

Title: The Phase Diagram of QCD: Results & Challenges

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Abstract: I will review the progress made in our understanding of the QCD phase diagram within an RG approach to QCD and effective QCD models. In particular this includes a discussion of the confinement-deconfinement phase transition/cross-over, the chiral phase transition/cross-over, as well as their interrelation.

# The phase diagram of QCD

## Results & Challenges

Jan M. Pawłowski  
Universität Heidelberg & ExtreMe Matter Institute

Perimeter Institute, March 5th 2010



# Outline

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- QCD phase diagram
- Quark confinement & chiral symmetry breaking
- Chiral phase structure at finite density
- Summary and outlook



## QCD phase diagram

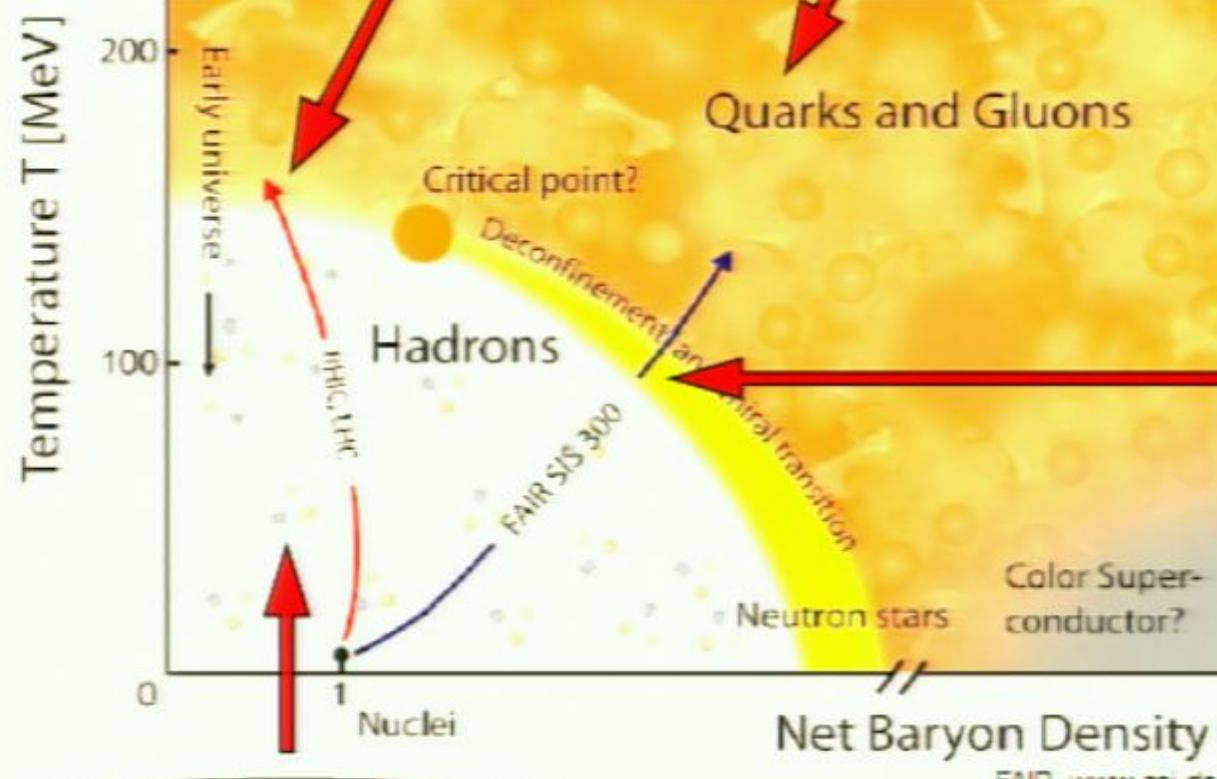
# Phase diagram of QCD

Strongly correlated quark-gluon-plasma

'RHIC serves the perfect fluid'

massless quarks (chiral symmetry)

deconfinement



hadronic phase

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confinement & chiral symmetry breaking

quarkyonic:

confinement & chiral symmetry?



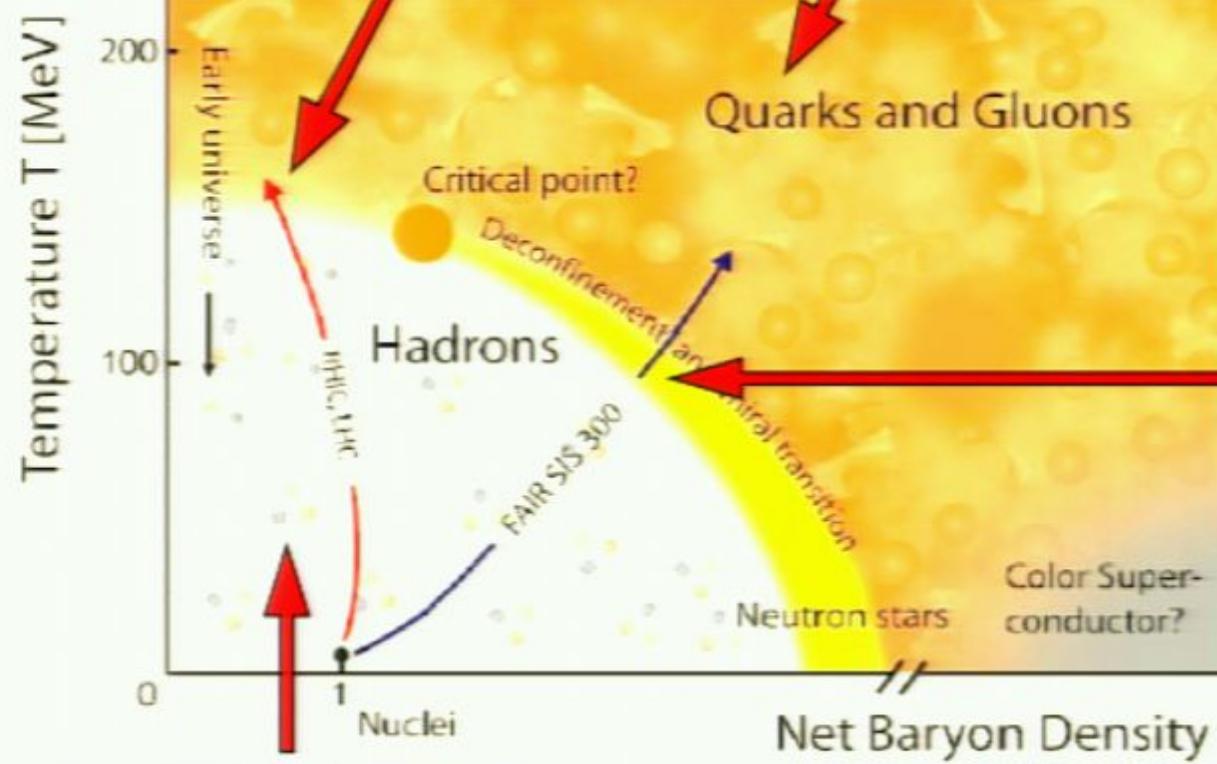
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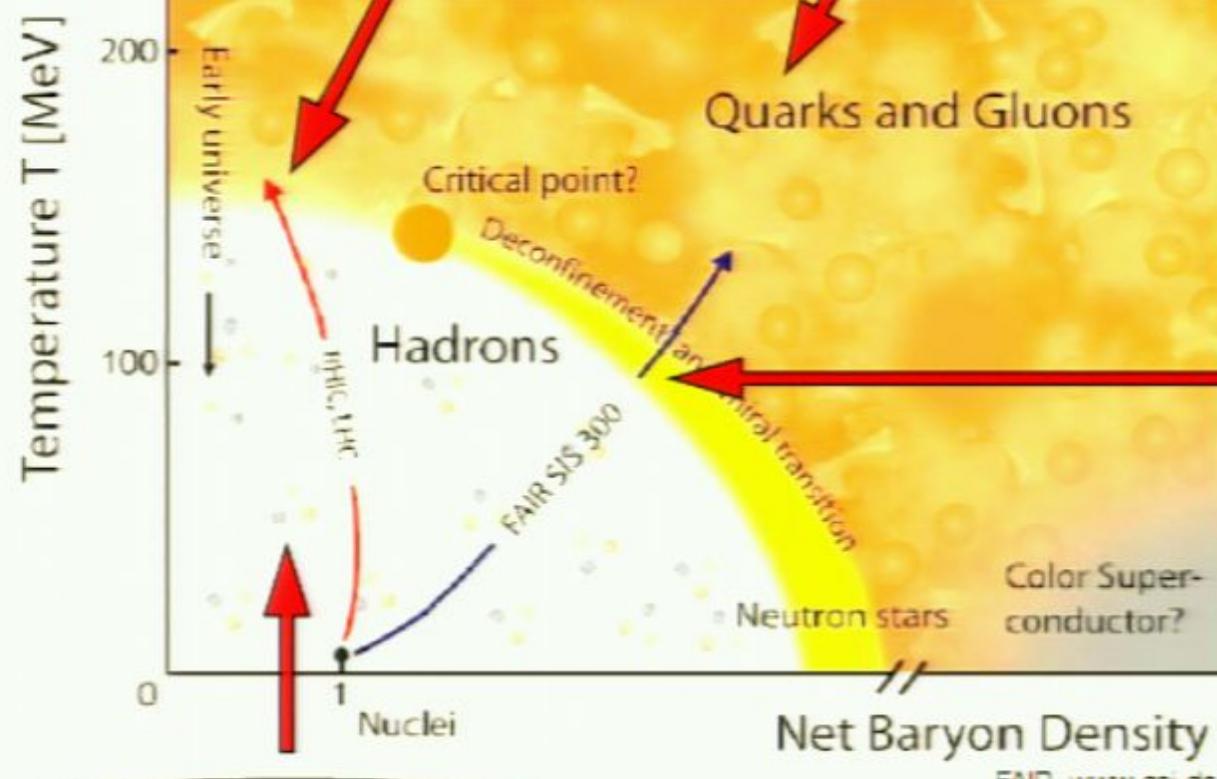
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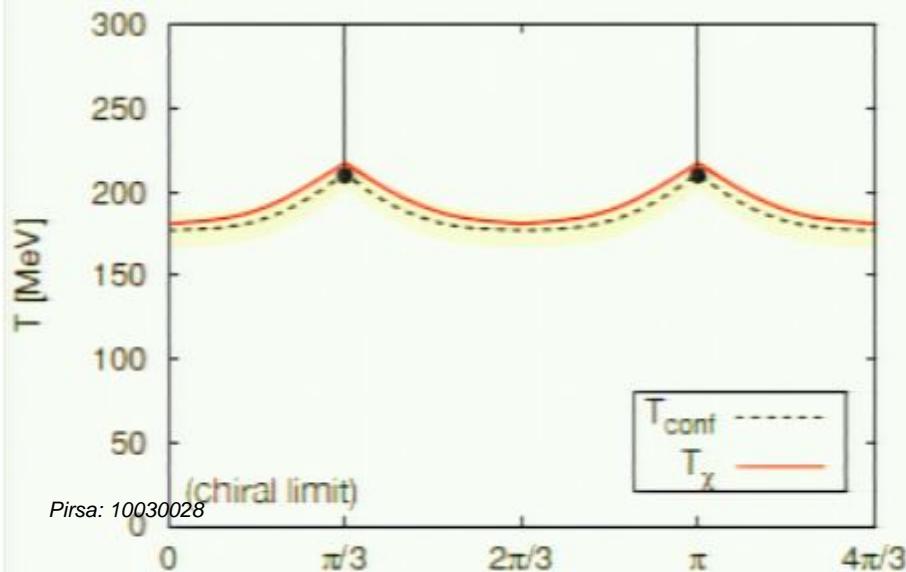
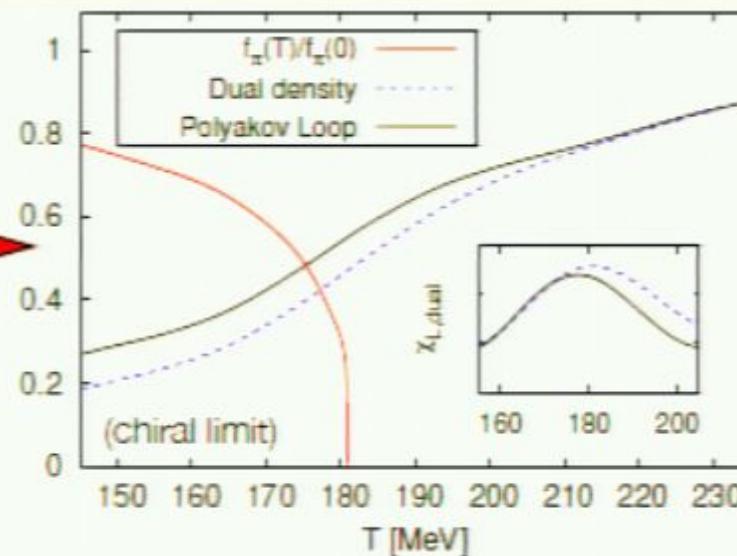
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# Phase diagram of two flavour QCD

Continuum methods

RG-flows in QCD

Braun, Haas, Marhauser, JMP '09



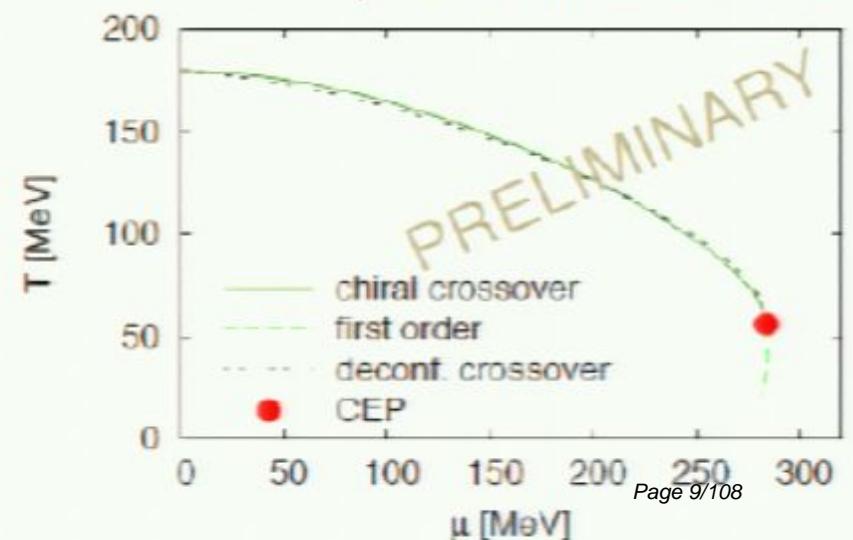
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PNJL & PQM model

Fukushima '03  
Ratti, Thaler, Weise '06

Back-coupling of matter fluctuations to YM-sector

Schaefer, JMP, Wambach '07  
Herbst, JMP, Schaefer, in prep



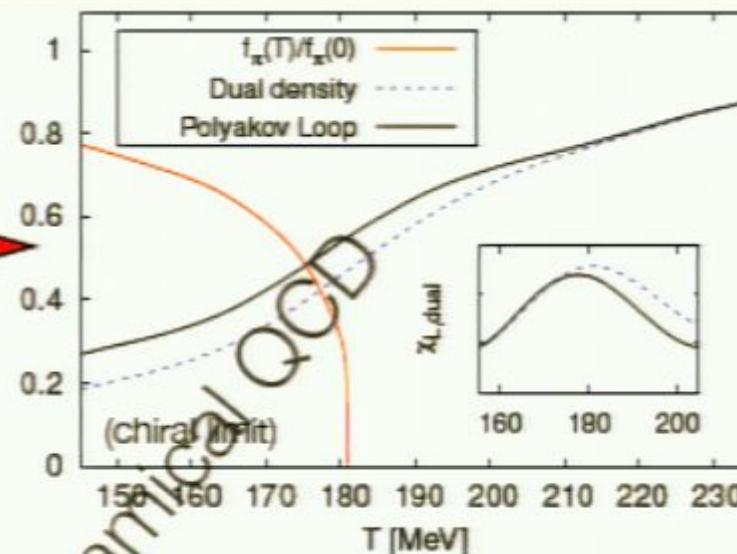
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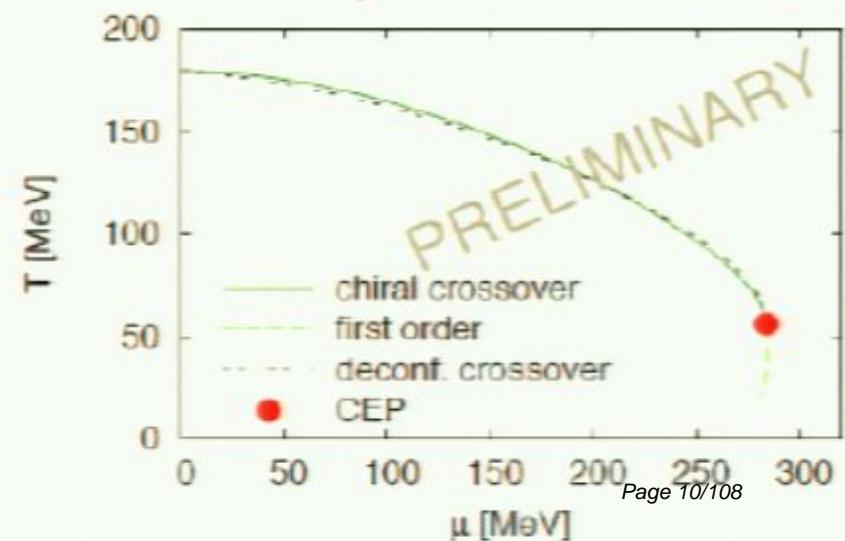
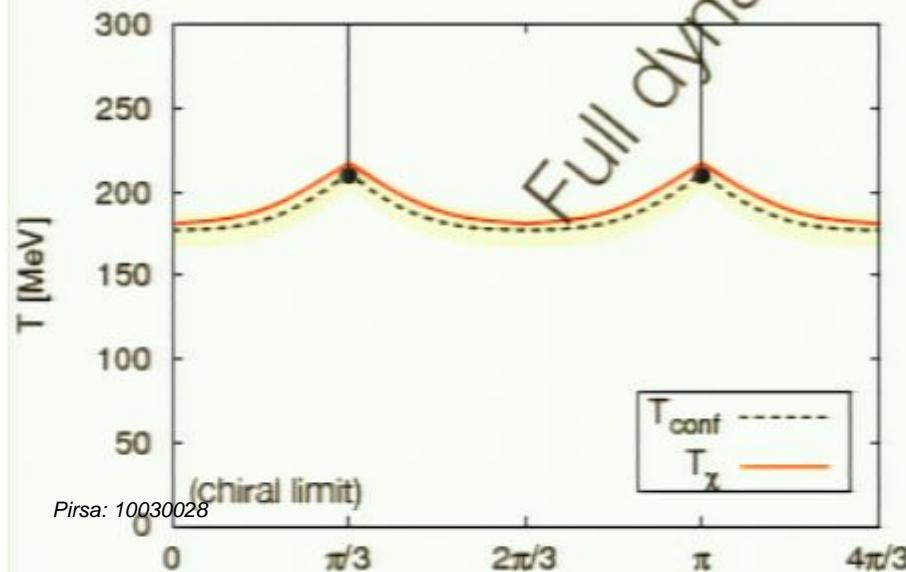


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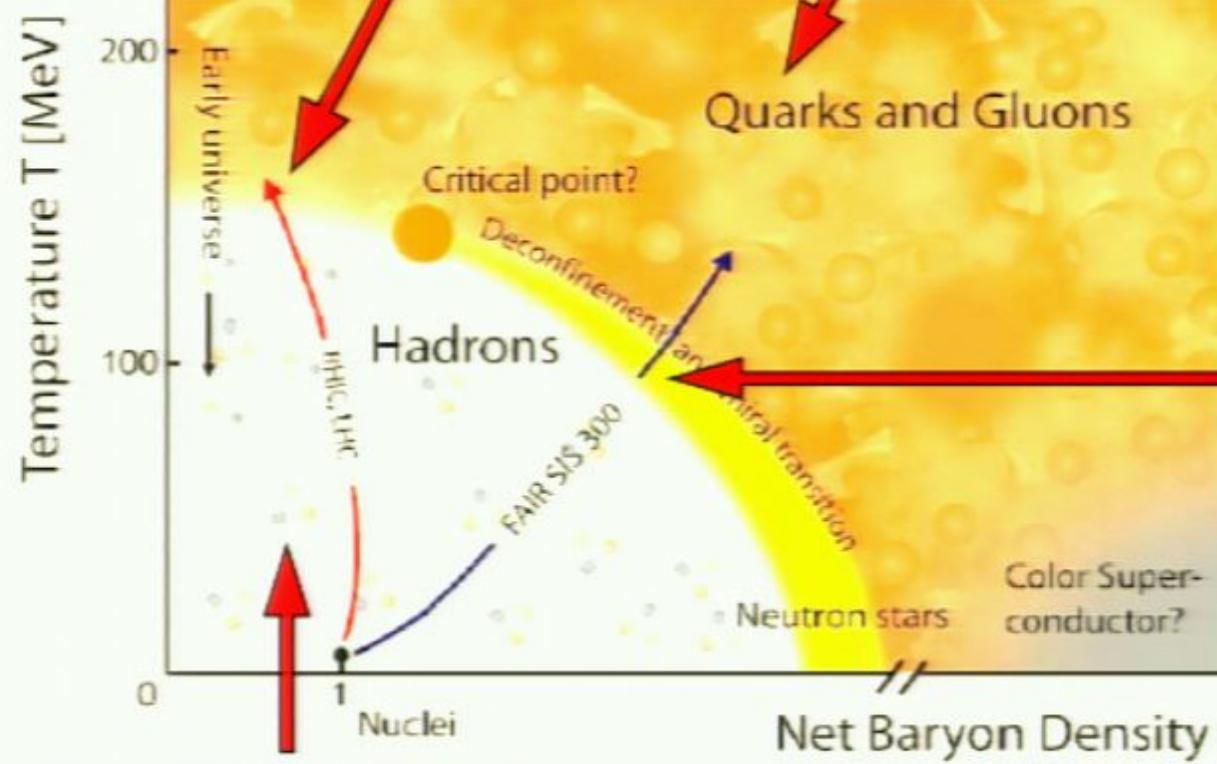
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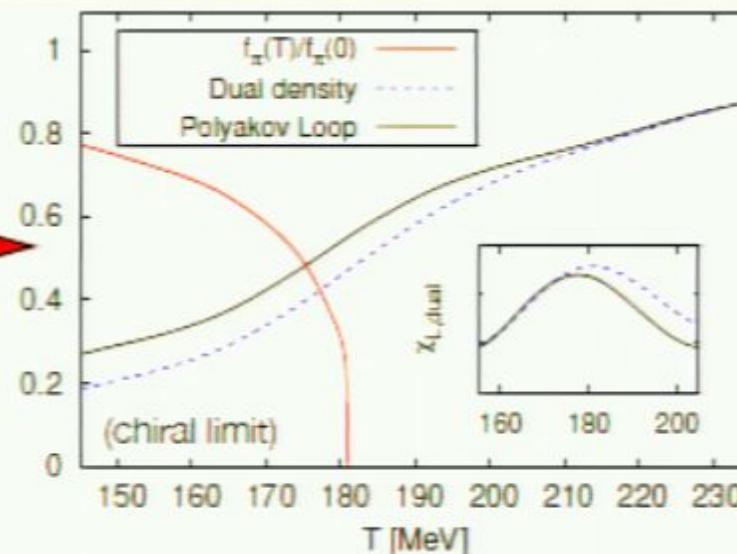
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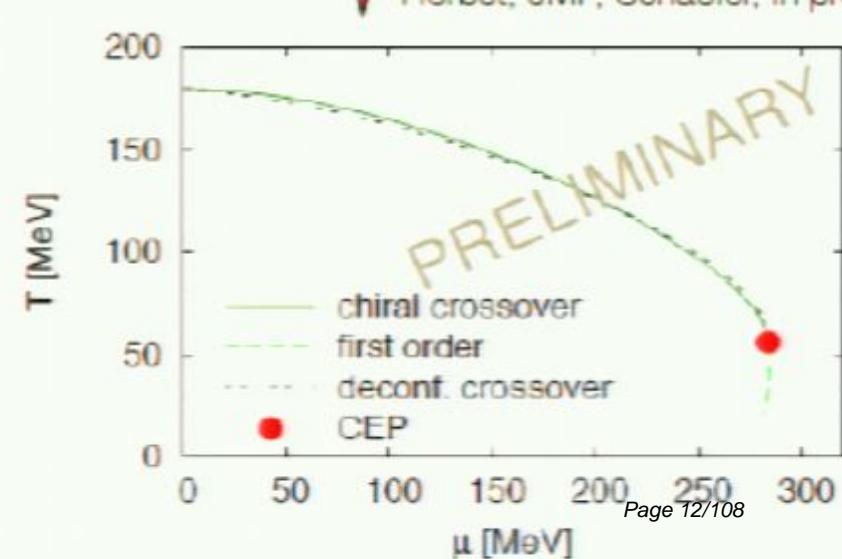
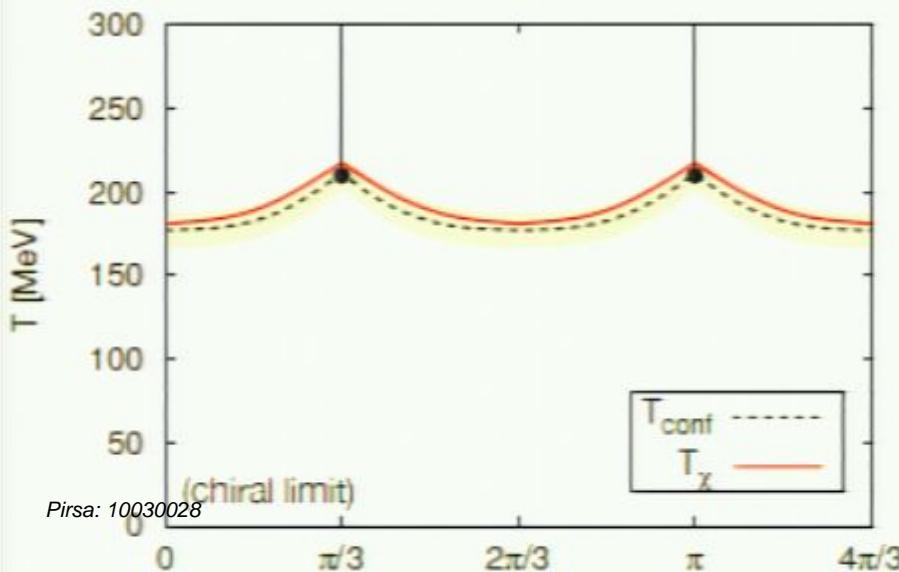


