

Title: Effects of decoherence on IR divergence

Date: Oct 30, 2010 11:30 AM

URL: <http://pirsa.org/10100082>

Abstract: We propose one way to regularize the fluctuations generated during inflation. We show that, as long as we consider the case that the non-linear interaction acts for a finite duration, observable fluctuations are free from IR divergences not only in the single field models but also in the multi field model. In contrast to the single field model, to discuss observables, we need to take into account the effects of quantum decoherence which pick up a unique history of the universe from various possibilities contained in initial quantum state set naturally in the early stage of the universe.

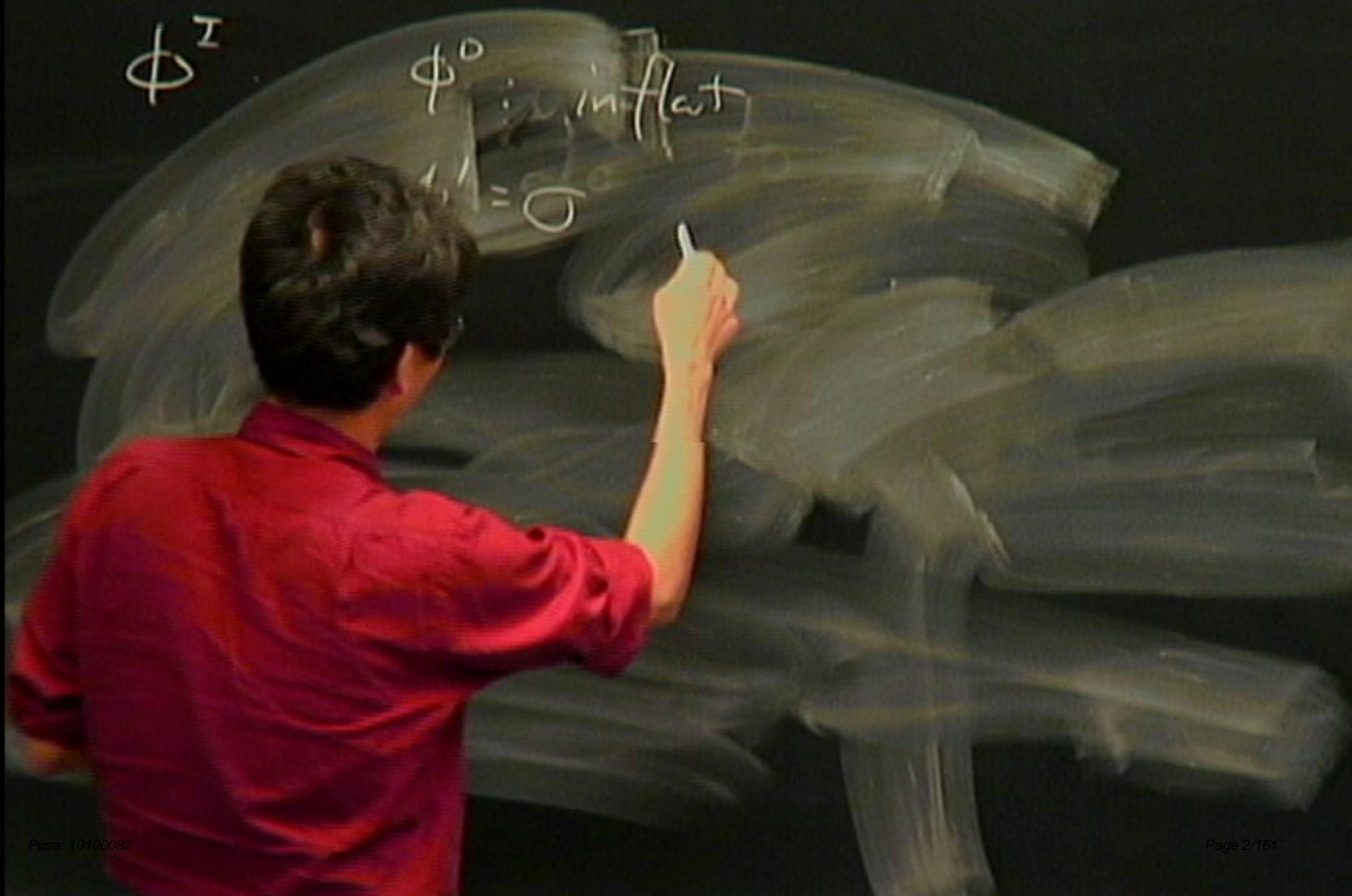
0904.8415

with Yuko.Urakawa

ϕ^I

ϕ^0

: inflat
= 0



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with Yuko. Urakawa

ϕ^I

ϕ^0

: inflat

$\phi^1 = \phi^0$

is curvature.

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with Yuko, Urakawa

ϕ^I

ϕ^0 : inflat

$\phi^1 = \sigma$ is curvature



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with Yuko. Urakawa

ϕ^I

ϕ^0 : inflat

$\phi^I = 0$

is curvature

$$\langle \sigma^2 \rangle \approx \int_{k_{min}}^{k_{out}}$$

$$k_{out} = \epsilon Q H_1$$

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with Yuko, Urakawa

ϕ^I

ϕ^0 : inflat

$\phi^1 = \sigma$

is curvature

$$\langle \sigma^2 \rangle \approx \int_{k_{\min}}^{k_{\text{cut}}}$$

$$k_{\text{cut}} = \epsilon_Q H$$

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with Yuko Urakawa

ϕ^I

ϕ^0 is inflat

$\phi^1 = 0$

nature

$$\langle \sigma^2 \rangle \approx \int_{k_{\min}}^{k_{\text{cut}}} d$$

k

large $k_{\min} \rightarrow 0$
 ∞ on
 $t \rightarrow \infty$

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ϕ^I

ϕ^0 is inflat

$\phi^1 = 0$ iso curvature

$$\langle \sigma^2 \rangle \approx \int_{k_{\min}}^{k_{\text{cut}}} d^3 k P_0(k)$$

$$k_{\text{cut}} = \epsilon Q H$$

large $k_{\min} \rightarrow 0$
 ∞ or
 $t \rightarrow \infty$

0904.8415

with Yuko Urakawa

ϕ^I

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$$\langle \sigma^2 \rangle \approx \int_{k_{\min}}^{k_{\text{cut}}} d^3 k P_0(k)$$

$$k_{\text{cut}} = \epsilon Q H$$

large $k_{\min} \rightarrow 0$
 ∞ on
 $t \rightarrow \infty$

Stock

$$k_{out} = CQH$$

view

$$\bar{J} = \left(\begin{array}{c} \end{array} \right)$$

$k_{out} = \epsilon Q H$

Stochastic view

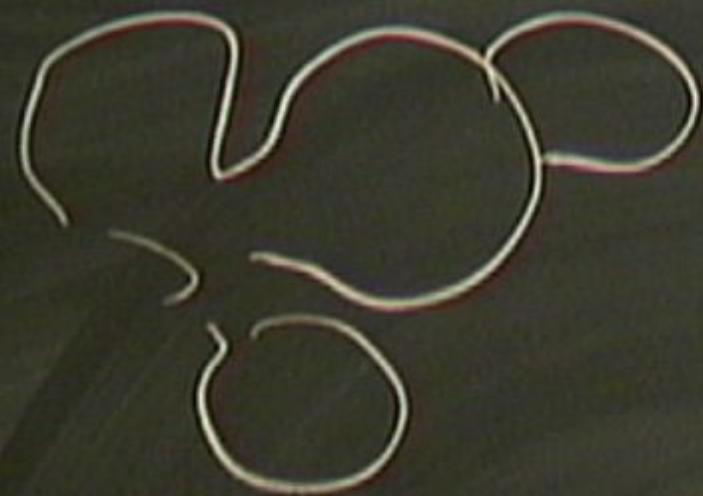
$$\bar{J} = \int d^3k$$

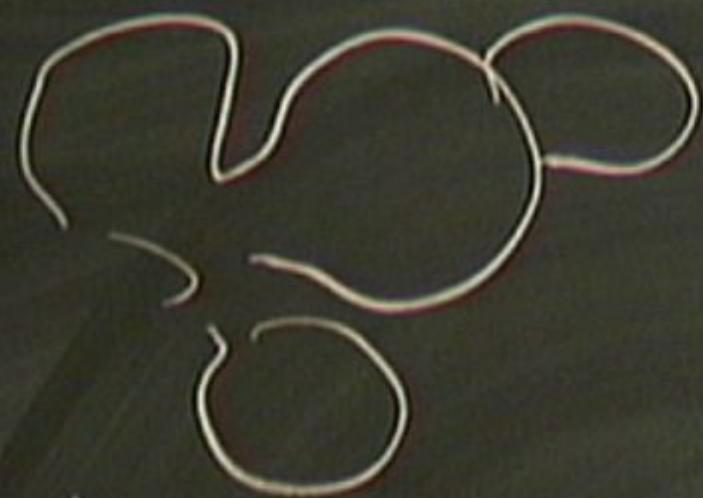
$$\langle \sigma^2 \rangle \approx \int_{k_{\min}}^{\infty} d^3 k P_0(k) \quad \begin{array}{l} \text{large } k_{\min} \rightarrow 0 \\ \text{or} \\ t \rightarrow \infty \end{array}$$

$$k_{\text{out}} = \epsilon Q H$$

Stochastic view

$$\bar{\sigma} = \int d^3 k \delta(k_{\text{out}} - k) \sigma_k e^{ik \cdot r}$$





k_{min}



k_{\min}

$$\frac{k_{\min} \rightarrow 0}{}$$

k_{\min} -depen



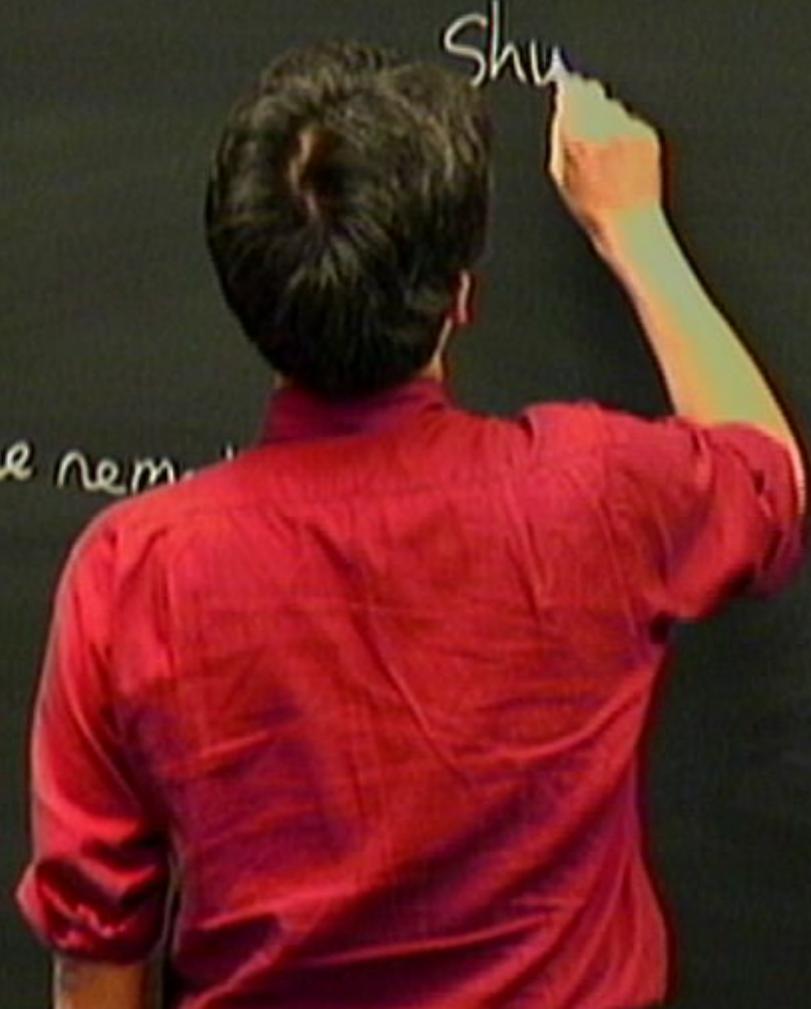
k_{\min}

$\underline{k_{\min} \rightarrow 0}$

k_{\min} -dependence removed

"True observable"

Shu





"True observable"

Should not depend on k_{\min}

k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains



k_{\min}

$\frac{k_{\min}}{h}$

ice remains

"True observable"

Should not depend on k_{\min}

$\langle \sigma_{V_0} \sigma_{H_2} \rangle$



"True observable"

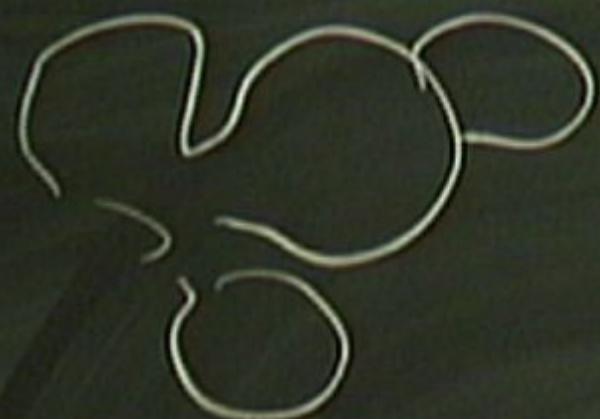
Should not depend on k_{\min}

k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains

$\langle \sigma_{lk} \sigma_{lh} \rangle$



"True observable"

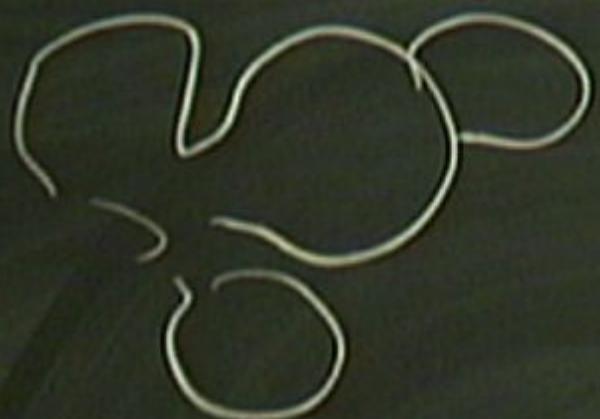
Should not depend on k_{\min}

k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains

$\langle \sigma_{ik} \sigma_{ihz} \rangle$



k_{\min}

$\frac{k_{\min} \rightarrow 0}{}$

k_{\min} -dependence remains

"True observable"

Should not depend on k_{\min}

$\langle \underline{\sigma}_{\text{LR}} \underline{\sigma}_{\text{Hz}} \rangle$



k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains

"True observable"

Should not depend on k_{\min}

$\langle \underline{\sigma}_{ik} \underline{\sigma}_{hk} \rangle$



"True observable"

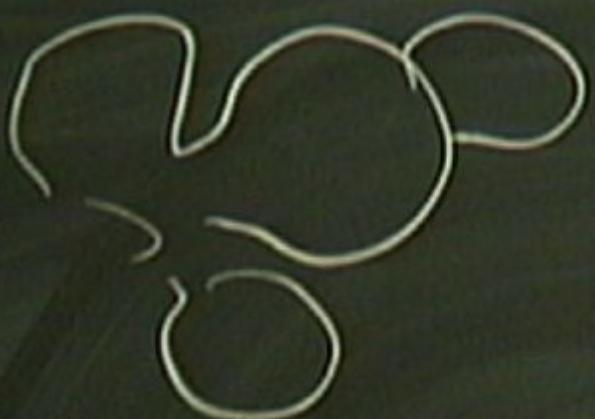
Should not depend on k_{\min}

k_{\min}

$\underline{k_{\min}} \rightarrow 0$

k_{\min} -dependence remains

$\langle \underline{\sigma}_{\text{LR}} \underline{\sigma}_{\text{LH}} \rangle$



"True observable"

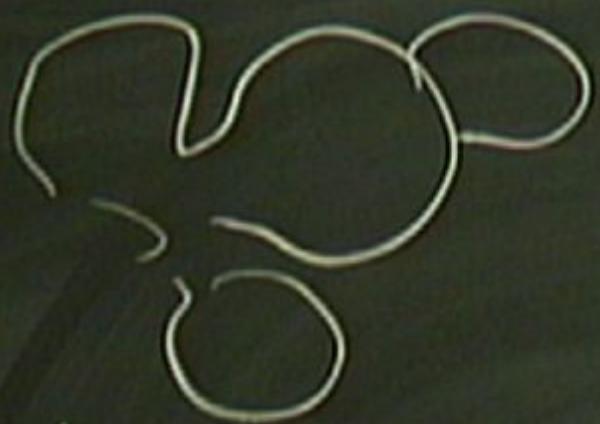
Should not depend on k_{\min}

k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains

$\langle \underline{\sigma}_{\text{LR}} \underline{\sigma}_{\text{Hz}} \rangle$



k_{\min}

$k_{\min} \rightarrow 0$

k_{\min} -dependence remains

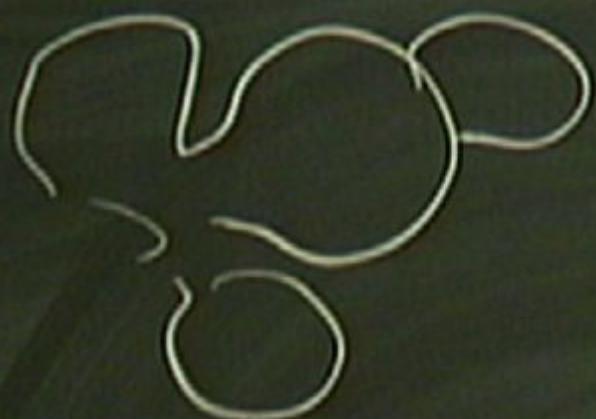
"True observable"

Should not depend on k_{\min}

$\langle \sigma_{ik} \sigma_{ij} \rangle$

$\langle \sigma(x) \sigma(y) \rangle$

$x, y \subset \odot$



k_{\min}

$\frac{k_{\min} \rightarrow 0}{k_{\min}}$

k_{\min} -dependence remains

"True observable"

Should not depend on k_{\min}

$\langle \sigma_{\text{R}} \sigma_{\text{Lz}} \rangle$

$\langle \sigma(x) \sigma(y) \rangle$

$x, y \subset \odot$



k_{\min}

$\underline{k_{\min}} \rightarrow 0$

k_{\min} -dependence remains

"True observable"

Should not depend on k_{\min}

$\langle \sigma_{ik} \sigma_{ik} \rangle$

$\langle \sigma(x) \sigma(y) \rangle$

$x, y \subset \odot$
can be divergent



k_{\min} $\frac{k_{\min} \rightarrow 0}{}$

k_{\min} -dependence remains

$$\zeta = \zeta(\delta\phi^x)$$

"True observable"

Should not depend on k_{\min}

$$\langle \sigma_{ik} \sigma_{ih} \rangle$$

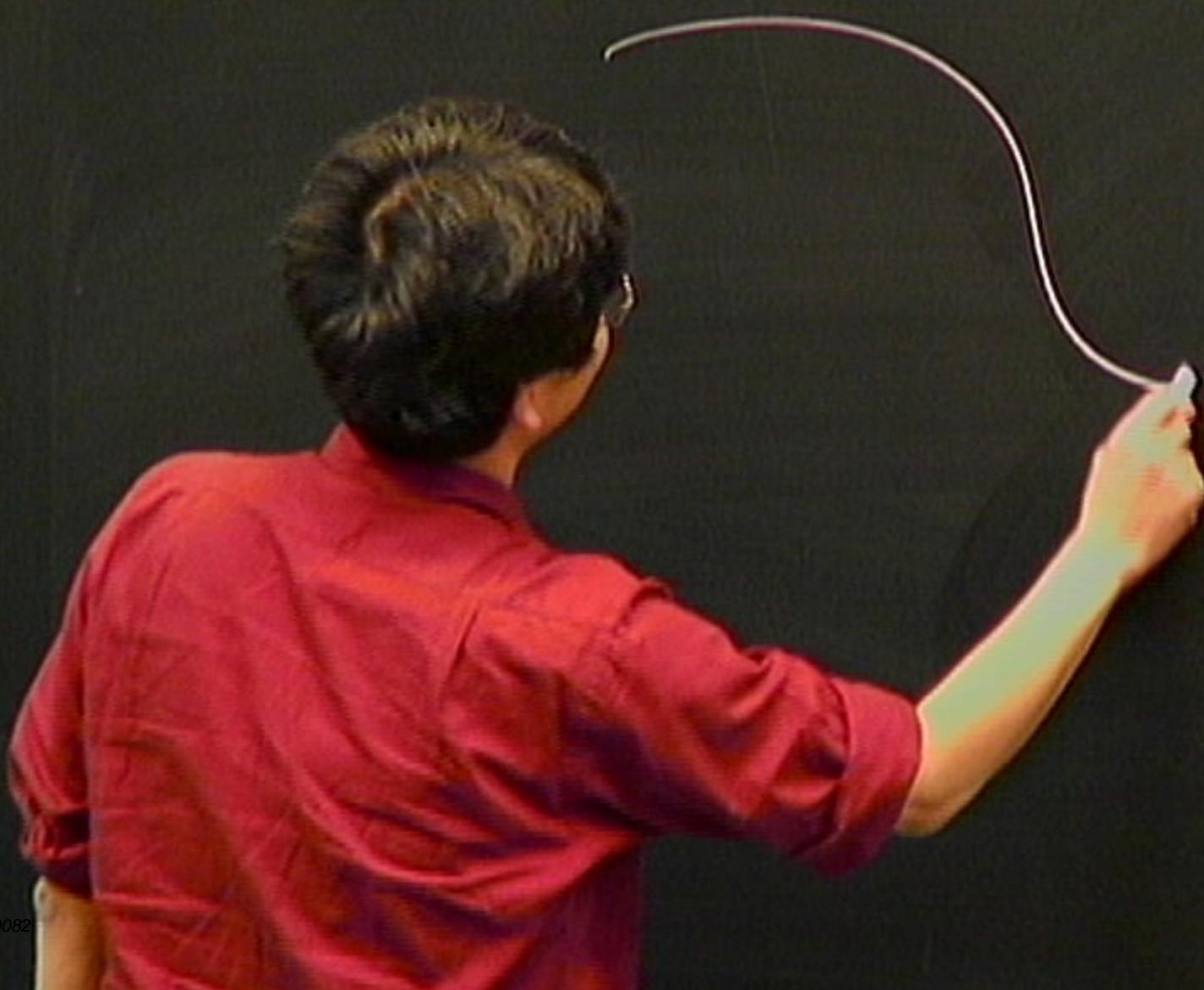
$$\langle \sigma(x) \sigma(y) \rangle$$

can be $x, y \subset \odot$
divergent

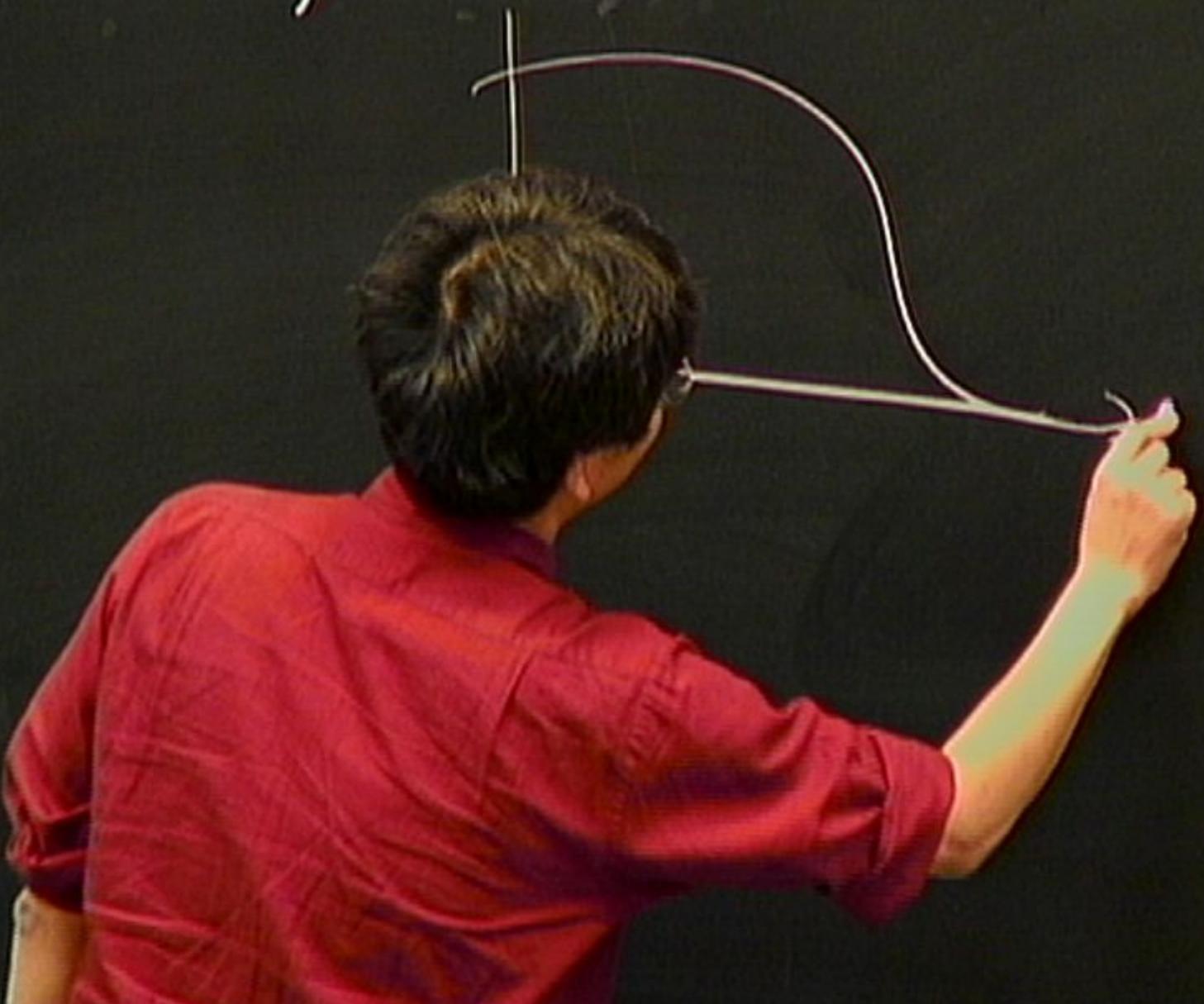
$$\bar{\sigma} = \int d^3x W(x)$$

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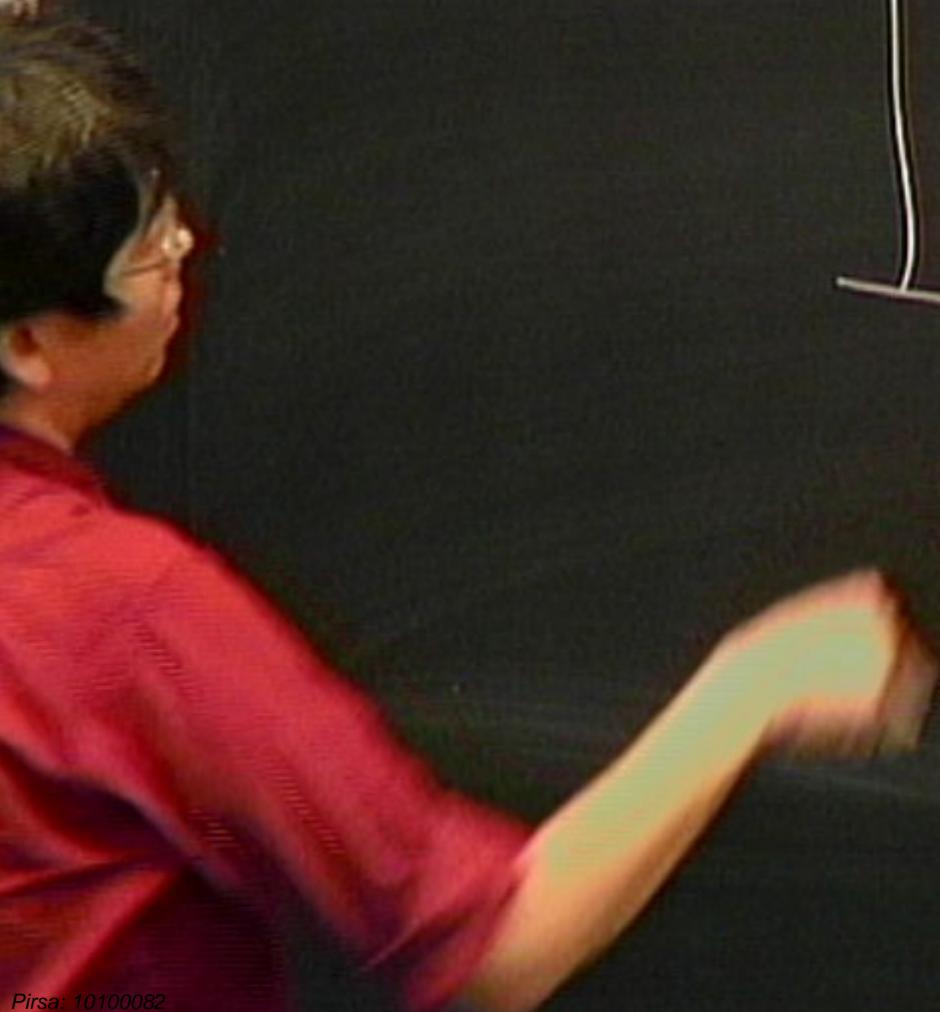
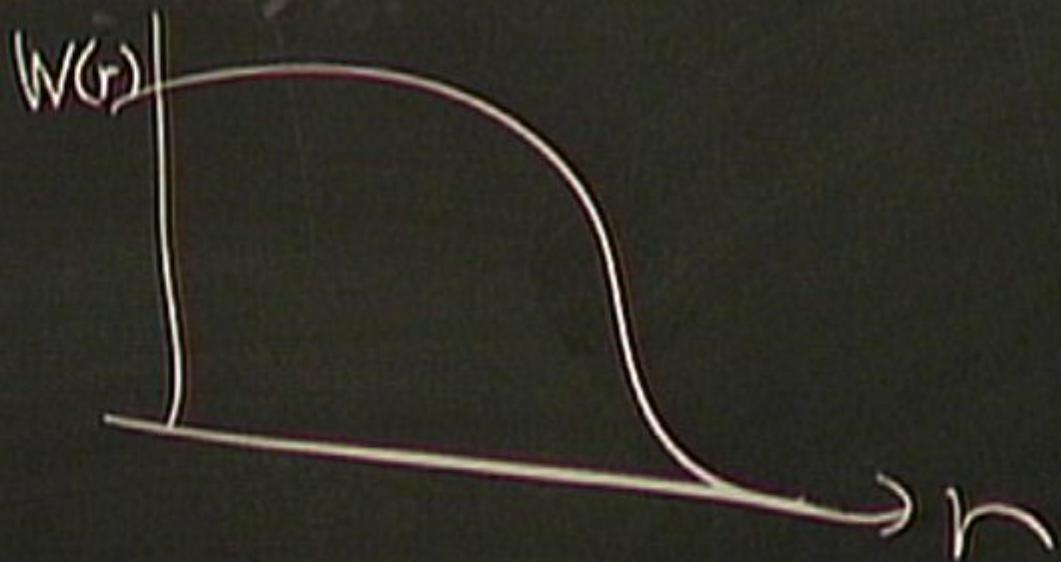
$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$



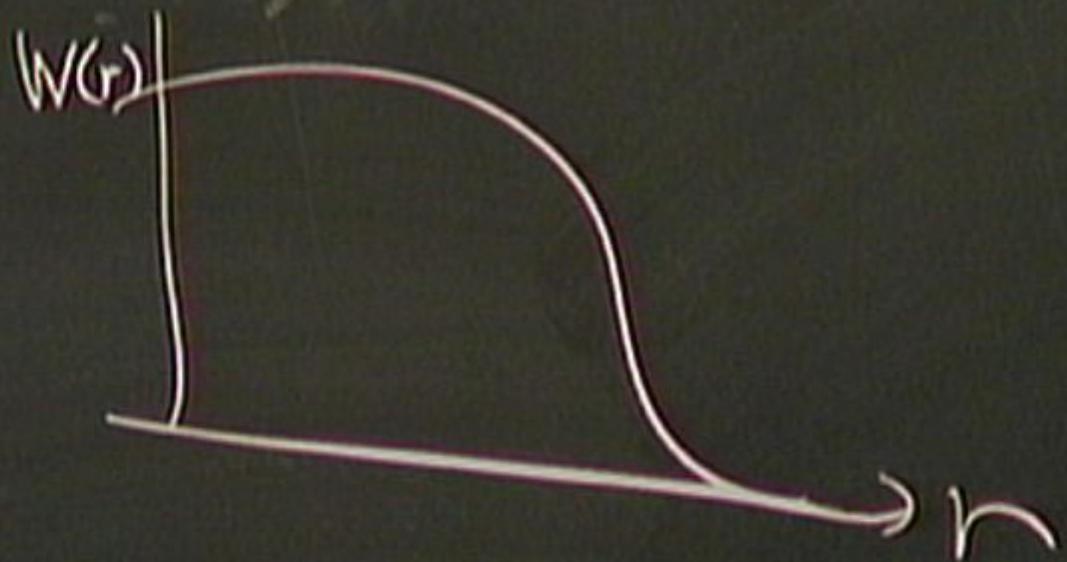
$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$



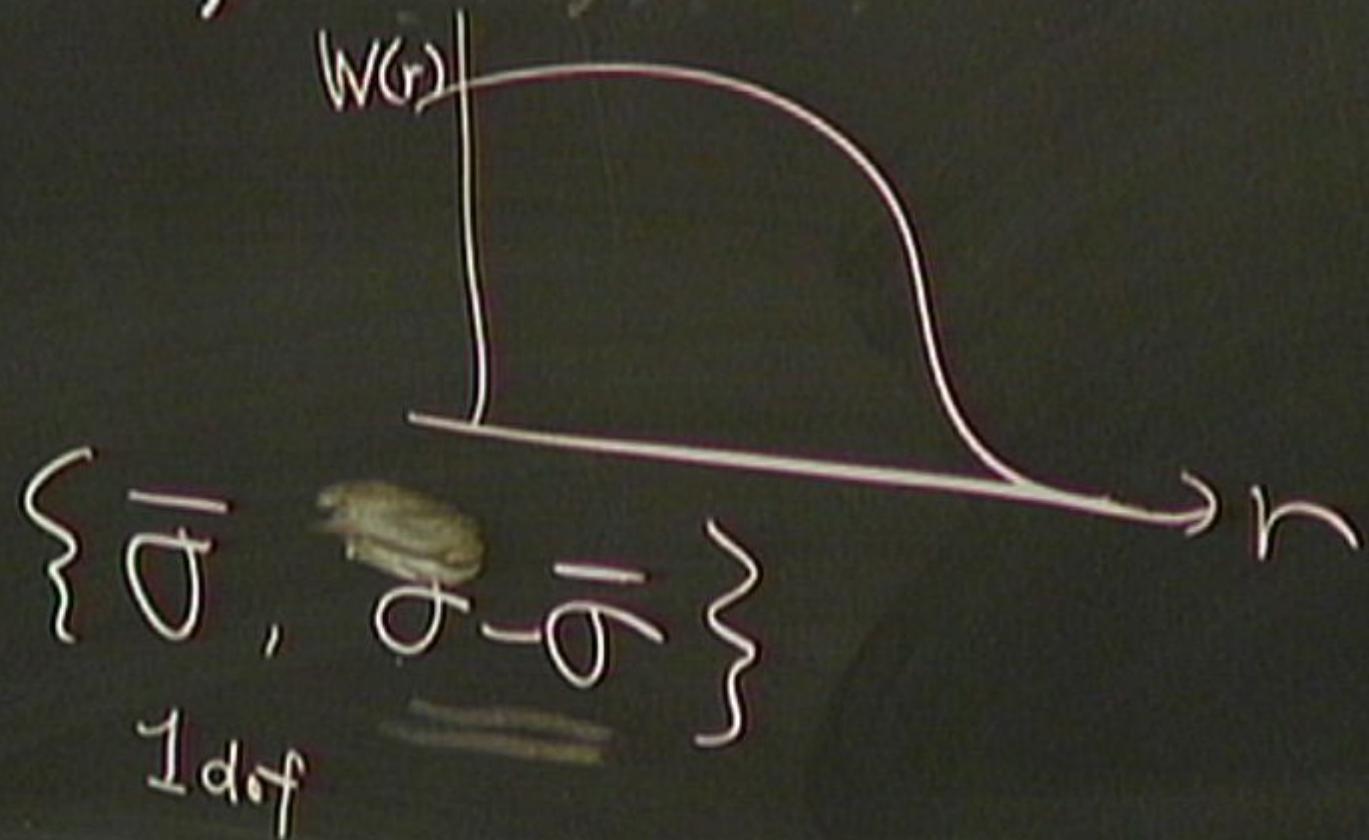
$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$



$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$

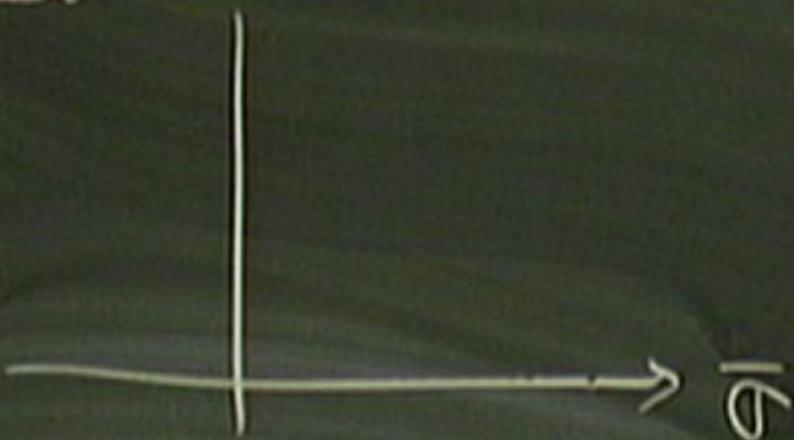


$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$



Do we really observe large fluctuation of $\bar{\sigma}$

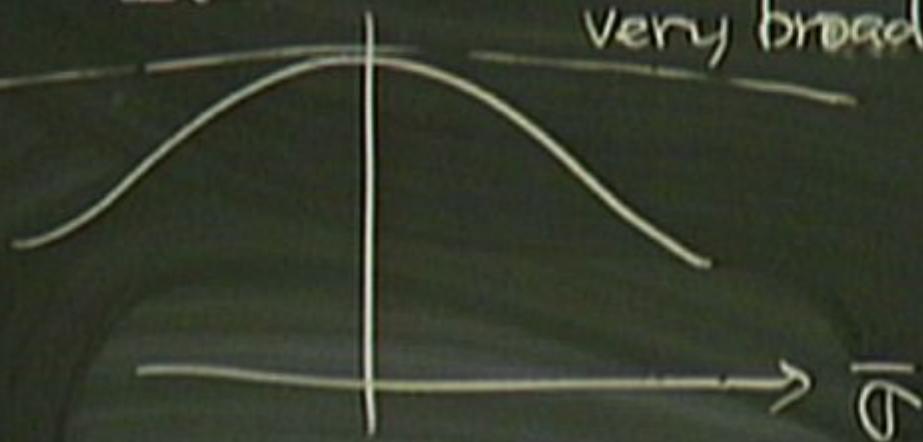
$$\bar{\Psi}(\bar{\sigma})$$



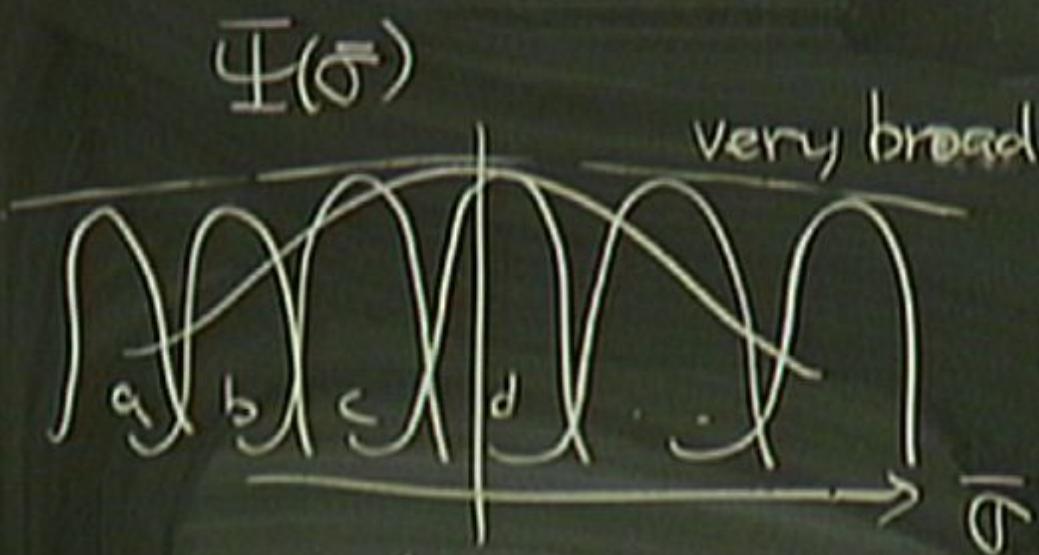
Do we really observe large fluctuation of $\bar{\sigma}$

$$\Phi(\bar{\sigma})$$

very broad.



