

Title: Foundations of Quantum Mechanics - Lecture 13

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URL: <http://pirsa.org/11010054>

Abstract:

The deBroglie-Bohm interpretation



Louis deBroglie
(1892-1987)



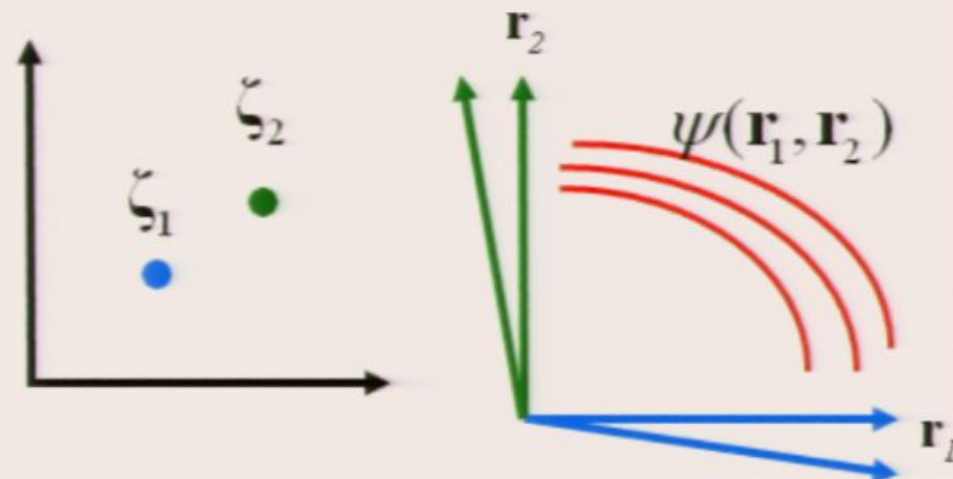
David Bohm
(1917-1992)

The deBroglie-Bohm interpretation for many particles

The ontic state: $(\psi(\mathbf{r}_1, \mathbf{r}_2), \zeta_1, \zeta_2)$

\nearrow
 Wavefunction on
 configuration space

\uparrow
 Particle
 positions



The evolution equations:

Schrödinger's equation

$$i\hbar \frac{\partial \psi(\mathbf{r}_1, \mathbf{r}_2, t)}{\partial t} = -\frac{\hbar^2}{2m_1} \nabla_1^2 \psi(\mathbf{r}_1, \mathbf{r}_2, t) - \frac{\hbar^2}{2m_2} \nabla_2^2 \psi(\mathbf{r}_1, \mathbf{r}_2, t) + V(\mathbf{r}_1, \mathbf{r}_2) \psi(\mathbf{r}_1, \mathbf{r}_2, t)$$

$$\left. \begin{aligned} \frac{d\zeta_1(t)}{dt} &= \frac{1}{m_1} [\nabla_1 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\zeta_1(t), \mathbf{r}_2=\zeta_2(t)} \\ \frac{d\zeta_2(t)}{dt} &= \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\zeta_1(t), \mathbf{r}_2=\zeta_2(t)} \end{aligned} \right\} \text{The guidance equation}$$

$$\psi = \sum_j c_j \psi_j$$

"waves" of the decomposition

$\zeta \in$ Spatial support of ψ_j j th wave is **occupied**

$\zeta \notin$ Spatial support of ψ_j j th wave is **empty**

If only the k th wave is occupied

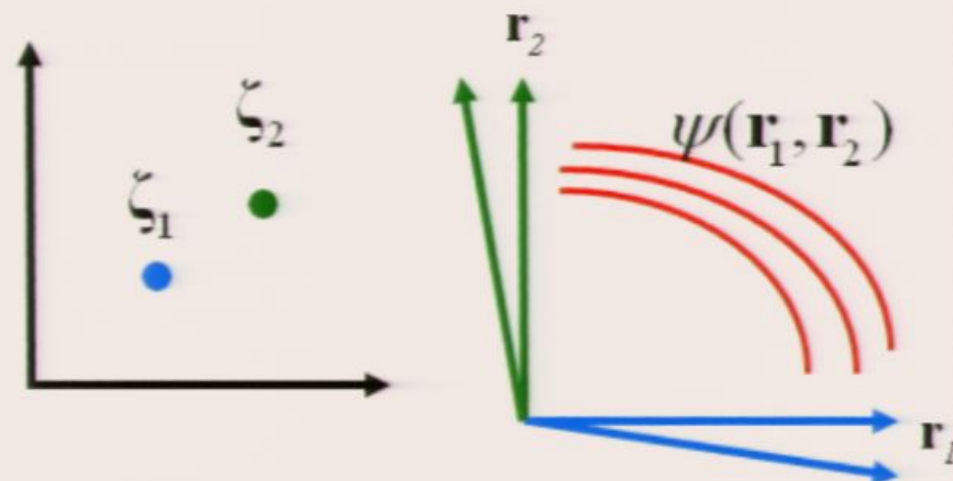
Then the guidance equation depends only on the k th wave

