Title: Two Higgs doublet models/Baryogenesis and LHC Date: Dec 07, 2010 03:40 PM URL: http://pirsa.org/10120030 Abstract: TBA

DØ dimuon anomaly and electroweak baryogenesis



Jim Cline (McGill U. and PI) with Kimmo Kainulainen (Jyväskylä) and Mike Trott (PI) PI-ATLAS LHC day, 7 Dec. 2010

New source of CP violation?

DØ observes 3.2σ deviation from SM prediction of



$$a_{SL}^b = \frac{N_b^{++} - N_b^{--}}{N_b^{++} + N_b^{--}}$$

from semileptonic B decays

 $B_{d,s} \to \mu^{\pm} X$

 $(a_{sl}^d \text{ and } a_{sl}^s \text{ are respective}$ contributions from $B_{d,s}$)

 $B_s \rightarrow J/\psi \phi$ and $B^- \rightarrow \tau \nu$ also deviate from SM. Evidence for new CP violation beyond SM

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A simple model

Two Higgs doublets (*H*, *S*) with minimal flavor violation (MFV) Trott, Wise 1009.2813

New Higgs S^0 FCNC couplings to b (and t) are CKM-suppressed:

 $y_b \overline{b}_L \left(H^0 \delta_{bi} + (\eta_D \delta_{bi} + \eta'_D V_{tb} V_{ti}^*) S^0 \right) q_R^i$



$$\sim {\eta'_D}^2 y_b^2 \frac{m_{S_R}^2 - m_{S_I}^2}{m_S^4} (V_{tb} V_{ti}^*)^2$$

same CKM structure as SM box diagram contribution

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New **CP** couplings

In addition to \mathcal{CP} Yukawa couplings η_U , η'_U , η_D , η'_D , scalar potential has new \mathcal{CP} couplings,

$$V = \lambda \left(H^{\dagger i} H_i - \frac{1}{2} v^2 \right)^2 + m_1^2 \left(S^{\dagger i} S_i \right)$$

+ $\left(\underline{m_2^2} H^{\dagger i} S_i + \text{h.c.} \right) + \lambda_1 \left(H^{\dagger i} H_i \right) \left(S^{\dagger j} S_j \right)$
+ $\lambda_2 \left(H^{\dagger i} H_j \right) \left(S^{\dagger j} S_i \right) + \left[\underline{\lambda_3} H^{\dagger i} H^{\dagger j} S_i S_j + \text{h.c.} \right]$
+ $\left[\underline{\lambda_4} H^{\dagger i} S^{\dagger j} S_i S_j + \underline{\lambda_5} S^{\dagger i} H^{\dagger j} H_i H_j + \text{h.c.} \right]$
+ $\lambda_6 \left(S^{\dagger i} S_i \right)^2$

of which 2 can be removed by field redefinitions.

Note: *H* is the "real Higgs," only $\langle H \rangle \neq 0$. New Higgses are $S_{R,I}^0$, S_{\pm} . There is no "tan β "; *H* and *S* both couple to all quarks.

CP and baryogenesis

People like to say that new *CP* is exciting because of baryogenesis. Electroweak baryogenesis might be testable at LHC. Can we put these together?

Necessary ingredients:

- new CP 🗸
- baryon violation B



SM provides anomalous **B** at high temperature through sphalerons

• getting out of thermal equilibrium ?

Electroweak Baryogenesis

- Sakharov: must go out of thermal equilibrium to make baryon asymmetry.
- Getting out of equilibrium can be achieved in a first order electroweak phase transition.



Phase transition is too weak in Standard Model, need new Higgs physics

How it works

At critical temperature T_c ~ 100 GeV, bubbles of true vacuum (⟨H⟩ ≠ 0) form and start expanding.
Particles reflect off wall in a CP violating way.
Baryon asymmetry forms inside the bubble.



- Sphalerons eat excess left-handed antiquarks in front of wall
- baryon asymmetry is created
- baryons diffuse back inside expanding bubble
- sphalerons must be ineffective inside bubble, otherwise baryons decay away: need $\langle H \rangle > T_c$



Electroweak baryogenesis has been previously considered in 2 Higgs doublet models. But our analysis is different:

 MFV couplings of quarks to Higgses instead of discrete symmetries to avoid FCNC's

 $\mathcal{L}_{\bar{q}Hq} = \bar{u}_R y_u QH + \bar{d}_R y_d QH^{\dagger}$ $+ \bar{u}_R (\eta_u y_u + \eta'_u y_u y_u^{\dagger} y_u) QS + \bar{d}_R (\eta_d y_d + \eta'_d y_d y_u^{\dagger} y_u) QS^{\dagger}$

• Explain dimuon + B decay anomalies

• Respect numerous particle physics constraints: R_b , EWPO, LEP/Tevatron mass limits, $b \rightarrow s\gamma$, neutron EDM, vacuum stability, Landau poles

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Pheno constraints are stringent

We find it hard to satisfy particle physics constraints and get a strong phase transition.

