

Title: Foundations and Interpretation of Quantum Theory - Lecture 19

Date: Mar 25, 2010 02:30 PM

URL: <http://pirsa.org/10030025>

Abstract:

MESO-MACROSCOPIC TESTS OF QM: MOTIVATION

At microlevel: (a) $|\uparrow\rangle + |\downarrow\rangle$ quantum superposition ✓
= (b) $|\uparrow\rangle$ OR $|\downarrow\rangle$ classical mixture *

how do we know? Interference

At macrolevel: (a)  +  quantum superposition }
OR (b)  OR  macrorealism }

* Decoherence DOES NOT reduce (a) to (b)!

Can we tell whether (a) or (b) is correct?

Yes, if and only if they give different experimental predictions. But if decoherence $\rightarrow$ no interference predictions of (a) and (b) identical.

$\Rightarrow$ must look for QIMDS

quantum interference of macroscopically distinct

What is "macroscopically distinct"?

(a) "extensive difference" A

(b) "disconnectivity" D

↑
-large number of particles behaving differently in two branches

Initial aim of program: interpret new test (a) vs (b).

("Electron-phonon" syst. $\Delta E \sim \hbar \omega \sim \text{eV} \sim \text{H} \sim \text{meV}$
 so just need any sort of ΔE to displace)

1. Matrix elements of $S-E$ interaction couple only a
 very restricted set of levels of S .

2. "Adiabatic" ("slow") displacements:

Ex.: spin-boson model

$$\hat{H} = \hat{H}_S + \hat{H}_B + \hat{H}_{S-B}$$

$$\hat{H}_S = \Delta \hat{\sigma}_x$$

$$\hat{H}_B = \text{set of SHO's with lower frequency cutoff } \omega_{\text{cut}} \gg 0$$

$$\hat{H}_{S-B} = \hat{\sigma}_z \sum_n C_n \hat{q}_n \text{ spin-boson coupling}$$

$$|\Psi_{\pm}(t=0)\rangle = |+\rangle |X_0\rangle \text{ or displaced state of BO}$$

$$\hat{p}_S(t=0) = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \text{ (non-dif.)}$$

$\Downarrow$

$$|\Psi_{\pm}(t=\hbar/\Delta_{\text{min}})\rangle \approx \frac{1}{\sqrt{2}}(|+\rangle |X_0\rangle + |-\rangle |X_{\pm 2}\rangle)$$

$$\langle X_0 | X_{\pm 2} \rangle = \exp(-F) \neq 0$$

$$|\Psi_{\pm}(t=\hbar/\Delta_{\text{min}})\rangle \approx \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 \\ 0 & \frac{1}{\sqrt{2}} \end{pmatrix}$$

1. Matrix elements of $S-E$ interaction depend only on very restricted set of levels of S .

2. "Adiabatic" (slow) approximation:

Ex.: spin-boson model

$$\hat{H} = \hat{H}_S + \hat{H}_B + \hat{H}_{S-B}$$

$$\hat{H}_S = \Delta \hat{\sigma}_x$$

$\hat{H}_B$ = set of SHO's with lower frequency cutoff $\omega_m \gg \Delta$

$$\hat{H}_{S-B} = \hat{\sigma}_z \sum_k c_k \hat{a}_k$$

$$\hat{\mathbb{I}}_m(t=0) = |+\rangle |X_0\rangle \leftarrow \text{discrete part of env.}$$

$$\hat{\rho}_r(t=0) = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & 0 \end{pmatrix} \quad (\text{unphysical})$$

or

$$\hat{\mathbb{I}}_m(t=0/\omega_m) \approx \frac{1}{\sqrt{2}} (|+\rangle |X_0\rangle + |+\rangle |X_-\rangle)$$

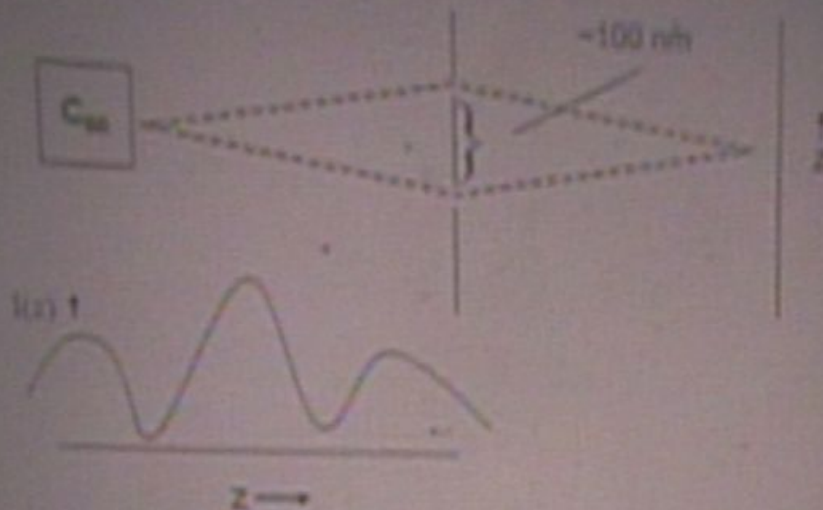
$$\langle X_+ | X_- \rangle = \exp(-F) \approx 0$$

