Title: A Computational Grand-Unified Theory

Date: Feb 03, 2010 04:00 PM

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Abstract: Are Quantum Mechanics and Special Relativity unrelated theories? Is Quantum Field Theory an additional theoretical layer over them? Where the quantization rules and the Plank constant come from? All these questions can find answer in the computational paradigm: "the universe is a huge quantum computer".

In my talk I'll take the computational-universe paradigm as genuine theoretical framework, and analyze some relevant implications. A new kind of quantum field theory emerges: "Quantum-Computational Field Theory" (QCFT). I will show how in QCFT Special Relativity unfolds from the fabric of the computational network, which also naturally embeds gauge-invariance, and even the quantization rule and the Planck constant, which thus resort to being properties of the underlying causal tapestry of space-time. In this way Quantum Mechanics remains the only theory needed to describe the computational-universe. I will analyze few simple toy-models in order to explore the mathematical structure of QCFT.

The new QCFT has many advantages versus the customary field theoretical framework, solving a number of logical and mathematical problems that plague quantum field theory. One further advantage of QCFT is the possibility of changing the computational engine without changing the field-theoretical framework. One can thus consider different kind of engines, e.g. classical, quantum, super-quantum, and even input-output networks with no pre-established causal relations, which are very interesting for addressing the problem of Quantum Gravity.

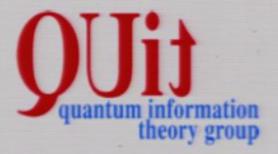
QCFT opens a large research line: I argue that this program should be addressed soon in the particle physics domain, before entering Quantum Gravity, notwithstanding the experimental success of the usual quantum field theory. It will also be the first test of the Lucien Hardy's program on Quantum Gravity.

Reference: arXiv:1001.1088 (http://arxiv.org/abs/1001.1088)

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A COMPUTATIONAL GUT

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Dipartimento di Fisica "A. Volta", Università di Pavia

WHY QUANTUM?

WHY QUANTUM? WHY RELATIVITY?



WHY QUANTUM? WHY RELATIVITY?

* Are Quantum Theory and Special Relativity unrelated theories?

* Is Quantum Field Theory (QFT) an additional theoretical layer over QT and SR?

* Where the quantization rules and the Planck constant come from?





HOWRELATIVITY EMERGES FROM THE COMPUTATION?

Lorentz transformations from Galileo principle

Lorentz transformations from Galileo principle

- * Galileo principle includes homogeneity and isotropy of space and homogeneity of time.
- *On the assumption of isotropy and homogeneity of space and homogeneity of time along with symmetry between the two references, the most general transformations of reference system are the Lorentz transformations with a parameter Ω with the dimensions of a velocity, which is independent on the relative velocity of frames.

Special Relativity from computational network

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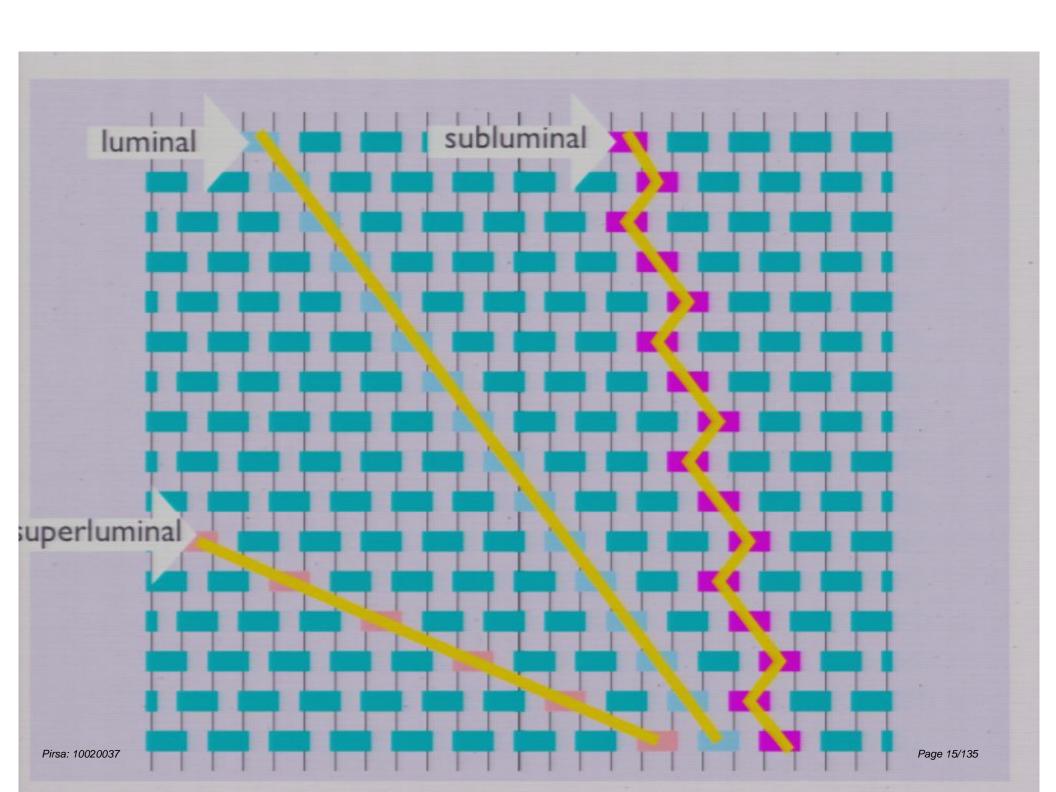
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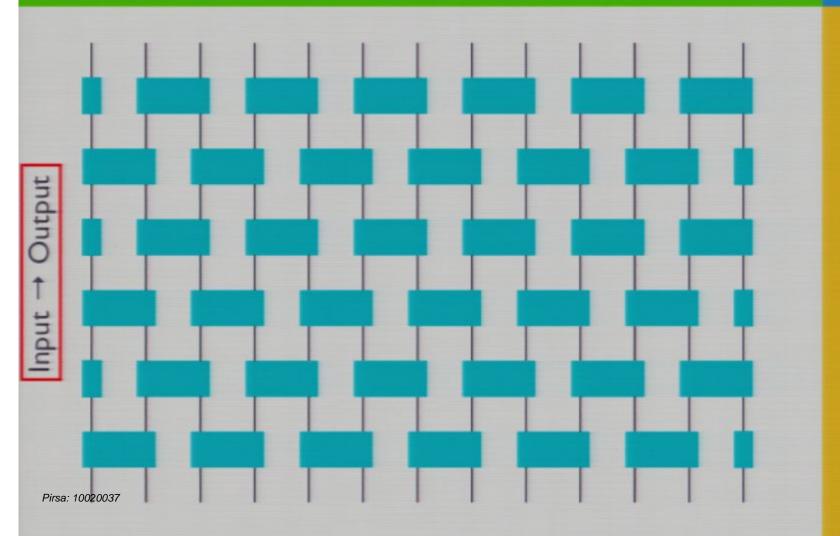
Special Relativity from computational network

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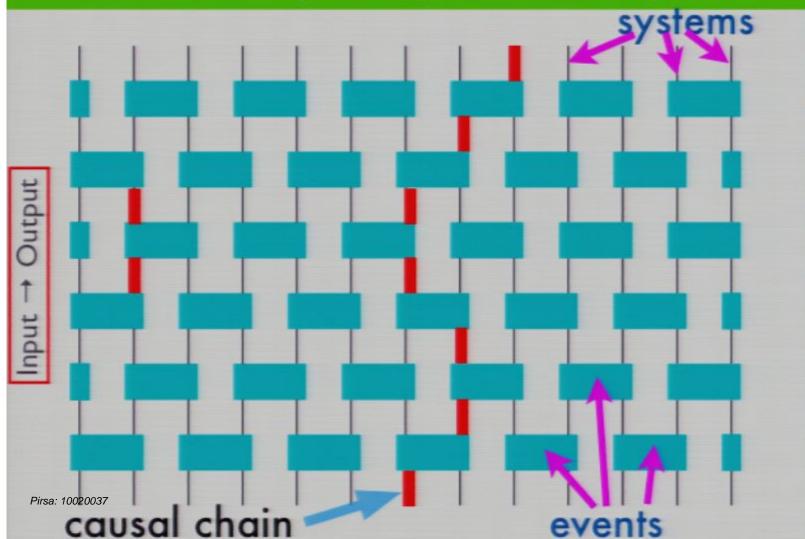
- *Take a computational circuit which is uniform and isotropic.
- *Take the "continuum limit" → space-time.



Relativity from QT (more generally from causality)



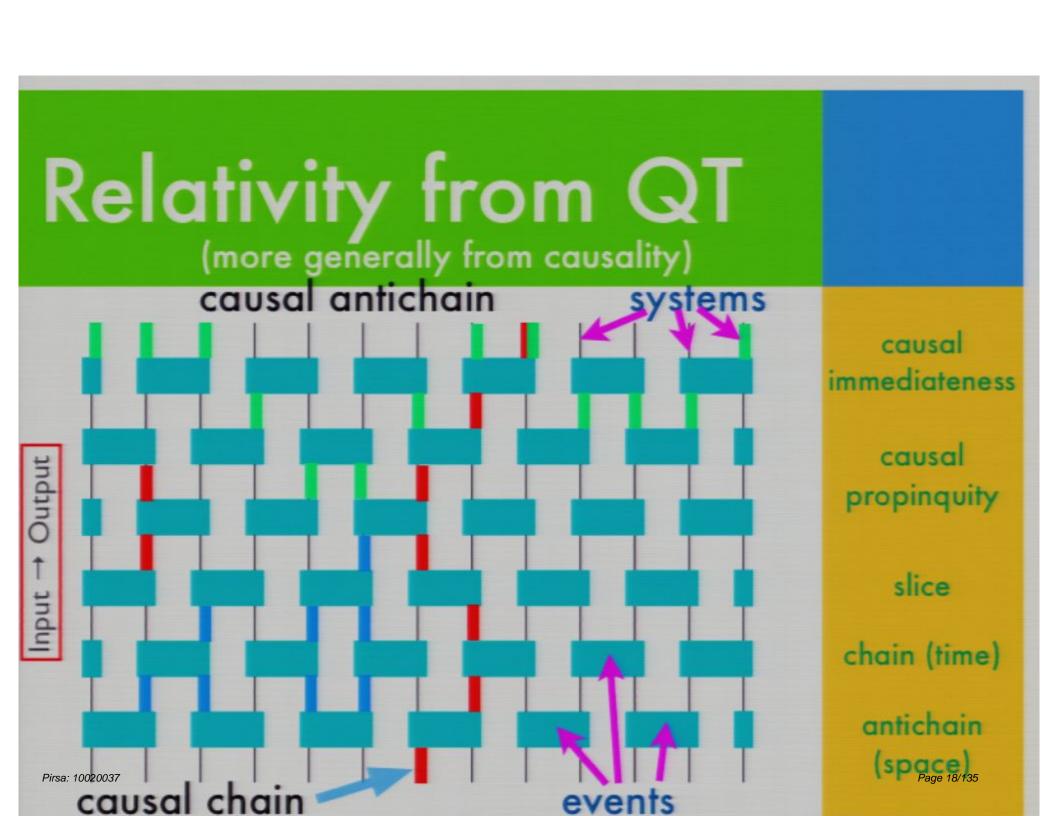
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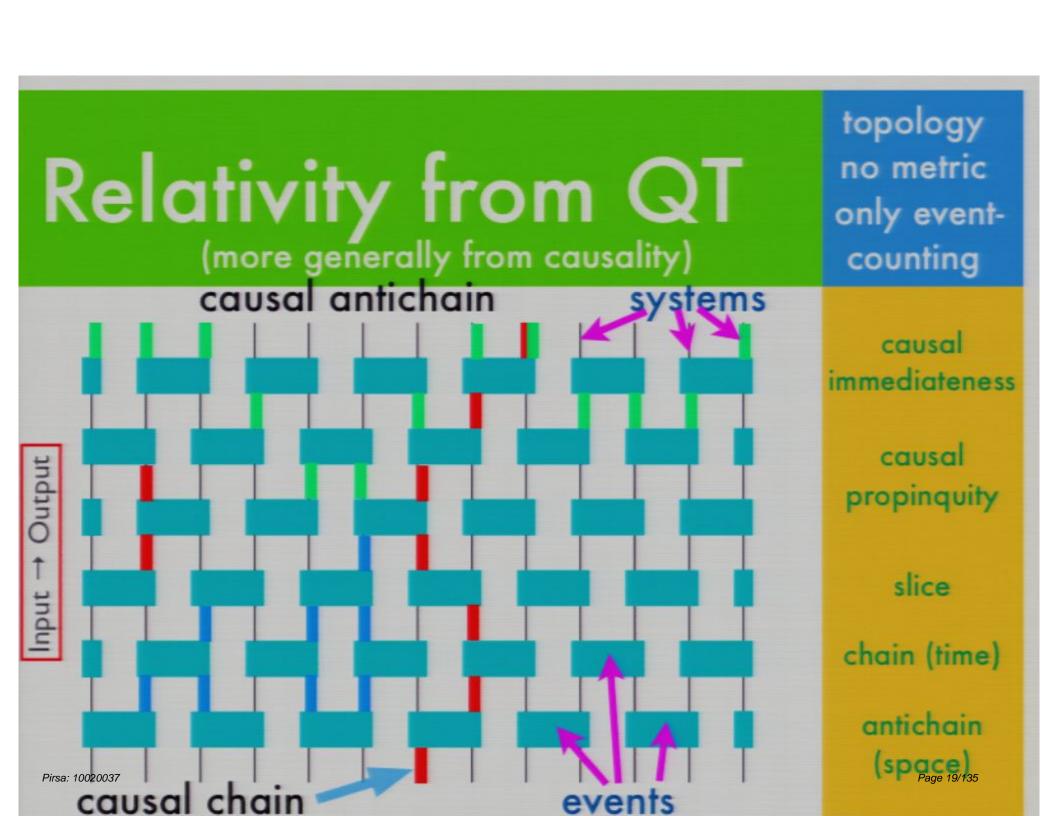


