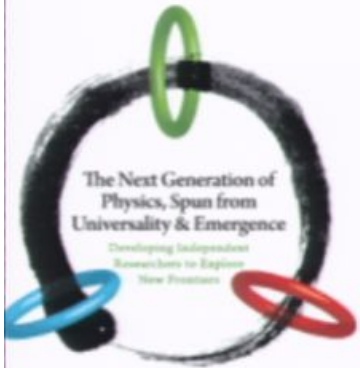


Title: Gravitational self-force effect on the periapsis advance in Schwarzschild spacetime

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Abstract: TBA



Gravitational self-force effect on the periapsis advance in Schwarzschild spacetime

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in collaboration with
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Outline of this talk

- Self-force calculation code
- Periapsis advance (test particle case)
- Conservative SF correction in periapsis advance
- Dissipative effect on periapsis advance
- Summary & Future work

Self-force calculation code

[Barack and NS, 2007, 2010]

Orbits

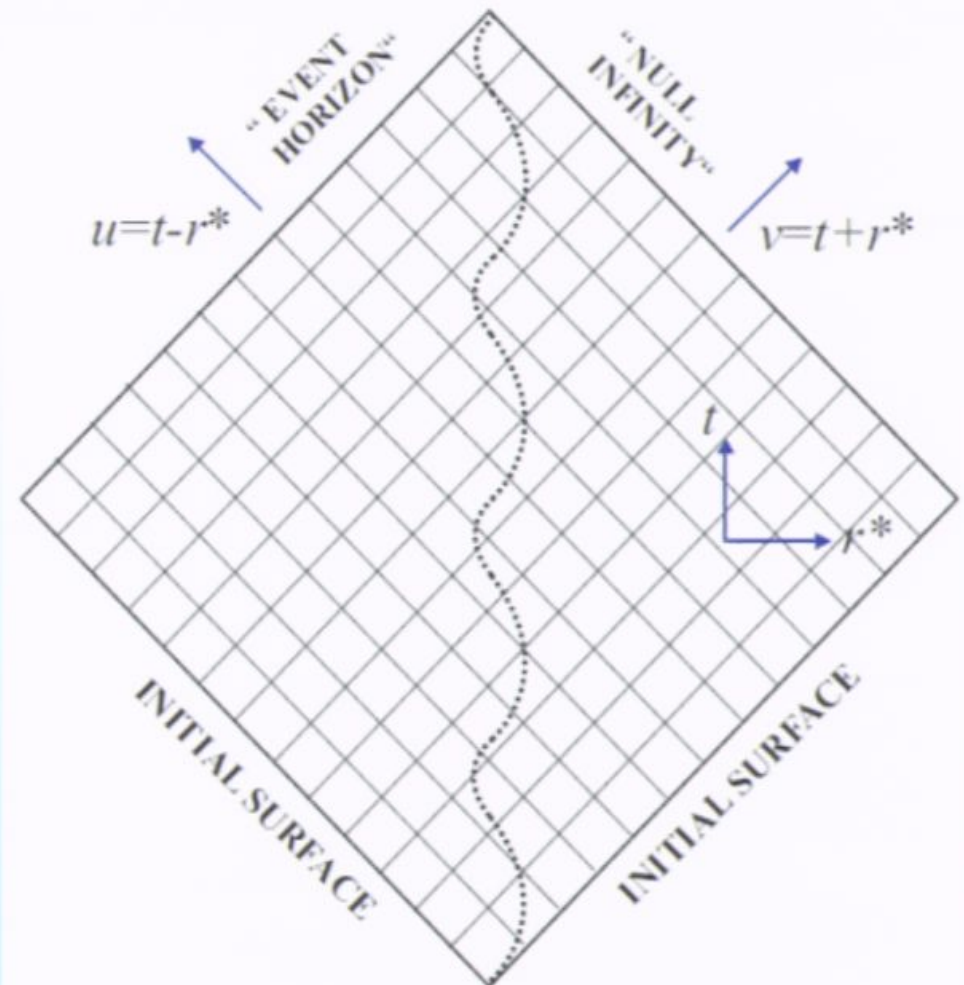
- Eccentric orbits

Metric perturbation

- Lorenz gauge
- (1+1)-dim. time domain
- Double null coordinates
- Uniform gridspace
- 4th order scheme

