

Title: A fast frequency-domain calculator for the gravitational self force

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Abstract: TBA

# A Fast Frequency-Domain Calculator for Gravitational Self-Force in Schwarzschild Spacetime

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June 20, 2010

# Ultimate Goal: Gravitational Wave Detection

- To see GWs (from EMRIs) in the LISA output, we need **very accurate GW** templates.
- To get these templates, we need a very **good modeling of orbits** in strong gravity.
- In short, we need very **precise, fast Self-Force** computations.
- F-domain SF computations: **fast(er)**  $\Rightarrow$  cover parameter space  $(p, e)$  in Schwarzschild.
- **Schwarzschild**: Solved (T-domain: Barack, Sago (full GSF); F-domain: Poisson, Detweiler). **Next**: Cover parameter space (F- vs. T-domains). **Kerr**: Eccentric scalar SF (Warburton); GSF: to come. Parameter space is huge for generic orbits.

# Background

- A point particle of mass  $\mu$  in Schwarzschild (Sch.) background with mass  $M$ .
- Particle induces a **perturbation** on the background spacetime. Geodesics are perturbed.
- We solve perturbed Einstein's equation in non-flat background.

$$G_{\mu\nu}(g_{\mu\nu} + h_{\mu\nu}) = 8\pi T_{\mu\nu}$$

- Retain only **linear order**  $\mathcal{O}(h)$  perturbation. Recall  $G(g) = 0$  (vacuum).

Further simplifications:

- Use  $\bar{h}_{\mu\nu} = h_{\mu\nu} - \frac{1}{2}g_{\mu\nu}h$

- Lorenz gauge where  $\nabla_\mu \bar{h}^{\mu\nu} = 0$  and with  $\nabla_\mu T^{\mu\nu} = 0$

$$\square \bar{h}_{\mu\nu} + 2R^{\alpha\beta} \bar{h}_{\alpha\beta} = -16\pi T_{\mu\nu}.$$

10 2nd order coupled wave-like equations sourced by  $T_{\mu\nu} \propto \delta(r - r_p(t))$ .

- The GSF is constructed from the gradient of  $\bar{h}_{\mu\nu}$  projected in the direction orthogonal to 4-velocity  $u^\mu$ .

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# Separation of Variables: Tensor Spherical Harmonics

- Separate the angular part. Use **tensor spherical harmonics** 10-dim. basis for any rank-2, symmetric 4-dim. tensor. In Kerr, spheroidal harmonics.

$$\bar{h}_{\mu\nu}^{\ell m} = \frac{\mu}{r} \sum_{\ell, m} \sum_{i=1}^{10} a^{(i)\ell} \bar{h}^{(i)\ell m}(t, r) Y_{\mu\nu}^{\ell m}(\theta, \phi; r).$$

- Angular variables decouple. Einstein equation becomes (at each  $\ell, m$ )

$$\square_{sc} \bar{h}^{(i)} + \mathcal{M}_{(j)}^{(i)} h^{(j)} = \mathcal{S}^{(i)}. \quad (1)$$

$\square_{sc}$  is the usual scalar field wave operator:

$$\square_{sc} = \frac{1}{4} \left[ \partial_t^2 - \partial_{r_*}^2 + f(r) \left( \frac{2M}{r^3} + \frac{\ell(\ell+1)}{r^2} \right) \right]$$

Pirsa: 10060035 where  $f(r) = 1 - \frac{2M}{r}$ .

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# Field Equations in Frequency Domain

- Bound orbits, up to  $e = e_{max}(?)$ , frequency domain **faster**. Determine  $e_{max}$ .
- For now: **Circular orbits** i.e.  $r_p(t) = r_0$

$$\bar{h}_{\ell m}^{(i)}(t, r) = \sum_m R_{\ell m}^{(i)}(r) \exp^{-i\omega_m t} \quad \omega_m = m\Omega_\phi$$

$$\tilde{\square}_{sc} R^{(i)}(r) + \mathcal{M}_{(j)}^{(i)} R^{(j)}(r) = \mathcal{S}^{(i)}(r).$$

- **10 2<sup>nd</sup> order coupled ODEs**.  $\mathcal{M}_{(j)}^{(i)} R^{(j)}$  contain the couplings.

$$\mathcal{M}_{(j)}^{(1)} R^{(j)} = \frac{1}{2} f^2 f' R_{,r}^{(3)} + \frac{f}{2r^2} (1 - 4M/r) \left( R^{(1)} - R^{(5)} - fR^{(3)} \right) - \frac{f^2}{2r^2} (1 - 6M/r) R^{(6)}$$

- 4 gauge equations from  $\nabla_\mu \bar{h}^{\mu\nu} = 0 \Rightarrow$  **4 1<sup>st</sup> order coupled ODEs**.

