

Title: Quantum Gravity - Review (PHYS 638) - Lecture 12

Date: Feb 09, 2010 10:00 AM

URL: <http://pirsa.org/10020066>

Abstract:

interpretation of WdW equation $\hat{H}_L \psi = 0$

+ interpretation of WdW equation $\hat{H}_L \psi = 0$

where is "time"?

+ interpretation of WdW equation $\hat{H}_1 \psi = 0$

where is "time"?

it

+ interpretation of WdW equation $\hat{H}_L \psi = 0$

where is "time"?

$$i\hbar \frac{d}{dt} \Psi(t) = \hat{H}[N, \bar{N}](t) \Psi(t)$$

+ interpretation of WdW equation $\hat{H}_\perp \psi = 0$
where is "time"?

$$i\hbar \frac{d}{dt} \psi(t) = \hat{H}[N, \vec{N}](t) \psi(t)$$

$$\text{if } \psi = \psi^{\text{phys}} \Rightarrow = 0$$

Interpretation of WdW equation $\hat{H}_\perp \psi = 0$

where is "time"?

$$i\hbar \frac{d}{dt} \psi(t) = \hat{H}[N, \vec{N}](t) \psi(t)$$

if $\psi = \psi^{\text{phys}} \Rightarrow = 0 \Rightarrow$ "frozen time"

interpretation of WdW equation $\hat{H}_\perp \psi = 0$

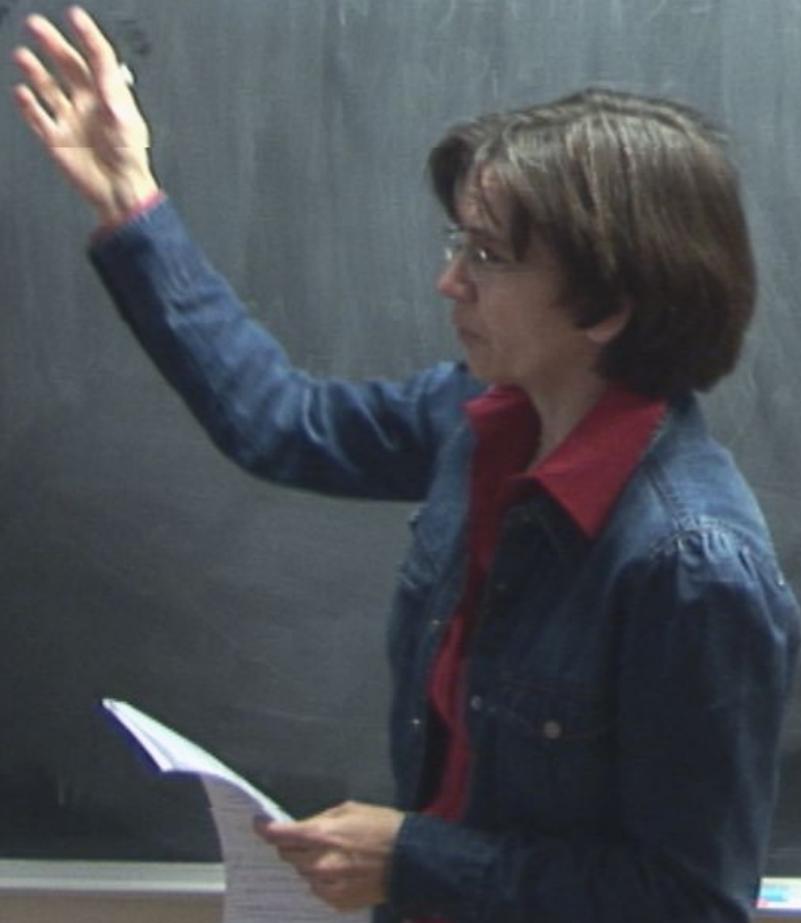
where is "time"?

$$i\hbar \frac{d}{dt} \psi(t) = \hat{H}[N, \dot{N}](t) \psi(t)$$

$\psi = \psi^{\text{phys}} \Rightarrow = 0 \Rightarrow$ "frozen time"
 \hookrightarrow "time" must be recovered from among the dynamical variables!

