

Title: Explorations in Cosmology - Lecture 2

Date: Apr 05, 2011 09:00 AM

URL: <http://pirsa.org/11040004>

Abstract:

least time: $ds^2 = a^2(-d\eta^2 + dx^2)$

$$\phi = \frac{u}{a}$$

$$S = \int d\eta d^3x \left[\frac{1}{2} Q \eta u^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

$$m_{\text{eff}}^2 = m^2 a^2 - \frac{a''}{a}$$

least time: $ds^2 = dx^2$

$$\phi = \frac{u}{a}$$

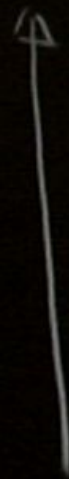
$$S = \int \left[\frac{1}{2} Q \eta u^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}} u^2 \right]$$

best time :

$$\phi = \frac{u}{a}$$

$$d\eta^c + dx^c$$

$$\int d^3x \left[\frac{1}{2} Q \eta u^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{eff}^c u^2 \right]$$



next time = $d(\eta^c + dx^c)$

$$\phi = \frac{u}{a}$$

$$\int d^3x \left[\frac{1}{2} Q |\eta u|^2 - \frac{1}{2} (\vec{\nabla} |\eta u|^2)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

+m_{eff} ↑

here time = $ds^2 = a^2(-dt^2 + dx^2)$

$$\phi = \frac{u}{a}$$

$$S = \int dt d^3x \left[\dots - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

m^2



here +

$$ds^2 = a^c (-dt^c + dx^c)$$

$$S = \int dt d^3x \left[\frac{1}{2} \rho u^c u^c - \frac{1}{2} (\vec{\nabla} u^c)^2 - \frac{1}{2} m_{eff}^c u^c \right]$$

$$m_{eff}^c = m^c a^c - \frac{a''}{a} + m_{ie}$$

next time :

$$ds^2 = a'(t) \dots$$

$$\phi = \frac{u}{a}$$

$$S = \int \left[\frac{1}{2} m'_{\text{eff}} |\dot{\mathbf{u}}|^2 - \frac{1}{2} (|\vec{\nabla} \phi|^2 - \frac{1}{2} m'_{\text{eff}} u^2) \right]$$



least time: $ds^2 = a^2(-dt^2 + dx^2)$

$$\phi = \frac{u}{a}$$

$$S = \int dt d^3x \left[|u|^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

m^2

here

$$ds^2 = a^2(-d\eta^2 + dx^2)$$

$$S = \int d\eta d^3x \left[\frac{1}{2} Q_{\eta\eta} |\dot{u}|^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

$$m_{\text{eff}}^2 = m^2 a^2 - \frac{a''}{a}$$

$$k^2 + m_{\text{eff}}^2$$

+ noise \uparrow

$\omega' \ll \omega^2$

WKB good

$\omega' > \omega^2$

WKB bad

WKB

well defined positive energy

$\omega' \ll \omega^2$

next time: $ds^2 = a^2(-d\eta^2 + dx^2)$

$$\phi = \frac{u}{a}$$

$$S = \int d\eta d^3x \left[\frac{1}{2} (\partial_\eta u)^2 - \frac{1}{2} (\vec{\nabla} u)^2 - \frac{1}{2} m_{\text{eff}}^2 u^2 \right]$$

$$m_{\text{eff}}^2 = m^2 a^2 - \frac{a''}{a}$$

$$\omega^2 = k^2 + m_{\text{eff}}^2$$

time

\uparrow

$$\omega' \ll \omega^2$$

WKB good

$$\omega' > \omega^2$$

WKB bad

WKB

well defined positive energy

$$\omega' \ll \omega^2$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} \alpha_k u_k + \alpha_k^\dagger u_k^\dagger$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k \cdot x} + a_k^\dagger u_k^\dagger e^{-i k \cdot x}$$

$$\hat{U}(\eta, \mathbf{x}) = \int \frac{d^3 k}{(2\pi)^3} a_{\mathbf{k}} u_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* e^{-i\mathbf{k} \cdot \mathbf{x}}$$

$$u_{\mathbf{k}}'' = -\omega_{\mathbf{k}}^2 u_{\mathbf{k}} \quad u_{\mathbf{k}}(\eta)$$

$$\hat{U}(\eta, \mathbf{x}) = \int \frac{d^3 \mathbf{k}}{(2\pi)^3} a_{\mathbf{k}} u_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* e^{-i\mathbf{k} \cdot \mathbf{x}}$$

$$u_{\mathbf{k}}'' = -\omega_{\mathbf{k}}^2(\eta) u_{\mathbf{k}} \quad u_{\mathbf{k}}(\eta)$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^\dagger e^{-i k x}$$

$$U_k'' = -\omega_k^2(\eta) U_k \quad U_k(\eta)$$

Normalized modes

$$\int \frac{d^3k}{(2\pi)^3} a_k u_k e^{i\mathbf{k}\cdot\mathbf{r}} + a_k^\dagger u_k^* e^{-i\mathbf{k}\cdot\mathbf{r}}$$

$$u_k'' = -\omega_k^2(\eta) u_k$$

$$u_k(\eta)$$

$$' = \frac{d}{d\eta}$$

$$W(u_k, u_k') = u_k^* u_k' - u_k'^* u_k$$

$$\frac{dW}{d\eta} = u_k'^* u_k' + u_k^* u_k'' - u_k''^* u_k - u_k'^* u_k$$

$$\int \frac{d^3k}{(2\pi)^3} \alpha_k u_k e^{i\mathbf{k}\cdot\mathbf{x}} + \alpha_k^\dagger u_k^* e^{-i\mathbf{k}\cdot\mathbf{x}}$$

$$-\omega_k^2 u_k$$

$$u_k(\eta)$$

$$' = \frac{d}{d\eta}$$

Wronskian

$$W(u_k, u_k') = u_k^* u_k' - u_k'^* u_k$$

$$= \cancel{u_k'^* u_k'} + u_k^* u_k'' - u_k''^* u_k - \cancel{u_k'^* u_k'}$$

$$= -\omega_k^2 u_k^* u_k + \omega_k^2 u_k'^* u_k = 0$$

$$\int \frac{d^3k}{(2\pi)^3} \alpha_k u_k e^{ikx} + \alpha_k^\dagger u_k^\dagger e^{-ikx}$$

$$u_k'' = -\omega_k^2(\eta) u_k \quad u_k(\eta) \quad ' = \frac{d}{d\eta}$$

Wronskian

adiabatic

$$W(u_k, u_k') = u_k^* u_k' - u_k'^* u_k$$

$$\begin{aligned} \frac{dW}{d\eta} &= \cancel{u_k'^*} u_k'' + u_k^* u_k'' - u_k'^* u_k - \cancel{u_k'^*} u_k' \\ &= -\omega_k^2 u_k^* u_k + \omega_k^2 u_k^* u_k = 0 \end{aligned}$$

$$iW = 1$$

$$W = -i$$

$$u_k^{i*} u_k$$

$$u_k^{i*} u_k - u_k^{i*} / u_k^i$$

$$u_k^{i*} u_k = 0$$

$$iW = 1$$

$$W = -i$$

$$u_k = \frac{1}{\sqrt{2\omega}} e^{-i\omega\eta}$$

u_k

$$u_k - u_k^{i\eta} / u_k^i$$

$$k = 0$$

$$iW = 1$$

$$W = -i$$

$$u_k = \frac{1}{\sqrt{2\omega}} e^{-i\omega\eta}$$

$$W = \frac{1}{\sqrt{2\omega}} e^{+i\omega\eta} - i\omega \frac{1}{\sqrt{2\omega}} e^{-i\omega\eta} \times 2$$

$$= \frac{1}{2\omega} \times (2\omega - i\omega) = -i$$

u_k

$$u_k - u_k' / u_k'$$

$$k = 0$$

$$iW = 1$$

$$W = -i$$

$$u_k = \frac{1}{\sqrt{2\omega}} e^{-i\omega\eta}$$

$$W = \frac{1}{\sqrt{2\omega}} e^{+i\omega\eta} - i\omega \frac{1}{\sqrt{2\omega}} e^{-i\omega\eta} \times 2$$

$$= \frac{1}{2\omega} \times (2\omega - i\omega) = -i$$

u_k

$$u_k - u_k' / u_k'$$

$$k = 0$$

wavefunction of a single particle -

Hilbert space

- inner product

$$\langle \alpha | \beta \rangle$$

Wavefunction of a single particle ~

Hilbert space

~ inner product

$$(\langle \alpha | \beta \rangle)^* = \langle \beta | \alpha \rangle$$

Wavefunction of a single particle ~

Hilbert space .