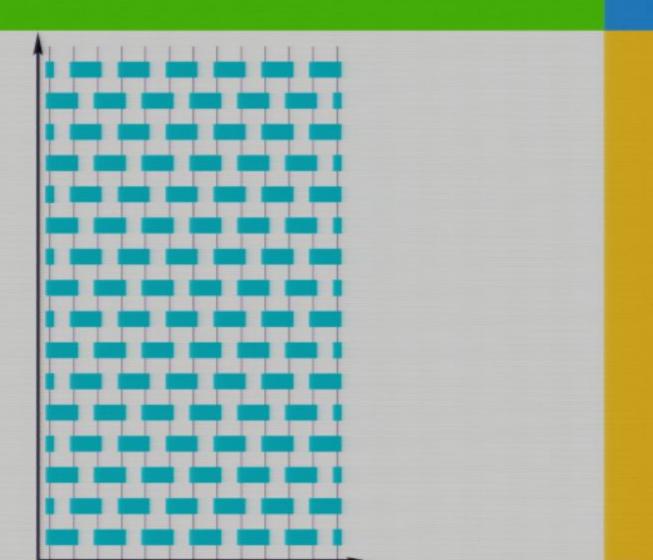
causal immediateness

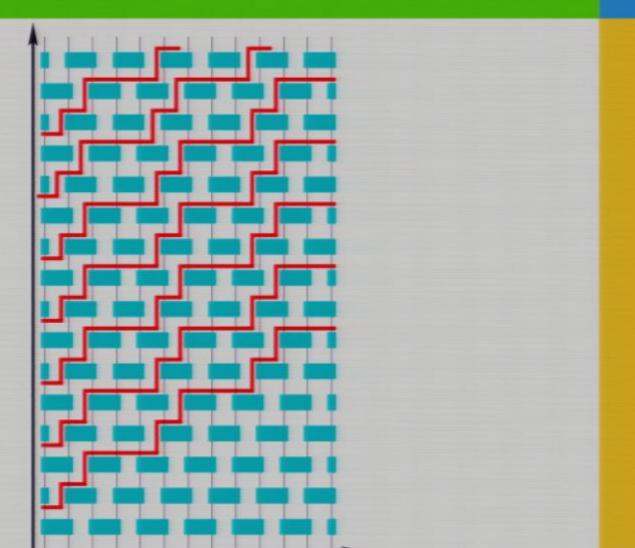
chain (time)

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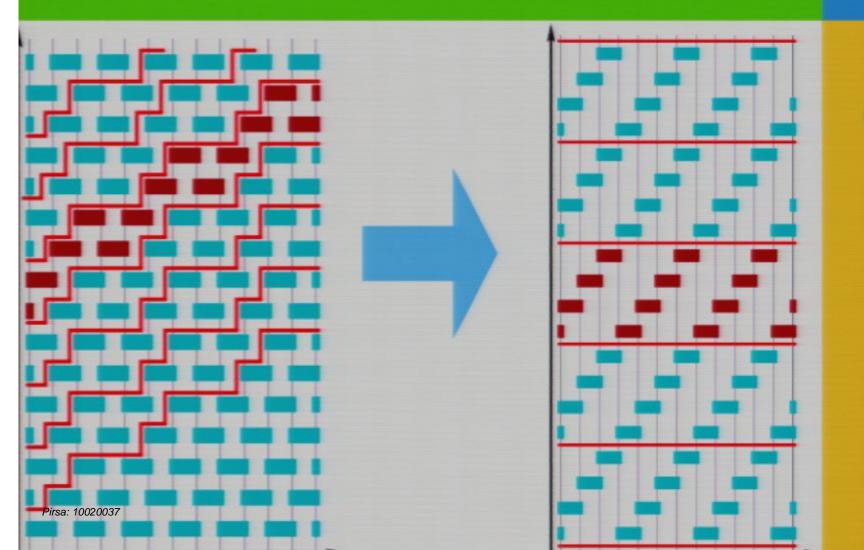




