Title: Constraining primordial non-Gaussianity with large-scale structure

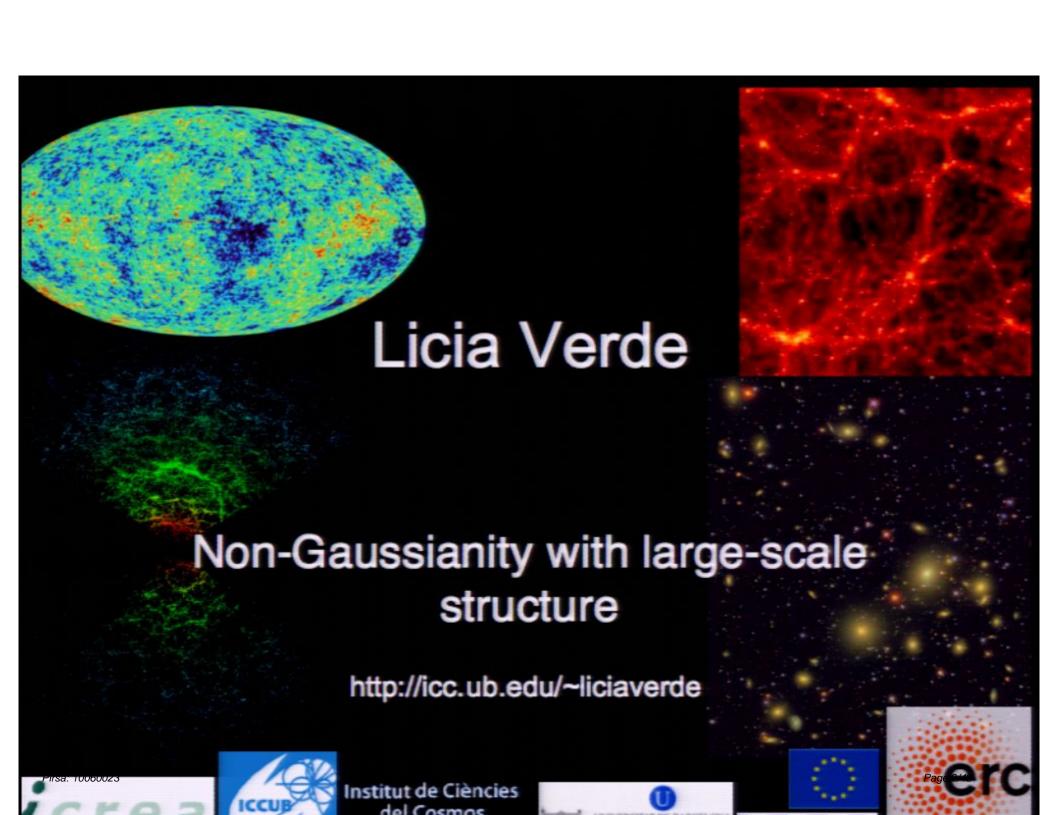
Date: Jun 17, 2010 12:45 PM

URL: http://pirsa.org/10060023

Abstract: Constraining primordial non-Gaussianity can offer a window into the early universe, and into testing the inflationary paradigm, which is fully complementary to the approach offered by Cosmic Microwave Background polarization.

Large-scale structure and galaxy surveys have recently received renewed attention as a way to constrain primordial non-Gaussianity. I will review the potential and the limitations of this approach and highlight its complementarity to Cosmic Microwave Background observations.

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## References

Verde, Matarrese, 2009 "Detectability of the effect of inflationary nongaussian halo bias", ApJL, 706, 91

Reid, Verde, Dolag, Matarrese, Moscardini, 2010, "Non-Gaussian halo assembly bias", arXiv:1004.1637

Mangilli, Verde, 2009 "Non-Gaussianity and the CMB Bispectrum: confusion between Primordial and Lensing-Rees Sciama contribution?" Phys. Rev. D 80, 123007

Carbone, Mena, Verde, 2010. "Cosmological Parameters Degeneracies and Non-Gaussian Halo Bias", arXiv:1003.0456, in press

Jimenez, Verde, 2010, "Implications for Primordial Non-Gaussianity (fNL) from weak lensing masses of high-z galaxy clusters", PRD 80.127302

Xia, Viel, Baccigalupi, De Zotti, Matarrese, Verde, 2010 "Primordial Pirsa: 10060023"

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Non-Gaussianity and the NRAO VLA Sky Survey" ApJLett, 717 (2010) L17-L21, 2010.

### **BASIC MOTIVATION**

#### see M. Zaldarriaga talk

Simplest inflationary models predict SMALL deviations from Gaussian initial conditions

How small is small? (How simple is simple?)
Can in some models "small" can be "detectable"?

There can always be non-standard models (strings, defects etc. yielding larger primordial non-Gaussianity)

Fully complementary approach to looking for r (primordial tensor modes) in the CMB.

But for large-scale structure dedicated telescopes/surveys are not needed: the data will be gathered "anyway".

