

Title: Constraints on dark matter annihilation in Milky Way halo from spherical harmonics of Fermi gamma-rays

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Abstract: Gamma-ray production by dark matter annihilation is one of the most universal indirect dark matter signals. In order to avoid intensive astrophysical background, one can study the gamma-rays away from the Galactic plane. The problem is that the dark matter annihilation signal at high latitudes is smooth and most probably subdominant to Galactic and extragalactic fluxes. I will discuss the use of spherical harmonics decomposition as a tool to distinguish a large scale small amplitude dark matter signal from astrophysical fluxes. The sensitivity of this method for currently available Fermi data is similar to the signal from thermal WIMP dark matter annihilation into  $W+W-$  or  $b\bar{b}$ .

## Outline

### I. Motivation

DM annihilation at high latitudes is a **large scale small amplitude** signal: ``Fourier transform'' is a tool of choice for such signals

### 2. Illustration of the idea: Fourier transform of a signal on an interval

### 3. Fitting the large scale structure in the space of spherical harmonics

## Sources of gamma-ray flux

### 1. Extragalactic sources:

1. Point sources
2. Isotropic flux (very faint point sources)

### 2. Galactic sources:

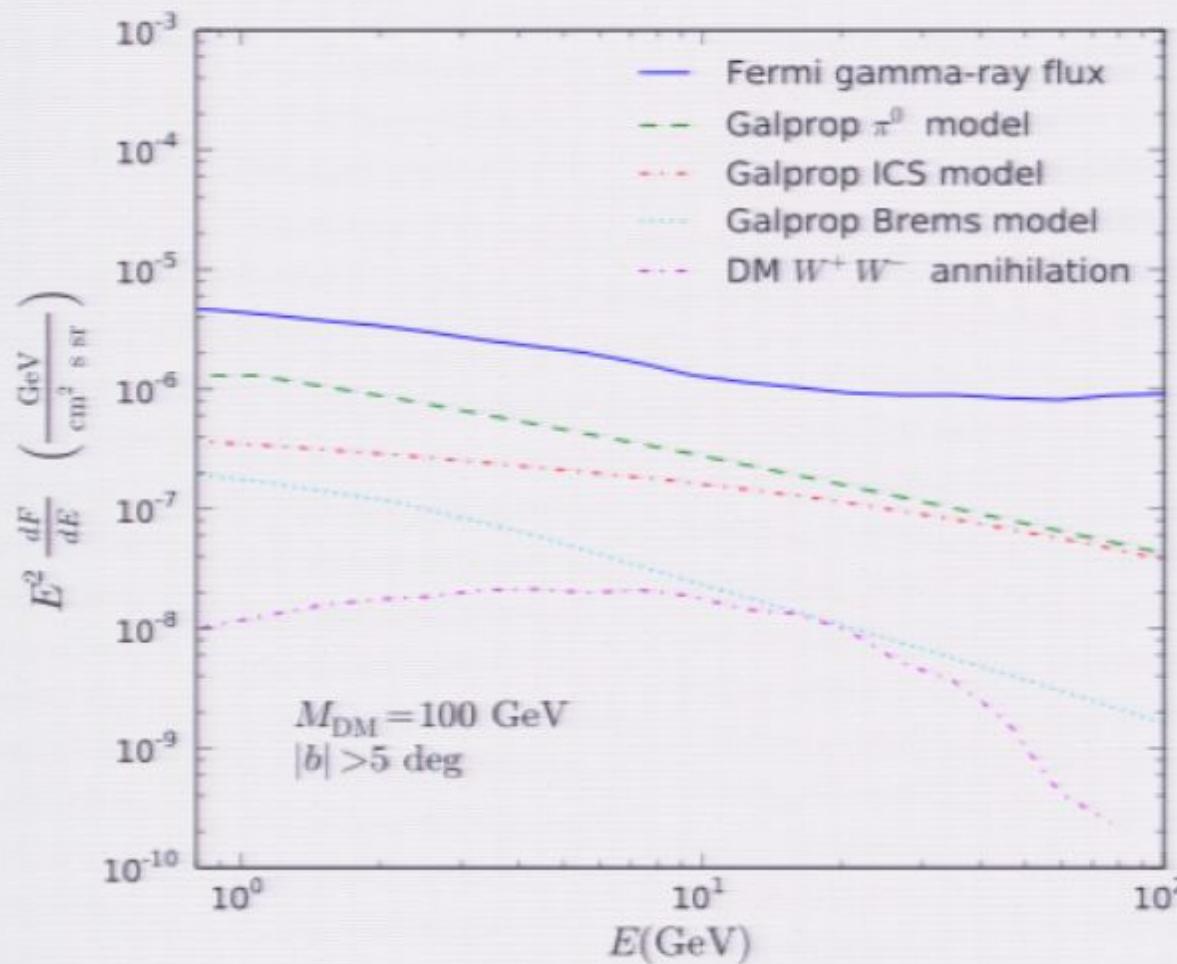
1. Point sources
2. Pion production
3. Inverse Compton scattering
4. Bremsstrahlung

### 3. Exotic sources:

1. New astrophysics, e.g., past AGN activity
2. Dark matter annihilation or decay

## Sources of gamma-ray flux

(Galactic plane and point sources are masked)



Expected DM annihilation signal is about 1% from the total gamma-ray flux

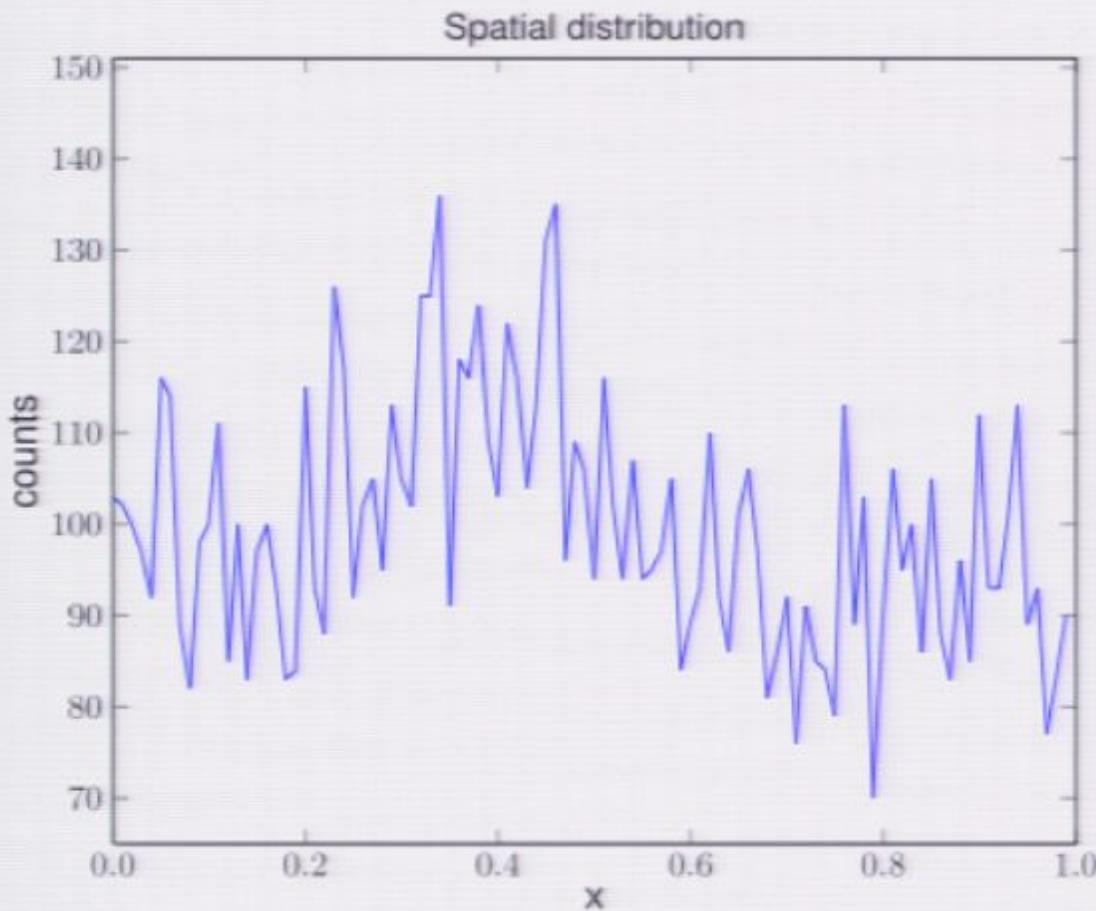
DM mass and annihilation cross section:

$$M_{\text{DM}} = 100 \text{ GeV}$$

$$\sigma v = 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

DM distribution:  
NFW profile

# One dimensional random map with a small signal

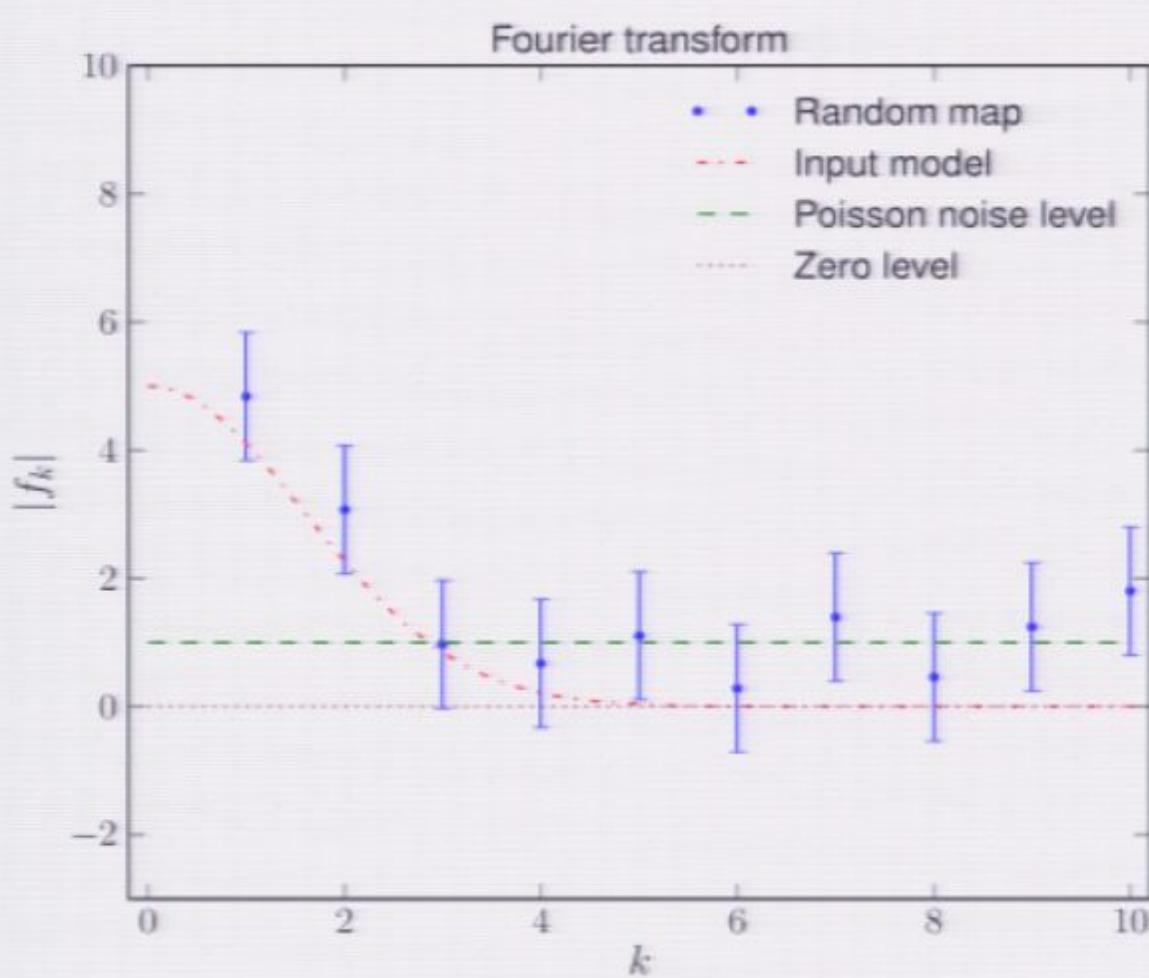


This map has a significant homogeneous component and a small Gaussian distribution of points.

There is a clear excess around  $x = 0.4$ , but how significant is it and what is its amplitude?

Number of events:  $N = 10^4$

Number of bins:  $N_{\text{bins}} = 10^2$



$$f_k = \frac{1}{\sqrt{N}} \int_0^1 dx f(x) e^{2\pi i k x}$$

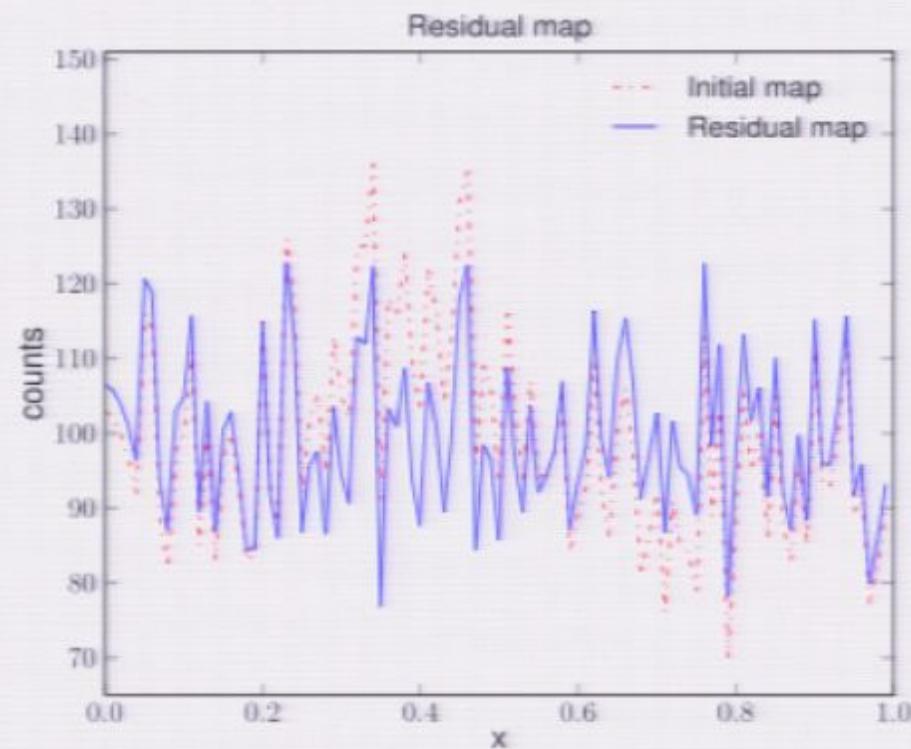
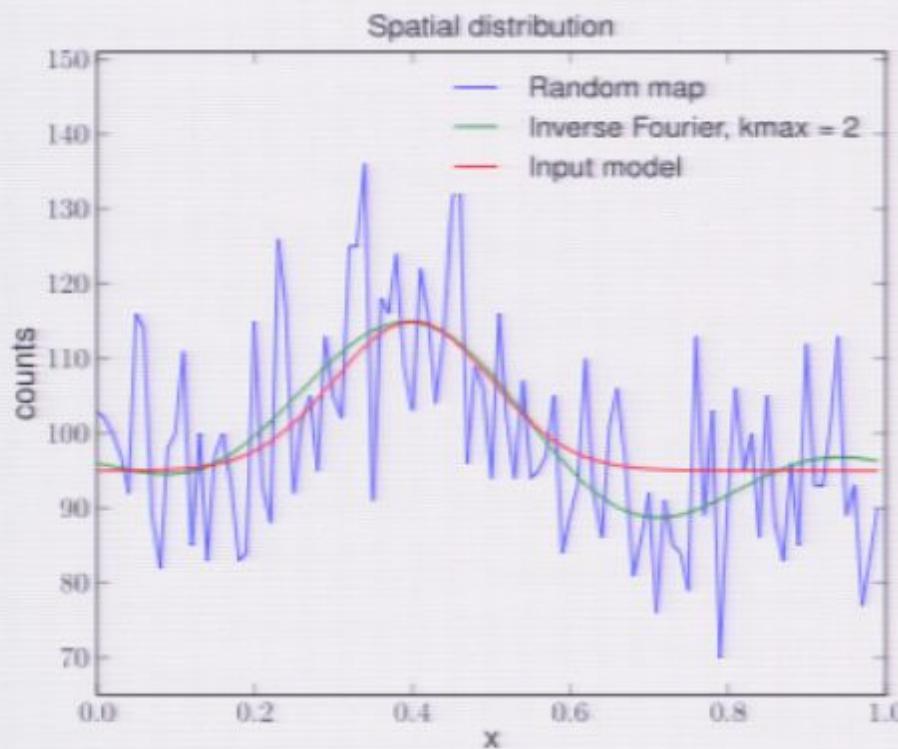
The normalization is chosen such that for the Poisson noise  $\langle |f_k|^2 \rangle = 1$

