

Title: A Matter Bounce By Means of Ghost Condensation

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Abstract: TBA

Ghost Bounce

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A matter bounce by means of ghost condensation

R. Brandenberger, L. Levasseur and C. Lin

arXiv:1007.2654

Chunshan Lin

McGill

USTC

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Outline

- I. Alternative inflation models
 - › Necessity
 - › Matter bounce
- II. Ghost condensation
 - › Basic philosophy
 - › Applications
 - › Interesting features
 - › Instability
- III. Matter bounce by means of ghost condensation
 - › Several advantages:
ghost free, stable against radiation and anisotropic stress...
 - › Perturbation
 - › Cut off issue

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Part I

Alternative inflation models

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- ◆ Inflation suffers from some conceptual problems

Flatness problem

Amplitude problem $\frac{V(\varphi)}{\Delta\varphi^4} \leq 10^{-12}$

Trans-Planckian problem

Singularity problem

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- ◆ Some other attempts

Matter bounce, Ekpyrotic,

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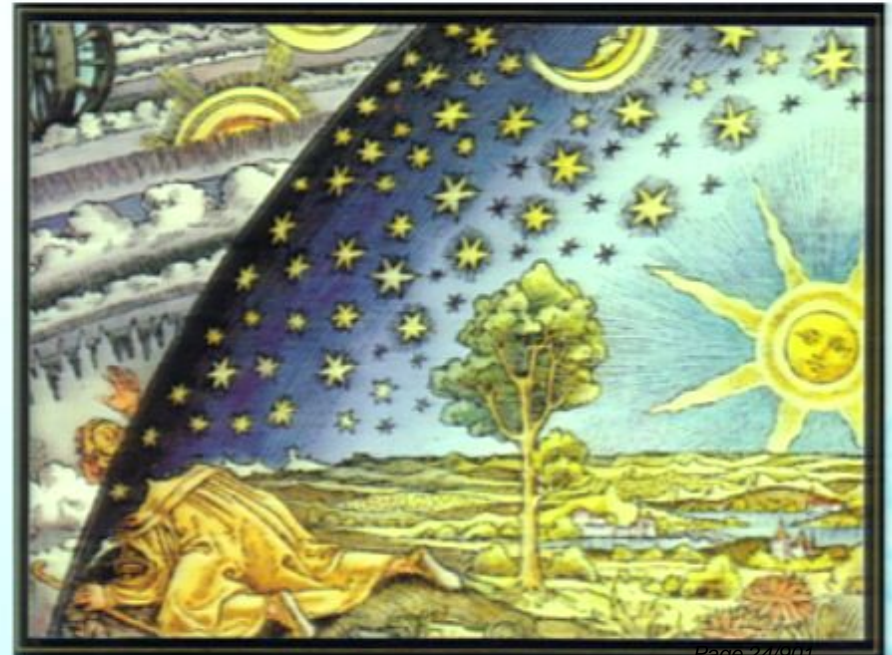
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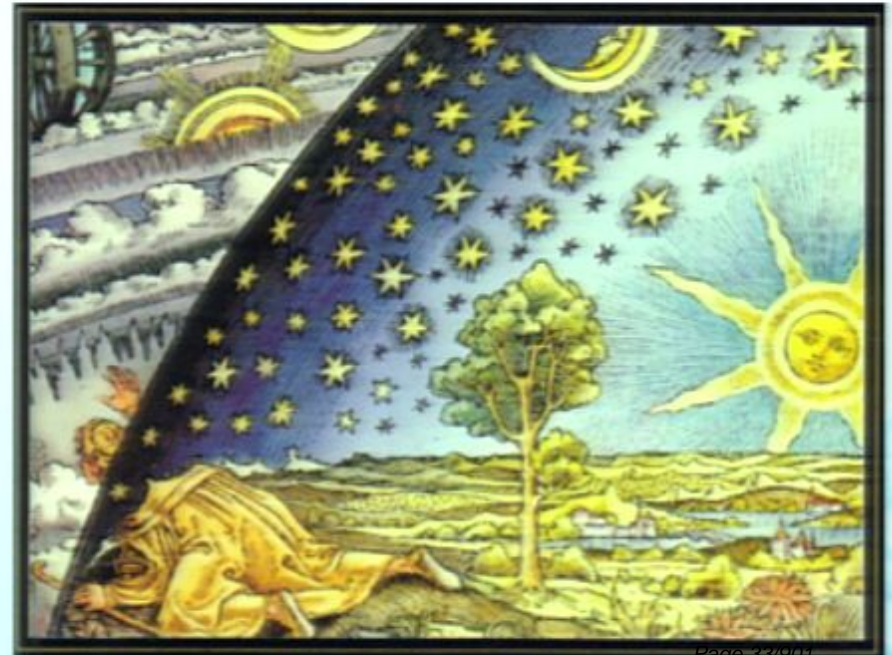
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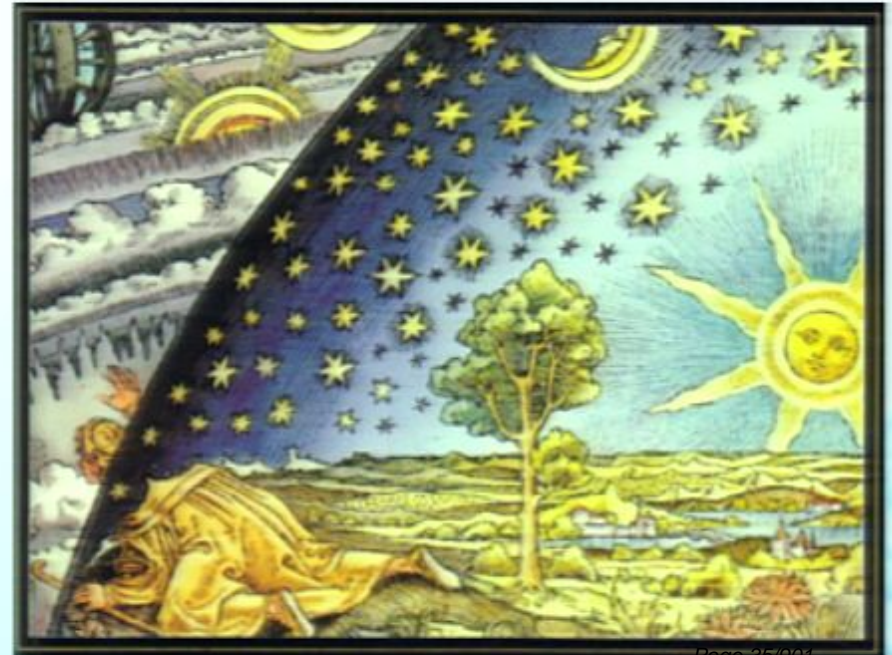
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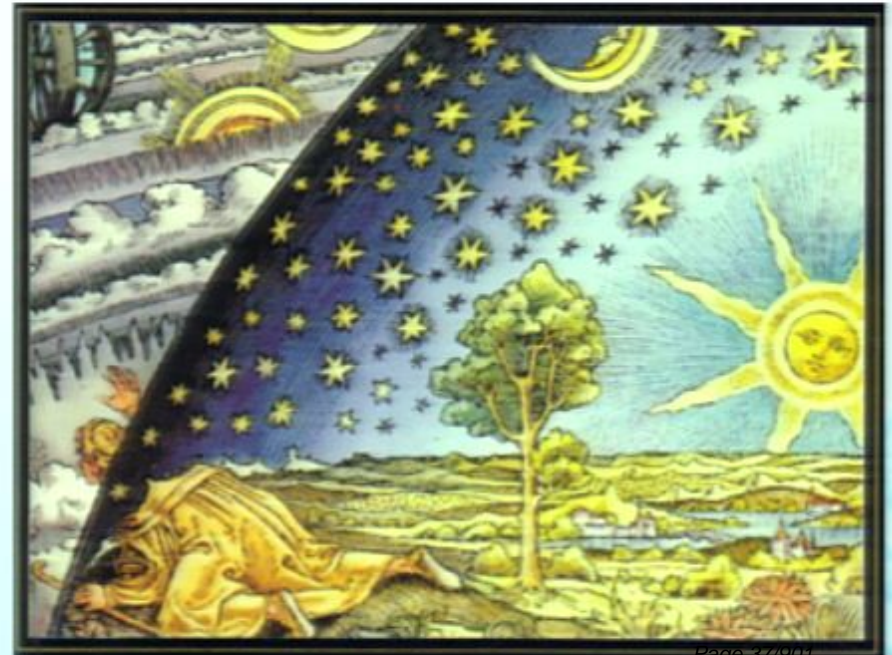
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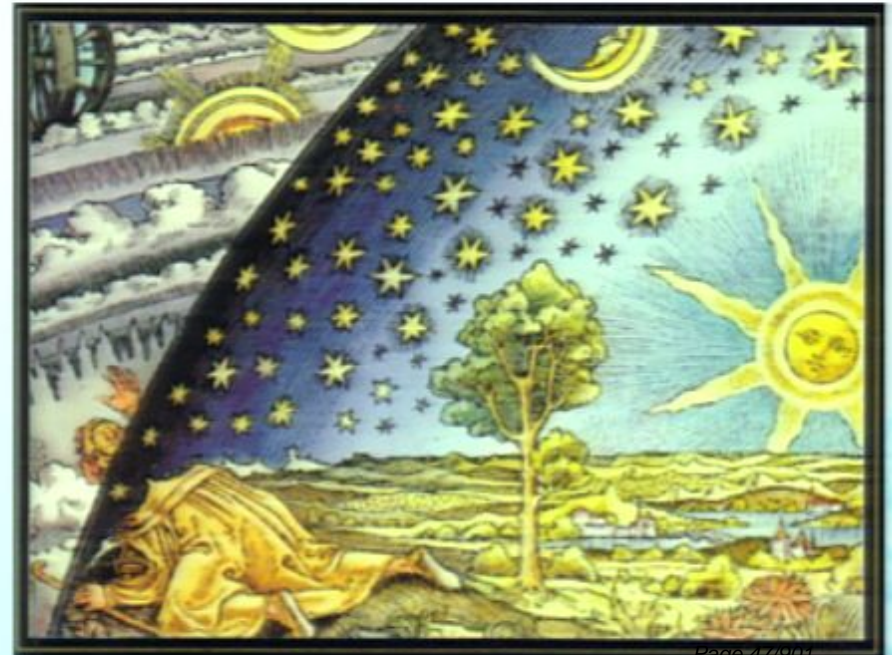
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- ◆ Cold pressureless matter
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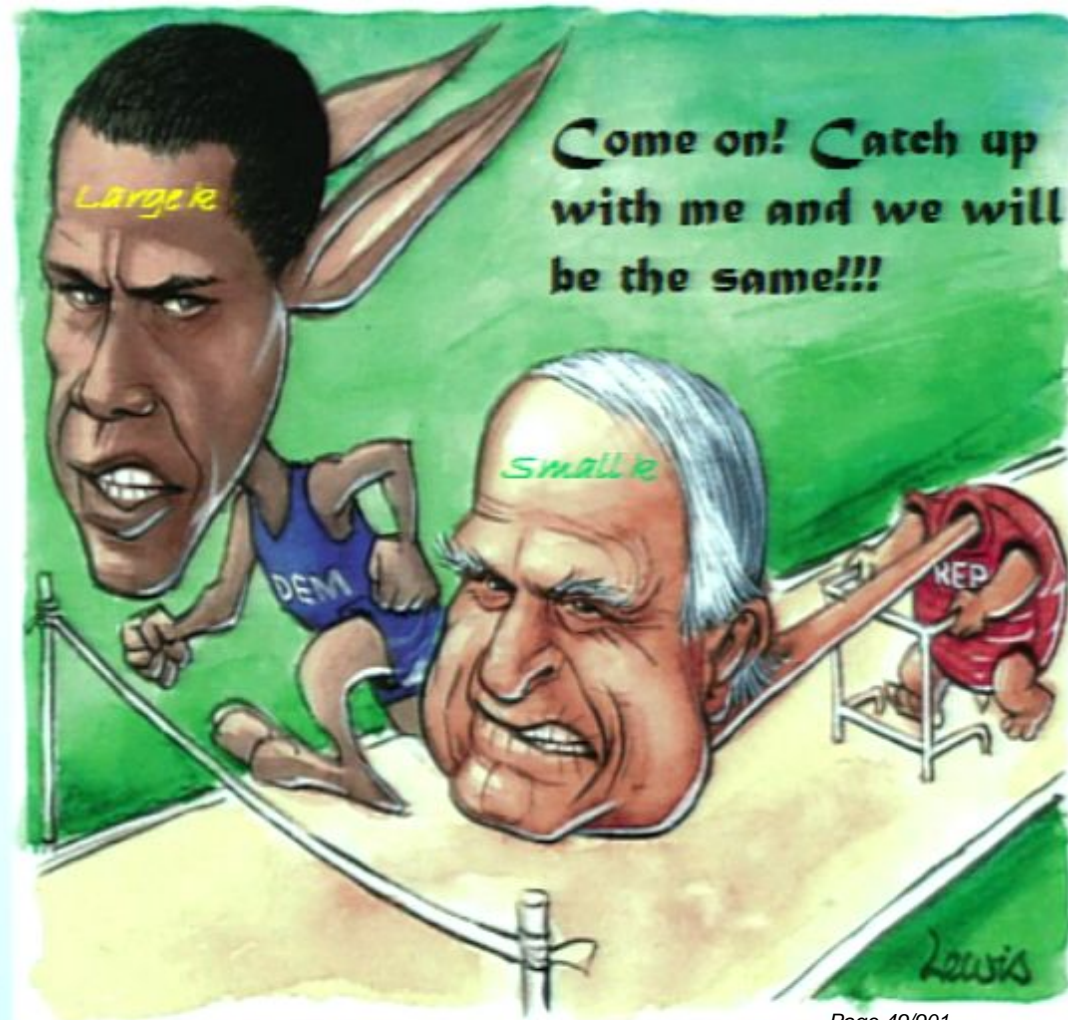
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$$\delta\varphi_* \propto H_* \propto t_*^{-1}$$

super horizon growing

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Amplitude of the larger scale perturbation mode will catch up with the smaller scale perturbation mode.



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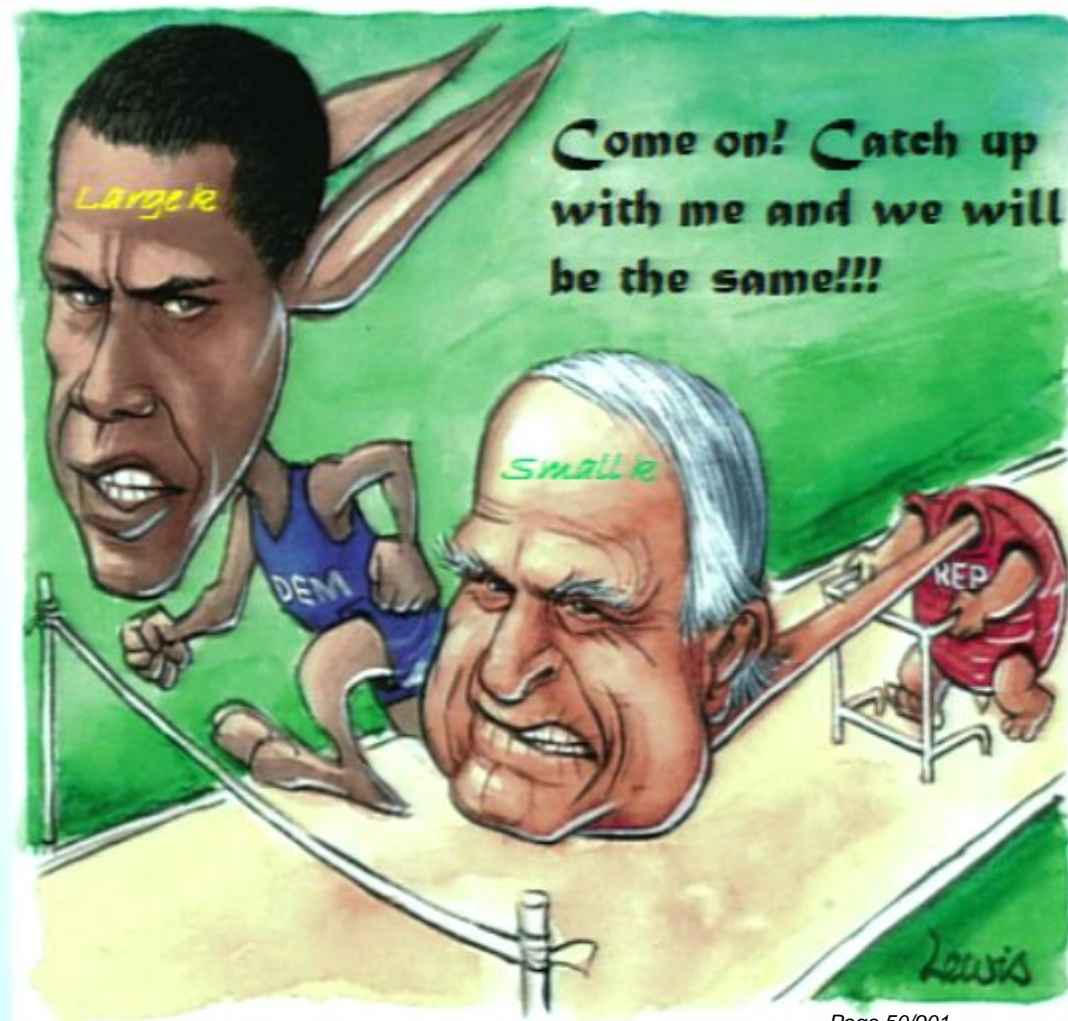
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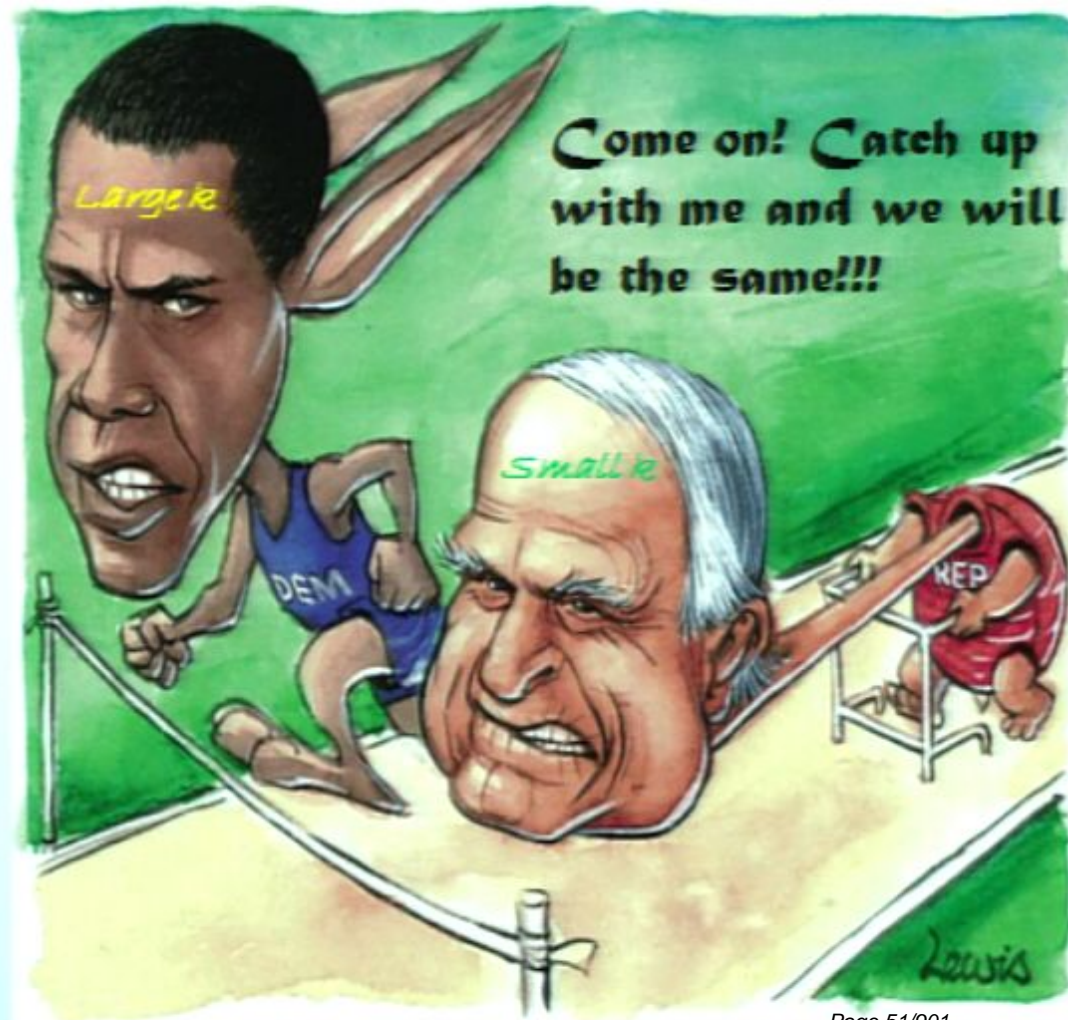
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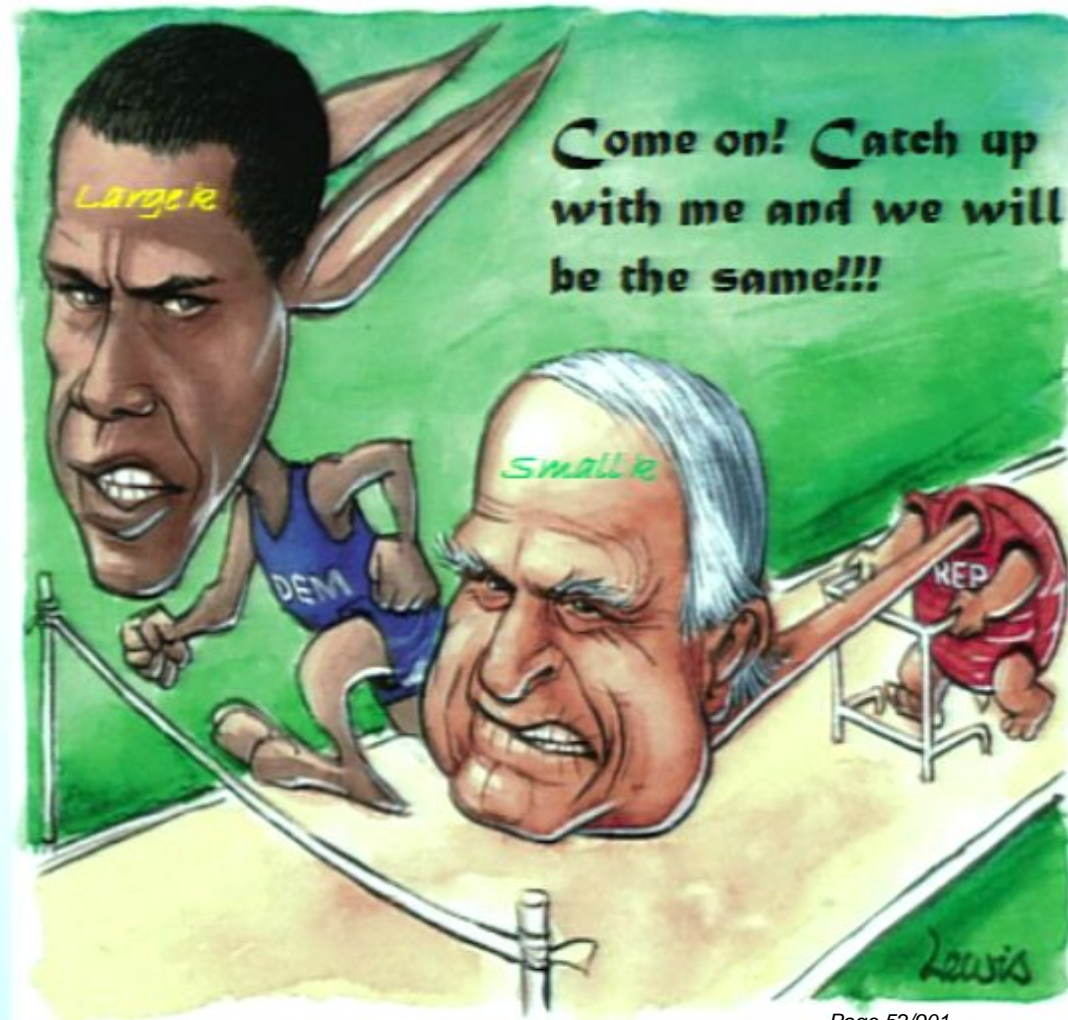
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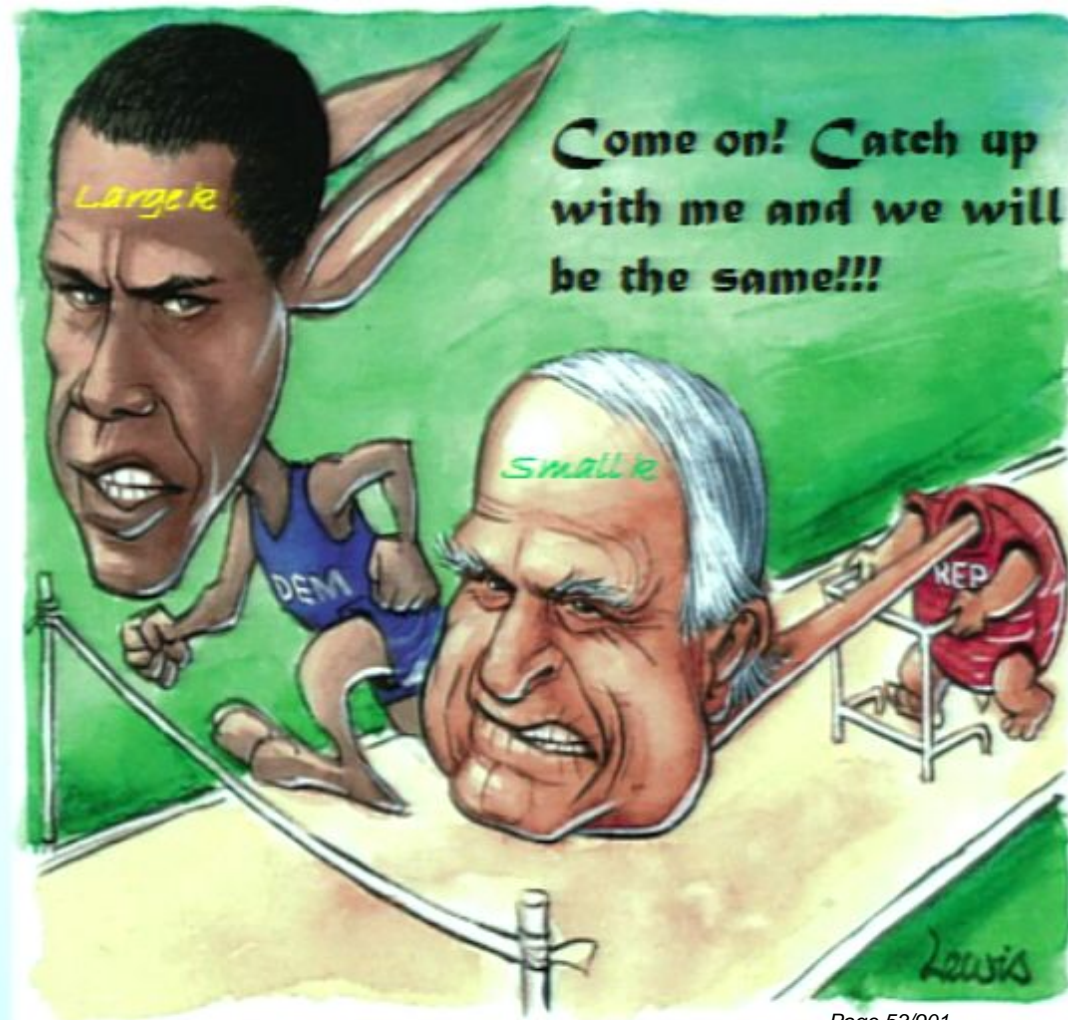
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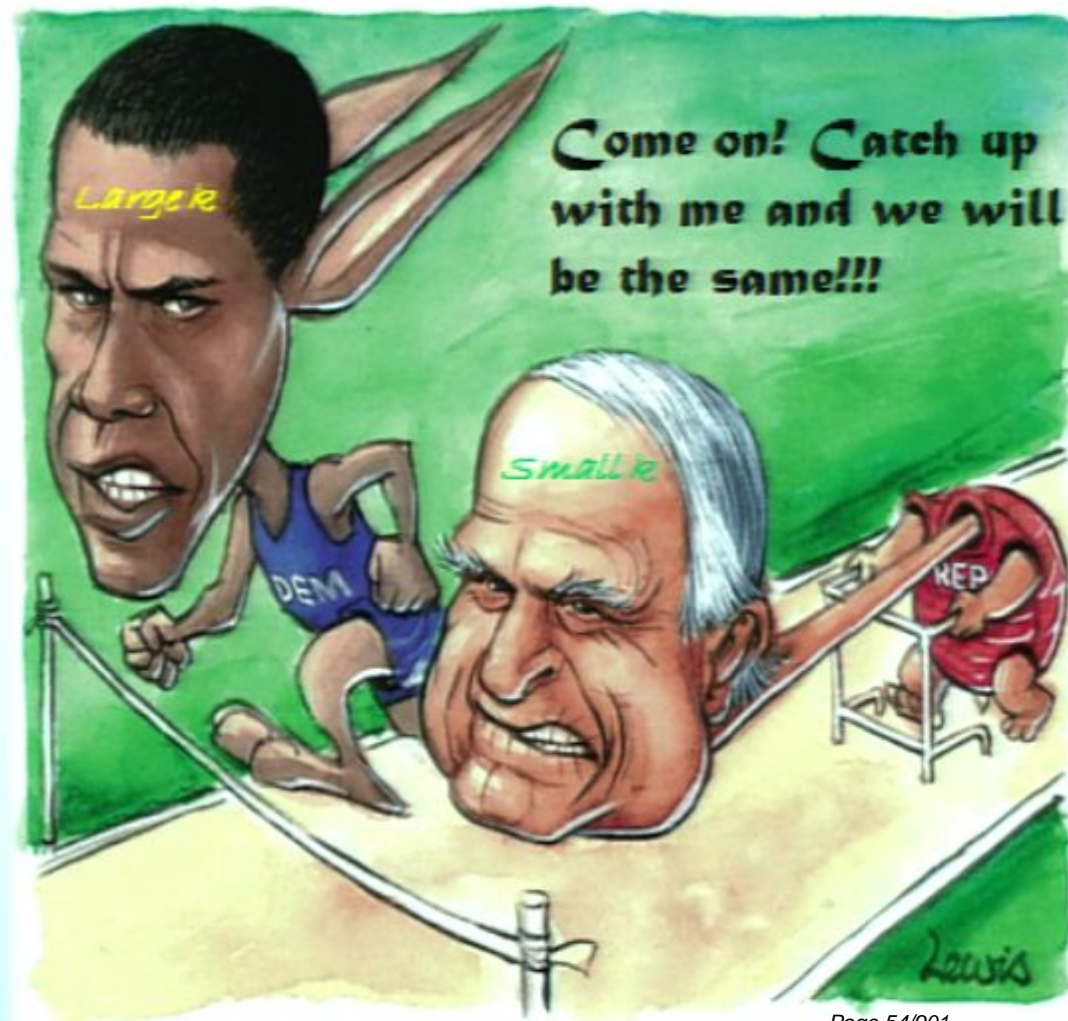
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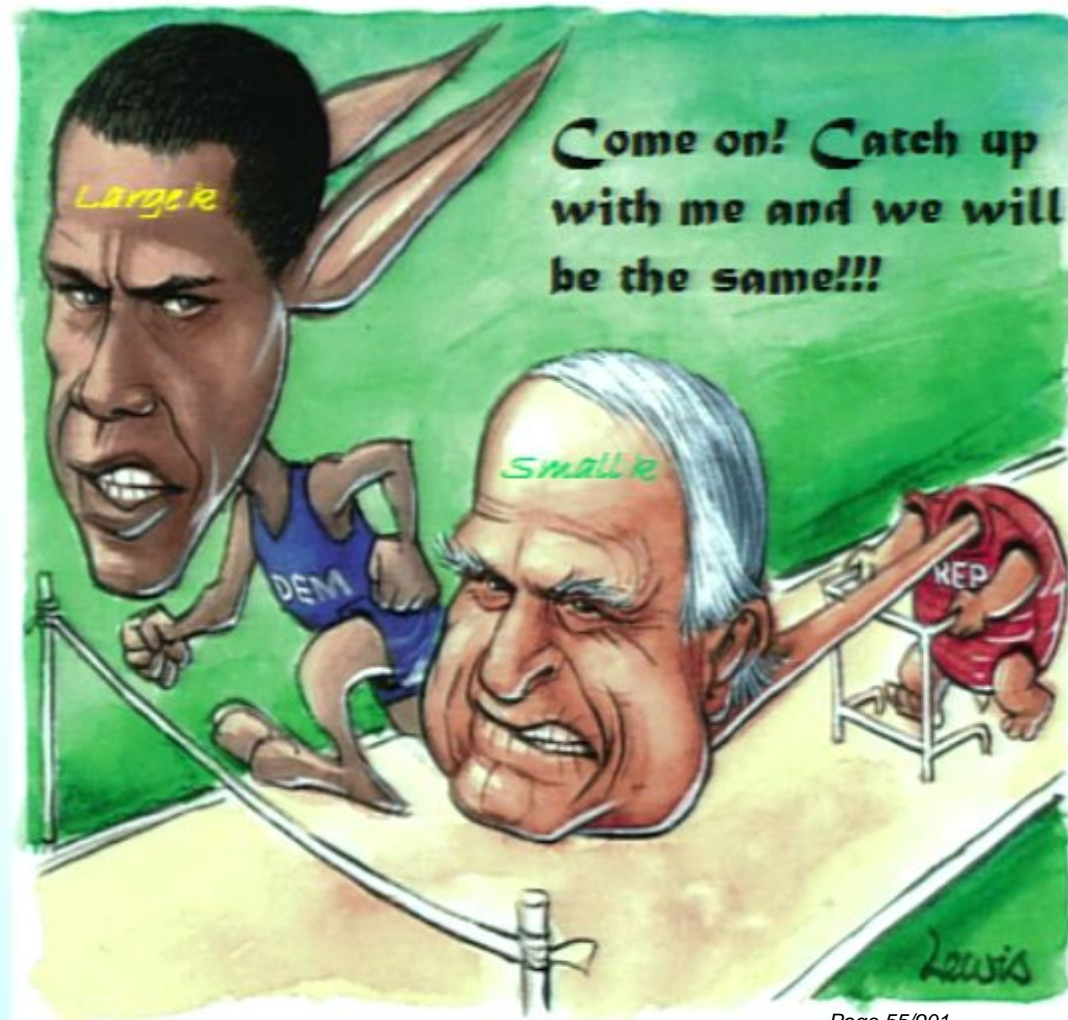
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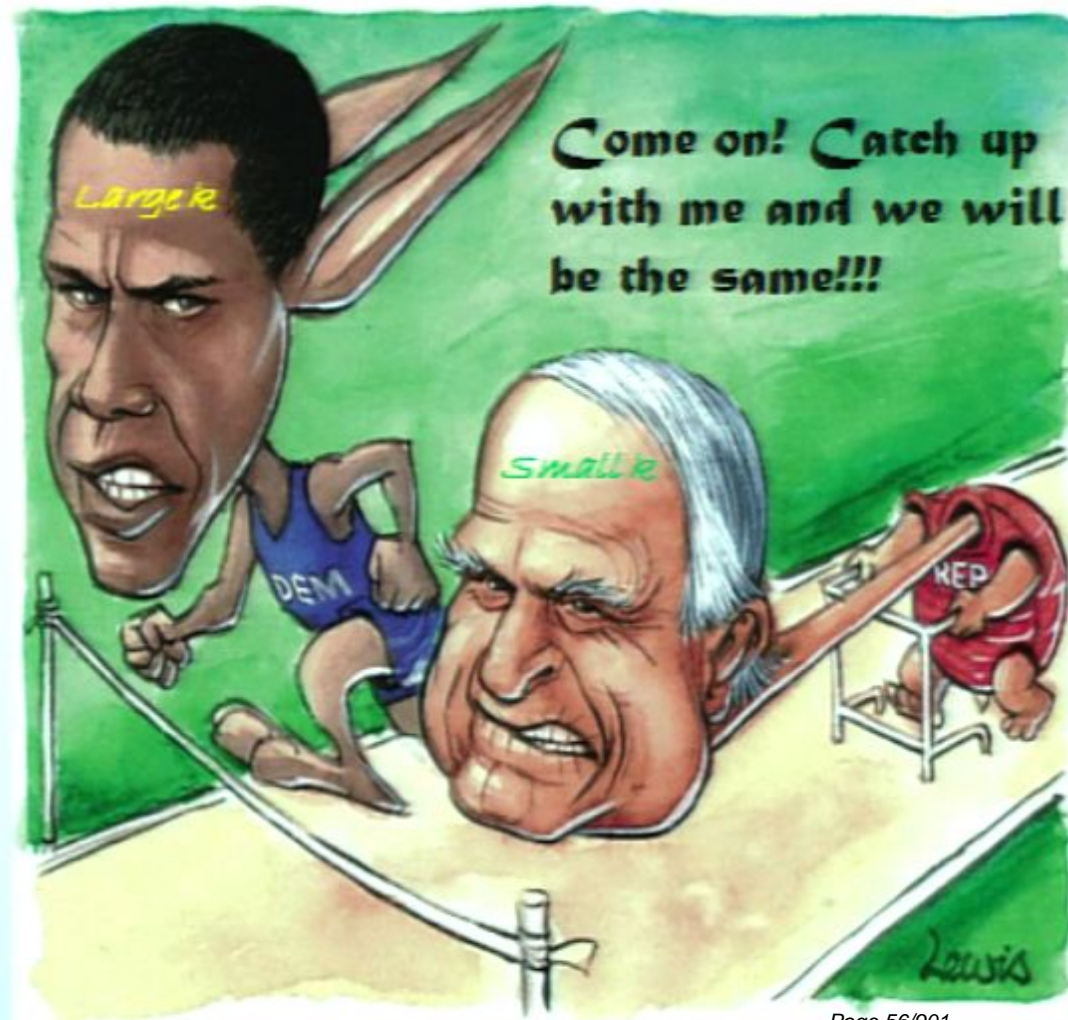
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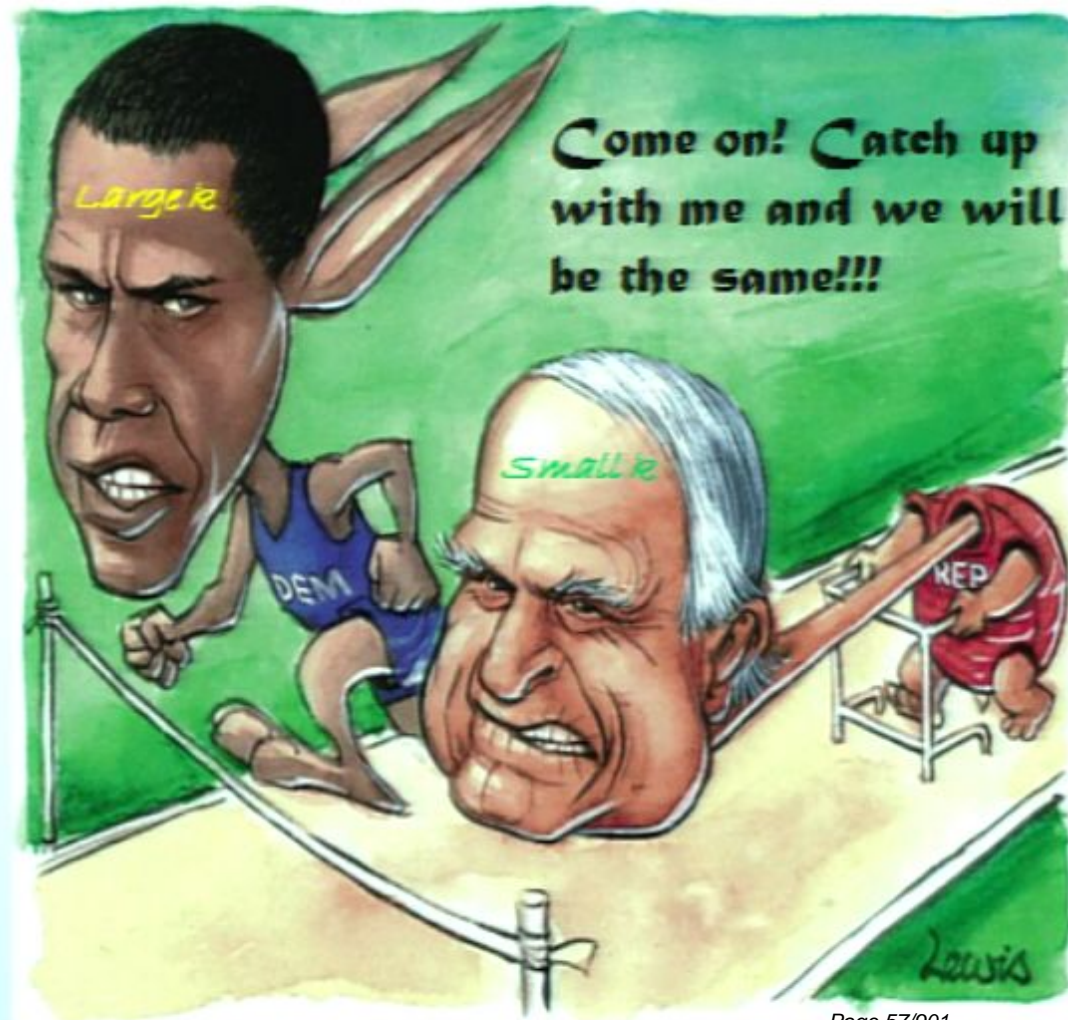
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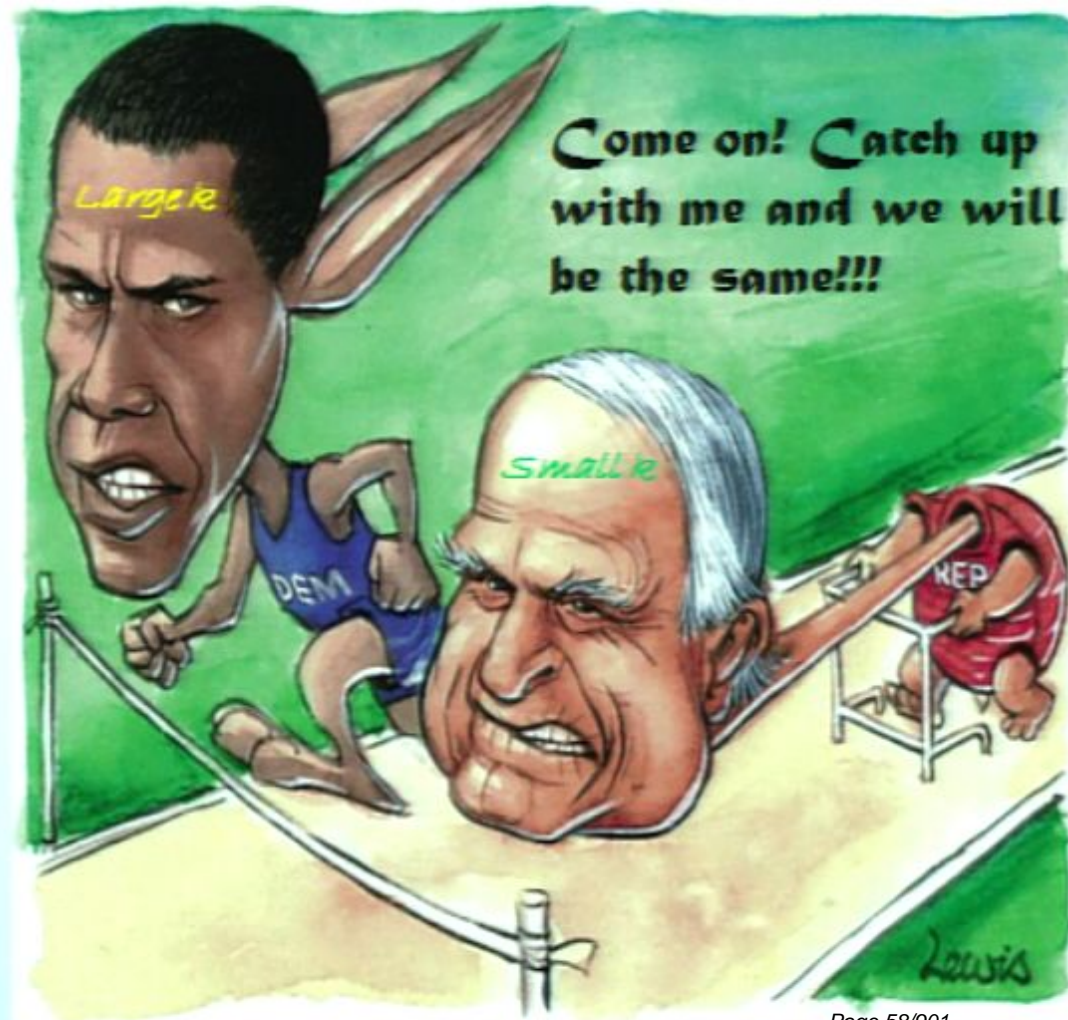
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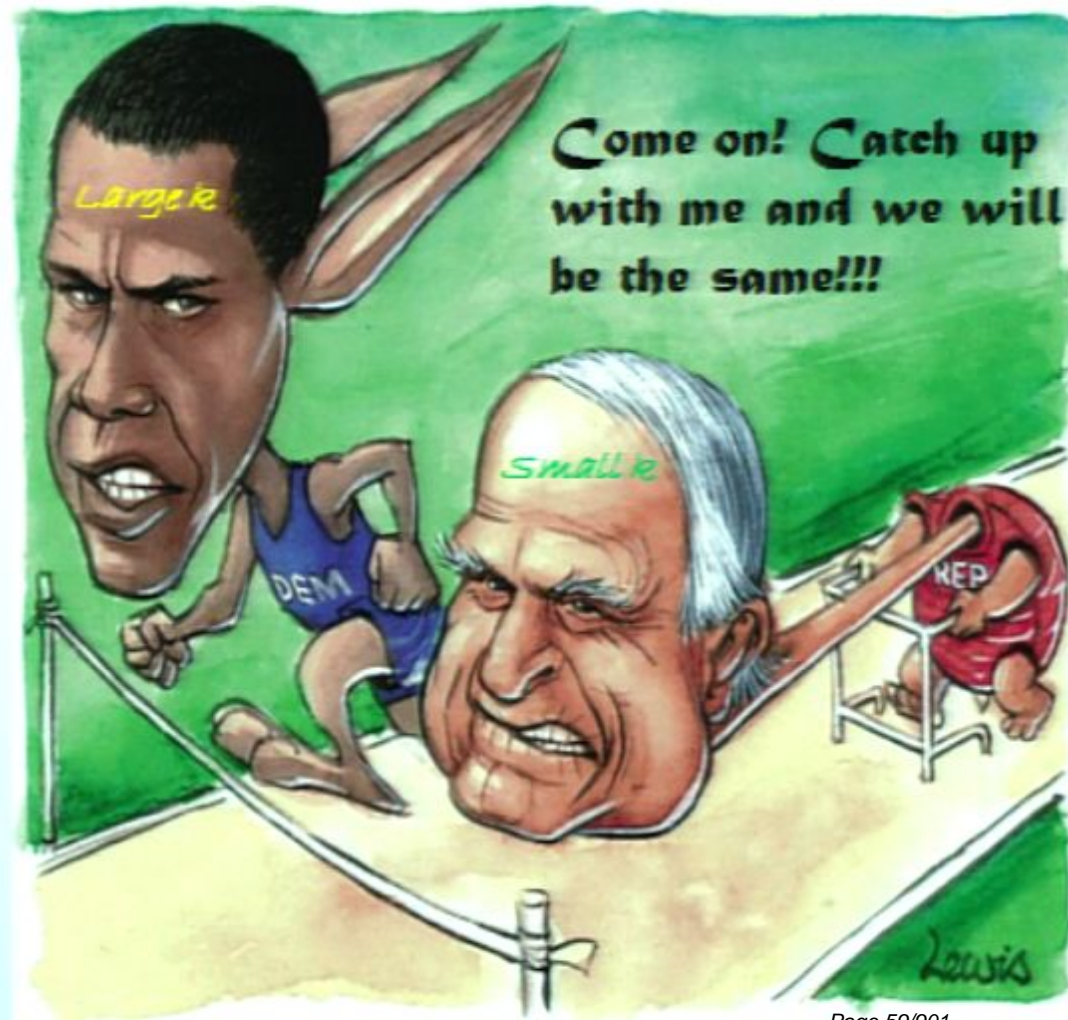
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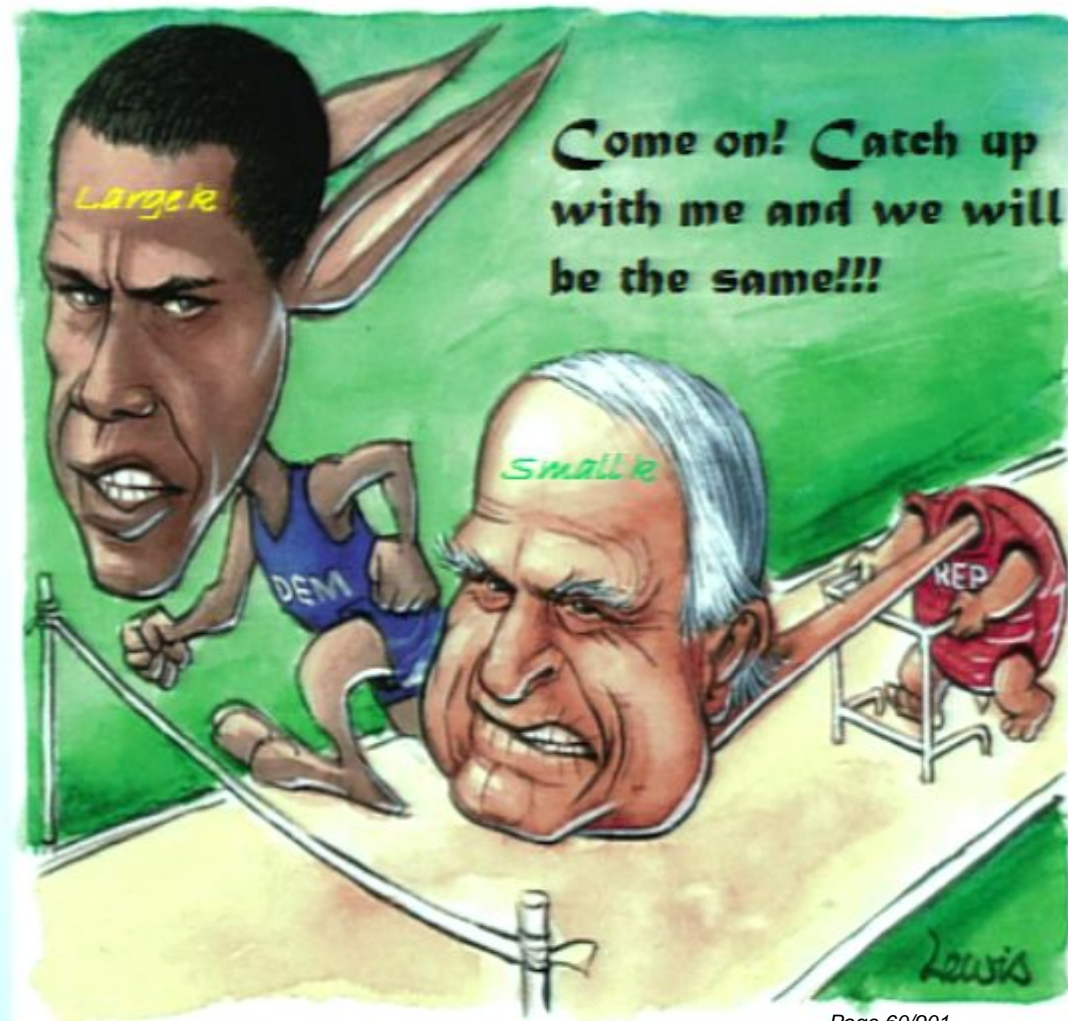
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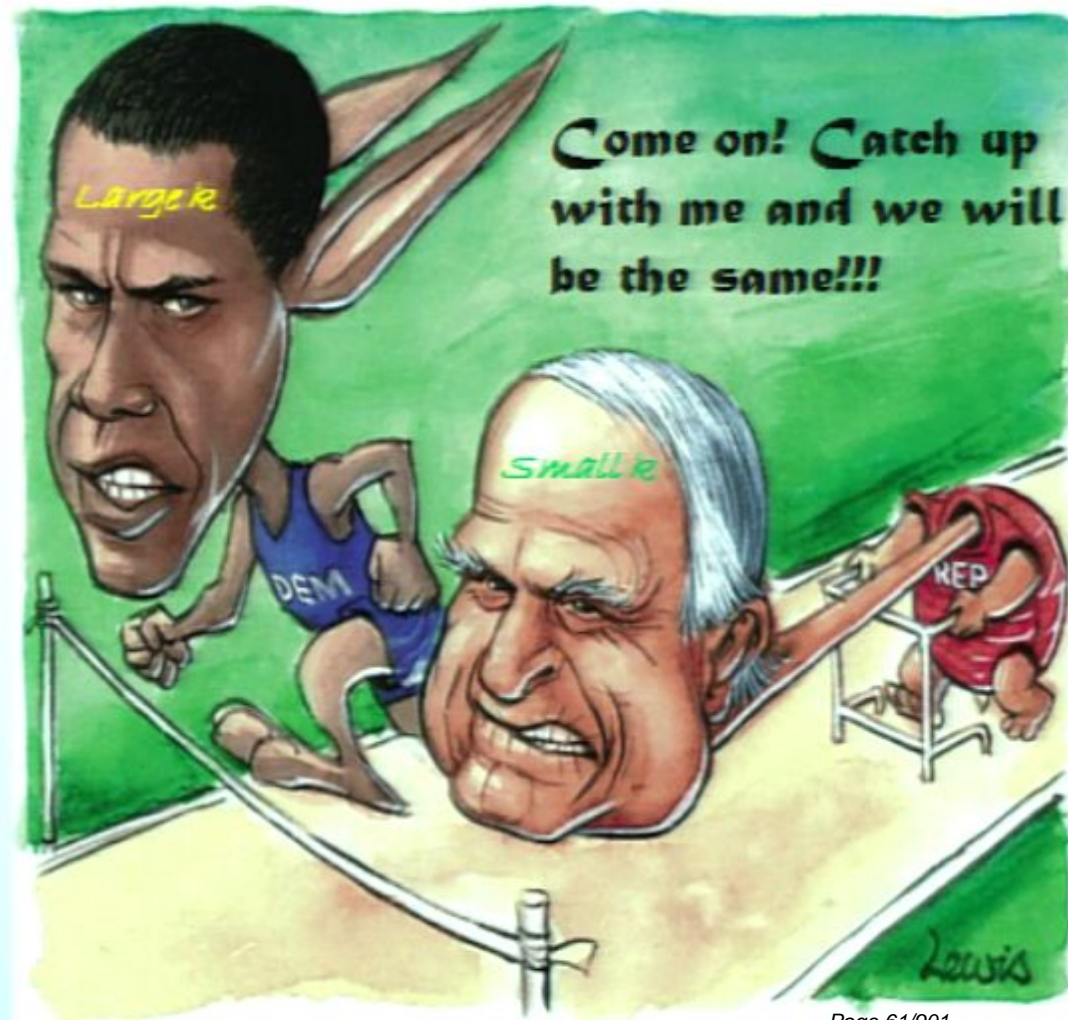
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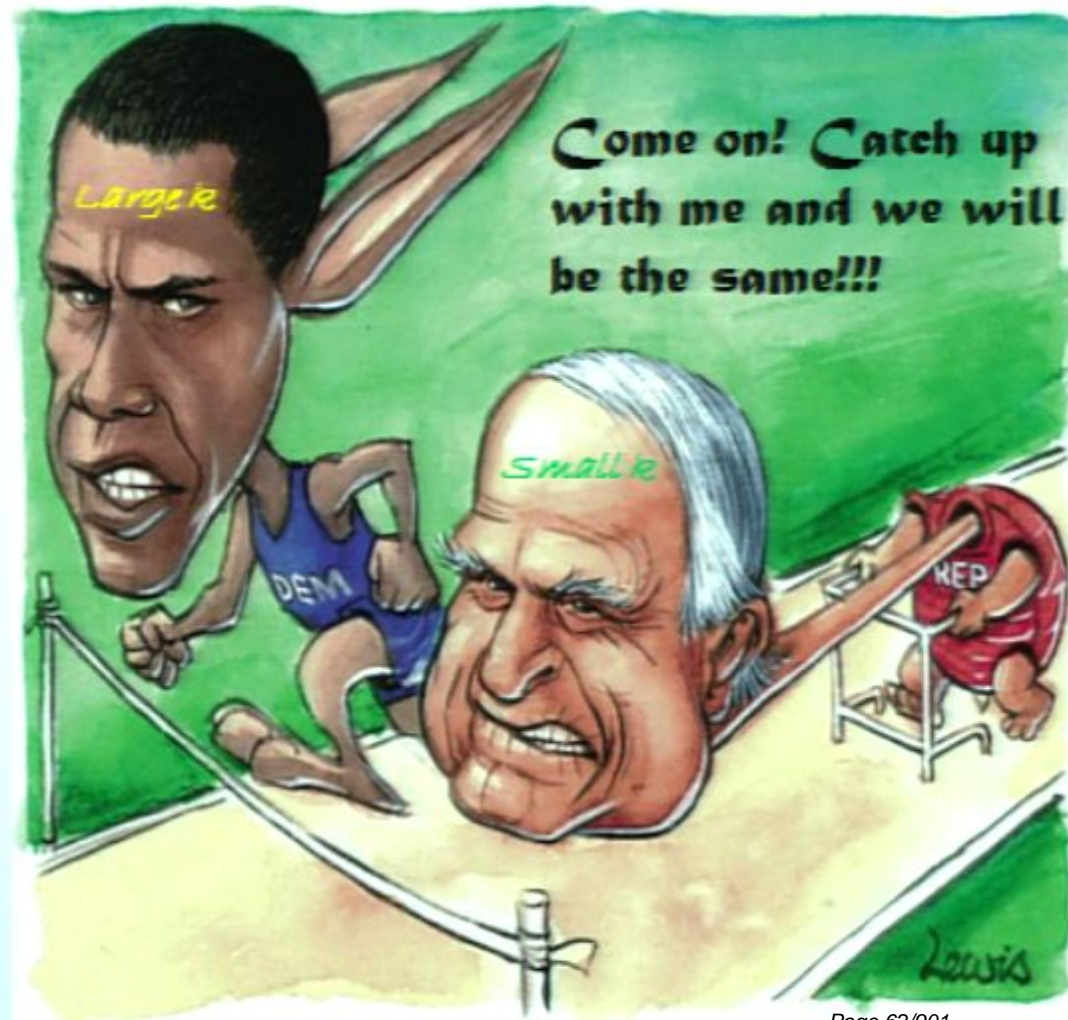
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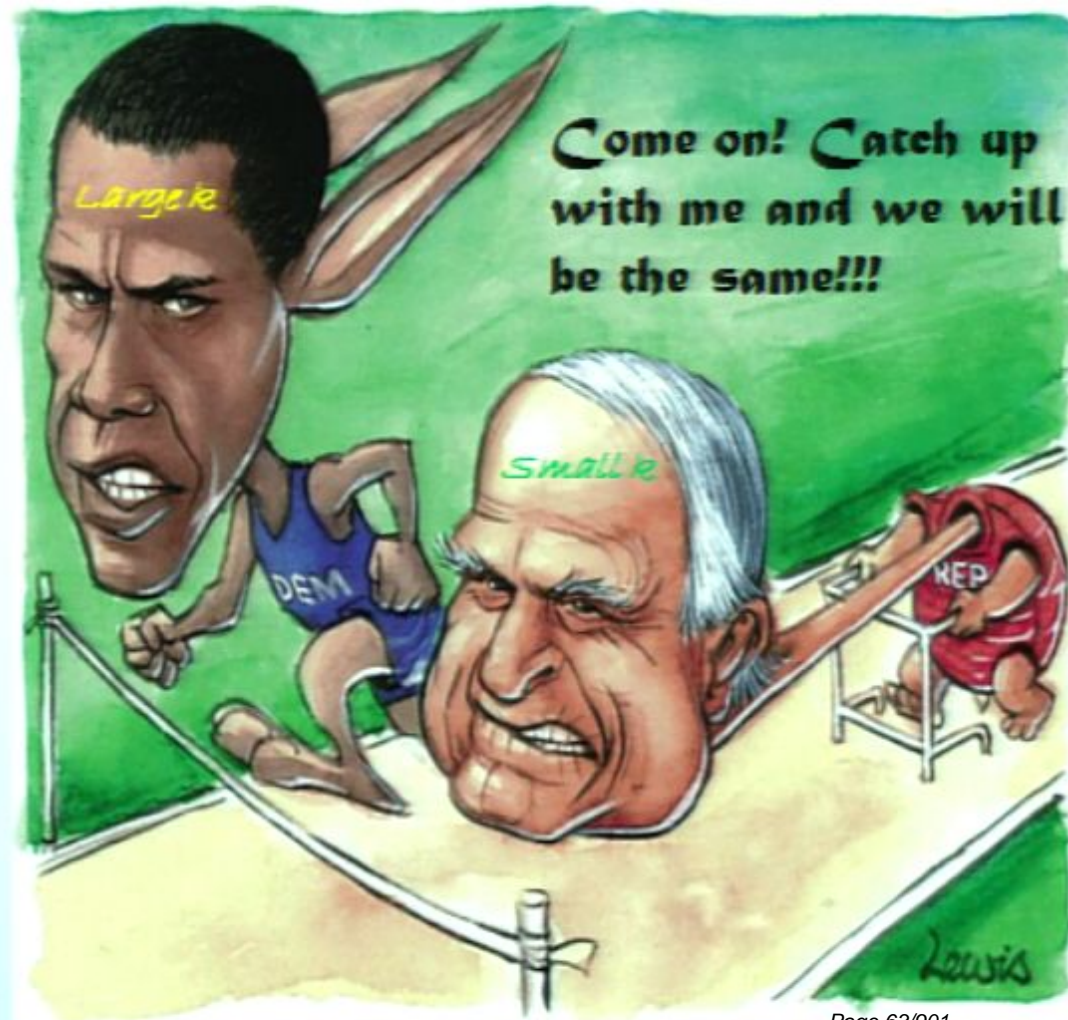
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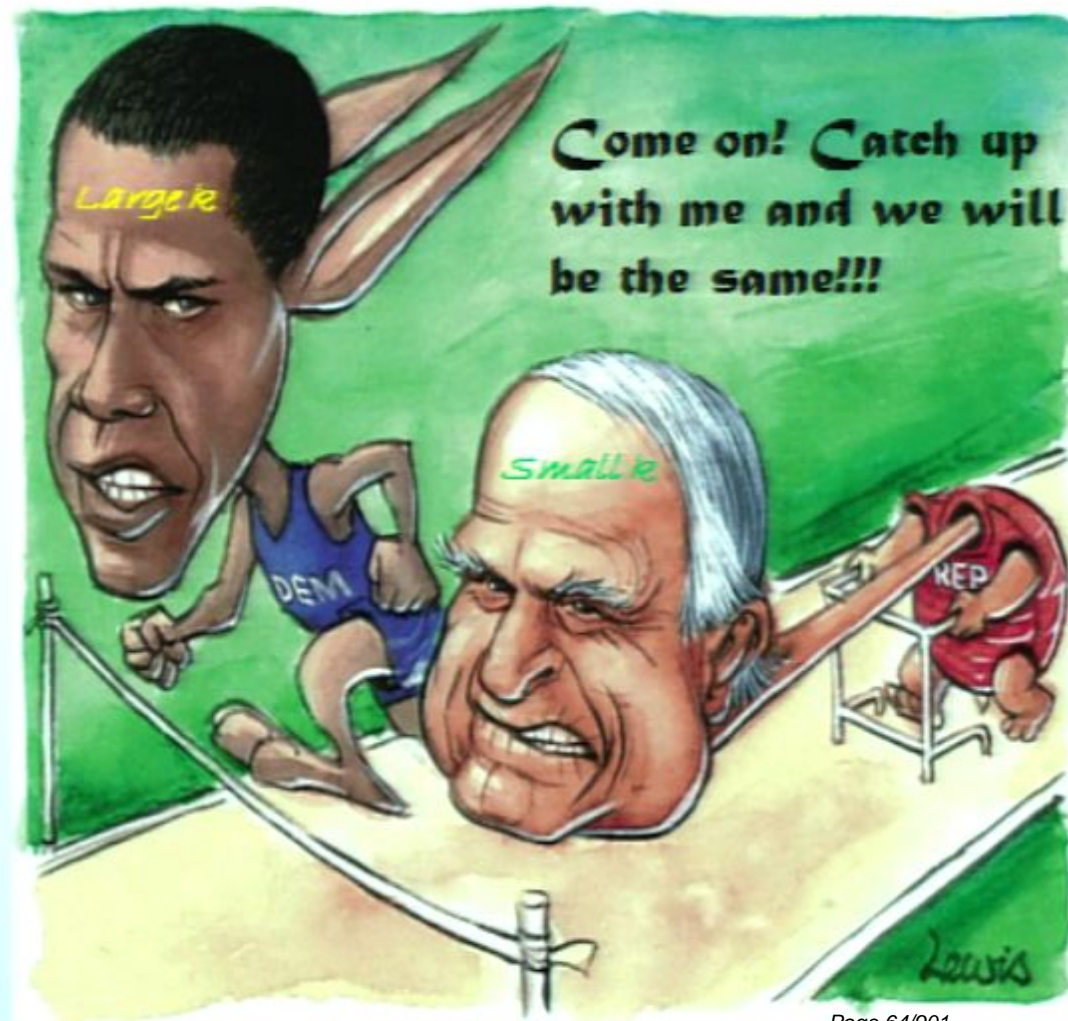
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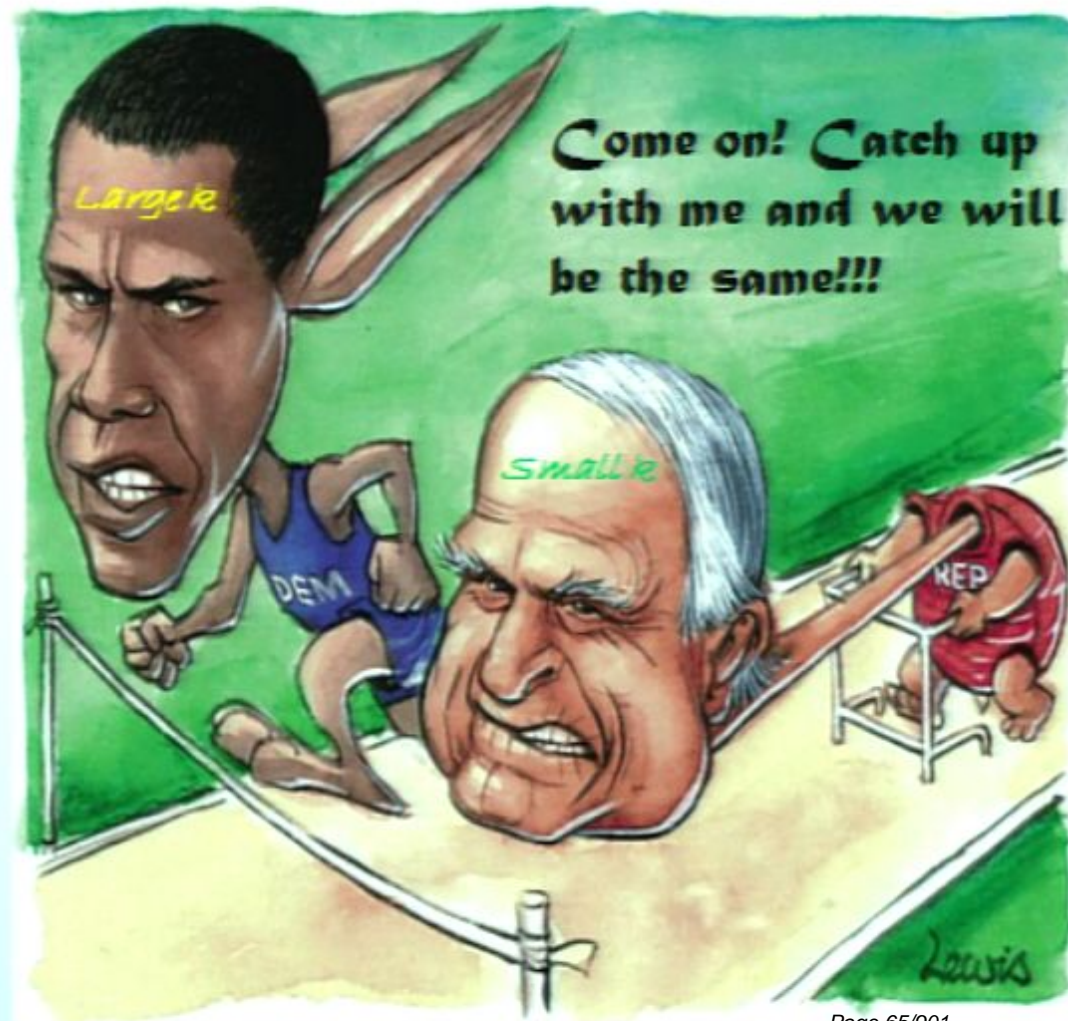
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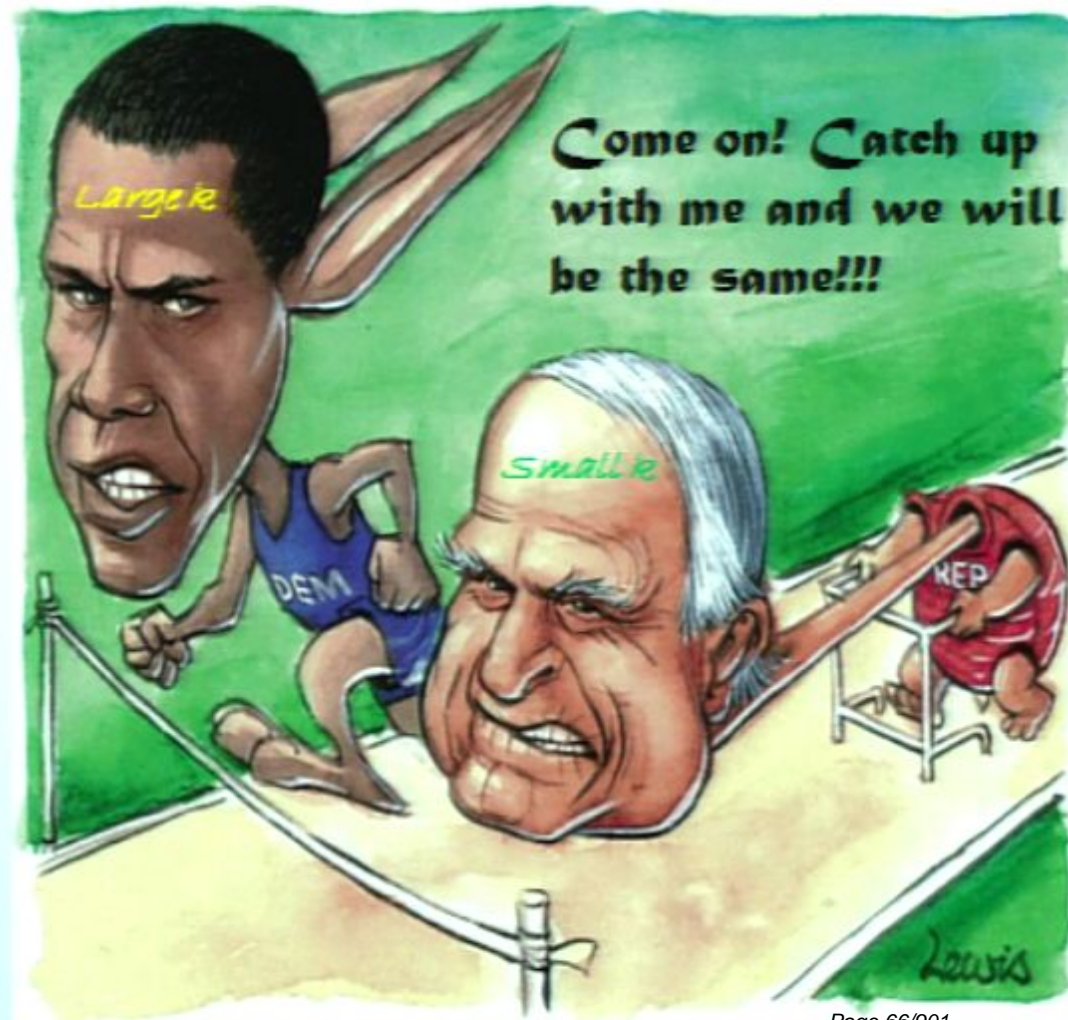
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- ◆ Contracting universe before big bang
- ◆ Cold pressureless matter
- ◆ Scale invariant spectrum

