

Title: Does entanglement persist at the macroscopic level?

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Abstract: The quantum states postulated to occur in situations of the 'Schroedinger's Cat' type are essentially N-particle GHZ states with N very large compared to 1, and their observation would thus be particularly compelling evidence for the ubiquity of the phenomenon of entanglement. However, in the traditional quantum measurement literature considerable scepticism has been expressed about the observability of this kind of 'macroscopically entangled' state, primarily because of the putatively disastrous effect on it of decoherence. In this talk I first examine why much of the literature has grossly overestimated the effects of decoherence, and then review the current experimental situation with respect to such states, as they (may) occur in fullerene diffraction, magnetic biomolecules, quantum-optical systems and Josephson devices; I also consider the prospects for their observation in nanomechanical systems. I conclude by reviewing and the theoretical implications of the experiments of the last decade in this area.

WHERE CAN WE FIND EVIDENCE  
IN EXPERIMENT FOR A HIGH DEGREE  
OF ENTANGLEMENT?

ANS: IN HIGH-N GHZ STATES!

$$\bar{\Psi} = \frac{1}{\sqrt{2}} \left\{ \underbrace{|\uparrow\uparrow\uparrow\uparrow\dots\uparrow\uparrow\uparrow\rangle}_N + \underbrace{|\downarrow\downarrow\downarrow\downarrow\dots\downarrow\downarrow\downarrow\rangle}_N \right\}$$

N PARTICLES "BEHAVING DIFFERENTLY"  
IN 2 BRANCHES OF SUPERPOS<sup>N</sup>,  $N \gg 1$

THESE ARE JUST "SCHRÖDINGER'S CAT"  
STATES!

⇒ FURTHER EFFORT TO BUILD SCHRÖDINGER'S  
CAT IN THE LABORATORY IS A PROBE OF  
ENTANGLEMENT.

TESTING QUANTUM MECHANICS  
TOWARDS THE LEVEL OF EVERYDAY LIFE:  
RECENT PROGRESS AND  
CURRENT PROSPECTS





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## MESO/MACROSCOPIC TESTS OF QM: MOTIVATION

At microlevel: (a)  $|\uparrow\rangle + |\downarrow\rangle$  quantum superposition ✓  
 $\neq$  (b)  $|\uparrow\rangle$  OR  $|\downarrow\rangle$  classical mixture ✗

how do we know? Interference

At macrolevel: (a)  +  quantum superposition }  
 OR (b)  OR  macrorealism }

⚡ Decoherence DOES NOT reduce (a) to (b)!

Can we tell whether (a) or (b) is correct?

Yes, if and only if they give different experimental predictions. But if decoherence  $\rightarrow$  no interference, then predictions of (a) and (b) identical.

$\Rightarrow$  must look for QIMDS

↑  
 quantum interference of macroscopically distinct states

What is “macroscopically distinct”?

(a) “extensive difference”  $\Lambda$

(b) “disconnectivity” D

↑  
 ~large number of particles behave differently in two branches

Initial aim of program: interpret raw data in terms of QM.  
 test (a) vs (b).

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## WHY HAS (MUCH OF) THE QUANTUM MEASUREMENT LITERATURE SEVERELY OVERESTIMATED DECOHERENCE?

("electron-on-Sirius" argument:  $\Delta\epsilon \sim a^{-N} \sim \exp - N \leftarrow \sim 10^{23}$ )

$\Rightarrow$  Just about any perturbation  $\gg \Delta\epsilon \Rightarrow$  decoherence)

1. Matrix elements of S-E interaction couple only a very restricted set of levels of S.
2. "Adiabatic" ("false") decoherence:

Ex.: spin-boson model

$$\hat{H} = \hat{H}_S + \hat{H}_E + \hat{H}_{S-E}$$

$$\hat{H}_S = \Delta\sigma_x$$

$$\hat{H}_E = \text{set of SHO's with lower frequency cutoff } \omega_{\text{min}} \gg \Delta$$

$$\hat{H}_{S-E} = \hat{\sigma}_z \sum_{\alpha} C_{\alpha} \hat{x}_{\alpha} \leftarrow \text{oscillator coords.}$$

$$\Psi_{\text{in}}(t=0) = |+\rangle |\chi_{+}\rangle \leftarrow \text{displaced state of oscillation}$$

$$\hat{\rho}_S(t=0) = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \text{ (trivially)}$$

