

Title: Specker's parable of the overprotective seer: Implications for Contextuality, Nonlocality and Complementarity

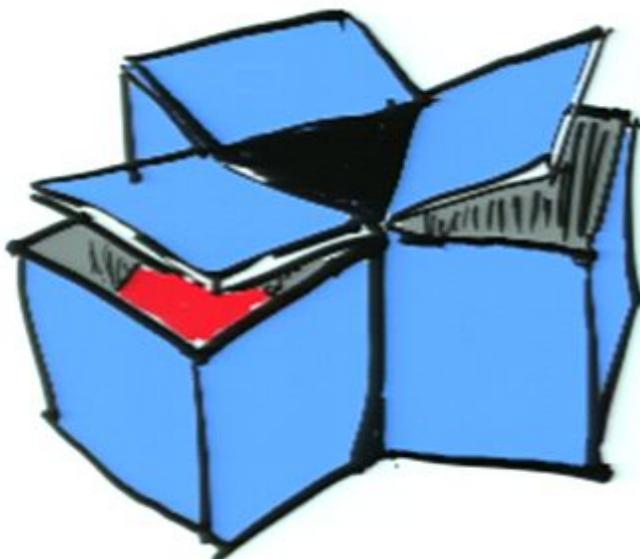
Date: Oct 12, 2010 04:00 PM

URL: <http://pirsa.org/10100060>

Abstract: I revisit an example of stronger-than-quantum correlations that was discovered by Ernst Specker in 1960. The example was introduced as a parable wherein an over-protective seer sets a simple prediction task to his daughter's suitors. The challenge cannot be met because the seer asks the suitors for a noncontextual assignment of values but measures a system for which the statistics are inconsistent with such an assignment. I will show how by generalizing these sorts of correlations, one is led naturally to some well-known proofs of nonlocality and contextuality, and to some new ones. Specker's parable involves a kind of complementarity that does not arise in quantum theory - three measurements that can be implemented jointly pairwise but not triplewise -- and therefore prompts the question of what sorts of foundational principles might rule out this kind of complementarity. This is joint work with Howard Wiseman and Yeong-Cherng Liang.

Specker's parable of the overprotective seer: Implications for Contextuality, Nonlocality and Complementarity

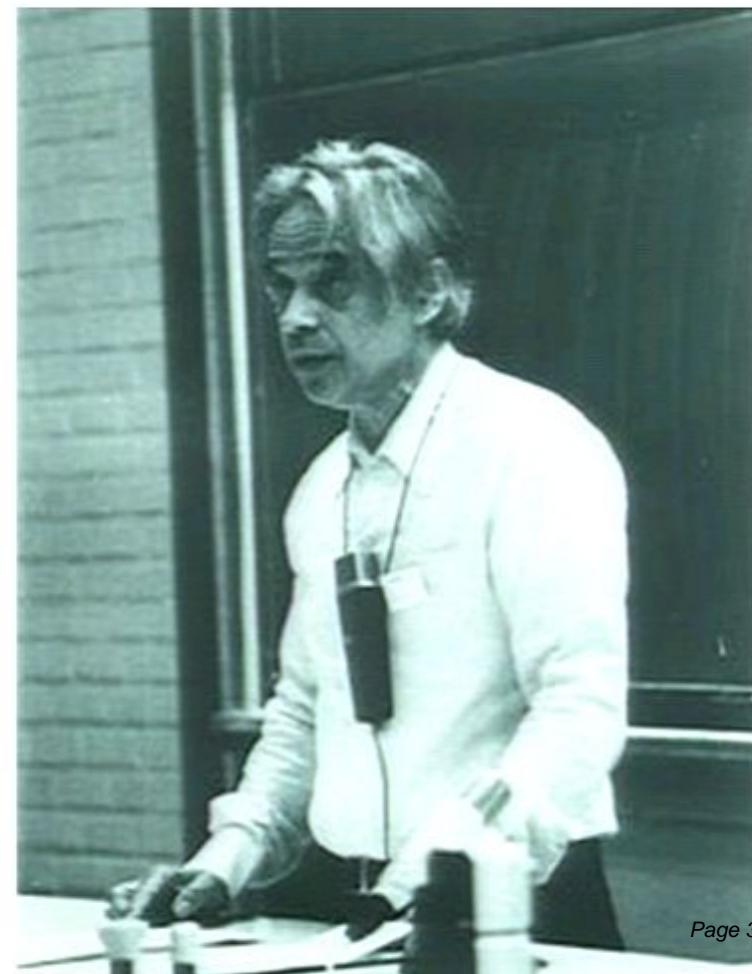
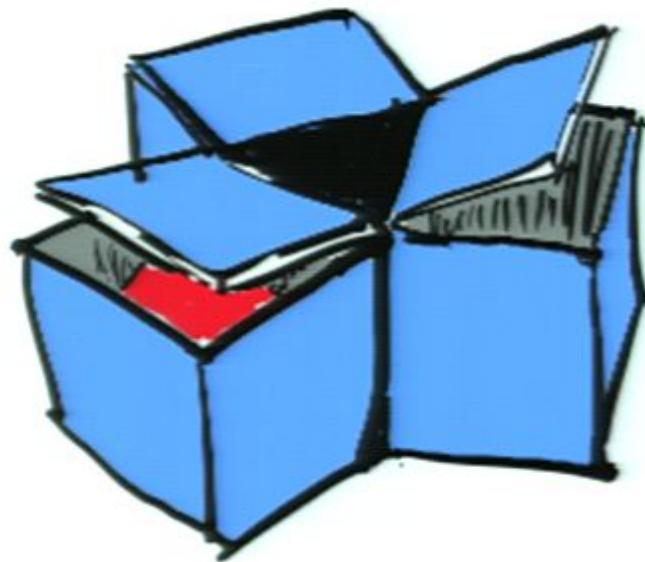
Robert Spekkens



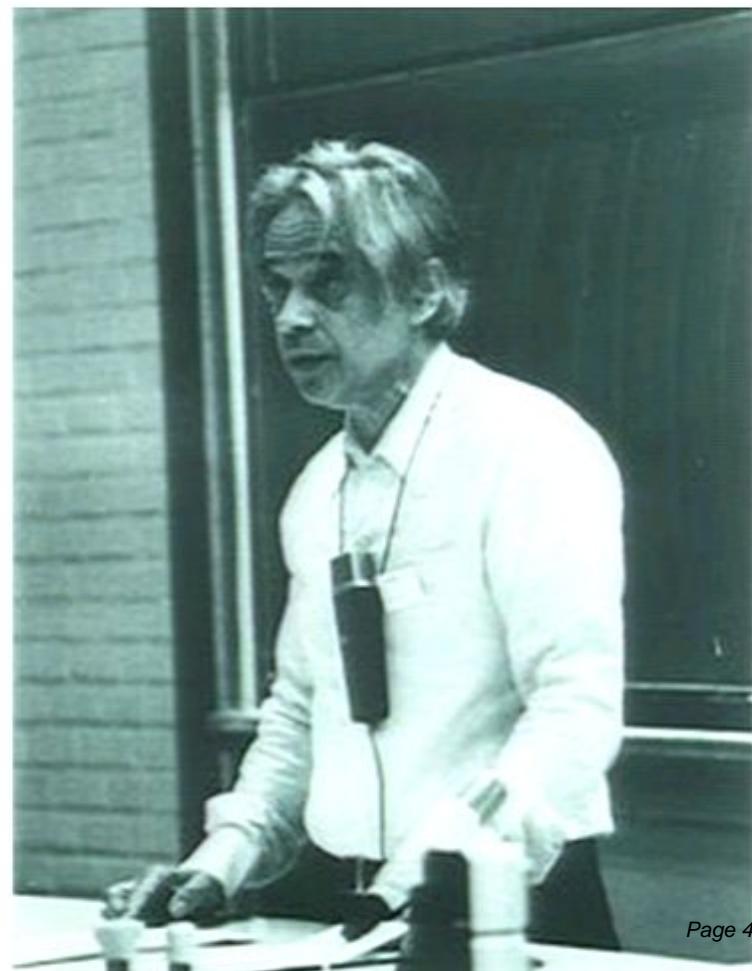
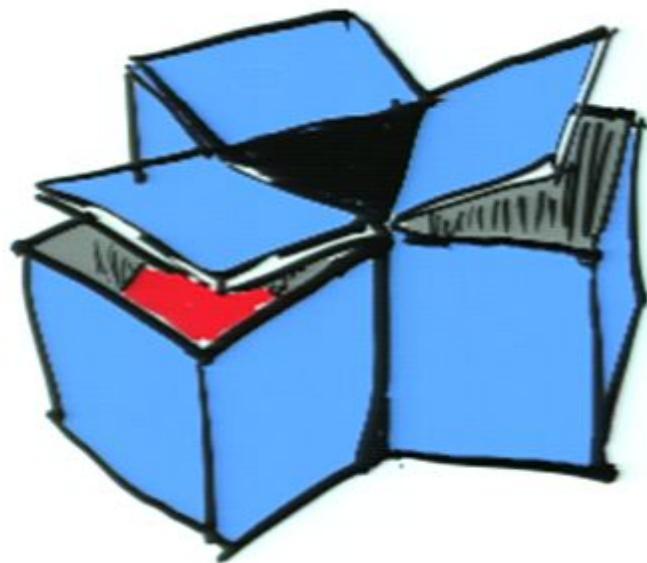
Joint work with:
Howard Wiseman
(Griffith University, Brisbane)
Yeong-Cherng Liang
(University of Sydney, now Geneva)

Funding by: the PIAF collaboration

Ernst Specker, "Die Logik nicht gleichzeitig entscheidbarer Aussagen",
Dialectica 14, 239 (1960)
("The logic of propositions which are not simultaneously decidable")



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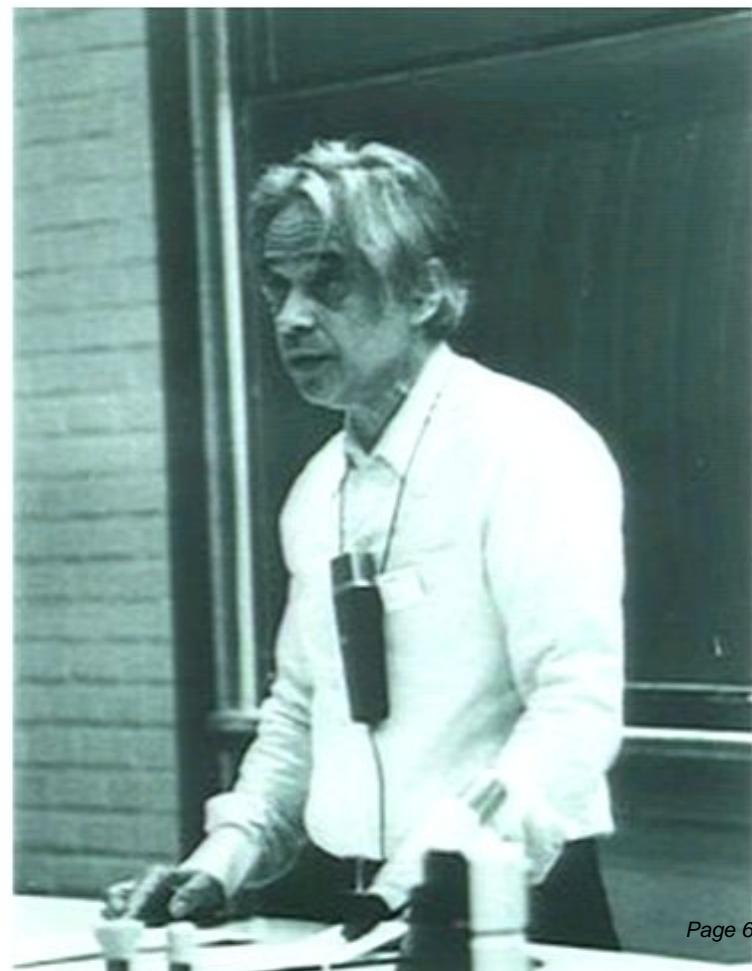
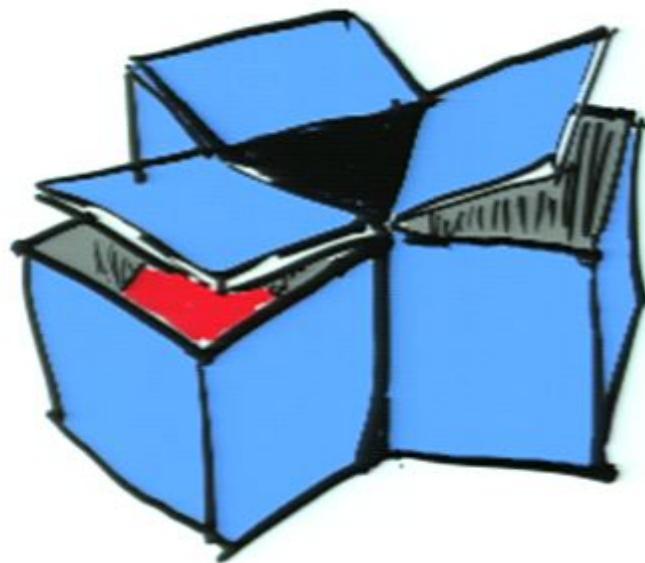
Joint measurability of POVMs

Noisy z-spin $\{E_+^z, E_-^z\}$ $E_{\pm}^z = \eta |\pm z\rangle\langle \pm z| + (1 - \eta) \frac{I}{2}$

Noisy x-spin $\{E_+^x, E_-^x\}$ η = sharpness factor

Noisy y-spin $\{E_+^y, E_-^y\}$

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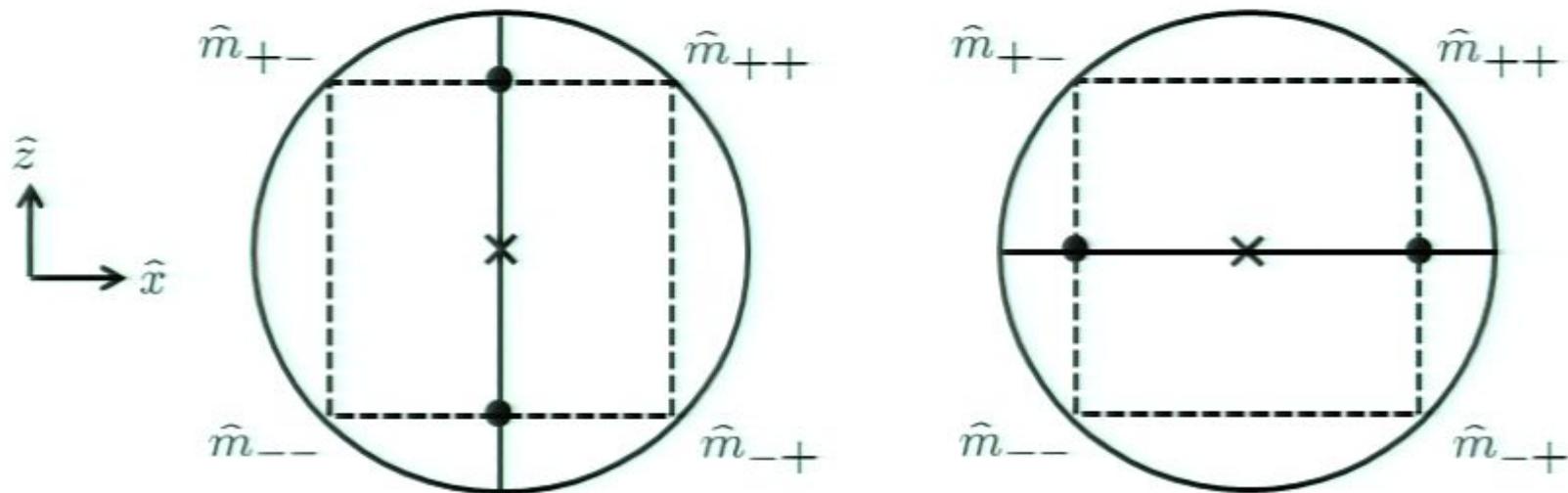
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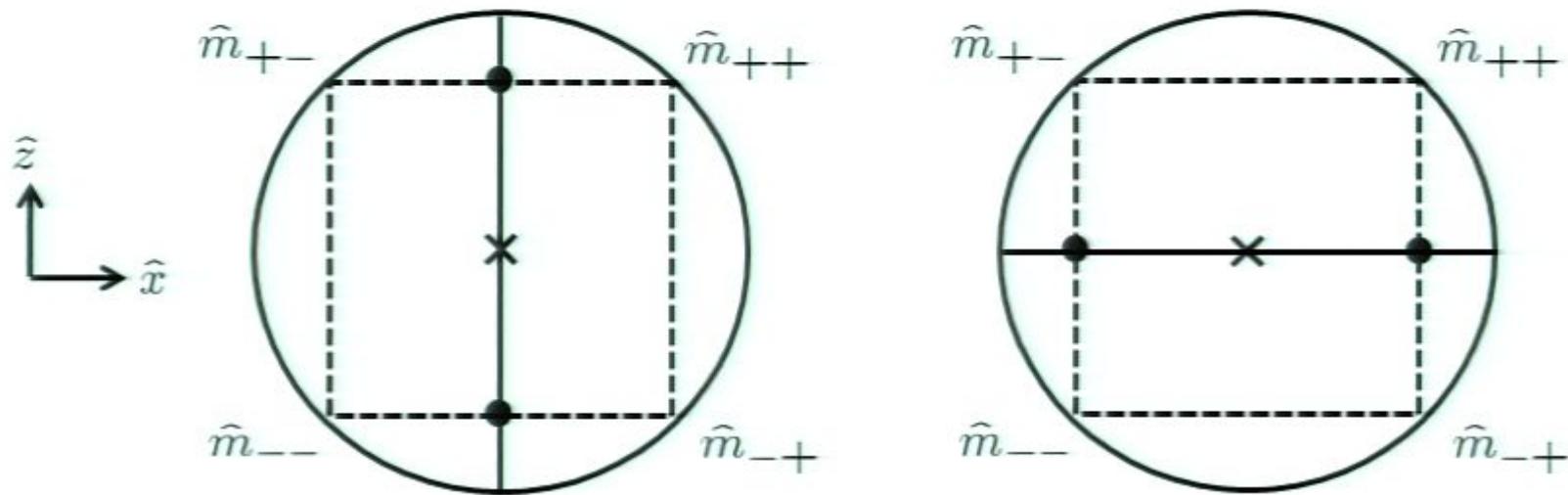


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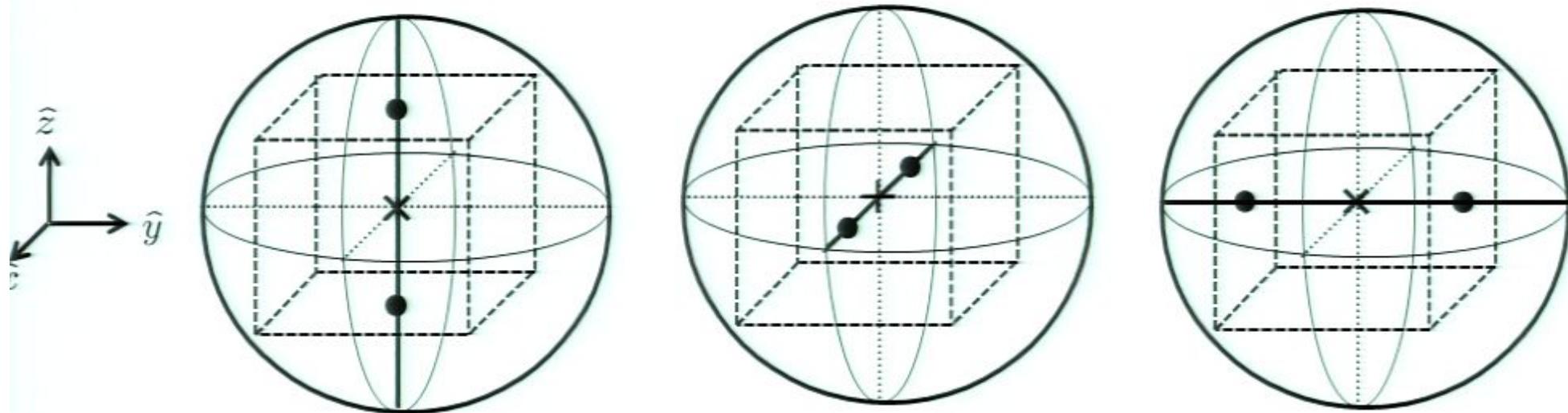
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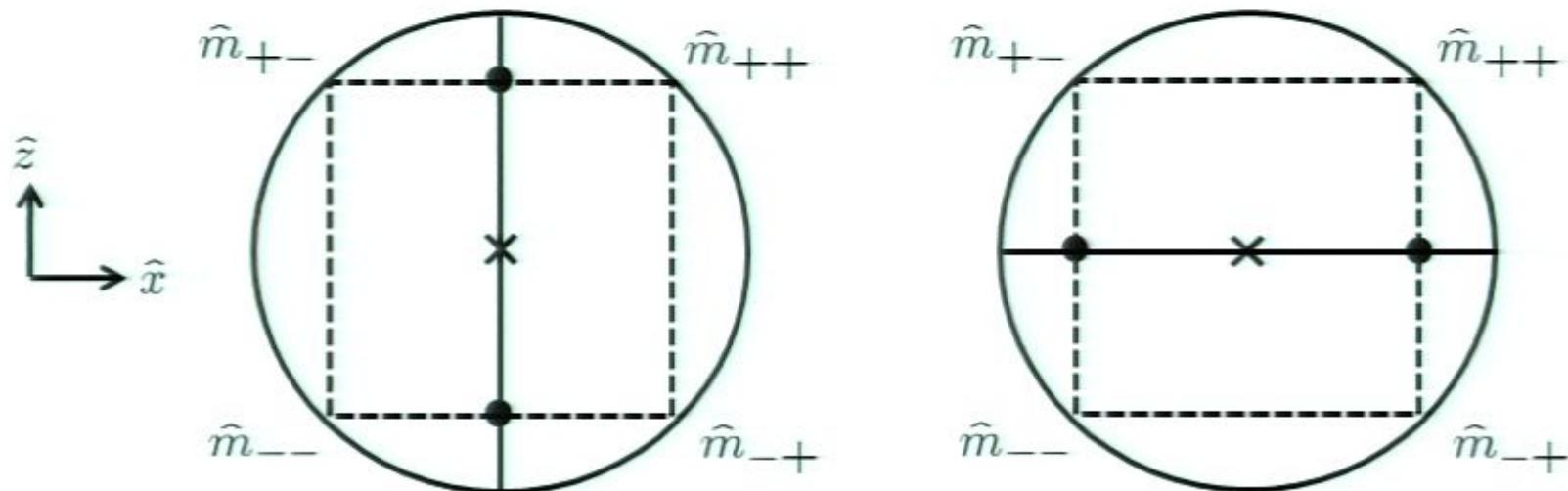
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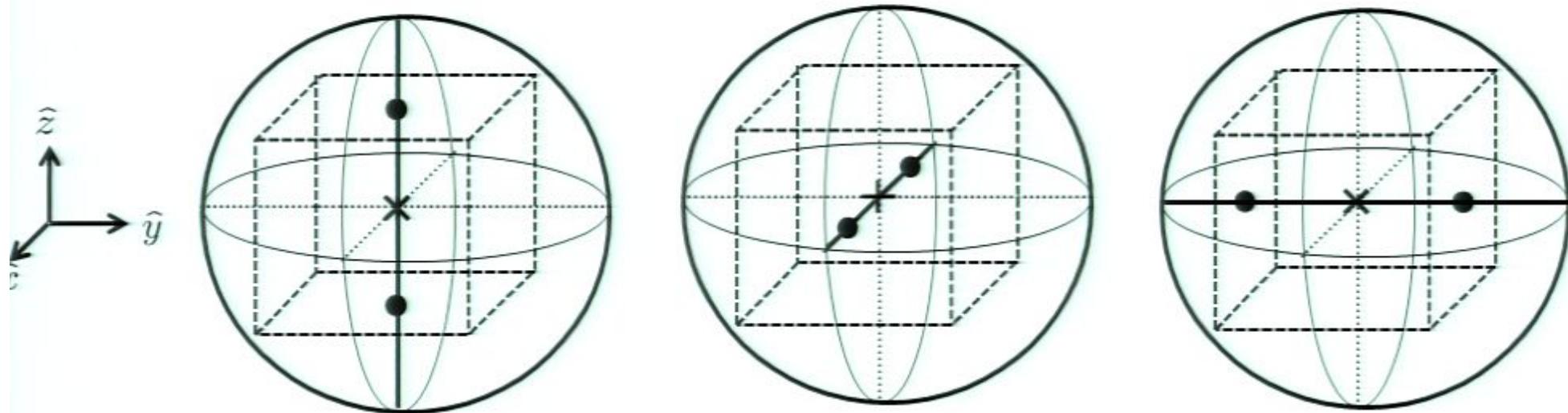
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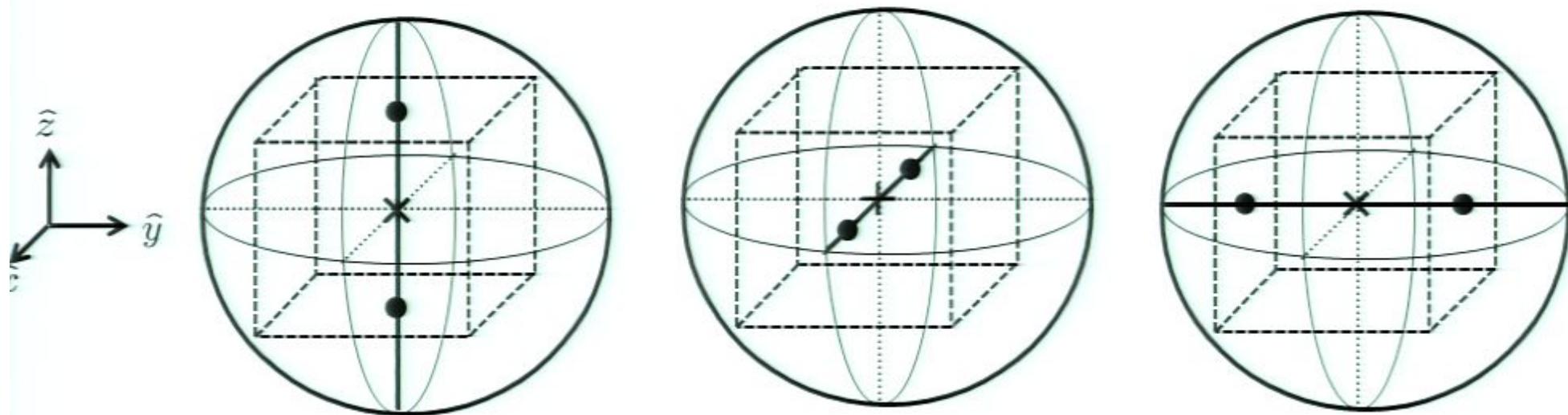
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Pairwise JM but not triplewise JM iff

$$\frac{1}{\sqrt{2}} \leq \eta \leq \frac{1}{\sqrt{3}}$$

Using a notion of noncontextuality for POVMs in
RWS, PRA 71, 052108 (2005)

we find that these correlations are
consistent with a noncontextual model

Therefore, we move on to projective measurements

Frustrated Networks

Nodes are binary variables

Edges imply joint measurability

- —○ Perfect correlation
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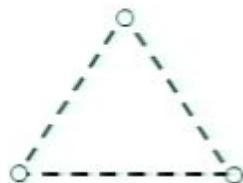
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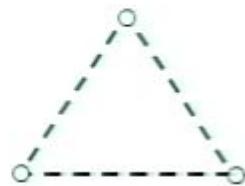
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Frustration = no valuation satisfying all correlations