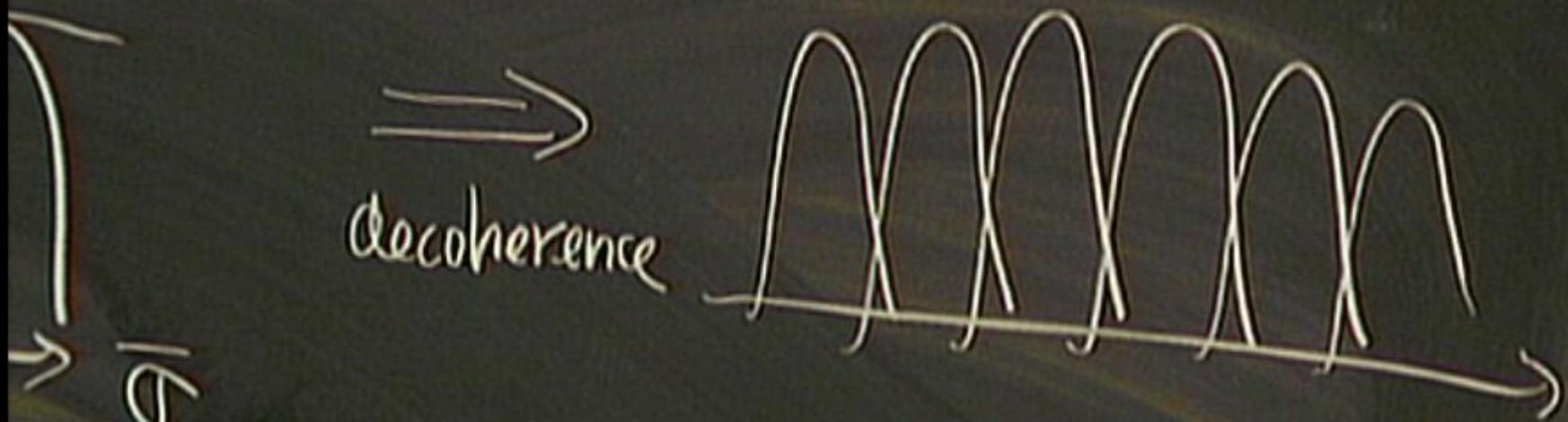
Do we really observe large fluctuation of $\bar{\sigma}$



$$P = \langle a > + \langle b > + \langle c > + \dots \rangle (\langle a | + \langle b | + \langle c | + \dots)$$

serve large fluctuation of $\bar{\sigma}$

broad.



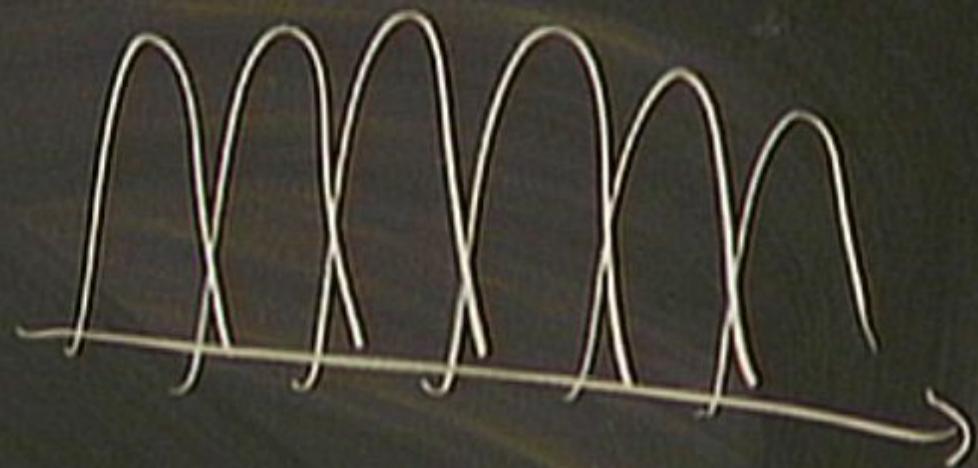
$$\dots (|a\rangle\langle a| + |b\rangle\langle b| + |c\rangle\langle c| + \dots) \Rightarrow \hat{\rho} = |a\rangle\langle a| + \boxed{|b\rangle\langle b|} + \dots$$

serve large fluctuation of $\bar{\sigma}$

broad.



\Rightarrow
Decoherence

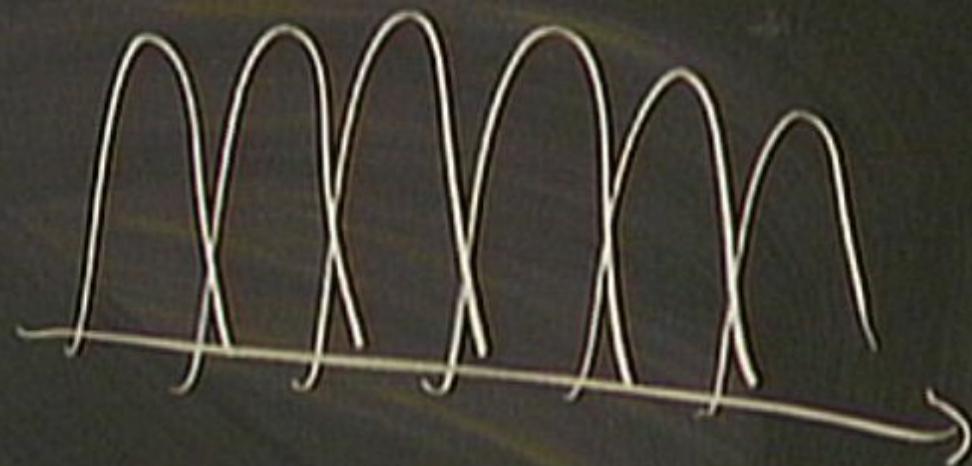


$$\dots (|a\rangle\langle a| + |b\rangle\langle b| + |c\rangle\langle c| + \dots) \Rightarrow \hat{\rho} = |a\rangle\langle a| + \boxed{|b\rangle\langle b|} + \dots$$

serve large fluctuation of $\bar{\sigma}$

broad.

\Rightarrow
decoherence

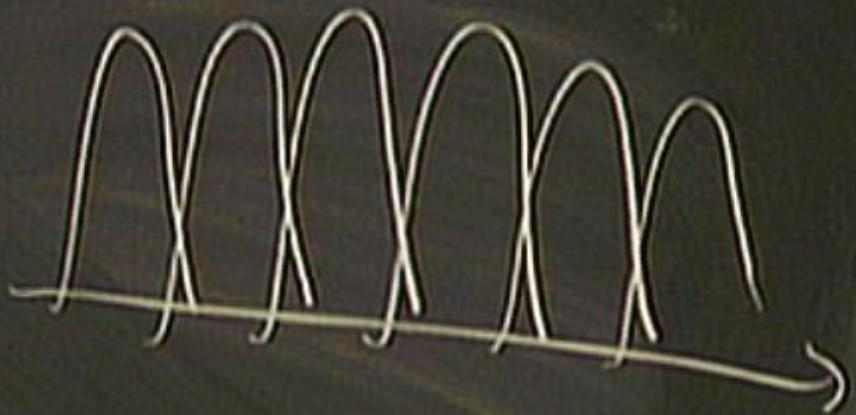


$$\dots (|a\rangle\langle a| + |b\rangle\langle b| + |c\rangle\langle c| + \dots) \Rightarrow \hat{\rho} = |a\rangle\langle a| + |b\rangle\langle b| + \dots$$

really observe large fluctuation of $\bar{\sigma}$

very broad.

\Rightarrow
Decoherence

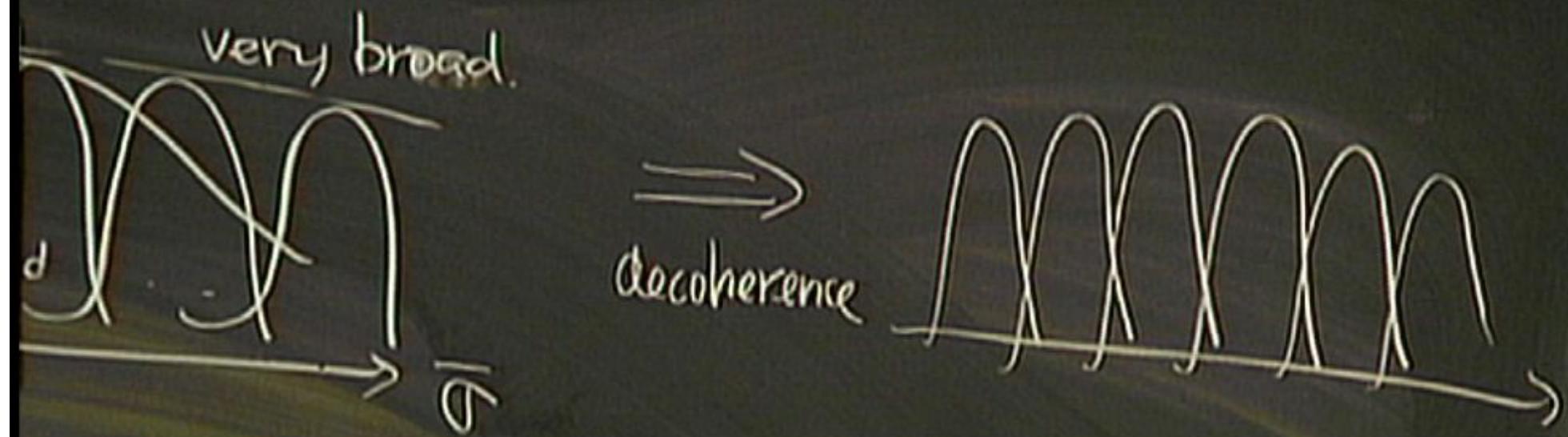


$|b\rangle + |c\rangle$

$$(\langle a| + \langle b| + \langle c| + \dots) \Rightarrow \rho = |\alpha\rangle\langle\alpha| + |\beta\rangle\langle\beta| + \dots$$

what we observe.

really observe large fluctuation of $\bar{\sigma}$

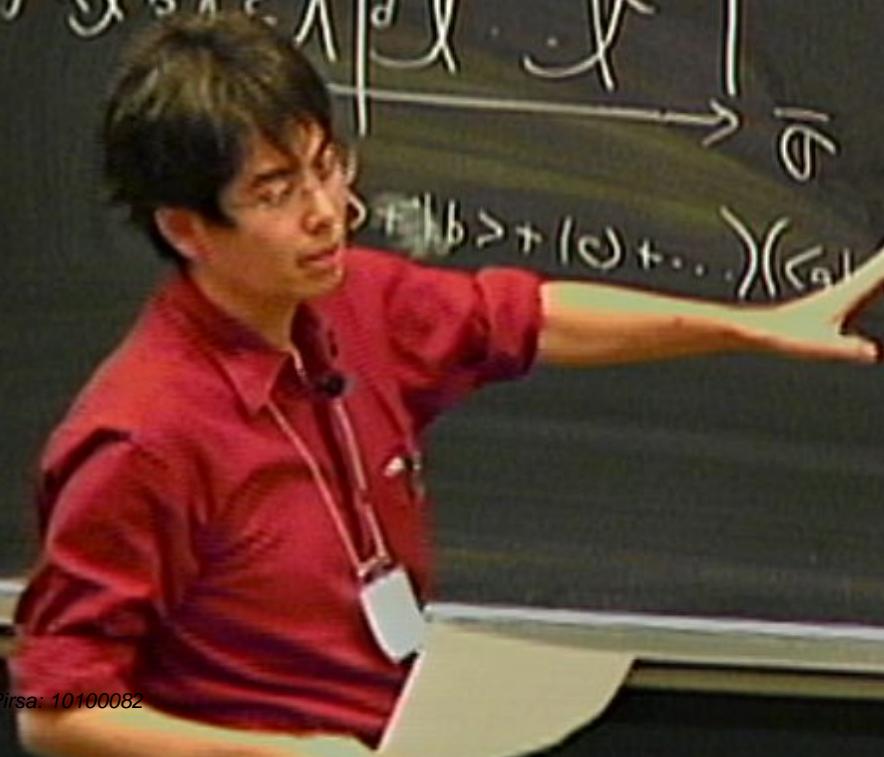
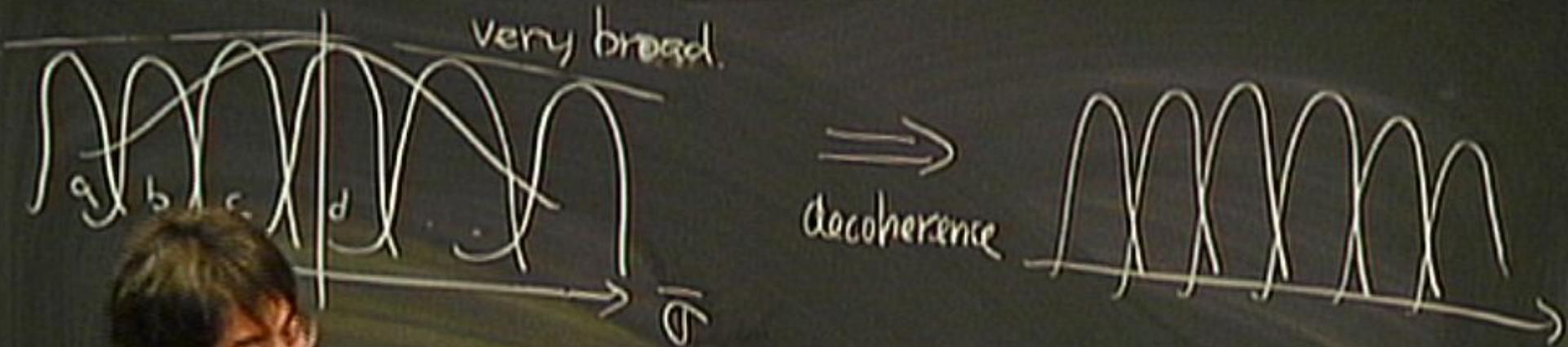


$$|\psi\rangle = |a\rangle + |b\rangle + |c\rangle + \dots \quad (\langle a| + \langle b| + \langle c| + \dots) \Rightarrow \rho = |\alpha\rangle\langle\alpha| + |\beta\rangle\langle\beta| + \dots$$

$\rho = |\alpha\rangle\langle\alpha| + |\beta\rangle\langle\beta| + \dots$

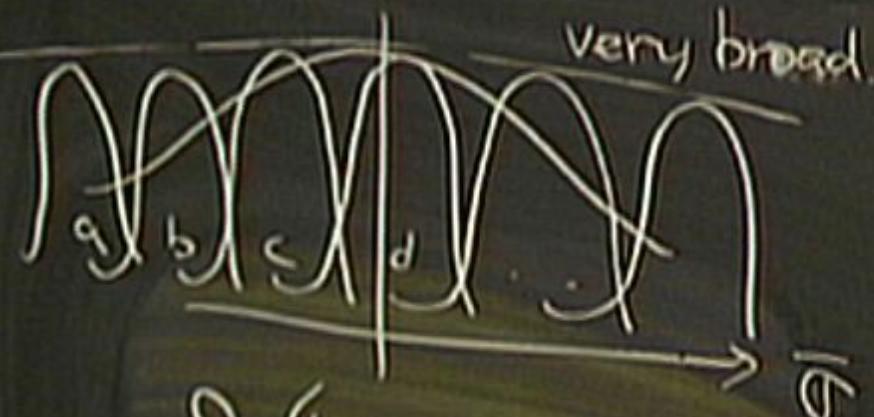
what we observe.

Do we really observe large fluctuation of $\bar{\sigma}$
 $\Psi(\bar{\sigma})$

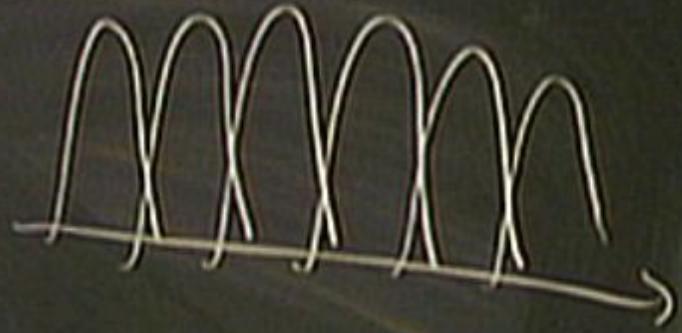


Do we really observe large fluctuation of $\bar{\sigma}$

$\Psi(\bar{\sigma})$



\Rightarrow
Decoherence



$$\rho = (|a\rangle + |b\rangle + |c\rangle + \dots)(\langle a| + \langle b| + \langle c| + \dots) \Rightarrow \rho = |a\rangle \langle a| + \boxed{|b\rangle \langle b|} + \dots$$

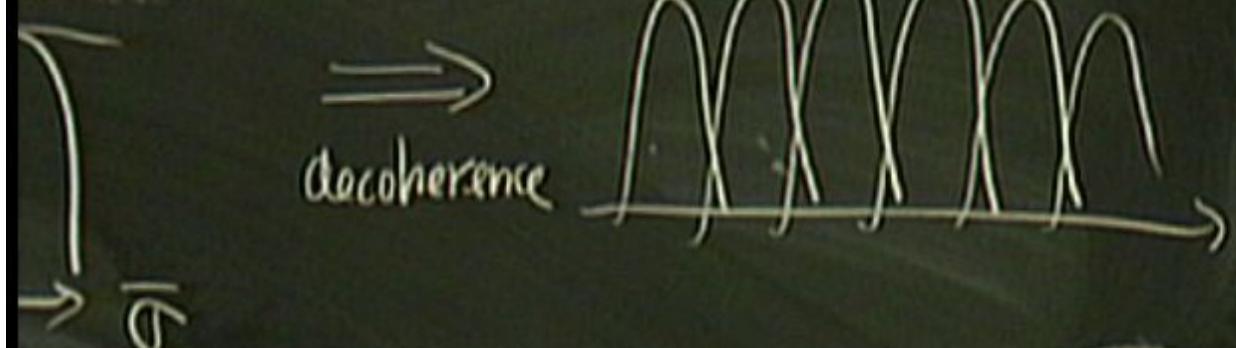
what we observe.

$$(a\dot{\varphi}_k)' + \left(k^2 - \frac{a''}{a} \right) (a\dot{\varphi}_k) = 0$$

serve large fluctuation of $\bar{\sigma}$

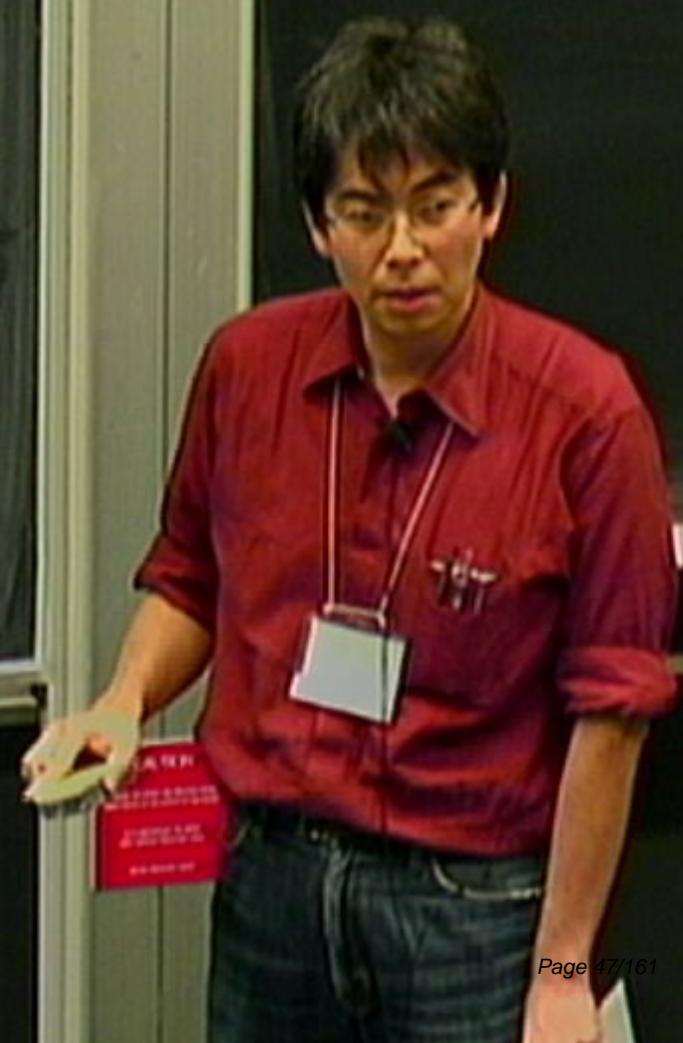
(a)

broad.



$$\dots (|a\rangle\langle a| + |b\rangle\langle b| + |c\rangle\langle c| + \dots) \Rightarrow \rho = |a\rangle\langle a| + |b\rangle\langle b| + \dots$$

what we observe.



Picking up one decohered wave packet is difficult.

T

start

decohered



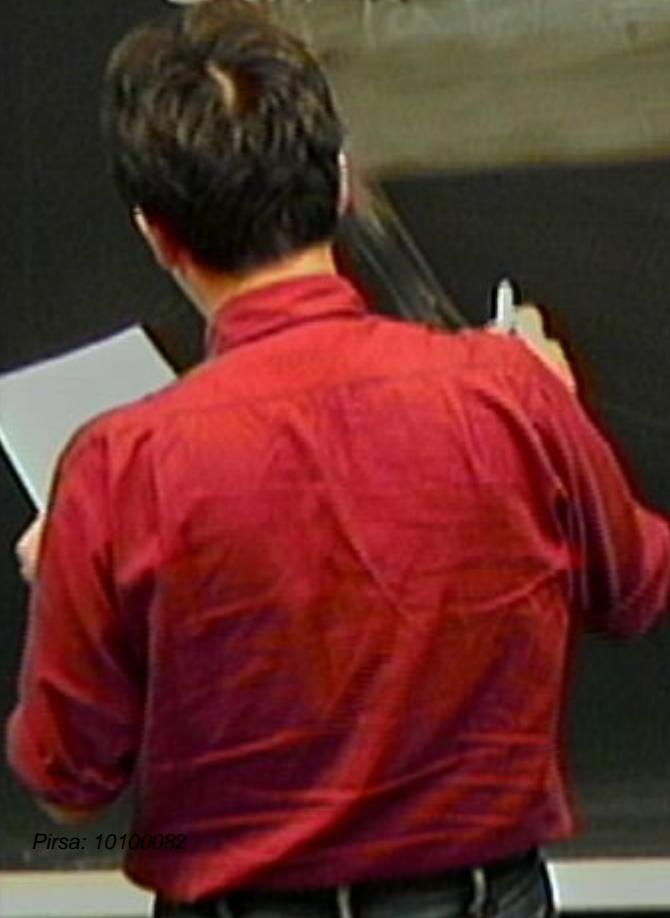
Picking up one decohered wave packet is difficult.

Instead,



Picking up one decohered wave packet is difficult.

Instead



Picking up one decohered wave packet is difficult.

Instead

$$P = \exp\left(-\frac{\sigma^2}{2(\Delta\sigma)^2}\right)$$

Picking up one decohered wave packet is difficult.

Instead

$$P = \exp\left(-\frac{\sigma^2}{2(\Delta\sigma)^2}\right)$$

$$\langle P \sigma(x) \sigma(y) \rangle$$

Picking up one decohered wave packet is difficult.

Instead

$$P = \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\bar{\sigma})^2}\right)$$

$$\underline{\langle P(\tau) \sigma(\psi) \rangle}$$

Picking up one decohered wave packet is difficult.

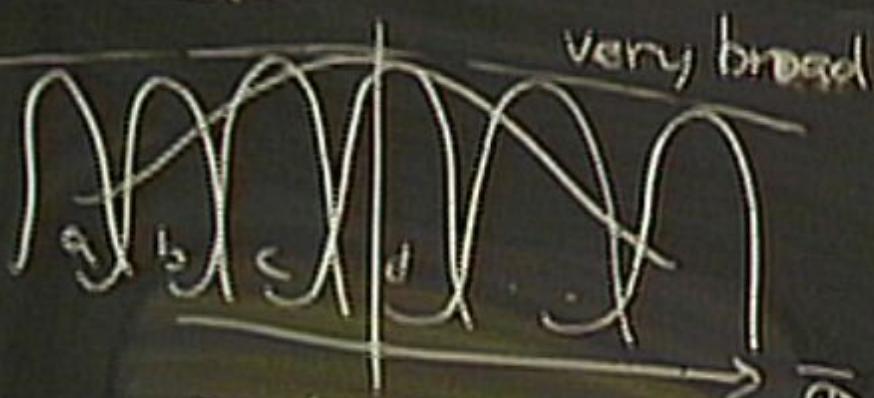
Instead

$$P = \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\bar{\sigma})^2}\right)$$

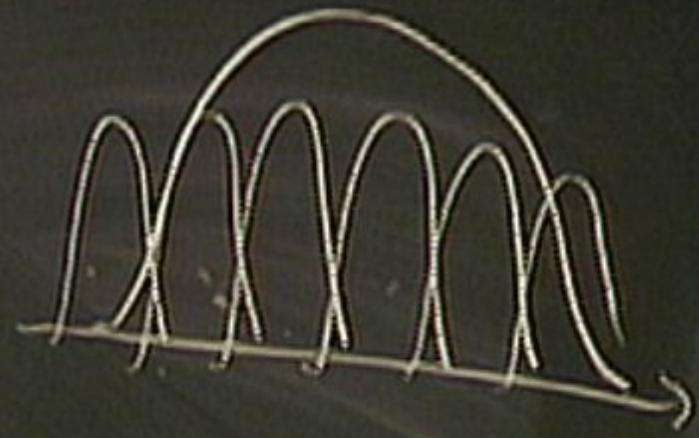
$$\frac{\langle P | \sigma(x) \sigma(y) \rangle}{\langle P \rangle}$$

Do we really observe large fluctuation of $\bar{\sigma}$

$$\Psi(\bar{\sigma})$$



\Rightarrow
decoherence



$$\rho = \sqrt{a|>}_a<_a| + b|>_b<_b| + \dots$$

$$\rho = |a><a| + \boxed{|b><b|} + \dots$$

what we observe.

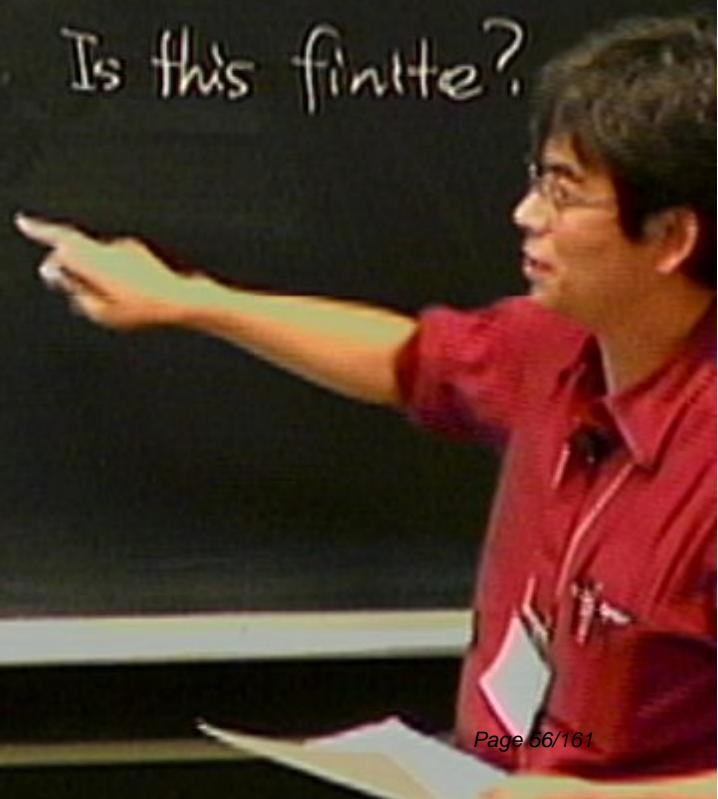
Picking up one decohered wave packet is difficult.

Instead

$$P = \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\bar{\sigma})^2}\right)$$

$$\frac{\langle P \cdot \sigma(x) \sigma(y) \rangle}{\langle P \rangle}$$

Is this finite?



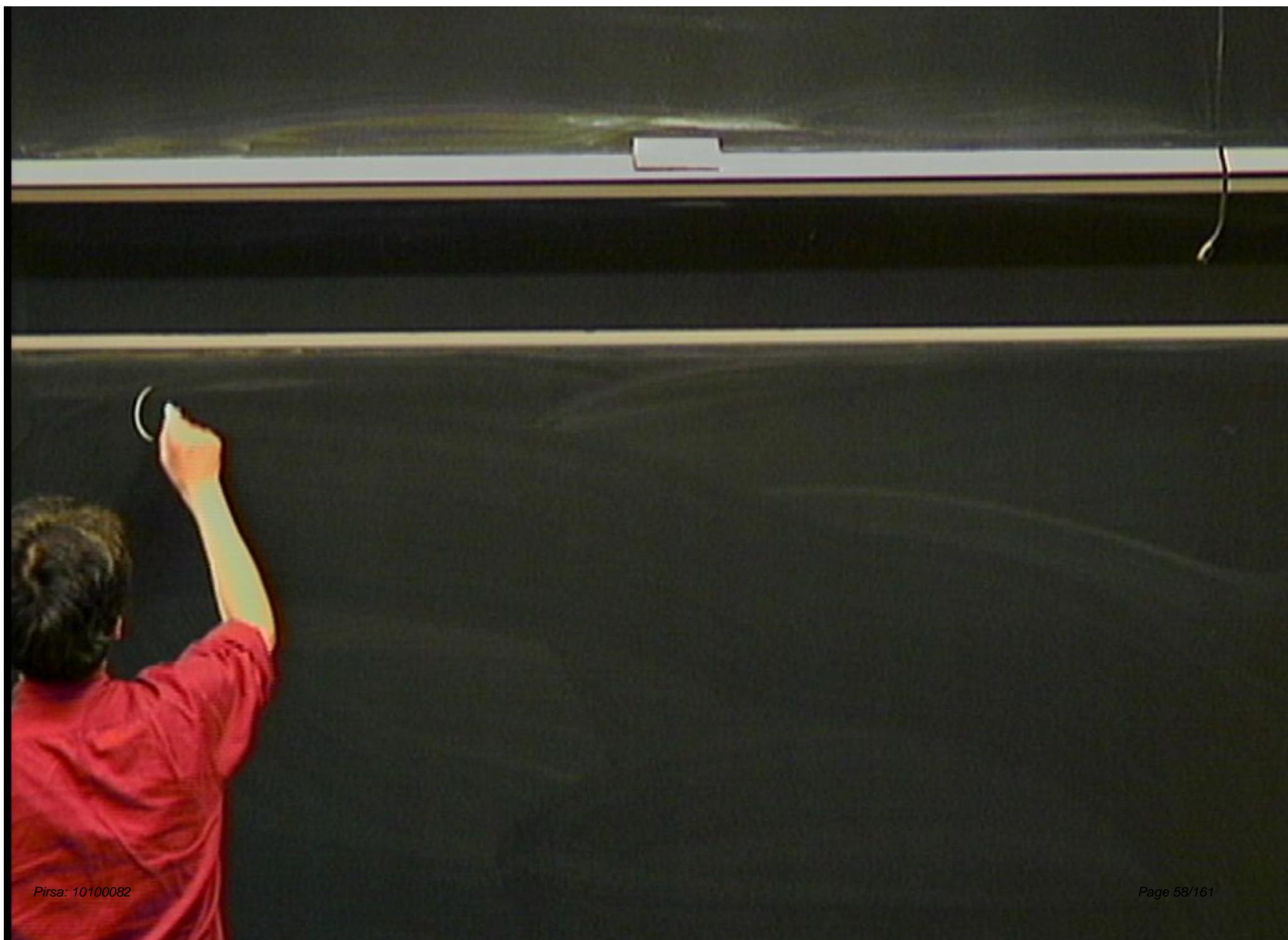
Picking up one decohered wave packet is difficult.

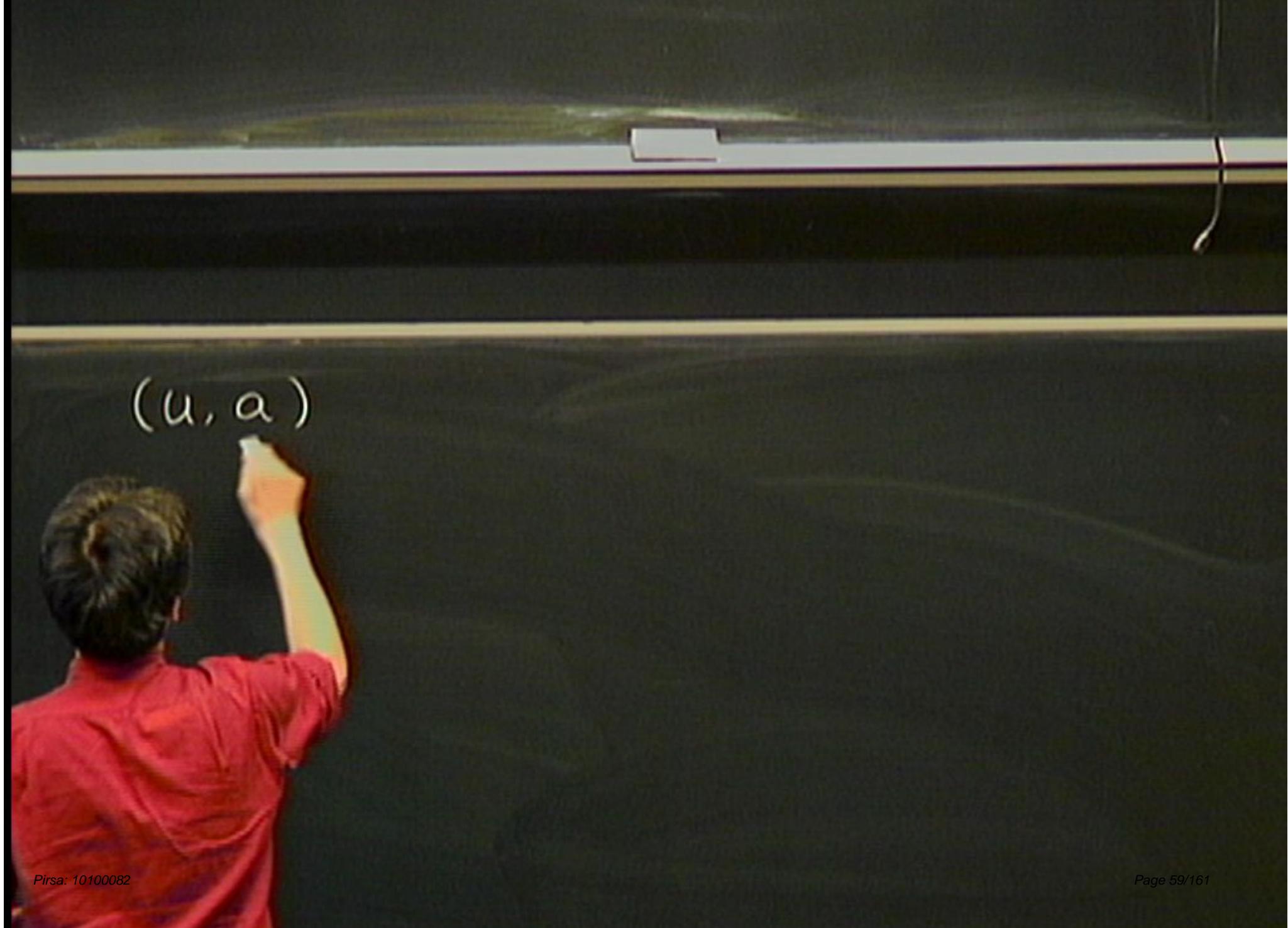
Instead

$$P = \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\bar{\sigma})^2}\right)$$

$$\frac{\langle P | \sigma(x) \sigma(y) \rangle}{\langle P \rangle}$$

Is this finite?



