

Title: Explorations in String Theory - Lecture 2

Date: Mar 15, 2011 11:30 AM

URL: <http://pirsa.org/11030049>

Abstract:

for Atoms

Molecule?

$$e^F = \int \mathcal{D}M \exp\left(-\frac{1}{2g_s} \text{Tr}(M^2 + \sum_P \frac{t_P}{P} M^P)\right)$$

$N \rightarrow \infty$
 $g_s \rightarrow 0$

$M \rightarrow \sqrt{g_s} M$

$$e^F = \int \mathcal{D}M \exp\left(-\frac{1}{2} \text{Tr}(M^2 + \sum_P g_s^{\frac{P}{2}-1} \frac{t_P}{P} M^P)\right)$$

propagator
 vertices



$$F = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} = \sum_{g=0}^{\infty} g_s^{2g-2} \int \mathcal{D}\lambda \left(\lambda_i \frac{t_i}{P}\right)$$

for Atoms

Molecule!

$$e^F = \int \mathcal{D}M \exp\left(-\frac{1}{2g_s} \text{Tr}(M^2 + \sum_P \frac{t_P}{P} M^P)\right)$$


$N \times N$

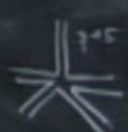
$$M \rightarrow \sqrt{g_s} M$$

$$e^F = \int \mathcal{D}M \exp\left(-\frac{1}{2} \text{Tr}(M^2 + \sum_P g_s^{\frac{P}{2}-1} \frac{t_P}{P} M^P)\right)$$

$$N \rightarrow \infty$$

$$g_s \rightarrow 0$$

propagator 

vertices 

$$F = \text{[circle with line]} + \text{[circle with loop]} + \text{[circle with two loops]} = \sum_{g=0}^{\infty} g_s^{2g-2} \int \mathcal{D}(\lambda_i, t_P)$$

$$\int \mathcal{D} = \sum_{\text{all planar diagrams}}$$

① E
 g_s
propagators

$$\begin{array}{ccc}
 \textcircled{1} & \begin{array}{c} E \\ \int_S \\ \text{propagators} \end{array} & \begin{array}{c} -V \\ \int_S \\ \text{vertices} \end{array} & \begin{array}{c} F \\ N \\ \text{color loops} \end{array}
 \end{array}$$

$$\begin{array}{ccc}
 \textcircled{2} & 1 & \prod_P \int_S^{(P_2-1)V_P}
 \end{array}$$

$$\begin{array}{ccc}
 \textcircled{1} & \begin{array}{c} E \\ g_s \end{array} & \begin{array}{c} -V \\ g_s \end{array} & \begin{array}{c} F \\ N \end{array} \\
 & \underbrace{\hspace{2cm}} & \underbrace{\hspace{2cm}} & \underbrace{\hspace{2cm}} \\
 & \text{propagators} & \text{vertices} & \text{color loops}
 \end{array}$$

$$\begin{array}{ccc}
 \textcircled{2} & 1 & \prod_P g_s^{(P/2 - 1) V_P} & N^F
 \end{array}$$

$$\begin{array}{cc}
 -\sum_P V_P & \sum_P \frac{P V_P}{2} \\
 g_s & g_s
 \end{array}$$

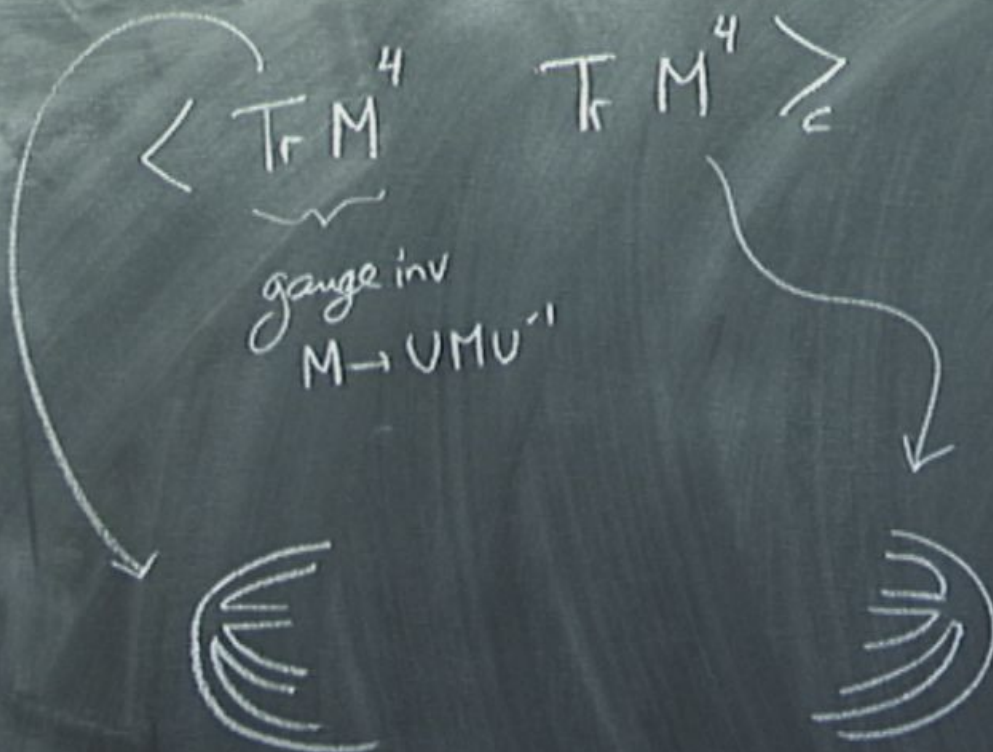


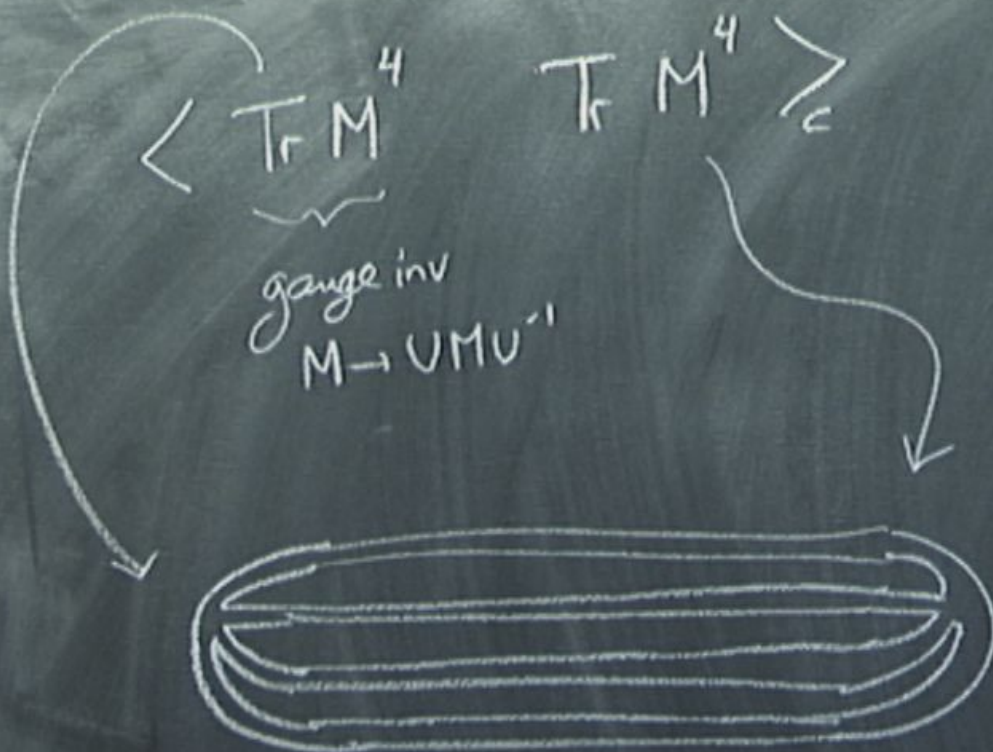
$$g_s^{-V}$$

$$\sum_P V_P P = 2E$$

$$\langle \text{Tr } M^4 \quad \text{Tr } M^4 \rangle$$

gauge inv
 $M \rightarrow U M U^{-1}$





$$\langle \text{Tr } M^4 \quad \text{Tr } M^4 \rangle$$

gauge inv
 $M \rightarrow U M U^{-1}$



change



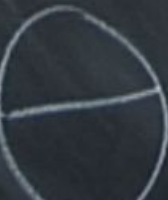
$$\left\langle \frac{1}{N^2} \text{Tr} M^4 \quad \frac{1}{N^2} \text{Tr} M^4 \right\rangle$$

gauge inv
 $M \rightarrow U M U^{-1}$



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

change



$$\frac{1}{N^2} T F M^4$$

free

μ^{-1}



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

$\langle 000 \rangle$

||



$$\frac{1}{N^2} T M^4 \left. \begin{array}{l} \text{free} \\ c \end{array} \right\}$$

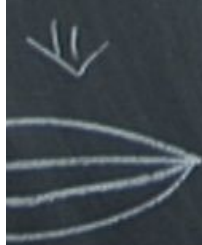
MU^{-1}



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

$\langle 000 \rangle$

\parallel



$$\langle 00 \rangle = \underbrace{\mathbb{F}(\lambda)}_{\lambda \leftrightarrow \text{string tension}}$$



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

$$\langle 00 \rangle = \underbrace{F(\lambda)}_{\lambda \leftrightarrow \text{string tension}}$$



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

$$\langle 0000 \rangle$$

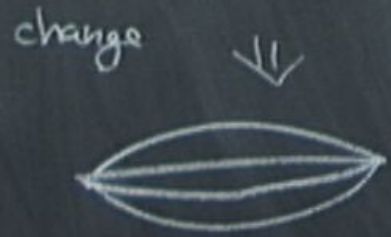
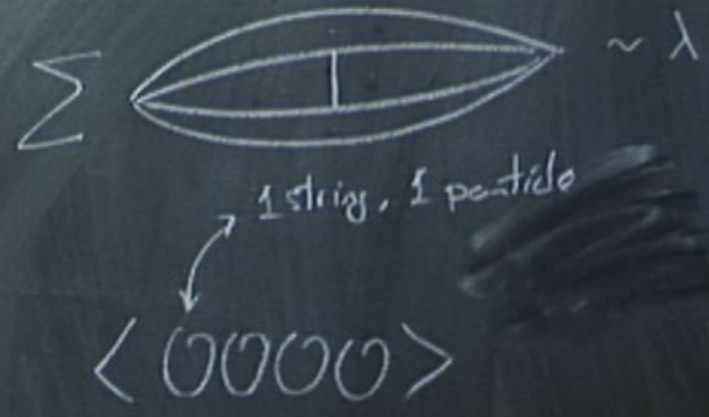
$$\left\langle \frac{1}{N^2} \text{Tr} M^4 \right\rangle \quad \left\langle \frac{1}{N^2} \text{Tr} M^4 \right\rangle$$