~ inner product

$$(\langle \alpha | \beta \rangle)^* = \langle \beta | \alpha \rangle$$

Wavefunction of a single particle ~

Hilbert space

inner product

$$(\langle \alpha | \beta \rangle)^* = \langle \beta | \alpha \rangle$$

Original scalar on arbitrary spacetime:

$$\square \phi = m^2 \phi$$

$$\square = \frac{1}{\sqrt{|g|}} \partial_\mu (\sqrt{|g|} g^{\mu\nu} \partial_\nu \phi)$$

Wavefunction of a single particle ~

Hilbert space

inner product

$$(\langle \alpha | \beta \rangle)^* = \langle \beta | \alpha \rangle$$

Original scalar on arbitrary spacetime:

$$\square \phi = m^2 \phi$$

$$\square = \frac{1}{\sqrt{|g|}} \partial_\mu (\sqrt{|g|} g^{\mu\nu} \partial_\nu \phi)$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_{\mathbf{k}} u_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* e^{-i\mathbf{k} \cdot \mathbf{x}}$$

$$\nabla_\mu J^\mu = 0$$

$$J_\mu(\phi_1, \phi_2) = -i \left[\phi_2 \partial_\mu \phi_1 - (\partial_\mu \phi_2) \phi_1 \right]$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\nabla_\mu J^\mu = 0$$

$$J_\mu(\phi_1, \phi_2) = -i \left[\phi_2 \partial_\mu \phi_1 - (\partial_\mu \phi_2) \phi_1 \right]$$

Klein-Gordon inner product

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\nabla_\mu J^\mu = 0$$

$$J_\mu(\phi_1, \phi_2) = -i \left[\phi_2 \partial_\mu \phi_1 - (\partial_\mu \phi_2) \phi_1 \right]$$

Klein-Gordon inner product

$$\nabla^\mu J_\mu = -i$$

$$W = -i$$

$$u = \frac{ie}{\hbar\omega}$$

$$\nabla^\mu J_\mu = -i \left[\nabla^\mu \phi_2 \nabla_\mu \phi_1 + \phi_2 \square \phi_1 - \dots \right]$$

$$\nabla^\mu J_\mu = -i \left[\nabla^\mu \phi_2 \nabla_\mu \phi_1 + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \nabla_\mu \phi_2 \nabla^\mu \phi_1 \right]$$

$$\nabla^\mu J_\mu = -i \left[\cancel{\nabla^\mu \phi_2 \nabla_\mu \phi_1} + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \cancel{\nabla_\mu \phi_2 \nabla^\mu \phi_1} \right]$$

$$= -i \left[\right]$$

$$u = \frac{1}{\omega} e$$

$$\nabla^\mu J_\mu = -i \left[\cancel{\nabla^\mu \phi_2 \nabla_\mu \phi_1} + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \cancel{\nabla_\mu \phi_2 \nabla^\mu \phi_1} \right]$$

$$= -i \left[m^2 \phi_2 \phi_1 - m^2 \phi_1 \phi_1 \right] = 0$$

$$\nabla^\mu J_\mu = -i \left[\cancel{\nabla^\mu \phi_2 \nabla_\mu \phi_1} + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \cancel{\nabla_\mu \phi_2 \nabla^\mu \phi_1} \right]$$

$$= -i \left[m^2 \phi_2 \phi_1 - m^2 \phi_1 \phi_1 \right] = 0$$

$$J_\mu(\phi_1, \phi_2) = -J_\mu(\phi_2, \phi_1)$$

$$\nabla^\mu J_\mu = -i \left[\cancel{\nabla^\mu \phi_2 \nabla_\mu \phi_1} + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \cancel{\nabla_\mu \phi_2 \nabla^\mu \phi_1} \right]$$

$$= -i \left[m^2 \phi_2 \phi_1 - m^2 \phi_1 \phi_1 \right] = 0$$

$$J_\mu(\phi_1, \phi_2) = -J_\mu(\phi_2, \phi_1)$$

Problem

Defn inner product on Hilbert space

$$\langle \alpha | \alpha \rangle \geq 0$$

$$\nabla^\mu J_\mu = -i \left[\cancel{\nabla^\mu \phi_2 \nabla_\mu \phi_1} + \phi_2 \square \phi_1 - \square \phi_2 \phi_1 - \cancel{\nabla_\mu \phi_2 \nabla^\mu \phi_1} \right]$$

$$= -i \left[m^2 \phi_2 \phi_1 - m^2 \phi_1 \phi_1 \right] = 0$$

$$J_\mu(\phi_1, \phi_2) = -J_\mu(\phi_2, \phi_1)$$

Problem

Defn inner product on Hilbert space

$$\langle \alpha | \alpha \rangle \geq 0$$

$$J_\mu(\phi_1, \phi_1) = 0$$

$$\square \phi = m^2 \phi$$

ϕ real field

$$\square \phi = m^2 \phi$$

ϕ real field

$$\phi = \phi_+ + \phi_-$$

positive norm

negative norm

$$\phi_- = (\phi_+)^{\dagger} S$$

$$V^+ \oplus V^-$$

$$\square\phi = m^2\phi$$

ϕ real field

$$\phi = \phi_+ + \phi_-$$

$$\phi_- = (\phi_+)^* S = V^+ \oplus V^-$$

positive norm

negative norm

freedom to split into

complex solutions

\equiv

freedom in choice of
vacuum state

\equiv choice of left particles.

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} a_k u_k e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^* e^{-i\mathbf{k}\cdot\mathbf{x}}$$

$$\nabla_\mu J^\mu = 0$$

$$(\phi_1, \phi_2) = -i \left[\phi_2 \partial_\mu \phi_1 - (\partial_\mu \phi_2) \phi_1 \right]$$

$$G_\mu(\phi_1, \phi_2) = -i \left[\phi_2^- \partial_\mu \phi_1^+ - \partial_\mu \phi_2^- \phi_1^+ \right]$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\nabla_\mu J^\mu = 0$$

$$J_\mu(\phi_1, \phi_2) = -i \left[\phi_2 \partial_\mu \phi_1 - (\partial_\mu \phi_2) \phi_1 \right]$$

$$G_\mu(\phi_1, \phi_2) = -i \left[\phi_2^- \partial_\mu \phi_1^+ - \partial_\mu \phi_2^- \phi_1^+ \right]$$

ϕ_1^+ is really wavefunction of single particle state

$$\phi_1^- = (\phi_1^+)^*$$

ϕ_1^+ is really wavefunction of single particle state

$$\phi_1^- = (\phi_1^+)^*$$

~ complex conjugate of wavefunction

4

ψ^*

ϕ_1^+ is really wavefunction of single particle state ψ

$\phi_1^- = (\phi_1^+)^*$ \sim complex conjugate of wavefunction ψ^*

$$G_{11} = -i \left[\psi_2^* \partial_x \psi_1 - \partial_x \psi_2^* \psi_1 \right]$$

ϕ_1^+ is really wavefunction of single particle state ψ

$\phi_1^- = (\phi_1^+)^*$ \sim complex conjugate of wavefunction ψ^*

$$G_{\mu\nu} = -i \left[\psi_2^* \partial_\mu \psi_1 - \partial_\mu \psi_2^* \psi_1 \right]$$

$$\nabla_\mu J^\mu = 0$$

ϕ_1^+ is really wavefunction of single particle state ψ

$\phi_1^- = (\phi_1^+)^*$ \sim complex conjugate of wavefunction ψ^*

$$G_{\mu\nu} = -i \left[\psi_2^* \partial_\mu \psi_1 - \partial_\mu \psi_2^* \psi_1 \right]$$

$$\nabla_\mu J^\mu = 0 \quad \frac{1}{\sqrt{-g}} \partial_\mu (\sqrt{-g} g^{\mu\nu} J_\nu) = 0$$

ϕ_1^+ is really wavefunction of single particle state ψ

$\phi_1^- = (\phi_1^+)^*$ \sim complex conjugate of wavefunction ψ^*

$$G_{\mu\nu} = -i \left[\psi_2^* \partial_\mu \psi_1 - \partial_\mu \psi_2^* \psi_1 \right]$$

$$\nabla_\mu J^\mu = 0$$

$$\partial_\mu (\sqrt{-g} g^{\mu\nu} J_\nu) = 0$$

$$\int d^3x \partial_\mu \left[\sqrt{g} g^{\mu\nu} T_{\nu\sigma} \right] = 0$$

$$\partial_\mu [\sqrt{g} g^{\mu\nu} T_{\nu\sigma}] = 0$$

$$\partial = - \int d^3x \partial_i (\sqrt{g} g^{ir} g_r) = 0$$

$$\int d^3x \partial_\mu [\sqrt{g} g^{\mu\nu} T_{\nu\sigma}] = 0$$

$$\frac{\partial Q}{\partial t} = - \int d^3x \partial_i (\sqrt{g} g^{ir} G_r) = 0$$

$$\int d^3x \partial_r [\sqrt{g} g^{rv} T_{rv}] = 0$$

$$\frac{\partial Q}{\partial t} = - \int d^3x \partial_i (\sqrt{g} g^{ir} g_{rj}) = 0$$

$$Q = \sqrt{g} g^{0r} g_{rn}$$

$$\int d^3x \partial_\mu [\sqrt{-g} g^{\mu\nu} T_{\nu\sigma}] = 0$$

$$\frac{\partial Q}{\partial t} = - \int d^3x \partial_i (\sqrt{-g} g^{ir} G_r) = 0$$

$$Q = \sqrt{-g} g^{0\mu} G_\mu$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i\vec{k} \cdot \vec{x}} + a_k^\dagger u_k^* e^{-i\vec{k} \cdot \vec{x}}$$

$$\langle \gamma_2 | \gamma_1 \rangle$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\langle \psi_2 | \psi_1 \rangle = -i \int d^3 x \sqrt{-g} g^{0\mu} \left[\psi_2^* \partial_\mu \psi_1 - (\partial_\mu \psi_2^*) \psi_1 \right]$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\langle \psi_2 | \psi_1 \rangle = -i \int d^3 x \sqrt{-g} g^{\mu\nu} \left[\psi_2^* \partial_\mu \psi_1 - (\partial_\mu \psi_2^*) \psi_1 \right]$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\langle \psi_2 | \psi_1 \rangle = -i \int d^3 x \sqrt{-g} g^{0\mu} \left[\psi_2^* \partial_\mu \psi_1 - (\partial_\mu \psi_2^*) \psi_1 \right]$$

not guaranteed positive definite

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i\vec{k}\cdot\vec{x}} + a_k^\dagger u_k^* e^{-i\vec{k}\cdot\vec{x}}$$

$$\langle \psi_2 | \psi_1 \rangle = -i \int d^3 x \sqrt{-g} g^{\mu\nu} \left[\psi_2^* \partial_\mu \psi_1 - (\partial_\mu \psi_2^*) \psi_1 \right]$$

not guaranteed positive definite

$$\phi = \phi^+ + \phi^- = \psi + \psi^*$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ikx} + a_k^\dagger u_k^* e^{-ikx}$$

$$\langle \psi_2 | \psi_1 \rangle = -i \int d^3 x \sqrt{-g} g^{\mu\nu} \left[\psi_2^* \partial_\mu \psi_1 - (\partial_\mu \psi_2^*) \psi_1 \right]$$

not guaranteed positive definite

$$\phi = \phi^+ + \phi^- = \psi + \psi^*$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

