

Title: Explorations in Quantum Information - Lecture 5

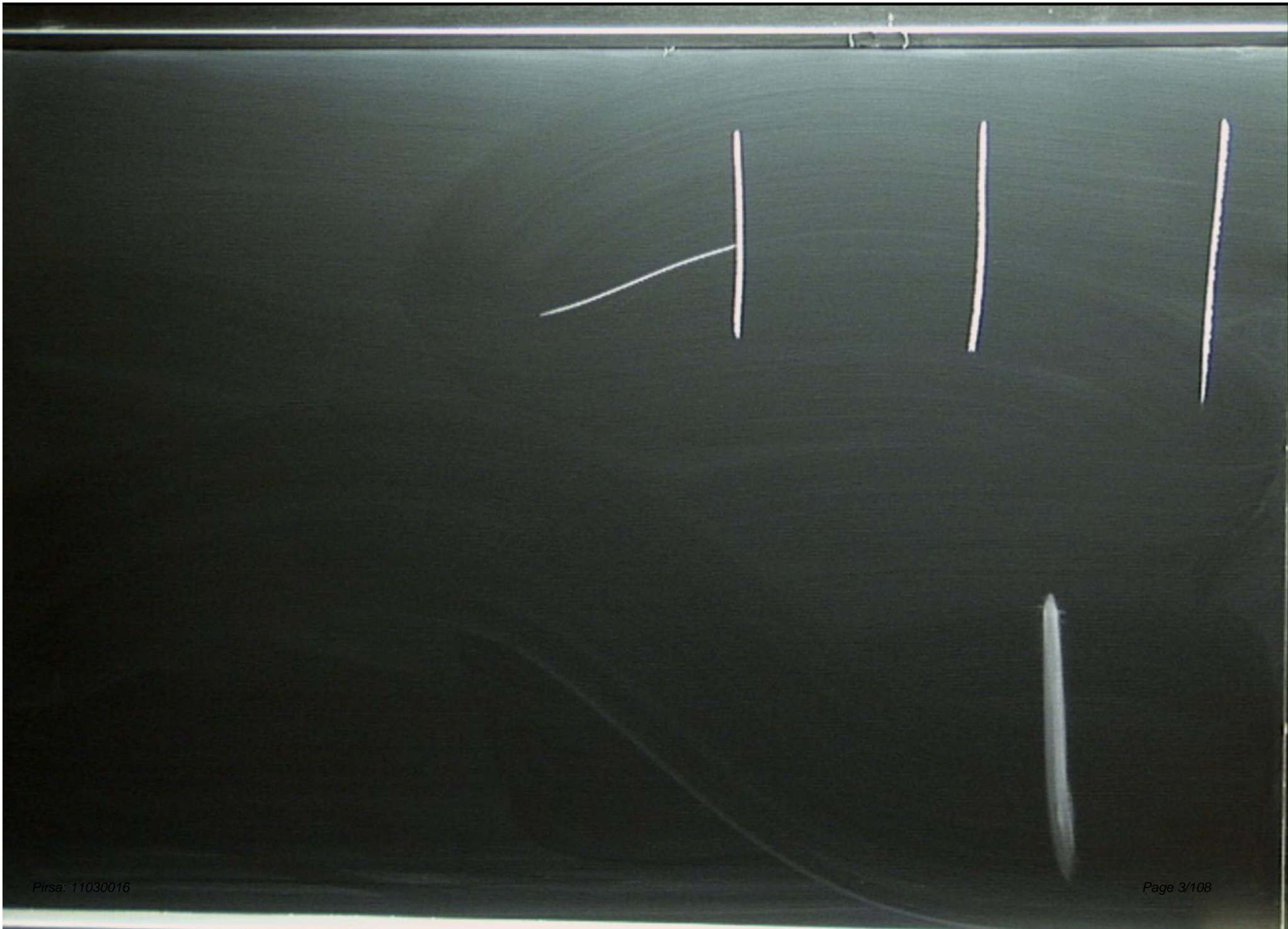
Date: Mar 18, 2011 09:00 AM

URL: <http://pirsa.org/11030016>

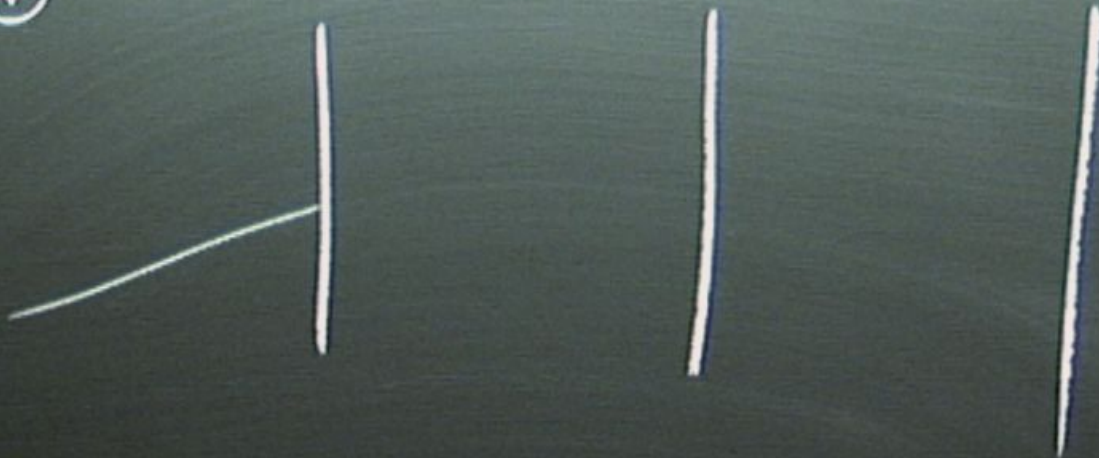
Abstract:



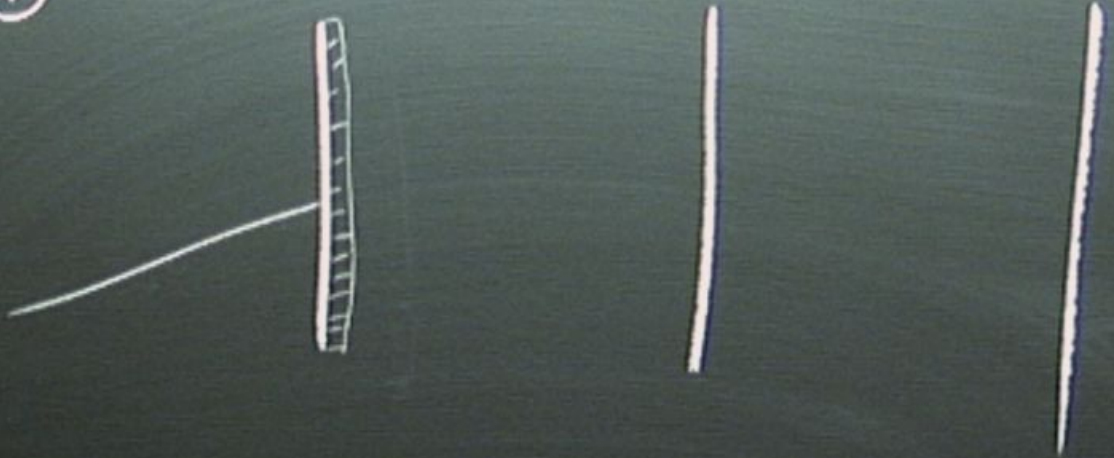
perimeter scholars
INTERNATIONAL



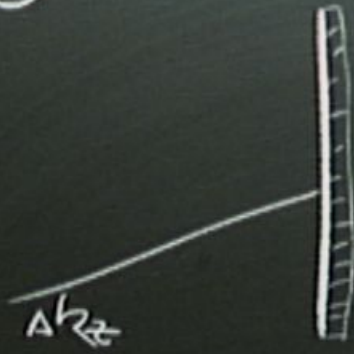
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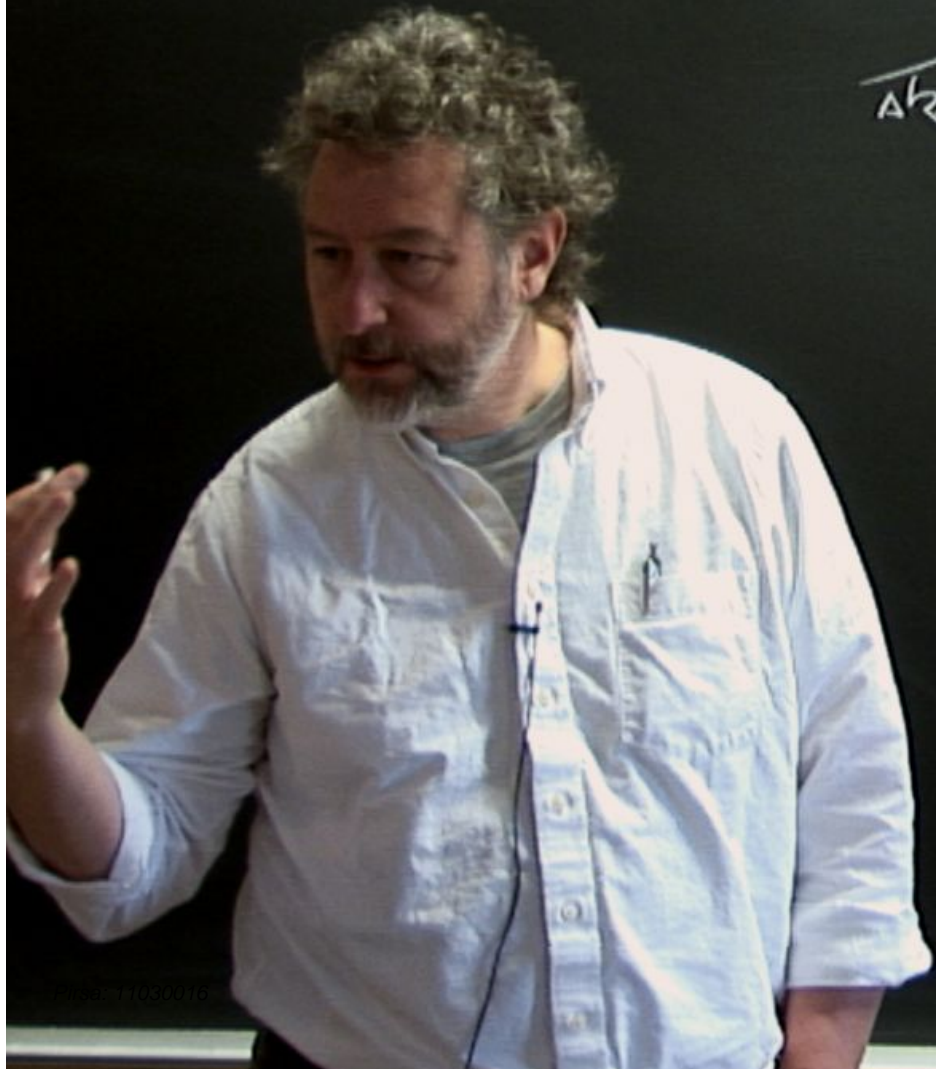
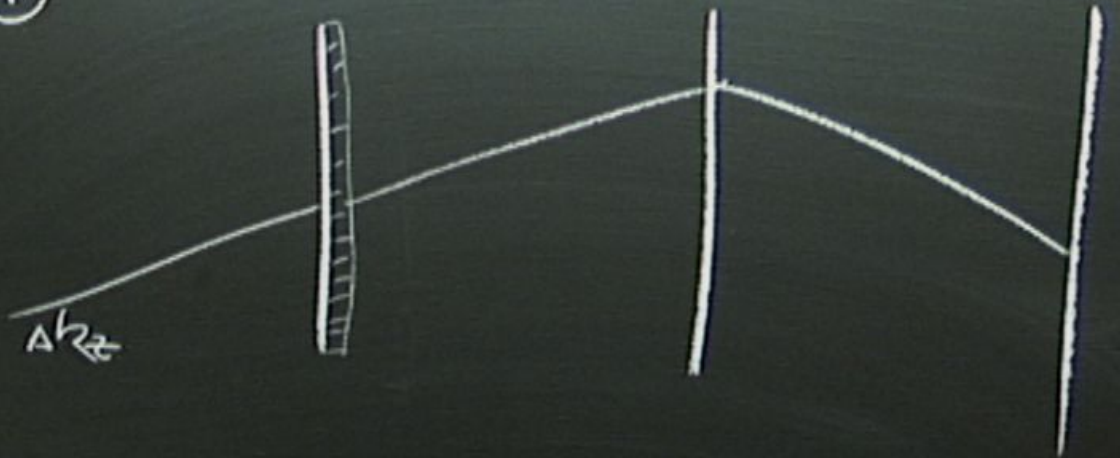
20



z①



20



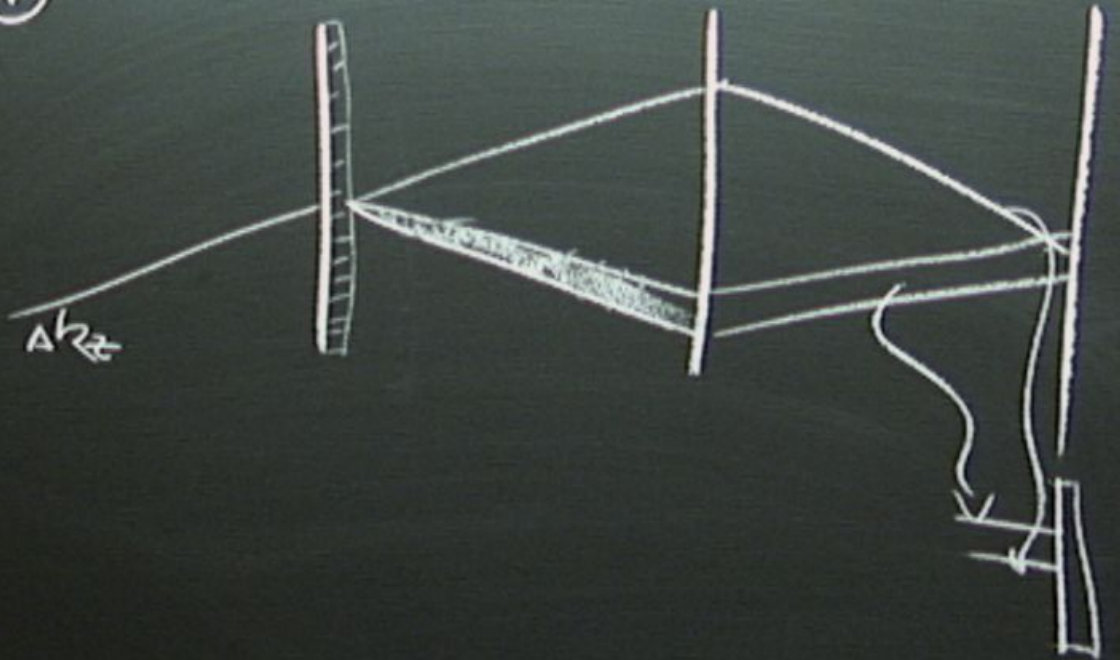
$z \rightarrow$



Δh_{22}



20

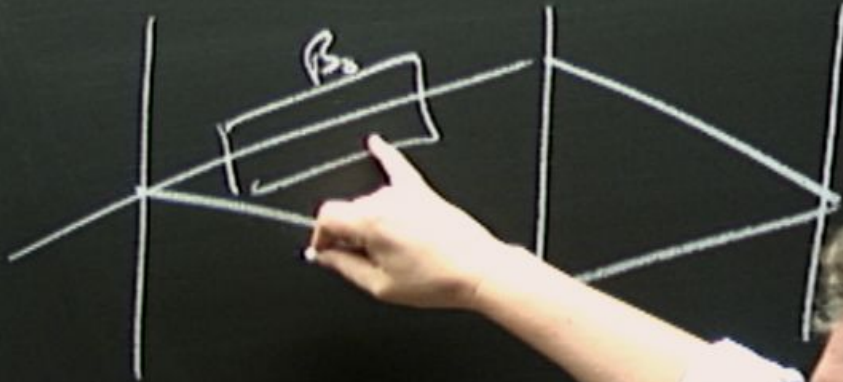


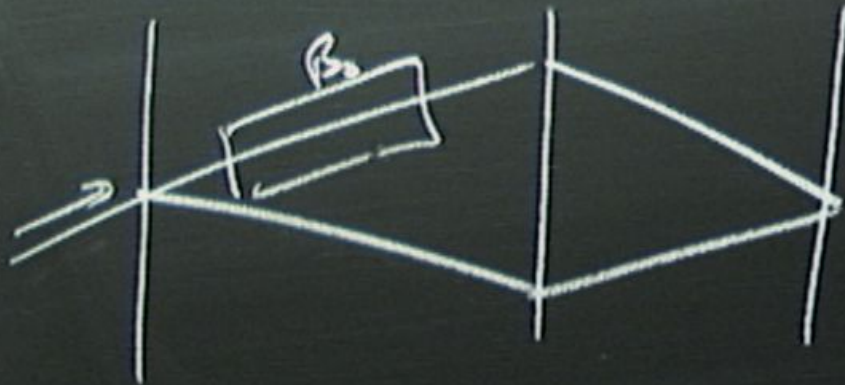
$$\nabla^2 \psi + K^2 \psi = 0$$

K

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{\sqrt{2m}}{\hbar} [E - V(x)]$$





$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$

$$\lambda = \frac{2\pi}{K}$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k}$$

↑
index of refraction

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

↑
index of refraction

$$\eta = 1 - \frac{\pi^2 B_c W}{2\pi}$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

↑ index of refraction

$$\eta = 1 - \frac{\pi^2 \beta_c W}{2\pi}$$

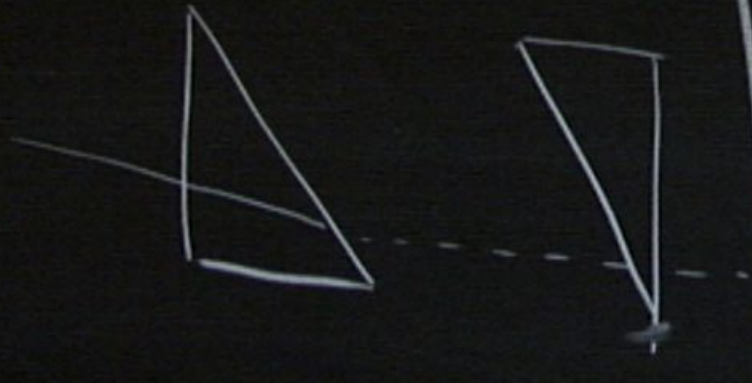
$\sim 10^{-5}$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$

$n = \frac{K}{k}$ ← Free space
 ↑
 index of refraction

$$\eta = 1 - \frac{\pi^2 B_c W}{2\pi} \sim 10^{-5}$$



$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

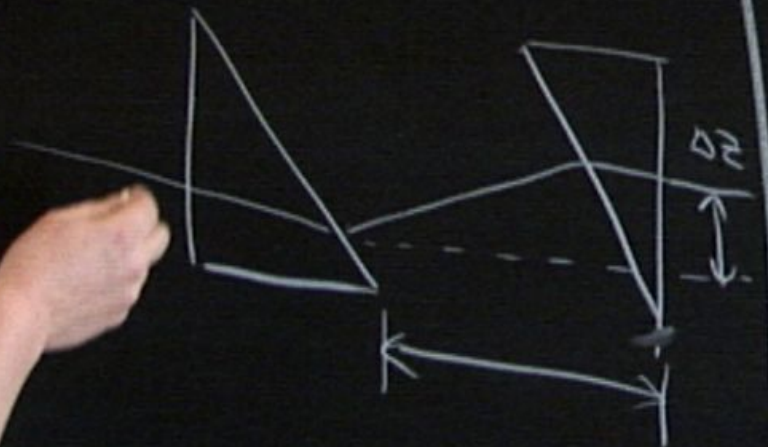
$$\eta = 1 - \frac{\pi^2 B_c W}{2\pi^2} \approx 10^{-5}$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

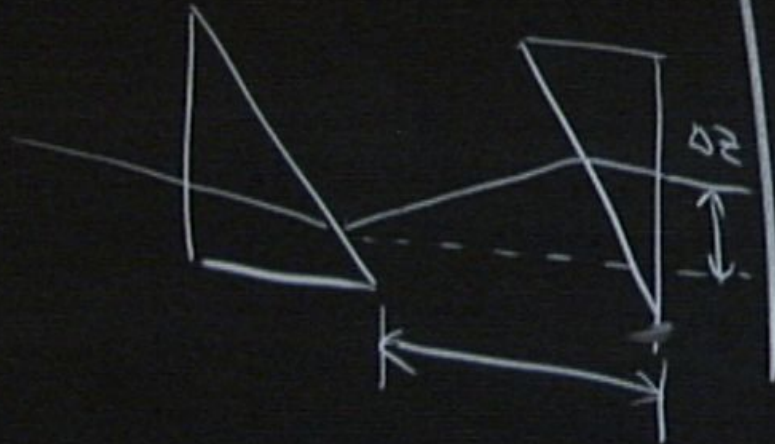


$$n = 1 - \frac{\pi^2 B_c W}{2\pi^2}$$

$\sim 10^{-5}$

$$\nabla^2 \psi + K^2 \psi = 0$$

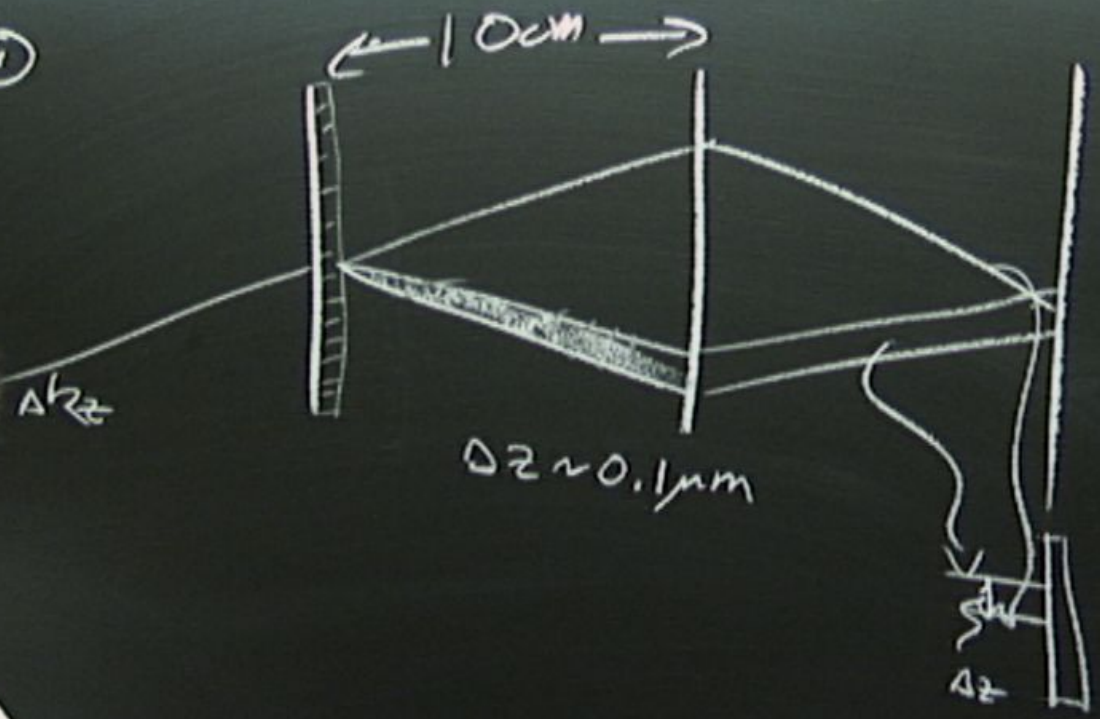
$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$



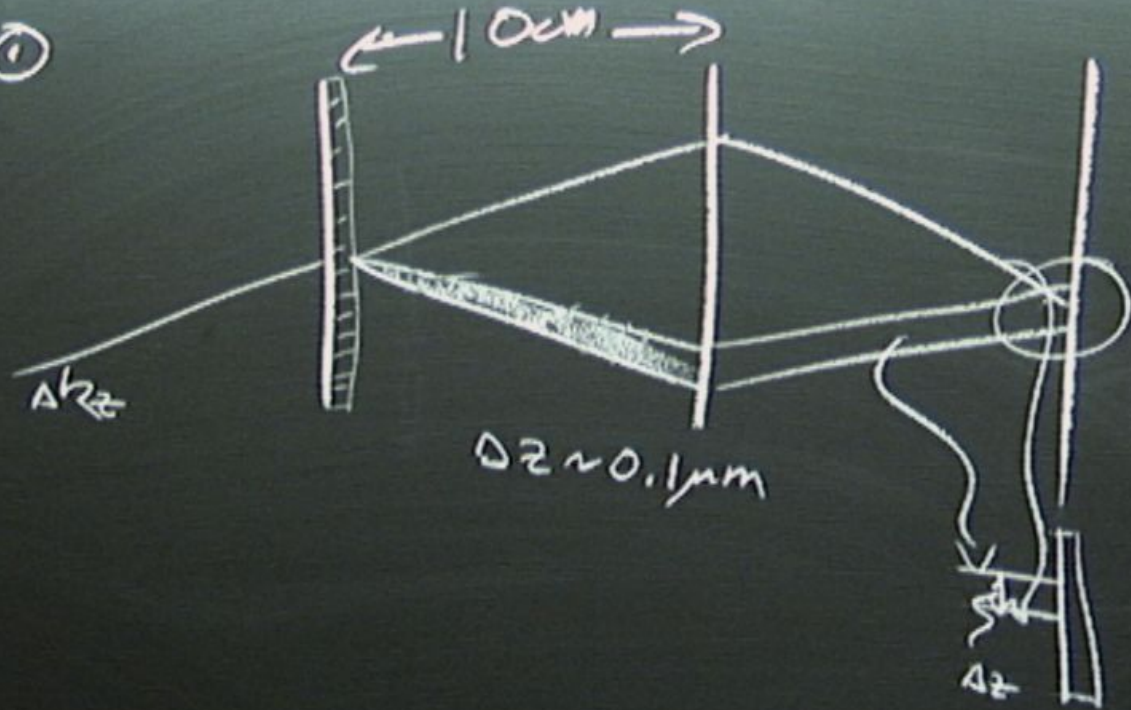
$$n = \frac{K}{k}$$

Free space
index of refraction

z ⊙

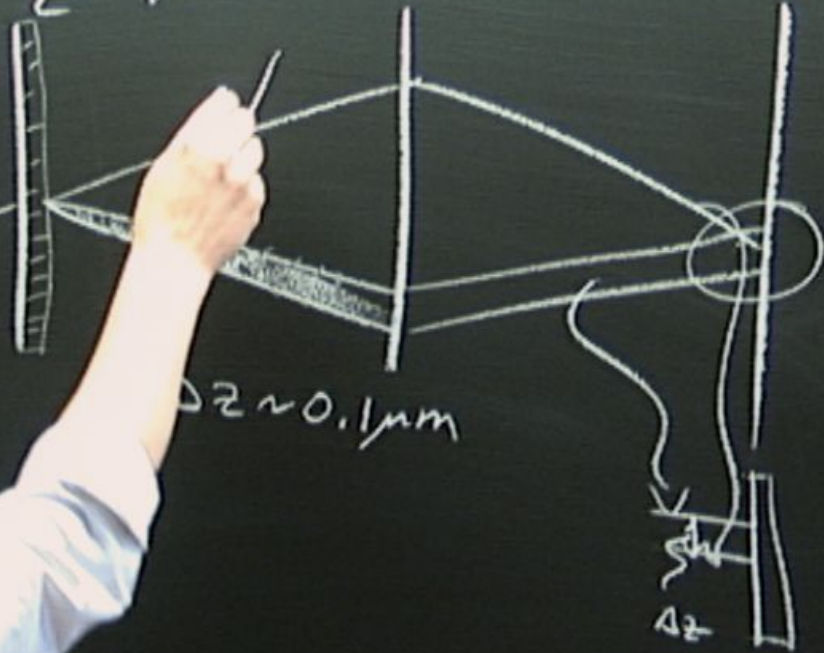


z ①

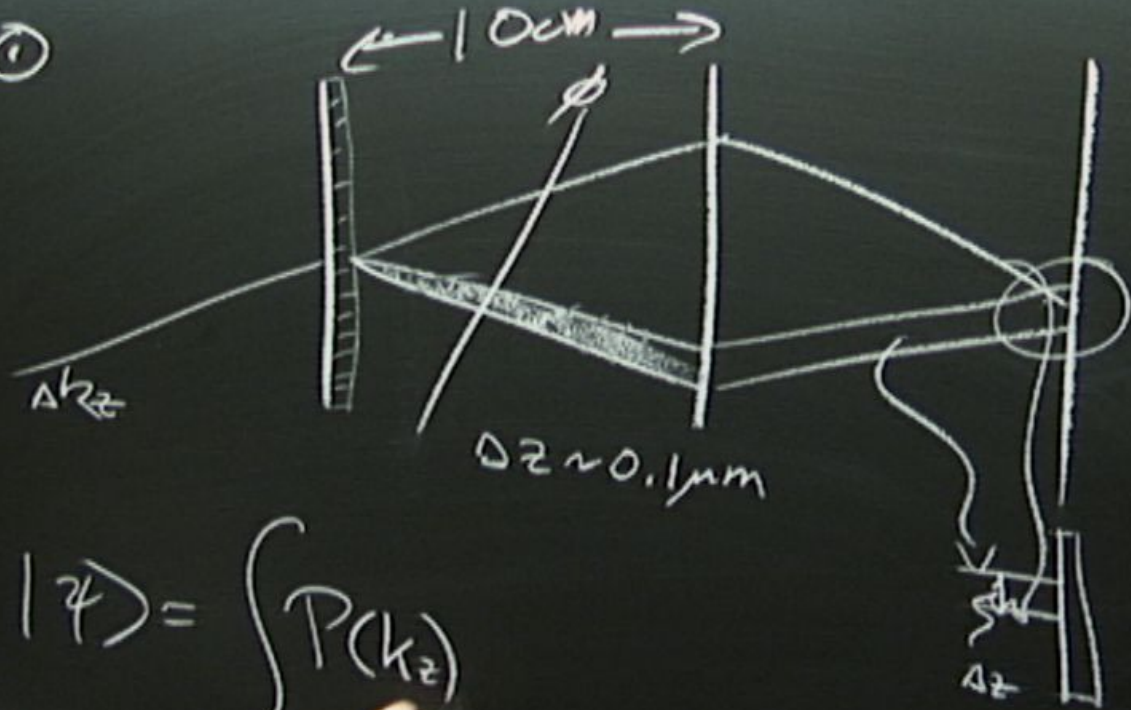


z ⊙

← 10cm →

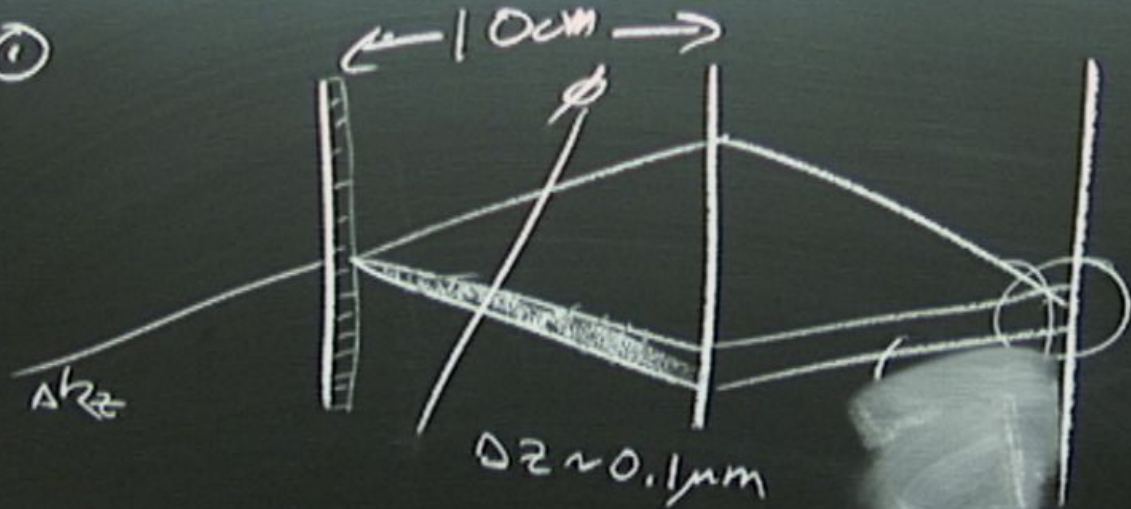


$z \rightarrow$



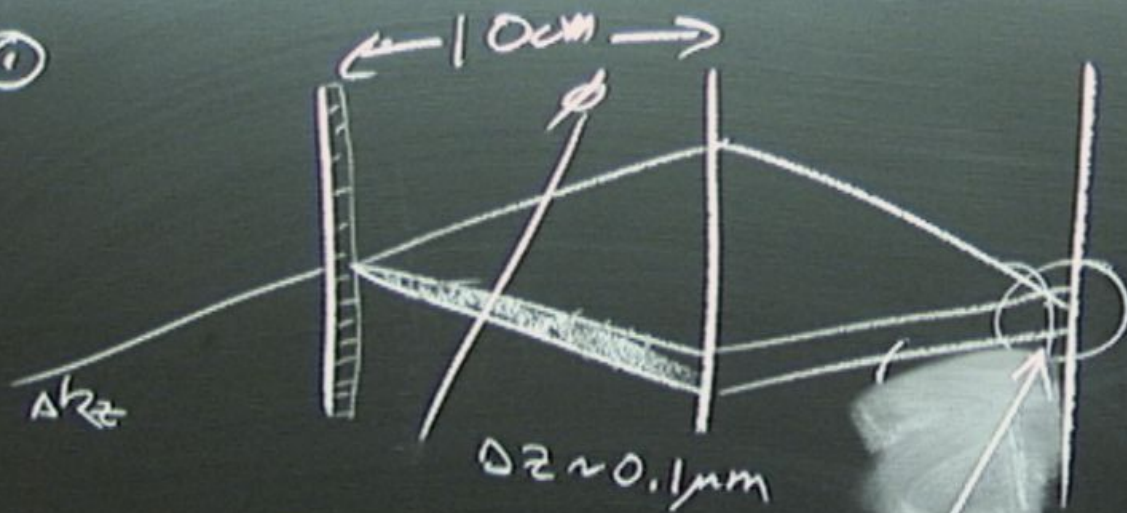
$$|\psi\rangle = \int P(k_z)$$

$z \rightarrow$



$$|\psi\rangle = \int P(k_z) \left[e^{i\phi} e^{-ik_z z_0} |0\rangle + e^{ik_z z_1} \right]$$