0

free

gauge inv
 $M \rightarrow U M U^{-1}$

$$\langle 00 \rangle_c^{\text{int}} = \underbrace{F(\lambda)}_{\lambda}$$



$\langle \frac{1}{N^2} \text{Tr} M^4 \rangle$

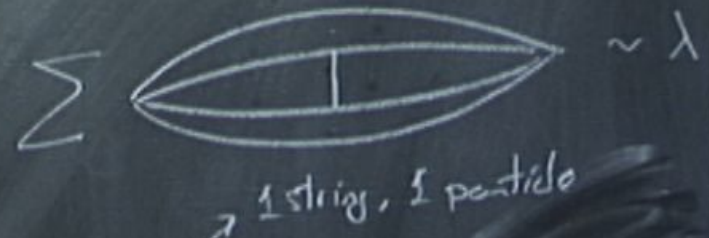
gauge inv
 $M \rightarrow U M U^{-1}$

free

$\langle 00 \rangle_c^{\text{int}} = \mathcal{F}(\lambda)$

\parallel

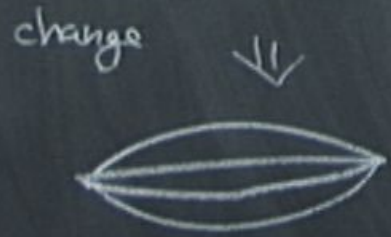
$\lambda \leftarrow$



$\langle 0000 \rangle =$

\Uparrow

\Uparrow



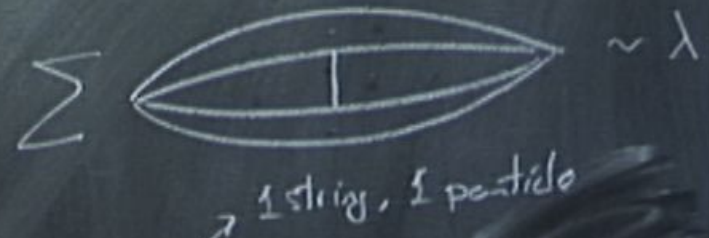
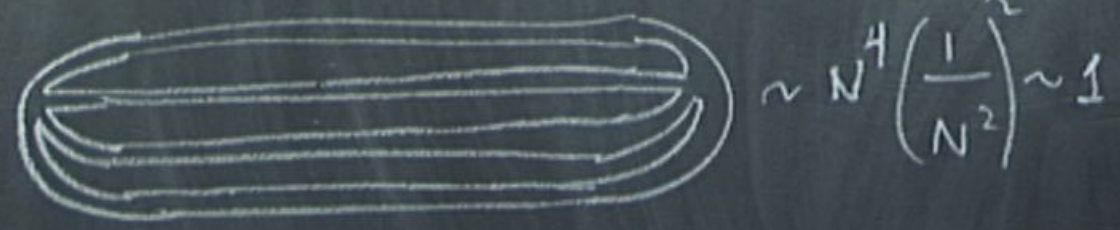
$$\langle \frac{1}{N^2} \text{Tr} M^4 \rangle \quad \langle \frac{1}{N^2} \text{Tr} M^4 \rangle$$

gauge inv
 $M \rightarrow U M U^{-1}$

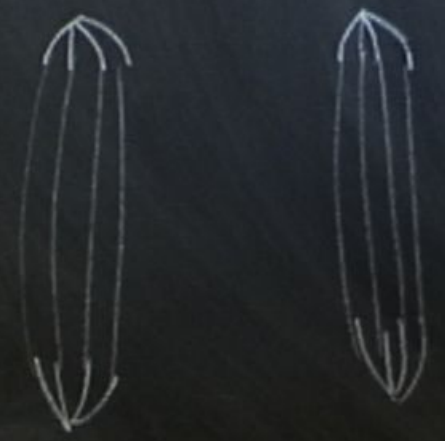
free

$$\langle \text{tr} \text{tr} \rangle_c^{\text{int}} = \mathcal{F}(\lambda)$$

$\lambda \leftarrow$



$$\langle \text{tr} \text{tr} \text{tr} \text{tr} \rangle =$$

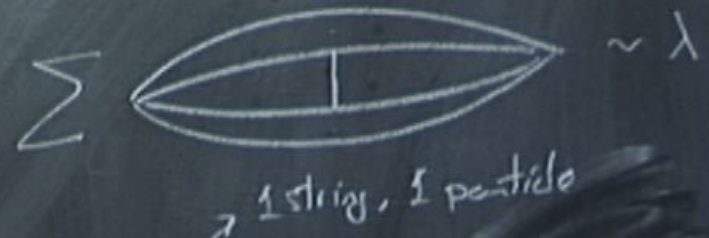


$$\left\langle \frac{1}{N^2} \text{Tr} M^4 \right\rangle \quad \left\langle \frac{1}{N^2} \text{Tr} M^4 \right\rangle$$

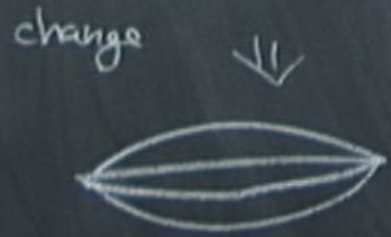
gauge inv
 $M \rightarrow U M U^{-1}$

free

$$\langle 00 \rangle_c^{\text{int}} = \underbrace{F(\lambda)}_{\lambda}$$

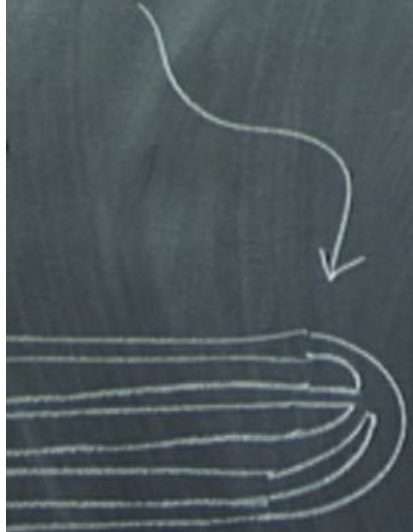


$$\langle 0000 \rangle =$$



$T M^4 \xrightarrow{c} \text{free}$

$$\langle 00 \rangle_c^{\text{int}} = \underbrace{F(\lambda)}_{\lambda \leftrightarrow \text{string tension}}$$

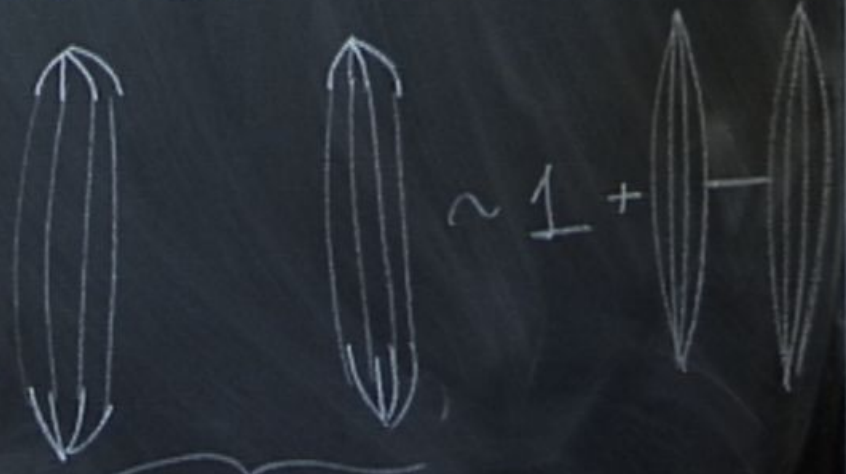


$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$



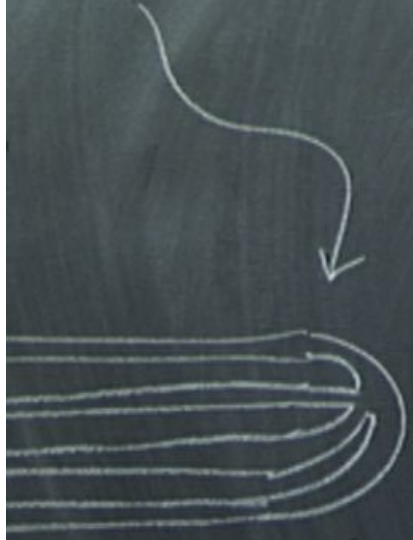
1 string, 1 particle

$$\langle 0000 \rangle =$$



$T M^4$ free

$$\langle 00 \rangle_c^{int} = \underbrace{F(\lambda)}_{\lambda \leftrightarrow \text{string tension}}$$



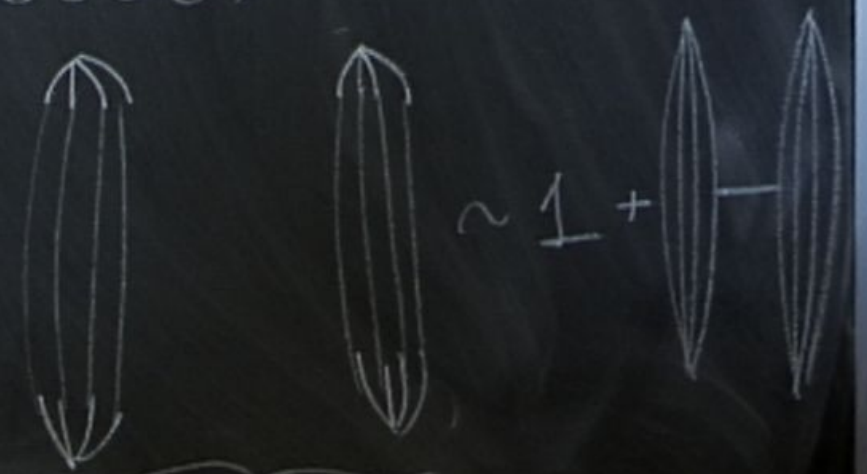
$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$



1 string, 1 particle

$$\langle 0000 \rangle =$$

suppressed @ large N



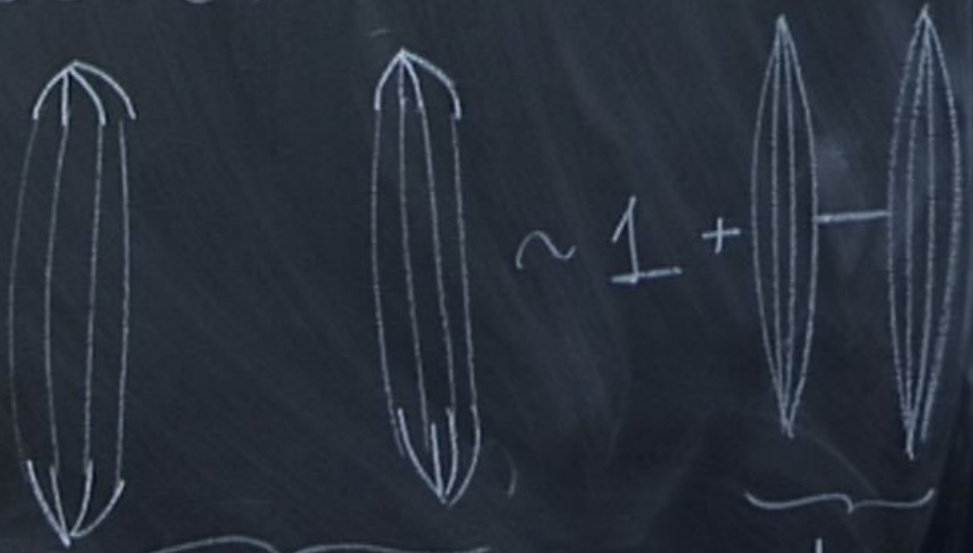
$\lambda \leftrightarrow$ string tension



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

1 string, 1 particle
 $\langle 0000 \rangle =$

suppressed @ large N



$\langle 00 \rangle \langle 00 \rangle$

$\frac{1}{N^2}$



①

g_s
 ───
 propagators

g_s
 ───
 vertices

N
 ───
 Ed

②

┌───┐
 1

$\prod_p g_s^{(P/2 - 1) V_p}$

$\text{Tr } \Pi^2$

$\text{Tr } \Pi^2$



$\text{Tr } \Pi^2$

$\text{Tr } \Pi^2$

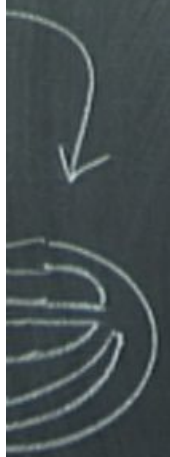
$g_s^{-\sum_p V_p}$ $g_s^{\sum_p V_p}$

g_s^{-V}

ee

$$\langle 00 \rangle_c^{\text{int}} = \underbrace{F(\lambda)}_{\lambda \leftrightarrow \text{string tension}}$$