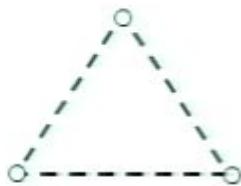
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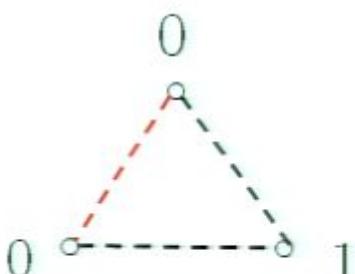
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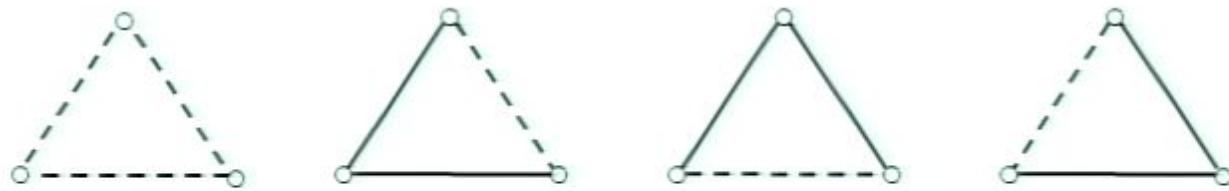


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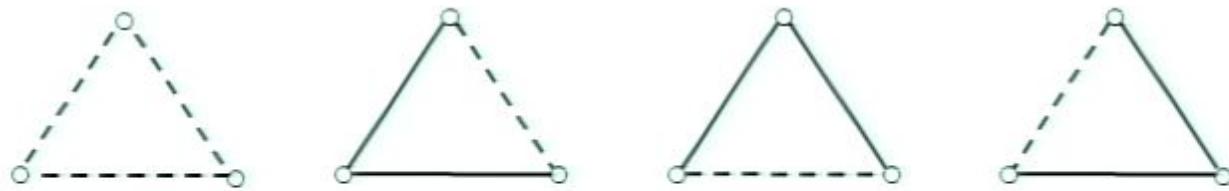


$$R \leq \frac{2}{3}$$

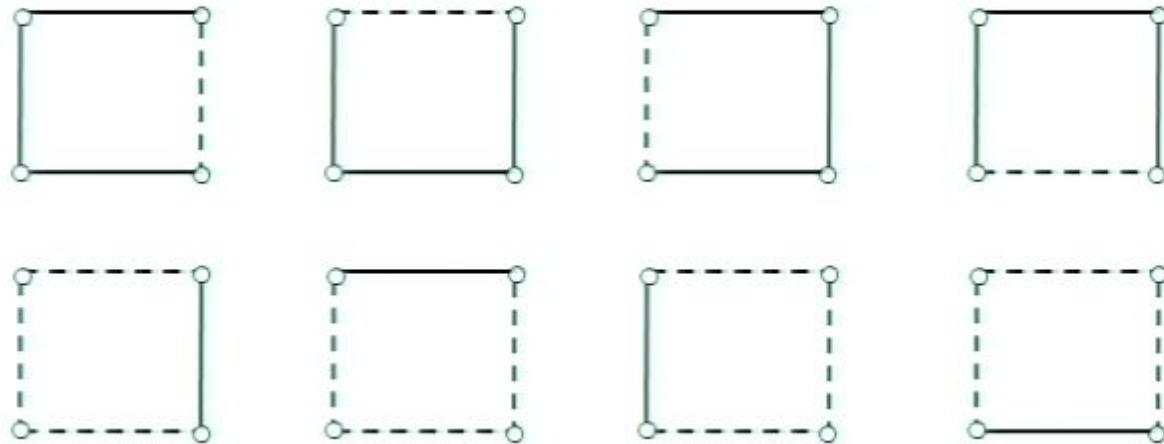
Kochen-Specker inequality



All 3-node frustrated networks
Equivalent under relabelling of the outcomes

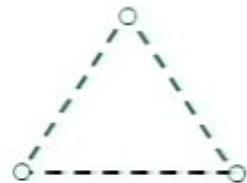


All 3-node frustrated networks
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All 4-node frustrated networks
Equivalent under relabelling of the outcomes

Kochen-Specker
bound



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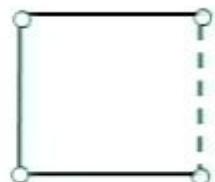
Quantum
bound

N/A

Maximum
value

$$R = 1$$

Overprotective
Seer (OS)
correlations

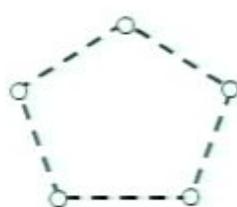


$$R \leq \frac{3}{4}$$

$$R \leq \frac{1}{2} + \frac{1}{\sqrt{2}} \\ \simeq 0.85$$

$$R = 1$$

CHSH
correlations
& PR box



$$R \leq \frac{4}{5}$$

$$R \leq \frac{2}{\sqrt{5}} \\ \simeq 0.89$$

$$R = 1$$

Klyachko
correlations



$$R \leq \frac{5}{6}$$

$$R \leq \frac{1}{2} + \frac{\sqrt{3}}{4} \\ \simeq 0.93$$

$$R = 1$$

Vaidman's
necklace
correlations



$$R \leq \frac{n-1}{n}$$

$$= 1 - \frac{1}{n}$$

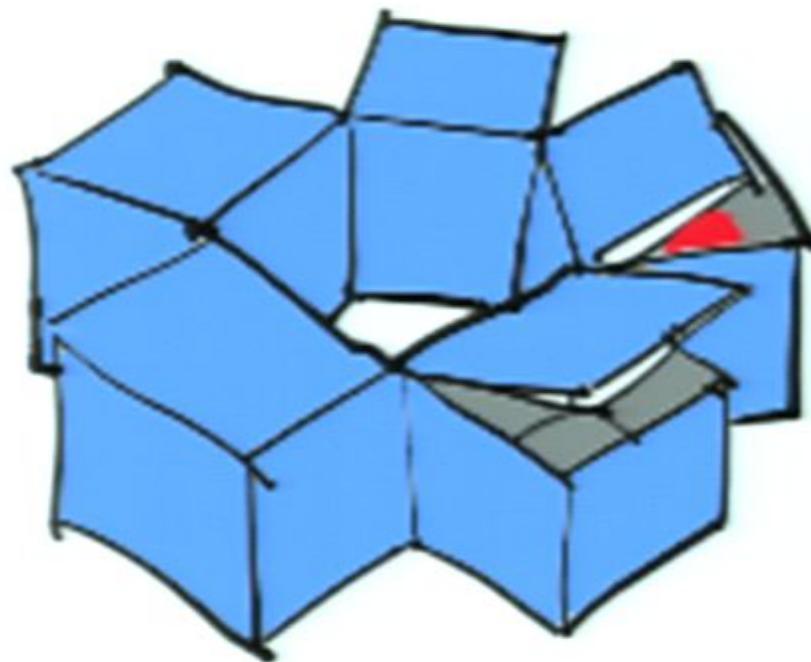
$$R \leq \frac{2 \cos(\frac{\pi}{n})}{1 + \cos(\frac{\pi}{n})}$$

$$= 1 - \frac{\pi^2}{n^2}$$

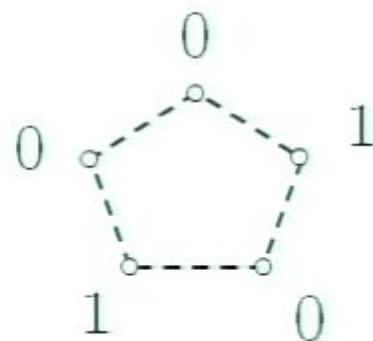
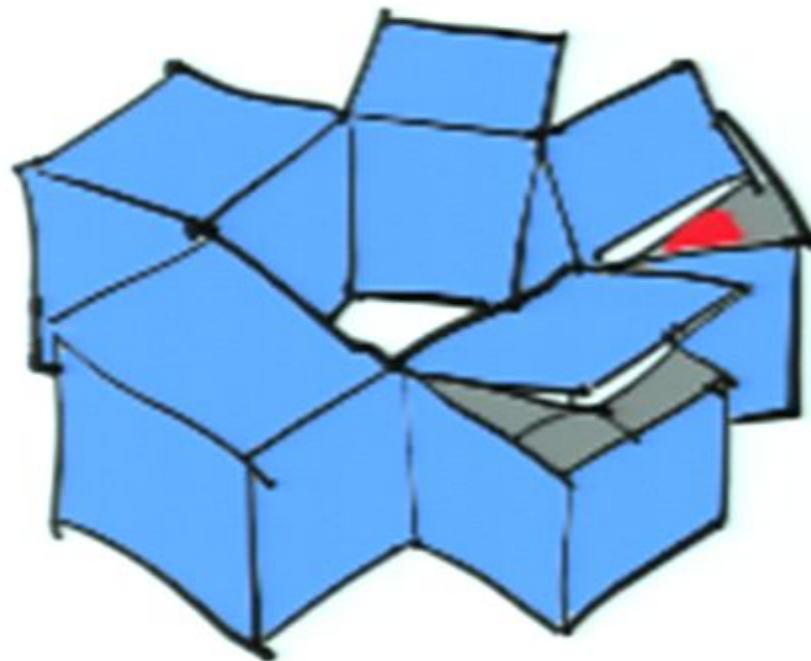
$$R = 1$$

Analogue of
Chained Bell
correlations

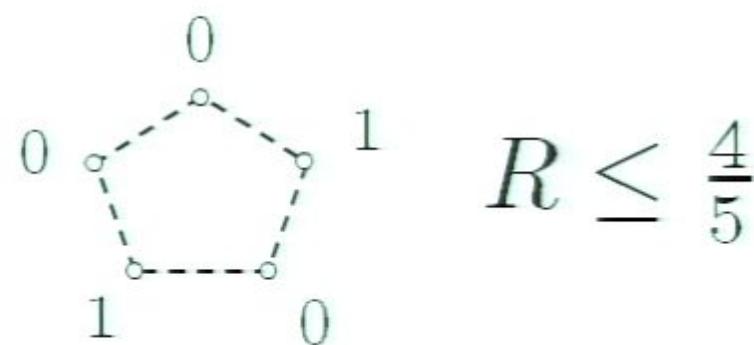
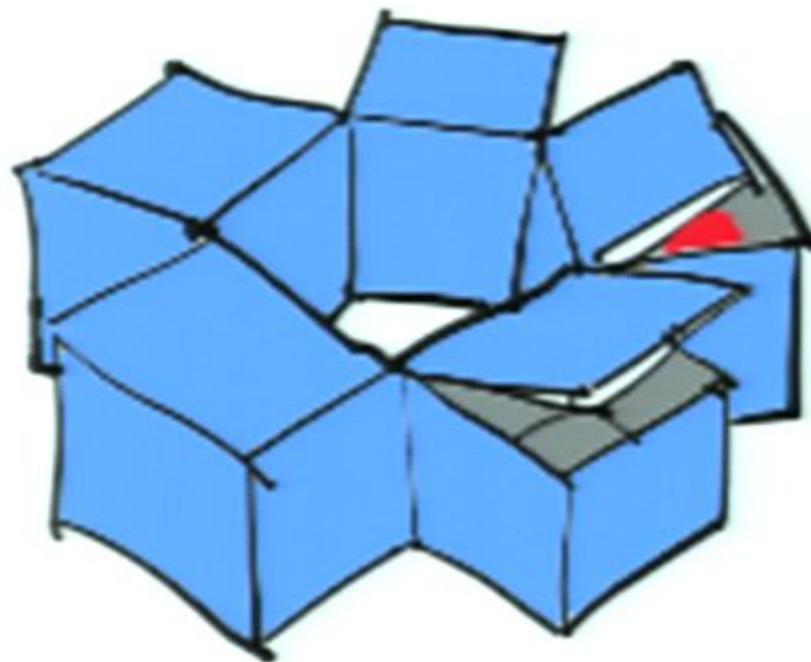
Double-query 5-box system allowing only **adjacent** queries



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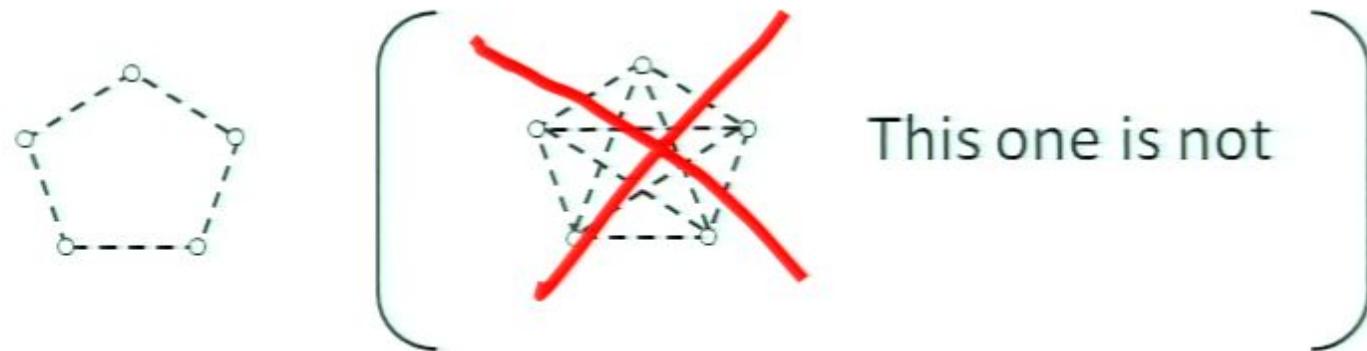
Double-query 5-box system allowing only **adjacent** queries



This pattern of joint measurability is **possible** in Quantum Theory



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E.g. 5 projective mmts:

$$\{\Pi_1, I - \Pi_1\}$$

$$\{\Pi_2, I - \Pi_2\}$$

$$\{\Pi_3, I - \Pi_3\}$$

$$\{\Pi_4, I - \Pi_4\}$$

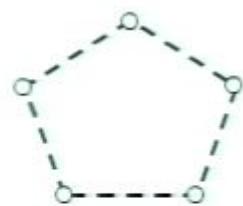
$$\{\Pi_5, I - \Pi_5\}$$

where $[\Pi_i, \Pi_{i+1}] = 0 \quad i \in \{1, \dots, 5\}$

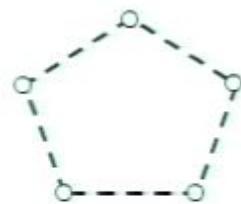
Klyachko's proof of the Kochen-Specker theorem

Klyachko, arXiv:quant-ph/0206012

Klyachko's proof of the Kochen-Specker theorem



Klyachko's proof of the Kochen-Specker theorem



5 projective mmts:

$$\{|l_1\rangle\langle l_1|, I - |l_1\rangle\langle l_1|\}$$

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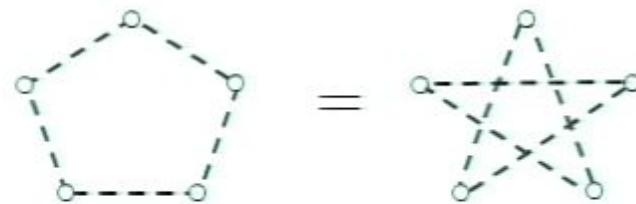
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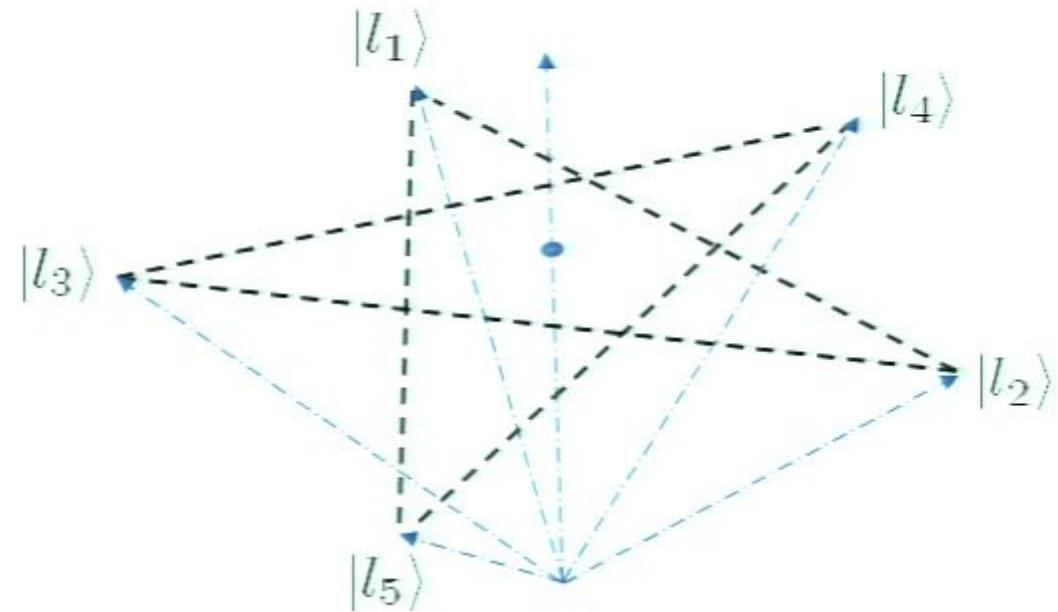
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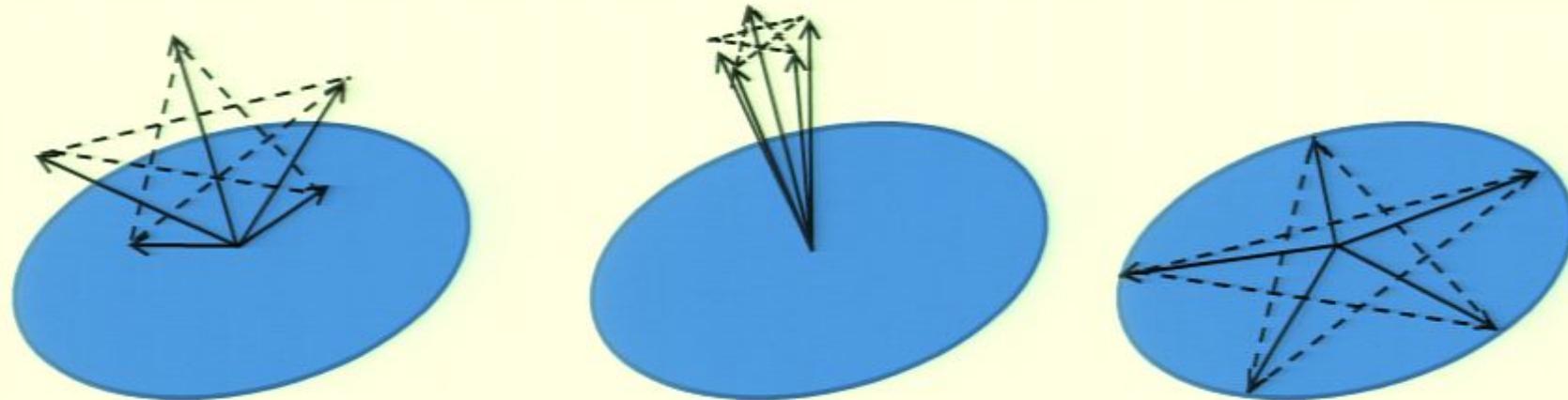
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where $\langle l_i | l_{i \oplus 1} \rangle = 0$ $i \in \{1, \dots, 5\}$



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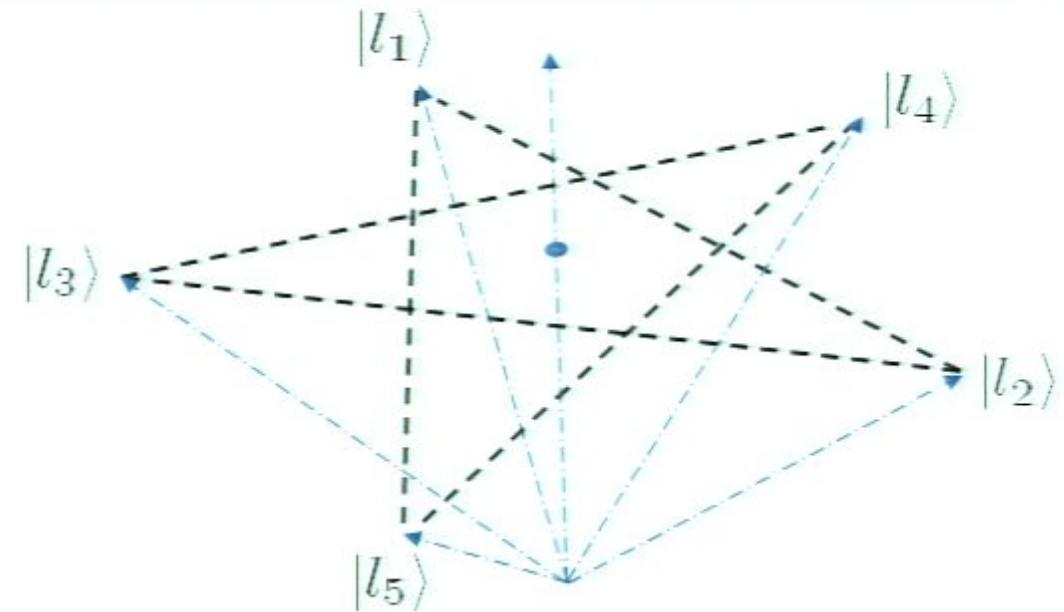
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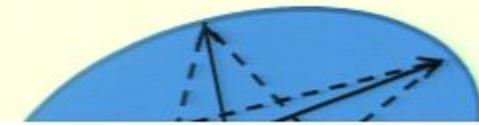
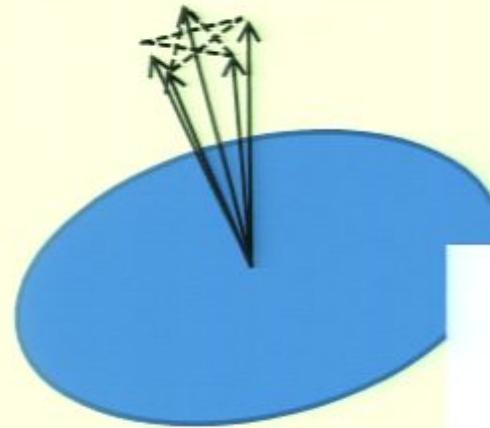
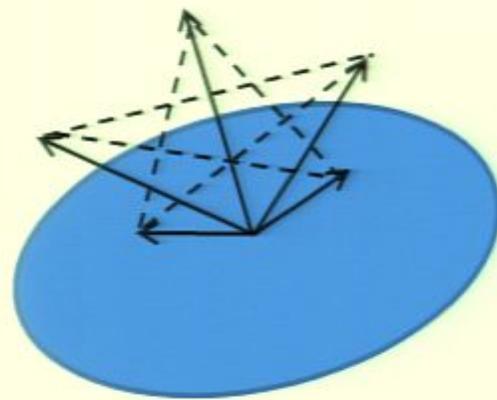
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$$\cos^2 \theta = \frac{1}{\sqrt{5}}$$

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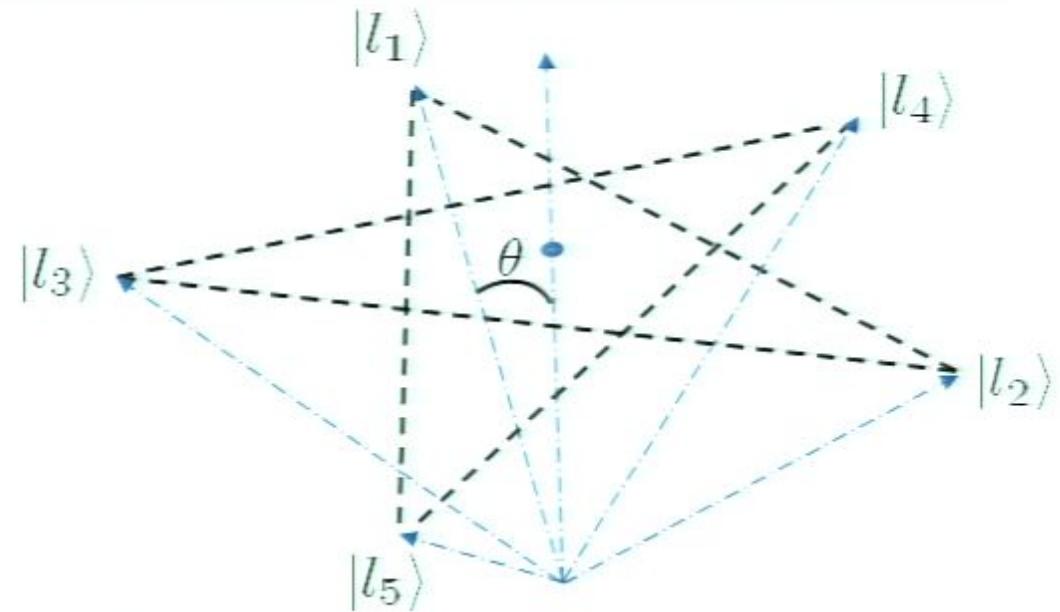
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Klyachko's proof of the Kochen-Specker theorem