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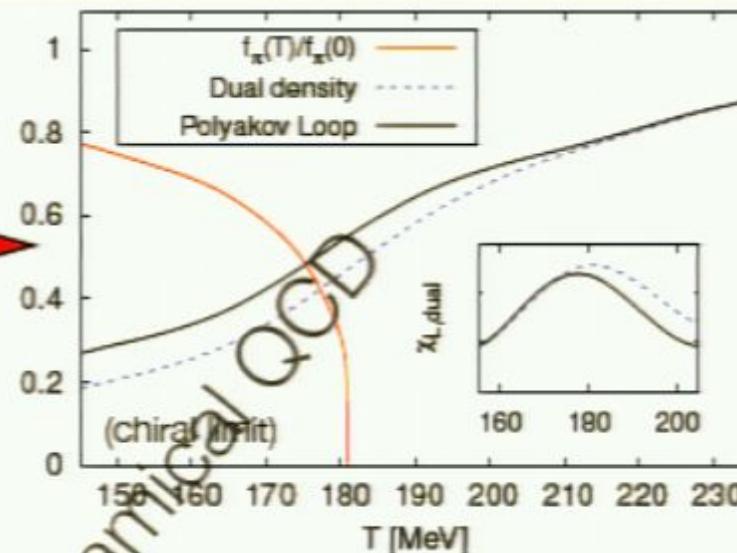


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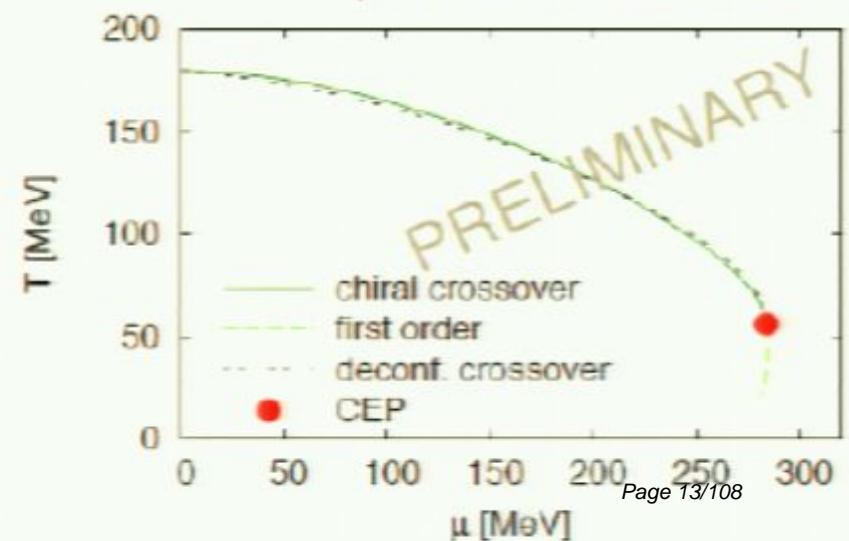
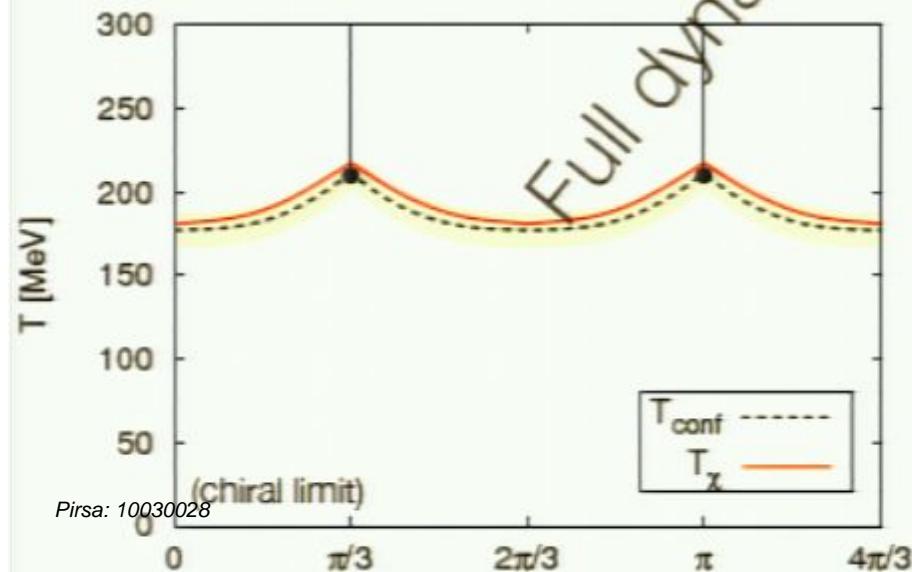


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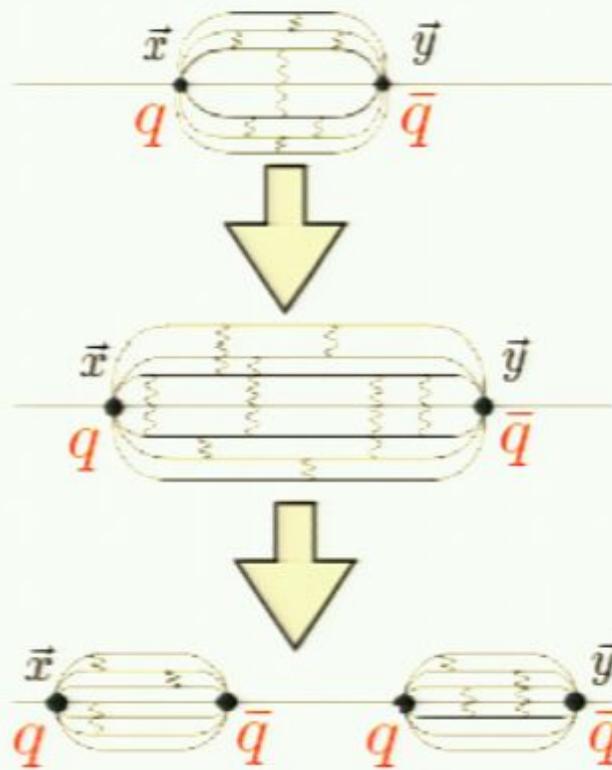


# Quark confinement & chiral symmetry breaking

# Confinement

$$r = |\vec{x} - \vec{y}|$$

Order parameter  $\sim \langle q \rangle'$



$$\Phi = e^{-\frac{1}{2}\beta F_{q\bar{q}}(\infty)}$$

- Confinement:  $\Phi = 0$
- Deconfinement:  $\Phi \neq 0$

$\Phi$  Polyakov loop

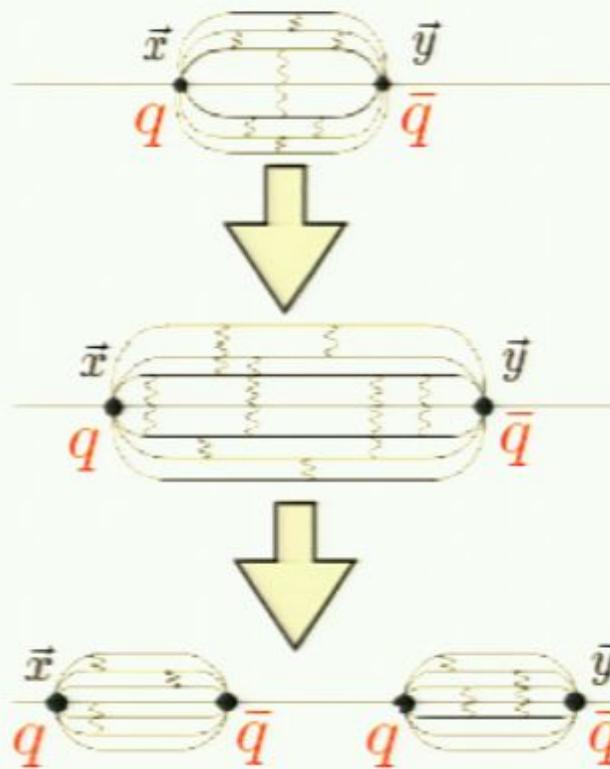
$$\Phi = \frac{1}{3} \langle \text{Tr } \mathcal{P} \exp\{ig \int_0^{1/T} dx_0 A_0\} \rangle$$



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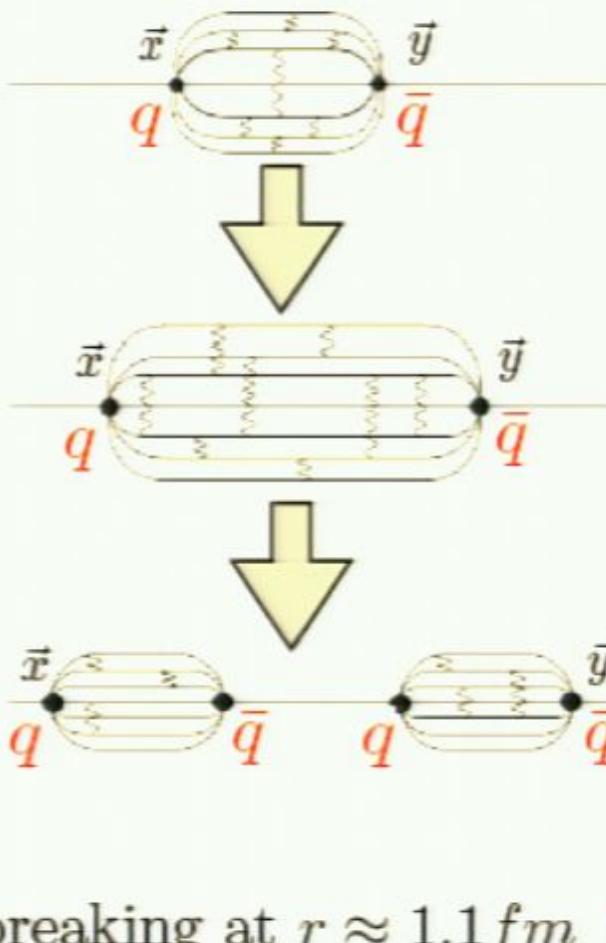
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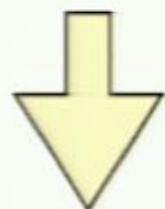
Symmetry

- $Z_3$  - symmetry:  $q \rightarrow zq$
- broken by dynamical quarks

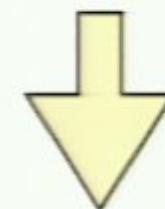
# Chiral symmetry breaking

chiral symmetry

Generation	first	second	third	Charge
Mass [MeV]	1.5-4	1150-1350	$170 \times 10^3$	
Quark	u	c	t	$\frac{2}{3}$
Quark	d	s	b	$-\frac{1}{3}$
Mass [MeV]	4-8	80-130	$(4.1-4.4) \times 10^3$	



chiral symmetry breaking:  $\Delta m \approx 400 \text{ MeV}$



2 light flavours, one heavy flavour 2 + 1

chiral symmetry breaking

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# Chiral symmetry breaking

$$\text{Diagram: Four gluons } g \text{ interact via a vertex with coupling constant } \lambda_s \text{ to form a cross symbol } \propto \alpha_s^2$$



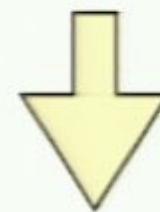
$$\int d^4x \lambda_\psi [(\bar{q}q)^2 - (\bar{q}\gamma_5 q)^2]$$

$$\langle \bar{q}q \rangle \neq 0$$

mass term:  $\langle \bar{q}q \rangle \bar{q}q$

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# Chiral symmetry breaking

$$\text{Fermion loop diagram with four external gluon lines and internal gluon lines, followed by a red arrow pointing to a quark-gluon vertex diagram with a coupling } \lambda_\psi.$$

$$\propto \alpha_s^2$$

Order parameter

$$\sigma = \langle \bar{q}q \rangle$$

chiral condensate

- chiral symmetry:  $\sigma = 0$

- symmetry breaking:  $\sigma \neq 0$

$$\int d^4x \lambda_\psi [(\bar{q}q)^2 - (\bar{q}\gamma_5 q)^2]$$

$$\langle \bar{q}q \rangle \neq 0$$



mass term:  $\langle \bar{q}q \rangle \bar{q}q$

Symmetry

- $SU_L(N_f) \times SU_R(N_f)$

- broken to  $SU(N_f)$

# Functional RG

# Functional RG

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- Introduction to Functional RG flows & some results in QCD (talks & lit)
  - Integrals from differential equations: The FRG-idea in 0+0-dimensions
  - Confinement & chiral symmetry breaking from Functional Methods
  - Aspects of the Functional RG

# Functional RG

## Functional RG

$$k\partial_k \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k(p)} k\partial_k R_k(p)$$

Wetterich '93

- Yang Mills Theory:  $\phi = (A, C, \bar{C})$  RG-scale  $k$ :  $t = \ln k$

$$\partial_t \Gamma_k[\phi] = -\frac{1}{2} \text{ (diagram)} - \text{ (diagram)}$$

- Fermions are straightforward though ‘physically’ complicated

- no sign problem numerics as in scalar theories
  - chiral fermions reminder: Ginsparg-Wilson fermions from RG arguments
  - bound states via dynamical hadronisation effective field theory techniques applicable

PI  
↓

$$m \partial_m \Gamma[\phi] = \frac{1}{2} T^{\mu}_{\mu} <\phi(x)\phi(y)> m \partial_m m^2$$



$$\downarrow$$
$$m \partial_m \Gamma[\phi] = \frac{1}{2} \bar{m} \langle \phi(x) \phi(y) \rangle m \partial_m m^2$$

CS-equ.



$$\partial_m \Gamma[\phi] = \frac{1}{2} T_w \langle \phi(x) \phi(y) \rangle \partial_m m^2$$

CS-equ.

170

$$m^2 \rightarrow R_\mu(p^2)$$



## Functional RG

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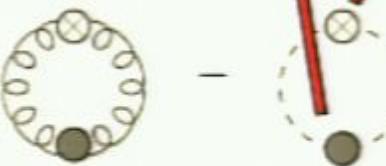
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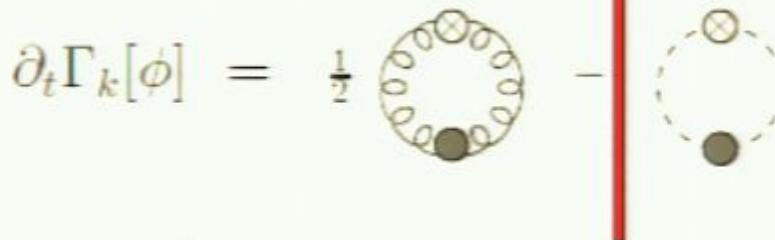
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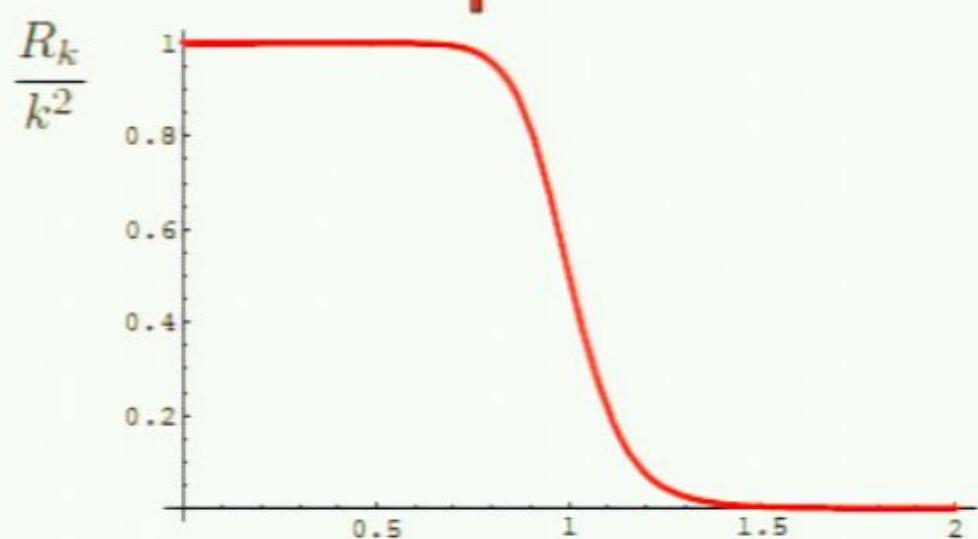
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RG-scale  $k$ :  $t = \ln k$



- Flow infrared finite

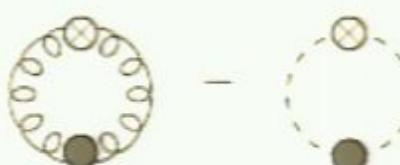


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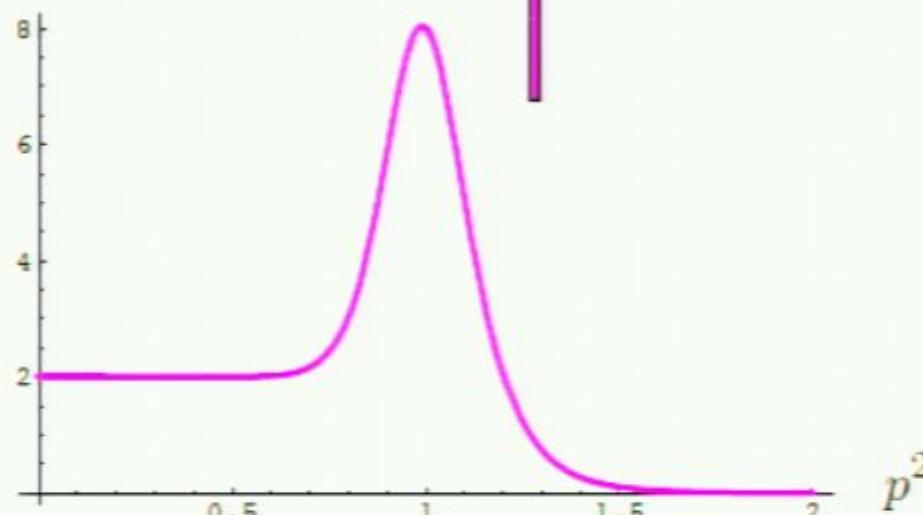
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- Flow **ultraviolet finite**

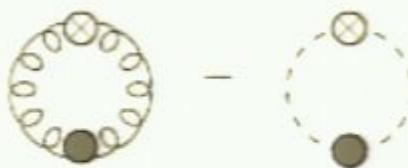
$$\frac{\partial_t R_k}{k^2}$$



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$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left( \text{Diagram A} - \text{Diagram B} \right)$$


- Perturbation theory

$$\partial_t \Gamma_k[\phi] = \partial_t \frac{1}{2} \text{Tr} \log \left( S_{\text{cl}}^{(2)}[\phi] + R_k(p) \right)$$

# Functional RG

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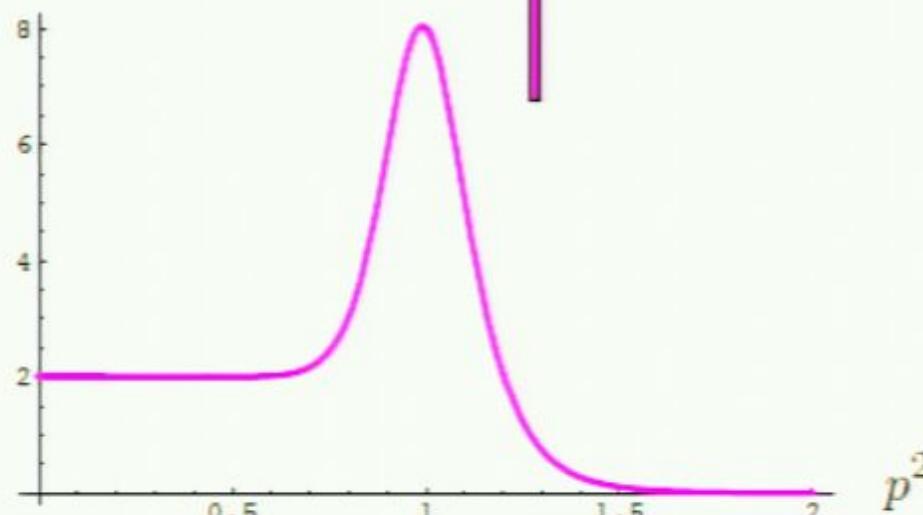
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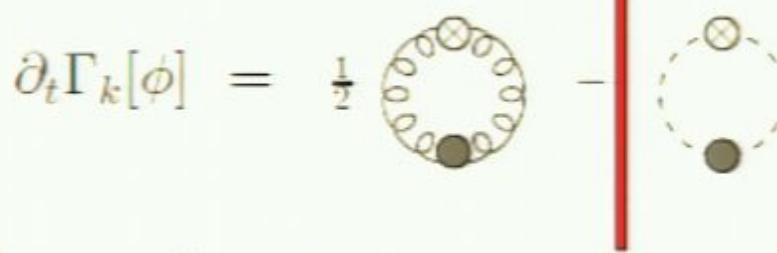


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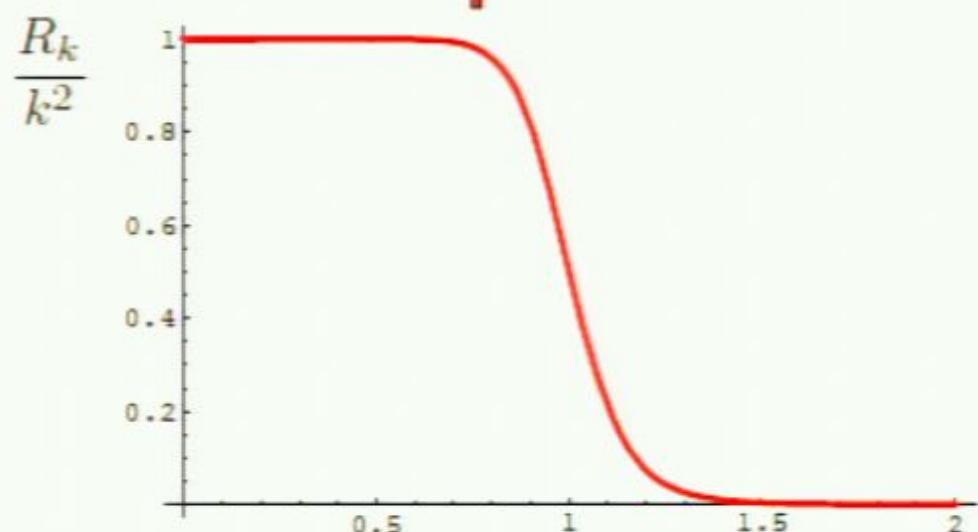
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- Flow infrared finite



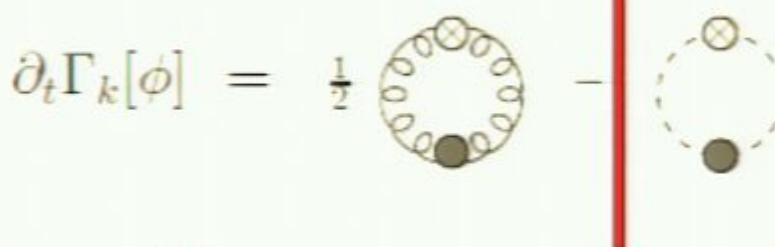


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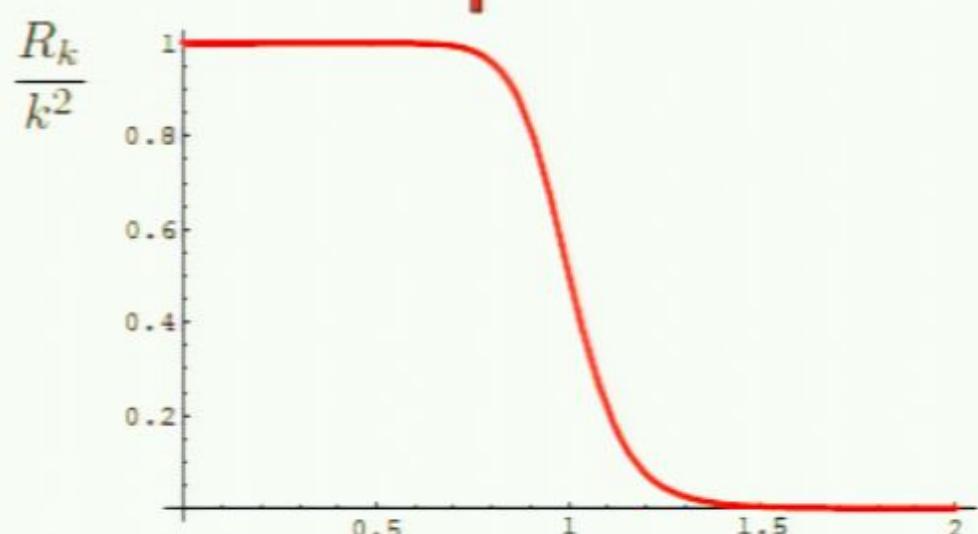
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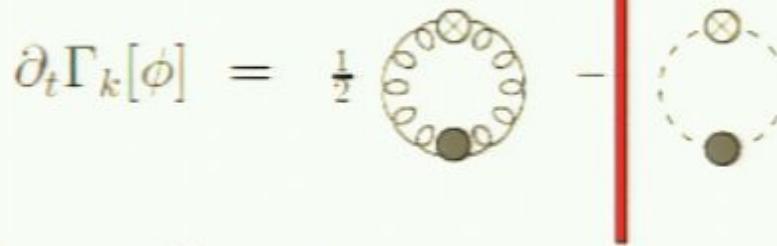


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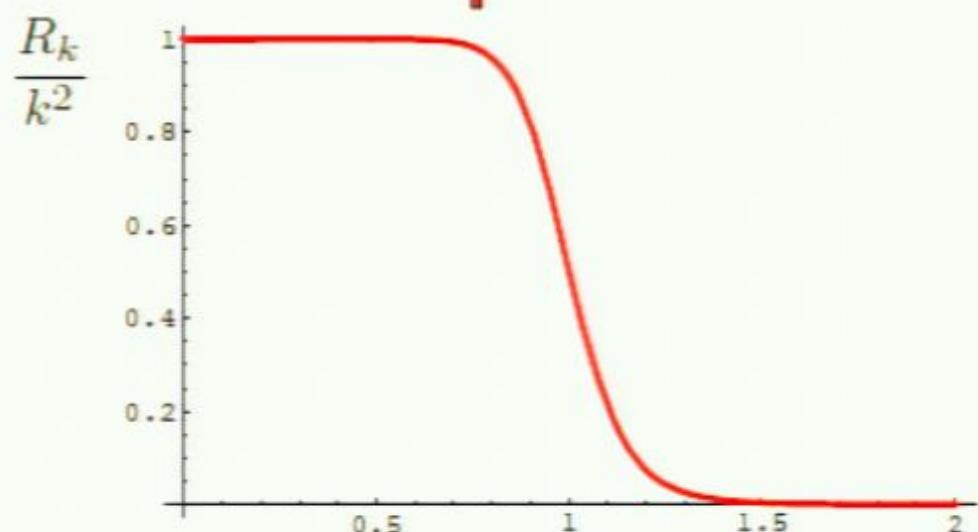
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RG-scale  $k$ :  $t = \ln k$



