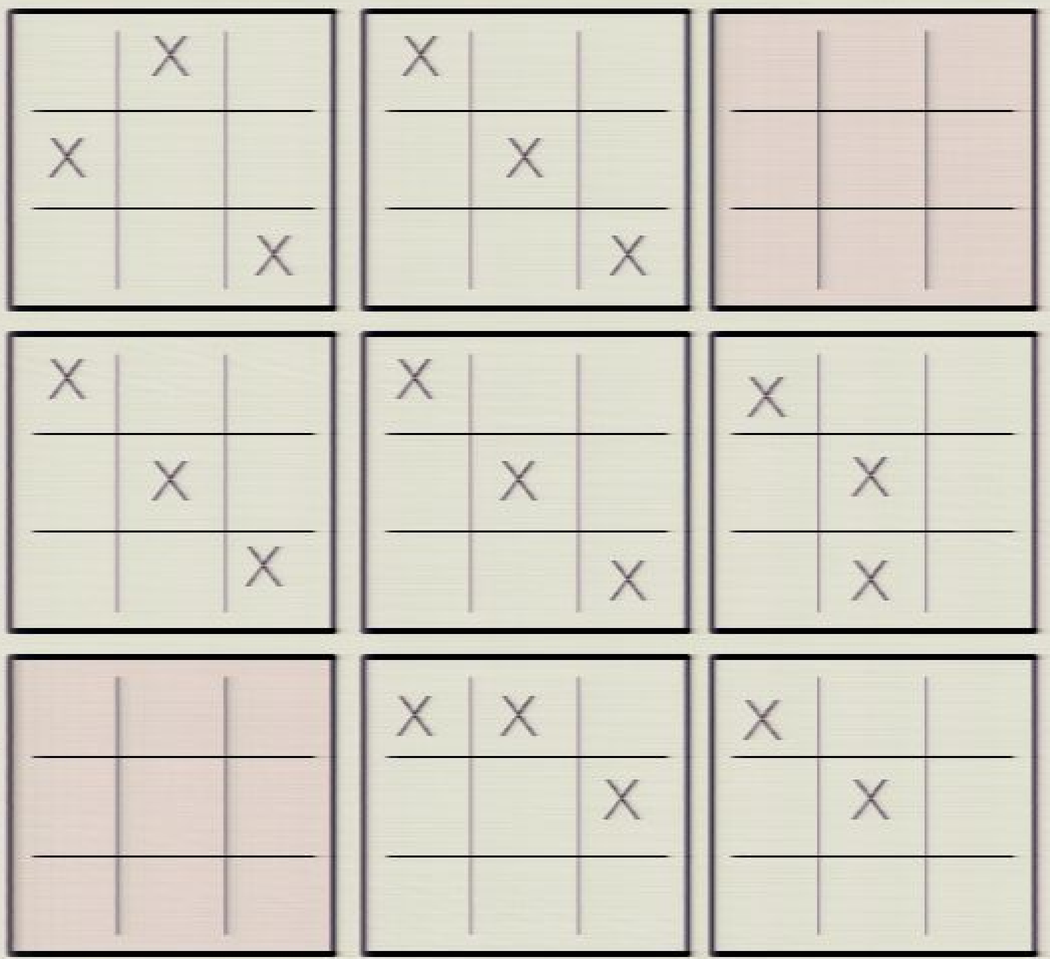
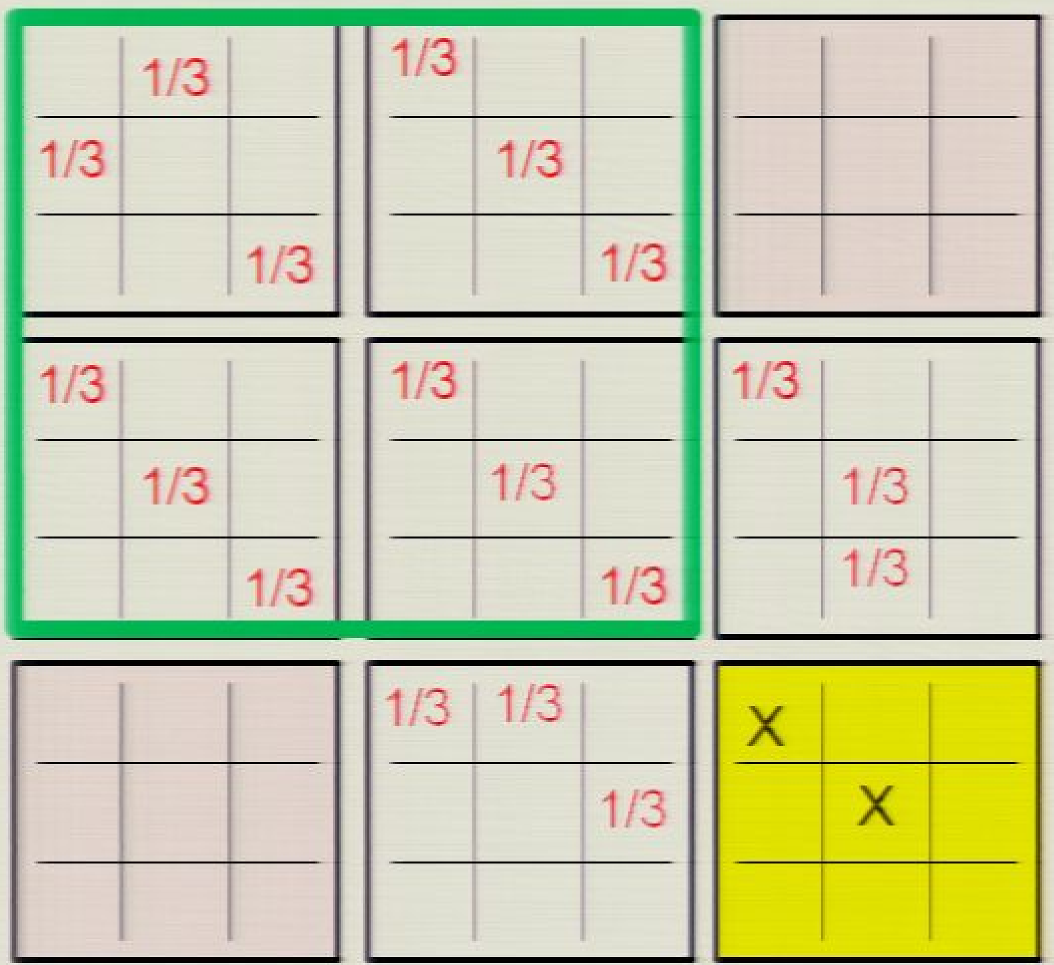


From possibility to probability?