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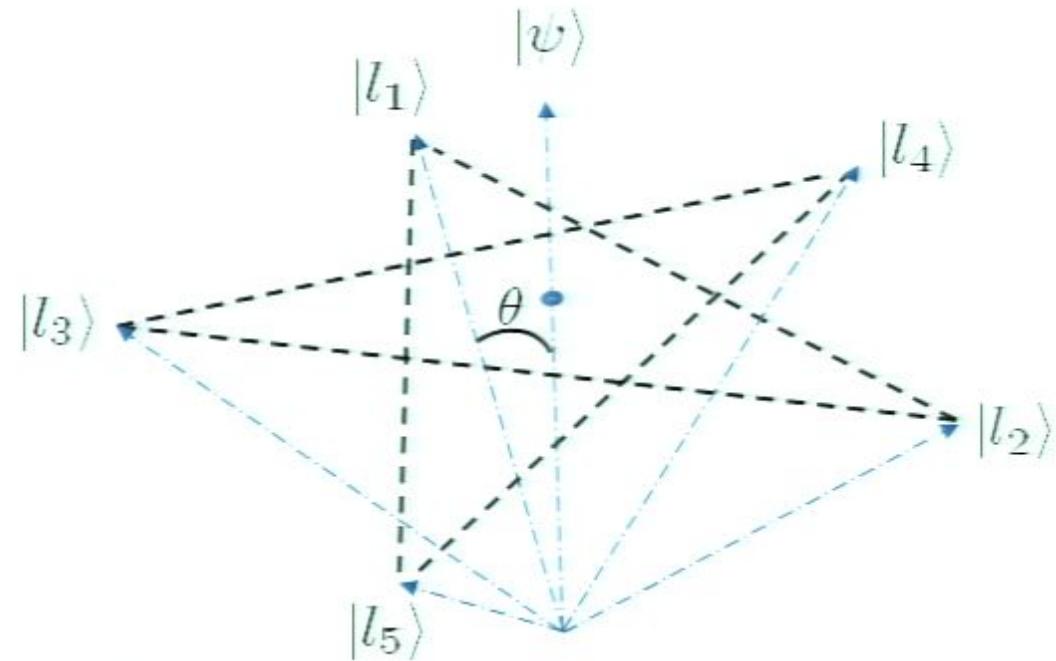
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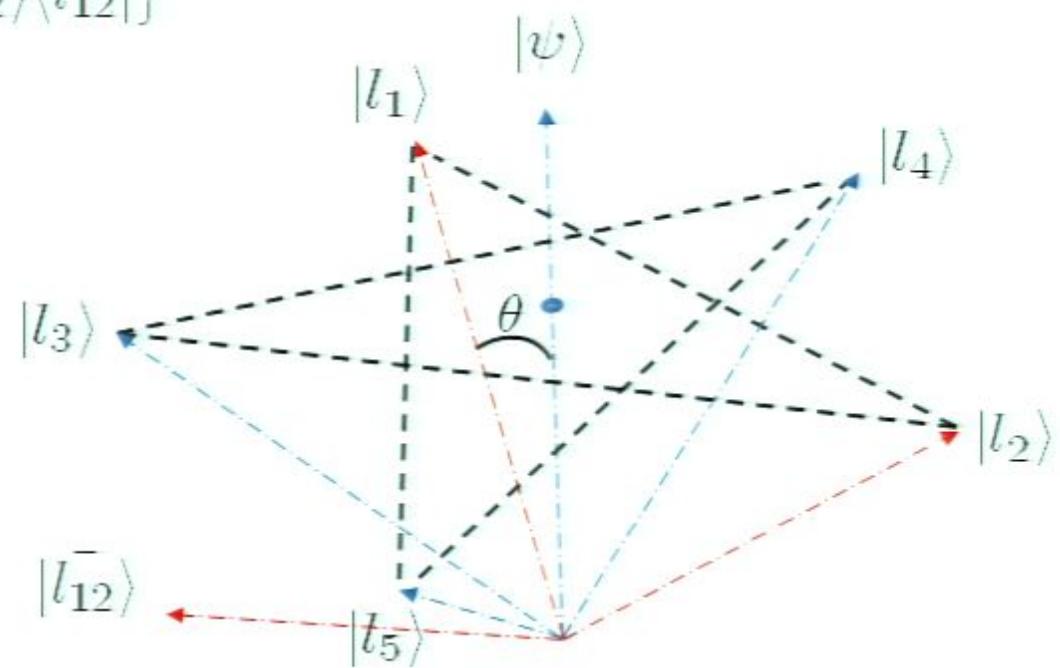


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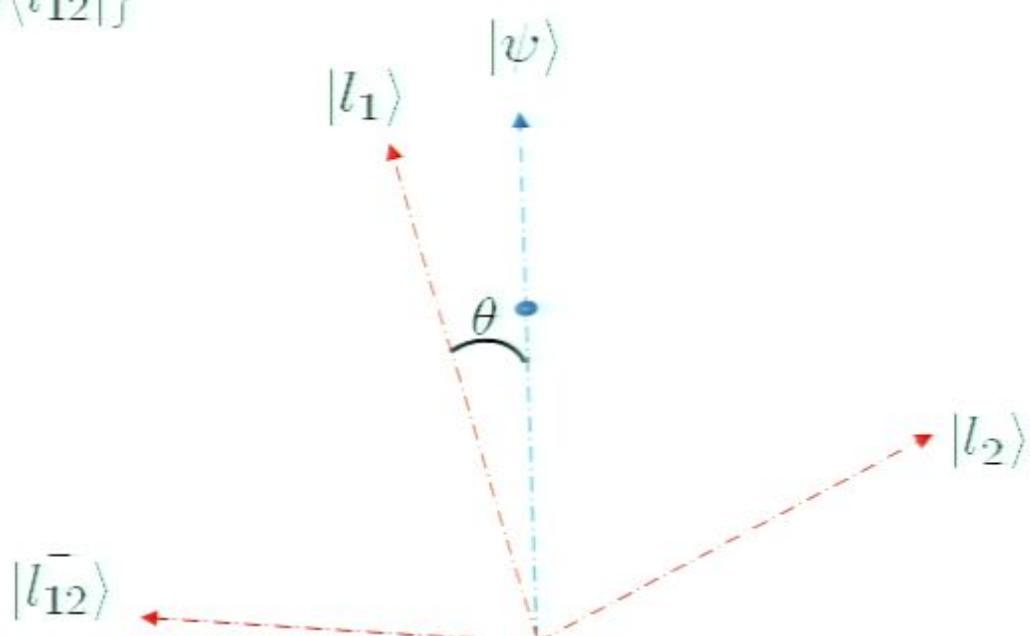
Consider: $\{|l_1\rangle\langle l_1|, |l_2\rangle\langle l_2|, |\bar{l_{12}}\rangle\langle \bar{l_{12}}|\}$



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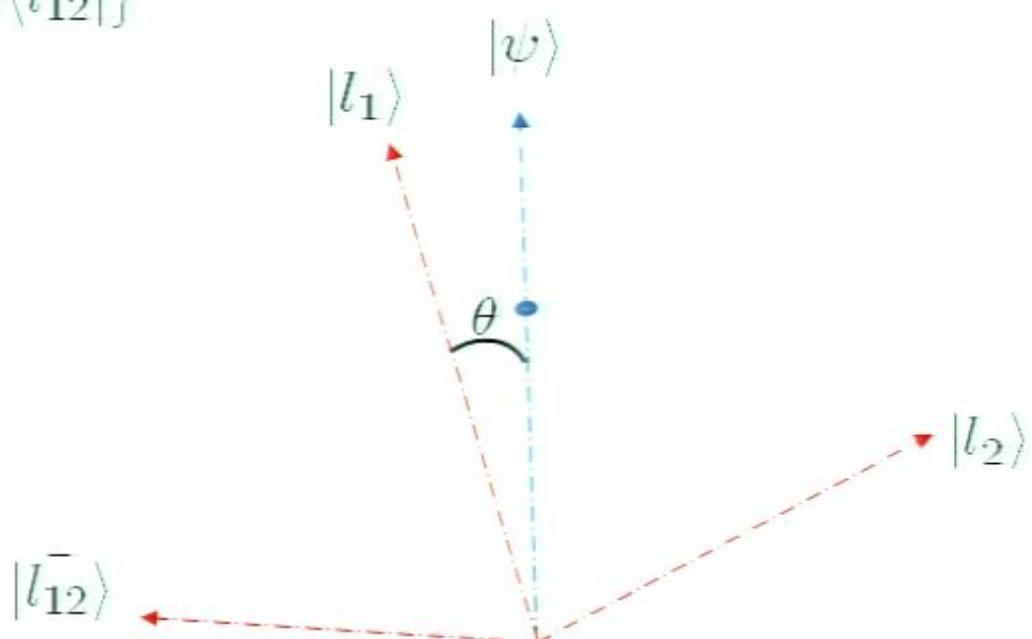
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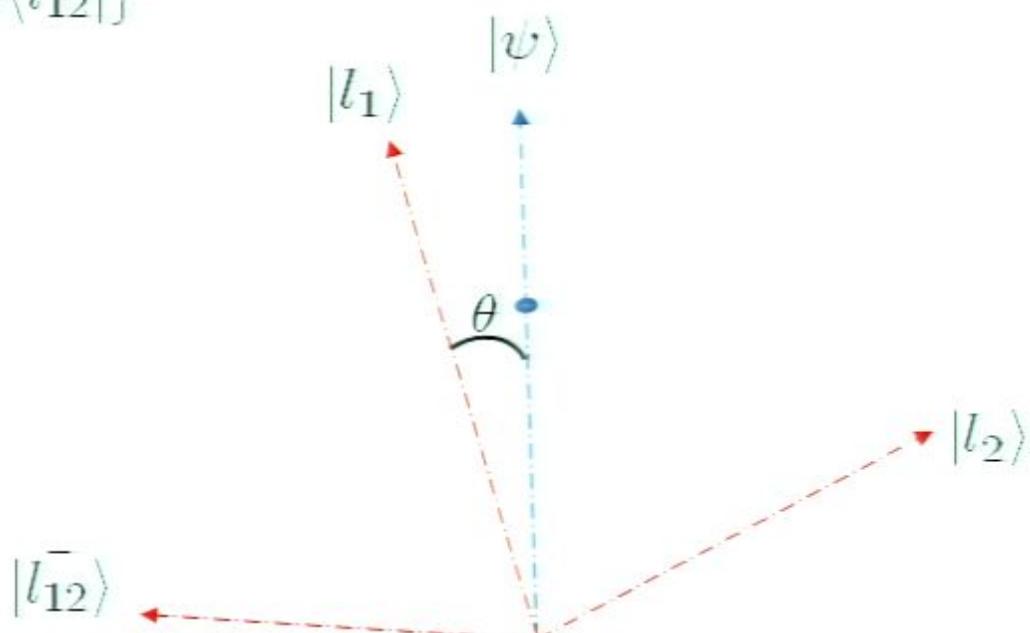
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Klyachko's proof of the Kochen-Specker theorem

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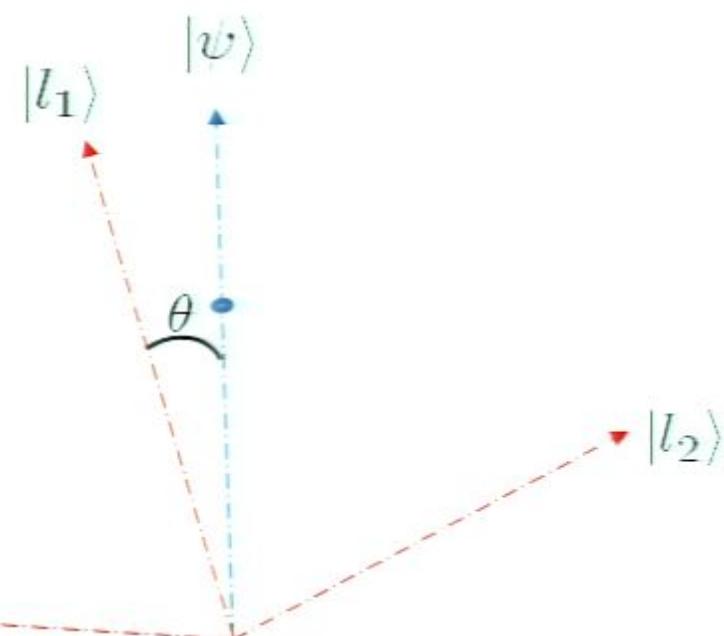
Consider: $\{|l_1\rangle\langle l_1|, |l_2\rangle\langle l_2|, |\bar{l_{12}}\rangle\langle \bar{l_{12}}|\}$

$$\left. \begin{array}{l} v(|l_1\rangle\langle l_1|) = 1 \\ v(|l_2\rangle\langle l_2|) = 0 \\ v(|\bar{l_{12}}\rangle\langle \bar{l_{12}}|) = 0 \\ \\ v(|l_1\rangle\langle l_1|) = 0 \\ v(|l_2\rangle\langle l_2|) = 1 \\ v(|\bar{l_{12}}\rangle\langle \bar{l_{12}}|) = 0 \\ \\ v(|l_1\rangle\langle l_1|) = 0 \\ v(|l_2\rangle\langle l_2|) = 0 \\ v(|\bar{l_{12}}\rangle\langle \bar{l_{12}}|) = 1 \end{array} \right\}$$

**Probability of
anticorrelated outcomes**

$$|\langle \psi | l_1 \rangle|^2 + |\langle \psi | l_2 \rangle|^2 = \frac{2}{\sqrt{5}}$$

$$|\bar{l_{12}}\rangle$$



$$\text{prob.} |\langle \psi | \bar{l_{12}} \rangle|^2 = 1 - \frac{2}{\sqrt{5}}$$

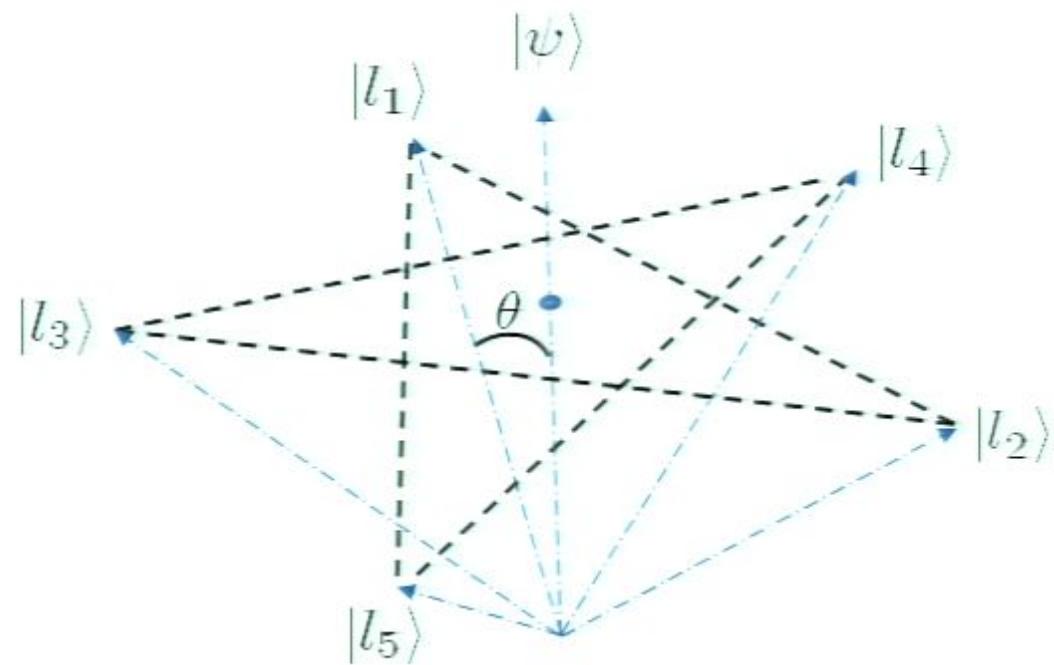
Klyachko's proof of the Kochen-Specker theorem

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Similarly for any pair of measurements...

Probability of anticorrelated outcomes

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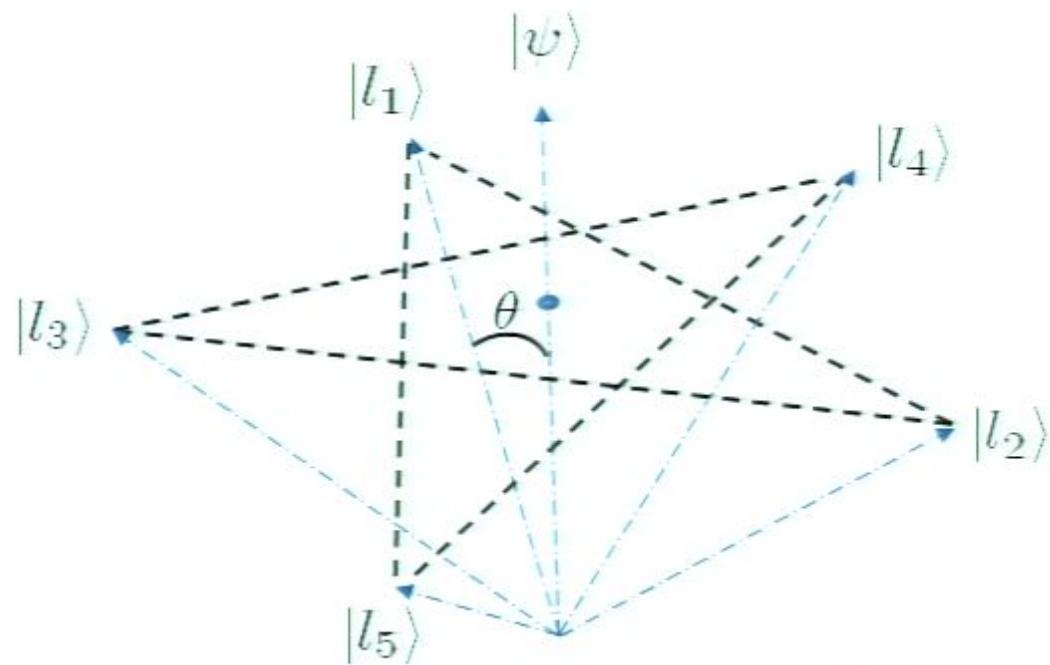
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Quantum violation of Kochen-Specker inequality

$$R = \frac{2}{\sqrt{5}} \simeq 0.89 > \frac{1}{\sqrt{2}}$$

Double-query n-box system allowing only adjacent queries



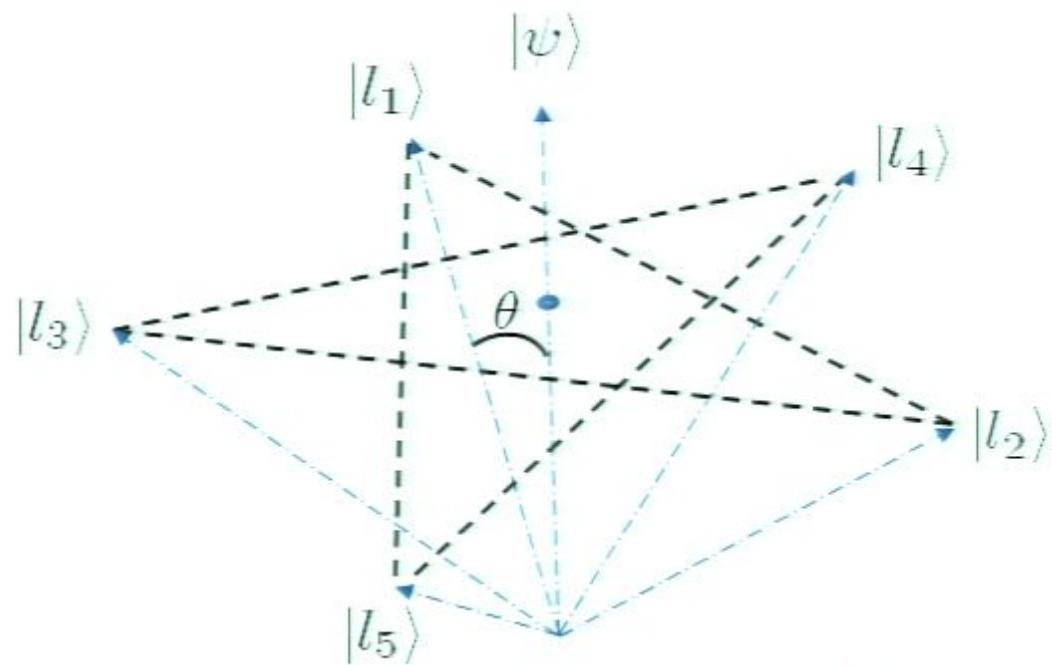
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Double-query n -box system allowing only adjacent queries



Double-query n-box system allowing only adjacent queries



$$R \leq \frac{n-1}{n} = 1 - \frac{1}{n}$$

Kochen-Specker
inequality

Double-query n-box system allowing only adjacent queries



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Generalization of Klyachko's proof

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Kochen-Specker
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Generalization of Klyachko's proof

n projective mmts:

$$\{|l_1\rangle\langle l_1|, I - |l_1\rangle\langle l_1|\}$$

$$\{|l_2\rangle\langle l_2|, I - |l_2\rangle\langle l_2|\}$$

...

$$\{|l_n\rangle\langle l_n|, I - |l_n\rangle\langle l_n|\}$$

where $\langle l_i | l_{i+1} \rangle = 0 \quad i \in \{1, \dots, n\}$

Double-query n-box system allowing only adjacent queries



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$\{n/\frac{n-1}{2}\}$ Star polygons

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Quantum violation of Kochen-Specker inequality

$$R = \frac{2\cos(\frac{\pi}{n})}{1 + \frac{2\cos(\frac{\pi}{n})}{n}} > 1 - \frac{1}{n}$$

How the seer can achieve his ends in a quantum world

	Kochen-Specker bound	Quantum bound
	$R \leq \frac{n-1}{n} = 1 - \frac{1}{n}$	$R \leq \frac{2 \cos(\frac{\pi}{n})}{1 + \cos(\frac{\pi}{n})} \simeq 1 - \frac{\pi^2}{4n^2}$

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Suppose $n \ll \text{no. of suitors} \ll n^2$

Suitors believe it is very likely that **one of them will win**

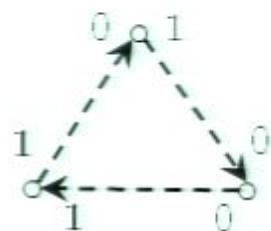
In fact it is very likely that **none of them will win**

The failure of transitivity of implication

Note: noncontextuality implies transitivity of implication

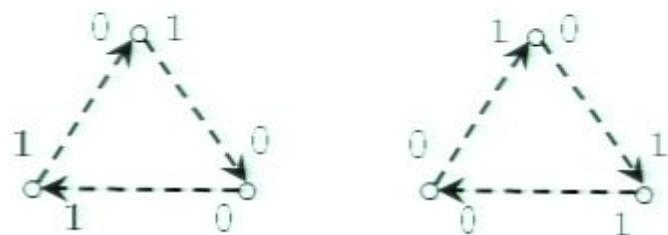
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CONTRADICTION!

The failure of transitivity of implication

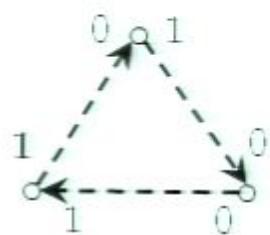
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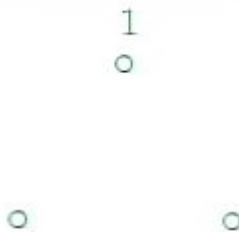
CONTRADICTION!

Alternatively:

Always



Sometimes



CONTRADICTION!

The failure of transitivity of implication

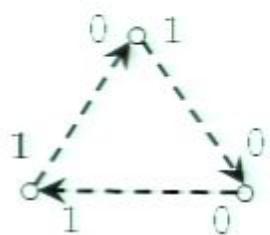
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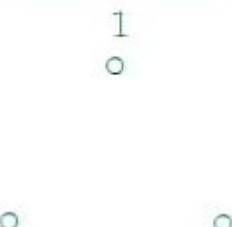
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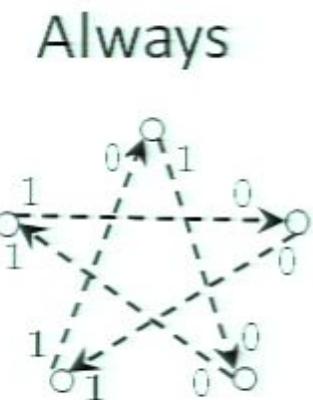


CONTRADICTION!