- Flow infrared finite



$$\frac{R_u(\rho)}{\kappa^2}$$



1

$$\kappa^2 = (10)$$

CS -

m<sup>2</sup> -

# Functional RG

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- Full flow

$$\partial_t \Gamma_k[\phi] = \partial_t \frac{1}{2} \text{Tr} \log \left( \Gamma_k^{(2)}[\phi] + R_k(p) \right) + \text{Diagram C} - \text{Diagram D}$$

Diagram C is circled in red, and a red arrow points from the text "RG-improvement" below to the term  $\partial_t \Gamma_k^{(2)}[\phi]$ .

$$\phi[A_0] = \frac{1}{N_C} \text{Tr} P e^{i \int_0^P A_0(t) dt}$$

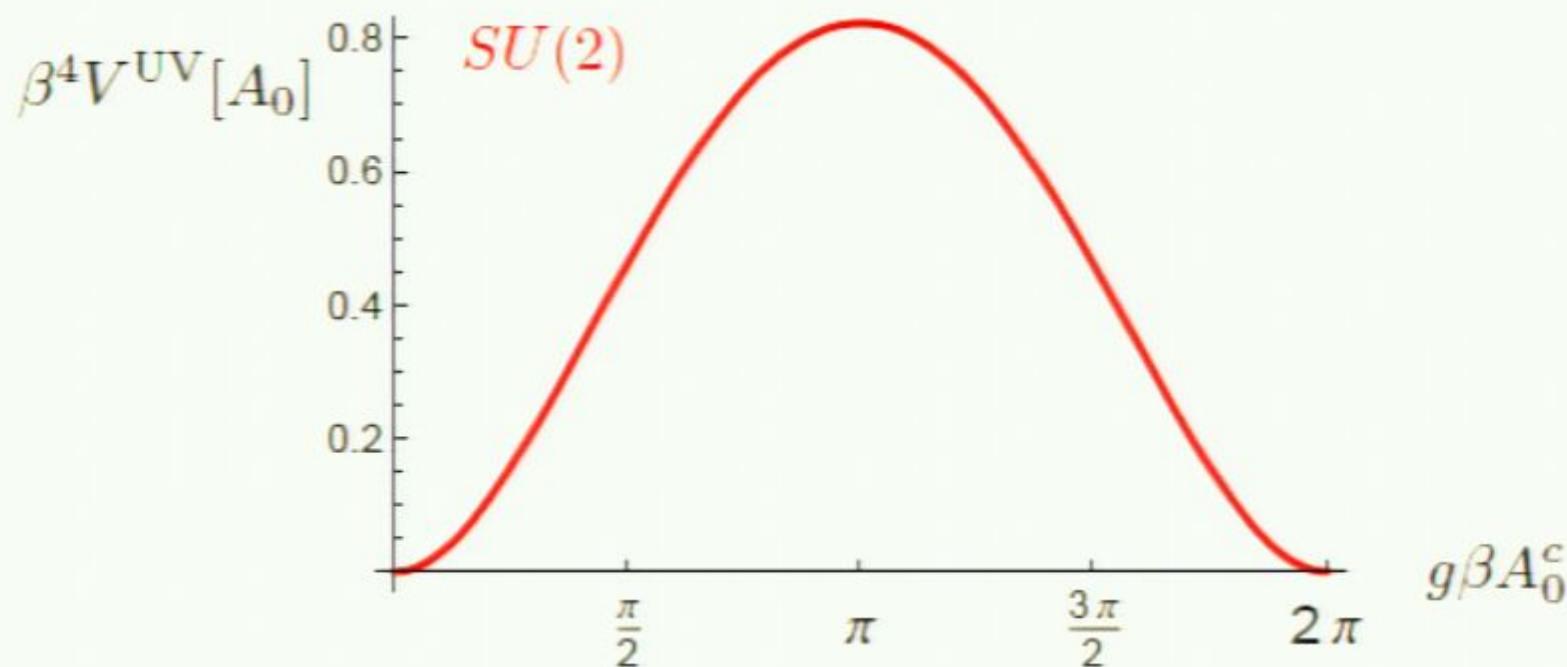
examples of  $\langle WIC \rangle_{CFT}$  calculations in AdS<sub>5</sub> flns

$$\gamma(\sigma), \theta(\tau)$$

# Confinement

Perturbation theory

$$V^{\text{UV}}[A_0] = \frac{1}{2\Omega} \text{Tr} \log S_{AA}^{(2)}[A_0] - \frac{1}{\Omega} \text{Tr} \log S_{CC}^{(2)}[A_0]$$

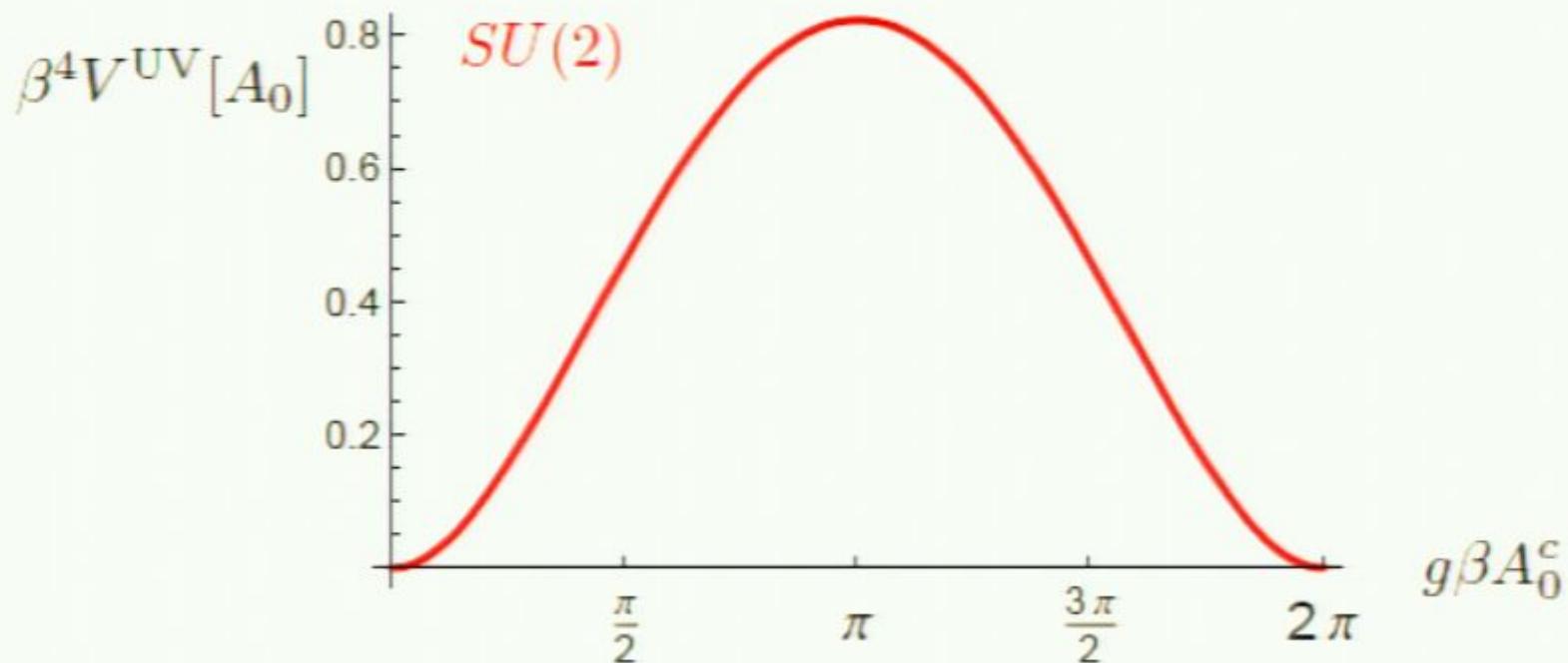




# Confinement

Perturbation theory

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# Confinement

Continuum methods  (Functional RG-flows)

Braun, Gies, JMP '07

$$V[A_0] = -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) - \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle) + O(V''[A_0])$$



- Yang-Mills propagators in Landau gauge ('96 - today)

- DSE, FRG, Lattice, Stochastic Quantisation

von Smekal, Hauck, Alkofer '96

Alkofer, Aguilar, Binosi, Bicudo, Boucaud, Bogolubsky, Bowman, Braun, Cucchieri, De Soto, Dudal, Fischer, Gies, Gracey, Ilgenfritz, Langfeld, Leinweber, Leroy, Litim, Llanes-Estrada, Natale, Mendes, Michel, Müller-Preußker, Oliveira, Papavassilio, JMP, Quandt, Reinhardt, Rodriguez-Quintero, Schwenzer, Skullerud, Sorella, Sternbeck, Verschelde, von Smekal, Williams, Zwanziger, ....

- Numerical solutions

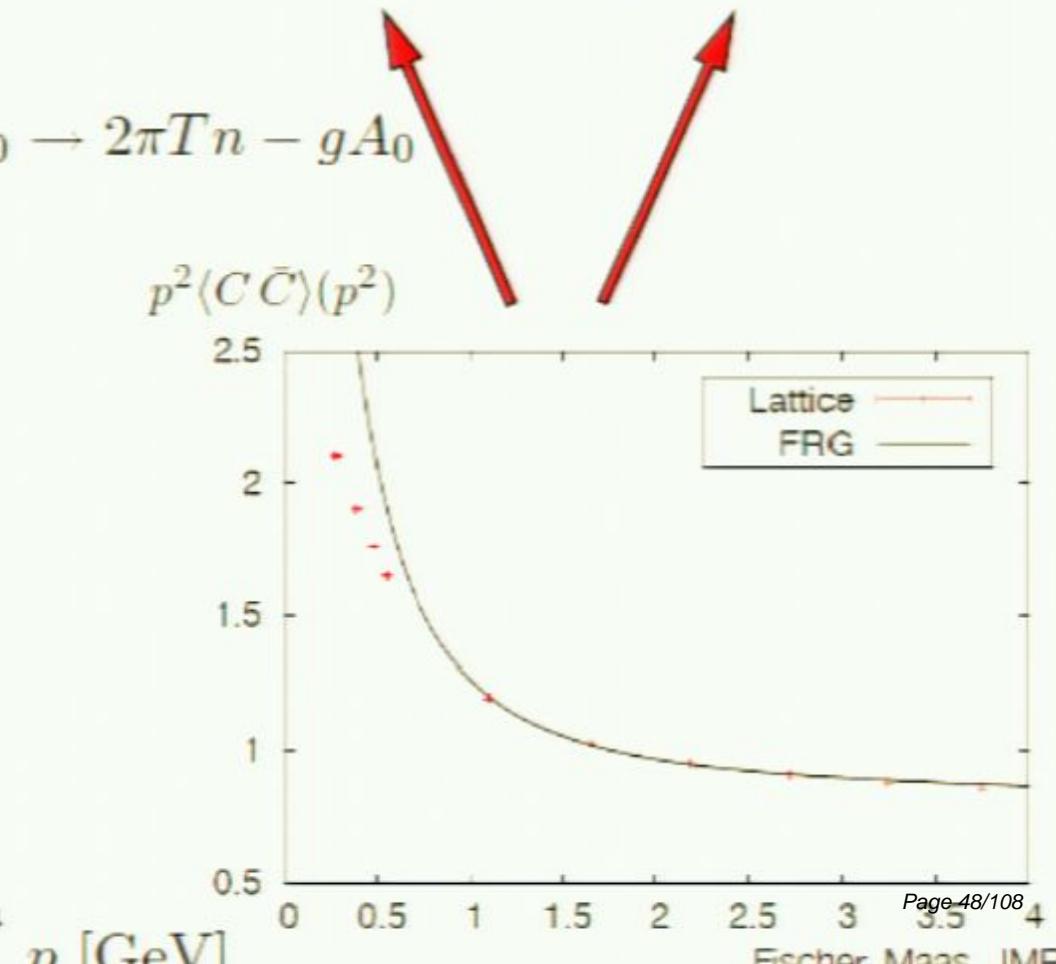
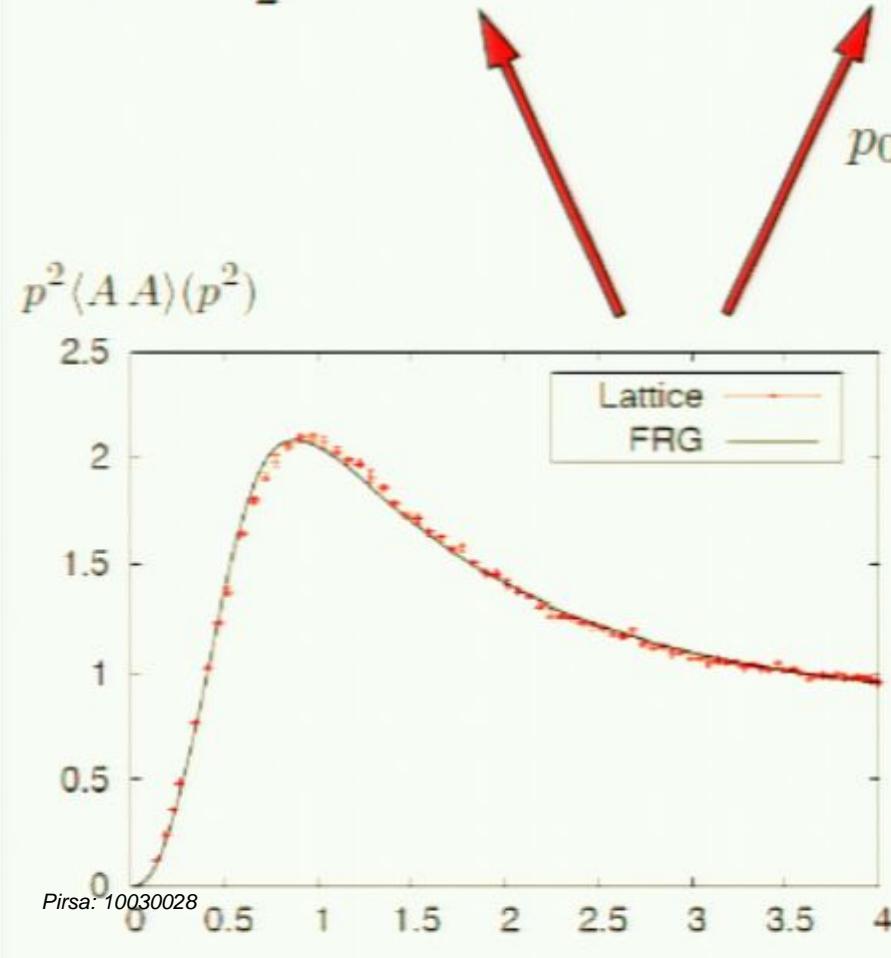
- Analytic IR-asymptotics IR-scaling & Gribov ambiguity

# Confinement

Continuum methods (Functional RG-flows)

Braun, Gies, JMP '07

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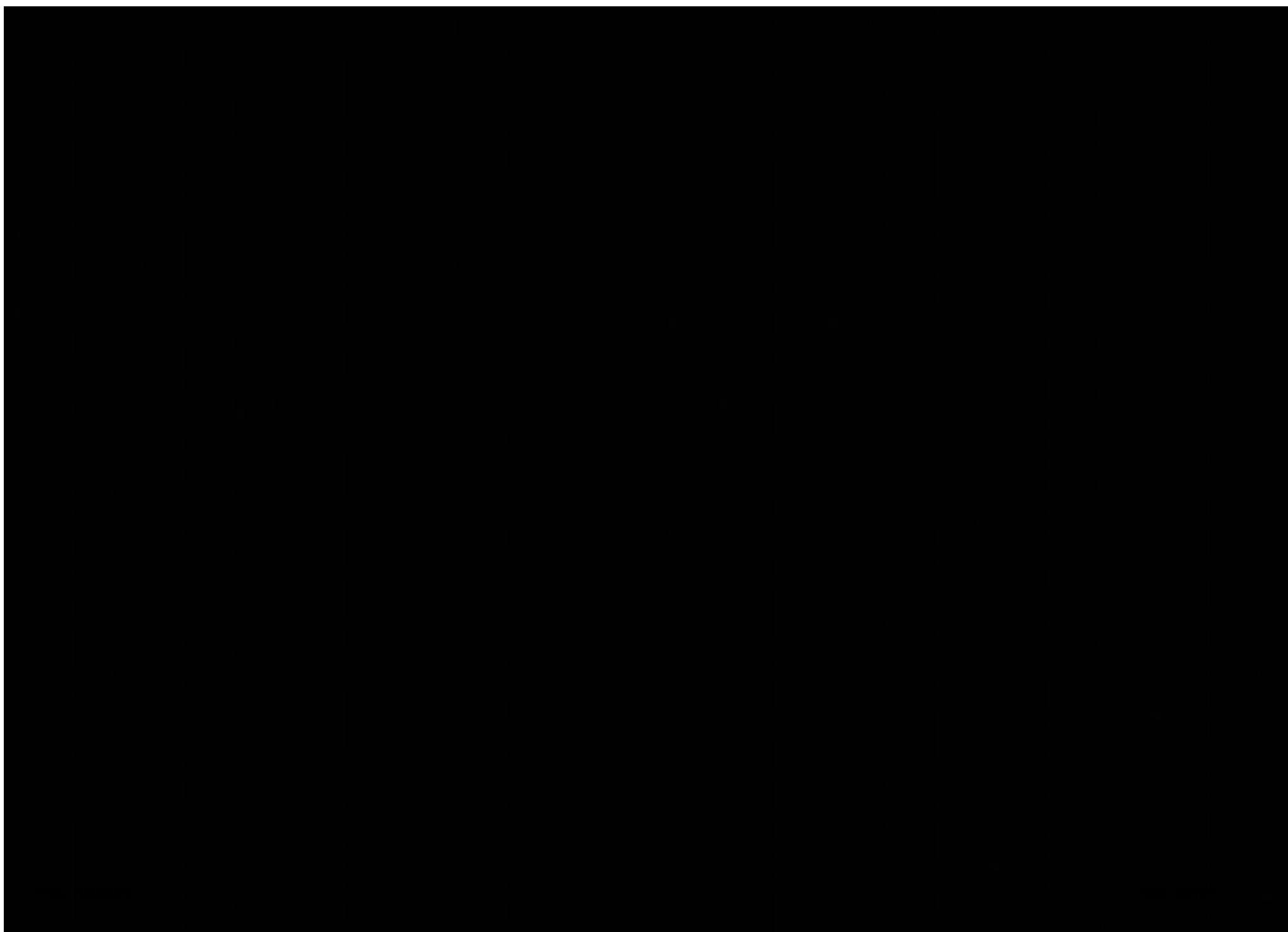
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- Numerical solutions

- Analytic IR-asymptotics IR-scaling & Gribov ambiguity



$$\phi[A_0] = \frac{1}{N_c} + Pe^{i \int_0^{\beta} A_0(t) dt}$$

$$\partial_\mu A_\nu = 0$$

Examples of  $\langle WIC \rangle_{CFT}$  calculations in AdS

$$\tau(\sigma), \theta(\tau)$$

$$ds^2 = \frac{L^2}{z^2} \left( (\alpha^2 - z^2) d\vartheta^2 + \frac{\alpha^2}{\alpha^2 - z^2} dz^2 \right)$$

# Confinement

Continuum methods  (Functional RG-flows)

Braun, Gies, JMP '07

$$V[A_0] = -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) - \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle) + O(V''[A_0])$$



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- Numerical solutions

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$$\subset \partial_\nu A_\nu \subset \phi[A_0] = \frac{1}{N_C} \text{tr } P e^{i \int_0^T A_0(t) dt}$$

$\partial \bar{c} A_\nu \subset \partial_\nu A_\nu = \emptyset$

Examples of  $\langle W(c) \rangle_{CFT}$  calculations in AdS

$$z(\sigma), \theta(\tau)$$

$$ds^2 = \frac{r^2}{a^2} \left( (a^2 - z^2) d\theta^2 + \frac{a^2}{z^2} dz^2 \right)$$

## Confinement

## Computation of propagators

$$k \partial_k \langle \text{---} \bullet \text{---} \rangle^{-1} = - \langle \text{---} \circ \text{---} \rangle + \langle \text{---} \bullet \text{---} \rangle - \langle \text{---} \circ \text{---} \rangle + \frac{1}{2} \langle \text{---} \bullet \text{---} \rangle + \frac{1}{2} \langle \text{---} \circ \text{---} \rangle - \frac{1}{2} \langle \text{---} \bullet \text{---} \rangle + \langle \text{---} \circ \text{---} \rangle$$

$$k \partial_k \left( \dots \right)^{-1} = \dots + \dots + \dots - \frac{1}{2} \dots + \dots$$

The diagram shows a series of Feynman-like diagrams representing a perturbative expansion. The first term is a single vertex labeled  $k \partial_k$ . Subsequent terms involve loops and vertices. One term shows a loop with two vertices, one solid black dot and one dashed circle. Another term shows a loop with three vertices, one solid black dot and two dashed circles. A third term shows a loop with four vertices, all dashed circles. Below these, a term with a dashed circle and a solid black dot is multiplied by  $-\frac{1}{2}$ , followed by another term with a dashed circle and a solid black dot.

# Confinement

Computation of propagators

- full momentum dependence of propagators
- vertices momentum-dependent RG-dressing
- functional optimisation JMP'05
- functional relations between diagrams: Flow=Flow(DSE)

$$\Rightarrow k\partial_k \langle A(p) A(-p) \rangle = \text{Flow}_A[\langle AA \rangle, \langle C\bar{C} \rangle]$$

$$k\partial_k \langle C(p) \bar{C}(-p) \rangle = \text{Flow}_C[\langle AA \rangle, \langle C\bar{C} \rangle]$$

- scaling/decoupling via boundary conditions at  $p^2 = 0$

# Confinement

Computation of propagators

$$k \partial_k \sim \bullet \sim^{-1} = - \sim \circ \circ \sim - \sim \circ \circ \sim + \frac{1}{2} \sim \circ \circ \sim + \frac{1}{2} \sim \circ \circ \sim - \frac{1}{2} \sim \circ \circ \sim + \sim \circ \circ \sim$$

$$k \partial_k \sim \bullet \sim^{-1} = \sim \circ \circ \sim + \sim \circ \circ \sim - \frac{1}{2} \sim \circ \circ \sim + \sim \circ \circ \sim$$

# Chiral symmetry breaking

$$\text{Diagram: Four gluons } g \text{ interacting via a vertex with coupling constant } \alpha_s^2.$$

Order parameter

$$\sigma = \langle \bar{q}q \rangle$$

chiral condensate

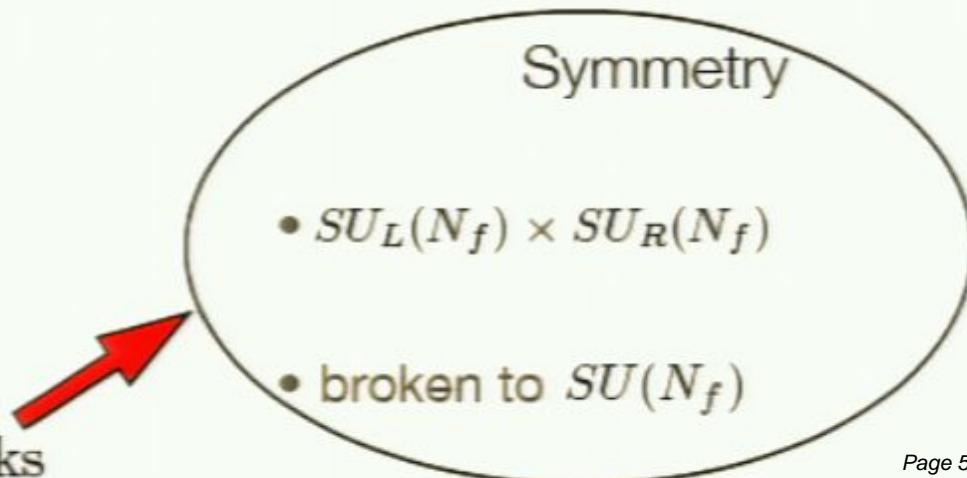
- chiral symmetry:  $\sigma = 0$

- symmetry breaking:  $\sigma \neq 0$

$$\int d^4x \lambda_\psi [(\bar{q}q)^2 - (\bar{q}\gamma_5 q)^2]$$

$$\langle \bar{q}q \rangle \neq 0$$

mass term:  $\langle \bar{q}q \rangle \bar{q}q$



explicitly broken by massive quarks

# Functional RG

---

- Introduction to Functional RG flows & some results in QCD (talks & lit)
  - Integrals from differential equations: The FRG-idea in 0+0-dimensions
  - Confinement & chiral symmetry breaking from Functional Methods
  - Aspects of the Functional RG

# Confinement

Continuum methods  (Functional RG-flows)

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$$V[A_0] = -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) - \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle) + O(V''[A_0])$$



