

Title: Entanglement entropy and infinite randomness fixed points in disordered magnetic and non-abelian quasi-particle chains

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Abstract: Many one dimensional random quantum systems exhibit infinite randomness phases, such as the random singlet phase of the spin-1/2 Heisenberg model. These phases are typically the result of destabilizing systems described by a conformal field theory with disorder. Interestingly, entanglement entropy in 1d infinite randomness phases also exhibits a universal log scaling with length. In my talk I will touch upon calculating the entanglement entropy for infinite-randomness phases, as well as describe the exotic infinite randomness phases realized in chains of non-abelian anyon chains. It was speculated that the entanglement entropy of an infinite-randomness phase is associated with the direction of RG flow, just as the c-theorem dictates the direction of RG flows for CFT's. I will also show that the entanglement entropy in disordered non-abelian chains provide the only known counter example.

Random spin chains: Abelian, non-Abelian, and non-linear

Gil Refael (Caltech)

Non-Abelian collaborators:

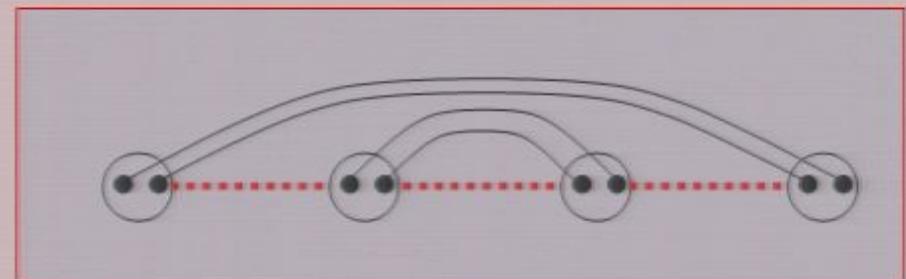
Lukasz Fidkowski (Caltech)

Joel Moore (UC Berkeley)

Nick Bonsteel (FSU)

Undergraduate forced labor:

Paraj Titum (IIT Kanpur), Han Hsuan Lin (MIT)



Pirsa: 10050071

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Sloan Foundation
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Outline

- Disordered Heisenberg models.

Random singlet and infinite randomness fixed points.

- Entanglement entropy in disordered spin chains

- Random Non-Abelian spin chains:

- Fibonacci anyons
- $SU(2)^k$ anyons

Disordered Heisenberg Chains

Heisenberg spin chains

Heisenberg spin chains

$$H = J \sum_i \vec{S}_i \cdot \vec{S}_{i+1}$$

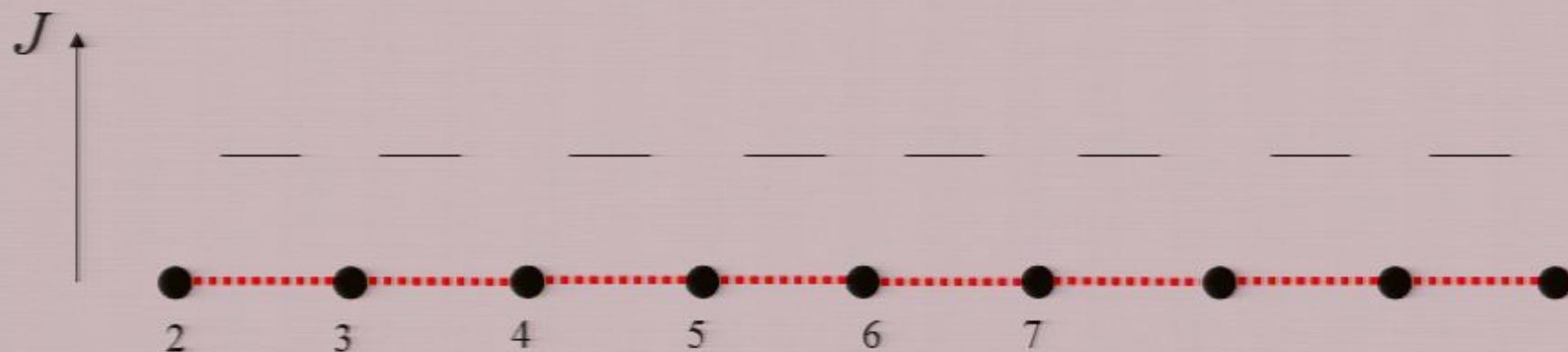
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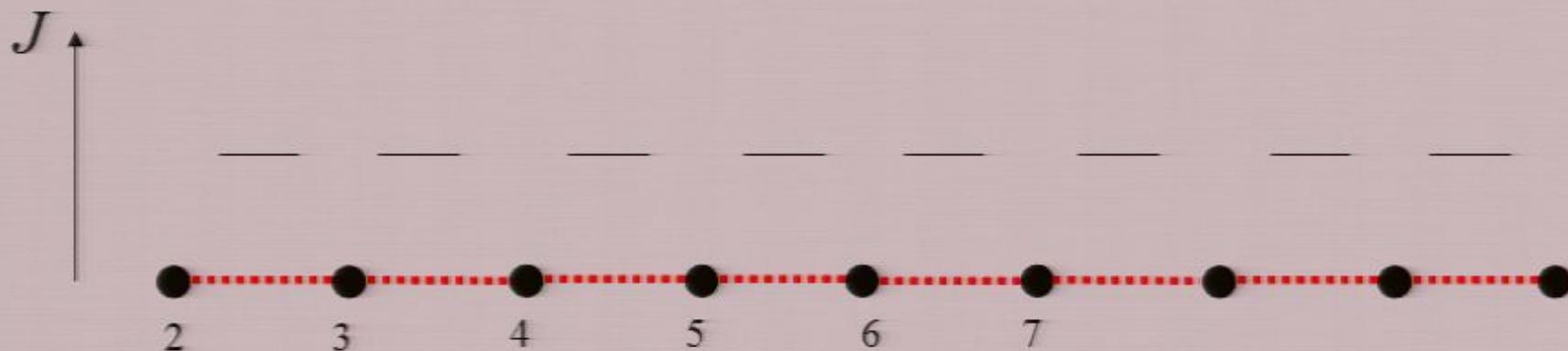
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Solution methods rely on translational invariance:

Heisenberg spin chains

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Solution methods rely on translational invariance:

- Bethe ansatz
- Bosonization

Spin chains with randomness

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Spin chains with randomness

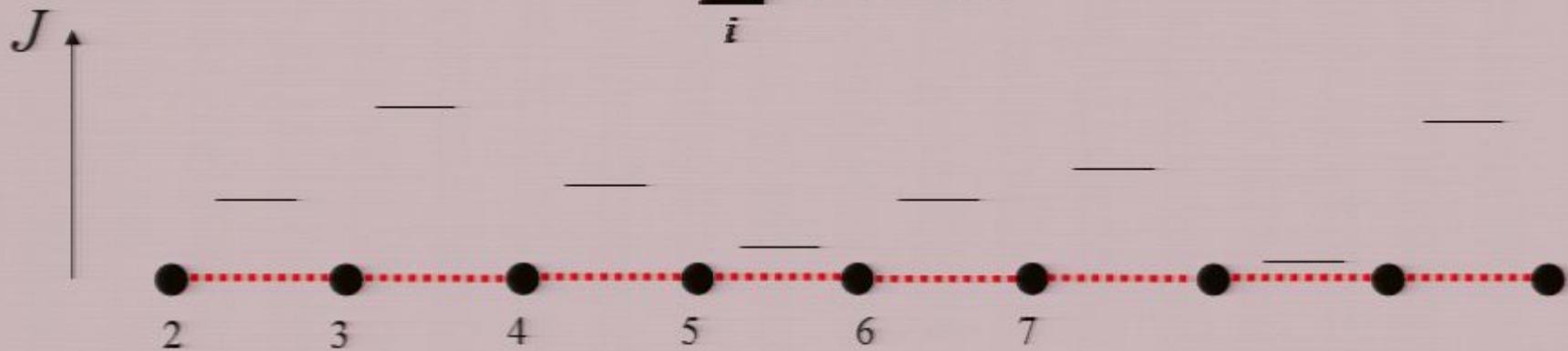
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$$H = \dots + J_2 \vec{S}_2 \cdot \vec{S}_3 + J_3 \vec{S}_3 \cdot \vec{S}_4 + J_4 \vec{S}_4 \cdot \vec{S}_5 \dots$$

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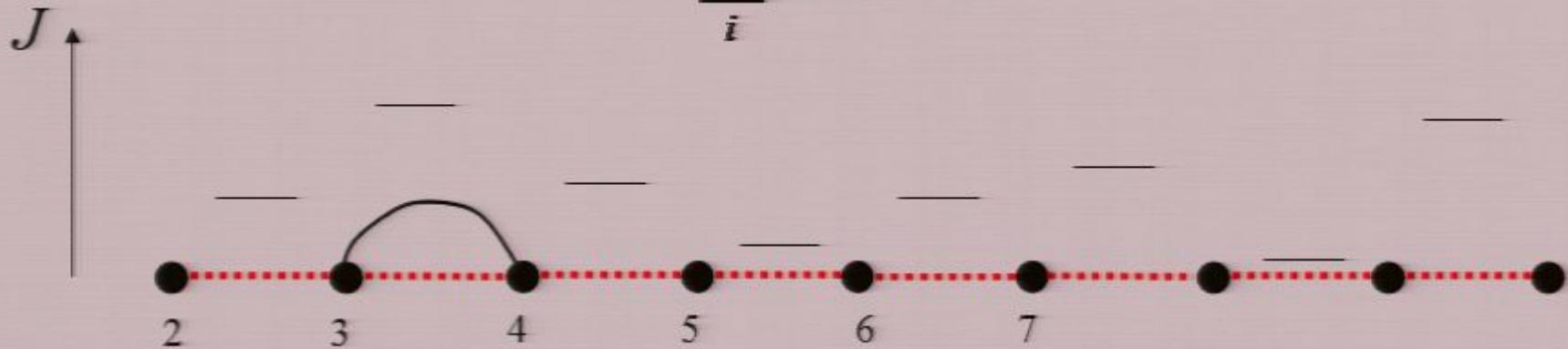


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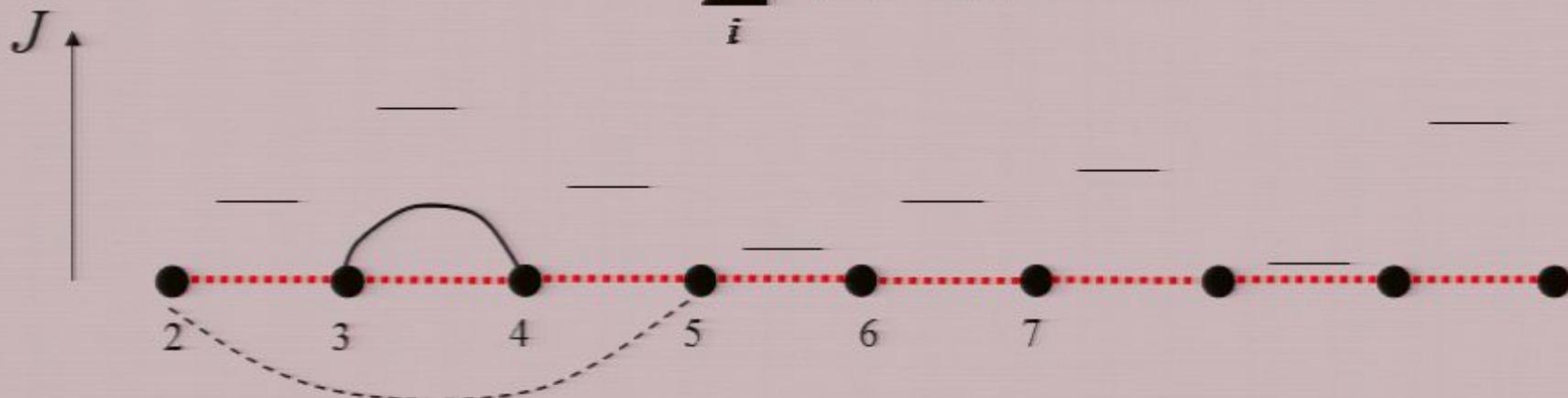


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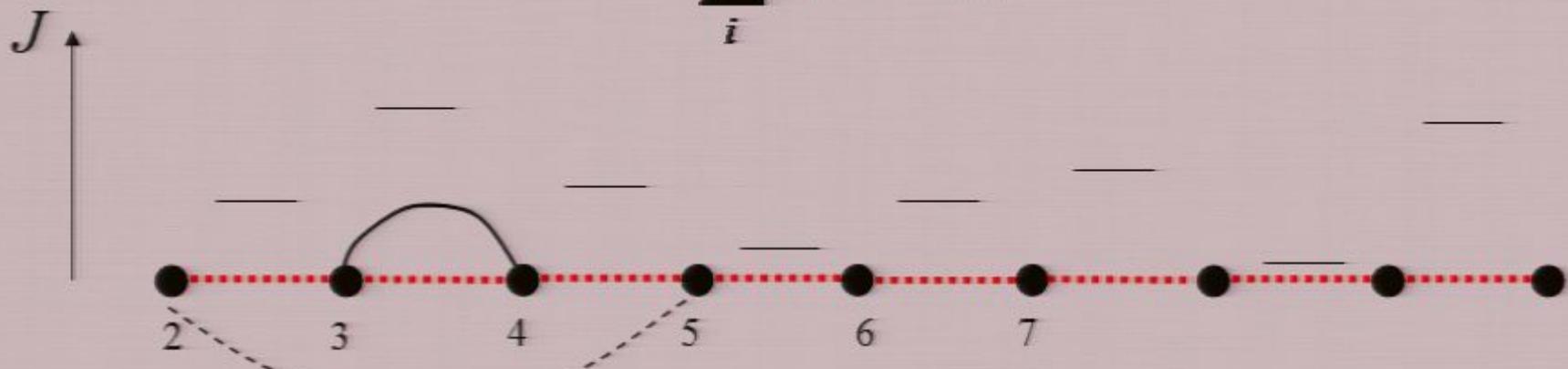
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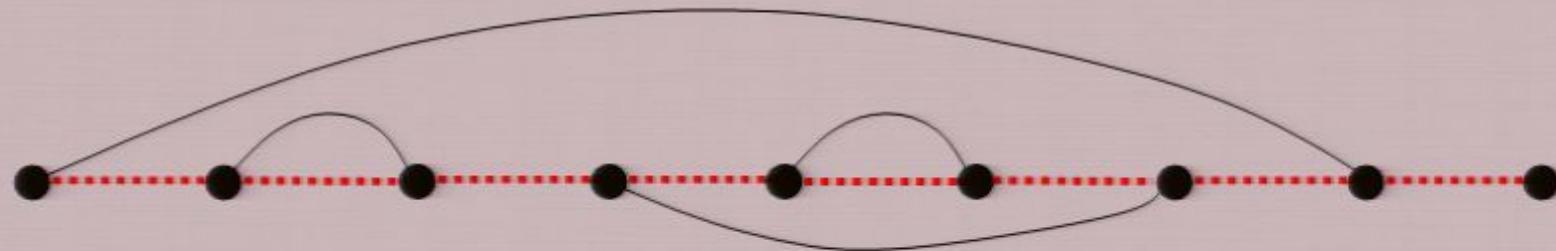
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[D.S. Fisher (1994), Bhatt, Lee (1982)]



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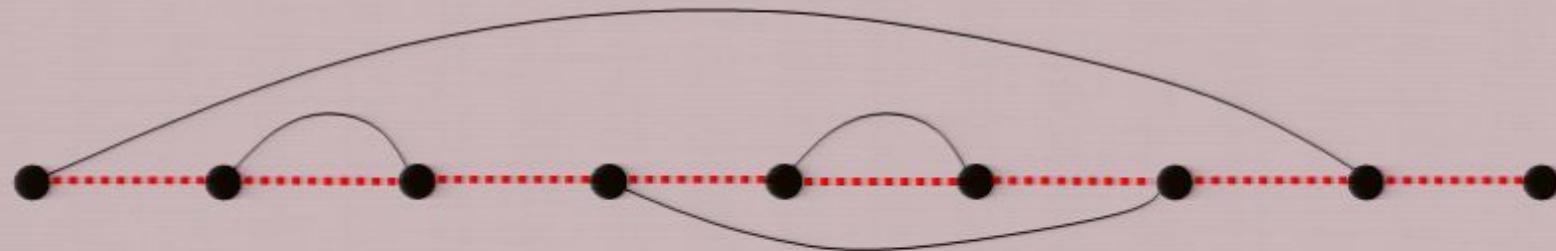
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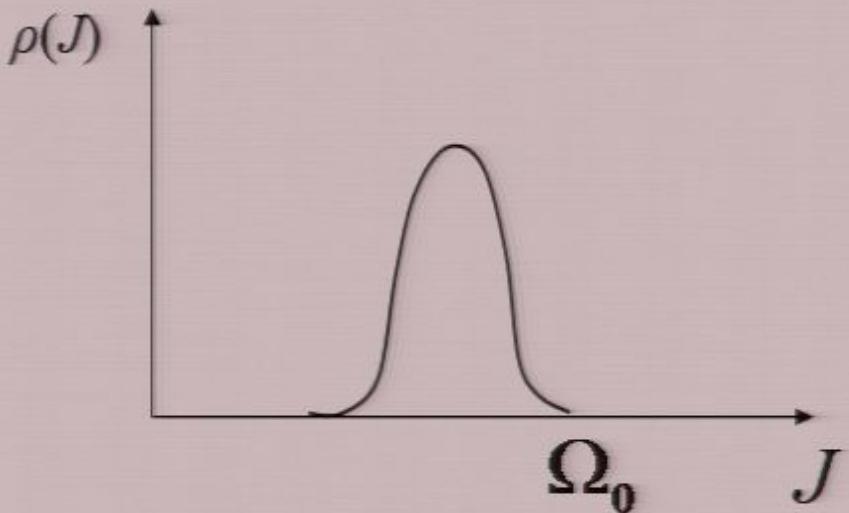
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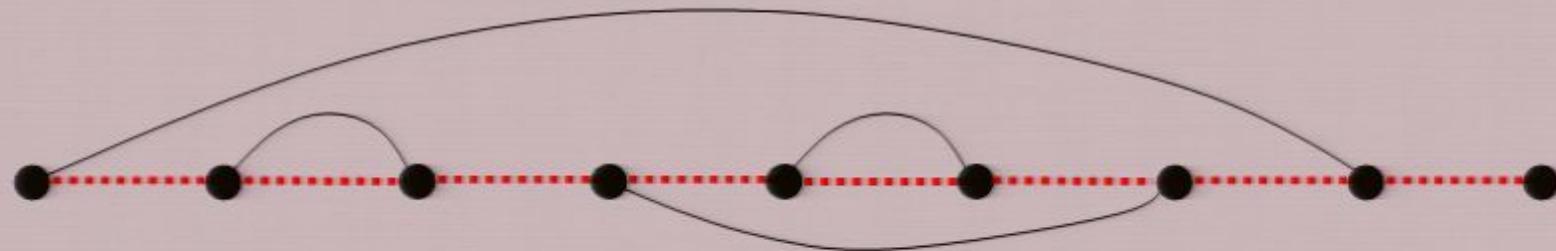


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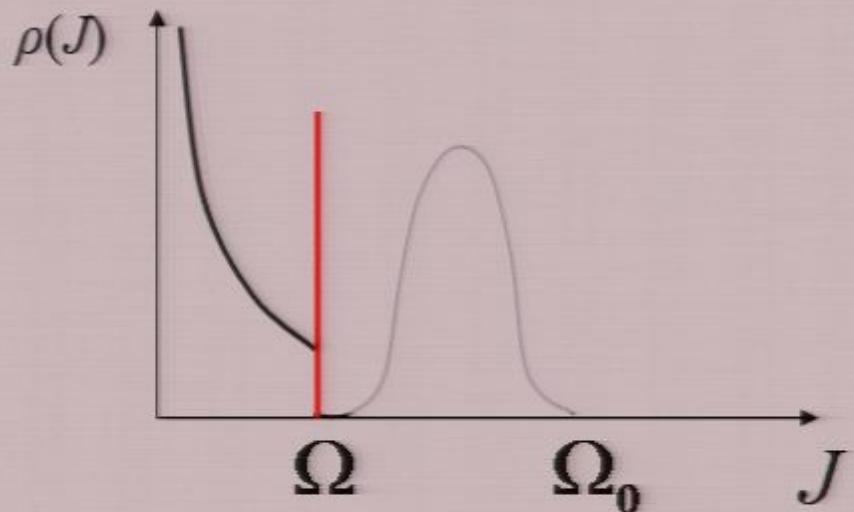
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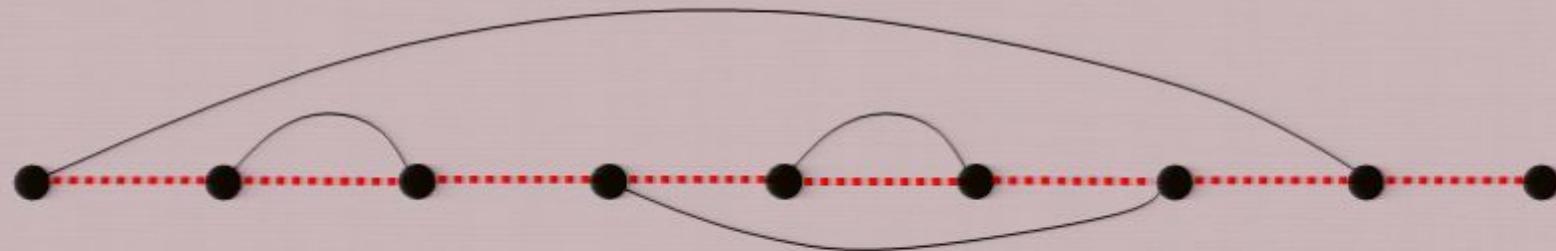
$$\rho(J) = \frac{\chi}{\Gamma} \cdot \frac{1}{\Omega} \left(\frac{\Omega}{J} \right)^{1-\chi/\Gamma}$$

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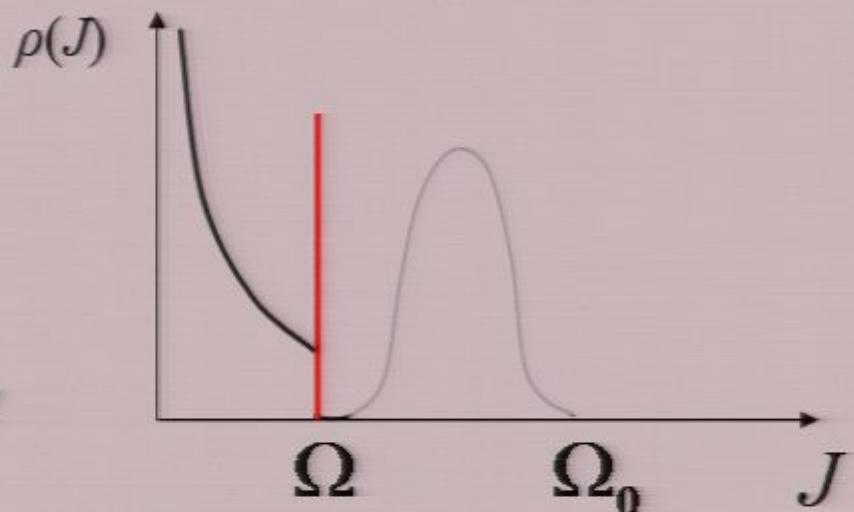


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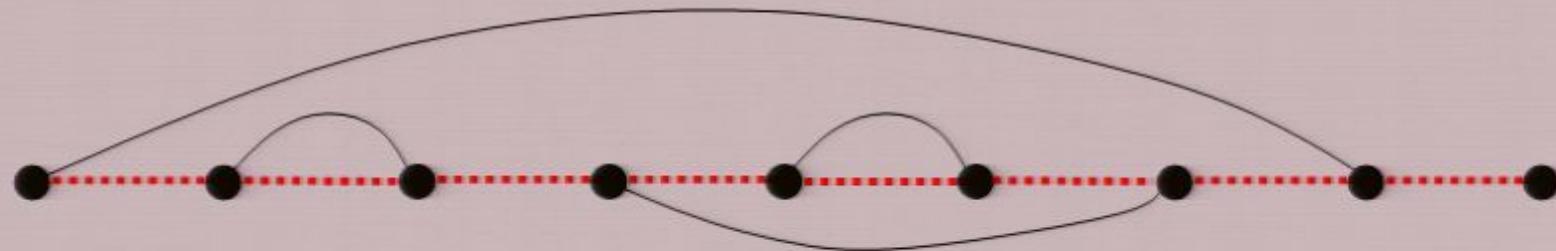
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For all initial distributions of the disorder!



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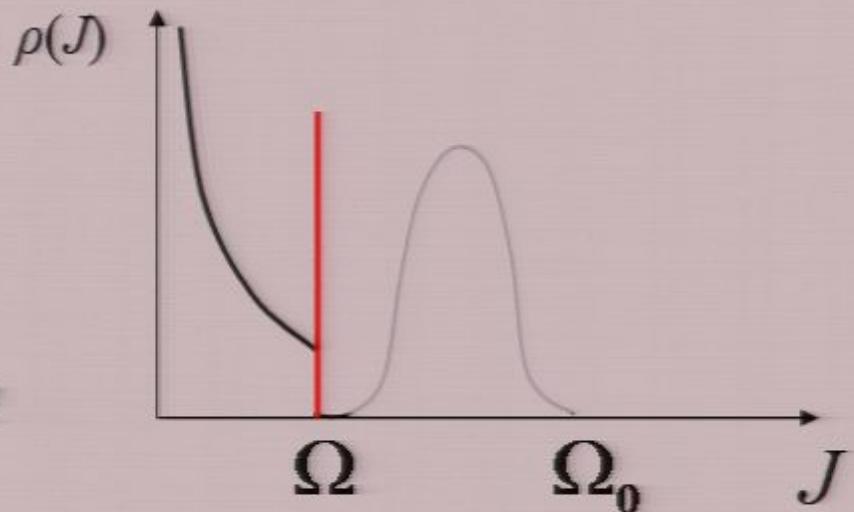


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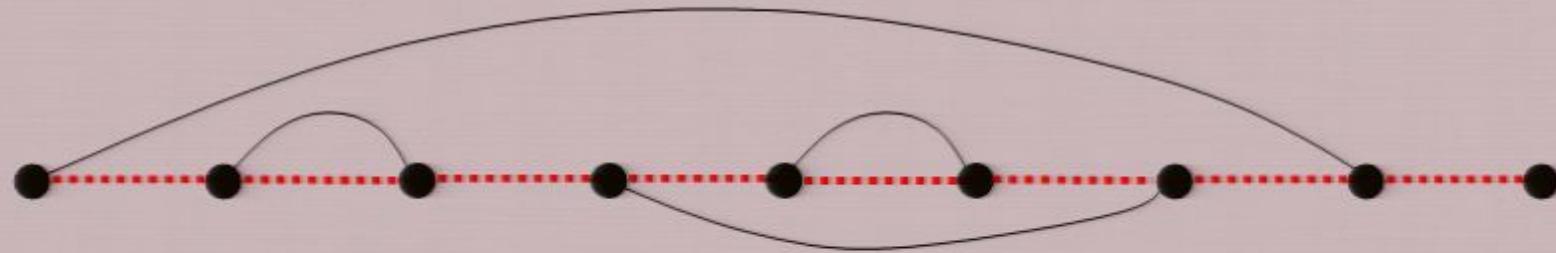
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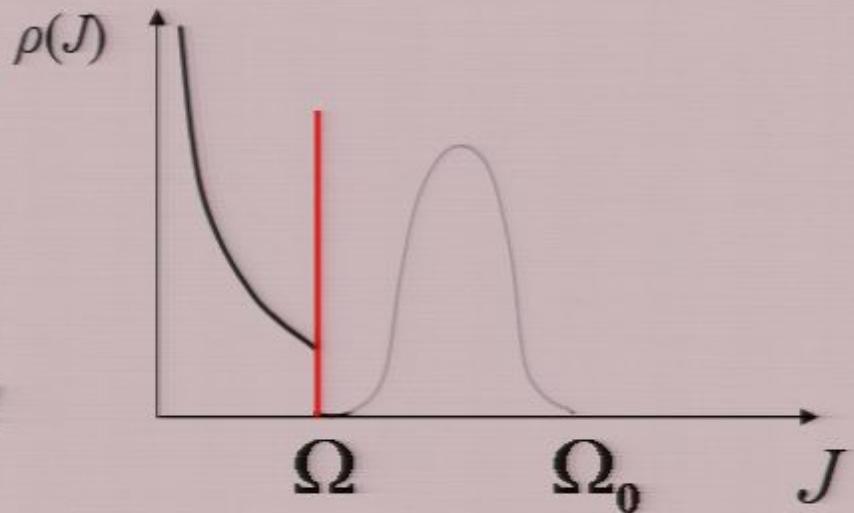


