

Title: First results from the T2K experiment

Date: Apr 19, 2011 12:30 PM

URL: <http://pirsa.org/11040088>

Abstract: Neutrino oscillations has been observed and confirmed at two mass splittings (Δm^2), which is consistent with three generations of neutrinos and an unitary mixing matrix. Despite the rapid progress in understanding neutrino oscillations in the last decade, two large questions remain about neutrino oscillation parameters at $\Delta m^2 \sim 0.001 \text{ eV}^2$. Is θ_{23} exactly 45 degrees, indicating an additional symmetry in neutrino mixing? Is θ_{13} non-zero, which would mean there could be CP violation in the neutrino sector. If θ_{13} is large enough, then such CP violation could be studied with future high intensity experiments such as the proposed Long Baseline Neutrino experiment in the US (LBNE). The Tokai-To-Kamioka (T2K) long baseline neutrino experiment is designed to precisely measure ν_{μ} disappearance (Δm^2_{23} , θ_{23}) and search for ν_e appearance (θ_{13}). A beam of muon neutrinos is generated at the J-PARC facility in Tokai-mura, Japan, and is sampled by two near detectors, ND280 and INGRID, before reaching the Super-Kamiokande detector, 295km away. In this talk, a first look at ν_{μ} disappearance and ν_e appearance will be shown from T2K, from the inaugural 6 month run ending in June 2010 (3.23×10^{19} protons on target, at $15.5 \text{ kW} \times 10^7 \text{s}$).

Earthquake in Japan

On March 11th, 2011, Japan experienced a severe earthquake followed by a tsunami

No reported injuries to members of the T2K collaboration or JPARC employees

All foreign collaborators have returned home safely

The tsunami did not reach JPARC

Inspection of the lab is ongoing

Priority is to restore water, power, and gas systems



Introduction

A reminder about neutrinos

In the Standard Model, there are three neutrinos: ν_e , ν_μ , ν_τ
paired with an associated charged lepton partner: e, μ , τ

Neutrinos interact via the
weak force (W, Z bosons)

To detect neutrinos:
Need "nothing" coming in

Detect the outgoing lepton to
determine neutrino flavor (CC only)

Nucleus can be excited or
additional particles emitted
(NC or CC)

Charged-Current (CC) Neutral-Current (NC)
Interactions Interactions
Neutrinos



Anti-Neutrinos



Quarks



Neutrino mixing

The flavor state of the neutrino, ν_α , is related to the mass states, ν_i , by a non unity mixing matrix, $U_{\alpha i}$

$$|\nu_i\rangle = \sum U_{\alpha i} |\nu_\alpha\rangle$$

Since there are three observed flavors of neutrinos (ν_e , ν_μ , ν_τ), U contains three mixing angles ($\theta_{12}, \theta_{23}, \theta_{13}$) and a CP violating phase δ .

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$c_{ij} = \cos \theta_{ij}, s_{ij} = \sin \theta_{ij}$$

$$U_{\alpha i} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

$$\begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix}$$

$$\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

"Atmospheric": $\theta_{23} \sim 37^\circ - 45^\circ$

"Reactor": $\theta_{13} < 11^\circ$

"Solar": $\theta_{12} \sim 34^\circ$

Quark mixing: unitary matrix, small angles: $\theta_{CKM}^{CKM}_{12} \sim 13.0^\circ$, $\theta_{CKM}^{CKM}_{23} \sim 2.3^\circ$, $\theta_{CKM}^{CKM}_{13} \sim 0.2^\circ$

▪ Is θ_{23} exactly 45 degrees, or not?

▪ Is θ_{13} zero, or just small?

▪ Is there CP violation in the neutrino sector? Is it large?

Neutrino oscillation

As the neutrinos propagate, the mass states interfere:

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right) + 2 \sum_{i>j} \operatorname{Im} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin \left(\frac{\Delta m_{ij}^2 L}{2E} \right)$$

Depends on:

L (km): Distance the neutrino has travelled

E (GeV): Energy of the neutrino

Δm^2 (eV 2): Difference of the square of the mass eigenvalues

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

"Atmospheric": $\Delta m_{23}^2 = 2.4 \times 10^{-3}$ eV 2

"Solar": $\Delta m_{12}^2 = 7.59 \times 10^{-5}$ eV 2

Probability for ν_μ oscillating into ν_x :
(ν_μ disappearance)

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Probability for ν_μ oscillating into ν_e
(ν_e appearance)

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right)$$

Neutrino mixing

The flavor state of the neutrino, ν_α , is related to the mass states, ν_i , by a non unity mixing matrix, $U_{\alpha i}$

$$|\nu_i\rangle = \sum U_{\alpha i} |\nu_\alpha\rangle$$

Since there are three observed flavors of neutrinos (ν_e , ν_μ , ν_τ), U contains three mixing angles ($\theta_{12}, \theta_{23}, \theta_{13}$) and a CP violating phase δ .

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$c_{ij} = \cos \theta_{ij}, s_{ij} = \sin \theta_{ij}$$

$$U_{\alpha i} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

$$\begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix}$$

$$\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

"Atmospheric": $\theta_{23} \sim 37^\circ - 45^\circ$

"Reactor": $\theta_{13} < 11^\circ$

"Solar": $\theta_{12} \sim 34^\circ$

Quark mixing: unitary matrix, small angles: $\theta_{CKM}^{CKM}{}_{12} \sim 13.0^\circ$, $\theta_{CKM}^{CKM}{}_{23} \sim 2.3^\circ$, $\theta_{CKM}^{CKM}{}_{13} \sim 0.2^\circ$

▪ Is θ_{23} exactly 45 degrees, or not?

11040088 ▪ Is θ_{13} zero, or just small?

▪ Is there CP violation in the neutrino sector? Is it large?

Neutrino oscillation

As the neutrinos propagate, the mass states interfere:

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right) + 2 \sum_{i>j} \operatorname{Im} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin \left(\frac{\Delta m_{ij}^2 L}{2E} \right)$$

Depends on:

L (km): Distance the neutrino has travelled

E (GeV): Energy of the neutrino

Δm^2 (eV 2): Difference of the square of the mass eigenvalues

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

"Atmospheric": $\Delta m_{23}^2 = 2.4 \times 10^{-3}$ eV 2
"Solar": $\Delta m_{12}^2 = 7.59 \times 10^{-5}$ eV 2

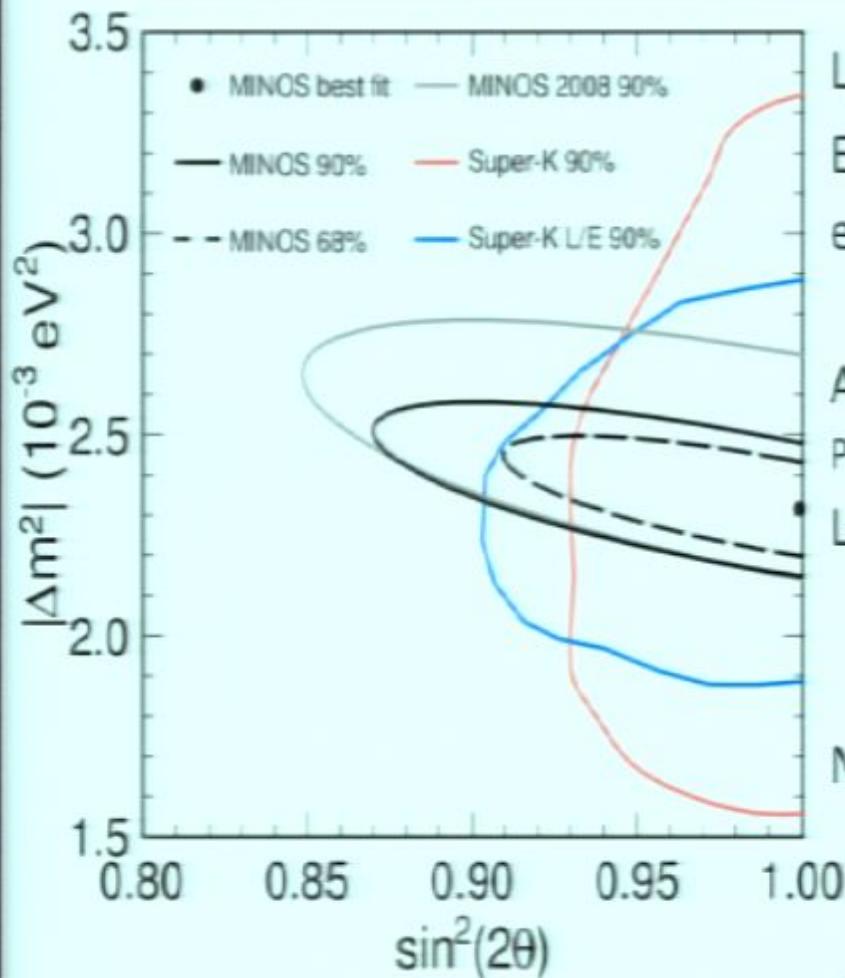
Probability for ν_μ oscillating into ν_x :
(ν_μ disappearance)

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Probability for ν_μ oscillating into ν_e
(ν_e appearance)

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right)$$

Existing measurements of ν_μ disappearance



L, E are fixed based on neutrino source
Extract $|\Delta m^2|$, $\sin^2 2\theta$ based on rate,
energy spectrum after oscillation

Accelerator-produced neutrino beam (MINOS)
Phys. Rev. Lett. 101, 131802 (2008)
 $L=735\text{ km}$, $E(\text{peak}) \sim 3 \text{ GeV}$

$$\Delta m_{23}^2 = 2.43 \pm 0.13 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{23} > 0.90 \text{ (90\% CL)}$$

New result (2011) 1103.0340 [hep-ex]

$$\Delta m_{23}^2 = 2.32^{+0.12}_{-0.08} \times 10^{-3} \text{ eV}^2 \text{ (68\% CL)}$$

Atmospheric neutrinos (Super-Kamiokande)

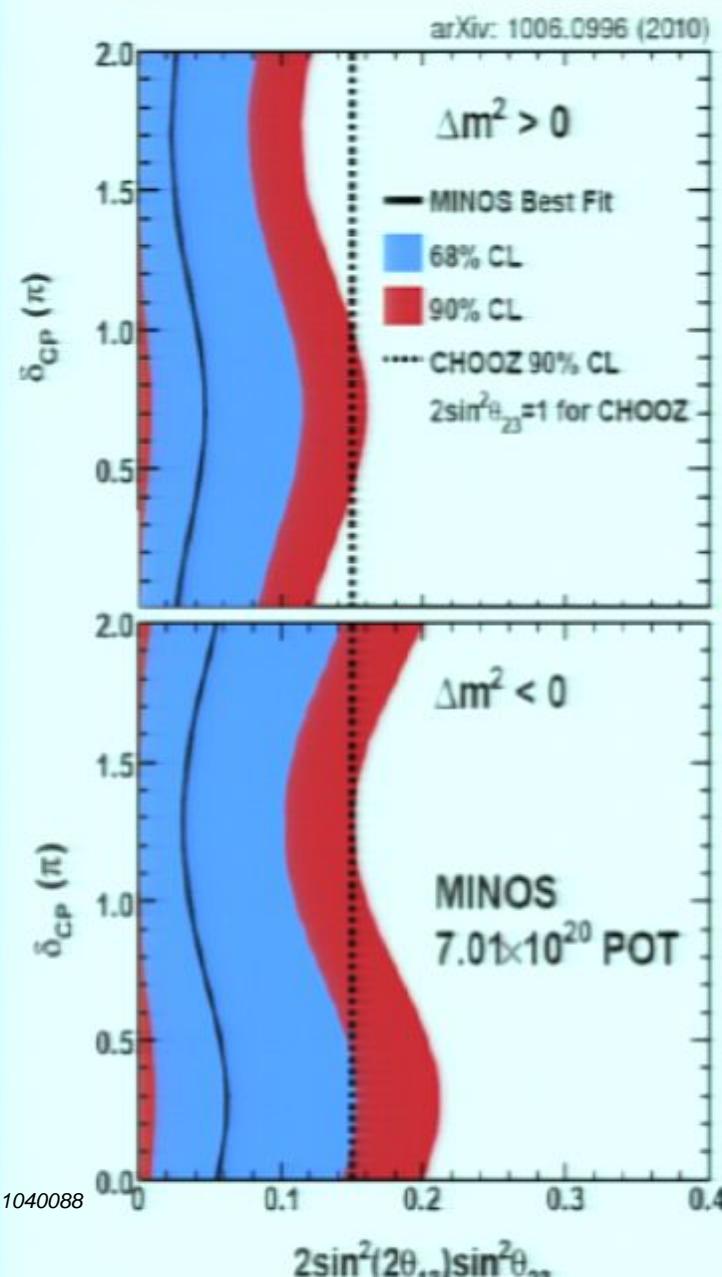
Phys. Rev. D71:112005, 2005; hep-ex/0501064

$L \sim 15\text{ km} - 13,000\text{ km}$, $E \sim 100 \text{ MeV} - 10 \text{ TeV}$

$$1.5 < \Delta m_{23}^2 < 3.4 \times 10^{-3} \text{ eV}^2$$

$$B(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Existing measurements of ν_e appearance



Accelerator-produced neutrino beam (MINOS)

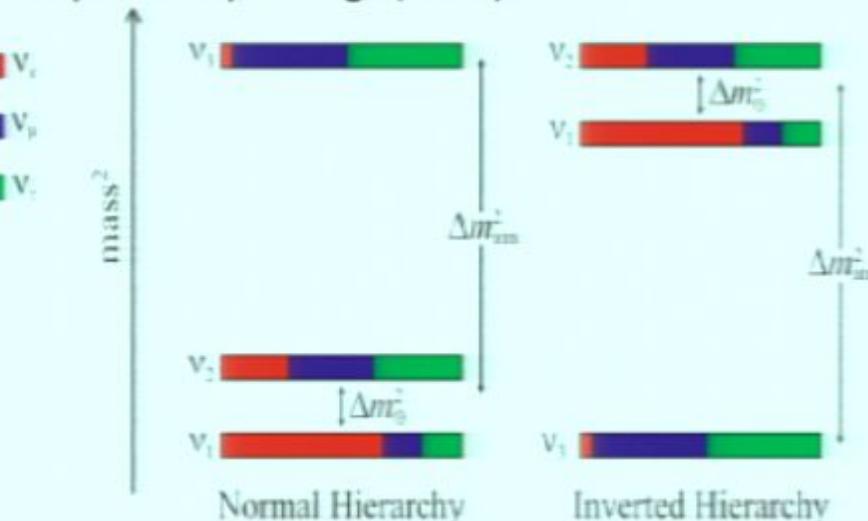
Phys.Rev.D82:051102,2010; hep-ex/1006.0996

L=735km, E(peak) ~3 GeV

$$2\sin^2\theta_{23}\sin^22\theta_{13} < 0.12 \text{ (90% CL)}$$

ν_e and ν_μ interact differently in matter vs vacuum, altering oscillation probability

Depends upon sign(Δm^2)

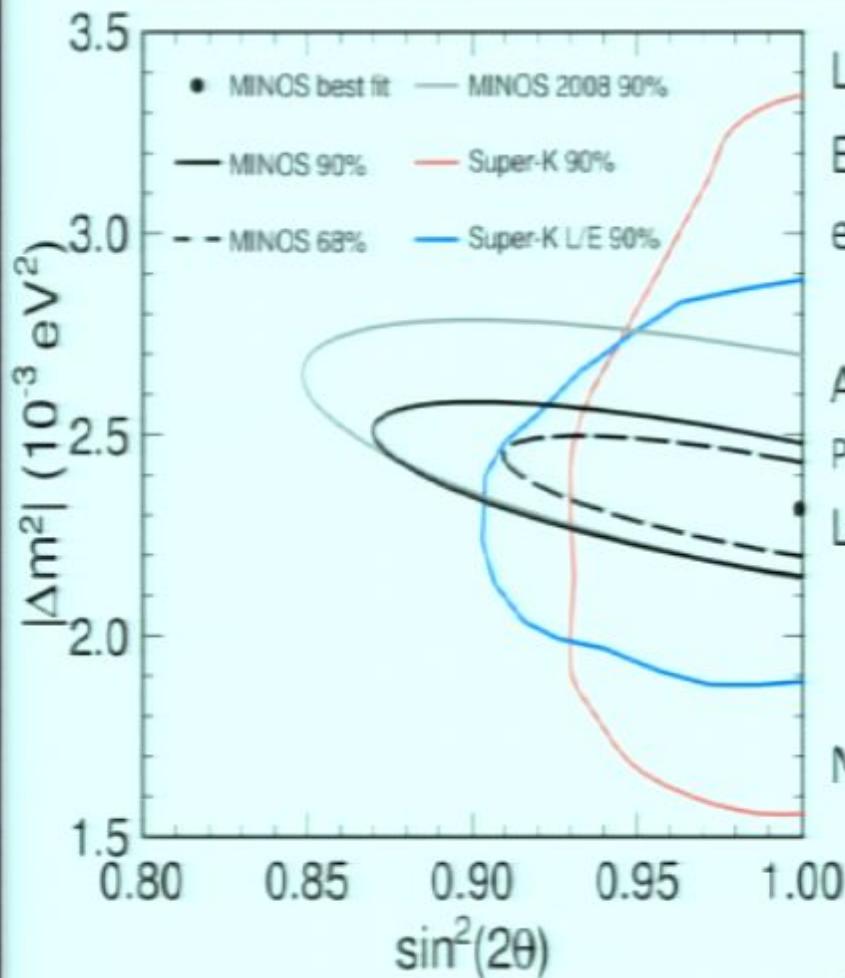


Reactor ν_e disappearance (CHOOZ)

Eur.Phys.J.C27:331-374,2003; hep-ex/0301017

L~1 km, E~3 MeV

Existing measurements of ν_μ disappearance



L, E are fixed based on neutrino source
Extract $|\Delta m^2|$, $\sin^2 2\theta$ based on rate,
energy spectrum after oscillation

Accelerator-produced neutrino beam (MINOS)
Phys. Rev. Lett. 101, 131802 (2008)

$L=735\text{ km}$, $E(\text{peak}) \sim 3 \text{ GeV}$

$$\Delta m_{23}^2 = 2.43 \pm 0.13 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{23} > 0.90 \text{ (90\% CL)}$$

New result (2011) 1103.0340 [hep-ex]

$$\Delta m_{23}^2 = 2.32^{+0.12}_{-0.08} \times 10^{-3} \text{ eV}^2 \text{ (68\% CL)}$$

Atmospheric neutrinos (Super-Kamiokande)

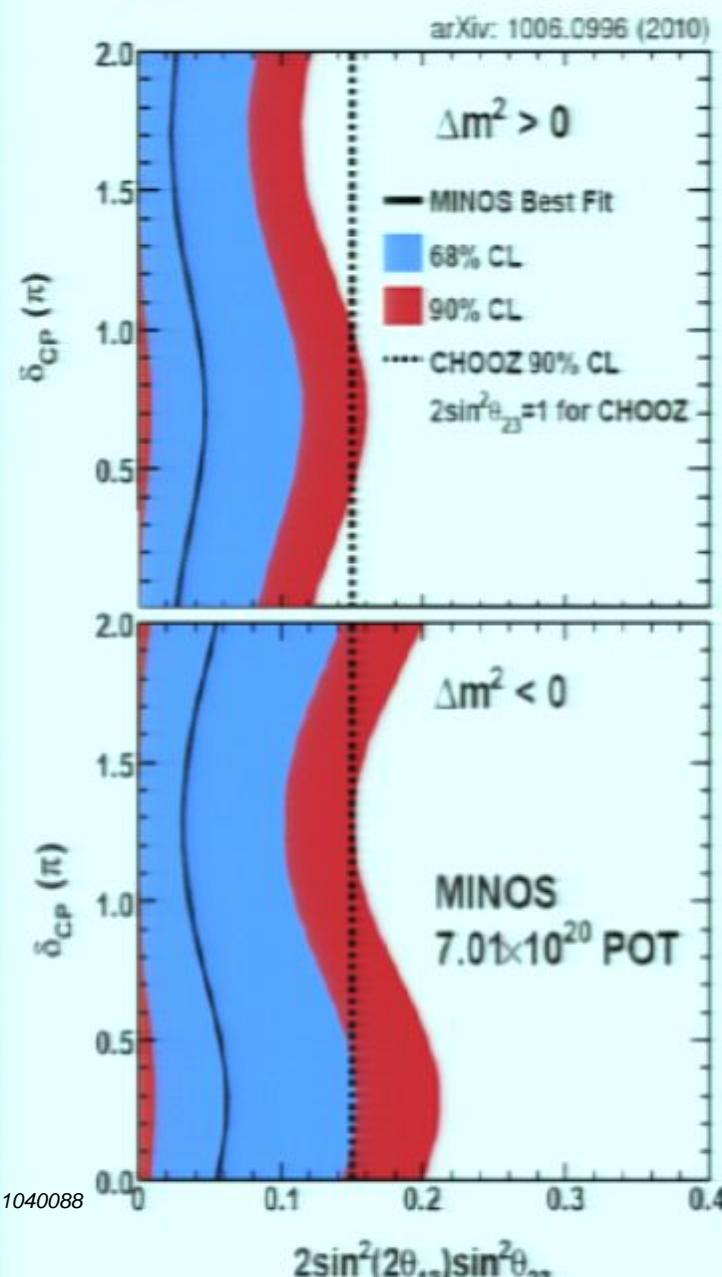
Phys. Rev. D71:112005, 2005; hep-ex/0501064

$L \sim 15\text{ km} - 13,000\text{ km}$, $E \sim 100 \text{ MeV} - 10 \text{ TeV}$

$$1.5 < \Delta m_{23}^2 < 3.4 \times 10^{-3} \text{ eV}^2$$

$$B(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Existing measurements of ν_e appearance



Accelerator-produced neutrino beam (MINOS)

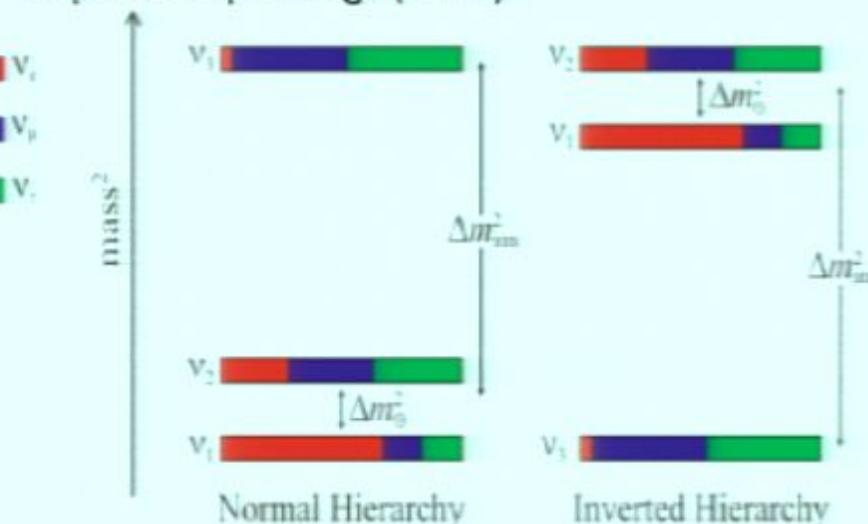
Phys.Rev.D82:051102,2010; hep-ex/1006.0996

$L=735\text{km}$, $E(\text{peak}) \sim 3\text{ GeV}$

$$2\sin^2\theta_{23}\sin^2\theta_{13} < 0.12 \text{ (90% CL)}$$

ν_e and ν_μ interact differently in matter vs vacuum, altering oscillation probability

Depends upon $\text{sign}(\Delta m^2)$



Reactor ν_e disappearance (CHOOZ)

Eur.Phys.J.C27:331-374,2003; hep-ex/0301017

$L \sim 1\text{ km}$, $E \sim 3\text{ MeV}$

Neutrino oscillation

As the neutrinos propagate, the mass states interfere:

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right) + 2 \sum_{i>j} \operatorname{Im} [U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j}] \sin \left(\frac{\Delta m_{ij}^2 L}{2E} \right)$$

Depends on:

L (km): Distance the neutrino has travelled

E (GeV): Energy of the neutrino

Δm^2 (eV²): Difference of the square of the mass eigenvalues

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

"Atmospheric": $\Delta m_{23}^2 = 2.4 \times 10^{-3}$ eV²
"Solar": $\Delta m_{12}^2 = 7.59 \times 10^{-5}$ eV²

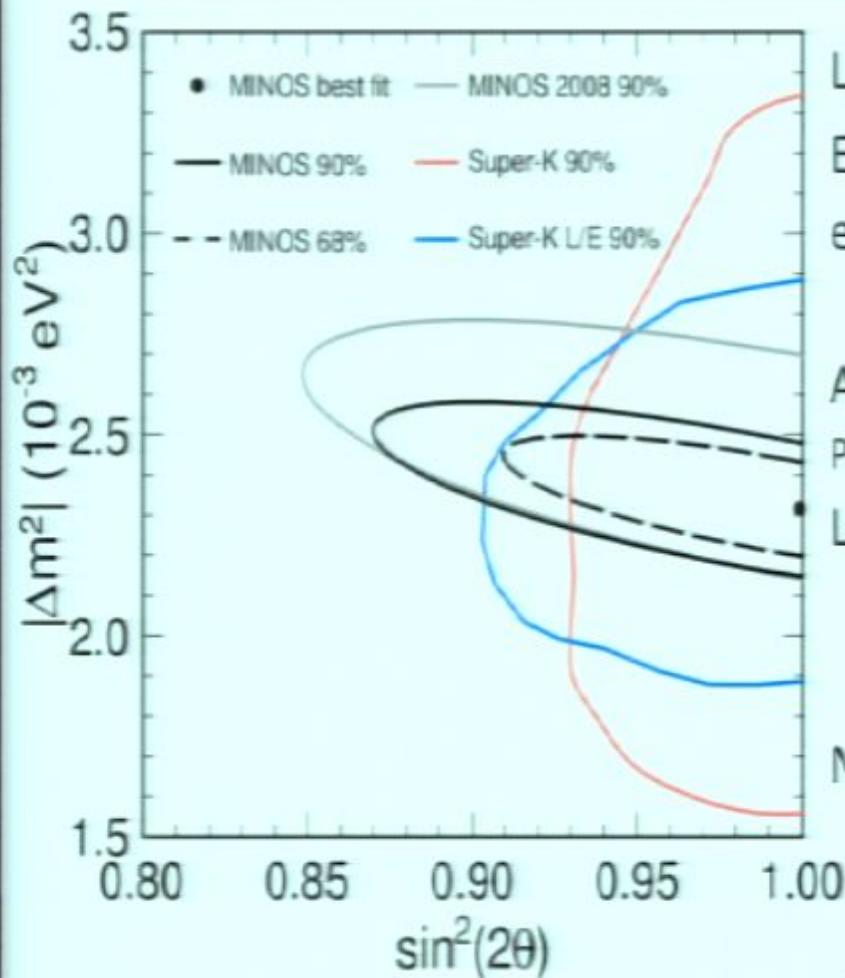
Probability for ν_μ oscillating into ν_x :
(ν_μ disappearance)

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Probability for ν_μ oscillating into ν_e
(ν_e appearance)

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right)$$

Existing measurements of ν_μ disappearance



L, E are fixed based on neutrino source
Extract $|\Delta m^2|, \sin^2 2\theta$ based on rate,
energy spectrum after oscillation

Accelerator-produced neutrino beam (MINOS)
Phys. Rev. Lett. 101, 131802 (2008)

$L=735\text{km}, E(\text{peak}) \sim 3\text{ GeV}$

$$\Delta m_{23}^2 = 2.43 \pm 0.13 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{23} > 0.90 \text{ (90\% CL)}$$

New result (2011) 1103.0340 [hep-ex]

$$\Delta m_{23}^2 = 2.32^{+0.12}_{-0.08} \times 10^{-3} \text{ eV}^2 \text{ (68\% CL)}$$

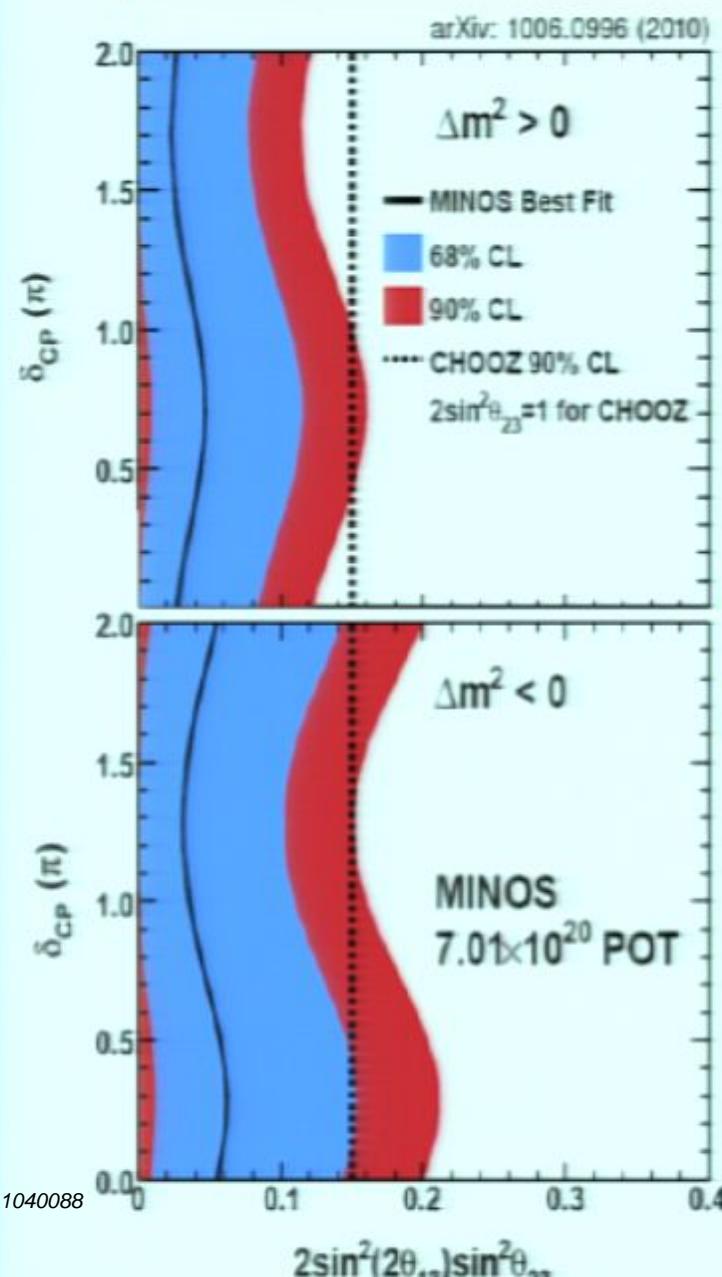
Atmospheric neutrinos (Super-Kamiokande)

Phys. Rev. D71:112005, 2005; hep-ex/0501064

$L \sim 15\text{km}-13,000\text{km}, E \sim 100\text{ MeV}-10\text{ TeV}$

$$1.5 < \Delta m_{23}^2 < 3.4 \times 10^{-3} \text{ eV}^2$$

Existing measurements of ν_e appearance



Accelerator-produced neutrino beam (MINOS)

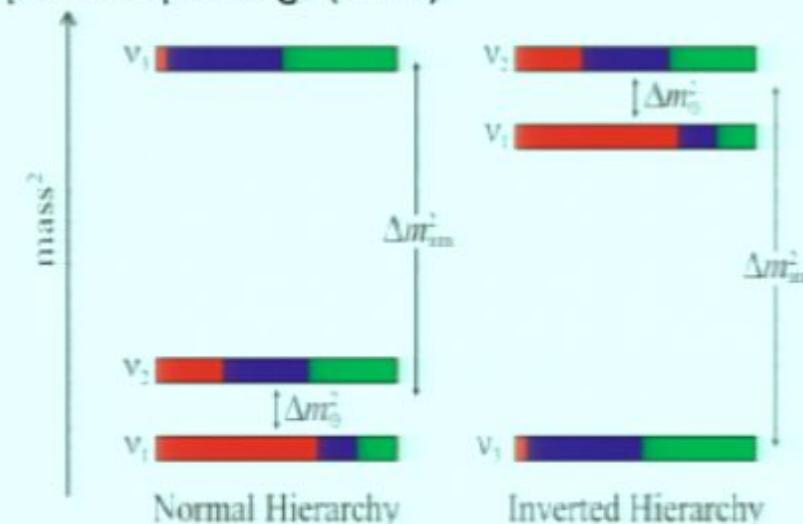
Phys.Rev.D82:051102,2010; hep-ex/1006.0996

$L=735\text{ km}$, $E(\text{peak}) \sim 3 \text{ GeV}$

$$2\sin^2\theta_{23}\sin^22\theta_{13} < 0.12 \text{ (90\% CL)}$$

ν_e and ν_μ interact differently in matter vs vacuum, altering oscillation probability

Depends upon $\text{sign}(\Delta m^2)$



Reactor ν_e disappearance (CHOOZ)

Eur.Phys.J.C27:331-374,2003; hep-ex/0301017

$L \sim 1 \text{ km}$, $E \sim 3 \text{ MeV}$

The T2K experiment

The T2K experiment

T2K is designed to measure oscillations at the atmospheric Δm^2 :

Measure ν_μ disappearance (Δm^2_{23} , θ_{23})

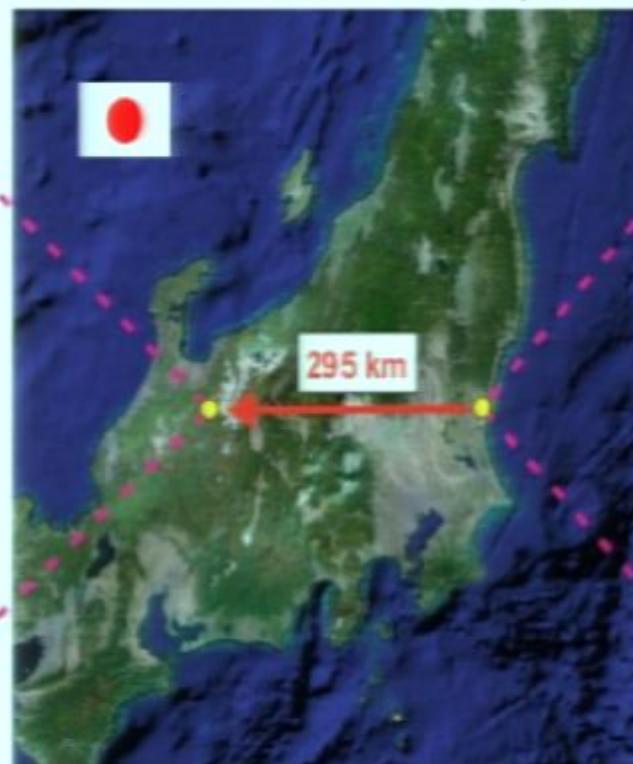
Discover ν_e appearance? (θ_{13})

Produce a beam of ν_μ on one side of Japan and detect it on the other

Super Kamiokande
50,000 tons of water
10,000 phototubes



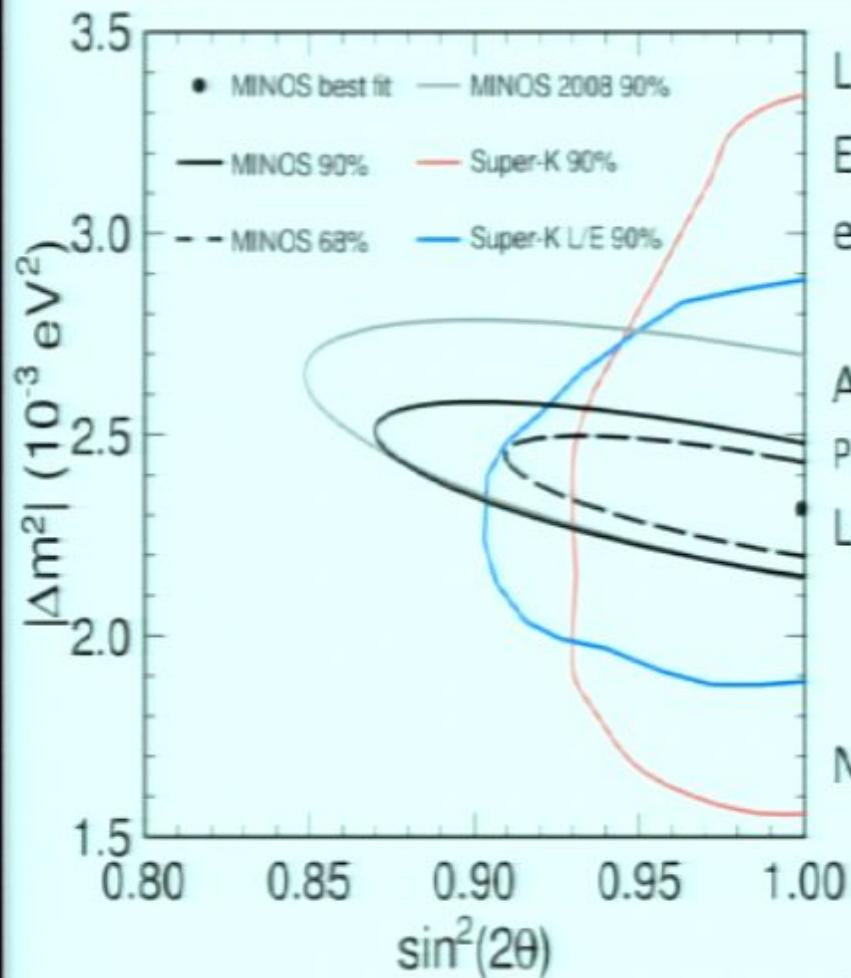
Neutrino beam directed across Japan



Tokai accelerator complex and location of near detector (ND280)



Existing measurements of ν_μ disappearance



L, E are fixed based on neutrino source
Extract $|\Delta m^2|$, $\sin^2 2\theta$ based on rate,
energy spectrum after oscillation

Accelerator-produced neutrino beam (MINOS)
Phys. Rev. Lett. 101, 131802 (2008)

$L=735\text{ km}$, $E(\text{peak}) \sim 3\text{ GeV}$

$$\Delta m_{23}^2 = 2.43 \pm 0.13 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{23} > 0.90 \text{ (90\% CL)}$$

New result (2011) 1103.0340 [hep-ex]

$$\Delta m_{23}^2 = 2.32^{+0.12}_{-0.08} \times 10^{-3} \text{ eV}^2 \text{ (68\% CL)}$$

Atmospheric neutrinos (Super-Kamiokande)

Phys. Rev. D71:112005, 2005; hep-ex/0501064

$L \sim 15\text{ km} - 13,000\text{ km}$, $E \sim 100\text{ MeV} - 10\text{ TeV}$

$$1.5 < \Delta m_{23}^2 < 3.4 \times 10^{-3} \text{ eV}^2$$

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

The T2K experiment

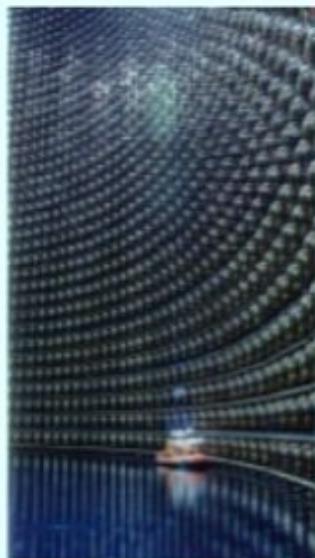
T2K is designed to measure oscillations at the atmospheric Δm^2 :

Measure ν_μ disappearance (Δm^2_{23} , θ_{23})

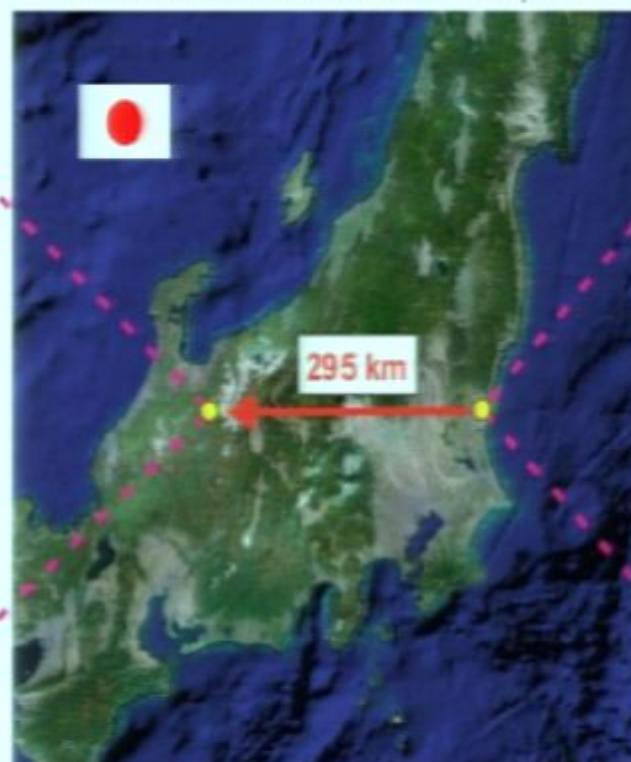
Discover ν_e appearance? (θ_{13})