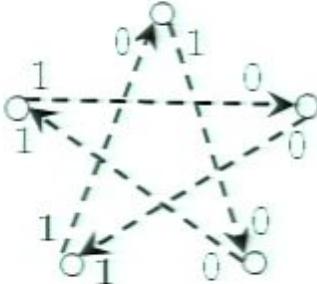
A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

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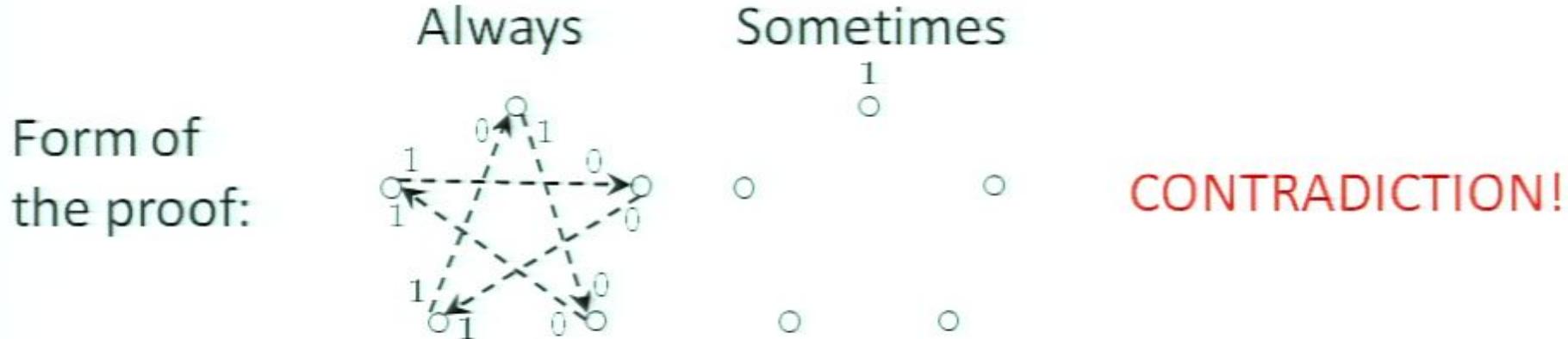
Form of
the proof:



A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

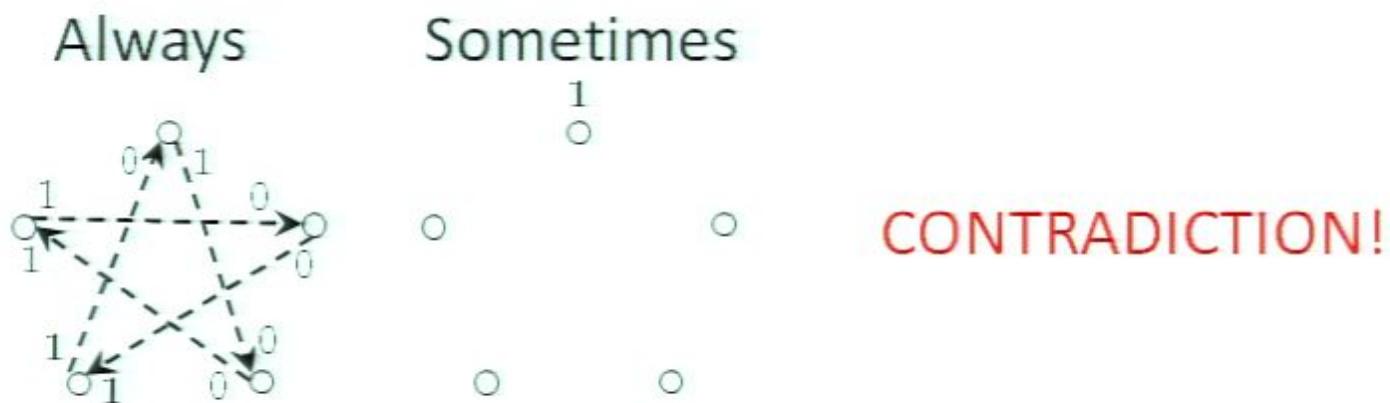
Form of the proof:	Always	Sometimes
		 

A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

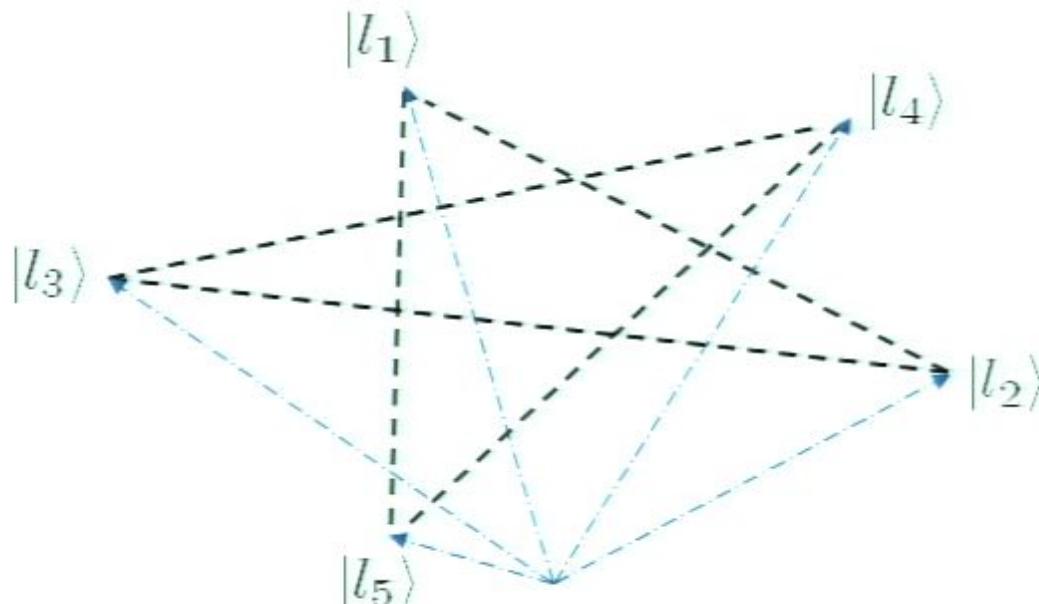


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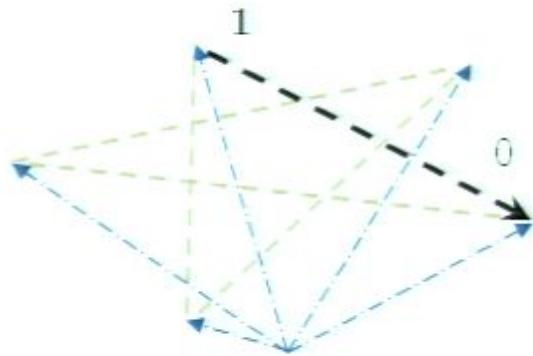
Form of the proof:



We use:



A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

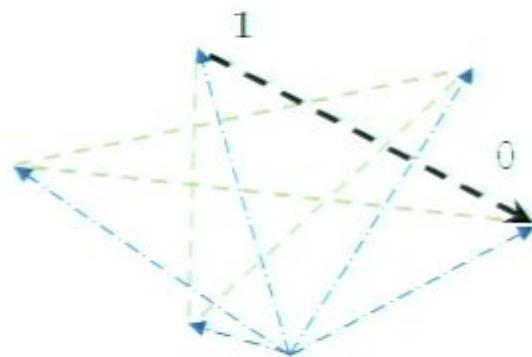


For all states

If $|l_1\rangle \perp |l_2\rangle$

Then $v(|l_1\rangle\langle l_1|) = 1 \implies v(|l_2\rangle\langle l_2|) = 0$

A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

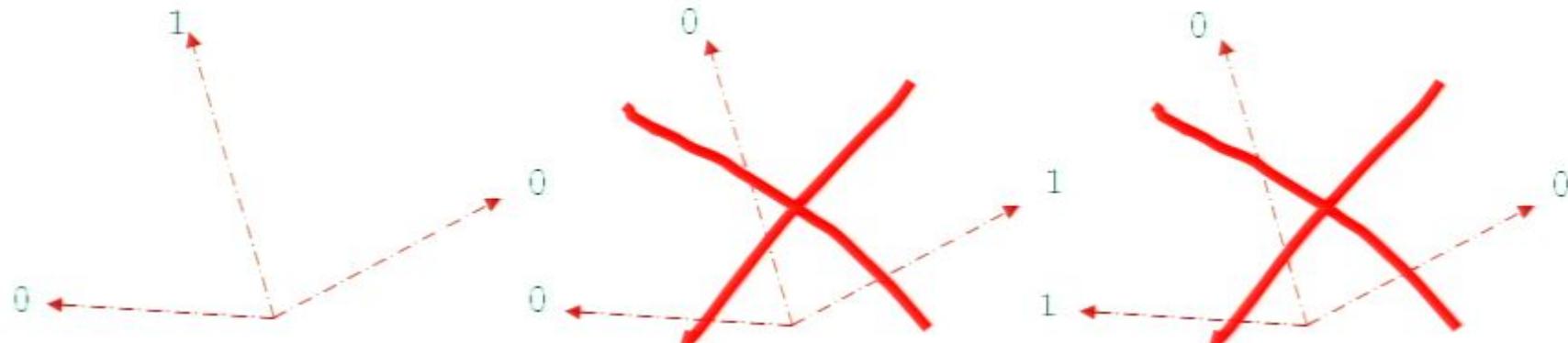


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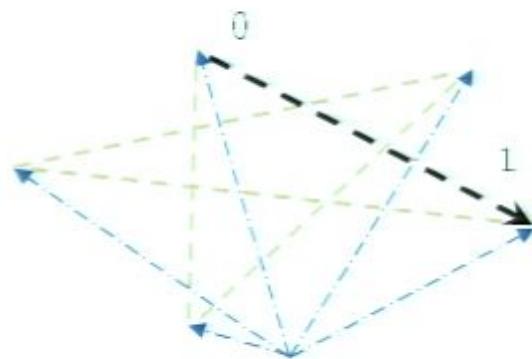
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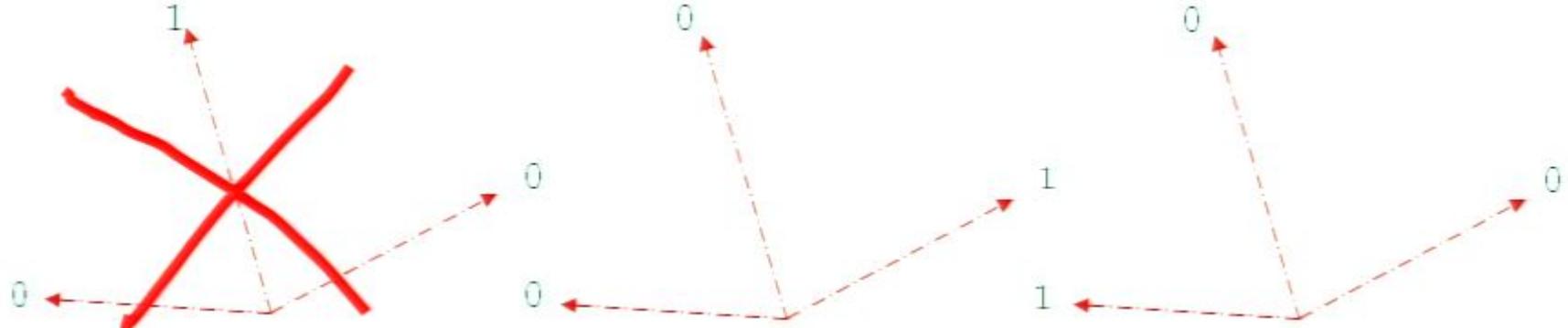


A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

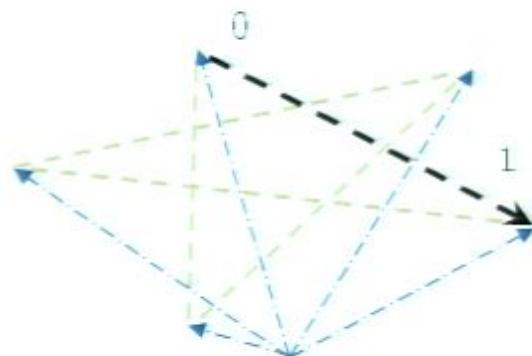


BUT NOT THE CASE THAT
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A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication



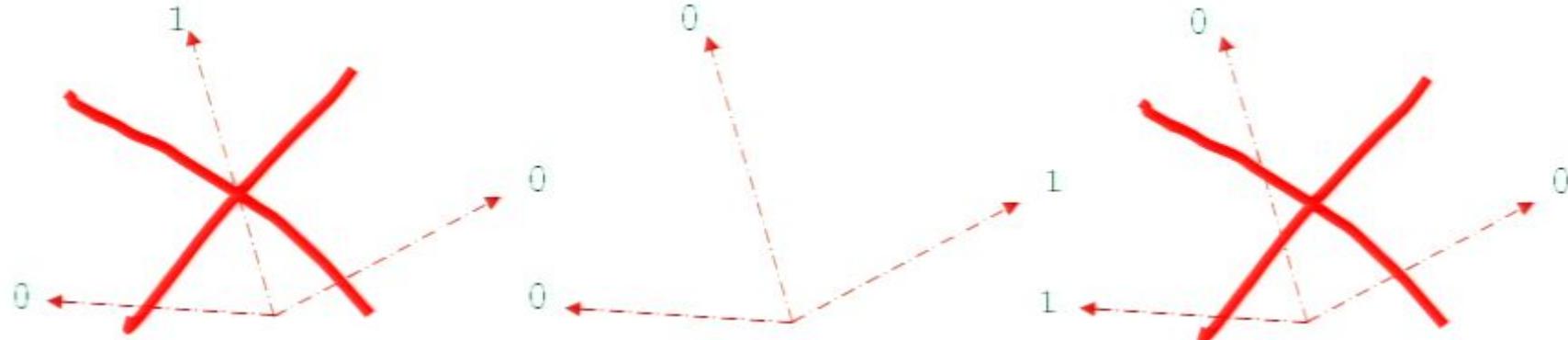
NONETHELESS

for all states $|\psi\rangle \in \text{span}(|l_1\rangle, |l_2\rangle)$

If $|l_1\rangle \perp |l_2\rangle$

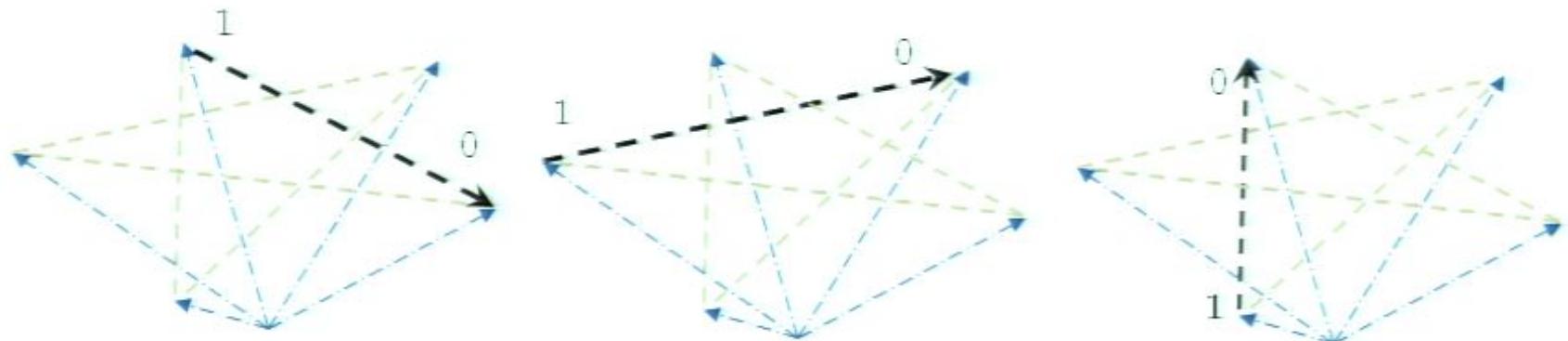
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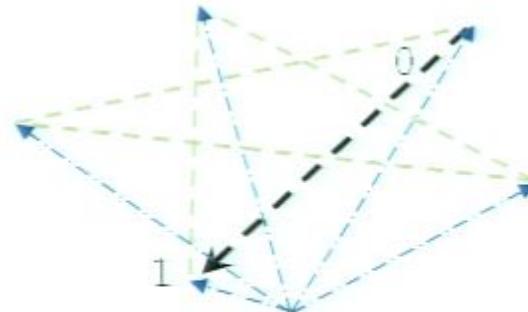
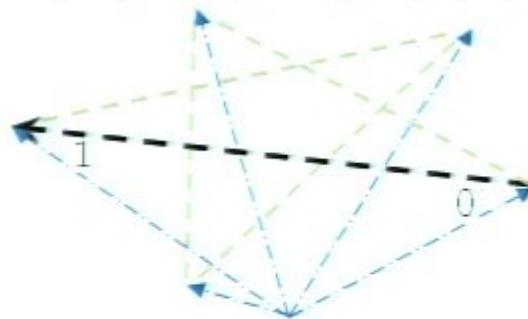
A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

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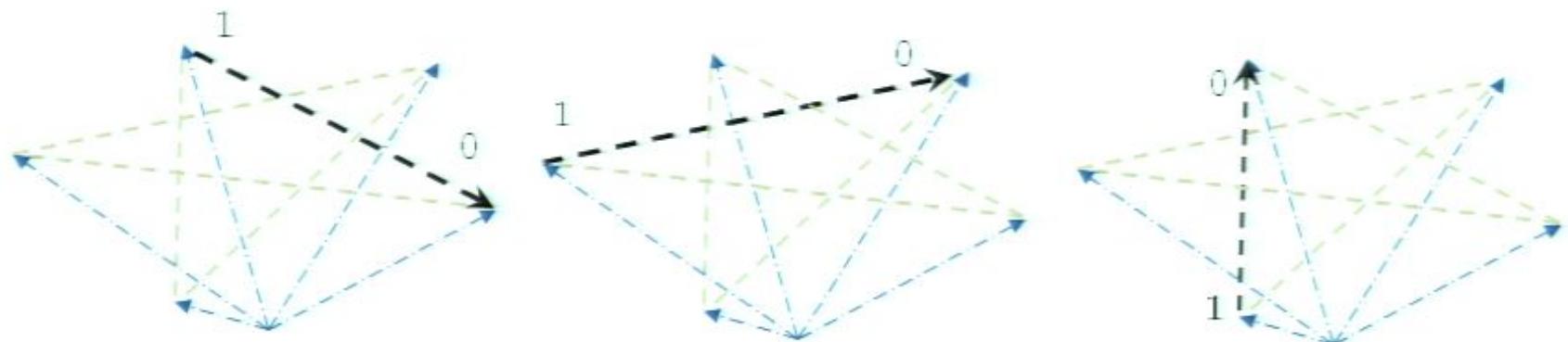
$$\forall |\psi\rangle \in \text{span}(|l_2\rangle, |l_3\rangle)$$

$$\forall |\psi\rangle \in \text{span}(|l_4\rangle, |l_5\rangle)$$



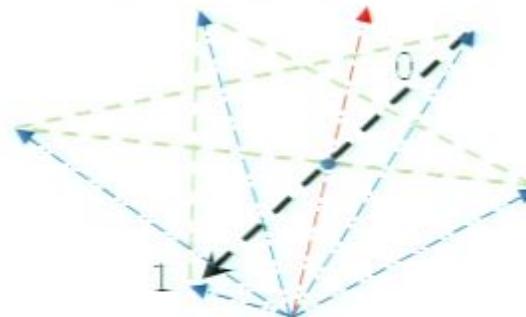
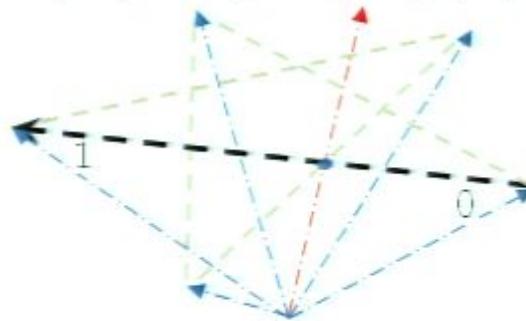
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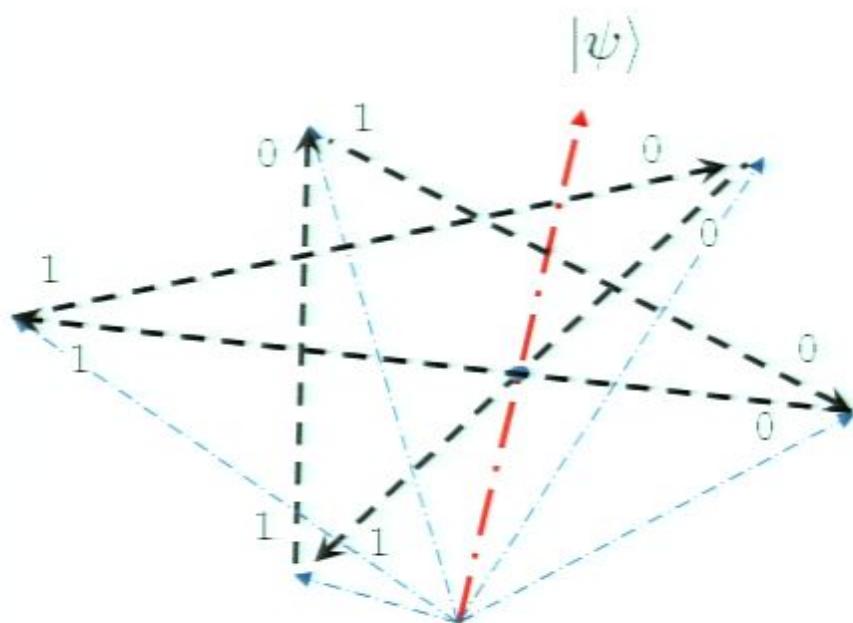
$$\forall |\psi\rangle \in \text{span}(|l_4\rangle, |l_5\rangle)$$



Therefore choose $|\psi\rangle \in \text{span}(|l_2\rangle, |l_3\rangle) \cap \text{span}(|l_4\rangle, |l_5\rangle)$

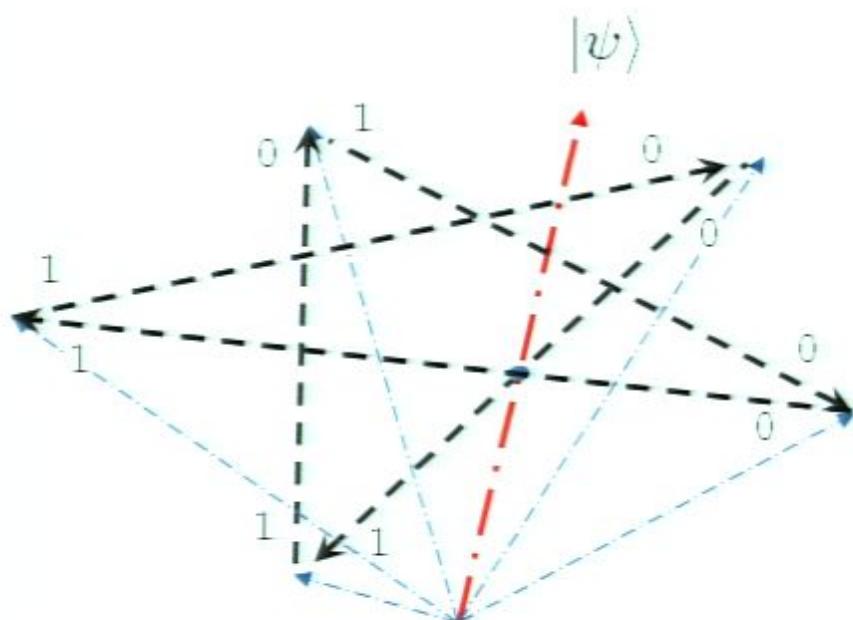
A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

Always

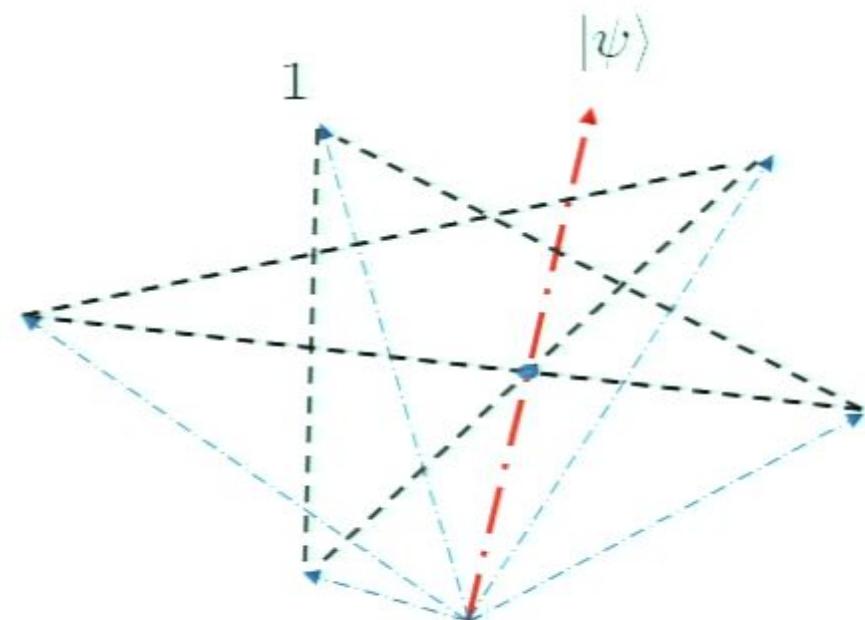


A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

Always



Sometimes

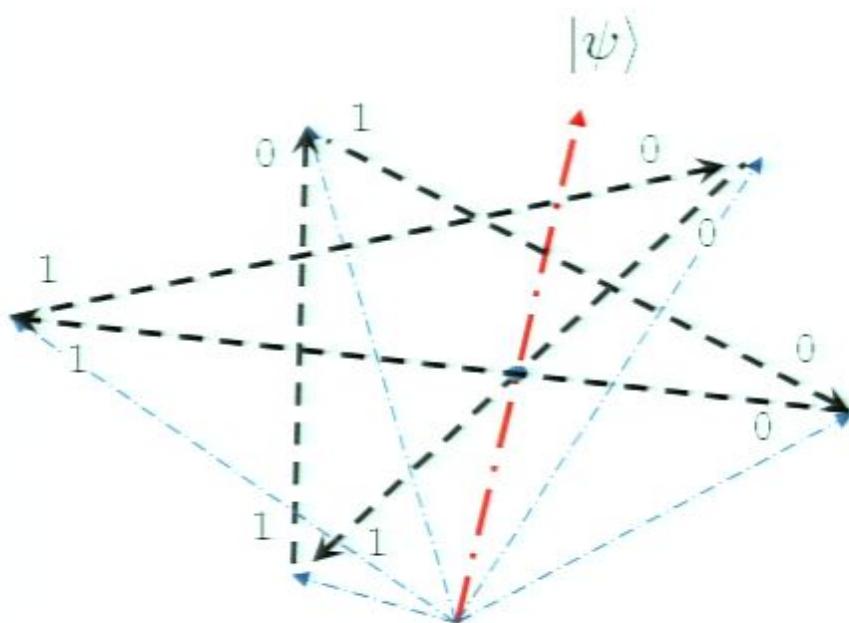


i.e. with probability

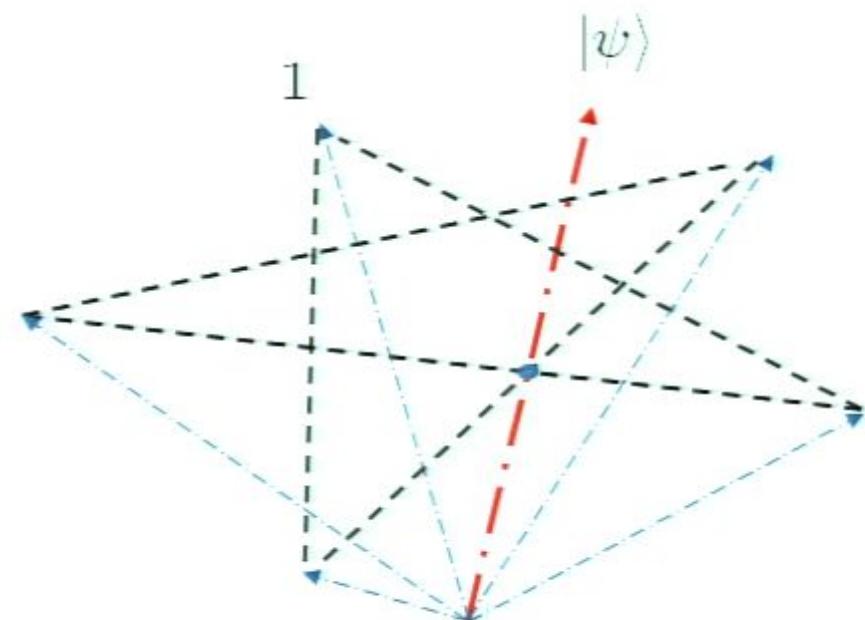
$$|\langle l_1 |\psi \rangle|^2 = 1 - \frac{2}{\sqrt{5}} \simeq 0.1056$$