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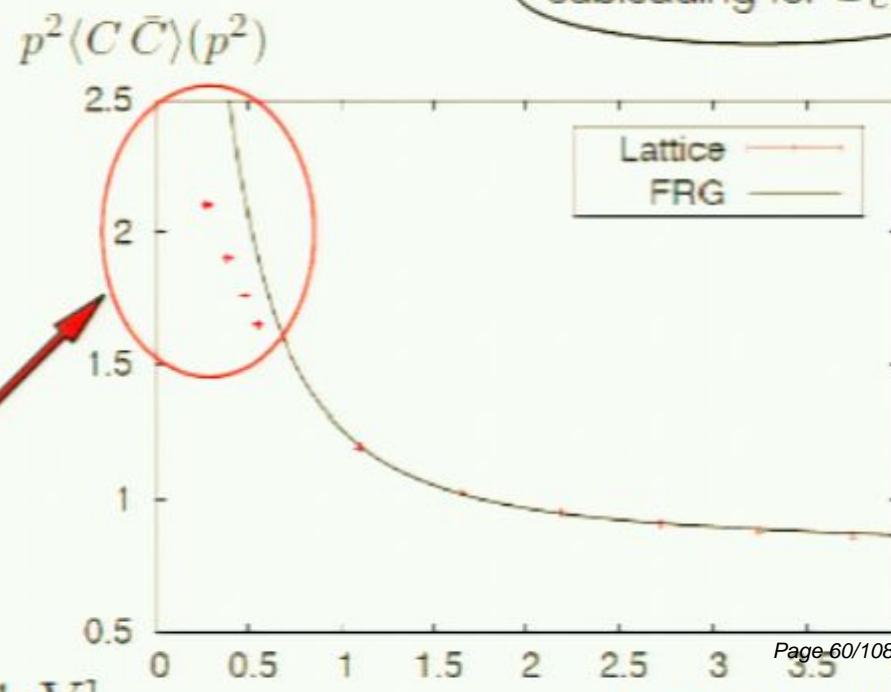
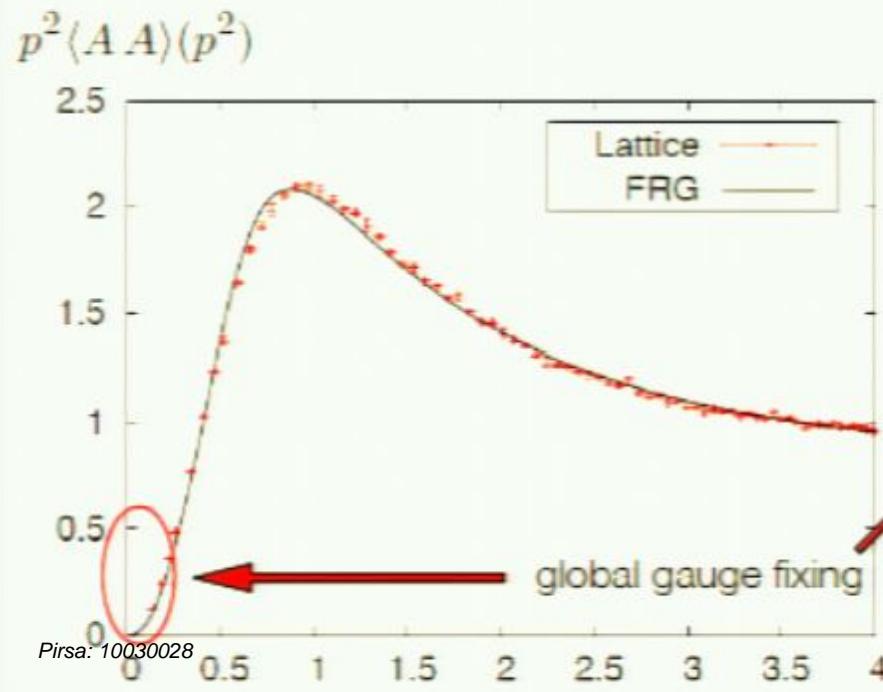
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'Polyakov loop potential'



subleading for  $T_{c,\text{conf}}$

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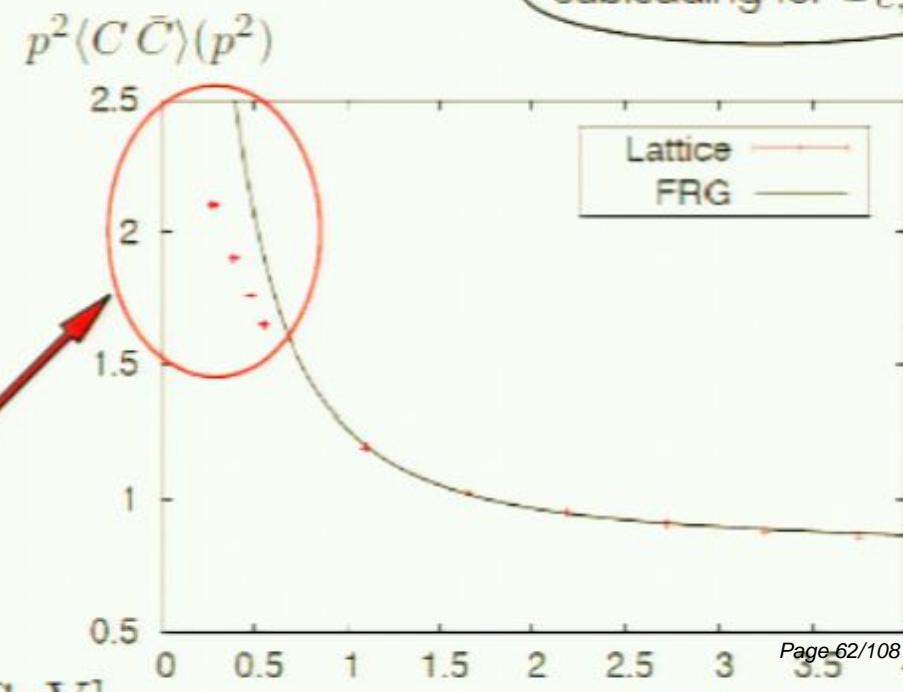
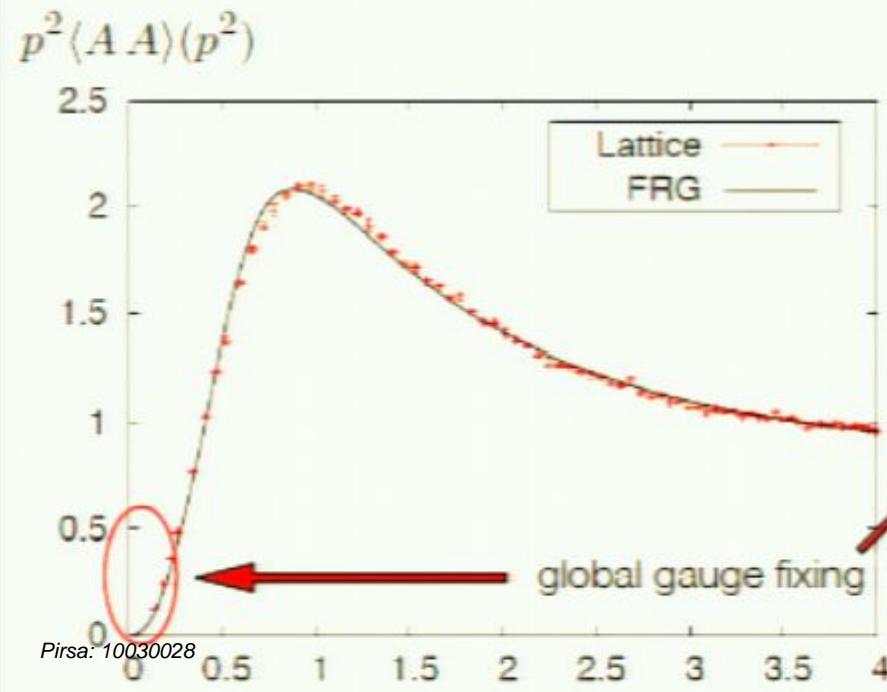
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# Confinement

Computation of propagators

$$k \partial_k \sim \bullet = - \sim \circlearrowleft \circlearrowright - \sim \circlearrowleft \circlearrowright + \frac{1}{2} \sim \circlearrowleft \circlearrowright + \frac{1}{2} \sim \circlearrowleft \circlearrowright - \frac{1}{2} \sim \circlearrowleft \circlearrowright + \sim \circlearrowleft \circlearrowright$$

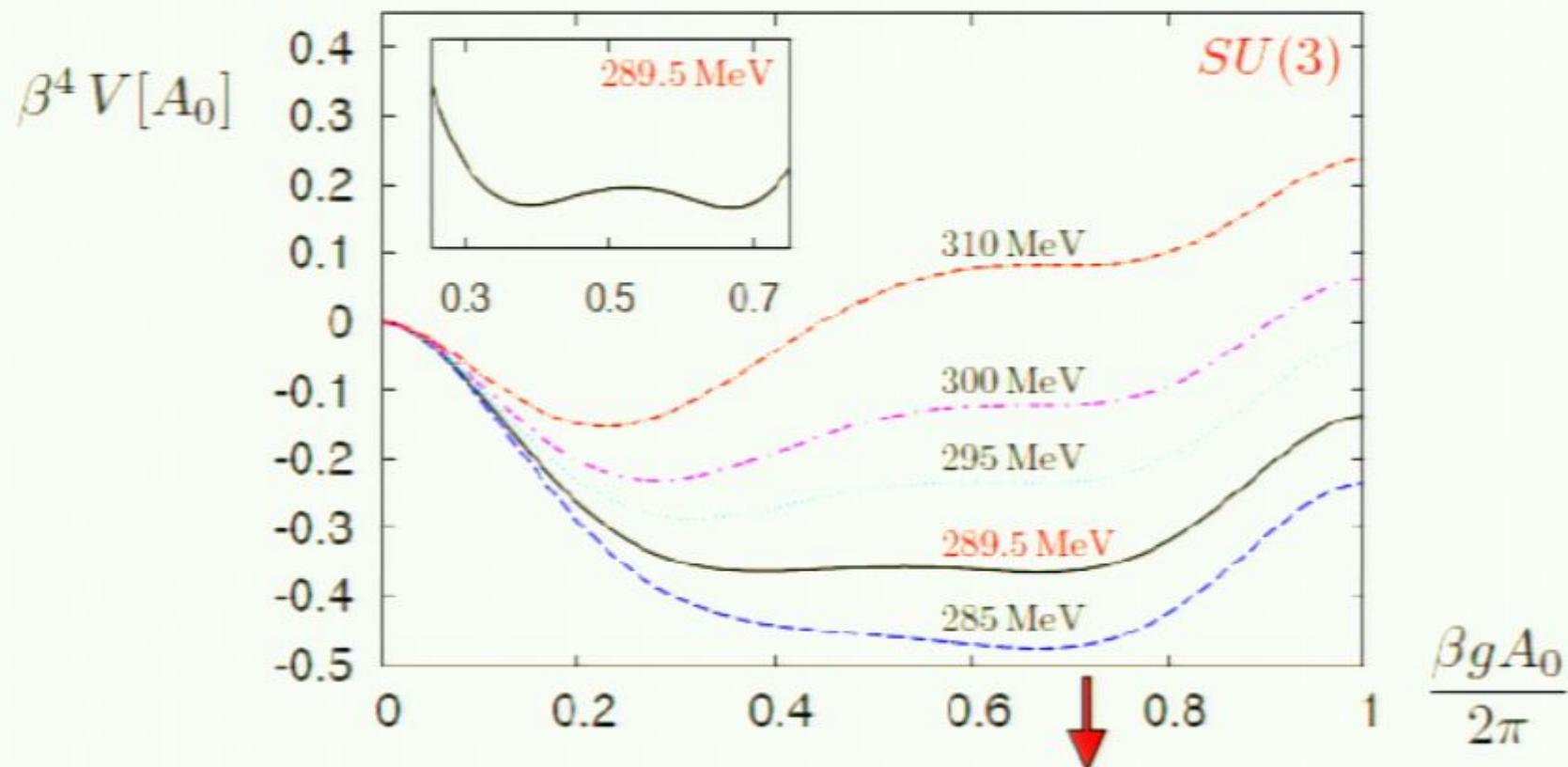
$$k \partial_k \sim \bullet = \sim \circlearrowleft \circlearrowright + \sim \circlearrowleft \circlearrowright - \frac{1}{2} \sim \circlearrowleft \circlearrowright + \sim \circlearrowleft \circlearrowright$$

# Confinement

$$T_c = 289.5 \pm 10 \text{ MeV}$$

$$T_c/\sqrt{\sigma} = 0.658 \pm 0.023$$

$$\text{lattice : } T_c/\sqrt{\sigma} = 0.646$$



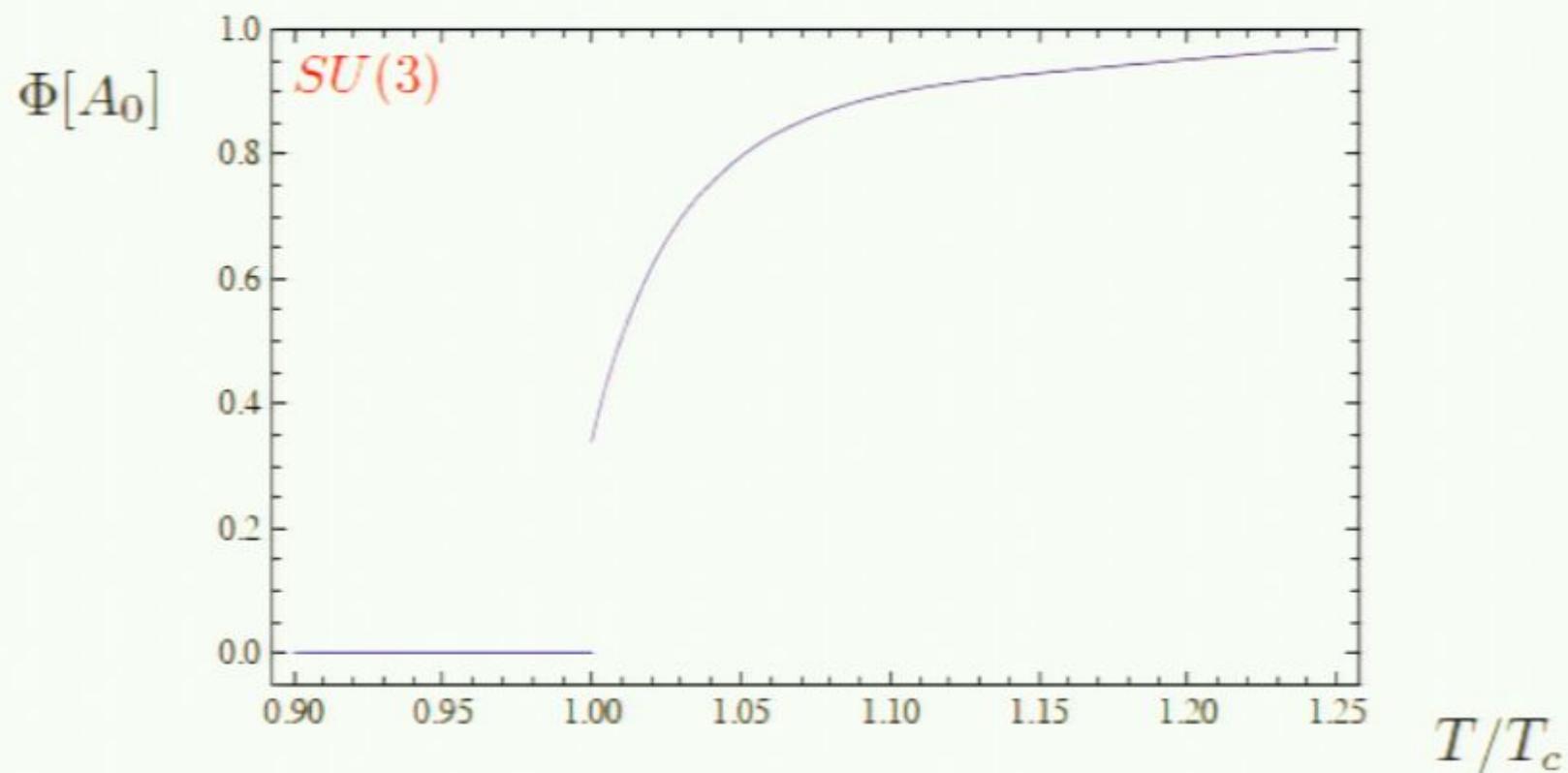
Pirsa: 10030028  $\Phi[A_0] = \frac{1}{3}(1 + 2 \cos \frac{1}{2}\beta g A_0) \longrightarrow \Phi[\frac{4}{3}\pi \frac{1}{\beta g}] = 0$

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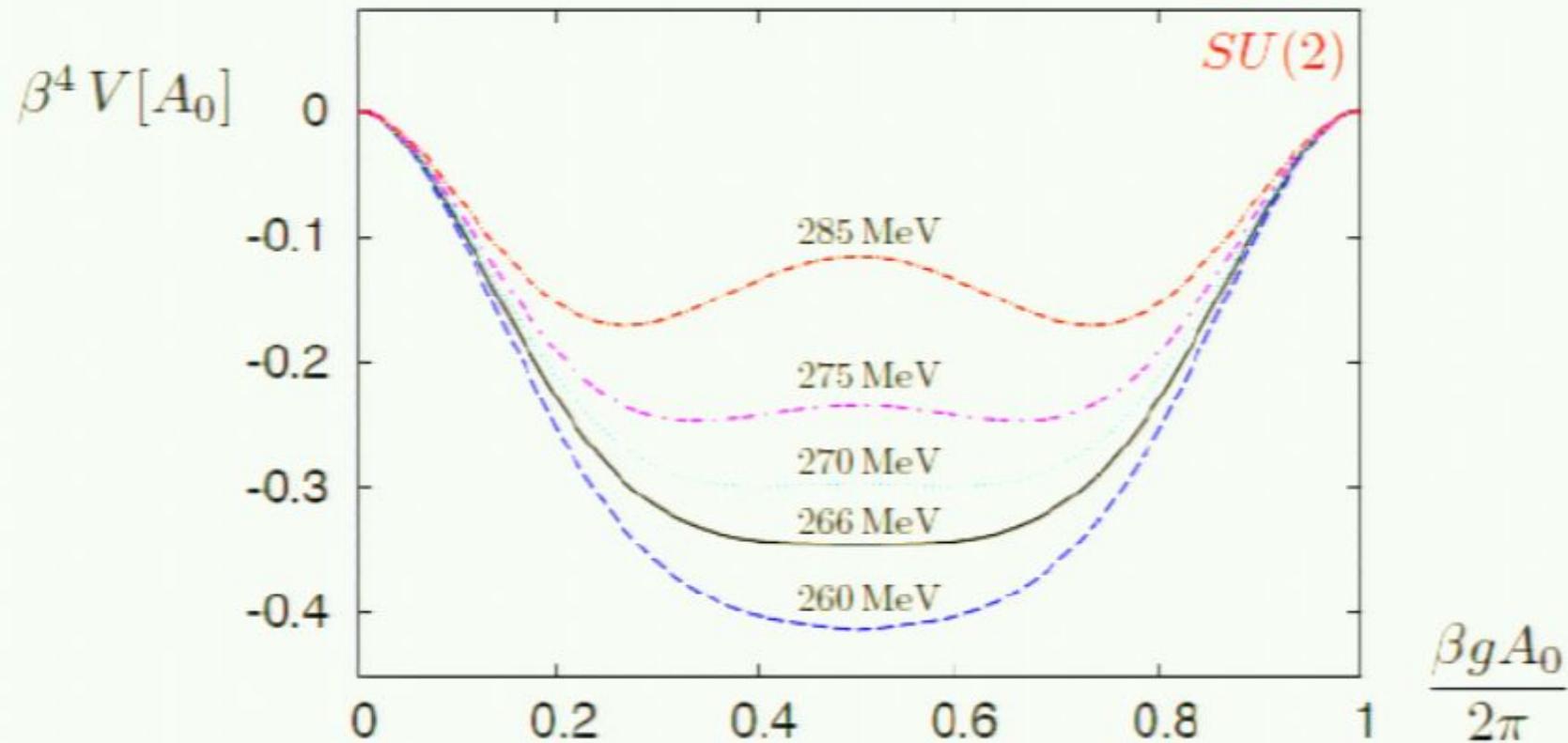
SU(N), G(2), Sp(2): Braun, Eichhorn, Gies, JMP, in preparation

# Confinement

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$$T_c/\sqrt{\sigma} = 0.605 \pm 0.023$$

$$\text{lattice : } T_c/\sqrt{\sigma} = 0.709$$



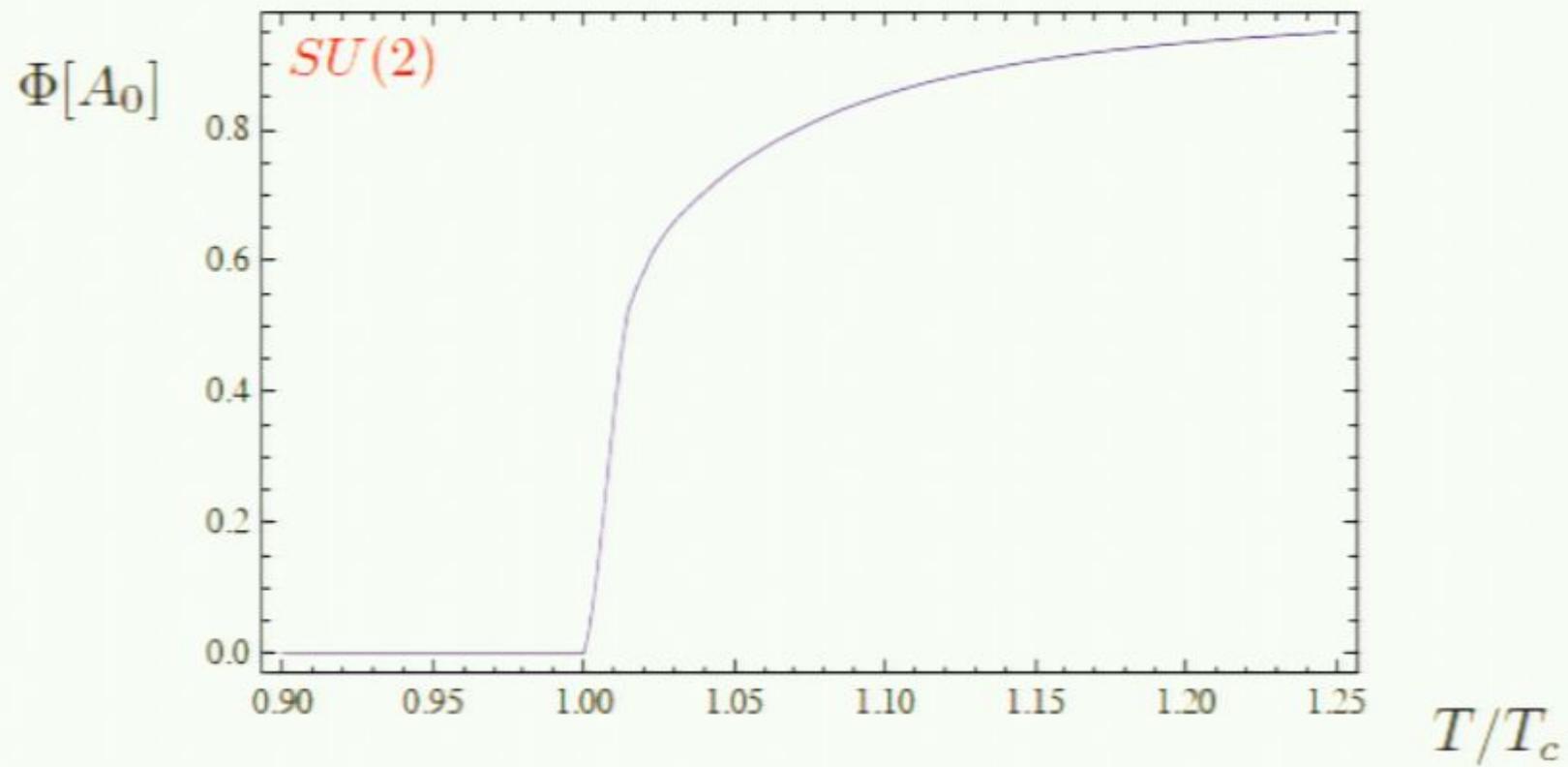
$$\Phi[A_0] = \cos \frac{1}{2} \beta g A_0 \longrightarrow \Phi[\pi/(\beta g)] = 0$$

# Confinement

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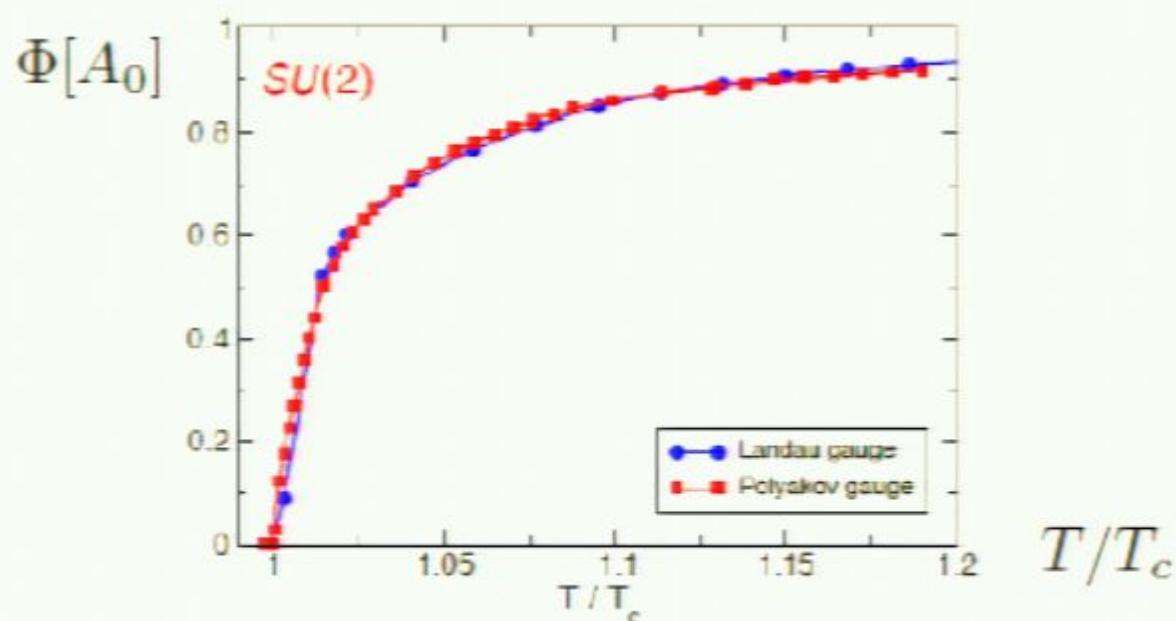
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# Universal properties & gauge independence