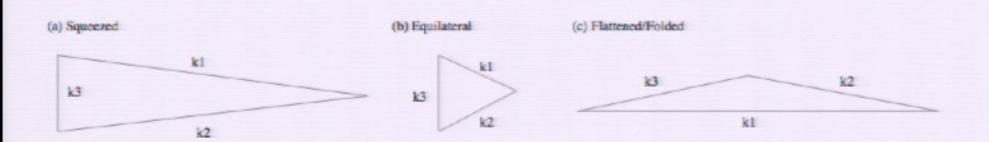
# shapes

#### Simple inflationary model:

One field, canonical kinetic energy, slow roll, Bunch-Davies vacuum

Small LOCAL non-Gaussianity  $\Phi = \phi + f_{NL}(\phi^2 - \langle \phi^2 \rangle)$ 

#### Look at bispectrum



Violation of each of the above conditions leaves a unique signal with specific **shape**From Komatsu et al. 2009, arxiv:0902.4759 & refs. there

## The non-Gaussianity parameter f<sub>NL</sub>:

f<sub>NL</sub> Defined in Fourier space, through the bispectrum, and in general with complex dependence on k (vectors)

But many just say:  $\Phi = \phi + \alpha (\phi^2 - \langle \phi^2 \rangle)$ 

Salopek Bond 1990; Gangui et al 1994; Verde et al 2000; Komatsu Spergel 2001

f<sub>NI</sub> Let's assume it is constant

This is called local model (Creminelli 03)

Typical of when non Gaussianity is generated outside the horizon

Defined on Gravitational potential (actually Bardeen potential, important for sign)

#### Measuring fNL allows us to constraint inflationary models

#### Very very simple example: single field

Remember slow-roll parameters

$$\epsilon_* = \frac{m_{\text{Pl}}^2}{16\pi} \left(\frac{V'}{V}\right)^2$$
, and  $\eta_* = \frac{m_{\text{Pl}}^2}{8\pi} \left[\frac{V''}{V} - \frac{1}{2} \left(\frac{V'}{V}\right)^2\right]$ 

The skewness is

$$S_{3,\Phi} = \frac{\langle \Phi^3 \rangle}{\langle \Phi^2 \rangle^2}$$
  
 $S_{3,\Phi} = 2f_{NL} \times 3[1 + \gamma(n)]$ 

## Measuring fNL allows us to determine the shape of the inflaton potential

Relating the skewnness to the slow-roll parameters

$$f_{NL} = (5/2)\epsilon_* - (5/3)\eta_*$$

But the primordial slope is

$$n=2\epsilon_*-6\eta_*+1$$

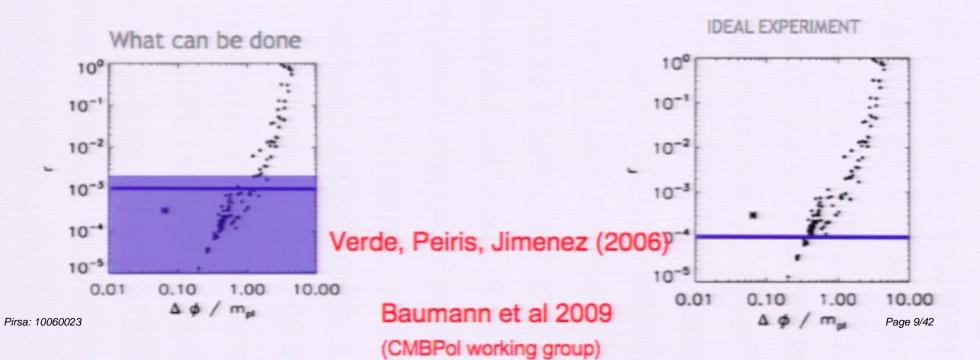
So a measurement of fNL gives you a measurement of the slow-roll parameters

$$\epsilon = 18/35 f_{NL} - n'/7$$
 $\eta = 6/35 f_{NL} - 3/14n'$ 
 $n' = n - 1$ 

#### And the scale of inflation...

$$\Delta_R^2 = \frac{V/M_{pl}^4}{24\pi^2\epsilon}$$

Recall that CMB polarization detection will be very challenging

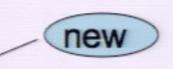


## Galaxy surveys:

Tools:

Bispectrum (or higher orders)

Clustering of peaks on large scales



Abundance of rare events (peaks, massive halos...)

## Searching for non-Gaussianity with LSS: COMPLEMENTARITY

#### Each probe is affected by different systematics

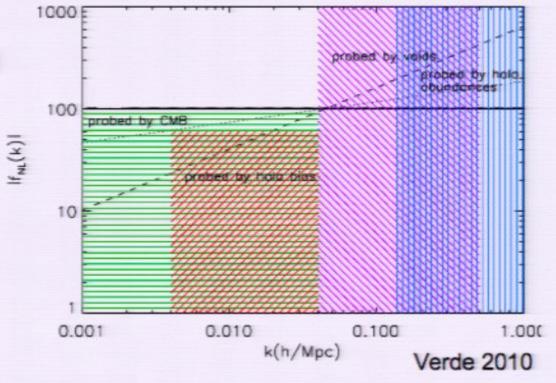
Clustering/spatial properties: Bispectrum, trispectrum, etc.

warning: gravity also generates NG

that's why trispectrum may be

interesting (LV & Heavens 2001)

Abundance of rare events: by looking at the tails of the halo mass function warning: what's a halo and what's its mass? What mass function?