In present case - split into positive and negative frequency solutions

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

in present case - split into positive and negative frequency solutions

$$\phi = \sum e^{-i\omega t} + e^{i\omega t}$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

In present case - split into positive and negative frequency solutions

$$\phi = \sum \underbrace{e^{-i\omega t}} + e^{i\omega t}$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

In present case - split into positive and negative frequency solutions

$$\phi = \underbrace{\sum e^{-i\omega t}}_{\psi} + \underbrace{\sum e^{i\omega t}}_{\psi^*}$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

in present case - split into positive and negative frequency solutions

$$\phi = \underbrace{\sum e^{-i\omega t}}_{\psi} + \underbrace{\sum e^{i\omega t}}_{\psi^*}$$

$$\langle \psi_2 | \psi_1 \rangle = 2\omega \int d^3x \psi_2^* \psi_1$$

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

in present case - split into positive and negative frequency solutions

$$\phi = \underbrace{\sum e^{-i\omega t}}_{\psi} + \underbrace{\sum e^{i\omega t}}_{\psi^*}$$

$$\langle \psi_2 | \psi_1 \rangle = 2\omega \int d^3x \psi_2^* \psi_1$$

Positive

In Minkowski space

$$\langle \psi_2 | \psi_1 \rangle = i \int d^3x \left[\psi_2^* \partial_t \psi_1 - (\partial_t \psi_2^*) \psi_1 \right]$$

In present case - split into positive and negative frequency solutions

$$\phi = \underbrace{\sum e^{-i\omega t}}_{\psi} + \underbrace{\sum e^{i\omega t}}_{\psi^*}$$

$$\langle \psi_2 | \psi_1 \rangle = 2\omega \int d^3x \psi_2^* \psi_1$$

Positive

$$\langle \psi_2 | \psi_1 \rangle \geq 0$$

$$\int d^3x \partial_\mu \left[\sqrt{-g} g^{\mu\nu} \right] = 0$$

$$\square \phi = m^2 \phi$$

$$\square \phi^\pm = m^2 \phi^\pm$$

$$\frac{\partial Q}{\partial t} = -$$

$$= 0$$

$$Q = \int \sqrt{-g} g^{\mu\nu}$$



FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$$\phi = \frac{u}{a(\eta)}$$

FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$$\phi = \frac{u}{a(\eta)}$$

$$u = \psi + \psi^*$$

<

FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$$\phi = \frac{u}{a(\eta)}$$

$$u = \psi + \psi^*$$

$$\langle \psi_2 | \psi_1 \rangle = +i \int d^3x \left(\frac{a^4}{a^2} \right) \left[\dots \right]$$

FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$\phi =$

$$u = \psi + \psi^*$$

$$\langle \chi_2 | + i \int d^3x \left(\frac{a^4}{a^2} \right) \left[\frac{\chi_2^*}{a} \partial_x \left(\frac{\chi_2}{a} \right) - \partial_\eta \left(\frac{\chi_2^*}{a} \right) \frac{\chi_2}{a} \right]$$

FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$$\phi = \frac{u}{a(\eta)}$$

$$u = \psi + \psi^*$$

$$\begin{aligned} \langle \psi_2 | \psi_1 \rangle &= +i \int d^3x \left(\frac{a^4}{a^2} \right) \left[\frac{\psi_2^*}{a} \partial_\eta \left(\frac{\psi_1}{a} \right) - \partial_\eta \left(\frac{\psi_2^*}{a} \right) \frac{\psi_1}{a} \right] \\ &= i \int d^3x \left[\psi_2^* \partial_\eta \psi_1 - \partial_\eta \psi_2^* \psi_1 \right] \end{aligned}$$

FRW

$$ds^2 = a^2(-d\eta^2 + d\vec{x}^2)$$

$$\phi = \frac{u}{a(\eta)}$$

$$u = \psi + \psi^*$$

$$\begin{aligned} \langle \psi_2 | \psi_1 \rangle &= +i \int d^3x \left(\frac{a^4}{a^2} \right) \left[\frac{\psi_2^*}{a} \partial_x \left(\frac{\psi_1}{a} \right) - \partial_x \left(\frac{\psi_2^*}{a} \right) \frac{\psi_1}{a} \right] \\ &= \int d^3x \left[\psi_2^* \partial_x \psi_1 - \partial_x \psi_2^* \psi_1 \right] \end{aligned}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k \cdot x} + a_k^\dagger u_k^* e^{-i k \cdot x}$$

$$\psi_{1,2} = u_{1,2 k} e^{i k \cdot x}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k \cdot x} + a_k^\dagger u_k^* e^{-i k \cdot x}$$

$$\psi_{1,2} = u_{1,2,k} e^{i k \cdot x}$$

$$\langle k_2 | k_1 \rangle = i \int d^3 x e^{i k_1 \cdot x}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ik \cdot x} + a_k^\dagger u_k^* e^{-ik \cdot x}$$

$$u_{1,2k} e^{ik \cdot x}$$

$$\langle \dots \rangle = i \int d^3 x e^{-ik_2 \cdot x} e^{ik_1 \cdot x} \left[u_{2k_2}^* \partial_\eta u_{1k_1} - \partial_\eta u_{2k_2}^\dagger u_{1k_1} \right]$$

= (

$$\langle \eta, x \rangle = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k \cdot x} + a_k^\dagger u_k^* e^{-i k \cdot x}$$

$$\psi_{1,2} = u_{1,2 k} e^{i k \cdot x}$$

$$\langle \psi_1, \psi_2 \rangle = i \int d^3 x e^{-i k_2 \cdot x} e^{i k_1 \cdot x} \left[u_{2 k_2}^* \partial_\mu u_{1 k_1} - \partial_\mu u_{2 k_2}^* u_{1 k_1} \right]$$

$$= (2\pi)^3 \delta^{(3)}(k_1 - k_2) i W$$

$$\langle \eta, x \rangle = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ik \cdot x} + a_k^\dagger u_k^* e^{-ik \cdot x}$$

$$\psi_{1,2} = u_{1,2}(k) e^{ik \cdot x}$$

$$\langle k_2 | k_1 \rangle = i \int d^3 x e^{-ik_2 \cdot x} e^{ik_1 \cdot x} \left[u_{2k_2}^\dagger \partial_\mu u_{1k_1} - \partial_\mu u_{2k_2}^\dagger u_{1k_1} \right]$$

$$= (2\pi)^3 \delta^{(3)}(k_1 - k_2) i W$$

$$\boxed{\langle k_2 | k_1 \rangle = (2\pi)^3 \delta^{(3)}(k_1 - k_2)}$$

$$\langle \eta, x \rangle = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k \cdot x} + a_k^\dagger u_k^\dagger e^{-i k \cdot x}$$

$$\psi_{1,2} = u_{1,2 k} e^{i k \cdot x}$$

$$\langle k_2 | k_1 \rangle = i \int d^3 x e^{-i k_2 \cdot x} e^{i k_1 \cdot x} \left[u_{2 k_2}^\dagger \partial_\mu u_{1 k_1} - \partial_\mu u_{2 k_2}^\dagger u_{1 k_1} \right]$$

$$= (2\pi)^3 \delta^{(3)}(k_1 - k_2) i W$$

$$\boxed{\langle k_2 | k_1 \rangle = (2\pi)^3 \delta^{(3)}(k_1 - k_2)}$$

$$\langle \eta, x \rangle = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ikx} + a_k^\dagger u_k^* e^{-ikx}$$

$$\psi_{1,2} = u_{1,2}(k) e^{ikx}$$

$$\langle k_2 | k_1 \rangle = i \int d^3 x e^{-ik_2 \cdot x} e^{ik_1 \cdot x} \left[u_{2k_2}^* \partial_\mu u_{1\mu} - \partial_\mu u_{2k_2}^* u_{1\mu} \right]$$

$$= (2\pi)^3 \delta^{(3)}(k_1 - k_2) i W$$

$$\boxed{\langle k_2 | k_1 \rangle = (2\pi)^3 \delta^{(3)}(k_1 - k_2)}$$

$$[a_k, a_n] = 0$$

$$[a_k^\dagger, a_n^\dagger] = 0$$

$$[a_k^\dagger, a_n] = (2\pi)^3 \delta^{(3)}(\mathbf{k} - \mathbf{n})$$

$$\left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \omega_{\mathbf{k}}$$

$$[a_k, a_{k'}] = 0$$

$$[a_k^+, a_{k'}^+] = 0$$

$$[a_k^+, a_{k'}] = (2\pi)^3 \delta^{(3)}(k - k')$$

$u_{\mathbf{k}}$

$$[a_k, a_{k'}] = 0$$

$$[a_k^\dagger, a_{k'}^\dagger] = 0$$

$$[a_k^\dagger, a_{k'}] = (2\pi)^3 \delta^{(3)}(k - k')$$

$$[u(\eta, x), u(\eta, x')] = 0$$

$$[p, \dots]$$

$$[a_k, x_{k'}] = 0$$

$$[a^{\dagger}_k, a^{\dagger}_{k'}] = 0$$

$$[a^{\dagger}_k, a_{k'}] = (\pi)^3 \delta^{(3)}(k - k')$$

$$[u(\eta, x), u(\eta, x')] = 0$$

$$[P(\eta, x), P(\eta, x')] = 0$$

$$P = u'$$

$$[u$$

$u \geq k,$

$$[a_k, x_{k'}] = 0$$

$$[a_k^+, a_{k'}^+] = 0$$

$$[a_k^+, a_{k'}] = (2\pi)^3 \delta^{(3)}(k - k')$$

$$[u(\eta, x), u(\eta, x')] = 0$$

$$[P(\eta, x), P(\eta, x')] = 0 \quad P = u'$$

$$[u(\eta, x), P(\eta, x')] = i \delta^{(3)}(x - x')$$

$$[a_k, x_{k'}] = 0$$

$$[a_k^+, a_{k'}^+] = 0$$

$$[a_k^+, a_{k'}] = (2\pi)^3 \delta^{(3)}(k - k')$$

$$[u(\eta, x), u(\eta, x')] = 0$$

$$[P(\eta, x), P(\eta, x')] = 0 \quad P = u'$$

$$[u(\eta, x), P(\eta, x')] = i \delta^{(3)}(x - x')$$

$$[W(x), P(y)] = \int \frac{d^3 k_1}{(2\pi)^3} \frac{d^3 k_2}{(2\pi)^3}$$

$$\left[\begin{array}{c} \gamma_1 \\ -\gamma_2 \\ \gamma_2 \\ \gamma_1 \end{array} \right]$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ikx} + a_k^\dagger u_k^* e^{-ikx}$$

$$\hat{\phi}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{ikx} + a_k^\dagger u_k^* e^{-ikx}$$

