Title: The Galileon as a local modification of gravity

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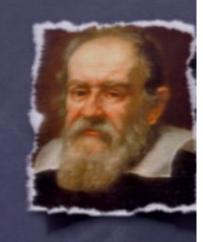
URL: http://pirsa.org/10040100

Abstract: TBA

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Alberto Nicolis Columbia University

The Galileon as a local modification of gravity



with R.Rattazzi and E.Trincherini

"local"

local 4D QFT

"local"

local 4D QFT

"local"

d << Hubble

"Beautiful" (?)

- "Beautiful" (?)
- Unique EFT for low-energy gravitons

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... Still:

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the C.C. problem

- "Beautiful" (?)
- Unique EFT for low-energy gravitons

... Still:

- the C.C. problem
- the present acceleration

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Difficult!

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Difficult!

Assuming:

Difficult!

Assuming:

spin 2

Difficult!

Assuming:

spin 2

massless

Difficult!

Assuming:

- spin 2
- massless
- Lorentz invariance

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(ca. 1965)

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Fierz-Pauli (bad)

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Fierz-Pauli

(bad)

ghost condensate

(good)

Difficult!

Assuming:

- spin 2
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 - 201011

->

Fierz-Pauli

ghost

condensate

(bad)

(good)

locality

Difficult!

Assuming:

- spin 2
- massiess
- Lorentz invariance

- @ locality
- DGP (5D)

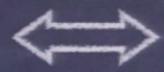
Fierz-Pauli (bad)

ghost (good) condensate

(good and bad)

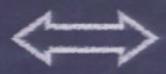
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Deviations from GR d << Hubble



effective light scalar dof mixed with metric

Deviations from GR d << Hubble



effective light scalar dof mixed with metric

$$h_{\mu\nu} \ll 1$$

$$\mathcal{L} \sim (1 - 2\pi)R + \mathcal{L}_{\pi} + h_{\mu\nu}T^{\mu\nu}$$

Deviations from GR d << Hubble



effective light scalar dof mixed with metric

$$h_{\mu\nu} \ll 1$$

$$\mathcal{L} \sim (1 - 2\pi)R + \mathcal{L}_{\pi} + h_{\mu\nu}T^{\mu\nu}$$

demix

$$\Rightarrow \mathcal{L} \sim \mathcal{L}_{GR} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$

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$$\mathcal{L}_{DGP} = \mathcal{L}_{GR} + (\partial \pi)^2 + (\partial \pi)^2 \Box \pi + \pi T^{\mu}_{\mu}$$

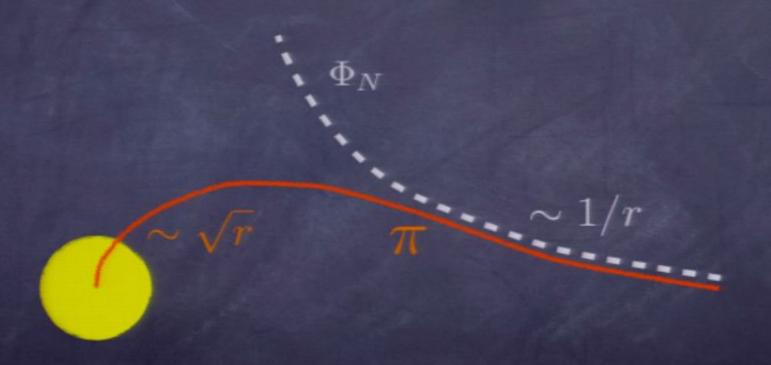
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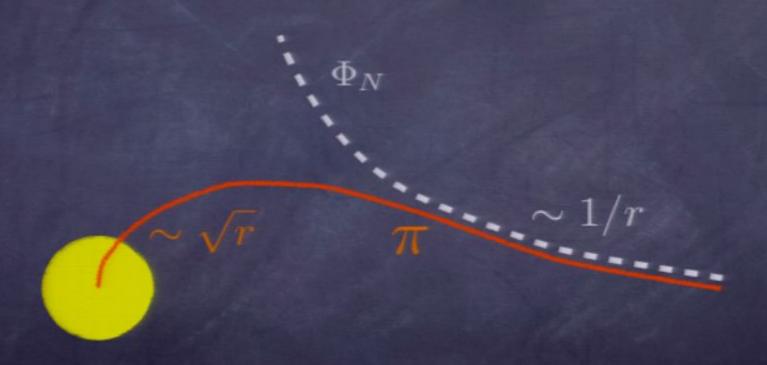
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"Vainshtein effect"