Can we always replace the X's with probabilities that satisfy the no-signaling principle?

All of these probabilities are forced to 1/3.



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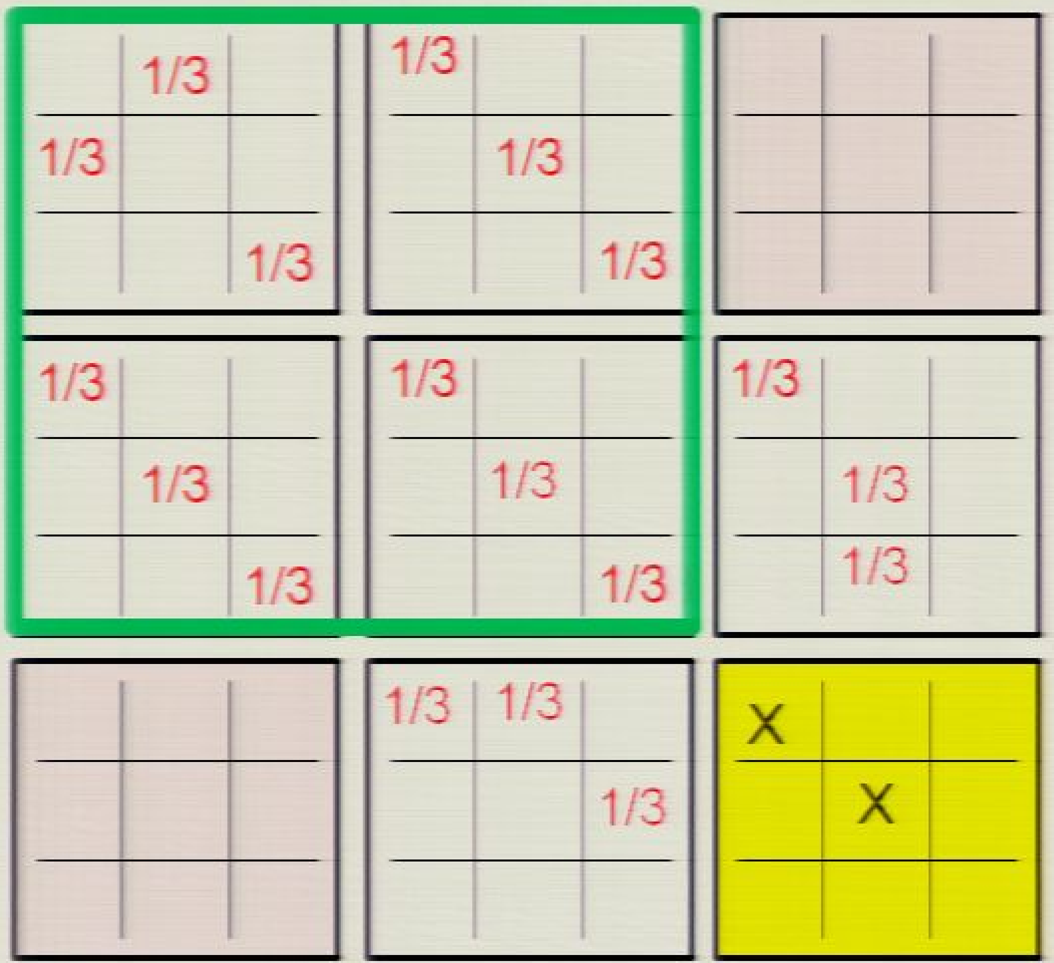
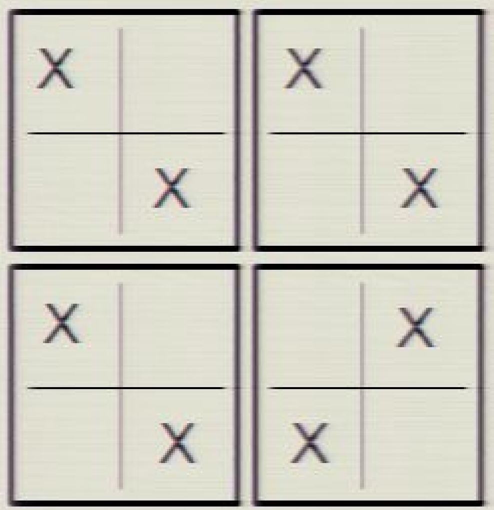


We cannot match this pattern with any probability assignment.