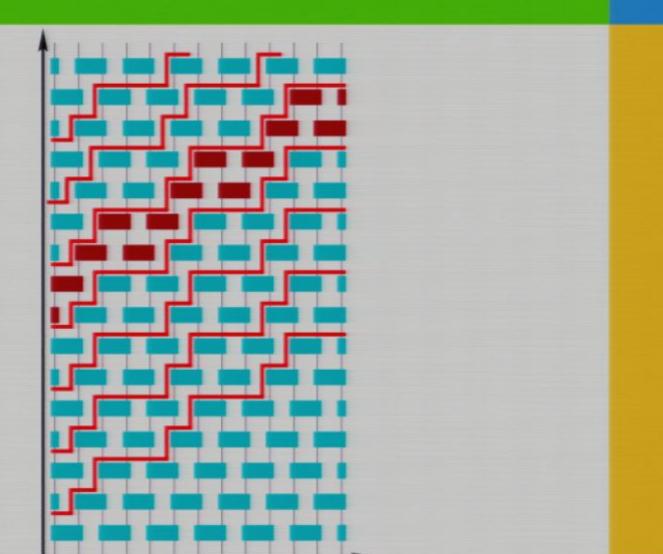


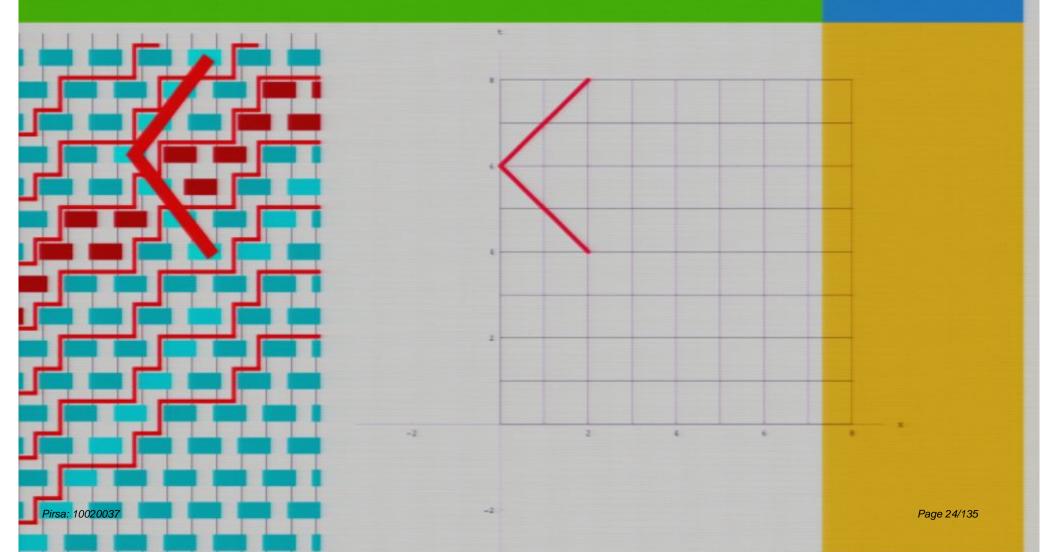


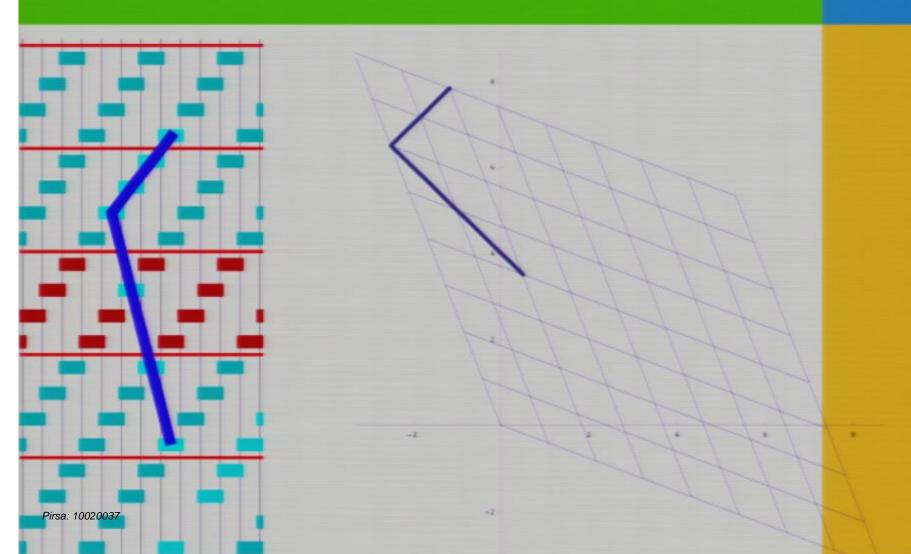
build a uniform foliation



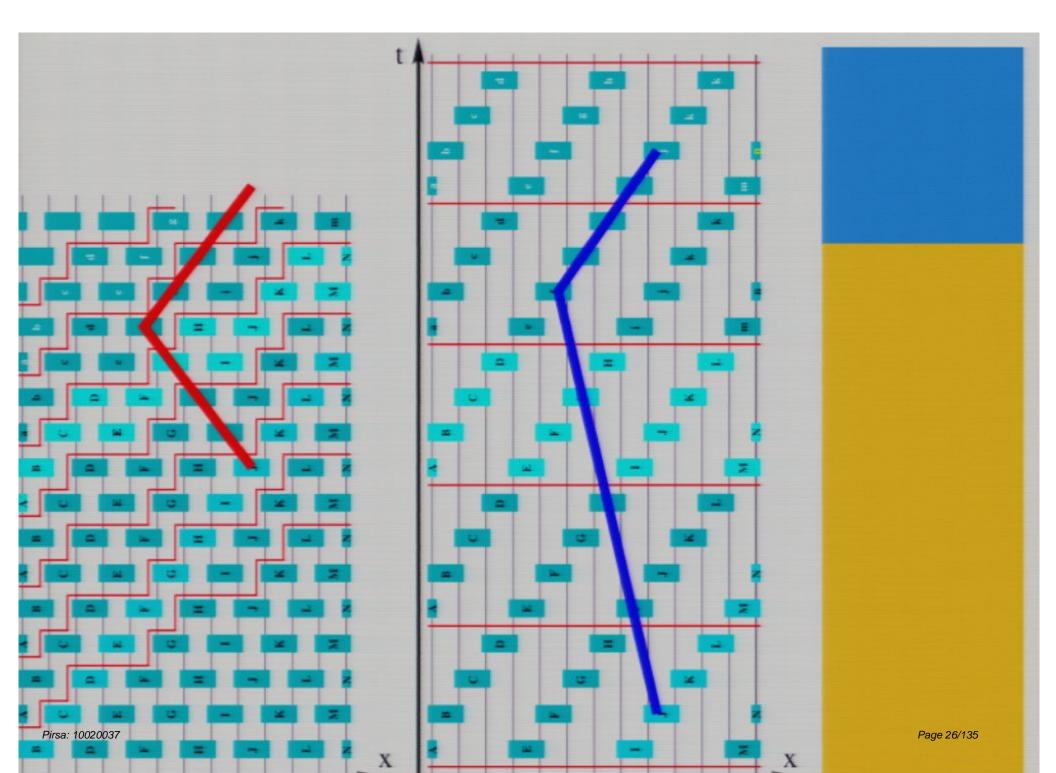
change reference

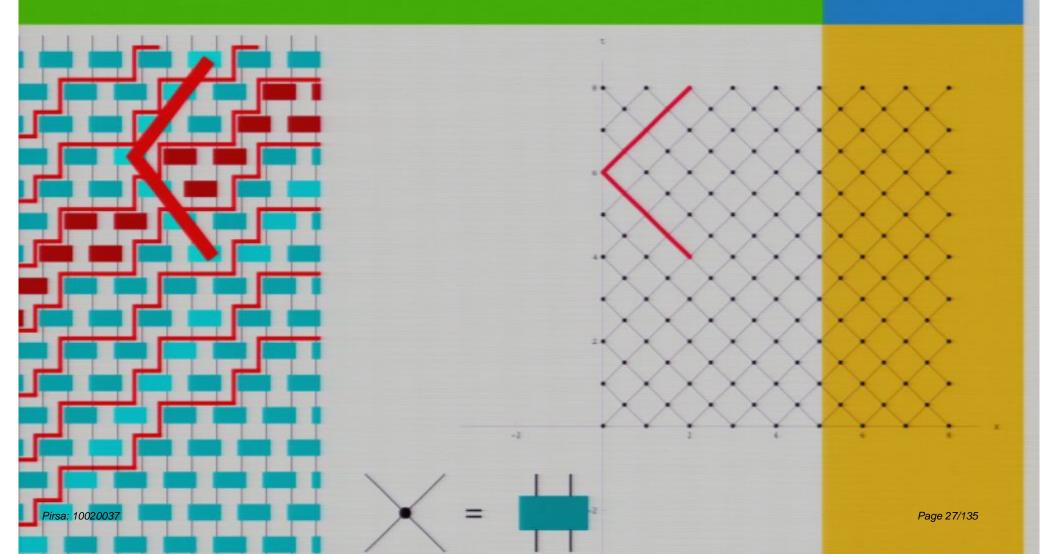


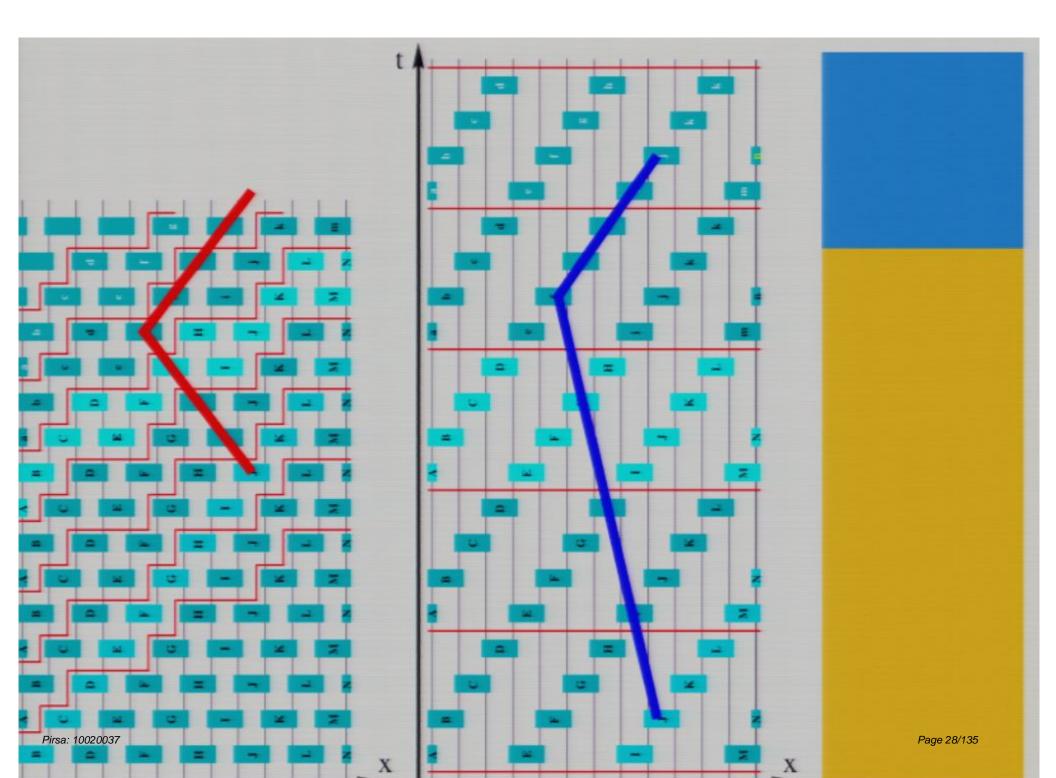




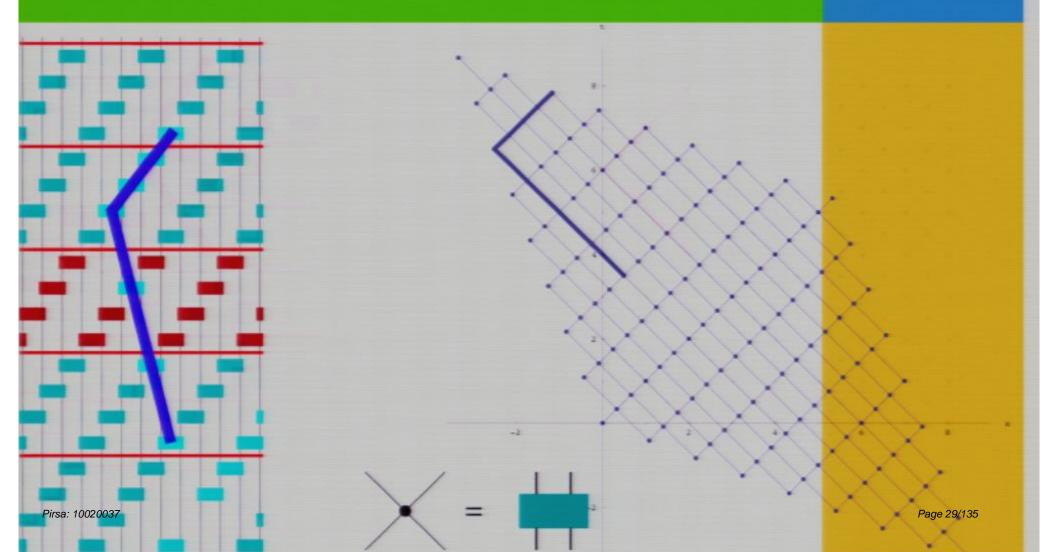
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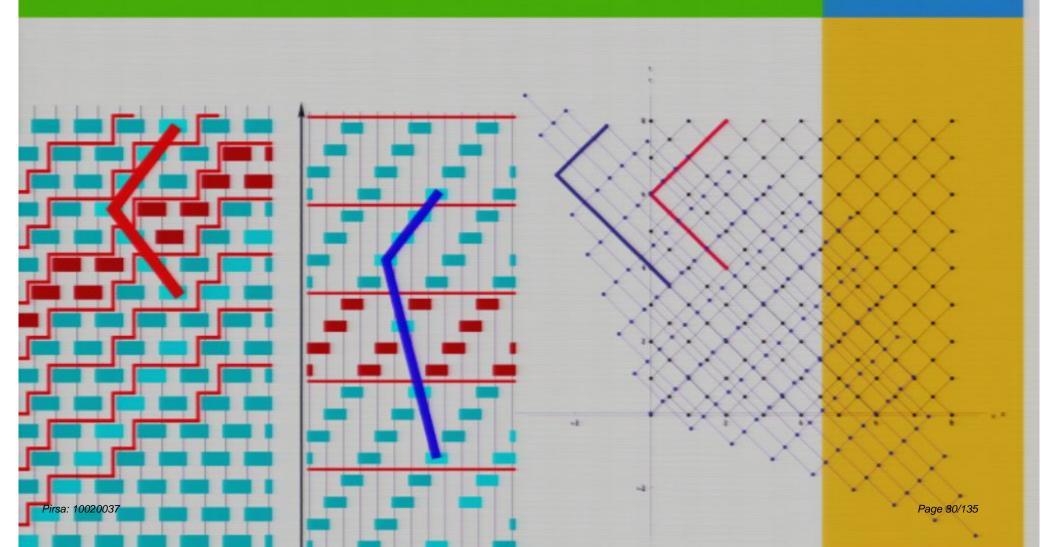