Polyakov gauge:  $A_0 = A_0^c(\vec{x})\sigma_3$

$$\text{RG-flow : } V[A_0] = - \int dt \text{flow}[V''[A_0], \alpha_s]$$



- $\text{---} \circ$ : Polyakov gauge: crit. exp.  $\nu = 0.65$

$t_{\text{tiny}} = 0.63$

- $\text{---} \bullet$ : Landau gauge propagators

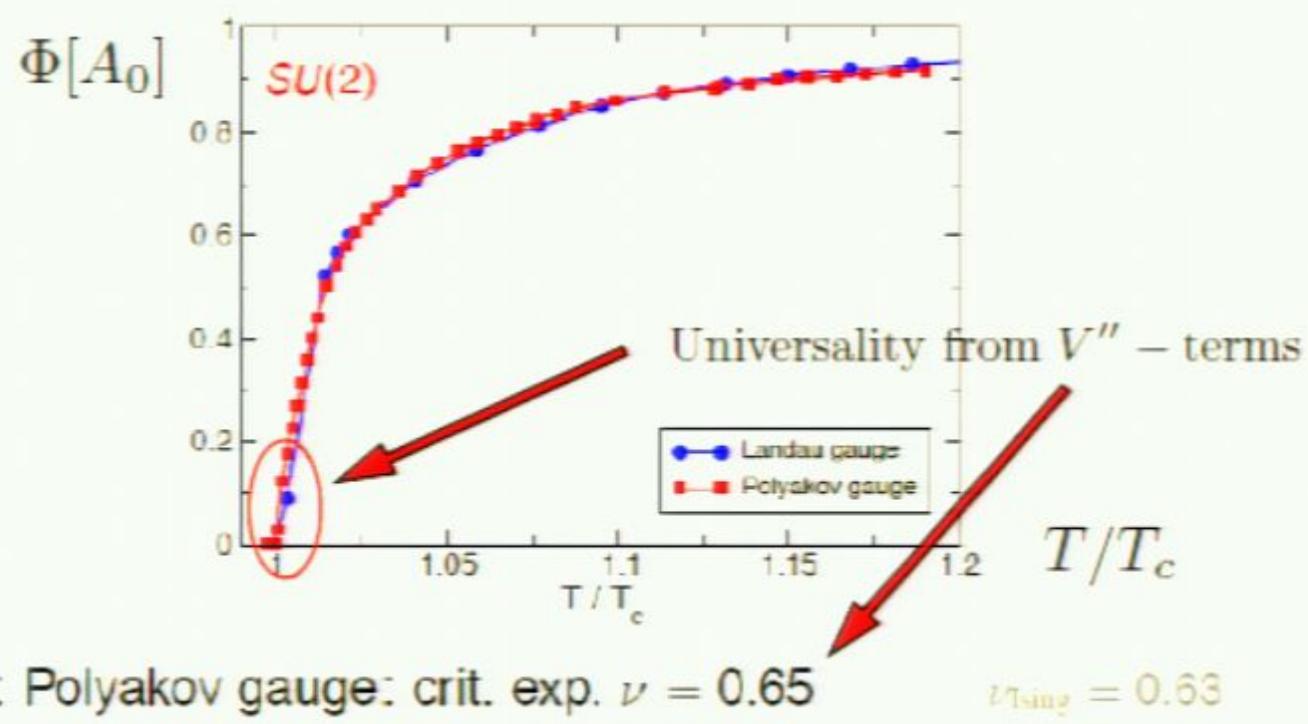
JMP, Marhauser '08  
Page 68/108

# Phase structure at vanishing density

# Universal properties & gauge independence

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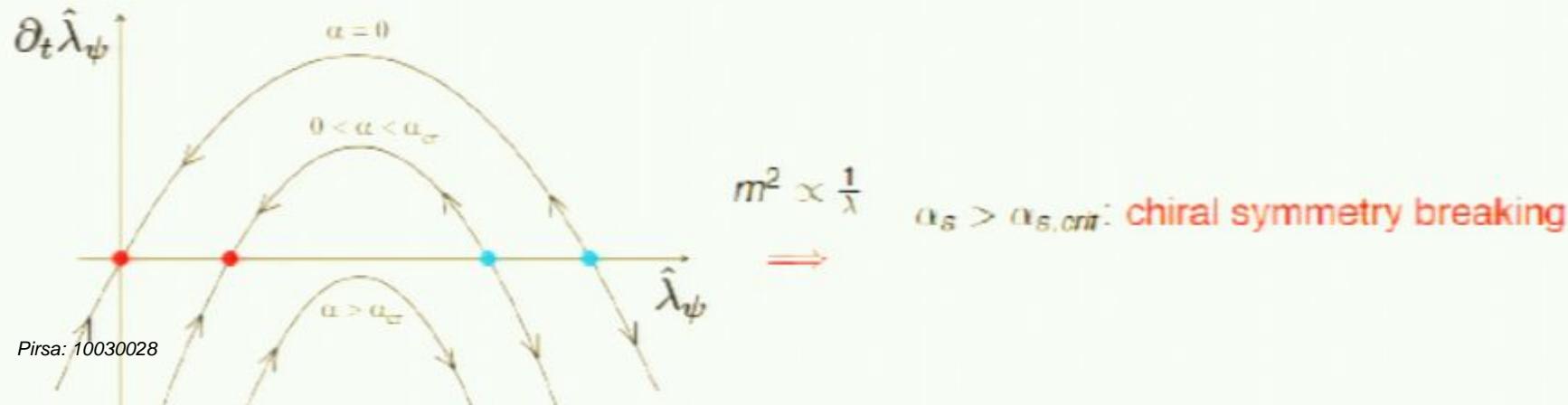
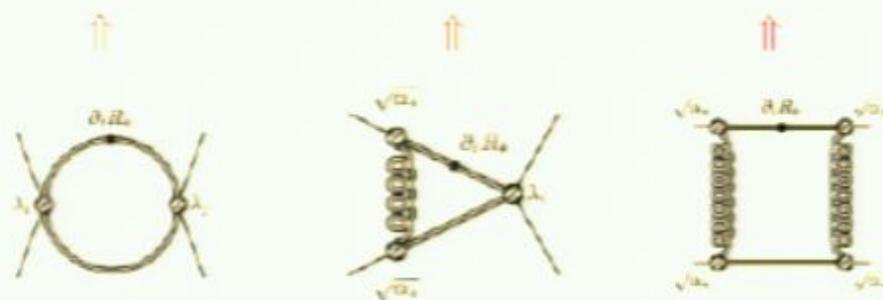
# Phase structure at vanishing density

# Chiral symmetry breaking

## A glimpse at chiral symmetry breaking

Flow of four-fermion coupling  $\hat{\lambda}_\psi = \lambda_\psi k^2$  with infrared scale  $k$

$$k \partial_k \hat{\lambda}_\psi = 2\hat{\lambda}_\psi - A\left(\frac{T}{k}\right) \hat{\lambda}_\psi^2 - B\left(\frac{T}{k}\right) \hat{\lambda}_\psi \alpha_s - C\left(\frac{T}{k}\right) \alpha_s^2$$



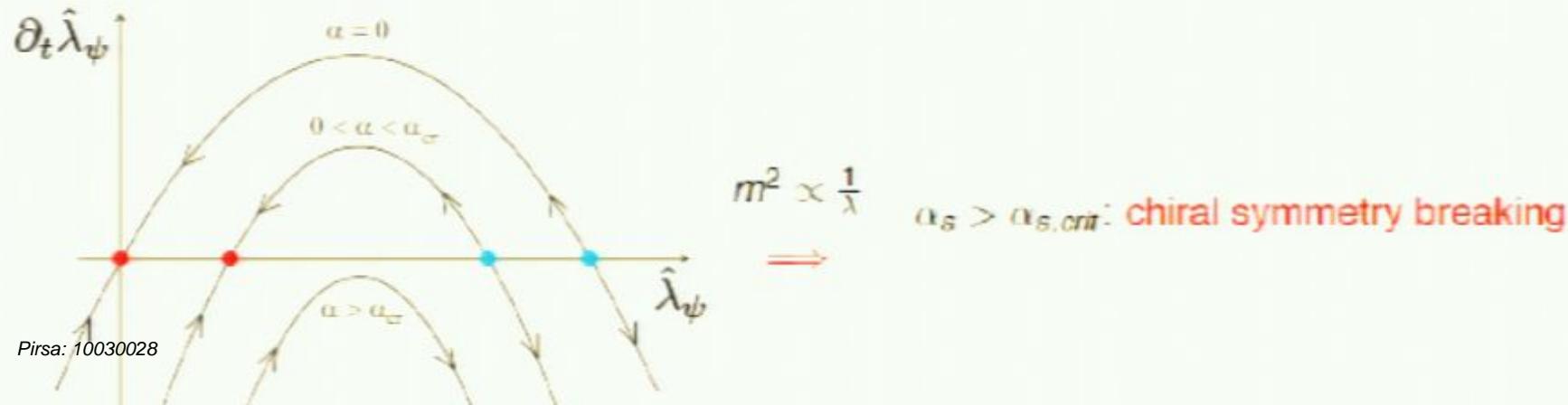
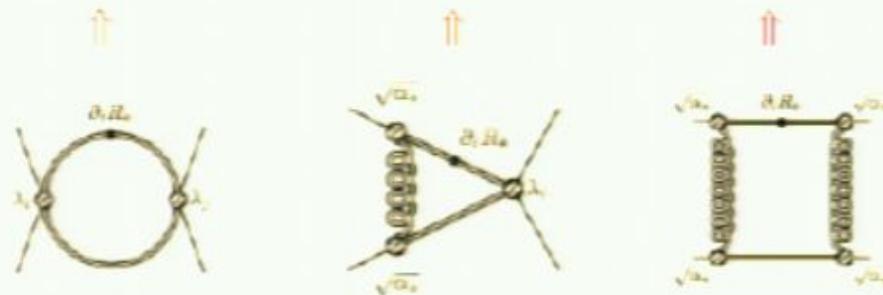


# Chiral symmetry breaking

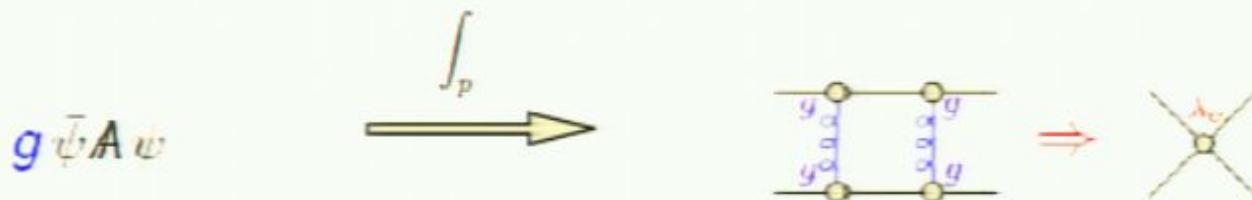
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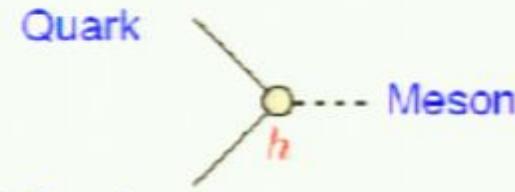
# Chiral symmetry breaking



## Hubbard-Stratonovitch

$$\lambda_\psi (\bar{\psi} \psi)^2 = h \bar{\psi} \psi \sigma - \frac{1}{2} m^2 \sigma^2$$

with  $m^2 = -\frac{h^2}{2\lambda_\psi}$  and EoM( $\sigma$ )



+Baryons and Glueballs

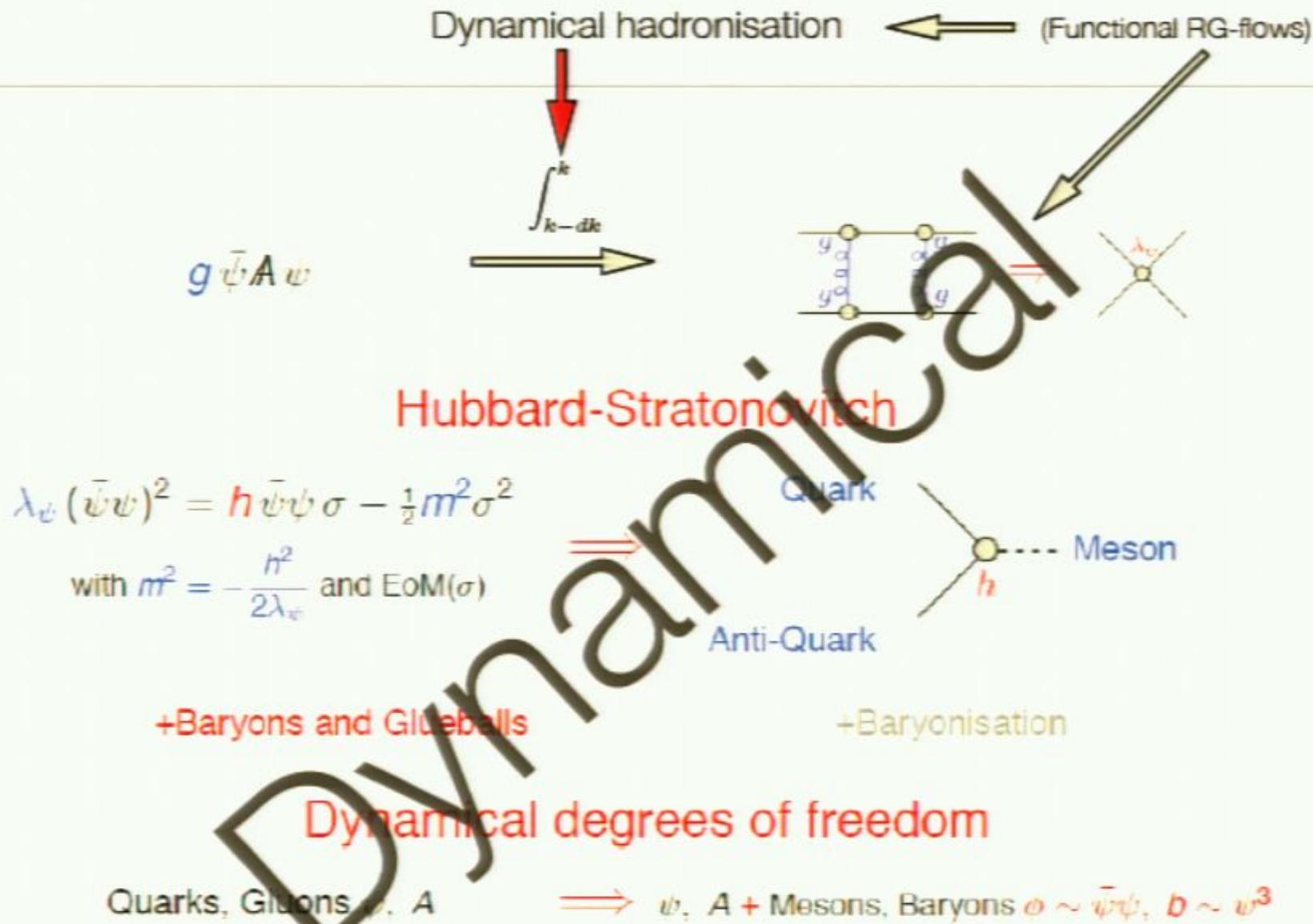
+Baryonisation

## Dynamical degrees of freedom

Quarks, Gluons  $\psi, A$

⇒  $\psi, A +$  Mesons, Baryons  $\phi \sim \bar{\psi} \psi, b \sim \psi^3$

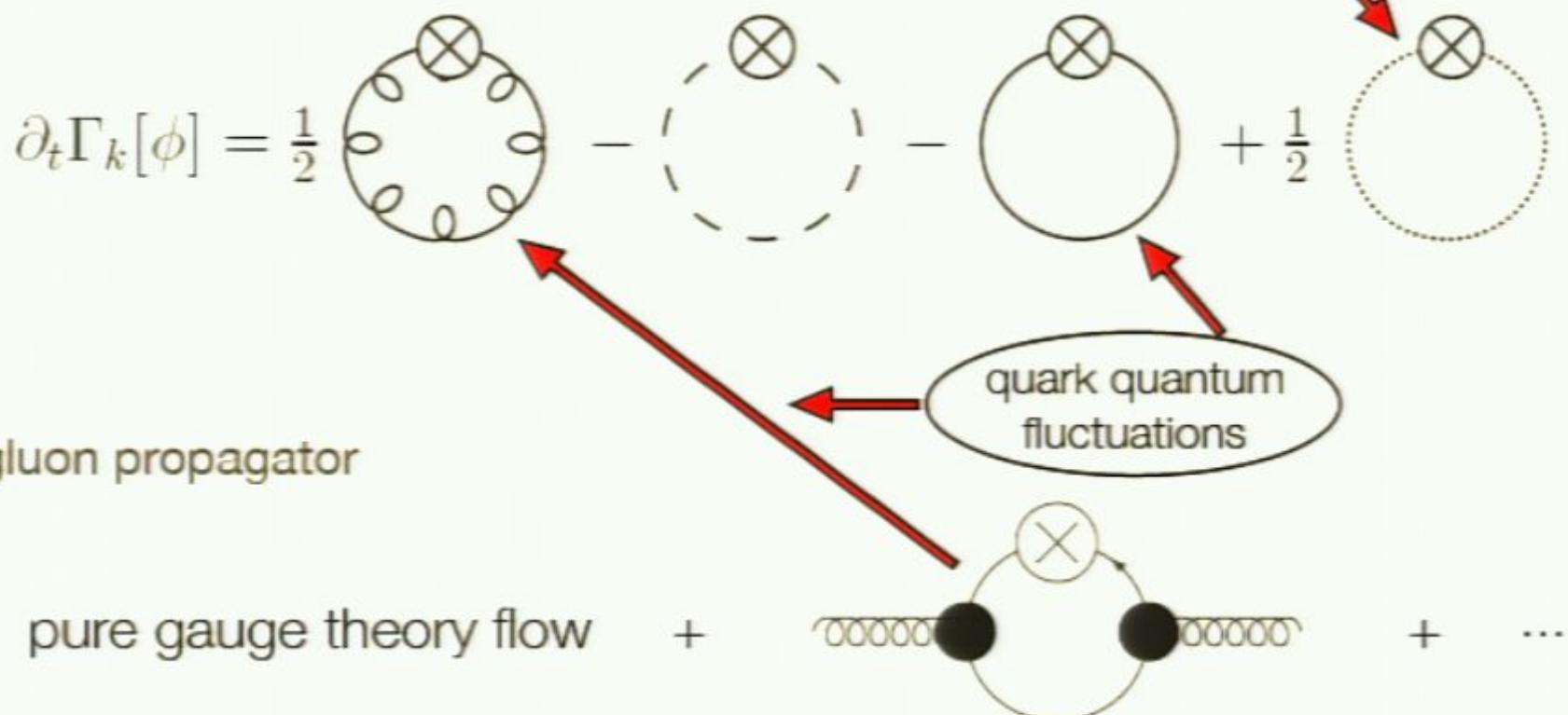
# Chiral symmetry breaking



# Full dynamical QCD: $N_f = 2$ & chiral limit

Continuum methods  $\leftarrow$  (Functional RG-flows)

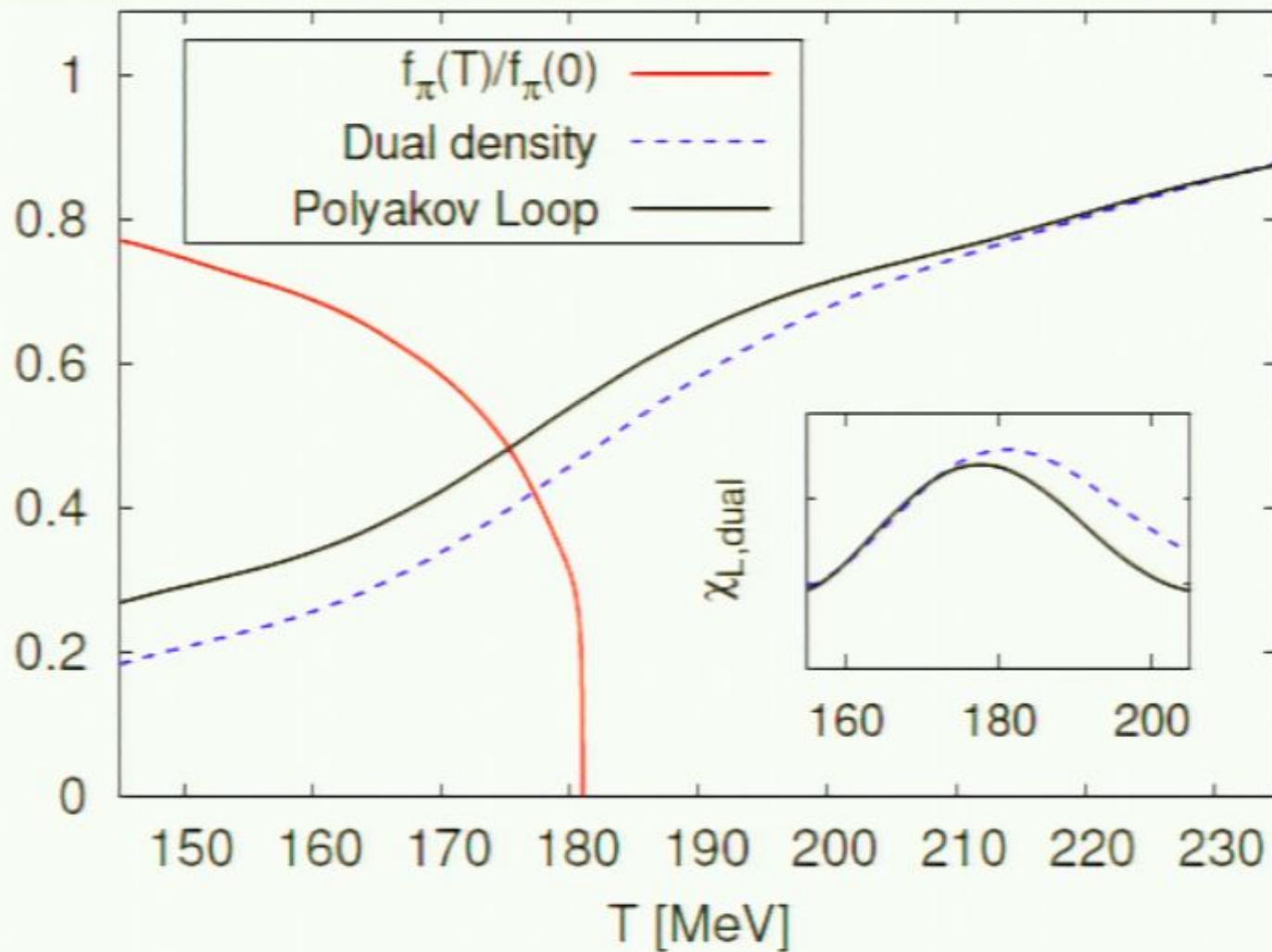
- RG-flow of Effective Action (Effective Potential)



- flow of gluon propagator

# Full dynamical QCD: $N_f = 2$ & chiral limit

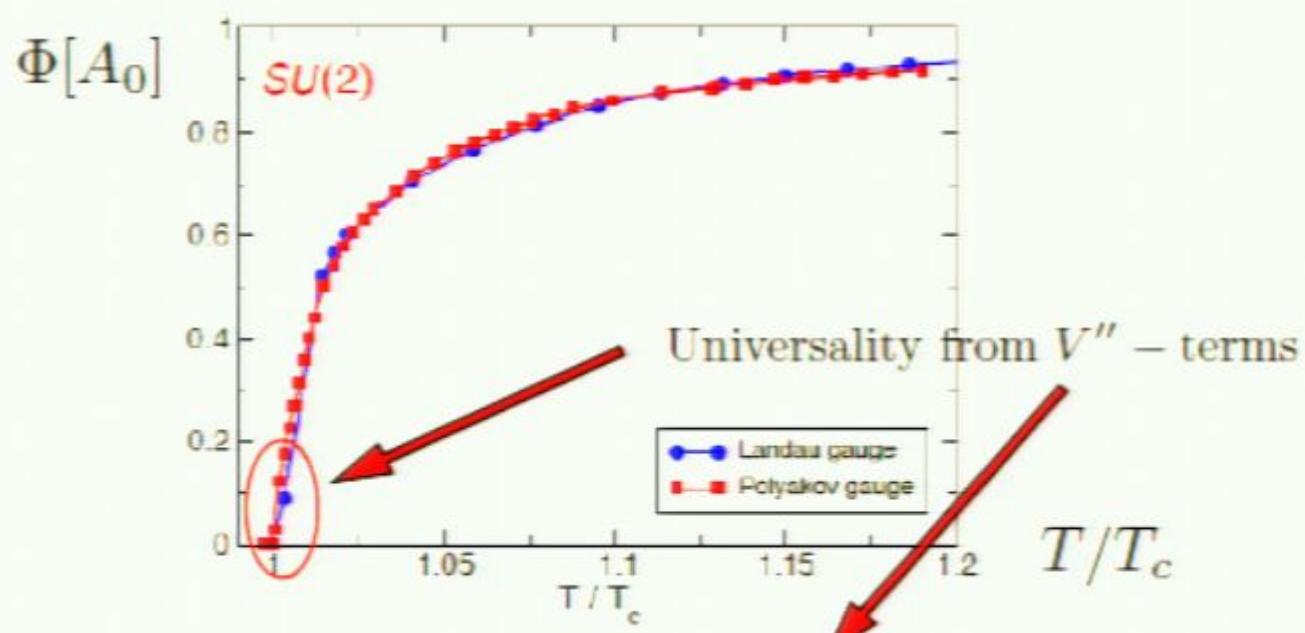
Continuum methods



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- $\text{---}$ : Polyakov gauge: crit. exp.  $\nu = 0.65$

$$\nu_{\text{Tang}} = 0.63$$

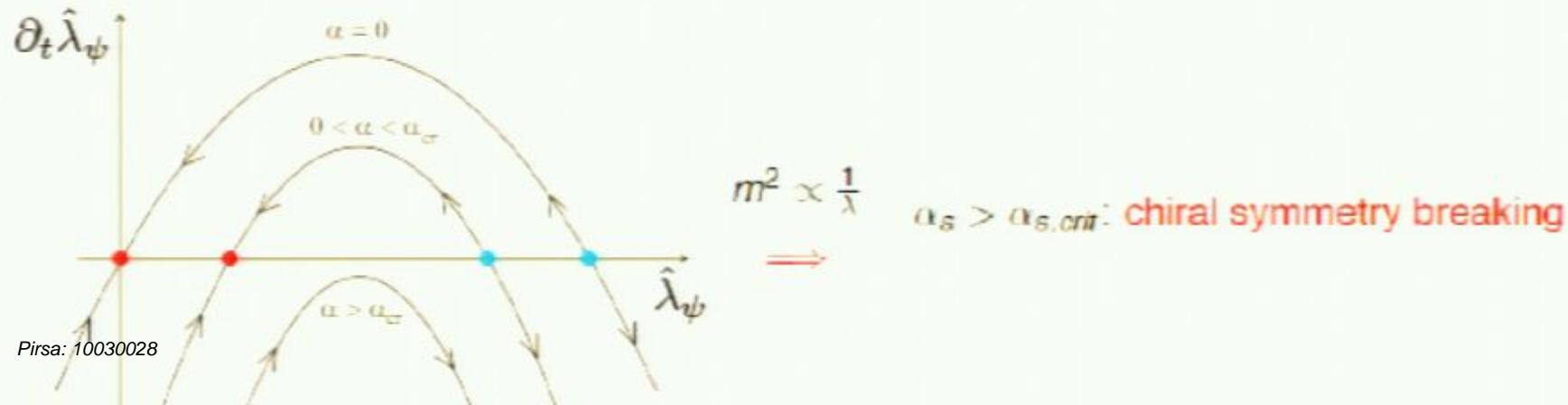
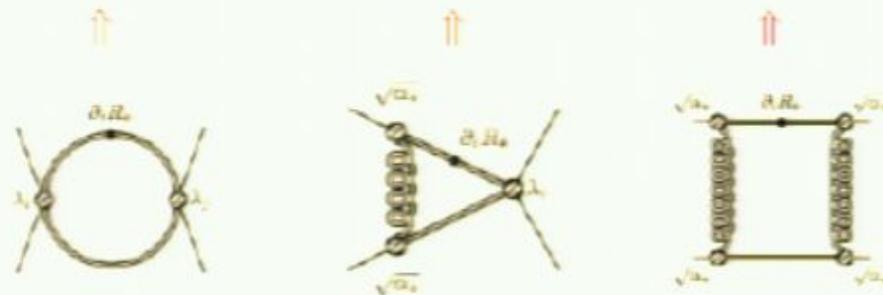
- $\text{---}$ : Landau gauge propagators

# Chiral symmetry breaking

## A glimpse at chiral symmetry breaking

Flow of four-fermion coupling  $\hat{\lambda}_\psi = \lambda_\psi k^2$  with infrared scale  $k$

