

Title: Condensed Matter Review - Lecture 13

Date: Jan 19, 2011 10:15 AM

URL: <http://pirsa.org/11010095>

Abstract:

$$\Delta(T=0) = 2\omega_D e^{-\frac{1}{N(0)V}}$$

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For $T > 0$ $\Delta(T) \neq 0$

$$1 = N(0)V \int_0^{\omega_D} \frac{d\epsilon}{\epsilon^2 + \Delta^2} \tanh \frac{\beta(\epsilon^2 + \Delta^2)}{2}$$

$$\Delta(T=0) = 2\omega_D e^{-\frac{1}{N(0)V}}$$

For $T > 0$ $\Delta(T) \neq 0$

$$1 = N(0)V \int_0^{\omega_D} \frac{d\epsilon}{\sqrt{\epsilon^2 + \Delta^2}} \tanh \frac{\beta(\sqrt{\epsilon^2 + \Delta^2}}{2}$$

$$\Delta(T=0) = 2.4 \omega_D e^{-\frac{1}{N(0)V}}$$

For $T > 0$ $\Delta(T) \neq 0$

$$1 = N(0)V \int_0^{\omega_D} \frac{d\epsilon}{\sqrt{\epsilon^2 + \Delta^2}} \tanh\left(\frac{\beta\sqrt{\epsilon^2 + \Delta^2}}{2}\right) \rightarrow \Delta(T)$$

As $T \rightarrow T_c$

$$\Delta(T=0) = 2.4 \omega_p e^{-\frac{1}{N(0)V}}$$

For $T > 0$ $\Delta(T) \neq 0$

$$1 = N(0)V \int_0^{\omega_n} \frac{d\epsilon}{\sqrt{\epsilon^2 + \Delta^2}} \tanh\left(\frac{\beta\sqrt{\epsilon^2 + \Delta^2}}{2}\right) \rightarrow \Delta(T)$$

As $T \rightarrow T_c$ $\Delta(T) \rightarrow 0$

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$$\Delta(T) \rightarrow 0$$

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As $T \rightarrow T_c$ $\Delta(T) \rightarrow 0$

$$\frac{1}{N(0)V} = \int_0^{\omega_D} \tanh\left(\frac{\epsilon}{2T_c}\right) \frac{d\epsilon}{\epsilon}$$

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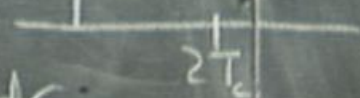
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Integral
 $\frac{1}{2T_c}$



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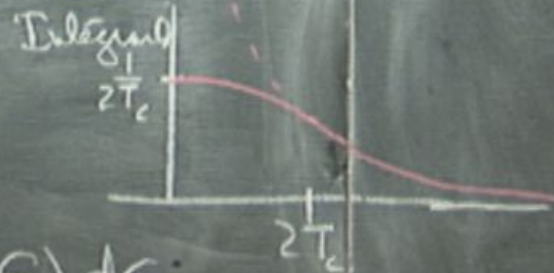
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As $\Delta(T) \rightarrow 0$

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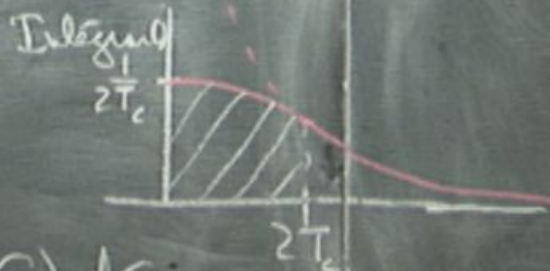
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$$\frac{1}{N(0)V} = \int_0^{\omega_D} \tanh\left(\frac{\epsilon}{2T_c}\right) \frac{d\epsilon}{\epsilon}$$

$$\approx \ln \frac{\omega_D}{2T_c} + 0.82$$

$$T_c = 1.13 \omega_D e^{-\frac{1}{N(0)V}}$$



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Weak Coupling, BCS Result

$\Delta(T)$



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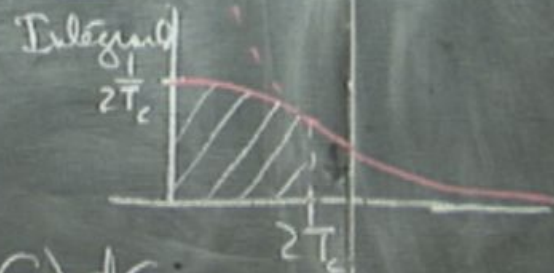
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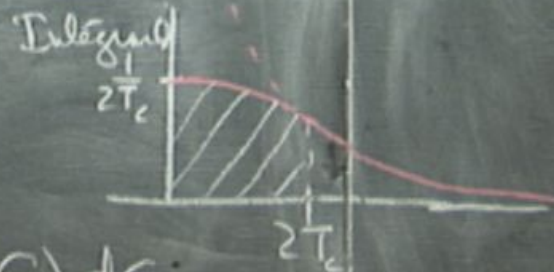
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$$T_c = 1.13 \omega_D e^{-\frac{1}{N(0)V}}$$