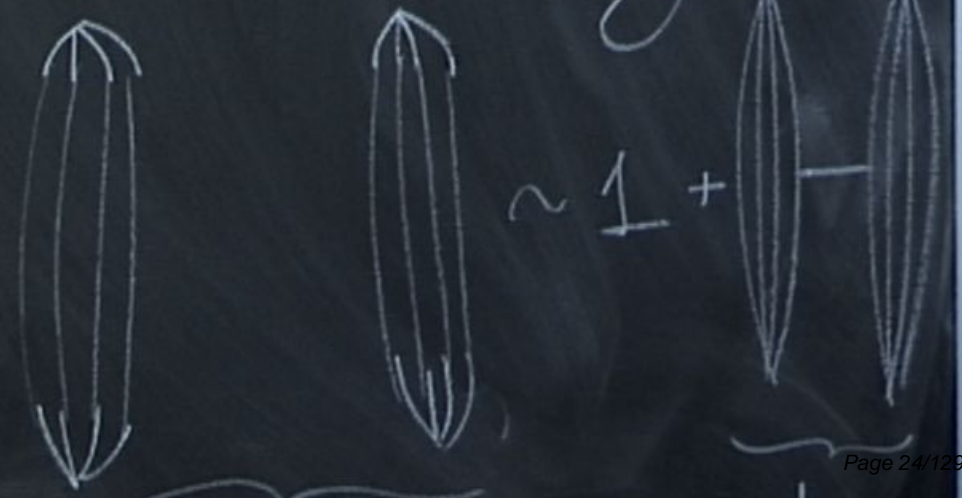
string tension



$$\sim N^4 \left(\frac{1}{N^2} \right)^2 \sim 1$$

1 string, 1 particle

$$\langle 0000 \rangle = \underbrace{\langle 00 \rangle \langle 00 \rangle}_{\text{suppressed @ large } N} + \dots$$





$$\lambda \sim g_s N \frac{1}{N}$$

E
 g_s
 propagators

$$g_s \rightarrow 0$$

$$N \rightarrow \infty$$

$g_s N$ fixed

$$-V$$

$$F$$

g_s N
 vertices edge loops

$$1$$

$$\prod_P g_s^{(P-1)V_P}$$

N^F

$$- \sum_P V_P$$

$$g_s$$

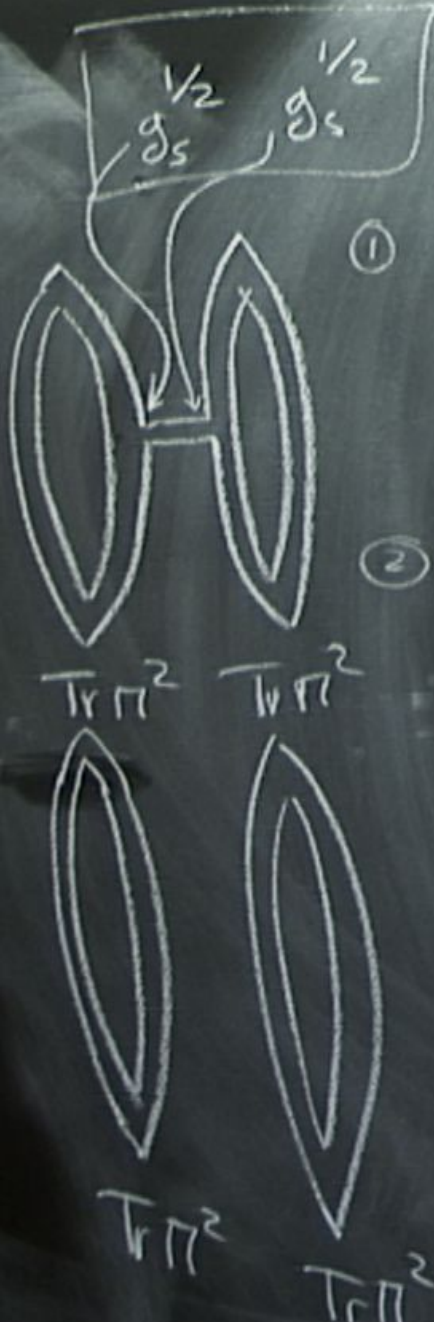
$$\sum_P \frac{P V_P}{2}$$

$$g_s$$



$$g_s^{-V}$$

$$\sum_P V_P P = 2E$$



$$\lambda \sim \frac{1}{g_s N} \sim \frac{1}{N}$$

$$E \sim g_s$$

propagators

$$g_s \rightarrow 0$$

$$N \rightarrow \infty$$

$$g_s N \text{ fixed}$$

$$-V$$

$$g_s$$

vertices

$$N^F$$

edge loops

$$1$$

$$\prod_P g_s^{(P-1)V_P} N^F$$

$$- \sum_P V_P$$

$$g_s$$

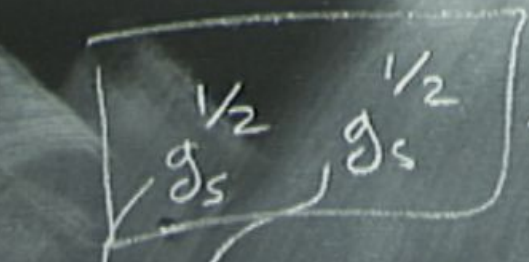
$$\sum_P \frac{P V_P}{2}$$

$$g_s$$



$$g_s^{-V}$$

$$\sum_P V_P P = 2E$$



①



②

$\text{Tr } \Pi^2$ $\text{Tr } \Pi^2$



$$\lambda \sim g_s N \frac{1}{N}$$

$$E \sim g_s$$

propagators

$$g_s \rightarrow 0$$

$$N \rightarrow \infty$$

$$g_s N \text{ fixed}$$

$$-V \quad F$$

$$g_s \quad N$$

vertices Edg loops

$$1 \quad \prod_P g_s^{(P/2 - 1) V_P} N^F$$

$$g_s^{-\sum_P V_P} g_s^{\sum_P \frac{P V_P}{2}}$$



$$\sum_P V_P P = 2E$$

2D QG

$$Z = \sum_{\text{genus}} \int \mathcal{D}g \exp(-\beta A + \gamma X) \quad \chi$$

$$\sum_{\text{genus}} \int \mathcal{D}g \exp(-\beta A + \gamma X)$$

$$X = \frac{1}{4\pi} \int \sqrt{g} R$$

$$A = \int \sqrt{g}$$

2D QG

$$Z = \sum_{\text{genus}} \int \mathcal{D}g \exp(-\beta A + \gamma \chi)$$

$$\chi = \frac{1}{4\pi} \int \sqrt{g} R$$

$$A = \int \sqrt{g}$$

↖ surface

$$(-\beta A + \gamma \chi)$$

$$\chi = 2 \text{ genus} = 2$$

$$\chi = \frac{1}{4\pi} \int_{\Sigma} \sqrt{g} R$$

$$A = \int_{\Sigma} \sqrt{g} \quad \leftarrow \text{cosmological constant}$$

surface

$$(-\beta A + \gamma \chi)$$

$$-\chi = 2 \text{ genus} - 2$$

$$\chi = \frac{1}{4\pi} \int_{\Sigma} \sqrt{g} R$$

$$A = \int_{\Sigma} \sqrt{g} \quad \leftarrow \text{cosmological constant}$$

surface

the free parameters

$$g \exp(-\beta A + \gamma \chi)$$

$$-\chi = 2 \text{ genus} - 2$$

$$\chi = \frac{1}{4\pi} \int_{\Sigma} \sqrt{g} R$$

$$A = \int_{\Sigma} \sqrt{g} \quad \leftarrow \text{cosmological constant}$$

surface

2D QG

$Z[\beta, \gamma]$

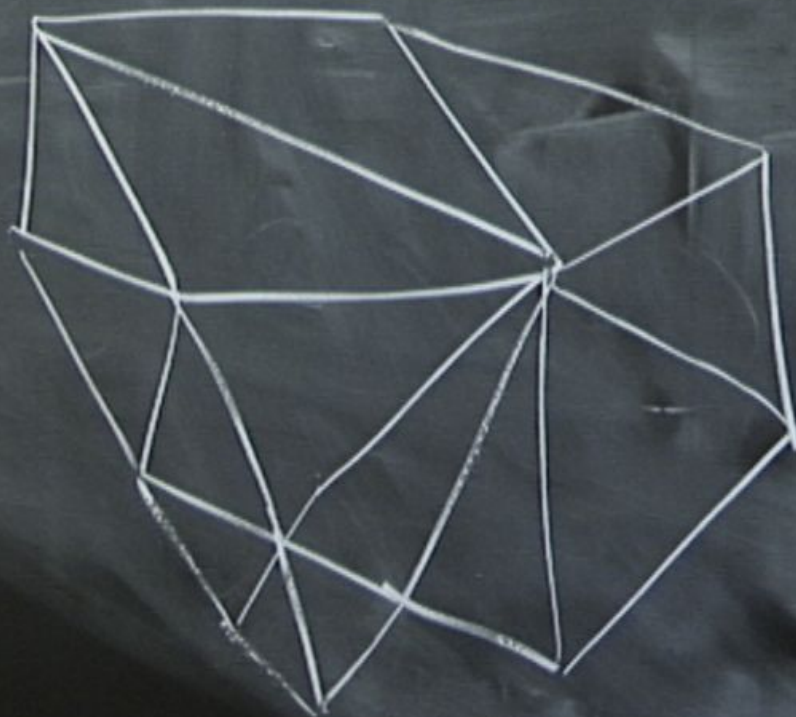
genus

$\int \mathcal{D}g$

the free part

$$\exp(-\beta A + \gamma \chi)$$

$$-\chi = 2 \text{ genus}$$



the free parameters

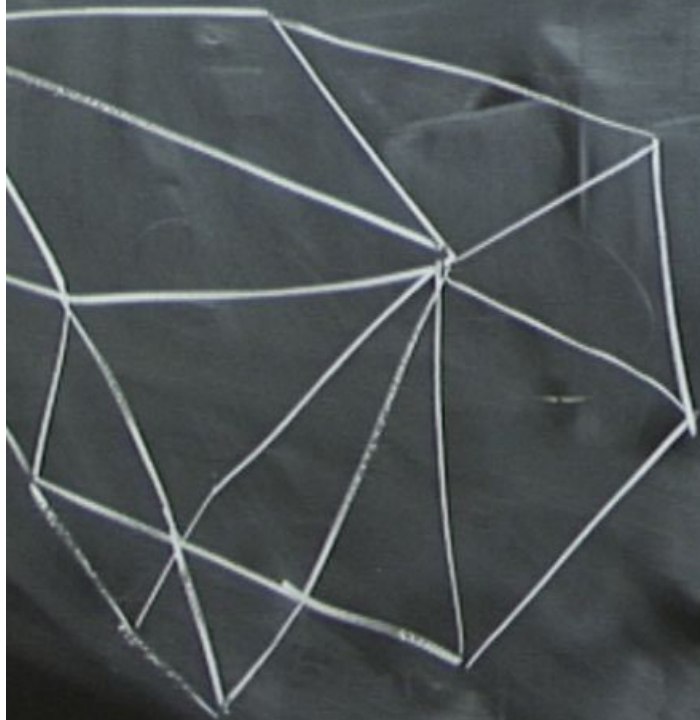
$$\sum_{\text{genus}} \int \mathcal{D}g \exp(-\beta A + \gamma \chi)$$

$$-\chi = 2 \text{ genus} - 2$$

$$\chi = \frac{1}{4\pi} \int$$

$$A = \int \sqrt{g}$$

\sum
 \curvearrowright Sur



$$\int \sqrt{g} R \rightarrow \sum_{i=4}$$

the free parameters

$$g \exp(-\beta A + \gamma \chi)$$

$$-\chi = 2 \text{ genus} - 2$$

$$\chi = \frac{1}{4\pi} \int_{\Sigma} \sqrt{g} R$$

$$A = \int_{\Sigma} \sqrt{g} \quad \leftarrow \text{cosmological constant}$$

surface



$$\int \sqrt{g} R \rightarrow \sum_{i=1}^V 4\pi \left(1 - \frac{N_i}{6} \right)$$