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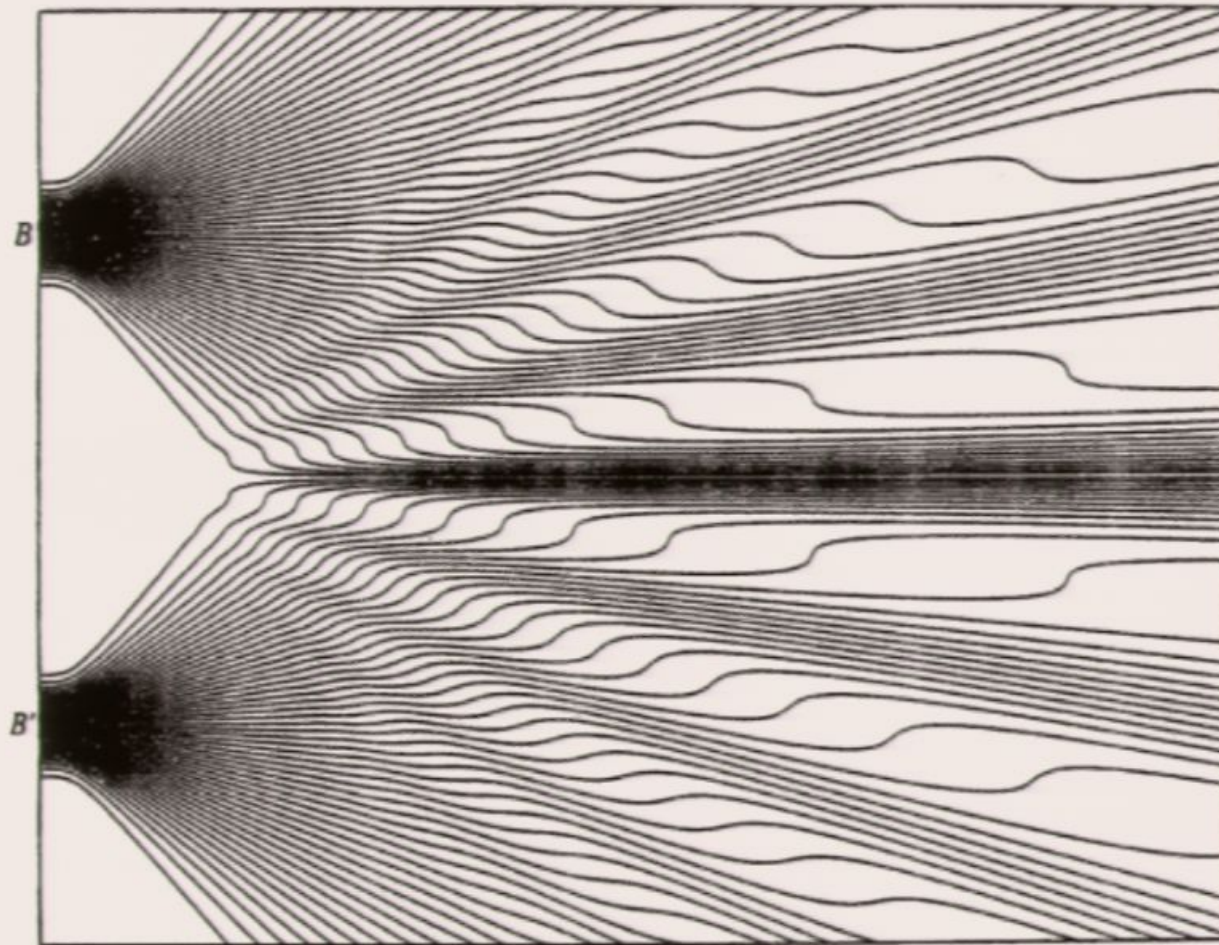
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Double slit experiment

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \phi^{(1)}(\mathbf{r}_1, t) \chi^{(2)}(\mathbf{r}_2, t) \quad \text{Product state}$$

$$= R_1(\mathbf{r}_1, t) e^{iS_1(\mathbf{r}_1, t)/\hbar} R_2(\mathbf{r}_2, t) e^{iS_2(\mathbf{r}_2, t)/\hbar}$$

$$S(\mathbf{r}_1, \mathbf{r}_2, t) = S_1(\mathbf{r}_1, t) + S_2(\mathbf{r}_2, t)$$

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The two particles evolve independently

Reproducing the operational predictions

Consider a measurement of A with eigenvectors $\phi_k(\mathbf{r})$

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Assumption: different outcomes of a measurement correspond to disjoint regions of the configuration space of the apparatus

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To have re-interference with the empty waves, it would be necessary to map all the η_k back to η

Do measurements reveal attributes of the particles?