PR boxes



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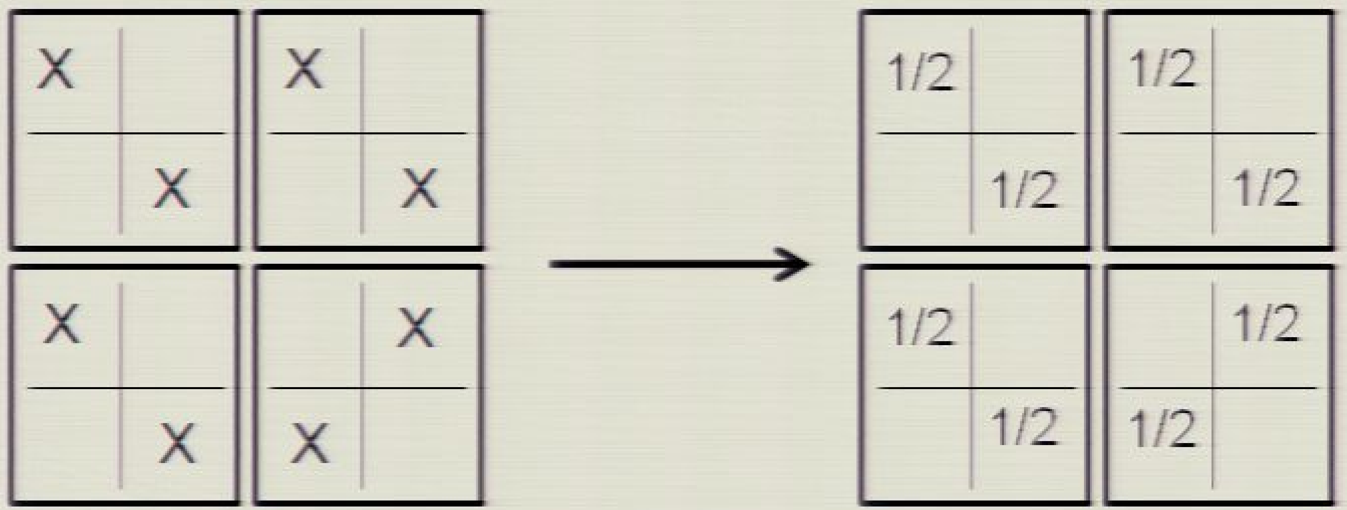
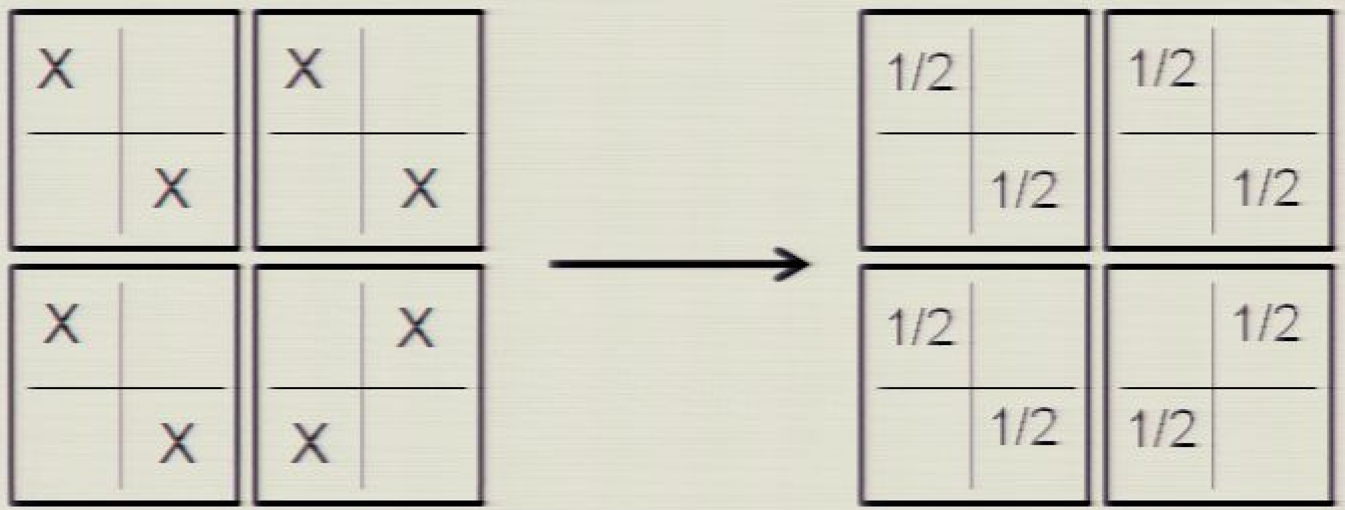


Popescu-Rohrlich
"nonlocal box"

PR boxes



Popescu-Rohrlich
"nonlocal box"

- This is not an allowed probability pattern for a pair of systems in AQT (Tsirelson bound).
- Is this an allowed possibility pattern for a pair of systems in MQT?

PR boxes in MQT?

WLOG, we need consider only pure states and "non-overlapping" measurements.

$$E_0^{(1)} \cap E_1^{(1)} = \langle 0 \rangle, \text{ etc.}$$

	$E_0^{(2)}$	$E_1^{(2)}$	$F_0^{(2)}$	$F_1^{(2)}$
$E_0^{(1)}$	X		X	
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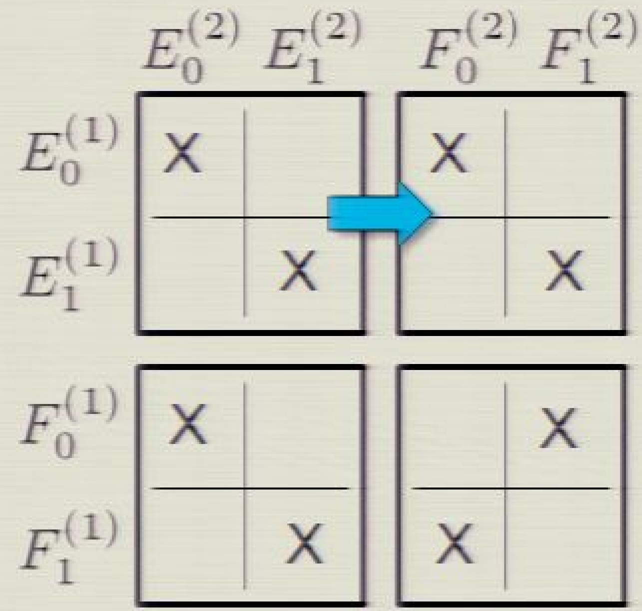

