

Title: Probing primordial non-Gaussianity with today's universe

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Abstract: Primordial non-Gaussianity has been traditionally constrained using three-point function of the cosmic microwave background. Two years ago, however, Dalal et al have shown that non-Gaussianity of the local type induces a scale dependent bias for biased tracers of the underlying dark matter structure. This allows constraining of the primordial non-Gaussianity from measurements of large-scale structure provided by redshift surveys. I will discuss the technique, its theoretical aspects, its surprising resilience towards systematics and current results from the real data. I will also show some preliminary new results: extension to the two field inflationary models and the analogue of the Dalal effect in the Lyman alpha forest.

Primordial non-Gaussianity with today's Universe

PI, March 2010

Luž e Slosar,

BNL

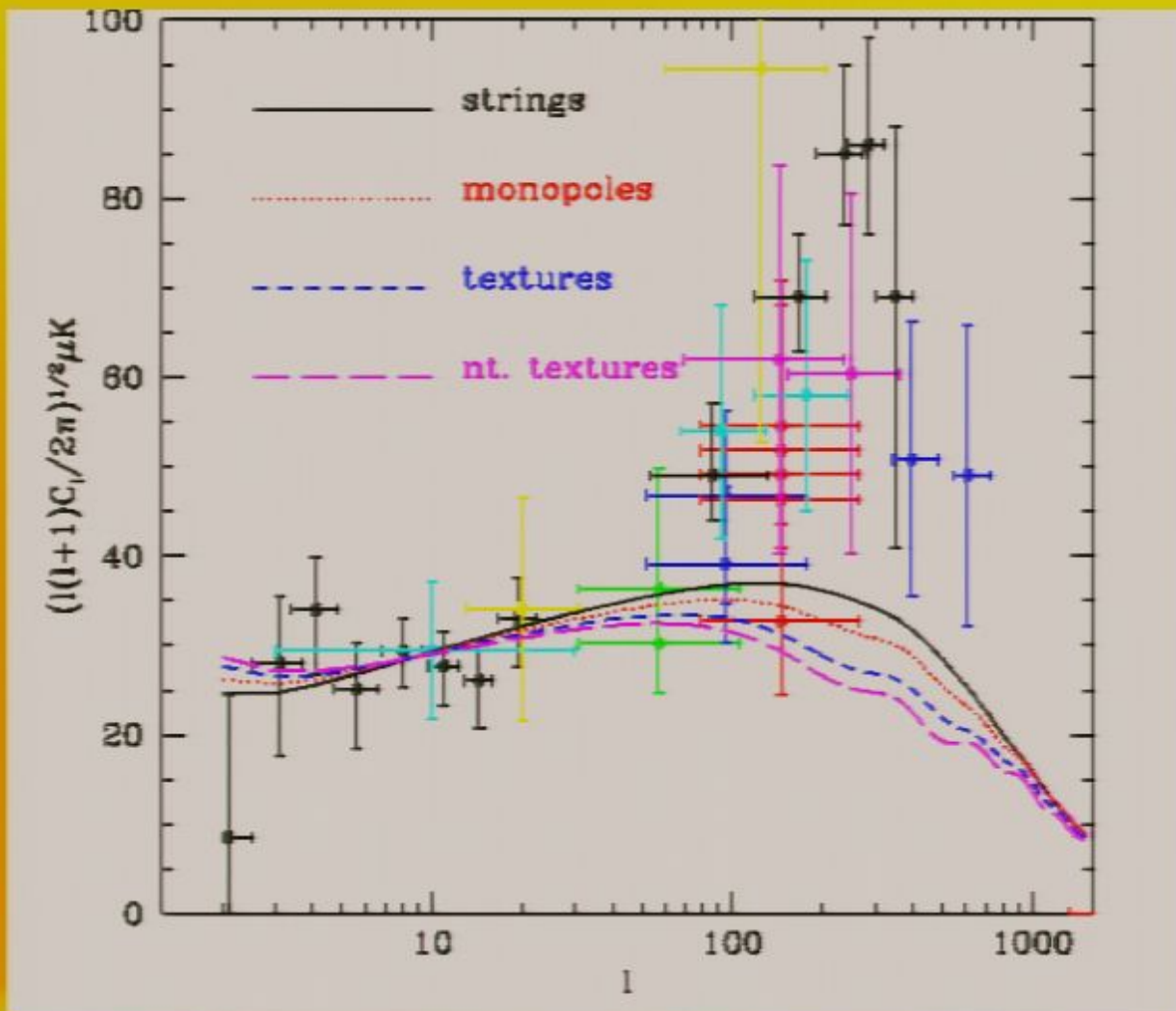


Historical note:

- In 90s, the source of primordial fluctuations was unknown. There were two options:
 - Inflation – generic predictions of the nearly scale invariant primordial power spectrum
 - Topological defects: model dependent and very different power spectrum



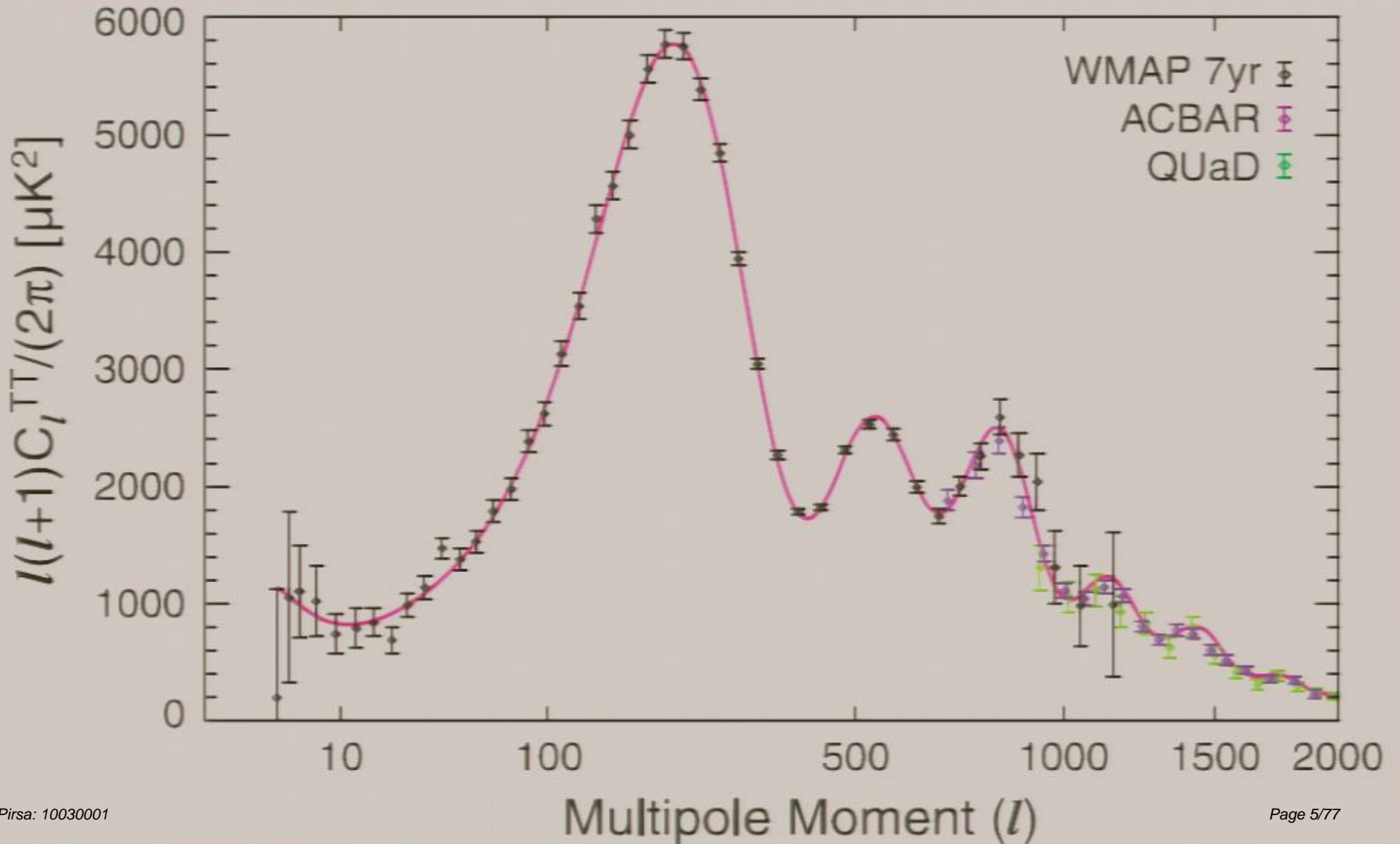
Inflation vs defects:



- Data in the older days (late 90s)
- HAHHAHA!!!
- Inflation clearly won
- The strength of inflation is its curse:

- predictions are very model independent

2010:



Inflation:

How to constrain inflation? Not too many options:

- measure departures from a power-law spectrum
- measure B-modes of CMB polarization
- Measure iso-curvature modes
- measure non-Gaussianity



Gaussianity:

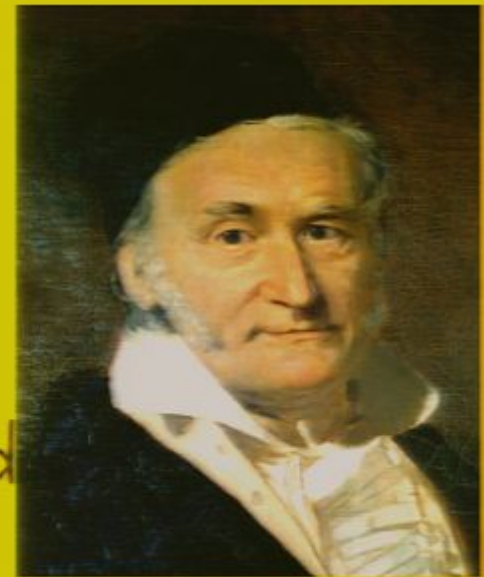
In standard inflationary models, primordial curvature perturbations and in turn other observables are Gaussian:

Every Fourier mode has a random phase and amplitude randomly taken from a Gaussian distribution with variance $P(k)$

All odd point statistics vanish, all even points statistics are completely determined by $P(k)$

Gaussianity can be falsified, non-

Gaussianity cannot



Non-Gaussianity:

In minimal inflationary models, curvature perturbation:

$$\mathcal{R} = \frac{H(\phi)}{\dot{\phi}} \delta\phi$$

• if ϕ is normally distributed, so is \mathcal{R} .

• One can break non-Gaussianity:

- scalar fields non-free

- non-linear corrections to above equation

- initial state non Bunch-Davis vacuum

Non-Gaussianity:

A generic way is to postulate that curvature perturbation in the matter domination obeys

$$\Phi = \Phi_L + f_{NL}^{\text{local}} \Phi_L^2.$$

- this is the so-called "local" non-Gaussianity:
 - local in real space and in potential!
- By far not the only possible non-Gaussianity
- If detected, we can check further, most non-Gaussianity will have a projected component
- Note the typical scales – these are very stringent

Non-Gaussianity:

What are the typical value of f_{NL} ?

2nd order perturbation theory of GR gives $f_{NL} \sim 1$

Simplest scalar field gives $f_{NL} \sim 0.01$

More complicated models up to $f_{NL} \sim 10$

$f_{NL} > 10$ implies very non-concordant physics: New
Ekpyrotic scenario, violent reheating, etc.

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How to measure f_{NL} ?

In principle, one can use 1-point distribution function

In practice, higher-order correlators give much better signal-to-noise: best to use bispectrum

CMB the cleanest probe (as always)

Difficult to measure, but special estimators have been constructed by Komatsu and company

f_{nl} from CMB

Komatsu et al in WMAP3:

$$-54 < f_{nl} < 114$$

Creminelli et al:

$$-36 < f_{nl} < 100$$

Komatsu et al in WMAP5:

$$-9 < f_{nl} < 111$$

Yadav & Wandelt:

$f_{nl} > 0$ at 2.9 sigma

Smith et al

$$-4 < f_{nl} < 80$$

Komatsu et al in WMAP7:

$$-10 < f_{nl} < 74$$

Hikage et al:

$$-71 < f_{nl} < 91$$

Curto et al (wavelets):

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Note on sigmas:



Linda Evangelista: "I don't go out of bed for less than \$10000."

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Andy Albrecht: "I don't go out of bed for less than 4 sigma."

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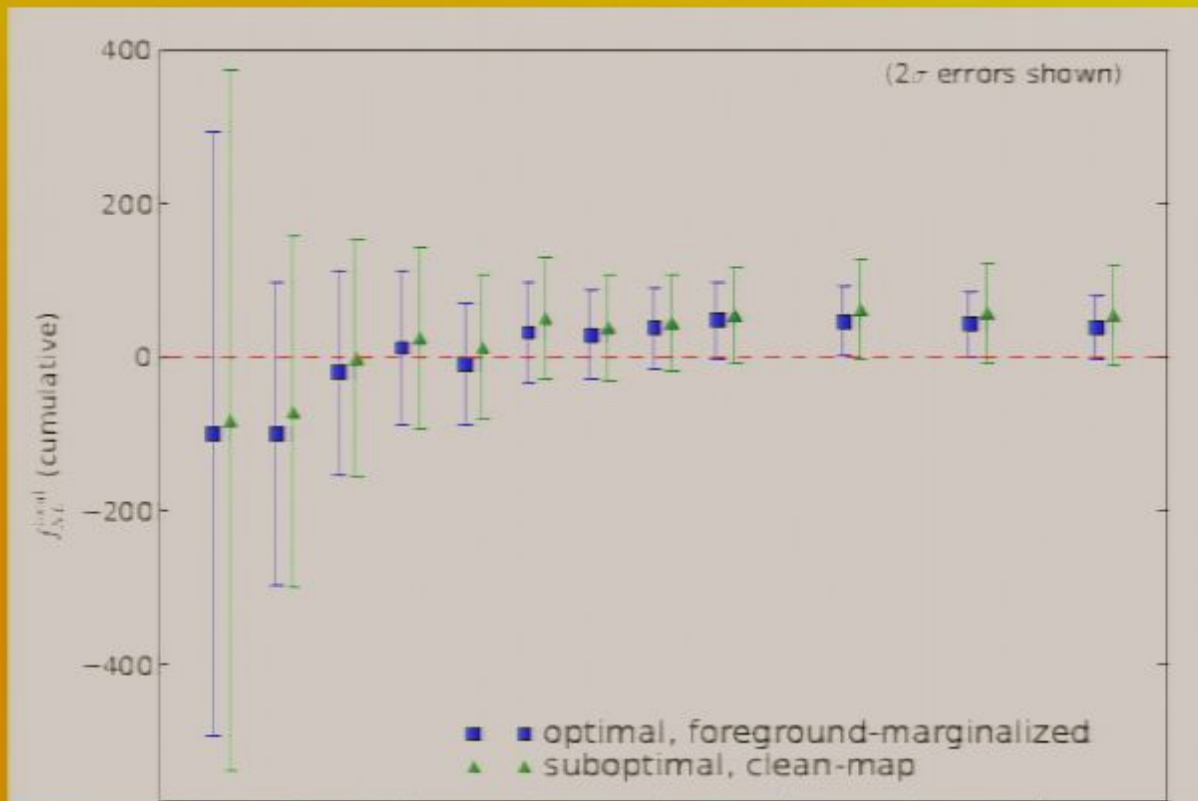
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WMAP 5:



- Smith et al
- Most optimal estimator
- The optimality is a compsci rather than physics problem
- Don't see anything exciting, some discrepancies remain

FWI from LSS?

If you have ability to probe the primordial density fields – see it direct (maybe 21cm at high z ?)

It is 3-dimensional so in principle much better

In practice this is never the case: non-linearities induce their own non-Gaussianity

Initial idea:

- Look at the mass function in the exponential tail
- rare peak probability affected a lot in even weakly

NG models

Galaxy power spectrum

- Assume that process of formation of galaxies is local on some scale.

