

Title: Kinetics of Microstructural Transitions in Block Copolymer Melts

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URL: <http://pirsa.org/10040084>

Abstract: Soft materials are dynamical by nature and the study of the dynamics of soft materials is an exciting, rich area of current interest. During macromolecular self-assembly, as occurs in block copolymers, long structural relaxation timescales due to collective molecular motion are often seen. How microstructure influences the dynamics, the existence and lifetime of metastable states, and the dynamics of long-lived non-equilibrium structures are all poorly-understood issues. I will discuss our dynamical simulations that address these questions in the context of the nucleation of one microstructured phase out of another in a block copolymer melt.

Kinetics of Microstructural Transitions in Block Copolymers

Robert Wickham

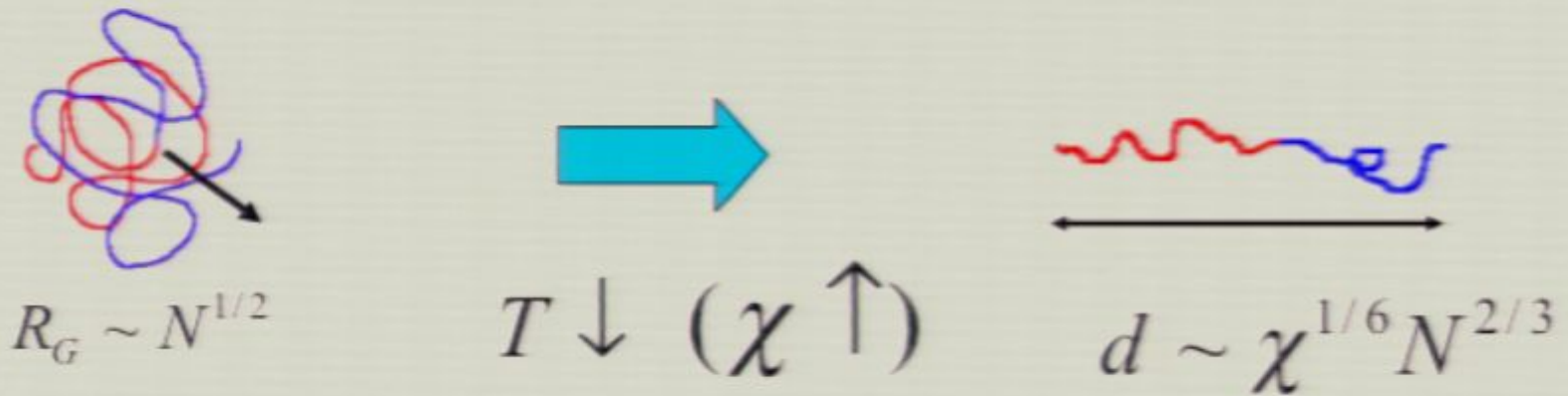
Department of Physics
University of Guelph



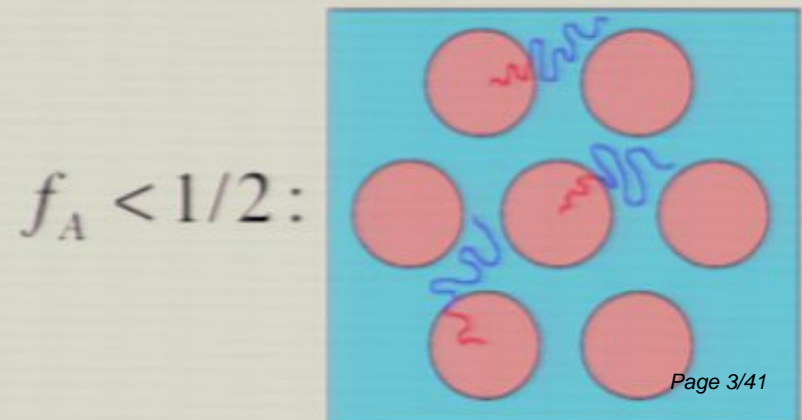
4-Corner Symposium
Perimeter Institute
April 22, 2010

Formation of Microstructures

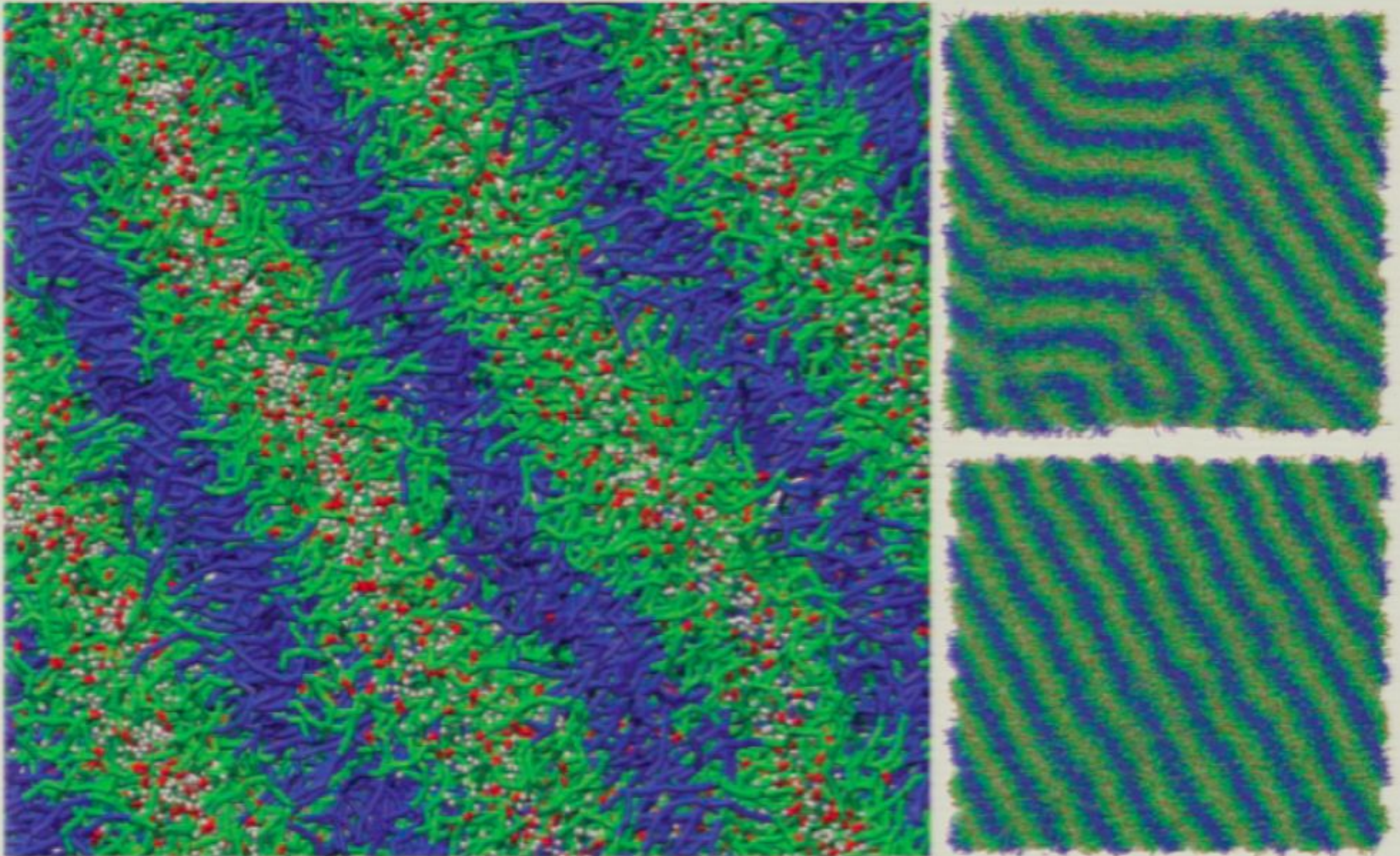
Competition between entropy (coil) and block repulsion (stretch)



Leads to microphase separation of diblock copolymers



From M. Klein and W. Shinoda, Science (2008)



Bulk diblock copolymer melt phase diagram

Self-consistent mean-field theory (SCMFT) - Matsen and Schick (1994):

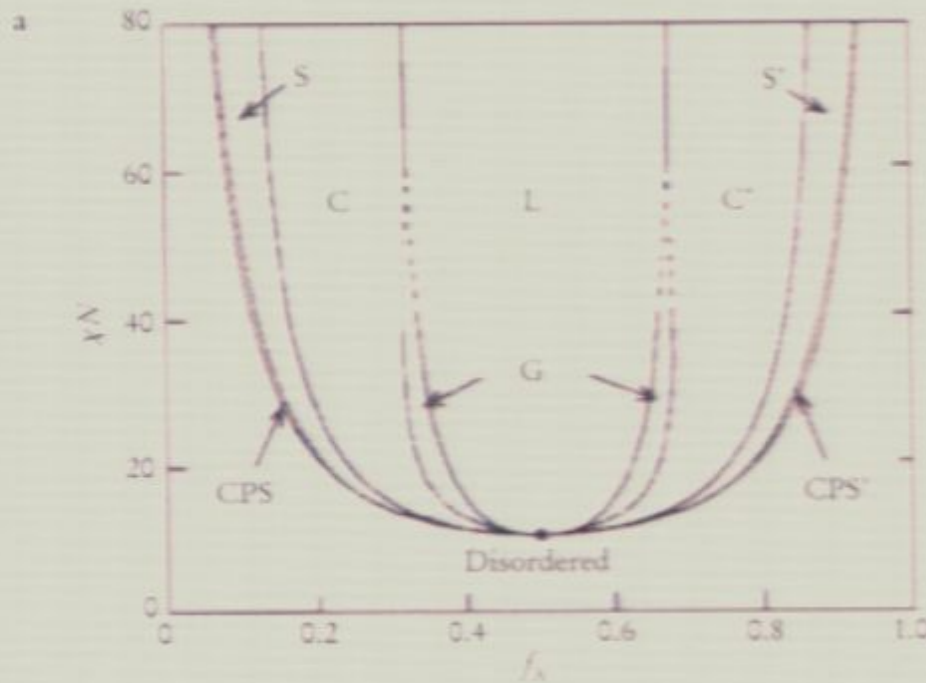
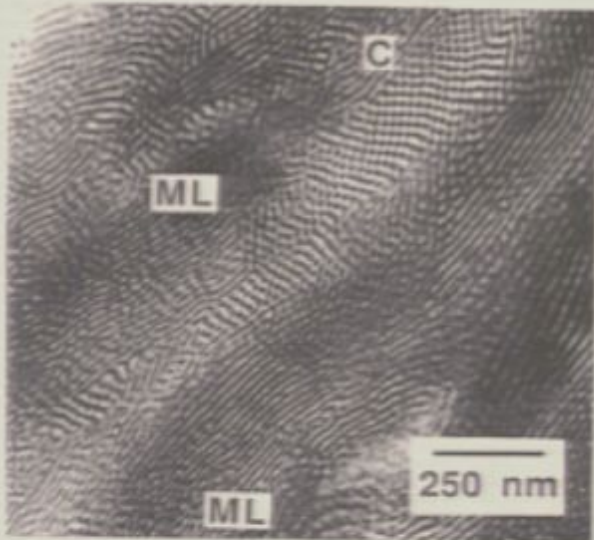


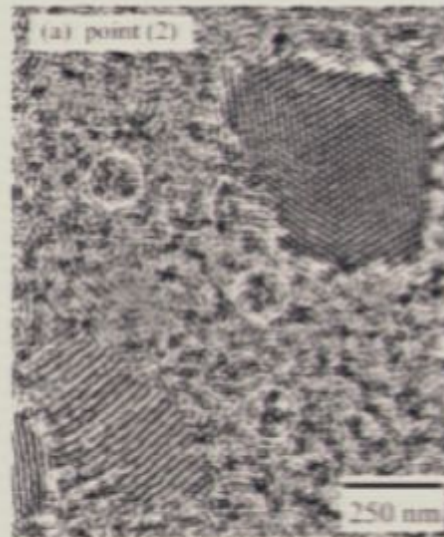
Diagram:
Fredrickson and
Bates (1999)



Reality is more complicated...