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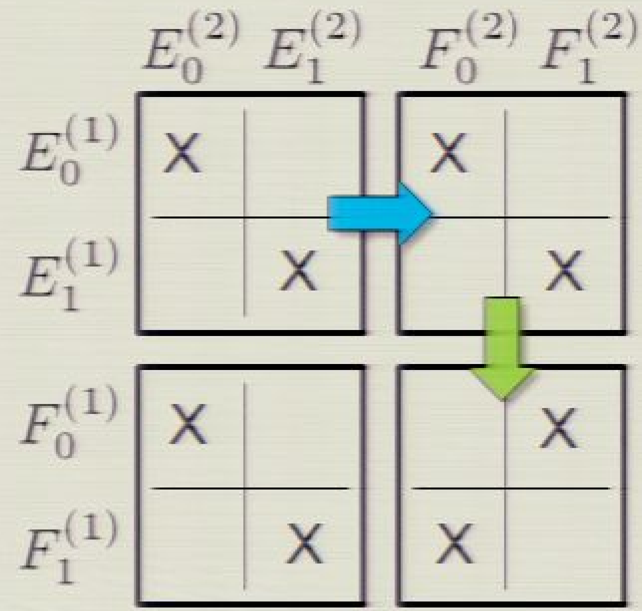

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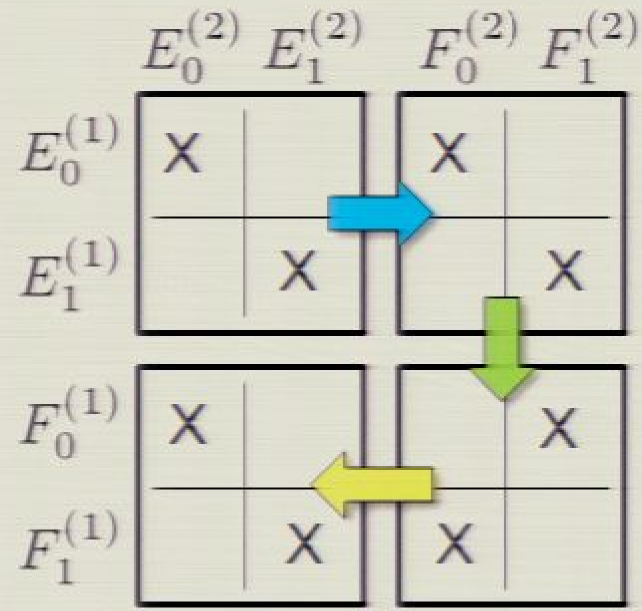

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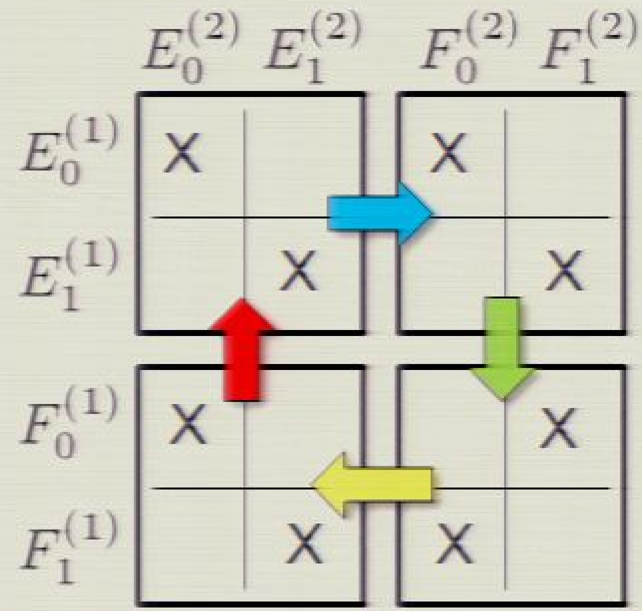

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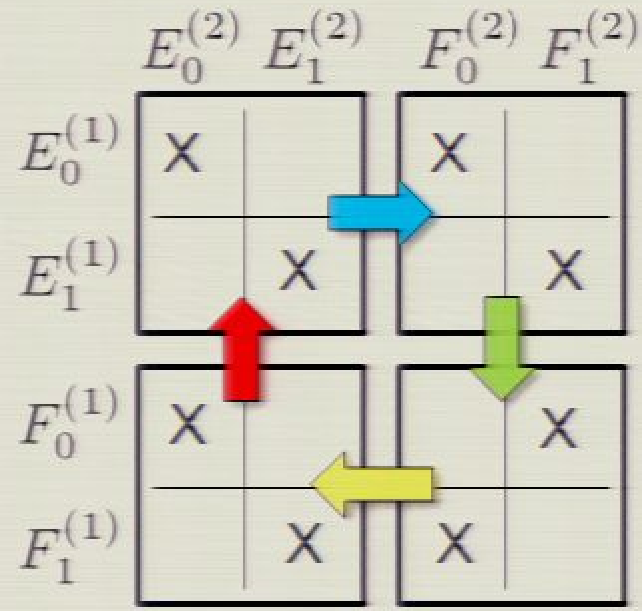

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Contradiction!

No PR boxes in MQT!