$u_{1,2} =$

$$\langle k_2 | k_1 \rangle = \int d^3 x e^{-ik_2 x} [u_{2k_2}^*]_{\eta}$$

$$(k_1 - k_2) \cdot W$$

$(k_1 - k_2)$

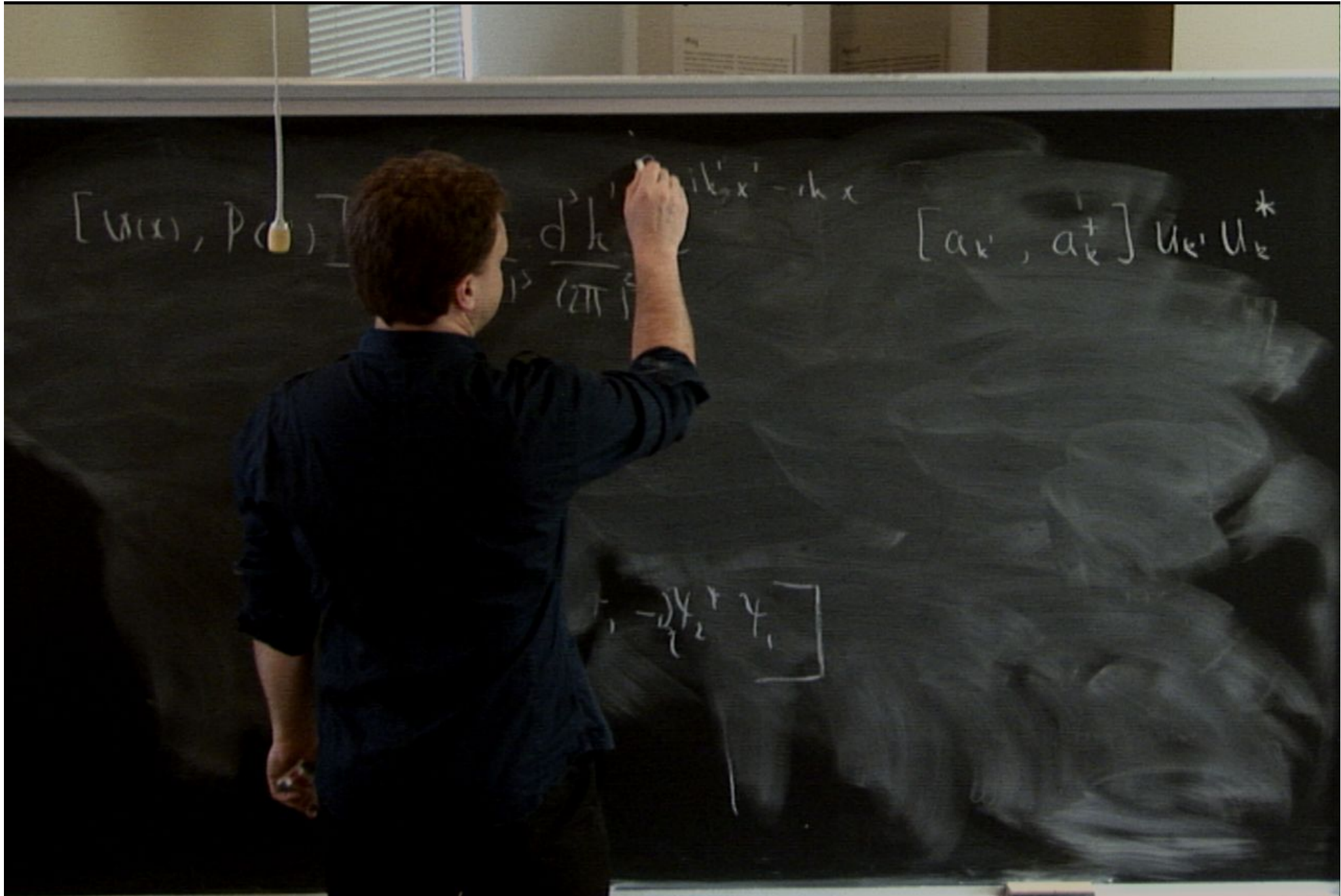
$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$\hat{\phi}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x}$$

$$= i \int d^3 x e^{-i k_2 x} e^{i k_1 x} \left[u_{2k_2}^* \partial_\eta \right]$$

$$= (2\pi)^3 \delta^{(3)}(k_1 - k_2) i W$$

$$\langle k_2 | k_1 \rangle = (2\pi)^3 \delta^{(3)}(k_1 - k_2)$$



$$[W(x), P(x)]$$

$$\frac{d}{dx} i\hbar x - \hbar x$$

$$[a_k, a_k^\dagger] U_k U_k^*$$

$$-\frac{1}{2} \psi_2 + \psi_1$$

$$[W(x), P(x')] = \int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} e^{ik'x' - ikx} [a_k, a_{k'}^*] u_{k'} u_k^*$$

$$= \left[-\frac{1}{2} \gamma_2 + \gamma_1 \right]$$

$[\psi(x), \psi(x)]$

$$\int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3}$$

$$\left[\begin{array}{l} e^{ik'x' - ikx} \\ + e^{-ik'x' + ikx} \end{array} \right]$$

$$[a_{k'}, a_{k'}^+] \psi_{k'} \psi_{k'}^*$$

$$[a_{k'}^+, a_{k'}] \psi_{k'}^* \psi_{k'}$$

$$\left[\begin{array}{l} -\psi_2^* \psi_1 \\ \psi_1 \end{array} \right]$$

$$\begin{aligned}
 [P(x)] = & \int \frac{d^3h}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \left[e^{ik'_x x' - ik_x x} [a_{k'}, a_{k'}^+] u_{k'} u_{k'}^* \right. \\
 & \left. + e^{-ik'_x x' + ik_x x} [a_{k'}^+, a_{k'}] u_{k'}^* u_{k'} \right]
 \end{aligned}$$

$$\begin{aligned}
 [u(x), p(x')] &= \int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \left[e^{ik'x' - ikx} [a_{k'}, a_k^\dagger] u_{k'} u_k^{1*} \right. \\
 &\quad \left. + e^{-ik'x' + ikx} [a_{k'}^\dagger, a_k] u_{k'}^* u_k' \right] \\
 &= \int \frac{d^3k}{(2\pi)^3} e^{ik(x'-x)} [u_k u_k^{1\dagger} - u_k^\dagger u_k']
 \end{aligned}$$

$$[u(x), p(x')] = \int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \left[e^{ik'x' - ikx} [a_{k'}, a_k^\dagger] u_{k'} u_k^{*\dagger} \right.$$

$$\left. + e^{-ik'x' + ikx} [a_{k'}^\dagger, a_k] u_{k'}^* u_k' \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{ik(x-x')} [u_k u_k^{*\dagger} - u_k^* u_k']$$

$$= \underline{i \delta^{(3)}(x-x')}$$

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} \left[a_k u_k e^{i\vec{k}\cdot\vec{x}} + a_k^\dagger u_k^* e^{-i\vec{k}\cdot\vec{x}} \right]$$

Well-defined
out state



Well-defined
in-state

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} \alpha_k u_k e^{i\vec{k}\cdot\vec{x}} + \alpha_k^\dagger u_k^* e^{-i\vec{k}\cdot\vec{x}}$$

Well-defined
out state



Well-defined
in-state

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} a_k u_k e^{i\vec{k}\cdot\vec{x}} + a_k^\dagger u_k^* e^{-i\vec{k}\cdot\vec{x}}$$

Well-defined
out state

$$\omega' \ll \omega^2$$

Vacuum

$$\omega' > \omega^2$$

Well-defined
in-state

$$\omega' \ll \omega^2$$

Vacuum

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} \alpha_k u_k e^{i\mathbf{k}\cdot\mathbf{x}} + \alpha_k^\dagger u_k^\dagger e^{-i\mathbf{k}\cdot\mathbf{x}}$$