$$\hat{\rho}_r(t=0/\omega_m) \approx \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$$

Eq. 77 (cf. random initial state)

The Search for QIMDS

1. Molecular diffraction*



Note: (a.) Beam does not have to be monochromated

$$f(\nu) = A\nu^3 \exp[-(\nu - \nu_0)^2 / \nu_m^2] \quad (\nu_0 \sim 1.5 \nu_m)$$

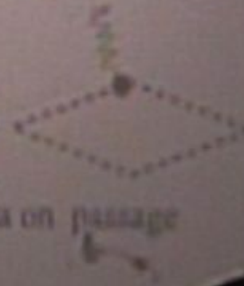
(b.) "Which-way" effects?

Oven is at 900–1000 K

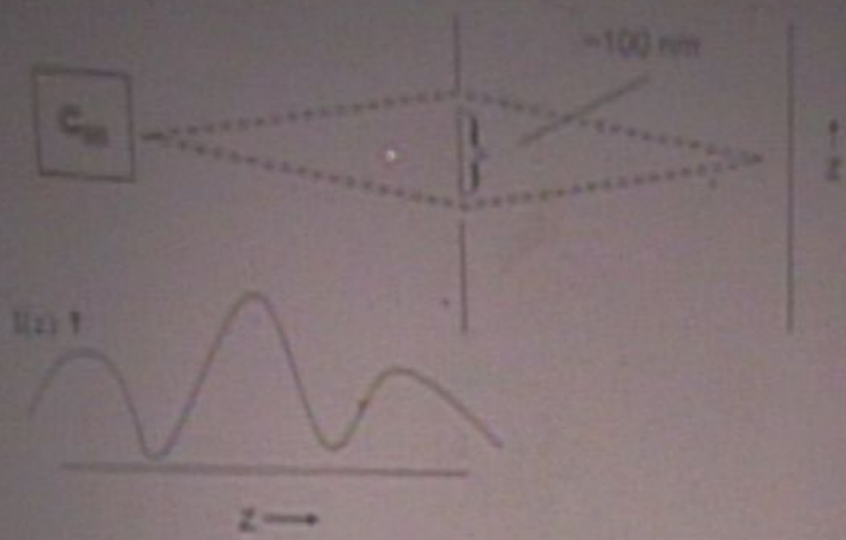
$\Rightarrow$ many vibrational modes excited

4 modes infrared active $\Rightarrow$

absorb/emit several radiation quanta on passage through apparatus!



Why doesn't this destroy interference?



Note: (a.) Beam does not have to be monochromated

$$f(\nu) = A\nu^3 \exp[-(\nu - \nu_0)^2 / \nu_m^2] \quad (\nu_0 \sim 1.8 \nu_m)$$

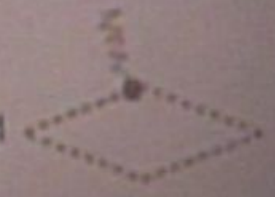
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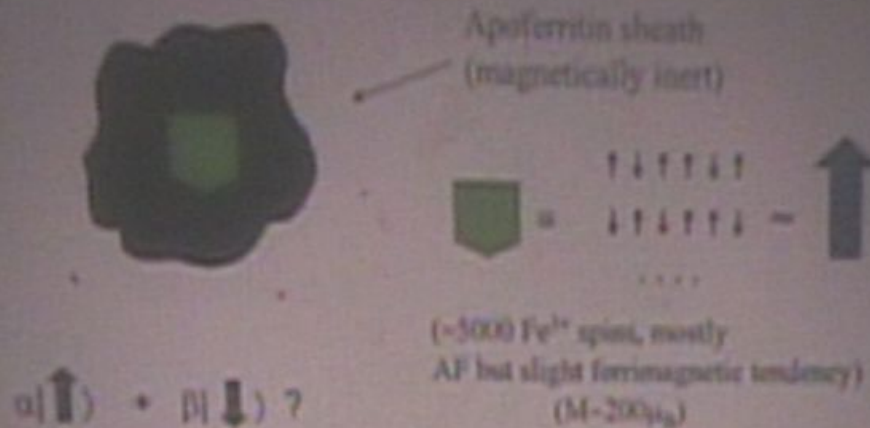
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Why doesn't this destroy interference?