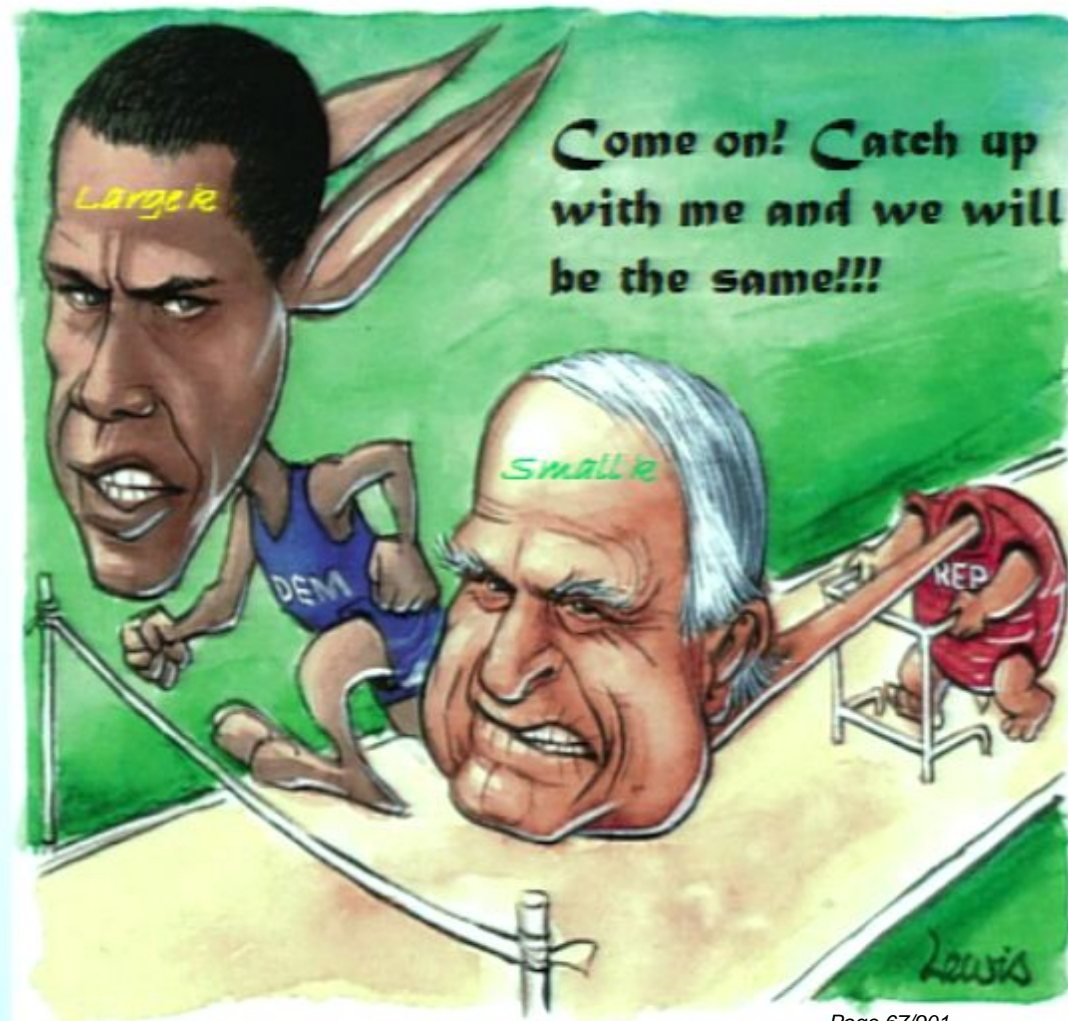
Horizon crossing

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super horizon growing

$$\zeta(t) \propto t^{-1}$$

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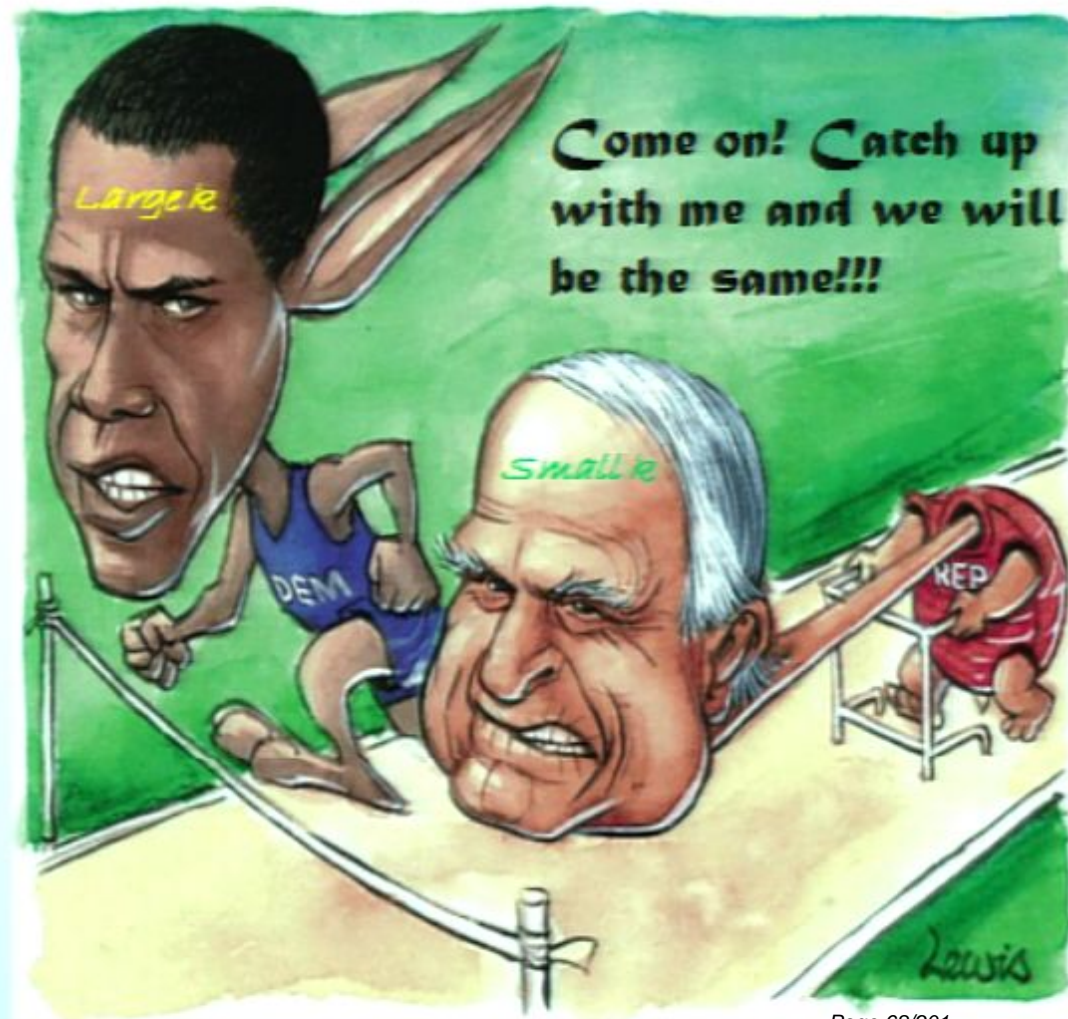
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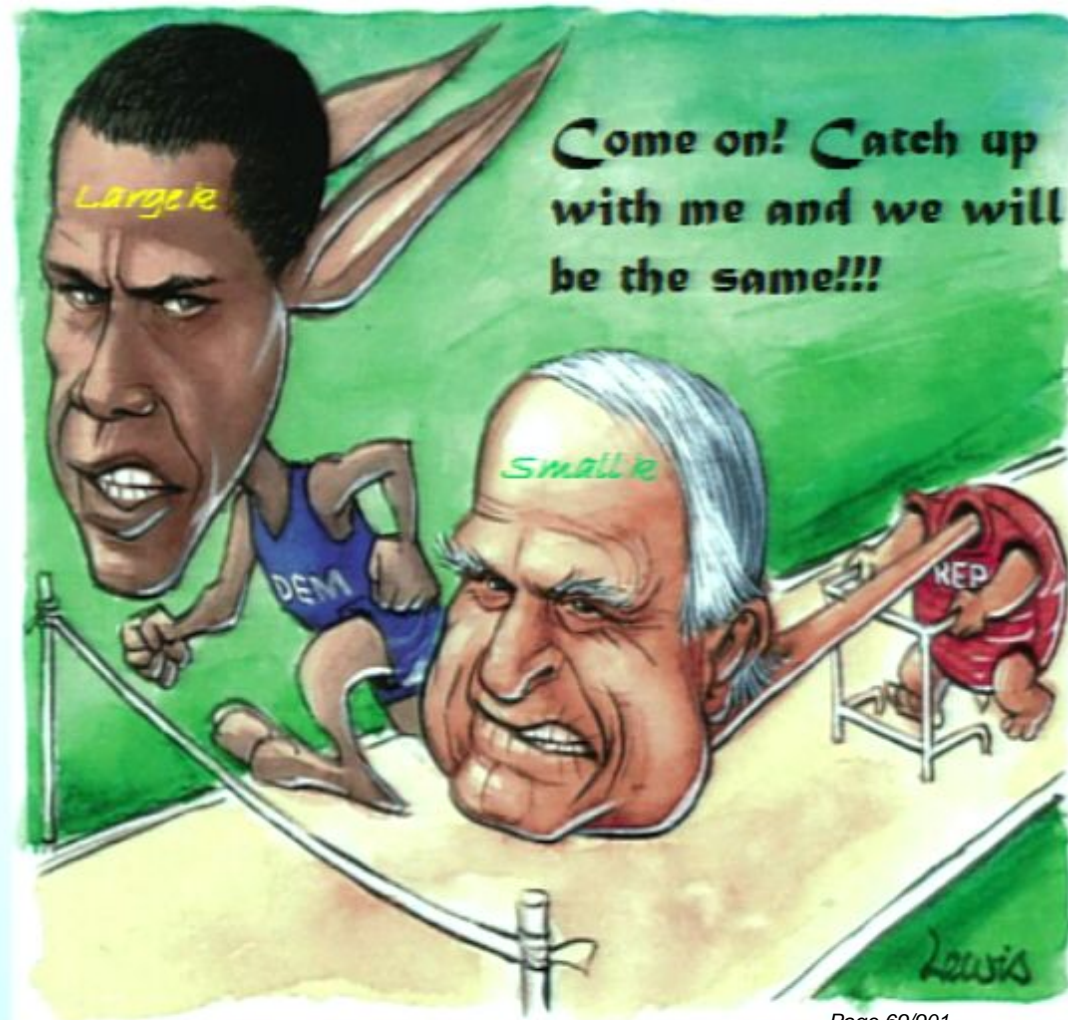
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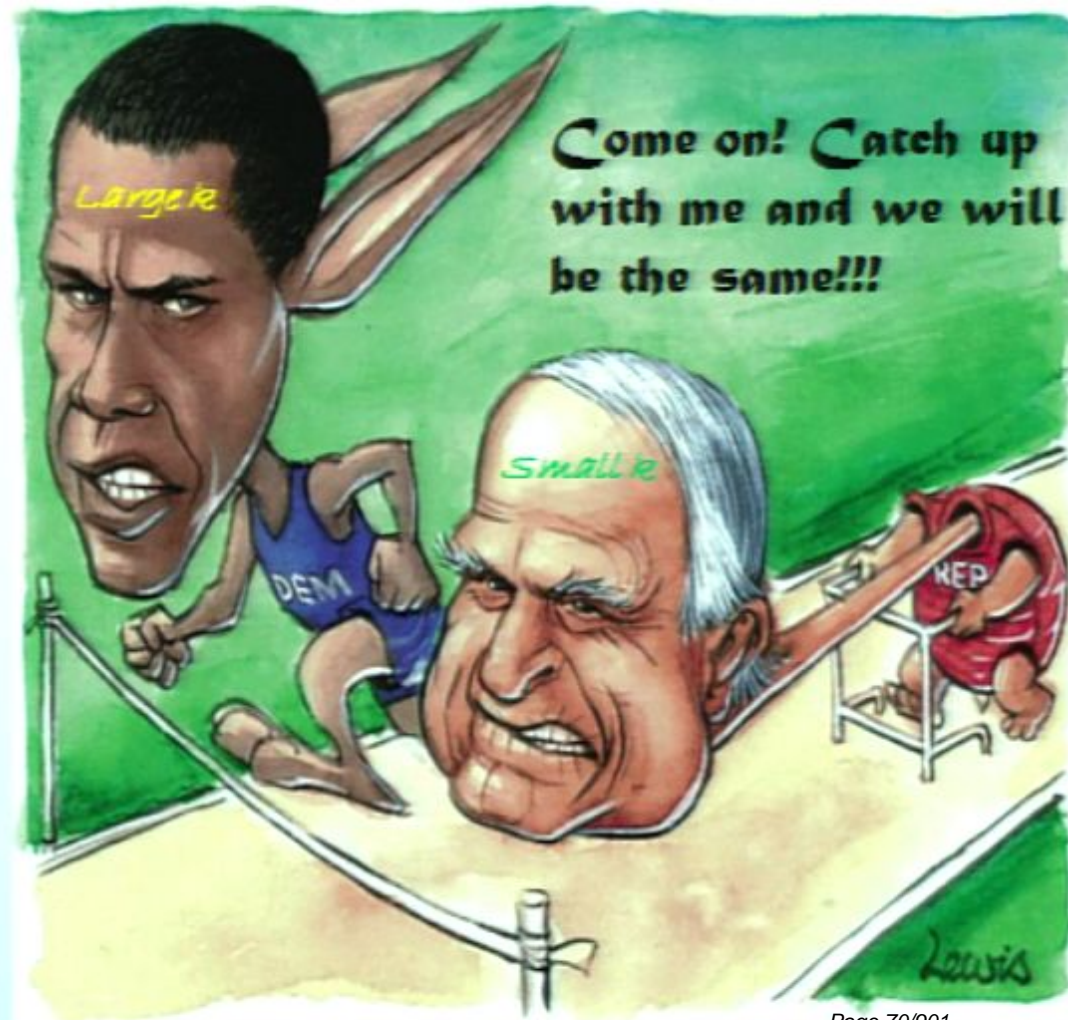
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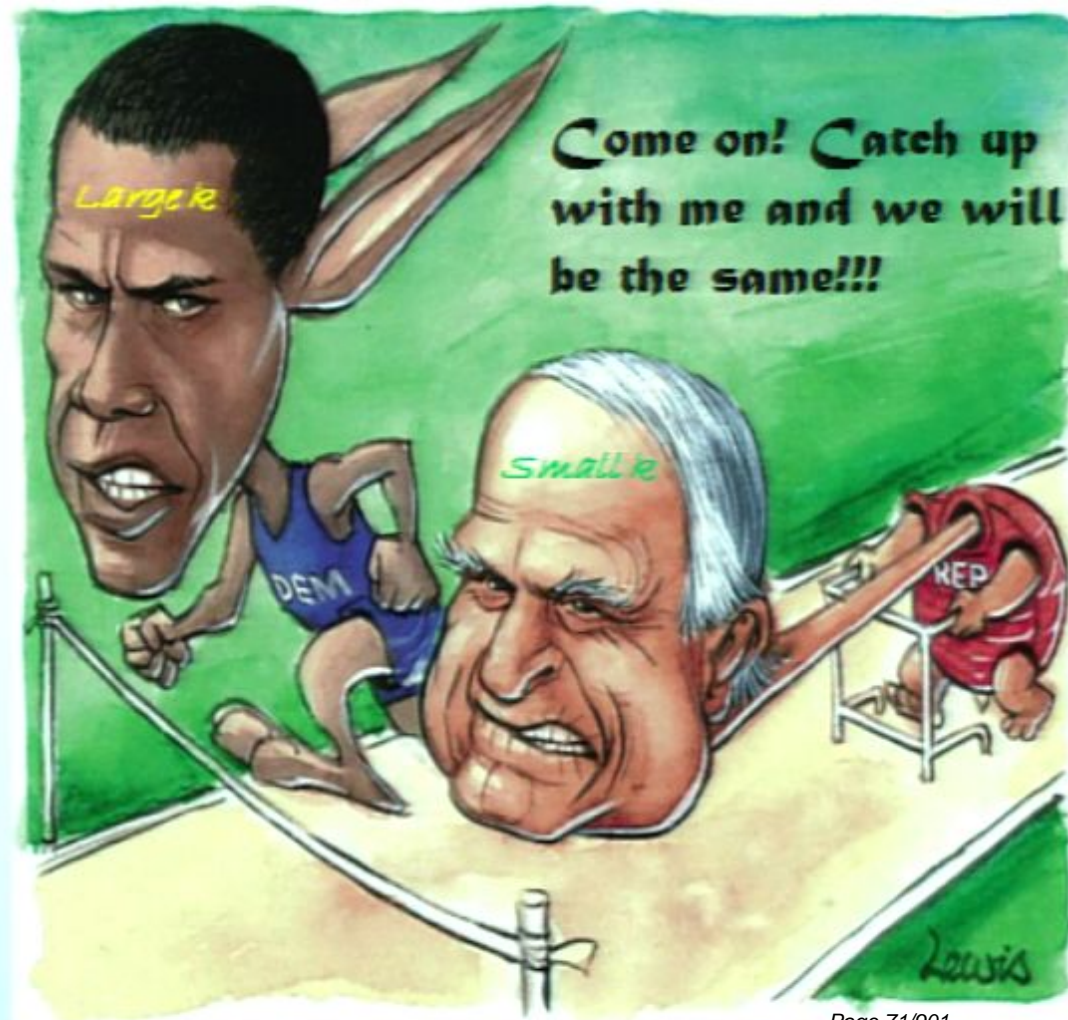
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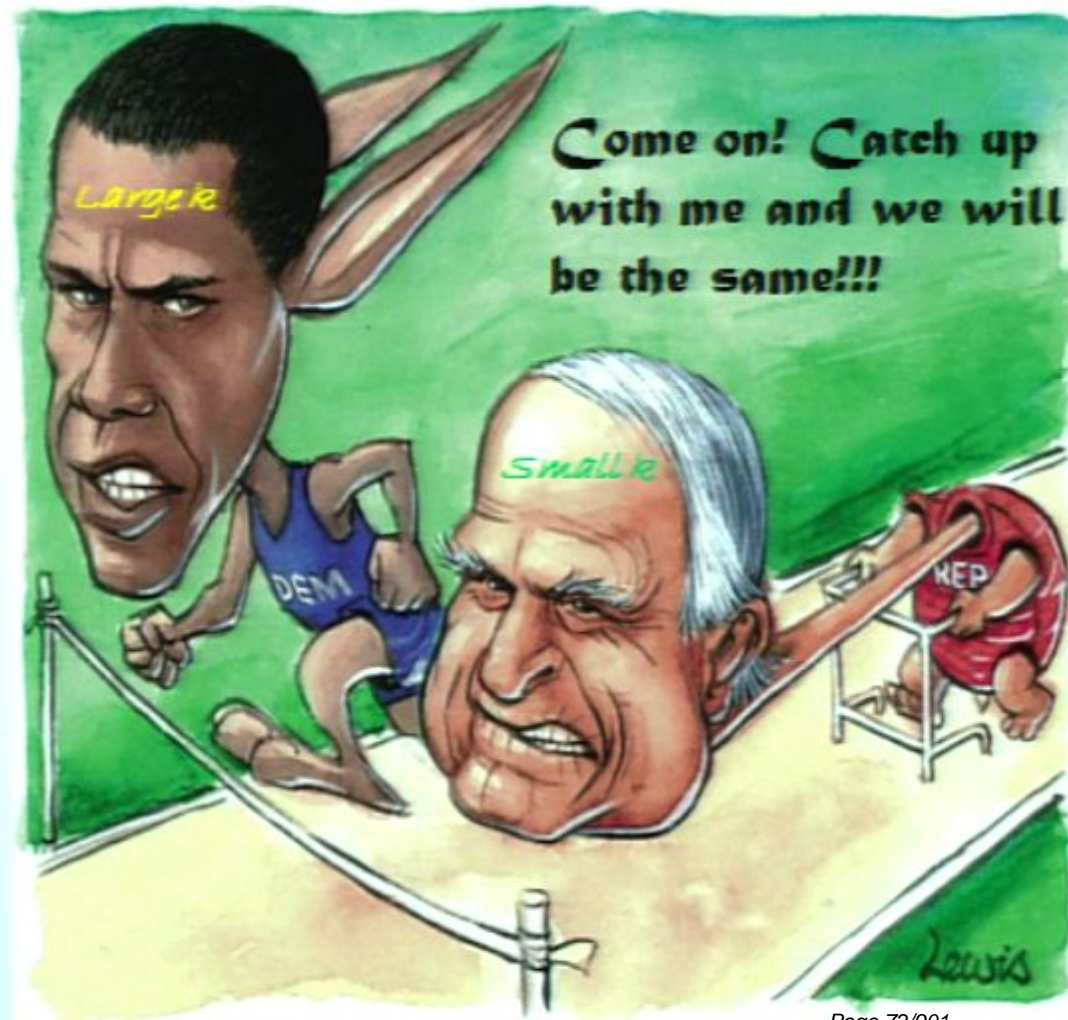
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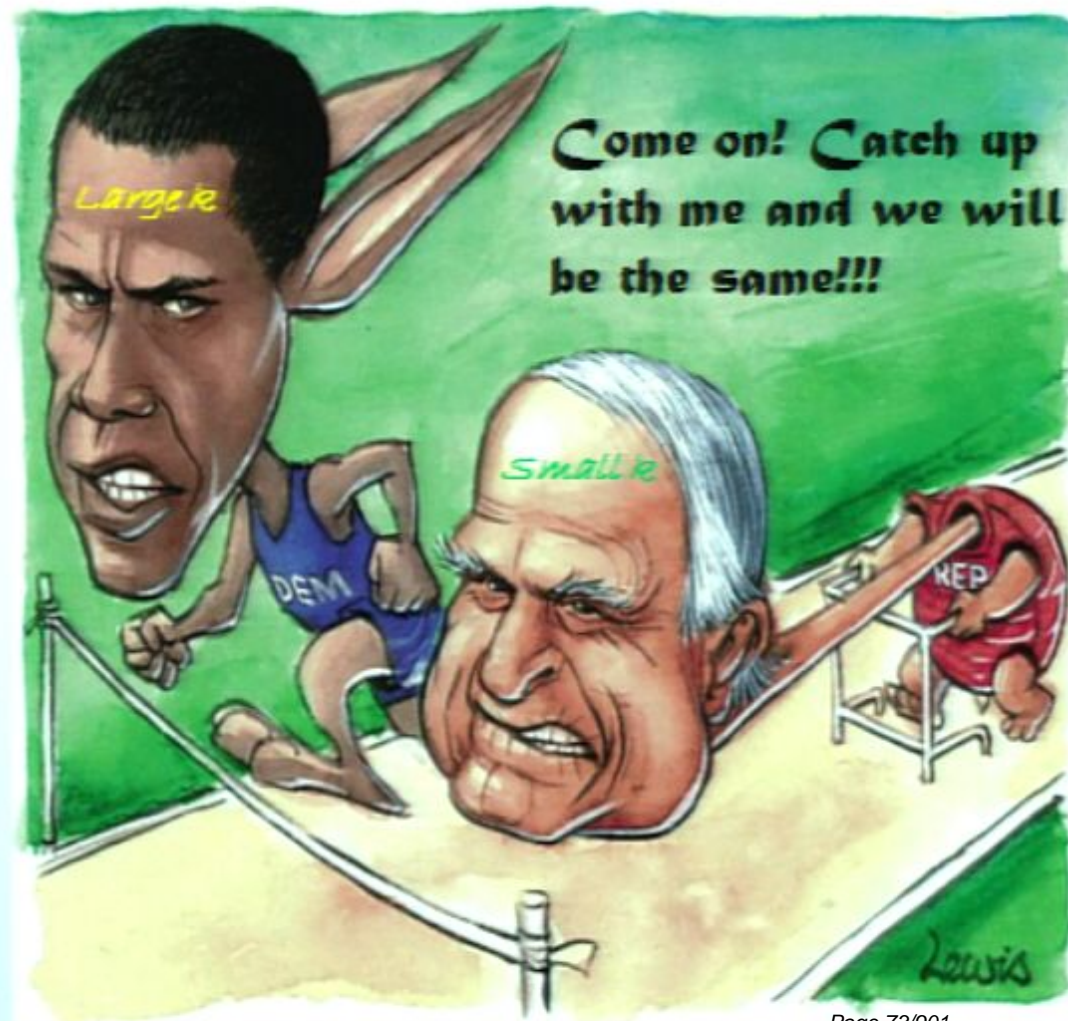
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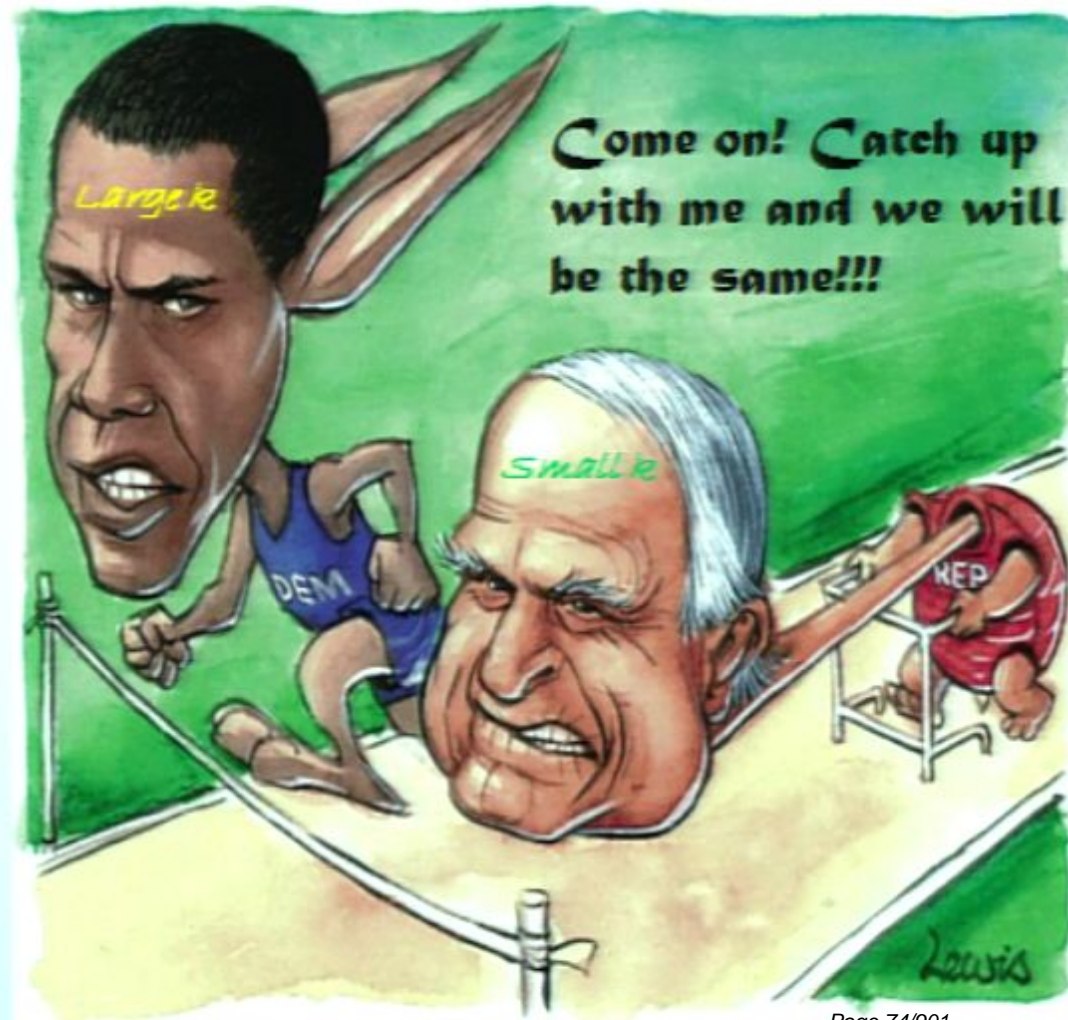
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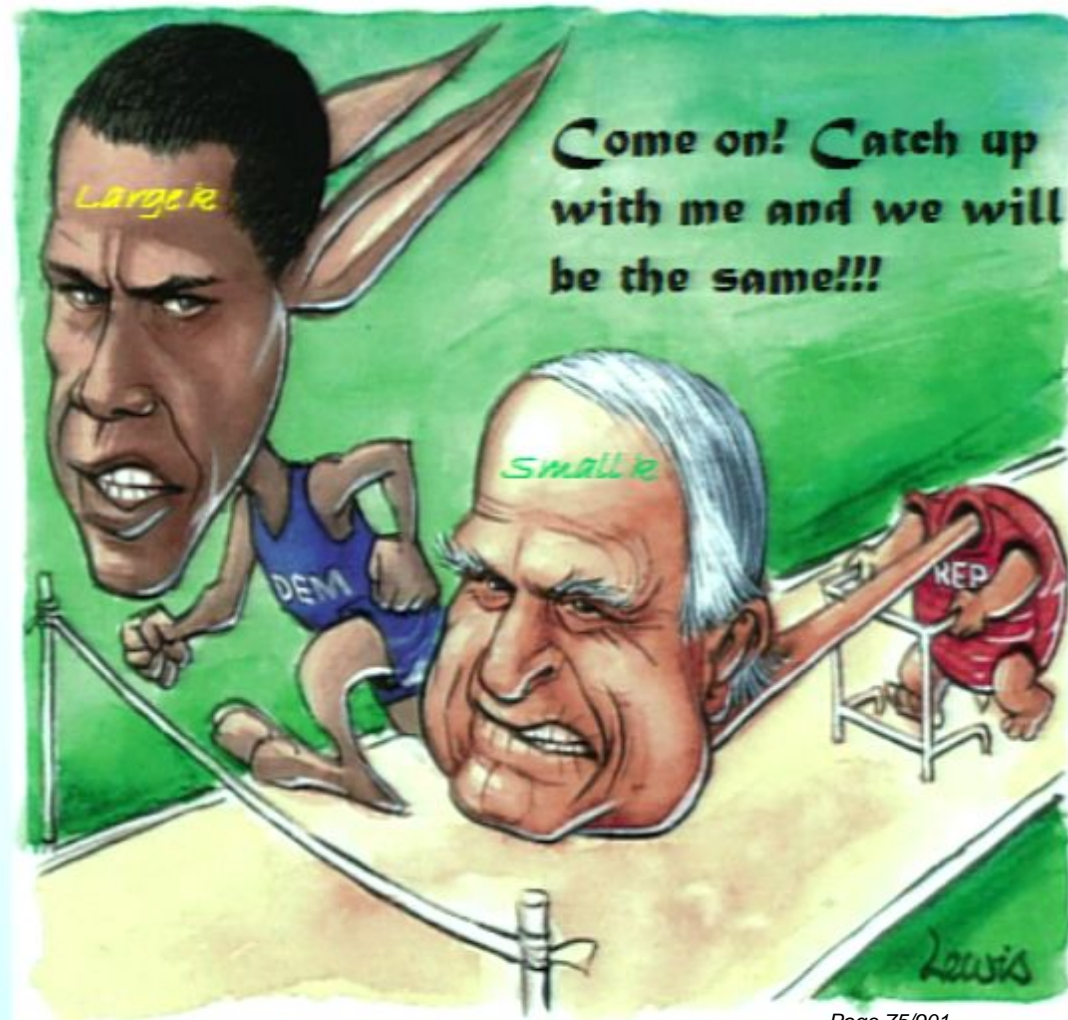
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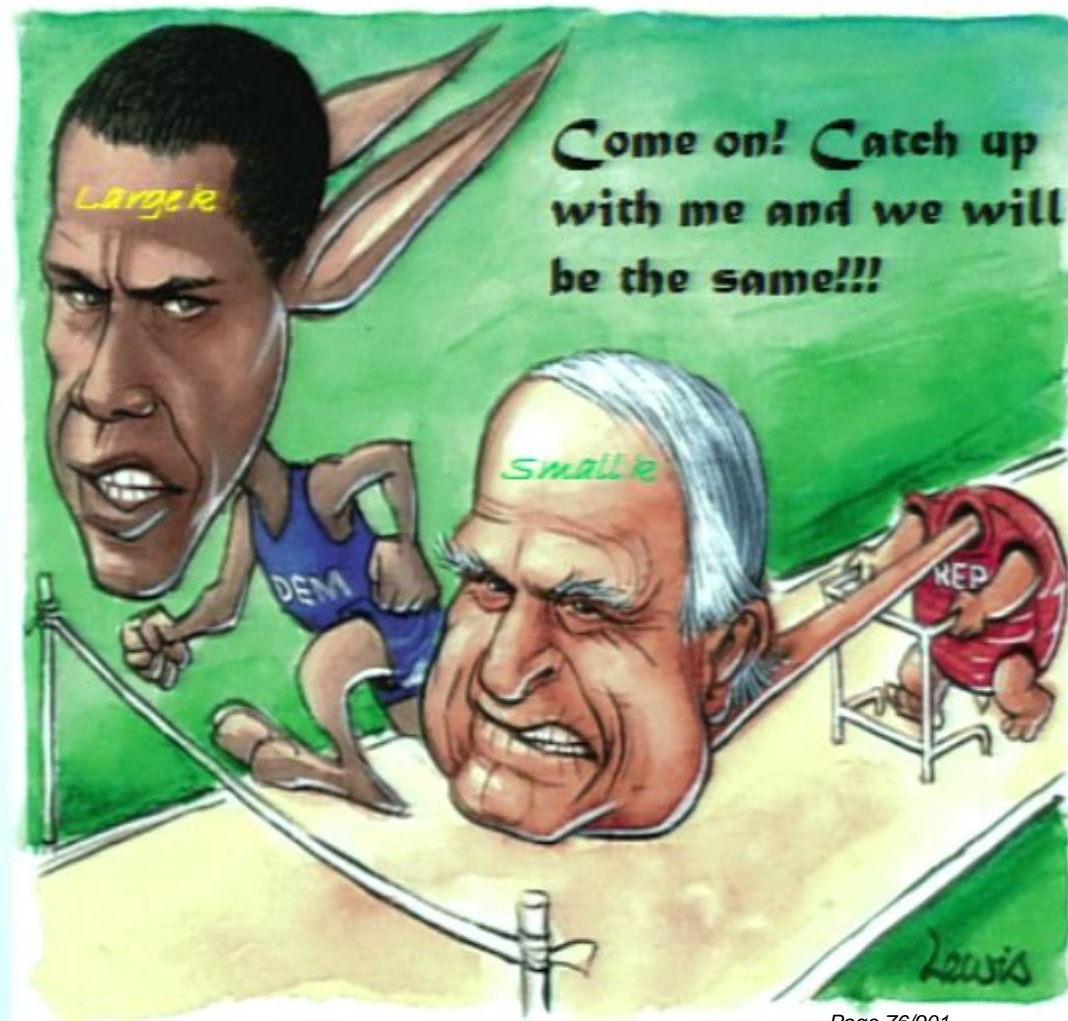
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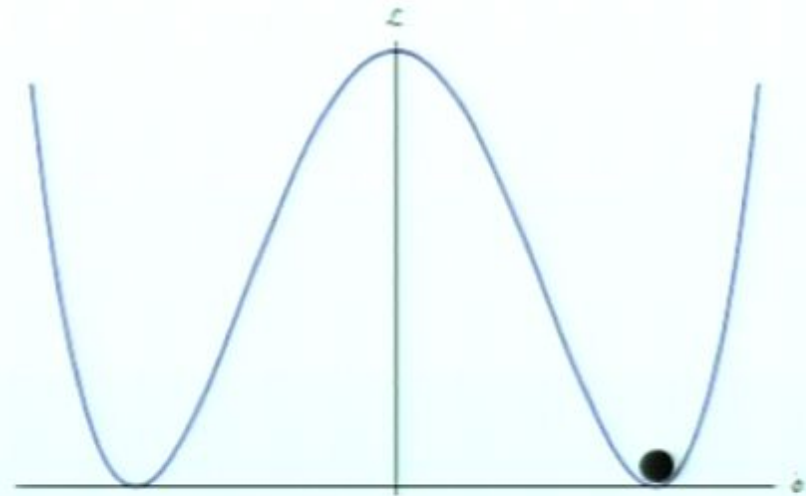
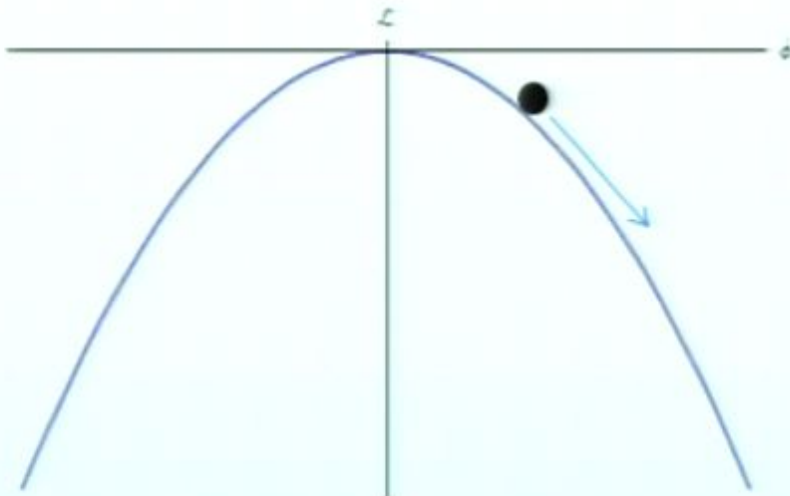
Part II

Ghost Condensation Theory

Ghost Condensation

$$L = -\frac{1}{2} \partial^\mu \phi \partial_\mu \phi + \dots$$

$$P = \frac{1}{8} (X - c^2)^2, \quad X = \partial^\mu \phi \partial_\mu \phi$$



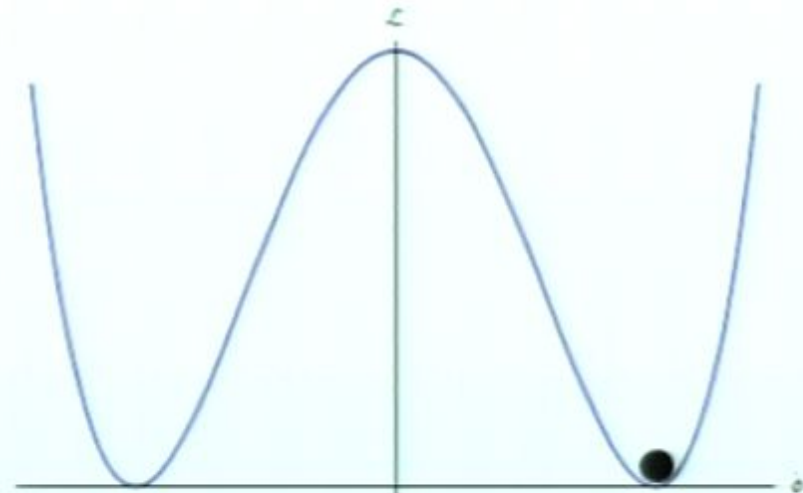
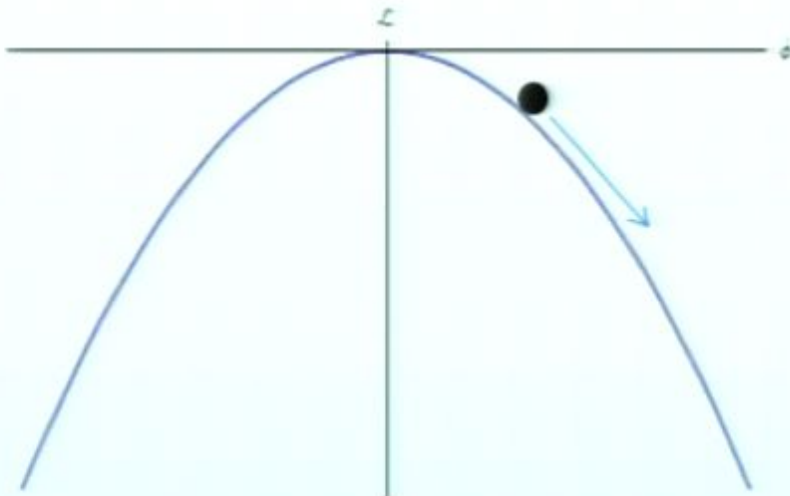
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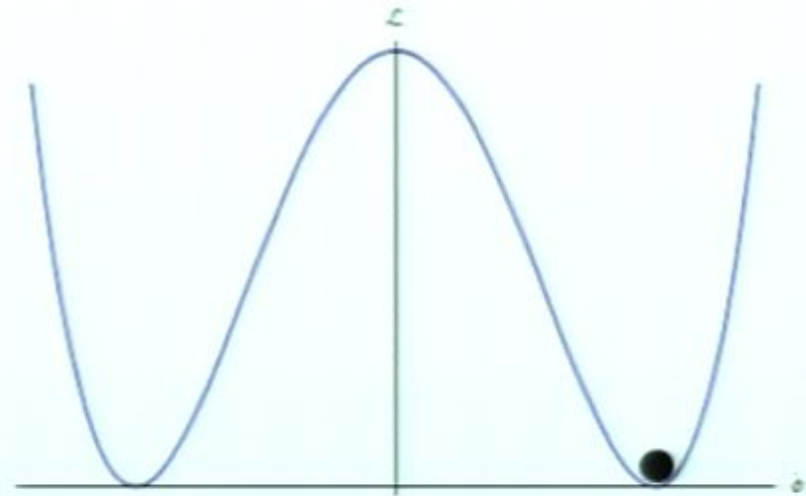
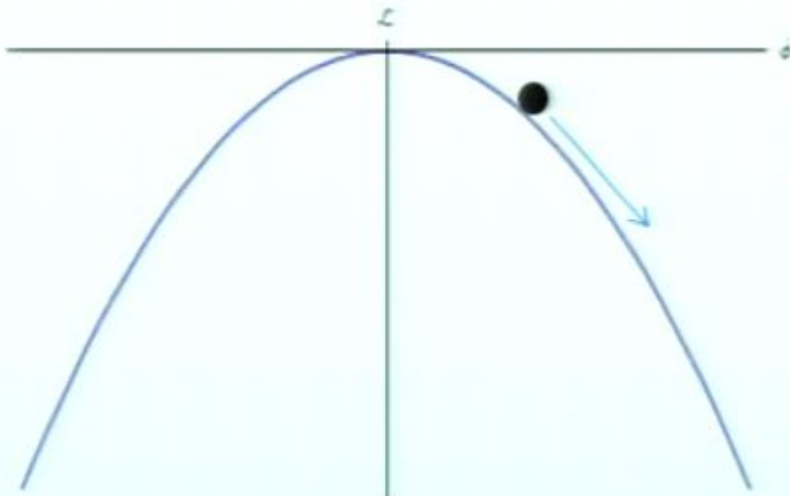
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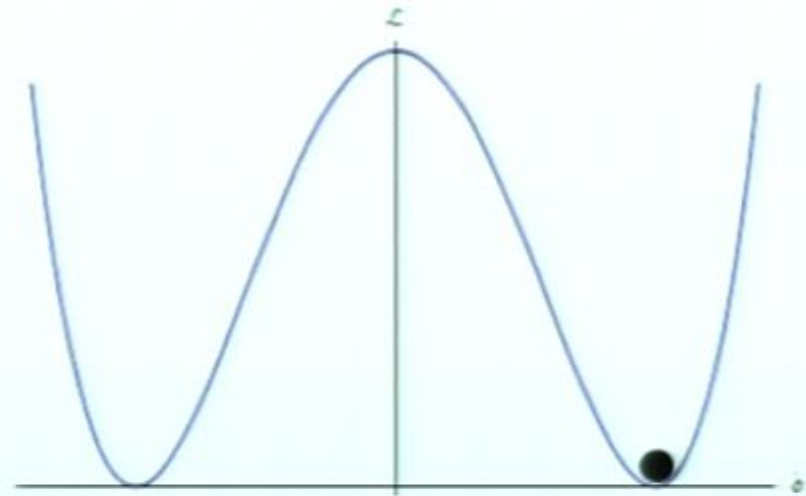
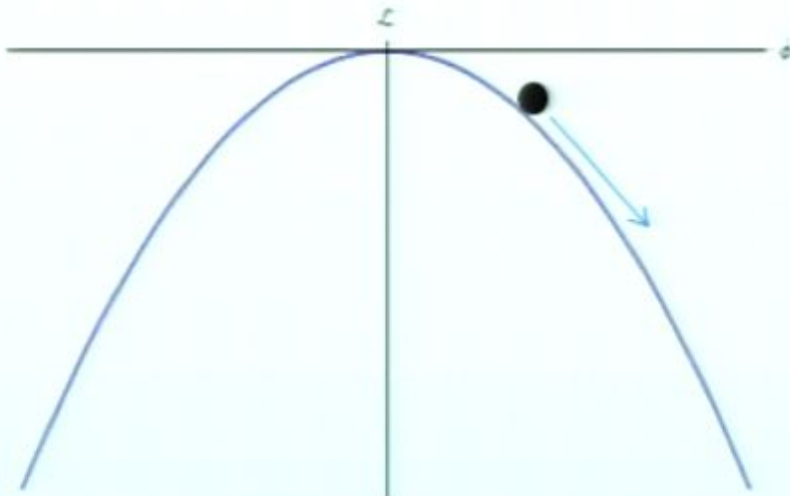
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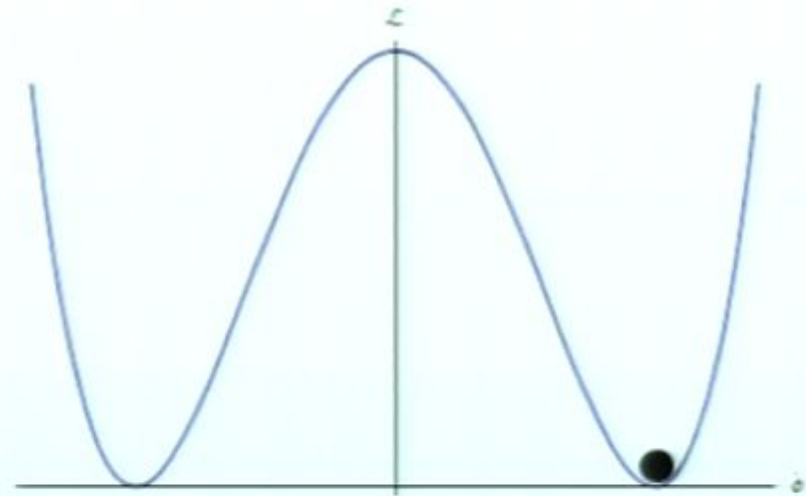
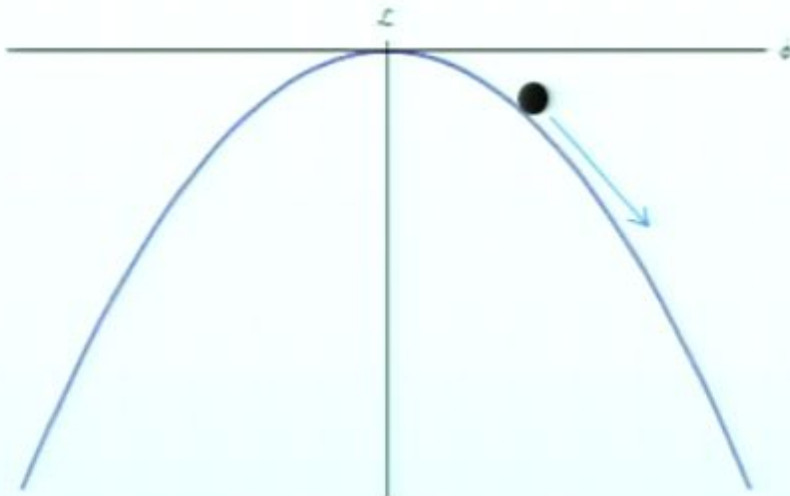
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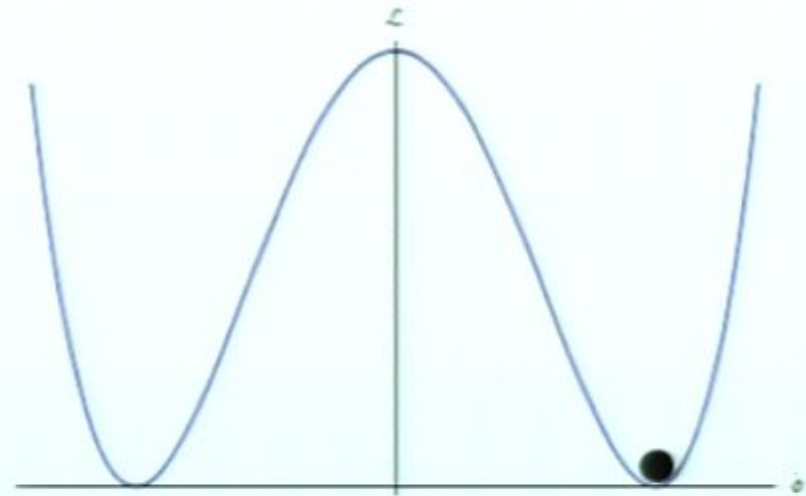
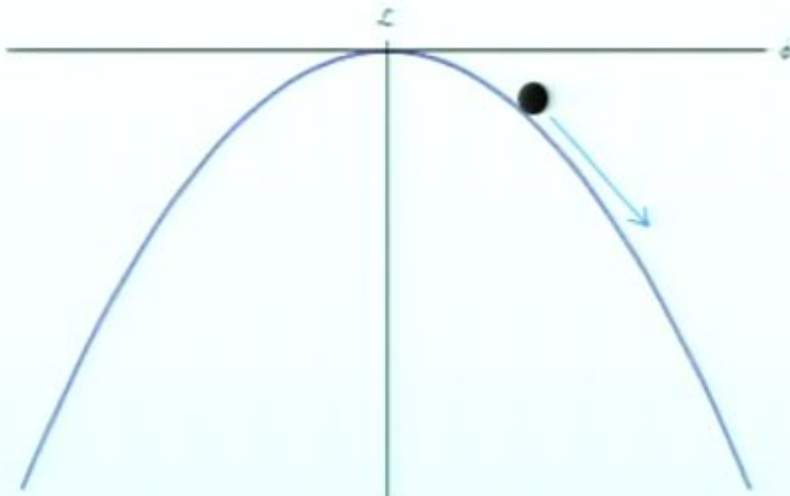
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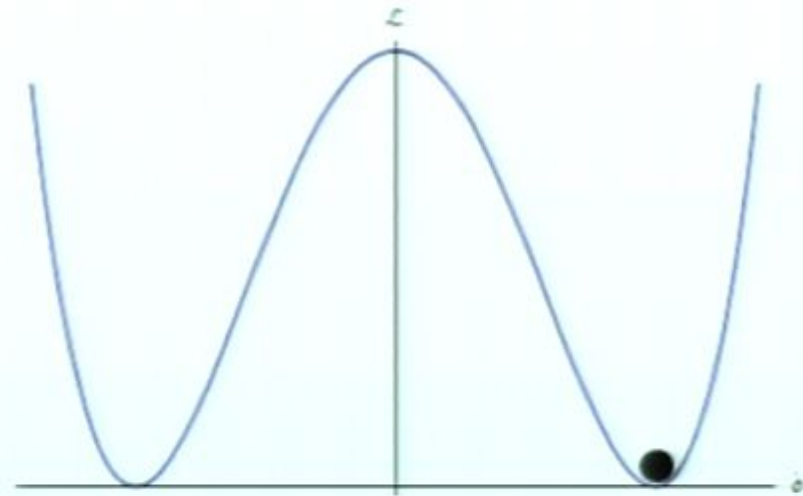
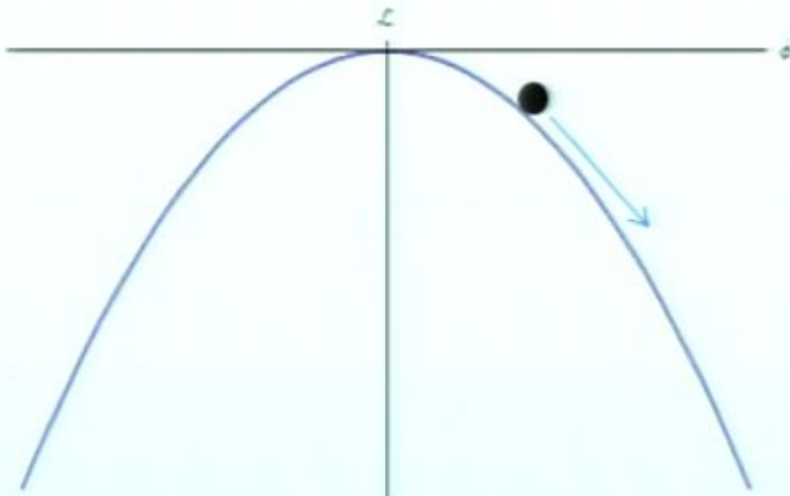
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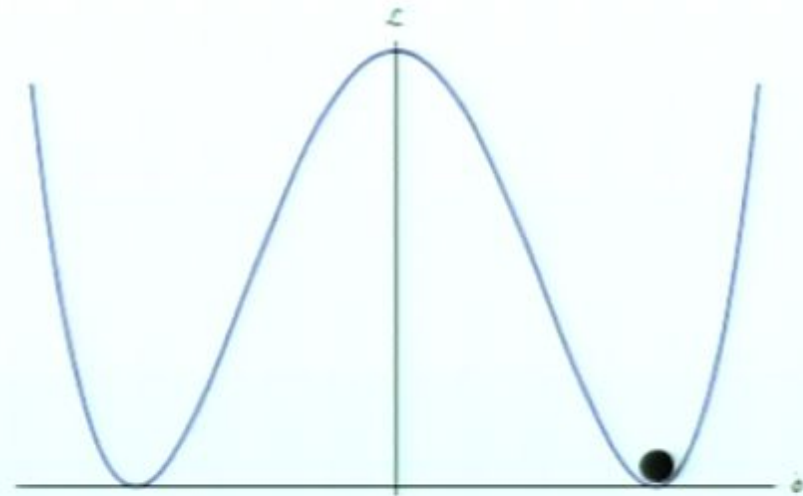
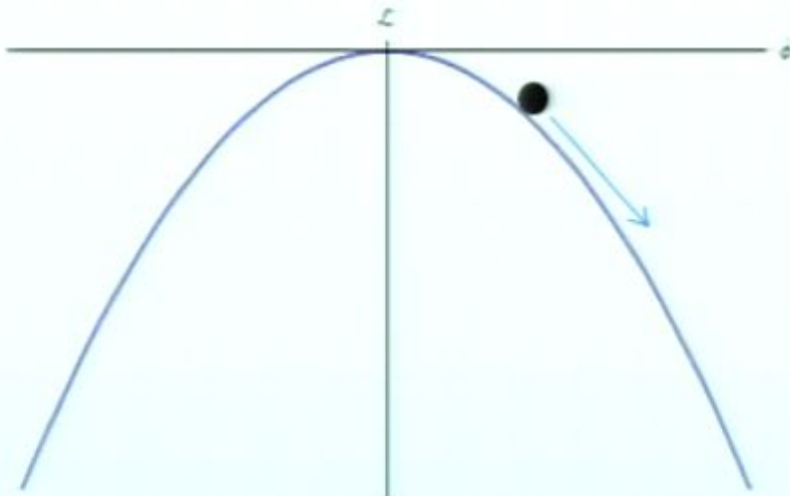
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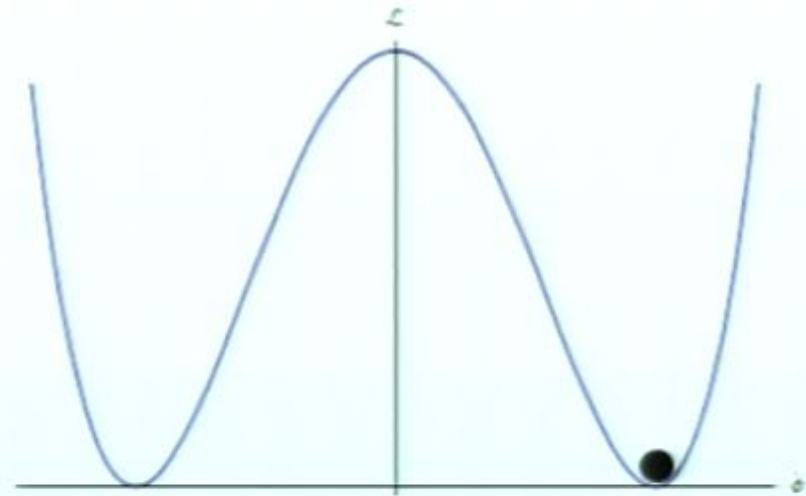
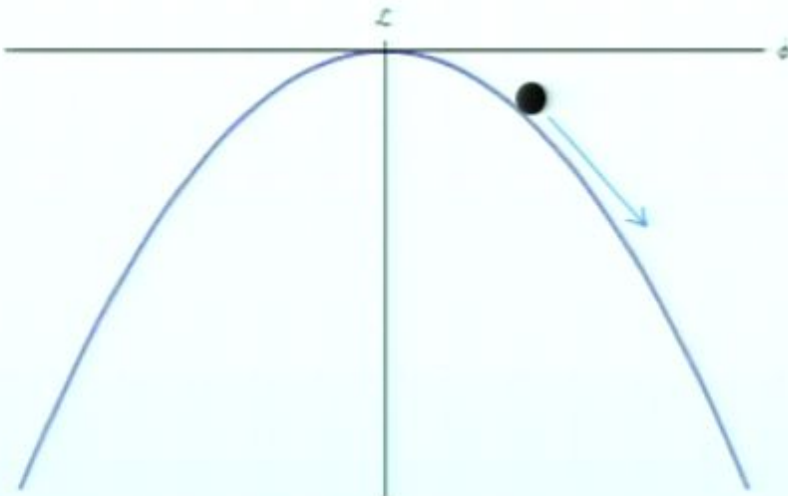
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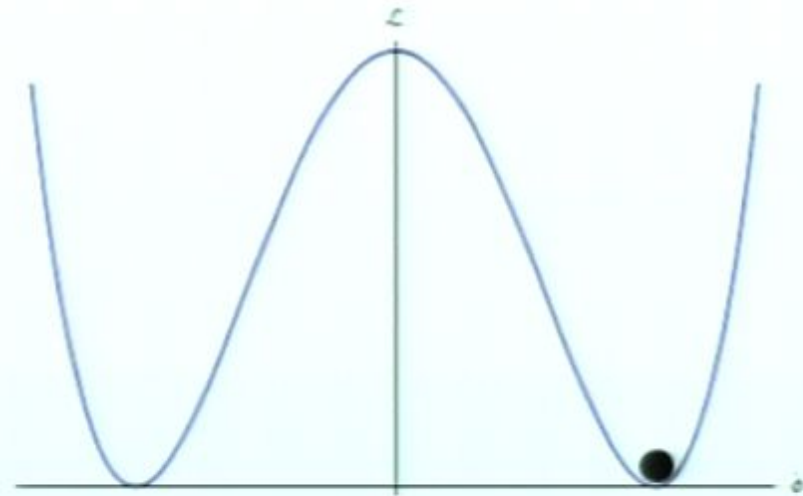
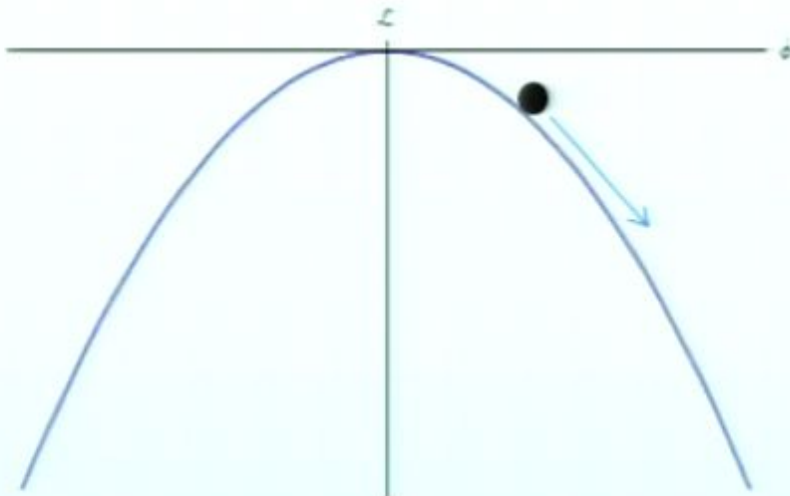
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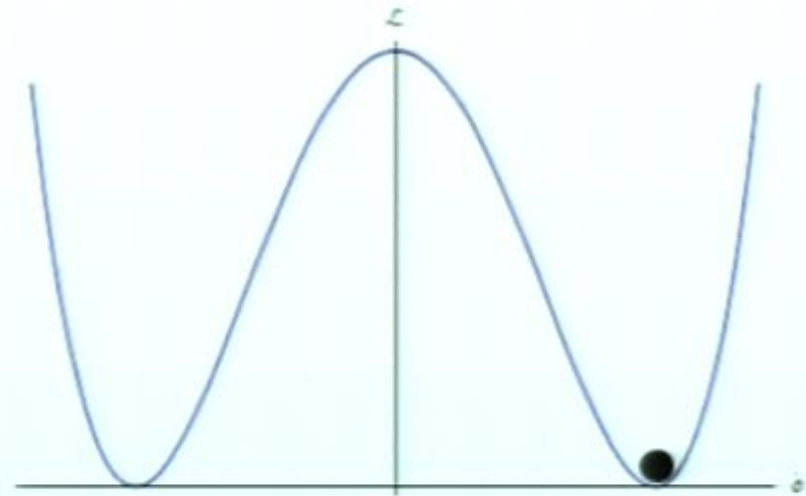
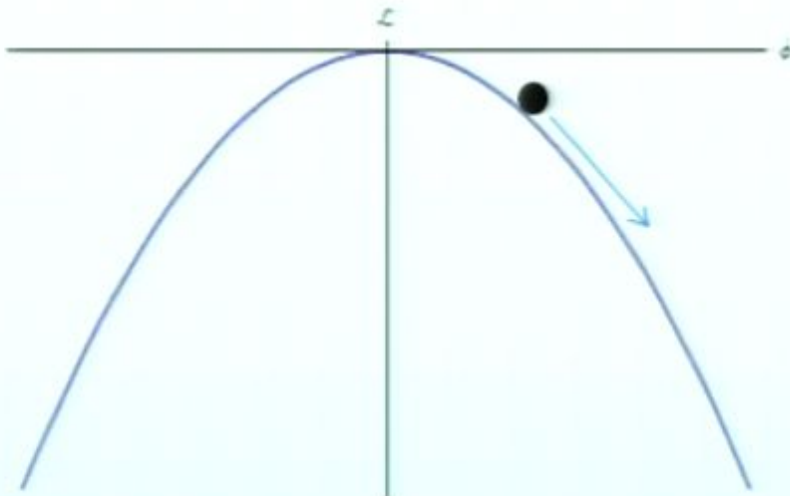
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Inflation, Dark Energy

Dark matter

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Ghost field locate at the minima, with scalar excitation

$$\phi = ct + \pi$$

Low energy effective action for π is

$$S \sim \int d^4x \left[\frac{1}{2} \dot{\pi}^2 - \frac{1}{2M^2} (\nabla^2 \pi)^2 + \dots \right],$$

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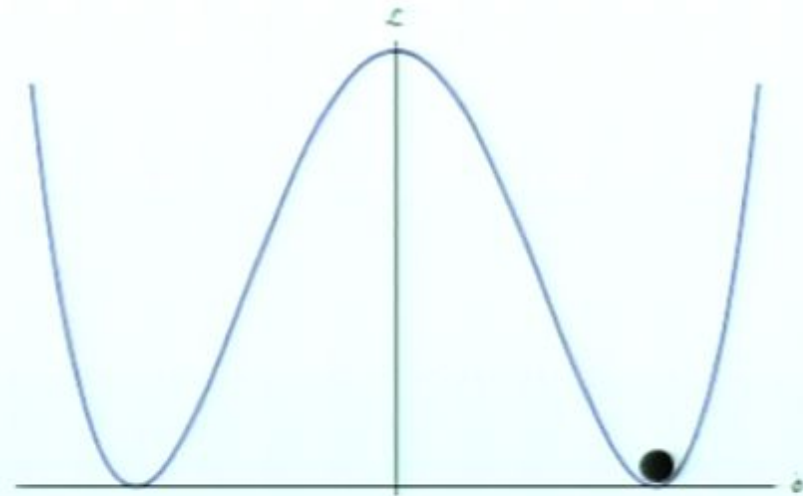
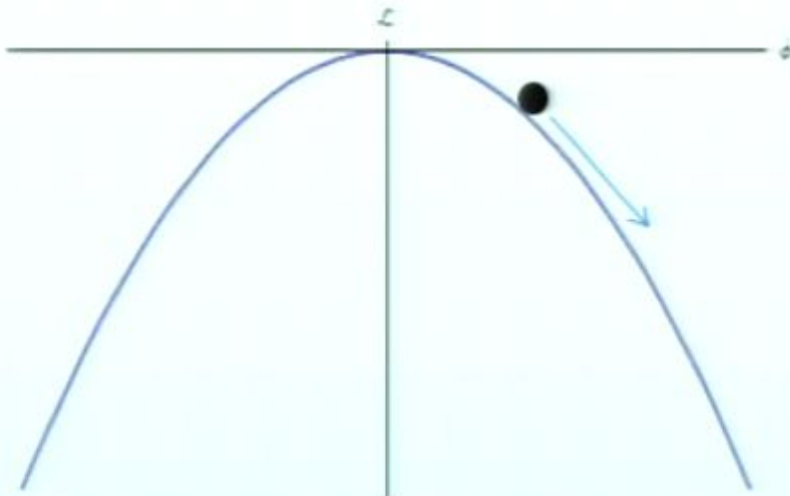