Produce a beam of ν_μ on one side of Japan and detect it on the other

Super Kamiokande
50,000 tons of water
10,000 phototubes



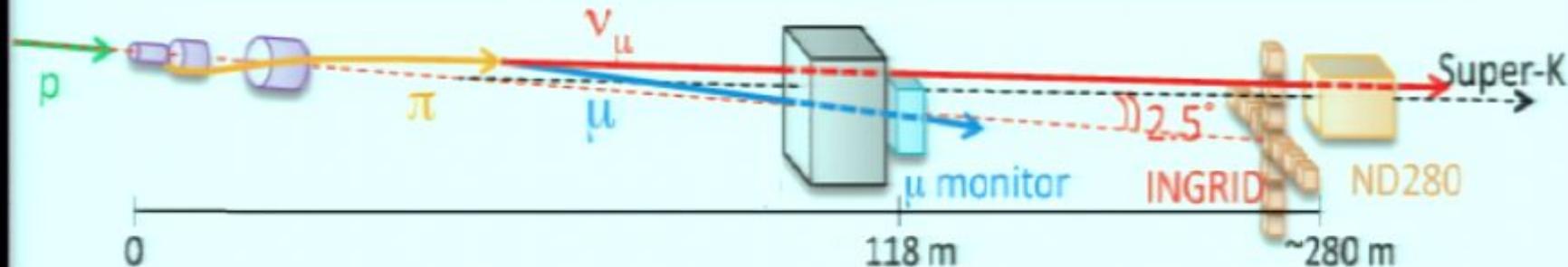
Neutrino beam directed across Japan



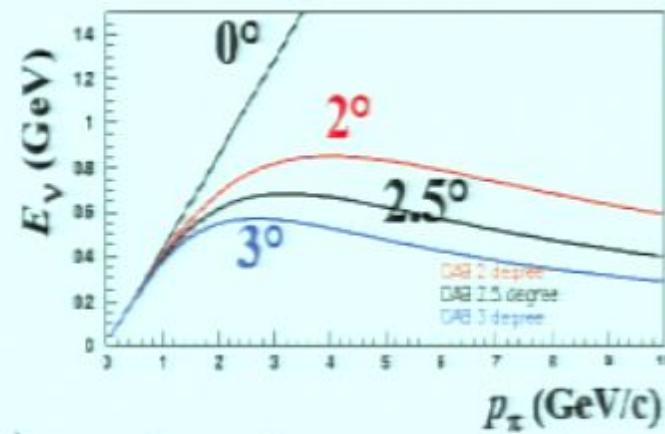
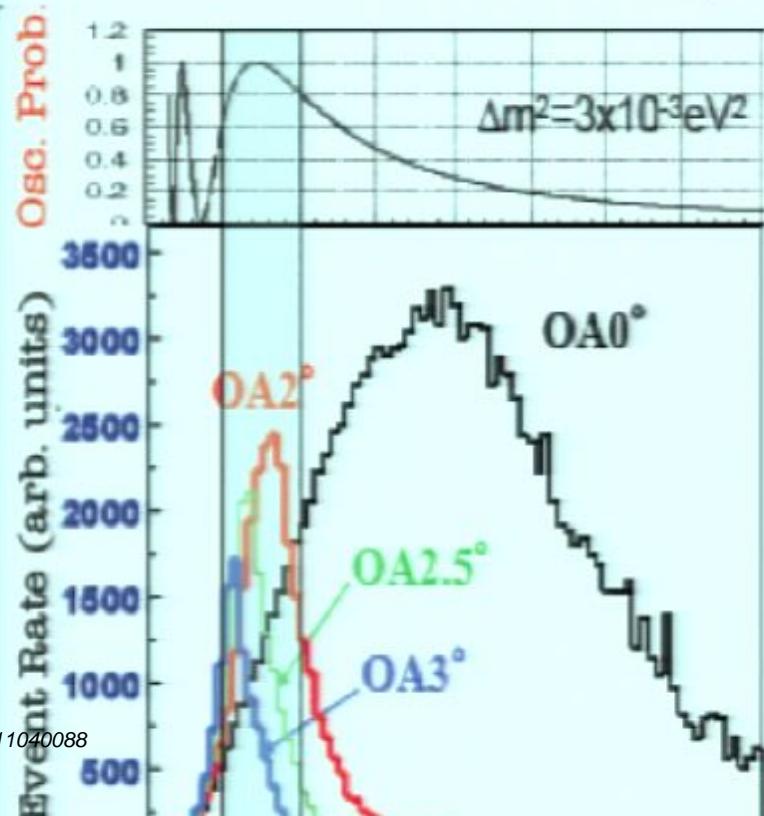
Tokai accelerator complex and location of near detector (ND280)



Creating an (offaxis) neutrino beam



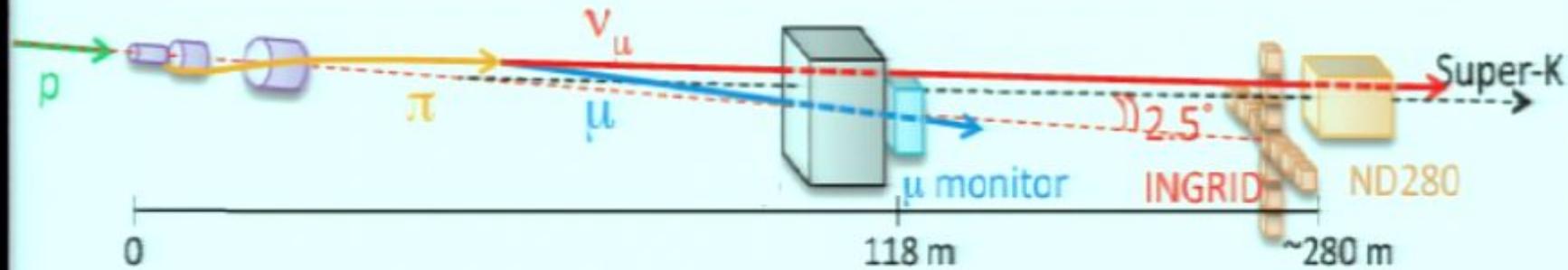
Primary protons hit a target (carbon) producing secondary unstable mesons (π, K) which decay to a tertiary ν_μ beam



At angles away from the parent pion's direction, the neutrino energy is independent of pion momentum, resulting in a narrower neutrino energy spectrum

Peak corresponds to oscillation maximum
Reduces backgrounds from higher energy

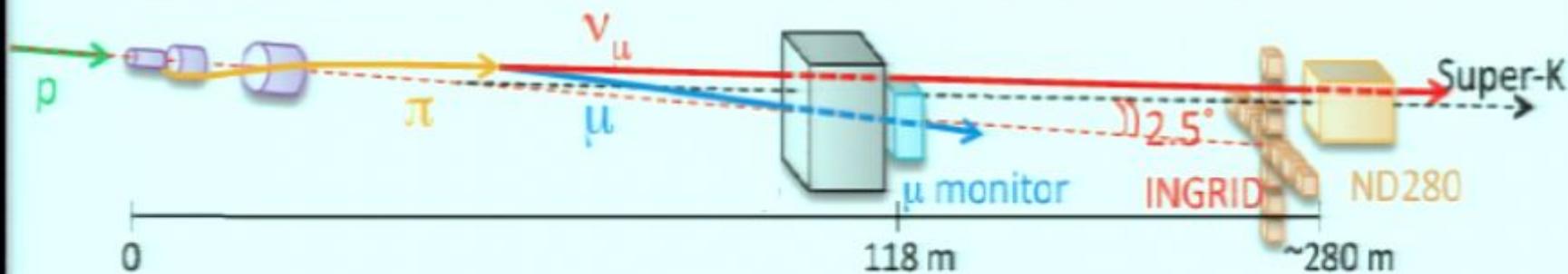
The proton beam



Start with 30 GeV protons, produced at JPARC in Tokai-mura, Japan

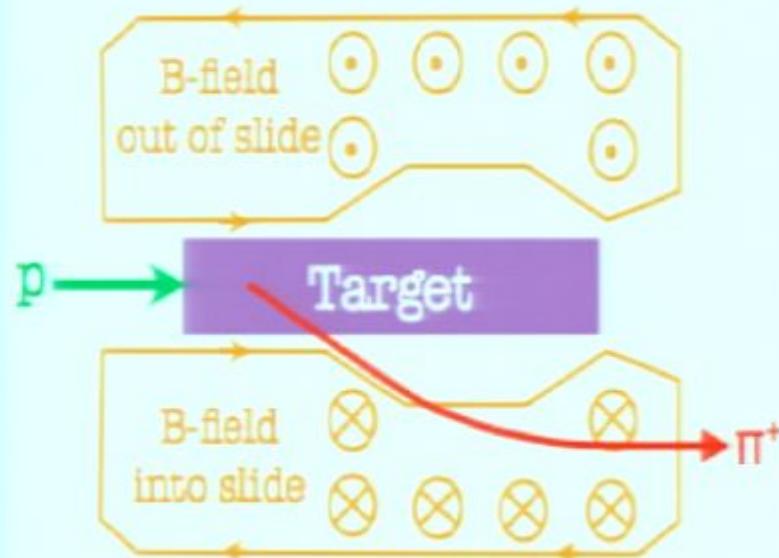


Secondary beam of mesons

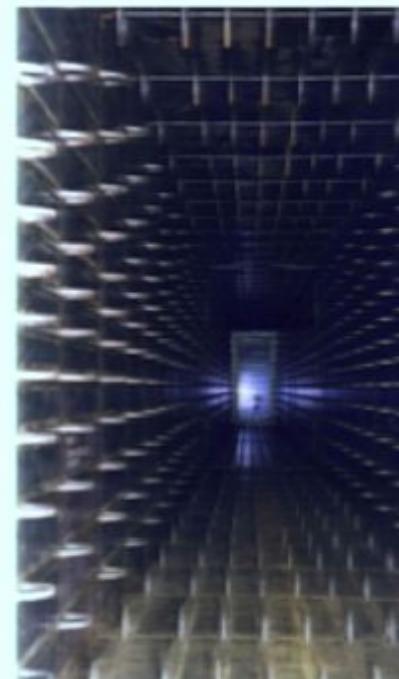


The protons hit a 91 cm long graphite target, producing pions and kaons

The mesons are focused by three magnetic "horns"



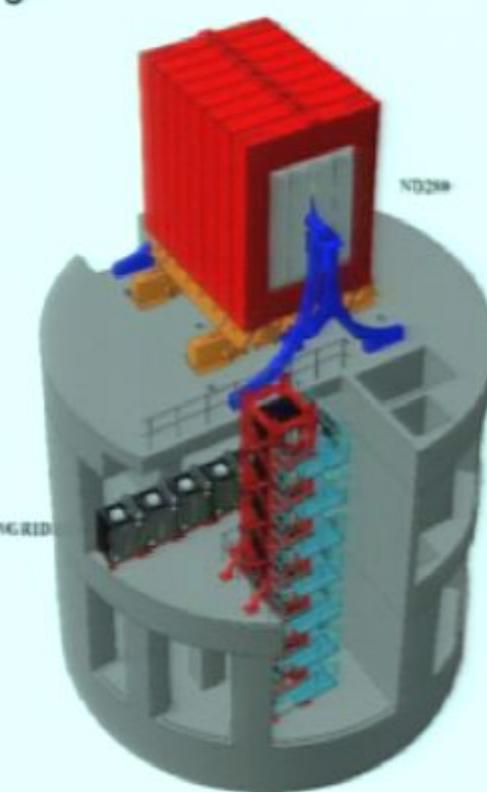
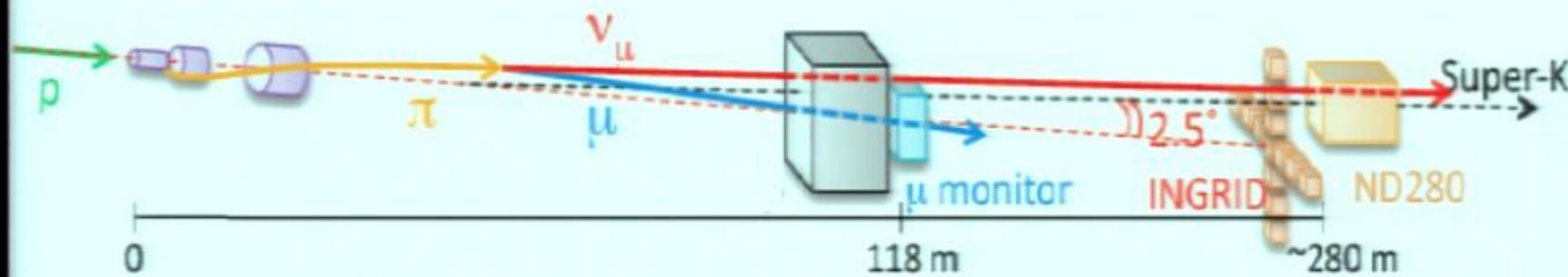
then decay in a 100m long decay volume



$$\pi \rightarrow \mu^- + \nu_\mu$$

The muons are sampled using a pair of muon monitors as a real-time neutrino beam monitor

Neutrino detectors



Off-axis with the ND280 detectors

Measure the unoscillated ν_μ rate

Constrain background processes

On-axis with the INGRID detector

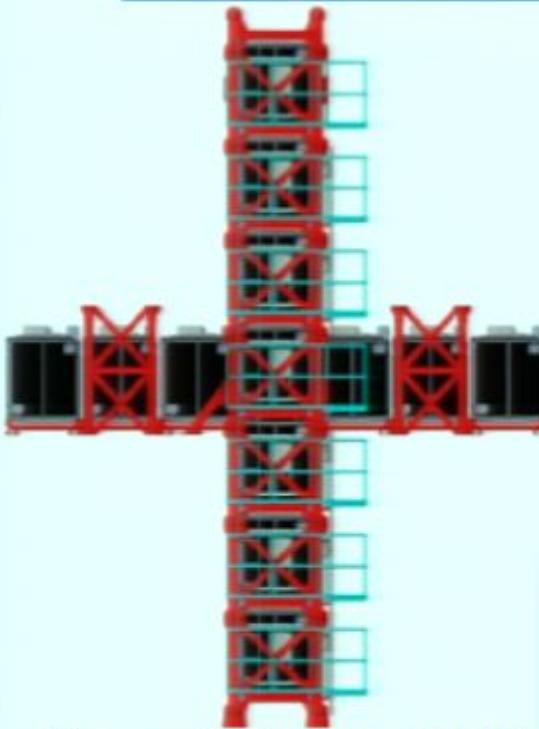
Determine neutrino beam direction



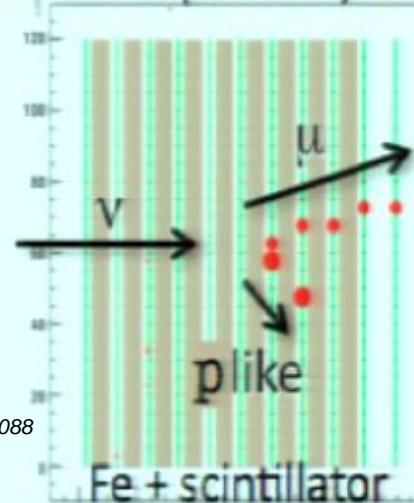
11040088 Off-axis@295km with Super-Kamiokande

Extract oscillation parameters

On-axis Interactive Neutrino GRID (INGRID)



1st event (Nov. 22, 2009)



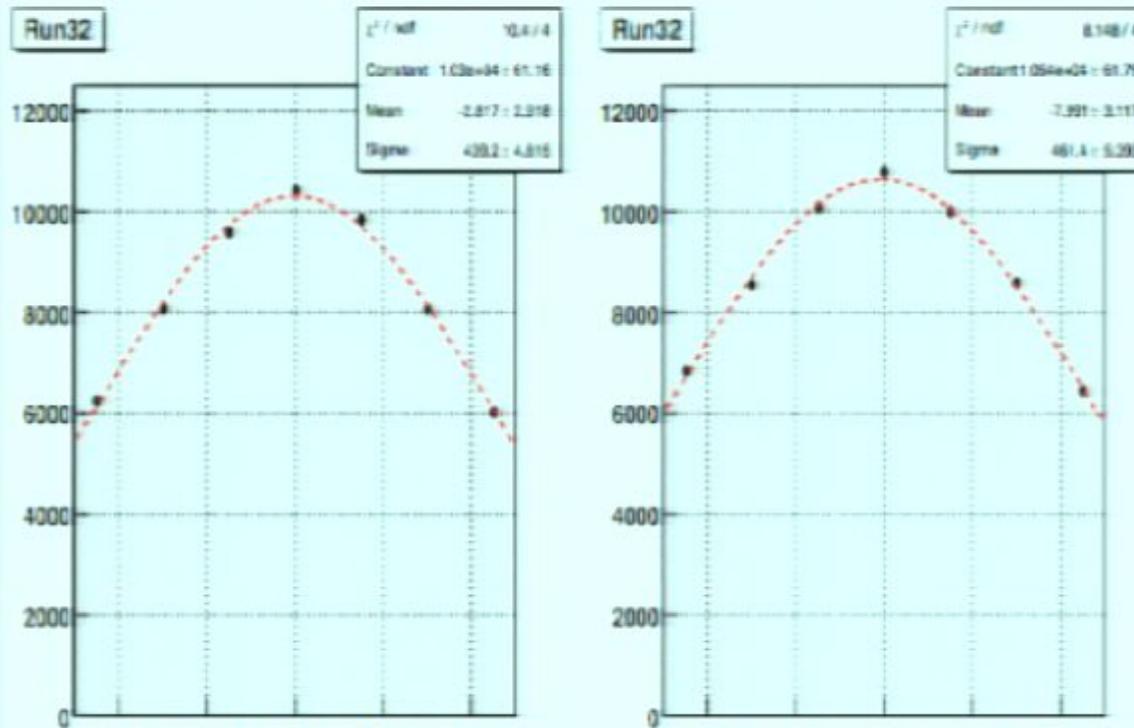
14 iron-scintillator modules arranged in a cross

X-Y scintillator layers

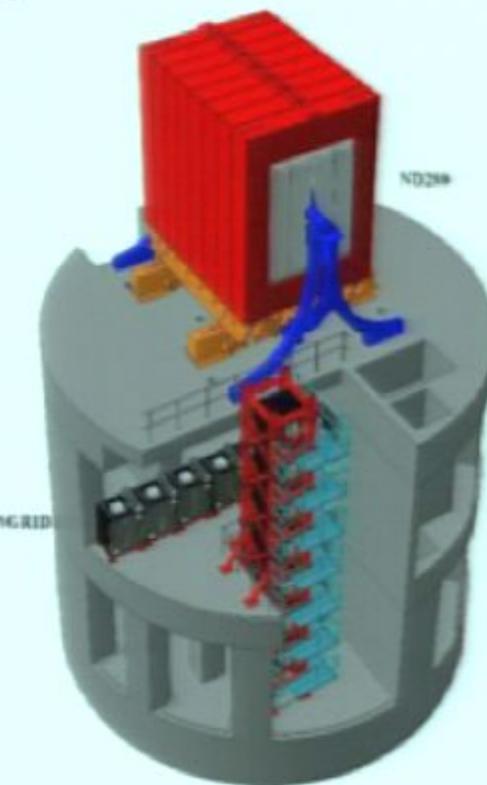
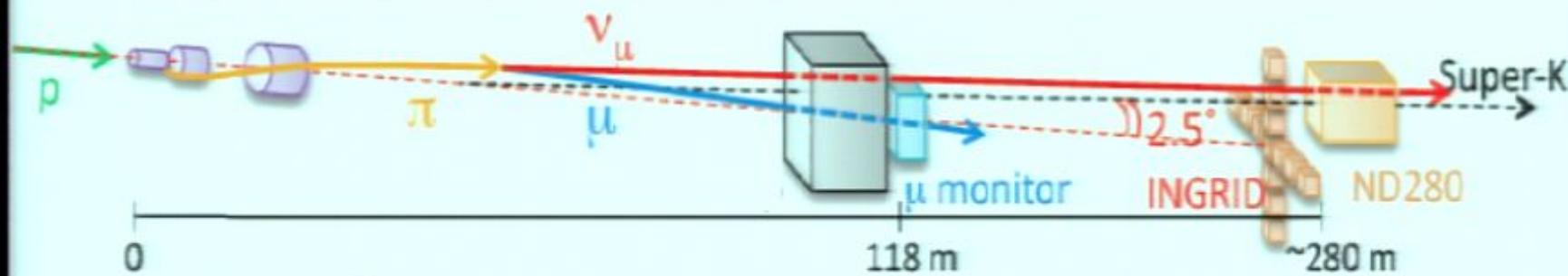
Count neutrino interactions in each module to determine neutrino rate vs. position

$\sim 1.5 \times 10^{14}$ protons on target
means 10,000 events / day

Extract beam direction better than 0.5 mrad



Neutrino detectors



Off-axis with the ND280 detectors

Measure the unoscillated ν_μ rate

Constrain background processes

On-axis with the INGRID detector

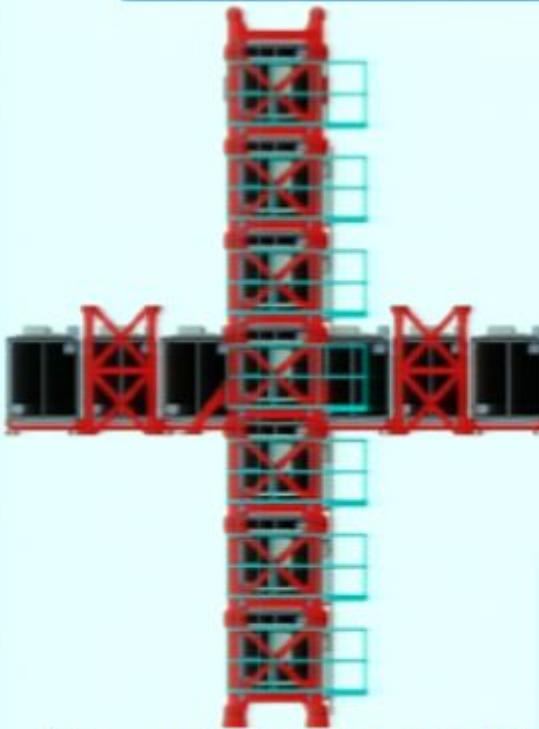
Determine neutrino beam direction



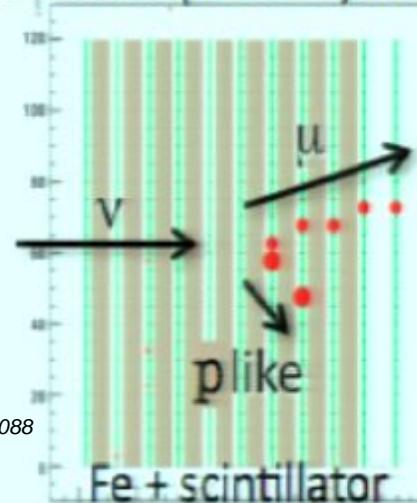
11040088 Off-axis@295km with Super-Kamiokande

Extract oscillation parameters

On-axis Interactive Neutrino GRID (INGRID)



1st event (Nov. 22, 2009)



11040088

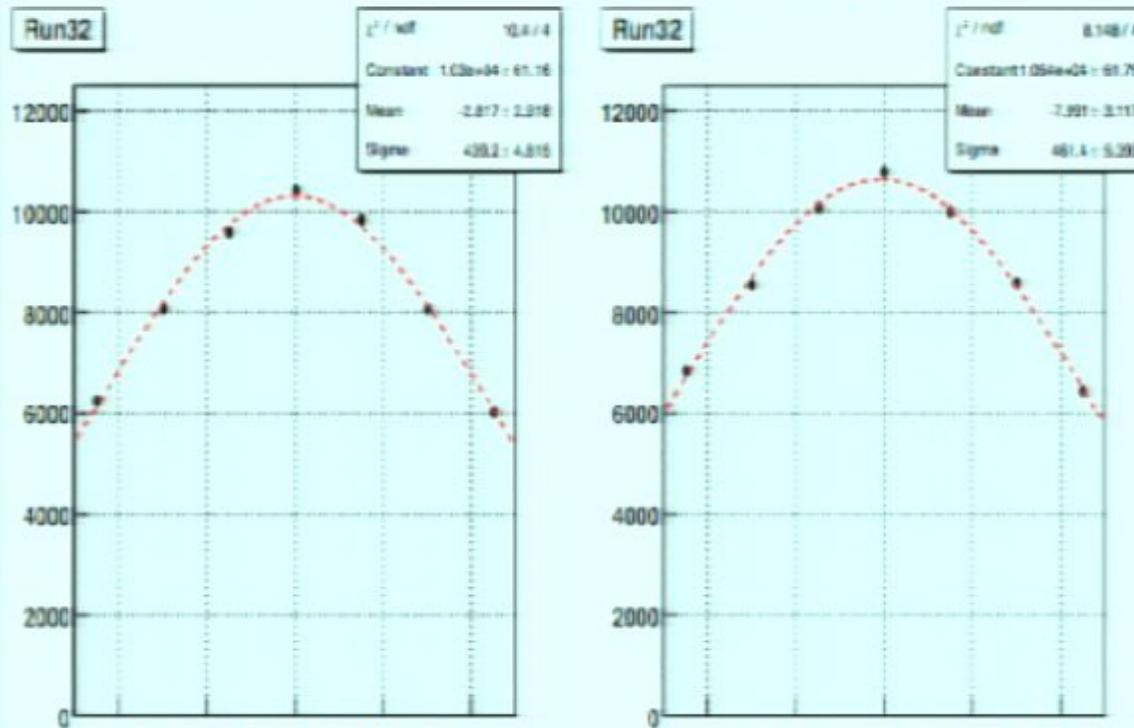
14 iron-scintillator modules arranged in a cross

X-Y scintillator layers

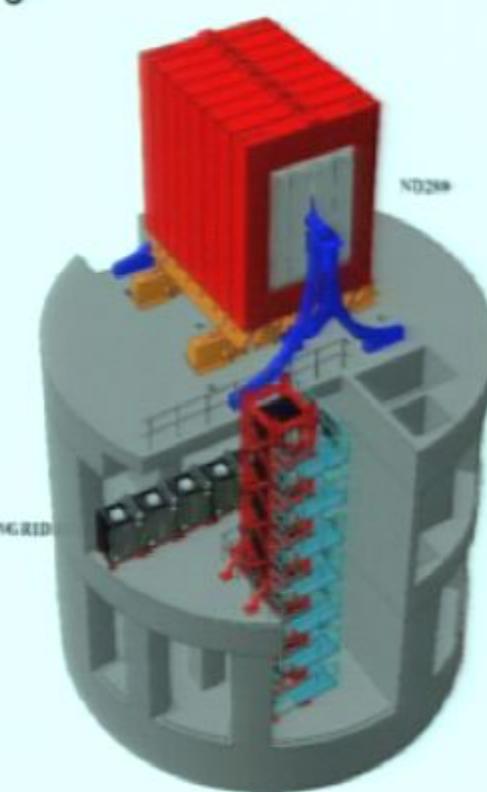
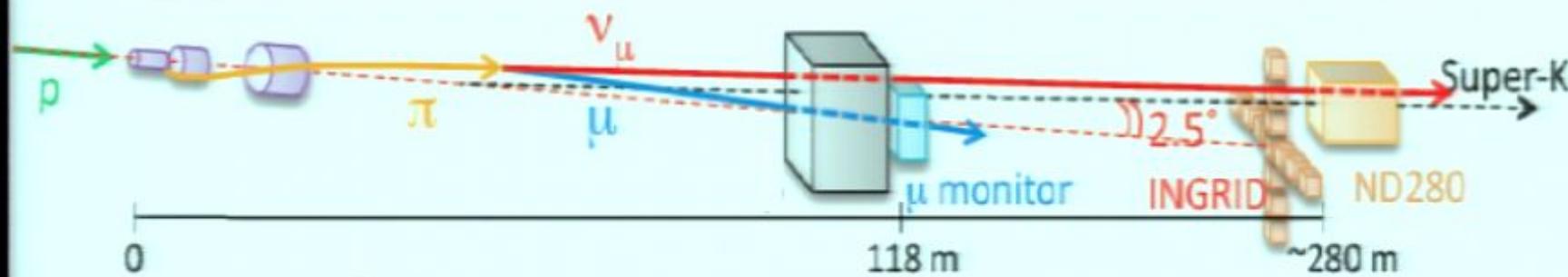
Count neutrino interactions in each module to determine neutrino rate vs. position

$\sim 1.5 \text{ v} / 10^{14}$ protons on target
means 10,000 events / day

Extract beam direction better than 0.5 mrad



Neutrino detectors



Off-axis with the ND280 detectors

Measure the unoscillated ν_μ rate

Constrain background processes

On-axis with the INGRID detector

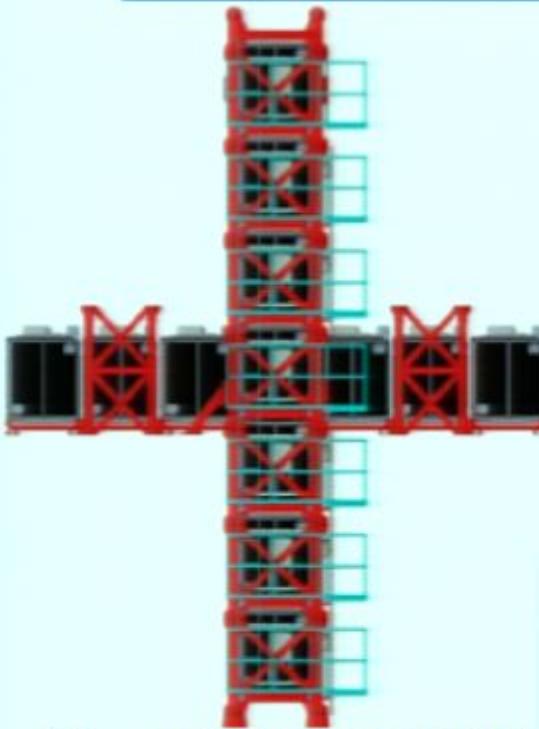
Determine neutrino beam direction



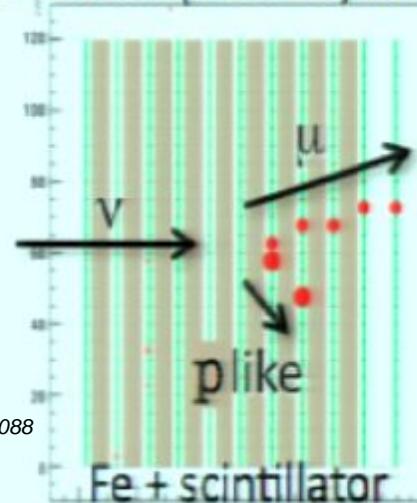
11040088 Off-axis@295km with Super-Kamiokande

Extract oscillation parameters

On-axis Interactive Neutrino GRID (INGRID)



1st event (Nov. 22, 2009)



14 iron-scintillator modules arranged in a cross

X-Y scintillator layers

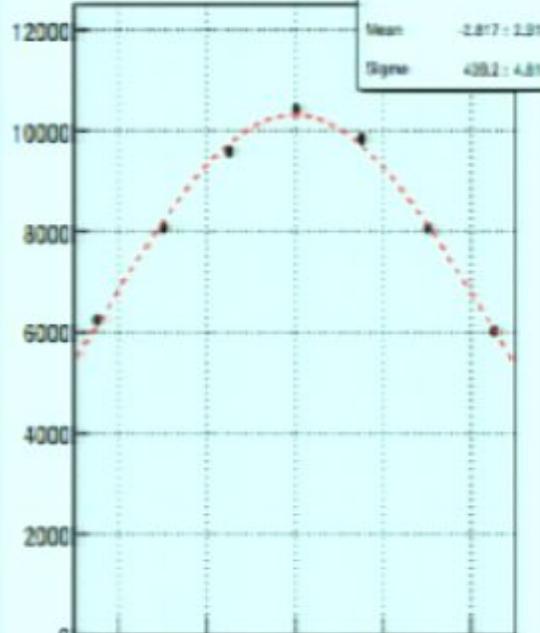
Count neutrino interactions in each module to determine neutrino rate vs. position

$\sim 1.5 \text{ v} / 10^{14}$ protons on target

means 10,000 events / day

Extract beam direction better than 0.5 mrad

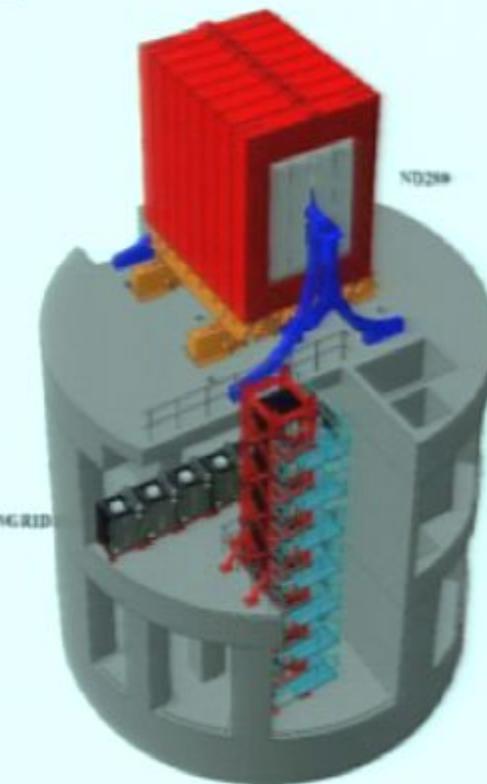
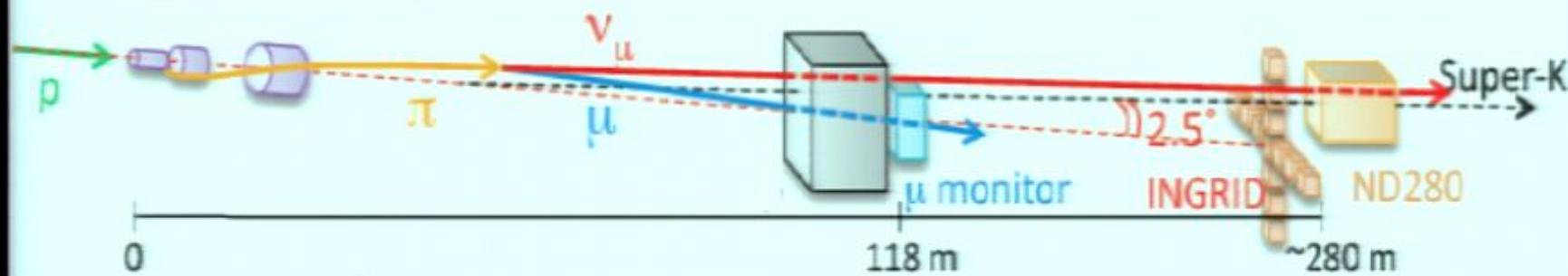
Run32



Run32



Neutrino detectors



Off-axis with the ND280 detectors

Measure the unoscillated ν_{μ} rate

Constrain background processes

On-axis with the INGRID detector

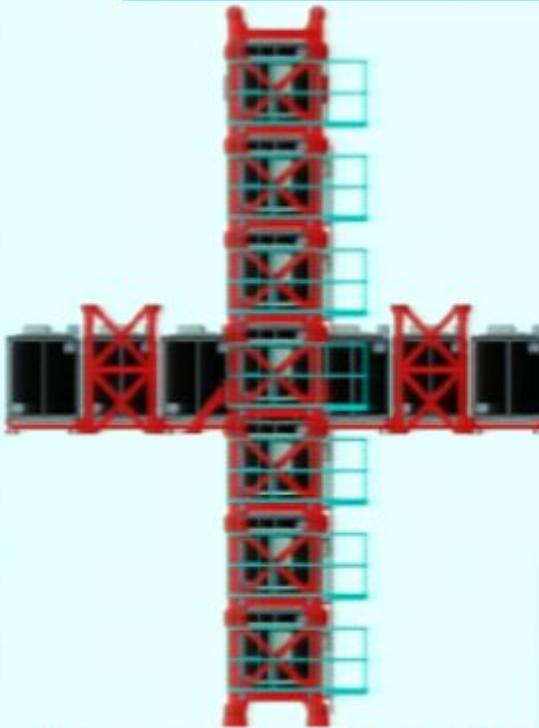
Determine neutrino beam direction



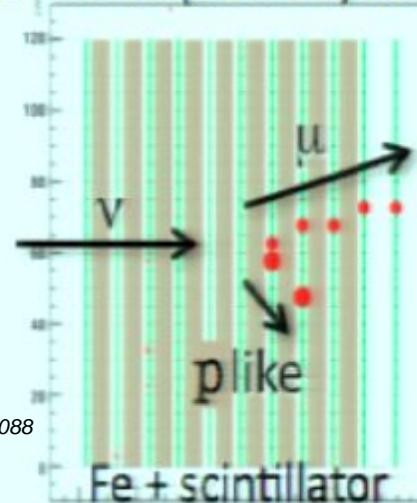
11040088 Off-axis@295km with Super-Kamiokande

Extract oscillation parameters

On-axis Interactive Neutrino GRID (INGRID)



1st event (Nov. 22, 2009)



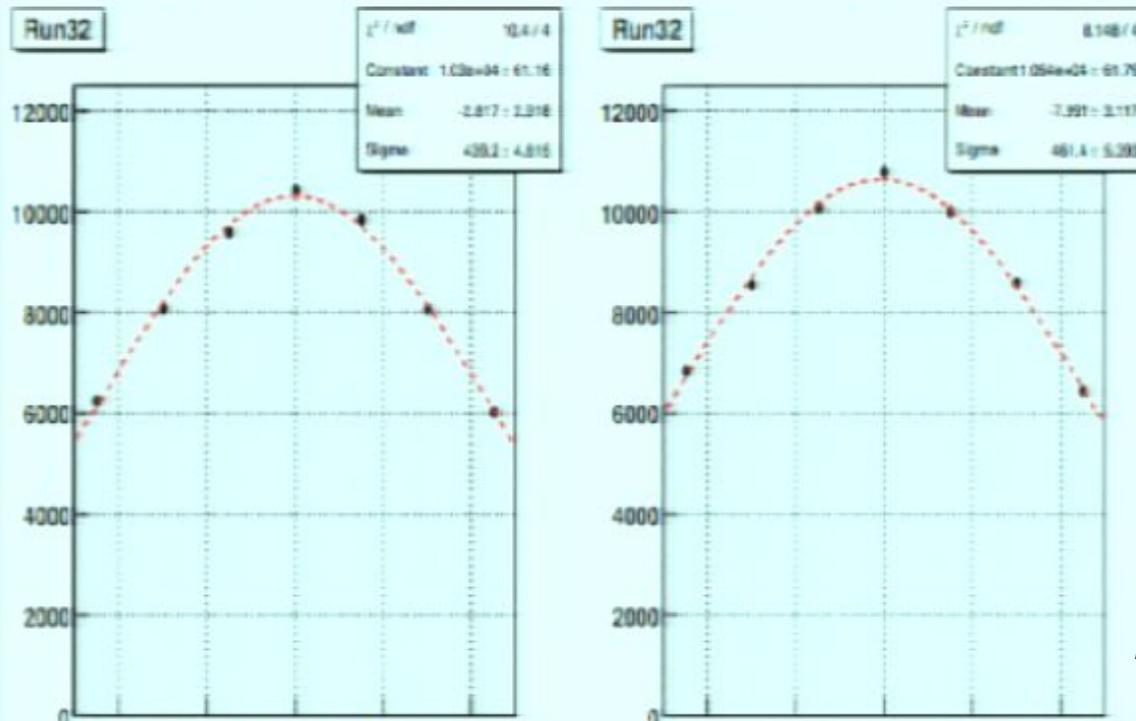
14 iron-scintillator modules arranged in a cross

X-Y scintillator layers

Count neutrino interactions in each module to determine neutrino rate vs. position

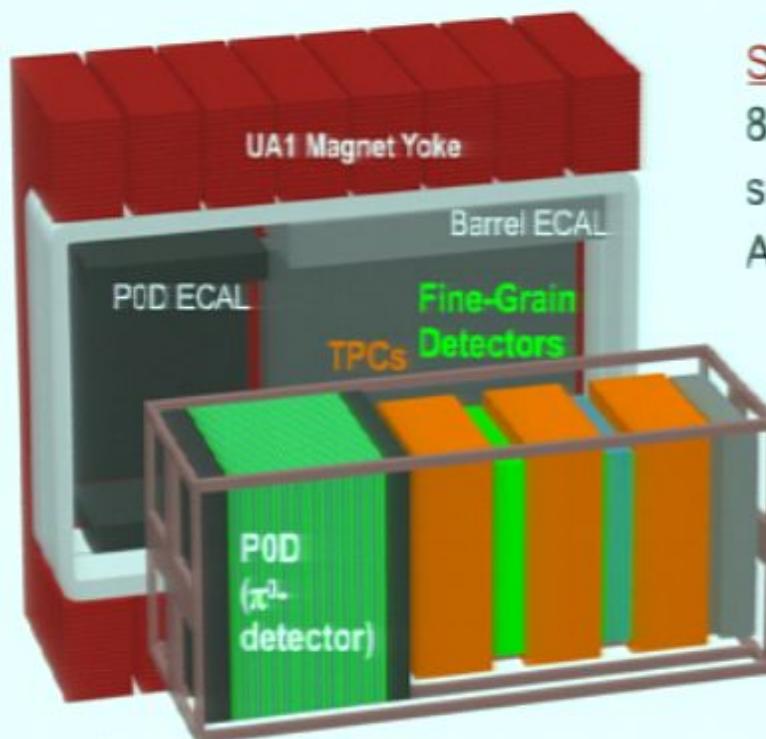
$\sim 1.5 \text{ v} / 10^{14}$ protons on target
means 10,000 events / day