The Poisson noise level also gives an estimate of uncertainty of  $f_k$ 's

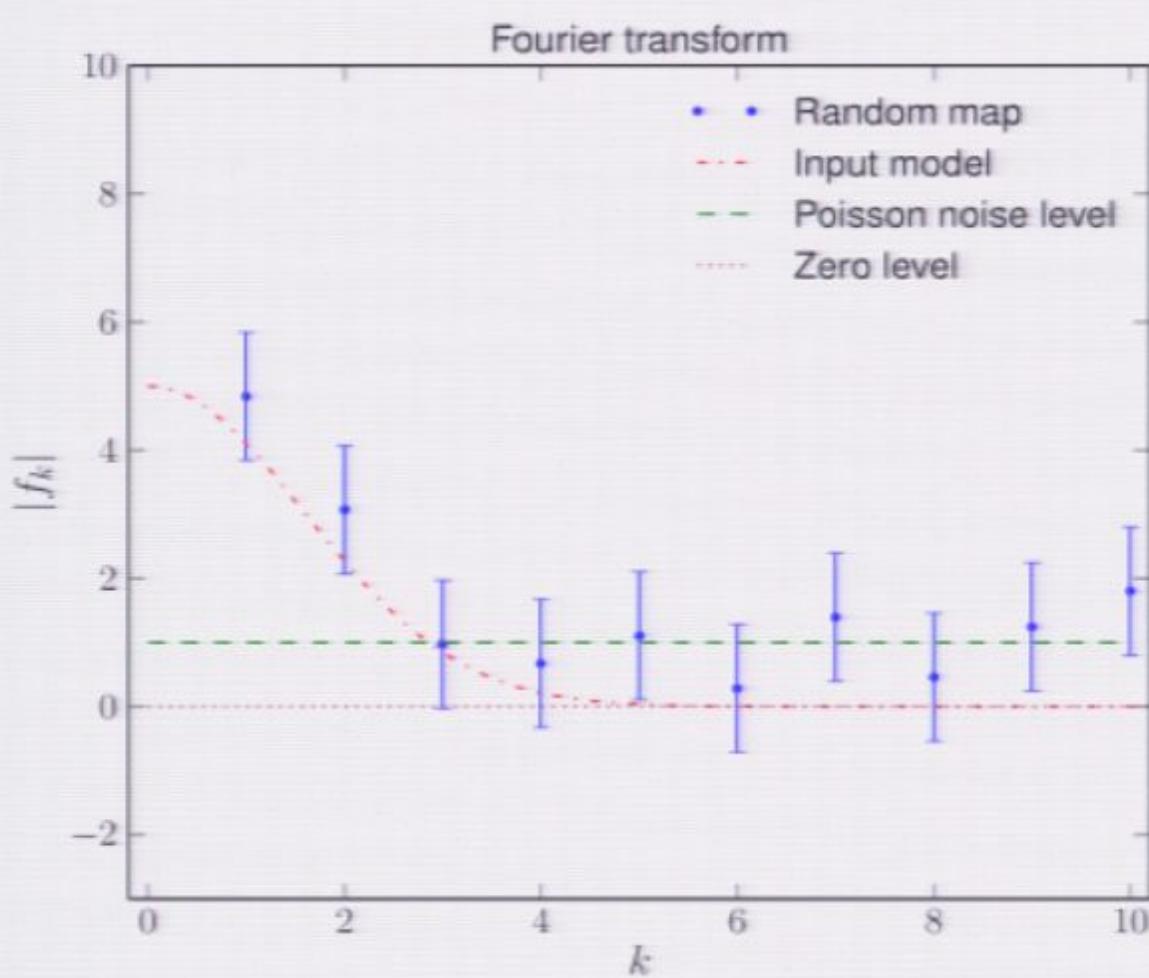
The significance of  $f_1$  is about 5 sigma, the significance of  $f_2$  is about 3 sigma.

The non-homogeneous contribution is  $\sqrt{N} f_0 = 500$  which is 5% in this case

## Inverse Fourier transform and the residual map



The inverse Fourier transform (**green line**)  
of the first non-trivial harmonics does an OK job  
in approximating the form of the excess (**red line**)



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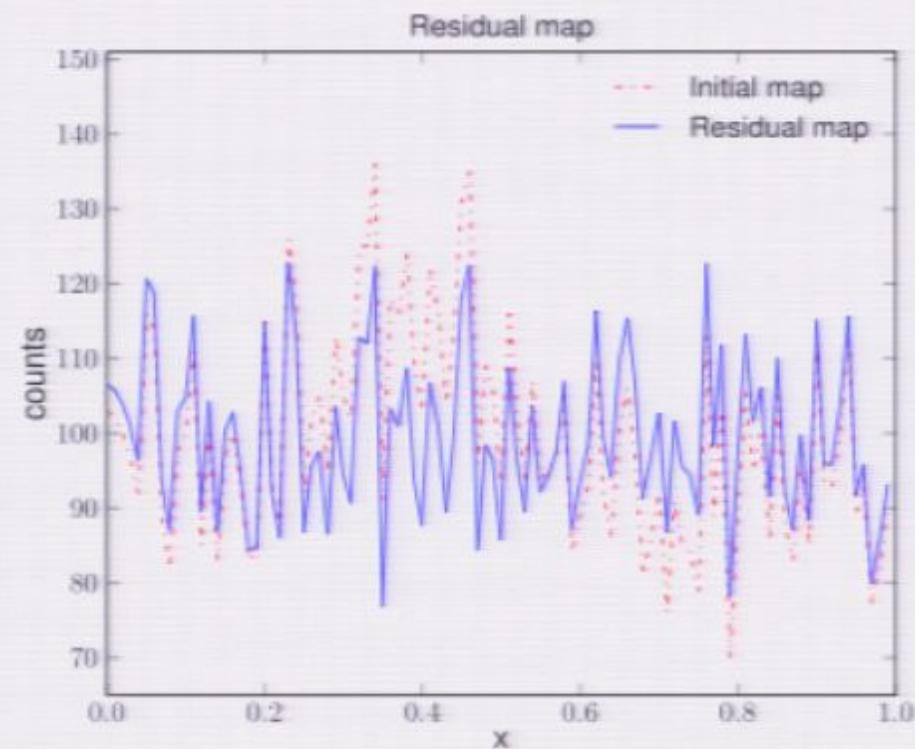
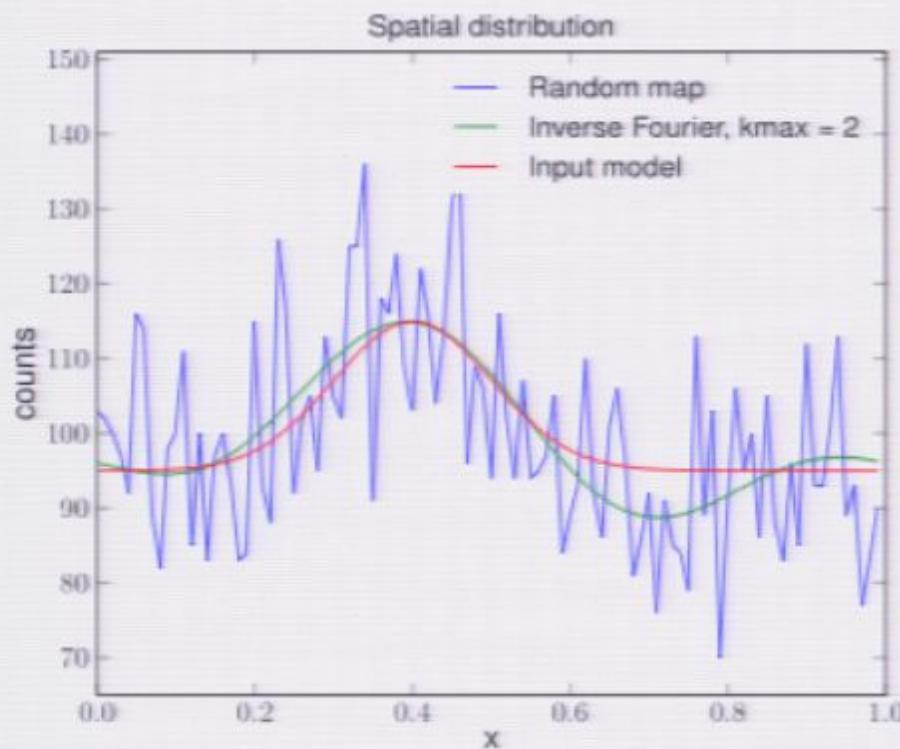
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## Fitting in Fourier space and in coordinate space

### Fourier space:

1. Find Fourier transform of the data and of the templates.
2. Determine the standard deviation of the Fourier modes (given by the Poisson noise level).
3. Find the maximum likelihood model by minimizing the chi^2:

$$\chi^2(\alpha) = \sum_k \frac{|f_k^{\text{data}} - f_k^{\text{model}}(\alpha)|^2}{\sigma_k^2}$$

### Coordinate space:

1. Determine the model in terms of the expected numbers of photons in every pixel  $x_p(\alpha)$ .
2. The likelihood is determined as a product over pixels of Poisson probabilities:

$$L(\alpha) = \prod_p \frac{x_p(\alpha)^{n_p}}{n_p!} e^{-x_p(\alpha)}$$

## Comparison of fitting a small gaussian excess on top of the homogeneous distribution

Quantity	Input	Fourier	Coordinate
Fraction	0.05	0.06 ±0.01	0.06 ±0.01
Position	0.40	0.38 ±0.01	0.38 ±0.01
Dispersion	0.10	0.10 ±0.02	0.10 ±0.01

The results are practically the same but in Fourier space I used only 4 Fourier modes instead of 100 pixels in coordinate space fitting

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## Fourier space fitting

1. Linear fitting procedure

2. Chi^2 doesn't depend on  
pixel size (for small k)

3. Stable wrt negative expected  
numbers of photons

## Coordinate space fitting

Non-linear likelihood function

Absolute likelihood depends  
on pixel size

Poisson probability for negative  
expectation is not defined

## DM annihilation via spherical harmonics of gamma-ray

The goal is to find a small amplitude large scale gamma-ray signal due to annihilation of dark matter in the main halo of our Galaxy using the decomposition in spherical harmonics.

The algorithm has the following steps:

1. Mask the Galactic plane and point sources
2. Choose the astrophysics template(s)
3. Decompose the data and the templates in  $a_{lm}$ 's
4. Calculate the variance of  $a_{lm}$ 's
5. The  $\chi^2$  is defined as

$$\chi^2 = \sum_{lm} \frac{(a_{lm}^{\text{data}} - a_{lm}^{\text{model}})^2}{\sigma^2(a_{lm})}$$

where  $a_{lm}^{\text{model}}$  is a linear sum of an isotropic component (mostly extragalactic), Galactic astrophysics, and any extra templates.

$$\int \gamma_{\ell^m}^* \gamma_{\ell^1 w} d$$

$$\int\limits_w \gamma_{e^w}^* \gamma_{e^{1/w}} d\mathcal{L}$$



$$\int_W Y_{e^m} \cdot Y_{e^{1/m}} \cdot d\mathcal{L} \neq \delta_{1,1} \cdot \delta_{m,m}$$



$$\int_w \gamma_{\ell^m}^* \gamma_{\ell^1 w} \cdot d\mathcal{L} \neq \delta_{\ell^1} \delta_{\ell^m}$$



$$\int_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell\ell'} \delta_{mm'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(\theta) f(\theta) d\Omega$$

$$\int\limits_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell\ell'} \delta_{mm'}$$

$$a_{\ell m} = \int\limits_W Y_{\ell m}^*(\theta) f(\theta) d\Omega_\theta$$

$$\Delta = \int\limits_W \| f(\theta) - a_{\ell m} Y_{\ell m}(\theta) \| d\Omega_\theta$$

$$\int_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell\ell'} \delta_{mm'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(\theta) f(\theta) d\Omega_\theta$$

$$\Delta = \int_W \| f(\theta) - a_{\ell m} Y_{\ell m}(\theta) \|$$

$$\int_W Y_{\ell m}^* Y_{\ell' m'} dR \neq \delta_{\ell \ell'} \delta_{m m'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(y) f(y) dR_y$$

$$\Delta = \int_W \left| f(y) - a_{\ell m} Y_{\ell m}(y) \right|^2 dR_y$$

$$\int_W Y_{e_m}^* Y_{e_n} \, d\mu = \delta_{e_m e_n}$$

$$a_{em} = \int_W Y_{e_m}^*(j) f(j) \, d\mu_j$$

$$\Delta = \int_W \|f(j) - a\|_2^2 \, d\mu_j$$

$$\frac{\partial \Delta}{\partial a_{in}} \Rightarrow \int_W$$

$$\int_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell\ell'} \delta_{mm'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(j) f(j) d\Omega_j$$

$$\Delta = \int_W \| f(j) - a_{\ell m} Y_{\ell m}(j) \|^2 d\Omega_j$$

$$\frac{\partial \Delta}{\partial a_{\ell m}^*} \Rightarrow \int_W Y_{\ell m}^* Y_{\ell m} d\Omega = \delta_{\ell\ell'} \delta_{mm'}$$

$$\int_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell \ell'} \delta_{m m'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(j) f(j) d\Omega_j$$

$$\Delta = \int_W \| f(j) - a_{\ell m} Y_{\ell m}(j) \|^2 d\Omega_j$$

$$\frac{\partial \Delta}{\partial a_{\ell m}^*} \Rightarrow \int_W \left\{ Y_{\ell m}^* Y_{\ell m} \right\} d\Omega_j = \delta_{\ell \ell'} \delta_{m m'} =$$

$$\int_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell\ell'} \delta_{mm'}$$

$$a_{\ell m} = \int_W Y_{\ell m}^*(y) f(y) d\Omega_y$$

$$\Delta = \int_W \| f(y) - a_{\ell m} Y_{\ell m}(y) \|^2 d\Omega_y$$

$$\frac{\partial \Delta}{\partial a_{\ell m}^*} \Rightarrow \int_W \left\{ Y_{\ell m}^* Y_{\ell' m'} d\Omega_y \right\} \partial a_{\ell m}^* = a_{\ell m}$$