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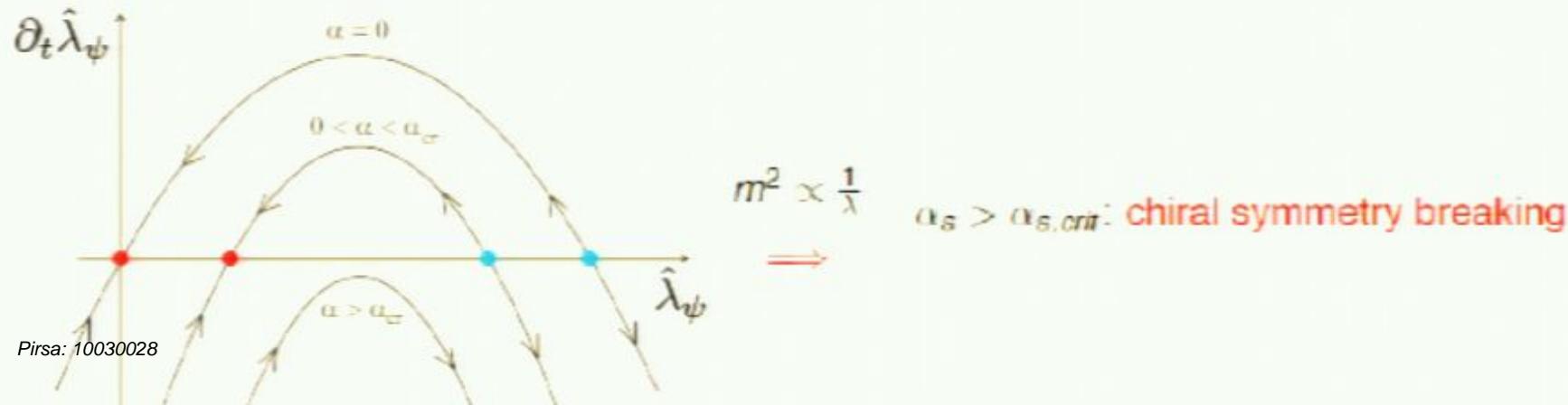
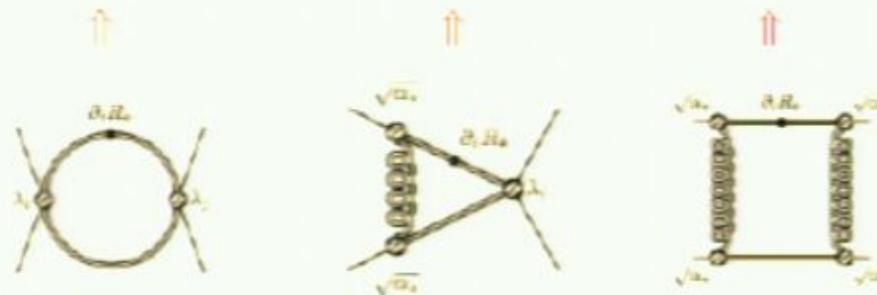


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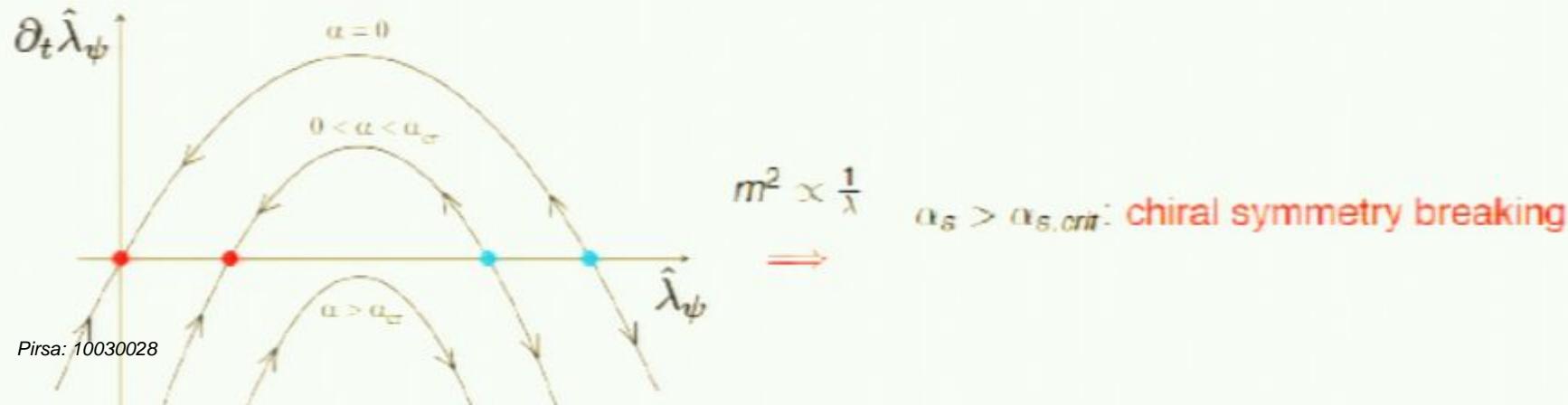
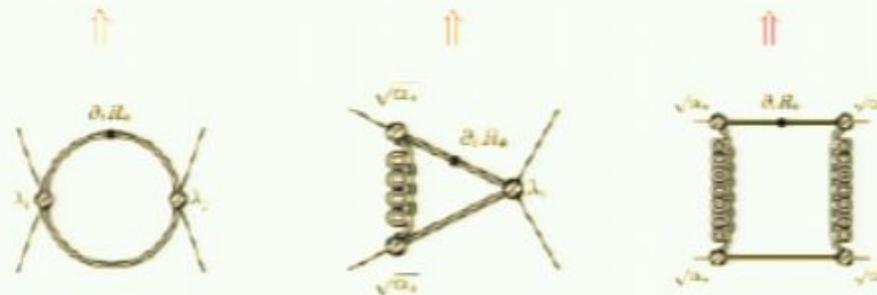


# Chiral symmetry breaking

## A glimpse at chiral symmetry breaking

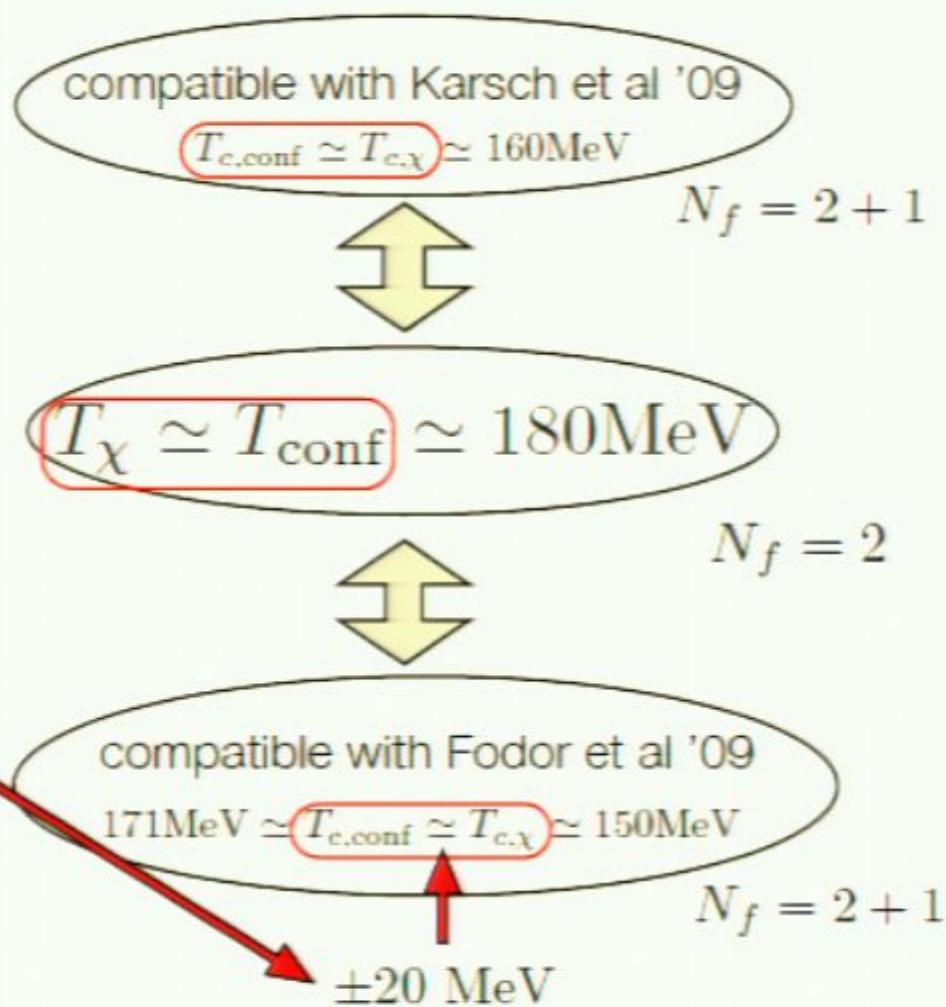
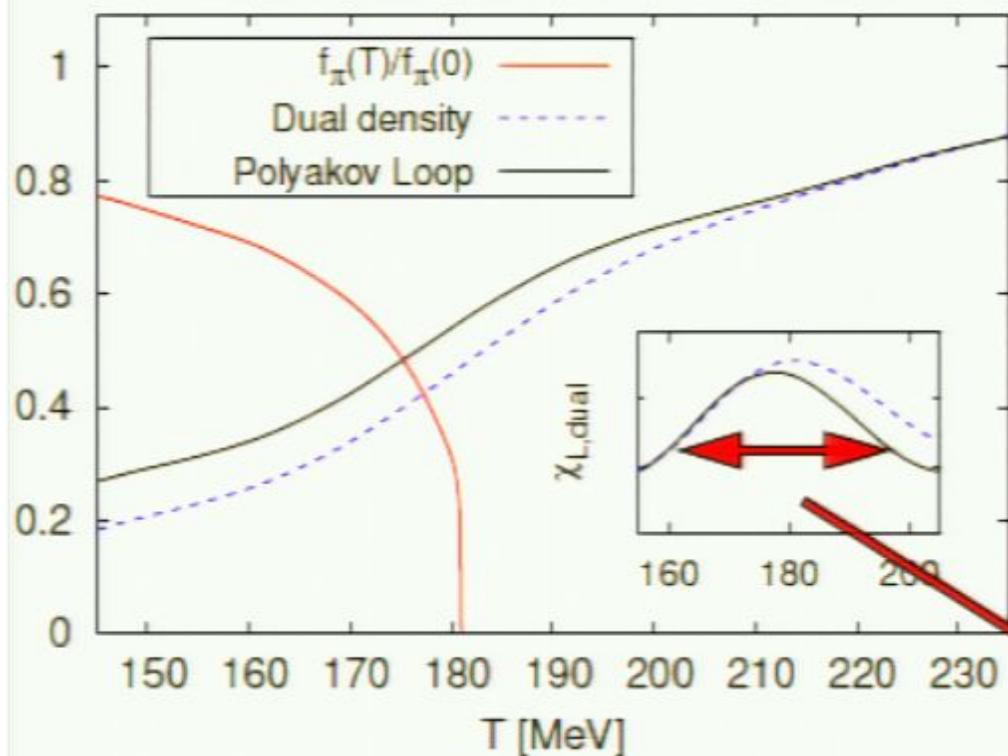
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# Full dynamical QCD: $N_f = 2$ & chiral limit

Continuum methods & lattice

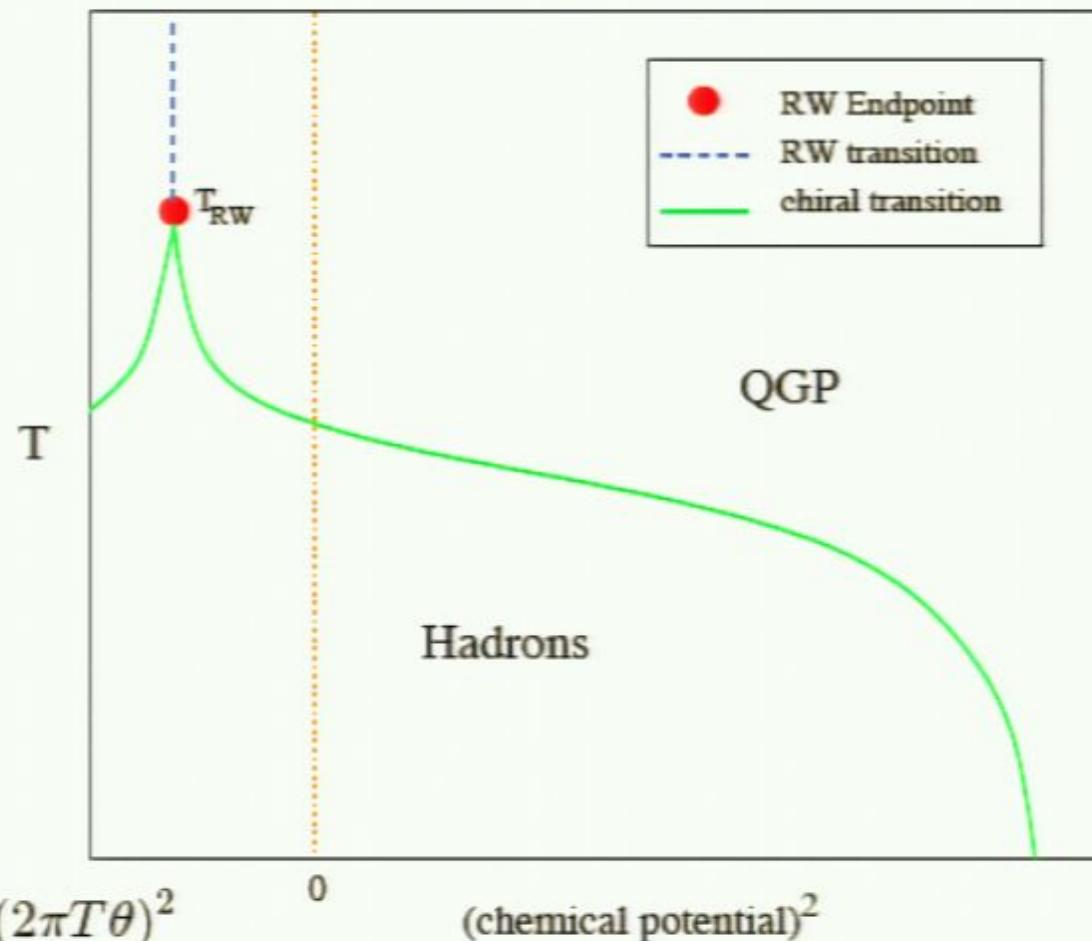


## Phase structure at finite density

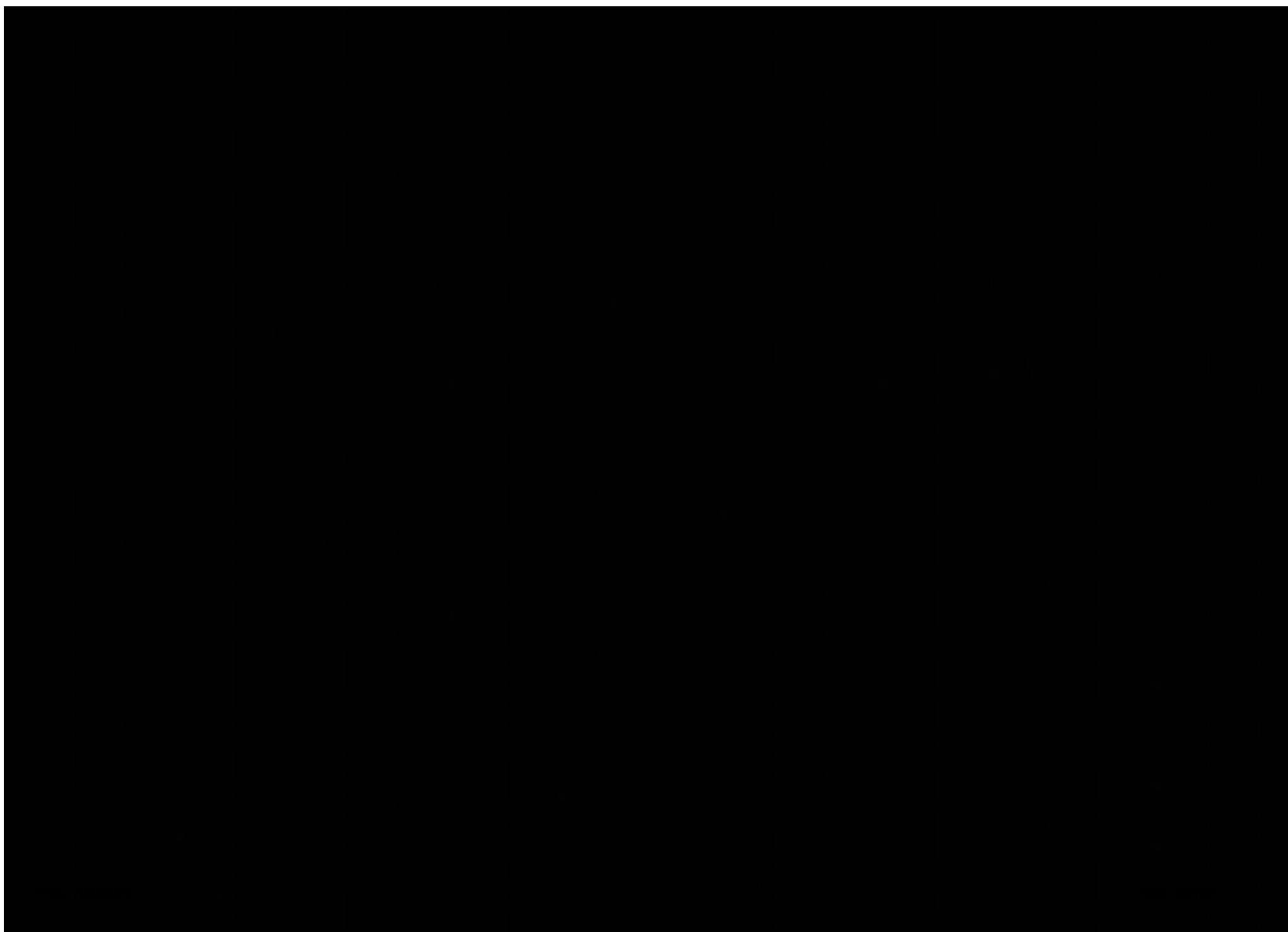
# Imaginary chemical potential

Lattice & Continuum QCD

$$\psi_\theta(t + \beta, \vec{x}) = -e^{2\pi i \theta} \psi_\theta(t, x) \quad \text{with} \quad \mu_I = 2\pi T \theta$$



## Phase structure at finite density



$$\bar{c} \subset A_\mu \subset \phi[A_0] = \frac{1}{N_c} \rightarrow P e^{i \int_0^P A_0(t) dt}$$

$\bar{c} \subset A_\mu \subset \partial A_\mu = 0$

examples of  $\langle N(c) \rangle_{\text{CFT}}$  calculations in AdS

$\varphi(\sigma), \theta(\tau)$

$$\ln \langle N(c) \rangle = \frac{L^2}{\alpha'}$$

$$ds^2 = \frac{L^2}{z^2} \left( (\alpha^2 - z^2) d\sigma^2 + \frac{\alpha}{\alpha^2 - z^2} dz^2 \right)$$

$$S_{NG} = \frac{\text{Area}}{2\pi\alpha'} = \frac{L}{2\pi\alpha'} \int_0^\pi d\sigma \int_0^{\alpha/\sqrt{1-z^2}} dz \frac{\alpha}{z^2} = \frac{L}{\alpha'} \left( \frac{\alpha}{e} - 1 \right)$$

$$\bar{c} \subset A_\mu \subset \phi[A_0] = \frac{1}{N_c} + Pe^{\int_0^\infty A_0(t) dt}$$

$\bar{c} \subset A_\mu \subset \partial_\mu A_\mu = 0$

Examples of  $\langle N(c) \rangle_{CFT}$  calculations in AdS

$\tau(0), \theta(\tau)$

$$\ln \langle N(c) \rangle = \frac{L^2}{\alpha'}$$

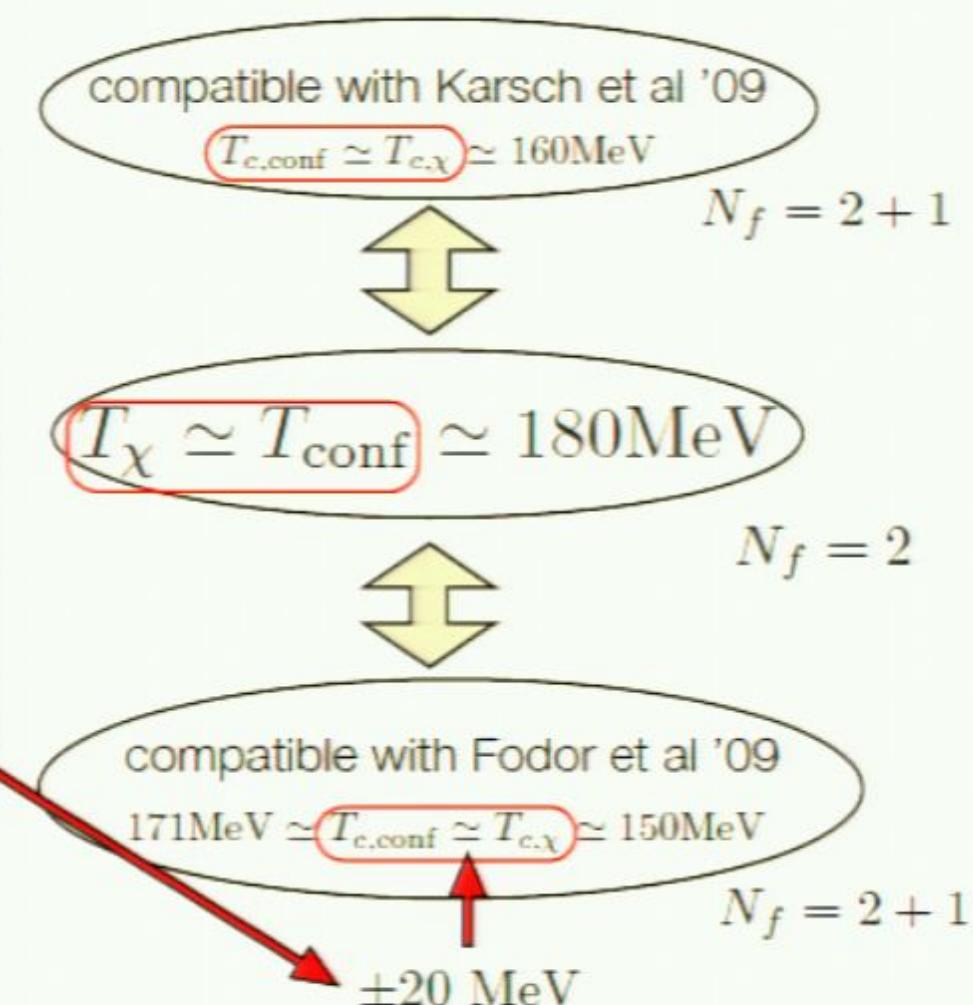
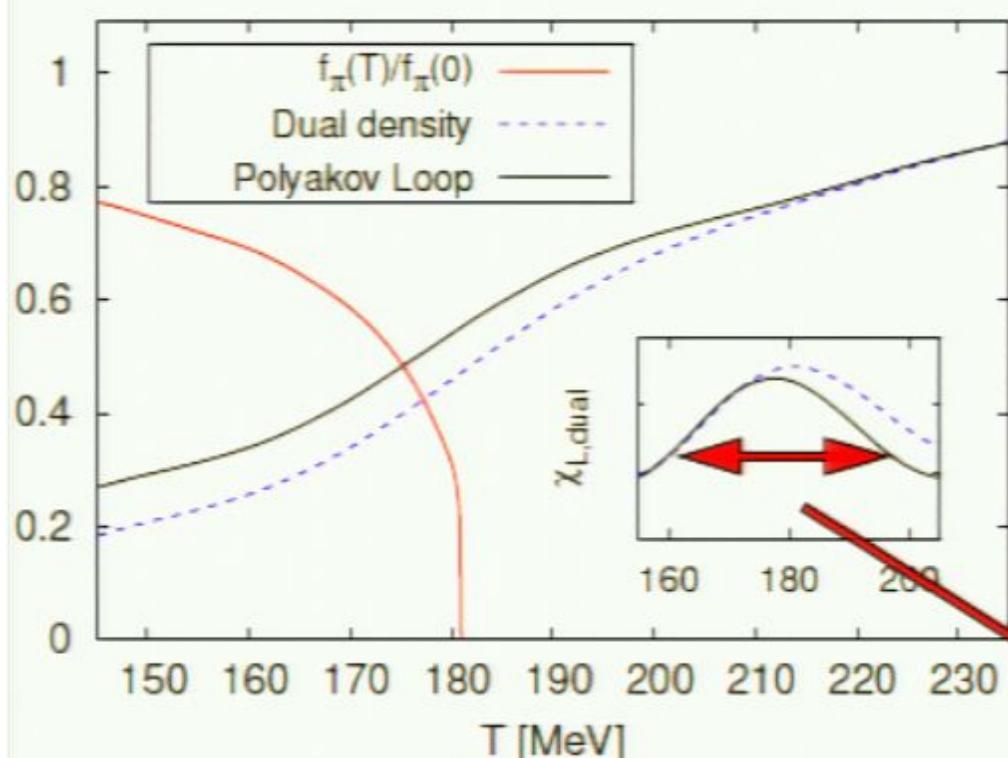
$$ds^2 = \frac{L^2}{z^2} \left( (a^2 - z^2) d\theta^2 + \frac{\alpha'}{a^2 - z^2} dz^2 \right)$$



$$S_{NG} = \frac{\text{Area}}{2\pi\alpha'} = \frac{L}{2\pi\alpha'} \int d\theta \int dz \frac{a}{z^2} = \frac{L}{\alpha'} \left( \frac{a}{e} - 1 \right)$$

