

Title: Explorations in Numerical Relativity (PHYS 642) - Lecture 5

Date: Mar 19, 2010 11:20 AM

URL: <http://pirsa.org/10030104>

Abstract:

Stability Analysis

- Will use the property captured by (95) as working definition of stability.
- In particular, if you believe (95) is true for the wave equation, then you believe the wave equation is stable.
- Fundamentally, if FDA approximation *converges*, then expect the same behaviour for the difference solution:

$$\|u_j^n\| \sim \|u_j^0\|. \quad (97)$$

- FD solution constructed by *iterating in time*, generating

$$u_j^0, u_j^1, u_j^2, u_j^3, u_j^4, \dots$$

in succession, using the FD equation

$$u_j^{n+1} = 2u_j^n - u_j^{n-1} + \lambda^2 \left(u_{j+1}^n - 2u_j^n + u_{j-1}^n \right).$$

Stability Analysis

- One of the most frustrating/fascinating features of FD solutions of time dependent problems: discrete solutions often “blow up”—e.g. floating-point overflows are generated at some point in the evolution
- ‘Blow-ups’ can sometimes be caused by legitimate (!) “bugs”—i.e. an incorrect implementation—at other times it is simply the *nature of the FD scheme* which causes problems.
- Are thus lead to consider the *stability* of solutions of difference equations
- Again consider the 1-d wave equation, $u_{tt} = u_{xx}$
- Note that it is a *linear, non-dispersive* wave equation
- Thus the “size” of the solution does *not* change with time:

$$\|u(x, t)\| \sim \|u(x, 0)\|, \quad (95)$$

where $\|\cdot\|$ is an suitable norm, such as the L_2 norm:

$$\|u(x, t)\| \equiv \left(\int_0^1 u(x, t)^2 dx \right)^{1/2}. \quad (96)$$

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Stability Analysis

- Not guaranteed that (97) holds for all values of $\lambda \equiv \Delta t / \Delta x$.

- For certain λ , have

$$\|u_j^n\| \gg \|u_j^0\|,$$

and for those λ , $\|u^n\|$ diverges from u , even (especially!) as $h \rightarrow 0$ —that is, the difference scheme is *unstable*.

- For many wave problems (including all linear problems), given that a FD scheme is *consistent* (i.e. so that $\hat{\tau} \rightarrow 0$ as $h \rightarrow 0$), *stability is the necessary and sufficient condition for convergence* (Lax's theorem).

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Heuristic Stability Analysis

- Write general time-dependent FDA in the form

$$\mathbf{u}^{n+1} = \mathbf{G}[\mathbf{u}^n], \quad (98)$$

- \mathbf{G} is some *update operator* (linear in our example problem)
- \mathbf{u} is a column vector containing sufficient unknowns to write the problem in first-order-in-time form.
- Example: introduce new, auxiliary set of unknowns, v_j^n , defined by

$$v_j^n = u_j^{n-1},$$

then can rewrite differenced-wave-equation (16) as

$$u_j^{n+1} = 2u_j^n - v_j^n + \lambda^2 (u_{j+1}^n - 2u_j^n + u_{j-1}^n), \quad (99)$$

$$v_j^{n+1} = u_j^n, \quad (100)$$

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Heuristic Stability Analysis

- Thus with

$$\mathbf{u}^n = [u_1^n, v_1^n, u_2^n, v_2^n, \dots, u_J^n, v_J^n],$$

(for example), (99-100) is of the form (98).

- Equation (98) provides compact way of describing the FDA solution.
- Given initial data, \mathbf{u}^0 , solution after n time-steps is

$$\mathbf{u}^n = \mathbf{G}^n \mathbf{u}^0, \quad (101)$$

where \mathbf{G}^n is the n -th power of the matrix \mathbf{G} .

- Assume that \mathbf{G} has a complete set of orthonormal eigenvectors

$$\mathbf{e}_k, \quad k = 1, 2, \dots, J,$$

and corresponding eigenvalues

$$\mu_k, \quad k = 1, 2, \dots, J,$$

Heuristic Stability Analysis

- Thus have

$$\mathbf{G} \mathbf{e}_k = \mu_k \mathbf{e}_k, \quad k = 1, 2, \dots, J.$$

- Can then write initial data as (spectral decomposition):

$$\mathbf{u}^0 = \sum_{k=1}^J c_k^0 \mathbf{e}_k,$$

where the c_k^0 are coefficients.

- Using (101), solution at time-step n is

$$\mathbf{u}^n = \mathbf{G}^n \left(\sum_{k=1}^J c_k^0 \mathbf{e}_k \right) \quad (102)$$

$$= \sum_{k=1}^J c_k^0 (\mu_k)^n \mathbf{e}_k. \quad (103)$$

Heuristic Stability Analysis

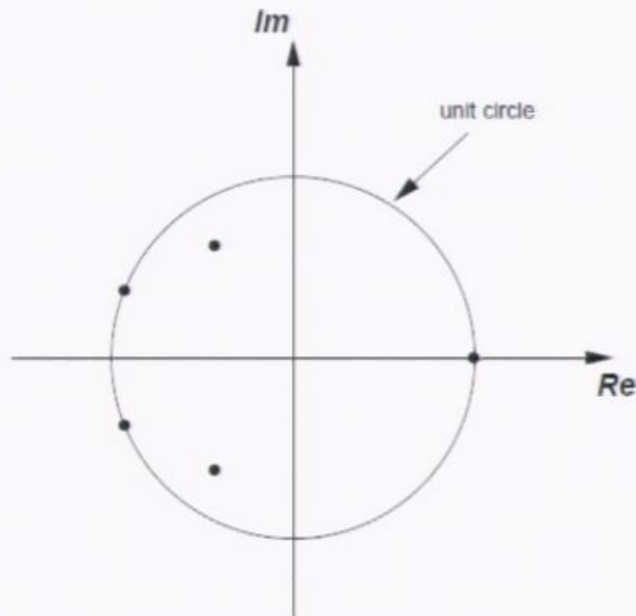
- If difference scheme is to be stable, must have

$$|\mu_k| \leq 1 \quad k = 1, 2, \dots, J \quad (104)$$

(Note: μ_k will be complex in general, so $|\mu|$ denotes the complex modulus, $|\mu| \equiv \sqrt{\mu\mu^*}$).

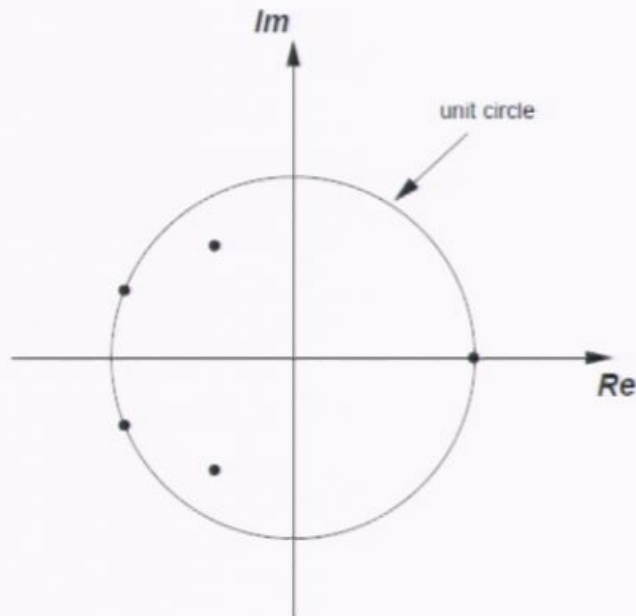
- Geometric interpretation: eigenvalues of the update matrix must lie on or within the unit circle

Heuristic Stability Analysis



- Schematic illustration of location in complex plane of eigenvalues of update matrix \mathbf{G} .
- In this case, all eigenvalues (dots) lie on or within the unit circle, indicating that the corresponding finite difference scheme is stable.

