

Title: Explorations in Quantum Info. (PHYS 641) - Lecture 5

Date: Feb 22, 2010 09:00 AM

URL: <http://pirsa.org/10020088>

Abstract:

spin $\frac{1}{2}$,

spin $\frac{1}{2}$,
control spin,
super mirror.
 $n = \{$

spin $\frac{1}{2}$,
control spin,
super mirror.
$$n = \left[1 - \frac{v(r)}{E} \right]^{1/2}$$

spin $1/2$,
control spin,
super mirror.

$$n = \left[1 - \frac{v(r)}{E} \right]^{1/2}$$

$$V_{\text{mag}} = \mu \cdot B$$

spin $\frac{1}{2}$,
control spin,
super mirror.

$$n = \left[1 - \frac{V(r)}{E} \right]^{1/2} \sim 1 - 10^{-5}$$

$$V_{\text{mag}} = \mu \cdot B$$

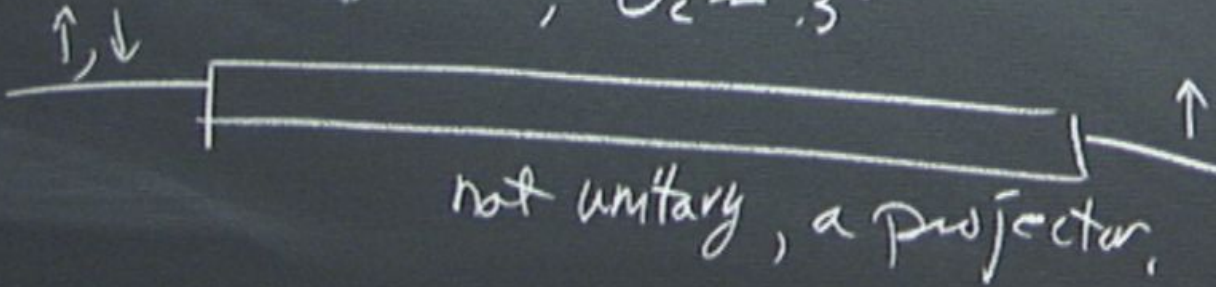
$$n_+ = n_-$$

spin $\frac{1}{2}$,
control spin,
super mirror.

$$n = \left[1 - \frac{V(\omega)}{E} \right]^{\frac{1}{2}} \sim 1 - 10^{-5}$$

$$V_{\text{mag}} = \mu \cdot B$$

$$n_+ = n_- ; \theta_c \sim 3^\circ$$



path & spin

$$|0\rangle = k_x > 0$$

$$|1\rangle = k_x < 0$$

$$|\uparrow\rangle = J_z = +\frac{1}{2}$$

$$|\downarrow\rangle = J_z = -\frac{1}{2}$$

$SU(2) \cong SU(2) - SU(2)$
path & spin

$$|0\rangle = k_x > 0$$

$$|1\rangle = 0 > k_x$$

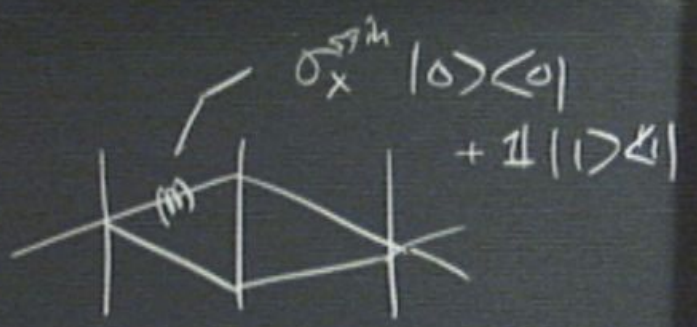
$$|\uparrow\rangle = J_z = +\frac{1}{2}$$

$$|\downarrow\rangle = J_z = -\frac{1}{2}$$

$SU(2) \cong \mathbb{S}^3$
path α

$|0\rangle = k_x > 0$
 $|1\rangle = 0 > 1$

$u(t)$
 $J_z = +1/2$
 $J_z = -1/2$



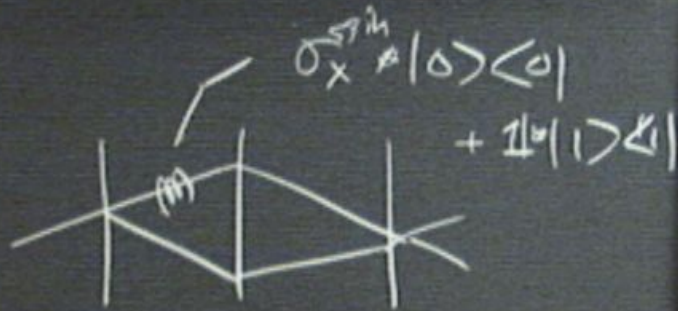
$SU(2) \cong SU(2) - SU(2)$
 path & spin

$|0\rangle = k_x > 0$

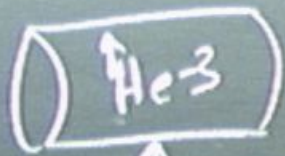
$|1\rangle = 0 > k_x$

$|\uparrow\rangle = I_z = +1/2$

$|\downarrow\rangle = I_z = -1/2$



\uparrow He-3



spin filter

Rotate Spm.

Adiabatic Rotator

$B(\frac{1}{2})$

Rotate Spm.

Adiabatic Rotator

$B(\theta)$



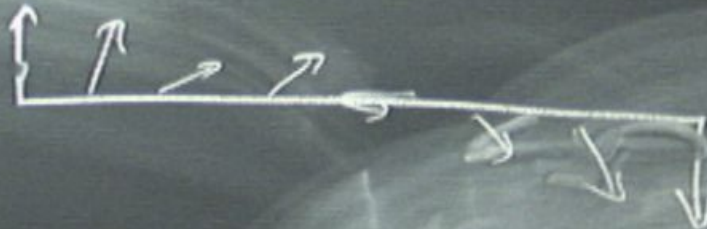
Rotate Spin.

Adiabatic Rotator

input state is
an eigenstate of
initial Field

Output state is
an eigenstate of
Final Field.

$B(\theta)$



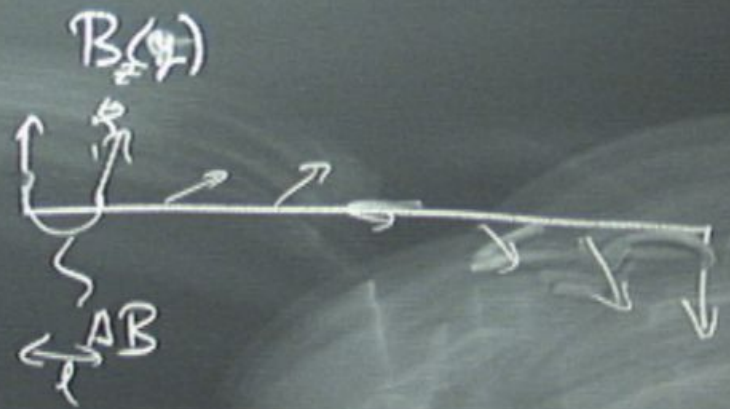
Rotate Spn.

Adiabatic Rotator

input state is
an eigenstate of
initial Field

Output state is
an eigenstate of
Final Field.

state follows
the field



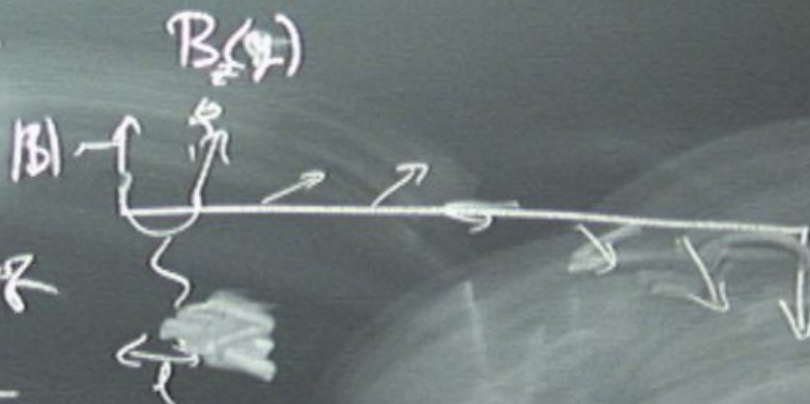
Rotate SpM.