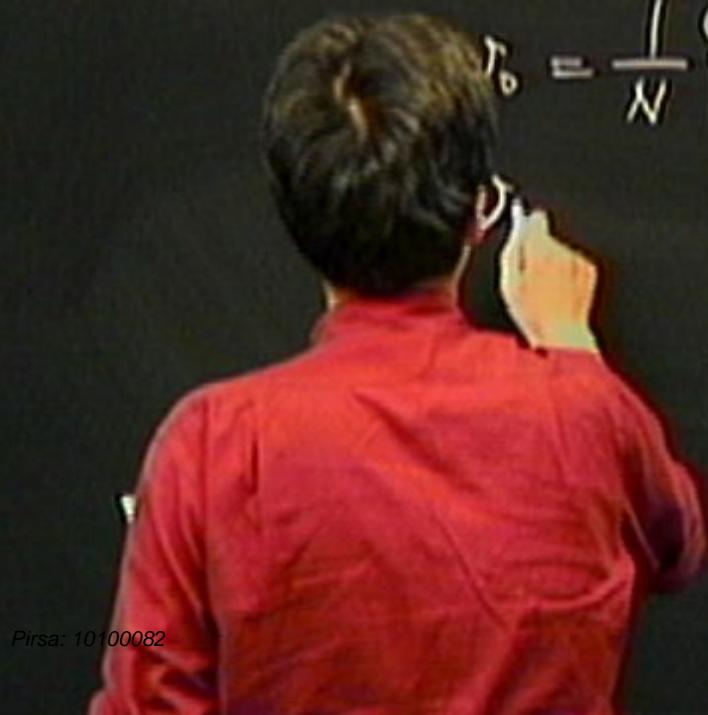
A person with dark hair tied back, wearing a red short-sleeved shirt, is standing in front of a chalkboard. They are facing away from the camera, writing with a piece of chalk. The chalkboard is dark, and the text they are writing is partially visible.

(u, α)

(u , α)

$(u, a) \xrightarrow{\text{unitary}} (v, b)$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$



$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v = \left\{ u_0 + \sum \dots \right\}$$

$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 \quad u_0 + \sum \dots \quad \}$$

$$v_p \quad p$$

$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$v_p = u_p e^{i p \cdot x} - \frac{w_p}{w_0} \frac{c(0)}{c(p)} u_0$$

$$\frac{1}{C} \sum_n u_p$$

$(u, a) \xrightarrow{\text{unitary}} (v, b)$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \lim_{n \rightarrow \infty} U_p$$

$$U_P = U_0 e^{i p \cdot x} - \frac{W_0}{W_0} \frac{C(0)}{C(p)} U_0$$

$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{c(p)} = \lim_{\eta \rightarrow 0} U_p$$

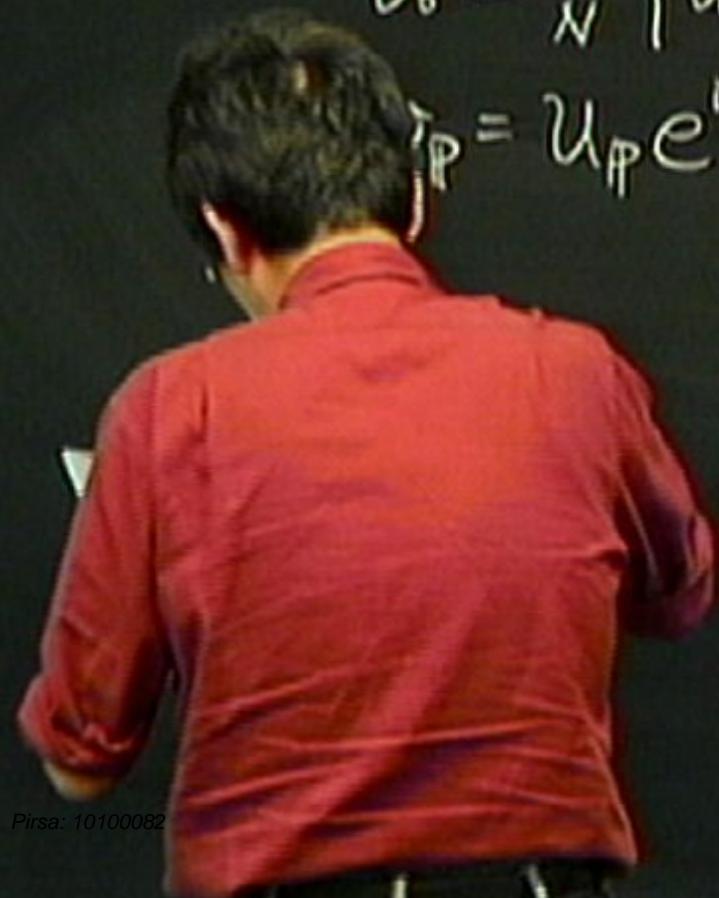
$$U_P = U_0 e^{i p \cdot x} - \frac{W_P}{W_0} \frac{C(0)}{C(P)} U_0$$

$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \lim_{\eta \rightarrow 0} U_p$$

$$j_p = U_p e^{ip \cdot x} - \frac{W_p}{W_0} \frac{C(0)}{C(p)} U_0$$



$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \lim_{\eta \rightarrow 0} U_p$$

$$U_P = U_0 e^{i P \cdot x} - \frac{W_P}{W_0} \frac{C(0)}{C(p)} U_0$$

$$\int W(x) U_P(x) d^3x = 0$$

$$(u, a) \xrightarrow{\text{unitary}} (v, b)$$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$v_P = u_P e^{i \int W(x) dx}$$

$W_0 \in \mathcal{L}(P)$

$$\int W(x) v_P(x) d^3x = D$$

v_0 contains information

$$\lim_{\eta \rightarrow 0} u_\eta$$

$(u, a) \xrightarrow{\text{unitary}} (v, b)$

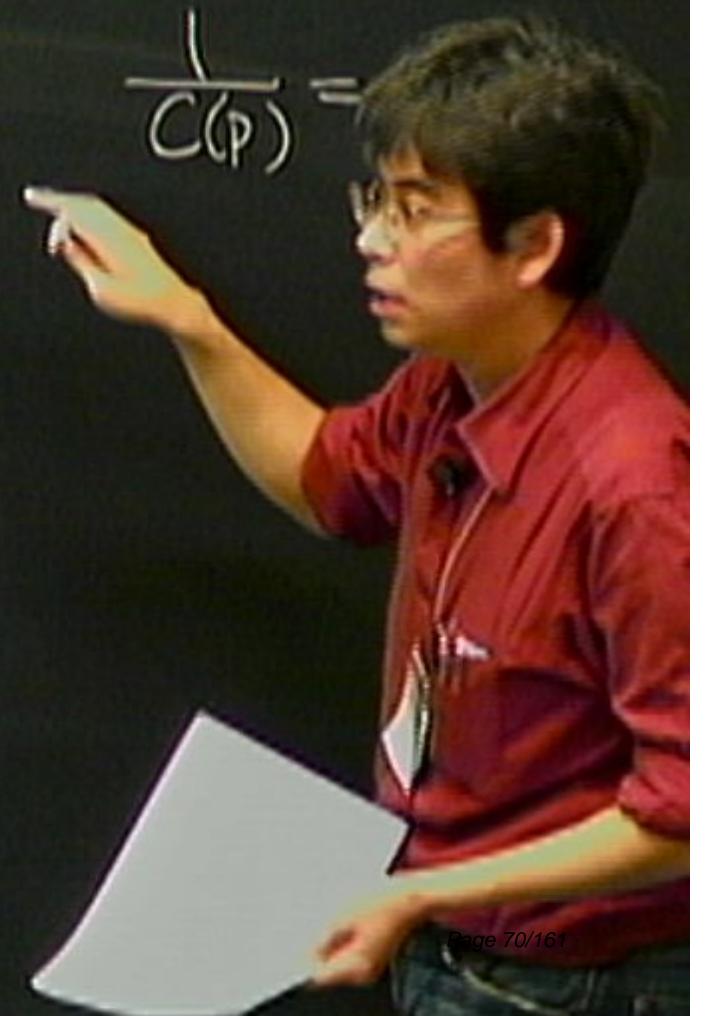
$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{C(p)} =$$

$$v_p = u_p e^{ipx} - \frac{W_p}{W_0} \frac{C(0)}{C(p)} u_0$$

$$\int W(x) v_p(x) dx = D$$

v_0 contains information about



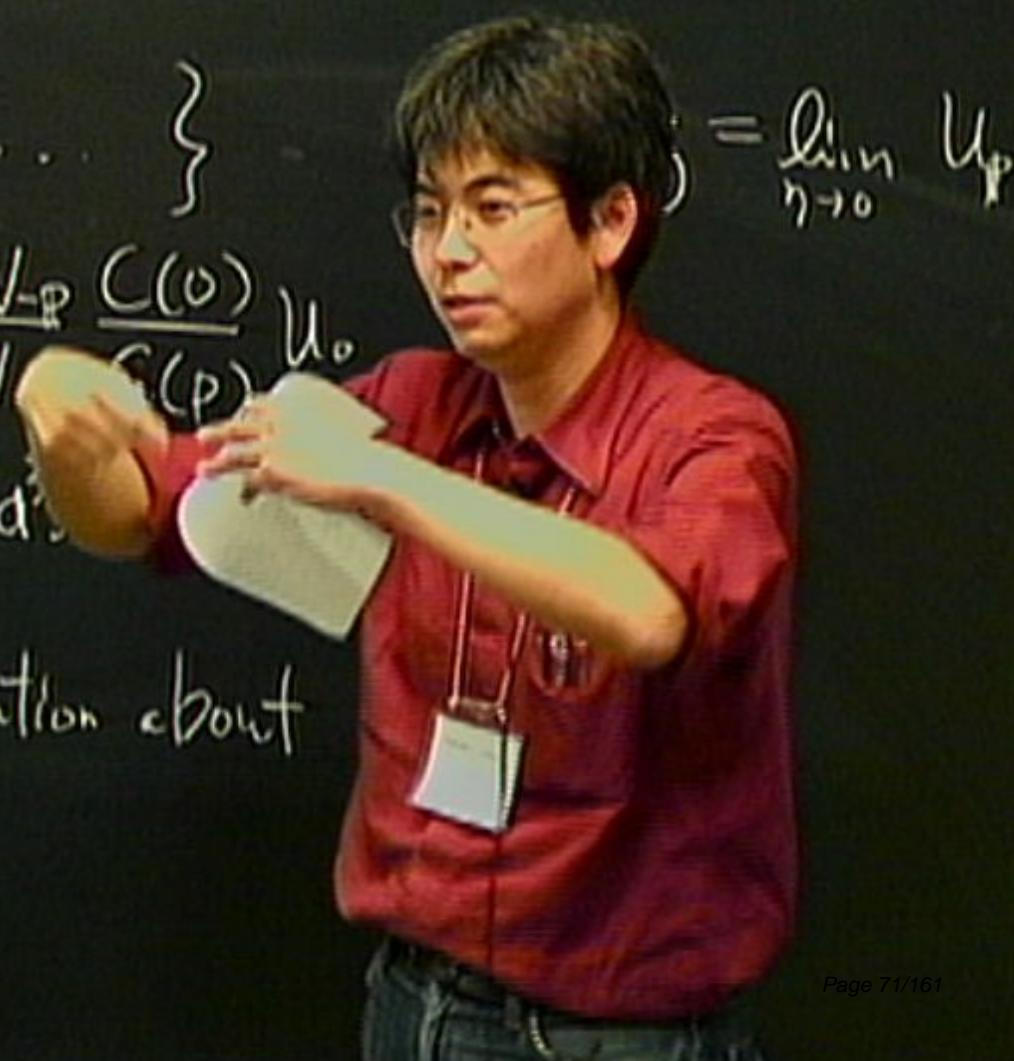
$(u, a) \xrightarrow{\text{unitary}} (v, b)$

$$U_0 = \frac{1}{N} \left\{ U_0 + \sum \dots \right\}$$

$$U_P = U_P e^{i P \cdot x} - \frac{W_P C(0)}{W C(P)} U_0$$

$$\int W(x) U_P(x) dx$$

U_0 contains information about



$(u, a) \xrightarrow{\text{unitary}} (v, b)$

$$v_o = \frac{1}{N} \left\{ u_o + \sum \dots \right\}$$

$$v_p = u_p e^{ip \cdot x} - \frac{w_p}{w_o} \frac{c(o)}{c(p)} u_o$$

$$\int w(x) v_p(x) d^3x = 0$$

v_o contains information about

$(u, a) \xrightarrow{\text{unitary}} (v, b)$

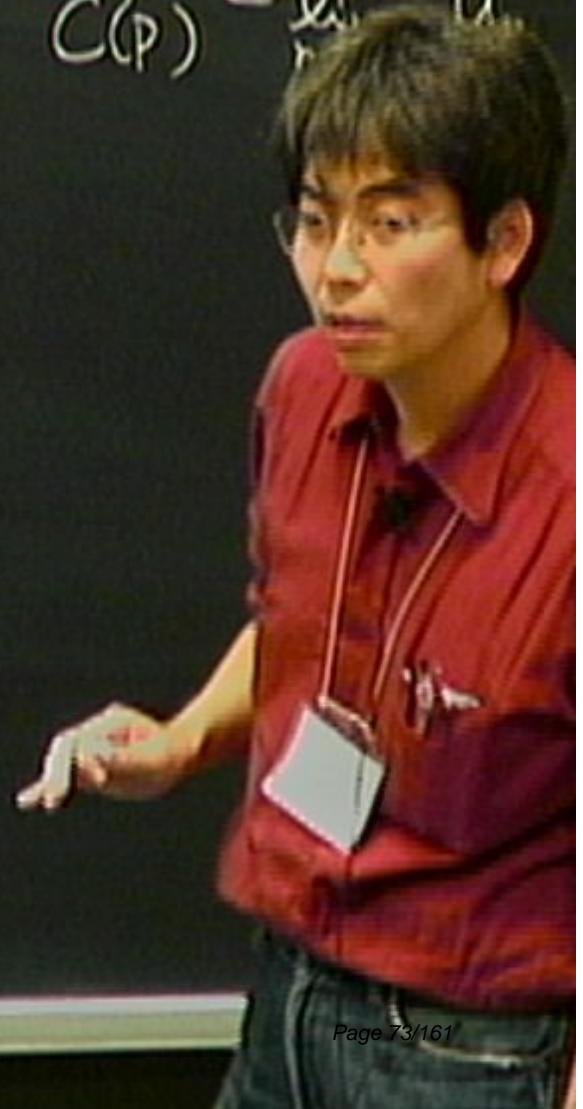
$$v_o = \frac{1}{N} \left\{ u_o + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \ell_1 \cdot u$$

$$v_p = u_p e^{ip \cdot x} - \frac{W_p}{W_o} \frac{C(o)}{C(p)} u_o$$

$$\int W(x) v_p(x) d^3x = 0$$

v_o contains information about \bar{J}



$(u, a) \xrightarrow{\text{unitary}} (v, b)$

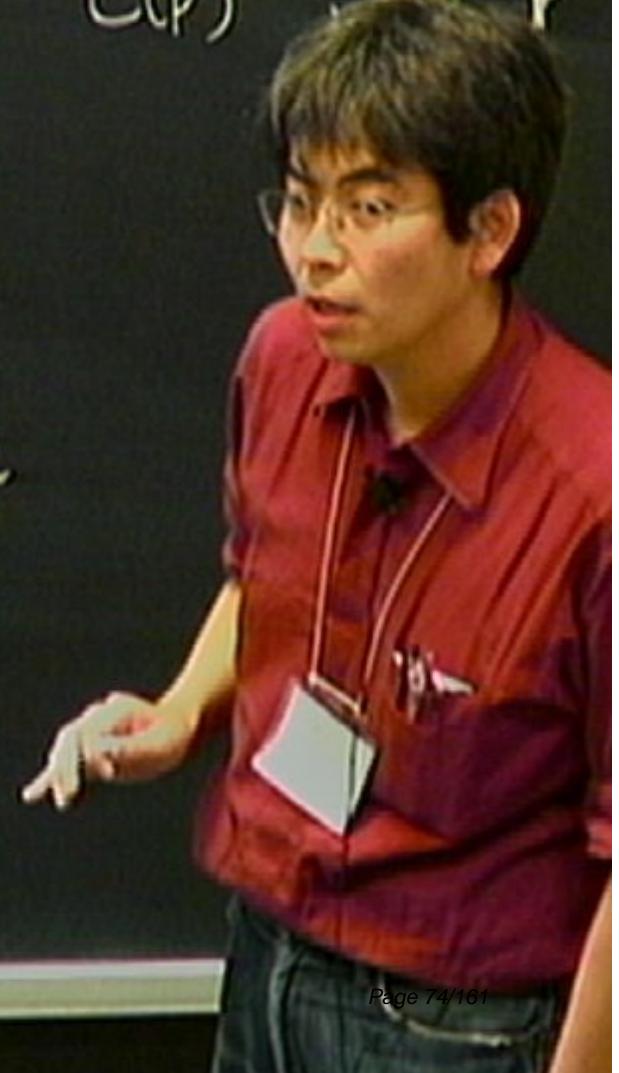
$$v_o = \frac{1}{N} \left\{ u_o + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \lim_{p \rightarrow \infty} U_p$$

$$U_p = U_p e^{ip \cdot x} - \frac{W_p}{W_o} \frac{C(o)}{C(p)} U_o$$

$$\int W(x) U_p(x) d^3x = 0$$

U_o contains information about \bar{J}



$(u, a) \xrightarrow{\text{unitary}} (v, b)$

$$v_0 = \frac{1}{N} \left\{ u_0 + \sum \dots \right\}$$

$$\frac{1}{C(p)} = \lim_{\eta \rightarrow 0} U_p$$

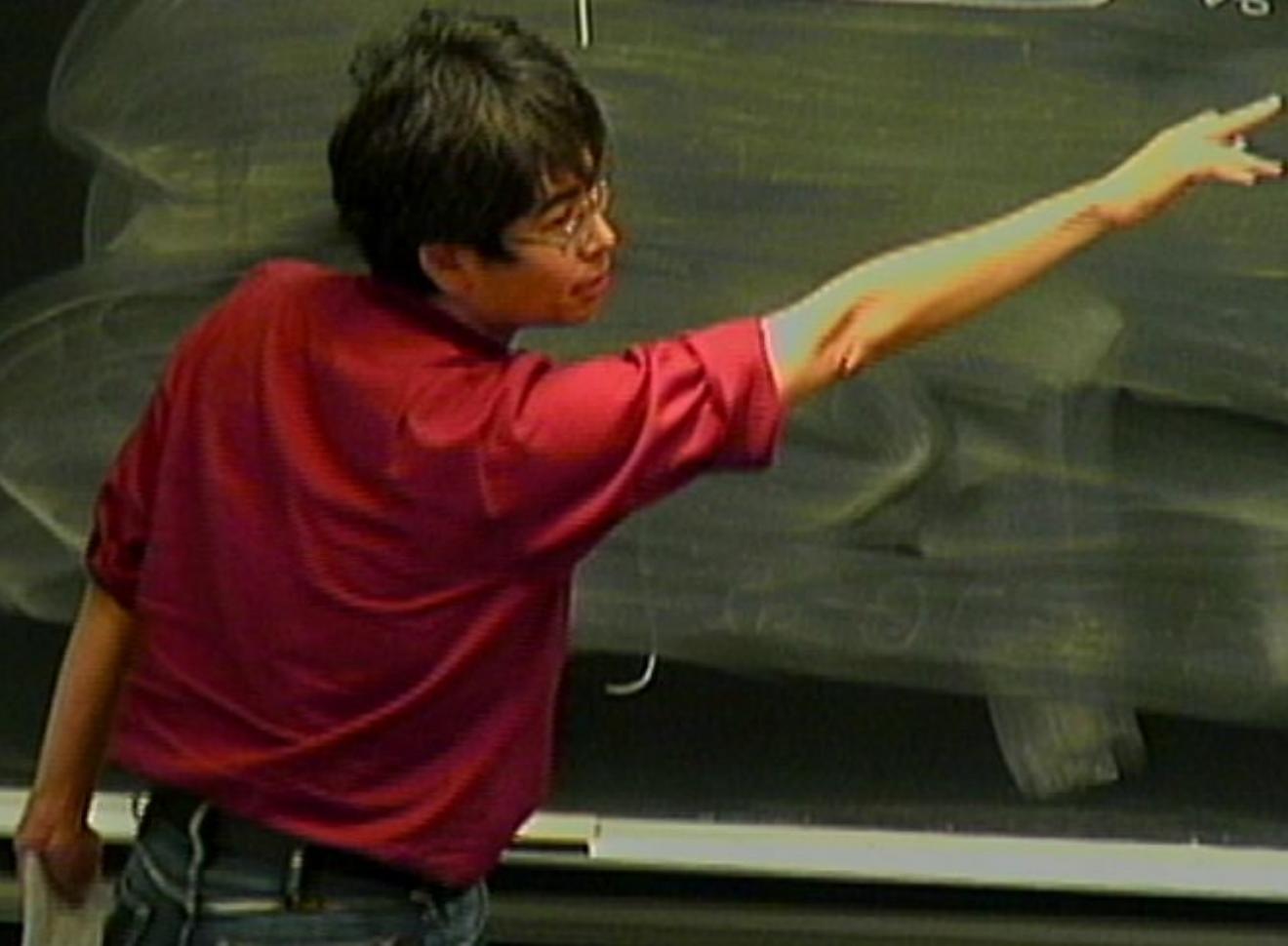
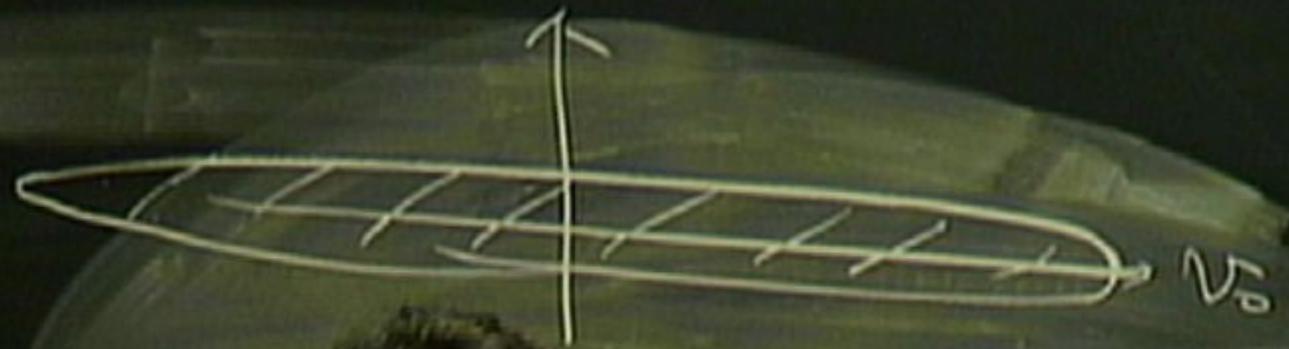
$$U_p = U_p e^{ip \cdot x} - \frac{W_p}{W_0} \frac{C(0)}{C(p)} u_0$$

$$\int W(x) U_p(x) d^3x = 0 \quad p \neq 0$$

U_0 contains information about \bar{J}

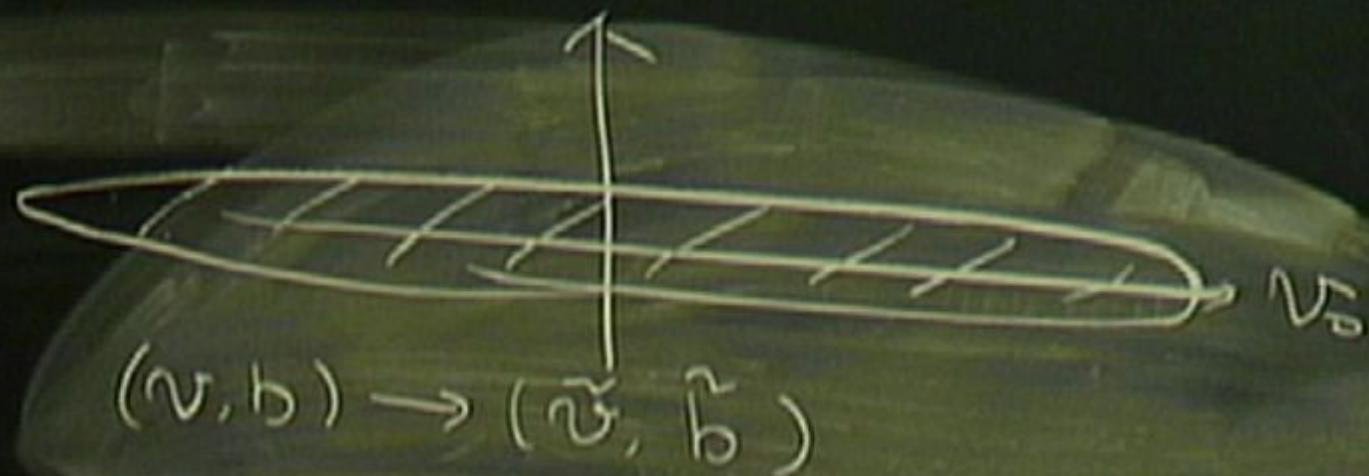
0904.8915

with Yuko.Urakawa



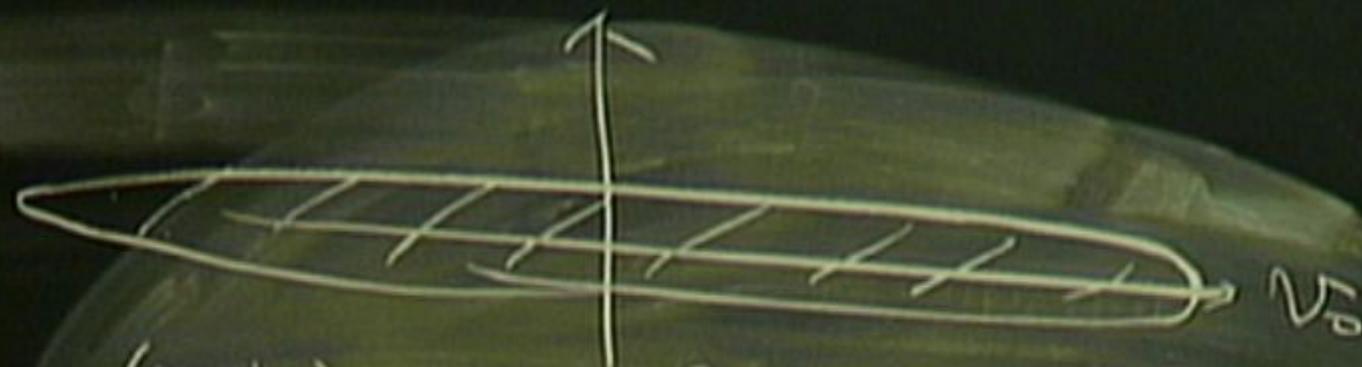
0904.8415

with Yuko.Urakawa



0904.8415

with Yuko, Urakawa

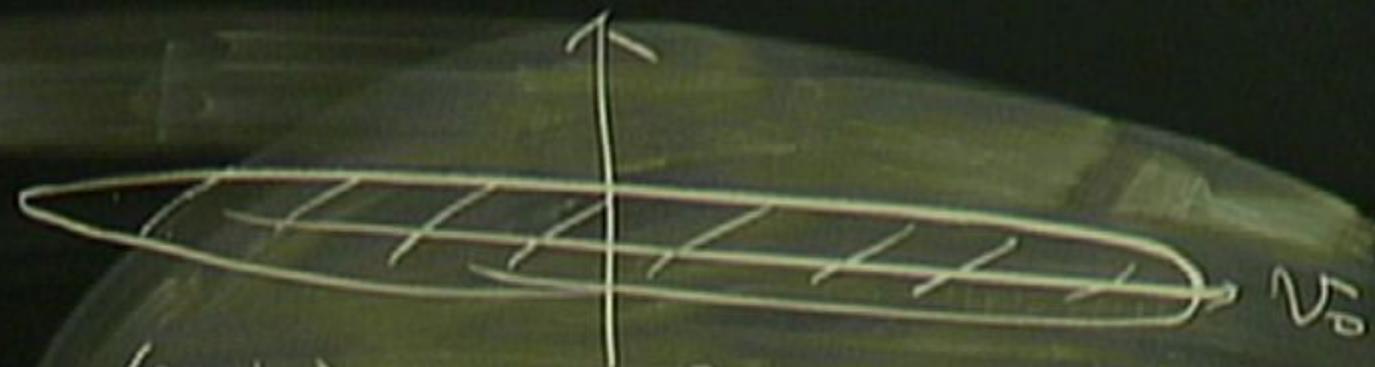


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \text{csh} r v_0 - \text{sh} r \gamma_0^*$$

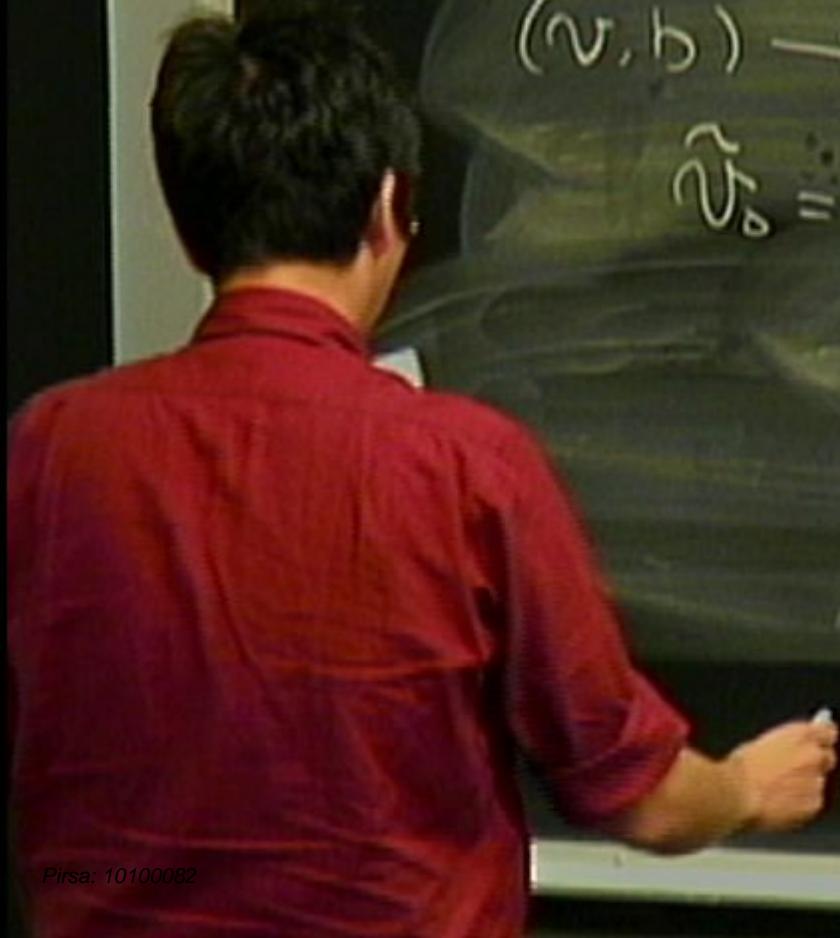
0904.8415

with Yuko, Urakawa



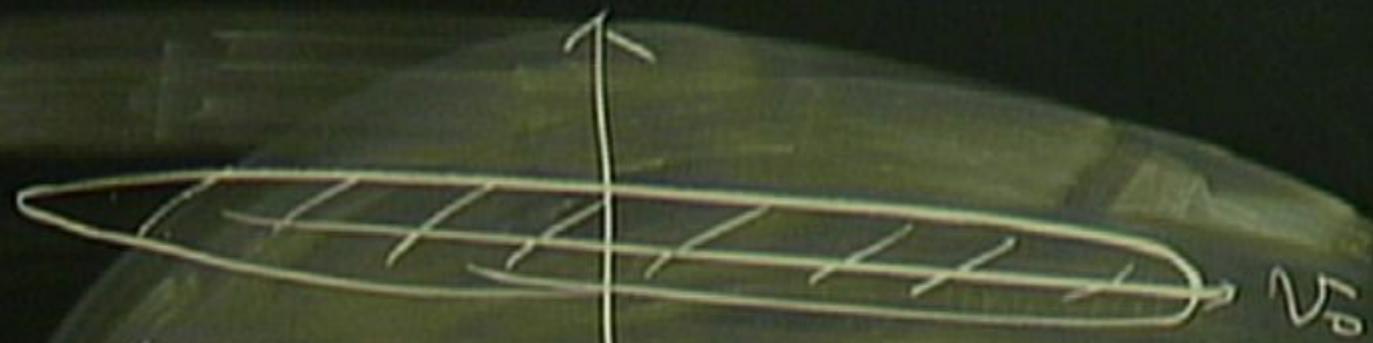
$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \text{cosh} r v_0 - \text{sh} r \gamma_0^*$$



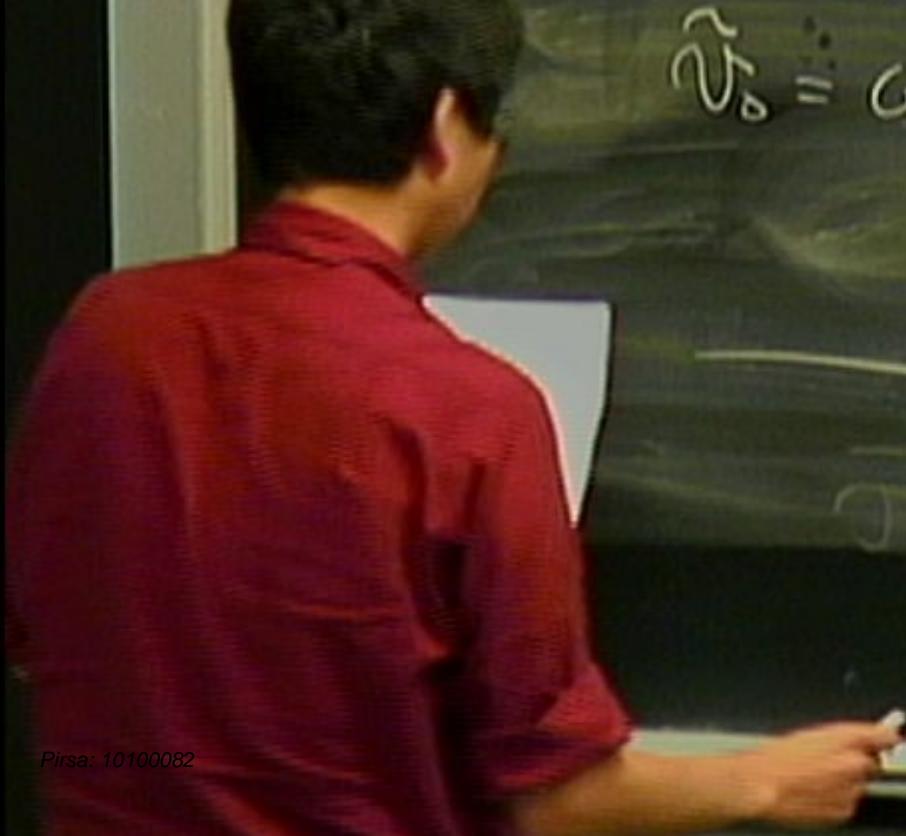
0904.8415

with Yuko, Urakawa



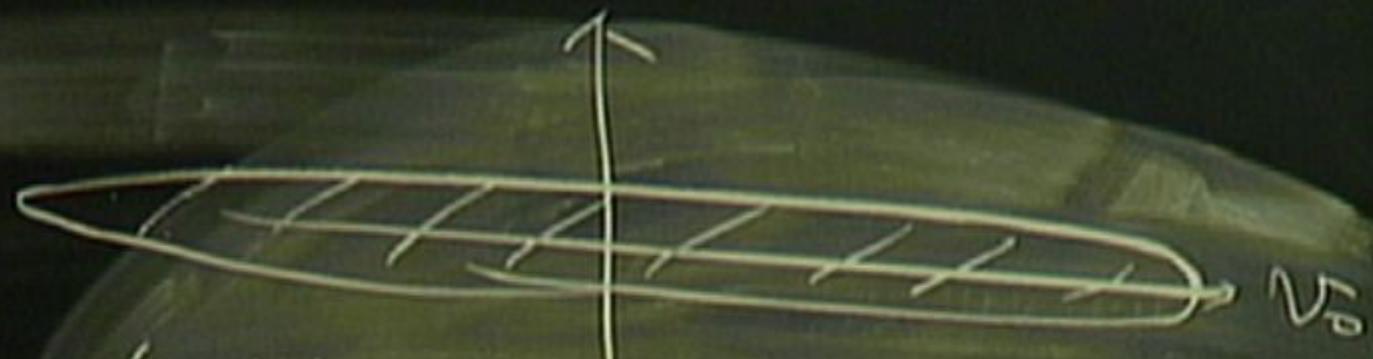
$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \text{cosh} r v_0 - \text{sh} r \gamma_0^*$$



0904.8415

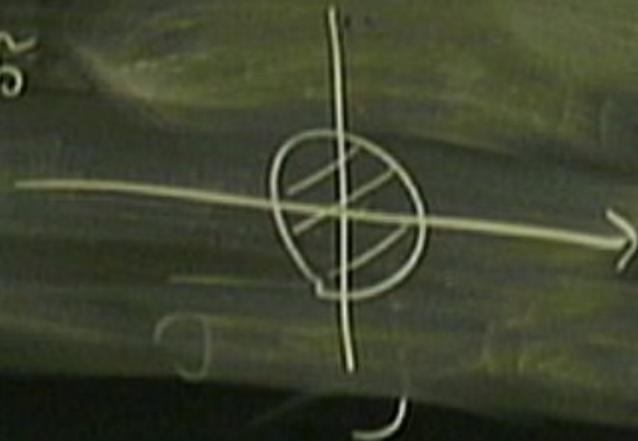
with Yuko, Urakawa



A diagram of a rotating frame with two axes. One axis is vertical and labeled v_0 . The other axis is horizontal and labeled b_0 . A grid of lines is drawn perpendicular to the v_0 axis.