Pirsa: 10060023 anyway interesting: can probe smaller scales than CMB2

# Searching for non-Gaussianity with LSS

Bispectrum, clustering; inflation-type

Verde et al. (1999) and Scoccimarro et al. (2004) showed that constraints on primordial NG in the gravitational potential from large redshift-surveys like 2dF and SDSS are not competitive with CMB ones: fnl has to be larger than 10<sup>2</sup>- 10<sup>3</sup> in order to be detected as a sort of non- linear bias in the galaxy-to-dark matter density relation. However LSS gives complementary constraints as it tests different scales than CMB.

Going to redshift z~2 can make LSS competitive (Sefusatti & Komatsu 2007). Going to higher z (e.g. through SZ cluster surveys or via 21-cm background anisotropies) helps, as the effective NG strength in the underlying CDM overdensity scales like (1+z) (LV et al 1999, Pillepich, Porciani & Matarrese 2006; Cooray 2006).

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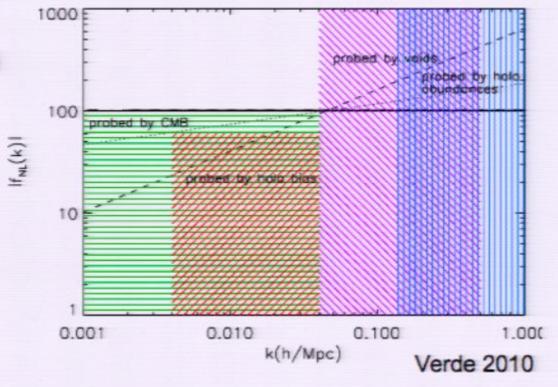
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# Searching for non-Gaussianity with LSS

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## Abundance of rare events

- Besides using standard statistical estimators, like bispectrum, trispectrum, three and four-point function, skewness, etc. ..., one can look at the tails of the distribution, i.e. at rare events.
- Rare events have the advantage that they often maximize deviations from what predicted by a Gaussian distribution, but have the obvious disadvantage of being ... rare!

Primordial non-Gaussianity also strongly affects the abundance of the first non-linear objects in the Universe, thereby modifying the reionization history.

- Matarrese LV & Jimenez (2000) and Verde, Jimenez, Kamionkowski & Matarrese showed that clusters at high redshift (z>1) can probe NG down to f<sub>NL</sub> ~ few 10 which is, however, not competitive with future CMB (Planck) constraints (but probe different scales)
- Voids (rare events) may be competitive (Kamionkowski, Verde, Jimenez 2008)

MVJ(2000) mass function & improved formulae obtained by LoVerde, Miller, Shandera, Verde 2007, Grossi et al 09, Maggiore, Riotto 2009, D'Amico et al

## Abundance of Rare events

$$\frac{dn}{dM} = f \frac{\rho_b}{M} \left| \frac{dP(>\delta_c|z,M)}{dM} \right| \quad \text{Press Schechter approach}$$

Non-Gaussianity changes especially the tails

Note: this was derived in the Press-Schechter framework. PS fails at some point (spherical collapse etc.).

Recommended: use the ratio NG/G (compare this to observations normalized to numerically calibrated Gaussian predictions)  $1+S_3 \delta_c^3/(6 \sigma_R^2)$ 

Must calibrate on N-body simulations. (e.g., Grossi et al. 09, Desjaques et al 09, Pillepich et al 10)

## Abundance of Rare events

A lightening fast history if the non-gaussian mass function!

#### Matarrese, LV & Jimenez (2001)

Derive the mass function for non-Gaussian fields using an approximation valid for rare events (MVJ) use a Press-Schecter-type approach

The resulting expressions should be used for fractional corrections

#### LoVerde, Miller, Shandera, Verde (2008)

Make different approximations, valid for not-so-rare events (LMSV, 1&2)

#### Grossi, Verde et al (2009)

Calibrate on simulations, LMSV, MVJ work, need a factor  $\sqrt{q}$  in front of  $\delta_c$ 

#### Riotto & Maggiore (2009)

Similar to LVMSV but improving over the Press-Schecter-type approach, giving the q factor (in the right place)

New: D'Amico, Musso, Norena, Paranjape

Get the best of three words: LVMSV2 not too shabby (in theory)!