Extract beam direction better than 0.5 mrad



Off-axis: ND280 detector complex

Suite of near detectors sit within UA1 ($B=0.2\text{T}$) magnet to measure the unoscillated neutrino rate, and relevant background processes to the oscillation analysis



Pi-zero Detector (POD)

Pb/brass/scintillator planes interspersed with water bags
Measure $\text{NC}\pi^0$ interactions
 π^0 decays to 2 photons, photons

Side Muon Range Detector

87x17x0.7cm instrumented scintillator in magnet yoke
Active veto, cosmic trigger

Electromagnetic Calorimeters

X-Y Pb/scintillator planes
P0D, Barrel, TPC3
Downstream ECAL Tag photons, e from Tracker ($\text{CC } v_e$) and P0D ($\text{NC}\pi^0$)

Tracker (FGD & TPC)

Measure $\text{CC } v_\mu, v_e$ interactions

Tracker: Fine Grained Detectors

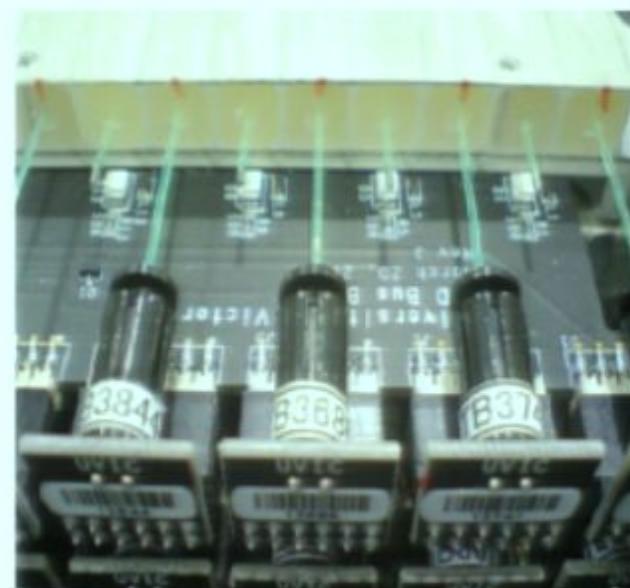
Tracking detector with planes in x or y of scintillator bars (8448 in total)

Read wavelength shifting fibers and
Multi-Pixel-Photon-Counters (MPPCs)

667 avalanche photodiodes in parallel

Functions in a magnetic field

Used for INGRID and ND280 tracking
detectors; first large scale use

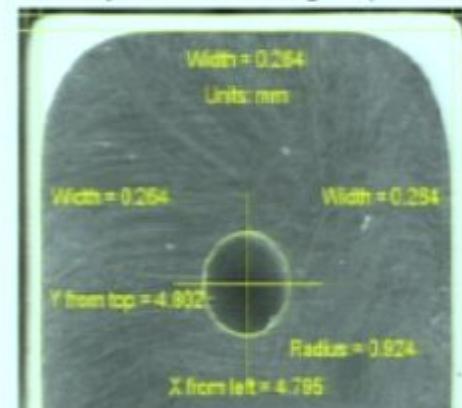


FGDs are neutrino interaction target with

2nd FGD has water planes to compare to far detector (water target)

1cm x 1cm bar granularity provides vertex information

Particle identification from energy loss in scintillator

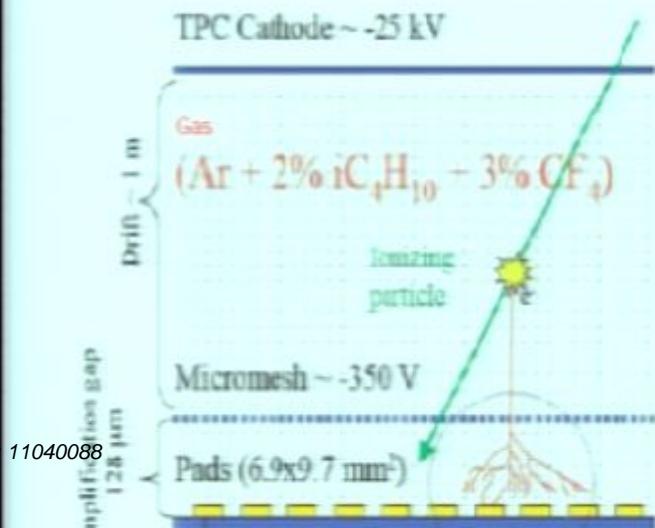
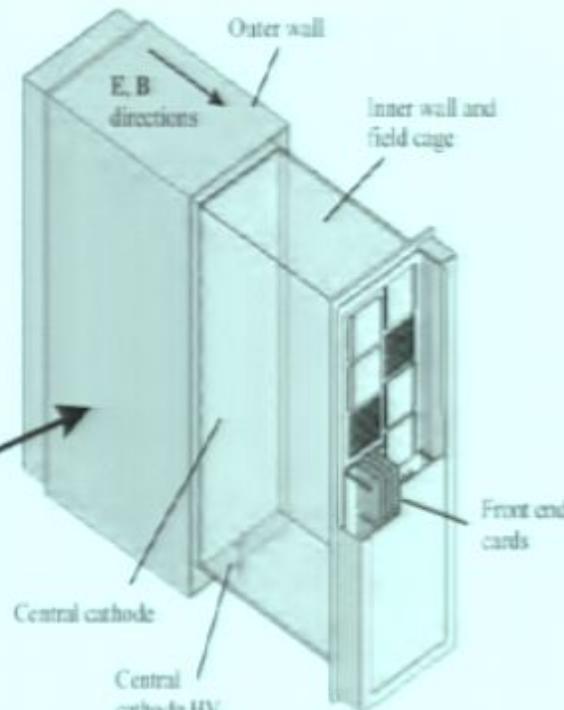


Tracker: Time Projection Chambers

Charged particles ionize the gas

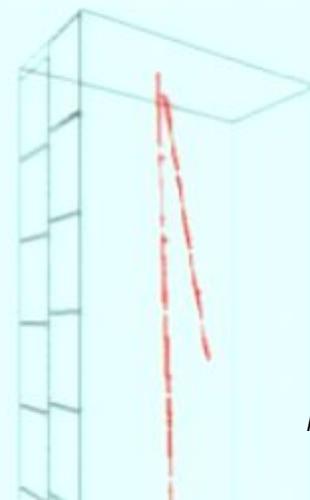
Electrons drift toward the readout plane due to the uniform E field (central cathode at -25kV)

"Wireless" TPC: Strong local field at the novel MicroMegas mesh creates a shower of electrons read out on $6.9 \times 9.7 \text{ mm}^2$ pads (1728 pads on each of 12 MicroMegas, 128 μm beyond the mesh)

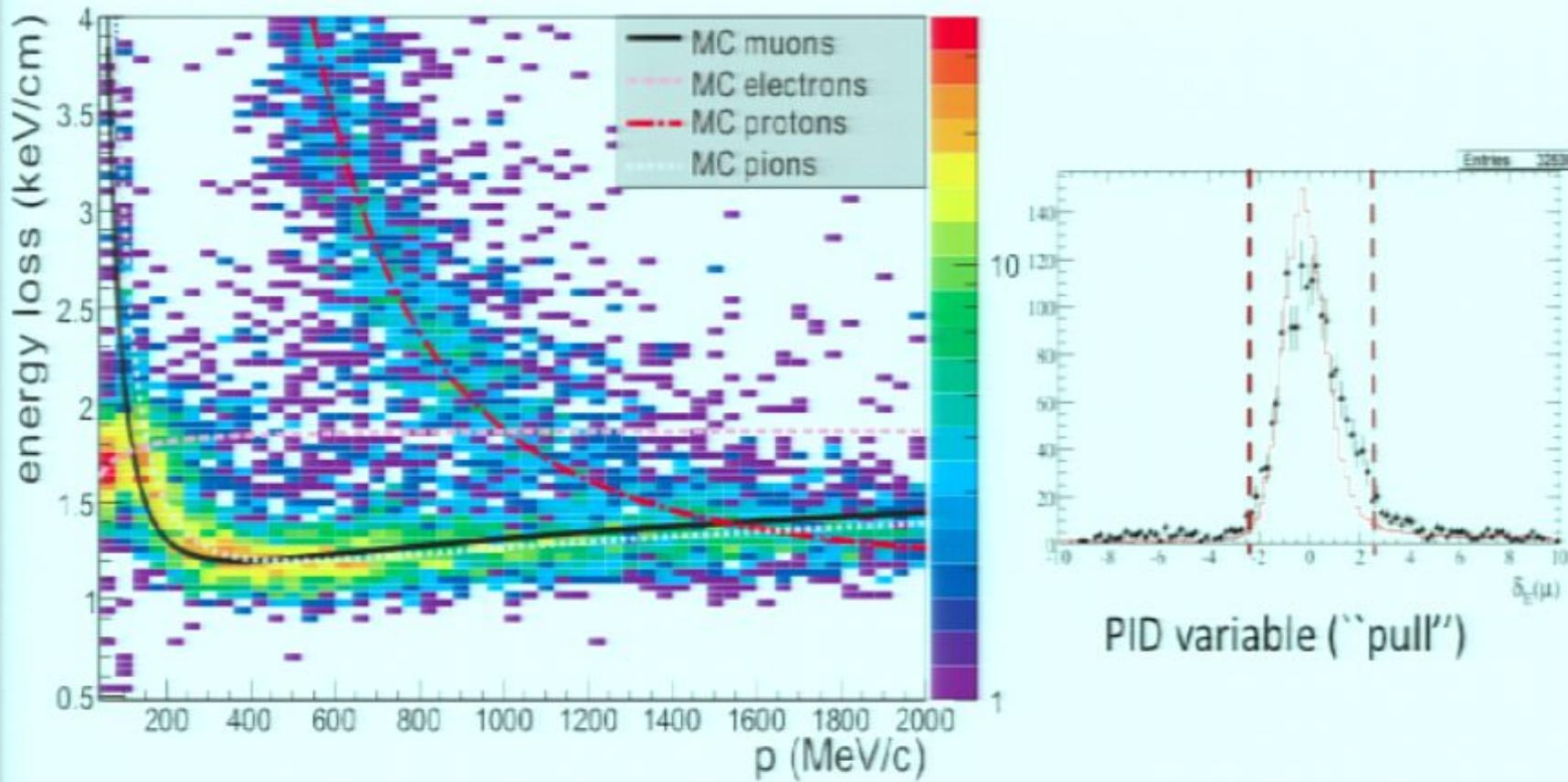


3D tracks are reconstructed provided drift velocity in the gas and timing of entry into TPC (from other subdetectors)

Spatial resolution:
600 μm @ 100cm drift distance



Particle ID with the TPC



Energy loss (dE/dx) in the TPC can be used to distinguish particle types
 dE/dx resolution ~8%

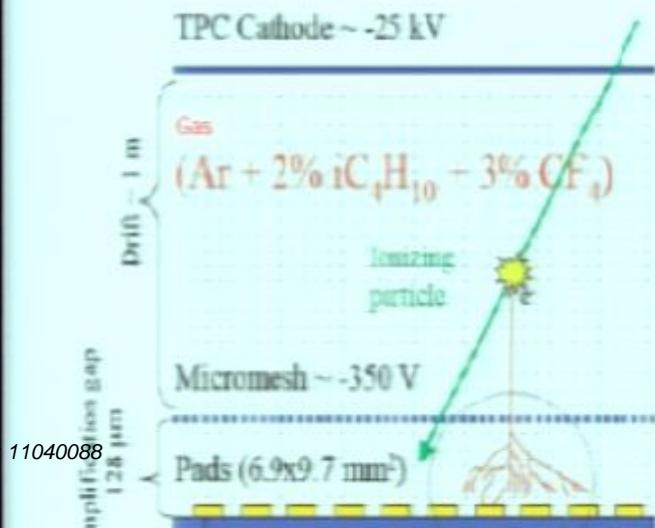
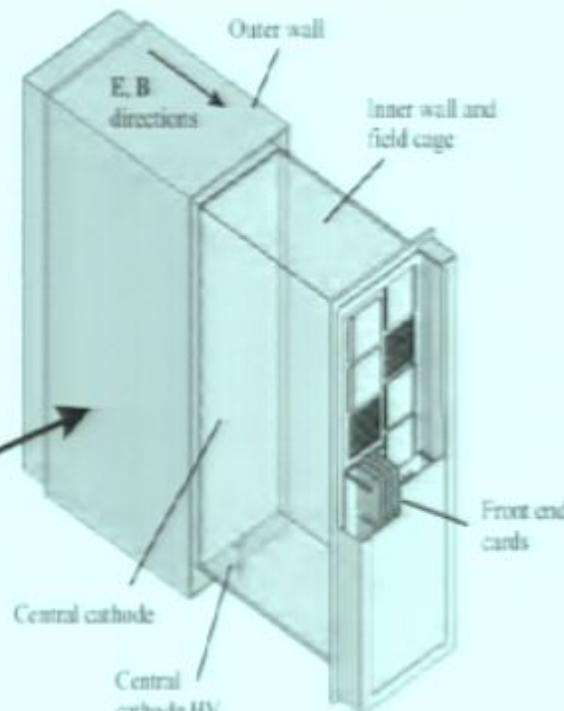
Select muon candidates for events consistent within 2.5 sigma of predicted muon energy loss, and more than 2 sigma away from electron energy loss

Tracker: Time Projection Chambers

Charged particles ionize the gas

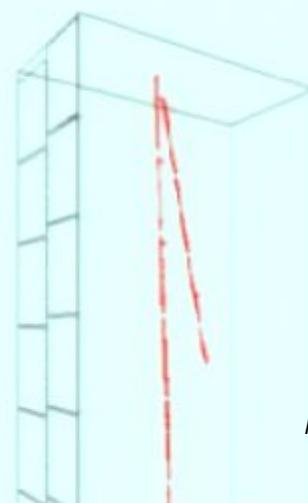
Electrons drift toward the readout plane due to the uniform E field (central cathode at -25kV)

"Wireless" TPC: Strong local field at the novel MicroMegas mesh creates a shower of electrons read out on $6.9 \times 9.7 \text{ mm}^2$ pads (1728 pads on each of 12 MicroMegas, 128 μm beyond the mesh)

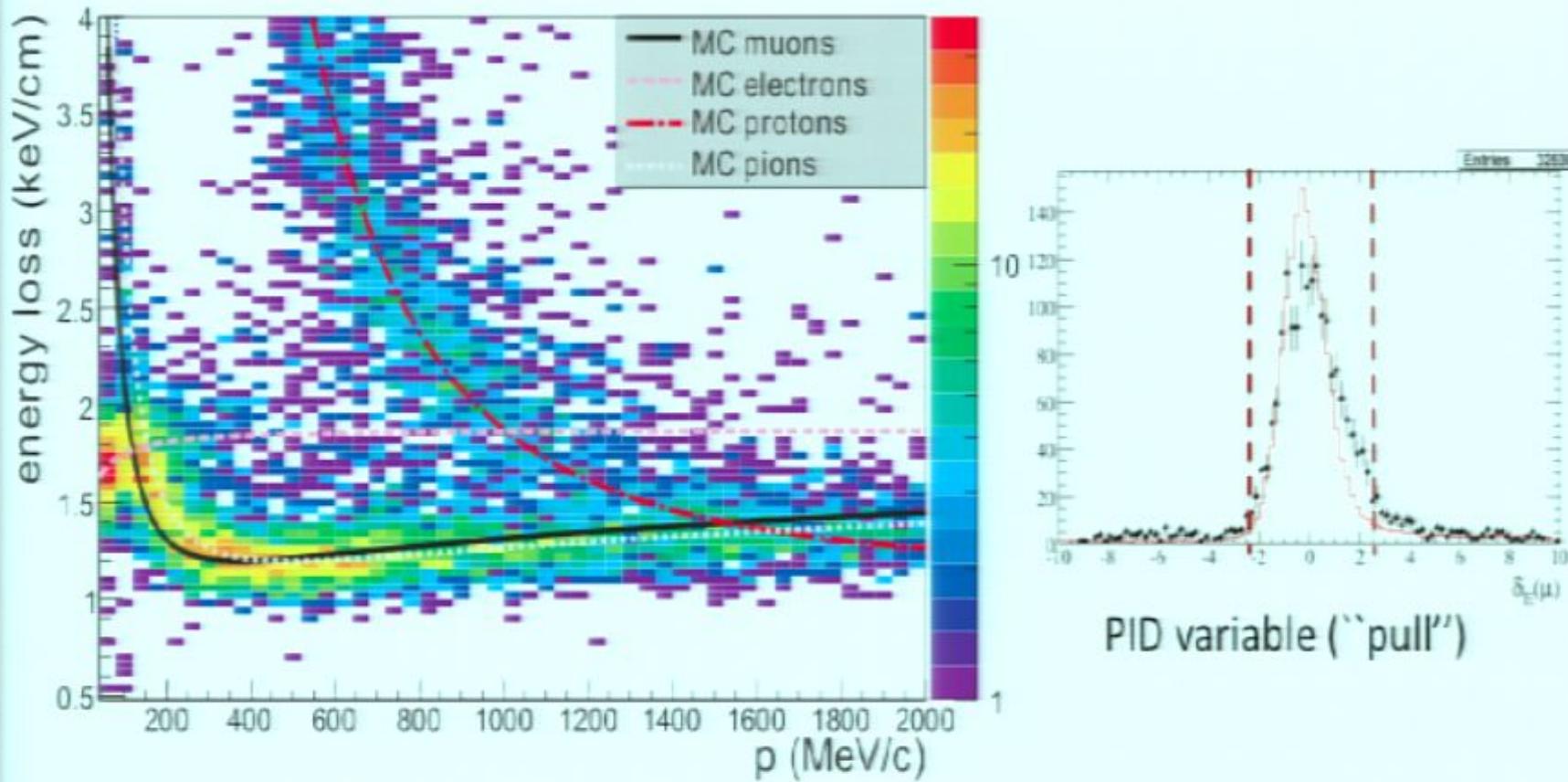


3D tracks are reconstructed provided drift velocity in the gas and timing of entry into TPC (from other subdetectors)

Spatial resolution:
600 μm @ 100cm drift distance



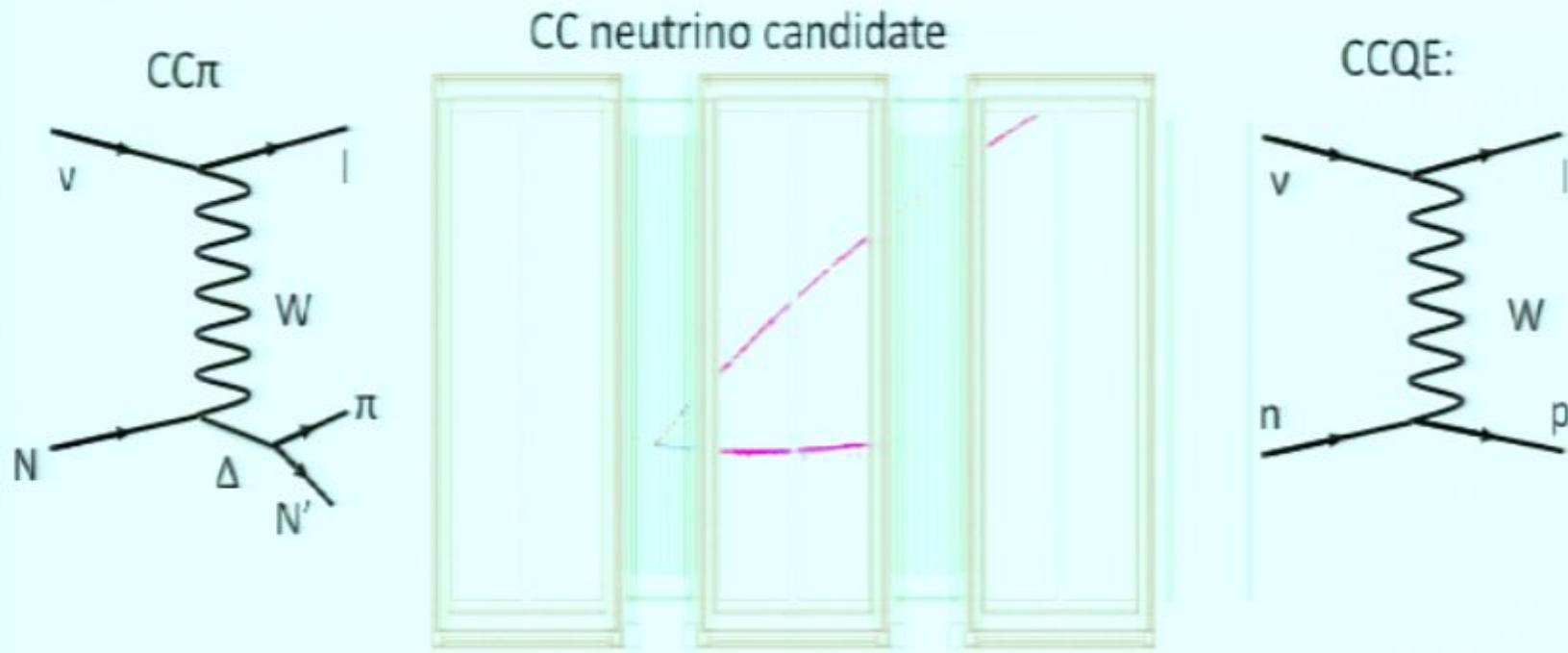
Particle ID with the TPC



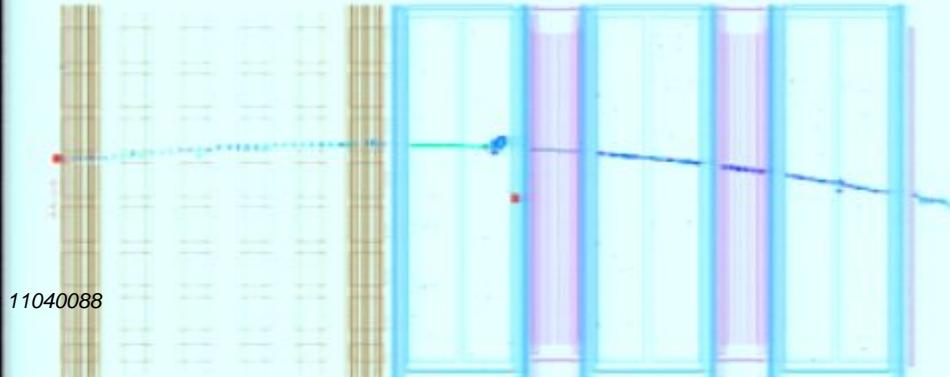
Energy loss (dE/dx) in the TPC can be used to distinguish particle types
 dE/dx resolution ~8%

Select muon candidates for events consistent within 2.5 sigma of predicted muon energy loss, and more than 2 sigma away from electron energy loss

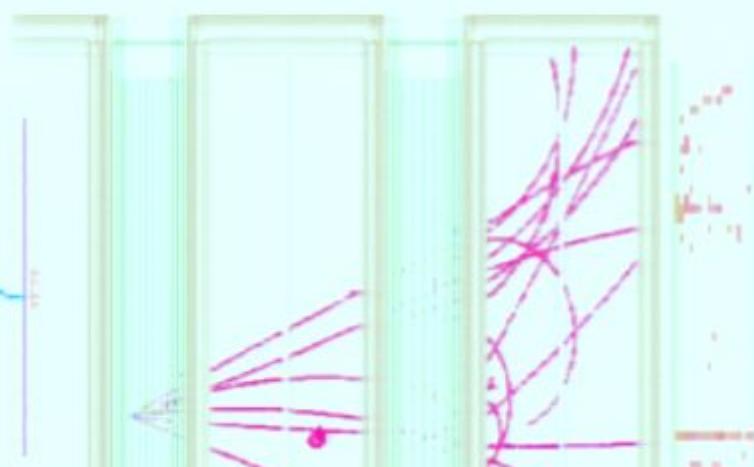
Example neutrino interactions in ND280



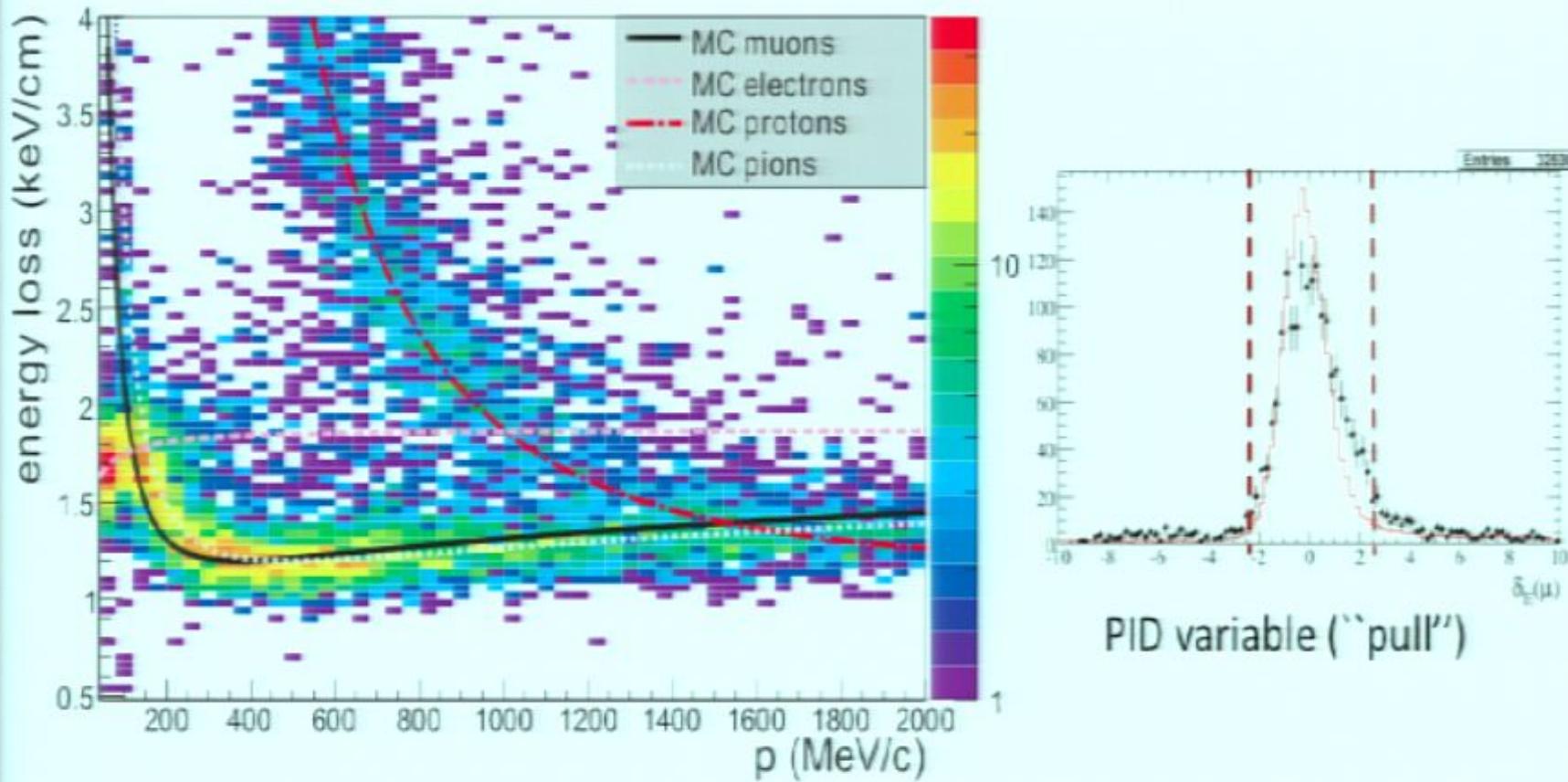
Neutrino interaction upstream of ND280
(sand muon)



More complex event (DIS?)



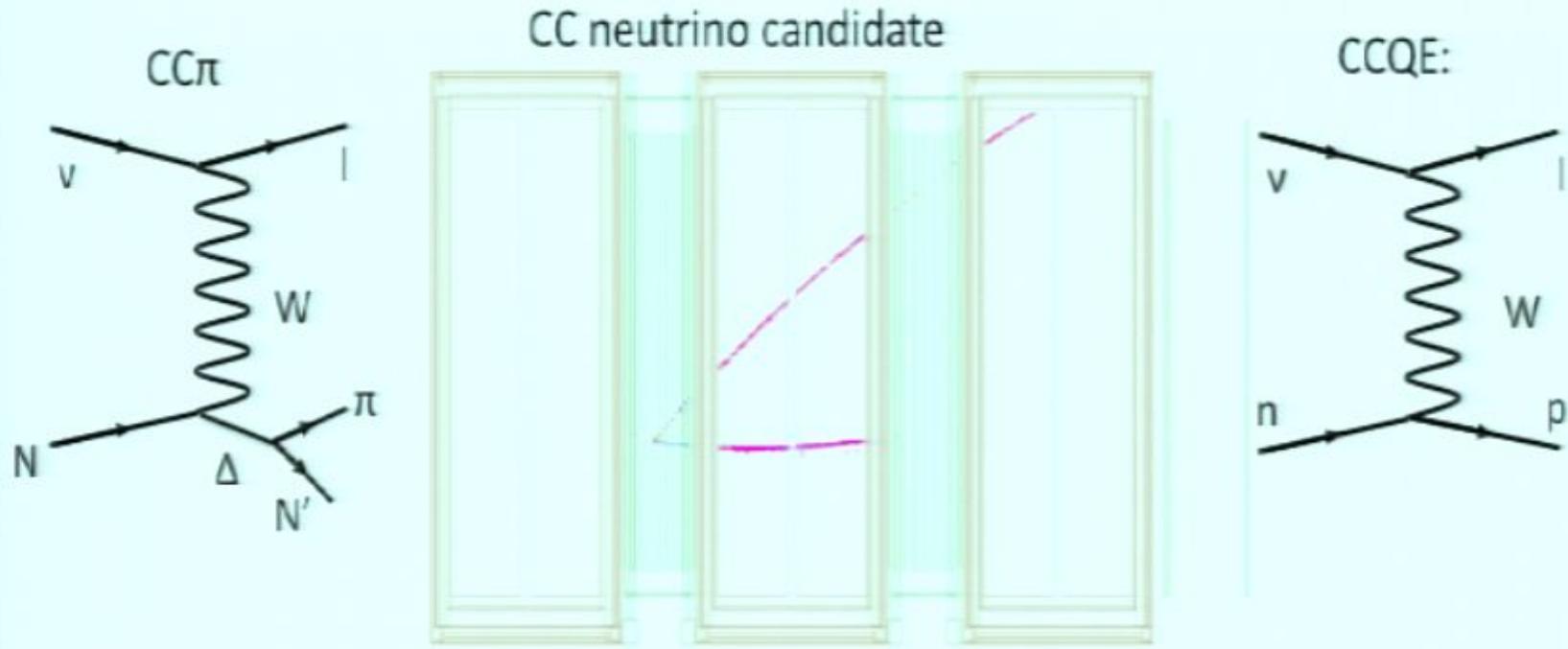
Particle ID with the TPC



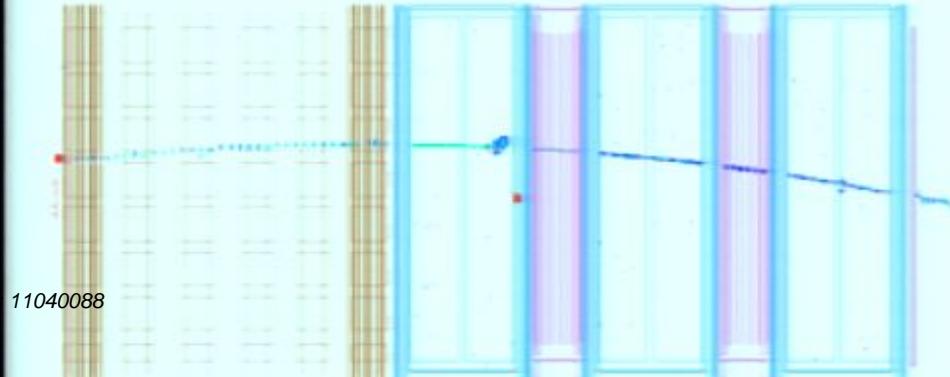
Energy loss (dE/dx) in the TPC can be used to distinguish particle types
 dE/dx resolution $\sim 8\%$

Select muon candidates for events consistent within 2.5 sigma of predicted muon energy loss, and more than 2 sigma away from electron energy loss

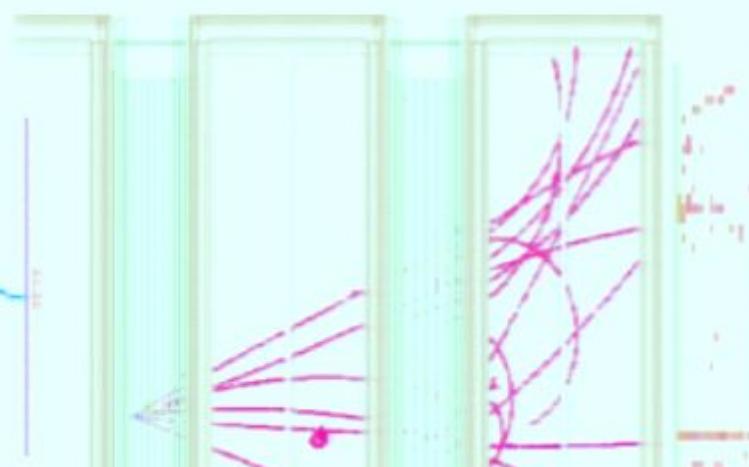
Example neutrino interactions in ND280



Neutrino interaction upstream of ND280
(sand muon)



More complex event (DIS?)



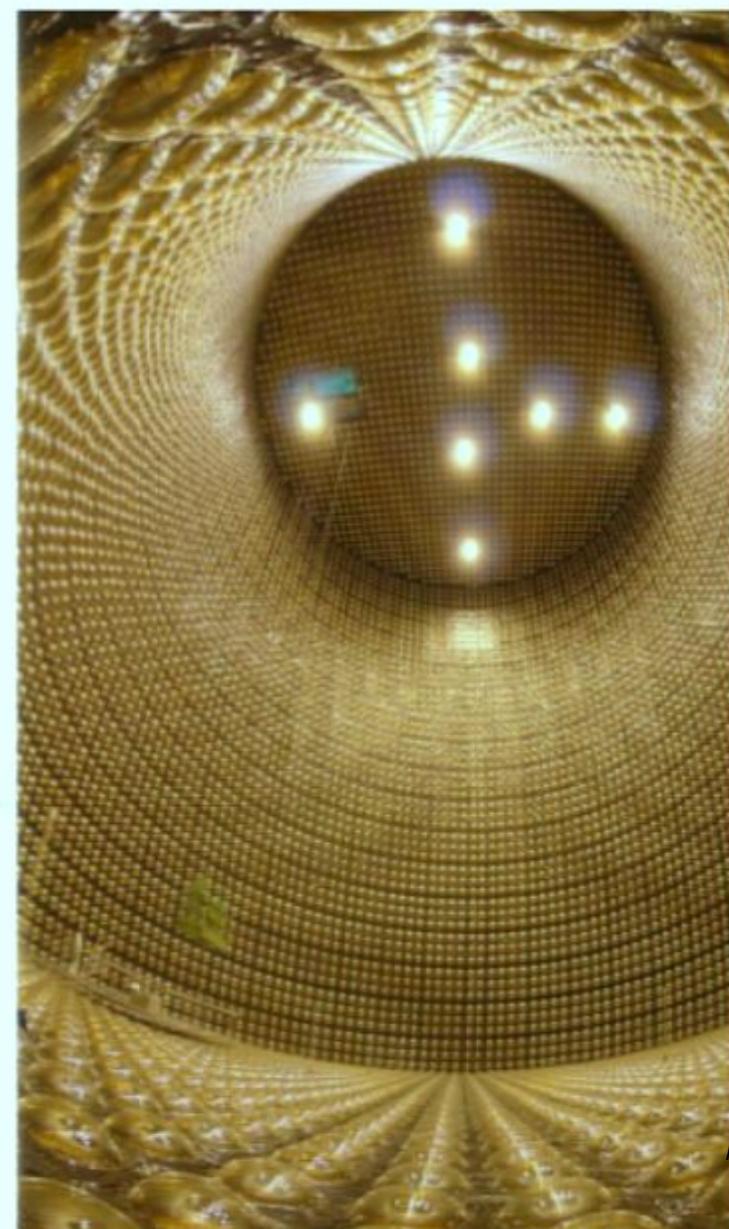
T2K far detector: Super-Kamiokande

50kton water Cherenkov detector
(22.5kton fiducial mass)

39.4 m diameter, 41.4 m tall
cylindrical tank lined with 11,129
photomultiplier tubes (PMTs)
40% photocathode coverage

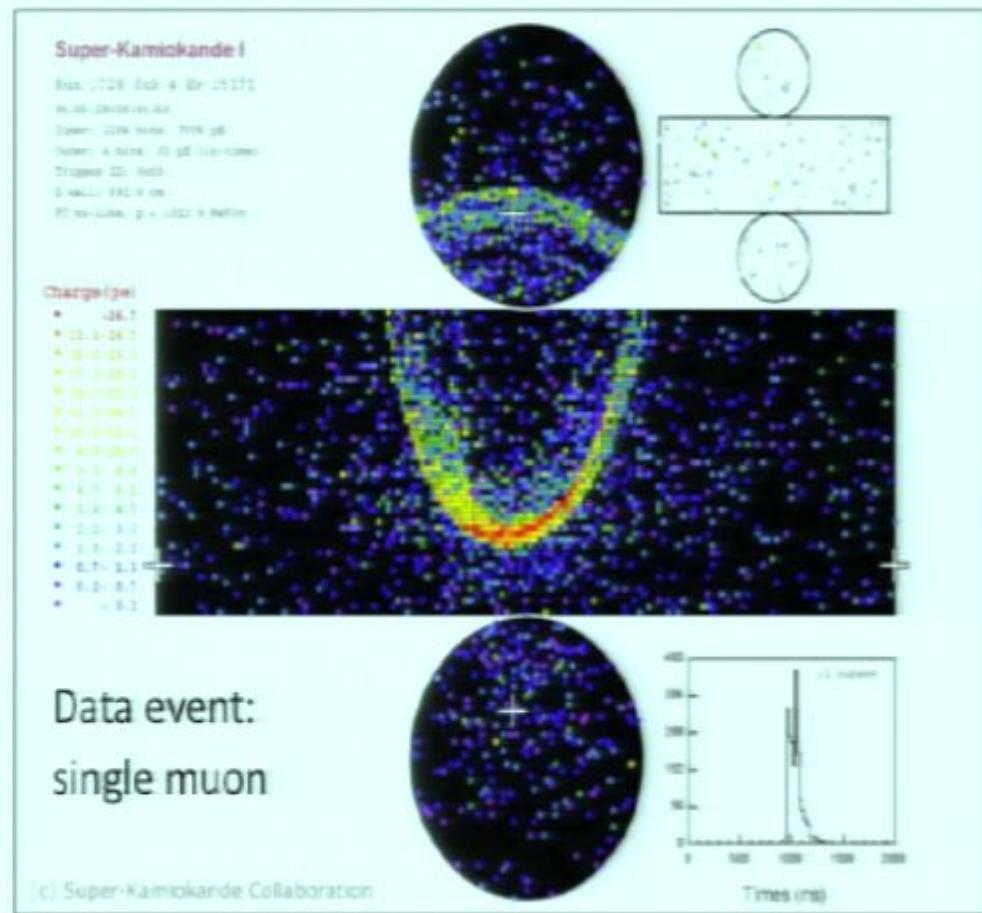
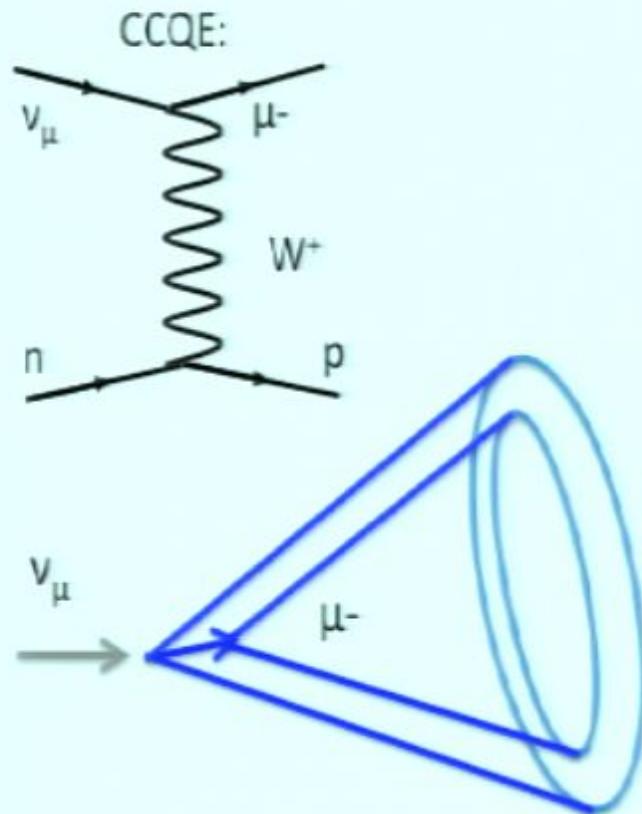
1885 veto PMTs located on the outside
of the tank reject events entering the tank,
such as cosmic rays