lines emanating from vertex i

$$\sum_{i=1}^V 4\pi = 4\pi V$$

$$- \sum_{i=1}^V \frac{4\pi}{6} N_i$$

now

$$2E = \sum_i N_i$$

$$\sum_{i=1}^V 4\pi = 4\pi V$$

$$- \sum_{i=1}^V \frac{4\pi}{6} N_i$$

now

$$2E = \sum_i N_i, \quad 2E = 3F$$

$$\sum_{i=1}^V 4\pi = 4\pi V$$

$$- \sum_{i=1}^V \frac{4\pi}{6} N_i = -4\pi \left(\frac{F}{2} \right)$$

now

$$2E = \sum_i N_i$$

$$2E = 3F$$

each face has 3
edges and each edge
is shared by 2
faces

$$\sum_{i=1}^V 4\pi = 4\pi V$$

$$- \sum_{i=1}^V \frac{4\pi}{6} N_i = -4\pi \left(\frac{F}{2} \right)$$

now

$$2E = \sum_i N_i, \quad 2E = 3F$$

each face has 3
edges and each edge
is shared by 2
faces

$$\# = 4\pi (V - E + F) = 4\pi \chi$$

2D QG

$Z[\beta, \gamma]$

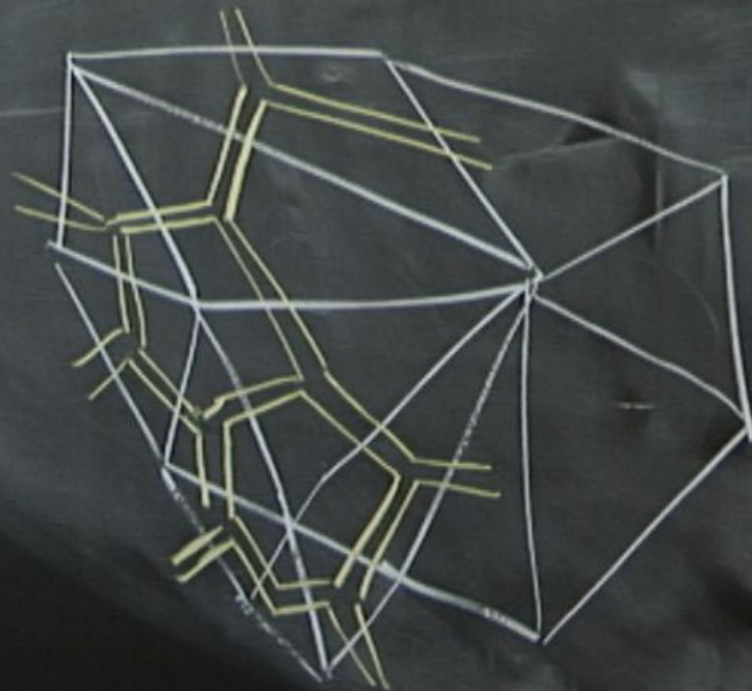
genus

$\int \mathcal{D}g$

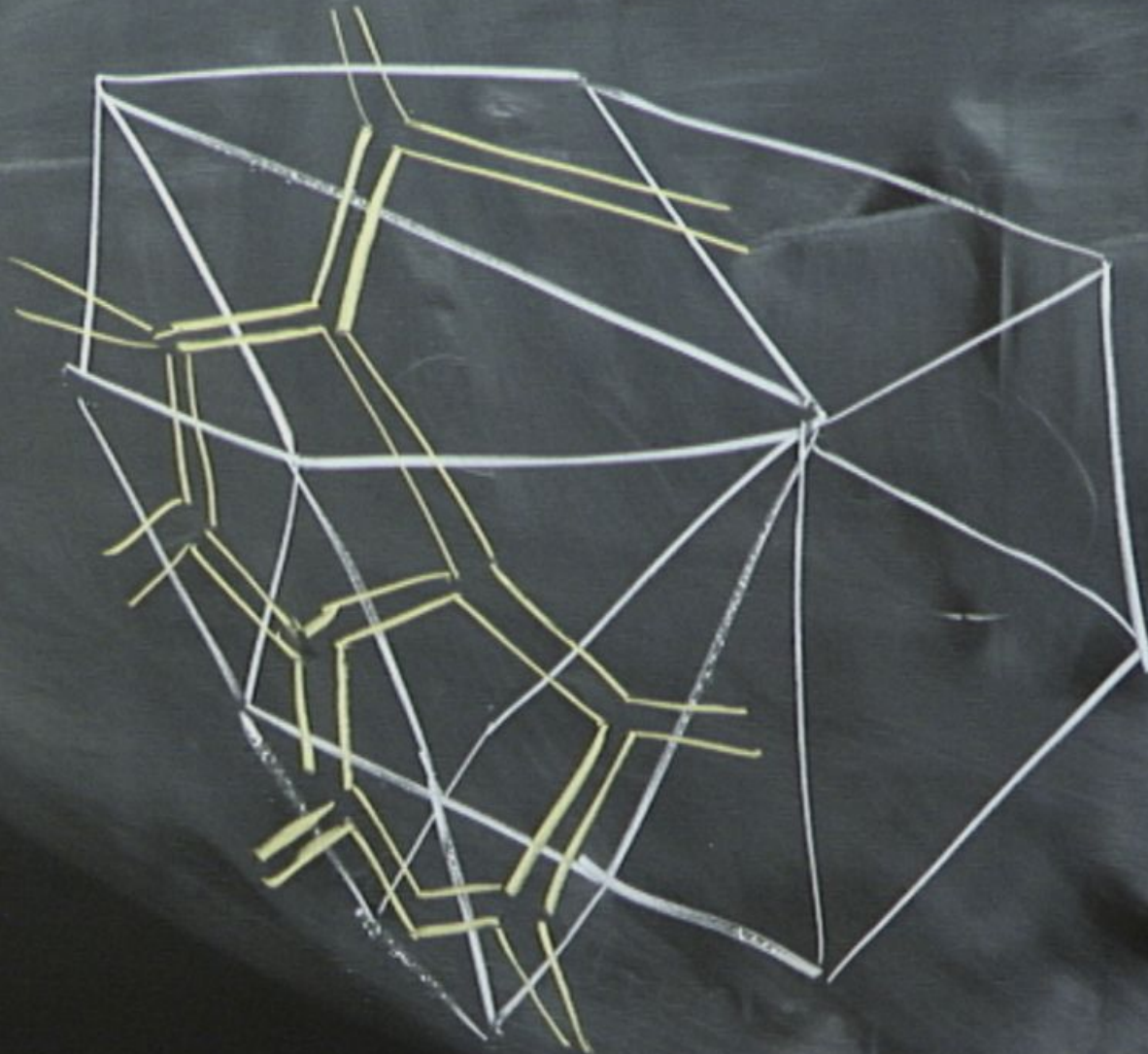
the free parameter

$$\exp(-\beta A + \gamma \chi)$$

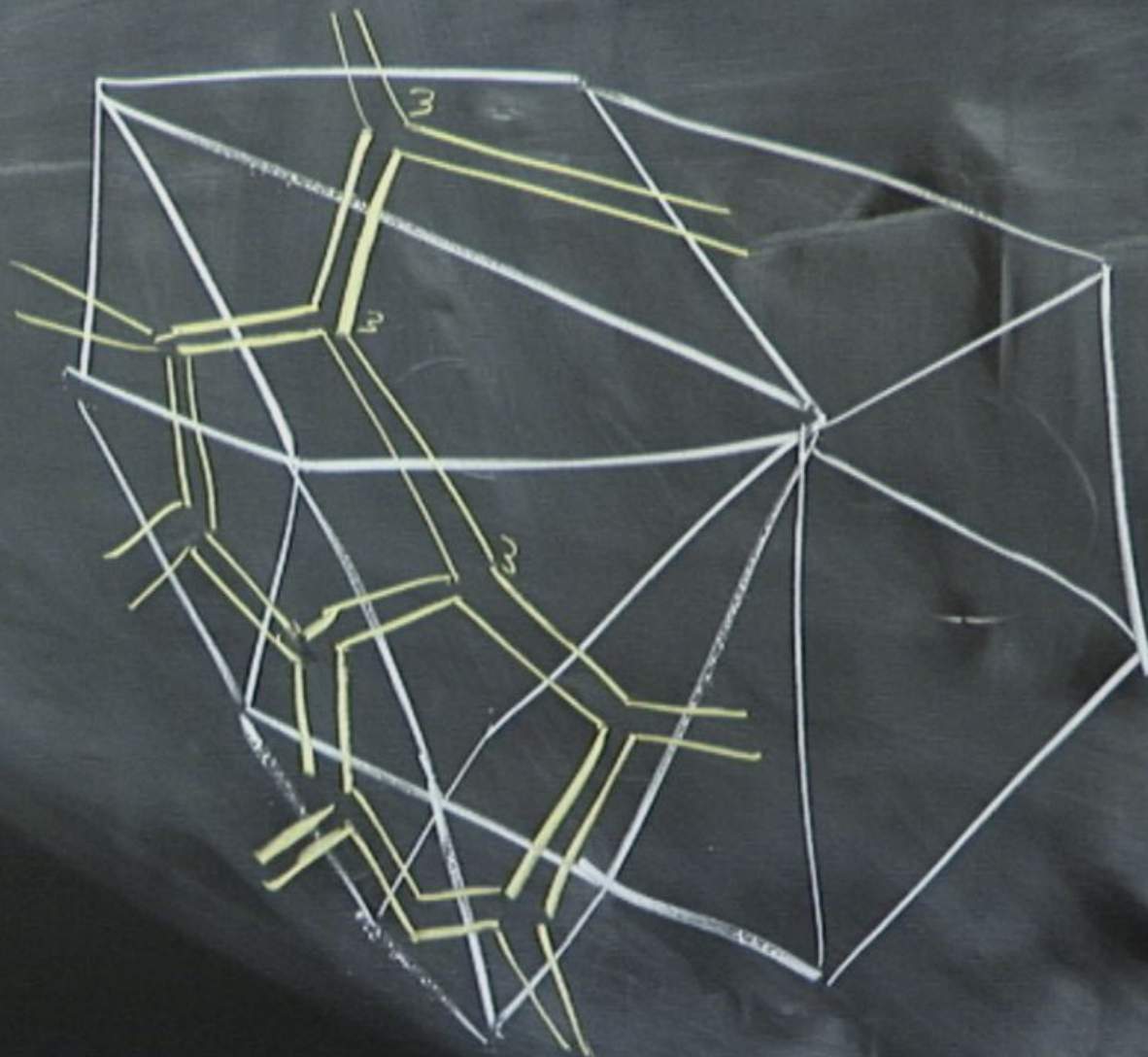
$$-\chi = 2 \text{ genus} - 2$$



$$\int \sqrt{g} R \rightarrow \sum_{i=1}^V 4\pi$$



$$\int \sqrt{g} R$$



$$\int \sqrt{g} R$$

$$e^Z = \int \mathcal{D}M e^{-\frac{1}{2} \text{Tr} M^2 + \frac{g}{\sqrt{N}} \text{tr} M^3}$$

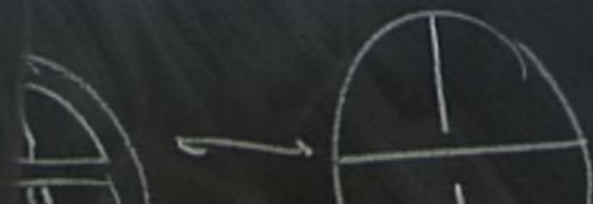
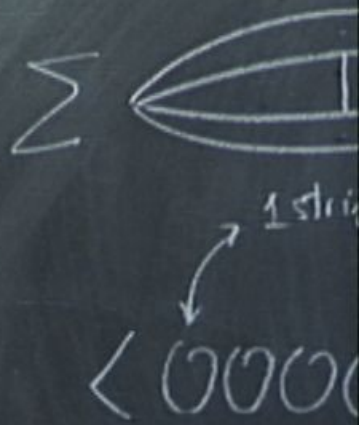
$\langle 00 \rangle$
 \parallel

$$N \equiv e^{\gamma}$$

$$g \equiv e^{-\beta}$$

EXERCISE:
 CHECK THE
 DETAILS

$$(N^2) \sim 1$$



$$e^{\mathcal{Z}} = \int \mathcal{D}M e^{-\frac{1}{2} \text{Tr} M^2 + \frac{g}{\sqrt{N}} \text{Tr} M^3}$$

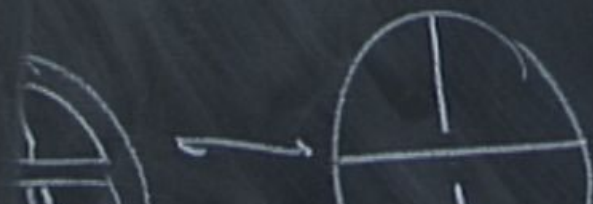
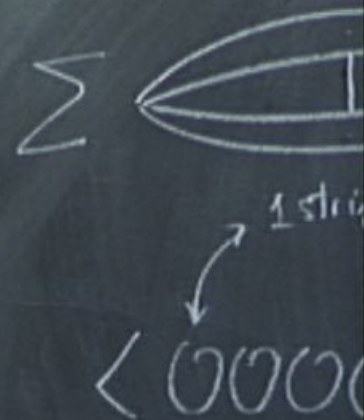
$\langle 00 \rangle$
 \parallel

$$N \equiv e^{\chi}$$

$$g \equiv e^{-\beta}$$

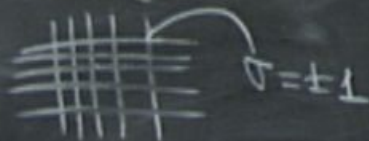
EXERCISE:
 CHECK THE
 DETAILS

$$\mathcal{Z} = N^2 \mathcal{Z}_0(g) + \mathcal{Z}_1(g) + \chi_2(g)/N^2 + \dots$$



2D Ising Model on a random surface

Z Ising model
in a fixed
lattice



$$= \sum_{\{\sigma_a = \pm 1\}}$$

$$e^{+ \beta \frac{J}{2} \sum_{\langle ab \rangle} (\sigma_a \sigma_b - 1) - \beta H \sum_a \sigma_a}$$

spin coupling

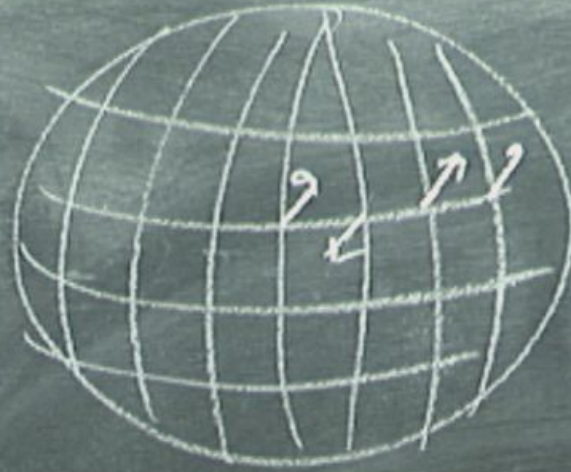
magnetic field

evidence
for Atoms

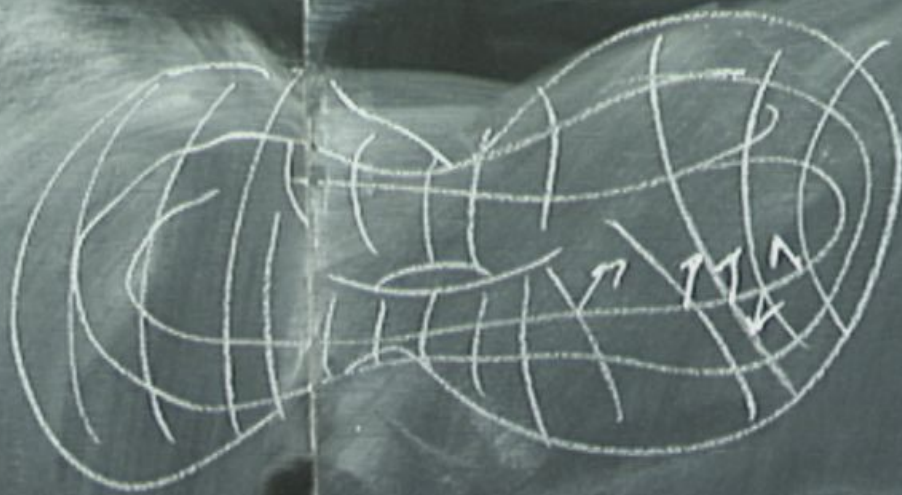
Big Is A
Molecule?