$\Downarrow$

$$\Psi_{\text{in}}(t \sim \hbar / \Delta_{\text{min}}) \cong \frac{1}{\sqrt{2}} (|+\rangle |\chi_{+}\rangle + |-\rangle |\chi_{-}\rangle),$$

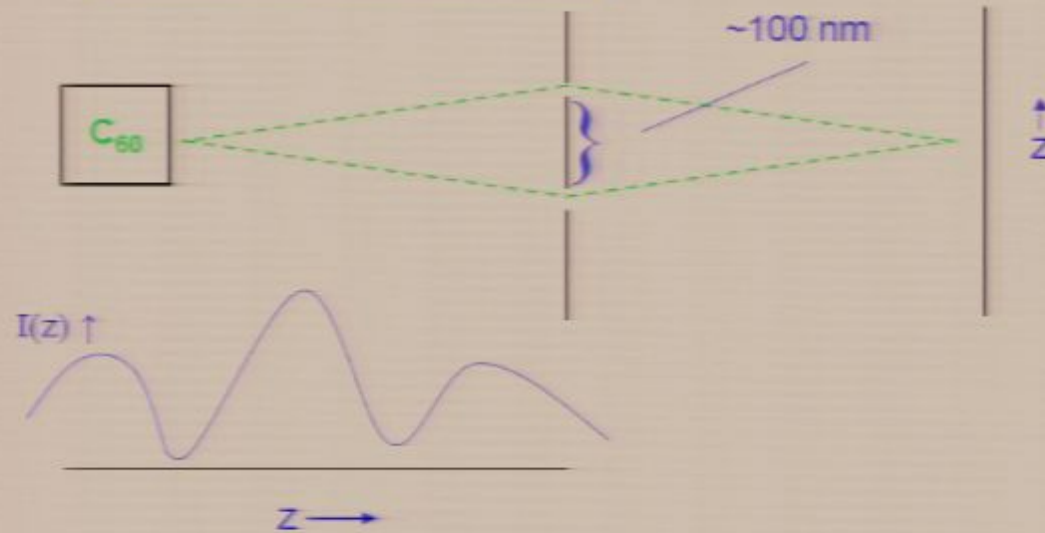
$$\langle \chi_{+} | \chi_{-} \rangle = \exp - F \cong 0 \quad \text{FC factor}$$

$$\Rightarrow \hat{\rho}_S(t \sim \hbar / \Delta_{\text{min}}) \cong \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$$

decohered?? (cf neutron interferometer)

## The Search for QIMDS

### I. Molecular diffraction\*



Note: (a.) Beam does not have to be monochromated

$$f(\nu) = A\nu^3 \exp\left[-(\nu - \nu_0)^2 / \nu_m^2\right] \quad (\nu_0 \sim 1.8 \nu_m)$$

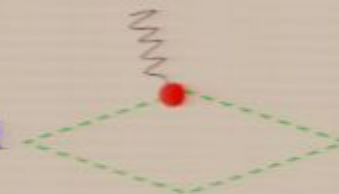
(b.) “Which-way” effects?