Well-defined
cut state

$$\omega' \ll \omega^2$$

Vacuum

filled with particles

$$\omega' > \omega^2$$

Well-defined
n-state

$$\omega' \ll \omega^2$$

Vacuum

Wavefunctions, complex solutions

Wavefunctions, complex solutions

$$U_k'' = -\omega^2 U_k$$

Wavefunctions, complex solutions

$$U_k'' = -\omega^2 U_k$$

Wavefunctions, complex solutions

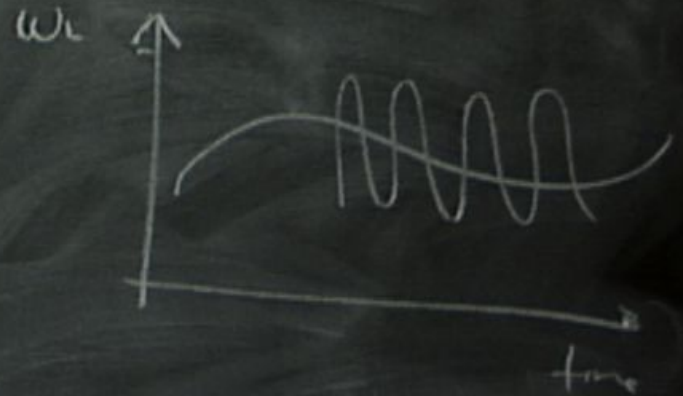
$$U_k'' = -\omega_k^2 U_k$$

U_k positive norm states

WKB

$$\omega_h^1 \ll \omega_k^2$$

adiabatic



Wavefunctions, complex solutions

$$U_k'' = -\omega_k^2 U_k$$

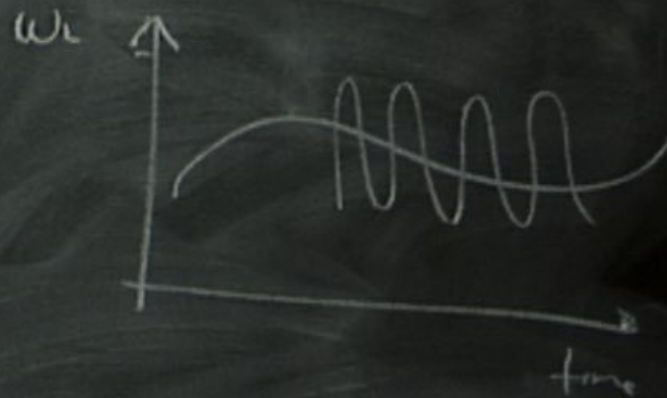
U_k positive norm states

WKB

$$\omega_h' \ll \omega_k^2$$

adiabatic

$$U_k = A e^{-i \int \omega_k(t) dt}$$



Wavefunctions, complex solutions

$$U_k'' = -\omega_k^2 U_k$$

U_k positive norm states

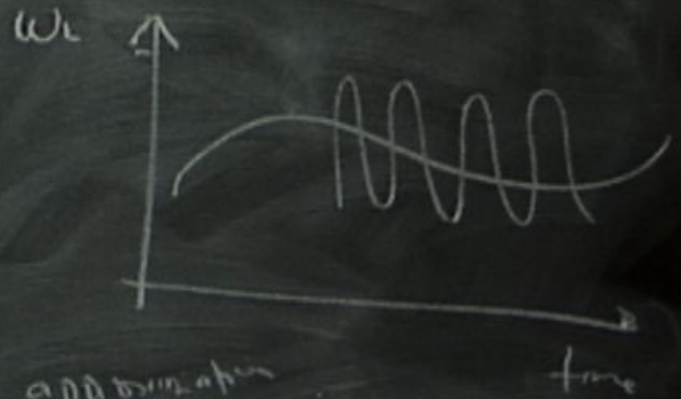
WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$U_k = A e^{-i \int \omega_k dt}$$

WKB approximation



Wavefunctions, complex solutions

$$U_k'' = -\omega_k^2 U_k$$

U_k positive norm states

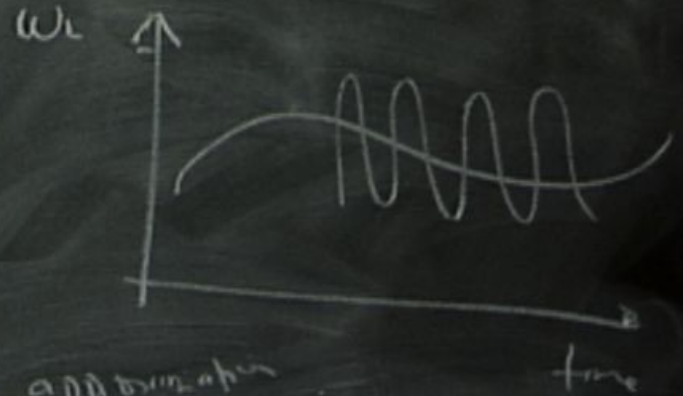
WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$U_k = A e^{-i \int \omega_k dt}$$

WKB approximation



Wavefunctions, complex solutions

$$U_k'' = -\omega_k^2 U_k$$

U_k positive norm states

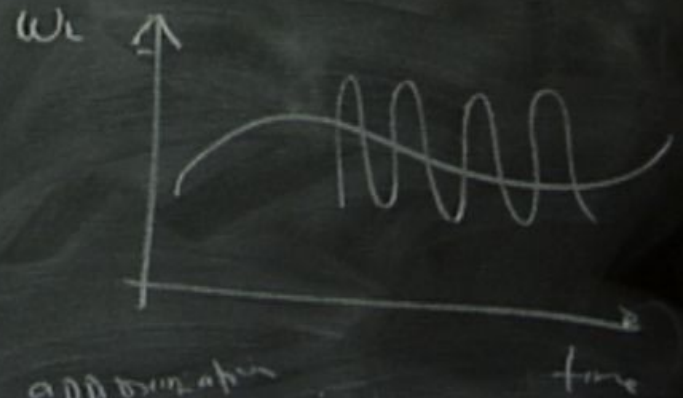
WKB

$$\omega_h' \ll \omega_k^2$$

adiabatic

$$U_k = A^{(1)} e^{-i \int \omega_k(t) dt}$$

WKB approximation



$$u_k^* e^{-ikx}$$

in particles ↓

Wronskian /
normalization

Wavefunctions, complex solutions

$$u_k'' = -\omega_k^2 u_k$$

u_k positive norm

WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$u_k = A^{(0)} e^{-i \int \omega_k(q) dq}$$

WKB

$$i [u_k^* u_k' - u_k'^* u_k] = 1$$

$$u_k^* e^{-ikx}$$

Wavefunctions, complex solutions

$$u_k'' = -\omega_k^2 u_k$$

u_k positive norm

WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$u_k = A e^{-i \int \omega_k dt}$$

WKB

Wronskian /
normalization

$$i [u_k^* u_k' - u_k'^* u_k] = 1$$

$$i |A|^2 2\omega_k$$

$$u_k^* e^{-ikx}$$

in particles ↓

Wronskian /
normalization

Wavefunctions, complex solutions

$$u_k'' = -\omega_k^2 u_k$$

u_k positive norm

WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$u_k = A e^{-i \int \omega_k dz}$$

WKB

$$i [u_k^* u_k' - u_k'^* u_k] = 1$$

$$-i |A|^2 2\omega_k = 1 \quad |A| = \frac{1}{\sqrt{2\omega_k}}$$

Wavefunctions, complex solutions

$$U_k'' = -\omega_k^2 U_k$$

U_k positive norm states

WKB

$$\omega_k' \ll \omega_k^2$$

adiabatic

$$U_k = \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt}$$

WKB approximation

Wronskian /
normalization

$$i [U_k^{*'} U_k' - U_k^{*'} U_k] = 1$$

$$-i |A|^2 2\omega_k = 1 \quad |A| = \frac{1}{\sqrt{2\omega_k}}$$

time



$$U_k = \frac{1}{\sqrt{1 - \beta^2}} \rho - \int \rho \alpha_k dz \quad \text{Vacuum.}$$

time



$$U_k = \frac{1}{\sqrt{v_{\text{avg}}}} \rho \int \sigma_{\text{avg}} dz \quad \text{Vacuum.}$$

time

↑

WKB approx is good

$$\omega' \ll \omega^2$$

WKB bad


$$U_k = \frac{1}{\sqrt{2\omega}} e^{-i\int \omega dz}$$

$$U_k = \frac{1}{\sqrt{2\omega}} e^{-i\int \omega dz}$$

for Atoms Molecule!

KB approx is good
 $c\omega^2$

KB bad

$$U_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i\int \omega_k dt} + \frac{\beta_k}{\sqrt{2\omega_k}} e^{+i\int \omega_k dt}$$


$$U_k = \frac{1}{\sqrt{2\omega_k}} e^{-i\int \omega_k dt} \text{ Vacuum.}$$

WKB approx is good
 $\omega' \ll \omega^2$

WKB

$$U_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \frac{\beta_k}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

created particles.

No particles created is proportional to $|\beta_k|^2$

$\int \omega_k dt$ Vacuum.

WKB approx is good
 $\omega' \ll \omega^2$

$$u_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

Bogoliubov co-efficients

created particles

No particles created is proportional to $|\beta_k|^2$

WKB bar

$$u_k = \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} \text{ Vacuum}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_{\mathbf{k}} u_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* e^{-i\mathbf{k} \cdot \mathbf{x}}$$

Norm/Normalization is preserved in time

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} \alpha_k U_k e^{i\vec{k}\cdot\vec{x}} + \alpha_k^* U_k^* e^{-i\vec{k}\cdot\vec{x}}$$

Norm/Normalization is preserved in time.

