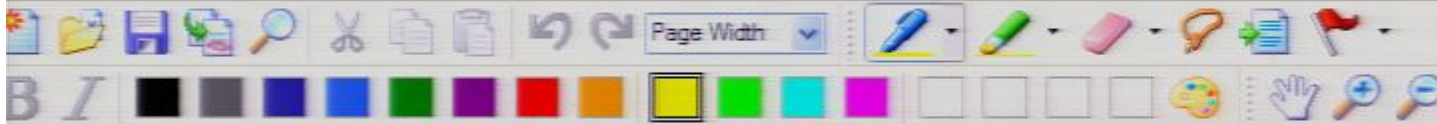


Title: Quantum Field Theory for Cosmology - Lecture 14

Date: Mar 04, 2010 05:00 PM

URL: <http://pirsa.org/10030014>

Abstract:



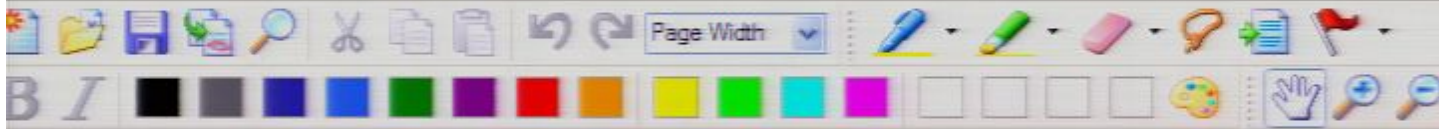
QFT for Cosmology, Achim Kempf, Winter 10, Lecture 14

2/27/2006

Quantum field theory on FRW spacetimes.

Friedmann Robertson Walker (FRW) spacetimes:

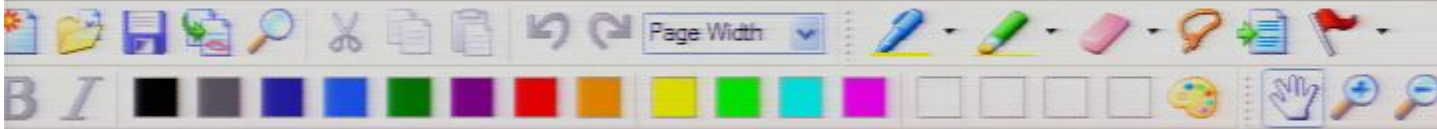
- While galaxies are causing some curvature locally, the universe is found to be spatially very flat on larger scales. (E.g., for triangles of billions of light years in size, Pythagoras' theorem holds with very good precision, the angles add up to 180° , etc)
- Thus, one often considers the simplifying approximation that spacetime has no spatial curvature at all.



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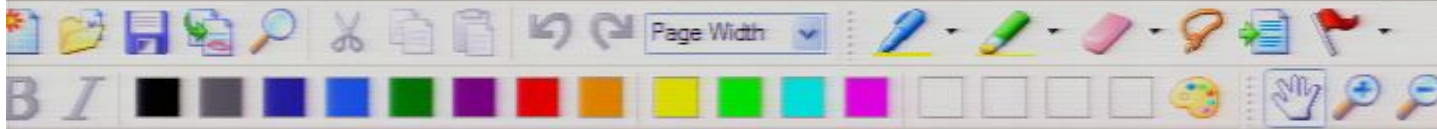
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- While galaxies are causing some curvature locally, the universe is found to be spatially very flat on larger scales. (E.g., for triangles of billions of light years in size, Pythagoras' theorem holds with very good precision, the angles add up to 180° , etc)
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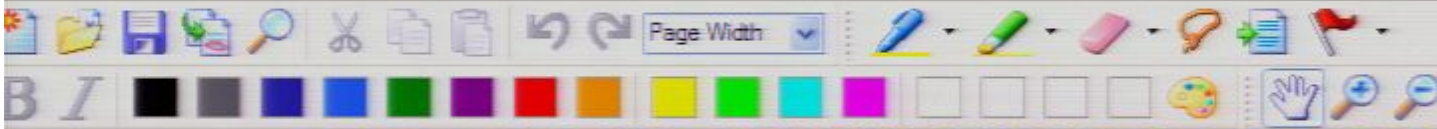
- This leaves the possibility that spacetime may contract or to expand, and potentially differently in different directions.
- The experimental evidence so far supports the simplifying assumption that the cosmic expansion is isotropic.

Remark: It is known that the Einstein equations allow for highly nontrivial evolutions of non-isotropic spacetimes, see, e.g., the text by Wainwright & Ellis.

- With these assumptions, let us choose these coordinates:

* Time:

- Galaxies mostly move away from another according to the Hubble flow.



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general spacetime expansion

o The peculiar velocity is the "small" extra random velocity that galaxies can possess relative to the general Hubble flow.

→ As the time coordinate, t , let us use the proper time, t , of a freely streaming observer who has no peculiar velocity.

(to a good approximation, you can use your wrist watch on earth)



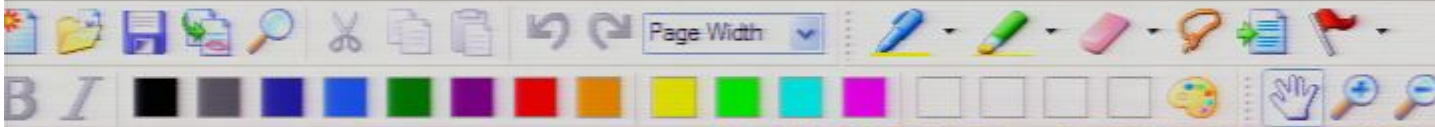
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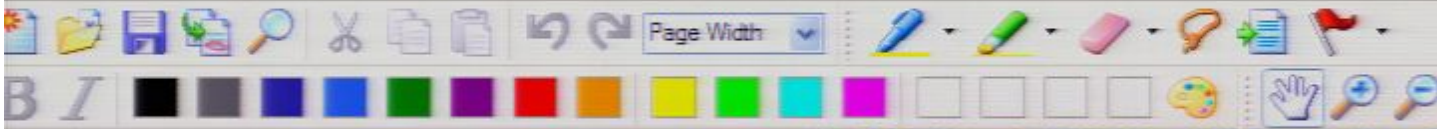
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* Space:

o It is convenient to use "comoving coordinates" x, y, z .

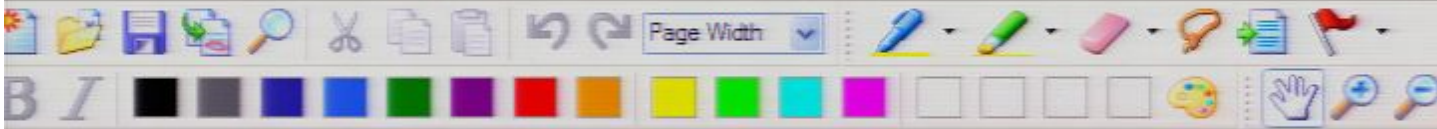


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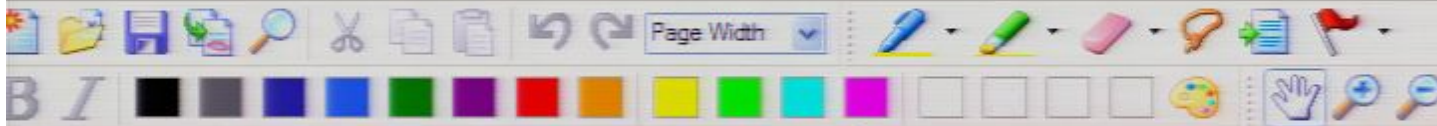
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- o It is convenient to use "comoving coordinates", x_1, x_2, x_3 :
- o At one time, t_0 , (say today) we measure the distances (say between all galaxies) and record those distances.
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- o If we let our spatial coordinate systems



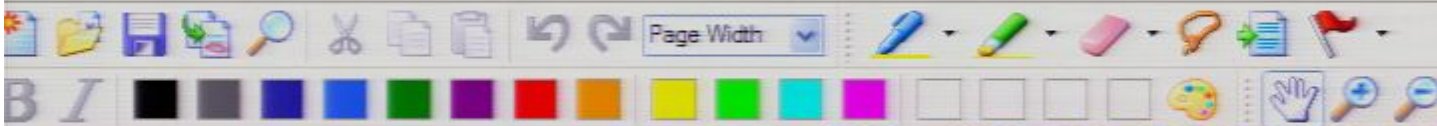
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□ The metric: In these coordinates, $g_{\mu\nu}(x)$ must read:

because we use what's called "proper" time

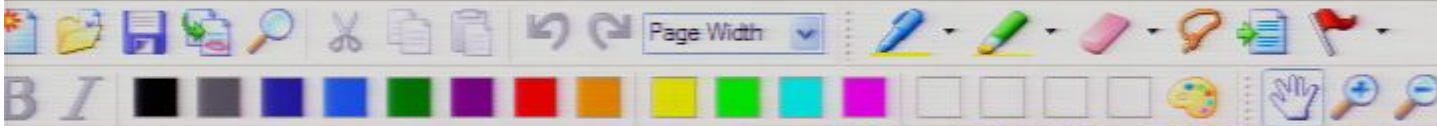
$$g_{\mu\nu}(t, \vec{x}) = \begin{pmatrix} 1 & & & \\ & -a^2(t) & & \\ & & -a^2(t) & \\ & & & -a^2(t) \end{pmatrix}$$

because our coordinate system's unit of length means over time a larger and larger proper length.

Recall: $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$
is "proper" distance.

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o The scale factor function $a(t)$ is



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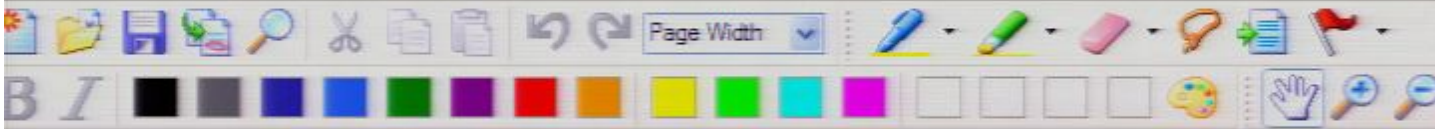
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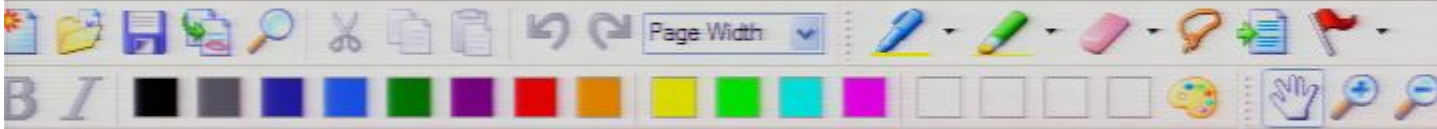
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- The scale factor function $a(t)$ is needed to take into account the expansion when calculating distances.