Adiabatic Rotator

input state is
an eigenstate of
initial Field

Output state is
an eigenstate of
Final Field.

state follows
the field



$$\frac{\gamma |B| \sin \theta}{\frac{\hbar}{n}} < \gamma |B|$$

|B| big enough to
dominate stray fields.

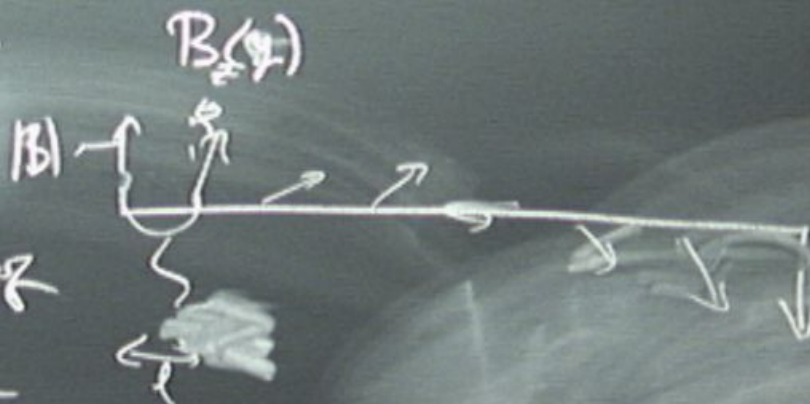
Rotate SpM.

Adiabatic Rotator

input state is
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state follows
the field



$$\frac{\gamma |B| \sin \theta}{\frac{\hbar}{\hbar}} < \gamma |B|$$

|B| big enough to
dominate stray fields.

$S(z) \approx \dots$
path α

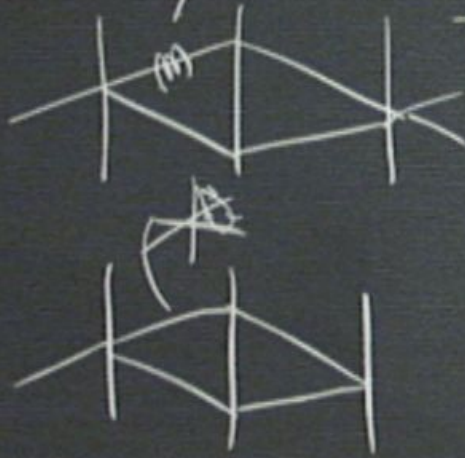
$$|0\rangle = \dots k_+ > 0$$

$$|1\rangle = \dots$$

$$I_2 = +1/2$$

$$I_2 = -1/2$$

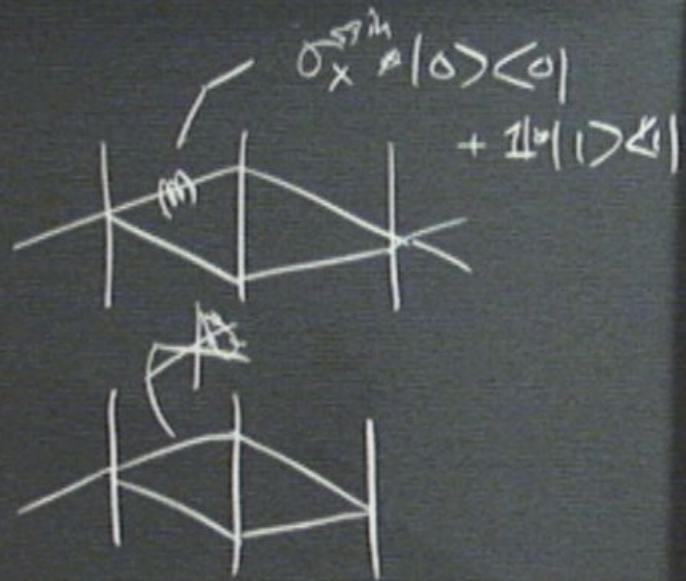
$$\sigma_x^m |0\rangle \langle 0| + 2|1\rangle \langle 1|$$



$SU(2) \cong SU(2) - SU(2)$
 path & spin

$|0\rangle = k_x > 0$
 $|1\rangle = 0 > k_x$

$|\uparrow\rangle = J_z = +1/2$
 $|\downarrow\rangle = J_z = -1/2$



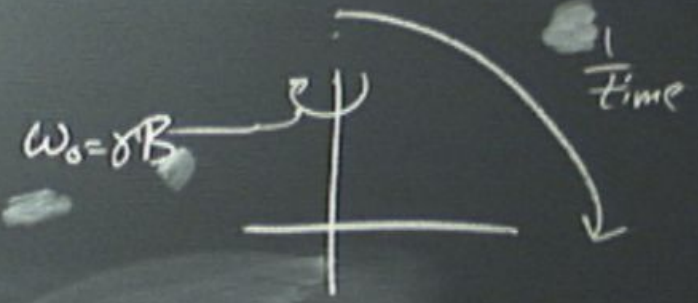
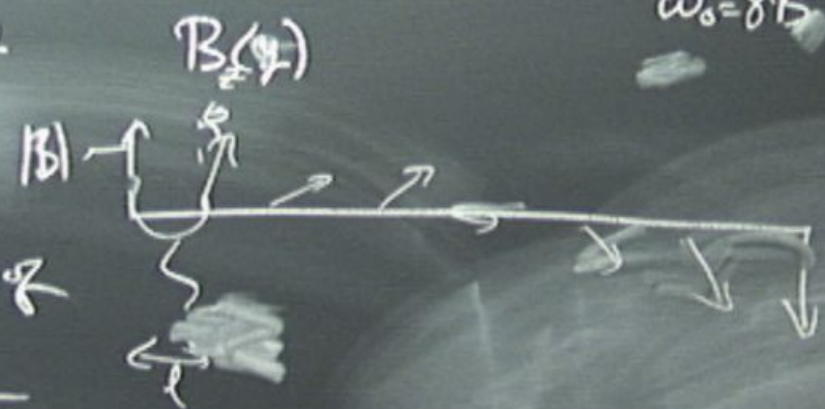
Rotate Spin

Adiabatic Rotator

input state is
an eigenstate of
initial field

Output state is
an eigenstate of
final field

state follows
the field



$$\frac{\gamma |B| \sin \theta}{\hbar} < \gamma |B|$$

|B| big enough to
dominate stray fields.

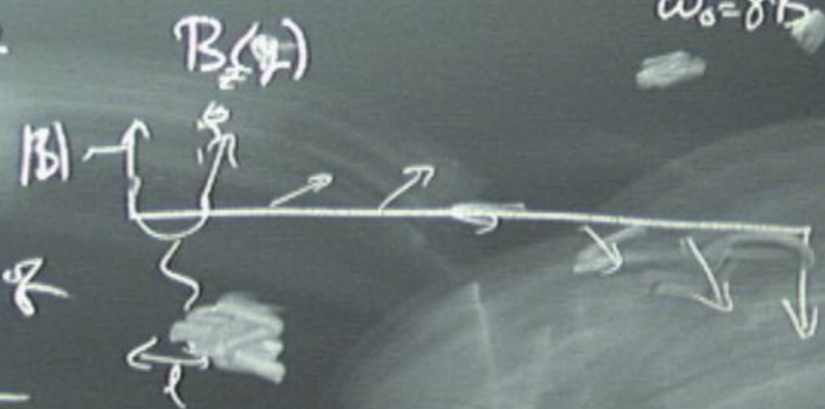
Rotate Spin

Adiabatic Rotator

input state is
an eigenstate of
initial field

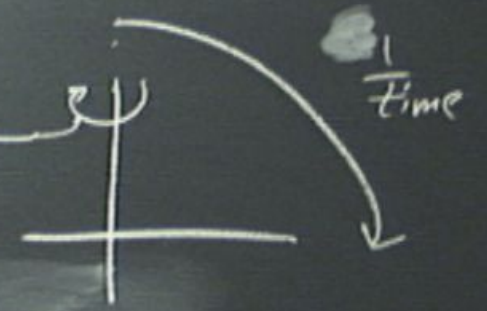
Output state is
an eigenstate of
final field

state follows
the field



$$\frac{\sin \theta}{\hbar} < \gamma |B|$$

$|B|$ big enough to
dominate stray fields



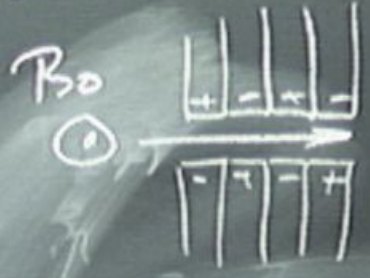
Rotate Spm.