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	X		X	X		X
$Y^{(1)}$	X	X		X	X	
		X	X		X	X
$Z^{(1)}$	X		X	X		X
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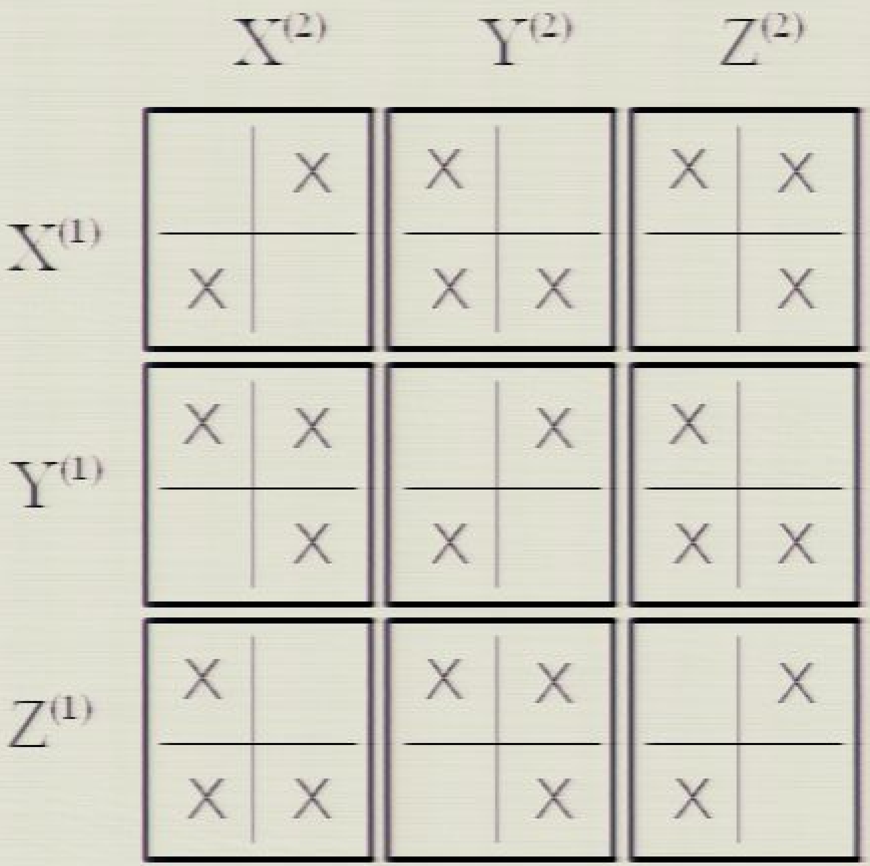

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	$X^{(2)}$		$Y^{(2)}$		$Z^{(3)}$	
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	1/2		0	1/2		1/2
$Y^{(1)}$	1/2	0		1/2	1/2	
		1/2	1/2		0	1/2
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We cannot simulate this MOT system using AOT.

Types of general modal theory (two systems)

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NSP = theories that satisfy the no-signaling principle

SPR = a "strong probabilistic resolution" exists ($p>0$ for each possibility)

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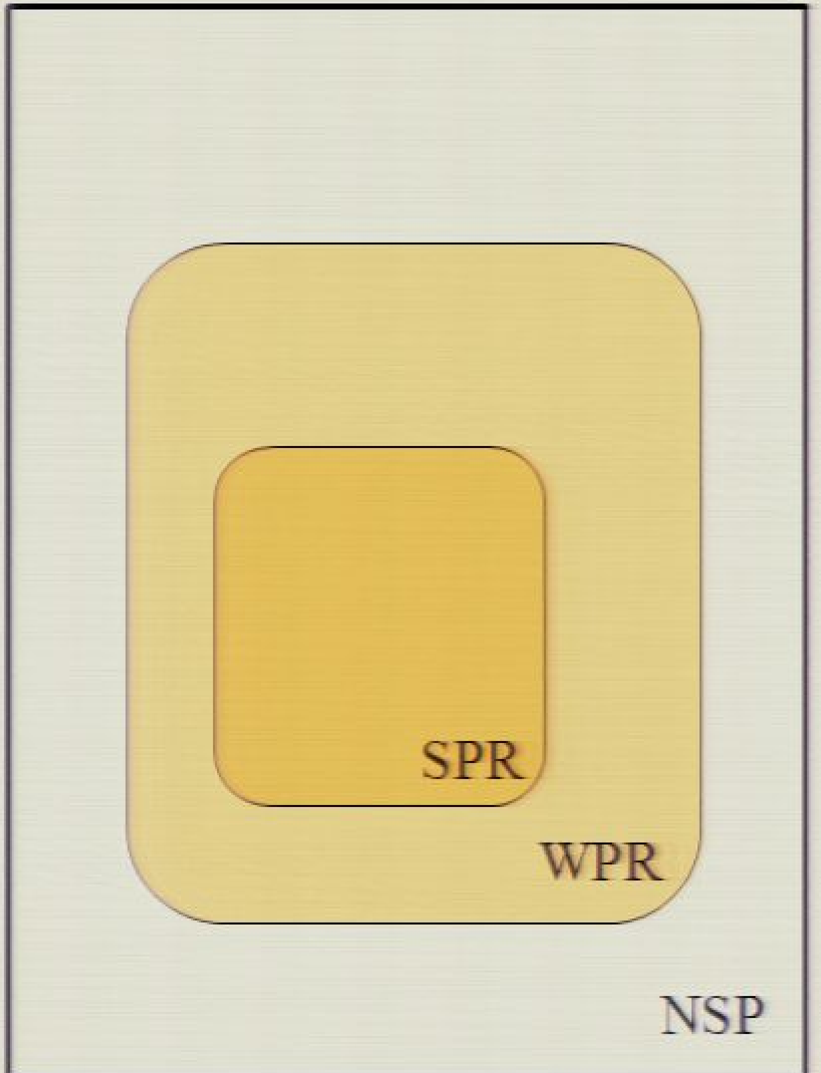