## Non-Gaussian halo bias

- A Gaussian field and a non-Gaussian field can have the same P(k)
- In a Gaussian field the P(k) of peaks is completely specified by the P(k)
- In a non-Gaussian field, however, the P(k) of the peaks, depends on all higher order correlations (i.e. fnl)

## Non-Gaussian halo bias

- Gaussian IC and a non-Gaussian IC can have the same P(k) for the dark matter
- For Gaussian IC the P(k) of massive halos is completely specified by the dark matter P(k)
- For Non Gaussian IC, however, the P(k) of the halos, depends on all higher order correlations (i.e. fnl)

## Non-Gaussian halo bias

### For Gaussian initial conditions (known since the '80)

$$\xi_{h,M}(r)=\exp\left[rac{
u^2}{\sigma_R^2}\xi_R(r)
ight]-1\simeqrac{
u^2}{\sigma_R^2}\xi_R(r)$$
  $b_E=1+b_L$  "The Kaiser formula"

In the '90 this was improved (e.g. Mo & White 1996, Catelan et al 1998)

#### For Non-Gaussian initial conditions

Dalal et al. PRD 2008 7713514

Matarrese, Verde, ApJLett, 2008, 77:L77

Slosar et al 08

McDonald 08

Afshordi & Tolley 08

Valageas 2009

A scale-dependent bias! (on top of the Gaussian one and proportional to it)

## The Effect

$$\frac{\Delta b}{b^{L,G}} = \beta(k) \frac{\Delta_c(z)}{D(z)}$$

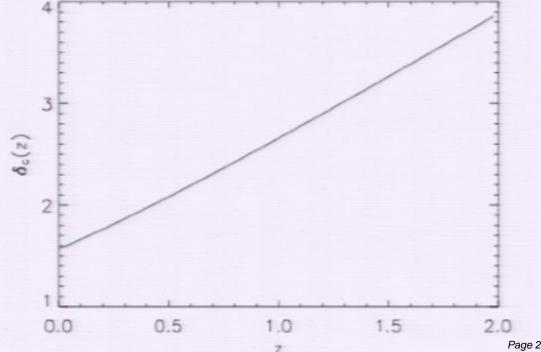
$$\delta(k) = \mathcal{M}_R(k)\Phi(k)$$

$$\alpha=k_1^2+k^2+2k_1k\mu$$

$$\beta(k) =$$

$$\frac{1}{8\pi^2\sigma_R^2\mathcal{M}_R(k)}\int dk_1k_1^2\mathcal{M}_R(k_1)\int_{-1}^1 d\mu\mathcal{M}_R\left(\sqrt{\alpha}\right)\frac{B_\phi(k_1,\sqrt{\alpha},k)}{P_\phi(k)}.$$

#### Redshift dependence



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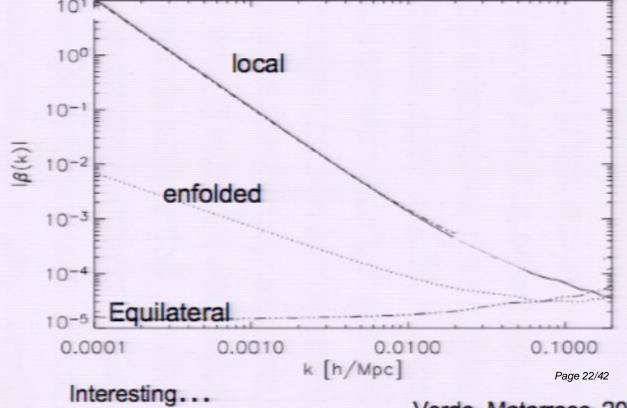
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## The Effect

$$\frac{\Delta b}{b^{L,G}} = \beta(k) \frac{\Delta_c(z)}{D(z)} \qquad \delta(k) = \mathcal{M}_R(k) \Phi(k) \qquad \alpha = k_1^2 + k^2 + 2k_1 k \mu$$

$$\beta(k) = \frac{1}{8\pi^2 \sigma_R^2 \mathcal{M}_R(k)} \int dk_1 k_1^2 \mathcal{M}_R(k_1) \int_{-1}^1 d\mu \mathcal{M}_R\left(\sqrt{\alpha}\right) \frac{B_{\phi}(k_1, \sqrt{\alpha}, k)}{P_{\phi}(k)}.$$