Position measurements are "statistically faithful":

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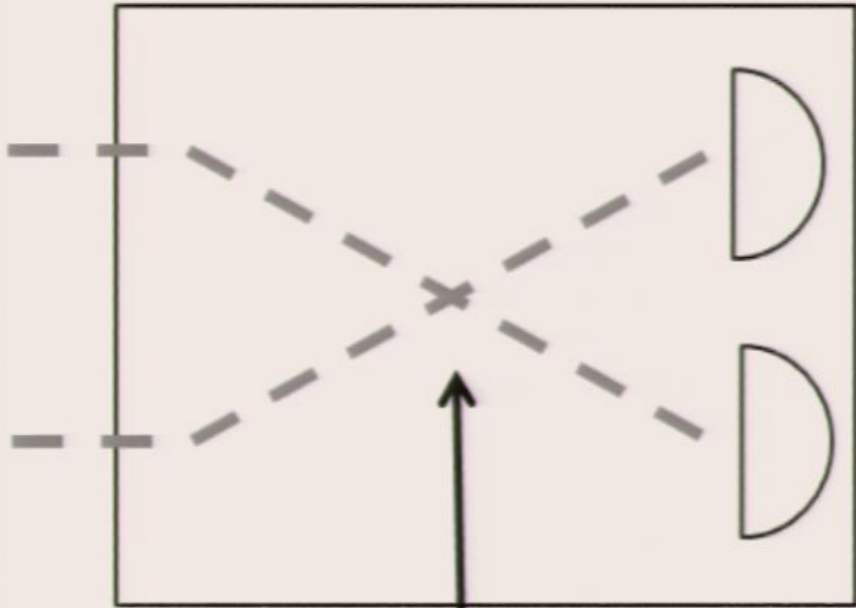
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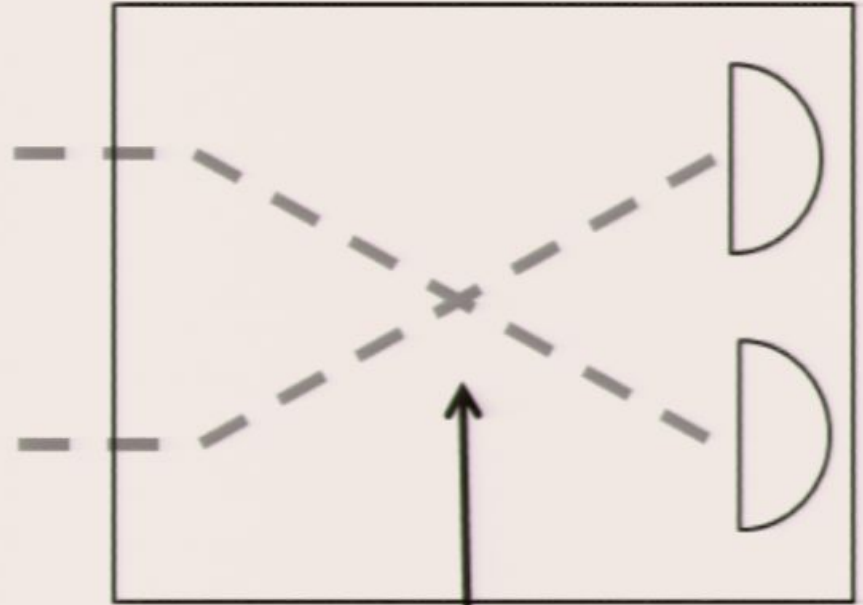
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Contextuality