$$\int\limits_W Y_{\rho m}^* Y_{\rho' n'} d\Omega \neq \delta_{\rho\rho'} \delta_{m'n'}$$

$$a_{\rho m} = \int\limits_W Y_{\rho m}^*(\theta) f(\theta) d\Omega_\theta$$

$$\Delta = \int\limits_W \| f(\theta) - a_{\rho m} Y_{\rho m}(\theta) \|^2 d\Omega_\theta$$

$$\frac{\partial \Delta}{\partial a_{\rho m}} \Rightarrow \int\limits_W \underbrace{Y_{\rho m}^* Y_{\rho m}}_{g_{\rho m}} d\Omega_\theta = a_{\rho m}$$

$$\int \limits_W Y_{\ell m}^* Y_{\ell' m'} d\Omega \neq \delta_{\ell \ell'} \delta_{mm'}$$

$$a_{\ell m} = \int \limits_W Y_{\ell m}^*(\theta) f(\theta) d\Omega_\theta$$

$$\Delta = \int \left| f(\theta) - a_{\ell m} Y_{\ell m}(\theta) \right|^2 d\Omega_\theta$$

$$\frac{\partial \Delta}{\partial a_{\ell m}} \Rightarrow \underbrace{\int \limits_W Y_{\ell m}^* Y_{\ell m} d\Omega_\theta}_{g_{\ell m}}$$

$$\neq \delta_{\ell\ell'} \delta_{mm'}$$

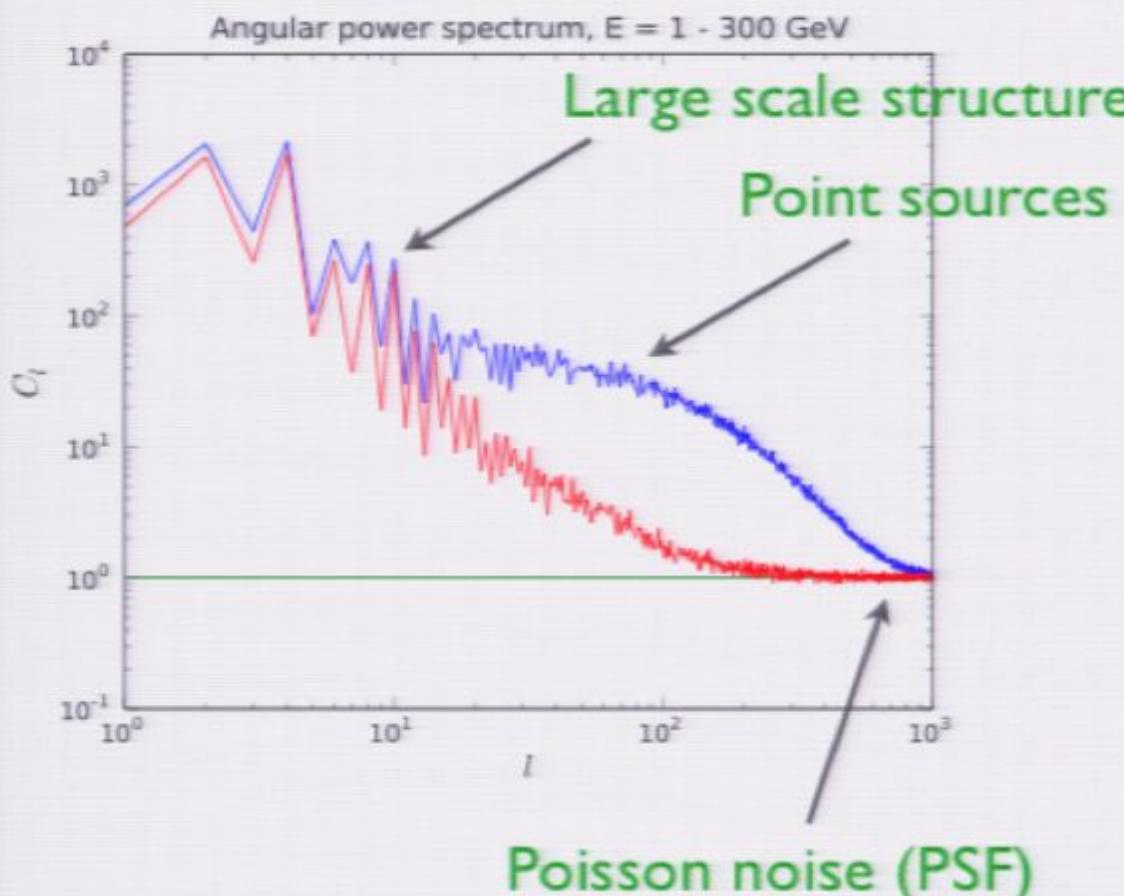
$$f(\gamma) = \sum_{\ell m} a^{\ell m} Y_{\ell m}(\gamma)$$

$$f(\gamma) d\gamma$$

$$|a^{\ell m} Y_{\ell m}(\gamma)|^2 d\gamma$$

$$a^{\ell m} \delta_{\ell\ell'} \delta_{mm'} = a_{\ell m}$$

# Angular power spectrum of Fermi data above 1 GeV



Blue curve - all Fermi data at  $|b| > 5 \text{ deg}$

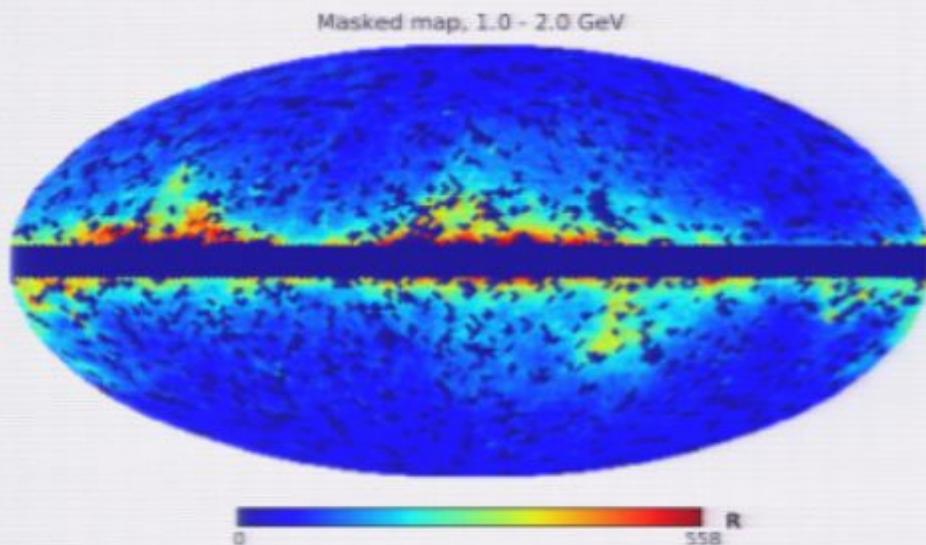
Red curve - gamma-ray point sources are masked

Point sources increase the level of Poisson noise.

In order to study the large scale distributions for  $\ell \gtrsim 10$  one needs to mask the point sources.

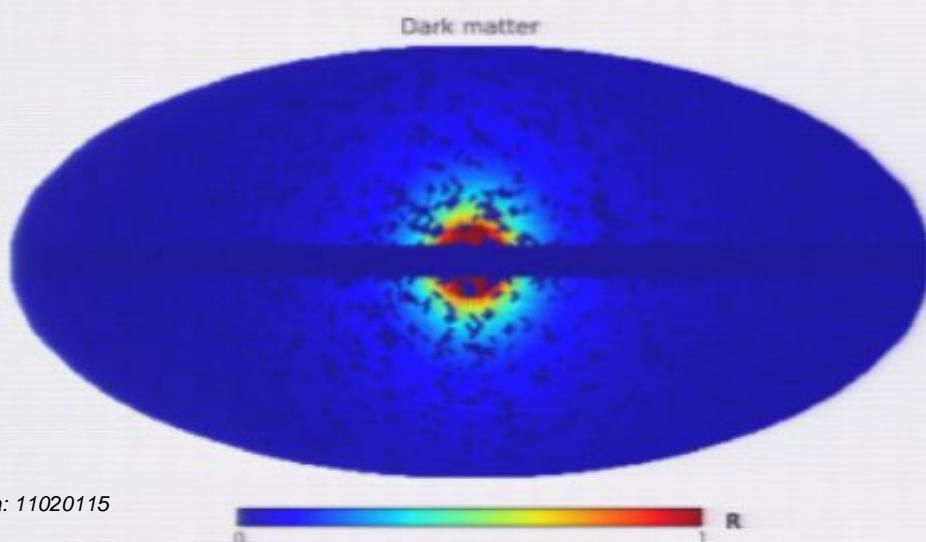
In the analysis we will use  
 $\ell_{\max} = 15$

## Templates



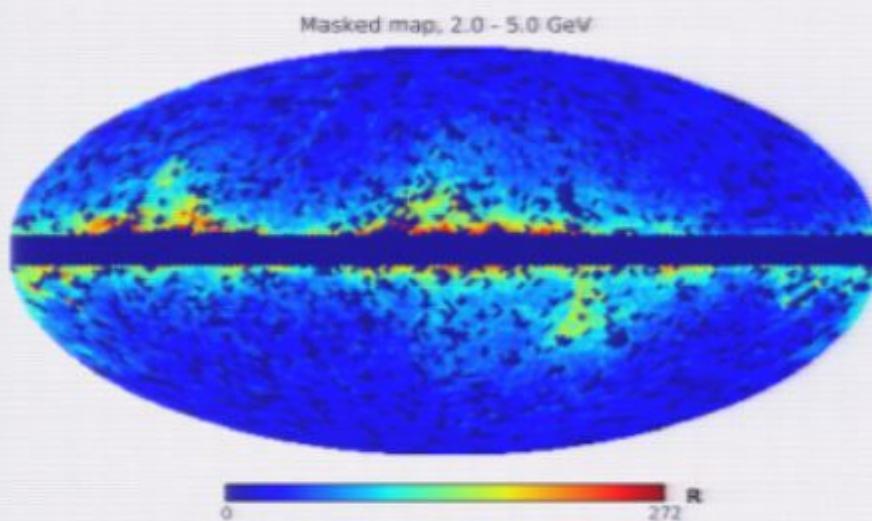
Masked Fermi data in 1-2 GeV  
energy bin.

Photon counts in pixels  
with pixel size about 2 deg



Dark matter annihilation in NF  
profile (arbitrary units)

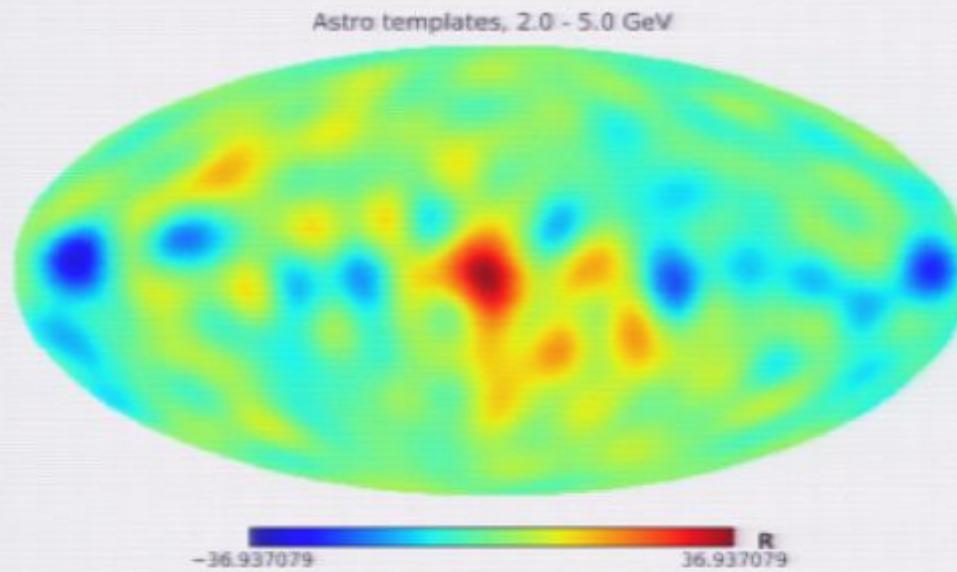
## Data



Consider the Fermi data in 2-5 GeV energy bin and make the spherical harmonics decomposition

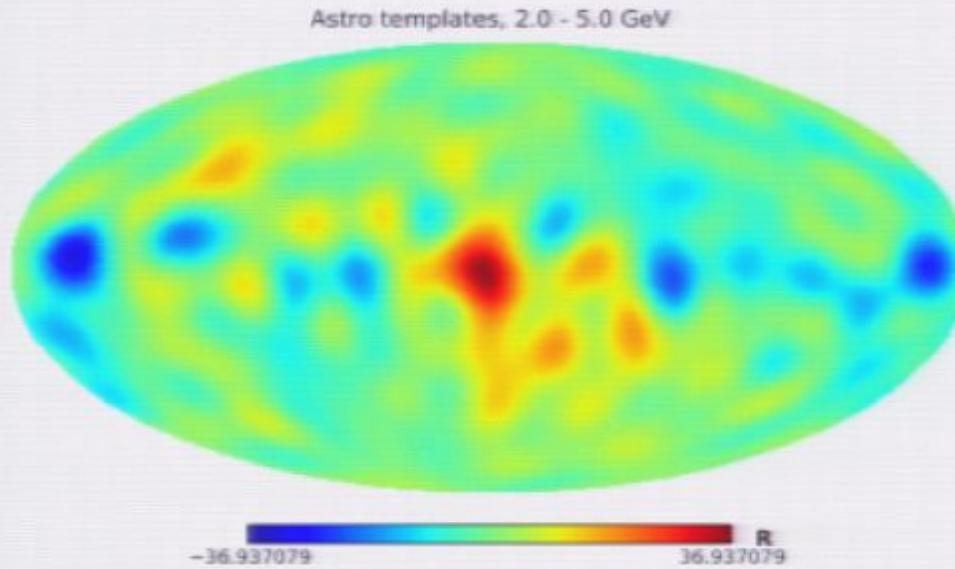
Inverse transform of the first 15 spherical harmonics