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Von-Neumann (Fourier) Stability Analysis (Summary)

- Von-Neumann (VN) stability analysis based on the ideas sketched above
- Assumes that difference equation is linear with constant coefficients, periodic boundary conditions boundary conditions are periodic
- Can then use Fourier analysis: difference operators in real-space variable $x \longrightarrow$ algebraic operations in Fourier-space variable k
- VN applied to wave-equation example shows that must have

$$\lambda \equiv \frac{\Delta t}{\Delta x} \leq 1,$$

for stability of scheme (16).

- Condition is often called the CFL condition—after Courant, Friedrichs and Lewy who derived it in 1928
- This type of instability has “physical” interpretation, often summarized by the statement *the numerical domain of dependence of an explicit difference scheme must contain the physical domain of dependence*

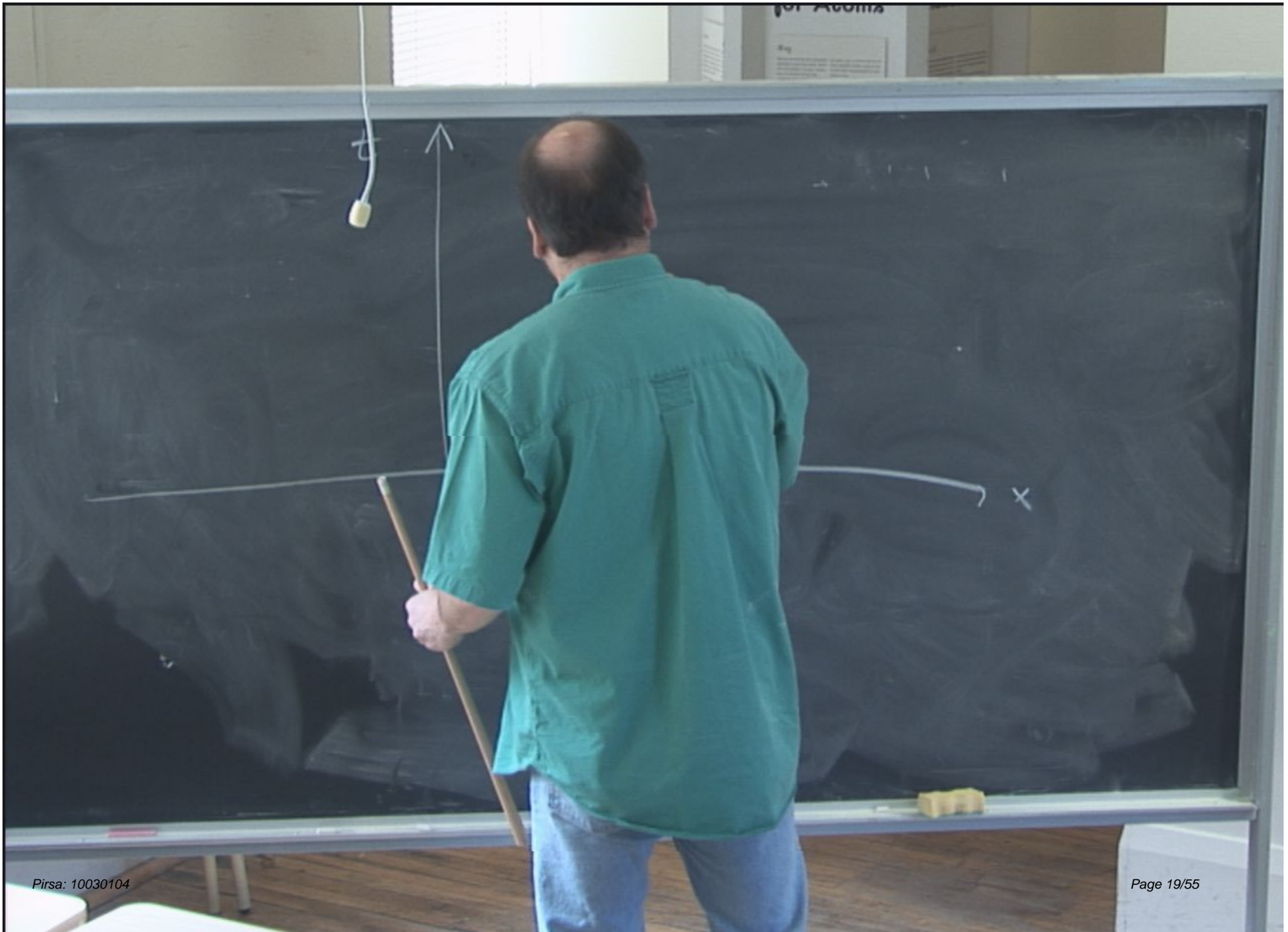