$$\delta_g(\vec{x}) = f(\delta_m(\vec{x}), \nabla \delta_m(\vec{x}), \dots)$$

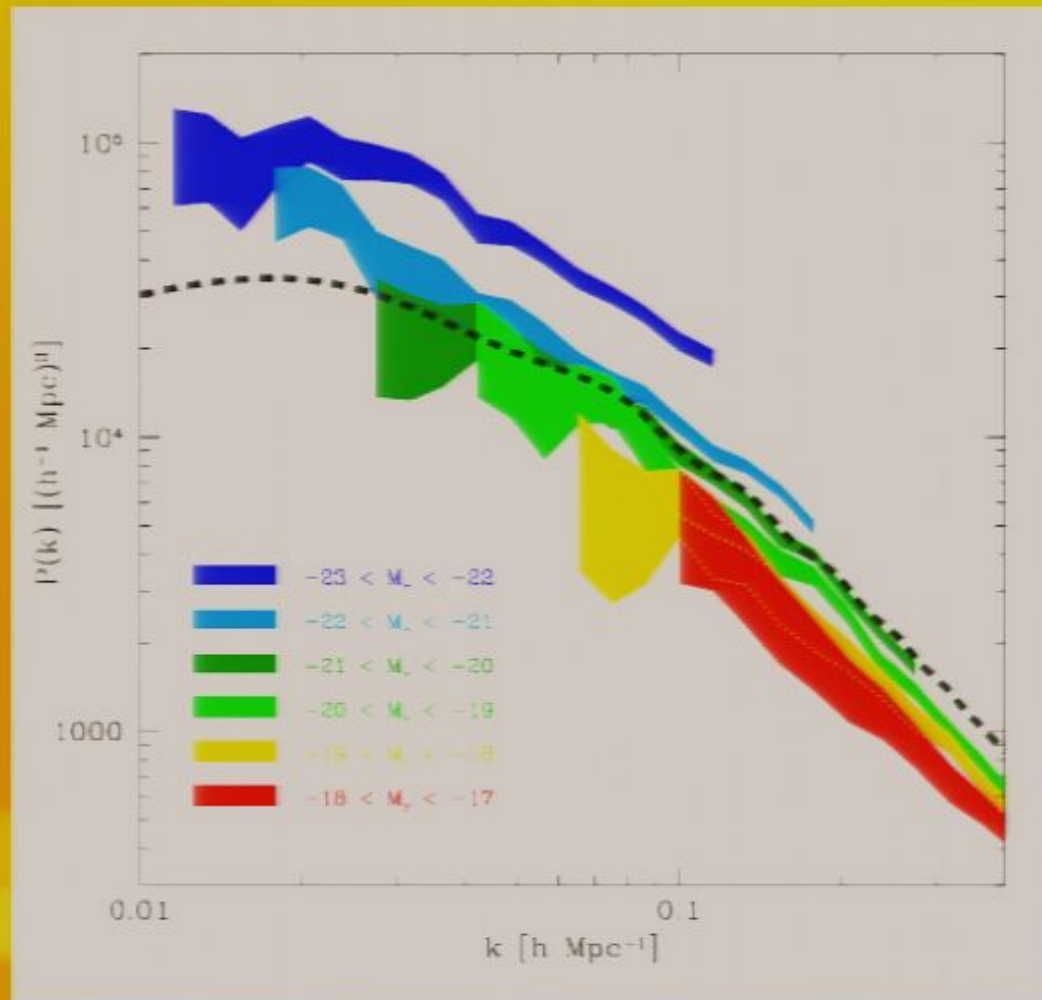
- On large scales, it follows that

$$\delta_g(\vec{k}) = b \delta_m(\vec{k}) + s(\vec{k})$$

- And hence

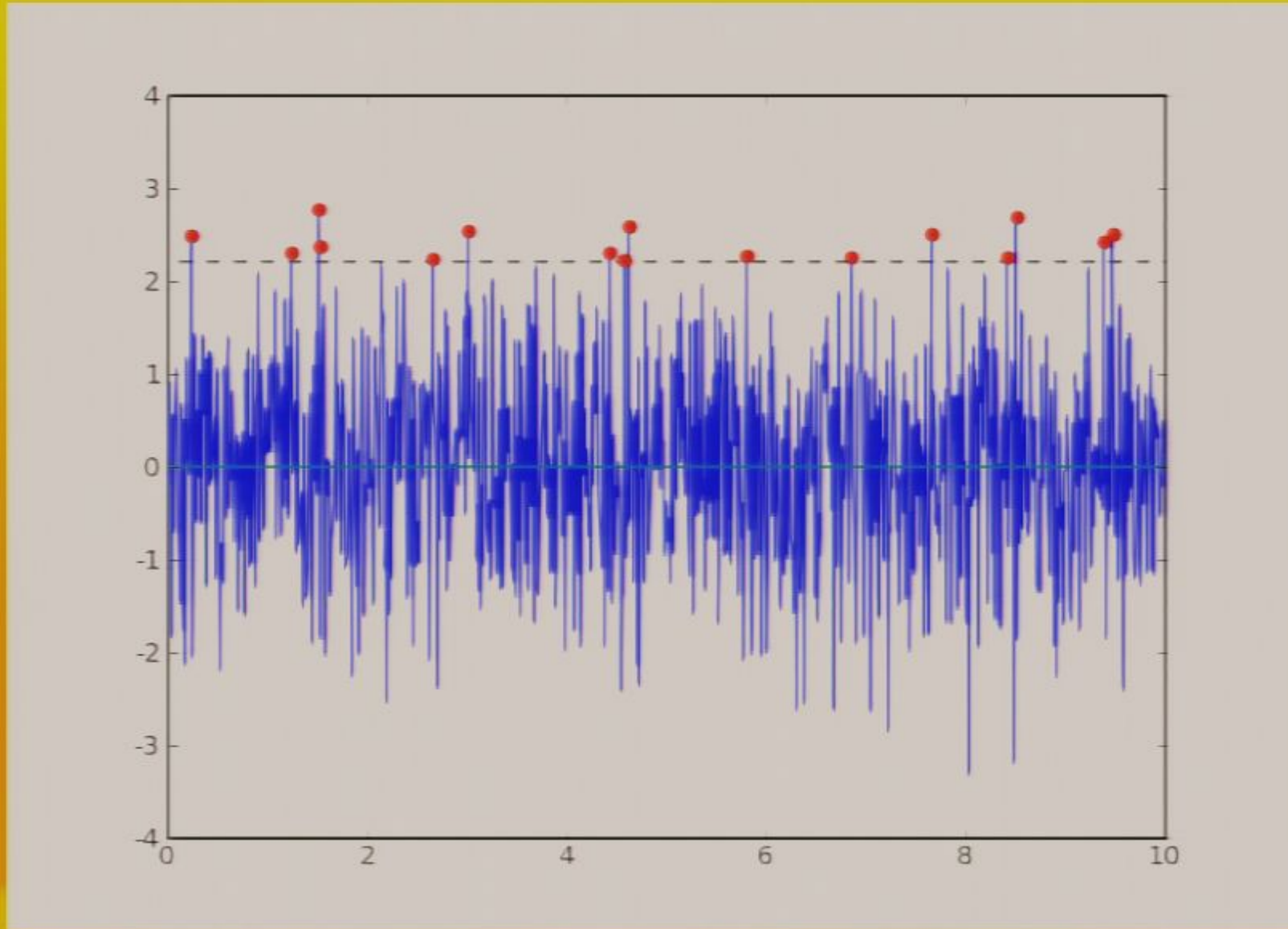
$$P_g(k) = b^2 P_m(k) + \text{const.}$$

Power spectrum of galaxies

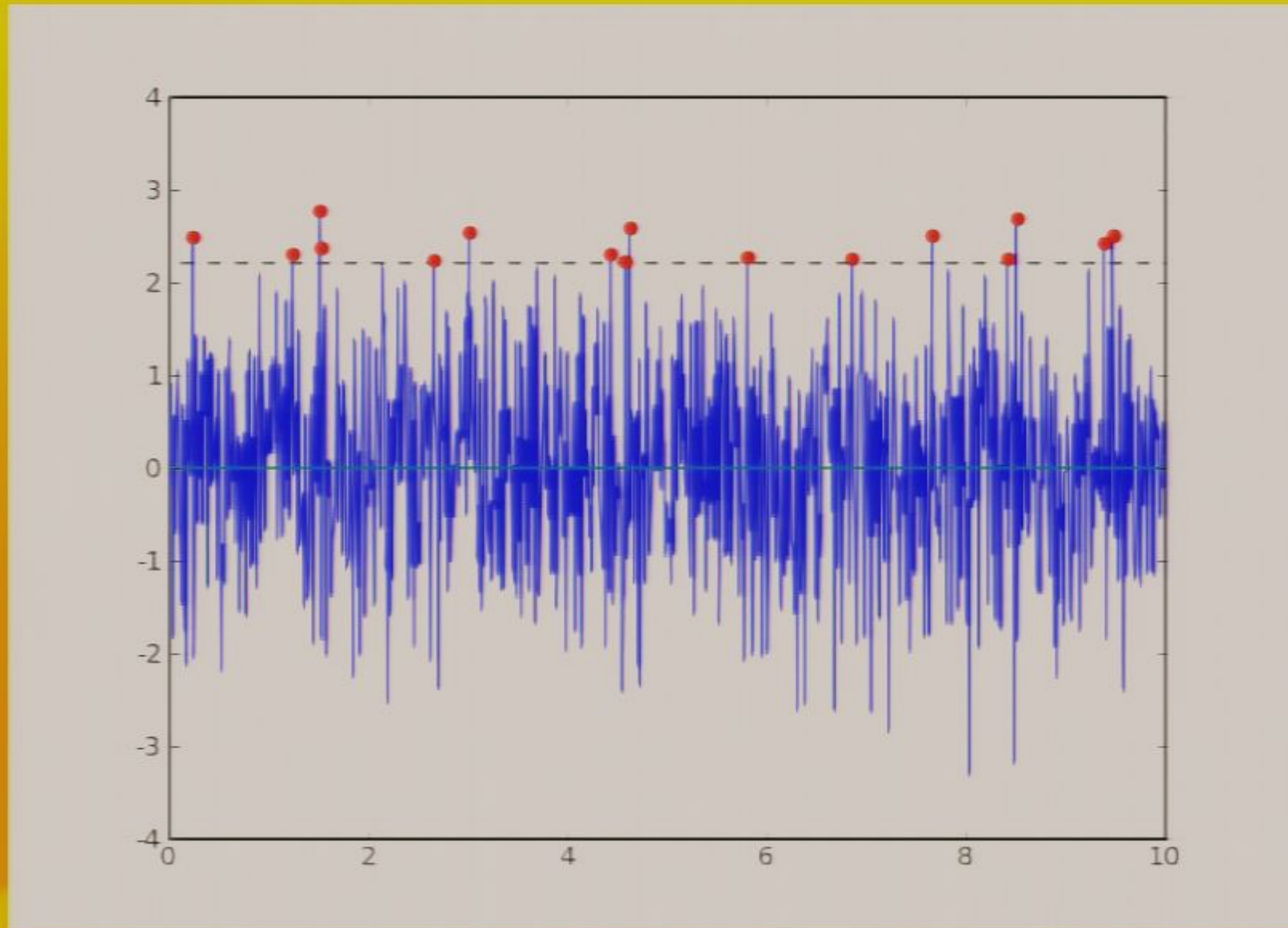


From Tegmark et al

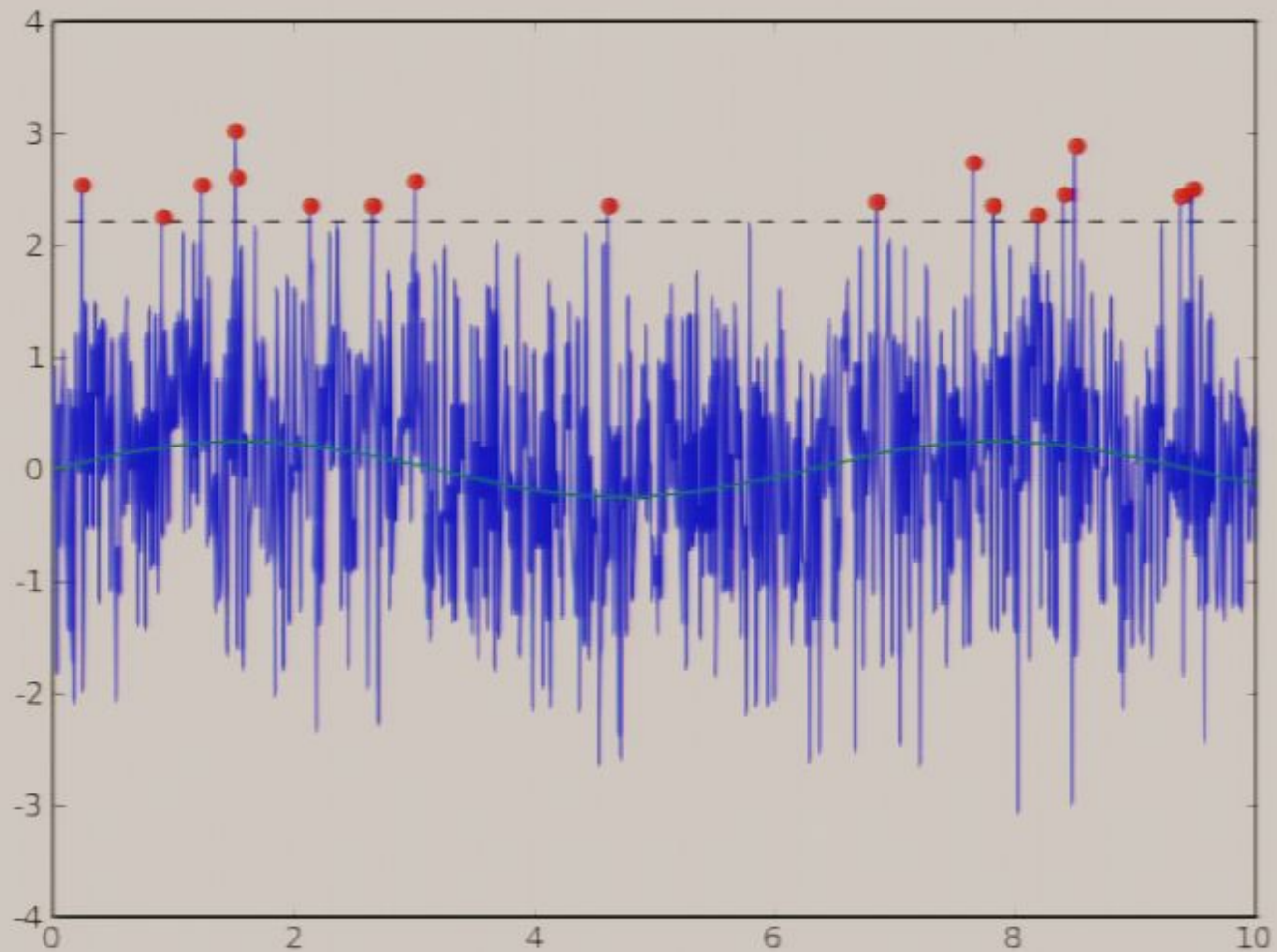
Detour: rare peaks



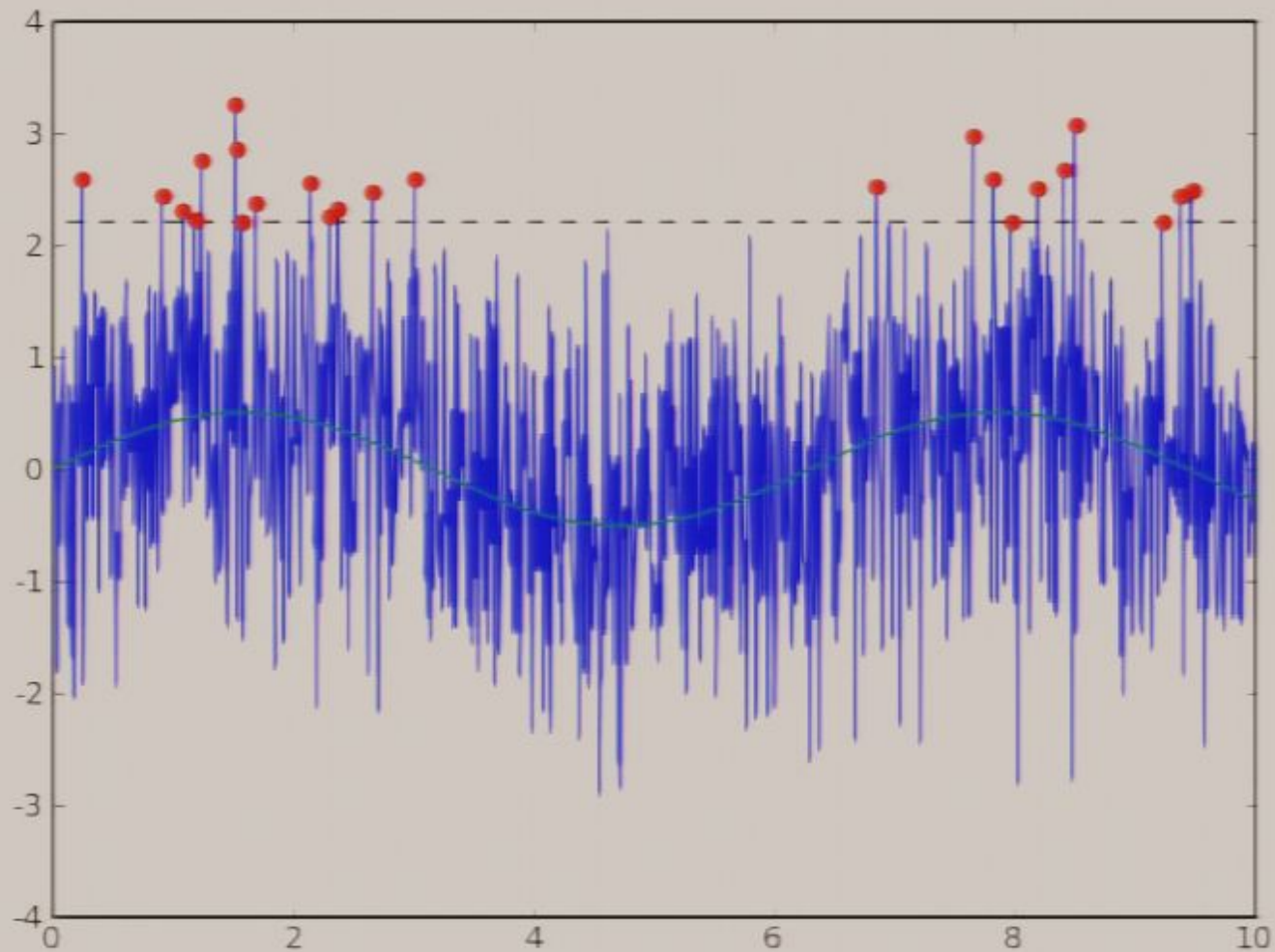
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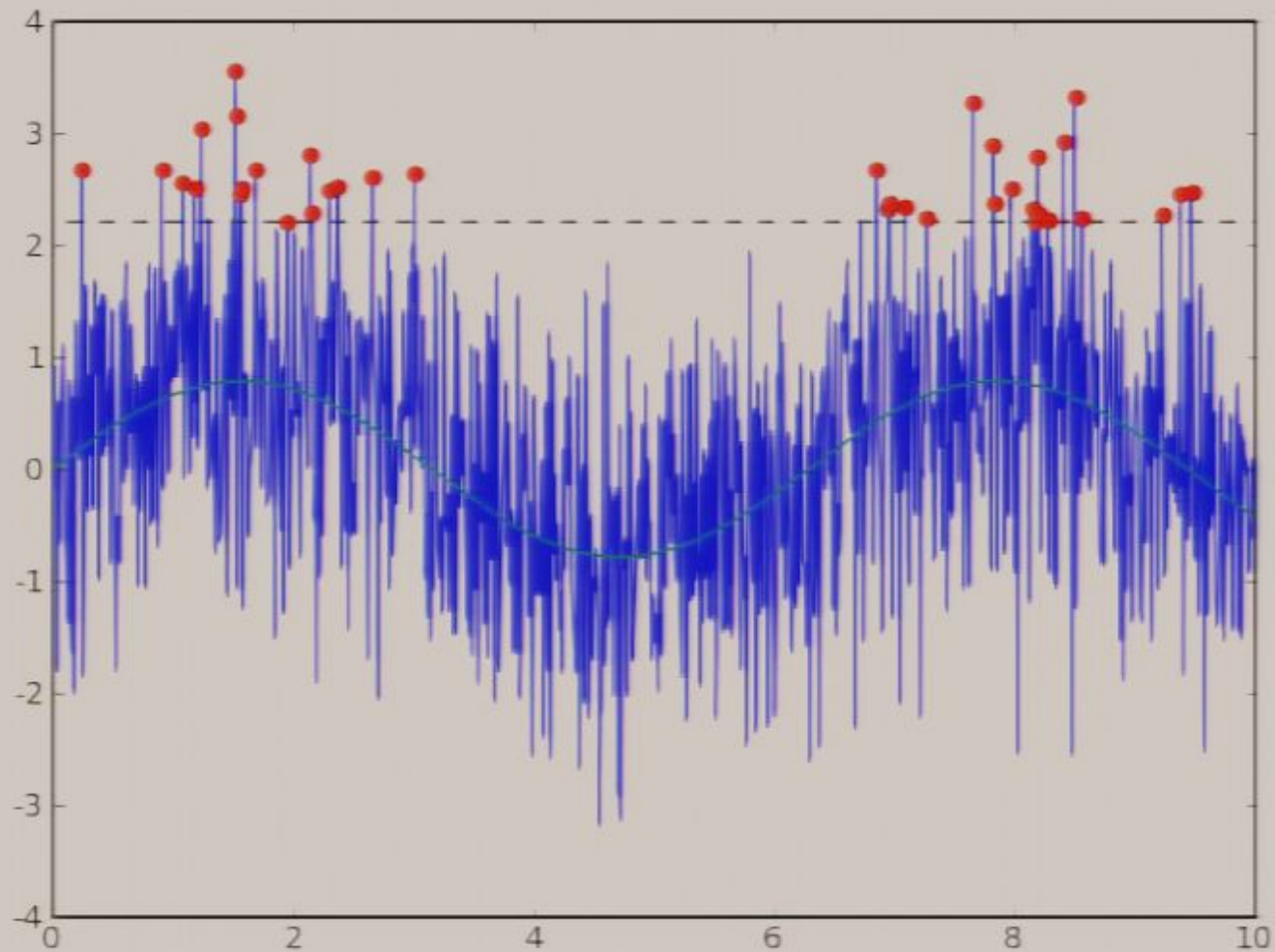
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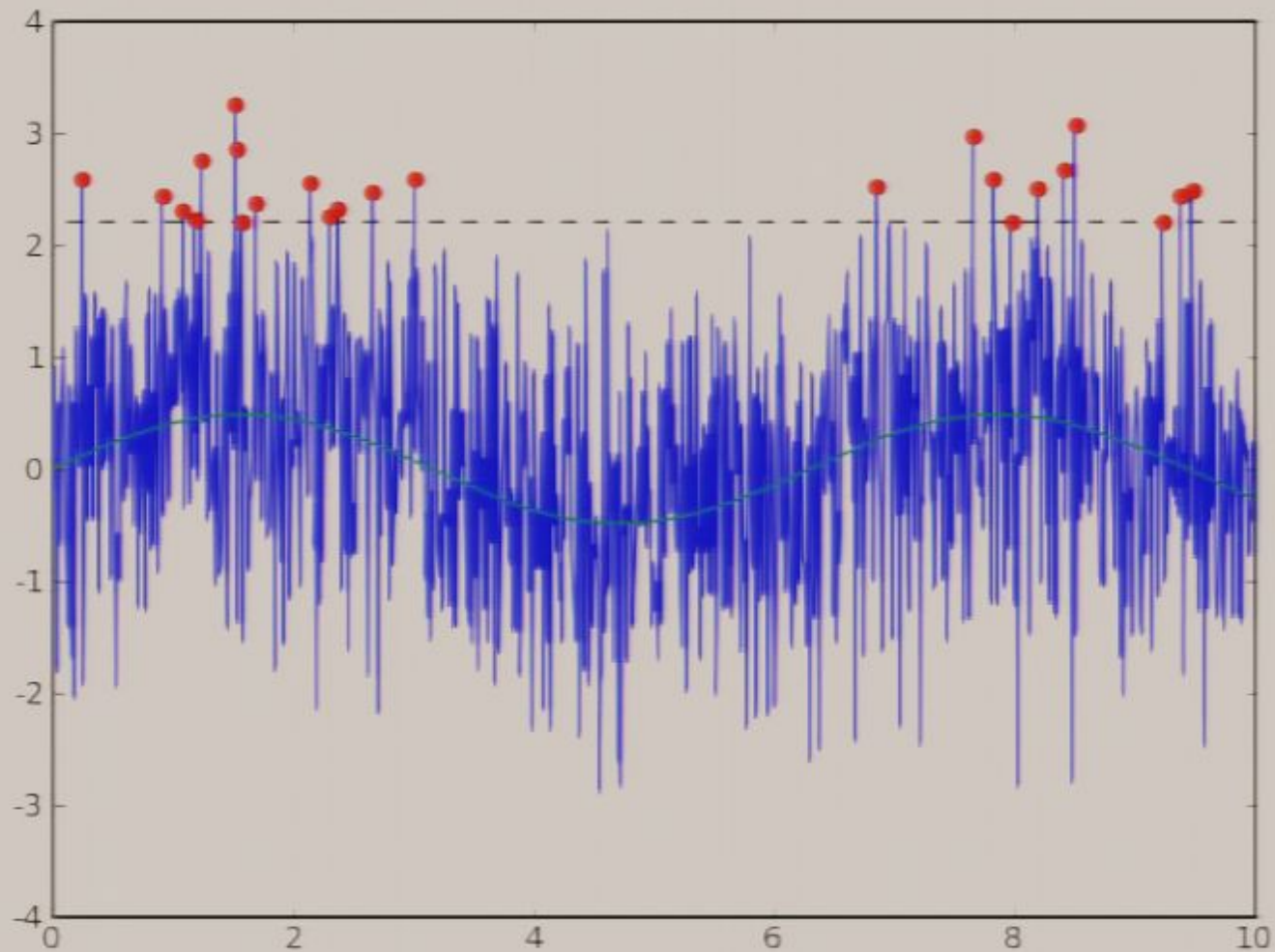
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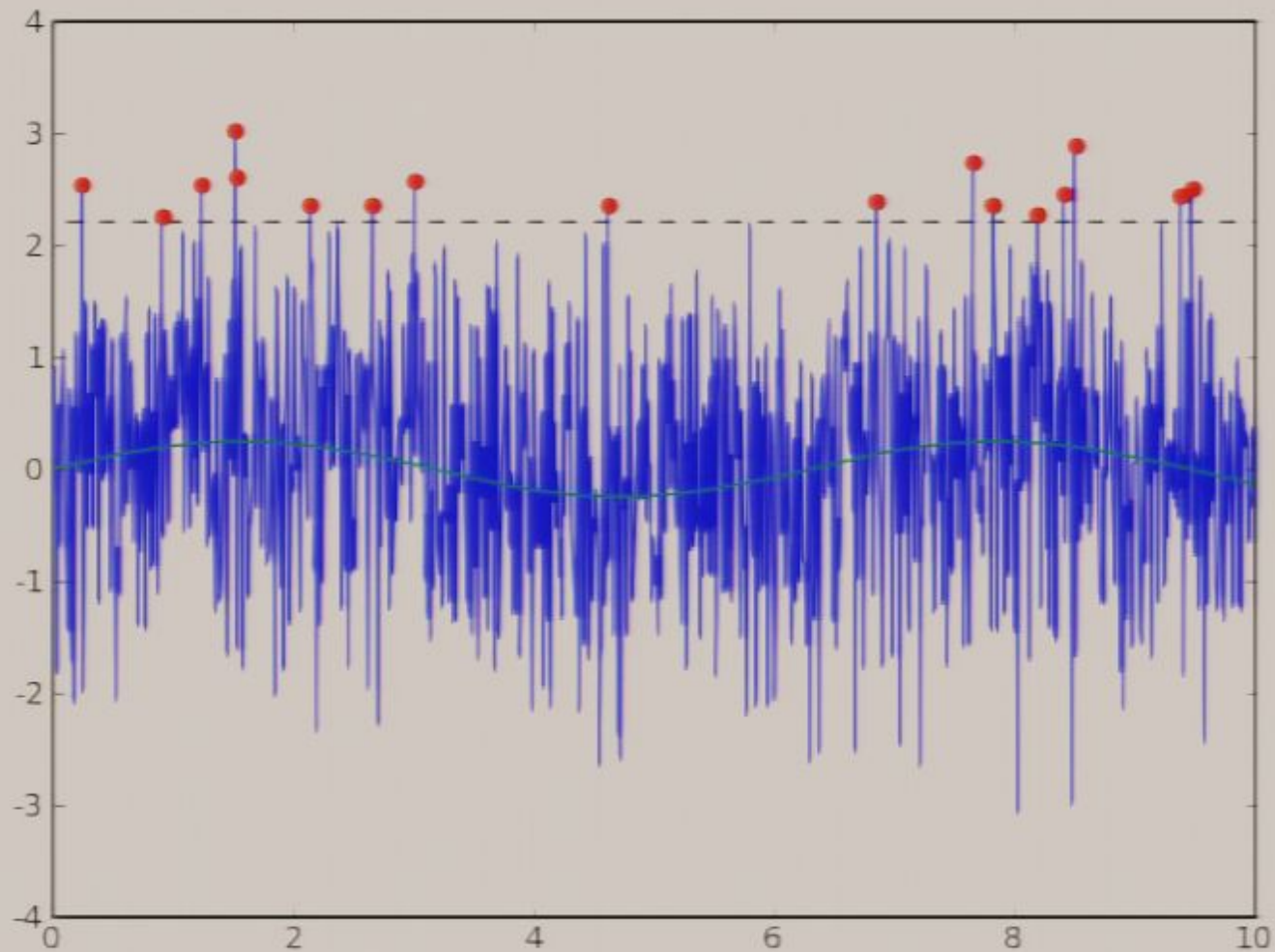
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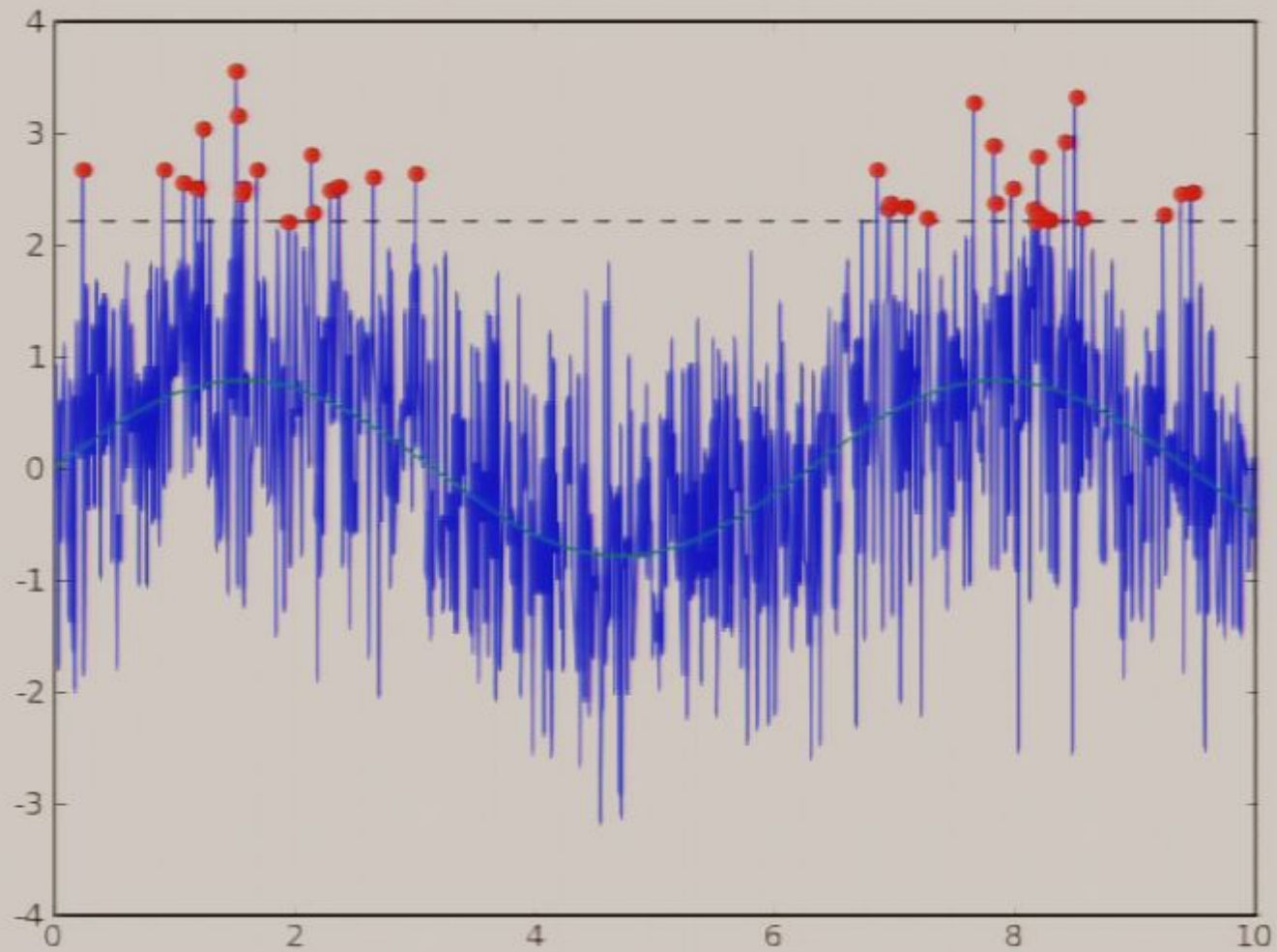
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Scale dependent bias:

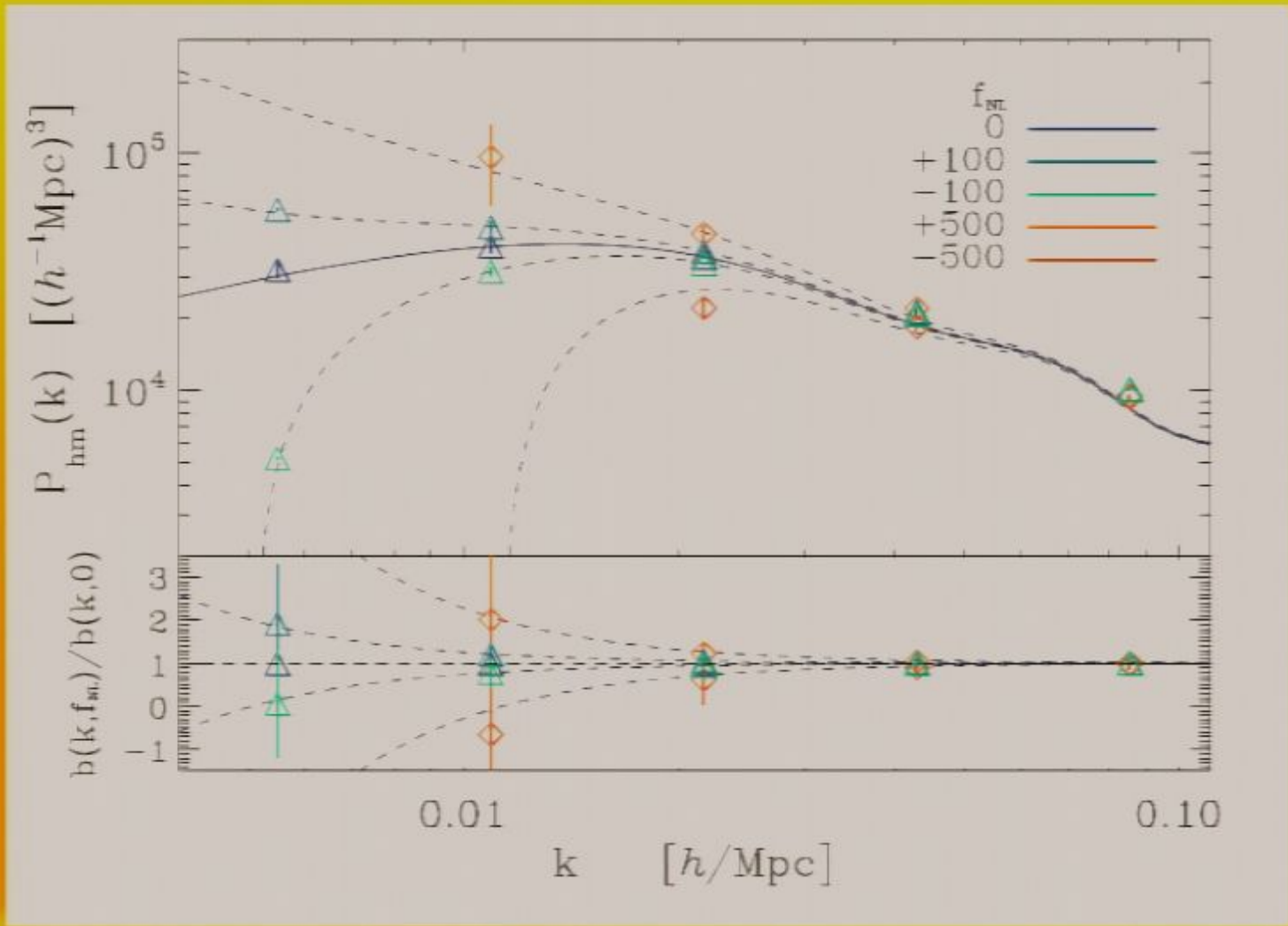
Dalal et al note a very different and surprising effect

- scale dependent bias in clustering of rare peaks in f_{NL} cosmologies.
- Get scale dependent bias in clustering of highly biased objects. Formula is:

$$\Delta b(M, k) = 3f_{\text{NL}}(b - 1)\delta_c \frac{\Omega_m}{k^2 T(k) D(z)} \left(\frac{H_0}{c} \right)^2$$

- Matarrese & Verde, Afshordi and Tolley, McDonald, Slosar et, Giannantonio & Porciani, etc derive an equivalent formula using a multitude of approaches.

- Note k and b dependence.



Hirata's magic:

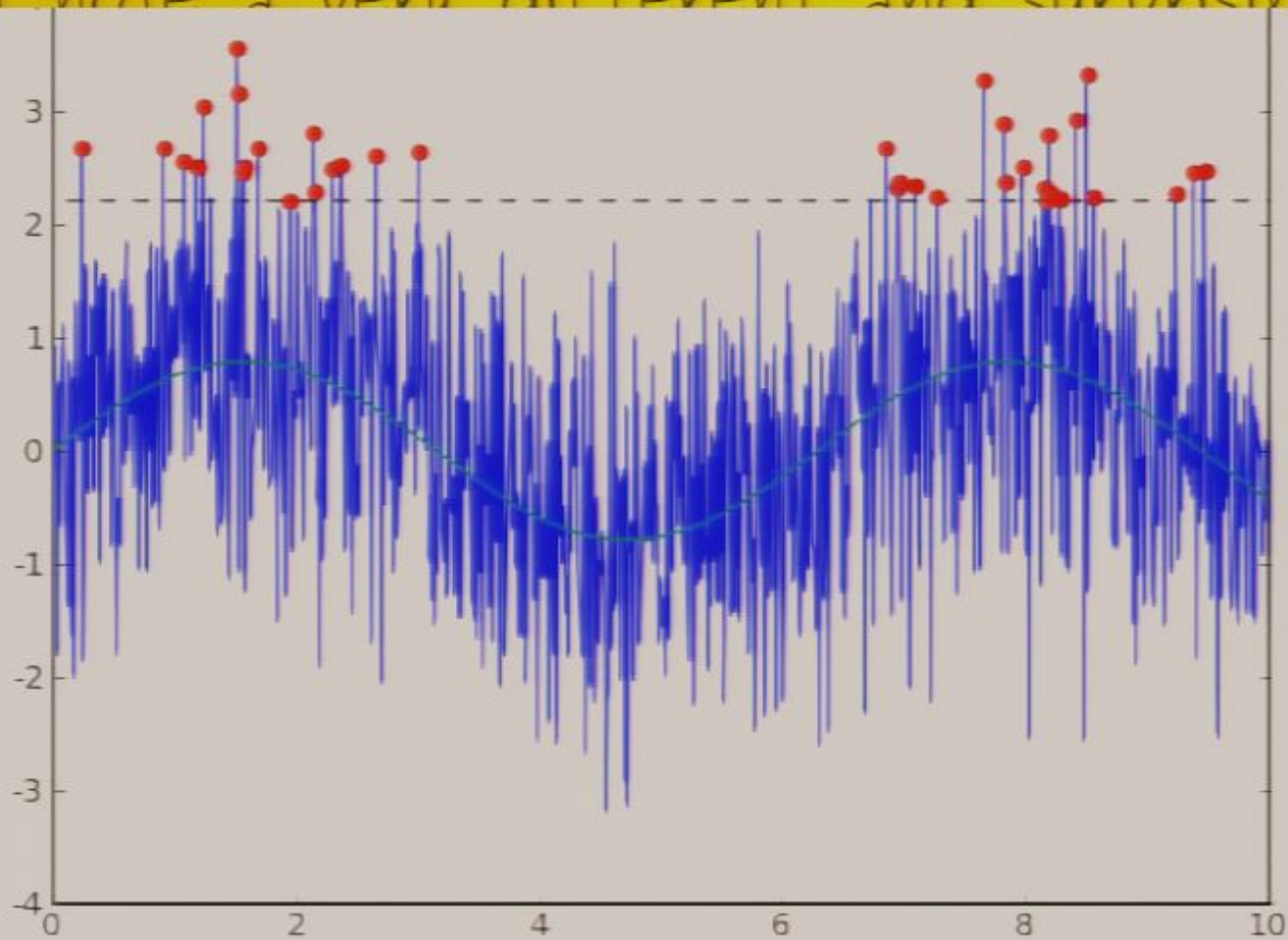
$$\Delta b(M, k) = \frac{3\Omega_m H_0^2}{c^2 k^2 T(k) D(z)} f_{\text{NL}} \frac{\partial \ln n}{\partial \ln \sigma_8}.$$

- It can be shown that all PS-like theories in which mass function is function of $\nu = \delta_c^2 / \sigma^2(M)$ only reduce to Dalal et al formula.
- But they need not be. For objects tracing recent mergers one gets:

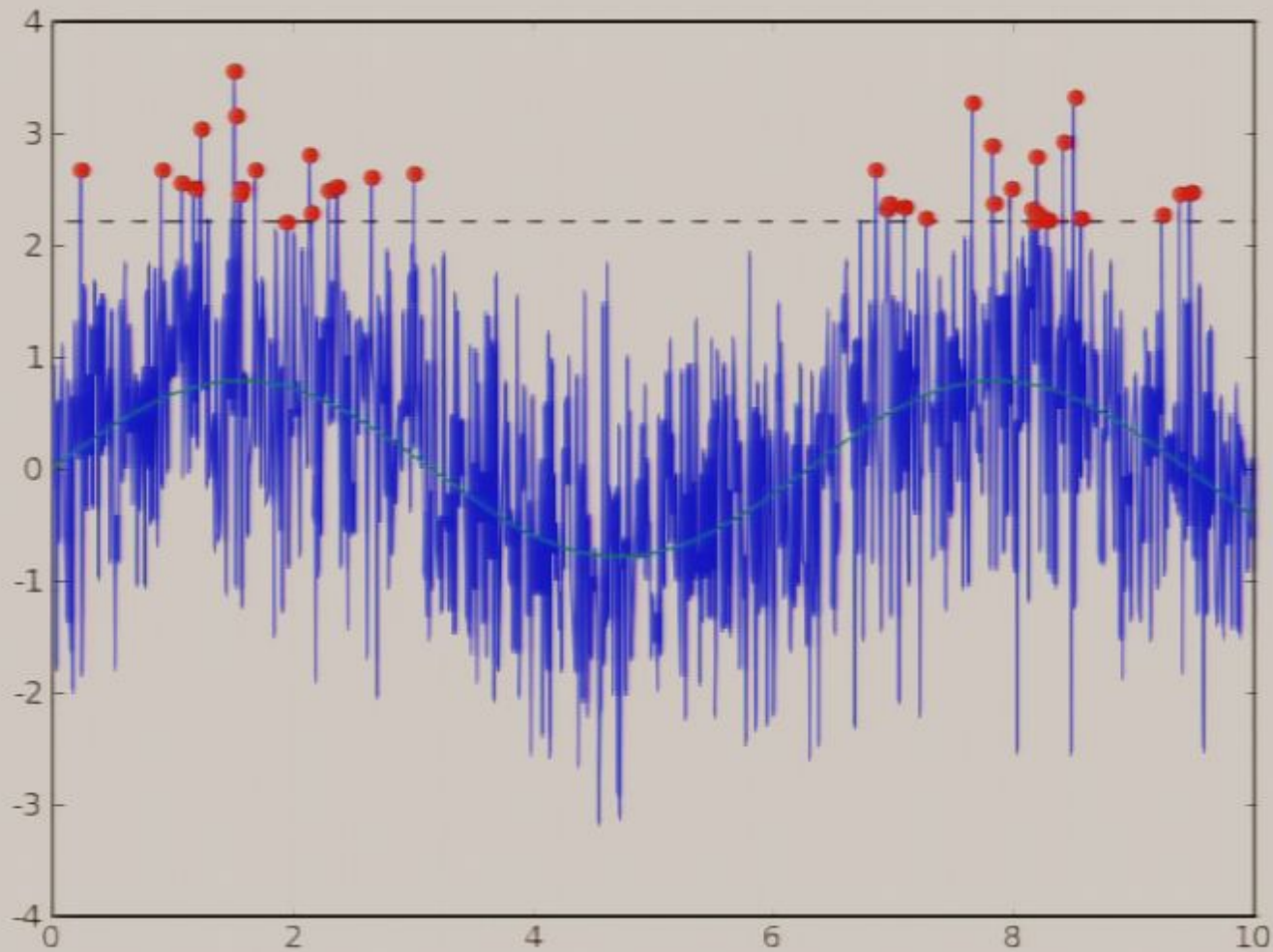
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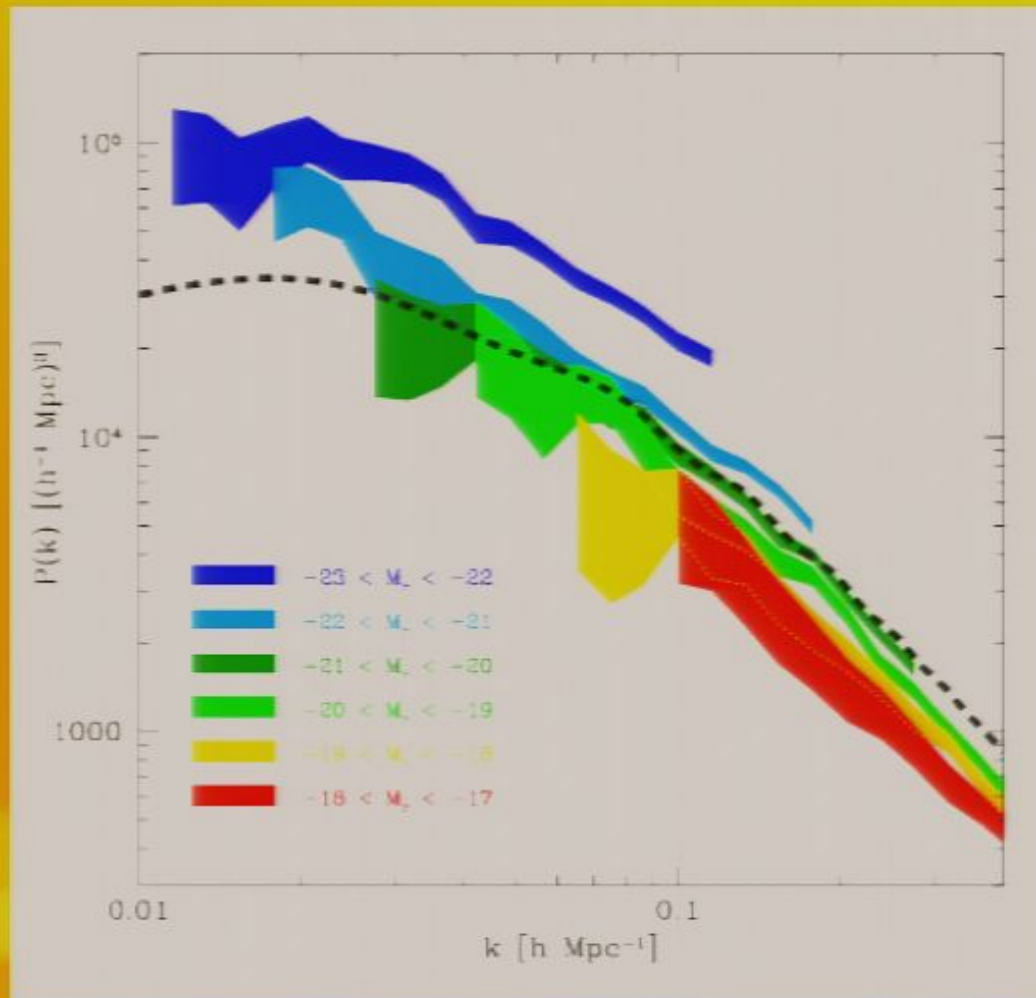
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Power spectrum of galaxies



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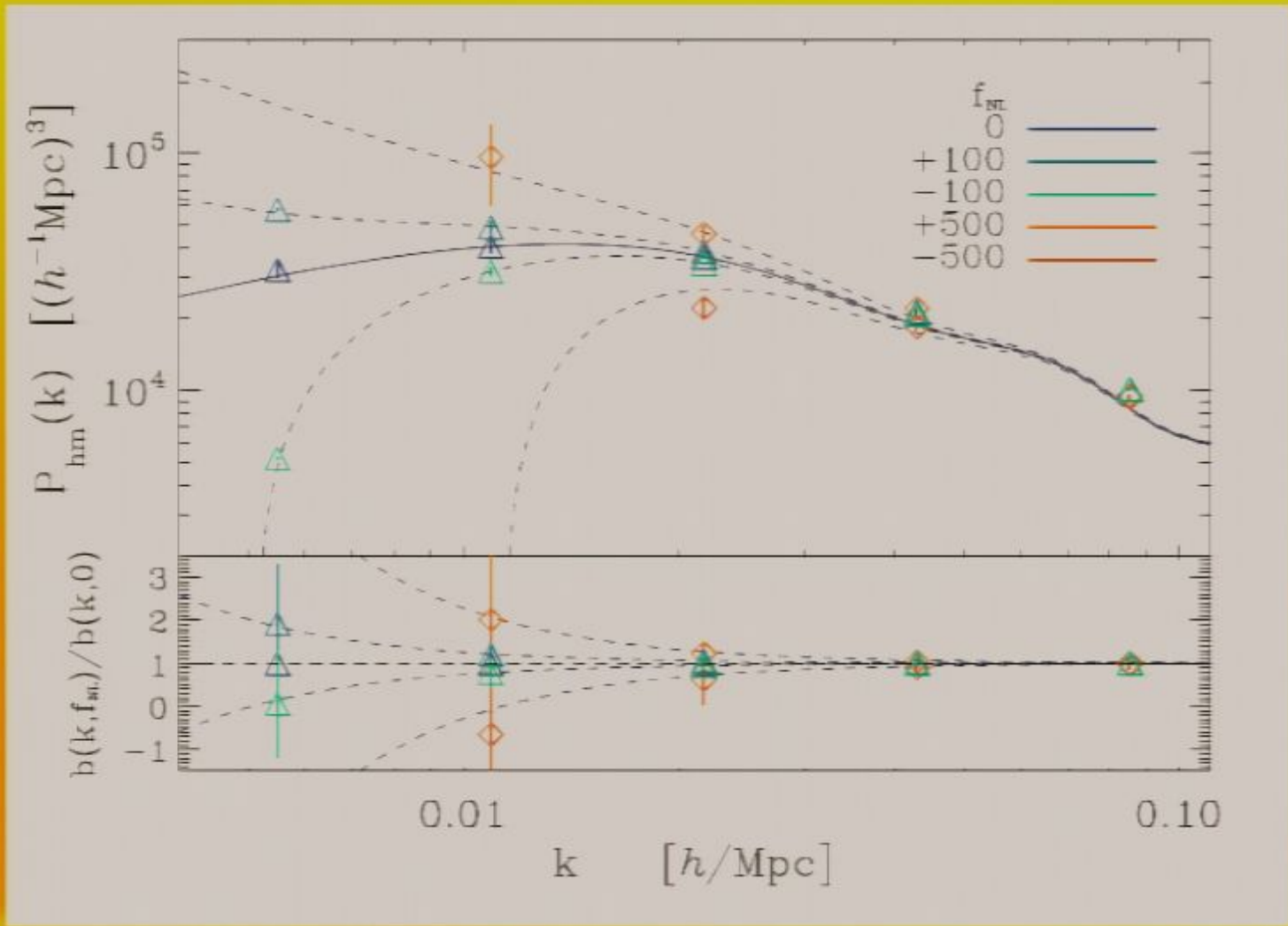
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Verde & Matarrese's magic:

- Two point function of peaks can be expressed as

$$\xi_{h,M}(|\mathbf{x}_1 - \mathbf{x}_2|) = -1 + \quad (\dagger)$$

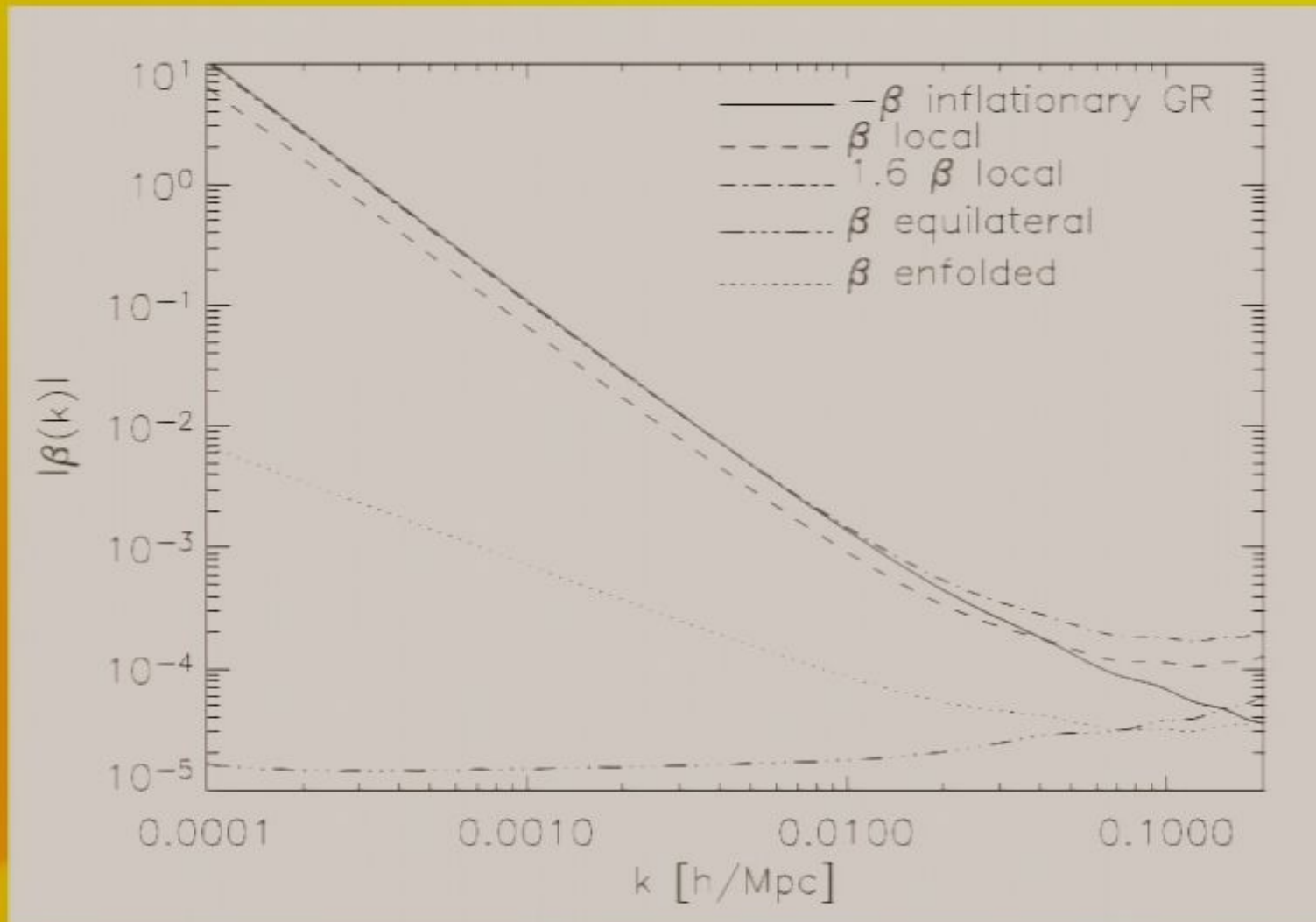
$$\exp \left\{ \sum_{N=2}^{\infty} \sum_{j=1}^{N-1} \frac{\nu^N \sigma_R^{-N}}{j!(N-j)!} \zeta^{(N)} \left[\begin{array}{c} \mathbf{x}_1, \dots, \mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_2 \\ j \text{ times} \quad (N-j) \text{ times} \end{array} \right] \right\}.$$

- They get:

$$P_h(k, z) = \frac{\delta_c^2(z) P_{\delta\delta}(k, z)}{\sigma_R^4 D^2(z)} \left[1 + \dagger f_{\text{NL}} \delta_c(z) \frac{P_{\phi\delta}(k) \mathcal{F}_R(k)}{P_{\delta\delta}(k)} \right]$$

- Can write two-point function of peaks for any non-Gaussianity as an integral over higher order correlators

