

Title: Condensed Matter II - Lecture 7

Date: Feb 24, 2010 10:10 AM

URL: <http://pirsa.org/10020103>

Abstract:

$$G_i = \frac{P_i}{m}$$

Non-int Bose Gas
Condensed in
state where

$$\vec{P}_i = \frac{\vec{P}}{N} \quad i=1, \dots, N$$

$$E_i = \frac{p_i^2}{2m}$$

Non-int Bose Gas
Condensed in
state where

$$\vec{p}_i = \frac{\vec{P}}{N} \quad i=1, \dots, N$$

$$E = \sum_i \frac{p_i^2}{2m} = \frac{P^2}{2mN}$$

$$\vec{v} = \frac{\vec{p}_c}{m}$$

Non-int Bose Gas
Condensed in
state where

$$\vec{p}_i = \frac{\vec{P}}{N} \quad i=1, \dots, N$$

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Consider an
excitation of one particle $\vec{p}_i \rightarrow \vec{p}_i + \Delta\vec{p}_i$

$$\vec{v} = \frac{\vec{p}_c}{m}$$

Non-int Bose Gas
Condensed in
state where

$$\vec{p}_i = \frac{\vec{P}}{N} \quad i=1, \dots, N$$

$$E = \sum_i \frac{p_i^2}{2m} = \frac{P^2}{2mN}$$

Consider an
excitation of one particle $\vec{p}_i \rightarrow \vec{p}_i + \Delta\vec{p}_i$

$$\therefore \Delta E = \frac{(\vec{p}_i + \Delta\vec{p}_i)^2}{2m} - \frac{p_i^2}{2m} = \frac{\vec{p}_i \cdot \Delta\vec{p}_i}{m} + \frac{(\Delta\vec{p}_i)^2}{2m}$$

(repulsive)
Weakly interacting Bose gas

(repulsive)

Weakly interacting Bose gas

$$\text{then } E(\vec{k}) = \hbar c k$$

$$c = \frac{\hbar}{m} \sqrt{4\pi a n}$$

S-wave scattering length

den of

(repulsive)

Weakly interacting Bose gas

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$$c = \frac{\hbar}{m} \sqrt{4\pi a n}$$

S-wave scattering length

density of particles

(repulsive)

Weakly interacting Bose gas

$$\text{then } \epsilon(\vec{k}) = \hbar c k$$

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S-wave scattering length

density of particles

(repulsive)

Weakly interacting Bose gas

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$$[c = \frac{\hbar}{m} \sqrt{4\pi a n}]$$

S-wave
scattering
length

density
of particles

For a uniformly flowing gas

$$E(\vec{k}) = \hbar c k + \hbar \vec{k} \cdot \vec{v}_s$$

(repulsive)

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For a uniformly flowing gas

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Critical velocity

$$|\vec{v}_s| = c$$

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S-wave scattering length

density of particles

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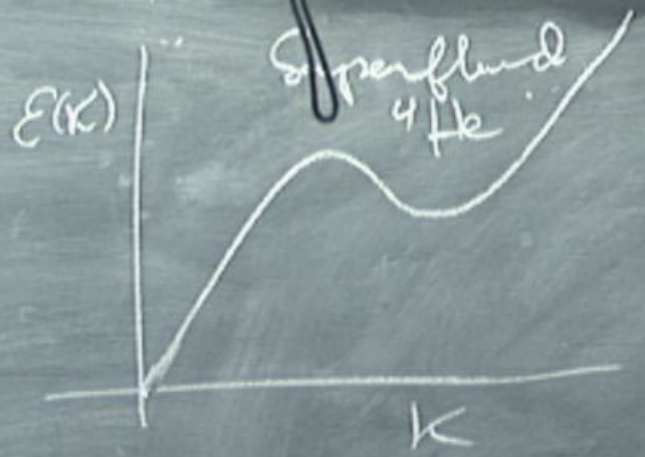
Critical velocity

$$|\vec{v}_s| = c$$

For

E

density
of particles



For a Superconductor

$$E(\vec{k}) = \hbar \vec{k} \cdot \vec{v}_S + \sqrt{\Delta^2 + \xi^2(\vec{k})}$$

electron
KE rel

density of particles



For a Superconductor

$$E(\vec{k}) = \hbar \vec{k} \cdot \vec{v}_s + \sqrt{\Delta^2 + \xi^2(\vec{k})}$$