Algebraic
Resonant
N-m-a



Rotate Spin

Applied in
Resonant
N.m. - commuting



Rotate Spin

Algebraic
Resonant
N m-commuting



Rotate Spin

Algebraic
Remnant
N m-commuting



$B_0 \Rightarrow B_1$

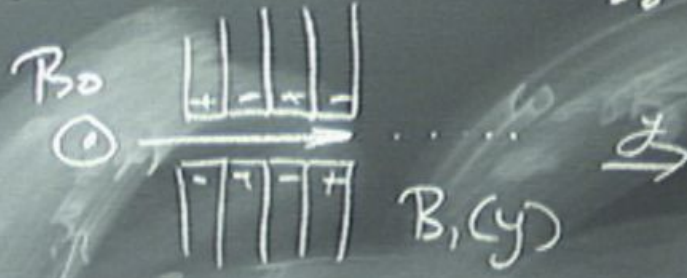
$B_0 \sigma_z$

is enough to
determine the

Rotate Spin

Algebraic
Resonant
Non-commuting

$$B_0 \Rightarrow B_1$$



$$H = \gamma B_0 \sigma_z ; \omega_0 = \gamma B_0$$

$$H_{trans} = \frac{\omega_1}{2} \begin{pmatrix} e^{i\omega_0 t \sigma_z} & \\ & e^{-i\omega_0 t \sigma_z} \end{pmatrix} \sigma_x$$

Rotate Spin

Algebraic
Resonant
Non-commuting

$$B_0 \Rightarrow B_1$$

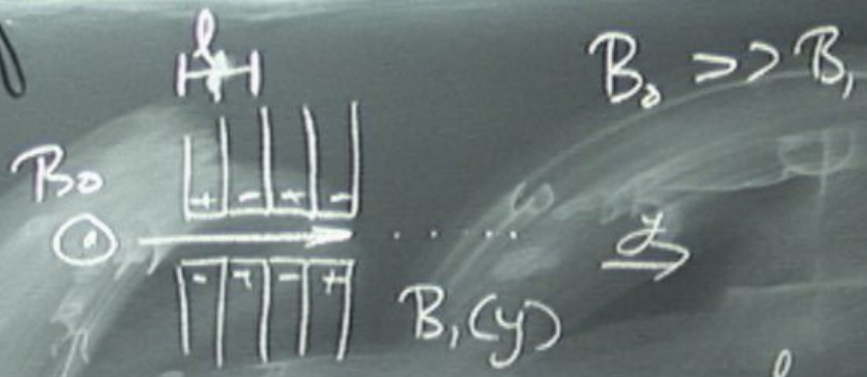


$$\mathcal{H} = \gamma B_0 \sigma_z ; \omega_0 = \gamma B_0$$

$$\mathcal{H}_{\text{trans}} = \frac{\omega_1}{2} \left[e^{i\omega_0 t \sigma_z} \sigma_x e^{-i\omega_0 t \sigma_z} + e^{-i\omega_0 t \sigma_z} \sigma_x e^{i\omega_0 t \sigma_z} \right]$$

Rotate Spin

Applied in
Resonant
NMR - commutation



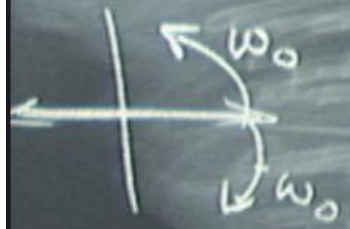
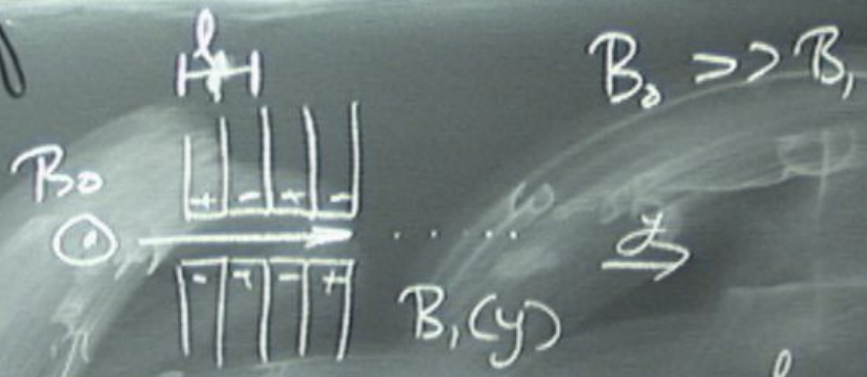
$$\mathcal{H} = \gamma B_0 \sigma_z \quad ; \quad \omega_0 = \gamma B_0$$

$$\frac{l}{v_n} = \frac{2\pi}{\omega_0}$$

$$\mathcal{H}_{trans} = \frac{\omega_1}{2} \left[e^{i\omega_0 t \sigma_z} \sigma_x e^{-i\omega_0 t \sigma_z} + e^{-i\omega_0 t \sigma_z} \sigma_x e^{i\omega_0 t \sigma_z} \right]$$

Rotate Spin

Applied in
Resonant
NMR-cooking



$$\mathcal{H} = \gamma B_0 \sigma_z \quad ; \quad \omega_0 = \gamma B_0$$

$$\frac{\hbar}{\gamma} = \frac{2\pi}{\omega_0}$$

$$\mathcal{H}_{\text{trans}} = \frac{\omega_1}{2} \left[e^{i\omega_0 t \sigma_z} \sigma_x e^{-i\omega_0 t \sigma_z} + e^{-i\omega_0 t \sigma_z} \sigma_x e^{i\omega_0 t \sigma_z} \right]$$

\vec{p} spin $1/2$,

control spin,

super-minor

$$\Pi = \int \dots \sim 1-10$$

$$V_{\text{max}} = \mu \cdot B$$

$$\uparrow \downarrow \dots \theta \leftarrow \rightarrow$$

not unitary, ...

Re



$$\vec{p} = \vec{u} \rho u' ; \quad u = e^{-i\hbar k x}$$

Conduct spin,

Super-minor

$$\Pi = \int \dots \sim 1-10$$

$$\nabla_{\text{avg}} = \langle \nabla \rangle$$

$$\Pi \sim \dots ; \theta \sim \dots$$

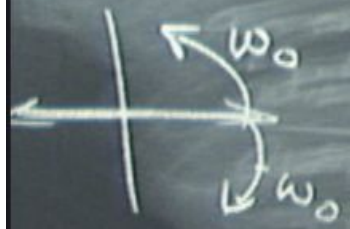
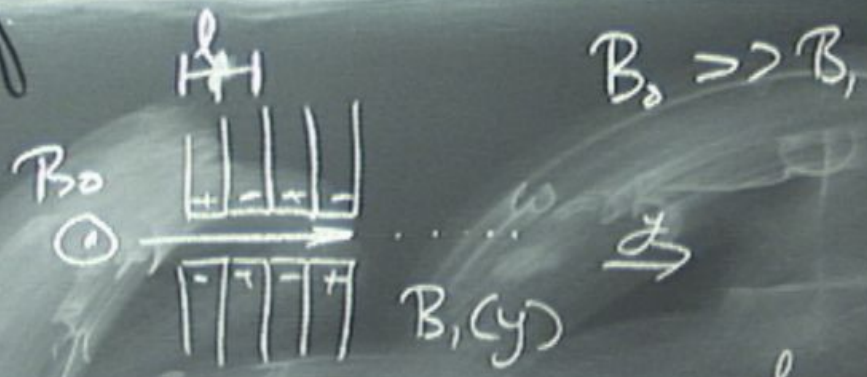
not analog, ...

