

Title: String Theory - Review (PHYS 623) - Lecture 11

Date: Feb 08, 2010 11:20 AM

URL: <http://pirsa.org/10020055>

Abstract:

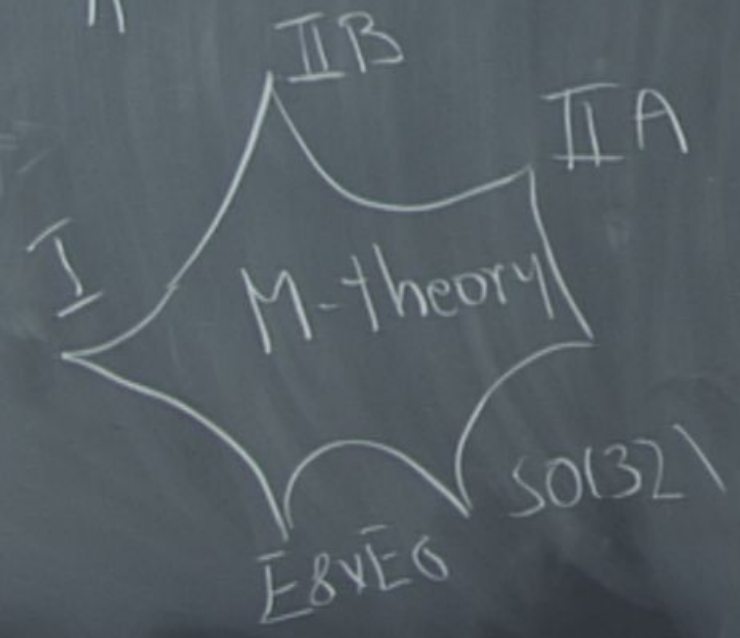


perimeter scholars
INTERNATIONAL

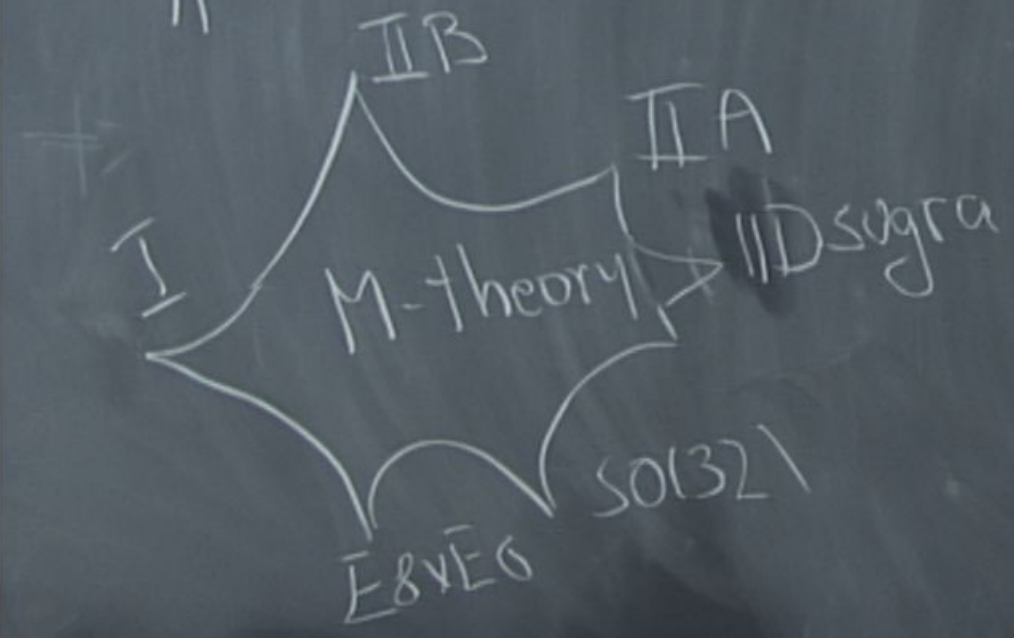
Type IIA, IIB, I

Type IIA, IIB, I heterotic $E_8 \times E_8$, $SO(32)$

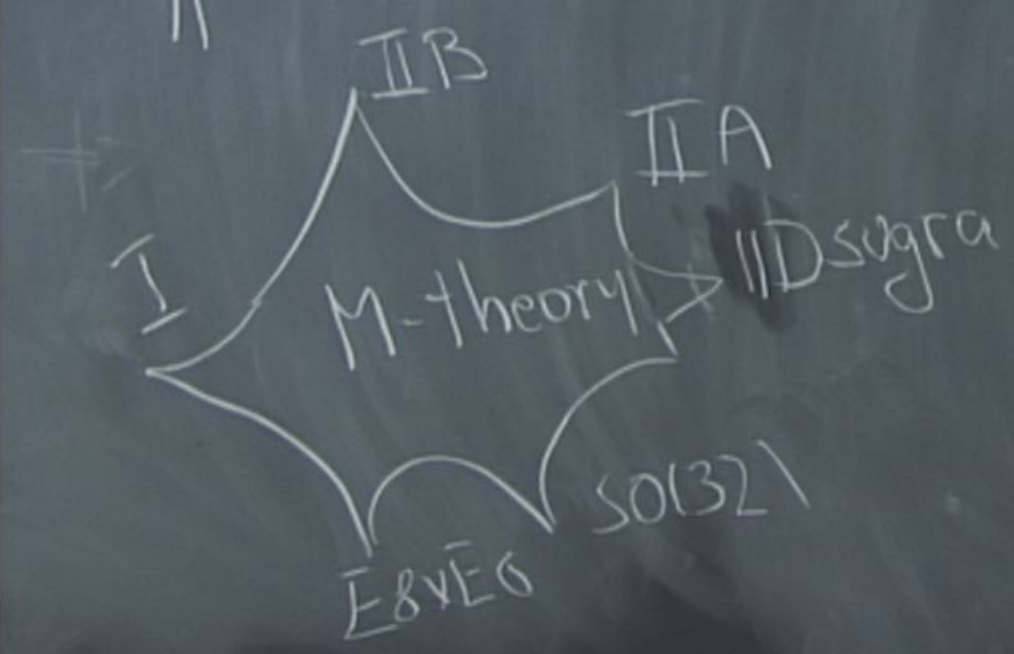
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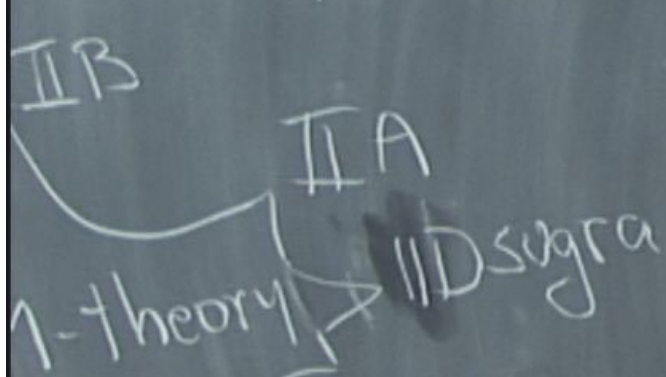