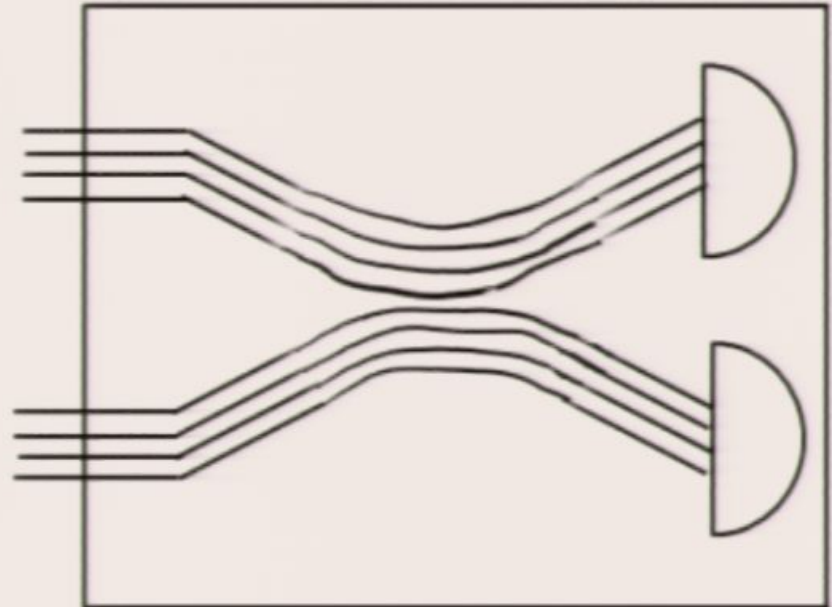
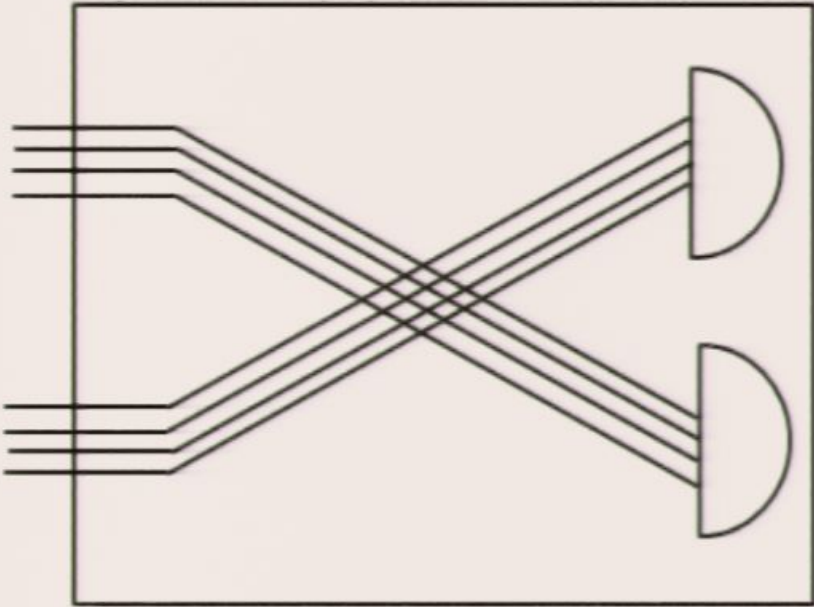