No conflict w/ solar system

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No conflict w/ solar system

Nice observational predictions (LLR)

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Nice observational predictions (LLR)

same non-linearity



self-accelerating cosmology

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self-accelerating cosmology

$$T_{\mu\nu} = 0 \quad \Rightarrow \quad \begin{cases} \pi = 0 \\ \pi = H^2 x_{\alpha} x^{\alpha} \end{cases}$$

- No conflict w/ solar system
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self-accelerating cosmology

$$T_{\mu\nu} = 0 \quad \Rightarrow \quad \left\{ \begin{array}{l} \pi = 0 \\ \pi = H^2 x_{\alpha} x^{\alpha} \end{array} \right.$$

$$h_{\mu\nu}^{
m phys} = h_{\mu\nu} + \pi\,\eta_{\mu\nu}$$
 \Longrightarrow deSitter (locally)

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$$\pi_{dS} = H^2 \, x_{\alpha} x^{\alpha}$$

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small fluctuations

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$$\mathcal{L}_{\pi}[\pi_{dS} + \varphi] \rightarrow -(\partial \varphi)^2$$

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solution unstable

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solution unstable

All this makes sense locally (d << Hubble)</p>

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At Hubble scales, the 4D description breaks down

All this makes sense locally (d << Hubble)</p>

At Hubble scales, the 4D description breaks down

Start seeing the 5th dimension (non-locality)

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Self-acceleration vs. stability

- Self-acceleration vs. stability
- For theories where:

Deviations from GR d << Hubble



effective light scalar dof mixed with metric

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- For theories where:

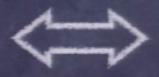
Deviations from GR d << Hubble



effective light scalar dof mixed with metric

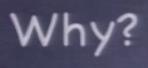
- Self-acceleration vs. stability
- For theories where:

Deviations from GR d << Hubble



effective light scalar dof mixed with metric

(i.e. universally coupled $\pi T^{\mu}{}_{\mu}$)



Examples above

- Examples above
- Simplest possibility

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- Examples above
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- Vector problematic:

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- Examples above
- Simplest possibility
- Vector problematic:

$$A_{\mu} \partial_{\nu} T^{\mu\nu} \rightarrow 0$$

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O(1) modifications to cosmology

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- o negligible effects in the solar system

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- stable self-accelerating solutions

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- O(1) modifications to cosmology
- negligible effects in the solar system
- stable self-accelerating solutions
- o do not assume global scalar field

Kinematics: FRW, locally = weak $h_{\mu\nu}$

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Convenient "conformal" gauge

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Convenient "conformal" gauge

$$ds^2 \simeq \left[1 - \frac{1}{2}H^2\vec{x}^2 + \frac{1}{2}(2\dot{H} + H^2)t^2\right](-dt^2 + d\vec{x}^2)$$

$$\mathcal{L}_{GR} \rightarrow \mathcal{L}_{GR} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$

$$\mathcal{L}_{\mathrm{GR}} \; o \; \mathcal{L}_{\mathrm{GR}} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$
 unspecified, so far

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Allow for generic non-linearities in \mathcal{L}_{π}

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 unspecified, so far

Allow for generic non-linearities in \mathcal{L}_{π} Keep lowest order in $h_{\mu\nu}$

Dynamics: Introduce the scalar

$$\mathcal{L}_{\mathrm{GR}}
ightarrow \mathcal{L}_{\mathrm{GR}} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$
 unspecified, so far