- Example: The proper distance d between two galaxies with comoving distance $(\Delta x_1, \Delta x_2, \Delta x_3)$ at proper time t_0 is:

$$d = \sqrt{|g_{\mu\nu}(t_0) \Delta x^\mu \Delta x^\nu|}$$

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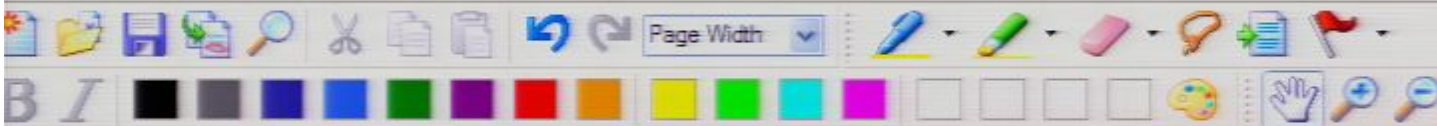


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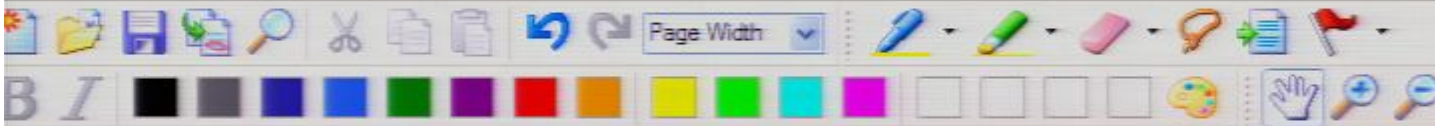
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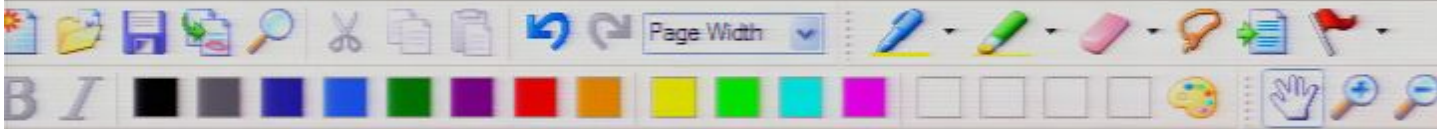
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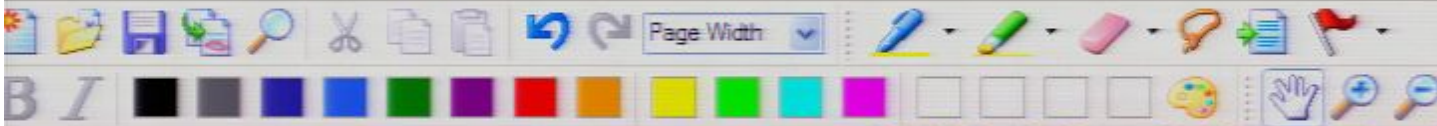
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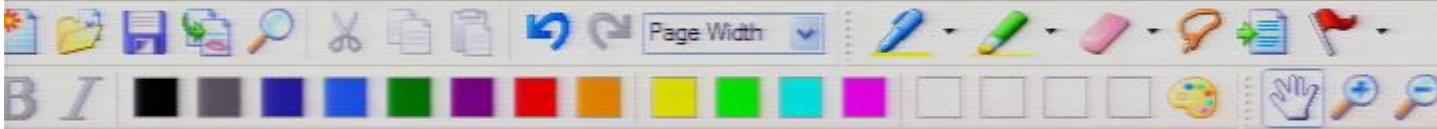
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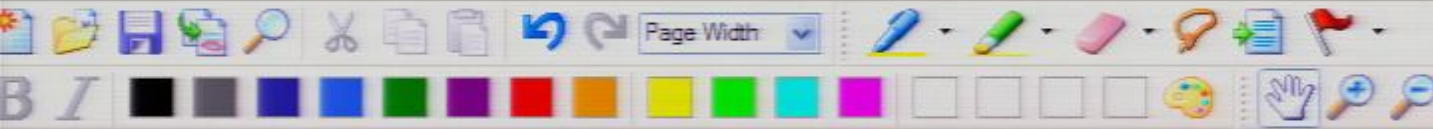
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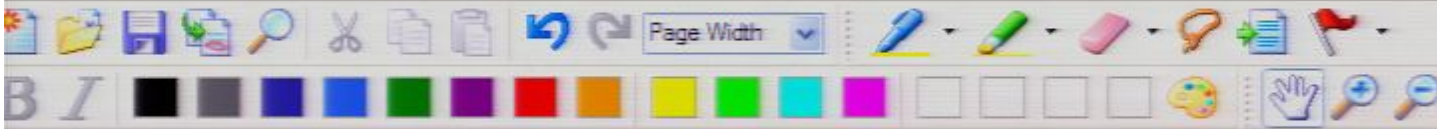
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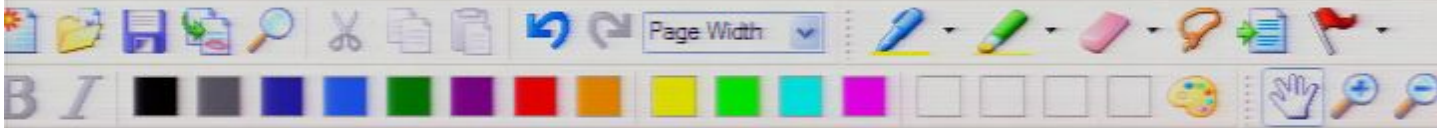
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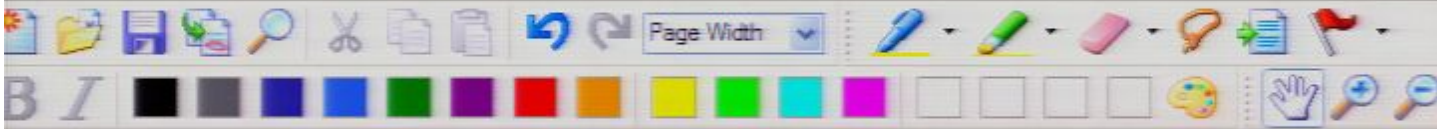
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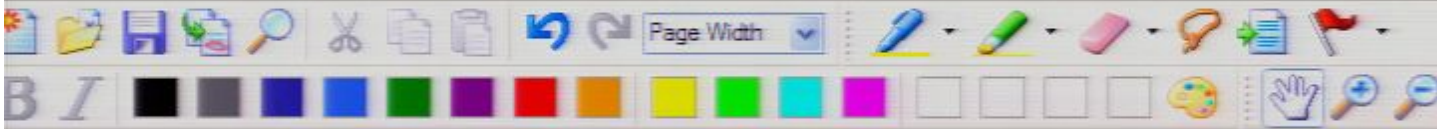
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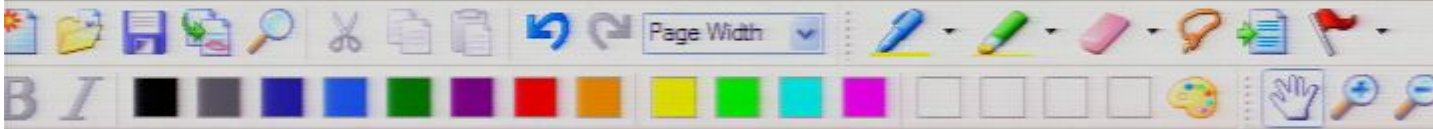
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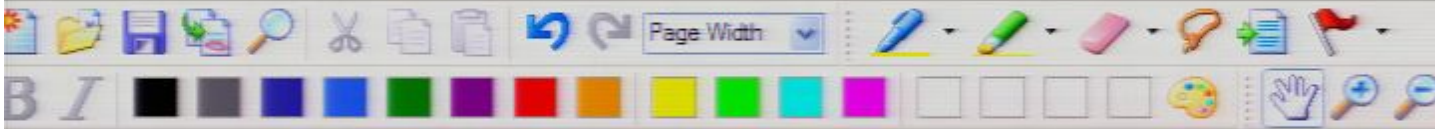
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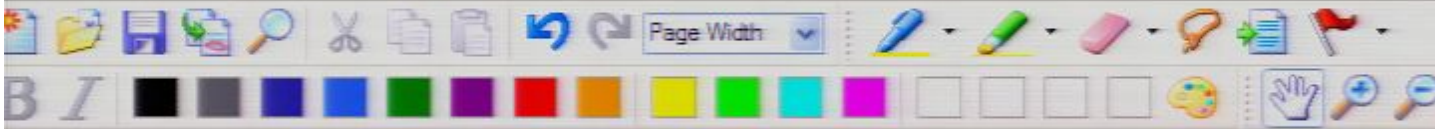
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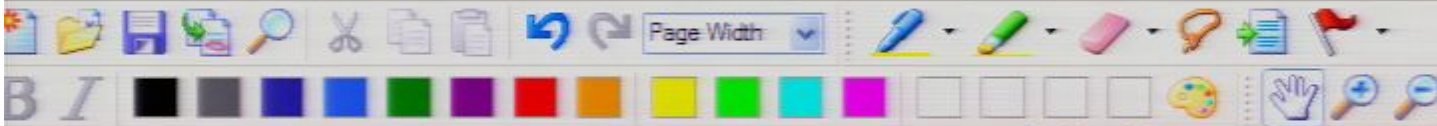
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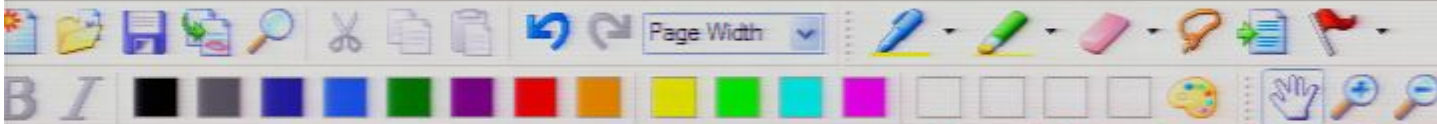
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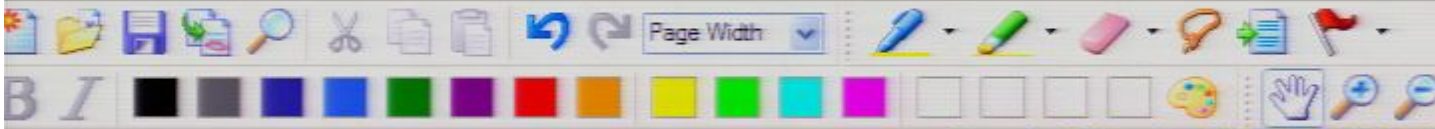
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while setting $g_{\mu\nu}(t, x) = \begin{pmatrix} 1 & & & \\ & -a^2 & & \\ & & -a^2 & \\ & & & 1 \end{pmatrix}$:

$$R_{\mu\nu}(x) - \frac{1}{2} g_{\mu\nu}(x) R(x) + \Lambda g_{\mu\nu}(x) = 8\pi G T_{\mu\nu}(x)$$

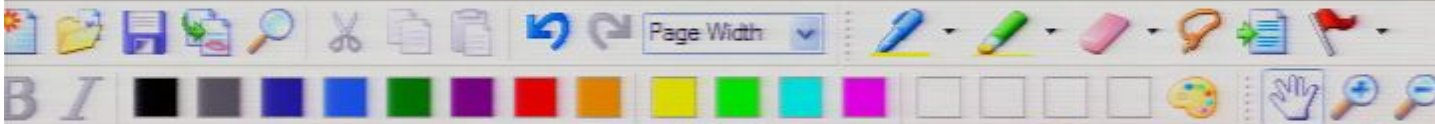
o We can do this classically, but not quantum mechanically:

* Can quantize only \mathcal{L}_i , not $g_{\mu\nu}$. So:
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"Semiclassical approximation" \rightarrow One uses $\bar{T}_{\mu\nu}(x) := \langle \Omega | T_{\mu\nu}(t, \vec{x}) | \Omega \rangle$.