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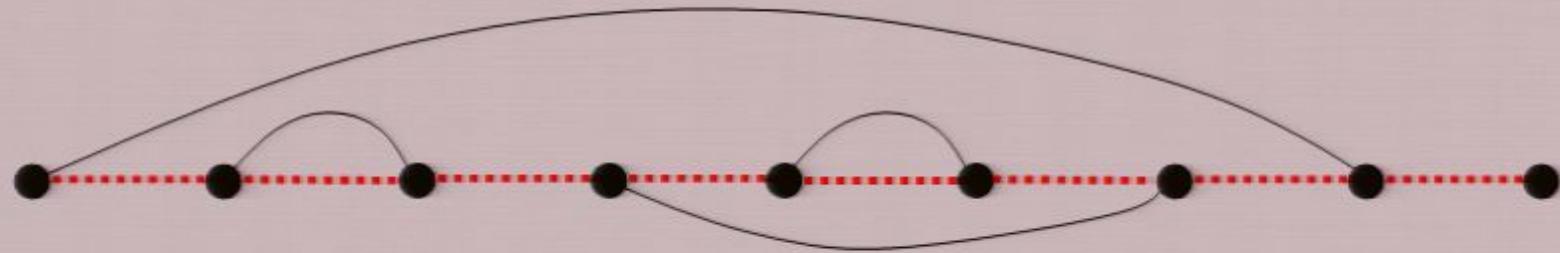
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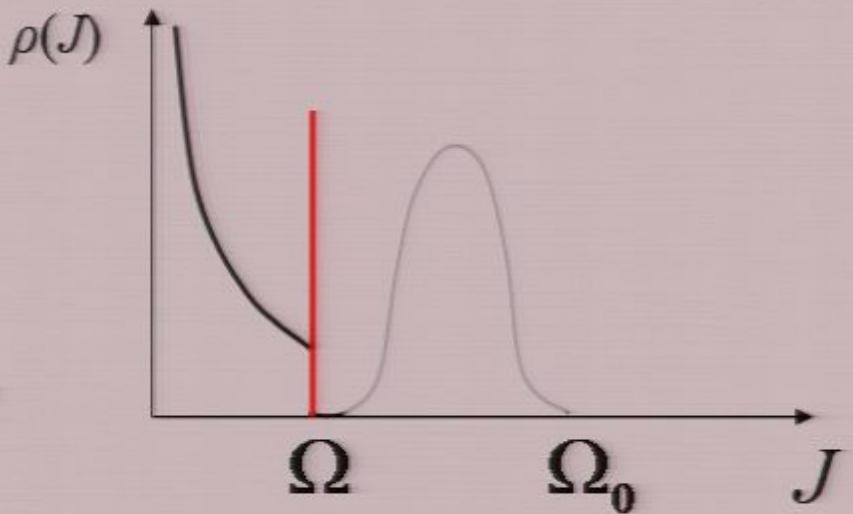


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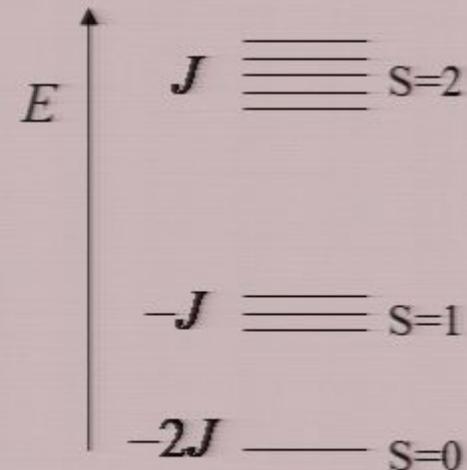
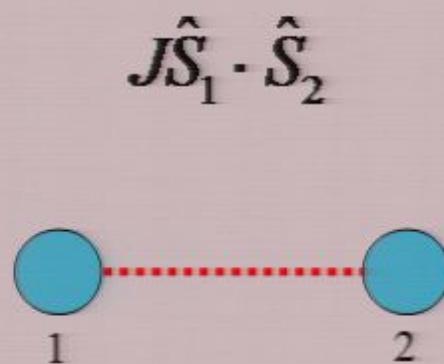
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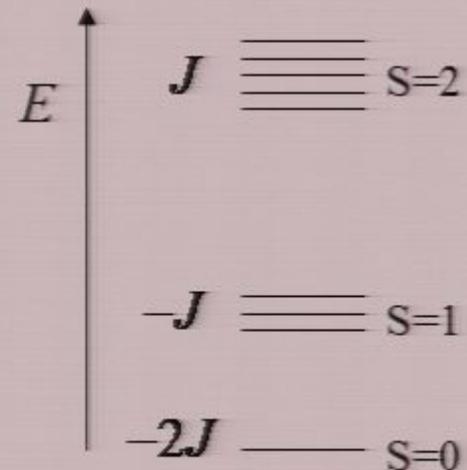
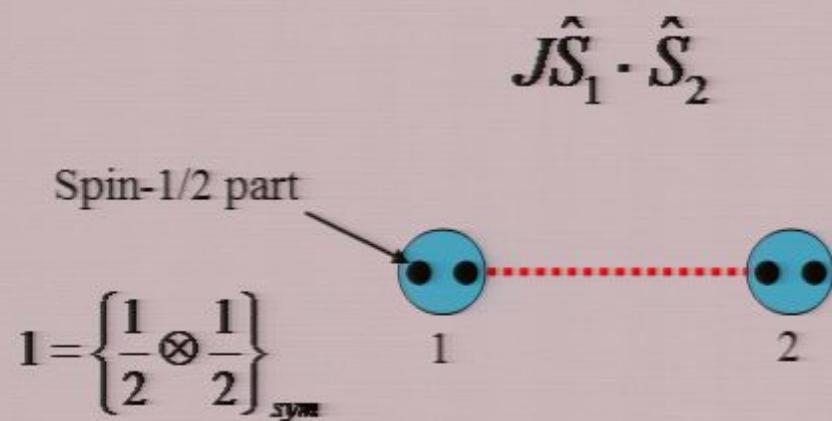
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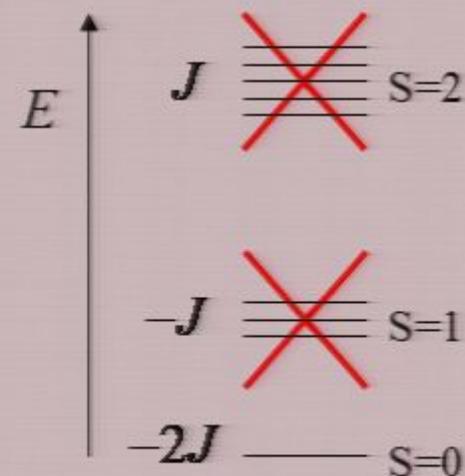
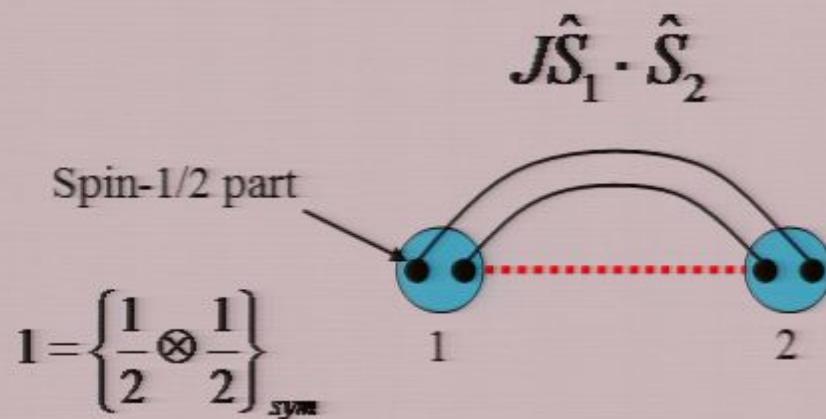
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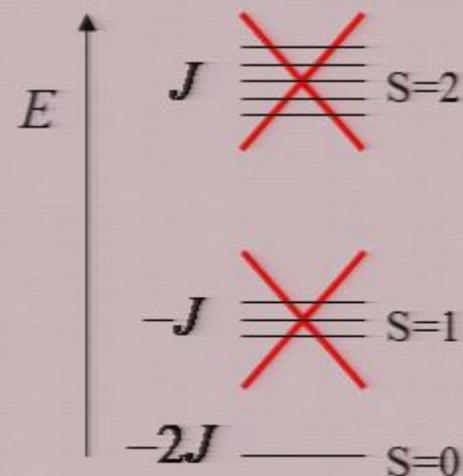
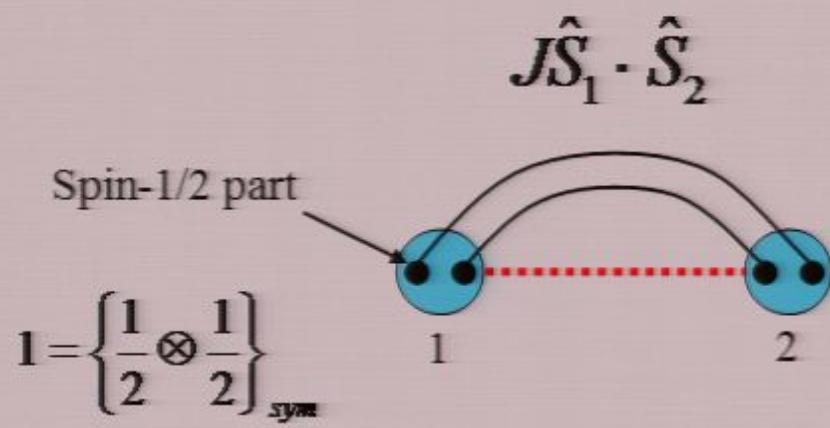
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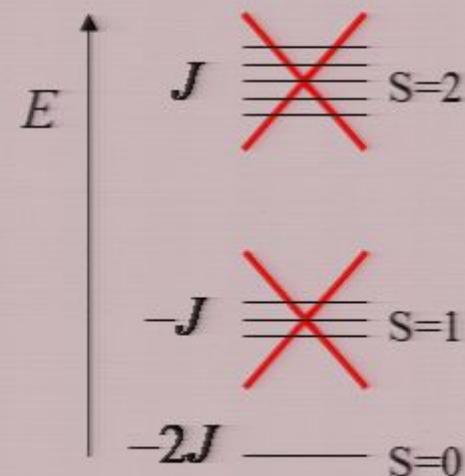
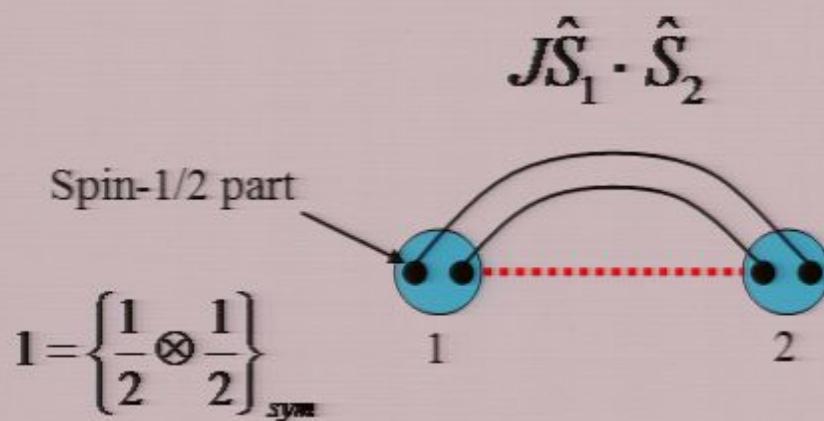
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- Weak randomness: Valence Bond Solid



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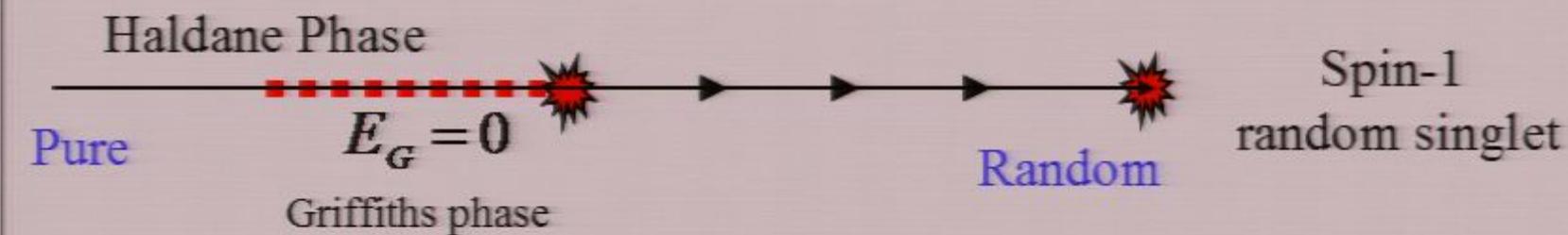
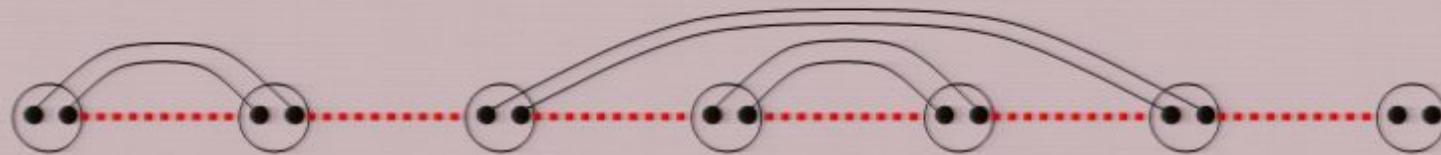


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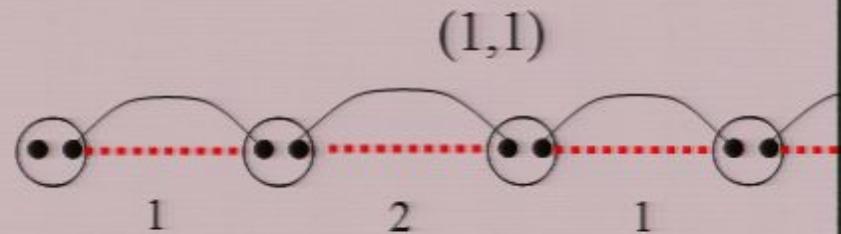
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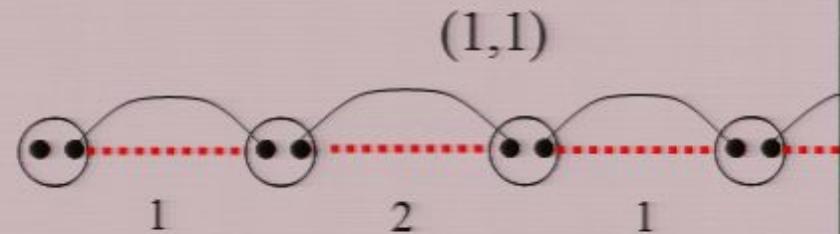
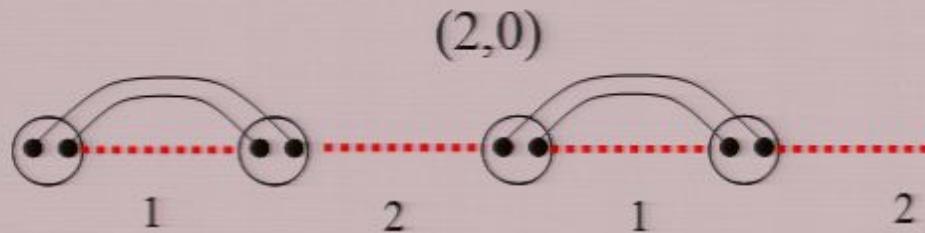
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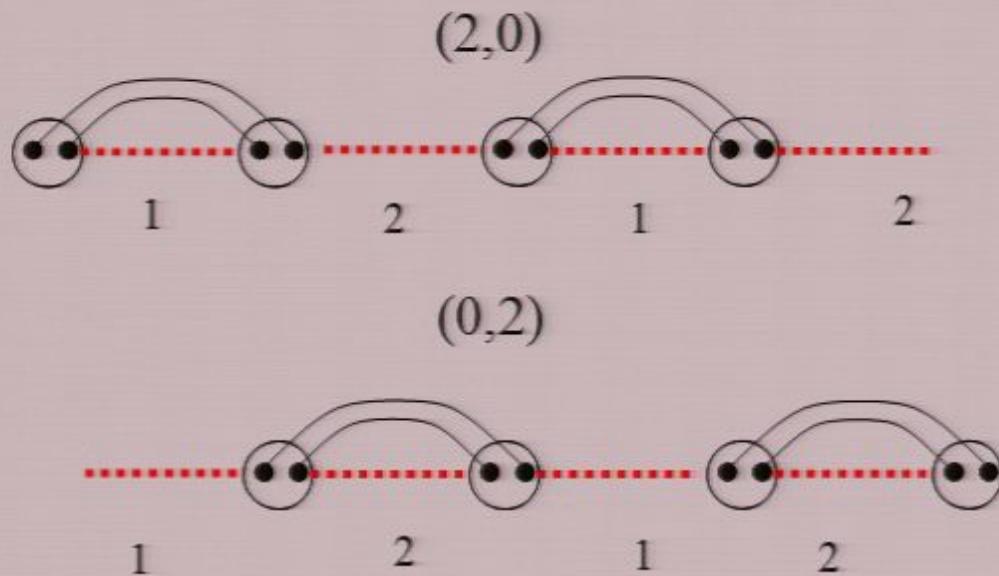
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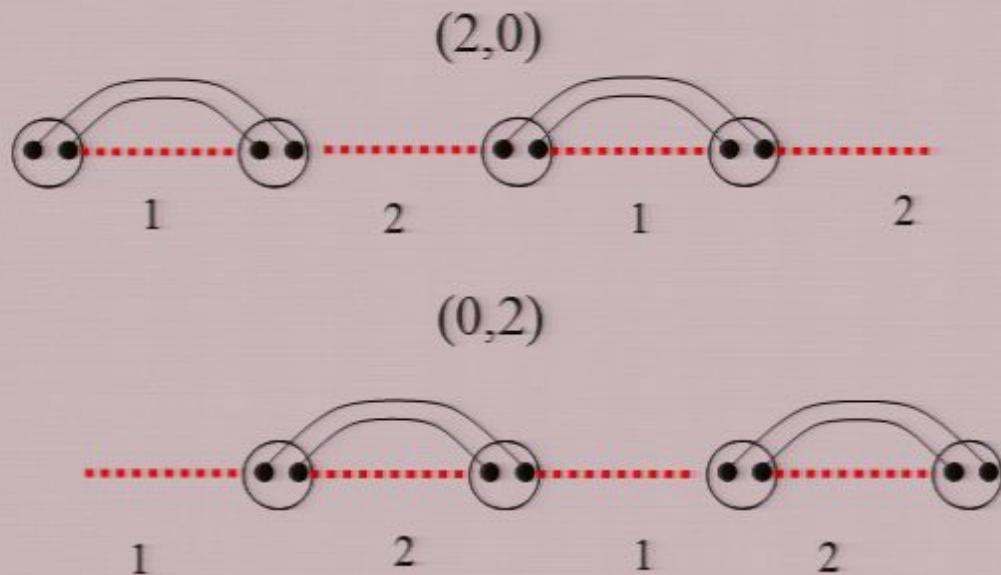
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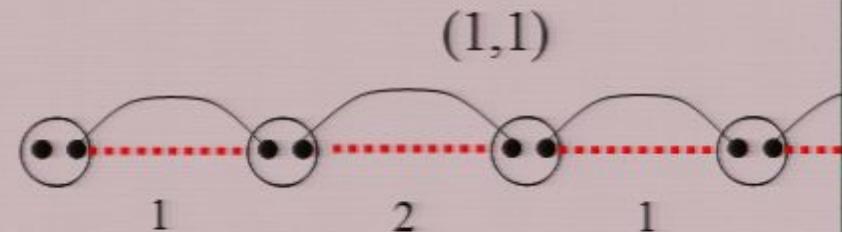
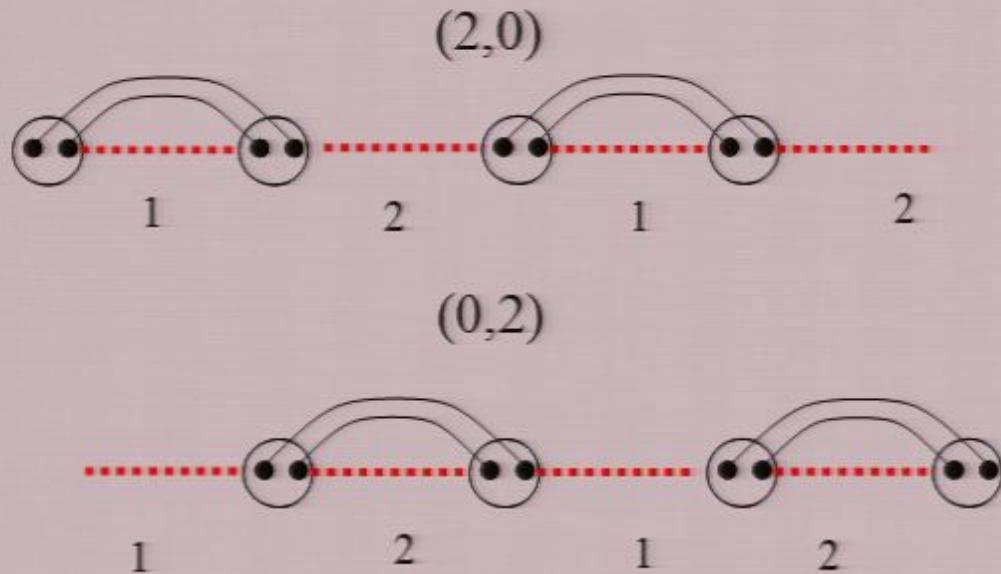


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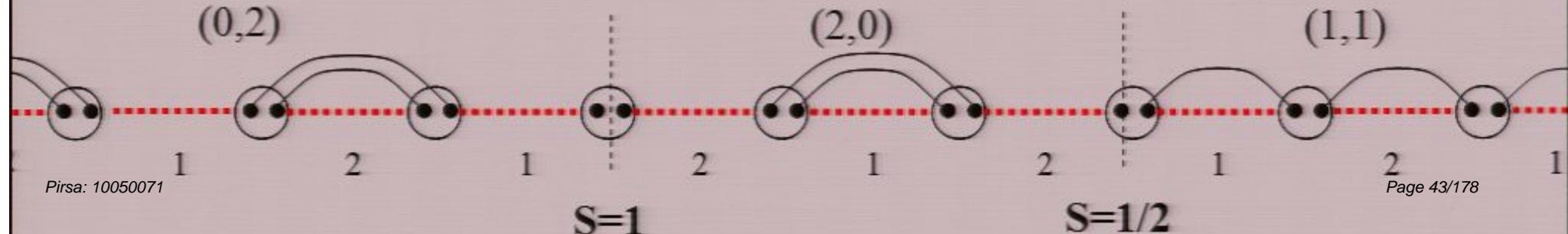
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- In general, spin S :

$$\psi = 1/(2S+1)$$

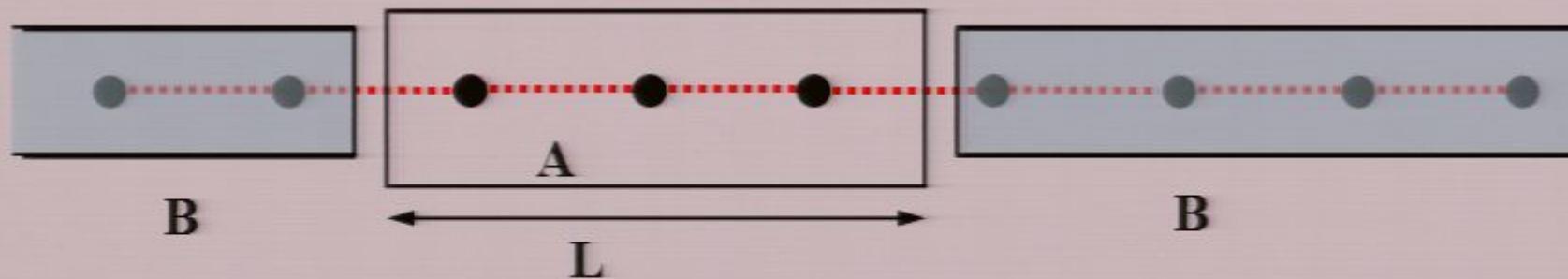
$$\chi = 2S$$

Entanglement entropy in Random Heisenberg chains

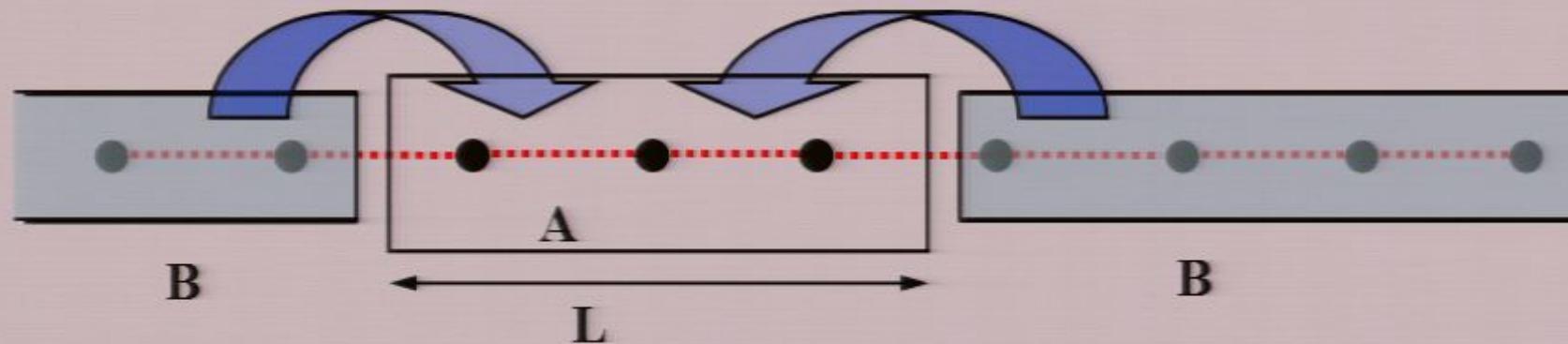
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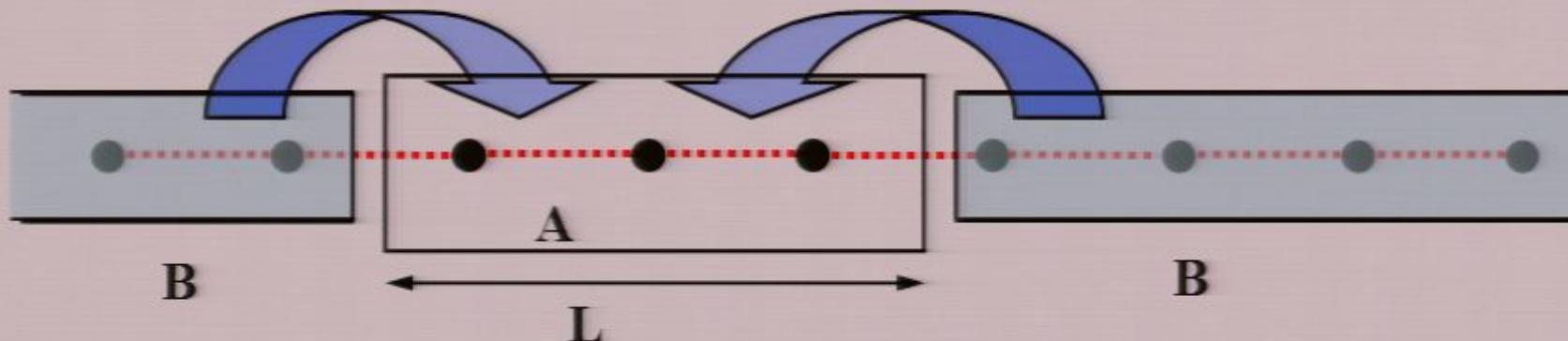
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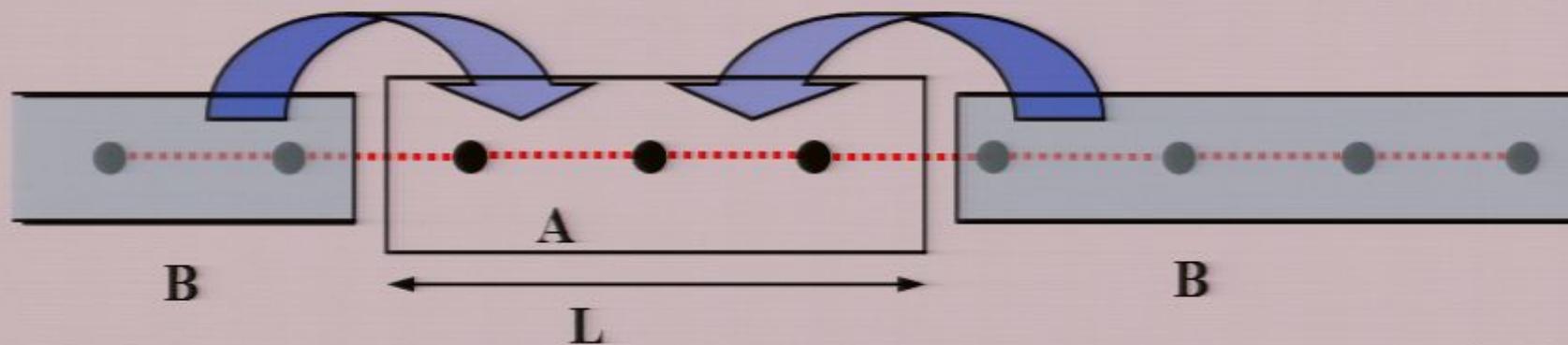
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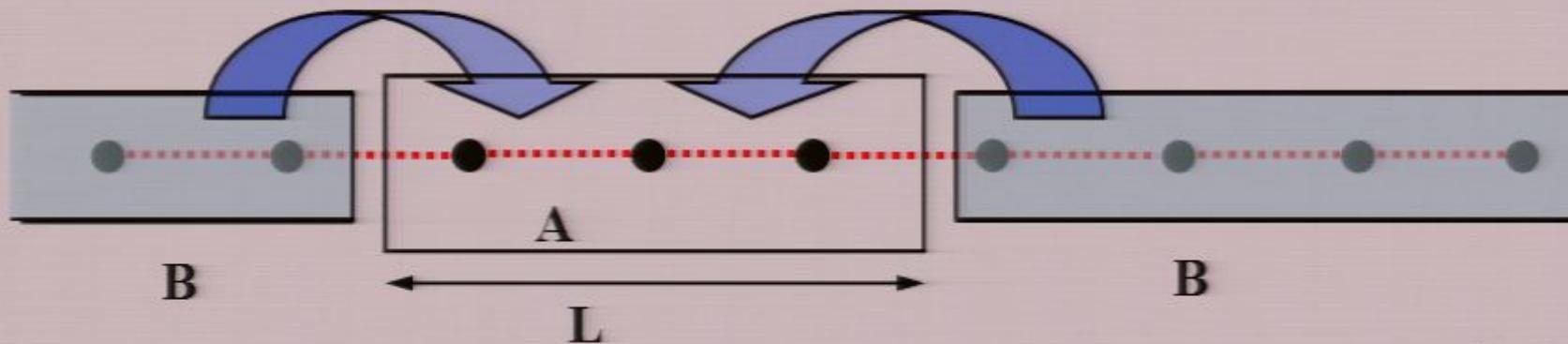


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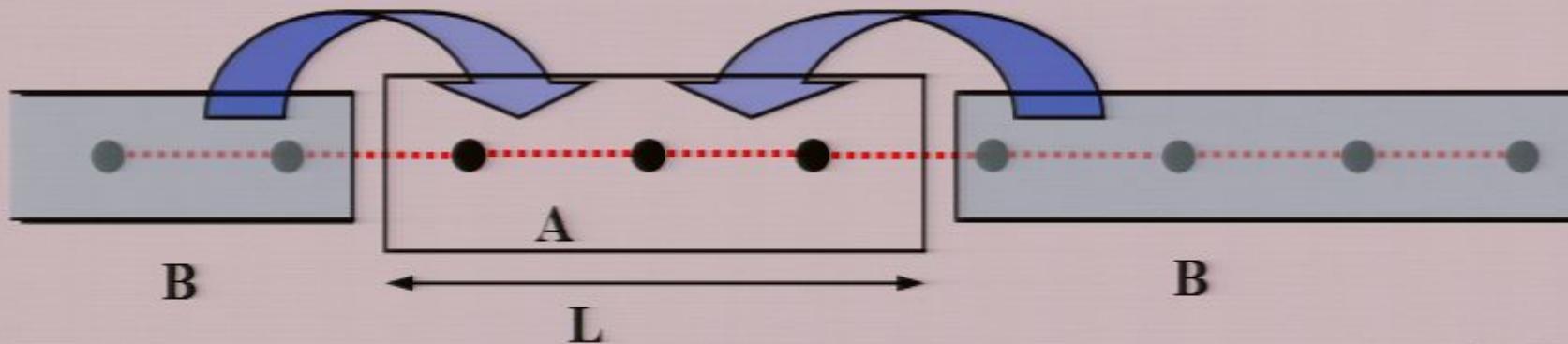
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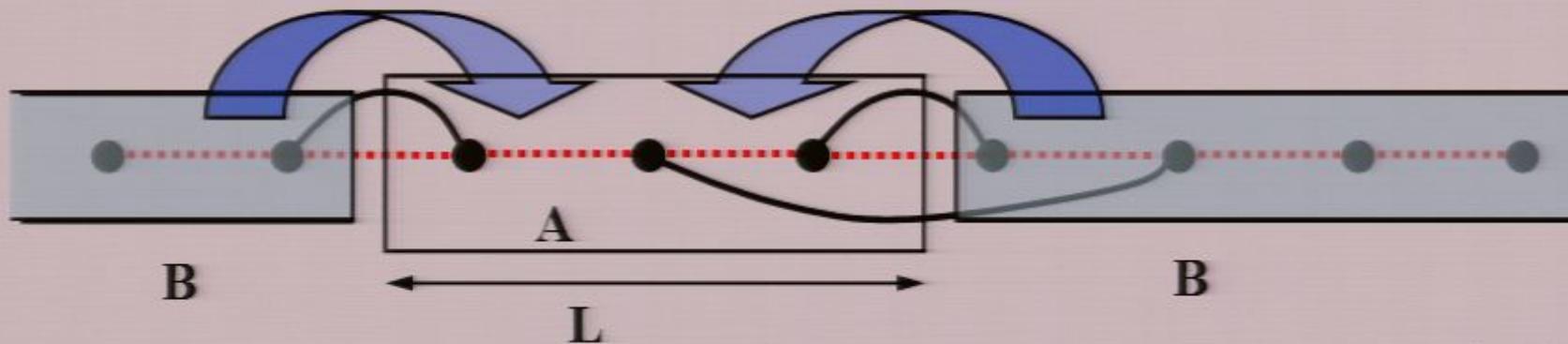
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Every singlet connecting A to B \rightarrow entanglement entropy 1.