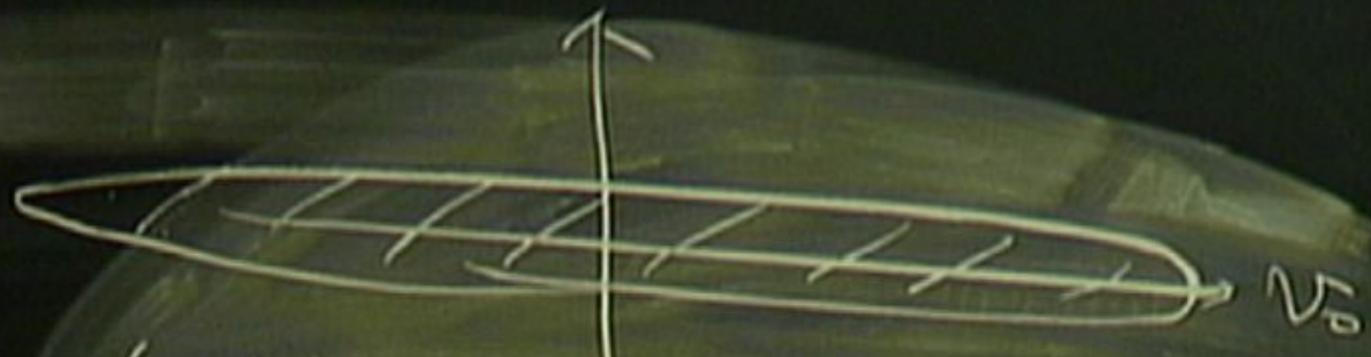
$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$
$$|0>_{\tilde{b}}$$



0904.8415

with Yuko, Urakawa


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

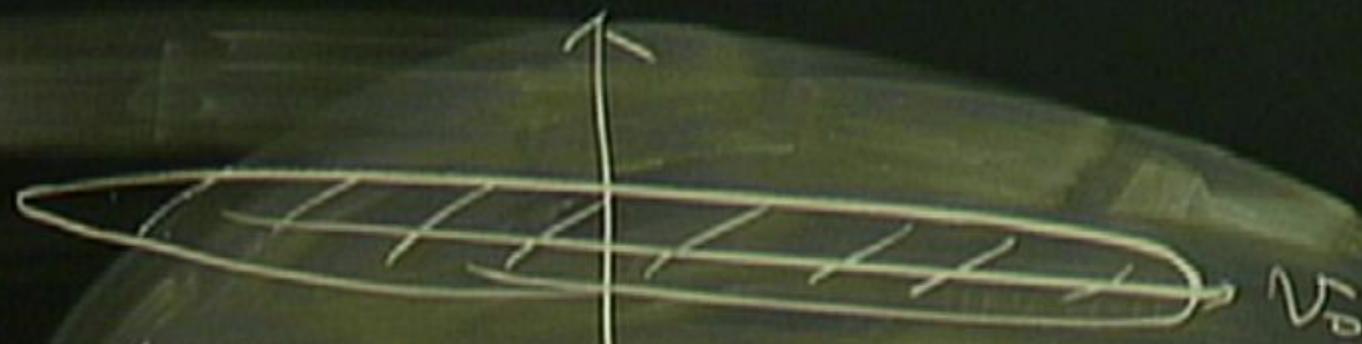
$$\tilde{v}_\delta = \text{csh} r v_\delta - \text{sh} r \gamma_\delta^*$$

$$|0>_{\tilde{b}}$$



0904.8415

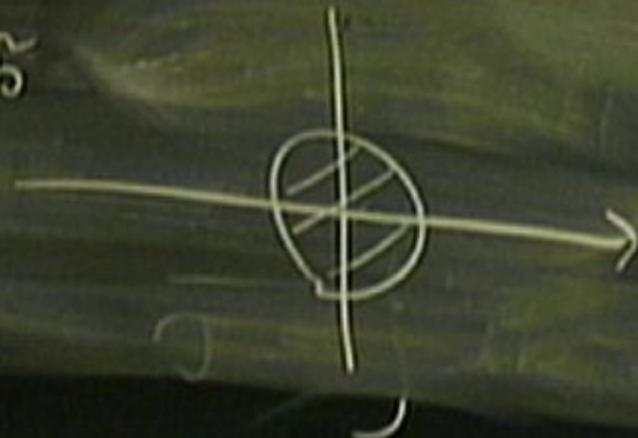
with Yuko, Urakawa



$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

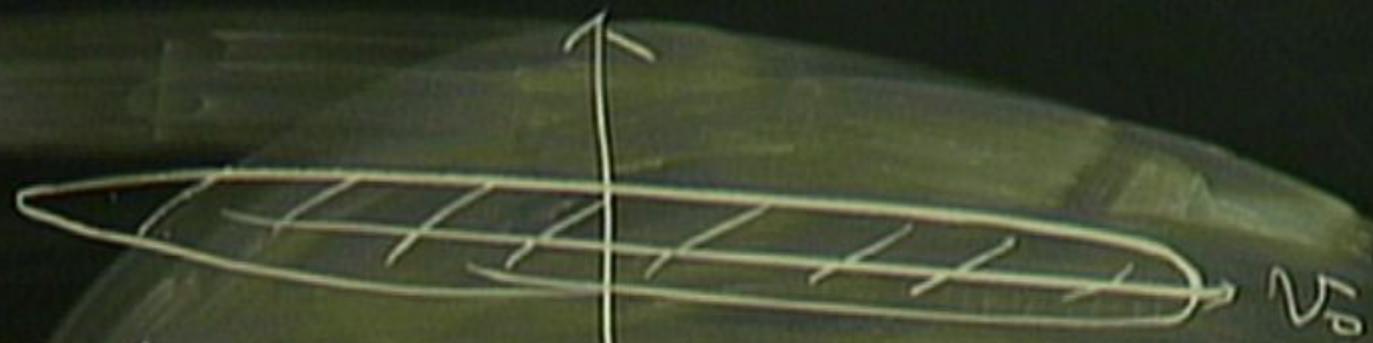
$$\tilde{v}_0 = \cosh r v_0 - \sinh r \gamma_0^*$$

$$|0>_{\tilde{b}}$$



0904.8415

with Yuko, Urakawa



$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

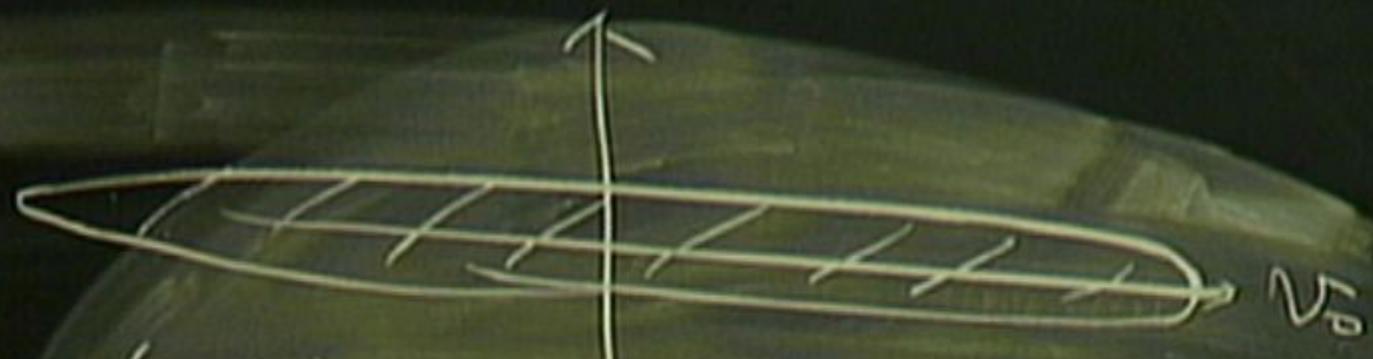
$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$

$$|0>_{\tilde{b}}$$



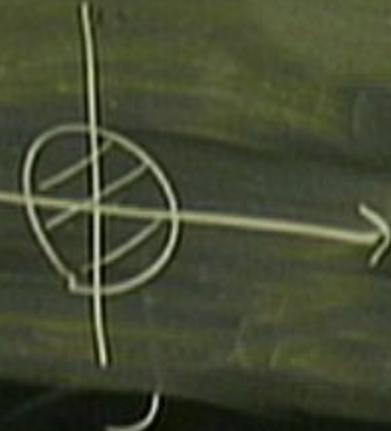
0904.8415

with Yuko, Urakawa


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$

$$|0>_{\tilde{b}}$$



0904.8415

with Yuko Urakawa

Coherent st

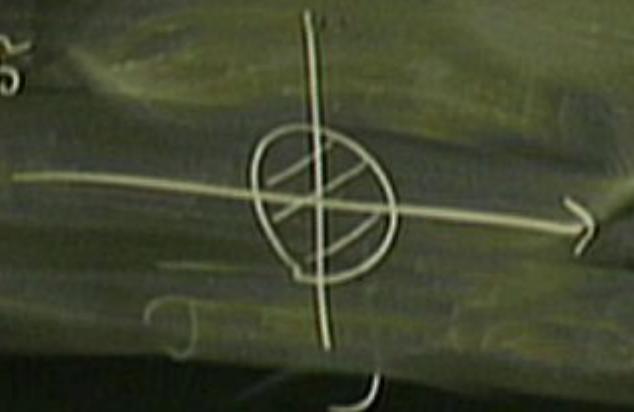
$$\tilde{b}|\beta\rangle =$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{...}$$

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$

$$|0\rangle_b$$



0904.8415

with Yuko.Urakawa

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

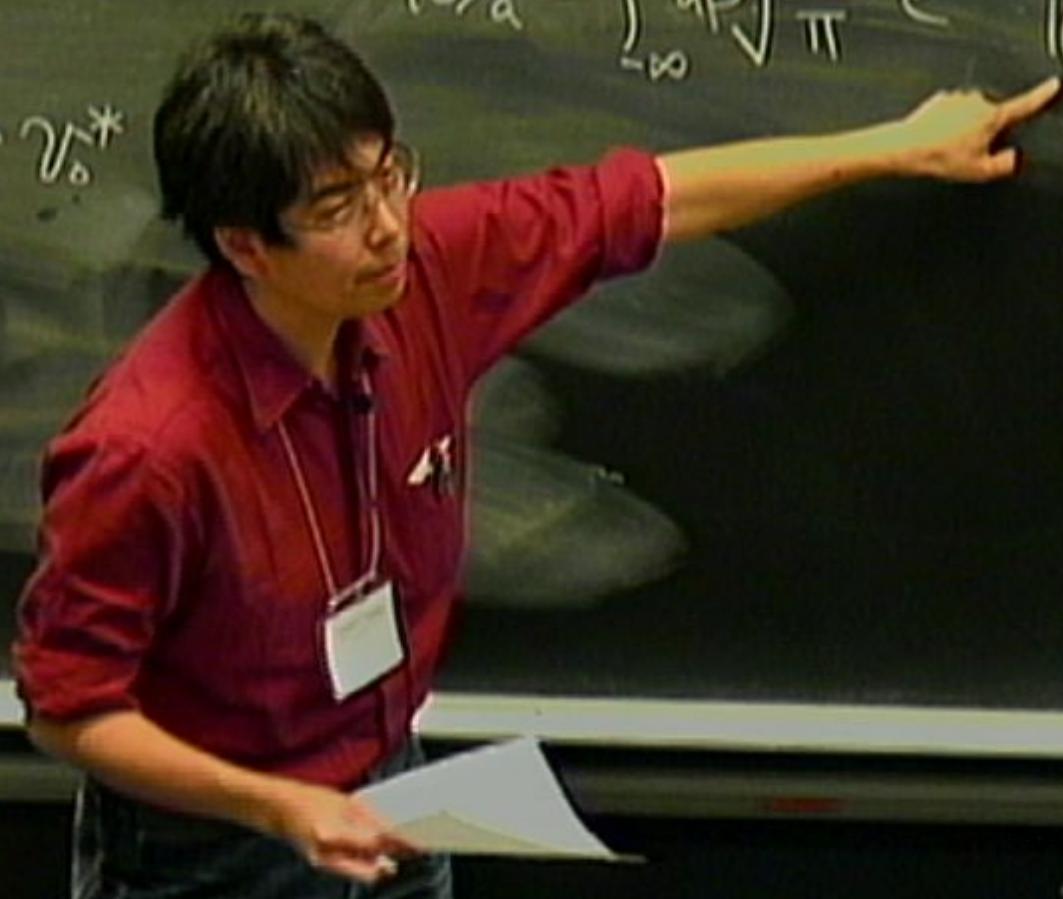
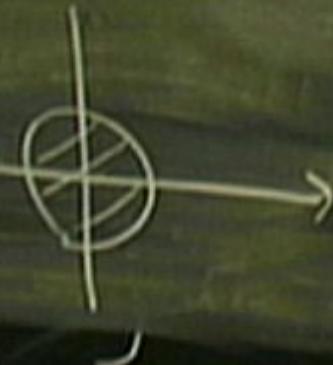
coherent state

$$\tilde{b} |\beta\rangle = \beta |\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{S}{\pi}} e^{-(S\beta)^2} |\beta\rangle_{\tilde{b}}$$

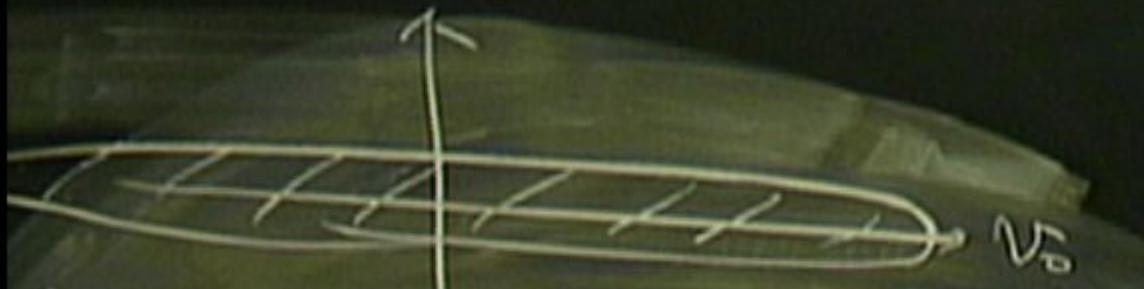
$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$

$$|0\rangle_b$$



0904.8415

with Yuko Urakawa



$$(\nu, b) \rightarrow (\tilde{\nu}, \tilde{b})$$

$$\tilde{\nu}_0 = \cosh r \nu_0 - \sinh r \nu_0^*$$

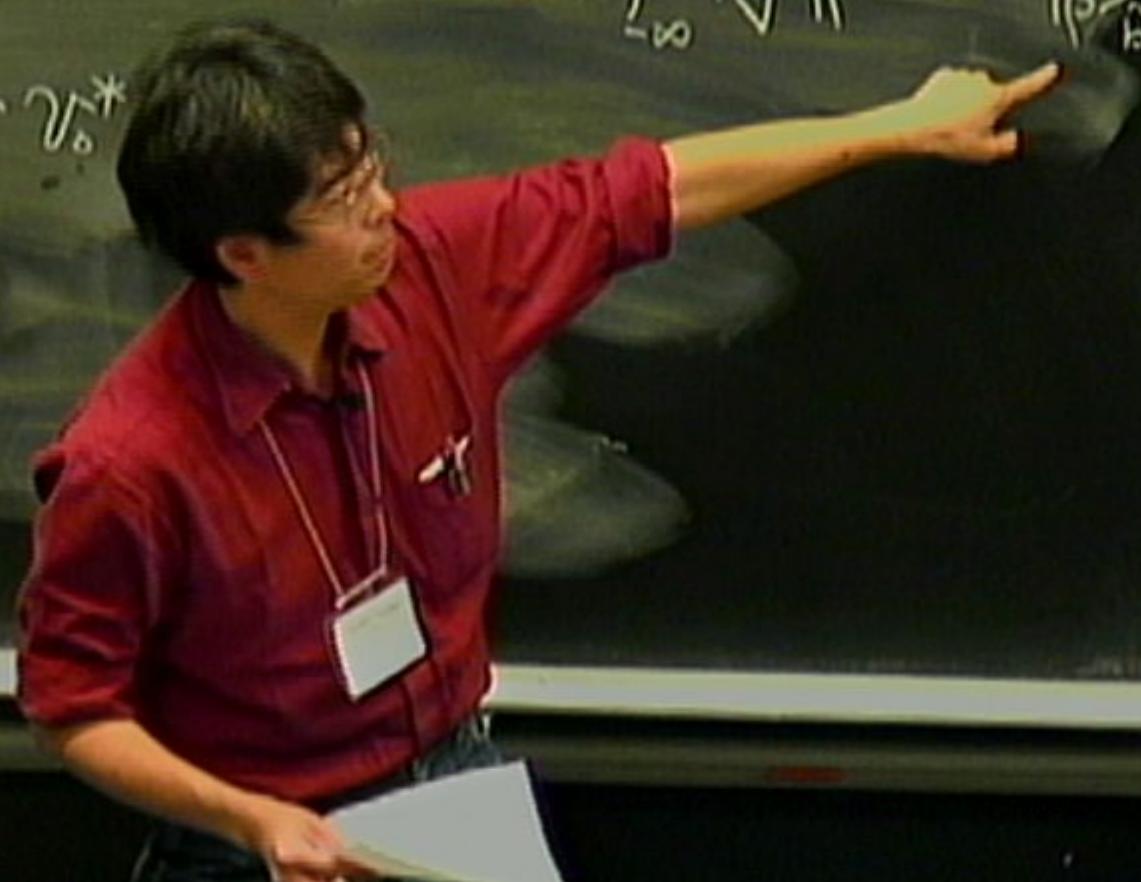
$$|0\rangle_{\tilde{b}}$$



coherent state

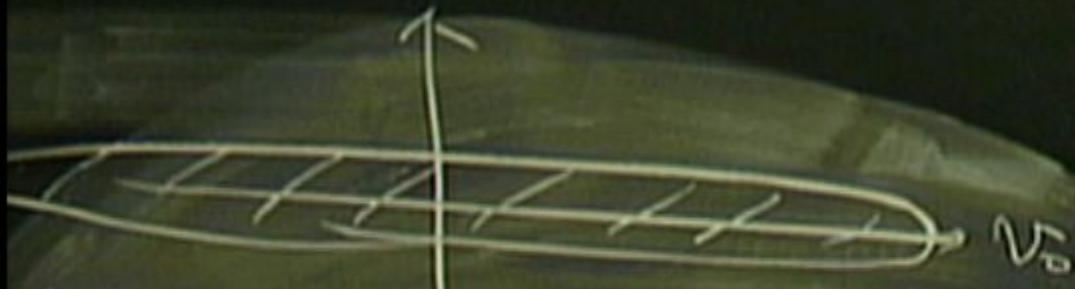
$$\hat{b} |\beta\rangle = \beta |\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{S}{\pi}} e^{-(S\beta)^2} |\beta\rangle_{\tilde{b}}$$



0904.8415

with Yuko, Urakawa



$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cos r v - \sin r \gamma$$

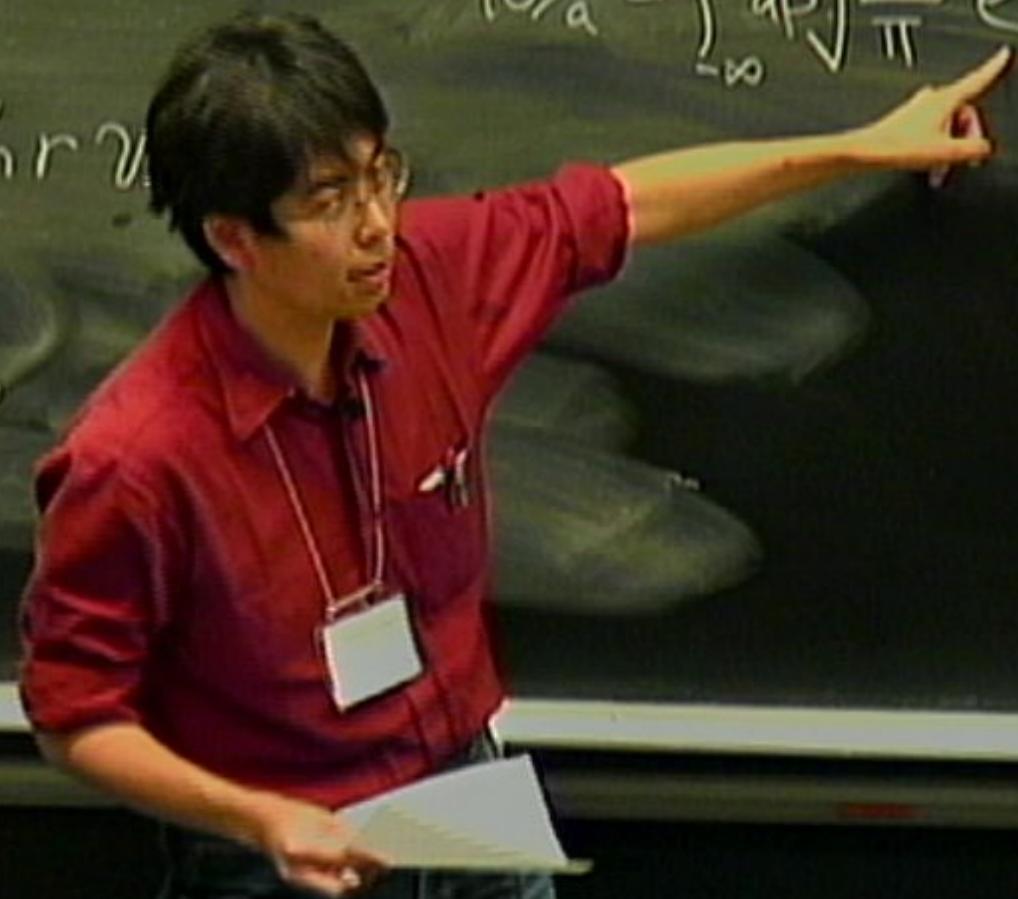
$$|0\rangle_{\tilde{b}}$$



coherent state

$$\hat{b} |\beta\rangle = \beta |\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{S}{\pi}} e^{-(S\beta)^2} |\beta\rangle_{\tilde{b}}$$



0904.8415

with Yuko Urakawa

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$
$$\tilde{v}_0 = \cosh r v_i - \sinh r v_o$$

$$|0\rangle_{\tilde{b}}$$



Coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|\beta\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{S}{\pi}} e^{-(S\beta)^2} |\beta\rangle_{\tilde{b}}$$



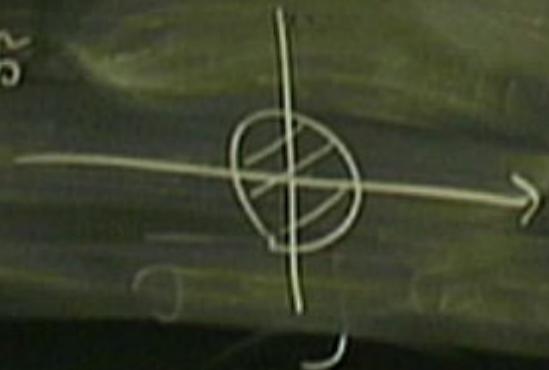
0904.8415

with Yuko Urakawa

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \text{cosh} r v_i - \text{sinh} r v_i^*$$

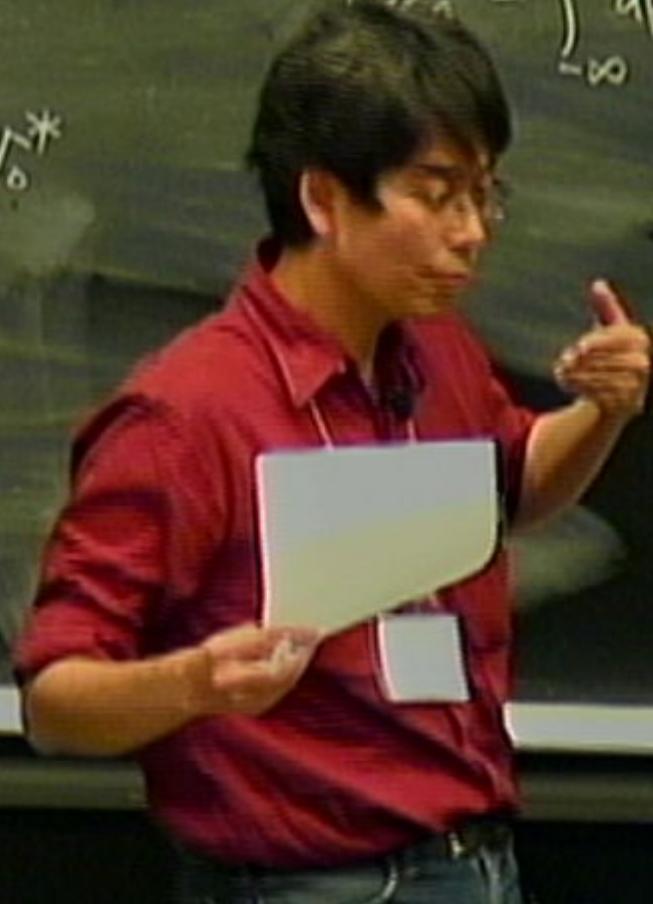
$$|0\rangle_{\tilde{b}}$$



coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|\beta\rangle = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{1}{\pi}} e^{-(\beta^2)} |\beta\rangle_{\tilde{b}}$$



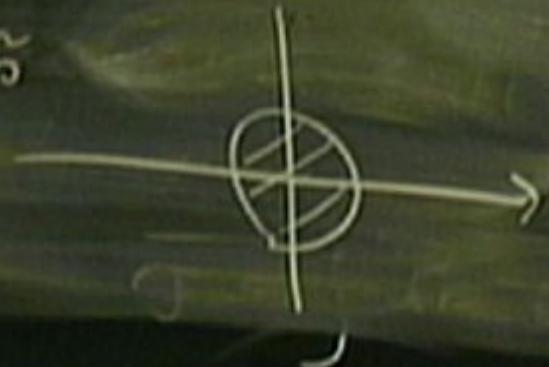
0904.8415

with Yuko Urakawa


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cosh r v_i - \sinh r v_i^*$$

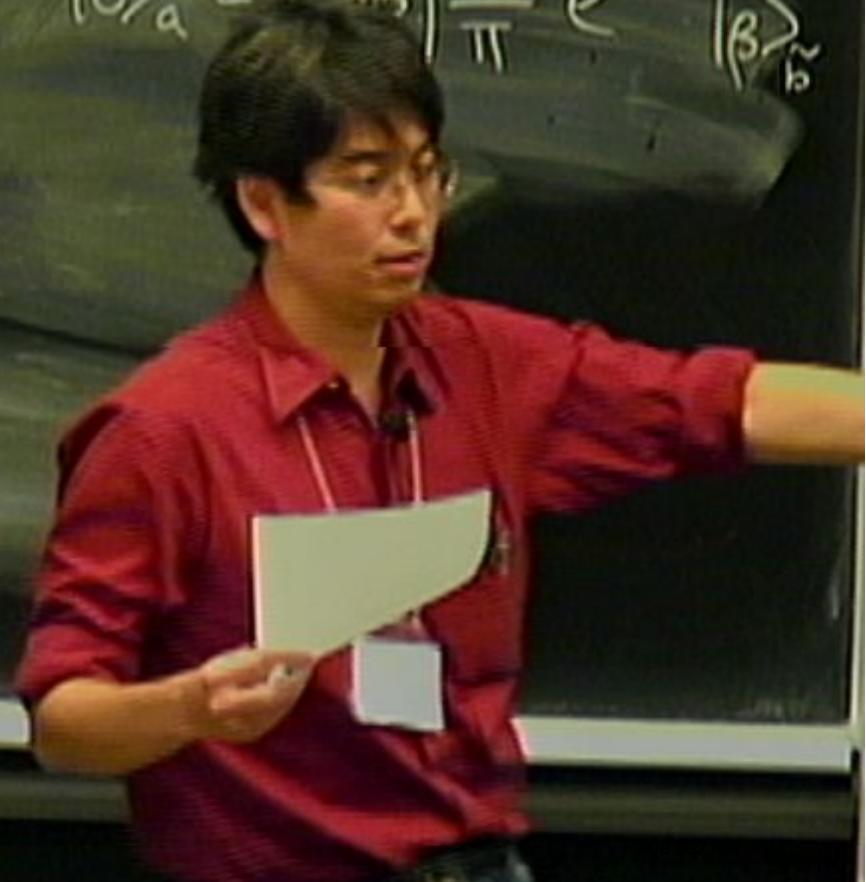
$$|0\rangle_{\tilde{b}}$$



coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{2}{\pi}} e^{-(\beta)^2} |\beta\rangle_{\tilde{b}}$$

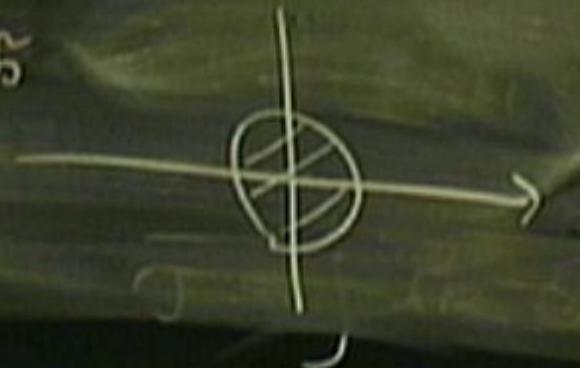


0904.8415

with Yuko Urakawa


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$
$$\tilde{v}_0 = \cosh r v_i - \sinh r v_0^*$$

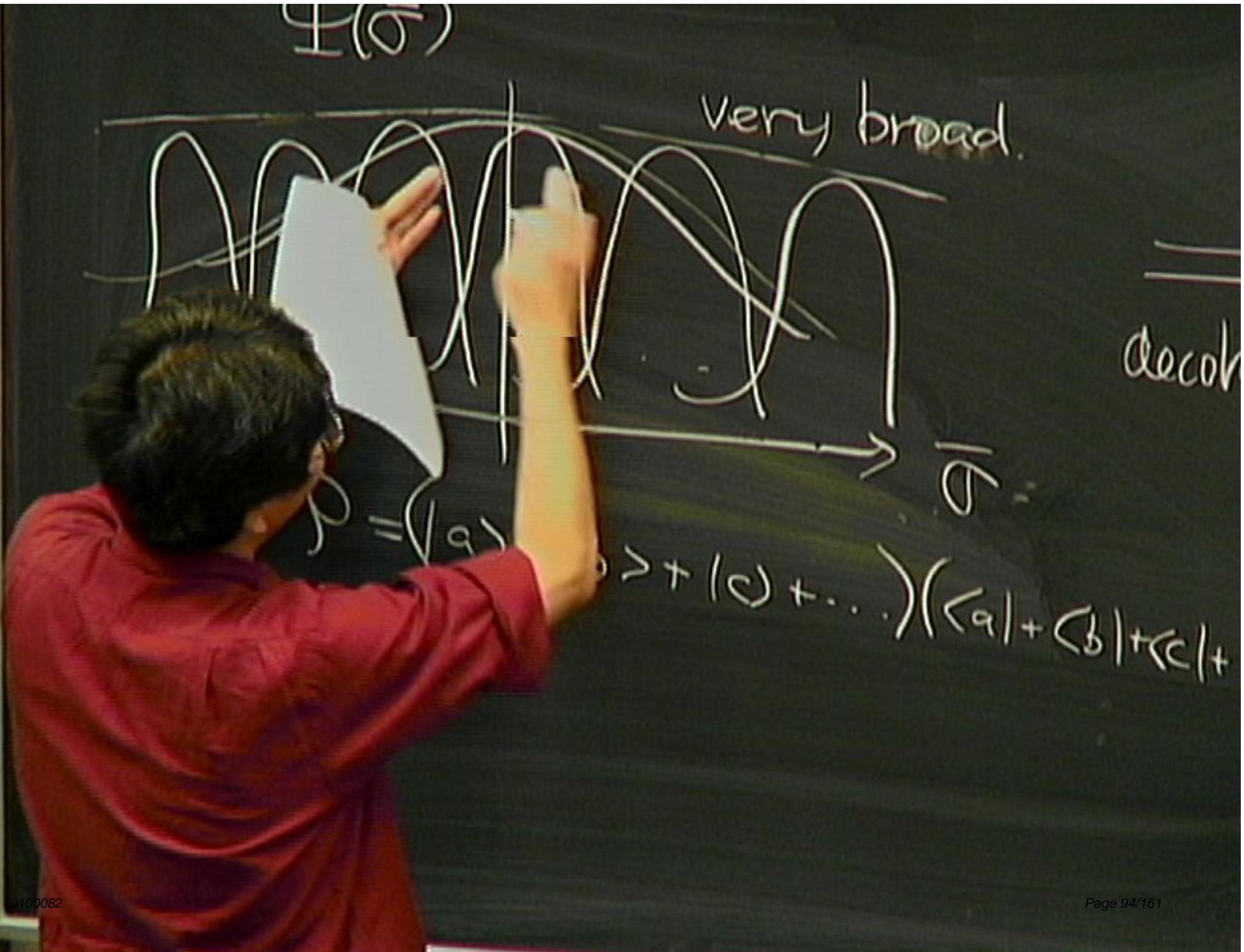
$|0\rangle_{\tilde{b}}$



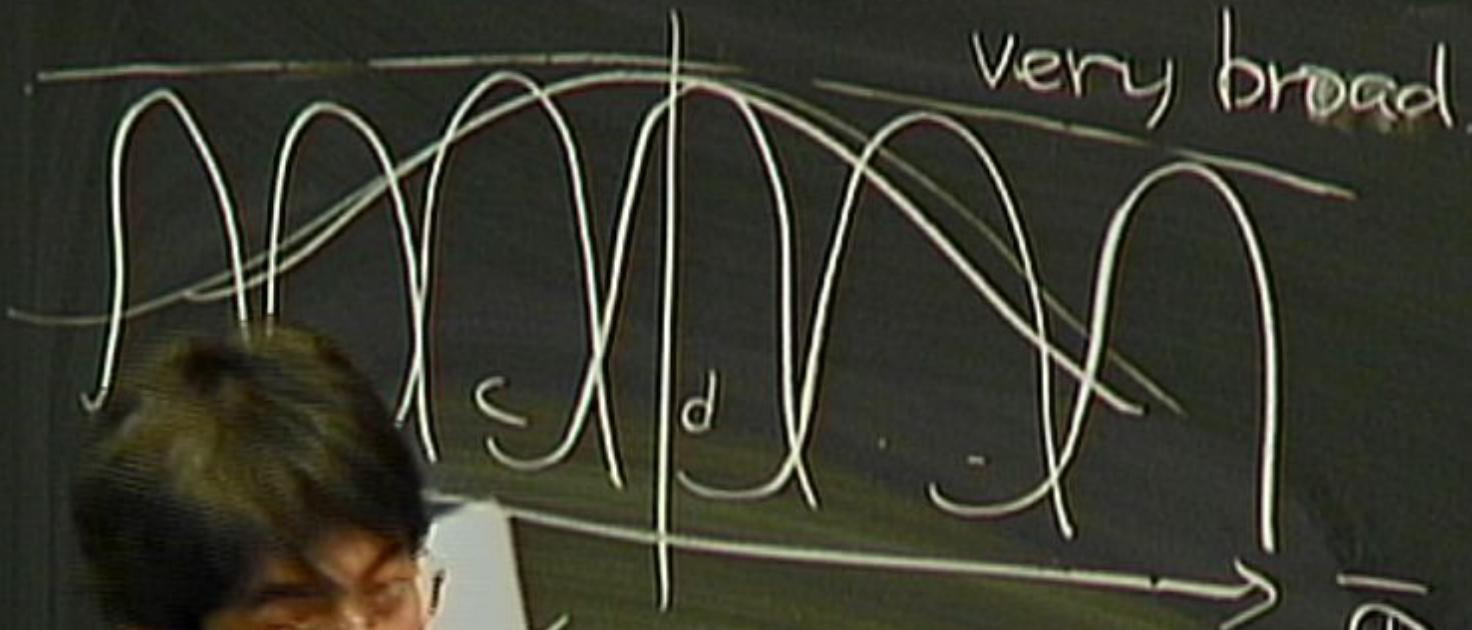
coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|\beta\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{S}{\pi}} e^{-(S\beta)^2}$$