Weak Coupling, BCS Result

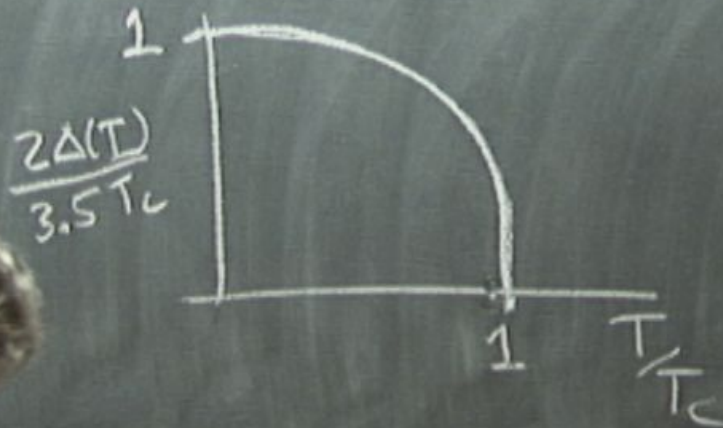
$$2\Delta(0) \approx 3.5 T_c$$

$\Delta(T)$

Weak Coupling, BCS Result

$$2\Delta(0) \approx 3.5 T_c$$

$\Delta(T)$



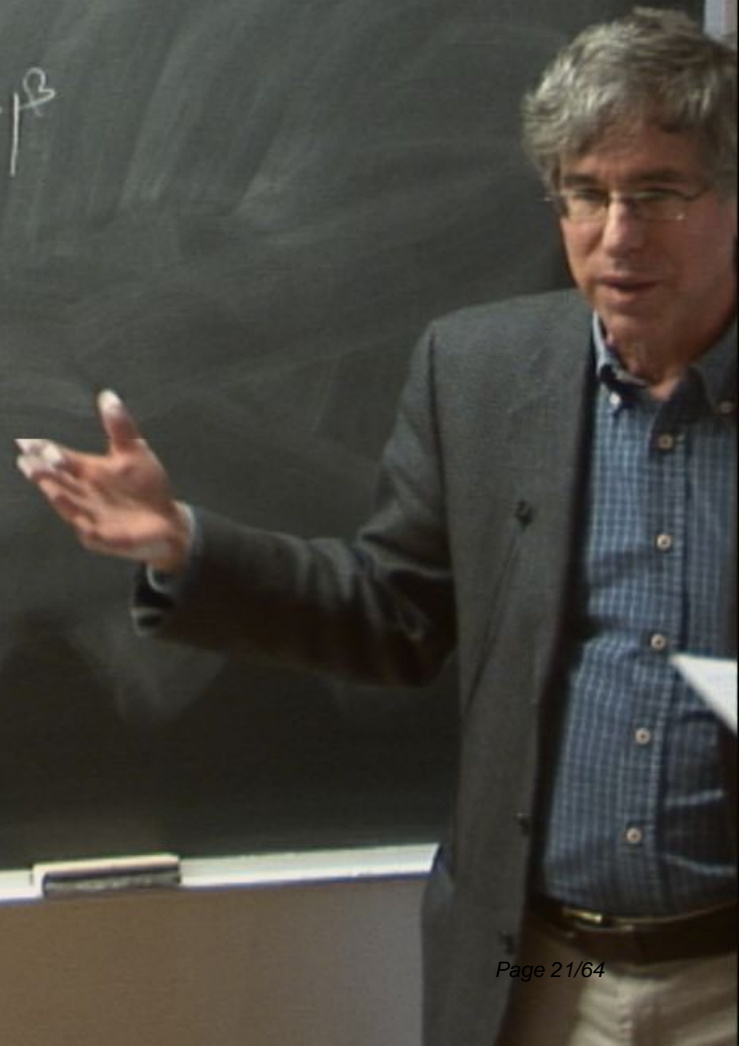
Weak Coupling, BCS Result

$$2\Delta(0) \approx 3.5 T_c$$

$$\sim 1 - a e^{-\Delta/T}$$

$\Delta(T)$

$$\frac{2\Delta(T)}{3.5 T_c}$$



Unconventional Metals and S.C.'s

Extended Hubbard Model

$$\mathcal{H} = -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + U \sum_i c_{i \uparrow}^\dagger c_{i \downarrow}^\dagger c_{i \downarrow} c_{i \uparrow}$$

Conventional Metals and S, C_i 's
 Extended Hubbard Model

$$-t \sum_{i, s, \alpha} c_{i+s, \alpha}^{\dagger} c_{i, \alpha} + U \sum_i c_{i \uparrow}^{\dagger} c_{i \downarrow}^{\dagger} c_{i \downarrow} c_{i \uparrow}$$

Weak Coupling

$$2 \Delta(0)$$

Hubbard 1

$$\frac{2\Delta(T)}{3.5 T_c}$$



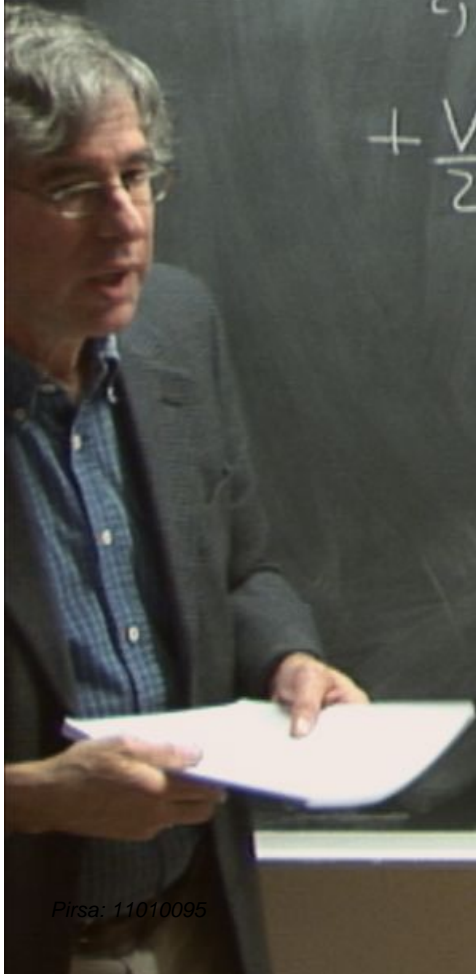
Unconventional Metals and S, C_i 's

Extended Hubbard Model

$$\begin{aligned}
 \mathcal{H} = & -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i\alpha} + U \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\downarrow} c_{i\uparrow} \\
 & + \frac{V}{2} \sum_{i, \delta, \alpha, \beta} c_{i\alpha}^\dagger c_{i+\delta, \beta}^\dagger c_{i+\delta, \beta} c_{i\alpha}
 \end{aligned}$$

