

Title: Gravitational Physics - Review (PHYS 636) - Lecture 9

Date: Jan 13, 2010 10:00 AM

URL: <http://pirsa.org/10010039>

Abstract:

$$Z \sim \text{tr} e^{-H/T}$$

$$Z \sim \text{tr} e^{-H/T}$$

$$t \rightarrow i\tau$$

$$H = \frac{1}{\beta} \int d\tau \int d^3x \mathcal{H}$$

$$Z \sim \text{tr} e^{-H/T}$$

$$t \rightarrow i\tau$$

$$H = \frac{1}{\beta} \int d\tau \int d^3x \mathcal{H} = \frac{1}{\beta} \int d^4x \mathcal{H}$$

$$Z \sim \text{tr} e^{-\frac{H}{T}} \quad I_E$$

$$t \rightarrow i\tau$$

$$\begin{aligned} H &= \frac{1}{\beta} \int d\tau \int d^3x \mathcal{H} = \frac{1}{\beta} \int d^4x \mathcal{H} \\ &= \frac{1}{\beta} I_E \end{aligned}$$

$$Z \sim \text{tr} e^{-\frac{H}{T}} \quad I_E$$

$$t \rightarrow i\tau$$

$$H = \frac{1}{\beta} \int d\tau \int d^3x \mathcal{H} = \frac{1}{\beta} \int d^4x \mathcal{H} \\ = \frac{1}{\beta} I_E$$

$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

cat.  
Schwarzschild

$$ds^2 = \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$



$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

at.  
Schwarzschild.

$$ds^2 = \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

cat.  
Schwarzschild.

$$ds^2 = \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

cat.  
Schwarzschild.

$$ds^2 = \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$r = 2GM$  is a co-ord singularity in Lorentz metric - the event horizon. Here it is a boundary of the Euclidean space.

$$I_E \stackrel{?}{=} \frac{1}{16\pi G} \int d^4x R(g) \sqrt{g}$$

deat.  
Schwarzschild.

$$ds^2 = \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$r = 2GM$  is a co-ord singularity in Lorentz metric - the event horizon. Here it is a boundary of the Euclidean space.

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} dt^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega^2$$

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} dt^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega^2$$

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} dt^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega^2$$

looking at  $dr^2$  suggests rewriting

$$\rho^* = 2\sqrt{2GM} \sqrt{r-2GM}$$

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} dt^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega_{II}^2$$

Looking at  $dr^2$  suggests rewriting

$$= 2\sqrt{2GM} \sqrt{r-2GM}$$

$$ds^2 = + dp^2 + (2GM)^2 d\Omega_{II}^2$$



$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} dt^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega_{II}^2$$

look at  $dr^2$  suggests rewriting

$$p^* = 2\sqrt{2GM} \sqrt{r-2GM}$$

$$+ dp^2 + (2GM)^2 d\Omega_{II}^2$$

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} d\tau^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega_{II}^2$$

looking at  $dr^2$  suggests rewriting

$$p^* = 2\sqrt{2GM} \sqrt{r-2GM}$$

$$ds^2 = \frac{p^{*2}}{(4GM)^2} d\tau^2 + dp^{*2} + (2GM)^2 d\Omega_{II}^2$$

$$r \rightarrow 2GM$$

$$ds^2 \sim \frac{(r-2GM)}{2GM} d\tau^2 + \frac{2GM}{r-2GM} dr^2 + (2GM)^2 d\Omega_{II}^2$$

looking at  $dr^2$  suggests rewriting

$$p^* = 2\sqrt{2GM} \sqrt{r-2GM}$$

$$ds^2 = \frac{p^{*2}}{(4GM)^2} d\tau^2 + dp^{*2} + (2GM)^2 d\Omega_{II}^2$$

Compare to  $dp^2 + p^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

Compare to  $dp^2 + p^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

↳ we make  $\tau$  periodic with period  $8\pi G M$ .

Compare to  $dp^2 + p^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

we make  $\tau$  periodic with period  $8\pi G M$ .

Euclidean-Sch is nonsingular at ZGM.