Oven is at 900–1000 K

$\Rightarrow$  many vibrational modes excited

4 modes infrared active  $\Rightarrow$

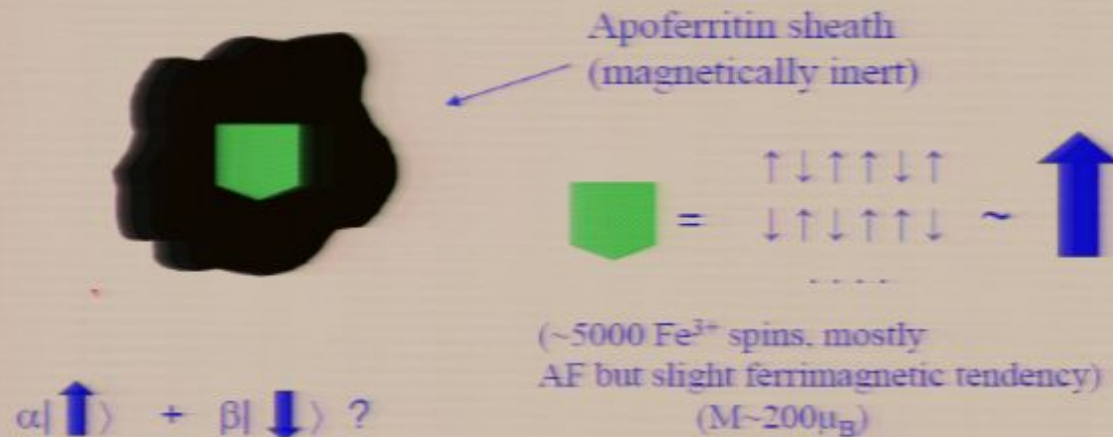
absorb/emit several radiation quanta on passage through apparatus!



**Why doesn't this destroy interference?**

## The Search for QIMDS (cont.)

### 2. Magnetic biomolecules\*



$$\text{AF: } \Delta \sim \hbar \omega_o \exp - N \sqrt{K/J}$$

no. of spins
uniaxial anisotropy
(isotropic)  
exchange en.

Raw data:  $\chi(\omega)$  and noise spectrum

above ~200 mK, featureless

below ~300 mK, sharp peak at ~ 1 MHz ( $\omega_{\text{res}}$ )

$$\omega_{\text{res}}^2 \cong \omega_o^2 + M^2 H^2$$

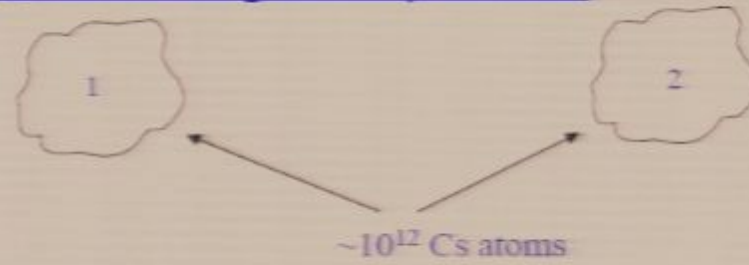
$$\ln \omega_o \sim a - bN \quad \leftarrow \text{no. of spins, exptly. adjustable}$$

Nb: data is on **physical** ensemble, i.e., only total magnetization measured.



## The Search for QIMDS (cont.)

### 3. Quantum-optical systems\*



for each sample separately, and also for total

$$\begin{aligned} & [J_x, J_y] = iJ_z \\ \Rightarrow & \langle \delta J_{x1} \delta J_{y1} \rangle \geq |J_{z1}| \\ & \langle \delta J_{x2} \delta J_{y2} \rangle \geq |J_{z2}| \\ & \langle \delta J_{x \text{ tot}} \delta J_{y \text{ tot}} \rangle \geq |J_{z \text{ tot}}| \end{aligned}$$

so, if set up a situation s.t.

$$J_{z1} = -J_{z2}$$

must have

$$\begin{aligned} \langle \delta J_{x1} \delta J_{y1} \rangle &> 0 \\ \langle \delta J_{x2} \delta J_{y2} \rangle &> 0 \end{aligned}$$

but may have

$$\langle \delta J_{x \text{ tot}} \delta J_{y \text{ tot}} \rangle = 0$$

(anal. of EPR)

Interpretation of idealized expt. of this type:

$$(\text{QM theory} \Rightarrow) \quad \langle \delta J_{x1} \delta J_{y1} \rangle \geq |J_{z1}| \sim N$$

$$\Rightarrow |\delta J_{x1}| \gtrsim N^{1/2}$$

But,

$$(\text{expt} \Rightarrow) \quad \langle \delta J_{x\text{tot}} \delta J_{y\text{tot}} \rangle \cong 0 \quad (\#)$$

$$\Rightarrow |\delta J_{x\text{tot}}| \sim 0$$

$$\Rightarrow \delta J_{x1} \text{ exactly anticorrelated with } \delta J_{x2}$$

$\Rightarrow$  state is either superposition or mixture of  $|n, -n\rangle$

but mixture will not give (#)

$\Rightarrow$  state must be of form

value of  $J_{x1}$       value of  $J_{x2}$

$$\sum_n c_n |n, -n\rangle$$

with appreciable weight for  $n \leq N^{1/2}$ .  $\Rightarrow$  high disconnectivity

Note:

(a) QM used essentially in argument

(b)  $D \sim N^{1/2}$  not  $\sim N$ .