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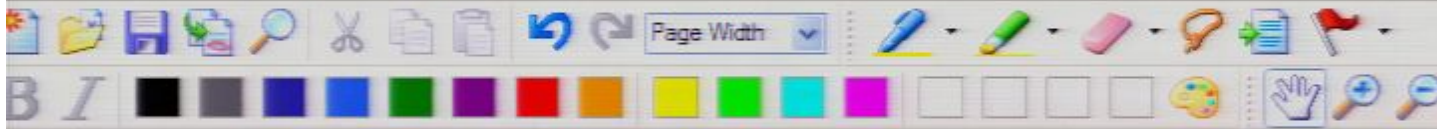
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Definition: The conformal time coordinate.



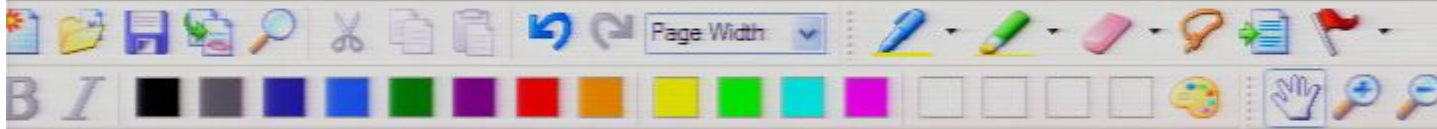
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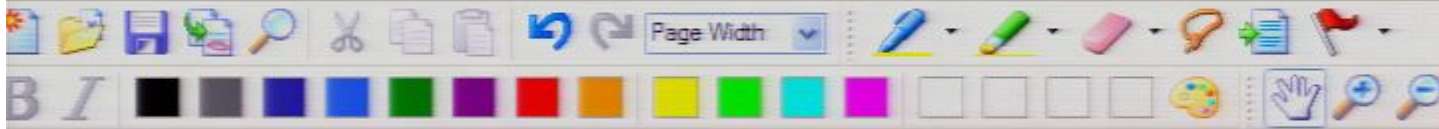
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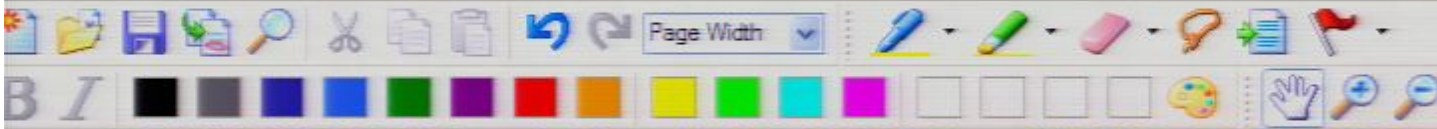
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(...because it shows that the FRW spacetime is equivalent to Minkowski space up to time-dependent conformal, i.e., angle-preserving, i.e. scale-factor-only transformations)

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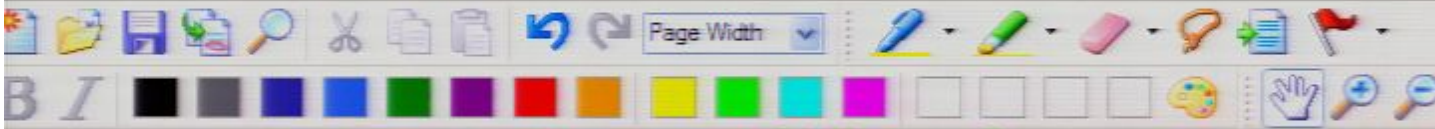
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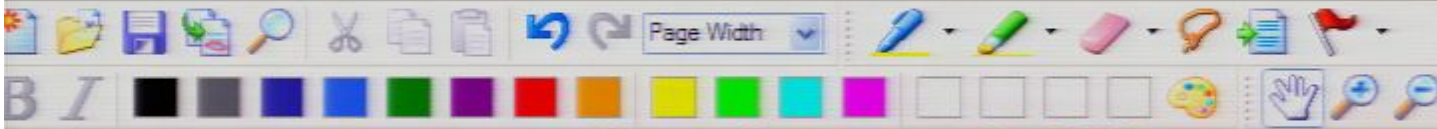
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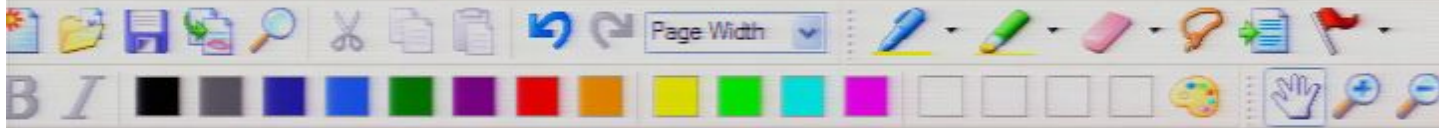
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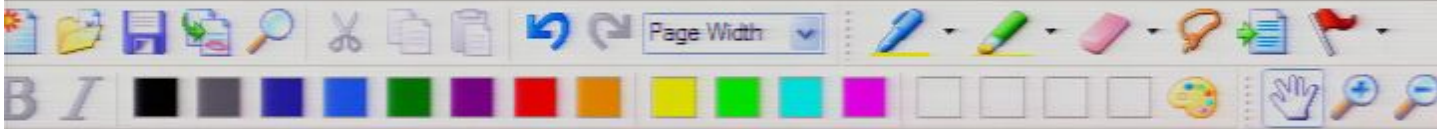
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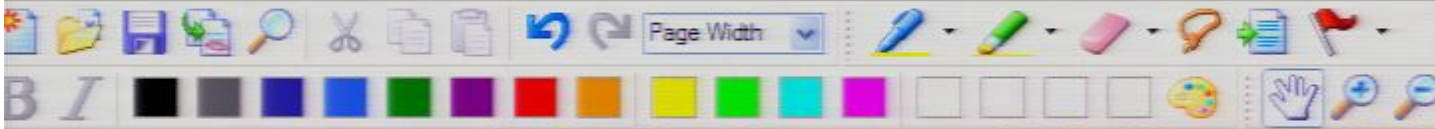
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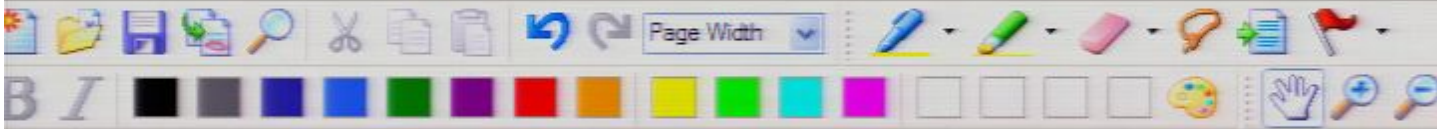
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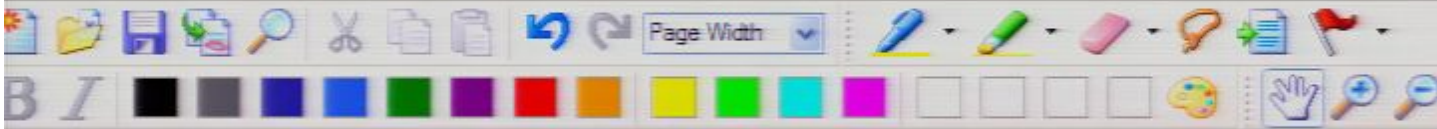


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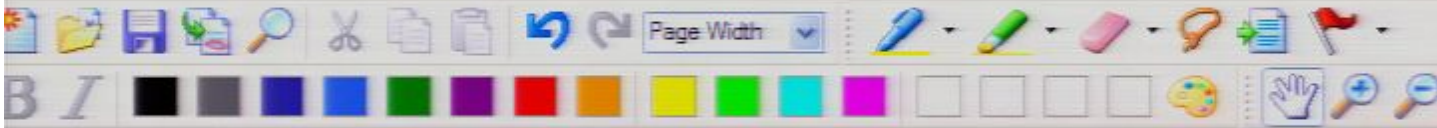
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The Klein Gordon field in FRW spacetimes