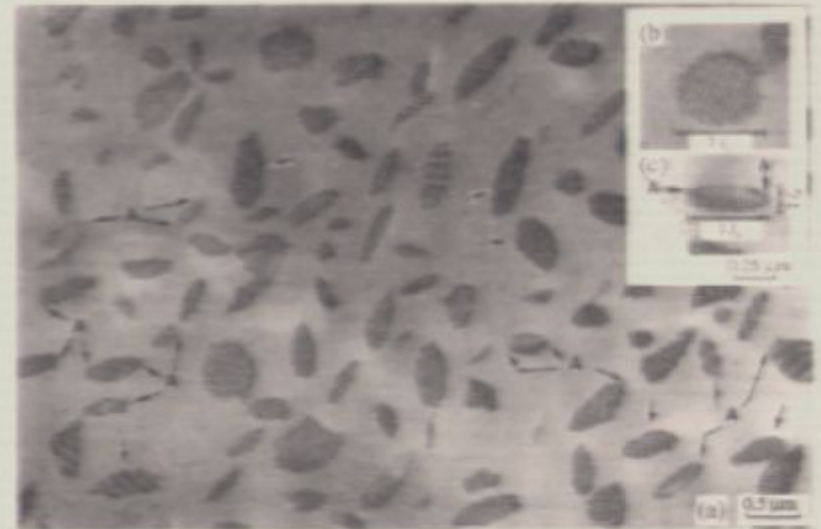
Hajduk et al. (1994)



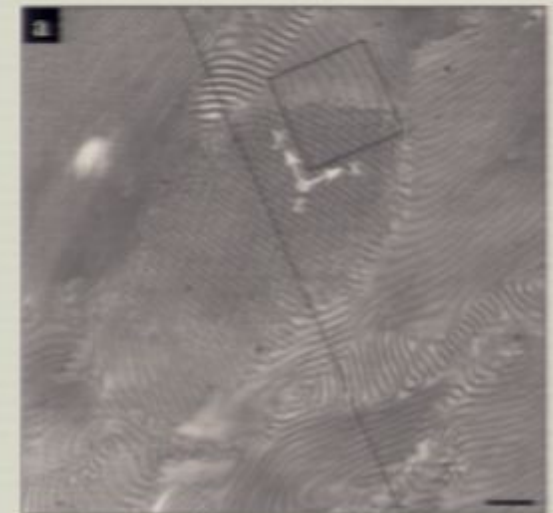
Sota et al. (2003)



Chastek and Lodge (2006)

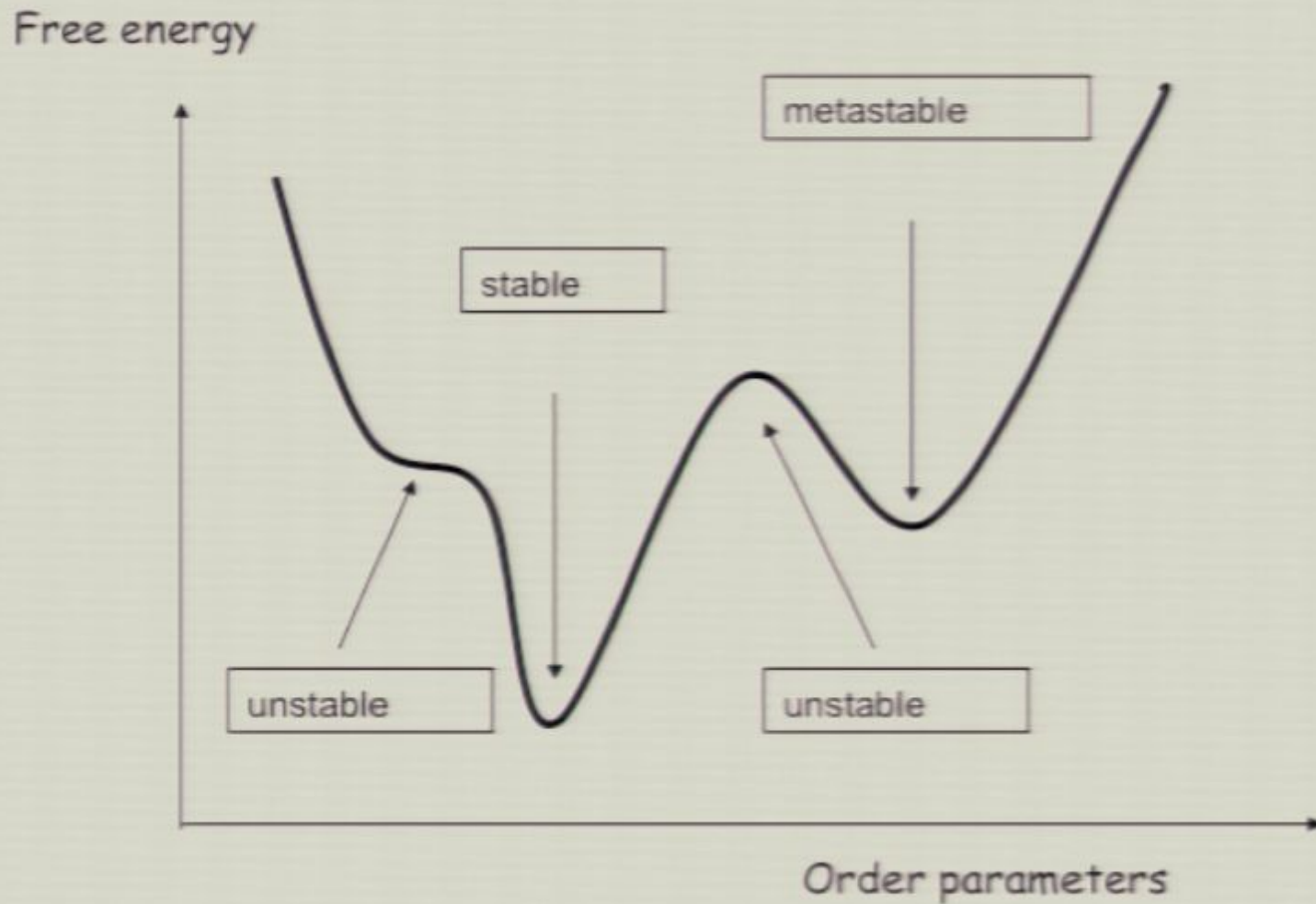


Koizumi et al. (1994)

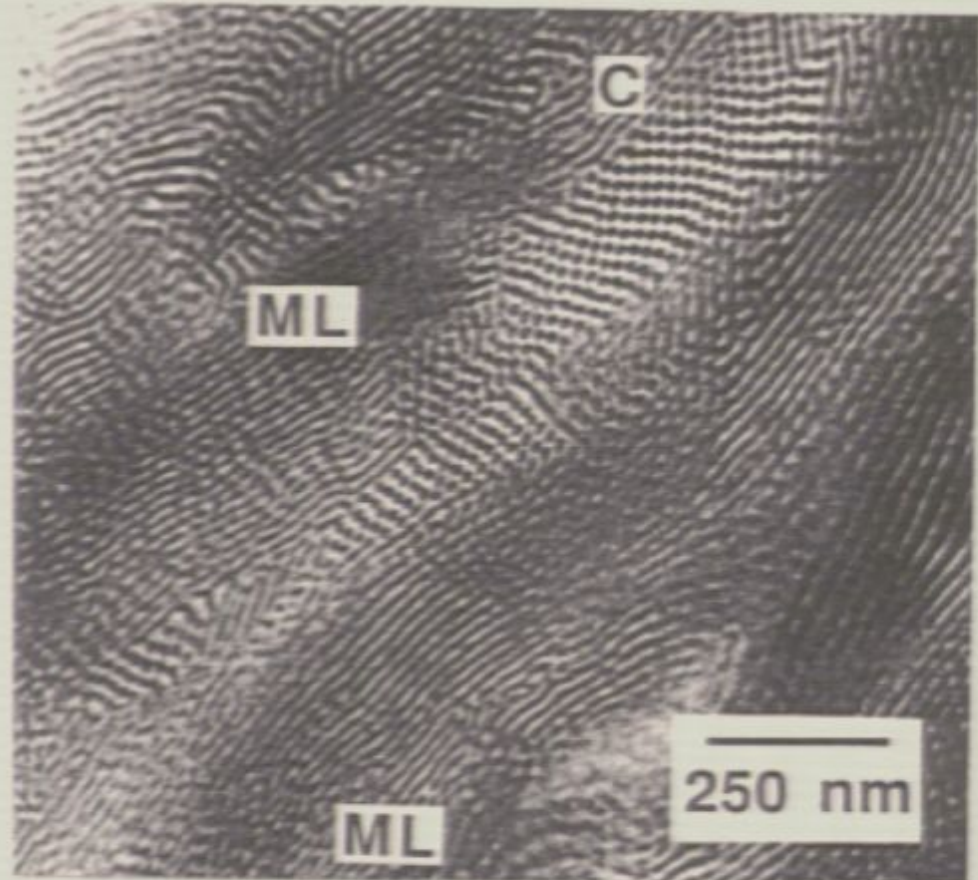
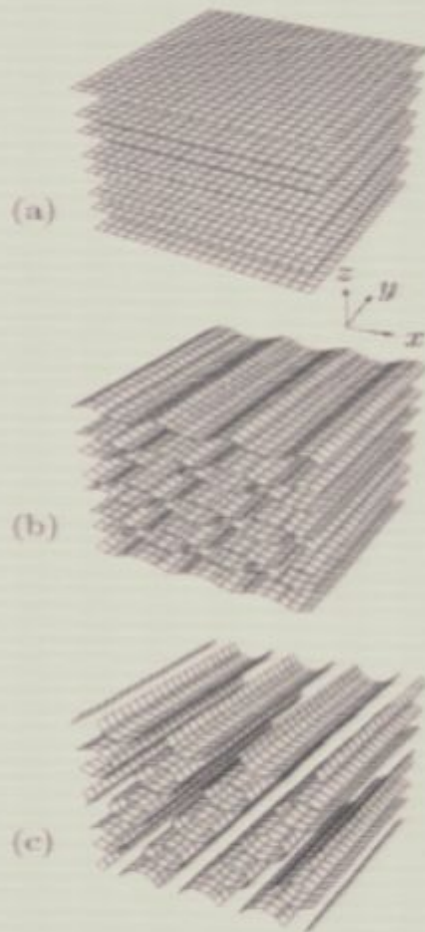


Mareau et al. (2007)

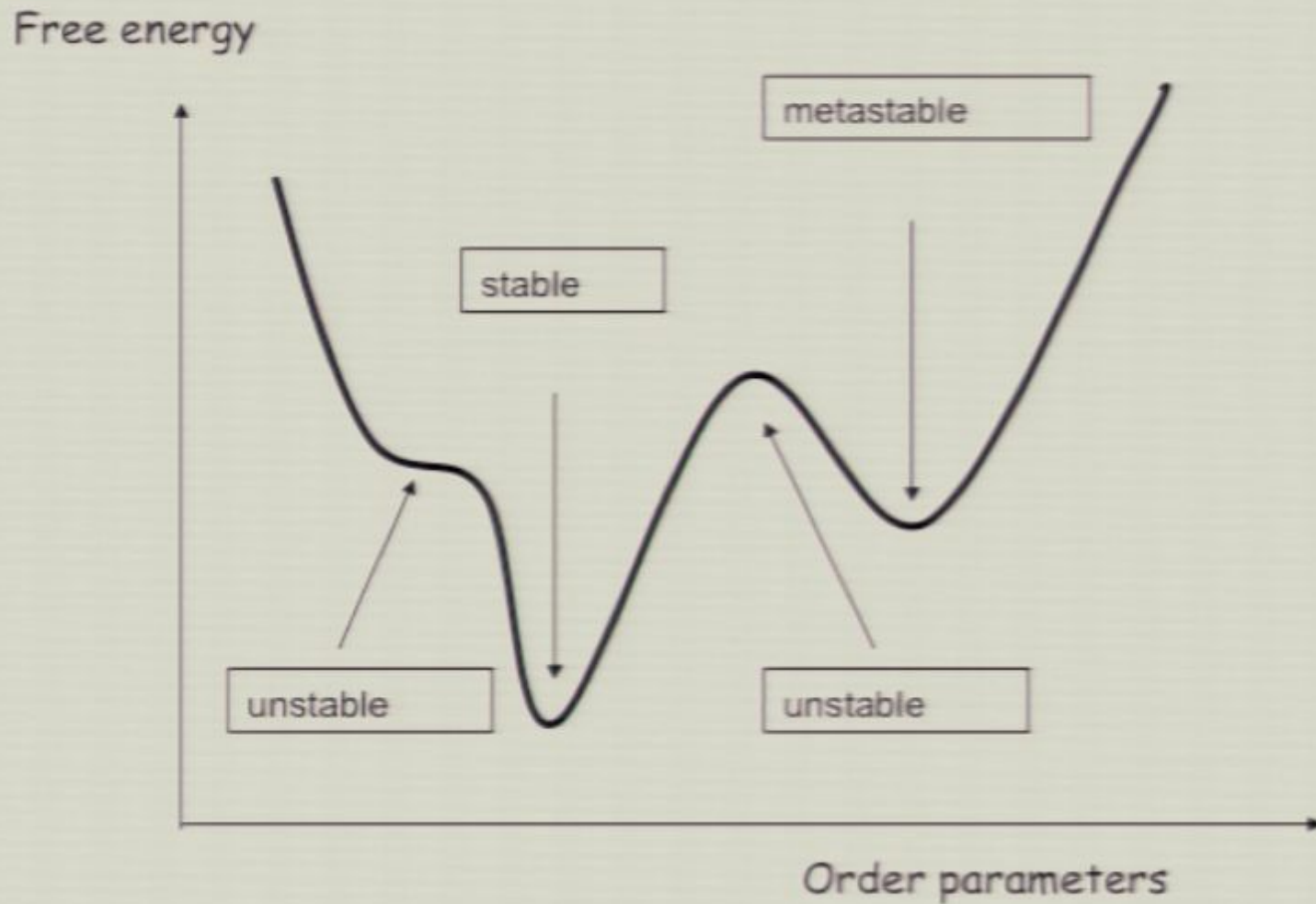
Free-energy landscape



Spinodal decomposition of unstable phase



Free-energy landscape



Classical Theory of Droplet Nucleation (CNT)