Z Ising model
on a random lattice

$$= \sum_{\uparrow\downarrow}$$



+



Z Ising model
on a random lattice

$$= \sum_{\mathbb{N}}$$



$$= \sum_g N^{2-2g}$$

Z Ising model
on a random lattice

$$= \sum_{\mathbb{N}}$$

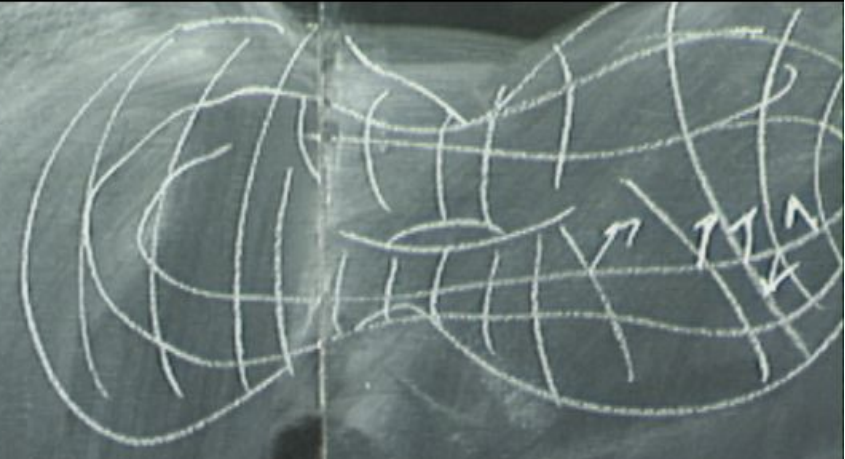


$$= \sum_g N^{2-2g} \sum_{n=1}^{\infty} e^{-\mu n}$$

$$= \sum_{\uparrow \downarrow}$$



+



$$N^{2-2g}$$

$$\sum_{n=1}^{\infty} e^{-\mu n \beta}$$

$$\sum_{\substack{n, g \\ \text{surfaces}}}$$

+



+ ...

$$\sum \sum_{n,g} \text{surfaces}$$

$$\sum \{ \sigma_a = \pm 1 \}$$

$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b)} - \beta H \sum \sigma_i$$

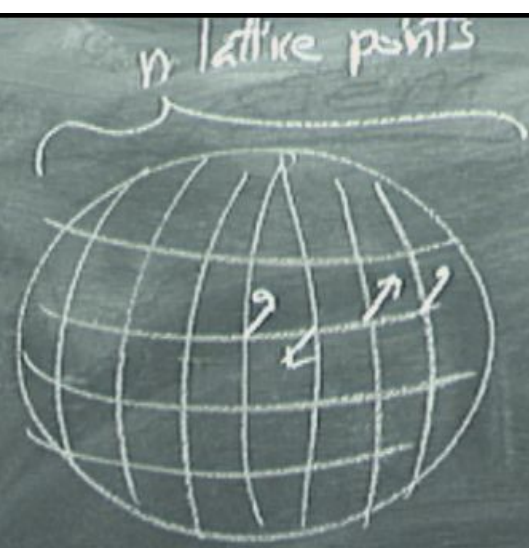


+

$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

Ising model
on a random lattice

$$= \sum_{\uparrow \downarrow}$$



$$Z \equiv \sum_g N^{2-2g} \sum_{n=1}^{\infty} e^{-\mu n \beta} \sum_{\text{surfaces}} \sum_{\{\sigma_a = \pm 1\}}$$

claim

$$e^Z = \int \mathcal{D}A \mathcal{D}B \exp \left(-\text{Tr} \left[\frac{A^2 + B^2}{2} \right] \right)$$

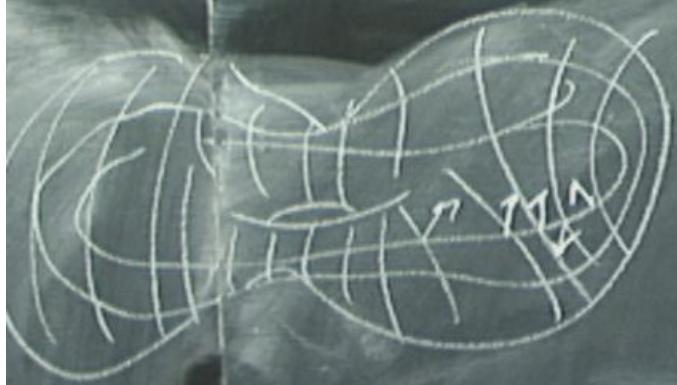


$$\sum_{\sum n, g} \sum_{\{\sigma_a = \pm 1\}}$$

(surfaces)

$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

$$\left(-\text{Tr} \left[\frac{A^2 + B^2 - 2e^{-\beta J} AB}{1 - e^{-2\beta J}} \right] - \frac{e^{-\beta \mu}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} \right) \right)$$



$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

\sum
 $\{\sigma_a = \pm 1\}$

$$\frac{1 + \beta^2 - 2e^{-\beta J} AB}{1 - e^{-2\beta J}} = \frac{e^{-\beta \mu}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$



+ -----

$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

$$\frac{2e^{-\beta J} AB}{1 - e^{-2\beta J}} = \frac{e^{-\beta \mu}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$

Z Ising model
on a random lattice
with coordination number 4

$$= \sum_{\uparrow \downarrow}$$



$$Z \equiv \sum_{g=0}^{\infty} N^{2-2g} \sum_{n=1}^{\infty} e^{-\mu n \beta} \sum_{\substack{n, g \\ \text{surfaces}}} \sum_{\{\sigma_a = \pm 1\}} e^{-\dots}$$

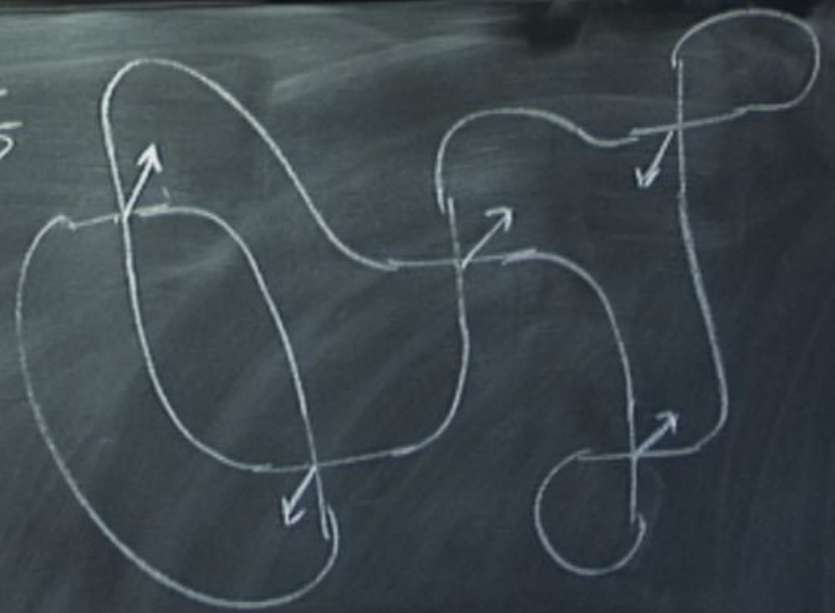
claim

$$e^Z = \int \mathcal{D}A \mathcal{D}B \exp \left(-\text{Tr} \left[\frac{A^2 + B^2}{1 - e^{-\dots}} \right] \right)$$



+ ...

example
for n=5



$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b)} - \beta H \sum_a \sigma_a$$

$$\frac{e^{-\beta J} AB}{e^{-2\beta J}} = \frac{e^{-\beta H}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$

2D Ising Model on a random surface

$$A \equiv \dot{A}$$

$$B \equiv B$$

$$A \equiv B$$

2D Ising Model on a random surface

$$A \equiv \dot{A} = 1$$

$$B \equiv B = 1$$

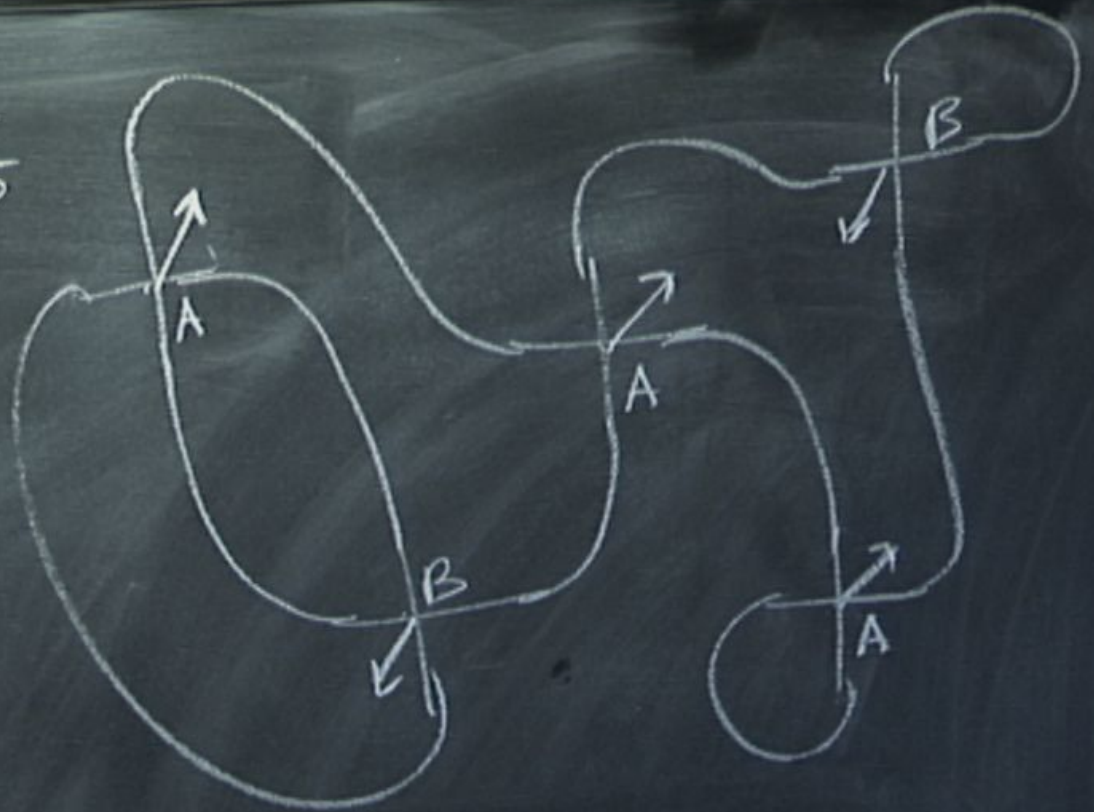
$$A \equiv B = e^{-\beta J}$$

2D Ising Model on a random surface

$$\begin{aligned} A & \equiv \dot{A} & = 1 \\ B & \equiv B & = 1 \\ A & \equiv B & = e^{-\beta J} \end{aligned}$$

this is
why we
chose that
kin. term

example
for $n=5$



$$- \beta H \sum_a \sigma_a$$

$$\left(e^{-\beta H A^4} + e^{\beta H B^4} \right)$$

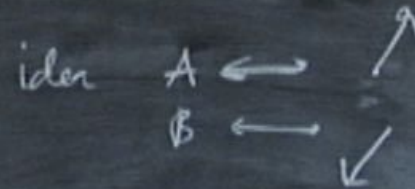
Ising Model on a random surface

$$A \equiv \bar{A} = 1$$

$$B \equiv \bar{B} = 1$$

$$A \equiv B = e^{-\beta J}$$

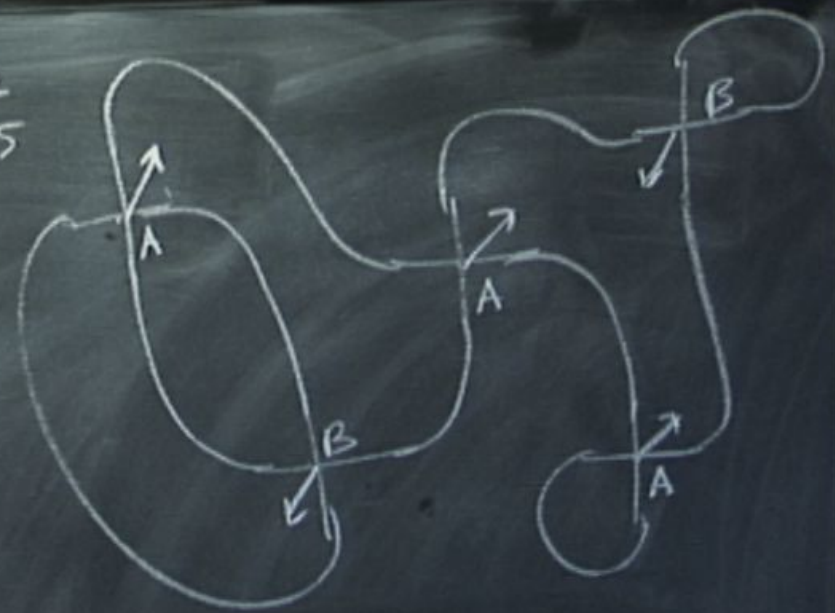
this is
why we
chose that
kin-term





+ ...