- **Parity**: Tensor spherical harmonics decouple:

10 ODEs  $\Rightarrow$  7 even  $\oplus$  3 odd, 4 gauge  $\Rightarrow$  3 even  $\oplus$  1 odd.

Even modes:  $l + m = \text{even}$

$(i) = 1 \dots 7$  eqtns. for  $R^{(1)} \dots R^{(7)}$  w/  $R^{(8)} = R^{(9)} = R^{(10)} = 0$

Odd modes:  $l + m = \text{odd}$

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$$10 \text{ ODEs} \Rightarrow 7 \text{ even} \oplus 3 \text{ odd}, \quad 4 \text{ gauge} \Rightarrow 3 \text{ even} \oplus 1 \text{ odd}.$$

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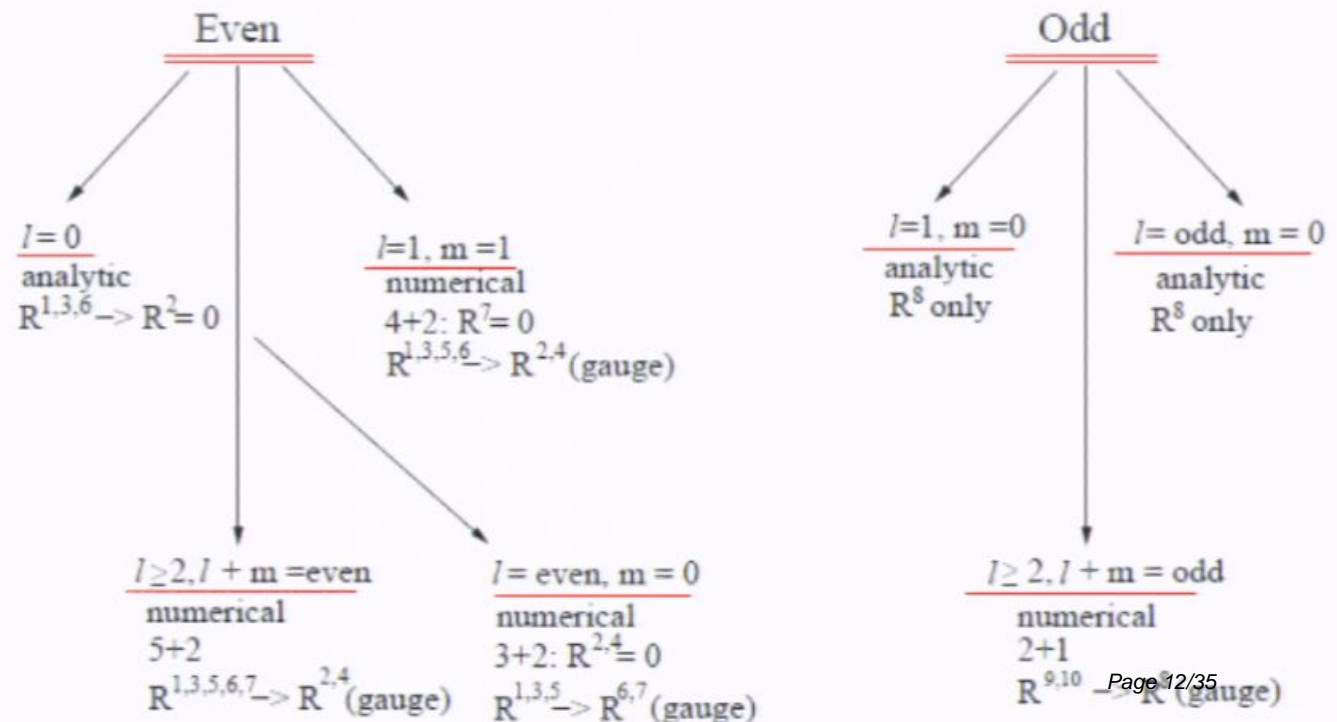
# Hierarchical Solving Scheme

- Not everything couples together at each  $\ell, m$ .  
Solve certain sets of ODEs first  $\Rightarrow$  Use solutions to obtain other fields.  
Can also use the gauge equations.