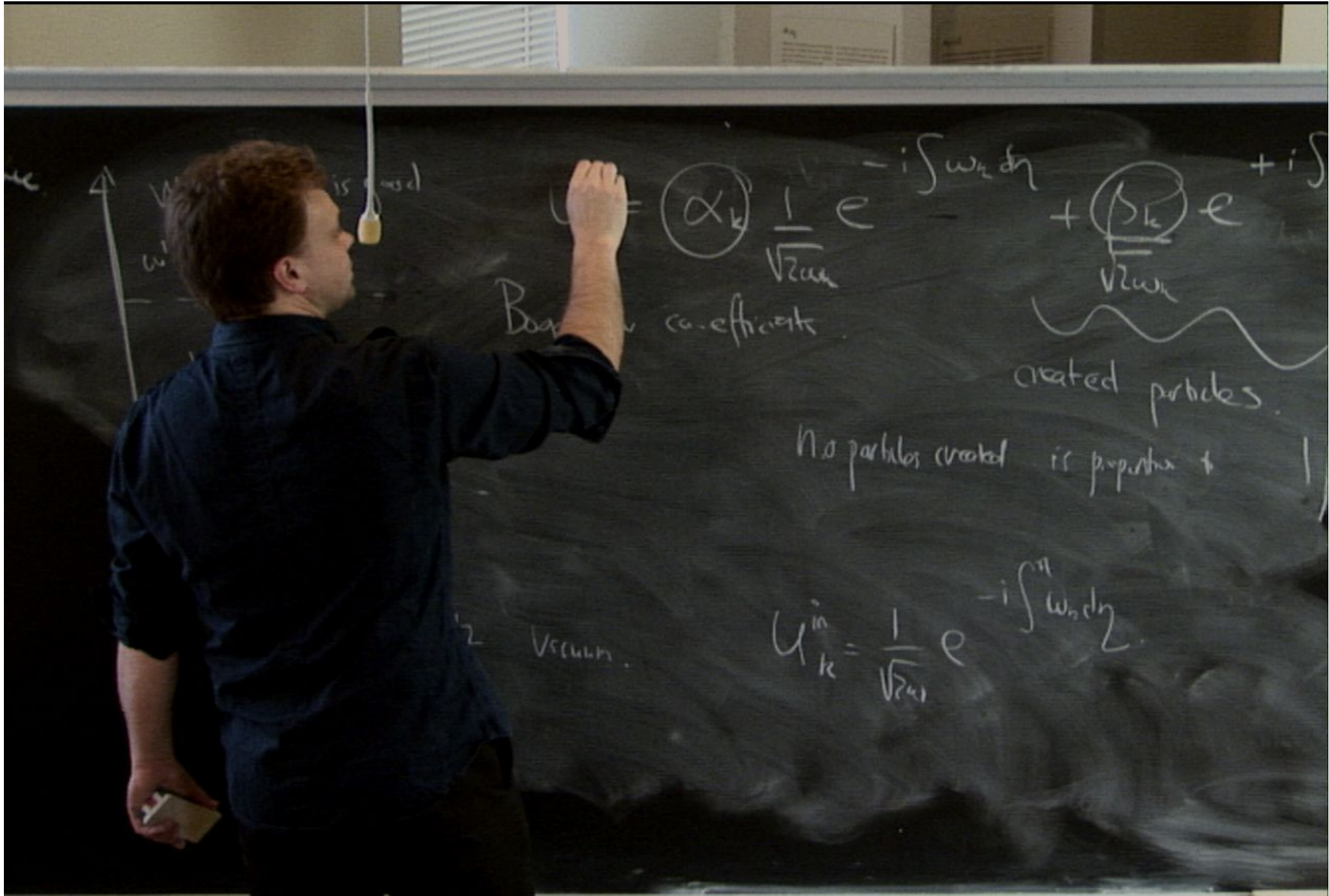
$$-i \left[U_k^* U_k' - U_k'^* U_k \right] = 1$$

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} a_k u_k e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^* e^{-i\mathbf{k}\cdot\mathbf{x}}$$

Norm/Normalization is preserved in time

$$-i \left[u_k^* u_k' - u_k'^* u_k \right] = 1$$

$$-i \left[\right]$$



$$U = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

α_k and β_k are Bogoliubov coefficients.
 β_k is related to created particles.

No particles created is proportional to $|\beta_k|^2$

vacuum

$$U_{in}^k = \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt}$$

Pollen to Evidence for Atoms

How Big Is A Molecule?

time ↑
 $\omega \ll \omega_c^2$
 WKB approx is good
 WKB bar

$$U_k^{\hat{}} = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

Boglubov coefficients

created particles

No particles created is proportional to $|\beta_k|^2$

$$U_k^{\hat{}} = \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} \quad \text{V vacuum}$$

$$U_k^{\text{in}} = \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt}$$

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} \alpha_k U_k e^{ikx} + \alpha_k^* U_k^* e^{-ikx}$$

Norm/Normalization is preserved in time.

$$-i \left[U_k^* U_k' - U_k'^* U_k \right] = 1$$

$$\lim_{\eta \rightarrow +\infty} U_k^{\text{in}}(\eta) = \alpha_k U_k^{\text{out}}(\eta) + \beta_k U_k^{\text{out}*}(\eta)$$

$$\hat{U}(\eta, x) = \int \frac{d^3k}{(2\pi)^3} \alpha_k U_k e^{i\vec{k}\cdot\vec{x}} + \alpha_k^* U_k^* e^{-i\vec{k}\cdot\vec{x}}$$

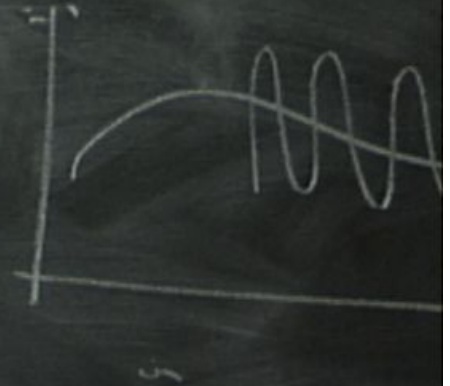
Norm/Normalization is preserved in time

$$-i \left[U_k^* U_k' - U_k'^* U_k \right] = 1$$

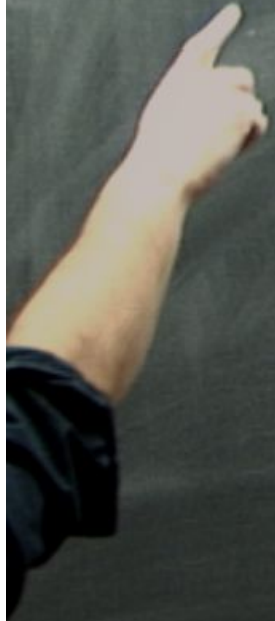
$$\lim_{\eta \rightarrow +\infty} U_k^{\text{in}}(\eta) = \alpha_k U_k^{\text{out}}(\eta) + \beta_k U_k^{\text{out}*}(\eta)$$

$$-i \left[(\alpha_k^* u_k^{out} + \beta_k^* u_k^{in}) (\alpha_k u_k^{in} + \beta_k u_k^{out}) \right]$$

$$-i \left[\left(\alpha_h^* U_h^{\text{in}+} + \beta_h^{\text{out}} U_h^{\text{out}+} \right) \left(\alpha_h U_h^{\text{in}+} + \beta_h U_h^{\text{out}+} \right) - \text{c.c} \right]$$

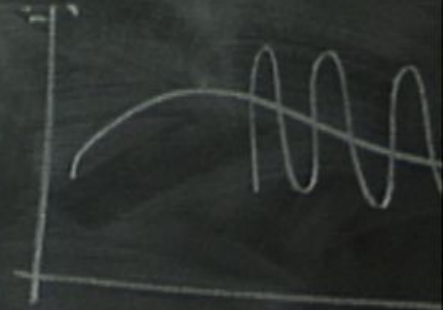


$$-i \left[\left(\alpha_k^* u_k^{in} + \beta_k^* u_k^{out} \right) \left(\alpha_k u_k^{in} + \beta_k u_k^{out} \right) - c.c \right]$$



$$-i \left[\left(\alpha_k^* U_k^{\text{in}+} + \beta_k^{\text{out}} U_k^{\text{out}+} \right) \left(\alpha_k U_k^{\text{in}+} + \beta_k U_k^{\text{out}+} \right) - \text{c.c} \right]$$

$$-i \left[|\alpha_k|^2 U_k^{\text{out}+} U_k^{\text{in}+} \right]$$



$$-i \left[\left(\alpha_k^* U_k^{out} + \beta_k^* U_k^{in} \right) \left(\alpha_k U_k^{in} + \beta_k U_k^{out} \right) - c.c. \right]$$

$$-i \left[|\alpha_k|^2 U_k^{out} U_k^{in} \right]$$



$$-i \left[\left(\alpha_k^* U_k^{out*} + \beta_k^* U_k^{out} \right) \left(\alpha_k U_k^{in} + \beta_k U_k^{in*} \right) - c.c \right]$$

$$-i \left[\begin{array}{l} |\alpha_k|^2 U_k^{out} U_k^{in} \\ -|\beta_k|^2 \end{array} \right]$$



$$-i \left[\left(\alpha_k^* U_k^{out} + \beta_k^* U_k^{in} \right) \left(\alpha_k U_k^{in} + \beta_k U_k^{out} \right) - c.c \right]$$

$$-i \left[\begin{pmatrix} |\alpha_k|^2 & \\ & -|\beta_k|^2 \end{pmatrix} \begin{pmatrix} U_k^{in} \\ U_k^{out} \end{pmatrix} - c.c \right]$$



$$-i \left[\left(\alpha_k^* U_k^{out*} + \beta_k^* U_k^{out} \right) \left(\alpha_k U_k^{in} + \beta_k U_k^{in*} \right) - c.c. \right]$$

$$-i \left[\begin{pmatrix} |\alpha_k|^2 & \\ & -|\beta_k|^2 \end{pmatrix} \begin{pmatrix} U_k^{out} \\ U_k^{in} \end{pmatrix} - c.c. \right]$$

$$= 1$$



$$-i \left[\left(\alpha_k^* U_k^{out*} + \beta_k^* U_k^{out} \right) \left(\alpha_k U_k^{in} + \beta_k U_k^{in*} \right) - c.c \right]$$

$$-i \left[\begin{pmatrix} |\alpha_k|^2 & U_k^{out} U_k^{in} \\ -|\beta_k|^2 & \end{pmatrix} - c.c \right]$$

$$= 1$$

$$|\alpha_k|^2 - |\beta_k|^2 = 1$$



$$-i \left[\left(\alpha_k^* U_k^{out} + \beta_k U_k^{in} \right) \left(\alpha_k U_k^{in} + \beta_k^* U_k^{out} \right) - c.c \right]$$

$$-i \left[\begin{pmatrix} |\alpha_k|^2 & \\ & -|\beta_k|^2 \end{pmatrix} U_k^{out} U_k^{in} - c.c \right]$$

$$= 1$$

$$|\alpha_k|^2 - |\beta_k|^2 = 1$$

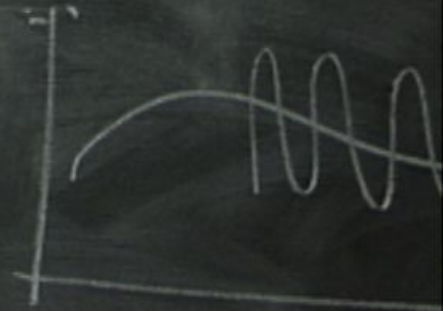
Unitarity / conservation of prob



$$-i \left[\left(\alpha_k^* U_k^{\text{out}^*} + \beta_k^{\text{out}} U_k^{\text{in}^*} \right) \left(\alpha_k U_k^{\text{out}} + \beta_k U_k^{\text{in}} \right) - \text{c.c} \right]$$

$$-i \left[\begin{pmatrix} |\alpha_k|^2 & \\ & -|\beta_k|^2 \end{pmatrix} U_k^{\text{out}^*} U_k^{\text{out}} - \text{c.c} \right]$$