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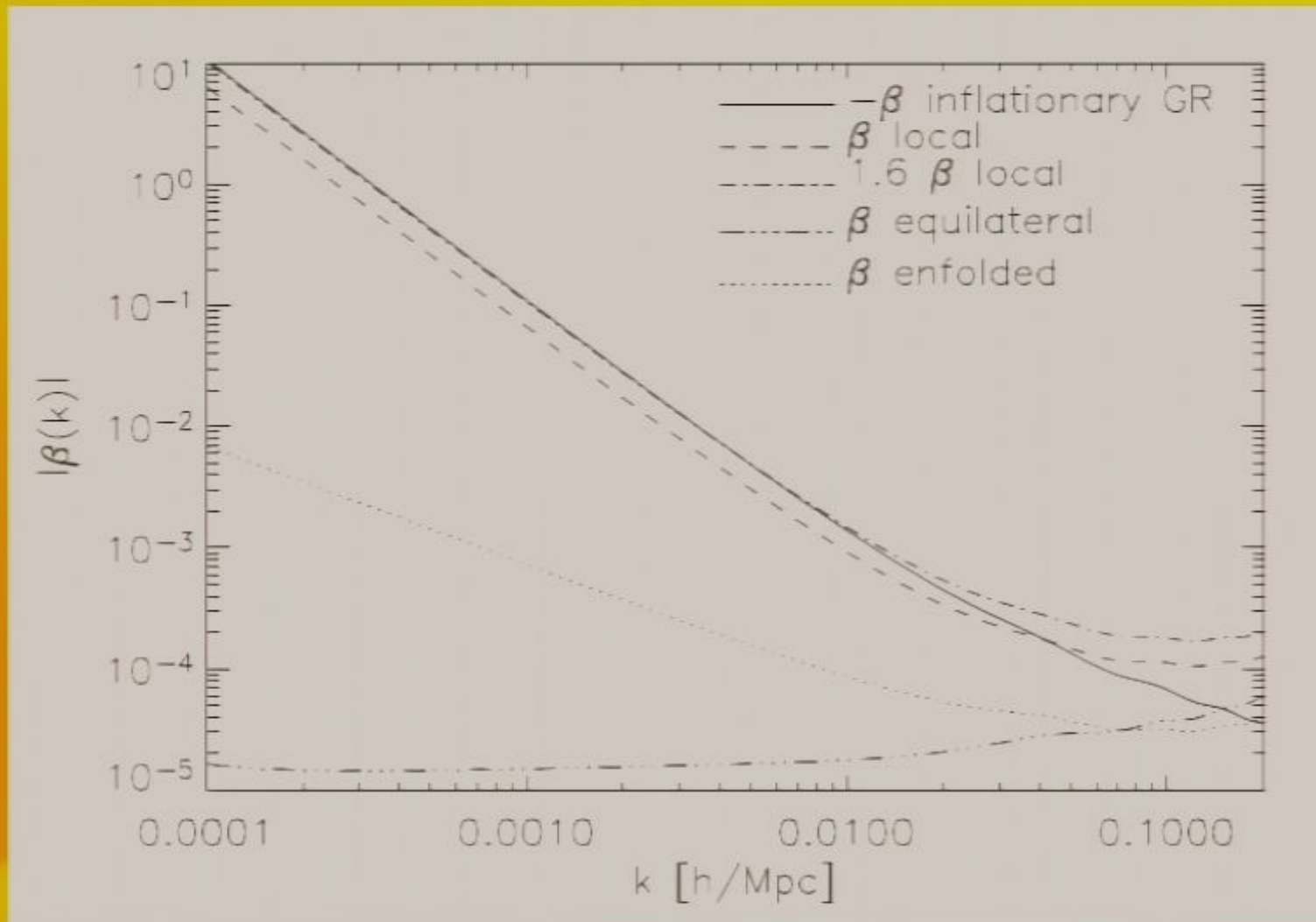
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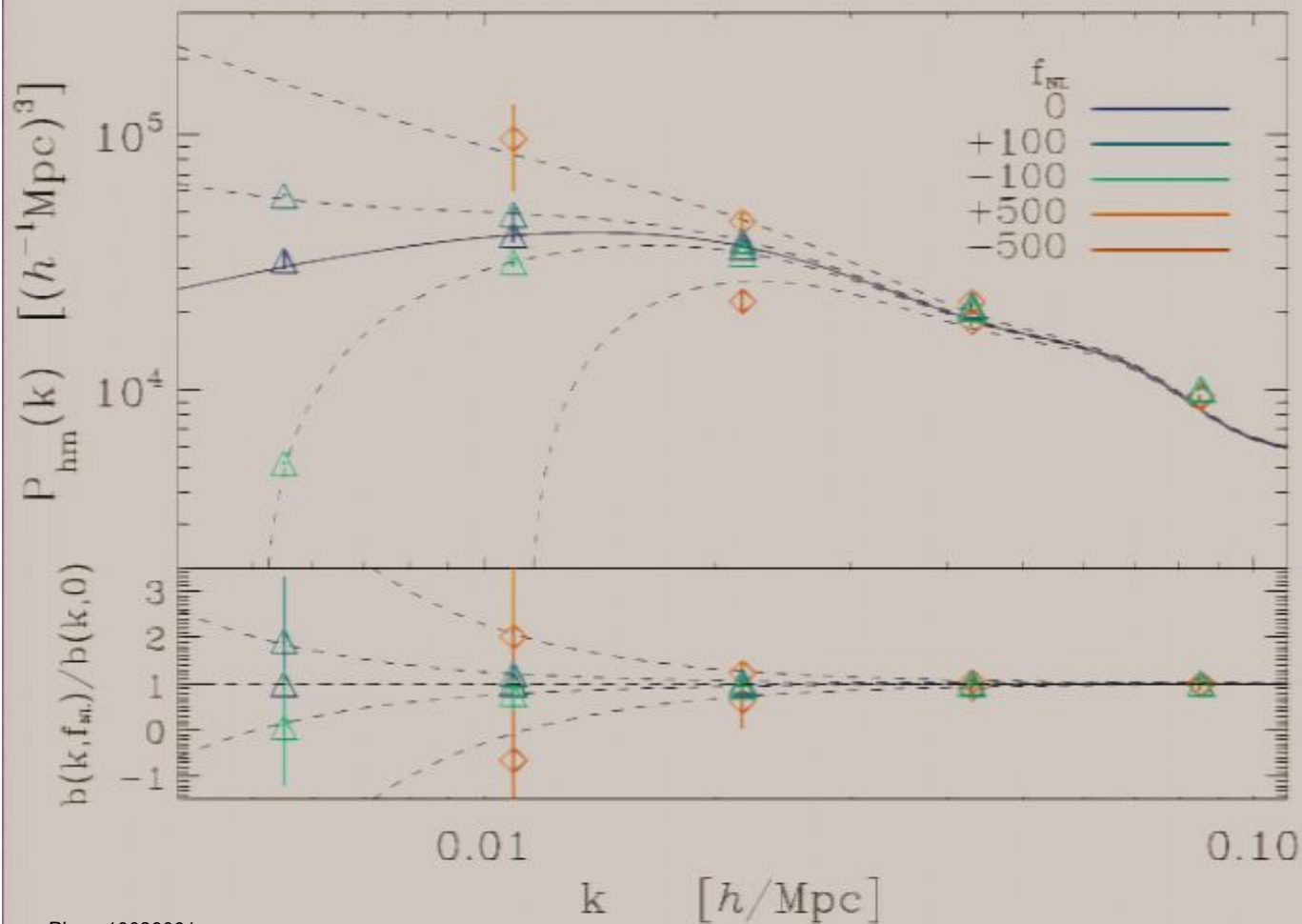
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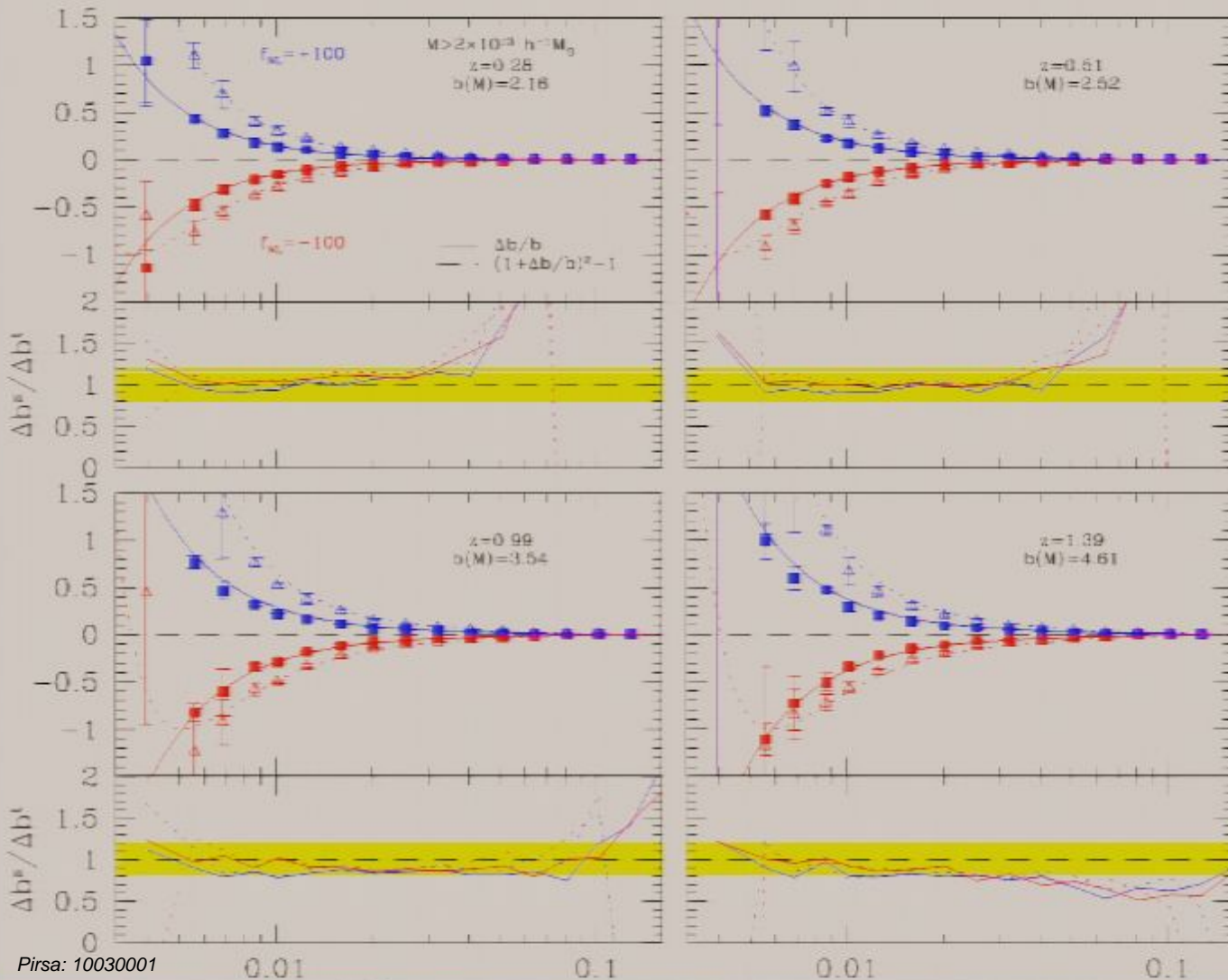


Simulations:



- Formula tested in cross-correlation!!
- seems to work like magic!
- (from Dalal et al)
- Note that one is safe from NL physics.

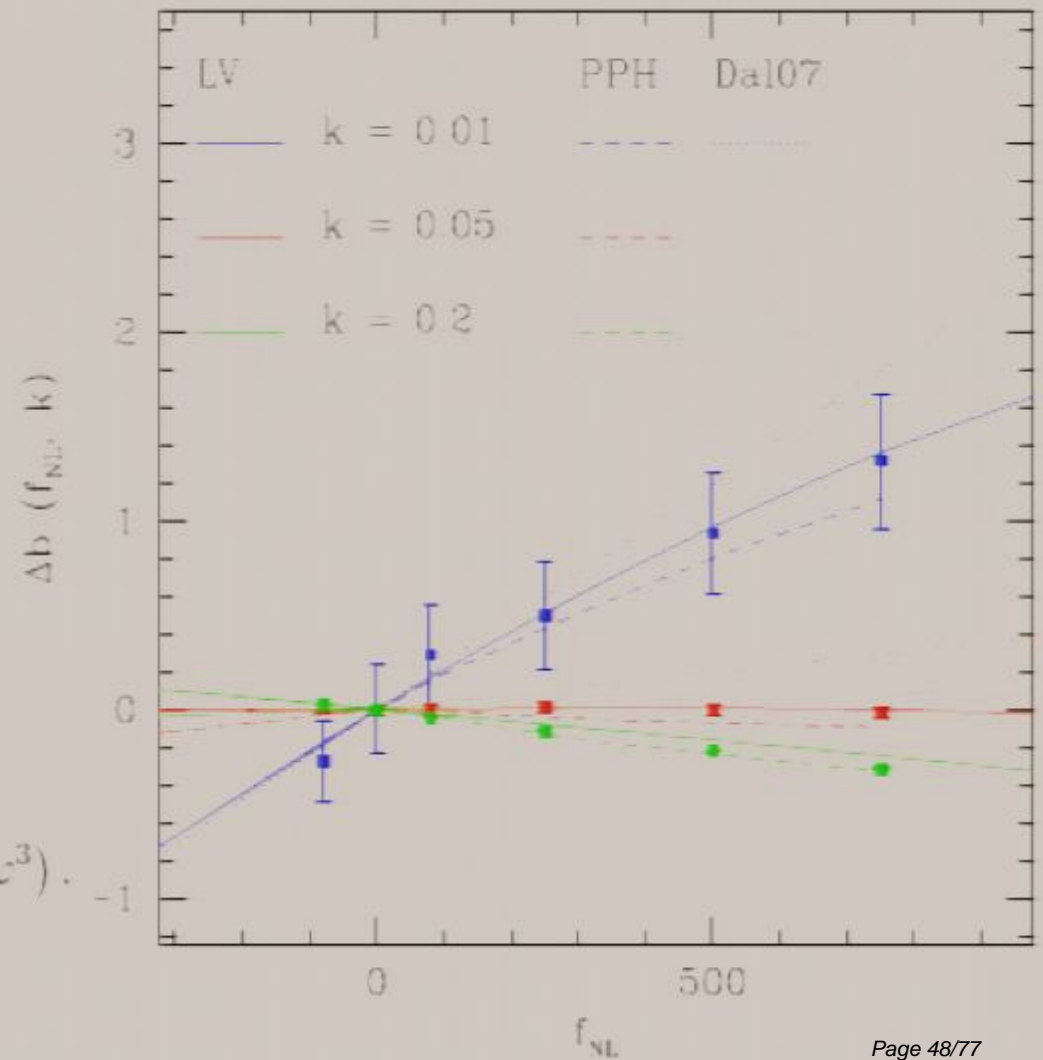
Simulations:



- seems to work like magic!
- from Desjacques et al
- they also see Taruya et al effect and scale indep. Bias.
- Licia Verde and collabs disagree?

Giannantonio & Porciani's magic

$$\begin{aligned} \delta_h(\mathbf{x}) = & b_{10} \delta + b_{01} \varphi + \\ & + \frac{1}{2!} (b_{20} \delta^2 + 2 b_{11} \delta \varphi + b_{02} \varphi^2) + \\ & + \frac{1}{3!} (b_{30} \delta^3 + 3 b_{21} \delta^2 \varphi + 3 b_{12} \delta \varphi^2 + b_{03} \varphi^3). \end{aligned}$$



How to measure it?

$$\Delta b(M, k) = 3f_{\text{NL}}(b - 1)\delta_c \frac{\Omega_m}{k^2 T(k) D(z)} \left(\frac{H_0}{c} \right)^2$$

- Need large volume surveys (small k) of highly biased tracers of structure.
- We use;
 - LRGs (photo, spectro)
 - QSO
 - ISW

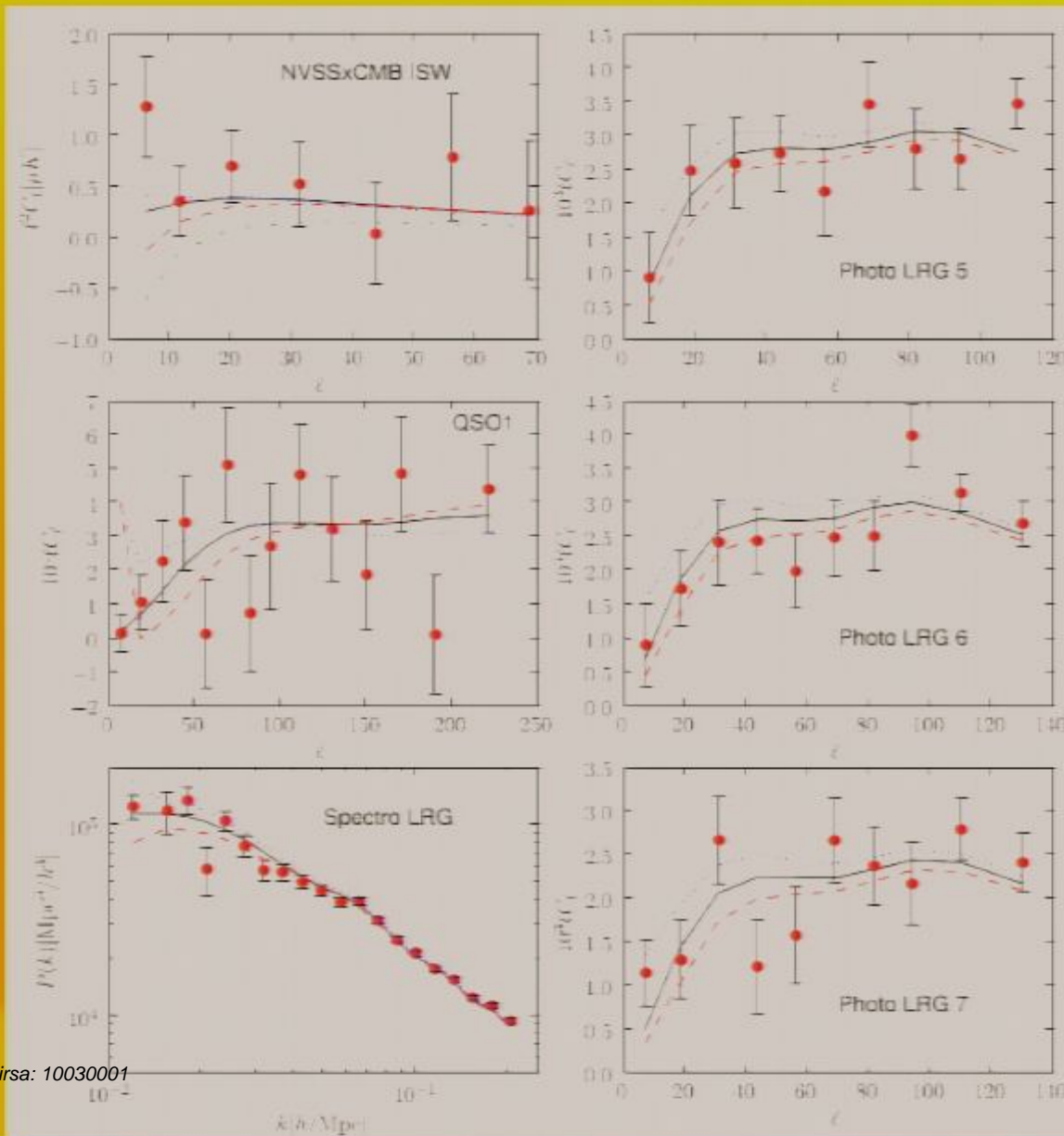
Datasets & Effects

ISW and QSO data same as in Hirata et al. & Ho et al.

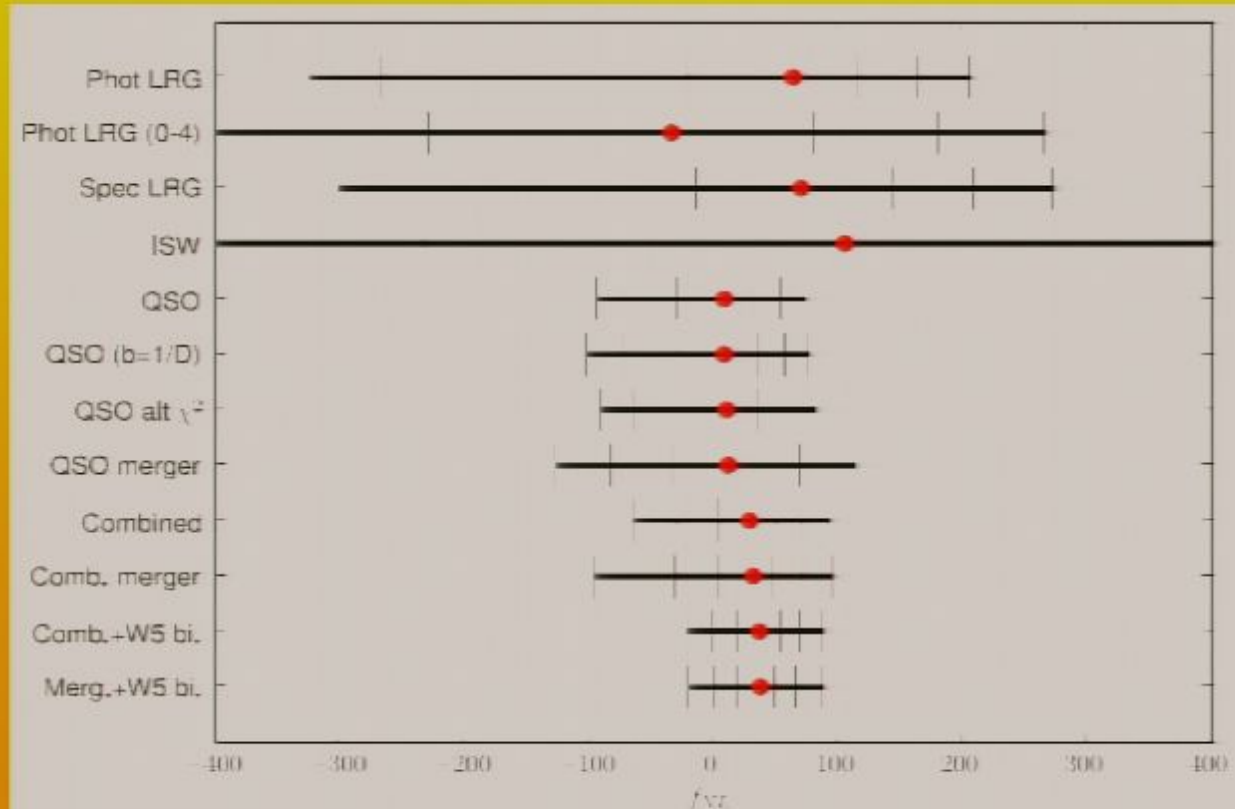
- Spectroscopic Tegmark's LRGs

- Photometric Padmanabhan's LRGs

- We don't find great sensitivity in ISW as opposed to Afshordi



Results (0805.3580):