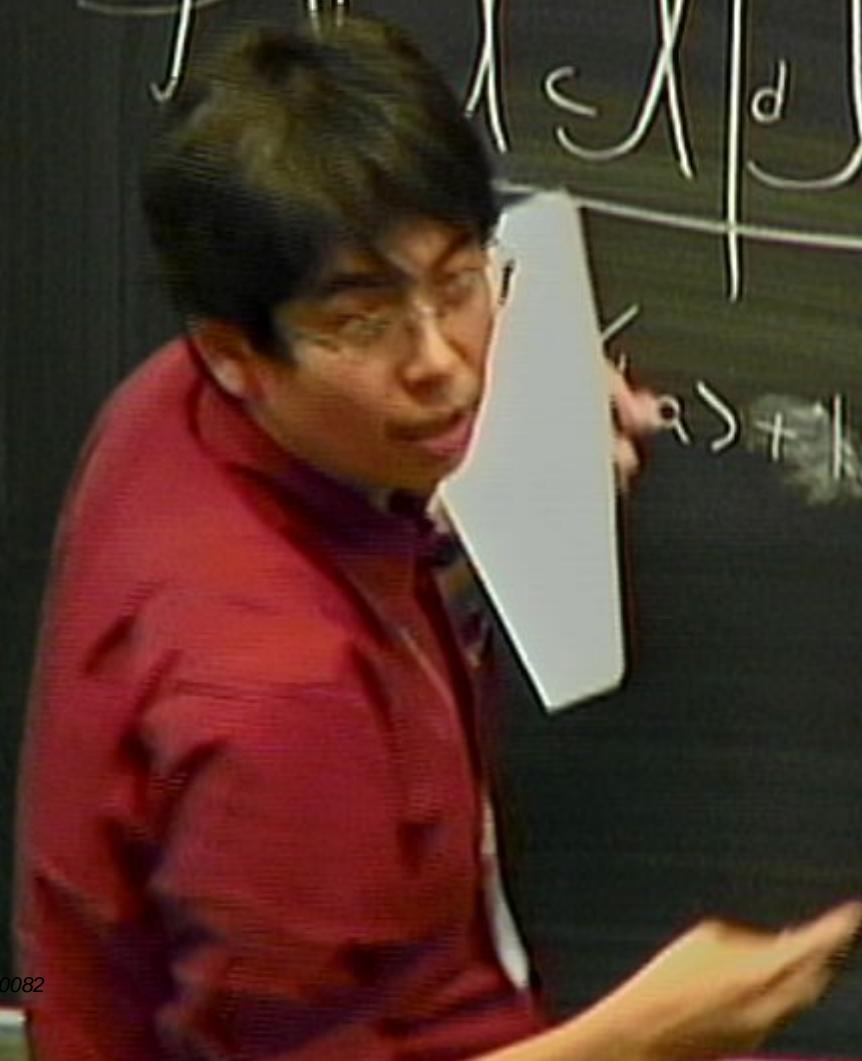


$\Psi(\sigma)$



$\bar{\Psi}$

$$|a\rangle + |b\rangle + |\psi + \dots\rangle (|a| + |b| + |\psi| + \dots)$$



0904.8415

with Yuko Urakawa

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$
$$\tilde{v}_0 = v_0 - s h r v_0^*$$

$$|0\rangle_{\tilde{b}}$$

coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|\beta\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{s}{\pi}} e^{-(s\beta)^2} |\beta\rangle_{\tilde{b}}$$
$$s \rightarrow 0$$

0904.8415

with Yuko.Umeda


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$
$$\tilde{v}_\delta = \cosh r v_i - \sinh r v_\delta^*$$

|0



coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{s}{\pi}} e^{-(s\beta)^2} |\beta\rangle_{\tilde{b}}$$
$$s \rightarrow 0$$

0904.24/5

with Yuko.Urakawa

$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$

$$\tilde{v}_0 = \cos r v_i - \sin r v_i^*$$

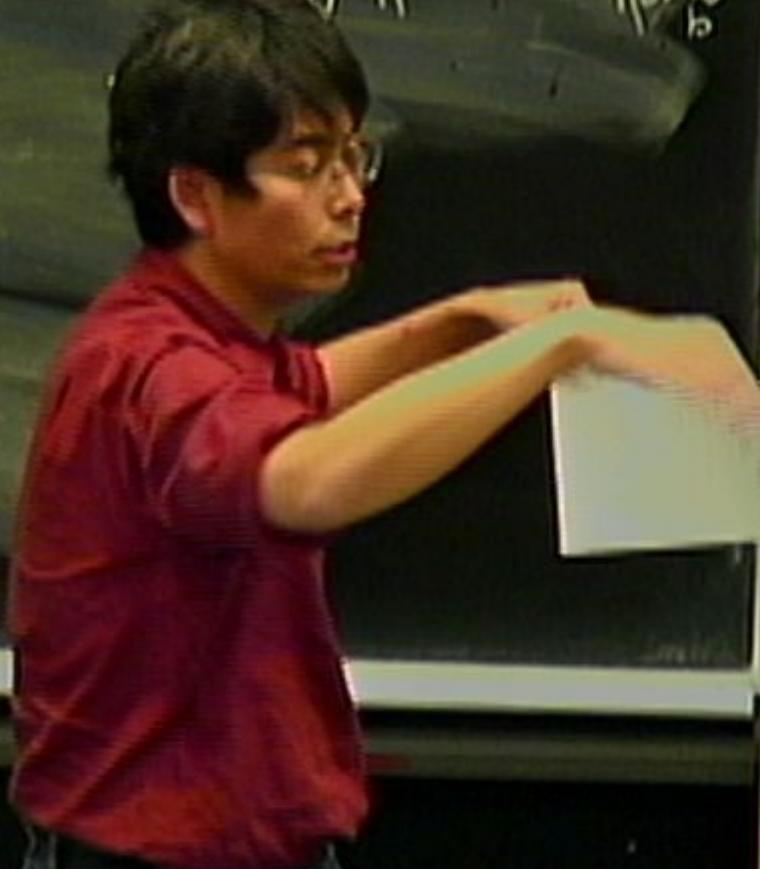
$$|0\rangle_{\tilde{b}}$$



coherent state

$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|\beta\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{2}{\pi}} e^{-(\beta)^2} |\beta\rangle_{\tilde{b}}$$

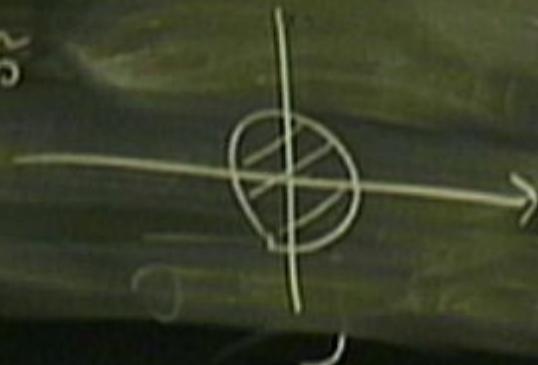


0904.24/5

with Yuko.Umekawa


$$(v, b) \rightarrow (\tilde{v}, \tilde{b})$$
$$\tilde{v}_0 = \cosh r v_0 - \sinh r v_0^*$$

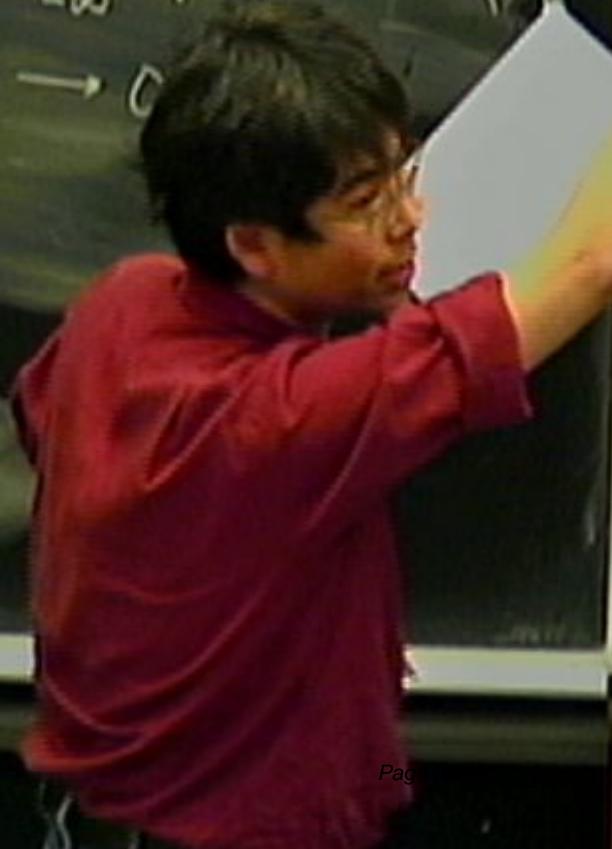
$|0\rangle_b$



coherent state

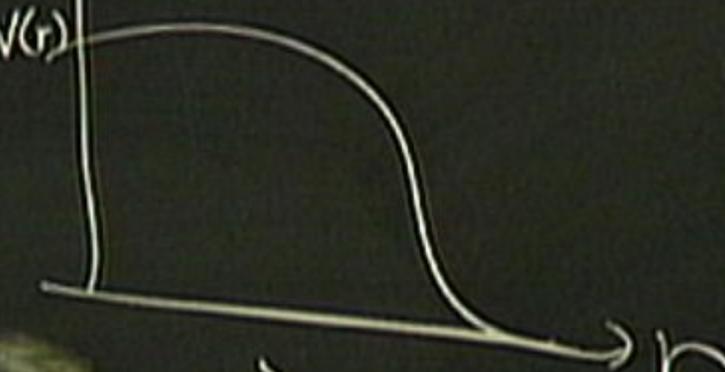
$$\tilde{b}|\beta\rangle = \beta|\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{s}{\pi}} e^{-(s\beta)^2} |\beta\rangle$$
$$s \rightarrow 0$$



$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$

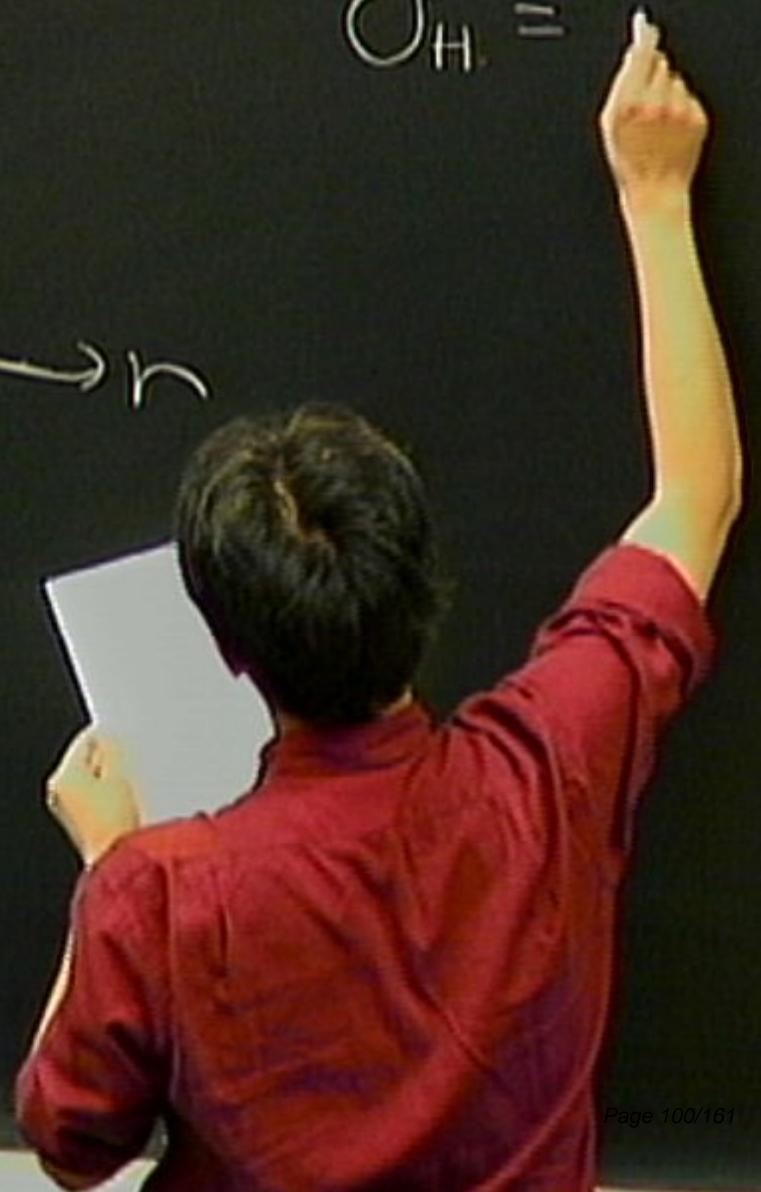
$W(r)$



$$\{\bar{\sigma}, \sigma - \bar{\sigma}\}$$

1_{def}

$$\hat{\sigma}_H =$$

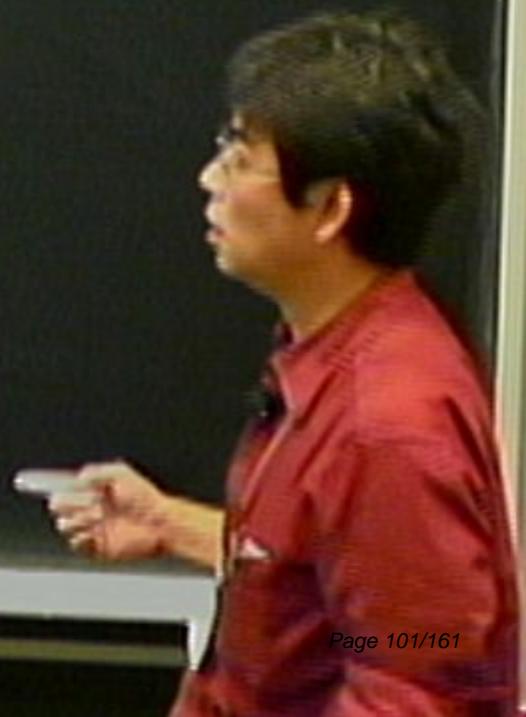


what we observe.

$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$
$$\left\{ \bar{\sigma}, \sigma - \bar{\sigma} \right\}$$

1 day

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(x') dx'$$

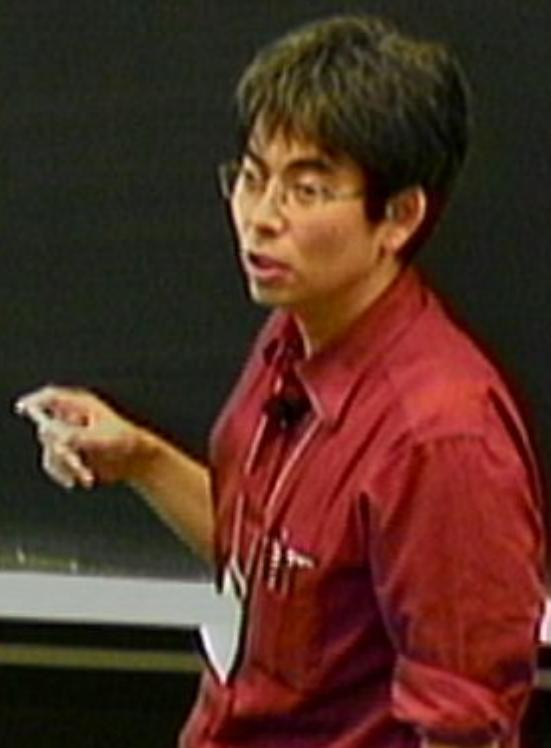


what we observe.

$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$
$$\{\bar{\sigma}, \sigma - \bar{\sigma}\}$$

1 day

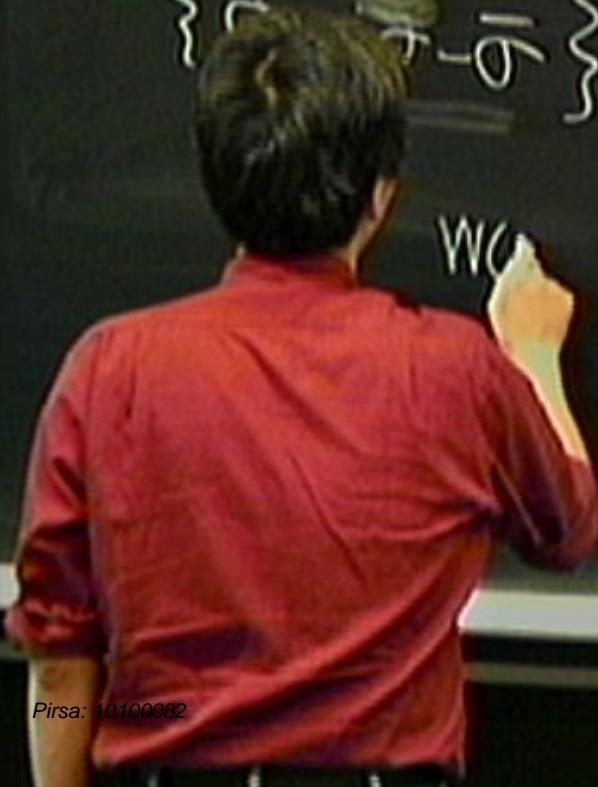
$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(z, x') \Gamma'(x') dx'$$



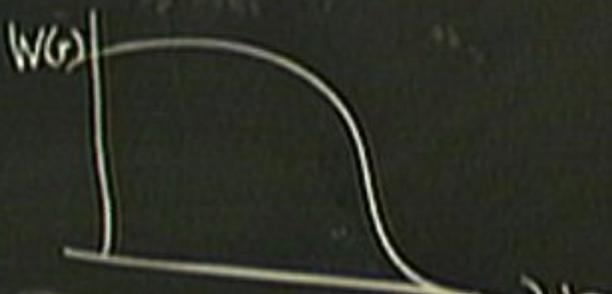
what we observe.

$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$
$$\{ \bar{r}, \bar{\tau}, \bar{\sigma} \}$$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma'(x') dx'$$



what we observe.

$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$


$$\{\bar{\sigma}, \dots\}$$

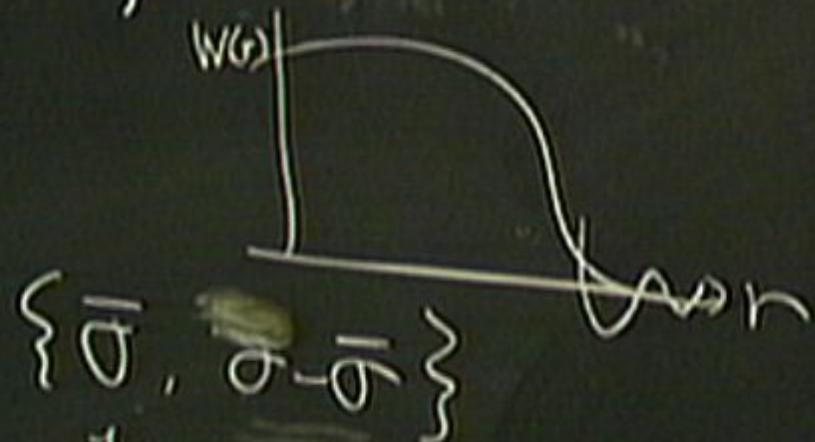
dag

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma'(x') dx'$$

$$\langle \dots \rangle = \int \dots$$

what we observe.

$$\bar{\sigma} = \int d^3x W(x) \sigma(x)$$

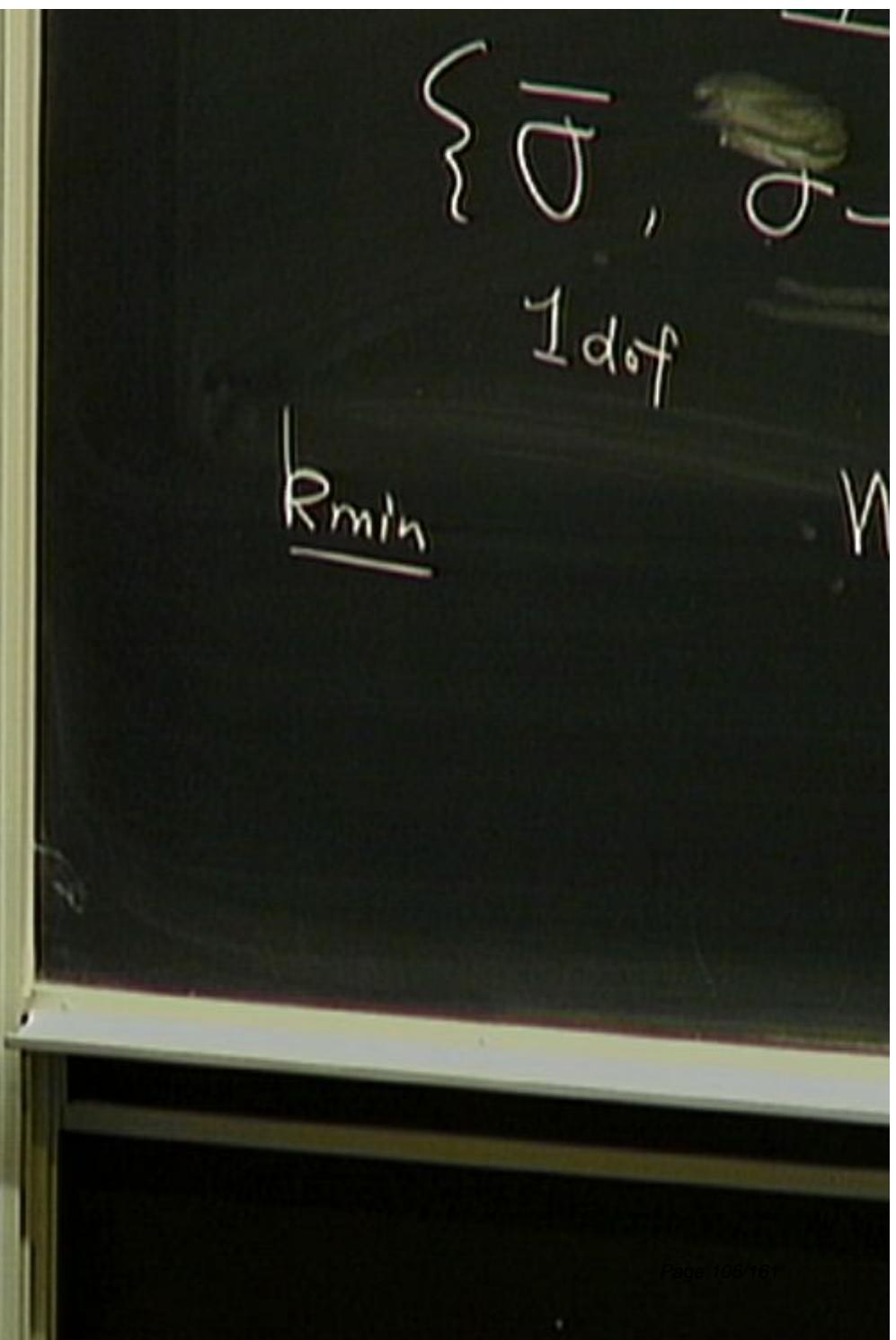


$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma'(x') dx'$$

$$\{\bar{\sigma}, \bar{\sigma} - \bar{\sigma}\}$$

1d_f

$$W(r) = \int d^3k e^{ikr} \Theta(k_{\text{cutt}} - k)$$



what we observe.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(z, z') T(\alpha') dz'$$



WHAT WE OBSERVE.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(z, z') P(\alpha') dz'$$
$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$\langle 0 | P(0) | 0 \rangle_a$

WHAT WE OBSERVE.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$
$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

and $\langle 0 | P(0) | 0 \rangle_a$

what we observe.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') P(\alpha') dx'$$
$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle 0 | P(0) | 0 \rangle_a$$

$$\approx \int d\alpha \int d\beta \langle \alpha | P(0) | \beta \rangle_b$$

what we observe.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$
$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle 0 | P(0) | 0 \rangle_a$$

$$\approx \int d\alpha \int d\beta \langle \alpha | P(0) | \beta \rangle_b$$

what we observe.

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') T(\alpha') dx'$$
$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle 0 | P(0) | 0 \rangle_a$$

$$\propto \int d\beta \frac{\langle \alpha | P(0) | \beta \rangle_b}{\langle \beta | \beta \rangle_b}$$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle 0 | P(0) | 0 \rangle_a$$

$$\frac{\int d\beta \langle \alpha | P(0) | \beta \rangle_b}{\int d\beta}$$



$$\langle \varphi |$$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle 0 | P(0) | 0 \rangle_a$$

$$\frac{i \alpha \int d\beta \langle \alpha | P(0) | \beta \rangle_b}{\langle \beta | \alpha \rangle_b}$$



$$\langle \beta | \alpha \rangle_b$$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_J^2 + \dots$$

~~and~~ $\langle 0 | P(0) | 0 \rangle_a$

$$\frac{\int d\alpha \int d\beta \langle \alpha | P(0) | \beta \rangle_b}{\int d\beta}$$

↓

$$(e^{ikx} e^{-ikx}) \langle \alpha | \beta \rangle_b$$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

$$\langle O | P O | O \rangle_a$$

$$\approx \frac{\int d\alpha \int d\beta \langle \alpha | P O | \beta \rangle_b}{\int d\alpha \int d\beta}$$

$$\theta(k_{\text{cut}} - k) \downarrow \langle \alpha | \beta \rangle_b \times \langle O | (PO) \rangle_b$$

$$\hat{O}_H = \hat{O}_I + \int G_R(x, x') f(x') dx'$$

$$= \hat{O}_I + \hat{O}_I^2 + \dots$$

$$\langle O | P O | O \rangle_a$$

$$\approx \frac{\int d\alpha \int d\beta \langle \alpha | P O | \beta \rangle_b}{\int d\alpha \int d\beta}$$

$$\theta(k_{\text{cut}} - k) \downarrow \langle \alpha | \beta \rangle_b \times \langle O | (PO) \rangle_b$$

$\alpha) \sigma(\alpha)$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\alpha') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

bar

$\langle 0 | P O | 0 \rangle_a$

$$\approx \int d\alpha \int d\beta \frac{\langle \alpha | P O | \beta \rangle_b}{\langle \beta | \alpha \rangle_b}$$

$$\langle 0 | P O | 0 \rangle_a = \int d^3 k e^{i k x} \theta(k_{cut} - k) \langle \tilde{b} | \alpha | \beta \rangle_b \times \langle 0 | (P O) | 0 \rangle_b$$

\Downarrow

$\tilde{b} \rightarrow \tilde{b}^+ \alpha$
 $\tilde{b}^+ \rightarrow \tilde{b}^+ \beta$

$\alpha) \sigma(\alpha)$

$$\hat{\sigma}_H = \hat{\sigma}_I + \int G_R(x, x') \Gamma(\sigma') dx'$$

$$= \hat{\sigma}_I + \hat{\sigma}_I^2 + \dots$$

bar

$\langle 0 | P O | 0 \rangle_a$

$$\approx \int d\alpha \int d\beta \frac{\langle \alpha | P O | \beta \rangle_b}{\langle \beta | \alpha \rangle_b}$$

$$\langle \rangle = \int d^3 k e^{i k x} \theta(k_{cut} - k)$$

$$\langle \alpha | \beta \rangle_b \times \langle 0 | (PO) | 0 \rangle_b$$

\Downarrow

$$\begin{aligned} b &\rightarrow b + \alpha \\ \tilde{b}^+ &\rightarrow \tilde{b}^+ + \beta \end{aligned}$$

$$\frac{\int d\alpha \int d\beta \delta(\alpha - P_O) \delta(\beta - b)}{\int d\alpha \int d\beta \delta(\alpha - P_O) \delta(\beta - b)}$$



$$\langle \alpha | \beta \rangle \times \langle \delta(P_O) | \rangle$$
$$\tilde{b} \rightarrow \tilde{b} + \alpha$$
$$\tilde{b}^+ \rightarrow \tilde{b}^+ + \beta$$

$$\approx \frac{\int d\alpha \int d\beta \zeta(\alpha) |PO| \beta^{\gamma_b}}{\int d\beta \zeta(\beta) |PO|}$$



$$ut - k) \zeta(\alpha | \beta) \times \zeta(O | (PO) | \beta^{\gamma_b})$$

$$\tilde{b} \rightarrow \tilde{b} + \alpha$$

$$\tilde{b}^+ \rightarrow \tilde{b}^+ + \beta$$

$$\approx \frac{\int d\alpha \int d\beta \langle Q | PO | \beta \rangle_b}{\int d\alpha \int d\beta \langle Q | PO | \alpha \rangle_b}$$



$$ut - k) \langle Q | \beta \rangle_b \times \langle O | (PO) | \alpha \rangle_b$$

$\tilde{b} \rightarrow \tilde{b} + \alpha$
 $\tilde{b}^+ \rightarrow \tilde{b}^+ + \beta$

Urakawa

Coherent state

$$\hat{b}^\dagger |\beta\rangle = \beta |\beta\rangle$$

$$|0\rangle_a = \int_{-\infty}^{\infty} d\beta \sqrt{\frac{s}{\pi}} e^{-(s\beta)^2} |\beta\rangle_b$$

$$s \rightarrow 0$$

$$= \hat{G}_I + \hat{G}_J + \dots$$

$$\langle 0 | P \Theta | 0 \rangle_a$$

$$\approx \int d\alpha \int d\beta \frac{\langle \alpha | P \Theta | \beta \rangle_b}{b}$$

↓

$$_{cut-k} \langle \tilde{b} | \alpha | \beta \rangle_{\tilde{b}} \times \langle 0 | (P \Theta) | 0 \rangle_{\tilde{b}}$$

$\tilde{b} \rightarrow \tilde{b}^+ \alpha$
 $\tilde{b}^+ \rightarrow \tilde{b}^+ \beta$

$$P \propto \exp\left(-\frac{\sigma^2}{2(\Delta\sigma)^2}\right)$$



$$P \propto \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\sigma)^2}\right)$$

↳ broadened wave function

$$\tilde{\sigma} > b + b^\dagger$$

$$P \propto \exp\left(-\frac{\bar{\sigma}^2}{2(\Delta\sigma)^2}\right) \rightarrow \exp\left(-\frac{(\alpha+\beta)^2}{2(\Delta\sigma)^2}\right)$$

$$\tilde{\sigma} \rightarrow \tilde{b} + \tilde{b}^\dagger \approx \alpha + \beta$$

$$\langle \alpha | \beta \rangle \propto \exp\left(-\frac{(\alpha-\beta)^2}{2}\right)$$

$$P \propto \exp\left(-\frac{\tilde{\sigma}^2}{2(\Delta\sigma)^2}\right) \hookrightarrow \exp\left(-\frac{(\alpha+\beta)^2}{2(\Delta\sigma)^2}\right)$$

$$\tilde{G} \supset \tilde{b} + \tilde{b}^\dagger \propto \alpha + \beta$$

$$\langle \alpha | \beta \rangle \propto \exp\left(-\frac{(\alpha-\beta)^2}{2}\right)$$

$$P \propto \exp\left(-\frac{\tilde{\sigma}^2}{2(\Delta\sigma)^2}\right) \rightarrow \exp\left(-\frac{(\alpha+\beta)^2}{2(\Delta\sigma)^2}\right)$$

$$\tilde{G} \rightarrow \tilde{b} + \tilde{b}^+ \approx \alpha + \beta$$

$$\langle \alpha | \beta \rangle \propto \exp\left(-\frac{(\beta)^2}{2}\right)$$

$(u, a) \xrightarrow{\text{unitary}}$

b)

$$v = \sum_{p \neq 0} \dots$$

$$+ \sum_{p \neq 0} \dots$$

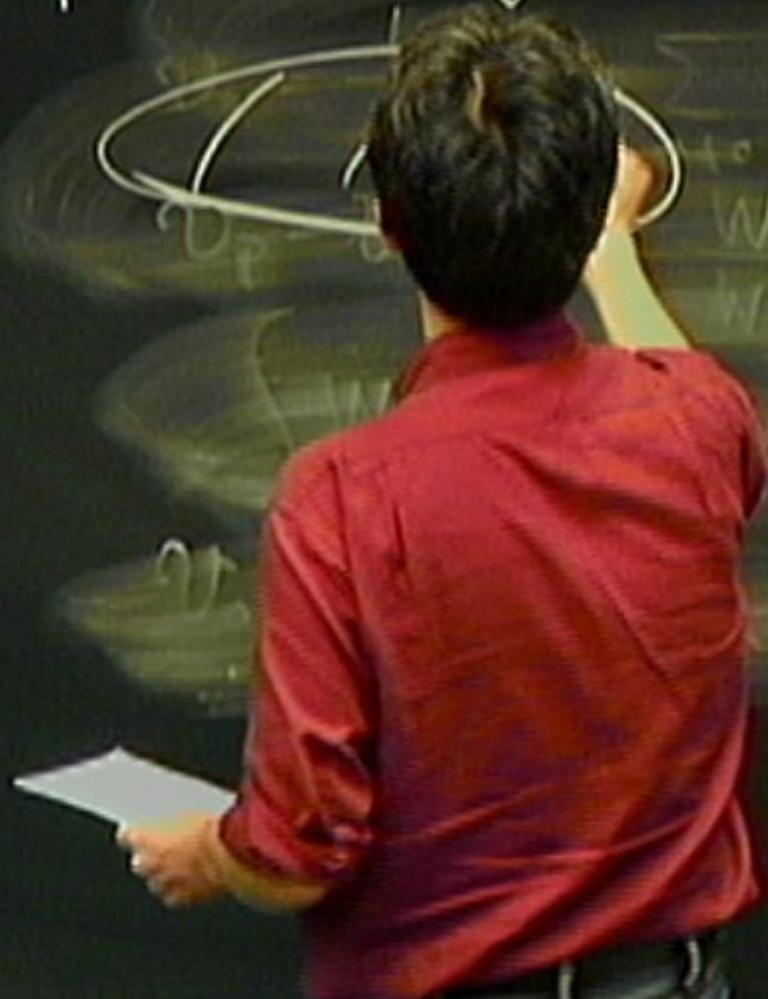
$$\frac{1}{C(p)} = \lim_{\eta \rightarrow 0} U_p$$

