Title: A new method to construct waveforms for Extreme-Mass-Ratio Inspirals

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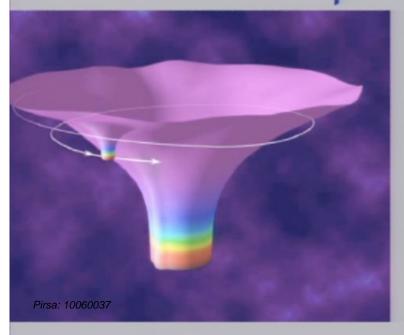
URL: http://pirsa.org/10060037

Abstract: TBA

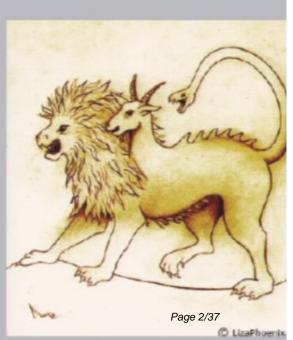
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4 New Method to Construct Waveforms for Extreme-Mass-Ratio Inspirals Carlos F. Sopuerta

Institute of Space Sciences (CSIC-IEEC) UAB Campus, Bellaterra, Barcelona



Work in collaboration with Nico Yunes (Princeton)



Outline

Motivation for a scheme to construct approximate EMRI waveforms.

A brief review of previous approximate methods

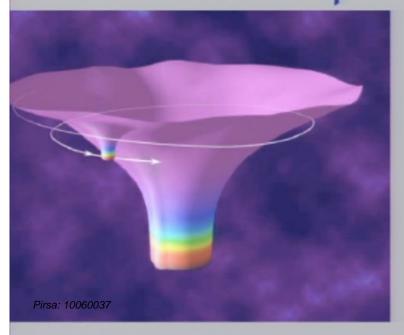
A new Method: The <u>Chimera</u> Scheme

Present Status and Future Prospects

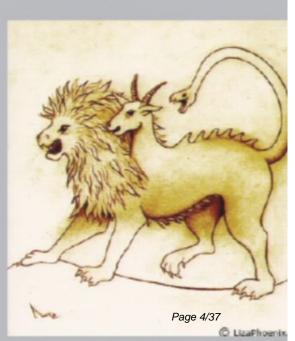
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Motivation

EMRIS are one of the main sources of GWS for LISA. The barameters of interest are:

$$\left. \begin{array}{lll} m & \sim & 1-30\,M_{\odot} \\ \\ M_{\bullet} & \sim & 10^5-10^7\,M_{\odot} \end{array} \right\} \implies \mu = \frac{m}{M_{\bullet}} \sim 10^{-7}-10^{-3}$$

The number of orbital cycles of an EMRI during the last year pefore plunge (inside the LISA band) is:

$$N_{\rm cycles} \sim 10^5 - 10^6$$

Then, it is clear that we cannot produce EMRI waveform emplate banks from the best methods presently available (SPISSE: 10060037) reukolsky + radiative approximation).

Motivation

Therefore, we need an approximate scheme that captures the essential characteristics of the EMRI dynamics (in particular coming from self-force calculations) in such a way that it can produce accurate waveforms in an efficient way.

Apart from their use for Data Analysis purposes, these approximate waveforms can also be used for parameterestimation purposes.

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Review of Approximate Methods

<u>Newtonian</u> [Peters & Mathews (1963)]: Newtonian trajectories + Quadrupolar waveforms.

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Teukolsky [Drasco & Hughes (2004)]: Kerr geodesics + Teukolsky fluxes + Teukolsky waveforms.

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Review of Approximate Methods

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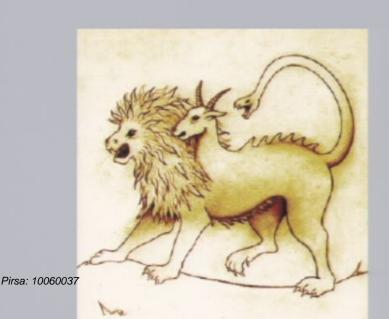
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Teukolsky [Drasco & Hughes (2004)]: Kerr geodesics + Teukolsky fluxes + Teukolsky waveforms.

<u>Numerical Kludge</u> [Babak et al (2007); Gair & Glampedakis (2006)]: Kerr geodesics + pN fluxes + Multipolar waveforms

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in Greek mythology, the <u>Chimera</u> was a monstrous ire-breathing creature of Lycia (in Asia Minor), omposed of the parts of multiple animals: upon the lody of a <u>lioness</u> with a tail that terminated in a <u>make</u>'s head, the head of a <u>goat</u> arose on her back at the enter of her spine.





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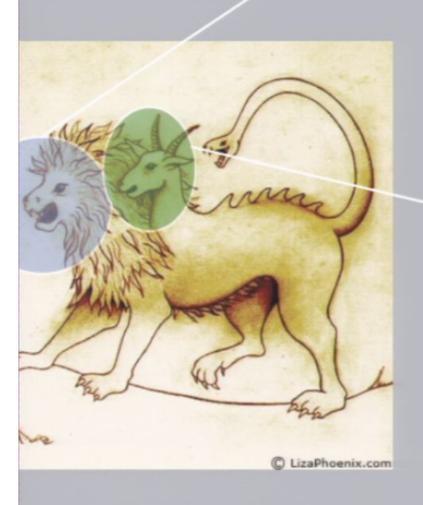
f The term <u>chimera</u> has also come to mean, more generally, an impossible or foolish fantasy, hard to pelieve... [Wikipedia]

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Kerr geodesics + Harmonic Coordinates



Radiative Self-Force from post-Minkowskian approximation +
Asymptotic matched expansions.

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Kerr geodesics + Harmonic Coordinates

Radiative Self-Force from post-Minkowskian approximation +
Asymptotic matched expansions.