electron KE referred to EF

superconducting gap $|\Delta|$

Specific Heat Jump

C_V



Specific Heat Jump

C_V



density
of particles



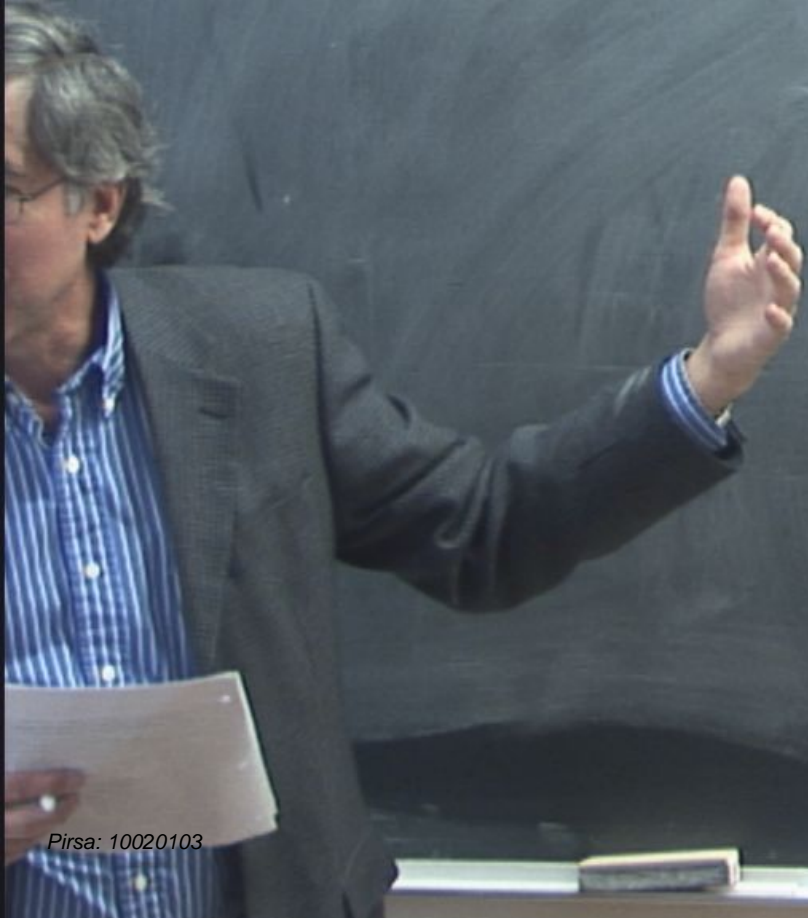
electron
bands

For a Superconductor

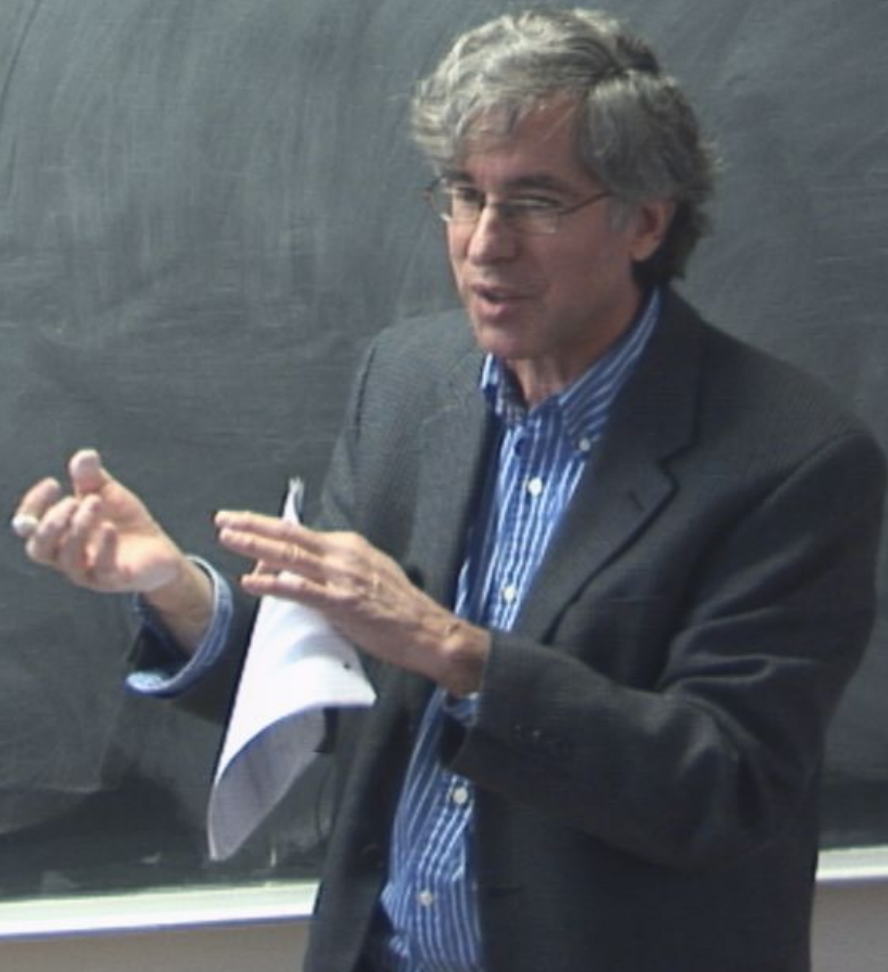
$$E(\vec{k}) = \hbar \vec{k} \cdot \vec{v}_s + \sqrt{\Delta^2 + \xi^2(\vec{k})}$$