SF calculation

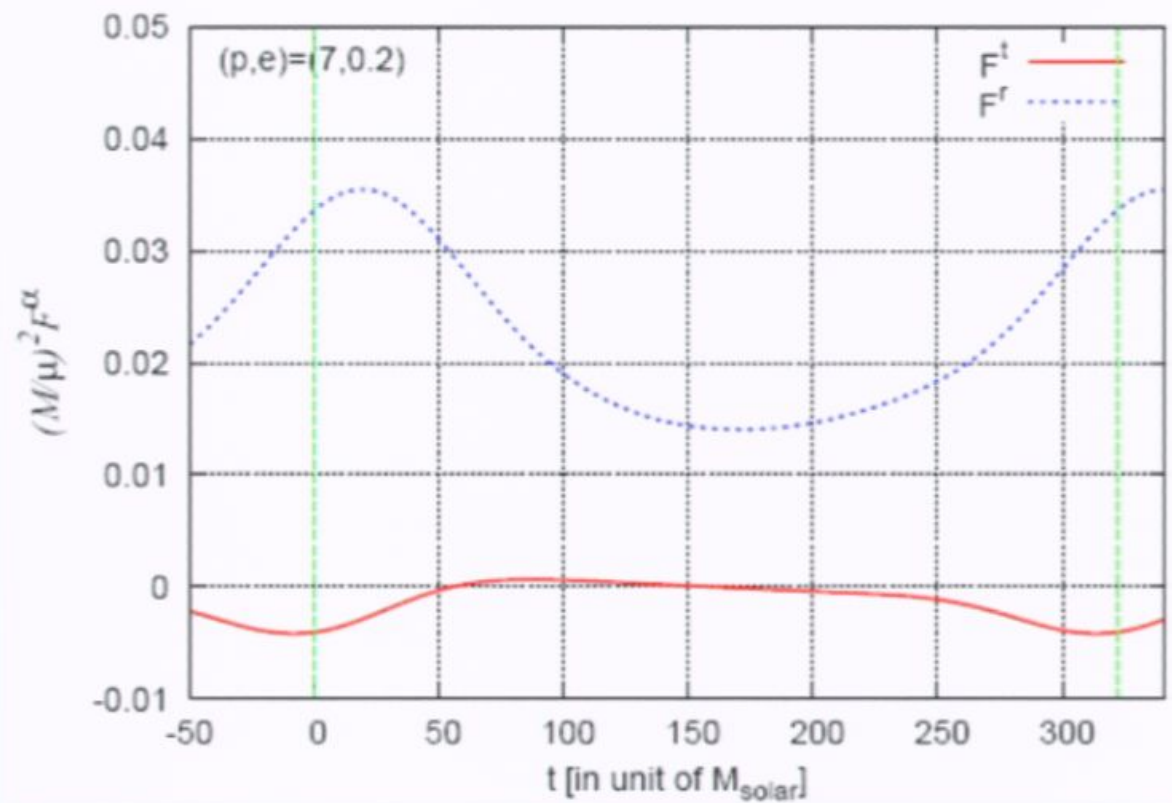
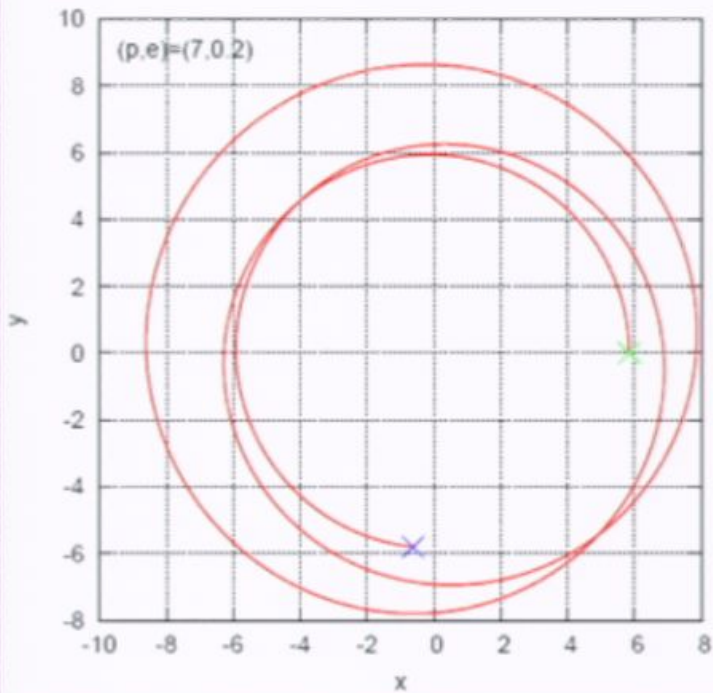
- Mode-sum scheme



$$F^\alpha(\tau) = \sum \left[F_{\text{full}}^{\alpha,l}(x_p) - A^\alpha L - B^\alpha \right]$$