Ro



Rotate Spin

Applied in
Resonant
NMR-cooking



$$H = \gamma B_0 \sigma_z ; \omega_0 = \gamma B_0$$

$$\frac{l}{\hbar} = \frac{\gamma}{\omega_0}$$

$$H_{trans} = \frac{\omega_1}{2} \left[e^{i\omega_0 t \sigma_z} \sigma_x e^{-i\omega_0 t \sigma_z} + e^{-i\omega_0 t \sigma_z} \sigma_x e^{i\omega_0 t \sigma_z} \right]$$

$$H_{total} = H + H_{trans}$$

$$\vec{P} = U S U^\dagger ; U = e^{-i H_0 t}$$

$$\frac{dP}{dt} = -i [H, P]$$

$$\frac{d\vec{P}}{dt} = -i [H, \vec{P}]$$

$$\vec{\rho} = U \rho U^{-1} ; U = e^{-i H_0 t}$$

$$\frac{d\rho}{dt} = -i [H, \rho] \quad \frac{dU}{dt} = -i H_0 U$$

$$\frac{d\vec{\rho}}{dt} = -i [H_{eff}, \vec{\rho}] = \frac{dU}{dt} U^{-1} \rho U^{-1} + U \frac{d\rho}{dt} U^{-1} + U \rho \frac{dU^{-1}}{dt}$$

$$\tilde{\rho} = U \rho U^{-1} ; U = e^{-i H_e t}$$

$$\frac{d\rho}{dt} = -i [H, \rho] \quad \frac{dU}{dt} = -i H_e U$$

$$\frac{d\tilde{\rho}}{dt} = -i [H_{\text{eff}}, \tilde{\rho}] = \frac{dU}{dt} U^{-1} \rho U^{-1} + U \frac{d\rho}{dt} U^{-1} + U \rho \frac{dU^{-1}}{dt}$$

$$= -i [\tilde{H}, \tilde{\rho}] - i [H_e, \rho]$$

$$\tilde{H} = U H_{\text{total}} U^{-1}$$

$$\tilde{\rho} = U \rho U^{-1} ; U = e^{-i H_R t}$$

$$\frac{d\rho}{dt} = -i [H, \rho] \quad \frac{dU}{dt} = -i H_R U$$

$$\frac{d\tilde{\rho}}{dt} = -i [H_{eff}, \tilde{\rho}] = \frac{dU}{dt} U^{-1} \rho U^{-1} + U \frac{d\rho}{dt} U^{-1} + U \rho \frac{dU^{-1}}{dt}$$

$$= -i [\tilde{H}, \tilde{\rho}] - i [H_R, \tilde{\rho}]$$

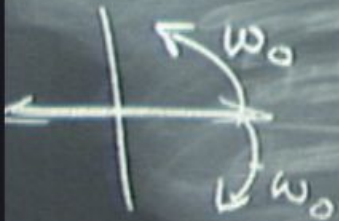
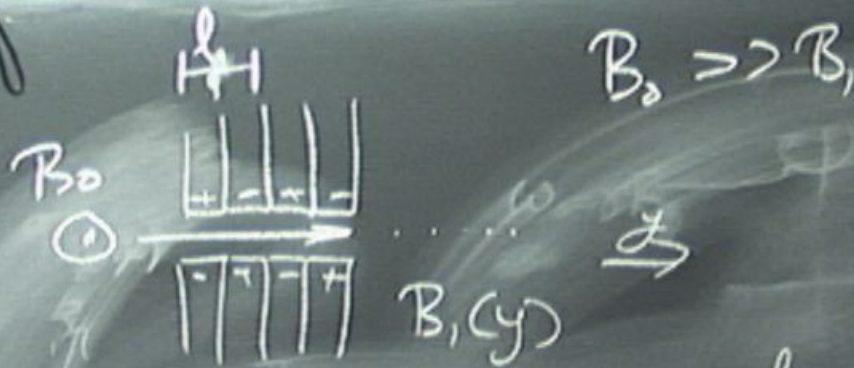
$$\tilde{H} = U H U^{-1}$$

$$H_{eff} = \tilde{H} - H_R$$

Ro

Rotate Spin

Applied in Resonant NMR-cooking



$$\mathcal{H} = \gamma B_0 \sigma_z \quad ; \quad \omega_0 = \gamma B_0$$

$$\frac{1}{\nu_n} = \frac{2\pi}{\omega_0}$$

$$\mathcal{H}_{trans} = \frac{\omega_1}{2} \left[e^{i\omega_0 t \sigma_z} \sigma_x e^{-i\omega_0 t \sigma_z} + e^{-i\omega_0 t \sigma_z} \sigma_x e^{i\omega_0 t \sigma_z} \right]$$

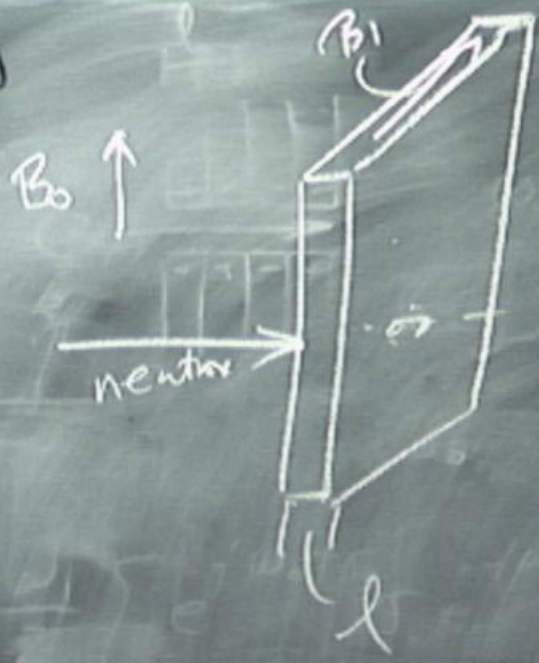
$$\mathcal{H}_{total} = \mathcal{H} + \mathcal{H}_{trans}$$

$$U = e^{-i\omega_0 t \sigma_z} \quad \mathcal{H} = \omega_0 \sigma_z$$

$$\tilde{\mathcal{H}} = \gamma B_0 \sigma_z + \frac{\omega_1}{2} \sigma_x + e^{-i\omega_0 t \sigma_z} \left[\frac{\omega_1}{2} \sigma_x \right] e^{i\omega_0 t \sigma_z}$$