Hubbard

$\frac{2}{3}$



Unconventional Metals and S, C, f 's

Extended Hubbard Model

$$\mathcal{H} = -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i\alpha} + U \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\downarrow} c_{i\uparrow}$$

$$+ \frac{V}{2} \sum_{i, \delta, \beta} c_{i\alpha}^\dagger c_{i+\delta, \beta}^\dagger c_{i+\delta, \beta} c_{i\alpha}$$

← "Extended" n.n.int

$$n_i = n_{i\uparrow} + n_{i\downarrow}$$

We

Hubbard

$\frac{2}{3}$

Unconventional Metals and S, C, i 's

Extended Hubbard Model

$$\mathcal{H} = -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + U \sum_i c_{i \uparrow}^\dagger c_{i \downarrow}^\dagger c_{i \downarrow} c_{i \uparrow}$$

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Hubbard

We

$\frac{2}{3}$

n.n. int

$+n_{i \downarrow}$

δ, α

Unconventional Metals and S, C_i 's

Extended Hubbard Model

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Consider $V=0$

← "Extended" n.n.i.

$$n_i = n_{i \uparrow} + n_{i \downarrow}$$

$$n_i n_{i+\delta} = 0, !$$

We

Hubbard

Unconventional Metals and S.C.'s

Extended Hubbard Model

$$\mathcal{H} = -t \sum_{i,j,\delta,\alpha} c_{i+\delta,\alpha}^\dagger c_{i\alpha} + U \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\downarrow} c_{i\uparrow}$$

$$+ \frac{V}{2} \sum_{i,j,\delta,\alpha,\beta} c_{i\alpha}^\dagger c_{i+\delta,\beta}^\dagger c_{i+\delta\beta} c_{i\alpha}$$

Consider $V=0$

1.) $U < 0 \rightarrow$ Superconductivity (S-wave singlet) $n_{i+\delta} = 0, 1, 2, \dots$

Hubbard

We

$\frac{2}{3}$

← "Extended" n.n. int

$$n_i = n_{i\uparrow} + n_{i\downarrow}$$

Unconventional Metals and S, C_i 's

Extended Hubbard Model

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Hubbard

$\frac{2}{3}$

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Unconventional Metals and S, C, i 's

Extended Hubbard Model

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Consider $V=0$

- 1.) $U < 0 \rightarrow$ Superconductivity (s-wave singlet) $n_{i+\delta} = 0, 1, 2, \dots$
- 2.) $U = 0$ and 1/2-filled ^{good} band metal

We

Hubbard

$\frac{2}{3}$

← "Extended" n.n. int

$$n_i = n_{i \uparrow} + n_{i \downarrow}$$

Unconventional Metals and S, C, I 's

Extended Hubbard Model

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Hubbard

$$+ V \sum_{i, \delta, \alpha, \beta} c_{i, \alpha}^\dagger c_{i+\delta, \beta}^\dagger c_{i+\delta, \beta} c_{i, \alpha}$$

← "Extended" n.n. int

$$n_i = n_{i \uparrow} + n_{i \downarrow}$$

when $V=0$

$t < 0 \rightarrow$ Superconductivity (s-wave singlet) $n_{i+\delta} = 0, 1, 2, \dots$

$t = 0$ and 1/2 site $\frac{1}{2}$ -filled ^{good} band metal

$\gg t$ and 1/2 site Mott Insulator

Unconventional Metals and S, C, I 's

Extended Hubbard Model

$$\mathcal{H} = -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + U \sum_i c_{i \uparrow}^\dagger c_{i \downarrow}^\dagger c_{i \downarrow} c_{i \uparrow}$$

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Consider $V=0$

- 1) $U < 0 \rightarrow$ Superconductivity (s-wave singlet)
- 2) $U = 0$ and 1 el/site $\frac{1}{2}$ -filled ^{good} band metal
- 3) $U \gg t$ and 1 el/site Mott Insulator

We

Hubbard

$\frac{2}{3}$

← "Extended" n.n. int

$$n_i = n_{i \uparrow} + n_{i \downarrow}$$

$$n_i n_{i+\delta} = 0, 1, 2, \dots$$

Unconventional Metals and S, C, I 's

Extended Hubbard Model

$$H = -t \sum_{i, \delta, \alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + U \sum_i c_{i \uparrow}^\dagger c_{i \downarrow}^\dagger c_{i \downarrow} c_{i \uparrow} \quad \text{Hubbard}$$

$$+ \frac{V}{2} \sum_{i, \delta, \alpha, \beta} c_{i, \alpha}^\dagger c_{i+\delta, \beta}^\dagger c_{i+\delta, \beta} c_{i, \alpha} \quad \leftarrow \text{"Extended" n.n. int}$$

Consider $V=0$

$$n_i = n_{i \uparrow} + n_{i \downarrow}$$

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Unconventional Metals and S.C.'s

Extended Hubbard Model

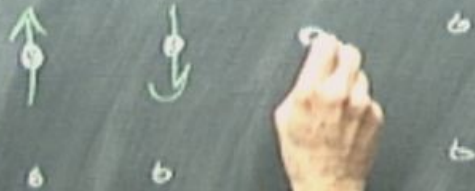
$$H = -t \sum_{i, s, \alpha} c_{i+s, \alpha}^\dagger c_{i, \alpha} + U \sum_i c_{i \uparrow}^\dagger c_{i \downarrow}^\dagger c_{i \downarrow} c_{i \uparrow} \quad \text{Hubbard}$$