Scale-dependence 2 10-2



Matarrese, Verde 08 Verde, Matarrese 09

Verde Matarrese 2009

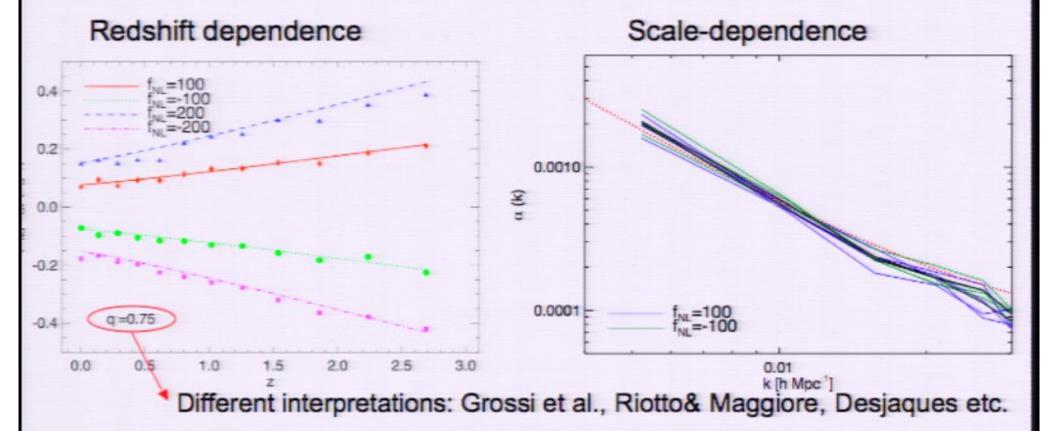
## Calibration on simulations

e.g., Grossi, Verde, Carbone et al. 2009, MNRAS, 398, 321

Volume: (1.2 Gpc/h)<sup>3</sup> 960<sup>3</sup> particles

 $m \approx 1.4 \times 10^{11} h^{-1} M_{\odot}$ 

Local non-gaussianity  $f_{NI} = 0$ , -100, +100, -200, +200



## Current constraints

Data/method	f <sub>NL</sub>	reference
Photometric LRG - bias	63 <sup>+54+101</sup> -85-331	Slosar et al. 2008
Spectroscopic LRG- bias	$70^{+74+139}_{-83-191}$	Slosar et al. 2008
QSO - bias	$8^{+26+47}_{-37-77}$	Slosar et al. 2008
combined	$28^{+23+42}_{-24-57}$	Slosar et al. 2008
NVSS-ISW	$105^{+647+755}_{-337-1157}$	Slosar et al. 2008
NVSS-ISW	$236 \pm 127(2 - \sigma)$	Afshordi&Tolley 2008

### Local-type only

### CMB bispectrum (WMAP7 Komatsu et al 2010)

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$$f_{\rm NL}^{\rm local} = 32 \pm 21 \; (68\% \; {\rm CL}) \qquad -10 \; < f_{\rm NL}^{\rm local} \; < \; 74.$$

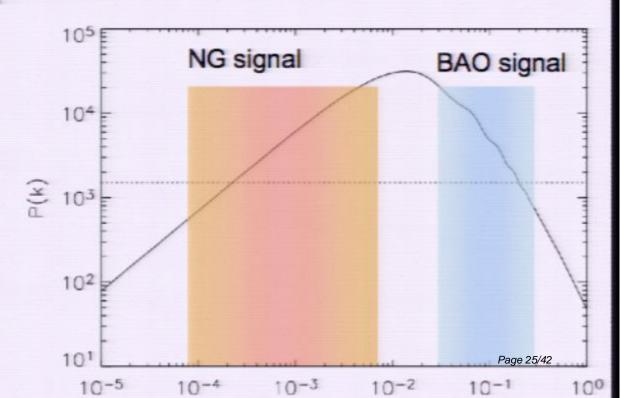
## Interesting features

Spectro- vs photo- z's Smooth feature in k (this is not BAO)
Smooth behavior in z
Standard photo-z accuracy will suffice

#### BAO surveys well suited!

Large volumes High-z Appropriate shot noise

If nP~1 at k=0.2 Then nP~1 at k=10-4



## How well can this do? Local

survey	z range	sq deg	mean galaxy density $(h/Mpc)^3$	$\Delta f_{ m NL}/q'$ LSS
SDSS LRG's	0.16 < z < 0.47	$7.6 \times 10^{3}$	$1.36 \times 10^{-4}$	40
BOSS	0 < z < 0.7	104	$2.66 \times 10^{-4}$	18
WFMOS low z	0.5 < z < 1.3	$2 \times 10^{3}$	$4.88 \times 10^{-4}$	15
WFMOS high z	2.3 < z < 3.3	$3 \times 10^{2}$	$4.55 \times 10^{-4}$	17
ADEPT	1 < z < 2	$2.8 \times 10^{4}$	$9.37 \times 10^{-4}$	1.5
EUCLID	0 < z < 2	$2 \times 10^{4}$	$1.56 \times 10^{-3}$	1.7
DES	0.2 < z < 1.3	$5 \times 10^{3}$	$1.85 \times 10^{-3}$	8
PanSTARRS	0 < z < 1.2	$3 \times 10^{4}$	$1.72 \times 10^{-3}$	3.5
LSST	0.3 < z < 3.6	$3 \times 10^{4}$	$2.77 \times 10^{-3}$	0.7

## How well can this do? Local

Data/method	$\Delta f_{\rm NL} \ (1-\sigma)$	reference	
BOSS-bias	18	Carbone et al 2008	
ADEPT/Euclid-bias	1.5	Carbone et al 2008	
PANNStarrs -bias	3.5	Carbone et al 2008	
LSST-bias	0.7	Carbone et al 2008	
LSST-ISW	7	Afshordi& Tolley 2008	
BOSS-bispectrum	35	Sefusatti & Komatsu 2008	
ADEPT/Euclid -bispectrum	3.6	Sefusatti & Komatsu 2008	
Planck-Bispectrum	3	Yadav et al . 2007	
BPOL-Bispectrum	2	Yadav et al . 2007	

Carbone, Verde, Matarrese 08

Carbone, Mena, Verde 2010:

there is no much degeneracy with cosmology!