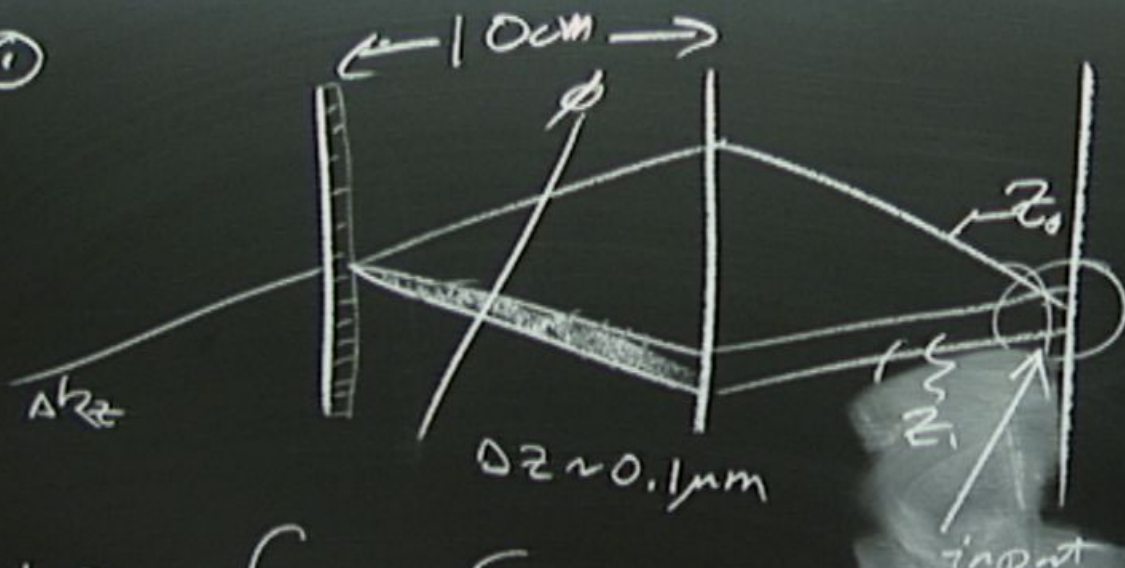
z ①



$$|\psi\rangle = \int P(k_z) \left[e^{i\varphi} e^{-ik_z z_0} |0\rangle + \right.$$

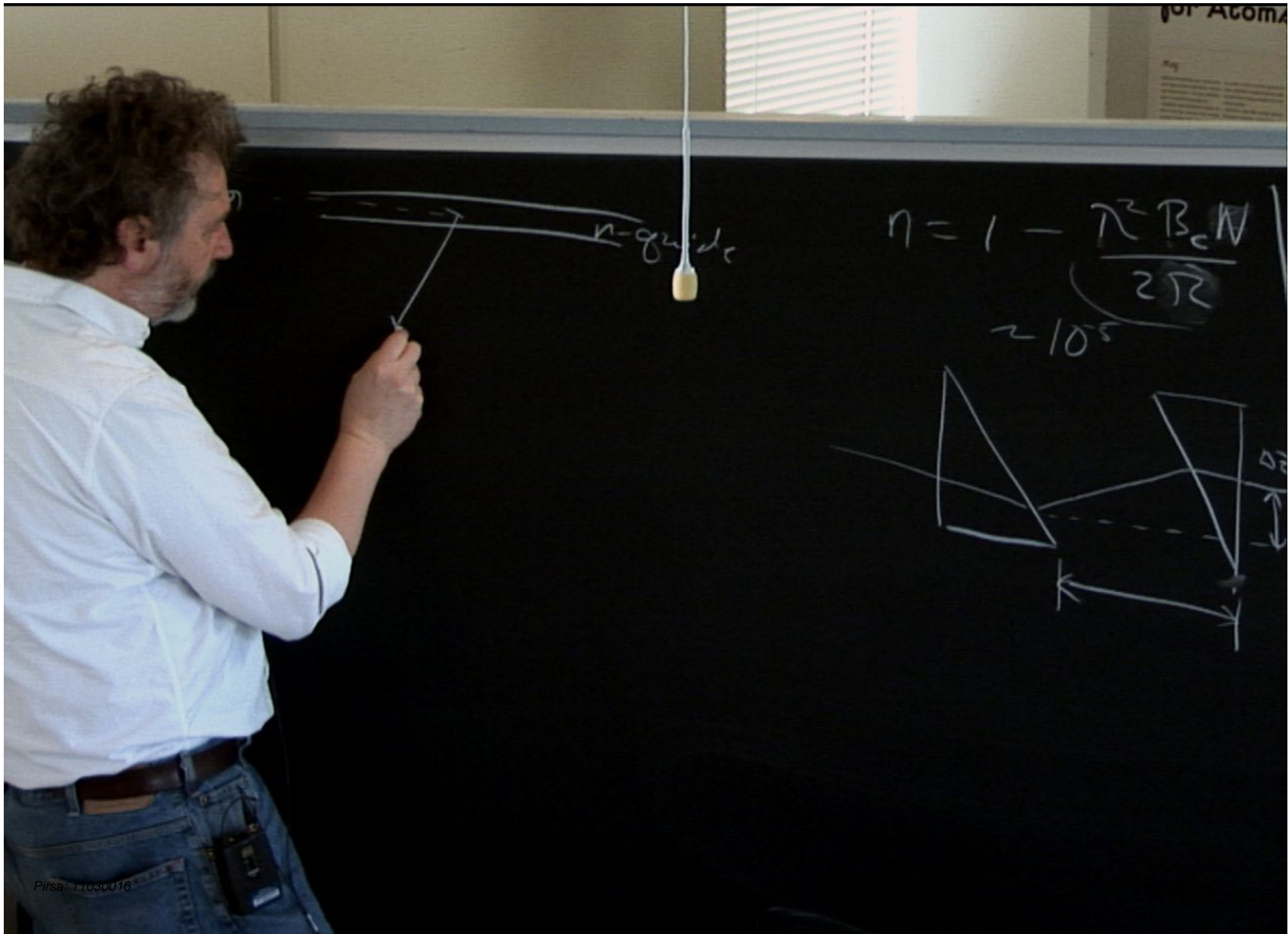
$$\left. e^{+ik_z z_1} |1\rangle \right] dk_z$$

z ①

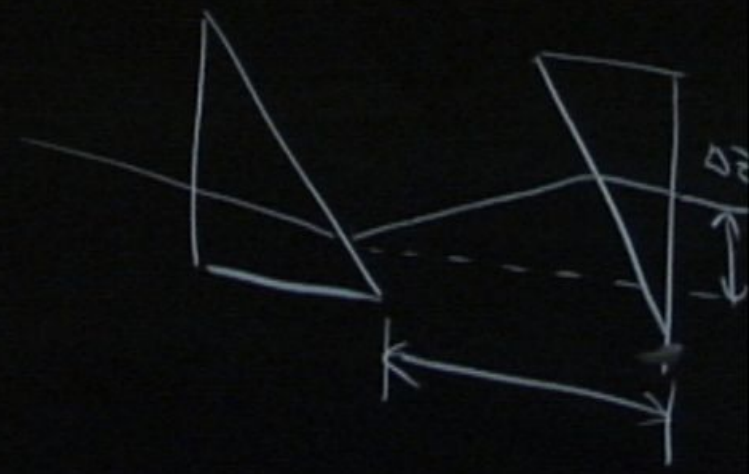


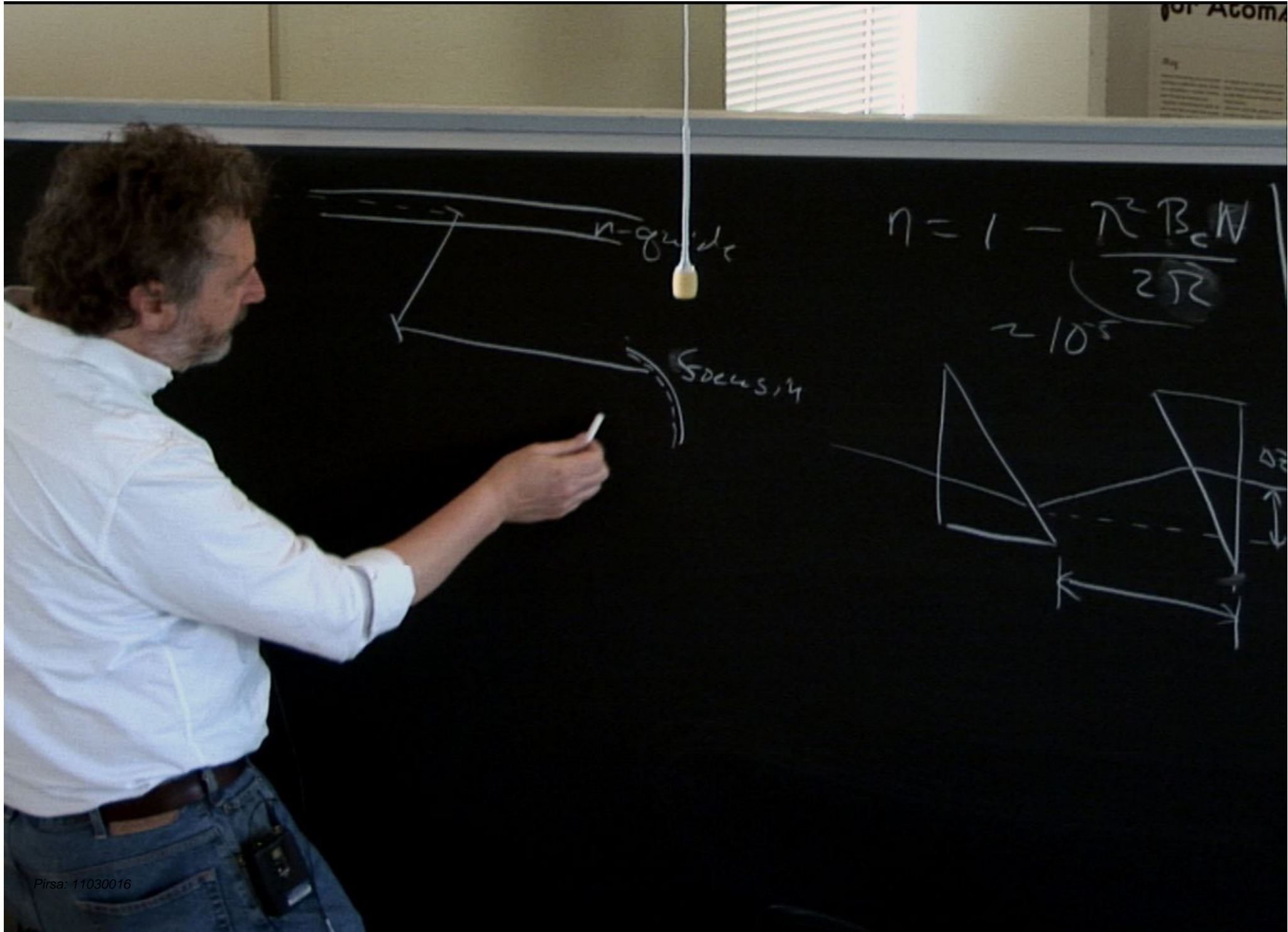
$$|\psi\rangle = \int P(k_z) \left[e^{i\varphi} e^{-ik_z z_0} |0\rangle + e^{ik_z z_1} |1\rangle \right] dk_z$$

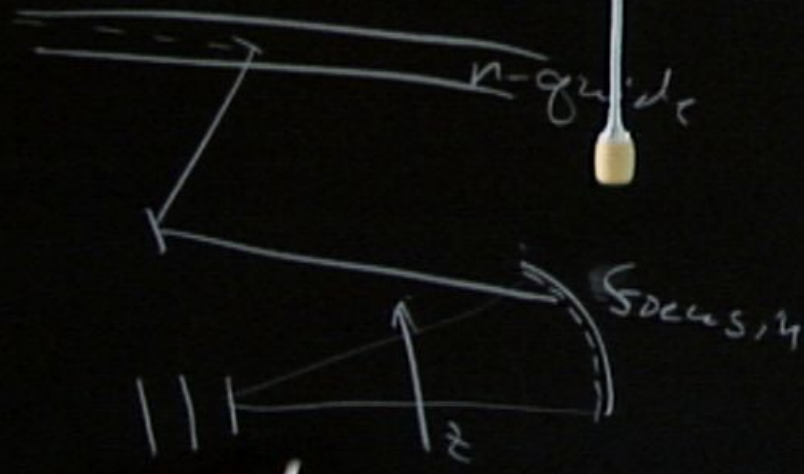
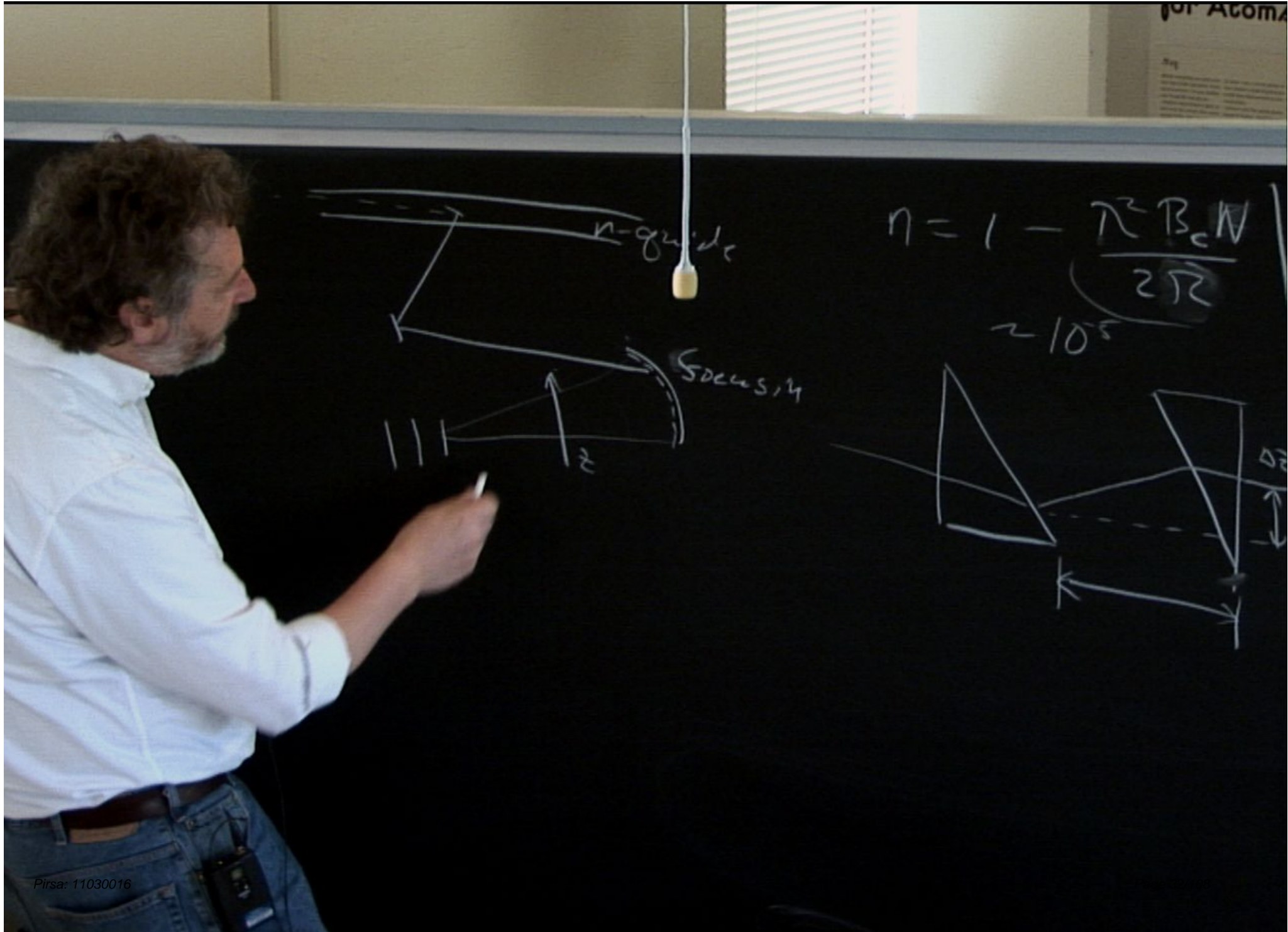
$$e^{ik_z z_1} |1\rangle dk_z$$



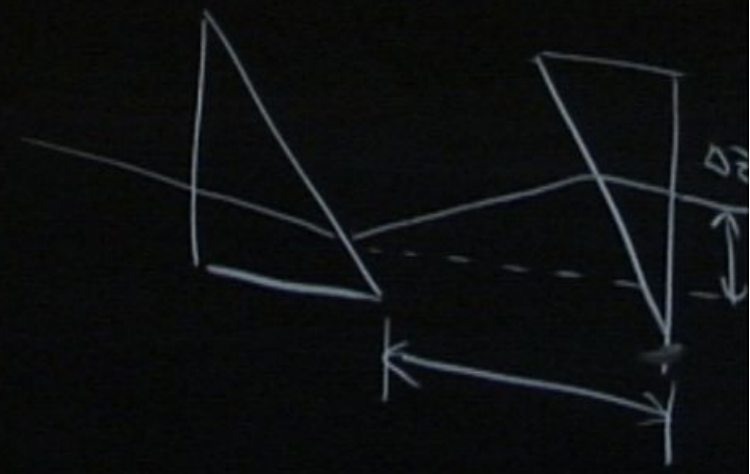
$$\eta = 1 - \frac{\pi^2 B_c W}{2L^2} \approx 10^{-5}$$



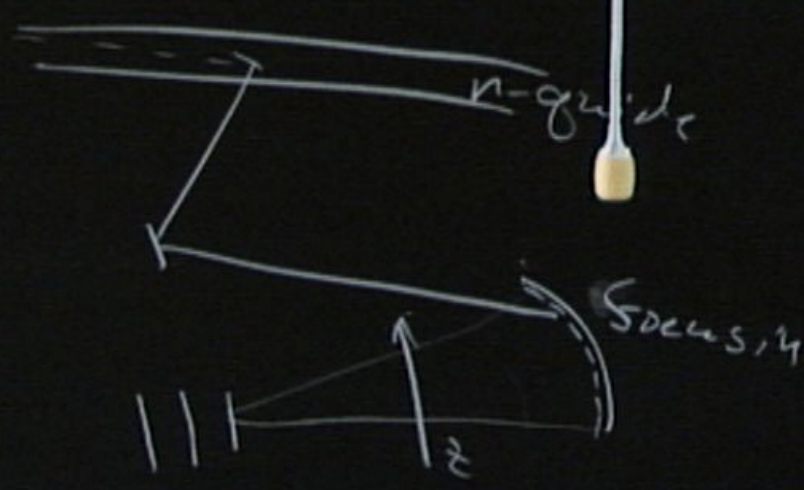




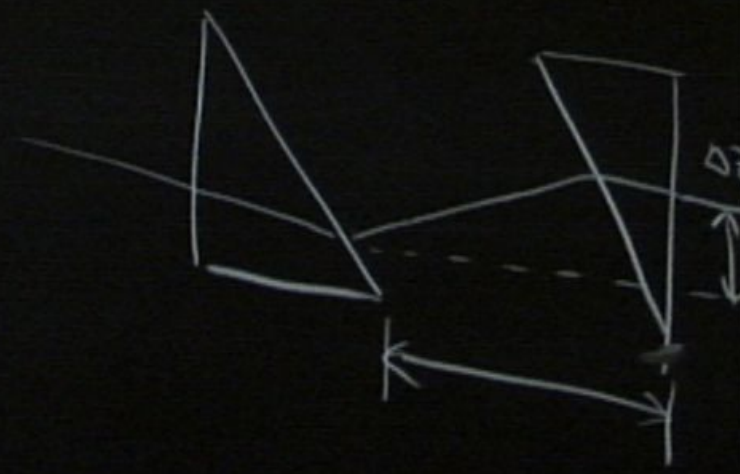
$$\eta = 1 - \frac{\pi^2 B_c W}{2L^2} \approx 10^{-5}$$



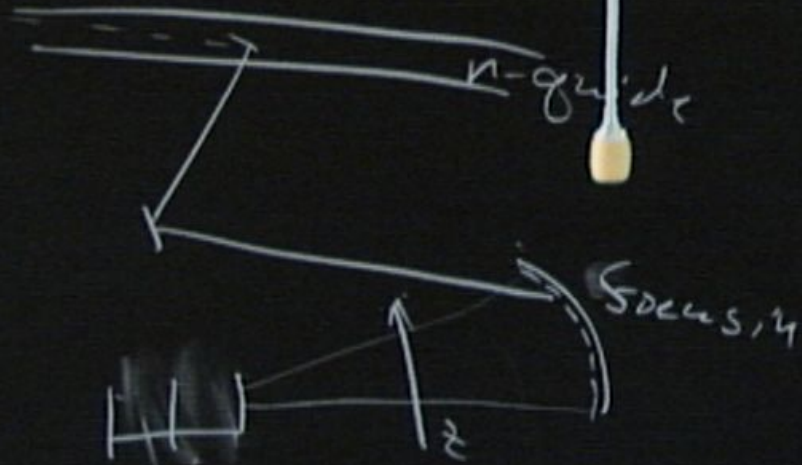
Results



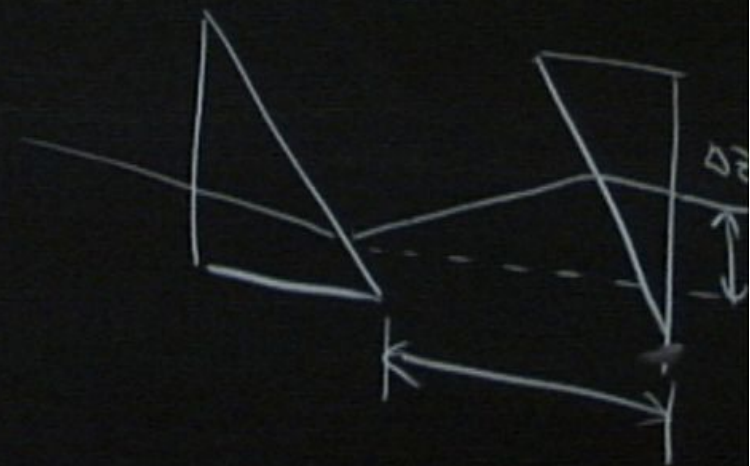
$$\eta = 1 - \frac{\pi^2 B_c W}{2\pi^2} \sim 10^{-5}$$

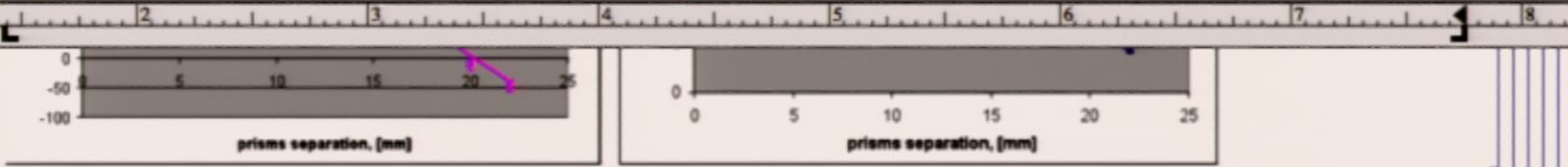


Resonator



$$\eta = 1 - \frac{\pi^2 B_c W}{2L^2} \approx 10^{-5}$$





point to remember here is that the contrast in the transmitted beam is lost due to a distribution of momenta. the neutron momenta in free space do not change, this can easily be recovered. Also note that we can compose any states of momenta through a series of such filters.

Experiment 16: coherence length with a focusing monochromator

distribution in k_z has two contributions: a set of focusing monochromator blades; and the thermal width of each

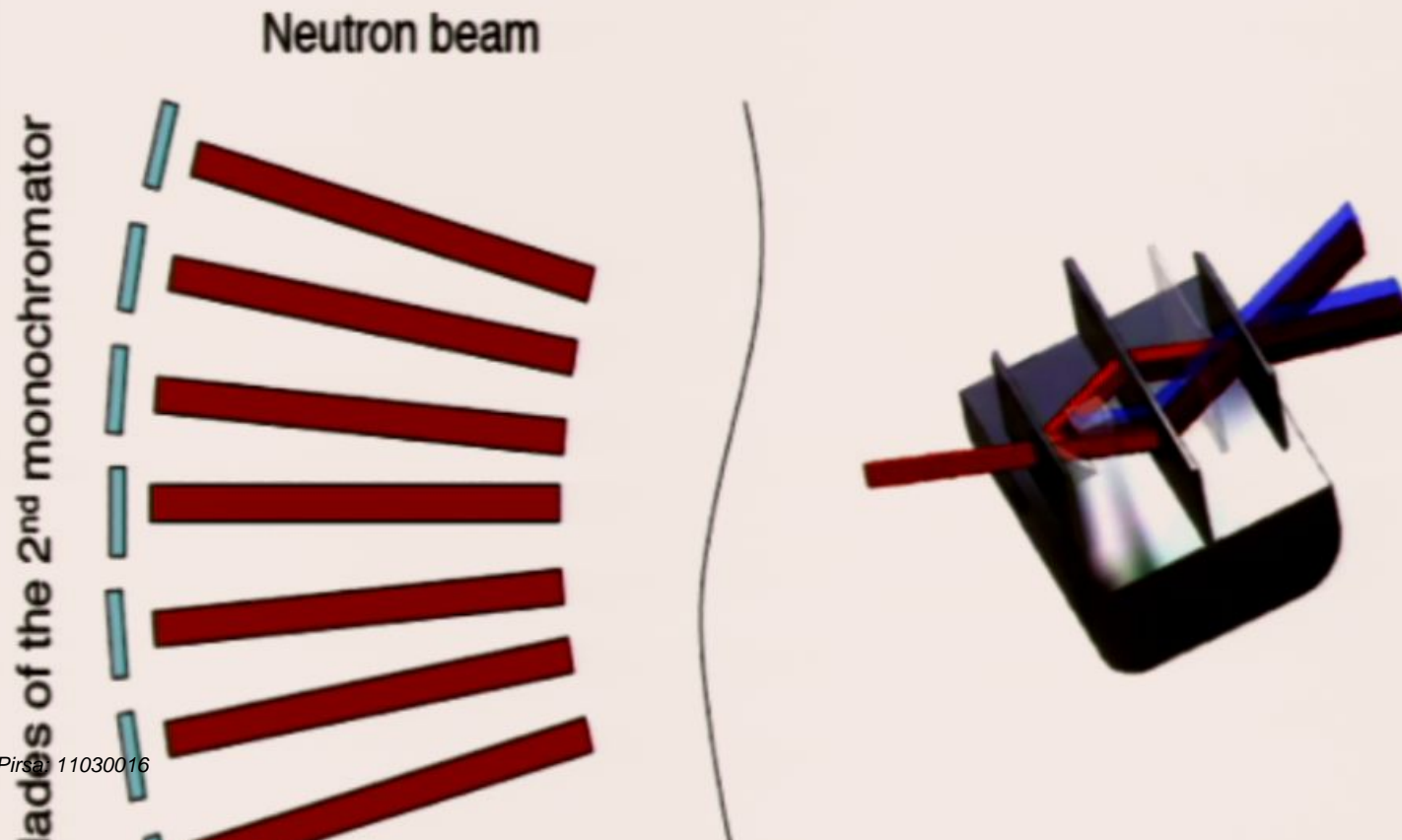
Neutron beam



the neutron momenta in free space do not change, this can easily be recovered. Also note that we can compose any states of momenta through a series of such filters.