## 2 - 5 GeV astro templates

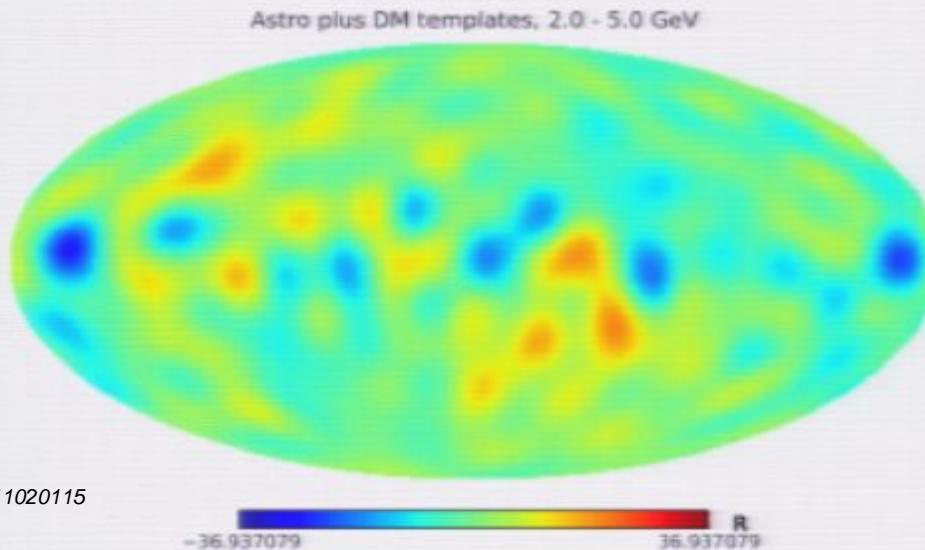


Inverse transform of the residual for low E bin  
and isotropic templates

## 2 - 5 GeV astro+DM templates

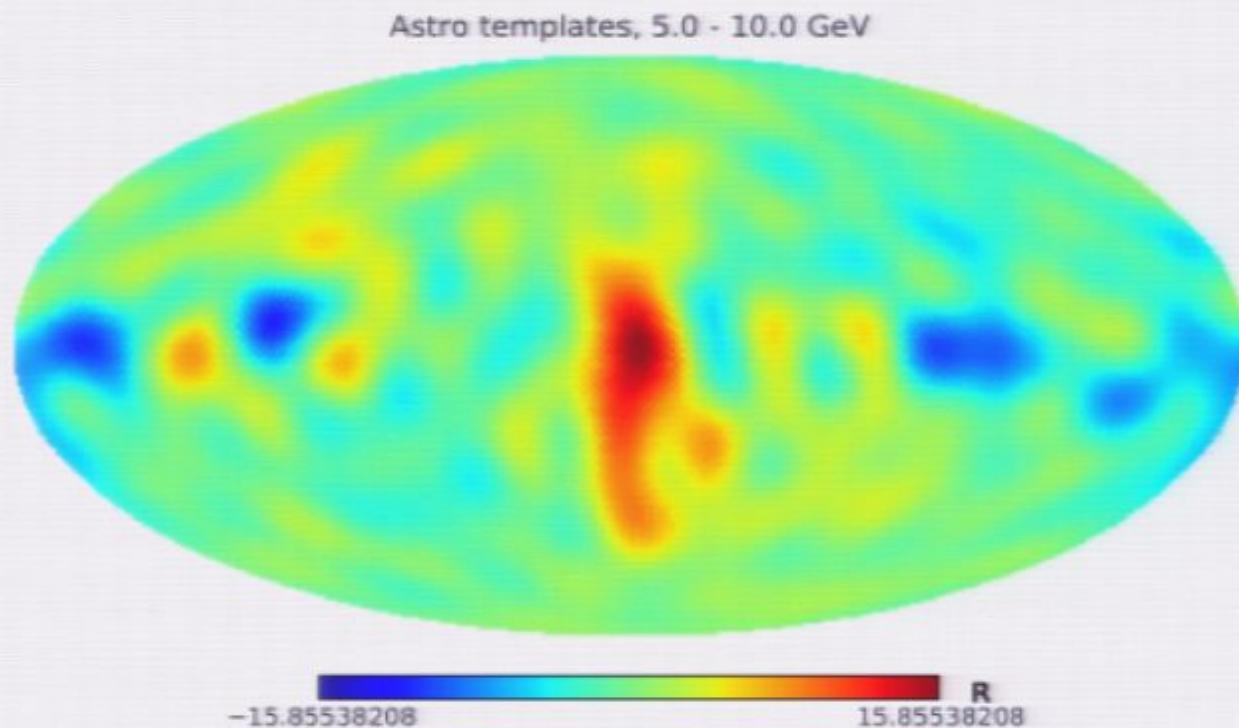


Inverse transform of the residual for low E bin and isotropic templates

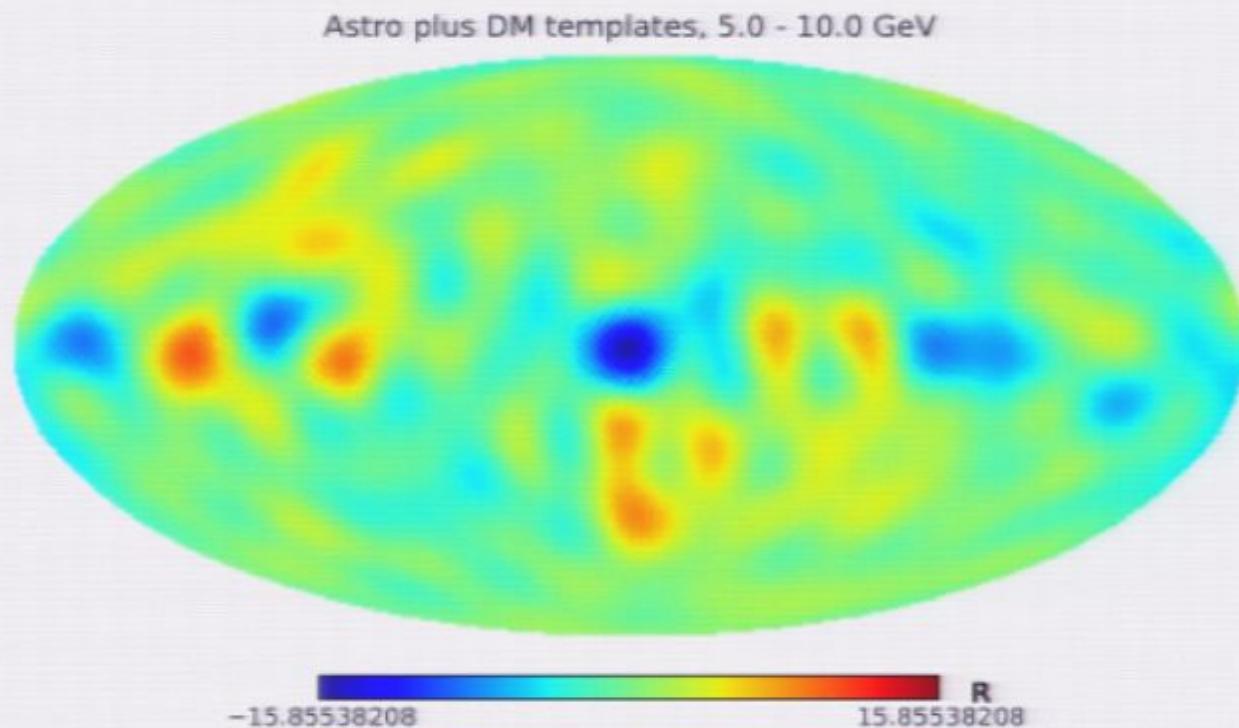


Inverse transform with DM template added

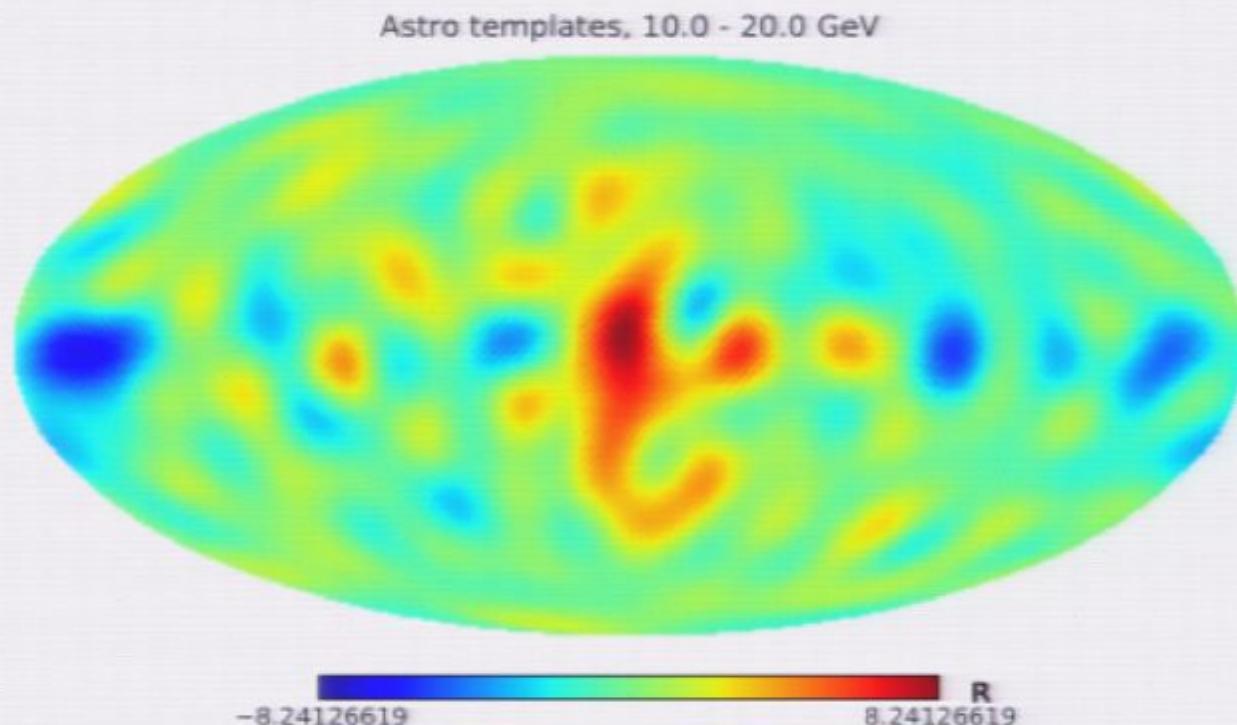
## 5 - 10 GeV astro templates



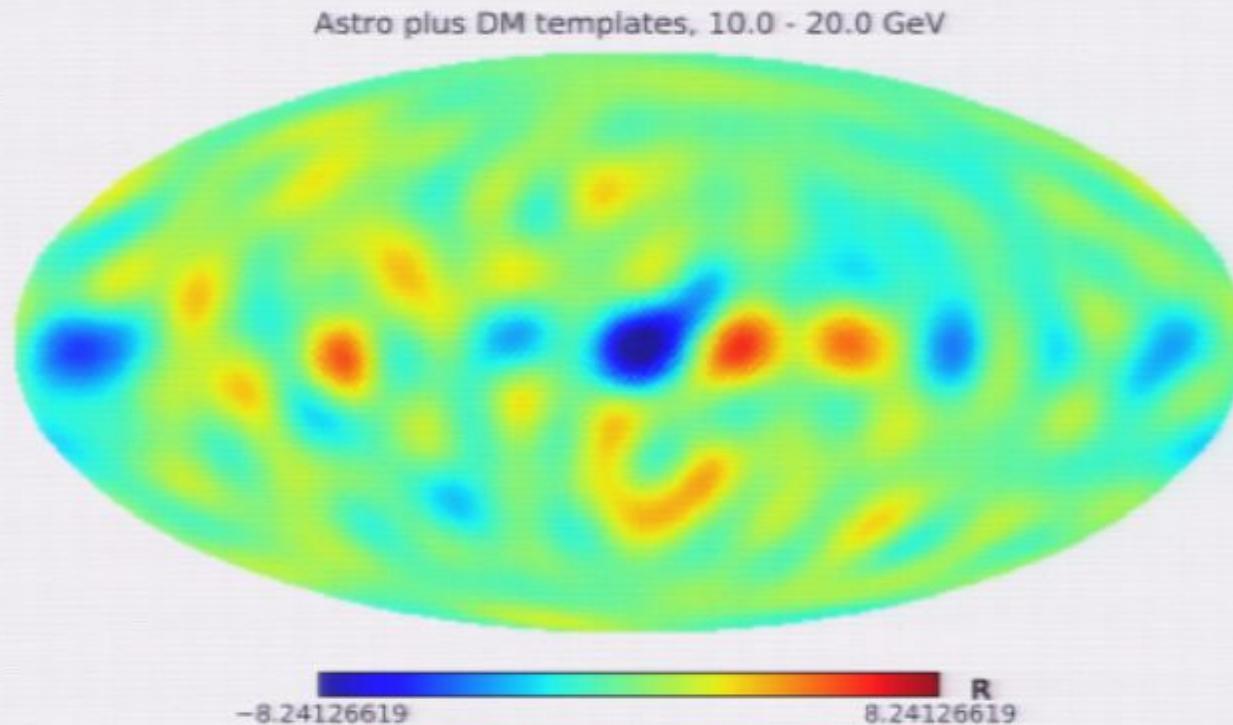