# Inflationary-GR Intrinsic to LSS

Bartolo, Matarrese, Riotto 2005, Bartolo et al 2006 Pillepich, Porciani, Matarrese, 2007

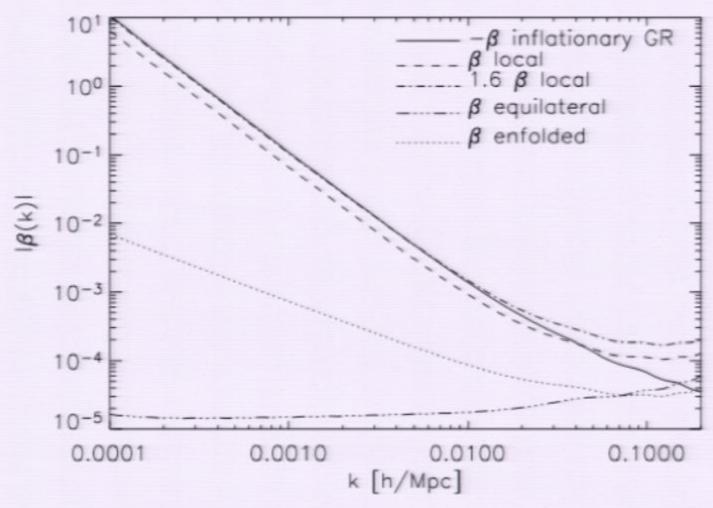
$$B_{\Phi}(k_1, k_2, k_3) = 2\left[\frac{5}{3}(a_{\rm NL} - 1) + f_{\rm NL}^{\rm infl, GR}(k_1, k_2, k_3)\right]P(k_1)P(k_2) + cyc.$$

$$f_{\mathrm{NL}}^{\mathrm{infl,GR}}(k_i,k_j,k_k) = -rac{5}{3}\left[1-rac{5}{2}rac{k_ik_jcos heta_{ij}}{k_k^2}
ight]$$

On horizon-scales Poisson equation gets quadratic corrections: Needs IC set up of inflation, parallels the TE anti-correlation.

Verde & Matarrese 2009, ApJL

# Inflationary-GR Intrinsic to LSS



n horizon-scales Poisson equation gets quadratic corrections: eeds IC set up of inflation, parallels the TE anti-correlation.

Verde & Matarrese 2009, ApJL

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## Selection effects?

Assembly bias! (discussion on the gaussian case amplitude of the effect, Non-gaussian case effect expected to be larger...)

$$b_L^G = ar{n}^{-1} rac{\partial n}{\partial \delta_l} = ar{n}^{-1} rac{\partial n}{\partial \delta_c},$$

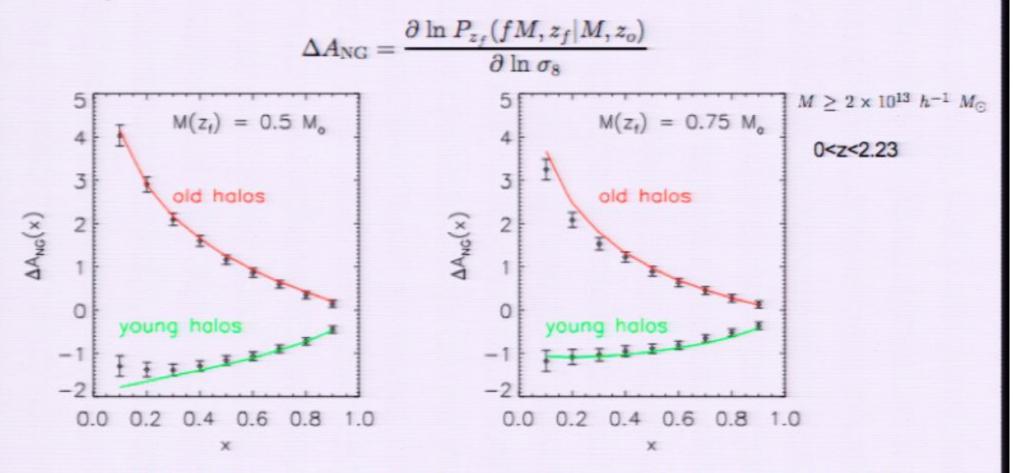
Peak Background split Slosar et al 2008

Local case only

$$\Delta b_{
m NG}(M,k,z_o) = rac{2f_{
m NL}}{D(z_o)\mathcal{M}(k)} rac{\partial \ln n(M,z_o)}{\partial \ln \sigma_8}$$

$$\Delta b_{
m NG}(M,k,z_o,z_f,f) =$$
 Reid et al. 2010, Extended PS approach:  $\frac{2f_{
m NL}}{D(z_o)\mathcal{M}(k)} \left( \frac{\partial \ln n(M,z_o)}{\partial \ln \sigma_8} \left( \frac{\partial \ln P_{z_f}(fM,z_f|M,z_0)}{\partial \ln \sigma_8} \right) \right)$ 

## Dependence on halo formation time



x is the fraction of halos with the highest (lowest) formation redshifts entering the sample

Points: simulations Grossi et al 2009

Line: ePS

$$ilde{\omega}_f = rac{\delta_c(z_f) - \delta(z_0)}{\sqrt{\sigma^2(fM) - rac{\sigma^2(M)}{
ho}}} \,.$$

## Is this worrying?

The effect is asymmetric between old and young halos. Even if the tracer population excludes only the 10% oldest halos, the value of the correction for the remaining 90% of the halos differs from the full sample by 0.44; at z = 0, this amounts to a change of 40%.

