

Title: Relativity (PHYS 604) - Lecture 3

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Abstract:

Stress

Stress-Energy Tensor

Stress-Energy Tensor

EM $T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$

Stress-Energy Tensor

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

Conserved $\frac{\partial}{\partial x^{\mu}}$ $\left(\partial_{\mu} T^{\mu\nu} = 0 \right)$

Stress-Energy Tensor

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

Conserved $\partial_{\mu} T^{\mu\nu} = 0$

$$\frac{\partial}{\partial x^{\mu}}$$

using

$$\partial_{\mu} F^{\mu\nu} = 0$$

$$\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$$

Stress-Energy Tensor

EM $T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$

Conserved $\partial_{\mu} T^{\mu\nu} = 0$

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$\partial_{\mu} T^{\mu\nu}$

Stress-Energy Tensor

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

Conserved $\partial_{\mu} T^{\mu\nu} = 0$

$$\frac{\partial}{\partial x^{\mu}}$$

using

$$\partial_{\mu} F^{\mu\nu} = 0$$

$$\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$$

$$\eta^{\mu\nu} \partial_{\mu} = \partial^{\nu} \quad \partial_{\mu} T^{\mu\nu} = \underbrace{(\partial_{\mu} F^{\alpha\mu})}_{=0} F^{\nu}_{\alpha} + F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha} - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta}$$

Stress-Energy Tensor

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$\frac{\partial}{\partial x^{\mu}}$$

using

$$\partial_{\mu} F^{\mu\nu} = 0$$

$$\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$$

$$\partial_{\mu} T^{\mu\nu} = \underbrace{(\partial_{\mu} F^{\alpha\mu})}_{=0} F^{\nu}_{\alpha} + F^{\nu\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta}$$

$$\eta^{\mu\nu} \partial_{\mu} = \partial^{\nu}$$

Stress-Energy Tensor

EM $T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$

Conserved $\partial_{\mu} T^{\mu\nu}$

$$\frac{\partial}{\partial x^{\mu}}$$

$$F^{\mu\nu} = 0$$

Bianchi id.

$$\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$$

$$\eta^{\mu\nu}$$

$$F^{\alpha\mu} \partial F^{\nu}_{\alpha} - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta} + \partial_{\alpha} F_{\beta\gamma} + \partial_{\alpha} F_{\gamma\beta}$$

-Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Bianchi Id

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} = \partial_{\beta} F_{\gamma\alpha} + \partial_{\gamma} F_{\alpha\beta}$

$$\partial_{\mu} T^{\mu\nu} = (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta}$$

$$= F^{\nu}_{\mu} \partial^{\mu} F^{\alpha\mu}$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu})$$

$$- F^{\mu\alpha}$$

- Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} =$$

Bianchi id.

using

$$= 0$$

$$\partial_\alpha F_{\beta\gamma} + \text{cycles} = 0$$

$$\partial_\alpha (F^\alpha_\nu + F^{\alpha\mu} \frac{1}{2} (\partial_\mu F^\nu_\alpha - \partial_\alpha F^\nu_\mu) + \partial_\alpha F^\nu_\mu)$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^\nu_\beta$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F^\nu_\alpha - \partial_\alpha F^\nu_\mu) - F^{\alpha\mu} \partial_\mu F^\nu_\alpha$$

-Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\begin{aligned} T^{\mu\nu} &= (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} + F^{\alpha\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) \\ &\quad - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta} \end{aligned}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$\begin{aligned} &F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu}) \\ &\quad \parallel \\ &F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha} \end{aligned}$$

- Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\begin{aligned} \partial_{\mu} T^{\mu\nu} &= (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} + F^{\alpha\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) \\ &\quad - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta} \end{aligned}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$\begin{aligned} &F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu}) \\ &\quad \parallel \\ &F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha} \end{aligned}$$

-Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\begin{aligned} \partial_{\mu} T^{\mu\nu} &= (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} + F^{\alpha\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) \\ &\quad - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta} \end{aligned}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$\begin{aligned} &F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu}) \\ &\quad \parallel \\ &F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha} \end{aligned}$$

$$\text{c.f. } \frac{1}{2} F^{\alpha\mu} (\partial_{\mu} F^{\nu}_{\alpha} + \partial_{\alpha} F^{\nu}_{\mu})$$

- Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\partial_{\mu} T^{\mu\nu} = (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} + F^{\alpha\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu})$$

$$\parallel$$

$$F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha}$$

c.f. $\frac{1}{2} F^{\alpha\mu} (\partial_{\mu} F^{\nu}_{\alpha} + \partial_{\alpha} F^{\nu}_{\mu} + \partial^{\nu} F_{\alpha\mu})$

- Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\begin{aligned} \partial_{\mu} T^{\mu\nu} &= (\partial_{\mu} F^{\alpha\mu}) F^{\nu}_{\alpha} + F^{\alpha\mu} \left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right) \\ &\quad - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta} \end{aligned}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$\begin{aligned} &F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu}) \\ &\quad \parallel \\ &F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha} \end{aligned}$$

$$\text{c.f. } \frac{1}{2} F^{\alpha\mu} (\partial_{\mu} F^{\nu}_{\alpha} + \partial_{\alpha} F^{\nu}_{\mu} + \partial^{\nu} F_{\mu\alpha})$$

- Energy Tensor

$$T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{\mu\nu} = 0$$

Biandhi id.

using $\partial_{\mu} F^{\mu\nu} = 0$ $\partial_{\alpha} F_{\beta\gamma} + \text{cycles} = 0$

$$\partial_{\mu} T^{\mu\nu} = \underbrace{(\partial_{\mu} F^{\alpha\mu})}_{0} F^{\nu}_{\alpha} + F^{\alpha\mu} \underbrace{\left(\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu} \right)}_{0} - \frac{1}{2} F^{\alpha\beta} \partial^{\nu} F_{\alpha\beta}$$

$$F_{\alpha\beta} \eta^{\alpha\nu} = F^{\nu}_{\beta}$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_{\mu} F^{\nu}_{\alpha} - \partial_{\alpha} F^{\nu}_{\mu})$$

$$\parallel$$

$$F^{\alpha\mu} \partial_{\mu} F^{\nu}_{\alpha}$$

c.f. $\frac{1}{2} F^{\alpha\mu} \left(\overset{①}{\partial_{\mu} F^{\nu}_{\alpha}} + \overset{②}{\partial_{\alpha} F^{\nu}_{\mu}} + \overset{③}{\partial^{\nu} F_{\mu\alpha}} \right)$

$\alpha\beta$

chi id.

cycles = 0

$$\textcircled{1} \quad \left(\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu \right)$$

$\textcircled{2}$

$$+ \partial_\alpha F_\mu^\nu$$

$$\text{c.f. } \frac{1}{2} F^{\alpha\mu} \left(\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial^\nu F_{\mu\alpha} \right)$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\parallel$$
$$F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = T^{00}$

$$\pi^i = T^{i0}$$



$$\frac{\partial}{\partial t} \int d^3x$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\equiv F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\begin{aligned} & \textcircled{2} \\ & - \partial_\alpha F_\mu^\nu \\ & + \partial_\alpha F_\mu^\nu \end{aligned}$$

$$\frac{1}{2} F^{\alpha\mu} \left(\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha} \right)$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = T^{00}$

$$\pi^i = T^{i0}$$



$$\frac{\partial}{\partial t} \int d^3x \rho(t)$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\equiv F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\begin{aligned} & \textcircled{2} \\ & \underbrace{-\partial_\alpha F_\mu^\nu}_{+ \partial_\alpha F_\mu^\nu} \end{aligned}$$

$$\frac{1}{2} F^{\alpha\mu} \left(\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha} \right)$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = T^{00}$

$$\pi^i = T^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$



$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

$$= - \int_S \pi^i dS^i$$

Gauss divergence th.

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\equiv F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\textcircled{2} \quad \left(-\partial_\alpha F_\mu^\nu + \partial_\alpha F_\mu^\nu \right)$$

$$\frac{1}{2} F^{\alpha\mu} \left(\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha} \right)$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = T^{00}$

$$\pi^i = T^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$



$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

$$= - \int_S \pi^i dS^i$$

Gauss divergence th.

If S' chosen so π^i is zero

$\rightarrow E = \int d^3x \rho$ is conserved.

