

Title: Gravitational Physics - Review (PHYS 636) - Lecture 7

Date: Jan 11, 2010 10:00 AM

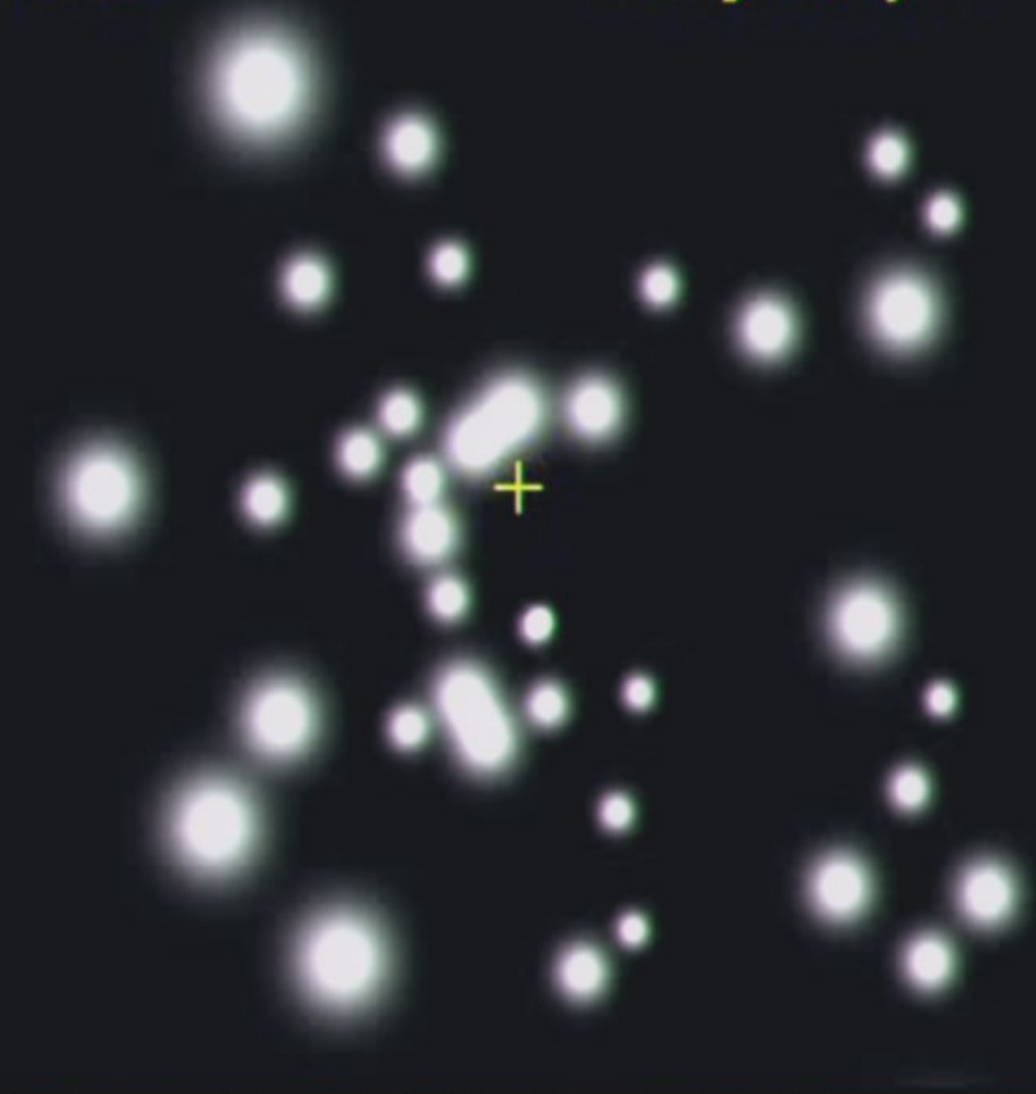
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Abstract:



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Recall our sph. symm. eqns.