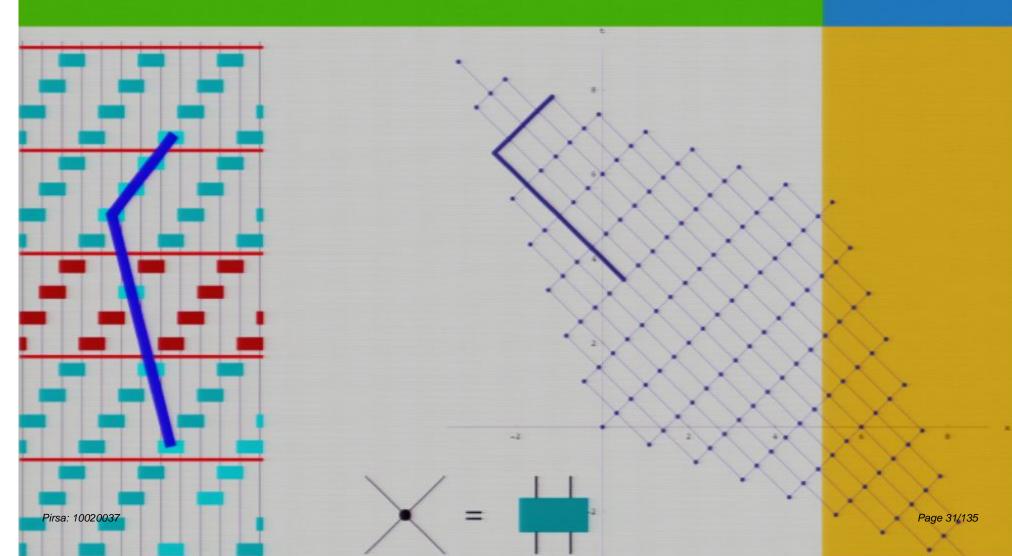


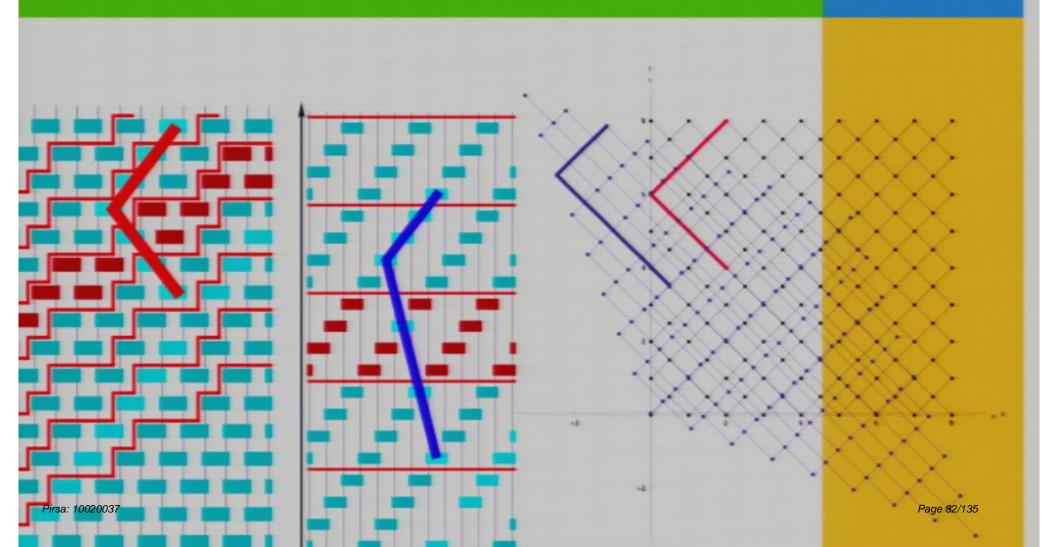


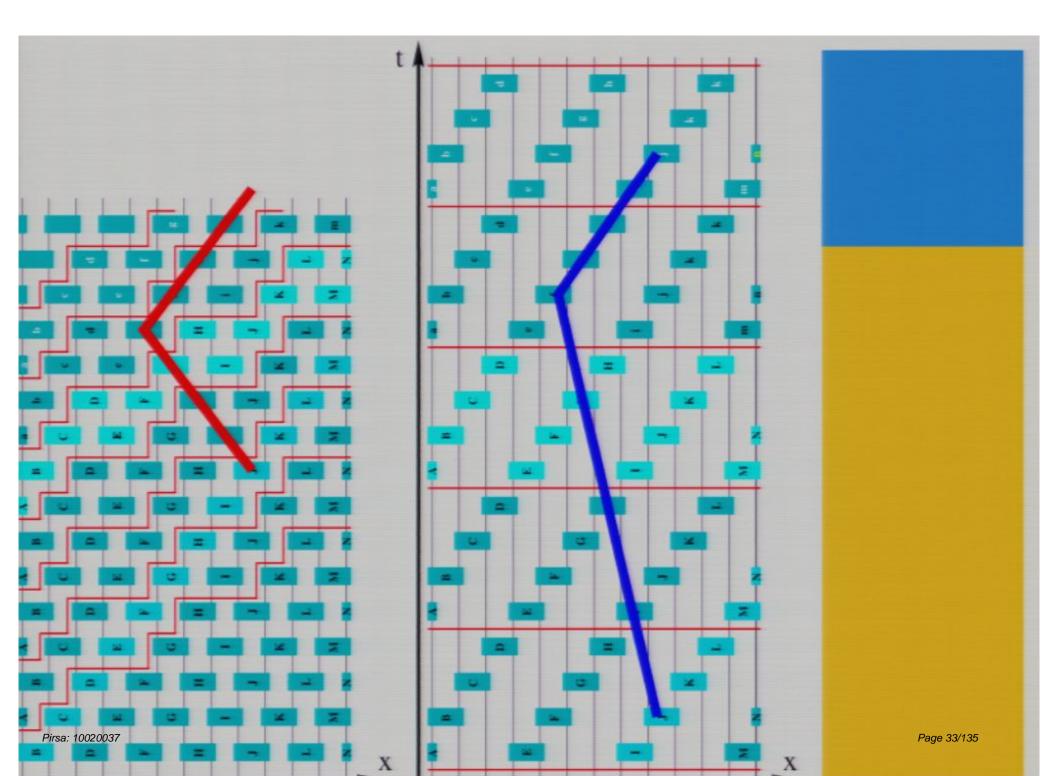


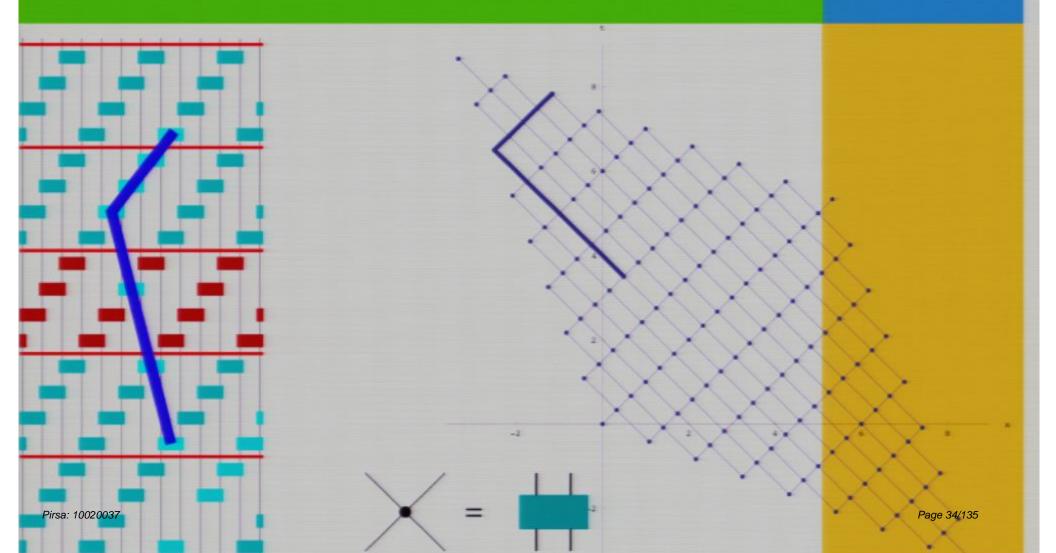


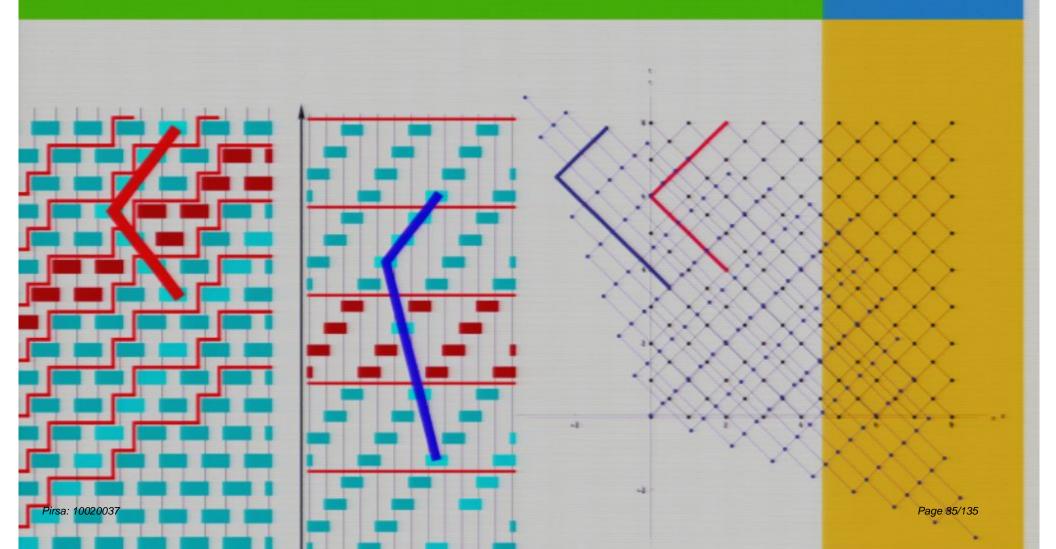










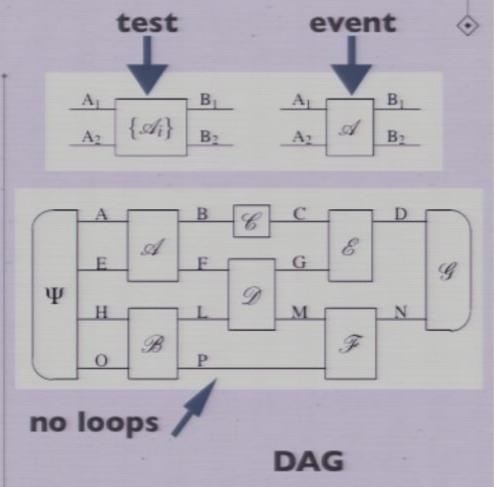


WE GOT SR FROM PURE CAUSALITY!

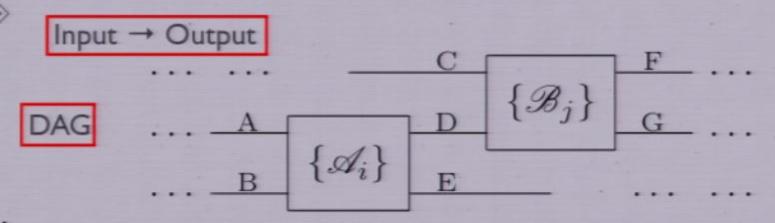
The Operational Framework

* Operational theory:
collection of systems closed
under parallel composition,
and of tests closed under
parallel/sequential
composition and under
randomization.

* Probabilistic operational theory: every test from the trivial system to the trivial system is associated to a probability distribution of outcomes.

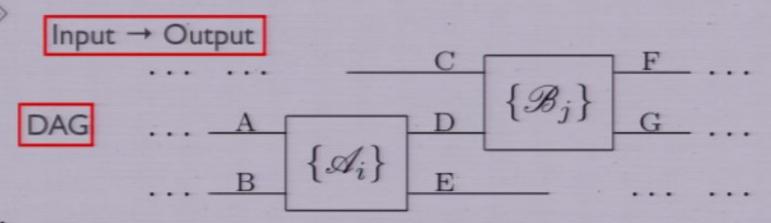


Causal probabilistic theories



A theory is causal, if for any two tests and that are connected (without loops) the marginal probability of the input event is independent on the choice of the output test, whereas, viceversa the marginal probability of the output event generally depends on the choice of the input test.

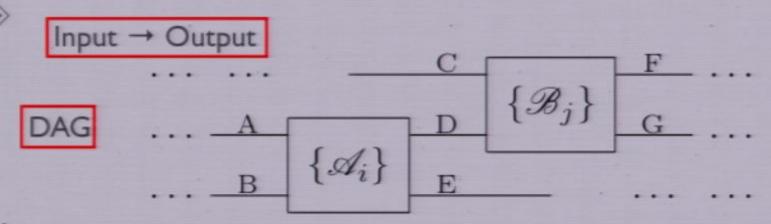
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Thm. A theory is causal iff the deterministic effect is unique for each system.

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Thm. A theory is causal iff the deterministic effect is unique for each system.

Thm. A theory is causal iff the normalization of states is a multiplication by

Wittgenstein-ism

1 The world is all that is the case.

1.1 The world is the totality of facts, not of things.

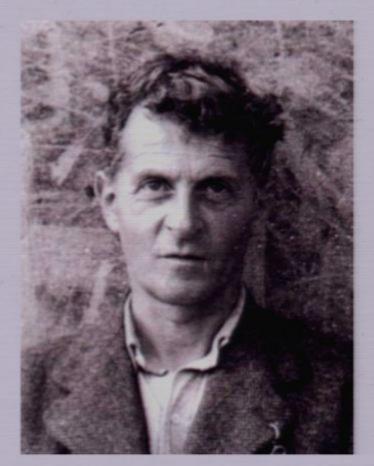
1.11 The world is determined by the facts, and by their being all the facts.

1.12 For the totality of facts determines what is the case, and also whatever is not the case.

1.13 The facts in logical space are the world.

1.2 The world divides into facts.

1.21 Each item can be the case or not the case while everything else remains the same.



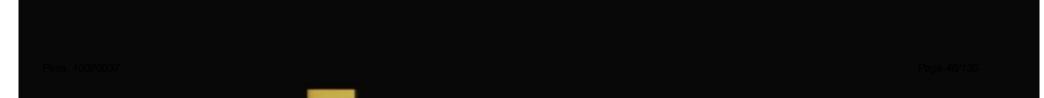
My Brief History of Space-Time

My Brief History of Space-Time

- *At the beginning there were only events ...
 - *Then the Man devised causal connections between events
 - *He modeled the causal connections in a unified framework which is space-

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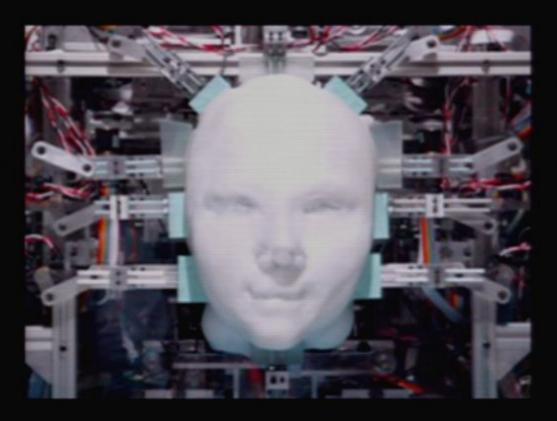
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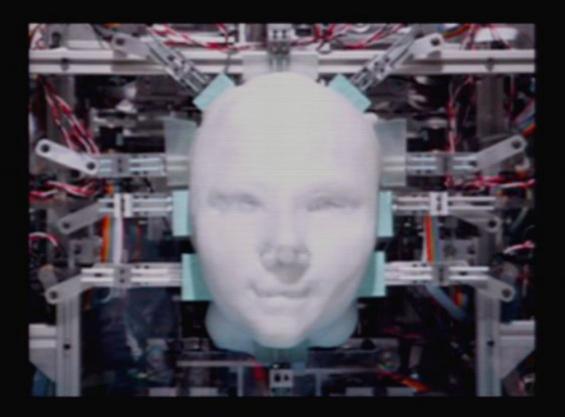
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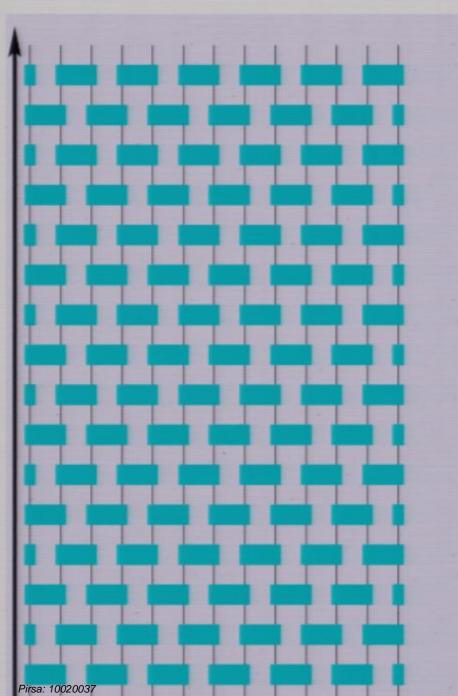


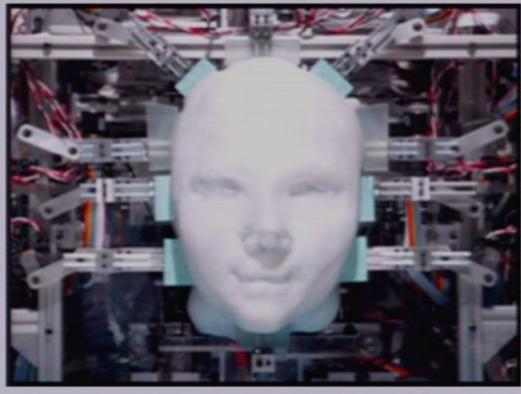


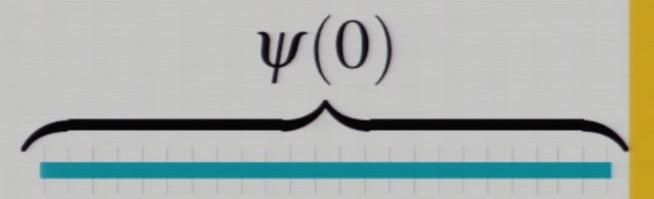


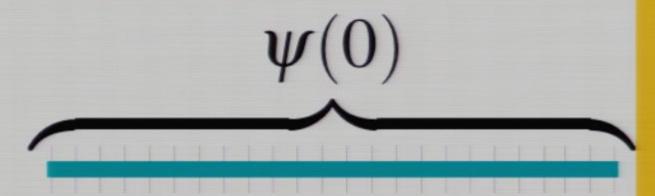


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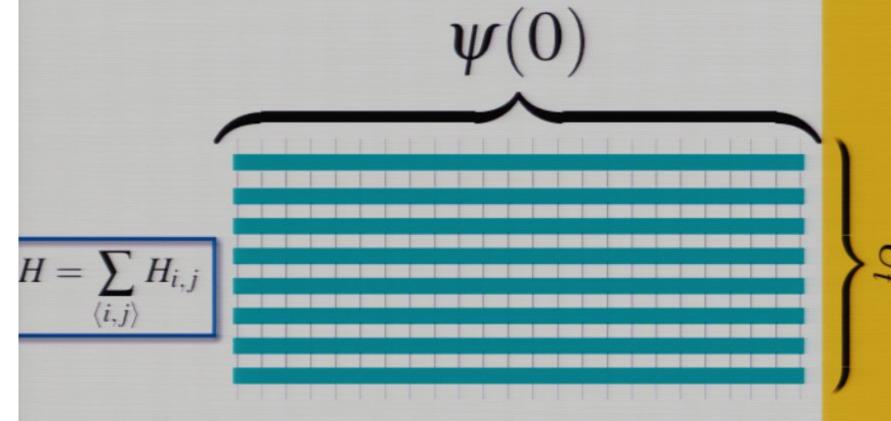