	Even ( $\ell + m = 2N$ )	Odd ( $\ell + m = 2N + 1$ )
$\ell = 0$	$(i) = 1, 3, 6 \rightarrow 2$	0
$\ell = 1$	$m = 1 : (i) = 1, 3, 5, 6 \rightarrow 2, 4$	$m = 0 : (i) = 8$ only
$\ell \geq 2$	$(i) = 1, 3, 5, 6, 7 \rightarrow 2, 4$	$(i) = 9, 10 \rightarrow 8$

## ODE Solving Tree:

Analytic solutions:  
monopole, odd  
dipole.



# The Crux: Boundary Conditions (BC)

- In **Time Domain**: No BC. Inaccurate IC cause spurious radiation.
- In **Freq. Domain**, we solve 1-dim. ODE so must specify correct BCs.
- **Ansatz**  $m \neq 0$  **dynamical modes**: in/outgoing waves at EH/infinity.

$$R_{in}^{(i)} = e^{-i\omega r_{in}^*} \sum_k b_k^{(i)} (r_{in} - 2M)^k, \quad R_{out}^{(i)} = e^{i\omega r_{out}^*} \sum_k \frac{a_k^{(i)}}{r_{out}^k}$$

- BC for  $m = 0$  **static modes** are determined by regularity.  
Series is infinite but truncate at some  $k_{max}$ .  $k = k_{start}$  are free.
- Substitute ansatz into each ODE. Obtain **recursion relations** for  $a_k^{(i)}, b_k^{(i)}$ .  
E.g.

$$4M^2 k(k - 4Mi\omega) b_k^3 = G_{k-1}^3 b_{k-1}^3 + 4M(b_{k-1}^1 - b_{k-1}^5 + b_{k-1}^6) \\ + H_{k-2}^3 - 2(b_{k-2}^1 - b_{k-2}^5 - b_{k-2}^6) - F_{k-3}^1 b_{k-3}^3$$

$$\text{where } G_k^3 = 2M [k(2k - 1) - l(l + 1) + 1 - 12Mi\omega k],$$

$$H_k^3 = k(k - 1) - l(l + 1) - 2 - 12Mi\omega k, \quad F_k^1 = 2i\omega k.$$

- 28 BCs in total.  $k_{max} = 5$  to 10 usually.

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# Matching the Inner and Outer Solutions

- 2 Junction Conditions at  $r = r_0$ :

(1)  $R_{in,out}^{(i)}$  are **continuous**, (2)  $\frac{dR_{in,out}^{(i)}}{dr}$  have a **jump discontinuity**  $\propto S^{(i)}$ .

- Single field: Undergrad. QM, like solving Schrodinger eqn. across  $\delta$ -potential.

- **Many coupled fields**: Follow Detweiler & Poisson 2003.

- $a_k^i, b_k^i$  that start the recursion relations for BC form a **vector space**.

**Example:** Even sector,  $\ell \geq 2, m > 1$ .

7 field and 3 gauge equations  $\Rightarrow$  **4 d.o.f** manifest in  $a_{k=0}^{1,5,6,7}, b_{k=0}^{1,5,6,7}$

**4-dim. vector space** for both inner (-) and outer (+) solutions.

Basis vectors:  $(1,0,0,0), (0,1,0,0), (0,0,1,0), (0,0,0,1)$

- Solve coupled ODEs with each basis vector as BC  $\rightarrow$   
4 times each for inner/outer solutions  $\rightarrow$  Solve ODEs 8 times.  
Actually 16 (for  $Re$  and  $Im$  parts).