A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication

Always



Sometimes



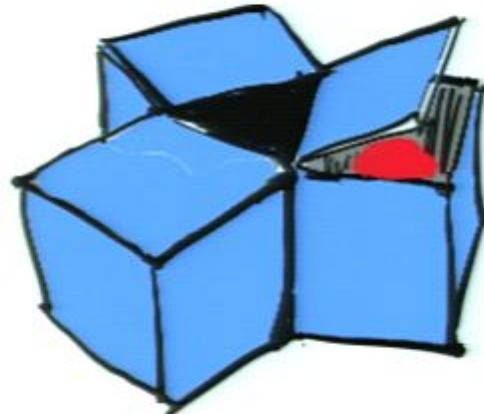
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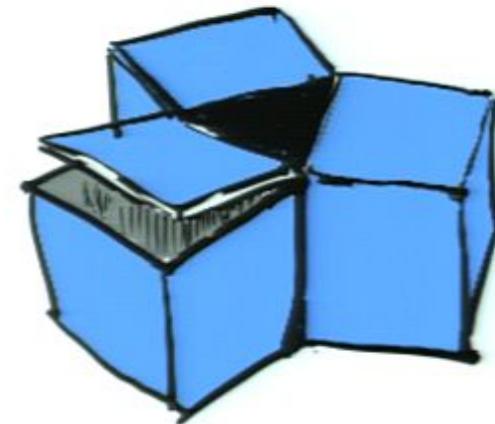
Page 74/138

A separated pair of single-query 3-box systems

Abydos

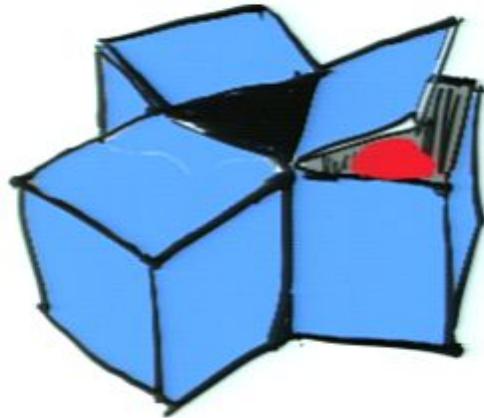


Babylon

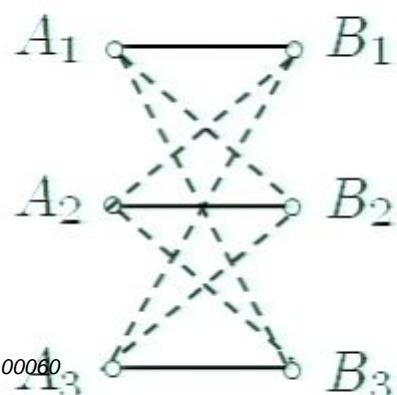
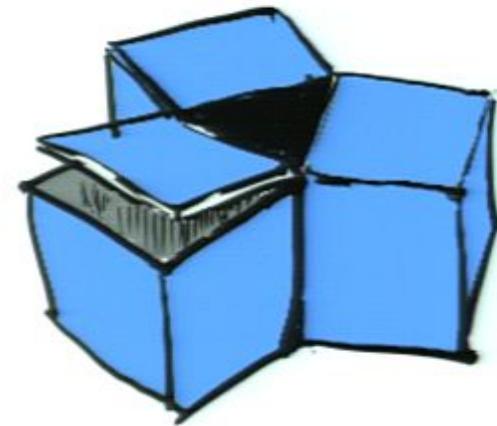


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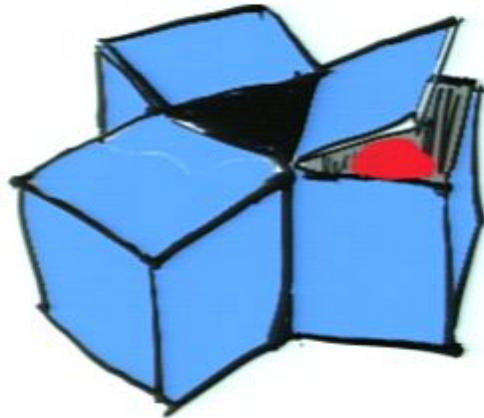


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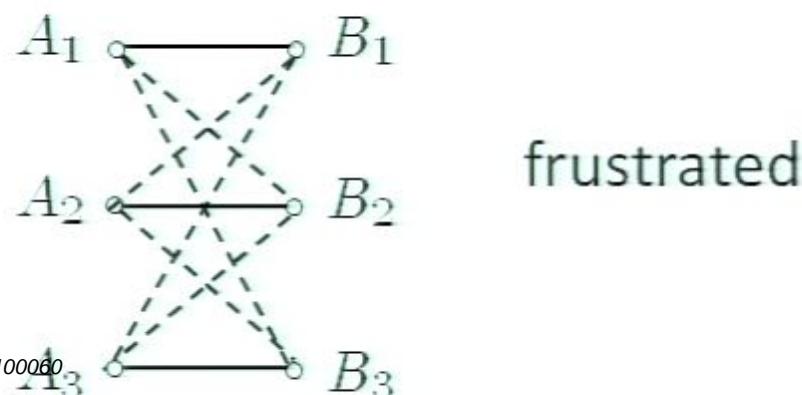
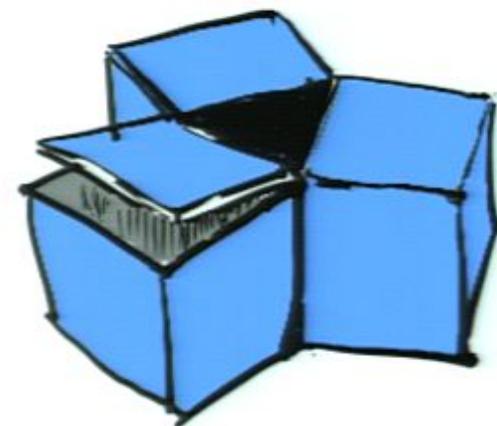


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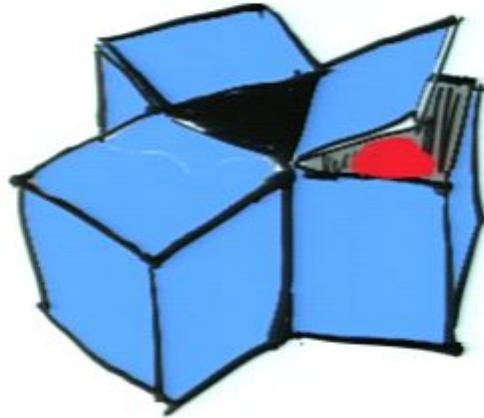


Babylon

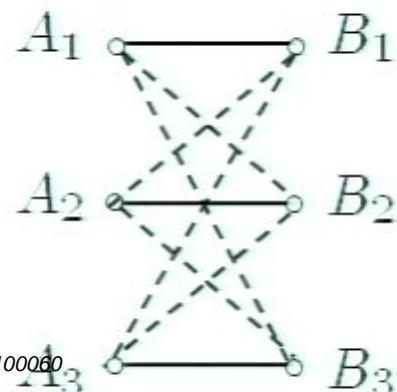
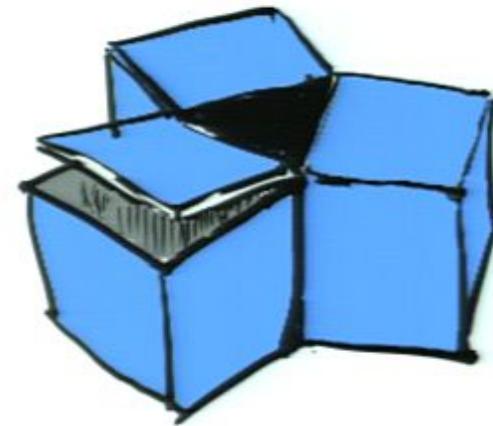


A separated pair of single-query 3-box systems

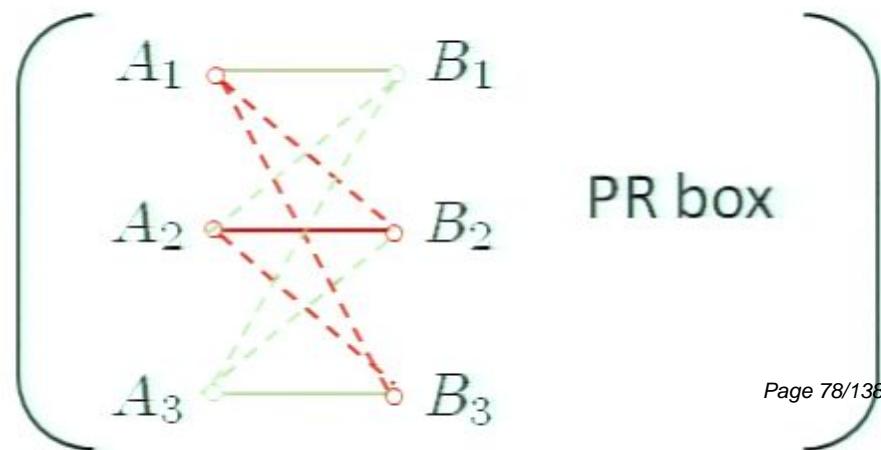
Abydos



Babylon



frustrated

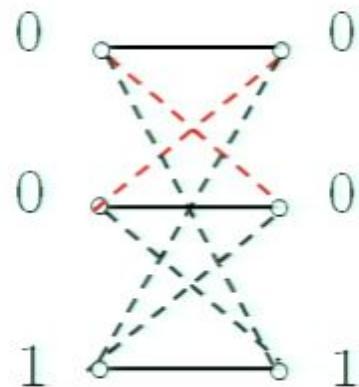


Bell inequality

Recall: Bell locality + perfect correlation
⇒ deterministic noncontextual values

Bell inequality

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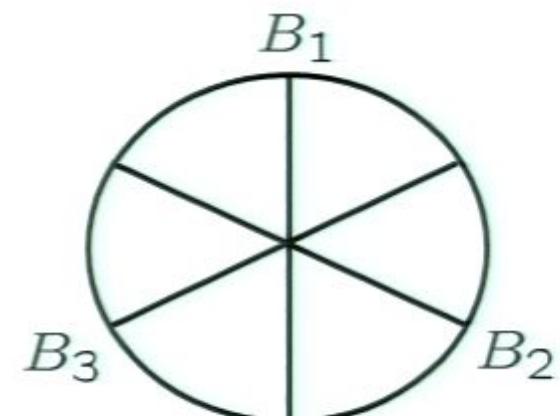
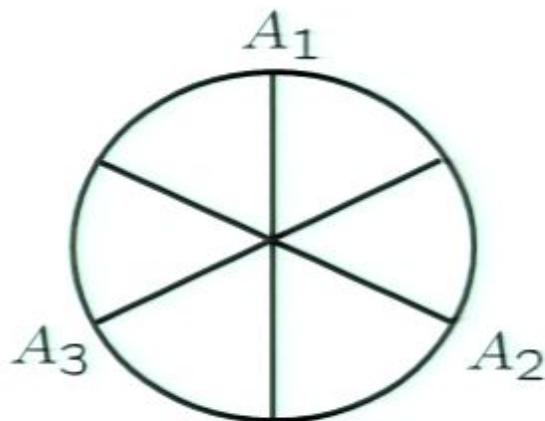
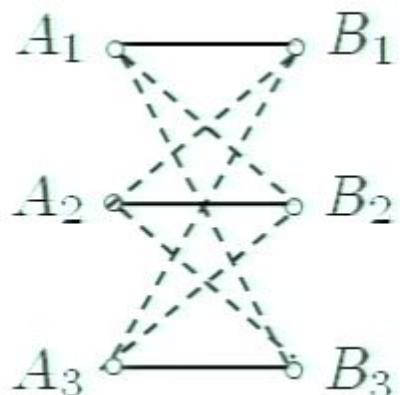


$$R \leq \frac{7}{9}$$

Local bound

Mermin's proof of Bell's theorem

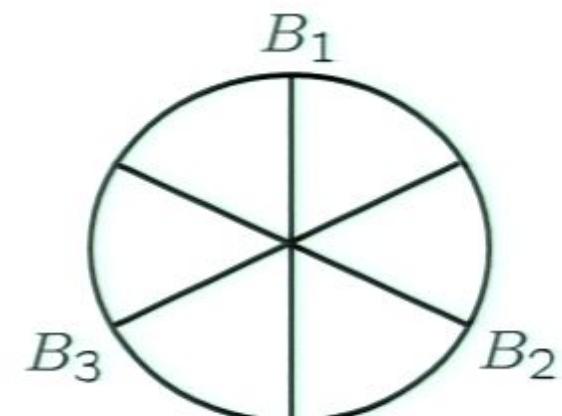
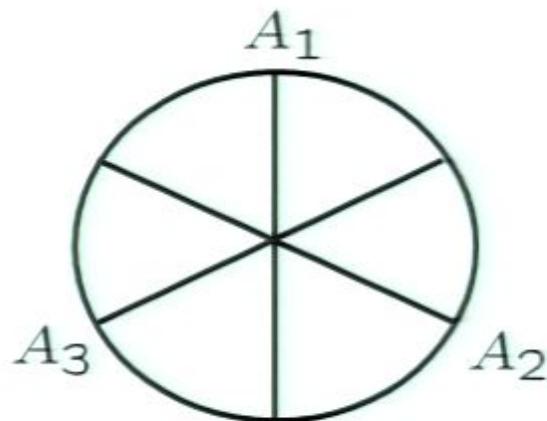
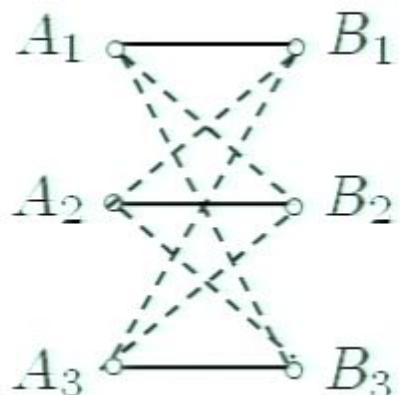
Mermin, Phys. Today 38 (4), 38 (1985)



$$|\psi\rangle_{AB} = \frac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$

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Mermin, Phys. Today 38 (4), 38 (1985)



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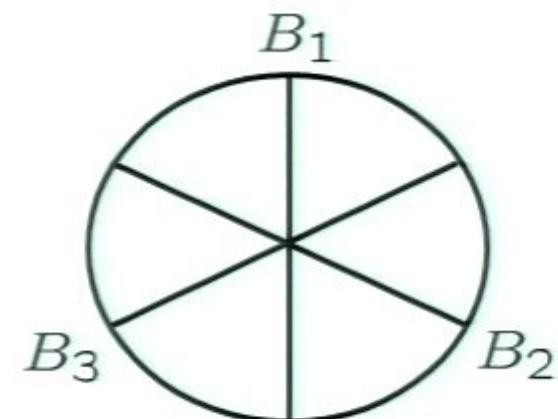
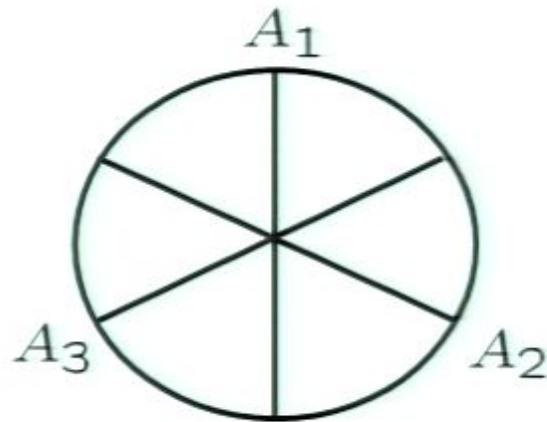
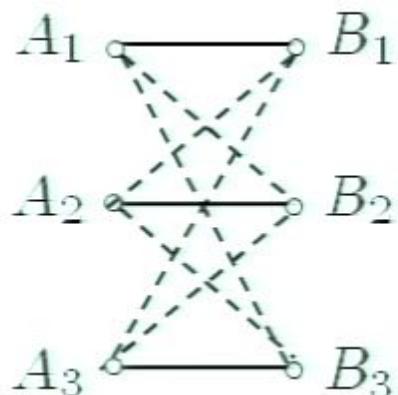
$$p(A_k = B_k) = 1$$

$$p(A_k \neq B_{k \oplus 1}) = \cos^2\left(\frac{\pi}{6}\right) = \frac{3}{4}$$

$$R = \frac{1}{3}(1) + \frac{2}{3}\left(\frac{3}{4}\right)$$

Mermin's proof of Bell's theorem

Mermin, Phys. Today 38 (4), 38 (1985)



$$|\psi\rangle_{AB} = \frac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$

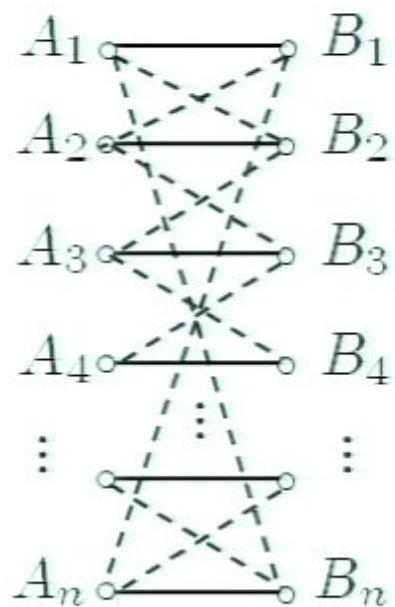
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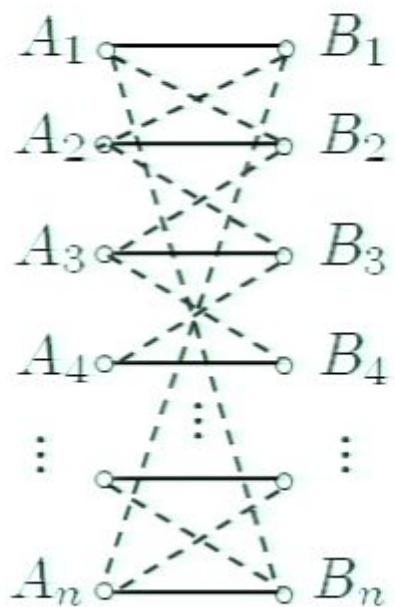
Quantum violation of Bell inequality

A separated pair of single-query n -box systems

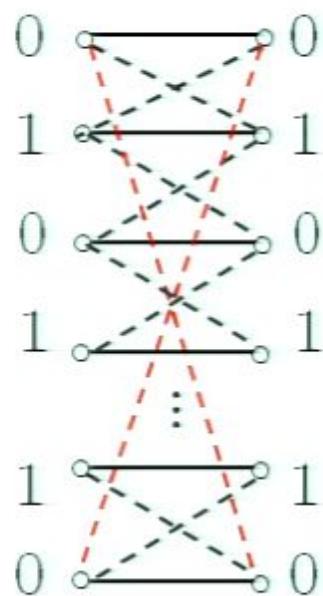


n odd

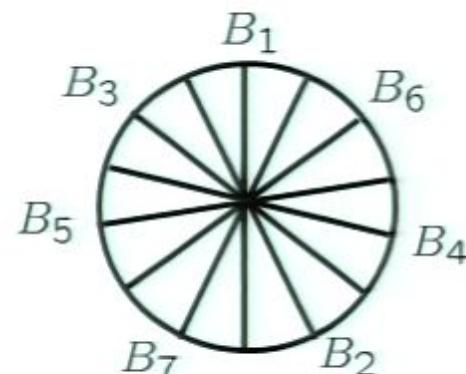
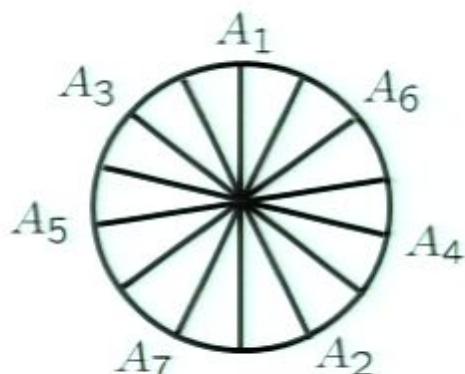
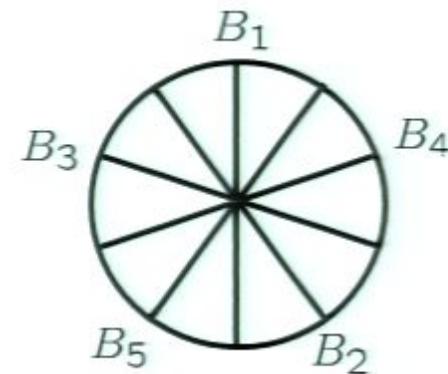
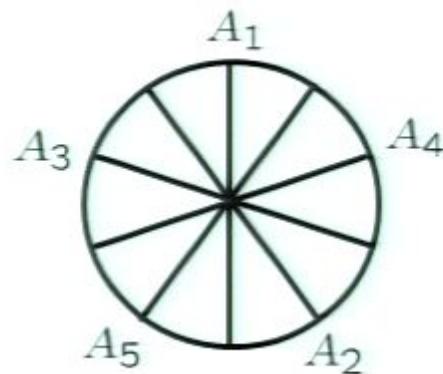
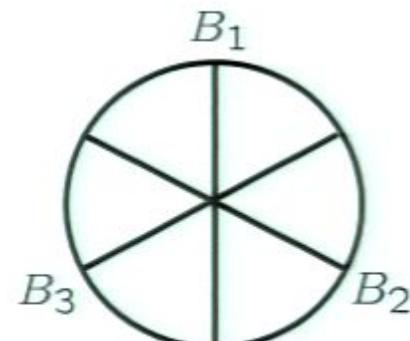
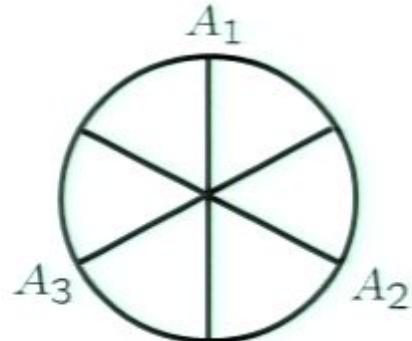
A separated pair of single-query n -box systems



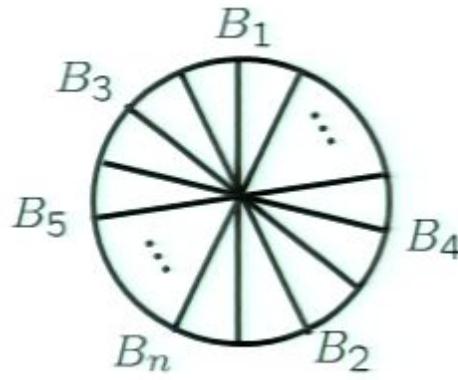
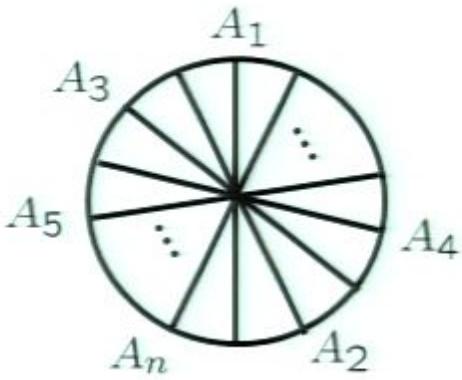
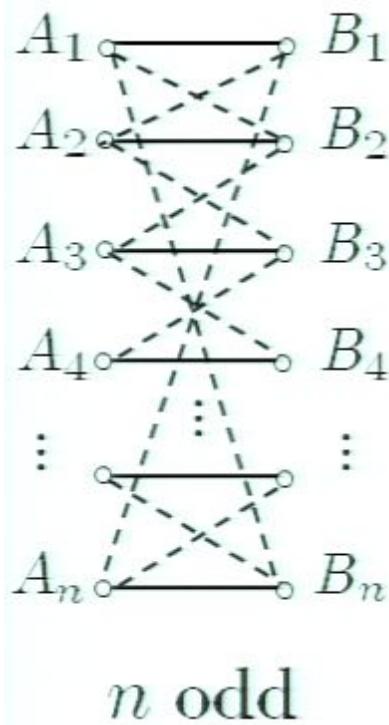
n odd



Generalization of Mermin's proof



Generalization of Mermin's proof



$$p(A_k = B_k) = 1$$

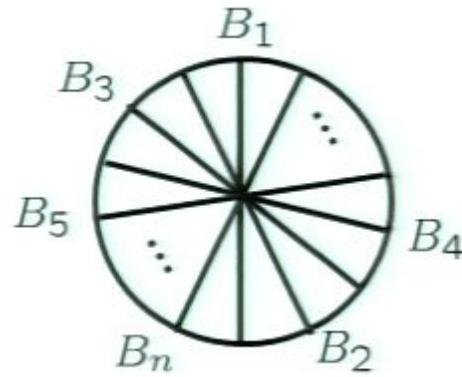
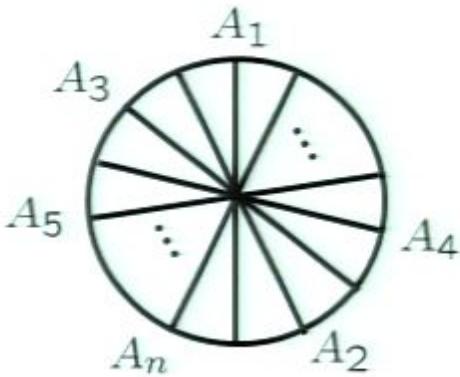
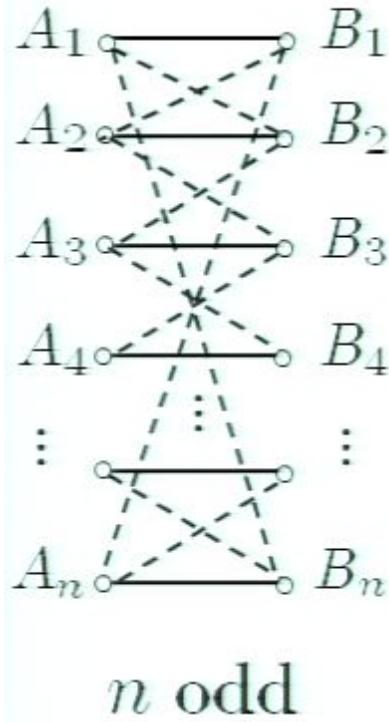
$$p(A_k \neq B_{k+1}) = \cos^2\left(\frac{\pi}{2n}\right)$$

$$\begin{aligned} R &= \frac{1}{3}(1) + \frac{2}{3} \cos^2\left(\frac{\pi}{2n}\right) \\ &\simeq 1 - \frac{\pi^2}{6n^2} \quad \text{as } n \rightarrow \infty \end{aligned}$$