Self-force calculation code

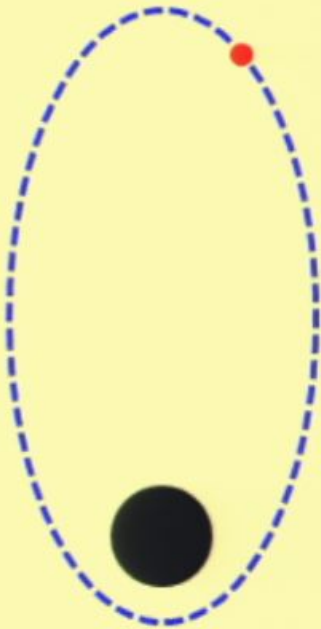
example: $(p,e)=(7,0.2)$



Main topic - SF correction in periaapsis advance -

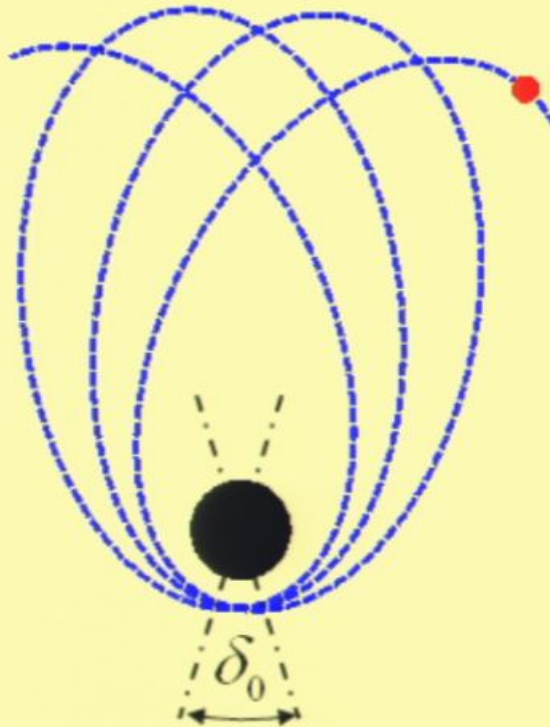
Newtonian

$$\frac{\Omega_\varphi}{\Omega_r} = 1$$



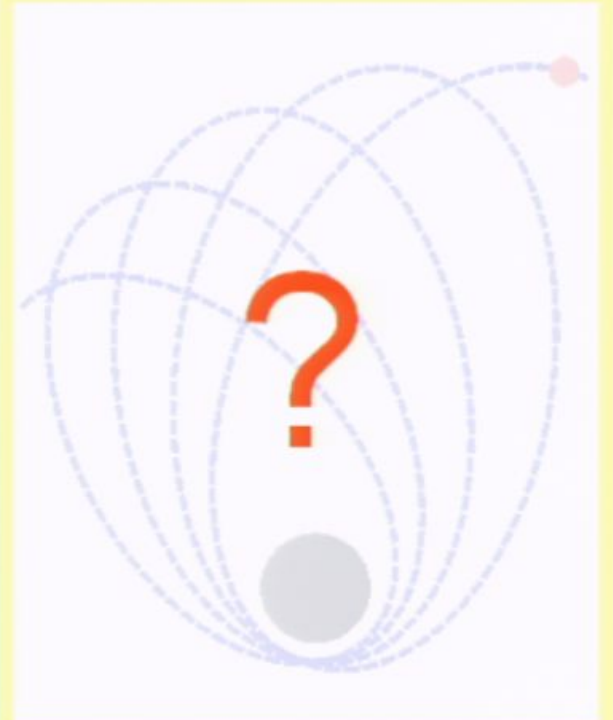
GR (geodesic)

$$\frac{\Omega_\varphi}{\Omega_r} = 1 + \frac{\delta_0}{2\pi}$$



GR (SF included)

$$\frac{\Omega_\varphi}{\Omega_r} = 1 + \frac{\delta_0 + \delta_{SF}}{2\pi}$$



Geodesic equations for eccentric orbit

Radial phase parameter χ

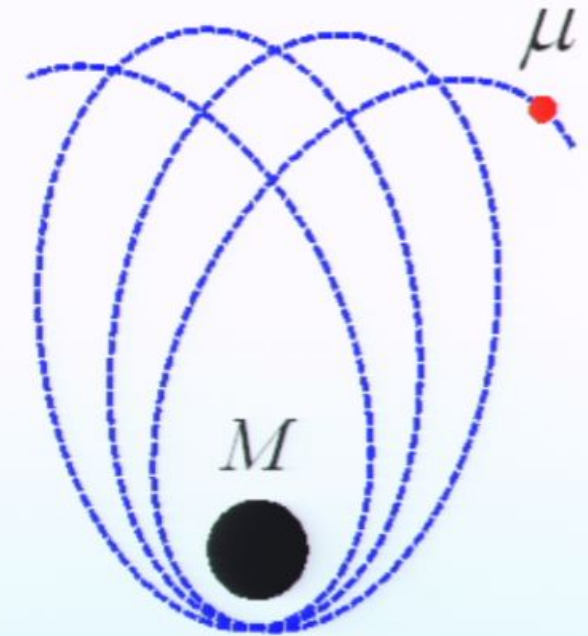
$$r_p(\chi) = \frac{pM}{1 + e \cos \chi} \quad (\chi=0 \text{ at periapsis})$$

$0 \leq \chi \leq 2\pi$ corresponds to one radial period.