1. Search grid in $\{m_h, m_1, \lambda_1, \lambda_2, \lambda_3\}$ space for strong phase transitions. Then filter with pheno constraints. All examples excluded by EWPO and dimuons $+R_b!$

2. Search grid for pheno-allowed points, then filter results on strong phase transition. Low m_h is favored.

m_h (GeV)	∦ start	# strong p.t.
115	210,000	92
120	195,000	49
130	171,000	25

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Distribution of S_I^0 - S_R^0 mass splittings:



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S_I^0 - S_R^0 mass splittings, pheno constraints only:



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Collider signatures

Parameter η'_D is large, $\sim 5 - 10$; appears in

$\eta'_D y_b \, \overline{b} Q S^\dagger$

Trott & Wise (1009.2813) suggest collinear gluon splitting + b quark fusion, followed by $S^0 \rightarrow b\overline{b}$, as main discovery channel at LHC.



Resulting 4b events have higher p_T , lower rapidity than ^{irse: 10120}SM background, plus resonance in one $b\overline{b}$ pair.

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S^0 production cross section

From Trott & Wise (1009.2813)



2b-2t events

We also predict production of $t\overline{b} + S_{-}$



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Is this as promising as 4b channel, or more so?Is production cross section greatly suppressed?

Production via W, Z

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Pirsa: 10120030

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The challenge of baryogenesis

So far we only included necessity of strong phase transition, $\langle h \rangle / T_c > 1$.

Getting enough baryogenesis is even more rare.

Case study that yields observed baryon asymmetry:

 $m_h = 115, \ m_{S_R} = 318, \ m_{S_T} = 201, \ m_{\pm} = 219 \text{ GeV}$

Baryogenesis is sensitive to values of

 $\lambda_4 = 0.23I, \quad \lambda_5 = 0, \quad \eta_U = 0.125$

which impact complex t quark mass in bubble wall,

$$m_t(z) \cong rac{y_t}{\sqrt{2}} \left(h(z) + \eta_U s(z)
ight)$$

 $\frac{d}{dz} \operatorname{Im}(m_t(z))$ must be nonnegligible in bubble wall.

Potential barrier + top mass phase

Phase of $m_t(z)$ comes from $\langle s_I(z) \rangle$ induced by Im (λ_4) in bubble wall: $\arg(m_t(z)) \sim \eta_U s_I(z)/h(z)$



We saturate $|\eta_U| \leq 0.125$, the 1- σ upper limit from R_b . Plise: 101205 Uning λ_4 , λ_5 , $\arg(\eta_U)$ gives barely large enough

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Why is it so difficult?

Naively, increasing Im(λ_5) increases Im($m_t(z)$) in bubble wall, good for baryogenesis.

Nonzero λ_5 makes it easier to explain dimuon anomaly with smaller η'_D , good for R_b

$$\begin{split} C^{\rm NP}(m_t) \ &= \ (\eta'_D)^2 \left(\frac{\sqrt{2}\,m_t}{v}\right)^4 \left[\frac{\lambda_3\,m_b^2}{m_s^4 - \lambda_3^2 v^4} + \frac{(\lambda_5^R)^2 \,v^2 \,m_b^2}{2 \,(m_s^2 + \lambda_3 \,v^2 - m_h^2) \,(m_s^4 - \lambda_3^2 \,v^4)} + \frac{(\lambda_5^R)^2 \,v^2 \,m_b^2}{2 \,(m_s^2 + \lambda_3 \,v^2 - m_h^2)^2 \,m_h^2}\right] \\ &+ \ (\eta'_D)^2 \left(\frac{\sqrt{2}\,m_t}{v}\right)^4 \left[-\frac{(\lambda_5^I)^2 \,v^2 \,m_b^2}{2 \,(m_s^2 - \lambda_3 \,v^2 - m_h^2) \,(m_s^4 - \lambda_3^2 \,v^4)} + \frac{(\lambda_5^I)^2 \,v^2 \,m_b^2}{2 \,(m_s^2 - \lambda_3 \,v^2 - m_h^2)^2 \,m_h^2}\right]. \end{split}$$