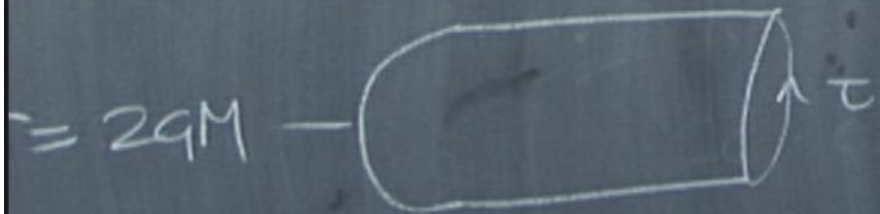
Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi GM$ .  
then Euclidean-Sch is nonsingular at ZGM.

Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

we make  $\tau$  periodic with period  $8\pi G M$ .

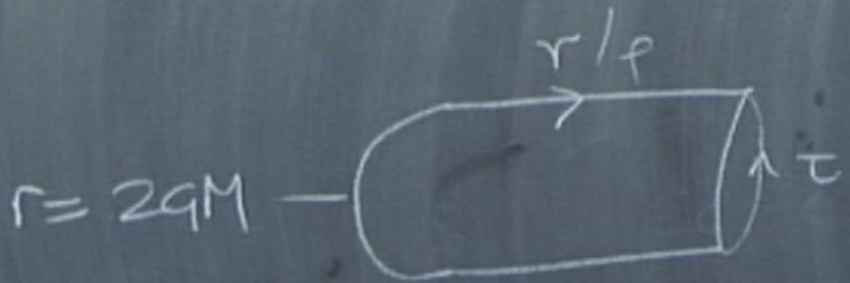
Euclidean-Sch is nonsingular at  $2GM$ .





Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi G M$ ,  
then Euclidean-Sch is nonsingular at  $2GM$ .

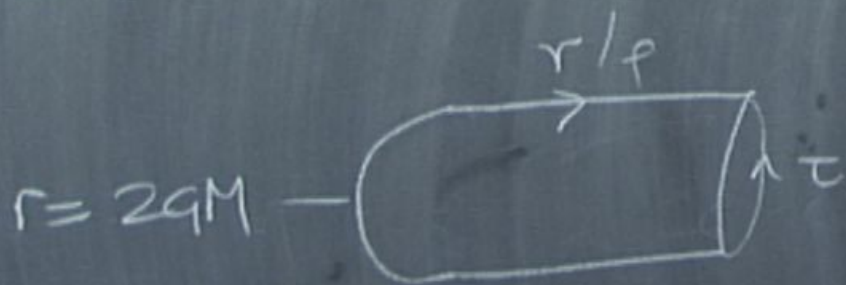


Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi G M$ .

then Euclidean - Sch is nonsingular at  $r = 2GM$ .

"black hole cigar"

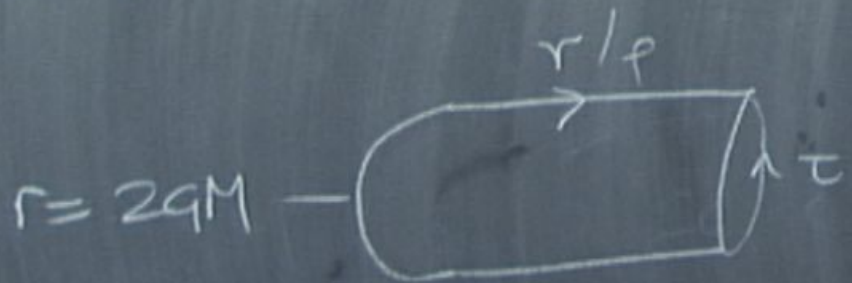


Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi GM$ .

then Euclidean-Sch is nonsingular at  $2GM$ .

"black hole cigar"



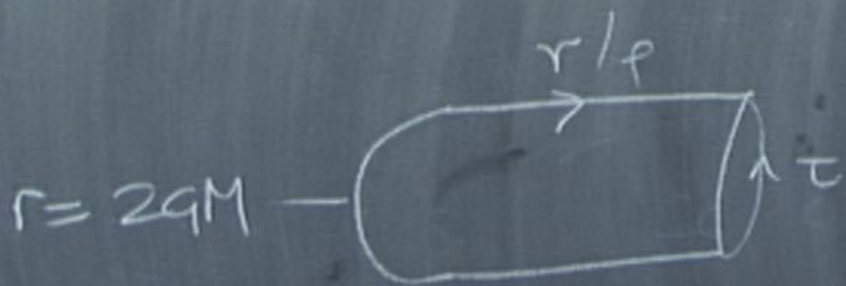
$$\Delta\tau = 8\pi GM = \beta.$$

Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi GM$ .

Then Euclidean-Sch is nonsingular at  $2GM$ .