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\equiv F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\textcircled{2} \quad \underbrace{-\partial_\alpha F_\mu^\nu}_{+ \partial_\alpha F_\mu^\nu}$$

$$\frac{1}{2} F^{\alpha\mu} \left(\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha} \right)$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = T^{00}$

$$\pi^i = T^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$



$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

Gauss divergence thm.

$$= - \int_S \pi^i dS^i$$

If S chosen so π^i is zero

$\rightarrow E = \int d^3x \rho$ is conserved.

$$\partial_\mu F^{\mu\nu}$$

$$\partial_\mu F_{\mu\alpha}$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

define $\rho = \kappa T^{00}$

$$\pi^i = \kappa T^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$

$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

$$= - \int_S \pi^i dS^i$$

then so π^i is zero

$$\Rightarrow \int d^3x \rho \text{ is conserved.}$$

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

$$\equiv F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\begin{matrix} \textcircled{1} & \textcircled{2} & \textcircled{3} \\ \partial_\mu F_\alpha^\nu & + \partial_\alpha F_{\mu}^\nu & + \partial_\mu F_{\alpha}^\nu \end{matrix}$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

$$\partial_0 T^{0i} + \partial_j T^{ji} = 0$$

define $\rho = kT^{00}$

$$\pi^i = kT^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$



$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

$$= - \int_S \pi^i dS^i$$

Gauss divergence theorem

If S' chosen so π^i is zero
 $\rightarrow E = \int d^3x \rho$ is conserved.

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

" $F^{\alpha\mu} \partial_\mu F_\alpha^\nu$

$$\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha}$$

$$\partial_\mu T^{\mu\nu} = 0 \Rightarrow \partial_0 T^{00} + \partial_i T^{i0} = 0$$

$$\partial_0 T^{0i} + \partial_j T^{ji} = 0$$

$$P^i = \int d^3x T^{0i}$$

define $\rho = kT^{00}$

$$\pi^i = kT^{i0}$$

$$\frac{\partial \rho}{\partial t} = -\partial_i \pi^i$$



$$\frac{\partial}{\partial t} \left(\int d^3x \rho(t, \mathbf{x}) \right)$$

$$= - \int d^3x \partial_i \pi^i$$

$$= - \int_S \pi^i dS^i$$

Gauss divergence th.

If S' chosen so π^i is zero

$\rightarrow E = \int d^3x \rho$ is conserved.

$$F^{\alpha\mu} \frac{1}{2} (\partial_\mu F_\alpha^\nu - \partial_\alpha F_\mu^\nu)$$

=

$$F^{\alpha\mu} \partial_\mu F_\alpha^\nu$$

$$\textcircled{1} \partial_\mu F_\alpha^\nu + \textcircled{2} \partial_\alpha F_\mu^\nu + \textcircled{3} \partial_\nu F_{\mu\alpha}$$

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$T^{00} = F^{i0} F^{i0} + \frac{1}{4} (-2 F^{0i} F^{0i} + F^{ij} F^{ij})$$

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$\begin{aligned} T^{00} &= F^{i0} F^{i0} + \frac{1}{4} (-2 F^{0i} F^{0i} + F^{ij} F^{ij}) \\ &= \frac{1}{2} \left(\frac{\mathbf{E}^2}{c^2} + \right. \end{aligned}$$

$$\text{EM} \quad T^{\mu\nu} = F^{\alpha\mu} F^{\nu}_{\alpha} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$$

$$\begin{aligned} T^{00} &= F^{i0} F^{i0} + \frac{1}{4} (-2 F^{0i} F^{0i} + F^{ij} F^{ij}) \\ &= \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2) \end{aligned}$$

$$\left(\begin{array}{l} \text{recall} \\ \text{from before} \end{array} \right) \rho = \frac{1}{\mu_0} T^{00}$$

EM $T^{\mu\nu} = F^{\alpha\mu} F_{\alpha}^{\nu} - \frac{1}{4} \eta^{\mu\nu} F^{\alpha\beta} F_{\alpha\beta}$

$T^{00} = F^{i0} F^{i0} + \frac{1}{4} (-2 F^{0i} F^{0i} + F^{ij} F^{ij})$
 $= \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$

(recall from before $\rho = \frac{1}{\mu_0} T^{00}$)

$$T_{0i} = F_{j0} F_{ji}$$

$$T_{0i} = F_{j0} F_{ji} = - \frac{1}{c} \sum_j \epsilon_{jik} E_j B_k$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \underline{E} \wedge \underline{B} \right)_i$$

$$T^{0i} = F^{j0} F_j^i = - \frac{1}{c} \epsilon_{jik} E_j B_k$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \underline{E} \wedge \underline{B} \right)_i$$

Poynting $S^i = \frac{1}{\mu_0} T^{0i}$

$$T_{0i} = F_{j0} F_{ji} = - \frac{1}{c} \sum_j \epsilon_{jik} E_j B_k$$

$$\epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \underline{E} \wedge \underline{B} \right)_i$$

Poynting vector: $\frac{1}{\mu_0} T_{0i}$

Ex: T_{ii}

$$- \frac{1}{2} E_i E_i$$

$$T^{0i} = F^{j0} F_j^i = - \frac{1}{c} \epsilon_{jik} E_j B_k$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \underline{E} \wedge \underline{B} \right)_i$$

Poynting $S^i = \frac{1}{\mu_0} T^{0i}$

$$T^{ij} = -B_i B_j - \frac{E_i E_j}{c^2} + \delta_{ij} \left(\frac{E^2}{2c^2} + \frac{1}{2} B^2 \right)$$

$$T^{0i} = F^{j0} F_j^i = -\frac{\mathbf{E}^j}{c} \epsilon_{jik} B^k$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \mathbf{E} \wedge \mathbf{B} \right)_i$$

Poynting $S^i = \frac{1}{\mu_0} T^{0i}$

Ex: $T^{ij} = -B_i B_j - \frac{E_i E_j}{c^2} + \delta_{ij} \left(\frac{E^2}{2c^2} + \frac{1}{2} B^2 \right)$

spatial stress in EM field.

$$T^{0i} = F^{j0} F_j^i = -\frac{\mathbf{E} \cdot \mathbf{B}}{c}$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \mathbf{E} \wedge \mathbf{B} \right)_i$$

Poynting $S^i = \frac{1}{\mu_0} T^{0i}$

Ex: $T^{ij} = -B_i B_j - \frac{E_i E_j}{c^2} + \delta_{ij} \left(\frac{E^2}{2c^2} + \frac{1}{2} B^2 \right)$

spatial stress in EM field.

$$\partial_j T^{0i}$$

$$= -\partial_j T^{ji}$$

$$T^{0i} = F^{j0} F_j^i = -\frac{\epsilon_0}{c} \epsilon_{jik} E_j B_k$$

$$= \epsilon_{ijk} \frac{E_j}{c} B_k = \left(\frac{1}{c} \underline{E} \wedge \underline{B} \right)_i$$

Poynting $S^i = \frac{1}{\mu_0} T^{0i}$

Ex: $T^{ij} = -B_i B_j - \frac{E_i E_j}{c^2} + \delta_{ij} \left(\frac{E^2}{2c^2} + \frac{1}{2} B^2 \right)$

spatial stress in EM field.

$$\partial_i \int T^{0i}$$

$$= - \int \partial_j T^{ji}$$

$$T^{00} + \nabla^2 \phi$$

$$T^{00} + \nabla^2 \phi$$

$$\partial_\alpha T^{\alpha\mu} = 0$$

$$x'^\mu = x^\mu + c^\mu$$

Emmy Noether

Lorentz Force Law

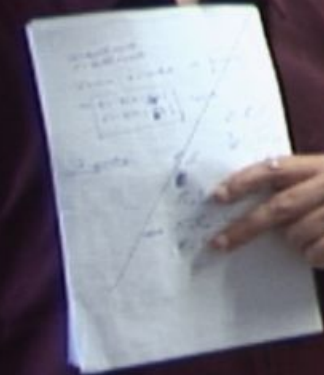
$$\underline{F} = m \ddot{\underline{x}} = q (\underline{E} + \underline{v} \wedge \underline{B})$$

4-acc:

$$m_0 \frac{dU^\mu}{d\tau} = q F^{\mu\nu} U_\nu$$

$$U^\mu(\tau) = \frac{dx^\mu}{d\tau}$$

$$x^\mu(\tau)$$



Lorentz Force Law

$$\mathbf{F} = m\mathbf{\ddot{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

4-acc?