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# Matching continued

- Output at each  $\ell, m$ : fields  $A_{in/out}^{1,2,3,4}, B_{in/out}^{1,2,3,4}, C_{in/out}^{1,2,3,4}, D_{in/out}^{1,2,3,4}$
- Similarly, for the derivatives:  $A_{in/out}^{\prime 1,2,3,4}, B_{in/out}^{\prime 1,2,3,4}, C_{in/out}^{\prime 1,2,3,4}, D_{in/out}^{\prime 1,2,3,4}$
- **e.g.**  $A_{out}^1$  is  $R_+^{(1)}$  obtained by using the vector  $(1, 0, 0, 0)$   
i.e.  $a_0^1 = 1, a_0^{5,6,7} = 0$ .
- The output fields and their derivatives form a  **$8 \times 8$  complex matrix**.
- To obtain the normalized, **inhomogeneous solutions**, we must solve for the complex amplitudes  $x_i$  in the following equation.

$$\begin{pmatrix} -A_{in}^1 & -A_{in}^2 & -A_{in}^3 & -A_{in}^4 & A_{out}^1 & A_{out}^2 & A_{out}^3 & A_{out}^4 \\ -B_{in}^1 & \dots & \dots & -B_{in}^4 & B_{out}^1 & \dots & \dots & B_{out}^4 \\ -C_{in}^1 & \dots & \dots & -C_{in}^4 & C_{out}^1 & \dots & \dots & C_{out}^4 \\ -D_{in}^1 & \dots & \dots & -D_{in}^4 & D_{out}^1 & \dots & \dots & D_{out}^4 \\ -A_{in}^{\prime 1} & -A_{in}^{\prime 2} & -A_{in}^{\prime 3} & -A_{in}^{\prime 4} & A_{out}^{\prime 1} & A_{out}^{\prime 2} & A_{out}^{\prime 3} & A_{out}^{\prime 4} \\ -B_{in}^{\prime 1} & \dots & \dots & -B_{in}^{\prime 4} & B_{out}^{\prime 1} & \dots & \dots & B_{out}^{\prime 4} \\ -C_{in}^{\prime 1} & \dots & \dots & -C_{in}^{\prime 4} & C_{out}^{\prime 1} & \dots & \dots & C_{out}^{\prime 4} \\ -D_{in}^{\prime 1} & \dots & \dots & -D_{in}^{\prime 4} & D_{out}^{\prime 1} & \dots & \dots & D_{out}^{\prime 4} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \\ x_5 \\ x_6 \\ x_7 \\ x_8 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ J_1 \\ 0 \\ J_6 \\ J_7 \end{pmatrix}$$

# Matching continued

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$$\begin{pmatrix} -A_{in}^1 & -A_{in}^2 & -A_{in}^3 & -A_{in}^4 & A_{out}^1 & A_{out}^2 & A_{out}^3 & A_{out}^4 \\ -B_{in}^1 & \dots & & -B_{in}^4 & B_{out}^1 & \dots & & B_{out}^4 \\ -C_{in}^1 & \dots & & -C_{in}^4 & C_{out}^1 & \dots & & C_{out}^4 \\ -D_{in}^1 & \dots & & -D_{in}^4 & D_{out}^1 & \dots & & D_{out}^4 \\ -A'_{in}{}^1 & -A'_{in}{}^2 & -A'_{in}{}^3 & -A'_{in}{}^4 & A'_{out}{}^1 & A'_{out}{}^2 & A'_{out}{}^3 & A'_{out}{}^4 \\ -B'_{in}{}^1 & \dots & & -B'_{in}{}^4 & B'_{out}{}^1 & \dots & & B'_{out}{}^4 \\ -C'_{in}{}^1 & \dots & & -C'_{in}{}^4 & C'_{out}{}^1 & \dots & & C'_{out}{}^4 \\ -D'_{in}{}^1 & \dots & & -D'_{in}{}^4 & D'_{out}{}^1 & \dots & & D'_{out}{}^4 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \\ x_5 \\ x_6 \\ x_7 \\ x_8 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ J_1 \\ 0 \\ J_6 \\ J_7 \end{pmatrix}$$

- **Inhomogeneous sol.:** e.g.  $R_{in}^{(1)} = x_1 A_{in}^1 + x_2 A_{in}^2 + x_3 A_{in}^3 + x_4 A_{in}^4$  etc.