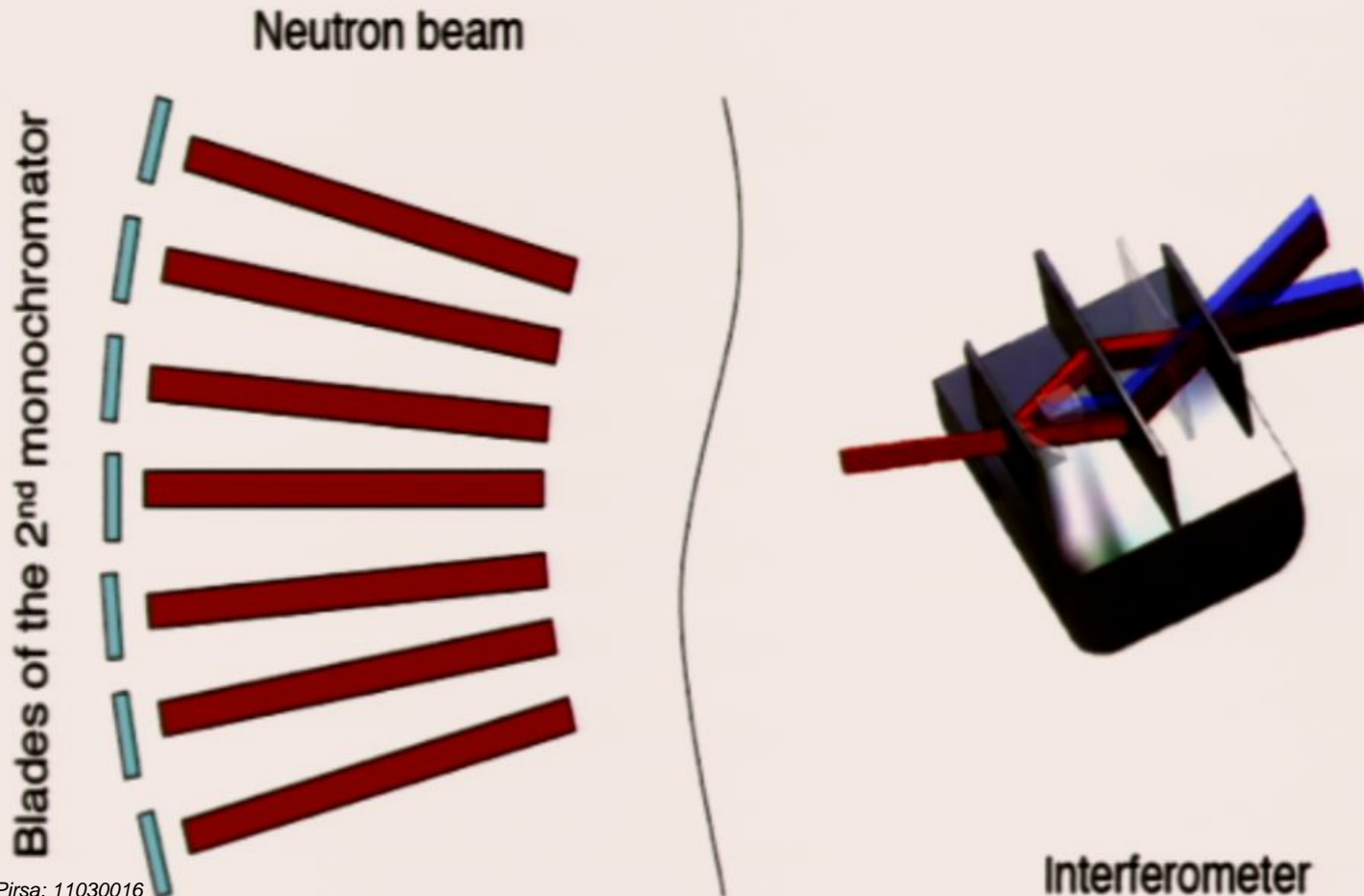
Experiment 16: coherence length with a focusing monochromator

distribution in k_z has two contributions: a set of focusing monochromator blades; and the thermal width of each



Experiment 16: coherence length with a focusing monochromator

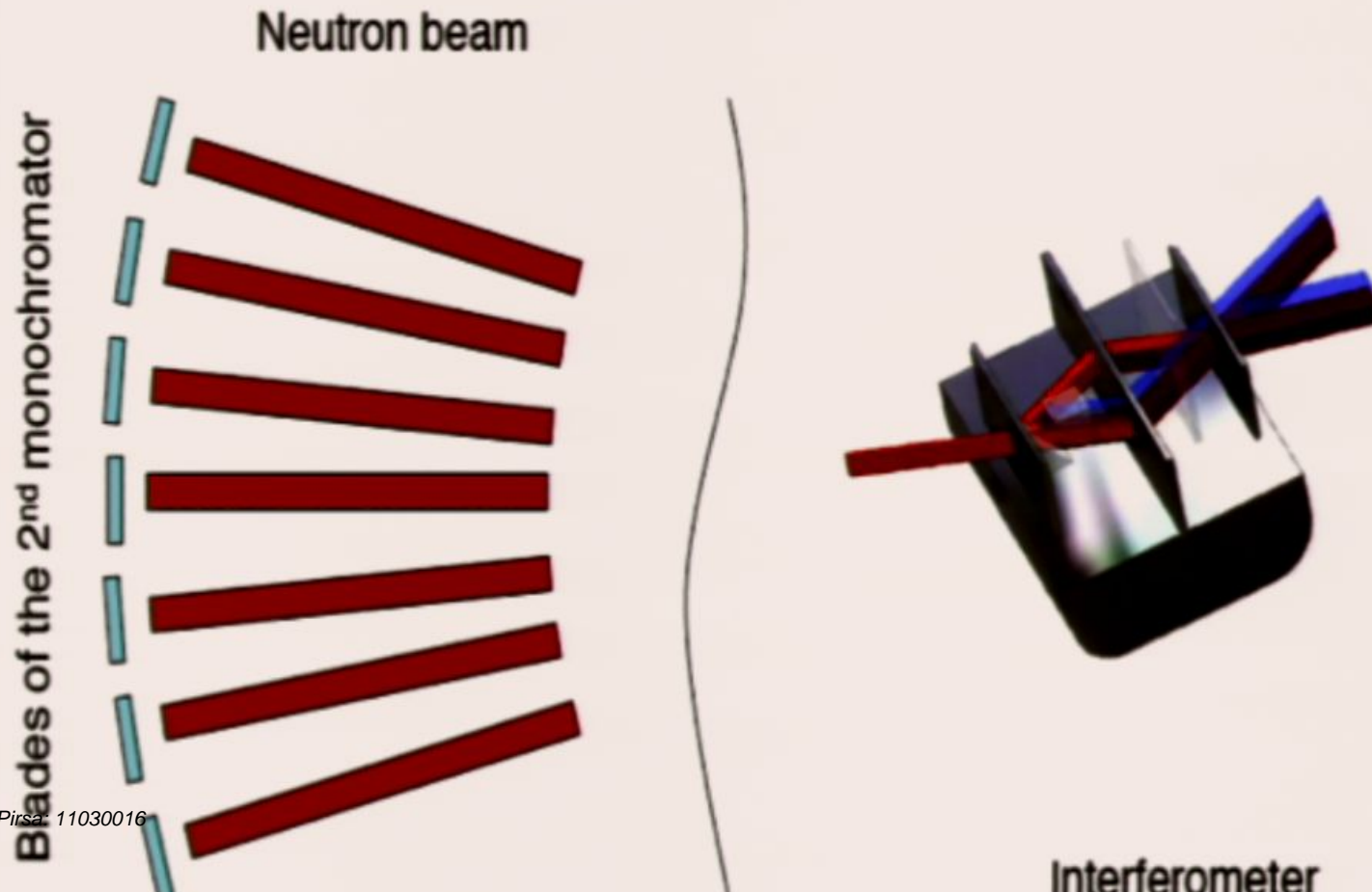
Distribution in k_z has two contributions: a set of focusing monochromator blades; and the thermal width of each

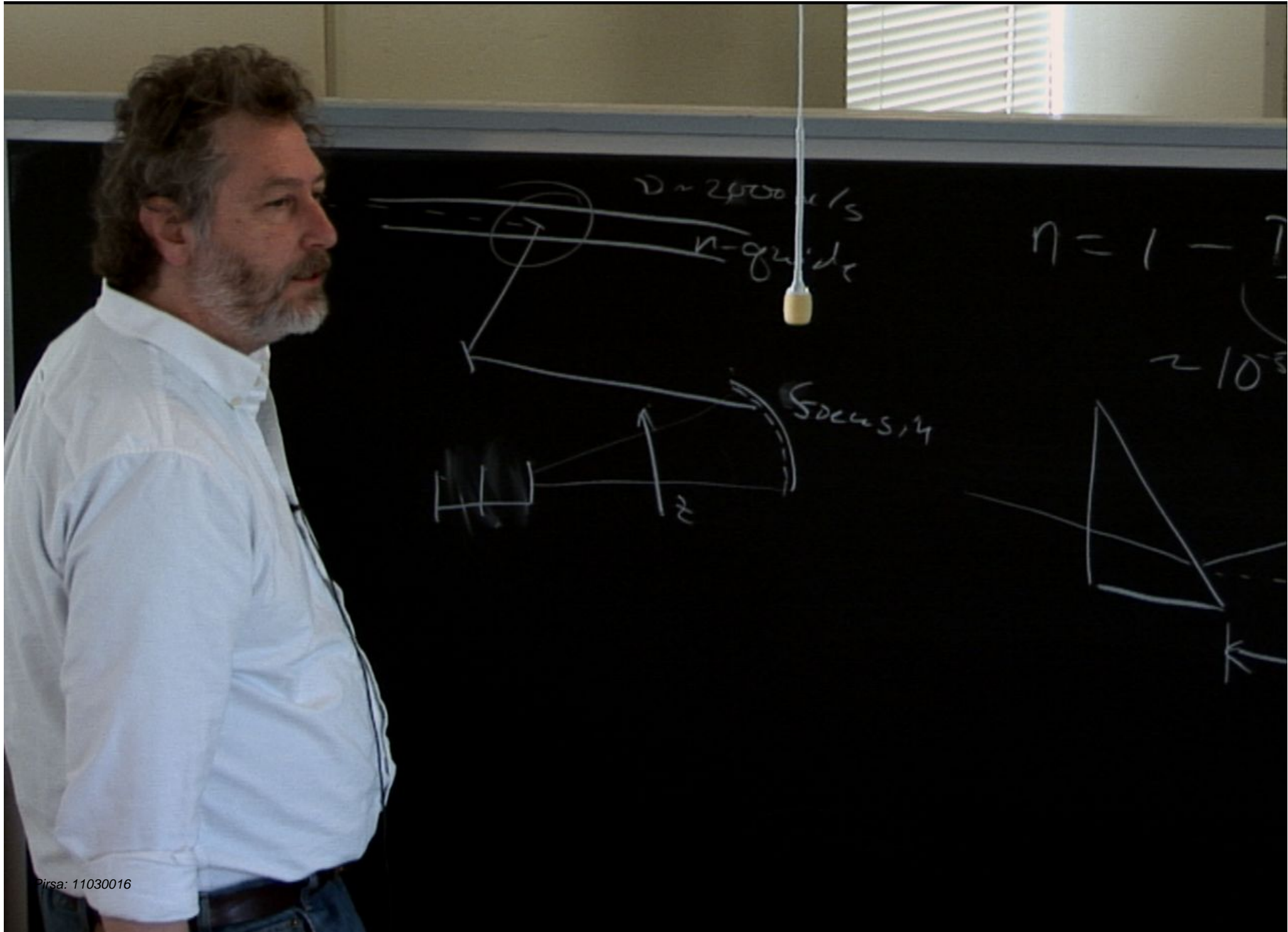


ary states of momentum through a series of such filters.

periment 16: coherence length with a focusing monochromator

istribution in k_z has two contributions: a set of focusing monochromator blades; and the thermal width of each



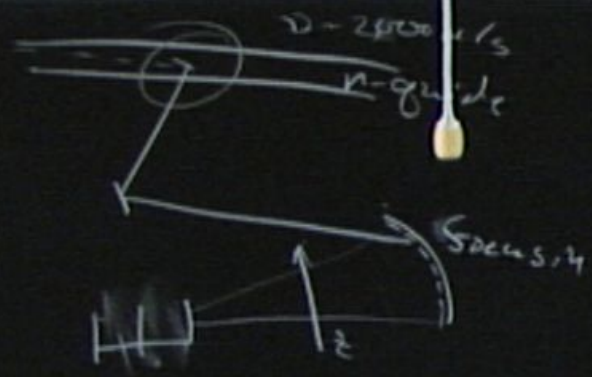


$v \sim 2000 \text{ m/s}$
n-guide

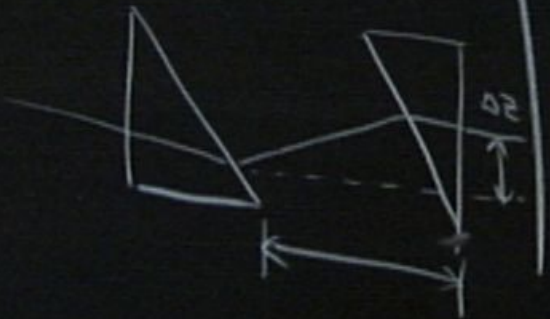
50000 m
z

$\eta = 1 - T$
 $\sim 10^{-5}$

electron



$$n = 1 - \frac{\pi^2 B_0^2 N}{2\epsilon^2} \sim 10^{-5}$$

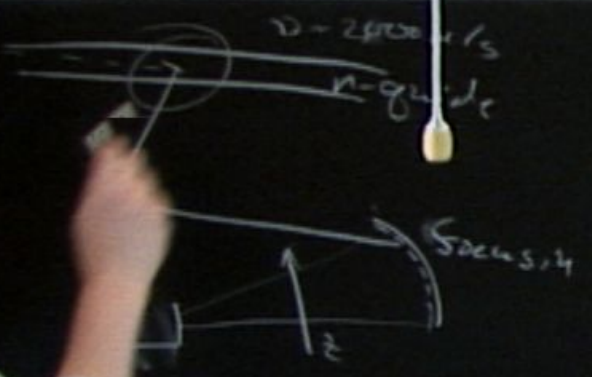


$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(r)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction



$$n = 1 - \frac{\mu^2 B_c N}{2R} \approx 10^{-5}$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

electron

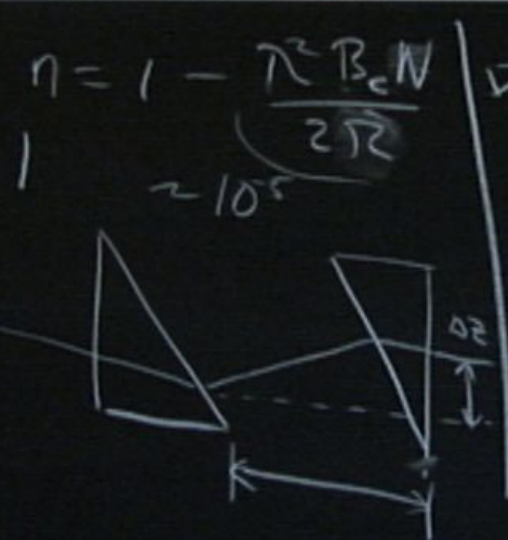
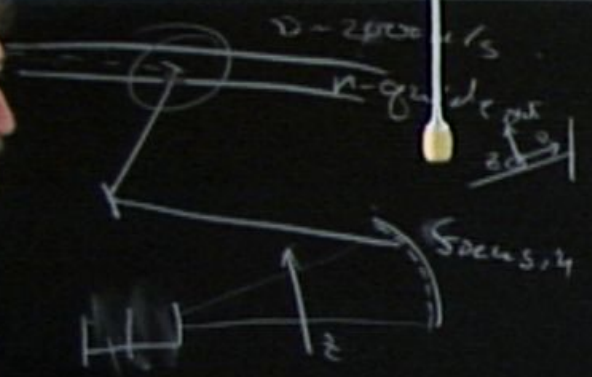
$V = 2000 \text{ eV}$
 $n\text{-guide}$
 500 nm
 z

$n = 1$

$\nabla^2 \psi + K^2 \psi = 0$

$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$

$n = \frac{K}{k} \leftarrow \text{Free space index of refraction}$



$$\eta = 1 - \frac{\hbar^2 B_c N}{2R} \approx 10^5$$

$$\nabla^2 \psi + K^2 \psi = 0$$
$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \begin{array}{l} \text{Free space} \\ \text{index of refraction} \end{array}$$

electron

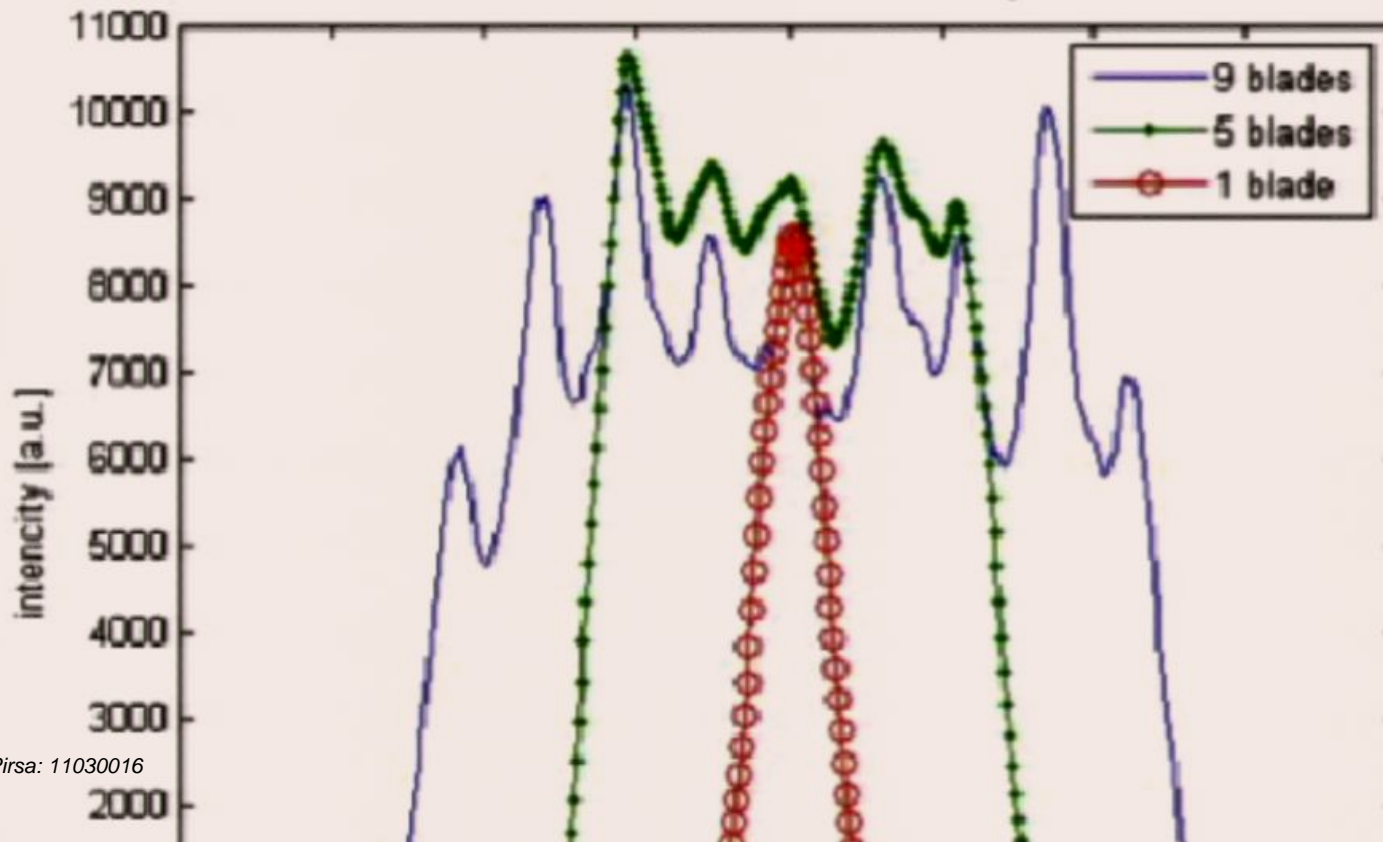
$v = 2\pi R / T$
 $r = \frac{mv}{eB}$
 $\omega = \frac{v}{r} = \frac{eB}{m}$

$n = 1 - \frac{\mu_0^2 B^2 N^2}{2R} \approx 10^5$

$\nabla^2 \psi + K^2 \psi = 0$
 $K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$
 $n = \frac{K}{k} \leftarrow$ Free space
 index of refraction

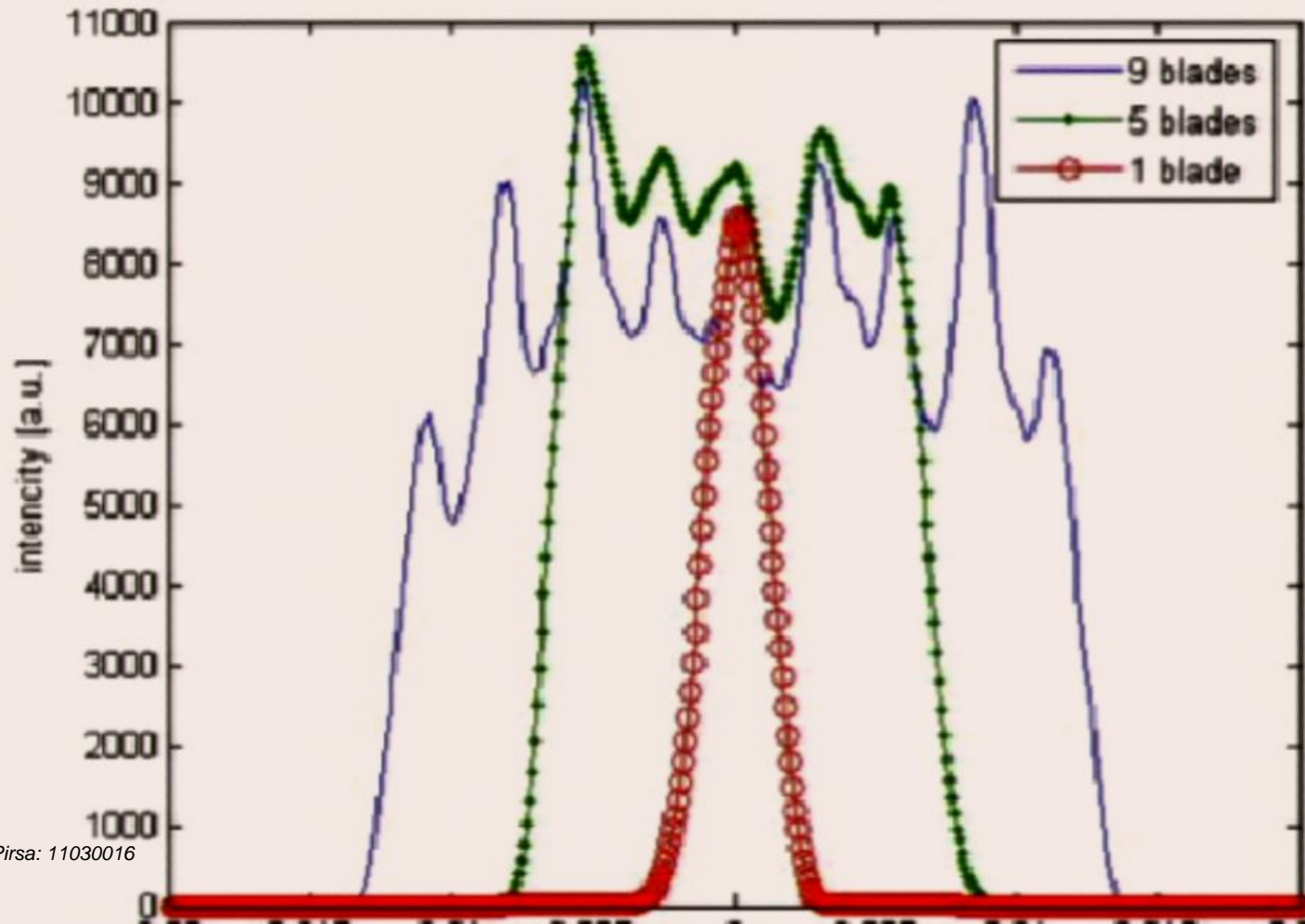


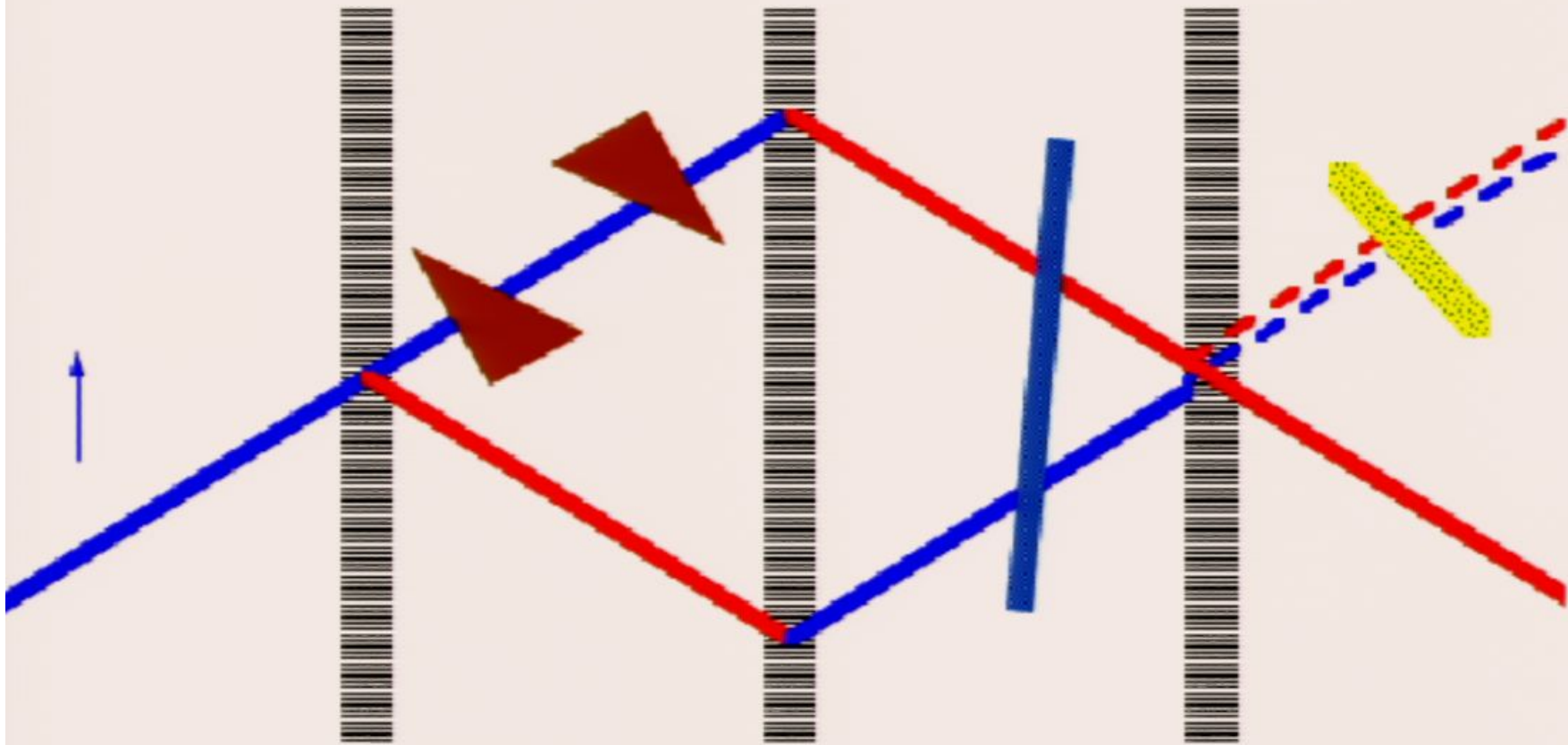
Vertical Neutron Momentum Distribution ($k_0 = 2.3 \times 10^{10} \text{ m}^{-1}$)





Vertical Neutron Momentum Distribution ($k_0 = 2.3 \times 10^{10} \text{ m}^{-1}$)

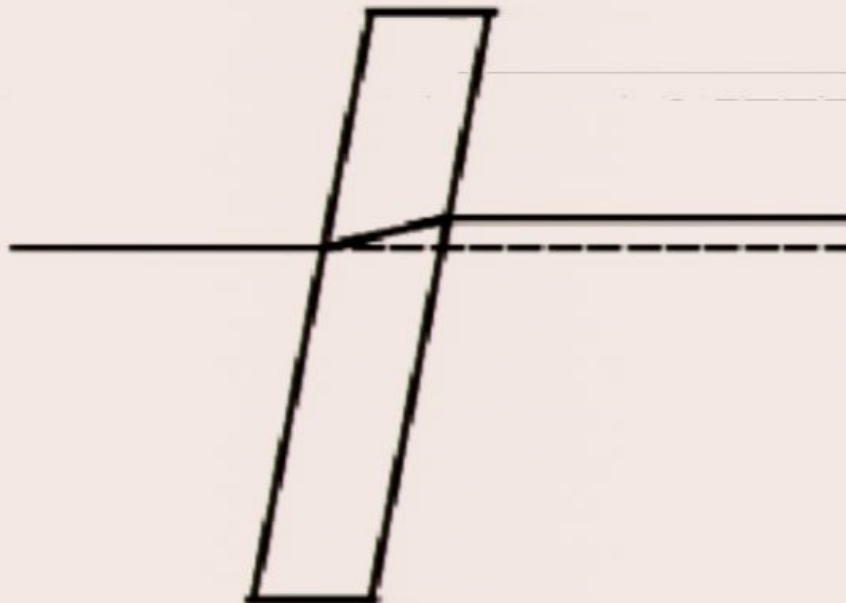




$$= 1 - \text{Cos}[\varphi(z) - \varphi(z + \Delta)]$$

that the outgoing beam is a coherent superposition of the beam and its displaced partner. We could use this to measure the spatial correlations of cross-sections.

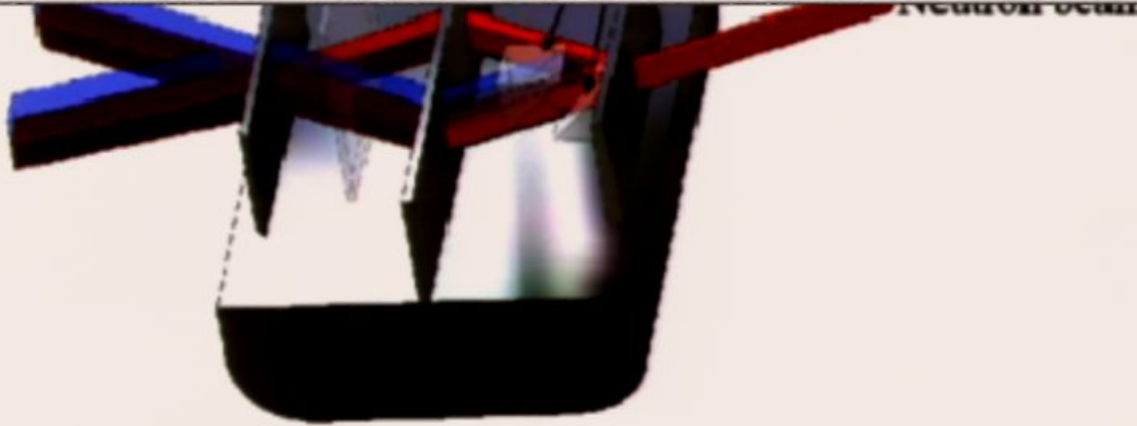
Now that contrast remains we can recombine the blades with a "thick crystal" momentum filter.



For most materials the index of refraction is dominated by the coherent scattering length and

$$n - 1 \approx \frac{2\pi N b_c}{\lambda^2}$$

To the detectors



Contrast measurements directly yields the coherence function

$$|\Psi_{k_z}\rangle = e^{i\phi_1} e^{ik_z z_1} |1\rangle + e^{i\phi_2} e^{ik_z z_2} |2\rangle$$

$$|\Psi_{k_z}\rangle = \int \rho(k_z) (e^{i\phi_1} e^{ik_z z_1} |1\rangle + e^{i\phi_2} e^{ik_z z_2} |2\rangle) dk_z$$

$$P_0 = \int \rho(k_z) \left| \langle \Psi_{k_z} | P_0 | \Psi_{k_z} \rangle \right| dk_z =$$

is a displacement z , then there is a phase shift of $k z$. Where k is the momentum in the z -direction.

```
Udisplacezdown[kz_, z_] := {{1, 0}, {0, Exp[I kz z]}};
```

```
Udisplacezdown[kz, z] // MatrixForm
```

$$\begin{pmatrix} 1 & 0 \\ 0 & e^{i k z z} \end{pmatrix}$$

```
Udisplacezinvdwn[kz_, z_] := {{1, 0}, {0, Exp[-I kz z]}};
```

```
res14[kz_, z_, a_] := TrigReduce[
```

```
ExpToTrig[Simplify[Ublade . Udisplacezdown[kz, z] . Um . Uphase[0, a] . Ublade . in .
Ubladeinv . Uphaseinv[0, a] . Uminv . Udisplacezinvdwn[kz, z] . Ubladeinv]]]
```

```
res14[kz, z, a] // MatrixForm
```

$$\begin{pmatrix} \frac{1}{2} (1 + \cos[a - kz z]) & \frac{1}{2} i \sin[a - kz z] \\ -\frac{1}{2} i \sin[a - kz z] & \frac{1}{2} (1 - \cos[a - kz z]) \end{pmatrix}$$

a Gaussian distribution of momenta along z .

```
nd[x_, sd_] := Exp[-x^2 / (2 sd^2)] / (sd Sqrt[2 π]);
```

```
M140[a_, sd_] :=
```

```
Sum[Sum[nd[kz, sd] Tr[Exp . res14[kz, z, a]] / 2, {kz, -0.1, 0.1, 0.001}],
{z, -.5, .5, .025}] / 40 / 500
```

```
Plot[M140[a, .01], {a, 0, 2 π}, {AxesLabel → {"phase (Radians)", "intensity"},
PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}, PlotRange → {0, 1}}]
```

```
$Aborted
```

```
Plot[M140[a, .02], {a, 0, 2 π}, {AxesLabel → {"phase (Radians)", "intensity"},
PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}, PlotRange → {0, 1}}]
```

$$\begin{pmatrix} \frac{1}{2} (1 + \cos[a - kz z]) & \frac{1}{2} i \sin[a - kz z] \\ -\frac{1}{2} i \sin[a - kz z] & \frac{1}{2} (1 - \cos[a - kz z]) \end{pmatrix}$$

a Gaussian distribution of momenta along z.