No overlap in 3D

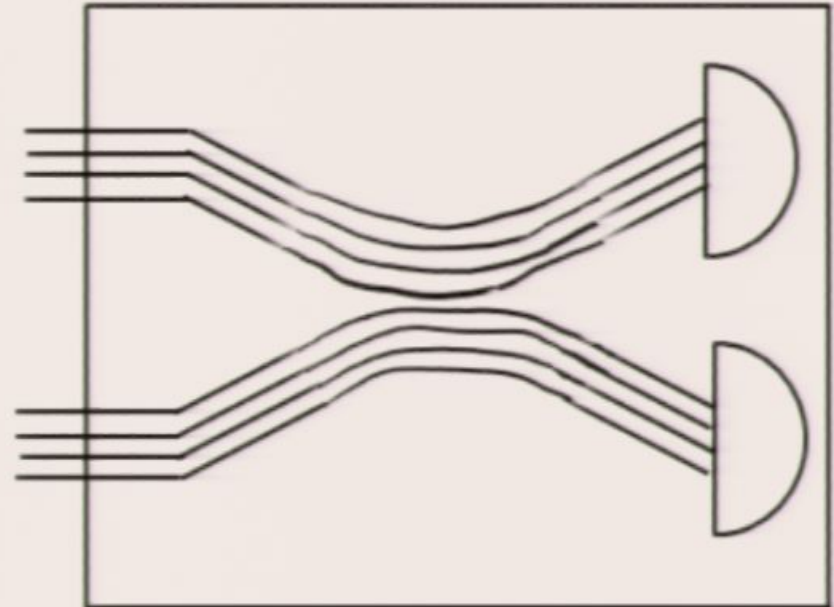
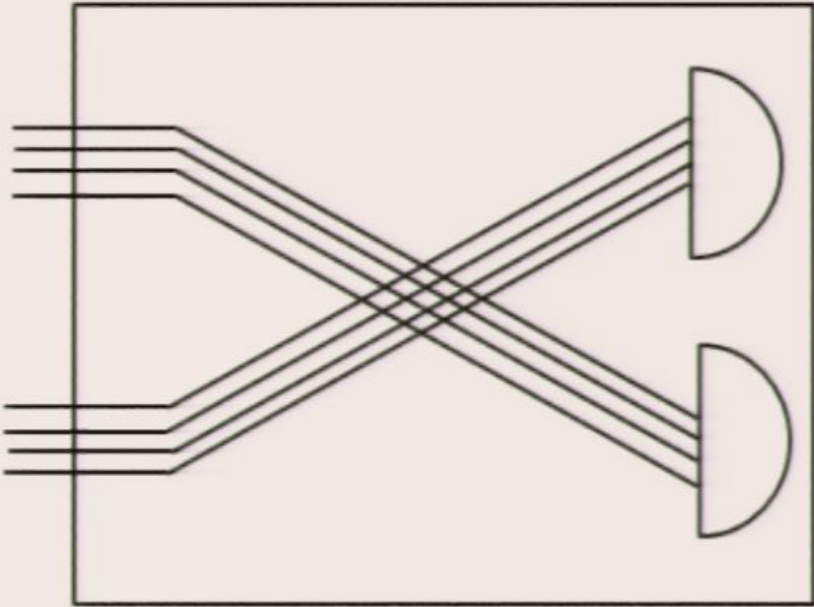


Overlap

Contextuality

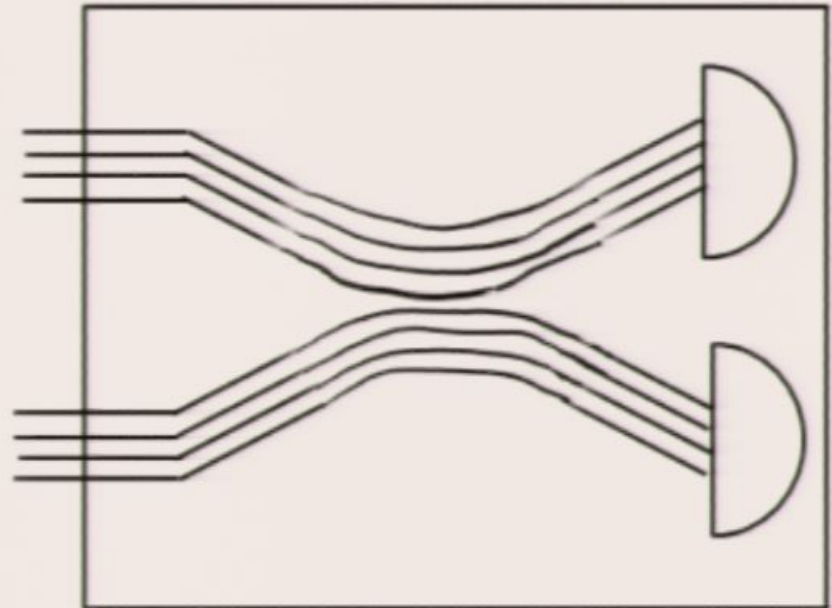
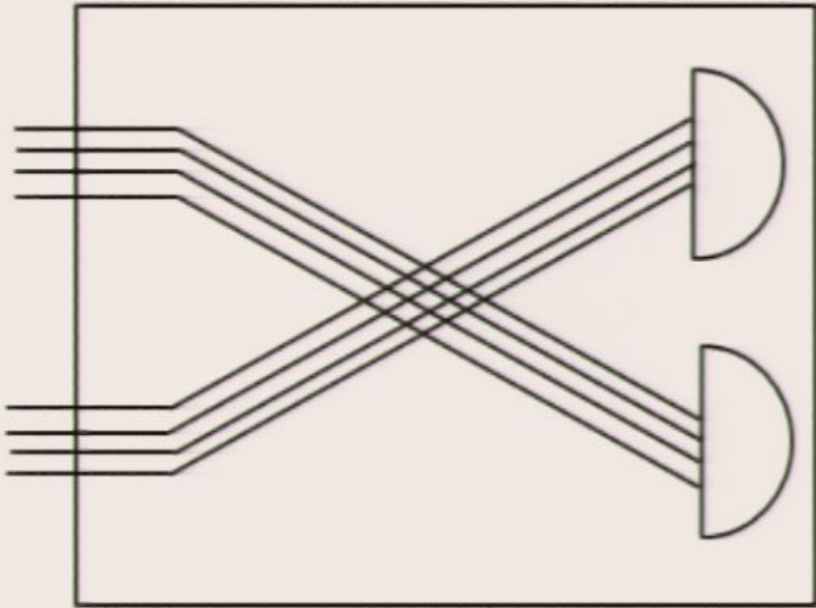


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"Surrealistic" trajectories?