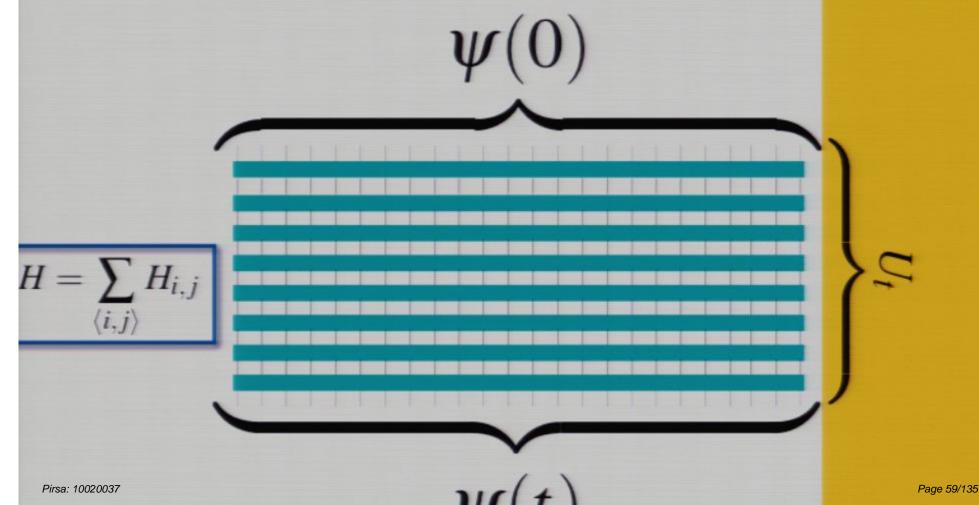






$$H = \sum_{\langle i,j \rangle} H_{i,j}$$

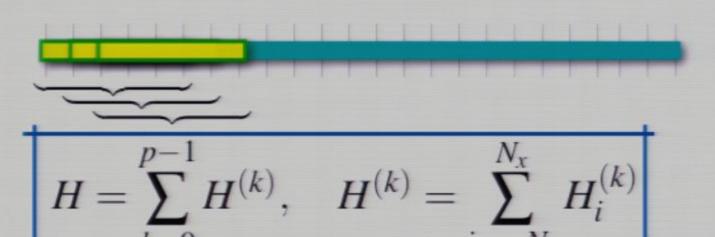




p nn translational-invariant "Hamiltonian"

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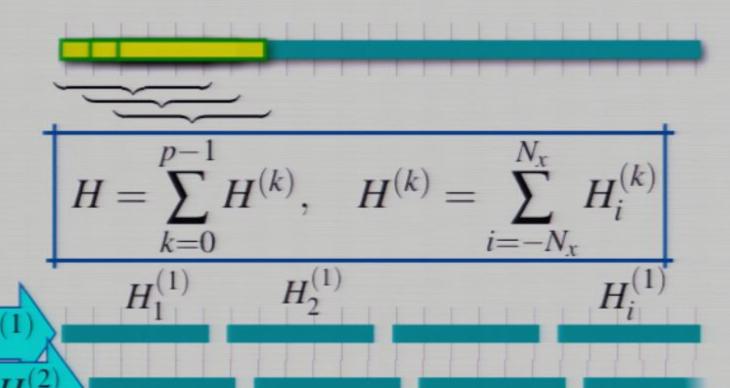
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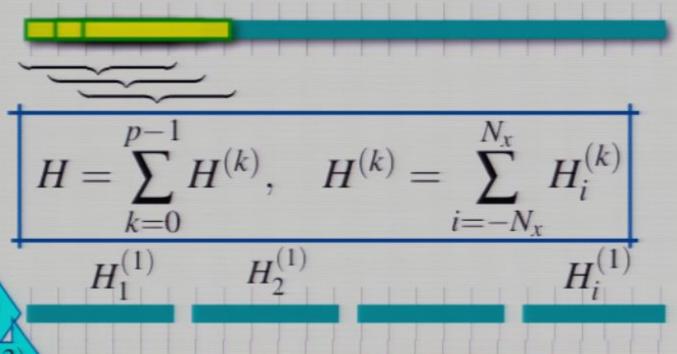




$$H = \sum_{k=0}^{p-1} H^{(k)}, \quad H^{(k)} = \sum_{i=-N_x}^{N_x} H_i^{(k)}$$







$\lim_{N\to\infty} \left(e^{-i\frac{\phi_t}{N}H^{(0)}} e^{-i\frac{\phi_t}{N}H^{(1)}} \right)^N = e^{-i\phi_t H}$

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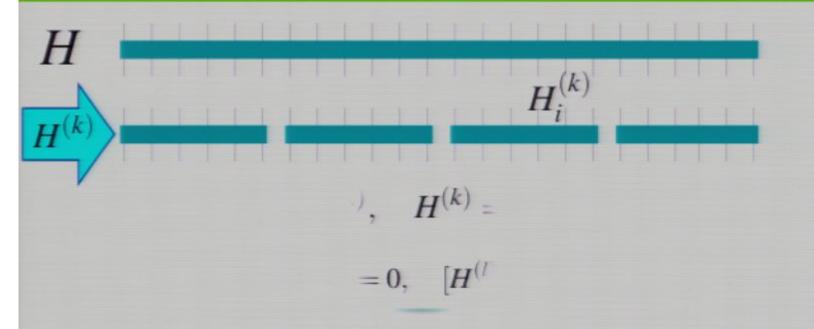
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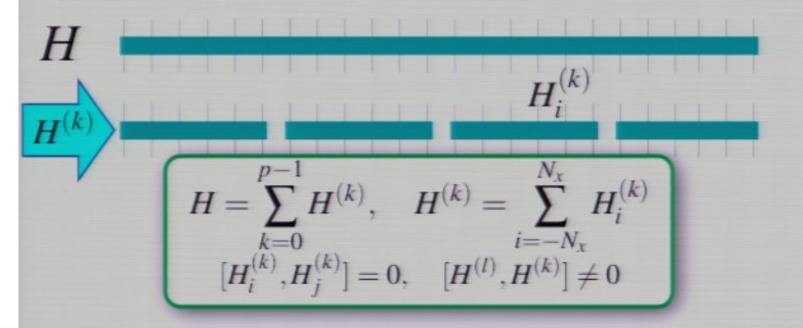
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H

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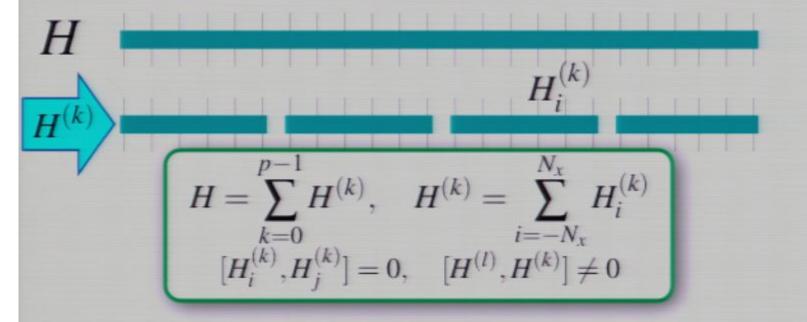




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Suzuki's bound

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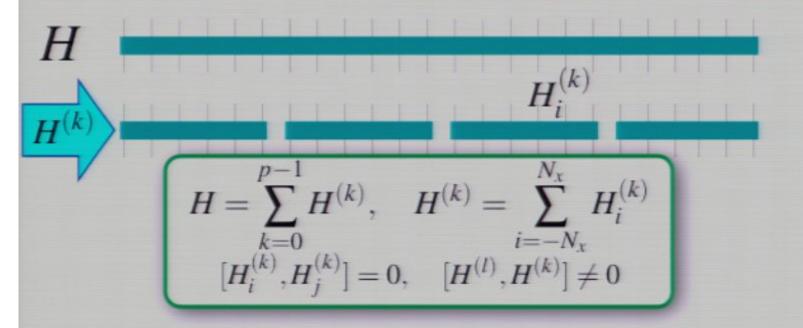


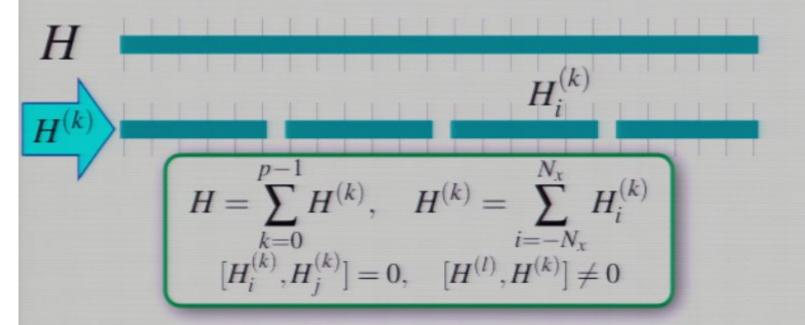
$$\exp\left(\sum_{k=0}^{p-1} H^{(k)}\right) - \left[\prod_{k=0}^{p-1} \exp\left(\frac{H^{(k)}}{n}\right)\right]^n \leqslant \frac{2}{n} \left(\sum_{i=0}^{p-1} \|H^{(k)}\|\right)^2 \exp\left(\frac{n+2}{n} \sum_{k=0}^{p-1} \|H^{(k)}\|\right)$$

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Suzuki's bound

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$$e^{-i\phi_t H} - \left(\prod_{i=1}^p e^{-i\frac{\phi_t}{N}H^{(k)}} \right)^N \bigg\| \leqslant \frac{2[p\phi_t(2N_x+1)\|H_0^{(0)}\|]^2}{N} \exp\left[\frac{N+2}{N} p\phi_t(2N_x+1)\|H_0^{(0)}\| \right]$$

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Simple scalar fields in 1 space dimension

Simple scalar fields in 1 space dimension

 infinitesimal space-granularity (minimal in principle discrimination between independent events);

 $\phi(x)$ field, operator function of space (evolving in time); we will describe it by the set of operators $\phi_n := a^{\frac{1}{2}}\phi(na)$

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Simplest equal-time microcausality:

Fermion:
$$\psi \left\{ \psi_n, \psi_m \right\} = \delta_{nm}$$
 (Dirac)

Boson:
$$\varphi \left[\varphi_n, \varphi_m \right] = \delta_{nm}$$
 (Newton-Wigner)

Simple scalar fields in 1 space dimension

Time evolution: $\phi(t) = U_t \phi(0) U_t^{\dagger}$

$$U_t = \exp\left(-\frac{i}{\hbar}t\hbar\omega H\right) = \exp(-2\pi i N_T H)$$

H: adimensional Hamiltonian

$$N_T = \frac{t}{T} = \frac{\omega t}{2\pi} = \frac{\phi_t}{2\pi}$$

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$$i\hbar\partial_t\phi=[\hbar\omega H,\phi]$$

Simple scalar fields in 1 space dimension

$$H = -\frac{i}{2} \sum_{n} (\phi_{n}^{\dagger} \phi_{n+1} - \phi_{n+1}^{\dagger} \phi_{n}) = \frac{a}{\hbar} P$$

$$\phi_n := a^{\frac{1}{2}}\phi(na)$$

$$P = -i\hbar \int \mathrm{d}x \phi^{\dagger}(x) \partial_x \phi(x)$$

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$$[H,\phi(x)] = [H,a^{-\frac{1}{2}}\phi_n] = a^{-\frac{1}{2}}\frac{i}{2}(\phi_{n+1} - \phi_{n-1}) = ia\partial_x\phi(x)$$

$$\omega a = c$$

$$\Box \phi = 0$$
 Klein-Gordon

Simple scalar fields in 1 space dimension

Trotter-ization

$$H = H^{(0)} + H^{(1)}, \quad H^{(k)} = -\frac{i}{2} \sum_{j=-N_x}^{N_x} H_j^{(k)} \quad H_n^{(k)} = -\frac{i}{2} (\phi_{2n+k}^{\dagger} \phi_{2n+k+1} - \phi_{2n+k+1}^{\dagger} \phi_{2n+k})$$

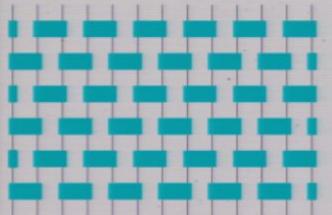
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Suzuki bound:

$$\left\| e^{-i\phi_t H} - \left(e^{-i\frac{\phi_t}{N}H^{(0)}} e^{-i\frac{\phi_t}{N}H^{(1)}} \right)^N \right\| \leq \frac{2\phi_t^2}{N} \left(\sum_{i=0}^1 \|H^{(k)}\| \right)^2 \exp\left(\phi_t \frac{N+2}{N} \sum_{k=0}^1 \|H^{(k)}\| \right)$$

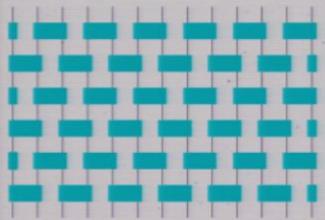


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$$\left\| e^{-i\phi_t H} - \left(e^{-i\frac{\phi_t}{N}H^{(0)}} e^{-i\frac{\phi_t}{N}H^{(1)}} \right)^N \right\| \leqslant \frac{8\|H_0^{(0)}\|^2 \phi_t^2 (2N_x + 1)^2}{N} \exp\left(2\phi_t (2N_x + 1) \frac{N+2}{N} \|H_0^{(0)}\| \right)$$

To guarantee convergence one needs a fixed!