1.77 Hz rate of cosmics for an
overburden of 2700 meters,
water equivalent in Kamioka mine



Neutrino events in Super-K

Cherenkov light emitted at a fixed angle produces ring(s) on the tank wall, recorded by PMTs



Muons produce well defined rings

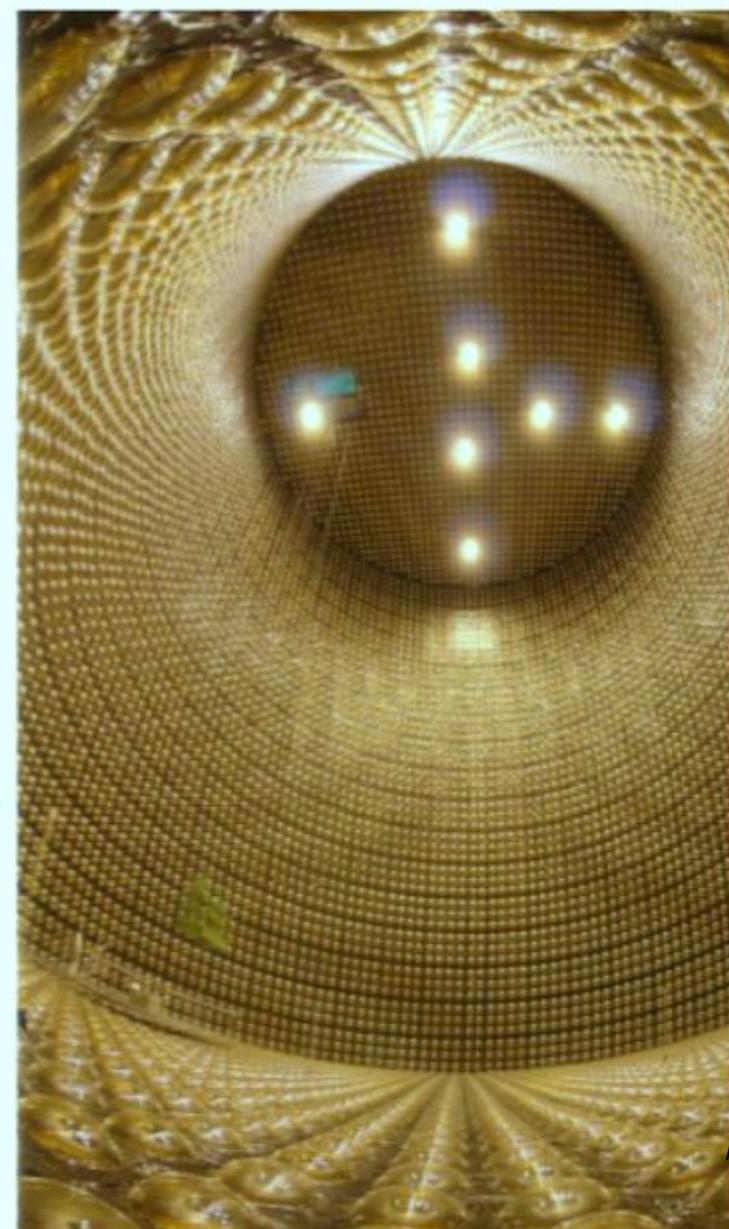
T2K far detector: Super-Kamiokande

50kton water Cherenkov detector
(22.5kton fiducial mass)

39.4 m diameter, 41.4 m tall
cylindrical tank lined with 11,129
photomultiplier tubes (PMTs)
40% photocathode coverage

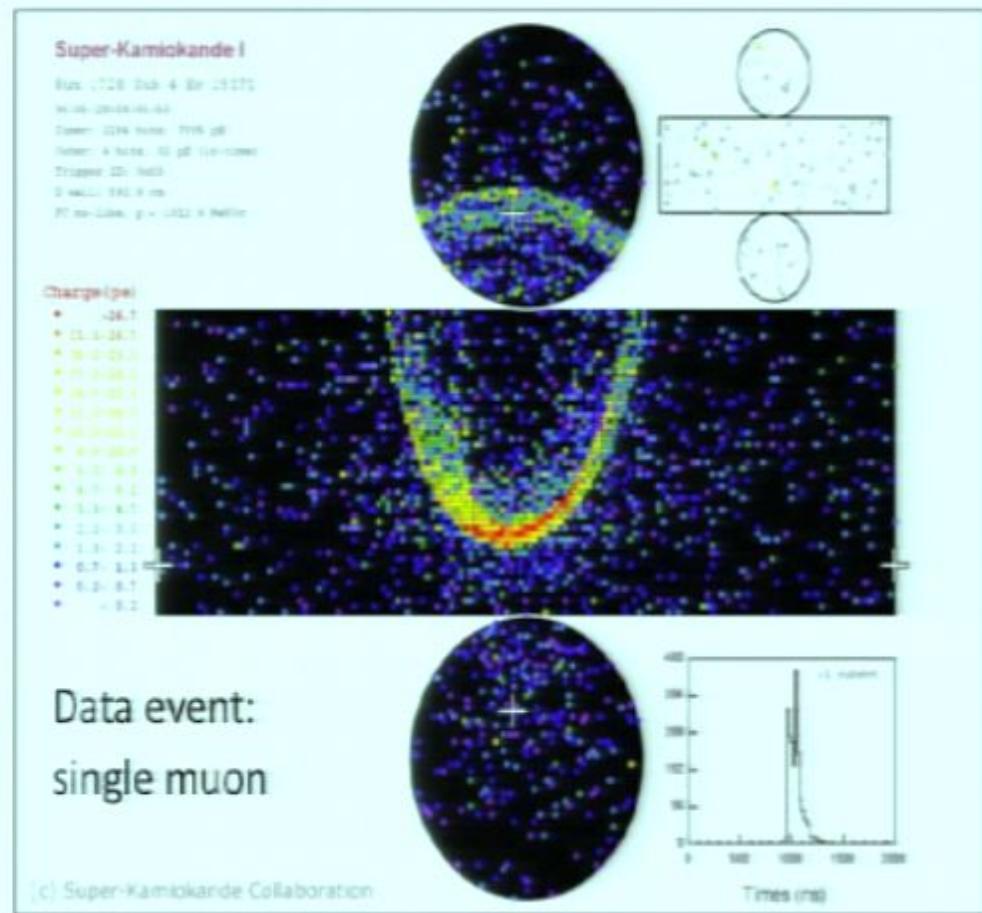
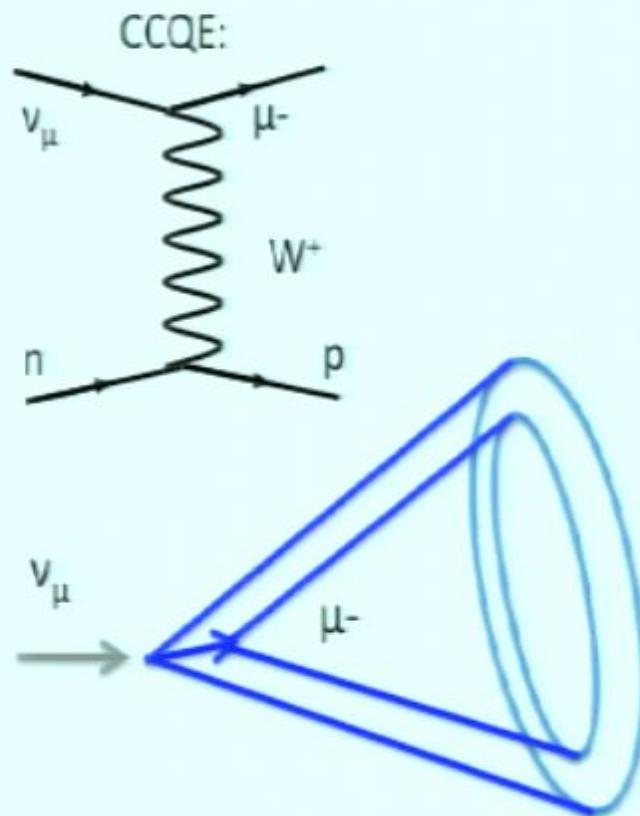
1885 veto PMTs located on the outside
of the tank reject events entering the tank,
such as cosmic rays

1.77 Hz rate of cosmics for an
overburden of 2700 meters,
water equivalent in Kamioka mine



Neutrino events in Super-K

Cherenkov light emitted at a fixed angle produces ring(s) on the tank wall, recorded by PMTs



Muons produce well defined rings

Neutrino events in Super-K

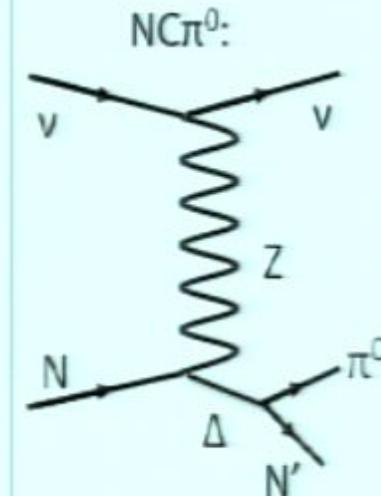
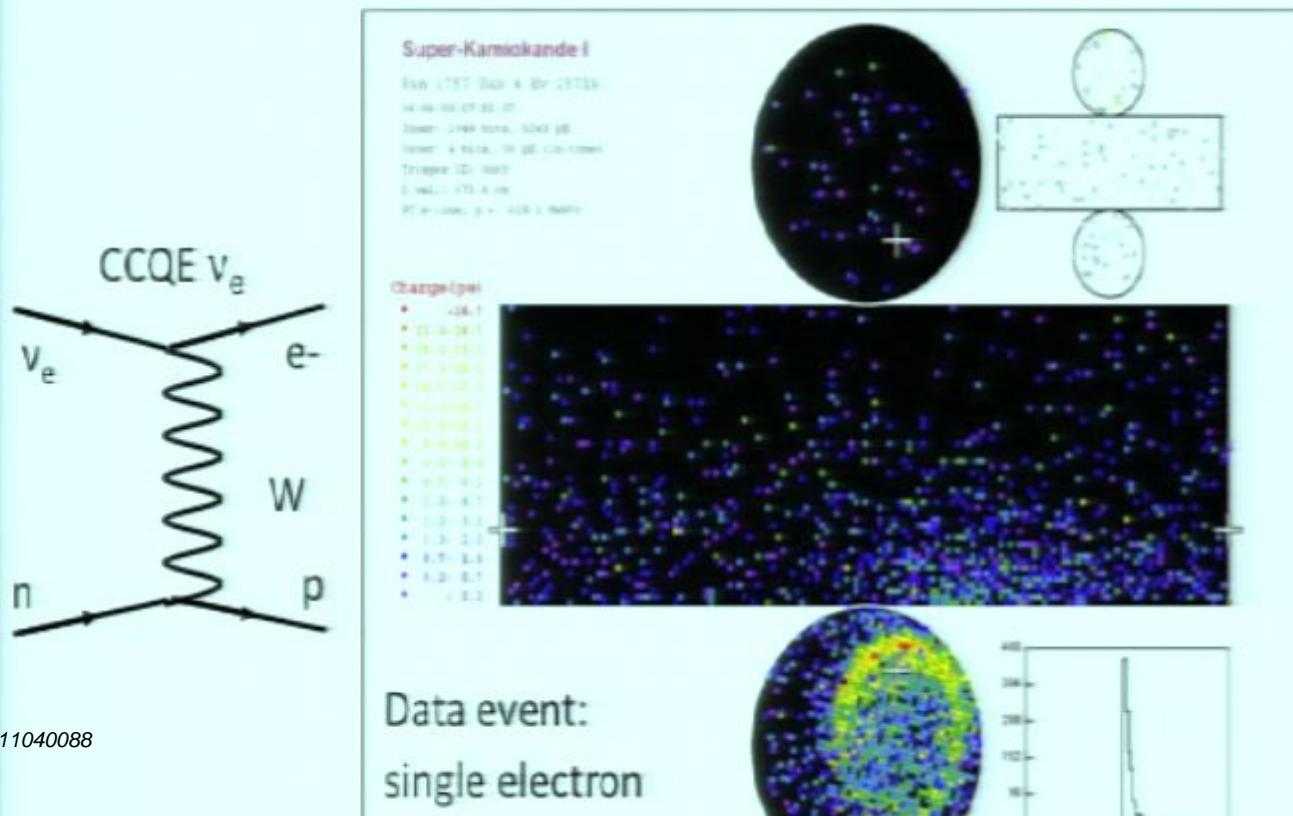
Electrons produce "fuzzy" rings, due to multiple scattering and showering

Electrons from CC ν_e interactions

Electrons from μ decay from CC μ (deadtime-less DAQ)

Neutral pions produce two electron-like rings from decay photons

If one ring is not reconstructed, mimics CCQE ν_e signal



NC π^0 background rejection

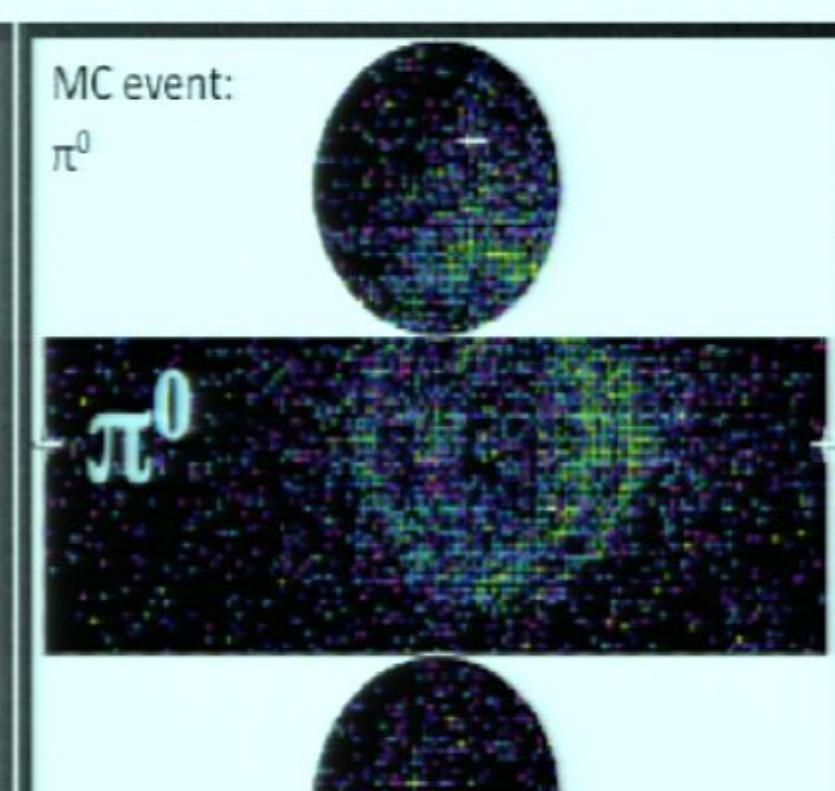
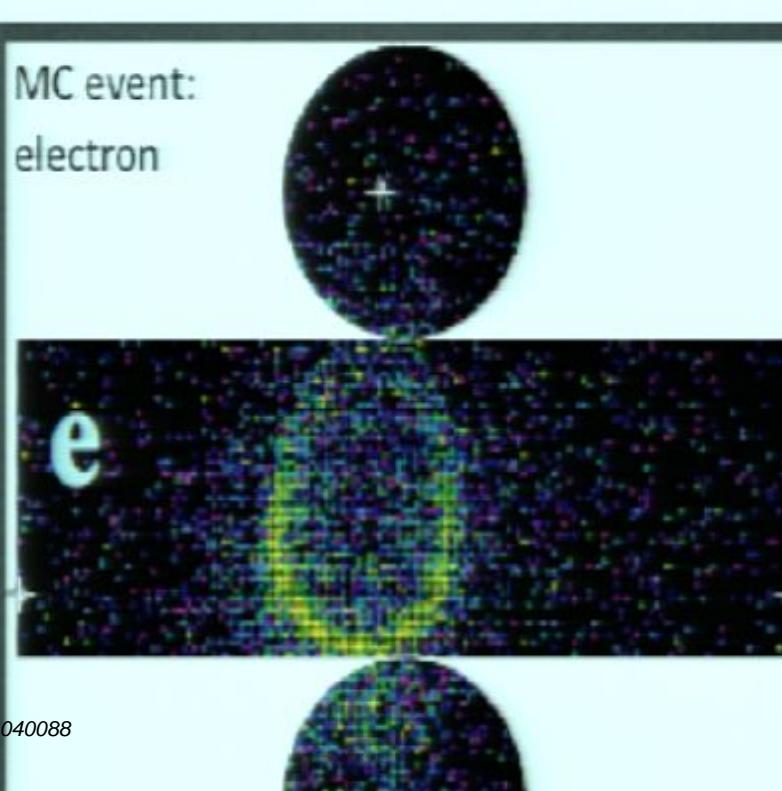
Simulate expected energy distribution for a faint second ring

Scan over possible second ring directions and energies

Select best match for observed light pattern

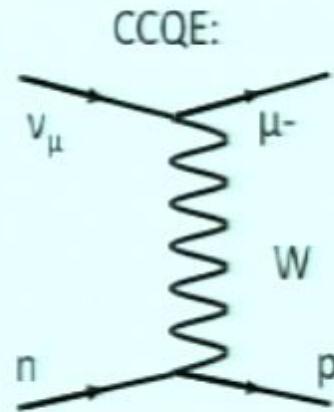
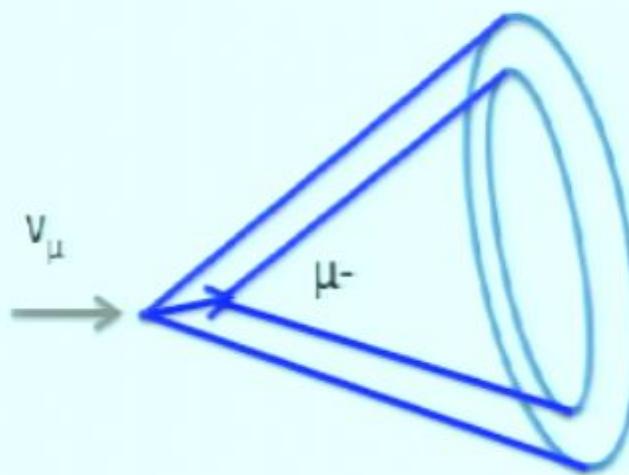
Cut on resulting invariant mass of two rings ($m(\pi^0)$) or $L_{2\gamma}/L_e$

Events with two true rings will have an invariant mass consistent with π^0



Determining neutrino energy

Energy of the outgoing lepton is determined from charge collected by the PMTs



For CCQE events, reconstruct neutrino energy from just the lepton kinematics provided:

- ✓ The neutrino direction is known
- ✓ The recoiling proton mass is known
- ✗ The target nucleon is at rest

$$E_{\nu}^{QE} = \frac{2M'_n E_\mu - [M'^2_n + m_\mu^2 - M_p^2]}{2[M'_n - E_\mu + p_\mu \cos\theta_\mu]}$$

Complications:

NC π^0 background rejection

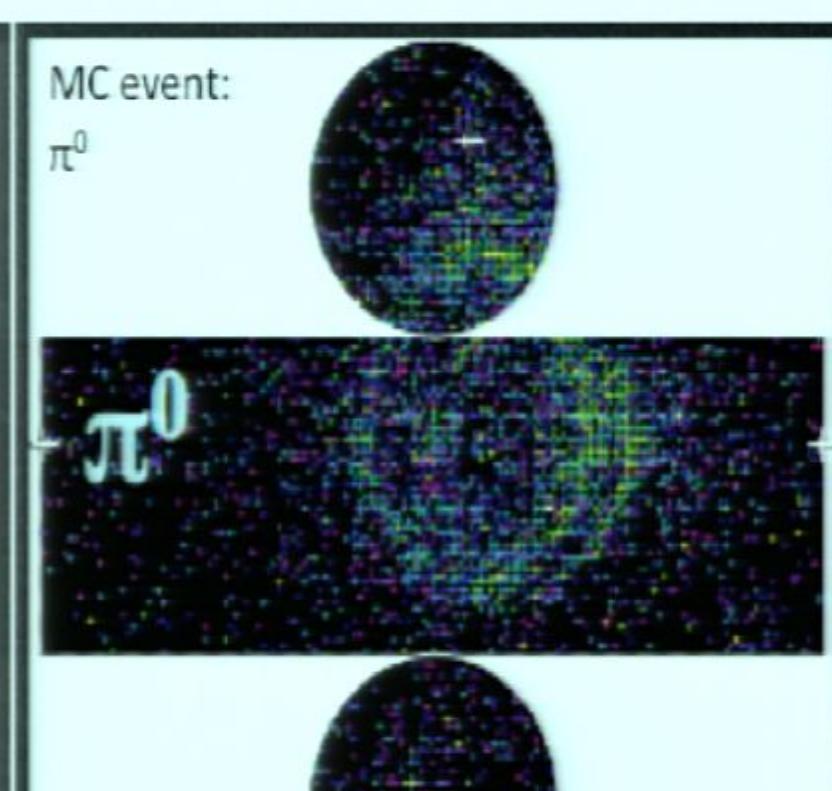
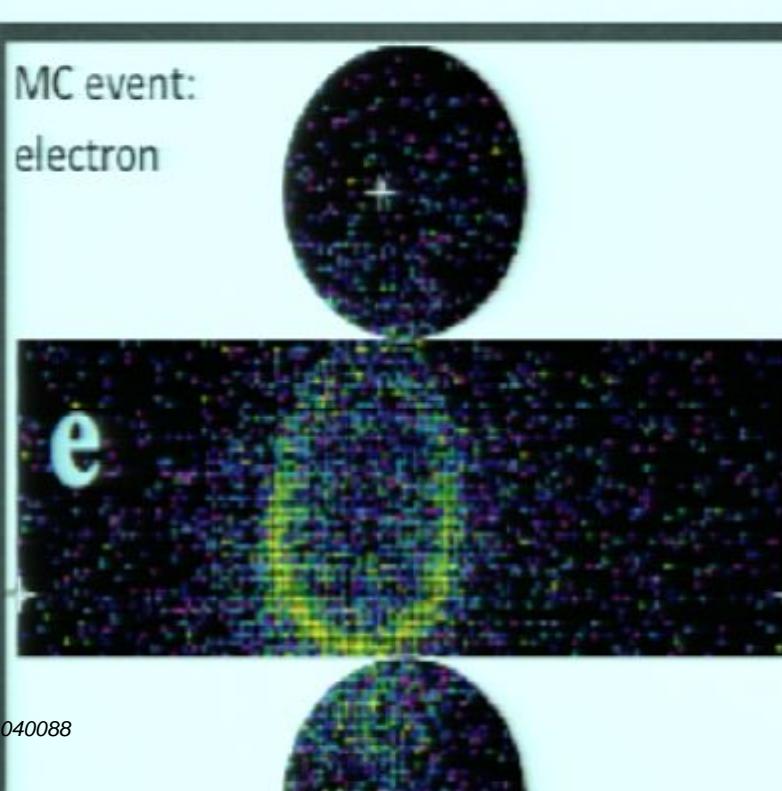
Simulate expected energy distribution for a faint second ring

Scan over possible second ring directions and energies

Select best match for observed light pattern

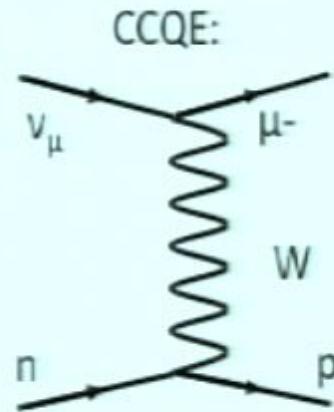
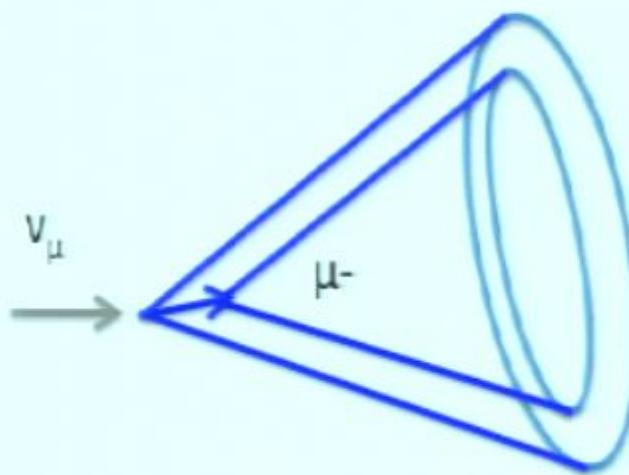
Cut on resulting invariant mass of two rings ($m(\pi^0)$) or $L_{2\gamma}/L_e$

Events with two true rings will have an invariant mass consistent with π^0



Determining neutrino energy

Energy of the outgoing lepton is determined from charge collected by the PMTs



For CCQE events, reconstruct neutrino energy from just the lepton kinematics provided:

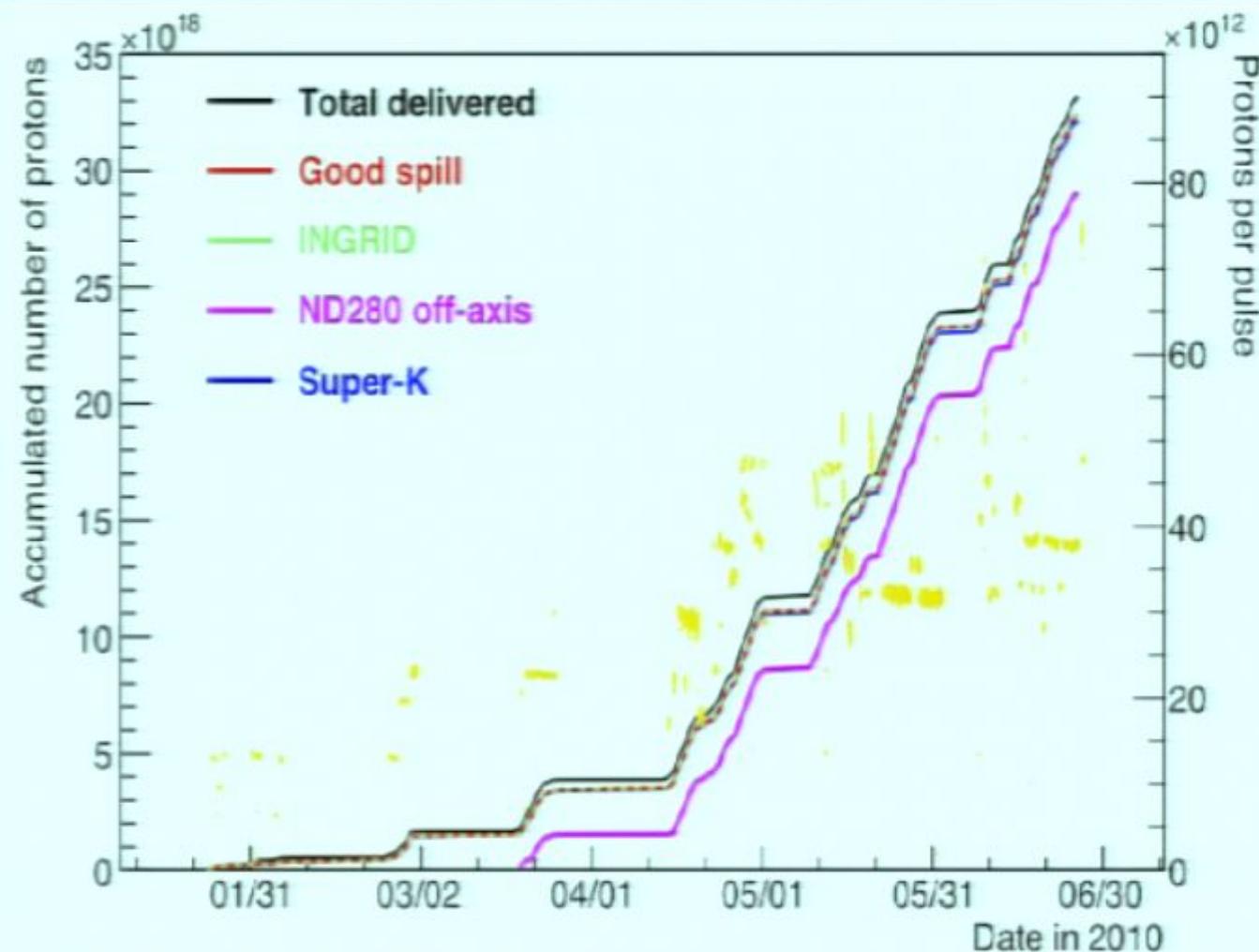
- ✓ The neutrino direction is known
- ✓ The recoiling proton mass is known
- ✗ The target nucleon is at rest

$$E_{\nu}^{QE} = \frac{2M'_n E_{\mu} - [M'^2_n + m_{\mu}^2 - M_p^2]}{2[M'_n - E_{\mu} + p_{\mu} \cos\theta_{\mu}]}$$

Complications:

Analysis of first dataset

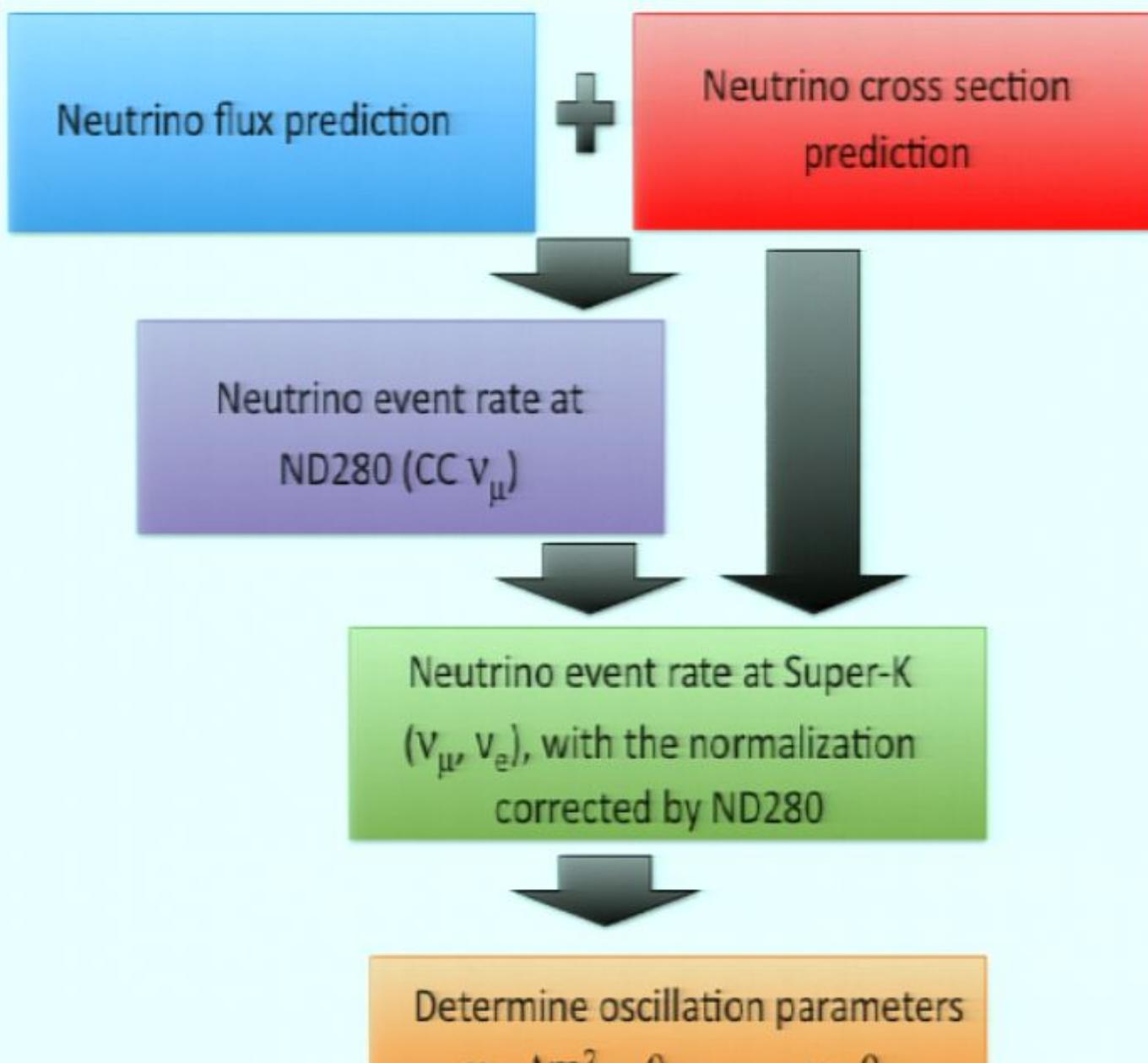
Run 1 dataset



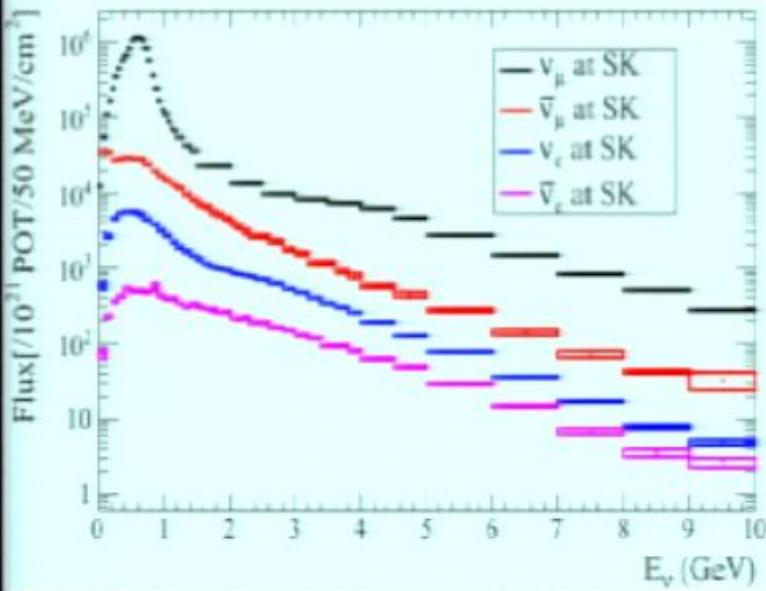
First run between Jan – Jun 2010:

6 bunches / spill / 3.54s

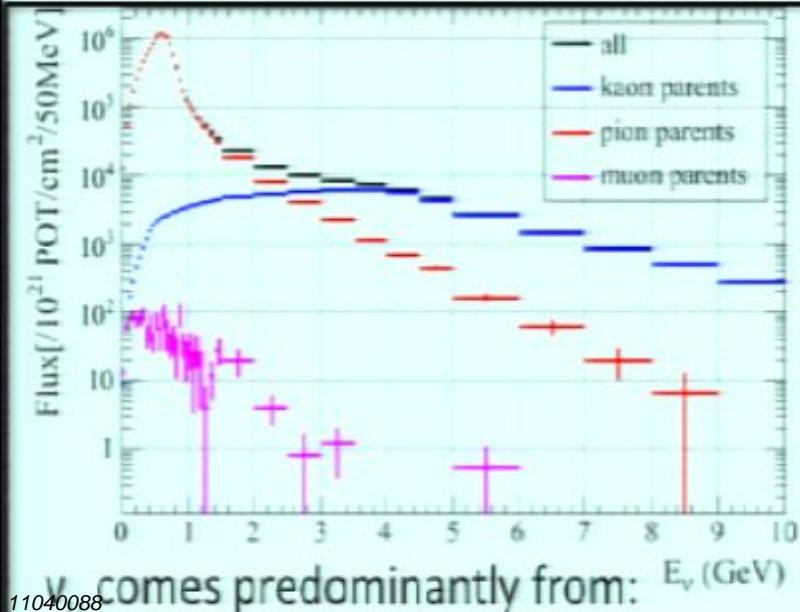
Analysis strategy



Neutrino flux prediction

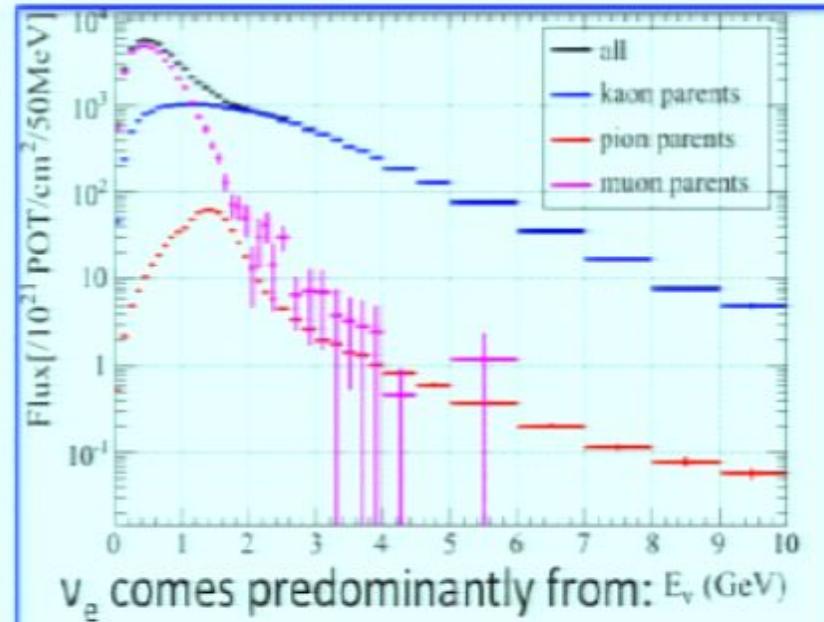


Majority of neutrino flux at Super-K is ν_μ
Also ν̄_μ (~6%) and
ν_e components (~1%)



11040088

π⁺ → ν_e decay



π → μ → ν_e decay chain (E_v < 1.5 GeV)

Neutrino flux uncertainties

Source	Uncertainty	Change at Peak	Max Change(<3GeV)
Pion Multiplicity	20%	16%	22%
Kaon Multiplicity	20-25%	1%	20%
Prod. Cross Sections	10-50%	7%	8%
Proton Beam	0.5mm, 0.3mrad	3%	9%
ν_μ Beam Direction	0.44mrad	1%	8%
Target Alignment	1.3mrad	<1%	1%
Horn Alignment	1mm	1%	3%
Horn Current	5kA	2%	2%
Horn Field Asymmetry	1.25%	0.5%	1%

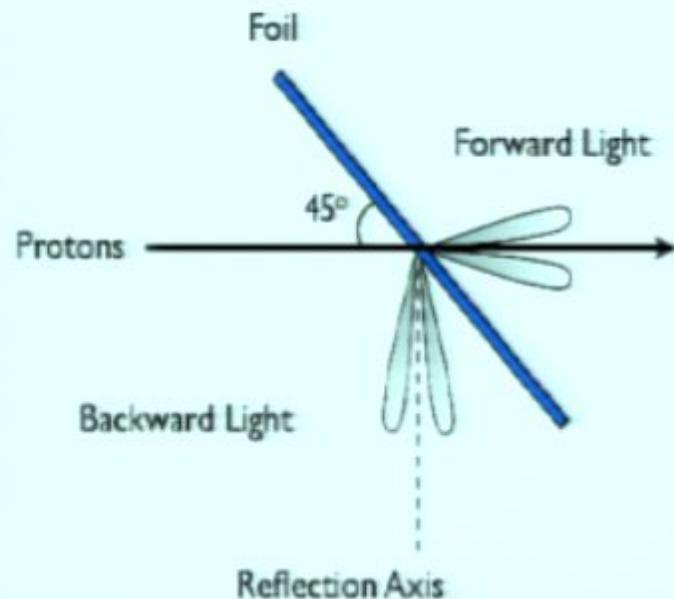
Constrained by in-situ measurements:

e.g. beamline monitors for the proton beam shape

Constrained by external experiments or data:

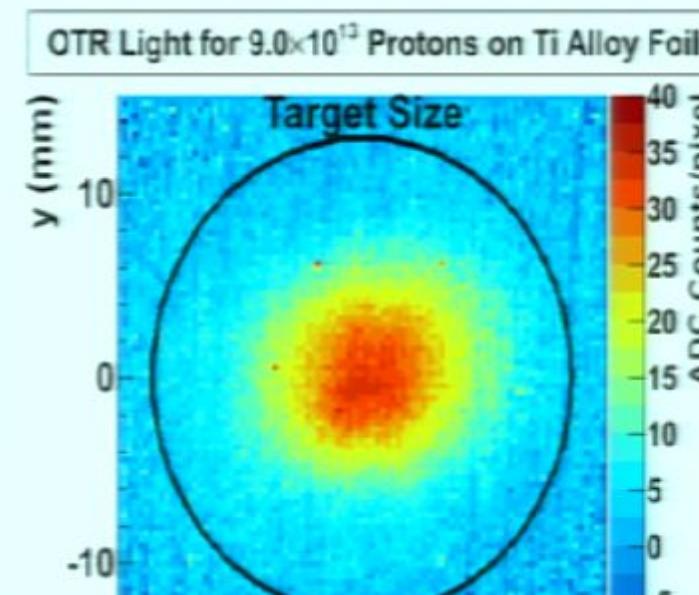
Proton beam monitoring

Multiple beam monitors measure the proton beam on the way to the neutrino target



Optical Transition Radiation is produced by the protons as they pass through a thin Ti foil in front of the neutrino target.