Arndt et al., *Nature* **401**, 680 (1999); Nairz et al., *Am. J. Phys.* **71**, 319 (2003).

2. Magnetic biomolecules*



$$\text{AF: } \Delta \sim \hbar \omega_s \exp - N \sqrt{K/J}$$

$\swarrow$ $\nwarrow$
 no. of spins uniaxial anisotropy

(isotropic) exchange en.

Raw data: $\chi(\omega)$ and noise spectrum
above ~200 mK, featureless

below ~300 mK, sharp peak at ~1 MHz (ω_m)

$$\omega_{\text{res}}^2 \approx \omega_0^2 + M^2 H^2$$

$$\ln \omega_s \sim a - bN \quad \leftarrow \text{no. of spins, exp. adjustable}$$

data is on physical ensemble, i.e., only total magnetization measured.

et al., Science 268, 77 (1995)

The Search for QIMDS (cont.)

2. Magnetic biomolecules*



Apo-ferritin sheath
(magnetically inert)



$\uparrow \downarrow \uparrow \uparrow \downarrow \uparrow$

$\downarrow \uparrow \downarrow \uparrow \uparrow \downarrow$

....



(~5000 Fe^{3+} spins, mostly
AF but slight ferrimagnetic tendency)
($M \sim 200 \mu_B$)

$\alpha|\uparrow\rangle + \beta|\downarrow\rangle ?$

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above ~200 mK, featureless

below ~300 mK, sharp peak at ~1 MHz (ω_{res})

$$\omega_{\text{res}}^2 \cong \omega_0^2 + M^2 H^2$$

$$\omega_0 \sim a - bN$$

no. of spins, exptly.
adjustable

The Search for QIMDS (cont.)

2. Magnetic biomolecules*



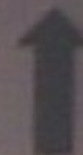
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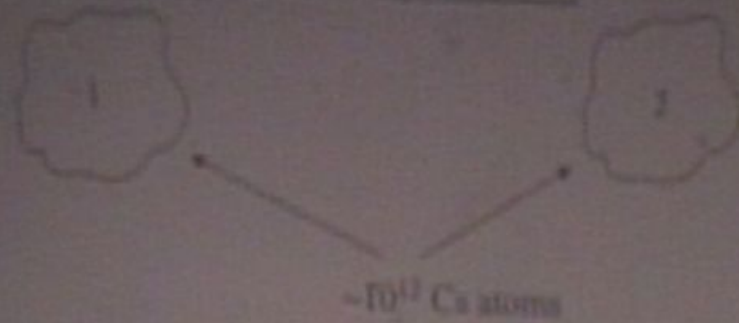
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The Search for QMDS (cont.)

3. Quantum-optical systems*



for each sample separately, and also for total

$$[J_x, J_y] = iJ_z$$

$$\Rightarrow \langle \delta J_{x1} \delta J_{y1} \rangle \geq |J_{z1}|$$

$$\langle \delta J_{x2} \delta J_{y2} \rangle \geq |J_{z2}|$$

$$\langle \delta J_{x\text{tot}} \delta J_{y\text{tot}} \rangle \geq |J_{z\text{tot}}|$$

so, if set up a situation s.t.

$$J_{z1} = -J_{z2}$$

must have

$$\langle \delta J_{x1} \delta J_{y1} \rangle > 0$$

$$\langle \delta J_{x2} \delta J_{y2} \rangle > 0$$

but may have

interpretation of idealized expt. of this type:

$$(\text{QM theory} \Rightarrow) \quad \langle \delta J_{+1} \delta J_{+1} \rangle \geq |J_{+1}| \sim N$$

$$\Rightarrow |\delta J_{+1}| \geq N^{1/2}$$

But

$$(\text{expt} \Rightarrow) \quad \langle \delta J_{+1} \delta J_{+1} \rangle \approx 0 \quad (ii)$$

$$\Rightarrow |\delta J_{+1}| \sim 0$$

$\Rightarrow \delta J_{+1}$ exactly anticorrelated with δJ_{-1}

$\Rightarrow$ state is either superposition or mixture of $|n, -n\rangle$

but mixture will not give (ii)

$\Rightarrow$ state must be of form

$$\sum_n c_n |n, -n\rangle$$

with appreciable weight for $n \lesssim N^{1/2}$, $\Rightarrow$ high disconnectivity

Note:

(a) QM used essentially in argument

(b) $D \sim N^{1/2}$ not $\sim N$.