$$m_0 \frac{dU^\mu}{d\tau} = q F^\mu{}_\nu U^\nu$$

$$P^\mu = m_0 U^\mu$$

$$\frac{dp^i}{d\tau} = q \left(F^i{}_0 \frac{d(ct)}{d\tau} + F^i{}_j \frac{dx^j}{d\tau} \right)$$

$$\times \frac{d\tau}{dt}$$

$$\frac{dp^i}{dt} = q \left(\frac{E^i}{c} + \epsilon^{ijk} B_k v_j \right)$$

Lorentz Force Law

$$\mathbf{F} = m\ddot{\mathbf{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

4-accⁿ

$$m_0 \frac{dU^\mu}{d\tau} = q F^\mu{}_\nu U^\nu$$

$$p^\mu = m_0 U^\mu$$

$$U^\mu(\tau) = \frac{dx^\mu}{d\tau}$$

$$x^\mu(\tau)$$

$$v^j = \frac{dx^j}{dt}$$

$$= \left(F^i{}_0 \frac{d(ct)}{d\tau} + F^i{}_j \frac{dx^j}{d\tau} \right)$$

$$= q \left(\frac{E^i}{c} + \epsilon^{ijk} B_k v^j \right)$$

Lorentz Force Law

$$F = m\ddot{x} = q(E + v \wedge B)$$

4-acc?

$$m_0 \frac{du^\mu}{d\tau} = q F^\mu{}_\nu u^\nu$$

$$p^\mu = m_0 u^\mu$$

$$\frac{dp^i}{d\tau} = q \left(F^i{}_0 \frac{d(ct)}{d\tau} + F^i{}_j \frac{dx^j}{d\tau} \right)$$

$$\begin{aligned} \frac{dp^i}{dt} &= q \left(\frac{E^i}{c} + \epsilon^{ijk} B_k v^j \right) \\ &= q \left(E^i + (v \wedge B)^i \right) \end{aligned}$$

$$\times \frac{d\tau}{dt}$$

$$u^\mu(\tau) = \frac{dx^\mu}{d\tau}$$

$$x^\mu(\tau)$$

$$v^j = \frac{dx^j}{dt}$$

Lorentz Force Law

$$\mathbf{F} = m\mathbf{\ddot{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

4-acc?

$$m_0 \frac{dU^\mu}{d\tau} = q F^\mu{}_\nu U^\nu$$

$$p^\mu = m_0 U^\mu$$

$$\frac{dp^i}{d\tau} = q (F^i{}_\nu U^\nu)$$

$$\frac{dp^i}{dt} = q \left(\mathbf{E} + \mathbf{v} \wedge \mathbf{B} \right)^i$$

$$\times \frac{d\tau}{dt}$$

$$U^\mu(\tau) = \frac{dx^\mu}{d\tau}$$

$$x^\mu(\tau)$$

$$v^j = \frac{dx^j}{dt}$$

Lorentz Force Law

$$\mathbf{F} = m\mathbf{\ddot{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

4-acc?

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$$U^\mu(\tau) = \frac{dx^\mu}{d\tau}$$

$$x^\mu(\tau)$$

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$$\frac{dp^i}{d\tau} = q \left(F^i{}_0 \frac{d(ct)}{d\tau} + F^i{}_j \frac{dx^j}{d\tau} \right)$$

$$v^j = \frac{dx^j}{dt}$$

$$\times \frac{d\tau}{dt}$$

$$\begin{aligned} \frac{dp^i}{dt} &= q \left(\frac{E^i}{c} + \epsilon^{ij} B_k v^j \right) \\ &= q \left(E^i + (\mathbf{v} \wedge \mathbf{B})^i \right) \end{aligned}$$

Lorentz Force Law

Note $\underline{F} = \frac{d\underline{p}}{dt}$

$$\underline{F} = m\ddot{\underline{x}} = q(\underline{E} + \underline{v} \wedge \underline{B})$$

4-acc?

$$m_0 \frac{du^\mu}{d\tau} = q F^\mu{}_\nu u^\nu$$

$$\frac{dp^i}{d\tau} = q \left(F^i{}_0 \frac{d(ct)}{d\tau} + F^i{}_j \frac{dx^j}{d\tau} \right)$$

$$\begin{aligned} \frac{dp^i}{d\tau} &= q \left(\frac{E^i}{c} c + \epsilon^{ijk} B_k v^j \right) \\ &= q \left(E^i + (\underline{v} \wedge \underline{B})^i \right) \end{aligned}$$

Force Law

$$= m\ddot{\mathbf{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

cc?

$$\frac{dU^M}{d\tau} = q F^M_{\nu} U^{\nu}$$

$$\frac{dp^i}{d\tau} = q \left(F^i_0 \frac{d(ct)}{d\tau} + F^i_j \frac{dx^j}{d\tau} \right)$$

$$\frac{dp^i}{d\tau} = q \left(\frac{E^i}{c} c + \epsilon^{ij} B_k V^j \right)$$

$$= q \left(E^i + (\mathbf{v} \wedge \mathbf{B})^i \right)$$

Note $\underline{F} = \frac{d\underline{p}}{d\underline{t}}$ not $\frac{d\underline{p}}{d\underline{\tau}}$

$$\frac{dp^0}{d\tau} = q F^0_i U^i$$

Force Law

$$= m\ddot{\mathbf{x}} = q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

cc:

$$\frac{dU^M}{d\tau} = q F^M_{\nu} U^{\nu}$$

$$\frac{dp^i}{d\tau} = q \left(F^i_0 \frac{d(ct)}{d\tau} + F^i_j \frac{dx^j}{d\tau} \right)$$

$$\frac{dp^i}{d\tau} = q \left(\frac{E^i}{c} + \epsilon^{ij} B_k v^j \right)$$

$$= q \left(E^i + (\mathbf{v} \wedge \mathbf{B})^i \right)$$

Note $\underline{F} = \frac{d\underline{p}}{d\underline{t}}$ not $\frac{d\underline{p}}{d\underline{\tau}}$

$$\frac{dp^0}{d\tau} = q F^0_i U^i = q \frac{E^i}{c}$$

the Law

$$= q(E + \mathbf{v} \wedge \mathbf{B})$$

$$= \mu \mathbf{v} \cdot \mathbf{u}^{\mathbf{v}}$$

$$\mathcal{L} \left(F^i_0 \frac{dx^0}{d\tau} + F^i_j \frac{dx^j}{d\tau} \right)$$

$$q \left(E^i + \epsilon^{ij} B_k v^j \right)$$

$$+ (\mathbf{v} \wedge \mathbf{B})^i$$

Note $\underline{F} = \frac{d\underline{p}}{d\tau}$ not $\frac{d\underline{p}}{dt}$

$$\frac{dp^0}{d\tau} = q F^0_i u^i = q \frac{E^i}{c} \cdot \frac{dx^i}{d\tau}$$

Law

$$= q(\mathbf{E} + \mathbf{v} \wedge \mathbf{B})$$

$$= q F^{\mu}_{\nu} u^{\nu}$$

$$\frac{dp^i}{dt} = q \left(F^i_0 \frac{dx^0}{dt} + F^i_j \frac{dx^j}{dt} \right)$$

$$\frac{dp^i}{dt} = q \left(\frac{E^i}{c} + \epsilon^{ijk} B_k v_j \right)$$

$$= q \left(\mathbf{E} + (\mathbf{v} \wedge \mathbf{B})^i \right)$$

Note $\underline{F} = \frac{d\underline{p}}{dt}$ not $\frac{d\underline{p}}{d\tau}$

$$\frac{dp^0}{d\tau} = q F^0_i u^i = q \frac{E^i}{c} \cdot \frac{dx^i}{d\tau}$$

$$\frac{dp^0}{dt} = q \mathbf{E} \cdot \mathbf{v}$$

the Law

$$= q(E + v \wedge B)$$

$$= q F^{\mu} v_{\mu}$$

$$\frac{dp^i}{dt} = q \left(E^i + \epsilon^{ijk} v_j B_k \right)$$

$$\frac{dp^i}{dt} = q$$

$$= q$$

Note $\underline{F} = \frac{d\underline{p}}{dt}$ not $\frac{dP}{dt}$

$$\frac{dP^0}{dt} = q F^0_i v^i = q \frac{E^i}{c} \cdot \frac{dx^i}{dt}$$

$$\frac{dP^0}{dt} = q \underbrace{E \cdot v}_{\text{Force}} = \text{Work done.}$$