Using the closest possible sample to recent major mergers get -1: the ePS prediction derived in Slosar et al. 2008

### Between theory and observations there is an ocean....

#### Go to Semi-analytic models

Non-Gaussian assembly bias in the Bertone et al. (2007) mock galaxy catalogs

$\bar{n} (h^{-1} \text{Mpc})^{-3}$	selection criteria	$\Delta A_{ m NG}^{ m gal}$	$b_G$	$\delta_c(b_G-1) \approx A_{ m NG}^{ m tot}$	% change
$4.5 \times 10^{-4}$	$\geq 8 \times 10^{10} \ h^{-1} \ M_{\odot}$	0.51	2.3	2.2	23%
$4.5 \times 10^{-4}$	$\geq 24~M_{\odot}/{\rm yr}$	0.24	1.3	0.51	48%

"observations" do not correlate well with halo formation time?????

# Tantalizing hints (this year only)

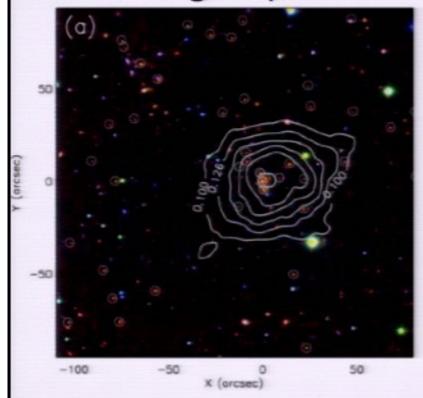
#### XMMUJ2235.3-2557

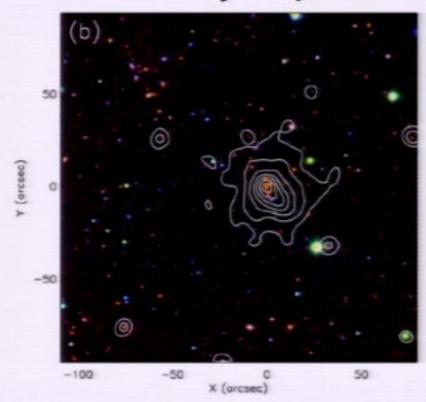
 $(8.5 \pm 1.7) \times 10^{14} M_{\odot}$ 

z=1.4

Lensing + optical

X-ray + optical





Declared survey area: 11 sq deg

Jee, et al., 2009, ApJ, 704, 672, arXiv:0908.3897

## XMMUJ2235.3-2557

œ

 $(8.5 \pm 1.7) \times 10^{14} M_{\odot}$  z=1.4

M> central estimate expect ZERO in the  $4\pi$ 

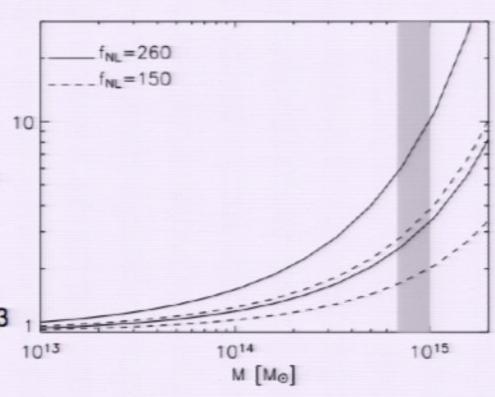
M> lower estimate expect 7 in the  $4\pi$ 

Jimenez, Verde, 2010 arXiv:0909.0403

Sartoris et al. arXiv:1003.0841

Holz, Perlmutter, arXiv:1004.5349

Cayon et al arXiv:1006.1950



NON-GAUSSIAN ENHANCEMENT

Weak lensing area 11 sq deg -> P=0.005 XMM serendipitous survey area

in 2006: 165 sq deg

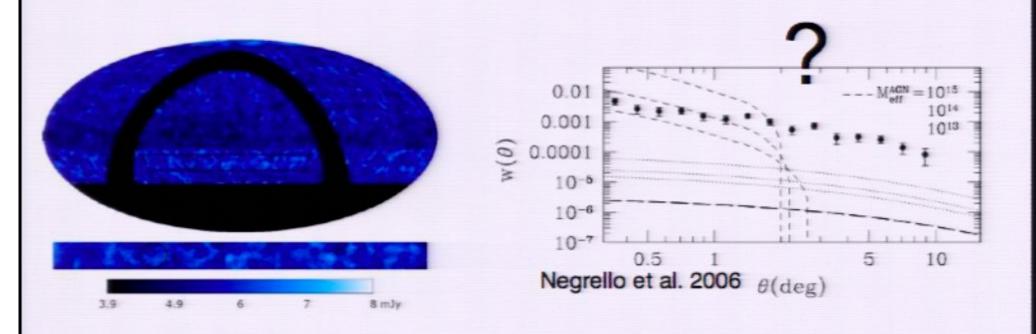
Now: 400 sq deg

P=0.07

-0 17

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## **NVSS** correlation function