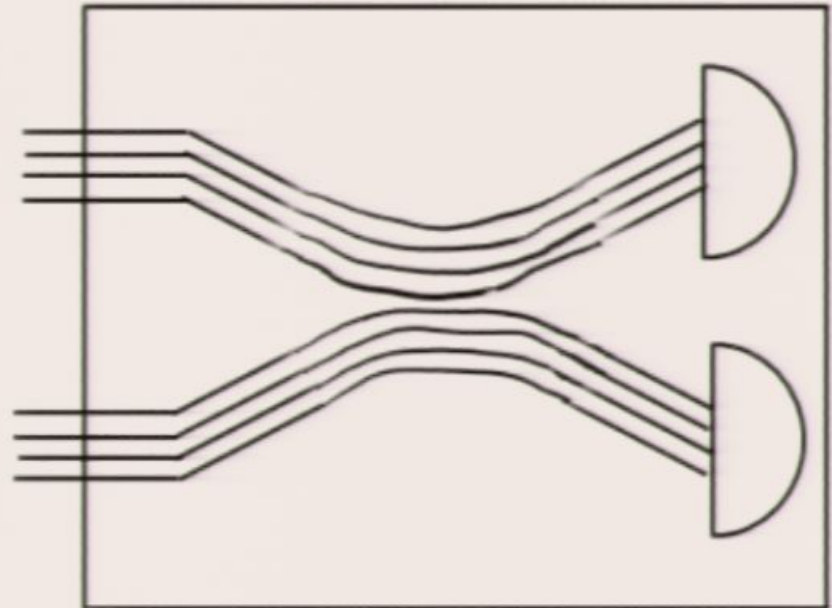
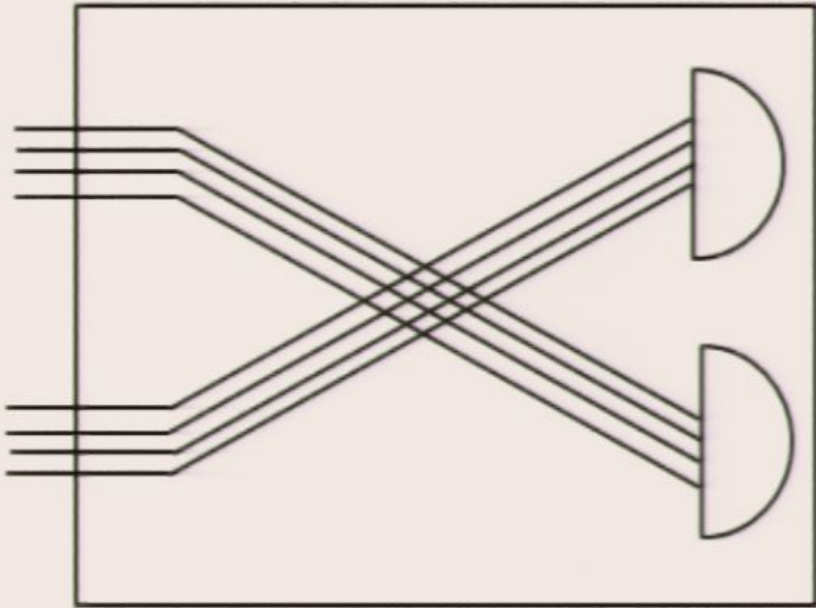
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Conclusion: it does not make sense to associate an attribute with an operator without also specifying the full experimental arrangement

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Operational correspondence vs. Ontological correspondence