(prob. generic to this kind of expt.)

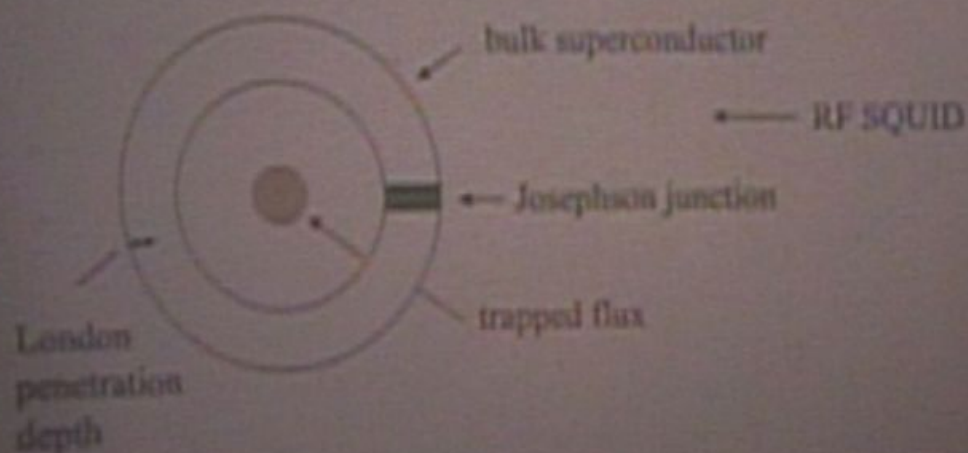
The Search for QIMDS (cont.)