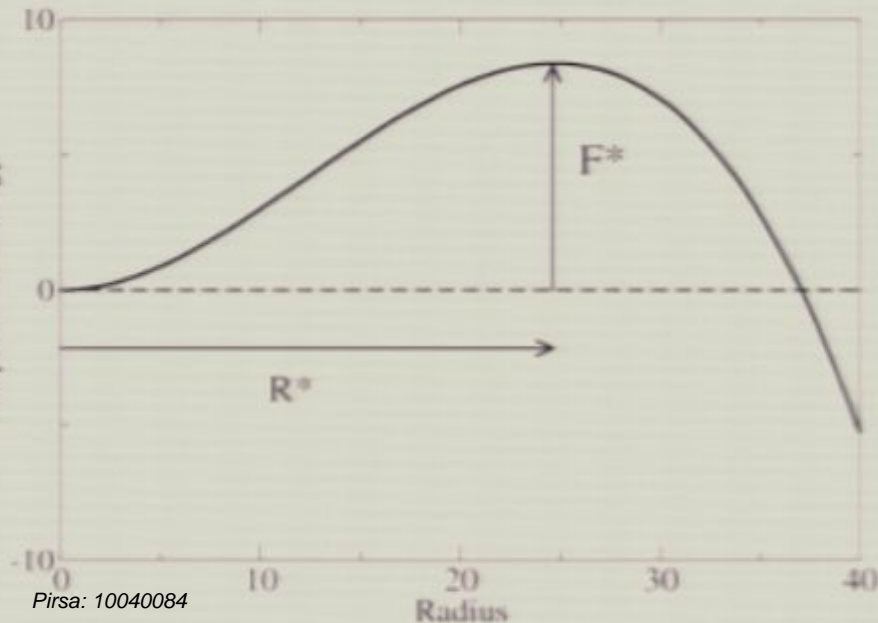
e.g. nucleation of liquid droplet from supersaturated vapour



$$F(r) = \frac{4}{3}\pi R^3 \times \Delta f + 4\pi R^2 \times \sigma$$

$$\Delta f = f_{\text{liquid}} - f_{\text{vapour}} < 0$$

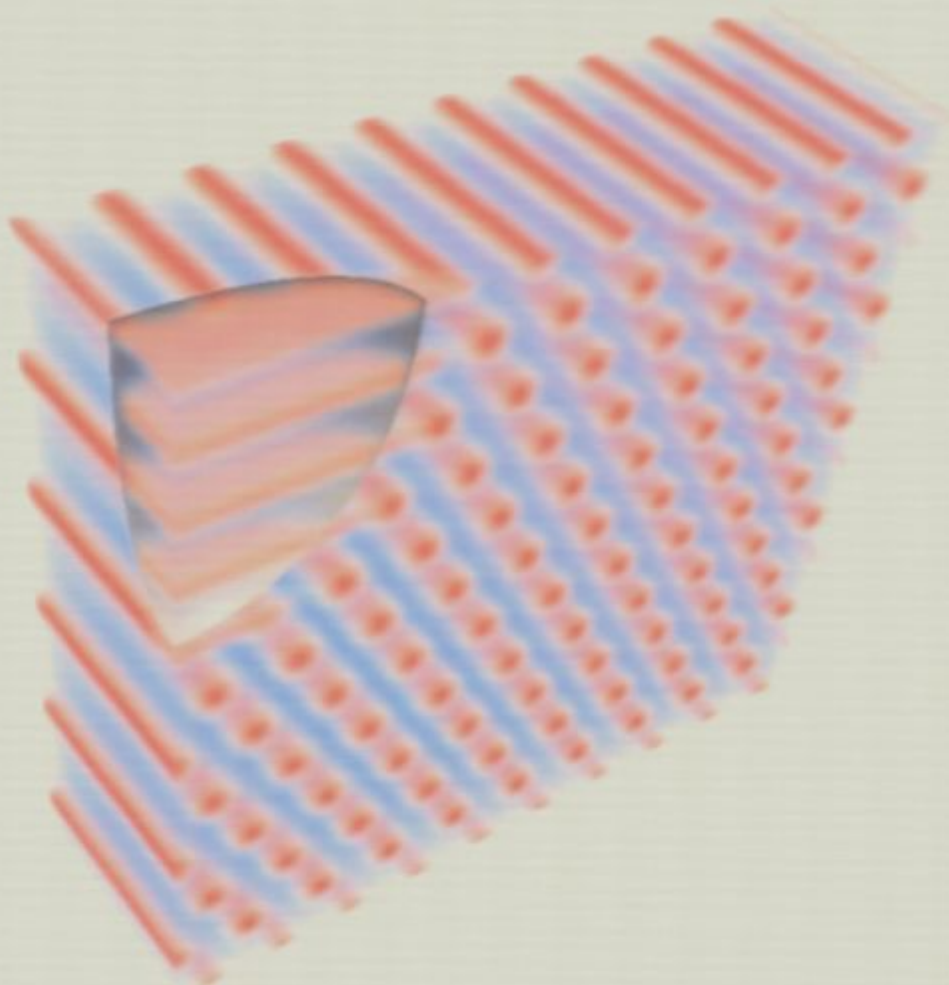
σ = liquid/vapour interfacial tension



$$R^* \sim \sigma / |\Delta f|$$

$$\text{Nucleation rate} \sim \exp(-F^*/kT)$$

Nucleation and growth from a metastable phase



- Inhomogeneous problem (droplet)

- Multiple length-scales

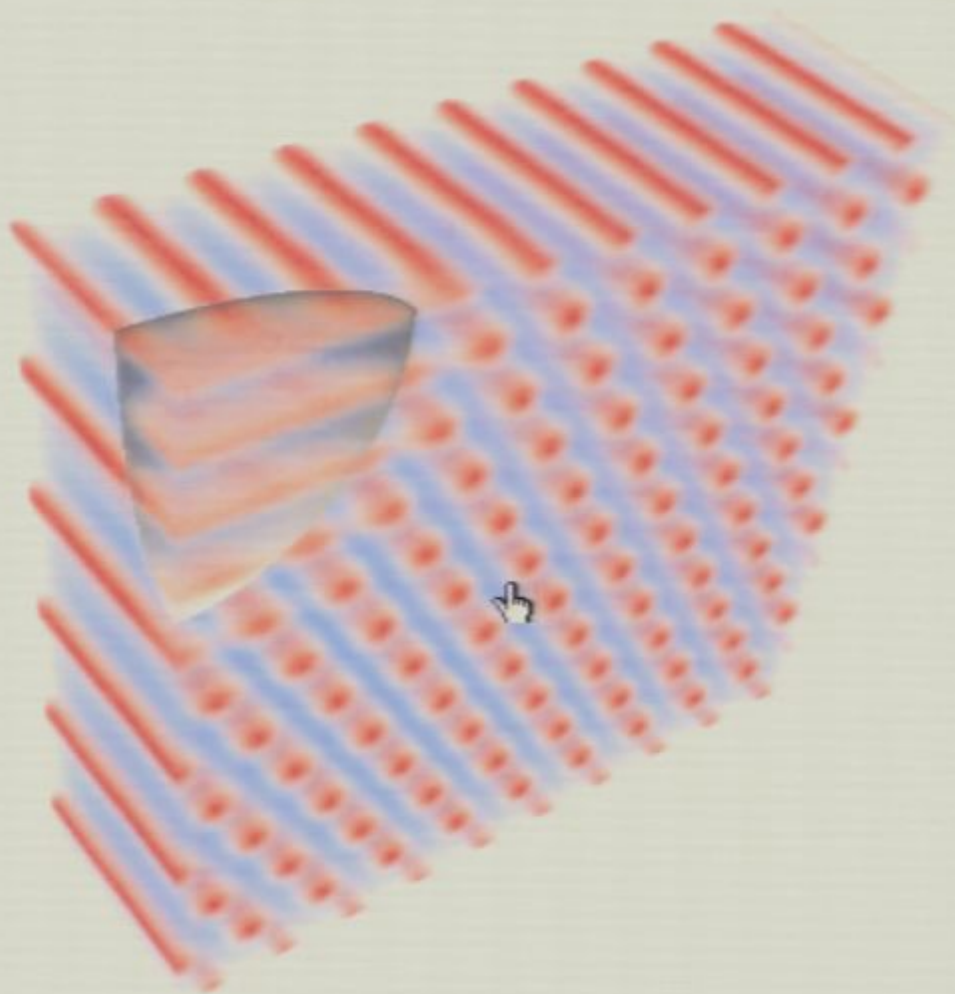
- Microstructure:

structure of interface?