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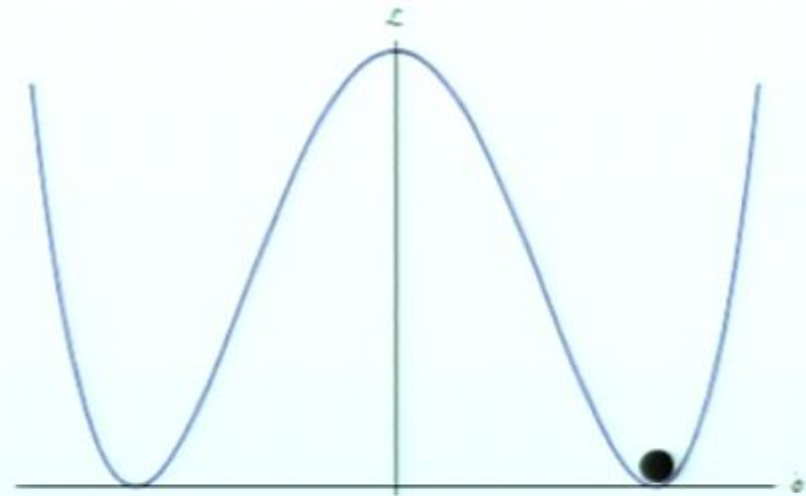
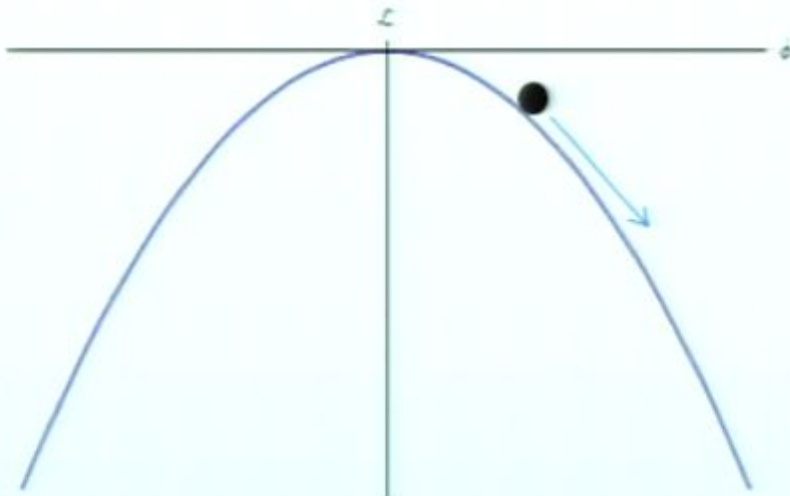
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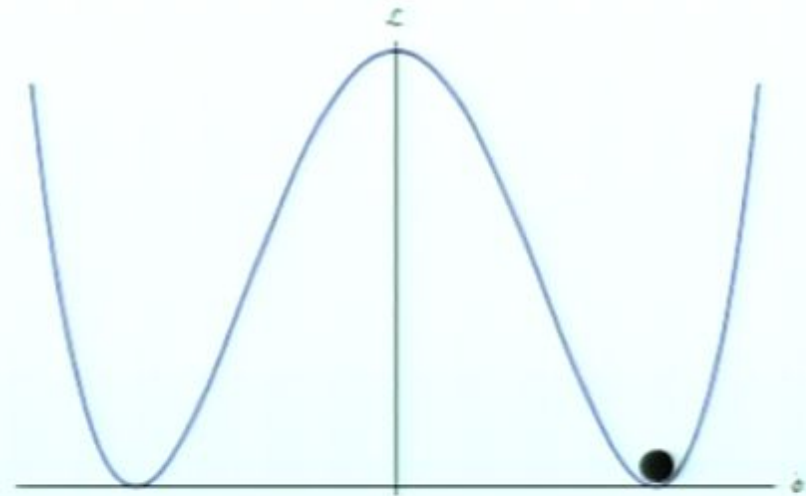
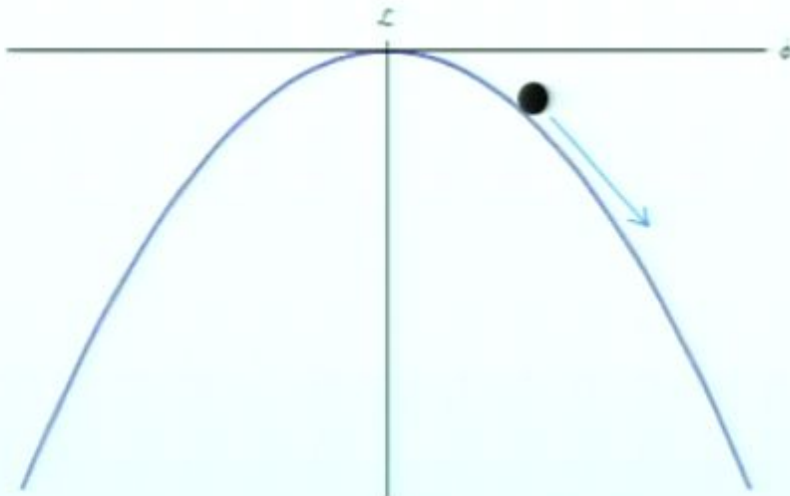
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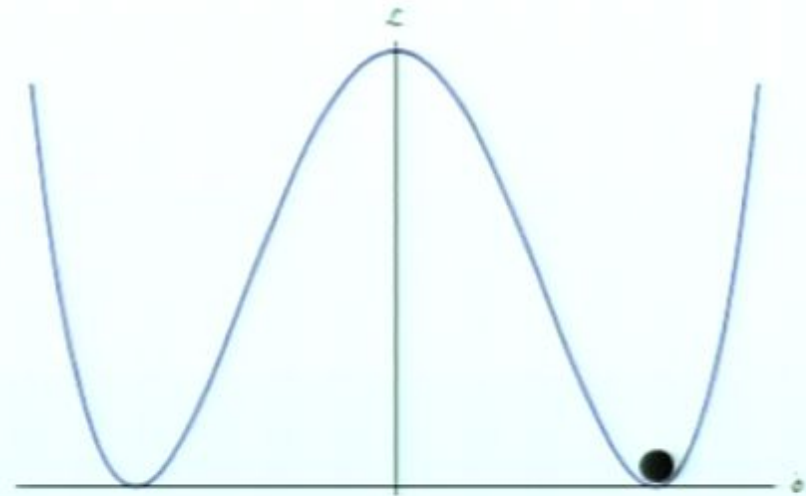
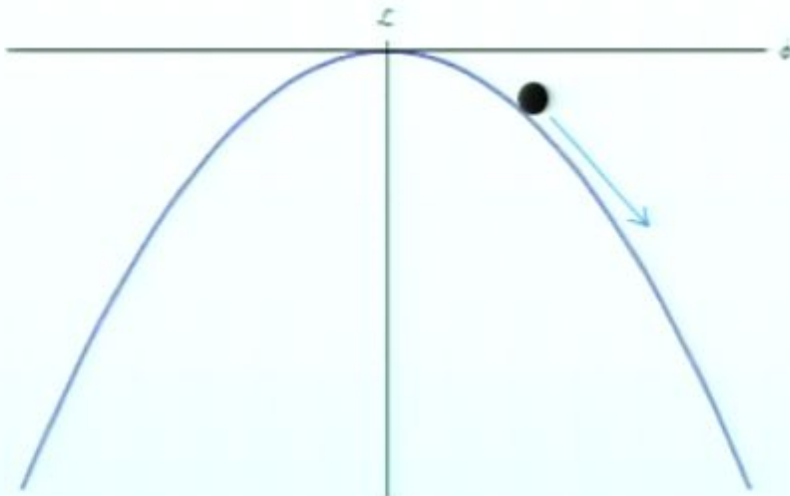
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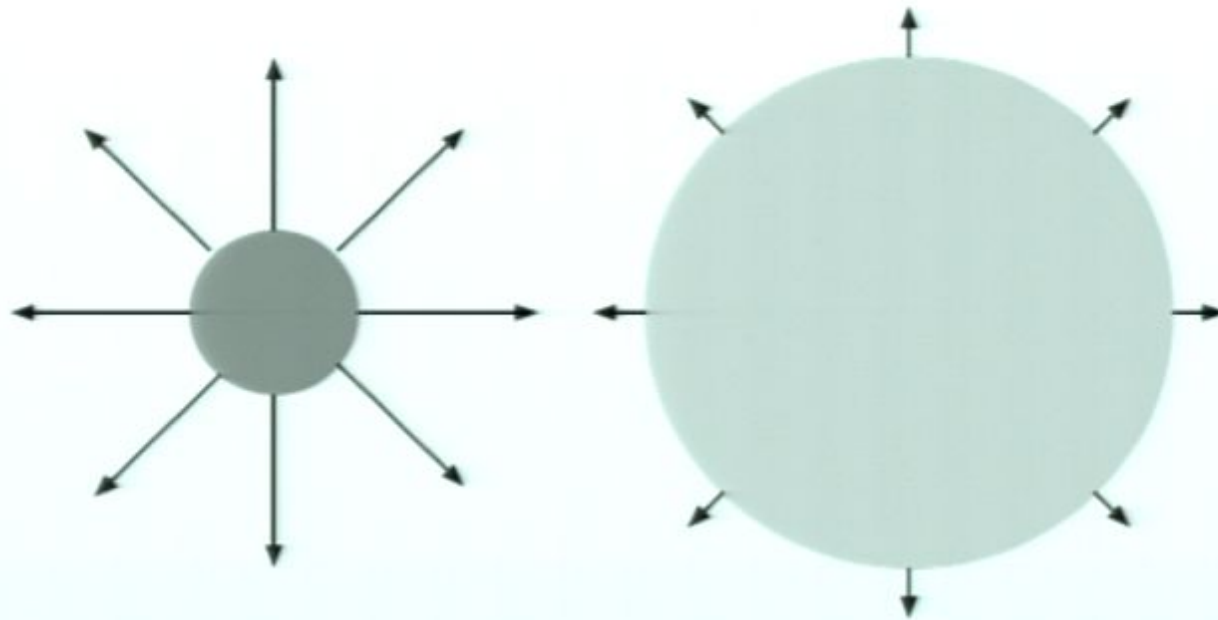
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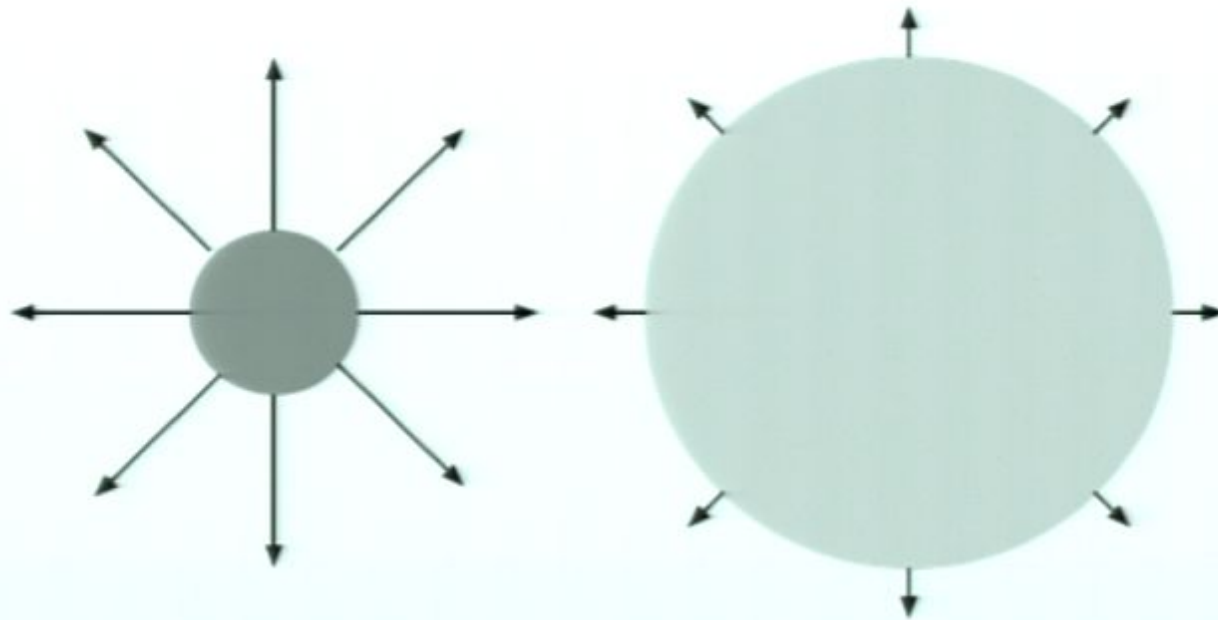
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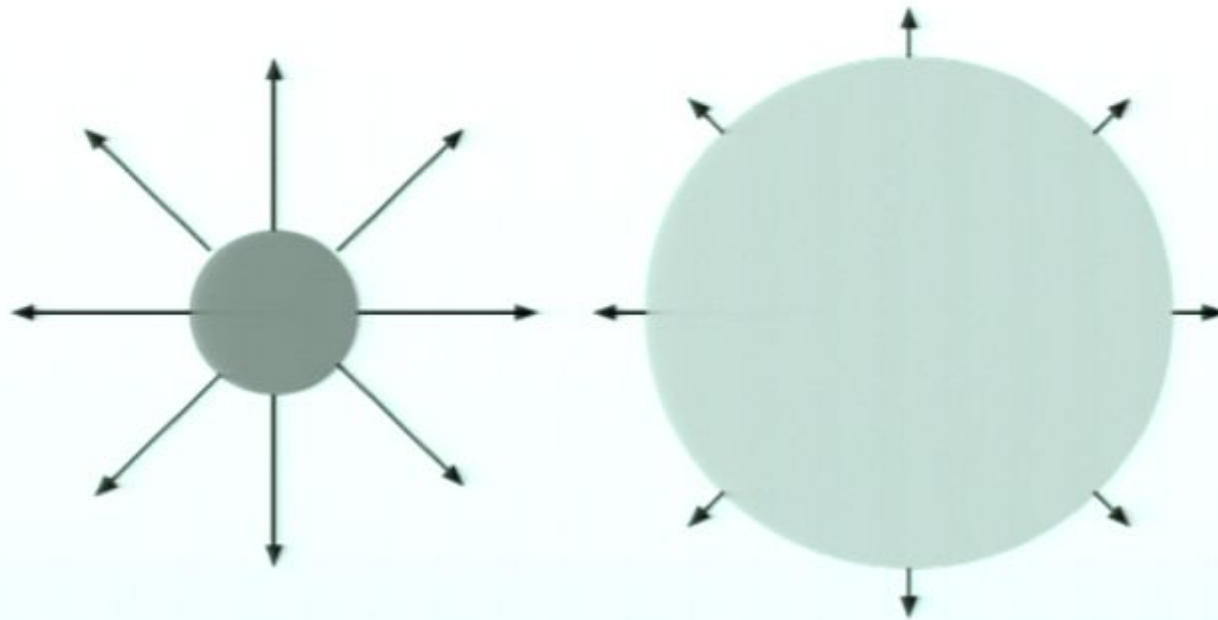
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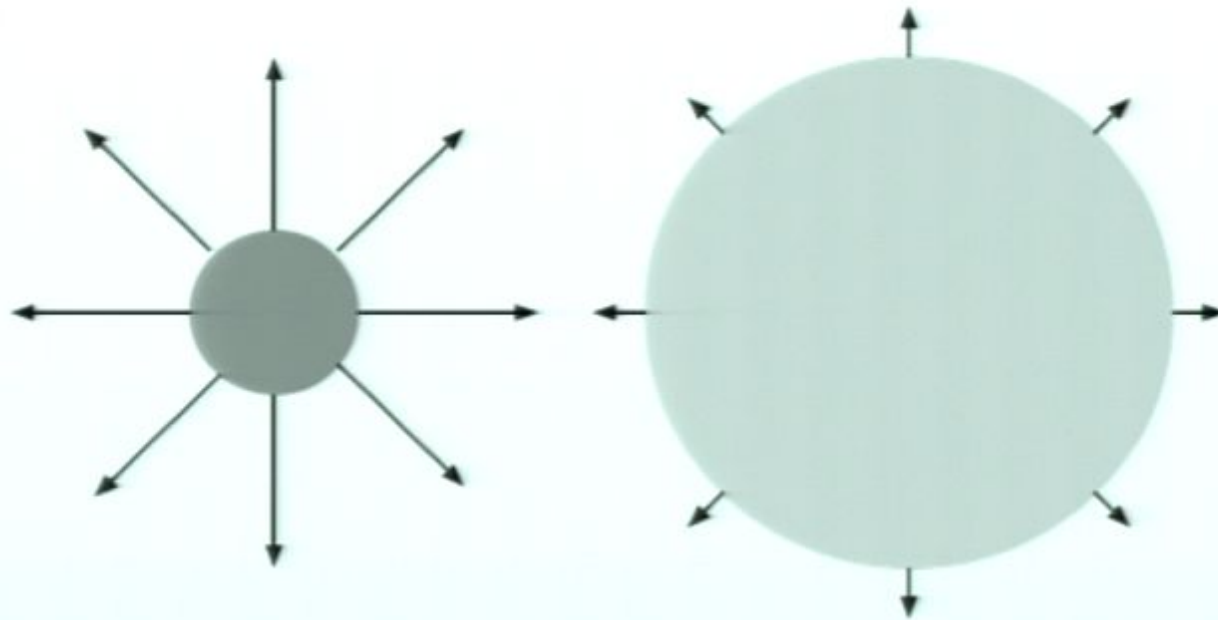
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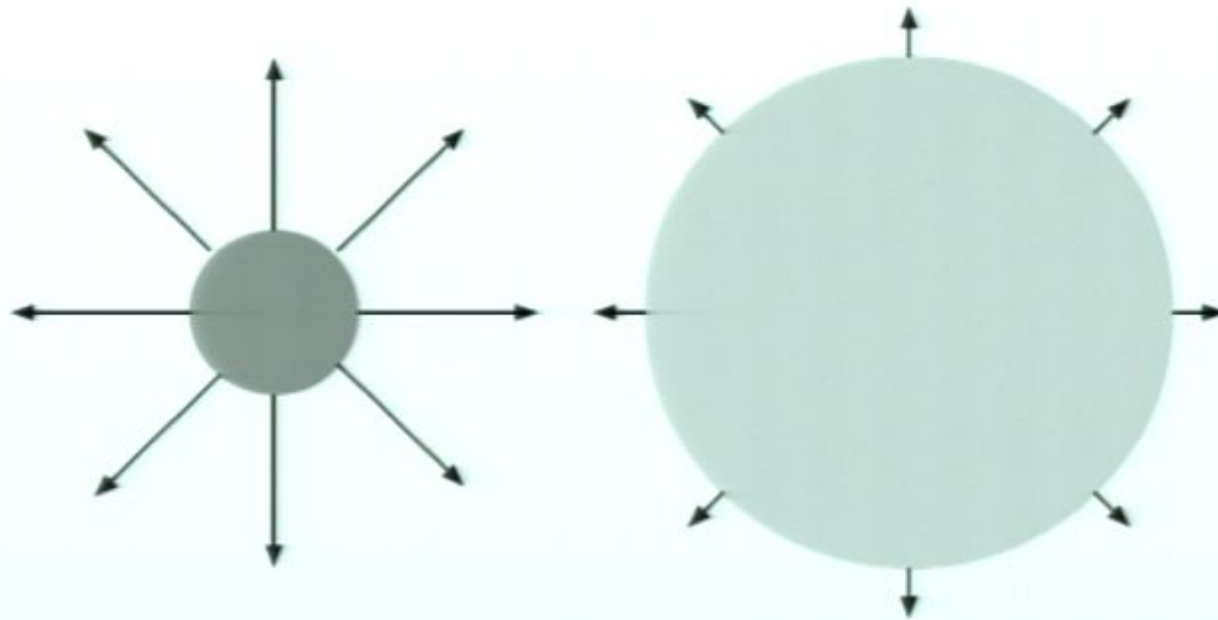
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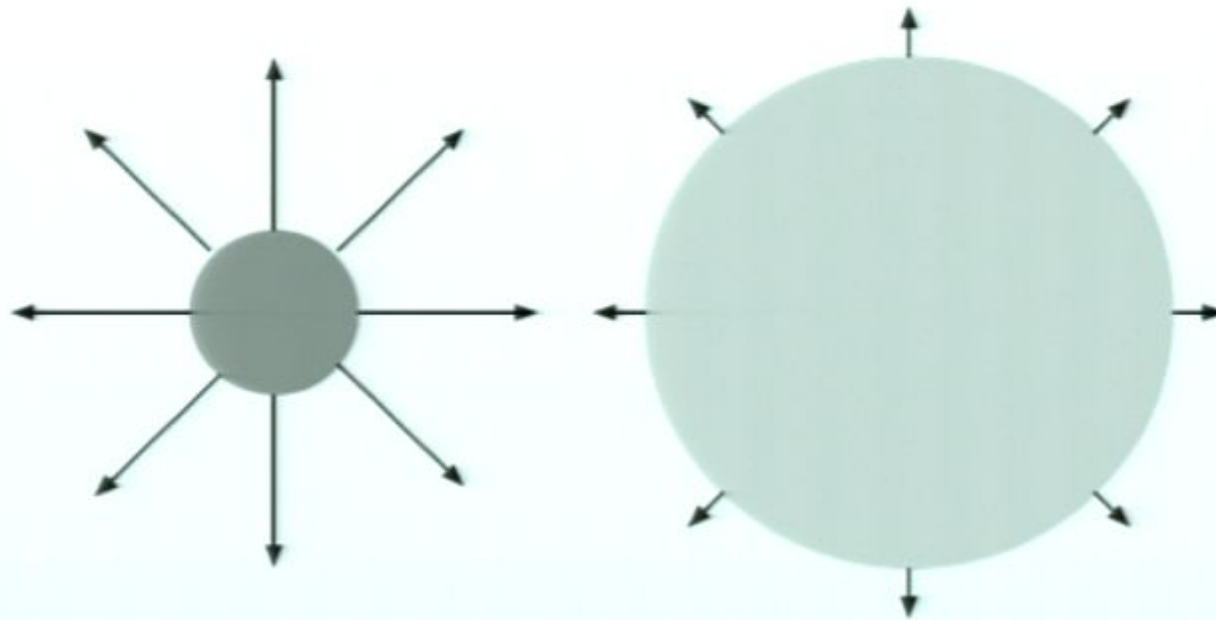
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Ghost field locate at the minima, with scalar excitation

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Inflation, Dark Energy

Dark matter



Matter Bounce

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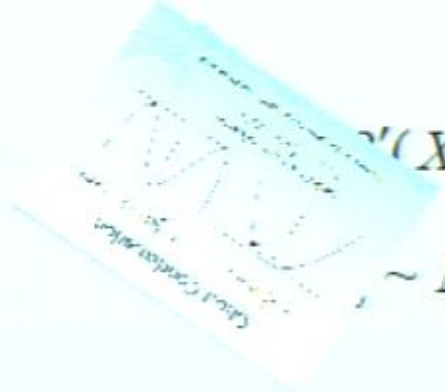
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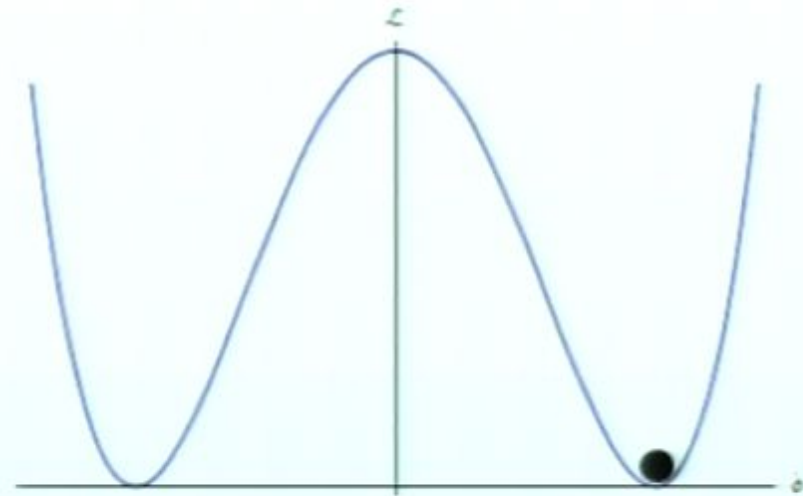
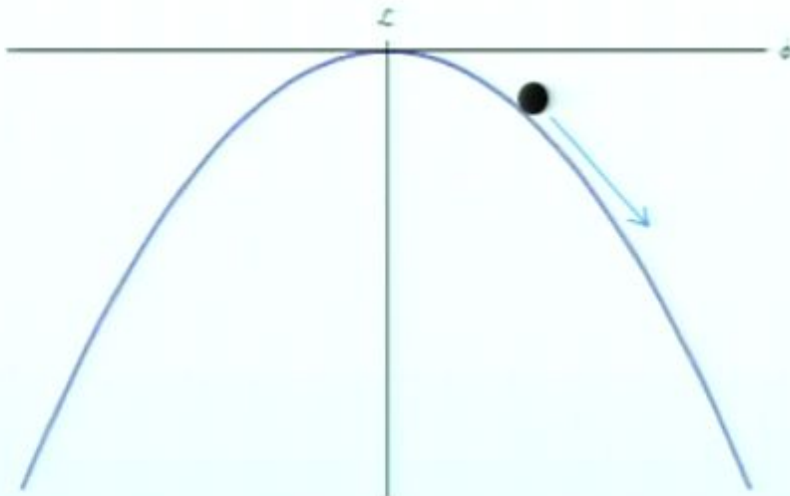
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Ghost Condensation

$$L = -\frac{1}{2} \partial^\mu \phi \partial_\mu \phi + \dots$$

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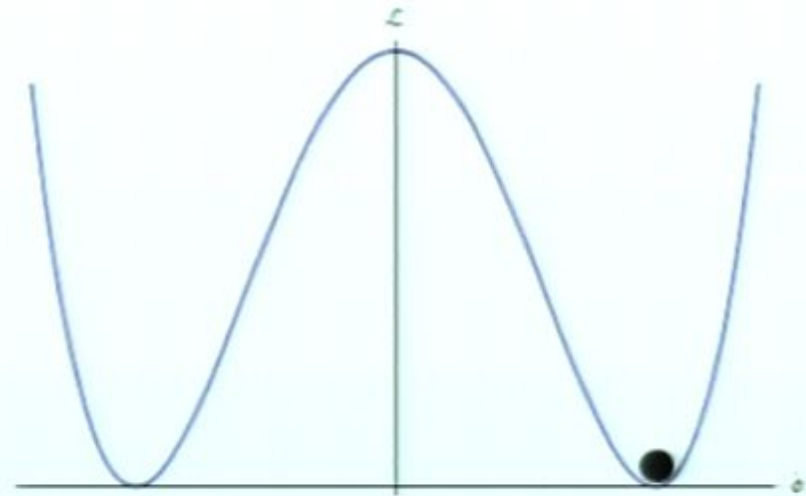
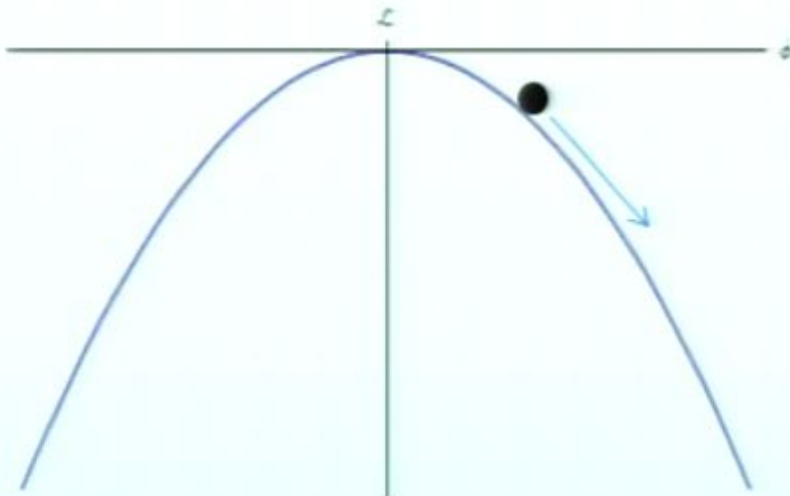
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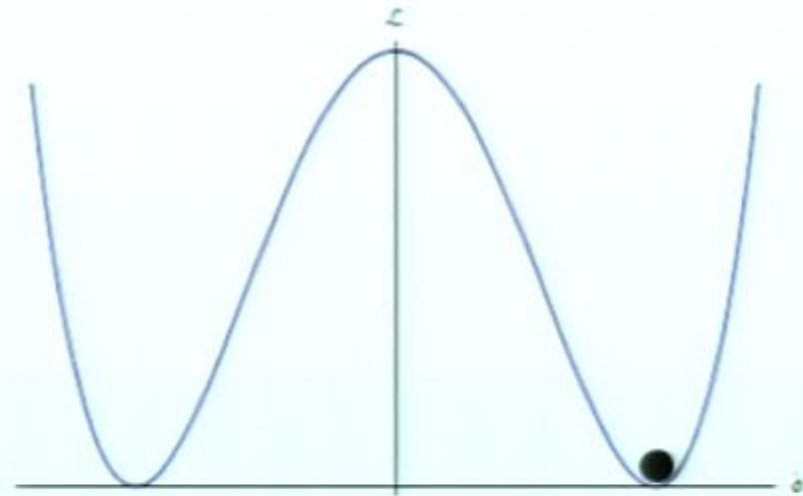
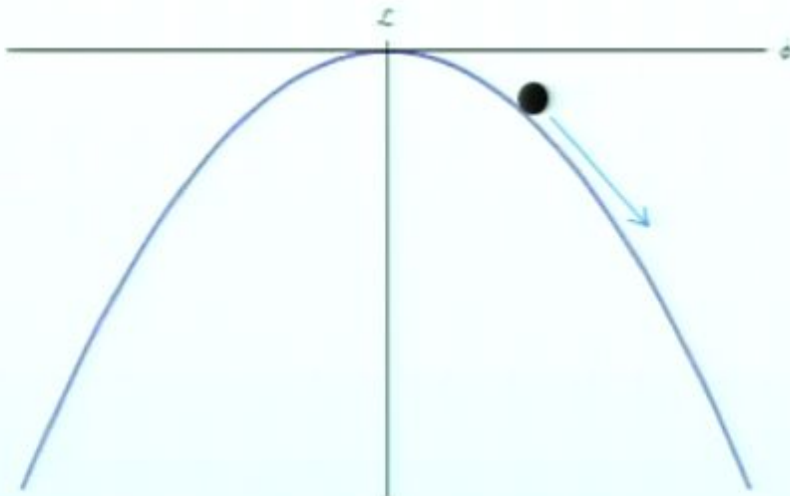
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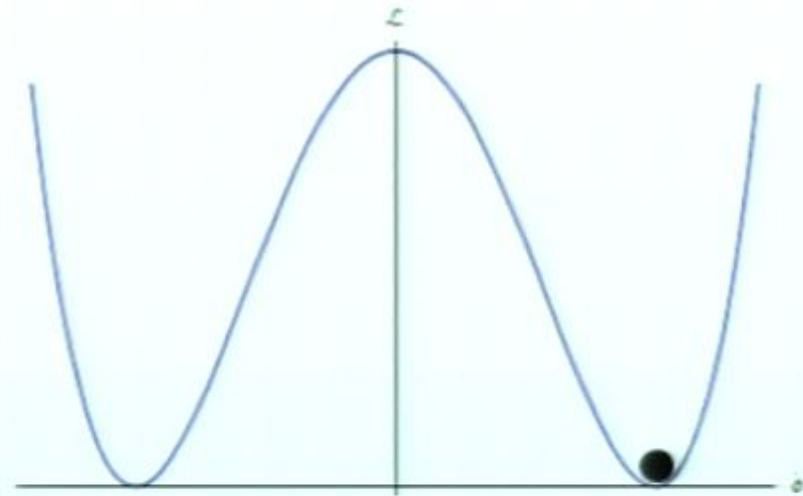
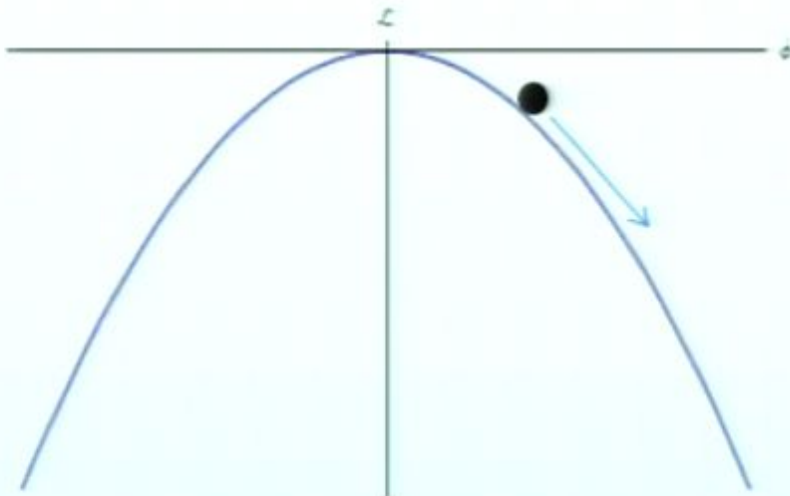
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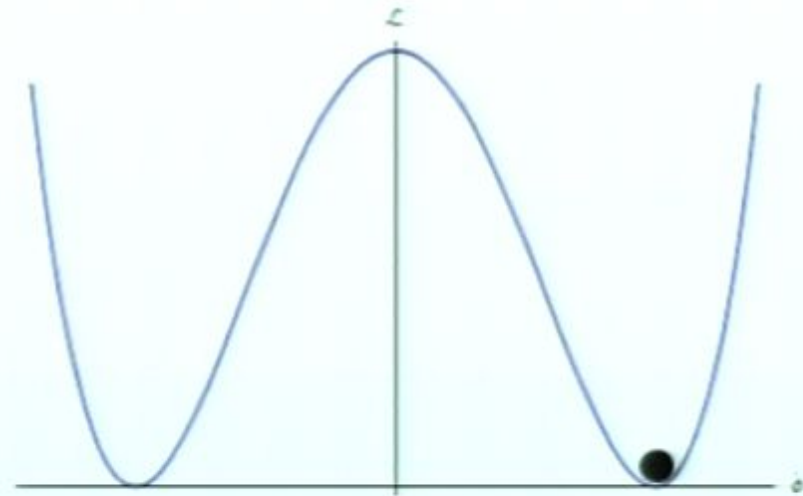
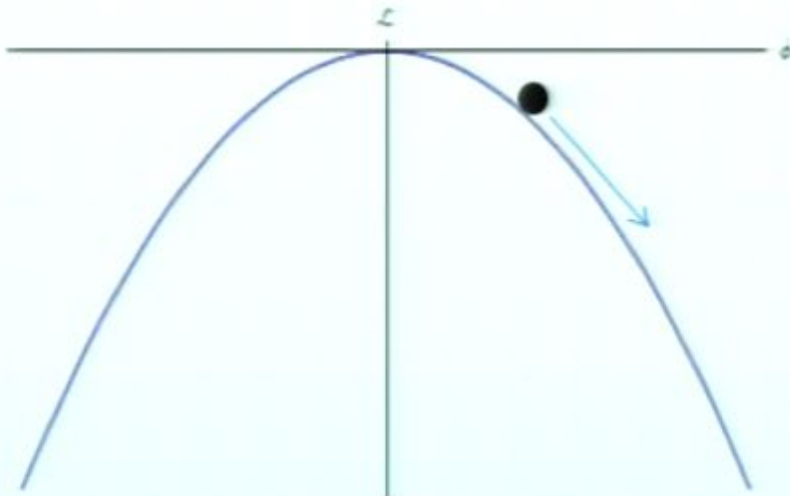
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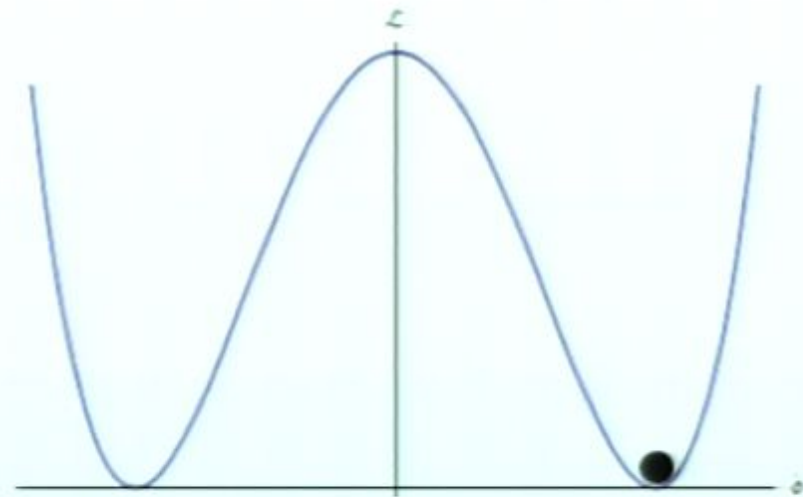
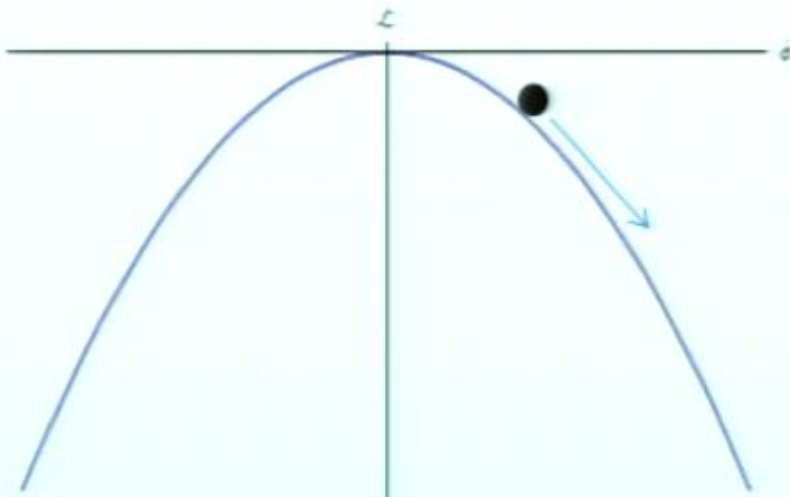
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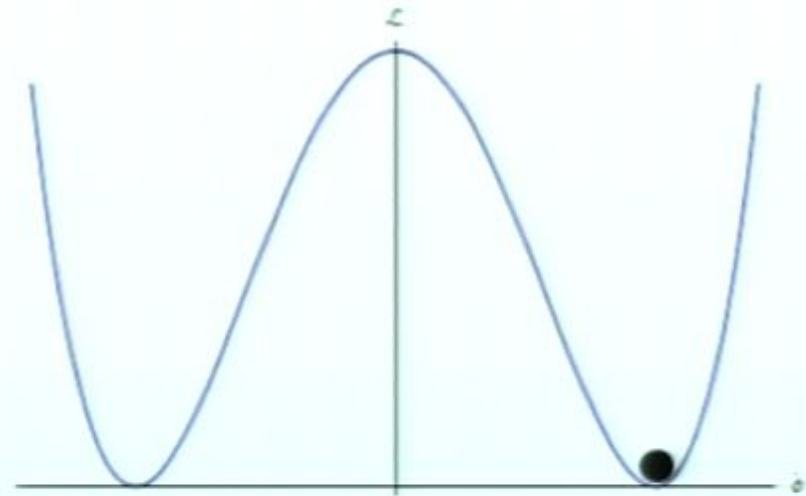
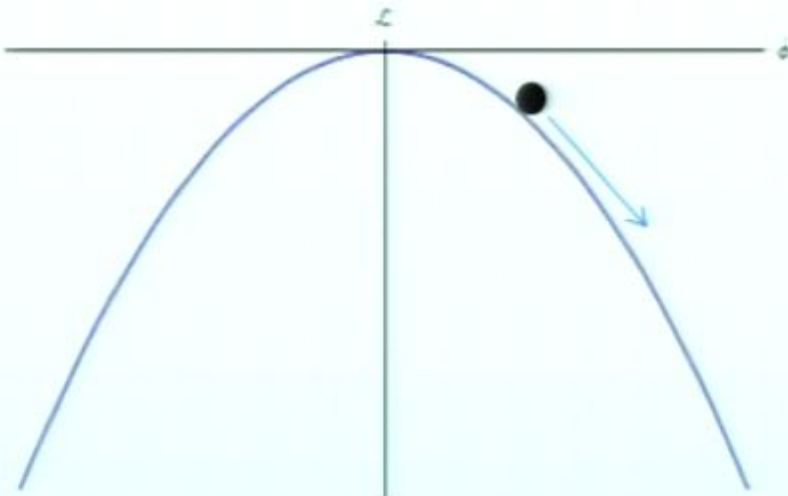
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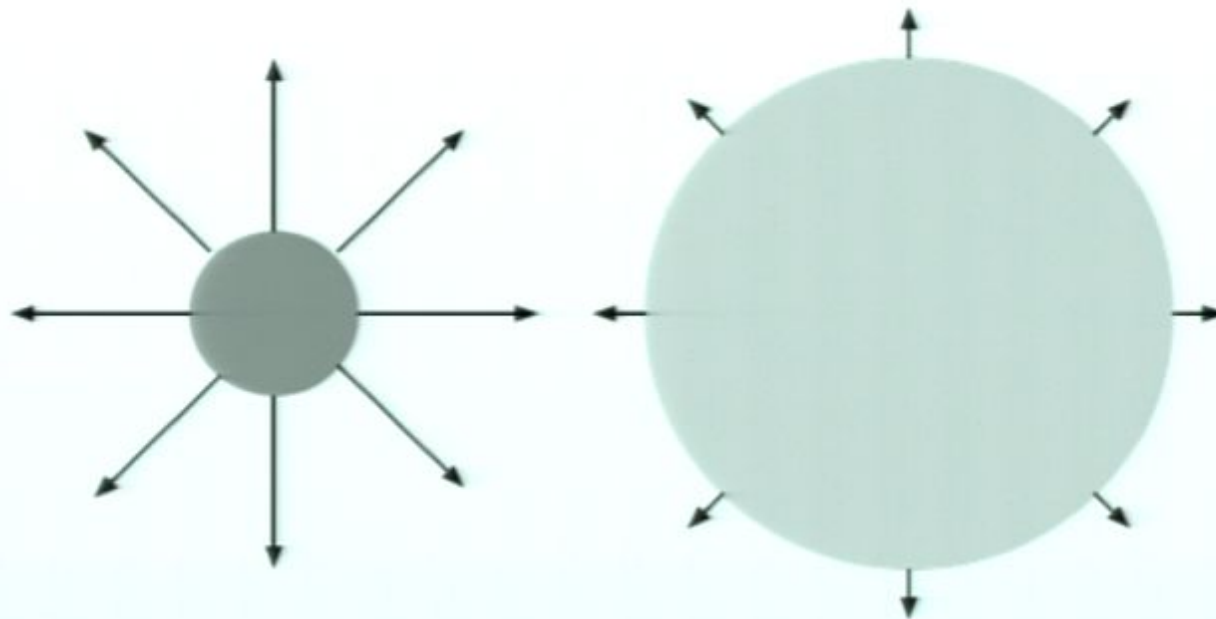
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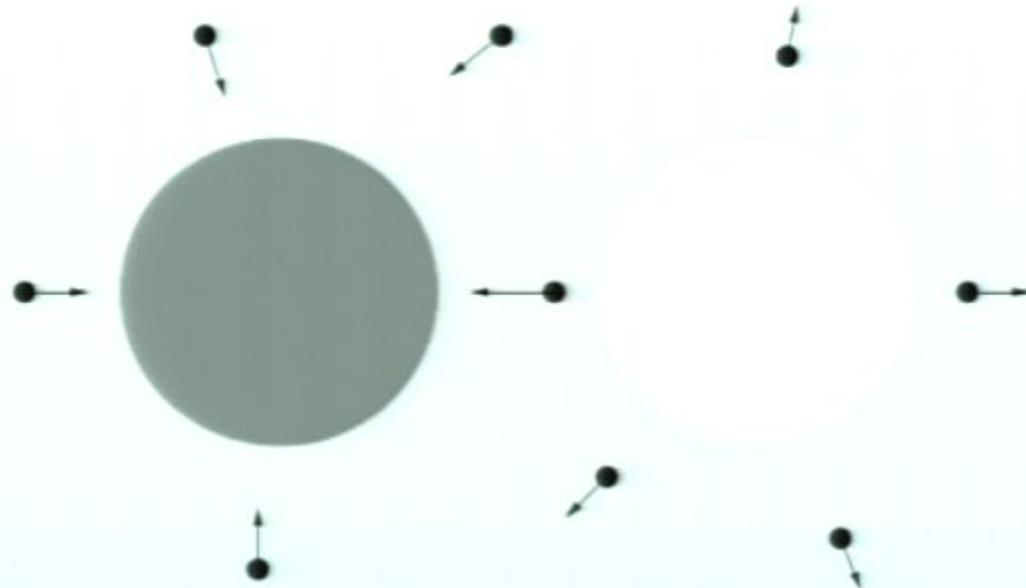
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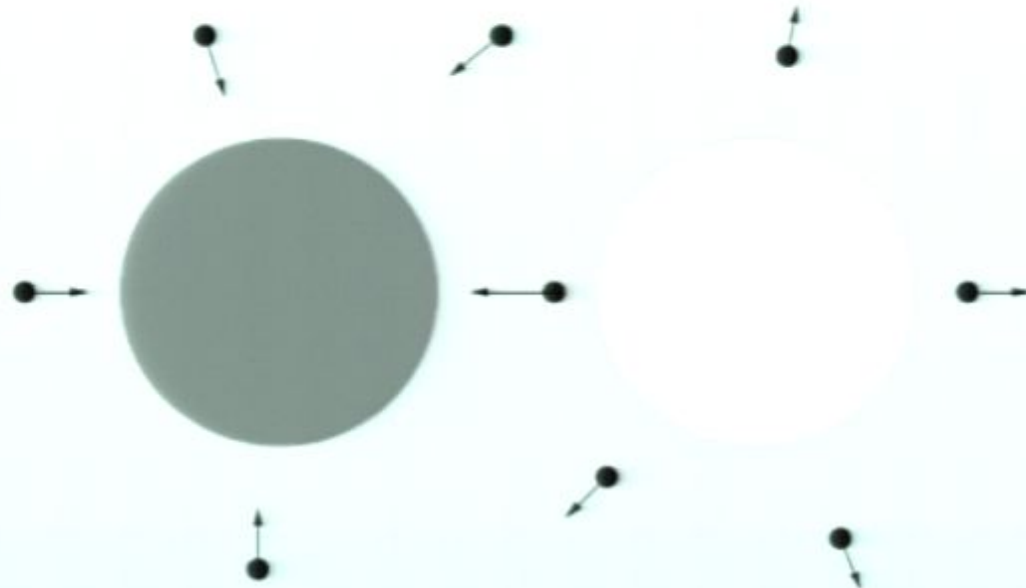
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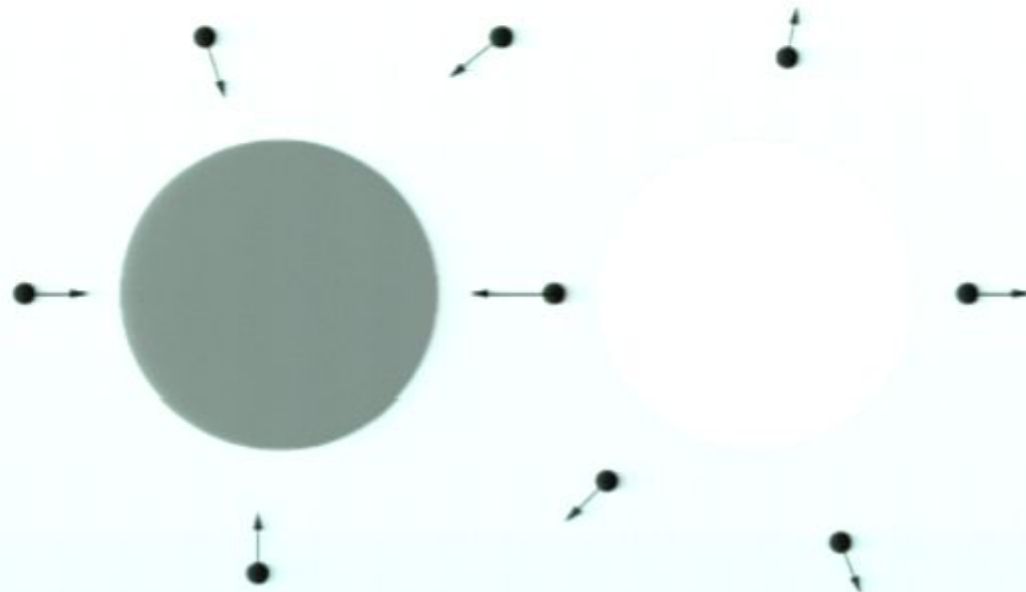
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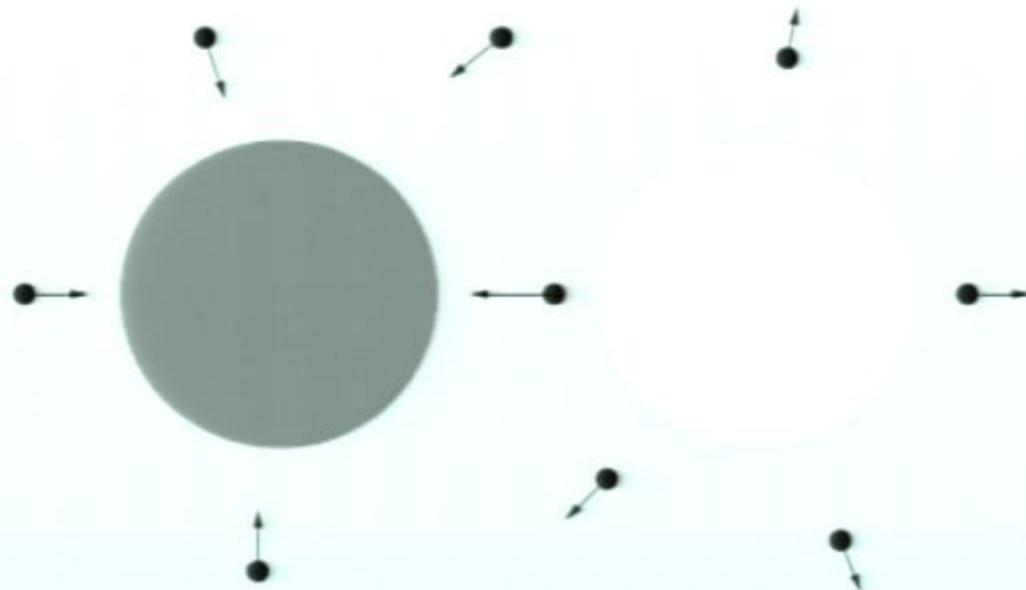


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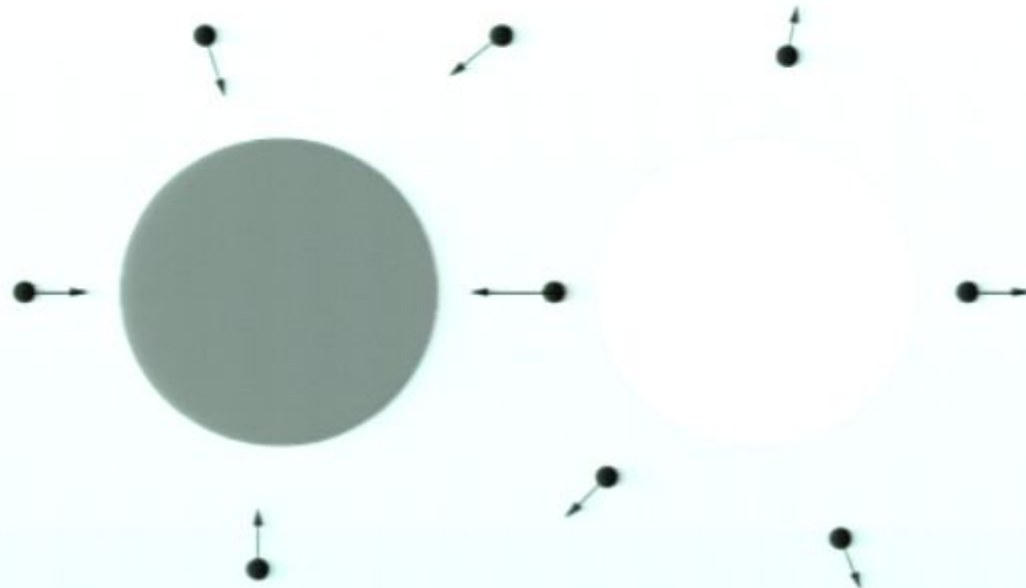


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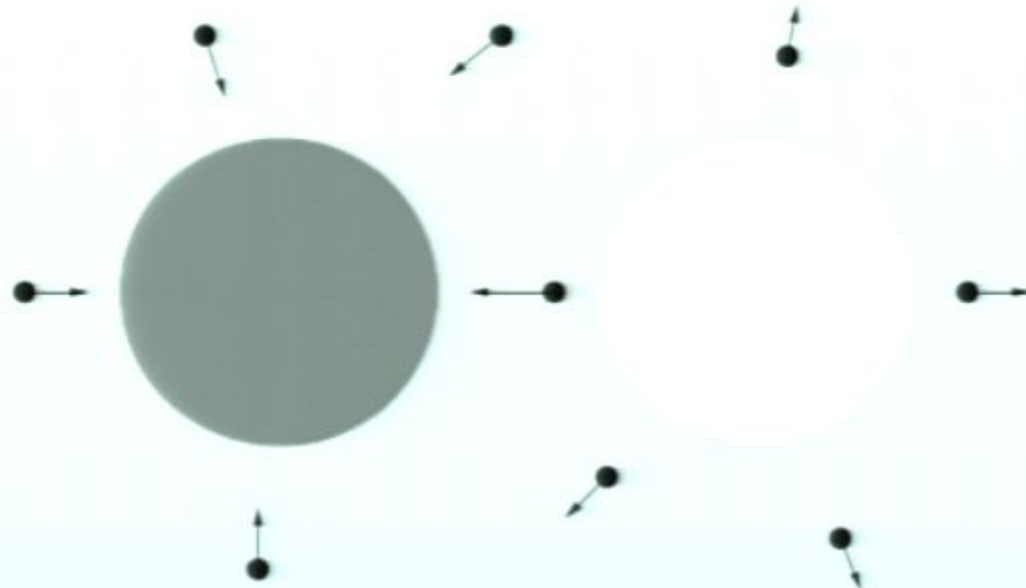


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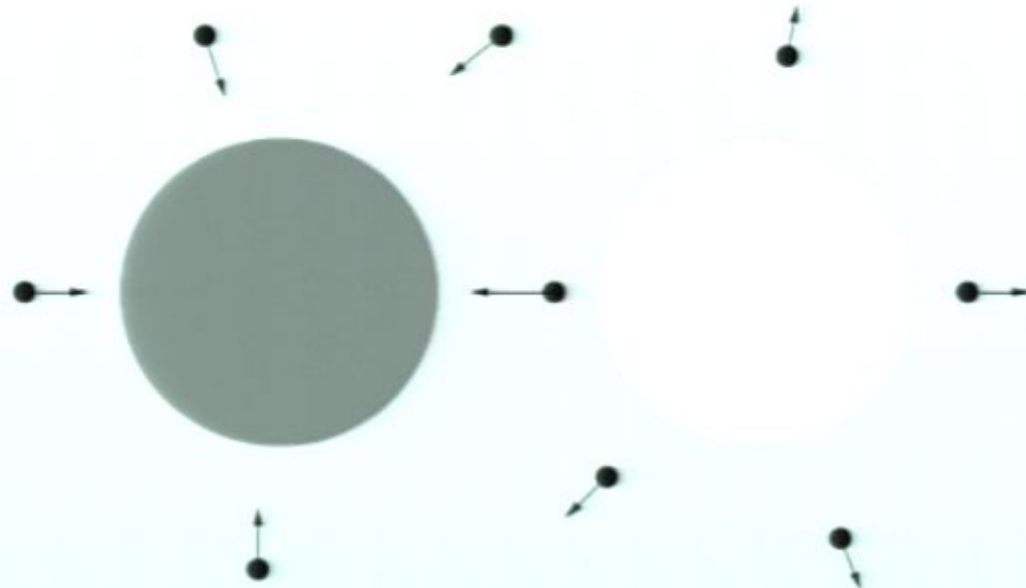
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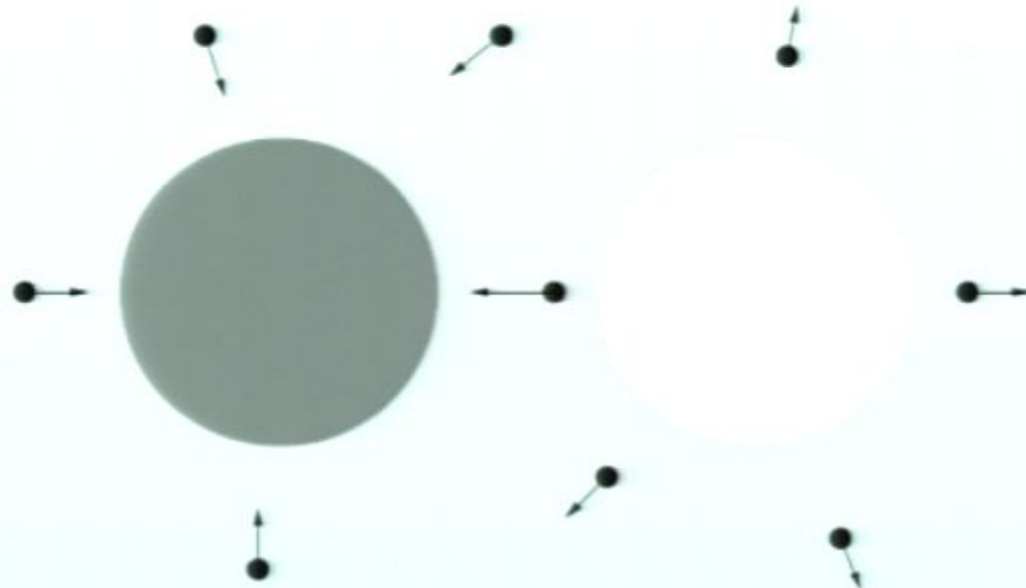


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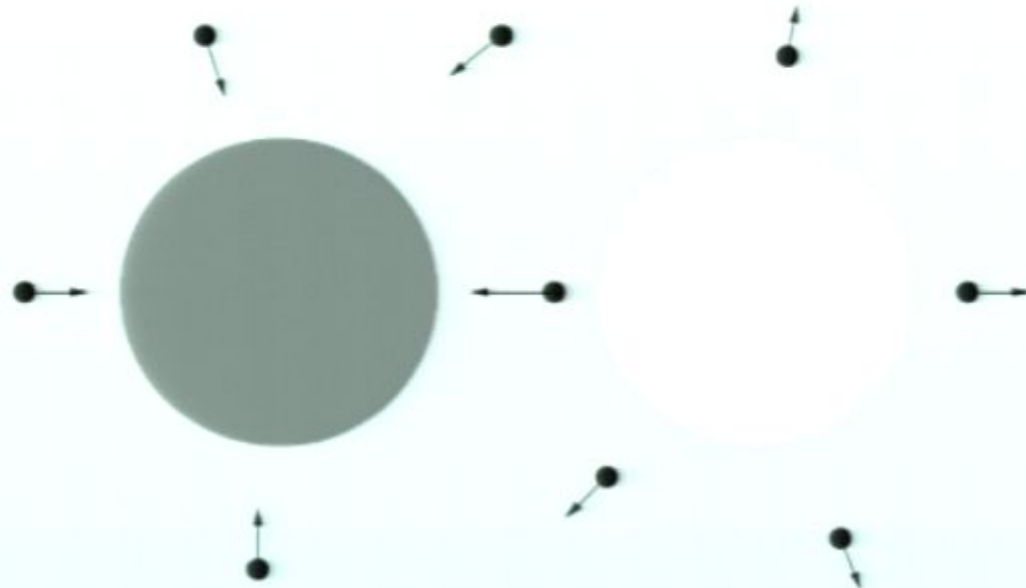
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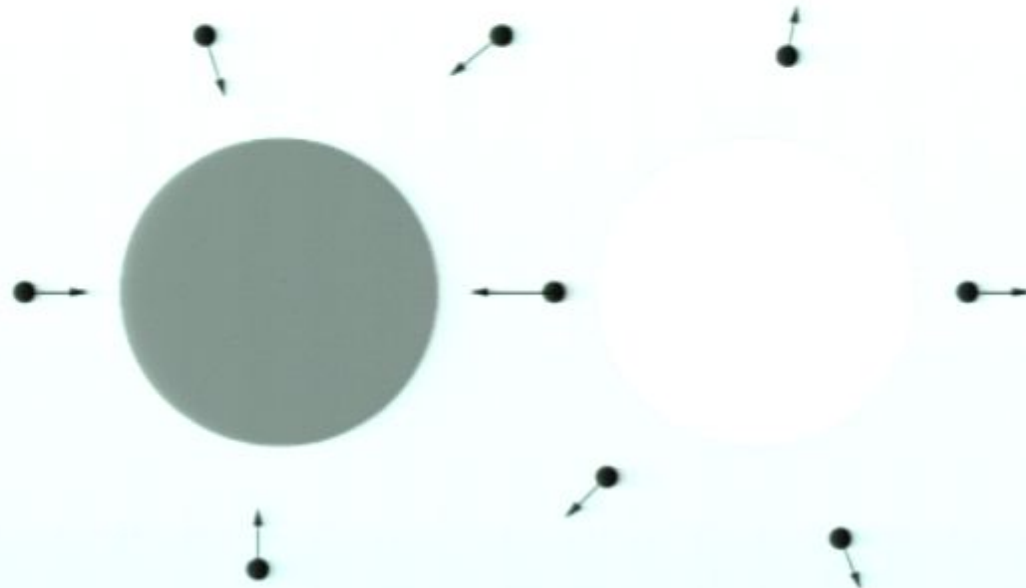


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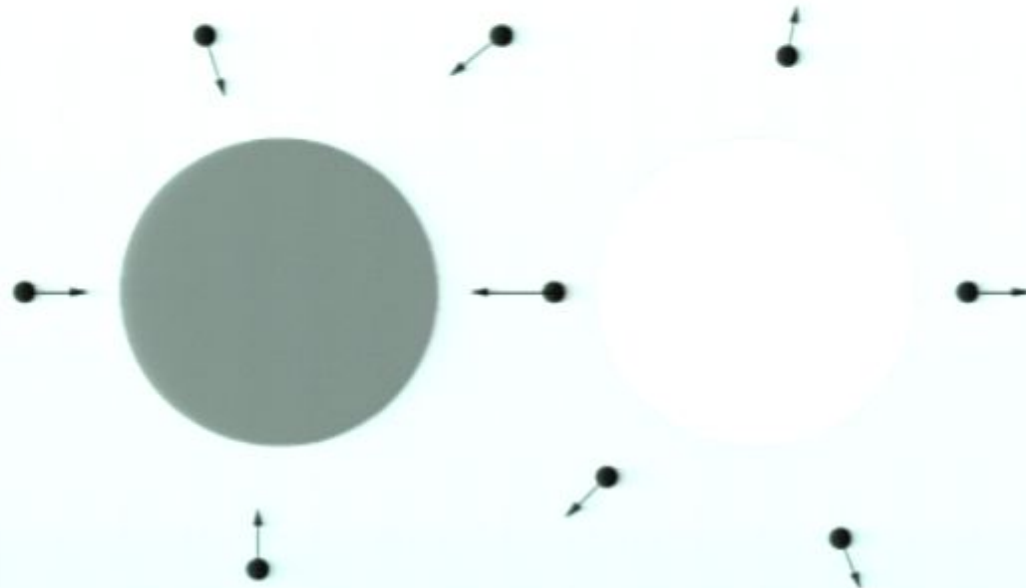


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Light lensing



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delay

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Light lensing



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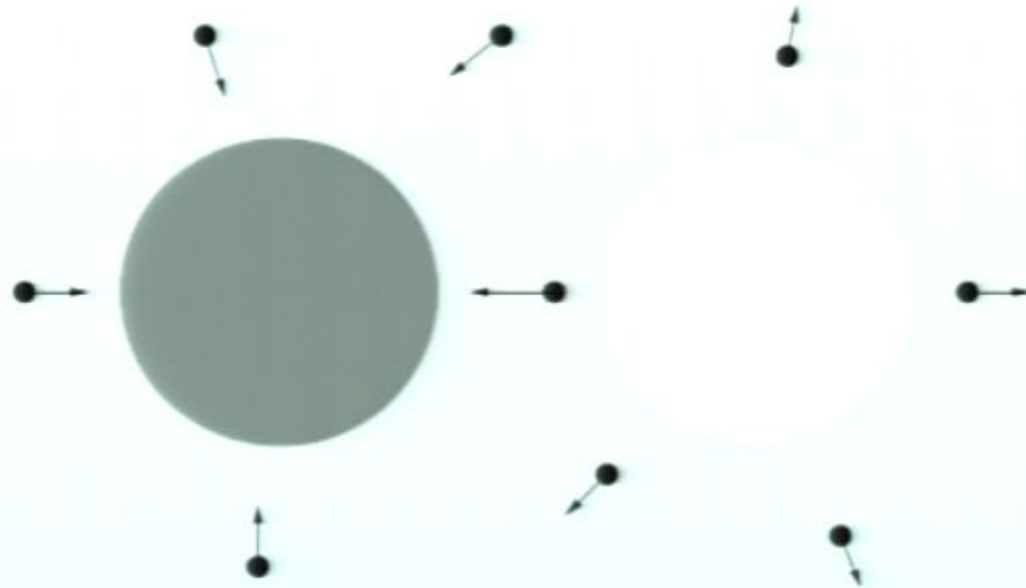
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Interesting features



Lumps come from scalar excitation, its energy density always positive in terms of “particle physics”, but the induced gravity can be either attractive or repulsive!

$\dot{\pi} > 0$ attractive

$\dot{\pi} < 0$ repulsive

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Ghost condensation locates at the minima of Lagrangian

$$P' = 0$$

$P' + 2P''\dot{\phi}^2 > 0$ is ghost free condition, so we get

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Ghost condensation stabilize vacuum on background and perturbation level !