Simple scalar fields in 1 space dimension

However, for a fixed one has maximal causal speed going to infinity!

$$\frac{a}{\tau} = \frac{x/2N_x}{t/2N}, \quad \lim_{N \to \infty} \frac{a}{\tau} = \infty$$

Simple scalar fields in 1 space dimension

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If we want to keep the maximal causal speed equal to c we need to rescale both time and space atoms in the same way, namely:

$$c = \frac{x}{t} = \frac{a2N_x}{\tau 2N} \Longrightarrow N_x = N$$

$$c = \omega a = \frac{a}{\tau} \Longrightarrow \omega = \frac{1}{\tau}$$

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but then Suzuki bound doesn't guarantee convergence of Trotter-ization since:

$$\left\| e^{-i\phi_t H} - \left(e^{-i\frac{\phi_t}{N}H^{(0)}} e^{-i\frac{\phi_t}{N}H^{(1)}} \right)^N \right\| \leqslant \frac{8\|H_0^{(0)}\|^2 \phi_t^2 (2N_x + 1)^2}{N} \exp\left(2\phi_t (2N_x + 1) \frac{N+2}{N} \|H_0^{(0)}\| \right)$$
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Since the phase for a swapping gate is π the time t must be discrete

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Question: Which kind of computation will simulate a first-quantized theory?

Answer: a classical computation! (Runge-Kutta integration)

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Question: Which kind of computation will simulate a first-quantized theory?

Answer: a classical computation! (Runge-Kutta integration)

 $\phi(x)$ will be described by a string of classical inflits $|\phi\rangle\langle\phi|$

$$\rho_{\vec{\phi}} = \otimes_n |\phi_n\rangle \langle \phi_n| \qquad \langle \phi |\phi'\rangle = \delta_{\phi\phi'}$$

$$\langle \phi | \phi' \rangle = \delta_{\phi \phi'}$$

$$\phi \in \mathbb{C}$$

A general classical processing will be described by a classical channel:

$$\mathcal{C}(\rho_{\vec{\phi}}) = \sum_{\vec{\phi}'} p(\vec{\phi}'|\vec{\phi}) \rho_{\vec{\phi}'}$$

Pisa: 10020037 ministic evolution:
$$p(ec{\phi}'|ec{\phi}) = \delta(ec{\phi}' - f(ec{\phi}))$$

1ST QUANTIZATION

classical channel:

$$\mathscr{C}(\rho_{\vec{\phi}}) = \sum_{\vec{\phi}'} p(\vec{\phi}'|\vec{\phi}) \rho_{\vec{\phi}'}$$

Deterministic evolution:
$$p(ec{\phi}'|ec{\phi}) = \delta(ec{\phi}' - f(ec{\phi}))$$

namely:

$$\mathscr{C}(\rho_{\vec{\phi}}) = \rho_{f(\vec{\phi})}$$

quantum evolution = unitary kernel:

$$\phi_i(t+\tau) = \sum_j U_{ij}\phi_j(t)$$

1ST QUANTIZATION BY QCFT₁

What is QCFT₁?

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QCFT₂ squanders the Hilbert space when simulating QFT₁!

It linearly combines the eigenvalues ϕ_n of the projectors $|\phi_n\rangle\langle\phi_n|$ without making superpositions of the kets $|\phi_n\rangle$ nor making entanglements between different systems

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QCFT₁ simulates QFT₁ more efficiently!

$$\phi = \begin{pmatrix} \dots \\ \phi_{n-1} \\ \phi_n \\ \phi_{n+1} \\ \dots \end{pmatrix}$$

 $\phi = \begin{pmatrix} \phi_{n-1} \\ \phi_n \\ \phi_{n+1} \end{pmatrix}$ Interactions between n quantum systems become many connecting n orthogonal states of a single quantum system.

1ST QUANTIZATION BY QCFT₁

Whatis $QCFT_1$?

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ight)$$

 $\phi = \begin{pmatrix} \phi_{n-1} \\ \phi_n \\ \phi_{n+1} \end{pmatrix}$ Interactions between n quantum systems become matrix blocks connecting n orthogonal states of a single quantum system

Time evolution:

 $\phi(t) = U_t \phi(0) := \exp(-i\omega t H^{age 98/135})$

Dirac Particle

$$i\hbar\partial_t\psi = \begin{pmatrix} ic\hbar\sigma_x\partial_x & mc^2\\ mc^2 & -ic\hbar\sigma_x\partial_x \end{pmatrix}\psi$$

$$\psi := \begin{pmatrix} \cdots \\ \psi_n \\ \psi_{n+1} \\ \cdots \end{pmatrix}, \qquad \psi_n = \begin{pmatrix} u_n \\ v_n \end{pmatrix} = \begin{pmatrix} u_n^1 \\ u_n^2 \\ v_n^1 \\ v_n^2 \\ v_n^2 \end{pmatrix}$$

$$H = \begin{pmatrix} \dots & 0 & 0 & -\frac{i}{2}\sigma_{x} & 0 & 0 & 0 & \dots \\ \dots & 0 & 0 & \frac{a}{\lambda}I & \frac{i}{2}\sigma_{x} & 0 & 0 & \dots \\ \dots & \frac{i}{2}\sigma_{x} & \frac{a}{\lambda}I & 0 & 0 & -\frac{i}{2}\sigma_{x} & 0 & \dots \\ \dots & 0 & -\frac{i}{2}\sigma_{x} & 0 & 0 & \frac{a}{\lambda}I & \frac{i}{2}\sigma_{x} & \dots \\ \dots & 0 & 0 & \frac{i}{2}\sigma_{x} & \frac{a}{\lambda}I & 0 & 0 & \dots \\ \dots & 0 & 0 & 0 & -\frac{i}{2}\sigma_{x} & \frac{a}{\lambda}I & 0 & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

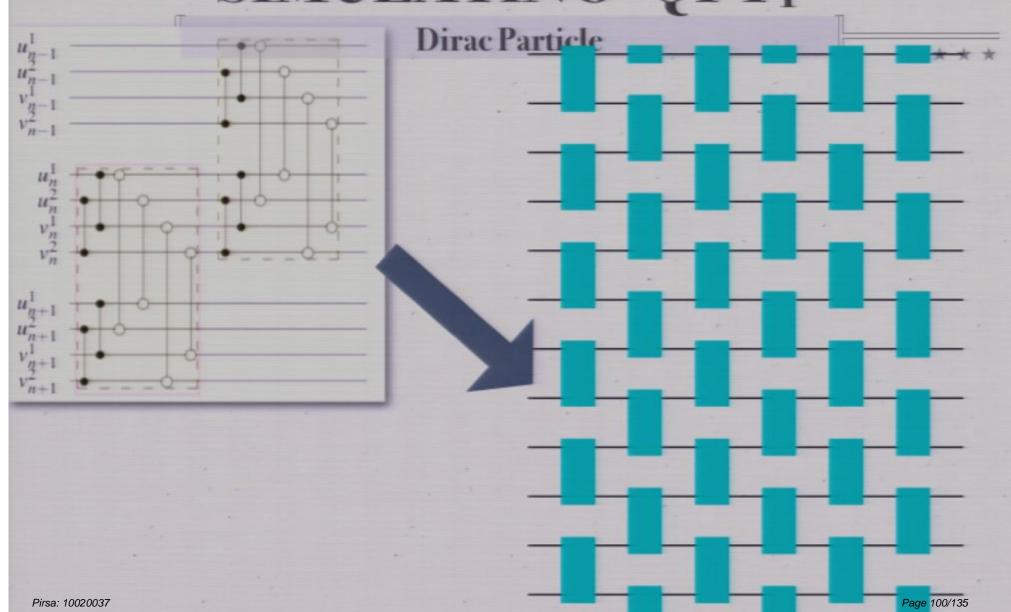
$$u_{n-1}^{2}$$
 v_{n-1}^{1}
 u_{n}^{1}
 u_{n}^{1}
 v_{n}^{2}
 v_{n}^{1}
 v_{n}^{2}
 v_{n}^{2}
 v_{n+1}^{1}
 v_{n+1}^{2}
 v_{n+1}^{2}

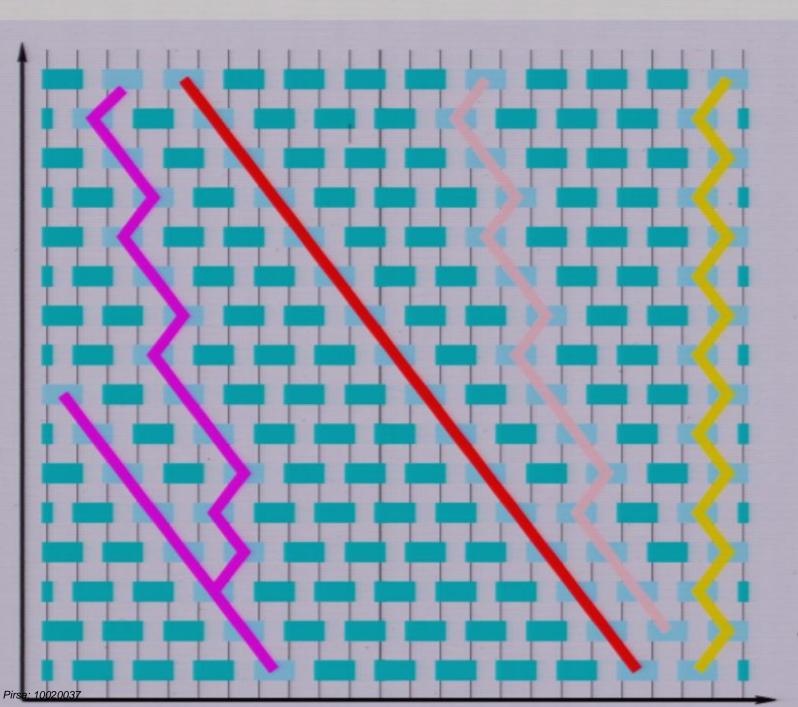
$$\lambda := \frac{\hbar}{mc} = 3.86159 * 10^{-13}$$

Reduced Compton wavelength

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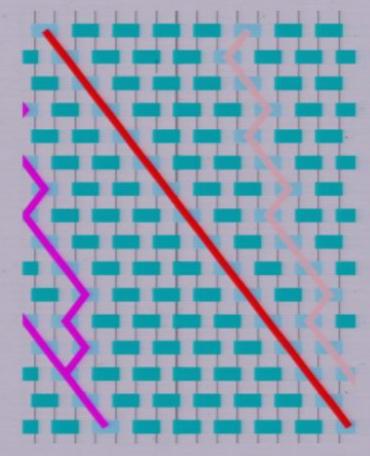
makes sense above the scale of





Mass: fequency of the zigzag

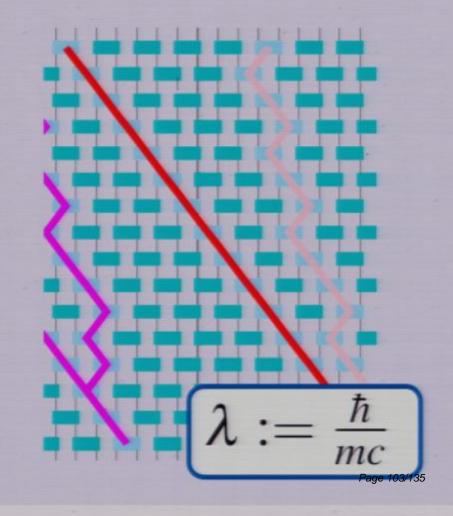
What is ħ?