De Witt metric g_{ij} has indefinite signature

$(-++++)$



The Witt metric g^{ijkl} has indefinite signature

$$(-++++), \quad (\square + m^2)\phi = 0 \quad (\text{KG equation})$$

The Witt metric g^{ijkl} has indefinite signature

$$(-++++), \quad (\square + m^2)\phi = 0 \quad (\text{KG equation})$$

Isham, $gr/qc - 9210011$

The Witt metric g^{ijkl} has indefinite signature

$$(-++++), \quad (\square + m^2)\phi = 0 \quad (\text{KG equation})$$

Isham, $gr/qc - 9210011$

\Rightarrow stuck?

The Witt metric g^{ijkl} has indefinite signature

$$(-++++), \quad (\square + m^2)\phi = 0 \quad (\text{KG equation})$$

Isham, $gr/qc - 9210011$

\Rightarrow stuck?

"new connection variables" (Ashtekar 1986)

Loop quantum gravity

- take classically equivalent, first-order formulation

Loop quantum gravity

- take classically equivalent, first-order formulation of GR in terms of orthonormal frame (tetrad) fields $e_{\mu}^A(x)$, $\mu=0, \dots, 3$, $A=0, \dots, 3$ internal $so(3,1)$ -index; and a spin connection $\omega_{\mu}^A{}_B(x)$, considered as independent variables.

Loop quantum gravity

- take classically equivalent, first-order formulation of GR in terms of orthonormal frame (tetrad) fields

$e_{\mu}^A(x)$, $\mu = 0, \dots, 3$, $A = 0, \dots, 3$ internal $so(3,1)$ -index;

and a spin connection $\omega_{\mu}^A{}_B(x)$, considered as independent variables.

$$g_{\mu\nu} = \eta_{AB} e_{\mu}^A e_{\nu}^B$$

internal Minkowski metric

Loop quantum gravity

- take classically equivalent, first-order formulation of GR in terms of orthonormal frame (tetrad) fields

$e_{\mu}^A(x)$, $\mu = 0, \dots, 3$, $A = 0, \dots, 3$ internal $so(3,1)$ -index

and a spin connection $\omega_{\mu}^A{}_B(x)$, considered as independent variables.

$$g_{\mu\nu} = \eta_{AB} e_{\mu}^A e_{\nu}^B$$

internal Minkowski metric

Loop quantum gravity

- take classically equivalent, first-order formulation of GR in terms of orthonormal frame (tetrad) fields $e_{\mu}^A(x)$, $\mu = 0, \dots, 3$, $A = 0, \dots, 3$ internal $so(3,1)$ -index; and a spin connection $\omega_{\mu}^A{}_B(x)$, considered as independent variables.

$$g_{\mu\nu} = \eta_{AB} e_{\mu}^A e_{\nu}^B$$

internal Minkowski metric

perform a 3+1 split and partially fix the gauge (s.t. $so(3,1)$)

\Rightarrow

$\rightarrow so(3)$

perform a 3+1 split and partially fix the gauge (s.t. $so(3,1)$)

\Rightarrow Yang-Mills like, conjugate variables

$$\{A_i^a(\vec{x}), E_b^j(\vec{y})\} = 2\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

$\rightarrow so(3)$

perform a 3+1 split and partially fix the gauge (s.t. $so(3,1)$)

\Rightarrow Yang-Mills like, conjugate variables

$$\{A_i^a(\vec{x}), E_b^j(\vec{y})\} = 2\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

i, j - spatial

perform a 3+1 split and partially fix the gauge (s.t. $so(3,1)$)

\Rightarrow Yang-Mills like, conjugate variables

$$\{A_i^a(\vec{x}), E_b^j(\vec{y})\} = 2\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

i, j - spatial, $a, b = 1, 2, 3$ - $so(3)$ -index

$$A \Rightarrow A^a(x)$$

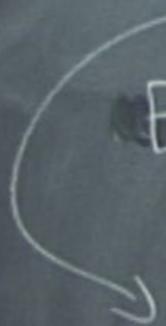
$A \Rightarrow A_i^a(x) \sim \text{so}(3) - \text{connection}$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$$E_b^j = \hbar e_i^a$$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$E_b^j = \sqrt{h} e_b^j$ (densitized triad)