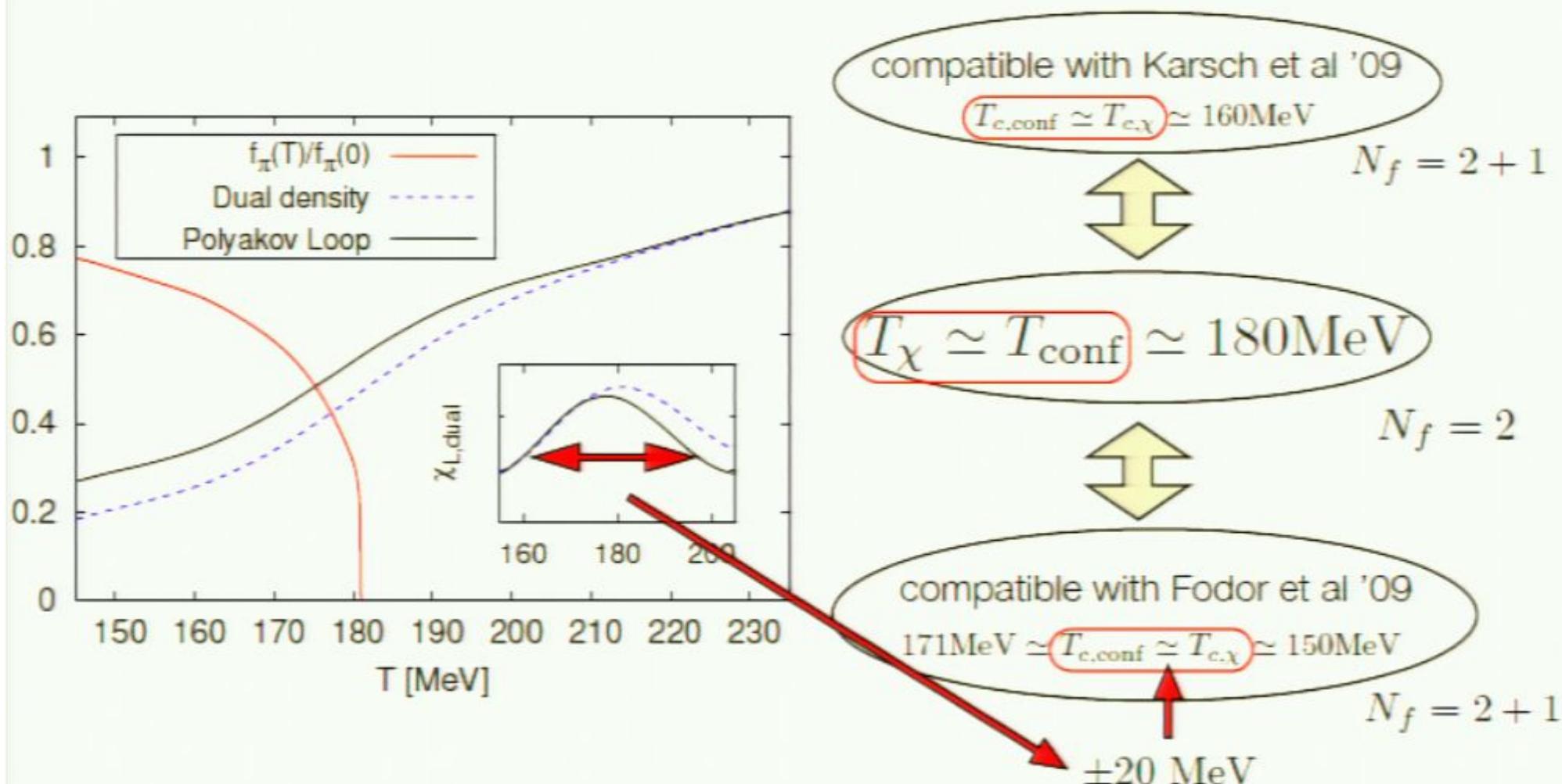
# Full dynamical QCD: $N_f = 2$ & chiral limit

Continuum methods & lattice



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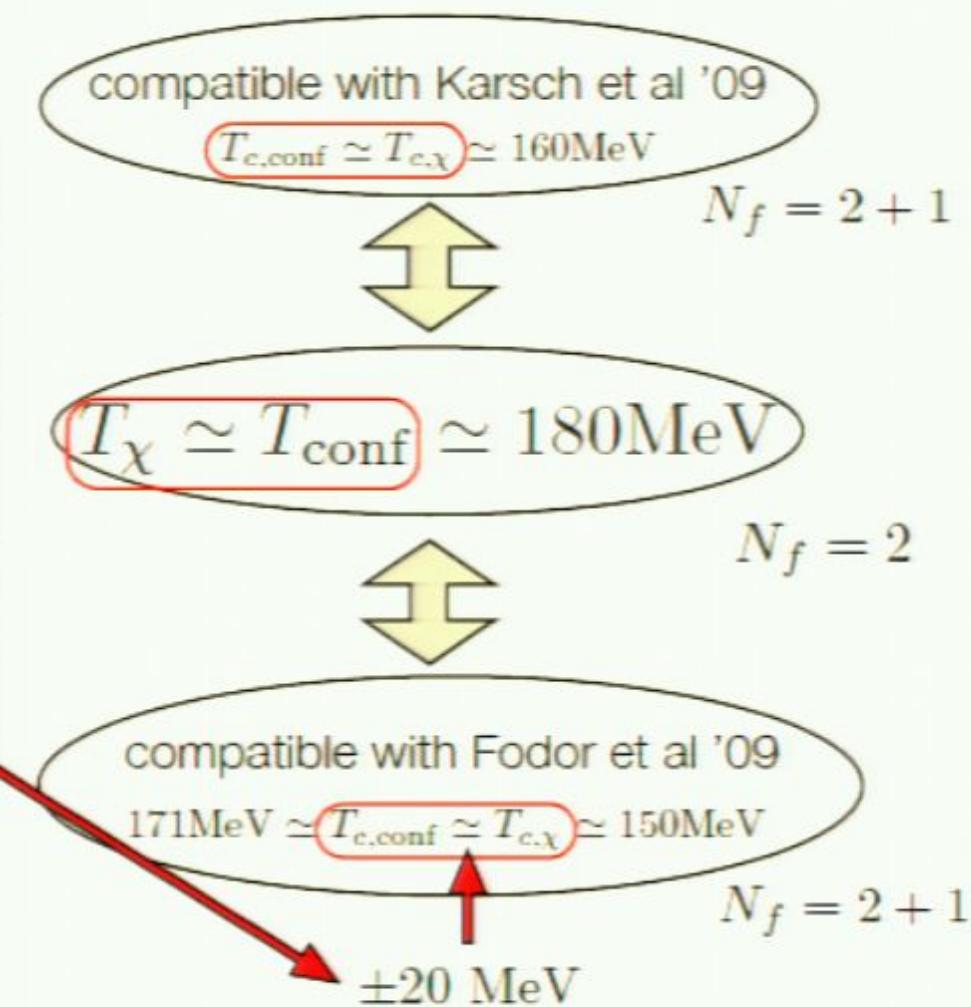
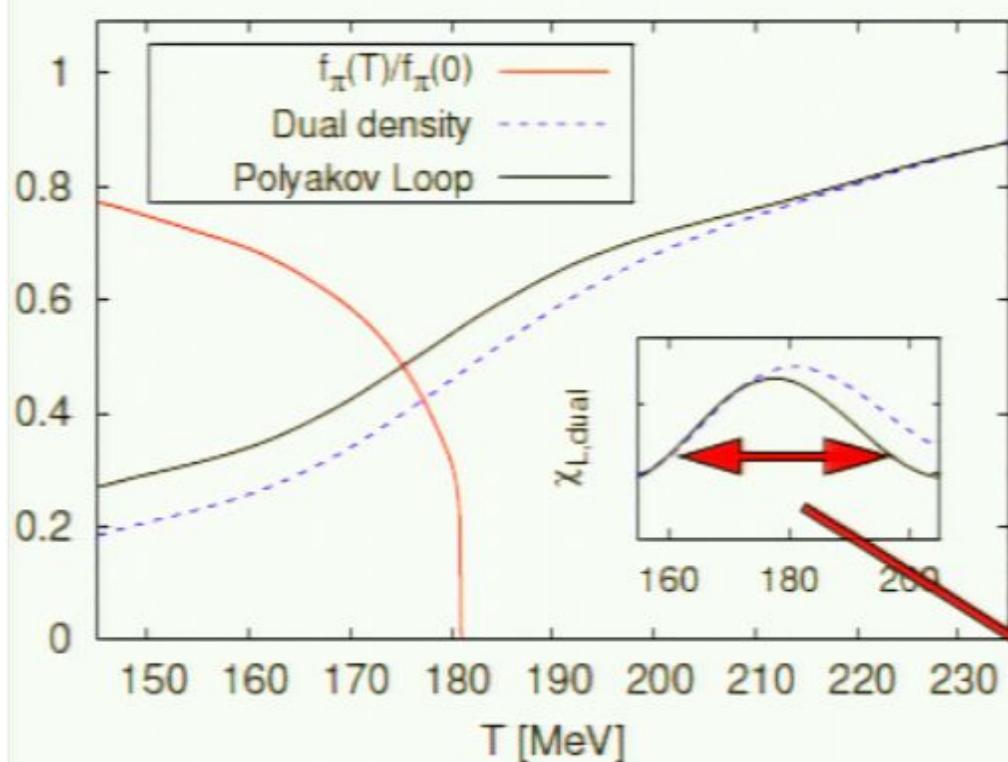
Continuum methods & lattice





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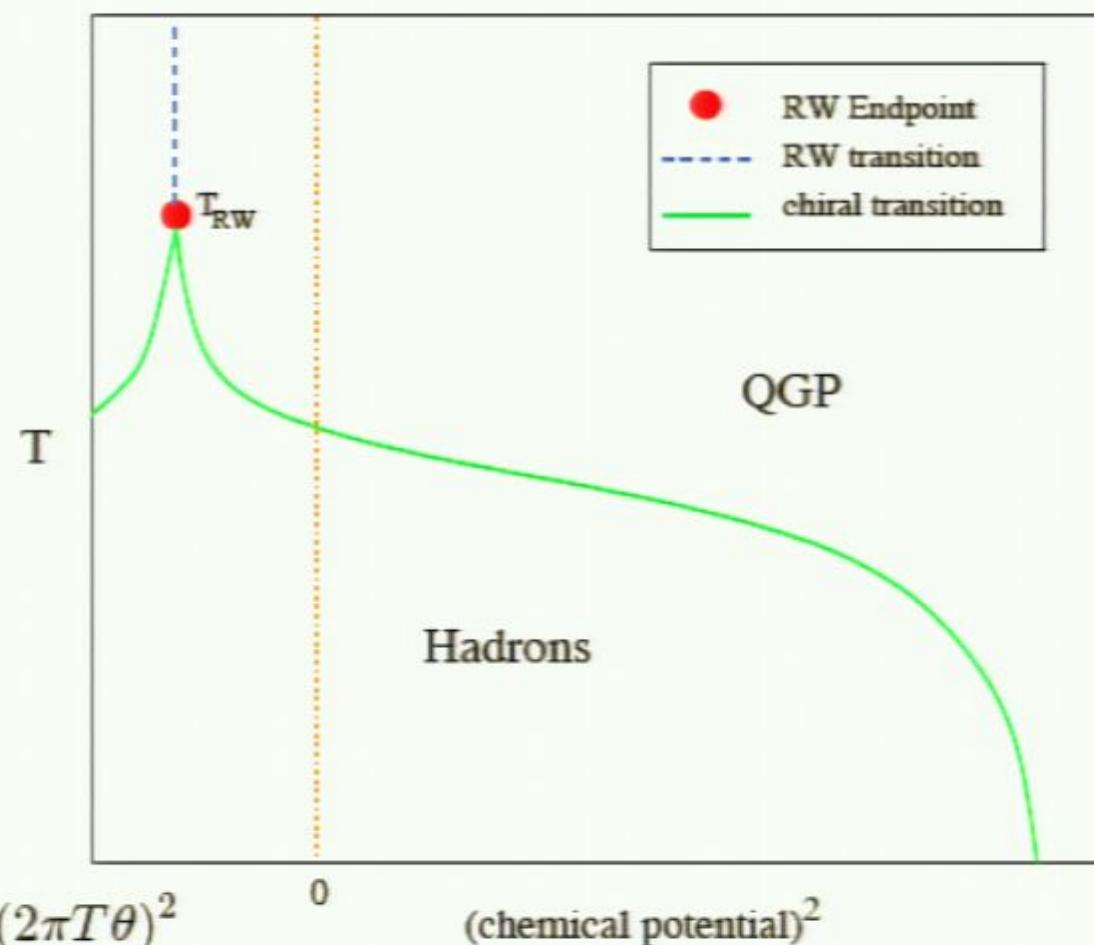
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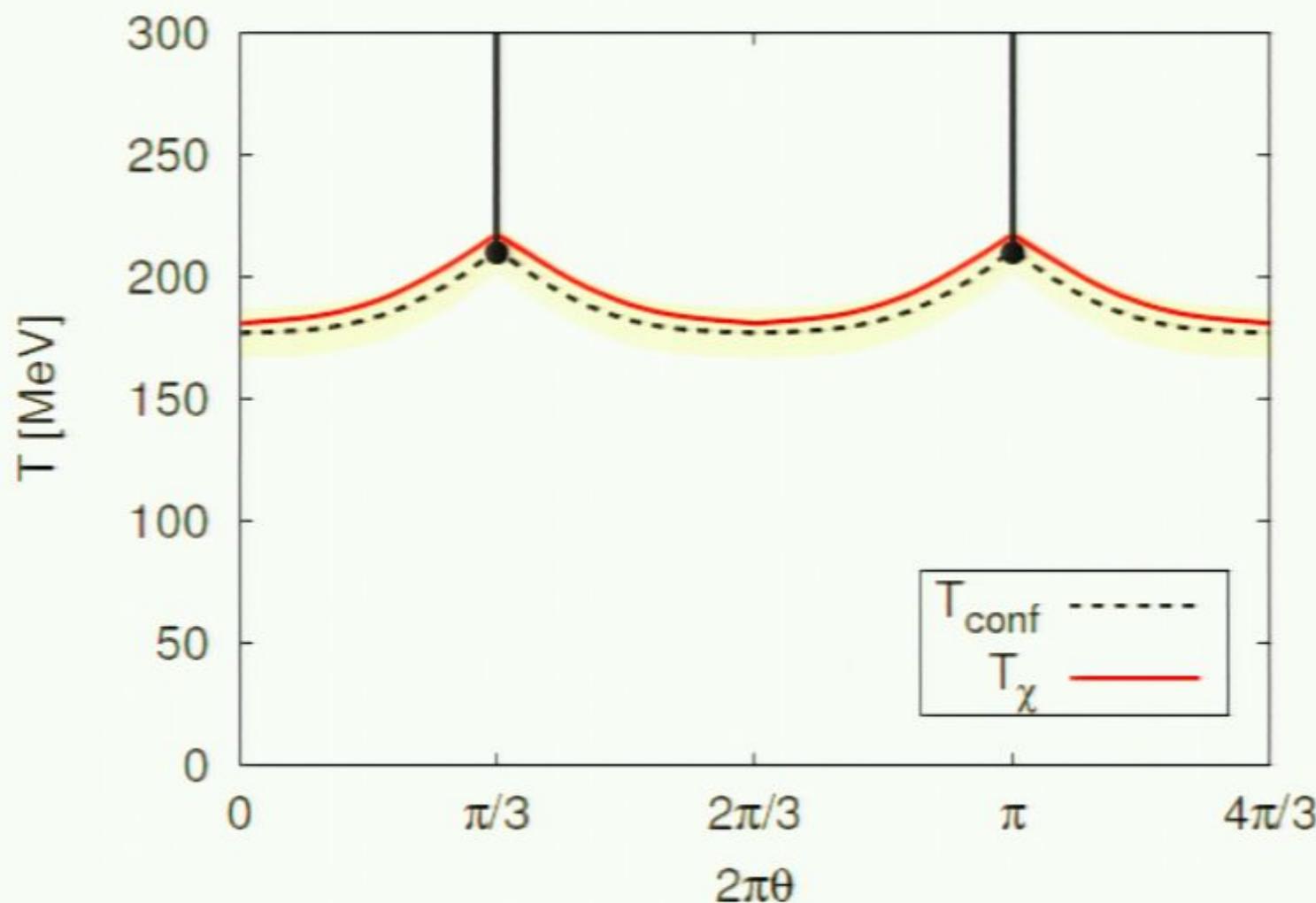
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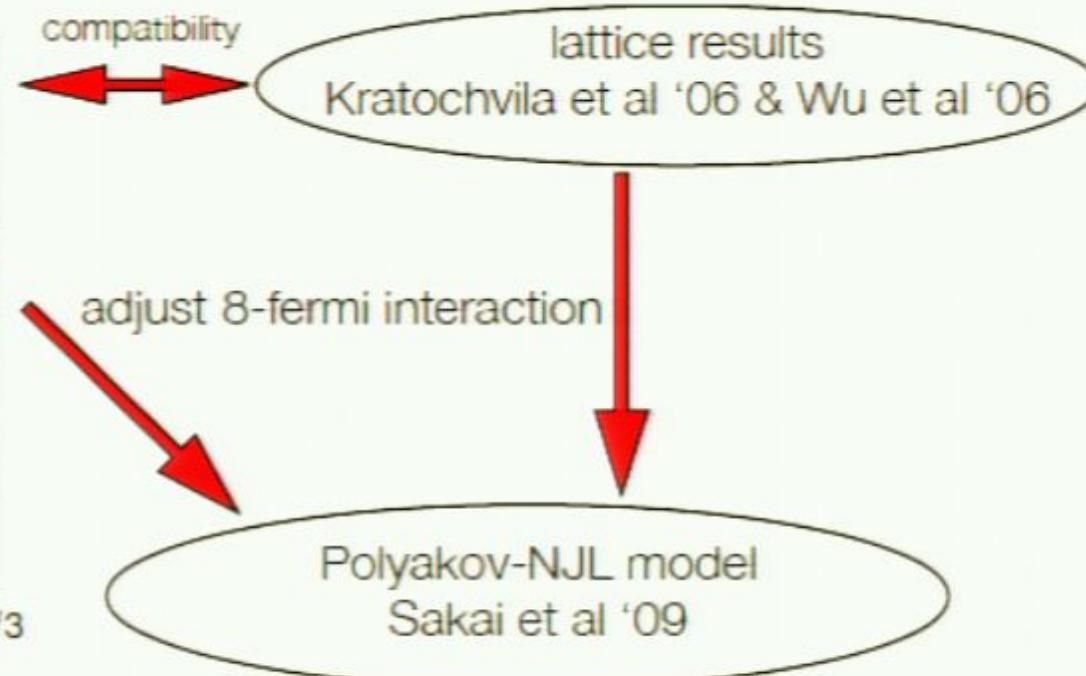
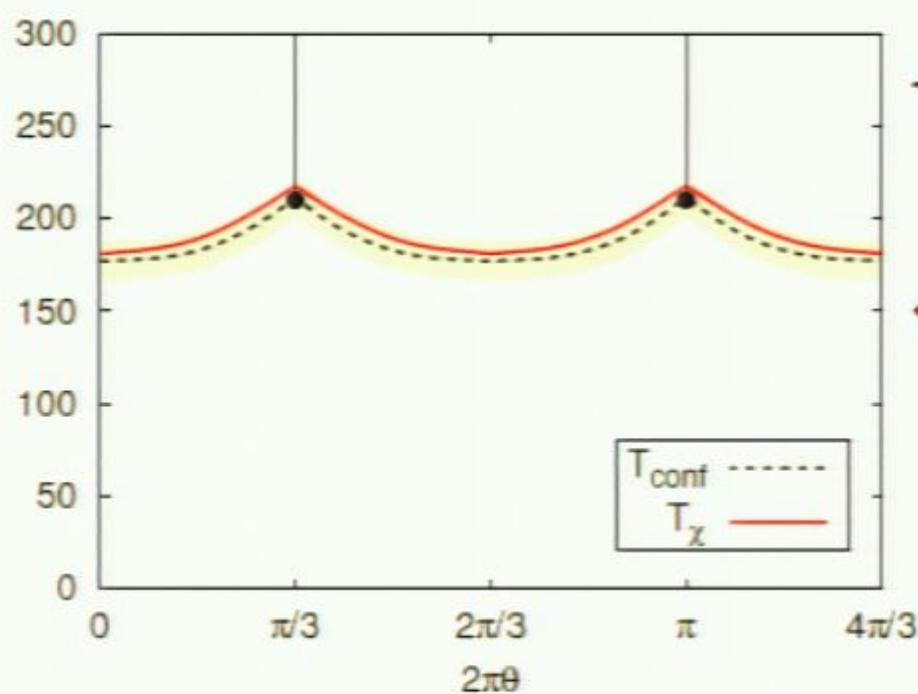
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Continuum methods

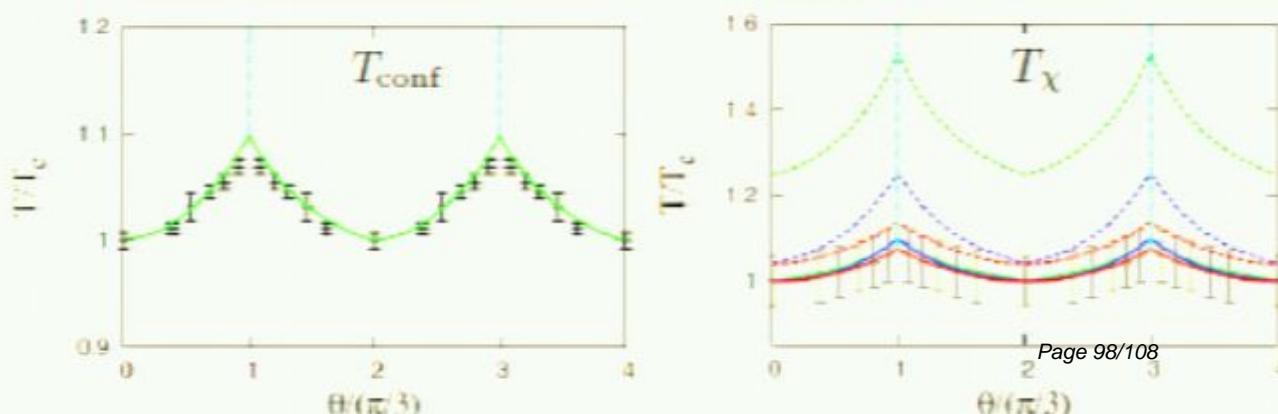


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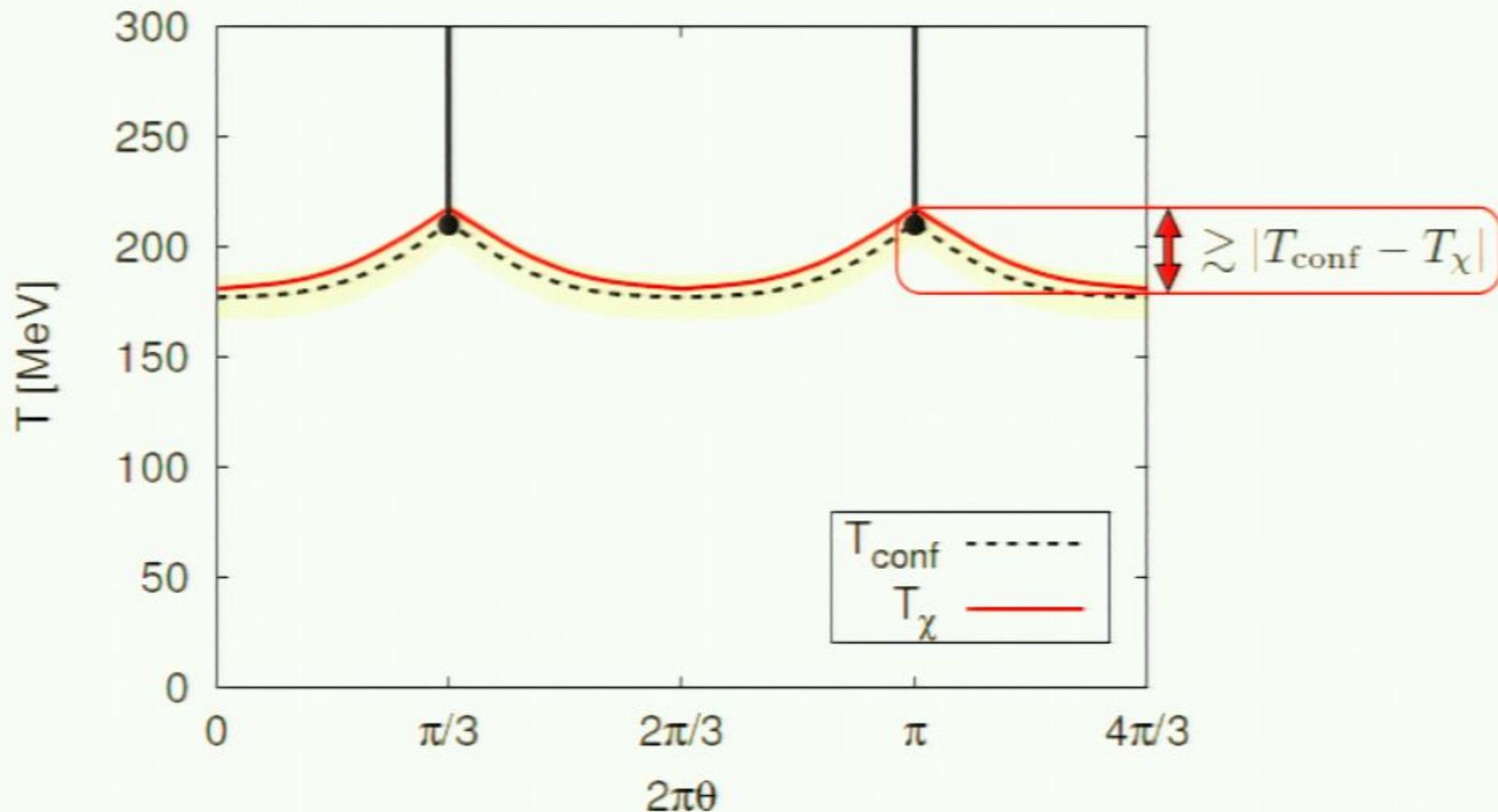