Type IIA, IIB, I heterotic $E_8 \times E_8$, $SO(32)$



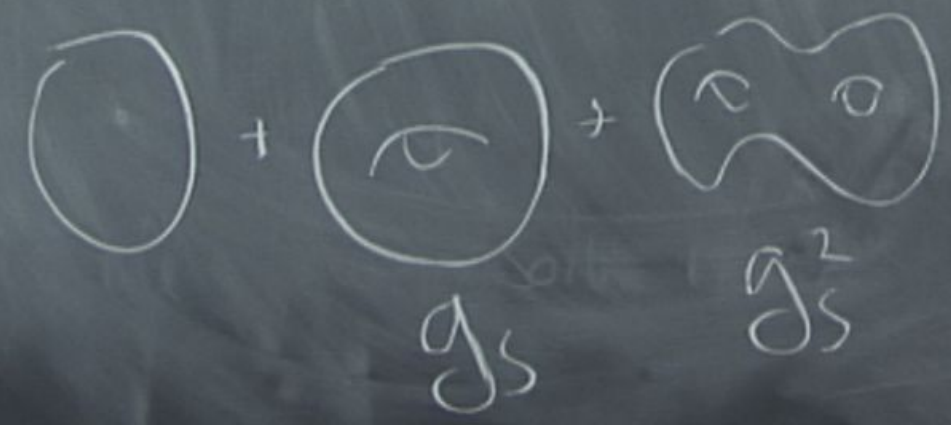
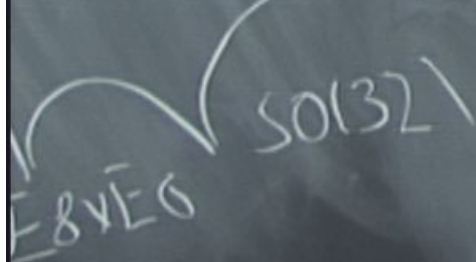
T-duality $R \sim \frac{1}{R}$
S-duality

A, IIB, I heterotic $E_8 \times E_8, SO(32)$



T-duality $R \sim \frac{1}{R}$ (perturbative)

S-duality



String Duality

① Low-energy effective actions
 $\alpha' \rightarrow 0$ (massless dof).