Equation of motion

$$\frac{dt_p}{d\chi}(r_p; E_0, L_0) = \frac{E_0}{f_p} \left[E_0^2 - f_p \left(1 + \frac{L_0^2}{r_p^2} \right) \right]^{-1/2} \left(\frac{dr_p}{d\chi} \right)$$

$$\frac{d\varphi_p}{d\chi}(r_p; E_0, L_0) = \frac{L_0}{r_p^2} \left[E_0^2 - f_p \left(1 + \frac{L_0^2}{r_p^2} \right) \right]^{-1/2} \left(\frac{dr_p}{d\chi} \right)$$



- μ : Mass of particle
- M : Mass of central BH
- p : Semi-latus rectum
- e : Eccentricity

where E_0 and L_0 are the specific energy and angular momentum:

$$r_p^2 = \frac{(p-2-2e)(p-2+2e)}{r_p^2} = \frac{p^2 M^2}{r_p^2}$$

Periapsis advance (test particle case)

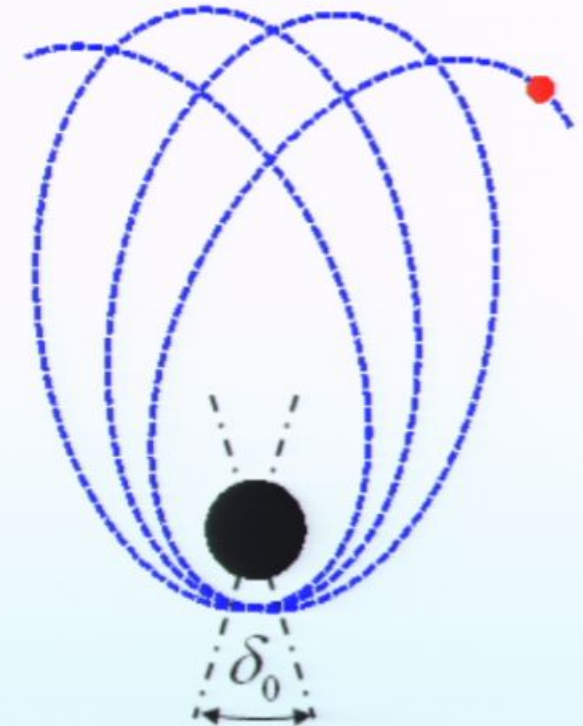
Periapsis advance

$$T_r \equiv \int_0^{2\pi} \left(\frac{dt_p}{d\chi} \right) (\chi; E_0, L_0) d\chi$$

$$\Delta\varphi \equiv \int_0^{2\pi} \left(\frac{d\varphi_p}{d\chi} \right) (\chi; E_0, L_0) d\chi$$

The periapsis advance can be obtained by:

$$\delta_0(p, e) = \Delta\varphi - 2\pi$$



Asymptotic form

For large p limit

$$\delta_0 = \frac{\pi}{c} \left[6 + \left(27 + \frac{3}{2} e^2 \right)^{1/2} \right] \frac{1}{p} + O(p^{-2})$$

For small e limit

$$\delta_0 = 2\pi \left[\frac{3}{c} \sqrt{\frac{p}{c} - 1} \right] + O(e^2)$$

Conservative part of self-force

[Hinderer-Flanagan '08]

Here we define the conservative piece of SF as that derived from the symmetric field (half retarded + half advanced):

$$F_{\text{cons}}^{\alpha} = \frac{1}{2} \left[F_{\text{ret}}^{\alpha} + F_{\text{adv}}^{\alpha} \right]$$

Using the symmetry of geodesics with respect to the transformation of $(\chi, t, r, \theta, \varphi) \rightarrow (-\chi, -t, r, \theta, -\varphi)$, we can obtain the relation between the retarded and advanced SF as

$$F_{\text{ret}}^r(-\chi) = F_{\text{adv}}^r(\chi), \quad F_{\text{ret}}^{t/\varphi}(-\chi) = -F_{\text{adv}}^{t/\varphi}(\chi)$$

$$F_{\text{cons}}^r(\chi) \equiv \frac{1}{2} \left[F^r(\chi) + F^r(-\chi) \right]$$

$$F_{\text{cons}}^{t/\varphi}(\chi) \equiv \frac{1}{2} \left[F^{t/\varphi}(\chi) - F^{t/\varphi}(-\chi) \right]$$

Conservative SF correction in periapsis advance

t - and φ - components of EOM including the conservative SF

$$\frac{d}{d\tau} E = -F_t^{\text{cons}}, \quad \frac{d}{d\tau} L = -F_\varphi^{\text{cons}}$$

Integrating these equations, we obtain:

$$E_0 \rightarrow E(\chi) = E_0 + \delta E(0) - \int_0^\chi F_t^{\text{cons}}(\chi') (d\tau/d\chi) d\chi'$$

$$L_0 \rightarrow L(\chi) = L_0 + \delta L(0) + \int_0^\chi F_\varphi^{\text{cons}}(\chi') (d\tau/d\chi) d\chi'$$

E, L are not constant along the orbit, but periodic.