anisotropic surface
free-energy

non-spherical nucleus

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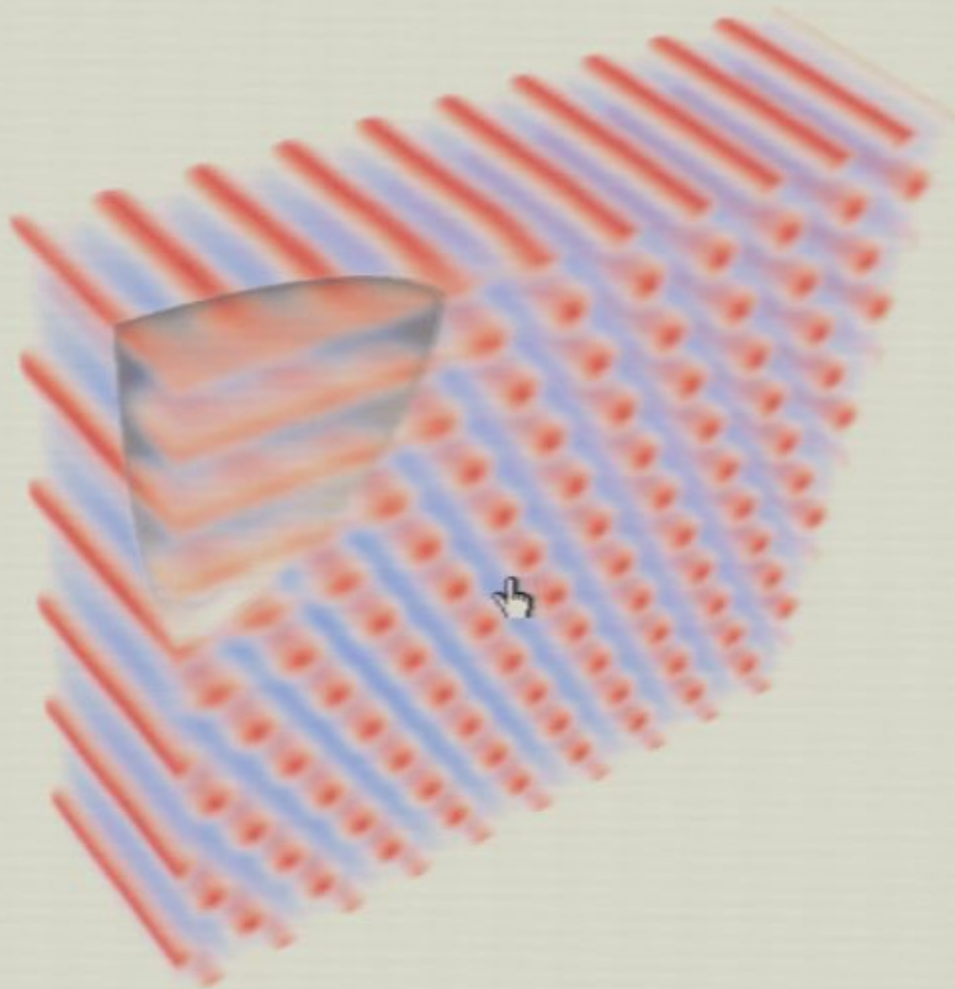
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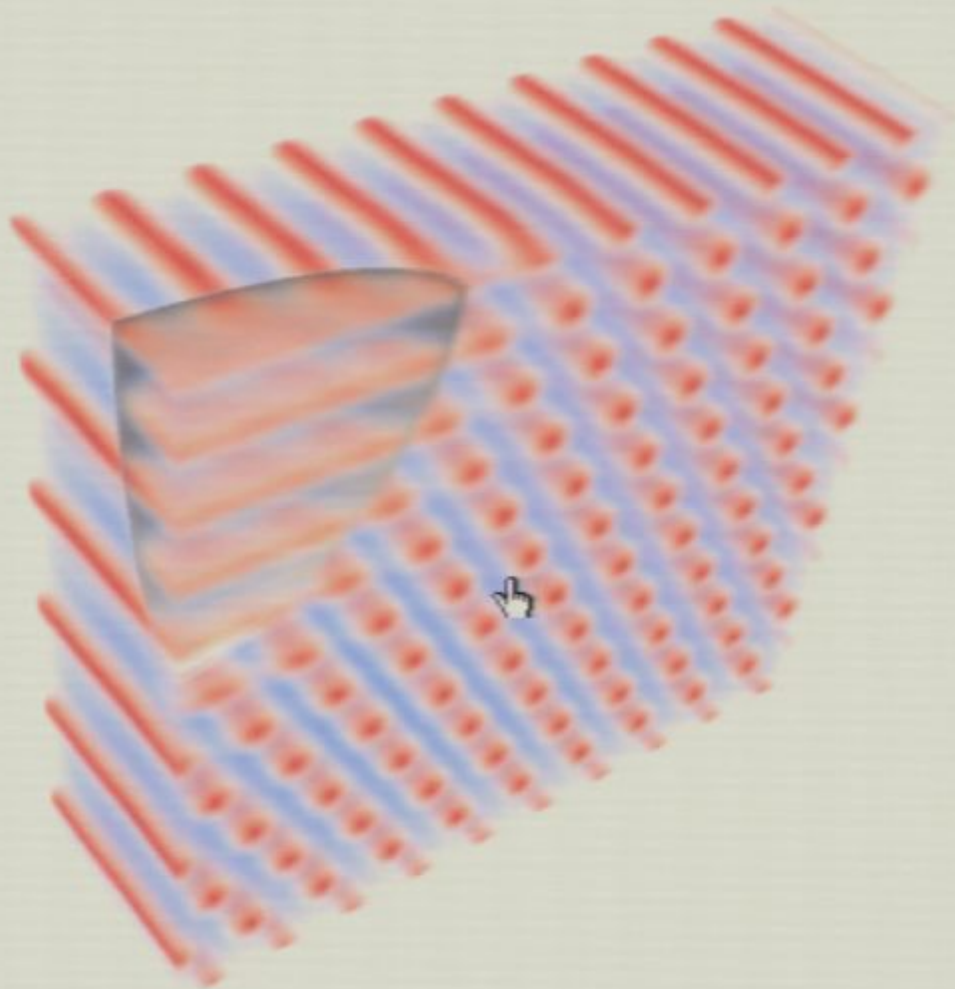
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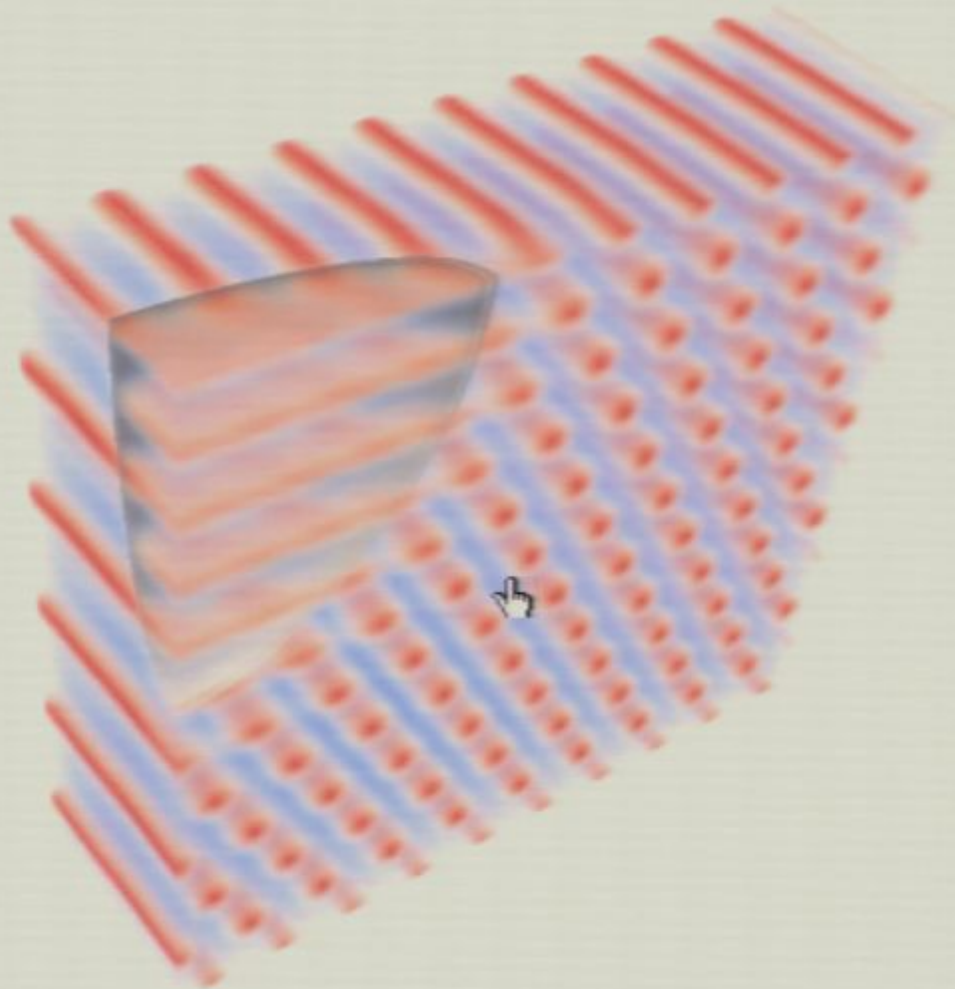
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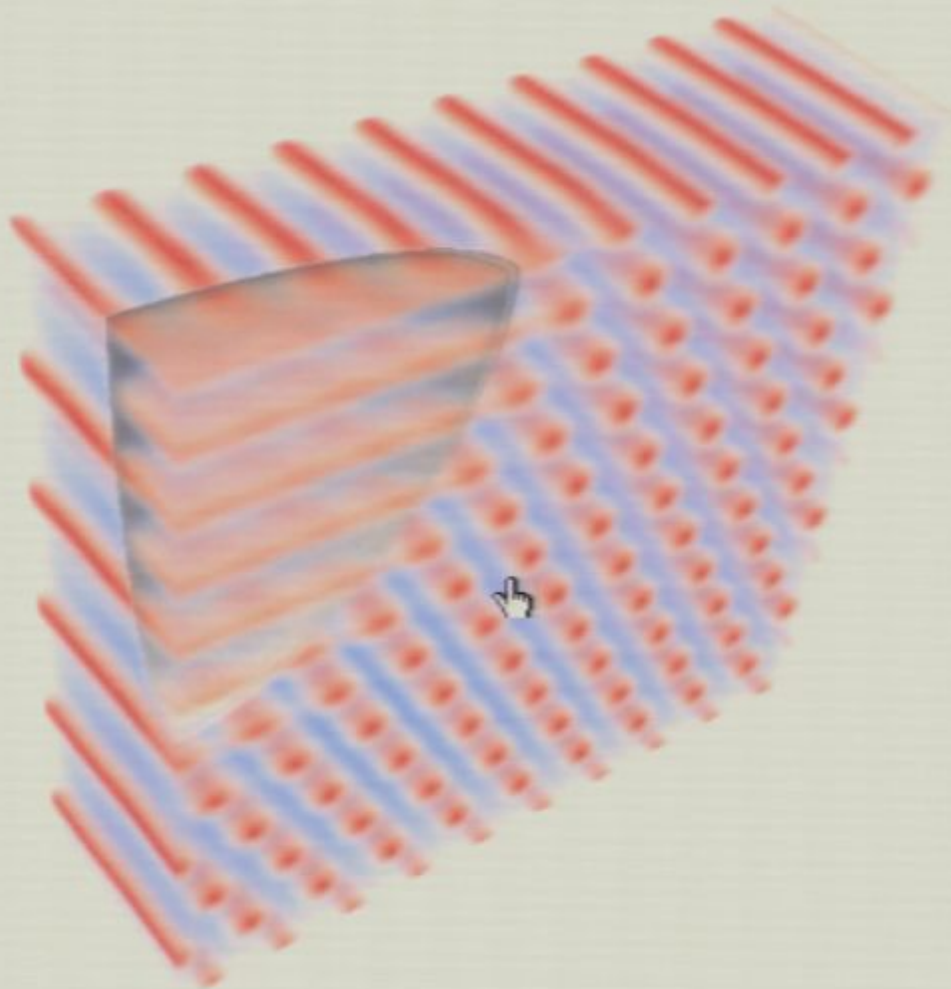
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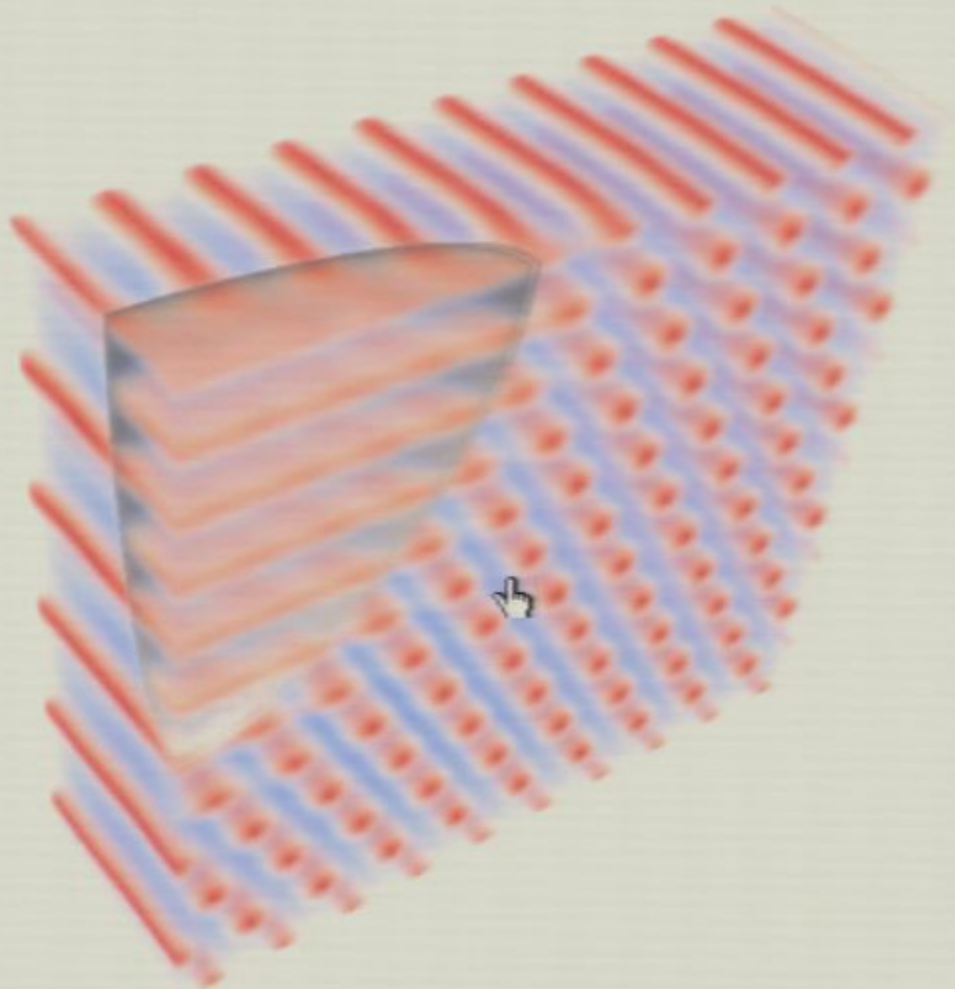
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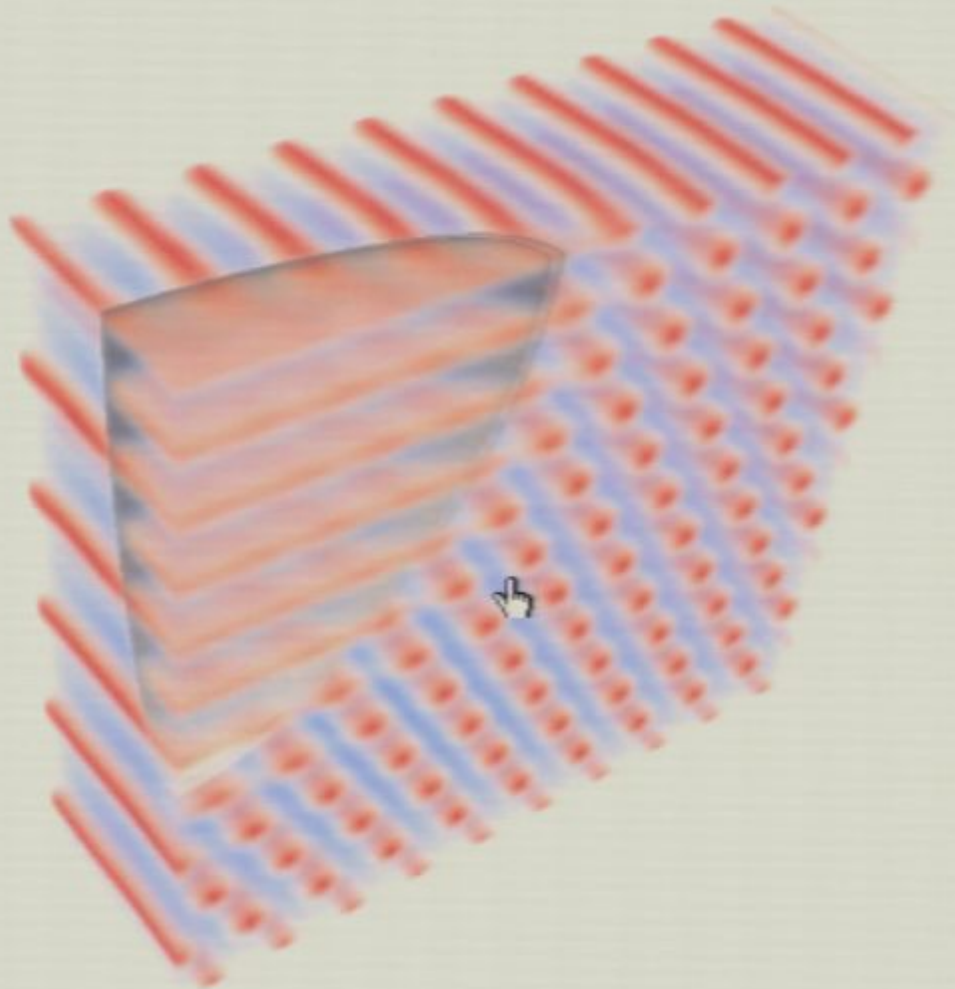
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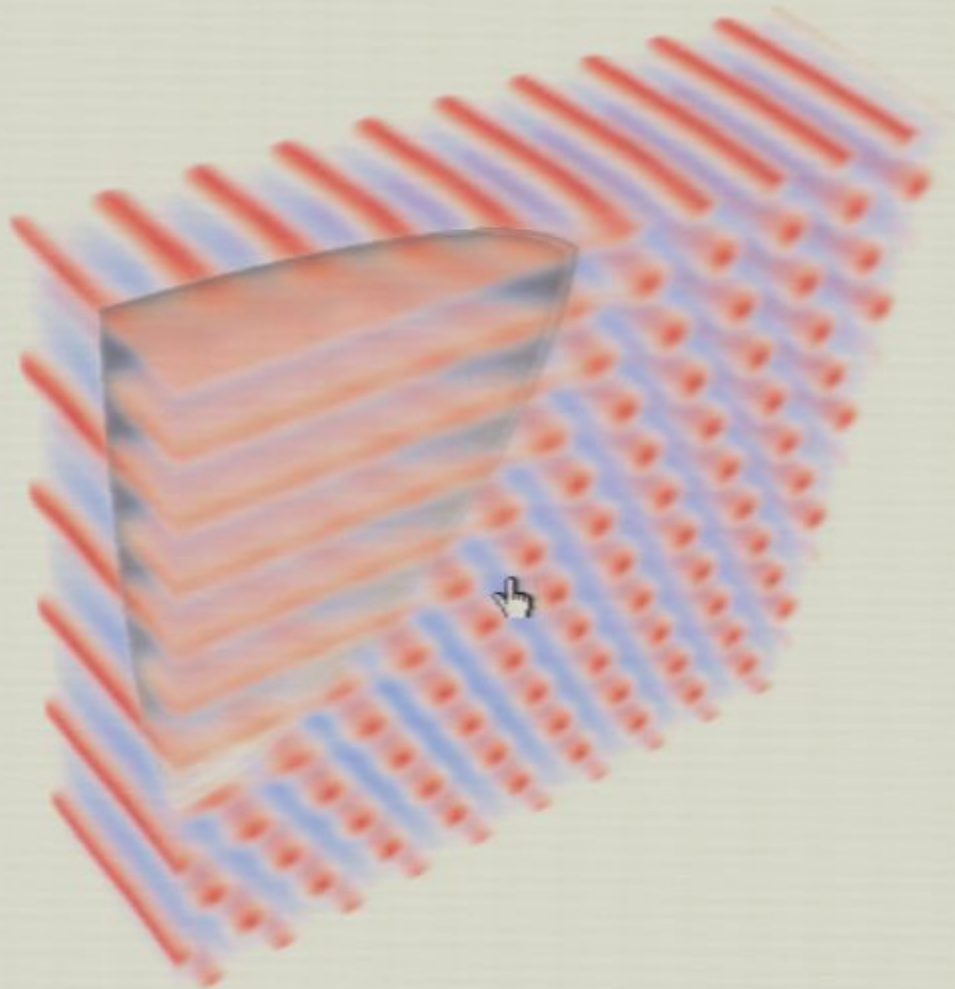
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Theoretical approach to incorporate anisotropy in surface free-energy