"black hole cigar"

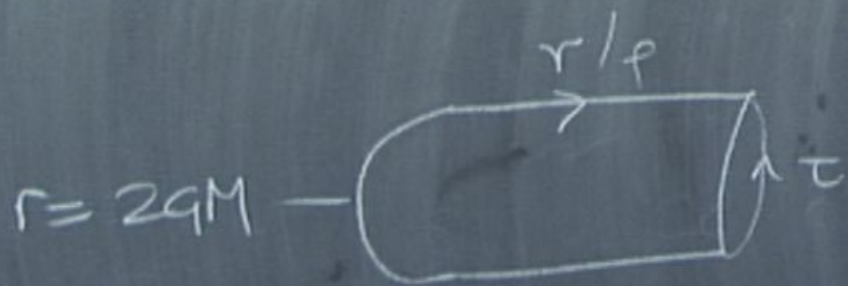


$$\Delta\tau = 8\pi GM = \beta.$$

Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi GM$ .

then Euclidean-Sch is nonsingular at  $2GM$ .



"black hole cigar"

$$\Delta\tau = 8\pi GM = \beta.$$

The black hole should induce a temperature

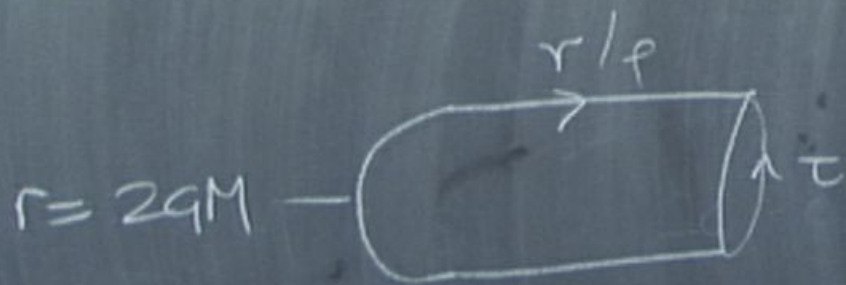
$$T = \frac{1}{8\pi GM} \rightarrow \frac{\hbar c^3}{8\pi GM k_B}$$

Compare to  $dr^2 + r^2 d\theta^2 \leftarrow \mathbb{R}^2$  in polar co-ords

If we make  $\tau$  periodic with period  $8\pi GM$ .

then Euclidean-Sch is nonsingular at  $2GM$ .

"black hole cigar"