Entanglement entropy in the Heisenberg model



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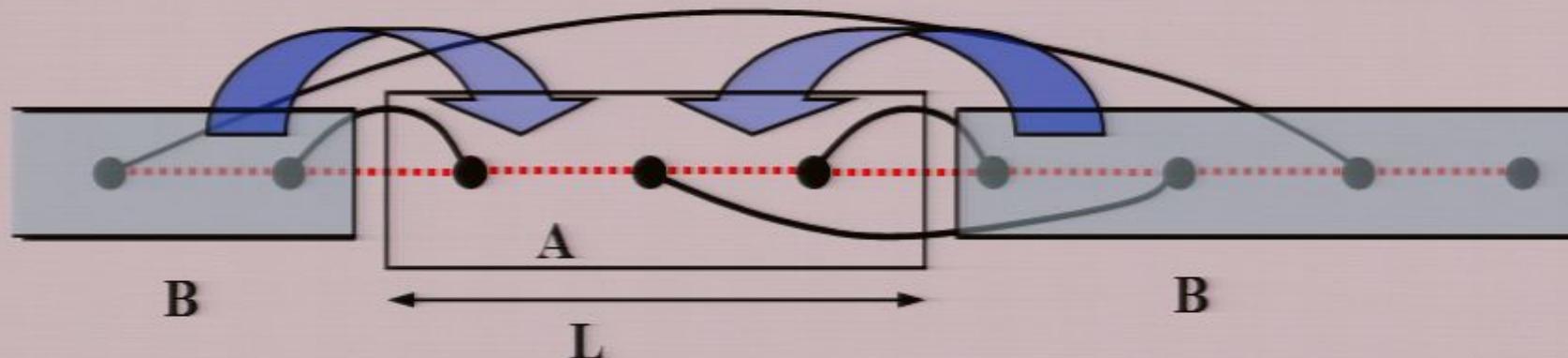
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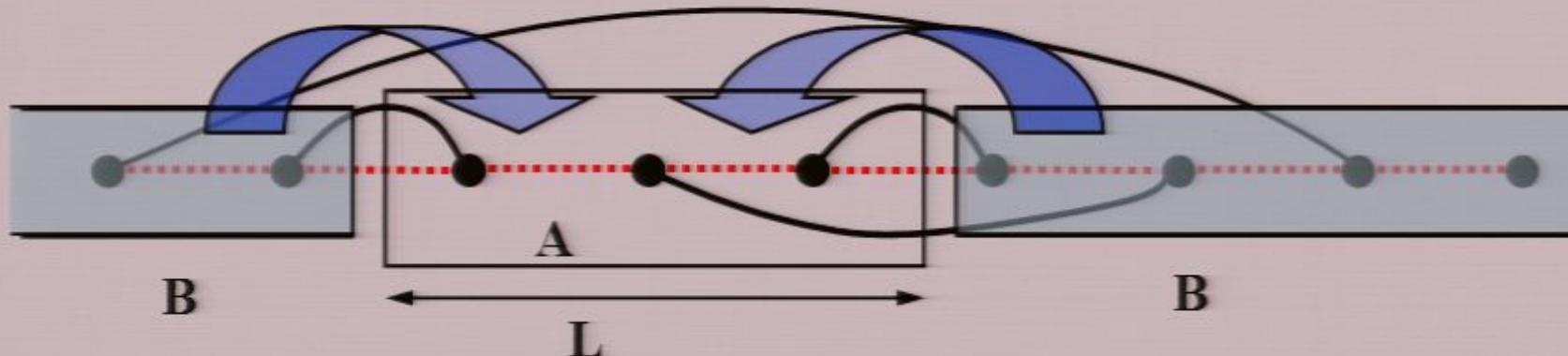
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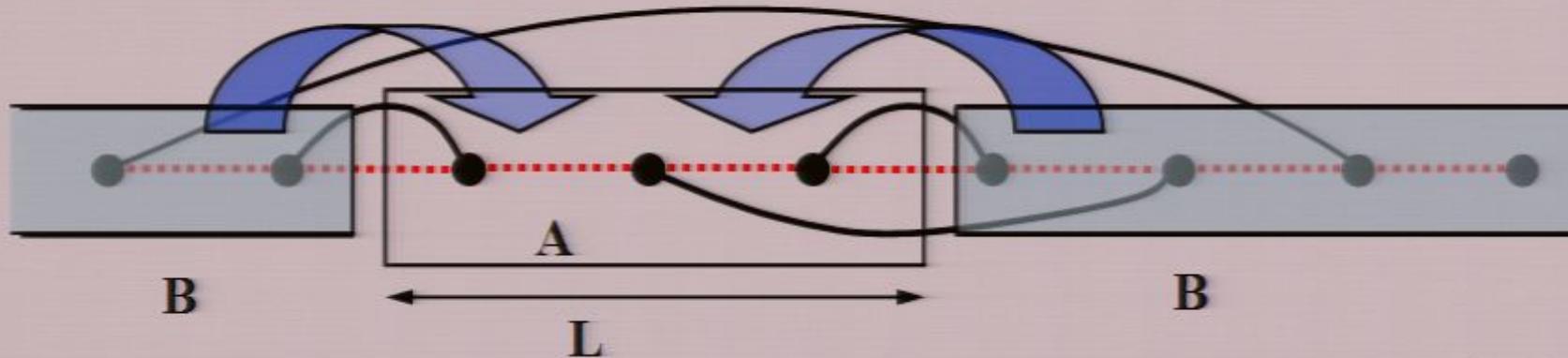
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Entanglement entropy in the Heisenberg model



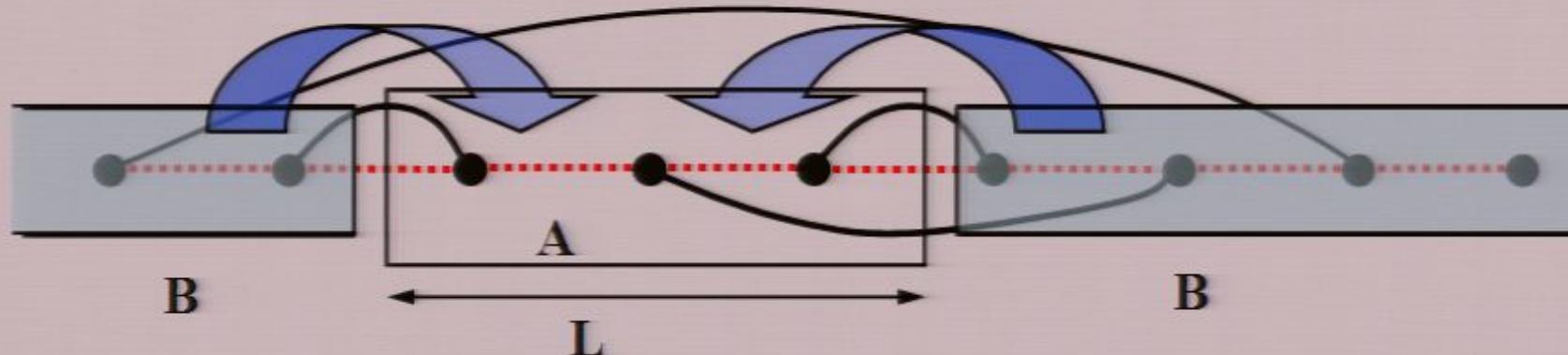
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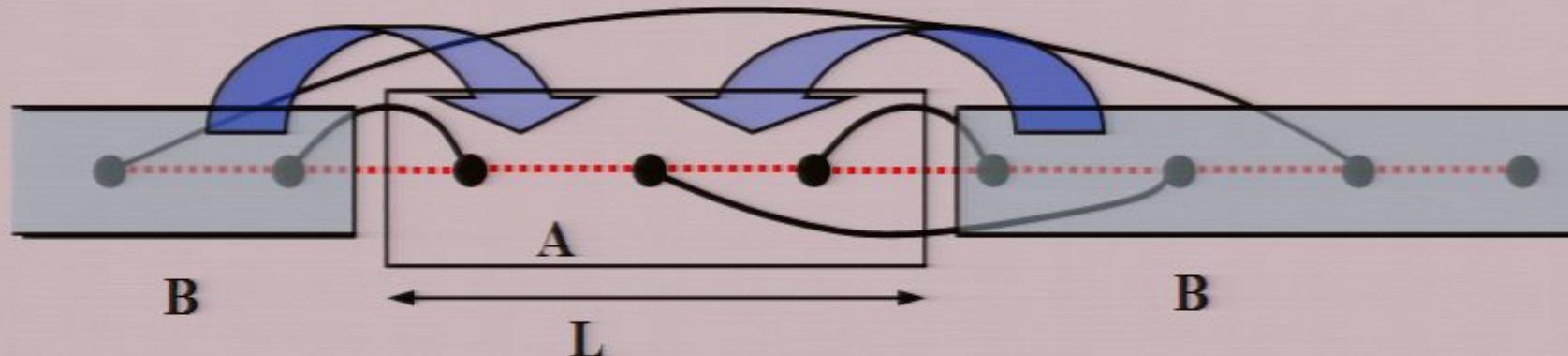
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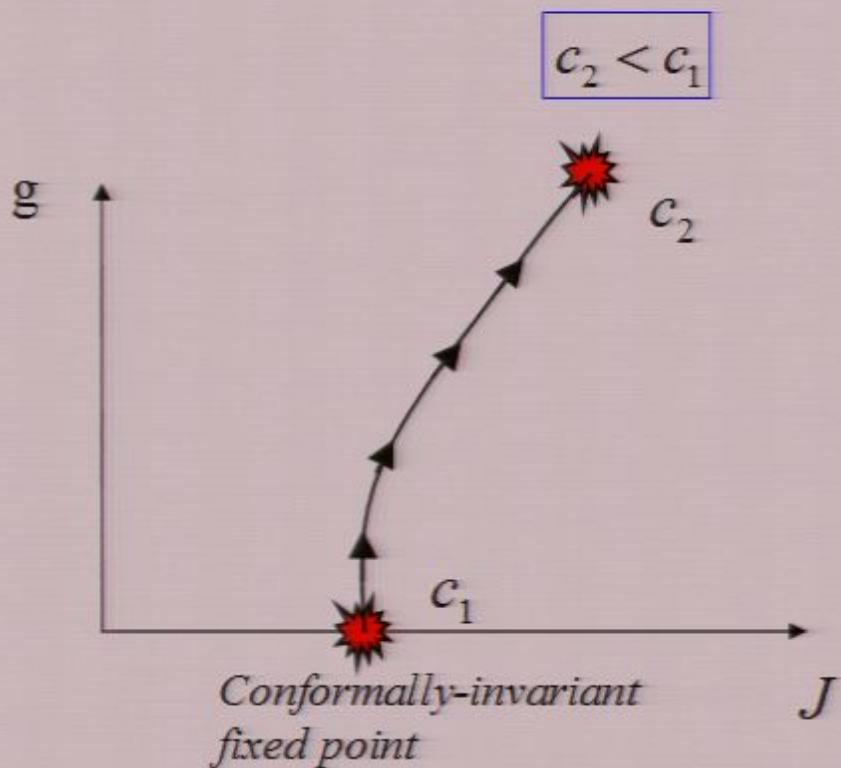
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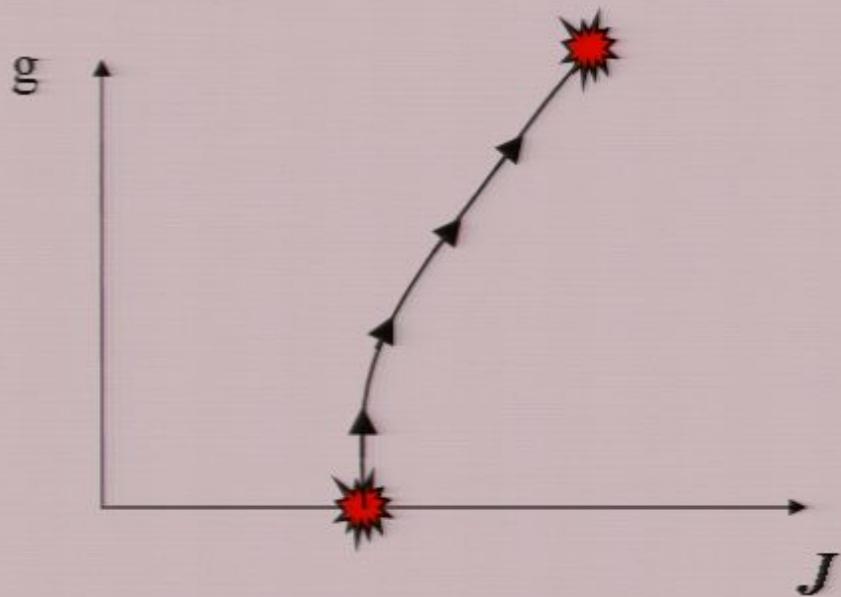
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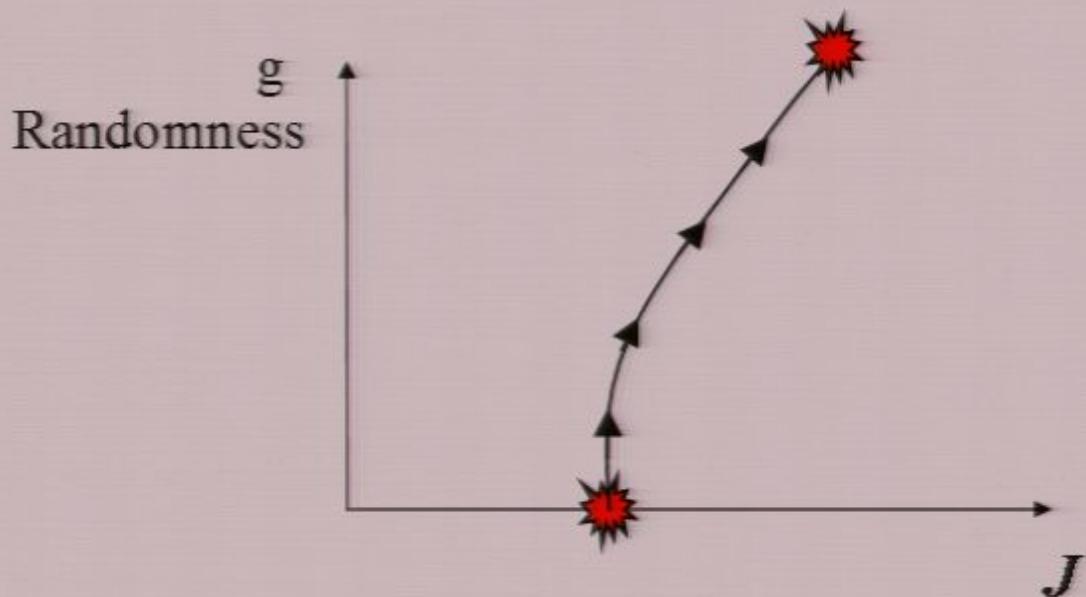
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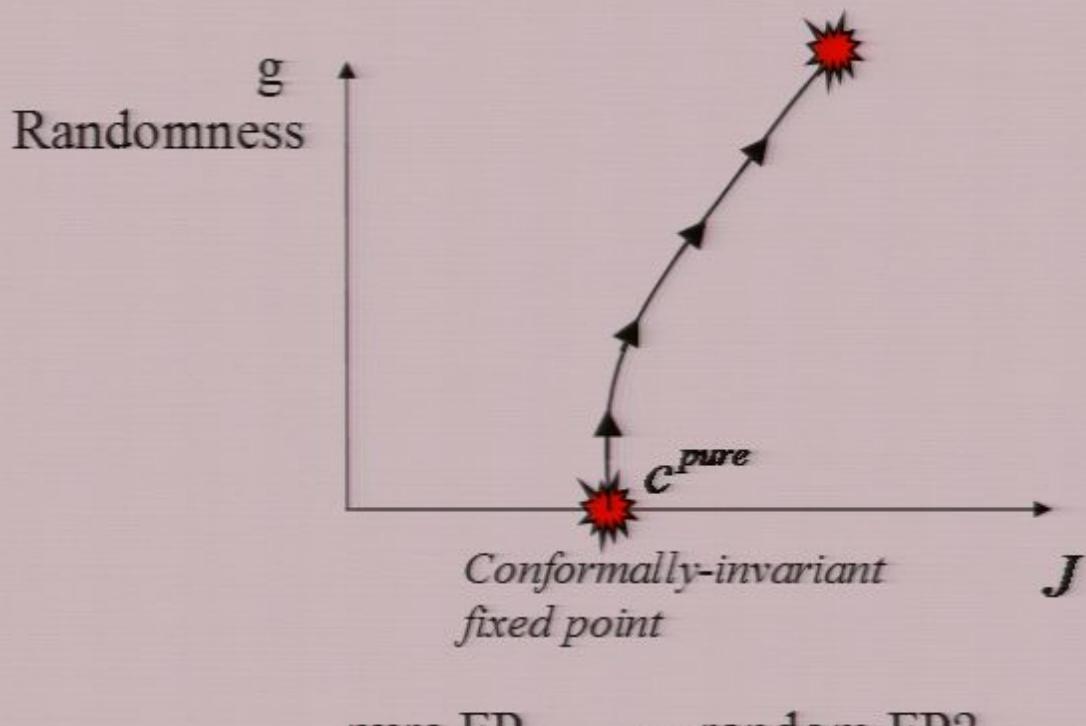
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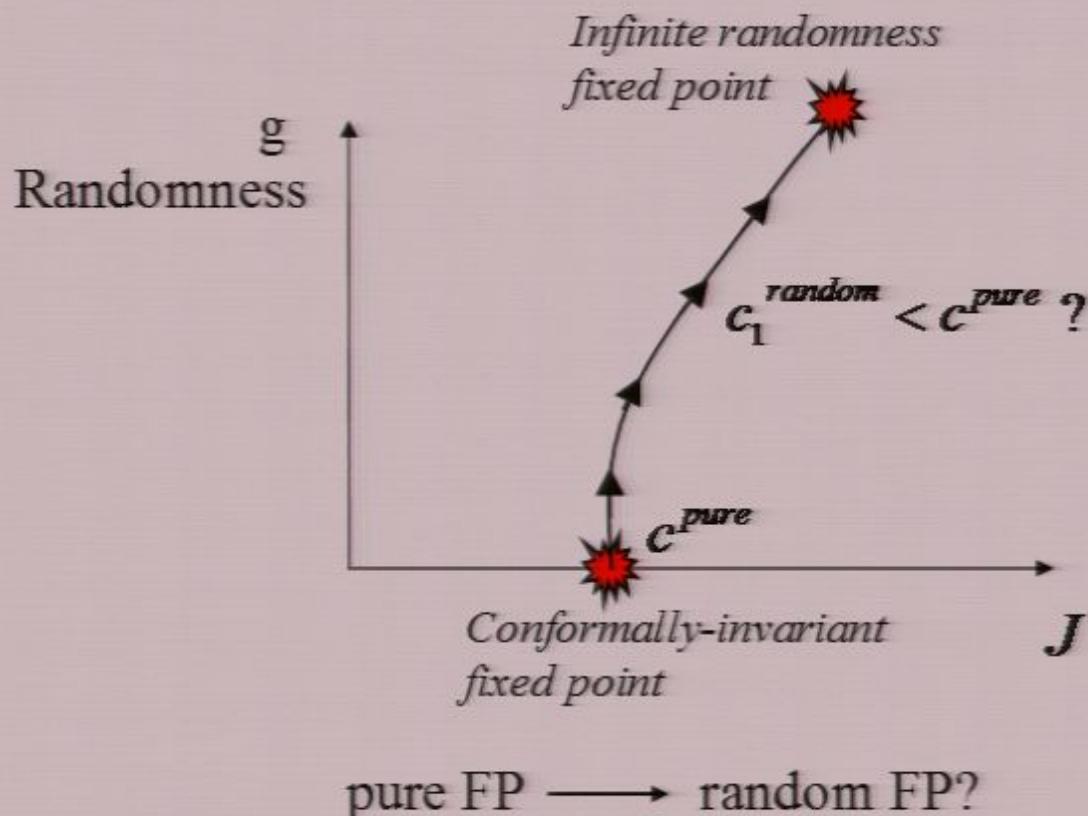
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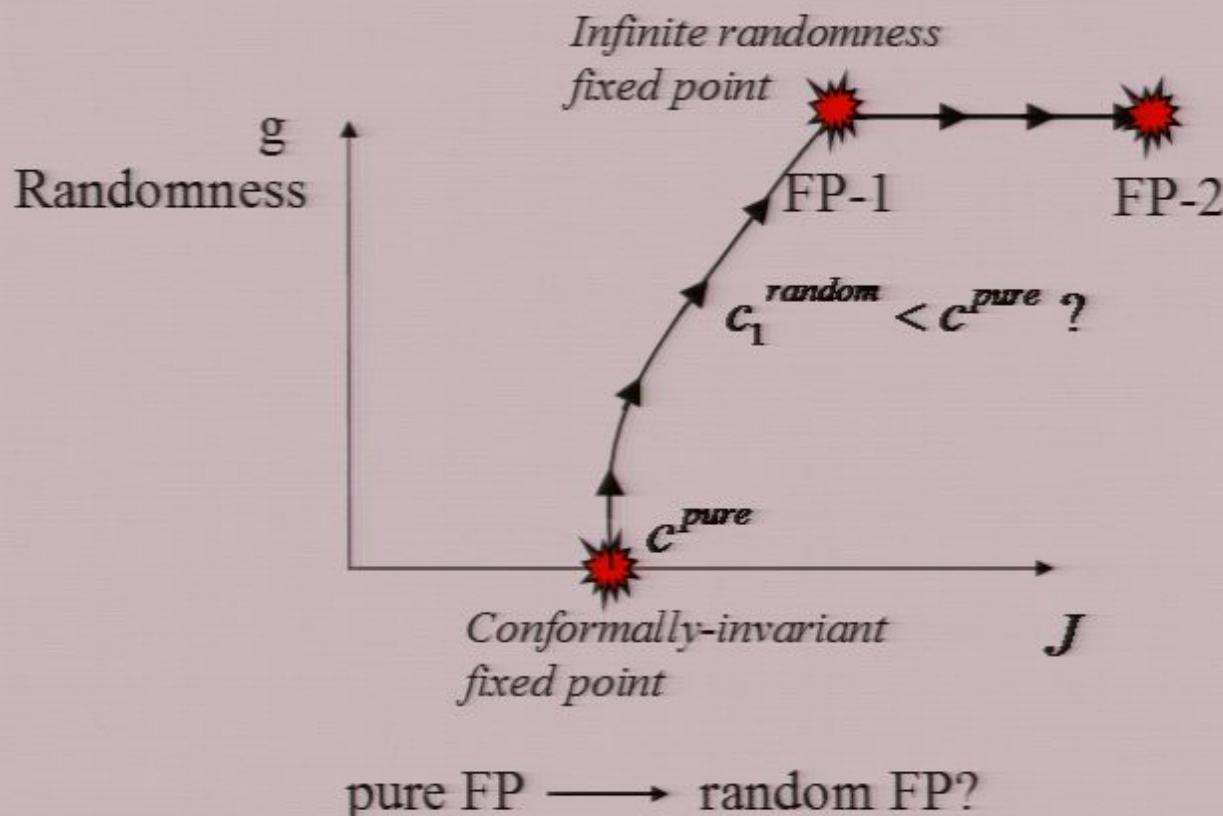
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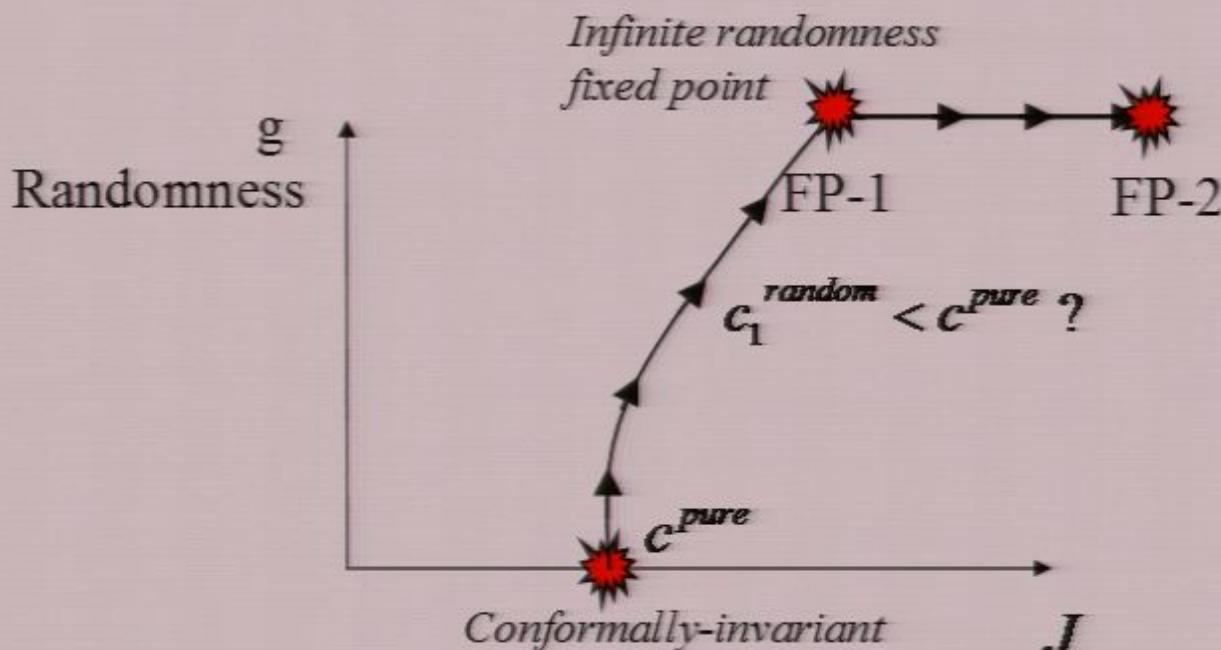
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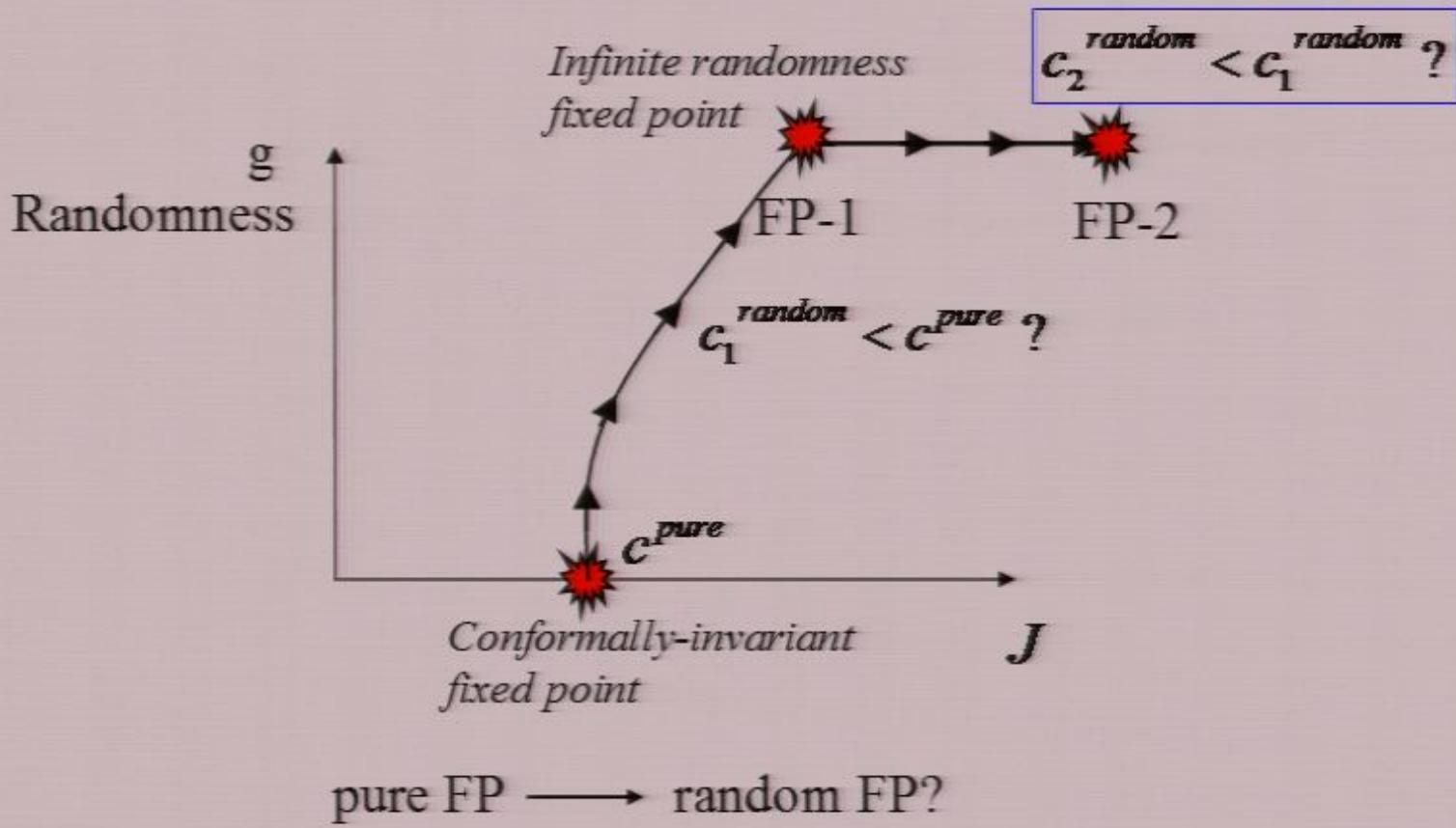