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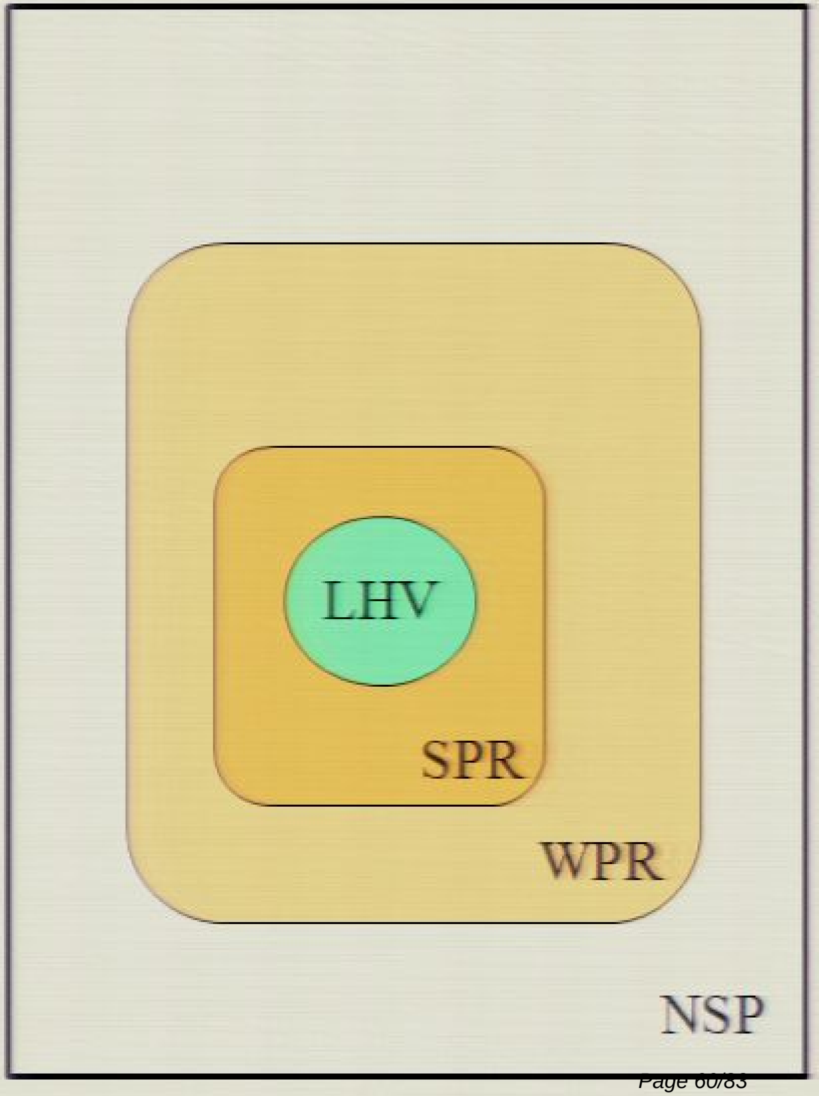

Things we know about MQT

Everything with an LHV model is also in MQT.

Some things in SPR (e.g., PR boxes) are not in MQT.
Some things in MQT (e.g., $|S\rangle$ state) are not in SPR.

Key question: Is everything in MQT also in WPR?

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Two simplifications of the problem:

- We only need to consider pure states and effects. (More X's in the tables can only make the WPR problem easier!)
- We only need to consider basic measurements and entangled states with Schmidt number = dim V.

Possibility table for an entangled state

	X	
X	X	
		X

X		X
X		
	X	X

	X	X
X		X
X	X	

X		X
	X	
X		X

X	X	
	X	X
	X	X

X		
	X	X
	X	

	X	
X	X	X
	X	X

X	X	
X		X
	X	

X	X	
	X	
	X	X

$d = \dim V$
N distinct measurements
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	X	
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	X	
X		X

X	X	
	X	X
	X	X

X		
	X	X
	X	

	X	
X	X	X
	X	X

X	X	
X		X
	X	

X	X	
	X	
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N distinct measurements on each system

Our problem: Devise a probability assignment for this table, respecting the NSP.

We may assign $p = 0$ to some of the X's if need be.

What we know: This table corresponds to basic measurements made on a "maximally entangled" MQT state.

Hall's marriage theorem

Two sets of d elements: $W = \{\text{Alice, Beth, Connie, ...}\}$

$M = \{\text{Adam, Bob, Carl, ...}\}$

A "compatibility" relation between W and M (subset of $W \times M$)

"Marriage condition": For any n, any subset of n elements of W is compatible with at least n elements of M.

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Each $d \times d$ sub-table satisfies the marriage condition. Thus, it includes a "permutation" sub-table on d elements.