$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$E_b^j = \sqrt{h} e_b^j$ (densitized triad)

$$F_{ij}^a = 2 \left(\partial_i A_j^a - \partial_j A_i^a \right) + G$$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$$E_b^j = \sqrt{h} e_b^j \quad (\text{densitized triad})$$

$$F_{ij}^a = 2G_N \partial_{[i} A_{j]}^a + G_N^2 \epsilon_{abc} A_i^b A_j^c$$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$E_b^j = \sqrt{h} e_b^j$ (densitized triad)

$F_{ij}^a = 2G_{ab} \partial_{[i} A_{j]}^a + G_{ab}^2 \epsilon_{abc} A_i^b A_j^c$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$E_b^j = \sqrt{h} e_b^j$ (densitized triad)

$$F_{ij}^a = 2G_{ab} \partial_{[i} A_{j]}^a + G_{ab}^2 \epsilon_{abc} A_i^b A_j^c$$

$\Rightarrow 7$

$A \ni A_i^a(x) \sim \text{so}(3)\text{-connection}$,

$$E_b^j = \sqrt{h} e_b^j \quad (\text{densitized triad})$$

$$F_{ij}^a = 2G_{ab} \partial_{[i} A_{j]}^a + G_{ab}^2 \epsilon_{abc} A_i^b A_j^c$$

$\Rightarrow 7 \times \infty^3$ first-class constraints

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \psi$$

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$C_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$\mathcal{H}_i = F_{ij}^a E_a^i = 0 \quad \text{momentum constraints}$$

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$\mathcal{H}_i = F_{ij}^a E_a^i = 0 \quad \text{momentum constraints}$$

$$\mathcal{H}_\perp = \frac{1}{2} \frac{\epsilon^{abc} F_{ijc}}{\sqrt{|\det E_a^i|}} E_a^i E_b^j - \frac{\beta^2 + 1}{\beta^2 \sqrt{|\det E_a^i|}} E_{[a}^i E_{b]}^j \cdot (G_N A_i^a - T_i^a) \cdot (G_N A_i^b - T_i^b)$$

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$\mathcal{H}_i = F_{ij}^a E_a^i = 0 \quad \text{momentum constraints}$$

$$\mathcal{H}_\perp = \frac{1}{2} \frac{\epsilon^{abc} F_{ijc}}{\sqrt{|\det E_a^i|}} E_a^i E_b^j - \frac{\beta^2 + 1}{\beta^2 \sqrt{|\det E_a^i|}} E_{[a}^i E_{b]}^j \cdot (G_N A_i^a - \Gamma_i^a) \cdot (G_N A_i^b - \Gamma_i^b)$$

$$\Gamma = \Gamma(E) \sim \mathfrak{so}(3)\text{-connection}$$

perform a 3+1 split and partially fix the gauge (s.t. $so(3)$

\Rightarrow Yang-Mills like, conjugate variables $\rightarrow so(3)$

$$\{A_i^a(\vec{x}), E_b^j(\vec{y})\} = 2\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

i, j - spatial, $a, b = 1, 2, 3$ - $so(3)$ -index

\mathcal{A}

Immirzi-Barbero param

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$\mathcal{H}_i = F_{ij}^a E_a^i = 0 \quad \text{momentum constraints}$$

$$\mathcal{H}_\perp = \frac{1}{2} \frac{\epsilon^{abc} F_{ijc}}{\sqrt{|\det E_a^i|}} E_a^i E_b^j - \frac{\beta^2 + 1}{\beta^2 \sqrt{|\det E_a^i|}} E_{[a}^i E_{b]}^j \cdot (G_N A_i^a - \Gamma_i^a) \cdot (G_N A_i^b - \Gamma_i^b)$$