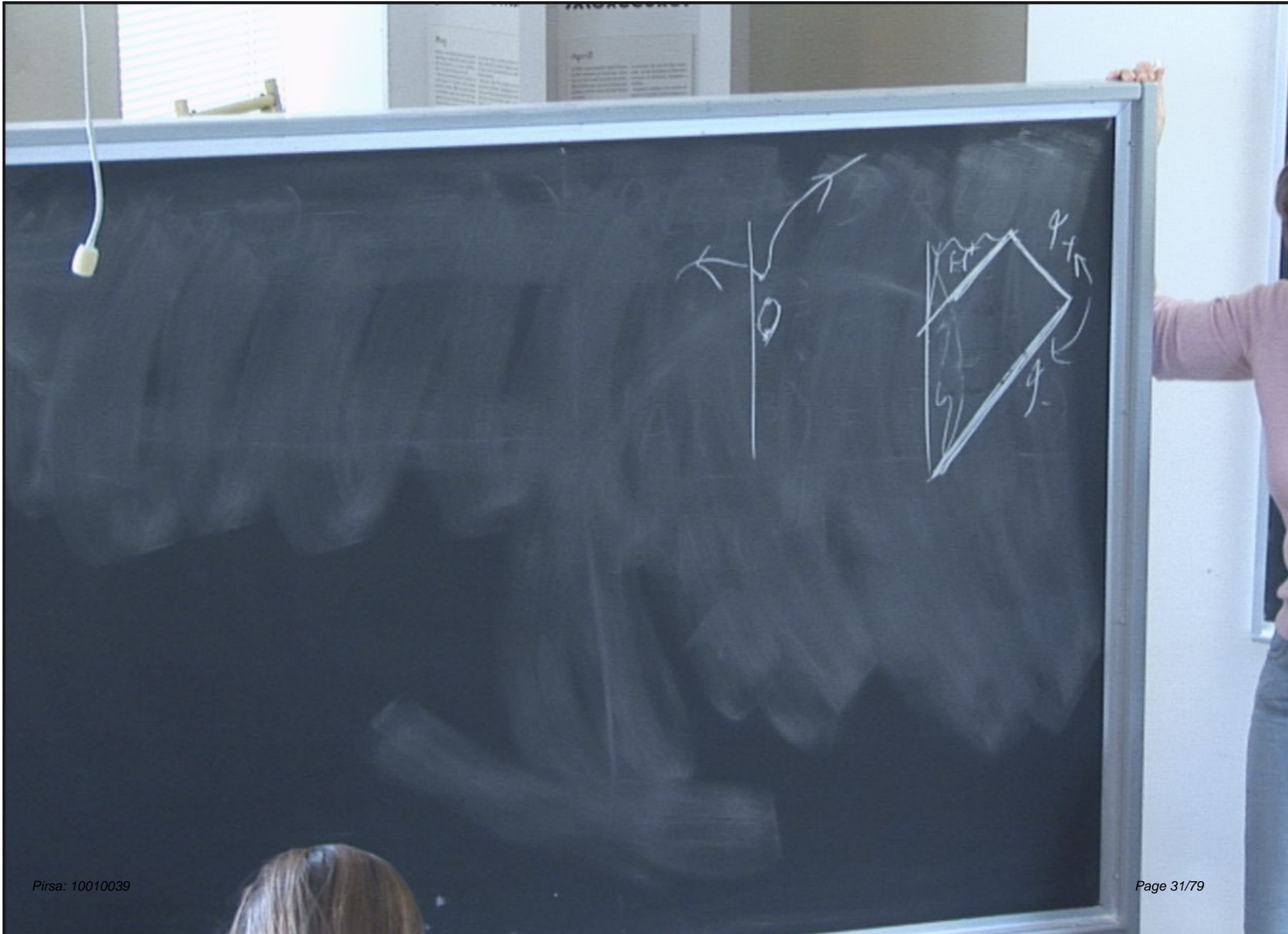


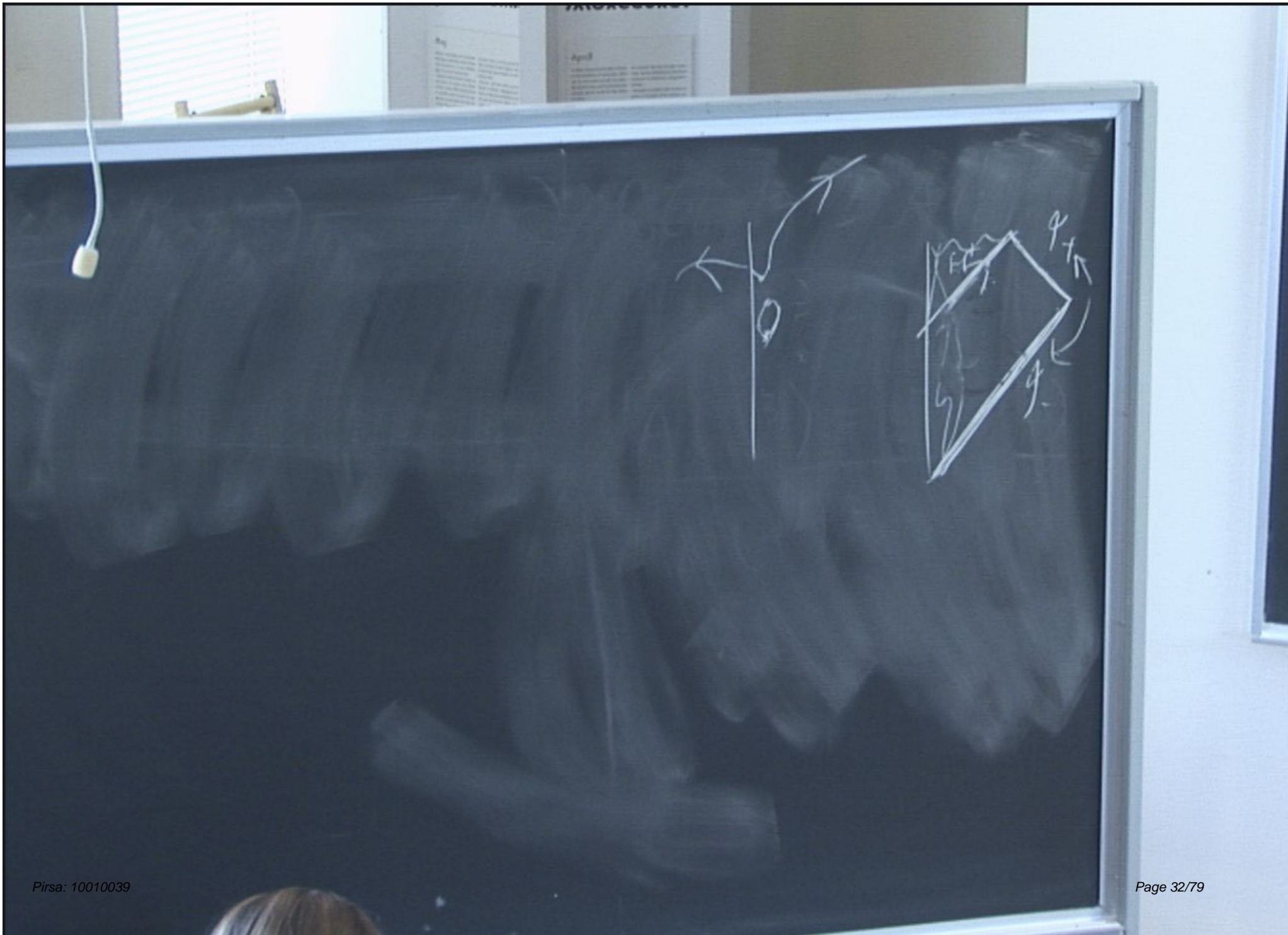
$$\Delta\tau = 8\pi GM = \beta.$$