example
for n=5



$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

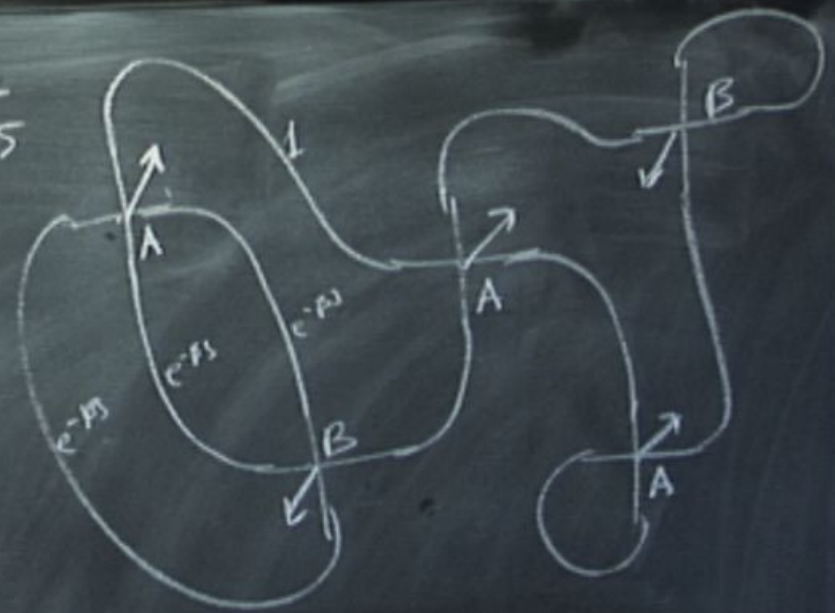
$e^{-\beta J}$ for each \downarrow connected to \uparrow

$$\frac{2e^{-\beta J} AB}{1 - e^{-2\beta J}} = \frac{e^{-\beta \mu}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$



+ ...

example
for n=5



$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

$e^{-\beta J}$ for each \downarrow connected to \uparrow

$$\frac{2e^{-\beta J} AB}{1 - e^{-2\beta J}} = \frac{e^{-\beta \mu}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$

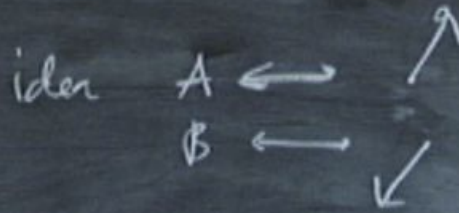
on a random surface

$$\text{---} \dot{A} = 1$$

$$\text{---} B = 1$$

$$\text{---} B = e^{-\beta J}$$

this is
why we
choose that
kin. term



$$+ e^{-\beta \mu} e^{-\beta H}$$

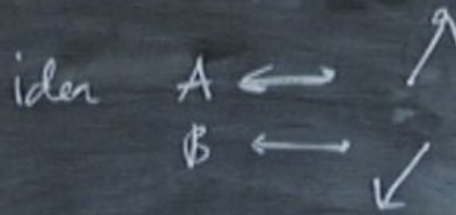
random surface

$$= 1$$

$$B = 1$$

$$B = e^{-\beta J}$$

this is why we chose that kin. term



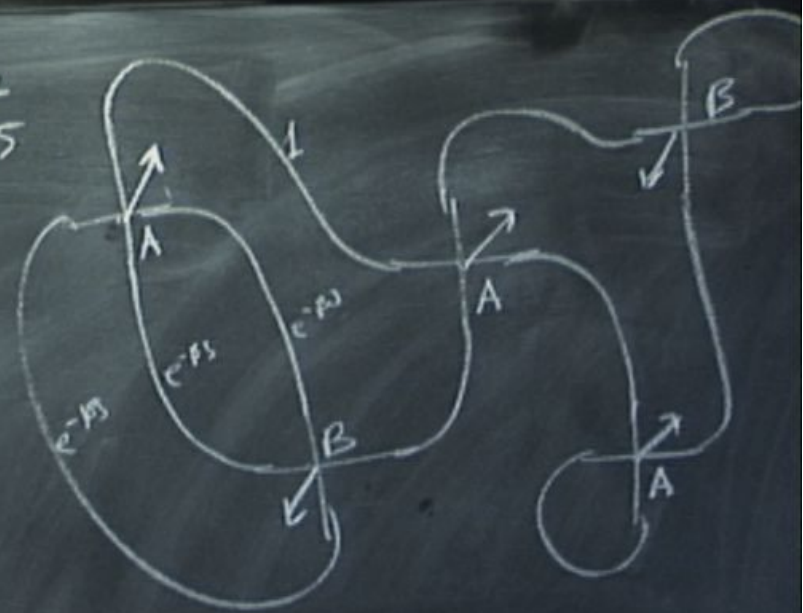
$$A \left| \begin{array}{l} e^{-\beta H} \\ e^{-\beta H} \end{array} \right.$$

$$B \left| \begin{array}{l} e^{-\beta H} \\ e^{+\beta H} \end{array} \right.$$



+ ...

example for n=5



$$e^{-\frac{\beta}{2} \sum_{\langle ab \rangle} (1 - \sigma_a \sigma_b) - \beta H \sum_a \sigma_a}$$

$e^{-\beta J}$ for each \downarrow connected to \uparrow

$$\frac{2e^{-\beta J} AB}{1 - e^{-2\beta J}} = \frac{e^{-\beta H}}{N} \text{Tr} \left(e^{-\beta H} A^4 + e^{\beta H} B^4 \right)$$

$$\frac{1}{Z} \int \mathcal{D}M \left(\frac{1}{N} \text{Tr} e^M \right) e^{-\frac{Z}{g^2} \text{Tr} M^2}$$

$$\frac{1}{Z} \int \mathcal{D}M \left(\frac{1}{N} \text{Tr} e^M \right) e^{-\frac{Z}{g^2} \text{Tr} M^2}$$

$N \times N$ matrices

$$\int \mathcal{D}M \left(\frac{1}{N} \text{Tr} e^M \right) e^{-\frac{2}{g^2} \text{Tr} M^2} = \frac{1}{N} \underbrace{L_{N-1}^1 \left(-\frac{\lambda}{4N} \right)}_{\text{Laguerre Polynomial}} e^{\frac{\lambda}{8N}}$$

$N \times N$ matrices $\lambda = g^2 N$

$$\int_{\mathcal{M}} \int_{\mathcal{M}} \left(\frac{1}{N} \text{Tr} e^M \right) e^{-\frac{2}{g^2} \text{Tr} M^2} = \frac{1}{N} \underbrace{L_{N-1}^{1} \left(-\frac{\lambda}{4N} \right)}_{\text{Laguerre Polynomial}} e^{\frac{\lambda}{8N}}$$

all torus diags!

$$= \frac{2}{\sqrt{\lambda}} I_1(\sqrt{\lambda}) + \frac{\lambda}{8\pi N^2} I_2(\sqrt{\lambda}) + \dots$$

all planar diagrams!
Bessel functions

$\lambda = g^2 N$

$$\int \mathcal{D}M \left(\frac{1}{N} \text{Tr} e^M \right) e^{-\frac{2}{g^2} \text{Tr} M^2} = \frac{1}{N} \underbrace{L_{N-1}^1 \left(-\frac{\lambda}{4N} \right)}_{\text{Laguerre Polynomial}} e^{\frac{\lambda}{8N}}$$

$N \times N$ matrices

all torus diags!

$$= \frac{2}{\sqrt{\lambda}} I_1(\sqrt{\lambda}) + \frac{\lambda}{8\pi N^2} I_2(\sqrt{\lambda}) + \dots$$

all planar diagrams!

Bessel functions

$$L_n^m \equiv \frac{e^x}{n!} x^{-m} \frac{d^n}{dx^n} \left(e^{-x} x^{n+m} \right)$$

$$\lambda = g^2 N$$

DERIVATION

$$M \rightarrow U M U^{-1} \quad \text{Lorentz Symmetry}$$

DERIVATION

$$M \rightarrow U M U^{-1} \quad \begin{array}{l} \text{Unitary} \\ \text{Symmetry} \end{array}$$

eigenvalues of M not transformed

DERIVATION

$$M \rightarrow U M U^{-1} \quad \begin{array}{l} \text{Huga} \\ \text{Symmetry} \end{array}$$

eigenvalues of M not transformed

$$M \rightarrow \{\lambda_i\}$$

N variables

N variables

$$S[M] = e^{-\text{Tr} M^2} \frac{1}{Zg^2} = e^{-\frac{1}{2g^2} \sum_i \lambda_i^2}$$

DERIVATION

$$M \rightarrow U M U^{-1} \quad \text{Hugo Symmetry}$$

eigenvalues of M not transformed

$$M \rightarrow \{\lambda_i\}$$

N^2 variables

N variables

$$S[M] = e^{-\frac{(\text{Tr } M^2)}{2g^2}} = e^{-\frac{1}{2g^2} \sum_i \lambda_i^2}$$

$$U \begin{pmatrix} \lambda_1 & & \\ & \lambda_2 & \\ & & \ddots \end{pmatrix} U^{-1}$$

$$\text{in } \# \quad \mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \quad \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\text{in } \# \\ \mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{D}M f(M)$$

in #

$$\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{D}M f(M)$$

$$1 = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^+ - \Lambda)$$

in #

$$\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{D}M f(M)$$

$$1 \equiv \int d^N \lambda_i dU \int_{S^{(N^2)}} (UMU^\dagger - \Lambda) \Delta^2(\{\lambda\})$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

in #

$$\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{D}M f(M) \mathbb{1}$$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \int_{\mathcal{S}^{(N^2)}} (UMU^\dagger - \Lambda) \Delta^2(\{\lambda\})$$

$$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \dots & \\ & & \lambda_N \end{pmatrix}$$

def of Δ^2

in #

$$\mathcal{DM} \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{DM} f(M) \mathbb{1} = \int dU \int \mathcal{DM} \int d\lambda f(\lambda) \Delta^2(\lambda) \int \mathcal{D}^{(N^2)}(UMU^{-1})$$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \int \mathcal{D}^{(N^2)}(UMU^{\dagger} - \Lambda) \Delta^2(\{\lambda\})$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$\int \mathcal{D}^{(N^2)}(UMU^{\dagger} - \Lambda) \Delta^2(\{\lambda\})$

in #

$$\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}}$$

$$\int \mathcal{D}M f(M) \mathbb{1} = \int dU \int \mathcal{D}\pi \int d^4\lambda f(\lambda) \Delta^2(\lambda) \int \mathcal{D}M \delta^{(N^2)}(UMU^{-1} - \Lambda)$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \delta^{(N^2)}(UMU^+ - \Lambda) \Delta^2(\{\lambda\})$$

def of Δ^2

in #
 $\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N$

$$\Delta^2(\{\lambda\})$$

Jacobian

(no measure!)