$\beta = 1/\alpha$

$\Gamma = \Gamma(E) \sim \text{so}(3)$ -connection

$$G_a = \mathcal{D}_i E_a^i + G_N \epsilon_{abc} A_i^b E^{ci} = \mathcal{D}_i E_a^i = 0 \quad \text{"Gauss constraints"}$$

$$\mathcal{H}_i = F_{ij}^a E_a^i = 0 \quad \text{momentum constraints}$$

$$\mathcal{H}_\perp = \frac{1}{2} \frac{\epsilon^{abc} F_{ijc}}{\sqrt{|\det E_a^i|}} E_a^i E_b^j - \frac{\beta^2 + 1}{\beta^2 \sqrt{|\det E_a^i|}} E_{[a}^i E_{b]}^j \cdot (G_N A_i^a - \Gamma_i^a) \cdot (G_N A_i^b - \Gamma_i^b)$$

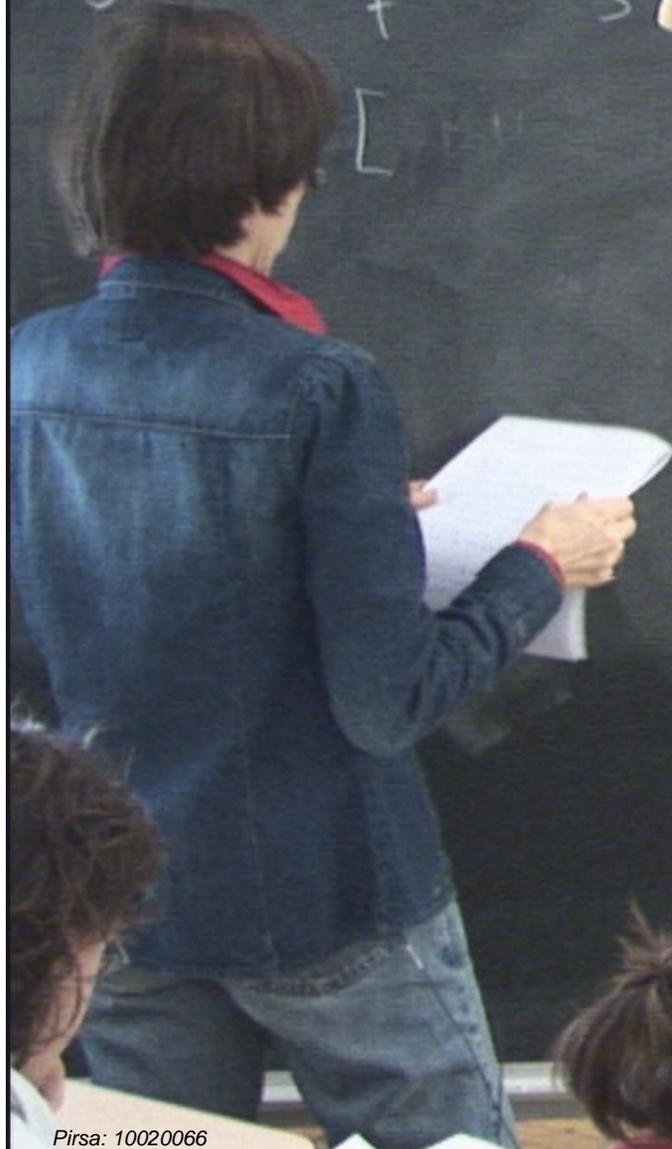
PB algebra closes

$$\beta = 1$$

$$\Gamma = \Gamma(E) \sim \text{so}(3)\text{-connection}$$

Dirac quantization! $\int_{\Sigma} \omega$ for the gauge (s.t. sol)

[



Dirac quantization! "fix the gauge (s.t.)"

$$[\hat{A}_i^a(\vec{x}), \hat{E}_j^b(\vec{y})] = i\hbar \delta\pi\beta \delta_b^a \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

formally $\psi(\vec{x}) \sim \psi[A] = A_i^a(\vec{x}) \psi[A]$

Dirac quantization! "fix the gauge (s.t. ...)

$$[\hat{A}_i^a(\vec{x}), \hat{E}_b^j(\vec{y})] = i\hbar \delta\pi\beta \delta_b^a \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

formally: $\hat{A}_i^a(\vec{x})\psi[A] = A_i^a(\vec{x})\psi[A]$

$$\hat{E}_b^j(\vec{x})\psi[A] = \frac{\hbar}{i} \delta\pi\beta \frac{\delta}{\delta A_i^b(\vec{x})} \psi[A]$$