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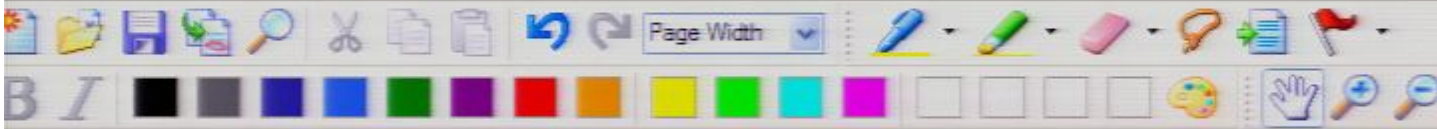
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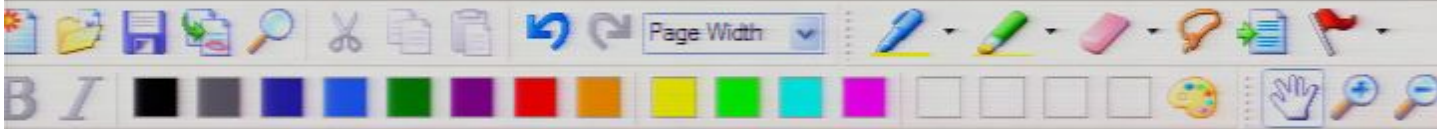
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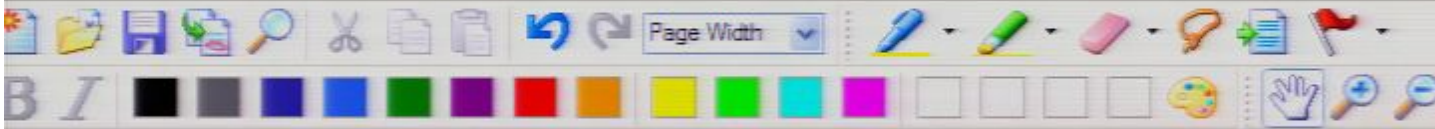
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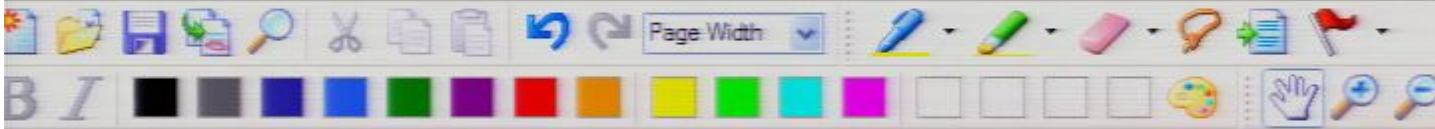
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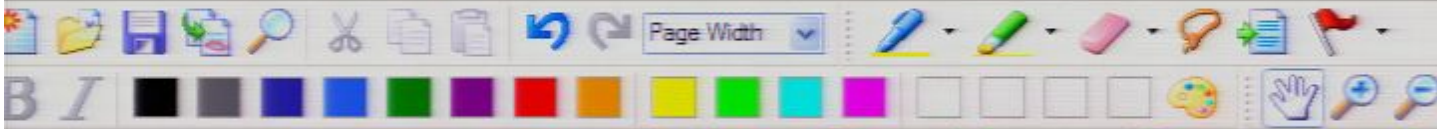
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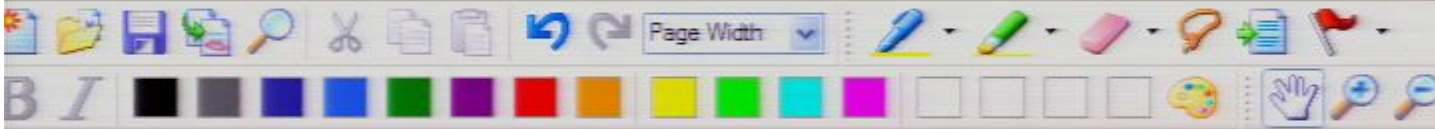
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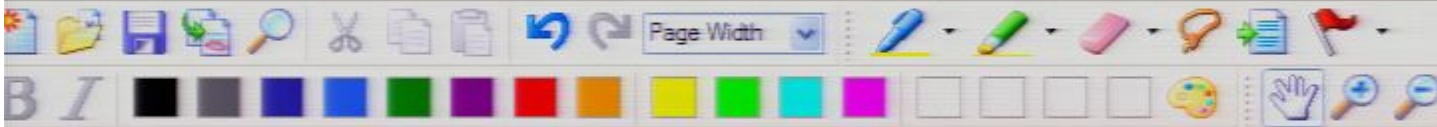
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a term that also occurs in the usual harmonic oscillator. Notice though that it is now time-dependent.

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We try to change from $\phi(\eta, \vec{x})$ to a new field variable, say $\chi(\eta, \vec{x})$, so that the equation of motion for χ has no "friction"-type term.



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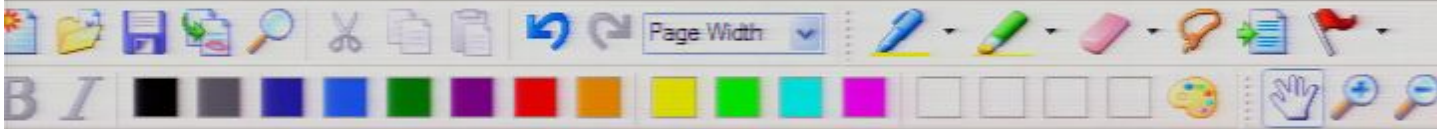
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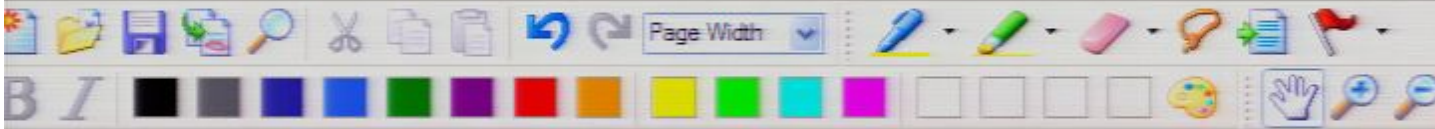
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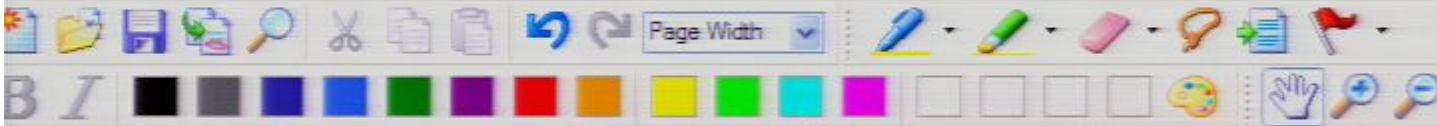
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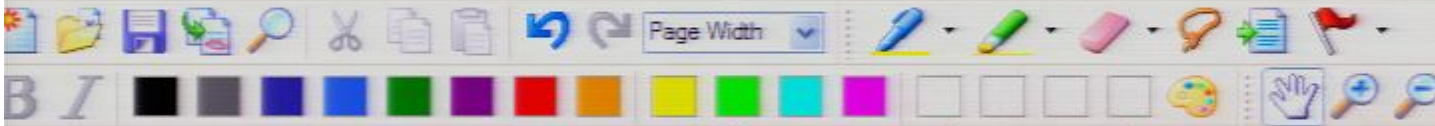
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$$\phi'' + \cancel{\gamma} \phi' + \cancel{\omega} \phi = 0$$


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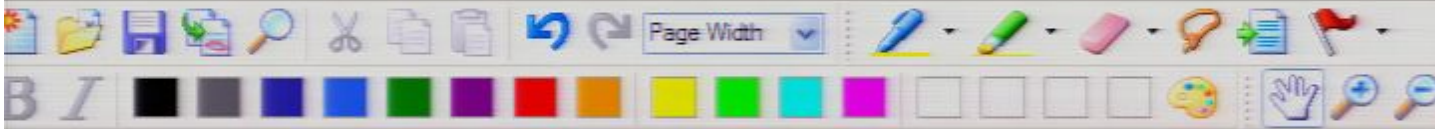
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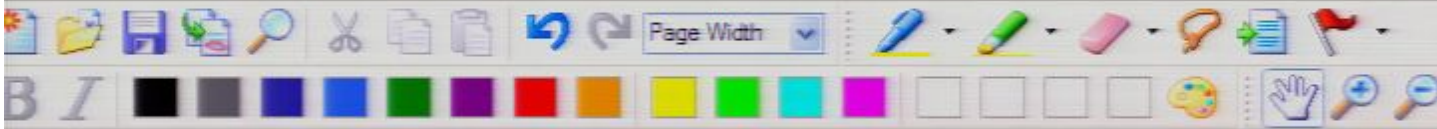
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and: $\phi_{,i} = \frac{\partial}{\partial x^i} \frac{1}{a} \chi(\eta, \vec{x}) = \frac{1}{a} \chi_{,i}$ for $i=1,2,3$

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$$S_{KG} = \int \frac{1}{2} \left(\chi'^2 - \sum_{i=1}^3 \chi_{,i}^2 - \underbrace{\left(m^2 a^2 - \frac{a''}{a} \right)} \chi^2 \right) d\eta d^3x$$



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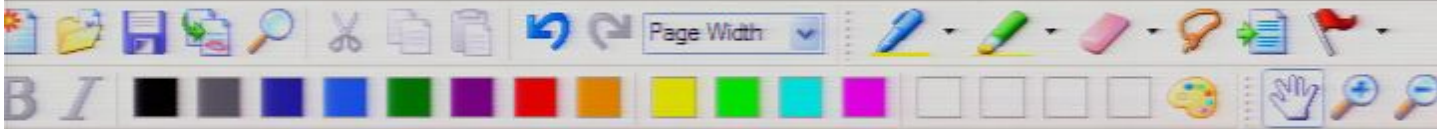
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Exercise: verify



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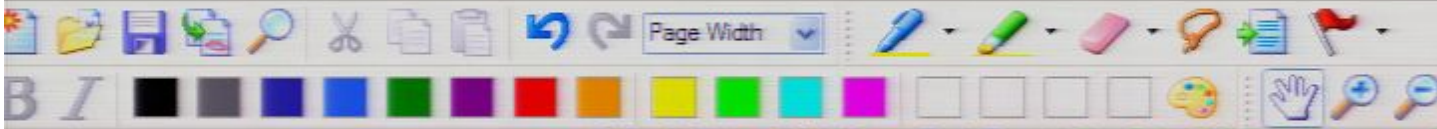
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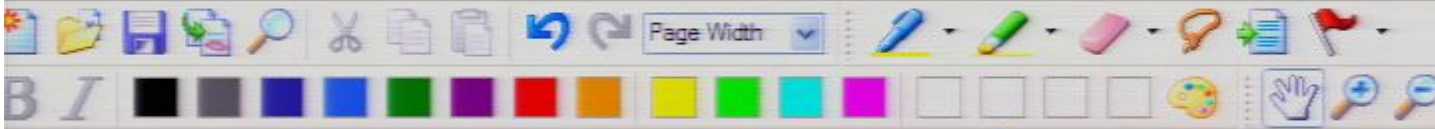
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□ Strategy:

We try to change from $\phi(\eta, \vec{x})$ to a new field
variable say $\chi(\eta, \vec{x})$ so that the equation of



$$\phi''(\eta, \vec{x}) + 2 \frac{a'(\eta)}{a(\eta)} \phi'(\eta, \vec{x}) - \Delta \phi(\eta, \vec{x}) + a^2(\eta) m^2 \phi(\eta, \vec{x}) = 0$$

Here, we use the notation: $f' := \frac{\partial f}{\partial \eta}$

Exercise: verify.

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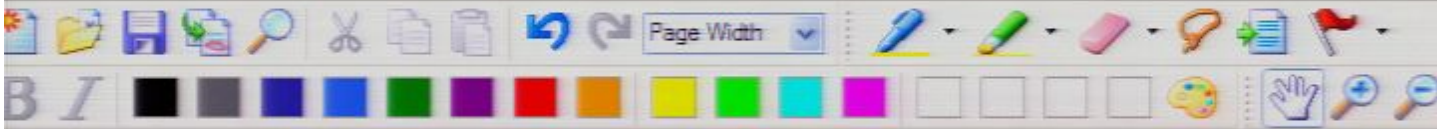
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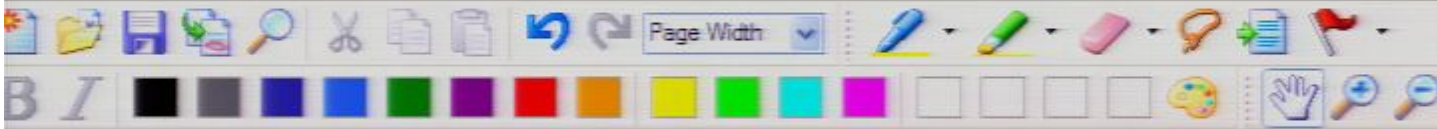
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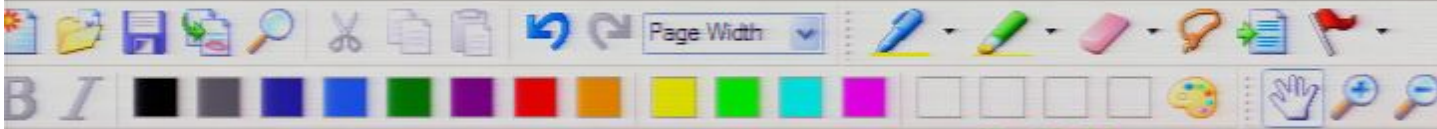
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Note that this term is like a time-dependent mass term $m_{\text{eff}}^2(\eta)$

Exercise: verify

□ Equation of motion:



* Using these, the action in terms of \mathcal{X} reads:

$$S_{\text{KG}} = \int \frac{1}{2} \left(\dot{\mathcal{X}}^2 - \sum_{i=1}^3 \mathcal{X}_{,i}^2 - \underbrace{\left(m^2 a^2 - \frac{a''}{a} \right)}_{\text{mass term } m_{\text{eff}}^2(\eta)} \mathcal{X}^2 \right) d\eta d^3x$$

Note that this term is like a time-dependent mass term $m_{\text{eff}}^2(\eta)$

Exercise: verify

□ Equation of motion:

* Do

$$\frac{\delta S'}{\delta \phi(\eta, \vec{x})} = 0 \quad \text{and} \quad \frac{\delta S'}{\delta \mathcal{X}(\eta, \vec{x})} = 0$$

would equivalent equations of motion?