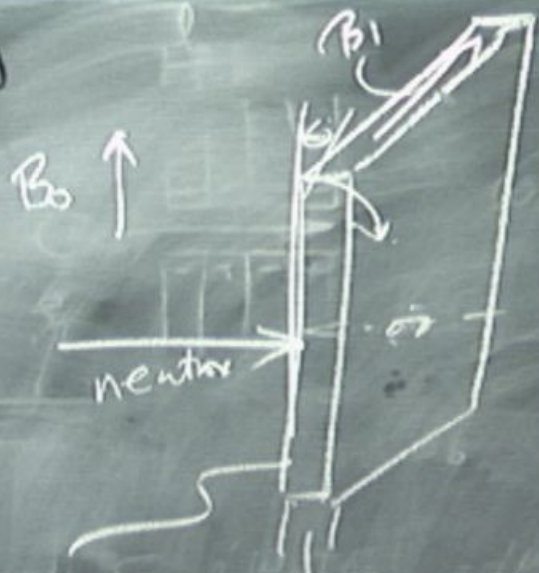
Rotated Spin

Adiabatic
Resonant
Non-commuting



Rotate Spin

Adiabatic
Resonant
Non-commuting



$$\chi = e \frac{\gamma B_1}{2} \sigma_x \frac{l}{v_n}$$

1
case

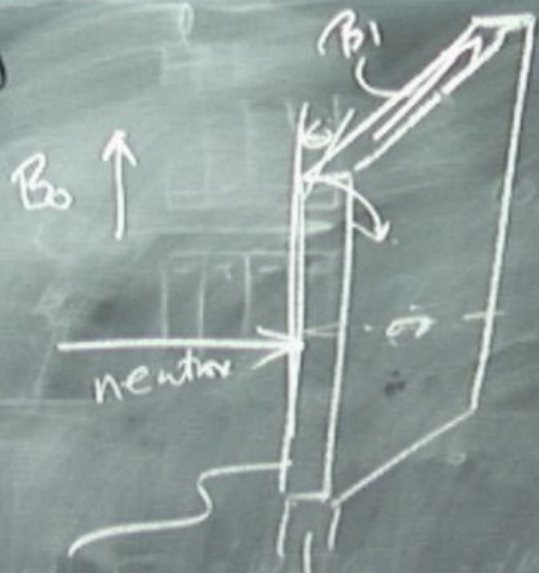
$$B_1 \sim 2T$$

$$\gamma = 1.76 \text{ H}^{-1}/g$$

$$v = 2,000 \text{ m/s}$$

Rotated Spinn.

Alindetie
Resmant
Nm-commutia



$$\alpha = e \frac{\gamma B_1}{2} \sigma_x \frac{l}{v_n} l$$

1
case

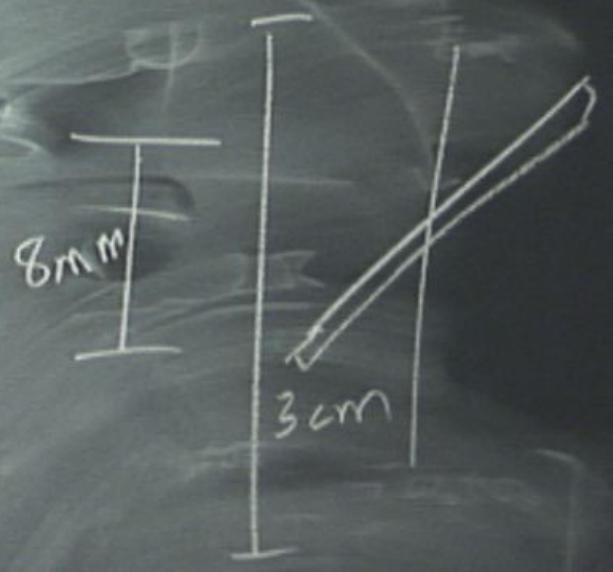
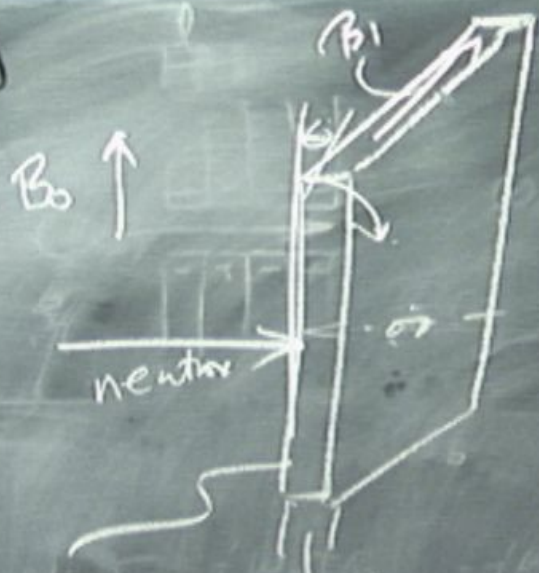
$$B_1 = 2T$$

$$\gamma = 1.76 \text{ H} \cdot 10^6$$

$$v = 2,000 \text{ m/s}$$

Rotated Spinn.

Alindetik
Resmanit
Nm-koordinat



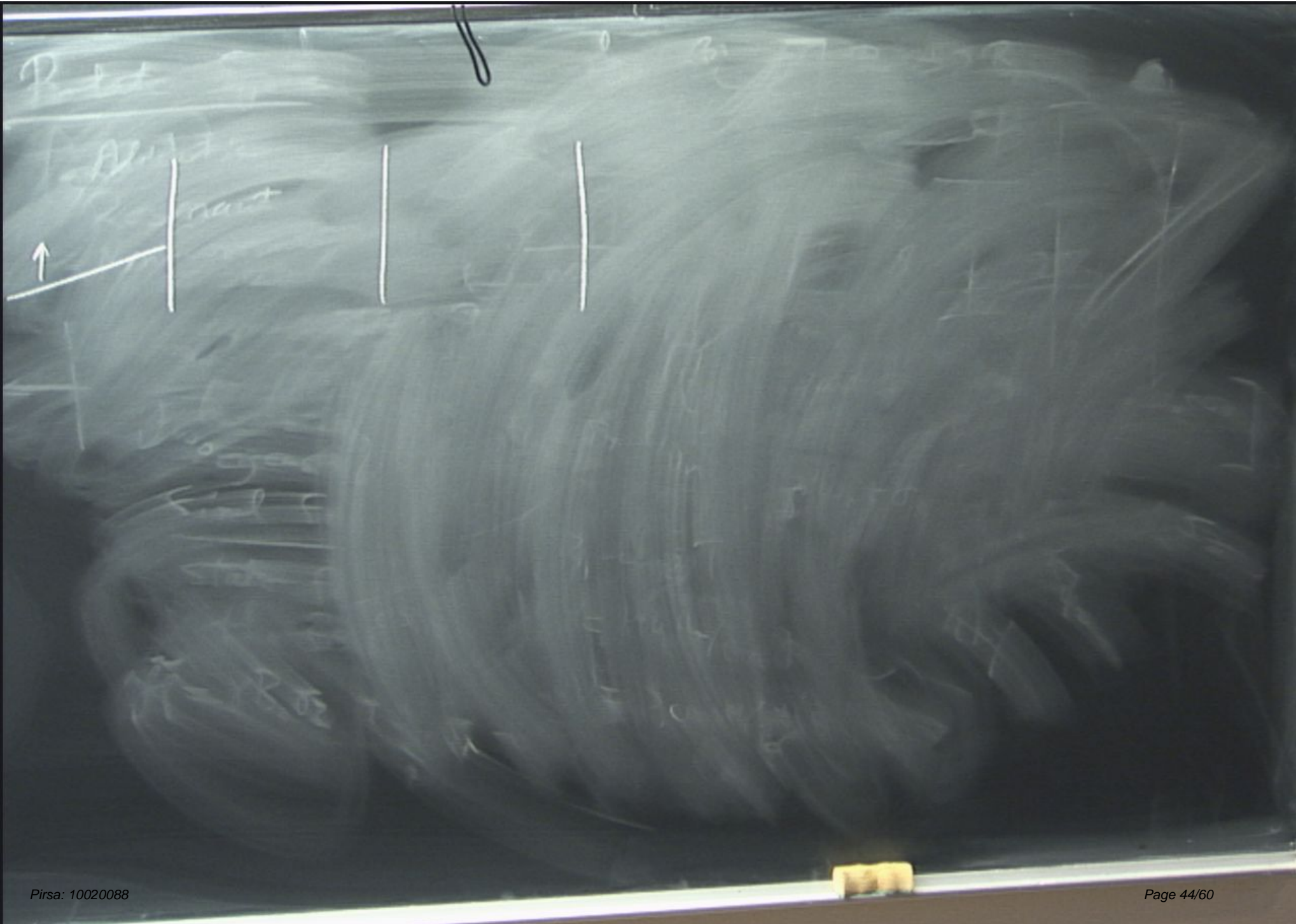
$$\chi = e^{-\frac{\gamma B_1}{2}} \sigma_x \frac{l}{l_0} l$$