What is ħ?

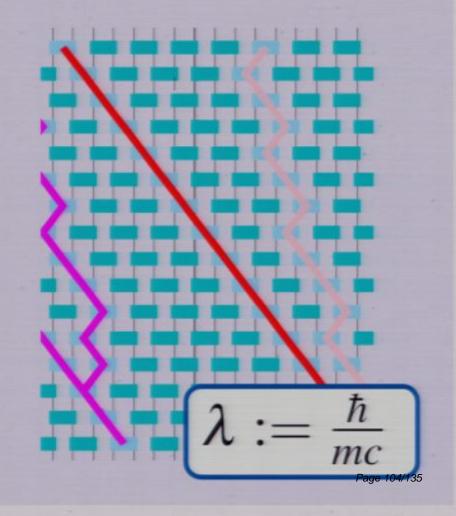
**m* is the ZBW frequency

*c is the max causal speed



What is ħ?

- **m* is the ZBW frequency
- *c is the max causal speed
- * ħ is the scale above which one can use ∂, i.e. at which one sees the gates!
- *quantization rules
 "emergent" (P is a
 swappiness, X is where
 the qubit is 1)



MULATING QFT₁

Schrödinger equation

$$\partial_t \phi = i \frac{\hbar}{2m} \partial_x^2 \phi$$

$$\omega = \frac{\hbar}{2ma^2}$$

$$H = \sum_{j} e_{j+1,j} - 2e_{j,j} + e_{j,j+1} =$$

$$H = H^{(0)} + H^{(1)}, \quad H^{(0)} = \sum_{j} H_{2j,2j+1}, \quad H^{(1)} = \sum_{j} H_{2j+1,2j+2}$$
 Pirsa: 1002003 $H_{j,j+1} = \sum_{j} e_{j+1,j} - e_{j,j} - e_{j+1,j+1} + e_{j,j+1}$

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Schrödinger equation

"Trotterize" the Hamiltonian

$$H = H^{(0)} + H^{(1)}, \quad H^{(0)} = \sum_{j} H_{2j,2j+1}, \quad H^{(1)} = \sum_{j} H_{2j+1,2j+2}$$

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Schrödinger equation

"Trotterize" the Hamiltonian

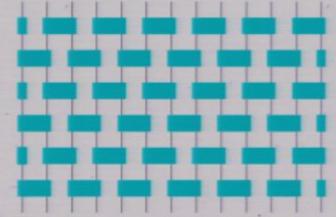
$$H = H^{(0)} + H^{(1)}, \quad H^{(0)} = \sum_{j} H_{2j,2j+1}, \quad H^{(1)} = \sum_{j} H_{2j+1,2j+2}$$

$$H_{j,j+1} = \sum_{j} e_{j+1,j} - e_{j,j} - e_{j+1,j+1} + e_{j,j+1}$$

By taking the maximal causal speed equal to C namely $a \propto N^{-1}$ one obtains:

$$\omega = \frac{\hbar}{2ma^2} \propto N^2$$

a contradiction!



Dirac Field Theory

Dirac Field Theory

Clifford algebra
$$\Gamma_k := \left(\prod_{j=-\infty}^{k-1} \sigma_j^z\right) \sigma_k^ \{\Gamma_k, \Gamma_h\} = \delta_{kh}$$

Fermi scalar field:
$$\psi_n = \Gamma_{4n+\alpha}$$
 $\{\psi(x), \psi^{\dagger}(x)\} = \delta(x-y)$ $\{\psi(x), \psi(y)\} = 0$

$$-\frac{i}{2}(\sigma_{l}^{+}\sigma_{l+1}^{-}-\sigma_{l}^{-}\sigma_{l+1}^{+})$$

Dirac Field Theory

Clifford algebra
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$$-\frac{i}{2}(\sigma_{l}^{+}\sigma_{l+1}^{-}-\sigma_{l}^{-}\sigma_{l+1}^{+})$$

Dirac field:

$$\{\psi_{\alpha}(x), \psi_{\beta}^{\dagger}(y)\} = \delta_{\alpha\beta}\delta(x - y) \{\psi_{\alpha}(x), \psi_{\beta}(y)\} = 0 \{u_{1}^{1} \\ v_{1}^{1} \\ v_{n-1}^{1} \\ v_{n-1}^{1} \\ v_{n-1}^{1} \\ v_{n}^{1} \\$$

$$H = \begin{pmatrix} \frac{i}{2}\sigma_{x}(\delta_{+} - \delta_{-}) & \frac{a}{\lambda}I \\ \frac{a}{\lambda}I & -\frac{i}{2}\sigma_{x}(\delta_{+} - \delta_{-}) \end{pmatrix}$$

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$$\mathcal{H} = -i\hbar c \psi^{\dagger}(x) \partial_x \psi(x)$$

Dirac Field Theory

Clifford algebra
$$\Gamma_k := \left(\prod_{j=-\infty}^{k-1} \sigma_j^z\right) \sigma_k^ \{\Gamma_k, \Gamma_h\} = \delta_{kh}$$

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$$-\frac{i}{2}(\sigma_l^+\sigma_{l+1}^- - \sigma_l^-\sigma_{l+1}^+)$$

 $\mathcal{H} = -i\hbar c \psi^{\dagger}(x) \partial_x \psi(x)$

Dirac field:

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$$\{\psi_{\alpha}(x), \psi_{\beta}(y)\} = 0$$

$$\downarrow_{\frac{u_{1}^{1}-1}{v_{n-1}^{2}}}$$

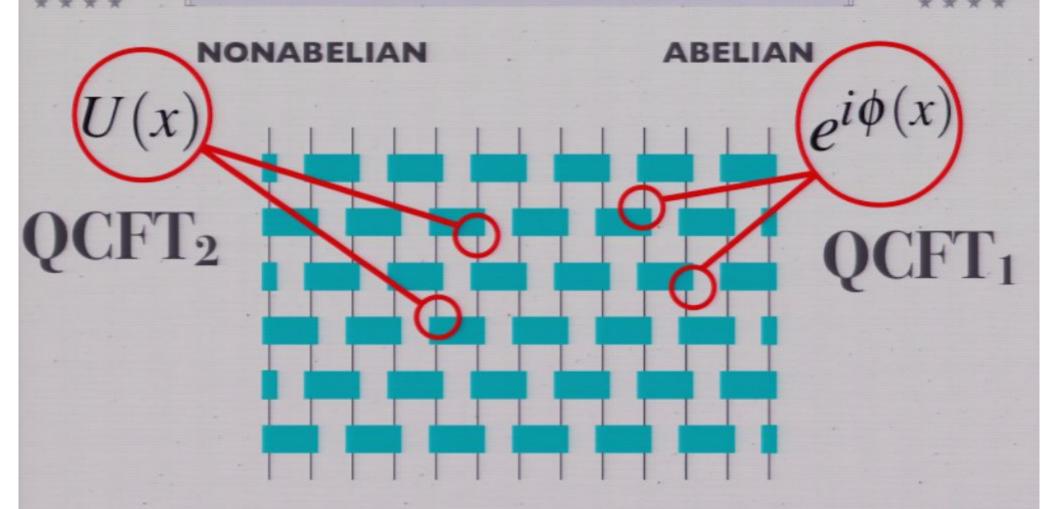
$$\downarrow_{\frac{u_{n}^{1}}{v_{n}^{2}}}$$

$$H = \begin{pmatrix} \frac{i}{2}\sigma_{x}(\delta_{+} - \delta_{-}) & \frac{a}{\lambda}I \\ \frac{a}{\lambda}I & -\frac{i}{2}\sigma_{x}(\delta_{+} - \delta_{-}) \end{pmatrix}$$

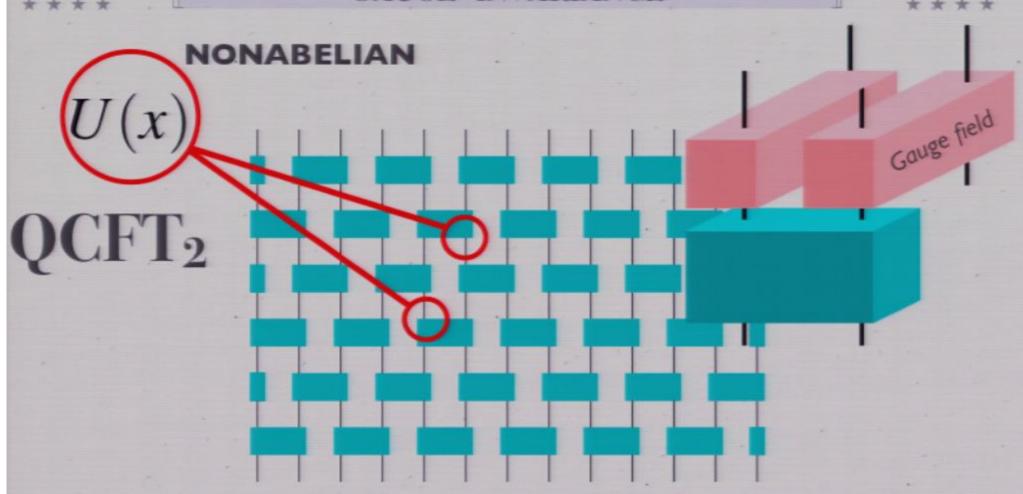
$$\sigma_l^+ \sigma_{l+1}^z \dots \sigma_{l+k-1}^z \sigma_{l+k}^- + h.c.$$

OFT cure: Grassmann variables!

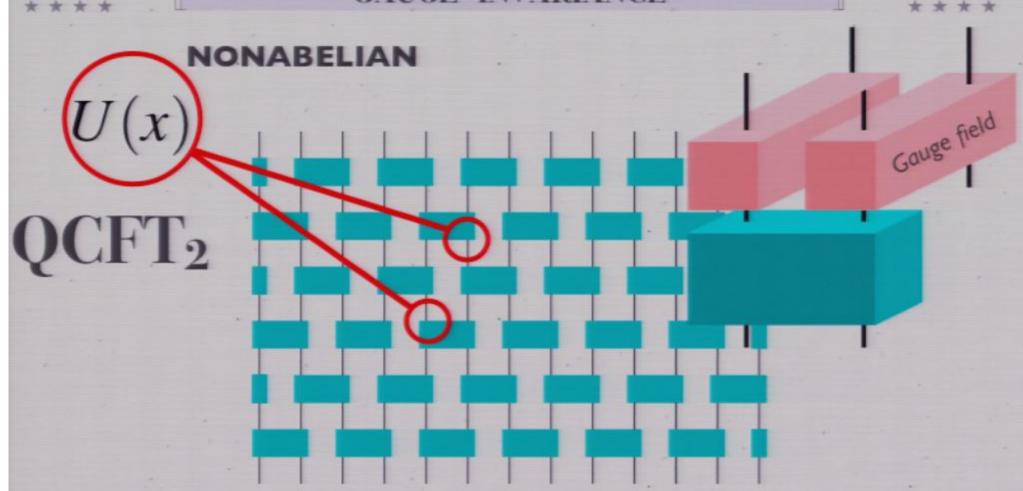
GAUGE INVARIANCE



GAUGE INVARIANCE



GAUGE INVARIANCE



Natively nonabelian Gauge theory!

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and on ... foliation !!!

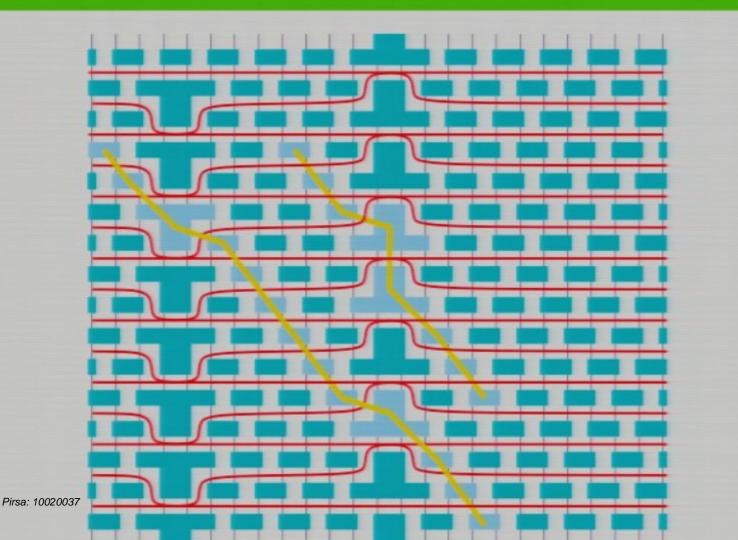


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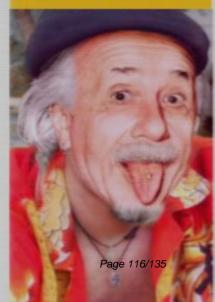
PLAY GOD WITH QCFT

or else: Einstein demystified

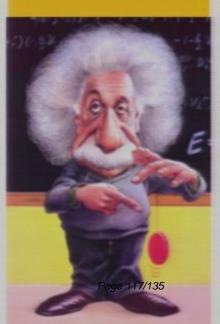




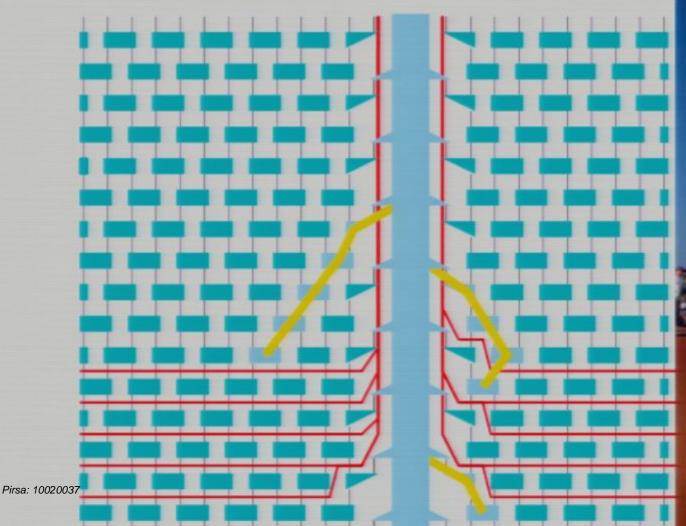
positive and negative masses

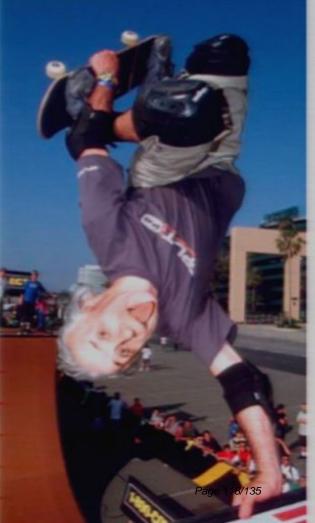


a worm hole!