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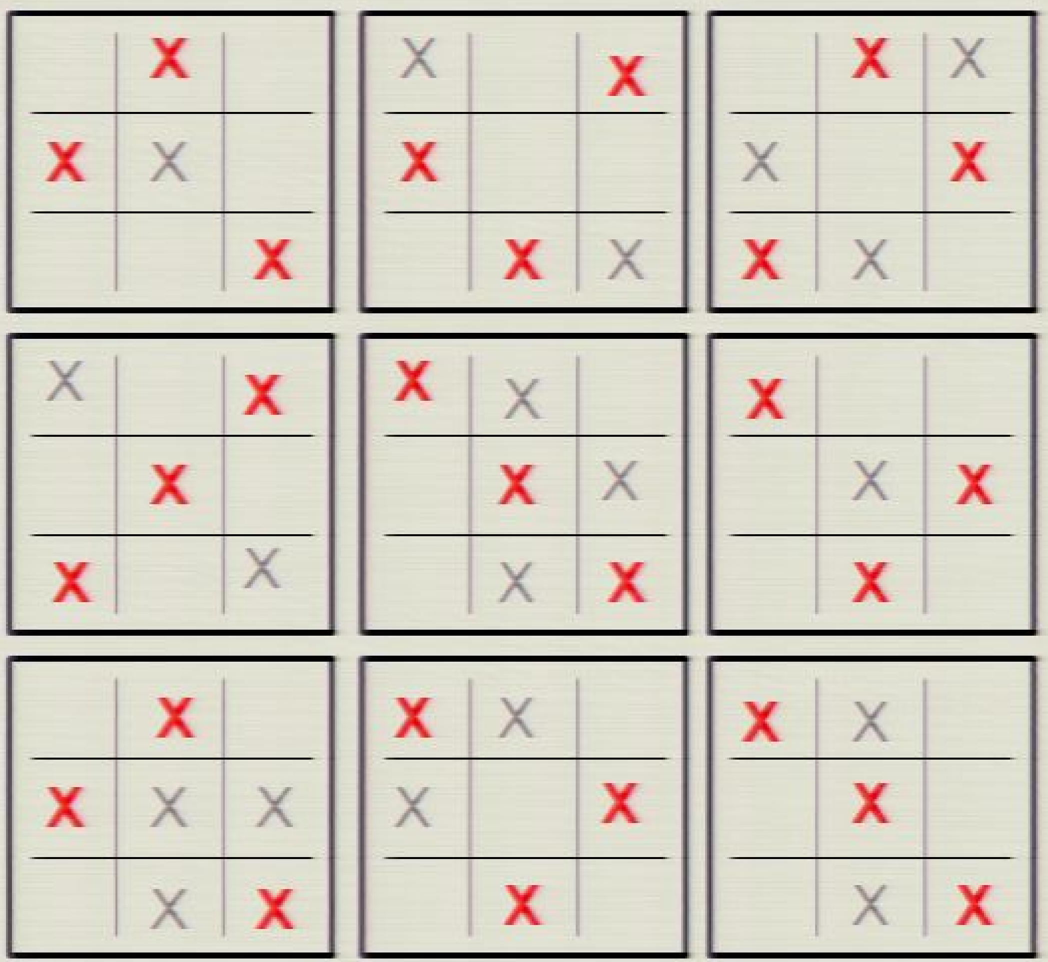

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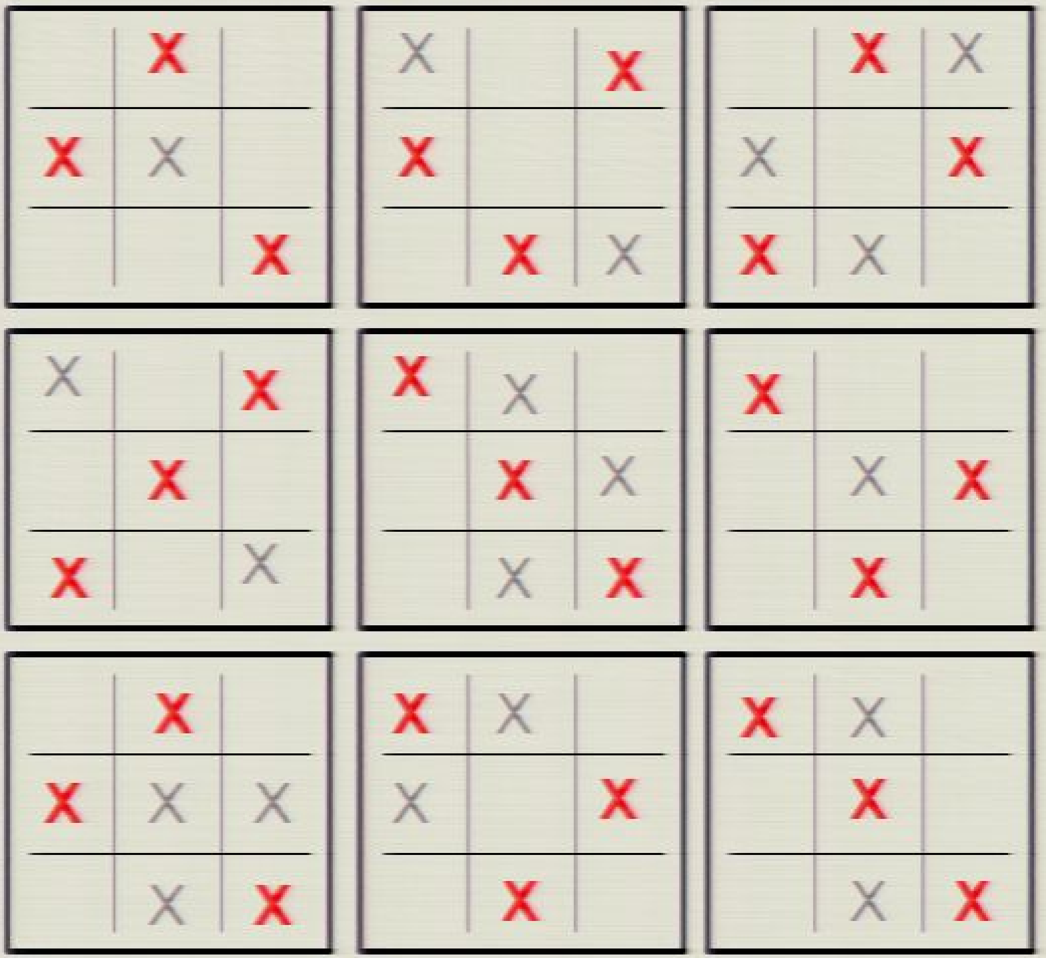

Do the "marriage trick" on each sub-table in the table.

Assign $p = 1/d$ to each marriage, $p = 0$ to everything else.

In each sub-table, each row and each column sums to $p = 1/d$. Thus, this assignment automatically satisfies the NSP.