* Using these, the action in terms of x reads:

$$S_{\text{KG}} = \int \frac{1}{2} \left(\dot{x}^2 - \sum_{i=1}^3 x_{,i}^2 - \underbrace{\left(m^2 a^2 - \frac{a''}{a} \right)}_{\text{mass term } m_{\text{eff}}^2(\eta)} x^2 \right) d\eta d^3x$$

Note that this term is like a time-dependent mass term $m_{\text{eff}}^2(\eta)$

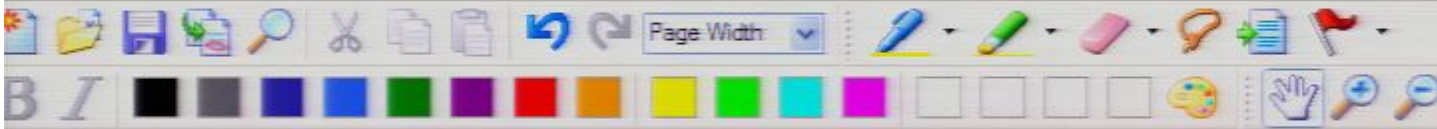
Exercise: verify

□ Equation of motion:

* Do

$$\frac{\delta S}{\delta \phi(\eta, \vec{x})} = 0 \quad \text{and} \quad \frac{\delta S}{\delta x(\eta, \vec{x})} = 0$$

yield equivalent equations of motion?



□ Equation of motion:

* Do

$$\frac{\delta S'}{\delta \phi(\gamma, \bar{x})} = 0 \quad \text{and} \quad \frac{\delta S'}{\delta x(\gamma, \bar{x})} = 0$$

yield equivalent equations of motion?

* Yes, because:

$$0 = \frac{\delta S'}{\delta \phi} = \frac{\delta S'}{\delta x} \frac{\delta x}{\delta \phi}$$

↳ if $\delta S'/\delta x$ vanishes then also $\delta S'/\delta \phi$ vanishes.

* Thus, we may calculate the equation of motion



$$S_{\text{KG}} = \int \frac{1}{2} \left(\dot{x}^2 - \sum_{i=1}^3 x_{,i}^2 - \underbrace{\left(m^2 a^2 - \frac{a''}{a} \right)}_{\text{mass term } m_{\text{eff}}^2(\eta)} x^2 \right) d\eta d^3x$$

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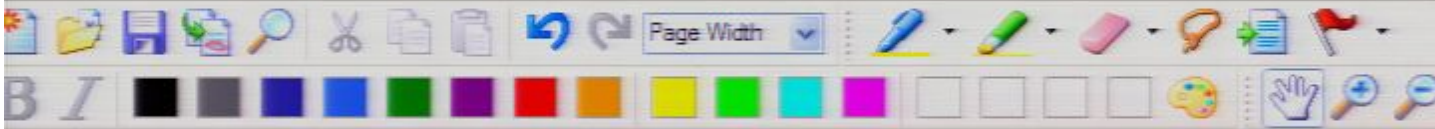
Exercise: verify

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$$S_{\text{KG}} = \int \frac{1}{2} \left(\dot{x}^{\prime 2} - \sum_{i=1}^3 x_{,i}^2 - \underbrace{\left(m^2 a^2 - \frac{a''}{a} \right)}_{\text{like a time-dependent mass term } m_{\text{eff}}^2(\eta)} x^2 \right) d\eta d^3x$$

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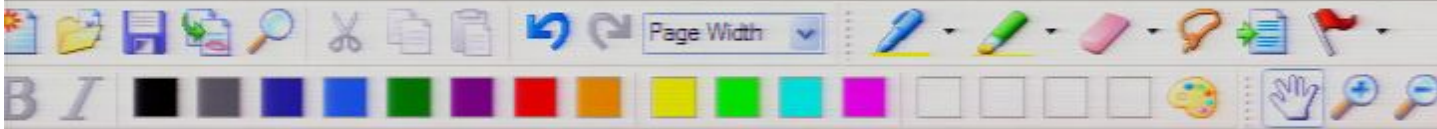
* Do

$$\frac{\delta S'}{\delta \phi(\eta, \vec{x})} = 0 \quad \text{and} \quad \frac{\delta S'}{\delta X(\eta, \vec{x})} = 0$$

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* Yes, because:



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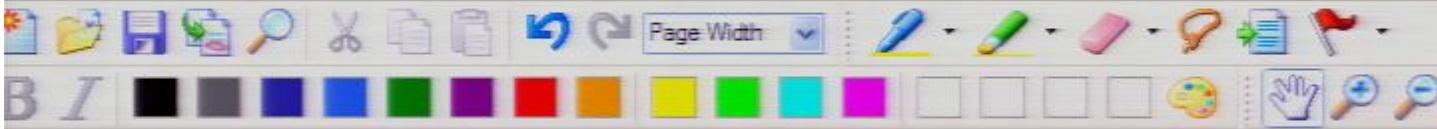
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* Thus, we may calculate the equation of motion directly in terms of x from $S[x]$, to obtain:

$$x'' = \Delta x + (m^2 - a'')x = 0 \quad (\text{Eq. 14/23})$$



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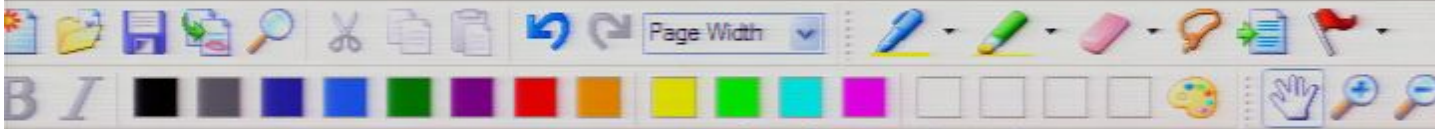
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$$x'' - \Delta x + \left(m^2 a^2 - \frac{a''}{a}\right)x = 0 \quad (\text{EOM!})$$



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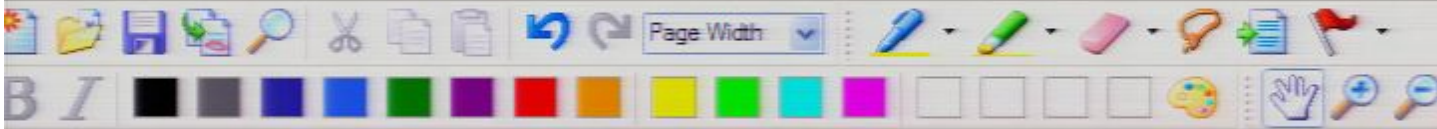
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Remark:



We could have obtained this equation of motion directly from that of ϕ by change of variable. But finding the action for x was still worthwhile, namely to get the conjugate to x !



* Thus, we may calculate the equation of motion directly in terms of χ from $S[\chi]$, to obtain:

Exercise: verify!

$$\chi'' - \Delta \chi + \left(m^2 a^2 - \frac{a''}{a}\right) \chi = 0 \quad (\text{EOM!})$$

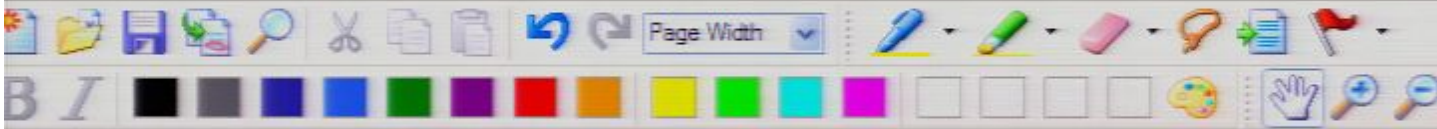
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□ Preparation for quantization:

* We need the canonically conjugate field



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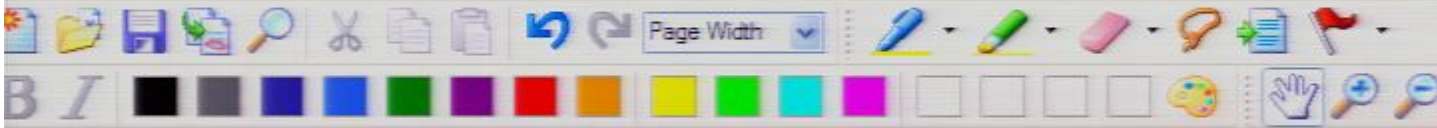
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Preparation for quantization:

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$$\pi^{(\chi)}(\eta, \vec{x})$$



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to the field $\mathcal{X}(\eta, \vec{x})$, i.e., the Legendre transform of \mathcal{X} :

* To this end, we consider the Lagrangian:

$$L = \int d^3x \left(\frac{1}{2} (\dot{\mathcal{X}}^2 - \sum_{i=1}^3 \mathcal{X}_{,i}^2 - (m^2 a^2 - \frac{a''}{a}) \mathcal{X}) \right)$$

* Thus, the Legendre transformed variable reads:

$$\pi^{(x)}(\eta, \vec{x}) := \frac{\delta L}{\delta \dot{\mathcal{X}}(\eta, \vec{x})} = \dot{\mathcal{X}}(\eta, \vec{x})$$