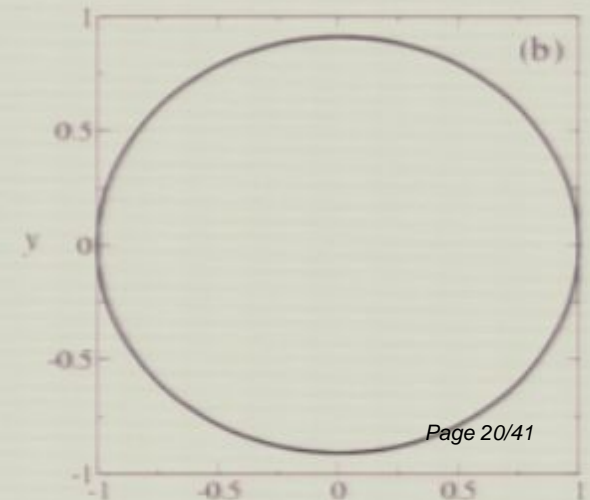
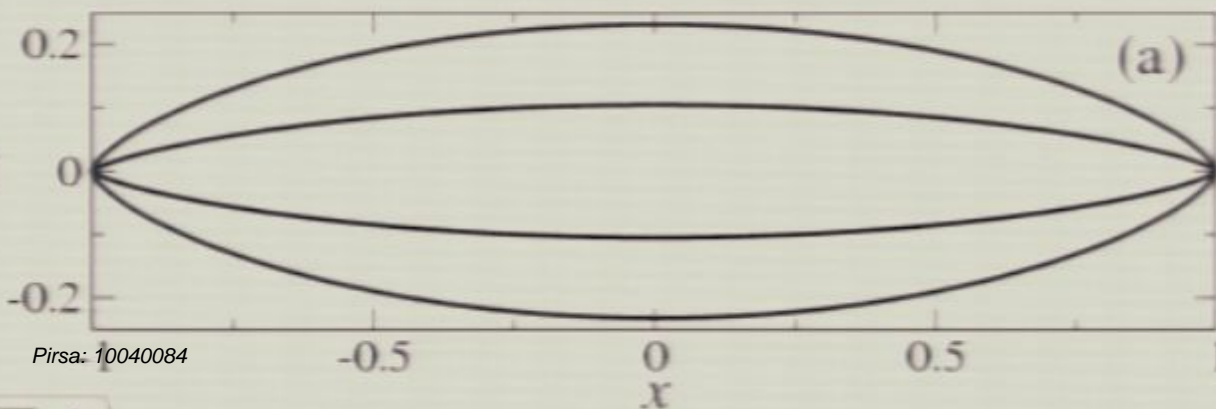
Wickham, Shi and Wang (2003)

Landau-Brazovskii free-energy in terms of order parameter $\phi(\vec{r}) = \rho_A(\vec{r}) - \langle \rho_A \rangle$:

$$F = \int d^3 r \left\{ \frac{\xi^2}{2} [(\nabla^2 + q_0^2)\phi]^2 + \frac{\tau}{2}\phi^2 - \frac{\gamma}{3}\phi^3 + \frac{1}{4}\phi^4 \right\}$$

Exhibits all equilibrium phases for diblock copolymers

Basis of classical nucleation theory for anisotropic droplets
(Wickham, Shi, Wang, 2003)



Time-dependent Landau-Brazovskii model

Spencer, Wickham (2010)

$$F = \int d^3r \left\{ \frac{\xi^2}{2} [(\nabla^2 + q_0^2)\phi]^2 + \frac{\tau}{2}\phi^2 - \frac{\gamma}{3}\phi^3 + \frac{1}{4}\phi^4 \right\}$$

Add dynamics:

$$\frac{\partial\phi}{\partial t} = \Gamma \nabla^2 \frac{\delta F}{\delta\phi} \quad (\text{conserved order-parameter, model B dynamics})$$

$$\frac{\partial\phi}{\partial t} = \Gamma \nabla^2 [\xi^2 (\nabla^2 + q_0^2)^2 \phi + \tau\phi - \gamma\phi^2 + \phi^3]$$

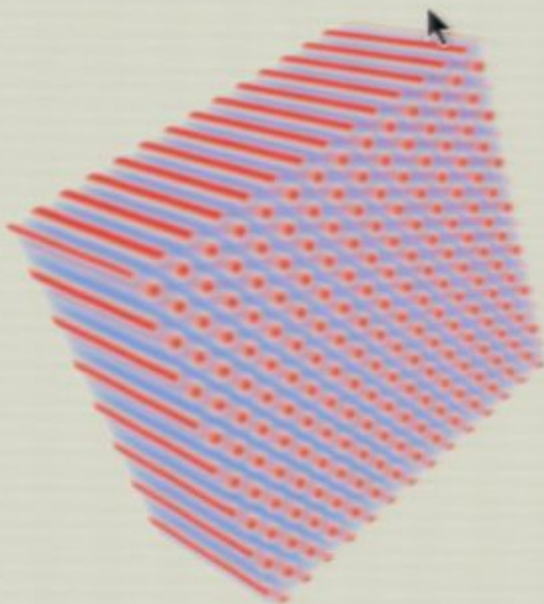
Note that there is no noise term in the model

Numerical Implementation

- Finite difference scheme on cubic lattice
- Largest run: 512x580x200 on 64 core (using MPI) for 1 week (1 million time steps)

Initial state:

Single nucleus of stable lamellar phase in background of metastable cylindrical phase

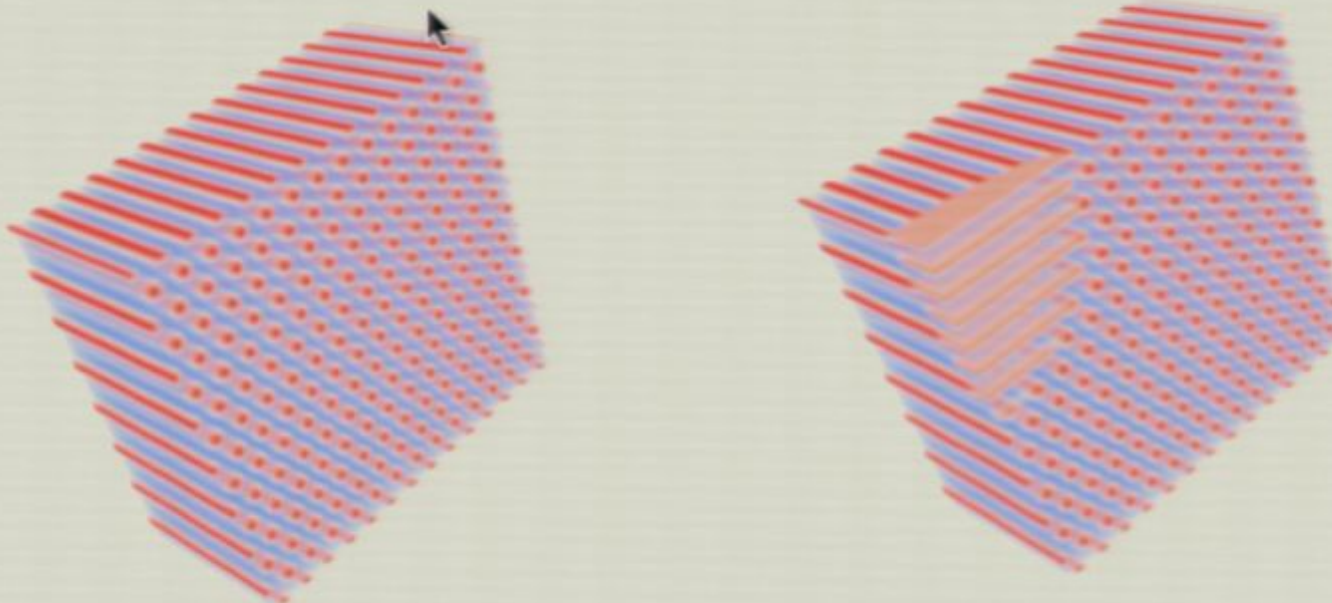


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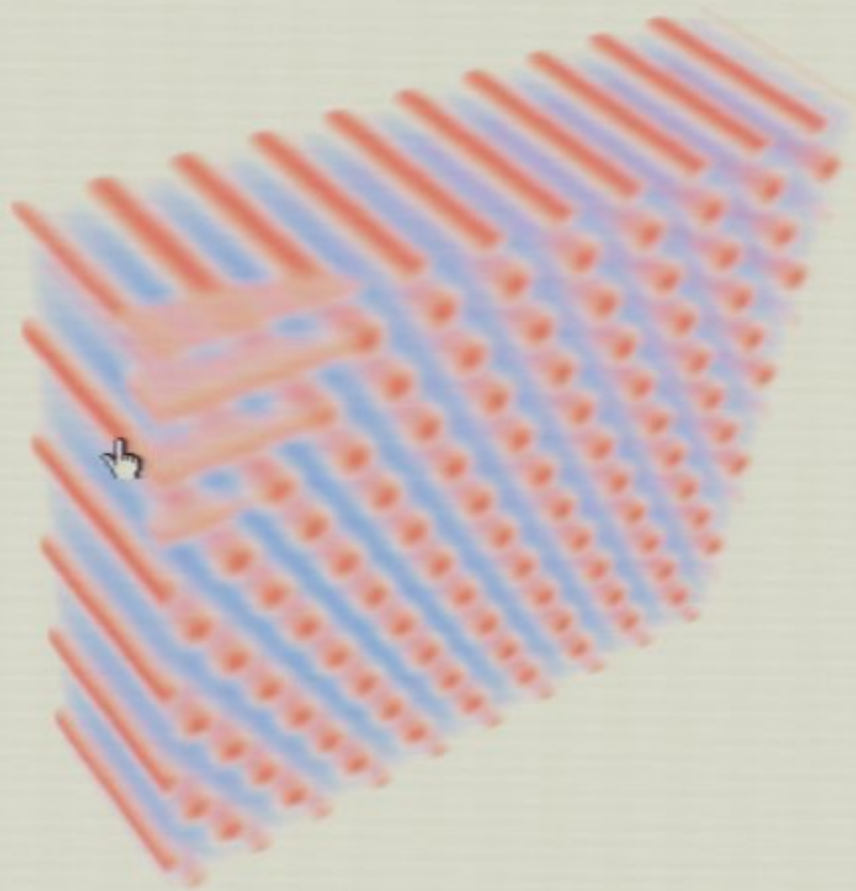
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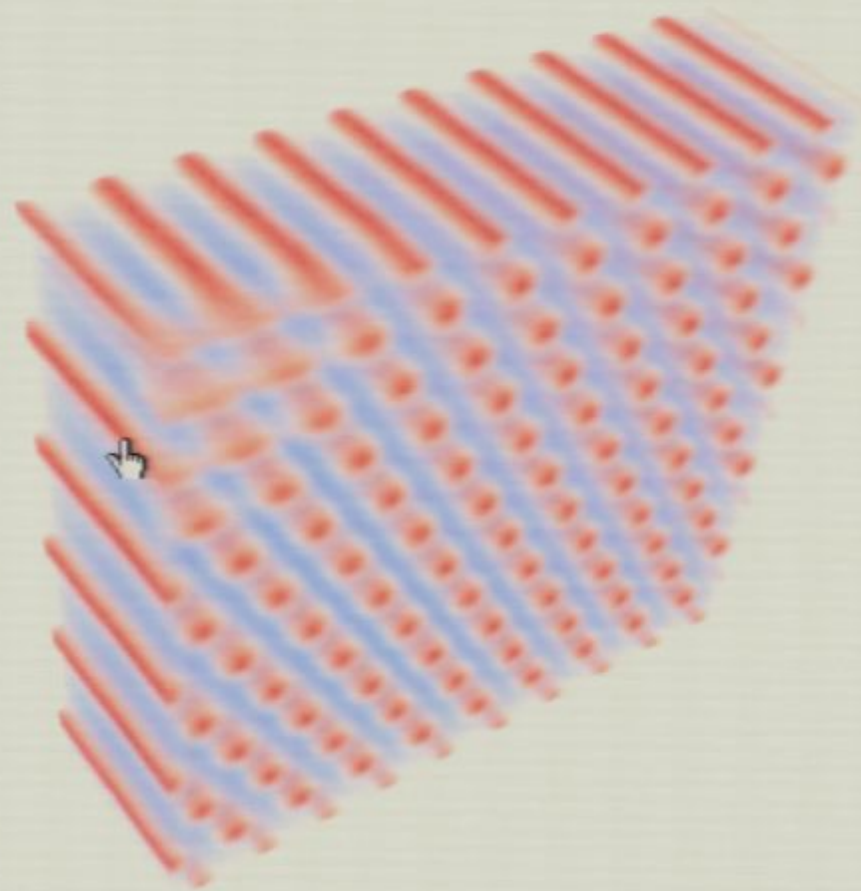
Identifying the critical nucleus size