perco

(non-linear) $\psi = \psi_0 + \psi_1$



$$H_{ph} = \sum_{\vec{s}, \lambda} \omega_{\vec{s}, \lambda} \left(b_{\vec{s}, \lambda}^\dagger b_{\vec{s}, \lambda} + \frac{1}{2} \right)$$



$$H_{ph} = \sum_{\vec{\xi}, \lambda} \omega_{\vec{\xi}, \lambda} (b_{\vec{\xi}, \lambda}^{\dagger} b_{\vec{\xi}, \lambda} + \frac{1}{2})$$

$$H_{el-ph} = \frac{1}{N} \sum_{\vec{k}, \sigma} \sum_{\vec{\xi}, \lambda} g_{\vec{\xi}, \lambda} C_{\vec{k} + \vec{\xi}, \sigma}^{\dagger} C_{\vec{k}, \sigma} (b_{\vec{\xi}, \lambda}^{\dagger} + b_{-\vec{\xi}, \lambda})$$

$$H_{ph} = \sum_{\vec{\xi}, \lambda} \omega_{\vec{\xi}, \lambda} (b_{\vec{\xi}, \lambda}^{\dagger} b_{\vec{\xi}, \lambda} + \frac{1}{2})$$

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$$H_{el-ph} = \frac{1}{\Omega} \sum_{\vec{k}, \sigma} \sum_{\vec{\xi}, \lambda} g_{\vec{\xi}, \lambda} C_{\vec{k}-\vec{\xi}, \sigma}^\dagger C_{\vec{k}, \sigma} (b_{\vec{\xi}, \lambda}^\dagger + b_{-\vec{\xi}, \lambda})$$



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Effective Int

$$V_{\text{eff}}(\vec{g}) = |g(\vec{g})|^2 \left\{ \frac{1}{\epsilon(\mathbf{k}+\vec{g}) - \epsilon(\mathbf{k}) - \omega_g} \right.$$

$$\left. - \frac{1}{\epsilon(\mathbf{k}+\vec{g}) - \epsilon(\mathbf{k}) + \omega_g} \right\}$$

$$= |g(\vec{g})|^2 \frac{2\omega_g}{\omega^2 - \omega_g^2} \quad \omega \equiv \epsilon(\mathbf{k}+\vec{g})$$

Effective Int

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$$= |g(\vec{g})|^2 \frac{2\omega_{\vec{g}}}{\omega^2 - \omega_{\vec{g}}^2} \quad \omega \equiv \epsilon(\vec{k}+\vec{g}) - \epsilon(\vec{k})$$