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 $\alpha' \rightarrow 0$ (massless dof).

Enough susy \Rightarrow Entire string

② BPS states (compare branes)

$$S = \frac{1}{2\pi G_D} \int d^D x \sqrt{-g} R$$

$[G_D] = \text{length}^{D-2}$

HID SUGRA

H D SUGRA 1978

Non-renormalizable

HID SUGRA 1978

Non-renormalizable

Not chiral

HID SUGRA 1978

Non-renormalizable

Not chiral : M2, M5, M9

11D SUGRA 1978

Non-renormalizable

(Not chiral : $M2, M5, M9 \Rightarrow$ Heterotic M-theory!

H D SUGRA 1978

Non-renormalizable

(Not chiral : M2, M5, M9 \Rightarrow Heterotic theory!

Field content

X Graviton

H D SUGRA 1978

Non-renormalizable

(Not chiral : M2, M5, M9 \Rightarrow Heterotic M-theory!

Field content

X Graviton $(D-2)(D-2)$ Matrix
 $\frac{1}{2}(D-1)(D-2) = 44$ bosonic

(2) Antisymmetric tensor A_3 $F_4 = dA_3$.

(1) Low-energy effective actions

dA_3 (massless dof)

$F_4 \rightarrow$ Antisymmetric

F_4 (massive dof)

② Antisymmetric tensor A_3 $F_4 = dA_3$.

$A_{MNP} = 84$ bosonic dof.

$$\frac{1 \cdot 9 \cdot 8 \cdot 7}{2}$$

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$$\frac{1 \cdot 9 \cdot 8 \cdot 7}{2}$$

$\Rightarrow 128$ bosonic dof

③ Gravitino
128 dof

$\psi_M \leftarrow$ vector index 32-component Majorana
 $\chi_M \leftarrow$ spinor index

Action

$$\kappa_{11}^2 S = \int dx \sqrt{-g} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

[60] = length

Action

$$\kappa_{11}^2 S = \int d^4x \sqrt{-g} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

$$\uparrow \quad \uparrow$$
$$(6\pi G_{11}) = 2\kappa_{11}^2 =$$

\uparrow Gravitational constant

Newton's const.

Action

$$\alpha_{11}^2 S = \int d^4x \sqrt{-g} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

$$\frac{1}{16\pi} G_{11} = 2\alpha_{11}^2 = \frac{1}{2\pi} (2\pi \ell_P^g) \leftarrow \text{11D Planck length}$$

\uparrow Gravitational constant

\uparrow Newton's const

$$|F_m|^2 = \frac{1}{m!} G^{M_1 N_1} \dots G^{M_n N_n} F_{M_1 \dots M_n} F_{N_1 \dots N_n}$$

Action

$$\kappa_{11}^2 S = \int d^4x \sqrt{-g} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

$$\uparrow \quad \uparrow \quad \uparrow \quad \leftarrow \text{11D Planck length}$$
$$6\pi G_{11} = 2\kappa_{11}^2 = \frac{1}{2\pi} (2\pi \ell_P^g)$$

\uparrow Gravitational constant

Newton's const

\downarrow Chern-Simons term

$$|F_m|^2 = \frac{1}{m!} G^{M_1 N_1} \dots G^{M_n N_n} F_{M_1 M_n} F_{N_1 \dots N_n} \\ F_{1234} F^{1234} + \dots$$

$$\Psi_M = \nabla_M \xi + \frac{1}{12} (\Gamma_M F^{(4)} - 3 F_M^{(4)}) \xi + (\text{fermi})^2$$

$$\Psi_M = \nabla_M \mathcal{E} + \frac{1}{12} (\Gamma_M F^{(4)} - 3 F_M^{(4)}) \mathcal{E} + (\text{fermi})^2$$

$$F^{(4)} = \frac{1}{4!} F_{MNPQ} \Gamma^{MNPQ}$$

$$+ \frac{1}{12} (\Gamma_M F^{(4)} - 3 F_M^{(4)}) \epsilon + (\text{fermi})^2$$

$$F^{(4)} = \frac{1}{4!} F_{MNPQ} \Gamma^{MNPQ} \quad \Gamma^{0123} \quad \text{antisymm}$$

$$F_M^{(4)} = \frac{1}{3!} F_{MNPQ} \Gamma^{NPO}$$

$$\Psi_M = \nabla_M \epsilon + \frac{1}{12} (\Gamma_M F^{(4)} - 3 F_M^{(4)}) \epsilon + (\text{fermi})^2$$

(check!)