Quantum violation of Bell inequality

$$R = \frac{1}{3} + \frac{2}{3} \cos^2\left(\frac{\pi}{2n}\right) \geq 1 - \frac{2}{3}$$

Generalization of Mermin's proof



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Compare to:
Chained Bell
inequalities

Quantum violation of Bell inequality

$$R = \frac{1}{3} + \frac{2}{3} \cos^2\left(\frac{\pi}{2}\right) \geq 1 - \frac{2}{3}$$

Hardy's proof of Bell's theorem

Hardy, PRL 71, 1665 (1993)

2 settings at each wing

2 outcomes for each measurement

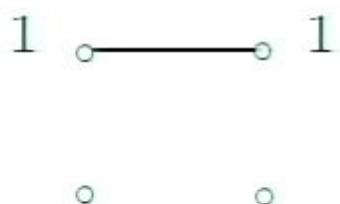
$$A_1 \circ \quad \circ B_1$$

$$A_2 \circ \quad \circ B_2$$

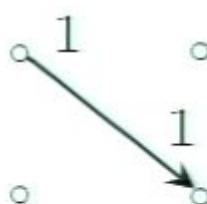
Hardy's proof of Bell's theorem

It is possible to find A_1, A_2, B_1 and B_2 and a state such that

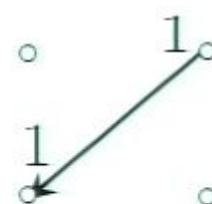
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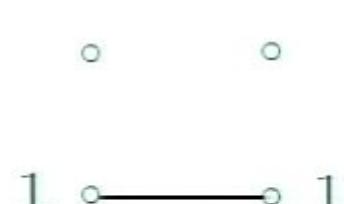
Always



Always

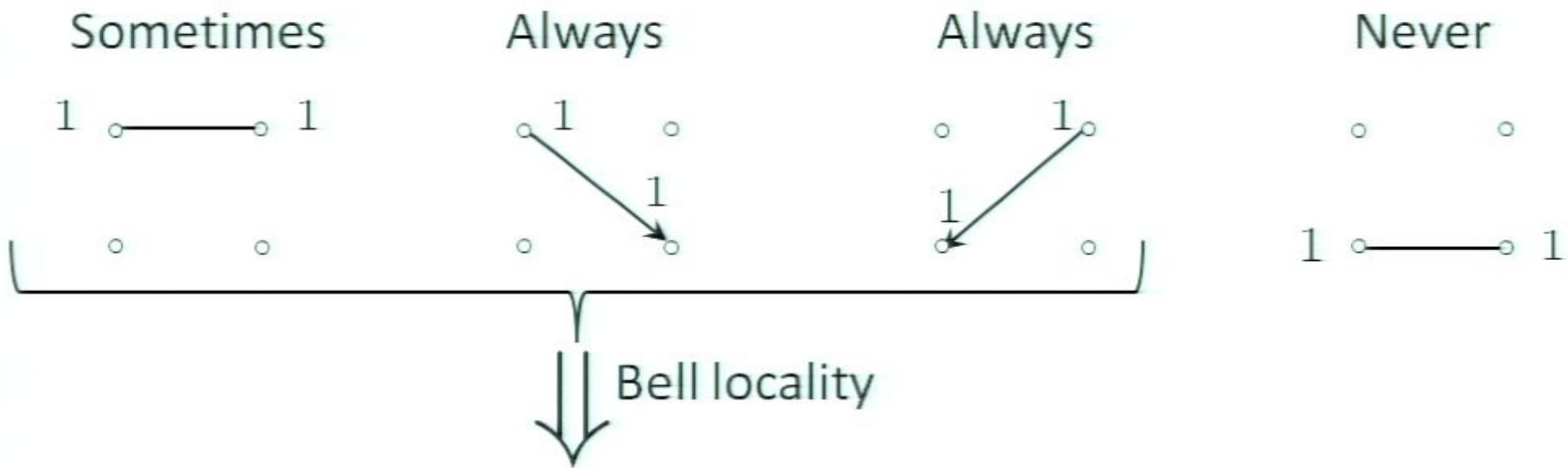


Never



Hardy's proof of Bell's theorem

It is possible to find A_1, A_2, B_1 and B_2 and a state such that



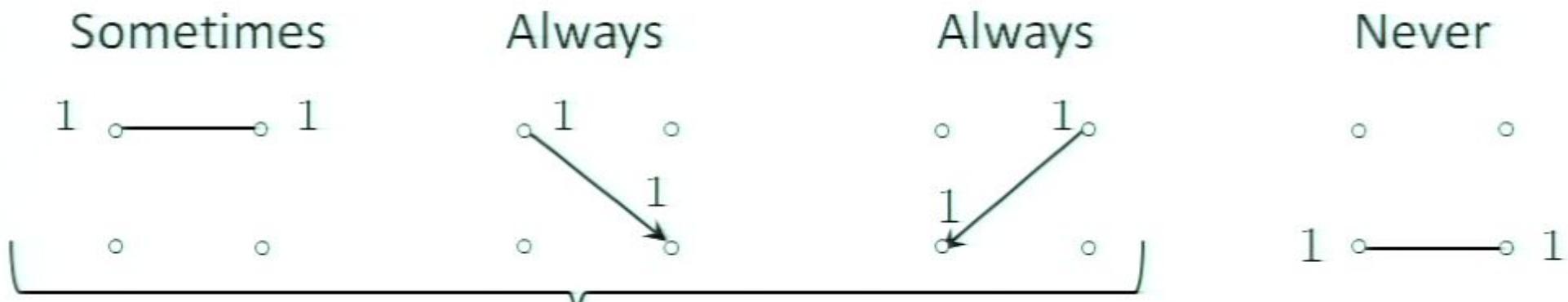
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○ ○

1 ○ — ○ 1

Hardy's proof of Bell's theorem

It is possible to find A_1, A_2, B_1 and B_2 and a state such that



↓
Bell locality

CONTRADICTION!

Sometimes

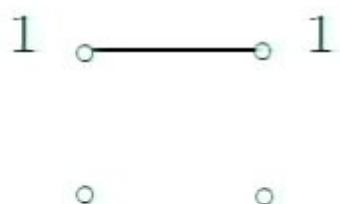
$\circ \quad \circ$

$1 \circ - \circ 1$

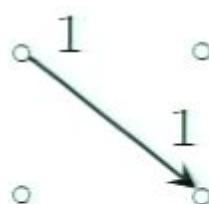
Recasting Hardy's proof as failure of transitivity

It is possible to find A_1, A_2, B_1 and B_2 and a state such that

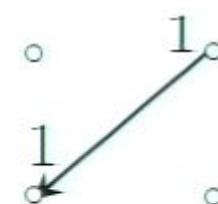
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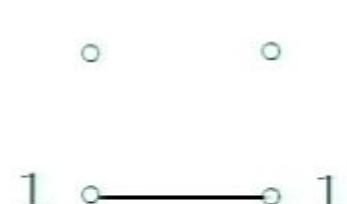
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Always



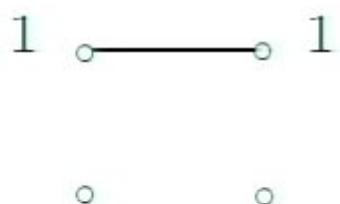
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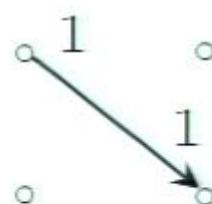
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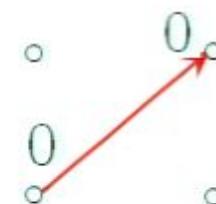
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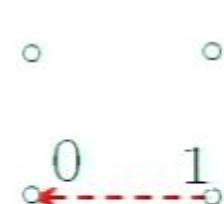
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Always



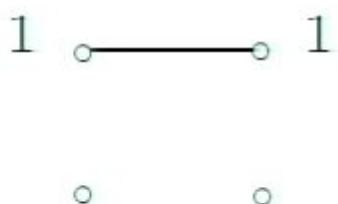
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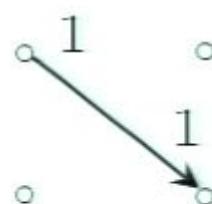
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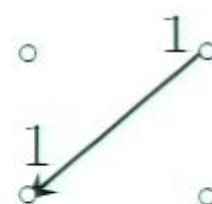
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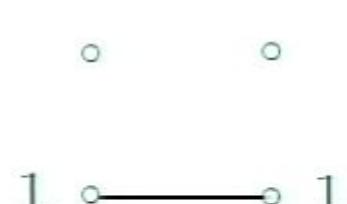
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Always



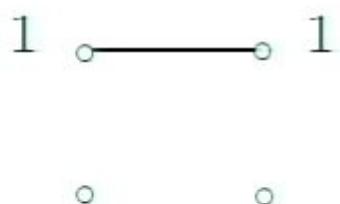
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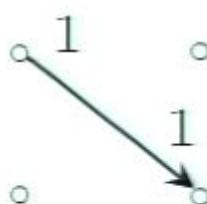
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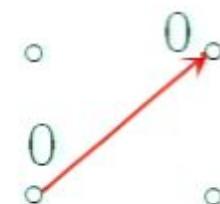
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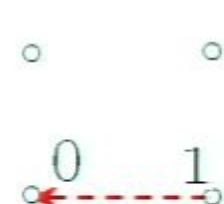
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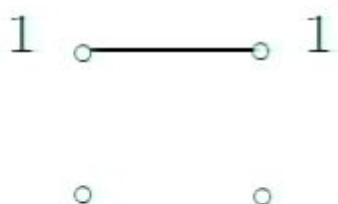
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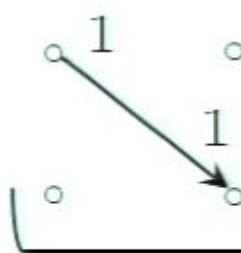
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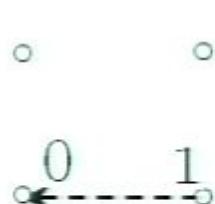
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Always



Always



Transitivity

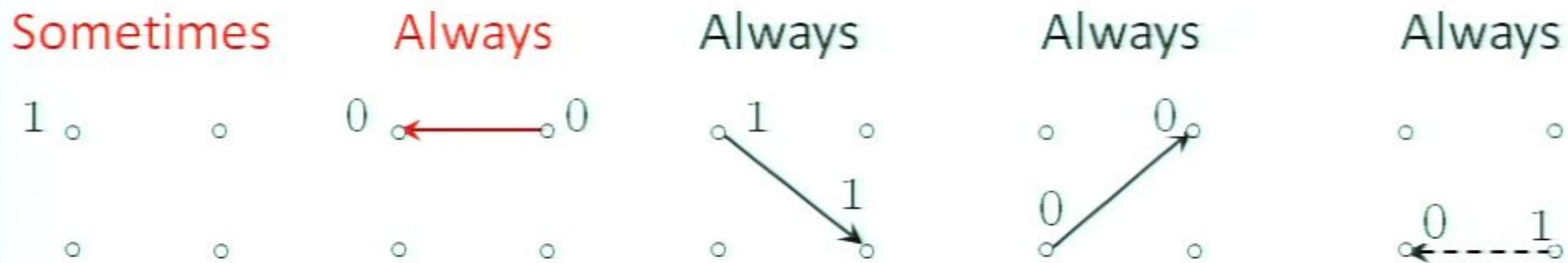
CONTRADICTION!

Always



How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that



How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that

Sometimes

1 o

o

Always

0 o ← 0

o

o

Always

1 o ← 1
o → 1

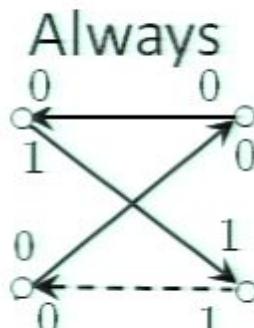
Transitivity

Always

0 o → 0
0 → 0

Always

o → 0
0 → 1



How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that

Sometimes

1 o

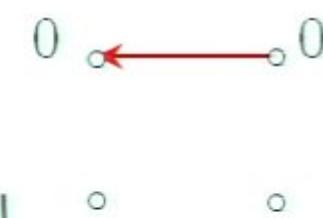
o

o

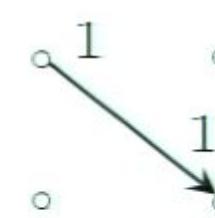
Always
0 o ← 0

o

o



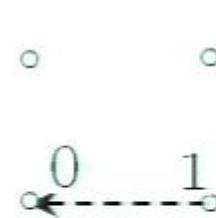
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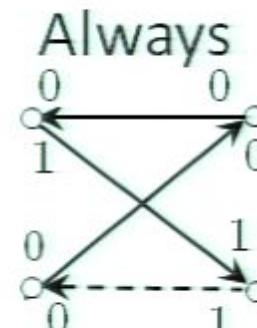


Always



Transitivity

CONTRADICTION!



How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that

Sometimes

1

0

0

Always

0

0

Always

1

0

0

Always

0

0

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Always

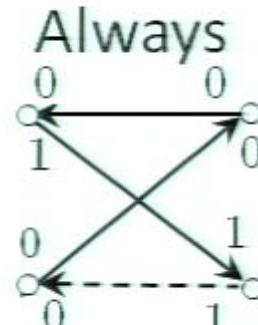
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Transitivity

CONTRADICTION!



Answer: NO!

How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that

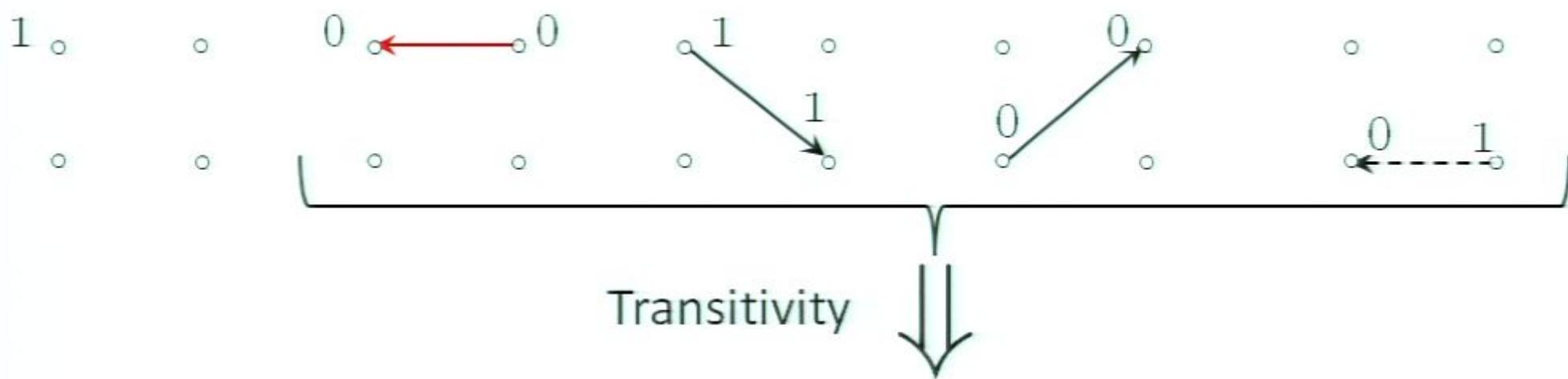
Sometimes

Always

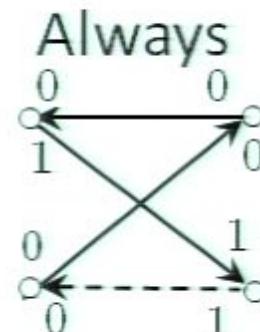
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Always

Always



CONTRADICTION!



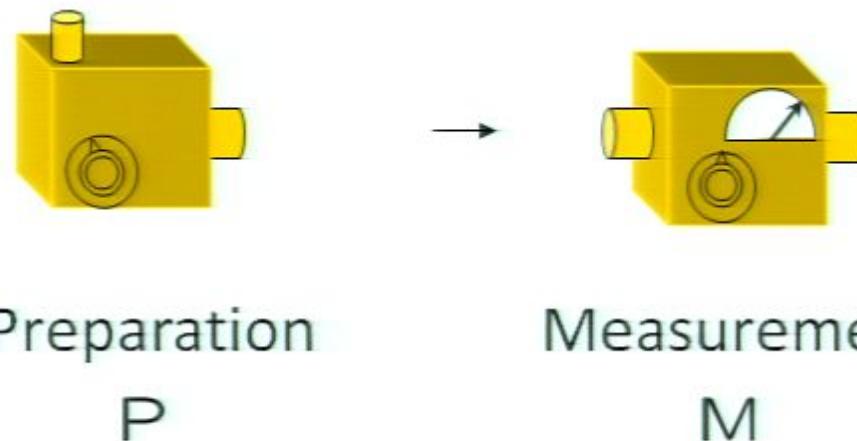
Answer: NO!

Theorem: No bipartite proof of nonlocality can have the form of a chain of implications where the final consequent denies the initial antecedent and the initial antecedent has nonzero probability.

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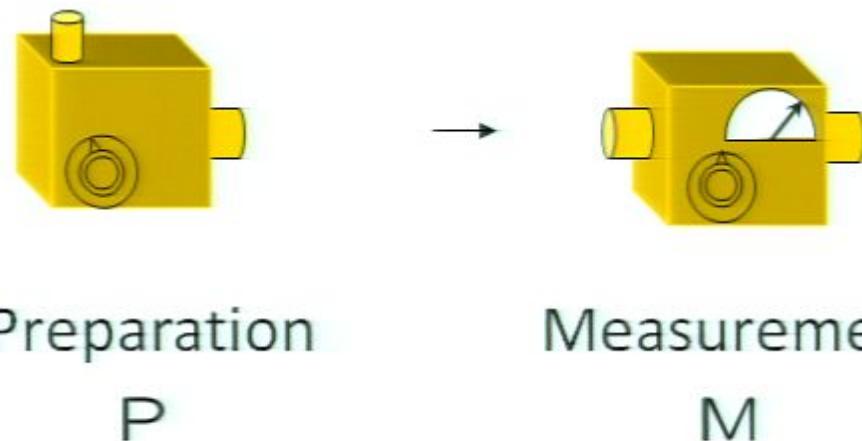
Why do there exist “striking” proofs (via the failure of transitivity) for the Kochen-Specker theorem, but not Bell’s theorem?

Operational probabilistic theories



These are defined as lists of [instructions](#)

Operational probabilistic theories

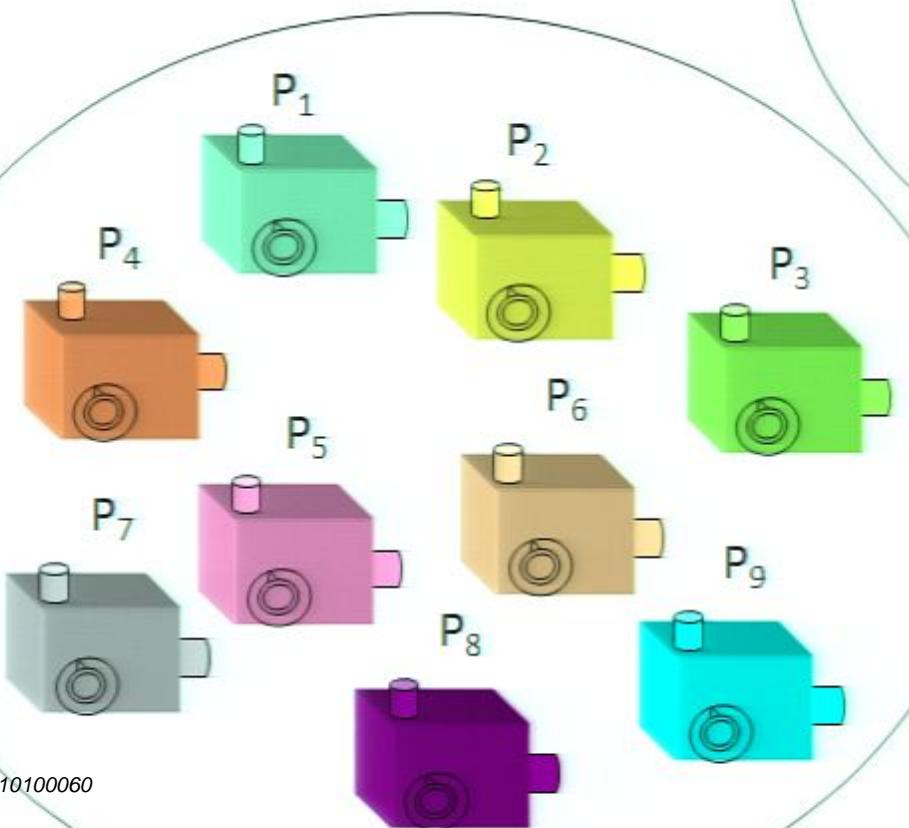
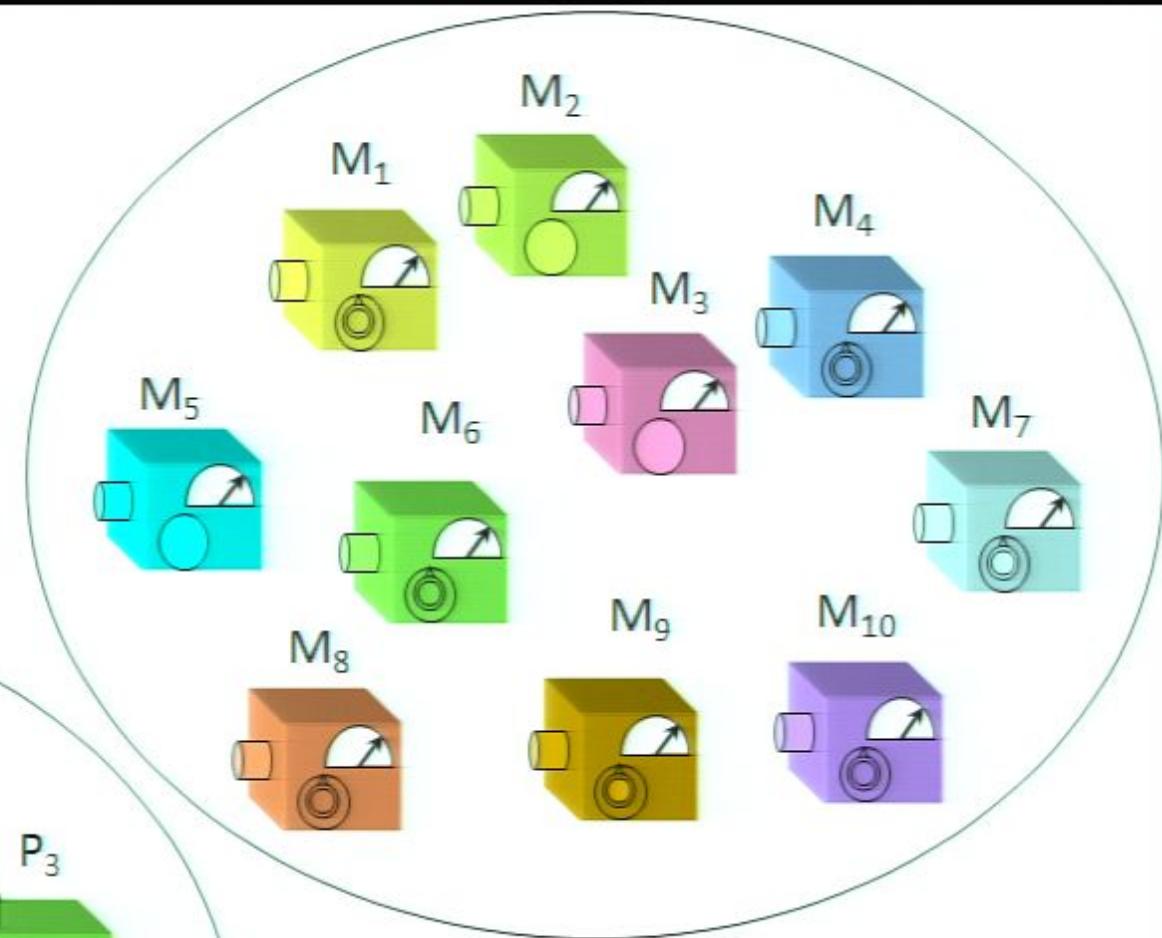


These are defined as lists of [instructions](#)