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Consider $V=0$

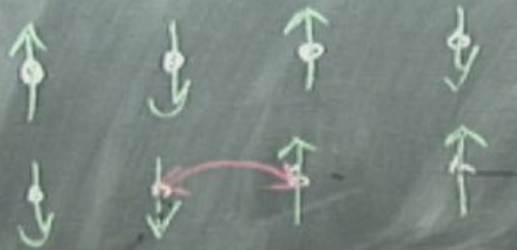
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- 3) $U \gg t$ and 1 el/site Mott Insulator

2^N -fold Spind degeneracy



\downarrow C_1 \uparrow

2^N -fold Spind degeneracy



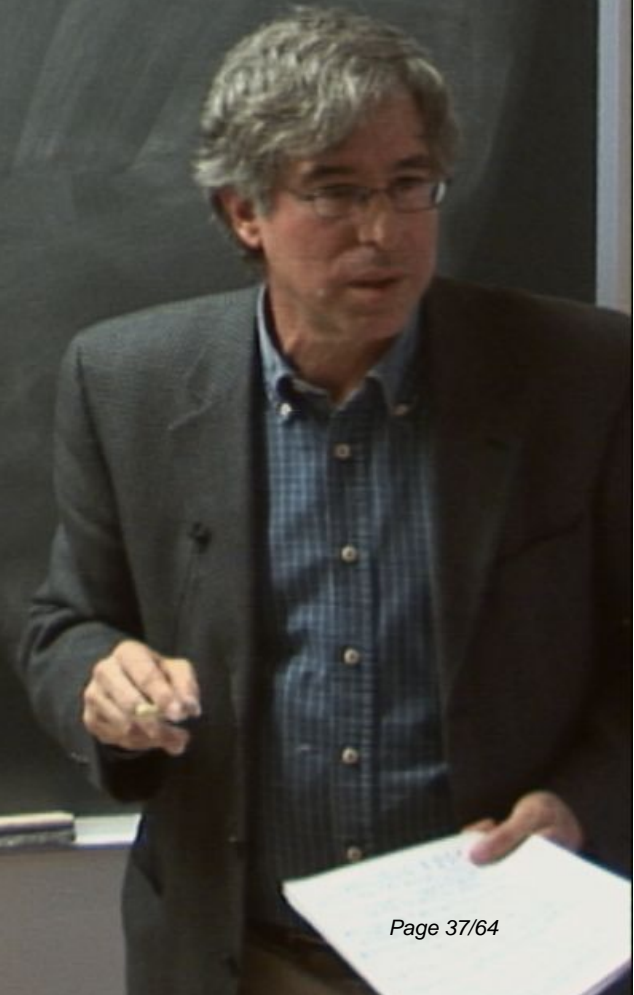
$\downarrow C_i \uparrow$

2^N -fold Spind degeneracy



$C_{1\uparrow}$

ingly
atal
the



2^N -fold Spind degeneracy

Antiferromagnetism

Consider 2 sites



$C_1 \uparrow$

2^N -fold Spin degeneracy

Antiferromagnetism

Consider 2 sites

$\downarrow \uparrow$



$\uparrow \downarrow$

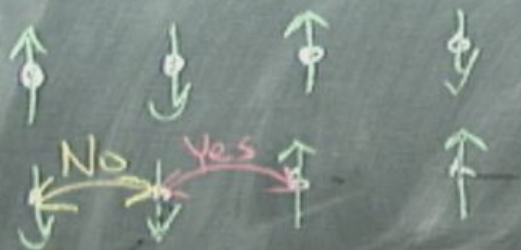
Energy lowered by $-\frac{4t^2}{u}$ for a singlet

" " by 0 for a triplet per

2^N -fold Spin degeneracy

Antiferromagnetism

Consider 2 sites

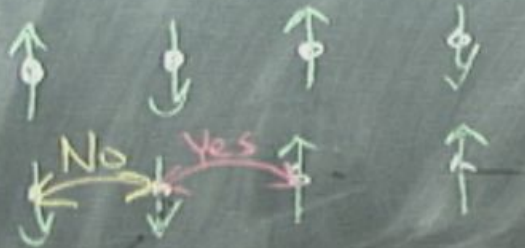


Energy lowered by $-\frac{4t^2}{U}$ for a singlet pair
 " " by 0 for a triplet pair

2^N -fold Spin degeneracy

Antiferromagnetism

Consider 2 sites



Energy lowered by $-\frac{4t^2}{U}$ for a singlet pair

$$\Delta E = \frac{4t^2}{U} \left[\frac{(\vec{S}_1 + \vec{S}_2)^2 - 2}{2} \right] (\vec{S}_1 + \vec{S}_2)^2 = 2 \text{ for triplet pair}$$

$$= 0 \text{ for singlet pair}$$

and S, C's
Model

$$-U \sum_i C_{i\uparrow}^+ C_{i\downarrow}^+ C_{i\downarrow} C_{i\uparrow}$$

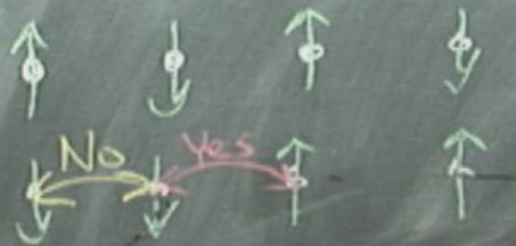
$$C_{i\beta} C_{i+\delta\beta} C_{i\alpha}$$

conductivity (S-wave singlet)

site $\frac{1}{2}$ -filled, ^{good} band metal

site Mott Insulator

2^N -fold Spind degeneracy



Antiferromagnetic

Consider



Energy lower

$$\Delta E = \frac{4t^2}{U} \left[\frac{(\overline{S_1 + S_2})^2 - 2}{2} \right] (\overline{S_1 + S_2})^2 = 2 = 0$$

idea

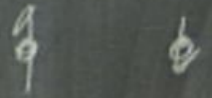
and S, C's
Model

$$-U \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\downarrow} c_{i\uparrow}$$