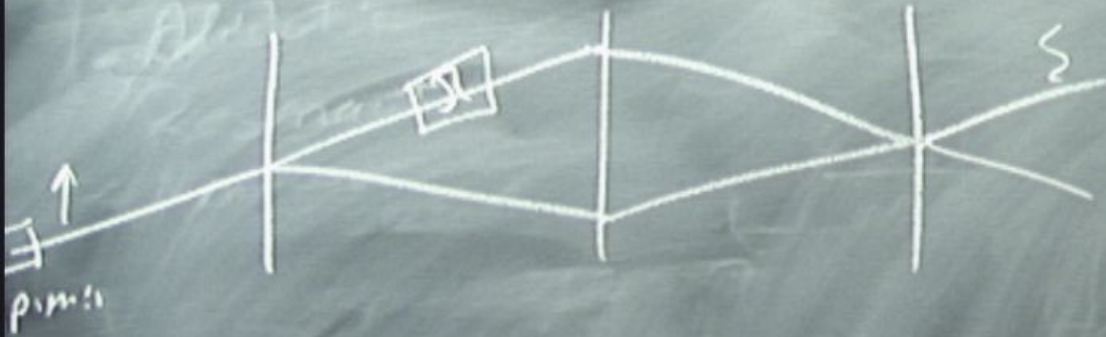
1
case

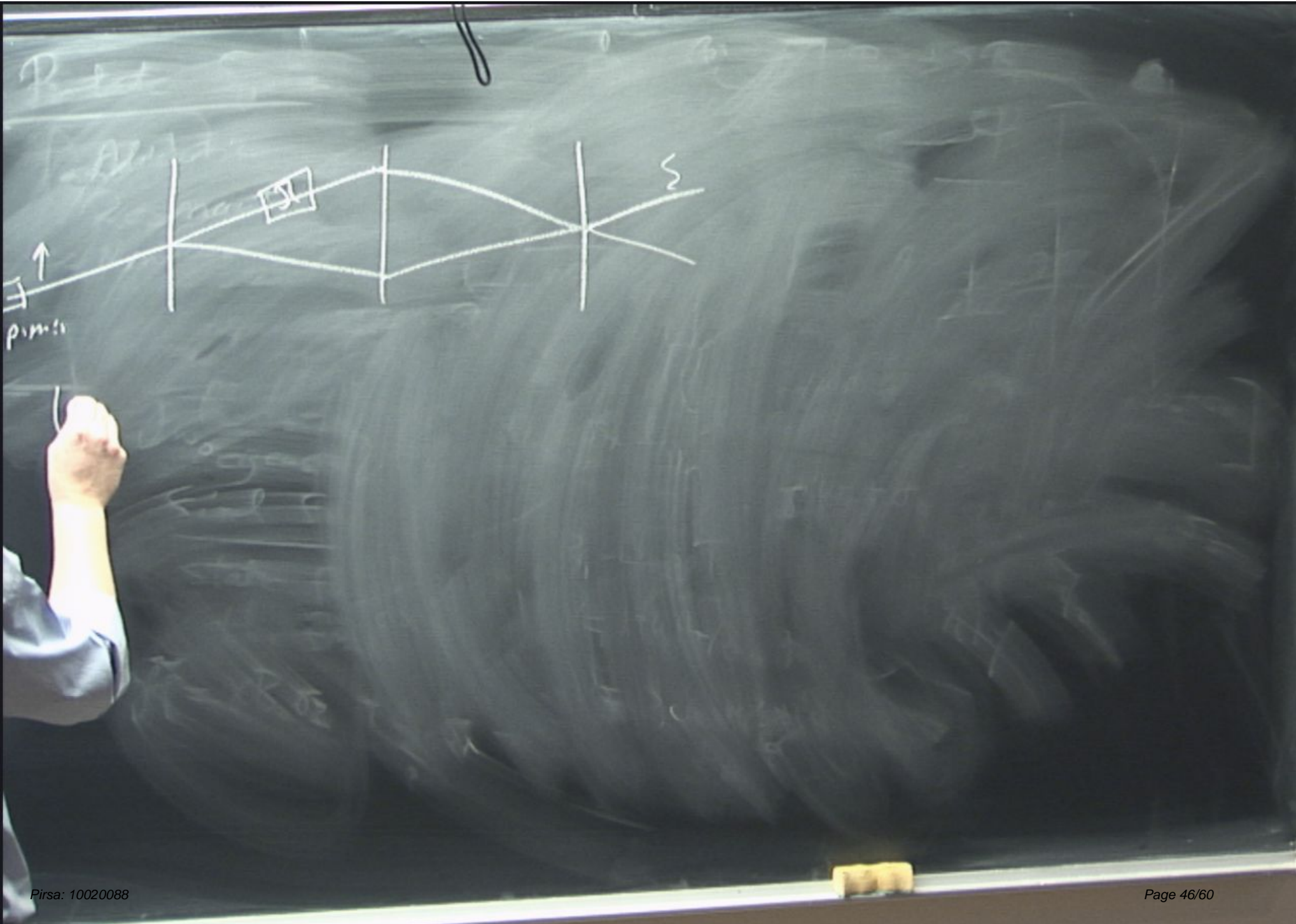
$$B_1 = 2T$$

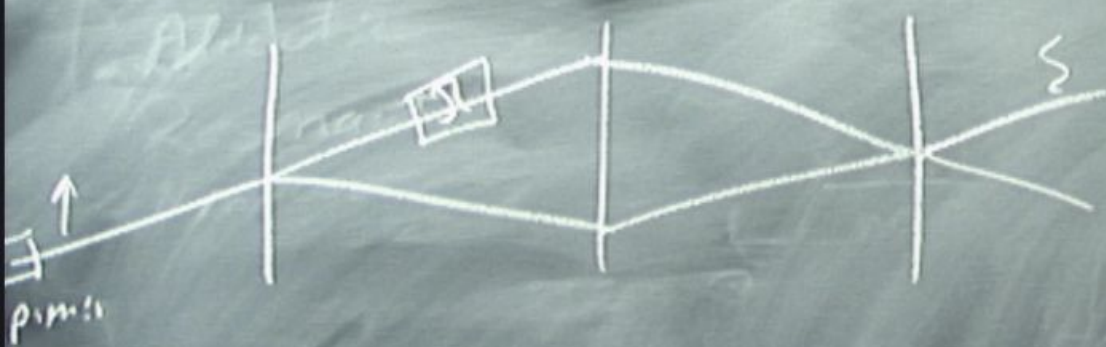
$$\gamma = 196 \text{ H} \cdot \text{m} / \text{A}$$