The black hole should induce a temperature

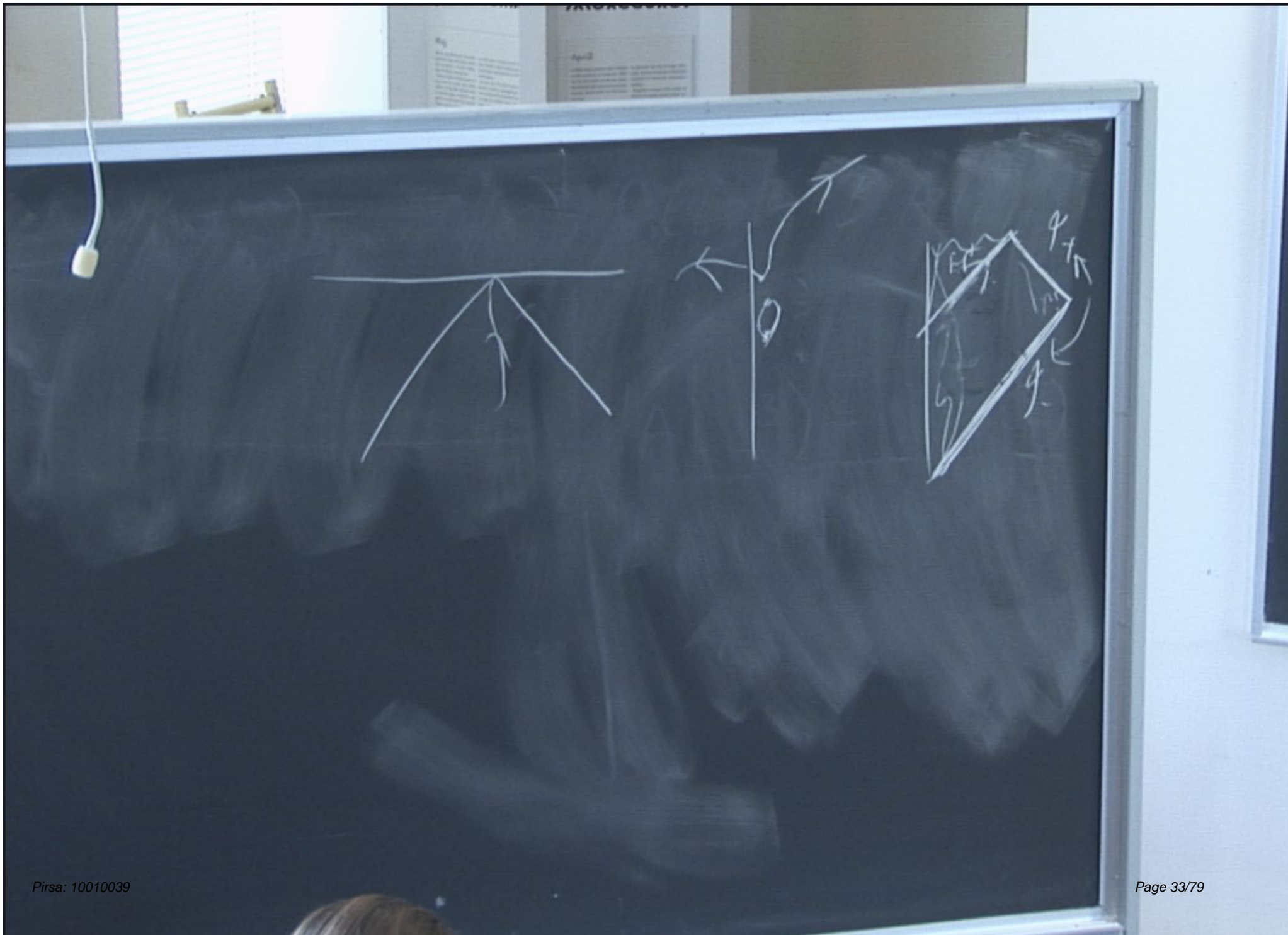
$$T = \frac{1}{8\pi GM} \rightarrow \frac{\hbar c^3}{8\pi GM k_B}$$