$$= \frac{W_p}{W_0} \frac{C(0)}{C(p)} U_0$$

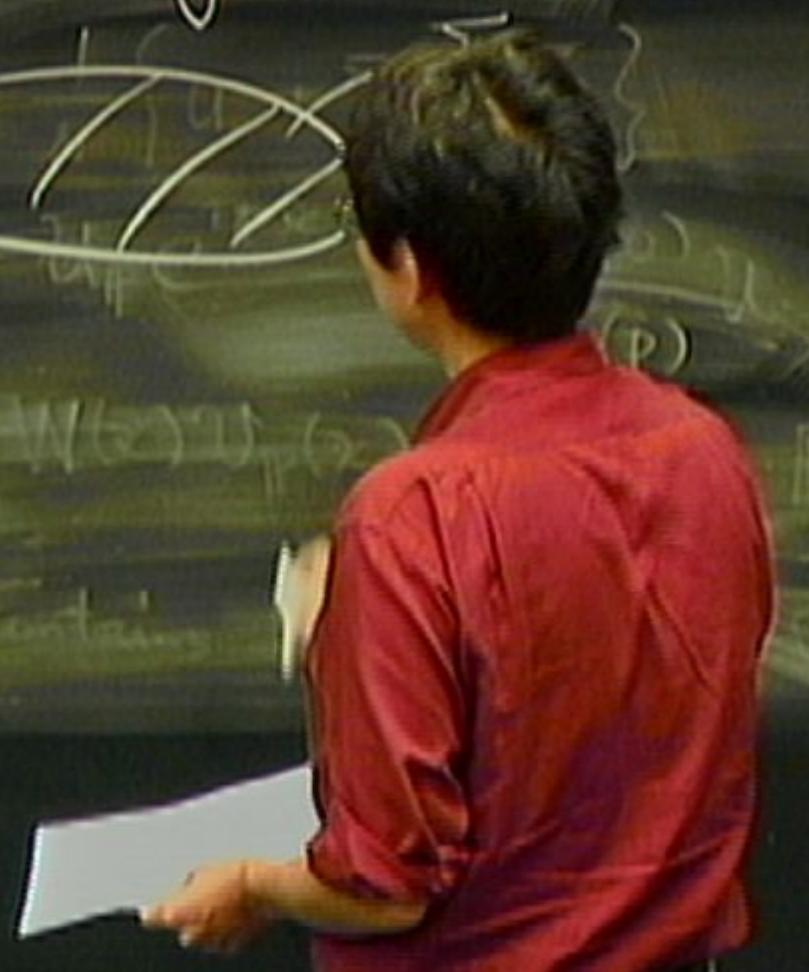
$$P(x) d^3x = D P \neq 0$$

information about \bar{J}

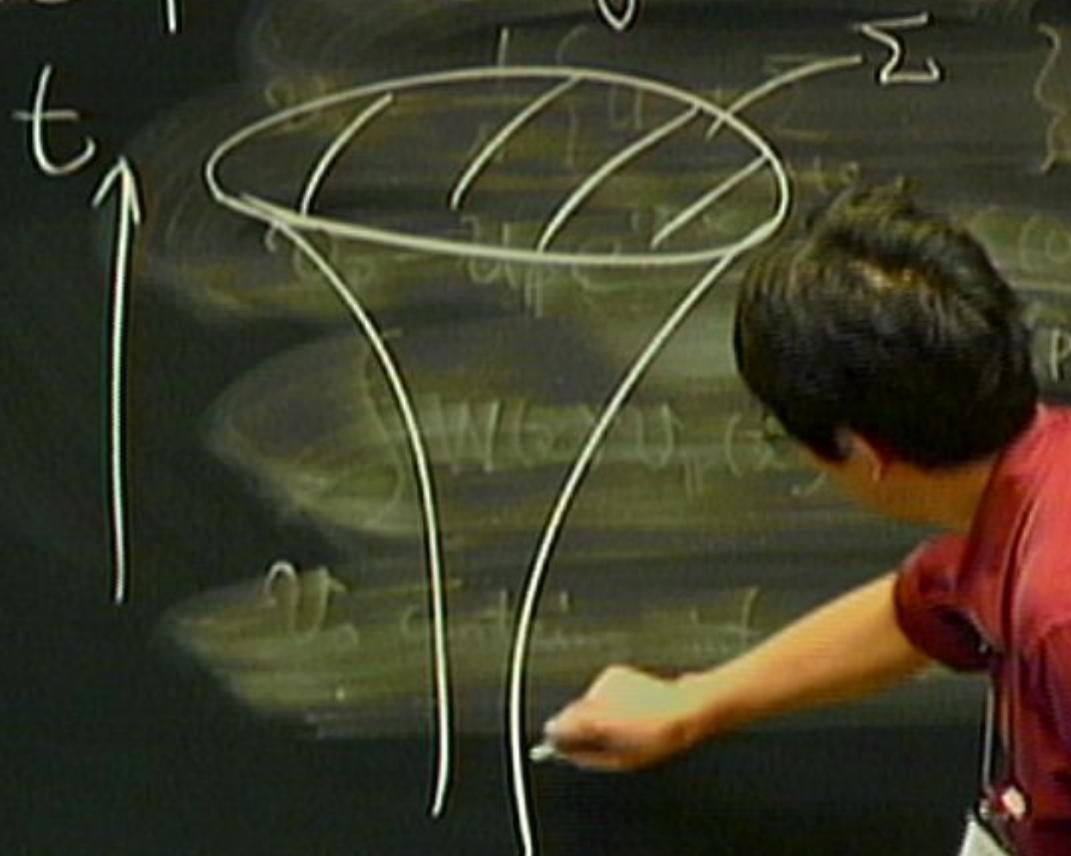
Temporal integral



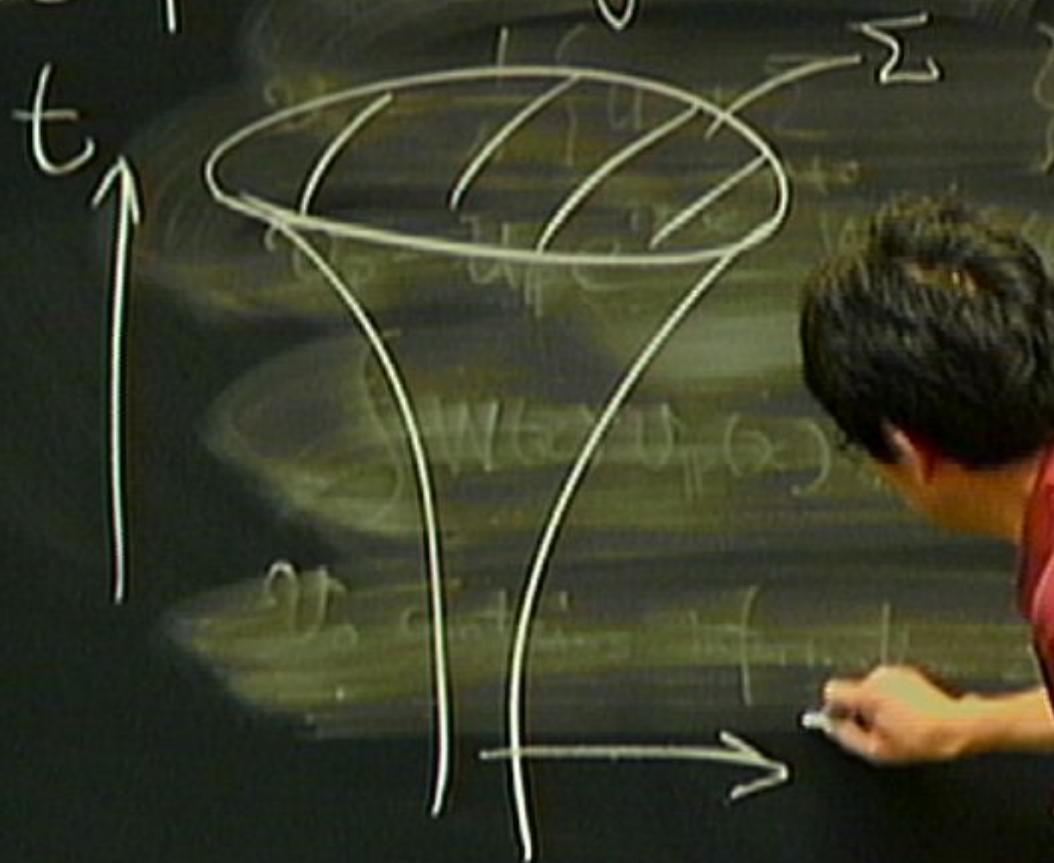
temporal integral



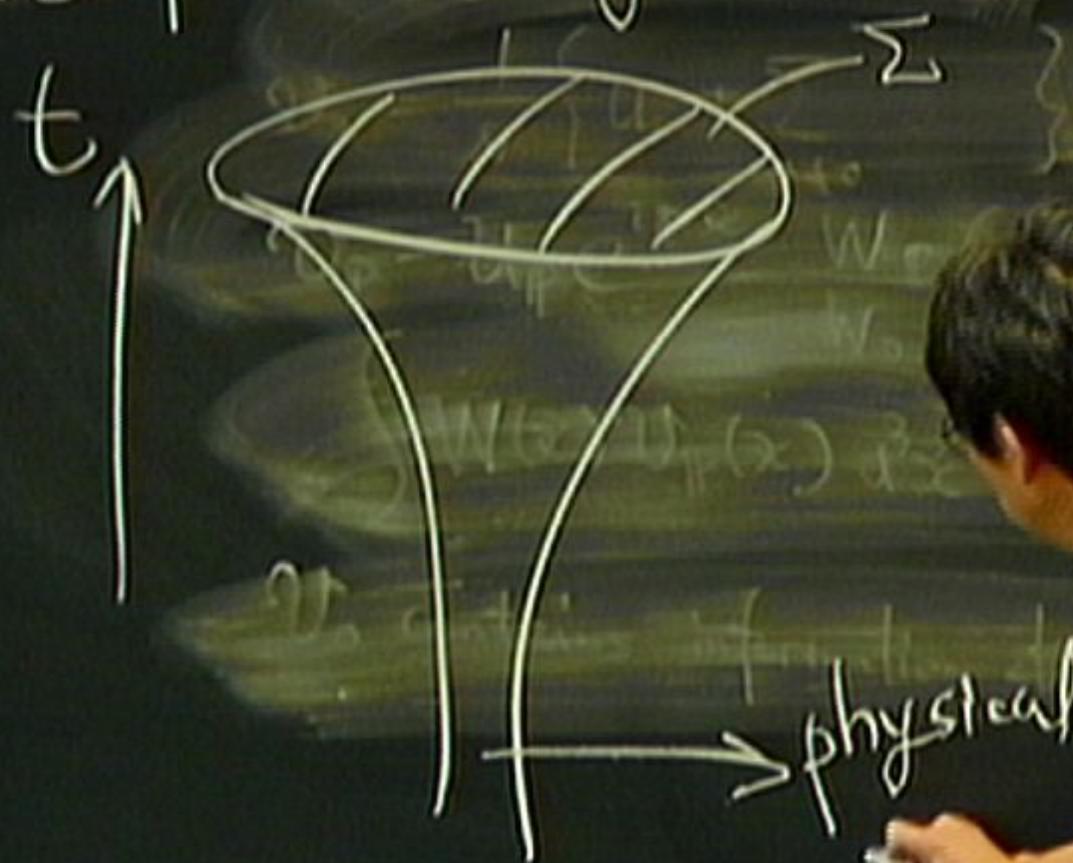
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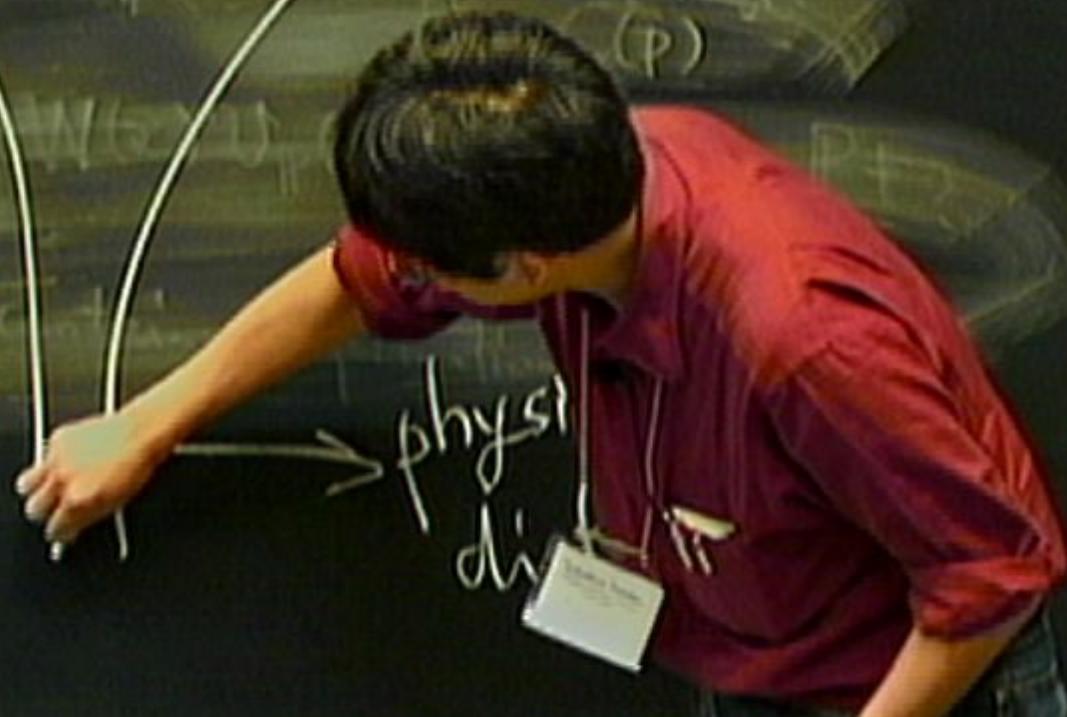
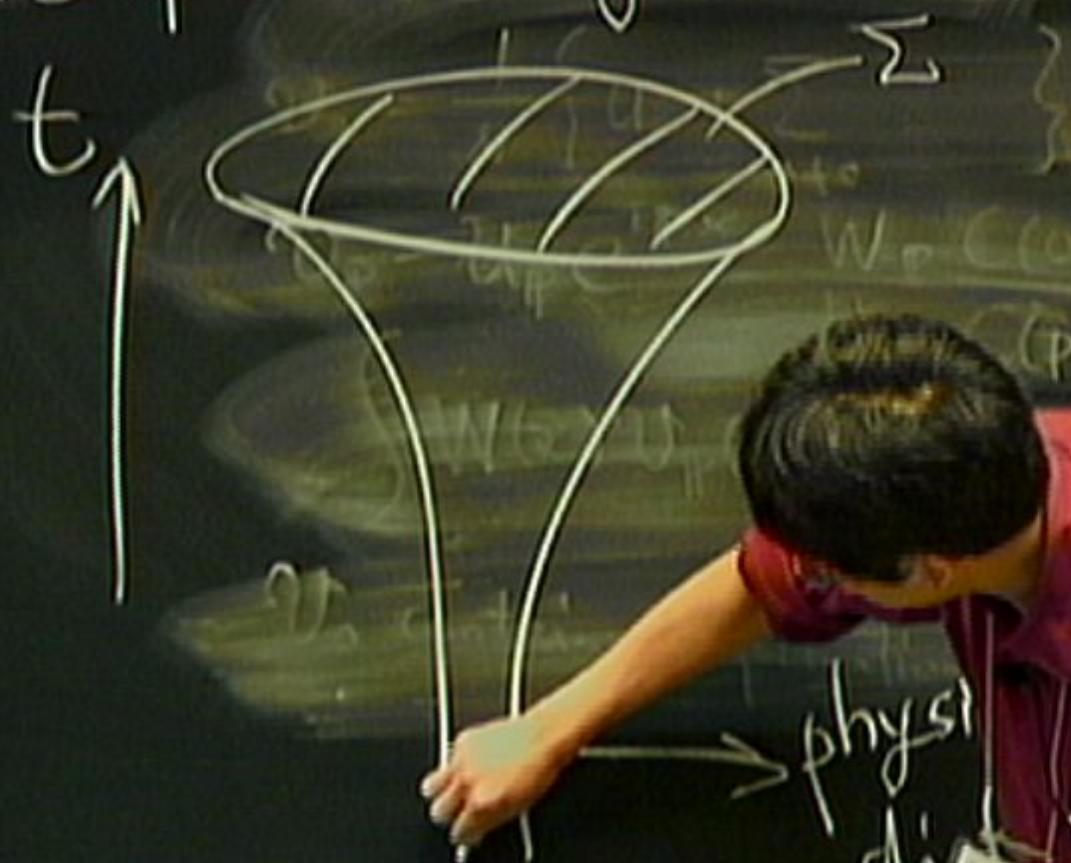
temporal integral