• Limits:

$$-29 (-65) < f_{NL} < +70 (+93)$$

$$-31 (-96) < f_{NL} < +70 (+96)$$

$$0 (-21) < f_{NL} < +69 (+88).$$

$$-1 < f_{NL}^{local} < 63$$

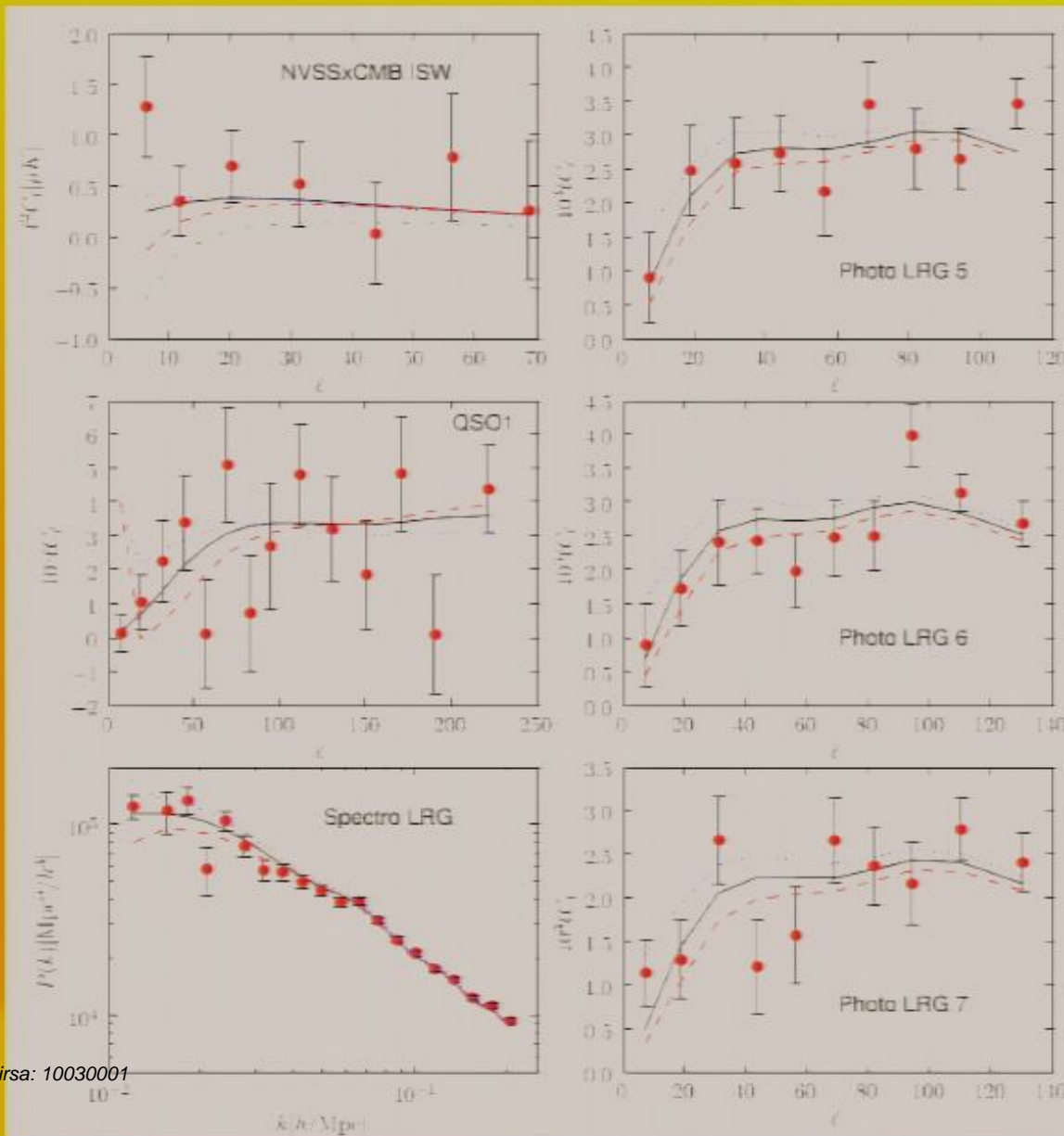
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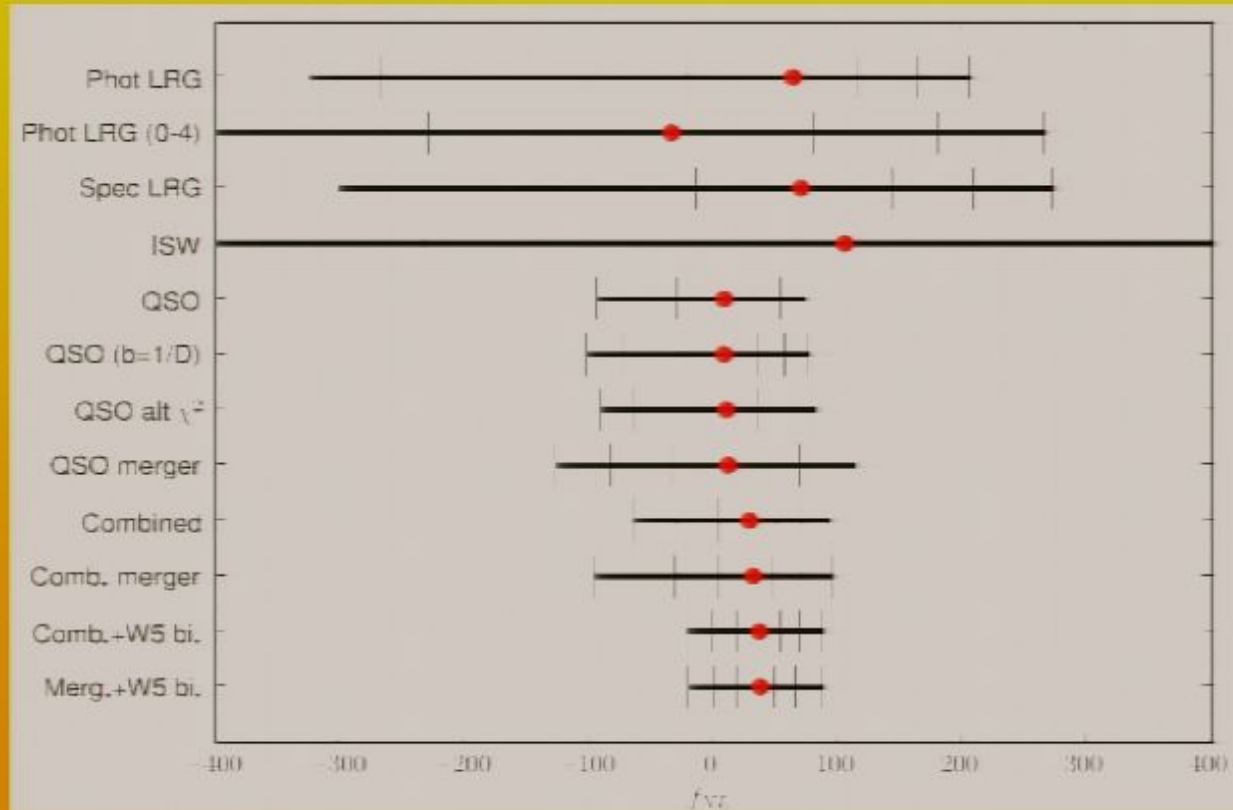
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Gen.: Bispectrum of a tracer

- Jeong + Komatsu et al, 0904.0497:

$$B_h(k_1, k_2, k_3, z) = h_1^3(z) D^4(z) \left[B_m^G(k_1, k_2, k_3) + \frac{b_2(z)}{b_1(z)} \{P_R(k_1)P_R(k_2) + (\text{cyclic})\} + f_{\text{NL}} \frac{B_{f_{\text{NL}}}^{\text{nG}}(k_1, k_2, k_3)}{D(z)} + f_{\text{NL}}^2 \frac{B_{f_{\text{NL}}^2}^{\text{nG}}(k_1, k_2, k_3)}{D^2(z)} + g_{\text{NL}} \frac{B_{g_{\text{NL}}}^{\text{nG}}(k_1, k_2, k_3)}{D^2(z)} \right].$$

- Different "triangle" dependence
- At higher redshift NG goes up, NL goes down

Infrared divergence?

- For small k , $P(k)$ goes as k^{-4} and correlation function diverges. Is this a problem?

$$\xi(r) = \frac{1}{(2\pi)^3} \int d^3k P(k) e^{i\mathbf{k}\cdot\mathbf{r}}.$$

- Measured correlation function always Ok:

$$\langle \tilde{\xi}(r) \rangle = \xi(r) - \sigma_W^2.$$

Infrared divergence?

- Any other GR corrections on large scales?
- Two issues:
 - What do you really measure on super-horizon scales?
 - Yoo et al working on it, sounds promising
 - What is the correct form of the Poisson equation to use when dealing with collapsing objects.
 - Somewhat trivially no effects within spherical collapse (Wands & Slosar), the implicit gauge is Newtonian

Two Fields

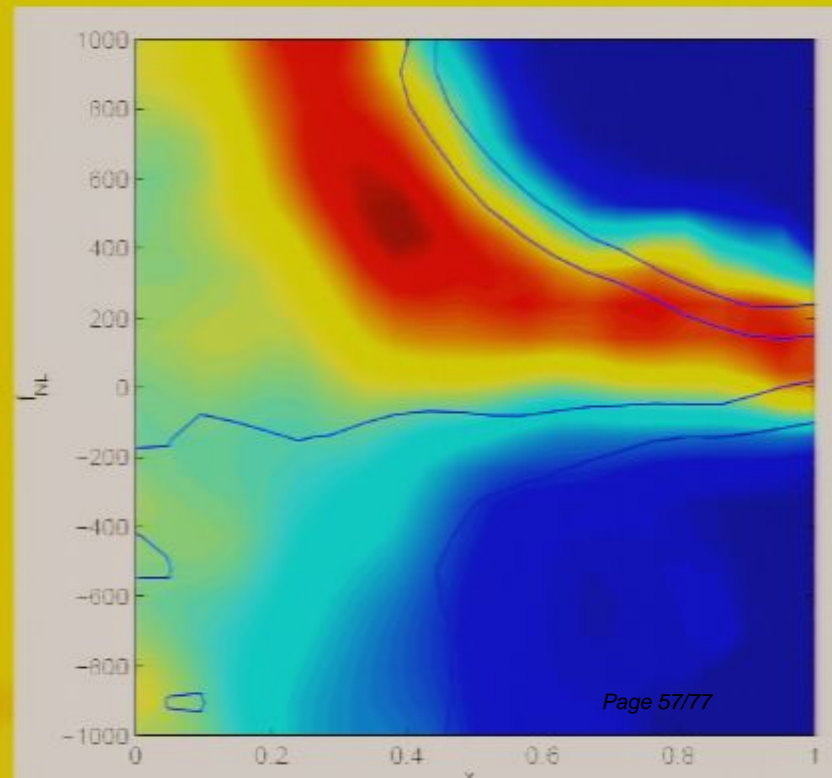
- With Tseliakovich & Hirata

$$\Phi = o_1 + o_2 + \tilde{f}_{\text{NL}} o_2^2.$$

$$P_L^h = b_g^2 \frac{P^{\text{lin}}}{1 + \xi^2} \left(\xi^2 + \left(1 + \frac{2\tilde{f}_{\text{NL}} \alpha^{-1} \delta_c}{1 + \xi^2} \right)^2 \right)$$

$$\xi = \sigma_1 / \sigma_2$$

$$x_2 = 1 / (1 + \xi^2).$$

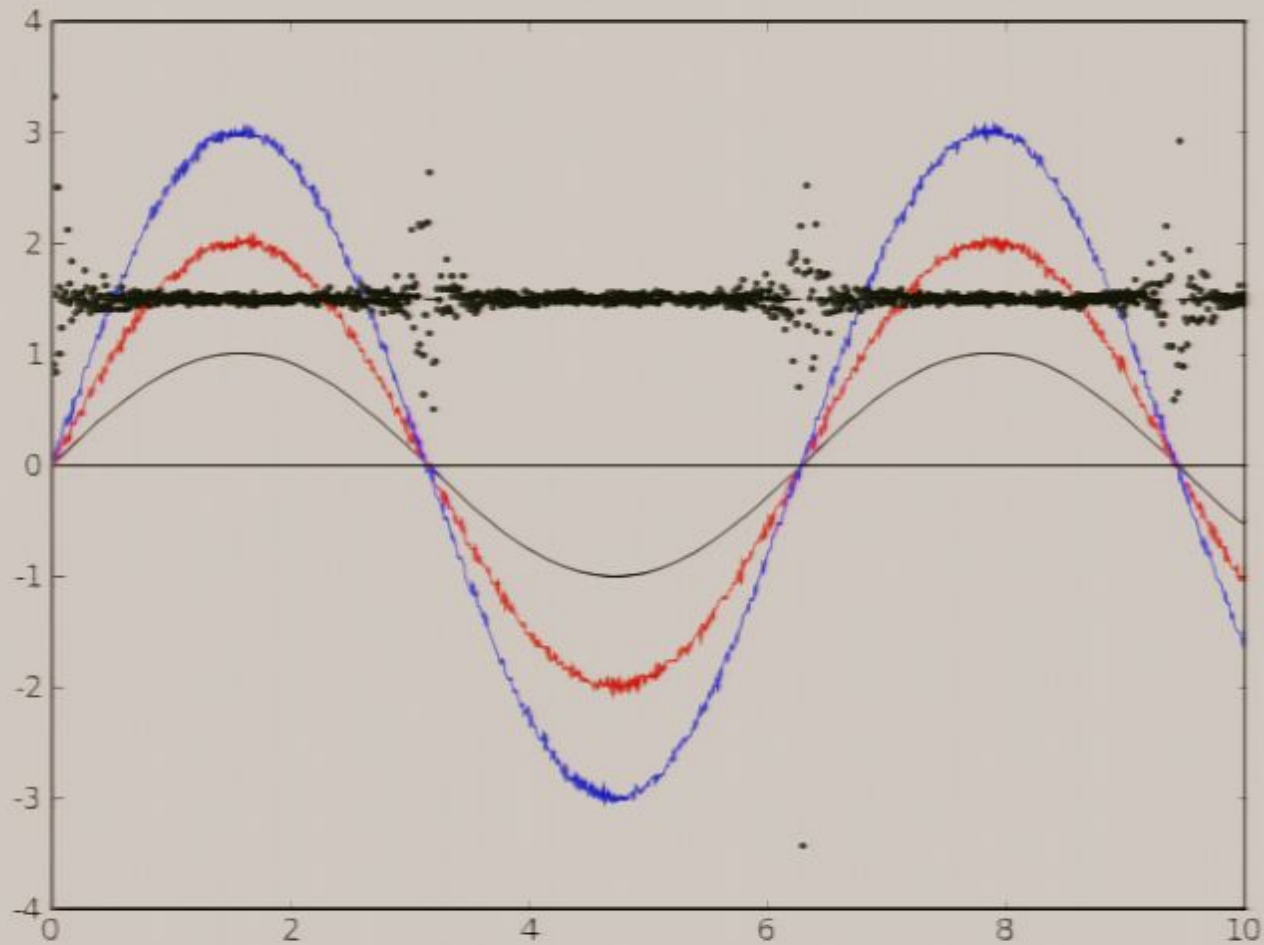


Avoiding sample variance.

- The above method relied on the scale dependent bias detected through power spectrum.
- Large scales \rightarrow fewer modes \rightarrow large sample variance
- Seljak (arxiv:0807.1770) proposed way around it
- Measure the same mode using two tracers with different intrinsic bias
- In the ratio, the mode amplitude cancels and one remains with bias ratio

• (useless at the moment as we're still Poisson limited)

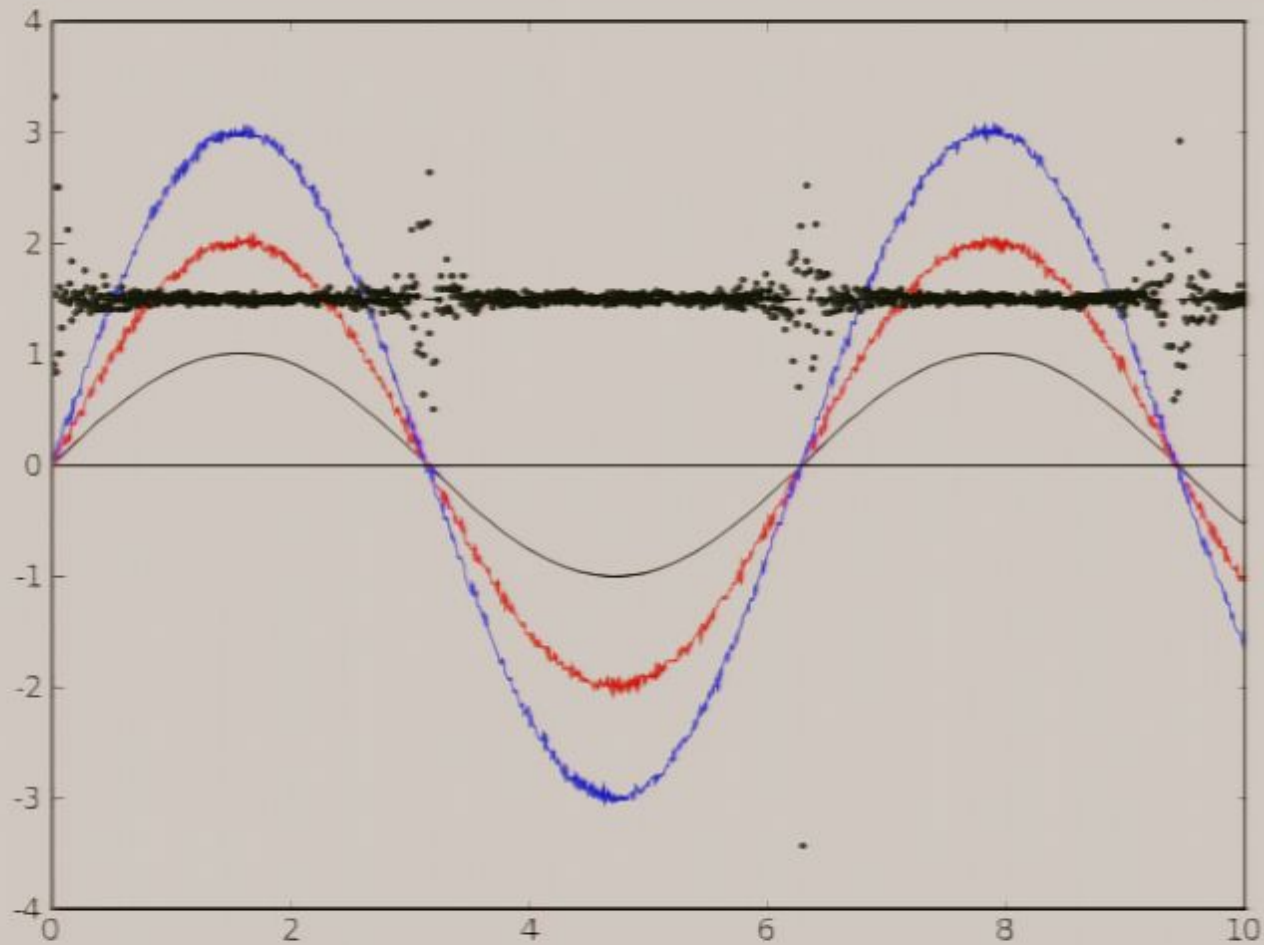
Avoiding sample variance.



FNL with Lyman-alpha forest

- Dalal's effect works, because the mass function is exponentially sensitive to PS normalisation.
- In Lyman-alpha forest, flux is also exponentially sensitive to delta under (over?) simplifying assumptions

Avoiding sample variance.

