

Title: Gravitational Physics - Review (PHYS 636) - Lecture 2

Date: Jan 04, 2010 11:00 AM

URL: <http://pirsa.org/10010049>

Abstract:



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Lie Derivative

Uses a vector field to take derivative of other vectors (tensors).

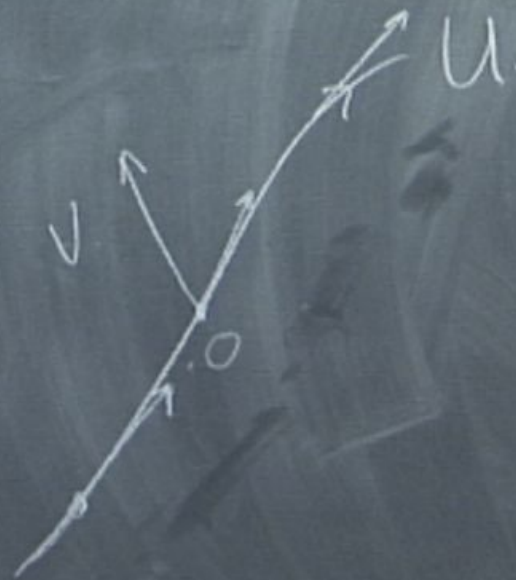
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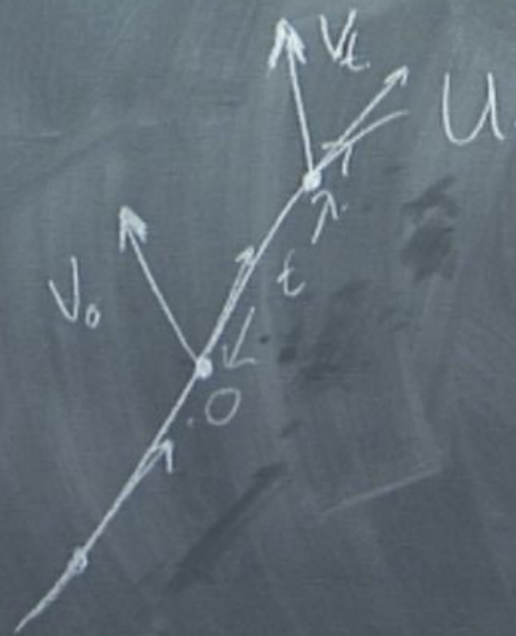
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Lie Derivative

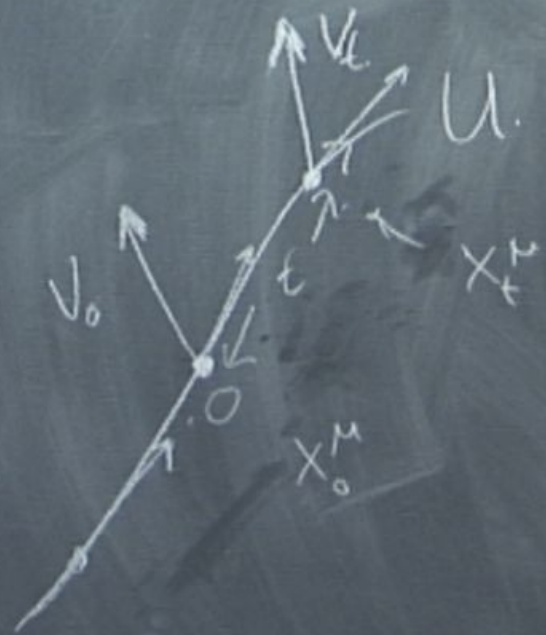
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Lie Derivative

Uses a vector field to take derivative of other vectors (tensors).

$$X_t^M = X_0^M + tU^M + O(t^2)$$



$$\left(L_u V \right)^n = \lim_{t \rightarrow 0} \frac{1}{t} \left[V^H (X_0^\alpha + tU^\alpha) - \frac{\partial X^H V_0}{\partial X_0^\alpha} \right]$$



$$\begin{aligned}
 \left(L_u V \right)^M &= \lim_{t \rightarrow 0} \frac{1}{t} \left[V^M (X_0^\alpha + t U^\alpha) - \frac{\partial X^M}{\partial X_0^\nu} V_0^\nu \right] \\
 &= \lim_{t \rightarrow 0} \frac{1}{t} \left[V_0^M + t V_{,\nu}^M U^\nu \right]
 \end{aligned}$$