Braun, Haas, Marhauser, JMP '09



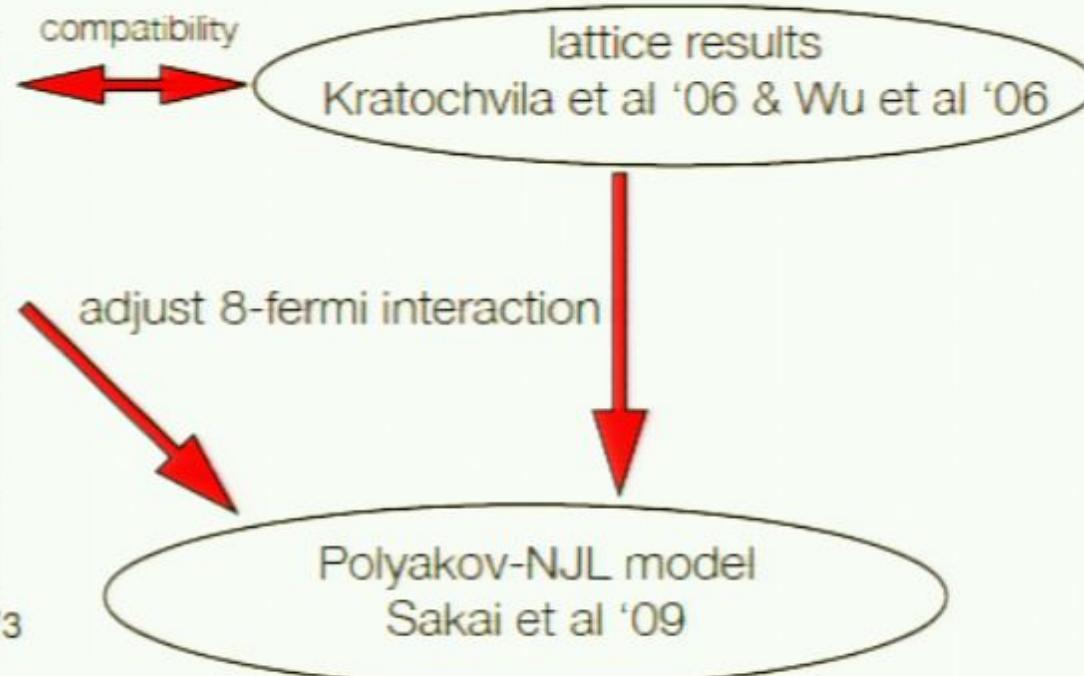
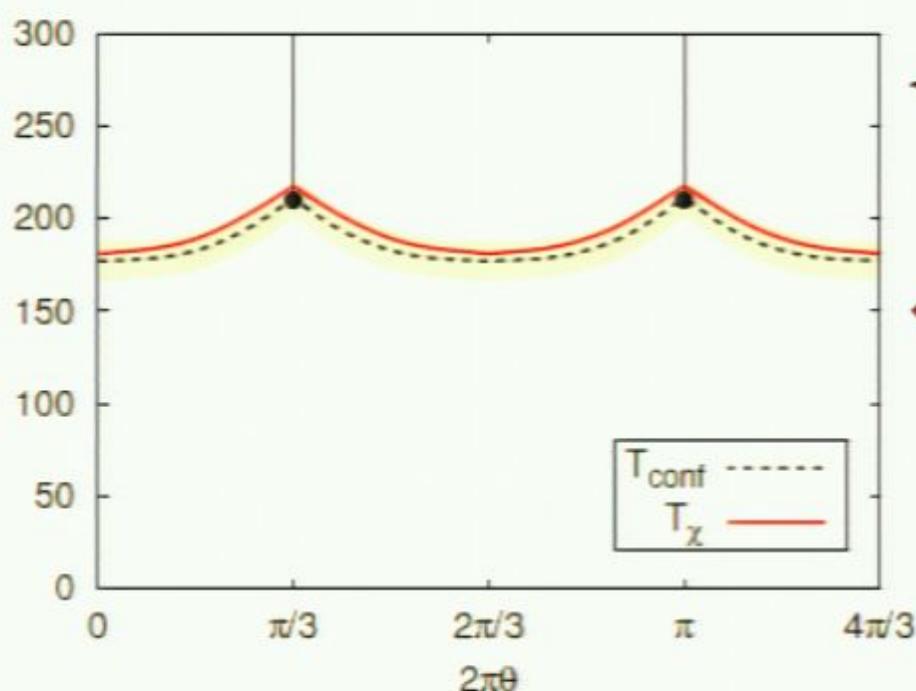
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Remark on dual order parameters for confinement

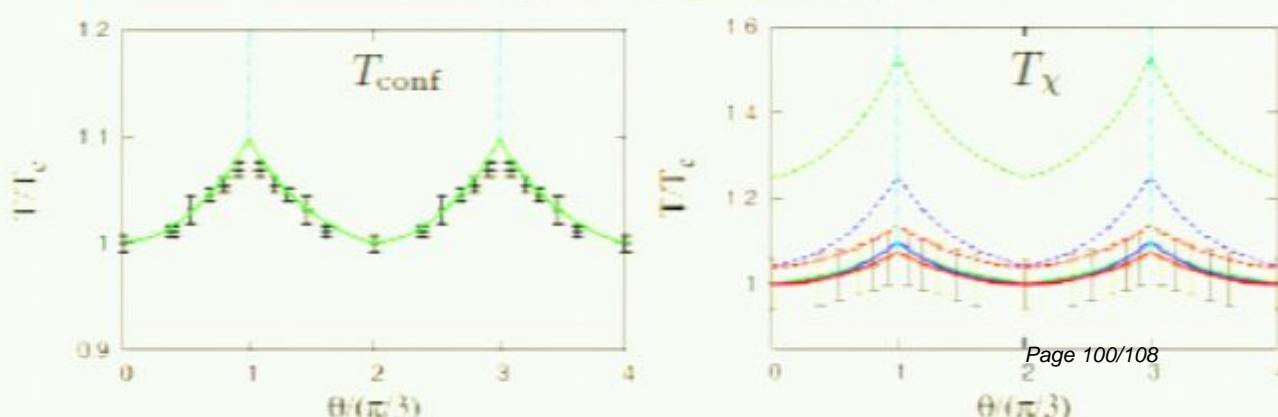


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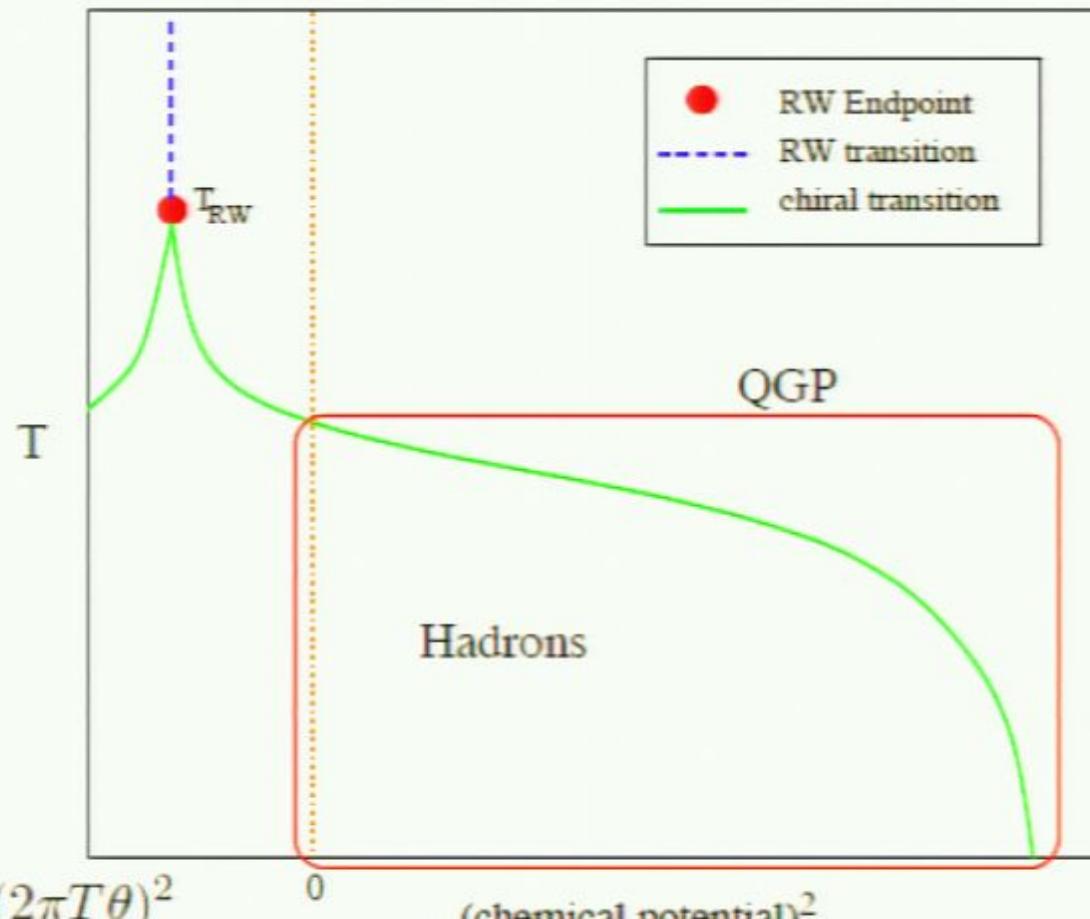


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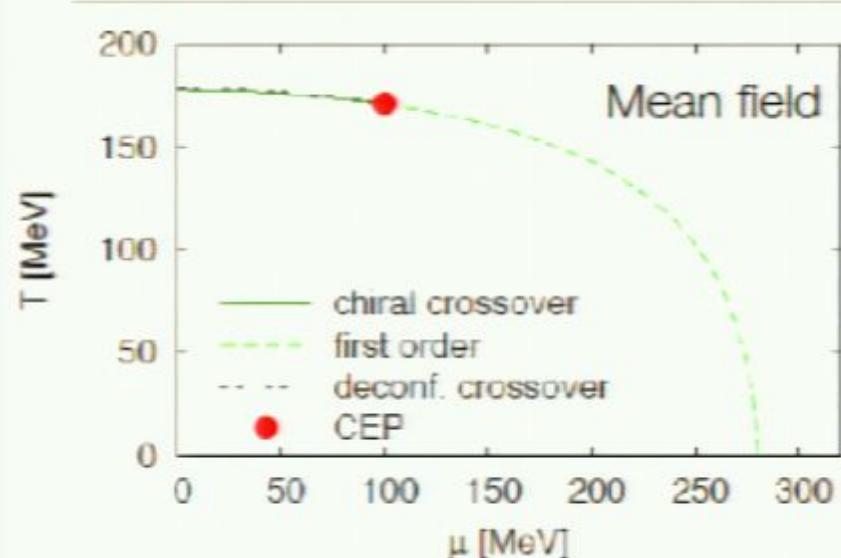
# Real chemical potential

$$\psi_\theta(t + \beta, \vec{x}) = -\psi(t, x)$$



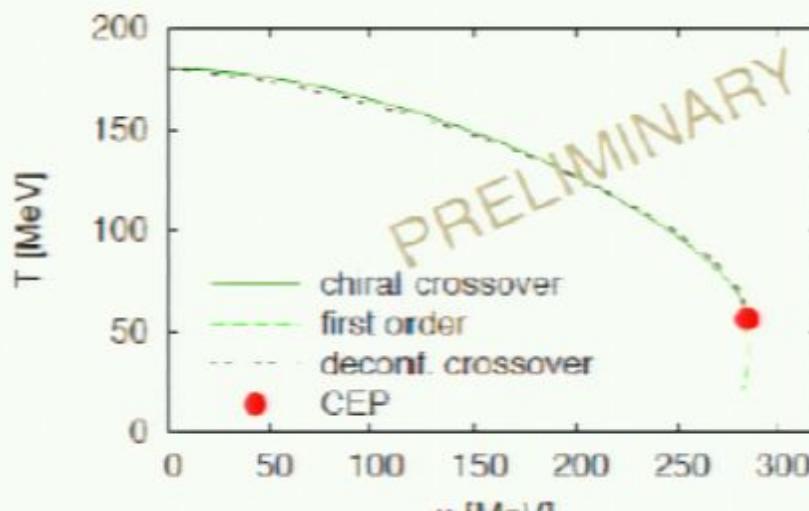
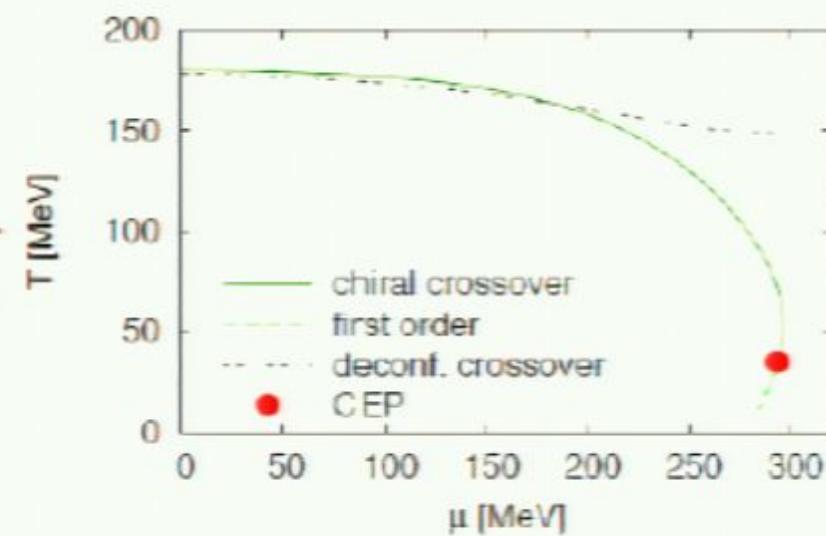
# A glimpse at real chemical potential

Polyakov - Quark-Meson model



RG  
quark-meson  
fluctuations

$N_f = 2$



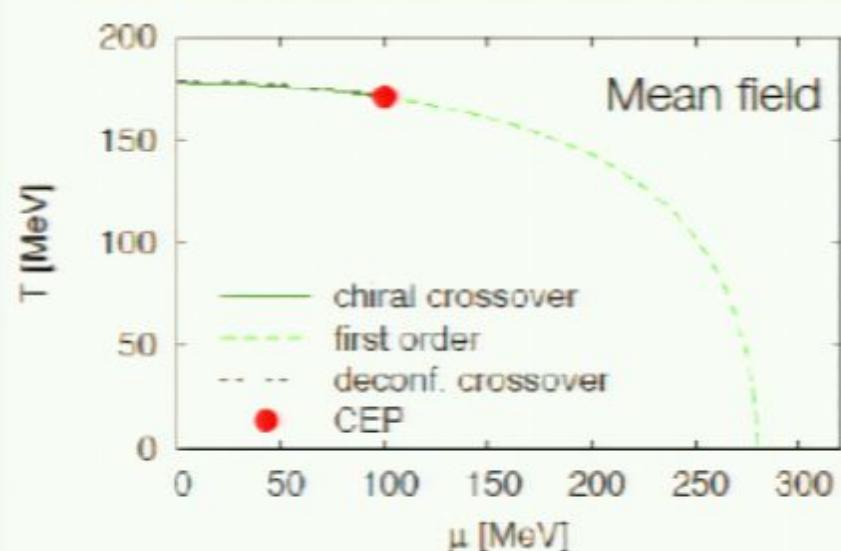
quark fluctuations  
in YM sector

Schaefer, JMP, Wambach '07

## Summary & Outlook

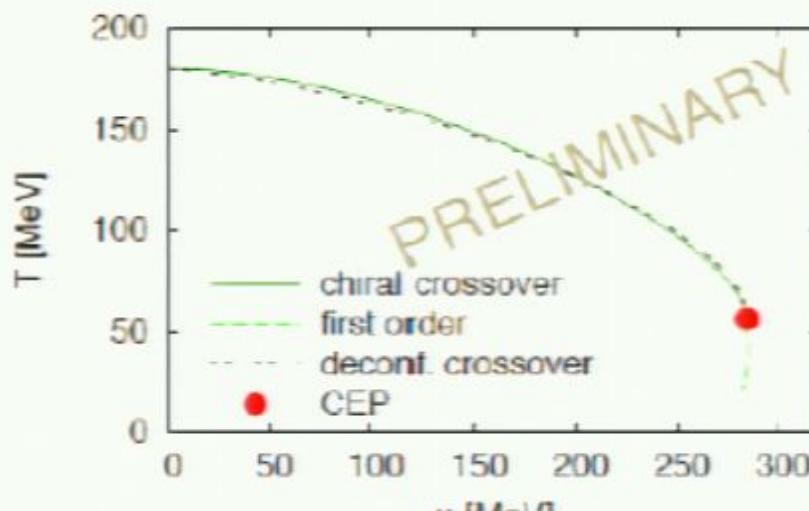
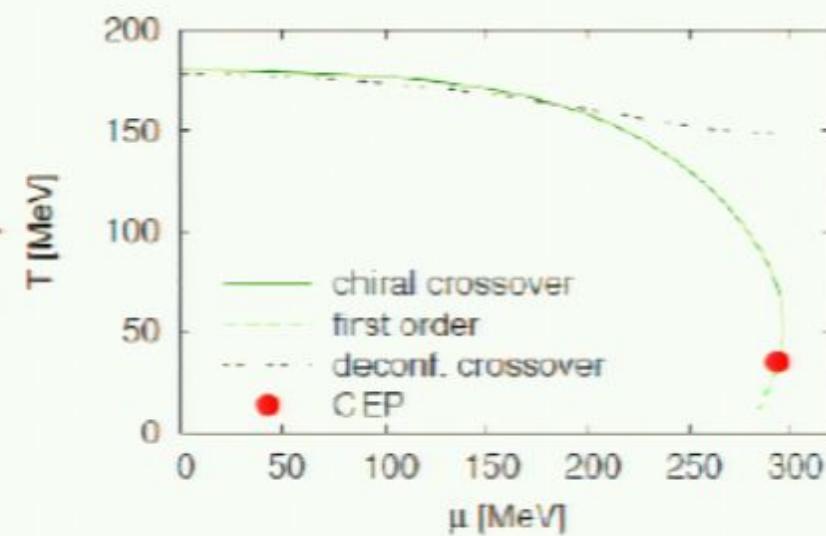
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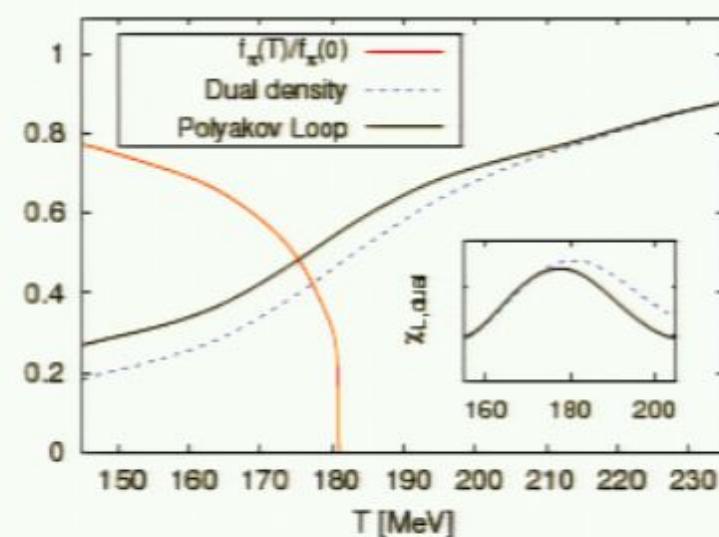
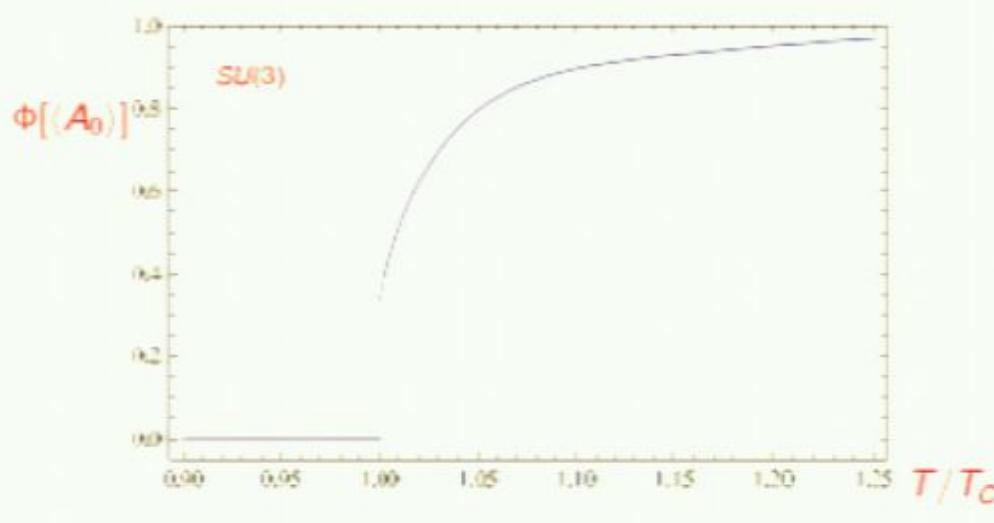


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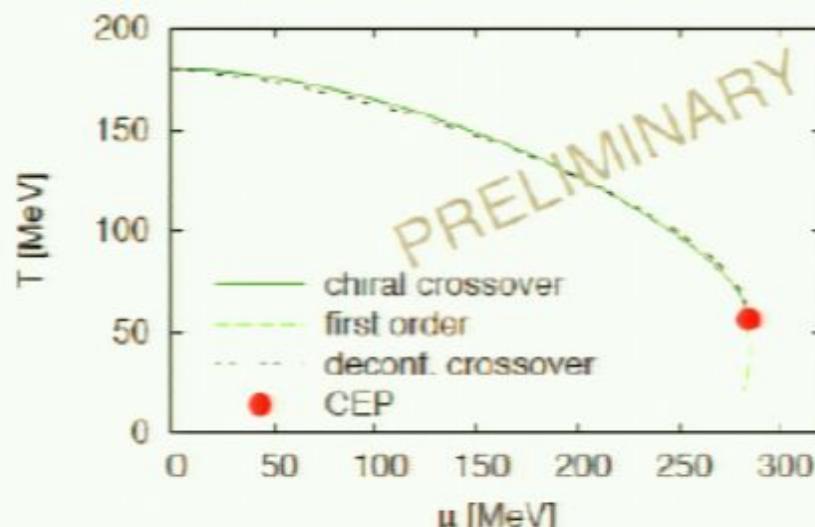
# Summary & outlook

- Phase diagram of QCD
  - Confinement & chiral symmetry breaking at finite temperature



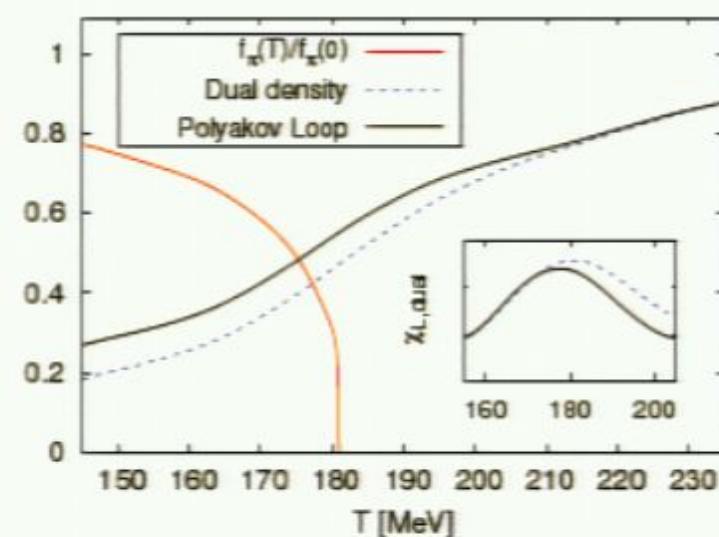
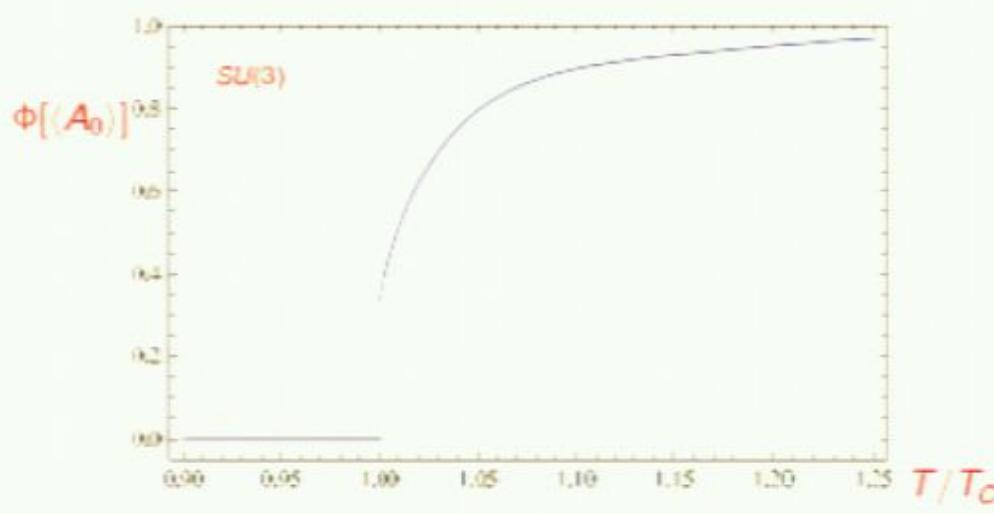
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- Phase diagram of QCD
  - Confinement & chiral symmetry breaking at finite temperature
    - Dynamical hadronisation
  - critical point and phase lines in effective theories



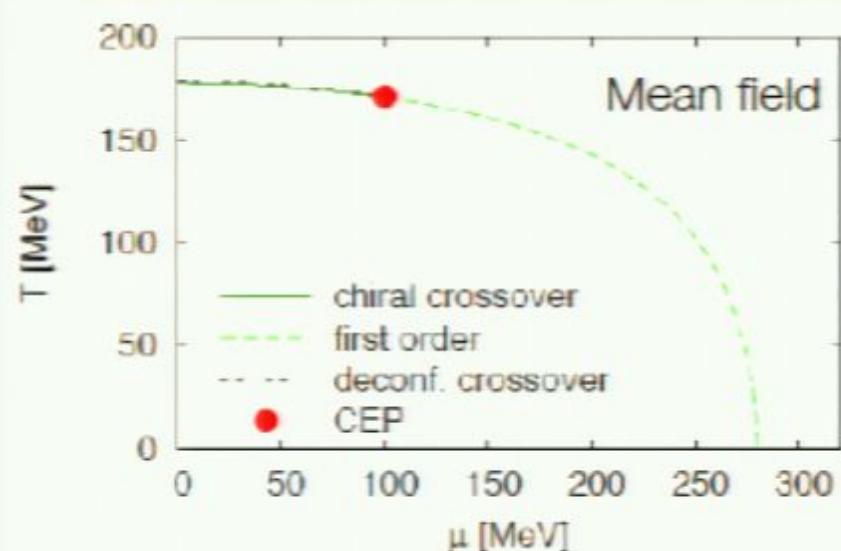
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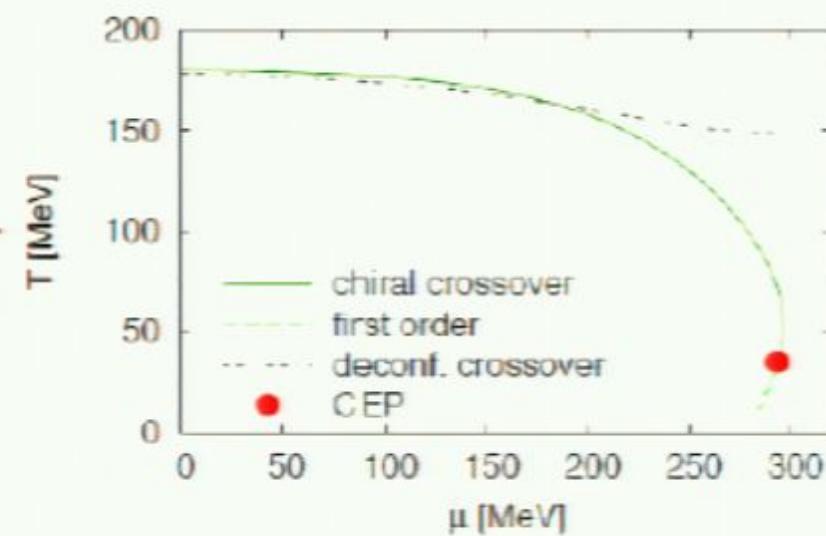
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HTL/HDL  
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