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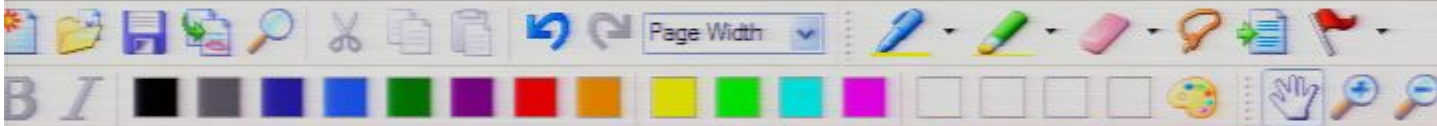
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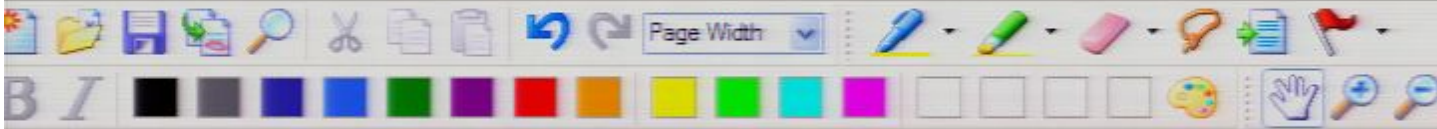
$$\pi^{(x)}(\eta, \vec{x}) := \frac{\delta L}{\delta x'(\eta, \vec{x})} = x'(\eta, \vec{x}) \quad (\text{EOM 2})$$

* Which is the field that is conjugate to ϕ ?

$$S_{x.o.} = \int \left(\frac{1}{2} a^{-2}(\eta) \eta^{\alpha\beta} \phi_{,\alpha} \phi_{,\beta} - \frac{1}{2} m^2 \phi^2 \right) a^4 d\eta d^3x$$

\Rightarrow The field $\pi^{(\phi)}$ which is conjugate to ϕ reads:

$$\pi^{(\phi)} := \frac{\delta L}{\delta \phi'} = a^2 \phi'$$



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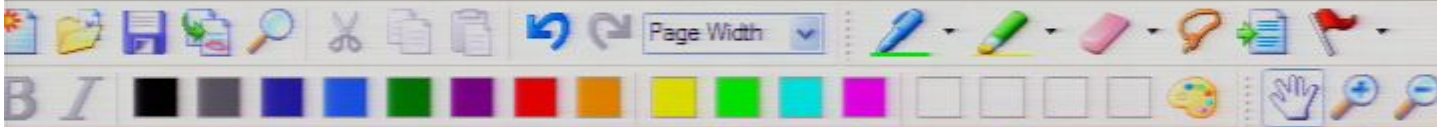
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* Compare:



$$\pi^{(x)}(\eta, \vec{x}) := \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \dot{\phi} \quad (\text{EOM 2})$$

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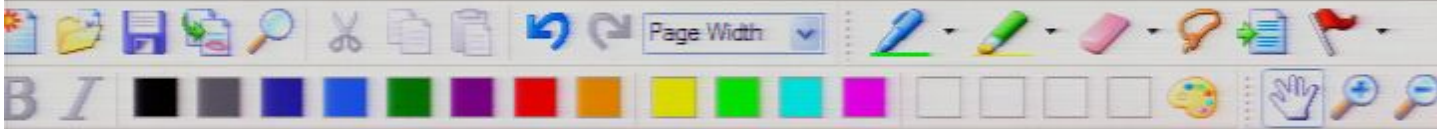
$$\mathcal{L}_{x.6} = \int \left(\frac{1}{2} \bar{a}^{-2}(\eta) \eta^{\alpha\beta} \phi_{,\alpha} \phi_{,\beta} - \frac{1}{2} m^2 \phi^2 \right) a^4 d\eta d^3x$$

\Rightarrow The field $\pi^{(\phi)}$ which is conjugate to ϕ reads:

$$\pi^{(\phi)} := \frac{\delta \mathcal{L}}{\delta \dot{\phi}} = \bar{a}^{-2} \dot{\phi}$$

* Compare:

$$\begin{aligned} \pi^{(x)} &= \dot{\phi} \\ &= (a\phi)' \end{aligned}$$



* Which is the field that is conjugate to ϕ ?

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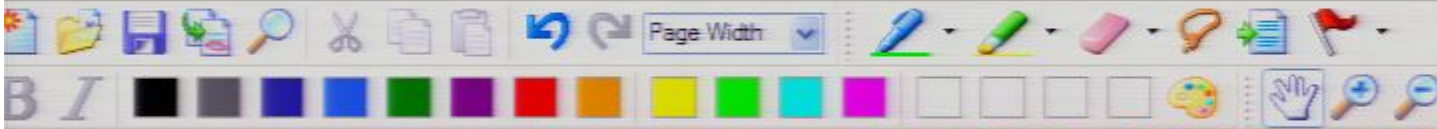
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* Compare:

$$\begin{aligned} \pi^{(x)} &= \mathcal{X}' \\ &= (a\phi)' \\ &= a\phi' + a'\phi \\ &= \frac{1}{a} \pi^{(\phi)} + a'\phi \end{aligned}$$

, i.e., $\pi^{(\phi)}$, $\pi^{(x)}$ are different!



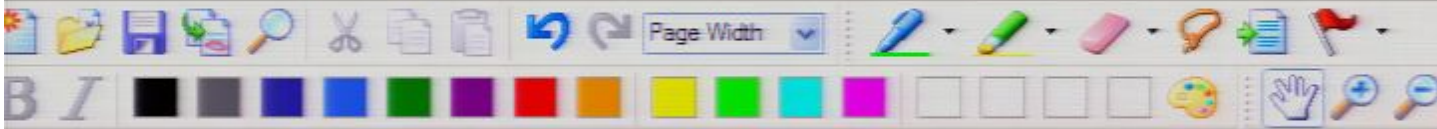
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Quantization:



⇒ The field $\pi^{(\phi)}$ which is conjugate to ϕ reads:

$$\pi^{(\phi)} := \frac{\delta \mathcal{L}}{\delta \phi'} = a^2 \phi'$$

* Compare:

$$\pi^{(\chi)} = \chi'$$

$$= (a\phi)'$$

$$= a\phi' + a'\phi$$

$$= \frac{1}{a} \pi^{(\phi)} + a'\phi, \text{ i.e., } \pi^{(\phi)}, \pi^{(\chi)} \text{ are different!}$$

□ Quantization:

$$[\phi(\eta, \vec{x}), \pi^{(\phi)}(\eta, \vec{x}')] = i \delta^3(\vec{x} - \vec{x}')$$



$= \frac{1}{a} \pi + a \psi$, i.e., π^+ , π^- are different!

□ Quantization:

$$[\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = i \delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\phi}(\eta, \vec{x}), \hat{\phi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(\phi)}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = 0$$

□ Proposition:

□ In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{x}'$, these commutation relations become:

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□ Proposition:

In terms of the fields $\hat{\chi} := a\hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}^{(\phi)}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(\chi)}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = 0$$



$$[\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(b)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\phi}(\eta, \vec{x}), \hat{\phi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(b)}(\eta, \vec{x}), \hat{\pi}^{(b)}(\eta, \vec{x}')] = 0$$

□ Proposition:

In terms of the fields $\hat{\chi} := a\hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(\chi)}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = 0$$

□ Proof. Only the first CCR is non-trivial to check

Quantization:

$$[\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\phi}(\eta, \vec{x}), \hat{\phi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(\phi)}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = 0$$

Proposition:

In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}^{(\phi)}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

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Proposition:

In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

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Quantization:

$$[\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(b)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\phi}(\eta, \vec{x}), \hat{\phi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(b)}(\eta, \vec{x}), \hat{\pi}^{(b)}(\eta, \vec{x}')] = 0$$

Proposition:

In terms of the fields $\hat{\chi} := a\hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

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$$[\hat{\pi}^{(\phi)}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = 0$$

□ Proposition:

In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{x}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

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$$[\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

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$$[\hat{\pi}^{(\phi)}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] = 0$$

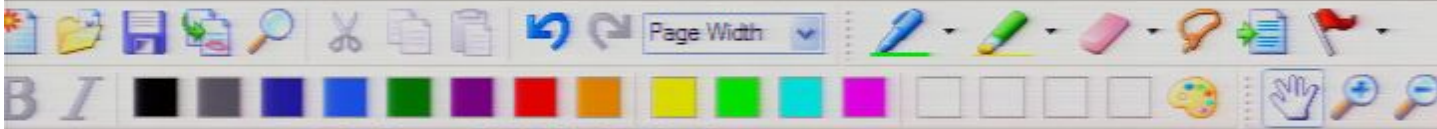
□ Proposition:

In terms of the fields $\hat{\chi} := a\hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\pi}^{(\phi)}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

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In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\chi}'$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

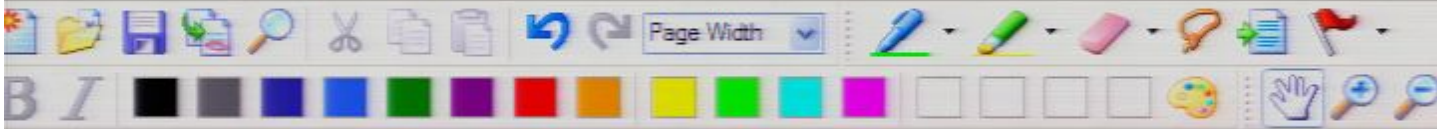
$$[\hat{\pi}^{(\chi)}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = 0$$

□ Proof Only the first CCR is nontrivial to check:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = [a(\eta) \hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)} \hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta) \hat{\phi}(\eta, \vec{x}')] =$$

$$= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] =$$

$$= i\delta^3(\vec{x} - \vec{x}')$$



In terms of the fields $\hat{\chi} := a \hat{\phi}$, $\hat{\pi}^{(\chi)} := \hat{\chi}'$, these commutation relations become:

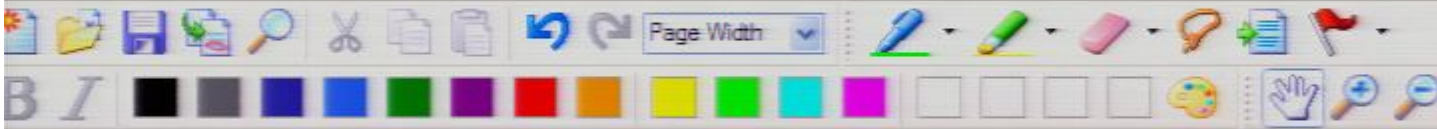
$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

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$$\begin{aligned} [\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(\chi)}(\eta, \vec{x}')] &= [a(\eta) \hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)} \hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta) \hat{\phi}(\eta, \vec{x}')] \\ &= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] \\ &= i\delta^3(\vec{x} - \vec{x}') \end{aligned}$$



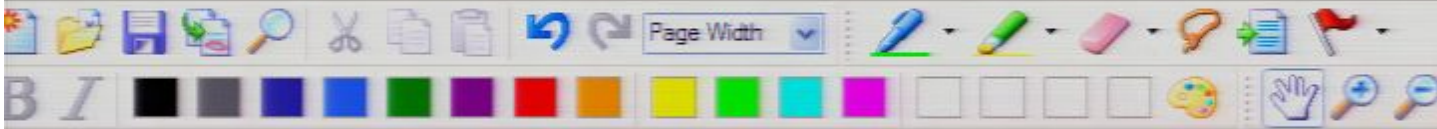
□ Proof Only the first CCR is nontrivial to check:

$$\begin{aligned} [\hat{x}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] &= [a(\eta) \hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)} \hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta) \hat{\phi}(\eta, \vec{x}')] \\ &= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] \\ &= i \delta^3(\vec{x} - \vec{x}') \end{aligned}$$

□ Thus, the change from ϕ to x is fairly trivial.