# Numerical Details

- Code in **C**, using **GNU Scientific Library (GSL)** repository.
  - Rewrote 2nd order ODEs as 1st order ODEs.
  - Used Embedded Runge-Kutta Prince-Dormand (**rk8pd**) integrator
  - Empirically shown to be the fastest integrator (the coarsest).
  - Typical ODE solver error threshold:  $(10^{-10}, 10^{-10})$ .
- At each  $\ell, m$ , same set of ODEs is solved many times:
  - 4 times for inner/outer, real/imag solutions
  - Another  $N$  times to cover the  $N$ -dim. vector space
    - e.g. Even sector  $\ell \geq 2, m > 1$ : 4-d vector space  $\Rightarrow$  Solve 16 times!
    - Similarly in Odd sector  $\ell \geq 2$ : 2-d vector space  $\Rightarrow$  Solve 8 times
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 $(12 + 7 \times 12 + 56 \times 8 + 63 \times 16) = 1552$  times!

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# Numerics contd.

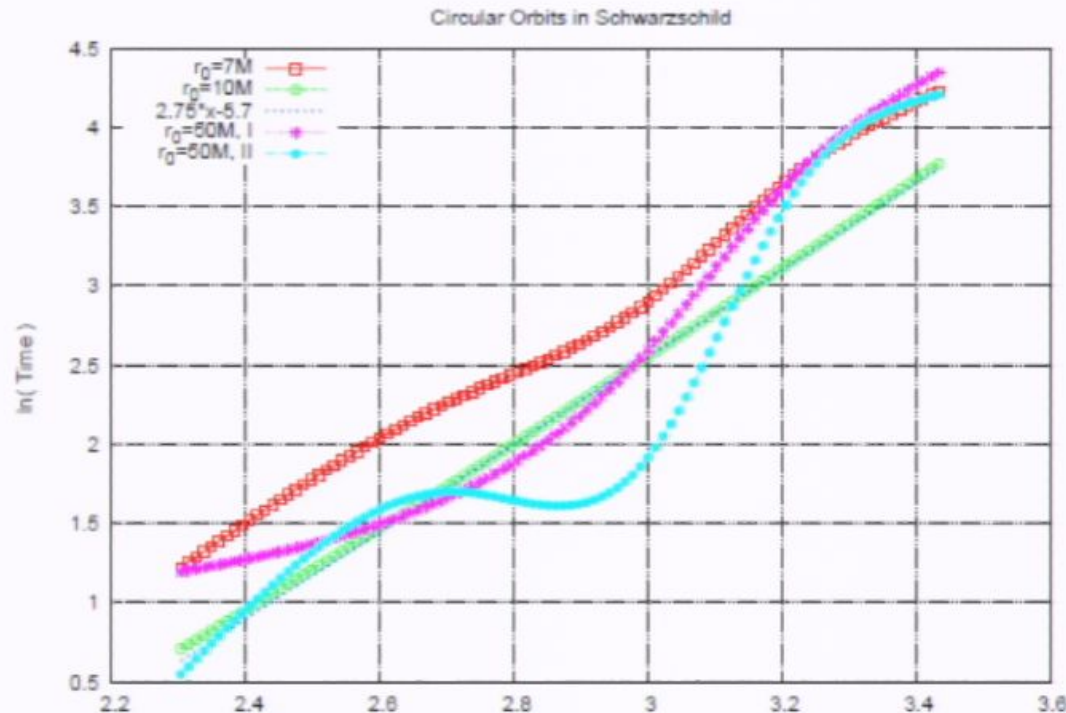
## Speed tests:

	$r_0 = 50M^{(I)}$	$r_0 = 50M^{(II)}$	$r_0 = 10M$	$r_0 = 7M$
$\ell$	time (m:s)	time (m:s)	time (m:s)	time (m:s)
10	3:18	1:44	2:02	3:32
15	5:22	5:29	5:48	9:45
20	13:10	6:39	12:43	17:58
25	40:20	36:31	23:43	40:40
31	1:17:31	1:07:21	43:30	1:08:43