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Part III

Realization of Matter Bounce

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Ghost bounce

- Matter sector + ghost condensation

$$\rho_m(t) \sim a(t)^{-3(1+w_m)} \quad \rho_X \sim a(t)^{-p}$$

minimal requirement $p > 3$

against radiation $p > 4$

against anisotropic stress $p > 6$

- Lagrangian of GC takes the following general form

$$\mathcal{L} = M^4 P(X) - V(\phi)$$

$P(X)$ takes the prototypical form

$$P(X) = \frac{1}{8}(X - c^2)^2$$

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$$P(X) = \frac{1}{8}(X - c^2)^2$$

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- Matter sector + ghost condensation

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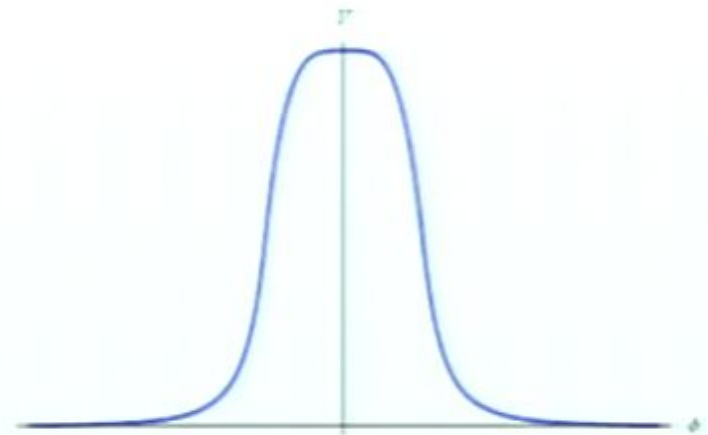
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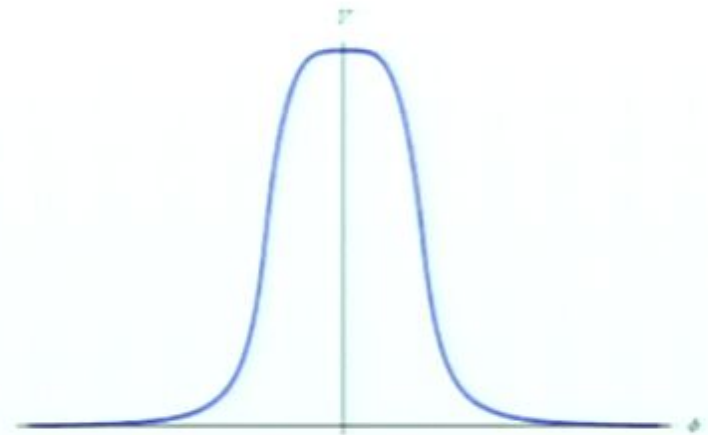
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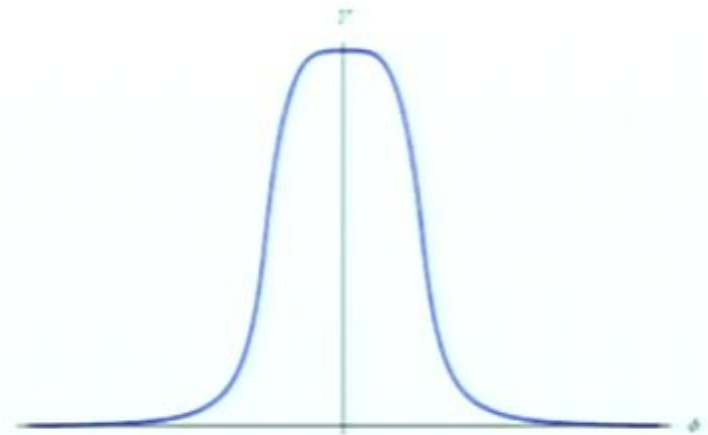
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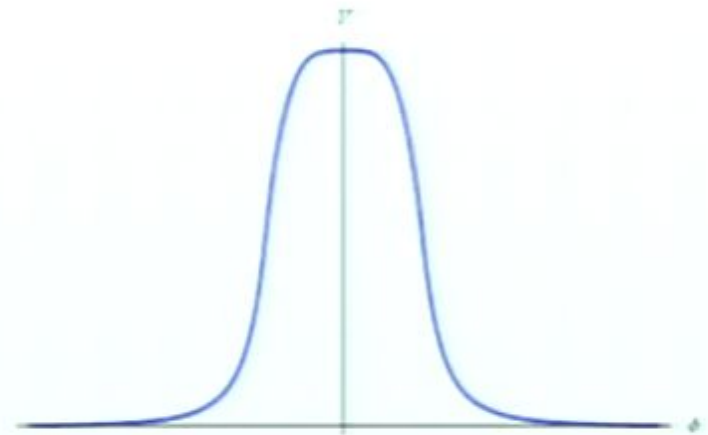
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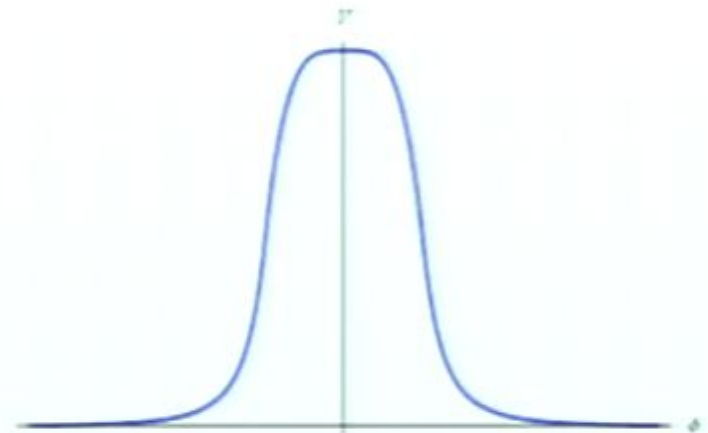
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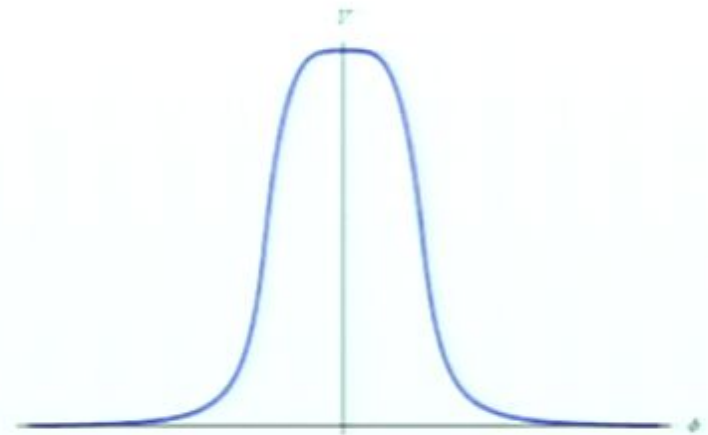
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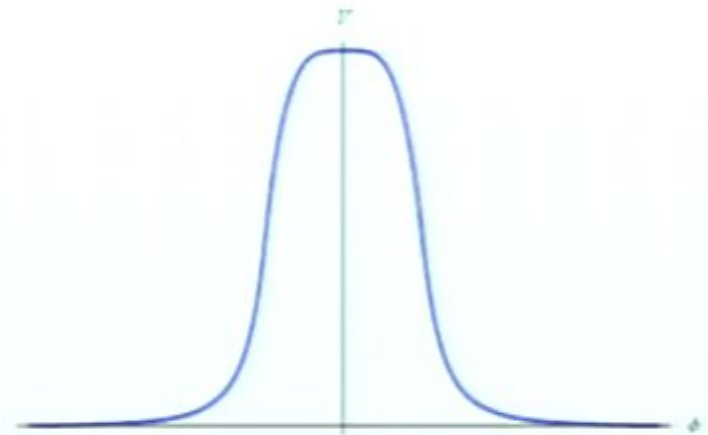
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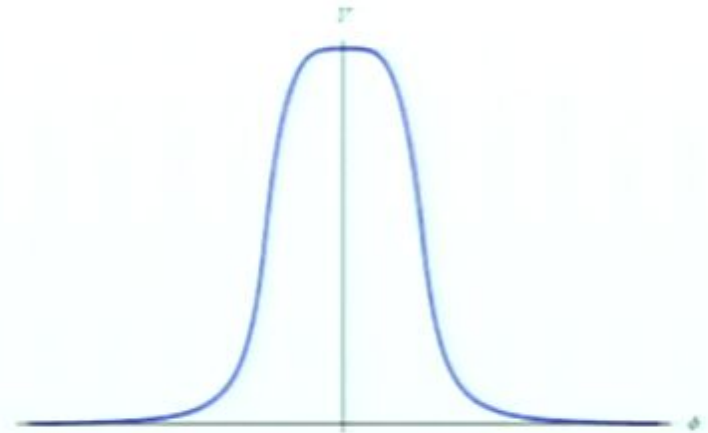
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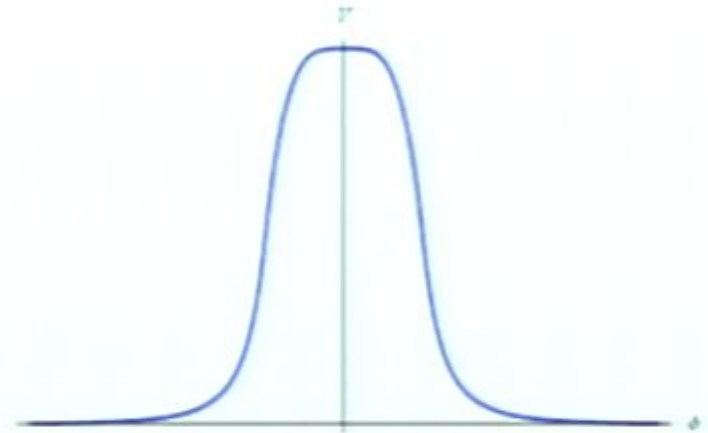
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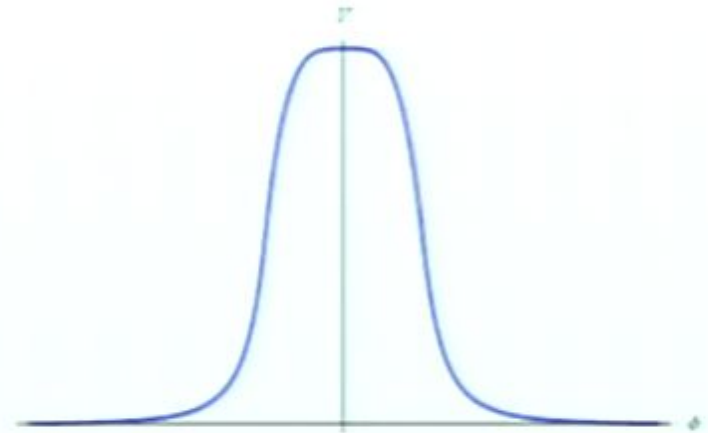
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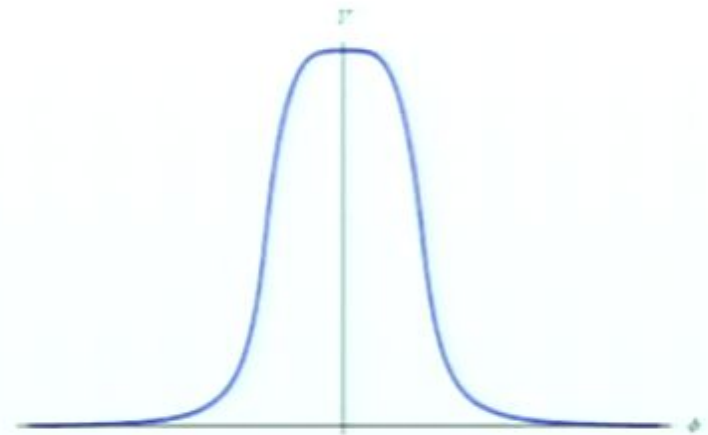
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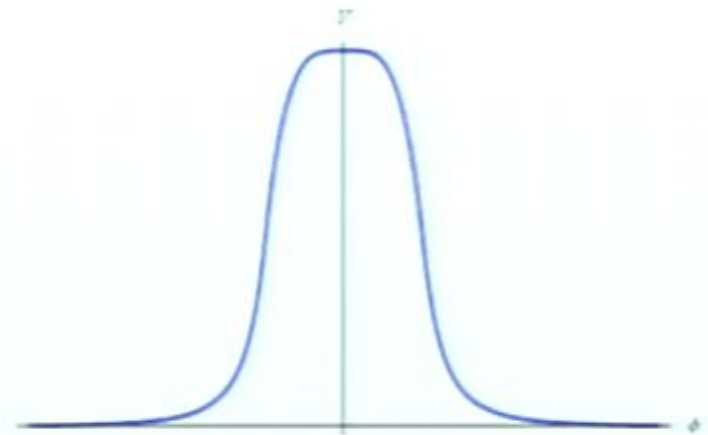
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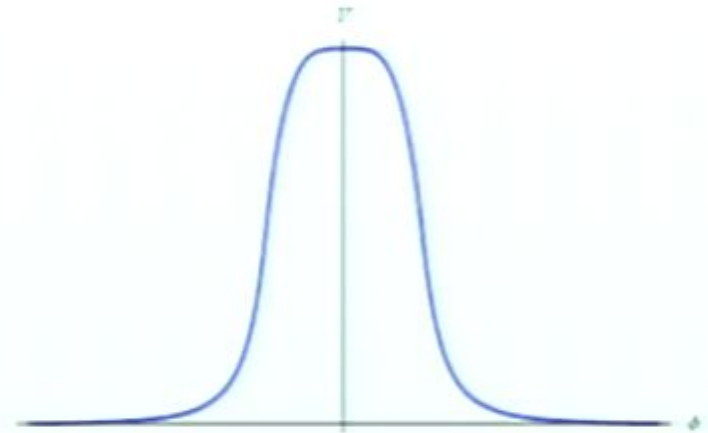
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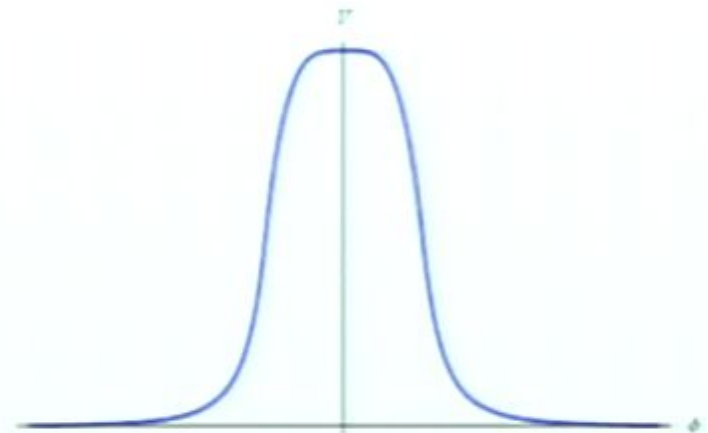
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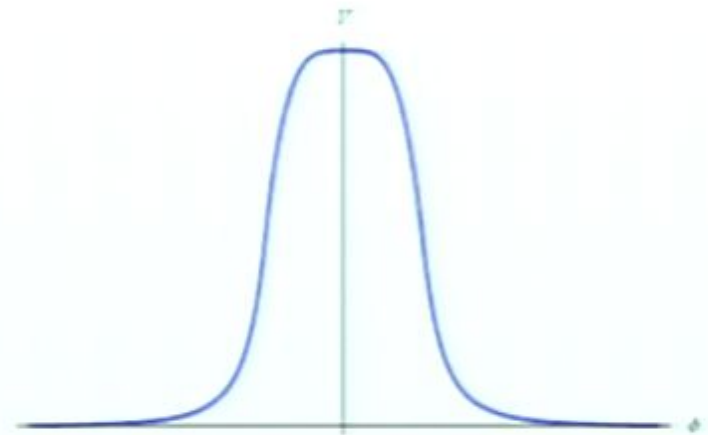
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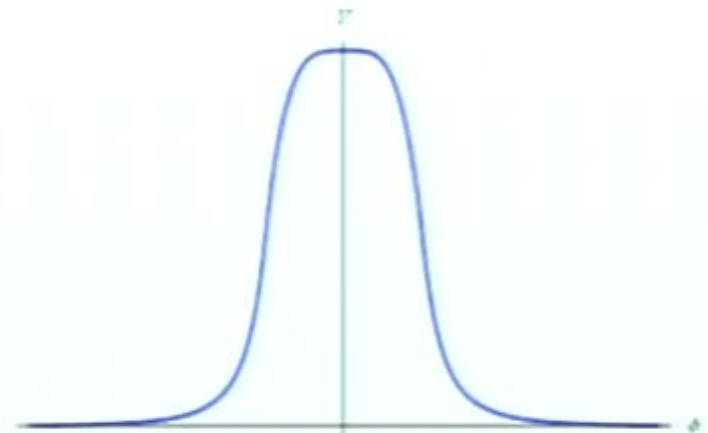
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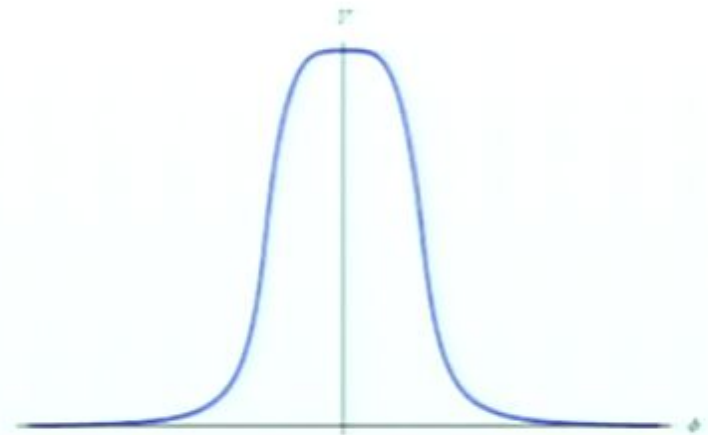
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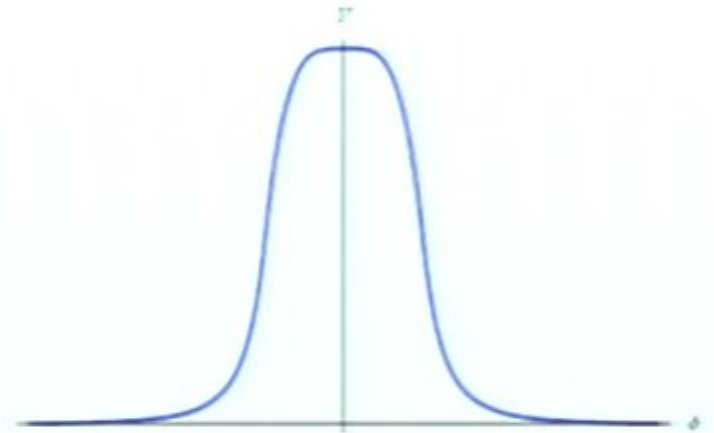
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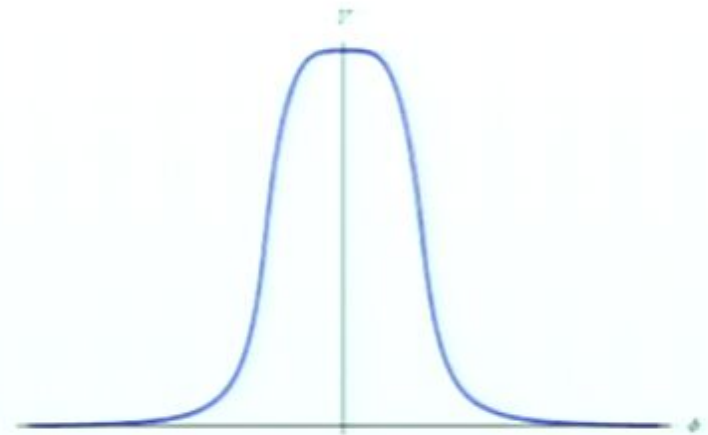
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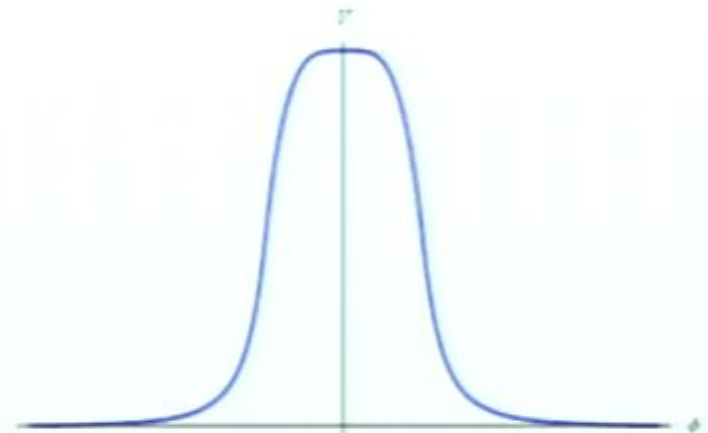
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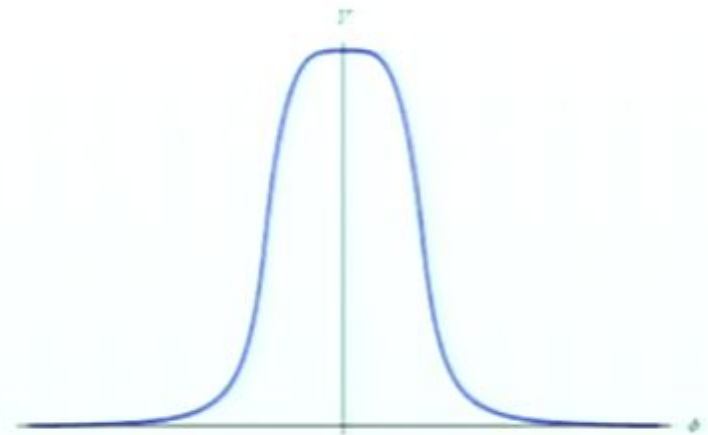
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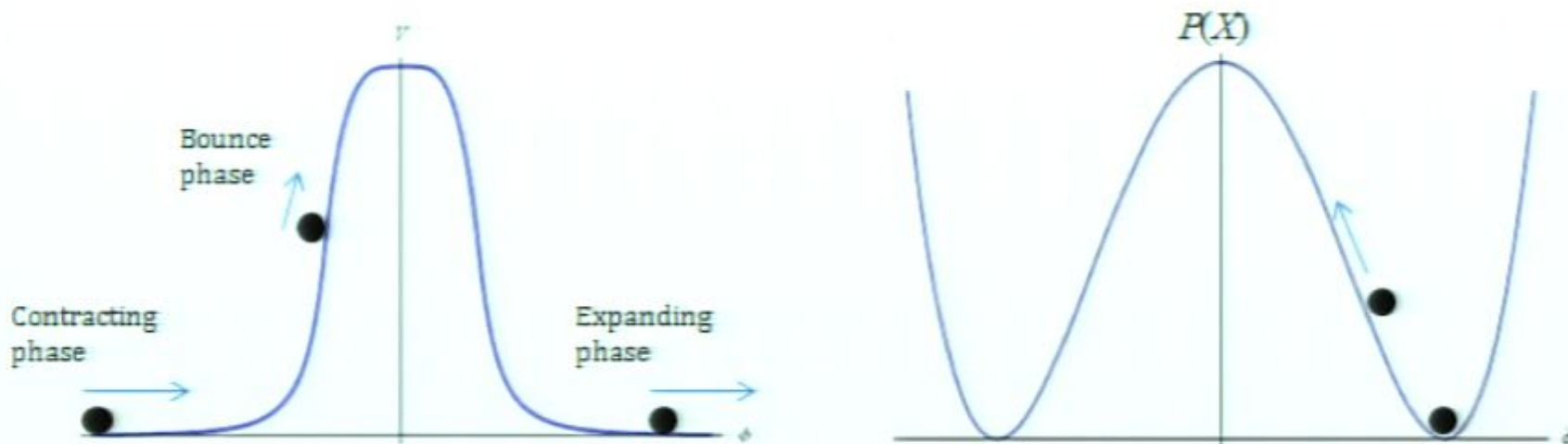
It yields $\rho_X \sim \dot{\pi} \sim t^{-\alpha}$.

$\alpha = 4$ Marginally stable against anisotropic stress

$\alpha = 6$ stable



Ghost bounce

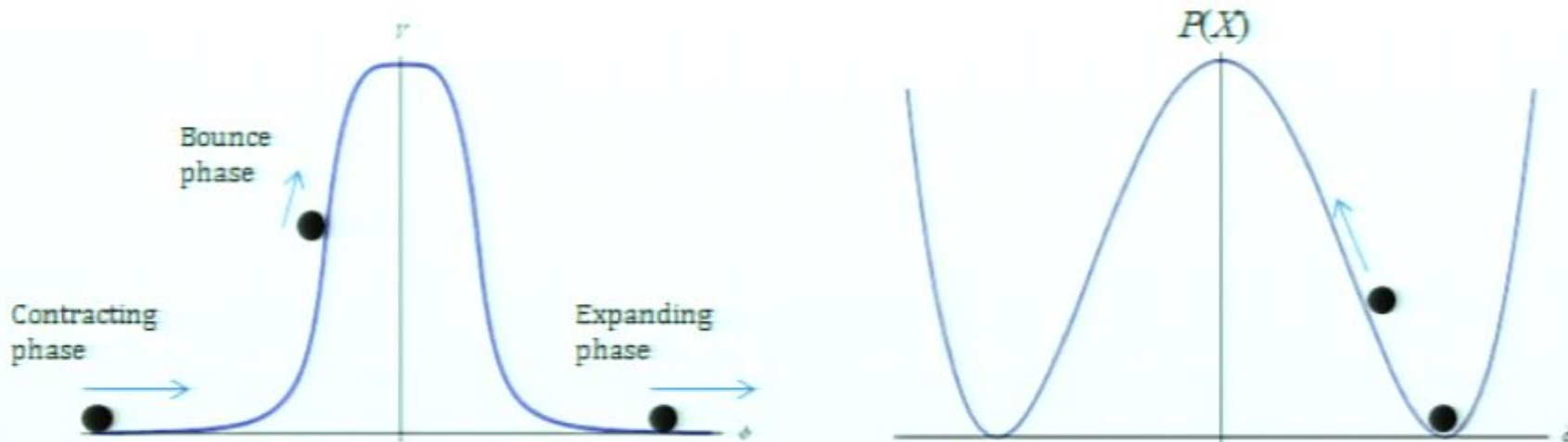


$$2M_p^2 \dot{H} = -2M^4 X P' - (1 + w_m) \rho_m$$

Realization of NEC violation!



Ghost bounce

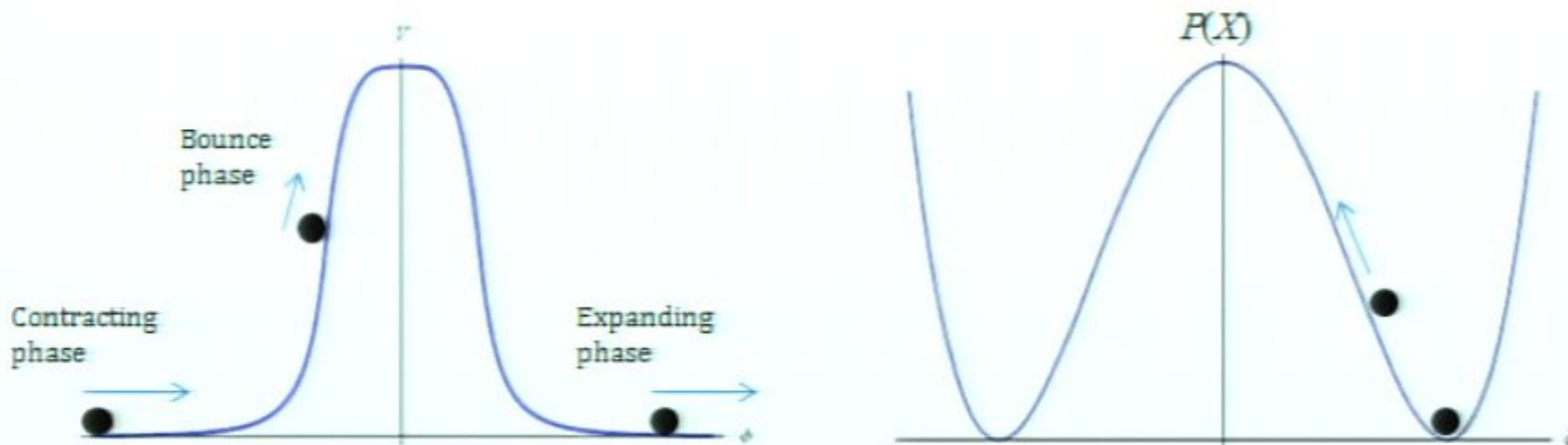


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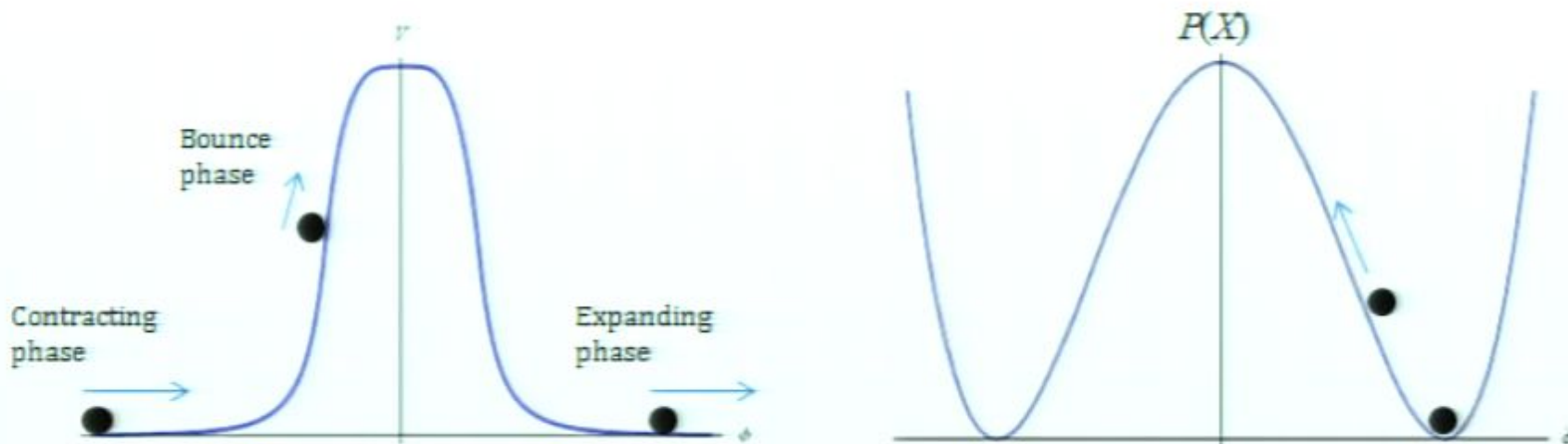


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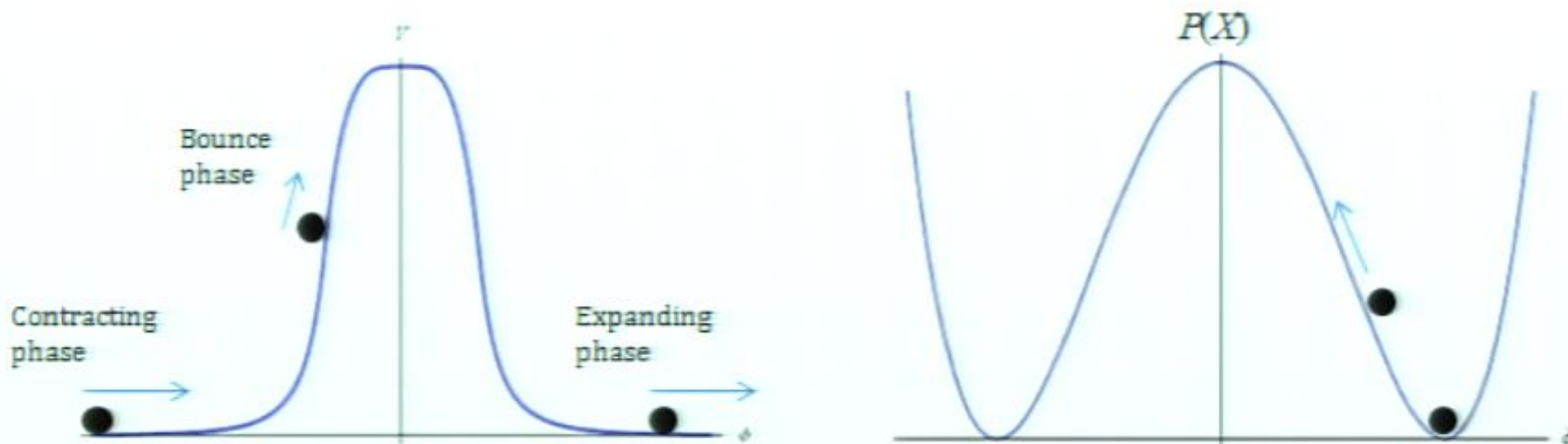


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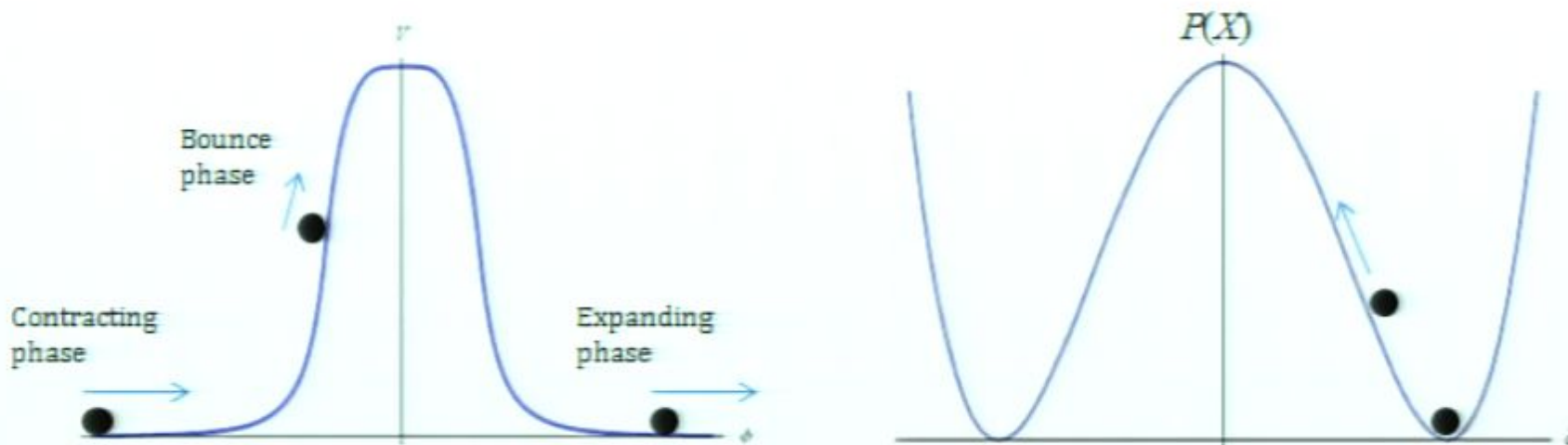


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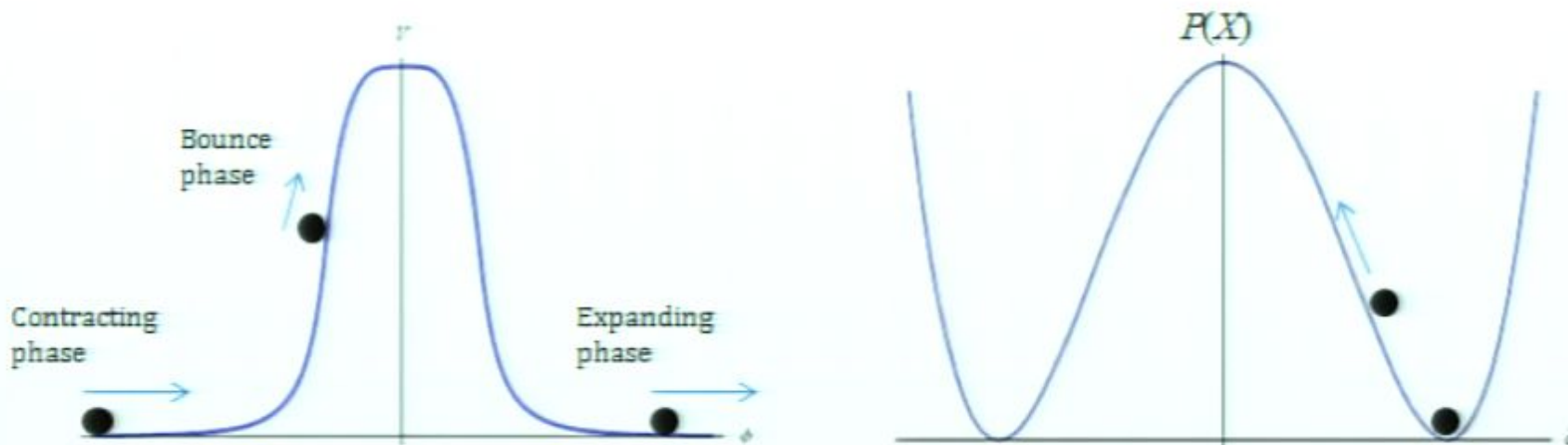


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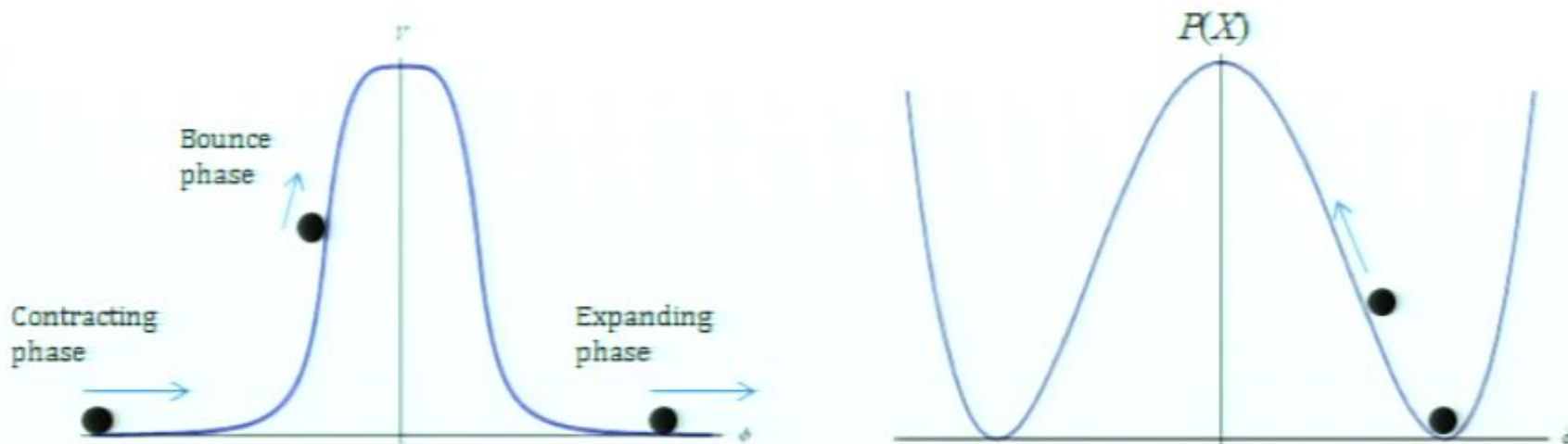


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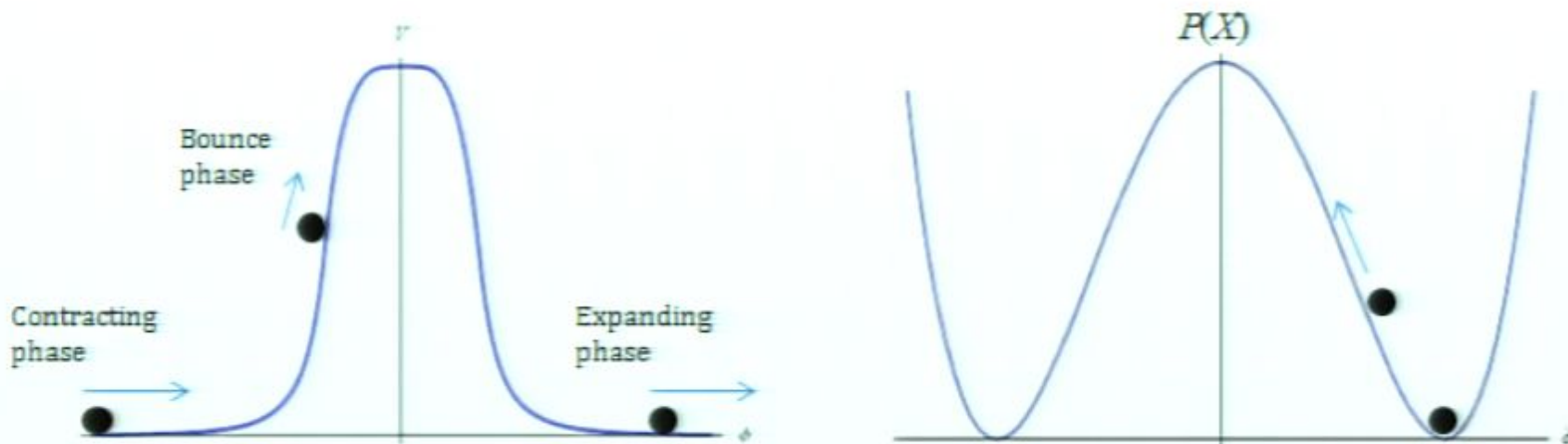


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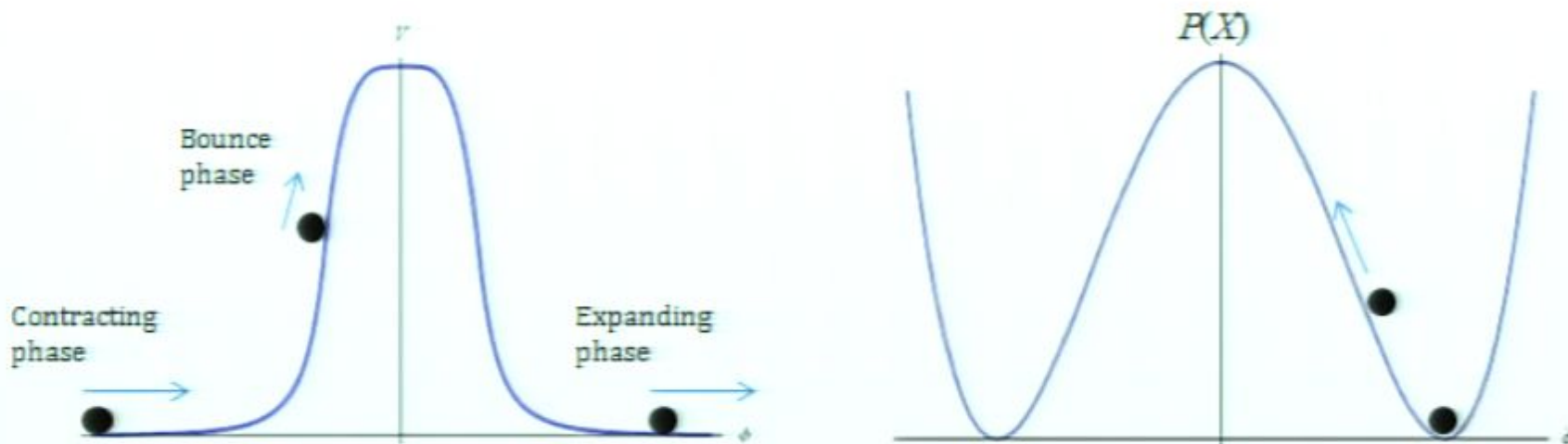


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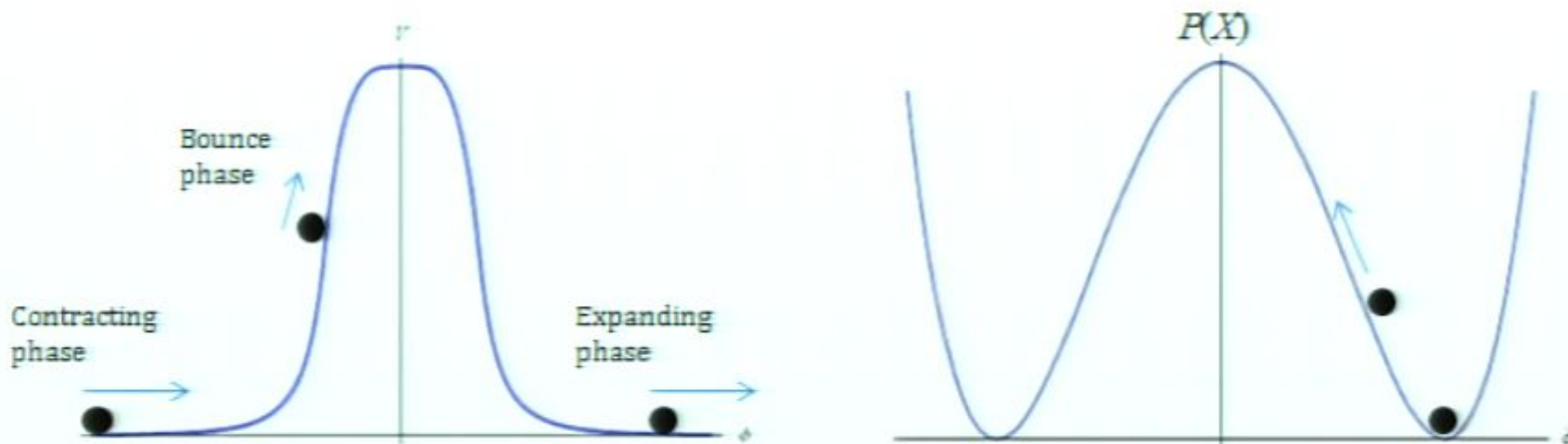


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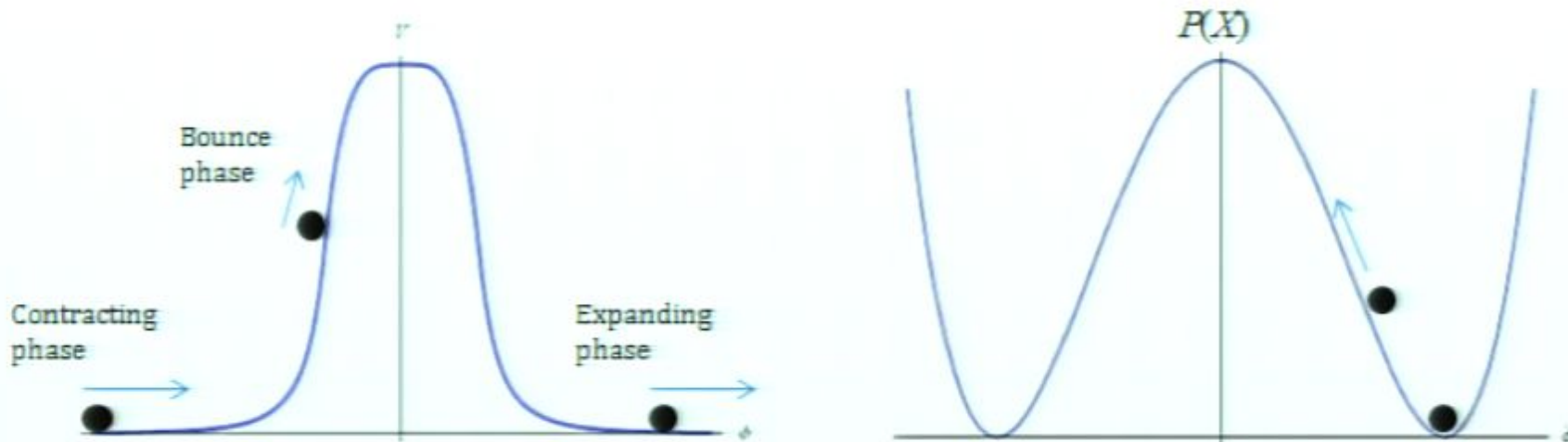


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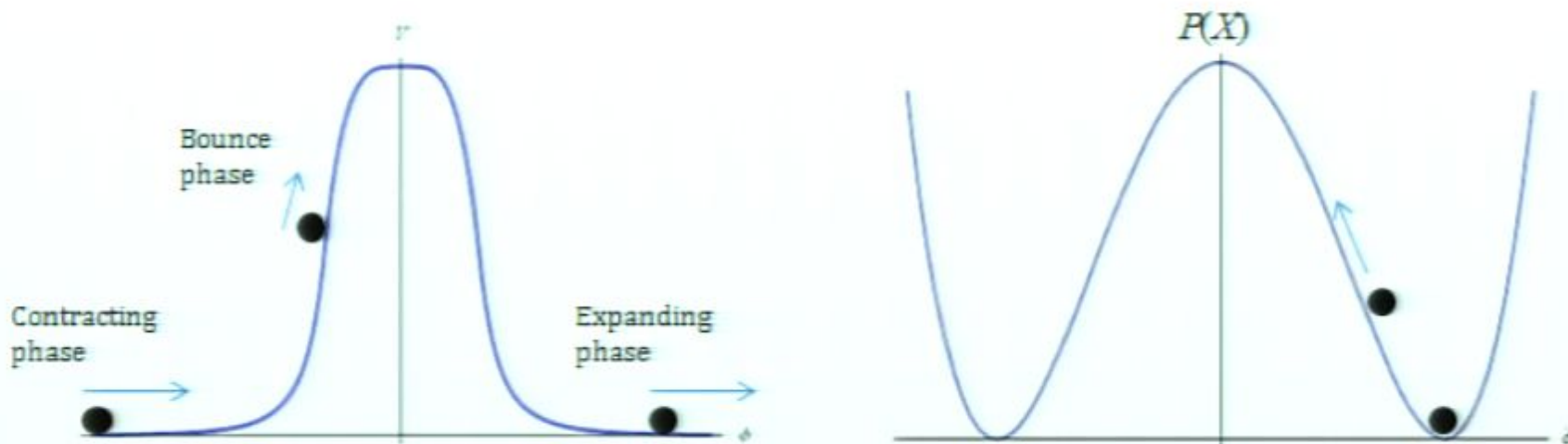


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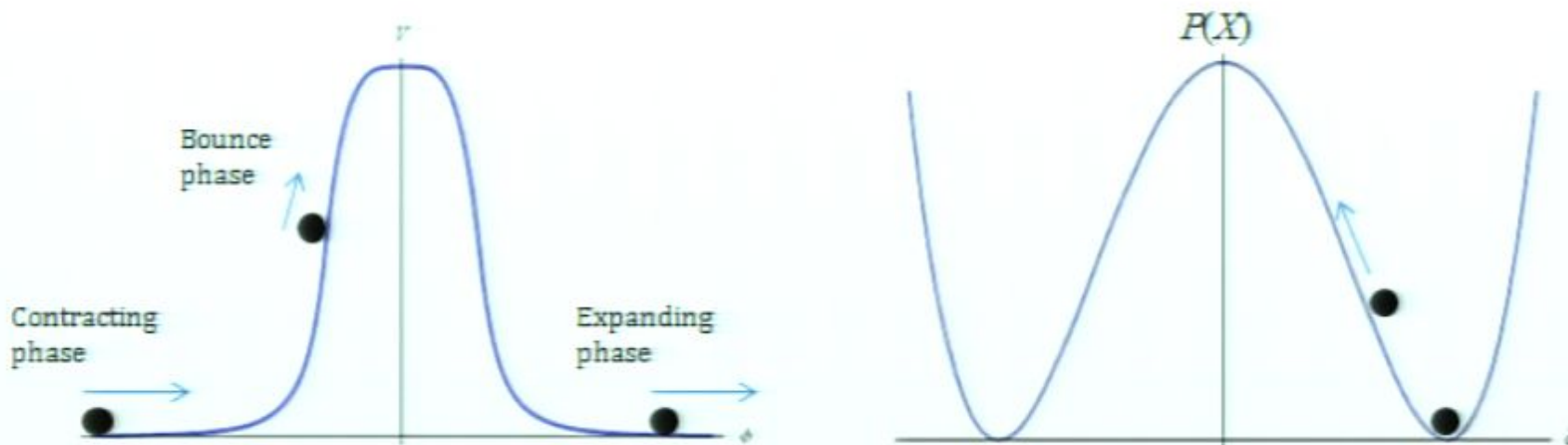


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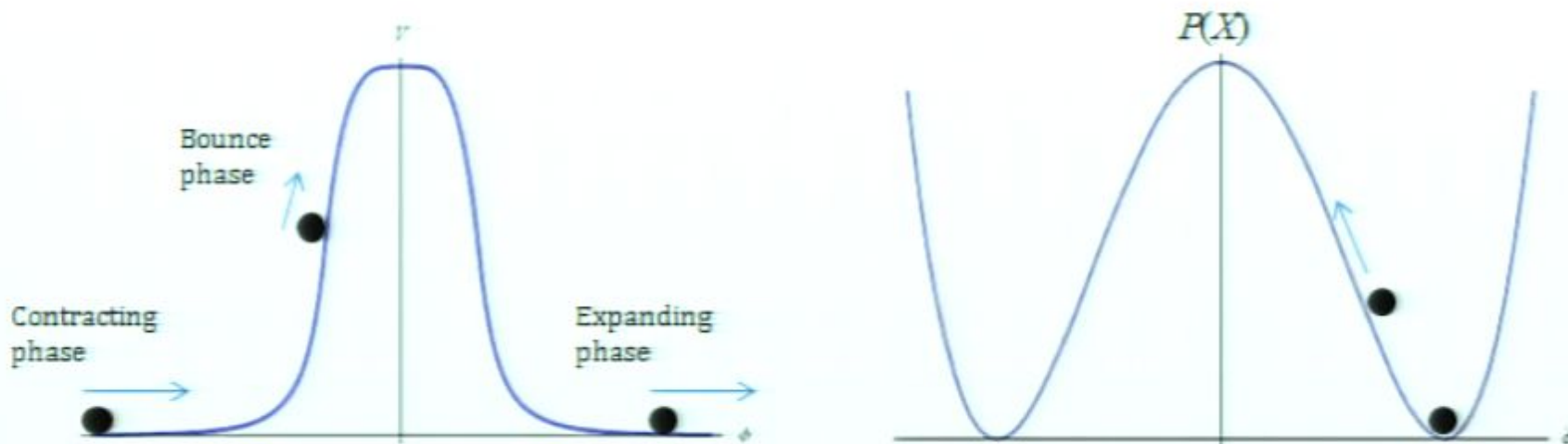


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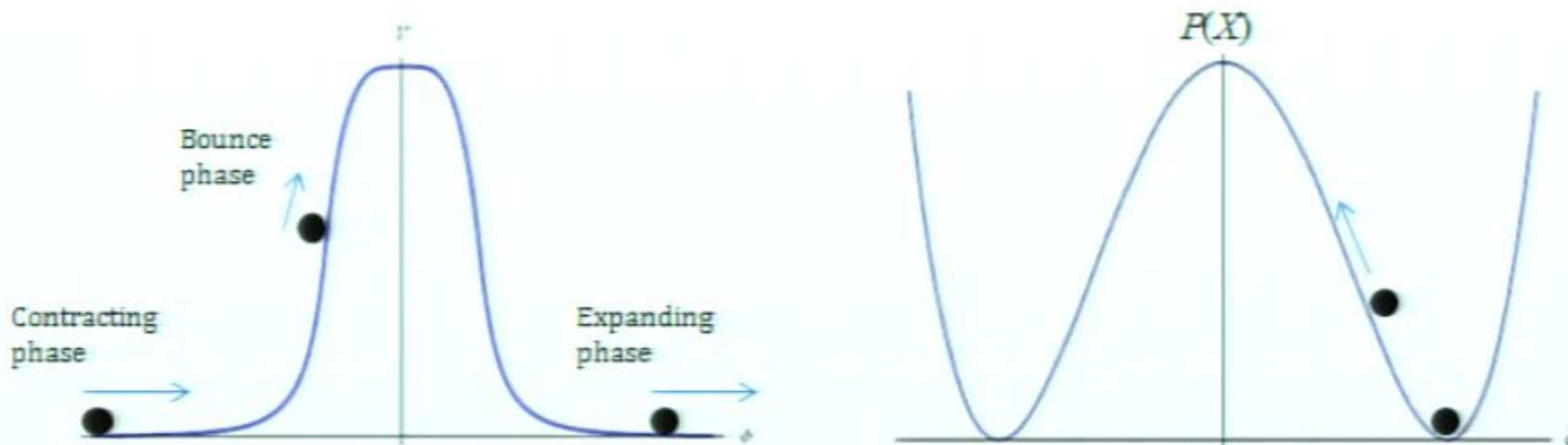


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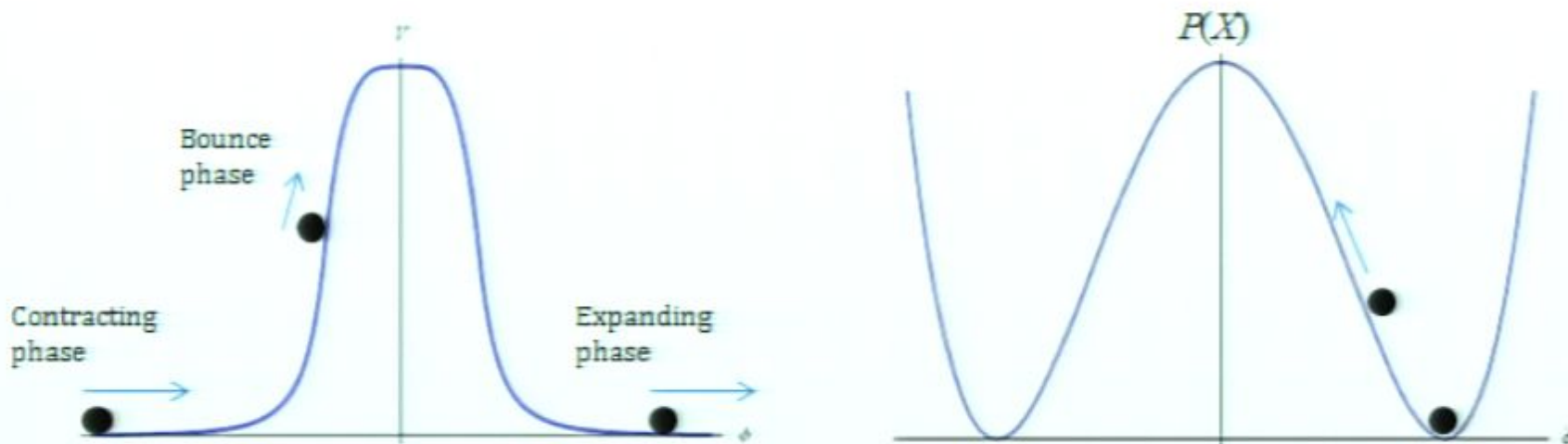


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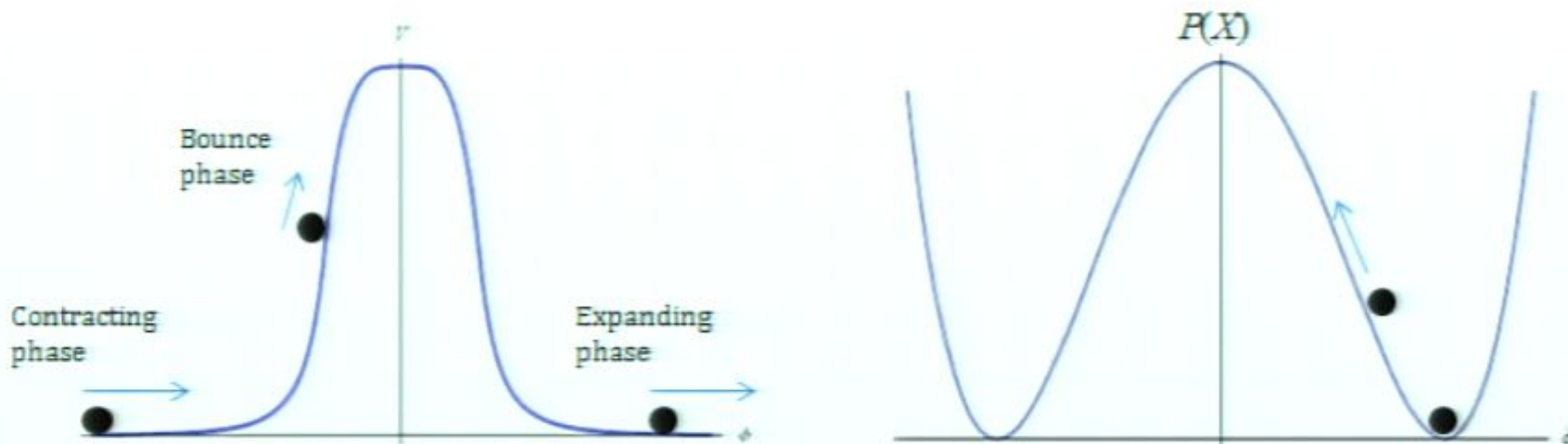


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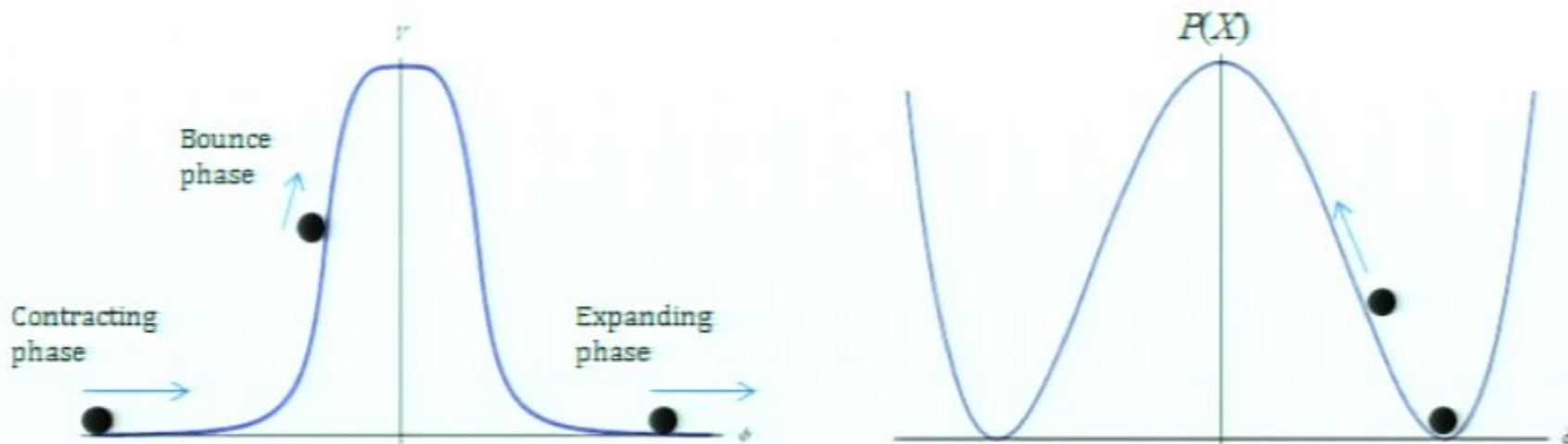


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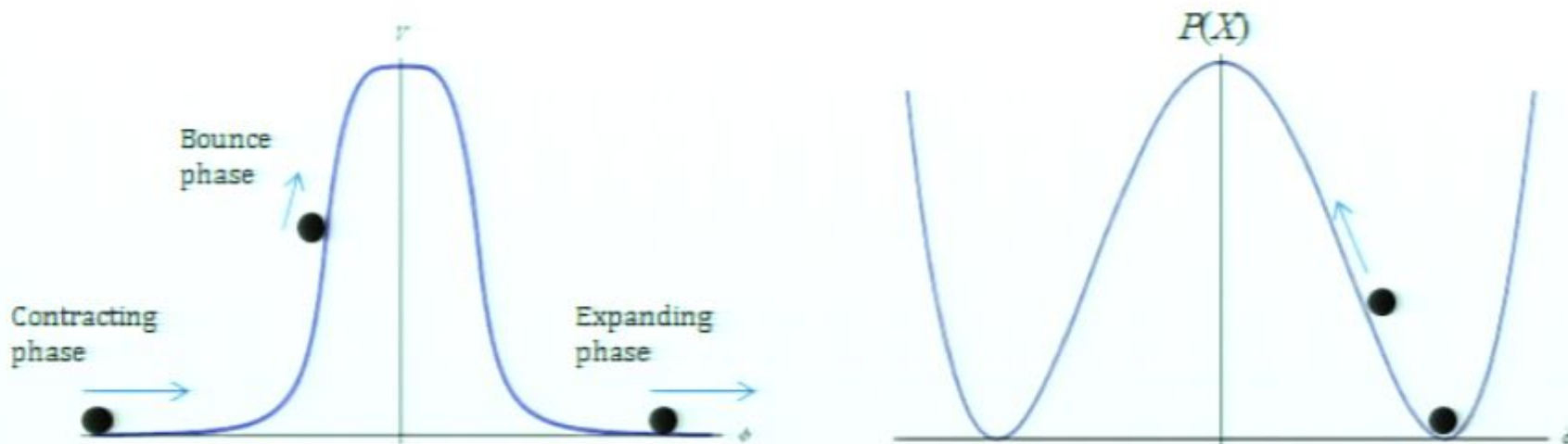


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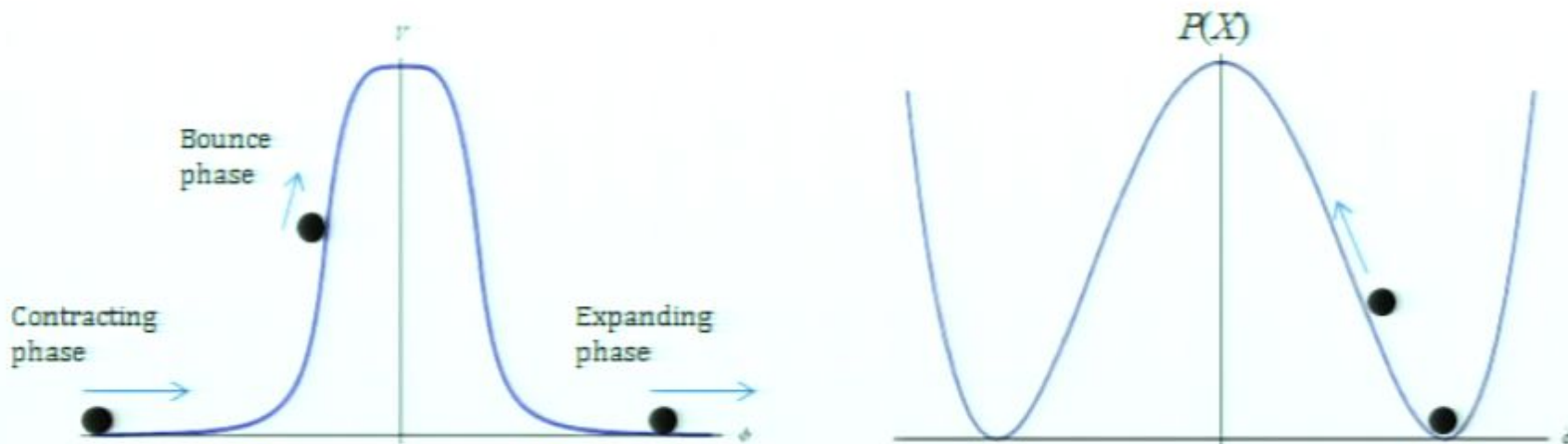


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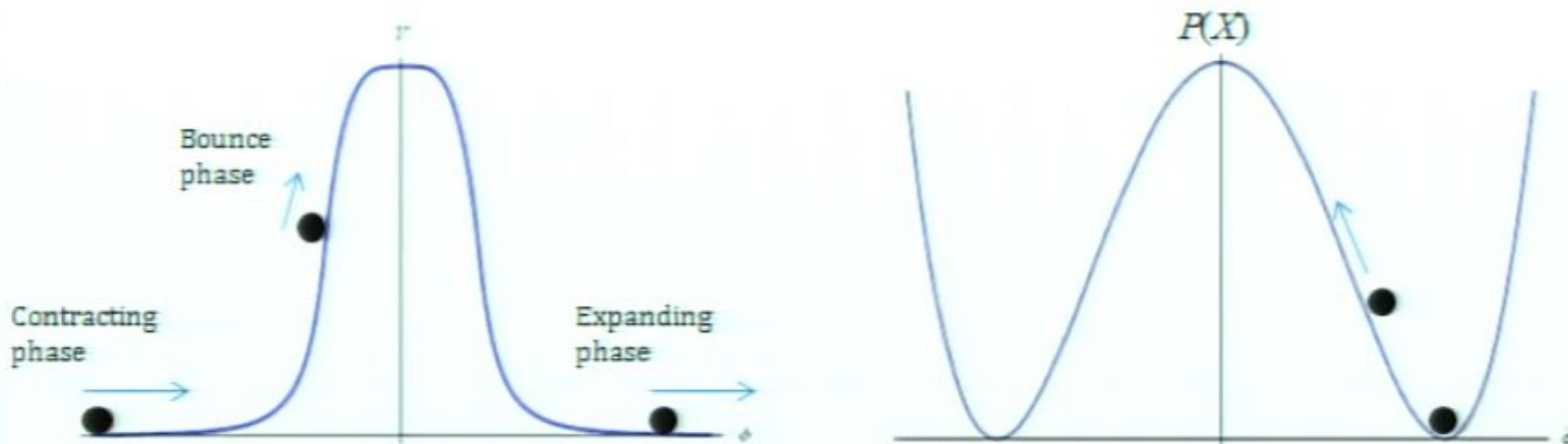


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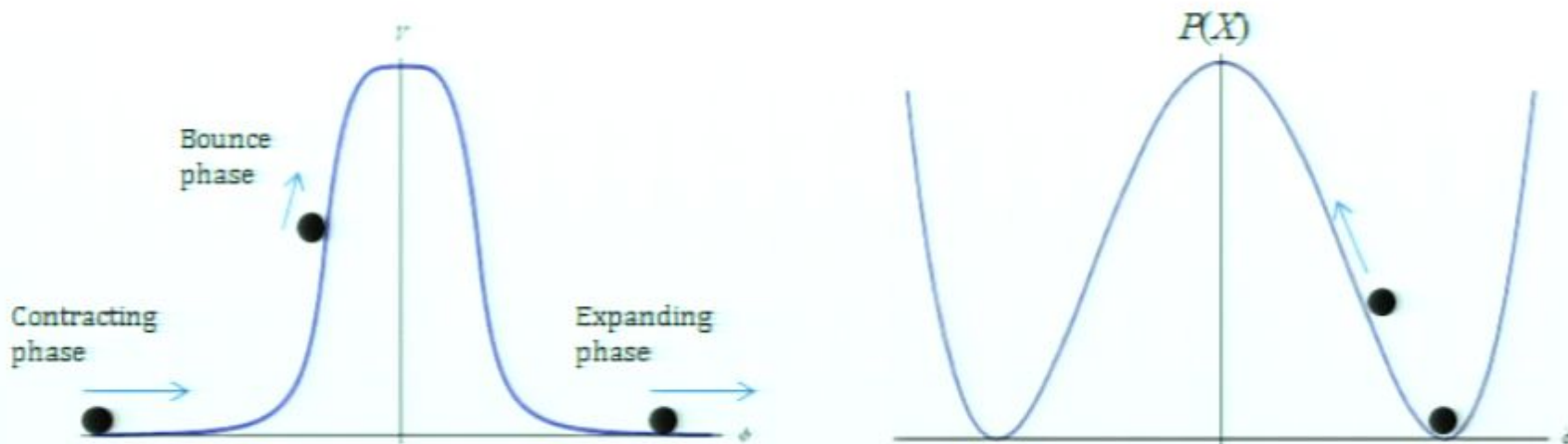


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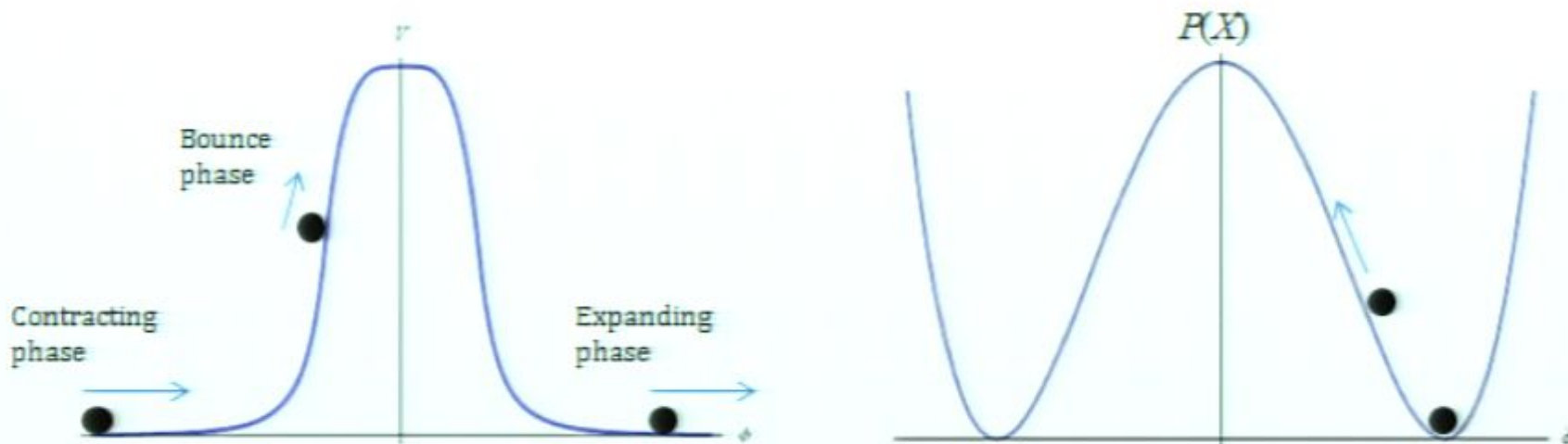


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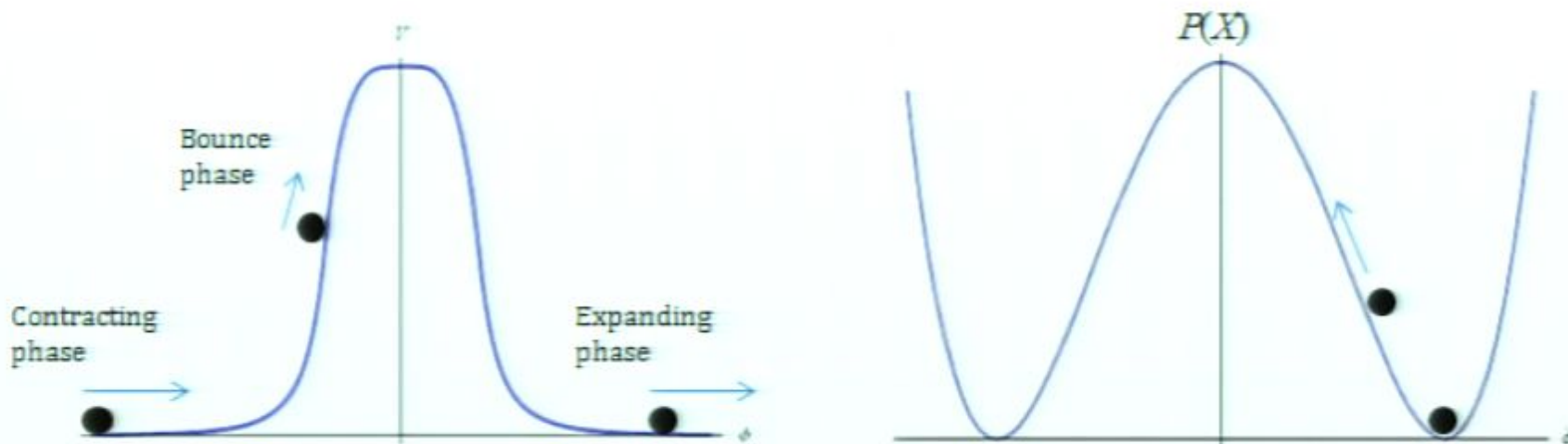


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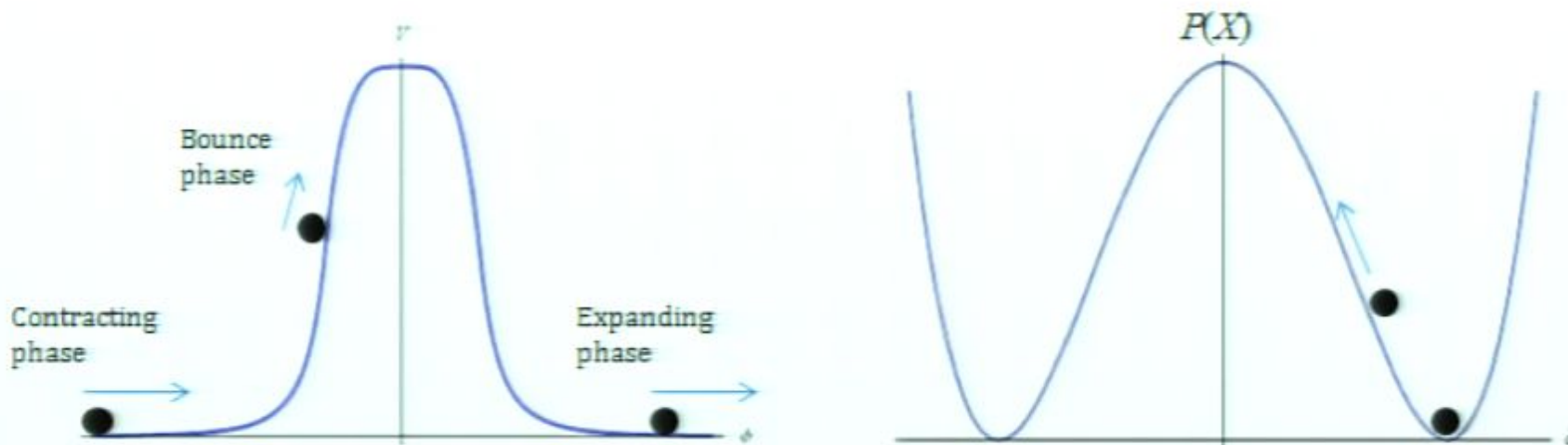


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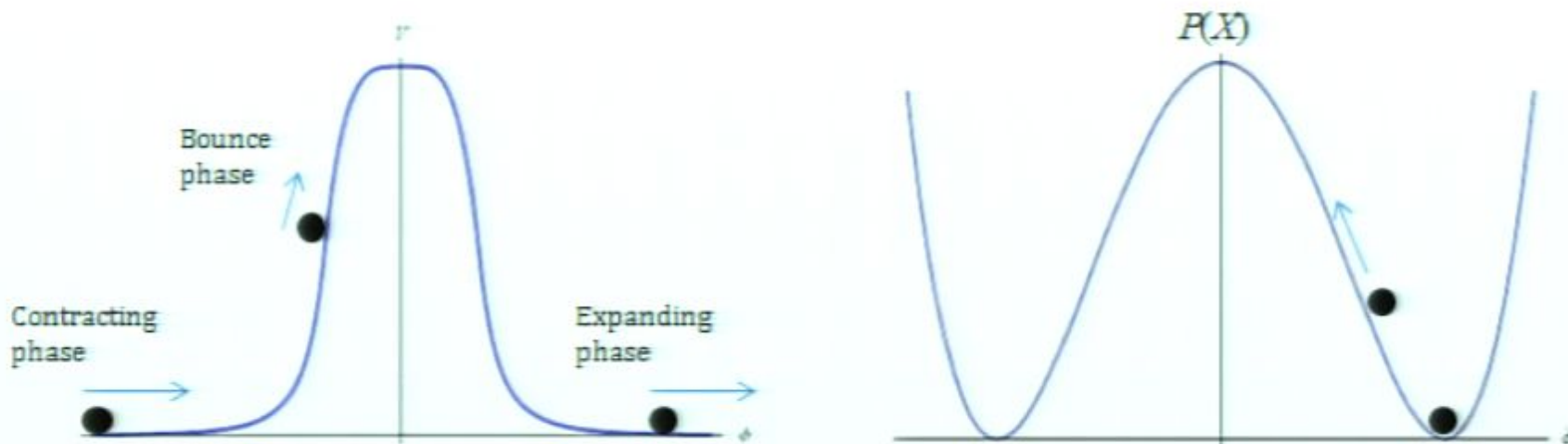


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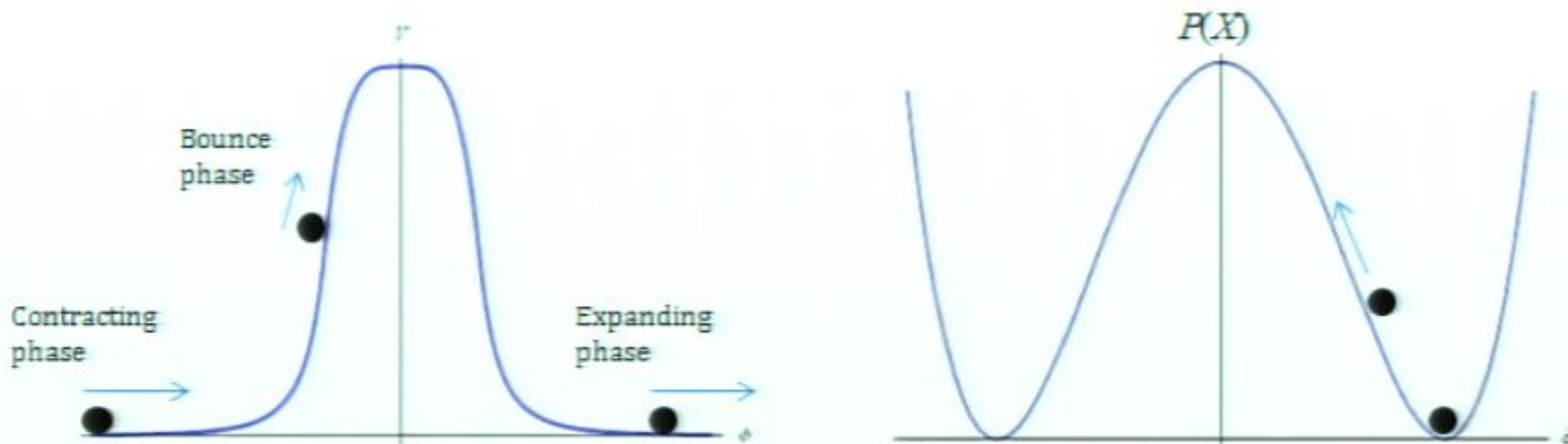