Notice, however:

$$\begin{array}{l} L \xrightarrow{\text{L.T. } \phi \text{ replaced by } \pi^{\phi}} H^{(\phi)} := \int \phi' \pi^{(\phi)} d^3x - L \\ \quad \searrow \text{L.T. } x' \text{ replaced by } \pi^x \\ \quad \quad H^{(\pi)} := \int x' \pi^{(x)} d^3x - L \end{array} \left. \begin{array}{l} \leftarrow \\ \leftarrow \end{array} \right\} \begin{array}{l} \text{they have no reason} \\ \text{to be the same!} \end{array}$$



$$[\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\mathcal{X}}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(0)}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] = 0$$

□ Proof Only the first CCR is nontrivial to check:

$$\begin{aligned} [\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] &= [a(\eta)\hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)}\hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta)\hat{\phi}(\eta, \vec{x}')] \\ &= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] \\ &= i\delta^3(\vec{x} - \vec{x}') \end{aligned}$$

□ Thus, the change from ϕ to \mathcal{X} is fairly trivial.

Notice, however:



$$[\hat{x}(q, \vec{x}), \hat{\pi}^{(0)}(q, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{x}(q, \vec{x}), \hat{x}(q, \vec{x}')] = 0$$

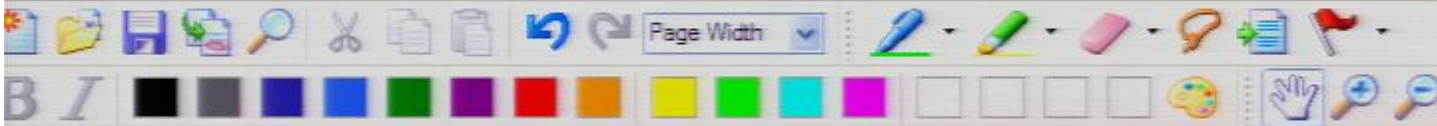
$$[\hat{\pi}^{(0)}(q, \vec{x}), \hat{\pi}^{(0)}(q, \vec{x}')] = 0$$

□ Proof Only the first CCR is nontrivial to check:

$$\begin{aligned} [\hat{x}(q, \vec{x}), \hat{\pi}^{(0)}(q, \vec{x}')] &= [a(q)\hat{\phi}(q, \vec{x}), \frac{1}{a(q)}\hat{\pi}^{(\phi)}(q, \vec{x}') + a'(q)\hat{\phi}(q, \vec{x}')] \\ &= [\hat{\phi}(q, \vec{x}), \hat{\pi}^{(\phi)}(q, \vec{x}')] \\ &= i\delta^3(\vec{x} - \vec{x}') \end{aligned}$$

□ Thus, the change from ϕ to x is fairly trivial.

Notice, however:



$$[\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

$$[\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\mathcal{X}}(\eta, \vec{x}')] = 0$$

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$$\begin{aligned} [\hat{\mathcal{X}}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] &= [a(\eta) \hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)} \hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta) \hat{\phi}(\eta, \vec{x}')] \\ &= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] \\ &= i\delta^3(\vec{x} - \vec{x}') \end{aligned}$$

□ Thus, the change from ϕ to \mathcal{X} is fairly trivial.

Notice, however:



In terms of the fields $\chi := a\phi$, $\pi := \dot{\chi}$, these commutation relations become:

$$[\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] = i\delta^3(\vec{x} - \vec{x}')$$

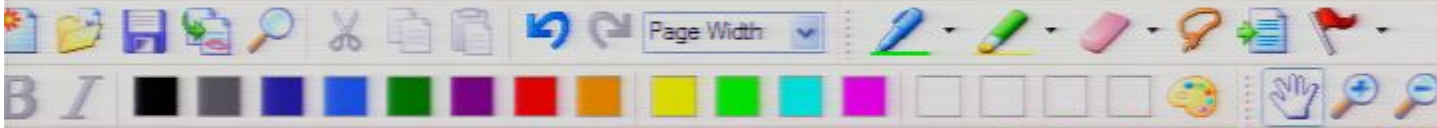
$$[\hat{\chi}(\eta, \vec{x}), \hat{\chi}(\eta, \vec{x}')] = 0$$

$$[\hat{\pi}^{(0)}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] = 0$$

□ Proof Only the first CCR is nontrivial to check:

$$\begin{aligned} [\hat{\chi}(\eta, \vec{x}), \hat{\pi}^{(0)}(\eta, \vec{x}')] &= [a(\eta)\hat{\phi}(\eta, \vec{x}), \frac{1}{a(\eta)}\hat{\pi}^{(\phi)}(\eta, \vec{x}') + a'(\eta)\hat{\phi}(\eta, \vec{x}')] \\ &= [\hat{\phi}(\eta, \vec{x}), \hat{\pi}^{(\phi)}(\eta, \vec{x}')] \\ &= i\delta^3(\vec{x} - \vec{x}') \end{aligned}$$

□ Thus, the change from ϕ to χ is fairly trivial.



$$= i\delta^3(\vec{x} - \vec{x}')$$

□ Thus, the change from ϕ to χ is fairly trivial.

Notice, however:

$$L \xrightarrow{\text{L.T. } \phi \text{ replaced by } \pi^\dagger} H^{(\phi)} := \int \phi' \pi^{(\phi)} d^3x - L \quad \left. \vphantom{H^{(\phi)}} \right\} \leftarrow$$

$$L \xrightarrow{\text{L.T. } \chi \text{ replaced by } \pi^\dagger} H^{(\pi)} := \int \chi' \pi^{(\chi)} d^3x - L \quad \left. \vphantom{H^{(\pi)}} \right\} \swarrow$$

they have no reason
to be the same!

□ Question:

How can both be valid generators of time evolution,



□ Thus, the change from ϕ to x is fairly trivial.

Notice, however:

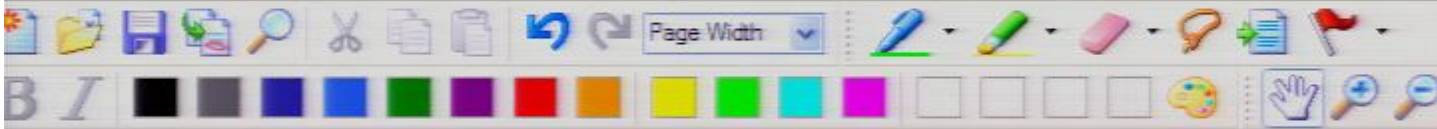
$$L \begin{cases} \xrightarrow{\text{L.T. } \phi \text{ replaced by } \pi^\dagger} H^{(\phi)} := \int \phi' \pi^{(\phi)} d^3x - L \\ \searrow \text{L.T. } x' \text{ replaced by } \pi^{xc} H^{(\pi)} := \int x' \pi^{(xc)} d^3x - L \end{cases} \left\{ \begin{array}{l} \leftarrow \\ \downarrow \end{array} \right.$$

they have no reason to be the same!

□ Question:

How can both be valid generators of time evolution,

i.e., how can we have:



Question:

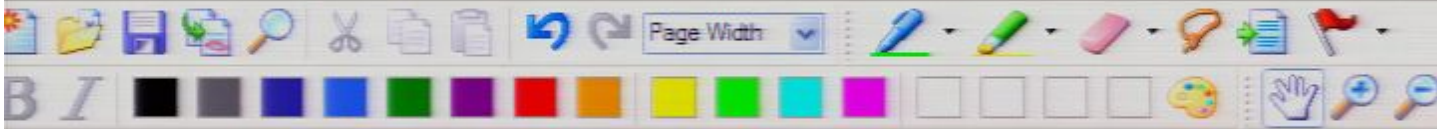
How can both be valid generators of time evolution,
i.e., how can we have:

$$i\dot{\hat{\phi}}' = [\hat{\phi}', \hat{H}^{(H)}] \quad \text{and} \quad i\dot{\hat{x}}' = [\hat{x}', \hat{H}^{(X)}]$$

and yet $\hat{H}^{(H)} \neq \hat{H}^{(X)}$?

Should there not be one Hamiltonian for all variables?

Answer: Yes, and it is, of course $\hat{H}^{(H)}$.



□ Question:

How can both be valid generators of time evolution,

i.e., how can we have:

$$i\dot{\hat{\phi}}' = [\hat{\phi}', \hat{H}^{(0)}] \quad \text{and} \quad i\dot{\hat{x}}' = [\hat{x}', \hat{H}^{(0)}]$$

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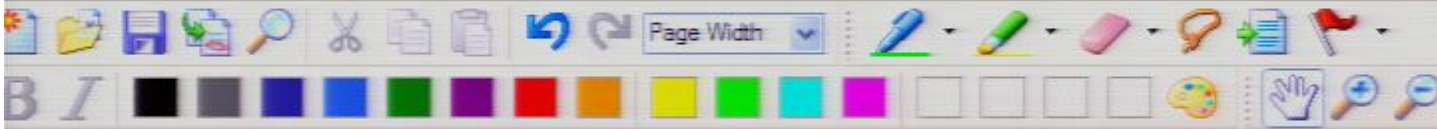
This extra term is there if the variable \hat{Q} has also explicit time-dependence, e.g., $\hat{Q} = \cos(\omega t + \phi) \hat{q} + c \hat{p}$, or here: $\hat{x}' = \frac{1}{\omega} \dot{\hat{\phi}}$.

Recall that in QM: $i\dot{\hat{Q}}' = [\hat{Q}', \hat{H}] + i\frac{\partial}{\partial t} \hat{Q}'$

$$i\dot{f} = \hat{f}, \hat{f}^2$$

$$i\dot{f} = [\hat{f}, \hat{H}]$$





□ Question:

How can both be valid generators of time evolution,

i.e., how can we have:

$$i\dot{\hat{\phi}}' = [\hat{\phi}', \hat{H}^{(0)}] \quad \text{and} \quad i\dot{\hat{x}}' = [\hat{x}', \hat{H}^{(0)}]$$

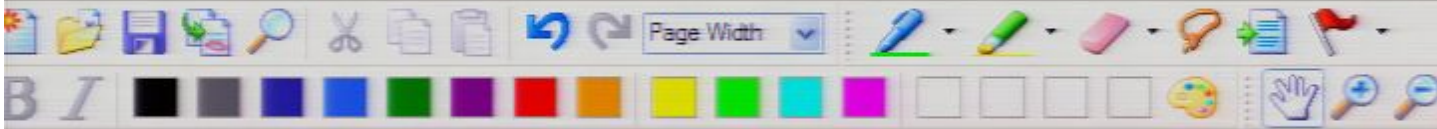
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$$i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(H)}] \quad \text{and} \quad i\hat{x}' = [\hat{x}, \hat{H}^{(H)}]$$

and yet $\hat{H}^{(H)} \neq \hat{H}^{(0)}$?

□ Should there not be one Hamiltonian for all variables?

□ **Answer:** Yes, and it is, of course $\hat{H}^{(H)}$.