General Relativity

General Relativity

Generalisation of Minkowski spacetime

General Relativity

Generalisation of Minkowski spacetime
← differentiable manifold

General Relativity

Generalisation of Minkowski spacetime

← differentiable manifold

topological space

General Relativity

Generalisation of Minkowski spacetime

← differentiable manifold

topological space
which is "locally $\mathbb{R}^n$ "

General Relativity

Generalisation of Minkowski spacetime
← differentiable manifold

manifold space
which is "locally $\mathbb{R}^n$ "

Note

$$\frac{dp}{dt}$$

General Relativity

Generalisation of Minkowski spacetime
→ differentiable manifold

topological space
which is "locally $\mathbb{R}^n$ "



Note

$$\frac{dp}{dt}$$

General Relativity

Generalisation of Minkowski spacetime

← differentiable manifold

topological space
which is "locally $\mathbb{R}^n$ "

eg. sphere, torus



cone



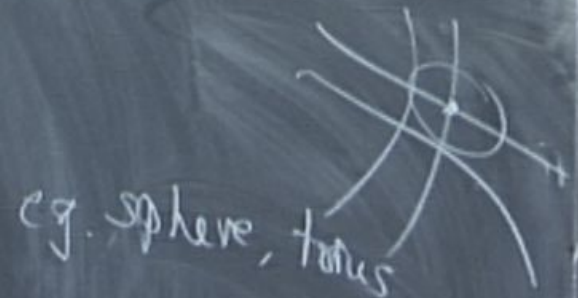
General Relativity

Generalisation of Minkowski spacetime

← differentiable manifold

usual operations
of calculus
should apply

topological space
which is "locally $\mathbb{R}^n$ "



General Relativity

Generalisation of Minkowski spacetime

← differentiable manifold

usual operations
of calculus
should apply

topological space
which is "locally $\mathbb{R}^n$ "

eg. sphere, torus



General Relativity

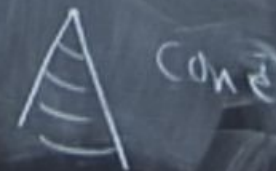
Generalisation of Minkowski spacetime

← differentiable manifold

usual operations
of calculus
should apply

topological space
which is "locally

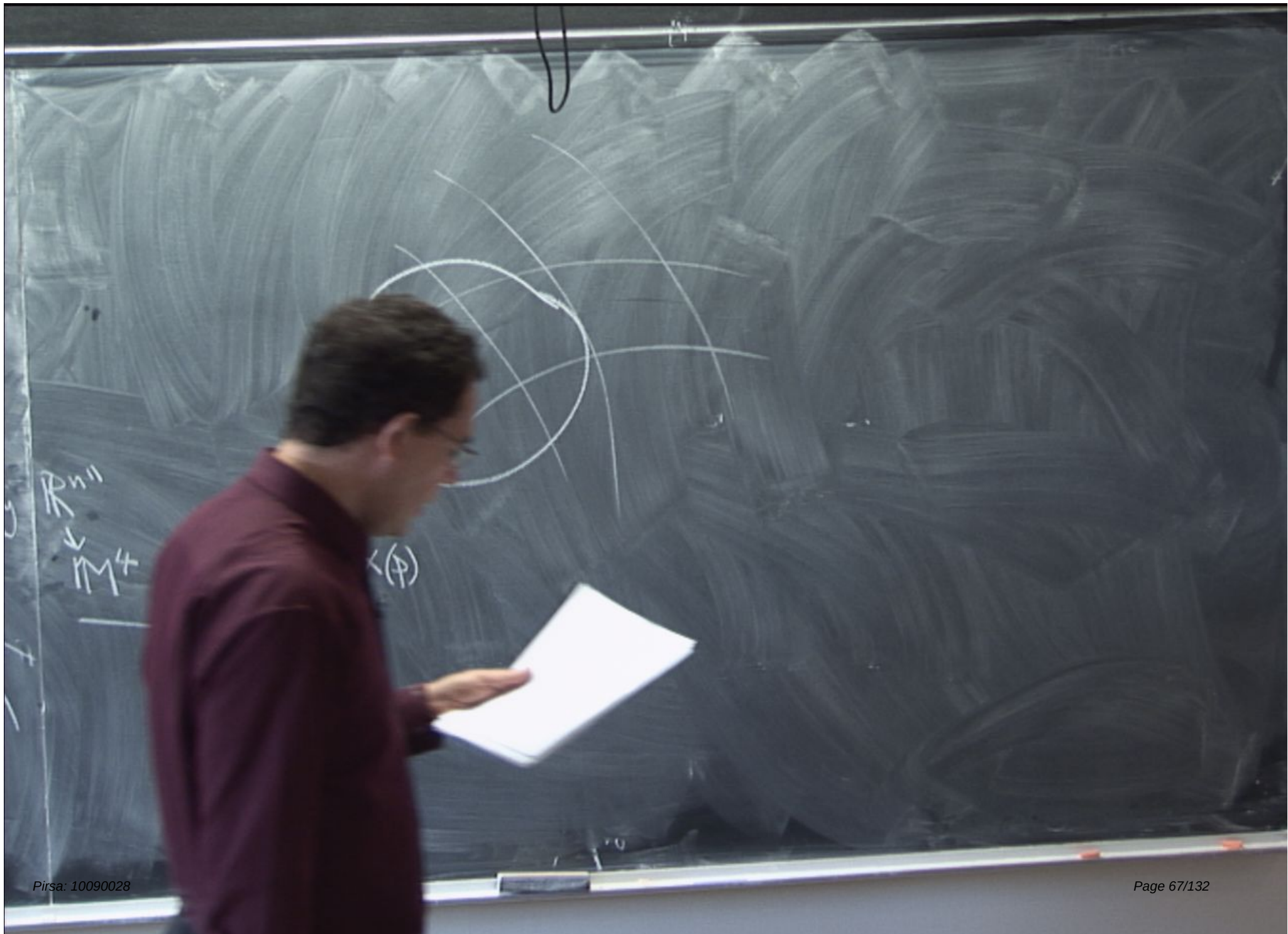
eg. sphere, torus

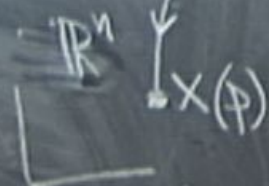


cone









every p has
unique $x(p)$

$\mathbb{R}^n$
 $\downarrow$
 M^4



$\mathbb{R}^n$
 $\downarrow$
 M^4

$\mathbb{R}^n$
 $\downarrow$
 $X(P)$

every p has
unique $X(p)$

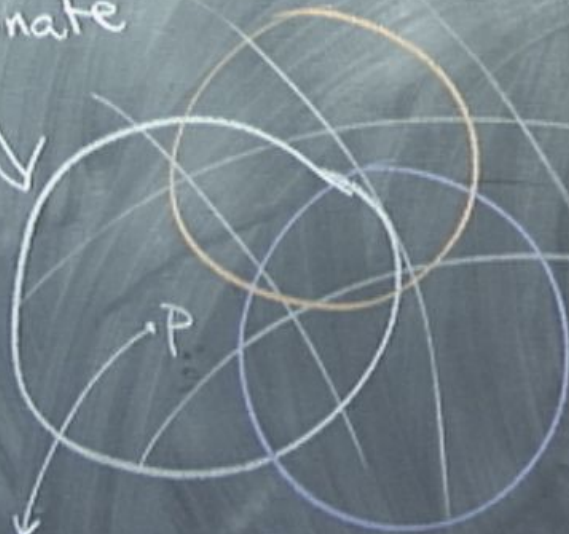
Coordinate
patch



$\mathbb{R}^n$
 $\downarrow$
 M^4

every p has
unique $x(p)$

Coordinate
patch



p

$\mathbb{R}^n$

$x(p)$

every p has
unique $x(p)$

$\mathbb{R}^n$

M^n

Coordinate
patch



$\mathbb{R}^n$
 $\downarrow$
 M^n

x
 $\mathbb{R}^n$
 $x(p)$

every p has
unique $x(p)$

Coordinate
patch



P

x

$\mathbb{R}^n$

$x(p)$

every p has
unique $x(p)$

$\mathbb{R}^n$

M^4

$\mathbb{R}^n$

Coordinate
patch

$$\begin{array}{c} x' \\ \uparrow \\ \mathbb{R}^n \\ \downarrow \\ x'(p) \end{array}$$

$$\begin{array}{c} x'' \\ \uparrow \\ \mathbb{R}^n \\ \downarrow \\ x''(p) \end{array}$$