2^N -fold Spind degeneracy

Antiferromagnetic

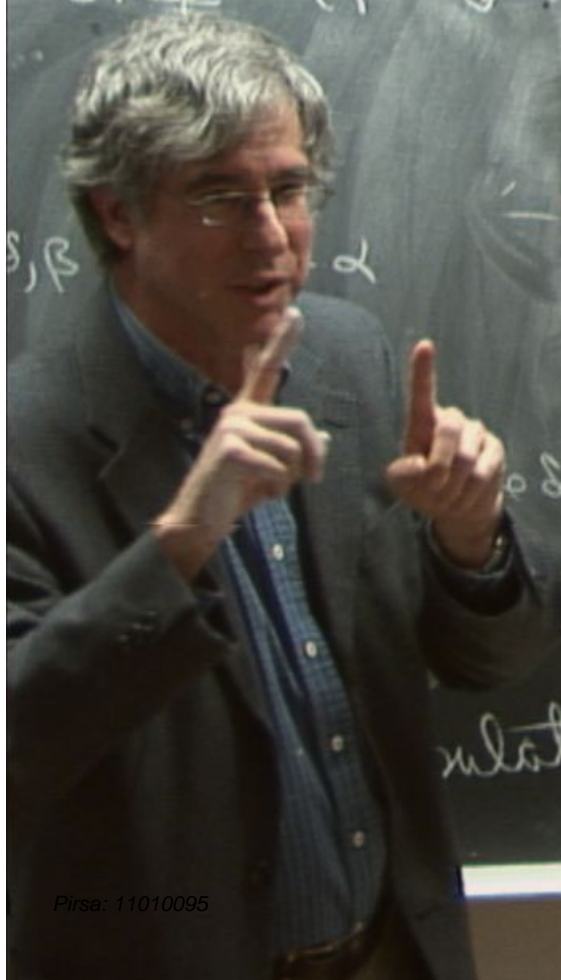
Consider



Energy lower



$$\Delta E = \frac{4t^2}{U} \left[\frac{(S_1 + S_2)^2 - 2}{2} \right] (S_1 + S_2)^2 = 2$$



of singlet
total
multiplicity

and S, C, S
Model

$$-U \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\downarrow} c_{i\uparrow} - U$$

2^N -fold Spin degeneracy

Antiferromagnetic

Consider

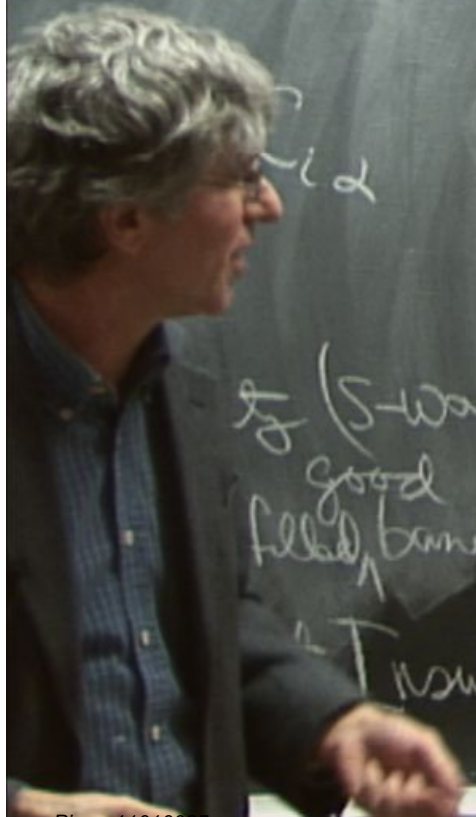


Energy lower

$$\Delta E = \frac{4t^2}{U} \left[\frac{(S_1 + S_2)^2 - 2}{2} \right] (S_1 + S_2)^2 = 2 = 0$$



(S-wave singlet)
good filled band metal
Insulator



and S, C's
Model

$$-U \sum_i C_{i\uparrow}^+ C_{i\downarrow}^+ C_{i\downarrow} C_{i\uparrow} - U$$

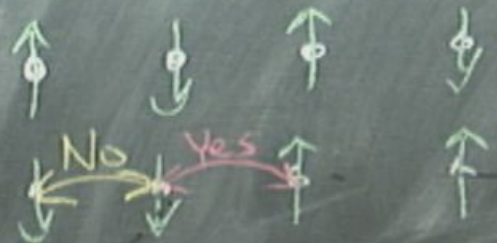
$$C_{i\uparrow} C_{i\downarrow}$$

conductivity (S-1)
site 1/2-filled,
site Mott J

2^N -fold Spind degeneracy

Antiferromagnetic

Consider



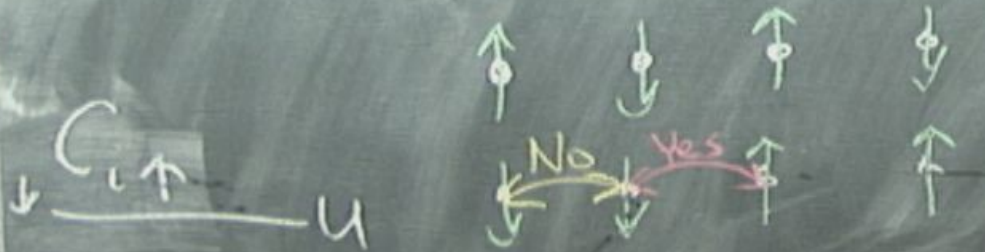
Energy lower

$$\Delta E = \frac{4t^2}{U} \left[\frac{(\overline{S_1 + S_2})^2 - 2}{2} \right] (\overline{S_1 + S_2})^2 = 2 = 0$$