in any quantum field.









Another soln at finite  $T$  is flat space



Another soln at finite  $T$  is flat space



Another soln at finite  $T$  is flat space



For both flat space & E-Sch,

$$R = 0$$

Another soln at finite  $T$  is flat space



For both flat space & E-Sch,

$$R = 0 \quad \text{so} \quad I_E \stackrel{?}{=} \int \frac{R}{16\pi G} = 0$$

Another soln at finite  $T$  is flat space



For both flat space & E-Sch,

$$R = 0 \quad \text{so} \quad I_E \stackrel{?}{=} \int \frac{R}{16\pi G} = 0$$

Recall 
$$g^{ab} \delta R_{ab} = -\nabla_a \nabla_b \delta g^{ab} + \square (g_{ab} \delta g^{ab})$$

and we had to fix  $\delta g^{ab}$  and  $\delta n g^{ab} = 0$  on  
boundary

and we had to fix  $\delta g^{ab}$  and  $\partial_n g^{ab} = 0$  on  
boundary. Suggests we need to add a boundary  
term to the action to take care of this normal  
derivative.



and we had to fix  $\delta g^{ab}$  and  $\partial_n g^{ab} = 0$  on  
boundary. Suggests we need to add a boundary  
term to the action to take care of this normal  
derivative. It is also useful to have the notion  
of a boundary as a physical object.

and we had to fix  $\delta g^{ab}$  and  $\partial_n g^{ab} = 0$  on boundary. Suggests we need to add a boundary term to the action to take care of this normal derivative. It is also useful to have the notion of a boundary as a physical object.

$$\mu_+ \rightarrow \frac{\delta M_+}{\delta M} = \mu_-$$

Gauss-Codazzi formalism



# Gauss-Codazzi formalism



Describes embeddings  
of submanifolds in spacetime

# Gauss-Codazzi formalism



Describes embeddings  
of submanifolds in spacetime

$\Sigma$  — submanifold

$n^a$  — normal (possibly several)

# Gauss-Codazzi formalism



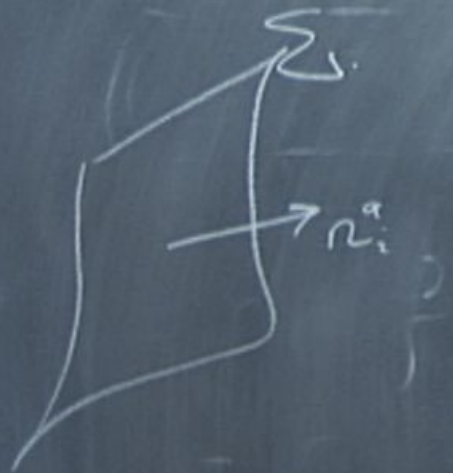
Describes embeddings  
of submanifolds in spacetime

$\Sigma$  - submanifold

$K \in n_i^a$  - normal (possibly several)

# lin indep normals = codimension of  $\Sigma$

# Gauss-Codazzi formalism



Describes embeddings  
of submanifolds in spacetime

$\Sigma$  - submanifold

$K \in n_i^a$  - normal (possibly several)