An [operational](#) theory specifies

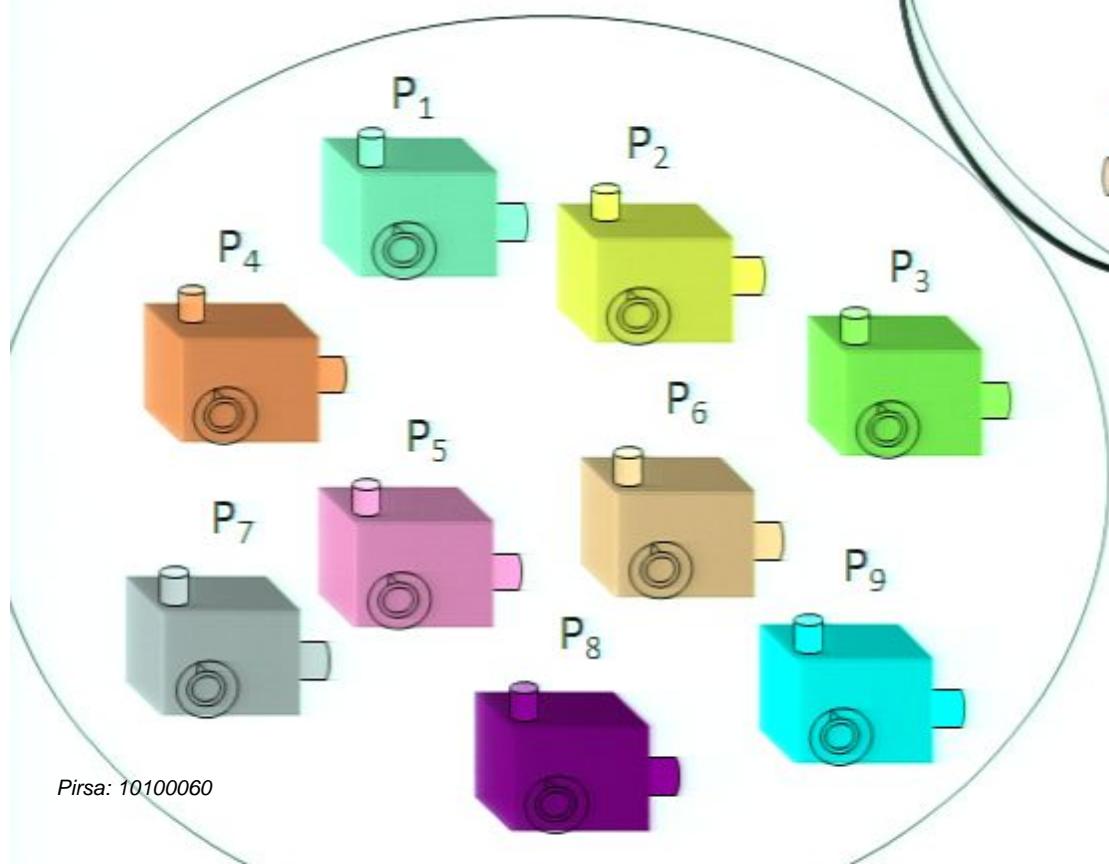
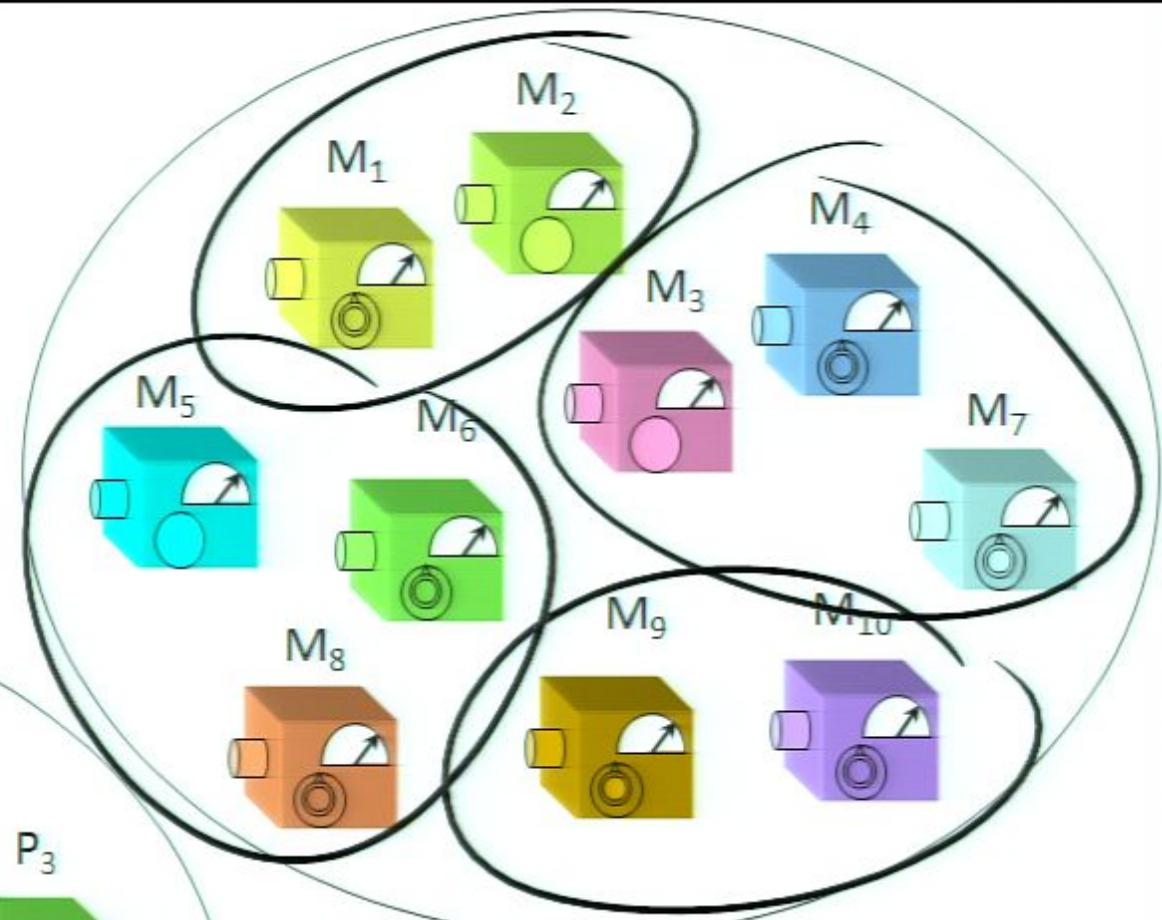
$$p(k|P, M) \equiv \begin{array}{l} \text{The probability of outcome } k \text{ of } M \\ \text{given } P \end{array}$$

Operational equivalence classes of measurements



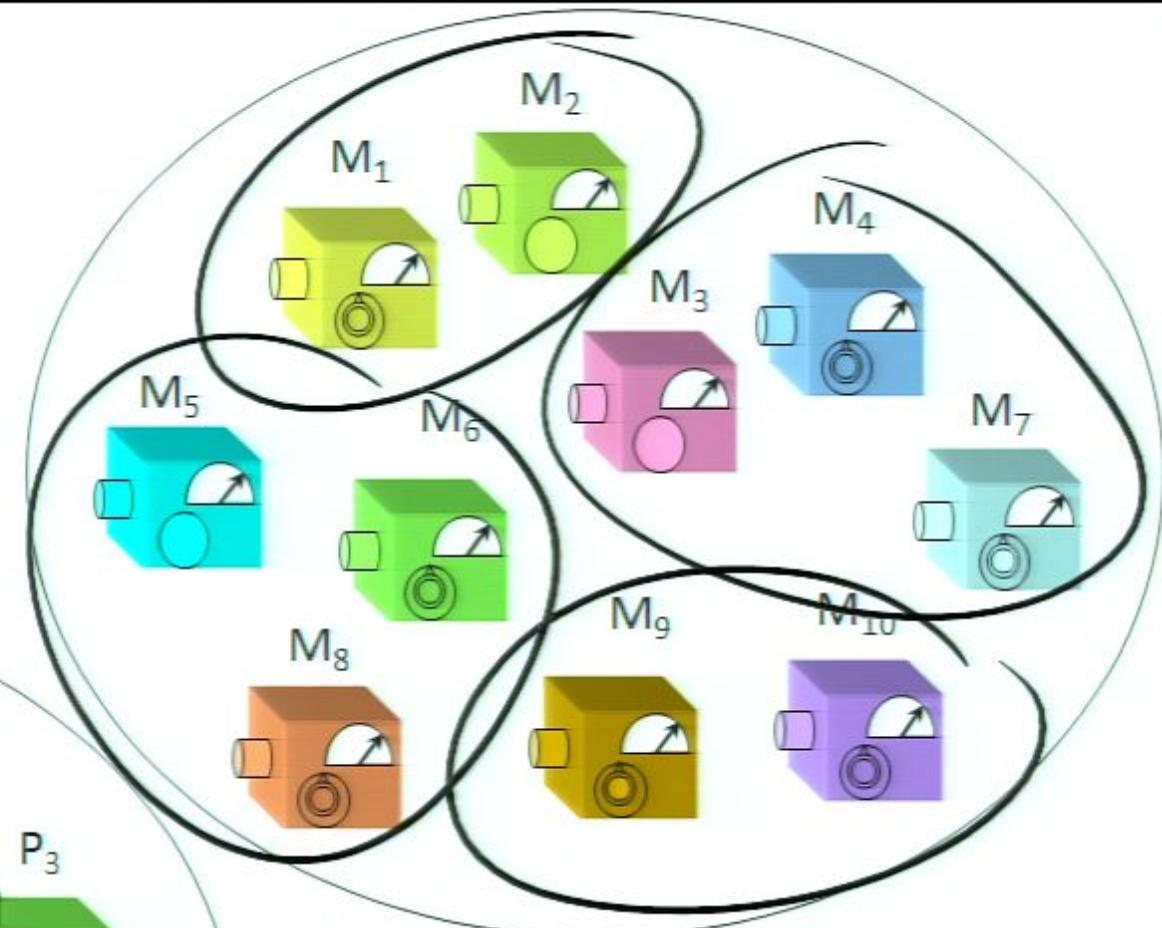
Operational equivalence classes of measurements

M is equivalent to M' if
 $\forall P : p(k|P, M) = p(k|P, M')$

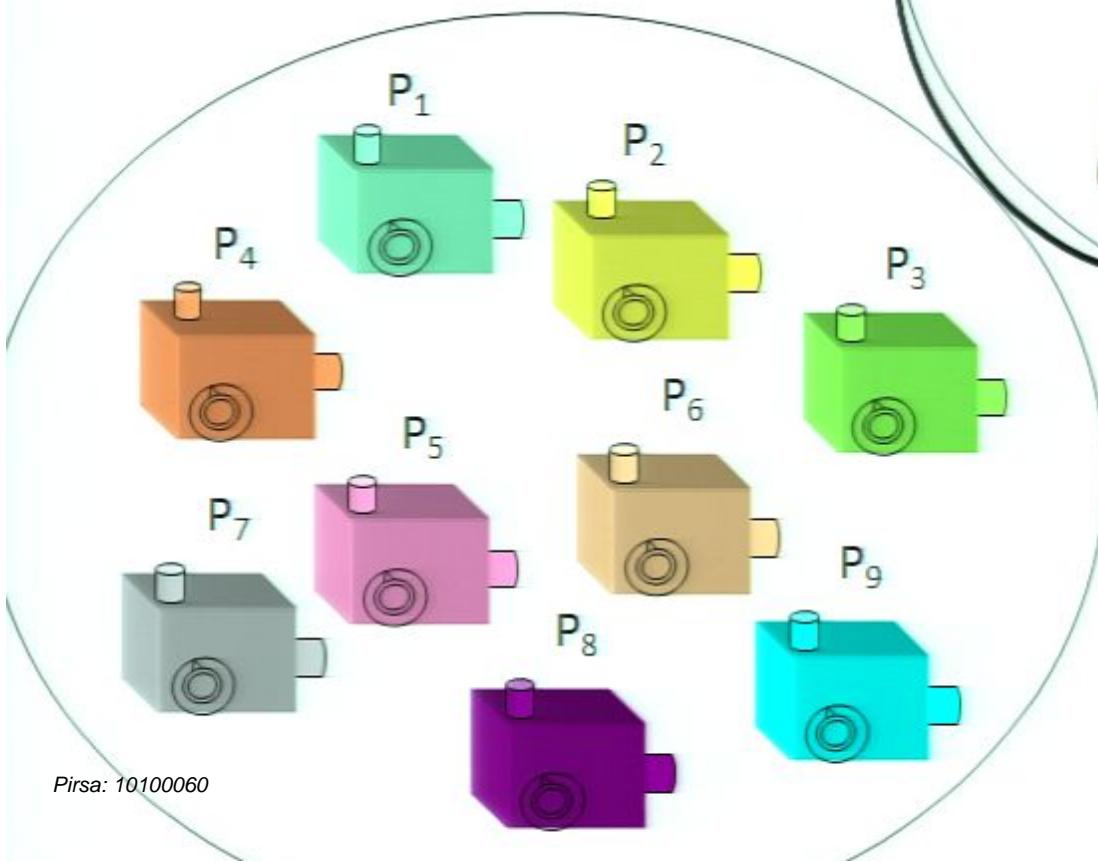


Operational equivalence classes of measurements

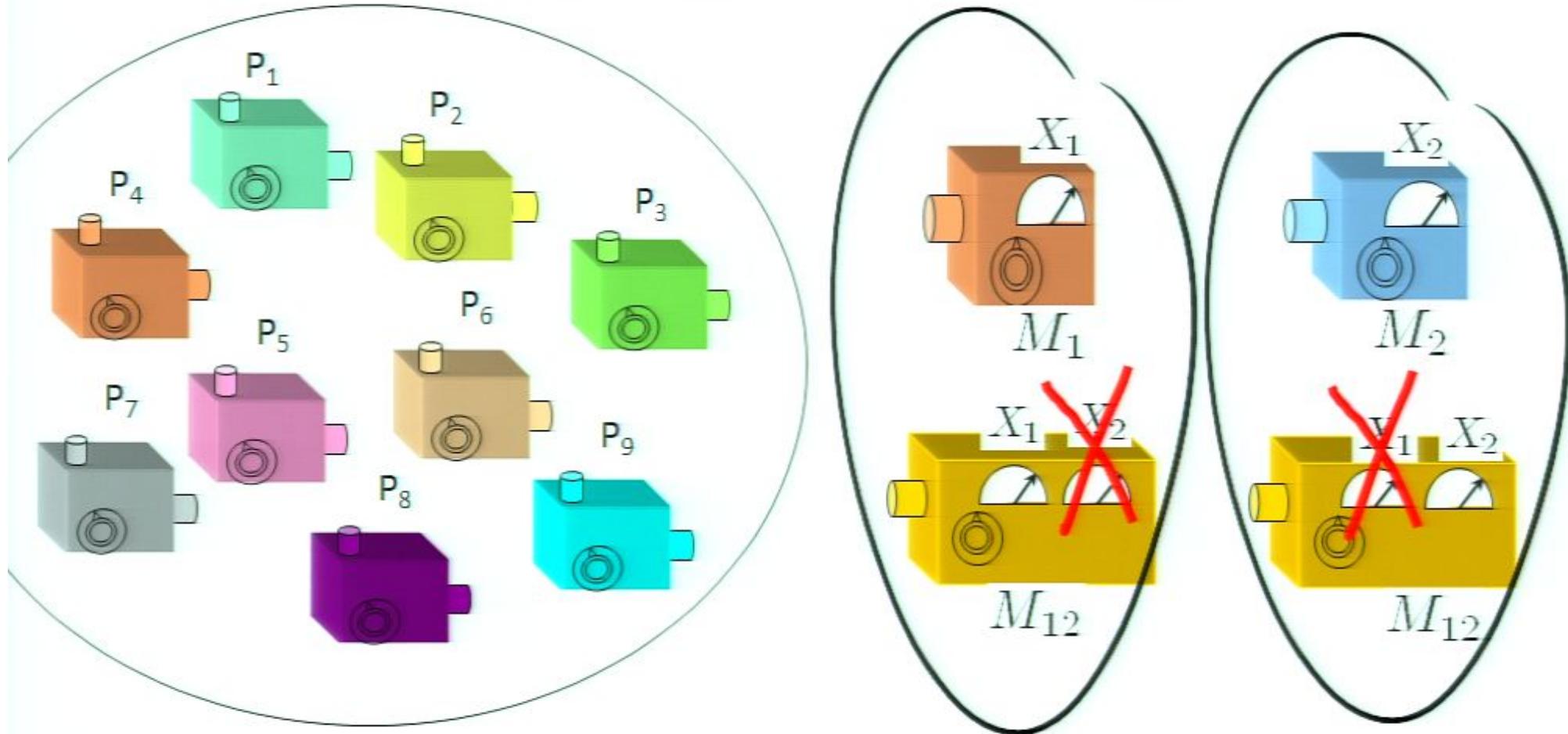
M is equivalent to M' if
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Denote the class of M by \mathcal{M}



Joint Measurability #1 (joint simulability)

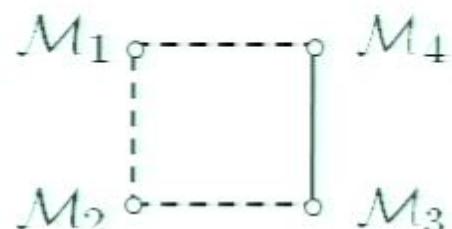


\mathcal{M}_1 and \mathcal{M}_2 are jointly measurable if $\exists \mathcal{M}_{12}$ such that

$$\forall \mathcal{F} : p(X_1 | \mathcal{M}_1, \mathcal{F}) = \sum_{X_2} p(X_1, X_2 | \mathcal{M}_{12}, \mathcal{F})$$

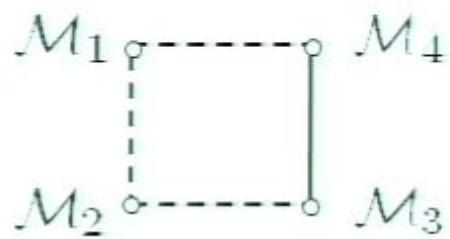
Ruling out extremal frustrated network correlations by ruling out OS correlations

Suppose a theory has

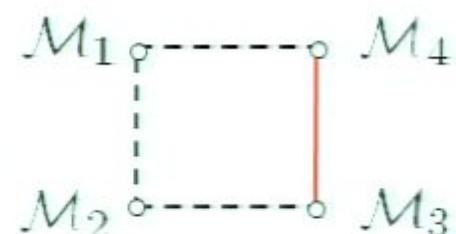


Ruling out extremal frustrated network correlations by ruling out OS correlations

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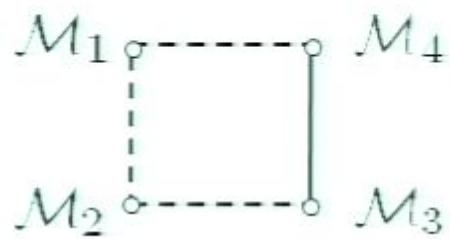


Consider the set of preparations for which

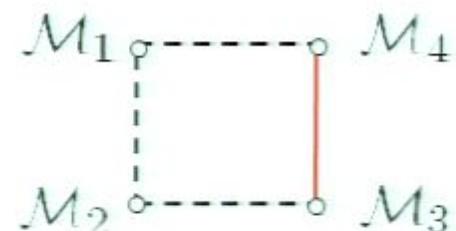


Ruling out extremal frustrated network correlations by ruling out OS correlations

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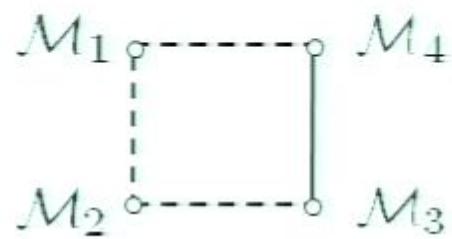
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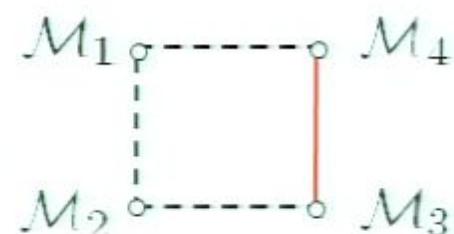
On this set, M_3 and M_4 are operationally equivalent

Ruling out extremal frustrated network correlations by ruling out OS correlations

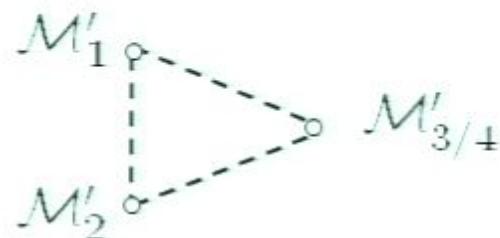
Suppose a theory has



Consider the set of preparations for which

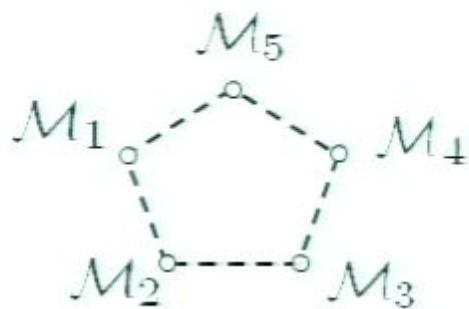


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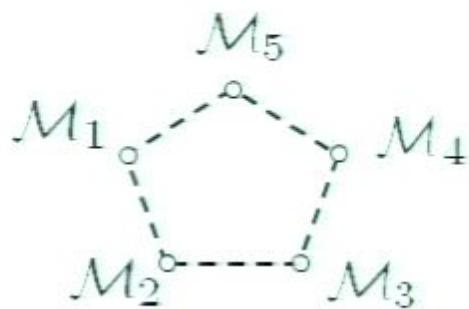
Ruling out extremal frustrated network correlations by ruling out OS correlations

Suppose a theory has

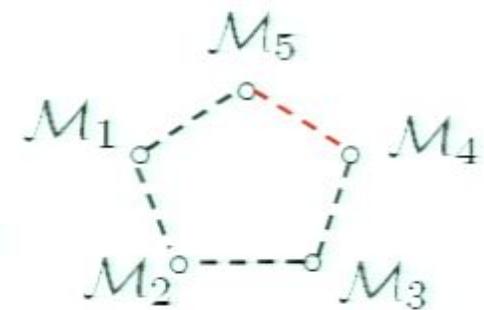


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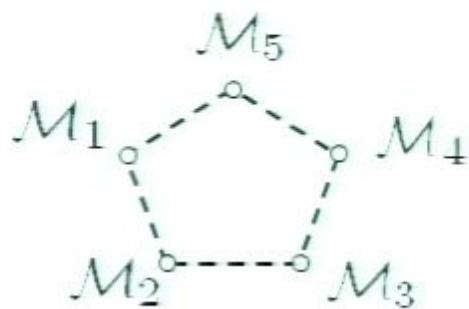


Consider the set of preparations for which

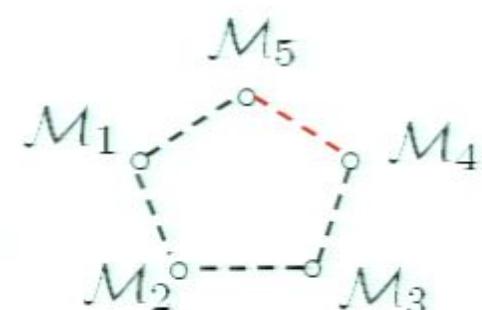


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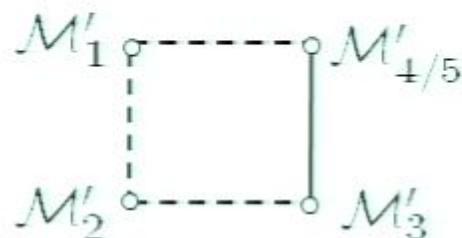
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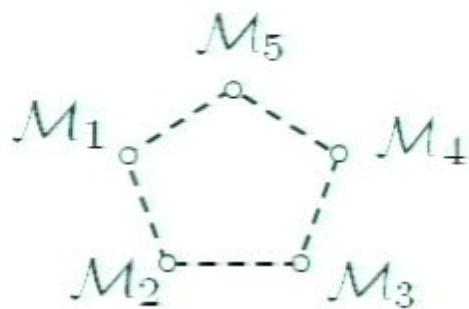


On this set, (a relabelling of) M_4 and M_5 are opt'ly equivalent

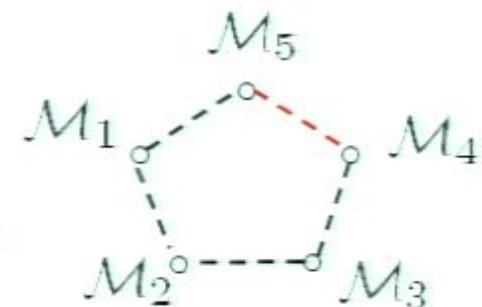


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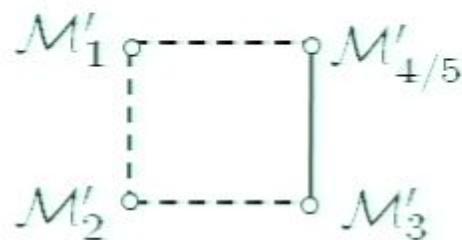
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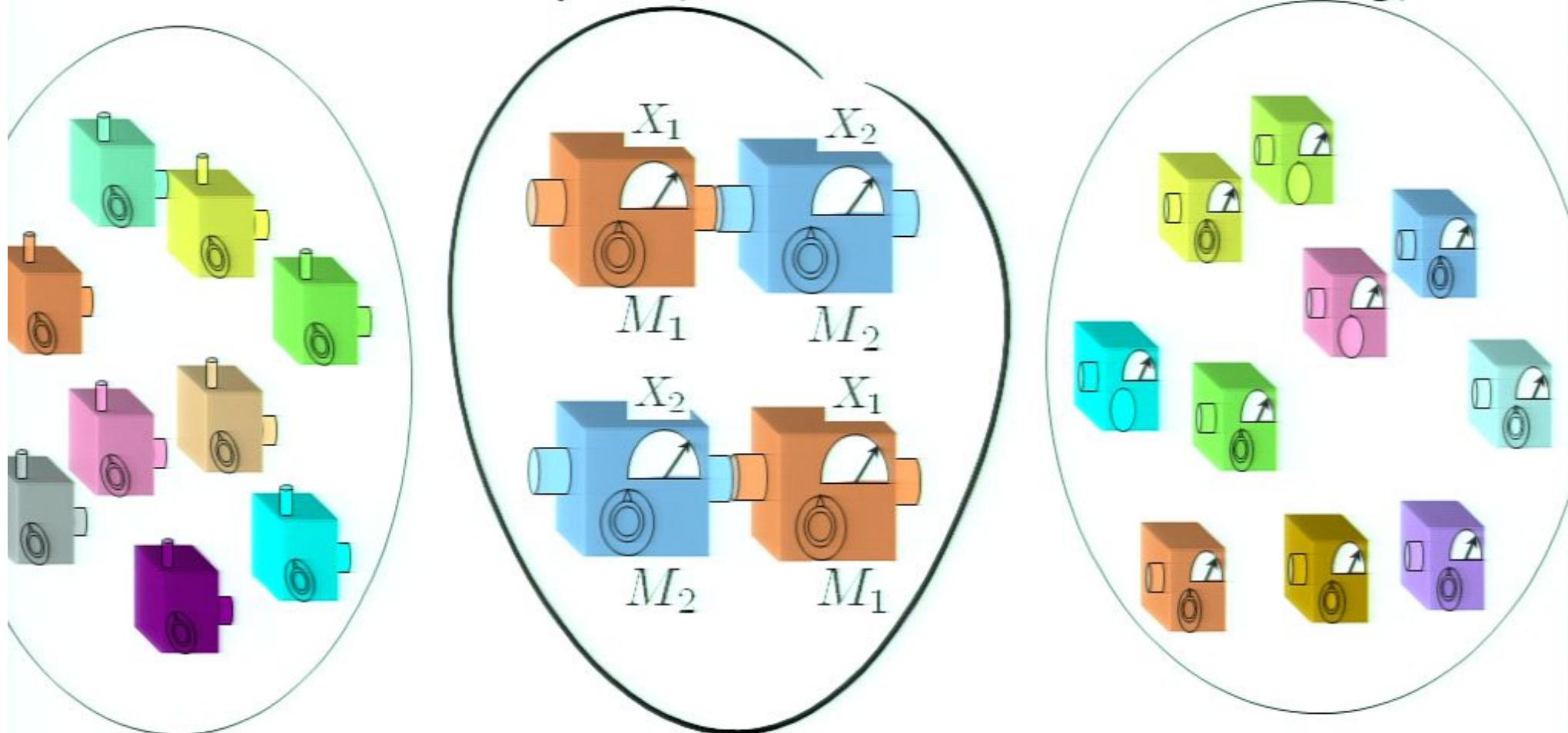