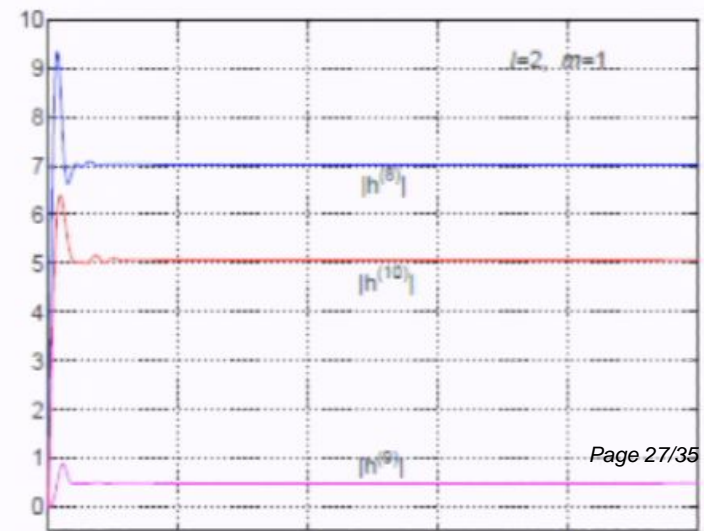
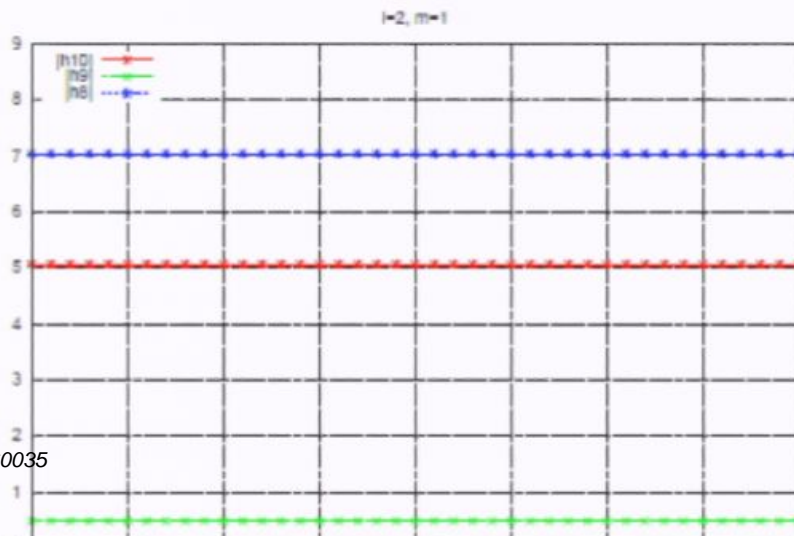
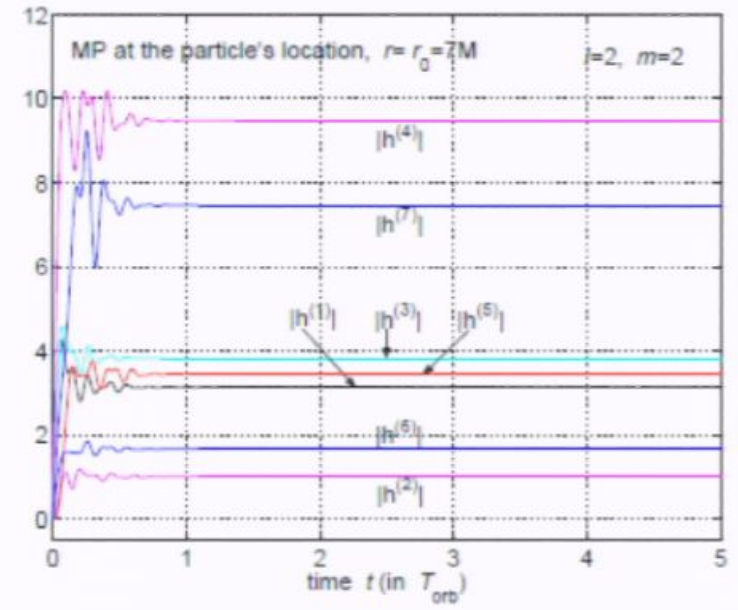
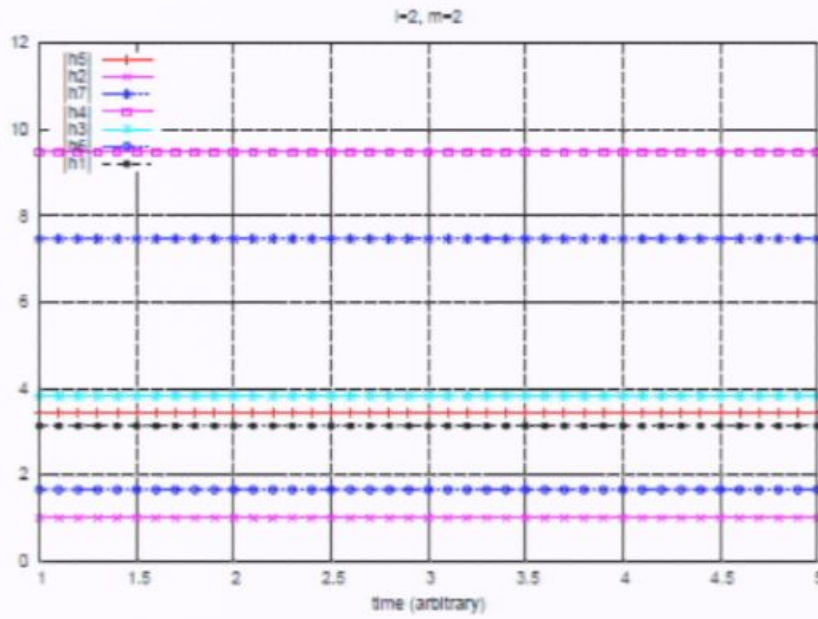
I:  $10^{-11}$  for rel. & abs. accuracies.