$$\begin{array}{c} \mathbb{R}^{n+1} \\ \downarrow \\ \mathbb{M}^n \end{array}$$

$$\begin{array}{c} x \\ \uparrow \\ \mathbb{R}^n \\ \downarrow \\ x(p) \end{array}$$

every p has
unique $x(p)$

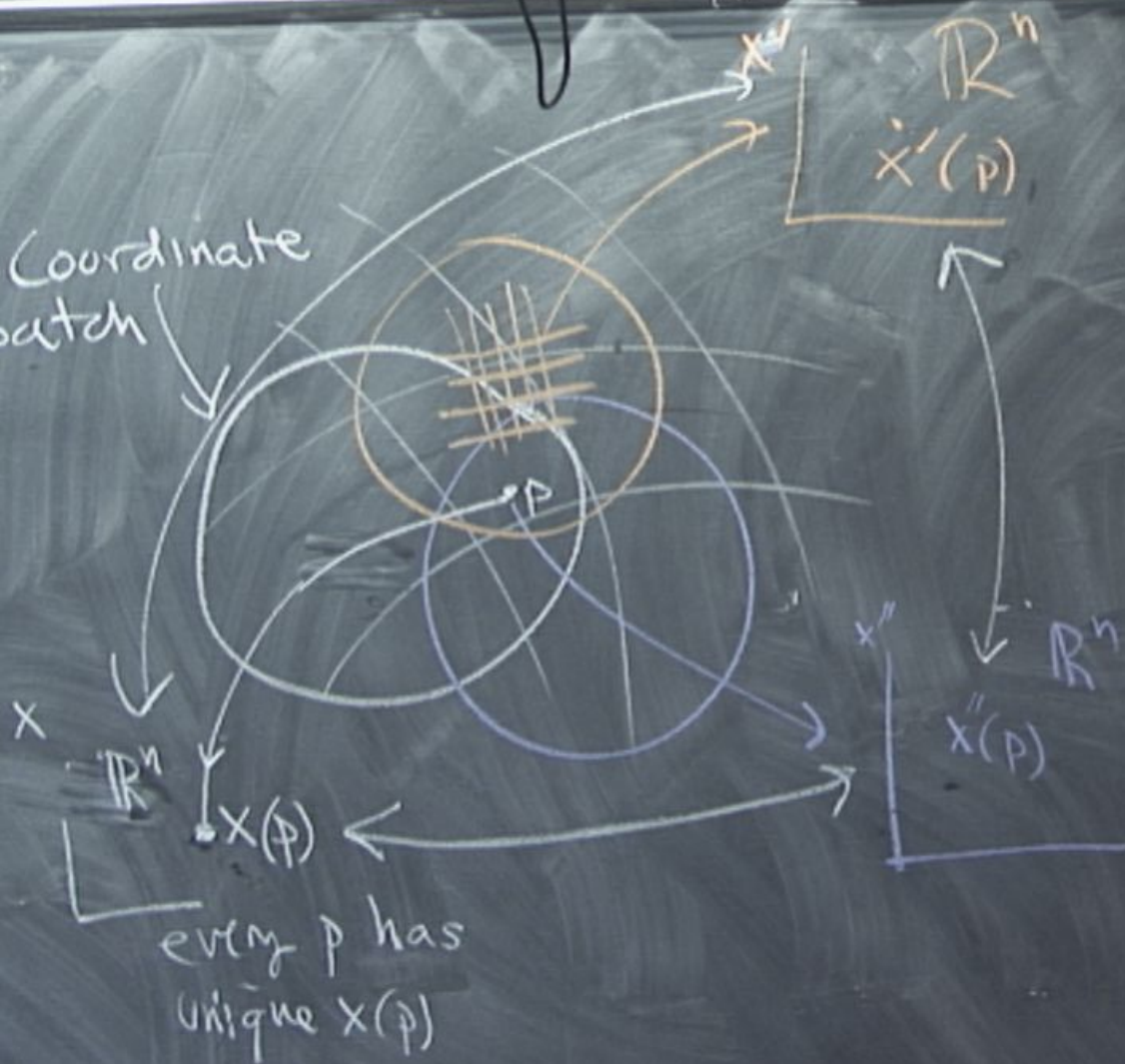
intersection regions: coords should
be compatible

$$x \longrightarrow x''$$

intersection regions: coords should
be compatible

$x \longrightarrow x''$

Coordinate
patch



every p has
unique $x(p)$

intersection regions: coords should
be compatible



must have $x(x) = x''(x'(x))$

in the overlap regions

intersection regions: coords should
be compatible



must have $x''(x) = x''(x'(x))$
in triple overlap regions

eg. S^2



eg. S^2



eg. S^2



$$x^2 + y^2 + z^2 = 1$$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

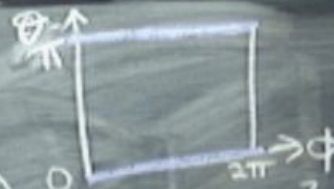
$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

- Singular at NP, SP because
some pt represented by any $0 < \phi \leq 2\pi$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

- Singular at NP, SP because
some pt represented by any $0 < \phi \leq 2\pi$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

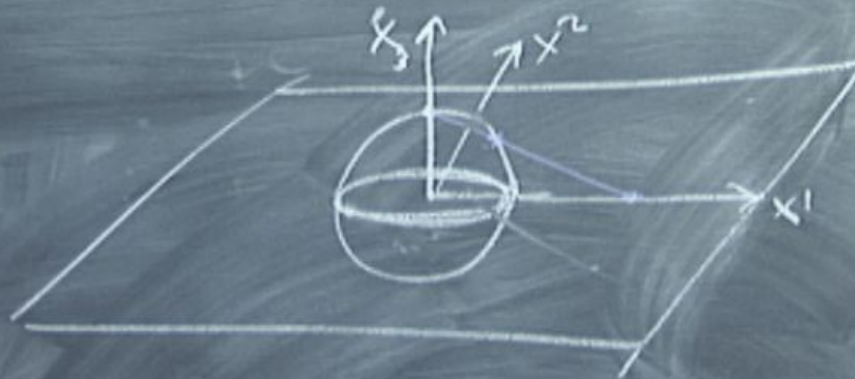
$$z = \cos \theta$$

- Singular at NP, SP because
some pt represented by any $0 < \phi \leq 2\pi$

Projective coords



Projective coords



Projective coords

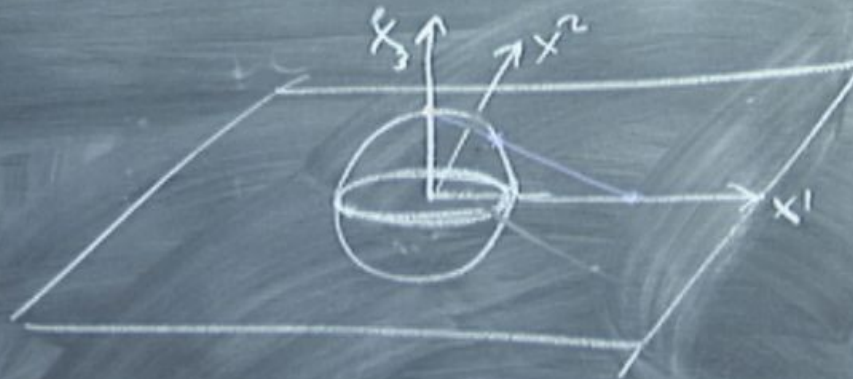


Projective coords



coords on plane (z_1, z_2)

Projective coords



coords on plane (z_1, z_2)

$$\left(\frac{x_1}{1-x_3}, \frac{x_2}{1-x_3} \right)$$

Projective coords



coords on plane (z_1, z_2)

NP proj coords $\left(\frac{x_1}{1-x_3}, \frac{x_2}{1-x_3} \right)$

Sp " " $\left(\frac{x_1}{1+x_3}, \frac{x_2}{1+x_3} \right)$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

- Singular at NP, SP because
some pt represented by any $0 < \phi \leq 2\pi$

eg. S^2



$$x^2 + y^2 + z^2 = 1$$

$$x = \sin \theta \cos \phi$$

$$y = \sin \theta \sin \phi$$

$$z = \cos \theta$$

- Singular at NP, SP because
some pt represented by any $0 < \phi \leq 2\pi$

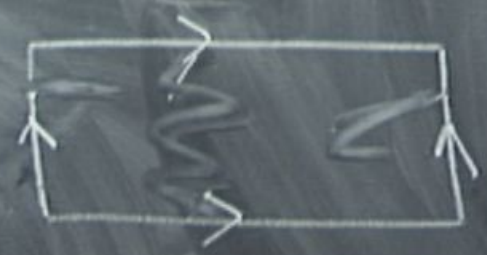
ords



me (z_1, z_2)