(prob. generic to this kind of expt.)

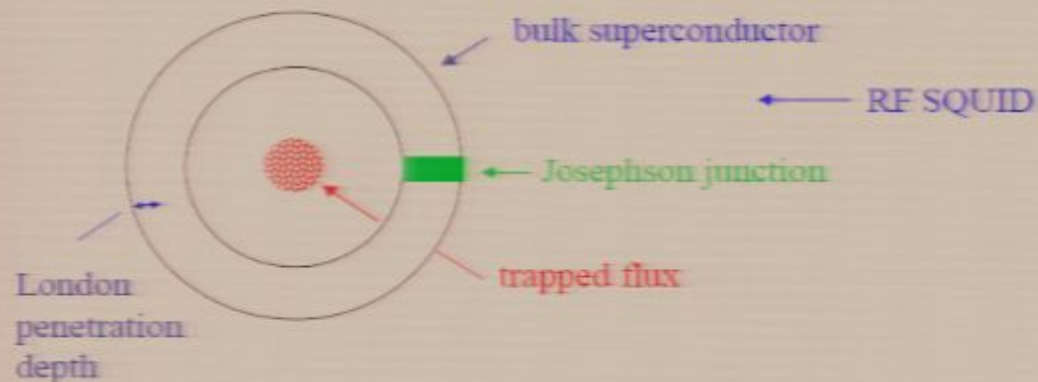
## The Search for QIMDS (cont.)

### 4. Superconducting devices

( $\nabla$ : not all devices which are of interest for quantum computing are of interest for QIMDS)

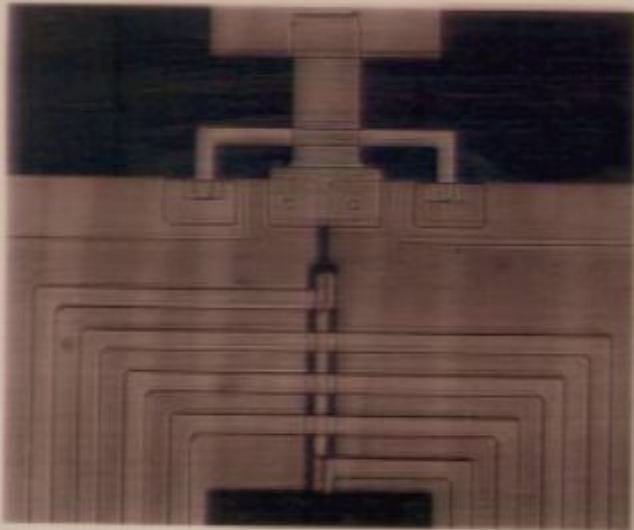
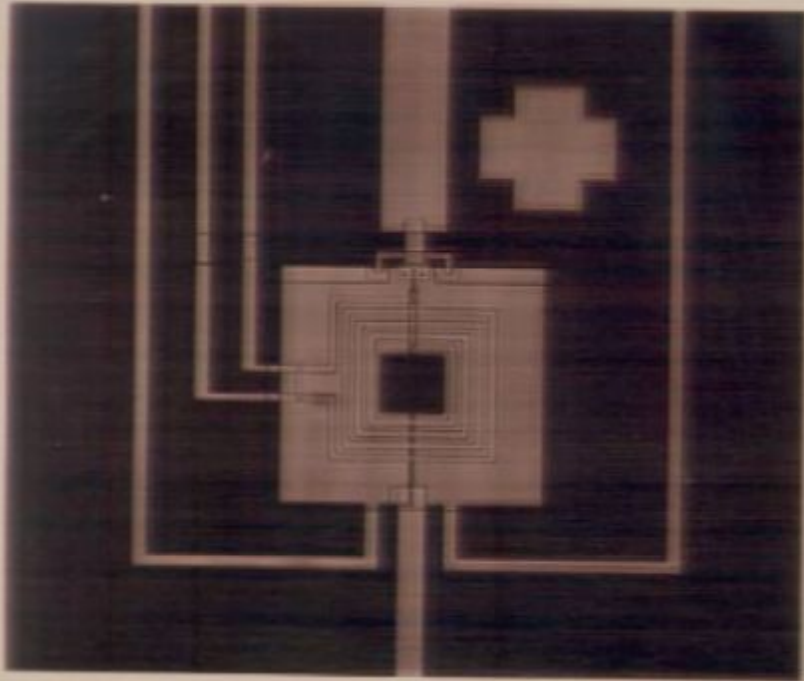
Advantages:

- classical dynamics of macrovariable v. well understood
- intrinsic dissipation (can be made) v. low
- well developed technology
- (non-) scaling of  $S$  (action) with  $D$ .



“Macroscopic variable” is trapped flux  $\Phi$   
[or circulating current  $I$ ]

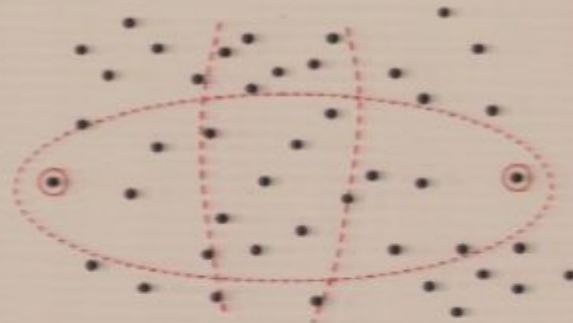
← 0.5 mm →



Pairing of electrons:



"di-electronic molecules"



Cooper Pairs

In simplest ("BCS") theory, Cooper pairs, once formed, must automatically undergo Bose condensation!

⇒ must all do exactly the same thing at the same time (also in nonequilibrium situation)

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## SUPERCONDUCTING RING IN EXTERNAL MAGNETIC FLUX:



$$E \propto K^2$$

Quantization condition for "particle" of charge  $2e$  (Cooper pair):

$$K \equiv \oint \mathbf{v} \cdot d\mathbf{l} = \frac{\hbar}{2m} (n - \Phi/\Phi_0)$$

integer  
"flux quantum"  
 $\hbar/2e$

A.  $\Phi = 0$ : groundstate unique ( $n = 0$ )

$\Rightarrow$  all pairs at rest.

B.  $\Phi = 1/2 \Phi_0$ : groundstate doubly degenerate:

( $n = 0$  or  $n = 1$ )



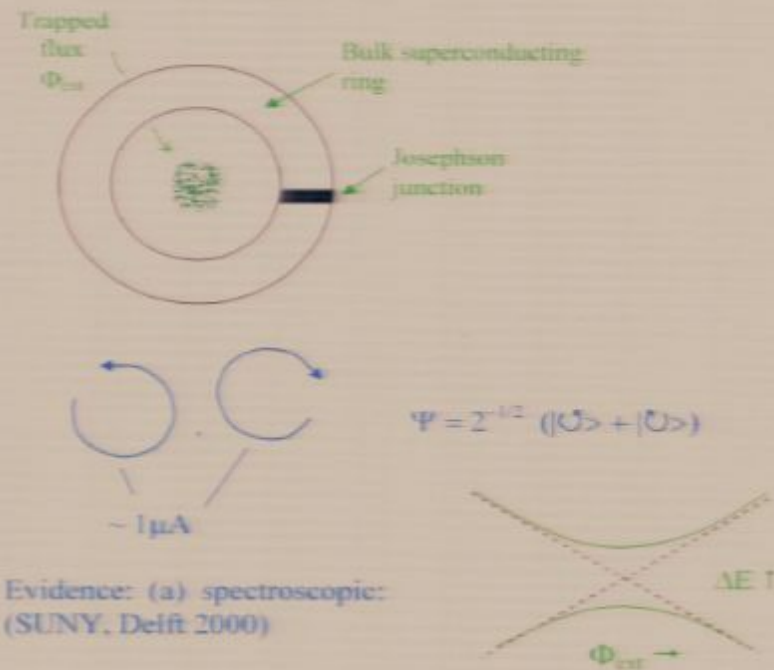
Either **all** pairs rotate **clockwise**:

Or **all** pairs rotate **anticlockwise**:

Note: state with 50% ↙ and 50% ↗

**strongly forbidden by energy considerations**

## Josephson circuits

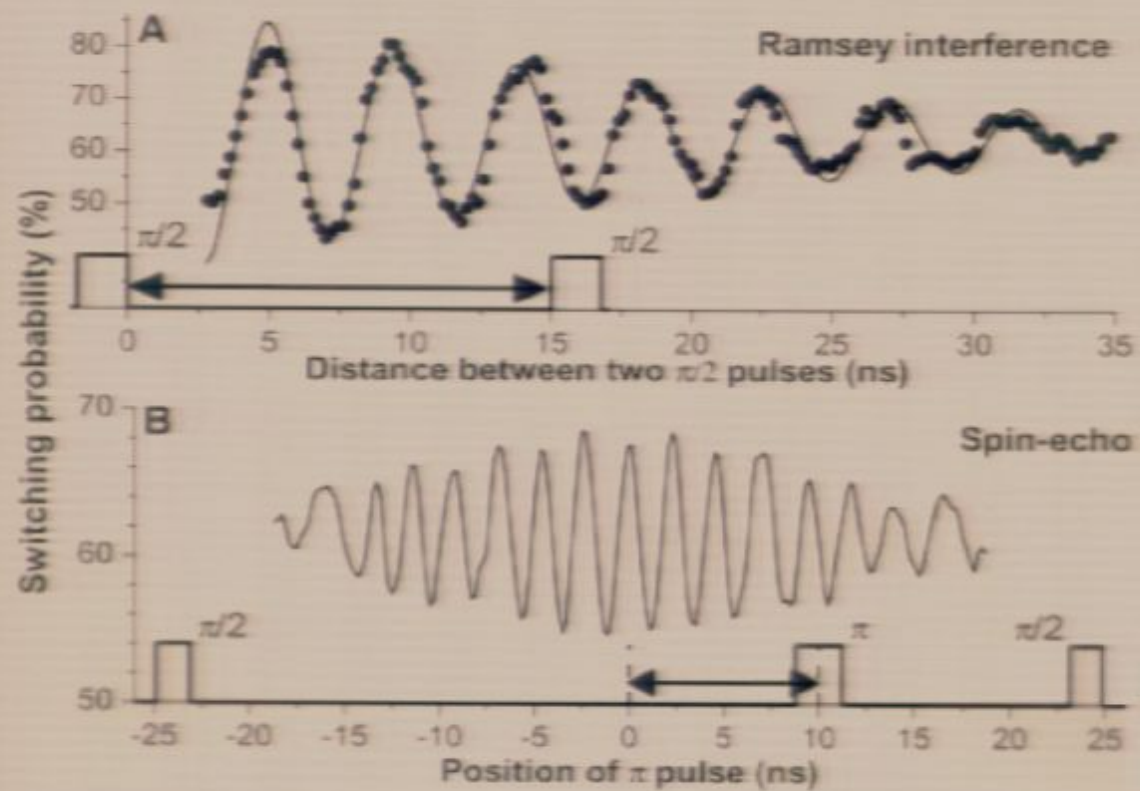


Evidence: (a) spectroscopic:  
(SUNY, Delft 2000)

(b) real-time oscillations (like  $NH_3$ )

between  $\downarrow$  and  $\uparrow$

(Saclay 2002, Delft 2003) ( $Q_0 \sim 50-350$ )



From I. Chiorescu, Y. Nakamura, C.J.P. Harmans, and J. E. Mooij, *Science*, **299**, 1869 (2003)



WHAT IS THE DISCONNECTIVITY “D” (“SCHRÖDINGER’S-CATTINESS”) OF THE STATES OBSERVED IN QIMDS EXPERIMENTS?

i.e., how many “microscopic” entities are “behaving differently” in the two branches of the superposition?