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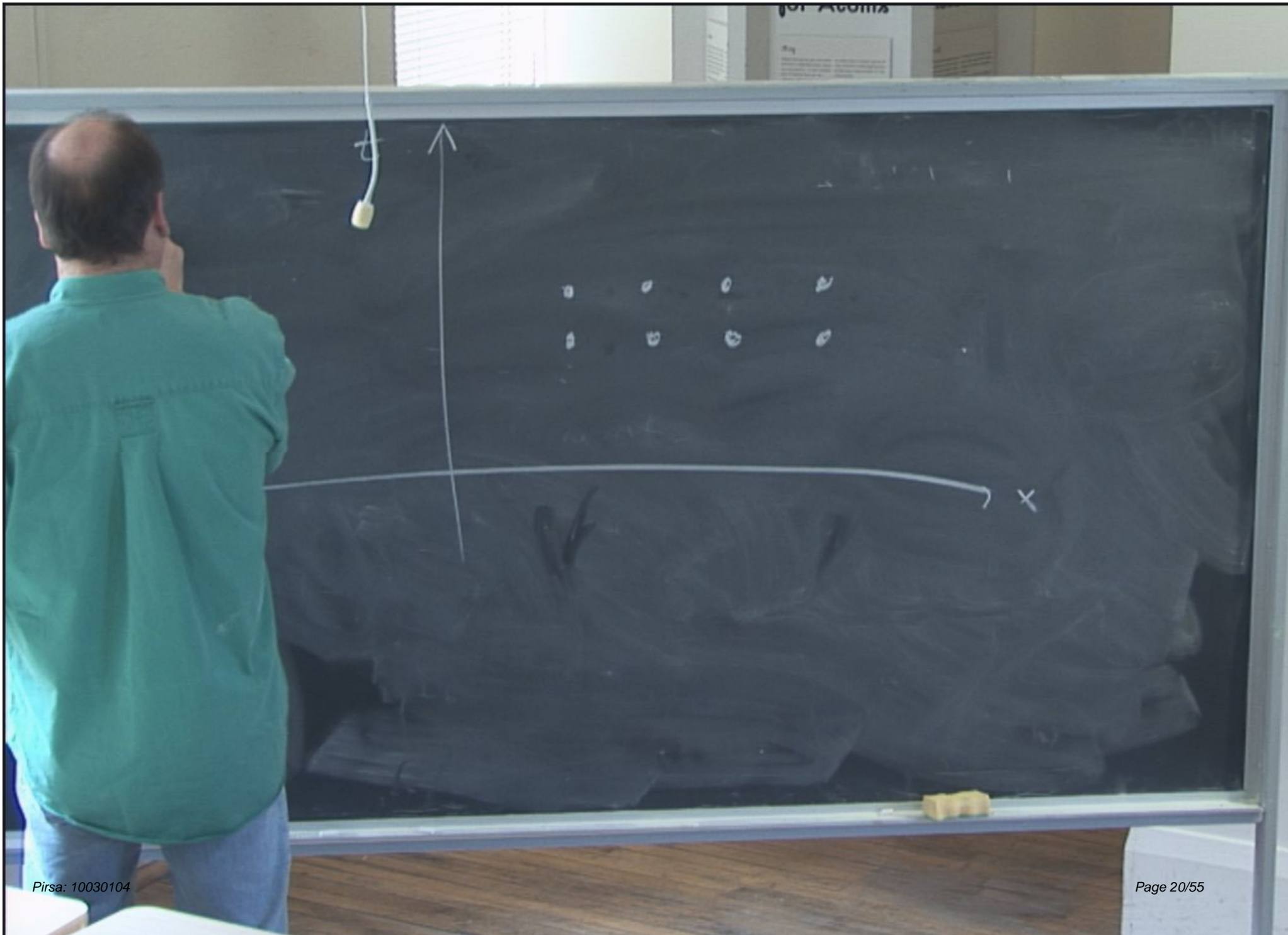
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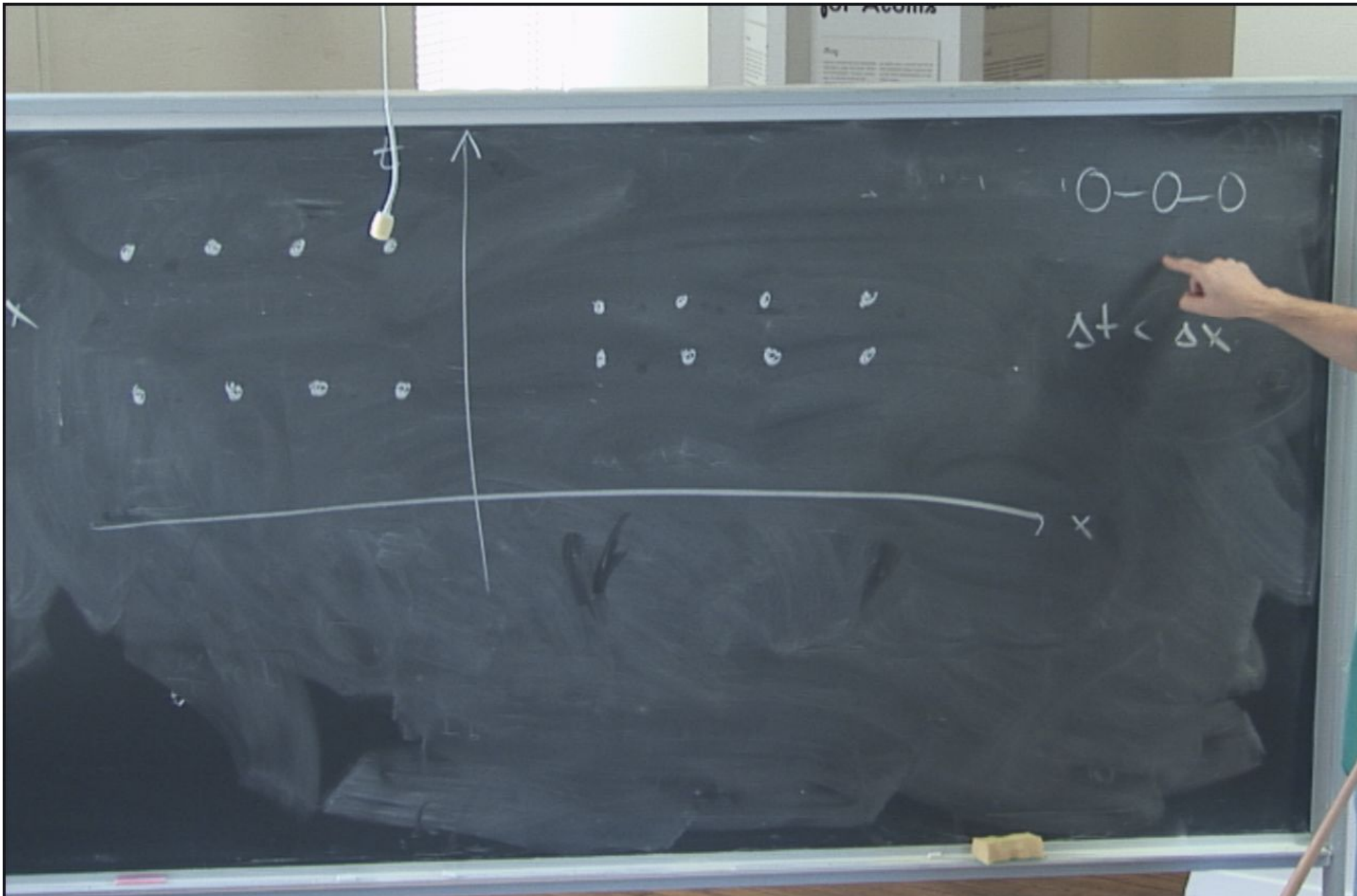
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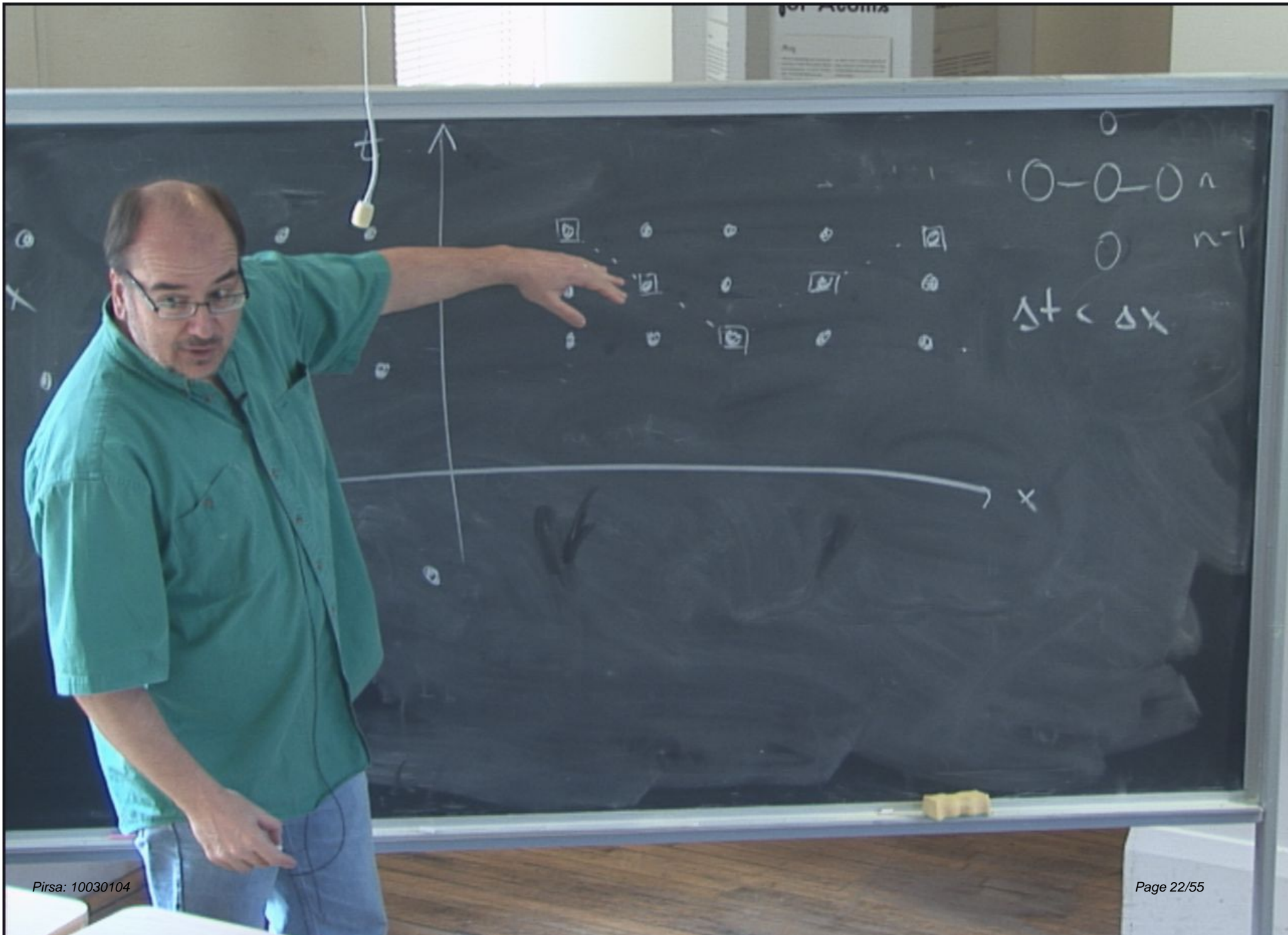
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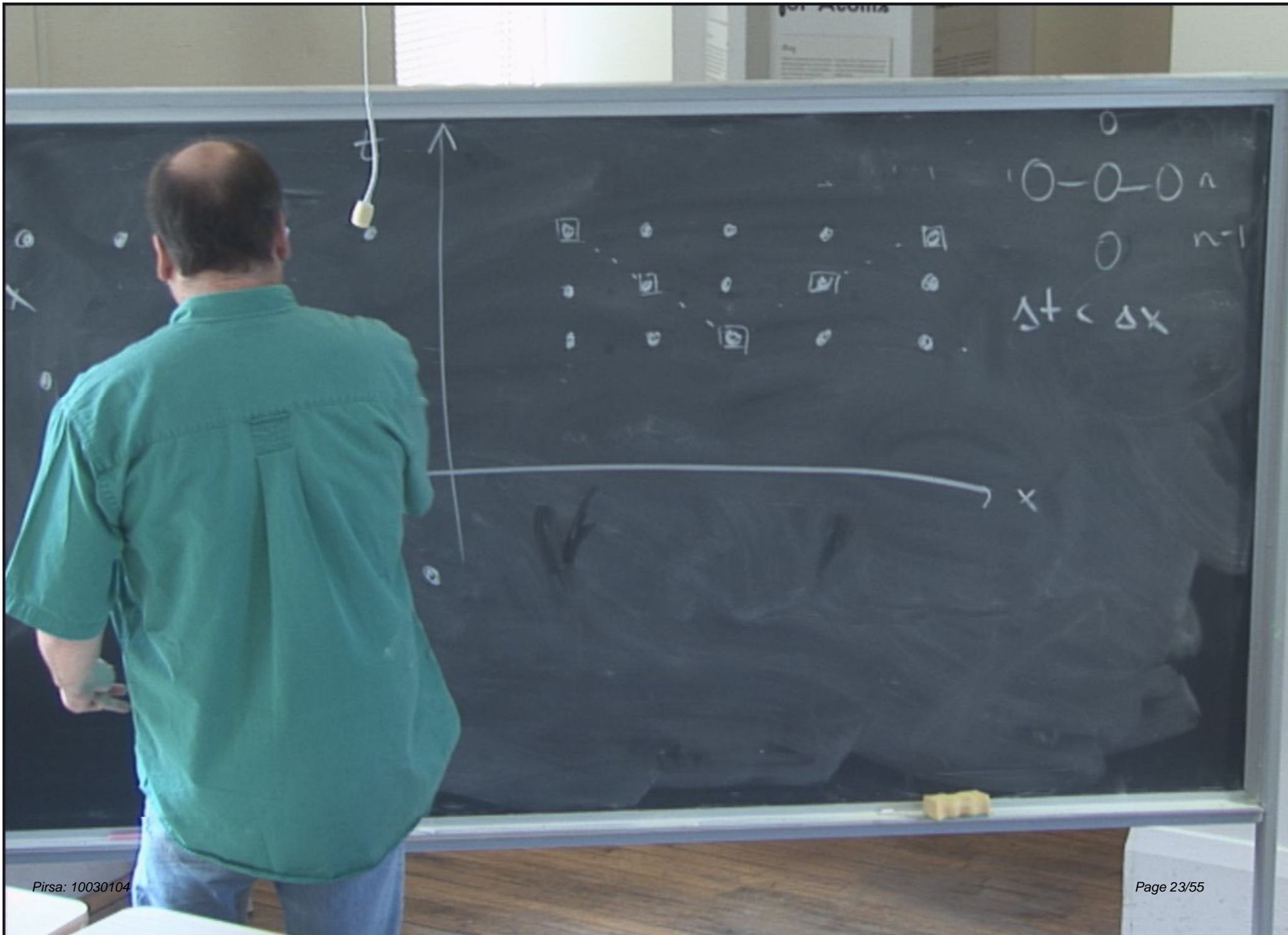
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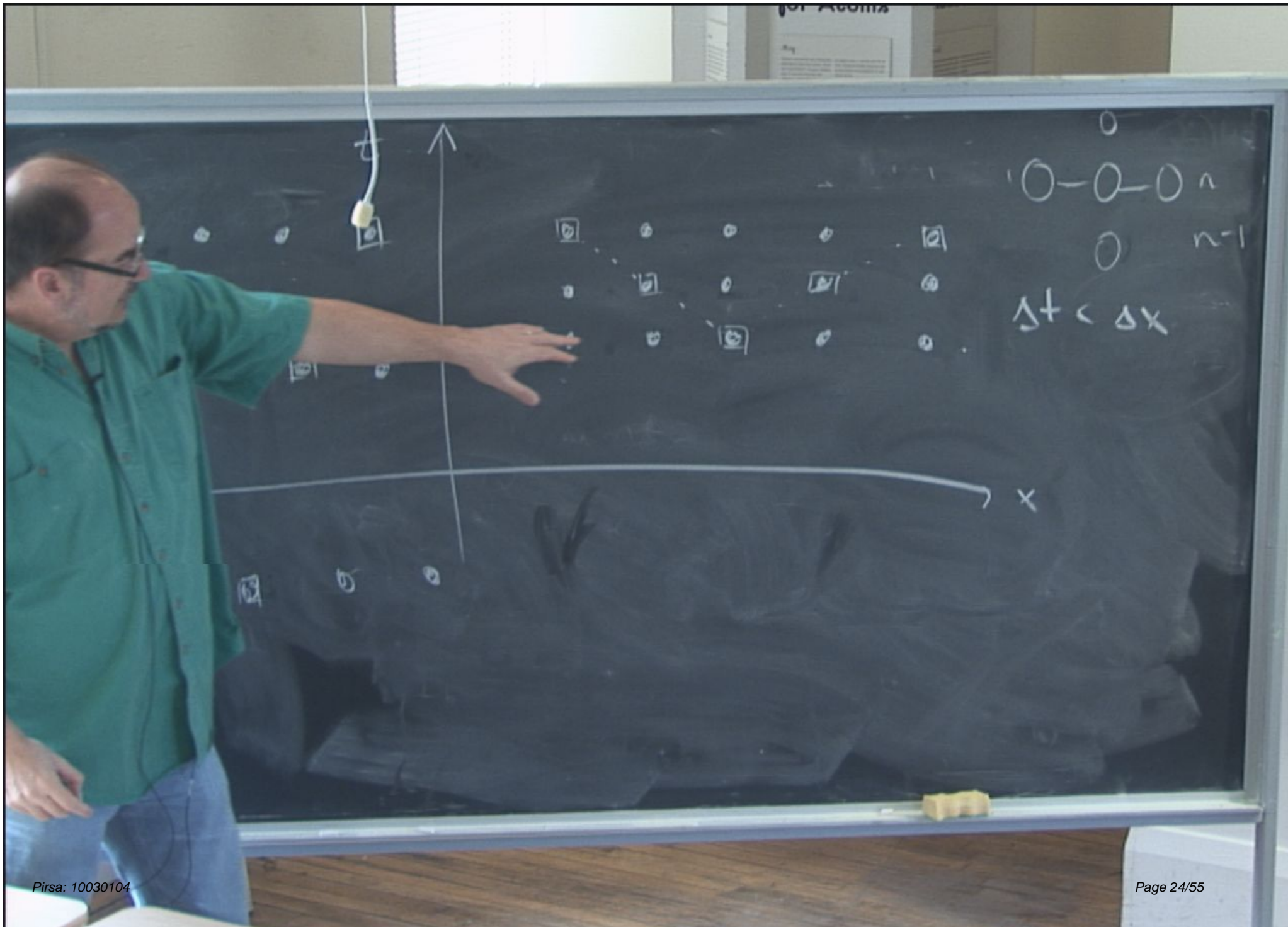