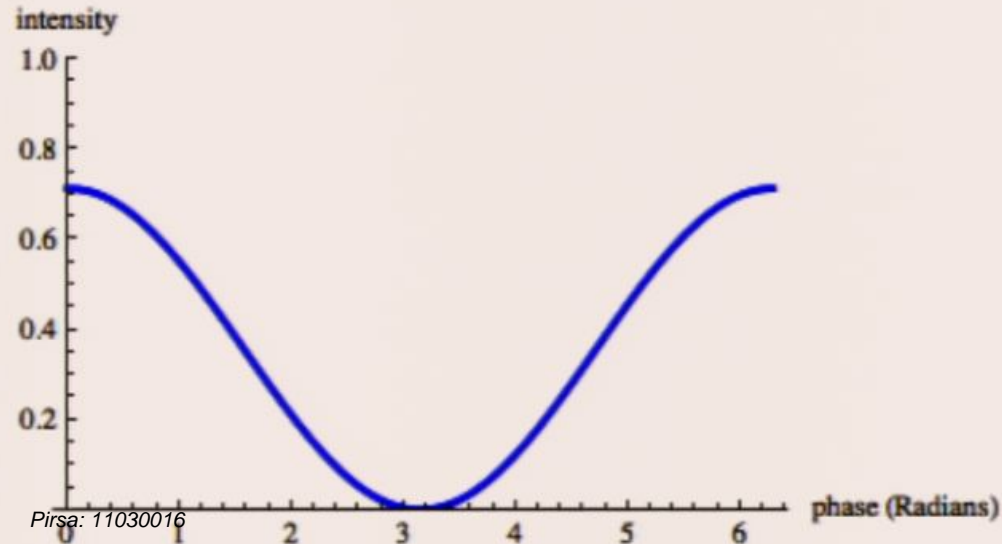
```
nd[x_, sd_] := Exp[-x^2 / (2 sd^2)] / (sd Sqrt[2 π]);
```

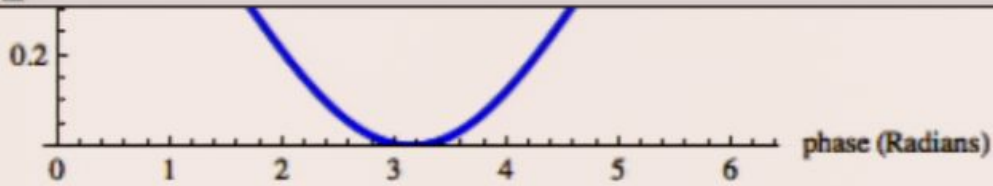
```
M140[a_, sd_] :=  
Sum[Sum[nd[kz, sd] Tr[Exp . res14[kz, z, a]] / 2, {kz, -0.1, 0.1, 0.001}],  
{z, -.5, .5, .025}] / 40 / 500
```

```
Plot[M140[a, .01], {a, 0, 2 π}, {AxesLabel → {"phase (Radians)", "intensity"},  
PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}, PlotRange → {0, 1}}]
```

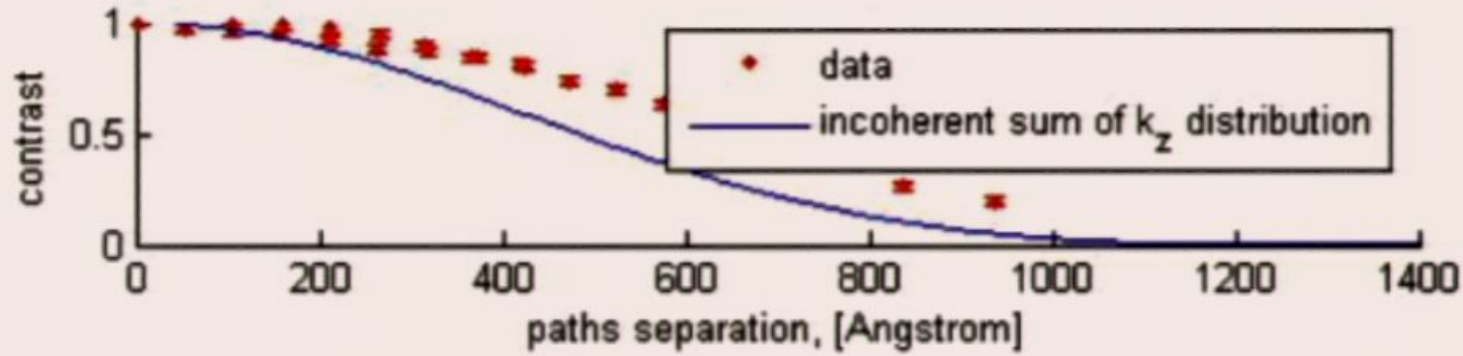
\$Aborted

```
Plot[M140[a, .02], {a, 0, 2 π}, {AxesLabel → {"phase (Radians)", "intensity"},  
PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}, PlotRange → {0, 1}}]
```



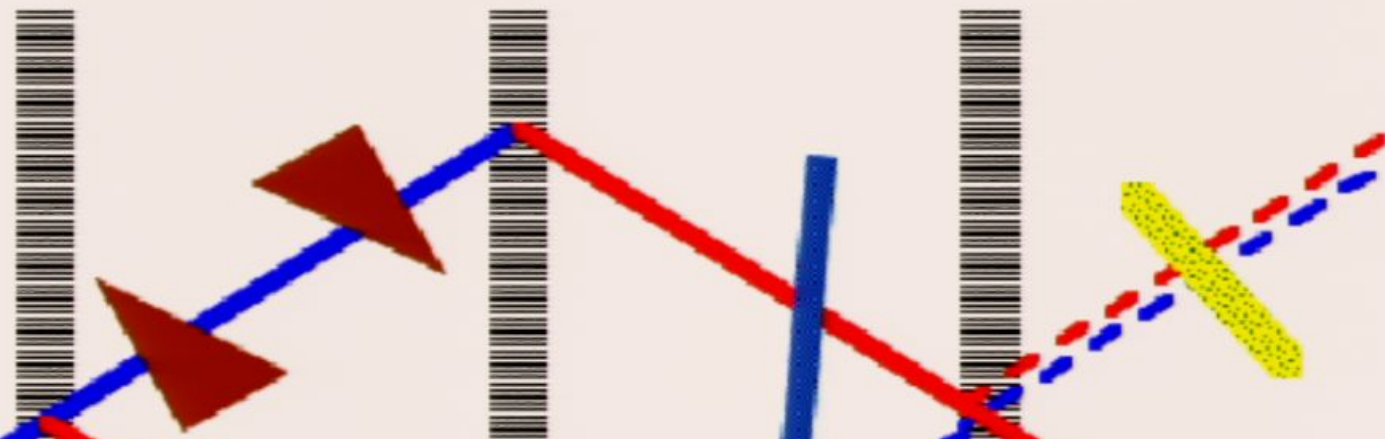


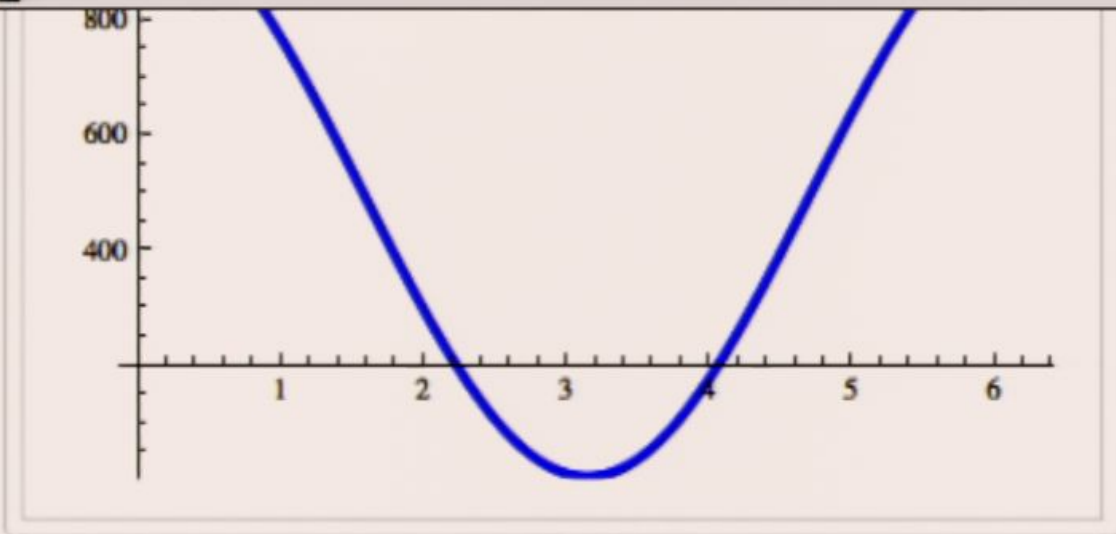
pected the contrast is a function of momentum spread.



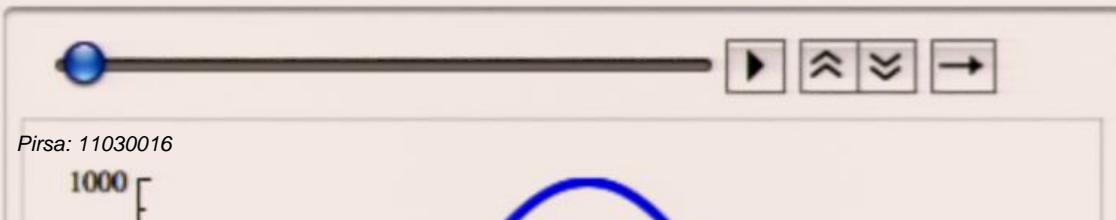
perimentally we see the same thing. Note that the contrast curve is not a Gaussian.

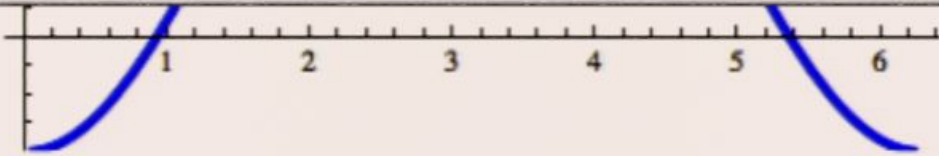
periment 15: outgoing wave



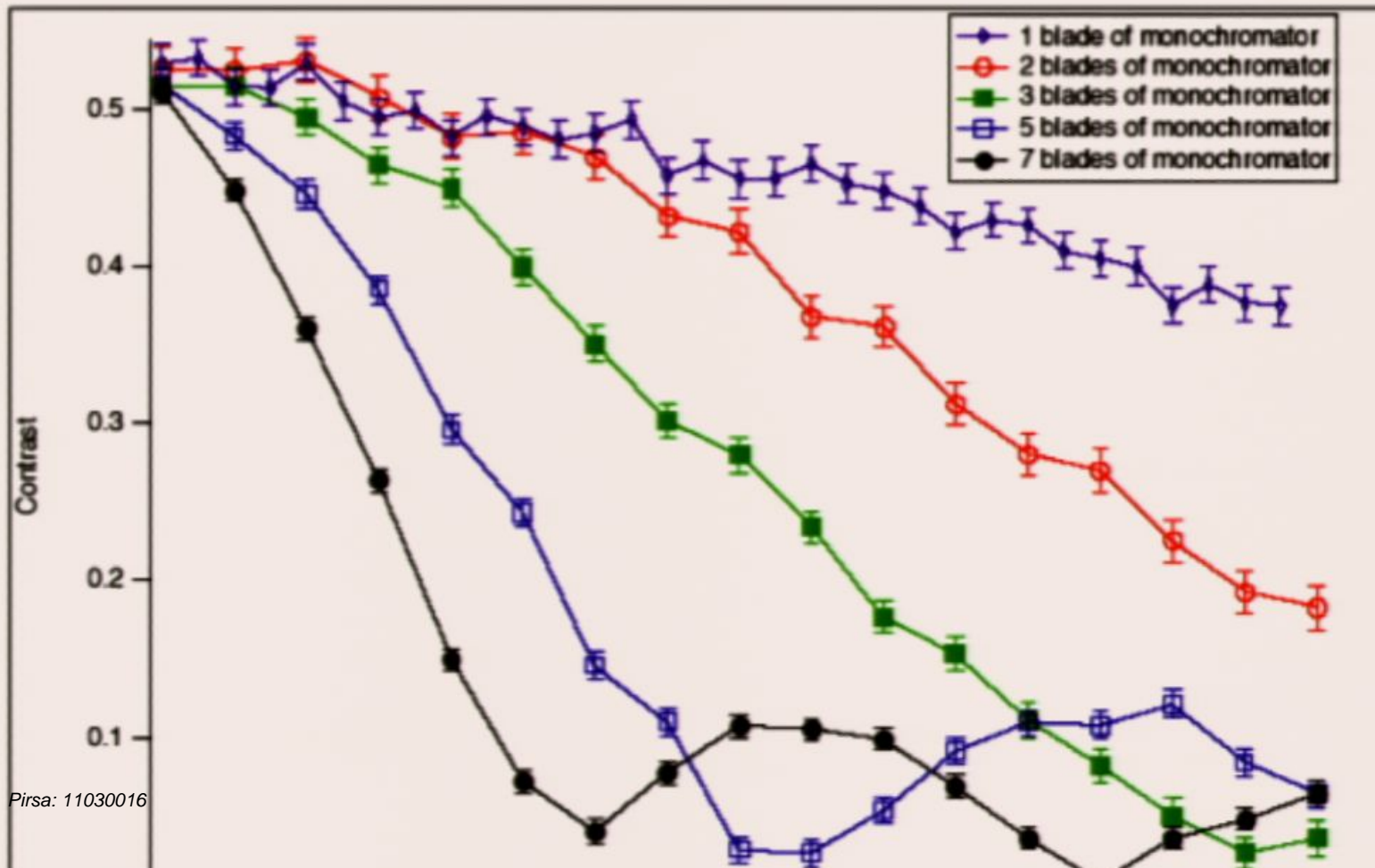


```
ListAnimate[{Plot[Re[M16H[0, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}],
Plot[Re[M16H[2 π, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}],
Plot[Re[M16H[2 π 3, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}],
Plot[Re[M16H[2 π 10, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}],
Plot[Re[M16H[2 π 30, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}],
Plot[Re[M16H[2 π 300, a]], {a, 0, 2 π},
  {PlotRange → {1, 1000}, PlotStyle → {RGBColor[0, 0, 1], Thickness[0.01]}}]]]
```

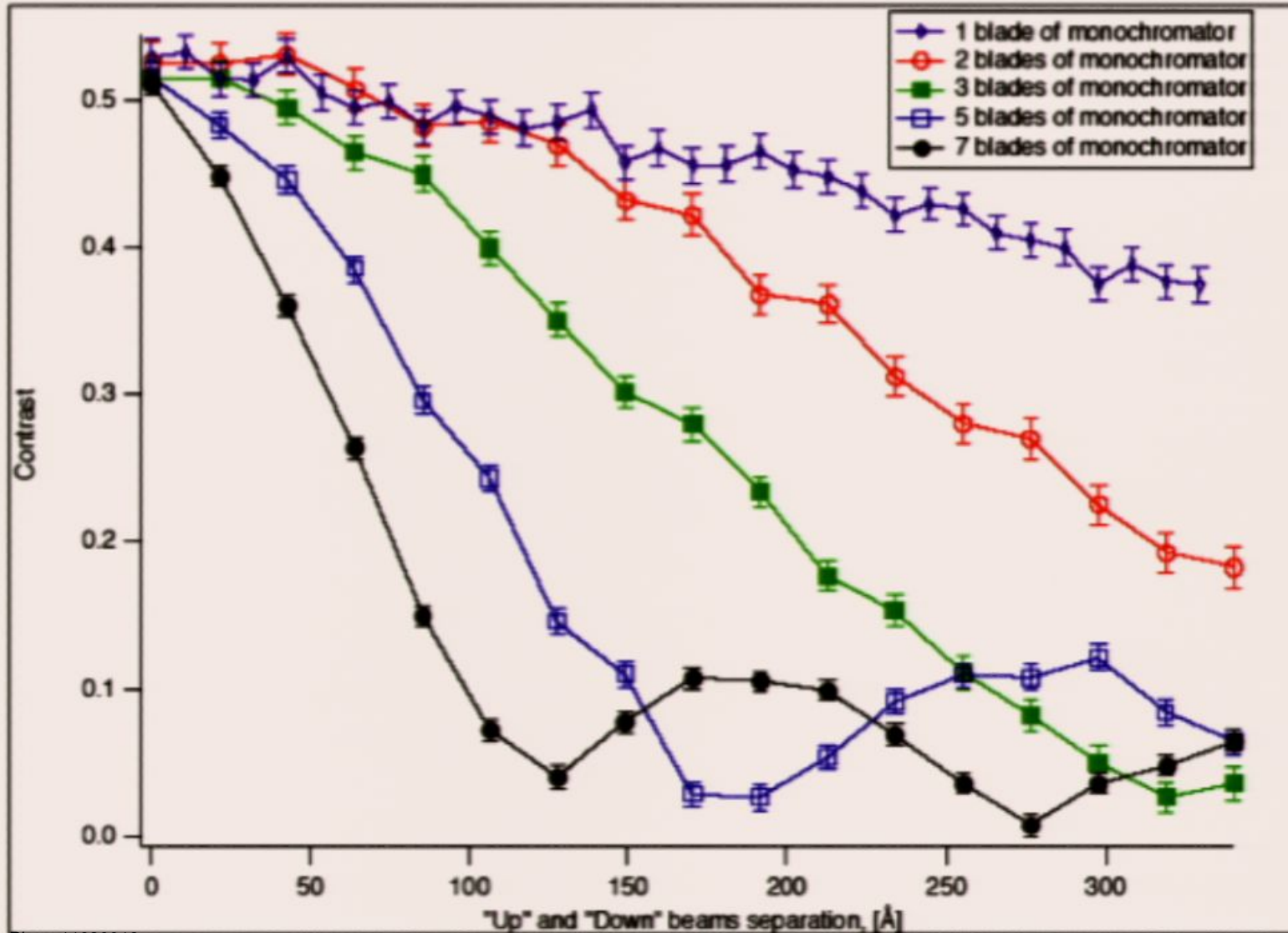




Experiments show the expected behavior.

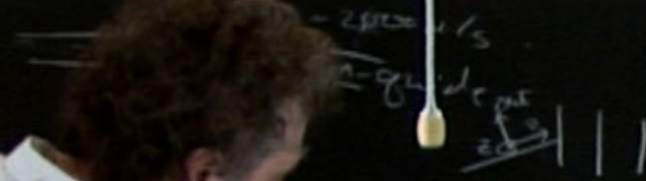


experiments show the expected behavior.



that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is

Resistor



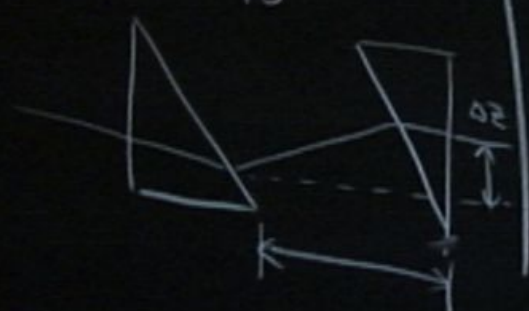
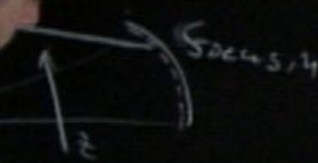
$$\eta = 1 - \frac{\pi^2 B_c N}{2\pi^2} \approx 10^5$$

$$\nabla^2 \psi + K^2 \psi = 0$$

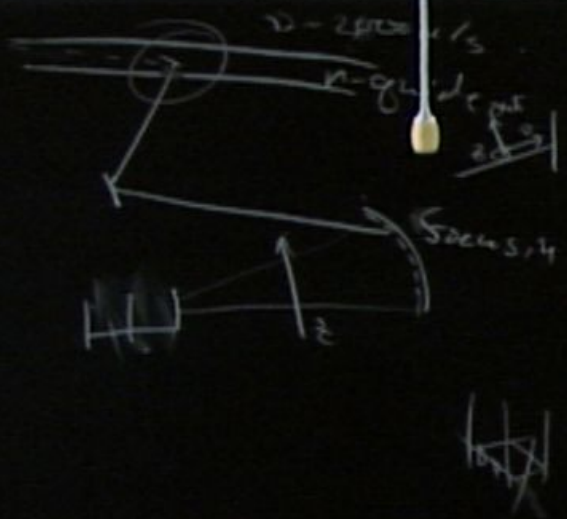
$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction



Resistor



$$\eta = 1 - \frac{\pi^2 B_c N}{2\pi^2} \approx 10^5$$

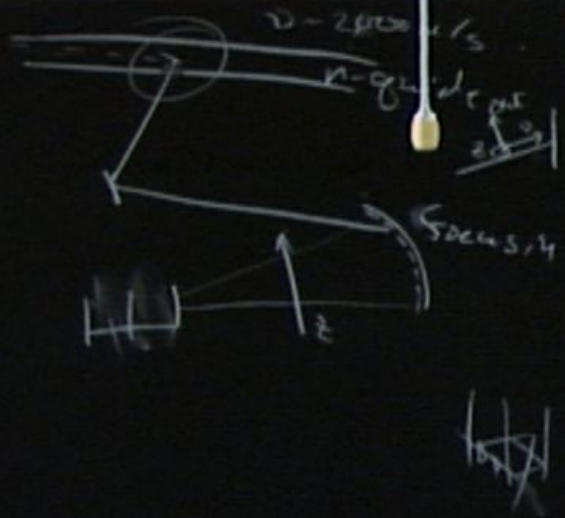
$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} E$$

$$N = \frac{K}{k} \leftarrow \text{index of refraction}$$



Resistor



$$n = 1 - \frac{\pi^2 B_c N}{2 \pi^2} \approx 10^5$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

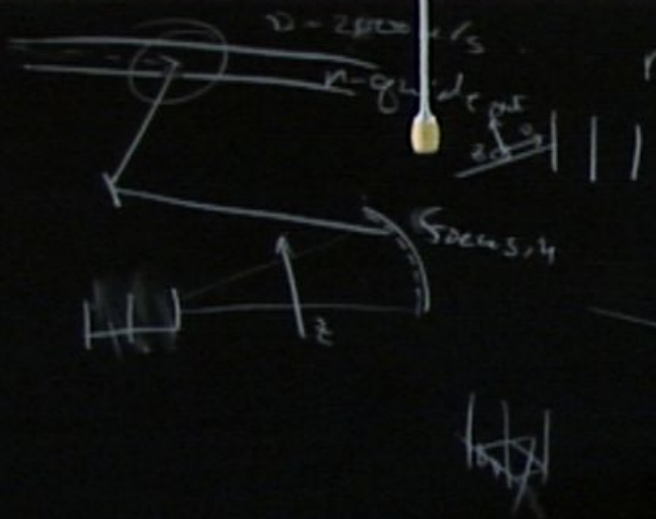
$$n = \frac{K}{k} \leftarrow \text{Free space index of refraction}$$

Handwritten notes on a piece of paper on the desk.

Chalk and a book on the desk.

A wooden board leaning against the chalkboard.

Resistor



$$\eta = 1 - \frac{\pi^2 B_c N}{2\pi^2} \sim 10^5$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

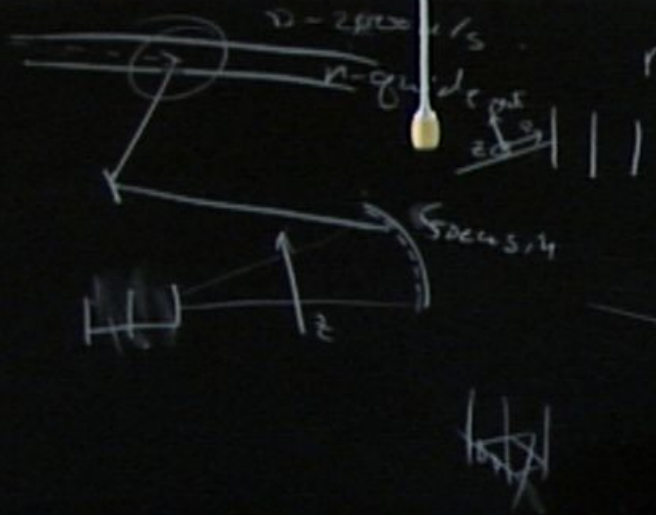
$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

$$\int P(k_z) e^{ik_z 0.2} dx$$



Resistor



$$n = 1 - \frac{\pi^2 B_c N}{2\pi} \approx 10^5$$

$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(z)]^{1/2}$$

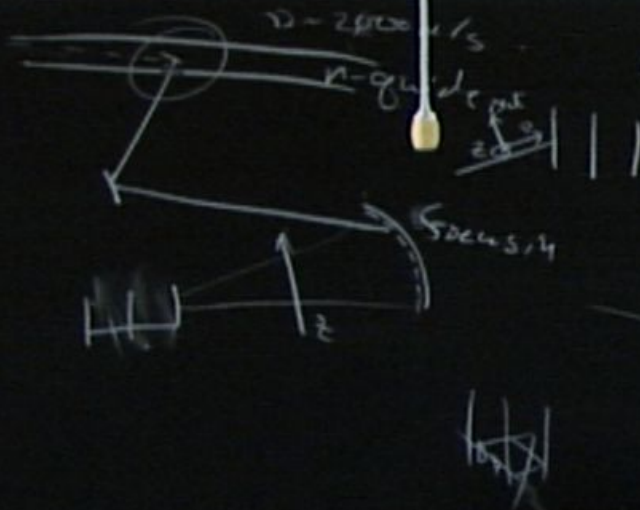
$$n = \frac{K}{k} \leftarrow \text{Free space}$$

index of refraction

$$\int P(k_z) e^{ik_z z} dk_z$$



Resistor



$$\eta = 1 - \frac{\mu^2 B_c N}{2\pi} \approx 10^5$$

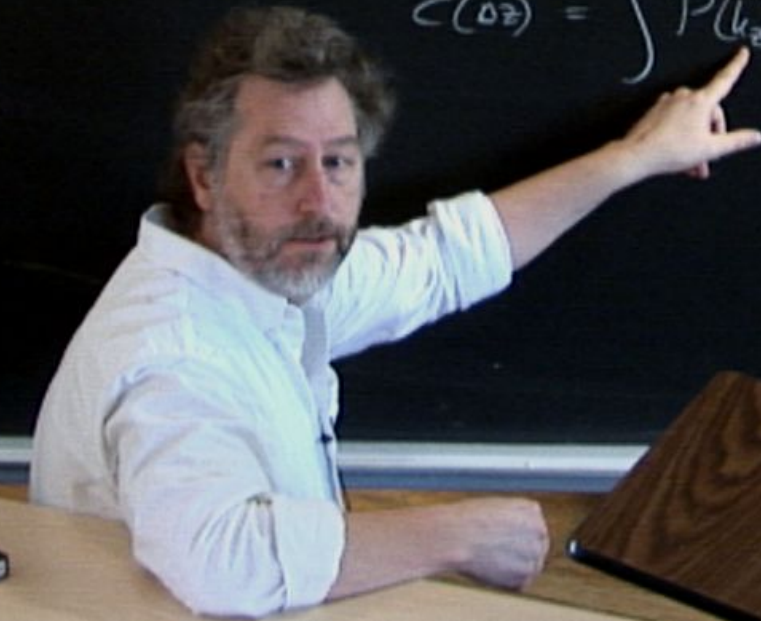
$$\nabla^2 \psi + K^2 \psi = 0$$

$$K = \frac{2m}{\hbar^2} [E - V(x)]^{1/2}$$

$$n = \frac{K}{k} \leftarrow \text{Free space}$$

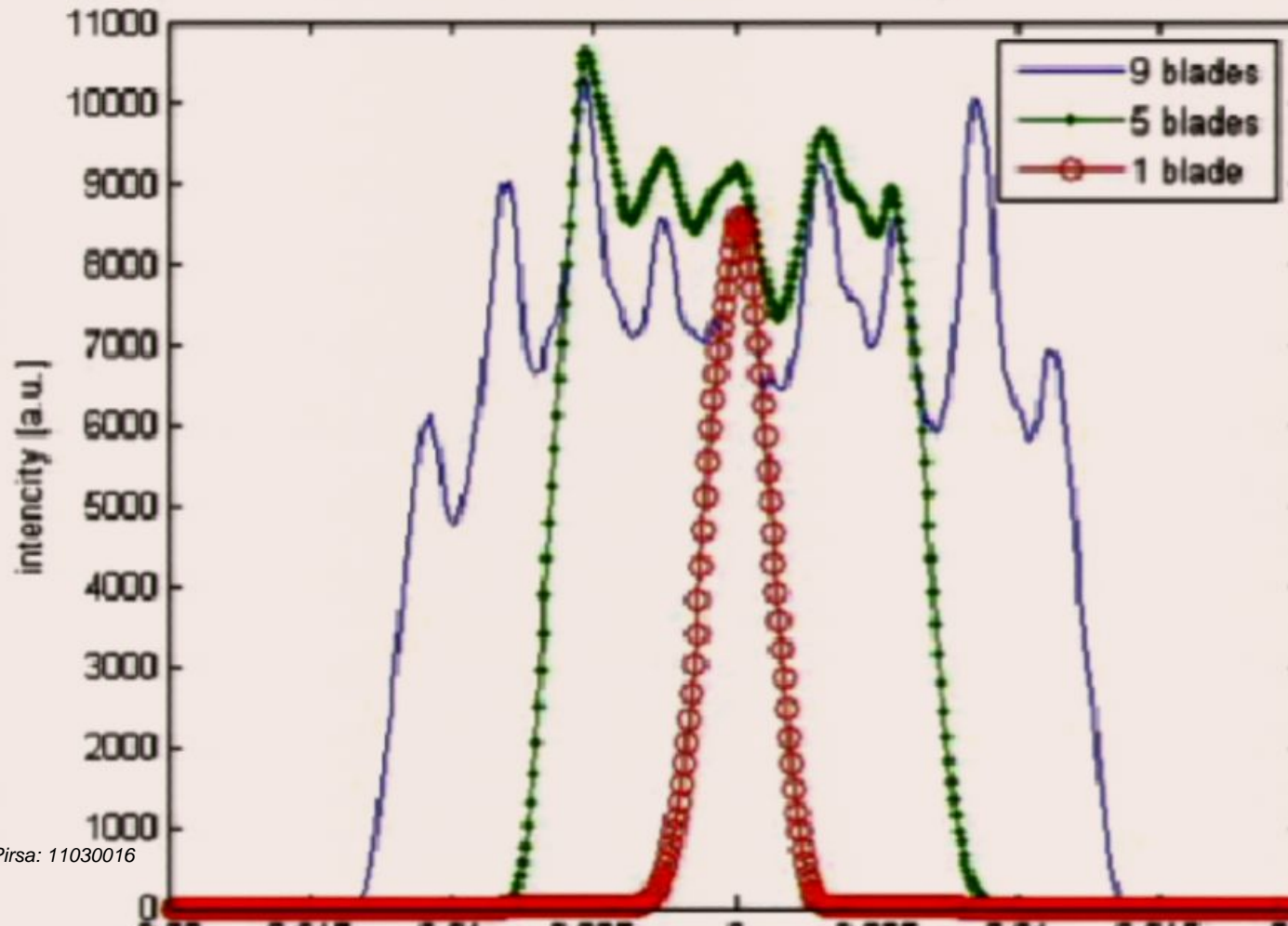
index of refraction

$$C(z) = \int P(k_z) e^{ik_z z} dk_z$$



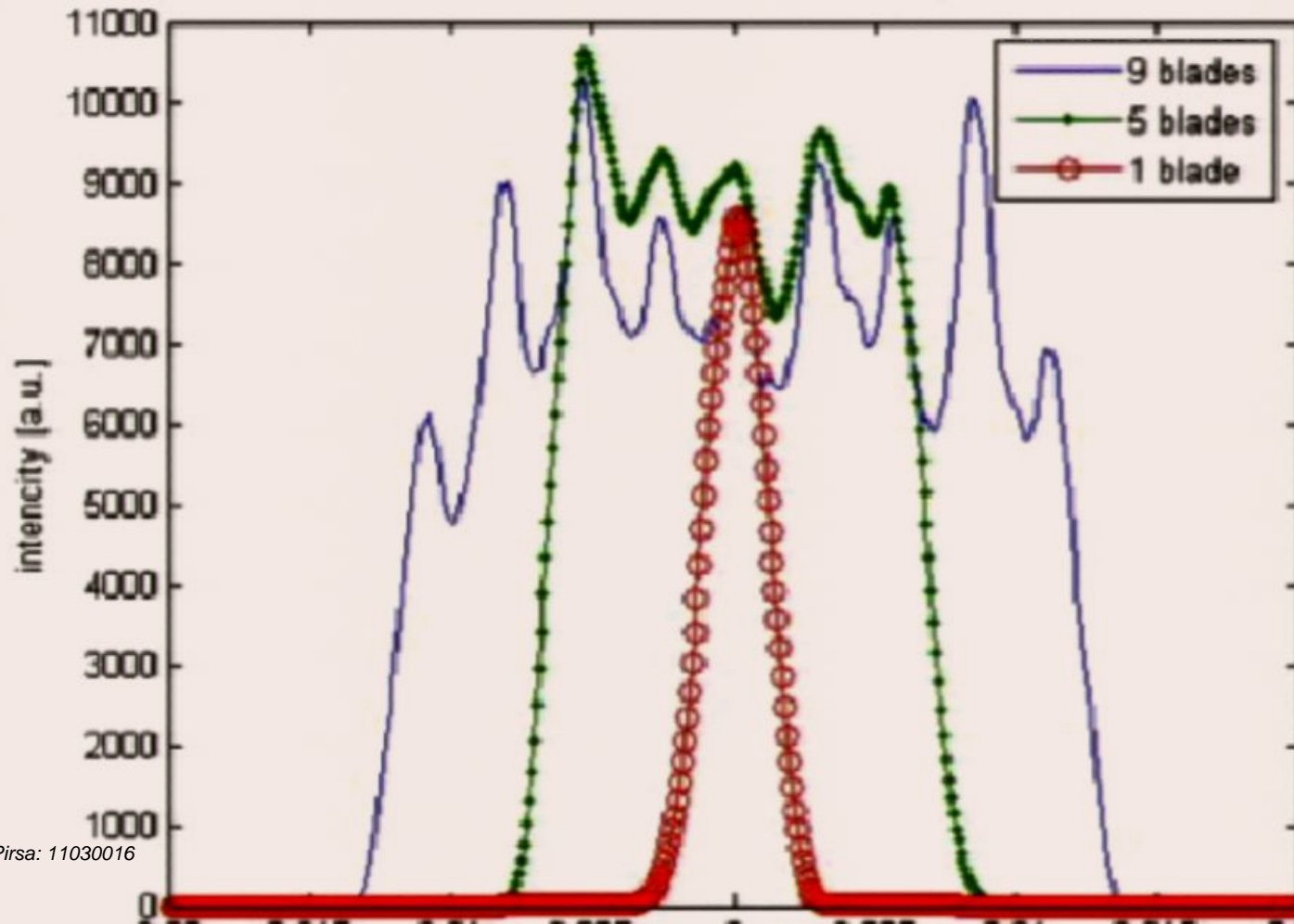


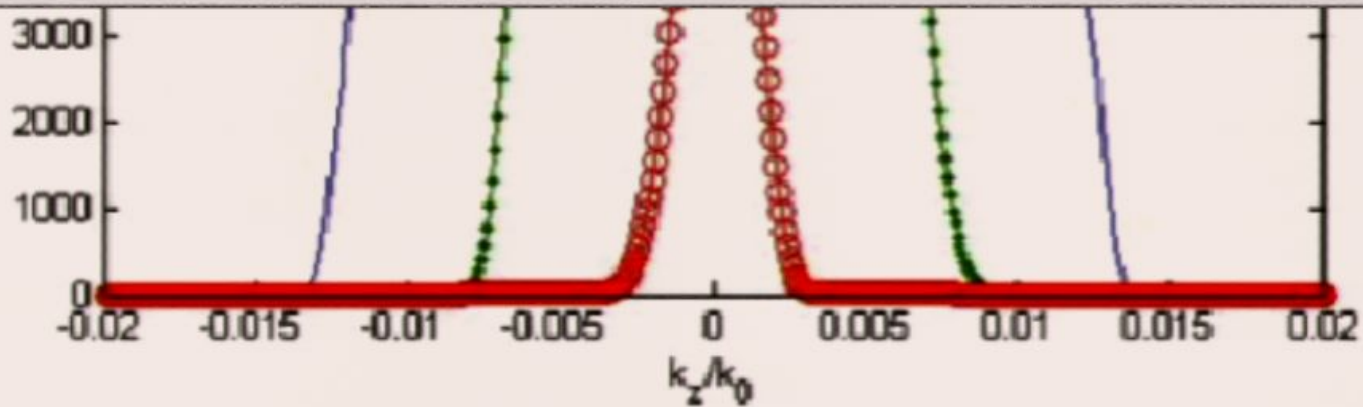
Vertical Neutron Momentum Distribution ($k_0 = 2.3 \times 10^{10} \text{ m}^{-1}$)