But λ_5 tends to kill delicate barrier in potential, making transition 2nd order:

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Integrating out S gives

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Pirsa: 10120030n anti-bump that cancels the positive barrier

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Also, $\text{Im}(m_t) \sim \eta_U$ must be small due to neutron EDM and R_b constraints

Further challenges

Baryon asymmetry depends on network of Boltzmann equations for all species of particles near bubble wall

(Fromme, Huber, Seniuch, hep-ph/0605242)

 $3v_{w}K_{1,i}\mu'_{t,2} + 3v_{w}K_{2,i}(m_{t}^{2})'\mu_{t,2} + 3u'_{t,2}$ $-3\Gamma_{y}(\mu_{t,2} + \mu_{t^{e},2} + \mu_{h,2}) - 6\Gamma_{w}(\mu_{t,2} + \mu_{t^{e},2}) - 3\Gamma_{W}(\mu_{t,2} - \mu_{h,2})$ $-3\Gamma_{ss}[(1 + 9K_{1,i})\mu_{t,2} + (1 + 9K_{1,b})\mu_{h,2} + (1 - 9K_{1,t})\mu_{t^{e},2}] = 0$

 $\begin{aligned} & 3v_{\rm w}K_{1,b}\mu_{b,2}' + 3u_{b,2}' \\ & -3\Gamma_s(\mu_{b,2} + \mu_{t^{\rm c},2} + \mu_{b,2}) - 3\Gamma_W(\mu_{b,2} - \mu_{t,2}) \\ & -3\Gamma_{ss}[(1+9K_{1,2})\mu_{t,2} + (1+9K_{1,2})\mu_{b,2} + (1-9K_{1,2})\mu_{t^{\rm c},2}] = 0 \end{aligned}$

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 $4v_w K_{1,\bar{a}} \mu'_{h,2} + 4u'_{h,2} \\ -3\Gamma_g (\mu_{t,2} + \mu_{b,2} + 2\mu_{F,2} + 2\mu_{h,2}) - 4\Gamma_h \mu_{h,2} = 0$

 $-3K_{4,t}\mu'_{t,2} + 3v_{w}\tilde{K}_{5,t}u'_{t,2} + 3v_{w}\tilde{K}_{6,t}(m_{t}^{2})'u_{t,2} + 3\Gamma_{t}^{tot}u_{t,2} = S_{t}$ $-3K_{4,b}\mu'_{b,2} + 3v_{w}\tilde{K}_{5,t}u'_{b,2} + 3\Gamma_{b}^{tot}u_{t,2} = 0$ $-3K_{4,t}\mu'_{t^{2},2} + 3v_{w}\tilde{K}_{5,t}u'_{t^{2},2} + 3v_{w}\tilde{K}_{6,t}(m_{t}^{2})'u_{t^{2},2} + 3\Gamma_{t}^{tot}u_{t^{2},2} = S_{t}$ $-4K_{4,b}\mu'_{b,2} + 4v_{w}\tilde{K}_{5,b}u'_{b,2} + 4\Gamma_{b}^{tot}u_{b,2} = 0.$ source term, $\theta = \text{Im}(m_{t})$:

Source term, $0 = \min(\operatorname{III}_{\mathfrak{t}})^{\prime}$ $S_{\mathfrak{t}} = -v_{\mathfrak{w}}K_{\mathfrak{s}}(m_{\mathfrak{t}}^{2}\theta_{\mathfrak{t}}^{\prime})^{\prime} + v_{\mathfrak{w}}K_{\mathfrak{s}}\theta_{\mathfrak{t}}^{\prime}m_{\mathfrak{t}}^{2}(m_{\mathfrak{t}}^{2})^{\prime}.$

To simplify network, FHS follow t_L , t_R , b_L , h (assuming all Higgses have same asymmetry).

Results depend upon how many species are explicitly Pirsa: 101260llowed; investigation in progress.