The light is emitted perpendicular to the beam direction, and is recorded with a 40mm camera. OTR light is used to determine the beam profile and position on the target.



π production from p+C

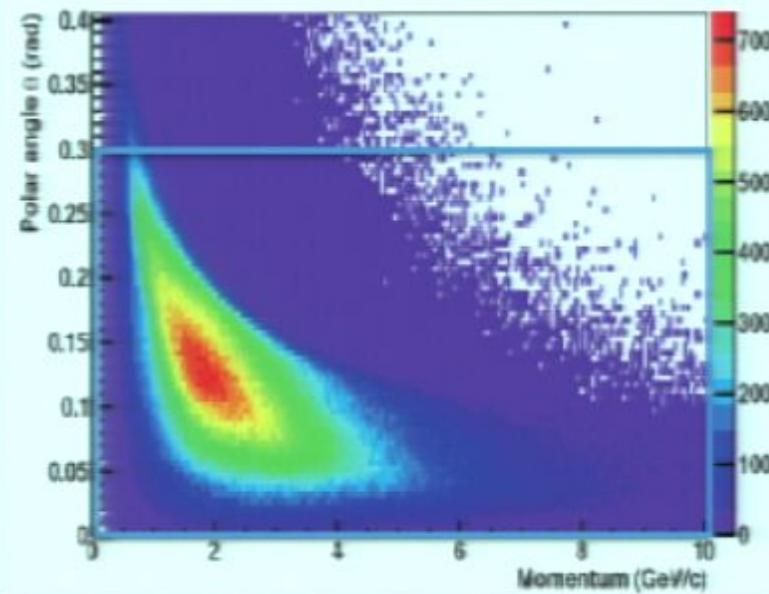
NA61/SHINE experiment at CERN

Designed to measure hadron production

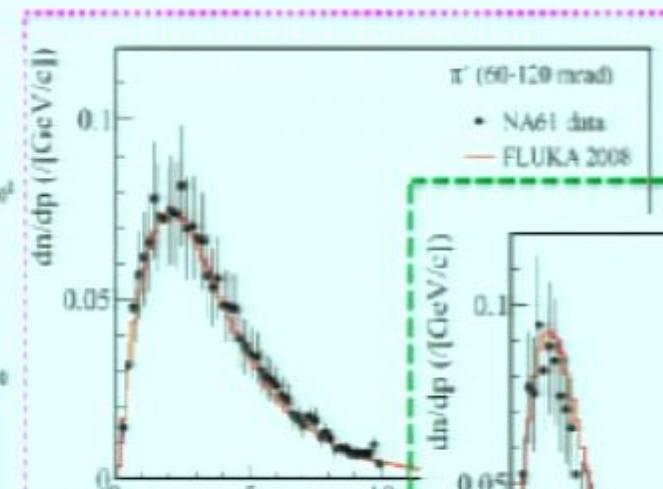
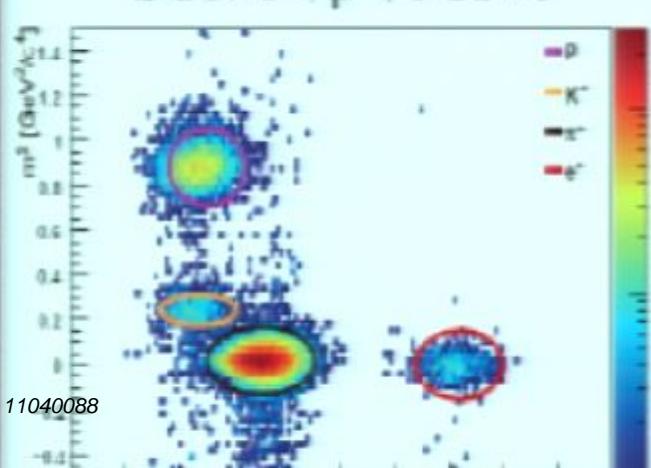
Used thin target (2cm) in 2007-2009 run,
T2K replica target (91cm, 1.9λ) in 2010 run

Use TOF and dE/dx to measure particle type; extract particle production cross section

i-p at production point of π^+ producing ν_μ @ SK



$2 \text{ GeV}/c < p < 3 \text{ GeV}/c$



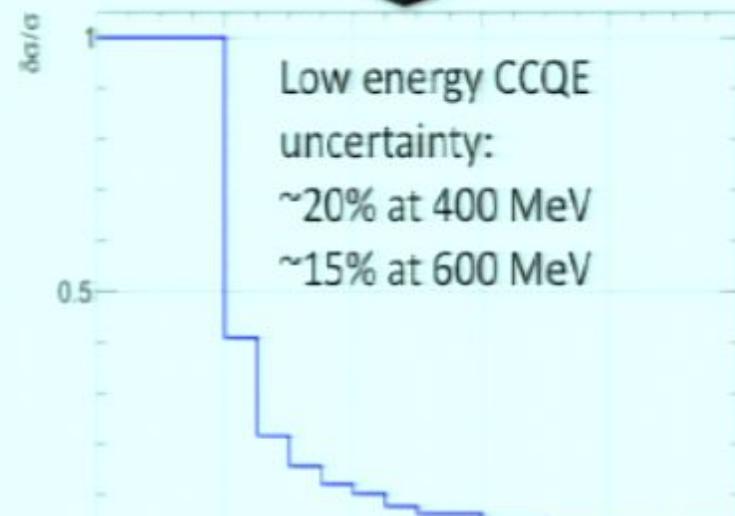
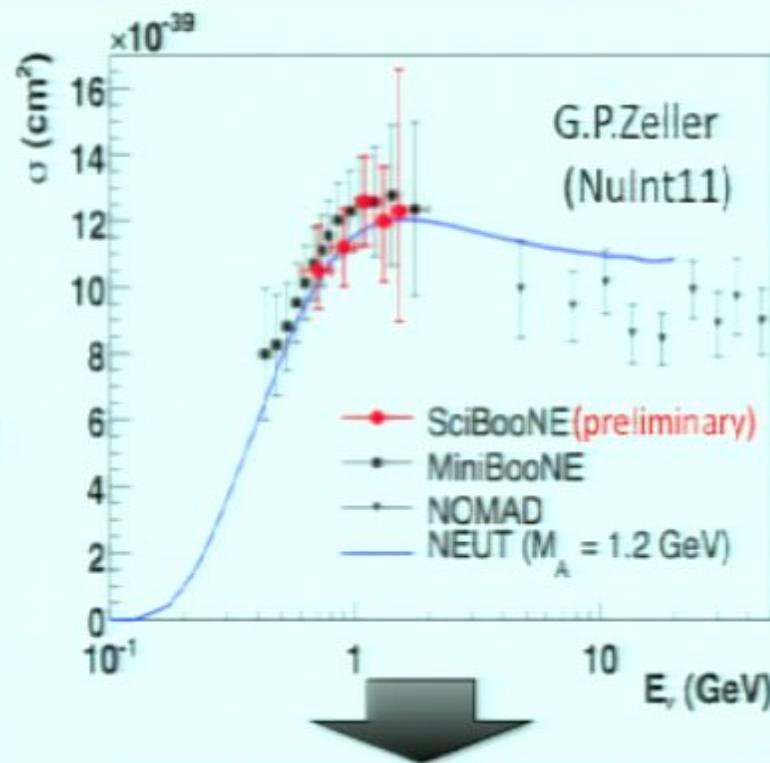
Neutrino cross section prediction, uncertainties

Cross section model set from external data

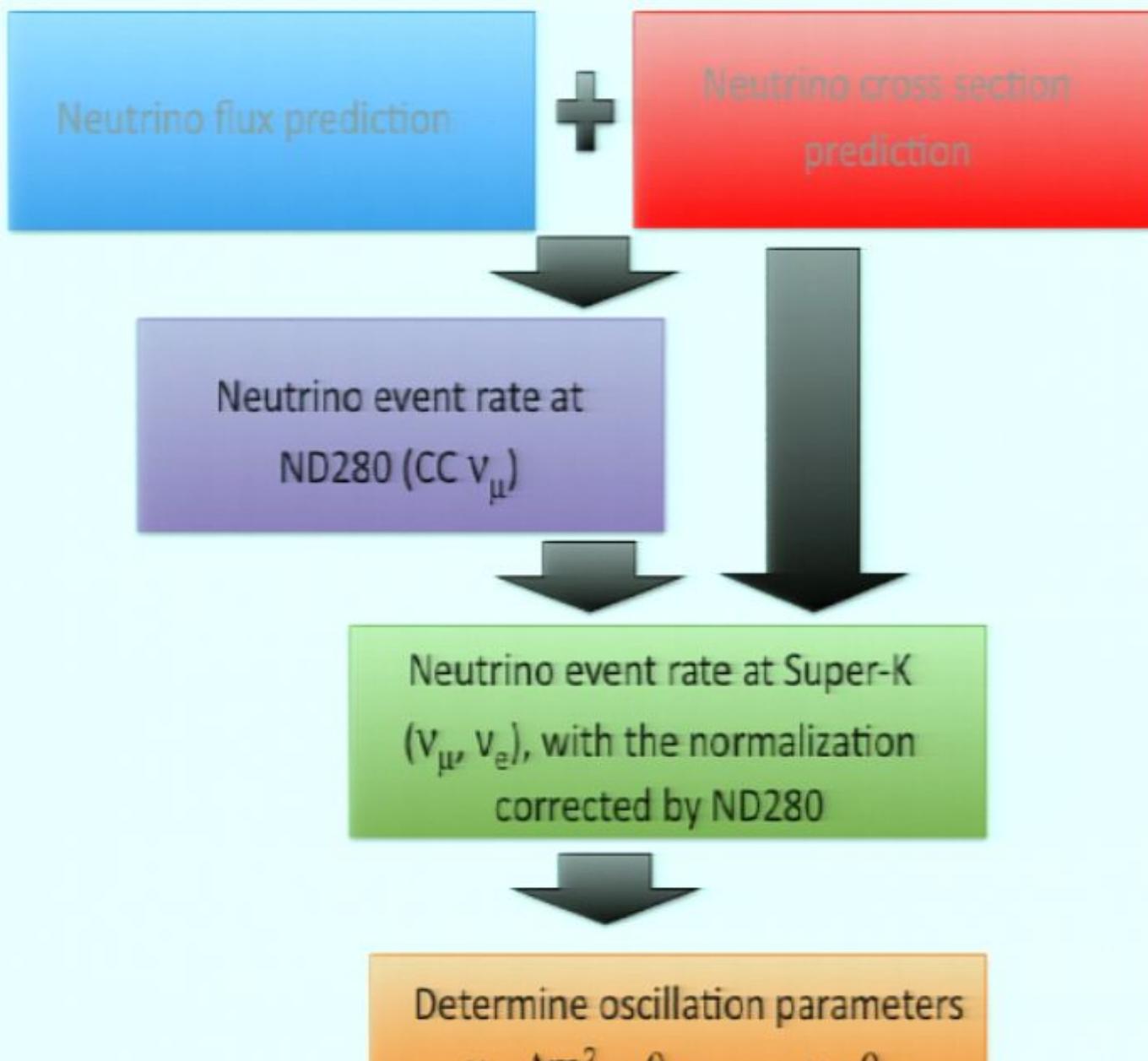
MiniBooNE, SciBooNE, K2K
Super-K atmospheric data

Uncertainties set from external data
constraints (includes variations to underlying
parameters and difference between models)

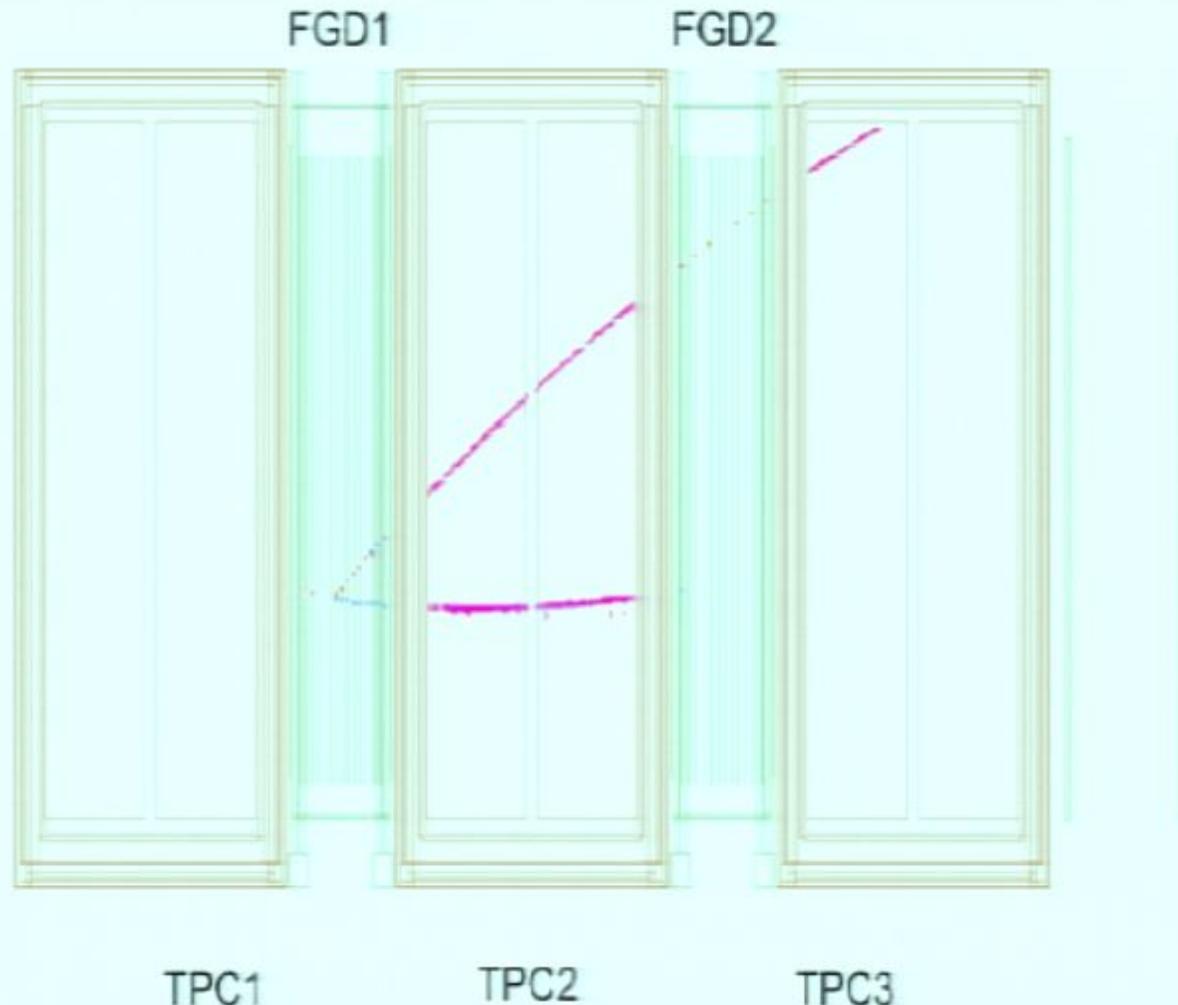
Source	Uncertainty < 2 GeV	Uncertainty > 2 GeV
CCQE	energy dependent	energy dependent
CC1 π	30%	20%
CC coherent	100%	100%
NC π^0	30%	25%
NC coherent/other	30%	30%



Analysis strategy

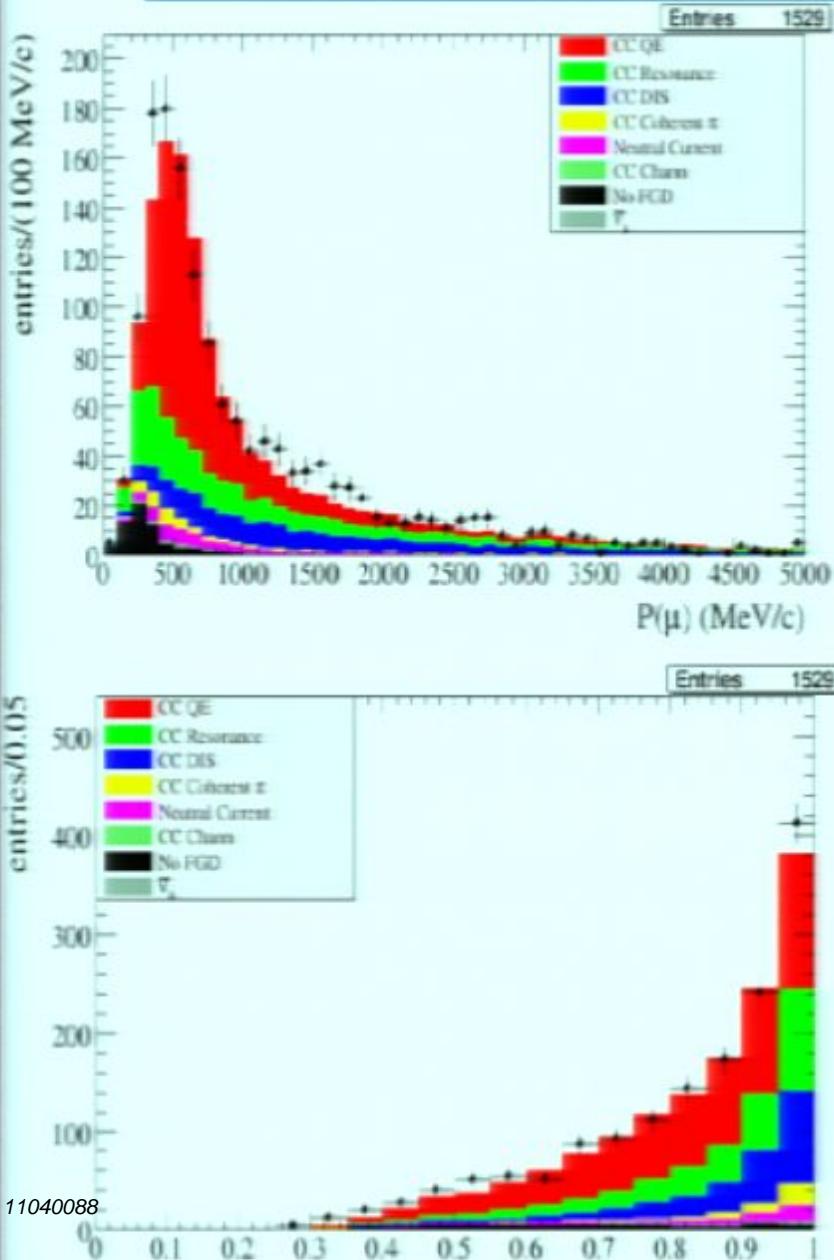


Basic CC selection in ND280



1. Select neutrino events: Use TPC1 as a veto (no tracks in TPC1)
2. Select events which originate in FGD1 or FGD2 fiducial volume
3. Use the highest momentum, negative TPC2 or TPC3 track

ND280 CC ν_μ sample



Reconstructed momentum and angle of the CC ν_μ candidates after selection

CC ν_μ purity: 91%

CCQE purity: 49%

No tuning to flux or cross section applied

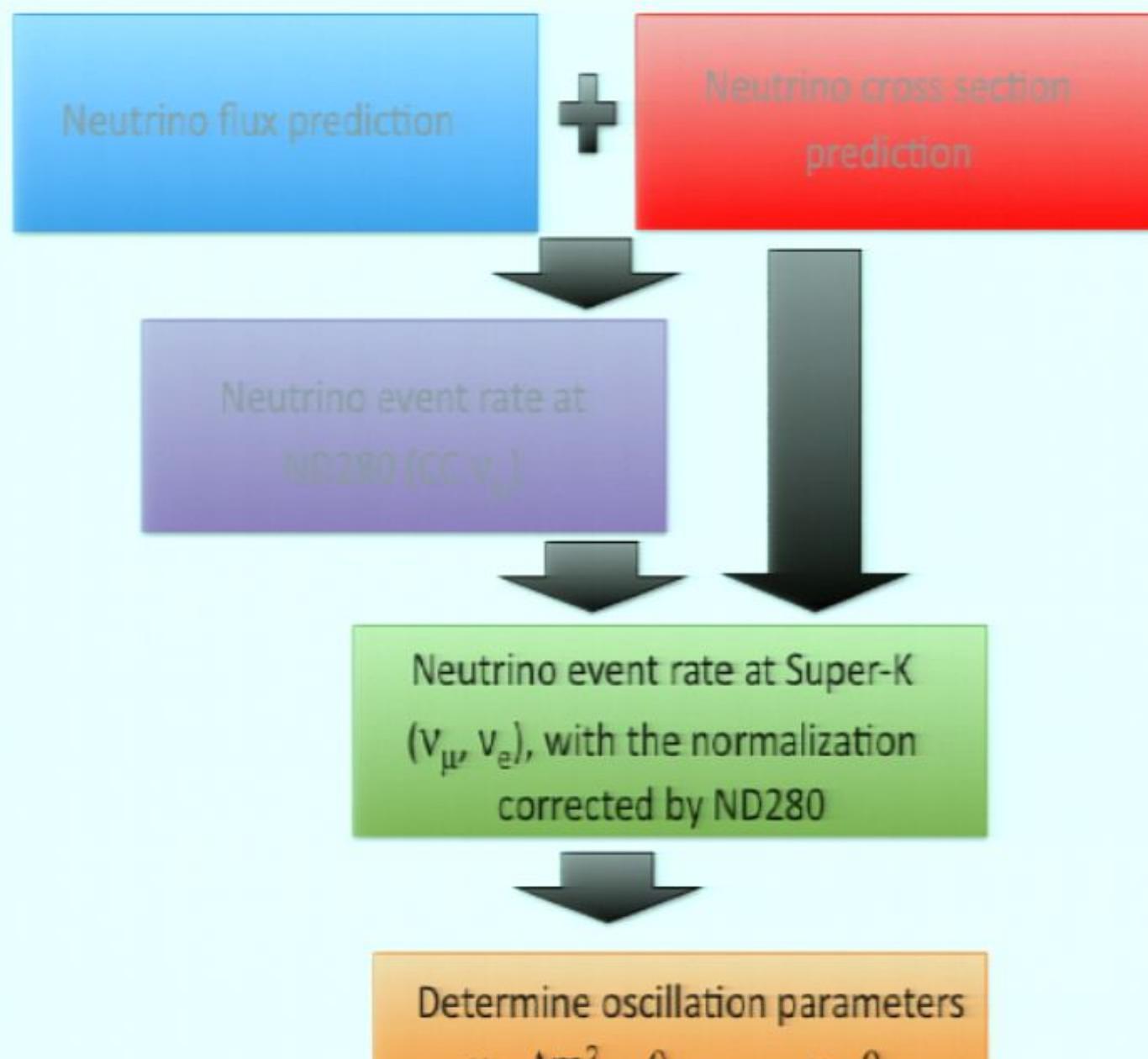
$$R(\text{data/MC}) = 1.061 \pm 0.028 \text{ (statistics)}$$
$$+0.044 \quad -0.038 \text{ (det systematics)}$$
$$\pm 0.039 \text{ (flux, xsec model)}$$

Dominant uncertainties:

TPC PID pull: 3.0%

TPC-FGD matching: 2.1%

Analysis strategy



Far detector neutrino event selection

Basic neutrino selection (precuts)

Event time within beam window

No activity in the veto

Reconstructed vertex $>2\text{m}$ from wall

Single reconstructed ring

ν_μ selection

Visible energy $> 30 \text{ MeV}$

μ -like ring

$E\mu > 200 \text{ MeV}$

decay electrons < 2

ν_e selection

Visible energy $> 100 \text{ MeV}$

e-like ring

No decay electron

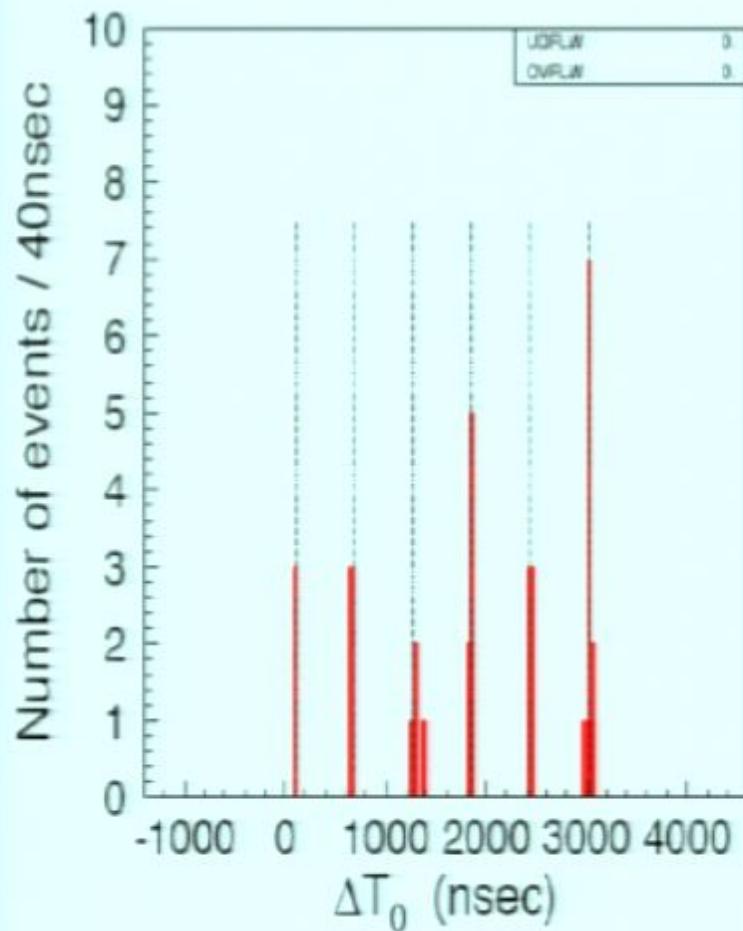
Invariant mass $< 105 \text{ MeV}/c^2$

$E\nu < 1250 \text{ MeV}$

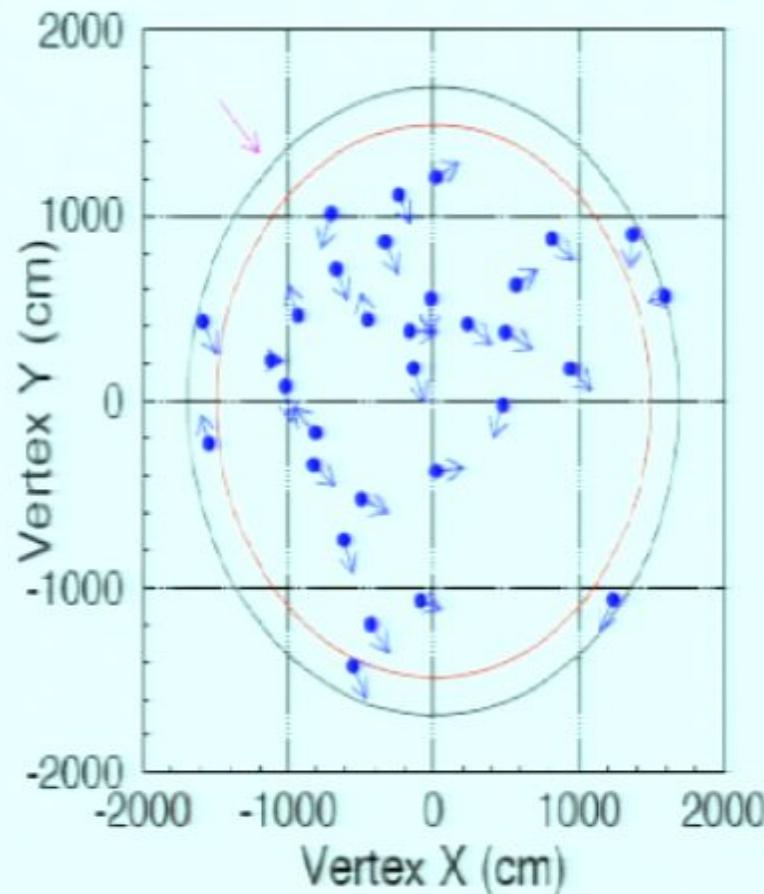
Cuts were set before looking at data from T2K

Feasible because of many years of experience with the Super-K detector
(atmospheric neutrino analyses)

Basic neutrino event selection

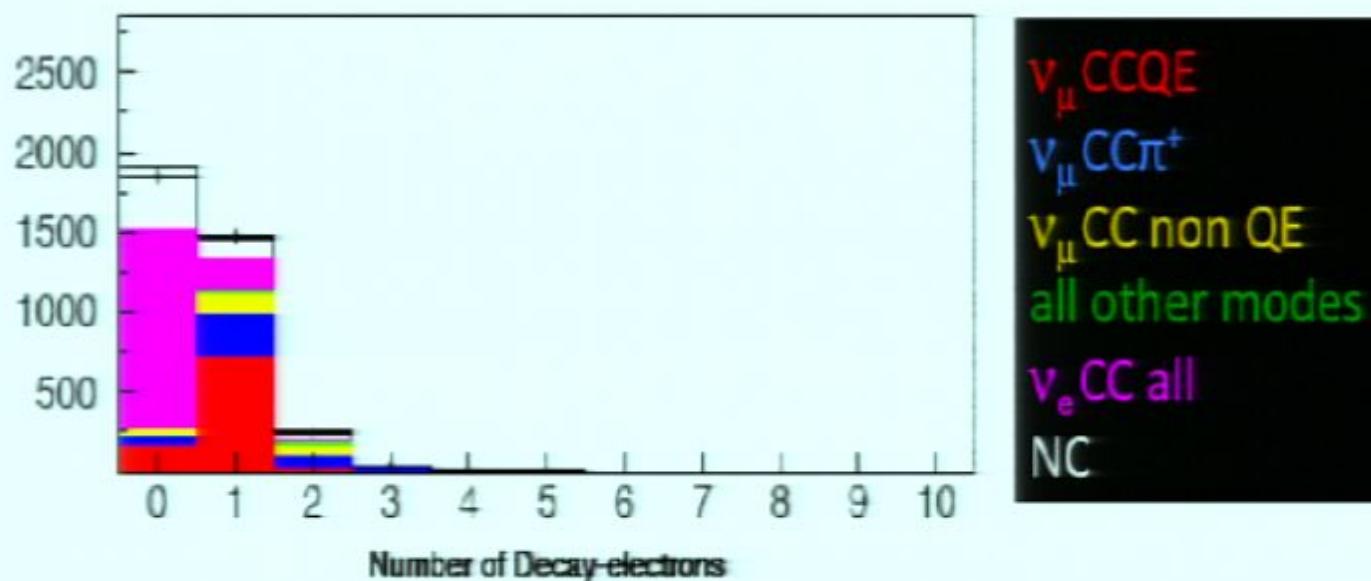


Events are consistent with beam timing,
established with GPS



Distribution within tank and
reconstructed event direction

Far detector uncertainties



ν_μ dominant uncertainties: ring counting, muon PID

ν_e dominant uncertainties: ring counting, electron PID, invariant mass selection

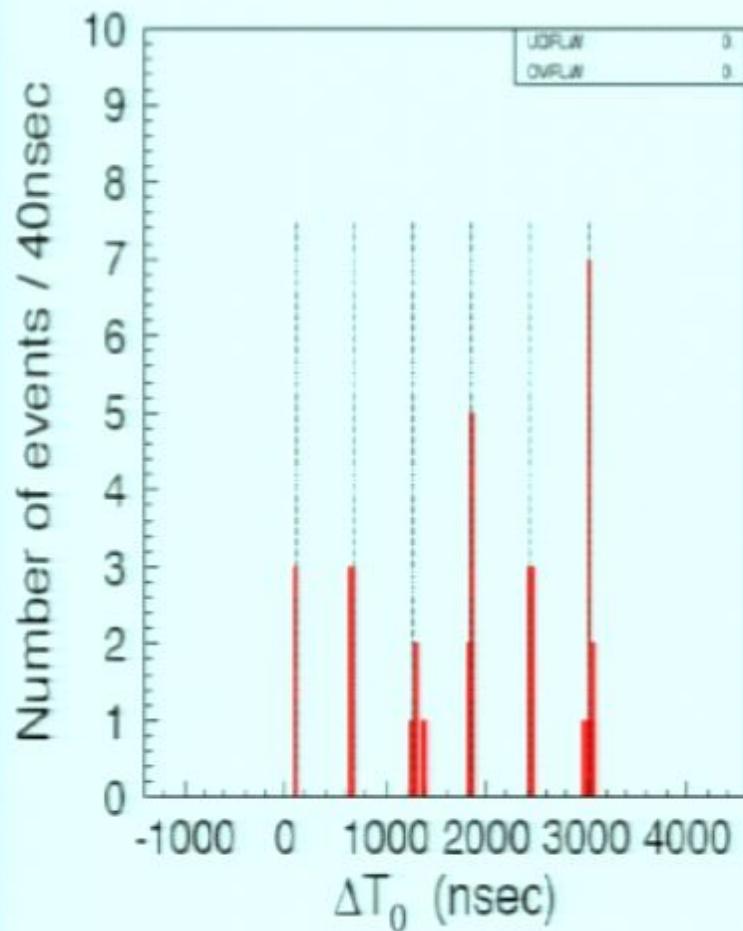
Use atmospheric ν sample, select unbiased sample using decay electron tagging

0 decay electrons: NC event, CCQE ν_e events

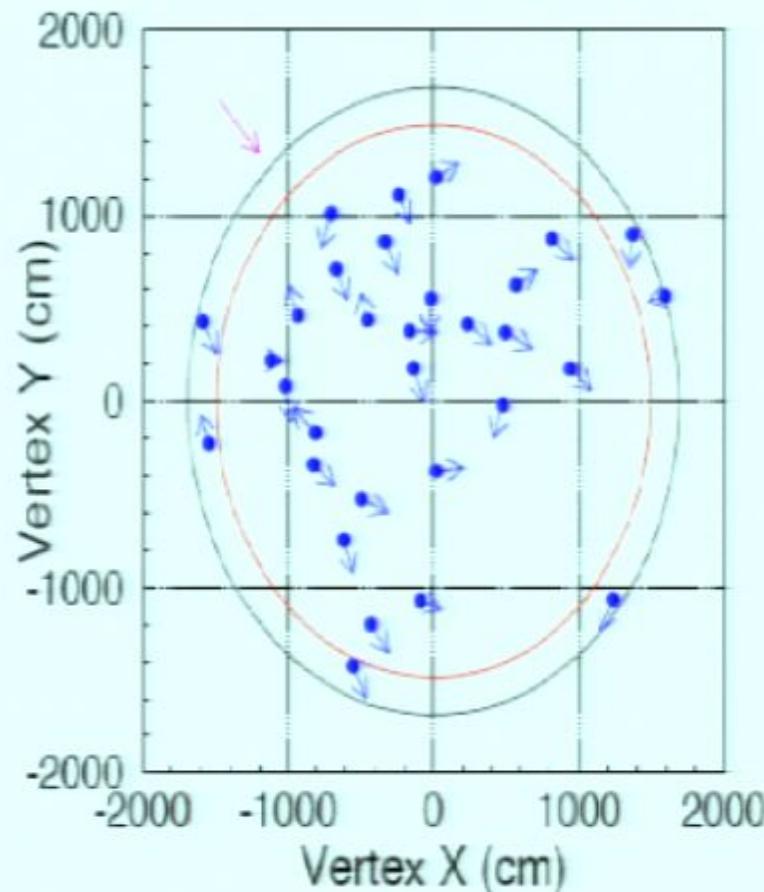
1 decay electrons: CCQE ν_μ

2 decay electrons: CC non QE (pion, multipion events)

Basic neutrino event selection

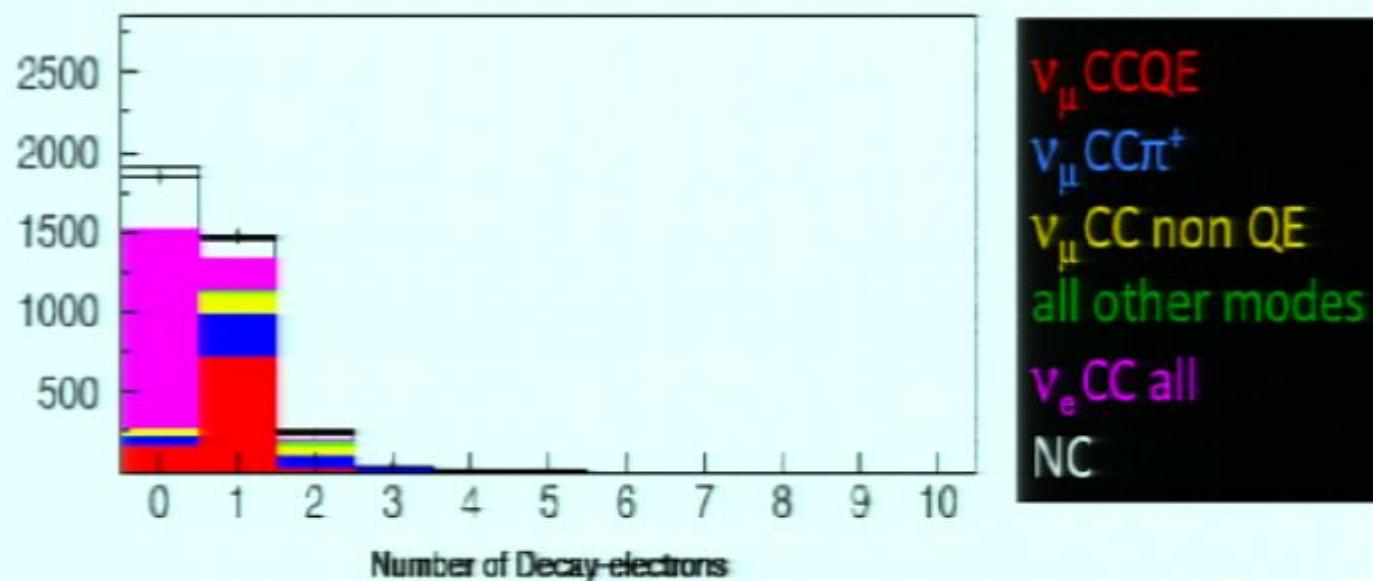


Events are consistent with beam timing,
established with GPS



Distribution within tank and
reconstructed event direction

Far detector uncertainties



ν_μ dominant uncertainties: ring counting, muon PID

ν_e dominant uncertainties: ring counting, electron PID, invariant mass selection

Use atmospheric ν sample, select unbiased sample using decay electron tagging

0 decay electrons: NC event, CCQE ν_e events

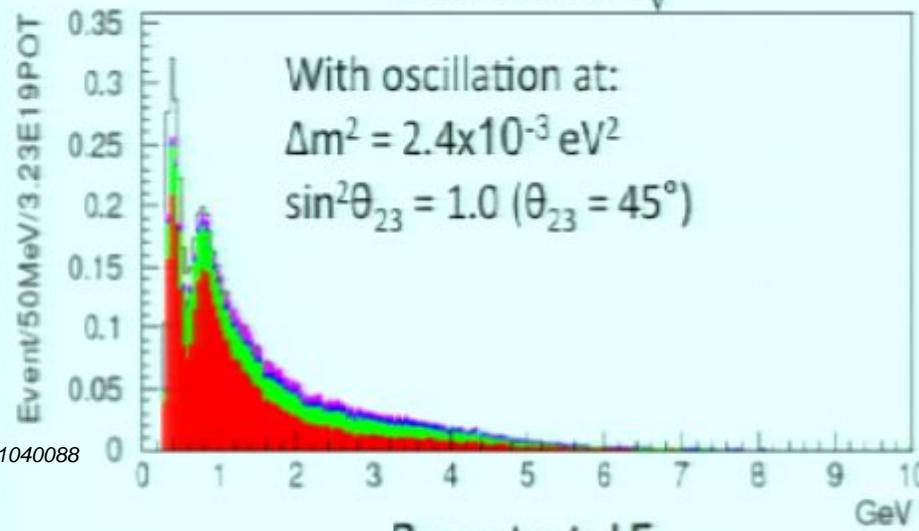
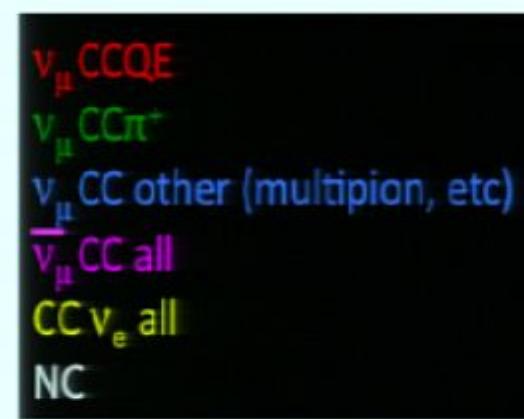
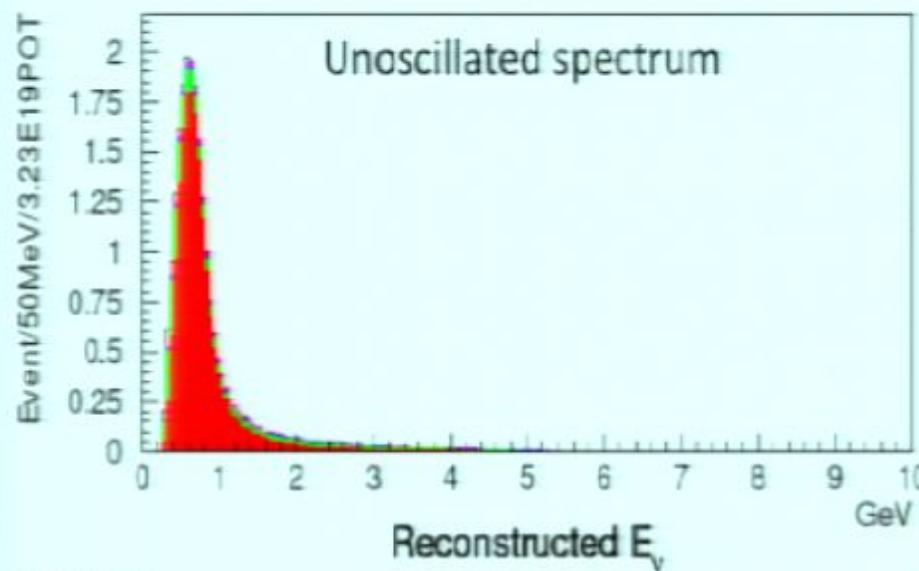
1 decay electrons: CCQE ν_μ