$$F^{(4)} = \frac{1}{4!} F_{MNPQ} \Gamma^{MNPQ}$$

$$F_M^{(4)} = \frac{1}{3!} F_{MNPQ} \Gamma^{NPO}$$

Γ^{1234}
+ ...
anti

$$3 F_M^{(4)} \epsilon + (\text{fermi})^2 = 0$$

$$F_{MN PQ} \Gamma^{MNPQ}$$

$$\frac{1}{3!} F_{MN PQ} \Gamma^{MNPQ}$$

$$\Gamma^{01234}$$

antisymm

$$+ \frac{1}{12} (\Gamma_M F^{(4)} - 3 F_M^{(4)}) \epsilon + (\text{fermi})^2 = 0$$

$\kappa!$

$$F^{(4)} = \frac{1}{4!} F_{MNPQ} \Gamma^{MNPQ}$$

$$F_M = \frac{1}{3!} F_{MNPQ} \Gamma^{MNPQ}$$

$$\pi_1 \pi_2 \pi_3 \pi_4$$

+ ... antisymm

\Rightarrow Classical susy solutions M2, M5 branes

$$A_3: \int d^3\sigma A_3.$$

↑ 3D extended object.

$F_4 : A_3 : \int d^3\sigma A_3$ M2-brane
↑ 3D extended object.

⊕
* $\bar{F}_4 = \bar{F}_7 ; \int d^6\sigma C_6$
" dC_6

$F_4 : A_3 : \int d^3\sigma A_3$ M2-brane
↑ 3D extended object.

$$*\overline{F}_4 = \overline{F}_7 ; \int d^6\sigma C_6$$

"
 dC_6

$F_4 : A_3 : \int d^3\sigma A_3$ M2-brane
↑ 3D extended object.

$*F_4 = F_7 ; \int d^6\sigma C_6$
" dC_6

$F_4 = dA_3$

$$F_4 : A_3 : T_2 \int d^3 \sigma A_3 \quad \text{M2-brane}$$

↑ 3D extended object.

$$*\bar{F}_4 = F_7 ; T_5 \int d^6 \sigma C_6 \quad \text{M5-brane}$$

$$dC_6$$

$$F_4 = dA_3$$

$$T_2 = 2\pi (2\pi l_p)^{-3}$$

$$T_5 = 2\pi (2\pi l_p)^{-6}$$

Action

$$\int_{\mathcal{M}} S = \int d^D x \sqrt{-g} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

\Rightarrow IIA supergravity.

Compactify on circle $X_{11} \rightarrow X_{11} + 2\pi R$.

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Metric Ansatz $\mu, \nu \dots$ $M, N \dots$
10D indices 11 D indices.

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Metric Ansatz $\mu, \nu \dots$ $M, N \dots$
10 D indices 11 D indices

Kaluza-Klein ansatz:

$$G_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

$$G_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

dilaton

$A, \phi, g^{\mu\nu}$ don't depend on X^9

Action

$$\int d^2 X \int d^4 X \sqrt{-G} \left(R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int A_3 \wedge F_4 \wedge F_4$$

\Rightarrow IIA supergravity.

$$\int d^4 X \sim R$$



$$G_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

dilaton

$A^{(1)}, \phi, g^{\mu\nu}$ don't depend on X^9

$$\textcircled{1} \Gamma_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

dilaton

$A^{(1)}, \phi, g^{\mu\nu}$ don't depend on X^9

$$\textcircled{2} F_{MNPQ}$$

$$\textcircled{1} \Gamma_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

dilaton

$A^{(1)}, \phi, g^{\mu\nu}$ don't depend on X^9

$$\textcircled{2} F_{MNPQ} \left\{ \begin{array}{l} F_{mmp11} = H_{mmp} \\ H_3 = dB_2 \end{array} \right.$$

$$\textcircled{1} \quad G_{MN} = e^{-2\phi/3} \begin{pmatrix} g_{\mu\nu} + e^{2\phi} A_\mu A_\nu & e^{2\phi} A_\mu \\ e^{2\phi} A_\nu & e^{2\phi} \end{pmatrix}$$

10D

dilaton

$A^{(1)}, \phi, g^{\mu\nu}$ don't depend on X^9

$$\textcircled{2} \quad F_{MNPQ} \begin{cases} F_{mmpq} = H_{mmp} & H_3 = dB_2 \\ F_{mmpq} \Rightarrow F_4 = dC_3 \end{cases}$$