Solution of Boltzmann equations

Chemical potentials for t_L , \overline{t}_R , b_L , h near bubble wall:



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 $\begin{aligned} & 3v_{w}K_{\mathbf{l},i}\mu_{t,2}' + 3v_{w}K_{2,i}(m_{t}^{2})'\mu_{t,2} + 3u_{t,2}'\\ & -3\Gamma_{y}(\mu_{t,2} + \mu_{te,2} + \mu_{h,2}) - 6\Gamma_{w}(\mu_{t,2} + \mu_{te,2}) - 3\Gamma_{W}(\mu_{t,2} - \mu_{h,2})\\ & -3\Gamma_{ss}[(1 + 9K_{\mathbf{l},i})\mu_{t,2} + (1 + 9K_{\mathbf{l},5})\mu_{b,2} + (1 - 9K_{\mathbf{l},i})\mu_{te,2}] = 0\\ & 3v_{w}K_{\mathbf{l},b}\mu_{b,2}' + 3u_{b,2}'\\ & -3\Gamma_{s}(\mu_{h,2} + \mu_{te,2} + \mu_{h,2}) - 3\Gamma_{W}(\mu_{h,2} - \mu_{1,2}) \end{aligned}$

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 $\begin{aligned} & 3v_{\rm w}K_{1,t}\mu_{t^{\rm e},2}' + 3v_{\rm w}K_{2,t}(m_t^2)'\mu_{t^{\rm e},2} + 3u_{t^{\rm e},2}' \\ & -3\Gamma_y(\mu_{t,2} + \mu_{b,2} + 2\mu_{t^{\rm e},2} + 2\mu_{h,2}) - 6\Gamma_m(\mu_{t,2} + \mu_{t^{\rm e},2}) \\ & -3\Gamma_{ss}[(1 + 9K_{1,t})\mu_{t,2} + (1 + 9K_{1,b})\mu_{b,2} + (1 - 9K_{1,t})\mu_{t^{\rm e},2}] = 0 \end{aligned}$

 $4v_w K_{1,5}\mu'_{h,2} + 4u'_{h,2} \\ -3\Gamma_y(\mu_{t,2} + \mu_{b,2} + 2\mu_{t',2} + 2\mu_{h,2}) - 4\Gamma_h\mu_{h,2} = 0$

 $\begin{aligned} -3K_{4,t}\mu'_{t,2} + 3v_{\mathbf{w}}\tilde{K}_{5,t}u'_{t,2} + 3v_{\mathbf{w}}\tilde{K}_{6,t}(m_{t}^{2})'u_{t,2} + 3\Gamma_{t}^{\mathrm{tot}}u_{t,2} &= S_{t} \\ -3K_{4,b}\mu'_{b,2} + 3v_{\mathbf{w}}\tilde{K}_{5,b}u'_{b,2} + 3\Gamma_{b}^{\mathrm{tot}}u_{b,2} &= 0 \\ -3K_{4,t}\mu'_{t^{c},2} + 3v_{\mathbf{w}}\tilde{K}_{5,t}u'_{t^{c},2} + 3v_{\mathbf{w}}\tilde{K}_{6,t}(m_{t}^{2})'u_{t^{c},2} + 3\Gamma_{b}^{\mathrm{tot}}u_{t^{c},2} &= S_{t} \\ -4K_{4,b}\mu'_{b,2} + 4v_{\mathbf{w}}\tilde{K}_{5,b}u'_{b,2} + 4\Gamma_{b}^{\mathrm{tot}}u_{b,2} &= 0. \end{aligned}$

source term, $\theta = \operatorname{Im}(m_t)$: $S_t = -v_w K_\theta (m_t^2 \theta'_t)' + v_w K_\theta \theta'_t m_t^2 (m_t^2)'.$

To simplify network, FHS follow t_L , t_R , b_L , h (assuming all Higgses have same asymmetry).

Results depend upon how many species are explicitly Pirse: 1012 foollowed; investigation in progress.

Solution of Boltzmann equations

Chemical potentials for t_L , \overline{t}_R , b_L , h near bubble wall:



Summary

- MFV 2HDM gives good account of new CP violation indicated in B mixing from DØ dimuons and deviations in $B_s \rightarrow J/\psi \phi$ and $B^- \rightarrow \tau \nu$.
- Same model can give baryogenesis; highly constrained.
- Unfortunately new phase in *B* mixing is not the one responsible for baryogenesis; model has 6 new *CP* phases.
- η_U should be as large as possible; more sensitive measurements of R_b might see a deviation
- Predict light Higgs and new Higgses $m_{S_I,S_R} \lesssim 450$ GeV, $m_{\pm} \lesssim 300$ GeV; these may become sharper
- 4b or 2b-2t discovery channels possible for LHC