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 \left(L_u V \right)^n &= \lim_{t \rightarrow 0} \frac{1}{t} \left[V^\mu (X_0^\alpha + t U^\alpha) - \frac{\partial X^\mu V_0^\nu}{\partial X_0^\nu} \right] \\
 &= \lim_{t \rightarrow 0} \frac{1}{t} \left[V_0^\mu + t V_{,\nu}^\mu U^\nu \right. \\
 &\quad \left. - (\delta_\nu^\mu + t U_{,\nu}^\mu) V_0^\nu \right]
 \end{aligned}$$

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 \left(\mathcal{L}_u V \right)^{\mu} &= \lim_{t \rightarrow 0} \frac{1}{t} \left[V^{\mu} (X_0^{\alpha} + t u^{\alpha}) - \frac{\partial X^{\mu}}{\partial X_0^{\nu}} V_0^{\nu} \right] \\
 &= \lim_{t \rightarrow 0} \frac{1}{t} \left[\cancel{V_0^{\mu}} + t V_{,\nu}^{\mu} u^{\nu} - \left(\cancel{\delta_{\nu}^{\mu}} + t u_{,\nu}^{\mu} \right) V_0^{\nu} \right] \\
 &= U^{\nu} V_{,\nu}^{\mu} - U_{,\nu}^{\mu} V^{\nu}
 \end{aligned}$$

For co-vector: $(L_V \omega)_\alpha = \omega_{\alpha, \beta} V^\beta + \omega_\beta V^\beta{}_{,\alpha}$

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+ natural extension for tensors

The Lie deriv is a map $T_p(M) \times T_p(M) \rightarrow T_p(M)$

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$$[v, w]f = v(wf) - w(vf)$$

$$\omega)_\alpha = \omega_{\alpha\beta} V^\beta + \omega_\beta \underline{\underline{V^\beta}}_{,\alpha}$$

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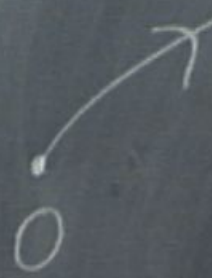
$$(\omega_\alpha u^\alpha)$$

Geometrical Significance



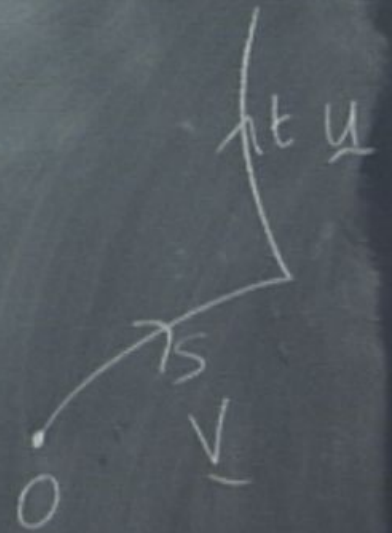
Geometrical Significance

The bracket tells us
how far a cunk is from "closing"