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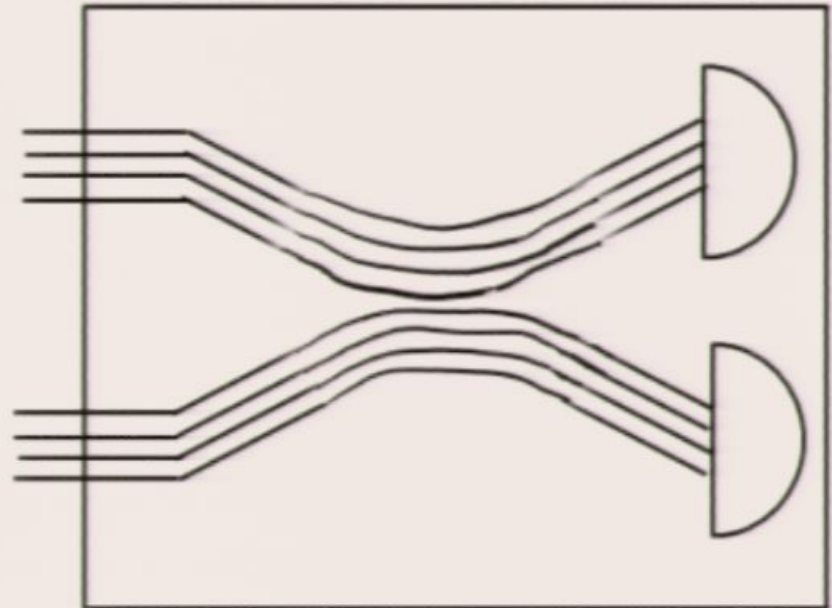
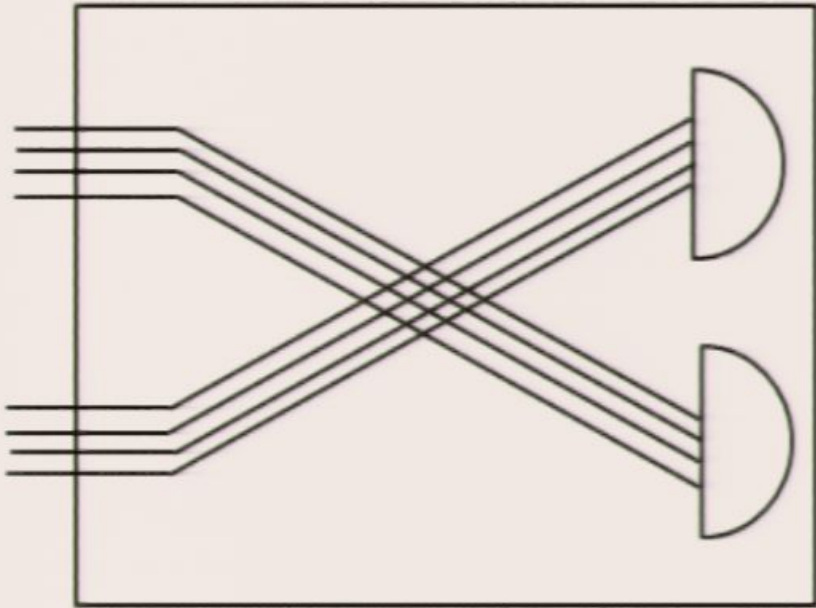
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Response:

Lorentz invariance is an emergent symmetry - a statistical consequence of quantum equilibrium

The underdetermination criticism

Underdetermination: when there are many possible choices of ontological structure that are consistent with observations

Underdetermination of the supplementary variables

Standard approach - Position preferred

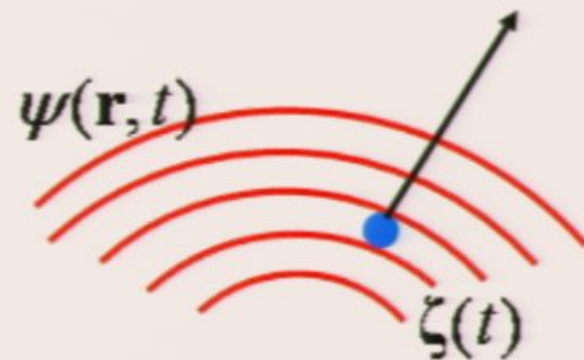
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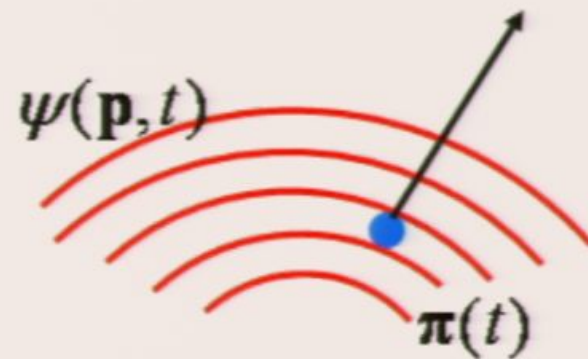
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The problem returns for the case of a deBroglie-Bohm theory of the electromagnetic field:

Supplementary variable: **Electric field** or **Magnetic field**?