Dirac quantization! ... for the gauge (s.t.)

$$[\hat{A}_i^a(\vec{x}), \hat{E}_b^j(\vec{y})] = i\hbar \delta\pi\beta \delta_b^a \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

formally: $\hat{A}_i^a(\vec{x})\Psi[A] = A_i^a(\vec{x})\Psi[A]$

$$\hat{E}_b^j(\vec{x})\Psi[A] = \frac{\hbar}{i} \delta\pi\beta \frac{\delta}{\delta A_i^b(\vec{x})} \Psi[A]$$

$$\psi[A] \in \mathcal{H}^{anx}$$

$$\psi[A] \in \mathcal{H}^{anx}$$

$$\psi[A] \in \mathcal{H}^{anx}$$

$$\hat{G}_a \psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a} \psi = 0 \quad \text{is}$$

$$\Psi[A] \in \mathcal{H}^{aux}$$

$$\hat{G}_a \Psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a} \Psi = 0 \quad \text{is solved}$$

functionals invariant under local gauge tr. s

$g(x)$

$$\Psi[A] \in \mathcal{H}^{aux}$$

$$\hat{G}_a \Psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a} \Psi = 0 \quad \text{is solved}$$

functionals invariant under local gauge tr/s

$$g(x) : A$$

$$\Psi[A] \in \mathcal{H}^{anx}$$

$$\hat{G}_a \Psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a_i} \Psi = 0 \quad \text{is solved}$$

by functionals invariant under local gauge tr's

$$g \Rightarrow g(x) : A_i \rightarrow A_i^g = g^{-1} A_i g - g^{-1} \mathcal{D}_i g$$

$$\Psi[A] \in \mathcal{H}^{anx}$$

$$\hat{G}_a \Psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a_i} \Psi = 0 \quad \text{is solved}$$

by functionals invariant under local gauge tr. s

$\mathcal{G}_{SO(3)}$

$$\Rightarrow g(x) : A_i \rightarrow A_i^g = g^{-1} A_i g - g^{-1} \partial_i g$$

$$\psi[A] \in \mathcal{H}^{anx}$$

$$\hat{G}_a \psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a_i} \psi = 0 \quad \text{is solved}$$

by functionals invariant under local gauge tr. s

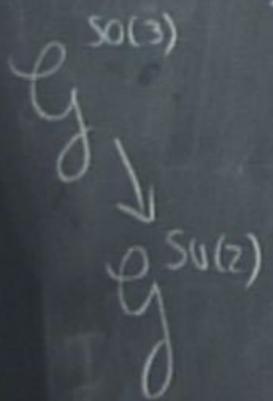
$g_{SO(3)}$

$$\Rightarrow g(x) : A_i \rightarrow A_i^g = g^{-1} A_i g - g^{-1} \mathcal{D}_i g$$

$$\Psi[A] \in \mathcal{H}^{anx}$$

$$\hat{G}_a \Psi \propto \mathcal{D}_i \frac{\delta}{\delta A^a_i} \Psi = 0 \quad \text{is solved}$$

by functionals invariant under local gauge tr. s



$$\Rightarrow g(x) : A_i \rightarrow A_i^g = g^{-1} A_i g - g^{-1} \mathcal{D}_i g$$

explicitly gauge-invariant functionals:

"Wilson loops"

explicitly gauge-invariant functionals:

"Wilson loops"

$$W_\gamma[A] = \text{Tr} \mathcal{P} \exp \oint_\gamma$$

explicitly gauge-invariant functionals:

"Wilson loops"

$$W_\gamma[A] = \text{Tr} \, \text{P exp} \oint_\gamma A = \text{Tr} \, \mathcal{H}(A; s_0, s_0)$$

explicitly gauge-invariant functionals:

"Wilson loops"

$$W_\gamma[A] = \text{Tr} \mathcal{P} \exp \oint_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

"path-ordered exponential of A along γ "

holonomy



explicitly gauge-invariant functionals:

"Wilson loops"

$$W_\gamma[A] = \text{Tr} \mathcal{P} \exp \oint_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