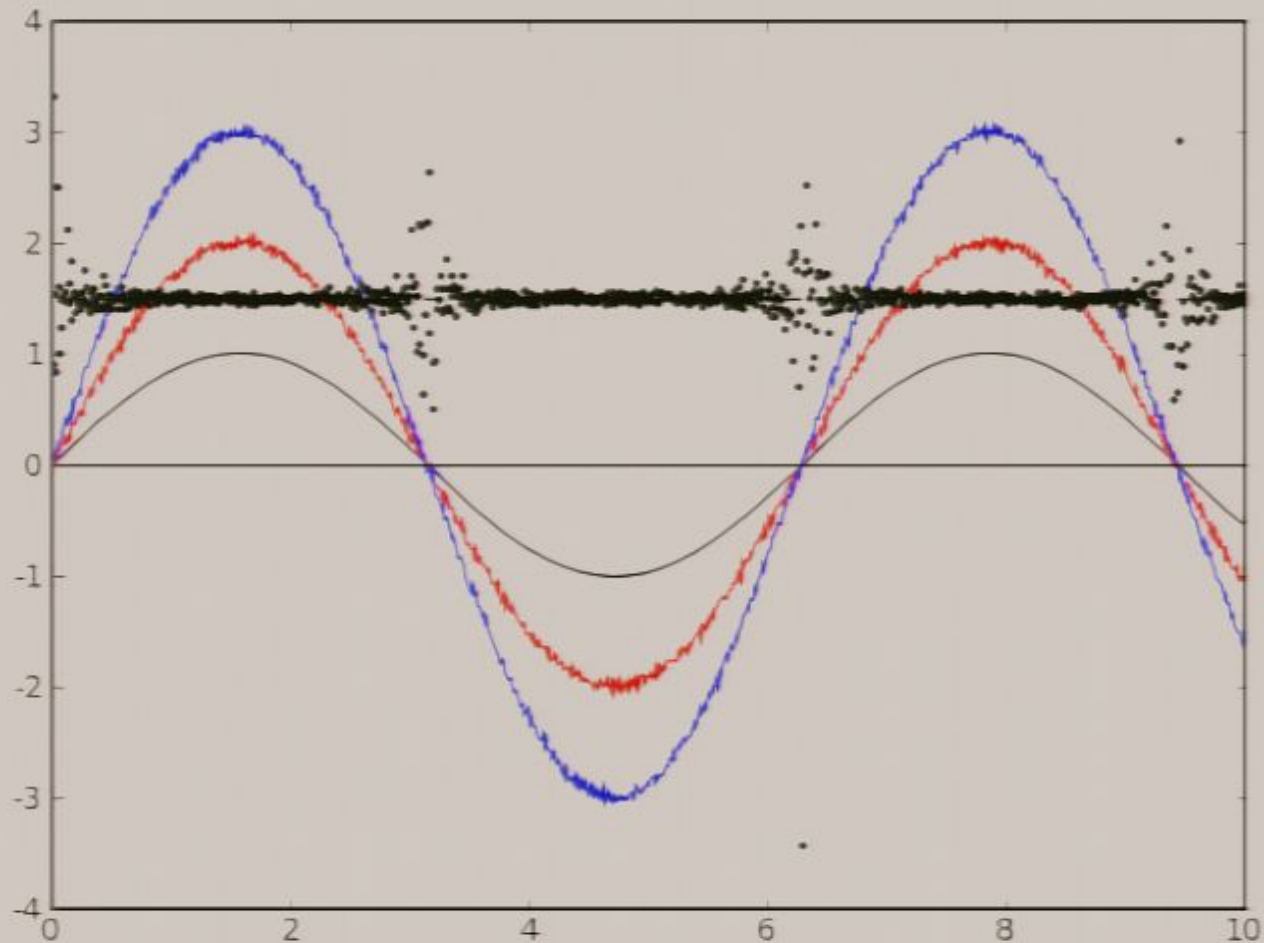


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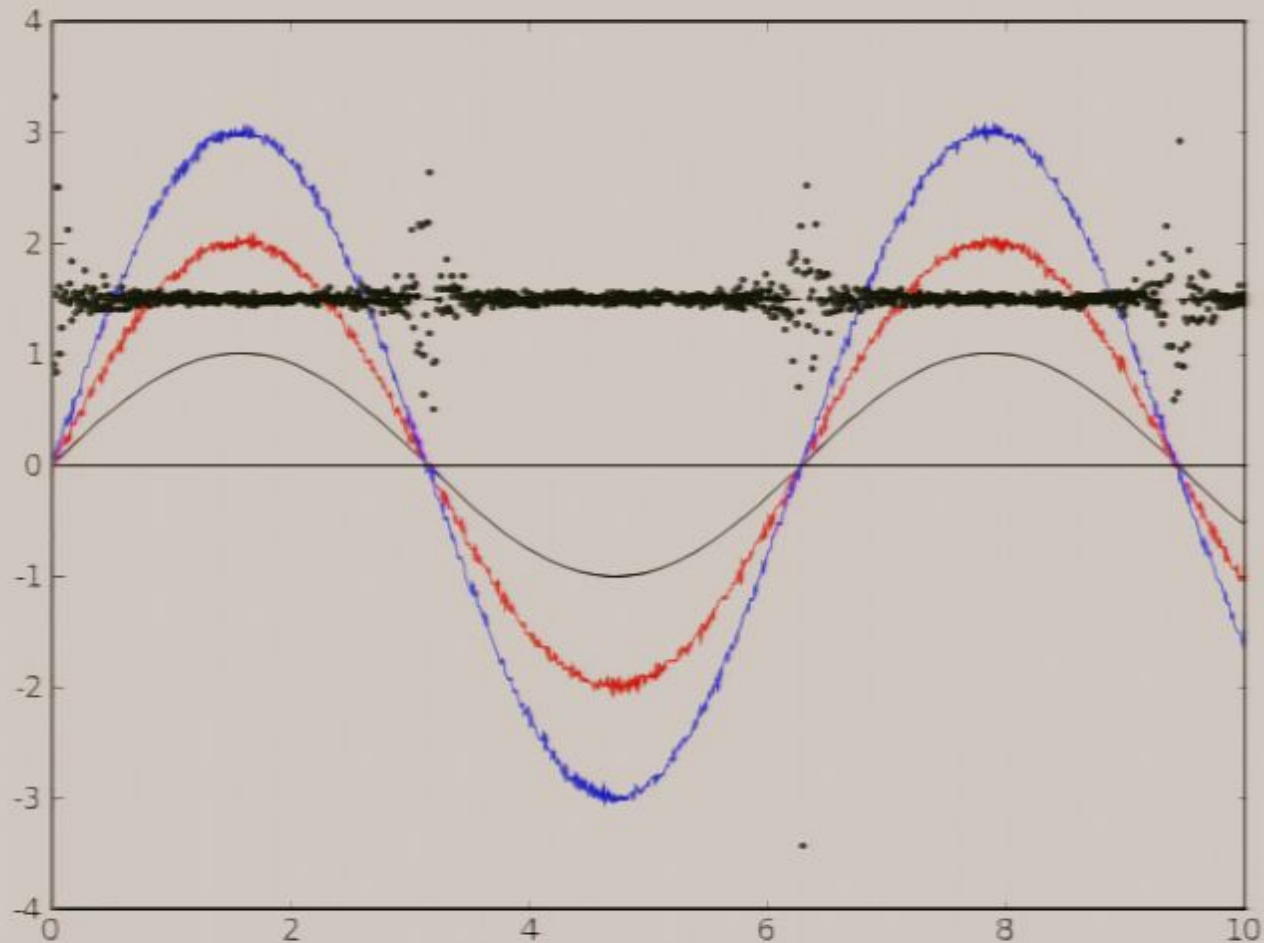
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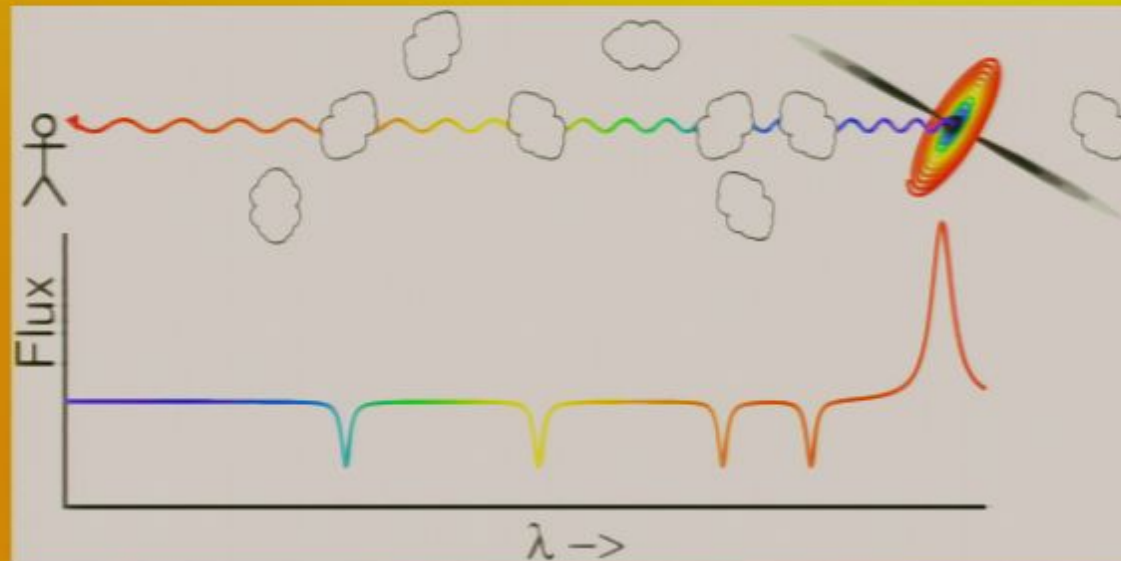
- Final errors now dominated by the Poisson error, that are subdominant on large scales.
- In other words, replace large scale mode counting with small scale mode counting – plenty of small scale modes.
- In principle a single mode is good enough. In practice, measure two power spectra and take ratio
- Degeneracies with other parameters also disappear
- Projections: with SNAP-like survey one gets to $f_{nl} \sim 1$

Avoiding sample variance.



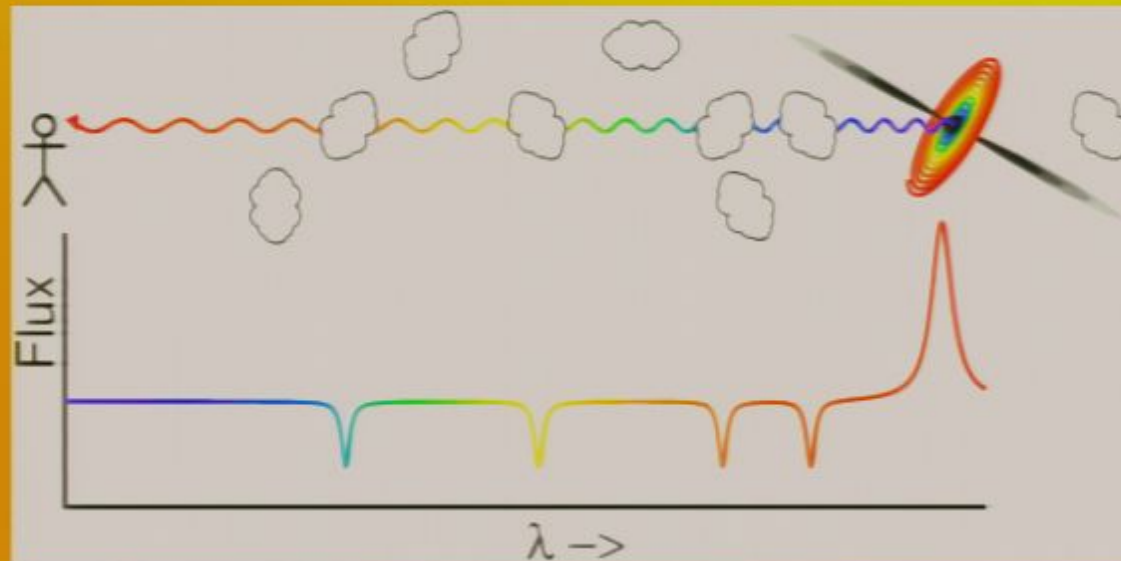
Lyman-alpha forest

- clouds of hydrogen absorb light from distant quasars, blueward of Lyman-alpha emission



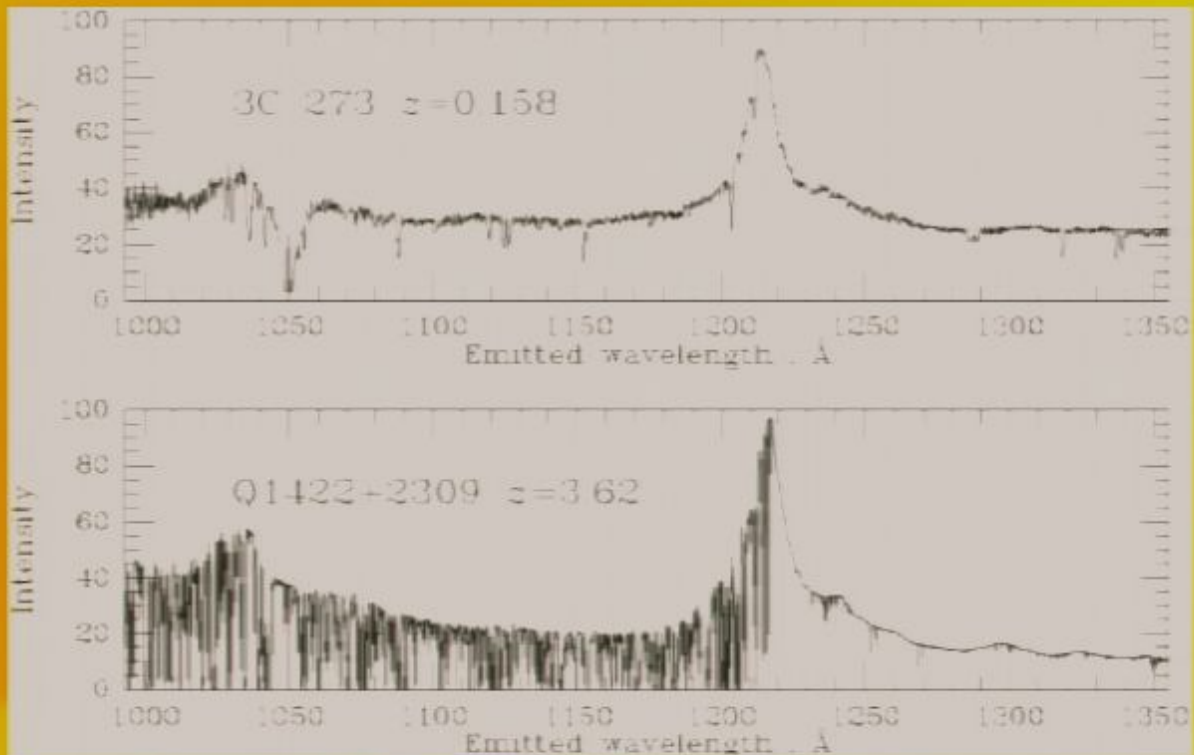
Lyman-alpha forest

- clouds of hydrogen absorb light from distant quasars, blueward of Lyman-alpha emission



Lyman-alpha forest

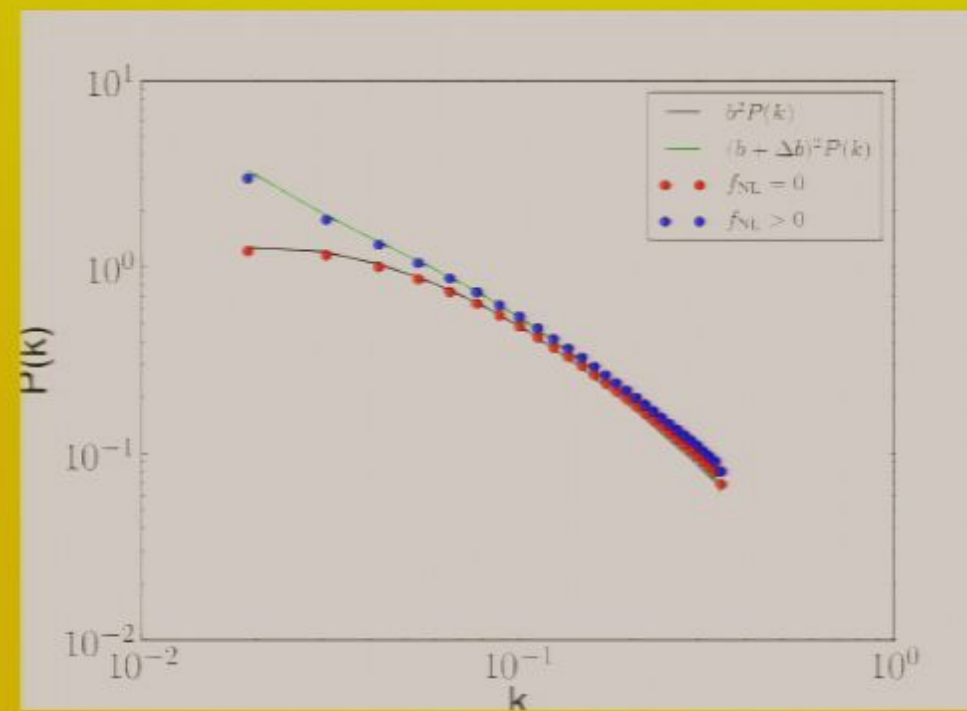
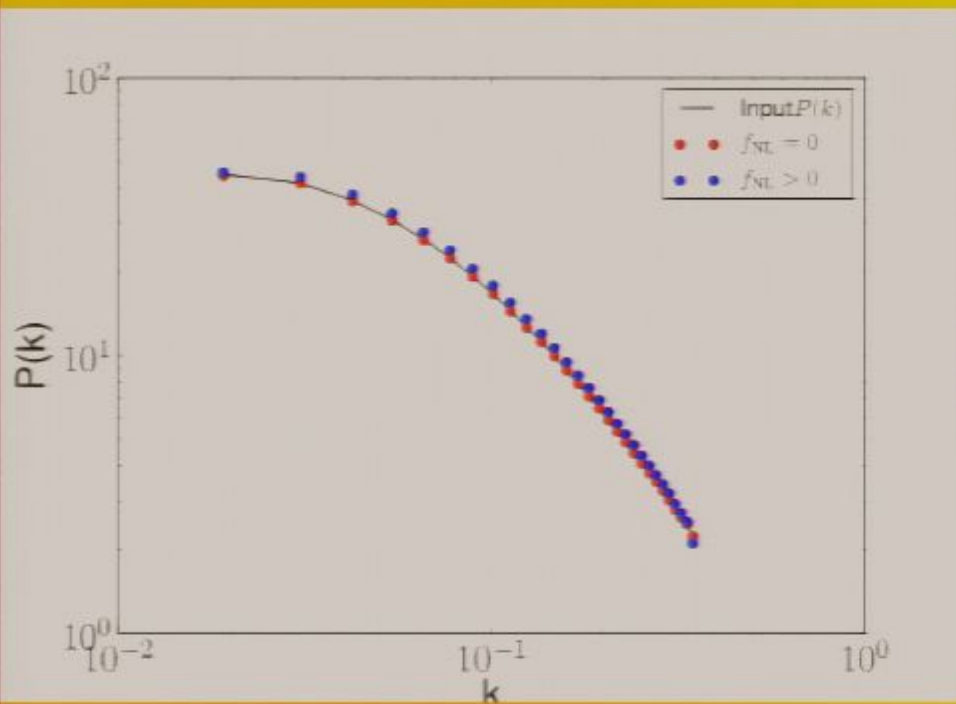
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Assuming photo-ionization equilibrium:

$$\tau \propto (1 + \delta)^\beta$$
$$f = e^{-\tau}$$

Results using Gaussian fields

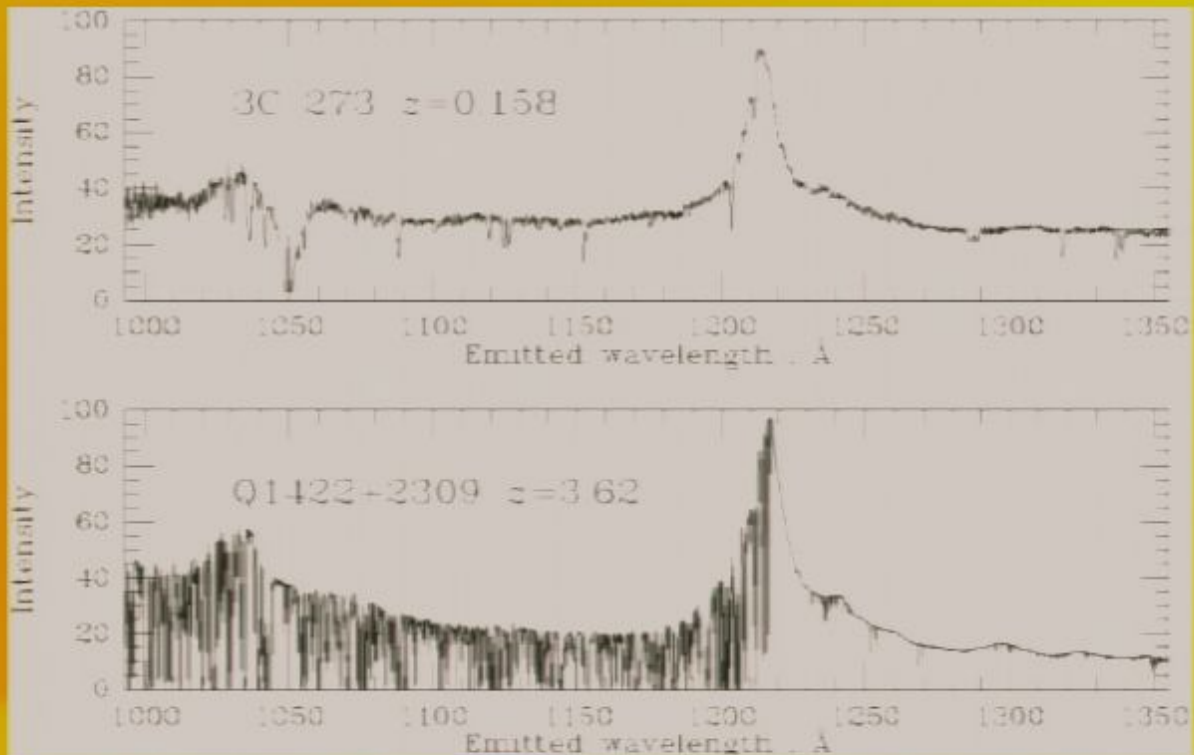


Gaussian field converted into fluxes using FGPA

Solid lines are simple analytical results using P -bg split

Lyman-alpha forest

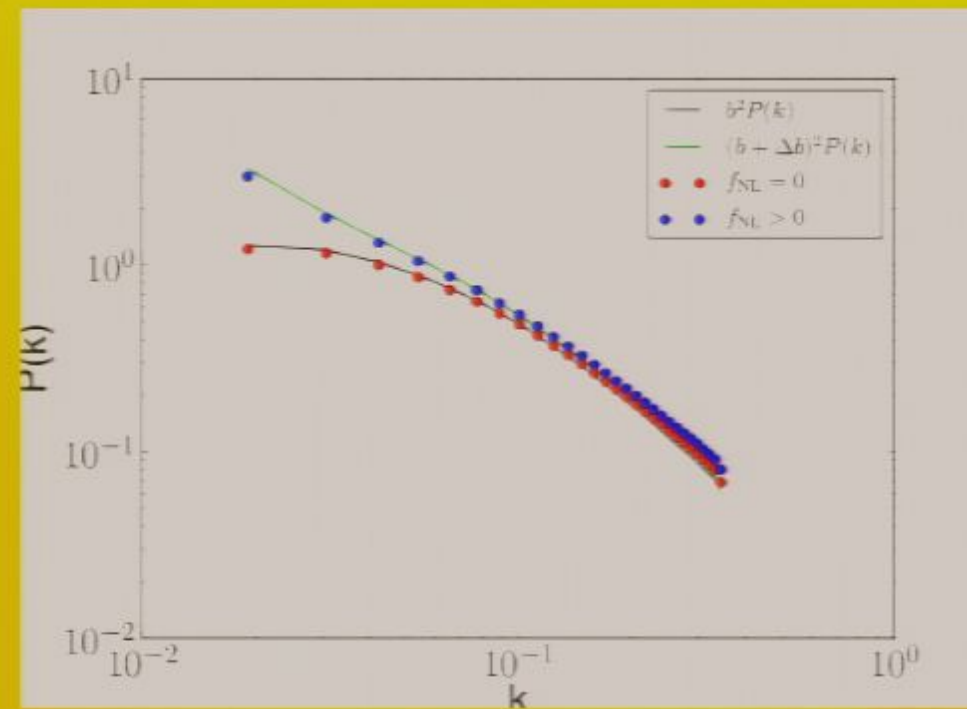
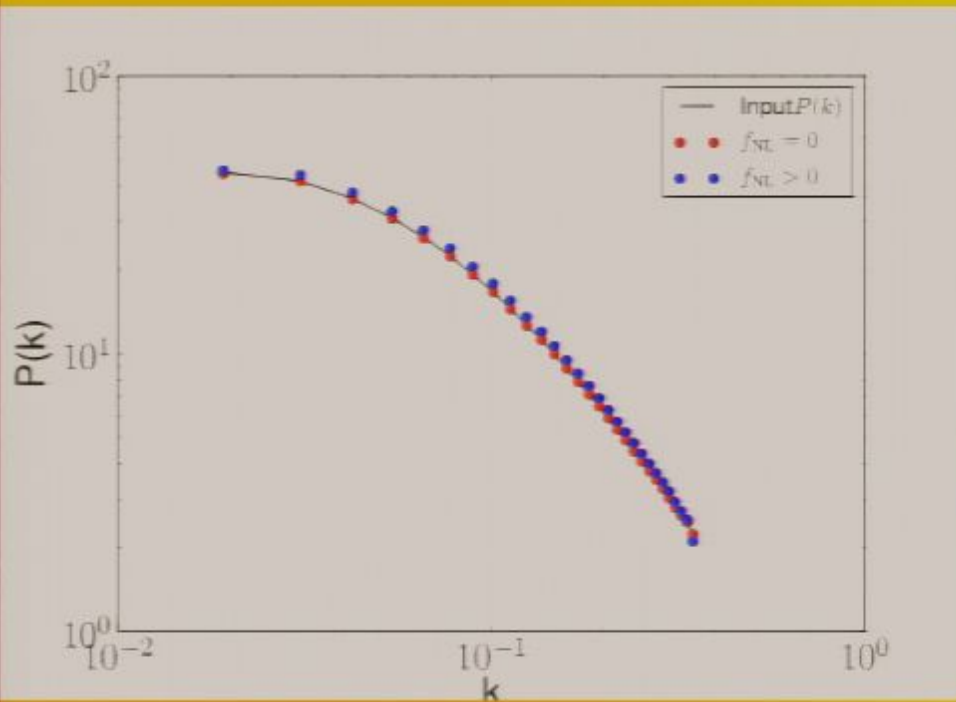
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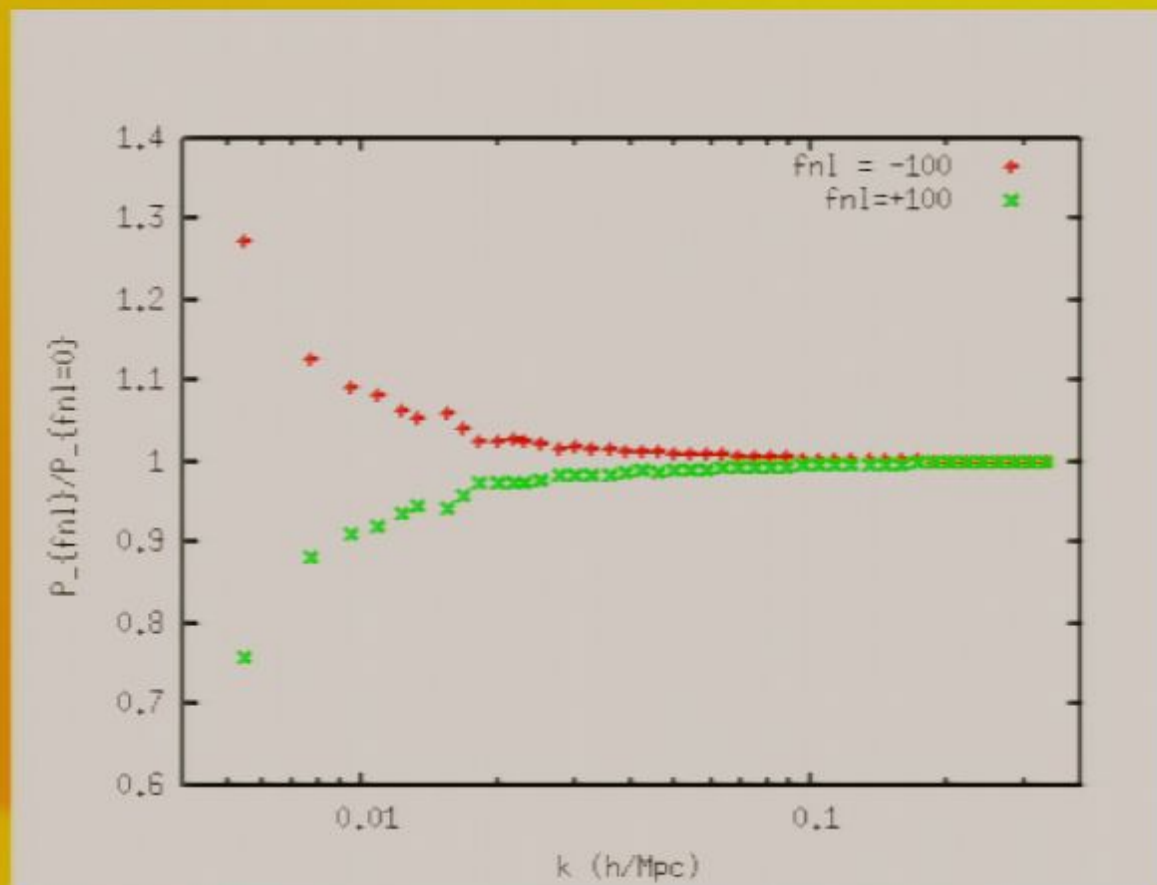
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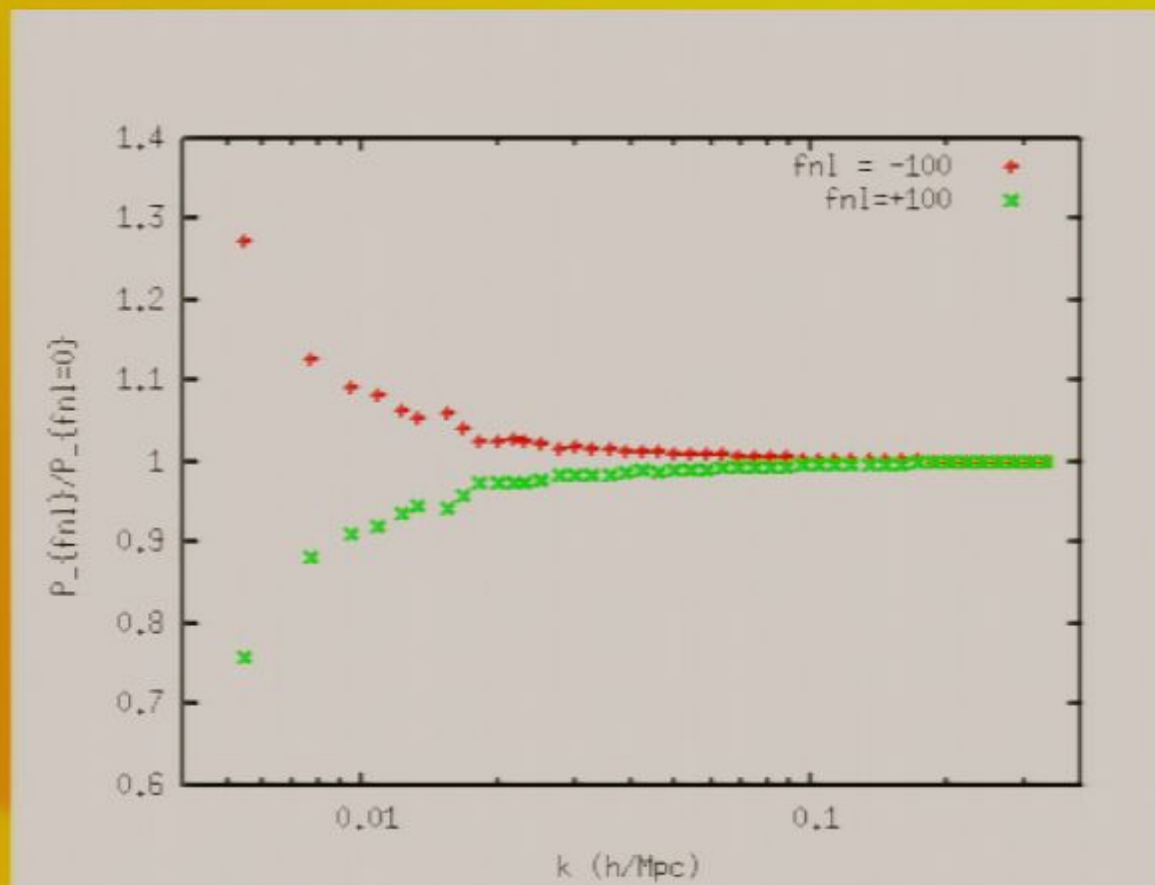
N-body sims (still FGPA)



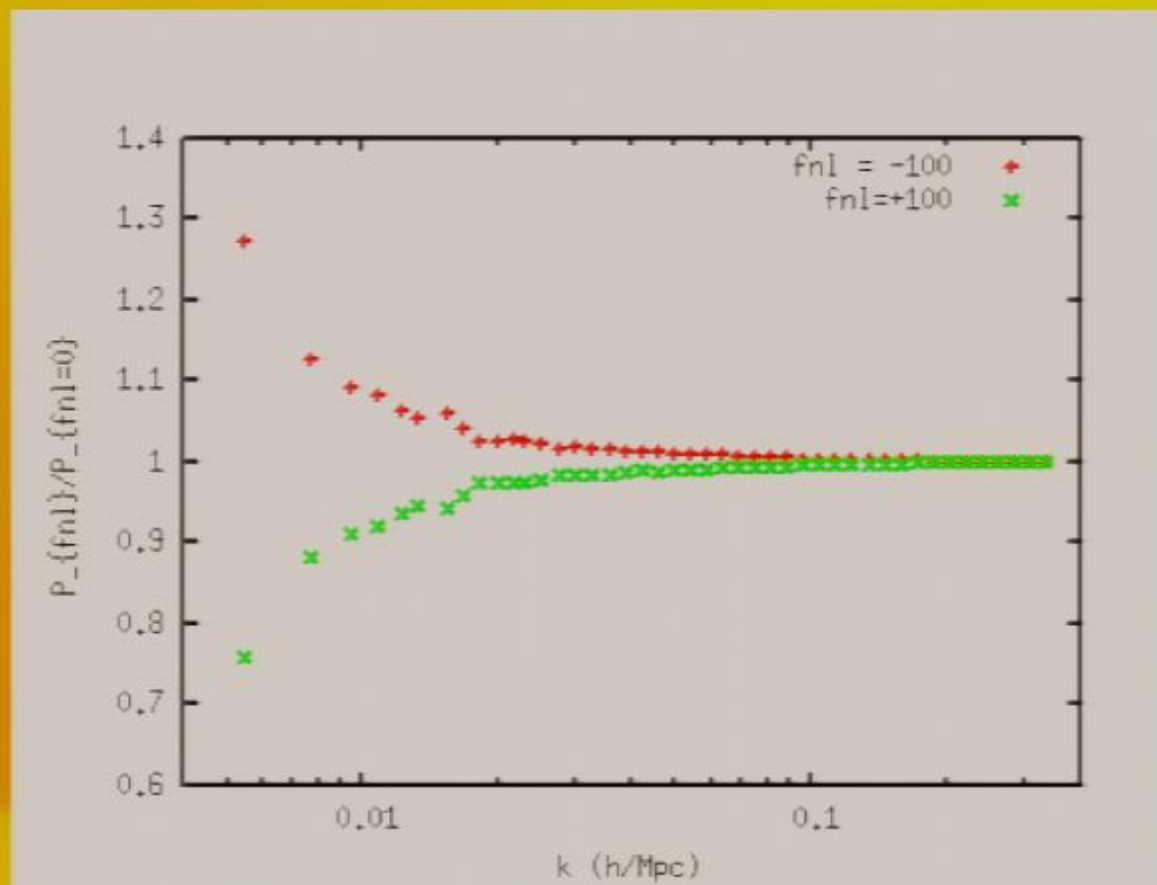
Conclusions.

- Still many checks to be done:
 - Currently even power spectrum is suboptimal with the present data: can do better.
 - simulations: seems to work very well
 - Future data such as LSST, BigBoss can push down a lot to fnl (a few)
 - Fnl from Lyman-alpha forest might or might not work
- Should remain very active field
- N_{eff} might be our best hope to constrain inflation

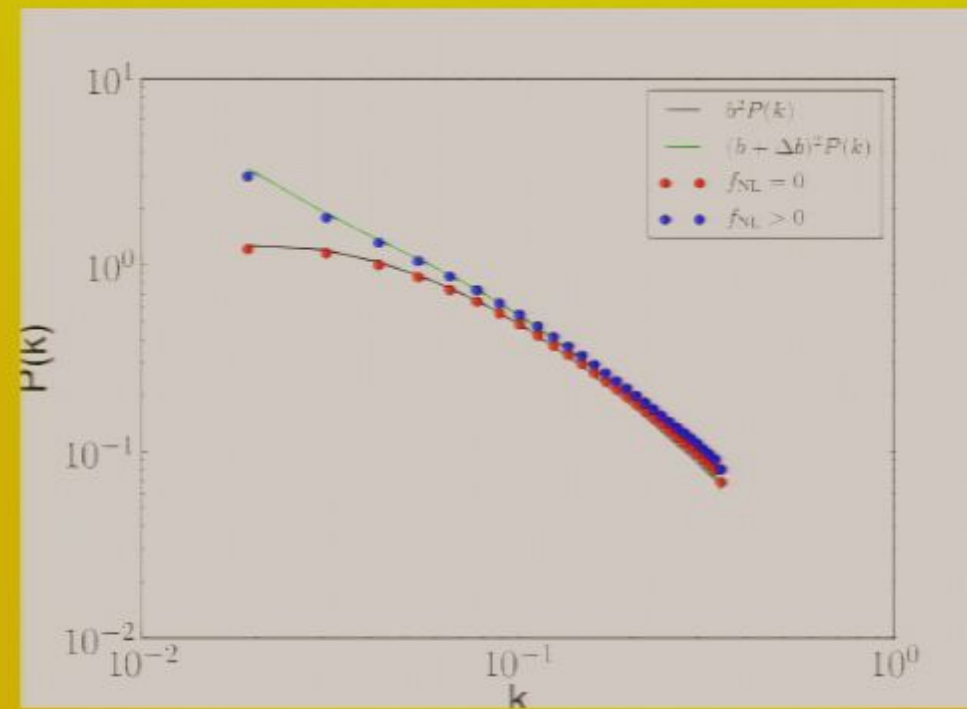
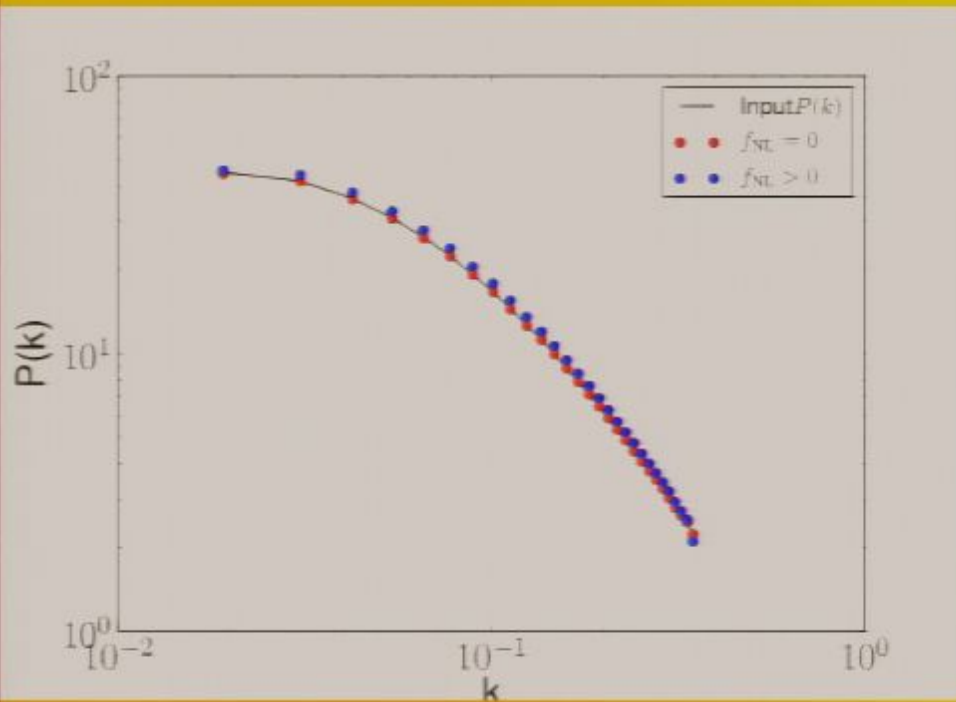
N-body sims (still FGPA)



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Results using Gaussian fields



Gaussian field converted into fluxes using FGPA

Solid lines are simple analytical results using P -bg split

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