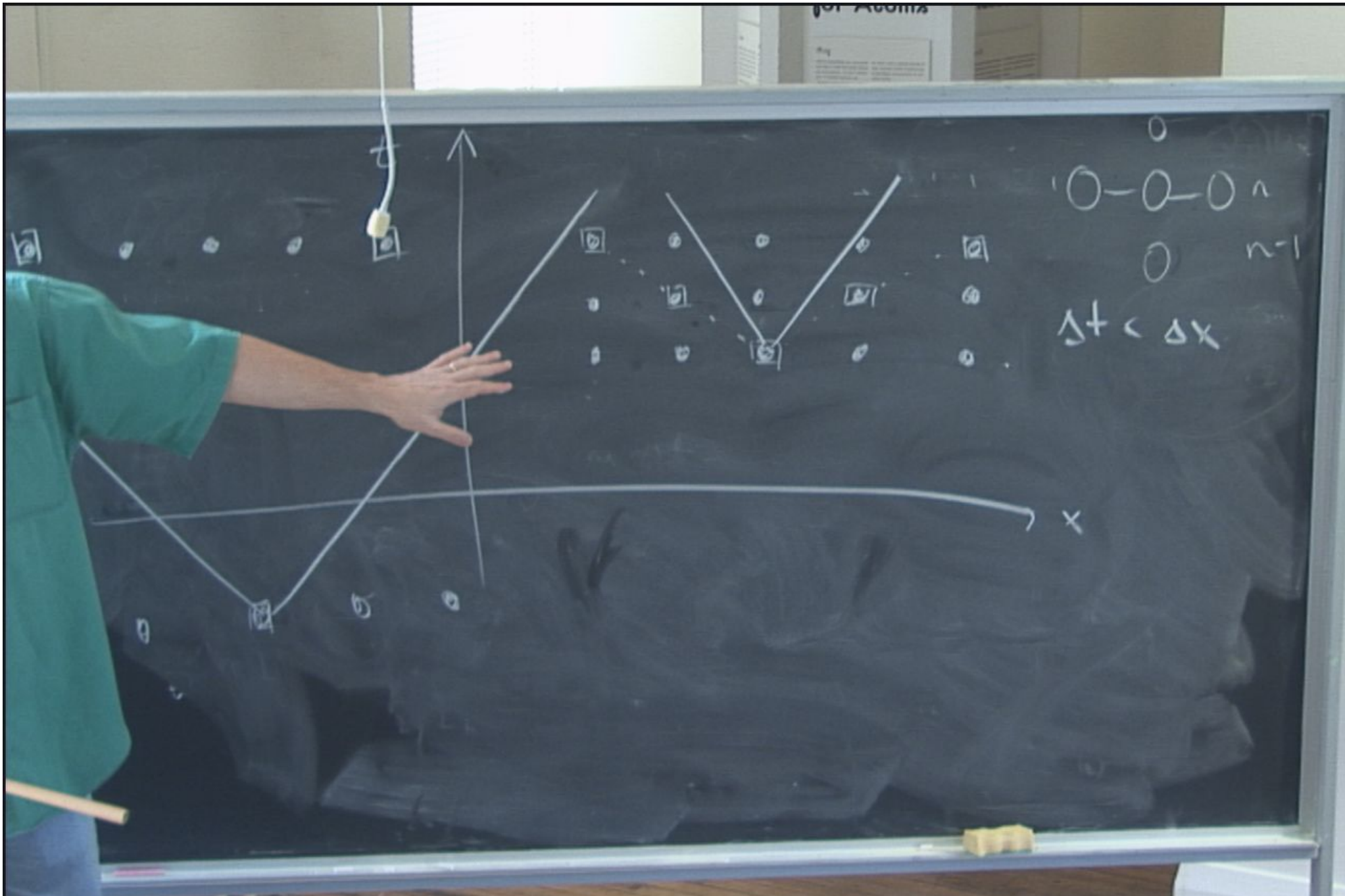


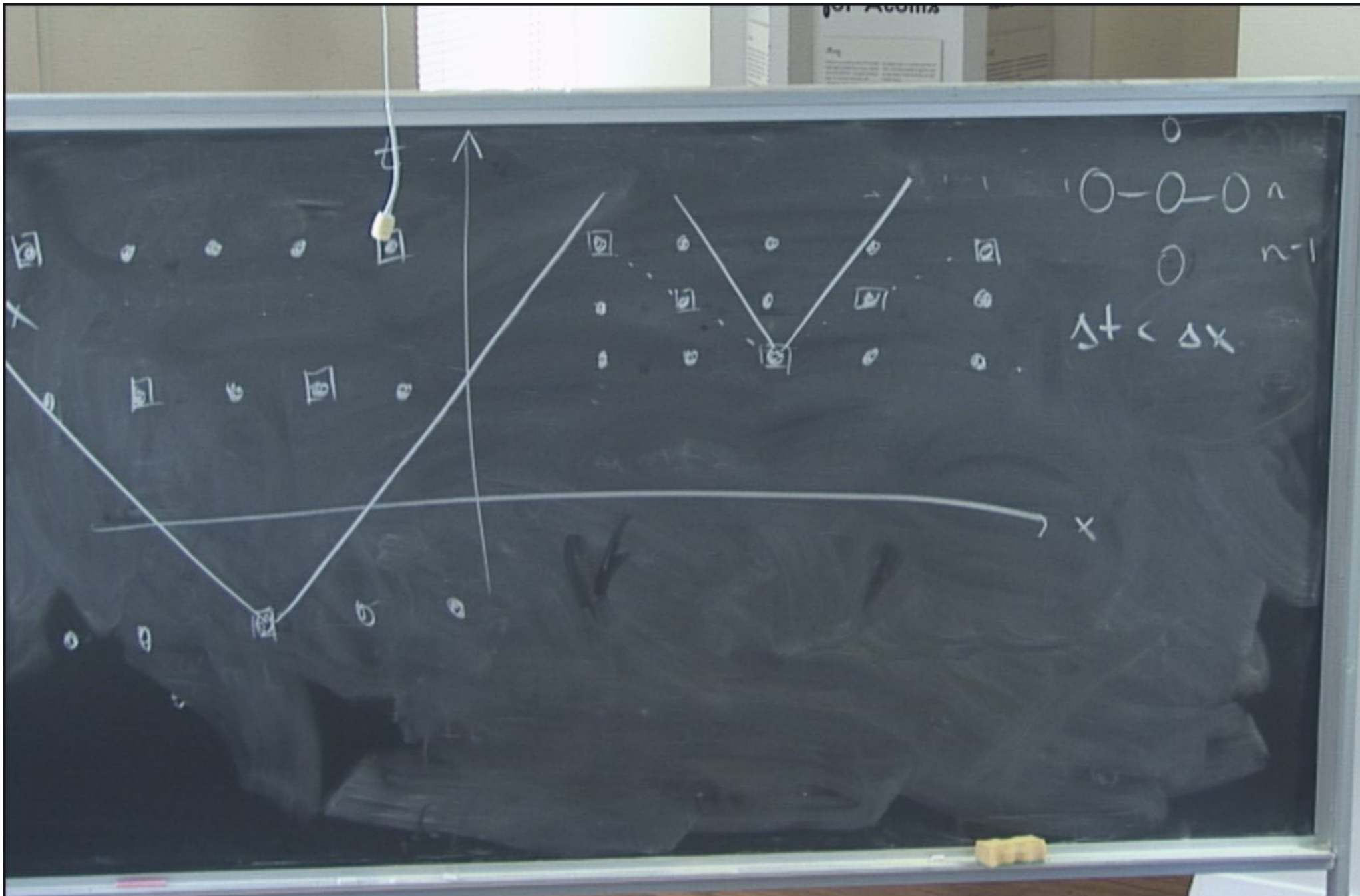


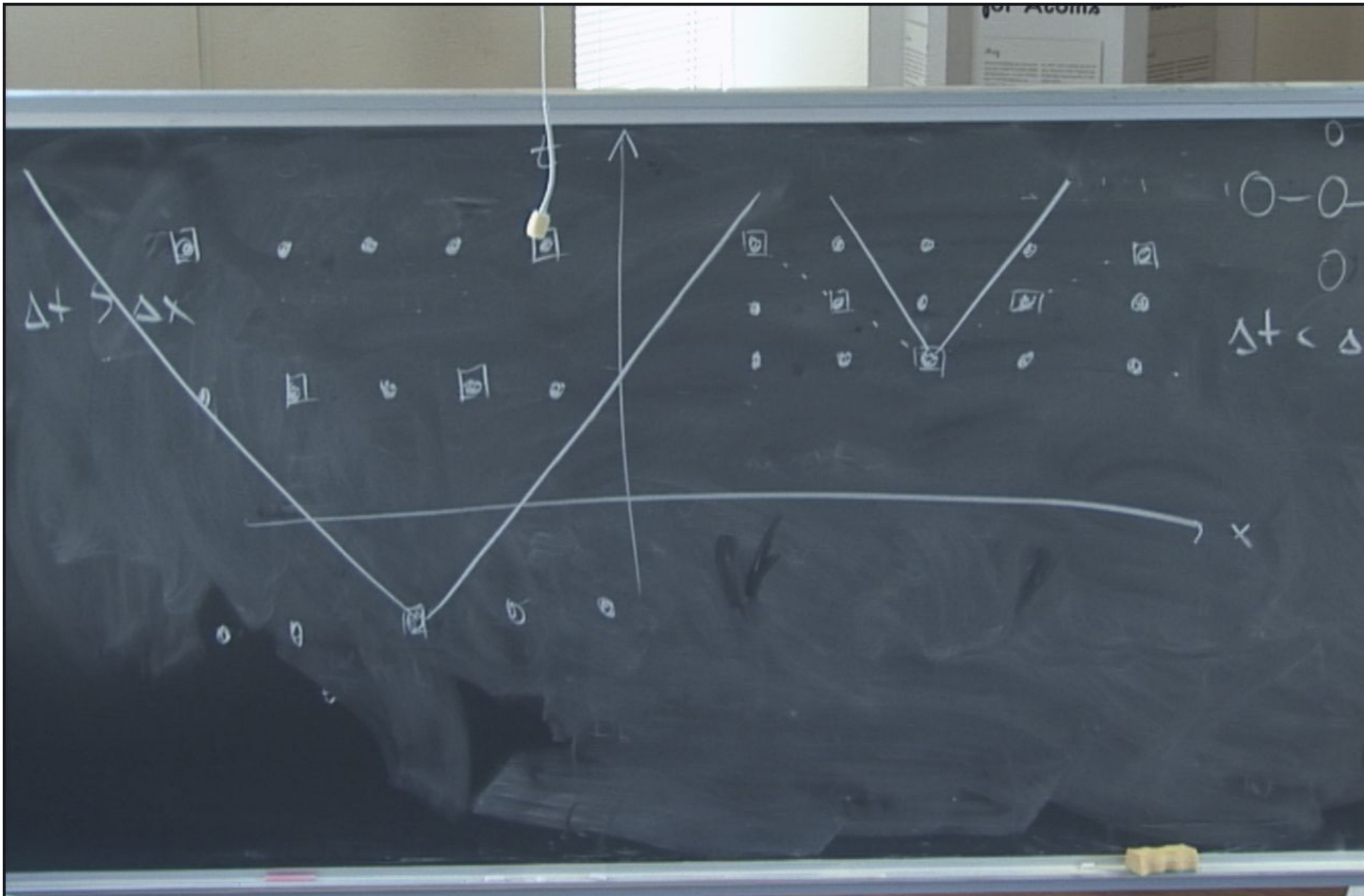












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$$\pi = \frac{\alpha}{2} (\dot{\phi} - \beta \phi')$$

$$\rightarrow \dot{\phi} = \frac{\alpha}{2} \pi + \beta \phi' = \frac{\alpha}{2} \pi + \beta \underline{\dot{\Phi}}$$

← BUT $\dot{\phi}' = \underline{\dot{\Phi}}$, so

$$\underline{\dot{\Phi}} = \left(\beta \underline{\dot{\Phi}} + \frac{\alpha}{2} \pi \right)' \quad (32)$$

• TO FIND π EQN. I, RECALL THAT

$$\square \phi = \frac{1}{\sqrt{-g}} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi)$$

$$\square \phi = 0 \rightarrow \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi) = 0$$

$$\rightarrow \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi) + \partial_\mu (\sqrt{-g} g^{\mu\rho} \partial_\rho \phi)$$

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PHYS 382M SPHERICAL SYMMETRY (6)

CHARACTERISTIC ANALYSIS OF THE SCALAR FIELD