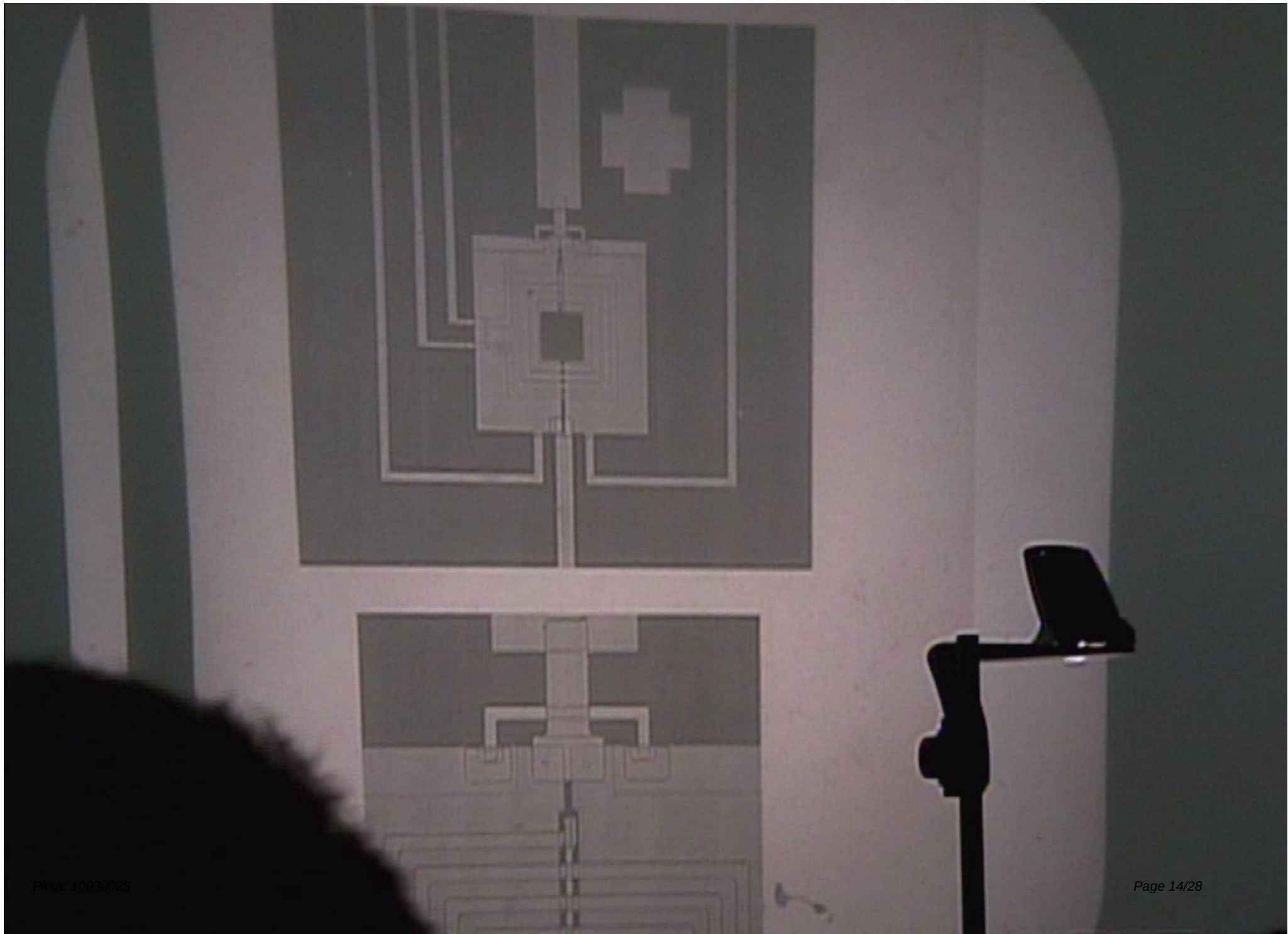
4. Superconducting devices

(Φ : not all devices which are of interest for quantum computing are of interest for QIMDS)

Advantages:

- classical dynamics of macrovariable v , well understood
- intrinsic dissipation (can be made) v , low
- well developed technology
- (non-) scaling of S (action) with D .

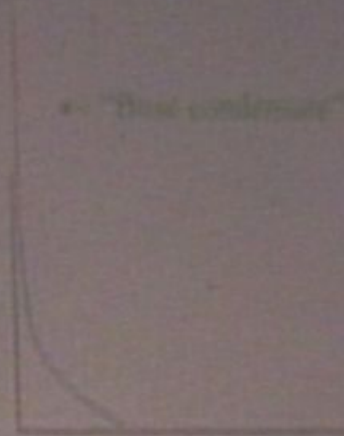
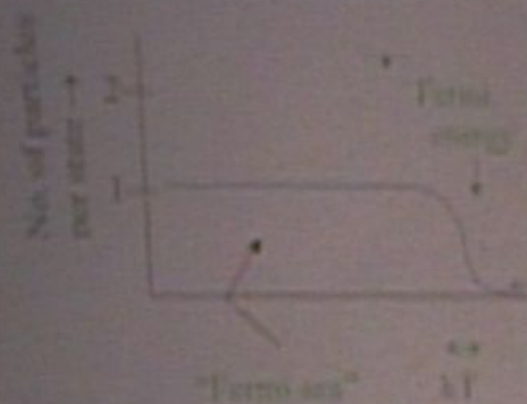




"Spin" of elementary particles $= \frac{n}{2} \hbar$

0, 1, 2, ... bosons
 $\frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots$ fermions

At low temperatures:



Electrons in metals: spin $\frac{1}{2} \Rightarrow$ fermions

But a compound object consisting of an even no.

of fermions has spin 0, 1, 2, ... $\Rightarrow$ boson.

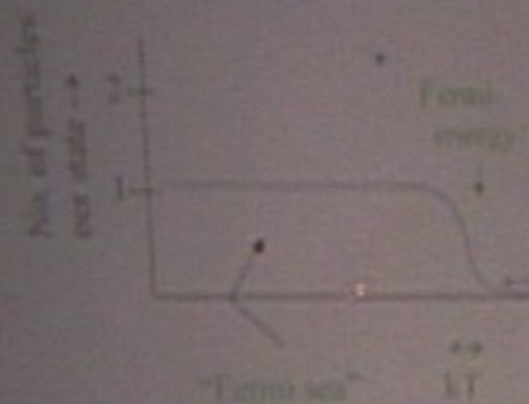
(e.g. $2p + 2n + 2e = {}^4\text{He atom}$)

... large Bose condensation

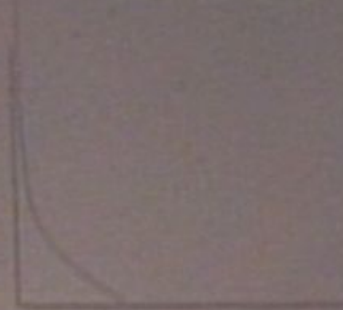
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* "Bose condensation"