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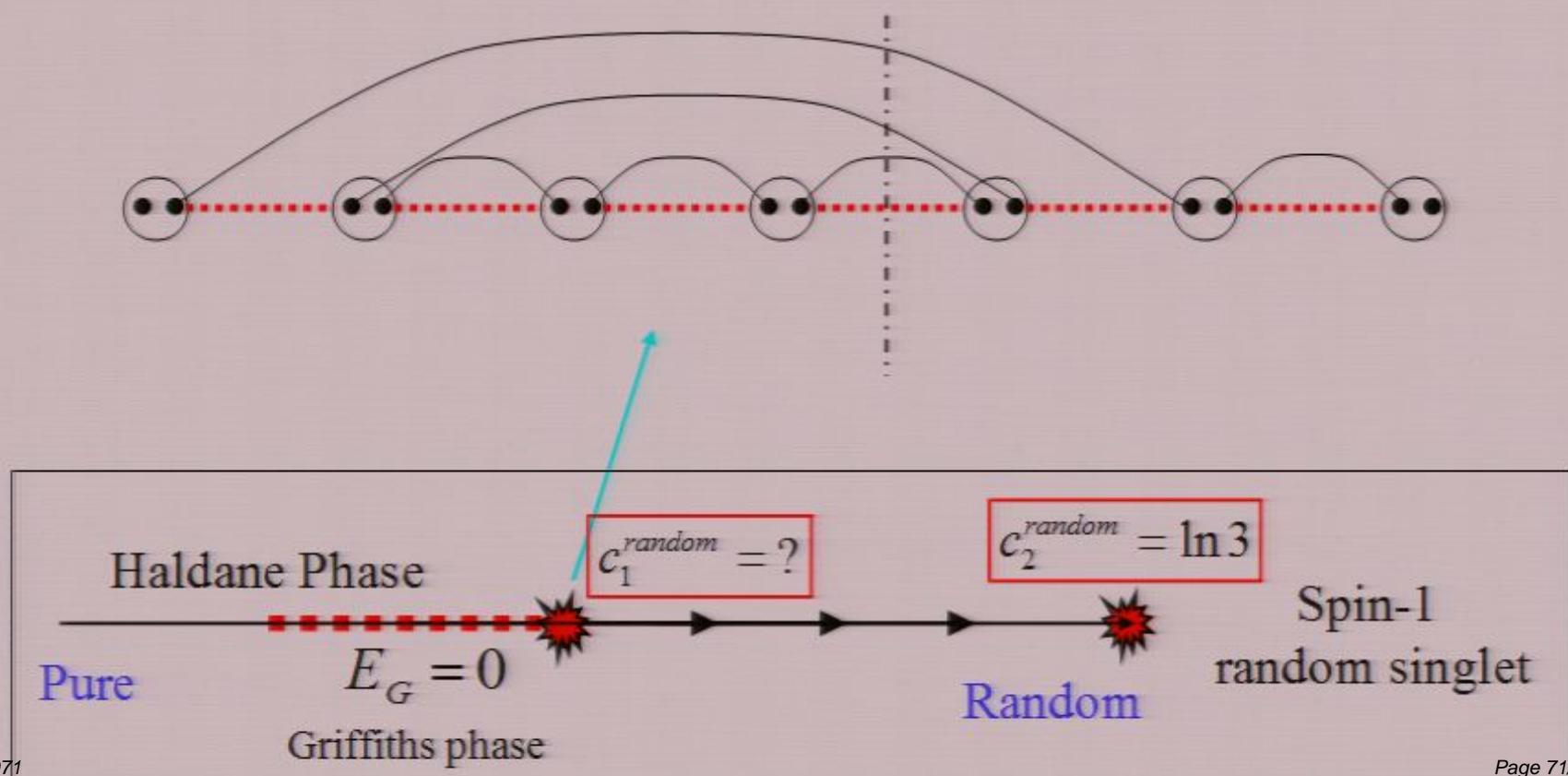
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Spin-1 “SU(2,2)” critical point

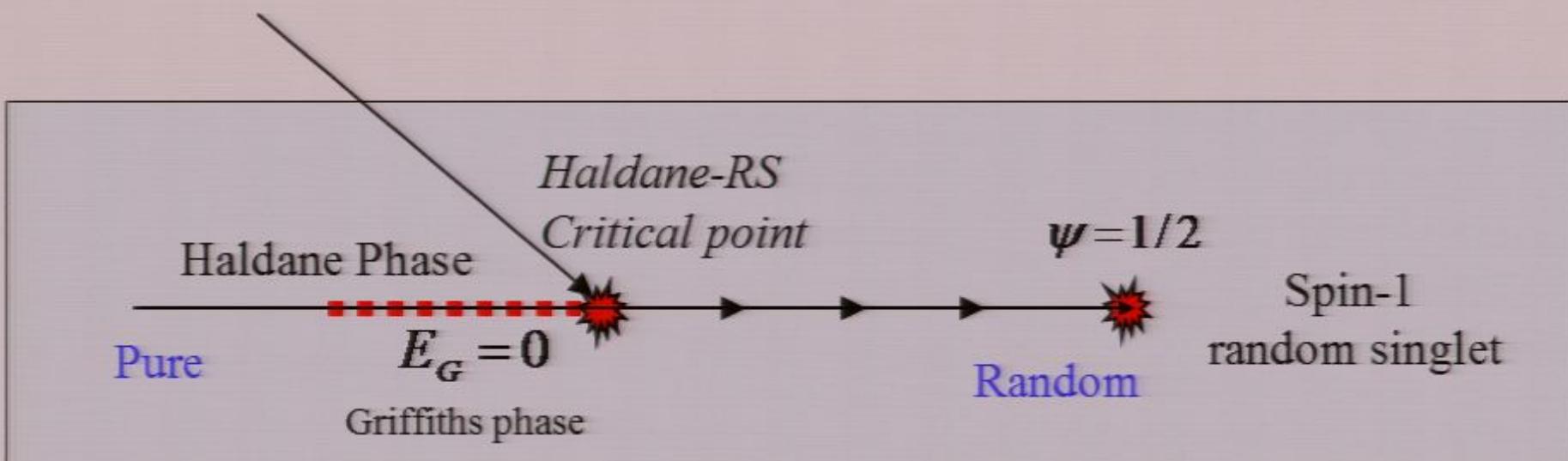
- Complex spin configurations:



Spin-1 Heisenberg results

- Lengthy calculation gives:

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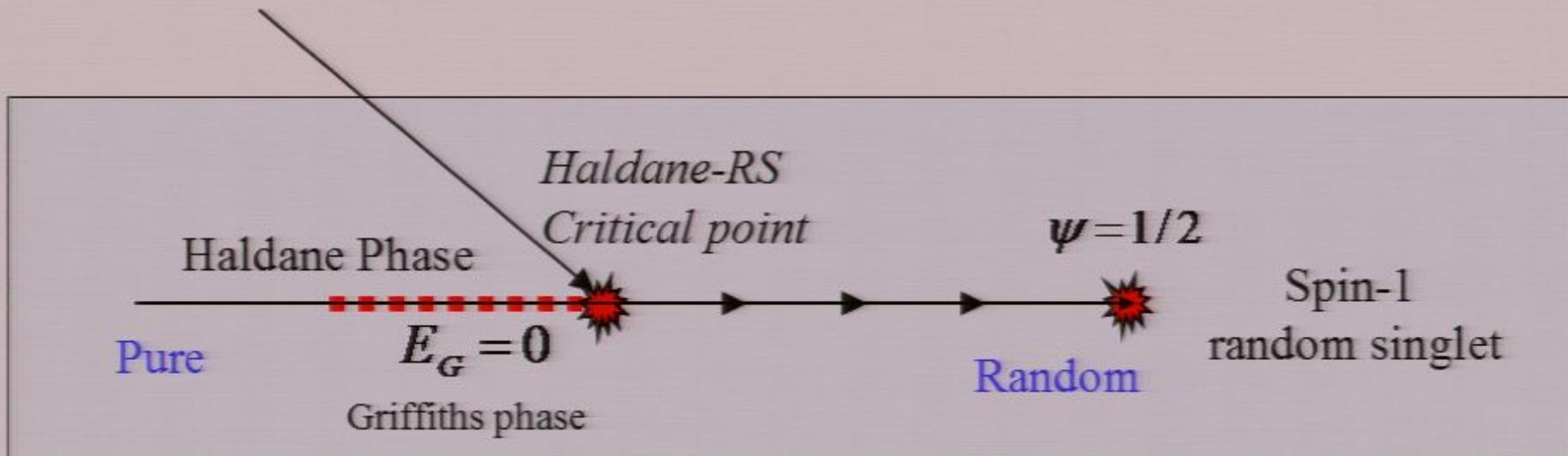


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$$E_{H-RS} \approx \frac{1}{3} \left(\frac{16}{9} \ln 2 \right) \log_2 L$$

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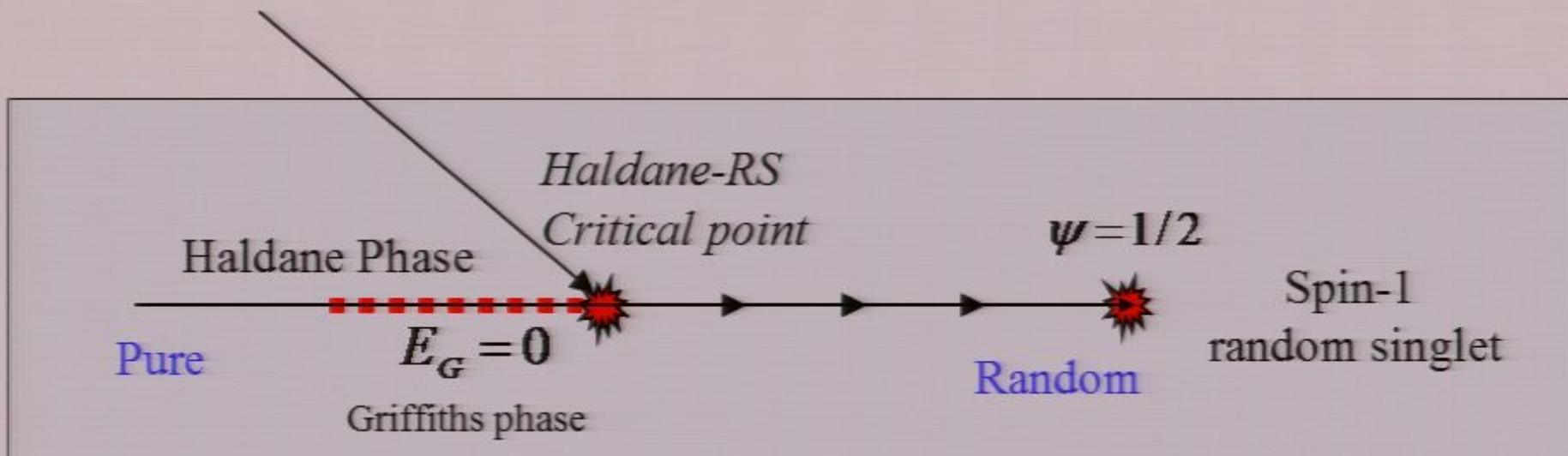


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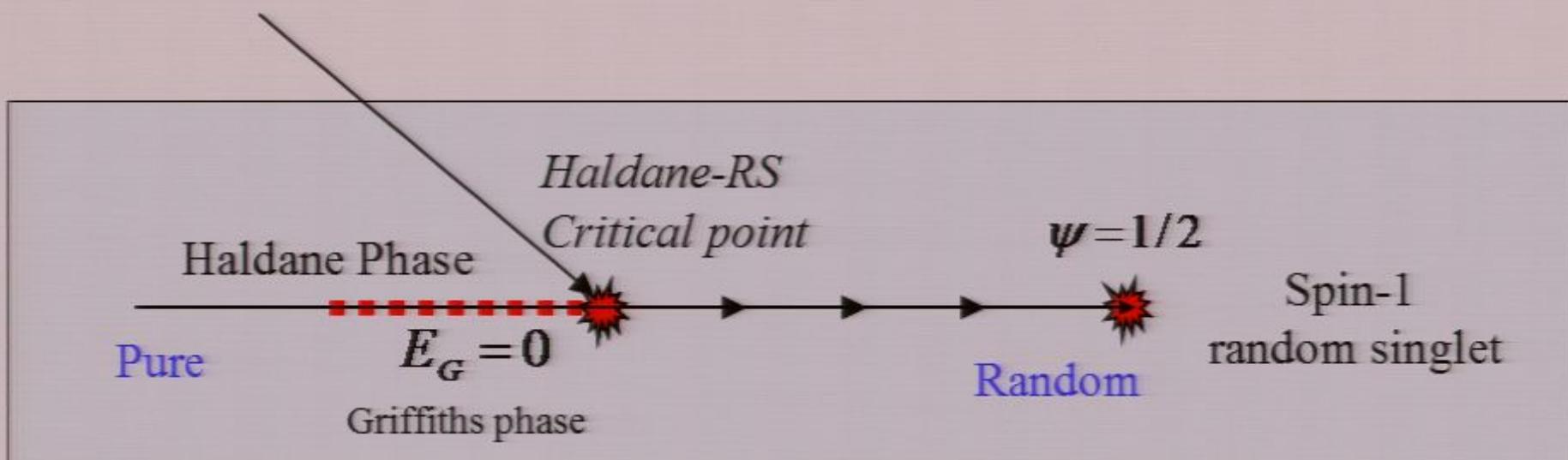
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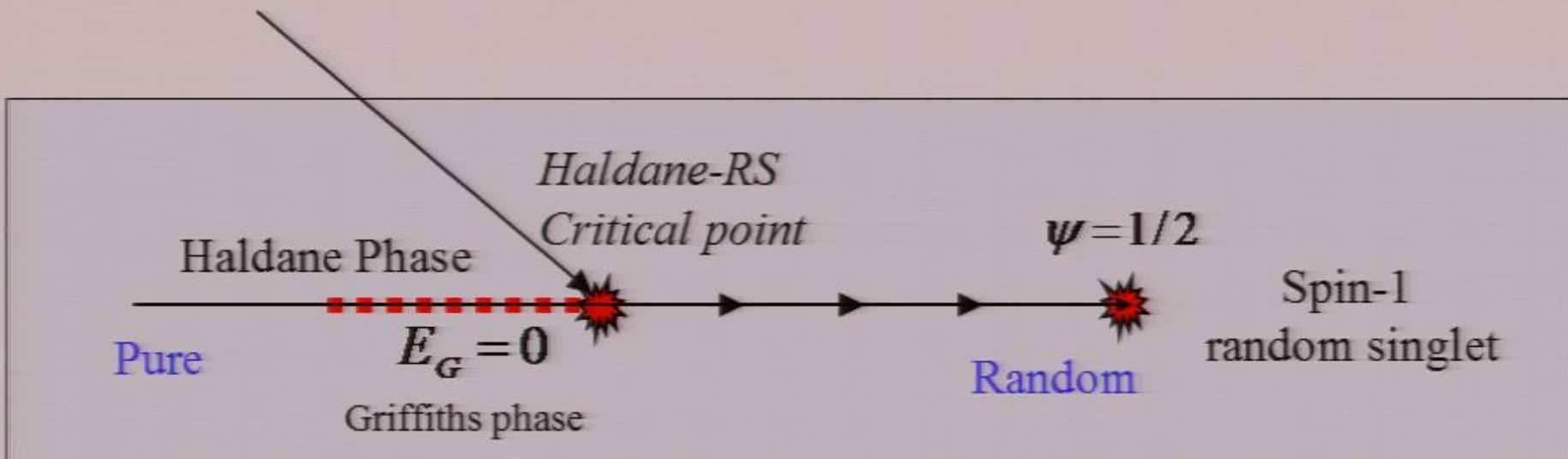
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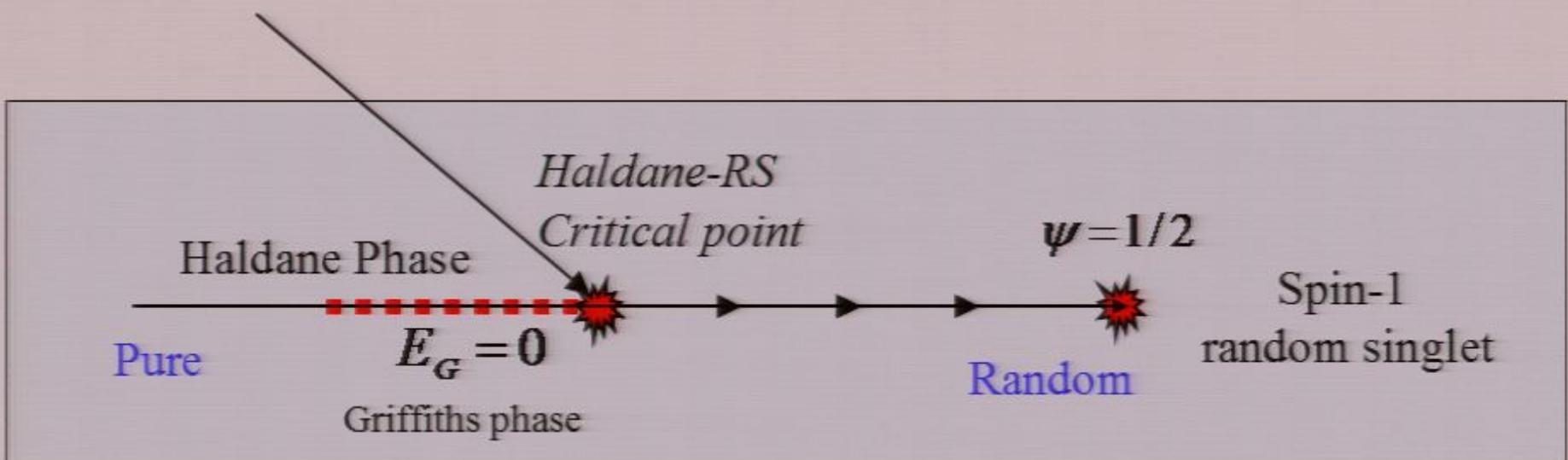
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Random Non-abelian Spin Chains