Action

$$S = \int d^4x \left(R - \frac{1}{2} |\tilde{F}_4|^2 \right) - \frac{1}{6} \int A_3 \wedge \tilde{F}_4 \wedge \tilde{F}_4$$

$$\tilde{F}_4 = dA_3 + A_1 \wedge H_3$$

↑ gauge invariant

Action

$$S = \int d^4x \left(R - \frac{1}{2} \underline{\underline{|\tilde{F}_4|^2}} \right) - \frac{1}{6} \int A_3 \wedge \tilde{F}_4 \wedge \tilde{F}_4$$

$$\tilde{F}_4 = dA_3 + A_1 \wedge H_3$$

↑ gauge invariant

$$A_3 \rightarrow A_3 + d\Lambda_2$$

Gauge transformations

Action

$$S = \int d^4x \left(R - \frac{1}{2} \|\tilde{F}_4\|^2 \right) - \frac{1}{6} \int A_3 \wedge \tilde{F}_4 \wedge \tilde{F}_4$$

$$H_3 = dA_3 + A_1 \wedge H_3$$

Gauge transformations

$$\delta A_1 = d\Lambda$$

$$\delta A_3 = d\Lambda \wedge B_2$$

Action

$$\int_{\mathbb{R}^4} S = \int d^4x \left(R - \frac{1}{2} |\tilde{F}_4|^2 \right) - \frac{1}{6} \int A_3 \wedge \tilde{F}_4 \wedge \tilde{F}_4$$

$$\tilde{F}_4 = dA_3 + A_1 \wedge H_3$$

Gauge transformations

$$\delta A_1 = d\Lambda$$

$$\delta \tilde{F}_4 = d(d\Lambda \wedge B_2) + d\Lambda \wedge H_3 = 0 \quad \delta A_3 = d\Lambda \wedge B_2$$

Action II A sugra (integrate over X'')
 $S = S_{NS} + S_{RT} + S_{CS}$.

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 $S = S_{NS} + S_{RT} + S_{CS}$.

$S_{NS} =$

Action IIA sugra (integrate over X'')

$$S = S_{NS} + S_R + S_{CS}$$

$$S_{NS} = \frac{1}{2\alpha'^2} \int d^{10}x \sqrt{-g}$$

Action IIA sugra (integrate over X'')

$$S = S_{NS} + S_R + S_{CS}$$

$$S_{NS} = \frac{1}{2\alpha'^2} \int d^{10}x \sqrt{-g} e^{-2\phi} \left(R + 4(\partial\phi)^2 - \frac{1}{2} H_3^2 \right)$$

Action IIA sugra (integrate over X'')

$$S = S_{NS} + S_R + S_{CS}$$

$$S_{NS} = \frac{1}{2\alpha_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} \left(R + 4(\partial\phi)^2 - \frac{1}{2} H_3^2 \right)$$

$$\alpha_{10}^2 \sim \ell_{10}^2$$

Action IIA sugra (integrate over X'')

$$S = S_{NS} + S_R + S_{CS}$$

$$S_{NS} = \frac{1}{2\alpha_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} \left(R + 4(\partial\phi)^2 - \frac{1}{2} H_3^2 \right)$$

$$\alpha_{10}^2 \sim G_{10} \Rightarrow G_{11} = 2\pi R G_{10}$$

Action IIA sugra (integrate over X'')

$$S = S_{NS} + S_R + S_{CS}$$

$$S_{NS} = \frac{1}{2\alpha_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} \left(R + 4(\partial\phi)^2 - \frac{1}{2} H_3^2 \right)$$

$$\alpha_{10}^2 \sim G_{10} \Rightarrow G_{11} = 2\pi R \int dx'' G_{10}$$

Action

String frame action

Einstein frame $\int d^{10}x \sqrt{-g} R$

Action

SR =

Action

$$S_R = -\frac{1}{4\alpha_{10}^2} \int d^{10}x \sqrt{-g} \left(|F_2|^2 + |\tilde{F}_4|^2 \right)$$

\uparrow
 $F_2 = dA_1$

Action

$$S_R = -\frac{1}{4\kappa_{10}^2} \int d^{10}x \sqrt{-g} \left(|F_2|^2 + |\tilde{F}_4|^2 \right)$$

\uparrow
 $F_2 = dA_1$

$$S_{CS} = -\frac{1}{4\kappa^2} \int B_2 \wedge F_4 \wedge F_4$$

Action

$$S_R = -\frac{1}{4\alpha_{10}^2} \int d^{10}x \sqrt{-g} \left(|F_2|^2 + |F_4|^2 \right)$$

↑
 $F_2 = dA_1$

$$S_{CS} = -\frac{1}{4\alpha^2} \int B_2 \wedge F_4 \wedge F_4$$

$\Psi_M \rightarrow \Psi_\mu + \lambda \rightarrow$ dilatino