2^N - fold Spin degeneracy

Antiferromagnetism

Consider 2 sites



Energy lowered by $-\frac{4t^2}{U}$ for a singlet pair

$$\Delta E = \frac{4t^2}{U} \left[\frac{(\vec{S}_1 + \vec{S}_2)^2 - 2}{2} \right]$$

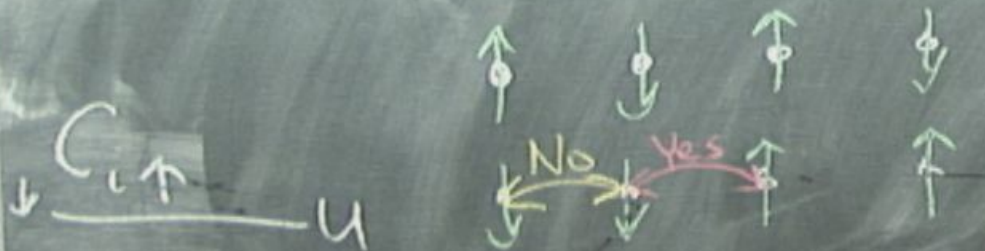
$(\vec{S}_1 + \vec{S}_2)^2 = 2$ for triplet
 $= 0$ for singlet

$$= \frac{4t^2}{U} \left[S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2 - 2 \right]$$

2^N -fold Spin degeneracy

Antiferromagnetism

Consider 2 sites



~~1 dozen~~

Energy lowered by $-\frac{4t^2}{u}$ for a singlet pair

$$\Delta E = \frac{4t^2}{u} \left[\frac{(\vec{S}_1 + \vec{S}_2)^2 - 2}{2} \right] \quad (\vec{S}_1 + \vec{S}_2)^2 = 2 \text{ for triplet } \begin{matrix} 0 \text{ for a} \\ \text{singlet pair} \end{matrix}$$

$$= 0 \text{ for singlet } \begin{matrix} 0 \text{ for a} \\ \text{triplet pair} \end{matrix}$$

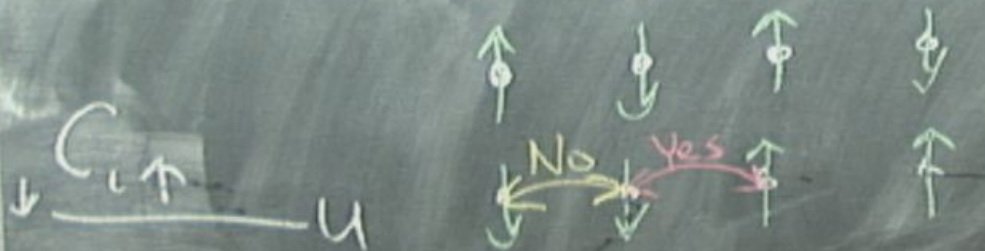
$$= \frac{4t^2}{u} \left[S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2 - 2 \right]$$

$$= \frac{4t^2}{u} \left[\vec{S}_1 \cdot \vec{S}_2 - \frac{1}{4} \right]$$

2^N - fold Spin degeneracy

Antiferromagnetism

Consider 2 sites



~~N dozen~~

Energy lowered by $-\frac{4t^2}{U}$ for a singlet pair

$$\Delta E = \frac{4t^2}{U} \left[\frac{(\vec{S}_1 + \vec{S}_2)^2 - 2}{2} \right] \quad (\vec{S}_1 + \vec{S}_2)^2 = 2 \text{ for triplet } \quad 0 \text{ for a triplet pair}$$

$$= 0 \text{ for singlet pair}$$

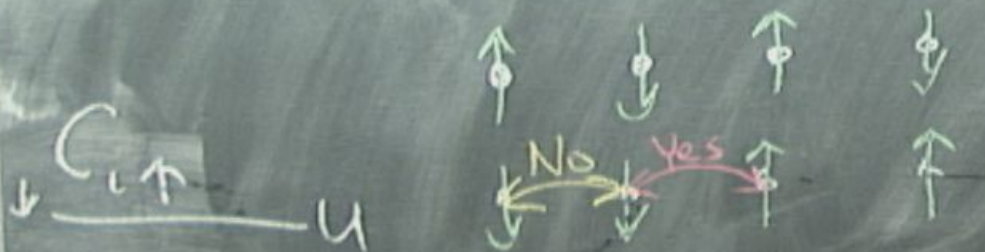
$$= \frac{4t^2}{U} \left[S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2 - 2 \right]$$

$$= \frac{4t^2}{U} \left[\vec{S}_1 \cdot \vec{S}_2 - \frac{1}{4} \right]$$

2^N -fold Spin degeneracy

Antiferromagnetism

Consider 2 sites



~~dozen~~

Energy lowered by $-\frac{4t^2}{u}$ for a singlet pair

$\Delta E = \frac{4t^2}{u} \left[\frac{(\vec{S}_1 + \vec{S}_2)^2 - 2}{2} \right]$

$(\vec{S}_1 + \vec{S}_2)^2 = 2$ for triplet $\Delta E = 0$ for a triplet pair

$= 0$ for singlet $\Delta E = 0$ for a singlet pair

$J = \frac{4t^2}{u} \left[S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2 - 2 \right]$

$= \frac{4t^2}{u} \left[\vec{S}_1 \cdot \vec{S}_2 - \frac{1}{4} \right]$

Do this for a lattice with ≤ 1 el/site

Ref: Web lectures of Jonathan Keeling
of Cambridge on "Quantum Magnetism"

$$H_{t-J} = -t \sum_{\langle i,j \rangle, \alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + J \sum_{\langle i,j \rangle} \left(\vec{S}_i \cdot \vec{S}_j - \frac{n_i n_j}{4} \right)$$

where $\vec{S}_i = \sum_{\alpha, \beta} c_{i, \alpha}^\dagger \frac{\vec{\sigma}_{\alpha\beta}}{2} c_{i, \beta}$