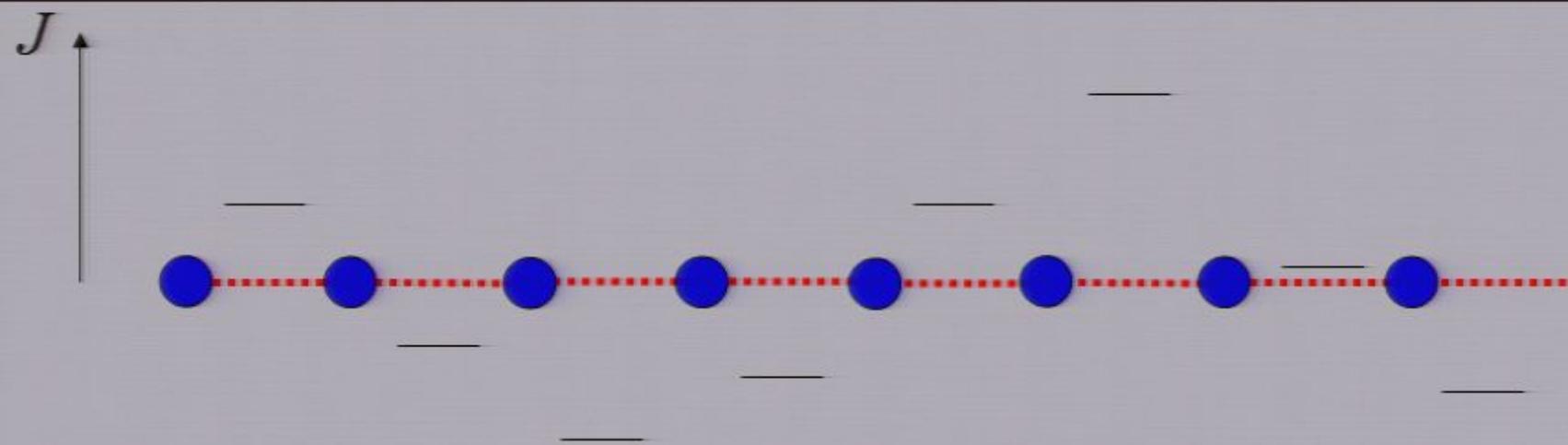
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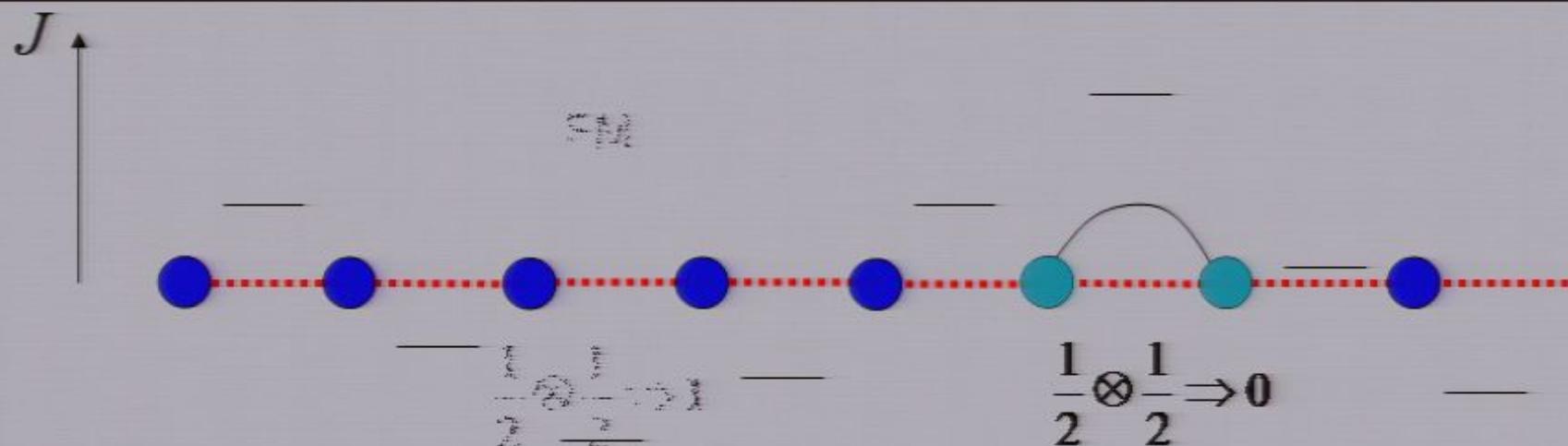
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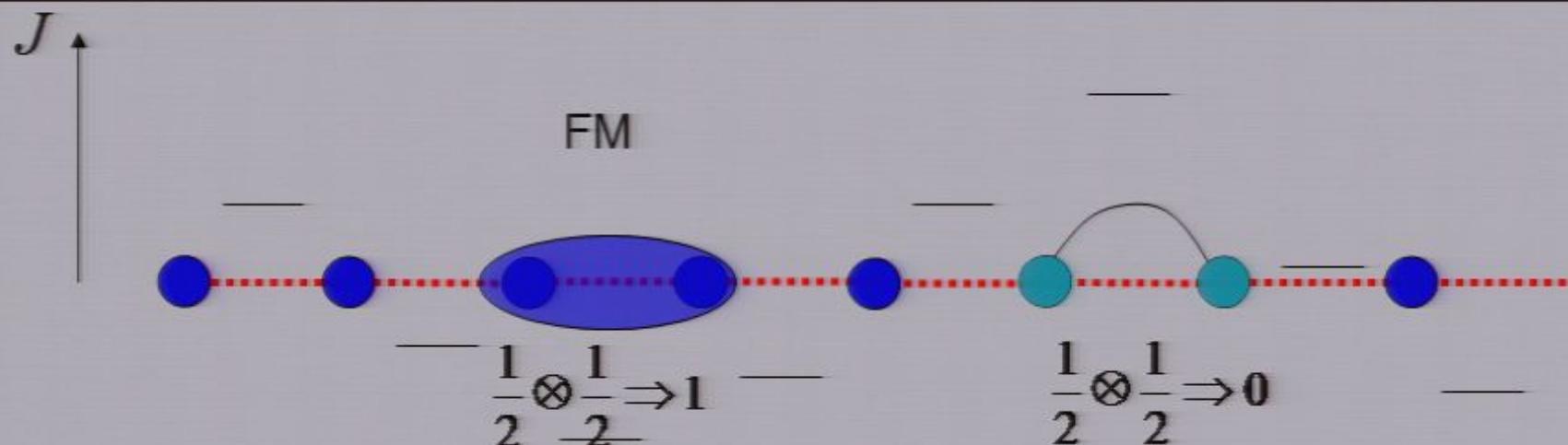
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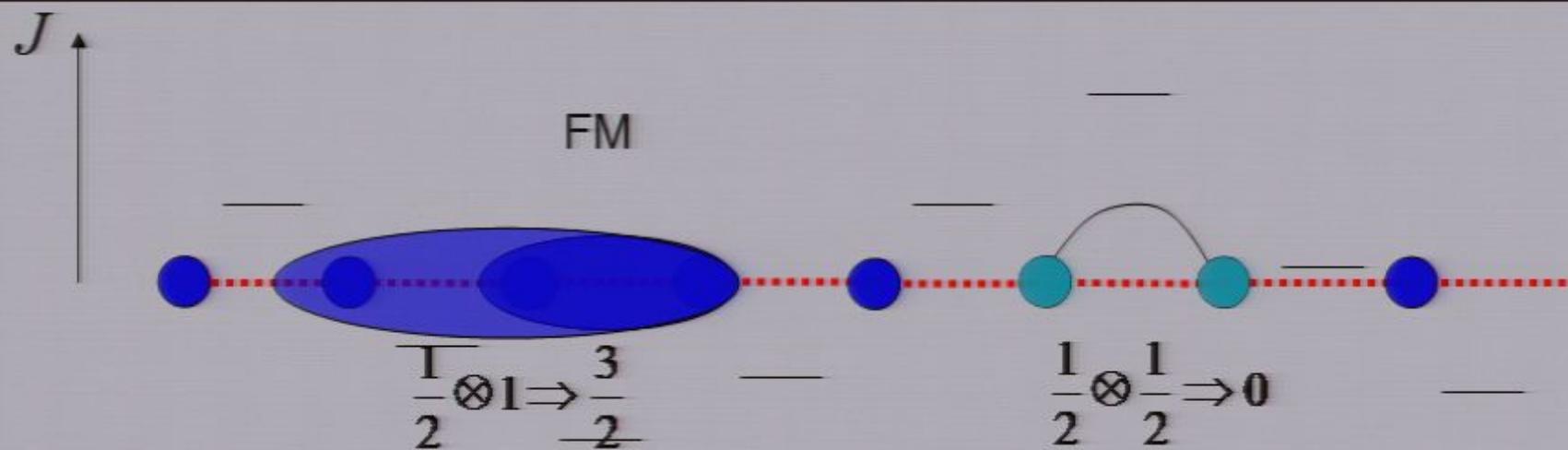
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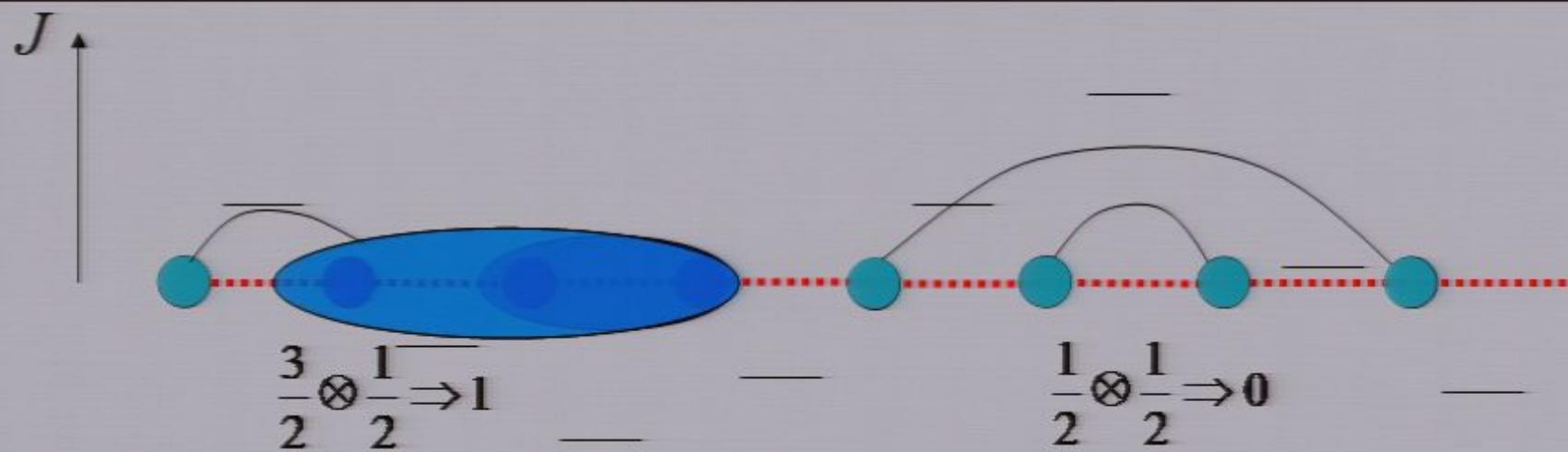
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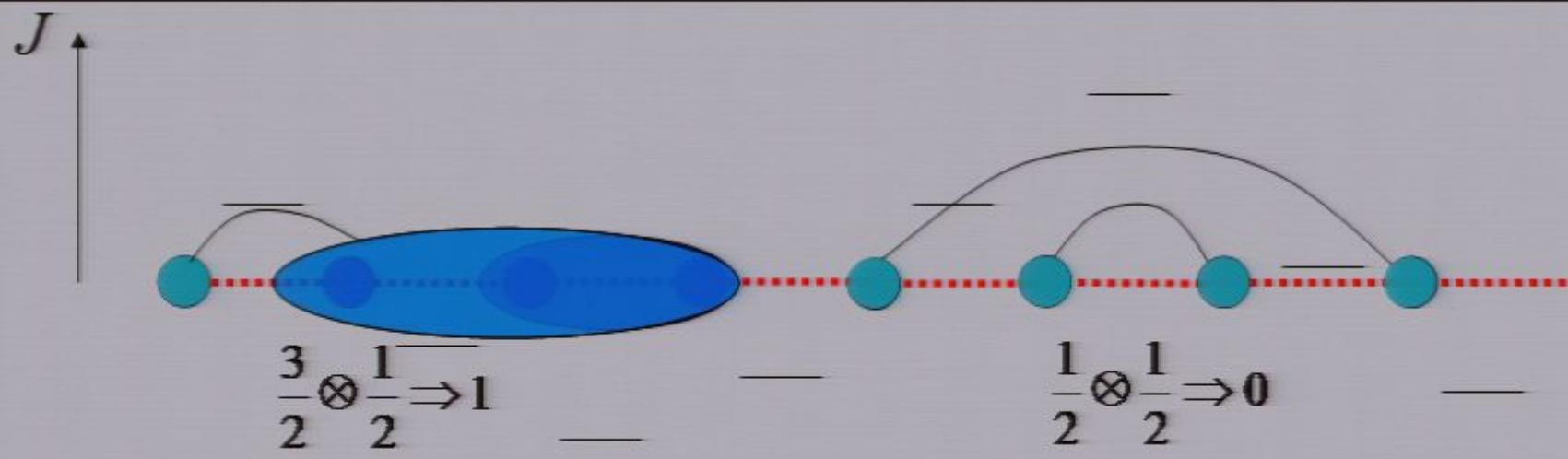
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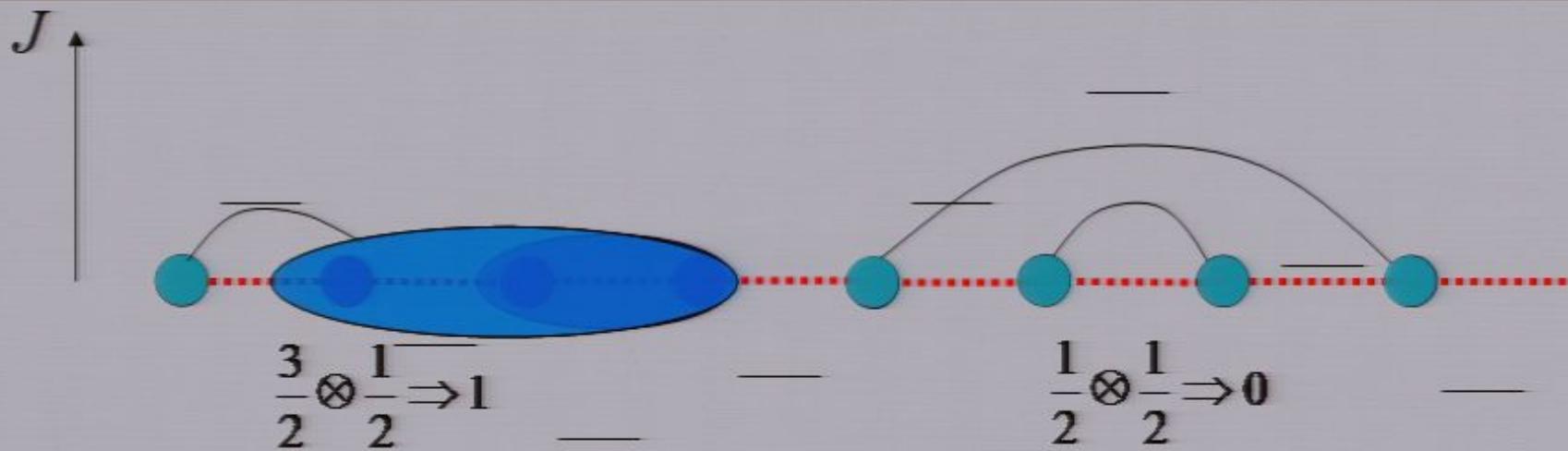
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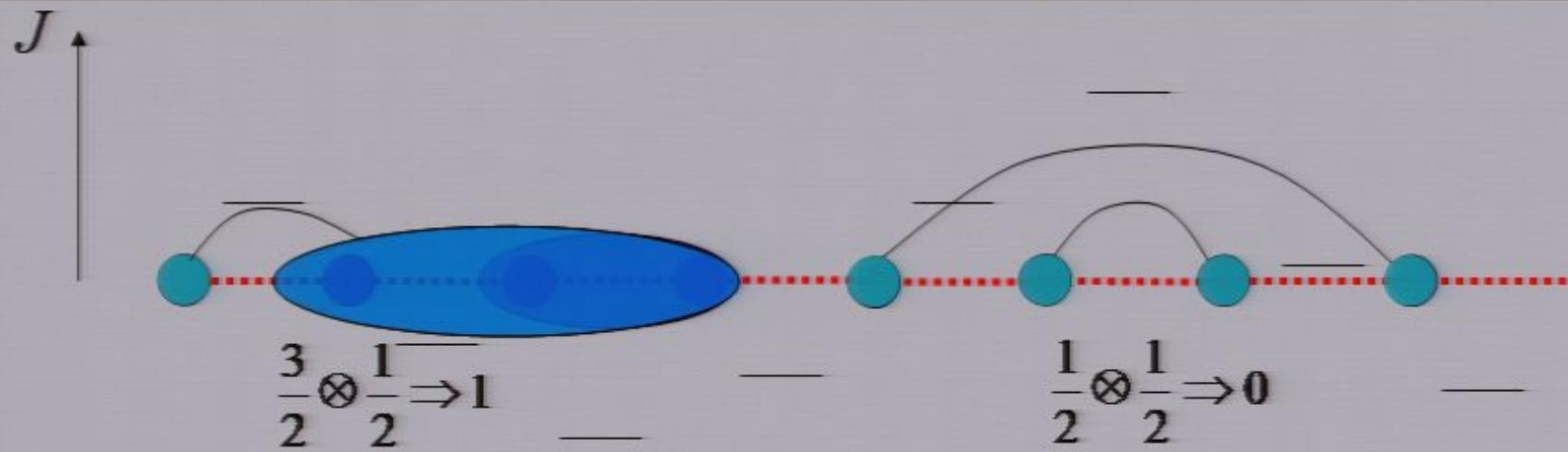
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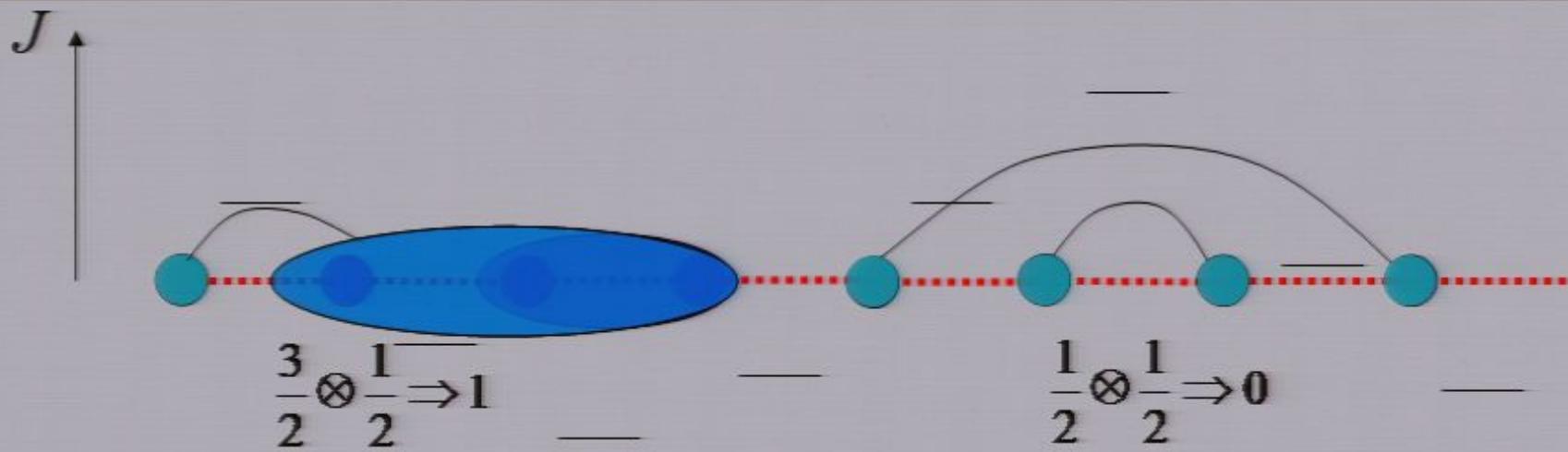


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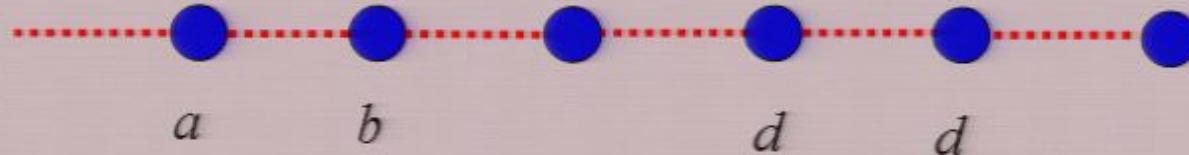
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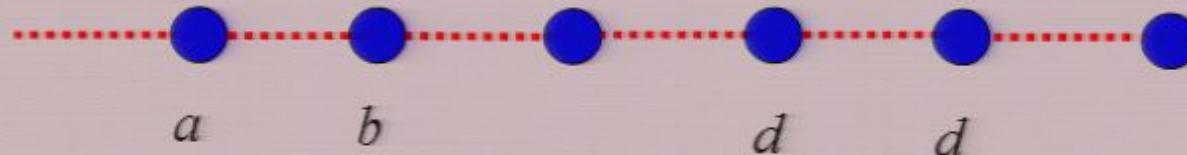


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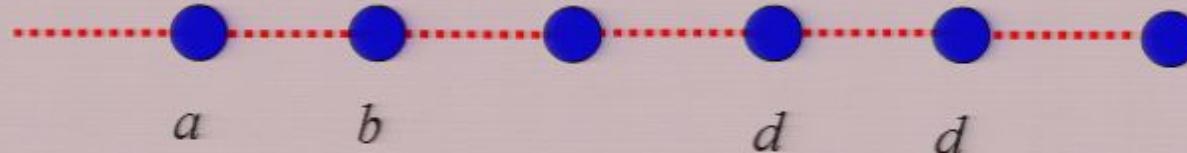
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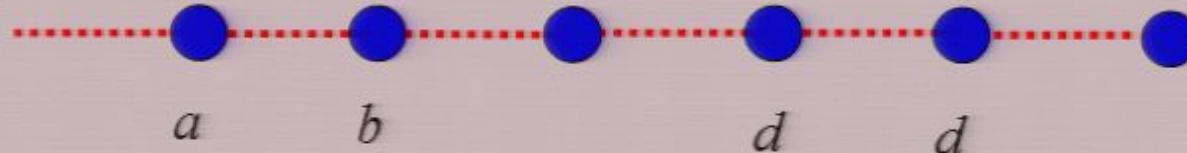
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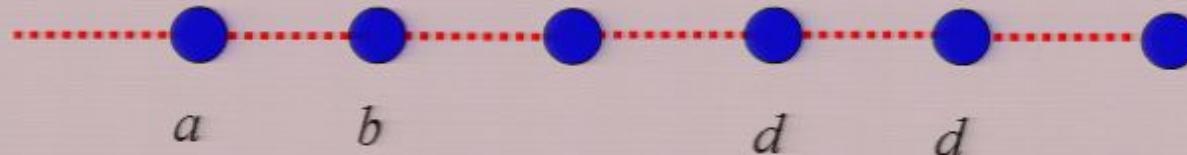
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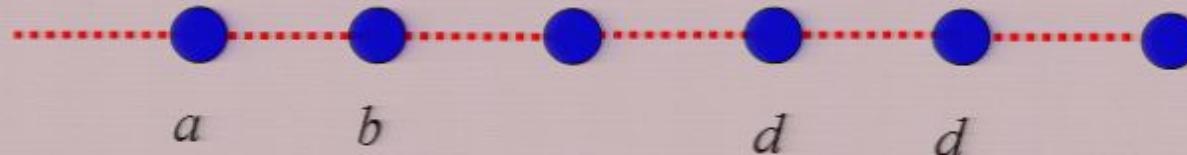
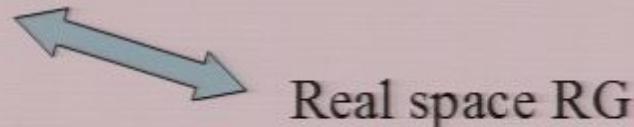
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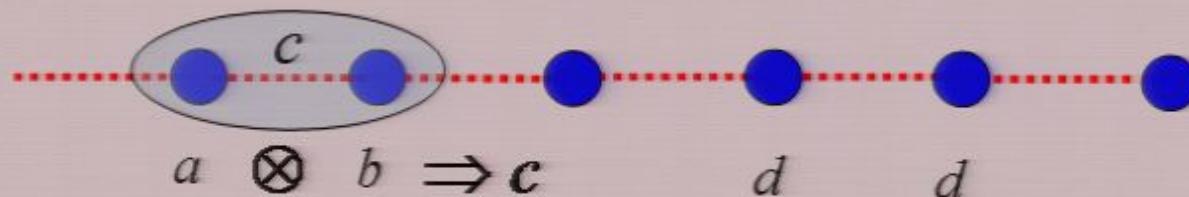
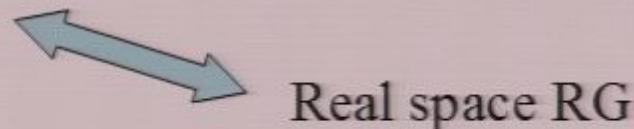
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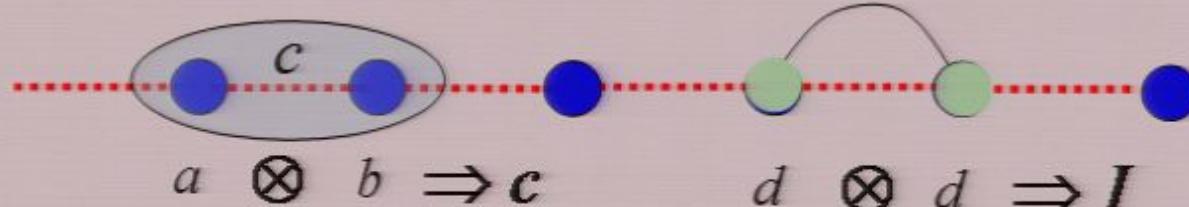
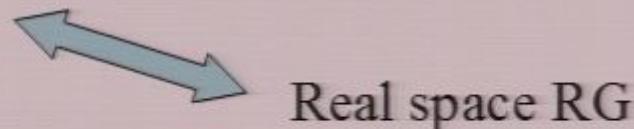
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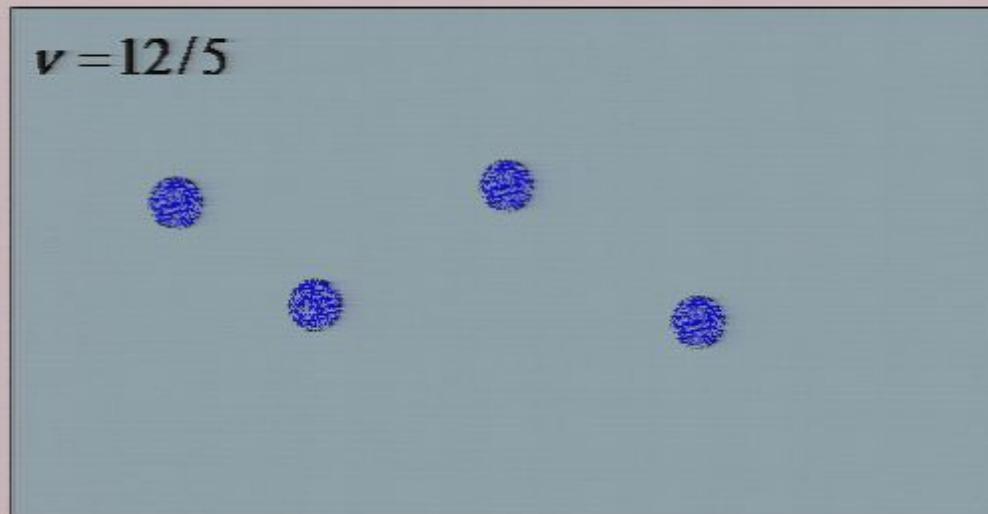
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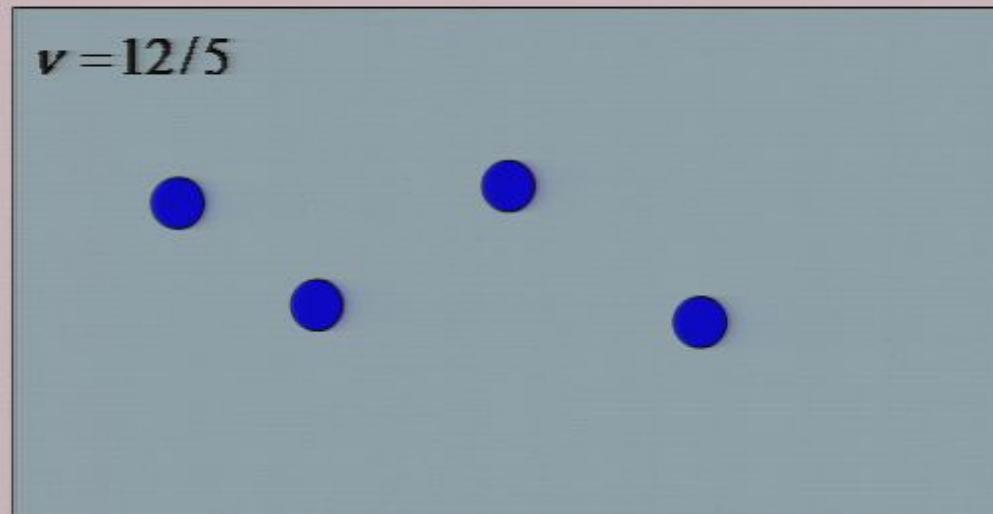
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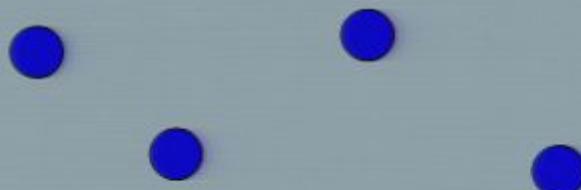
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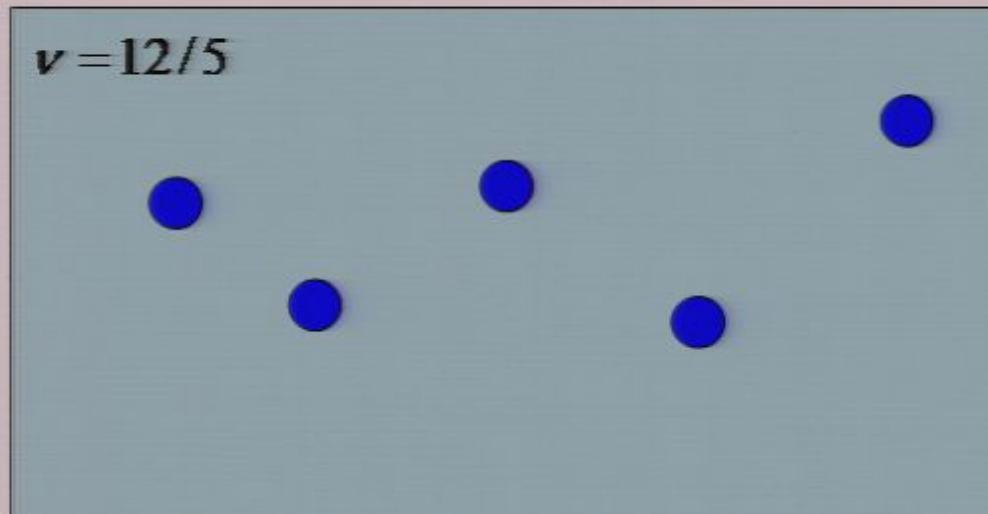


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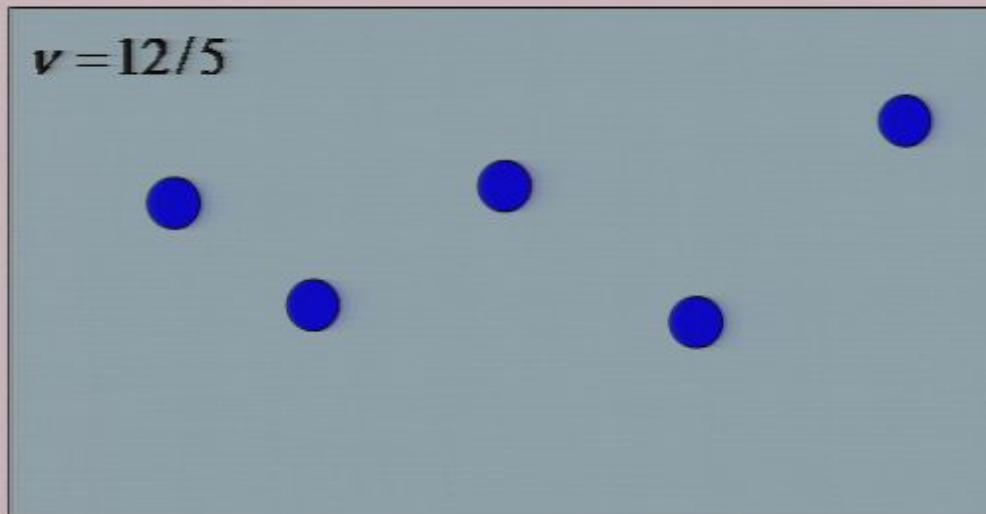
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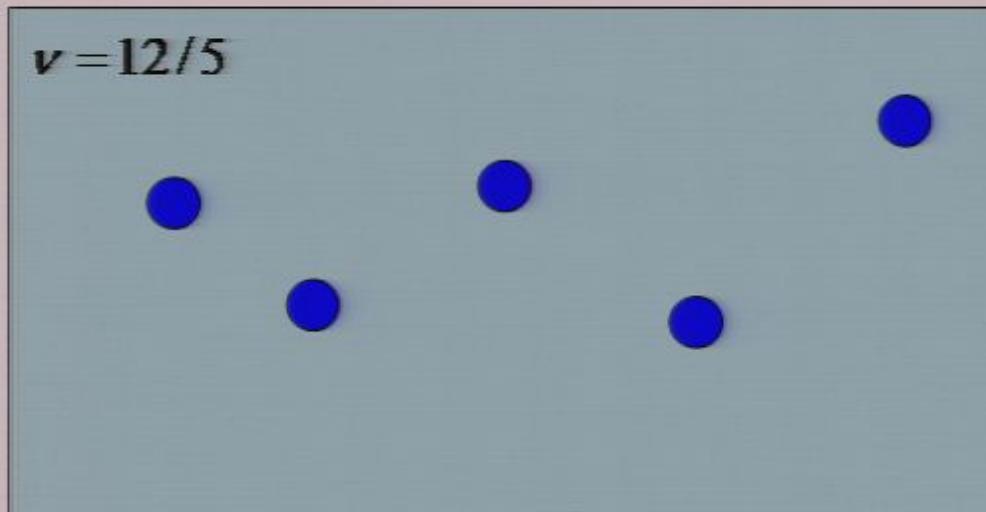
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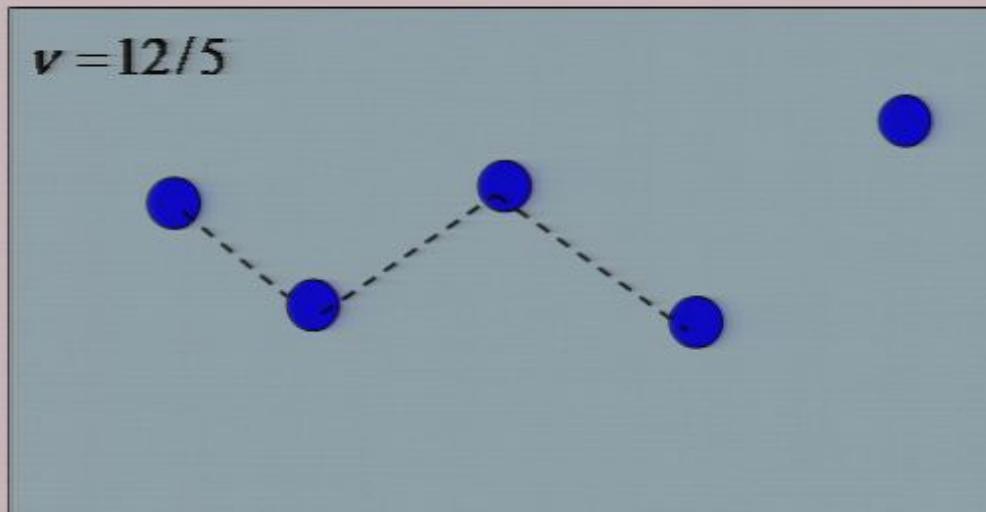
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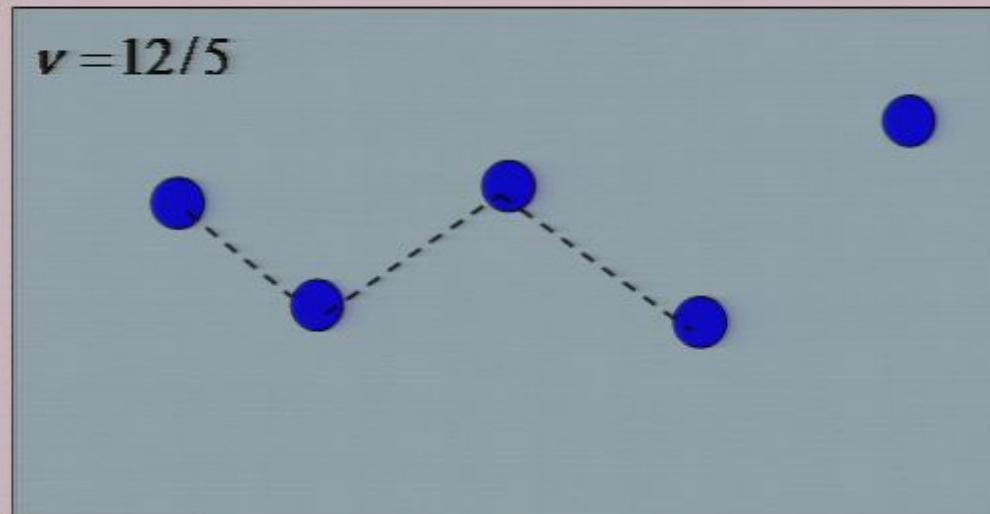
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$$\tau \otimes \tau = Id \oplus \tau$$

Compare to spin half:

$$(1/2) \otimes (1/2) = (Id) \oplus (1)$$

(golden mean)

$$\dim(H) = \left(\frac{1 + \sqrt{5}}{2} \right)^N = \phi^N = 1.618^N$$

- Non-interacting degeneracy:

- But quasi-particles interact with neighbors

Ground state?