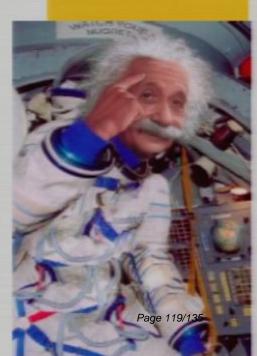


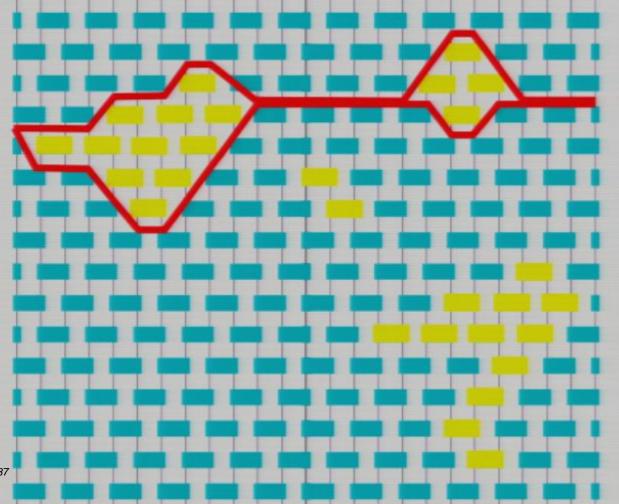
a black hole!



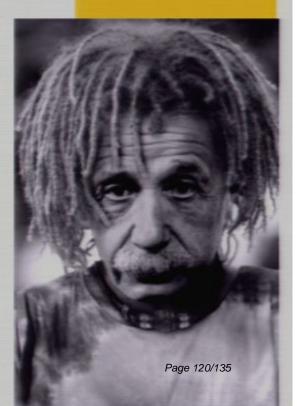


a time tunnel!





patterns?

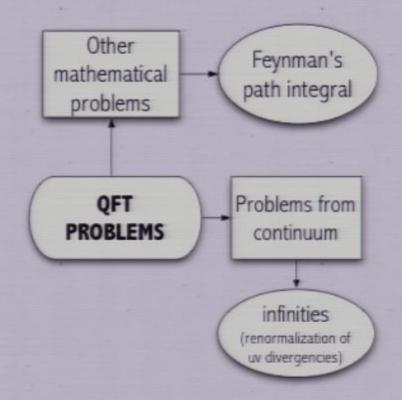


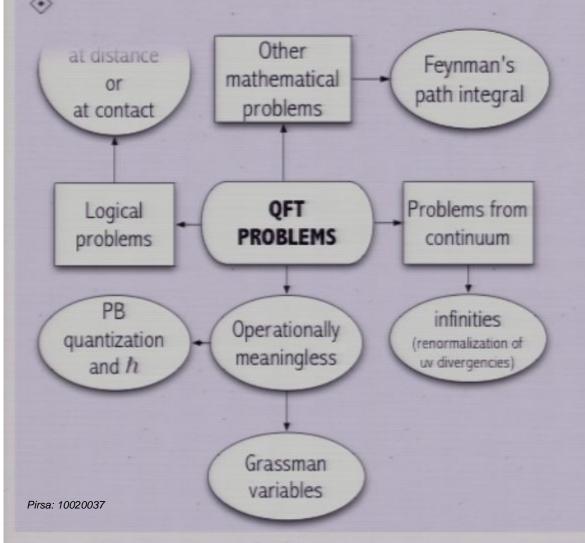
Einstein operationalist?

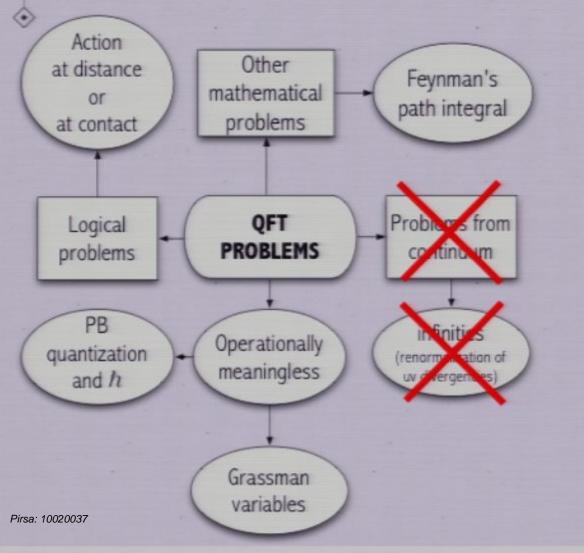
Einstein's disagreement with the operationalist approach was criticized by Bridgman[‡] as follows: "Einstein did not carry over into his general relativity theory the lessons and insights he himself has taught us in his special theory."

Bridgman, Einstein's Theories and the Operational Point of View, in: P.A. Schilpp, ed., Albert Einstein: Philosopher-Scientist, Open Pirsa: 10020037 COURT, La Salle, III., CUP, 1982, Vol. 2, p. 335-354.

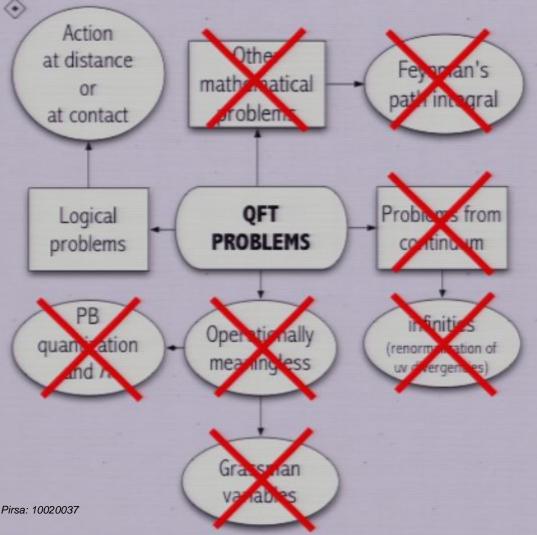


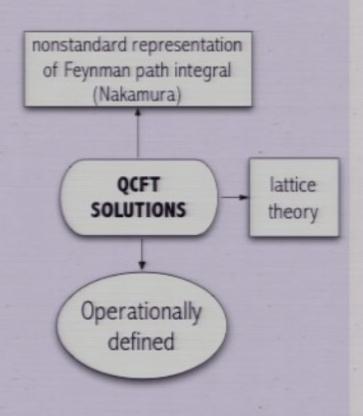


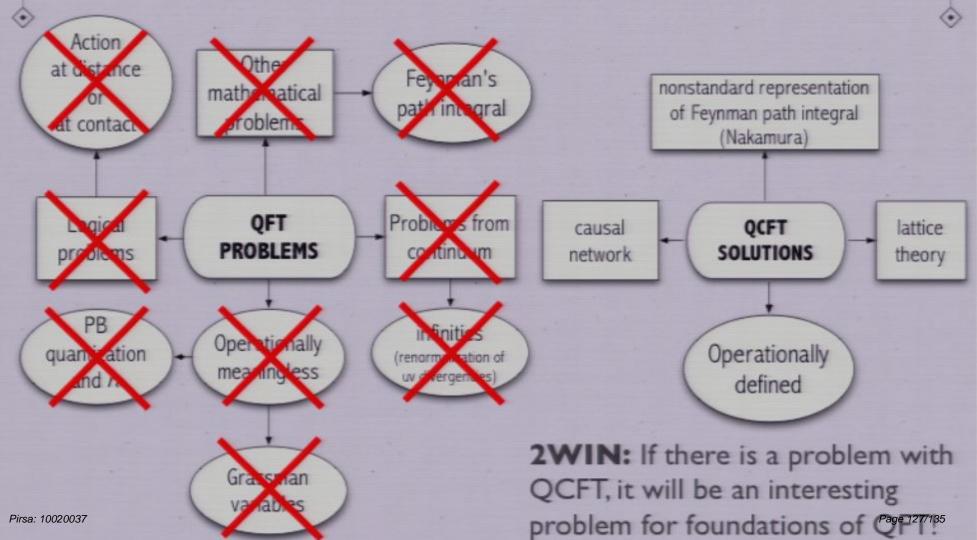




QCFT SOLUTIONS lattice theory

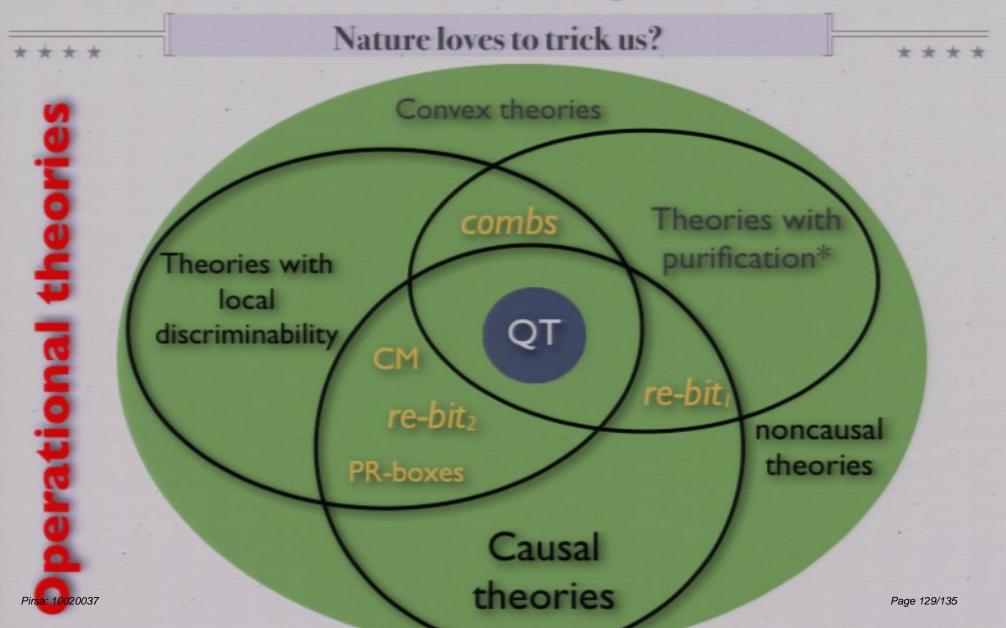






Moreover, you can change the computational engine from QT to super-QT, or even non-causal OpT, without changing the theoretical framework

THE PRINCIPLE OF THE QUANTUMNESS



"Emergent" Physics

- *Relativity
- ***Gravity**
- *Field Theory
- *Quantization rules and ħ

*...

TODO list

- * Improve Suzuki bound
- * Derive Lorentz covariance of field
- * Rederive quantization rules
- * Microcausality:
 - * Fermi (Grassman var.)
 - * Bose
 - * para-statistics?
- * Derive a 1dim toy nonabelian gauge theory

- * Derive Feynman path integral via Trotter
- * Connect Lagrangian density with a circuit tile
- Explore connections with lattice theories
- * Rederive GR Einstein's equation
- * Explore Penrose spin-networks, Regge calculus, etc.
- * Rederive gauge theories
- * Write a new Theory of ... Quantum Gravity!!!

Concludingremarks

Concluding remarks

- *QCFT seems to have many advantages versus QFT
- * It puts the finger on the foundational problems in QFT
- *It is QG-ready
- * It offers a first test of Lucien Hardy approach to QG
- * It's fun

TODO list

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No Signal

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