### Explanations include:

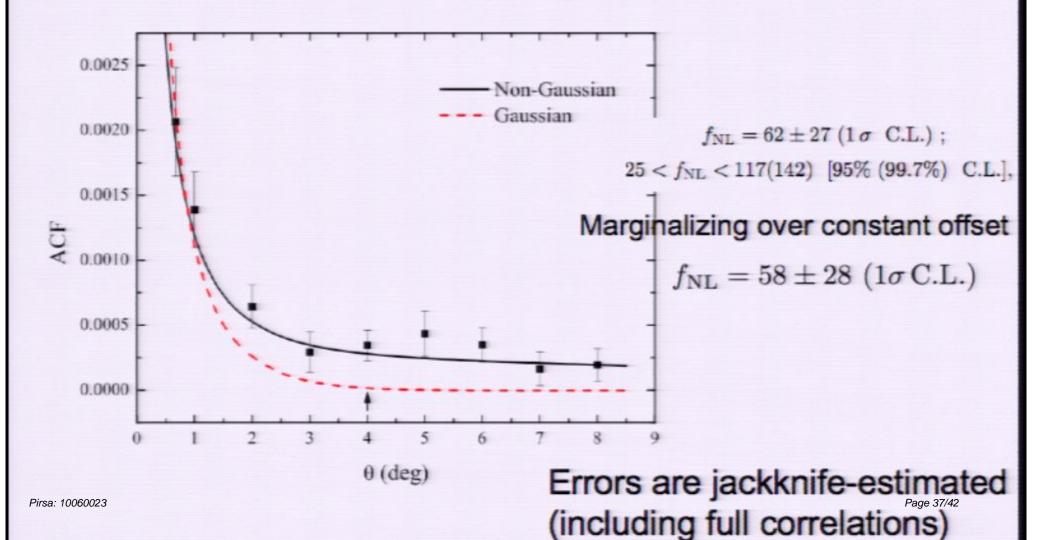
inker with redshift distribution of sources...

Evolution of bias factor radically different from that of optical QSOs e.g., Magliocchetti et al 1999, Negrello et al. 2006; Massardi et al. 2010....)

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# Could it be f<sub>NI</sub>?

Xia, Viel, Baccigalupi, DeZotti, Matarrese, LV (2010)



My point being: if there is an fnl local ~few x 10 out there, today we have the capability of seeing it.

Extraordinary claims need extraordinary proofs

# Complementarity

# Putting it all together

	CMB Bispectrum		Halo	bias
type NG	Planck	(CM)BPol	Euclid	LSST
	$1-\sigma$ errors			
Local	$3^{A)}$	$2^{A)}$	$1.5^{B}$ )	$0.7^{B)}$
Equilateral	$25^{C)}$	$14^{C}$	-	_
Enfolded	O10	O10	$39^{E)}$	$18^{E)}$
		#σ Detec	tion	
GR	N/A	N/A	$1^{E)}$	$2^{E)}$
Secondaries	$3^F)$	$5^F)$	N/A	N/A

YADAV, KOMATSU & WANDELT (2007) B)

CARBONE ET AL. (2008) C) BAUMANN ET AL. (2009);

SERUSATTI ET AL. (2009) E)Verde & Matarrese 2009

F) Mangilli & Verde 2009, Hanson et al. 2009

# Conclusions & future prospects

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\* Constraining/detecting non-Gaussianity is a powerful tool to discriminate among competing scenarios for perturbation generation (standard inflation, curvaton, modulated-reheating, DBI, ghost inflation, multi-field, etc. ...) some of which imply large non-Gaussianity. Non-Gaussianity will soon become the smoking-gun for (non?)-standard inflation models.

# Conclusions & future prospects

- Constraining/detecting non-Gaussianity is a powerful tool to discriminate among competing scenarios for perturbation generation (standard inflation, curvaton, modulated-reheating, DBI, ghost inflation, multi-field, etc. ...) some of which imply large non-Gaussianity. Non-Gaussianity will soon become the smoking-gun for (non?)-standard inflation models.
- Complementary approach to looking for primordial B models in CMB polarization
- Constraining non-Gaussianity in LSS allows to put independent limits on NG and on a different range of scales. Massive/high redshift objects (rare events) are most sensitive to primordial non-Gaussianity, both in their abundance and clustering (bias).
- Predicting/constraining non-Gaussianity is ready to become a branch of Precision Cosmology: this requires accurate analytical calculations, highresolution numerical simulations.

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