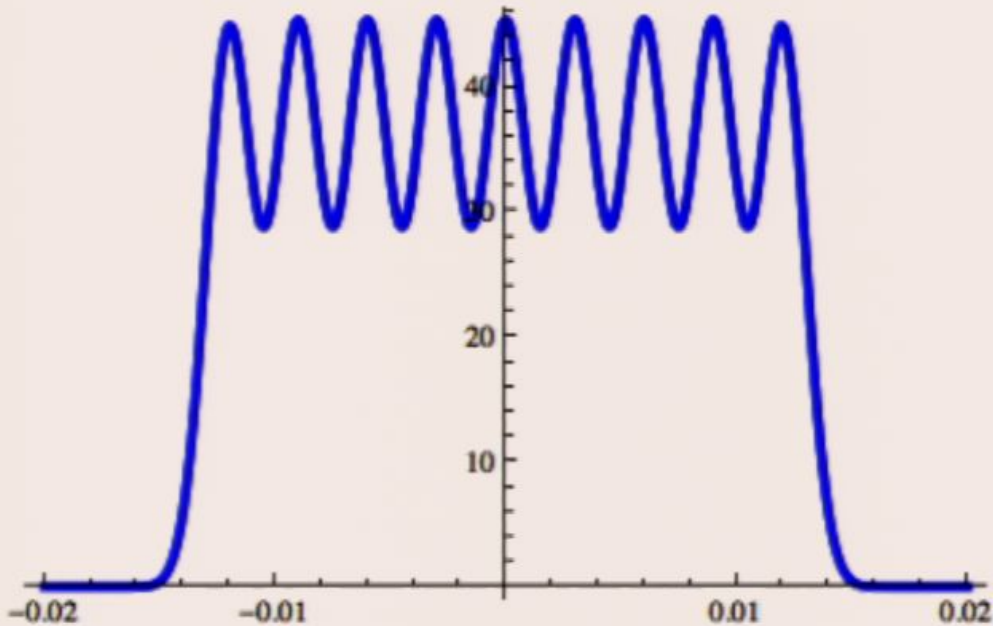
Vertical Neutron Momentum Distribution ($k_0 = 2.3 \times 10^{10} \text{ m}^{-1}$)



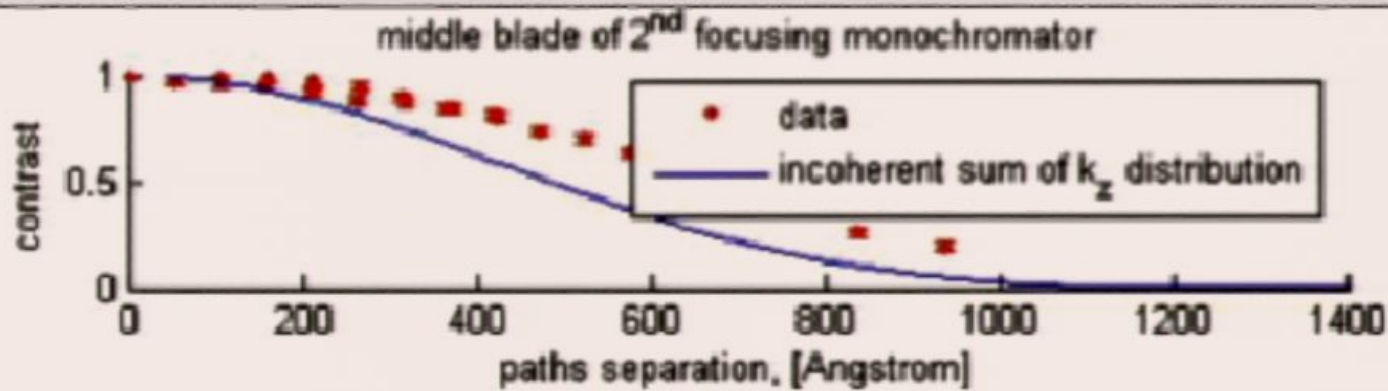


```
Pkz[kz_] := Sum[nd[kz - d, 0.001], {d, 0.012, -0.012, -0.003}]/9
```

```
Plot[Pkz[kz], {kz, -0.02, 0.02}, PlotStyle -> {RGBColor[0, 0, 1], Thickness[0.01]}
```

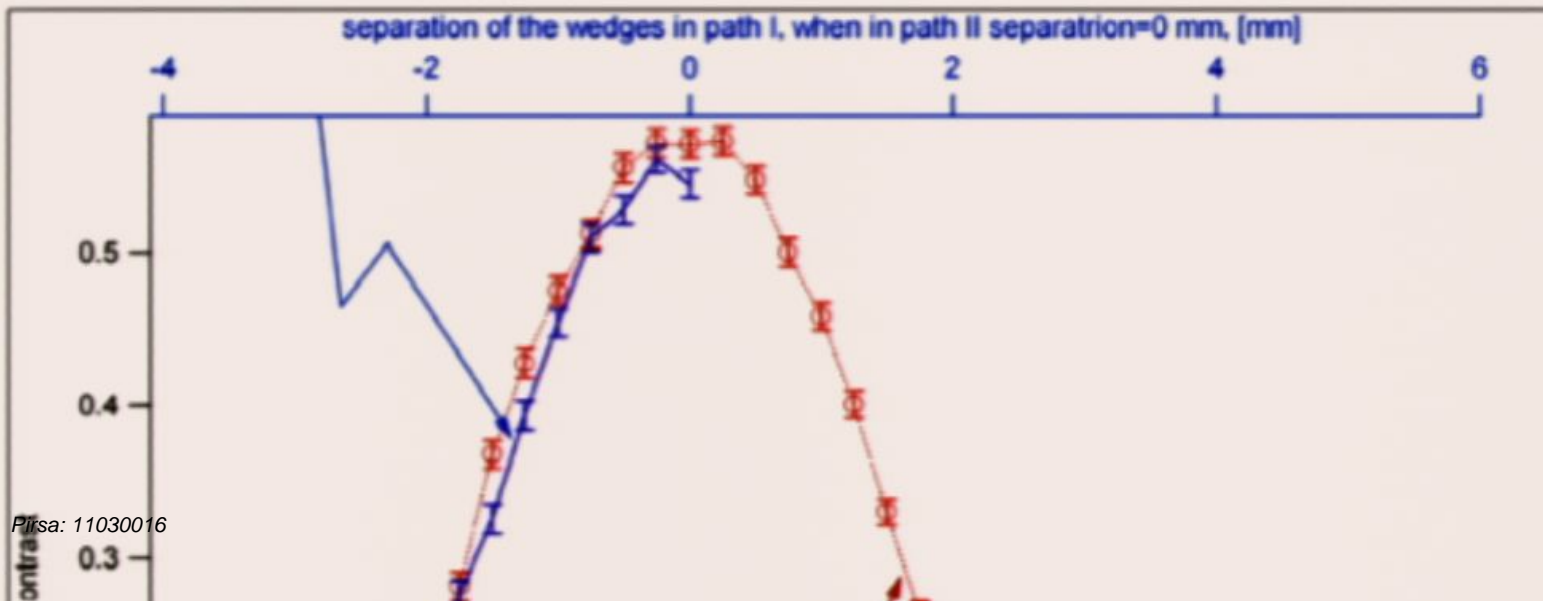


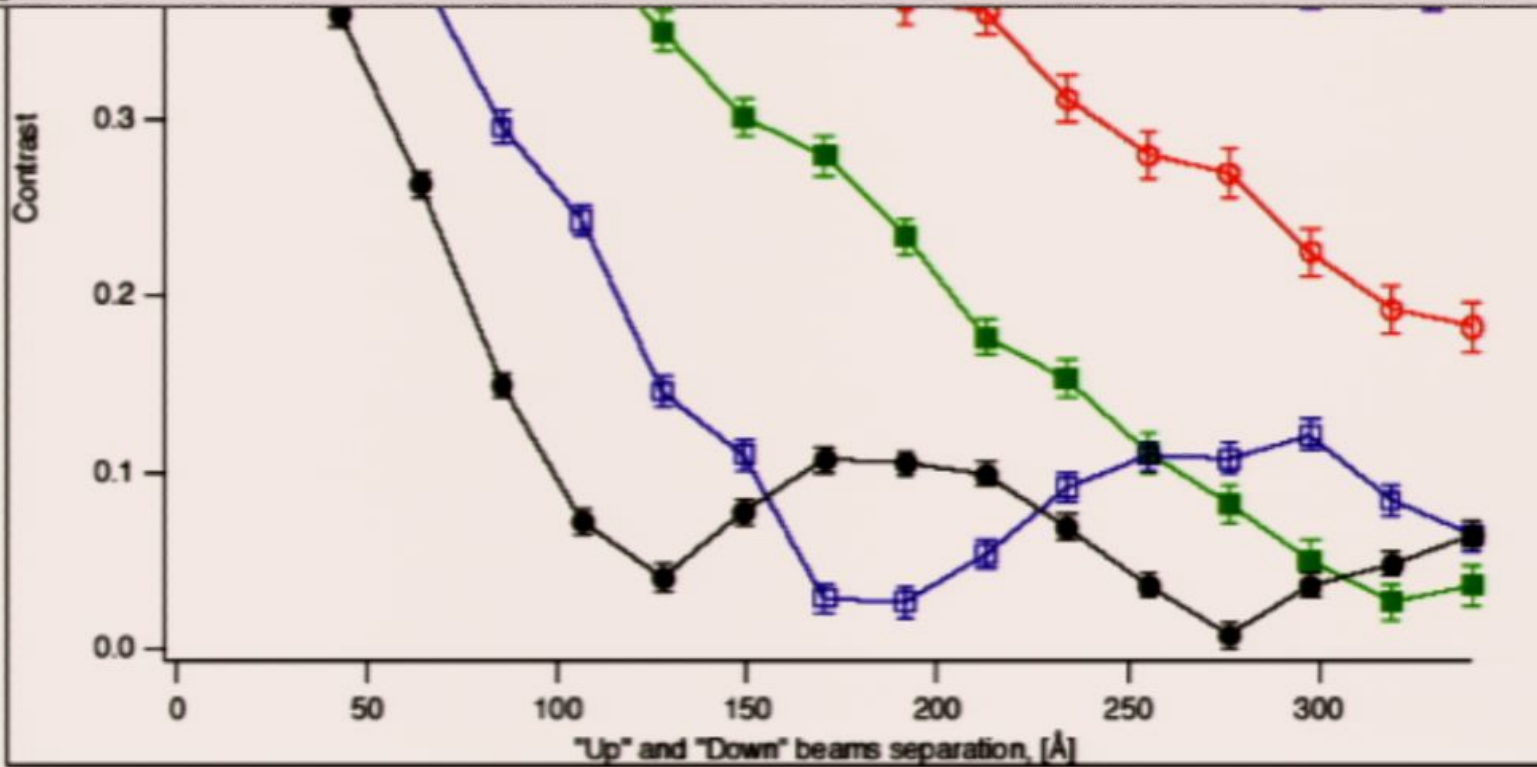
```
M160[z_, a_] := Sum[Pkz[kz] Tr[Exp . res14[kz, z, a]], {kz, -0.02, 0.02, 0.001}]
```

periment 17: Recombining the shifted beams

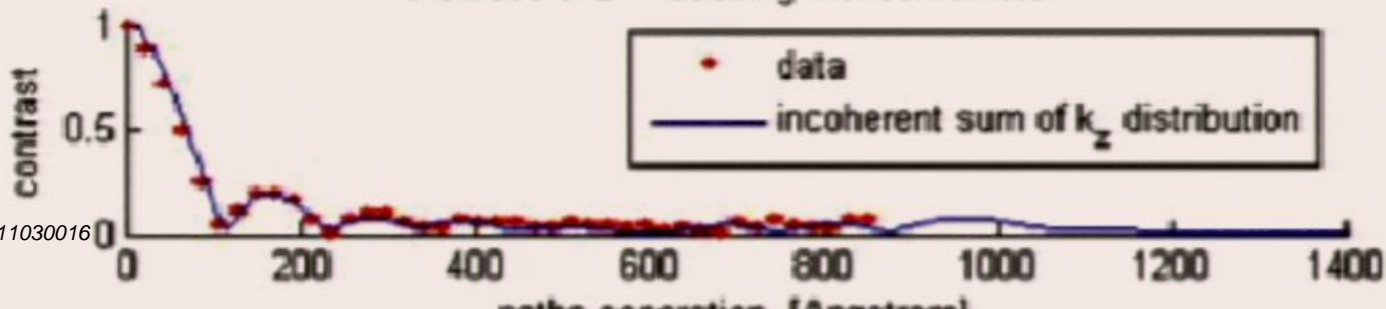
we put a set of prisms in each arm and show that the contrast is fully recovered as we displace the beam in both of the interferometer.

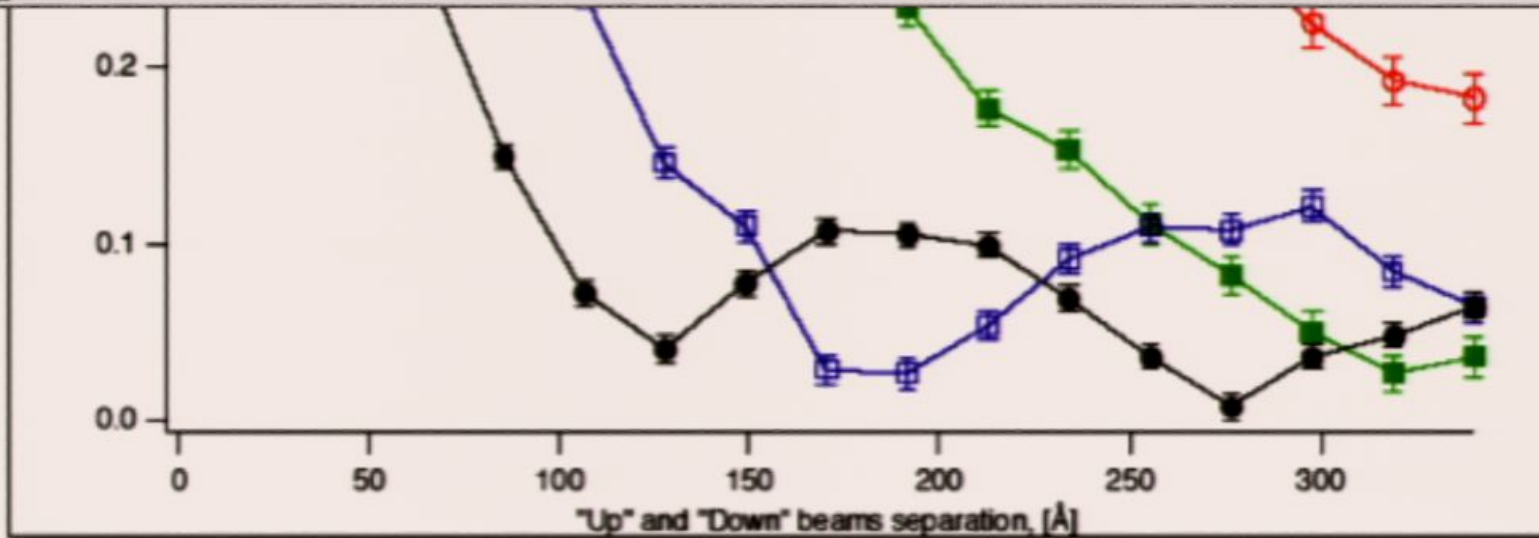




that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is the Fourier transform of the momentum distribution.

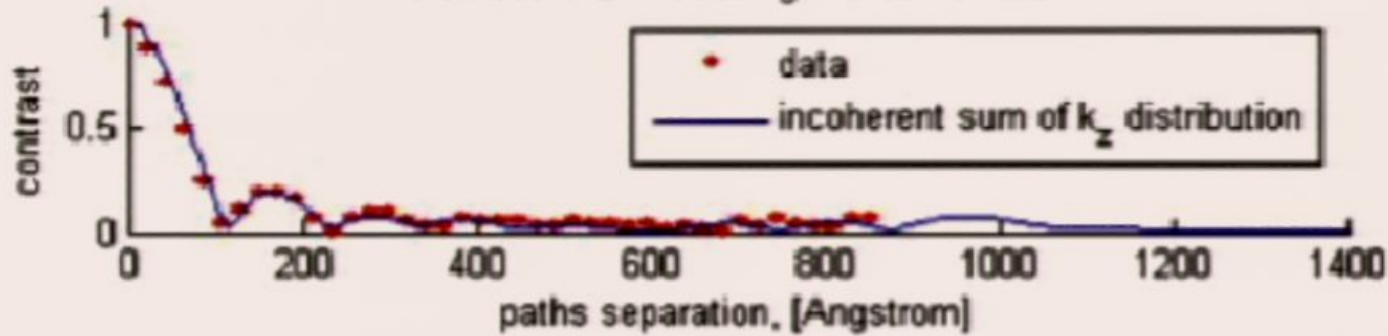
9 blades of 2nd focusing monochromator





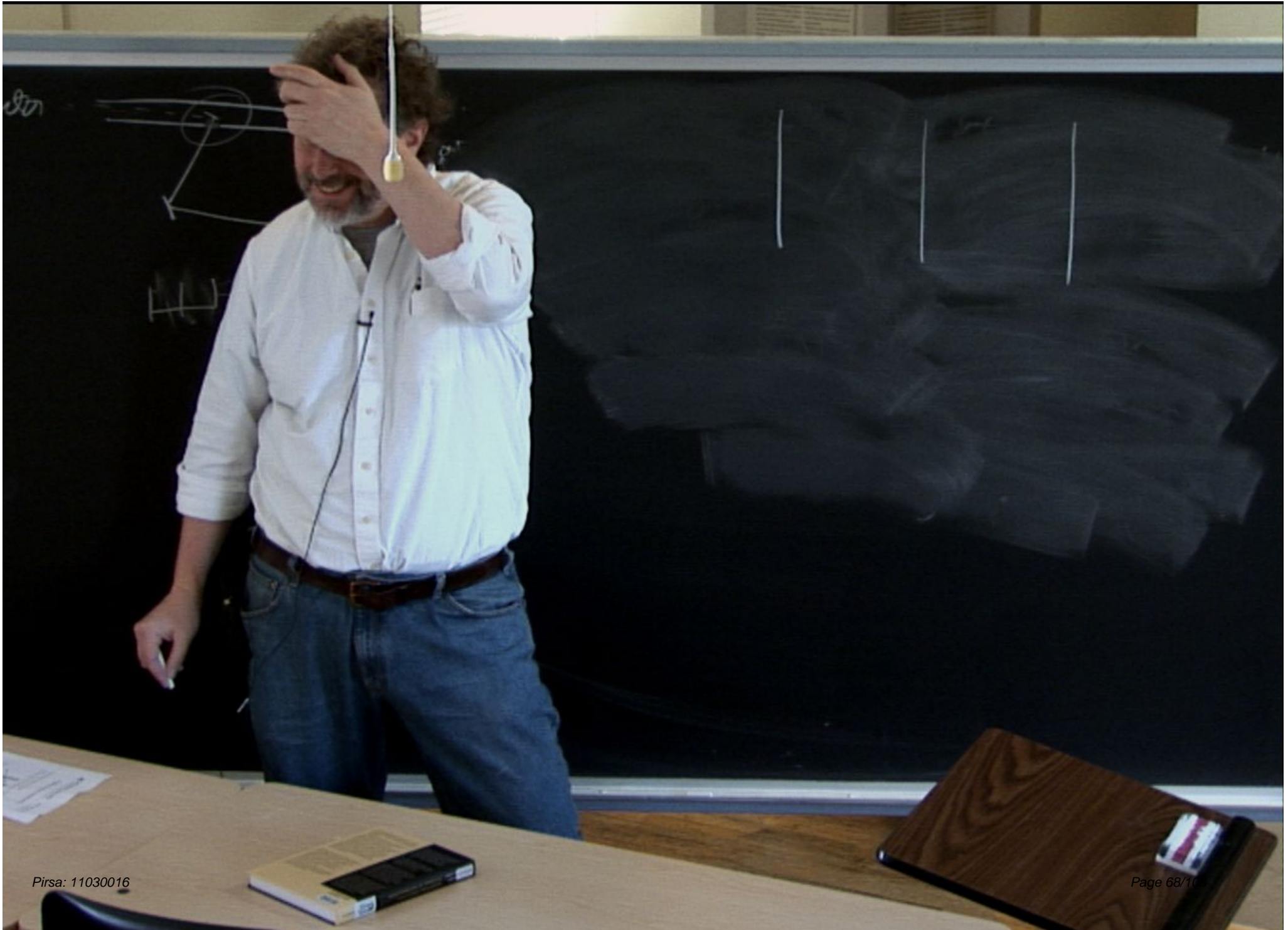
that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is the Fourier transform of the momentum distribution.