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Low energy properties?

Diagrammatic representation

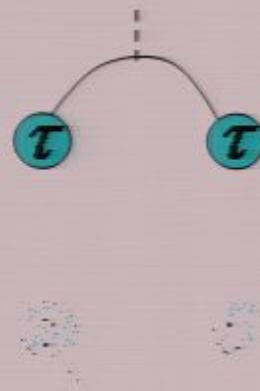
Diagrammatic representation

- Propagation lines:

Id (zero charge) ----- τ —————

- Expectation values:

$$\langle \tau \otimes \tau = Id | \tau \otimes \tau = Id \rangle =$$



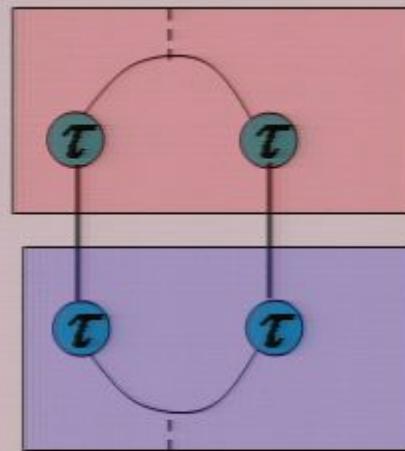
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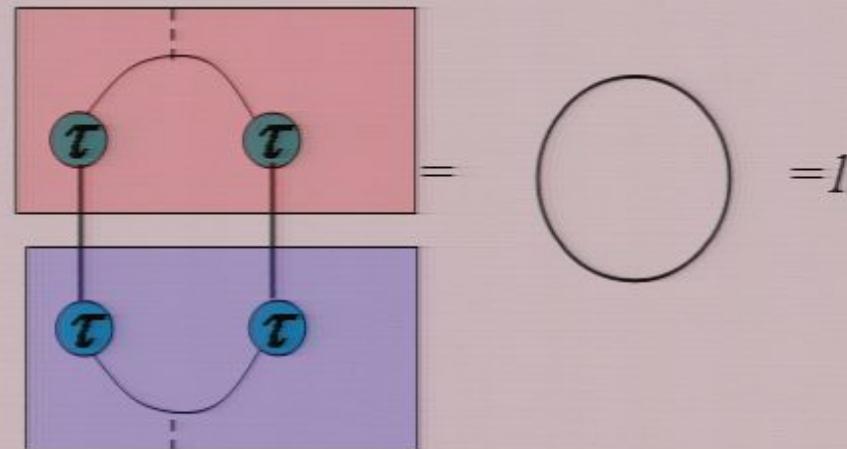
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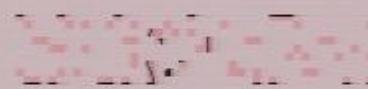
$$\langle \tau \otimes \tau = Id | \tau \otimes \tau = Id \rangle =$$



- Hamiltonian:

$$H = - \sum_i J_i \hat{P}_{i,i+1}^{(Id)} \quad (\text{sum of projections})$$

- Projections:



$$\langle \psi | P^{(Id)} | \chi \rangle =$$



Diagrammatic representation

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= 1

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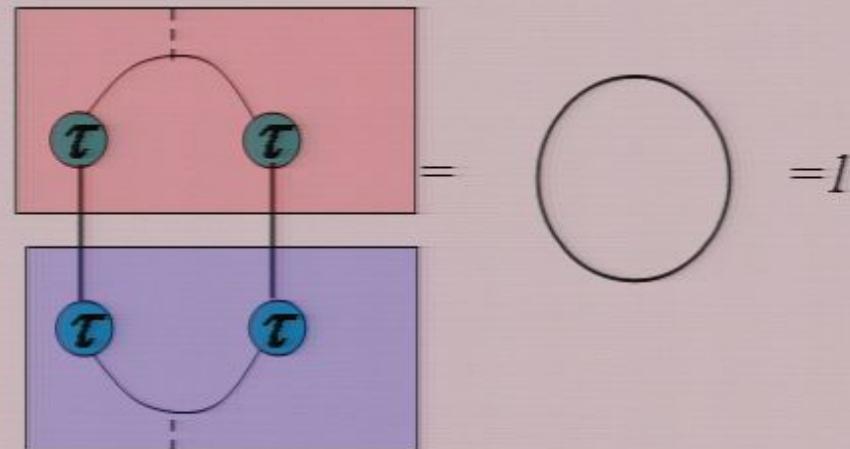
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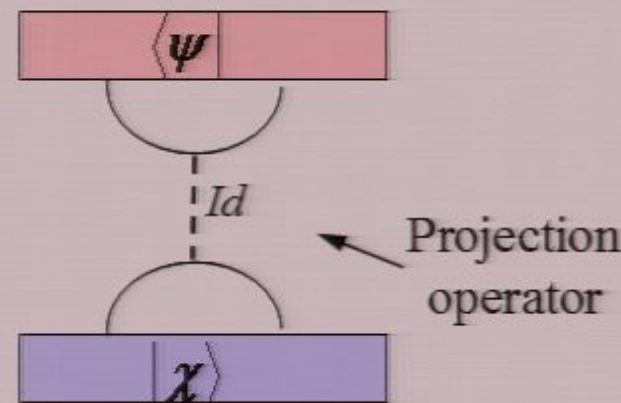


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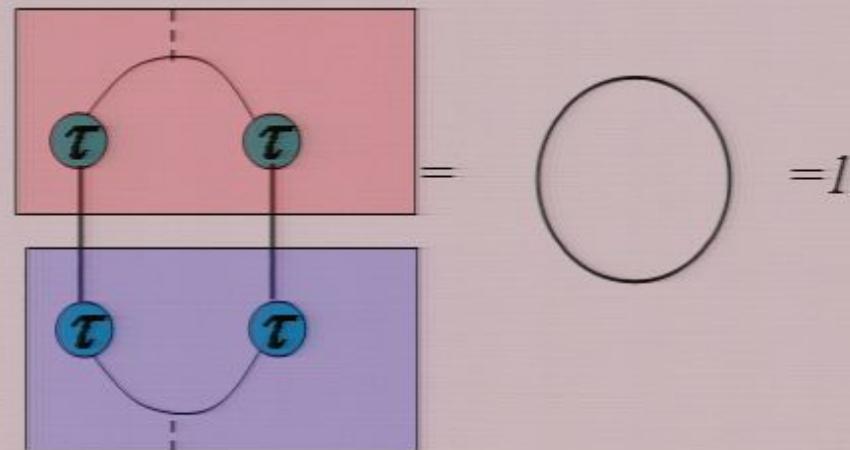
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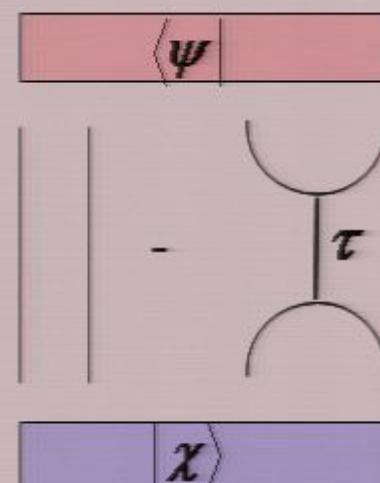
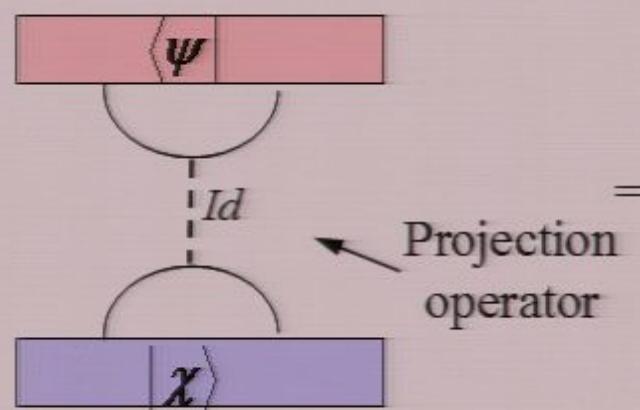


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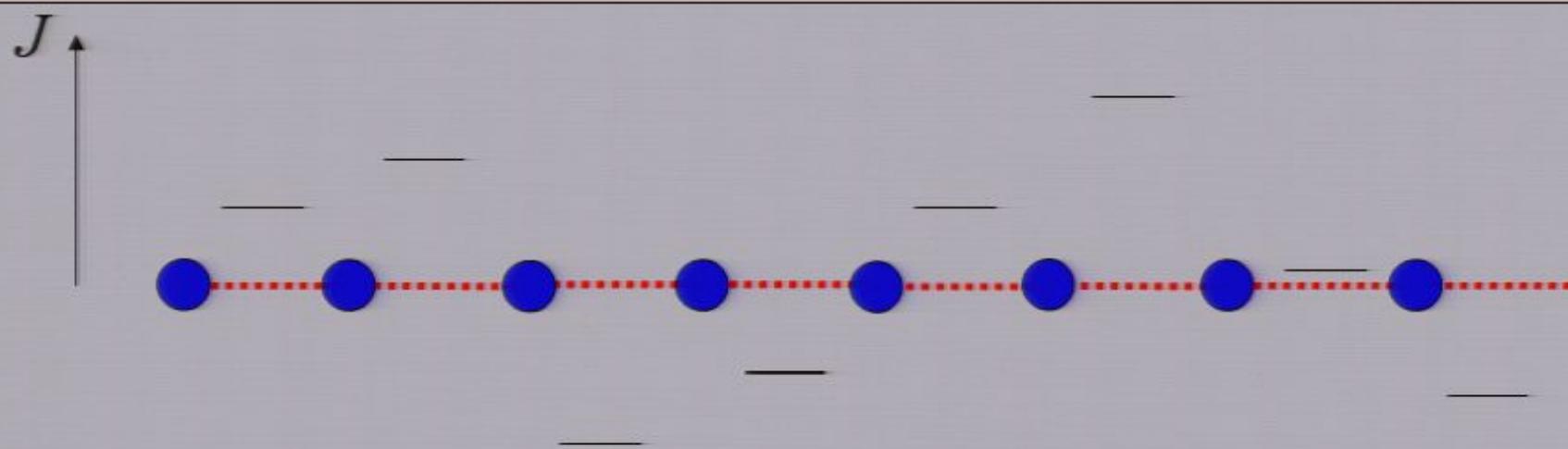
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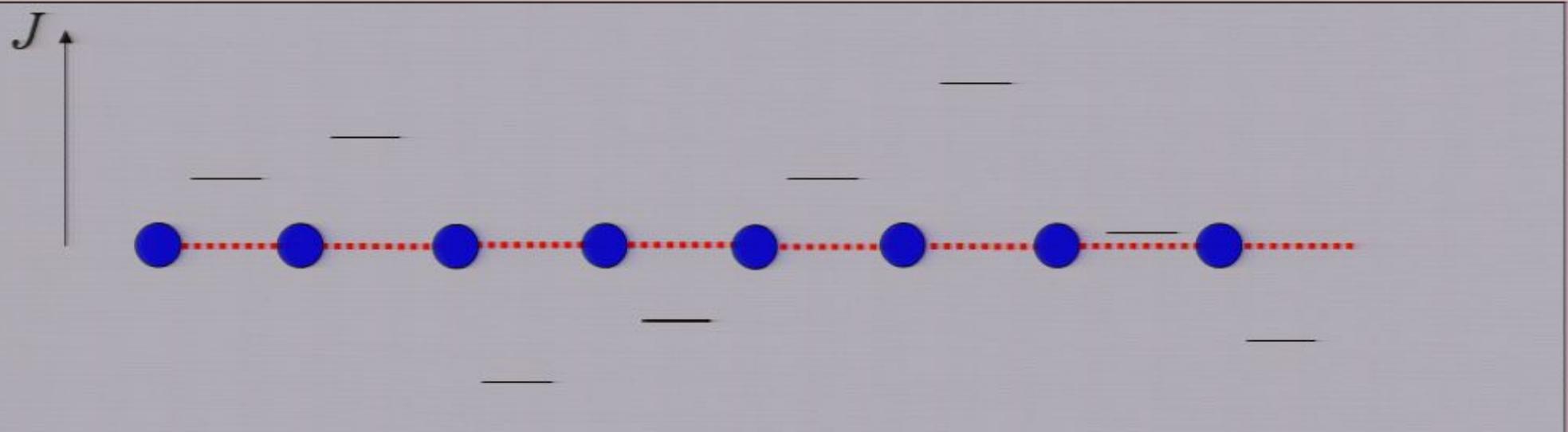
Random Fibonacci chain

$$H = -\sum_i J_i \hat{P}_{i,i+1}^{(Id)}$$



Random Fibonacci chain

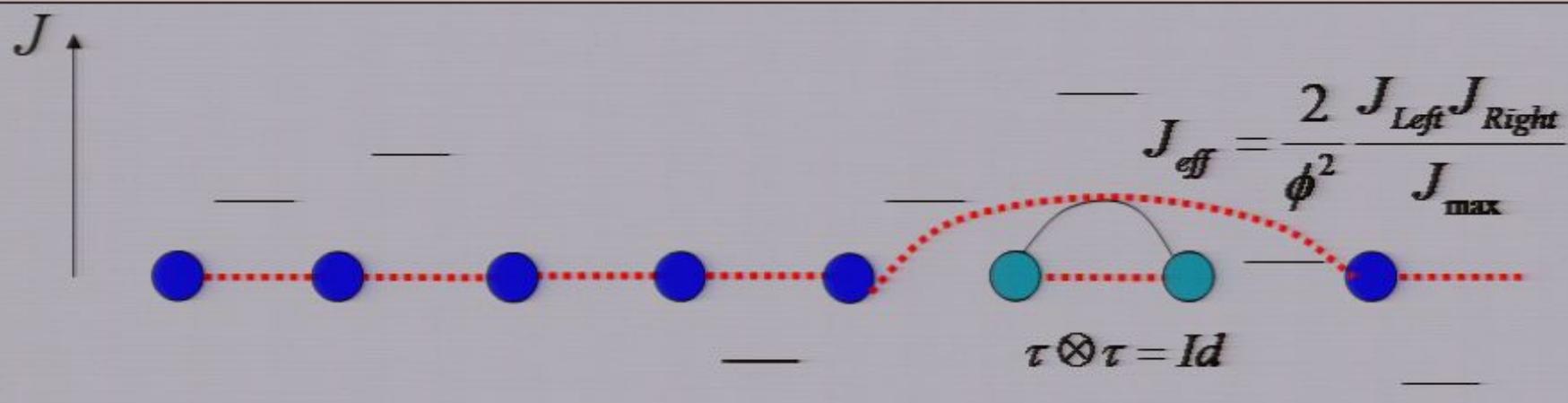
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- Ma-Dasgupta singlet channel decimation – same as before.

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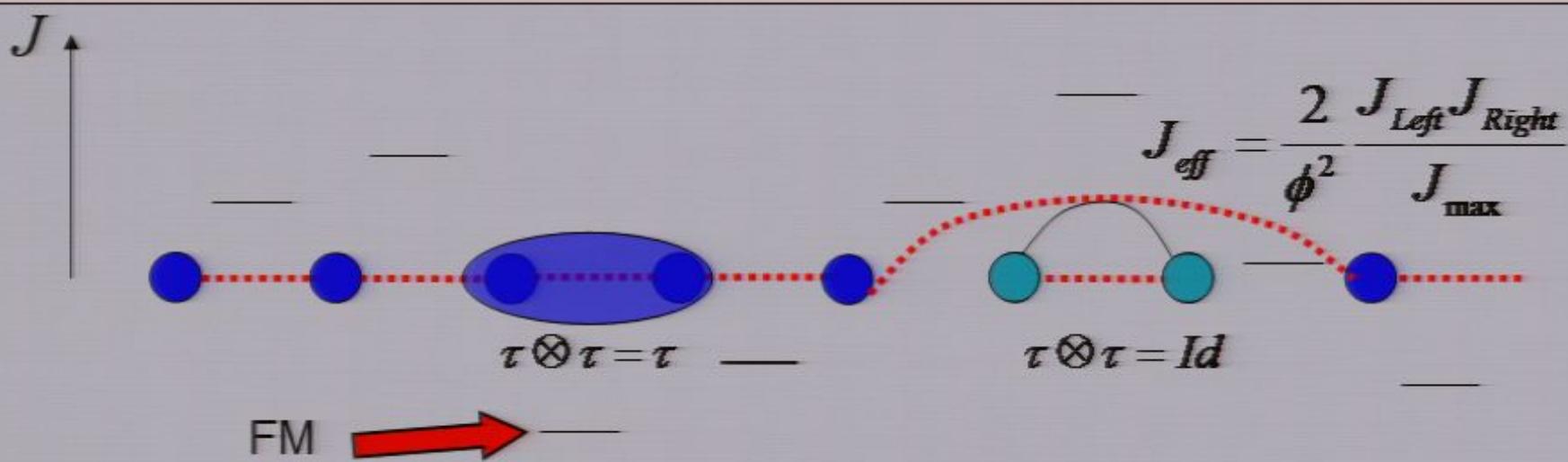
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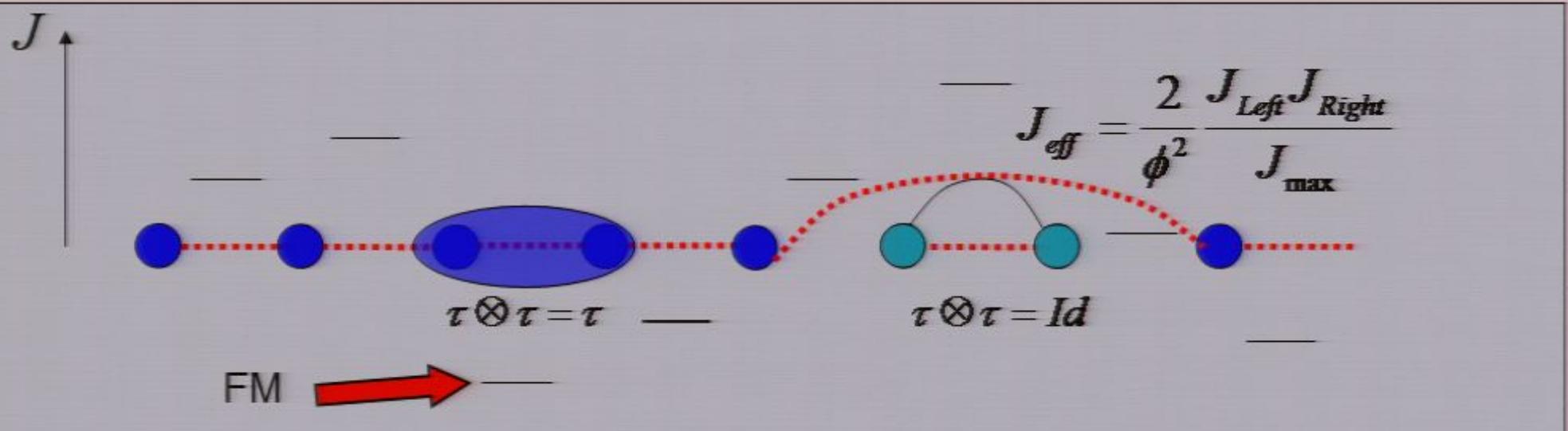
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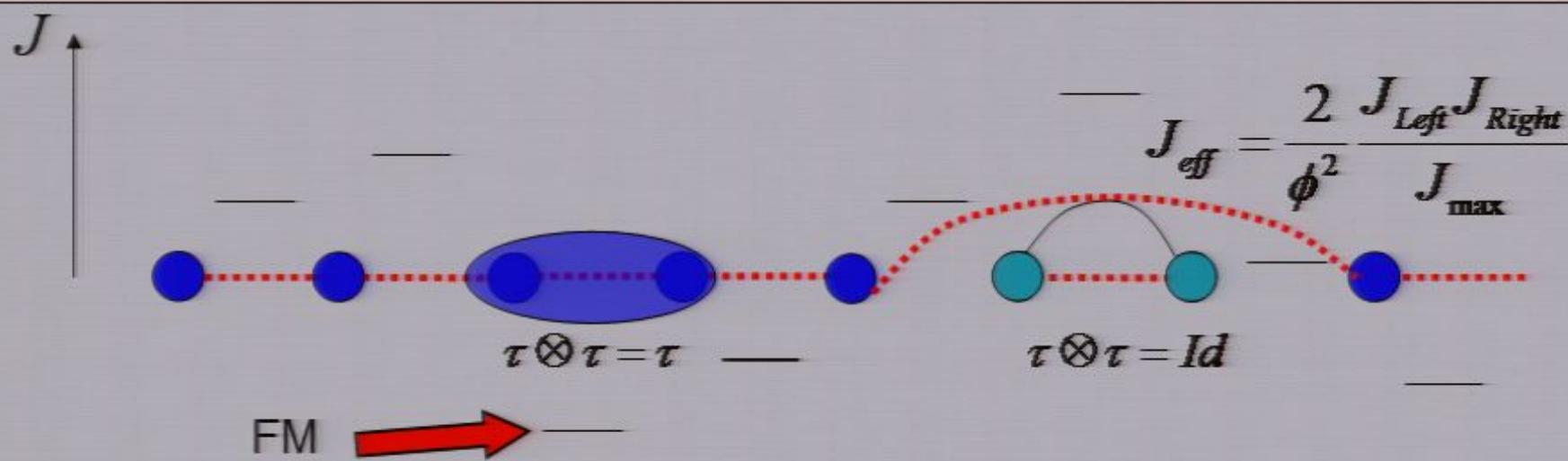
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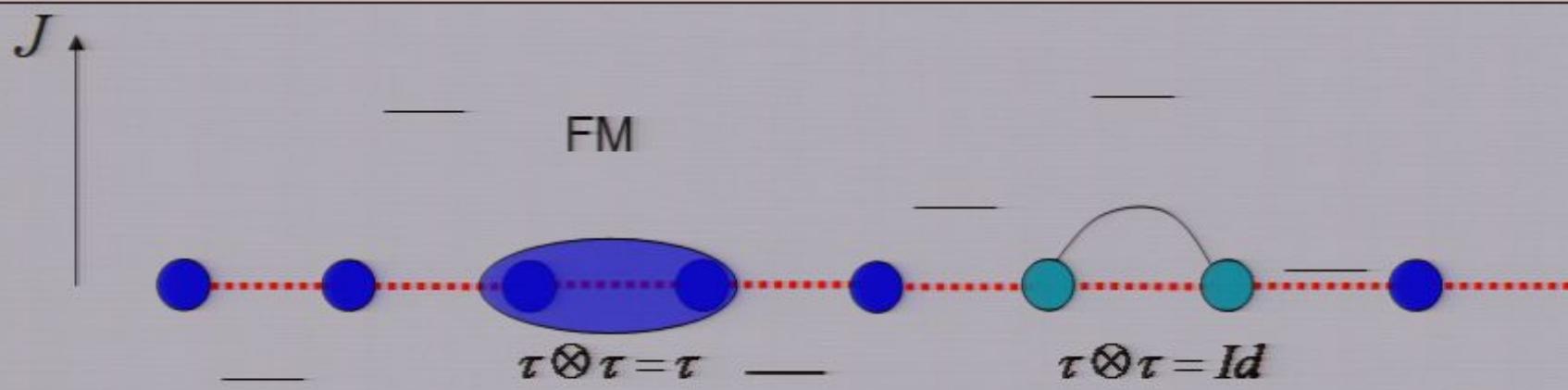


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→ $-J_{\text{eff}} \cdot P_{(12),3}^{(Id)}$

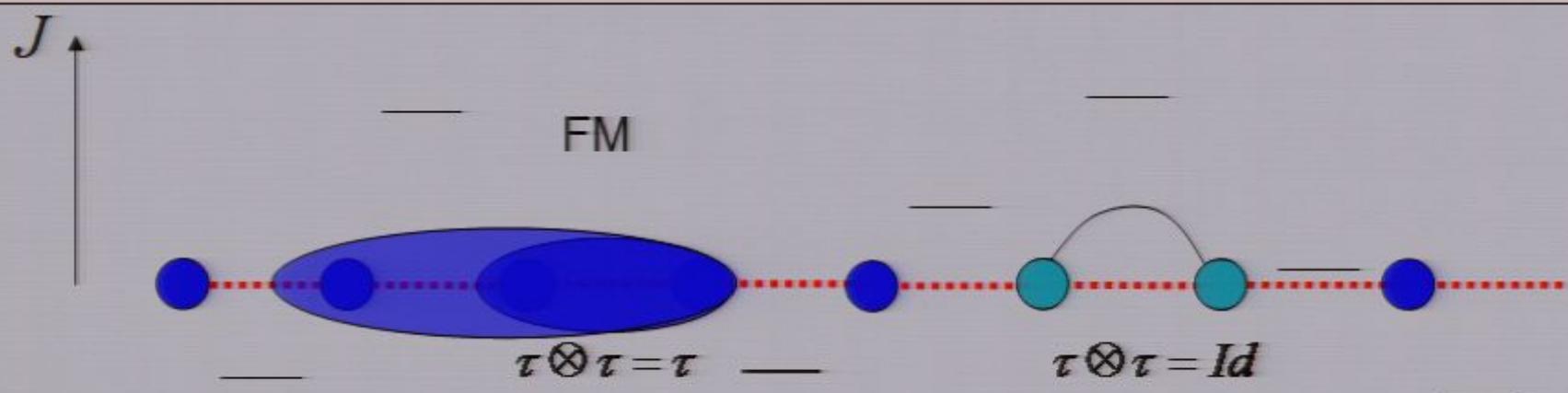
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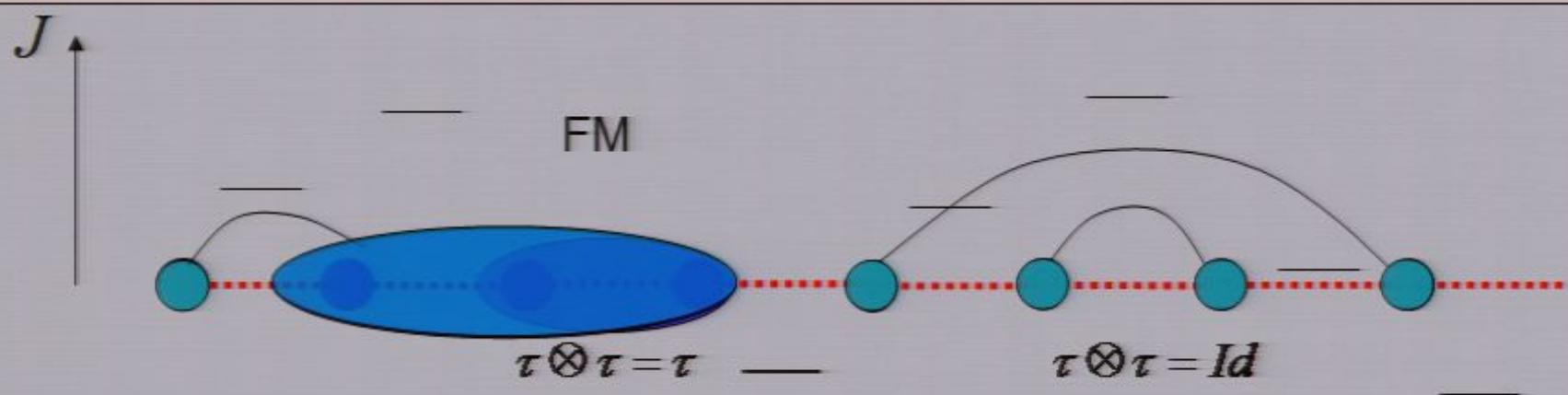
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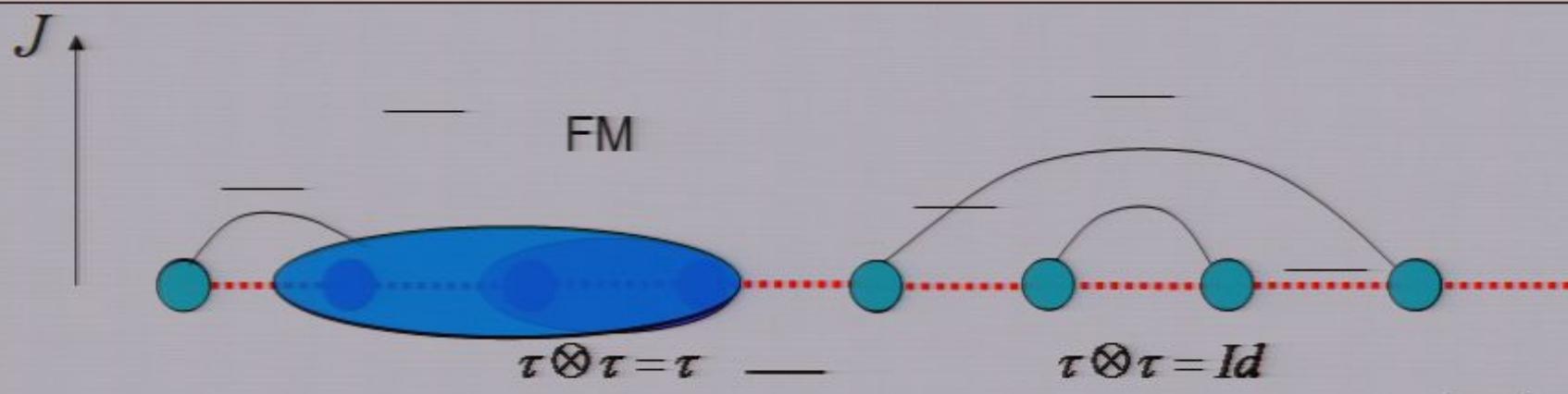
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- Phase diagram:

Random Fibonacci chain

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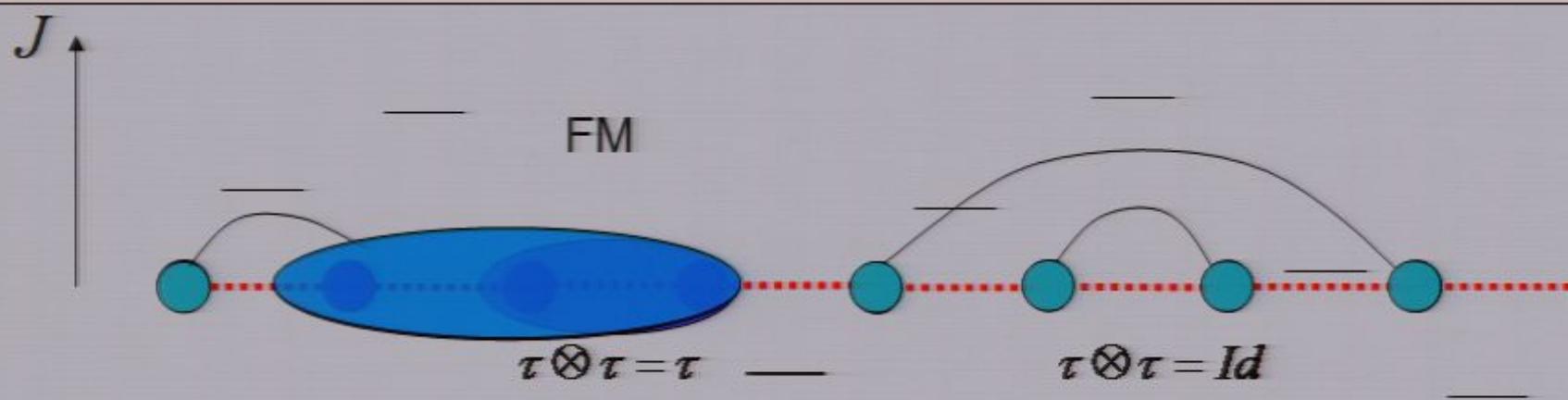
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$$N_{J>0}/N_{J<0} =$$

$$FM/AFM$$

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- Phase diagram:

Random singlet
phase

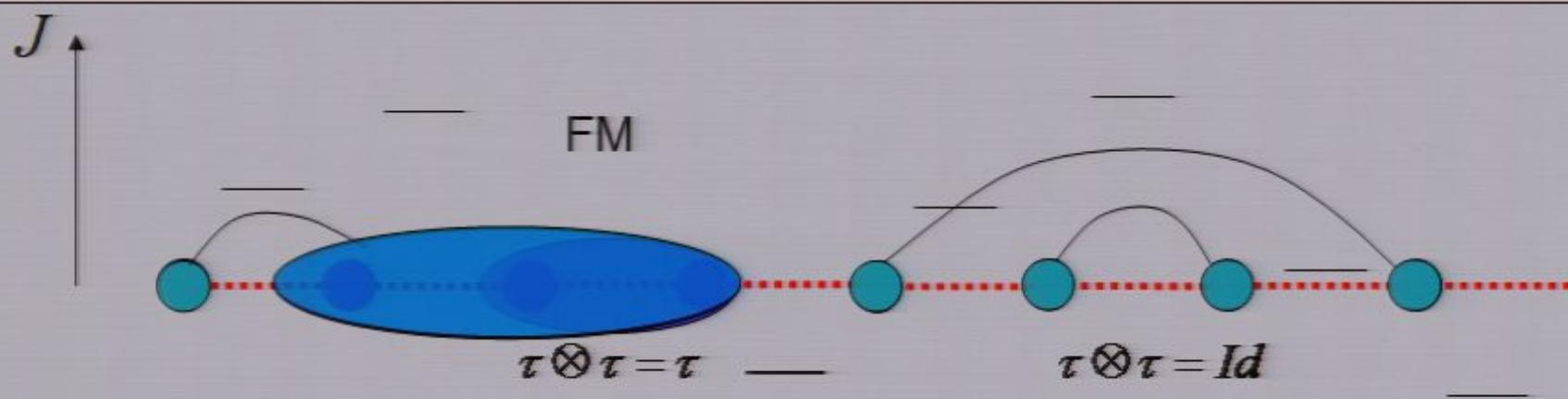


0
(AFM)
Bonesteel, Yang

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FM / AFM

Random Fibonacci chain

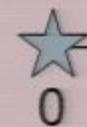
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- Phase diagram:

Random singlet
phase

Mixed phase



0



1

$N_{J>0} / N_{J<0} =$

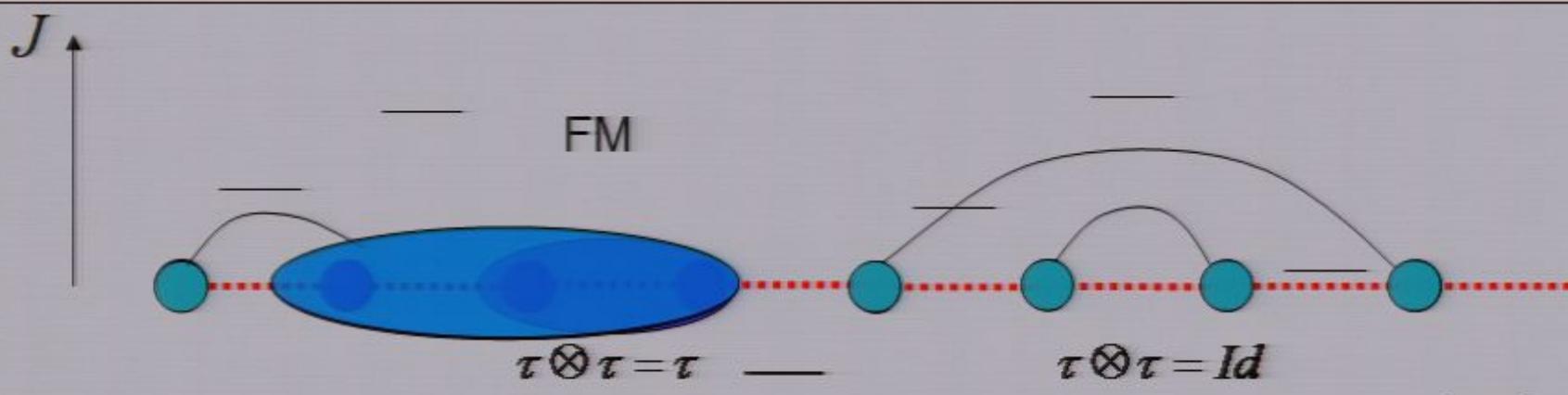
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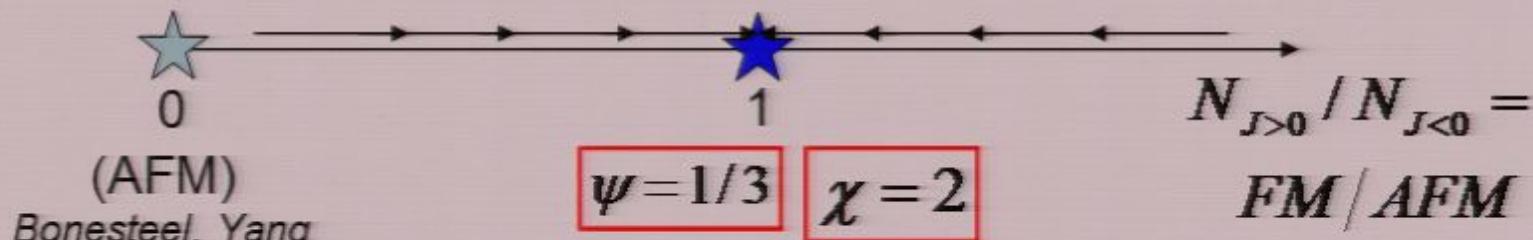
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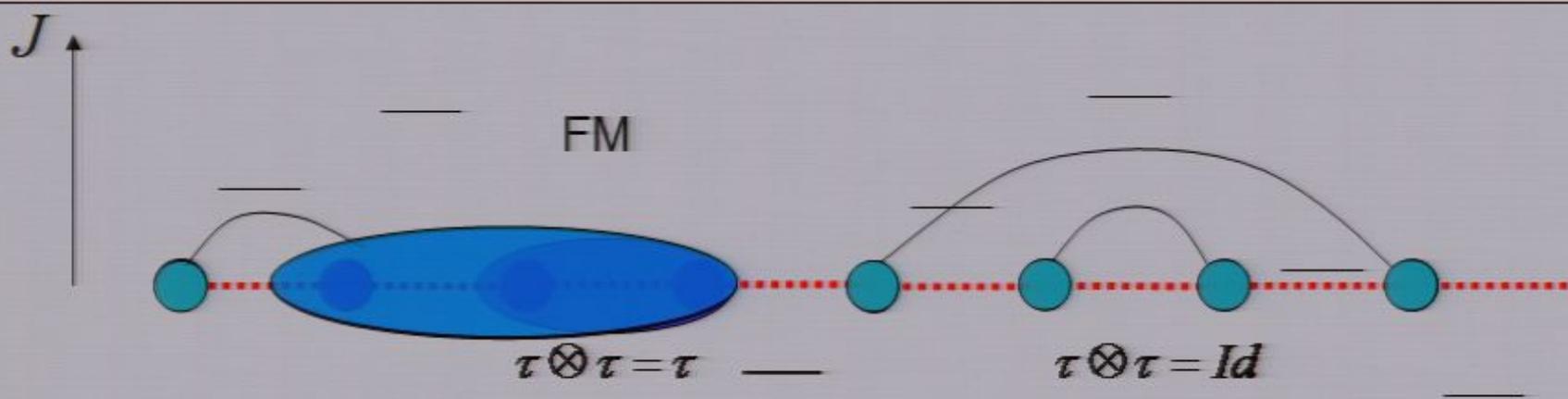
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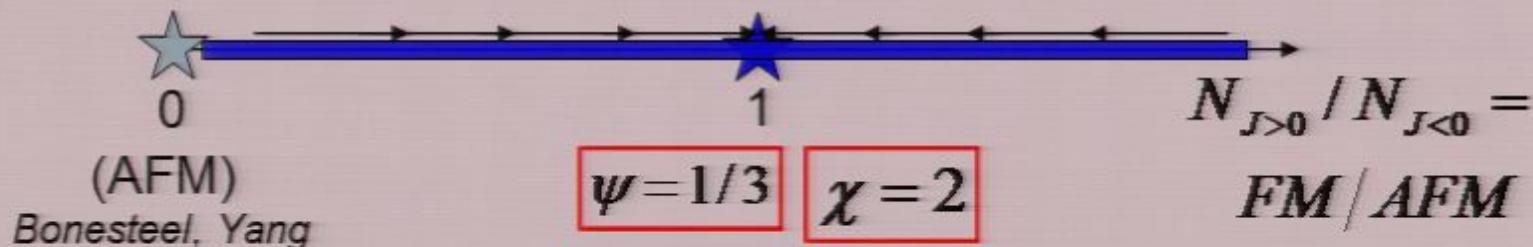
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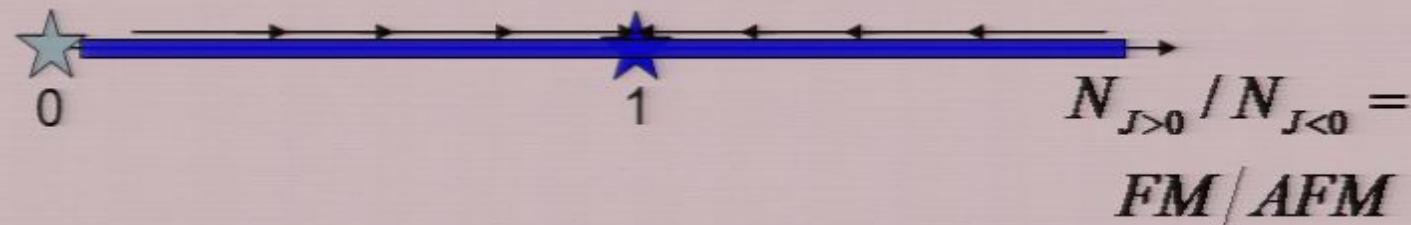
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$$E_L = \frac{1}{3} c_{\text{eff}} \cdot \log_2 L$$

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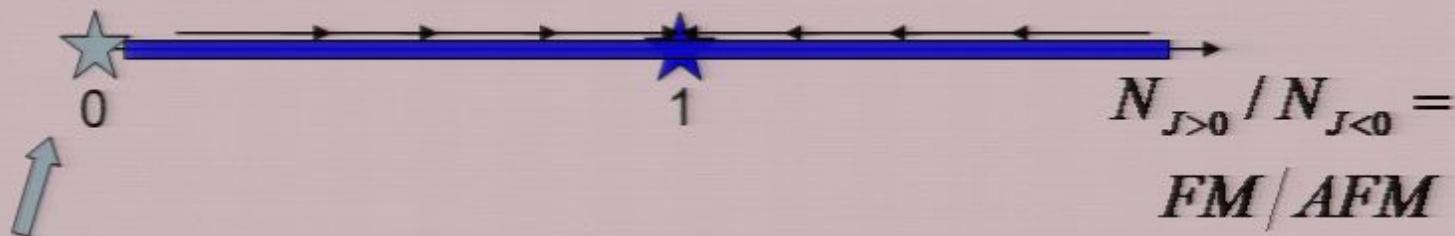
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$$\ln \phi = 0.481$$

(Yang, Bonsteel, 2006)

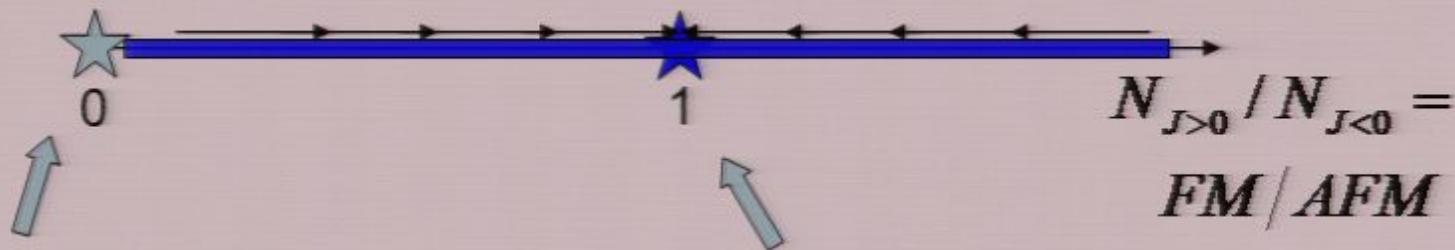
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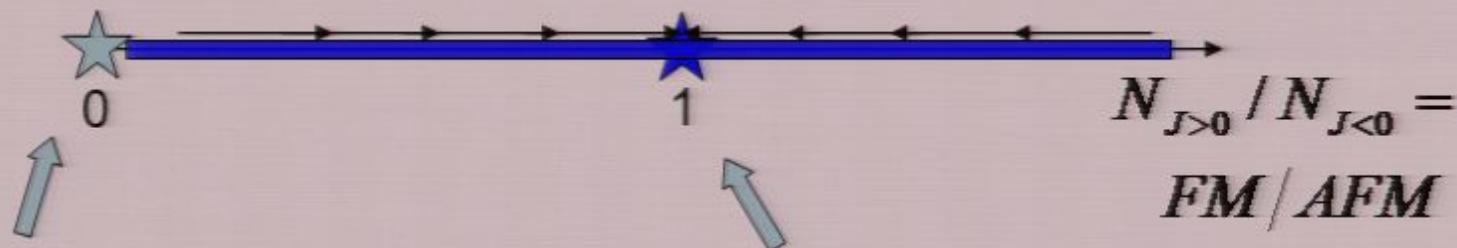
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Effective central charge grows along RG flow..

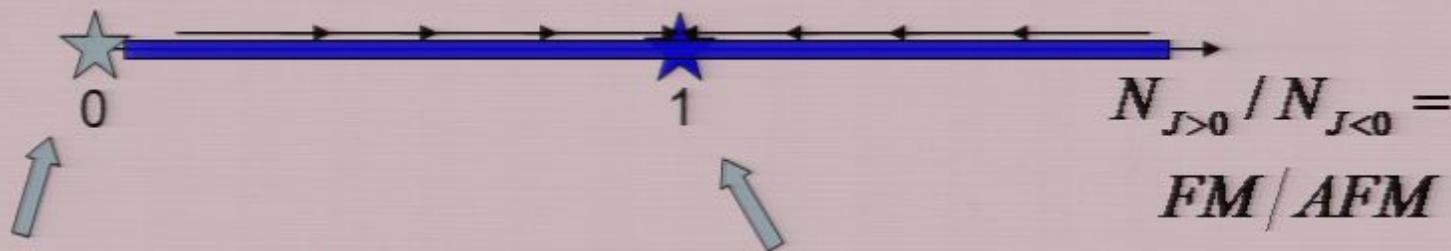
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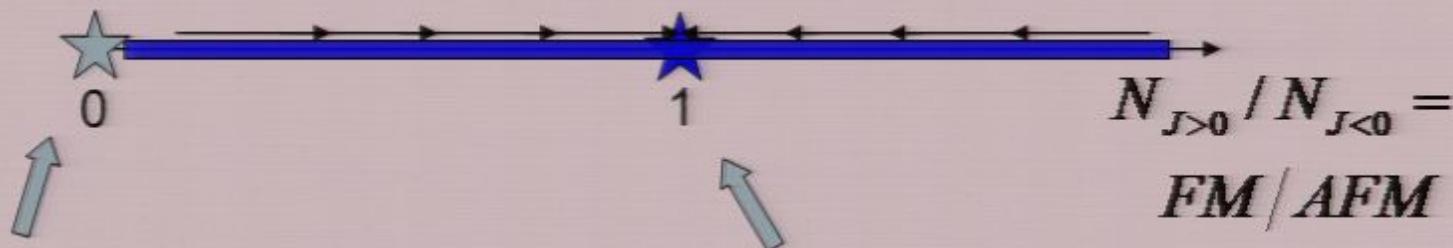
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Effective central charge grows along RG flow..

General class of non-Abelian anyons: $SU(2)k$

(w/ Lukasz Fidkowski,
and SURF students Paraj Titum, Han-Hsuan Lin)

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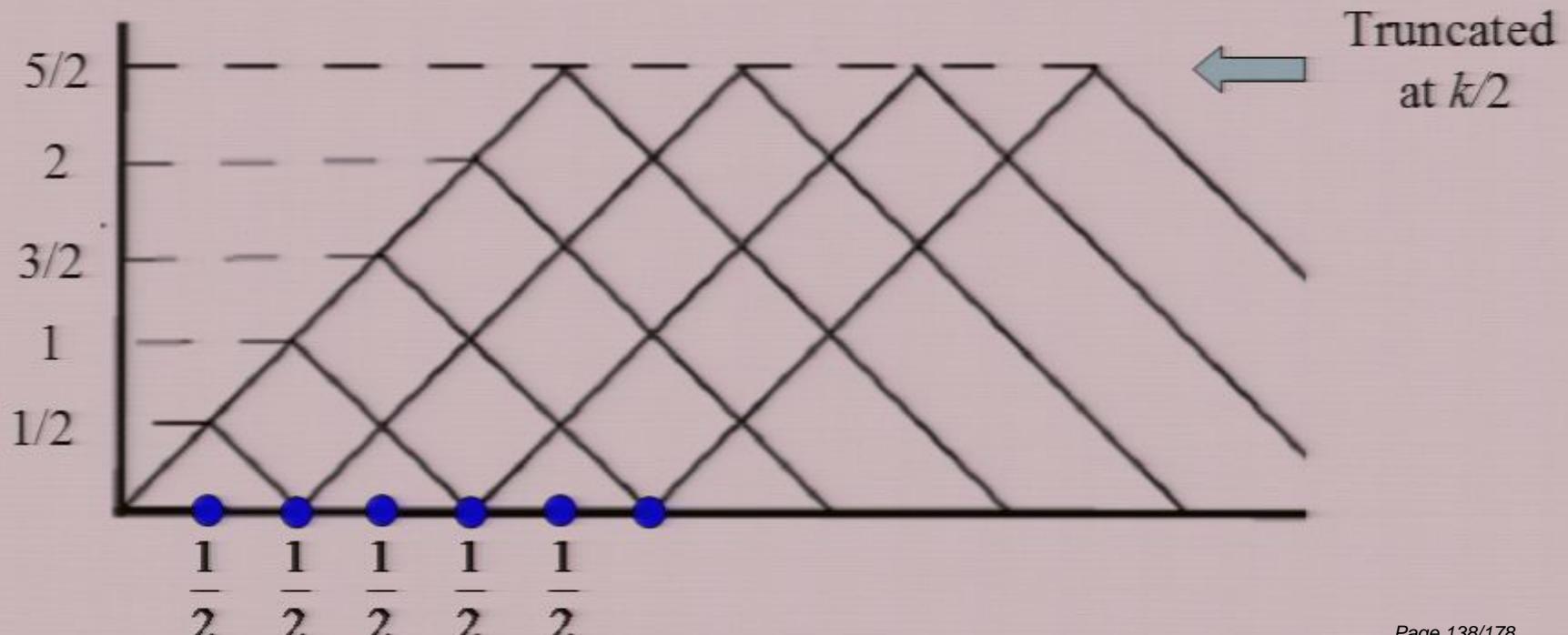
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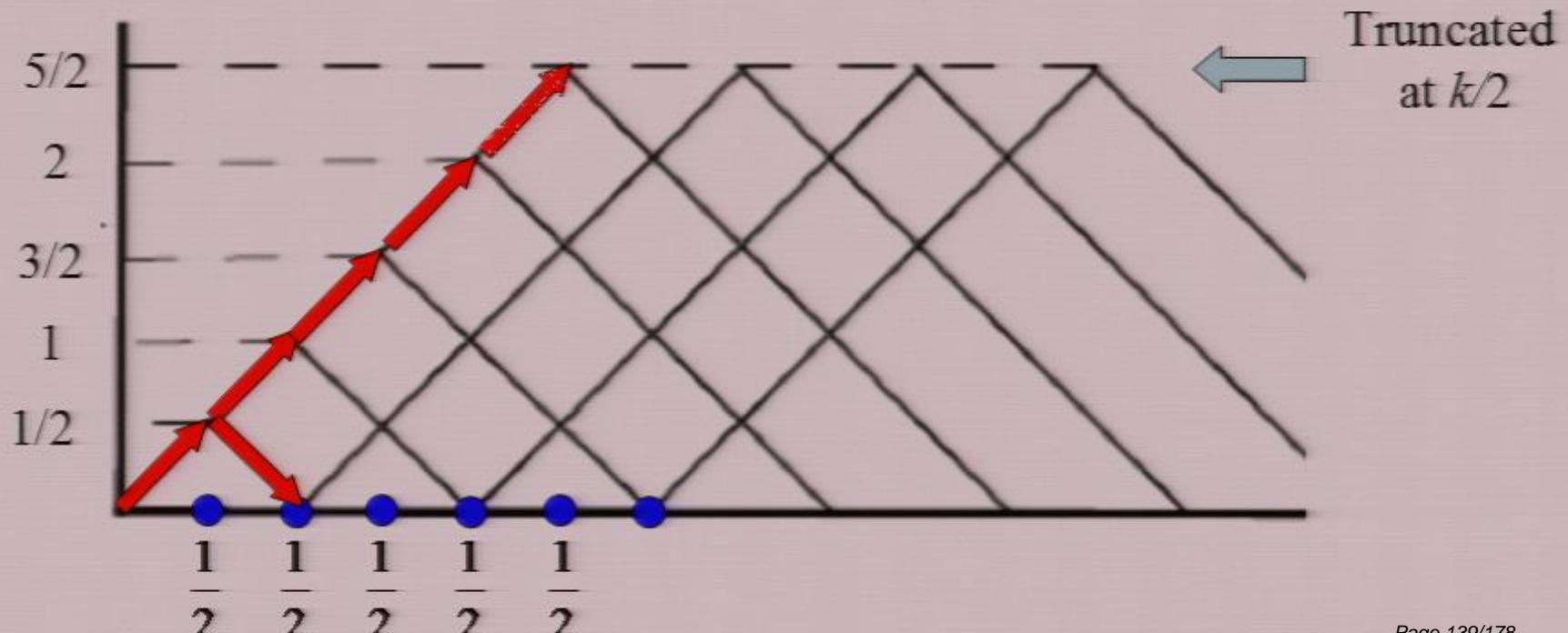
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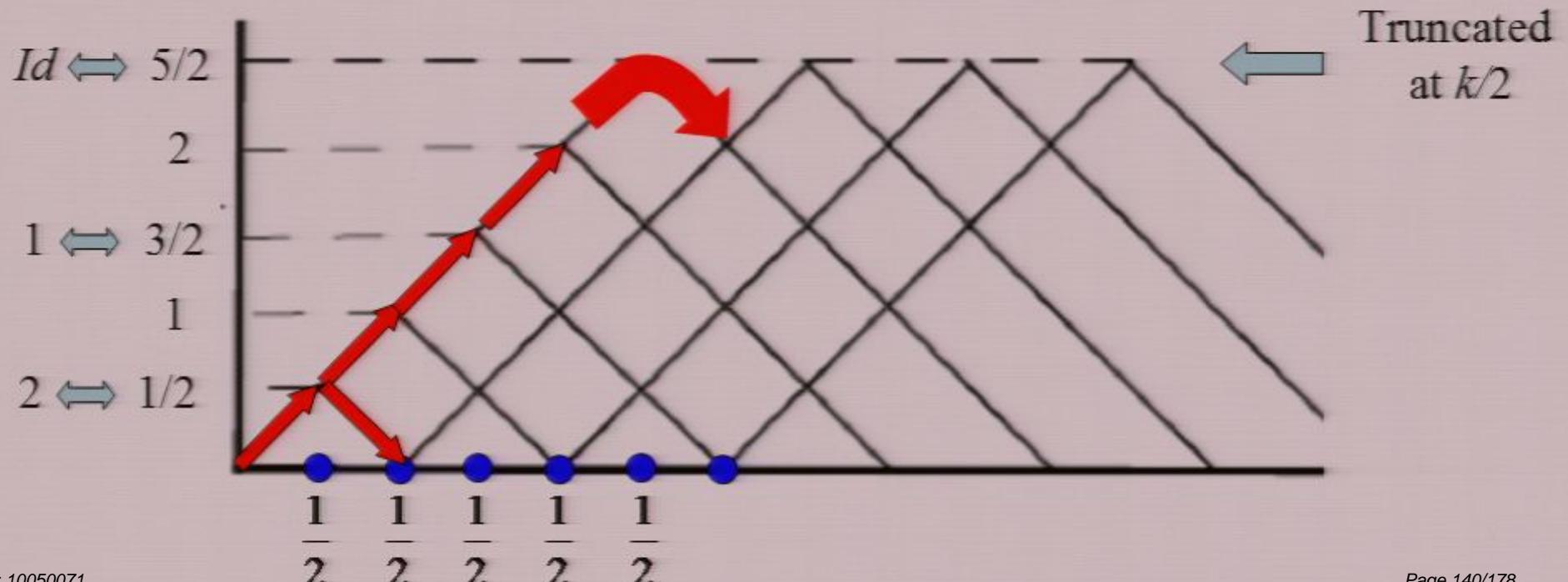
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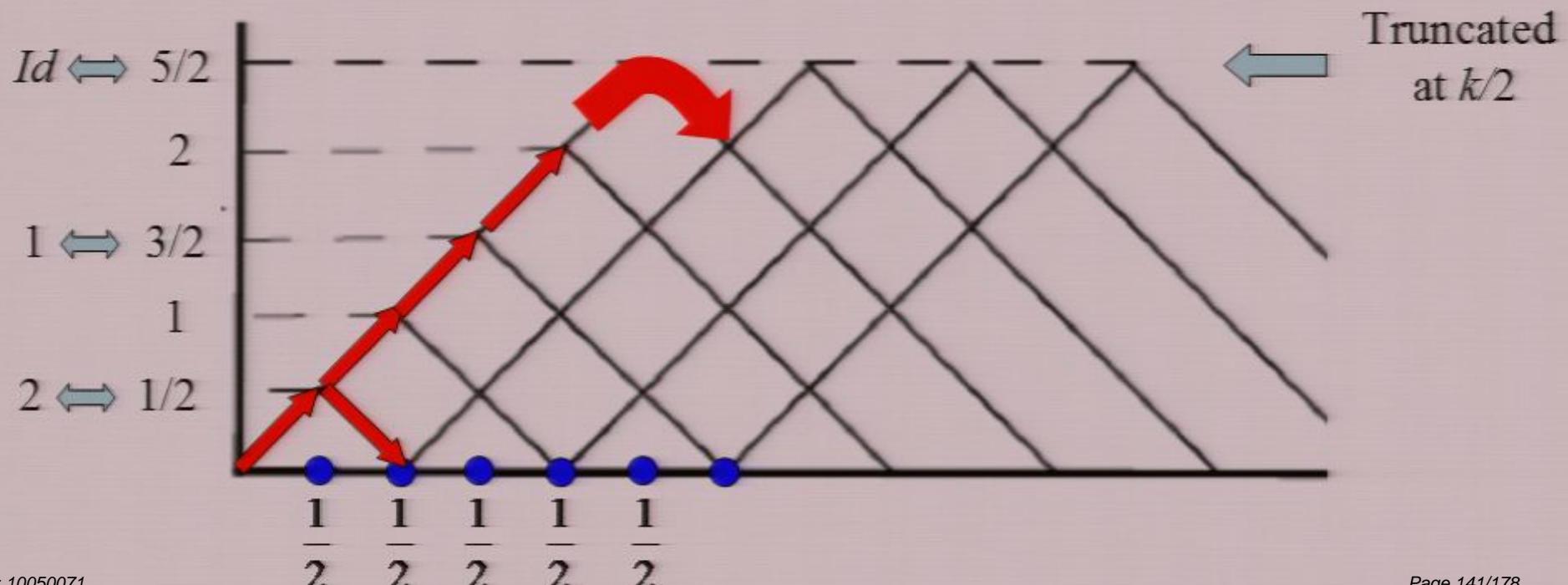
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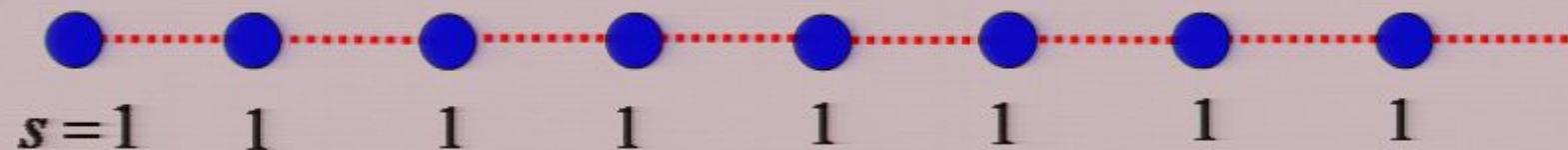
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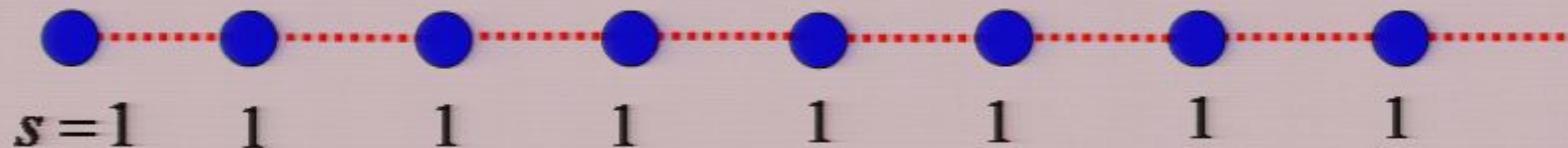
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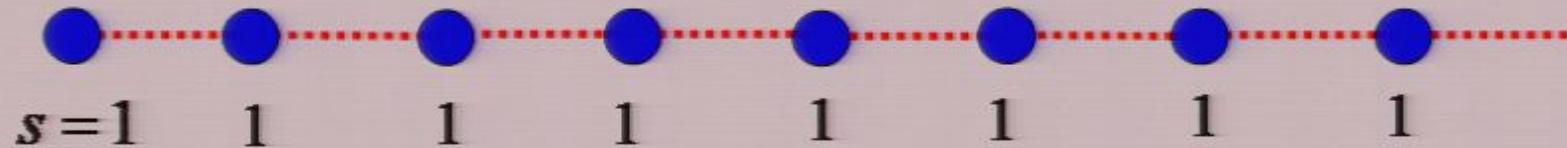
- Typical Hamiltonian:

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The diagram shows two sets of circles labeled i and $i+1$. The top set has two circles labeled i and $i+1$ connected by a curved line. A vertical line labeled s connects the bottom set of circles labeled i and $i+1$ to the top set. The bottom set also has two circles labeled i and $i+1$ connected by a curved line.