Allow for generic non-linearities in $\,\mathcal{L}_{\pi}\,$

Keep lowest order in $\,h_{\mu
u}$

"Physical" Metric:

$$\hat{h}_{\mu\nu} = h_{\mu\nu} + 2\pi \, \eta_{\mu\nu}$$

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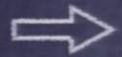
$$\pi$$
 's stress-energy

$$\pi T^{\mu}_{\mu}$$

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$$\pi$$
 's stress-energy

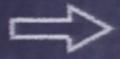
$$\pi T^{\mu}_{\mu}$$



Neglect π 's gravitational backreaction

$$\pi$$
 's stress-energy

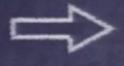
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Neglect π 's gravitational backreaction (check!)

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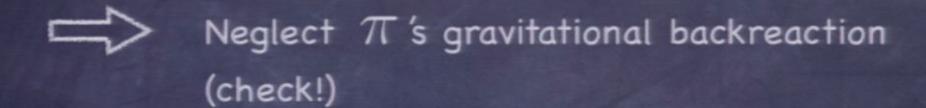


Neglect π 's gravitational backreaction (check!)

$$\mathcal{L} = \mathcal{L}_{GR} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$

$$\pi$$
 's stress-energy

$$\pi T^{\mu}_{\mu}$$



$$\mathcal{L} = \mathcal{L}_{GR} + \mathcal{L}_{\pi} + \pi T^{\mu}{}_{\mu}$$

For given $T_{\mu\nu}$ (and bdy. conditions)

$$h_{\mu\nu} = h_{\mu\nu}^{\rm GR}$$

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$$\pi \eta_{\mu\nu} = \frac{1}{2} (\hat{h}_{\mu\nu} - h_{\mu\nu}^{GR})$$

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full metric seen by matter

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$$\pi \to \pi + c + b_{\mu} x^{\mu}$$

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Compare curvature tensors

$$\pi\eta_{\mu\nu}=rac{1}{2}(\hat{h}_{\mu\nu}-h^{\mathrm{GR}}_{\mu\nu})_{\mathrm{conformal}}$$
 full metric seen by matter

- What gauge?
- Not unique: conformal transformations

$$\pi \to \pi + c + b_{\mu} x^{\mu}$$

Compare curvature tensors

$$\hat{R}_{\mu\nu} = R_{\mu\nu} - 2\partial_{\mu}\partial_{\nu}\pi - \eta_{\mu\nu}\Box\pi$$

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$$\pi = (\Phi_{\mathrm{DGP}} - \Phi_{\mathrm{GR}})_{\mathrm{conformal}}$$

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$$= -\frac{1}{4} \Delta H^2 \vec{x}^2 + \frac{1}{4} (2\Delta \dot{H} + \Delta H^2) t^2 + c + b t$$

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Claim: brane-bending mode of Luty, Porrati, Rattazzi

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Claim: brane-bending mode of Luty, Porrati, Rattazzi

Check: plug π into their EOM

$$\mathcal{L}_{\pi} = ?$$

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negligible effects in solar system

$$\mathcal{L}_{\pi} = ?$$

VS.

negligible effects in solar system



need non-linearities

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need non-linearities

IF:

$$\frac{\delta \mathcal{L}_{\pi}}{\delta \pi} + T^{\mu}{}_{\mu}$$

$$\mathcal{L}_{\pi} = ?$$

VS.

negligible effects in solar system



need non-linearities

IF:

$$\frac{\delta \mathcal{L}_{\pi}}{\delta \pi} + T^{\mu}{}_{\mu} = F(\Delta R_{\mu\nu}) + T^{\mu}{}_{\mu}$$

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VS.

negligible effects in solar system



need non-linearities

IF:

$$\frac{\delta \mathcal{L}_{\pi}}{\delta \pi} + T^{\mu}{}_{\mu} = F(\Delta R_{\mu\nu}) + T^{\mu}{}_{\mu}$$

local matter distribution



corrections to geometry

$$\Delta R_{\mu\nu} \sim \partial_{\mu}\partial_{\nu}\pi$$

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$$\Delta R_{\mu\nu} \sim \partial_{\mu}\partial_{\nu}\pi$$