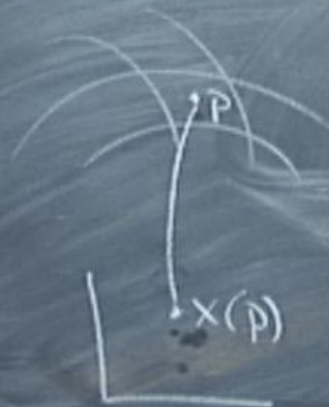
ords $\left(\frac{x_1}{1-x_3}, \frac{x_2}{1-x_3} \right)$
 " $\left(\frac{x_1}{1+x_3}, \frac{x_2}{1+x_3} \right)$

eg. S^2



- Singular at NP, SP
 some pt represented

Tensors on Manifolds



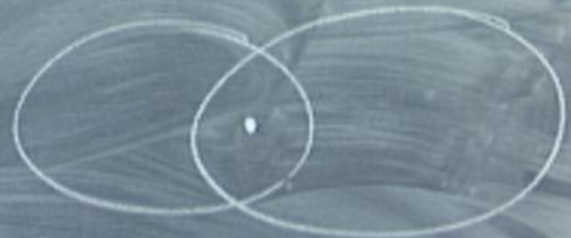
Tensors on Manifolds

Scalar
field

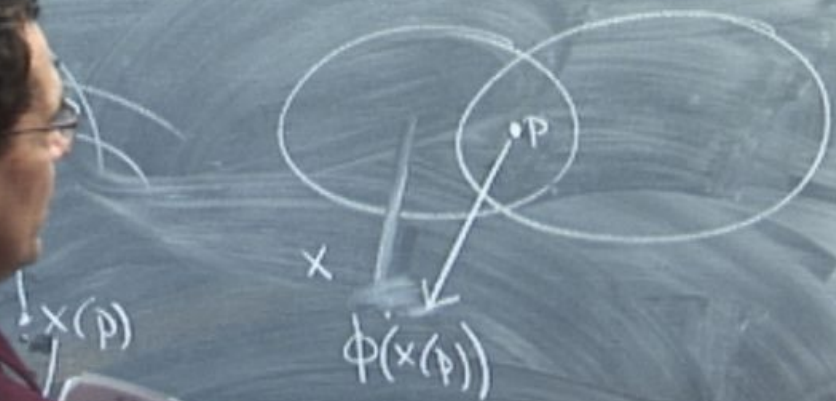
$$\phi(x(p))$$

$x(p)$

P



Tensors on Manifolds



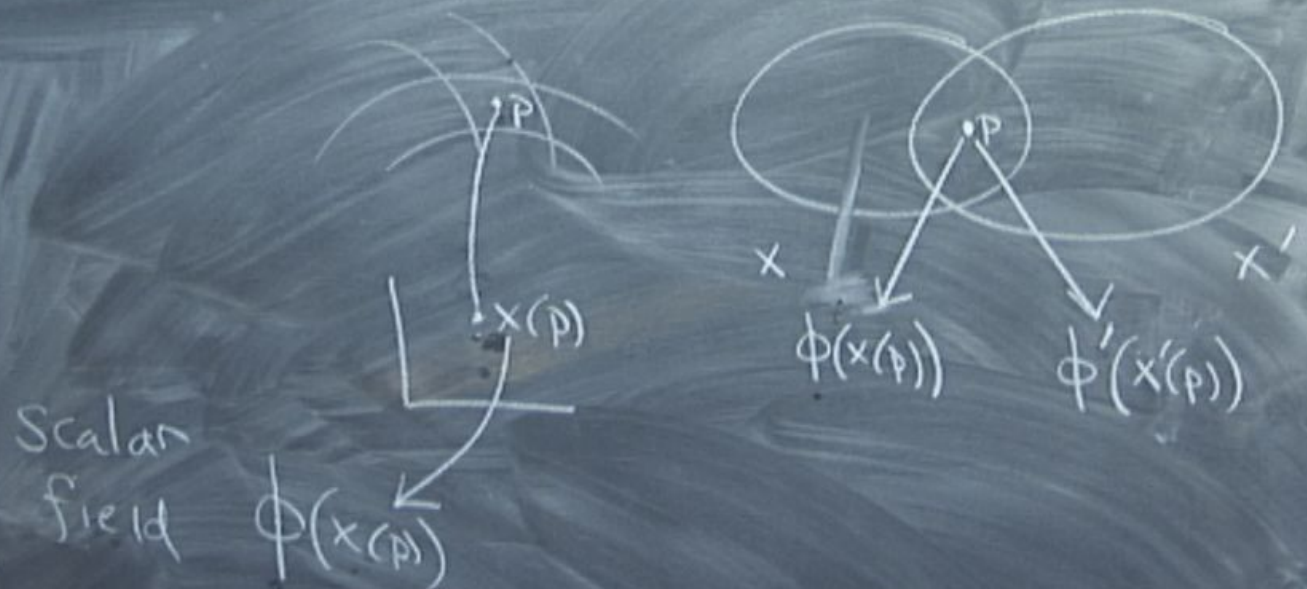
Tensors on Manifolds

Scalar
field

$$\phi(x(p))$$



Tensors on Manifolds



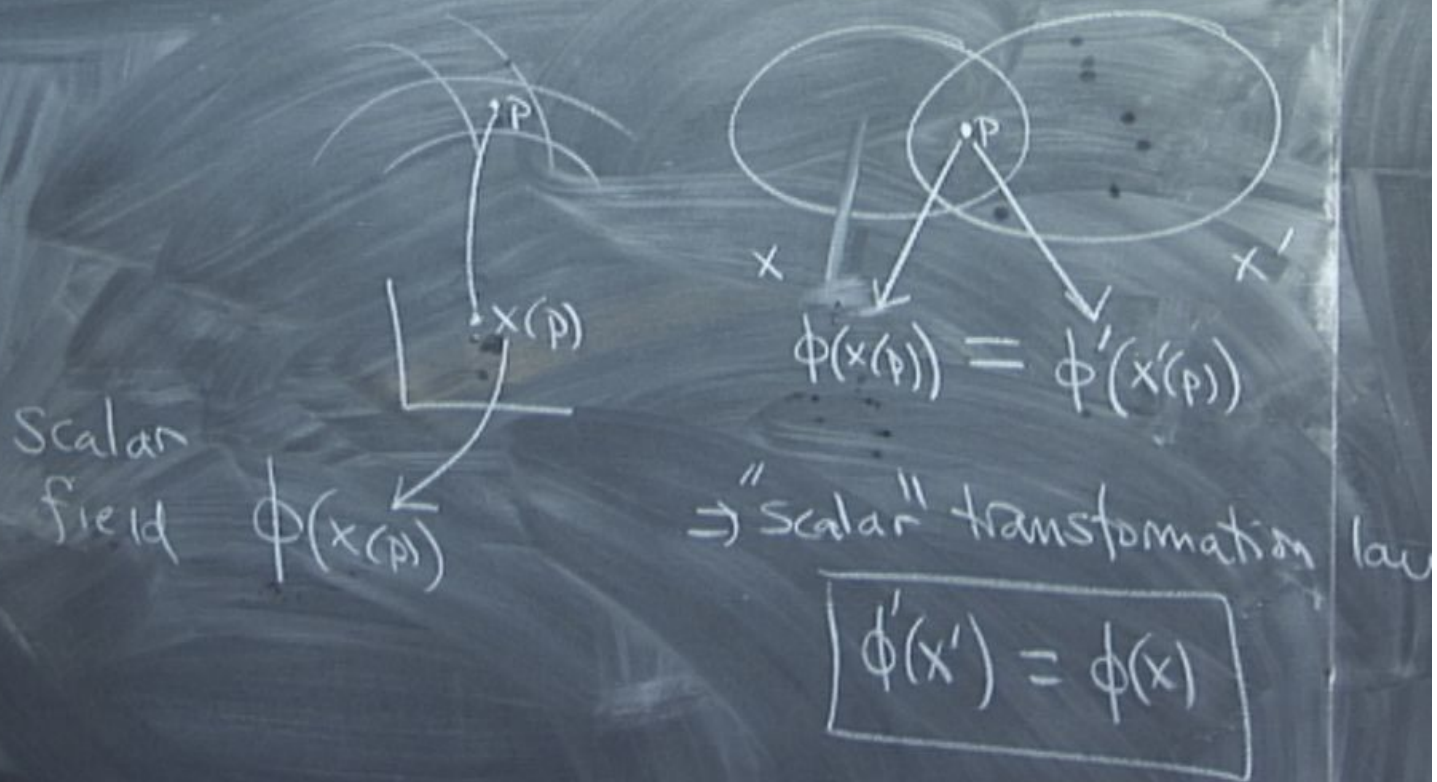
Scalar
field

$$\phi(x(p))$$

$$\phi(x(p))$$

$$\phi'(x'(p))$$

Tensors on Manifolds



Contravariant vectors

$$x^{\mu}(\lambda)$$

law

Contravariant vectors

$\mathbb{R}^n$

(x)