$\hbar \rho \hbar$

$\hbar \omega_{\vec{g}}$

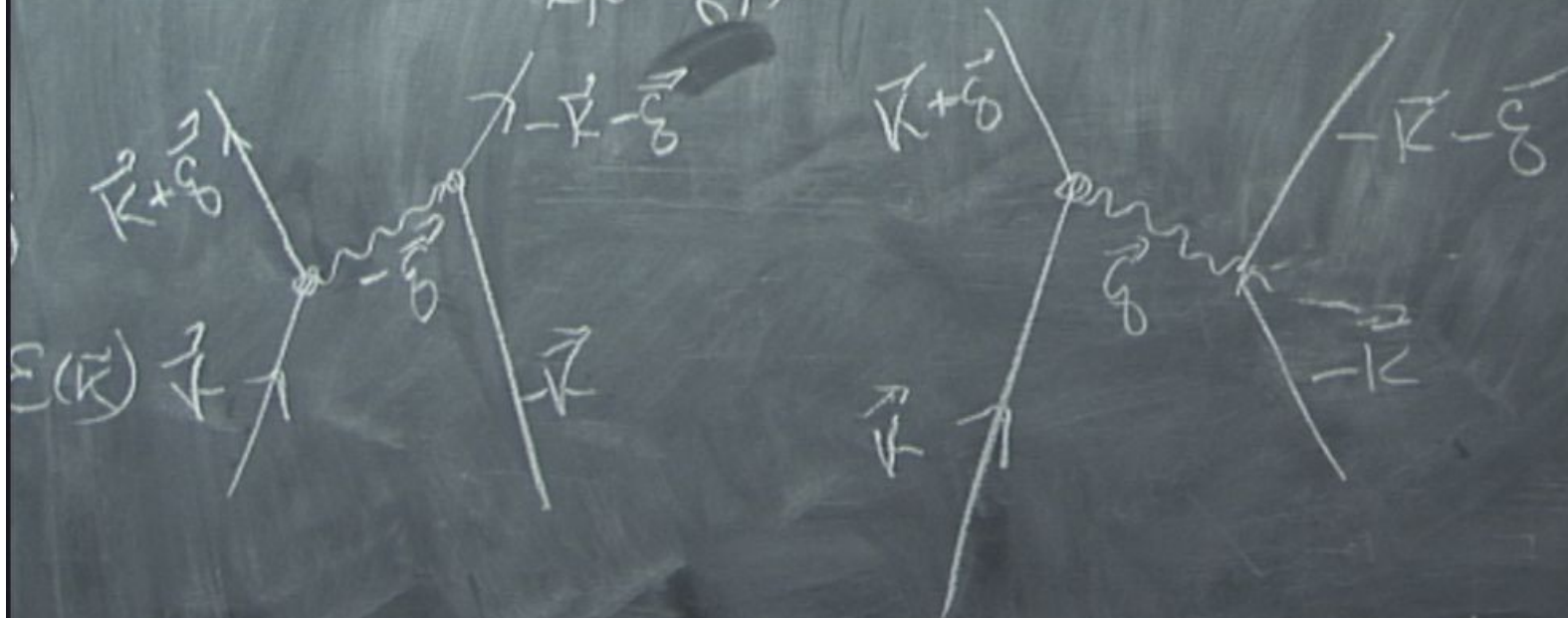
$\vec{k} + \vec{g}$

$\epsilon(\vec{k}+\vec{g}) - \epsilon(\vec{k})$

$$H_{ph} = \sum_{\vec{\xi}, \lambda} \omega_{\vec{\xi}, \lambda} \left(b_{\vec{\xi}, \lambda}^\dagger b_{\vec{\xi}, \lambda} + \frac{1}{2} \right)$$



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Effective Int

$$V_{\text{eff}}(\vec{g}) = |g(\vec{g})|^2 \left\{ \frac{1}{\epsilon(\kappa+\vec{g}) - \epsilon(\kappa) - \omega_{\vec{g}}} \right.$$

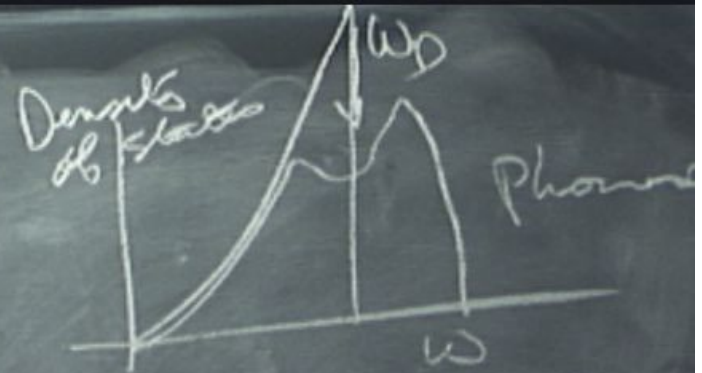
$$\left. - \frac{1}{\epsilon(\kappa+\vec{g}) - \epsilon(\kappa) + \omega_{\vec{g}}} \right\}$$

$$= |g(\vec{g})|^2 \frac{2\omega_{\vec{g}}}{\omega^2 - \omega_{\vec{g}}^2} \quad \omega \equiv \epsilon(\kappa+\vec{g})$$