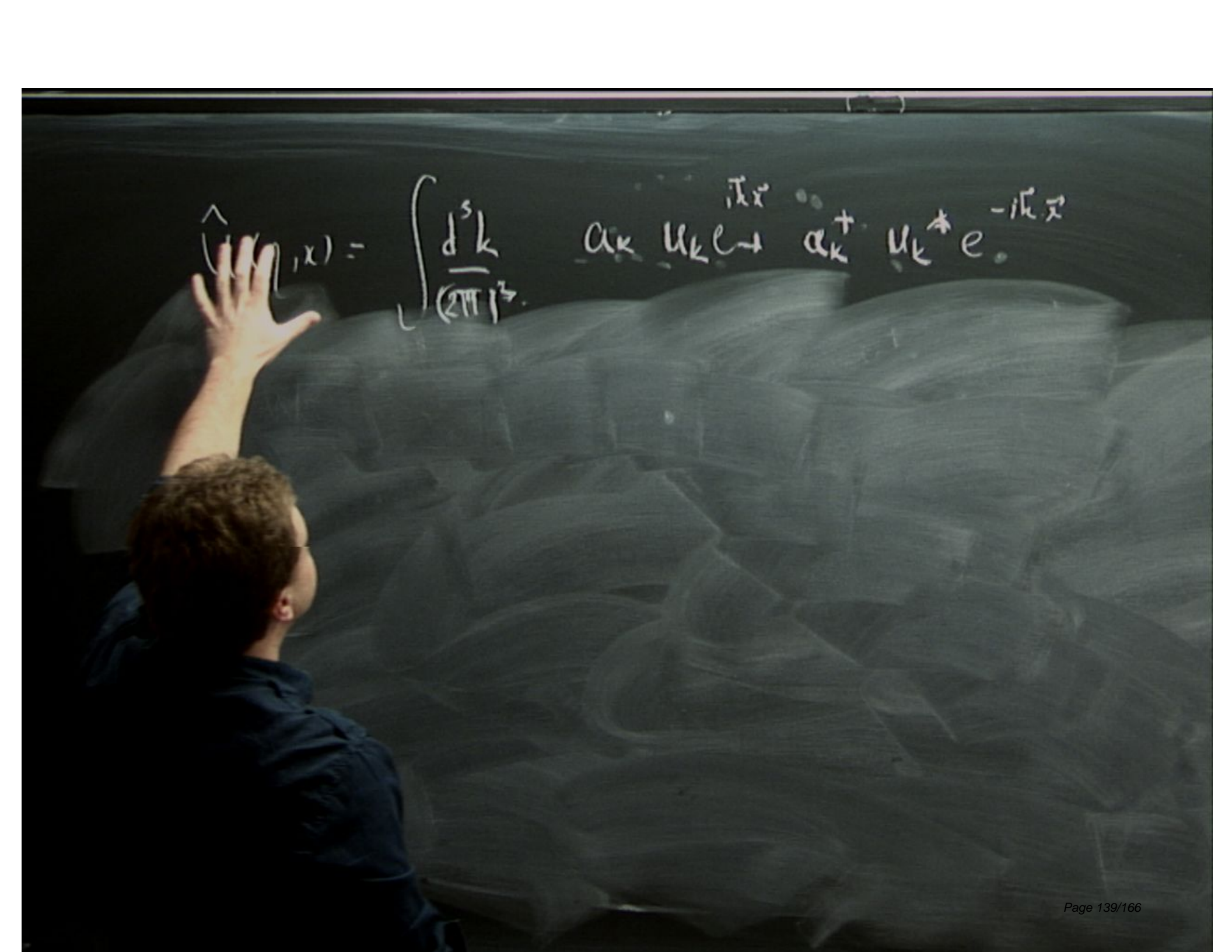
$$= 1$$



$$|\alpha_k|^2 - |\beta_k|^2 = 1$$

Unitarity / conservation of prob

$\alpha, \beta \sim \{ \}$ free constants function of k

A person is seen from the back, pointing their right hand towards a chalkboard. The chalkboard is dark and has a mathematical equation written in white chalk. The equation is for the expansion of a scalar field operator in terms of creation and annihilation operators. The person is wearing a dark blue shirt.
$$\hat{\psi}(t, \vec{x}) = \int \frac{d^3k}{(2\pi)^3} a_{\vec{k}} u_{\vec{k}} e^{-i\vec{k}\cdot\vec{x}} + a_{\vec{k}}^\dagger u_{\vec{k}}^* e^{-i\vec{k}\cdot\vec{x}}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_{\mathbf{k}} u_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x}} + a_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* e^{-i\mathbf{k}\cdot\mathbf{x}}$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{ikx} + a_k^\dagger u_k^{\text{in}*} e^{-ikx}$$

$$a_k | \text{in} \rangle = 0$$

in-vacuum

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\mathbf{k}\cdot\mathbf{x}}$$

$$a_k | \text{in} \rangle = 0 \quad \text{in-vacuum}$$

$$\lim_{\eta \rightarrow -\infty} u_k^{\text{in}}(\eta) =$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\mathbf{k}\cdot\mathbf{x}}$$

$$a_k | \text{in} \rangle = 0 \quad \text{in-vacuum}$$

$$\lim_{\eta \rightarrow +\infty} u_k^{\text{in}}(\eta) = \alpha_k u_k^{\text{out}}(\eta) + \beta_k u_k^{\text{out}*}(\eta)$$

$$* e^{-i\vec{k}\cdot\vec{x}}$$

=

$$\int \frac{d^3k}{(2\pi)^3} e^{i\vec{k}\cdot\vec{x}} \left[a_k u_k^{\text{in}} + a_{-k}^\dagger u_k^{\text{in}*} \right]$$

$$\left(a_k u_k^{\text{out}*} \right)$$

Unitary / conservation

constant function of k

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} u_{\mathbf{k}}^{\text{in}} + a_{-\mathbf{k}}^{\dagger} u_{\mathbf{k}}^{\text{in}*} \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} \left(\alpha_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + \beta_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}*} \right) + a_{-\mathbf{k}}^{\dagger} \left(\dots \right) \right]$$

Unitary

constant function of k

$$\int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} u_{\mathbf{k}}^{\text{in}} + a_{-\mathbf{k}}^{\dagger} u_{\mathbf{k}}^{\text{in}*} \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} \left(\alpha_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + \beta_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}*} \right) + a_{-\mathbf{k}}^{\dagger} \left(\alpha_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}} + \beta_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}*} \right) \right]$$

Unitary

constants function of k

$$\int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} u_{\mathbf{k}}^{\text{in}} + a_{-\mathbf{k}}^{\dagger} u_{\mathbf{k}}^{\text{in}*} \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} \left(\alpha_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + \beta_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}*} \right) + a_{-\mathbf{k}}^{\dagger} \left(\alpha_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}} + \beta_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}*} \right) \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[b_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + b_{-\mathbf{k}}^{\dagger} u_{\mathbf{k}}^{\text{out}*} \right]$$

$$\int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} u_{\mathbf{k}}^{\text{in}} + a_{-\mathbf{k}}^+ u_{\mathbf{k}}^{\text{in}*} \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[a_{\mathbf{k}} \left(\alpha_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + \beta_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}*} \right) + a_{-\mathbf{k}}^+ \left(\alpha_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}} + \beta_{-\mathbf{k}}^* u_{\mathbf{k}}^{\text{out}*} \right) \right]$$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} \left[b_{\mathbf{k}} u_{\mathbf{k}}^{\text{out}} + b_{-\mathbf{k}}^+ u_{\mathbf{k}}^{\text{out}*} \right]$$

$$b_{\mathbf{k}} = \alpha_{\mathbf{k}} a_{\mathbf{k}} + \beta_{\mathbf{k}}^* a_{-\mathbf{k}}^+$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\mathbf{k}\cdot\mathbf{x}} =$$

$$a_k | \text{in} \rangle = 0 \quad \text{in-vacuum}$$

$$\lim_{\eta \rightarrow +\infty} u_k^{\text{in}}(\eta) = \alpha_k u_k^{\text{out}}(\eta) + \beta_k u_k^{\text{out}*}(\eta)$$

$$b_k | \text{out} \rangle = 0 \quad b_k, b_k^\dagger \quad \text{annihilation + creation operators}$$

def particles $b_k^\dagger + 0$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\mathbf{k}\cdot\mathbf{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\mathbf{k}\cdot\mathbf{x}} =$$