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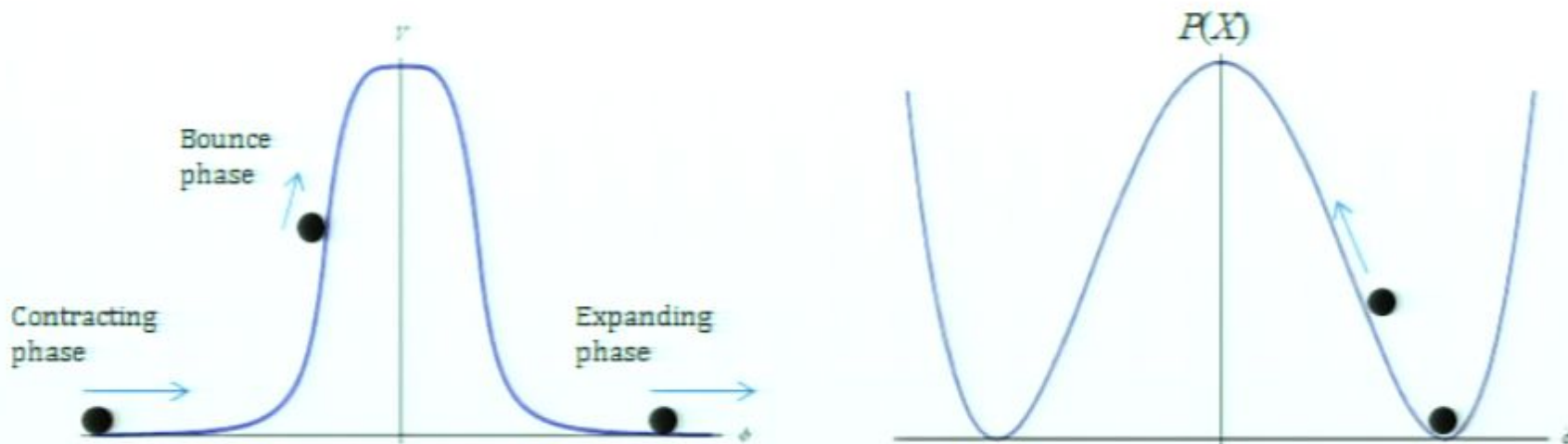


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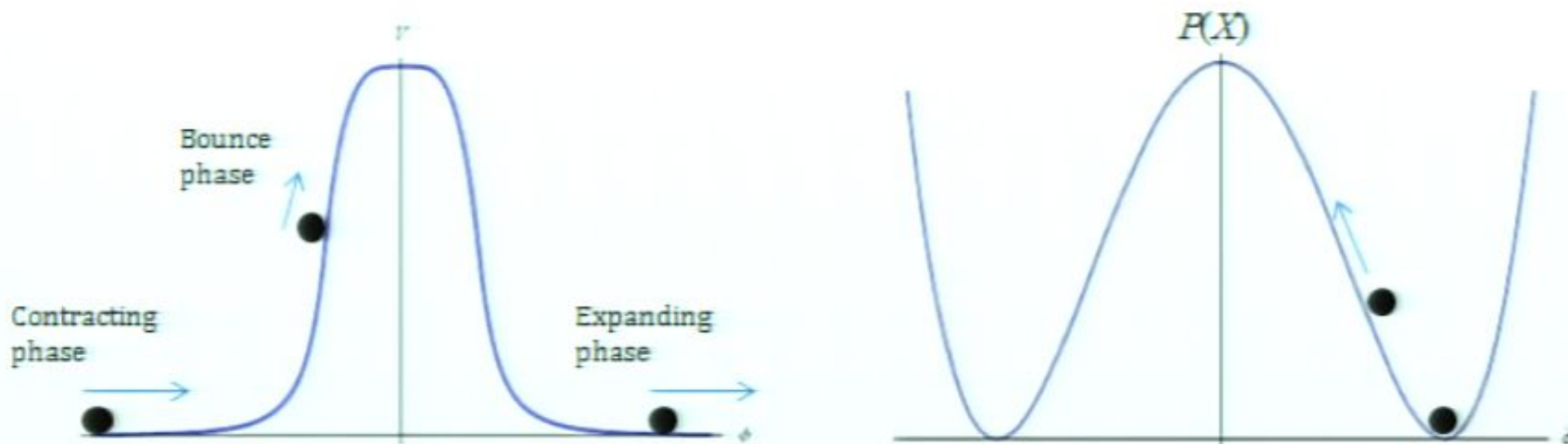


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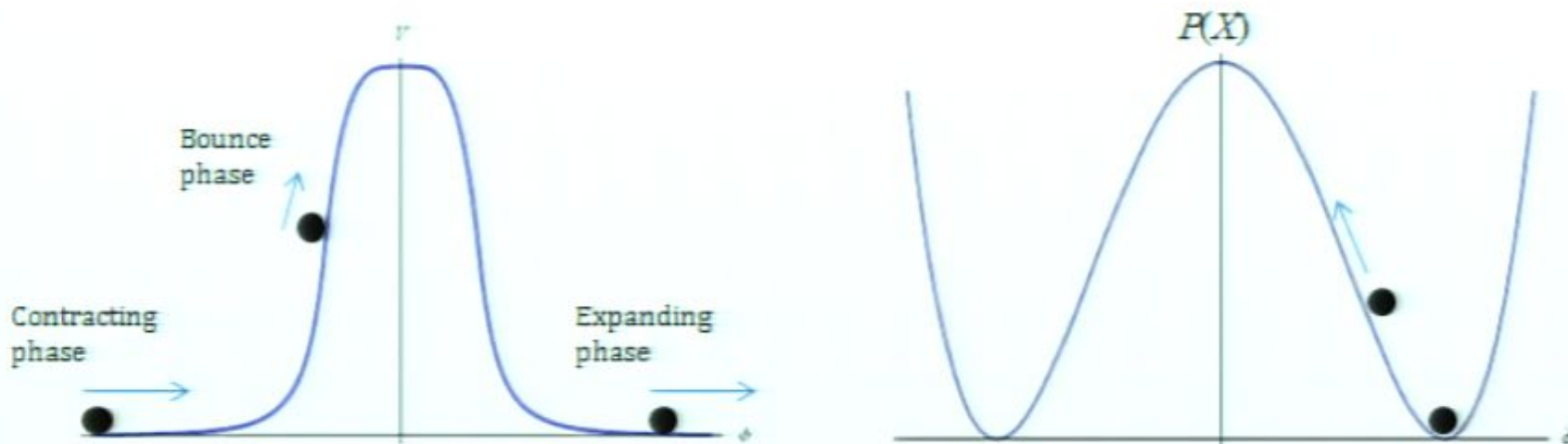


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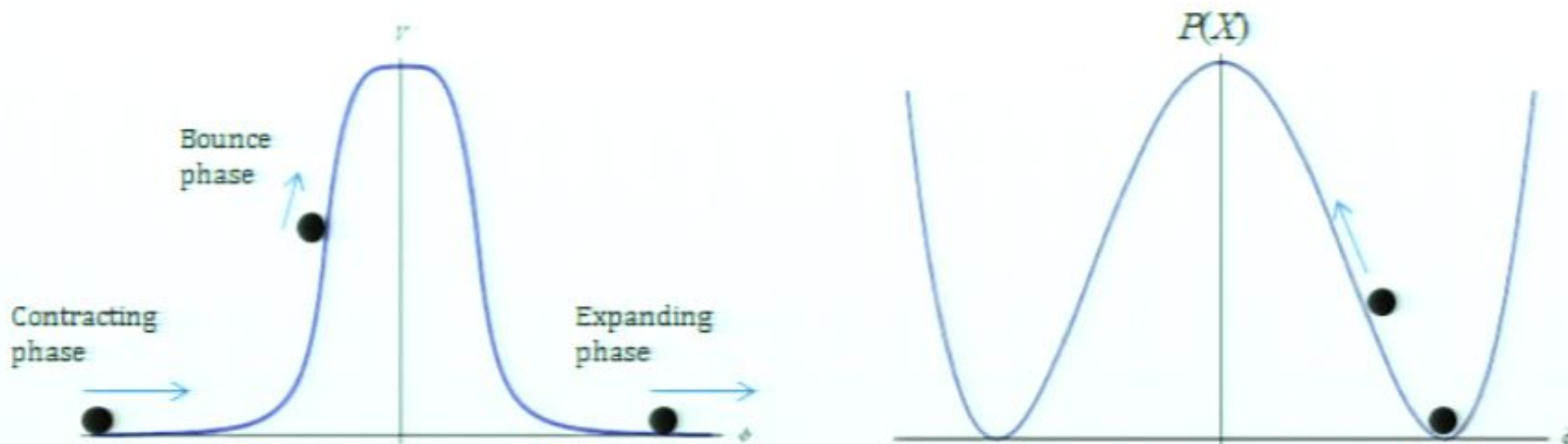


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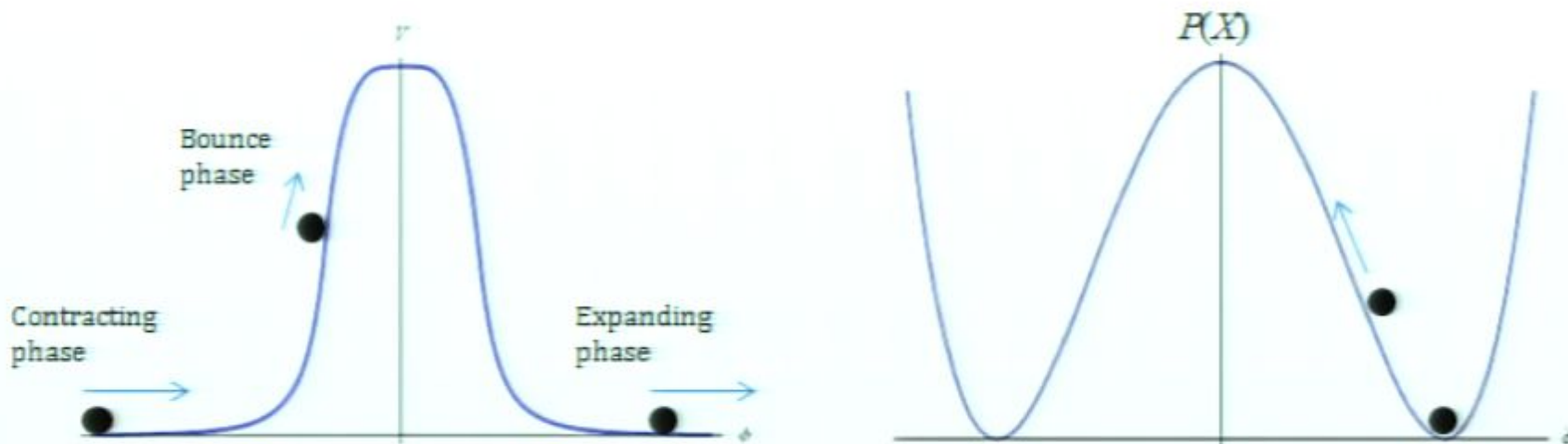


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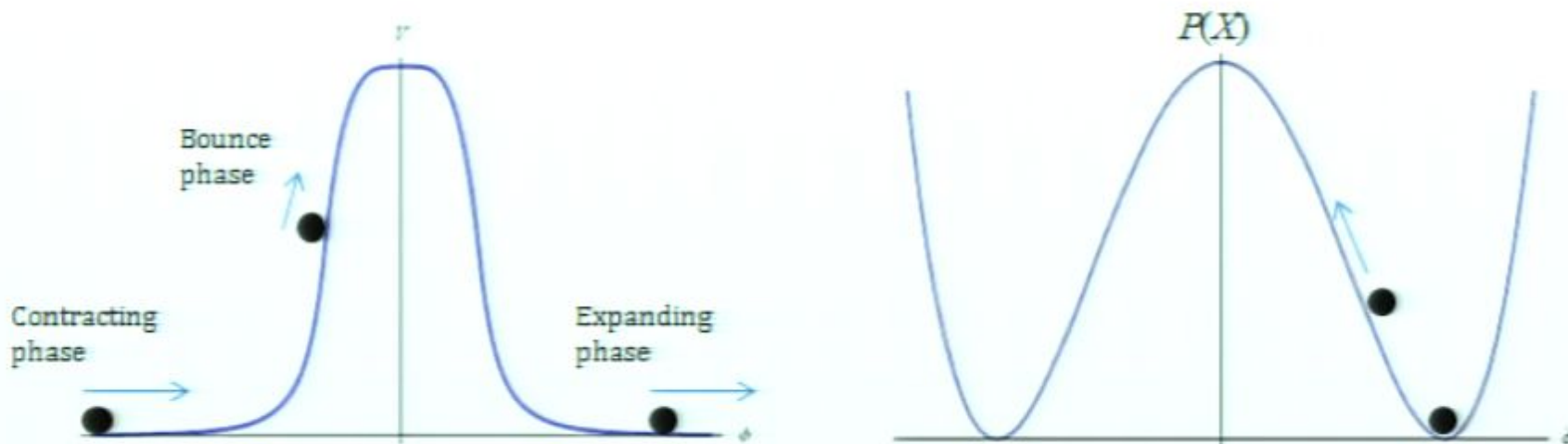


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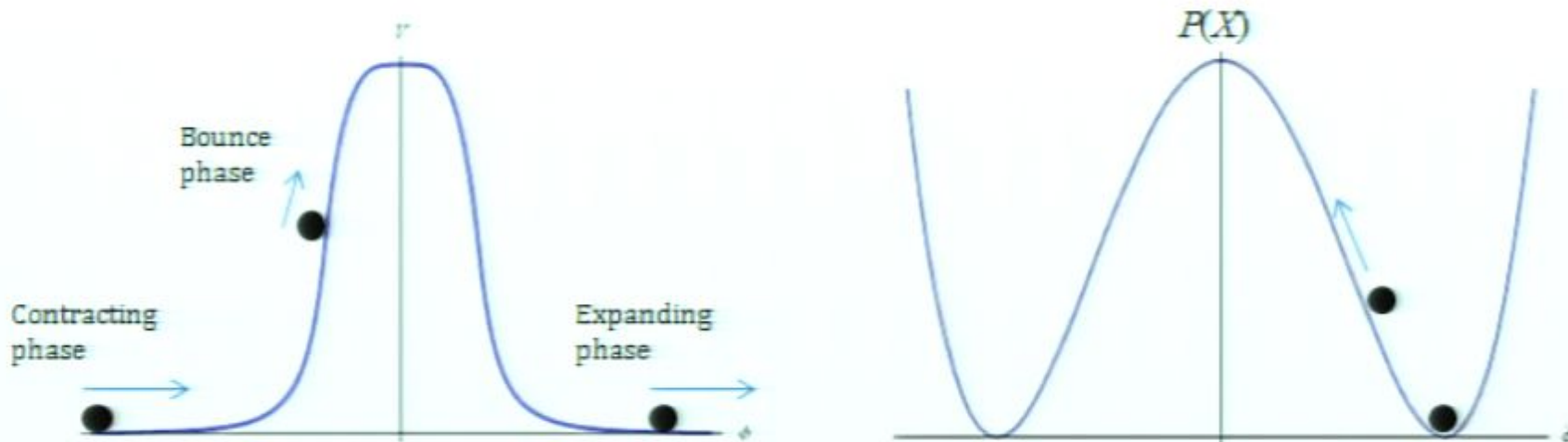


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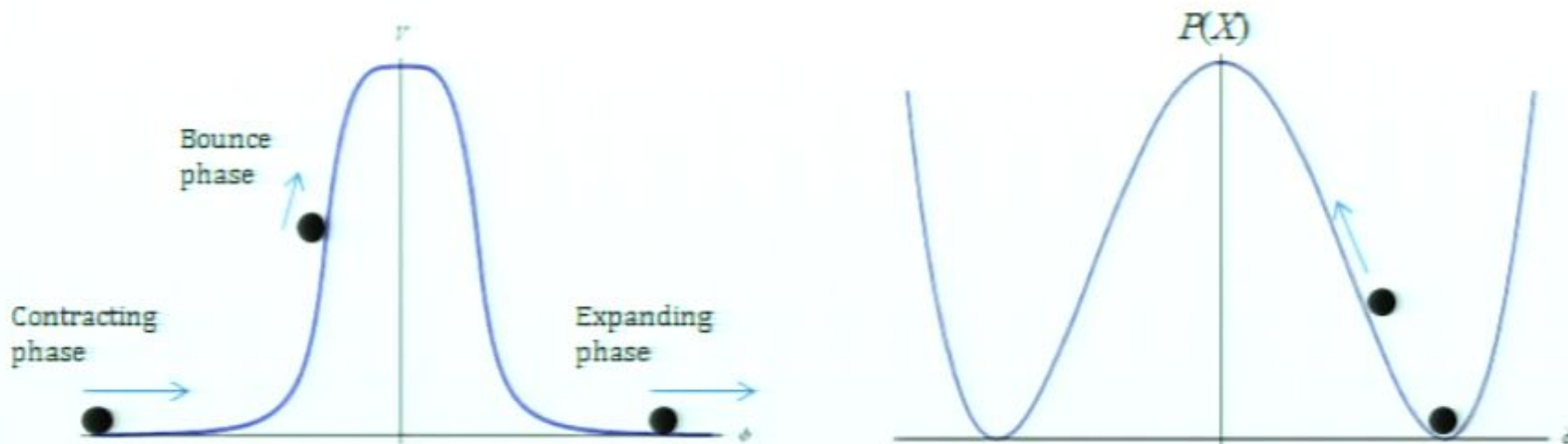


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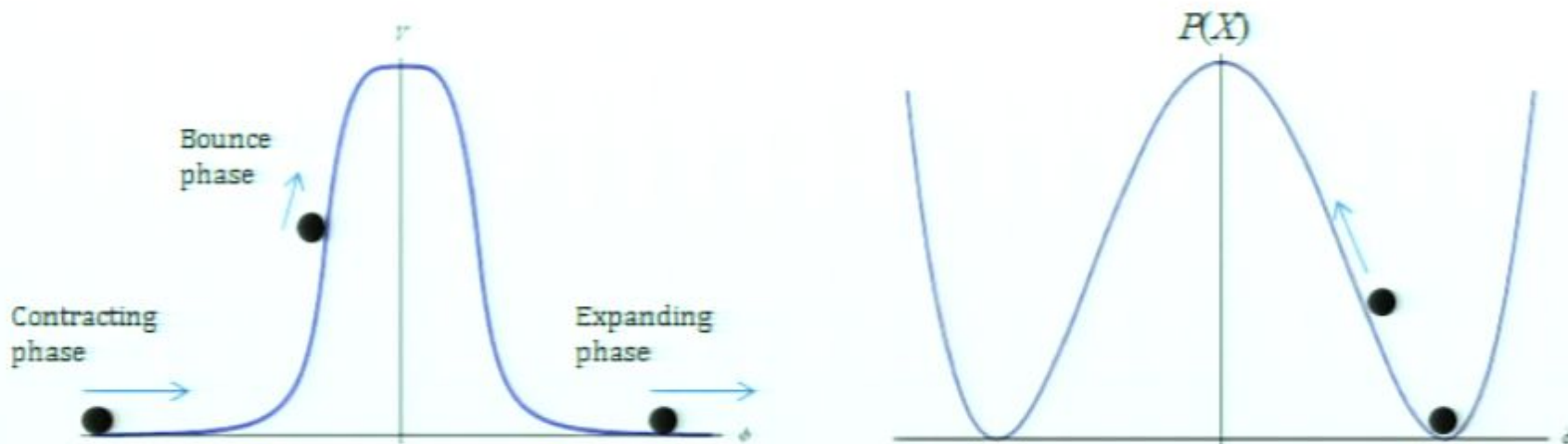


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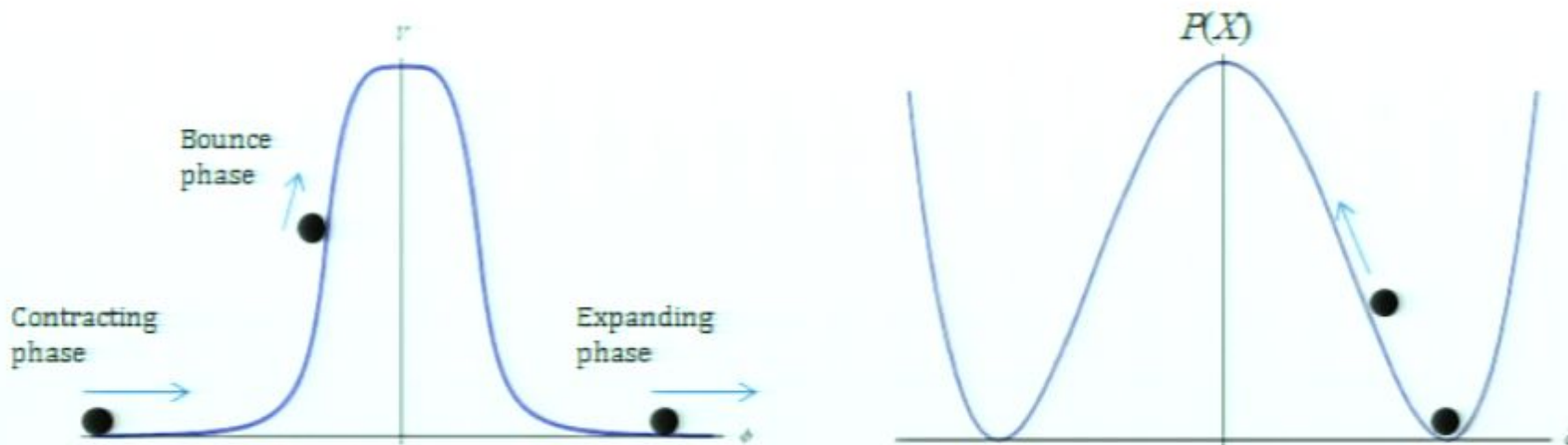


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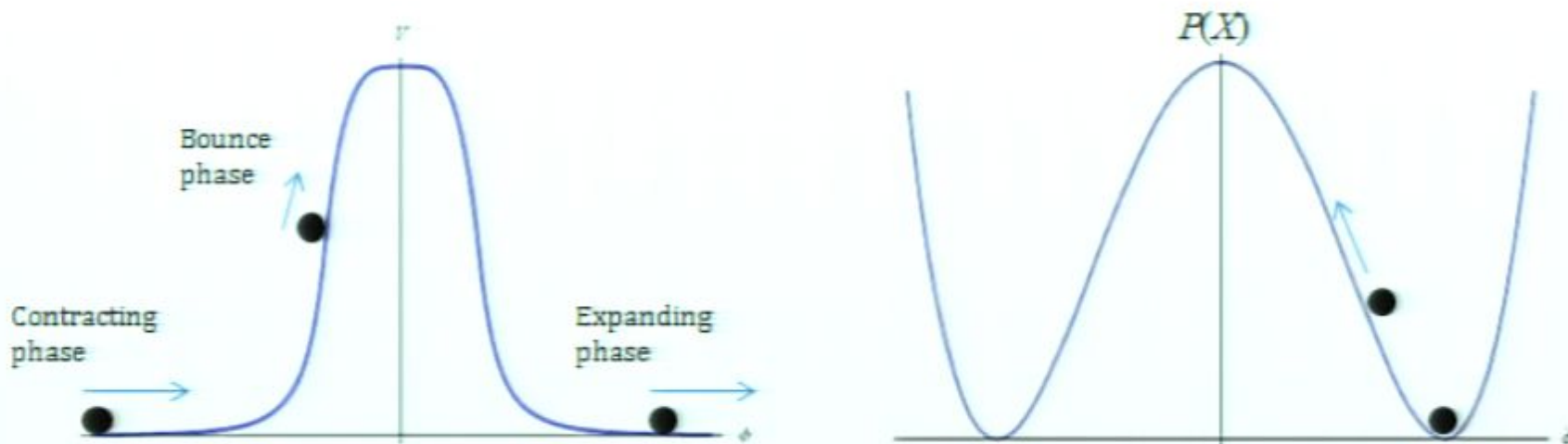


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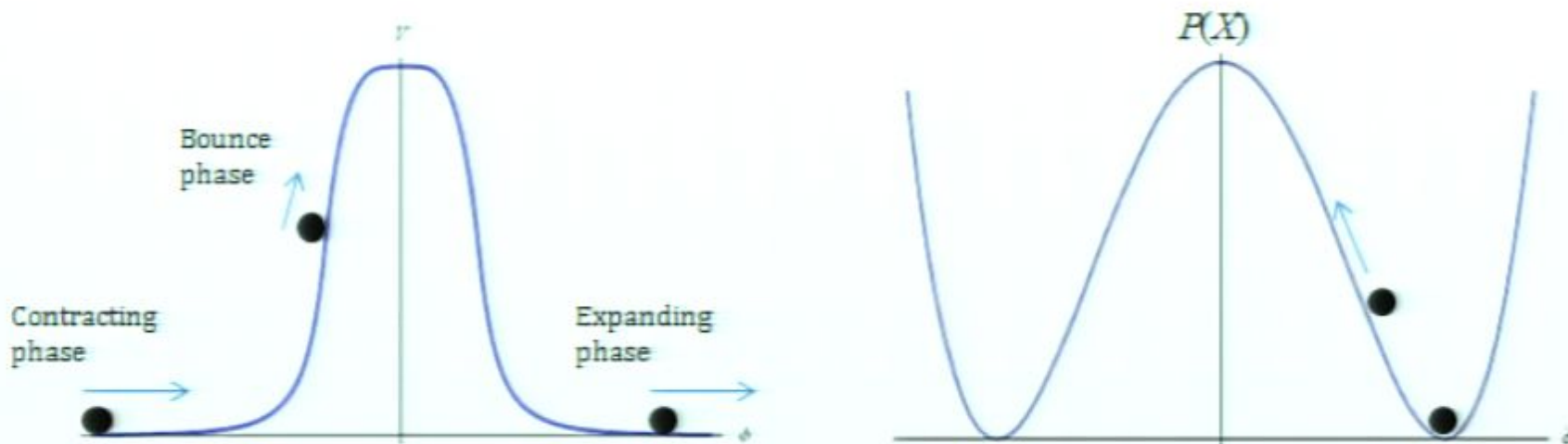


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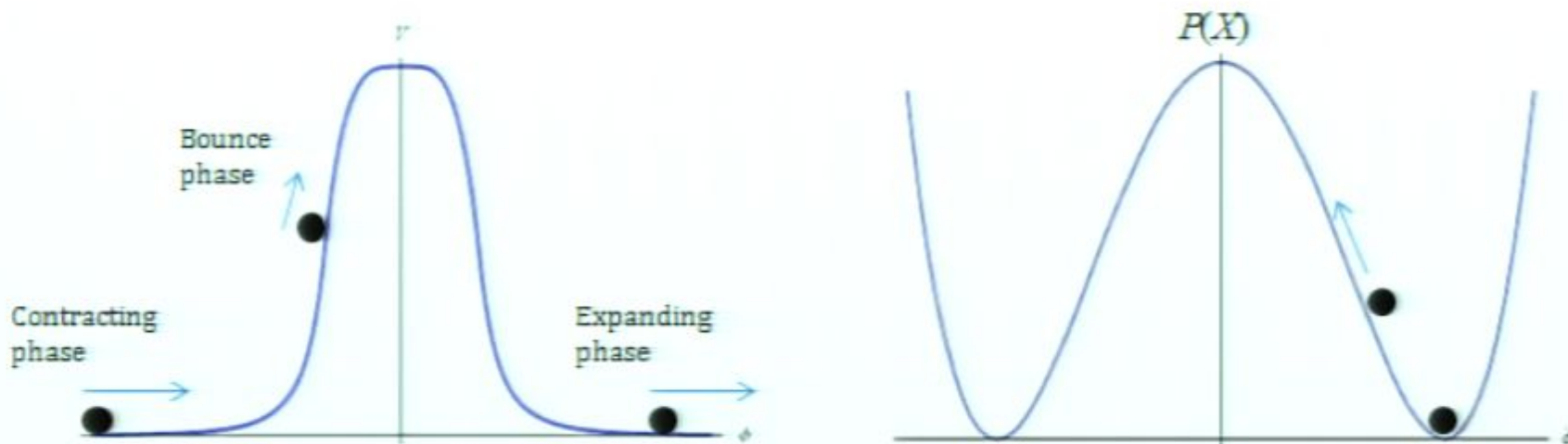


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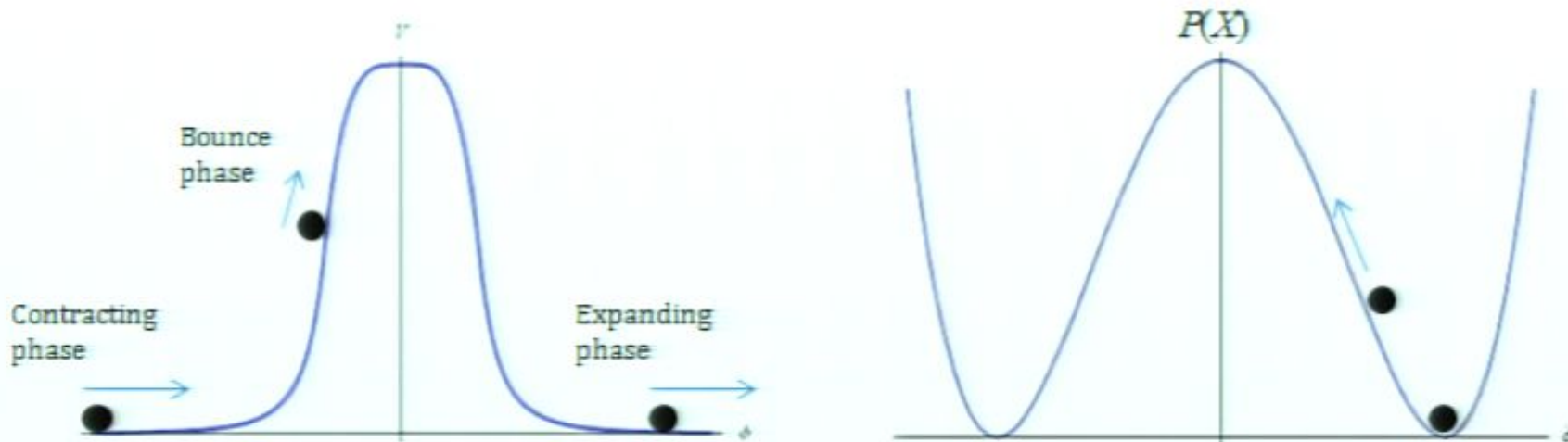


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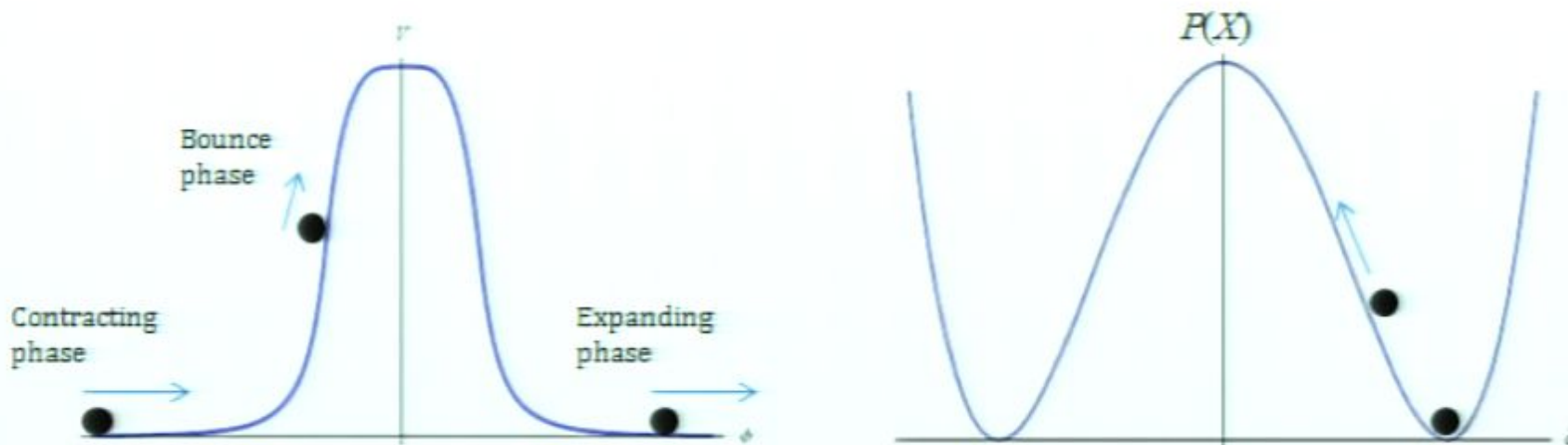


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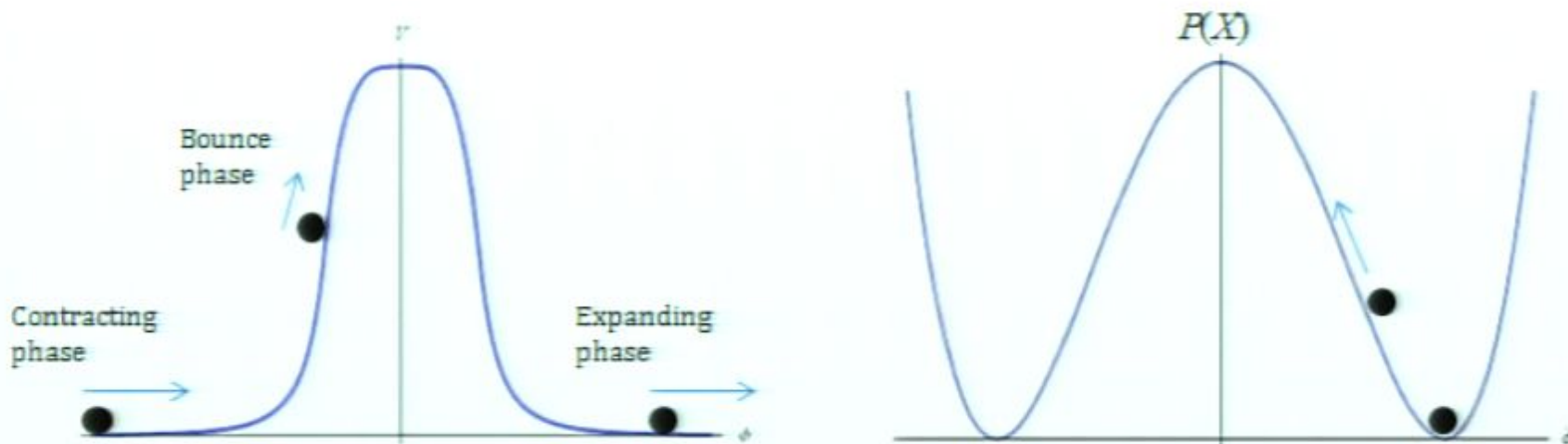


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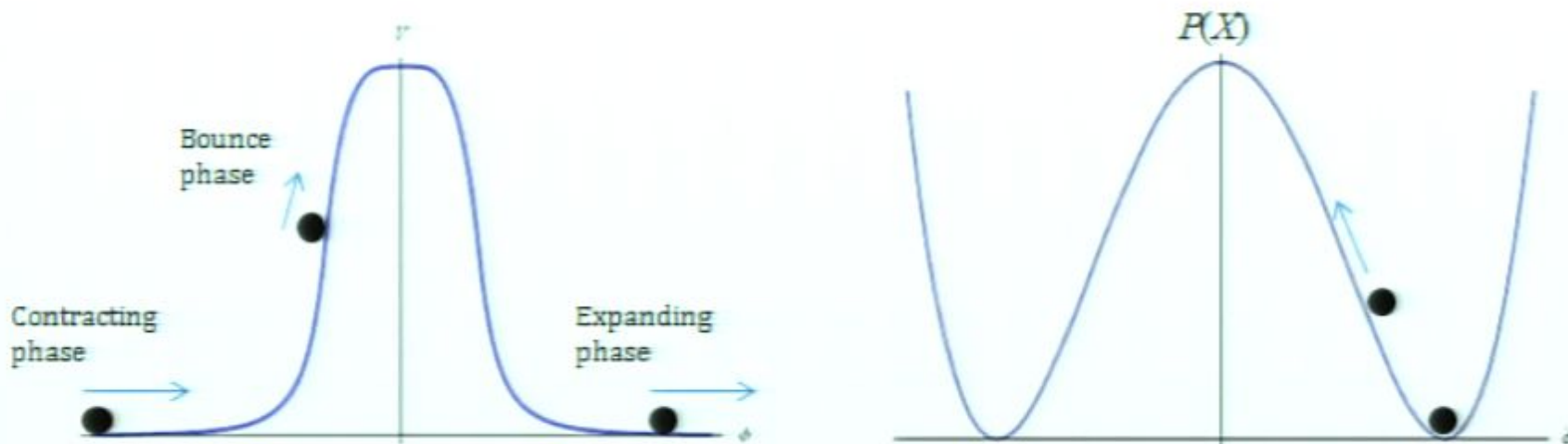


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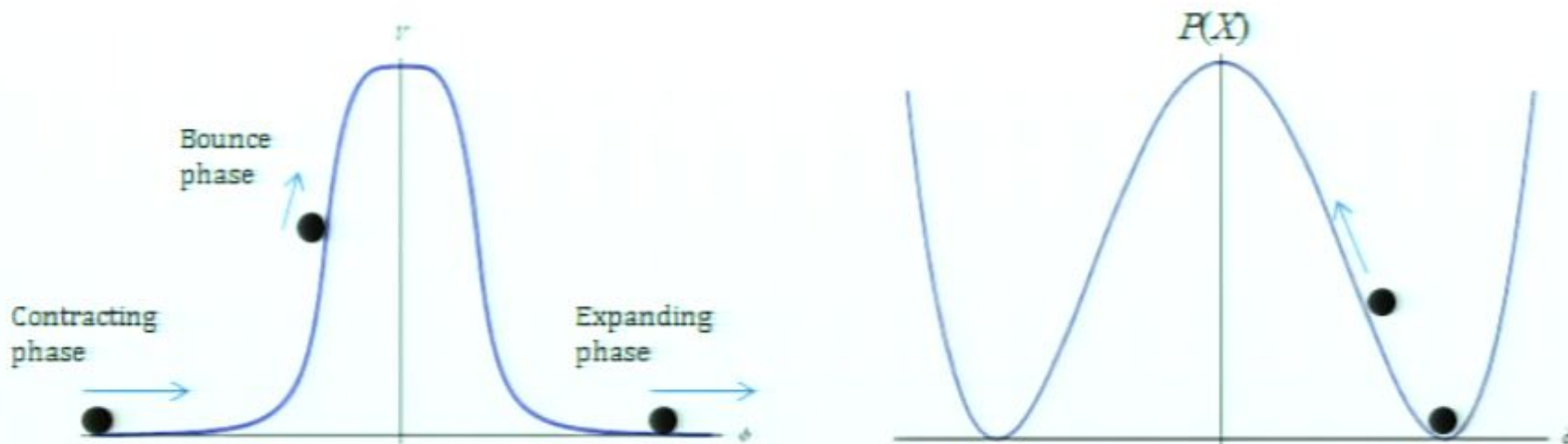


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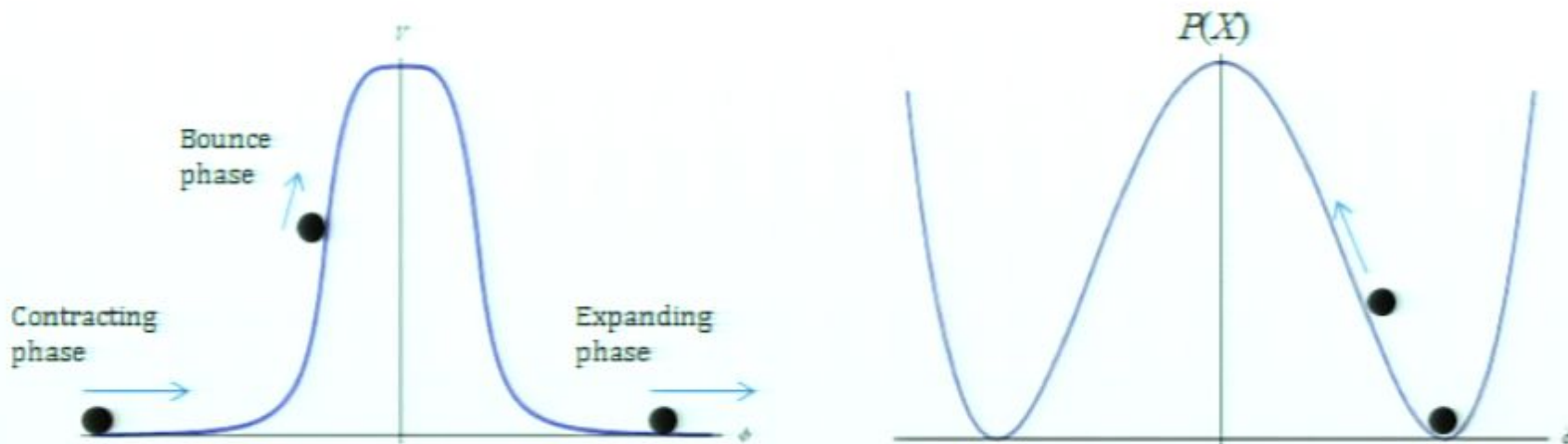


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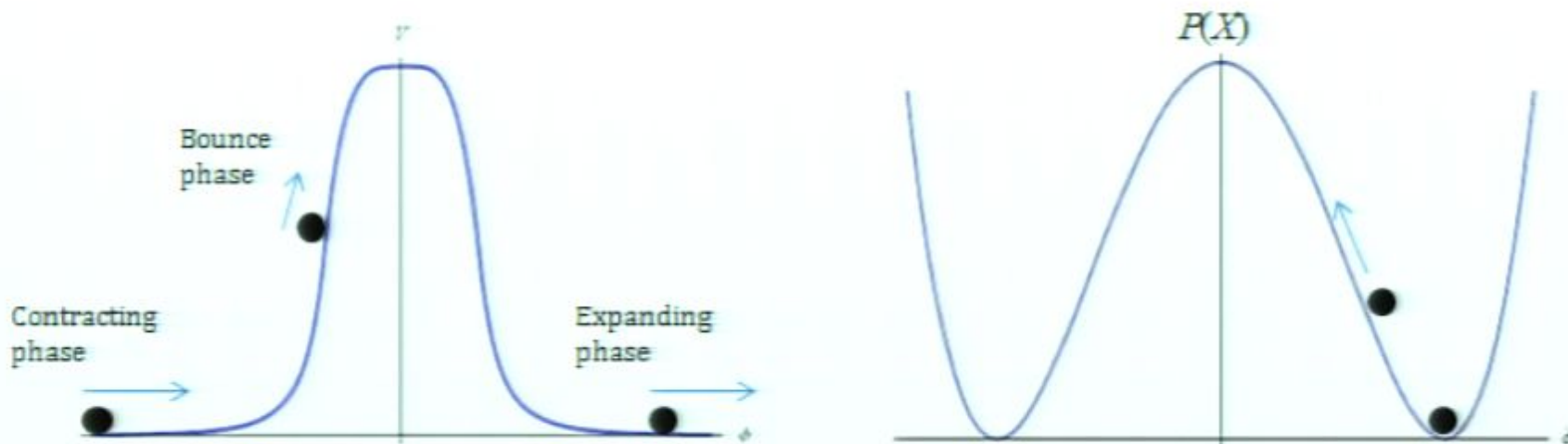


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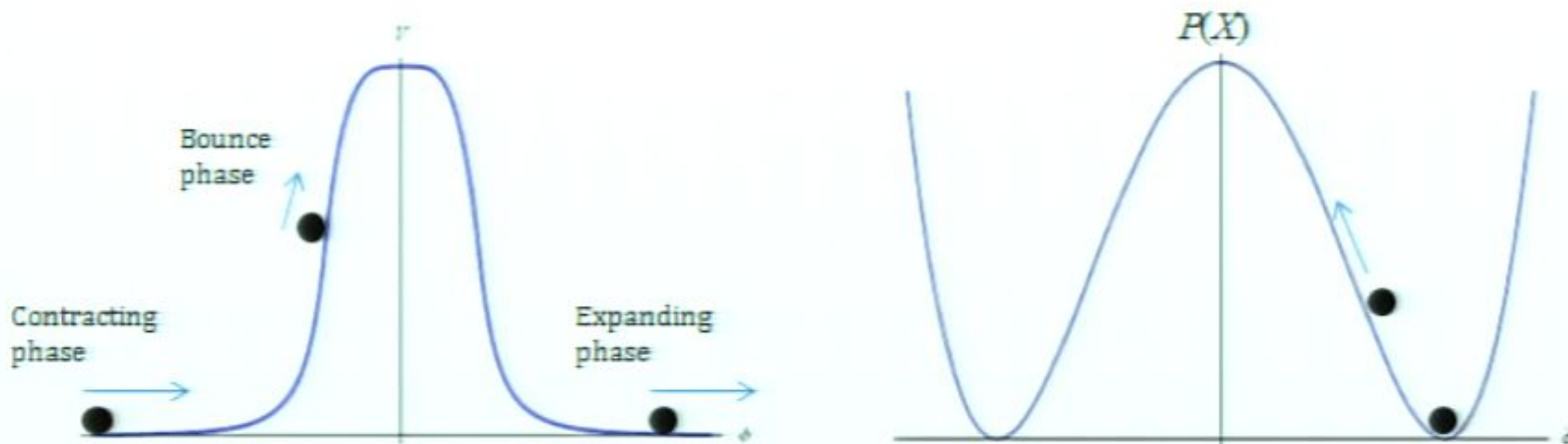


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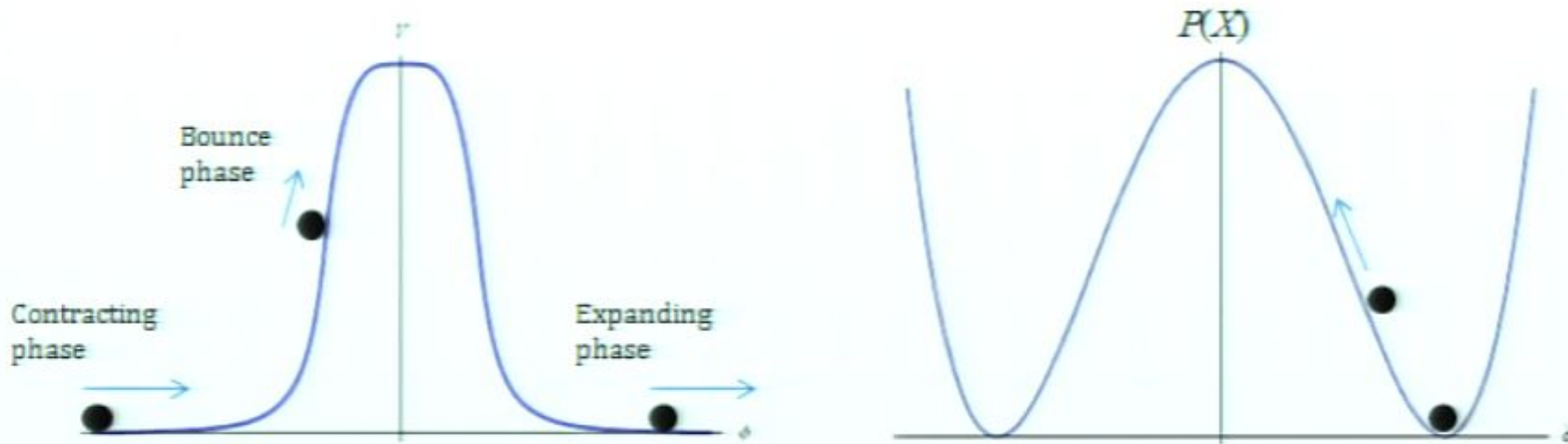


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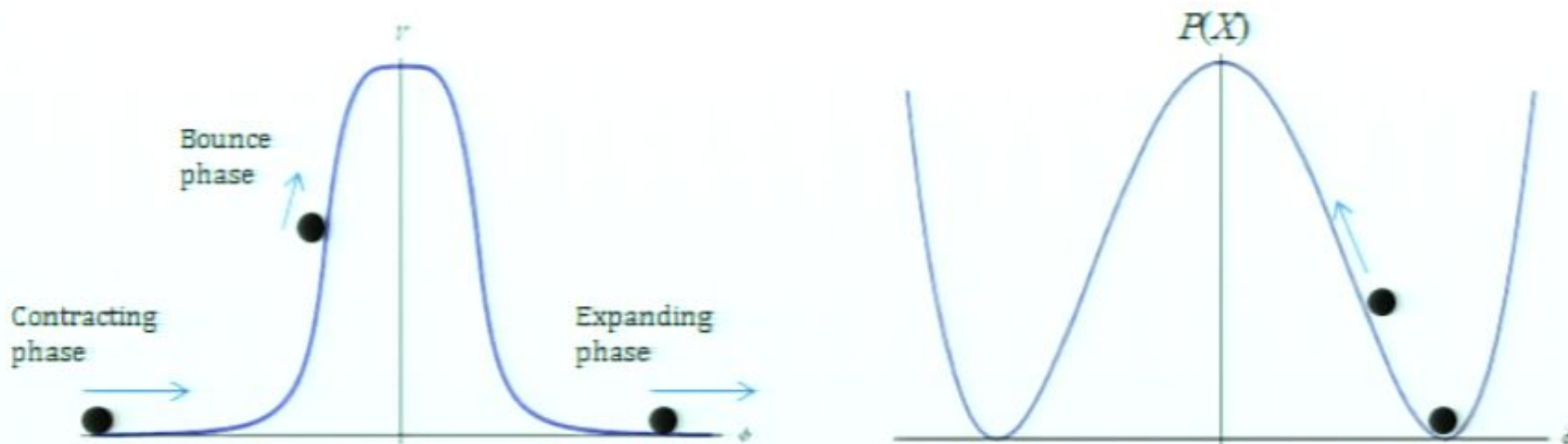


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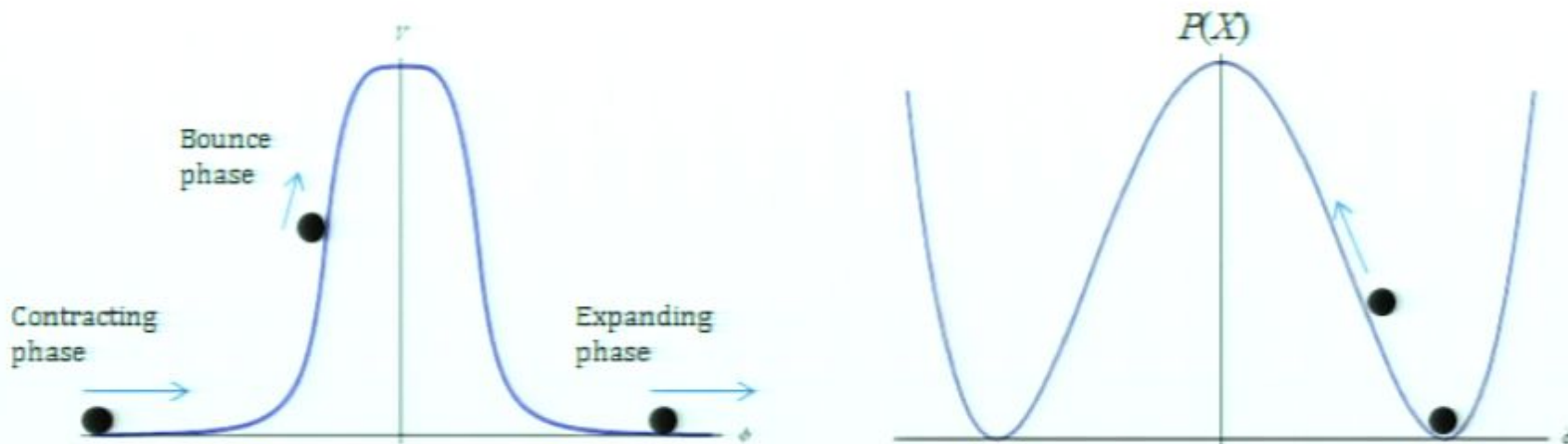


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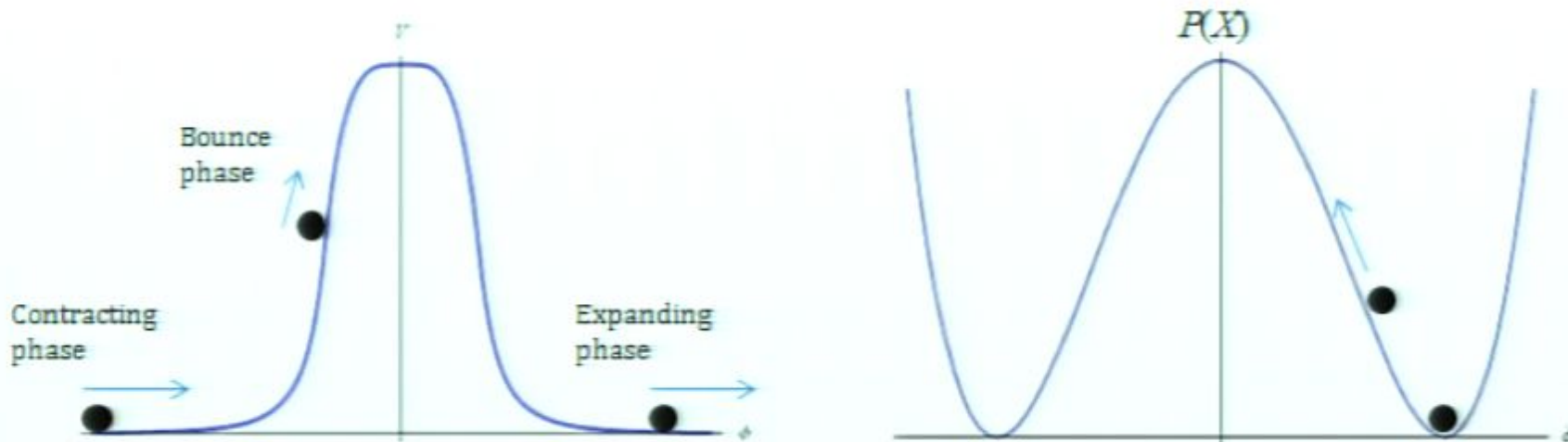


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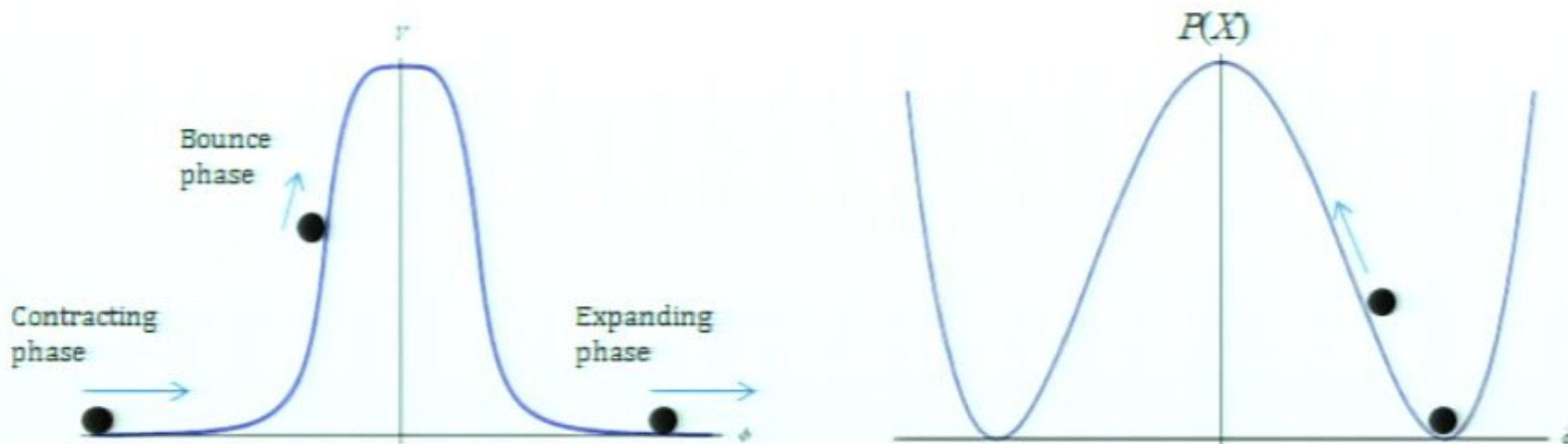


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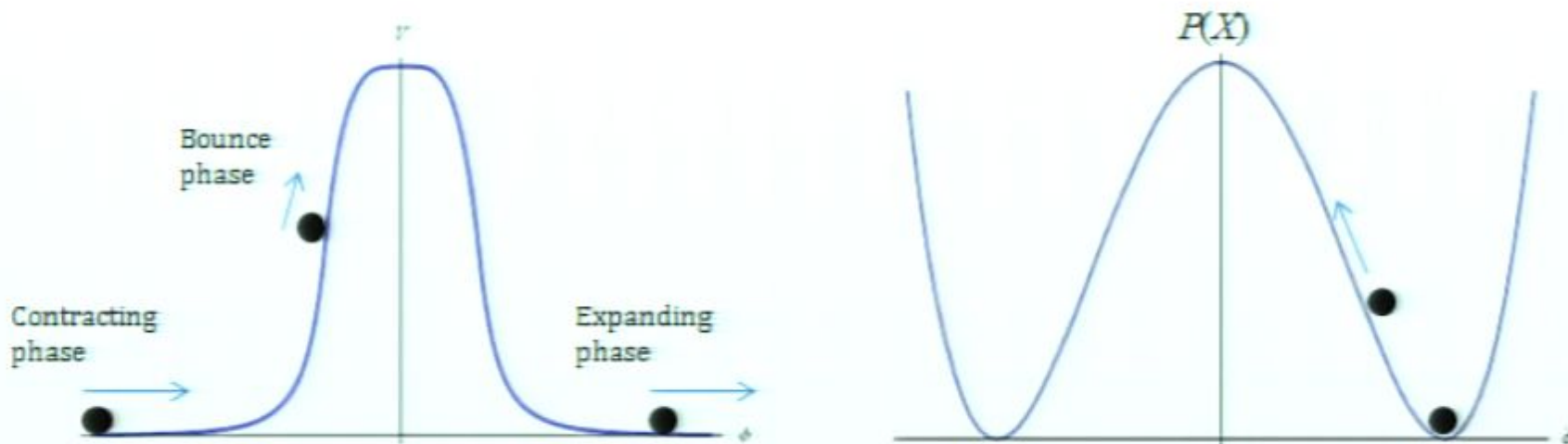


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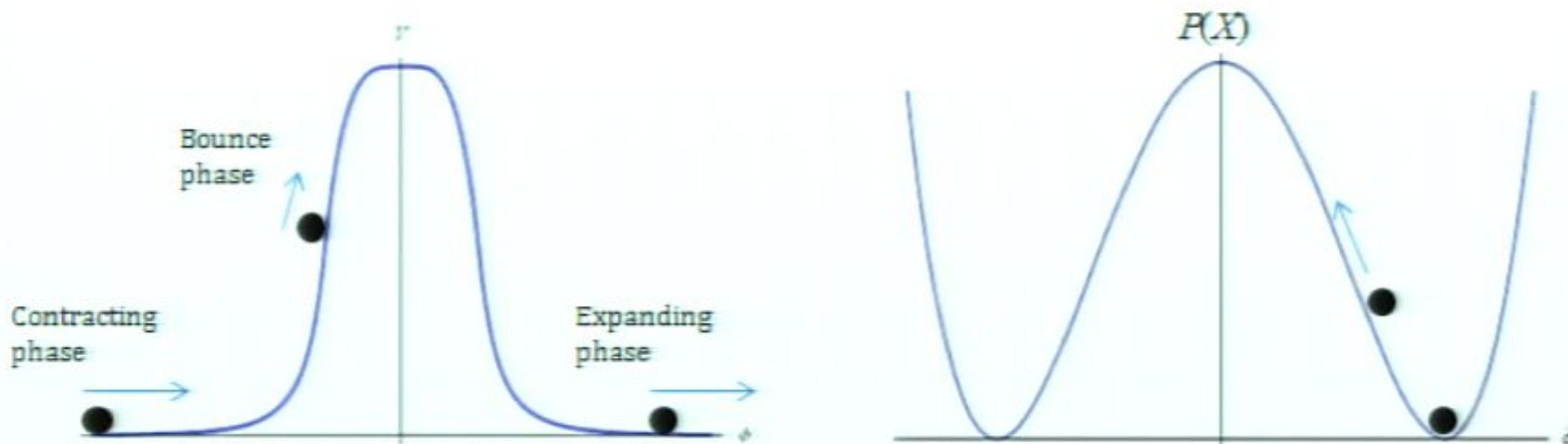


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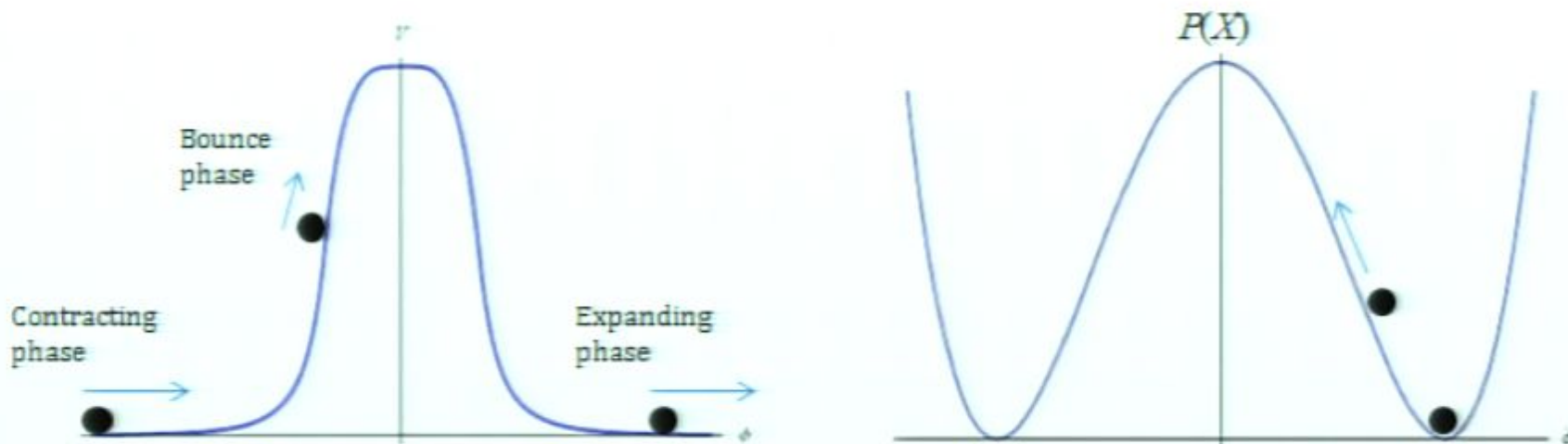


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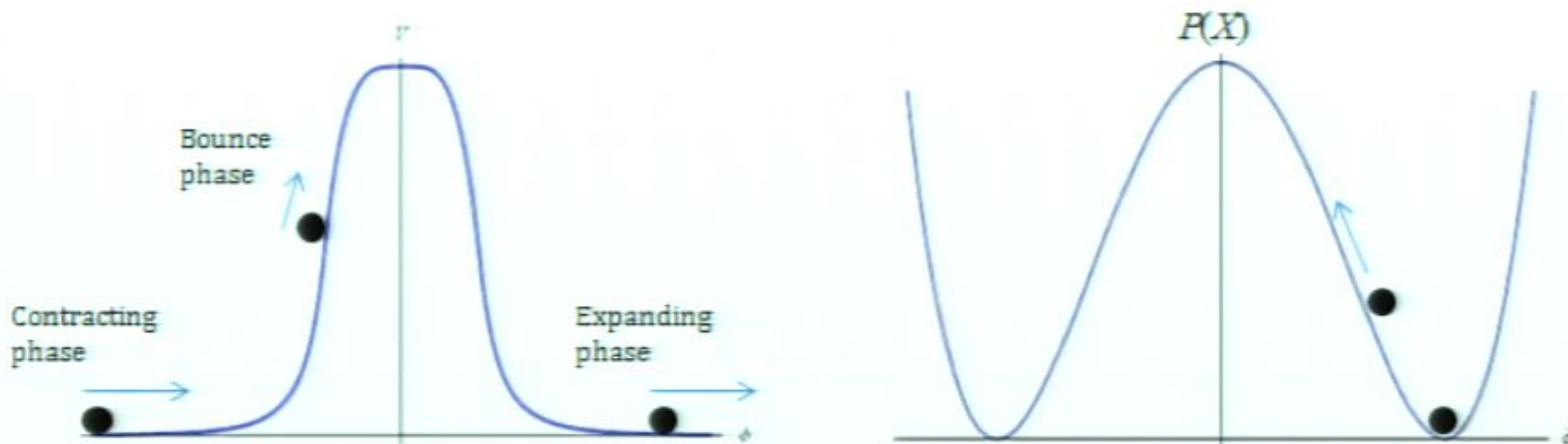


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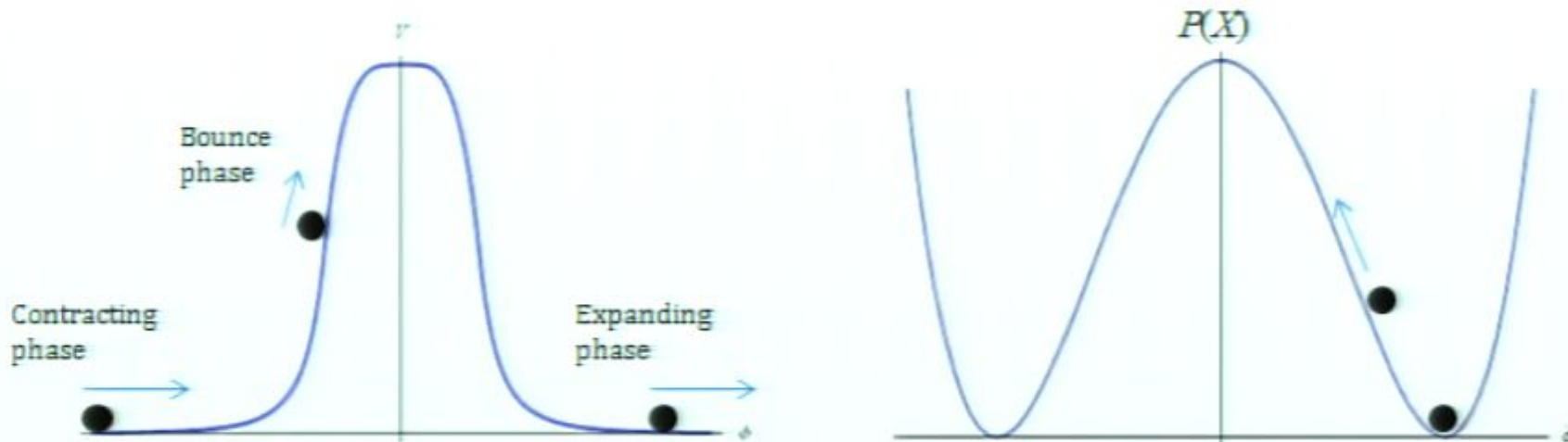


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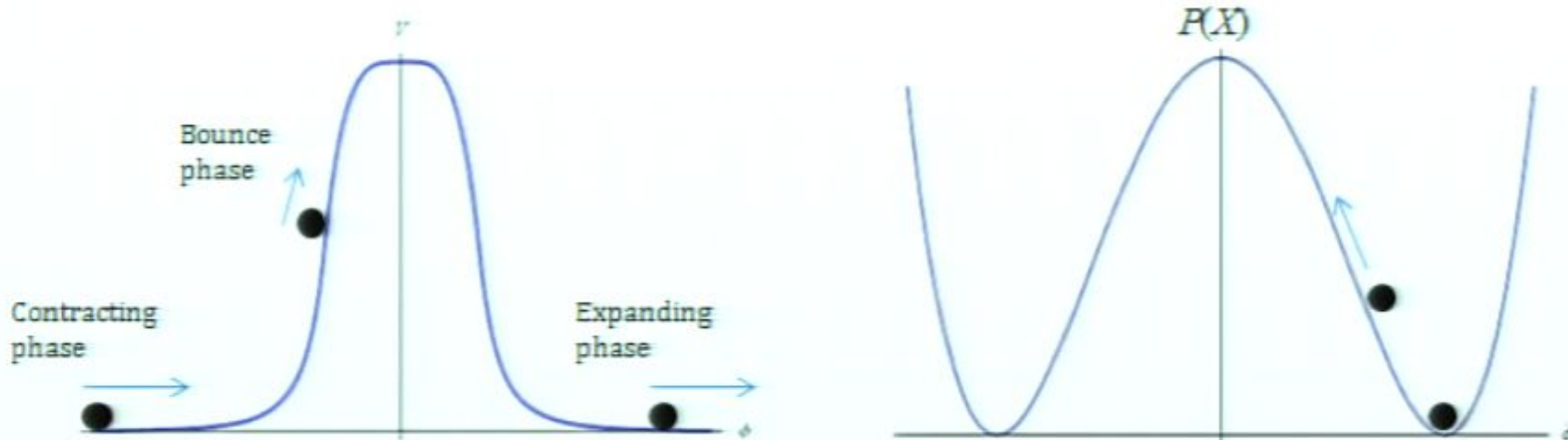