Fullerene (etc.) diffraction experiments: straightforward, number of “elementary” particles in  $C_{60}$  (etc.) ( $\sim 1200$ )

Magnetic biomolecules: number of spins which reverse between the two branches ( $\sim 5000$ )

Quantum-optical experiments }  
SQUIDS } matter of definition

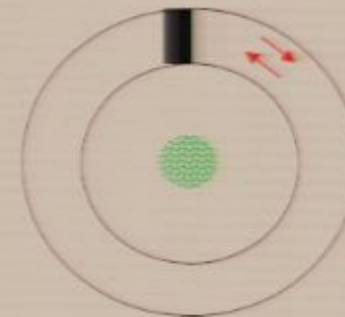
e.g. SQUIDS (SUNY experiment):

(a) naïve approach: no. of C. pairs

$$\Psi_{\sim} \sim \chi_{\sim}^{N/2}, \Psi_{\sigma} = \chi_{\sigma}^{N/2}$$

mutually orthogonal C. pair w.f.

$$\Rightarrow D \sim N \sim 10^9 - 10^{10} \quad \Delta: \text{Fermi statistics!}$$



(b) how many single electrons do we need to displace in momentum space to get from  $\Psi_{\sim}$  to  $\Psi_{\sigma}$ ? (Korsbakken et al., preprint, Nov. 08)

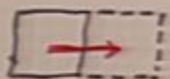
$$\Rightarrow D \sim N(v_x/v_F) \sim 10^3 - 10^4 \quad \Delta: \text{intuitively, severe underestimate in “BEC” limit (e.g. Fermi alkali gas)}$$

(c) macroscopic eigenvalue of 2-particle density matrix (corresponding to (fairly) orthogonal states in 2 branches):

$$\Rightarrow D \sim N(\Delta/\varepsilon_F) \sim 10^6 - 10^7$$



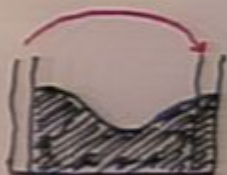
observation time  
 $\sim 1 \text{ sec}$



LiF



$v = 0$



$v = v_0$

$$W_{cl} \equiv \sum_r |\delta n_r| \sim v_0/v_F$$

What about nuclei?

Prime facie,

$$v_{nuc} \sim v_0/v_n \ll v_0/v_F$$

*rms nuclear velocity*

But - should think in terms of individual  
nucleons within nucleus! then

$$W_{nuc} \sim v_0/v_p \ll v_0/v_F!$$

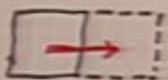
*velocity of proton (tc)*

If that's "cheating", is 1-electron calc?  
for SQWID.

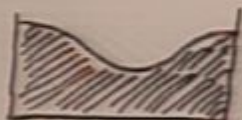
$$\text{Real FOM} \sim N(\Delta/\epsilon_F) \sim 10^6 - 10^7?$$



observation time  
 $\sim 1 \text{ sec}$



LiF



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$$W_{nuc} \sim v_0/v_p \ll v_0/v_F!$$

velocity of proton ( $\approx$ )

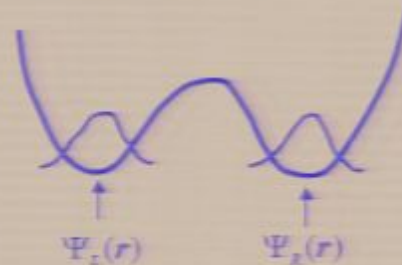
If that's "cheating", is 1-electron case?  
for squid.

$$\text{Real FOM} \sim N(\Delta/E_F) \sim 10^6 - 10^7?$$

More possibilities for QIMDS:

(a) BEC's of ultracold alkali gases:

Bose-Einstein condensates



(Gross-Pitaevskii)

Ordinary GP state:

$$\Psi_N = (a\psi_L(r) + b\psi_R(r))^N$$

“Schrödinger-cat” state (favored if interactions attractive):

$$\Psi_N = a(\psi_L(r))^N + b(\psi_R(r))^N$$

problems:

(a) extremely sensitive to well asymmetry  $\Delta E$   
 (energy stabilizing arg (a/b)  $\sim t^N \sim \exp - NB/\hbar$ )  
 so  $\Delta E$  needs to be  
**exp'ly** small in N

↑  
 single-particle tunnelling  
 matrix element

(b) detection: tomography unviable for  $N \gg 1$ ,  
 $\Rightarrow$  need to do time-sequence experiments (as in SQUIDS), but  
 period v. sensitive e.g. to exact value of N

## WHAT HAVE WE SEEN SO FAR?

1. If we interpret raw data in QM terms, then can conclude we have a **quantum superposition rather than a mixture** of meso/macroscopically distinct states.