:

dx^m

Contravariant vectors

$\mathbb{R}^n$

$x^M(\lambda)$

: $\frac{dx^M}{d\lambda}$ is a contravariant vector

$$V^{M'}(x') = \frac{\partial x'^M}{\partial x^\alpha} V^\alpha(x)$$

Contravariant vectors

$\mathbb{R}^n$
 $x^{\mu}(\lambda)$: $\frac{dx^{\mu}}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x'^{\mu}}{\partial x^{\alpha}} V^{\alpha}(x)$$

law

Contravariant vectors

$\mathbb{R}^n$
 $x^\mu(\lambda) : \frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^\alpha} V^\alpha(x)$$

Covariant vectors

Contravariant vectors

$\mathbb{R}^n$

$x^{\mu}(\lambda)$: $\frac{dx^{\mu}}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^{\alpha}} V^{\alpha}(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^{\mu}} = \frac{\partial x^{\alpha}}{\partial x^{\mu}} \frac{\partial}{\partial x^{\alpha}}$

Contravariant vectors

$\mathbb{R}^n$

$x^{\mu}(\lambda)$: $\frac{dx^{\mu}}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^{\alpha}} V^{\alpha}(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^{\mu}} = \frac{\partial x^{\alpha}}{\partial x^{\mu}} \frac{\partial}{\partial x^{\alpha}}$

Contravariant vectors

$\mathbb{R}^n$
 $x^\mu(\lambda) : \frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^\alpha} V^\alpha(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^\mu} = \frac{\partial x^\alpha}{\partial x^{\mu'}} \frac{\partial}{\partial x^\alpha}$

Contravariant vectors

$\mathbb{R}^n$

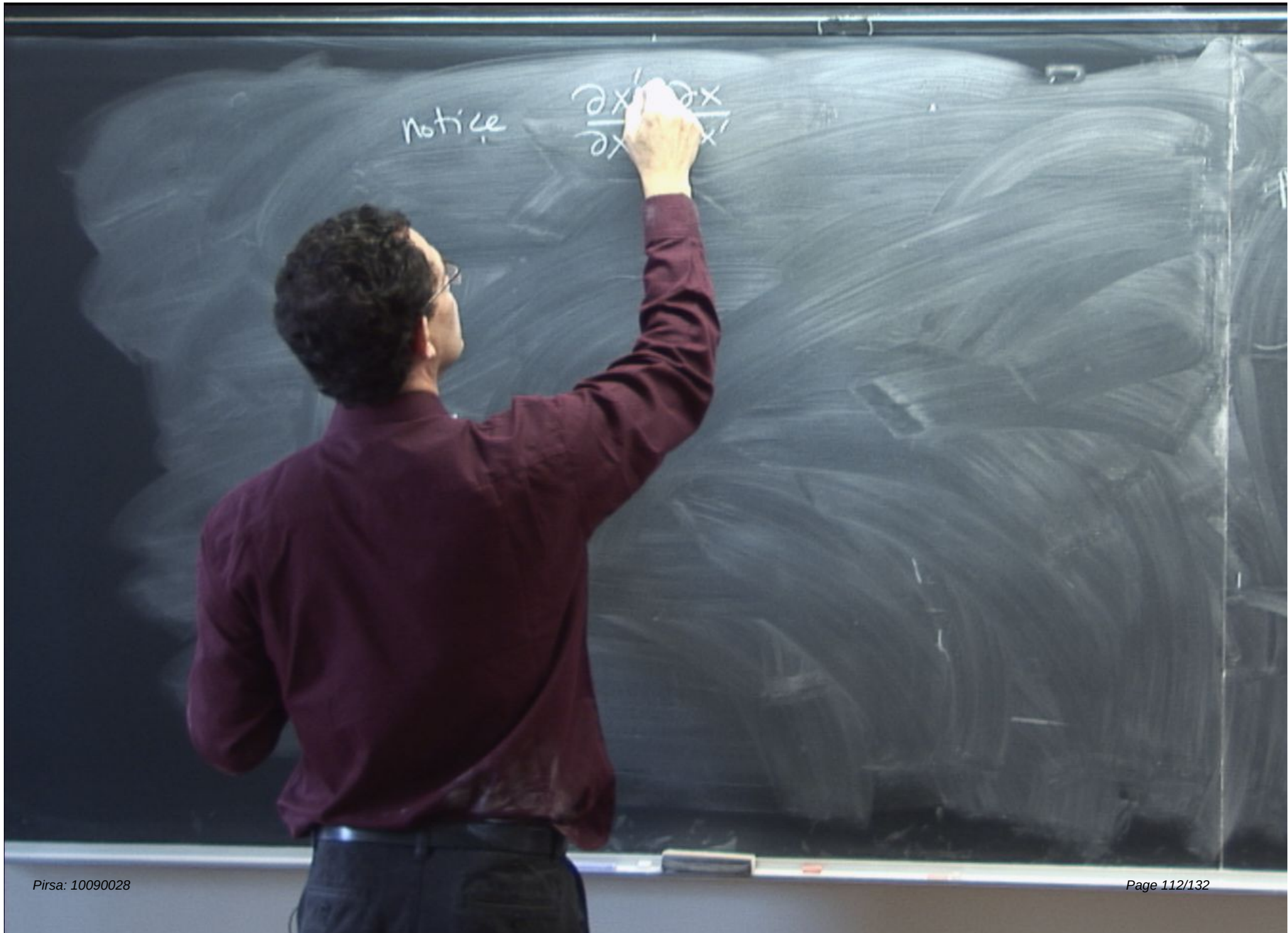
$x^\mu(\lambda)$: $\frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^\alpha} V^\alpha(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^{\mu'}} = \frac{\partial x^\alpha}{\partial x^{\mu'}} \frac{\partial}{\partial x^\alpha}$ in general

$$W_{\mu'}(x') = \frac{\partial x^\alpha}{\partial x^{\mu'}} W_\alpha(x)$$



notice

$$\frac{\partial x}{\partial x} \quad \frac{\partial x}{\partial x}$$

notice

$$\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$$

↔
inverses

Contravariant vectors

$\mathbb{R}^n$

$x^\mu(\lambda) : \frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V'^\mu = \Lambda^\mu_\alpha V^\alpha$$

$$V'^\mu(x') = \frac{\partial x'^\mu}{\partial x^\alpha} V^\alpha(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^\mu} = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial}{\partial x^\alpha}$ in general

$$W'_\mu(x') = \frac{\partial x^\alpha}{\partial x'^\mu} W_\alpha(x)$$

Contravariant vectors

$\mathbb{R}^n$

$x^\mu(\lambda)$: $\frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V'^\mu = \Lambda^\mu_\alpha V^\alpha$$

$$V'^\mu(x') = \frac{\partial x'^\mu}{\partial x^\alpha} V^\alpha(x)$$

$$W'_\mu = \tilde{\Lambda}^\alpha_\mu W_\alpha$$

Covariant vectors

e.g. $\frac{\partial}{\partial x'^\mu} = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial}{\partial x^\alpha}$ in general

$$W'_\mu(x') = \frac{\partial x^\alpha}{\partial x'^\mu} W_\alpha(x)$$

Contravariant vectors

$\mathbb{R}^n$

$x^\mu(\lambda)$: $\frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V^{\mu'}(x') = \frac{\partial x^{\mu'}}{\partial x^\alpha} V^\alpha(x)$$

$$V^{\mu'}(x') = \Lambda^\mu_{\alpha} V^\alpha$$

$$W_\mu = \tilde{\Lambda}^\alpha_\mu V^\alpha$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^{\mu'}} = \frac{\partial x^\alpha}{\partial x^{\mu'}} \frac{\partial}{\partial x^\alpha}$

$$W_\mu(x') = \frac{\partial x^\alpha}{\partial x^{\mu'}} W_\alpha(x)$$

Contravariant vectors

$$F'^{\alpha\beta}_{(x')} = \Lambda^\alpha_\mu \Lambda^\beta_\nu F^{\mu\nu}(x)$$

$\mathbb{R}^n$

$x^\mu(\lambda)$: $\frac{dx^\mu}{d\lambda}$ is a contravariant vector

$$V'^{\mu}(x') = \frac{\partial x'^{\mu}}{\partial x^{\alpha}} V^{\alpha}(x)$$

$$V'^{\mu}_{(x')} = \Lambda^{\mu}_{\alpha} V^{\alpha}(x)$$

$$W'_{\mu} = \tilde{\Lambda}^{\alpha}_{\mu} W_{\alpha}$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^{\mu}} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial}{\partial x^{\alpha}}$ in general

$$W'_{\mu}(x') = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} W_{\alpha}(x)$$