$\mathcal{D}M \rightarrow \mathcal{D}\tilde{M}$

$$\int \mathcal{D}M f(M) \mathbb{1}$$

$$= \int dU \int \mathcal{D}\tilde{M} \int d^N \lambda f(\lambda) \Delta^2(\lambda) \delta^{(N^2)}(U\tilde{M}U^{-1} - \Lambda)$$

$$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \dots & \\ & & \lambda_N \end{pmatrix}$$

$$\mathbb{1} \equiv \int d^N \lambda_i dU$$

def of Δ^2

$$\delta^{(N^2)}(UMU^+ - \Lambda) \Delta^2(\{\lambda\})$$

\tilde{M}

in #
 $\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N$ $\Delta^2(\{\lambda\})$
 Jacobian (no measure!)

$$\int \mathcal{D}M f(M) \mathbb{1} = \int dU \int \mathcal{D}\Pi \int d^4\lambda f(\lambda) \Delta^2(\lambda) \int \mathcal{D}M \delta^{(N^2)}(UMU^{-1} - \Lambda)$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\{\lambda\})$$

def of Δ^2

$$\mathbb{1} = \left(\int dU \right) \int d^4\lambda f(U) \Delta^2(\{\lambda\})$$

in #

$$\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N \quad \underbrace{\Delta^2(\{\lambda\})}_{\text{Jacobian}} \quad \mathcal{D}M \rightarrow \mathcal{D}\tilde{M} \quad (\text{no measure!})$$

$$\int \mathcal{D}M f(M) \mathbb{1} = \int dU \int \mathcal{D}\tilde{M} \int d^4\lambda f(\lambda) \Delta^2(\lambda) \delta^{(N^2)}(UMU^\dagger - \Lambda)$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \quad \delta^{(N^2)}(UMU^\dagger - \Lambda) \Delta^2(\{\lambda\})$$

def of Δ^2

$$\mathbb{1} = \left(\int dU \right) + \int d^4\lambda f(U) \Delta^2(\{\lambda\})$$

in # $\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N$ $\Delta^2(\{\lambda\})$
 Jacobian $\mathcal{D}\Pi \rightarrow \mathcal{D}\tilde{\Pi}$ (no measure!)

$$\int \mathcal{D}M f(M) \mathbb{1} = \int dU \int \mathcal{D}\Pi \int d^4\lambda f(\lambda) \Delta^2(\lambda) \delta^{(N^2)}(U\Pi U^{-1} - \Lambda)$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$$\mathbb{1} \equiv \int d^N \lambda_i dU \delta^{(N^2)}(U M U^+ - \Lambda) \Delta^2(\{\lambda\})$$

$$\mathbb{1} = \left(\int dU \right) \int d^4\lambda f(U) \Delta^2(\{\lambda\})$$

next use Δ to compute

in #
 $\mathcal{D}M \rightarrow d\lambda_1 \dots d\lambda_N$ $\Delta^2(\{\lambda\})$
 Jacobian (no measure!)

$$\int \mathcal{D}M f(M) \mathbb{1} = \int dU \int \mathcal{D}\Pi \int d^N \lambda f(\lambda) \Delta^2(\lambda) \delta^{(N^2)}(UMU^T - \Lambda)$$

$\Lambda = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$

$$\mathbb{1} \equiv \int d^N \lambda_i dU$$

def of Δ^2

$$\mathbb{1} = \left(\int dU \right) \int d^N \lambda f(U) \Delta^2(\{\lambda\})$$

next use Δ to compute

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(U M U^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda^2 U_0$

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UHU^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda U_0$

$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

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$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

U will not

give

$$UMU^{-1} - \Lambda = 0!$$

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$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

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U will not

give

$$UMU^{-1} - \Lambda = 0!$$

only if $U \simeq U_0$

$$U = (1 + T)U_0$$

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda U_0$

$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

U will not

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$$UMU^{-1} - \Lambda = 0!$$

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only $U = (1 + T)U_0$ matter

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda U_0$

$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

U will not

give

$$UMU^+ - \Lambda = 0!$$

$$U^+ = U^{-1}$$

only if $U \simeq U_0$

only $U = (1 + T)U_0$ matter

$$UMU^{-1} = \Lambda + [T, \Lambda]$$

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda U_0$

$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

U will not

give

$$UMU^+ - \Lambda = 0!$$

$$U^+ = U^{-1}$$

only if $U \simeq U_0$

only $U = (1 + T)U_0$ matter

$$(UMU^{-1})^{-1} = \Lambda^{-1} + [T, \Lambda^{-1}]$$

$$I = \int d^N \lambda_i dU \delta^{(N^2)}(UMU^{-1} - \Lambda) \Delta^2(\Lambda)$$

for a generic $M = U_0^{-1} \Lambda U_0$

$$\begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix}$$

U will not

give

$$UMU^+ - \Lambda = 0!$$

$$U^+ = U^{-1}$$

only if $U \simeq U_0$

only $U = (1 + T)U_0$ matter

$$UMU^{-1} = \Lambda + [T, \Lambda]$$

$$\delta^{(\lambda^2)}(\dots) = \delta(\lambda - \lambda') \delta([T, \lambda'])$$

Grains of
Pollen to
Evidence
for Atoms

How
Big Is A
Molecule?

$$1 = \int d^N \lambda \prod_{i < j} dT_{ij} \left[\delta^{(z)}(\dots) = \delta(\lambda - \lambda') \delta([T, \lambda']) \right] \Delta^2(\lambda)$$

$\delta(\text{Re}(\dots)) \delta(\text{Im}(\dots))$
(z)

$$1 = \Delta^2(\lambda') \int \prod_{i < j} dT_{ij} \prod_{i < j} \delta(T_{ij}(\lambda_j - \lambda_i)) \quad , \quad 1 = \frac{\Delta^2(\lambda')}{\prod_{i < j} (\lambda'_i - \lambda'_j)^2}$$

$d\text{Re} T_{ij} d\text{Im} T_{ij}$

$$\Delta(\lambda) = \prod_{i < j} (\lambda_i - \lambda_j)$$

Vandermonde
Det

$$\int \mathcal{D}M f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$N^2 \longrightarrow N \quad \Delta^2$

$$\int \mathcal{D}M f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$$N^2 \longrightarrow N \quad \Delta$$

$$\Delta(\lambda) = \det_{i,j} \lambda_i^{j-1}$$

$$\int_{\mathcal{D}M} f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$N^2 \longrightarrow N \quad \Delta$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$= \det \begin{vmatrix} 1 & \lambda_1 & \lambda_1^2 & \dots & \lambda_1^{N-1} \\ 1 & \lambda_2 & \lambda_2^2 & \dots & \lambda_2^{N-1} \\ 1 & \lambda_3 & \lambda_3^2 & \dots & \lambda_3^{N-1} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \lambda_N & \lambda_N^2 & \dots & \lambda_N^{N-1} \end{vmatrix}$$

$$\begin{vmatrix} \lambda_1^{N-1} \\ \lambda_2^{N-1} \\ \lambda_3^{N-1} \\ \vdots \\ \lambda_N^{N-1} \end{vmatrix}$$

$$\int_{\mathcal{D}M} f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$N^2 \longrightarrow N \quad \Delta$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$\equiv \prod_{i < j} (\lambda_i - \lambda_j)$$

$$= \det$$

| | | | |
|---|-------------|---------------|---------|
| 1 | λ_1 | λ_1^2 | \dots |
| 1 | λ_2 | λ_2^2 | \dots |
| 1 | λ_3 | λ_3^2 | \dots |
| ⋮ | ⋮ | ⋮ | ⋮ |
| 1 | λ_N | λ_N^2 | \dots |

| |
|-------------------|
| λ_1^{N-1} |
| λ_2^{N-1} |
| λ_3^{N-1} |
| ⋮ |
| λ_N^{N-1} |

$$\int_{\mathcal{D}M} f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$$N^2 \longrightarrow N \quad \Delta$$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$\prod_{i < j} (\lambda_i - \lambda_j)$$

$$= \det$$

$$\begin{vmatrix} 1 & \lambda_1 & \lambda_1^2 & \dots & \lambda_1^{N-1} \\ 1 & \lambda_2 & \lambda_2^2 & \dots & \lambda_2^{N-1} \\ 1 & \lambda_3 & \lambda_3^2 & \dots & \lambda_3^{N-1} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \lambda_N & \lambda_N^2 & \dots & \lambda_N^{N-1} \end{vmatrix}$$

Ex: $N=3$

$$\int_{\mathcal{M}} f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$$N^2 \longrightarrow N \quad \Delta$$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$\prod_{i < j} (\lambda_i - \lambda_j)$$

$$= \det$$

$$\begin{vmatrix} 1 & \lambda_1 & \lambda_1^2 & \dots & \lambda_1^{N-1} \\ 1 & \lambda_2 & \lambda_2^2 & \dots & \lambda_2^{N-1} \\ 1 & \lambda_3 & \lambda_3^2 & \dots & \lambda_3^{N-1} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \lambda_N & \lambda_N^2 & \dots & \lambda_N^{N-1} \end{vmatrix}$$

$$\begin{matrix} \lambda_1^{N-1} \\ \lambda_2^{N-1} \\ \lambda_3^{N-1} \\ \vdots \\ \lambda_N^{N-1} \end{matrix}$$

Ex: $N=3$

$$\prod_{i < j} (\lambda_i - \lambda_j) = \det \lambda_i^{j-1}$$

PROOF: LHS and RHS are both Polynomials with the same degree, zeros and large λ_i behavior. Hence they are the same.

$$\prod_{i < j} (\lambda_i - \lambda_j) = \det \lambda_i^{j-1}$$

PROOF: LHS and RHS are both Polynomials with the same degree, zeros and large λ_i behavior. Hence they are the same

det = 0 when $\lambda_i = \lambda_j$ (=) } 2 equal rows

$$\prod_{i < j} (\lambda_i - \lambda_j) = \det \lambda_i^{j-1}$$

PROOF: LHS and RHS are both Polynomials with the same degree, zeros and large λ_i behavior. Hence they are the same

det = 0 when $\lambda_i = \lambda_j$ (=) } 2 equal rows

$\lambda_1 \rightarrow \infty$ LHS

$$\prod_{i < j} (\lambda_i - \lambda_j) = \det \lambda_i^{j-1}$$

PROOF: LHS and RHS are both Polynomials with the same degree, zeros and large λ_i behavior. Hence they are the same

det = 0 when $\lambda_i = \lambda_j$ (=) } zeros

$$\lambda_1 \rightarrow \infty \quad \text{LHS} \quad \lambda_1^{N-1} \prod_{\substack{i < j \\ i \neq 1}} (\lambda_i - \lambda_j) = \Delta_{N-1}(\lambda_2, \dots, \lambda_N)$$

$$\prod_{i < j} (\lambda_i - \lambda_j) = \det \lambda_i^{j-1}$$

PROOF: LHS and RHS are both Polynomials with the same degree, zeros and large λ_i behavior. Hence they are the same.

det = 0 when $\lambda_i = \lambda_j$ (=) } zeros

$\lambda_1 \rightarrow \infty$ LHS $\lambda_1^{N-1} \prod_{i < j} (\lambda_i - \lambda_j)$

RHS $\begin{vmatrix} \lambda_1^{N-1} & \lambda_1^{N-2} & \dots & \lambda_1 \\ \lambda_2^{N-1} & \lambda_2^{N-2} & \dots & \lambda_2 \\ \vdots & \vdots & \ddots & \vdots \\ \lambda_N^{N-1} & \lambda_N^{N-2} & \dots & \lambda_N \end{vmatrix} = \lambda_1^{N-1} \Delta_{N-1}(\lambda_2, \dots, \lambda_N)$

QED

$$\left\langle \frac{1}{N^k} \text{tr}(e^M) \right\rangle = \frac{1}{Z} \int \prod_{i=1}^N d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 \underbrace{e^{\lambda_1} \dots e^{\lambda_N}}_{\frac{1}{N} \sum_i e^{\lambda_i}} e^{-\frac{2}{g^2} \sum_{i=1}^N \lambda_i^2}$$

for Atoms Molecules

Goal $\langle \frac{1}{N} \text{tr}(M) \rangle = \frac{1}{Z} \int \prod_{i=1}^N d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 e^{\lambda_1} e^{-\frac{2}{g^2} \sum_{i=1}^N \lambda_i^2}$