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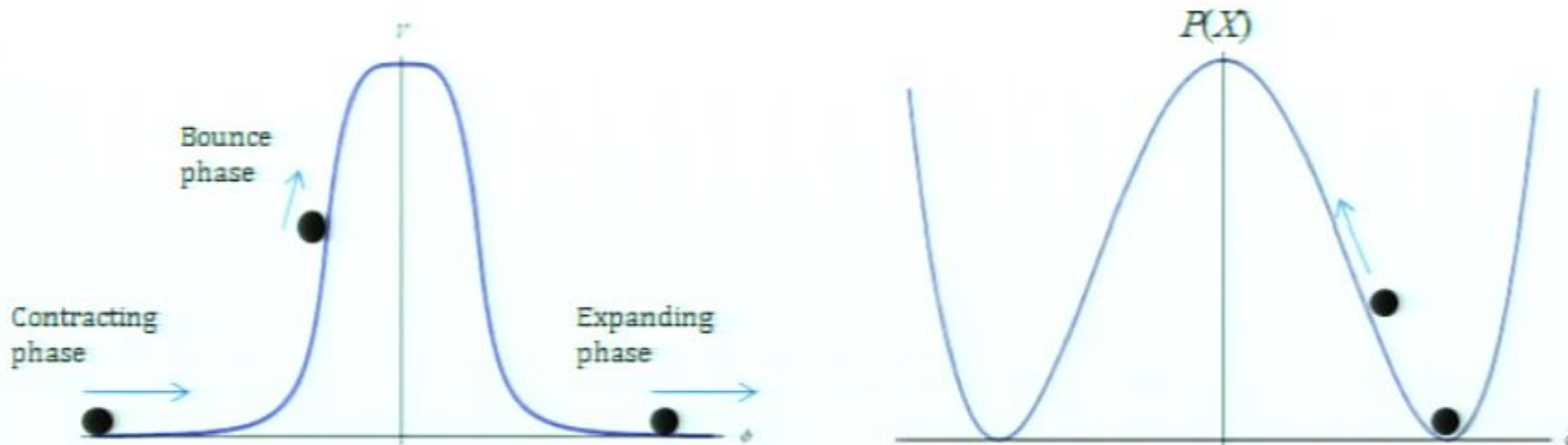


$$2M_p^2 \dot{H} = -2M^4 X P' - (1 + w_m) \rho_m$$

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Ghost bounce

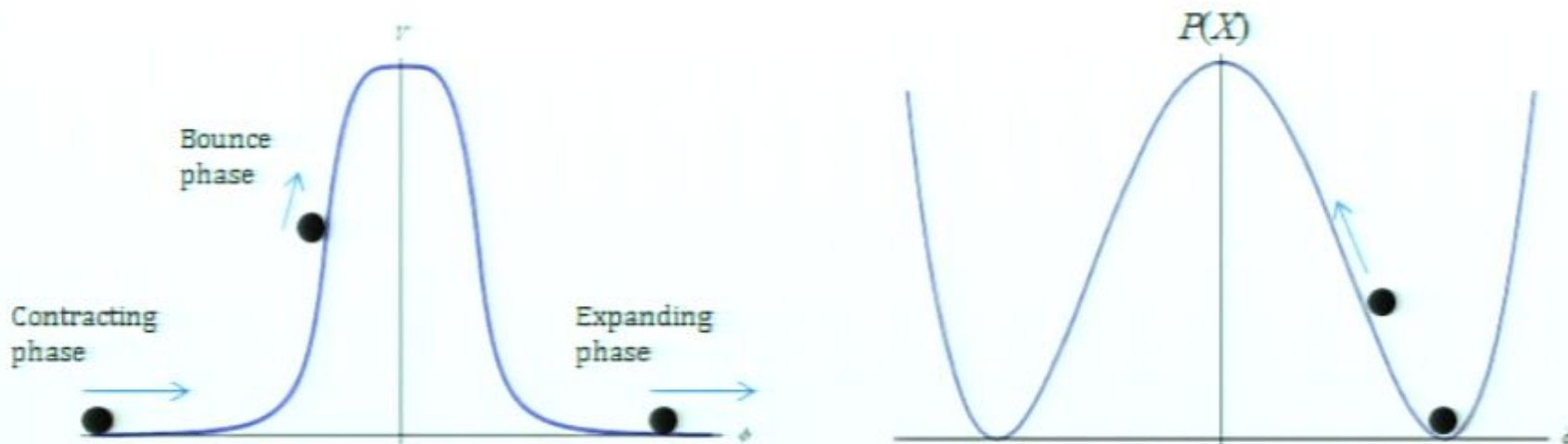


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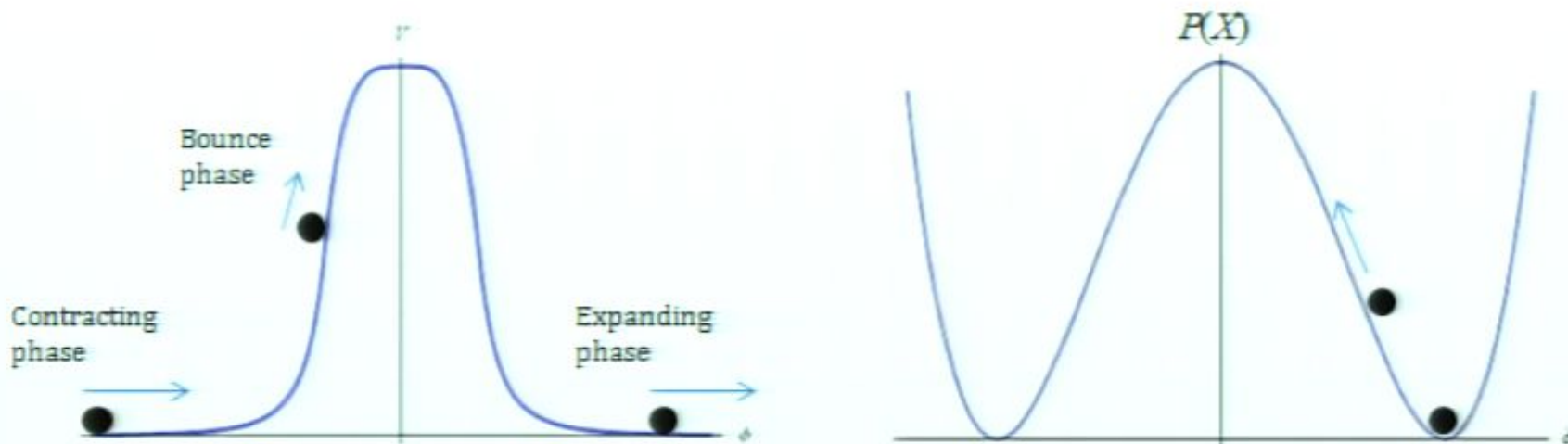


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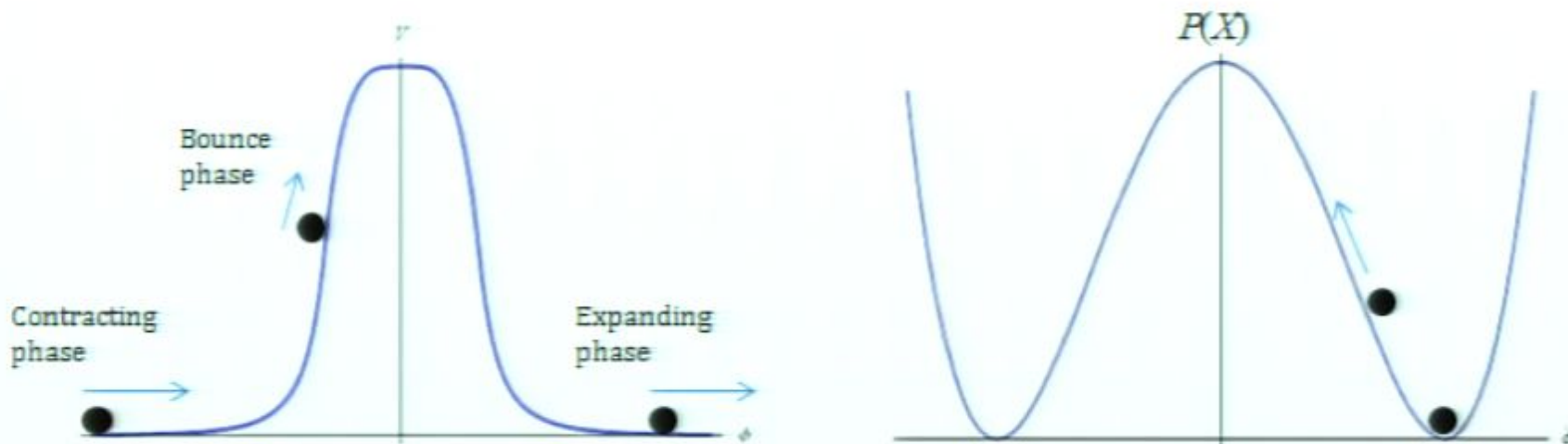


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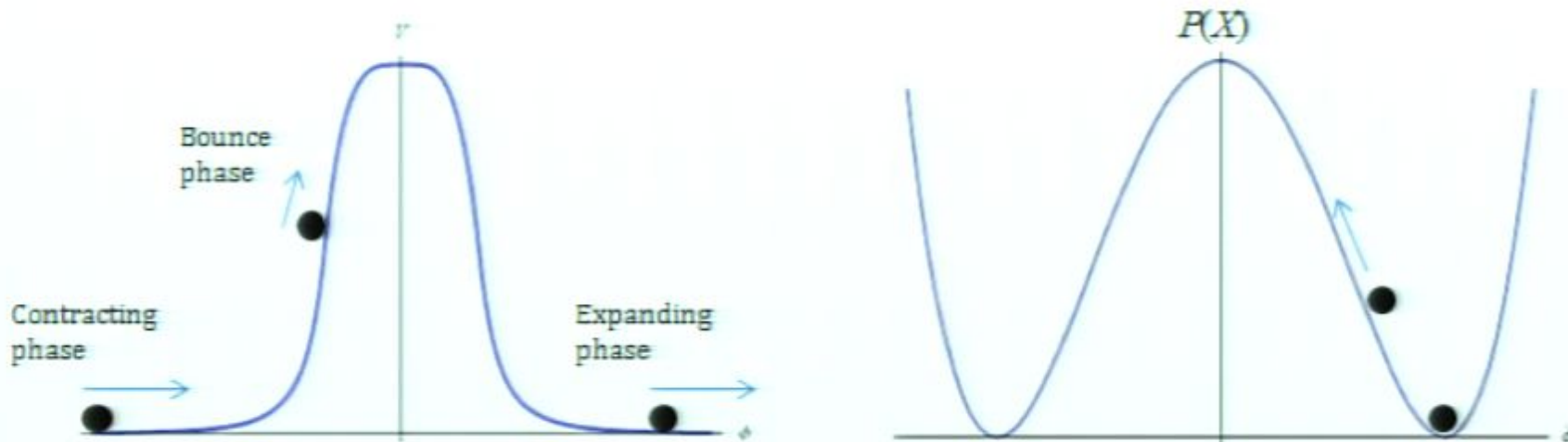


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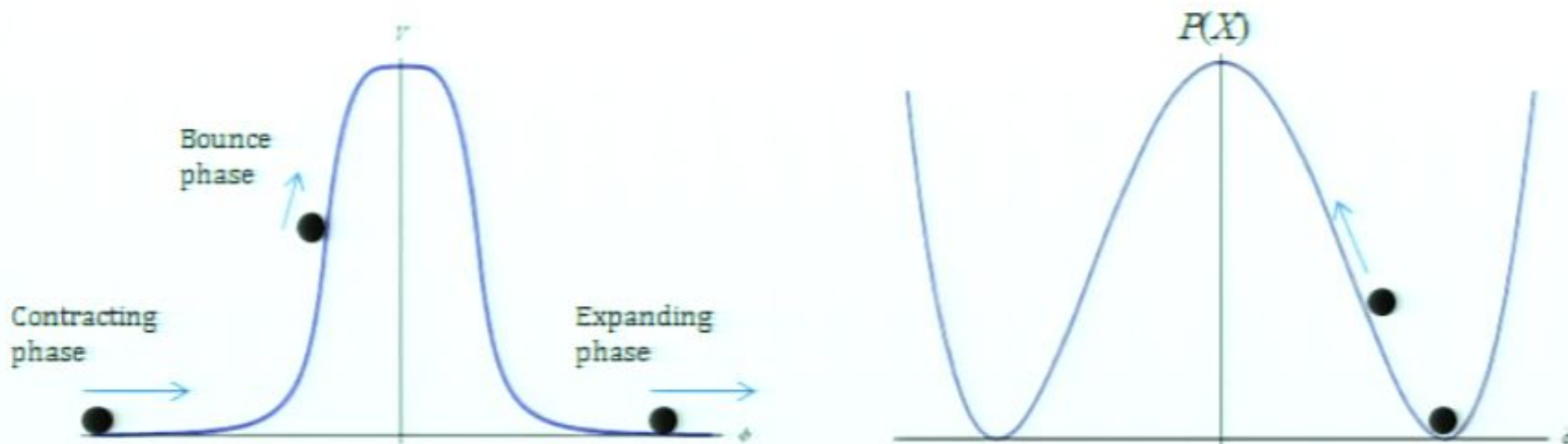


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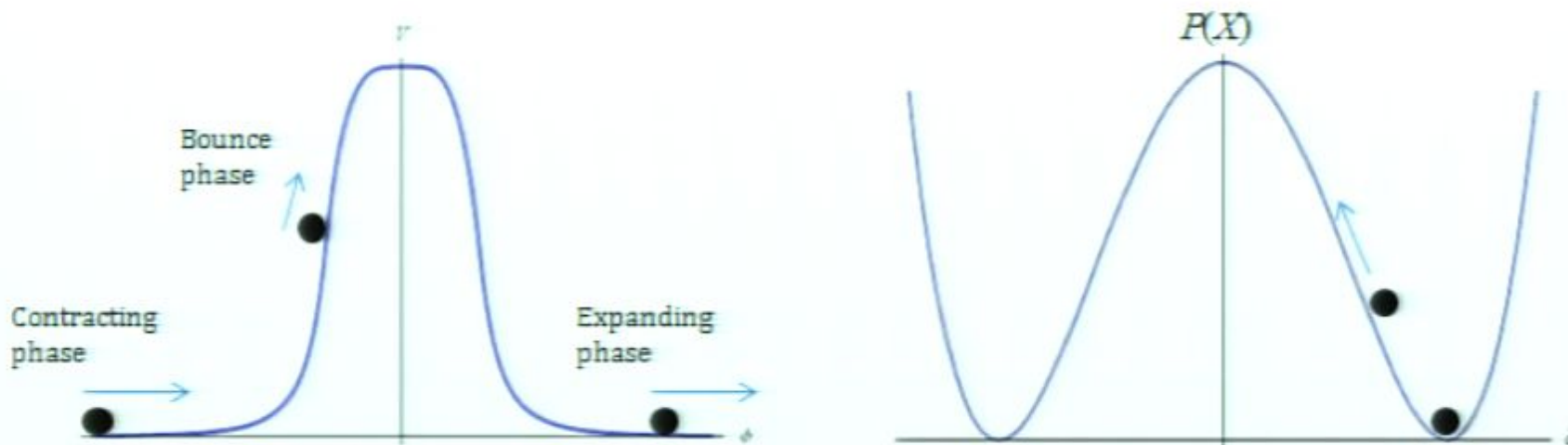


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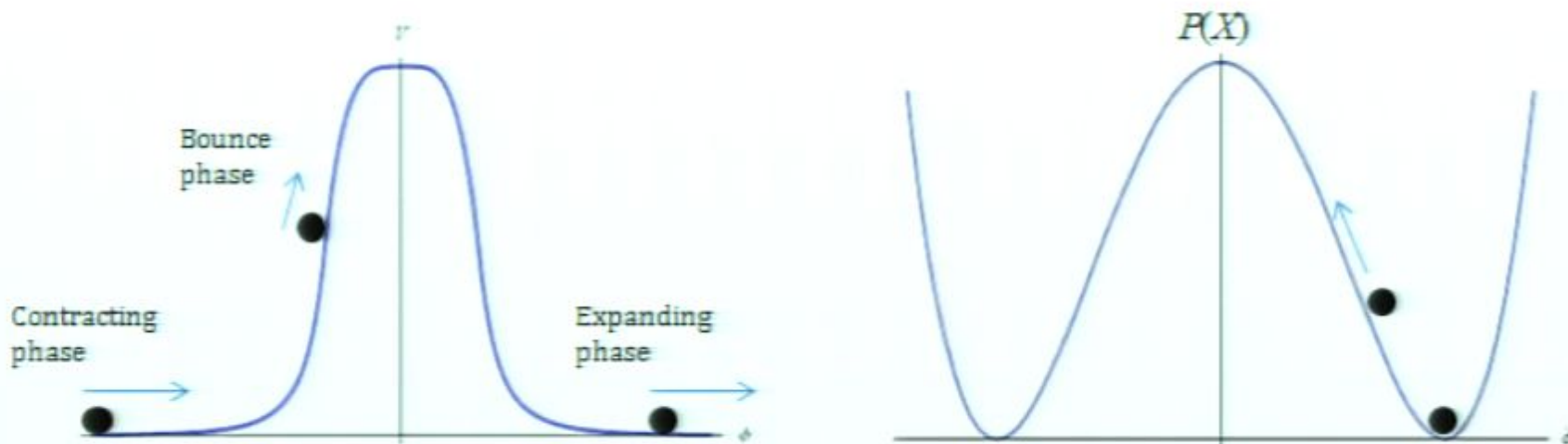


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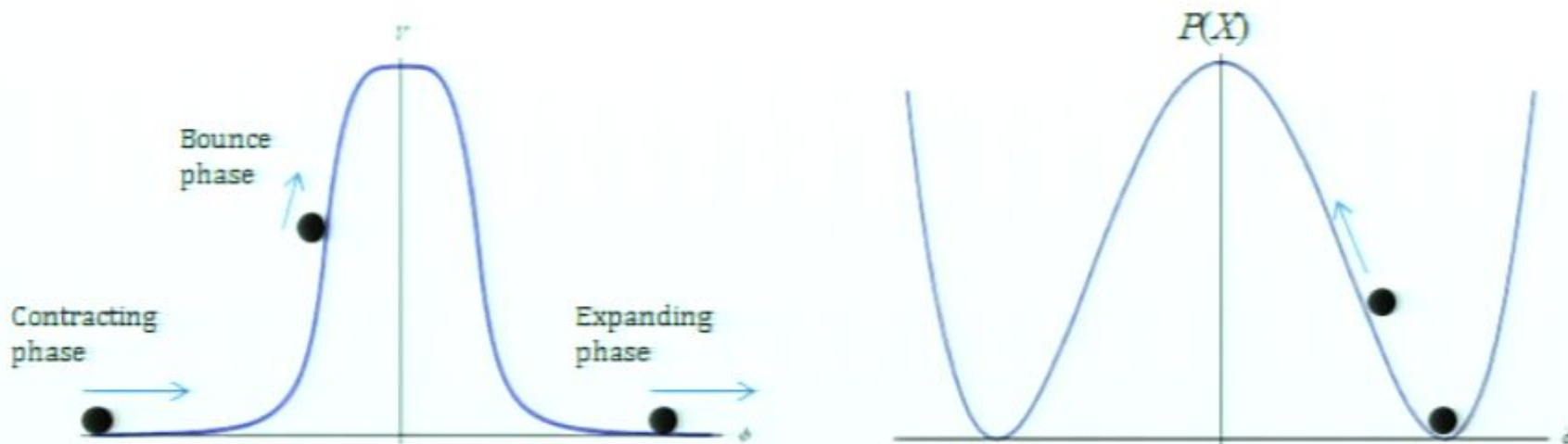


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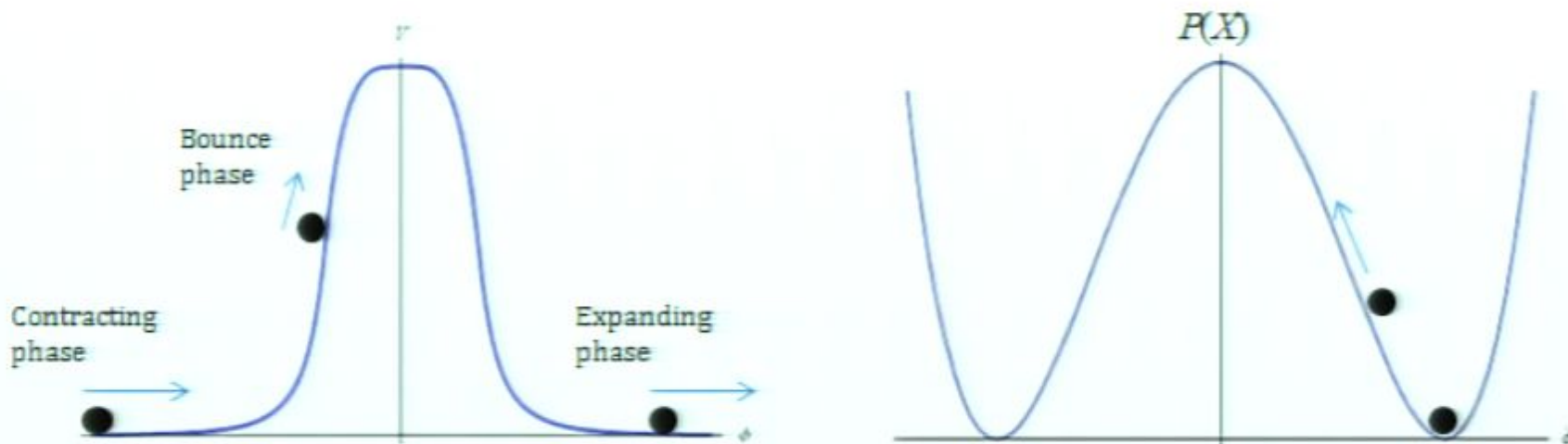


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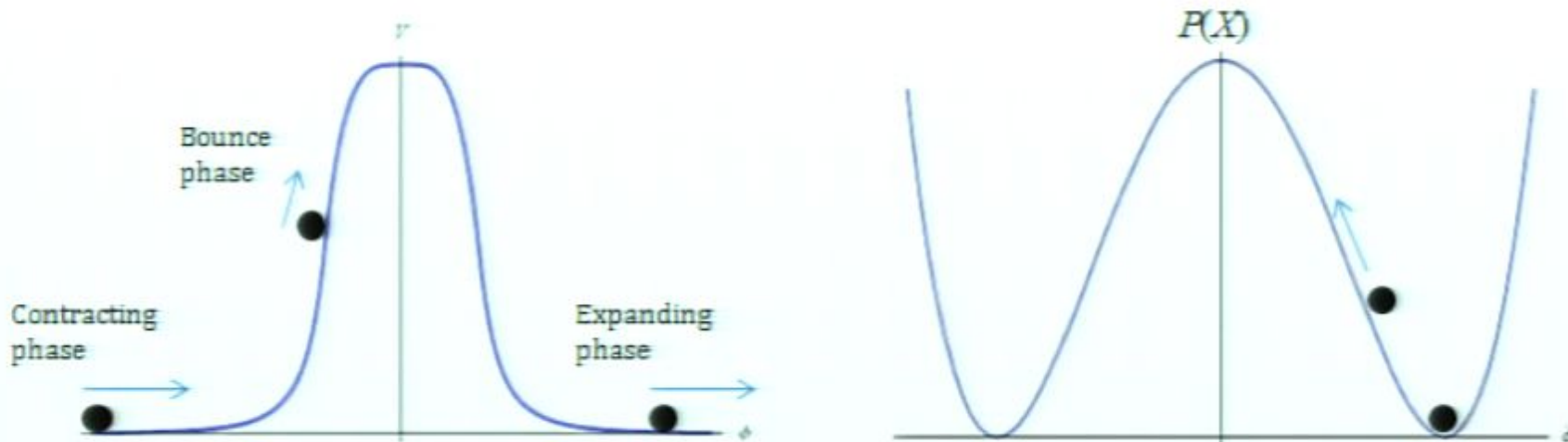


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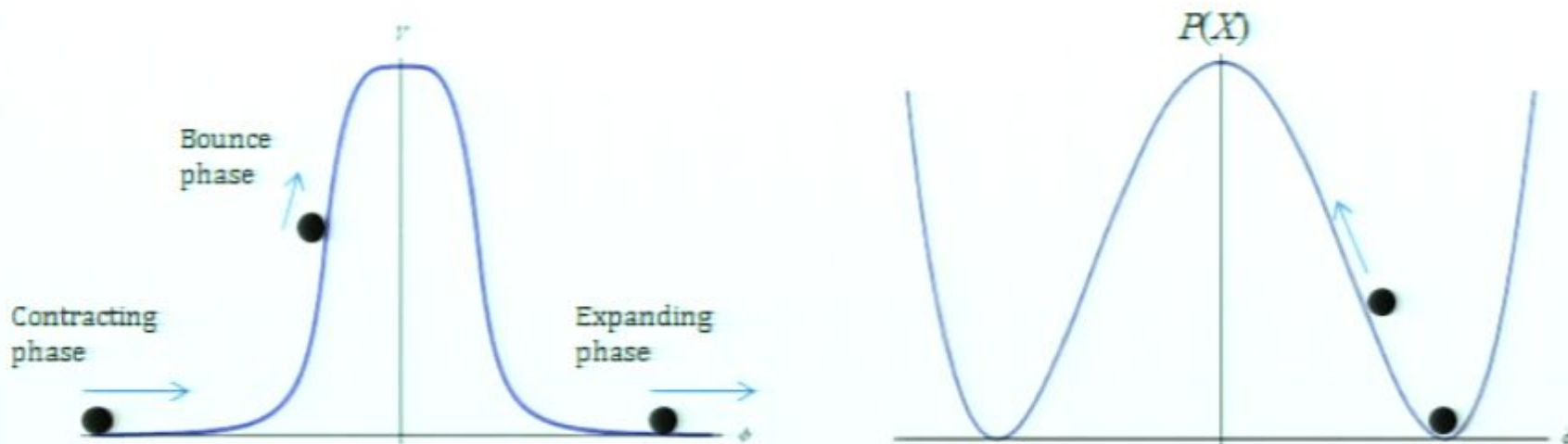


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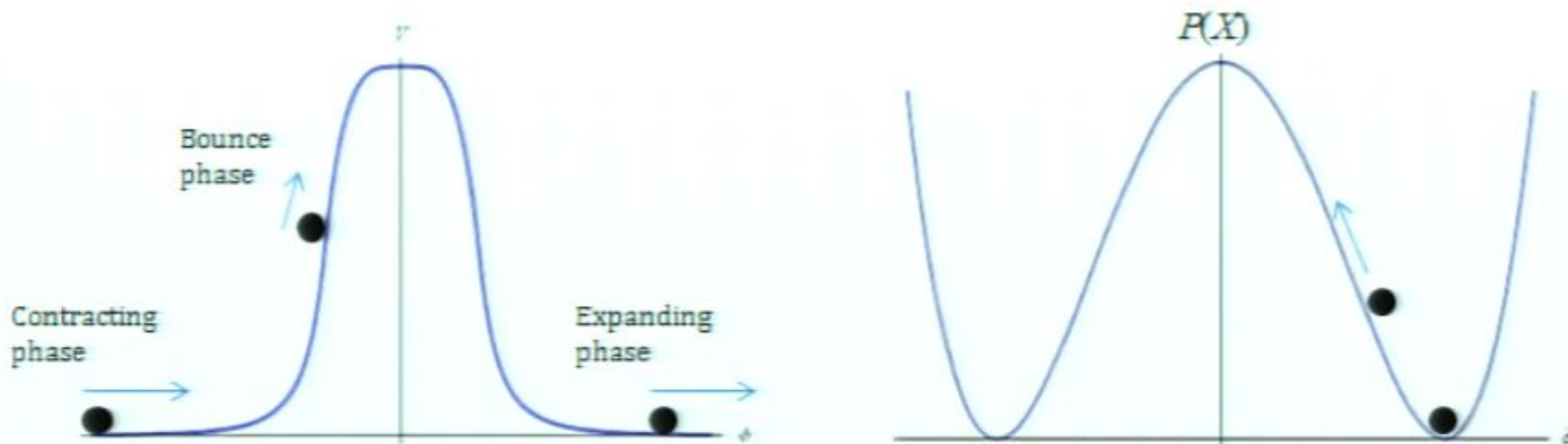


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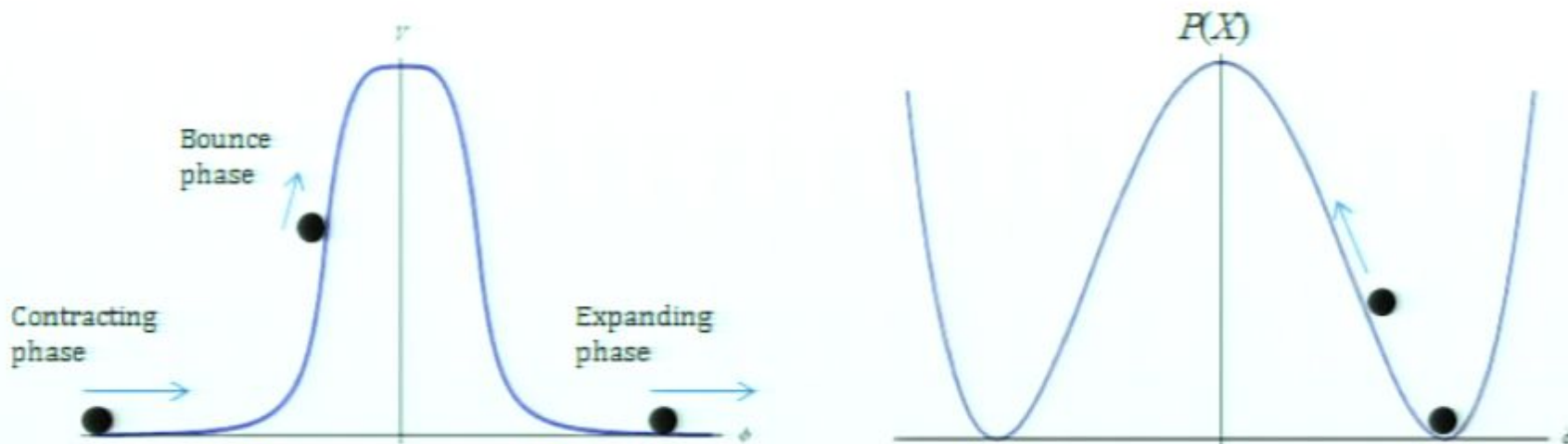


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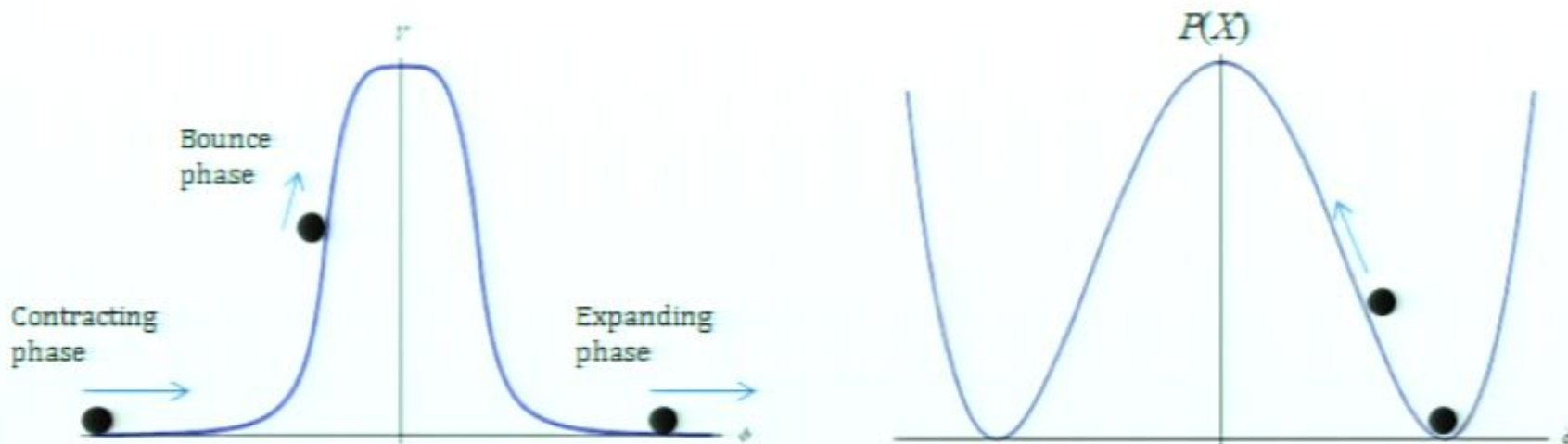


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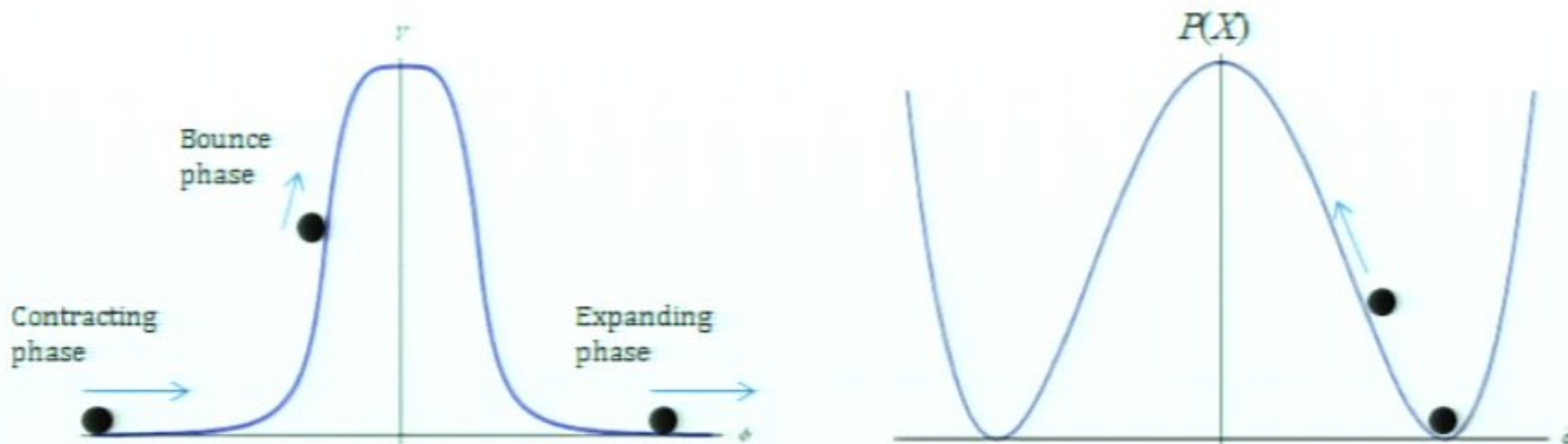


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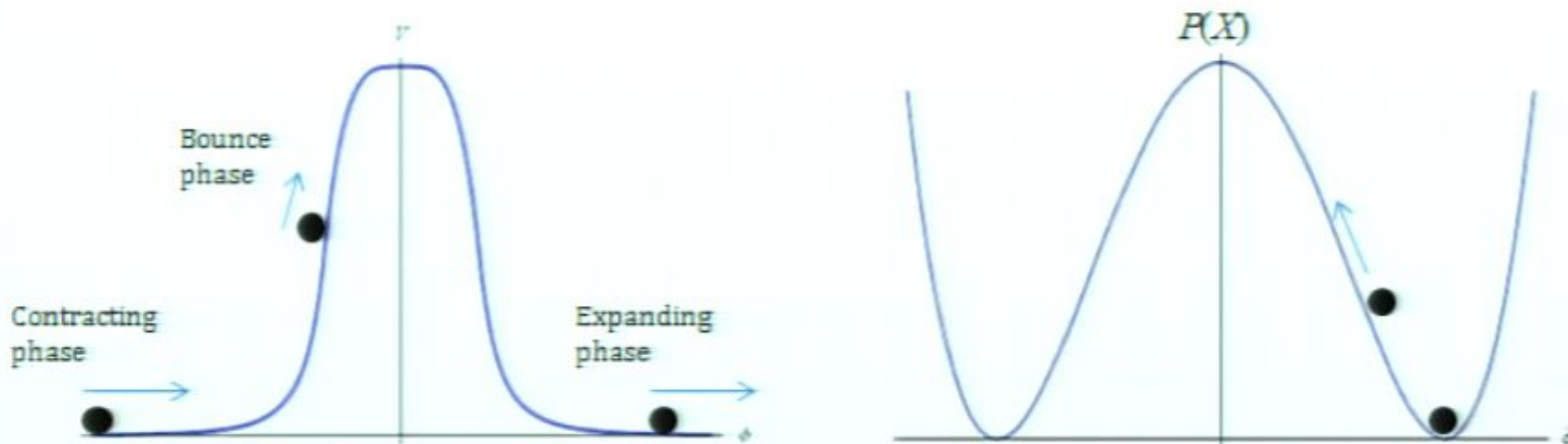


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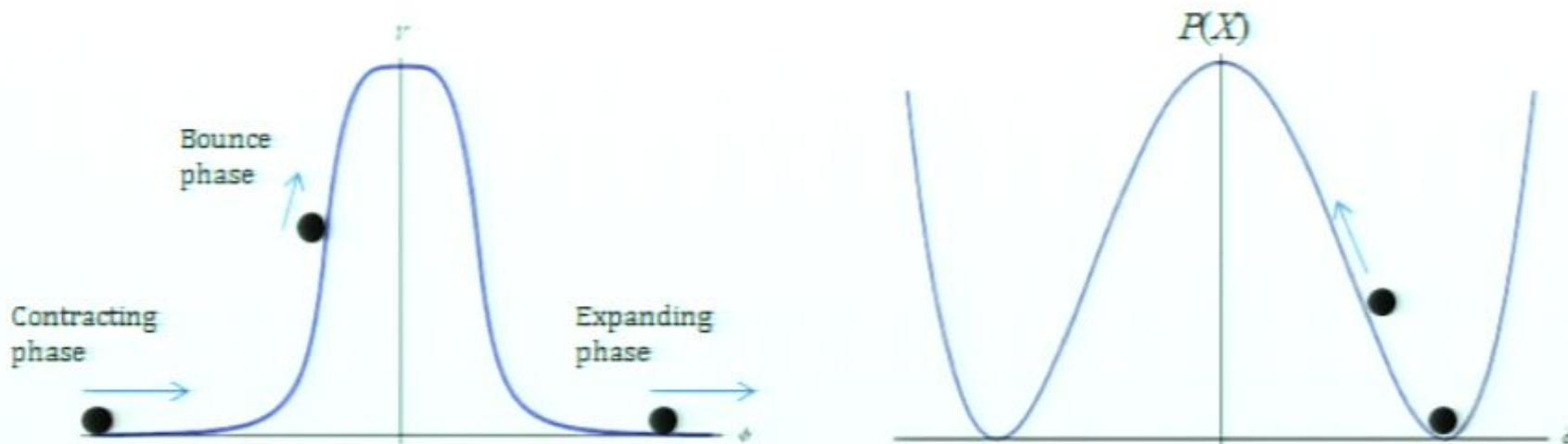


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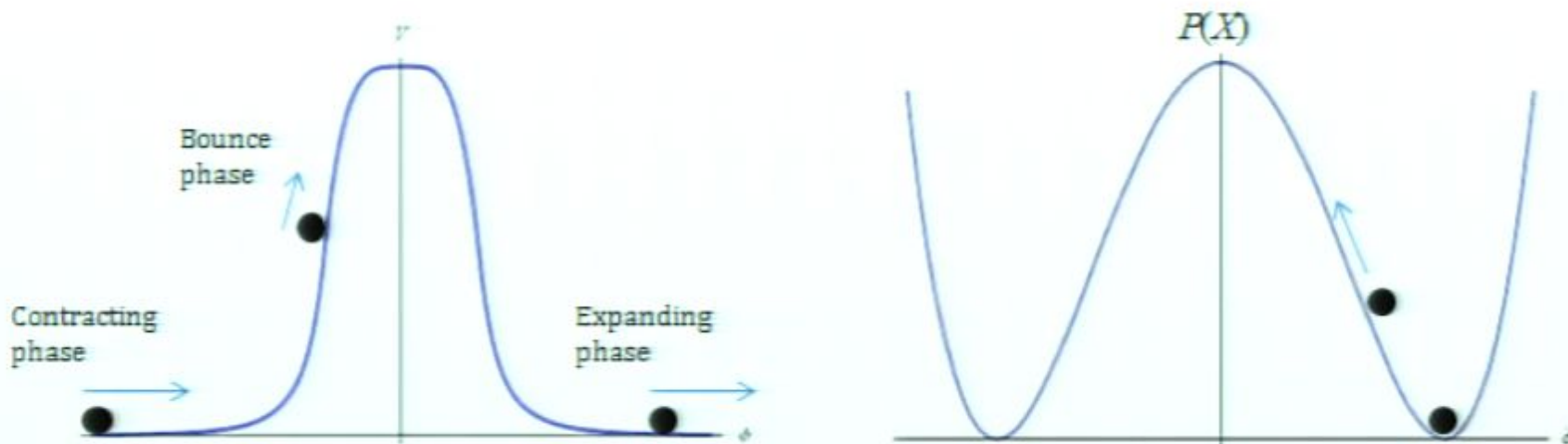


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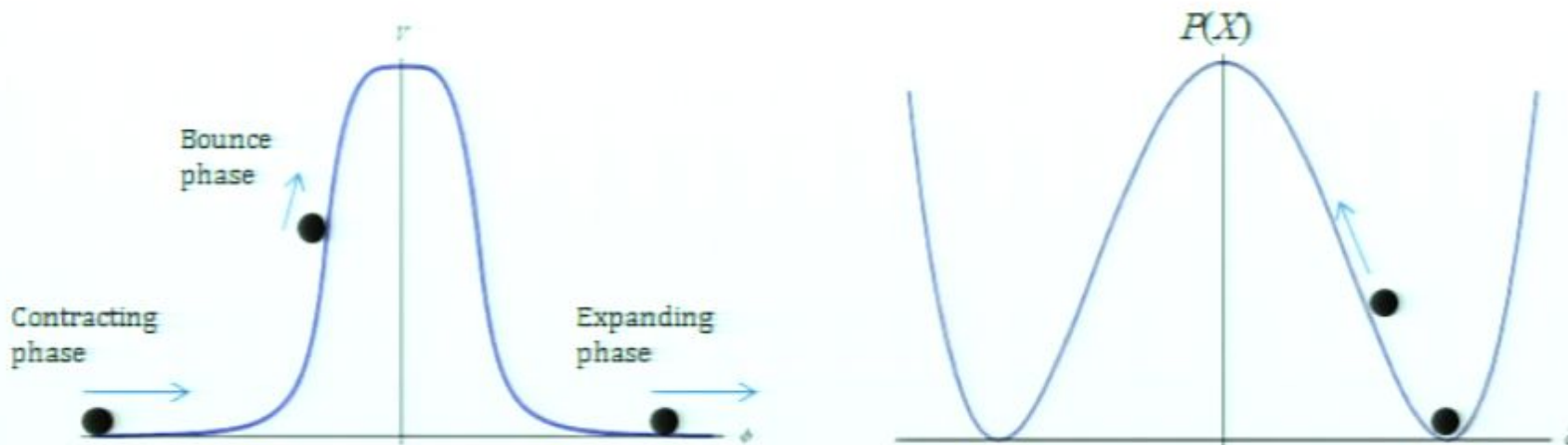


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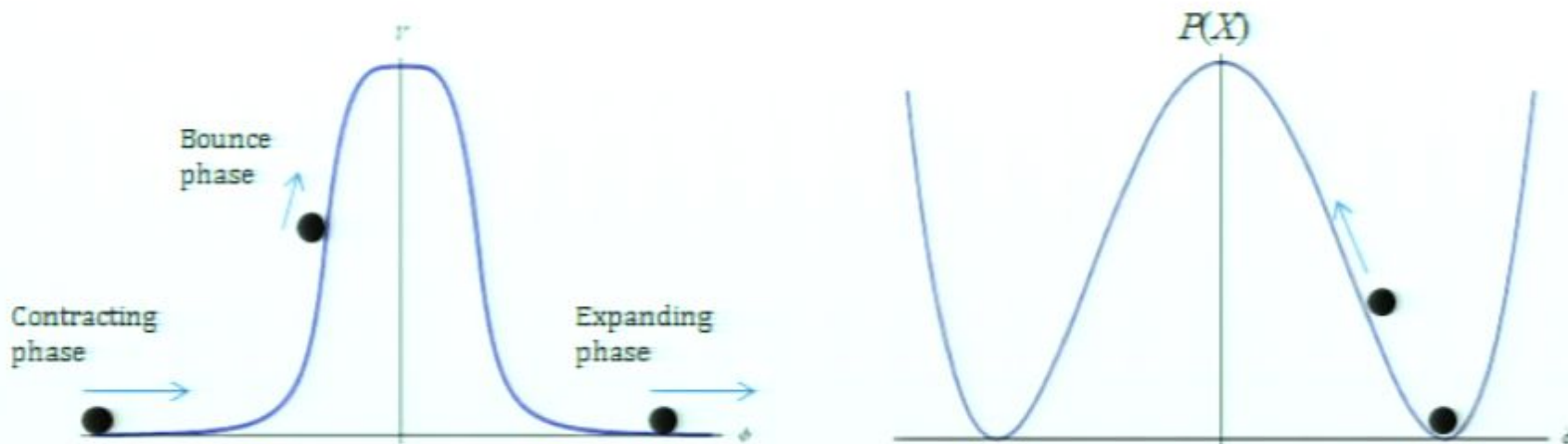


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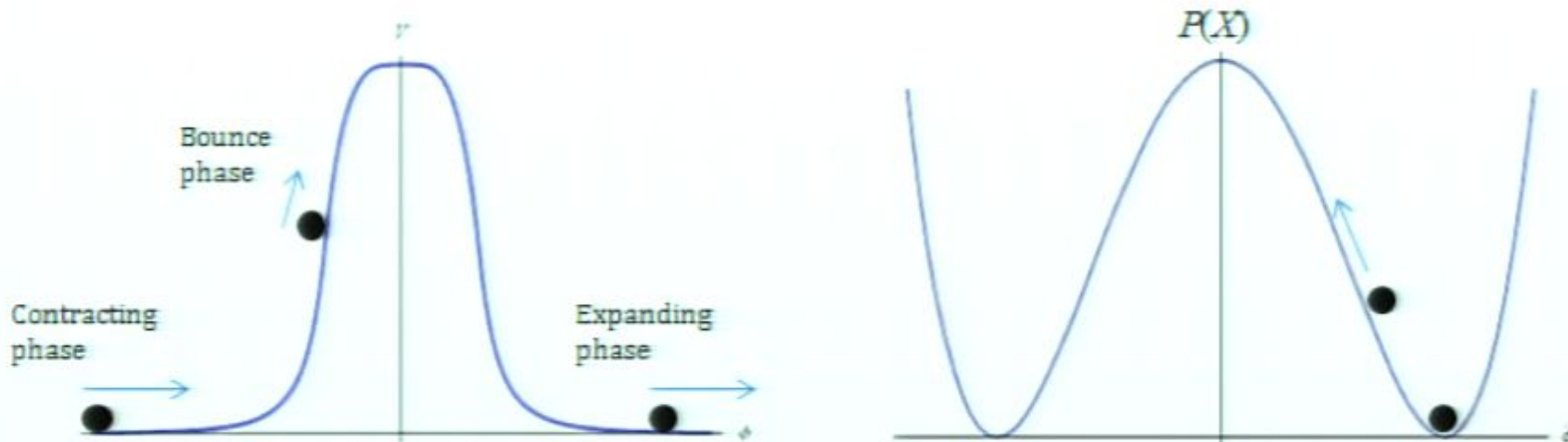


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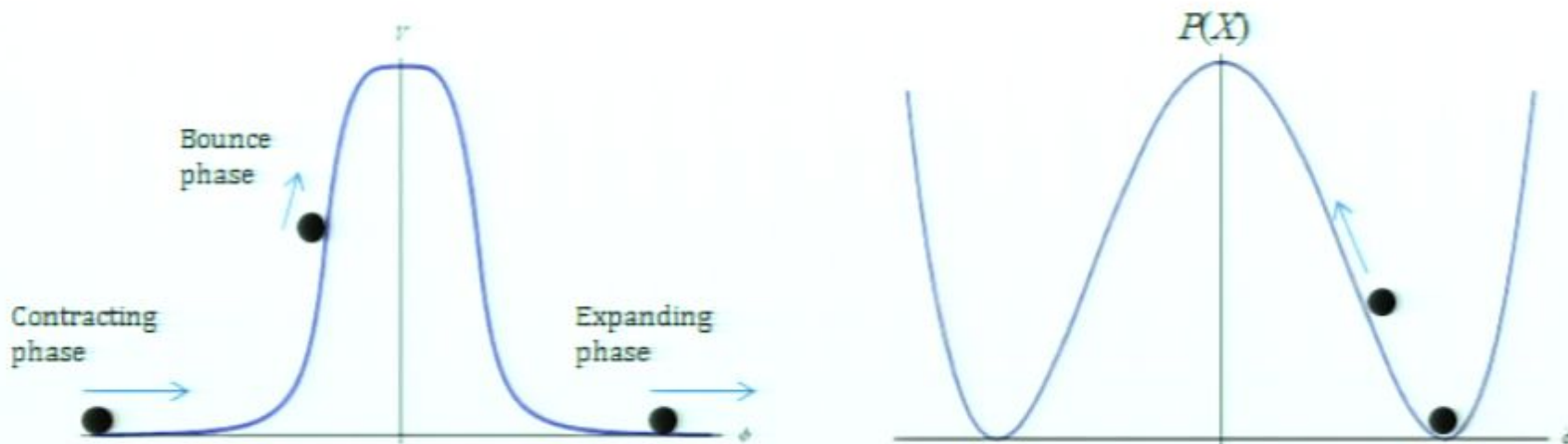


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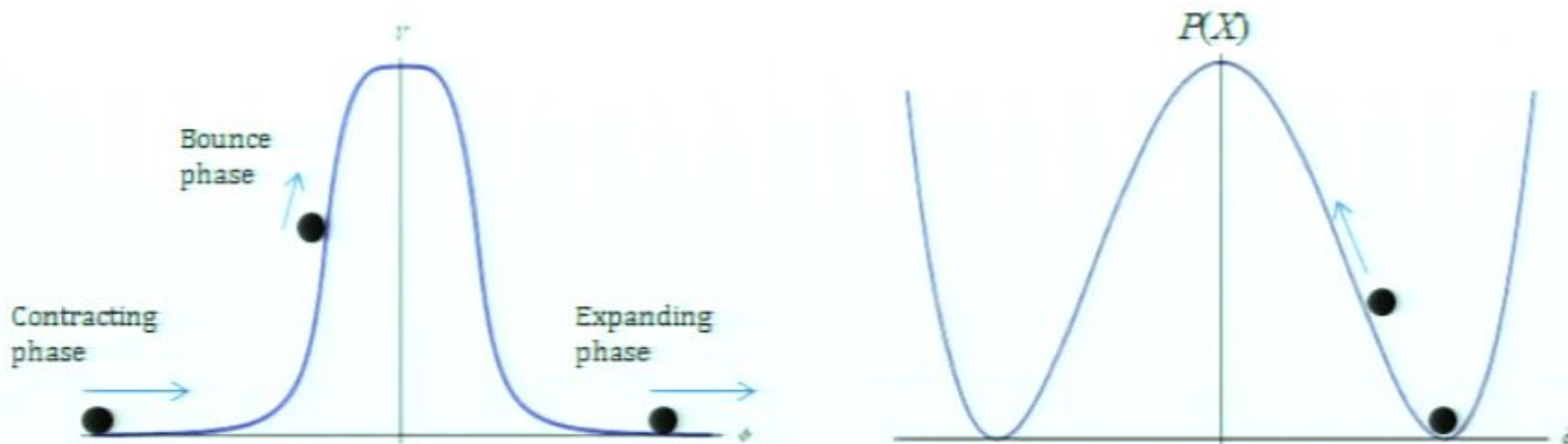


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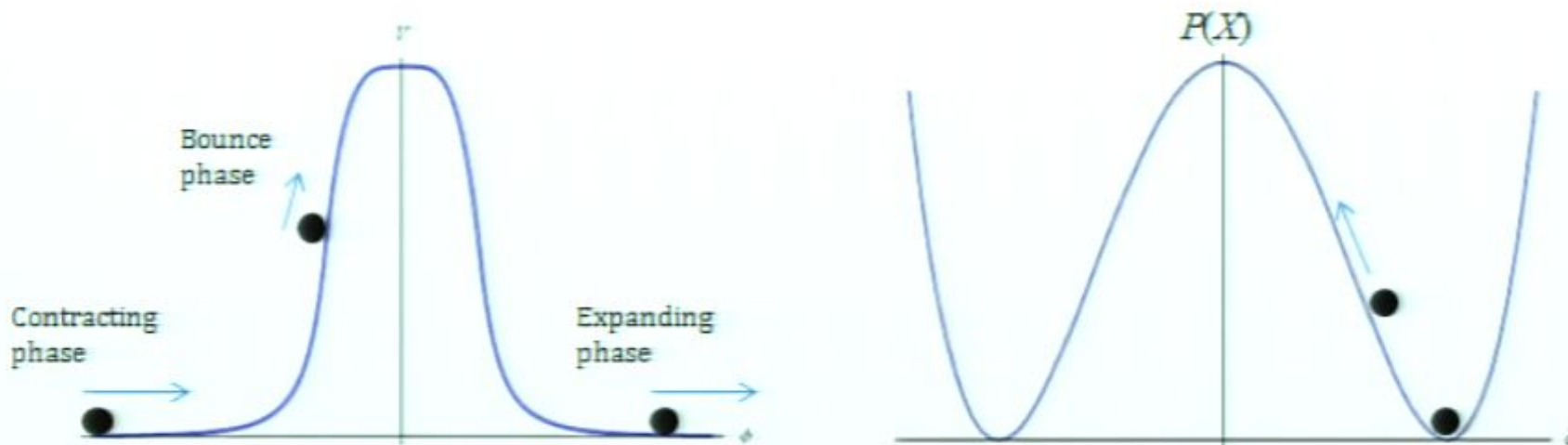


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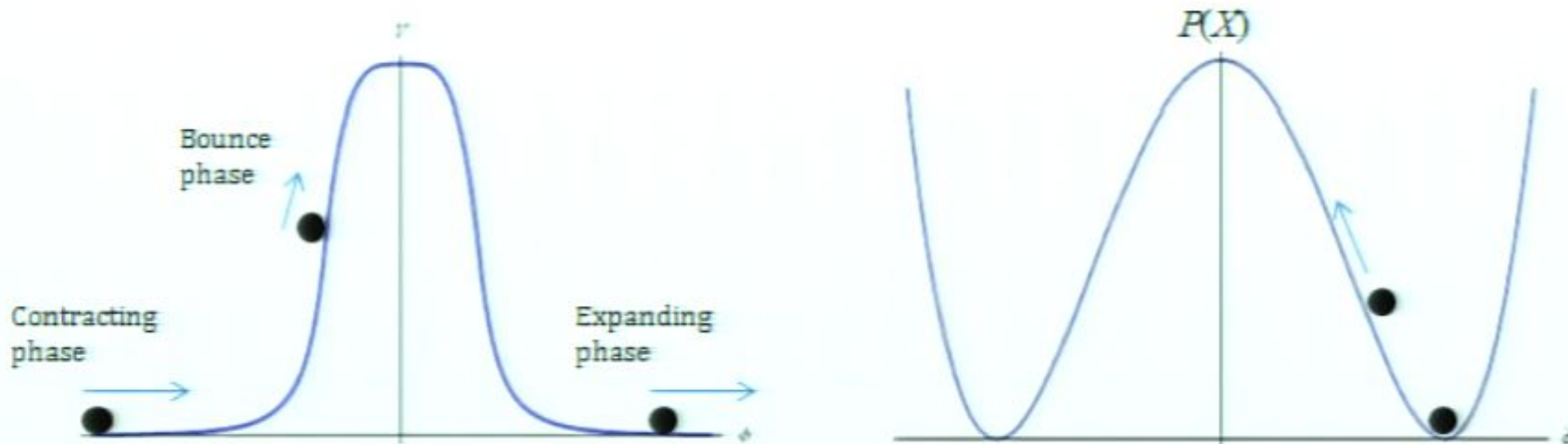


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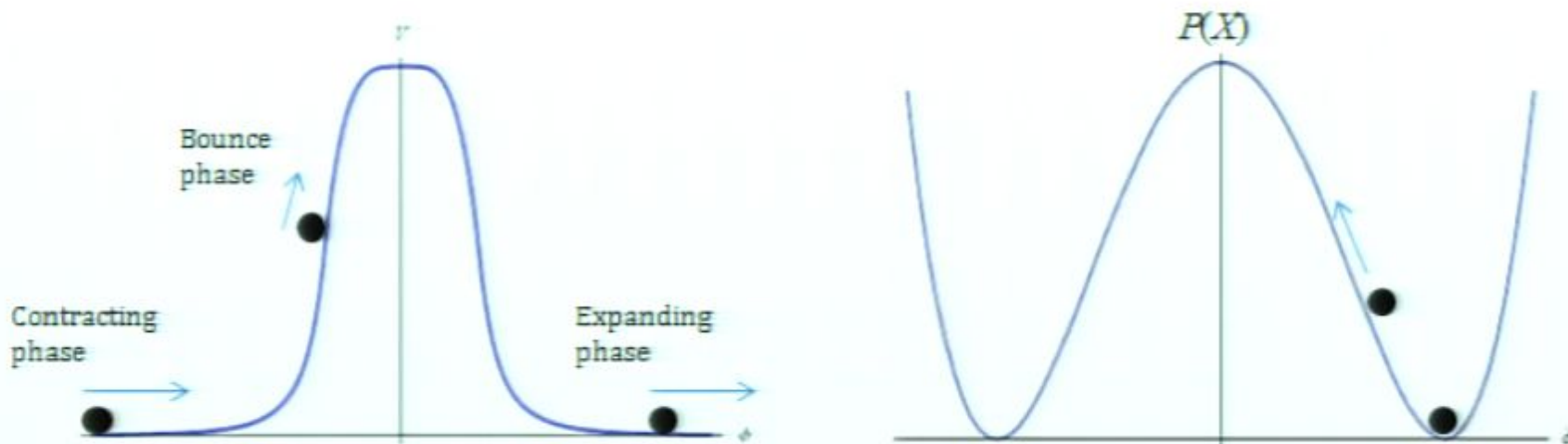


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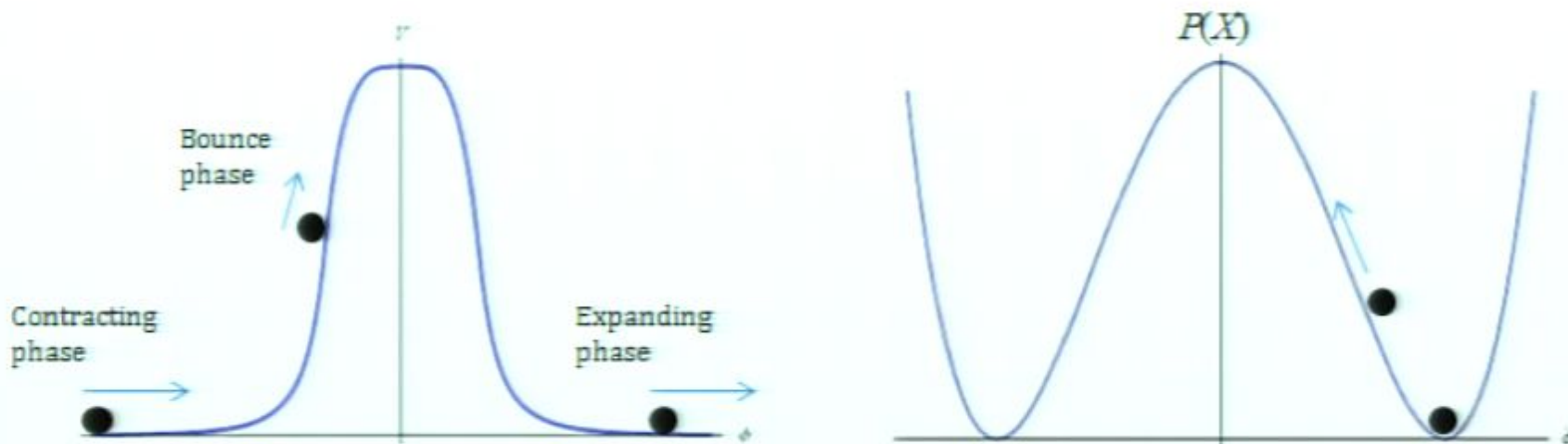


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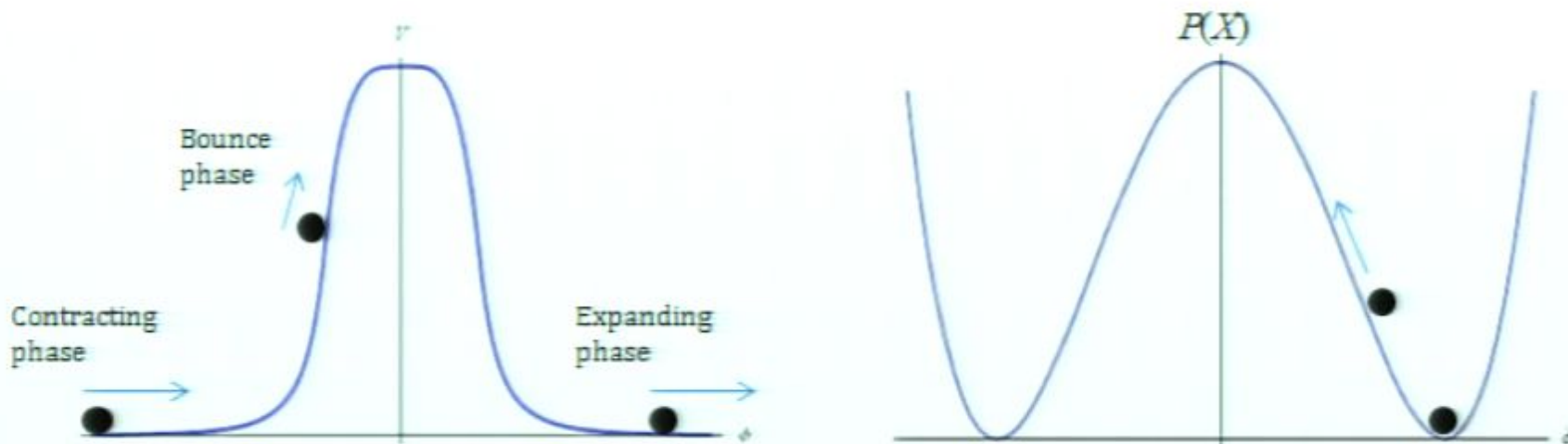


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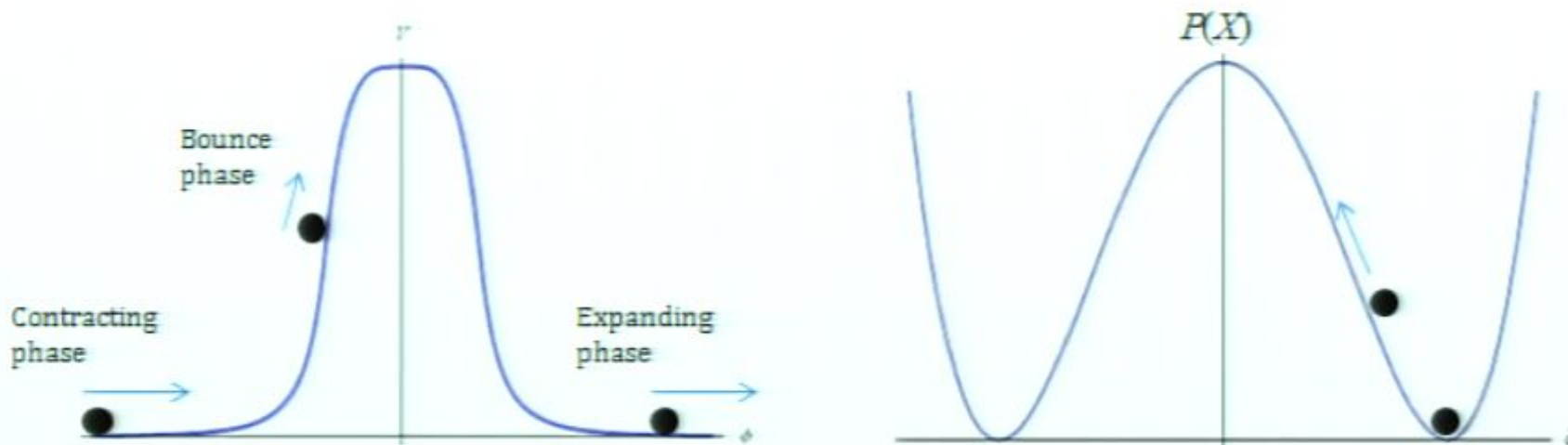


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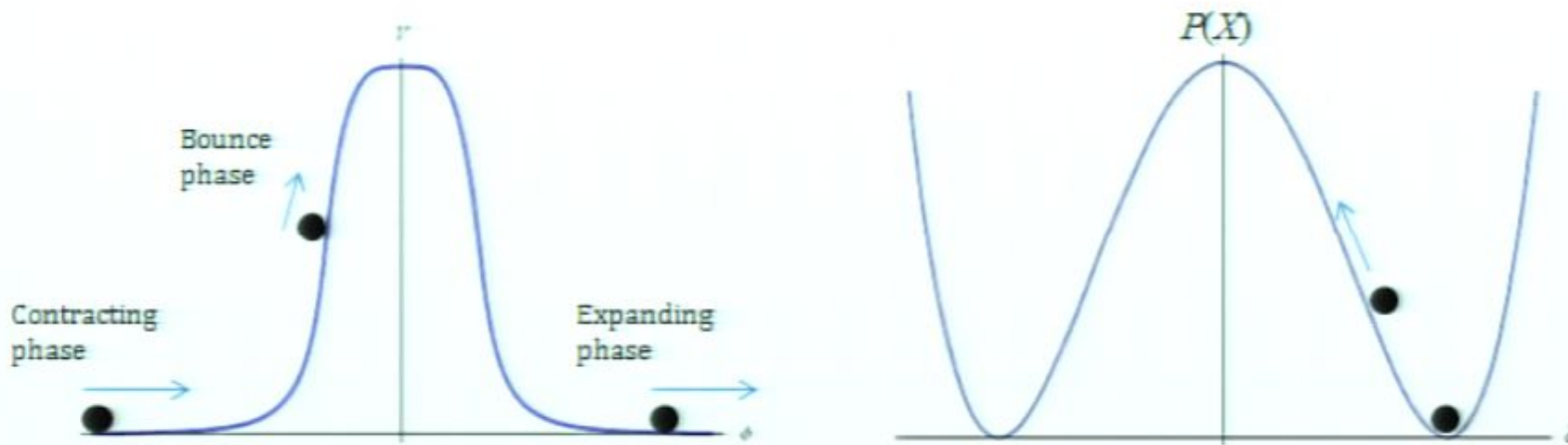


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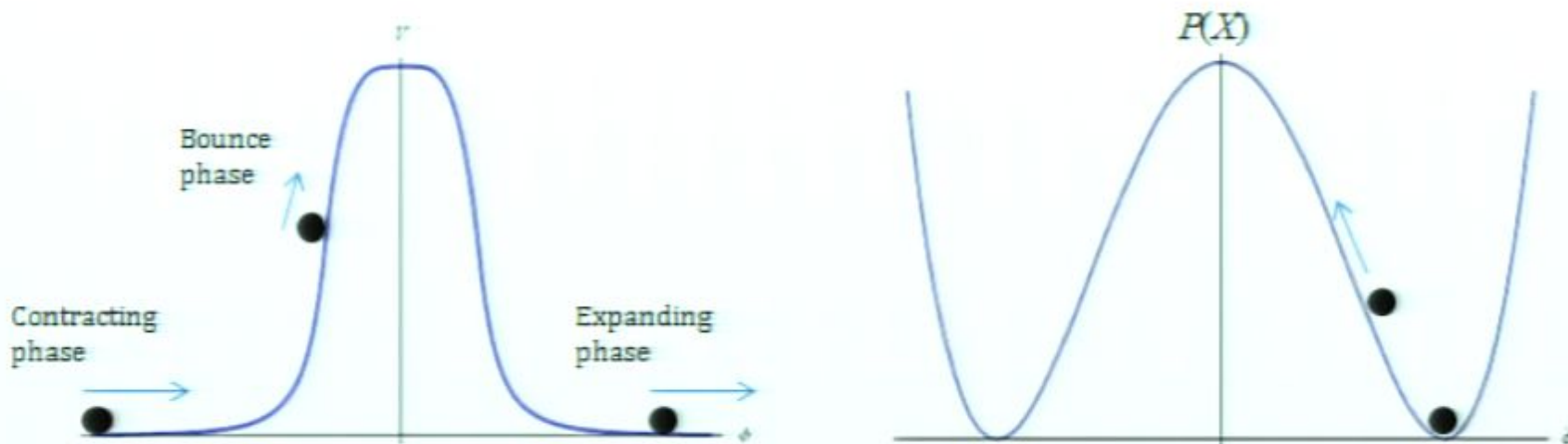


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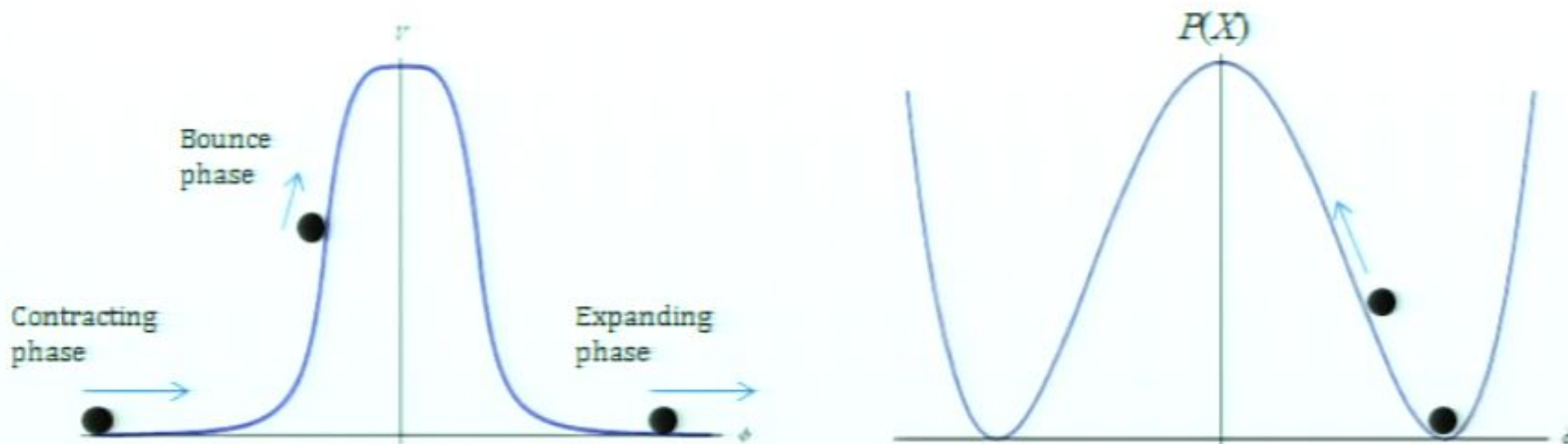


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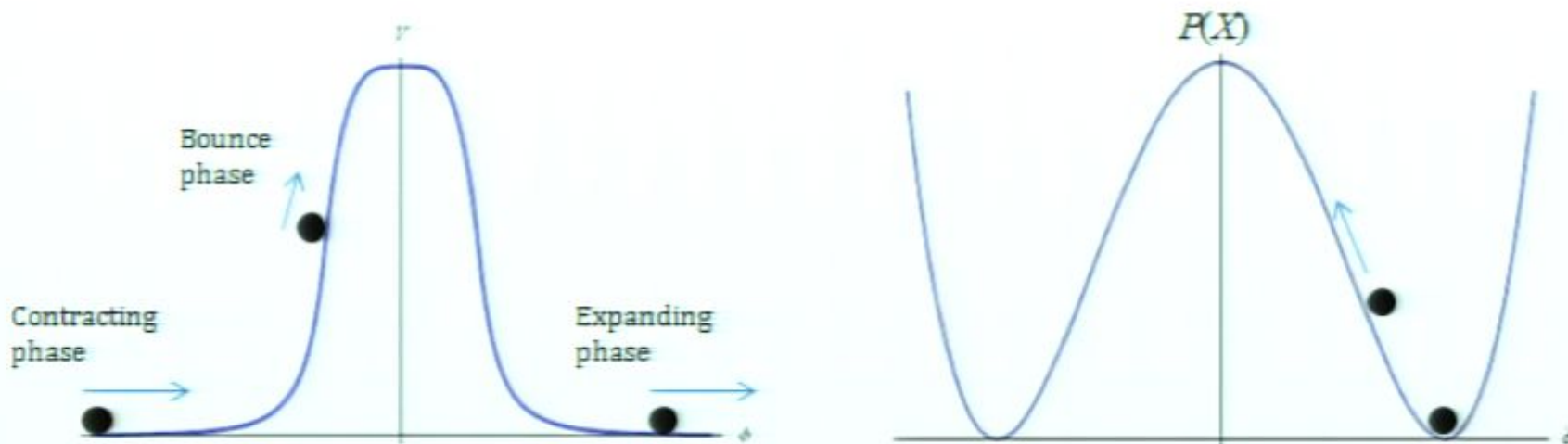