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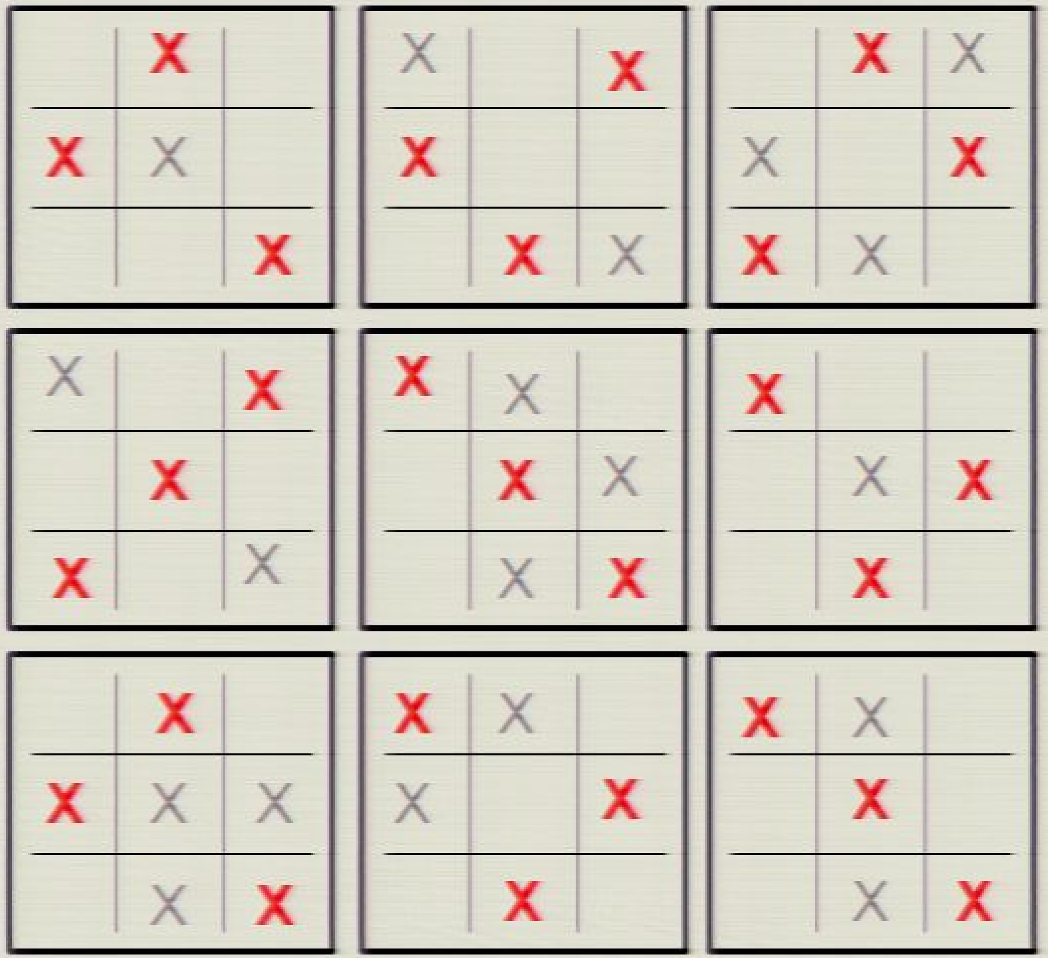

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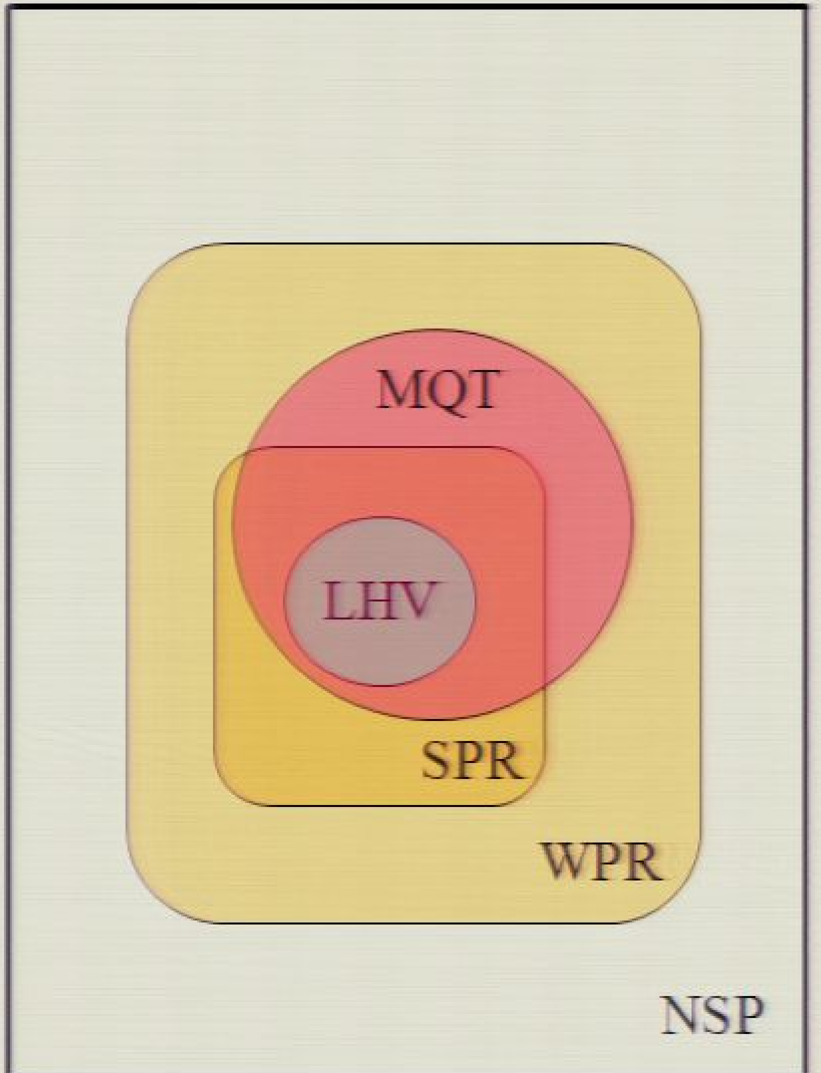

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arXiv: 1010.2929
arXiv: 1010.5452
PIRSA: 10090069 (BWS)
PIRSA: 10100050 (MDW)

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