## 5 - 10 GeV astro+DM templates



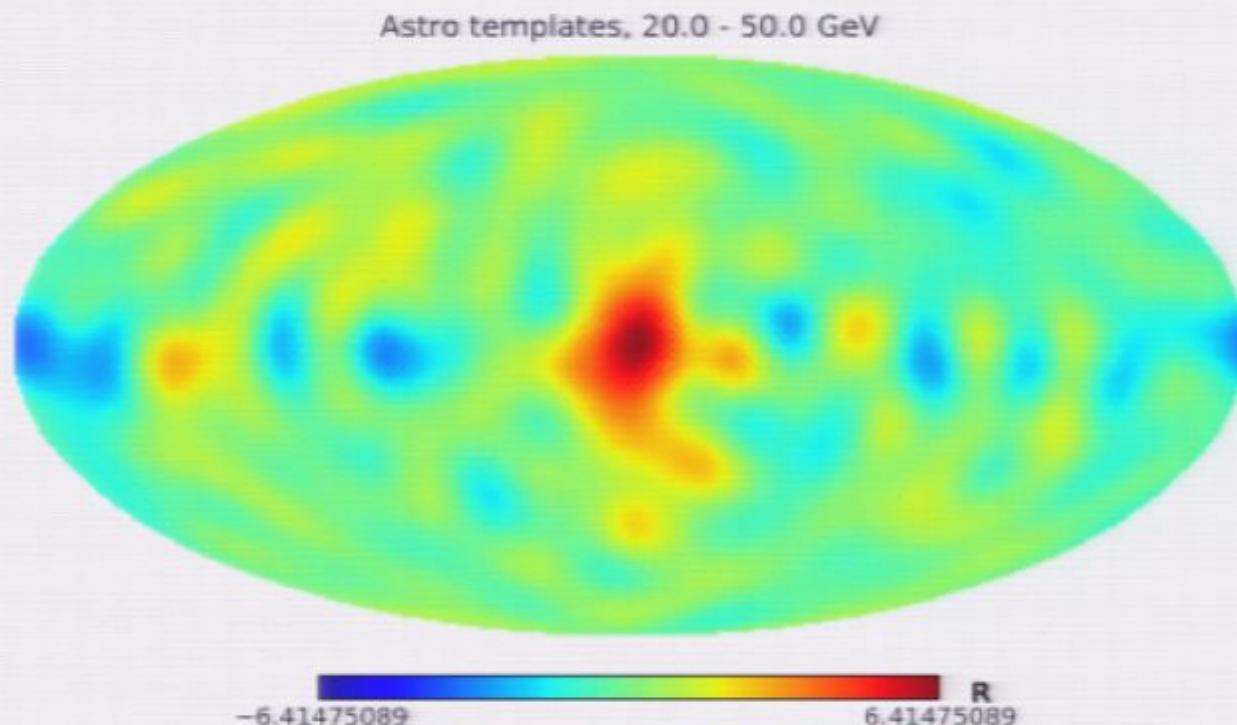
## 10 - 20 GeV astro templates



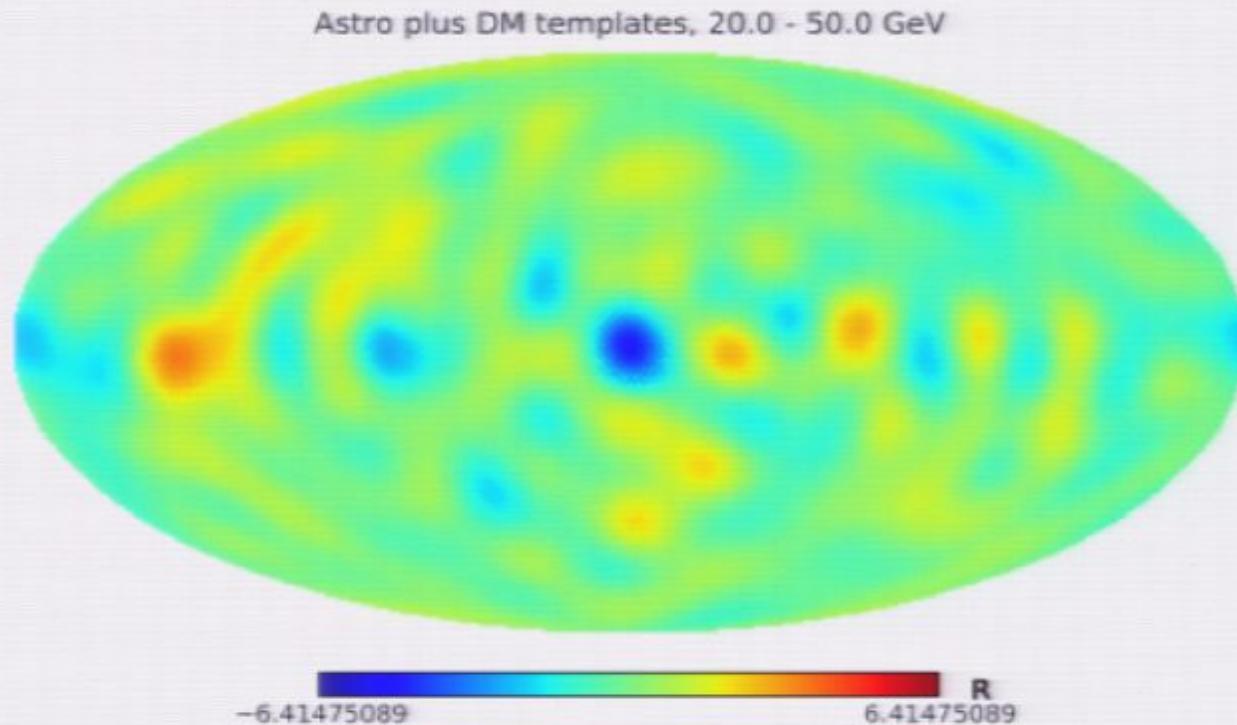
## 10 - 20 GeV astro+DM templates



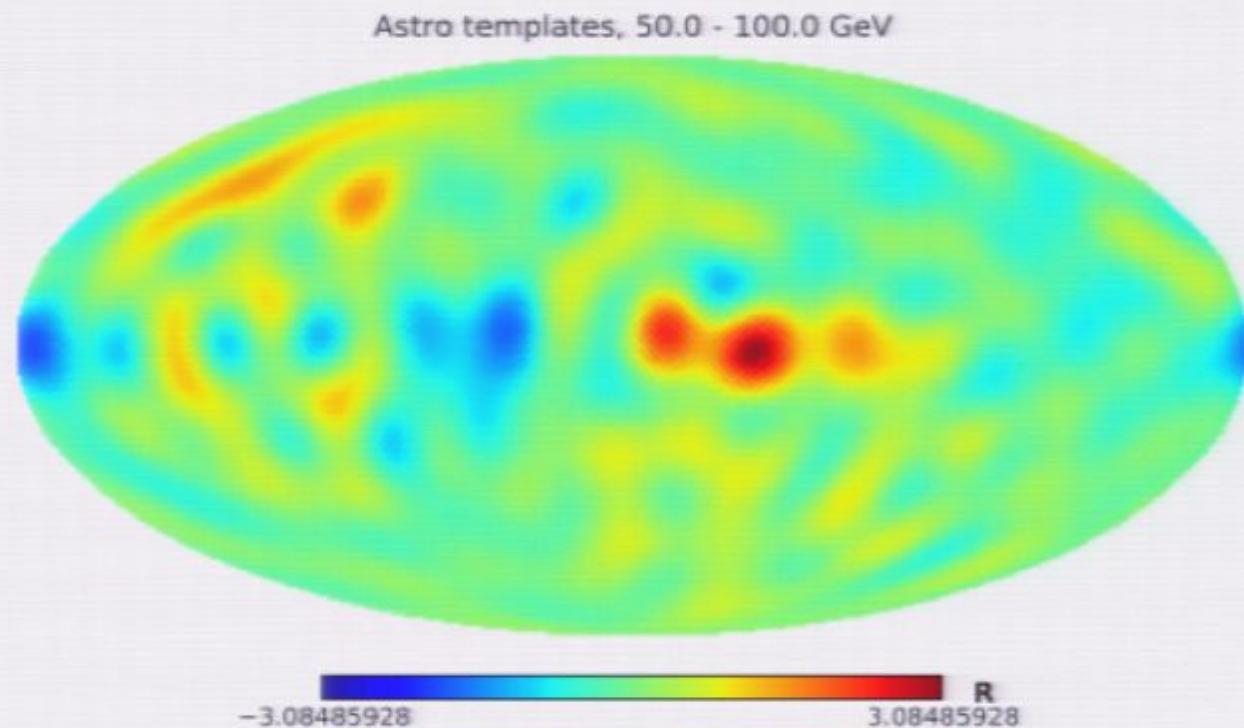
## 20 - 50 GeV astro templates



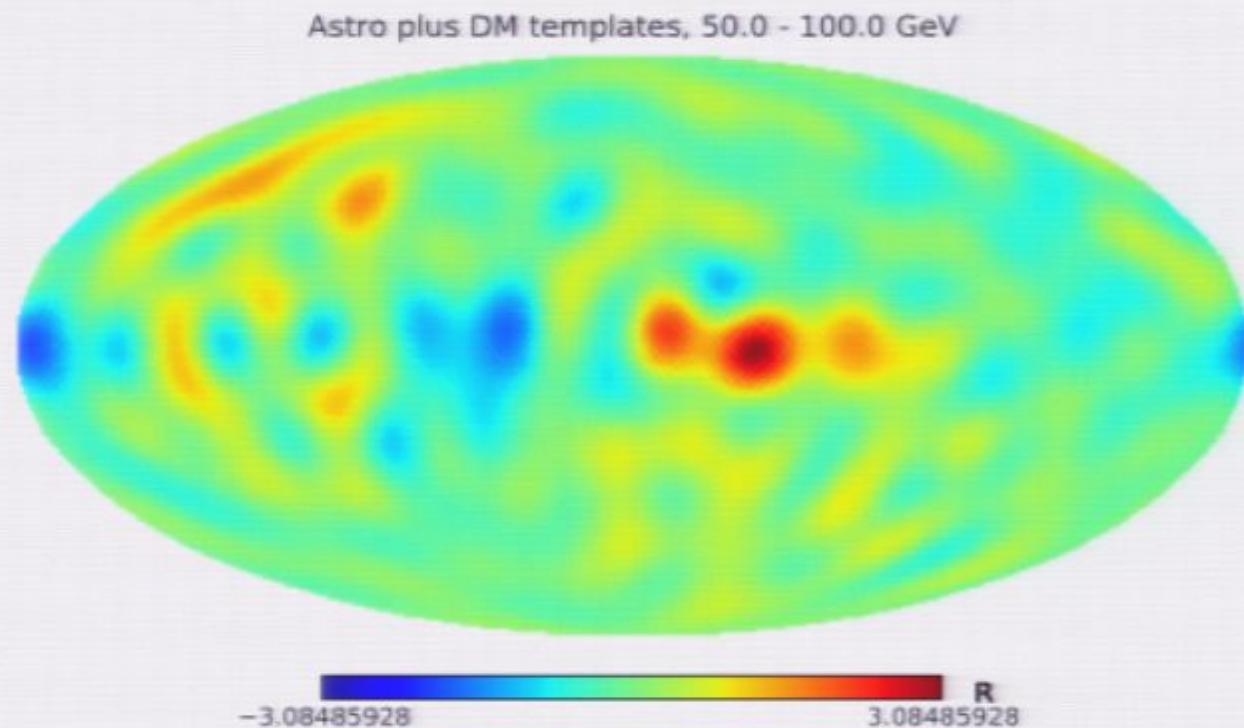
## 20 - 50 GeV astro+DM templates



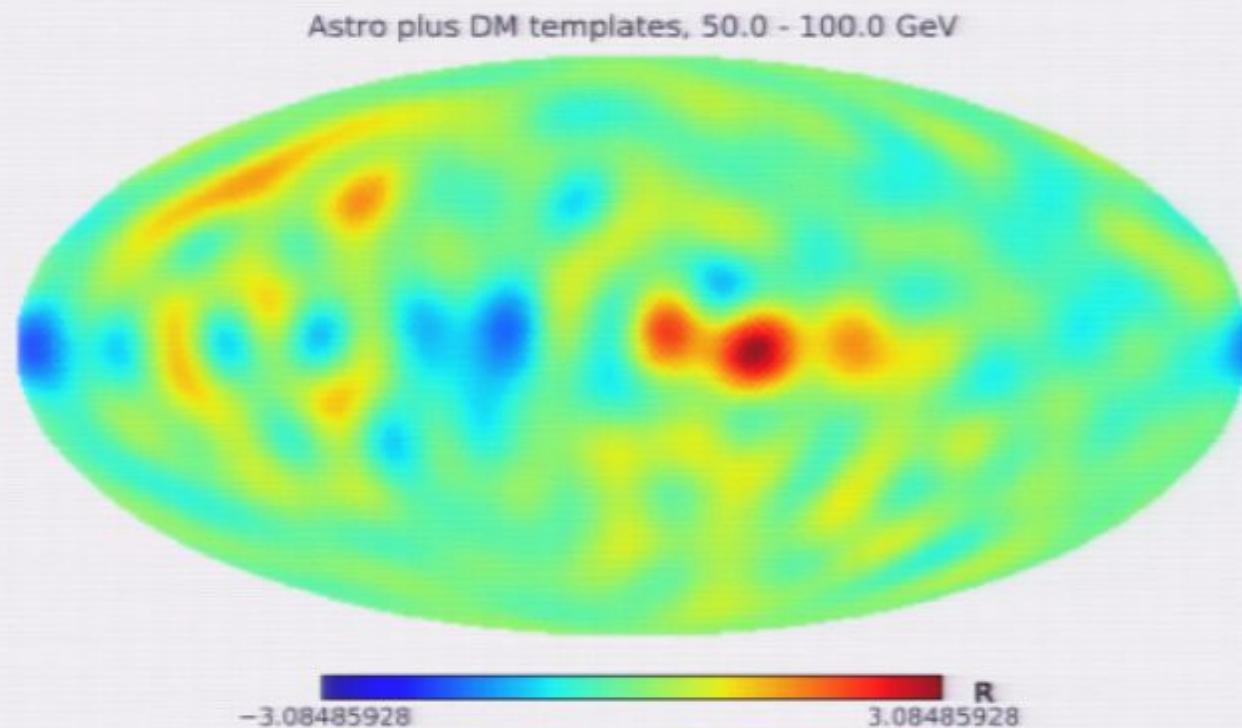
## 50 - 100 GeV astro templates



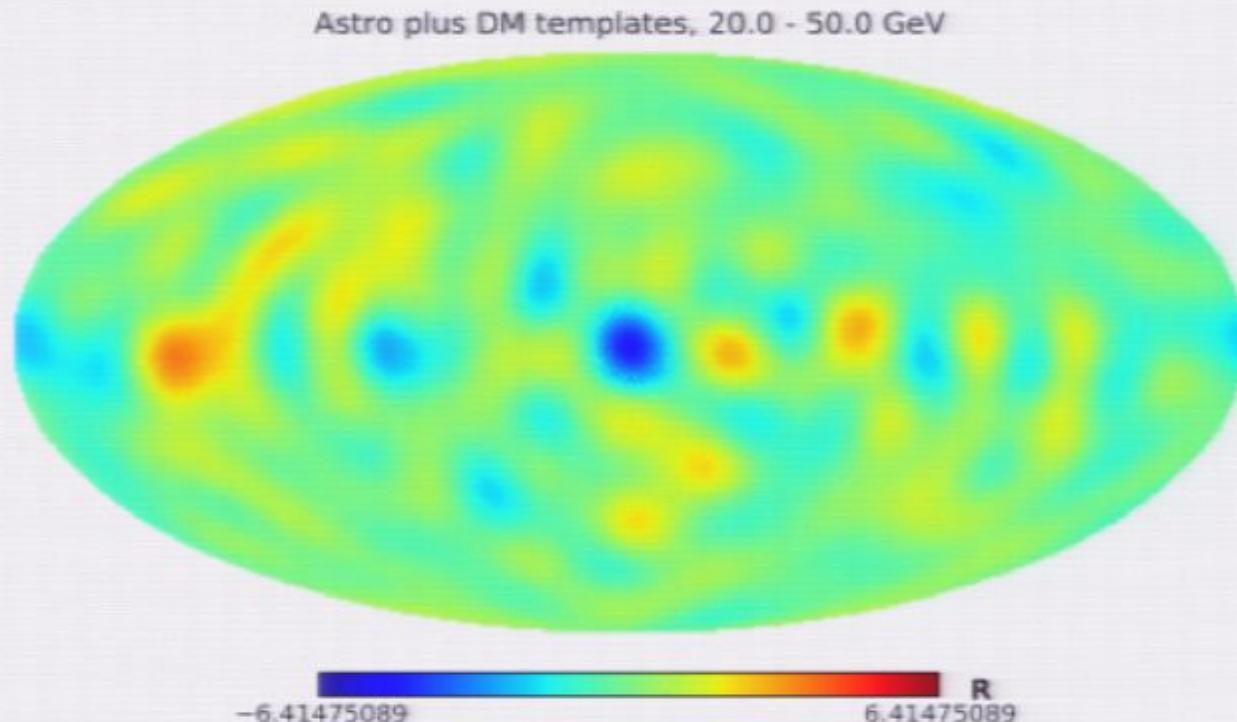
## 50 - 100 GeV astro+DM templates



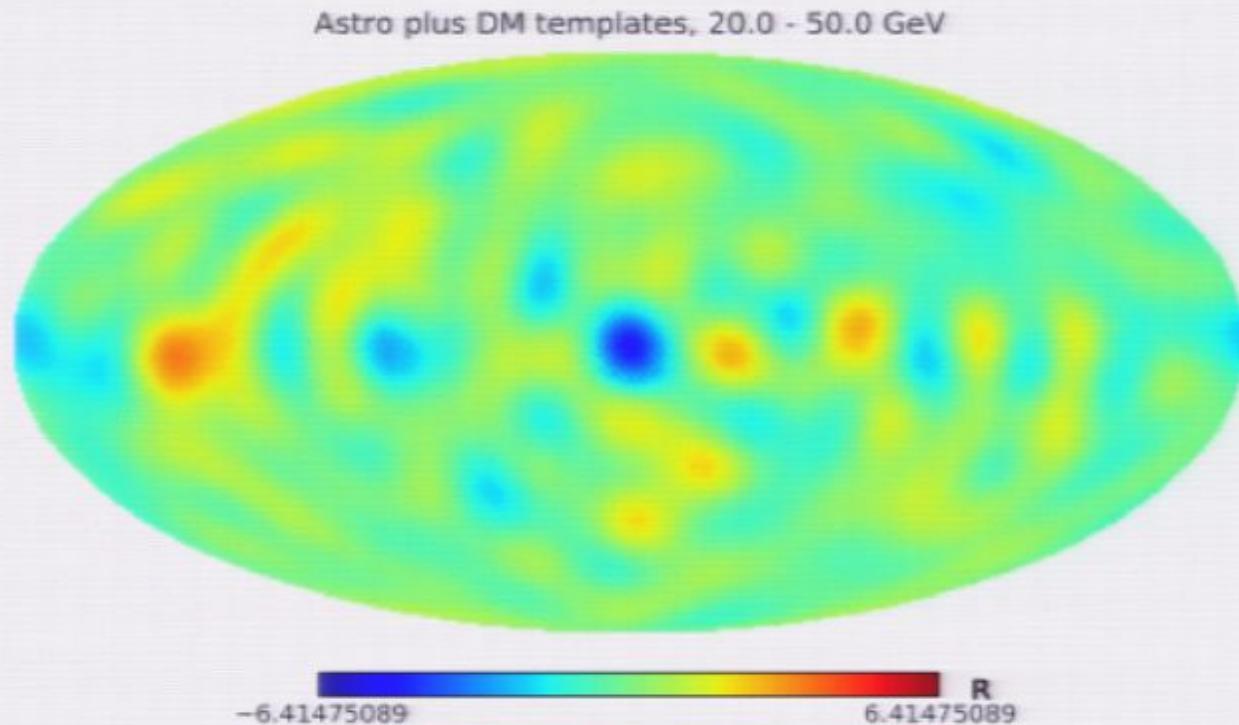