REF: COUPANT : HILBERT "METHODS OF MATHEM. PHYS",
VOL II, ch 5

• EQNS (37)-(38) ARE A 1ST-ORDER, QUASI-LINEAR SYS
FOR OUR RADIATION FIELD. DEFINING

$$u = (\mathbf{E}, \pi)^T$$

WE CAN WRITE

$$u_t + A u_x = B \tag{39}$$

$$A = - \begin{pmatrix} \beta & \chi/a \\ \chi/a & \beta \end{pmatrix} \tag{40}$$

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$$|A - \tau I| = 0$$

$$\rightarrow \tau = -\beta \pm \frac{\alpha}{a} \tag{41}$$

THESE ARE THE "LOCAL SIGNAL SPEEDS" FOR THE SCALAR

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• MASSLESS SCALAR FIELD - WEAK (NON-SELF INTERACTING)
 SIGNALS TRAVEL ALONG NULL GEODESICS → ALTERNATE
 DERIVATION of (41)

$$ds^2 = -\alpha^2 dt^2 + a^2 (dr + \beta dt)^2 = 0$$

REGULARITY / LOCAL FLATNESS AT $r=0$

• OUR CAUCHY PROBLEM FOR THE EINSTEIN MODEL IS TO BE
 SOLVED ON

$$t \geq 0, r \geq 0$$

• BOUNDARY CONDITIONS AS $r \rightarrow 0$ WILL FOLLOW FROM ASYM-
 PTOTIC FLATNESS, NO INCOMING RADIATION; $r=0$ NOT
 A REAL BOUNDARY, BUT COMPLETARILY (I.E. WHEN FINITE
 DIFFERENTIATING) IS EFFECTIVELY ONE