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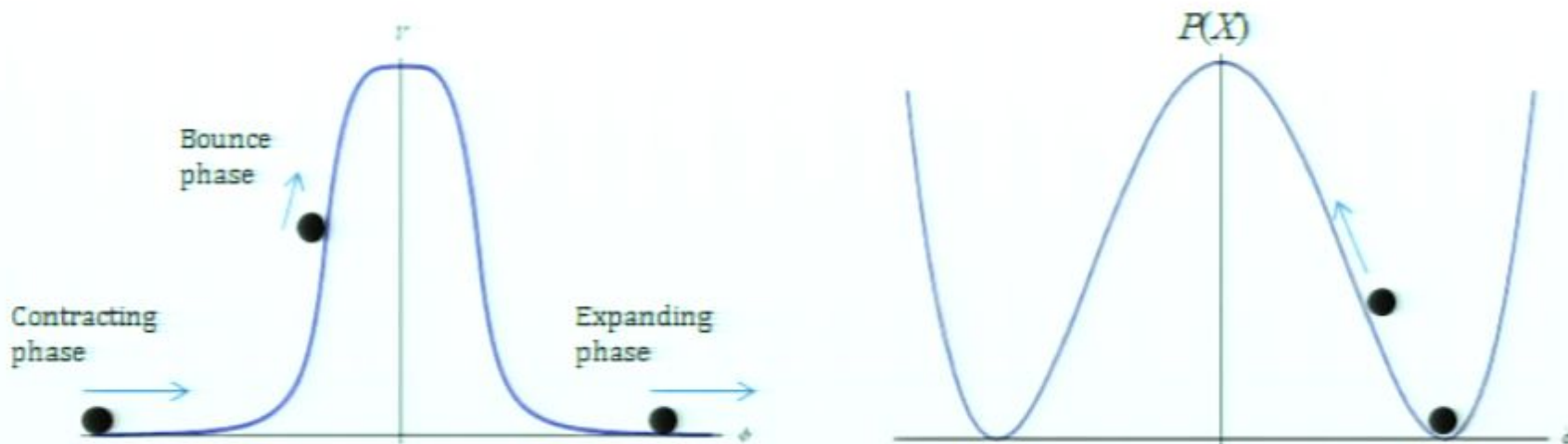


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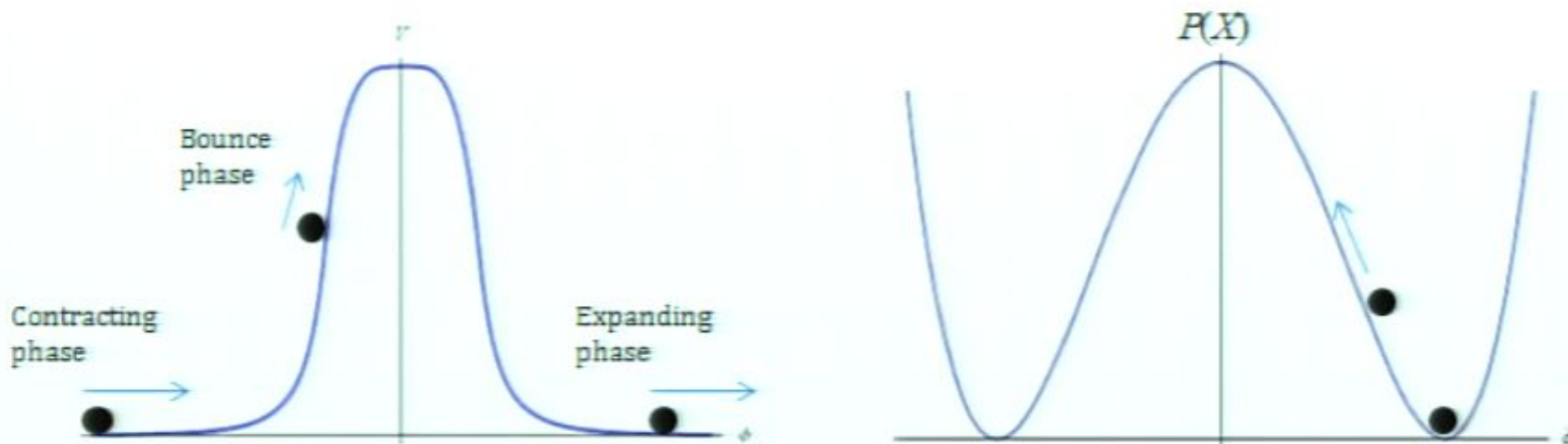


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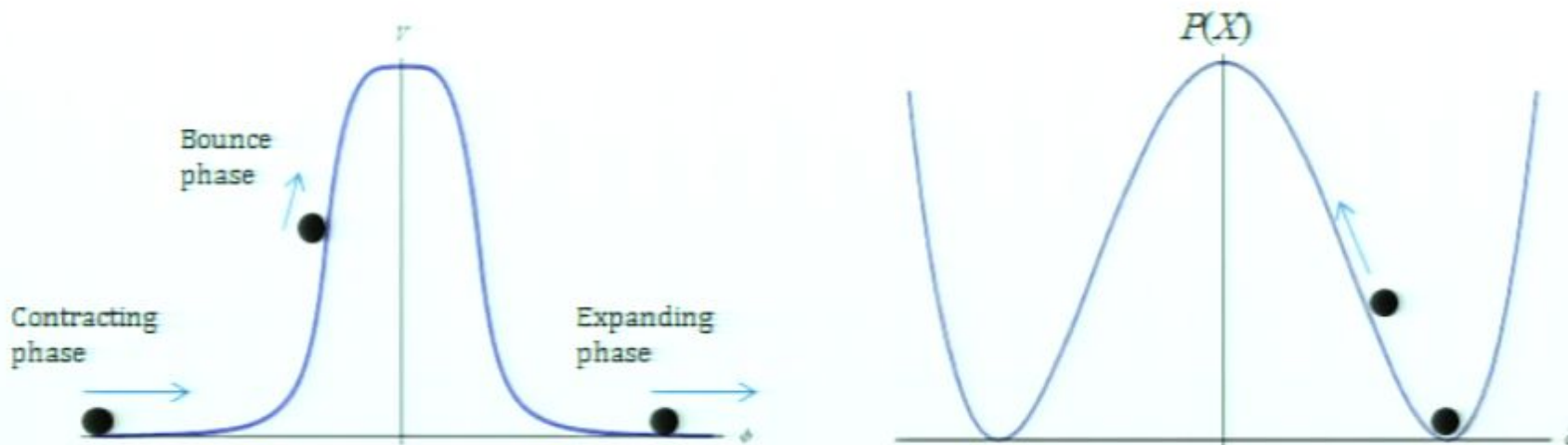


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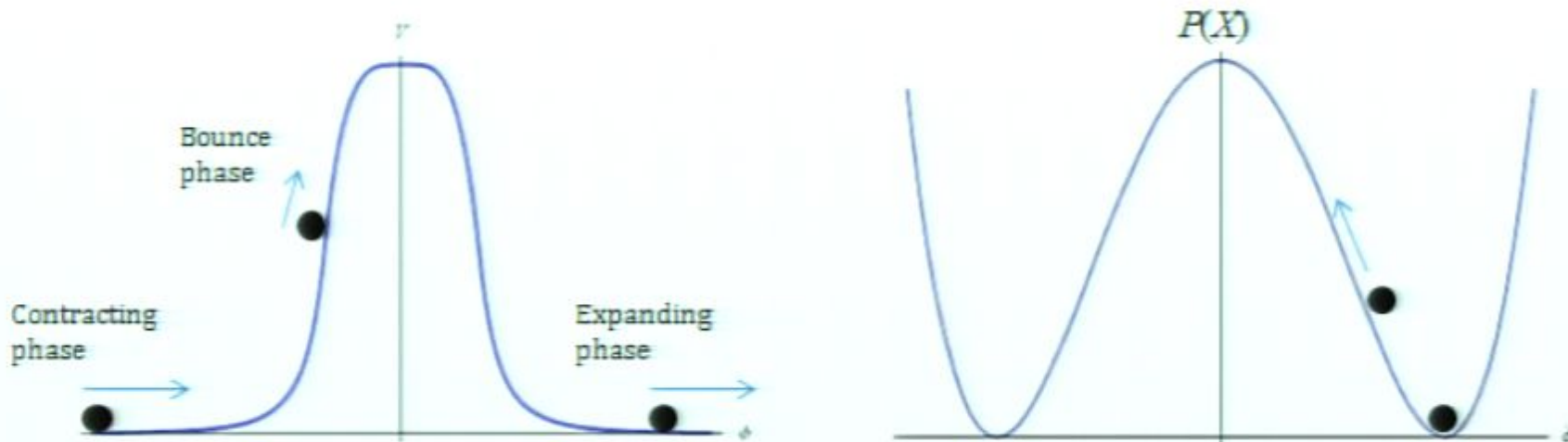


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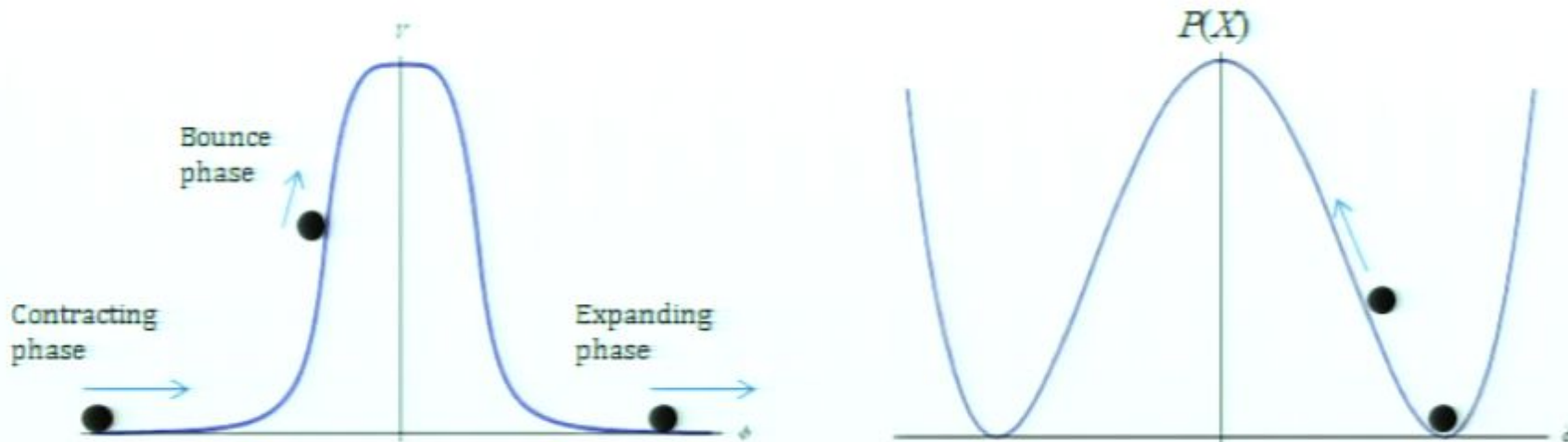


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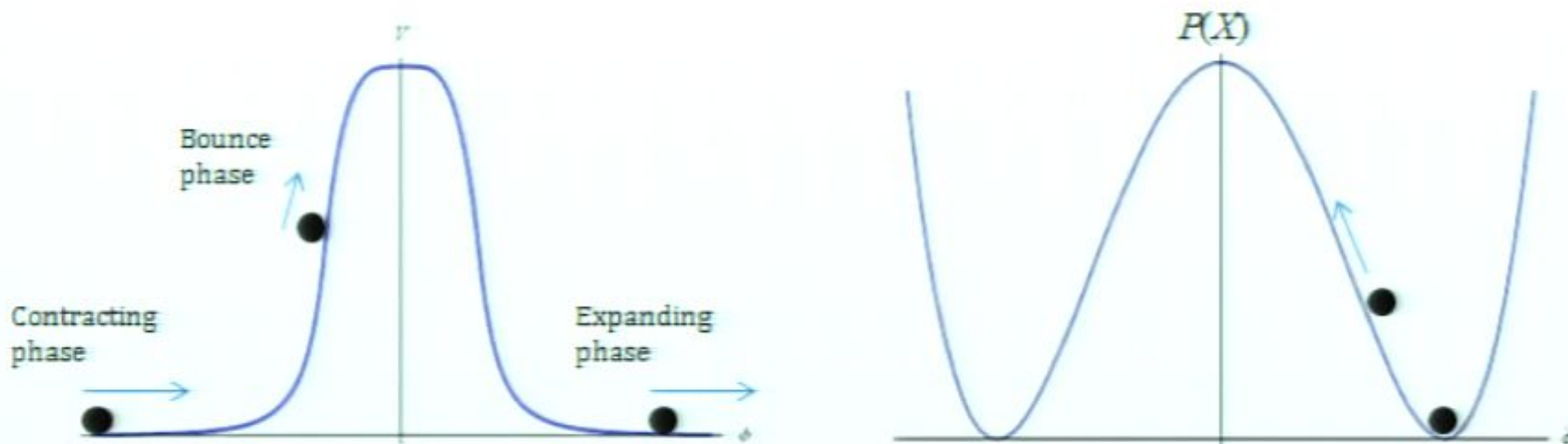


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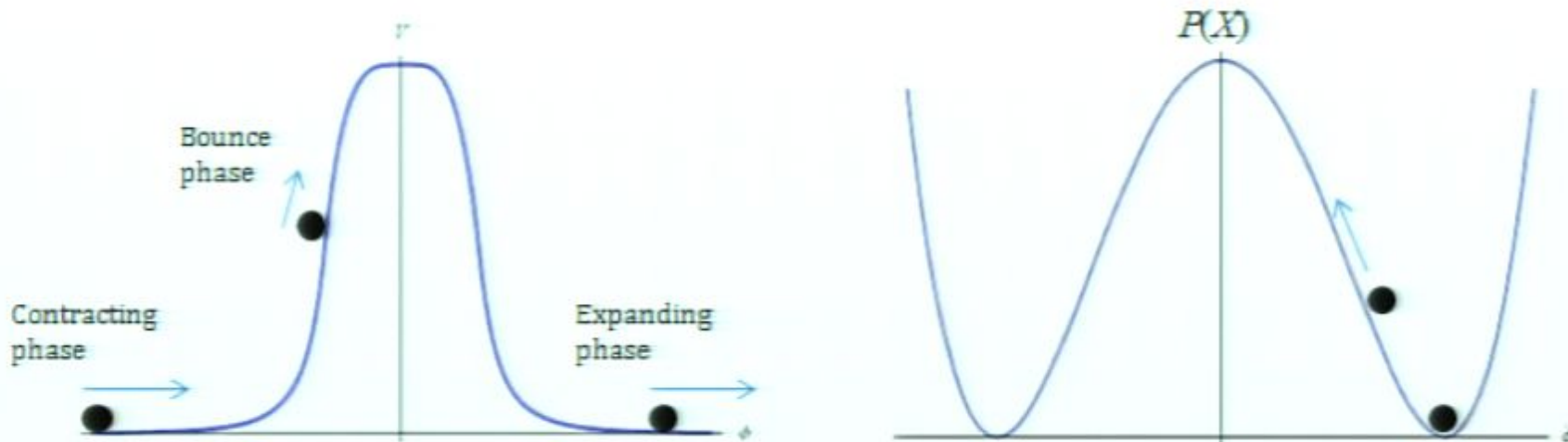


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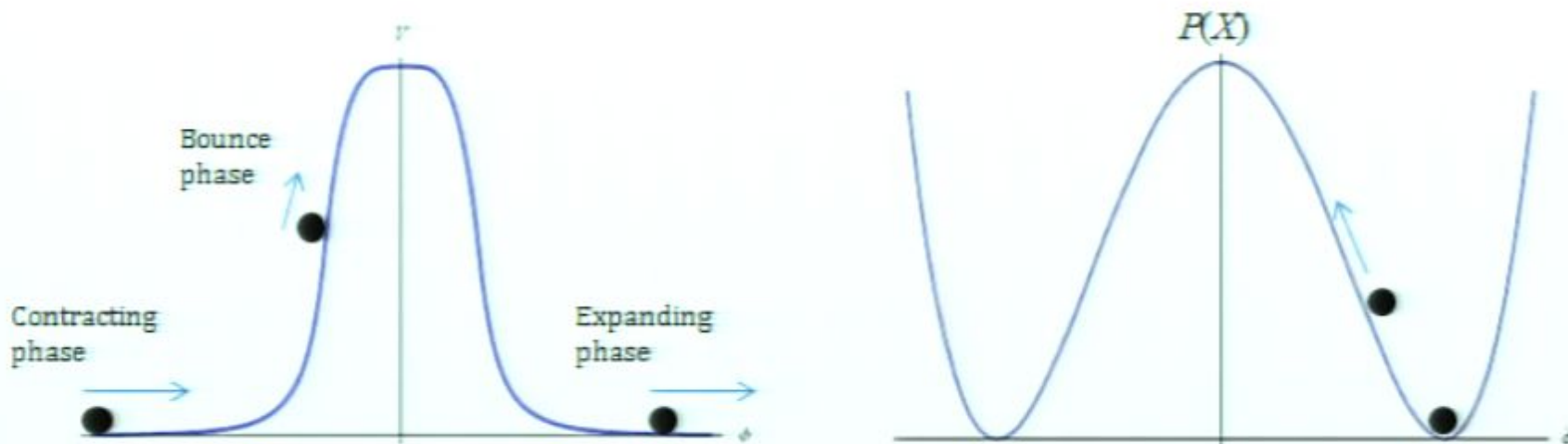


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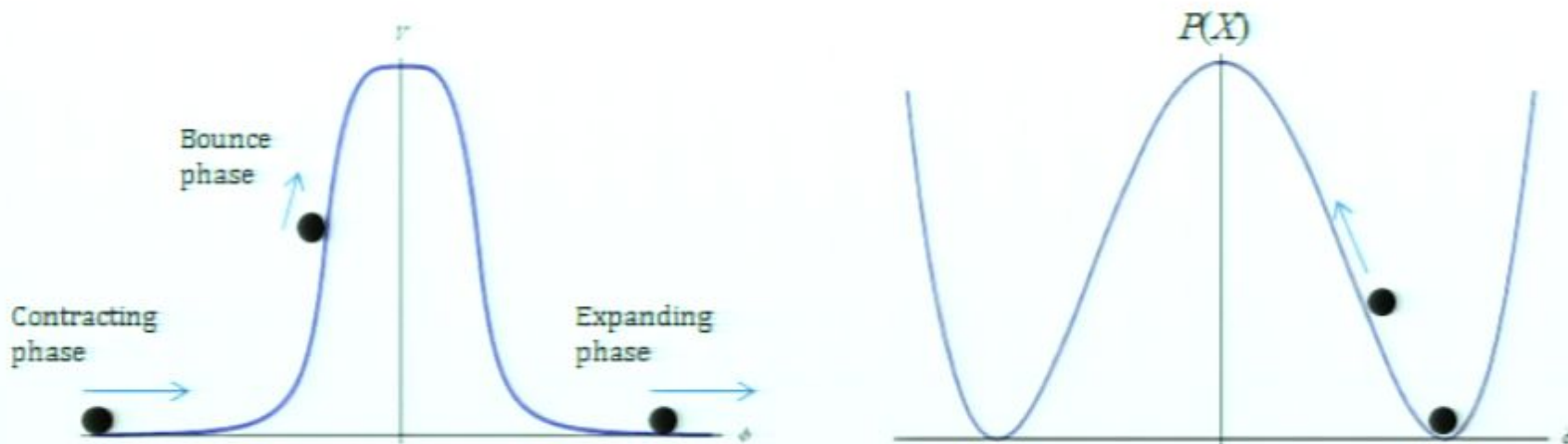


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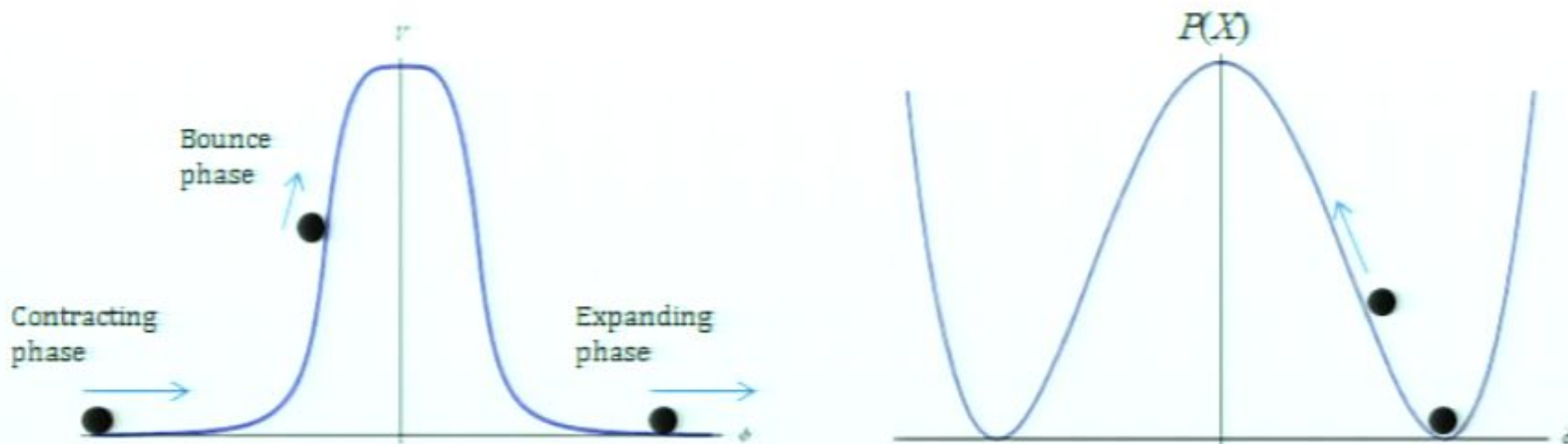


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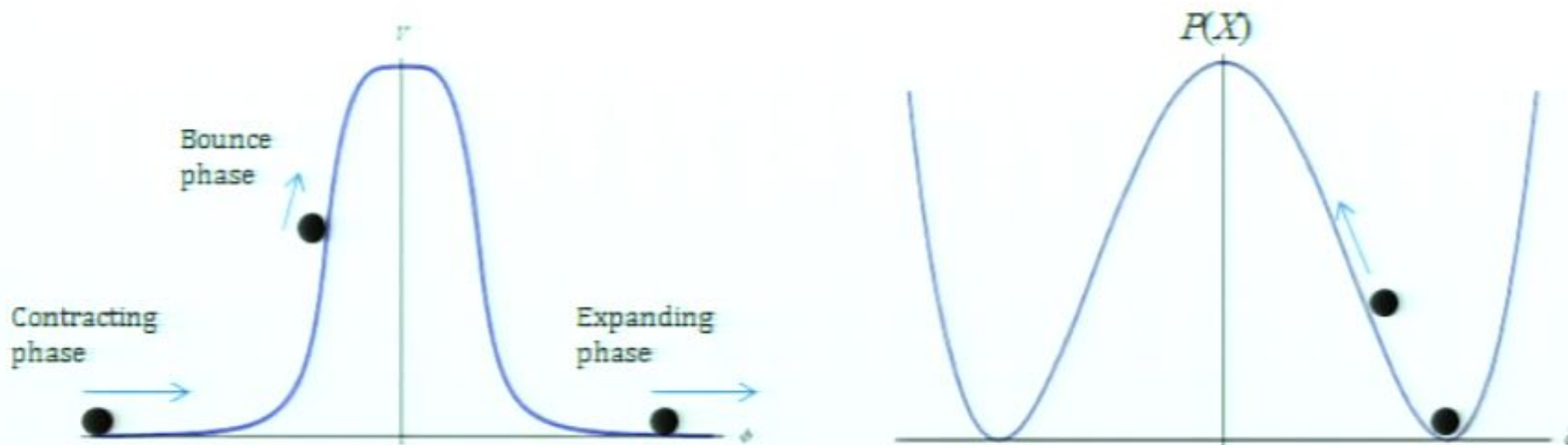


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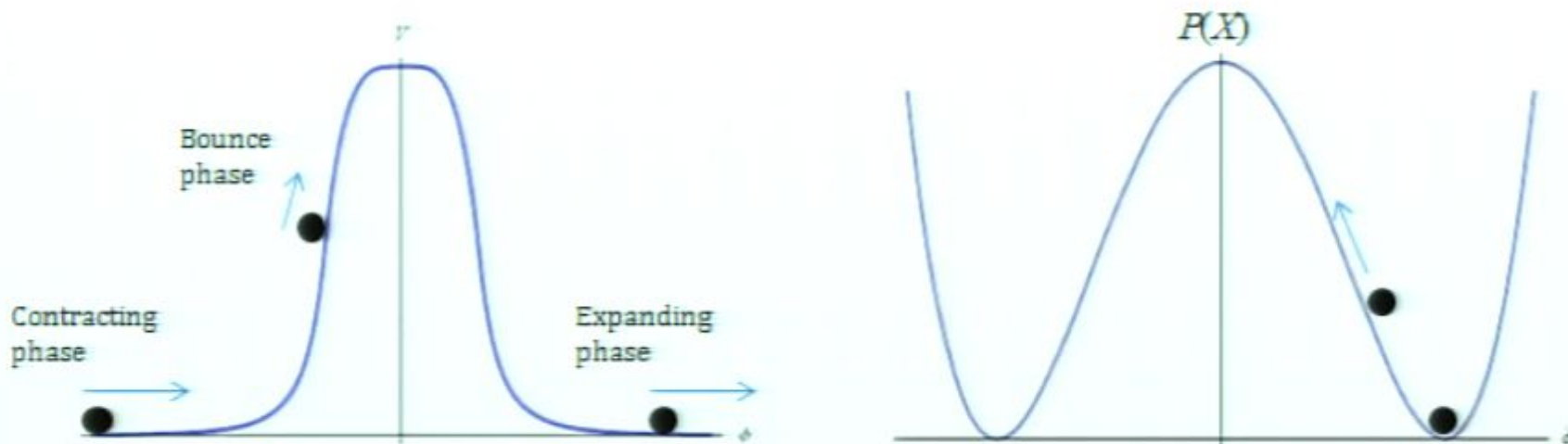


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Divergence is cut off at M^4

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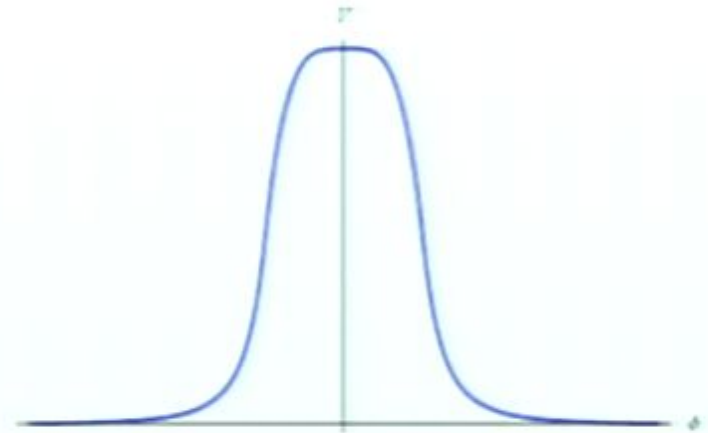
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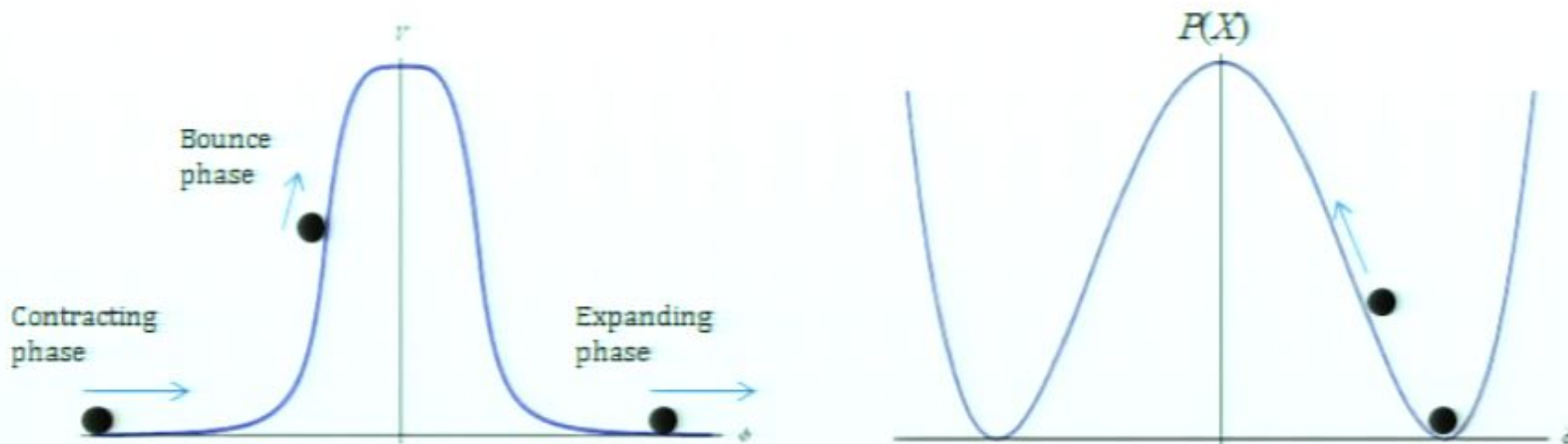
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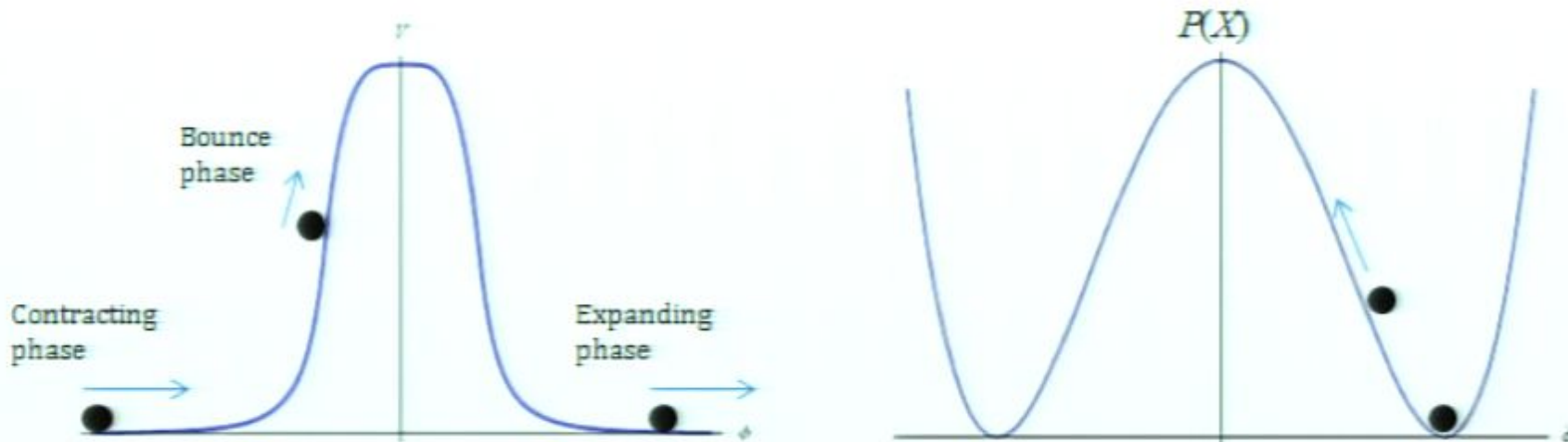


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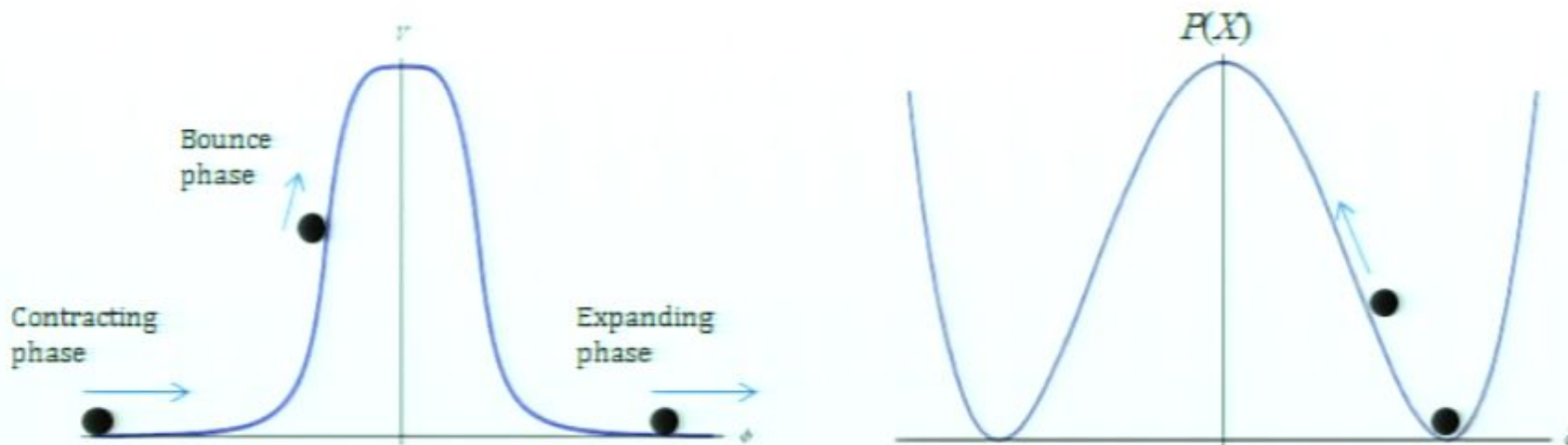


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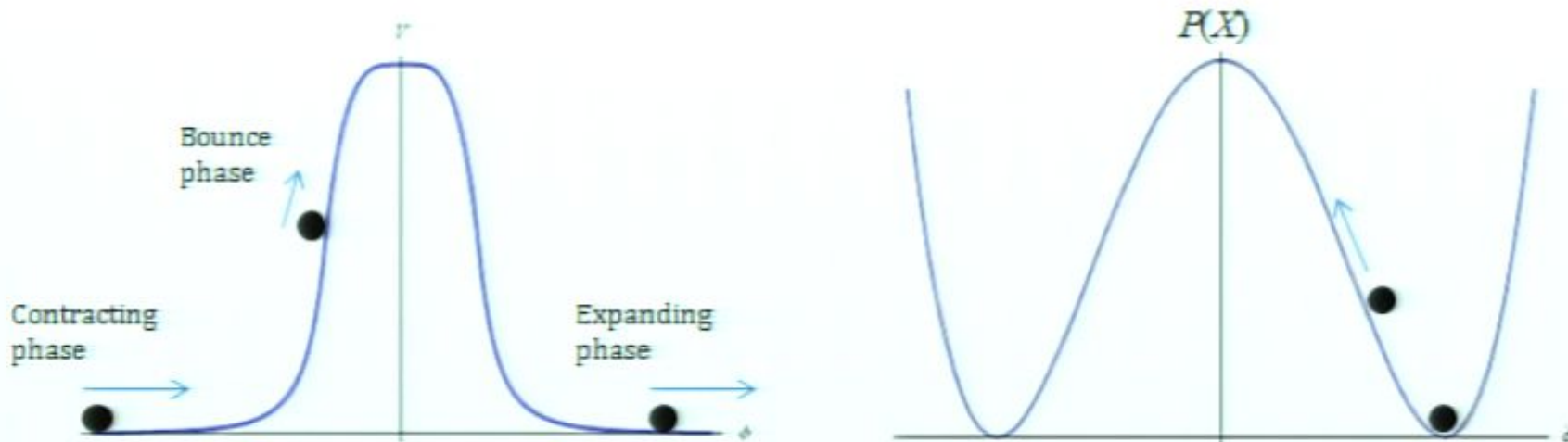


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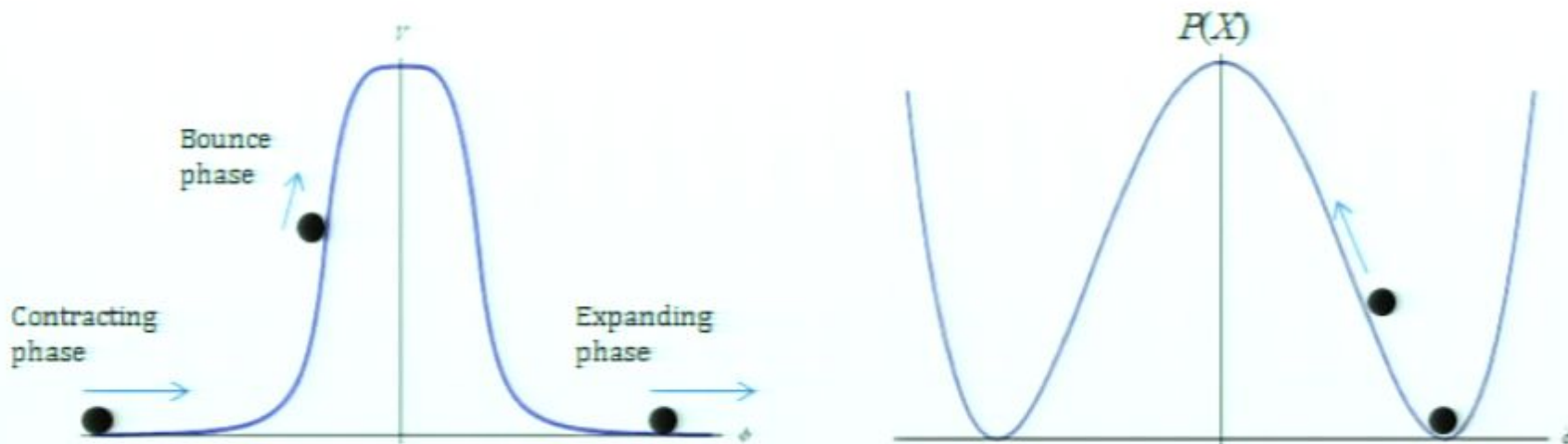


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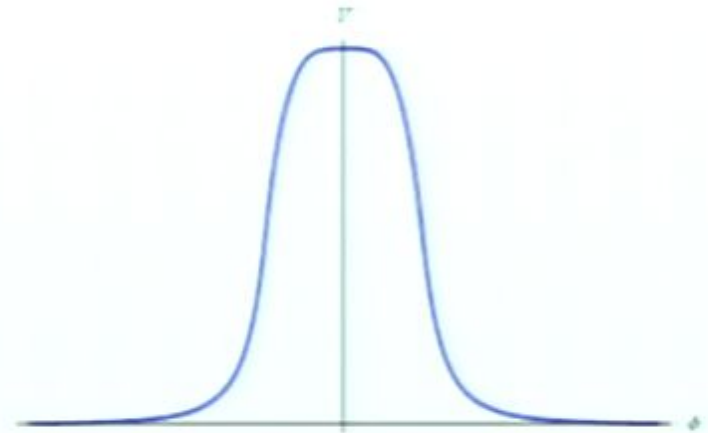
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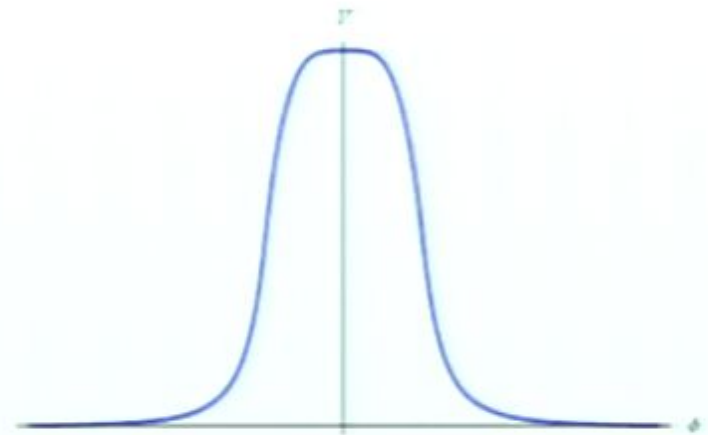
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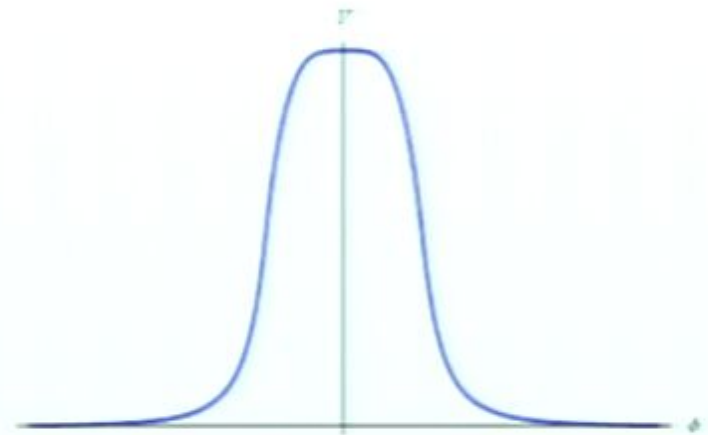
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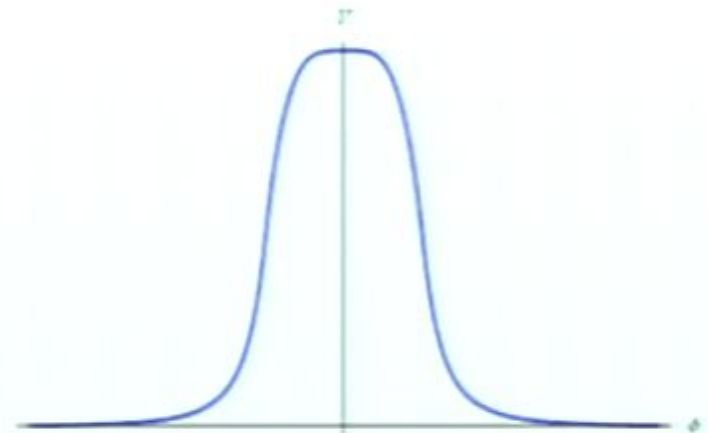
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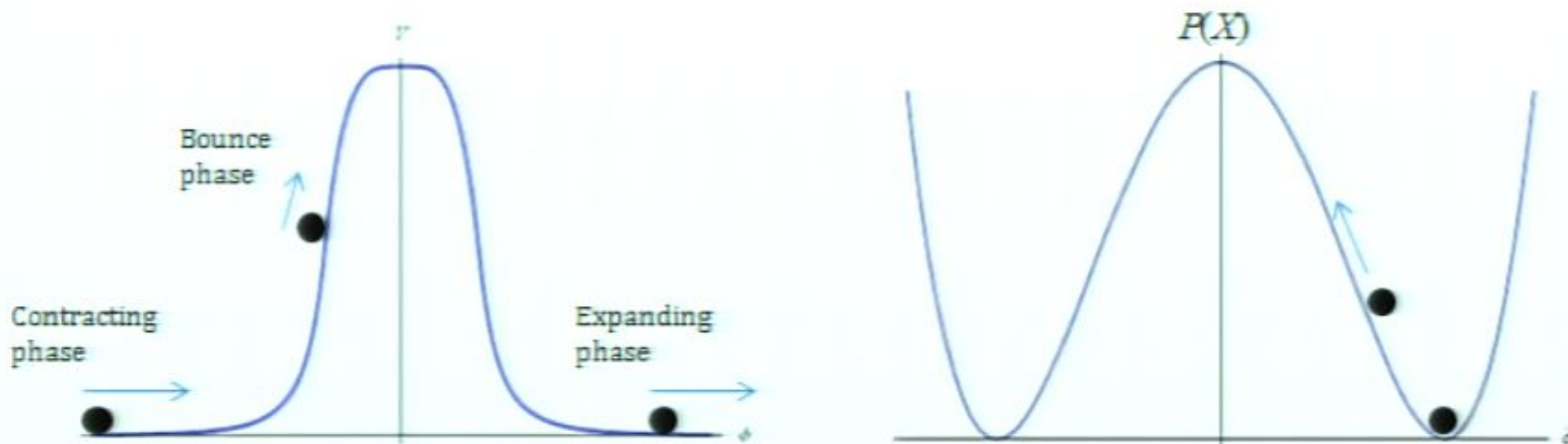
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$\alpha = 4$ Marginally stable against anisotropic stress

$\alpha = 6$ stable



Ghost bounce



$$2M_p^2 \dot{H} = -2M^4 X P' - (1 + w_m) \rho_m$$

Realization of NEC violation!



Ghost bounce

- Ansatz for potential

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Divergence is cut off at M^4

- Ghost field changes as

$$\phi(t) = ct + \pi(t)$$

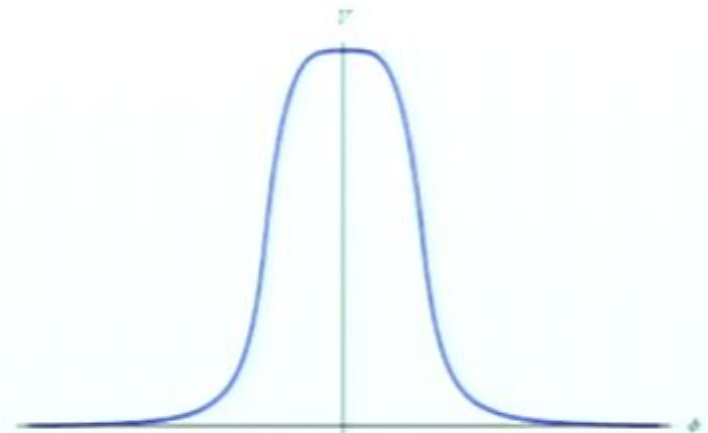
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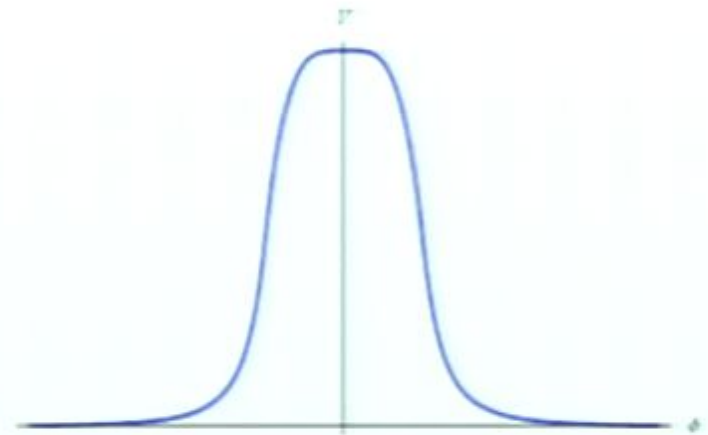
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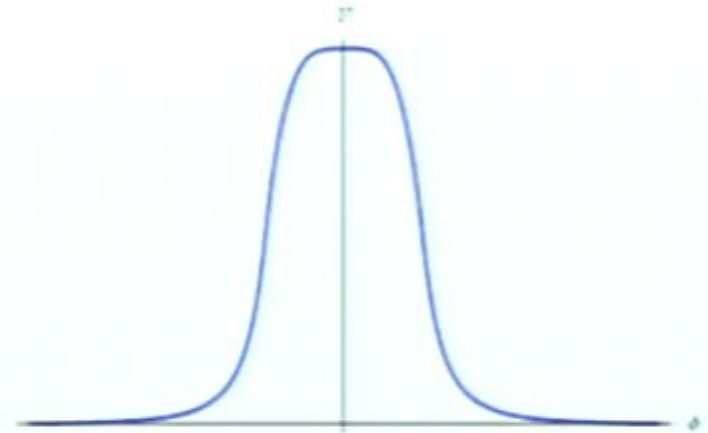
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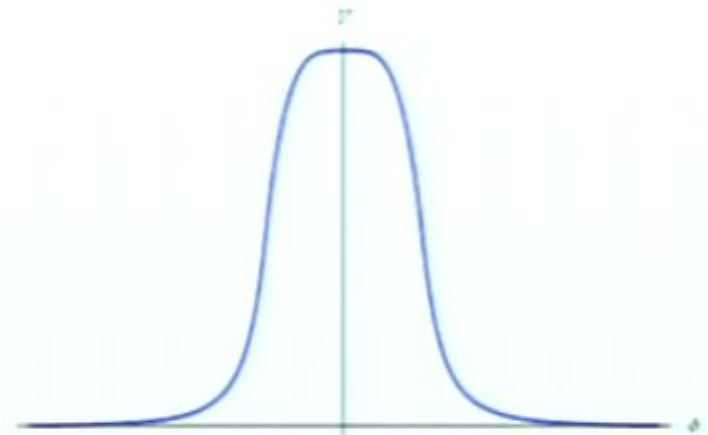
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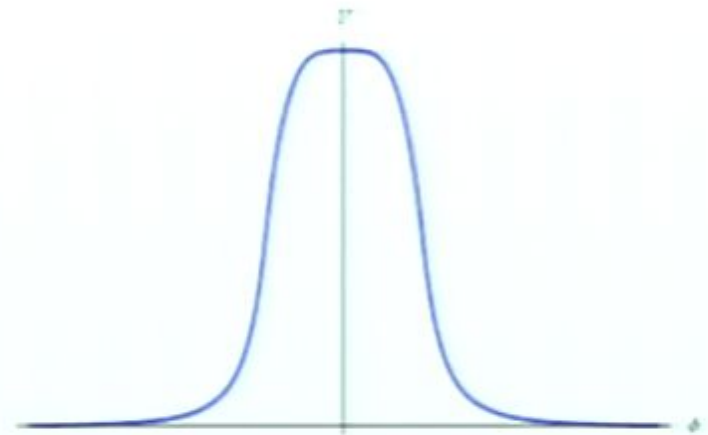
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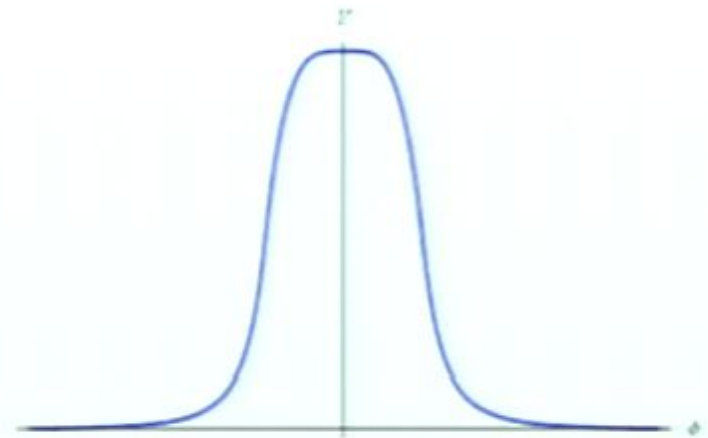
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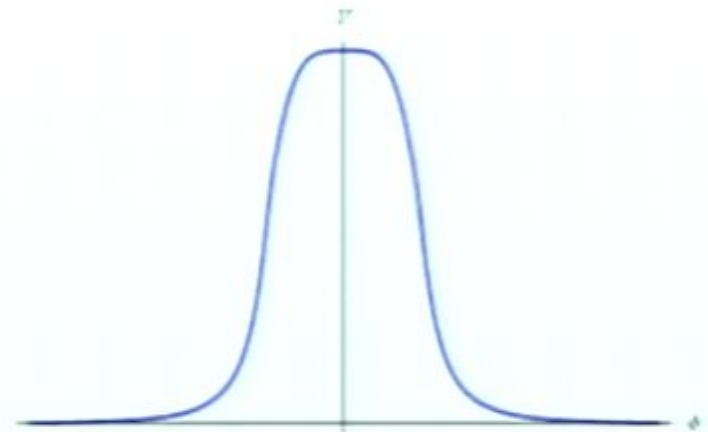
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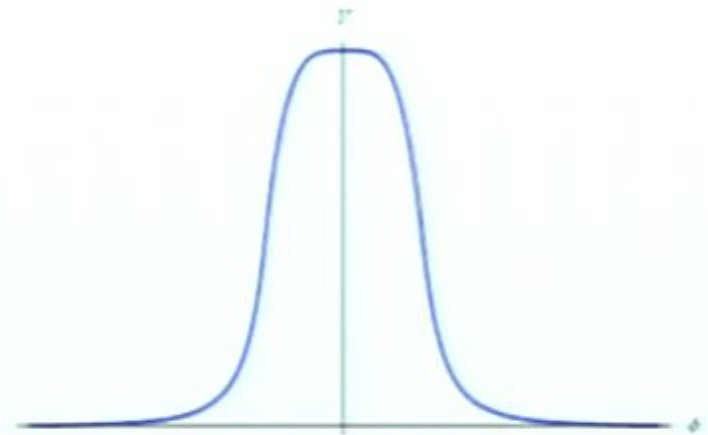
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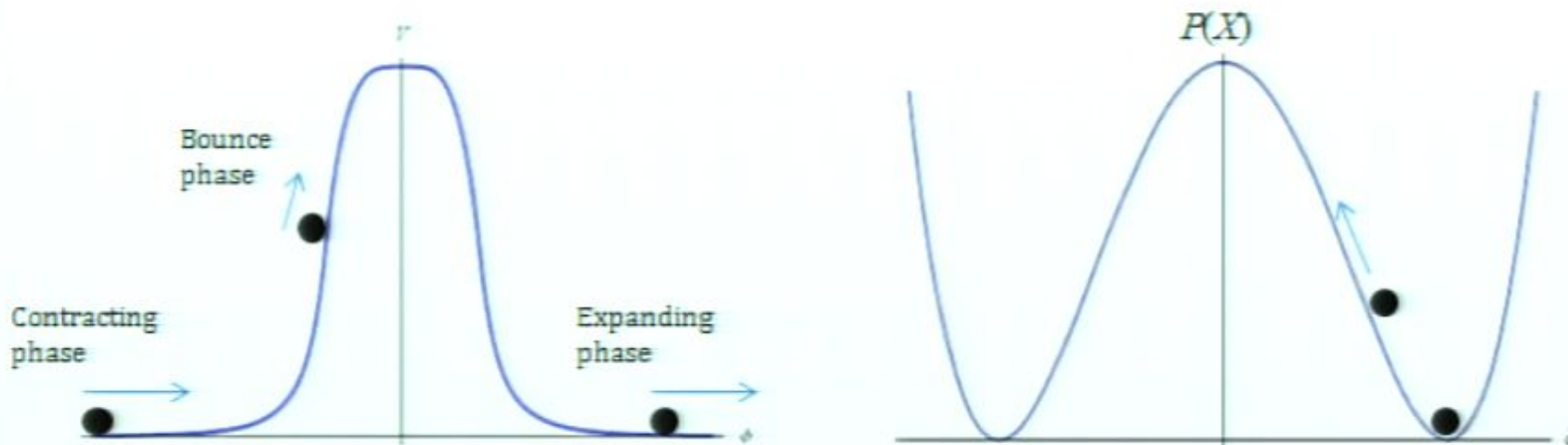
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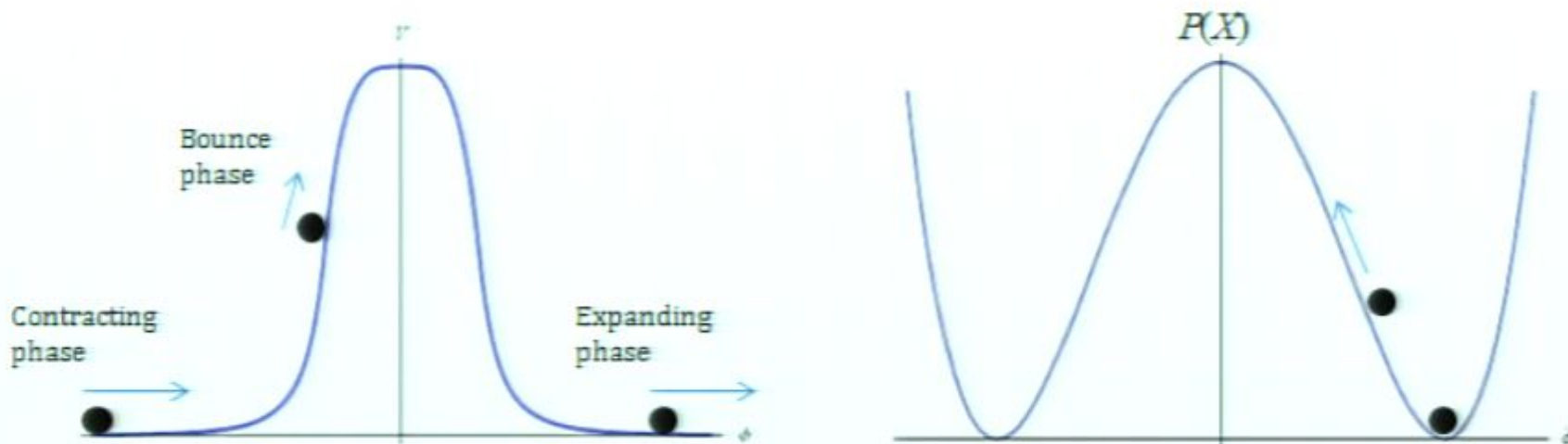


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Realization of NEC violation!



Ghost bounce

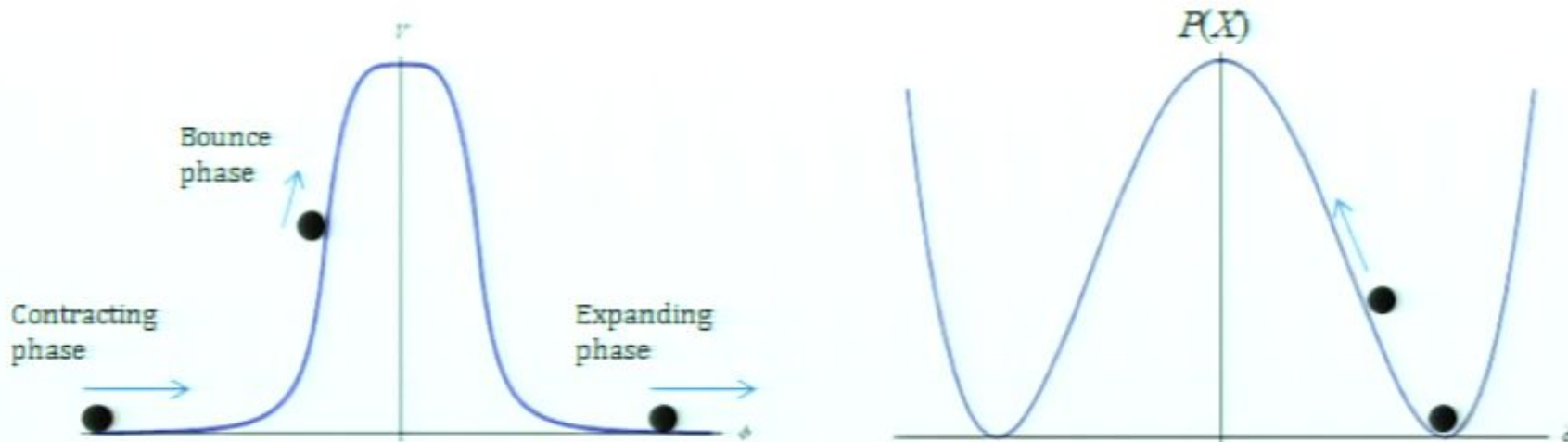


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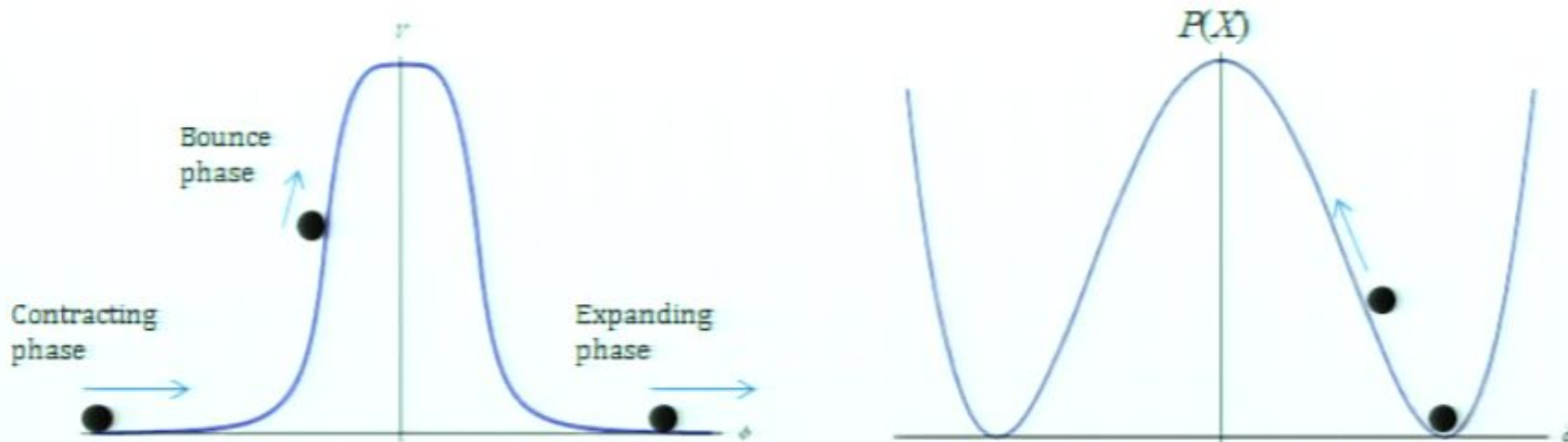


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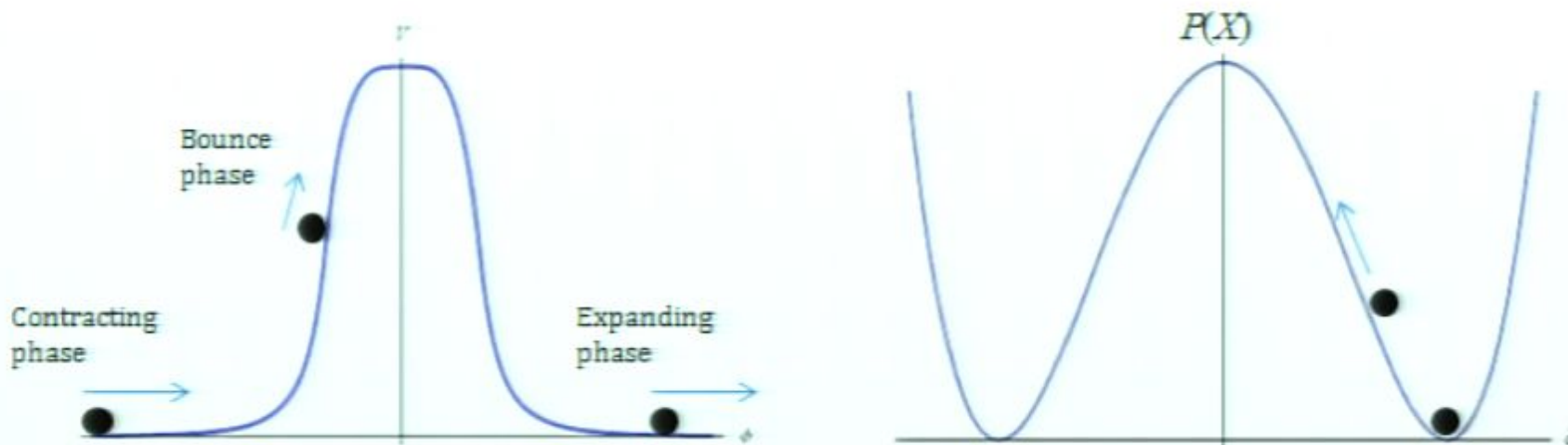


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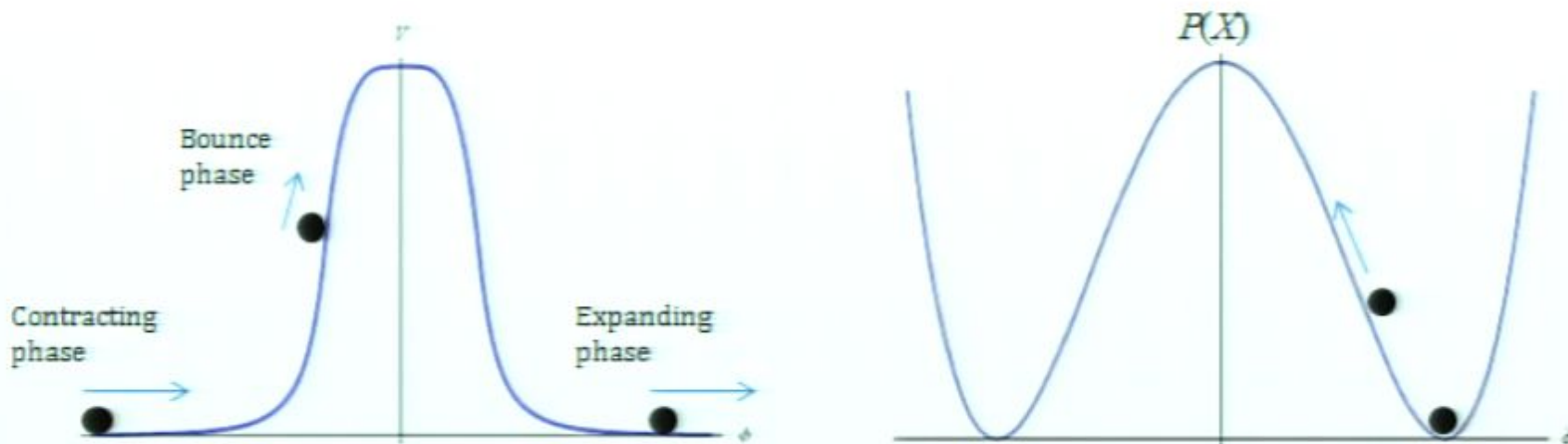


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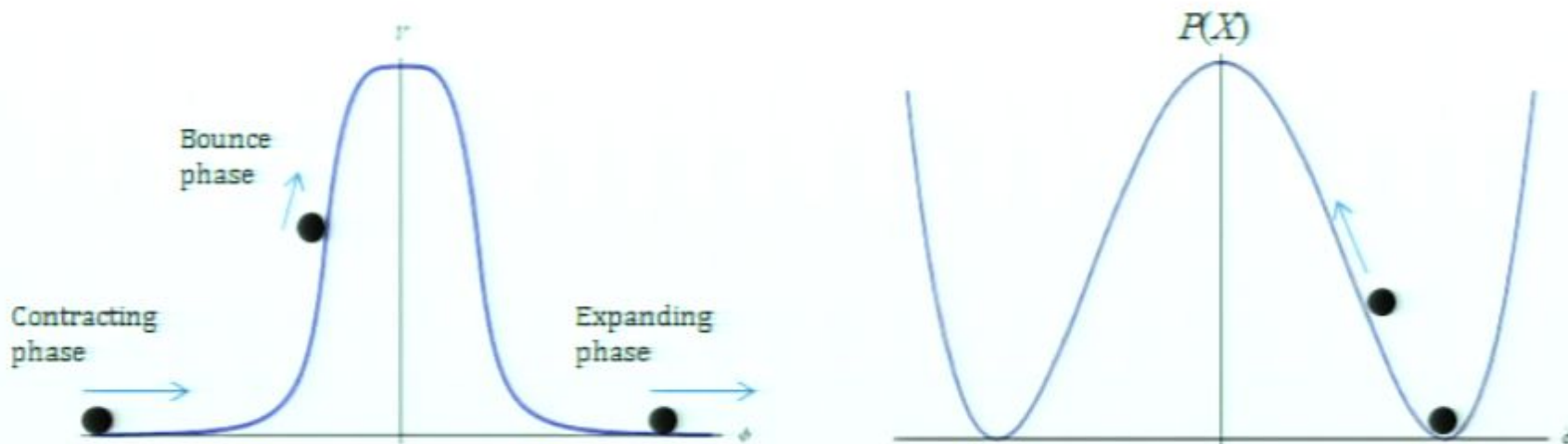


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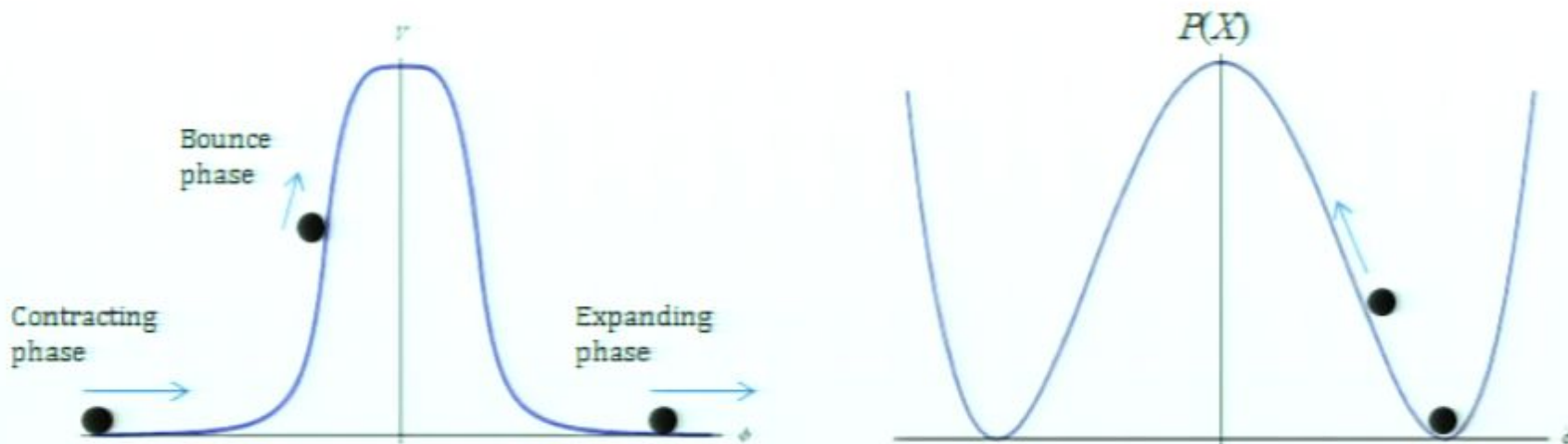


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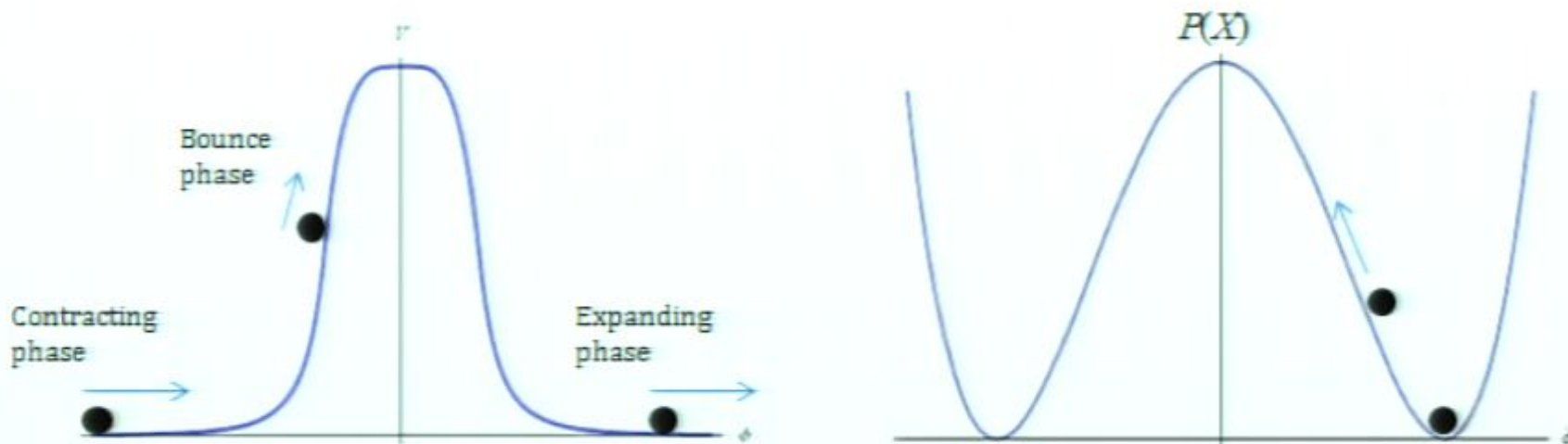


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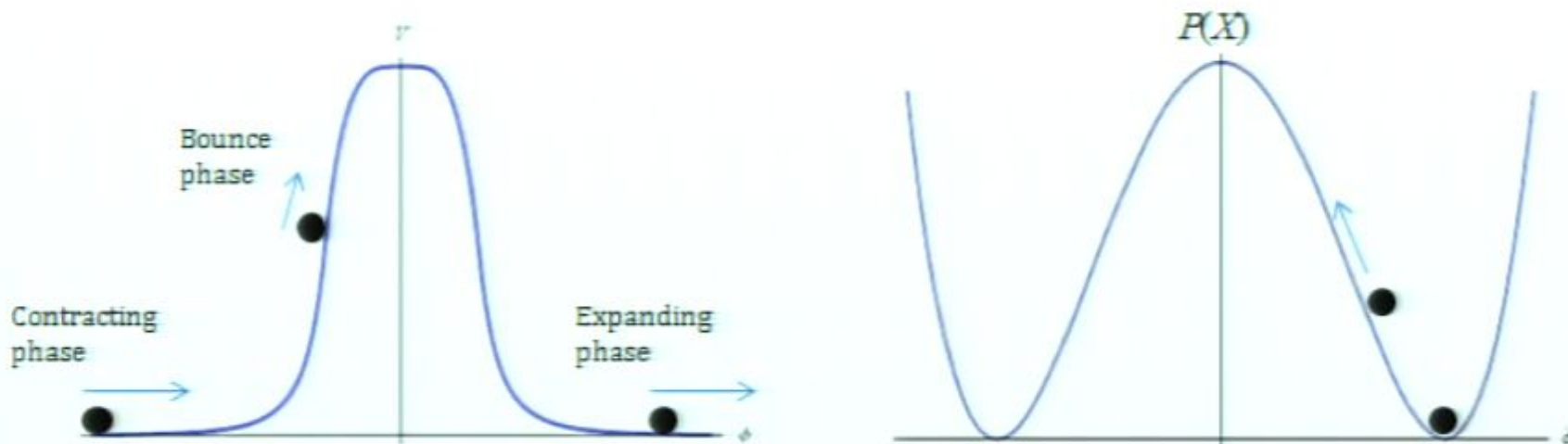


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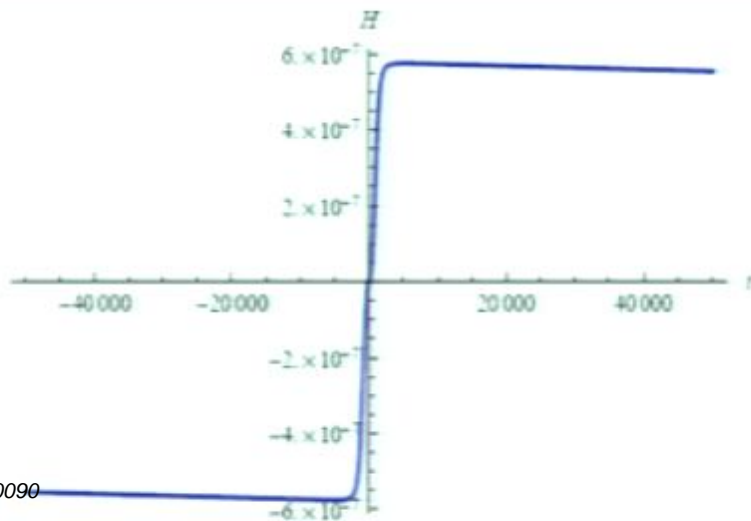
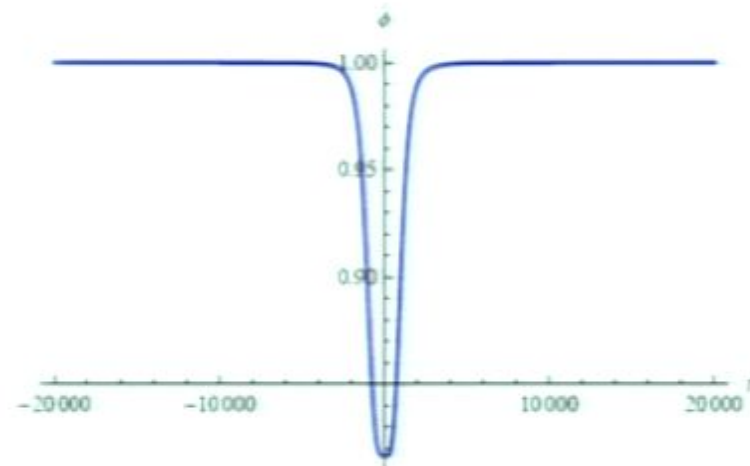
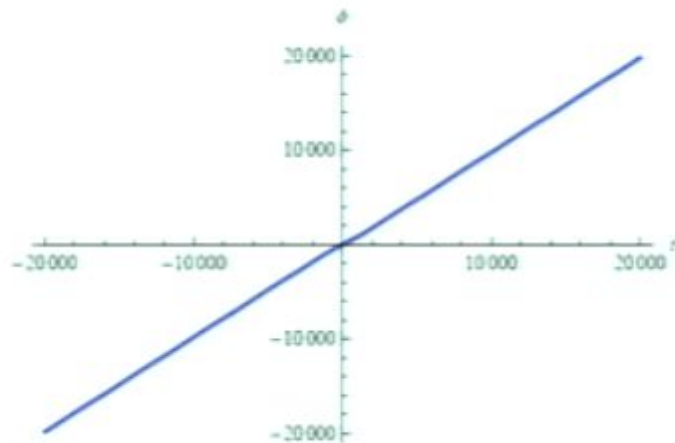


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Realization of NEC violation!