$$V = 2,000 \text{ m/s}$$









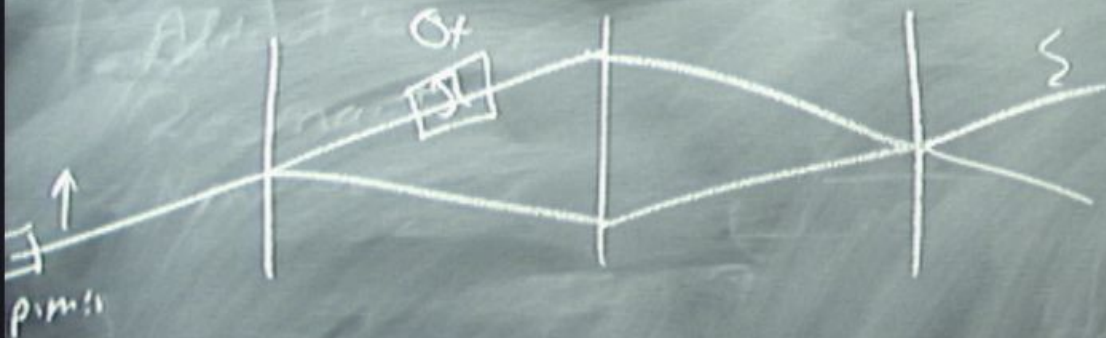
$$|0\rangle$$

$$|0\rangle$$

$$|1\rangle$$

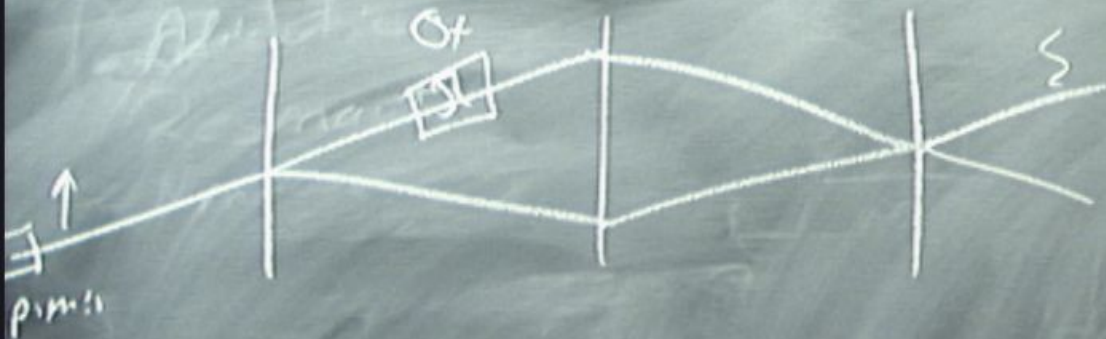
$$|0\rangle$$

$$|1\rangle$$



$$|0\rangle \xrightarrow[\text{H} \otimes \text{I}]{\text{Controlled-Not}}$$

- $|0\rangle$
- $|1\rangle$
- $|0\rangle$
- $|1\rangle$

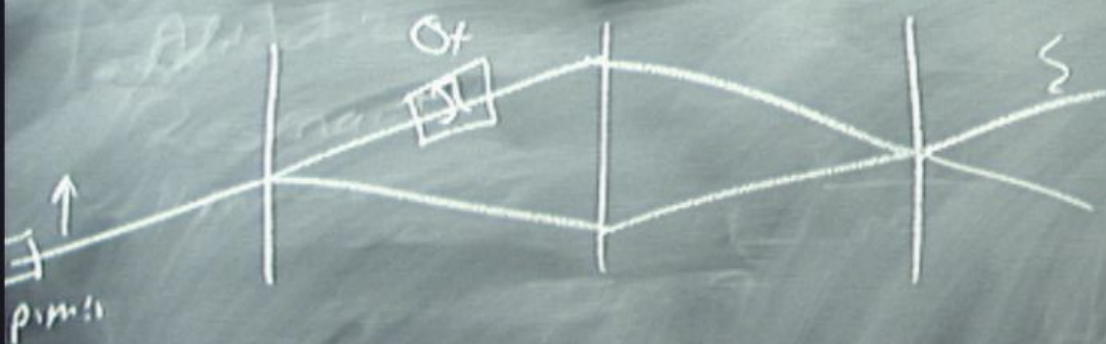


$$|0\uparrow\rangle \xrightarrow[\text{H} \otimes \text{II}]{\text{Sivertblu}} \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle)$$

$$H|\uparrow\rangle\langle\uparrow| + H|\downarrow\rangle\langle\downarrow|$$

- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & -1 & 1 \end{pmatrix}$$

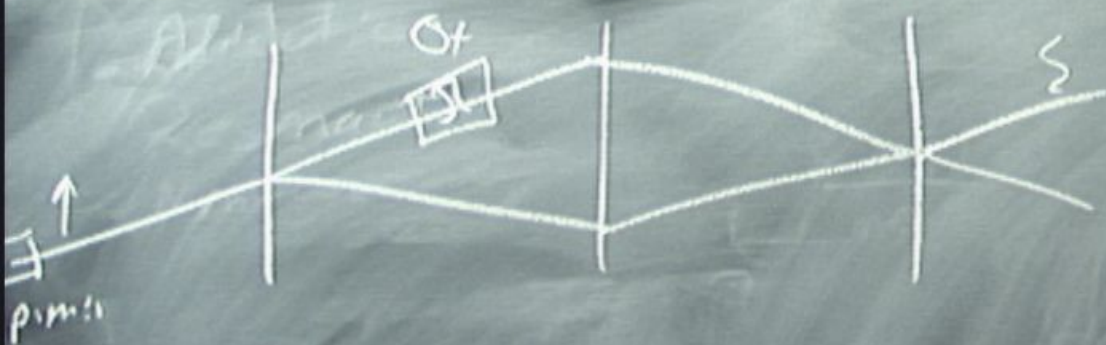


$$|0\uparrow\rangle \xrightarrow[\text{H} \otimes \mathbb{I}]{\text{Sivert bla}} \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |\uparrow\rangle$$

$$|1\uparrow\rangle \langle 1| + |1\downarrow\rangle \langle 1\downarrow| \quad \text{spin flip}$$

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & -1 & 1 \end{pmatrix}$$

- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$



$$|0\uparrow\rangle \xrightarrow[\text{H} \otimes \mathbb{I}]{\text{Sivertblin}} \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |\uparrow\rangle$$

$$|1\uparrow\rangle\langle 1| + |1\downarrow\rangle\langle 1| \quad \text{Spin flip}$$

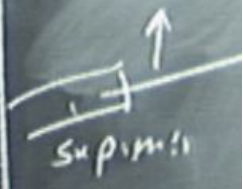
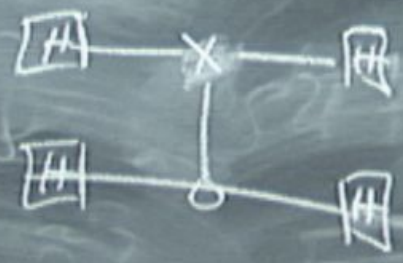
$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & -1 & 1 \end{pmatrix}$$

$$|0\rangle\langle 0| \otimes \sigma_x + |1\rangle\langle 1| \otimes \mathbb{I}$$

$$U = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$

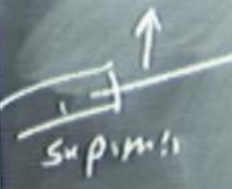
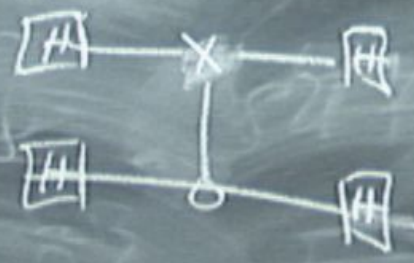
$CNOT \equiv \sigma_x |0\rangle\langle 0| + |1\rangle\langle 1|$



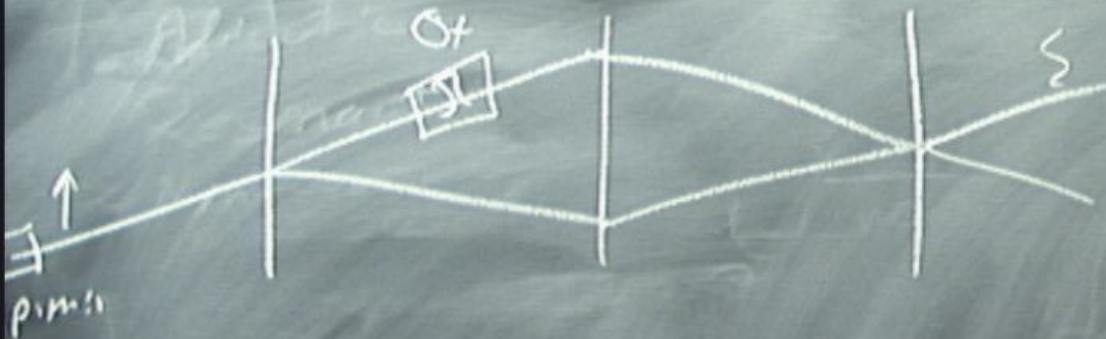
$|0\rangle$
 $|0\rangle$
 $|1\rangle$
 $|0\rangle$
 $|1\rangle$

$CNOT \equiv$

$\sigma_x \frac{|00\rangle + |11\rangle}{\sqrt{2}} + \frac{|01\rangle + |10\rangle}{\sqrt{2}}$