Fix undercooling: $\tau_{coexist} - \tau = 0.0026$



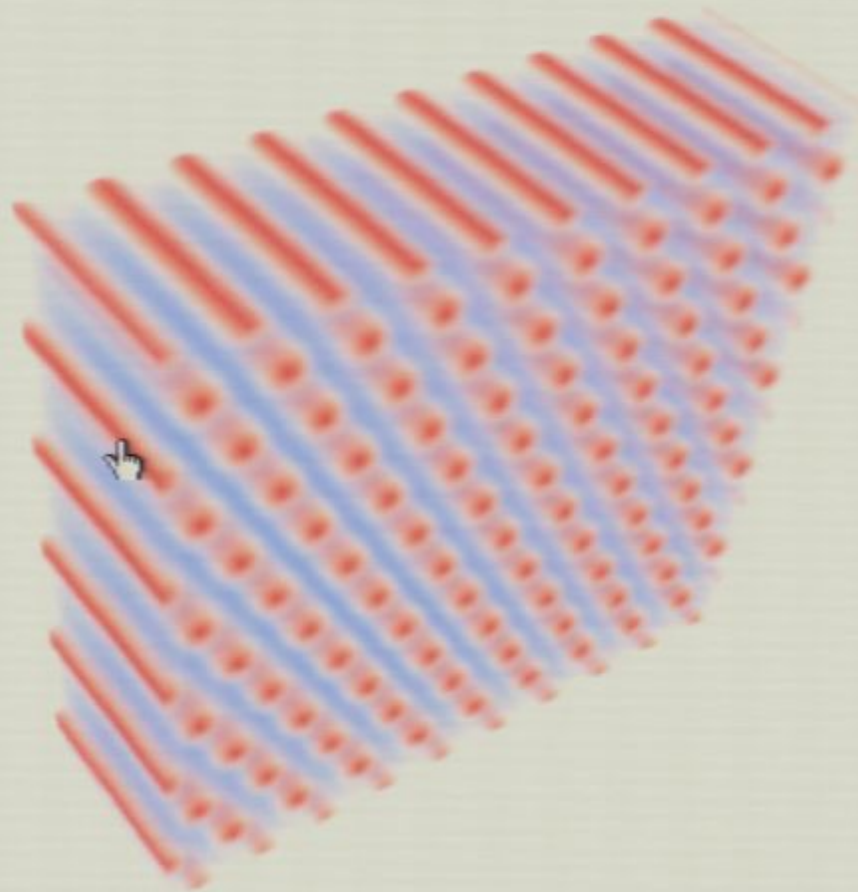
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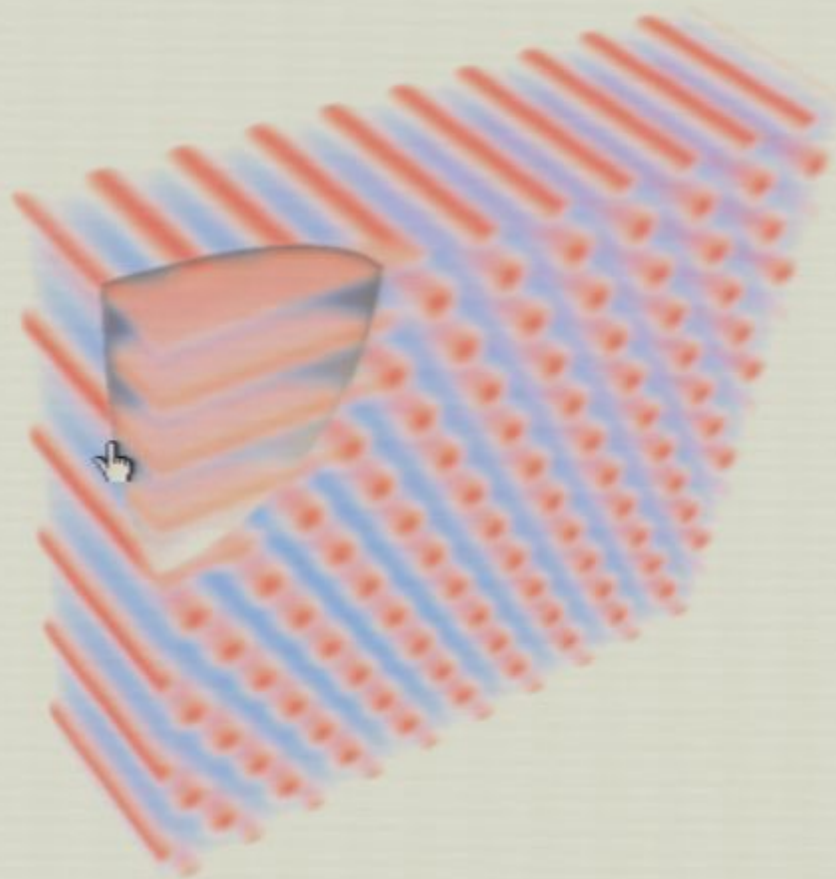
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Identifying the critical nucleus size

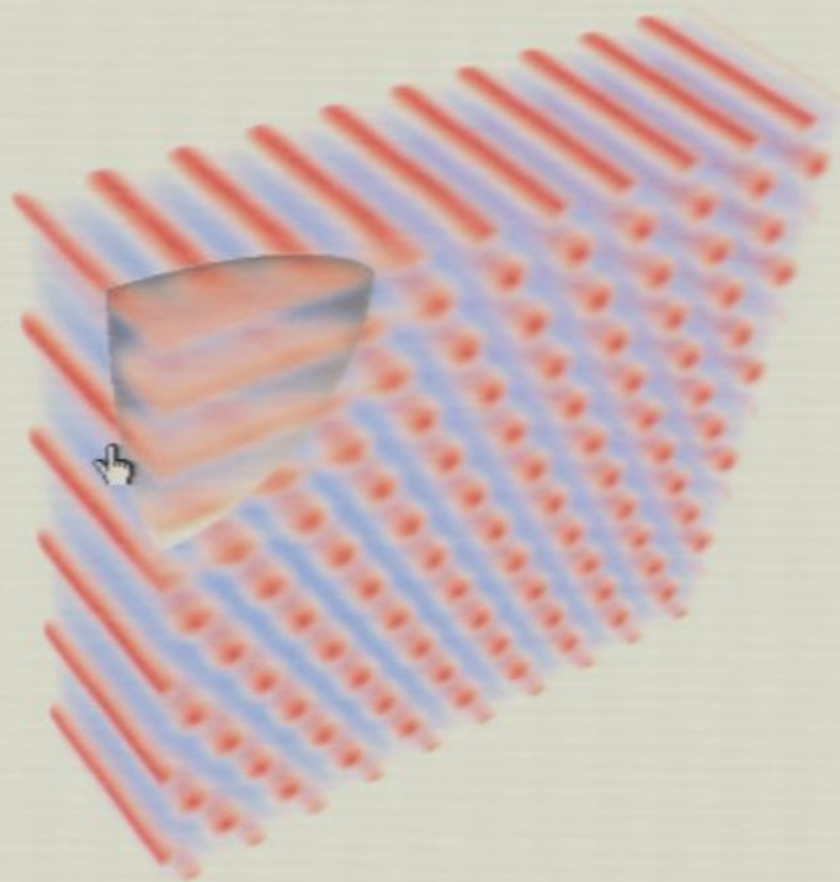
Nuclei with an initial size above a certain threshold will grow