Do this for a lattice with $\leq 1 \text{ el/site}$

Ref: Web lectures of Jonathan Keeling
of Cambridge on "Quantum Magnetism"

$$H_{t,J} = -t \sum'_{i,j,\alpha} c_{i+\delta, \alpha}^{\dagger} c_{i, \alpha} + J \sum_{i,j} \left(\vec{S}_i \cdot \vec{S}_j - \frac{n_i n_j}{4} \right)$$

where $\vec{S}_i = \sum_{\alpha\beta} c_{i\alpha}^{\dagger} \frac{\vec{\sigma}_{\alpha\beta}}{2} c_{i\beta}$

\sum' means "no double occupancy"

Do this for a lattice with $\leq 1 \text{ el/site}$

Ref: Web lectures of Jonathan Keeling
of Cambridge on "Quantum Magnetism"

$$H_{t,J} = -t \sum_{i,j,\alpha} c_{i+\delta, \alpha}^\dagger c_{i, \alpha} + J \sum_{i,j} \left(\vec{S}_i \cdot \vec{S}_{i+\delta} - \frac{n_i n_{i+\delta}}{4} \right)$$

where $\vec{S}_i = \sum_{\alpha\beta} c_{i\alpha}^\dagger \frac{\vec{\sigma}_{\alpha\beta}}{2} c_{i\beta}$

\sum' means "no double occupancy"

Do this for a lattice with $\leq 1 \text{ el/site}$

Ref: Web lectures of Jonathan Keeling
of Cambridge on "Quantum Magnetism"

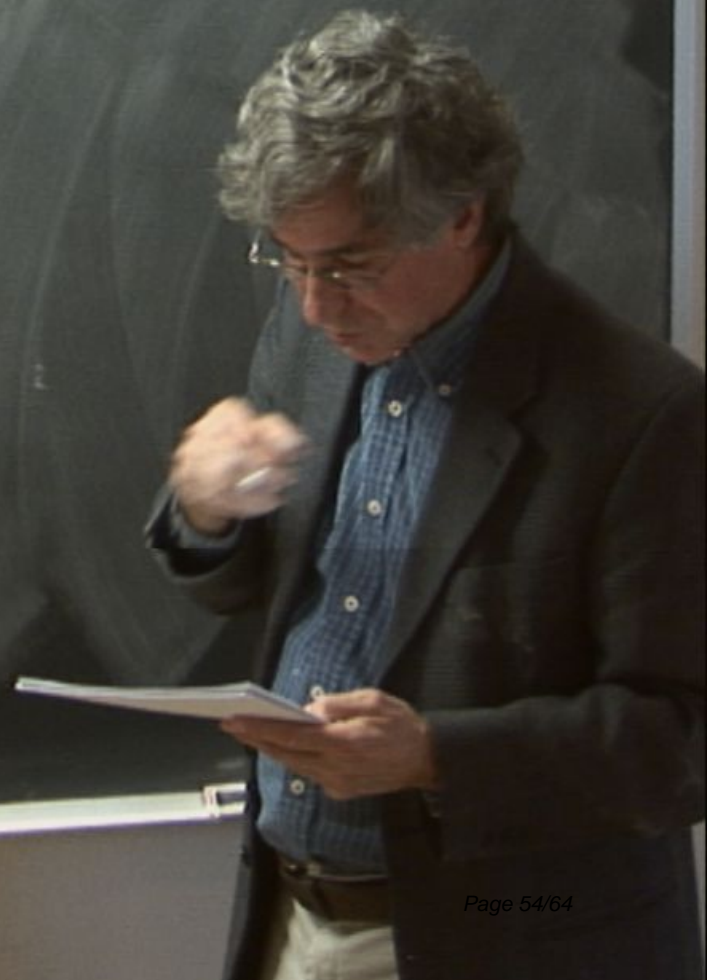
$$H_{t-J} = -t \sum_{i,j,\alpha} c_{i+\delta, \alpha}^\dagger c_{i\alpha} + J \sum_{i,j,\delta} \left(\vec{S}_i \cdot \vec{S}_{i+\delta} - \frac{n_i n_{i+\delta}}{4} \right)$$

$$\text{where } \vec{S}_i = \sum_{\alpha\beta} c_{i\alpha}^\dagger \frac{\vec{\sigma}_{\alpha\beta}}{2} c_{i\beta}$$

\sum' means "no double occupancy"

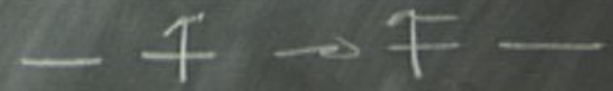
$$H_H = H_u + H_{t,D}^{\text{diagonal}} + H_{t,OD}^{\text{off-diag}}$$

"then"
 $\frac{N(4s)}{4}$



$$H_H = H_U + \overset{\text{diagonal}}{H_{t,D}} + \overset{\text{off-diag}}{H_{t,OD}}$$