## 50 - 100 GeV astro+DM templates



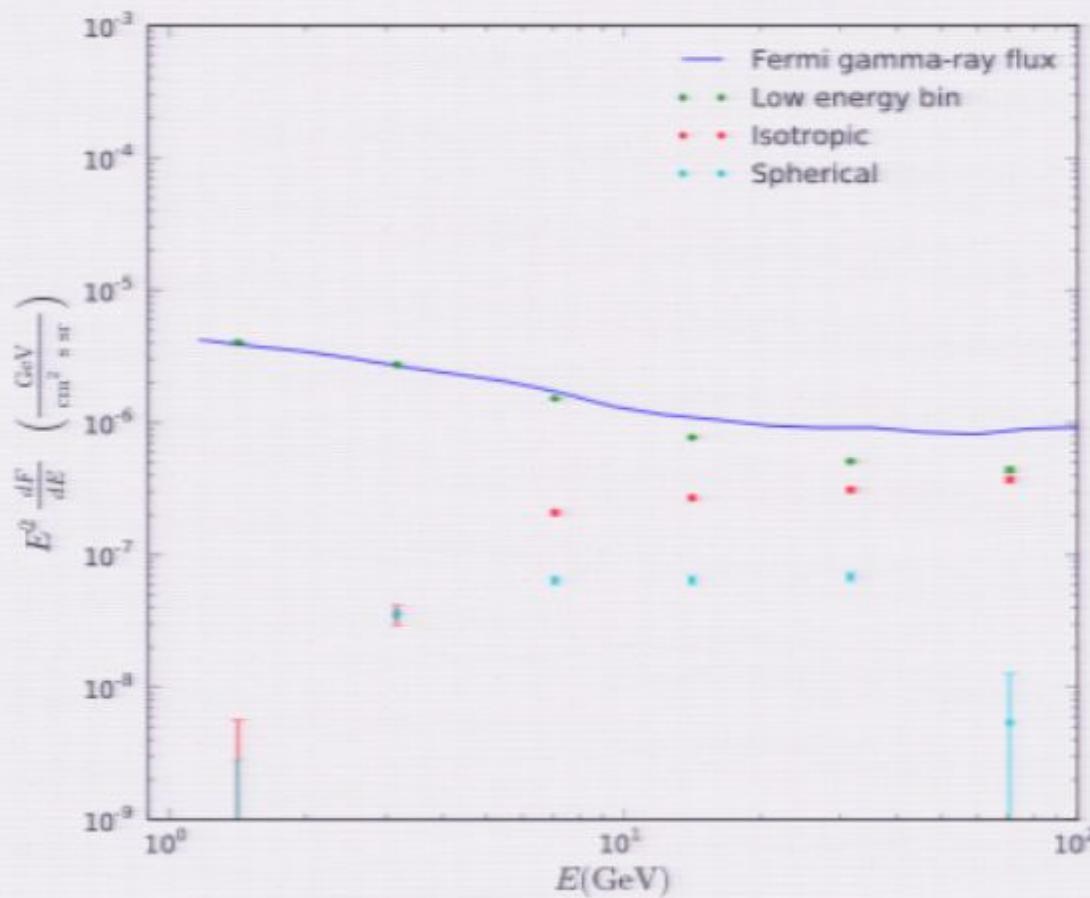
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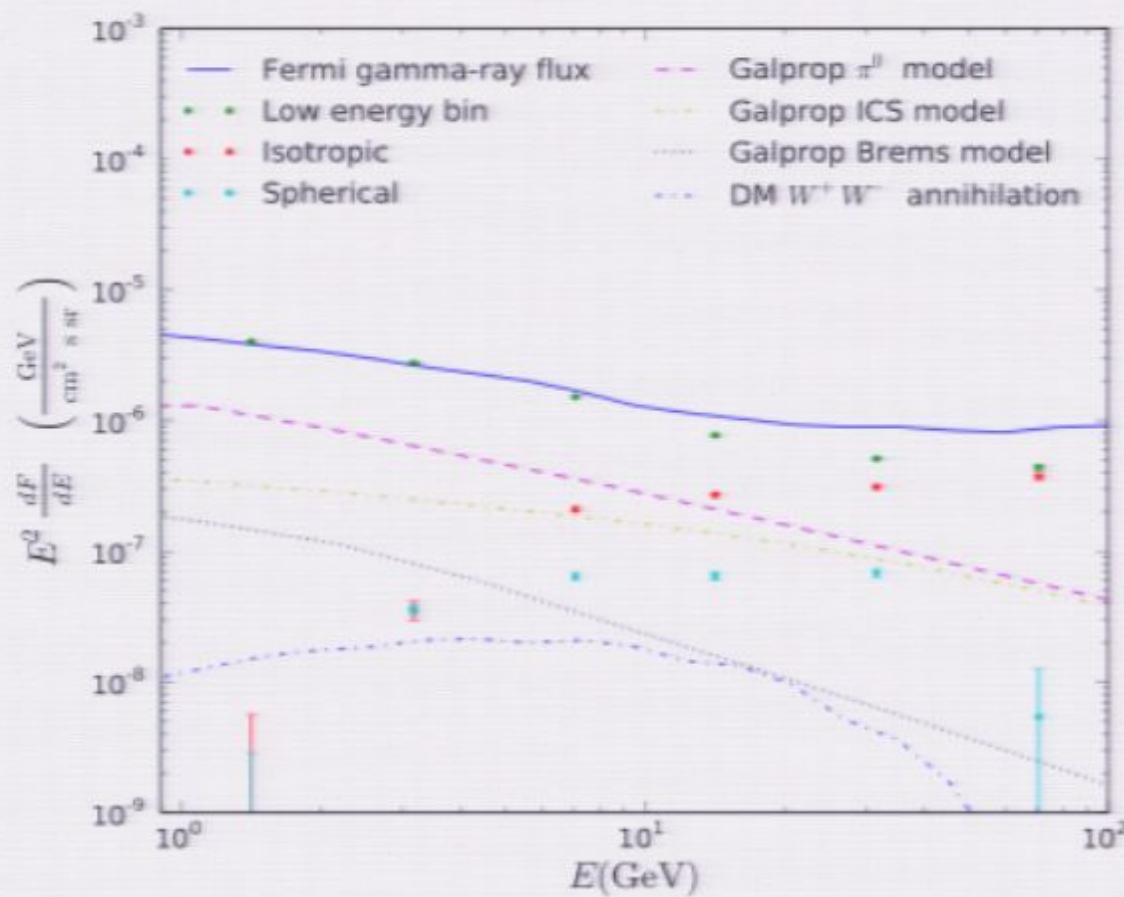
## 20 - 50 GeV astro+DM templates



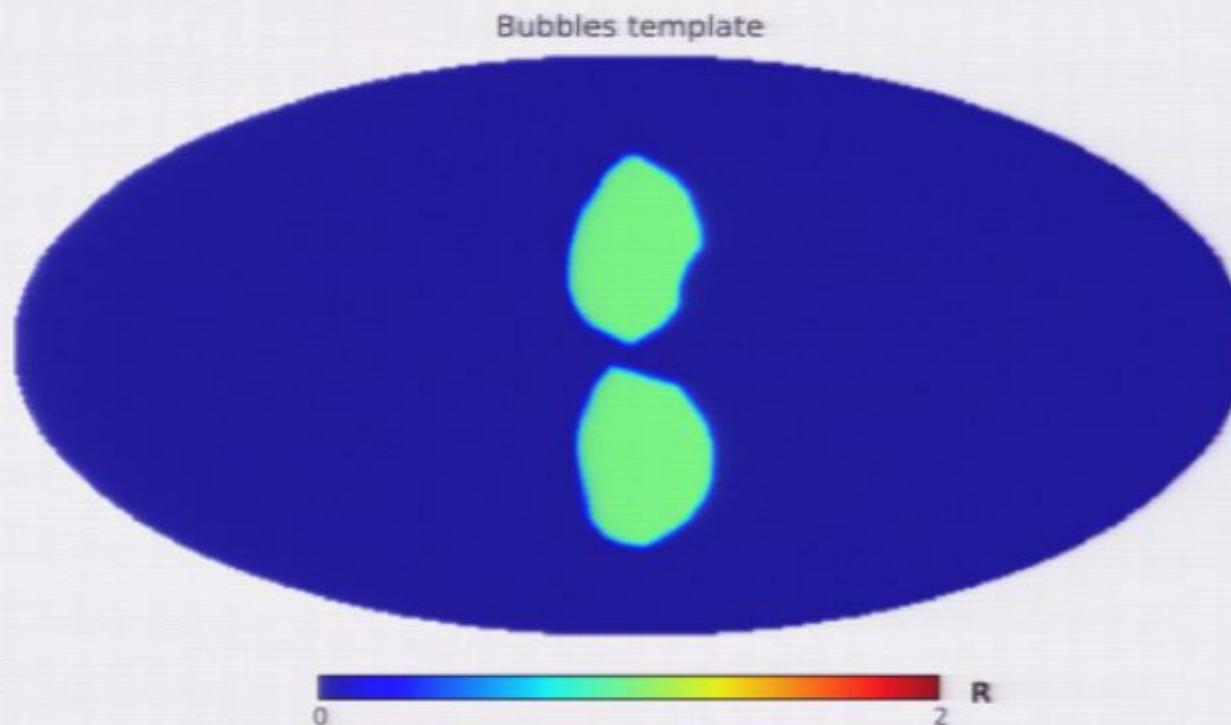
## Fluxes for astro and DM templates



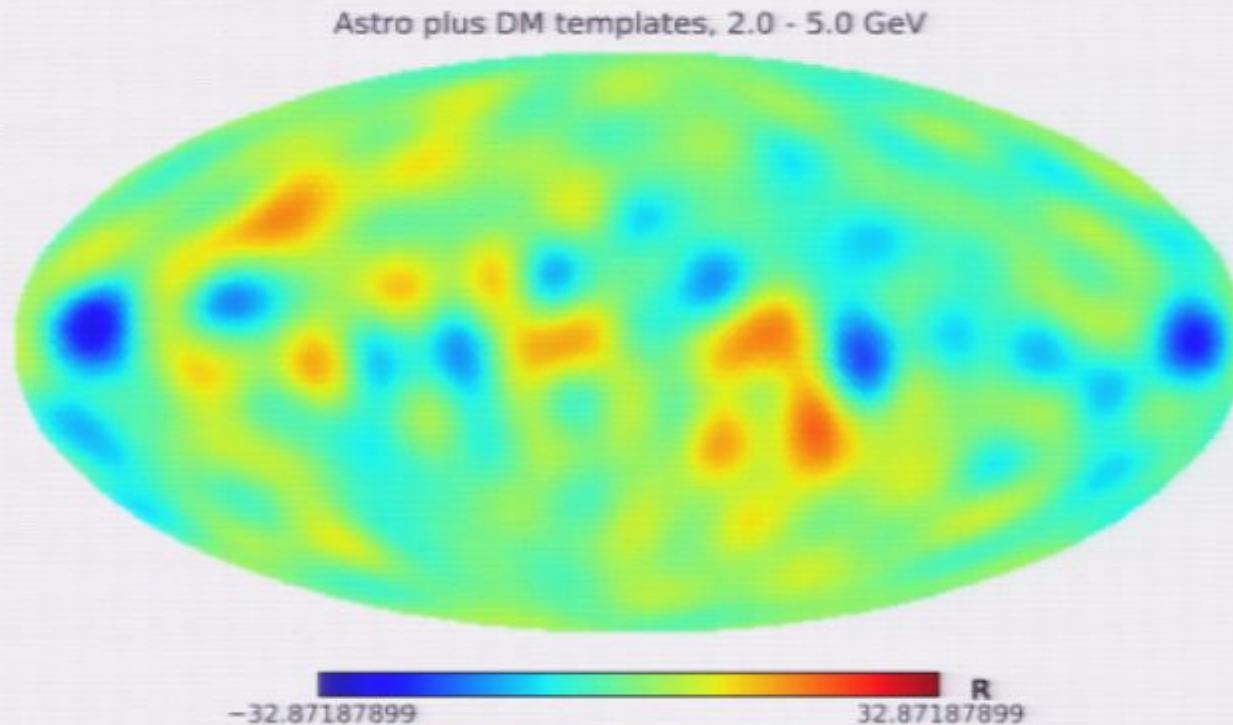
## Comparison with Galprop and DM annihilation



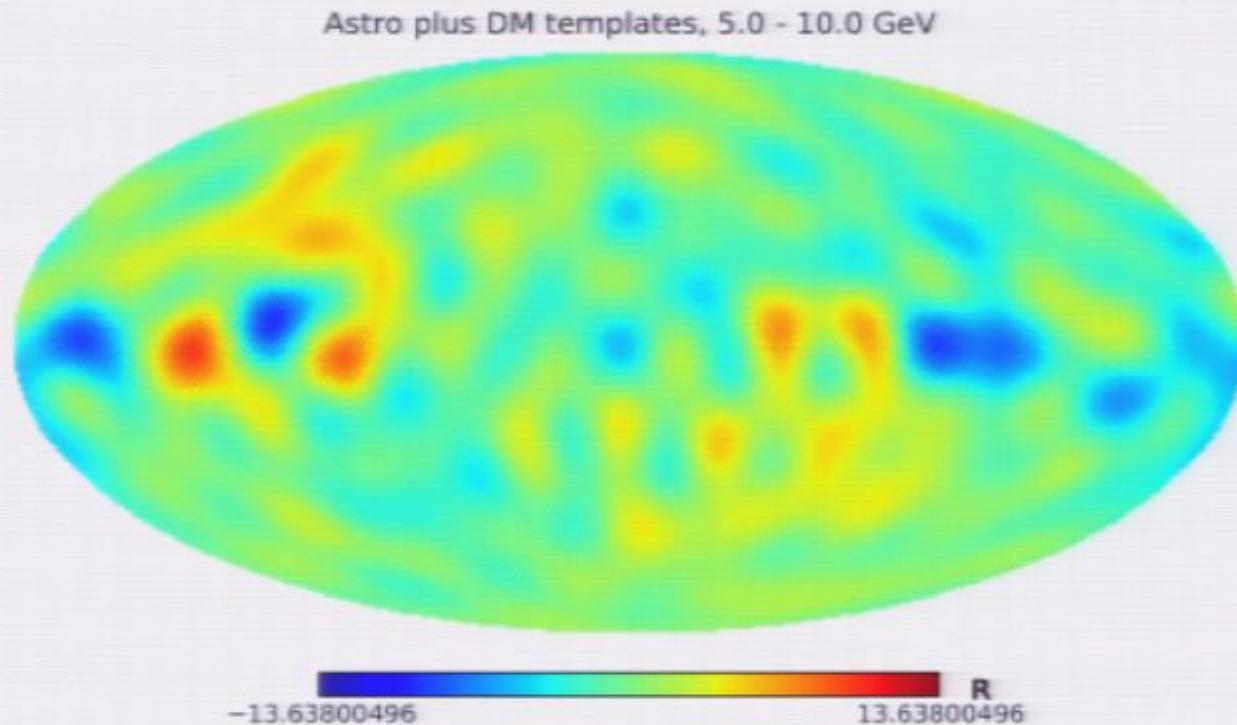
## Add bubbles template (Su, Slatyer, Finkbeiner 2010)