2 decay electrons: CC non QE (pion, multipion events)

ν_μ disappearance results

ν_μ disappearance at T2K

For a fixed baseline ($L=295\text{ km}$) neutrinos of energy E_1 and E_2 will have different probabilities for oscillation



$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Extract Δm^2 , $\sin^2 \theta_{23}$ from the change
in overall rate and spectrum
distortion observed

Backgrounds are events where pion
is unobserved (absorbed)

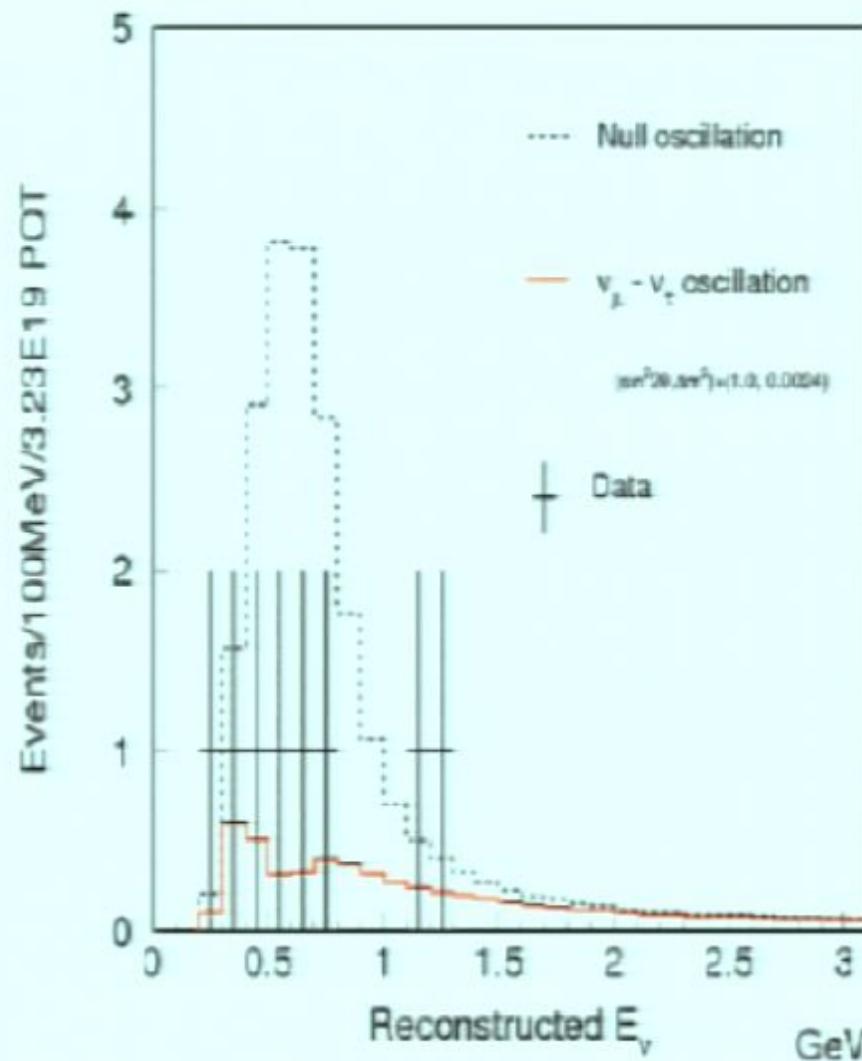
ν_μ disappearance results

8 events after selection

No-oscillation prediction:
 22.81 ± 3.19 (systematics)

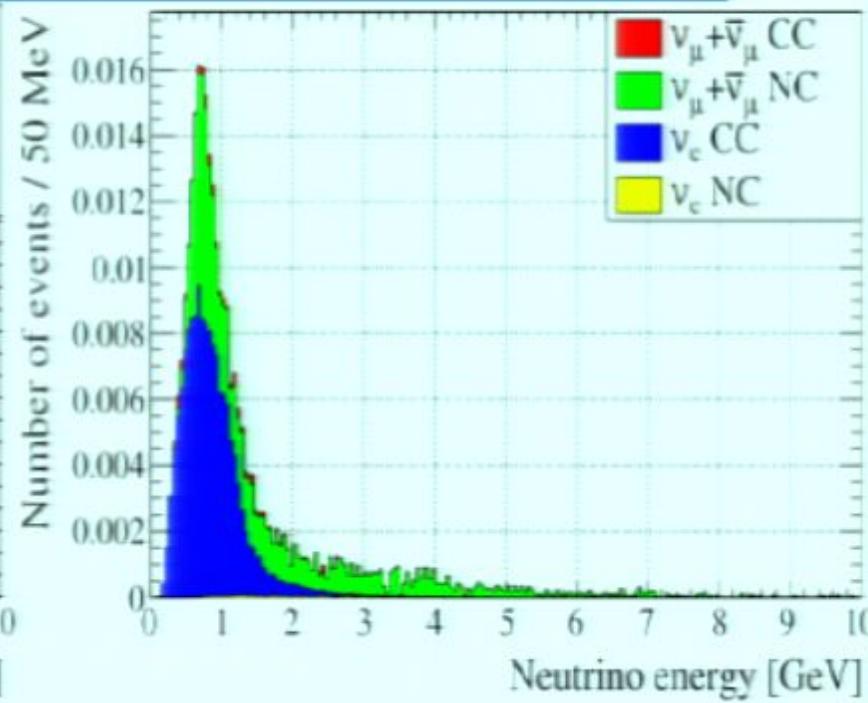
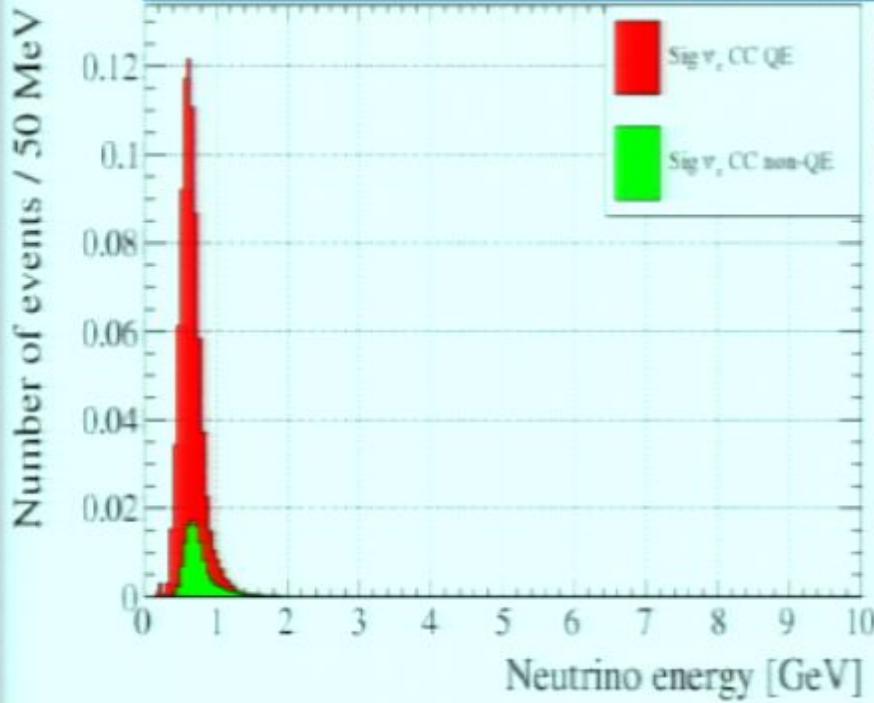
With oscillation:
 6.34 ± 1.04 (systematics)

at MINOS experimental best fit:
 $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 2\theta_{23} = 1.0$



v_e appearance results

ν_e appearance at T2K



Signal: CC interactions of ν_e from ν_μ to ν_e oscillations

Signal	# events
@ $\sin^2 2\theta_{13} = 0.1, \delta \text{cp} = 0$	0.9

Background	# events
intrinsic ν_e	0.16
ν_μ (70% NC π^0)	0.13
$\bar{\nu}_\mu$	0.01
total:	$0.30 \pm 0.07 \text{ (sys)}$

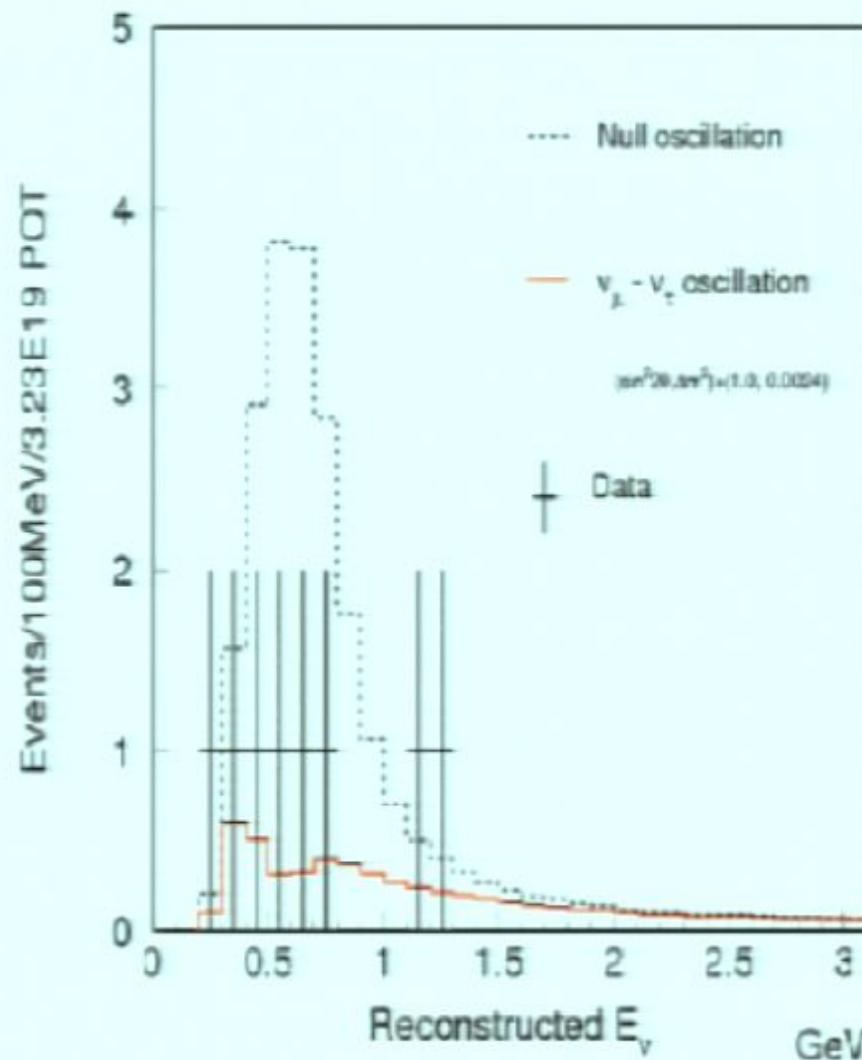
ν_μ disappearance results

8 events after selection

No-oscillation prediction:
 22.81 ± 3.19 (systematics)

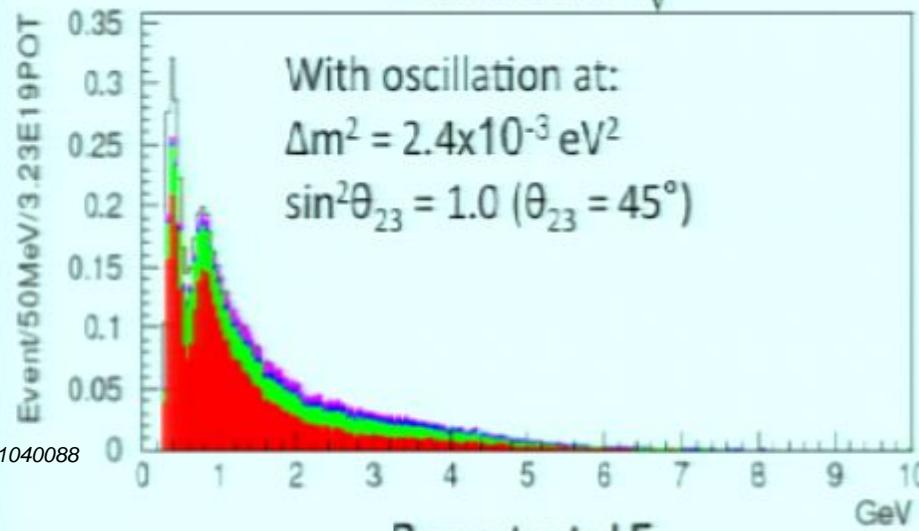
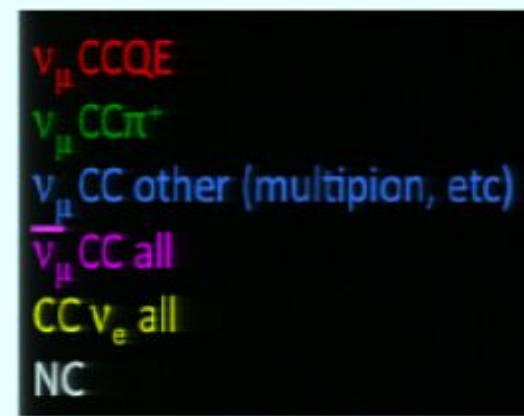
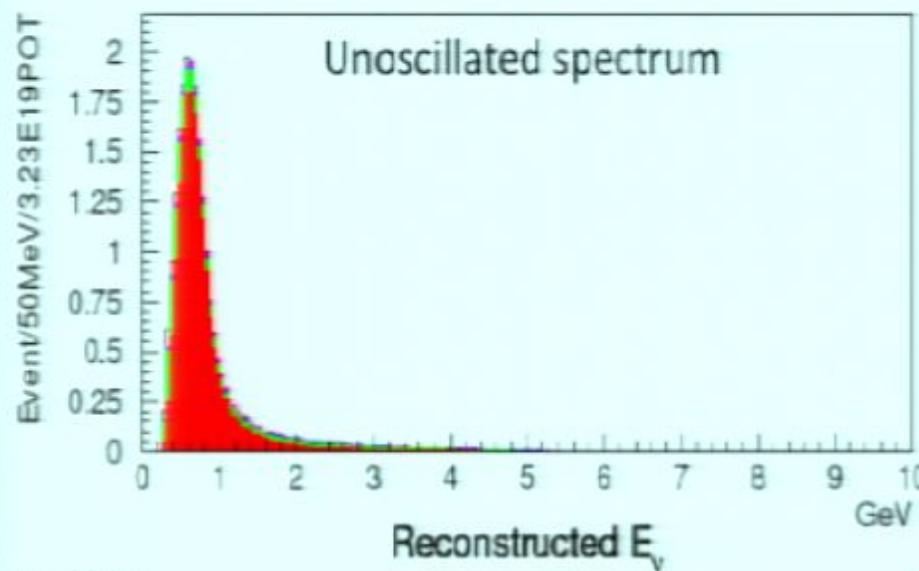
With oscillation:
 6.34 ± 1.04 (systematics)

at MINOS experimental best fit:
 $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 2\theta_{23} = 1.0$



ν_μ disappearance at T2K

For a fixed baseline ($L=295\text{ km}$) neutrinos of energy E_1 and E_2 will have different probabilities for oscillation



$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right)$$

Extract Δm^2 , $\sin^2 \theta_{23}$ from the change in overall rate and spectrum distortion observed

Backgrounds are events where pion is unobserved (absorbed)

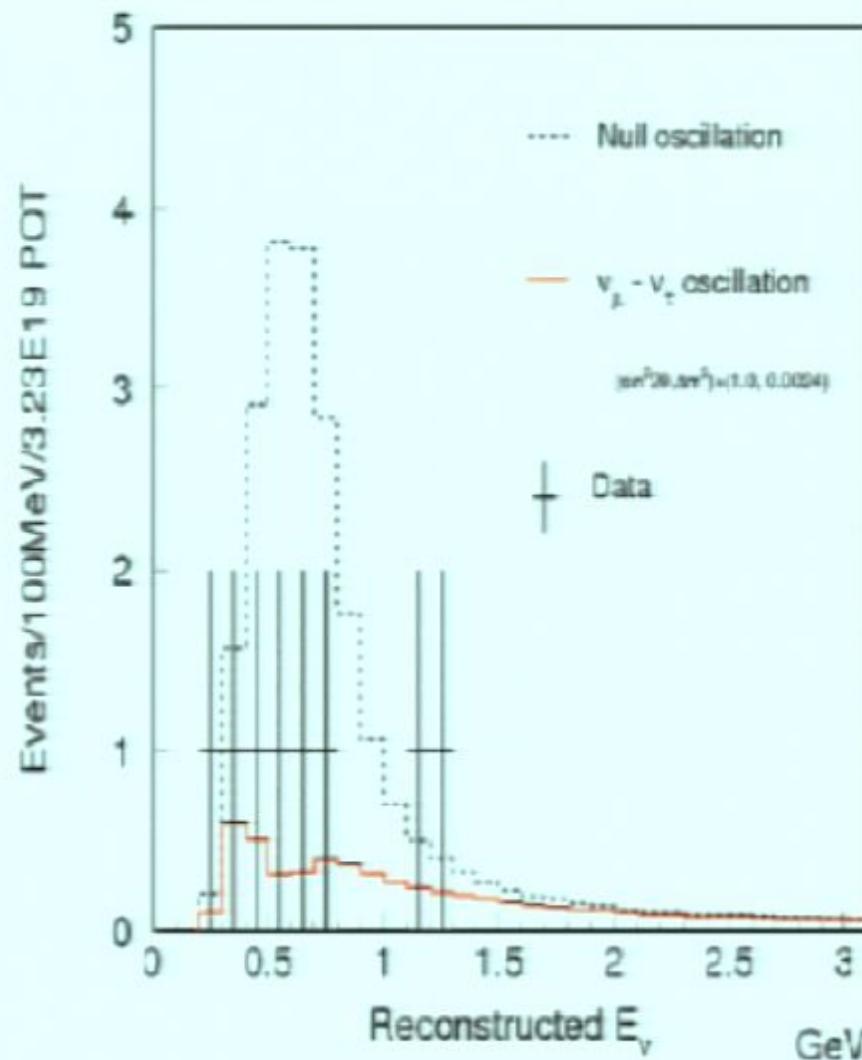
ν_μ disappearance results

8 events after selection

No-oscillation prediction:
 22.81 ± 3.19 (systematics)

With oscillation:
 6.34 ± 1.04 (systematics)

at MINOS experimental best fit:
 $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 2\theta_{23} = 1.0$



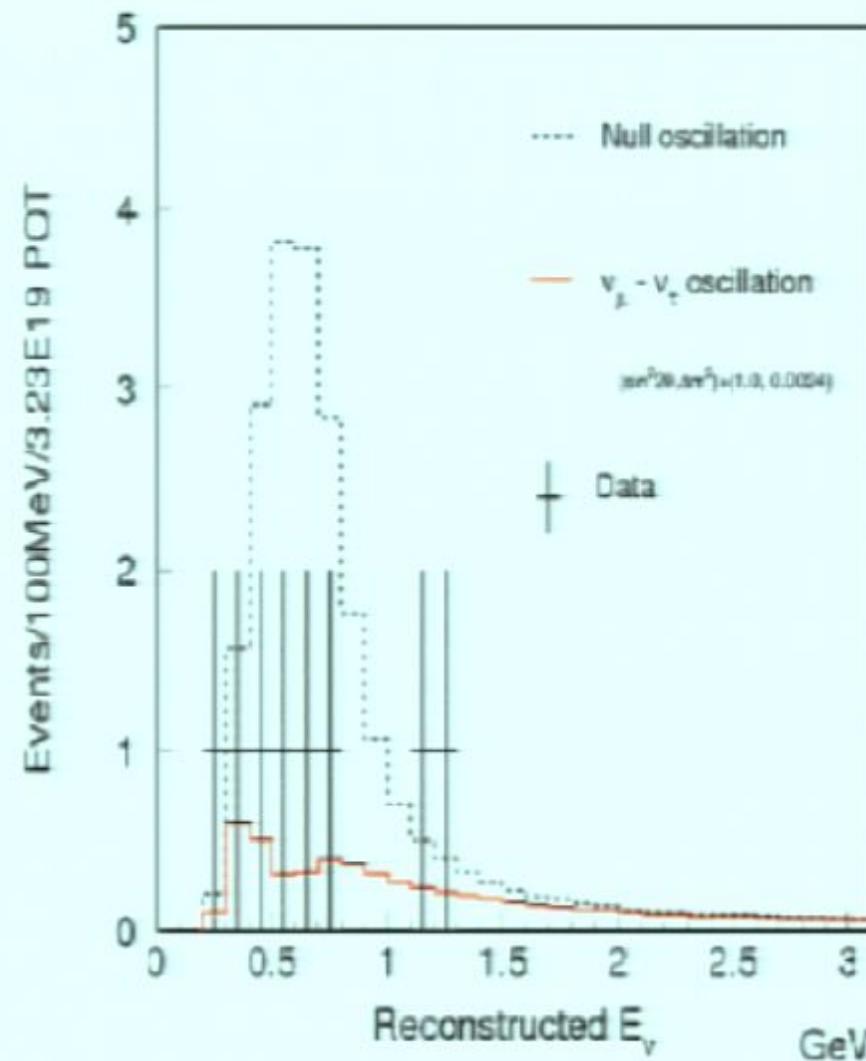
ν_μ disappearance results

8 events after selection

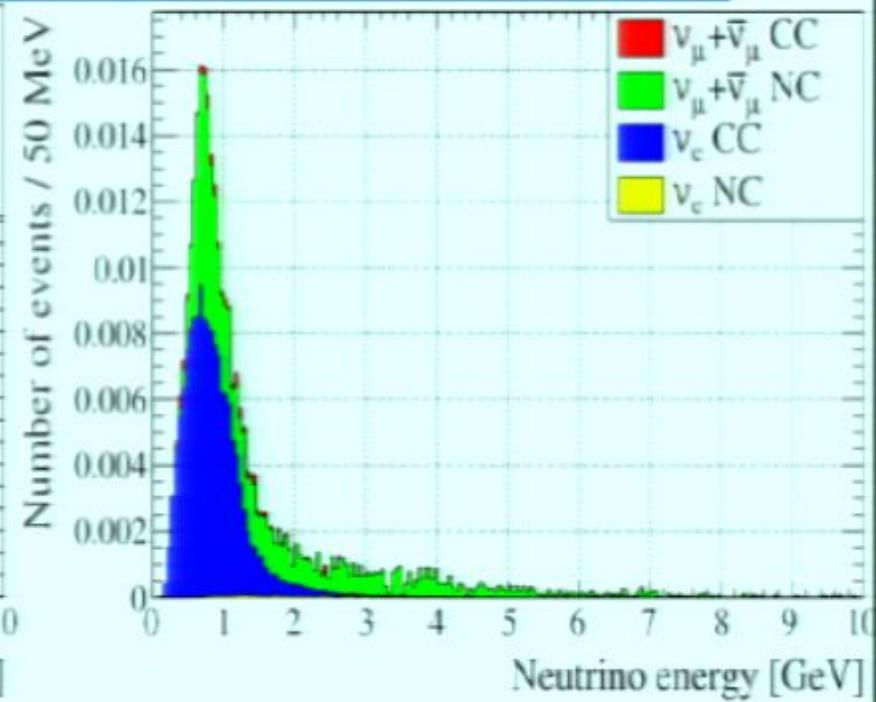
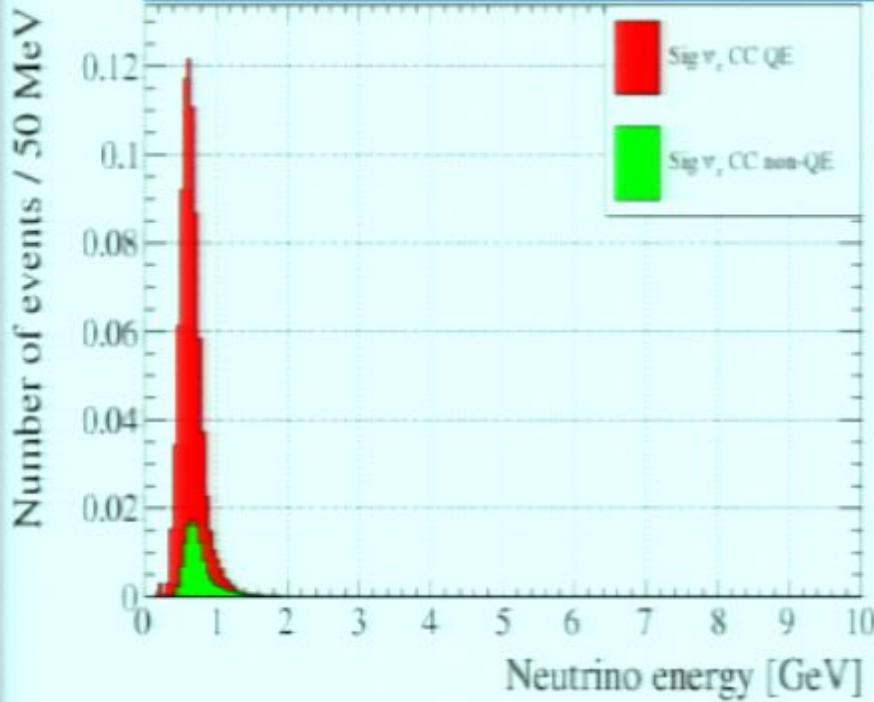
No-oscillation prediction:
 22.81 ± 3.19 (systematics)

With oscillation:
 6.34 ± 1.04 (systematics)

at MINOS experimental best fit:
 $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 2\theta_{23} = 1.0$



ν_e appearance at T2K



Signal: CC interactions of ν_e from ν_μ to ν_e oscillations

Signal	# events
@ $\sin^2 2\theta_{13} = 0.1, \delta \text{cp} = 0$	0.9

Background	# events
intrinsic ν_e	0.16
ν_μ (70% NC π^0)	0.13
$\bar{\nu}_\mu$	0.01
total:	$0.30 \pm 0.07 \text{ (sys)}$

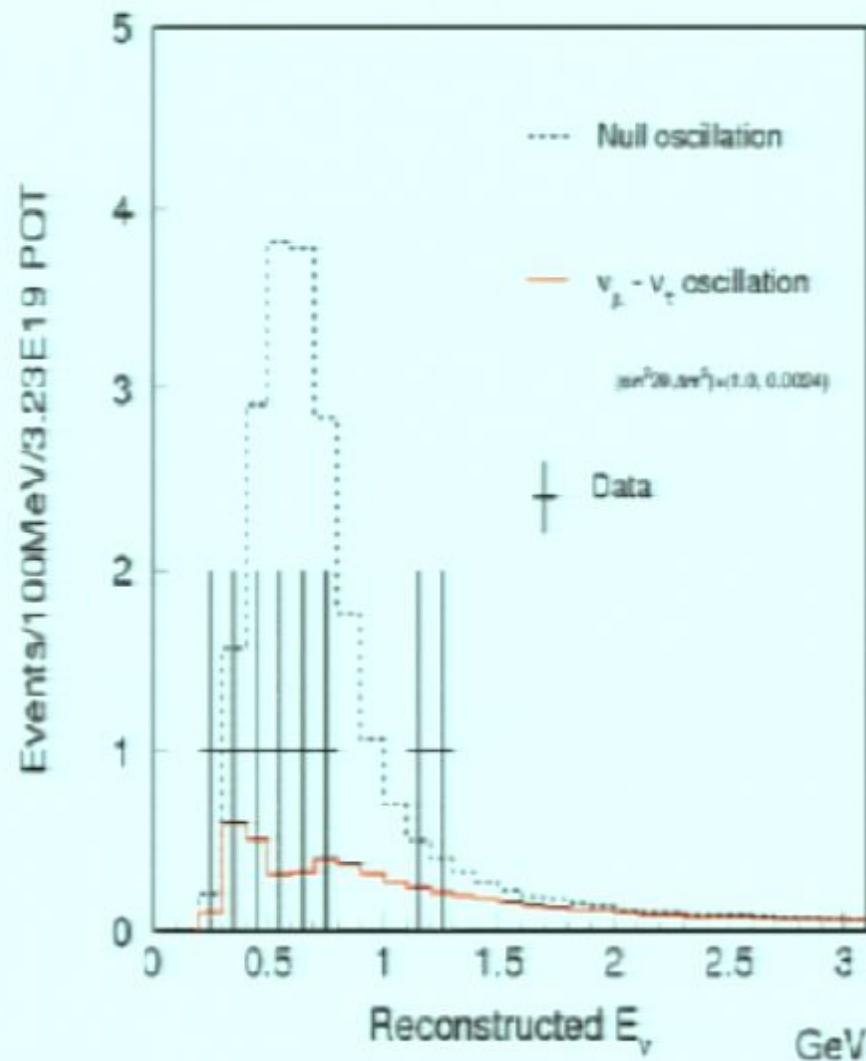
ν_μ disappearance results

8 events after selection

No-oscillation prediction:
 22.81 ± 3.19 (systematics)

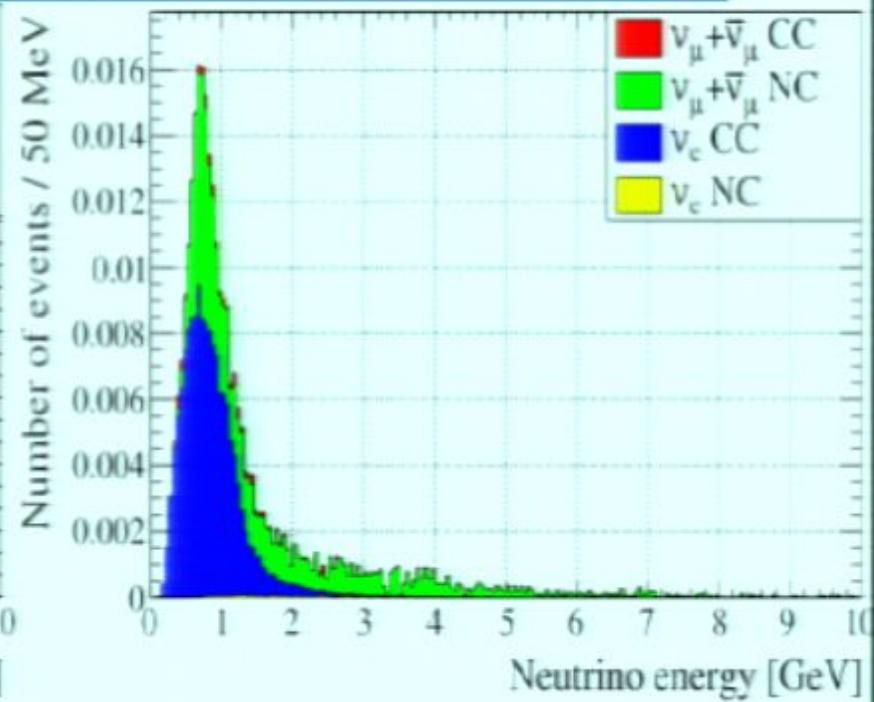
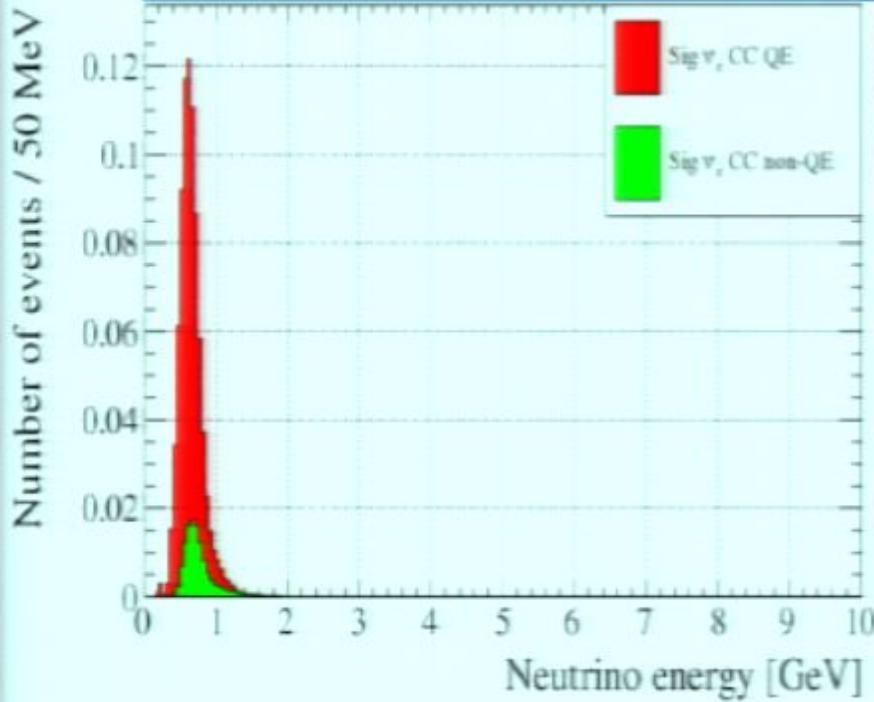
With oscillation:
 6.34 ± 1.04 (systematics)

at MINOS experimental best fit:
 $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 2\theta_{23} = 1.0$



v_e appearance results

ν_e appearance at T2K



Signal: CC interactions of ν_e from ν_μ to ν_e oscillations

Signal	# events
@ $\sin^2 2\theta_{13} = 0.1, \delta \text{cp} = 0$	0.9

Background	# events
intrinsic ν_e	0.16
ν_μ (70% NC π^0)	0.13
$\bar{\nu}_\mu$	0.01
total:	$0.30 \pm 0.07 \text{ (sys)}$

ν_e events after cuts

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
Selection cut

ν_e selection

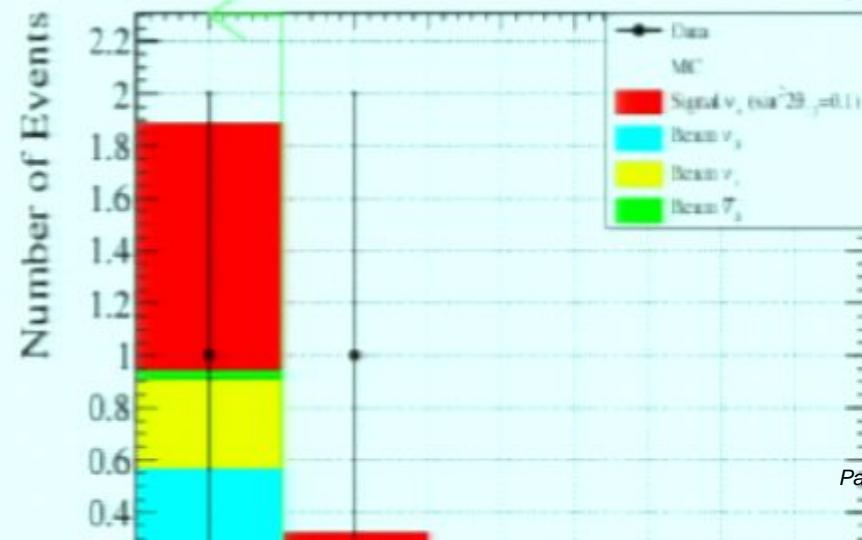
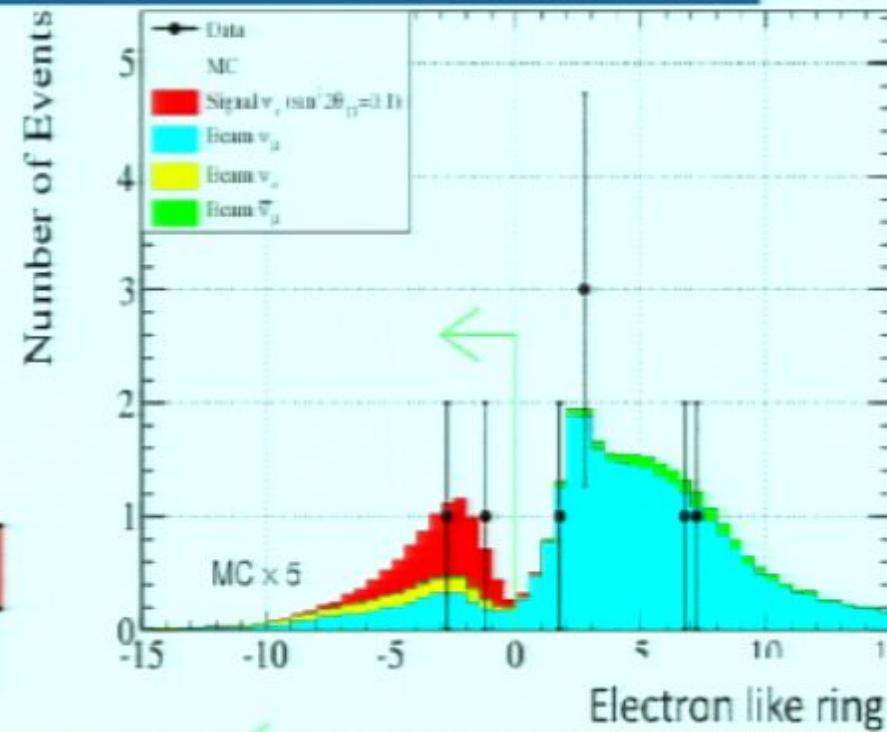
Visible energy > 100 MeV

e-like ring

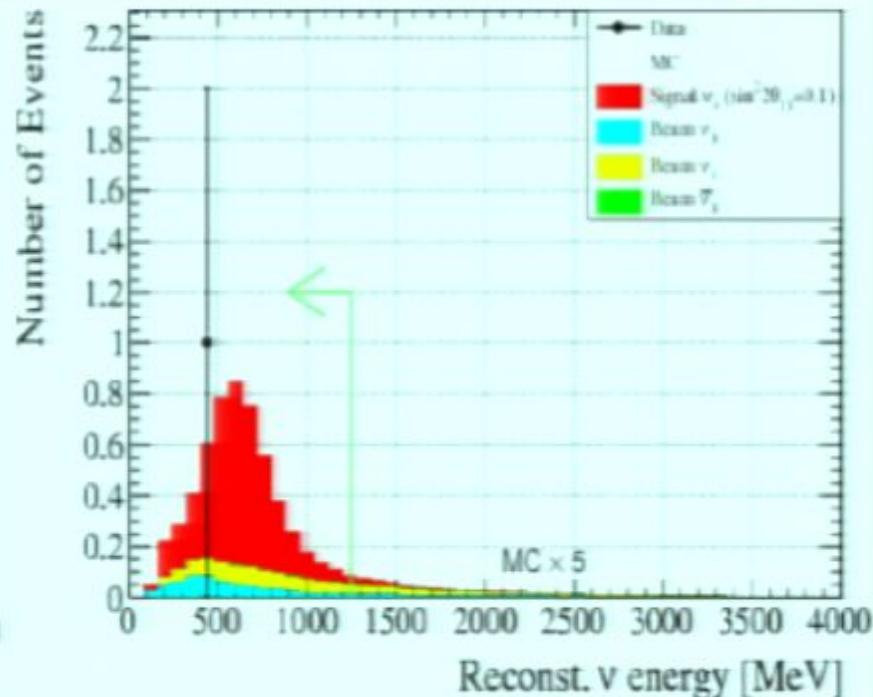
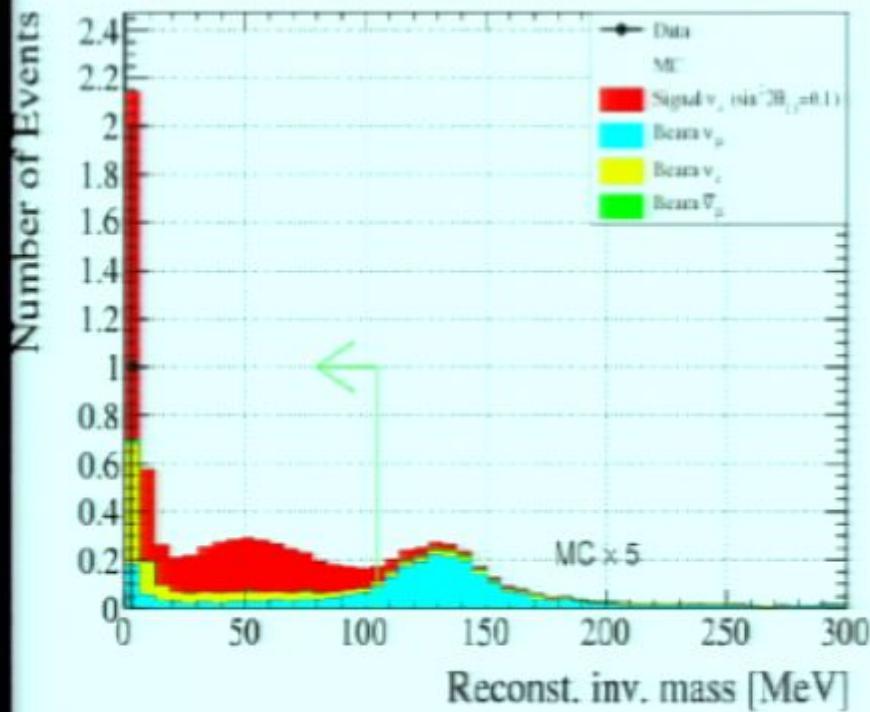
No decay electron