$\delta E(0)$: correction at $\chi=0$

$\delta L(0)$ (in terms of SF)

Conservative SF correction in periaapsis advance

Correction of periaapsis advance

$$\Delta\varphi = \int_0^{2\pi} \left(\frac{d\varphi_p}{d\chi} \right) (\chi; E, L) d\chi$$

$$= \int_0^{2\pi} \left[\left(\frac{d\varphi_p}{d\chi} \right) (\chi; E_0, L_0) + \delta \left(\frac{d\varphi_p}{d\chi} \right) (\chi; E_0, L_0) \right] d\chi$$

$$\delta_0(p, e) + 2\pi$$

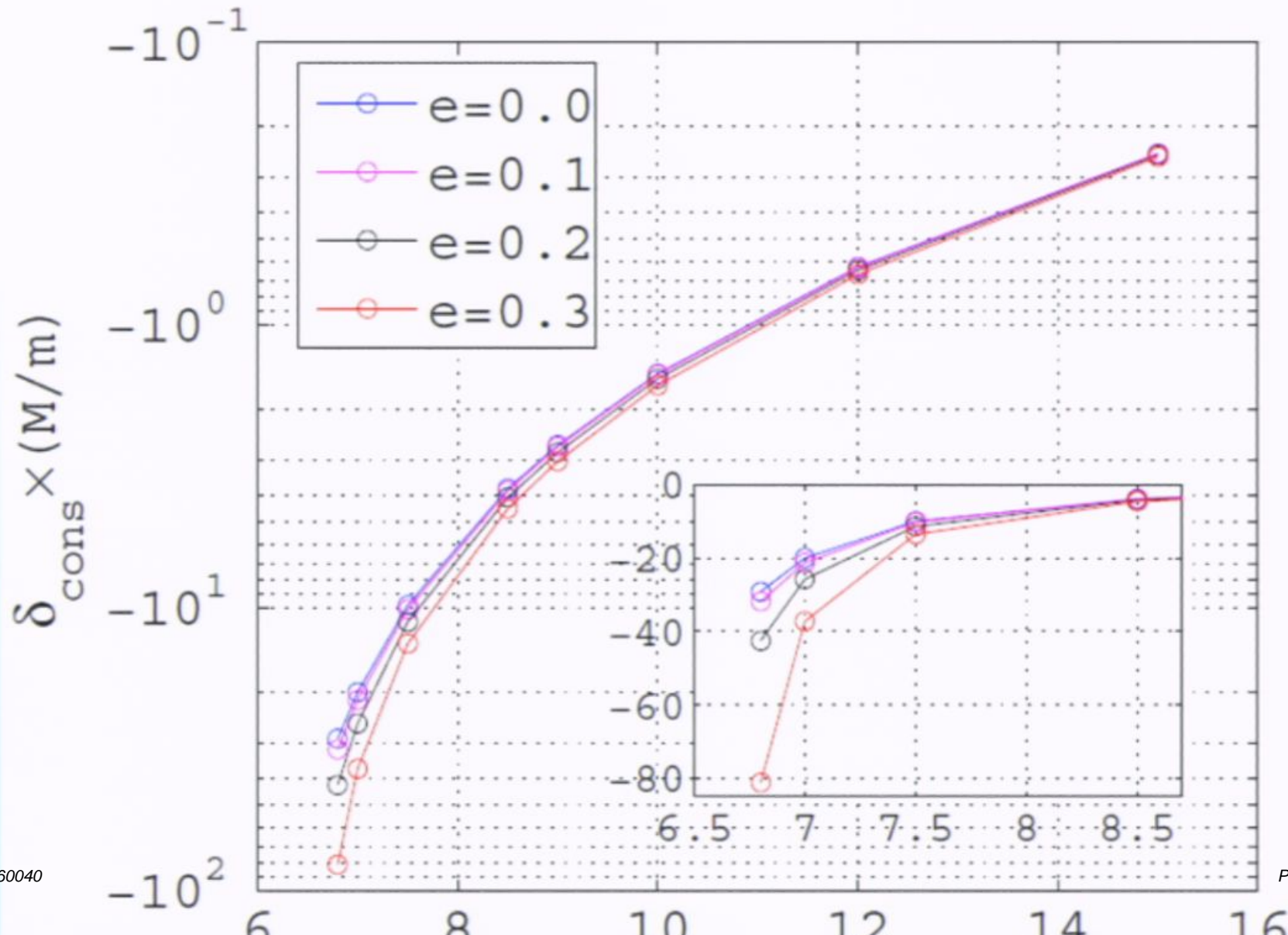
Leading (geodesic) term

$$\delta_{\text{cons}}(p, e)$$

Conservative correction term

$$\delta \left(\frac{d\varphi_p}{d\chi} \right) \equiv \frac{\partial}{\partial E} \left(\frac{d\varphi_p}{d\chi} \right) \Big|_0 \left[\delta E(0) - \int_0^x F_t^{\text{cons}}(\chi') (d\tau/d\chi) d\chi' \right] + \frac{\partial}{\partial L} \left(\frac{d\varphi_p}{d\chi} \right) \Big|_0 \left[\delta L(0) + \int_0^x F_\phi^{\text{cons}}(\chi') (d\tau/d\chi) d\chi' \right]$$

Conservative SF correction in periaapsis advance



Conservative SF correction in periaapsis advance

Correction of periaapsis advance

$$\Delta\varphi = \int_0^{2\pi} \left(\frac{d\varphi_p}{d\chi} \right) (\chi; E, L) d\chi$$

$$= \int_0^{2\pi} \left[\left(\frac{d\varphi_p}{d\chi} \right) (\chi; E_0, L_0) + \delta \left(\frac{d\varphi_p}{d\chi} \right) (\chi; E_0, L_0) \right] d\chi$$