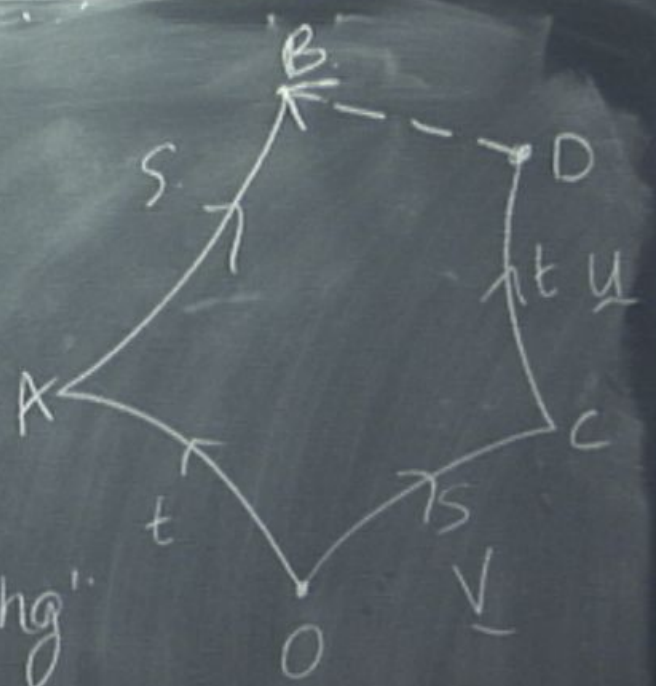
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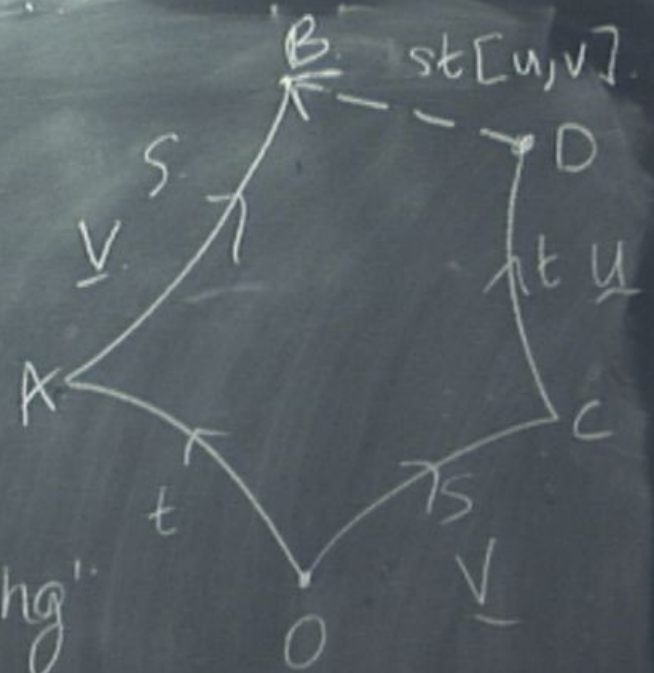
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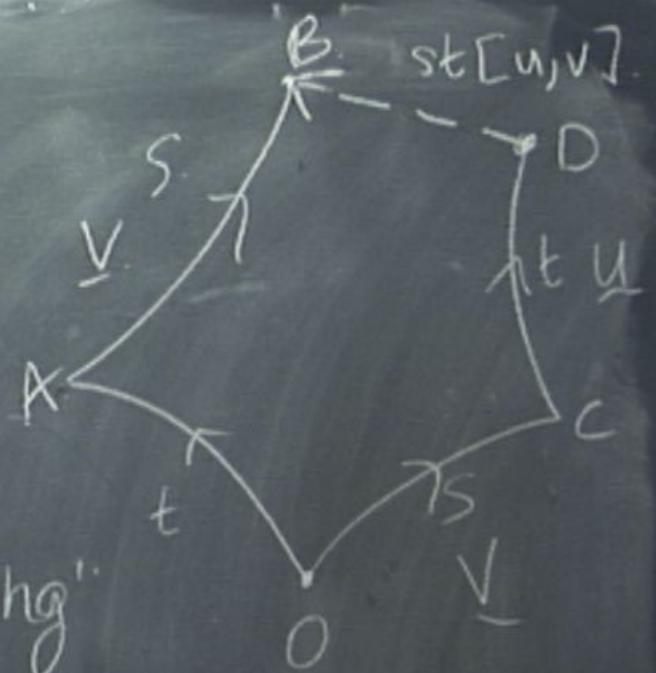
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$$X_B^\alpha = X_A^\alpha + S V_A^\alpha + \frac{1}{2} S^2 V^\alpha_{,\beta} V^\beta|_A$$

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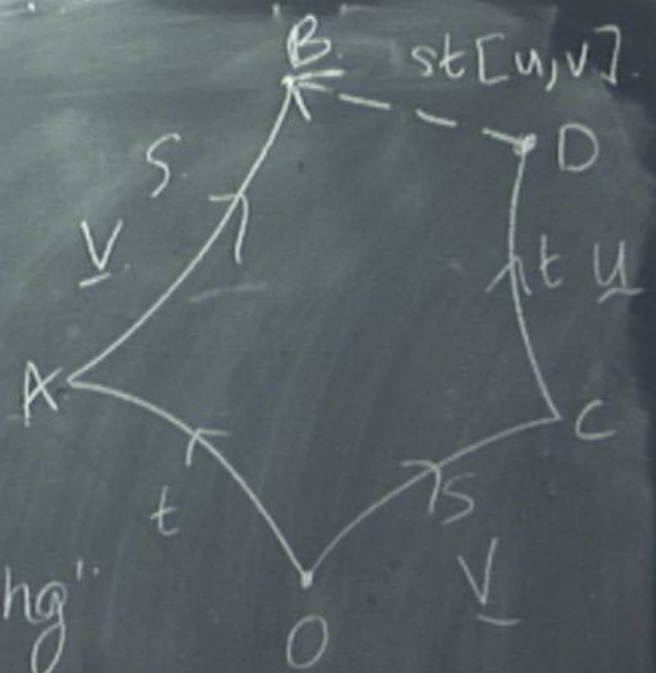
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$$\begin{aligned}
 X_B^\alpha &= X_A^\alpha + S V_A^\alpha + \frac{1}{2} S^2 V_{,\beta}^\alpha V^\beta|_A \\
 &= X_O^\alpha + t U_O^\alpha + \frac{1}{2} t^2 U_{,\beta}^\alpha U^\beta|_O
 \end{aligned}$$