Underdetermination of the supplementary variables

Multiple treatments of spin:

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Bohm, Schiller and Tiomno approach

Supplementary variables: particle position and orientation

The particle is taken to be an extended rigid object which makes a 'spin' contribution to the total angular momentum

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or

Bell's minimalist approach

Supplementary variables: **particle position**

The effect of spin is seen only in the dynamics of the particle positions

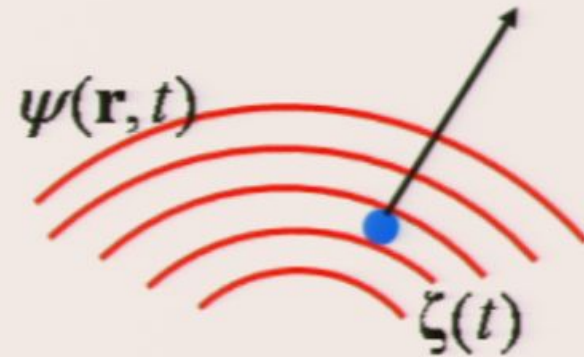
The operational predictions are reproduced by virtue of localization of pointers

Bell's minimalist approach to spin

Bell's minimalist approach to spin

The ontic state: $(\psi(\mathbf{r}), \zeta)$

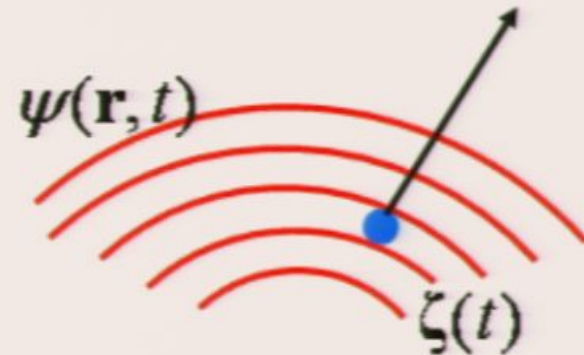
Two-component
wavefunction \nearrow
Particle
position \nwarrow



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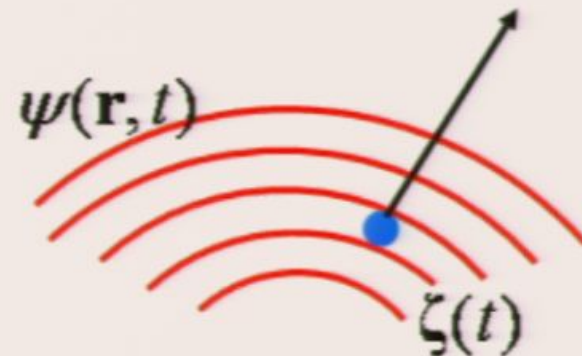
$$i\hbar \frac{\partial \psi(\mathbf{r}, t)}{\partial t} = \left[-\frac{\hbar^2}{2m} \left(\nabla - \frac{ie}{\hbar c} \mathbf{A} \right)^2 + \boldsymbol{\sigma} \cdot \mathbf{B} + V(\mathbf{r}) \right] \psi(\mathbf{r}, t) \quad \text{Pauli eq'n}$$

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$$\frac{d\zeta(t)}{dt} = \frac{\mathbf{j}(\mathbf{r}, t)}{R(\mathbf{r}, t)} \Big|_{\mathbf{r}=\zeta(t)} \quad \text{The guidance eq'n}$$

$$\text{where } R(\mathbf{r}, t) = \sum_s |\psi_s(\mathbf{r}, t)|^2$$

$$\mathbf{j}(\mathbf{r}, t) = \sum_s \left(\frac{\hbar}{2mi} (\psi_s^* \nabla \psi_s - \psi_s \nabla \psi_s^*) - \frac{e}{mc} \mathbf{A} \psi_s^* \psi_s \right) (\mathbf{r}, t)$$

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Bohm's model of the free electromagnetic field

Supplementary variables: electric field (or magnetic field)

combined with

Bell's model of fermions (indeterministic, discrete) or Colin's continuum version of it

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Supplementary variables: fermion number at each lattice point
(Note: field variables for fermions have been problematic)

or

Struyve and Westman's minimalist model of QED:

Supplementary variables : magnetic field (or electric field)
(the classical EM field carries an image of pointer positions)