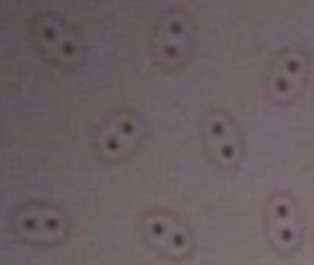
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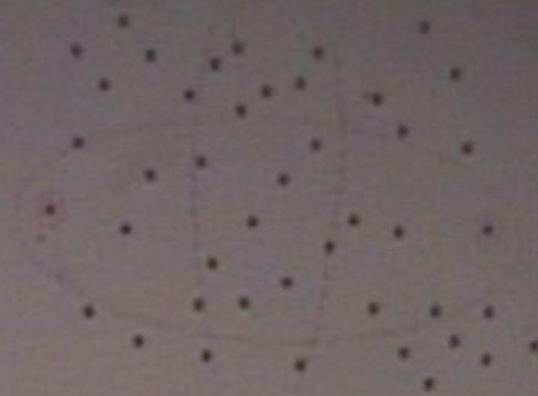
(Ex: $2p + 2n + 2e = {}^4\text{He atom}$)

$\Rightarrow$ can undergo Bose condensation

Pairing of electrons



"di-electronic molecules"



Cooper Pairs

In simplest ("BCS") theory, Cooper Pairs, once formed, must automatically undergo Bose condensation!

→ must all do exactly the same thing at the same time (also in many-particle situation)

SUPERCONDUCTING RING IN EXTERNAL MAGNETIC FLUX



Quantization condition for
"particle" of charge $2e$ (Cooper
pair)

$$K = \oint \mathbf{v} \cdot d\mathbf{l} = \frac{\hbar}{2m} (n \cdot \Phi - \Phi_0)$$

integer

"flux quantum"
 $\hbar/2e$

$$E \propto K^2$$

A. $\Phi = 0$: groundstate unique ($n = 0$)

$\Rightarrow$ all pairs at rest.

B. $\Phi = 1/2 \Phi_0$: groundstate doubly degenerate

($n = 0$ or $n = 1$)



pairs rotate clockwise

anticlockwise

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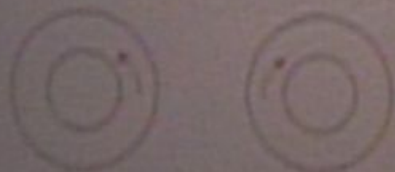
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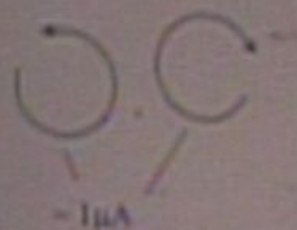
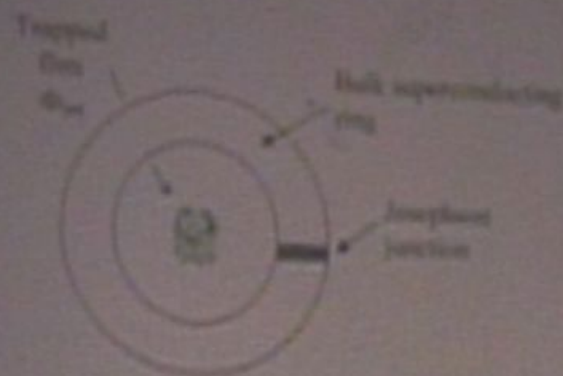
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Either all pairs rotate clockwise

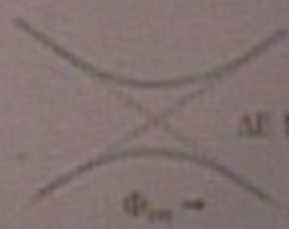
Or all pairs rotate anticlockwise

Josephson circuits



$$\Psi = 2^{-1/2} (|U\rangle + |L\rangle)$$

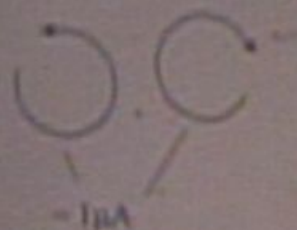
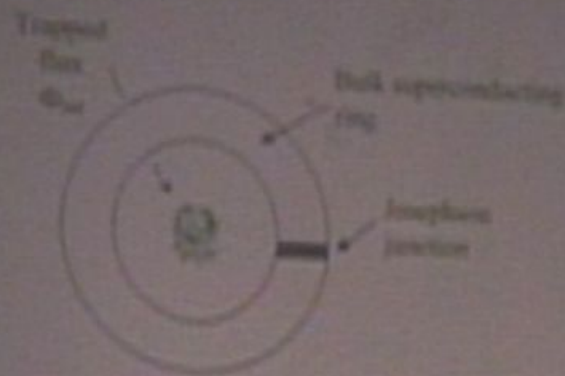
Evidence: (a) spectroscopic (SUNY, Delft 2000)