# lin indep normals = codimension of  $\Sigma$

e.g. in our universe, codim 1 — wall  
codim 2 — string  
3  
⋮

(domain)



e.g. in our universe, codim 1 - wall  
codim 2 - string  
codim 3 - point ptele (domain)

e.g. in our universe, codim 1 - wall  
codim 2 - string  
codim 3 - point ptele (domain)

$$\Sigma_d \subset \mathcal{M}_N$$

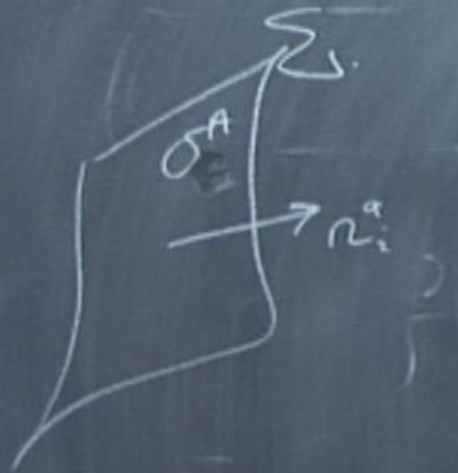
$$\text{codim } n = N - d$$

e.g. in our universe, codim 1 - wall  
codim 2 - string  
codim 3 - point<sup>2</sup> ptcle (domain)

$$\Sigma_d \subset \mathcal{M}_N$$

$$\text{codim } n = N - d$$

# Gauss-Codazzi formalism



Describes embeddings  
of submanifolds in spacetime

$\Sigma$  — submanifold

$K_{\mu}^{\nu}$  — normal (possibly several)

# lin indep normals = codimension of  $\Sigma$

e.g. in our universe, codim 1 - wall (domain)  
codim 2 - string  
codim 3 - point ptele

$$\Sigma_d \subset \mathcal{M}_N \quad \text{codim } n = N - d$$

Normals satisfy  $n(\sigma^A) = 0$

e.g. in our universe, codim 1 - wall (domain)  
 codim 2 - string  
 codim 3 - point ptele

$$\Sigma_d \subset \mathcal{M}_N$$

$$\text{codim } n = N - d$$

Normals satisfy  $n(\sigma^A) = 0$

$$n_{i\mu} \frac{\partial X^\mu}{\partial \sigma} = 0$$

e.g. in our universe, codim 1 - wall (domain)  
 codim 2 - string  
 codim 3 - point ptele

$$\Sigma_d \subset \mathcal{M}_N$$

$$\text{codim } n = N - d$$

Normals satisfy  $n(\sigma^A) = 0$

$$n_{i\mu} \frac{\partial X^\mu}{\partial \sigma^A} = 0$$

e.g. in our universe, codim 1 - wall (domain)  
 codim 2 - string  
 codim 3 - point ptele

$$\Sigma_d \subset \mathcal{M}_N \quad \text{codim } n = N - d$$

Normals satisfy  $n(\sigma^A) = 0$   $n_{i\mu} \frac{\partial X^\mu}{\partial \sigma^A} = 0$

(will take  $n_i{}^\mu n_{j\mu} = 0 \quad i \neq j$ )



The 1<sup>st</sup> fundamental form of  $\Sigma$  is

$$h_{ab} = g_{ab}$$

The 1<sup>st</sup> fundamental form of  $\Sigma$  is

$$h_{ab} = g_{ab} + \sum_{i=1}^n \eta_{ia} \eta_{ib} (-)^i$$

↑  
- timelike  
+ spacelike.

The 1<sup>st</sup> fundamental form of  $\Sigma$  is

$$h_{ab} = g_{ab} + \sum_{i=1}^n \eta_{ia} \eta_{ib} (-)^i$$

↑  
- timelike  
+

$h_{ab}$  is the projection of the metric of  $M$  onto  $\Sigma$  (still lies in  $T^*(M)^2$ )

$h_{ab}$  is the projection of the metric of  $M$  onto  $\Sigma$  (still lives in  $T^*(M)^2$ )

has in the projection of the metric of  $M$  onto  $\Sigma$  (still lies in  $T^*(M)^2$ )

Can also view  $\Sigma$  as a manifold,

$\frac{\partial X^M}{\partial \sigma^A}$  projects  $T$

T.