• GET CONDITIONS AT $r=0$ BY DEMANDING THAT SCALAR

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TRAPPED SURFACES / APPARENT HORIZONS

• WANT TO STUDY BLACK HOLE FORMATION; BH CHARACTERIZED BY AN EVENT HORIZON WHICH CAN ONLY BE DETERMINED ONCE COMPLETE S.T. HAS BEEN CONSTRUCTED

• USEFUL TO BE ABLE TO COMPUTE "INSTANTANEOUS" APPROX. (I.E. ON ANY HYPERSURFACE $\Sigma(t)$) TO EH. - PROVIDED BY APPARENT HORIZON \equiv OUTERMOST MARGINALLY TRAPPED SURFACE

PHI ZERO

SPHERICAL SYMMETRY

(13)

• TRAPPED SURFACE: 2-SURFACE WITH TOTALLY S^2 SUCH THAT DIVERGENCE OF OUTGOING NULL GEODESICS EMANATING FROM SURFACE < 0

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SPHERICAL SYMMETRY

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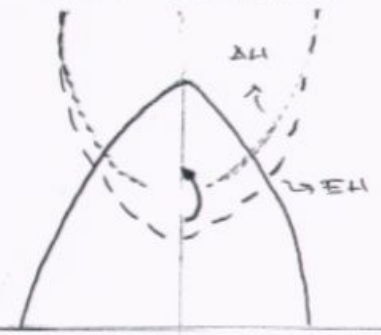
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◦ MARGINALLY TRAPPED SURFACE: " < 0 " \rightarrow " $= 0$ "

◦ MODULO COSTIC CENSORSHIP (NO NAKED SINGULARITIES)
 EXISTENCE OF $\Delta H \Rightarrow$ EXISTENCE OF EH; HOWEVER
 CAN HAVE EH WITHOUT ΔH



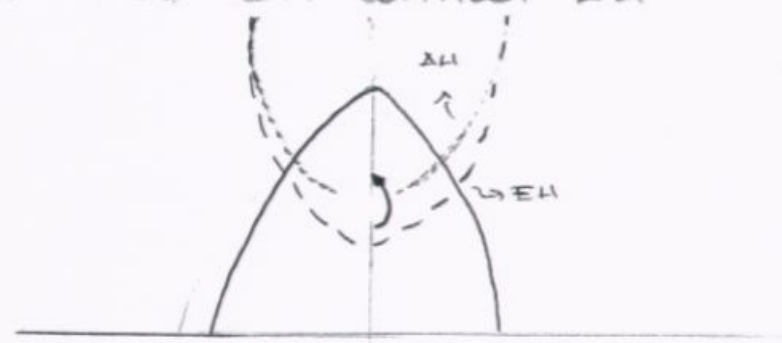
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◦ WILL TEND TO USE TERMS $\Delta H, IS$; BUT MATHS
 CHANGIABLY, BUT SHOULD BE ALWAYS Δ DISTINCTIONS

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(MARGINALLY) TRAPPED SURFACE EGM (ΔH EGM)

CONSIDER A 2-SURFACE WITH OUTGOING NULL DIRECTION u^a WHICH IS MARGINALLY TRAPPED, THEN

$$\nabla_a u^a = 0$$

NOW, CAN WRITE u^a AS

$$u^a = s^a + n^a$$

$\left\{ \begin{array}{l} \hookrightarrow \text{UNIT FUTURE-DIRECTED TIMELIKE NORMAL TO } \Sigma \\ \hookrightarrow \text{UNIT OUTWARDS-POINTING SPACELIKE NORMAL TO 2-SURF.} \end{array} \right.$

IN \neq TO 3+1 DECOMPOSITION, METRIC h_{ab} , h^{ab} IS INDUCED ON THE 2-SURFACE BY PROJECTION

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IN \neq TO 3+1 DECOMPOSITION, METRIC h_{ab} , h^{ab} IS INDUCED ON THE 2-SURFACE BY PROJECTION

$$h^{ab} = \gamma^{ab} - s^a s^b = g^{ab} + n^a n^b - s^a s^b$$

CAN SHOW (EXERCISE) THAT $\nabla_a u^a$ IS A "2-VECTOR";
 I.E. IS INTRINSIC TO 2-SURFACE; I.E. DOES NOT DEPEND ON
 HOW 2-SURFACE IS EMBEDDED IN Σ , THEN

$$\nabla_a u^a = g^{ab} \nabla_a u_b = h^{ab} \nabla_a u_b$$

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$$u^a = s^a + n^a$$

UNIT FUTURE-DIRECTED TIME-LIKE NORMAL TO Σ
 UNIT OUTWARDS-POINTING SPACE-LIKE NORMAL TO 2-SURF.

IN $3+1$ DECOMPOSITION, METRIC h_{ab} , h^{ab} IS INDUCED ON THE 2-SURFACE BY PROJECTION

$$h^{ab} = \gamma^{ab} - s^a s^b = g^{ab} + n^a n^b - s^a s^b$$

CAN SHOW (EXERCISE) THAT $\nabla_a u^a$ IS A "2-VECTOR"; I.E. IS INTRINSIC TO 2-SURFACE; I.E. DOES NOT DEPEND ON HOW 2-SURFACE IS EMBEDDED IN Σ , THEN

$$\begin{aligned} \nabla_a u^a &= g^{ab} \nabla_a u_b = h^{ab} \nabla_a u_b \\ &= h^{ab} \nabla_a (s_b + n_b) \end{aligned}$$

$$= h^{ab} \nabla_a (s_b + n_b)$$

h^{ab} PROJECTS OUT

$$h_{ab} = \gamma_{ab} - s^a s^b = g^{ab} + n^a n^b - s^a s^b$$

CAN SHOW (EXERCISE) THAT $\nabla_a u^a$ IS A "2-VECTOR";
 I.E. IS INTRINSIC TO 2-SURFACE; I.E. DOES NOT DEPEND ON
 HOW 2-SURFACE IS EMBEDDED IN Σ , THEN

$$\begin{aligned} \nabla_a u^a &= g^{ab} \nabla_a u_b = h^{ab} \nabla_a u_b \\ &= h^{ab} \nabla_a (s_b + n_b) \\ &= h^{ab} \perp \nabla_a (s_b + n_b) \\ &= h^{ab} (D_a s_b + \perp \nabla_a n_b) \\ &= h^{ab} (D_a s_b - K_{ab}) \\ &= (g^{ab} - s^a s^b) (D_a s_b - K_{ab}) \end{aligned}$$

h^{ab} PROJECTS INTO
 2-SURFACE, SO CAN
 FIRST PROJECT INTO Σ

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How 2-SURFACE IS EMBEDDED IN Σ , THEN