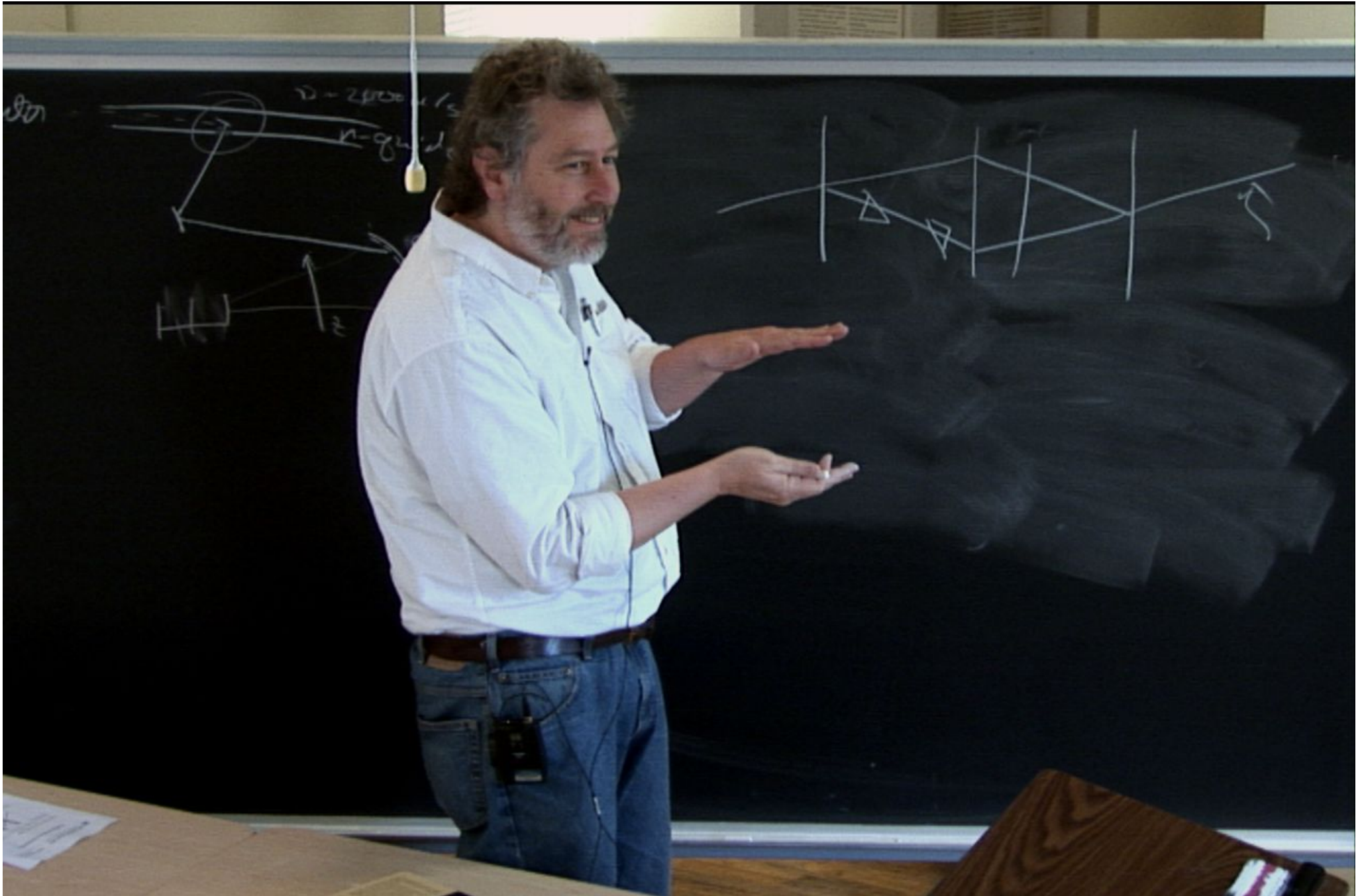
9 blades of 2nd focusing monochromator

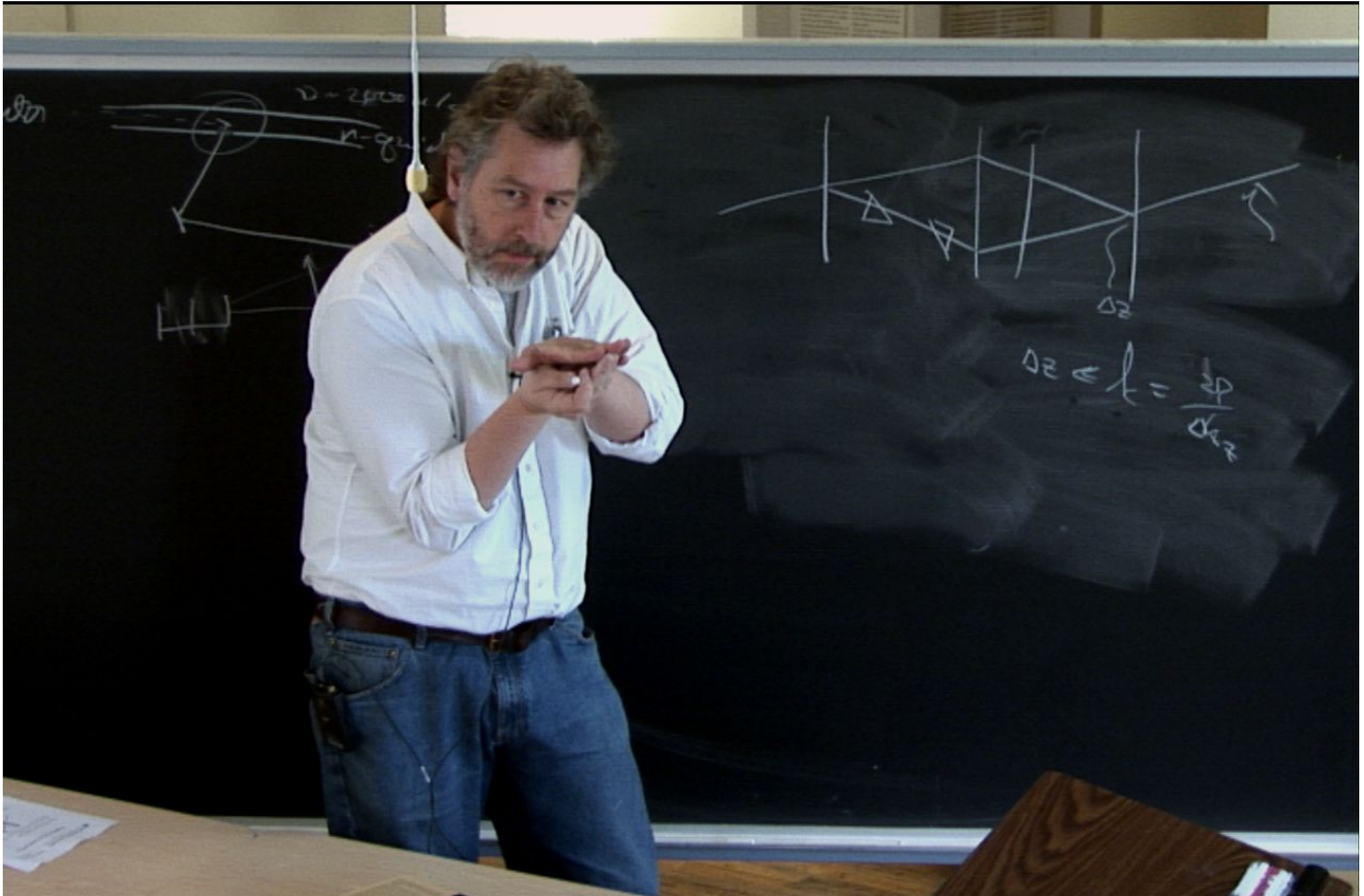


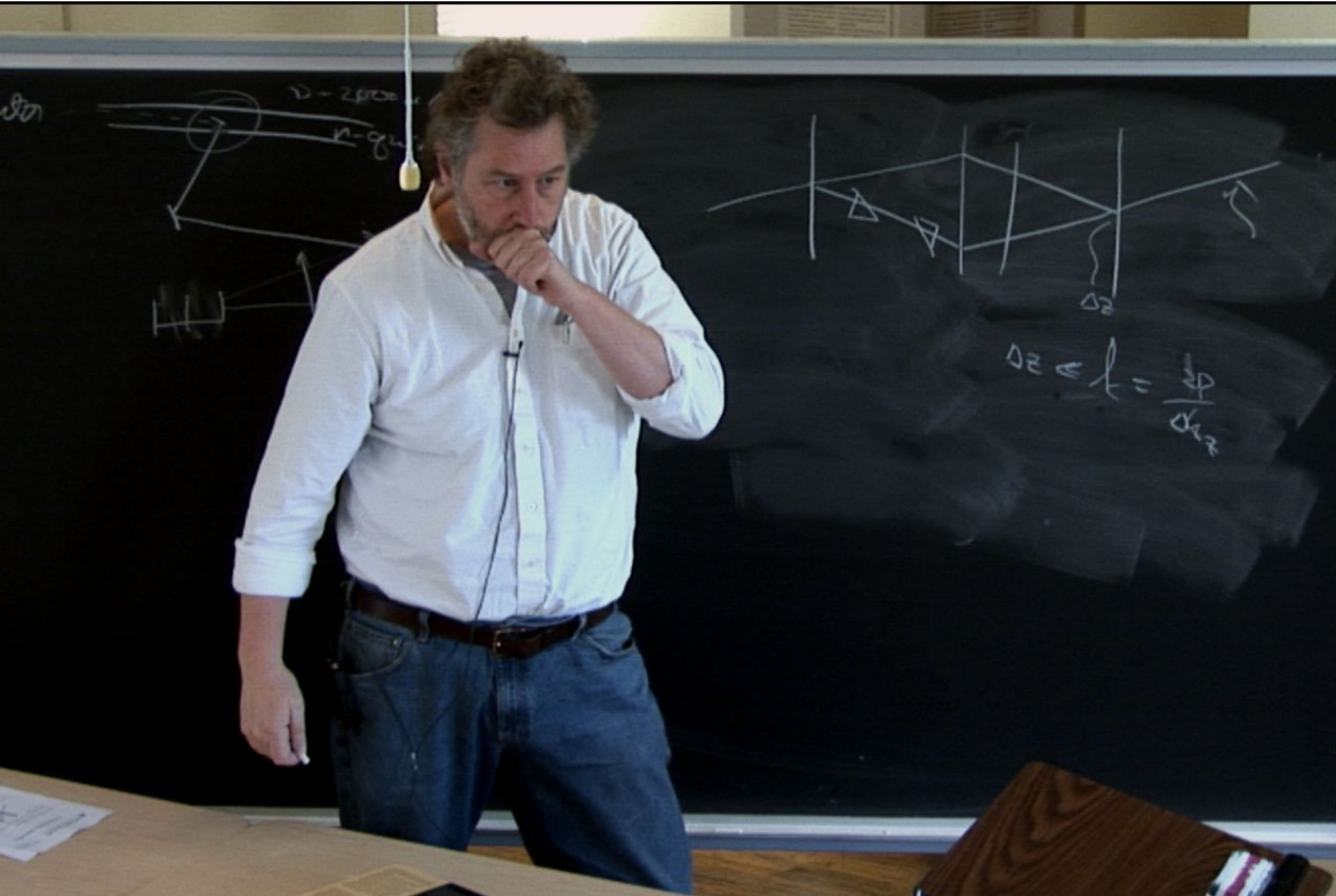
5 blades of 2nd focusing monochromator

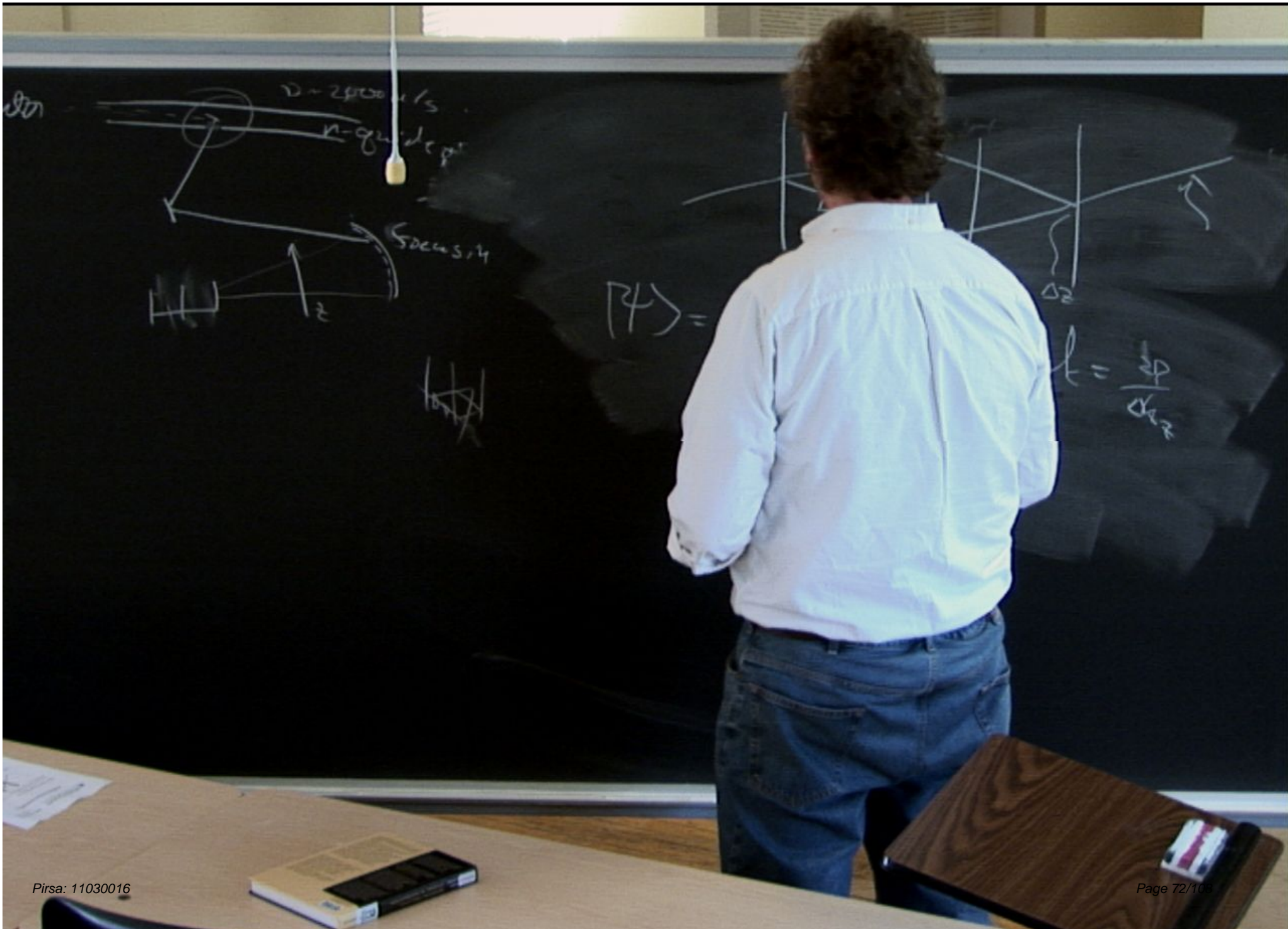






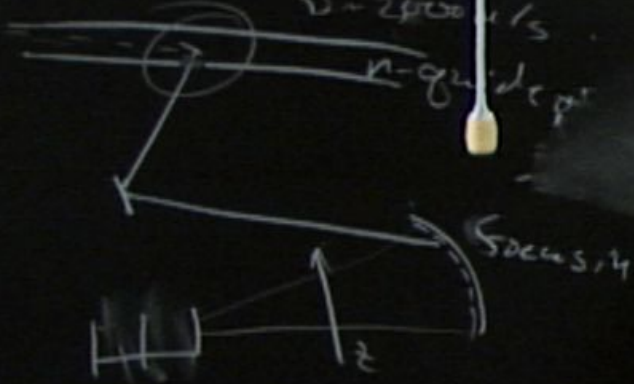






800

$v = 2000 \text{ m/s}$
 $n = 1.5$



focus f

$$f = \frac{zD}{\Delta z}$$

$$f = \frac{zD}{\Delta z}$$

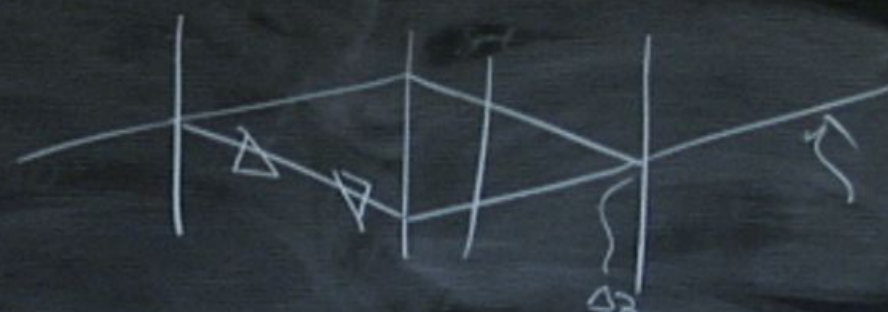
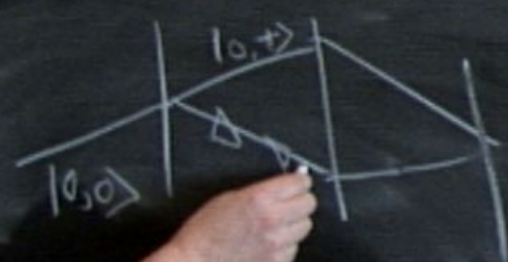
$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of Δz



$$|4\rangle =$$

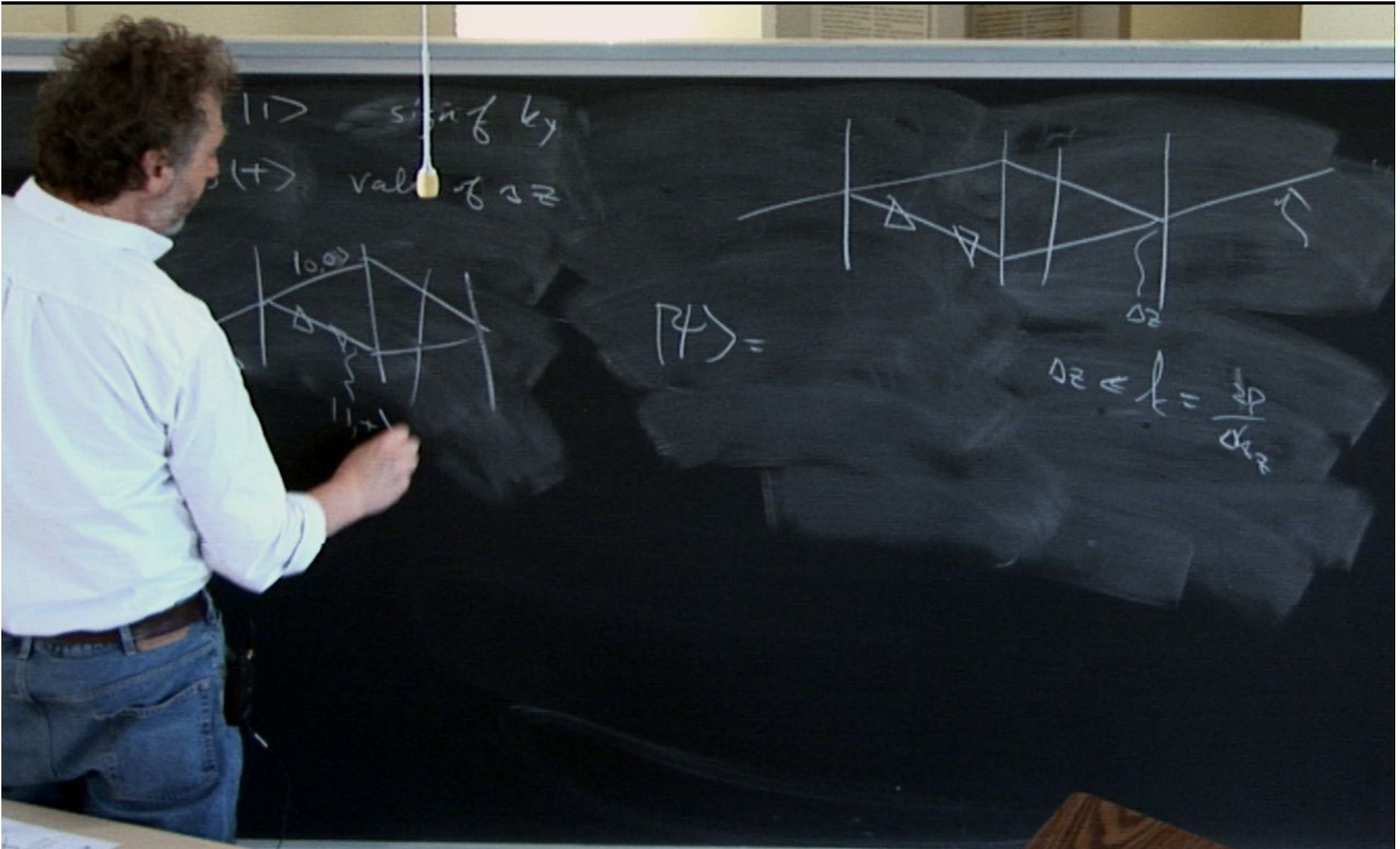
$$\Delta z \approx \frac{\partial \epsilon}{\partial k_z}$$

$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of Δz

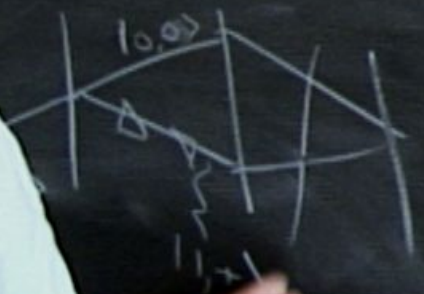


$|1\rangle =$

$$\Delta z \ll \hbar = \frac{\hbar \omega}{\omega_{Lz}}$$



$11 >$ sign of k_y
 $(+)$ value of z



$P4 =$



$$dz \approx \frac{\partial P}{\partial z}$$

$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of σ_z

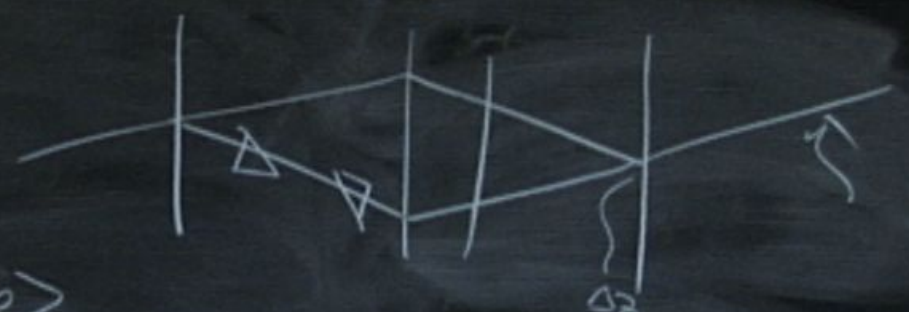
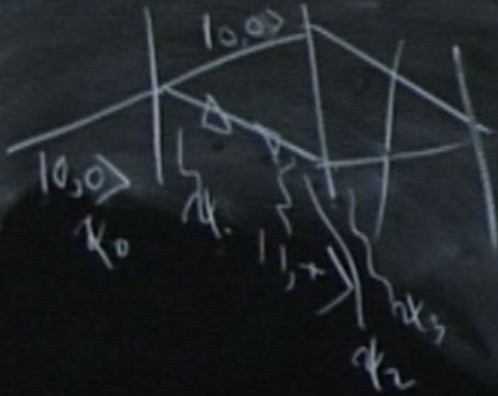


$$\begin{aligned}
 |4_0\rangle &= |0,0\rangle \\
 |4_1\rangle &= \frac{1}{\sqrt{2}}(|0,0\rangle + |1,0\rangle) \\
 |4_2\rangle &= \frac{1}{2\sqrt{2}}(|0,0\rangle + |1,+\rangle) \\
 &= \frac{1}{\sqrt{2}}(|1,0\rangle + |0,+\rangle) \\
 &= \frac{1}{\sqrt{2}}(e^{i\phi}|1,0\rangle + |0,+\rangle)
 \end{aligned}$$



$$\Delta z \ll \ell = \frac{\hbar^2 p}{2m}$$

$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of σ_z



$$|\psi_0\rangle = |0,0\rangle$$

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}(|0,0\rangle + |1,0\rangle)$$

$$|\psi_2\rangle = \frac{1}{2\sqrt{2}}(|0,0\rangle + e^{i\phi} |1,+\rangle)$$

$$|\psi_3\rangle = \frac{1}{\sqrt{2}}(|1,0\rangle + e^{i\phi} |0,+\rangle)$$

$$|\psi_4\rangle = \frac{1}{\sqrt{2}}(e^{-i\phi} |1,0\rangle + e^{i\phi} |0,+\rangle)$$

$$|\psi_5\rangle = \dots$$

$$\Delta z \ll \lambda = \frac{2\pi}{\alpha_{kz}}$$

$$\dots + e^{i\alpha_{kz} z} (|0,+\rangle + \dots)$$

$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of k_z



$$|\psi_0\rangle = |0,0\rangle$$

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}(|0,0\rangle + |1,0\rangle)$$

$$|\psi_2\rangle = \frac{1}{2\sqrt{2}}(|0,0\rangle + e^{i\phi_0} |1,+\rangle)$$

$$|\psi_3\rangle = \frac{1}{\sqrt{2}}(|1,0\rangle + e^{i\phi_0} |0,+\rangle)$$

$$|\psi_4\rangle = \frac{1}{2\sqrt{2}}(e^{i\phi} |1,0\rangle + e^{i\phi_0} |0,+\rangle)$$

$$|\psi_5\rangle = \frac{1}{2}(e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i\phi_0} (|0,+\rangle + |1,+\rangle))$$

$$\Delta z \ll \lambda = \frac{2\pi}{\alpha k_z}$$

$|0\rangle, |1\rangle$

$|0\rangle, |+\rangle$

van

$||K_H^0$

$|0,0\rangle$
 K_0

$$|\psi_0\rangle = |0,0\rangle$$

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}(|0,0\rangle + |1,0\rangle)$$

$$|\psi_2\rangle = \frac{1}{2\sqrt{2}}(|0,0\rangle + e^{i\phi_0} |1,+\rangle)$$

$$|\psi_3\rangle = \frac{1}{\sqrt{2}}(|1,0\rangle + e^{i\phi_0} |0,+\rangle)$$

$$|\psi_4\rangle = \frac{1}{\sqrt{2}}(e^{-i\phi} |1,0\rangle + e^{i\phi_0} |0,+\rangle)$$

$$|\psi_5\rangle = \frac{1}{2}(e^{-i\phi} (|1,0\rangle - |0,0\rangle) + e^{i\phi_0} (|0,+\rangle + |1,+\rangle))$$

$|0\rangle, |1\rangle$ sign \uparrow $|K_H^0\rangle$ $|2_0\rangle = (-e^{i\phi} |0,0\rangle$
 $|0\rangle, |+\rangle$ val \uparrow

$|K_0\rangle = |0,0\rangle$
 $|2_0\rangle = \frac{1}{\sqrt{2}}(|0,0\rangle + |1,0\rangle)$
 $= \frac{1}{2\sqrt{2}}(|0,0\rangle + e^{i\phi} |1,+\rangle)$
 $\frac{1}{\sqrt{2}}(|1,0\rangle + e^{i\phi} |0,+\rangle)$
 $|2_+$ $(e^{i\phi} |1,0\rangle + |0,+\rangle)$
 $|2_c\rangle = \frac{1}{\sqrt{2}}(|1,0\rangle - |0,0\rangle) + e^{i\phi}(|0,+\rangle + |1,+\rangle)$

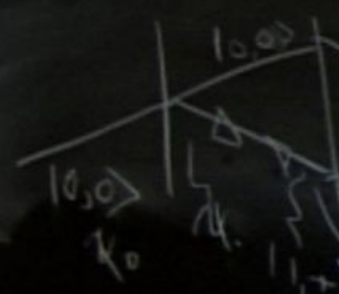
$|0\rangle, |1\rangle$
 $|0\rangle, |+\rangle$

$$|1\rangle_K = \frac{1}{\sqrt{2}} \left(-e^{i\phi} |0,0\rangle + e^{i(k_2 + \phi)z} |0,+\rangle \right)$$

$$\begin{aligned}
 |0\rangle_K &= |0,0\rangle \\
 &= \frac{1}{\sqrt{2}} (|0,0\rangle + |1,0\rangle) \\
 &= \frac{1}{\sqrt{2}} (|0,0\rangle + e^{i\phi} |1,+\rangle) \\
 &= \frac{1}{\sqrt{2}} (|1,0\rangle + e^{i\phi} |0,+\rangle) \\
 &= \frac{1}{\sqrt{2}} (e^{i\phi} |1,0\rangle + e^{i\phi} |0,+\rangle) \\
 |1\rangle_K &= \frac{1}{\sqrt{2}} (e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i(k_2 + \phi)z} (|0,+\rangle + |1,+\rangle))
 \end{aligned}$$

$|0\rangle, |1\rangle$ sign
 $|0\rangle, |+\rangle$ val
 K_H

$$|2\rangle = \frac{1}{\sqrt{2}} \left(-e^{-i\phi} |0,0\rangle + e^{i\phi} |0,+\rangle \right)$$



$$(1,0\rangle + |1,0\rangle)$$

$$\frac{1}{\sqrt{2}} \left(|0,0\rangle + e^{i\phi} |1,+\rangle \right)$$

$$\frac{1}{\sqrt{2}} \left(|1,0\rangle + e^{i\phi} |0,+\rangle \right)$$

$$= \left(e^{-i\phi} |1,0\rangle + e^{i\phi} |0,+\rangle \right)$$

$$\frac{1}{\sqrt{2}} \left(e^{-i\phi} (|1,0\rangle - |0,0\rangle) + e^{i\phi} (|0,+\rangle + |1,+\rangle) \right)$$

$|0\rangle, |1\rangle$ sign of k
 $|0\rangle, |+\rangle$ value of k

$$|2\rangle = \frac{1}{\sqrt{2}} \left(-e^{i\phi} |0,0\rangle + e^{i k_{z0} z} |0,+\rangle \right)$$

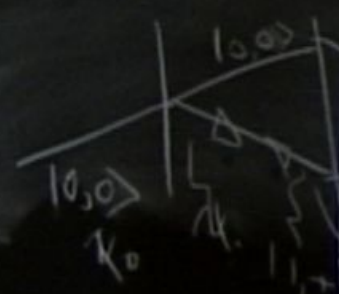
$$= \frac{1}{2} \left(e^{-i k_{z0} z} |+\rangle - e^{i\phi} |0\rangle \right) \otimes |0\rangle$$

$$= \frac{1}{\sqrt{2}} \left(|0,0\rangle + e^{i\phi} |1,+\rangle \right)$$

$$= \frac{1}{\sqrt{2}} \left(|1,0\rangle + e^{i\phi} |0,+\rangle \right)$$

$$= \frac{1}{2} \left(e^{i\phi} |1,0\rangle + e^{i\phi} |0,+\rangle \right)$$

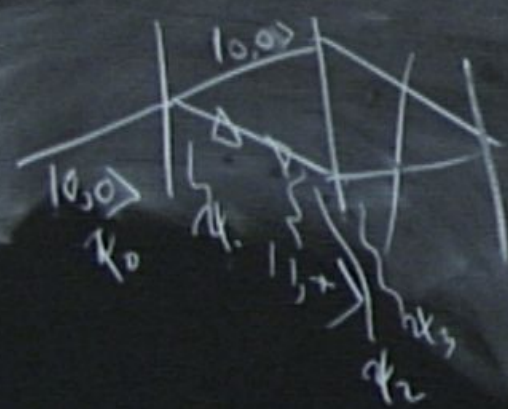
$$= \frac{1}{2} \left(e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i k_{z0} z} (|0,+\rangle + |1,+\rangle) \right)$$



$|0\rangle, |1\rangle$ sign of k_y
 $|0\rangle, |+\rangle$ value of σ_z

$||K_H$

$$\begin{aligned}
 |\psi_0\rangle &= \frac{1}{\sqrt{2}} \left(-e^{i\phi} |0,0\rangle + e^{i k_{z0} z} |0,+\rangle \right) \\
 &= \frac{1}{2} \left(e^{-i k_{z0} z} |+\rangle - e^{i\phi} |0\rangle \right) \otimes |0\rangle
 \end{aligned}$$



$$|\psi_0\rangle = |0,0\rangle$$

$$|\psi_1\rangle = \frac{1}{\sqrt{2}} (|0,0\rangle + |1,0\rangle)$$

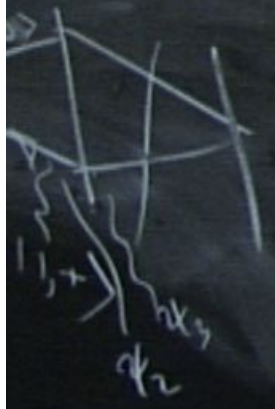
$$|\psi_2\rangle = \frac{1}{\sqrt{2}} (|0,0\rangle + e^{i\phi} |1,+\rangle)$$

$$|\psi_3\rangle = \frac{1}{\sqrt{2}} (|1,0\rangle + e^{i\phi} |0,+\rangle)$$

$$|\psi_4\rangle = \frac{1}{\sqrt{2}} (e^{i\phi} |1,0\rangle + e^{i k_{z0} z} |0,+\rangle)$$

$$|\psi_5\rangle = \frac{1}{2} (e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i k_{z0} z} (|0,+\rangle + |1,+\rangle))$$

sign of k_y
 value of σ_z



$$|1\rangle_{K_H} \quad |\psi_0\rangle = \frac{1}{\sqrt{2}} \left(-e^{i\phi} |0,0\rangle + e^{i k_{z0} z} |0,+ \rangle \right)$$

$$= \frac{1}{2} \left(e^{i k_{z0} z} |+\rangle - e^{i\phi} |0\rangle \right) \oplus |0\rangle$$

then
↑
rule

Variation in k_z

$$|\psi_0\rangle = |0,0\rangle$$

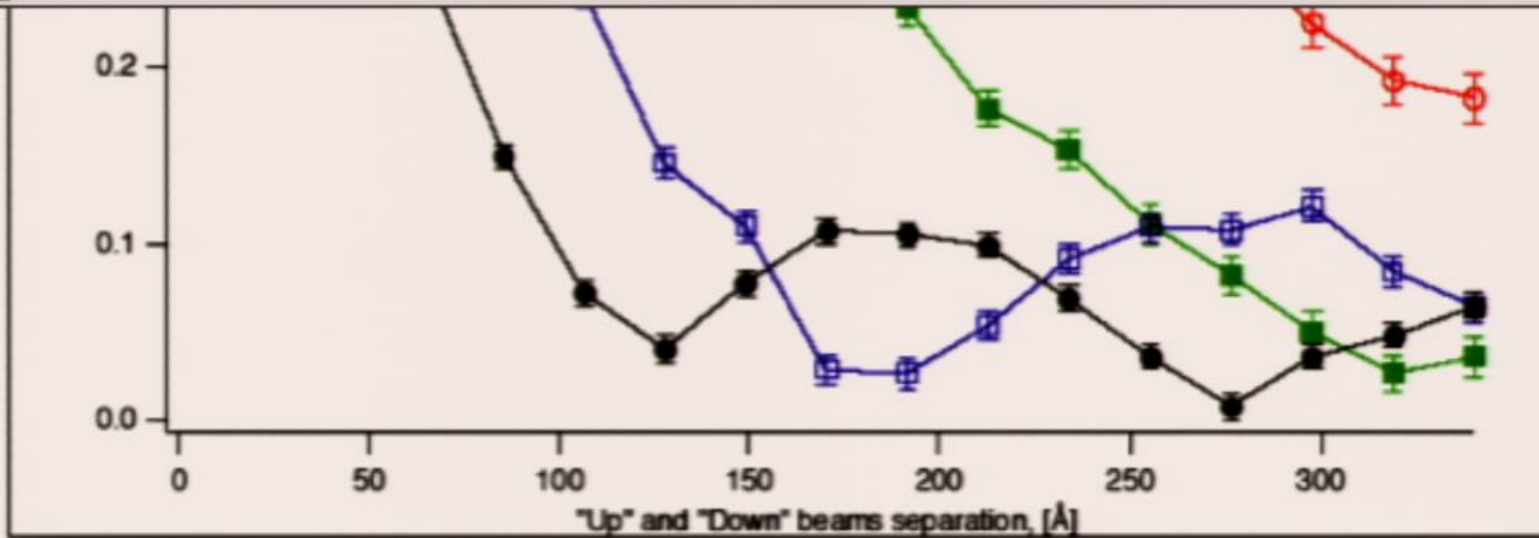
$$|\psi_1\rangle = \frac{1}{\sqrt{2}} (|0,0\rangle + |1,0\rangle)$$

$$|\psi_2\rangle = \frac{1}{\sqrt{2}} (|0,0\rangle + e^{i\phi} |1,+ \rangle)$$

$$|\psi_3\rangle = \frac{1}{\sqrt{2}} (|1,0\rangle + e^{i\phi} |0,+ \rangle)$$

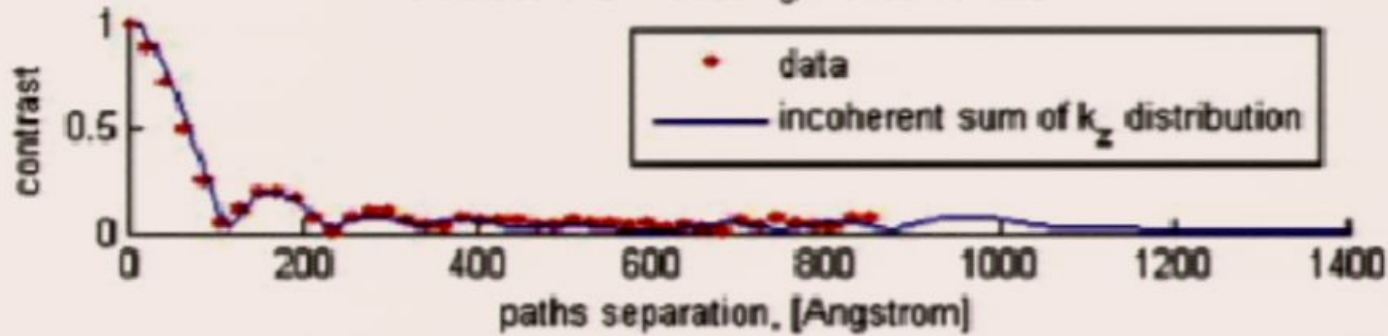
$$|\psi_4\rangle = \frac{1}{\sqrt{2}} (e^{i\phi} |1,0\rangle + e^{i\phi} |0,+ \rangle)$$

$$|\psi_5\rangle = \frac{1}{2} \left(e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i k_{z0} z} (|0,+ \rangle + |1,+ \rangle) \right)$$



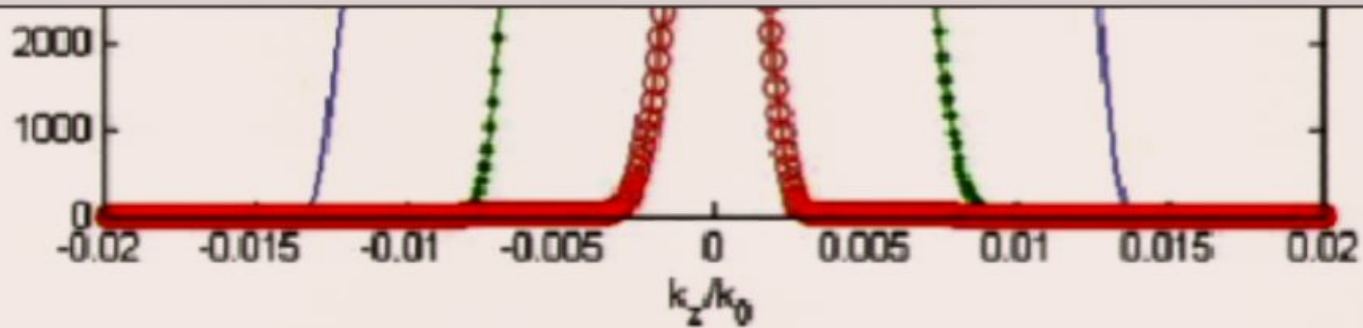
that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is the Fourier transform of the momentum distribution.