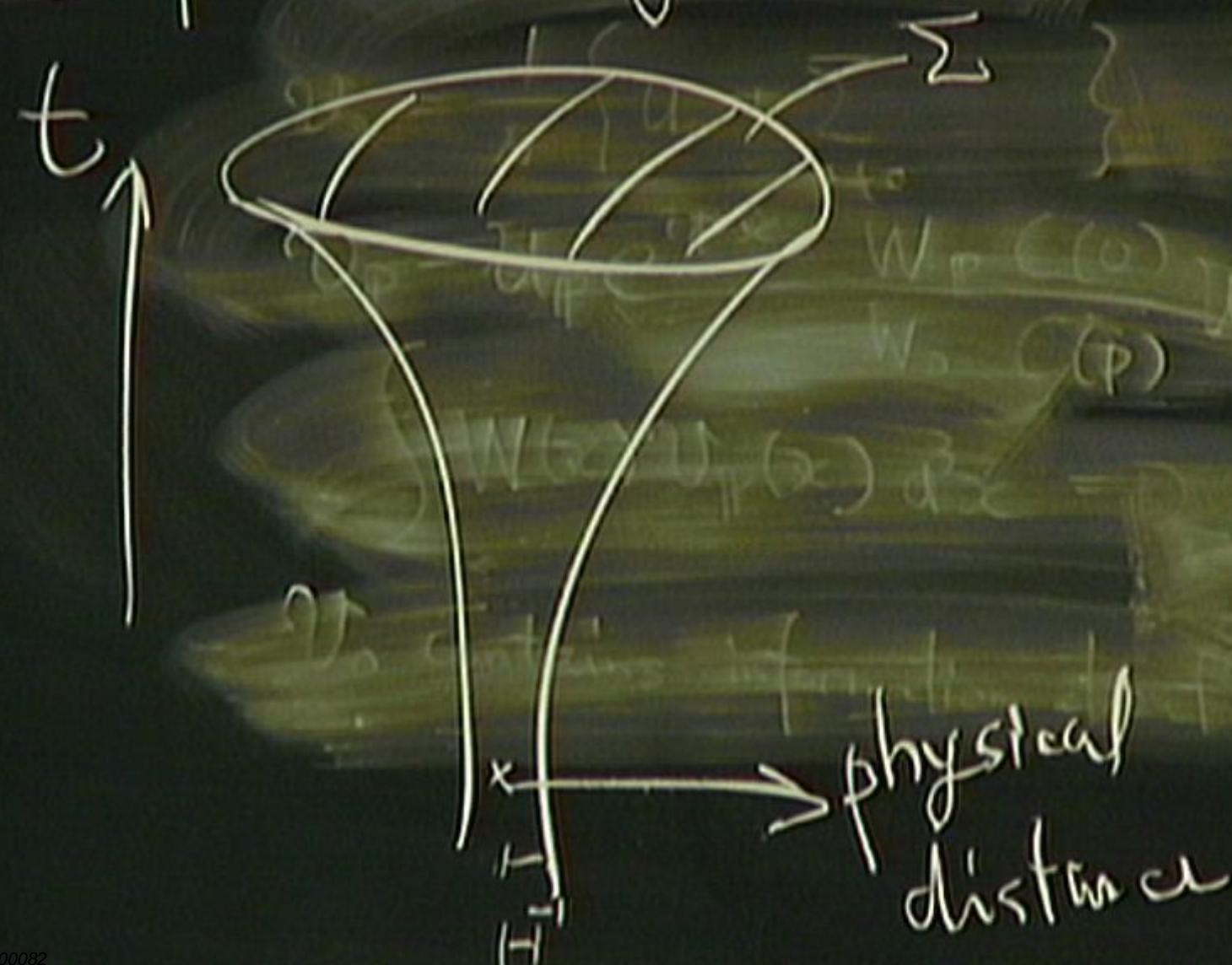
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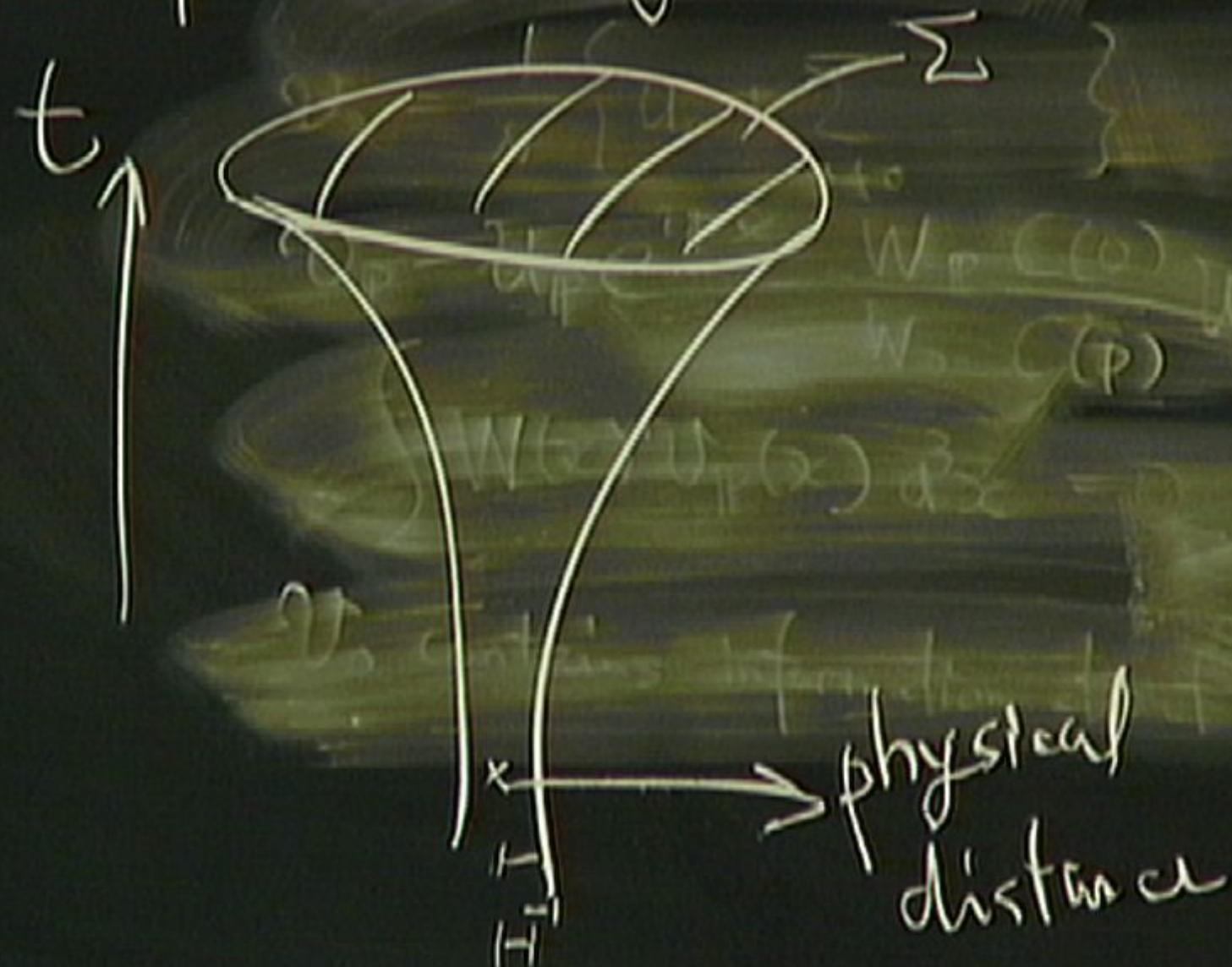
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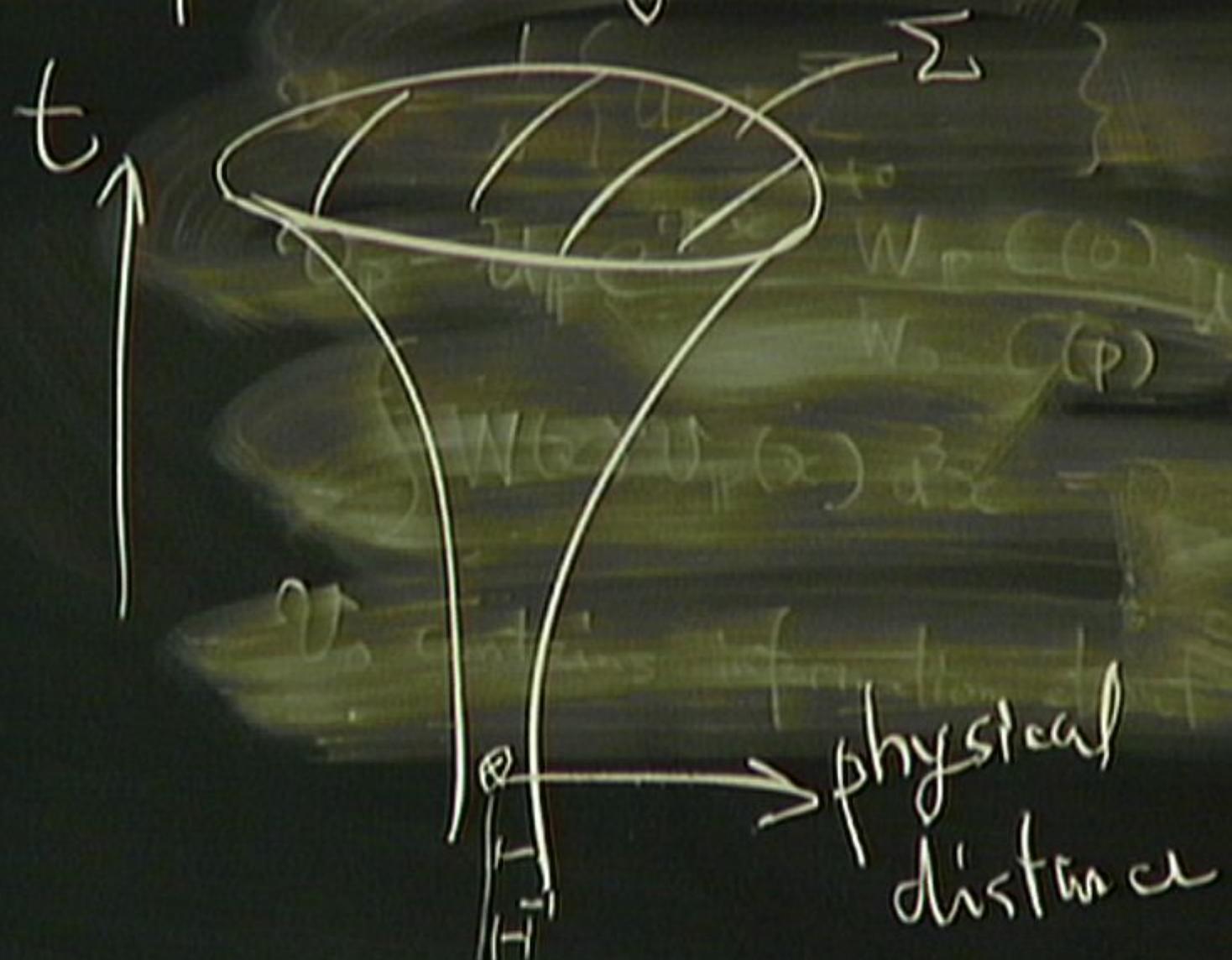
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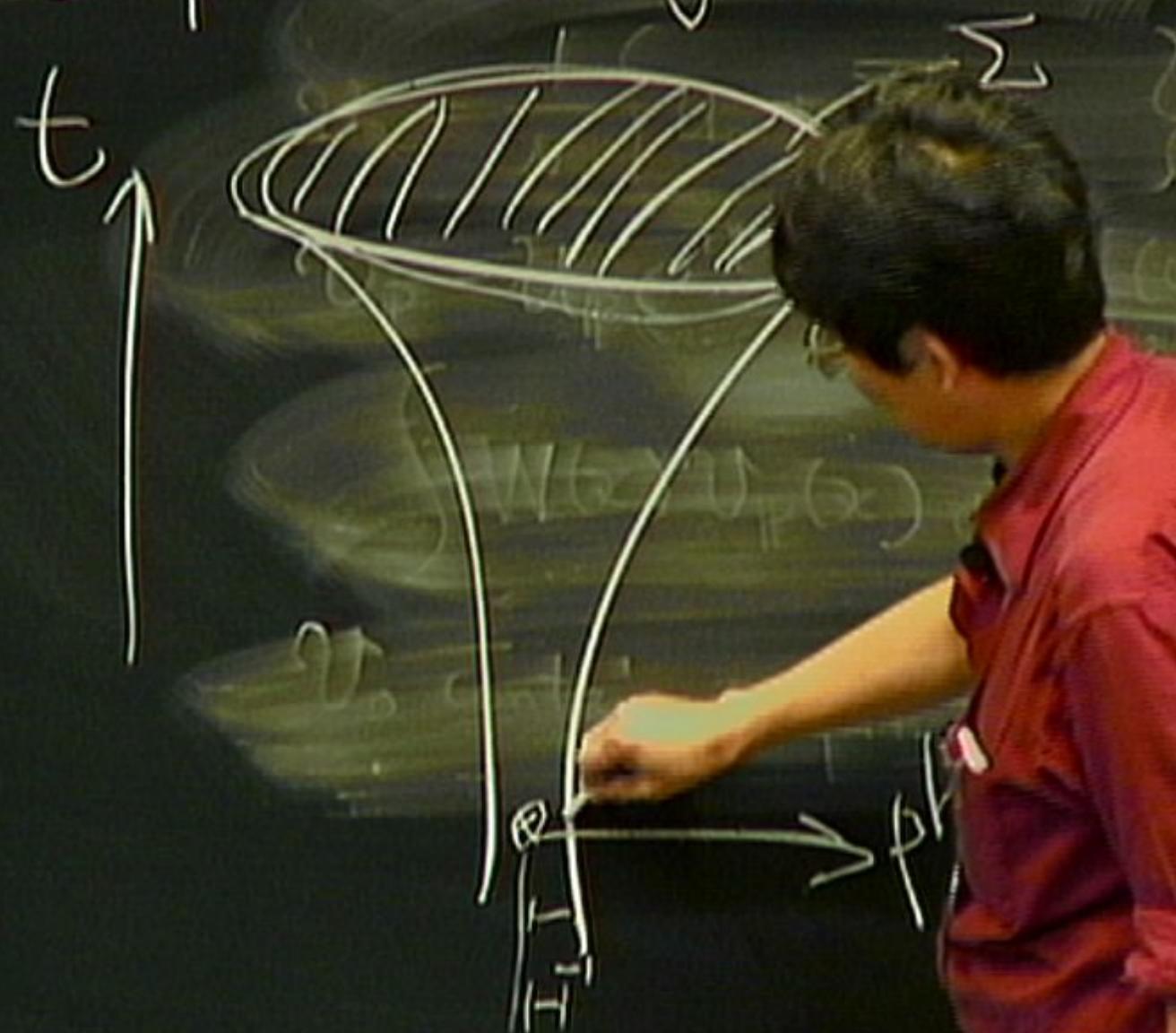
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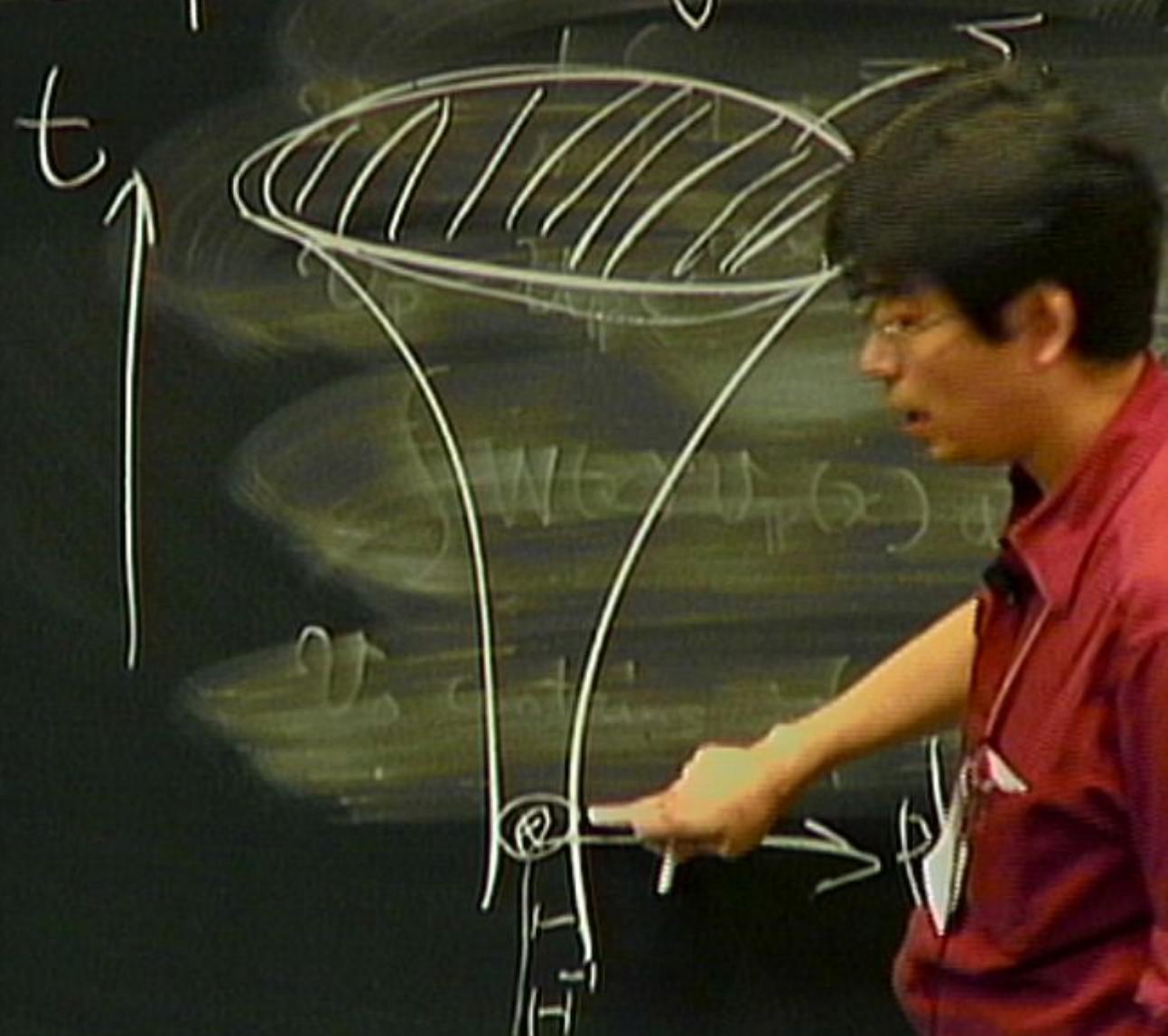
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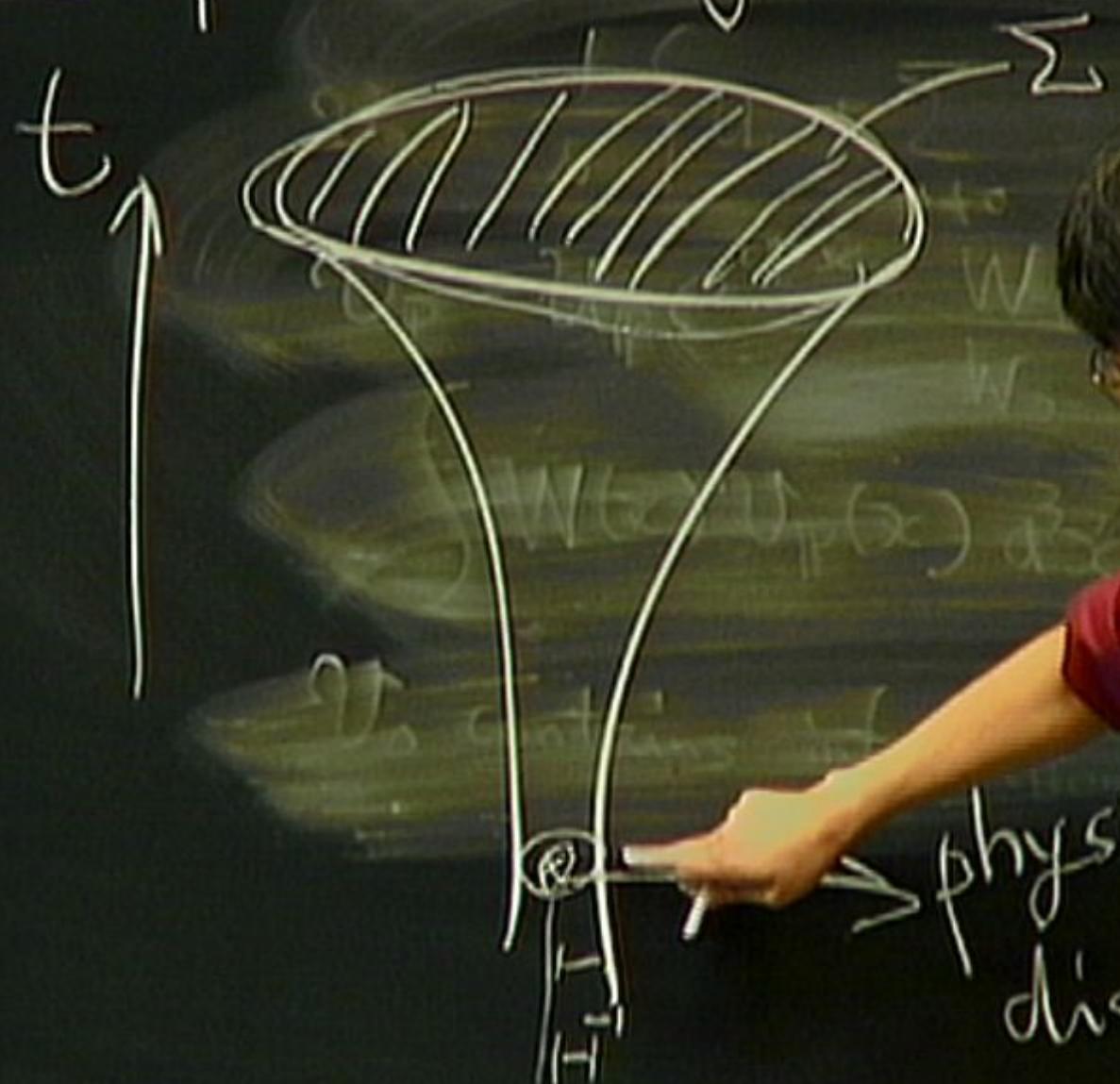
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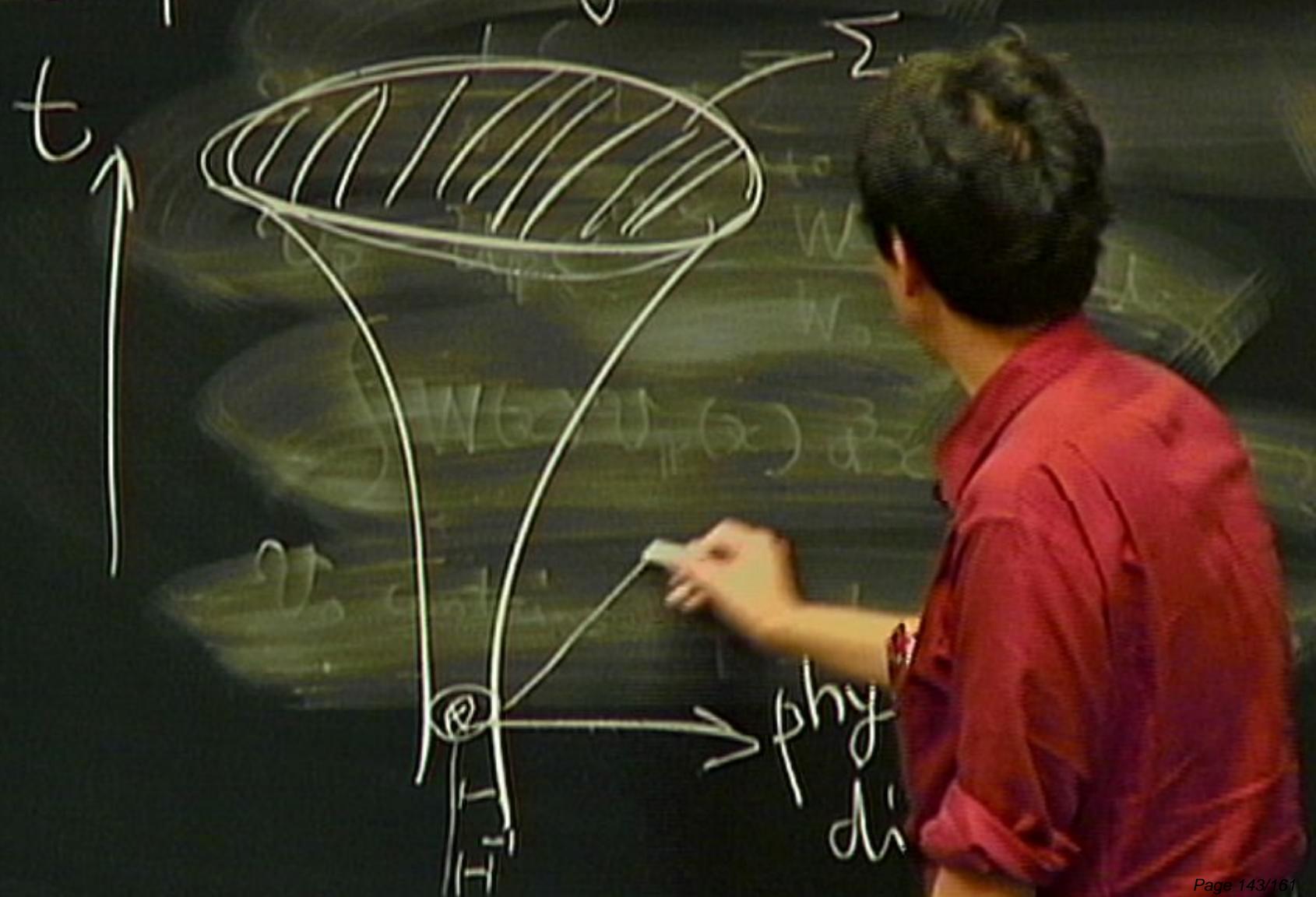
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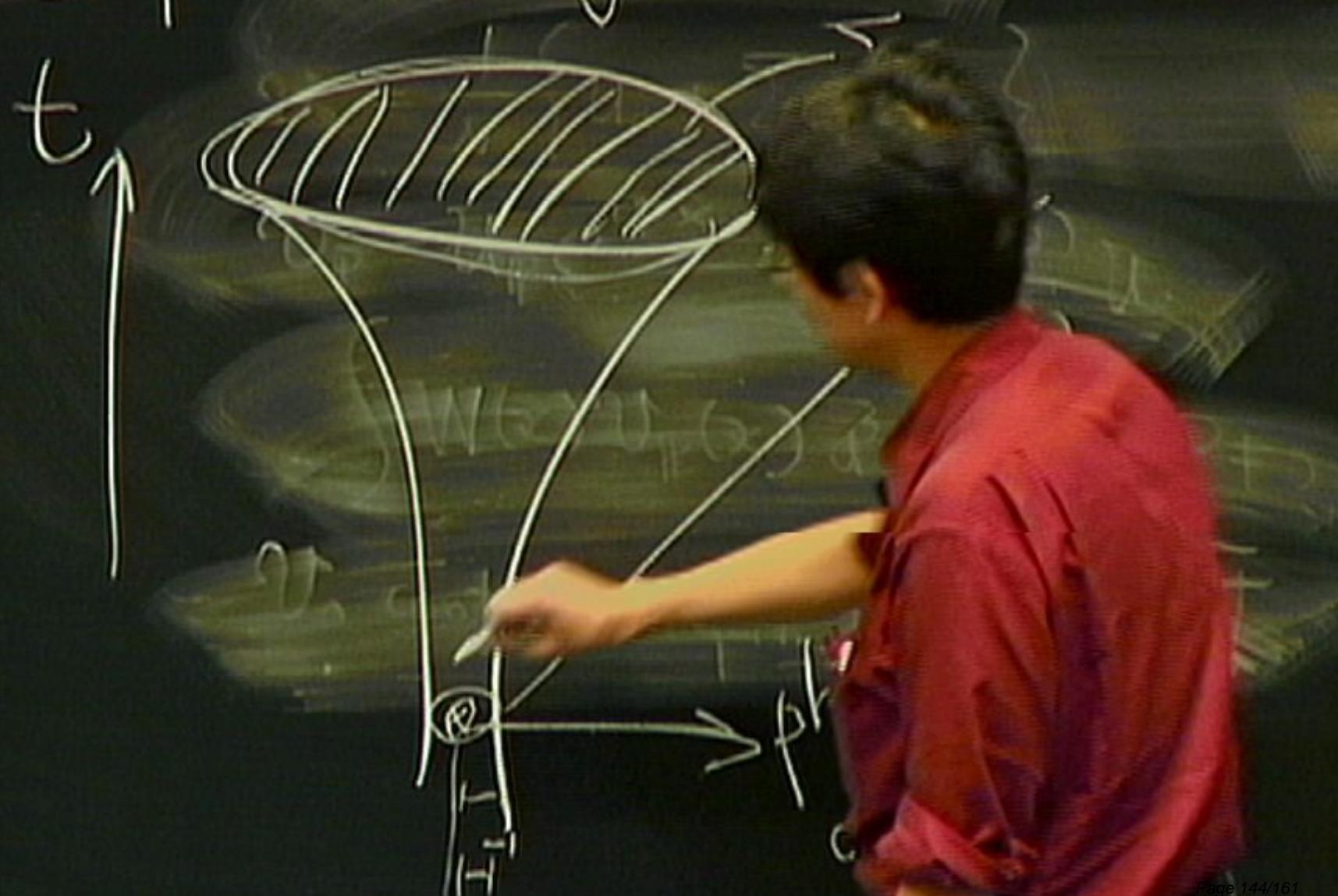
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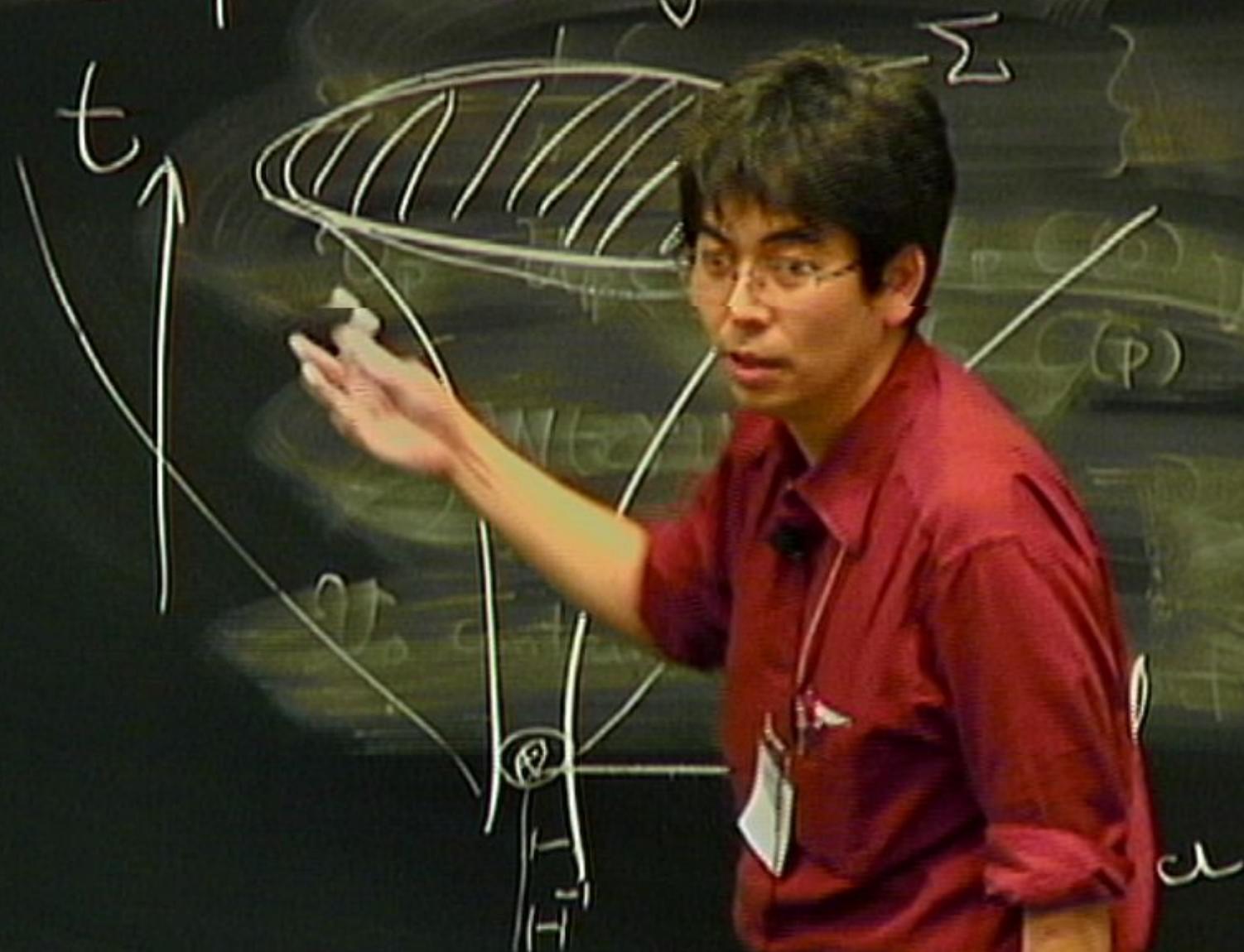
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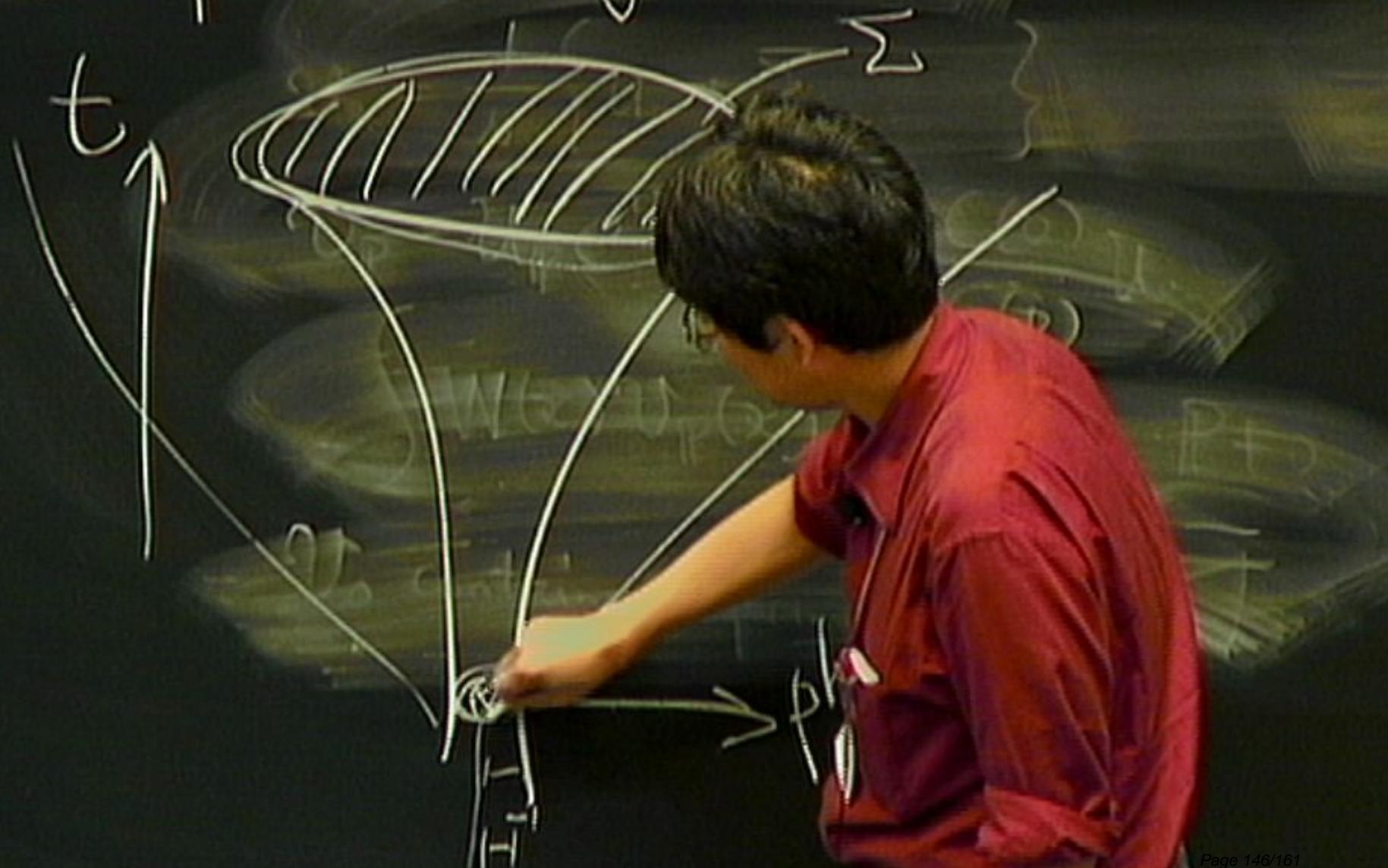
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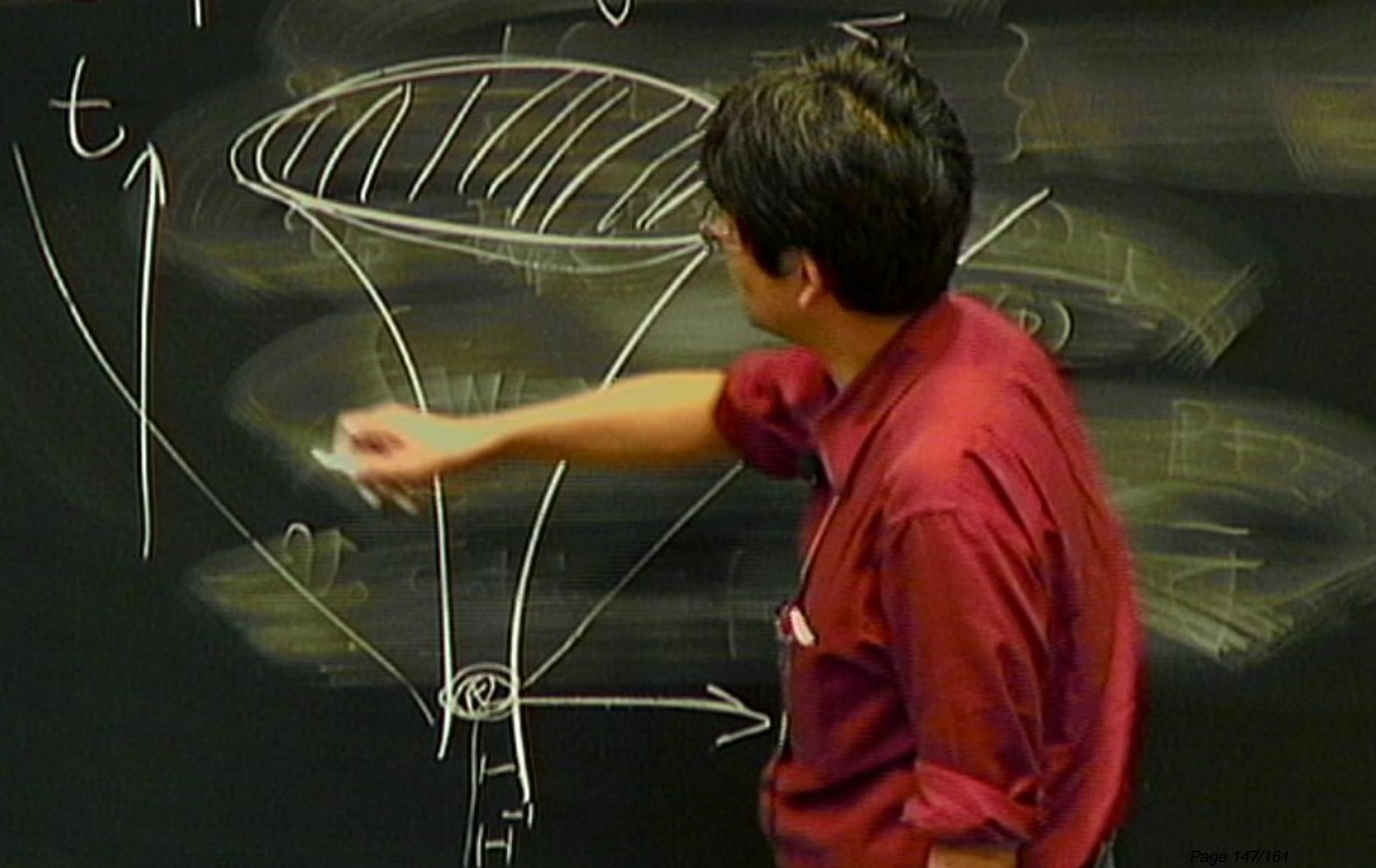
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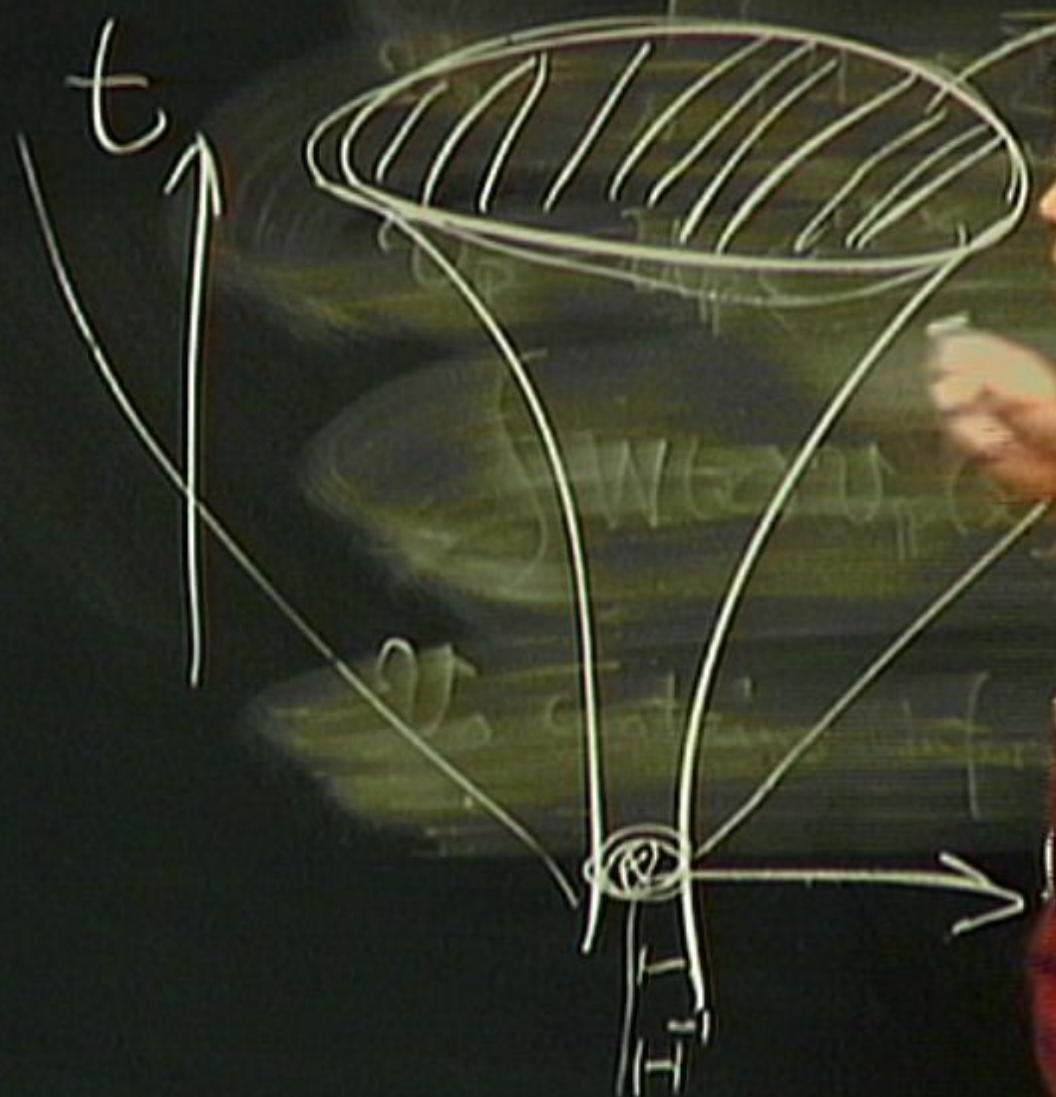
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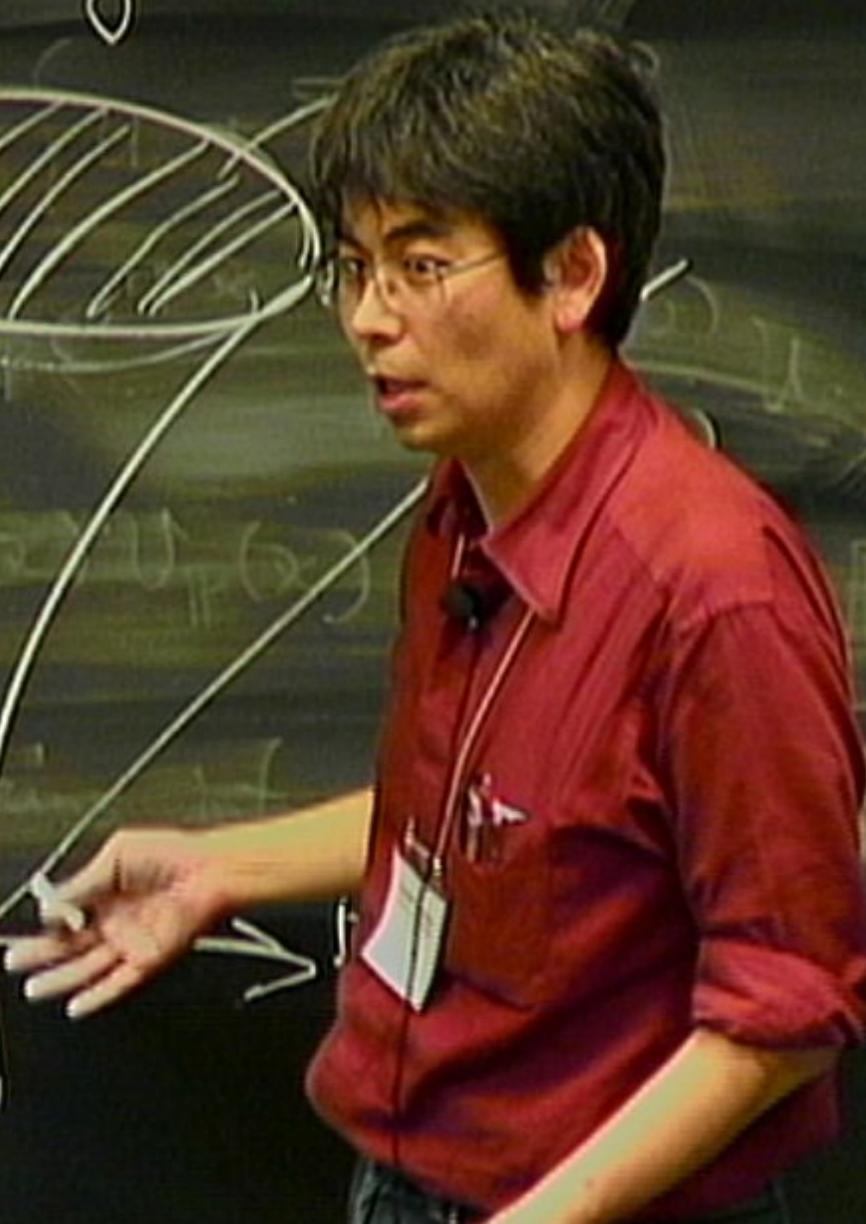
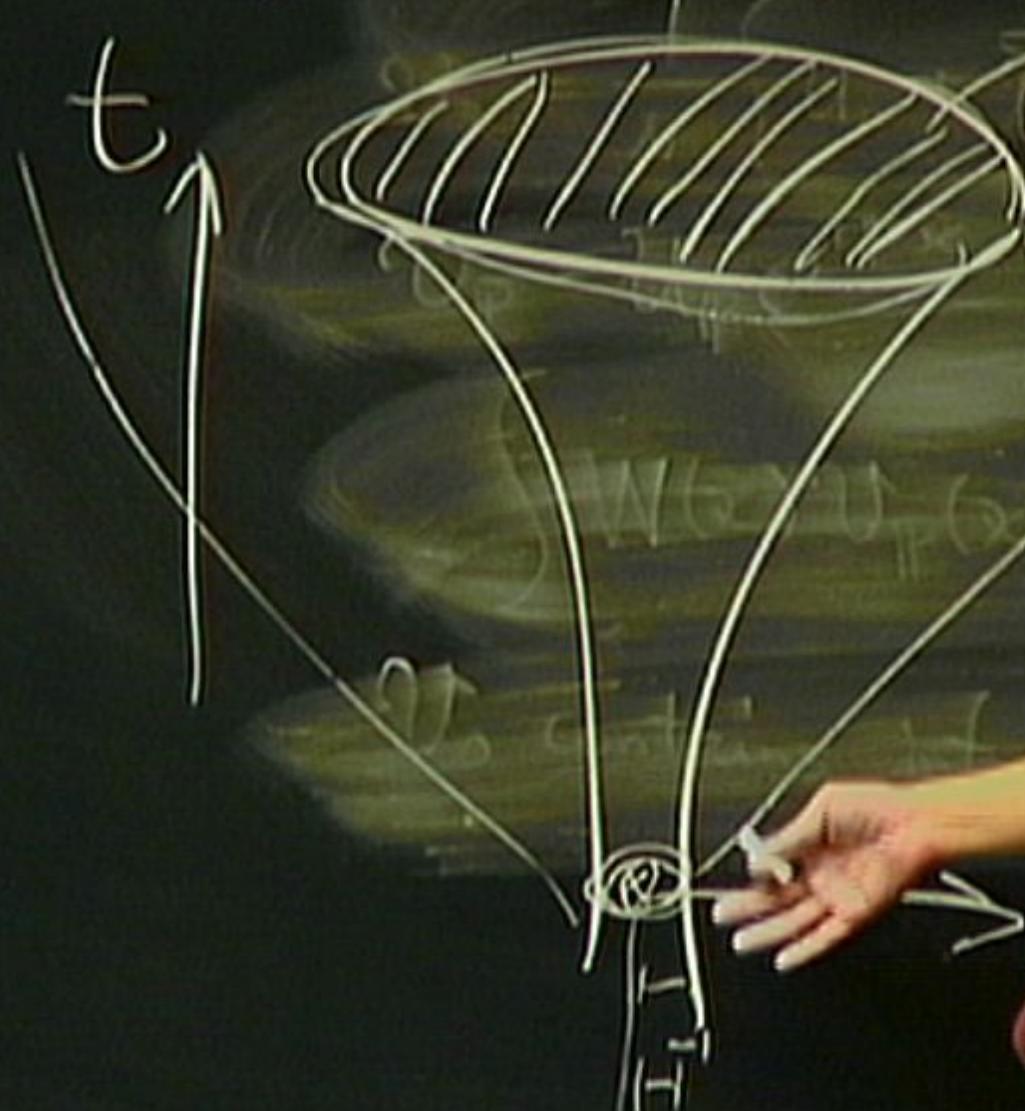