$h_{ab}$  is the projection of the metric of  $M$  onto  $\Sigma$  (still lives in  $T^*(M)^2$ )

Can also view  $\Sigma$  as a manifold,

$$\frac{\partial X^M}{\partial \sigma^A} \text{ projects } T_p^*(M) \rightarrow T_p^*(\Sigma)$$

$$g_{\mu\nu} \frac{\partial X^M}{\partial \sigma^A} \frac{\partial X^\nu}{\partial \sigma^B} = \gamma_{AB} \text{ - intrinsic metric}$$

$h_{ab}$  is the projection of the metric of  $M$  onto  $\Sigma$  (still lives in  $T^*(M)^2$ )

Can also view  $\Sigma$  as a manifold,

$$\frac{\partial X^M}{\partial \sigma^A} \text{ projects } T_p^*(M) \rightarrow T_p^*(\Sigma)$$

$$g_{\mu\nu} \frac{\partial X^M}{\partial \sigma^A} \frac{\partial X^\nu}{\partial \sigma^B} = \gamma_{AB} \text{ — intrinsic metric}$$



The 2<sup>nd</sup> fundamental form or extrinsic  
curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c n_{id}$$

The 2<sup>nd</sup> fundamental form or extrinsic  
curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c \Gamma_{id}$$

measures how  $\Sigma$  curves in  $M$

The 2<sup>nd</sup> fundamental form or extrinsic  
curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c \Pi_{id}$$

measures how  $\Sigma$  curves in  $M$



The 2<sup>nd</sup> fundamental form or extrinsic  
curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c \Pi_{id}$$

measures how  $\Sigma$  curves in  $M$



The 2<sup>nd</sup> fundamental form or extrinsic  
curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c n_d$$

measures how  $\Sigma$  curves in  $M$

$K_{iAB}$



The 2<sup>nd</sup> fundamental form or extrinsic curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c n_{id}$$

measures how  $\Sigma$  curves in  $M$



$$K_{iAB} = \frac{\partial X^c}{\partial \sigma^A} \frac{\partial X^d}{\partial \sigma^B} \nabla_c n_{id}$$

$$= -n_{id} \nabla_A \left( \frac{\partial X^d}{\partial \sigma^B} \right)$$

The 2<sup>nd</sup> fundamental form or extrinsic

curvature is

$$K_{iab} = h_a^c h_b^d \nabla_c n_{id}$$

measures how  $\Sigma$  curves in  $M$

$$K_{iAB} = \frac{\partial X^c}{\partial \sigma^A} \frac{\partial X^d}{\partial \sigma^B} \nabla_c n_{id}$$

$$= -n_{id} \nabla_A \left( \frac{\partial X^d}{\partial \sigma^B} \right)$$



Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



$$n_a = (\cos\theta, \sin\theta, 0)$$

$$h_{ab} = \delta_{ab} - n_a n_b$$

Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



$$n_a = (\cos\theta, \sin\theta, 0)$$

$$h_{ab} = \delta_{ab} - n_a n_b$$

$$= \begin{pmatrix} \sin^2\theta & -\sin\theta\cos\theta & 0 \\ -\sin\theta\cos\theta & \cos^2\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \text{ in Cartesian.}$$

Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



$$n_a = (\cos\theta, \sin\theta, 0)$$

$$h_{ab} = \delta_{ab} - n_a n_b$$

$$= \begin{pmatrix} \sin^2\theta & -\sin\theta\cos\theta & 0 \\ -\sin\theta\cos\theta & \cos^2\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \text{ in Cartesian.}$$

$$= \begin{pmatrix} 0 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & 1 \end{pmatrix} \text{ in polars.}$$

Easiest to see by example:

$$\Sigma = x^2 + y^2 = a^2 \text{ in } \mathbb{R}^3$$



$$n_a = (\cos\theta, \sin\theta, 0)$$

$$h_{ab} = \delta_{ab} - n_a n_b$$

$$= \begin{pmatrix} \sin^2\theta & -\sin\theta\cos\theta & 0 \\ -\sin\theta\cos\theta & \cos^2\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \text{ in Cartesian.}$$

$$= \begin{pmatrix} 0 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & 1 \end{pmatrix} \text{ in polars.}$$

In polars

$$K_{ab} = \nabla_a n_b = -T_{ab}^r$$

In polars

$$K_{ab} = \nabla_a n_b = -\Gamma_{ab}^r$$

$$K_{\theta\theta} = r = a$$

Taking  $(\theta, z)$  coords on cylinder

$$K_{AB} = \begin{pmatrix} a & 0 \\ 0 & 0 \end{pmatrix} \quad \gamma_{AB} = \begin{pmatrix} a^2 & 0 \\ 0 & 1 \end{pmatrix}$$