$$\nabla_a u^a = g^{ab} \nabla_a u_b = h^{ab} \nabla_a u_b$$

$$= h^{ab} \nabla_a (s_b + n_b)$$

$$= h^{ab} \perp \nabla_a (s_b + n_b)$$

$$= h^{ab} (D_a s_b + \perp \nabla_a n_b)$$

$$= h^{ab} (D_a s_b - K_{ab})$$

$$= (g^{ab} - s^a s^b) (D_a s_b - K_{ab})$$

$$= D^a s_a - K + s^a s^b K_{ab}$$

$$(s^b D_a s_b = \frac{1}{2} D_a (s^b s_b) = \frac{1}{2} D_a (1) = 0)$$

h^{ab} PROJECTS OUT
 2-SURFACE, SO CAN
 FIRST PROJECT INTO Σ

How 2-SURFACE IS EMBEDDED IN Σ , THEN

$$\nabla_a u^a = g^{ab} \nabla_a u_b = h^{ab} \nabla_a u_b$$

$$= h^{ab} \nabla_a (s_b + n_b)$$

$$= h^{ab} \perp \nabla_a (s_b + n_b)$$

$$= h^{ab} (D_a s_b + \perp \nabla_a n_b)$$

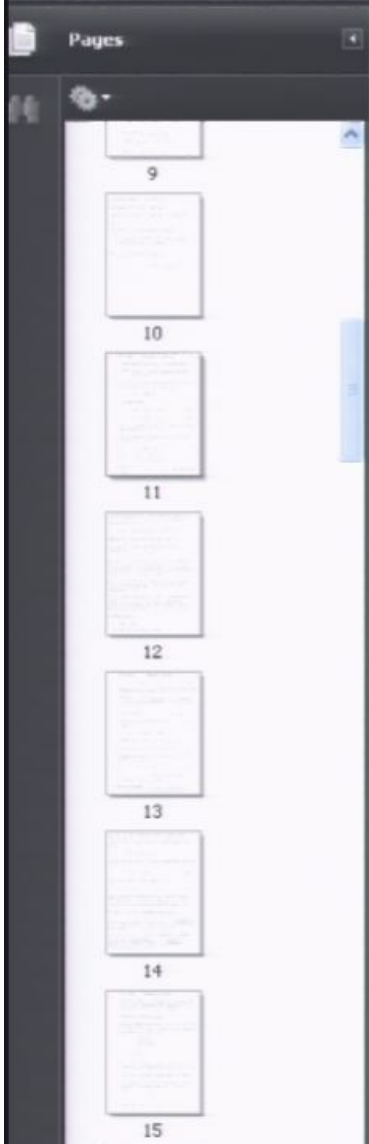
$$= h^{ab} (D_a s_b - K_{ab})$$

$$= (g^{ab} - s^a s^b) (D_a s_b - K_{ab})$$

$$= D^a s_a - K + s^a s^b K_{ab}$$

$$(s^b D_a s_b = \frac{1}{2} D_a (s^b s_b) = \frac{1}{2} D_a (1) = 0)$$

h^{ab} PROJECTS ONTO
 2-SURFACE, SO CAN
 FIRST PROJECT INTO Σ



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PHY 321

SPHERICAL SYMMETRY

⑤

THUS, OUR TRAPPED SURFACE (ΔH) EQUATION IS

$$D^a s_a - K + s^a s^b K_{ab} = 0 \quad (47)$$

AND ARGUING AS WE DID FOR THE 3+1 EGNS WE GET A VALID COMPONENT PART OF THIS EQUATION BY TAKING $a \rightarrow i, b \rightarrow j$

$$D^i s_i - K + s^i s^j K_{ij} = 0 \quad (48)$$

SPECIALIZING NOW TO SPHERICAL SYMMETRY

$$ds^2 = a^2 dr^2 + r^2 h^2 d\Omega^2$$

$$r_{ij} s^i s^j = 1 \rightarrow s^i = (a^{-2}, e, e)$$

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AND ARGUING AS WE DID FOR THE 3+1 EGNS WE GET A VALID COMPONENT PART OF THIS EQUATION BY TAKING $a \rightarrow i, b \rightarrow j$

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$$r_{ij} s^i s^j = 1 \rightarrow s^i = (a^{-2}, e, e)$$

$$D_i s^i = r^{-1/2} \partial_i (r^{1/2} s^i) \quad v^i = ar^2 b^2$$

$$= \frac{1}{ar^2 b^2} (r^2 b^2)' = \frac{2(rb)'}{arb}$$

THUS, (48) BECOMES

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$$r_{ij} s^j s^i = 1 \rightarrow s^i = (a^{-2}, e, e)$$

$$D_i s^i = r^{-\frac{1}{2}} \partial_i (r^{\frac{1}{2}} s^i) \quad v_i^2 = ar^2 b^2$$

$$= \frac{1}{ar^2 b^2} (r^2 b^2)' = \frac{2(rb)'}{arb}$$

THUS, (4E) BECOMES

$$\frac{2(rb)'}{arb} - (K^r_r + 2K^e_e) + a^{-2} K^m_m = 0$$

(K^r_r)

$$(rb)' = arb K^e_e$$

(49)

NOW, RECALL EVOL. EQN (34) FOR b

$$\dot{b} = -\alpha b K^e_e + \beta (rb)'$$

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$$r_{ij} s^j s^i = 1 \rightarrow s^i = (a^{-2}, e, e)$$

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THUS, (4E) BECOMES

$$\frac{2(rb)'}{arb} - (K^r_r + 2K^e_e) + a^{-2} K^m_m = 0$$

(K^r_r)

$$(rb)' = arb K^e_e$$

(49)

NOW, RECALL EVOL. EQN (3A) FOR b

$$\dot{b} = -\alpha b K^e_e + \beta (rb)'$$



$$v = \frac{1}{ar^2b^2} \quad (r^2b^2)' = 2(rb)'$$

$$= \frac{1}{ar^2b^2} (r^2b^2)' = \frac{2(rb)'}{arb}$$

I

THUS, (4E) BECOMES

$$\frac{2(rb)'}{arb} - (K^r_r + 2K^{\theta}_{\theta}) + a^{-2} K^m_m = 0$$

$$(rb)' = arb K^{\theta}_{\theta}$$

(49)

NOW, RECALL EVOL. EQU (3A) FOR b

$$b' = -\alpha b + \dots$$

$$\Rightarrow K^{\theta}_{\theta} = -(\alpha b)^{-1} (b' - \frac{1}{2} (rb)')$$

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NOW, RECALL EVOL. EQN (34) FOR b

$$\dot{b} = -\alpha b K_0^e + \frac{\beta}{r} (rb)'$$

$$\Rightarrow K_0^e = -(\alpha b)^{-1} (\dot{b} - \frac{\beta}{r} (rb)')$$

$$rK_0^e = -(\alpha b)^{-1} (r\dot{b} - \beta (rb)')$$

SO (49) CAN BE REWRITTEN AS

$$(rb) + \left(\frac{\alpha}{a} - \beta \right) (rb)' = 0 \quad (50)$$

WHICH SAYS THAT THE SURFACE OF CONSTANT AREAL RADIUS $R \equiv rb$ IS OUTGOING NULL AT THE MARGINALLY TRAPPED SURFACE (M) ACCORD WITH OUR PHYSICAL PICTURE

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NOW, RECALL EVOL. EQU (34) FOR b

$$\dot{b} = -\alpha b K^0 + \frac{\beta}{r} (rb)'$$

$$\rightarrow K^0 = -(\alpha b)^{-1} (\dot{b} - \frac{\beta}{r} (rb)')$$

$$rK^0 = -(\alpha b)^{-1} (r\dot{b} - \beta (rb)')$$

SO (49) CAN BE REWRITTEN AS

$$(rb) + \left(\frac{\alpha}{a} - \beta \right) (rb)' = 0 \quad (50)$$

WHICH SAYS THAT THE SURFACE OF CONSTANT AREAL RADIUS
 $R \equiv rb$ IS ORTHOGONAL NULL AT THE MARGINALLY TRAPPED
 SURFACE IN ACCORD WITH OUR PHYSICAL PICTURE