Invariant mass < 105 MeV/c²

E_V < 1250 MeV



ν_e events after cuts



ν_e selection

Visible energy > 100 MeV

e-like ring

No decay electron

Invariant mass < 105 MeV/c²

E_v < 1250 MeV

11040088

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
 Selection cut

Page 80/105

ν_e events after cuts

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
Selection cut

ν_e selection

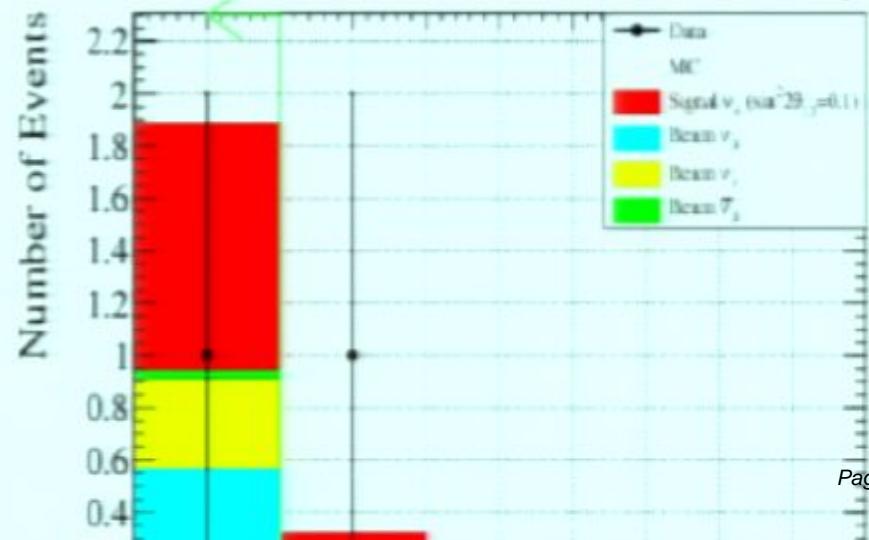
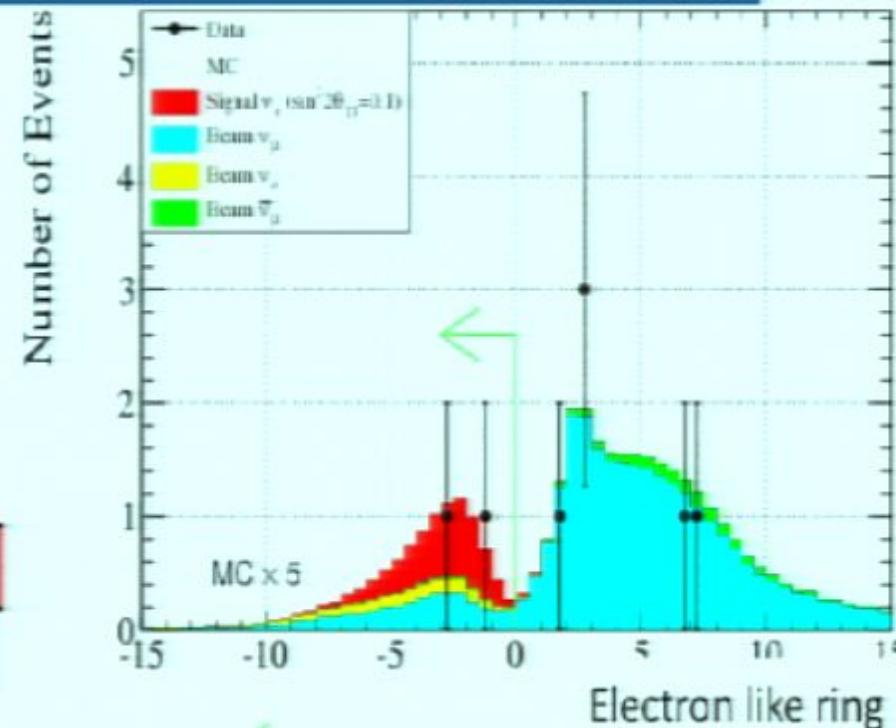
Visible energy > 100 MeV

e-like ring

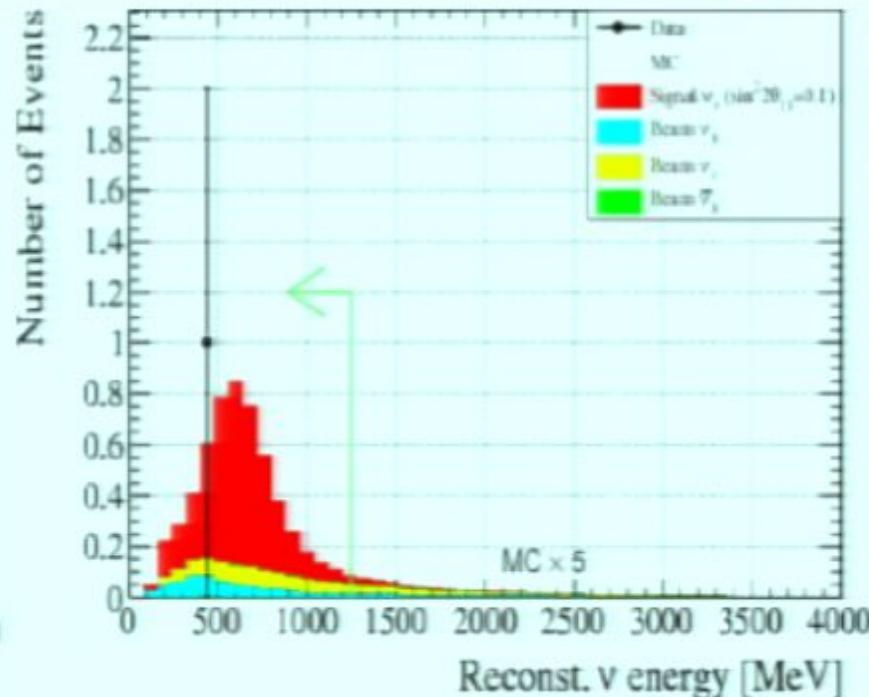
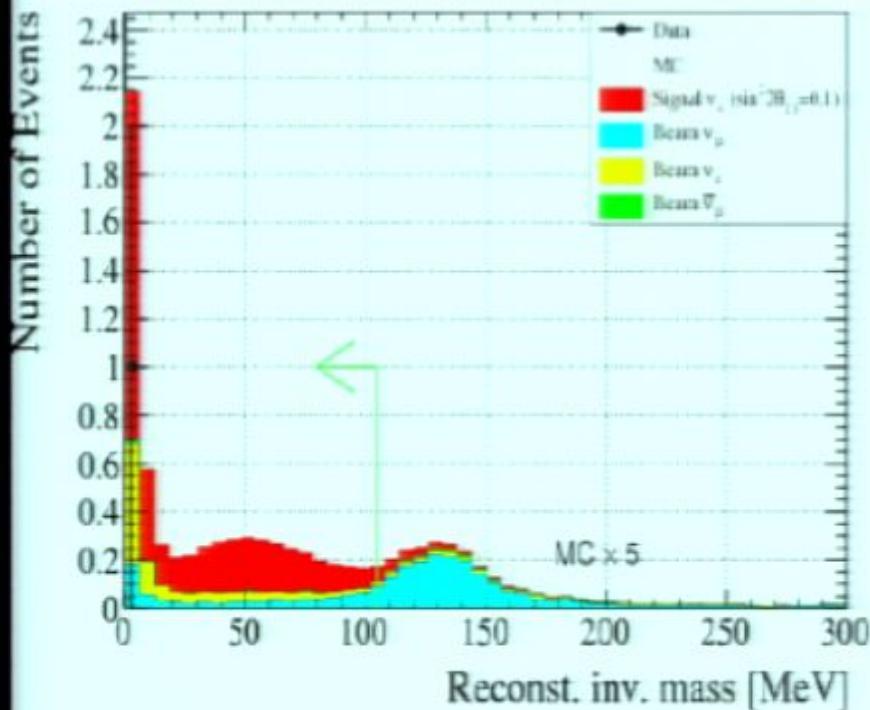
No decay electron

Invariant mass < 105 MeV/c²

E_V < 1250 MeV



ν_e events after cuts



ν_e selection

Visible energy > 100 MeV

e-like ring

No decay electron

Invariant mass < 105 MeV/c²

11040088

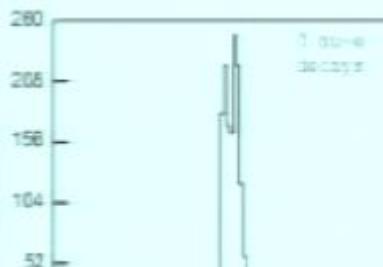
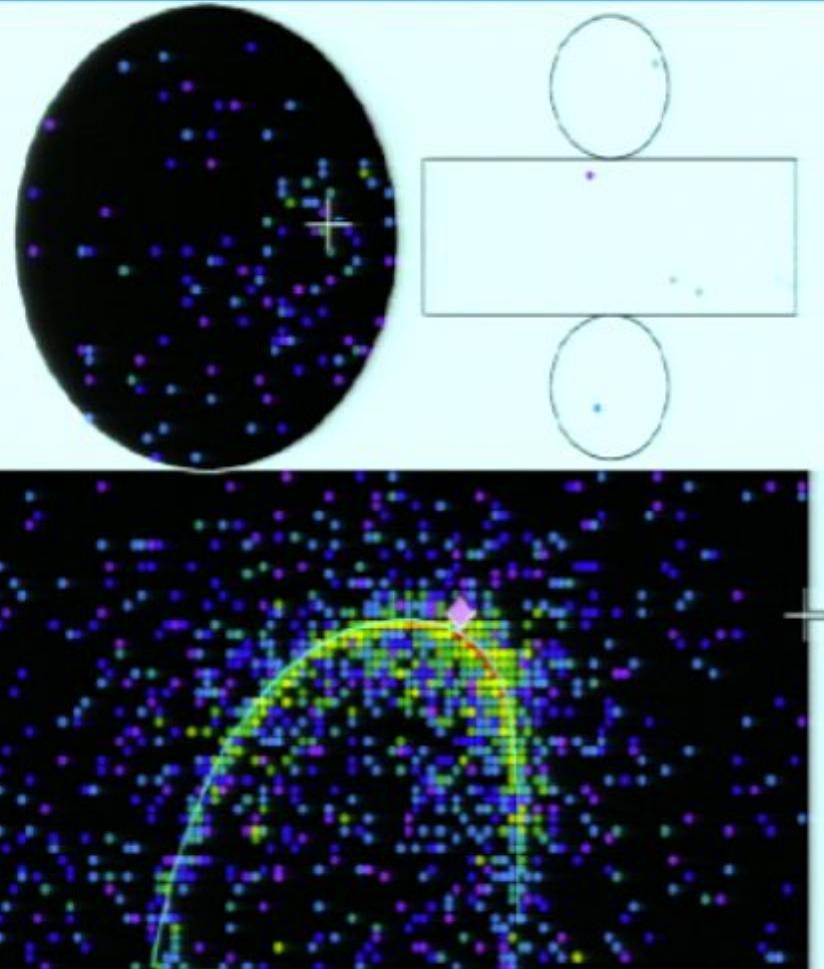
E_v < 1250 MeV

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
 Selection cut

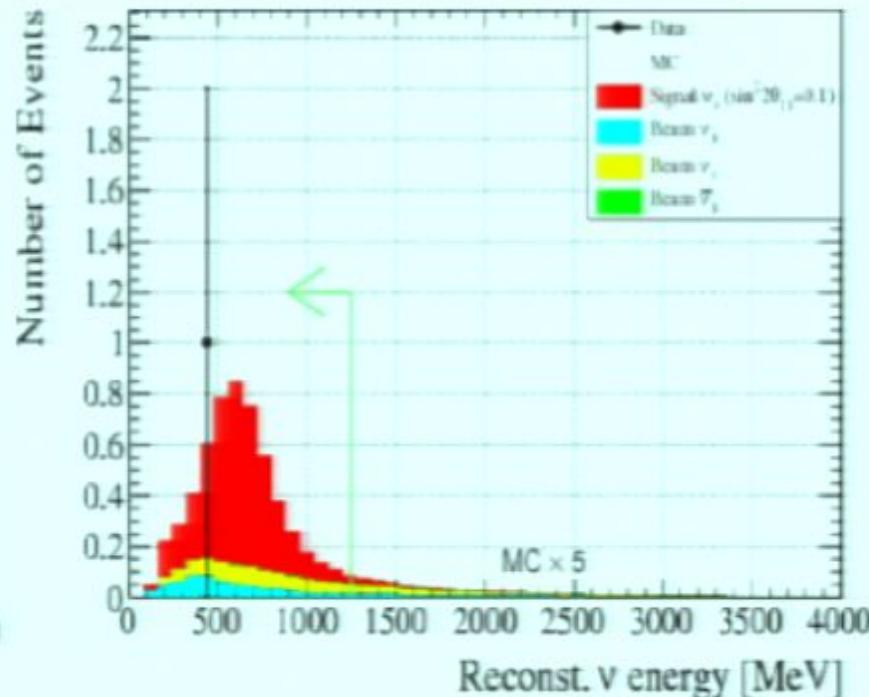
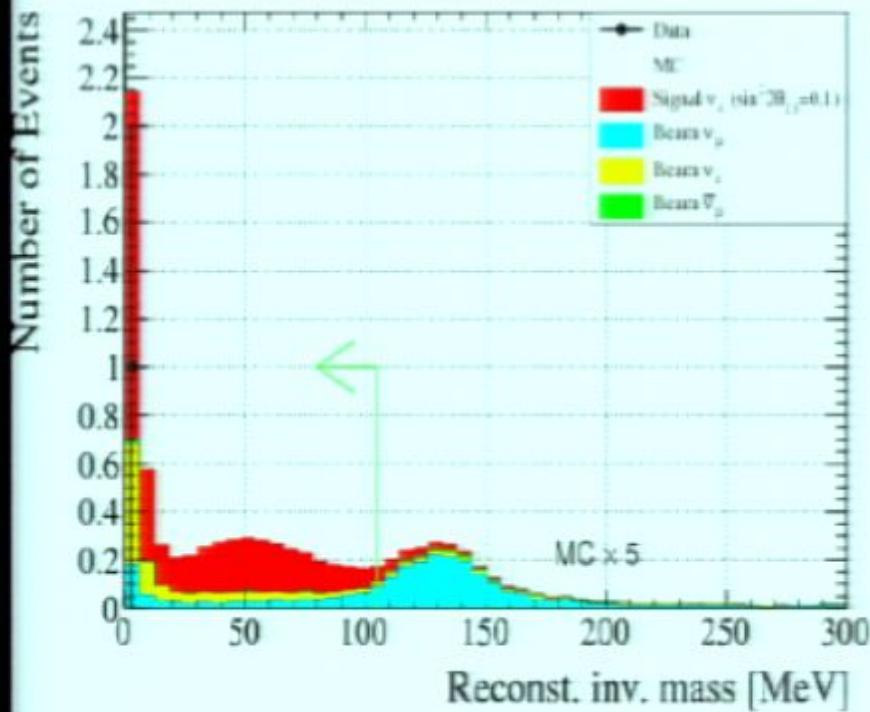
The candidate ν_e event

Super-Kamiokande IV

T2K Beam Run C Spill 822275
Run 66778 sub 585 Event 134229437
10-06-12:11:03:122
T2K beam fit = 1362.0 ns
Invert 1610 hits, 1681 ps
Invert 2 hits, 2 pe
Trigger: 1e0000000
D_mu11: 604.9 ns
alpha_0: p = 177.6 MeV/c



ν_e events after cuts



ν_e selection

Visible energy > 100 MeV

e-like ring

No decay electron

Invariant mass < 105 MeV/c²

11040088

E_v < 1250 MeV

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
 Selection cut

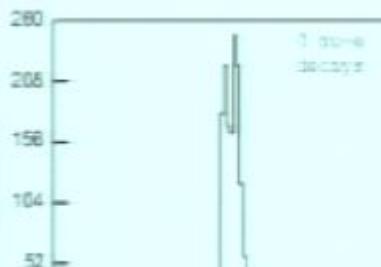
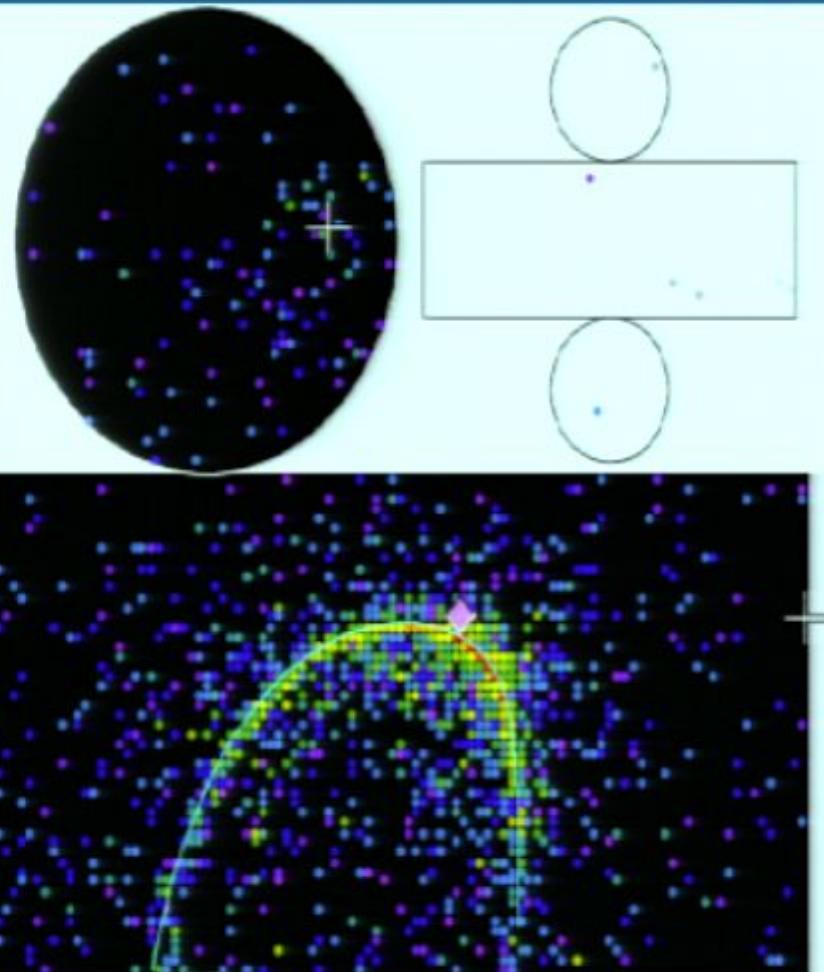
The candidate ν_e event

Super-Kamiokande IV

T2K Beam Run C Spill 822275
Run 66778 sub 585 Event 134229437
10-06-12:01:03:122
T2K beam fit = 1362.2 m
Invert 1810 hits, 1881 ps
Invert 2 hits, 2 pe
Trigger: 1x0000000
 $d_{\text{max}} = 604.4$ cm
 $\phi_{\text{angle}} = 377.6$ mrad/c

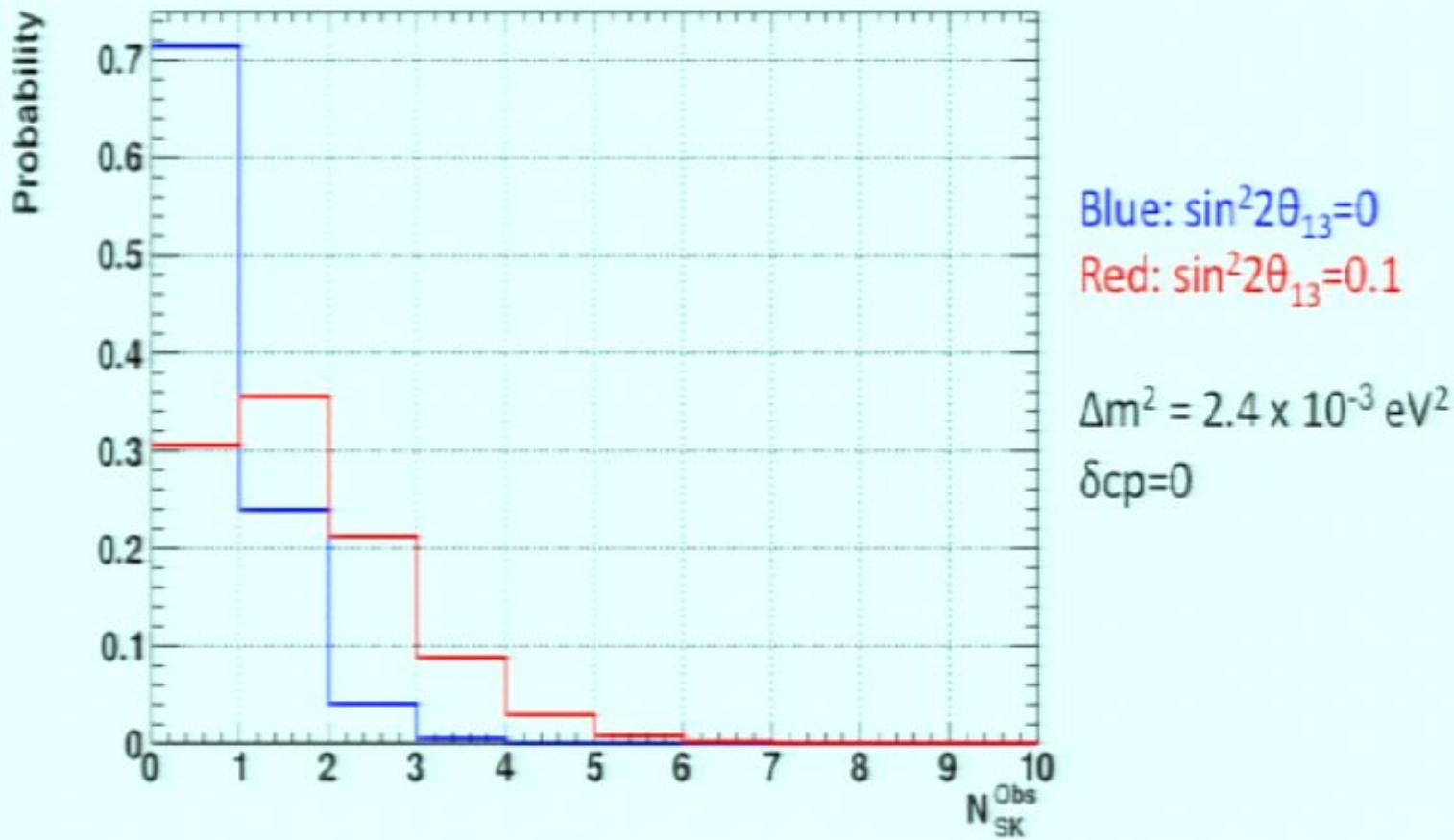
Charge (pe)

- >25.7
- 23.0-25.7
- 20.2-23.0
- 17.0-20.2
- 14.0-17.0
- 12.0-14.0
- 10.0-12.0
- 8.0-10.0
- 6.0-8.0
- 4.0-6.0
- 3.0-4.0
- 2.0-3.0
- 1.0-2.0
- 0.7-1.3
- 0.2-0.7
- < 0.2



ν_e appearance results

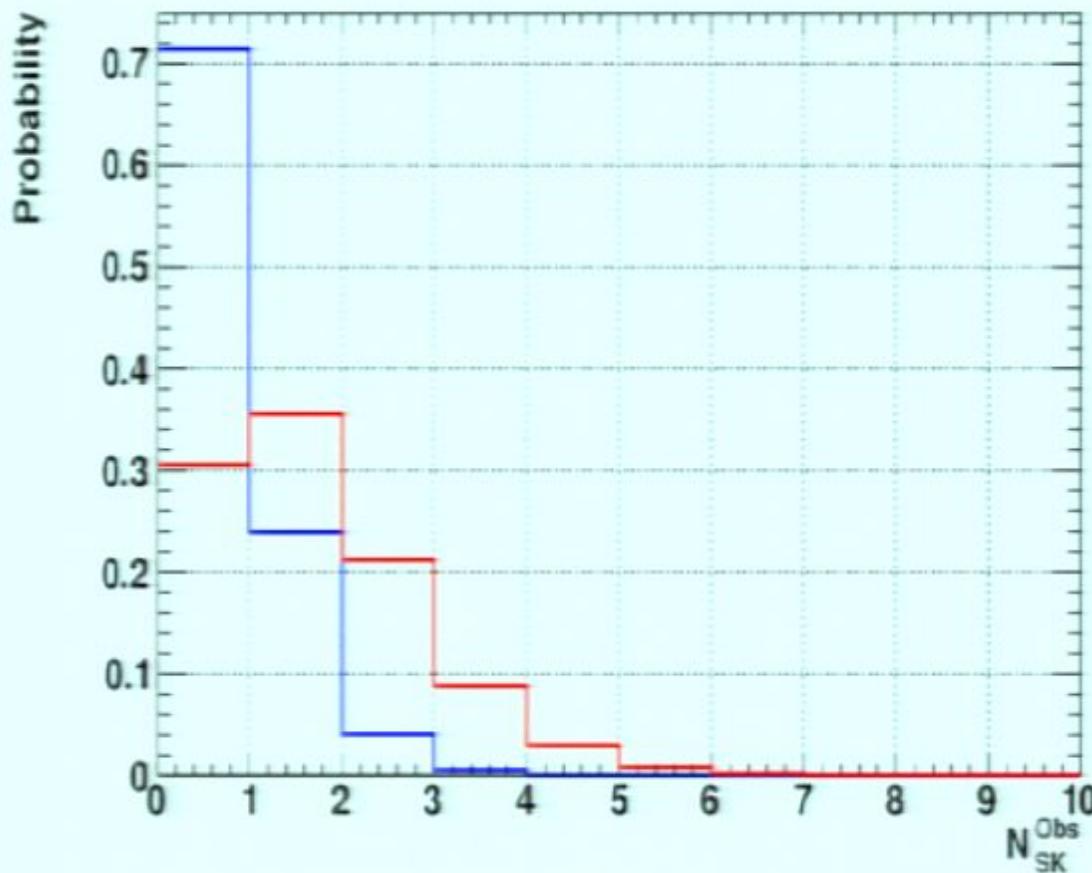
One candidate event remains after selection



PDF for expected number of Super-K ν_e events, including statistical and systematic errors

ν_e appearance results

One candidate event remains after selection



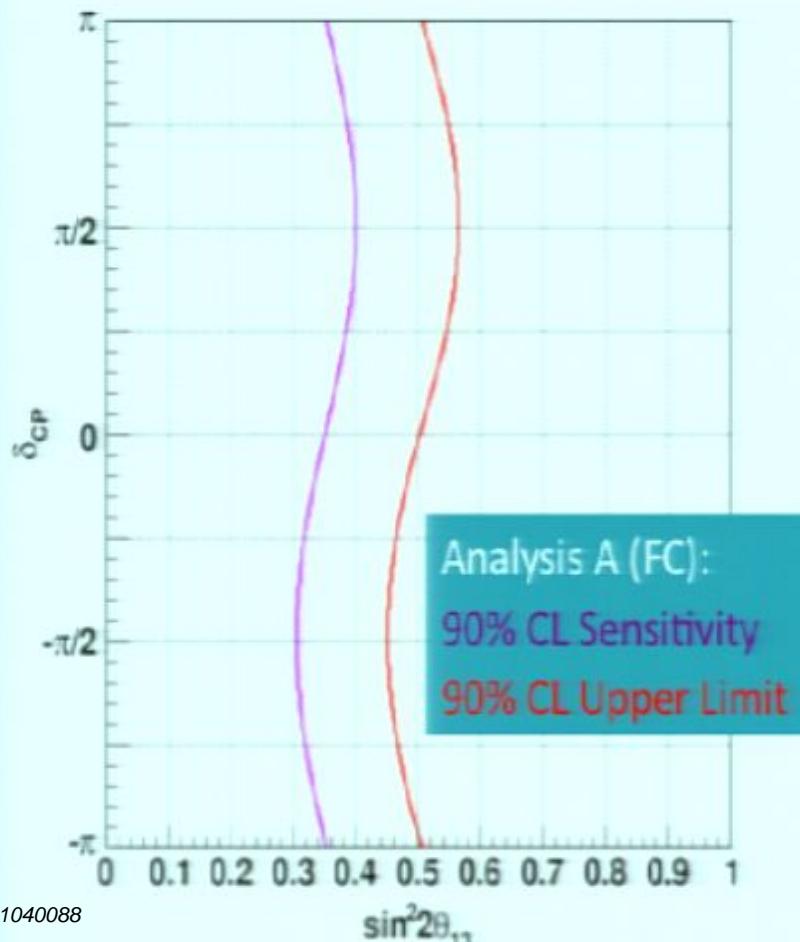
Blue: $\sin^2 2\theta_{13} = 0$
Red: $\sin^2 2\theta_{13} = 0.1$

$\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$
 $\delta \text{cp} = 0$

PDF for expected number of Super-K ν_e events, including statistical and systematic errors

ν_e oscillation sensitivity and limits

Two separate methods used to determine sensitivities and limits



Analysis A: Feldman-Cousins

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.50	0.35
Inverted ($\Delta m_{23}^2 < 0$)	0.59	0.42

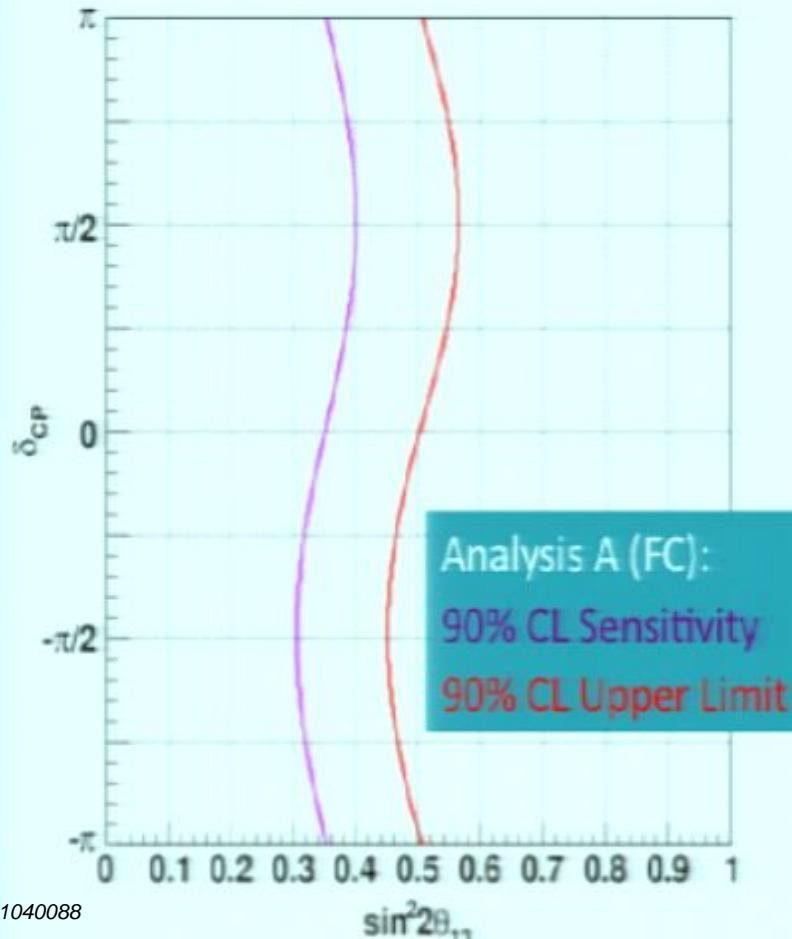
Analysis B: Classical one-sided limit

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.44	0.32
Inverted ($\Delta m_{23}^2 < 0$)	0.53	0.39

Future T2K analysis prospects

ν_e oscillation sensitivity and limits

Two separate methods used to determine sensitivities and limits



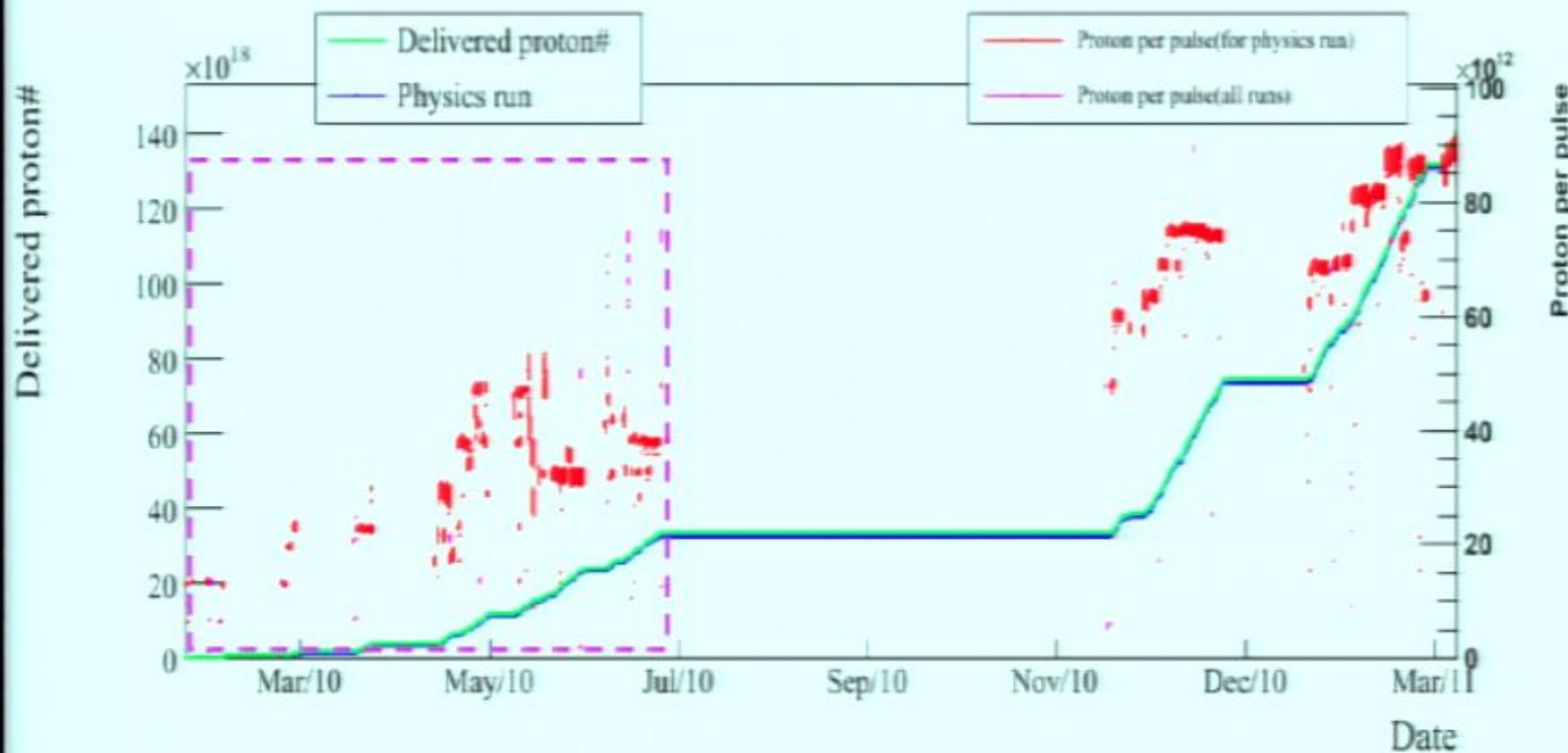
Analysis A: Feldman-Cousins

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.50	0.35
Inverted ($\Delta m_{23}^2 < 0$)	0.59	0.42

Analysis B: Classical one-sided limit

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.44	0.32
Inverted ($\Delta m_{23}^2 < 0$)	0.53	0.39

Run 2



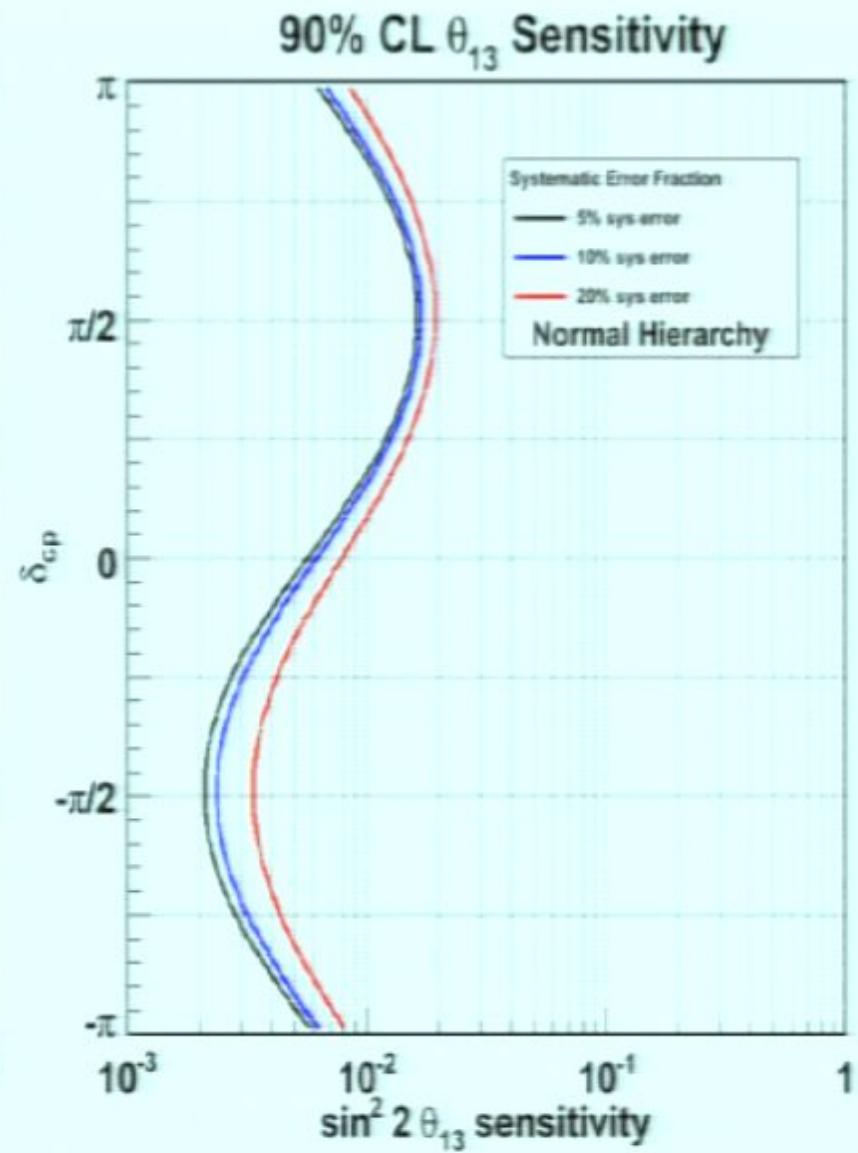
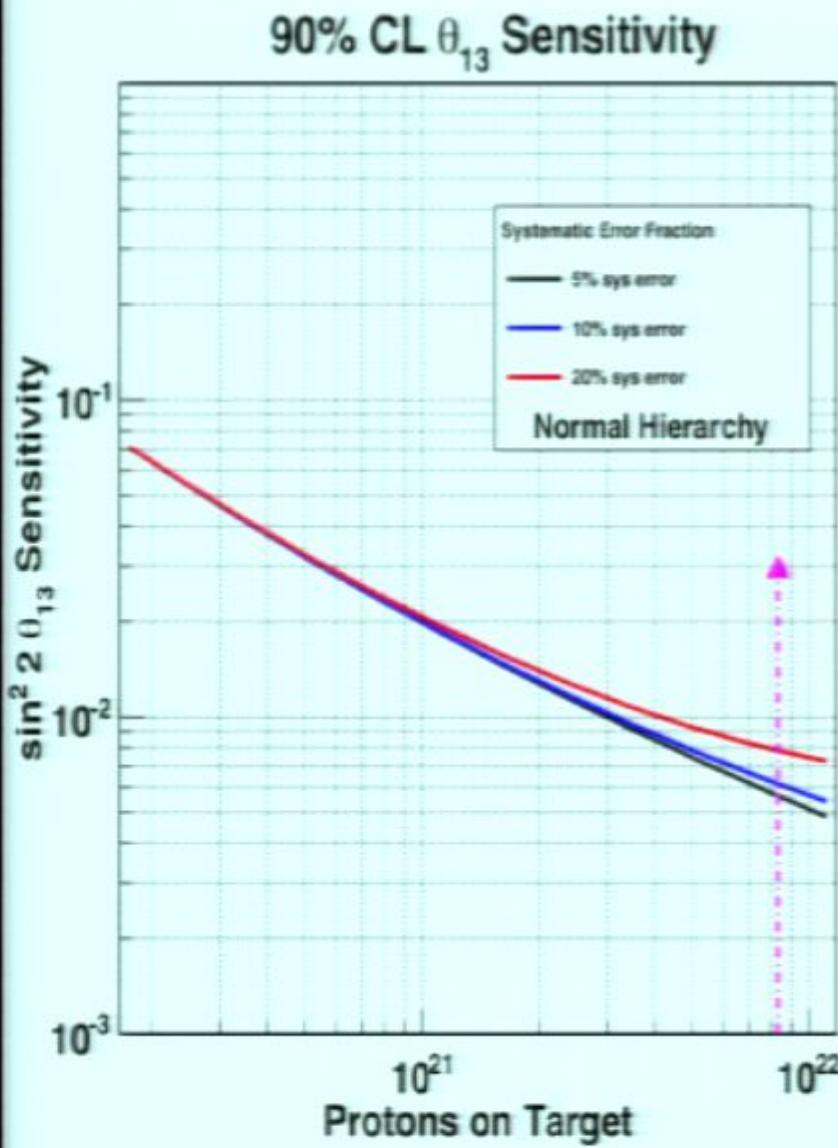
Achieved 145kW stable run in March prior to earthquake