$$\delta_0(p, e) + 2\pi$$

Leading (geodesic) term

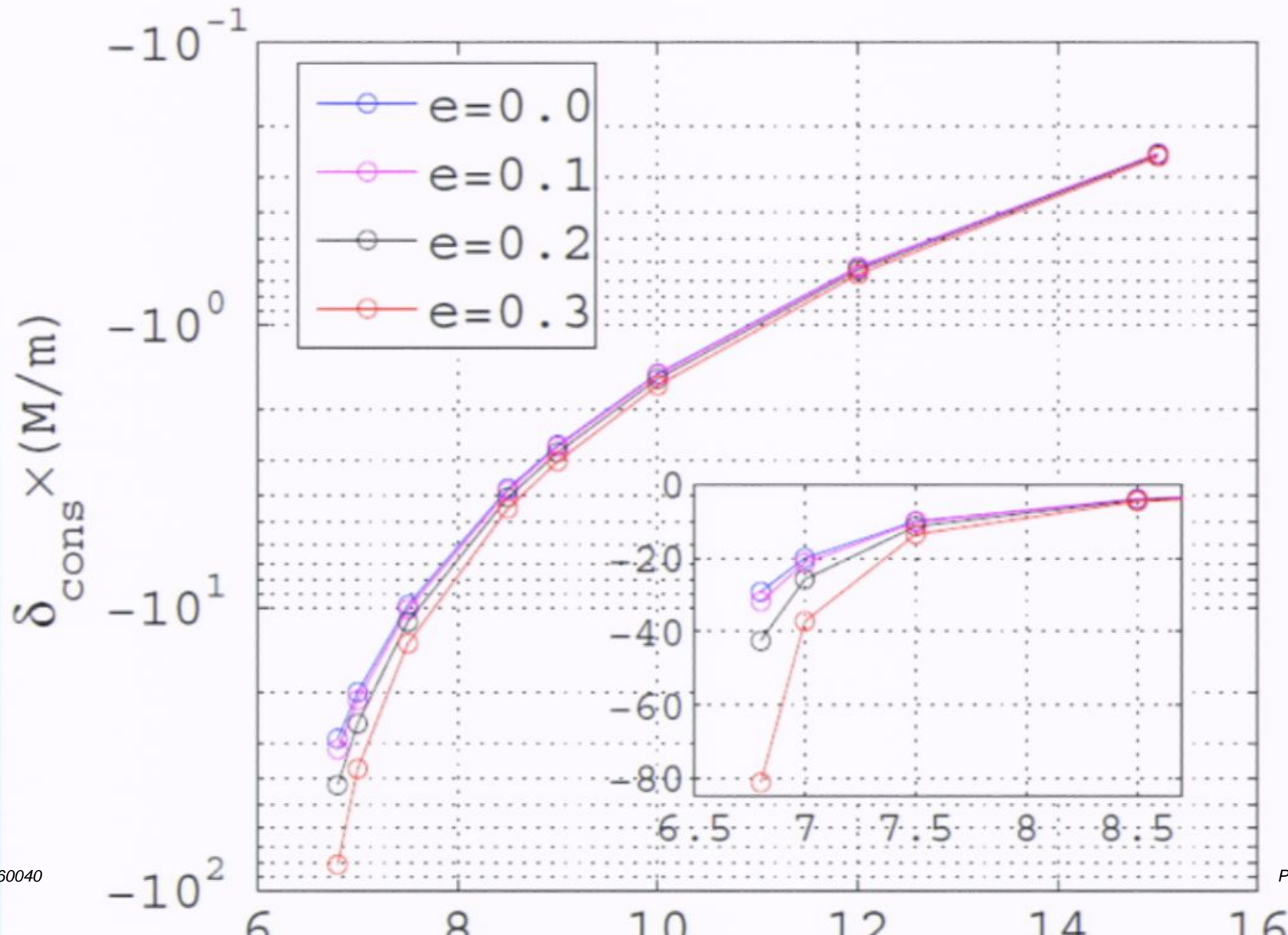
$$\delta_{\text{cons}}(p, e)$$

Conservative correction term

$$\delta \left(\frac{d\varphi_p}{d\chi} \right) \equiv \frac{\partial}{\partial E} \left(\frac{d\varphi_p}{d\chi} \right) \Big|_0 \left[\delta E(0) - \int_0^{\chi} F_t^{\text{cons}}(\chi') (d\tau/d\chi) d\chi' \right] + \frac{\partial}{\partial L} \left(\frac{d\varphi_p}{d\chi} \right) \Big|_0 \left[\delta L(0) + \int_0^{\chi} F_\phi^{\text{cons}}(\chi') (d\tau/d\chi) d\chi' \right]$$