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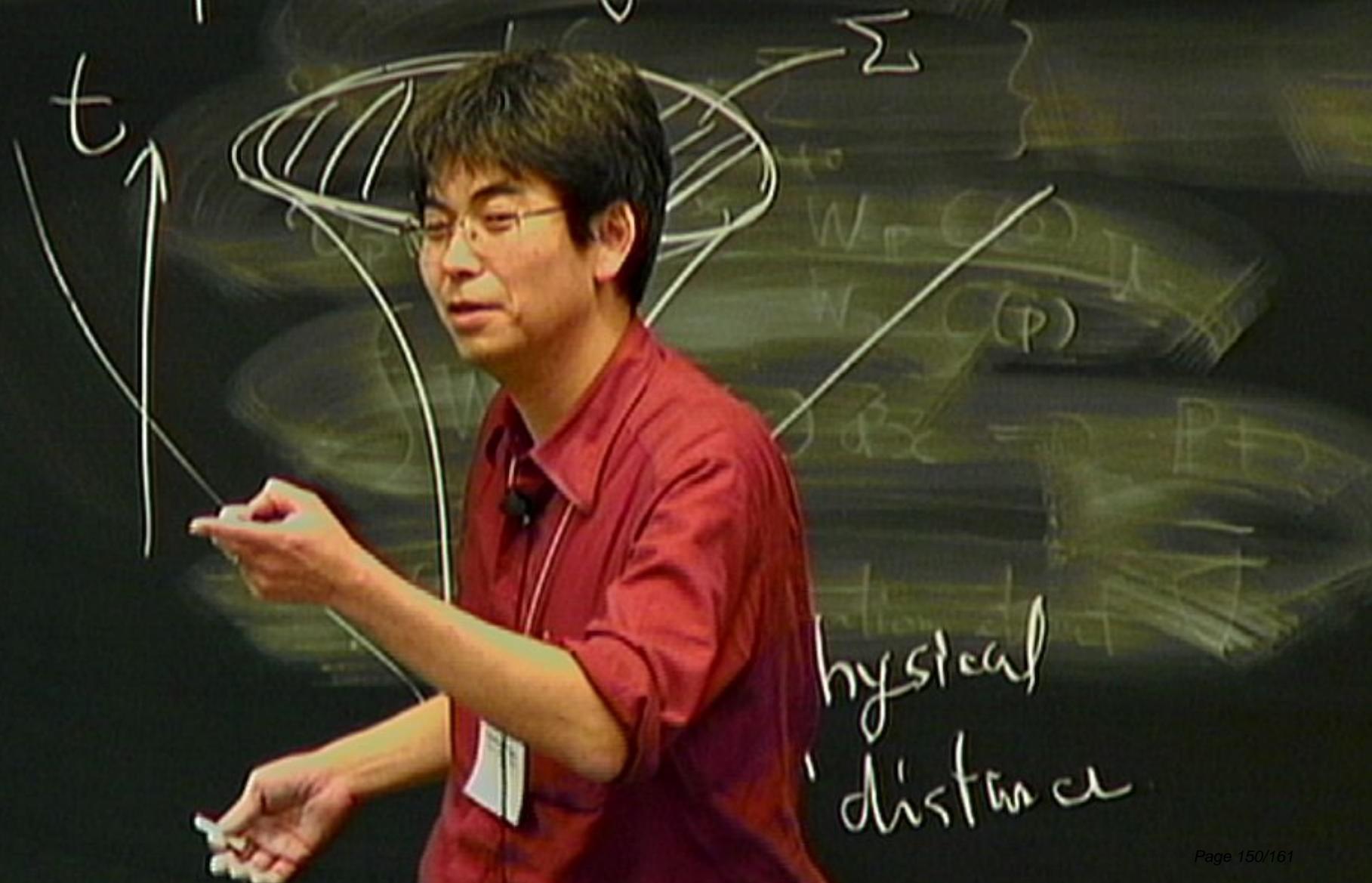
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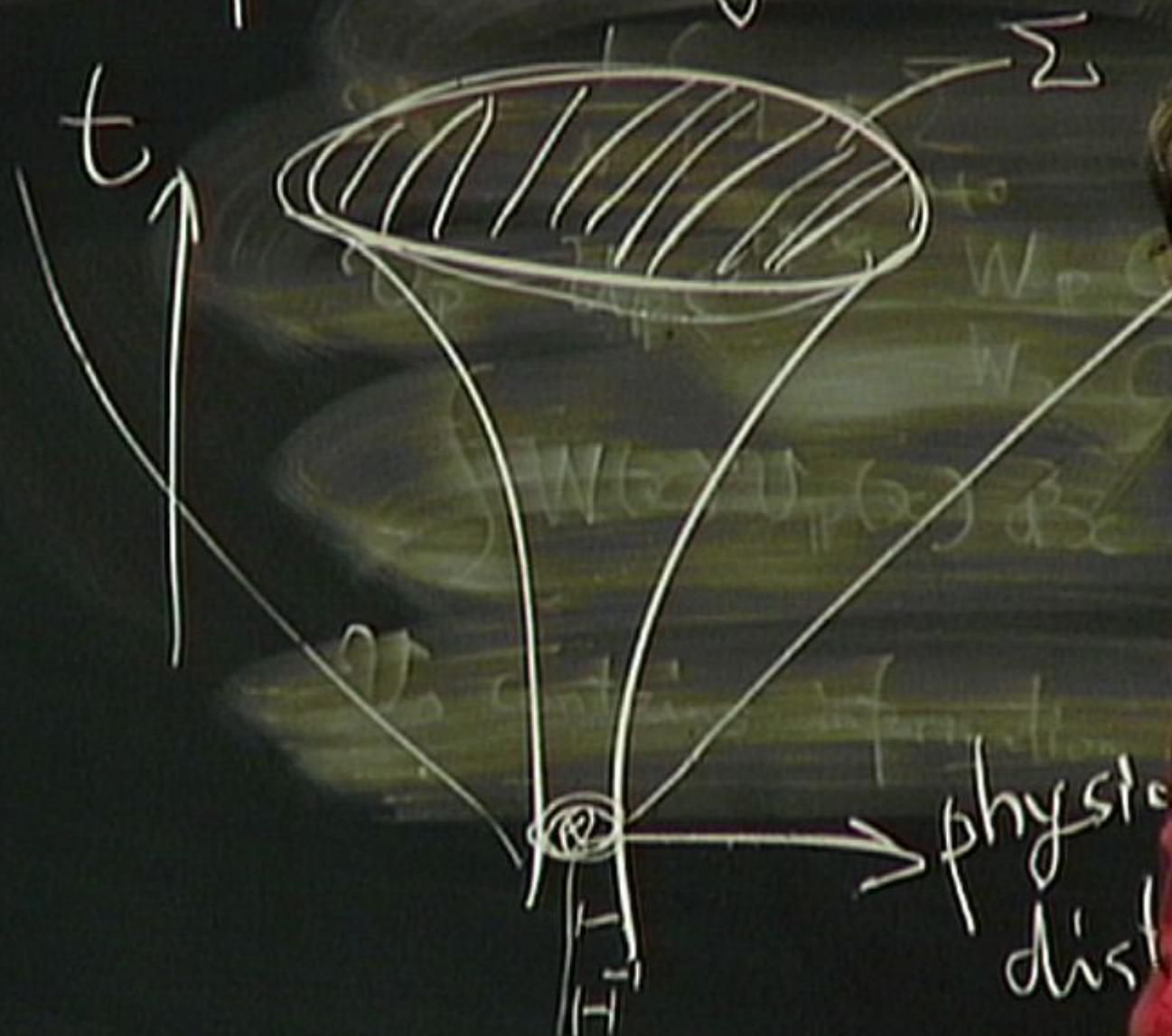
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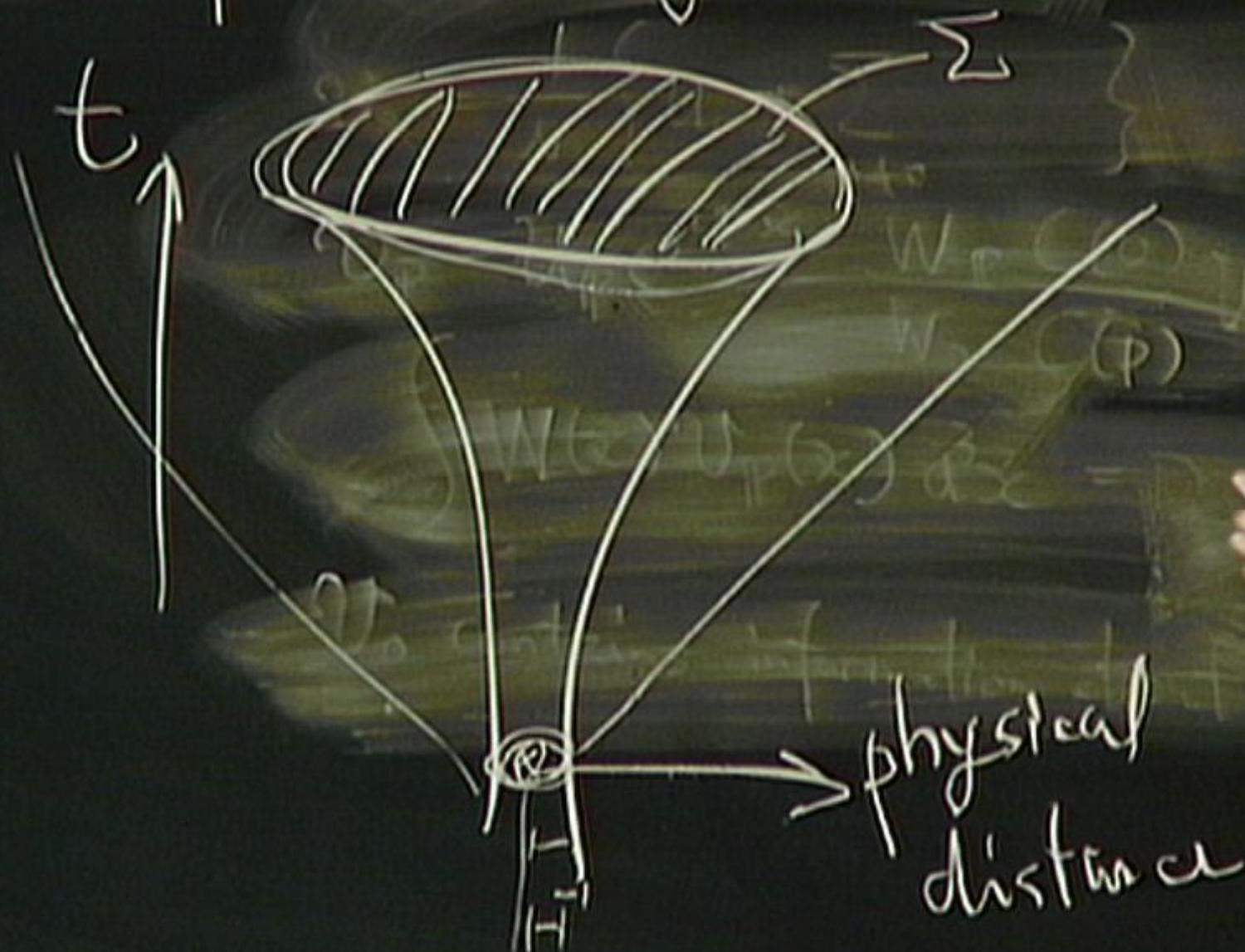
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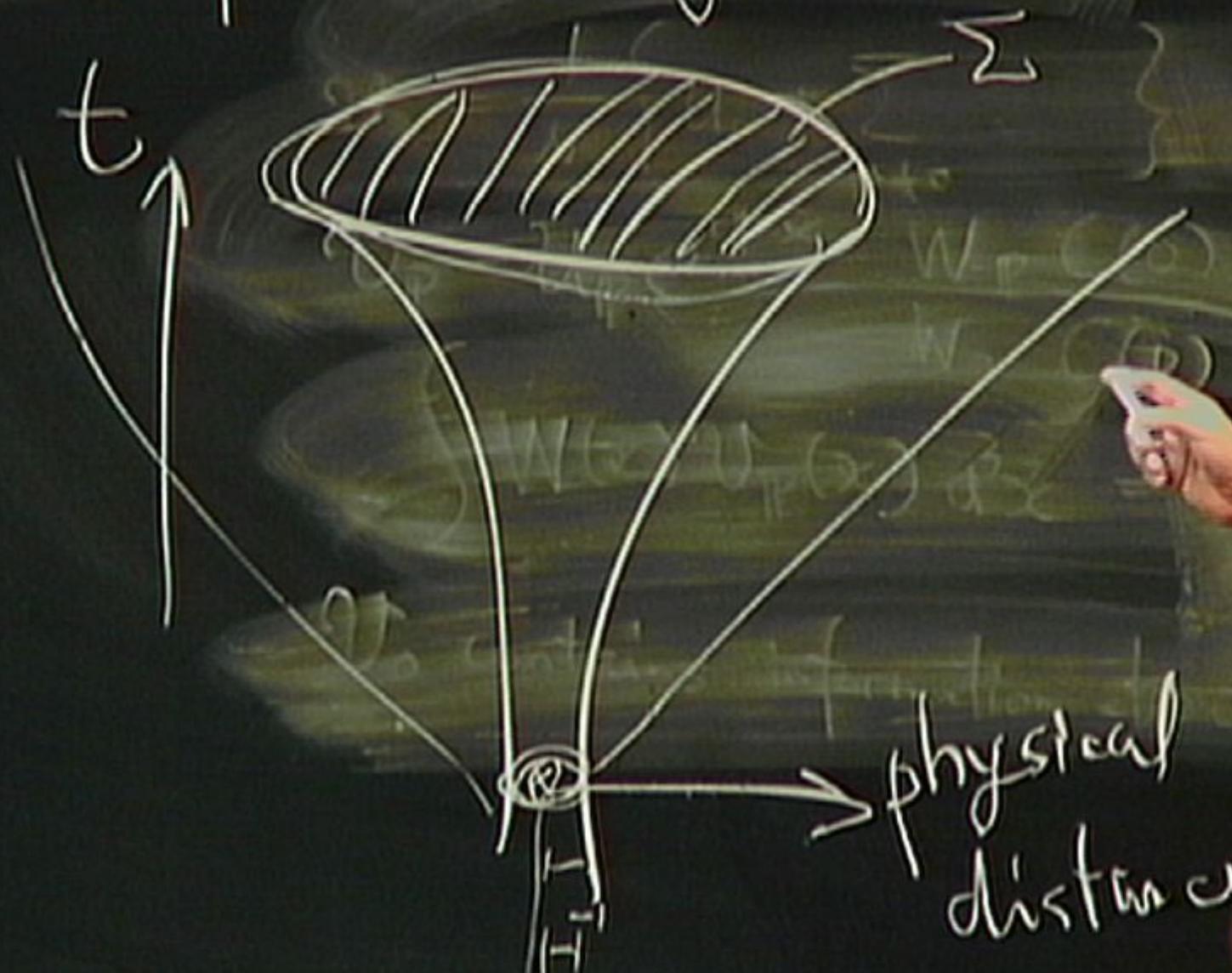
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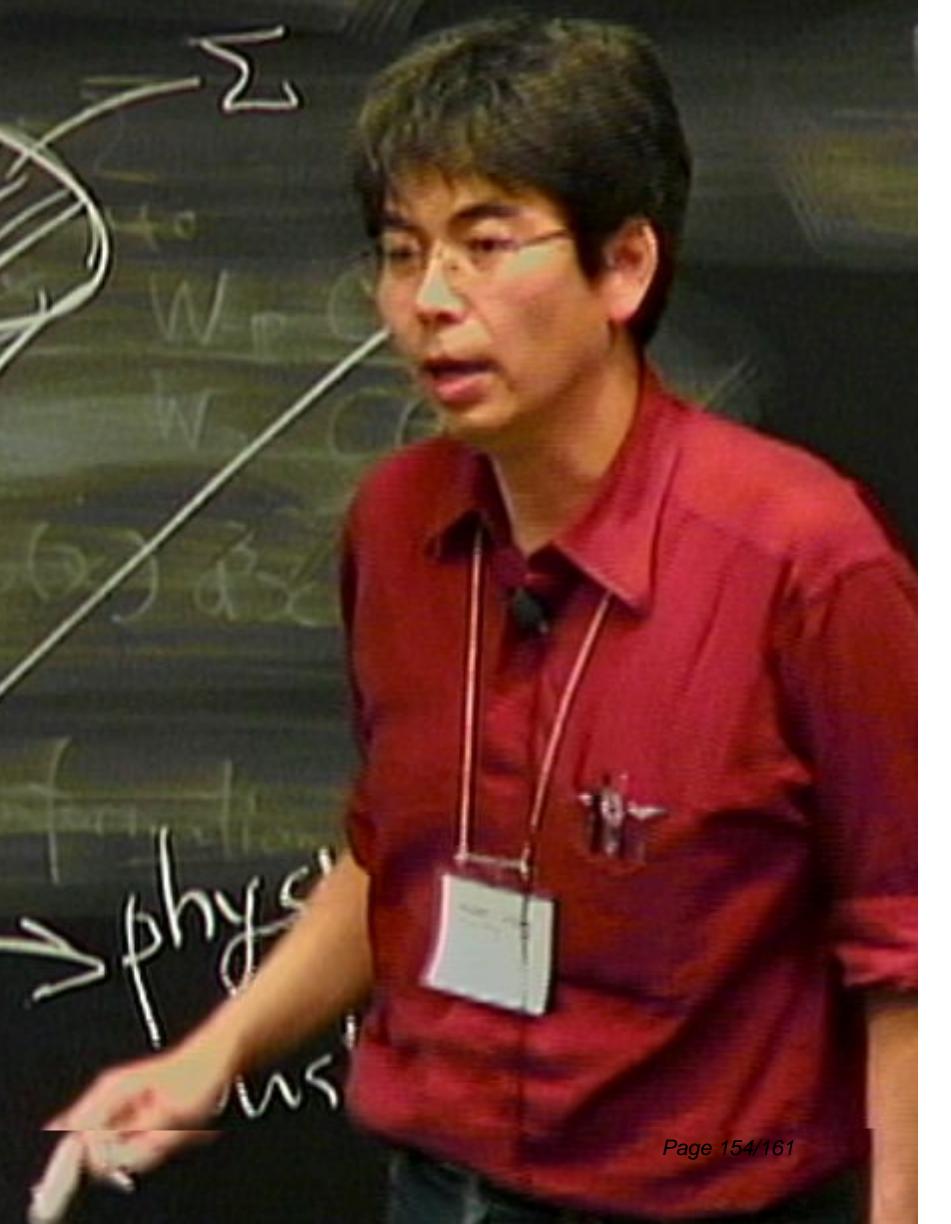
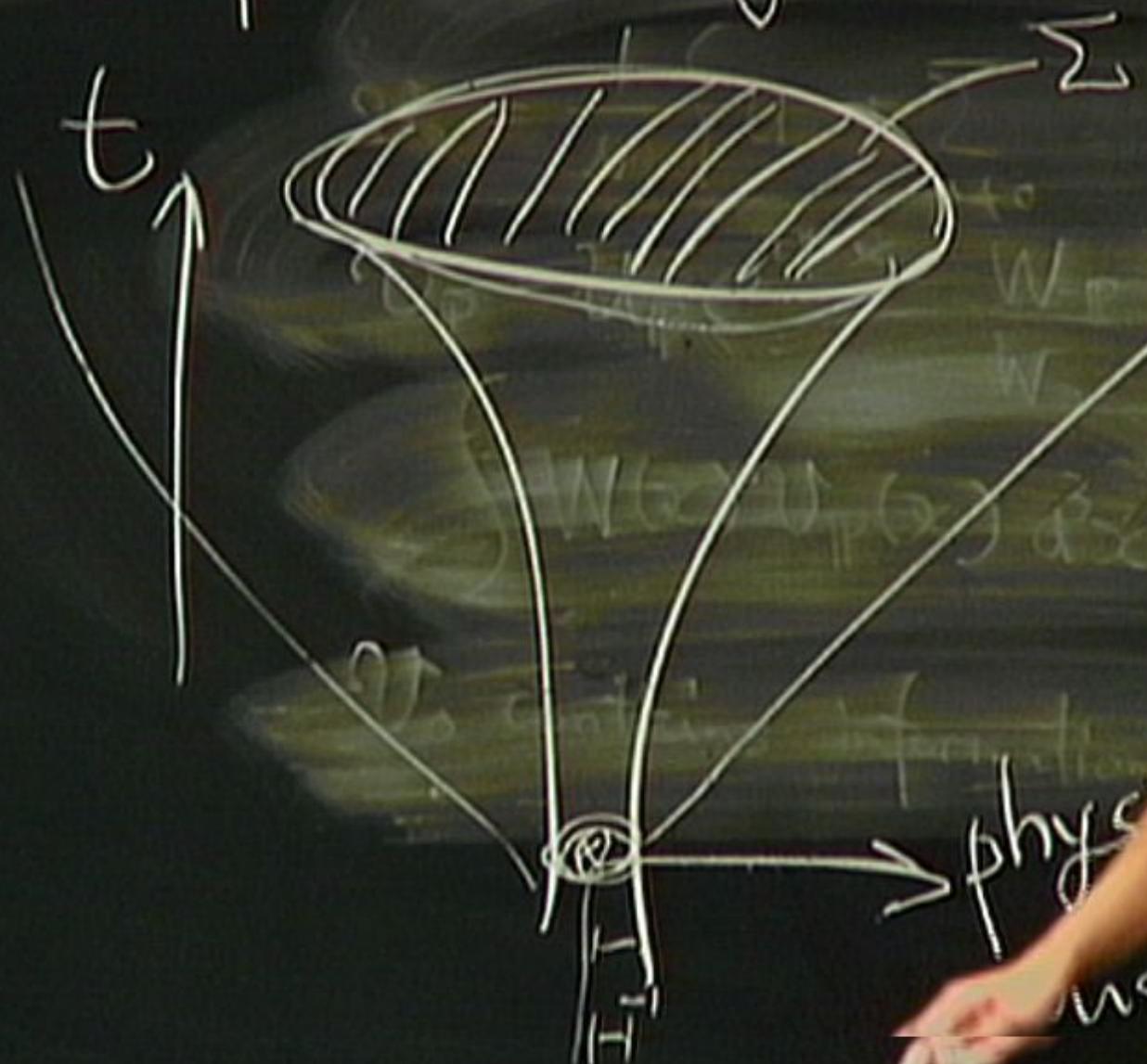
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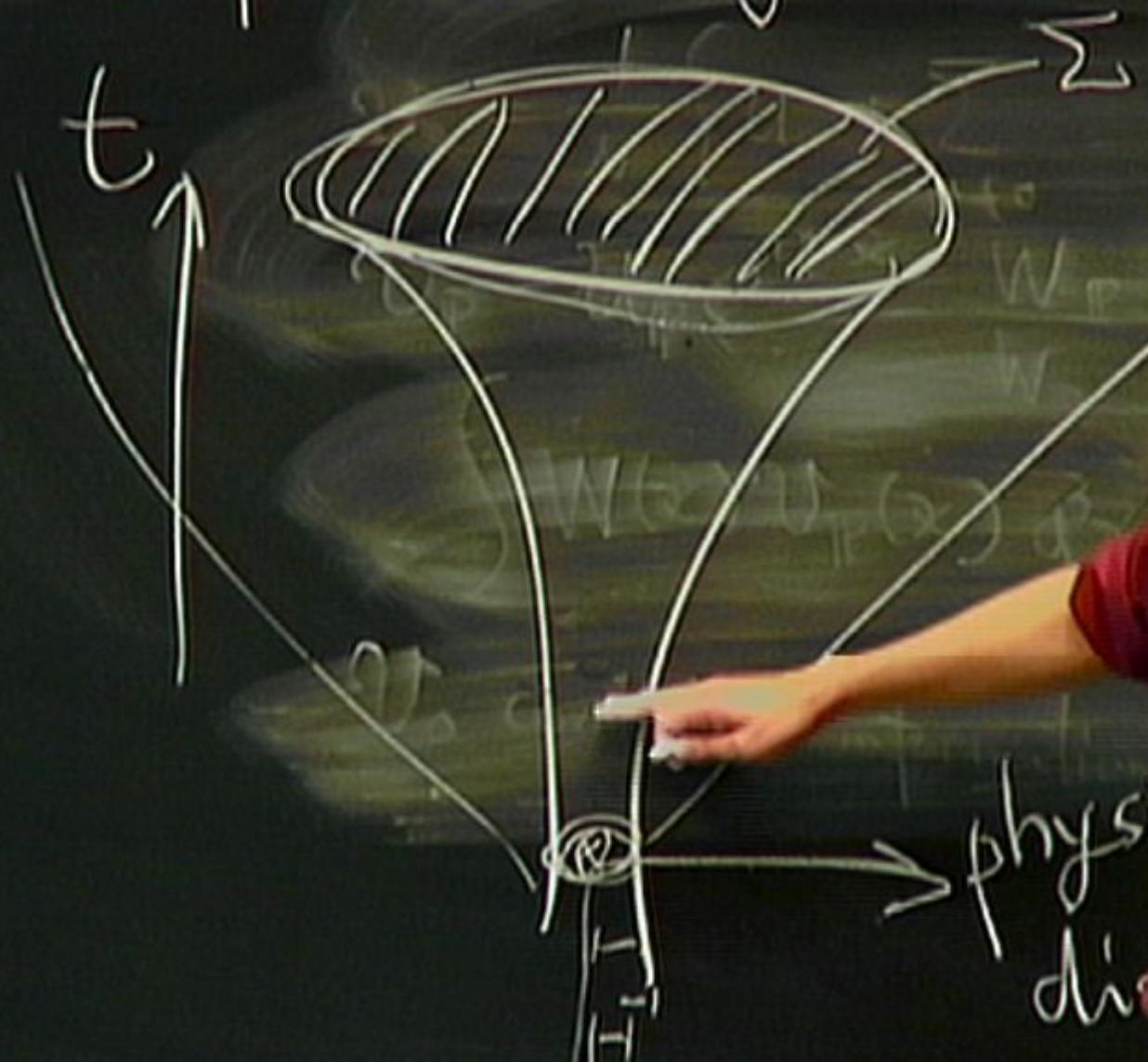
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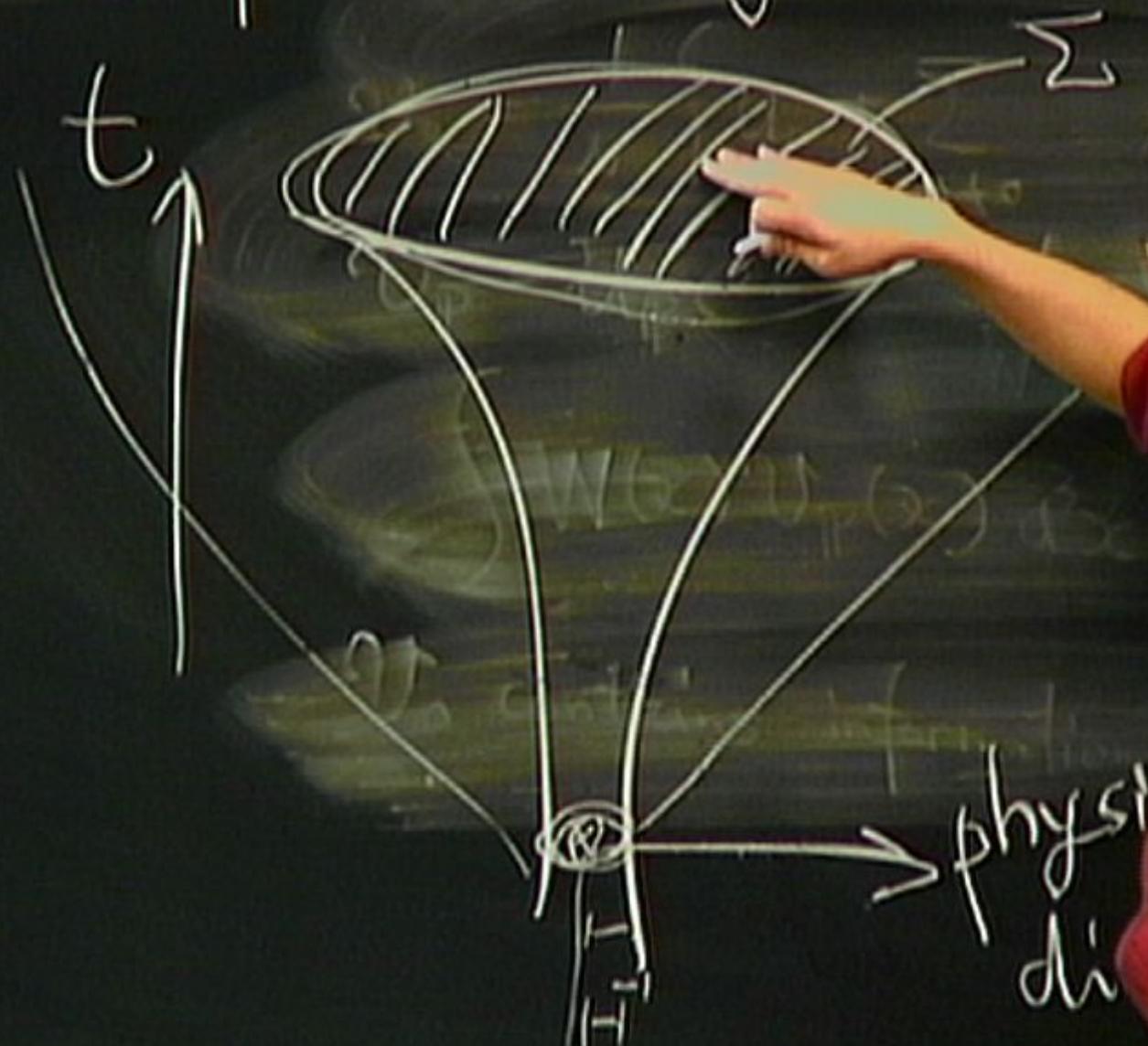
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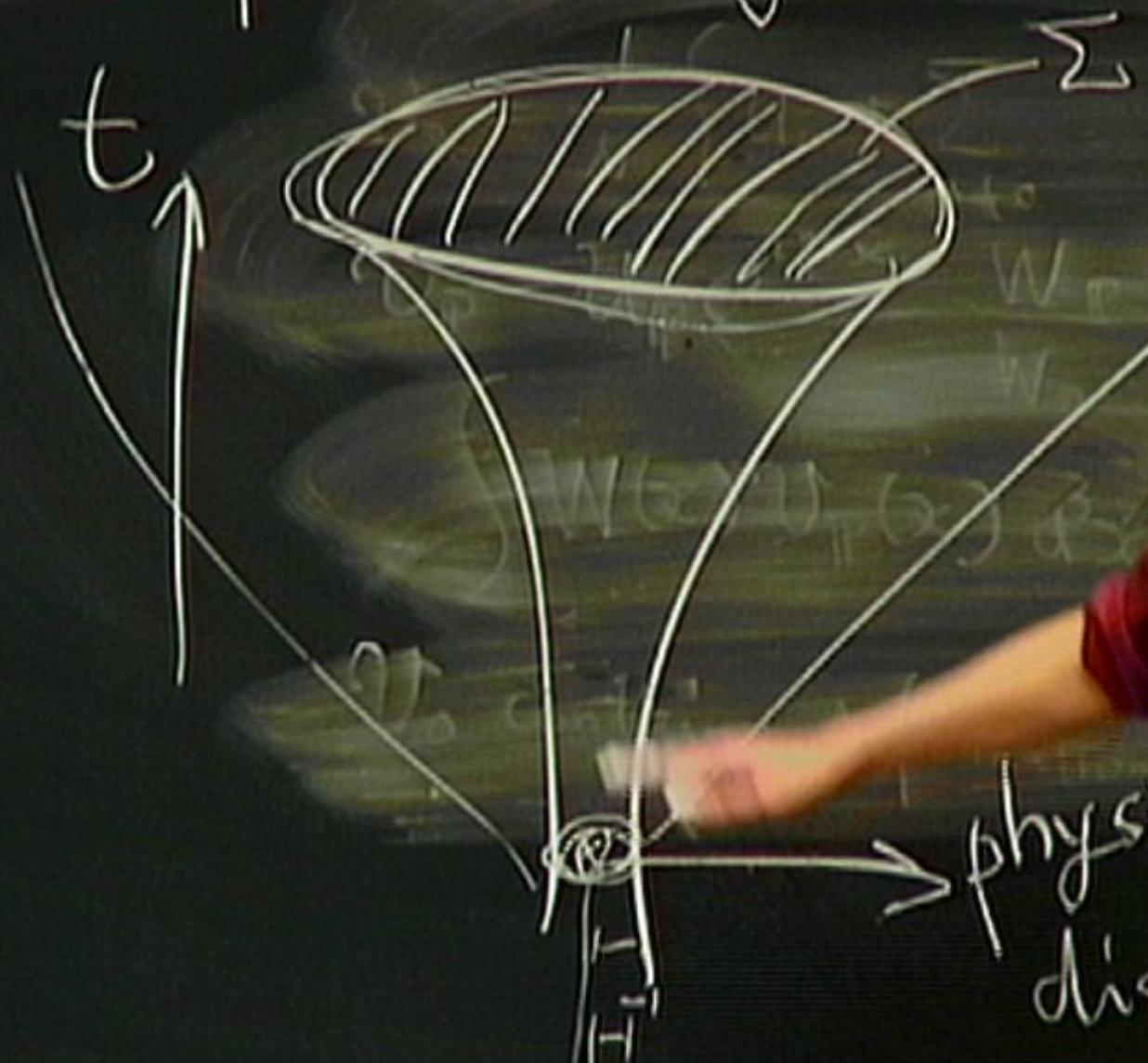
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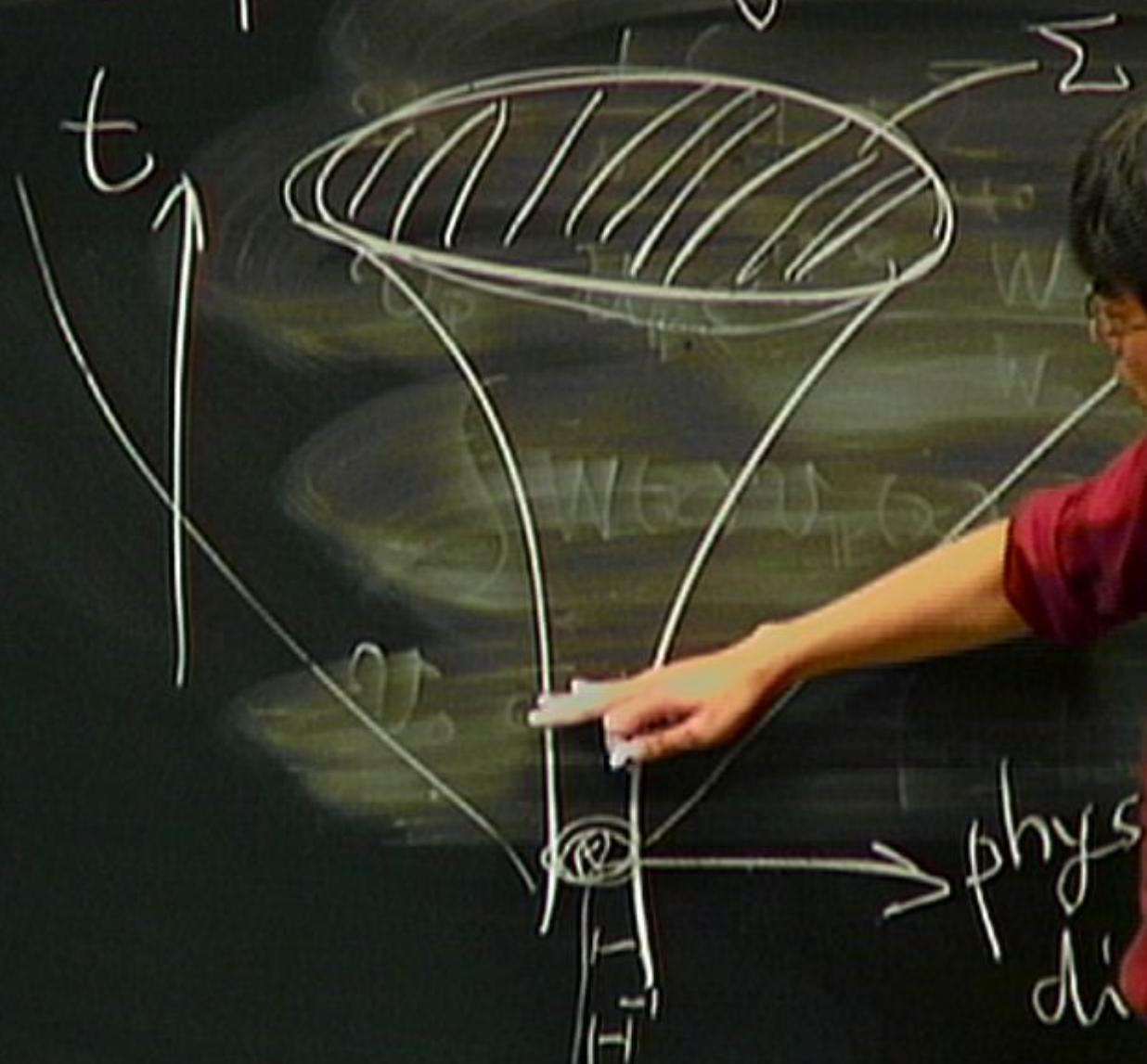
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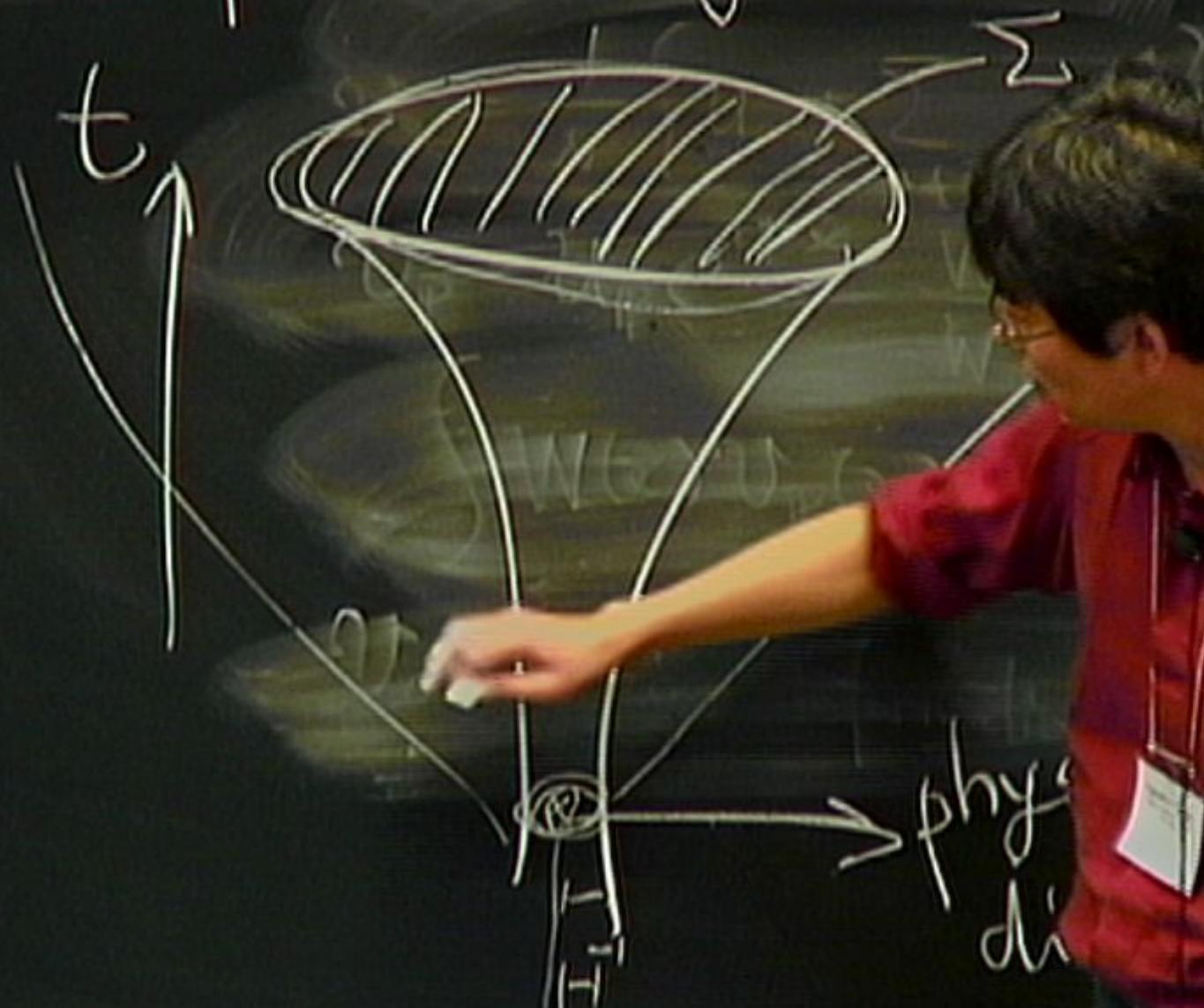
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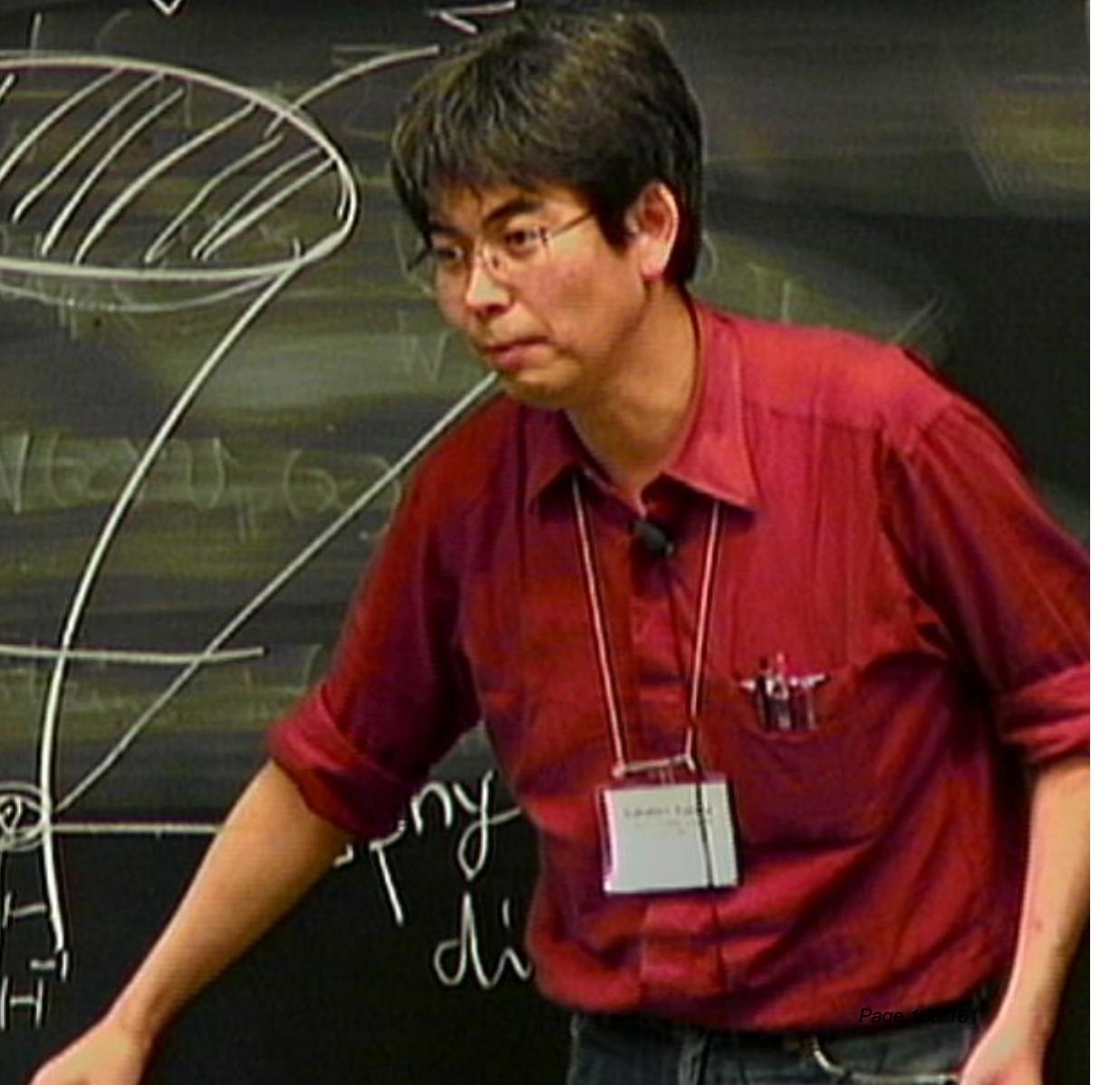
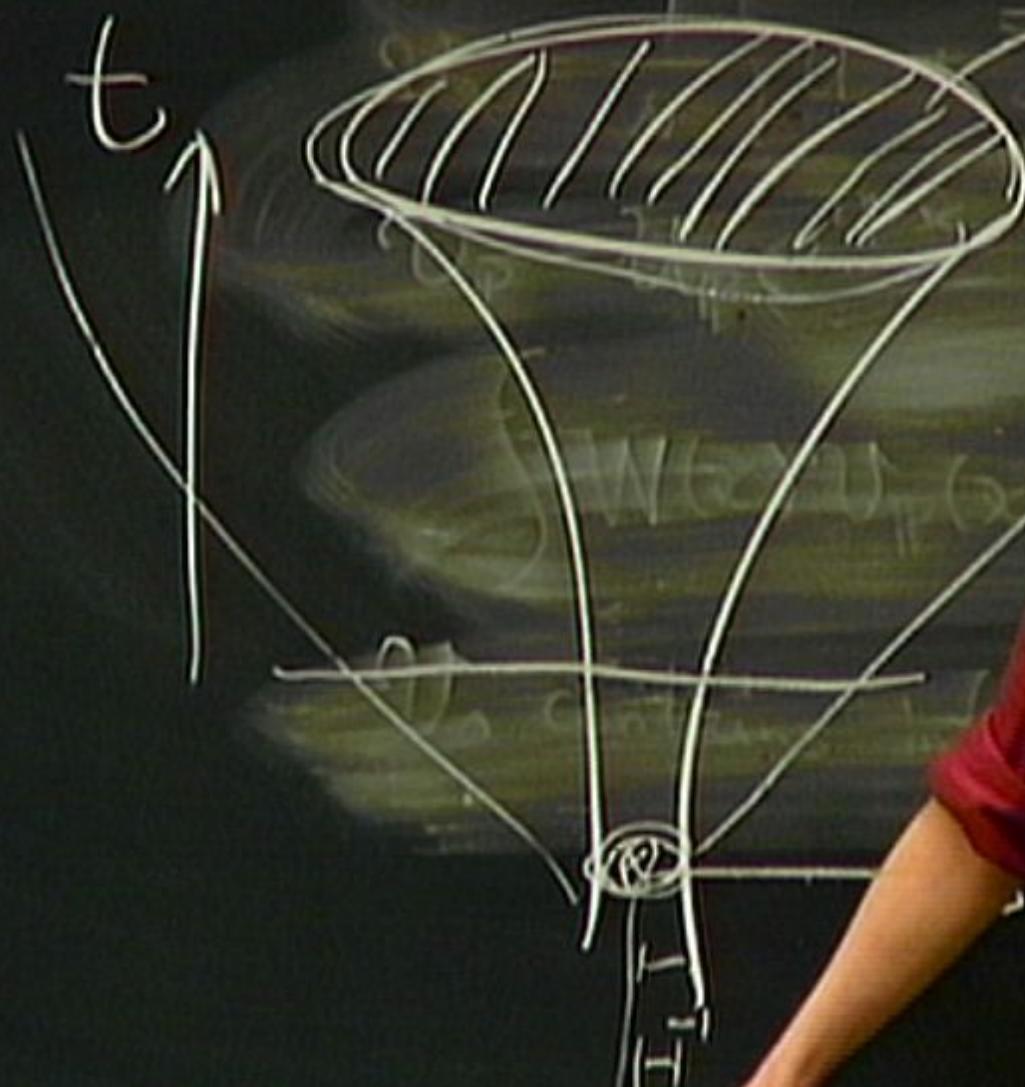
Temporal integral



Temporal integral



Temporal integral



Temporal integral