Find $\frac{L_{crit}}{L_0} = 6.0 \pm 0.5$ in this case

Identifying the critical nucleus size

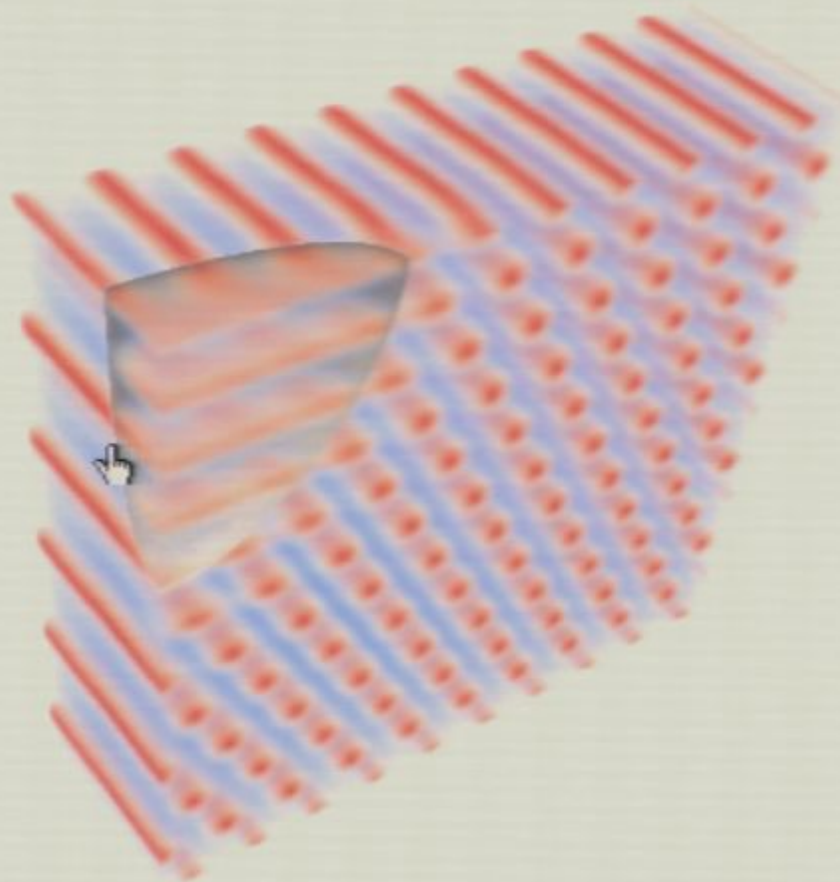
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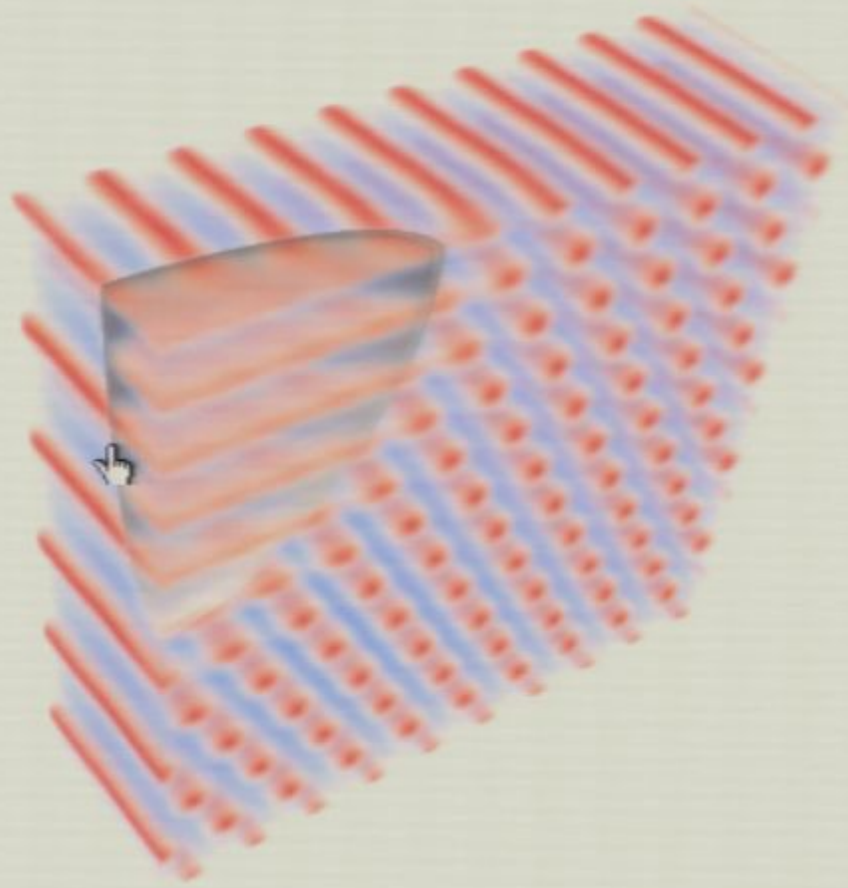
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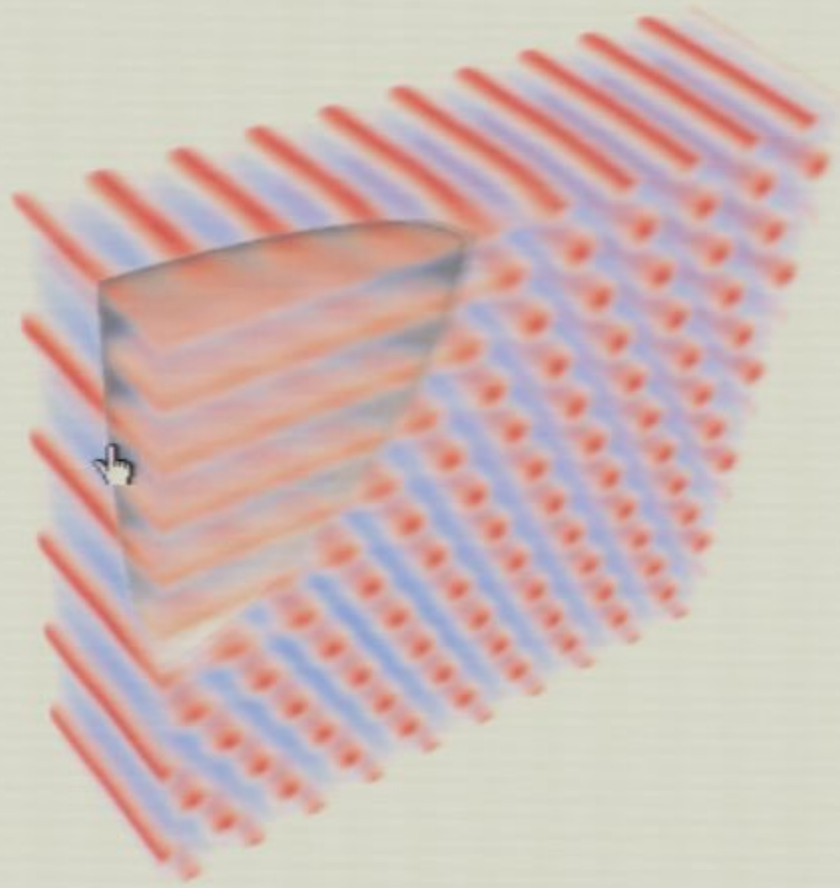
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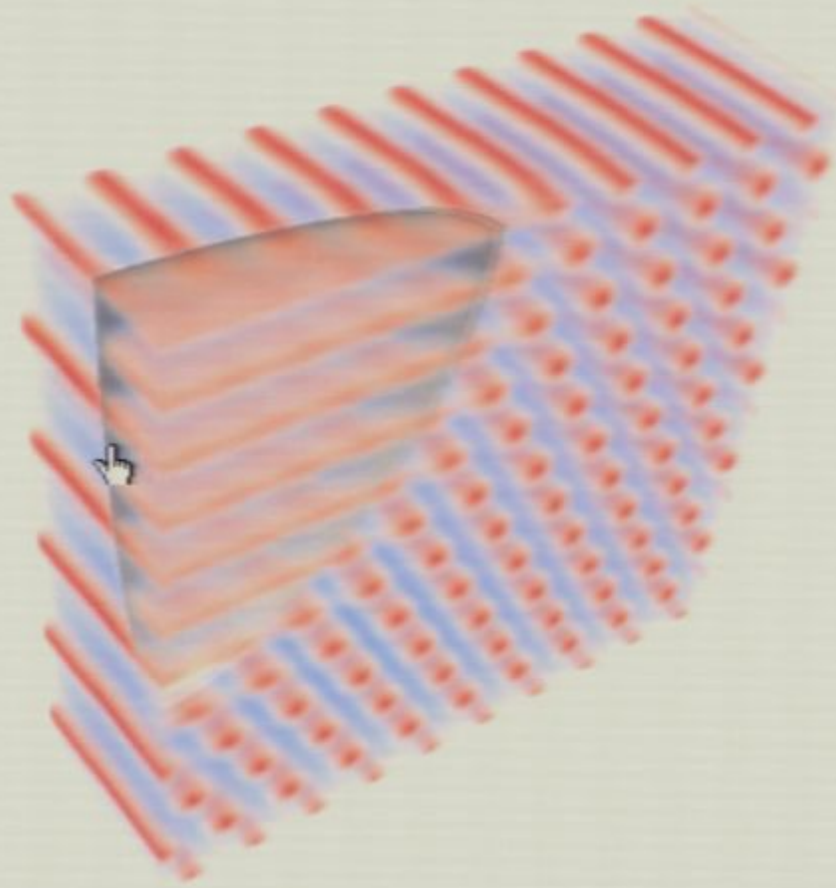
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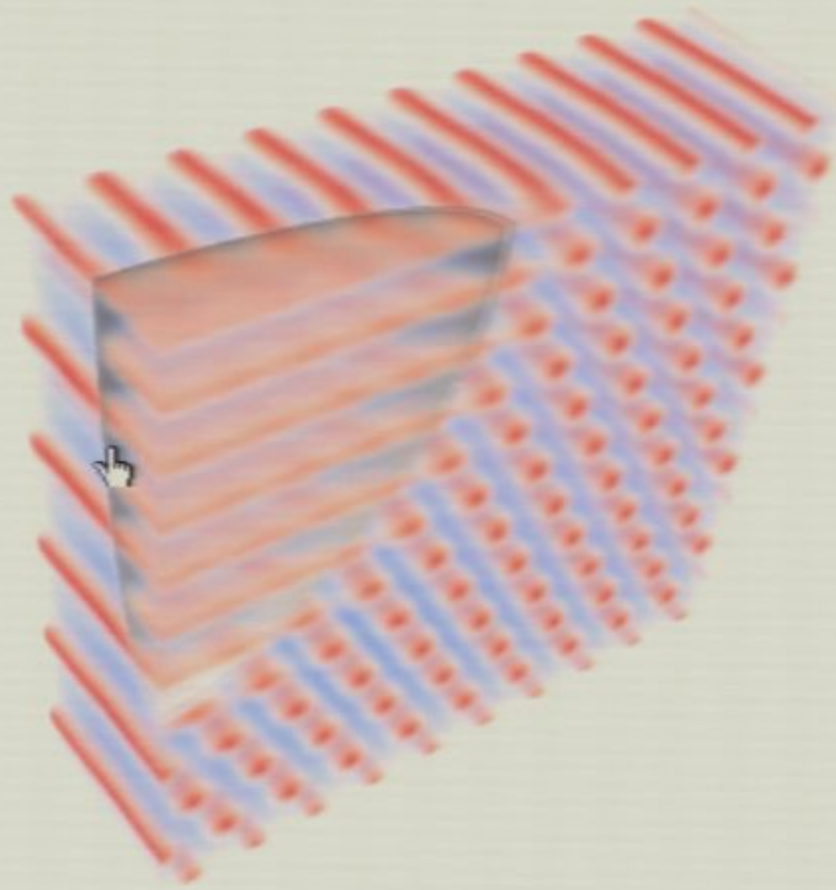
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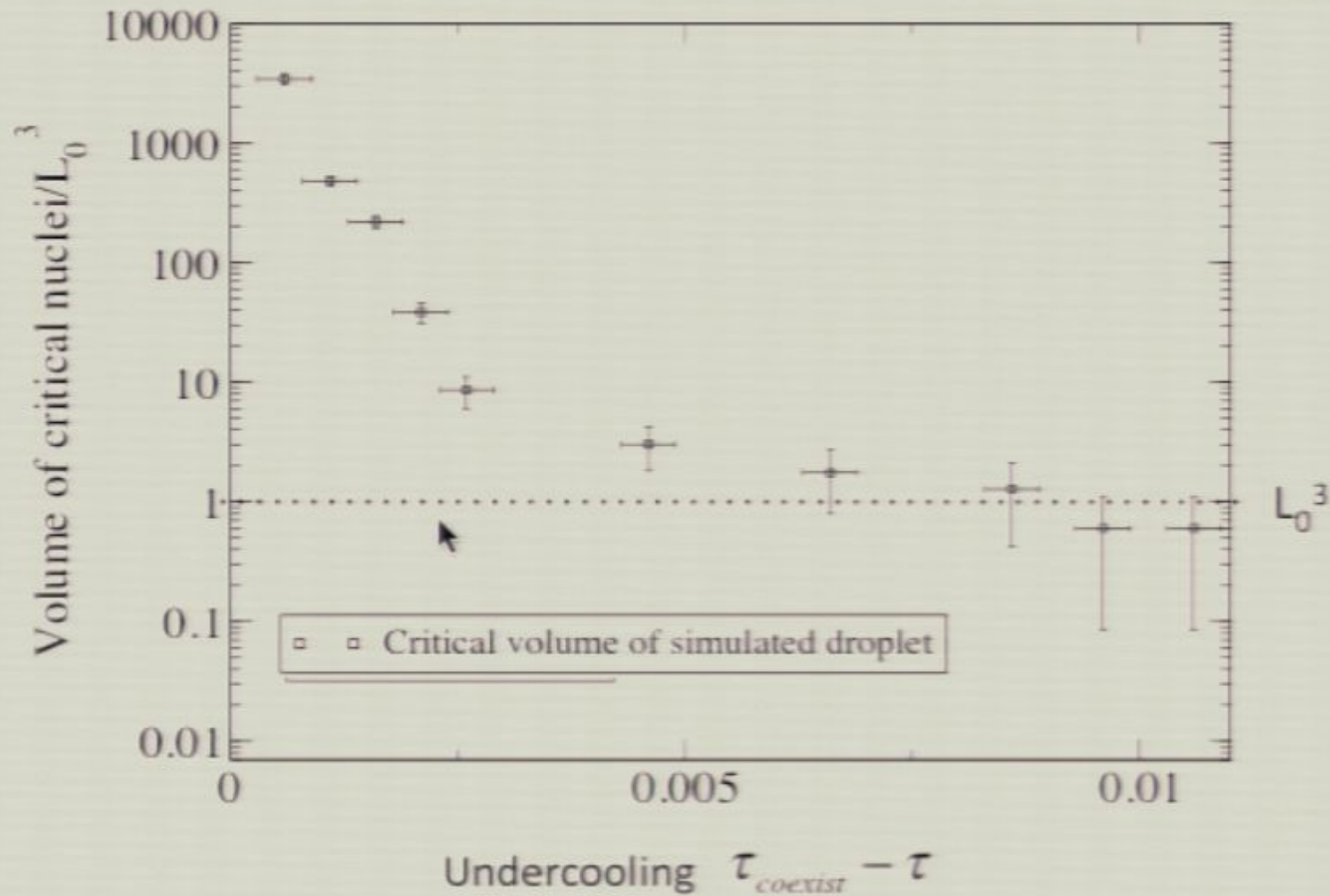
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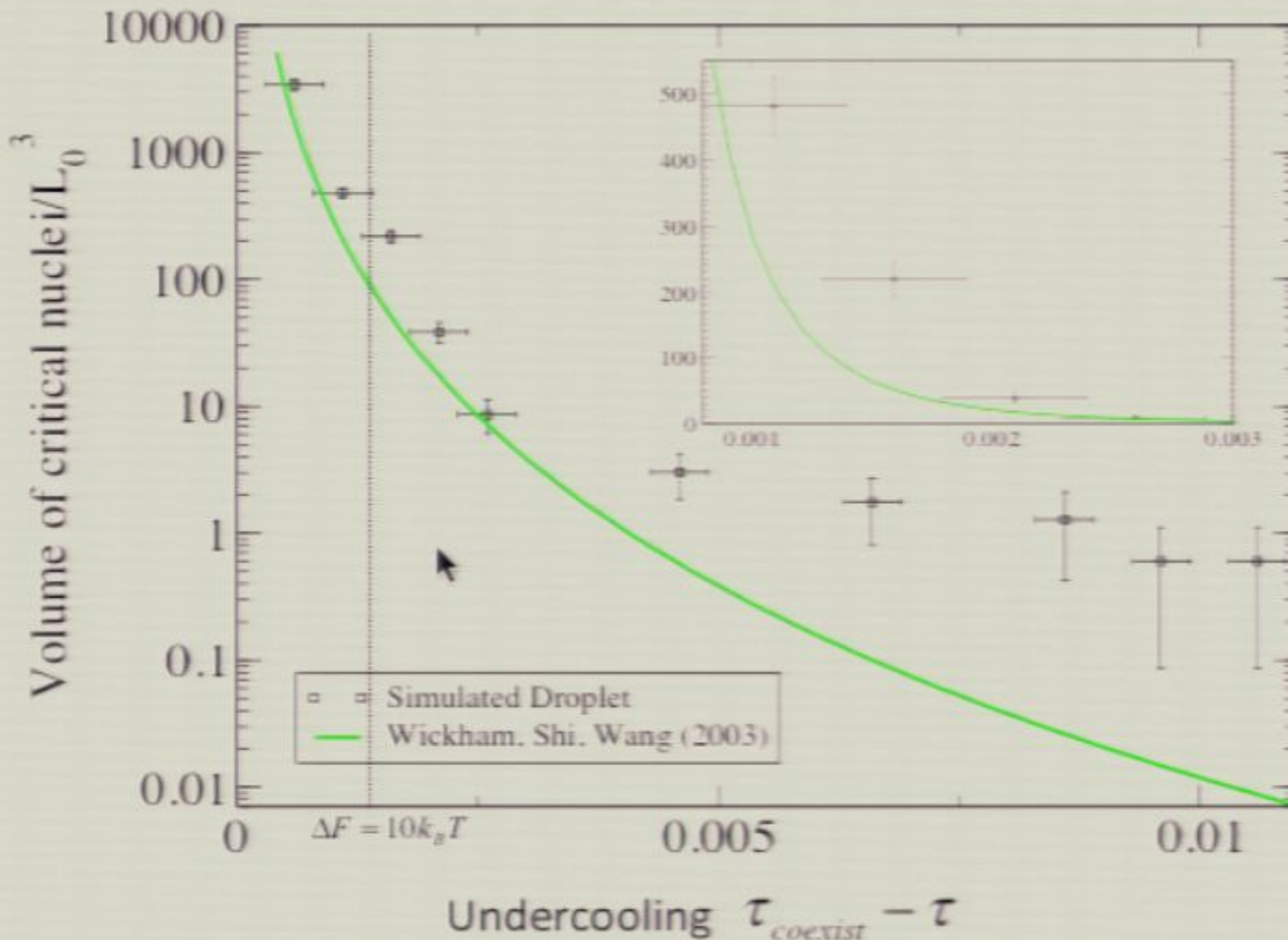
Find $\frac{L_{crit}}{L_0} = 6.0 \pm 0.5$ in this case

Critical nucleus volume vs. undercooling



Size of critical nucleus increases as coexistence is approached

Test of the approximate (2003) theory



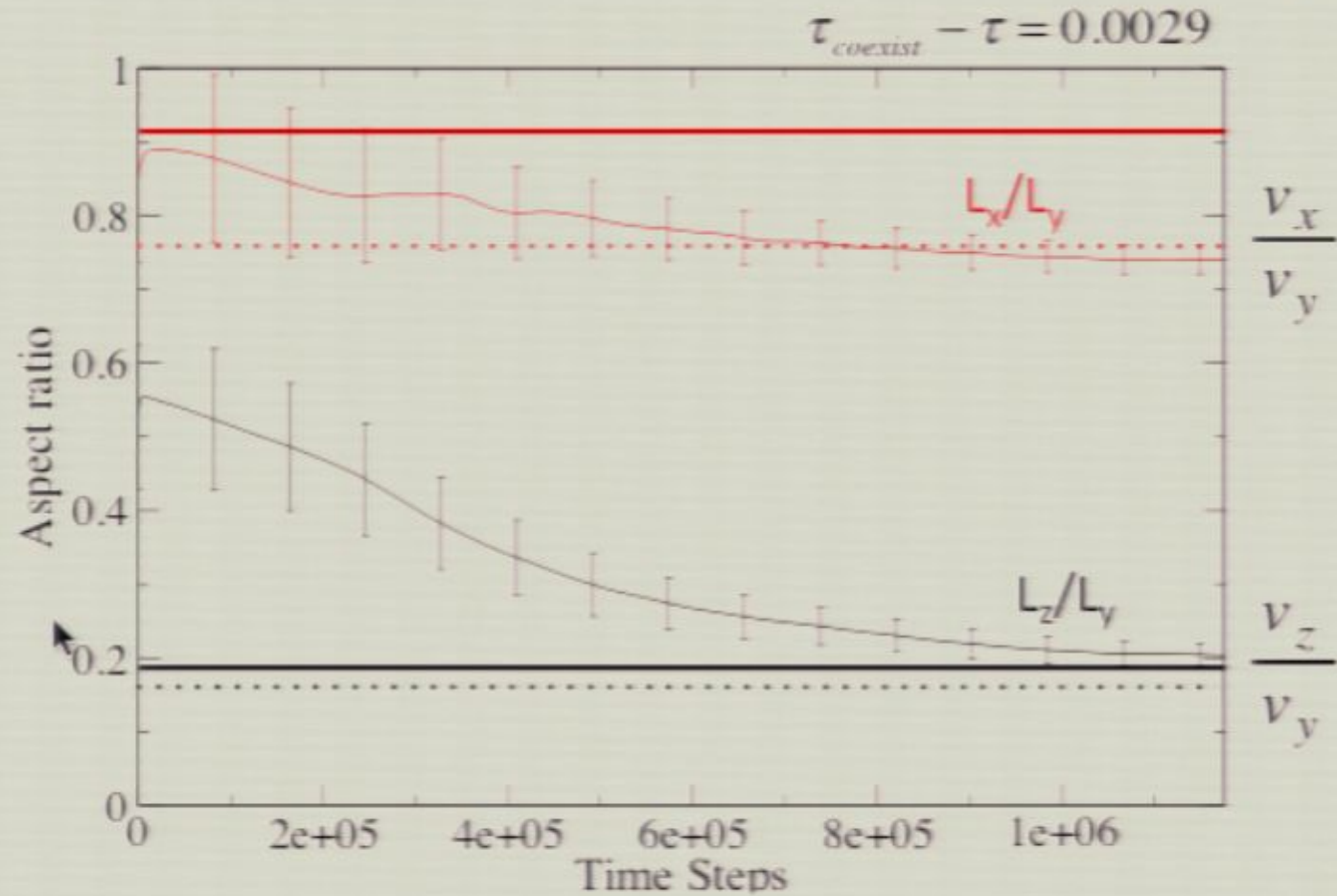
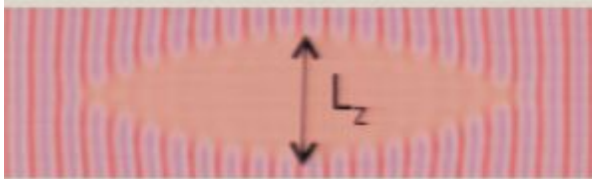
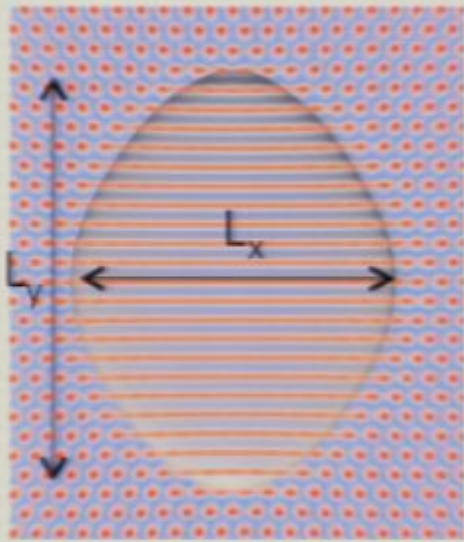
Theory:

$$V_{crit} \sim \left(\frac{\sigma}{|\Delta f|} \right)^3$$

No adjustable parameters

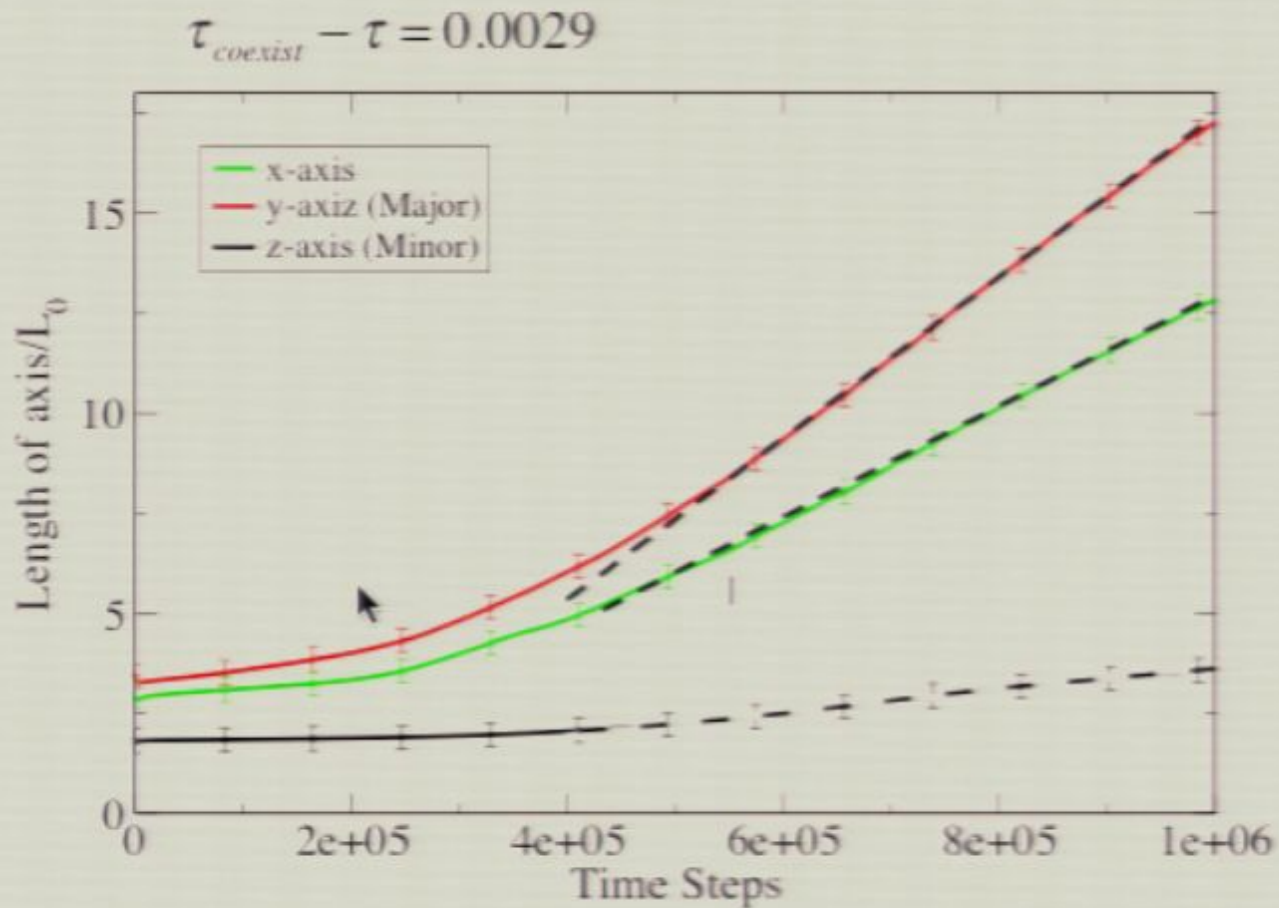
Approximate theory reasonably describes simulation near coexistence

Time-evolution of nucleus shape



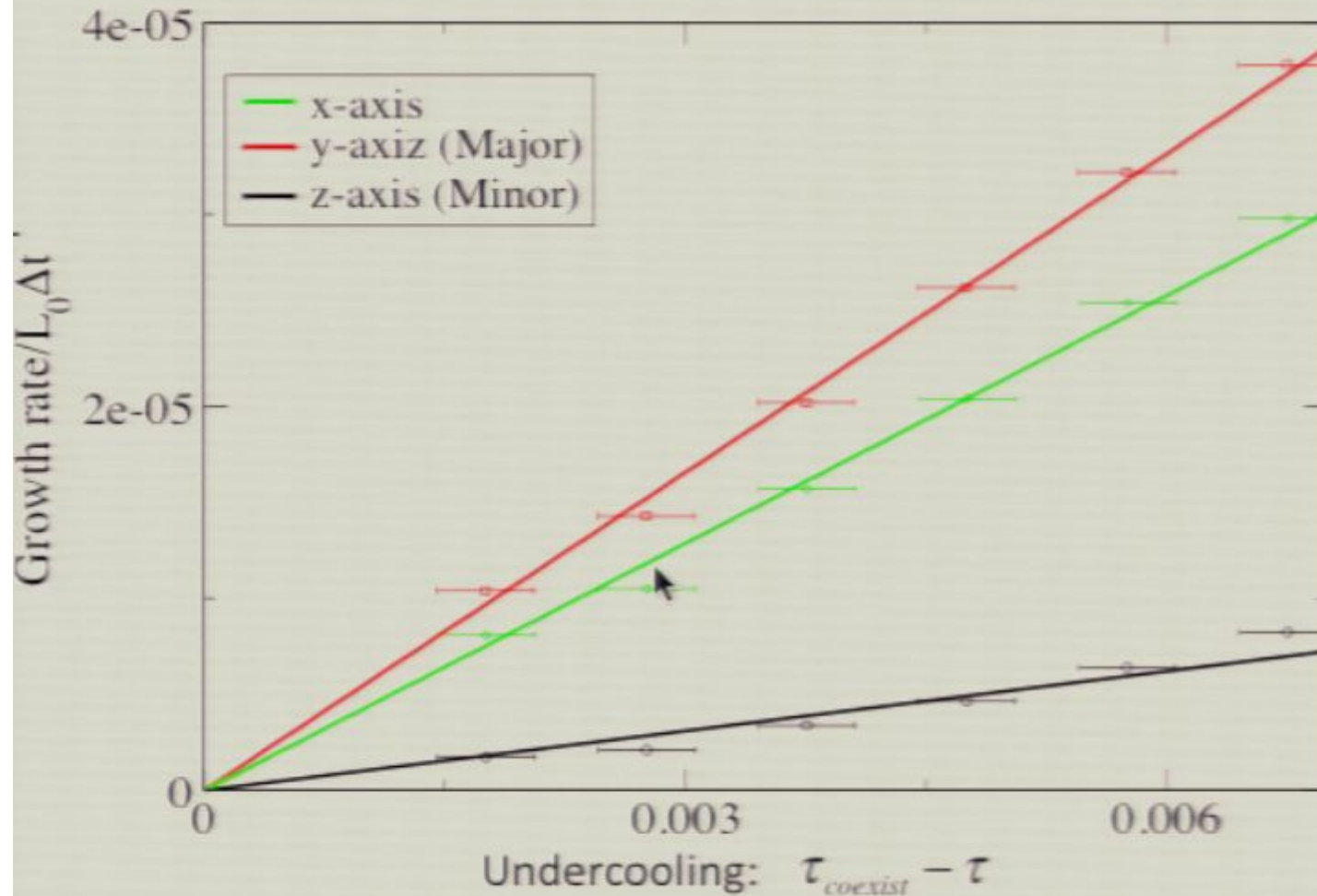
- Observe lens-like nucleus predicted by approximate theory
- Shape becomes constant

Kinetics: Growth of the droplet



Growth enters linear regime where interface velocity is constant

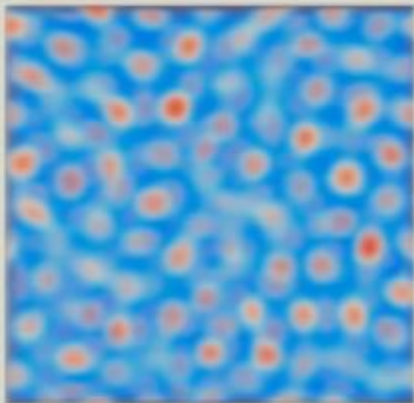
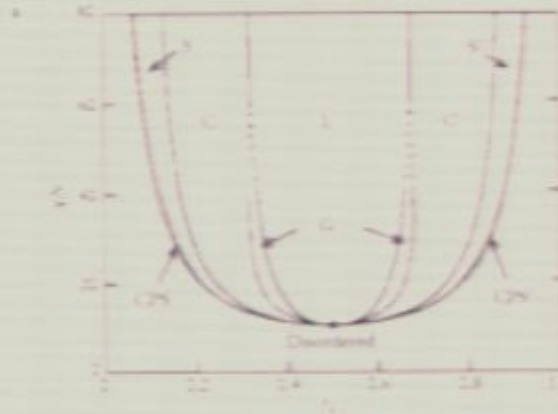
Interface velocity as a function of undercooling



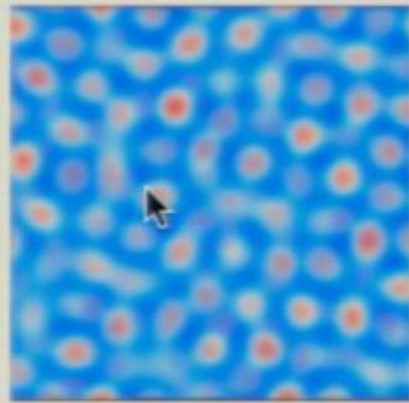
$$v \sim \frac{R_{\tau}}{T} (\tau_{coexist} - \tau)$$

Interface velocity is proportional to undercooling

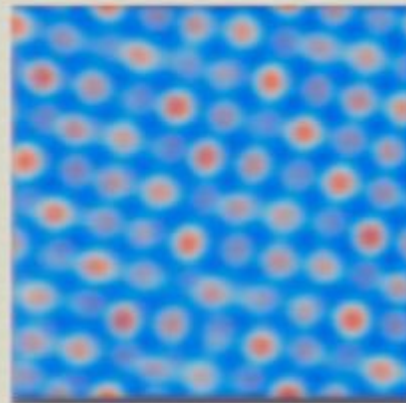
Other dynamical scenarios: Crystalization



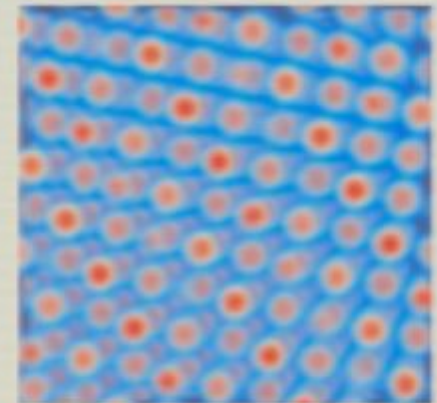
1.1 million
time steps



1.3 million
time steps



1.6 million
time steps



2 million
time steps

- Nature of disordered phase

- Kinetic pathways to forming the equilibrium BCC sphere phase

Summary

- Simulating nucleation and growth when microstructure is present is challenging
- We were able to identify a critical nucleus size
- The nucleus size increases as coexistence is approached, consistent with theory of Wickham, Shi and Wang (2003)
- The nucleus is lens-shaped, consistent with theory, but detailed comparison suggests kinetic effects are important
- Observed behaviour of interfacial velocity is consistent with theory of Goveas and Milner (1997)

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for Innovation

Fondation canadienne
pour l'innovation