Contravariant vectors

$$F'^{\alpha\beta}_{(x')} = \Lambda^\alpha_\mu \Lambda^\beta_\nu F^{\mu\nu}(x)$$

$\mathbb{R}^n$

$x^\mu(\lambda)$

$$\frac{dx^\mu}{d\lambda}$$

is a contravariant vector

$$x'^\mu = \Lambda^\mu_\alpha x^\alpha$$

$V^\alpha(x)$

$$V'^\mu(x') = \frac{\partial x'^\mu}{\partial x^\alpha} V^\alpha(x)$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^\mu} = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial}{\partial x^\alpha}$

in general

(x)

Contravariant vectors

$$F'^{\alpha\beta}(x') = \Lambda^\alpha_\mu \Lambda^\beta_\nu F^{\mu\nu}(x)$$

$\mathbb{R}^n$

$x^\mu(\lambda)$

$$\frac{dx^\mu}{d\lambda}$$

is a contravariant vector

$$x'^\mu = \Lambda^\mu_\alpha x^\alpha$$

$$V'^\mu(x') = \Lambda^\mu_\alpha V^\alpha(x)$$

$$V'^\mu(x') = \frac{\partial x'^\mu}{\partial x^\alpha} V^\alpha(x)$$

$$W'_\mu = \tilde{\Lambda}^\alpha_\mu W_\alpha$$

Covariant vectors

e.g. $\frac{\partial}{\partial x^\mu} = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial}{\partial x^\alpha}$

in general

$$W'_\mu(x') = \frac{\partial x^\alpha}{\partial x'^\mu} W_\alpha(x)$$

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

$\swarrow \quad \searrow$
 inverses

$x'(x)$ well-behaved provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

↔
inverses

$x'(x)$ well-behaved provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

Contravariant

$\mathbb{R}^n$ $x^{\mu}(\lambda) : \frac{dx^{\mu}}{d\lambda}$

$V^{\mu'}(x')$

Covariant vec

e.g. $\frac{\partial}{\partial x'^{\mu}} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial}{\partial x^{\alpha}}$

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

$\swarrow \searrow$
 inverses

$x'(x)$ well-behaved

$$\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$$

invertible.

similarly, tensor of

$$T^{\mu\nu}_{\alpha\beta}$$

$$\begin{bmatrix} m \\ n \end{bmatrix}$$

up
down

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

↔
inverses

$x'(x)$ well-defined provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

similarly, to valence $\begin{bmatrix} m \\ n \end{bmatrix}$ $\leftarrow \begin{matrix} \# \text{ up} \\ \# \text{ down} \end{matrix}$

$$= \left(\frac{\partial x'^{\alpha}}{\partial x^{\beta}} \right) \frac{\partial x^{\beta}}{\partial x'^{\alpha}}$$

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

$\swarrow \searrow$
 inverses

$x'(x)$ well-behaved provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

similarly, tensor of valence $\begin{bmatrix} m \\ n \end{bmatrix}$ $\begin{matrix} \leftarrow \# \text{ up} \\ \leftarrow \# \text{ down} \end{matrix}$

$$T^{\overbrace{\alpha\beta\gamma\dots}^m}_{\underbrace{\mu\nu\dots}_n}(x') = \left(\frac{\partial x'^{\alpha}}{\partial x^{\beta}} \dots \right) \left(\frac{\partial x^{\mu}}{\partial x'^{\nu}} \dots \right) T^{\alpha\beta\gamma\dots}_{\mu\nu\dots}$$

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

↔
inverses

$x'(x)$ well-behaved provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

similarly, tensor of valence $\begin{bmatrix} m \\ n \end{bmatrix}$ # up
down

$$T^{\overbrace{\alpha\beta\gamma\dots}^m}_{\underbrace{\mu\nu\dots}_n}(x') = \left(\frac{\partial x'^{\alpha}}{\partial x^{\beta}} \frac{\partial x'^{\beta}}{\partial x^{\gamma}} \right) \left(\frac{\partial x^{\gamma}}{\partial x'^{\mu}} \dots \right) T^{\rho\sigma\dots}_{\mu\dots}$$

notice $\frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} = \frac{\partial x^{\beta}}{\partial x^{\alpha}} = \delta_{\alpha}^{\beta}$

$\swarrow \quad \searrow$
 inverses

$x'(x)$ well-behaved provided $\frac{\partial x'^{\alpha}}{\partial x^{\beta}}$ invertible.

similarly, tensor of valence $\begin{bmatrix} m \\ n \end{bmatrix}$ $\begin{matrix} \nwarrow \text{up} \\ \swarrow \text{down} \end{matrix}$

$$T^{\overbrace{\alpha\beta\gamma\dots}^m}_{\underbrace{\mu\nu\dots}_n}(x') = \underbrace{\left(\frac{\partial x'^{\alpha}}{\partial x^{\rho}} \frac{\partial x'^{\beta}}{\partial x^{\sigma}} \dots \right)}_{m \text{ factors}} \underbrace{\left(\frac{\partial x^{\mu}}{\partial x'^{\nu}} \dots \right)}_{n \text{ factors}} T^{\rho\sigma\dots}_{\mu\dots}$$

rank of a tensor is $m+n$

Tensor algebra:

• S are tensors of same valence,

up
down
 $\left. \begin{matrix} \dots \\ \dots \\ \dots \end{matrix} \right) T^{\rho\sigma} \dots$
n factors

rank of a tensor is $m+n$

Tensor algebra:

if T, S are tensors of same valence,
 $T+S$ is also.

up

down

$\left. \begin{matrix} \dots \\ \mu_1 \end{matrix} \right) T^{\rho_1 \dots} \mu_1 \dots (x)$
n factors

rank of a tensor is $m+n$

Tensor algebra:

1) if T, S are tensors of same valence,
 $T+S$ is also.

2) outer multiplication

$$T^{\alpha\beta\dots\mu}_m S_{\mu\nu\dots\rho}_n$$

$$T^{\alpha\beta\dots\mu}_m(x)$$

up
 # down
 ...
 μ
 $\underbrace{\hspace{1cm}}_{n \text{ factors}}$

rank of a tensor is $m+n$

Tensor algebra:

1) if T, S are tensors of same valence,
 $T+S$ is also.

2) outer multiplication

$$T^{\overbrace{\alpha\beta\ldots}^m} S_{\underbrace{\mu\nu\ldots}_n} \rightarrow \begin{matrix} \text{valence} \\ \left[\begin{matrix} m \\ n \end{matrix} \right] \end{matrix}$$

3) contractions

$$T^{\overbrace{\alpha\beta\delta\ldots}^m} S_{\underbrace{\alpha\beta\delta\ldots}_n} \rightarrow \begin{bmatrix} m-2 \\ n-2 \end{bmatrix}$$

$$T^{\rho\delta\ldots}_{\mu\ldots}(x)$$

up
down
n factors

rank of a tensor is $m+n$

Tensor algebra:

1) if T, S are tensors of same valence,
 $T+S$ is also.

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3) contractions

$$T^{\overbrace{\alpha\beta\dots}^m} S_{\underbrace{\alpha\beta\dots}_n} \rightarrow \begin{bmatrix} m-2 \\ n-2 \end{bmatrix}$$

If a tensor vanishes in one coordinate system, then it vanishes in any coordinate system.

rank of a tensor is $m+n$

Tensor algebra:

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