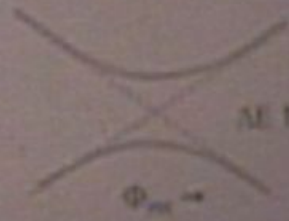
(b) real-time oscillations (like NH_3)

between $\downarrow$ and $\uparrow$

Josephson circuits



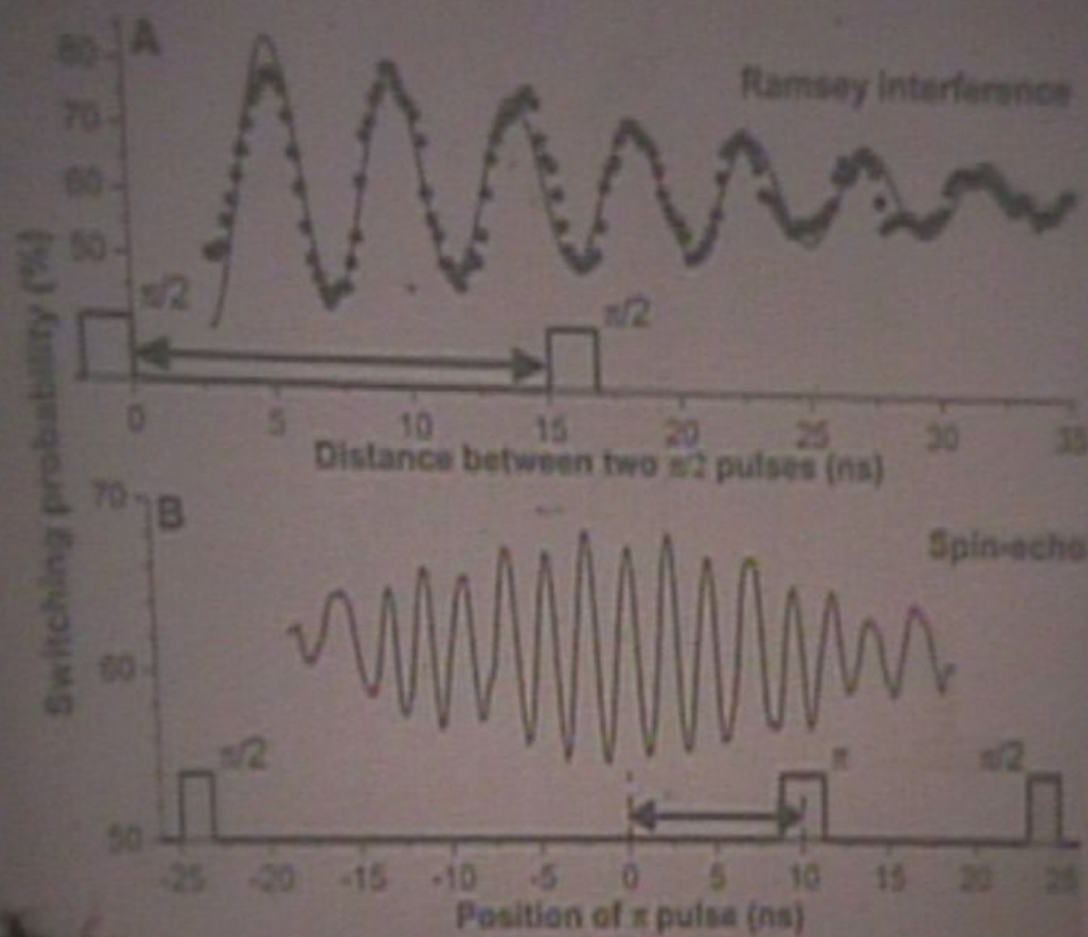
$$\Psi = 2\pi \Phi / \Phi_0$$



Evidence: (a) spectroscopic (SUNY, Delft 2000)

(b) real-time oscillations (like NMR)

between Ψ and Φ



WHAT IS THE DIMENSIONALITY "D" ("SCHRÖDINGER SCATTERNESS")
OF THE STATES OBSERVED IN QMIS EXPERIMENTS?

i.e., how many "microscopic" entities are "behaving differently"
in the two branches of the superposition?