9 blades of 2nd focusing monochromator

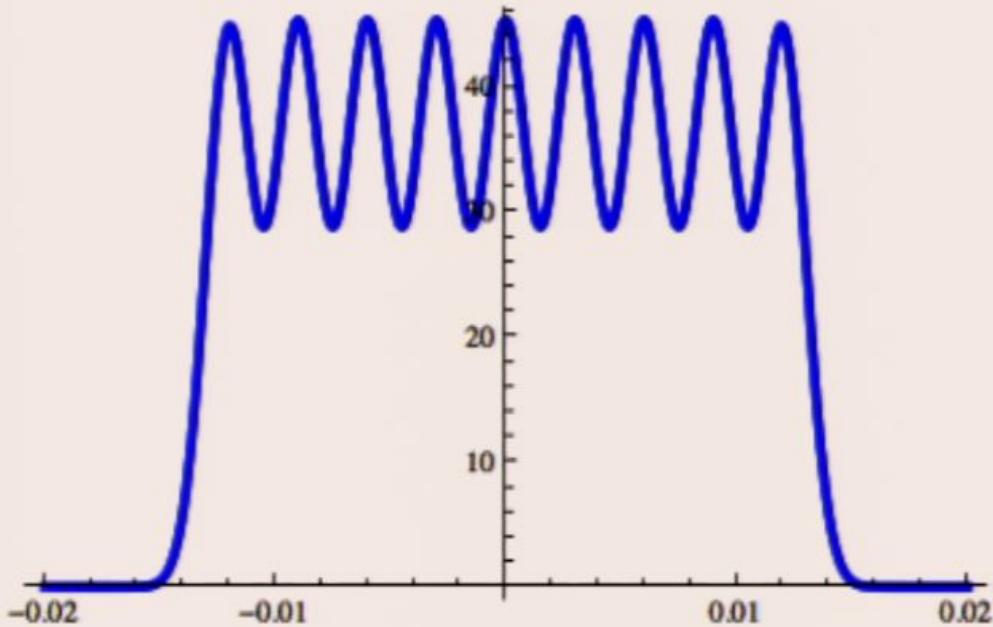


5 blades of 2nd focusing monochromator





```
Pkz[kz_] := Sum[nd[kz - d, 0.001], {d, 0.012, -0.012, -0.003}]/9
Plot[Pkz[kz], {kz, -0.02, 0.02}, PlotStyle -> {RGBColor[0, 0, 1], Thickness[0.01]}]
```

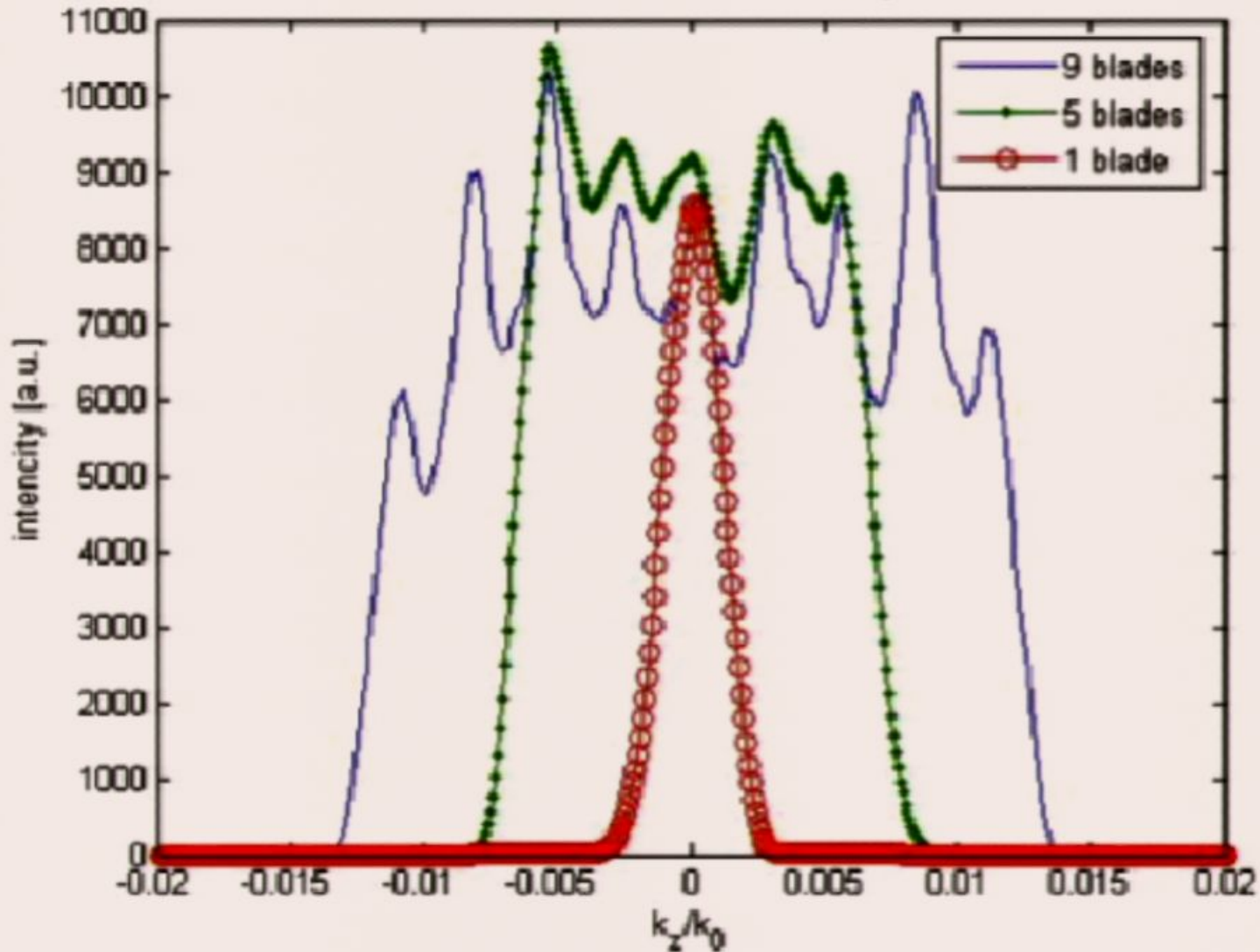


```
M160[z_, a_] := Sum[Pkz[kz] Tr[Exp . res14[kz, z, a]], {kz, -0.02, 0.02, 0.001}]
```

```
M16R[z_, a_] := Sum[Pkz[kz] Tr[Ezm . res14[kz, z, a]], {kz, -0.02, 0.02, 0.001}]
```

```
ListAnimate[{Plot[Re[M160[0, a]], {a, 0, 2 π},
```

Vertical Neutron Momentum Distribution ($k_0 = 2.3 \times 10^{10} \text{ m}^{-1}$)

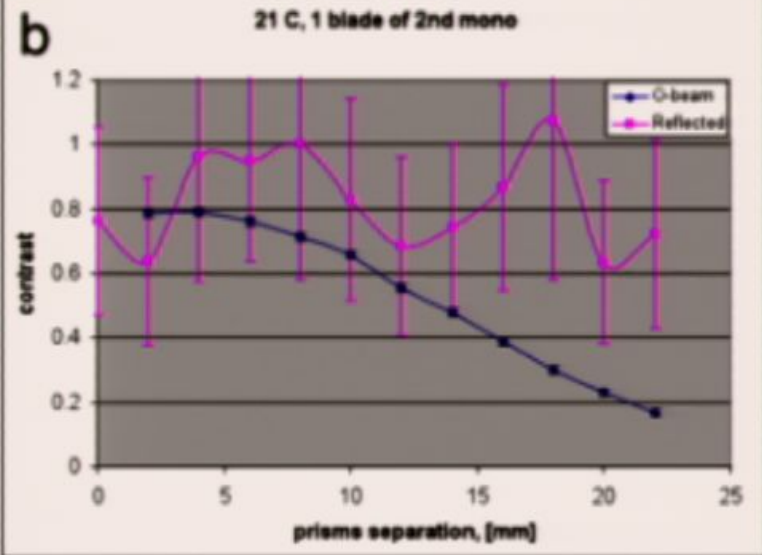
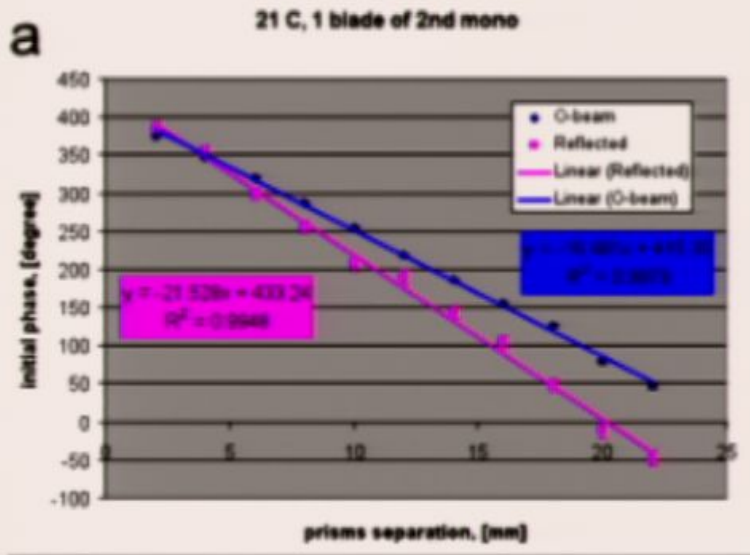


```
PKz[kz_] := Sum[nd[kz - d, 0.001], {d, 0.012, -0.012, -0.003}]/9
```

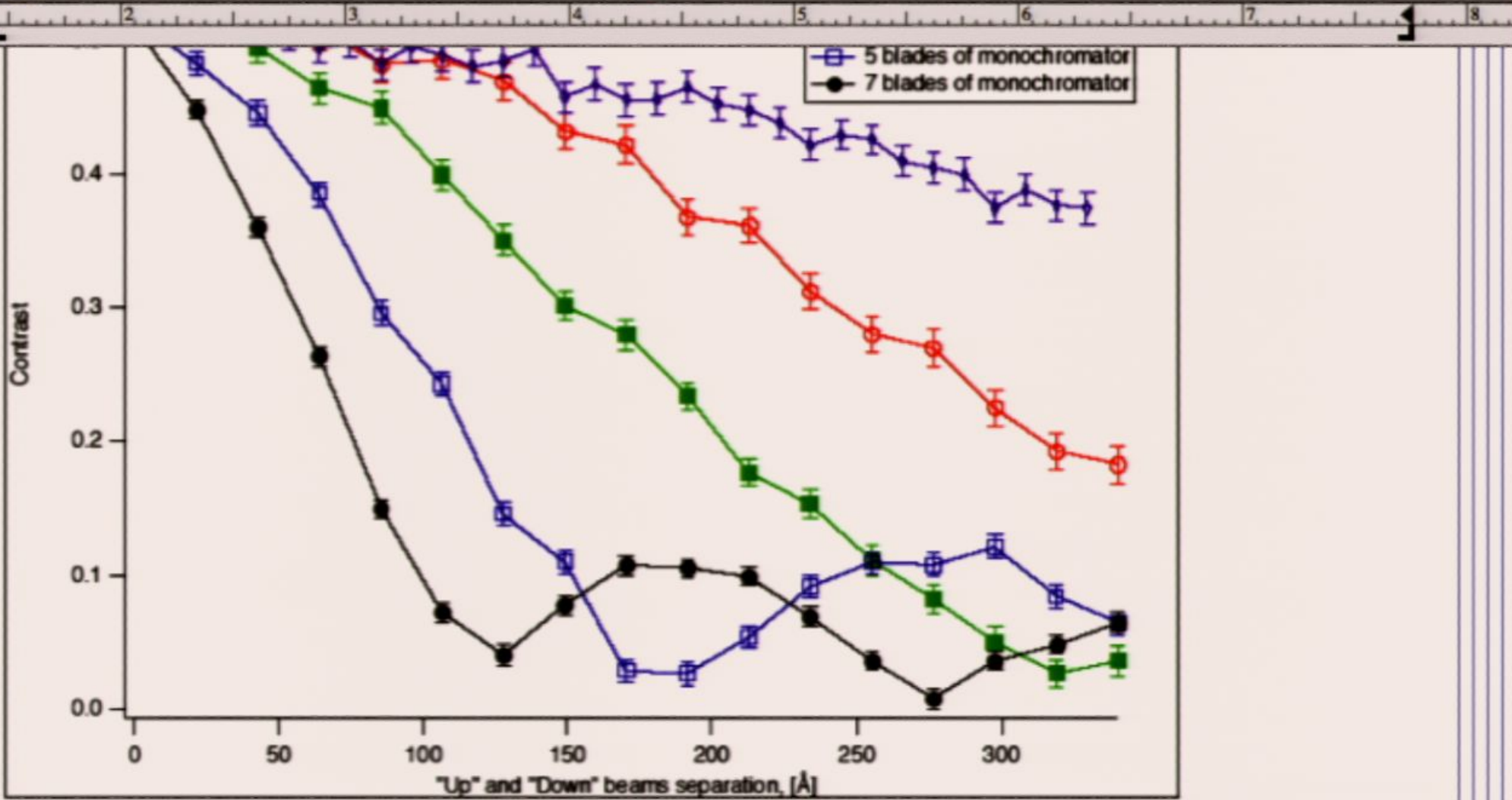
```
Plot[PKz[kz], {kz, -0.02, 0.02}, PlotStyle -> {RGBColor[0, 0, 1], Thickness[0.01]}]
```




data is noisy but shows that the scattered (recombined) beam retains contrast, while the direct beam loses contrast with increased separation of the beams.



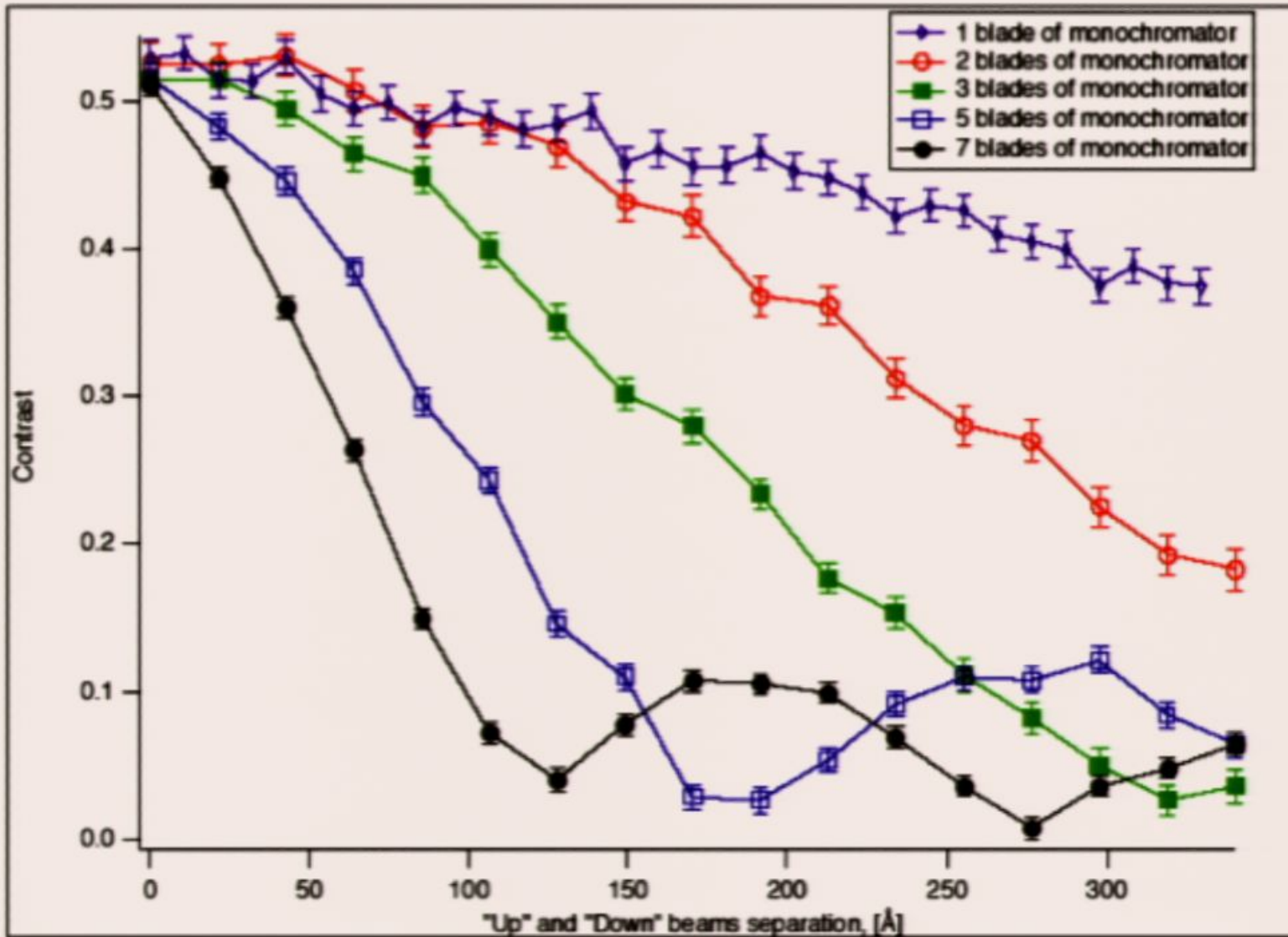
point to remember here is that the contrast in the transmitted beam is lost due to a distribution of momenta. the neutron momenta in free space do not change, this can easily be recovered. Also note that we can compose any state of momenta through a series of such filters.



that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is the Fourier transform of the momentum distribution.

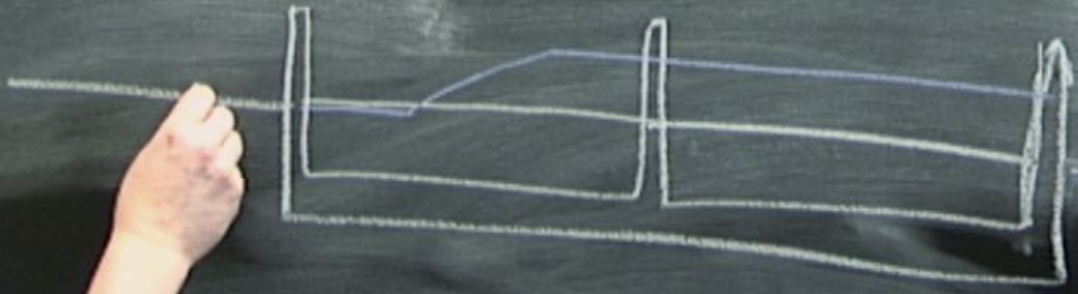
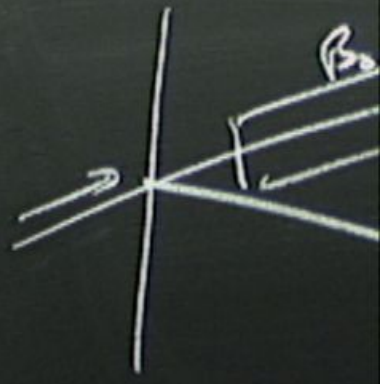
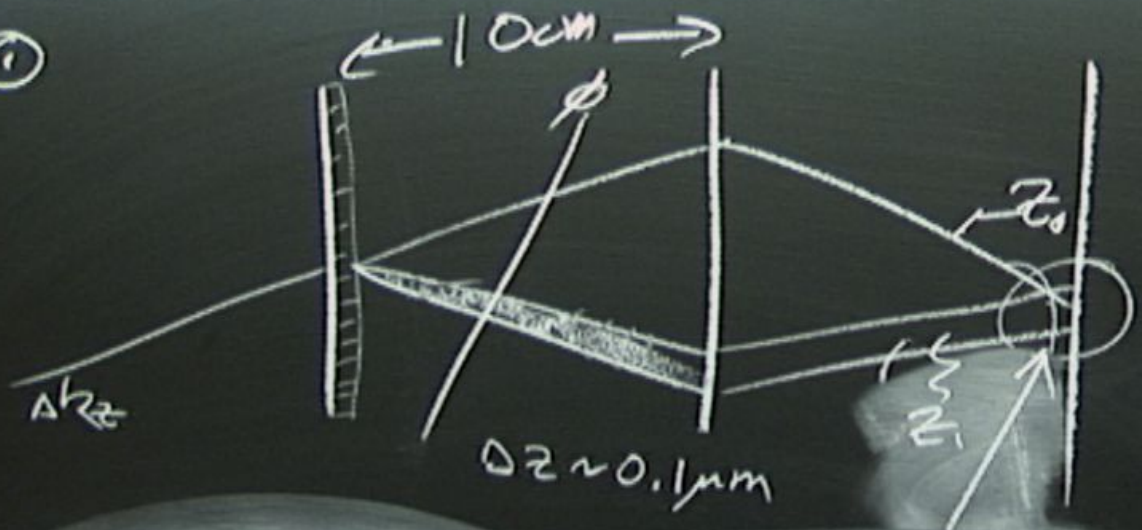
9 blades of 2nd focusing monochromator



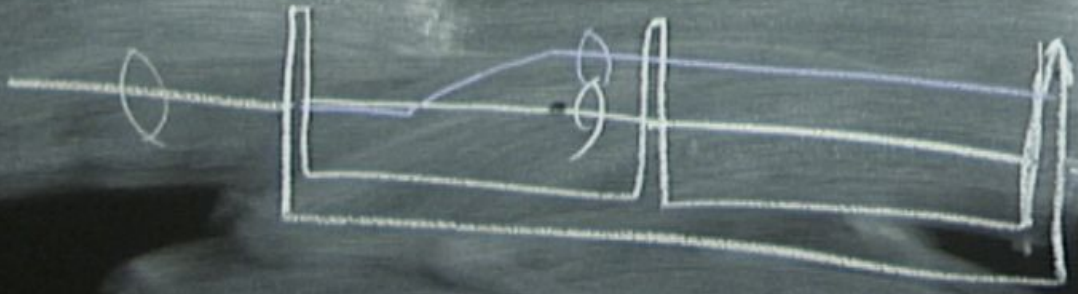
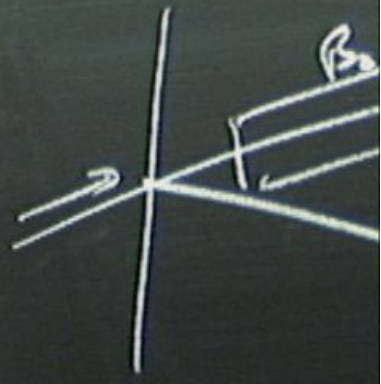
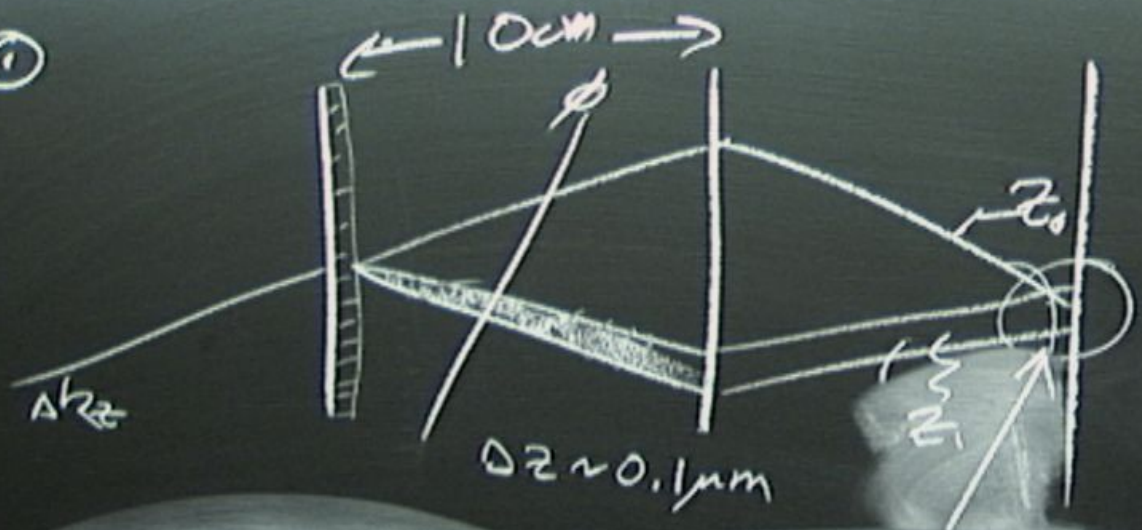


that the beats in the above data are predicted since the momentum spread has features. The Coherence curve is the Fourier transform of the momentum distribution.

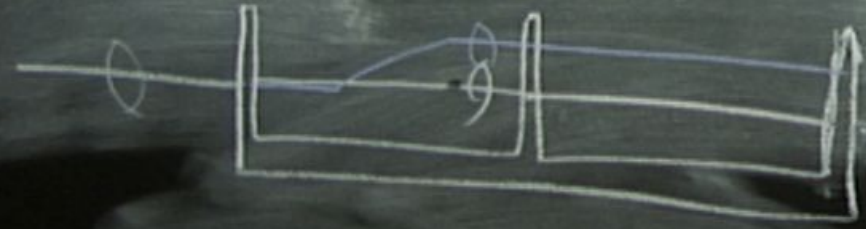
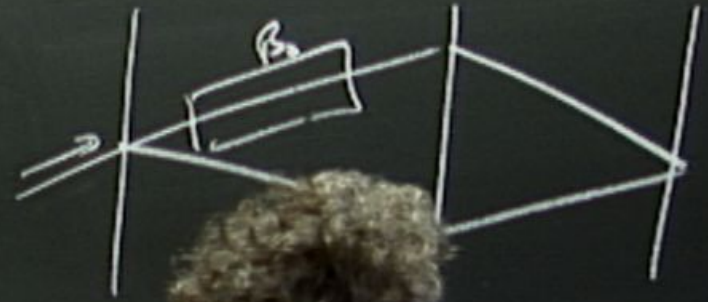
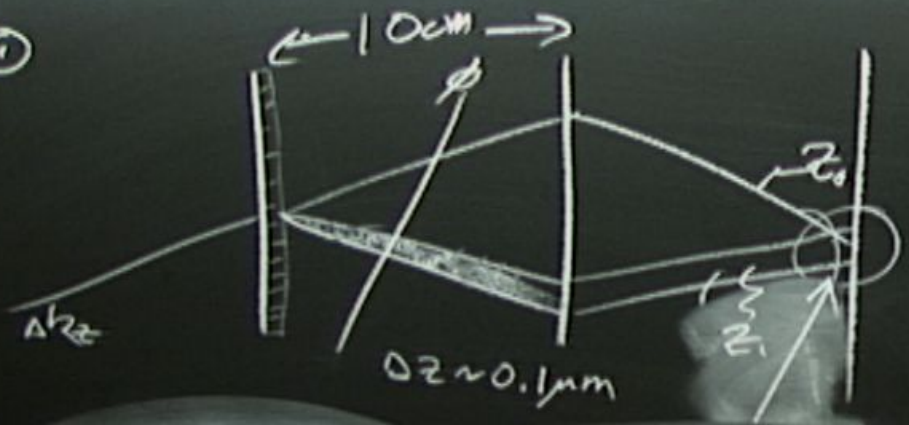
z ⊙



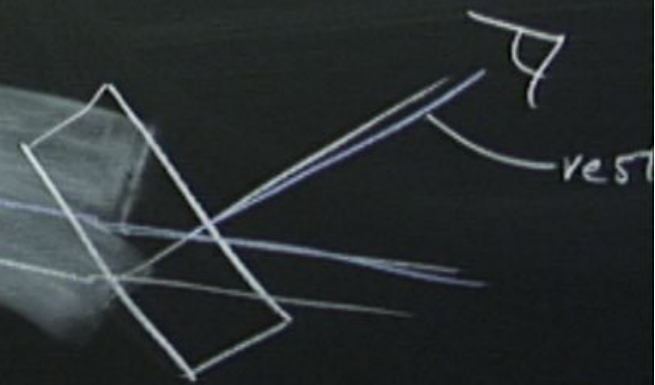
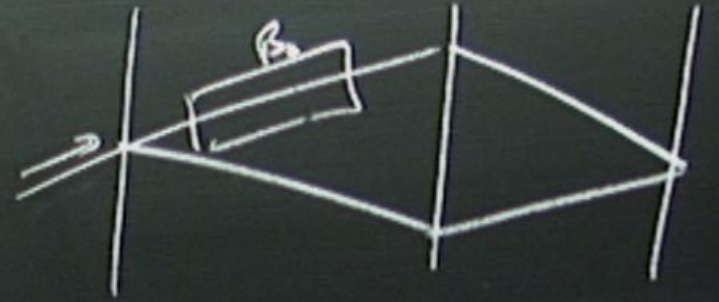
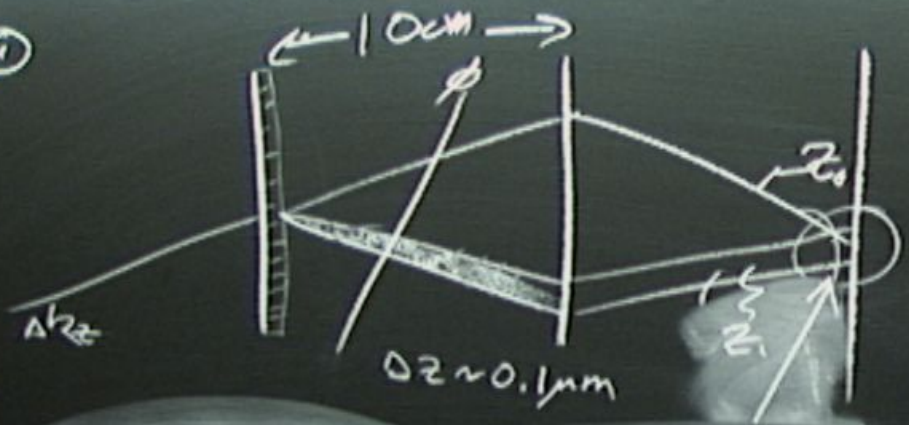
20