$$a_k | \text{in} \rangle = 0 \quad \text{in-vacuum}$$

$$\lim_{\eta \rightarrow +\infty} u_k^{\text{in}}(\eta) = \alpha_k u_k^{\text{out}}(\eta) + \beta_k u_k^{\text{out}*}(\eta)$$

$$b_k | \text{out} \rangle = 0 \quad b_k, b_k^\dagger \text{ annihilator + creator operators}$$

def particles at $t=0$

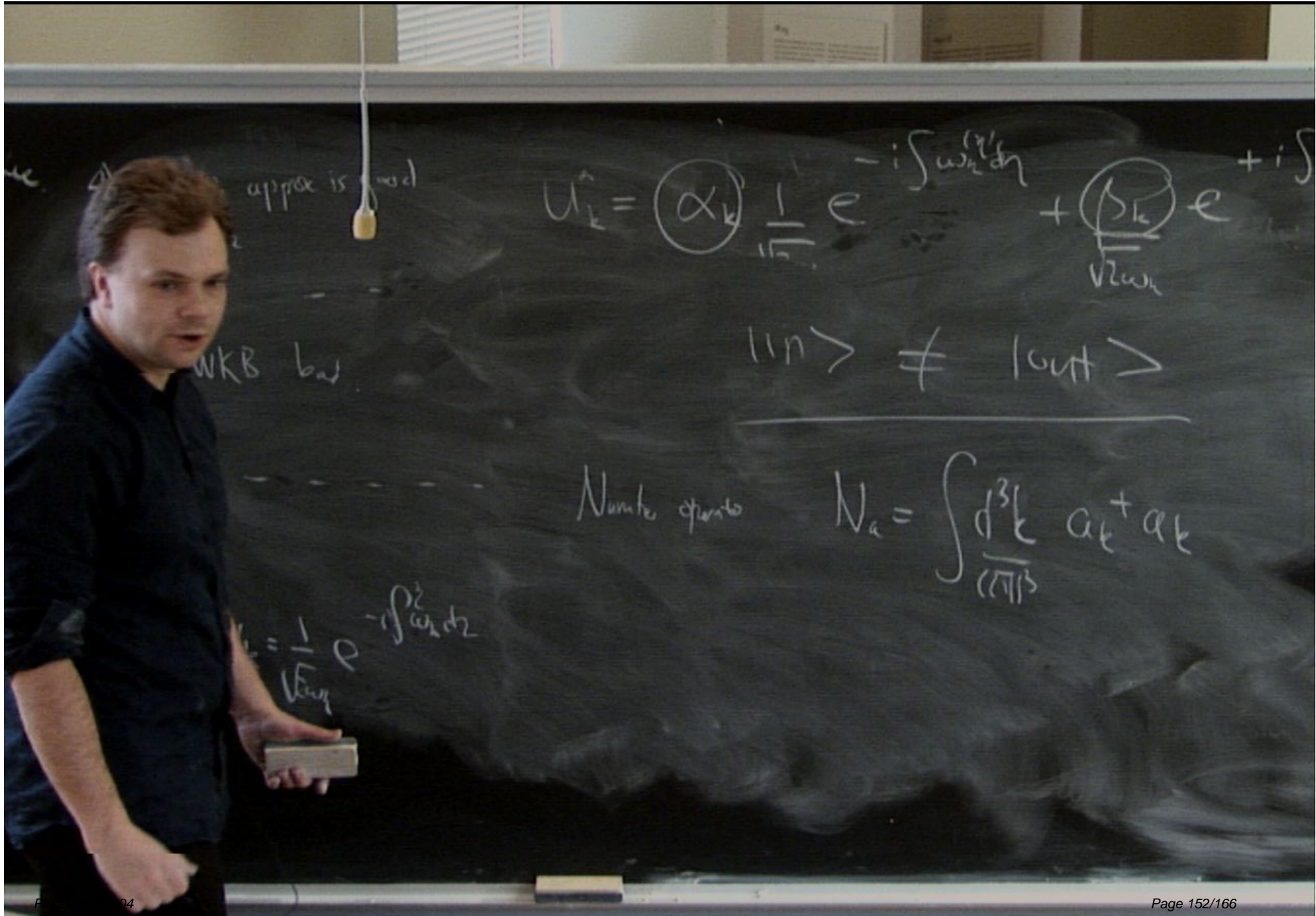
WKB approx is

$$\omega' \ll \omega^2$$

WKB 1

$$U_k^{\wedge} = \underbrace{(\alpha_k)}_{\frac{1}{\sqrt{v}} \frac{1}{\omega_k}} e^{-i \int \omega_k dt} + \underbrace{(\beta_k)}_{\frac{1}{\sqrt{v}} \frac{1}{\omega_k}} e^{+i \int \omega_k dt}$$

$$|in\rangle \neq |out\rangle$$



approx is good

$$U_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

WKB bad

$$|in\rangle \neq |out\rangle$$

Number operator

$$N_a = \int \frac{d^3k}{(2\pi)^3} a_k^\dagger a_k$$

$$= \frac{1}{\sqrt{2\omega_k}} p - i \int \omega_k dt$$

is good

$$U_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

$$|in\rangle \neq |out\rangle$$

Number operator

$$N_a = \int \frac{d^3k}{(2\pi)^3} a_k^\dagger a_k$$

$$[a_k, N_a] = a_k \quad N_b = \int \frac{d^3k}{(2\pi)^3} b_k^\dagger b_k$$

is good

$$U_k = \alpha_k \frac{1}{\sqrt{2\omega_k}} e^{-i \int \omega_k dt} + \beta_k \frac{1}{\sqrt{2\omega_k}} e^{+i \int \omega_k dt}$$

$$|in\rangle \neq |out\rangle$$

Number operator

$$N_a = \int \frac{d^3k}{(2\pi)^3} a_k^\dagger a_k$$

$$[a_k, N_a] = a_k \quad N_b = \int \frac{d^3k}{(2\pi)^3} b_k^\dagger b_k$$

$e^{-i \int \omega_k dt}$

$$\hat{U}(\eta, x) = \int \frac{d^s k}{(2\pi)^s} \alpha_k u_k^{\text{in}} e^{ikx} + \alpha_k^\dagger u_k^{\text{in}*} e^{-ikx} =$$

N_b

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} \alpha_k u_k^{\text{in}} e^{i\mathbf{k}\cdot\mathbf{x}} + \alpha_k^\dagger u_k^{\text{in}*} e^{-i\mathbf{k}\cdot\mathbf{x}} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\vec{k} \cdot \vec{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\vec{k} \cdot \vec{x}} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} \langle \text{in} | b_k^\dagger b_k | \text{in} \rangle$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{i\vec{k}\cdot\vec{x}} + a_k^\dagger u_k^{\text{in}*} e^{-i\vec{k}\cdot\vec{x}} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} \langle \text{in} | b_k^\dagger b_k | \text{in} \rangle \quad a_k | \text{in} \rangle = 0$$

$$b_k | \text{in} \rangle =$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k^{\text{in}} e^{ikx} + a_k^\dagger u_k^{\text{in}*} e^{-ikx} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$\int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle) b_k | \text{in} \rangle \quad a_k | \text{in} \rangle = 0$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\int \frac{d^3 k}{(2\pi)^3} e^{ikx}$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e^{ikx}$$

$$b_k$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle) b_k | \text{in} \rangle \quad a_k | \text{in} \rangle = 0$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k}$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

b_k

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle) b_k | \text{in} \rangle \quad a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k|^2 \langle \text{in} | a_{-k} a_{-k}^\dagger | \text{in} \rangle$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

$$b_k$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle b_k | \text{in} \rangle) \quad a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k|^2 \langle \text{in} | a_{-k} a_{-k}^\dagger | \text{in} \rangle$$

$$(2\pi)^3 \int^{(3)} (a)$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

$$b_k$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\delta^{(3)}(k) = \int d^3 x e^{i k x}$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$= \int \frac{d^3 k}{(2\pi)^3} \langle \text{in} | b_k | \text{in} \rangle \langle \text{in} | a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k| \langle \text{in} | b_k | \text{in} \rangle = \beta_k^\dagger \langle \text{in} | a_{-k} | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e^{i k x}$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e^{i k x}$$

$$| b_k \rangle$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle \quad (2\pi)^3 \delta^{(3)}(k) = \int d^3 x = V \rightarrow$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle b_k | \text{in} \rangle) \quad a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k|^2 \langle \text{in} | a_{-k} a_{-k}^\dagger | \text{in} \rangle$$

$$(2\pi)^3 \int^{(3)} \delta(a)$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

$$b_k$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$(2\pi)^3 \delta^{(3)}(k) = \int d^3 x = V$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle b_k | \text{in} \rangle) \quad a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k|^2 V$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k} | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

$$b_k$$

$$\hat{U}(\eta, x) = \int \frac{d^3 k}{(2\pi)^3} a_k u_k e^{i k x} + a_k^\dagger u_k^* e^{-i k x} =$$

$$\langle \text{in} | N_b | \text{in} \rangle$$

$$(2\pi)^3 \delta^{(3)}(k) = \int d^3 x = V$$

$$= \int \frac{d^3 k}{(2\pi)^3} (\langle \text{in} | b_k^\dagger \rangle b_k | \text{in} \rangle)$$

$$a_k | \text{in} \rangle = 0$$

$$= \int \frac{d^3 k}{(2\pi)^3} |\beta_k|^2 V$$

$$b_k | \text{in} \rangle = \beta_k^\dagger a_{-k}^\dagger | \text{in} \rangle$$

$$\langle \text{in} | b_k^\dagger = \langle \text{in} | a_{-k} \beta_k$$

$$\int \frac{d^3 k}{(2\pi)^3} e$$

$$= \int \frac{d^3 k}{(2\pi)^3}$$

$$= \int \frac{d^3 k}{(2\pi)^3} e$$

$$b_k$$