This extra term is there if the variable \hat{Q} has also explicit time-dependence, e.g., $\hat{Q} = c\cos(\omega t + k\epsilon)\hat{q} + c\hat{p}$, or here: $\hat{x} = \frac{1}{a}\hat{\phi}$.

Recall that in QM: $i\hat{Q}' = [\hat{Q}, \hat{H}] + i\frac{\partial}{\partial t}\hat{Q}$

□ Explicitly:

* From $\hat{x} = a\hat{\phi}$ and $i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(H)}]$ we obtain:

$$i\dot{f} = [\hat{f}, \hat{H}] + i\partial_t f$$

$$\hat{f} = \hat{x}, \quad \hat{f} = \hat{p}, \quad \hat{f} = \hat{x}$$

$$\hat{f} = \sin(t)\hat{x} + \cos^2(t)\hat{p}$$



$$i\dot{f} = [\hat{f}, \hat{H}] + i\partial_t f$$

$$\hat{f} = \hat{x}, \quad \hat{f} = \hat{p}, \quad \hat{f} = \hat{x} + 7\hat{p}^2$$

$$\hat{f} = \sin(t)\hat{x} + \cos^2(t)\hat{p}^2$$



i.e., how can we have:

$$i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(H)}] \quad \text{and} \quad i\hat{x}' = [\hat{x}, \hat{H}^{(0)}]$$

and yet $\hat{H}^{(H)} \neq \hat{H}^{(0)}$?

□ Should there not be one Hamiltonian for all variables?

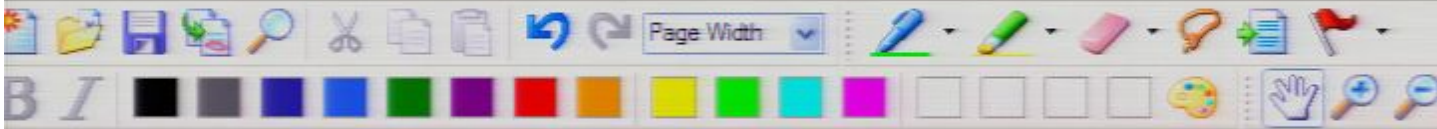
□ **Answer:** Yes, and it is, of course $\hat{H}^{(H)}$.

This extra term is there if the variable \hat{Q} has also explicit time-dependence, e.g., $\hat{Q} = \cos(\omega t + \phi)\hat{q} + c\hat{p}$, or here: $\hat{x} = \frac{1}{\omega}\hat{\phi}$.

Recall that in QM:
$$i\hat{Q}' = [\hat{Q}, \hat{H}] + i\frac{\partial}{\partial t}\hat{Q}$$

□ Explicitly:

* From $\hat{x} = a\hat{\phi}$ and $i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(H)}]$ we obtain:



□ Explicitly:

* From $\hat{x} = a\hat{\phi}$ and $i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(b)}]$ we obtain:

$$i\left(\frac{1}{a}\hat{x}\right)' = \frac{1}{a}[\hat{x}, \hat{H}^{(b)}]$$

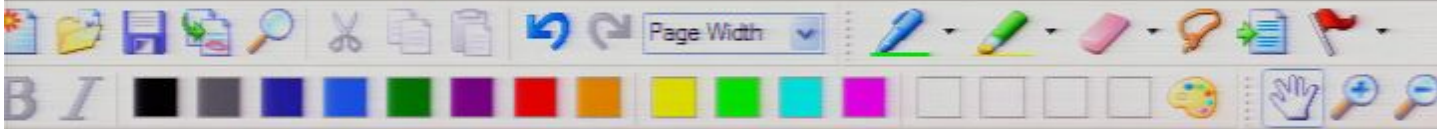
$$\Rightarrow i\frac{1}{a}\hat{x}' - i\frac{a'}{a^2}\hat{x} = \frac{1}{a}[\hat{x}, \hat{H}^{(b)}]$$

$$\Rightarrow i\hat{x}' = [\hat{x}, \hat{H}^{(b)}] + i\frac{a'}{a}\hat{x}$$

* But we also have:

$$i\hat{x}' = [\hat{x}, \hat{H}^{(a)}]$$

\Rightarrow We must have: $\hat{H}^{(a)} \neq \hat{H}^{(b)}$



□ Explicitly:

* From $\hat{x} = a\hat{\phi}$ and $i\hat{\phi}' = [\hat{\phi}, \hat{H}^{(a)}]$ we obtain:

$$i\left(\frac{1}{a}\hat{x}\right)' = \frac{1}{a}[\hat{x}, \hat{H}^{(a)}]$$

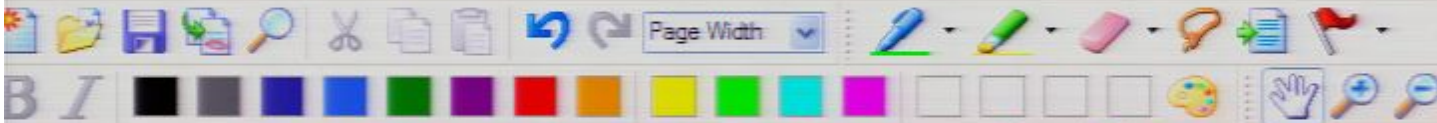
$$\Rightarrow i\frac{1}{a}\hat{x}' - i\frac{a'}{a^2}\hat{x} = \frac{1}{a}[\hat{x}, \hat{H}^{(a)}]$$

$$\Rightarrow i\hat{x}' = [\hat{x}, \hat{H}^{(a)}] + i\frac{a'}{a}\hat{x}$$

* But we also have:

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$$i\left(\frac{1}{a}\hat{x}\right)' = \frac{1}{a} [\hat{x}, \hat{H}^{(a)}]$$

$$\Rightarrow i\frac{1}{a}\hat{x}' - i\frac{a'}{a^2}\hat{x} = \frac{1}{a} [\hat{x}, \hat{H}^{(a)}]$$

$$\Rightarrow i\hat{x}' = [\hat{x}, \hat{H}^{(a)}] + i\frac{a'}{a}\hat{x}$$

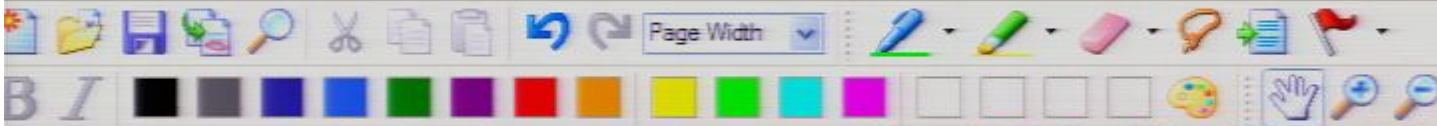
* But we also have:

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Since there are multiple Hamiltonians, which, if anyone, is the energy?

One usually defines the energy as the generator of time



$$\Rightarrow i \hat{x}' = [\hat{x}, \hat{H}^{(a)}] + i \frac{a'}{a} \hat{x}$$

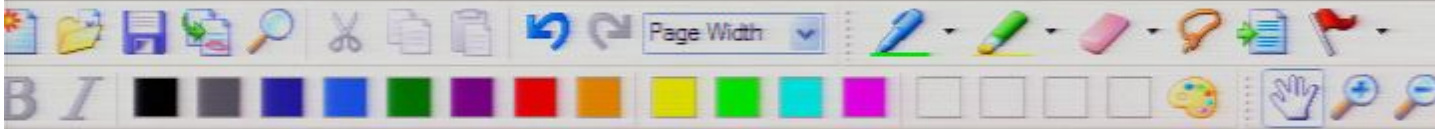
* But we also have:

$$i \hat{x}' = [\hat{x}, \hat{H}^{(b)}]$$

\Rightarrow We must have: $\hat{H}^{(a)} \neq \hat{H}^{(b)}$

Since there are multiple Hamiltonians, which, if anyone, is the energy?

△ One usually defines the energy as the generator of time evolution. We saw that in the presence of gravity this



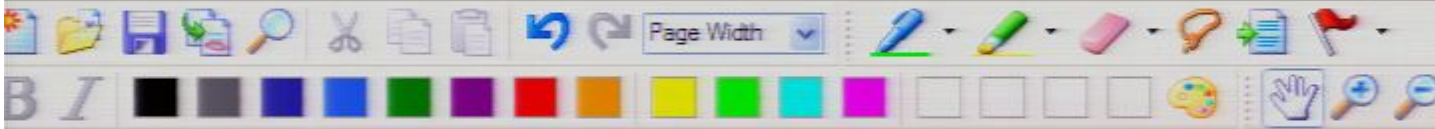
Since there are multiple Hamiltonians, which, if anyone, is the energy?

□ One usually defines the energy as the generator of time evolution. We saw that in the presence of gravity this is ambiguous: one can define many different Hamiltonians for the same theory (same action).

□ Therefore, with Einstein, we define the energy (density) not as the generator of time evolution but as a generator of curvature:

□ Recall: The Einstein equation

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□ Recall: The K.G. field's energy-momentum tensor

$$T_{\mu\nu}^{\text{K.G.}}(\eta, \vec{x}) = \frac{2}{\sqrt{|\eta|}} \frac{\delta S}{\delta g^{\mu\nu}} = \phi_{,\mu} \phi_{,\nu} - g_{\mu\nu} \left[\frac{1}{2} g^{\rho\sigma} \phi_{,\rho} \phi_{,\sigma} - \frac{1}{2} m^2 \phi^2 \right]$$

□ Consider $T_{00}(\eta, \vec{x})$, which is called the "energy density":

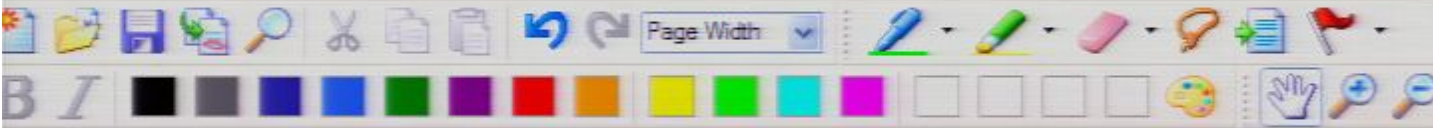
In differential geometry, there is also another use of the term "density":

For any tensor, say $A_{\mu\nu}$, there is a so-called "tensor density" $\tilde{A}_{\mu\nu}$, defined as $\tilde{A}_{\mu\nu} := A_{\mu\nu} \sqrt{g}$, which absorbs the obligatory measure factor in integrations.

$$T_{00}(\eta, \vec{x}) = a^{-4} \frac{1}{2} \pi^{(0)2} + \frac{1}{2} \sum_{i=1}^3 \phi_{,i}^2 + \frac{a^2}{2} m^2 \phi^2 \quad (T)$$

□ Exercises:

a) Verify (T).



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□ Exercises:

a) Verify (T).

b) Calculate $H^{(0)}$.

Notice that $H^{(0)}$ is not a scalar.

c) Show that $H^{(0)}(\eta) = \int_{\mathbb{R}^3} T_{00}^{\text{KG}}(\eta, \vec{x}) \sqrt{|g|} d^3x$.

d) Calculate $H^{(00)}(\eta)$.