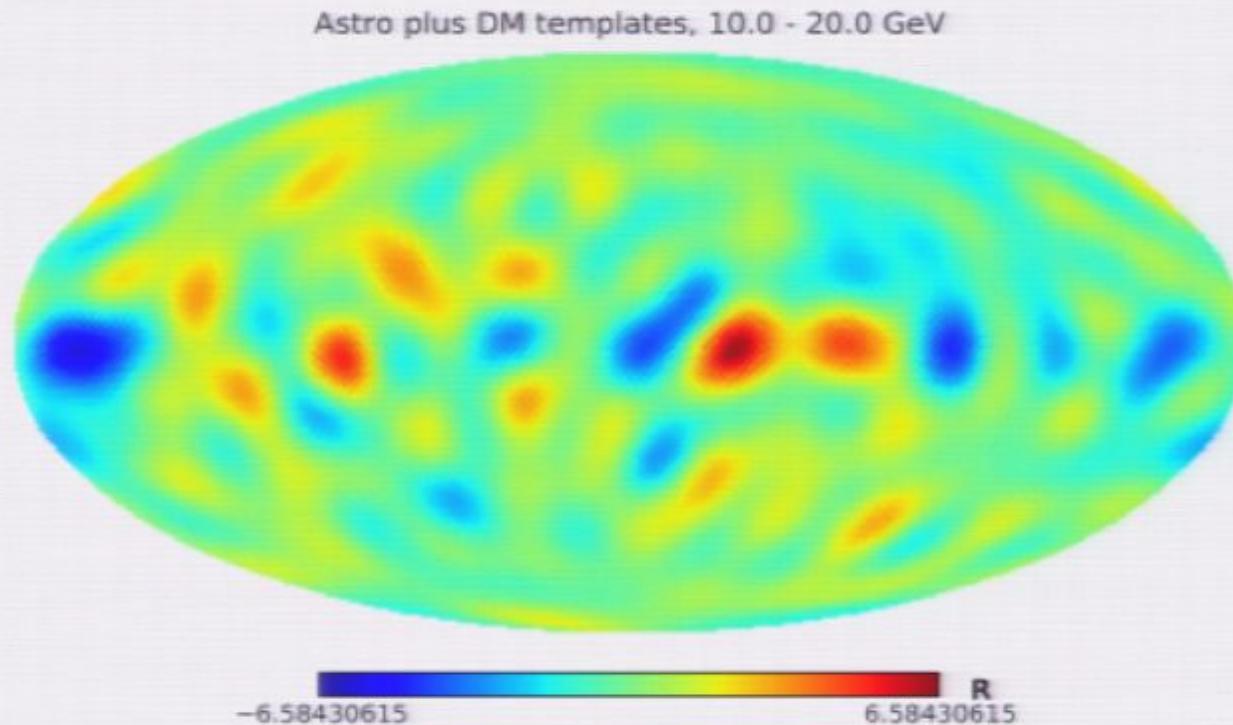
## 2 - 5 GeV astro+bubbles+DM templates



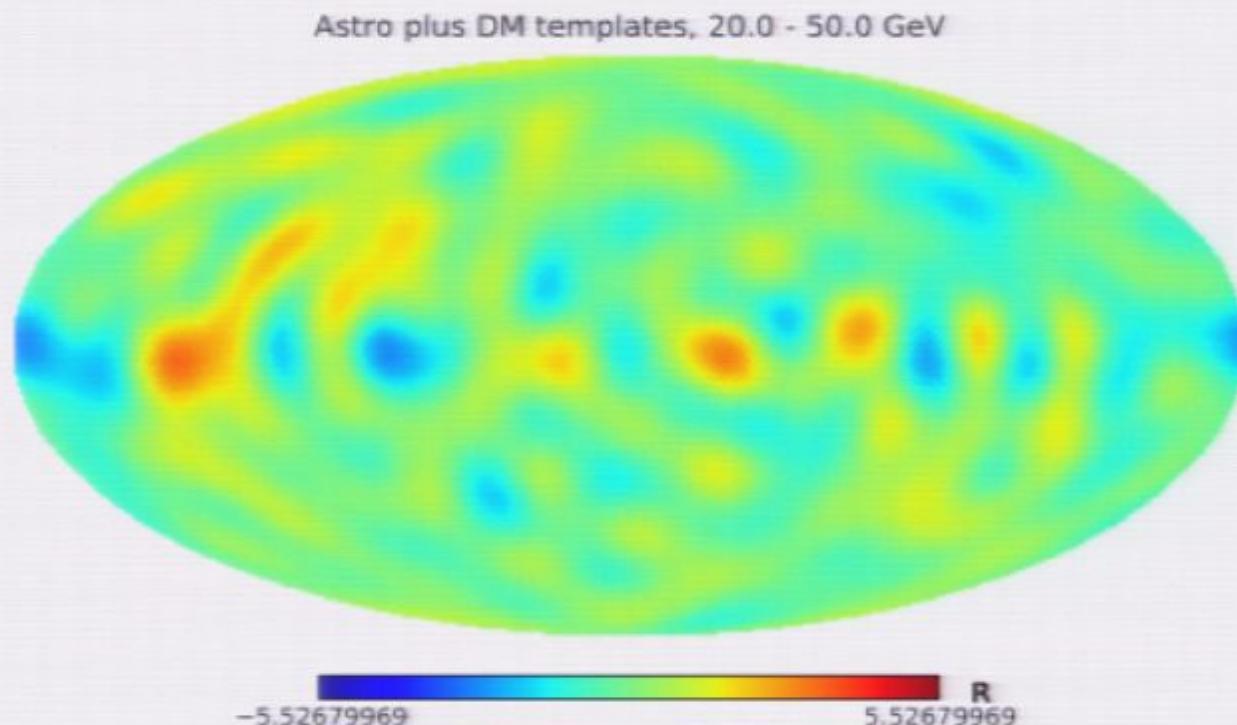
## 5 - 10 GeV astro+bubbles+DM templates



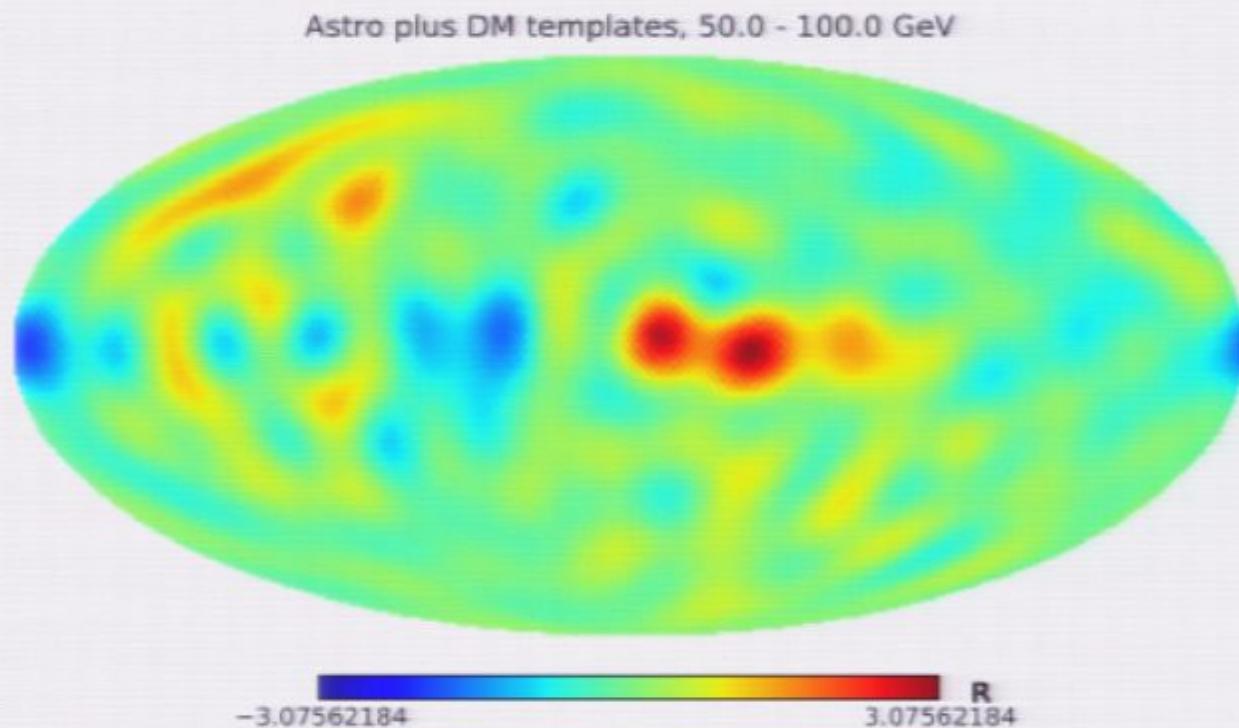
## 10 - 20 GeV astro+bubbles+DM templates



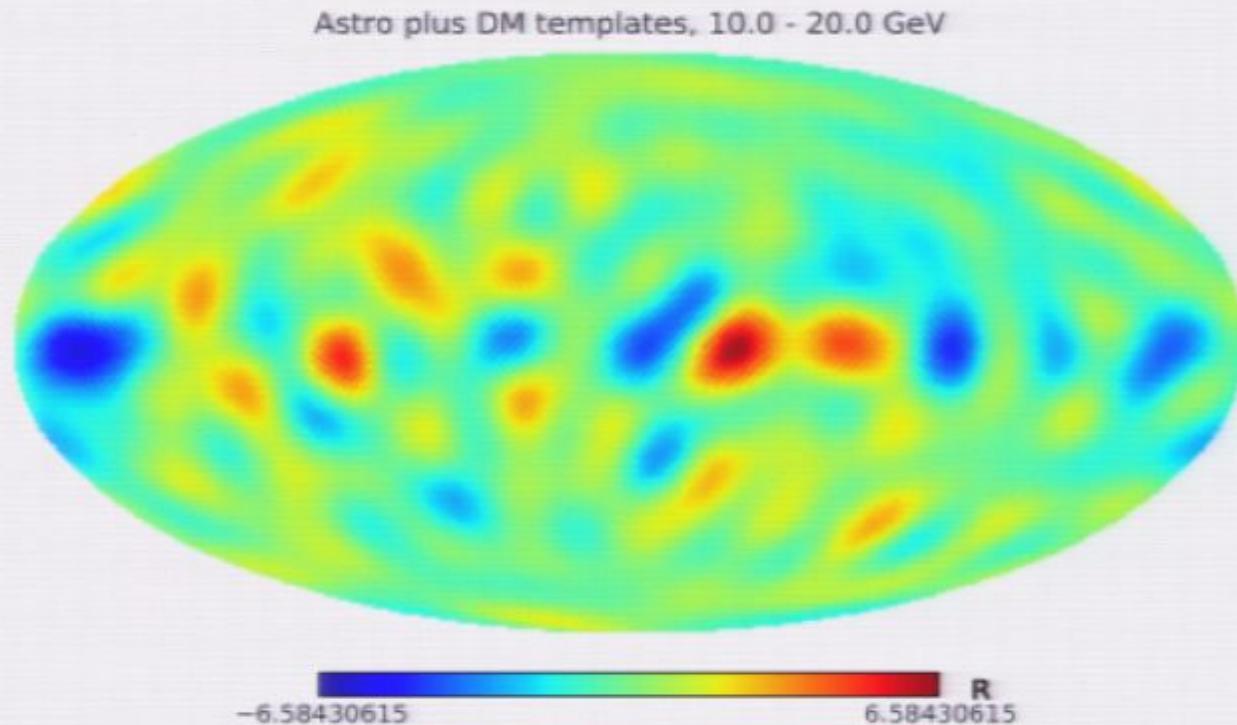
## 20 - 50 GeV astro+bubbles+DM templates



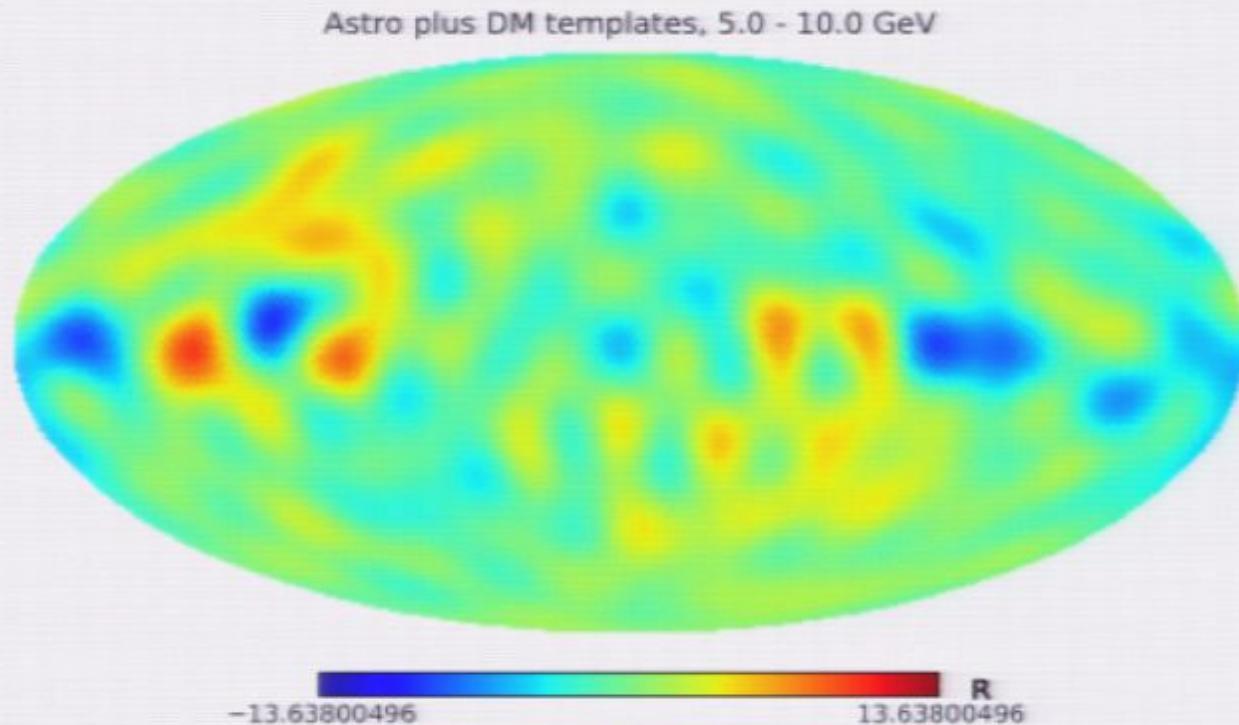
## 50 - 100 GeV astro+bubbles+DM templates