Entire run at 8 bunches / spill / 3.04s

1.45×10^{20} POT collected total

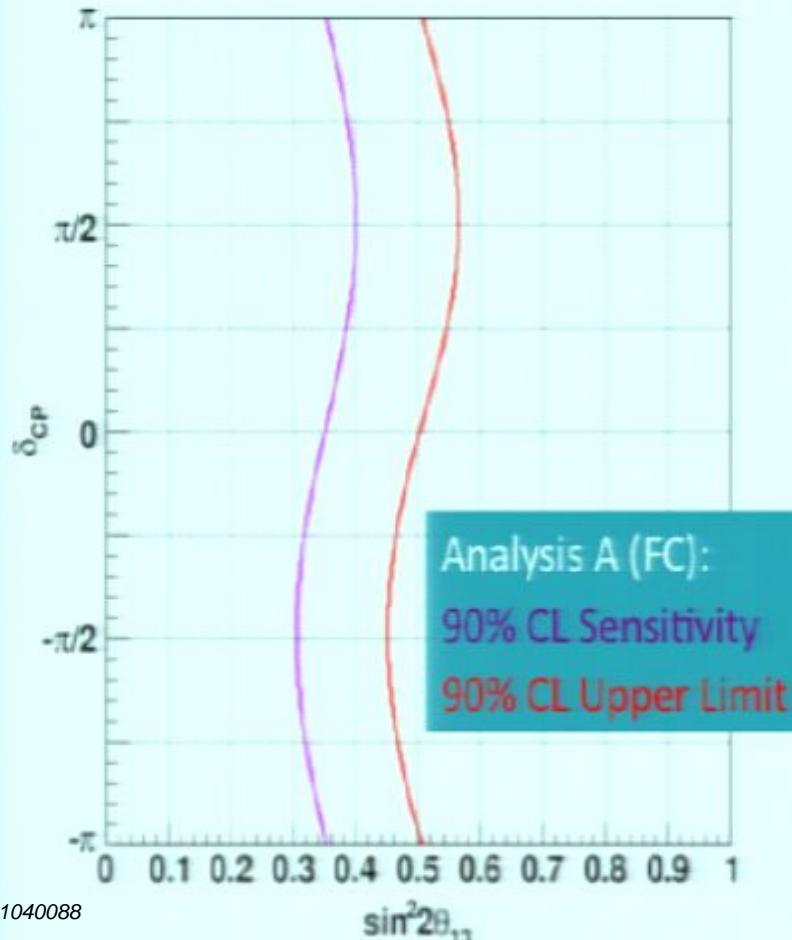
Extra POT is ~4x dataset presented today

Future ν_e appearance sensitivity



ν_e oscillation sensitivity and limits

Two separate methods used to determine sensitivities and limits



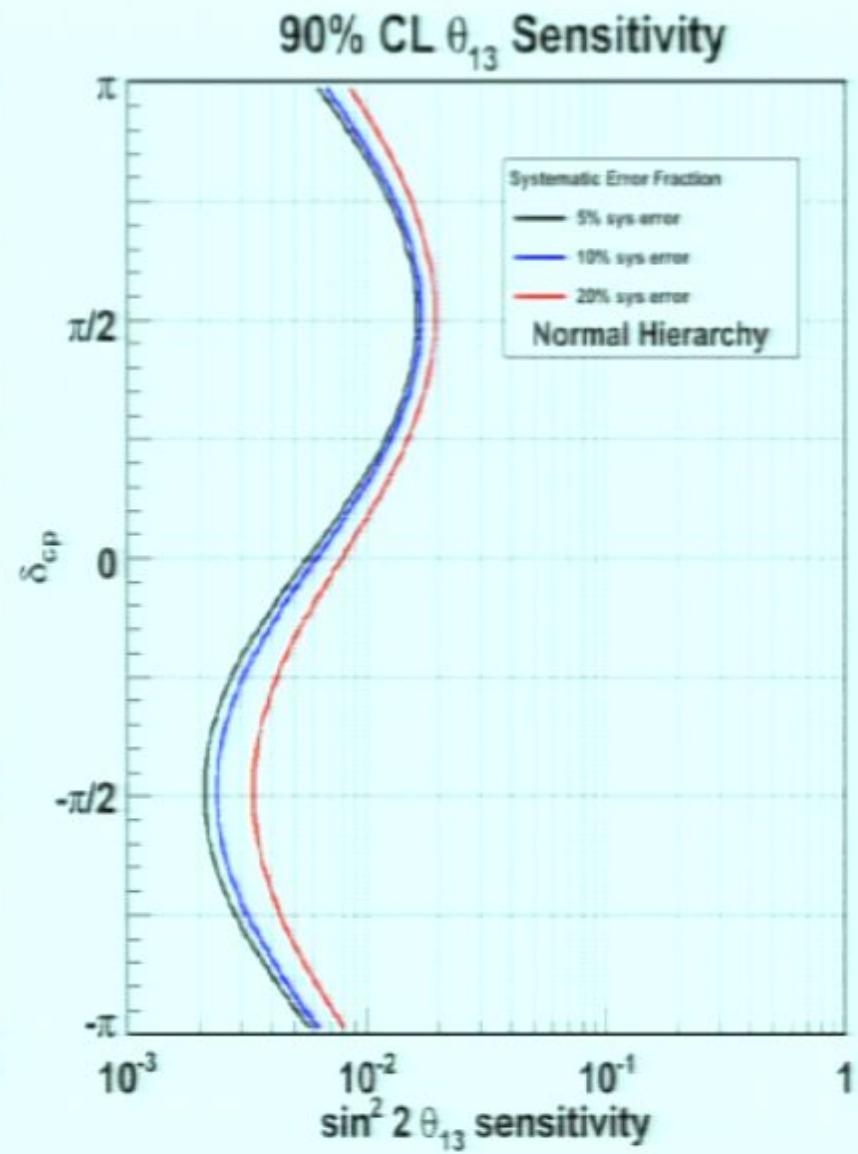
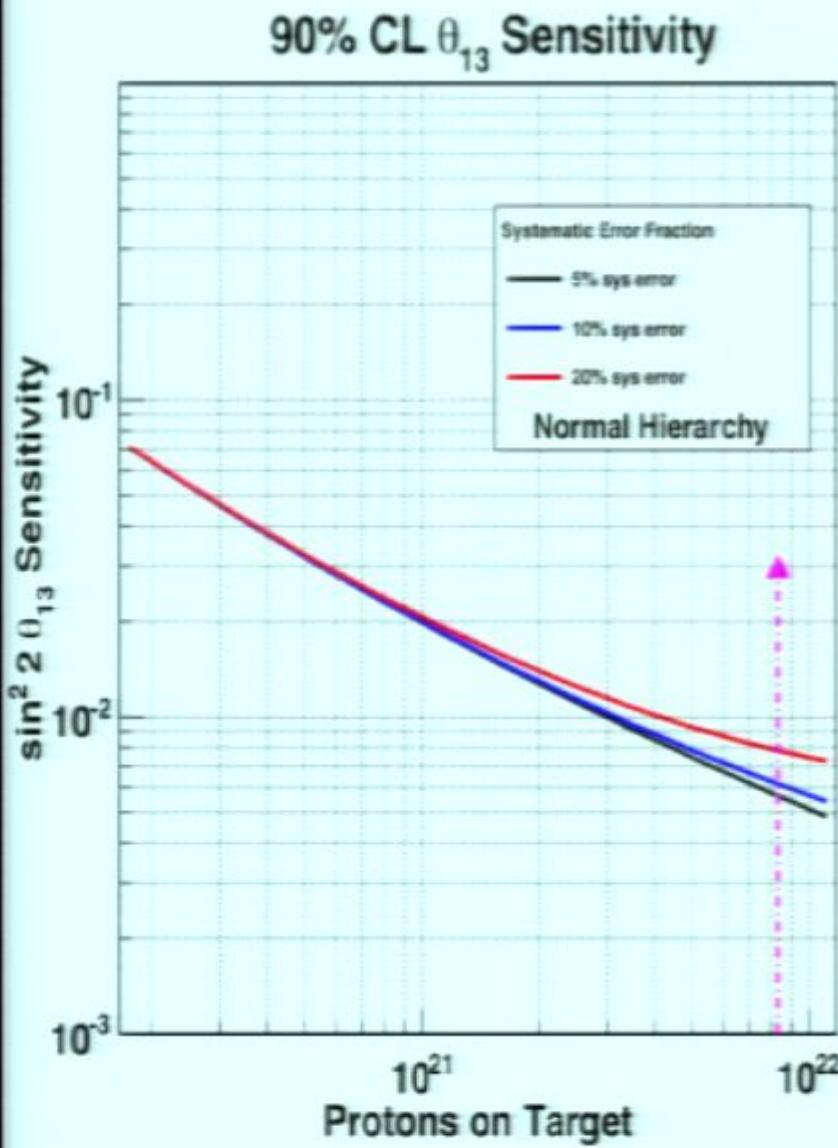
Analysis A: Feldman-Cousins

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.50	0.35
Inverted ($\Delta m_{23}^2 < 0$)	0.59	0.42

Analysis B: Classical one-sided limit

Hierarchy	Upper Limit	Sensitivity
Normal ($\Delta m_{23}^2 > 0$)	0.44	0.32
Inverted ($\Delta m_{23}^2 < 0$)	0.53	0.39

Future ν_e appearance sensitivity



Planned improvements to the analysis

Improvements to the flux uncertainties:

Reduced π production systematics from NA61

Addition of NA61 kaon data

Error (+)	N_{SK}^{Sig}	N_{SK}^{bkd}	$N_{SK}^{\text{S+B}}$	N_{ND}	$N_{SK}^{\text{bkd}}/N_{ND}$	$N_{SK}^{\text{S+B}}/N_{ND}$
Flux	21.97	18.12	20.49	19.83	9.17	11.88
CCQE	4.91	2.62	4.33		2.72	4.33
CC1 π	4.28	3.76	4.15	5.93	2.10	1.78
NC π^0	-	5.86	1.48	0.05	5.56	1.43
FSI	3.83	10.34	5.47	-	10.32	5.47
ND detector	-	-	-	5.60	5.60	5.60
ring counting	3.90	8.40	5.03	-	8.40	5.03
electron PID	3.80	8.10	4.88	-	8.10	4.88
invariant mass	5.10	8.70	6.01	-	7.70	6.01

Planned improvements to the analysis

Improvements to the constraints provided by ND280:

ν_μ spectrum measurement

CCQE, CC π rate measurements

Intrinsic ν_e measurement

Error (+)	N_{SK}^{Sig}	N_{SK}^{bkd}	N_{SK}^{S+B}	N_{ND}	N_{SK}^{bkd}/N_{ND}	N_{SK}^{S+B}/N_{ND}
Flux	21.97	18.12	20.49	19.83	9.17	11.88
CCQE	4.91	2.62	4.33		2.72	4.33
CC 1π	4.28	3.76	4.15	5.93	2.10	1.78
NC π^0	-	5.86	1.48	0.05	5.56	1.43
FSI	3.83	10.34	5.47	-	10.32	5.47
ND detector	-	-	-	5.60	5.60	5.60
ring counting	3.90	8.40	5.03	-	8.40	5.03
electron PID	3.80	8.10	4.88	-	8.10	4.88
invariant mass	5.10	8.70	6.01	-	7.70	6.01

Planned improvements to the analysis

Improved Super-K selection cuts

Improved Super-K detector uncertainties

Error (+)	N_{SK}^{Sig}	N_{SK}^{bkd}	N_{SK}^{S+B}	N_{ND}	N_{SK}^{bkd}/N_{ND}	N_{SK}^{S+B}/N_{ND}
Flux	21.97	18.12	20.49	19.83	9.17	11.88
CCQE	4.91	2.62	4.33		2.72	4.33
CC1 π	4.28	3.76	4.15	5.93	2.10	1.78
NC π^0	-	5.86	1.48	0.05	5.56	1.43
FSI	3.83	10.34	5.47	-	10.32	5.47
ND detector	-	-	-	5.60	5.60	5.60
ring counting	3.90	8.40	5.03	-	8.40	5.03
electron PID	3.80	8.10	4.88	-	8.10	4.88
invariant mass	5.10	8.70	6.01	-	7.70	6.01

Summary

The T2K experiment is designed to make precision measurements of:

ν_μ disappearance (Δm^2_{23} , θ_{23})

ν_e appearance (θ_{13})

With the initial 6 month run ending in Jun 2010:

the ν_μ dataset of 8 events is consistent with disappearance measured by previous experiments (MINOS, Super-K, K2K)

One candidate ν_e event was observed

Expected background is 0.30 ± 0.07

Analysis of the 1.45×10^{20} POT collected thus far is underway

Expected sensitivity is better than the current MINOS/CHOOZ limit

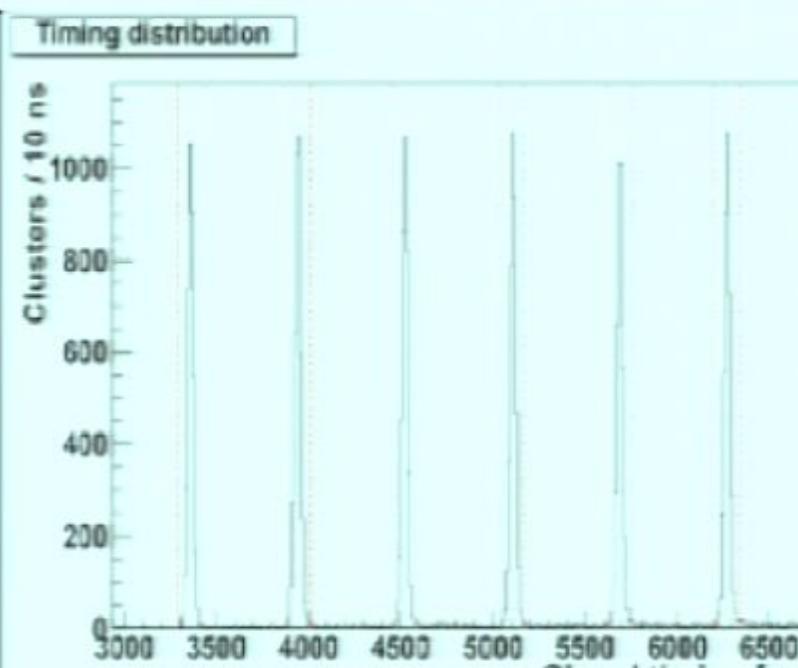
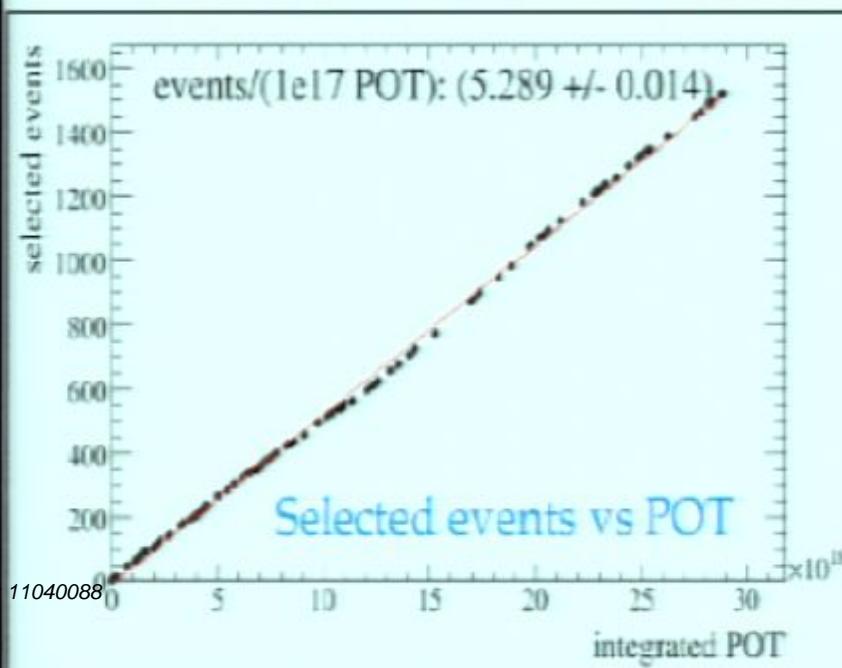
ND280 performance

Low rate of broken channels

Detector	Channels	Bad ch.	Bad fraction
ECAL (DSECAL)	22,336 (3,400)	35 (11)	0.16% (0.32%)
SMRD	4,016	7	0.17%
P0D	10,400	7	0.07%
FGD	8,448	20	0.24 %
INGRID	10,796	18	0.17 %
TPC	124,416	160	0.13 %

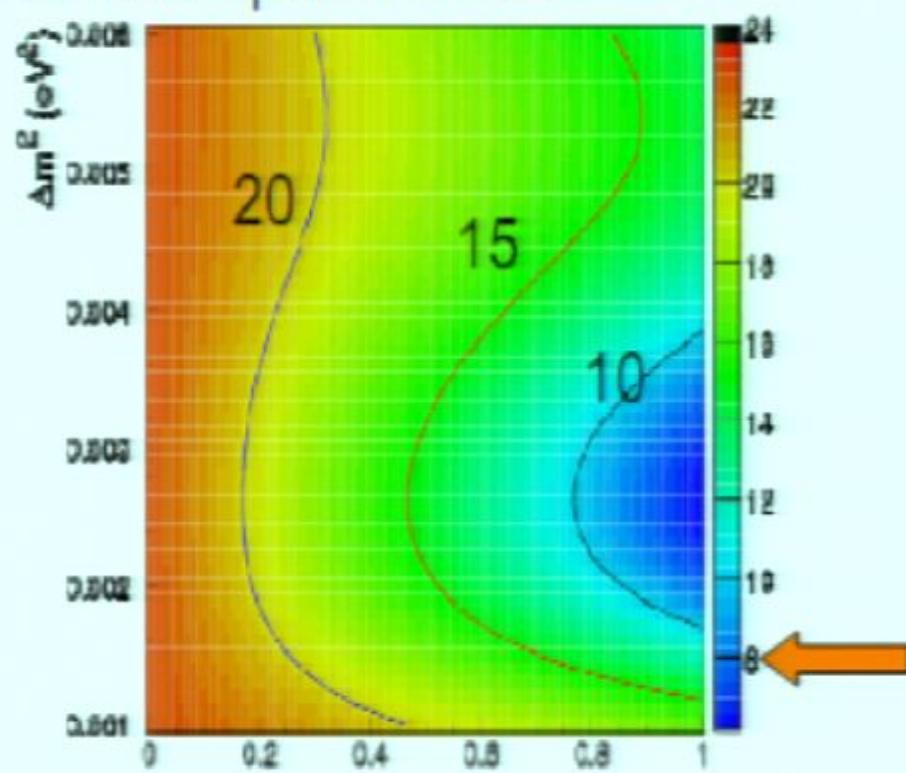
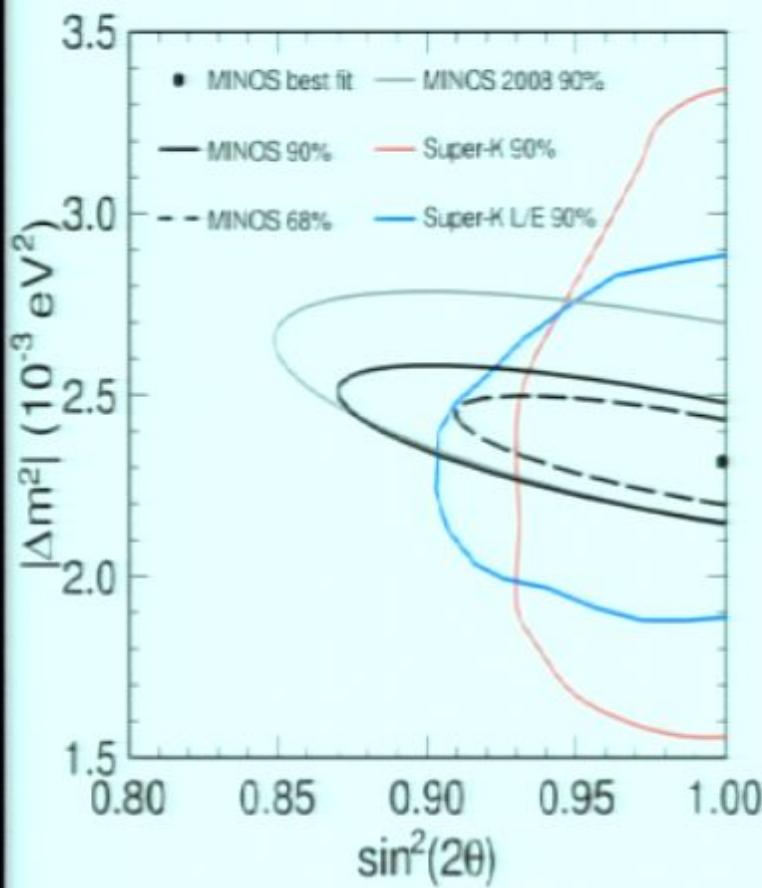
Events / POT stable

Timing consistent with beam (FGD)

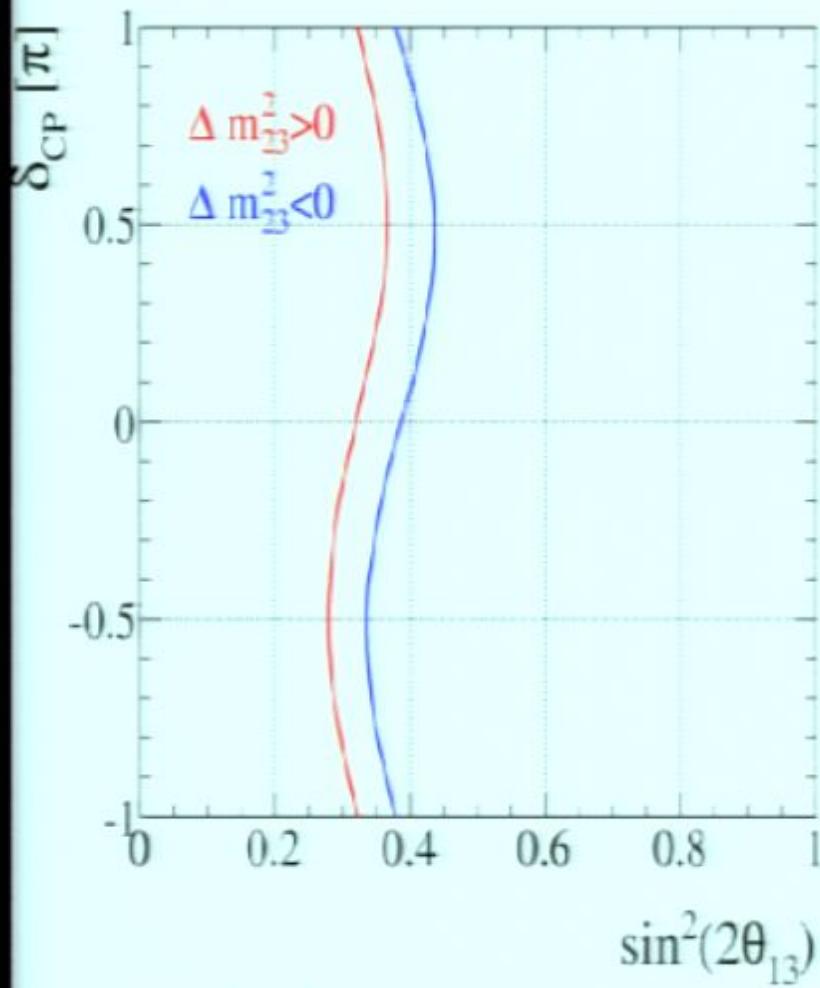


ν_μ disappearance consistency

Expected # of events as a function of Oscillation parameters.



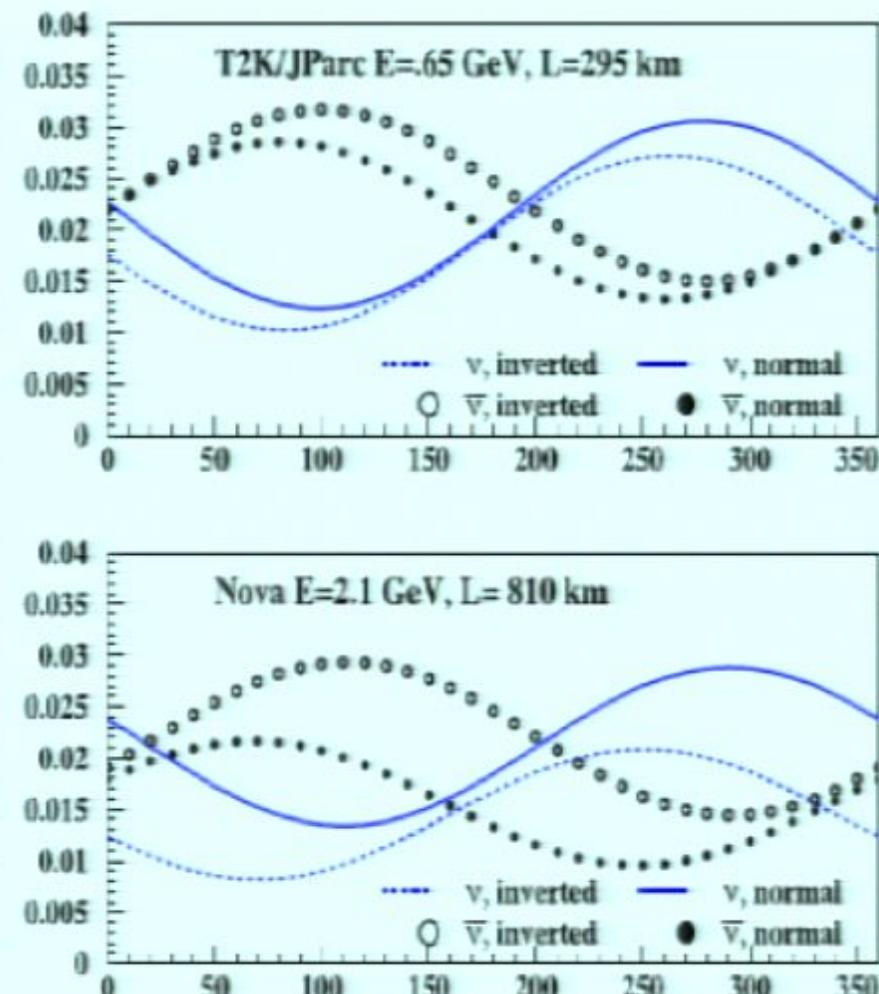
v_e results: normal vs. inverted hierarchy



Analysis B:

Normal hierarchy

Inverted hierarchy



Oscillation probability vs. dCP (deg)

1 pt

Slide Themes Slide Layouts Transitions Table Styles Charts SmartArt Graphics WordArt

5 | 4 | 3 | 2 | 1 | 0 | 1 | 2 | 3 | 4 | 5

The T2K Collaboration



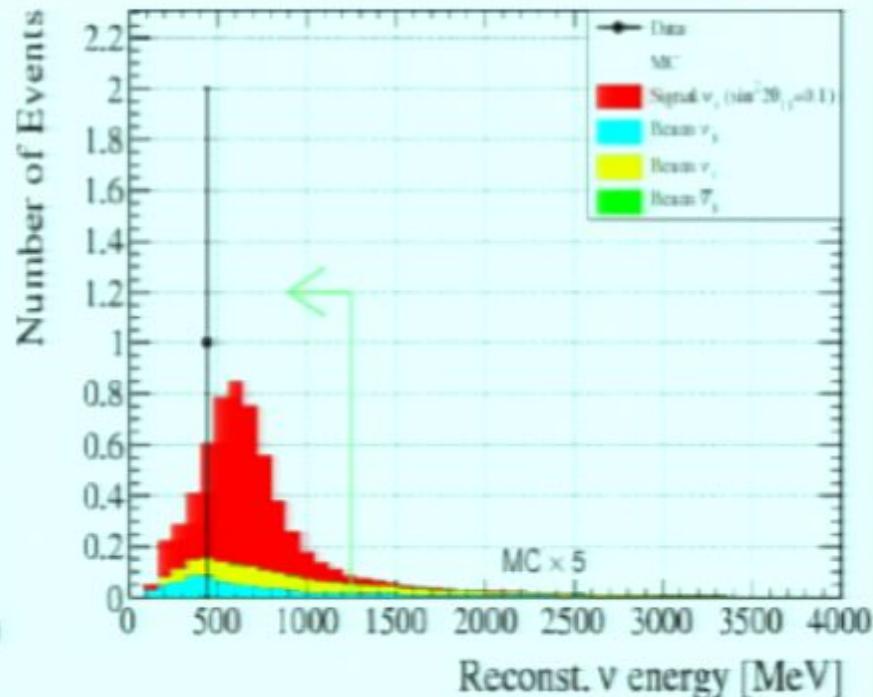
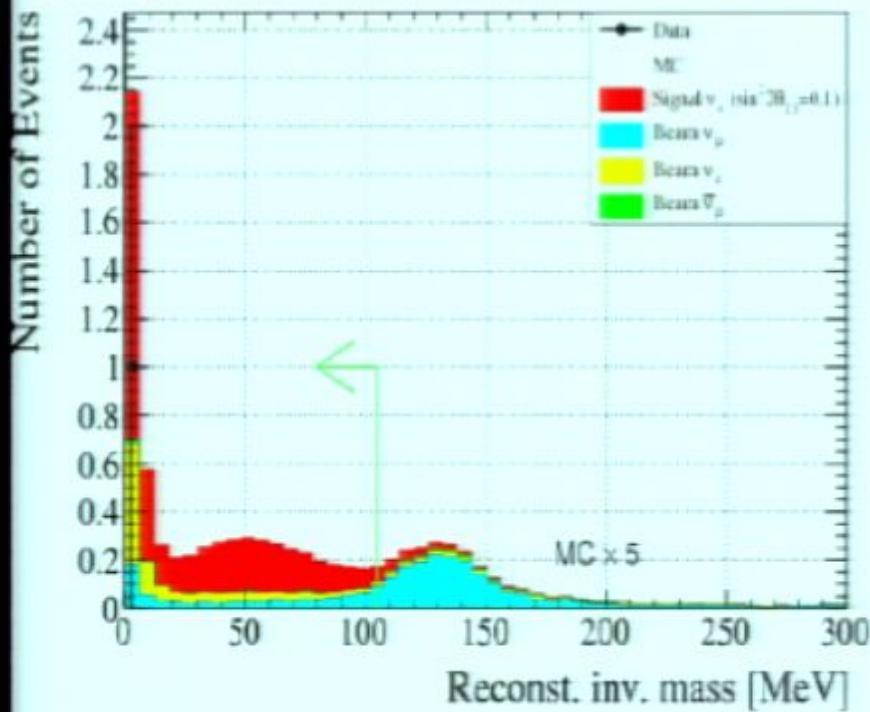
59 institutions in 12 countries

Canada	Korea	Switzerland	Japan	USA
TRUMF	Chonnam Nat'l U	Bern	ICRR Kamioka	Boston U
U of Alberta	Duisburg U	ETH Zurich	ICRR RICCN	BNL
U of B.C. Columbia	Seoul Nat'l U	U of Geneva	KDG	Colorado State U.
U of Regensburg	Spain	UK	Kanagawa U	Duke U.
U of Toronto	IFIC Valencia	U of Oxford	Kyoto U	Louisiana State U.
U of Victoria	U.A. Barcelona	Imperial Coll London	Meijo U of Ec.	Stanford Univ.
York U	Poland	Lancaster U	Osaka City U	U of California, Irvine
France	A.S. Soltan Warsaw	Queen Mary U of L.	U of Tokyo	U of Colorado
CEA Saclay	Hamburg University	Sheffield U	Italy	U of Pittsburgh
IPN Lyon	TU Warsaw	STFCRAL	NFN Genova	U of Rochester
LURE Paris	U of Silesia	STFC Daresbury	NFN Roma	U of Washington
LPNHE-Paris	Warsaw U.	U of Liverpool	Napoli U.	Germany
Russia	Wroclaw U.	U of Warwick	Pattaya U.	RWTH Aachen U.
INR				



I'm here on behalf of the T2K collaboration, approximately 500 people with members around the world.

ν_e events after cuts



ν_e selection

Visible energy > 100 MeV

e-like ring

No decay electron

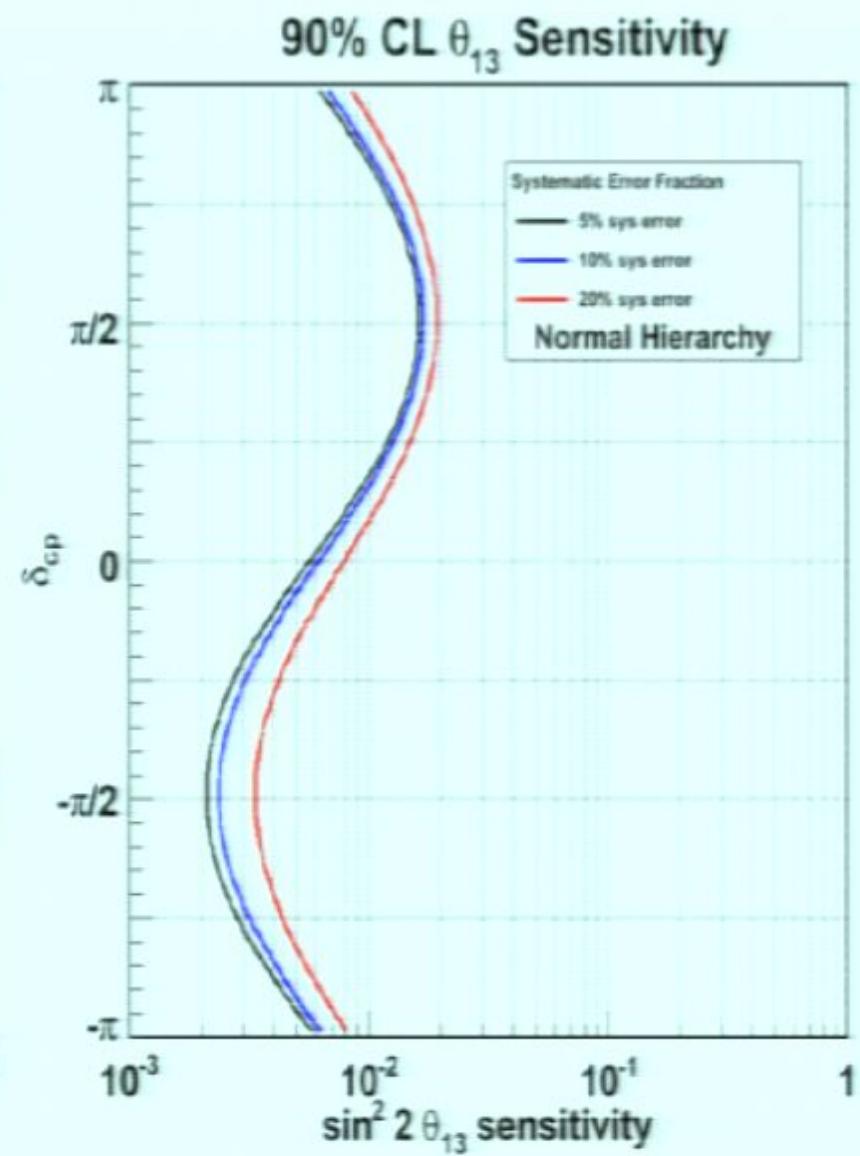
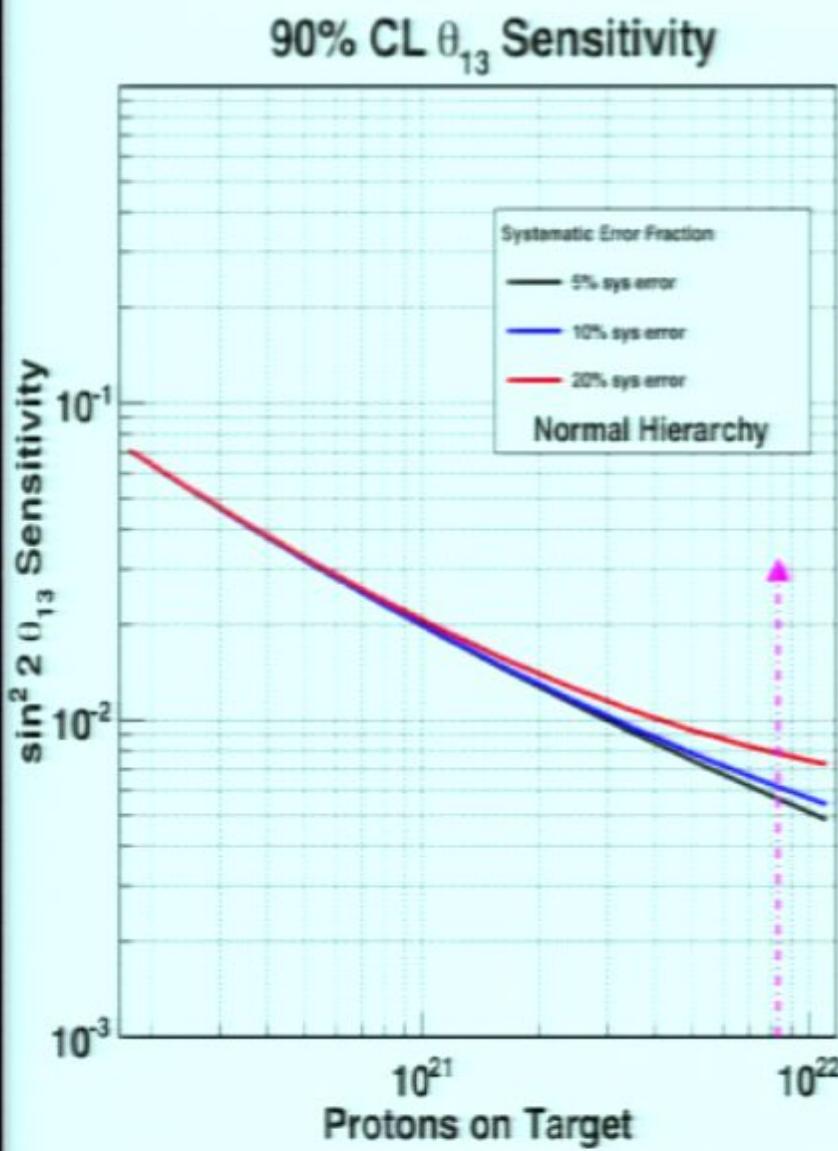
Invariant mass < 105 MeV/c²

11040088

E_v < 1250 MeV

Signal ν_e
 ν_e from beam
 ν_μ (NC/CC)
 Selection cut

Future ν_e appearance sensitivity



Future ν_e appearance sensitivity