$|0\rangle$
 $|1\rangle$
 $|0\rangle$
 $|1\rangle$



$$|0\uparrow\rangle \xrightarrow[\text{H} \otimes \mathbb{I}]{\text{Spin flip}} \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |\uparrow\rangle$$

$$|1\uparrow\rangle\langle 1| + |1\downarrow\rangle\langle 1|$$

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & -1 & -1 \end{pmatrix}$$

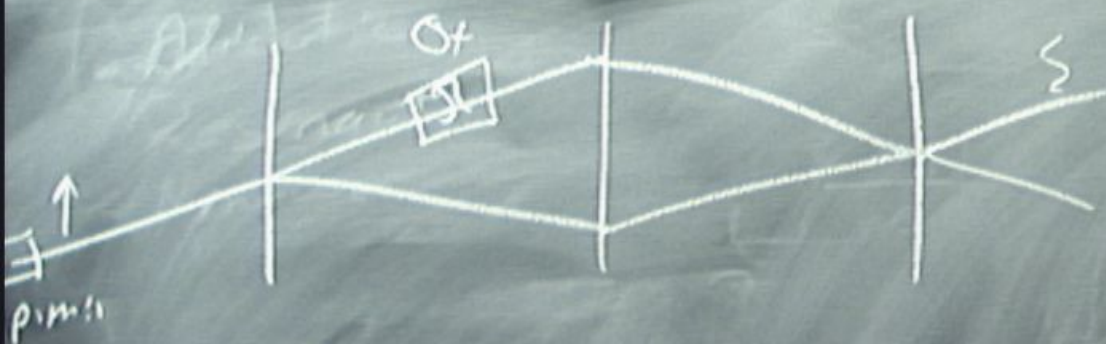
Spin flip

$$\frac{1}{\sqrt{2}} (|0\downarrow\rangle + |1\uparrow\rangle)$$

$$|0\rangle\langle 0| \otimes \sigma_x + |1\rangle\langle 1| \otimes \mathbb{I}$$

$$U = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$



$$|0\uparrow\rangle \xrightarrow[H \otimes II]{\text{Sierblu}} \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |\uparrow\rangle$$

$$|1\uparrow\rangle\langle 1| + |1\downarrow\rangle\langle 1|$$

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & -1 & -1 \end{pmatrix}$$

spin flip

$$\frac{1}{\sqrt{2}} (|0\downarrow\rangle + |1\uparrow\rangle)$$

$$|0\rangle\langle 0| \sigma_x + |1\rangle\langle 1| I$$

$$U = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

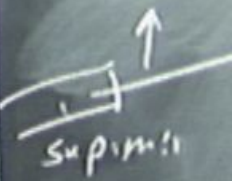
entangled state

- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$

$$\rightarrow \frac{1}{\sqrt{2}} (|1\downarrow\rangle + |0\uparrow\rangle)$$

$$\xrightarrow{\text{3rd}} \frac{1}{\sqrt{2}} \left[\frac{1}{\sqrt{2}} (|1\downarrow\rangle - |0\downarrow\rangle) + \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) \right]$$

D-beam



$|0\uparrow\rangle$

$|0\uparrow\rangle$

$|1\uparrow\rangle$

$|0\downarrow\rangle$

$|1\downarrow\rangle$

$$\longrightarrow \frac{1}{\sqrt{2}} (|1\downarrow\rangle + |0\uparrow\rangle)$$

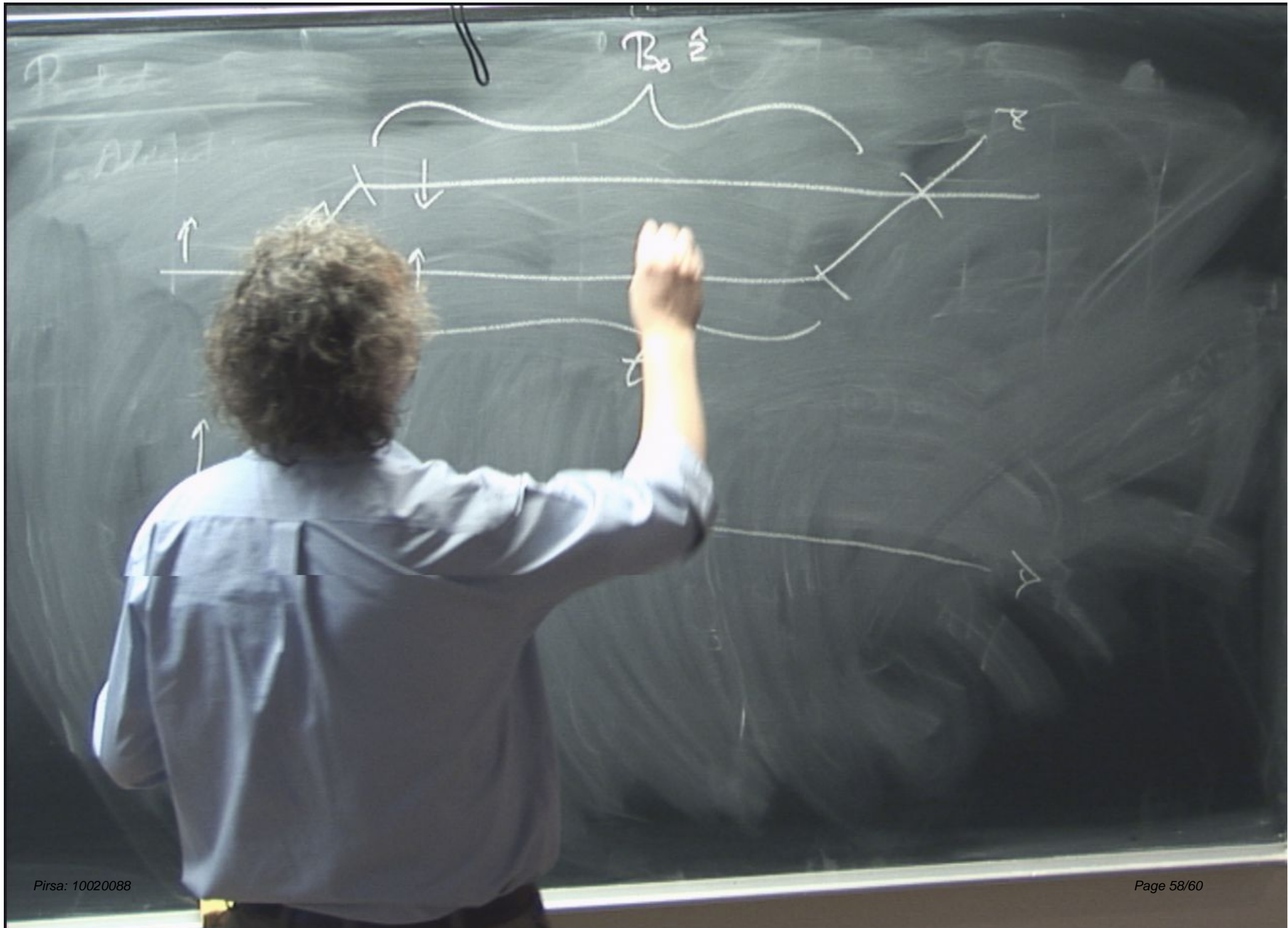
$$\begin{matrix} 3^{rd} \\ \text{blade} \end{matrix} \longrightarrow \frac{1}{\sqrt{2}} \left[\frac{1}{\sqrt{2}} (|1\downarrow\rangle - |0\downarrow\rangle) + \frac{1}{\sqrt{2}} (|0\uparrow\rangle + |1\uparrow\rangle) \right]$$

$$\text{O-beam} = \frac{1}{2} \left[|0\rangle (|1\uparrow\rangle - |1\downarrow\rangle) \right]$$

coherence



- $|0\uparrow\rangle$
- $|1\uparrow\rangle$
- $|0\downarrow\rangle$
- $|1\downarrow\rangle$



Product

$B_0 \hat{=} \text{Random}$



Product

$B_0 \approx$
Random