However, "only 1 degree of freedom involved."

2. Do data exclude **general** hypothesis of macrorealism?

**NO**

3. Do data exclude **specific** macrorealistic theories?

e.g. GRWP ← **Ghirardi, Rimini, Weber, Pearle**

**NO** (fullerene diffraction: N not large enough, SQUIDS:  
no displacement of COM between branches)

Would MEMS experiments (if in agreement with QM) exclude GRWP?

alas:  $\Gamma_{coll} \propto \Delta x$ ,  $\Gamma_{dec} \propto (\Delta x)^2$

collapse rate in GRWP theory      decoherence rate acc. to QM

⇒ do not gain by going to larger  $\Delta x$   
(and small  $\Delta x$  may not be enough to test GRWP)

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## HOW CONFIDENT ARE WE ABOUT (STANDARD QM'I) DECOHERENCE RATE?

Theory:

- (a) model environment by oscillator bath (may be questionable)
- (b) Eliminate environment by standard Feynman-Vernon type calculation (seems foolproof)

Result (for SHO):

$$\Gamma_{dec} \sim \Gamma \left( \frac{k_B T}{\hbar \Omega} \right) \cdot \left( \frac{\Delta x}{x_0} \right)^2$$

provided  $k_B T \gg \hbar \Omega$

energy relaxation rate ( $\Omega/Q$ )

zero-point rms displacement

ARE WE SURE THIS IS RIGHT?

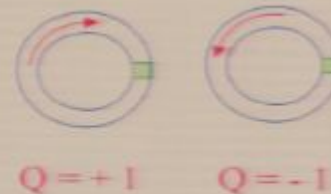
Tested (to an extent) in cavity QED: never tested (?) on MEMS.

Fairly urgent priority!

Where do we go from here?

1. Larger values of  $\Lambda$  and/or  $D$ ?  
(Diffraction of virus?)
2. Alternative Dfs. of "Measures" of Interest
  - More sophisticated forms of entanglement?\*
  - Biological functionality (e.g. superpose states of rhodopsin?)
  - Other (e.g. GR)
- 3. Exclude Macrorealism

Suppose: **Whenever observed,**  $Q = \pm 1$ .



Df. of "MACROREALISTIC" Theory:

- "COMMON SENSE"?
- I.  $Q(t) = \pm 1$  at (almost)  $\forall t$ ,  
whether or not observed.
  - II. Noninvasive measurability
  - III. Induction

Can test with existing SQUID Qubits!



Df:

$$K \equiv K(t_1, t_2, t_3, t_4) \equiv \langle \underline{Q}(t_1) \underline{Q}(t_2) \rangle + \langle \underline{Q}(t_2) \underline{Q}(t_3) \rangle \\ + \langle \underline{Q}(t_3) \underline{Q}(t_4) \rangle - \langle \underline{Q}(t_1) \underline{Q}(t_4) \rangle$$

Take  $t_2 - t_1 = t_3 - t_2 = t_4 - t_3 = \pi/4\Delta$  ← tunnelling frequency

Then,

- (a) Any macrorealistic theory:  $K < 2$
- (b) Quantum mechanics, ideal:  $K = 2.8$
- (c) Quantum mechanics, with all the real-life complications:  $K > 2$  (but  $< 2.8$ )

Thus: to extent analysis of (c) within quantum mechanics is reliable, **can force nature to choose between** macrorealism and quantum mechanics!

Possible outcomes:

- (1) Too much noise  $\Rightarrow K_{QM} < 2$
- (2)  $K > 2 \Rightarrow$  macrorealism refuted
- (3)  $K < 2$ : ?!