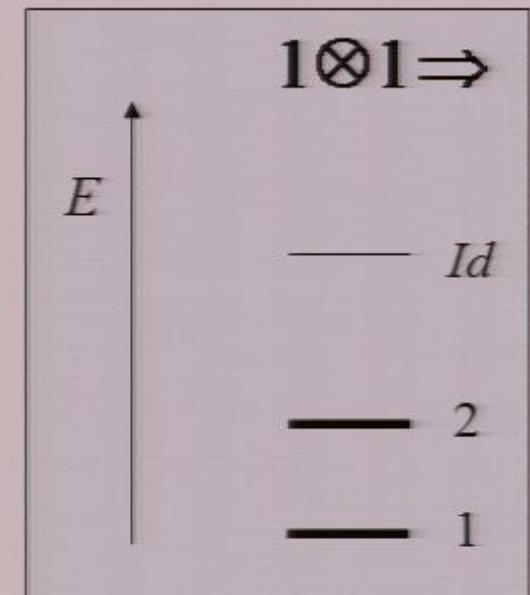
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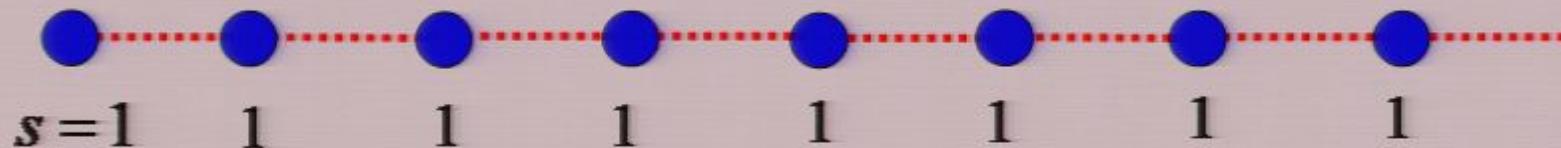
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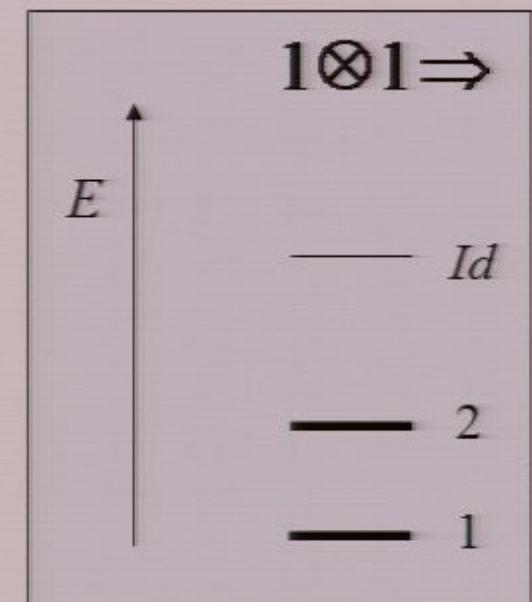
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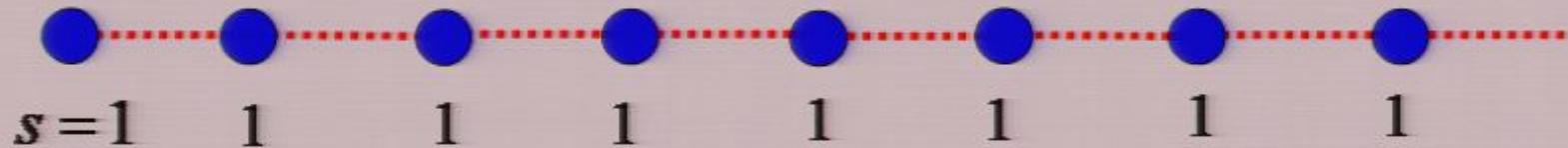
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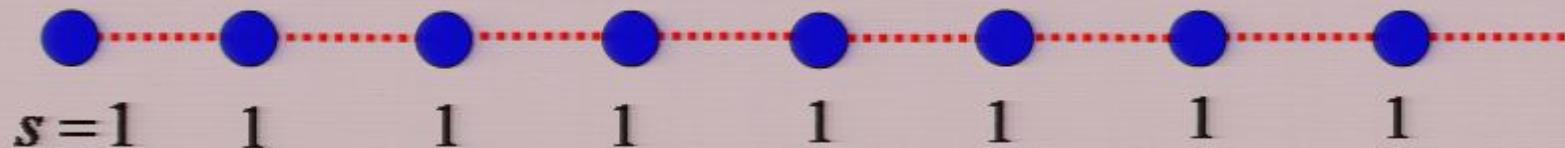
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A diagram showing two vertical teal circles labeled 'i' and 'i+1' connected by a horizontal line. A vertical line connects the top 'i' to the bottom 'i'. A vertical line connects the top 'i+1' to the bottom 'i+1'. A horizontal line connects the two bottom circles. The number '1' is placed above the horizontal line connecting the bottom circles.

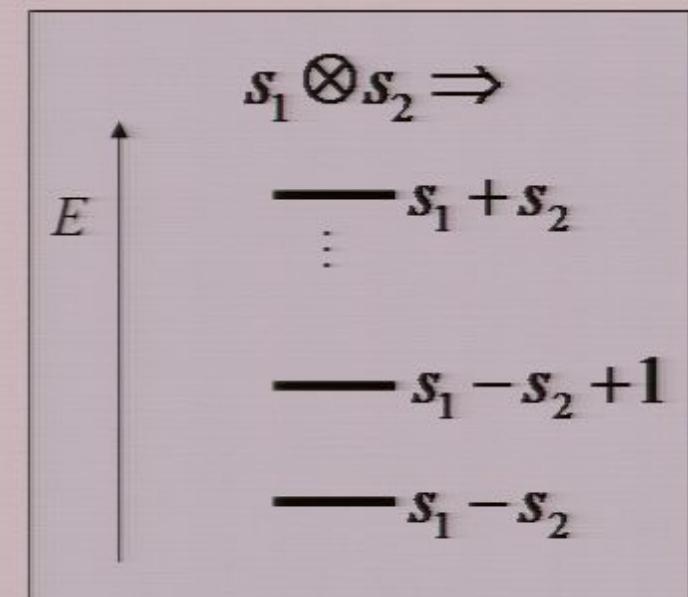
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$$s_1 \otimes s_2 = \sum_{s_3} \left(N_{s_1 s_2}^{s_3} \right) s_3 \quad s_n \in \{0, 1, \dots, \frac{k-1}{2}\}$$



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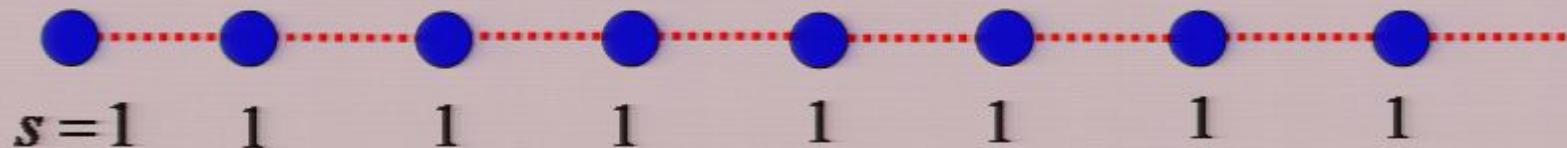


Fixed points of $SU(2)^k$ – Damle-Huse hierarchy

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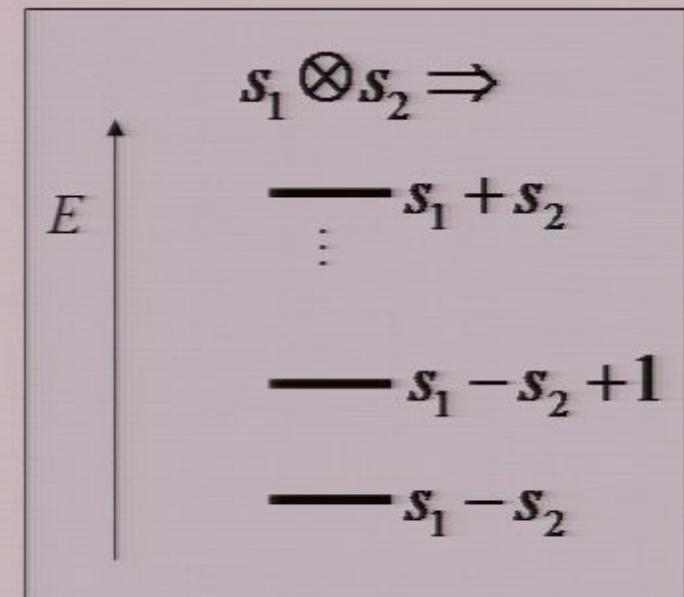
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- All Fixed Points are **Stable!**  Represent stable phases.

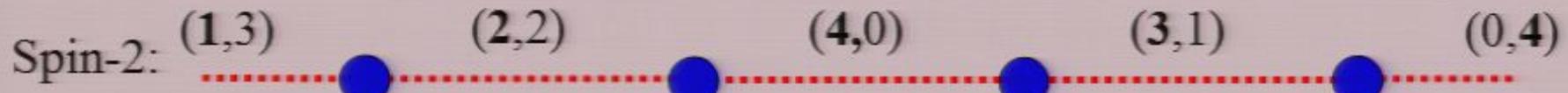
(Contrary to DH points in abelian spins)

Duality

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$k=5$: Domains 0-4



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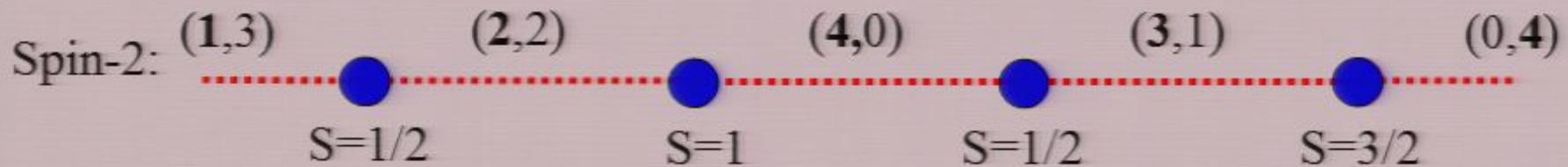


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Int spin
 $f(s) = s$



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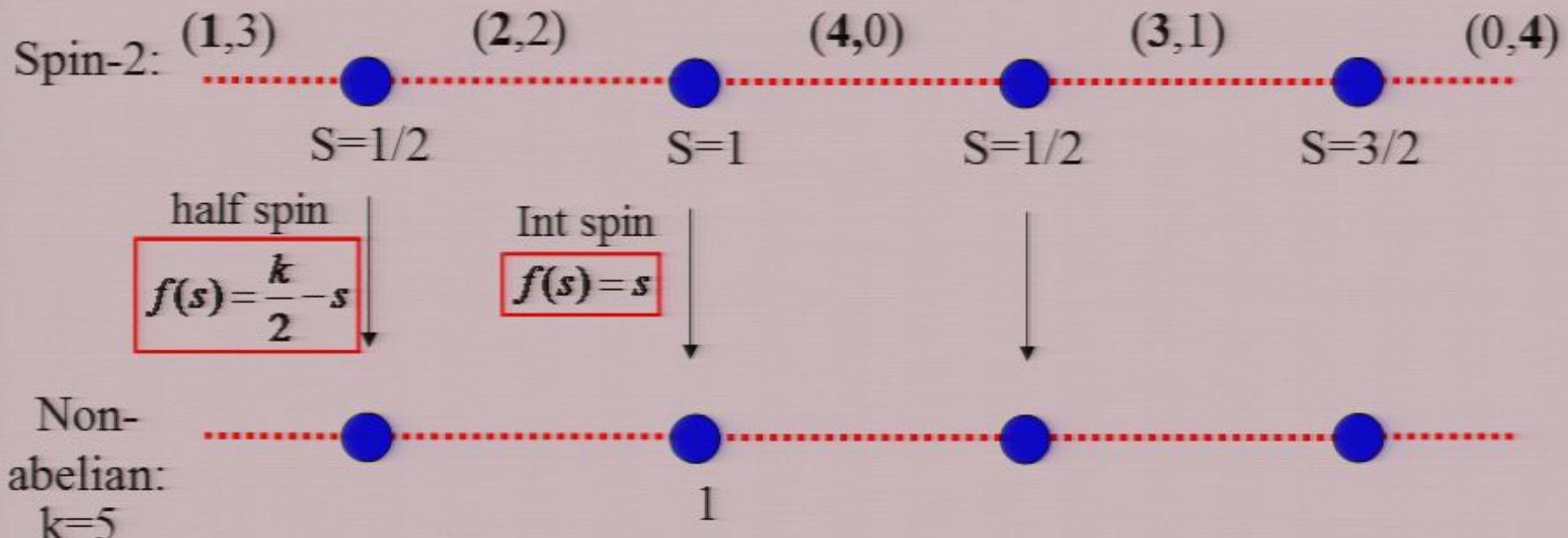


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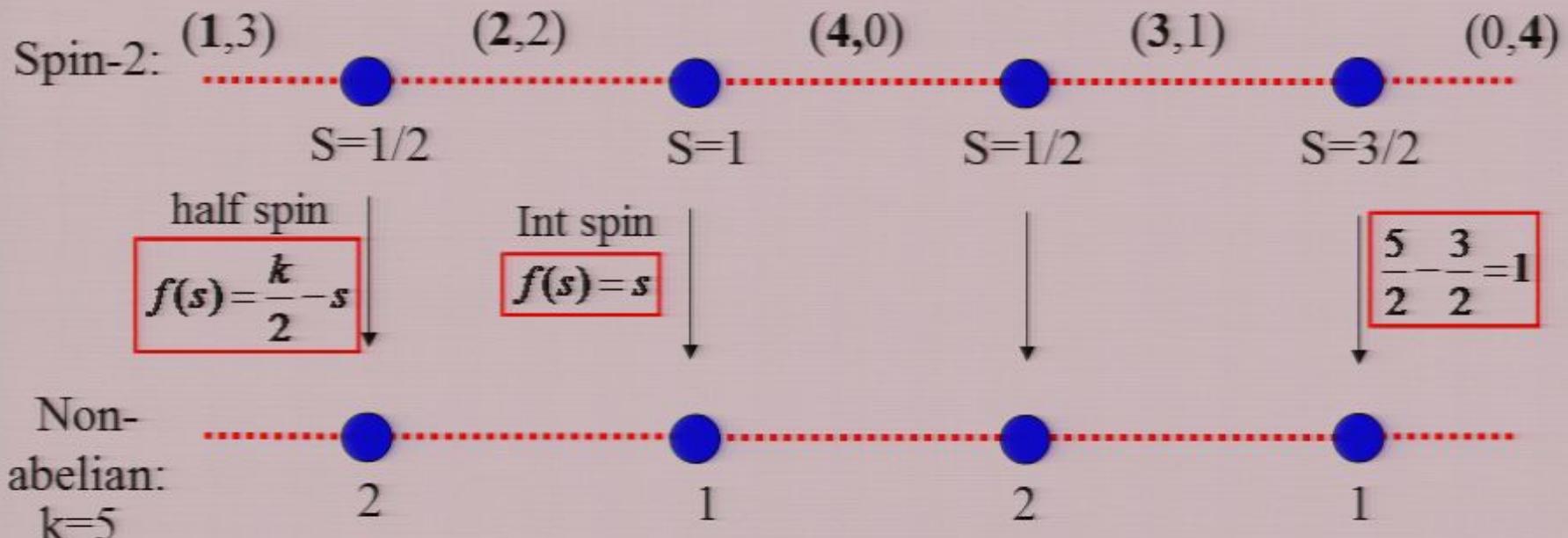


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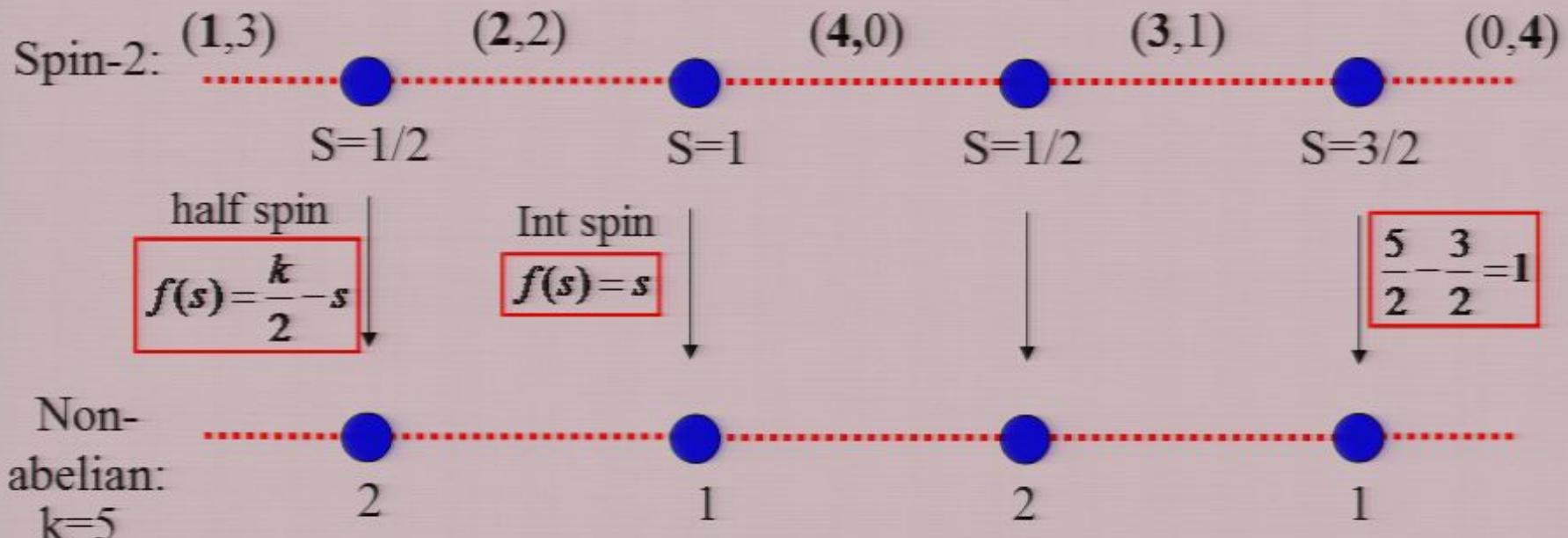


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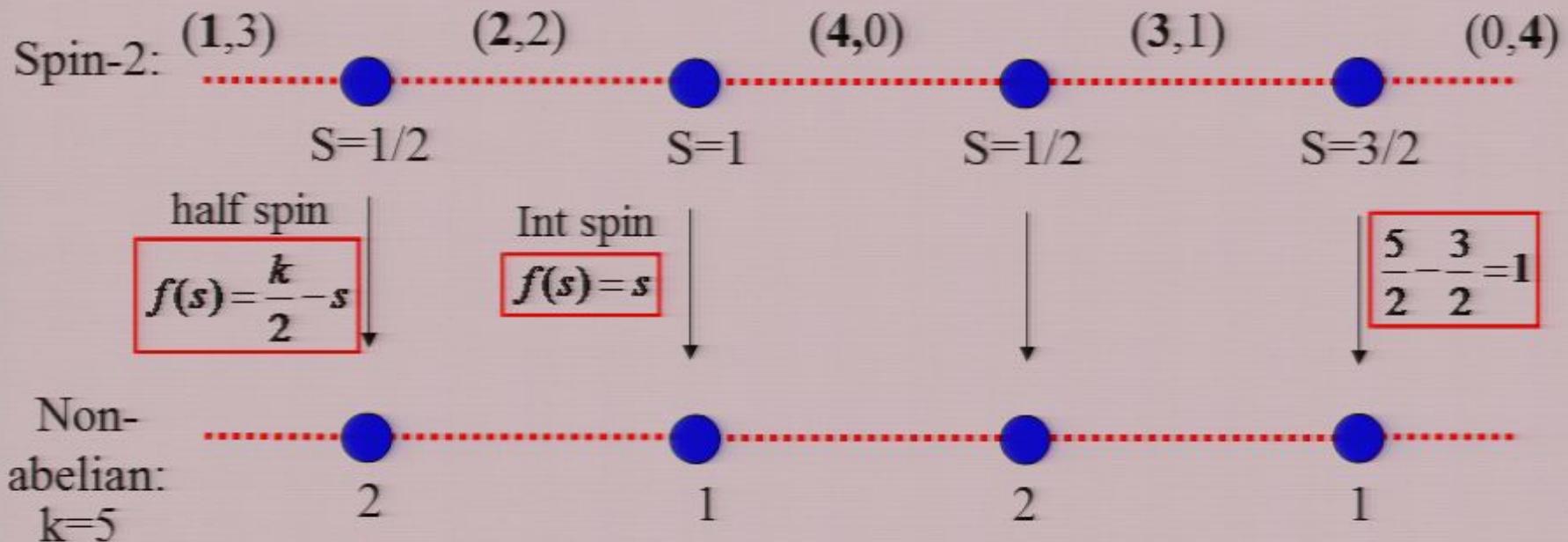
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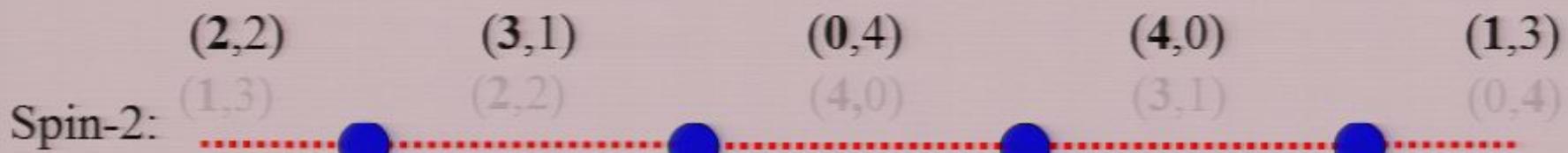
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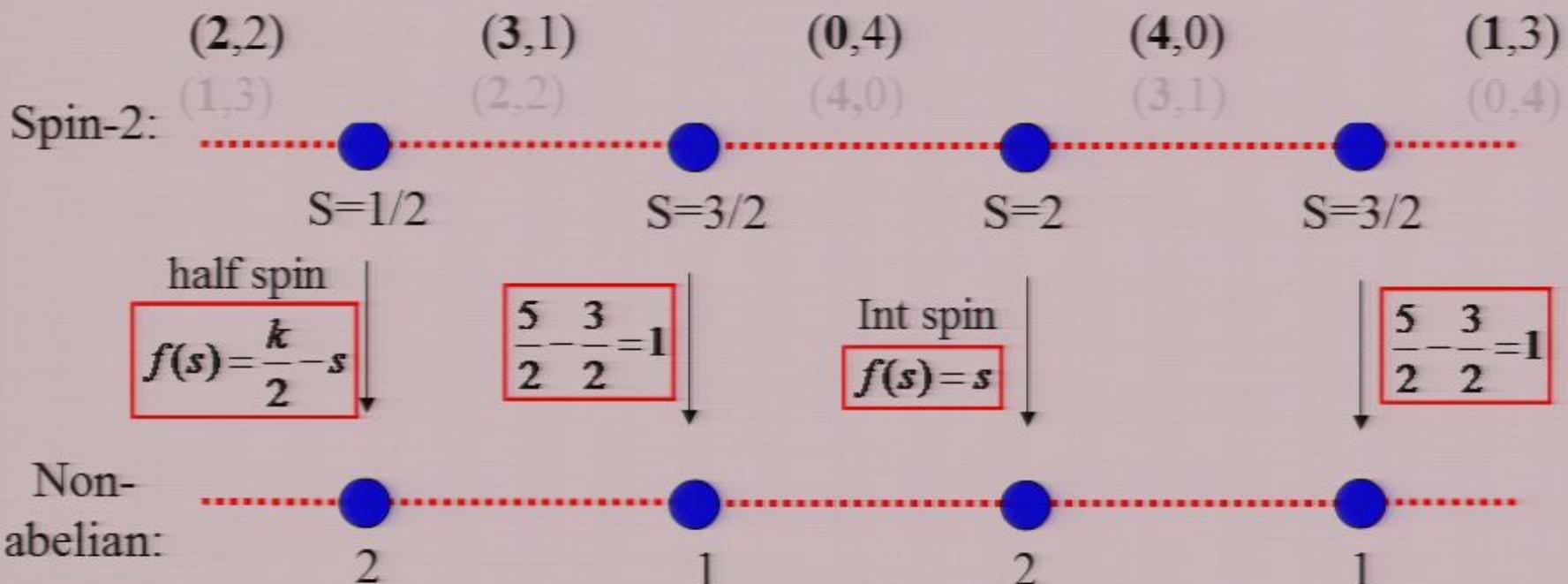
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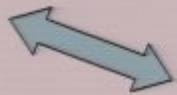
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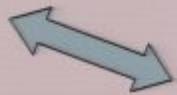
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No relevant perturbations of the non-Abelian permutation symmetric IR points!

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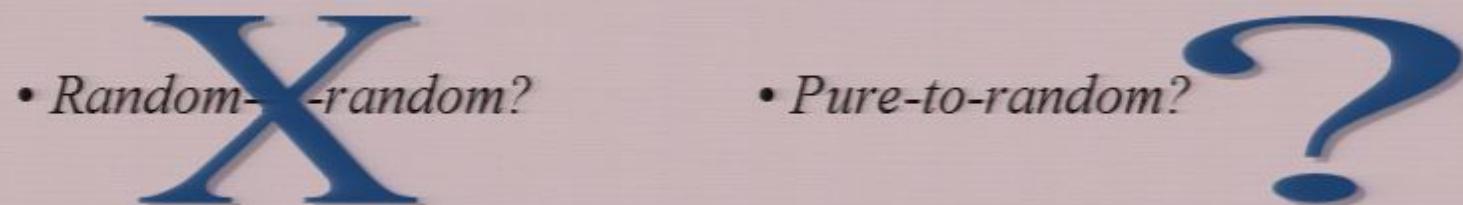
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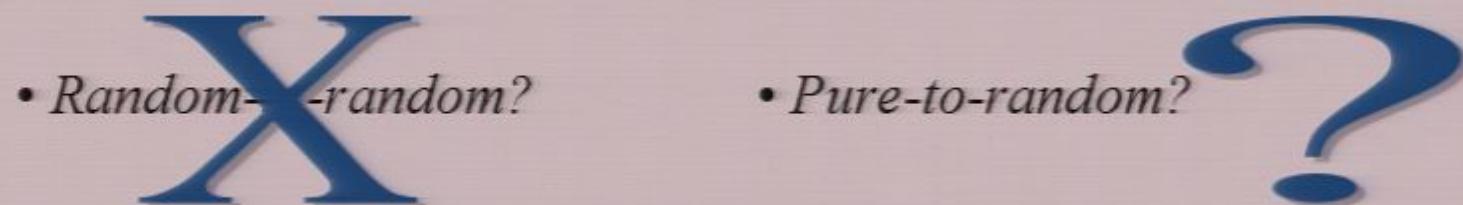
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- Mapping to the Abelian spin- $(k-1)/2$ chains, proves stability.

Entanglement and CFT's central charge