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$$M(r) = \int 4\pi r^2 T^0_0(r) dr$$

r-r eqn.

$$8\pi G\rho = \frac{2A'}{AB^2r} + \frac{1}{r^2}(B^{-2}-1)$$



r-r eqn

$$8\pi G\rho = \frac{2A'}{AB^2r} + \frac{1}{r^2}(B^{-2}-1)$$

$$\Rightarrow \frac{(A^2)'}{A^2} = \frac{2GM(r) + 8\pi G r^3 \rho}{r(r - 2GM(r))}$$

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TOLMAN  
OPPENHEIMER  
VOLKOFF

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$$\begin{aligned} \text{Then } p' &= -(\rho + p_0) \frac{(G M(r) + 4\pi G r^3 p)}{r(r - 2G M(r))} \\ &= -(\rho + p_0) \frac{4\pi G r (3p + \rho_0)}{3r - 8\pi G r^2 \rho_0} \end{aligned}$$



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solved by:

$$p = p_0$$

$$\left[ \frac{\sqrt{1 - \frac{2GM}{R}} - \sqrt{1 - \frac{2GM r^2}{R^3}}}{\sqrt{1 - \frac{2GM r^2}{R^3}} - 3 \sqrt{1 - \frac{2GM}{R}}} \right]$$



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$$p(0) = p_0 \left( \frac{\sqrt{1 - \frac{2GM}{R}} - 1}{1 - 3 \sqrt{1 - \frac{2GM}{R}}} \right)$$



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$$\rightarrow \infty \text{ when } R = \frac{9GM}{4}$$



ie.  $R_{\max} = \frac{9G}{4} \cdot \frac{4\pi}{3} R_{\max}^3 \rho_0$

$$\text{i.e. } R_{\max} = \frac{9G}{4} \cdot \frac{4\pi}{3} R_{\max}^3 \rho_0$$

$$\Rightarrow R_{\max} = \sqrt[3]{\frac{3\pi G \rho_0}{4}}$$



$$\text{i.e. } dR_{\max} = \frac{9G}{4} \cdot \frac{4\pi}{3} R_{\max}^3 \rho_0$$

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$$\Rightarrow R_{\text{max}} = \frac{1}{\sqrt{3\pi G \rho_0}}$$

$$\& M_{\text{max}} = \frac{4}{9G R_{\text{max}}}$$

- limit to mass of star - Chandrasekhar  
 $\sim 1.4 M_{\odot}$



How do we "observe" a black hole

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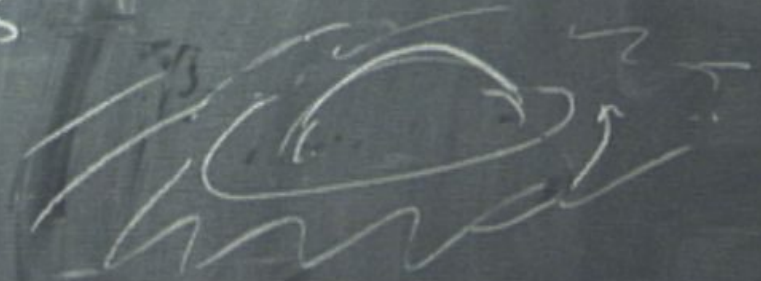
- lensing
- orbits of visible objects
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then 1<sup>st</sup> integral,  $g_{ab} \dot{X}^a \dot{X}^b = k \begin{cases} 1 & m^2 > 0 \\ 0 & \text{null} \end{cases}$



Equatorial geodesics:  $\Theta = \mathbb{P}^2$ .

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NEWTONIAN  
POT.



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$\uparrow$   
GR  
(ATT)

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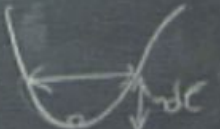


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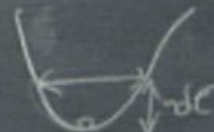
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Real" black holes rotate - Kerr Soln

$$ds^2 = dt^2 - \frac{\rho^2}{\Delta} (dr^2 + \Delta d\theta^2) - (r^2 + a^2) \sin^2\theta d\phi^2 - \frac{2GMr}{\rho^2} (dt - a \sin^2\theta d\phi)^2$$

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