The conservative correction can be described in terms double

Conservative SF correction in periaapsis advance



Comparison with post-Newtonian (circular limit)

Consider the ratio of two frequency at the circular limit:

$$\left. \begin{array}{l} \Omega_{\varphi} : \text{azimuthal frequency} \\ \Omega_r : \text{radial frequency} \end{array} \right\} W \equiv \left(\frac{\Omega_r}{\Omega_{\varphi}} \right)^2$$

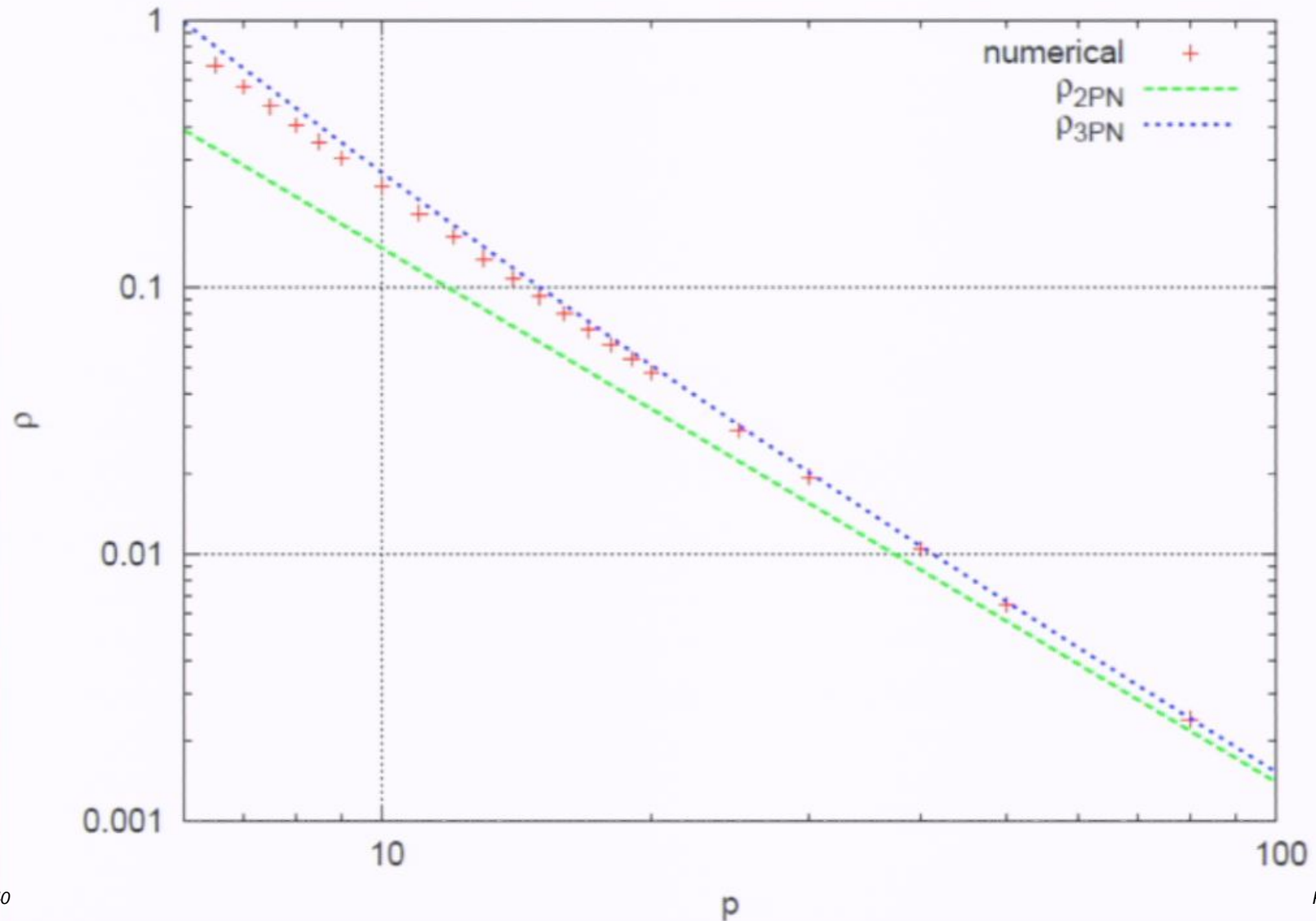
The SF correction in 'W' can be given by