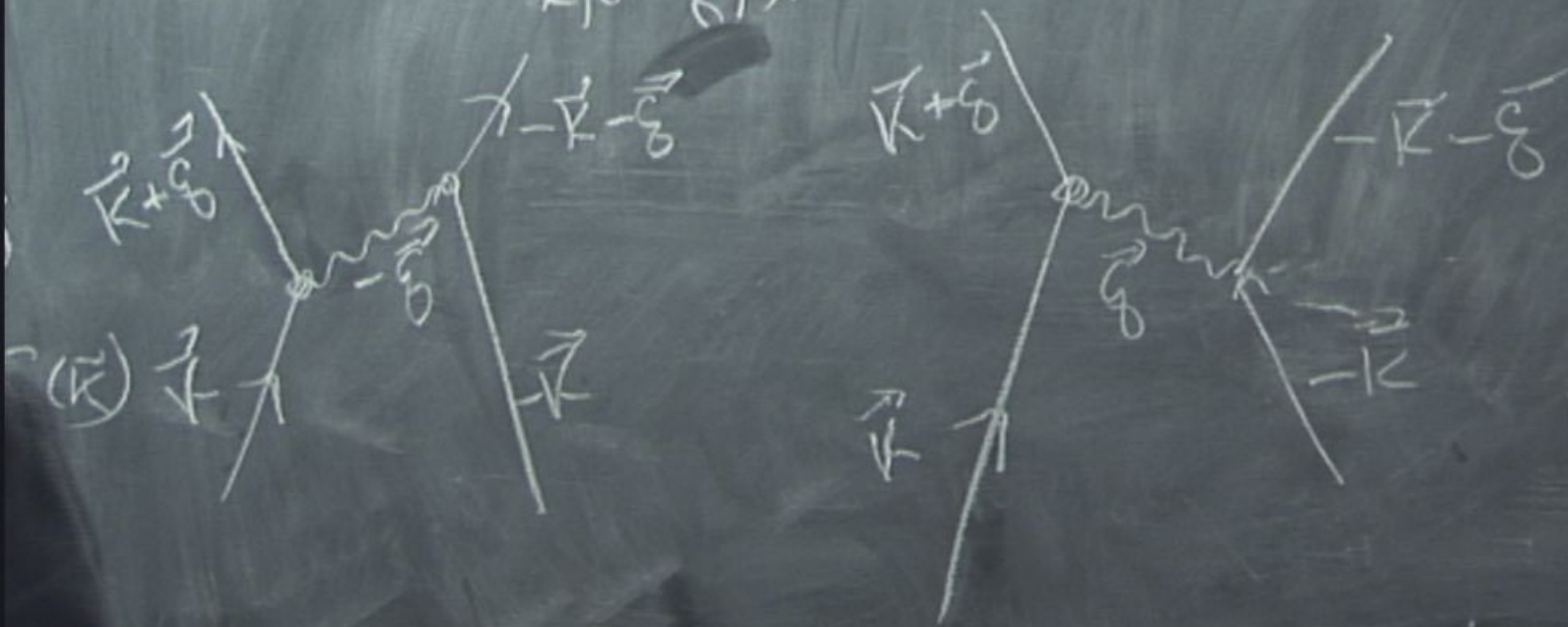
Simplifying assumption
of BCS

$$V_{\text{eff}} = -\frac{V}{N} \quad \text{if } \epsilon(\kappa) \text{ and } \epsilon(\kappa+\vec{g}) \text{ are within } \omega_D \text{ of } E_F \\ = 0 \text{ otherwise}$$

$$H_{ph} = \sum_{\vec{\xi}, \lambda} \omega_{\vec{\xi}, \lambda} (b_{\vec{\xi}, \lambda}^\dagger b_{\vec{\xi}, \lambda} + \frac{1}{2})$$



$$H_{el-ph} = \frac{1}{N} \sum_{\vec{k}, \sigma} \sum_{\vec{\xi}, \lambda} g_{\vec{\xi}, \lambda} C_{\vec{k}-\vec{\xi}, \sigma}^\dagger C_{\vec{k}, \sigma} (b_{\vec{\xi}, \lambda}^\dagger + b_{-\vec{\xi}, \lambda})$$



Specific Heat Jump

C_V



Cooper Pairing
Cooper considered the
problem