"path-ordered exponential of A along γ "

holonomy
↓

Σ



explicitly gauge-invariant functionals:

"Wilson loops"

holonomy

$$W_\gamma[A] = \text{Tr} \mathcal{P} \exp \oint_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

"path-ordered" exponential of A along γ



explicitly gauge-invariant functionals:

"Wilson loops"

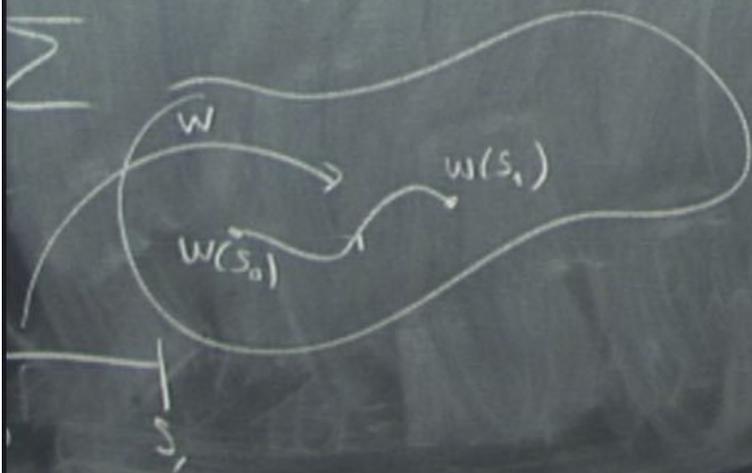
$$W_\gamma[A] = \text{Tr} \mathcal{P} \exp \oint_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

holonomy
↓

"path-ordered exponential of A along γ "

path in Σ

$$W: [s_1, s_2]$$



explicitly gauge-invariant functionals:

"Wilson loops"

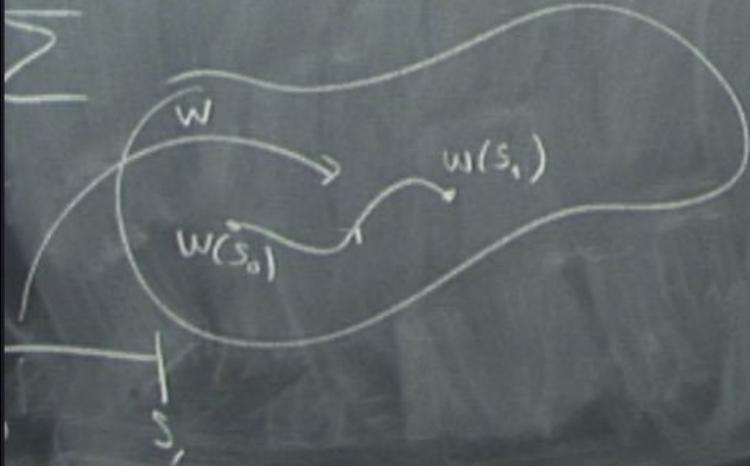
holonomy
↓

$$W_\gamma[A] = \text{Tr} \text{P exp} \int_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

"path-ordered exponential of A along γ "

path in Σ

$$W: [s_1, s_2] \rightarrow \Sigma, s \mapsto w^\mu(s)$$



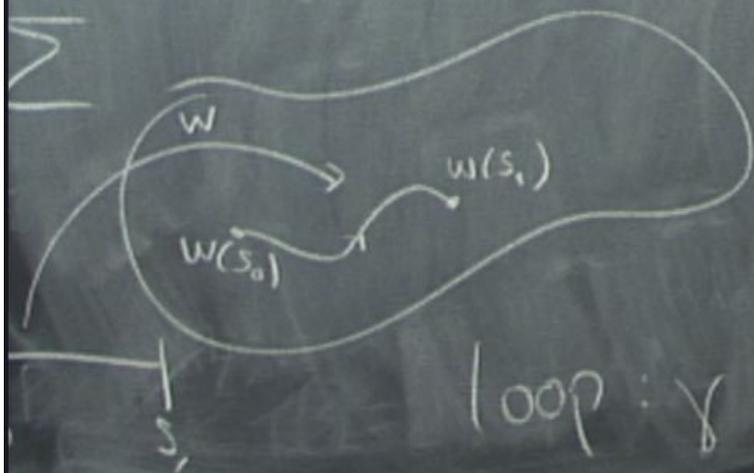
explicitly gauge-invariant functionals:

"Wilson loops"

$$W_\gamma[A] = \text{Tr} \text{P exp} \int_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

holonomy
↓

"path-ordered exponential of A along γ "



path in Σ

$$W: [s_1, s_2] \rightarrow \Sigma, s \mapsto w^\mu(s)$$

$$\text{loop } \gamma: [0, 1] \rightarrow \Sigma, s \mapsto \gamma^\mu(s), \gamma(0) = \gamma(1)$$

explicitly gauge-invariant functionals:

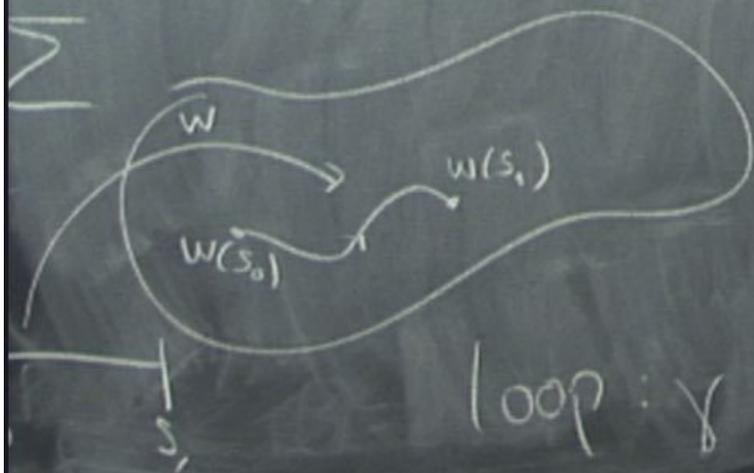
"Wilson loops"

holonomy
↓

$$W_\gamma[A] = \text{Tr} \text{P exp} \int_\gamma A = \text{Tr} U_\gamma[A; s_0, s_0]$$

"path-ordered" exponential of A along γ

path in Σ



$$W: [s_1, s_2] \rightarrow \Sigma, s \mapsto W^\mu(s)$$

$$\text{loop } \gamma: [0, 1] \rightarrow \Sigma, s \mapsto \gamma^\mu(s), \gamma(0) = \gamma(1)$$

$$A(x) = A_i(x) dx^i = A_i^a(x) X_a dx^i$$

3 algebra generators in the fund. repr. of $SU(2)$

$$A(x) = A_i(x) dx^i = A_i^a(x) X_a dx^i \quad \text{Fix the gauge } (s_0, s_1)$$

3 algebra generators in
the fund repr. of $SU(2)$

The holonomy path $w^M(s)$ with initial point s_0 and end-point

$$\frac{dU_w(s, s_0)}{ds} = A_i^a(x) \frac{dw^i}{ds} U_w(s, s_0), \quad s_0 \leq s \leq s_1$$

$$A(x) = A_i(x) dx^i = A_i^a(x) X_a dx^i \quad \text{Fix the gauge } (s, s_0)$$

3 algebra generators in
the fund repr. of $SU(2)$

$$U_w(s, s_0)$$

The holonomy of a path $w^M(s)$ with initial point s_0 and end

point s_1 is the solution of $\frac{dU_w(s, s_0)}{ds} = A_i(x) \frac{dw^i}{ds} U_w(s, s_0)$

where $x = w(s)$, with initial cond. $U_w(s_0, s_0) = 1$

$$A(x) = A_i(x) dx^i = A_i^a(x) X_a dx^i \quad \text{Fix the gauge } (s_0, s_1)$$

3 algebra generators in
the fund repr. of $SU(2)$

$$U_w(s, s_0)$$

The holonomy of a path $w^M(s)$ with initial point s_0 and end-

point s_1 is the solution of
$$\frac{dU_w(s, s_0)}{ds} = A_i(x) \frac{dw^i}{ds} U_w(s, s_0),$$

where $x = w(s)$, with initial cond. $U_w(s_0, s_0) = \mathbb{1}$
 ← unit in $SU(2)$