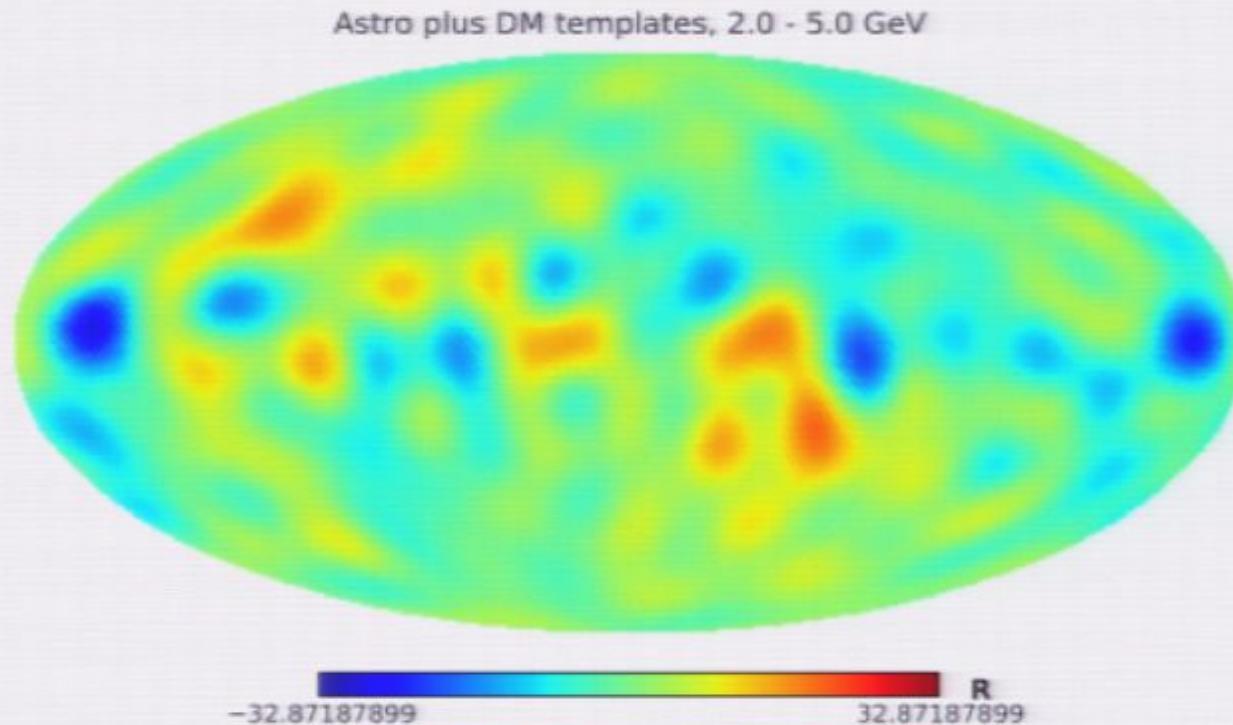
## 10 - 20 GeV astro+bubbles+DM templates



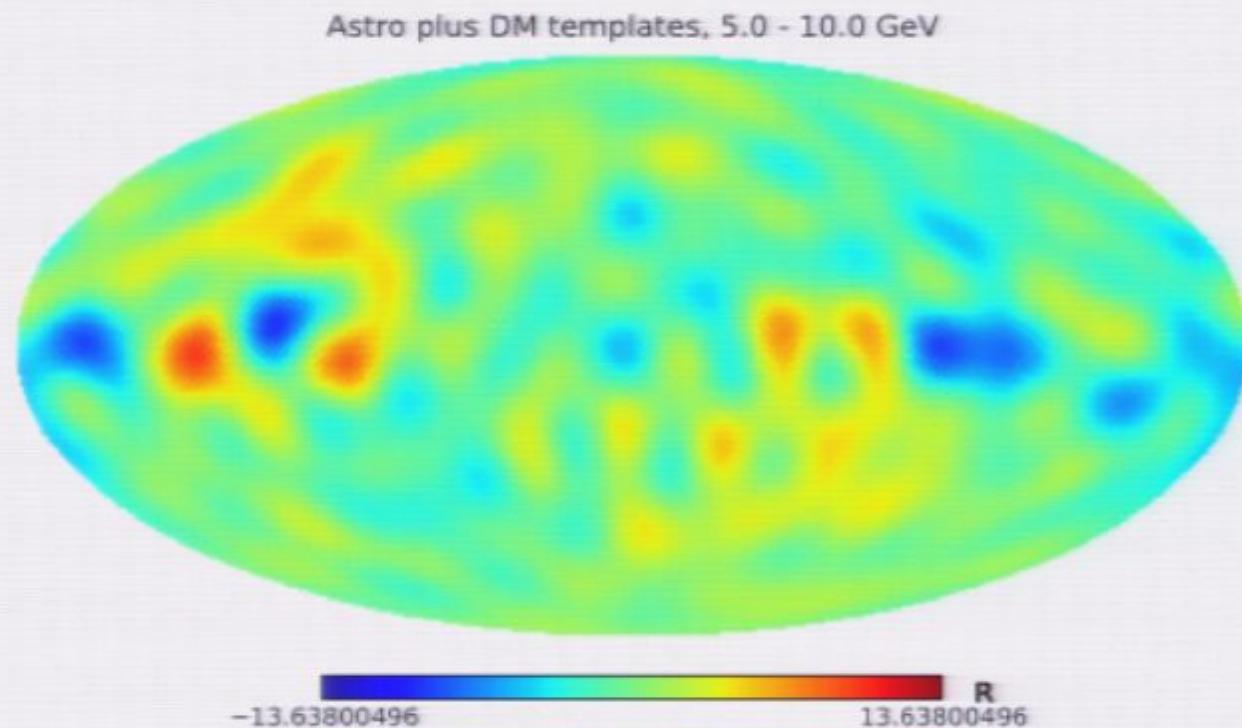
## 5 - 10 GeV astro+bubbles+DM templates



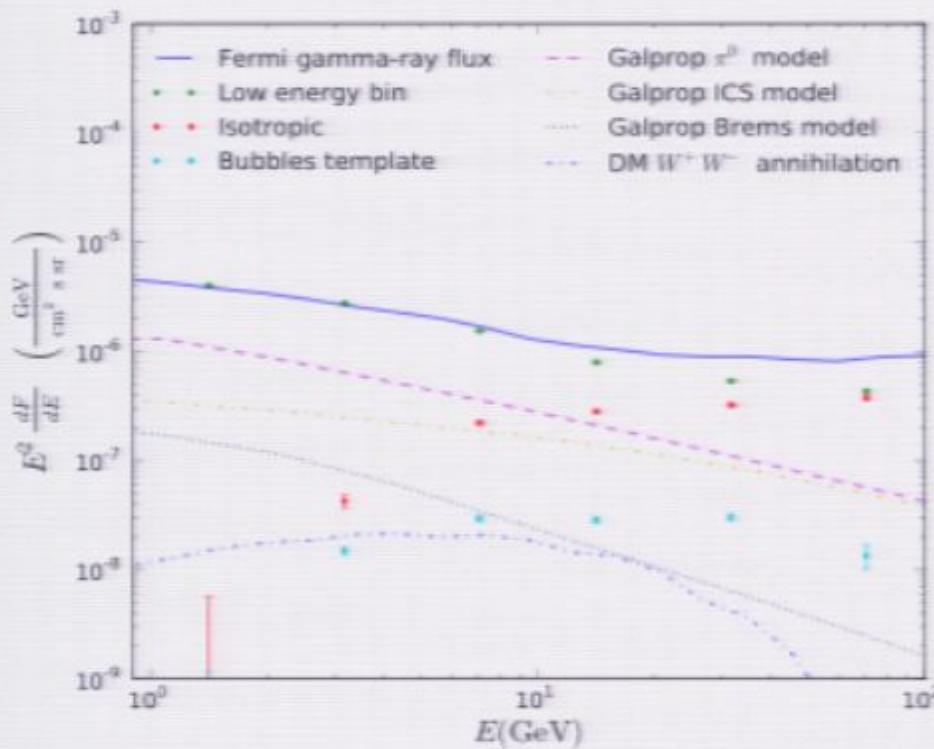
## 2 - 5 GeV astro+bubbles+DM templates



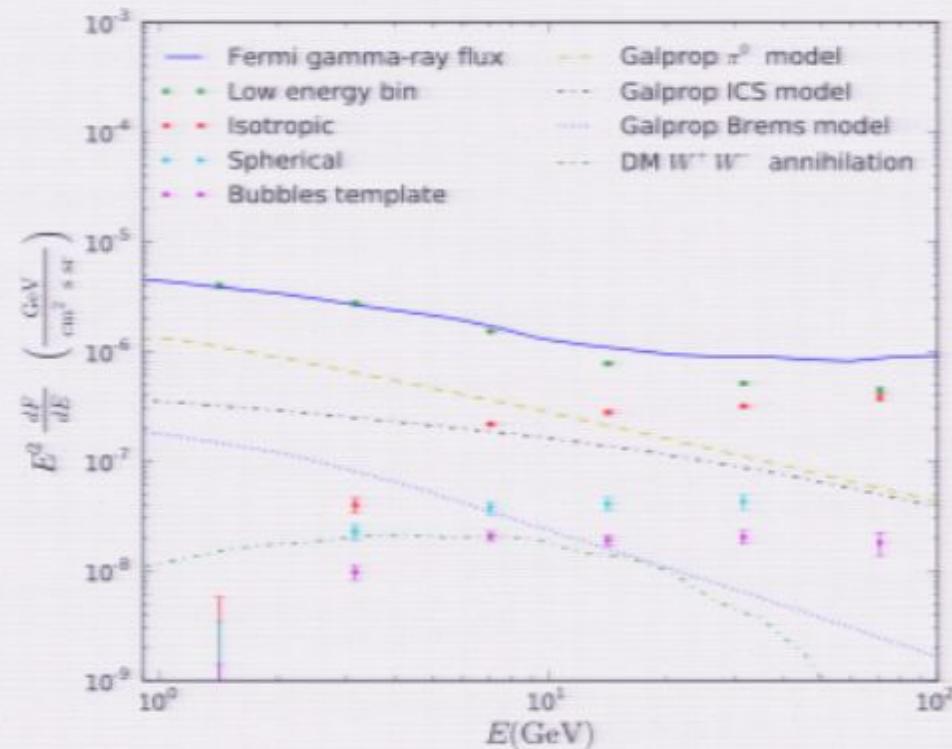
## 5 - 10 GeV astro+bubbles+DM templates



## Astro + bubbles



## Astro + DM + bubbles



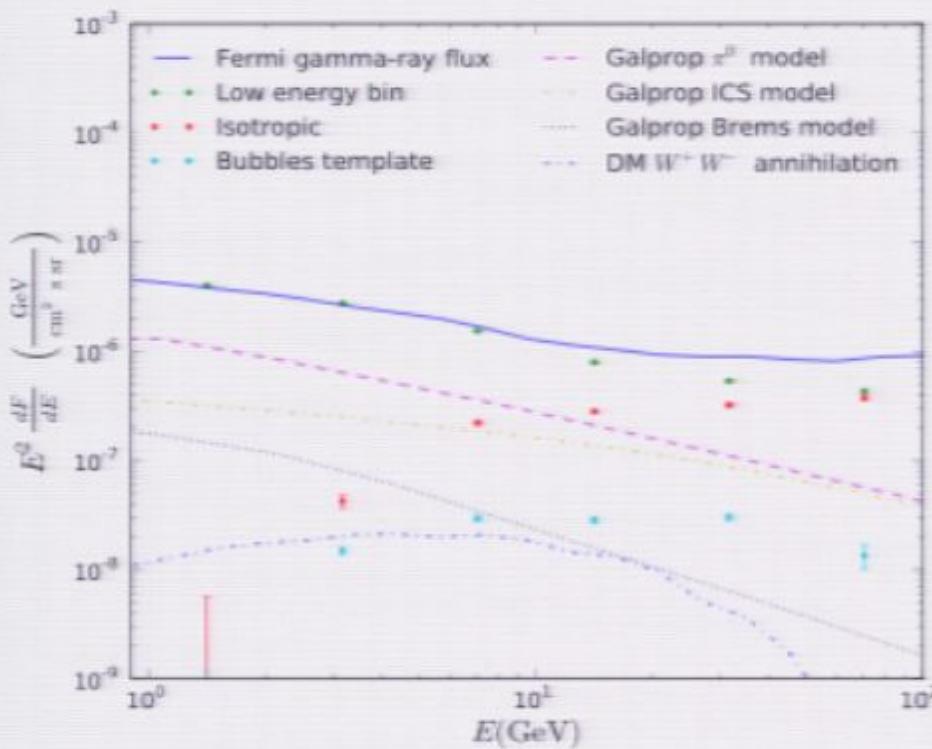
Adding the bubble template reduces the power in the spherical template.

The spherical template contribution remains significant and a little closer to the expected DM annihilation signal

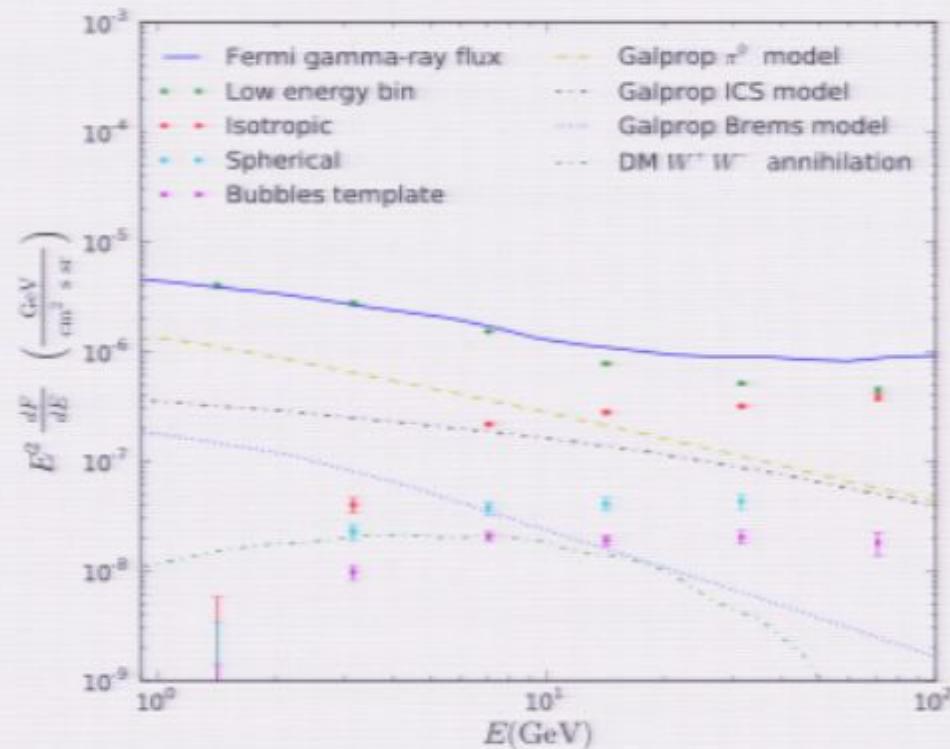
## Conclusions

1. Spherical harmonics are a natural choice for fitting the large scale structures in gamma-rays. It is a linear procedure, that avoids many difficulties associated with coordinate space fitting (e.g., upgrading / degrading of templates to fit to the same pixelation etc.)
2. The calculation of  $\chi^2$  and the fitting procedure are simple.  
The data is conveniently separated into low  $\ell$  modes (relevant for the fitting) and high  $\ell$  modes (dominated by the noise).
3. We observe that adding an extra spherical template significantly reduces the  $\chi^2$ . Adding the bubbles template reduces  $\chi^2$  but doesn't eliminate the DM-like spherical contribution.

## Astro + bubbles



## Astro + DM + bubbles



Adding the bubble template reduces the power in the spherical template.

The spherical template contribution remains significant and a little closer to the expected DM annihilation signal