$\underbrace{\quad}_{\frac{1}{N} \sum_i e^{\lambda_i}}$

$$\gg = \frac{1}{Z} \int \prod_{i=1}^N \pi d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 e^{\lambda_1} e^{-\frac{2}{g^2} \sum_{i=1}^N \lambda_i^2}$$

$$= \frac{1}{Z} \int \prod_{i=1}^N \pi d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 e^{\frac{1}{g^2} \sum_{i=1}^N \lambda_i} e^{-\sum_{i=1}^N \lambda_i^2}$$

$$\int \mathcal{D}M f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

$$N^2 \longrightarrow N \quad \text{!}$$

We can replace

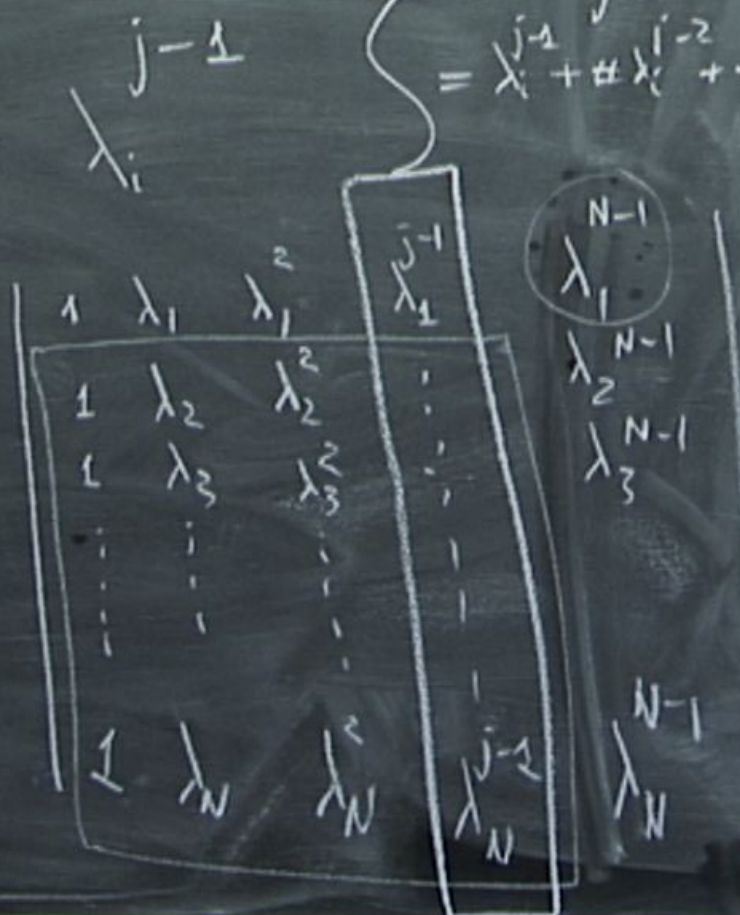
$$\begin{aligned} \lambda_i^{j-1} &\rightarrow P_{j-2}(\lambda_i) \\ &= \lambda_i^{j-1} + \# \lambda_i^{j-2} + \dots \end{aligned}$$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$\prod_{i < j} (\lambda_i - \lambda_j)$$

$$= \det$$

Ex: $N=3$



$\prod_{i < j} (\lambda_i - \lambda_j)$
 PROOF
 deg
 det
 λ_1

$$\int dM f(M) = \int d^N \lambda \Delta^2(\lambda) f(\lambda)$$

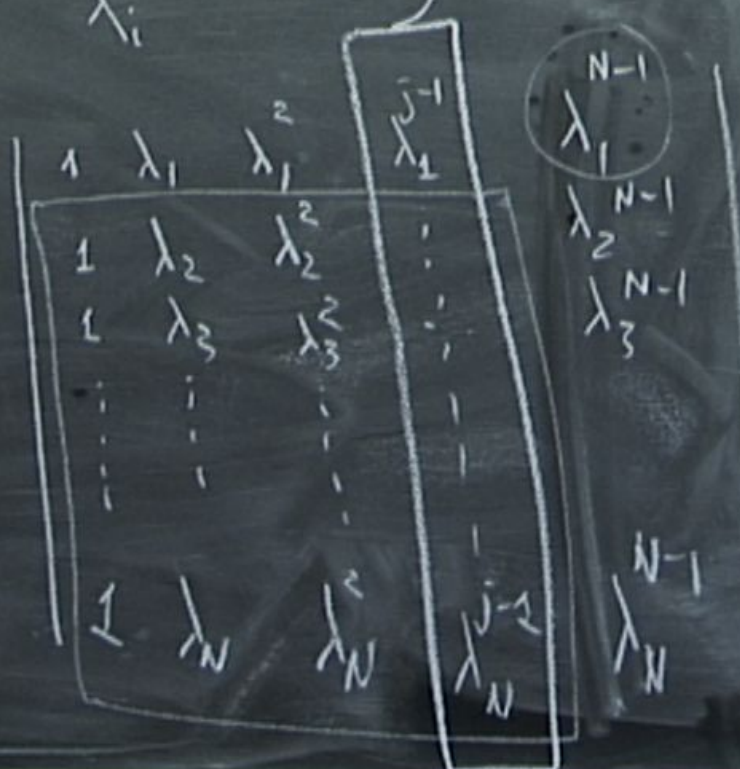
$$N^2 \longrightarrow N \quad \text{!}$$

$$\Delta(\lambda) = \det_{ij} \lambda_i^{j-1}$$

$$\prod_{i < j} (\lambda_i - \lambda_j)$$

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Ex: $N=3$



$\prod_{i < j} (\lambda_i - \lambda_j)$
 PROOF
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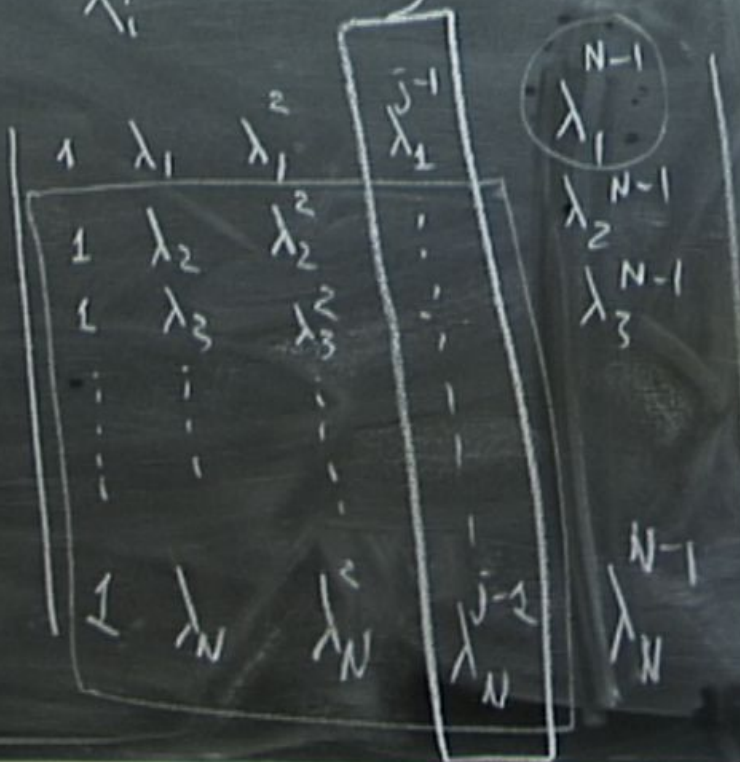
$$\begin{vmatrix} 1 & \lambda_1 & \lambda_1^2 \\ 1 & \lambda_2 & \lambda_2^2 \\ 1 & \lambda_3 & \lambda_3^2 \end{vmatrix} = \begin{vmatrix} 1 & \lambda_1 & \lambda_1^2 + \alpha\lambda_1 + \beta \\ 1 & \lambda_2 & \lambda_2^2 + \alpha\lambda_2 + \beta \\ 1 & \lambda_3 & \lambda_3^2 + \alpha\lambda_3 + \beta \end{vmatrix}$$

We can replace

$$\begin{aligned} \lambda_i^{j-1} &\rightarrow P_{j-2}(\lambda_i) \\ &= \lambda_i^{j-2} + \dots \end{aligned}$$

$$\begin{aligned} \Delta(\lambda) &= \det_{i,j} \lambda_i^{j-1} \\ &= \prod_{i < j} (\lambda_i - \lambda_j) \\ &= \det \end{aligned}$$

Ex: $N=3$



PROOF

deg

det

λ_1

Pollen to
Evidence
for Atoms

How
Big Is A
Molecule?

Goal $\langle \frac{1}{N!} P^M \rangle = \frac{1}{Z} \int \prod_{i=1}^N d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 e^{\lambda_1} e^{-\frac{2}{g^2} \sum_{i=1}^N \lambda_i^2}$

$= \frac{1}{Z} \int \prod_{i=1}^N d\lambda_i \left(\det_{i,j} \lambda_i^{j-1} \right)^2 e^{\frac{2}{g^2} \sum_{i=1}^N \lambda_i} e^{-\sum_{i=1}^N \lambda_i^2}$

$\left(\det P_{j-1}(\lambda_i) \right)^2$

$$\sqrt{2^n n! \sqrt{\pi}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$\left(\det_{i,j} P_i(\lambda_j) \right) \left(\det_{k,l} P_k(\lambda_l) \right)$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$\int dx_1 \dots dx_N \left(\det_{i,j} P_i(x_j) \right) \left(\det_{k,l} P_k(x_l) \right) \dots$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$f_{1, \dots, N} \left(\det_{i,j} P_i(\lambda_j) \right) \left(\det_{k,l} P_k(\lambda_l) \right) e^{-\lambda_1^2} e^{-\lambda_2^2} \dots e^{-\lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_2^2}{\sqrt{2}} g^2}$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$\int_{-\infty}^{+\infty} dx \dots \left(\det_{i,j} P_i(\lambda_j) \right) \left(\det_{k,l} P_k(\lambda_l) \right) e^{-\lambda_2^2} e^{-\lambda_3^2} \dots e^{-\lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_2^2}{\sqrt{2}} g^2}$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$\int_{-\infty}^{+\infty} dx_1 \dots dx_N \left(\det_{i,j} P_j(\lambda_i) \right) \left(\det_{k,l} P_k(\lambda_l) \right) e^{-\lambda_1^2 - \lambda_2^2 - \lambda_3^2 - \dots - \lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2}$$

$$= \sum_{j=0}^{N-1} \int dx_1 P_j(\lambda_1) e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2} \int \dots \int$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$= \sum_{j=0}^1 \frac{1}{N} e^{\frac{\lambda}{8N}} L_j^0 \left(-\frac{\lambda}{4N} \right)$$

$$= \frac{1}{N} L_{N-1}^1 \left(-\frac{\lambda}{4N} \right) e^{\frac{\lambda}{8N}}$$

$$\int \prod_{i,j} d\lambda_i P_j(\lambda_i) \left(\prod_{k,l} d\lambda_k P_l(\lambda_k) \right) e^{-\lambda_2^2 - \lambda_3^2 - \dots - \lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2}$$

$$= \sum_{j=0}^{N-1} \int d\lambda_j P_j(\lambda_j)^2 e^{-\lambda_j^2 + \frac{\lambda_j}{\sqrt{2}} g^2} \int \prod_{i=0}^{N-1} d\lambda_i$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} H_n(x)$$

$x^n + \dots$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$= \sum_{j=0}^N \frac{1}{N} e^{\frac{\lambda}{8N}} L_j^0 \left(-\frac{\lambda}{4N} \right)$$

$$= \frac{1}{N} L_{N-1}^1 \left(-\frac{\lambda}{4N} \right) e^{\frac{\lambda}{8N}}$$

$$\int_{-\infty}^{+\infty} dx_1 \dots dx_N \left(\det_{i,j} P_j(\lambda_i) \right) \left(\det_{k,l} P_k(\lambda_l) \right) e^{-\lambda_1^2 - \lambda_2^2 - \lambda_3^2 - \dots - \lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2}$$

$$= \sum_{j=0}^{N-1} \int dx_1 P_j(\lambda_1) e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2} \int dx_2 \dots \int dx_N$$

$$P_n(x) = \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} \underbrace{H_n(x)}_{x^n + \dots}$$

$$\int_{-\infty}^{+\infty} dx e^{-x^2} P_n(x) P_m(x) = \delta_{nm}$$

$$\begin{aligned} &= \sum_{j=0}^1 \frac{1}{N} e^{\frac{\lambda}{8N}} L_j^0 \left(-\frac{\lambda}{4N} \right) \\ &= \frac{1}{N} L_{N-1} \left(-\frac{\lambda}{4N} \right) e^{\frac{\lambda}{8N}} \\ &= \frac{2}{\sqrt{\lambda}} I_1(\sqrt{\lambda}) + \frac{1}{N^2} \frac{\lambda}{4\pi^2} I_2(\sqrt{\lambda}) + \dots \end{aligned}$$

$$\begin{aligned} & \int \dots \left(\det_{i,j} P_j(\lambda_i) \right) \left(\det_{k,l} P_k(\lambda_l) \right) e^{-\lambda_2^2 - \lambda_3^2 - \dots - \lambda_N^2} e^{-\lambda_1^2 + \frac{\lambda_1}{\sqrt{2}} g^2} \\ &= \sum_{j=0}^{N-1} \int d\lambda_j P_j(\lambda_j)^2 e^{-\lambda_j^2 + \frac{\lambda_j}{\sqrt{2}} g^2} \int \dots \int \dots \end{aligned}$$