Numerical results



Initial condition of numerical calculation

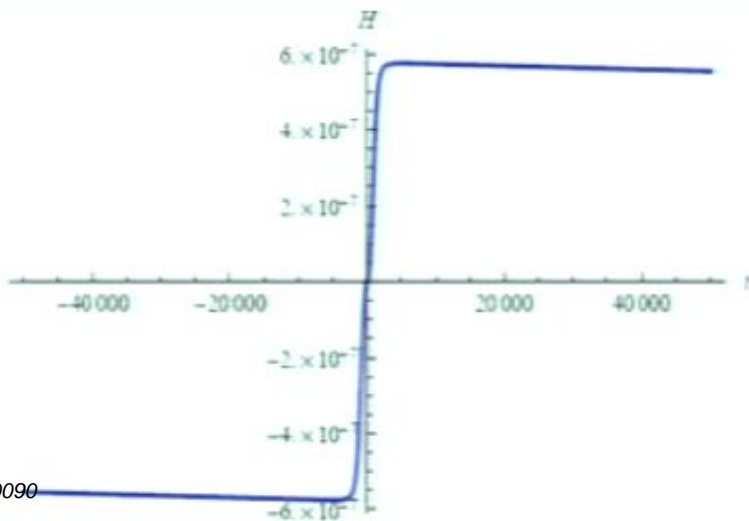
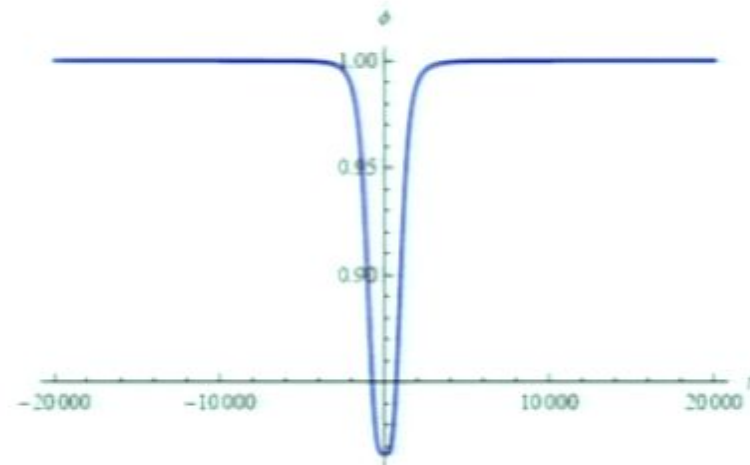
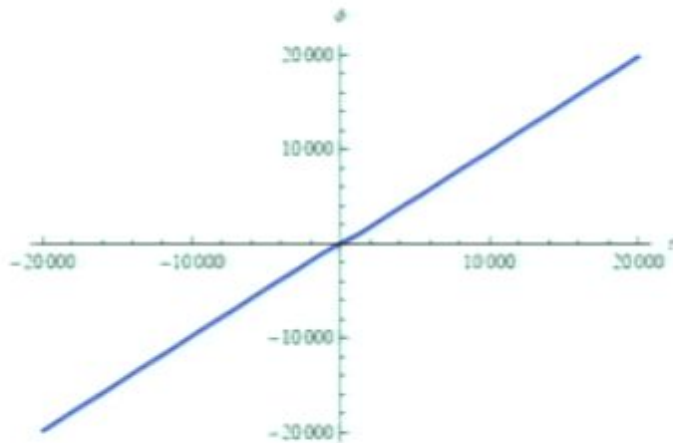
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Numerical results



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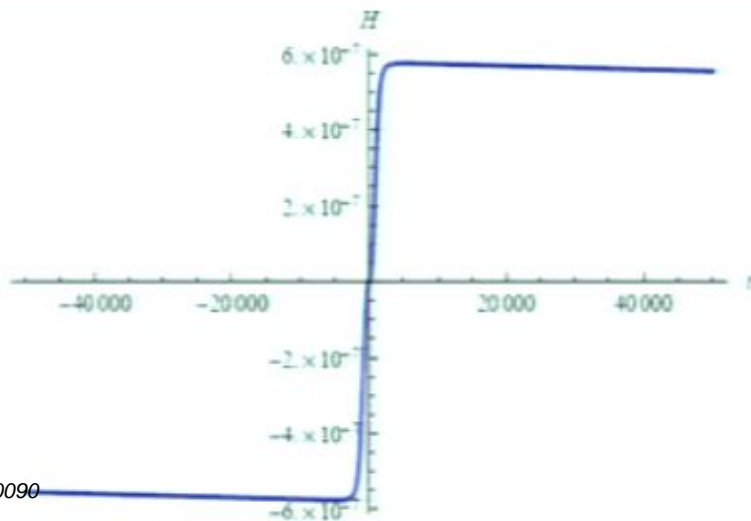
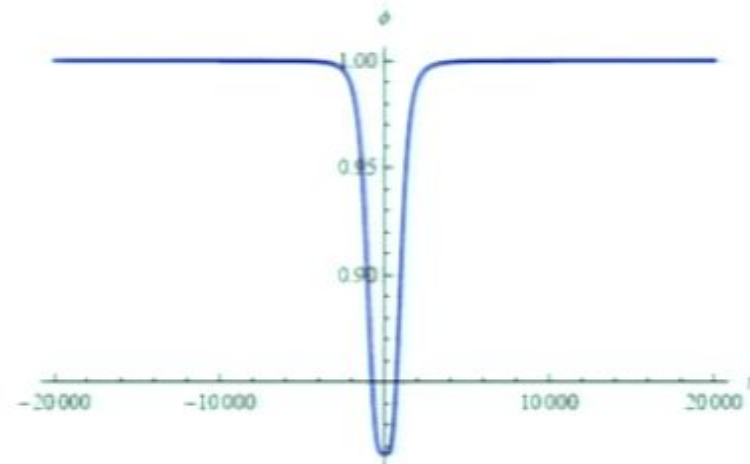
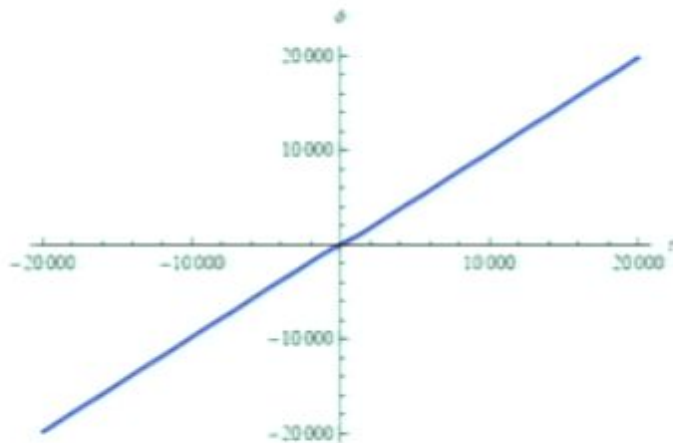
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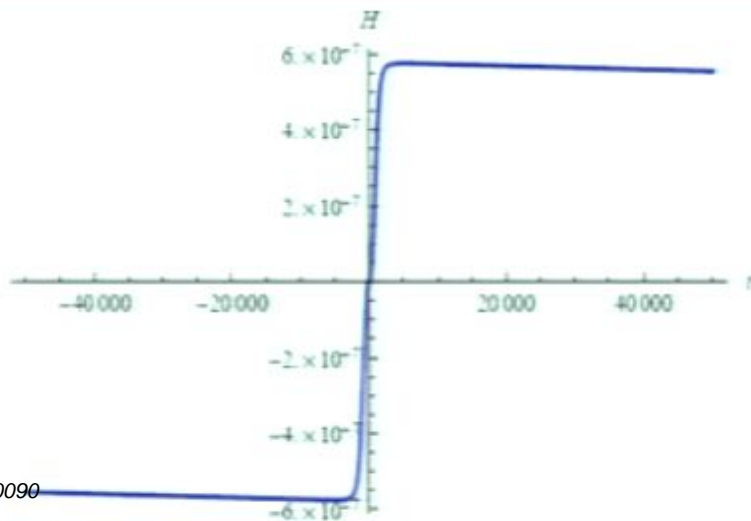
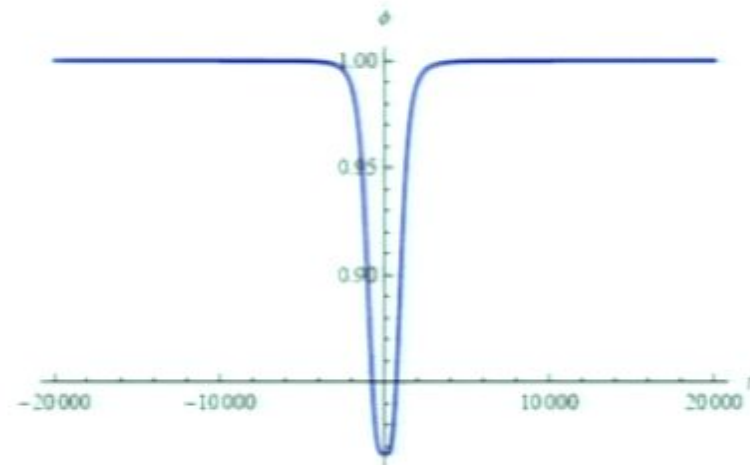
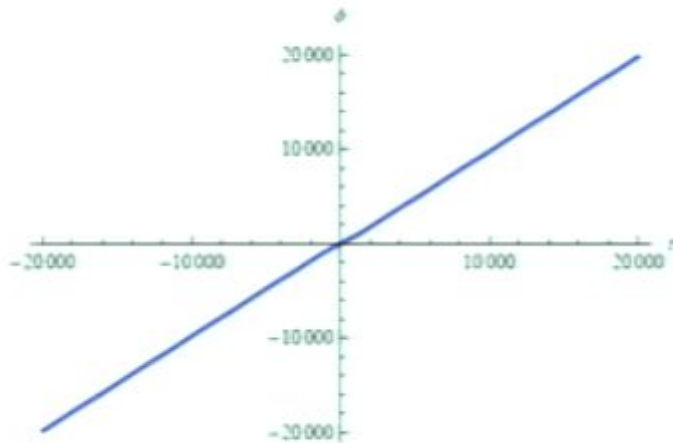
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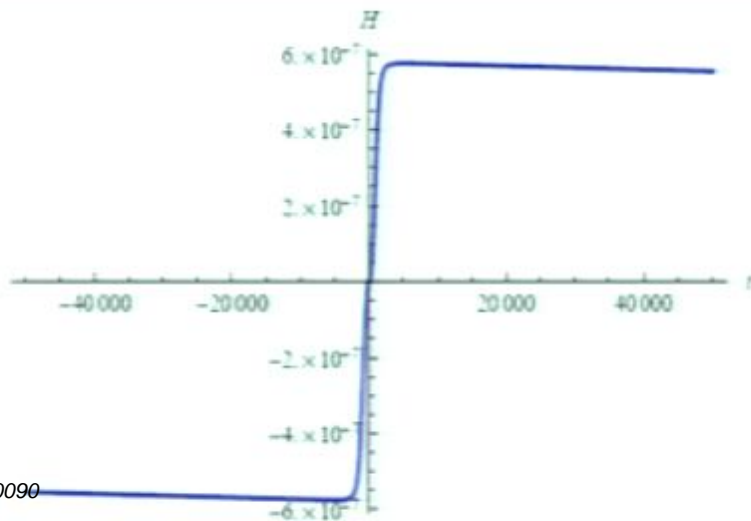
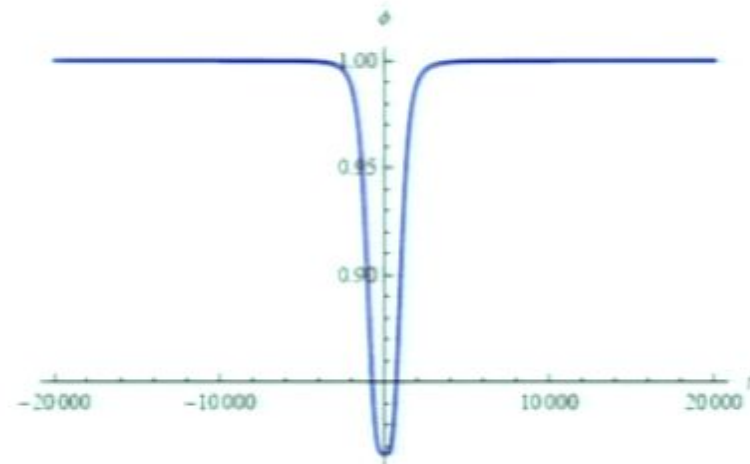
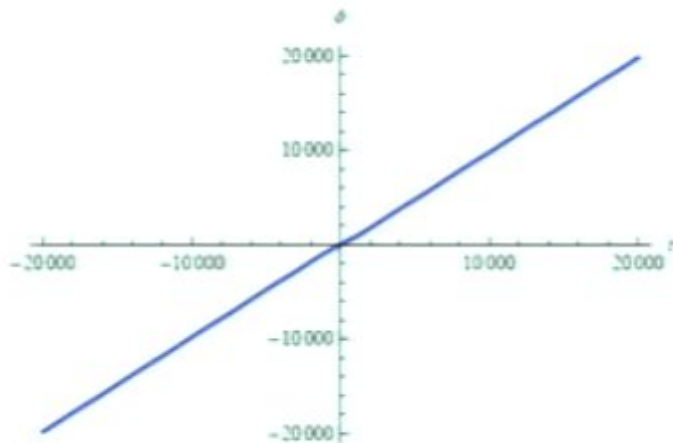
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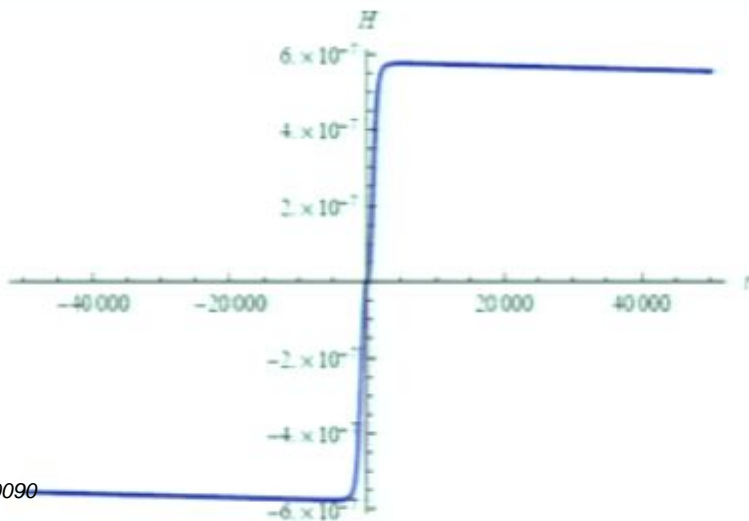
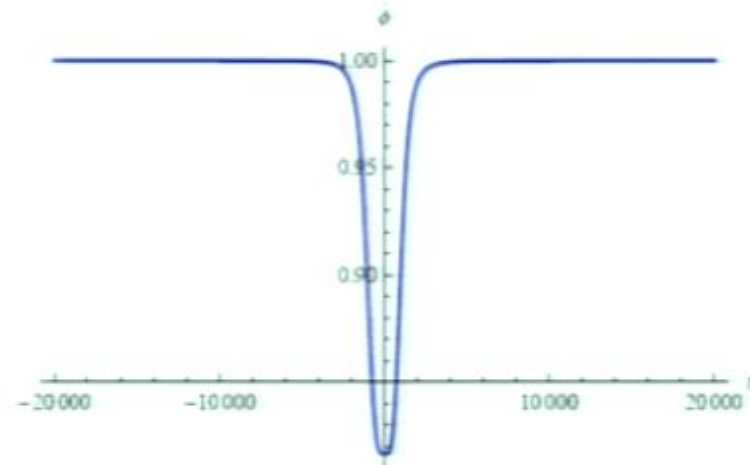
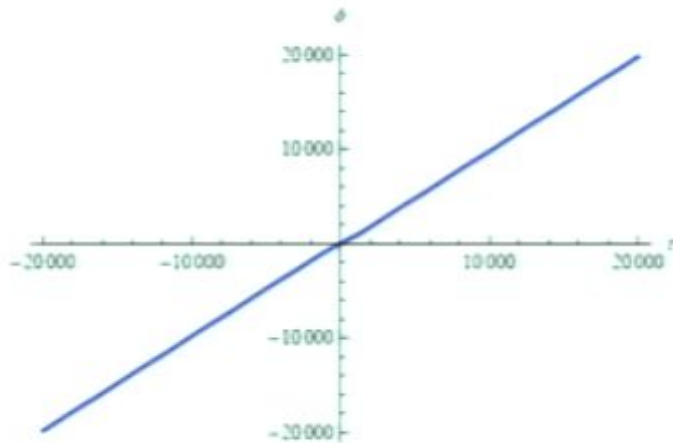
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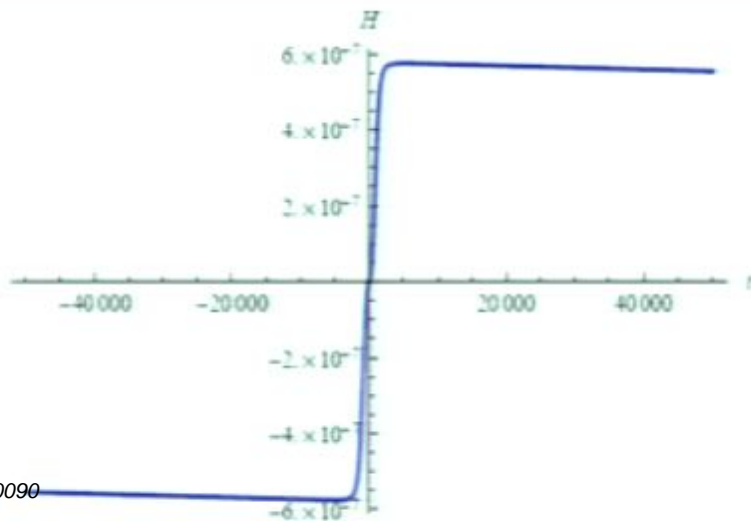
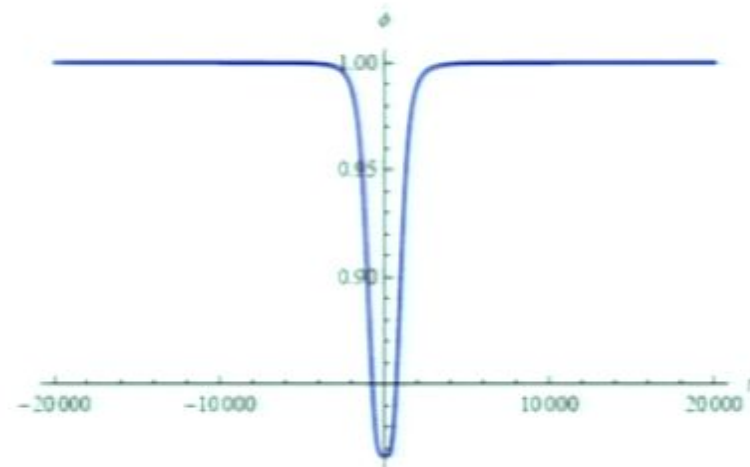
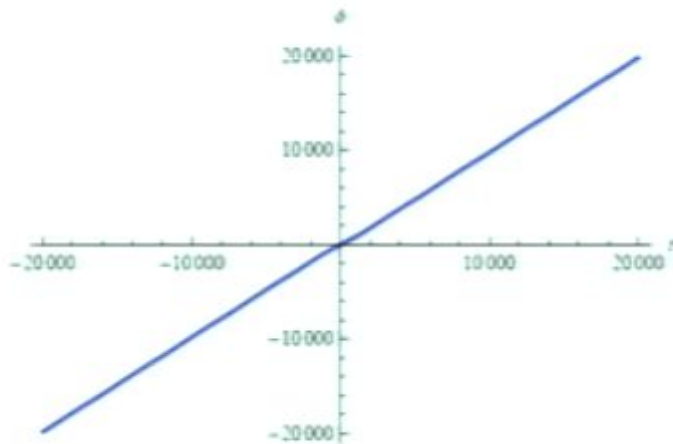
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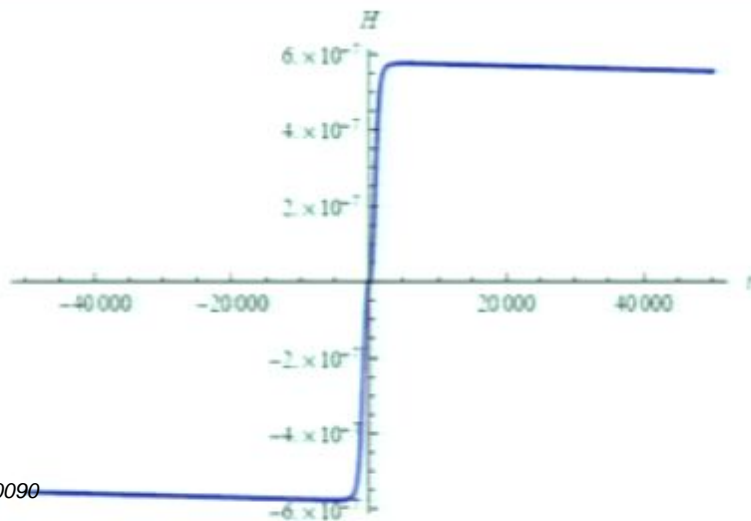
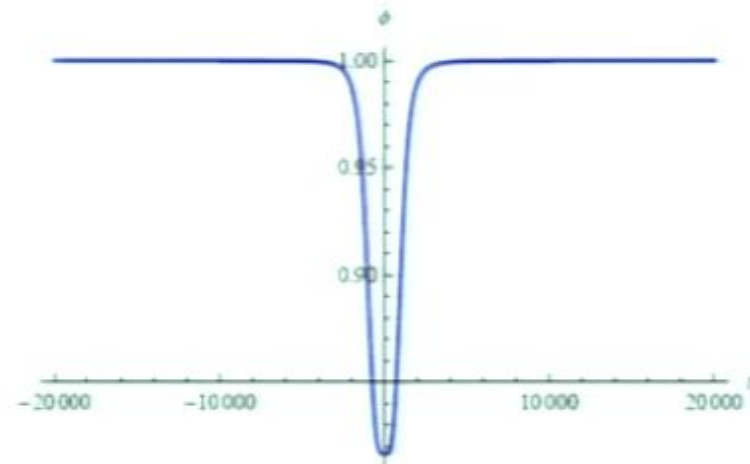
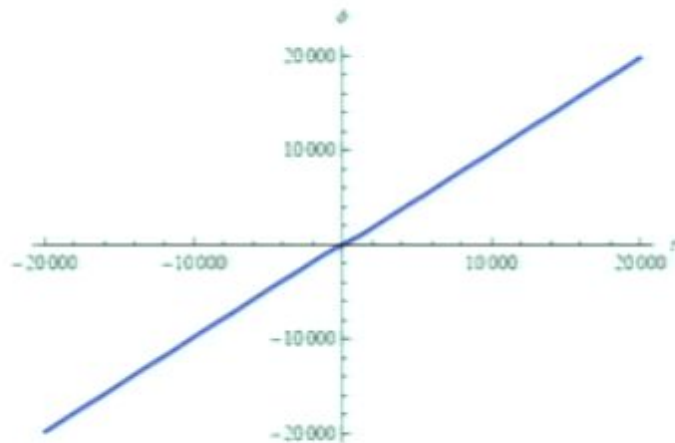
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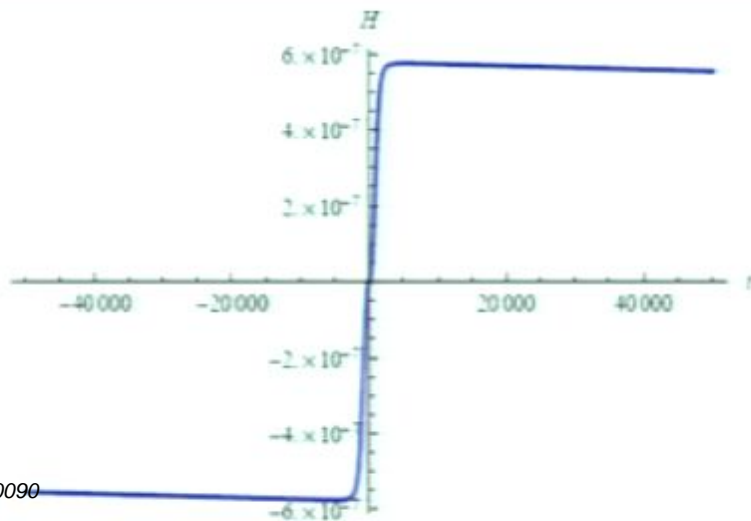
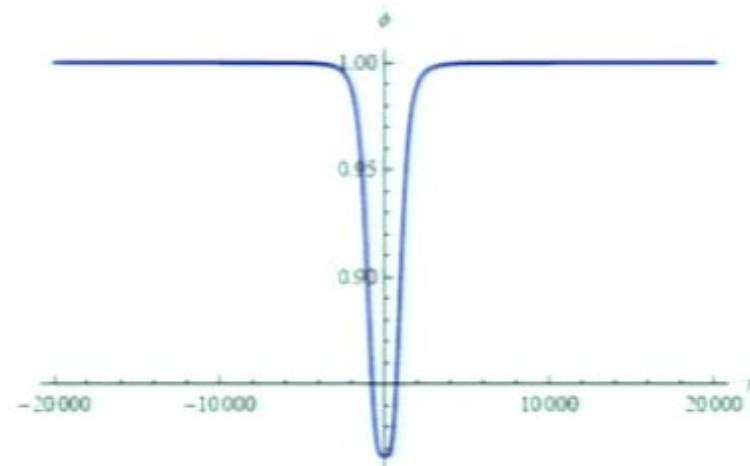
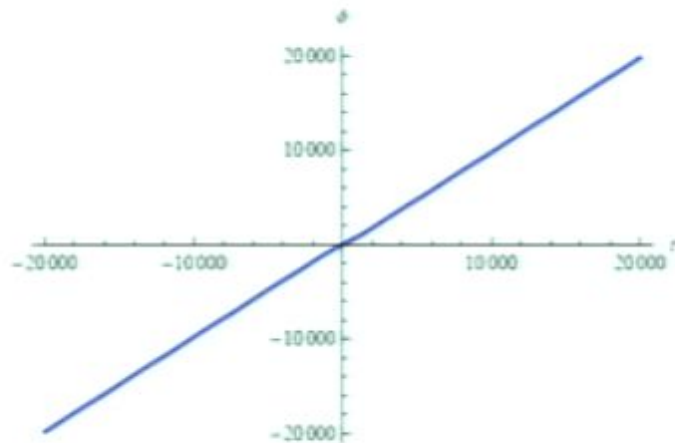
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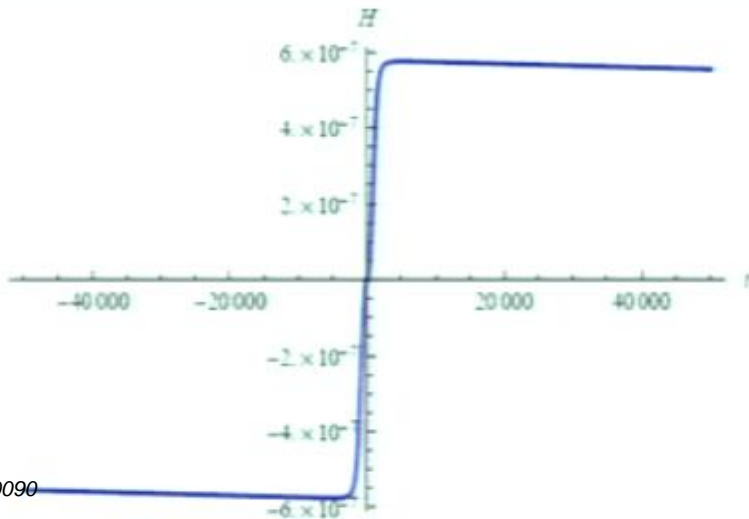
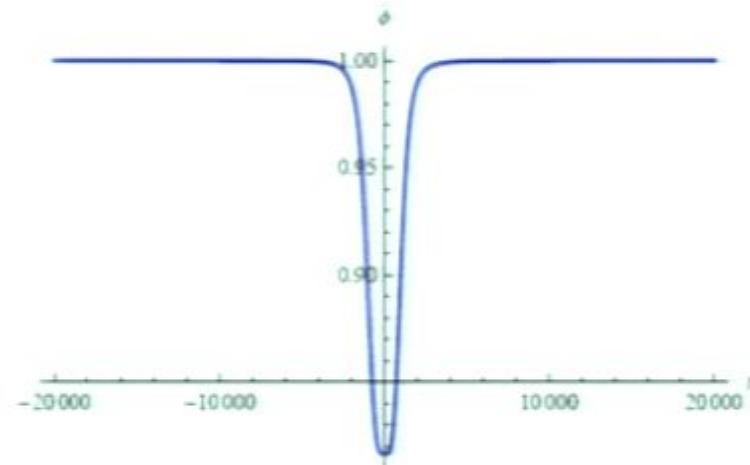
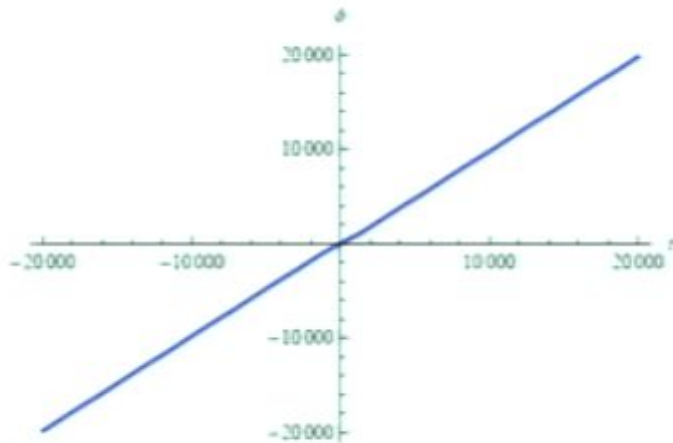
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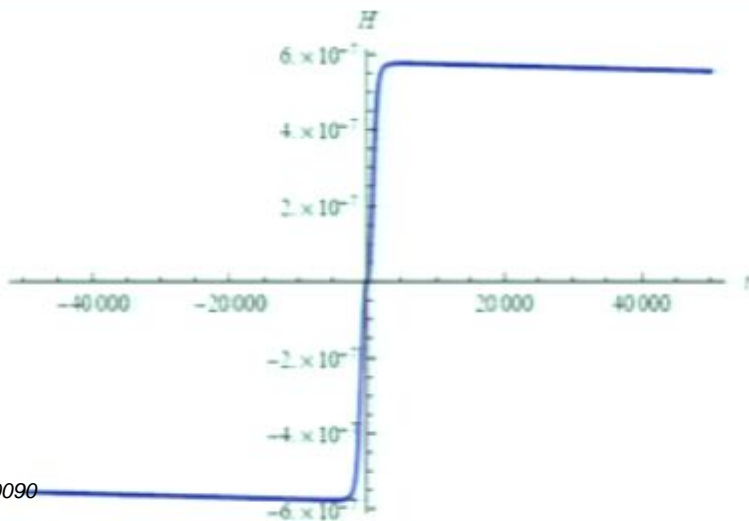
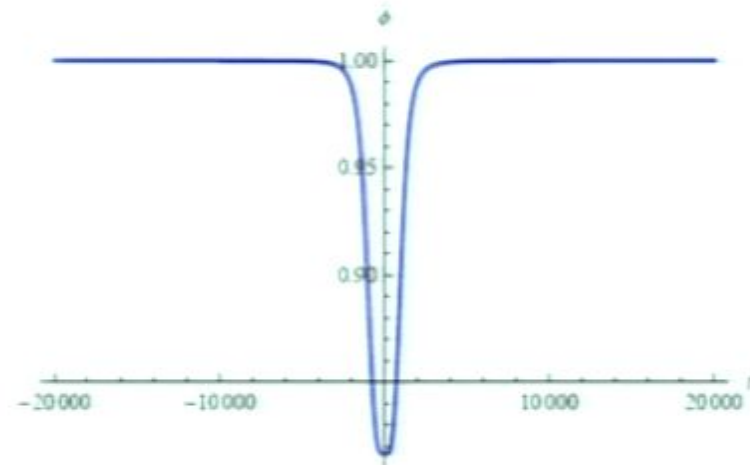
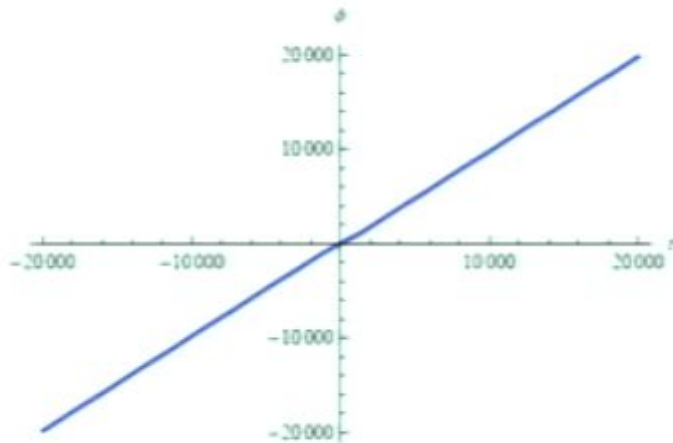
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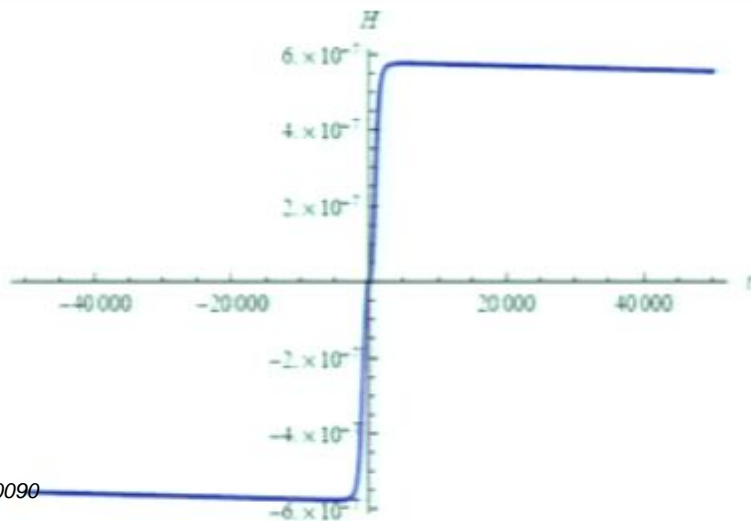
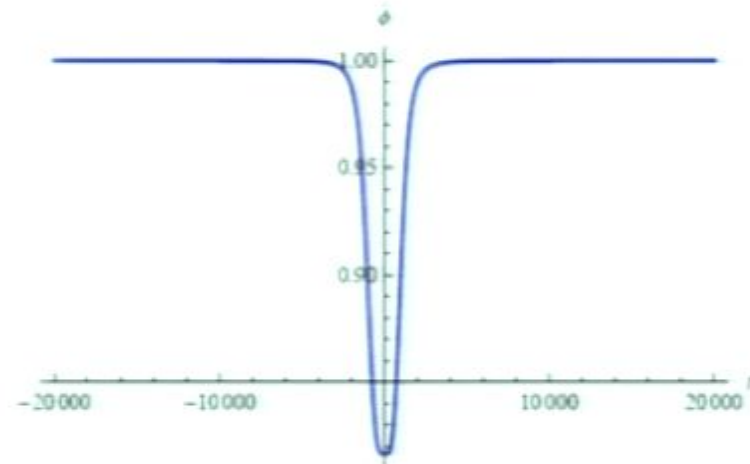
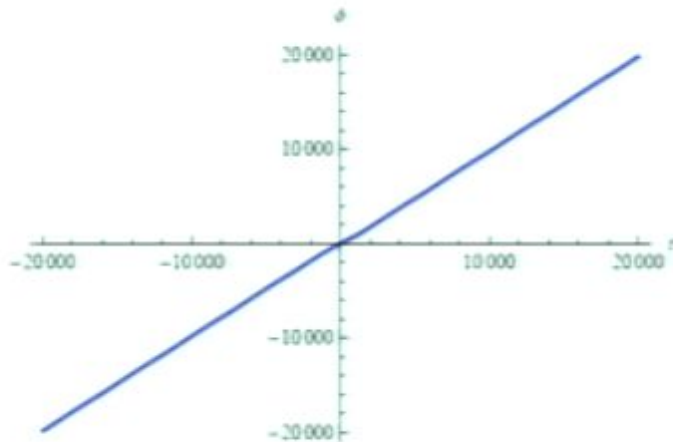
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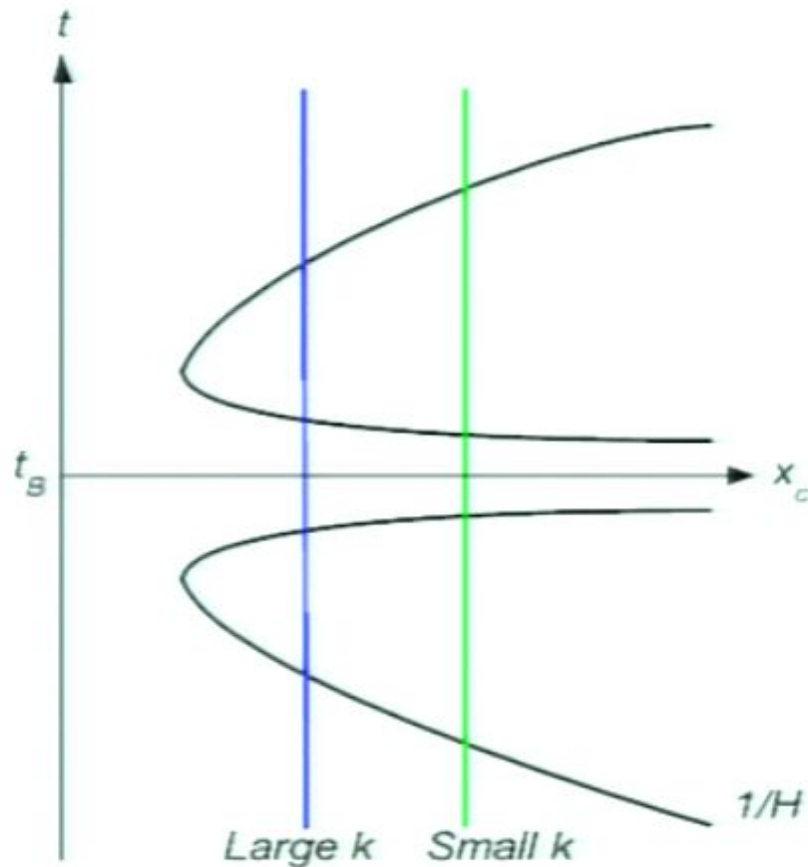
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Matter perturbation



The metric in Newtonian gauge

$$ds^2 = -(1 + 2\Phi)dt^2 + a(t)^2(1 - 2\Psi)dx^2$$

$$\varphi(\eta, \mathbf{x}) = \varphi_0(\eta) + \delta\varphi(\eta, \mathbf{x})$$

Introduce M-S Variable

$$v = a\left[\delta\varphi + \frac{\dot{z}}{a}\Phi\right]$$

$$S^{(2)} = \frac{1}{2} \int d^4x \left[v'^2 - v_{,i}v_{,i} + \frac{z''}{z}v^2 \right]$$

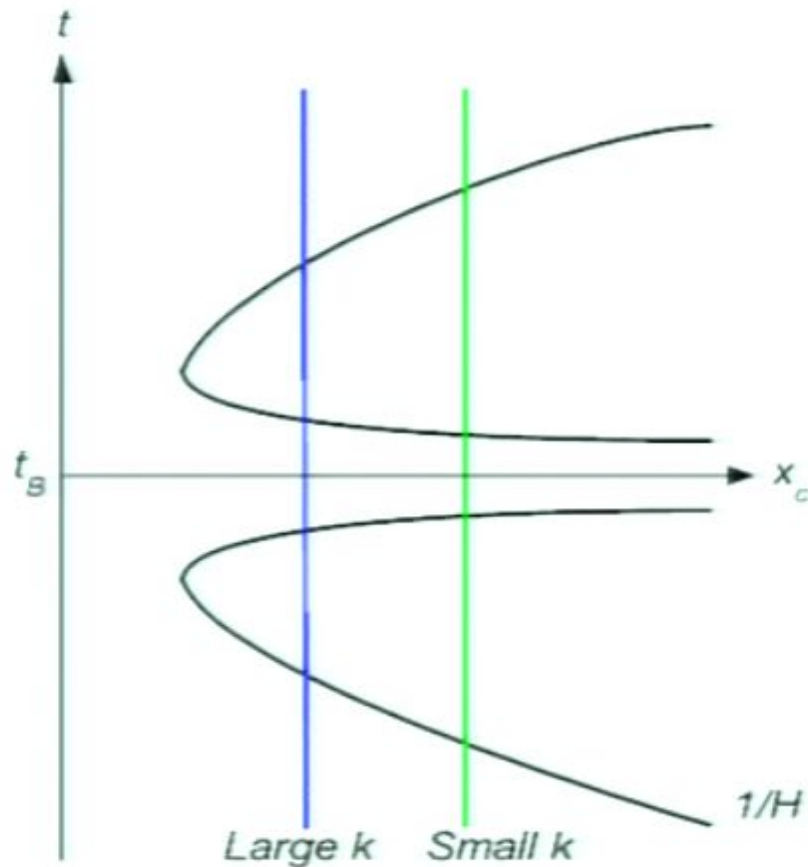
EoM:

$$v_k'' + k^2 v_k - \frac{z''}{z} v_k = 0$$

In matter contracting phase $z \sim a$

$$v(t) \sim t^{-1/3}$$

Matter perturbation



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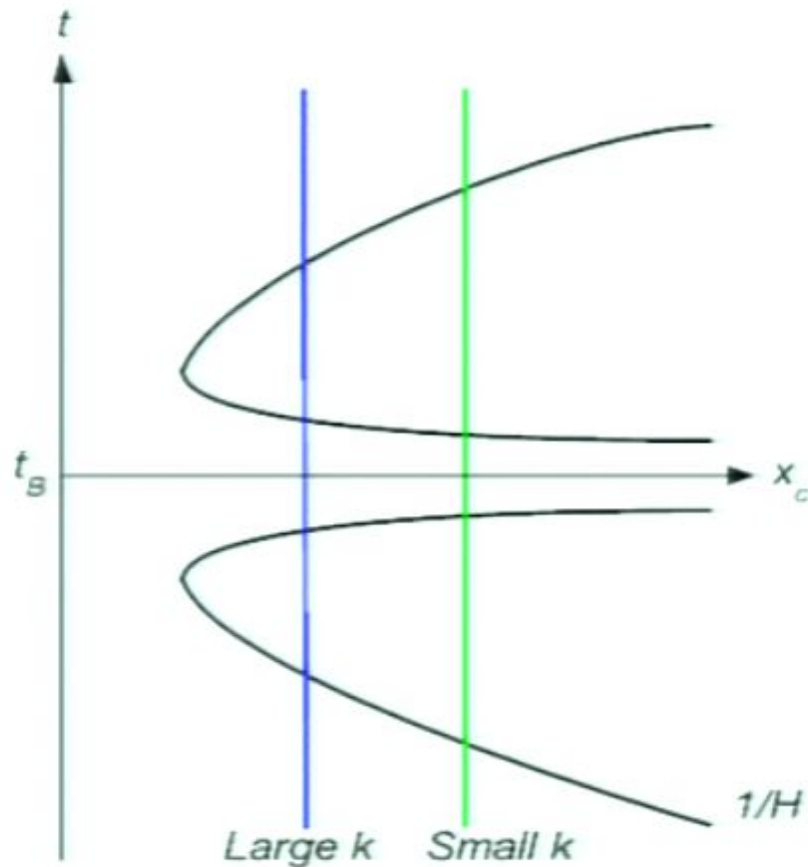
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$$v_k'' + k^2 v_k - \frac{z''}{z} v_k = 0$$

In matter contracting phase $z \sim a$

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Matter perturbation



The metric in Newtonian gauge

$$ds^2 = -(1 + 2\Phi)dt^2 + a(t)^2(1 - 2\Psi)dx^2$$

$$\varphi(\eta, \mathbf{x}) = \varphi_0(\eta) + \delta\varphi(\eta, \mathbf{x})$$

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$$v = a\left[\delta\varphi + \frac{\dot{z}}{a}\Phi\right]$$

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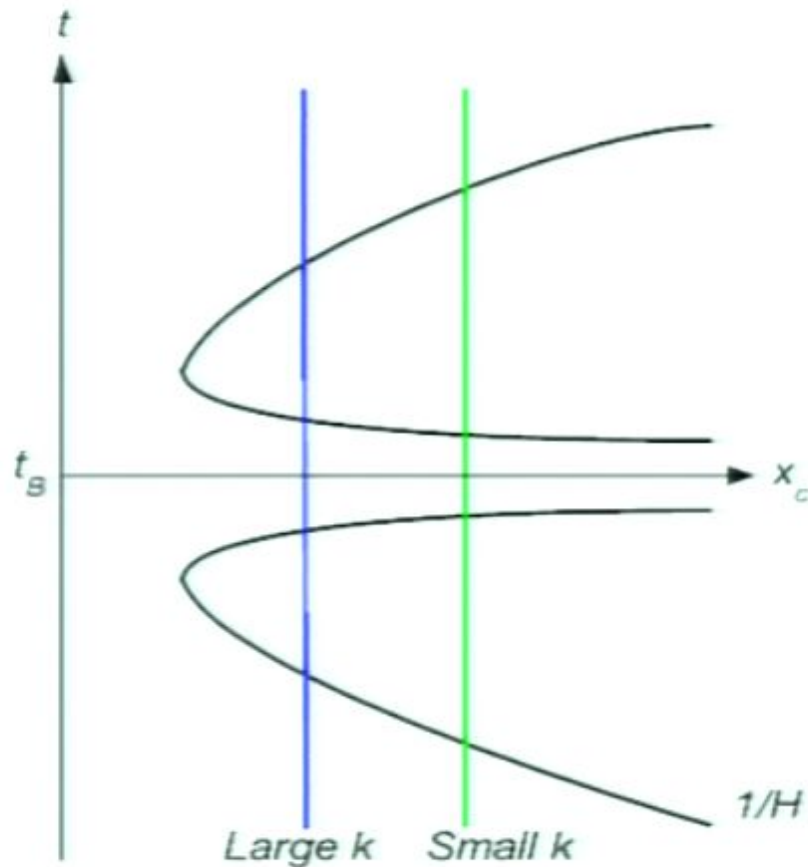
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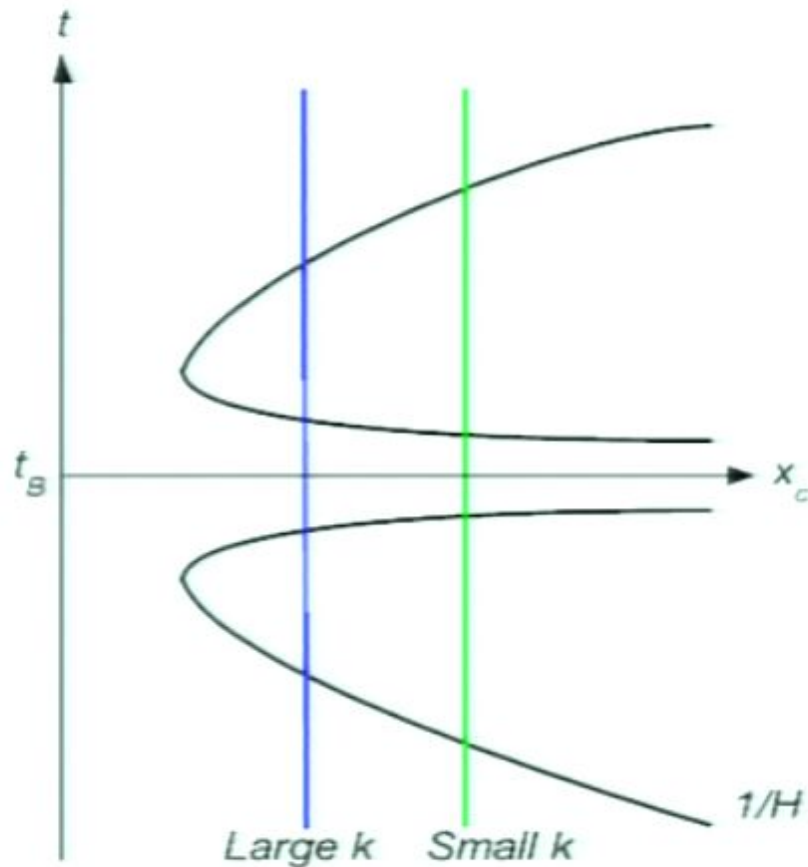
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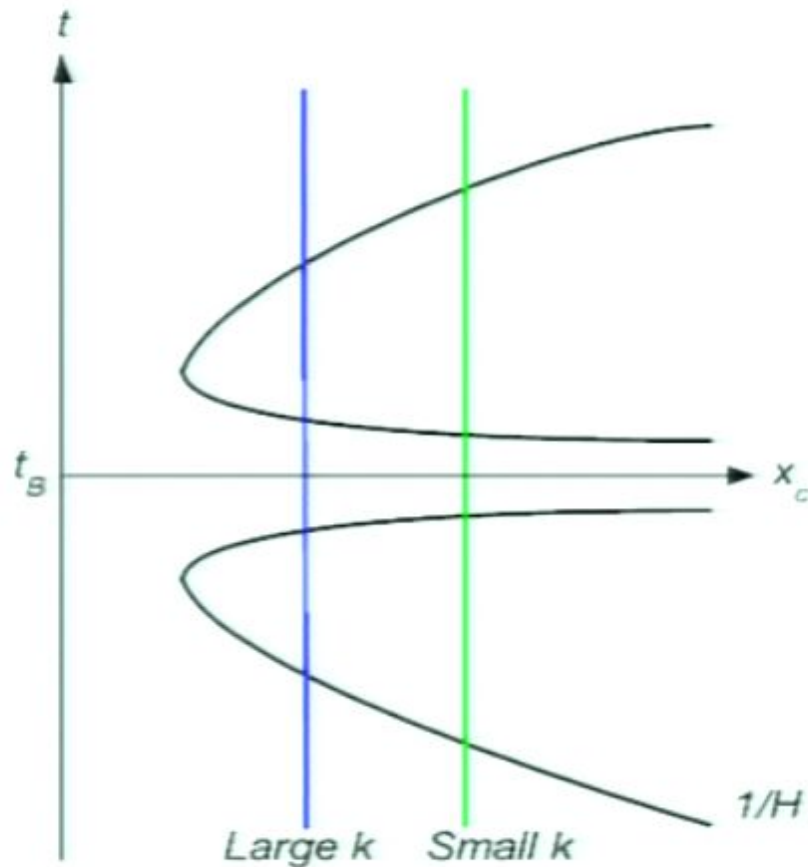
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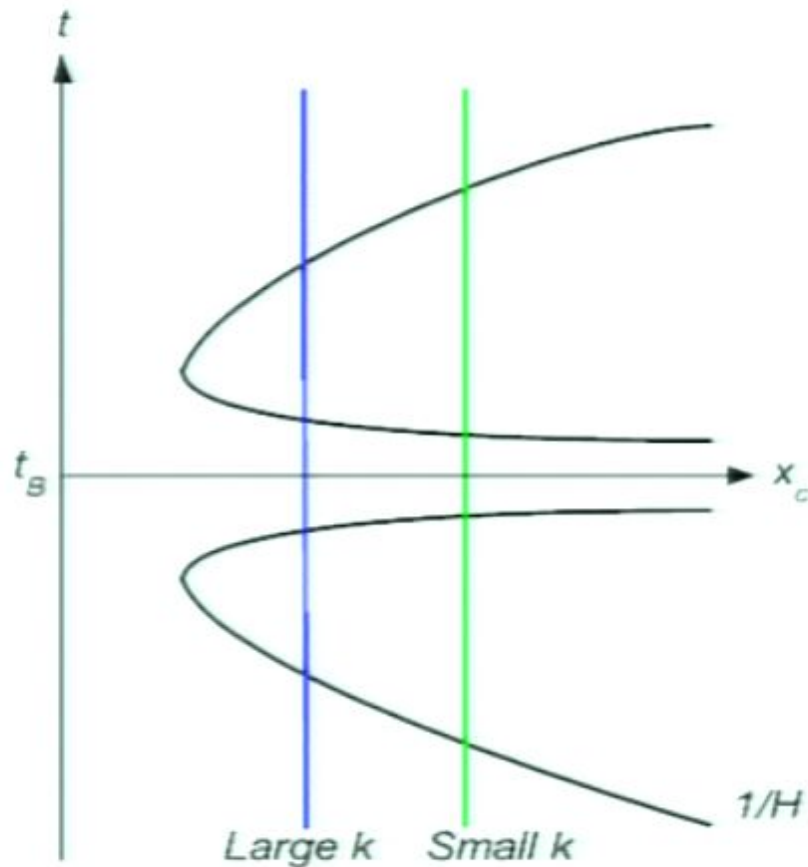
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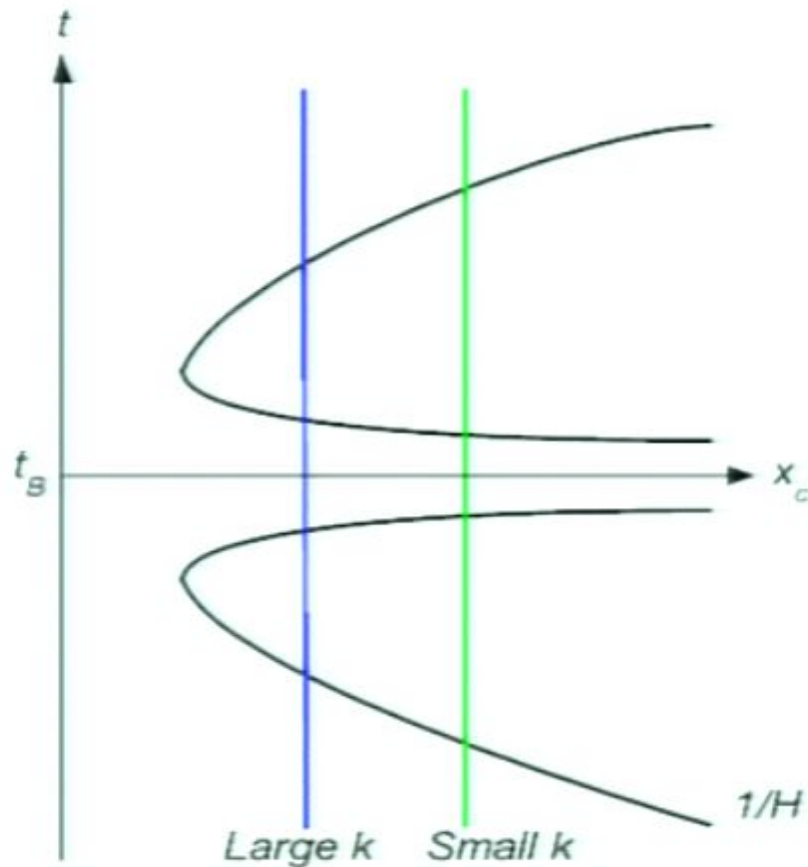
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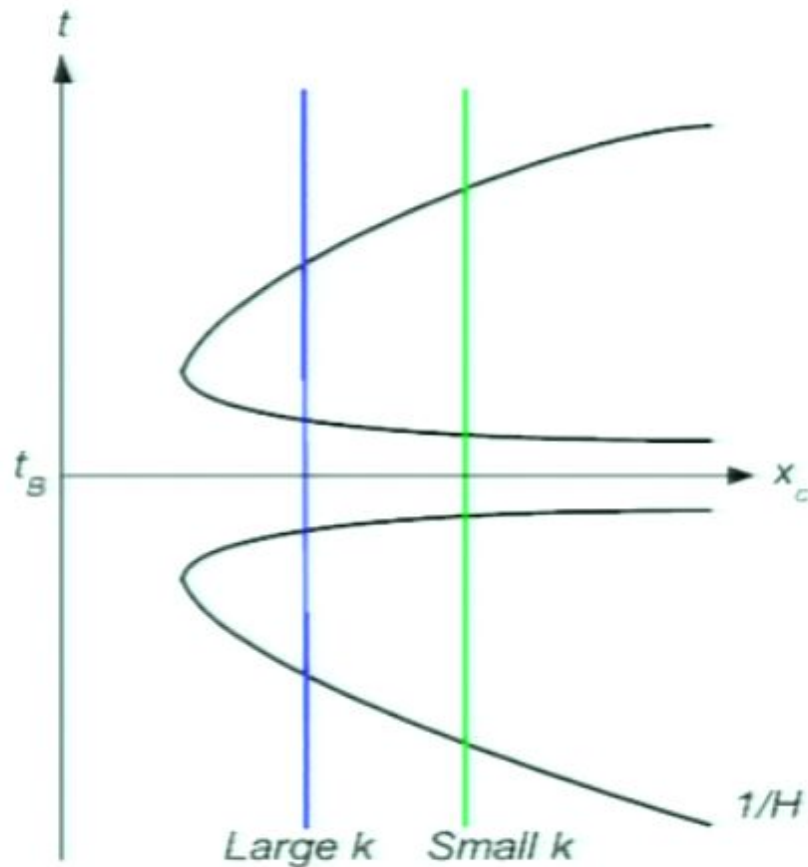
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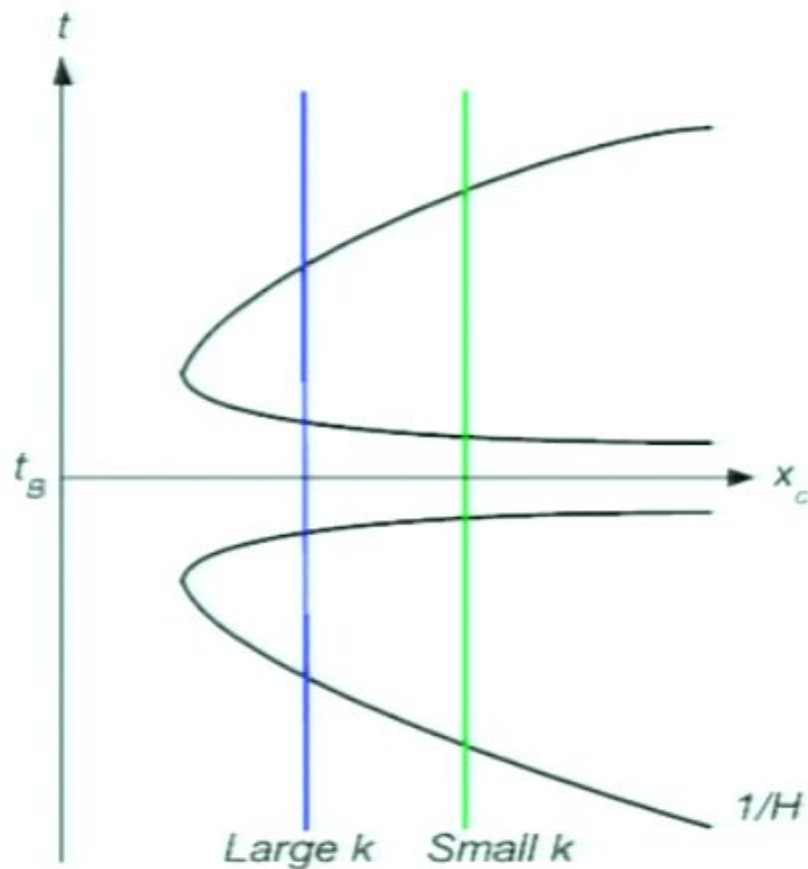
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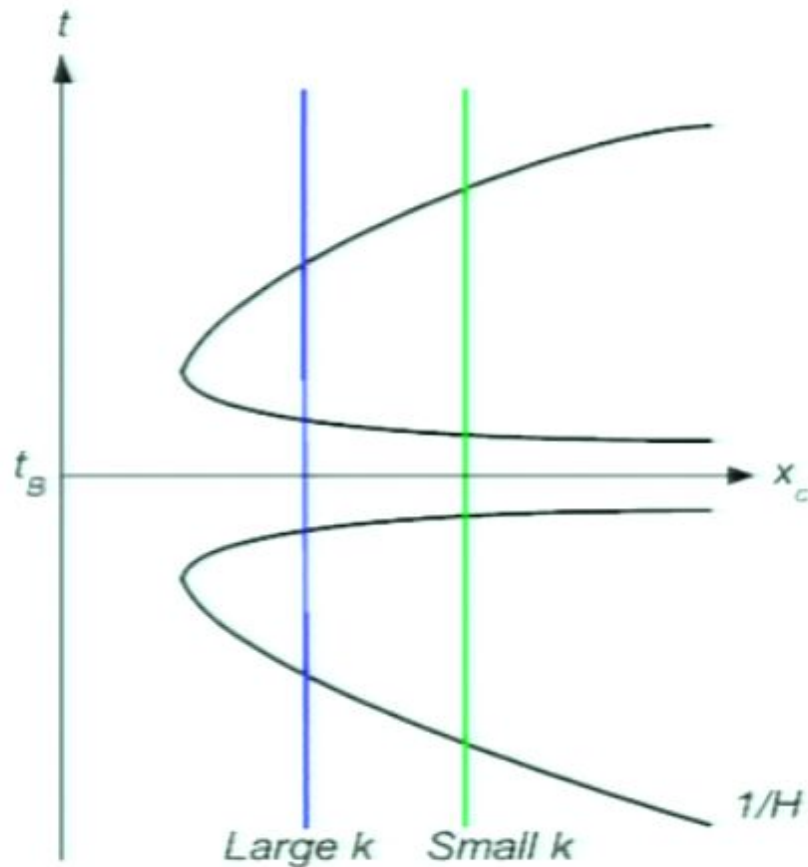
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 - a) In matter contracting phase, Ghost perturbation does NOT grow faster than matter perturbation;
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- ◆ Since the duration of bounce phase is short

$$H = \theta \cdot (t - t_B)$$

Where $\theta \gg H_c^2$ we interested in large scale perturbation

$$\partial_t^2 \Phi_g + \theta \Phi_g = 0$$

the solution is

$$\Phi_g = d_1 e^{i\sqrt{\theta}t} + d_2 e^{-i\sqrt{\theta}t}$$

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- ◆ Since the duration of bounce phase is short

$$H = \theta \cdot (t - t_B)$$

Where $\theta \gg H_c^2$ we interested in large scale perturbation

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So we need a very small M to protect the IR gravity.

A constraint condition has been given in hep-ph/0507120, where

$$M < 100 \text{Gev}$$

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$$\alpha = 4, \quad \rho_g \sim a^{-6}$$



Radiation epoch

δ_g



Logarithmic growing

Matter epoch

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δ_g



δ_m

Cut-off issue

Our ghost bounce model is free from cut-off upper bound

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Radiation epoch

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 - ◆ Interesting feature
 - ◆ Jeans instability \rightarrow low energy scale 100Gev
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Advantages:

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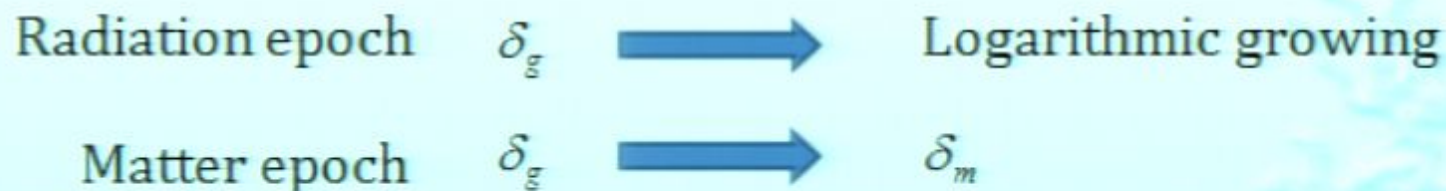


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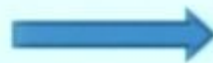
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- ◆ Ghost condensation theory
 - ◆ Stabilize vacuum
 - ◆ Interesting feature
 - ◆ Jeans instability \rightarrow low energy scale 100Gev
- ◆ We realize matter bounce by means of ghost condensation

Advantages:

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2. Background is stable against radiation and anisotropic stress;
3. We have a high energy scale bounce $\gg 100\text{Gev}$

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Stability during bounce

the dispersion relation

$$\omega^2 = \frac{-(\tilde{M}^2 M^4 + 4M_{pl}^4 \dot{H})k^2 + 2M_{pl}^2 \tilde{M}^2 k^4}{2M_{pl}^2 M^4}$$

the typical instability rate

$$\omega_c = \frac{1}{4} \frac{\tilde{M} M^2}{M_{pl}^2} + \dot{H} \frac{M_{pl}^2}{M^2 \tilde{M}}$$

Its growing rate during bounce phase

$$\Delta t \omega_c \sim \left(\frac{V_0}{M^4}\right)^{1/\alpha} \left[\frac{1}{4} \frac{\tilde{M} M}{M_{pl}^2} + \dot{H} \frac{M_{pl}^2}{M^3 \tilde{M}} \right]$$

Since $\dot{H} \sim \frac{M^4 \dot{\pi}}{M_{pl}^2}$ if $V_0 \ll M^4$ We get

$$\Delta t \omega_c \ll 1$$

Ghost perturbation

- ◆ Since the duration of bounce phase is short

$$H = \theta \cdot (t - t_B)$$

Where $\theta \gg H_c^2$ we interested in large scale perturbation

$$\partial_t^2 \Phi_g + \theta \Phi_g = 0$$

the solution is

$$\Phi_g = d_1 e^{i\sqrt{\theta}t} + d_2 e^{-i\sqrt{\theta}t}$$

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$$\Phi_g(t) \simeq \Phi_g^c$$

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$$H = \theta \cdot (t - t_B)$$

Where $\theta \gg H_c^2$ we interested in large scale perturbation

$$\partial_t^2 \Phi_g + \theta \Phi_g = 0$$

the solution is

$$\Phi_g = d_1 e^{i\sqrt{\theta}t} + d_2 e^{-i\sqrt{\theta}t}$$

Since the bounce phase is very short

$$\Phi_g(t) \simeq \Phi_g^c$$

This result is still true near bounce point since $k^2 \ll \theta$

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Interesting Features



- “Particle physics” energy density

$$\mathcal{E}_{PP} = \int d^3x T_{00} - c_* Q \sim \frac{1}{2} \dot{\pi}^2 + \frac{(\nabla^2 \pi)^2}{2M^2} + \dots$$

Inertial Mass!

- Gravitational energy density

$$\mathcal{E}_{grav} = T_{00} \sim M^2 \dot{\pi} + \dots$$

Gravitational Mass!

Interesting features

More generally,

$$\mathcal{L} = M^4 P(X) + M^2 S_1(X) (\Box \phi)^2 + M^2 S_2(X) \partial^\mu \partial^\nu \phi \partial_\mu \partial_\nu \phi + \dots$$

Ghost field locate at the minima, with scalar excitation

$$\phi = ct + \pi$$

Low energy effective action for π is

$$S \sim \int d^4x \left[\frac{1}{2} \dot{\pi}^2 - \frac{1}{2M^2} (\nabla^2 \pi)^2 + \dots \right],$$

The dispersion relation $\omega^2 \sim \frac{k^4}{M^2}$.

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