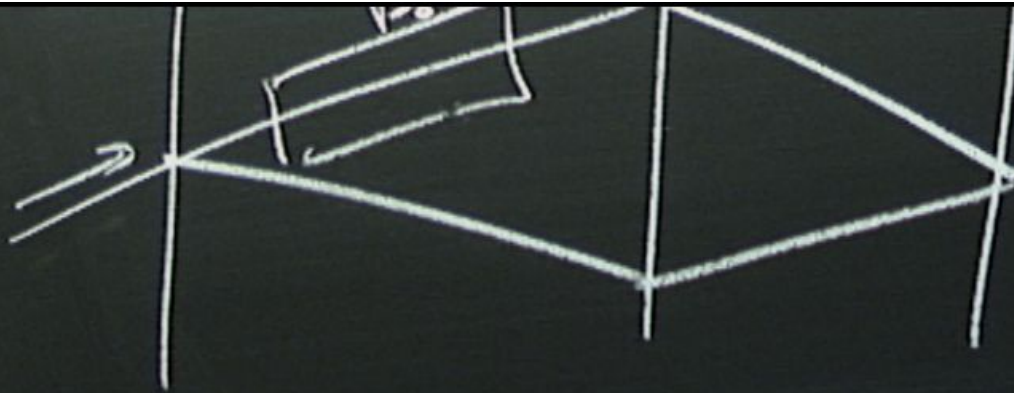


20

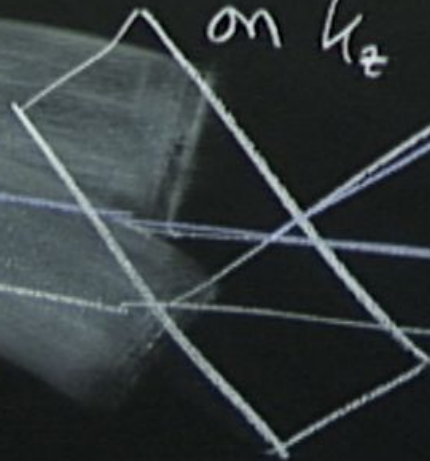


2 ①



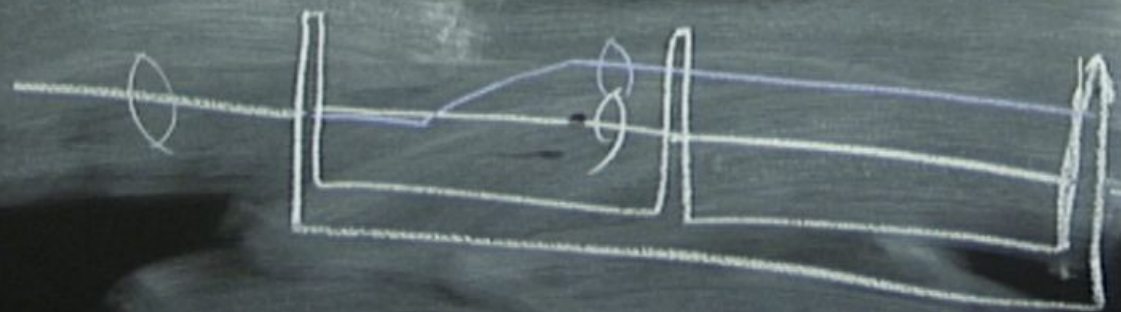
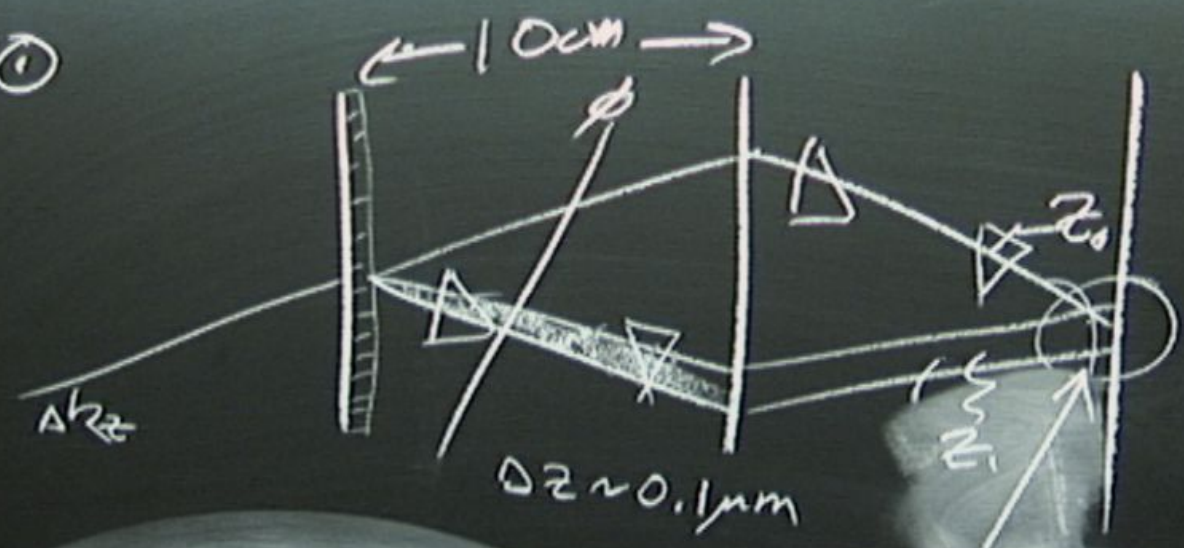


Post selection
on k_e

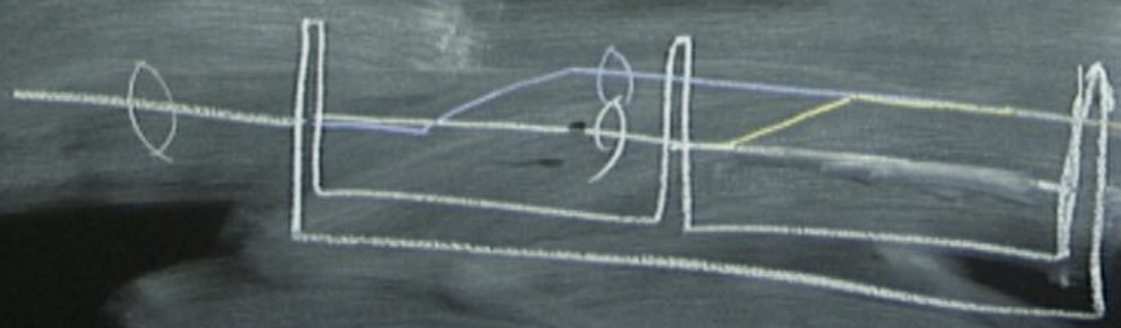
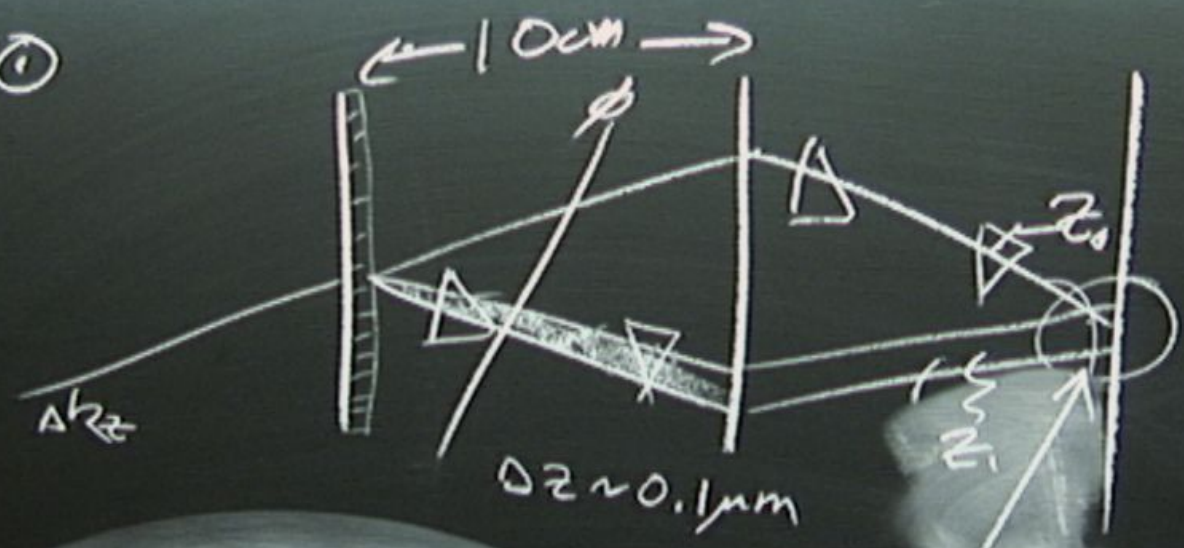


restoro content

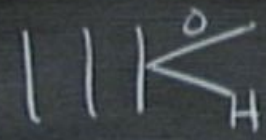
z ①



z ①



sign of k_y
 value of s_z



$$|\psi_0\rangle = \frac{1}{\sqrt{2}} \left(-e^{i\phi} |0,0\rangle + e^{i k_z a z} |0,+ \rangle \right)$$

$$= \frac{1}{2} \left(e^{-i k_z a z} |+\rangle - e^{i\phi} |0\rangle \right)$$

Variation in k_z

$$|\psi_0\rangle = |0,0\rangle$$

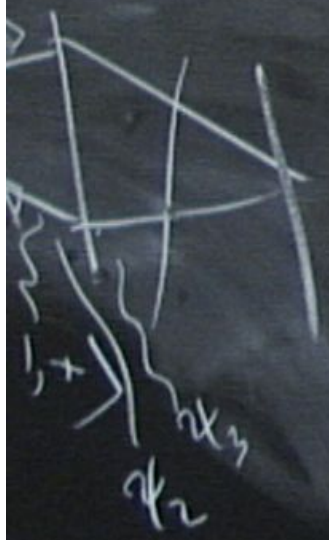
$$|\psi_1\rangle = \frac{1}{\sqrt{2}} (|1,0\rangle + |1,0\rangle)$$

$$|\psi_2\rangle = \frac{1}{\sqrt{2}} (|0,0\rangle + e^{i\phi} |1,+ \rangle)$$

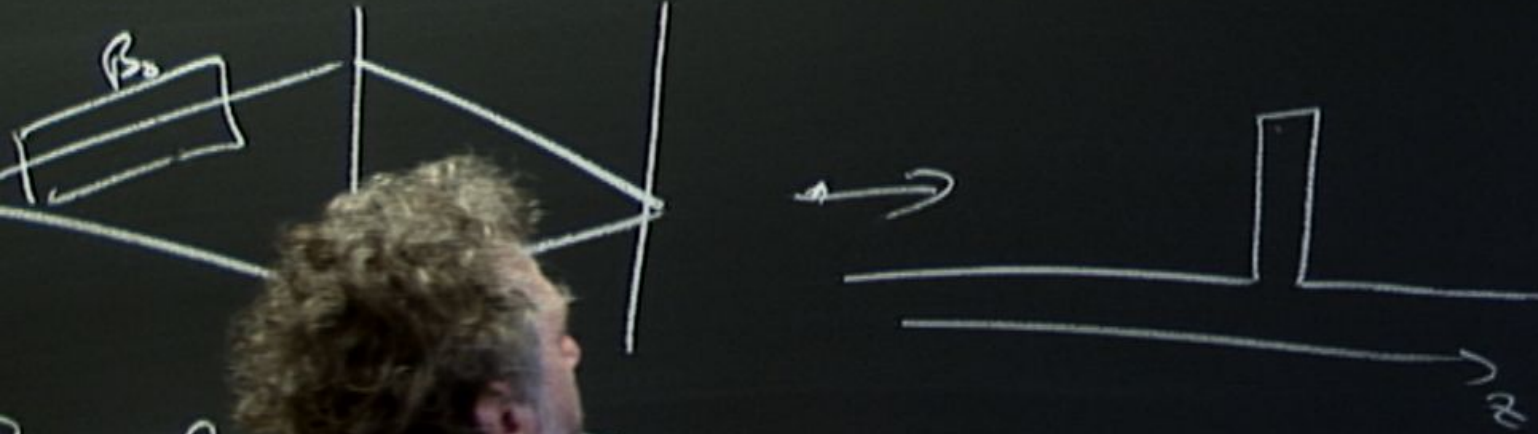
$$|\psi_3\rangle = \frac{1}{\sqrt{2}} (|1,0\rangle + e^{i\phi} |0,+ \rangle)$$

$$|\psi_4\rangle = \frac{1}{\sqrt{2}} (e^{i\phi} |1,0\rangle + e^{i\phi} |0,+ \rangle)$$

$$|\psi_5\rangle = \frac{1}{2} \left(e^{i\phi} (|1,0\rangle - |0,0\rangle) + e^{i k_z a z} (|0,+ \rangle + |1,+ \rangle) \right)$$



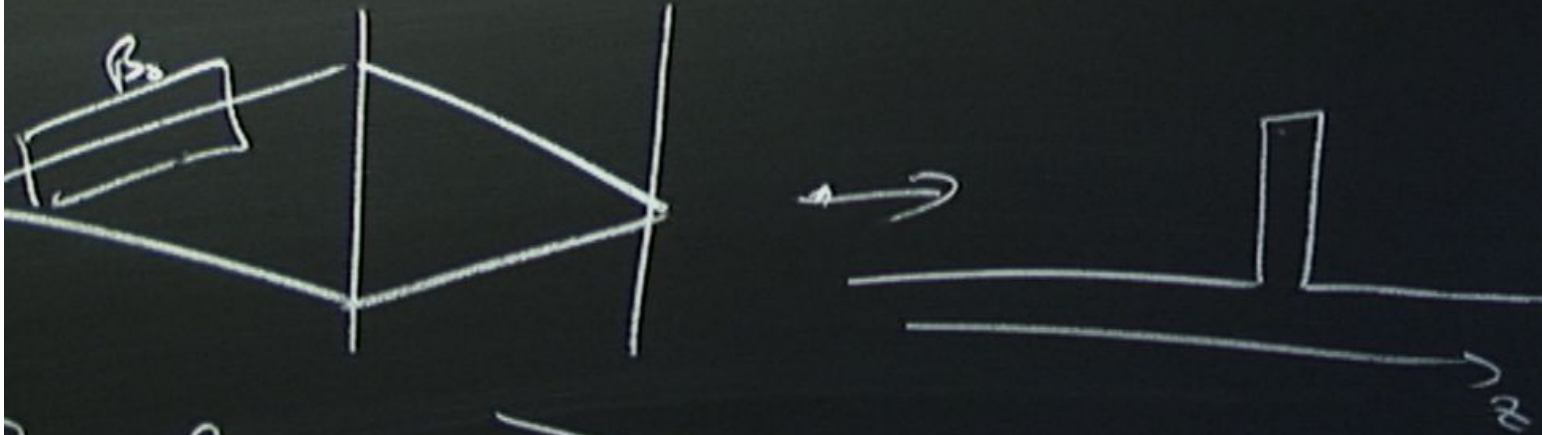
μB_0



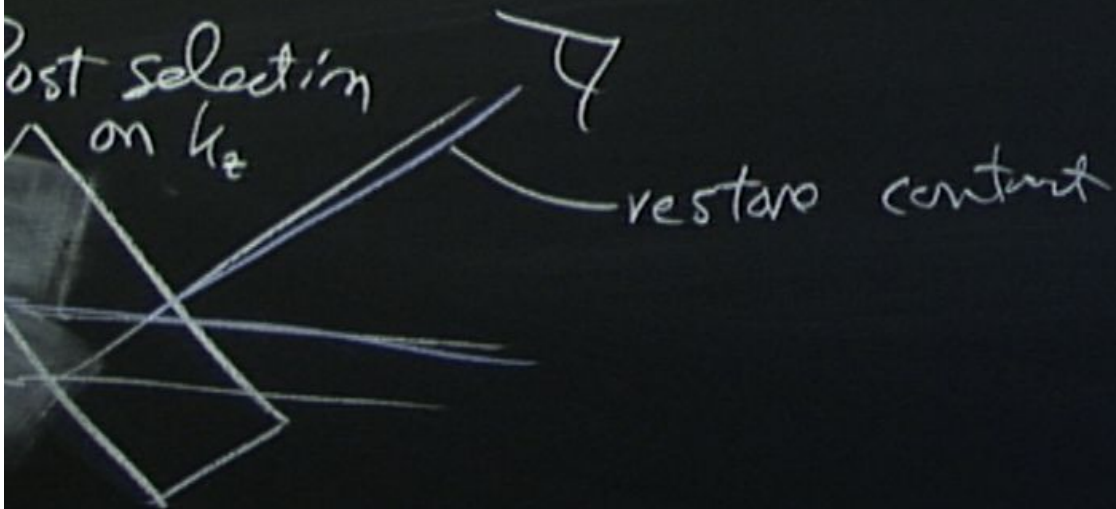
Post selection
on k_z

restored content

μB_0

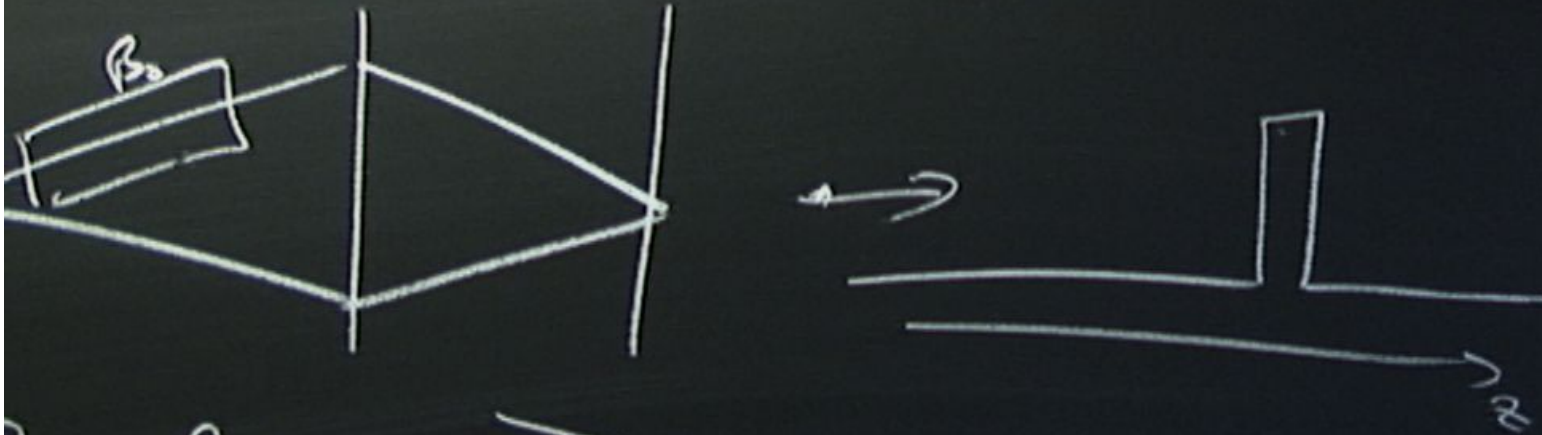


Post selection
on k_z



restoro content

$1. B_0$

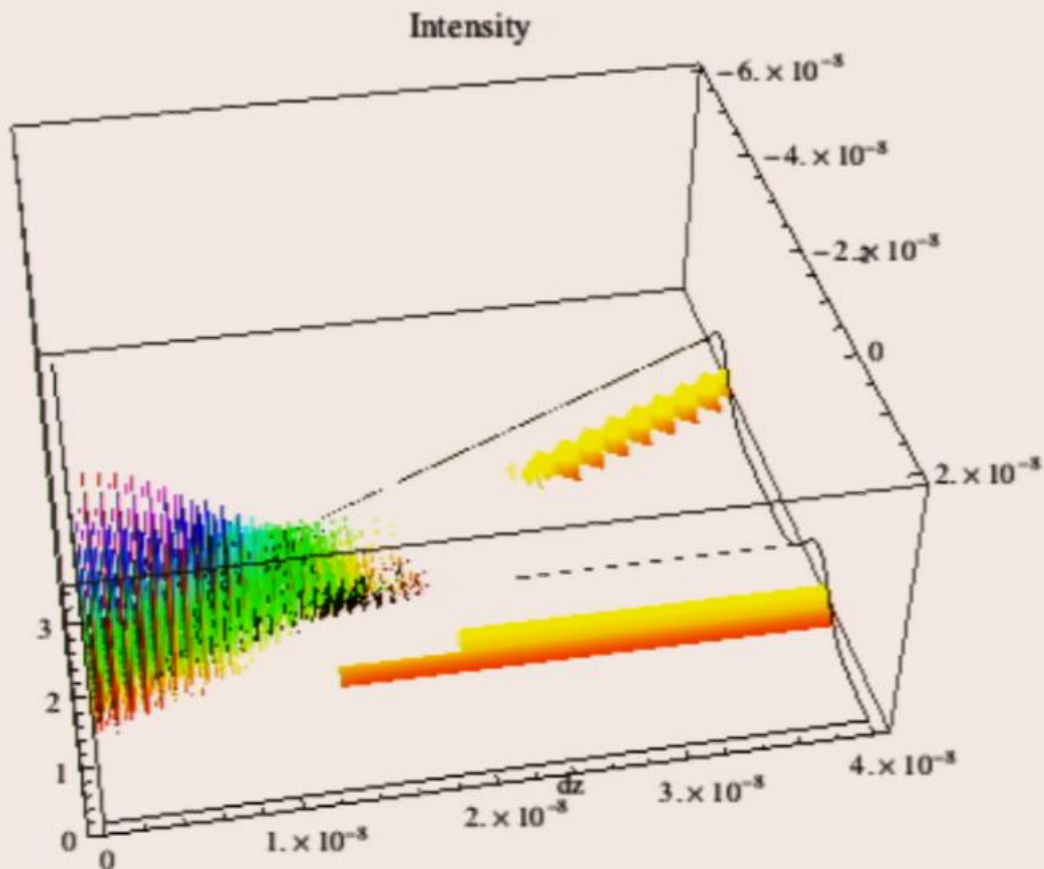


Post selection
on k_2



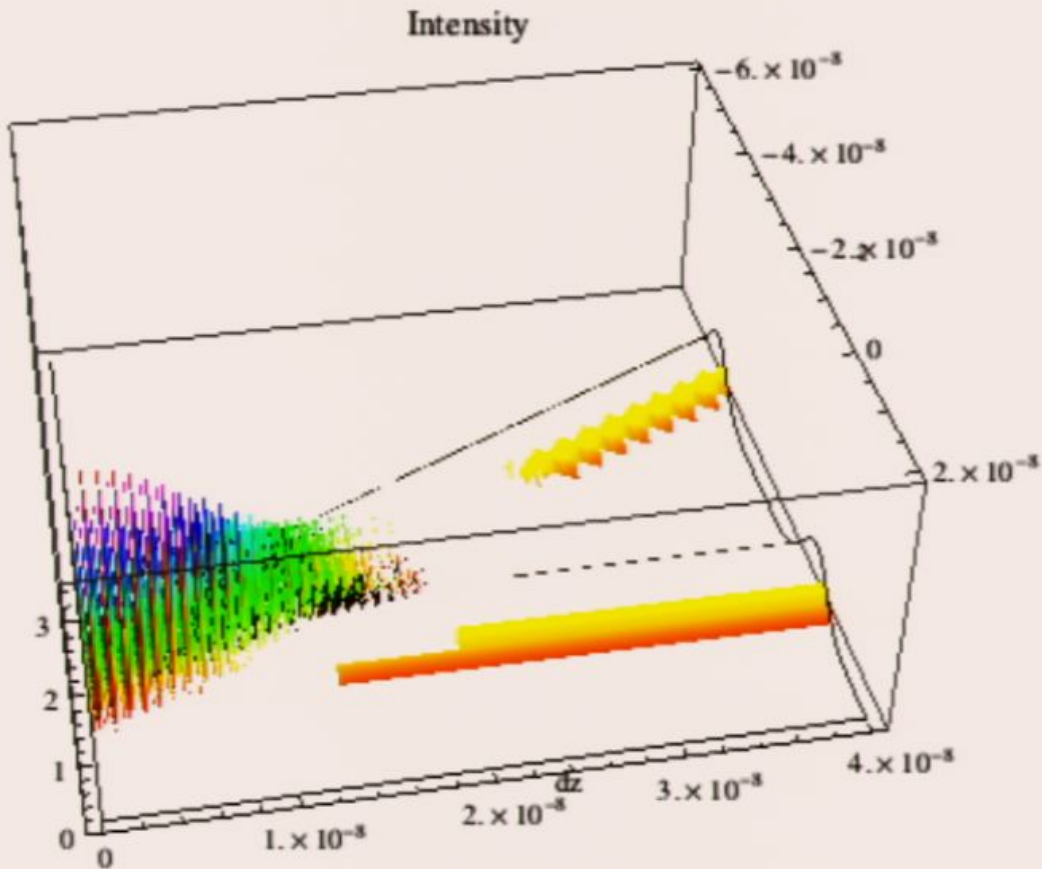
Γ
restoro content

```
Plot3D[Intensity[z, dz], {z, -600*10^(-10)}, 200*10^(-10)},
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5}, PlotPoints -> 50,
MaxRecursion -> 3, Mesh -> None, ColorFunction -> Function[{x, y, z}, Hue[z]],
PlotLabel -> "Intensity", AxesLabel -> Automatic]
```



```
ContourPlot[Intensity[z, dz], {z, -600*10^(-10)}, 200*10^(-10)},
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5},
Contours -> Function[{min, max}, Range[min, max, 0.05]],
```

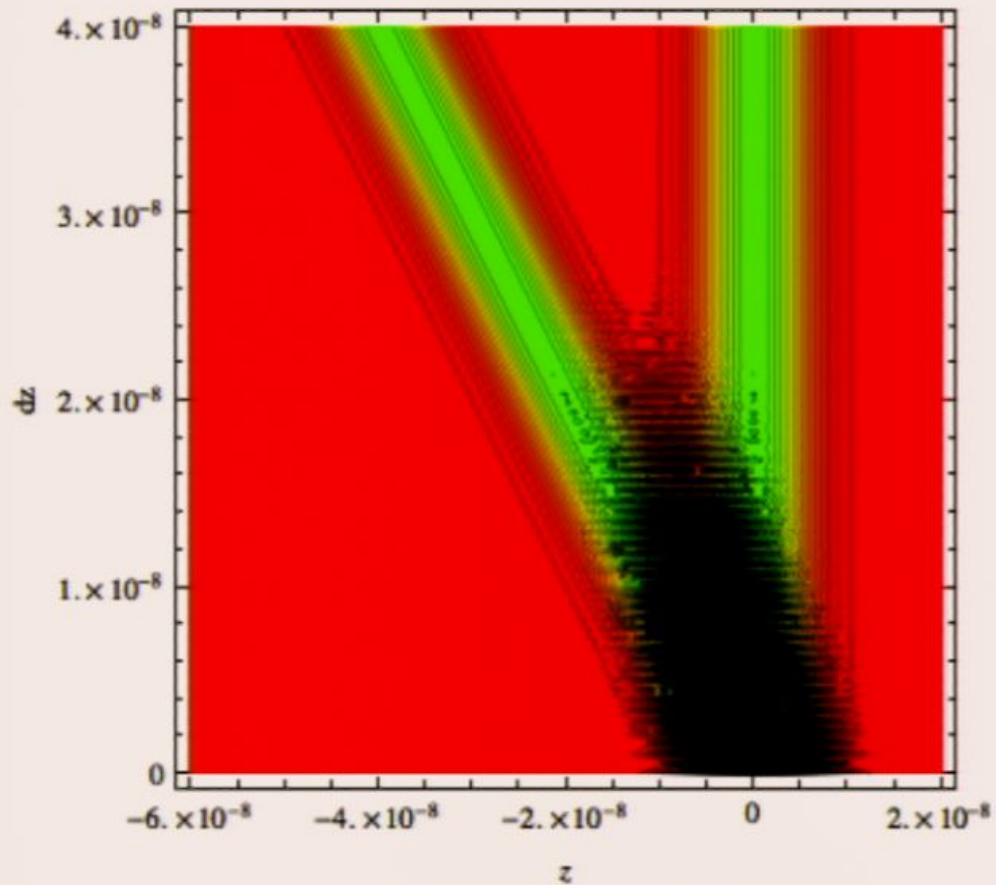
```
Plot3D[Intensity[z, dz], {z, -600*10^(-10)}, 200*10^(-10)},
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5}, PlotPoints -> 50,
MaxRecursion -> 3, Mesh -> None, ColorFunction -> Function[{x, y, z}, Hue[z]],
PlotLabel -> "Intensity", AxesLabel -> Automatic]
```



```
ContourPlot[Intensity[z, dz], {z, -600*10^(-10)}, 200*10^(-10)},
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5},
Contours -> Function[{min, max}, Range[min, max, 0.05]],
```

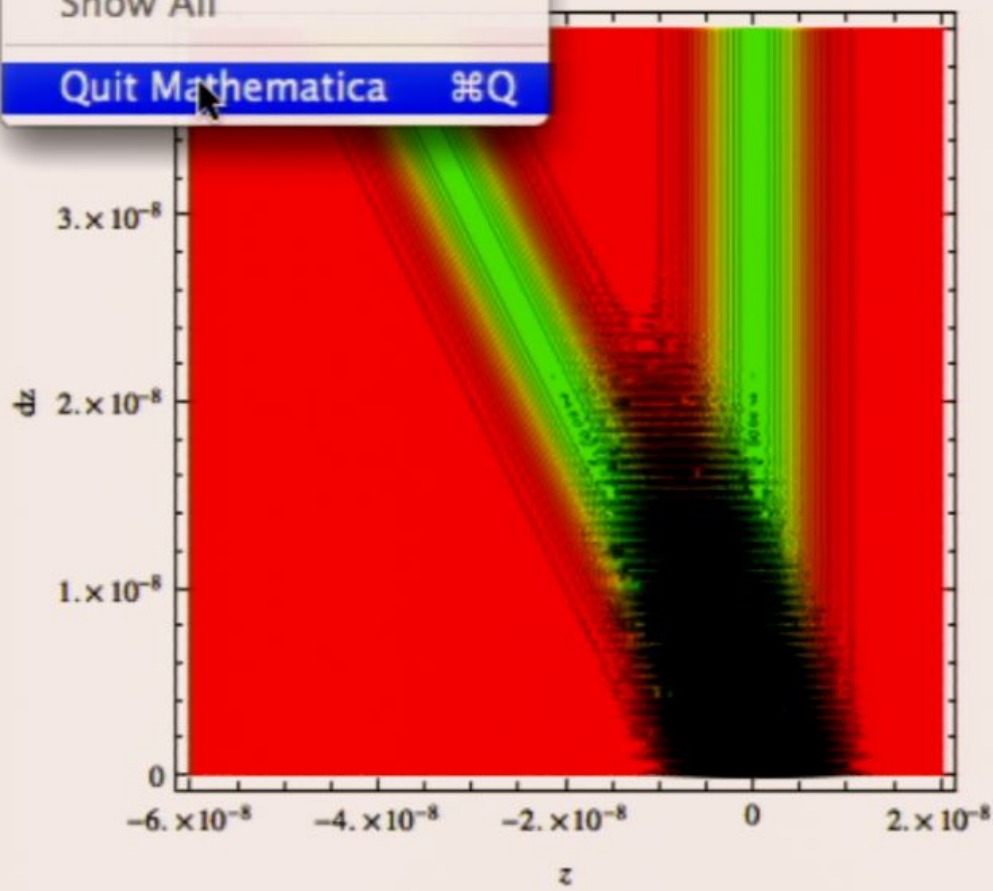


```
ContourPlot[Intensity[z, dz], {z, -600*10^(-10), 200*10^(-10)},  
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5},  
Contours -> Function[{min, max}, Range[min, max, 0.05]],  
ColorFunction -> Hue, FrameLabel -> {z, dz}]
```

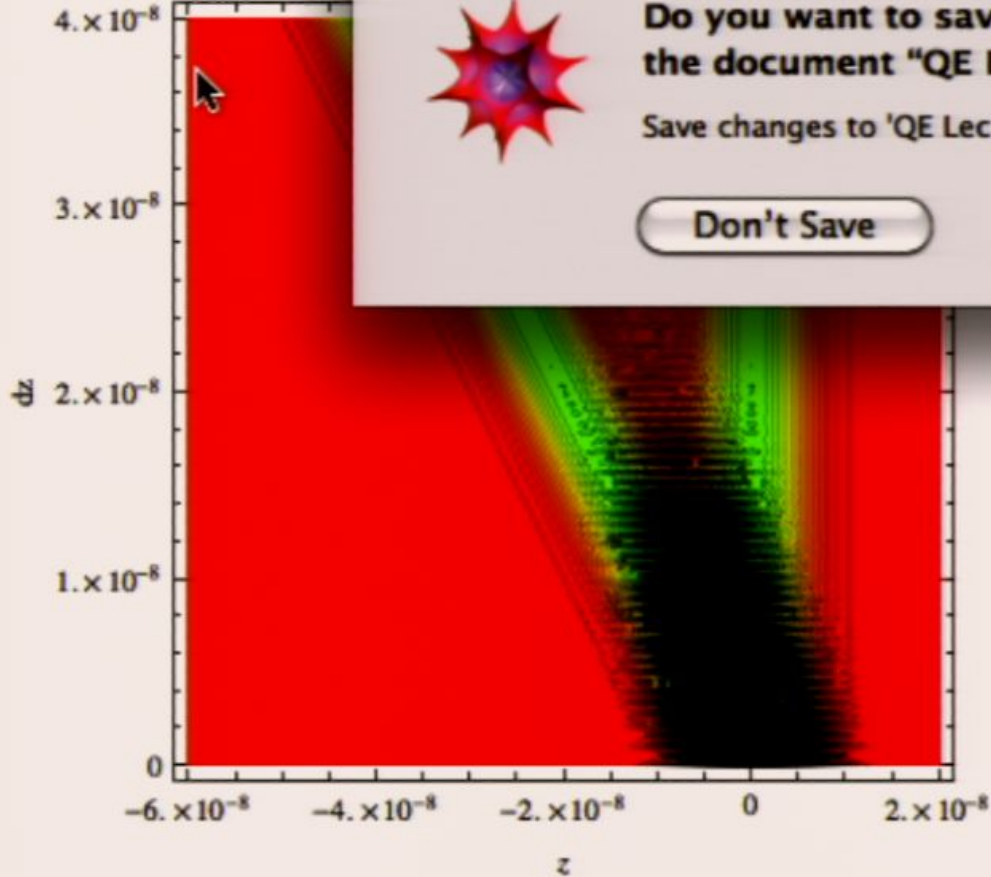


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```
{z, -600*10^(-10), 200*10^(-10)},  
PlotRange -> {0, 3.5},  
max}, Range[min, max, 0.05]],  
Label -> {z, dz}]
```



```
ContourPlot[Intensity[z, dz], {z, -600*10^(-10), 200*10^(-10)},  
{dz, 0, 400*10^(-10)}, PlotRange -> {0, 3.5},  
Contours -> Funct:  
ColorFunction ->
```



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