II:  $10^{-10}$  for rel. accuracy.

Some Power law scaling for run time as i.e.  $t \sim \ell^\alpha$  with  $\alpha \simeq 2.75$ .



# Results: Comparing $|h^{(i)}|$ with Barack & Sago 2007



# What about the Self-Force?

- Physical SF  $F_{SF}^\mu$  via “mode-sum” scheme:  $F_{SF}^\mu = \sum_{\ell=0}^{\infty} F_{\ell}^{\mu reg \pm}$
- Regularization:** Removes divergent parts of the field so  $\ell$ -modes scale as  $L^{-2}$  in where  $L = \ell + 1/2$  and

$$F_{\ell}^{\mu reg \pm} = F_{\ell}^{\mu \pm} - LA_{\pm}^{\mu} - B^{\mu}$$

No t-regularization in circular orbits, only r.

⇒ Only keeps the timelike support of the Green’s function: tail field  $\bar{h}_{\mu\nu}^{tail}$ .

- Complication: Regularization is done using  $Y^{\ell' m}$  but our decomposition is in  $Y_{\mu\nu}^{\ell m}$ .  
 ⇒ coupling between each tensor mode  $\ell$   
 and up to 7 scalar modes  $\ell - 3 \leq \ell' \leq \ell + 3$ .  
 Complicated final expression for unregularized  $\ell$ -modes:

Large  $\ell$ -tail: truncate at some  $\ell_{max}$ , extrapolate the rest by fitting to  $\frac{D_2}{\ell^2} + \frac{D_4}{\ell^4} \dots$

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# Progress, Current Status and Outlook

- **Currently:** Code computes  $F_{\ell}^{\mu reg \pm}, F_{SF}^{\mu}$  for  $\ell = 0 \dots 28$ .

For  $r_0 = 7M$ , time  $\gtrsim$  **1 hour**;  $r_0 = 10M$ , time  $\sim$  **40 minutes**.

Not there yet: (1)  $F_{\ell}^{\mu reg +} \neq F_{\ell}^{\mu reg -}$ , (2) Fluxes.

Soon: Large  $l$ -tail.

- Speed of progress so far:

Scalar		Gravitational		
Prep.	Coding	Prep	Coding	
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1 month	2 months	5 weeks	5 weeks	3 weeks

- **Future:** Eccentric orbits.

First: Revisit scalar field then gravitational SF.

Expect code to slowdown a lot depending on  $e$ .

Circular  $l = 15$  takes  $\lesssim$  10 minutes at  $r_0 = 7M$  so still not bad.

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