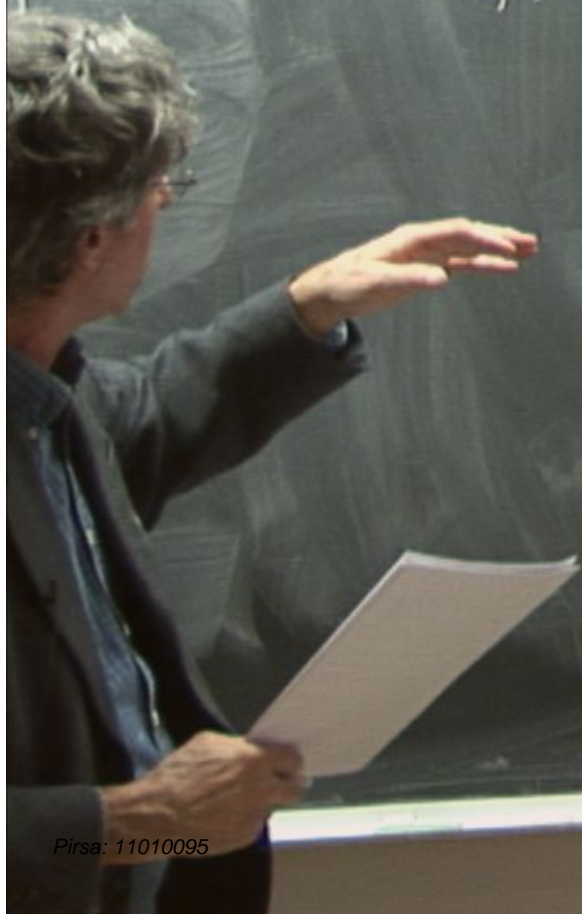
H_t terms
 - hopping of holes
 - hops



$$H_H = H_U + \overset{\text{diagonal}}{H_{t,D}} + \overset{\text{off-diag}}{H_{t,OD}}$$

$H_{t,D}$ has terms

- hopping of holes $\text{---} \uparrow \text{---} \rightarrow \uparrow \text{---}$
- hopping of singlet pairs $\uparrow \downarrow \text{---} \uparrow \text{---} \rightarrow \uparrow \text{---} \uparrow \downarrow$



Do a unitary transform

$$\tilde{H} = e^{iS} (H_u + H_{t,0} + H_{t,00}) e^{-iS}$$

$$= H_u + i[S, H_u] + H_{t,0} + H_{t,00}$$

$$+ i[S, H_{t,0} + H_{t,00}] - \frac{1}{2}[S, [S, H_u]]$$

Do a unitary transform

$$\tilde{H} = e^{iS} (H_u + H_{t,0} + H_{t,00}) e^{-iS}$$

$$= H_u + i[S, H_u] + H_{t,0} + H_{t,00}$$

$$+ i[S, H_{t,0} + H_{t,00}] - \frac{1}{2}[S, [S, H_u]]$$

choose S so that

$$i[S, H_u] + H_{t,00} = 0$$

Do a unitary transform

$$\tilde{H} = e^{iS} (H_u + H_{t,0} + H_{t,00}) e^{-iS}$$

$$= H_u + i[S, H_u] + H_{t,0} + H_{t,00}$$

$$+ i[S, H_{t,0} + H_{t,00}] - \frac{1}{2}[S, [S, H_u]]$$

Choose S so that

$$i[S, H_u] + H_{t,00} = 0 \rightarrow S \sim \frac{t}{\hbar} H_u$$

Do a unitary transform

$$\tilde{H} = e^{iS} (\mathcal{H}_u + \mathcal{H}_{t,0} + \mathcal{H}_{t,\infty}) e^{-iS}$$

$$= \mathcal{H}_u + i[S, \mathcal{H}_u] + \mathcal{H}_{t,0} + \mathcal{H}_{t,\infty}$$

$$+ i[S, \mathcal{H}_{t,0} + \mathcal{H}_{t,\infty}] - \frac{1}{2}[S, [S, \mathcal{H}_u]]$$

Choose S so that

$$i[S, \mathcal{H}_u] + \mathcal{H}_{t,0} = 0 \rightarrow S \sim \frac{t}{4}$$

Do a unitary transform

$$\tilde{H} = e^{iS} (\mathcal{H}_u + \mathcal{H}_{t,0} + \mathcal{H}_{t,00}) e^{-iS}$$

$$= \mathcal{H}_u + i[S, \mathcal{H}_u] + \mathcal{H}_{t,0} + \mathcal{H}_{t,00}$$

$$+ i[S, \mathcal{H}_{t,0} + \mathcal{H}_{t,00}] - \frac{1}{2}[S, [S, \mathcal{H}_u]]$$

Choose S so that

$$i[S, \mathcal{H}_u] + \mathcal{H}_{t,00} = 0 \rightarrow S \sim \frac{t}{4}$$

Eff. nn attr. from Hubbard Model

Suggests considering

$$\mathcal{H} = -t \sum_{i,j,\sigma} c_{i+\delta, \sigma}^{\dagger} c_{i, \sigma} - u \sum_{i, \alpha} c_{i, \alpha}^{\dagger} c_{i, \alpha}$$

$$+ \frac{1}{2} \sum_{i,j,\sigma} \left[\Delta_S c_{i \uparrow}^{\dagger} c_{i+\delta \downarrow}^{\dagger} + \Delta_S^* c_{i+\delta \downarrow} c_{i \uparrow} \right]$$

\mathcal{H}_u

$\rightarrow S \frac{t}{u}$

Eff. nn attr. from Hubbard Model

Suggests considering

$$\mathcal{H} = -t \sum_{\langle i, j \rangle} c_{i\sigma}^\dagger c_{j\sigma} - u \sum_i c_{i\uparrow}^\dagger c_{i\downarrow}^\dagger c_{i\uparrow} c_{i\downarrow}$$

$$+ \frac{1}{2} \sum_i \left[\Delta_S c_{i\uparrow}^\dagger c_{i+S\downarrow}^\dagger + \Delta_S^* c_{i+S\downarrow} c_{i\uparrow} \right]$$

Eff. nn attr. from Hubbard Model

Suggests considering

$$\mathcal{H} = -t \sum_{\langle i, j \rangle} c_{i+\delta}^{\dagger} c_j - u \sum_i c_i^{\dagger} c_i$$

$$+ \frac{1}{2} \sum_{\langle i, j \rangle} \left[\Delta_S c_{i\uparrow}^{\dagger} c_{j+\delta\downarrow} + \Delta_S^* c_{i+\delta\downarrow} c_{j\uparrow} \right]$$

\mathcal{H}_u

$\rightarrow \frac{t}{4}$