Multipolar Expansion of the Gravitational

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Mayoforma

KRadiative Self-From from post-Minkowskian pproximation + Asymptotic matched expansions Blanchet & Damour (1984); Iyer & Will (1995); slanchet (1997)]

KBasic Idea: To determine the gravitational field oth at the "near zone" and the "exterior zone" and to natch the two solutions at the overlapping region. These solutions for the gravitational fields are xpanded in G (post-Minkowskian expansion) and not here.

KRAdíatíve Self-From from post-Minkowskian pproximation + Asymptotic matched expansions Blanchet & Damour (1984); Iyer & Will (1995); lanchet (1997)]

(using the half retarded minus half advanced plution, the matching of the solutions provides an opression for a dissipative (radiative) self-force.

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K At IPN order, the radiative self-force is letermined from a scalar and a vector radiationeaction potentials:

$$V_{RR}(t, \mathbf{x}) = -\frac{1}{5} x^{ij} M_{ij}^{(5)}(t) + \frac{1}{189} x^{ijk} M_{ijk}^{(7)}(t) - \frac{1}{70} \mathbf{x}^2 x^{ij} M_{ij}^{(7)}(t),$$

$$V_{\text{Sai-PORT}}^{i}(t,\mathbf{x}) = \frac{1}{21}\hat{x}^{ijk}M_{jk}^{(6)}(t) - \frac{4}{45}\epsilon_{ijk}x^{jl}S_{kl}^{(5)}(t),$$

K Comment: All this assumes a harmonic gauge harmonic coordinates for the multipolar expansion).

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K The radiative self-force has two pieces:

i) Gradients of the potentials [this implies we need up to eight-order time derivatives of the trajectory].
ii) Interaction terms of the radiative potentials with he near-zone potentials (Kerr gravitational field).

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KKerr geodesics +Harmonic Coordinates[Ding 1983); Abe, Ichinose & Nakanishi (1987)] [See also zuiz (1986)]:

$$t_{\rm H} = t$$
,

$$x_{\rm H} = \sqrt{(r - M_{\bullet})^2 + a^2 \sin \theta \cos[\phi - \Phi(r)]},$$

$$y_{\rm H} = \sqrt{(r - M_{\bullet})^2 + a^2 \sin \theta \sin[\phi - \Phi(r)]},$$

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KKerr geodesics +Harmonic Coordinates: An Iternative to this is to use Asymptotically Cartesian Ind Mass Centered coordinates (ACMC) [Thorne 1980)]

KThis may be a convenient path to deal with the case n which the massive object is not described by the cerr metric.

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KMultípolar Expansion of the Gravitational Vaveforms [Thorne (1980)]

$$h_{ij}^{\rm TT} = \frac{4}{r} \sum_{\ell=2}^{\infty} \left[\frac{1}{\ell!} \mathfrak{M}_{ijI_{\ell-2}}^{(\ell)}(t-r) N_{I_{\ell-2}} \right]$$

$$+ \frac{2\ell}{(\ell+1)!} \varepsilon_{kl(i} \mathfrak{J}_{j)kI_{\ell-1}}^{(\ell)}(t-r) n_l N_{I_{\ell-2}} \bigg]^{\mathrm{TT}}$$

KHere, we also use harmonic coordinates (improving ver the procedure of the Numerical Kludge Mark (10000037 rms).

KGeneral Procedure, Step by Step: For initial data xo,yo,zo) & (p,e,i)/(E,Lz,C)

- I) Evolution of the Geodesic Equation from a time tep Dt (in harmonic coordinates).
- 11) Evaluation of the dissipative self-force.
- III) Estimation of the change in the constants of notion (E,Lz,C)

$$\dot{E} = -\Gamma^{-1} u^{\mu} \nabla_{\mu} \left(\xi_{\mu}^{(t)} u^{\mu} \right) = -\Gamma^{-1} \xi_{\mu}^{(t)} a_{RR}^{\mu} ,$$

$$\dot{L}_z \equiv \Gamma^{-1} u^{\mu} \nabla_{\mu} \left(\xi_{\mu}^{(\phi)} u^{\mu} \right) = \Gamma^{-1} \xi_{\mu}^{(\phi)} a_{RR}^{\mu} ,$$

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$$\dot{\Omega} = \Gamma^{-1} \cdot \mu \nabla \left(c \quad \alpha \cdot \beta \right) \quad \Omega \Gamma^{-1} c \quad \alpha \cdot \rho \cdot \mu$$

KGeneral Procedure, Step by Step: For initial data xo, yo, zo) & (p,e,i)/(E,Lz,C)

IV) Evaluation of the waveforms (can be done at ay ime).

V) Go back to (1).

K Note: Since our self-force is not averaged over a umber of orbital cycles, this procedure can be pplied, in principle, for an arbitrary time step.

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K We have constructed numerical codes that mplement of the steps described in this talk.

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K We have constructed numerical codes that mplement of the steps described in this talk.

KThe only difficulty is to evaluate the numerical ime derivatives that enter in the self-force: up to ight-order derivatives (six-order if we use analytical xpressions for the second-order derivatives of the uultipole moments). This requires high-order lifferentiation rules and the use of Richardson extrapolation (in part to check the accuracy).

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K We are currently generating waveforms with the loal of comparing with other approximate approaches Numerical Kludge, Teukolsky, EOB).

K We want to study the evolution of the "constants f motion" and compare with several approximations n the literature.

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K For the future, we would like to investigate the ossibility of using the Chimera waveforms for arameter-estimation analysis.

K It can clearly be improved in several ways: ntroducing conservative pieces of the self-force, etc.

t It would be interesting to see how this Scheme can extended to other geometries (non-Kerr) and even other gravitational theories.

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