$$\rho \equiv \Delta_{SF} W = -\frac{\delta_{SF}}{\pi} \left(1 - \frac{6}{p} \right)^{3/2}$$

Post-Newtonian formula of ρ for circular limit: [Damour '09]

$$\rho_{2PN} = \frac{14}{p^2} + \left(\frac{397}{p} - \frac{123}{p^2} \pi^2 \right) \frac{1}{p^3}$$

Comparison with post-Newtonian (circular limit)



Dissipative effect on periapsis advance

In reality, the parameters (p, e) changes in time due to the dissipation. Even in this case, we can define the periapsis advance:

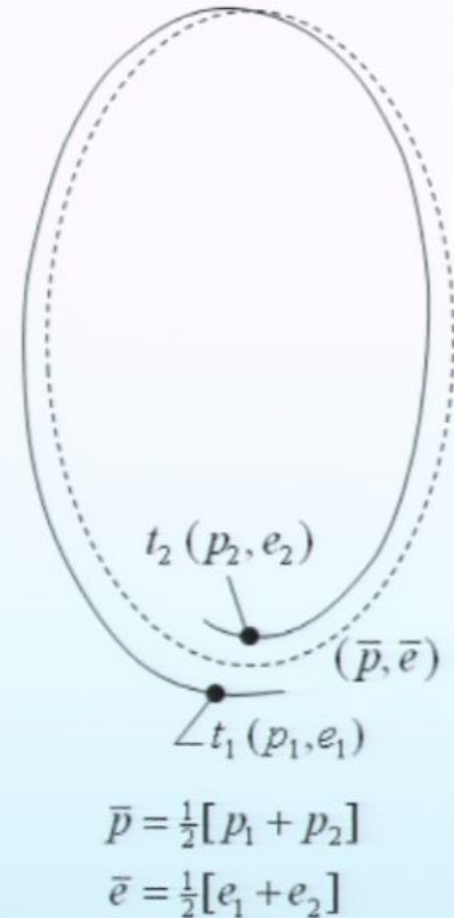
$$\Delta \varphi_{\text{true}}(t_1; t_2) = \int_{t_1}^{t_2} \dot{\varphi}(p(t), e(t)) dt$$

$$= \int_{t_1}^{t_2} \left[\dot{\varphi}(\bar{p}, \bar{e}) + \partial_{\bar{p}} \dot{\varphi}(\bar{p}, \bar{e})(p(t) - \bar{p}) + \partial_{\bar{e}} \dot{\varphi}(\bar{p}, \bar{e})(e(t) - \bar{e}) \right] dt$$

t_1, t_2 : consecutive radial turning points with $\dot{u}^r > 0$

[1st term] = $\Delta \varphi_{\text{cons}}(\bar{p}, \bar{e}) + O(\mu^2)$

[2nd, 3rd term] = $O(\mu^2)$



$$\Delta \varphi_{\text{true}}(t_1; t_2) = 2\pi + \delta_0(\bar{p}, \bar{e}) + \delta_{\text{cons}}(\bar{p}, \bar{e}) + O(\mu^2)$$

Summary & Future work

This work

Derive the SF correction in the periapsis advance

$$\Delta\varphi_{\text{true}}(t_1; t_2) = 2\pi + \delta_0(\bar{p}, \bar{e}) + \delta_{\text{cons}}(\bar{p}, \bar{e}) + O(\mu^2)$$

- Conservative effect is derived in terms of integrals of SF
- Dissipative effect is at the higher order $O(\mu^2)$.

Future work

Comparison to other methods (NR, PN, EOB, ...)

For this, we need some gauge invariant quantities.

- Circular limit [Barack, Damour, NS]
Described by only one orbital parameter (p).
We can use the angular frequency instead of p .
- Eccentric orbit [Barack and NS]
Described by two parameters (p, e).