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Any principle that rules out the OS correlations also rules out all other extremal correlations for frustrated networks

Joint Measurability #2 (invariance under reordering)



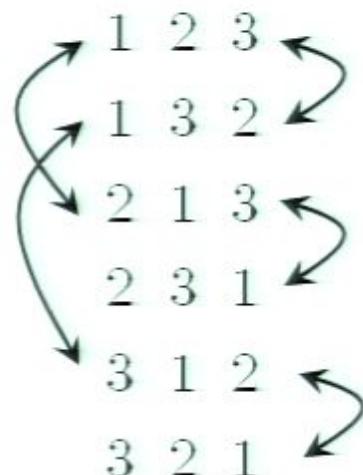
\mathcal{M}_1 and \mathcal{M}_2 are jointly measurable if $\exists M_1 \in \mathcal{M}_1$ and $\exists M_2 \in \mathcal{M}_2$ such that

$$\forall \mathcal{P} \forall \mathcal{M}_3 : p(X_1, X_2, X_3 | \mathcal{P} \text{ then } M_1 \text{ then } M_2 \text{ then } \mathcal{M}_3)$$

$$= p(X_1, X_2, X_3 | \mathcal{P} \text{ then } M_2 \text{ then } M_1 \text{ then } \mathcal{M}_3)$$

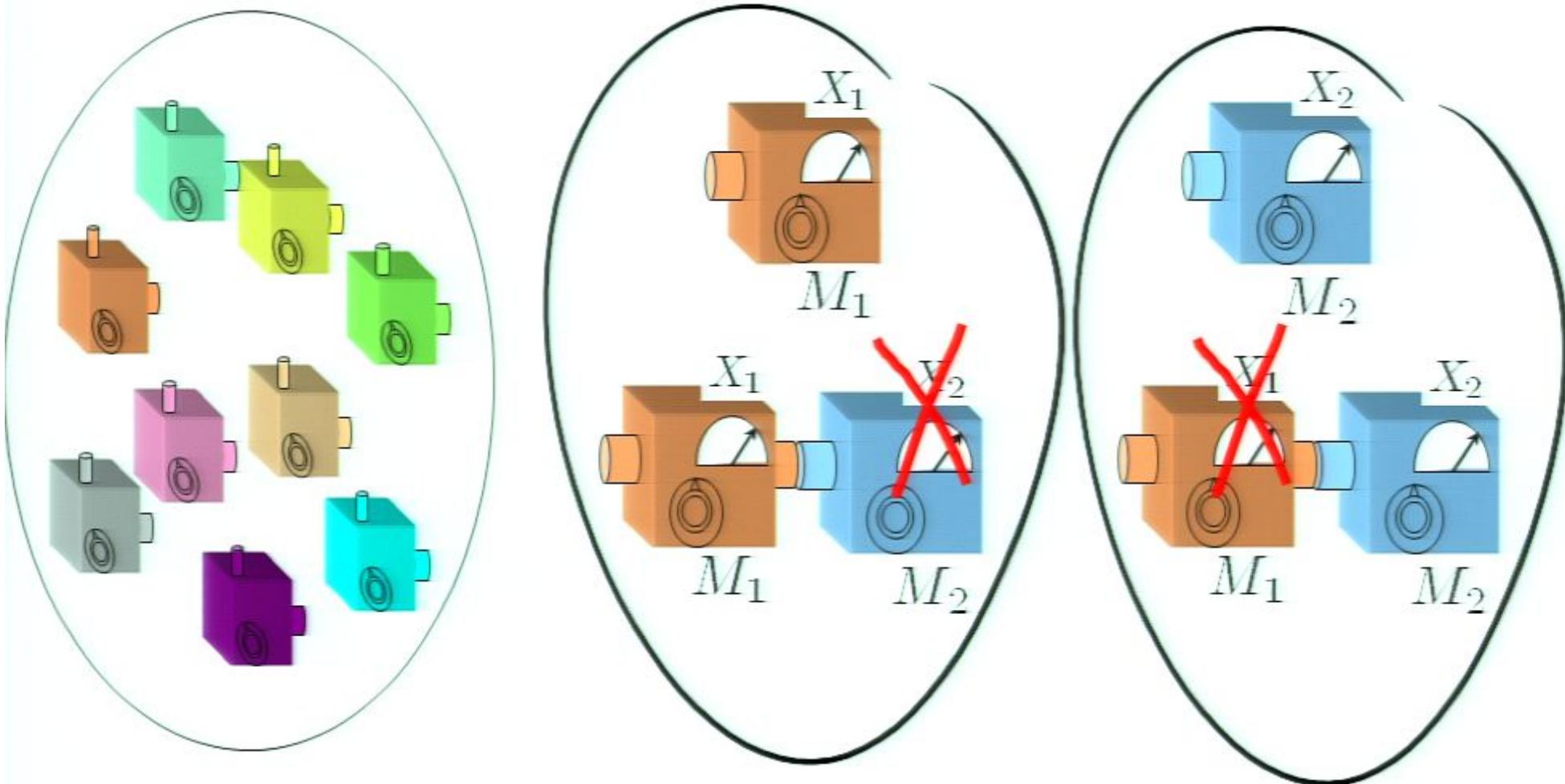
Lemma A: Pairwise JM #2 \Rightarrow Triplewise JM #2

Proof:

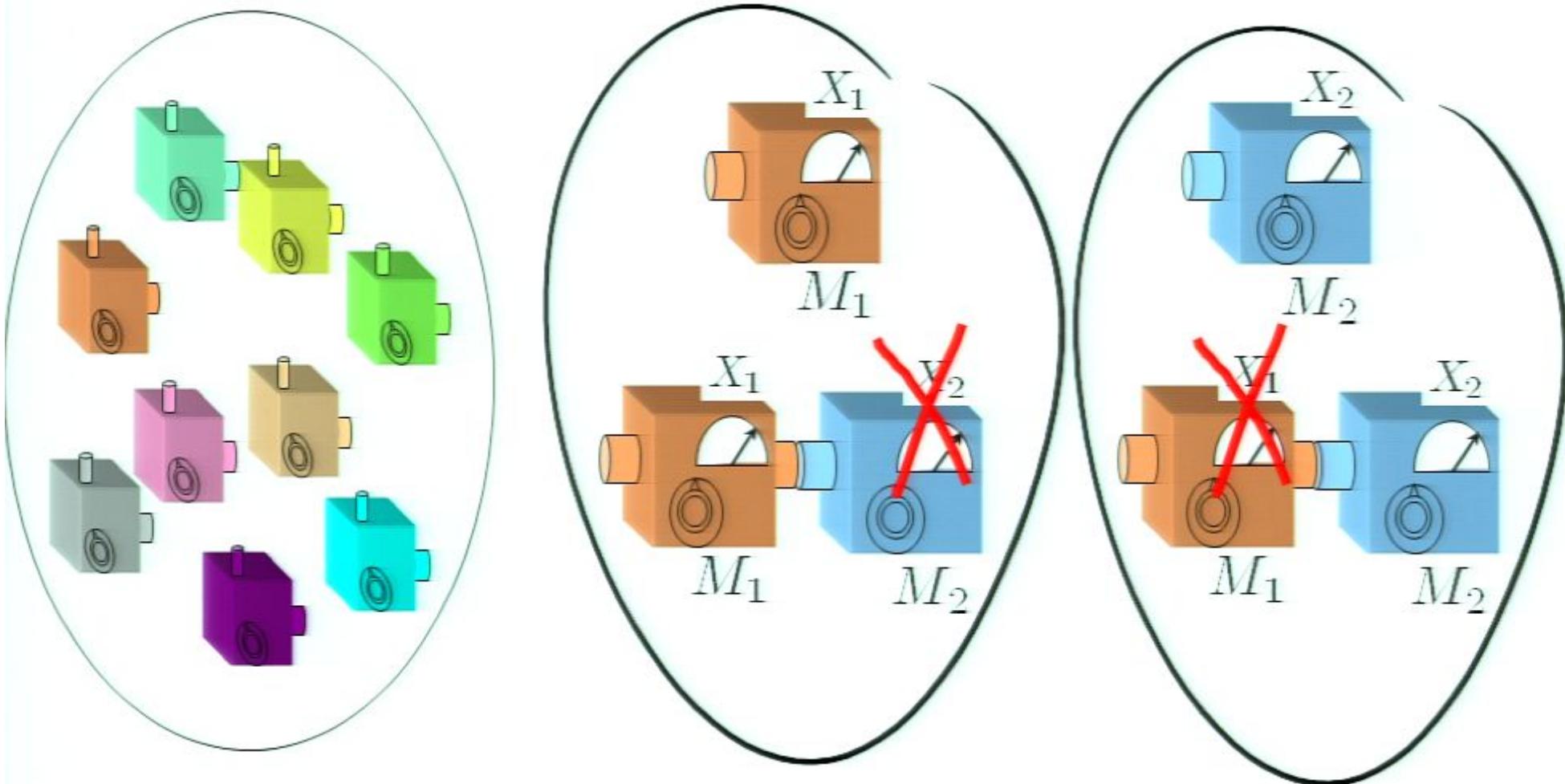


Lemma B: JM #2 (invariance under reordering)
⇒ JM #1 (joint simulability)

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Open question: Are there cases in which JM #1 ⇒ JM #2 ?

A possible argument against OS correlations:

Pairwise JM #1

↓ ???????

Pairwise JM #2

↓ lemma A

Triplewise JM #2

↓ lemma B

Triplewise JM #1

Research Directions

- In a general theory, for which measurements is it justified to assume deterministic outcomes in the ontological model?
- Can one find principles which disqualify the OS correlations?
- Counterfactual retrodiction (failure of joint measurability) vs. counterfactual prediction (standard uncertainty relations). Can we derive the precise amount of contextuality in quantum theory from time symmetry?

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Pairwise JM #2

↓ lemma A

Triplewise JM #2

↓ lemma B

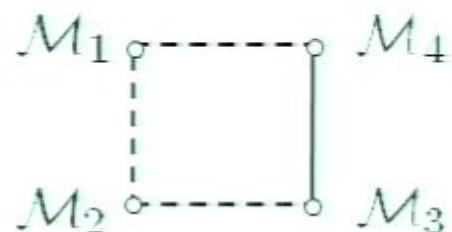
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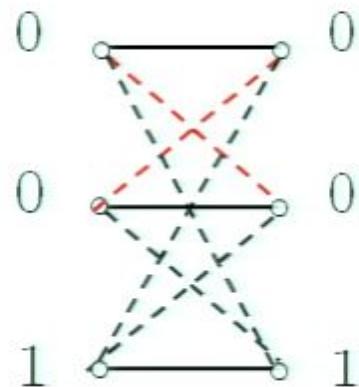
How the implications could be even more striking

Question: Can we find A_1, A_2, B_1 and B_2 and a state such that



Bell inequality

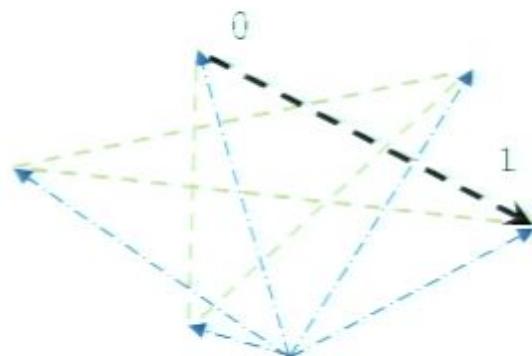
Recall: Bell locality + perfect correlation
⇒ deterministic noncontextual values



$$R \leq \frac{7}{9}$$

Local bound

A novel proof of the Kochen-Specker theorem based on the failure of transitivity of implication



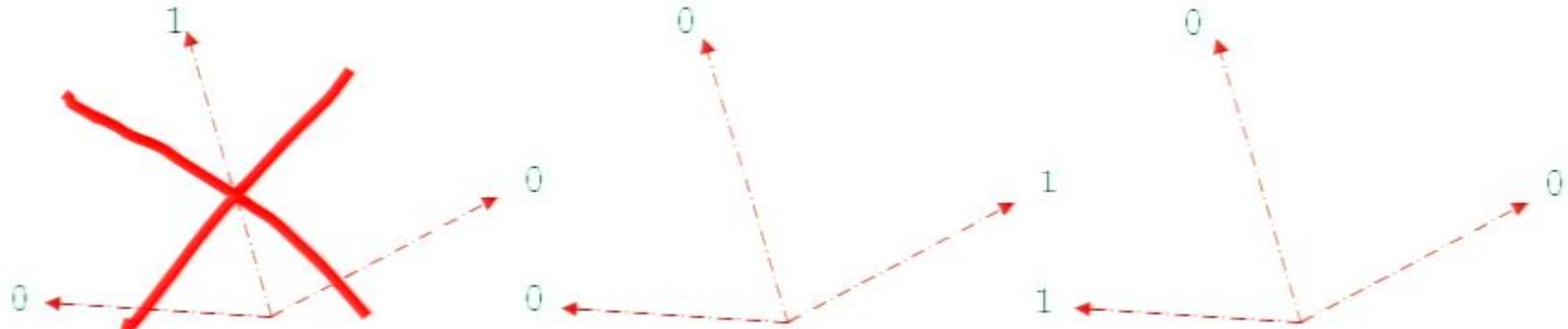
NONETHELESS

for all states $|\psi\rangle \in \text{span}(|l_1\rangle, |l_2\rangle)$

If $|l_1\rangle + |l_2\rangle$

Then $v(|l_1\rangle\langle l_1|) = 0 \implies v(|l_2\rangle\langle l_2|) = 1$

Because:



How the seer can achieve his ends in a quantum world

	Kochen-Specker bound	Quantum bound
	$R \leq \frac{n-1}{n}$ $= 1 - \frac{1}{n}$	$R \leq \frac{2 \cos(\frac{\pi}{n})}{1 + \cos(\frac{\pi}{n})}$ $\simeq 1 - \frac{\pi^2}{4n^2}$

The seer's challenge to the suitor:
identify a correlated pair of boxes

Suitor's expected prob. of winning: $\frac{1}{n}$

Suitor's actual prob. of winning: $\simeq O(\frac{1}{n^2})$

Suppose $n \ll \text{no. of suitors} \ll n^2$

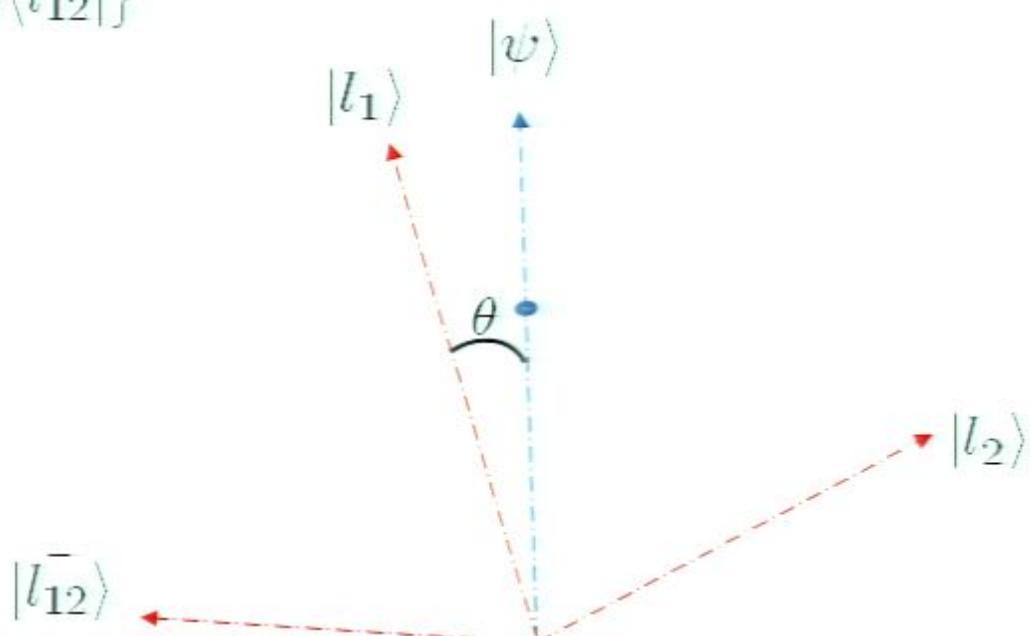
Suitors believe it is very likely that **one of them will win**

In fact it is very likely that **none of them will win**

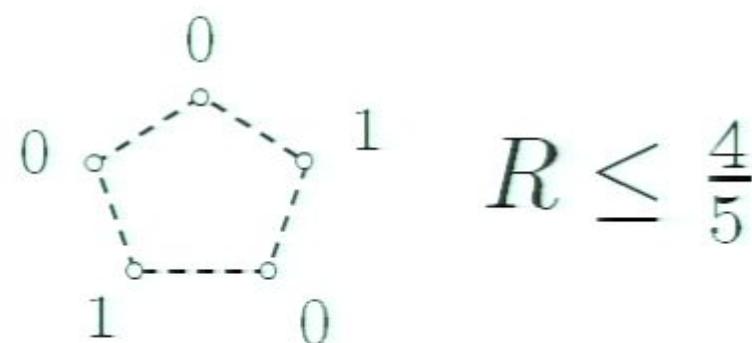
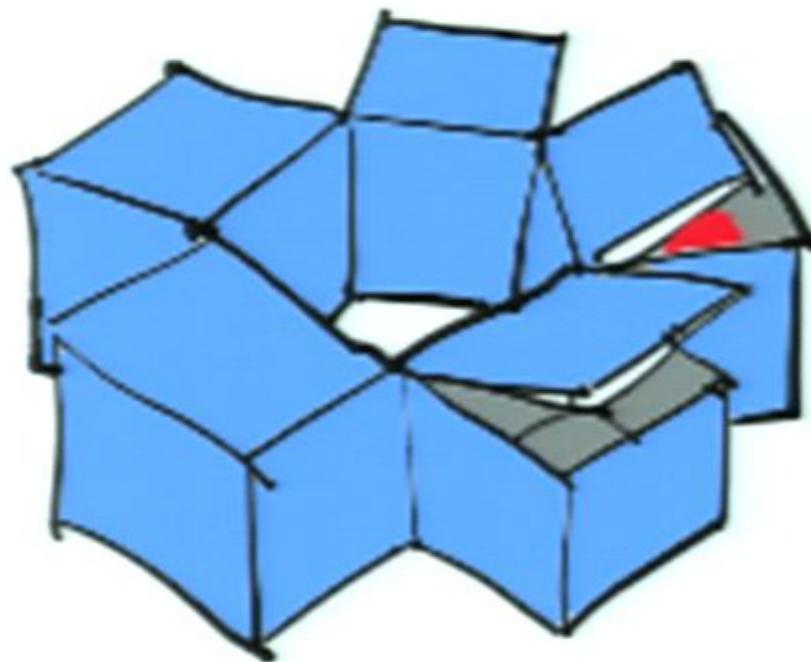
Klyachko's proof of the Kochen-Specker theorem

$$\cos^2 \theta = \frac{1}{\sqrt{5}}$$

Consider: $\{|l_1\rangle\langle l_1|, |l_2\rangle\langle l_2|, |\bar{l_{12}}\rangle\langle \bar{l_{12}}|\}$



Double-query 5-box system allowing only **adjacent** queries



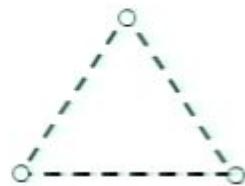
Frustrated Networks

Nodes are binary variables

Edges imply joint measurability

- Perfect correlation
- Perfect anti-correlation

E.g. Correlations in
the parable of the
overprotective seer:



Frustration = no valuation satisfying all correlations

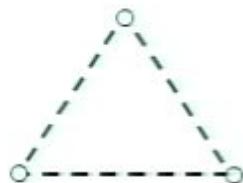
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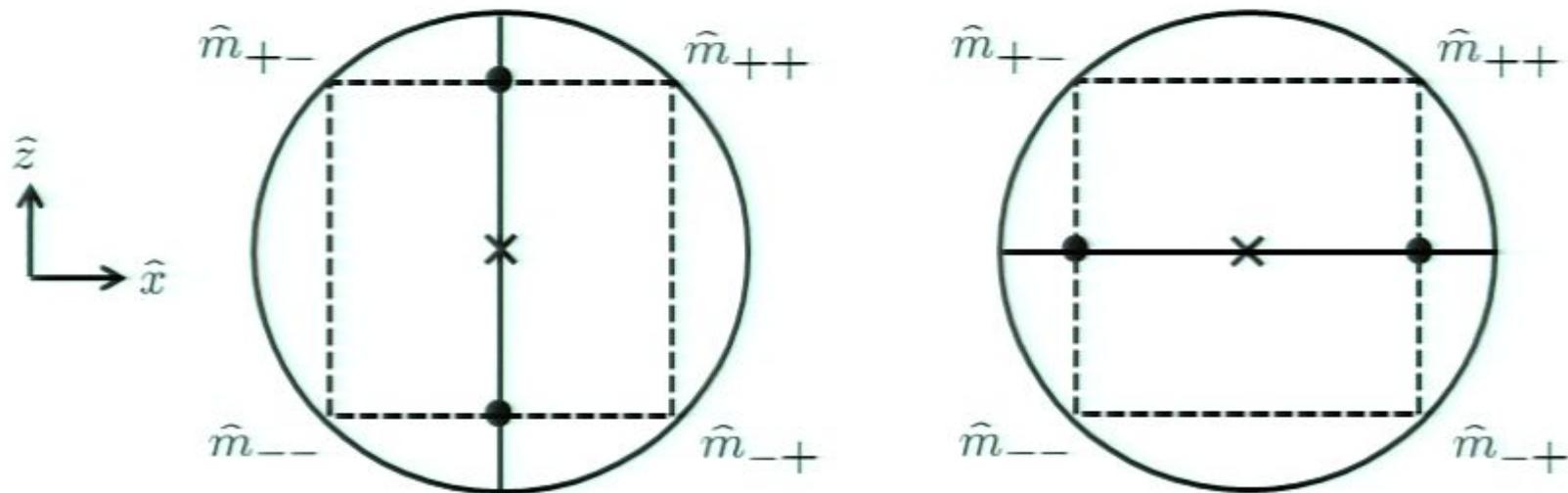


Joint measurability of POVMs

Noisy z-spin $\{E_+^z, E_-^z\}$ $E_\pm^z = \eta |\pm z\rangle\langle \pm z| + (1 - \eta) \frac{I}{2}$

Noisy x-spin $\{E_+^x, E_-^x\}$ η = sharpness factor

Noisy y-spin $\{E_+^y, E_-^y\}$



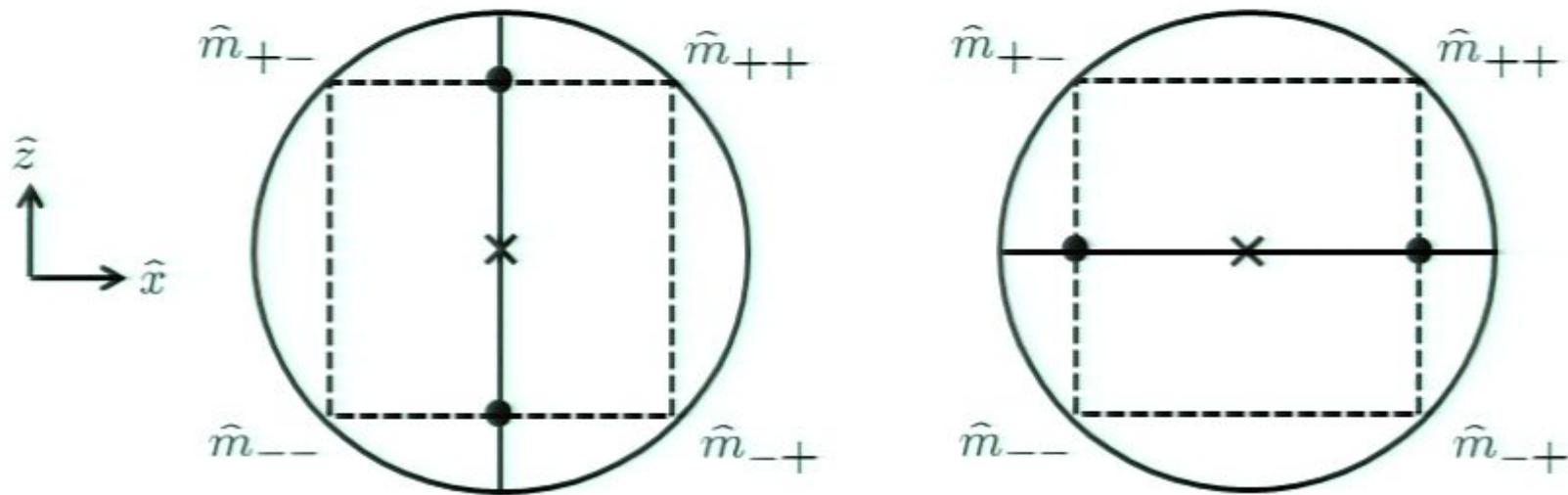
Pairwise JM iff $\eta \leq \frac{1}{\sqrt{2}}$

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