EOM:
$$\frac{\delta \mathcal{L}_{\pi}}{\delta \pi} = F(\partial_{\mu} \partial_{\nu} \pi) = -T^{\mu}{}_{\mu}$$

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$$\pi(x) \to \pi(x) + c + b_{\mu} x^{\mu}$$

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"Galilean invariance"

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(Analogous to
$$x(t)
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$$\pi(x) \to \pi(x) + c + b_{\mu} x^{\mu}$$

"Galilean invariance"

(Analogous to
$$x(t)
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2-derivative EOM

Galilean-invariant terms

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$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \, \pi^n$$

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Simplest:

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \, \pi^n$$

 \circ Simplest: $\mathcal{L}^{(1)}=\pi$

$$\mathcal{L}^{(1)} = \pi$$

Next to simplest:

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \, \pi^n$$

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$$\mathcal{L}^{(1)} = \pi$$

$$\mathcal{L}^{(2)} = (\partial \pi)^2$$

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \, \pi^n$$

$$\circ$$
 Simplest: $\mathcal{L}^{(1)}=\pi$

• Next to simplest:
$$\mathcal{L}^{(2)} = (\partial \pi)^2$$

Less trivial (DGP):

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \, \pi^n$$

$$\circ$$
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Invariance:

$$\mathcal{L}^{(n)} \sim \partial^{2n-2} \pi^n$$

$$\circ$$
 Simplest: $\mathcal{L}^{(1)}=\pi$

• Next to simplest:
$$\mathcal{L}^{(2)} = (\partial \pi)^2$$

• Less trivial (DGP):
$$\mathcal{L}^{(3)} = (\partial \pi)^2 \, \Box \pi$$

$$m{\mathcal{L}}^{(n)}
ightarrow \mathcal{L}^{(n)} + \partial_{\mu} G^{\mu}$$

Galilean Invariants

$$\mathcal{L}_{1} = \pi
\mathcal{L}_{2} = -\frac{1}{2} \partial \pi \cdot \partial \pi
\mathcal{L}_{3} = -\frac{1}{2} [\Pi] \partial \pi \cdot \partial \pi
\mathcal{L}_{4} = -\frac{1}{4} ([\Pi]^{2} \partial \pi \cdot \partial \pi - 2 [\Pi] \partial \pi \cdot \Pi \cdot \partial \pi - [\Pi^{2}] \partial \pi \cdot \partial \pi + 2 \partial \pi \cdot \Pi^{2} \cdot \partial \pi)$$

$$\mathcal{L}_{5} = -\frac{1}{5} ([\Pi]^{3} \partial \pi \cdot \partial \pi - 3 [\Pi]^{2} \partial \pi \cdot \Pi \cdot \partial \pi - 3 [\Pi] [\Pi^{2}] \partial \pi \cdot \partial \pi + 6 [\Pi] \partial \pi \cdot \Pi^{2} \cdot \partial \pi
+ 2 [\Pi^{3}] \partial \pi \cdot \partial \pi + 3 [\Pi^{2}] \partial \pi \cdot \Pi \cdot \partial \pi - 6 \partial \pi \cdot \Pi^{3} \cdot \partial \pi)$$
(38)

$$\Pi^{\mu}{}_{\nu} \equiv \partial^{\mu}\partial_{\nu}\pi \qquad [\cdots] \equiv \text{Tr}\{\cdots\}$$

In conclusion:

$$\mathcal{L}_{\pi} = \sum_{i=1}^{5} c_i \mathcal{L}_i$$

- Lagrangian very constrained: 5 coeffs.
- thorough analysis of highly symmetric solutions
- for dS and spherical sols.: algebraic problem

good news:

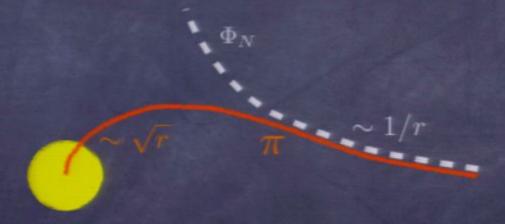
good news:

Stable dS self-accelerating solutions

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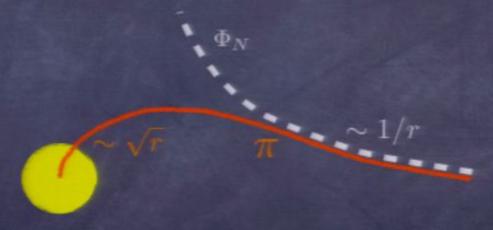
good news:

Stable dS self-accelerating solutions Stable Vainshtein-like spherical sols.



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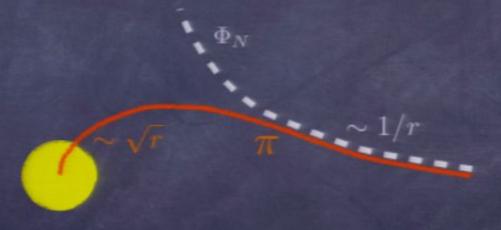
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bad news:

good news:

Stable dS self-accelerating solutions Stable Vainshtein-like spherical sols.



bad news:

Superluminal radial excitations

In fact superluminality very general

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ullet For linear π background soutions:

$$\delta c \sim \frac{1}{(LH_0)^2} \qquad L \gg H_0^{-1}$$

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$$\delta$$
 phase $\sim \frac{k}{H_0} \frac{1}{(LH_0)^2}$

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phase $\sim \frac{k}{H_0} \frac{1}{(LH_0)^2}$



Naively low strong-coupling scale

$$\frac{1}{\Lambda^{3n-6}}\partial\pi\partial\pi(\partial^2\pi)^{n-2}$$

$$\Lambda \sim (1000 \text{ km})^{-1}$$

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 \odot Consistent down to $\Lambda_{\rm eff} \sim {
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Naively low strong-coupling scale

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$$\Lambda \sim (1000 \text{ km})^{-1}$$

- ullet Consistent down to $\Lambda_{
 m eff}\sim {
 m cm}^{-1}$ on earth
- Problems w/ relativistic UV completion

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so far: d << Hubble

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so far: d << Hubble

ullet is it possible π \longrightarrow global scalar w/ local 4D action?

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DGP: no (physics is 5D)

so far: d << Hubble

ullet is it possible $\pi \to \mathrm{global}$ scalar w/ local 4D action?

DGP: no (physics is 5D)

here: still an open question

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 Galileon: local dynamics of general theories where

> Deviations from GR d << Hubble



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effective light scalar dof mixed with metric

Vainshtein effect demands Galilean invariance

 Galileon: local dynamics of general theories where

> Deviations from GR d << Hubble



- Vainshtein effect demands Galilean invariance
- 5 Lagrangian terms overall

 Galileon: local dynamics of general theories where

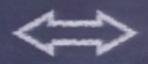
> Deviations from GR d << Hubble



- Vainshtein effect demands Galilean invariance
- 5 Lagrangian terms overall
- Stable, self-accelerating solutions

 Galileon: local dynamics of general theories where

> Deviations from GR d << Hubble



- Vainshtein effect demands Galilean invariance
- 5 Lagrangian terms overall
- Stable, self-accelerating solutions
- Stable Vainshtein-like spherical solutions

 Superluminal excitations = problems w/ relativistic UV completion

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 Superluminal excitations = problems w/ relativistic UV completion

low QM strong-coupling scale

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low QM strong-coupling scale

No satisfying global 4D completion yet

- Superluminal excitations = problems w/ relativistic UV completion
- low QM strong-coupling scale
- No satisfying global 4D completion yet
- Stable NEC violation (see next talk)