Geometrical Significance

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$$\begin{aligned}
 X_B^\alpha &= X_A^\alpha + S V_A^\alpha + \frac{1}{2} S^2 V^\alpha_{,\beta} V^\beta|_A \\
 &= X_O^\alpha + t U^\alpha + \frac{1}{2} t^2 U^\alpha_{,\beta} U^\beta|_O + S \left(V_0^\alpha + t V^\alpha_{,\beta} U^\beta \right) \\
 &\quad + \frac{1}{2} S^2 V^\alpha_{,\beta} V^\beta|_O
 \end{aligned}$$

To get X_D^α , swap $u \leftrightarrow v$, $s \leftrightarrow t$



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Then

$$\begin{aligned} X_B^\alpha - X_D^\alpha &= st [V^\alpha, \beta U^\beta - U^\alpha, \beta V^\beta] \\ &= st [u, v]^\alpha \end{aligned}$$

Lie derivatives are also used for symmetries.

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A Killing vector is a vector along which the Lie deriv of the metric vanishes.

$$(\mathcal{L}_V g)_{\mu\nu} = 0.$$

From
Grains of
Pollen to
Evidence
for Atoms

How
Big Is A
Molecule?

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$$(L_V g)_{\mu\nu} = 0 \quad \parallel$$
$$\nabla_\mu V_\nu + \nabla_\nu V_\mu = 0$$

Forms and Exterior Derivative

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d: Acting on a f_n , the ext. deriv.
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\underline{d} : Acting on a fn, the ext. deriv.
maps the function onto a covector,

or 1-form $\underline{d}f$.

$$\langle \underline{d}f | X \rangle = Xf$$

all $X \in T_p(M)$

For a co-ord basis:

$$\langle df | \left(\frac{\partial}{\partial x^m} \right) \rangle = \frac{\partial f}{\partial x^m}$$

iv.

$T_p(M)$

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A p-form is an antisymmetric tensor
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$T_p(M)$

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$$\langle df | \left(\frac{\partial}{\partial x^m} \right) \rangle = \frac{\partial f}{\partial x^m}$$

A p-form is an antisymmetric tensor of rank $(0, p)$ (p indices "downstairs"), write the space of p -forms as $\bigwedge_p T^*(M)$

Construct by taking antisymmetric
product " \wedge "

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product " \wedge "

$$A \wedge B = A \otimes B - B \otimes A$$

$$(A \wedge B)_{\mu\nu} = A_{\mu} B_{\nu} - A_{\nu} B_{\mu}$$

For a p -form and a q -form.

$$\left(A^{(p)} \wedge B^{(q)} \right)_{a_1 \dots a_{p+q}} = \frac{(p+q)!}{p! \cdot q!} A_{[a_1 \dots a_p} B_{a_{p+1} \dots a_{p+q}]}$$

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$$Q_{\{\mu\nu\}} = \frac{1}{2} (Q_{\mu\nu} - Q_{\nu\mu})$$

$$Q_{(\mu\nu)} = \frac{1}{2} (Q_{\mu\nu} + Q_{\nu\mu})$$

notation is taking all permutations of indices, weighted by sign of permutation.

" \wedge " is linear but not commutative.

$$A^{(p)} \wedge B^{(q)} = (-1)^{pq} B \wedge A$$

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d acts on forms, and takes a p -form to a $(p+1)$ -form.

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$$A_{a_1, \dots, a_p} \rightarrow (dA)_{a_1, \dots, a_{p+1}} = \frac{(p+1)!}{p!} \partial_{[a_1} A_{a_2, \dots, a_{p+1}]}$$

d does not need a metric or connection

$$d(A \wedge B) = (dA) \wedge B + (-)^p A \wedge dB$$

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$$\underline{d}(A \wedge B) = (\underline{d}A) \wedge B + (-)^p A \wedge \underline{d}B$$

relation between d and \mathcal{L} :

\underline{d} does not need a metric or connection

$$\underline{d}(A \wedge B) = (\underline{d}A) \wedge B + (-)^p A \wedge \underline{d}B$$

Relation between \underline{d} and $\langle _ | _ \rangle$:

$$\langle \underline{d}\omega | \underline{u}, \underline{v} \rangle = \underline{u} \langle \underline{\omega} | \underline{v} \rangle - \underline{v} \langle \underline{\omega} | \underline{u} \rangle - \langle \underline{\omega} | \underline{\epsilon}_{\underline{u}, \underline{v}} \rangle$$

$(\phi, \underline{A}) \sim A_\mu \Leftrightarrow$ e-m gauge pot.

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

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$$\left(\begin{array}{c|ccc} 0 & E_1 & E_2 & E_3 \\ \hline & 0 & B_3 & -B_2 \\ & & 0 & B_1 \\ & & & 0 \end{array} \right)$$



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