Fullerene (etc.) diffraction experiments: straightforward, number of
"elementary" particles in C_{60} (etc.) (~1200)

Magnetic biomolecules: number of spins which reverse between the
two branches (~5000)

Quantum-optical experiments } matter of definition
SQUIDS }

e.g. SQUIDS (SUNY experiment):

(a) naïve approach:

no. of C. pairs

$$\Psi_{\uparrow} \sim \chi_{\uparrow}^{N/2}, \Psi_{\downarrow} \sim \chi_{\downarrow}^{N/2}$$

mutually orthogonal C. pair w.f.

$$\Rightarrow D \sim N \sim 10^9 - 10^{10} \quad \Delta: \text{Fermi statistics!}$$



(b) how many single electrons do we need to displace in
momentum space to get from $\Psi_{\uparrow}$ to $\Psi_{\downarrow}$? (Korshakken et al.,
preprint, Nov. 08)

$$\Rightarrow D \sim N(v_f/v_f) \sim 10^1 - 10^4 \quad \Delta: \text{intuitively, severe underestimate in "BEC" limit (e.g. Fermi alkali gas)}$$

WHAT IS THE DISCRETELY "D" ("SCHRÖDINGER'S CAT")
OF THE STATES OBSERVED IN QMDS EXPERIMENTS?

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(b) how many single electrons do we need to displace in
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Nov. 08)

$$\sim 10^3 - 10^4 \quad \& \text{ intuitively, severe underestimate in "BEC" limit (e.g. Fermi alkali gas)}$$

	<u>DISTURBANCE</u>	<u>ESTIMATED</u>
Single e	1	1
Neutron in interferometer	$\sim 10^3$	1
QED cavity	~ 10	$\in 10$
Cooper-pair box	$\sim 10^3$	$\frac{1}{2}$
Cu	~ 1100	~ 1100
Fe	~ 5000 (7)	~ 5000
Ashton quantum-optics exp.	$\sim 10^2$ ($\approx 10^{2.5}$)	$\sim 10^2$
SLAC SQUID exp.	$\sim 10^2 - 10^{2.5}$ (≈ 10)	$(10^2 - 10^{2.5})$
Smallest visible dust particle	$\sim 10^{13}$	$(10^2 - 10^{2.5})$
	$\sim 10^{22}$	$\sim 10^{22}$

Where do we go from here?

1. Larger values of Λ and/or D ?
(Diffraction of virus?)

2. Alternative Dfs. of "Measures" of Interest

- More sophisticated forms of entanglement;
- Biological functionality (e.g. superpose states of rhodopsin?)
- Other (e.g. GR)

* 3. Exclude Macrorealism

Suppose: Whenever
observed, $Q = \pm 1$.



$Q = +1$



$Q = -1$

Def. of "MACROREALISTIC" Theory:

I. $Q(t) = \pm 1$ at (almost) $\forall t$,
whether or not observed.

II. Noninvasive measurability

III. Induction

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(Diffraction of virus?)

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Def of "MACROREALISTIC" Theory:

"COMMON
SENSE"

- I. $Q(t) = \pm 1$ at (almost) $\forall t$,
whether or not observed.
- II. Noninvasive measurability
- III. Induction

$$K(t_1, t_2) = \langle Q(t_1)Q(t_2) \rangle_{\text{qm}} + \langle Q(t_1)Q(t_2) \rangle_{\text{qm}} \\ + \langle Q(t_1)Q(t_2) \rangle_{\text{qm}} - \langle Q(t_1)Q(t_2) \rangle_{\text{qm}}$$

Take $t_1 - t_2 = t_2 - t_3 = t_3 - t_4 = \pi/4\Delta$ $\leftarrow$ tunnelling frequency

Then,

- (a) Any macrorealistic theory: $K \leq 2$
- (b) Quantum mechanics, ideal: $K = 2.8$
- (c) Quantum mechanics, with all the real-life complications: $K > 2$ (but < 2.8)

Thus: to extent analysis of (c) within quantum mechanics is reliable, can force nature to choose between macrorealism and quantum mechanics!

Possible outcomes:

Enough noise $\Rightarrow K_{\text{QM}} < 2$

Macrorealism refuted