

Title: Asymptotics of the Ponzano-Regge model for handlebodies

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URL: <http://pirsa.org/09120113>

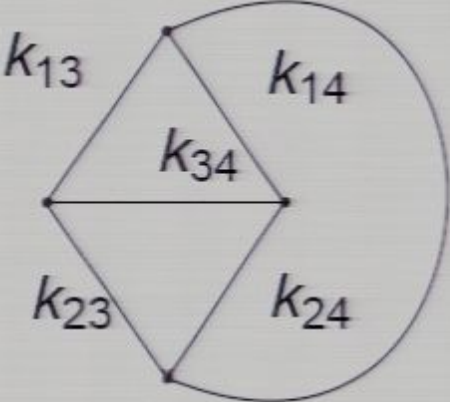
Abstract: The asymptotic formula for the Ponzano-Regge model amplitude is given for non-tardis triangulations of handlebodies in the limit of large boundary spins. The formula produces a sum over all possible immersions of the boundary triangulation in three dimensional Euclidean space weighted by the cosine of the Regge action evaluated on these immersions. Furthermore the asymptotic scaling registers the existence of flexible immersions.

Plan

- 6j symbol and the Ponzano-Regge model
- Coherent intertwiners
- Asymptotic formula
- Flexible polyhedra

Asymptotics of the 6j symbol

6j symbol given by the $SU(2)$ spin network

$$\left\{ \begin{array}{ccc} k_{12} & k_{13} & k_{14} \\ k_{34} & k_{24} & k_{23} \end{array} \right\} = \text{Diagram} \quad (1)$$


The diagram shows a tetrahedron with vertices at the top, bottom-left, bottom-right, and a central point. The edges are labeled as follows: top-left edge is k_{13} , top-right edge is k_{14} , bottom-left edge is k_{23} , bottom-right edge is k_{24} , and the central horizontal edge is k_{34} . The edge k_{12} is represented by a large circle on the right side of the tetrahedron, indicating it is the edge connecting the top and bottom vertices.

Ponzano and Regge gave an asymptotic formula

$$\left\{ \begin{array}{ccc} \lambda k_{12} & \lambda k_{13} & \lambda k_{14} \\ \lambda k_{34} & \lambda k_{24} & \lambda k_{23} \end{array} \right\} \rightarrow \frac{1}{\sqrt{12\pi \text{Vol}}} \cos \left(\sum_{a < b} (\lambda k_{ab} + \frac{1}{2}) \Theta_{ab} + \frac{\pi}{4} \right)$$

For $\lambda \rightarrow \infty$

- Vol - volume of tetrahedron with edge lengths $\lambda k_{ab} + \frac{1}{2}$
- Θ_{ab} - dihedral angles of the tetrahedron

The Ponzano-Regge model

3d quantum gravity

On the triangulation \mathcal{T} of 3-manifold-with-boundary Σ^3

$$\mathcal{Z}_{PR}(\Psi, \mathcal{T}) = \sum_{j_e} \prod_{\text{edges } e} \dim(j_e) \prod_{\text{tetrahedra}} \{6j\} \quad (2)$$

Ψ is the boundary state - depends on boundary spins

Sum over internal spin label

- Spin foam model for 3d gravity
- Fourier transform of path integral for discrete $SU(2)$ BF theory
- Topological invariant with regularization

One tetrahedron doesn't capture the dynamics (ie 4d BF)

What is the asymptotic formula for the whole model?

The Ponzano-Regge model

Spin network on the boundary

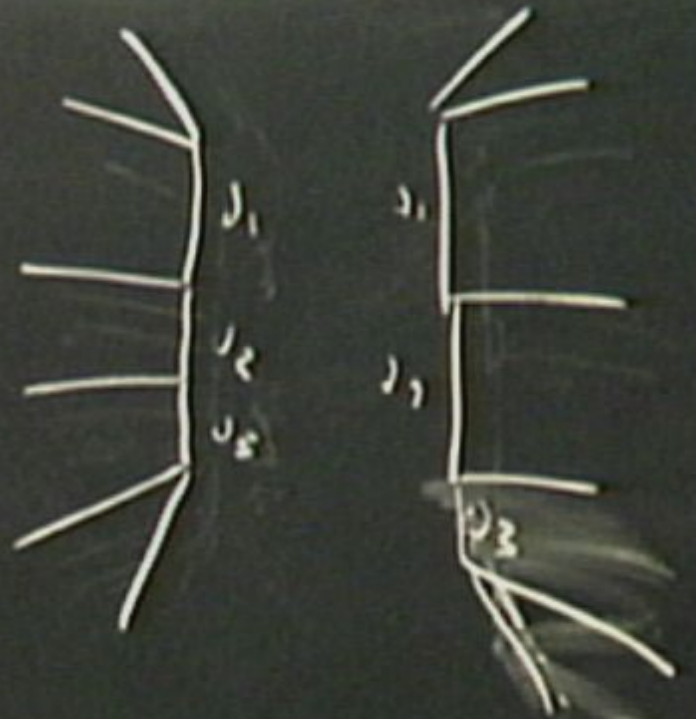
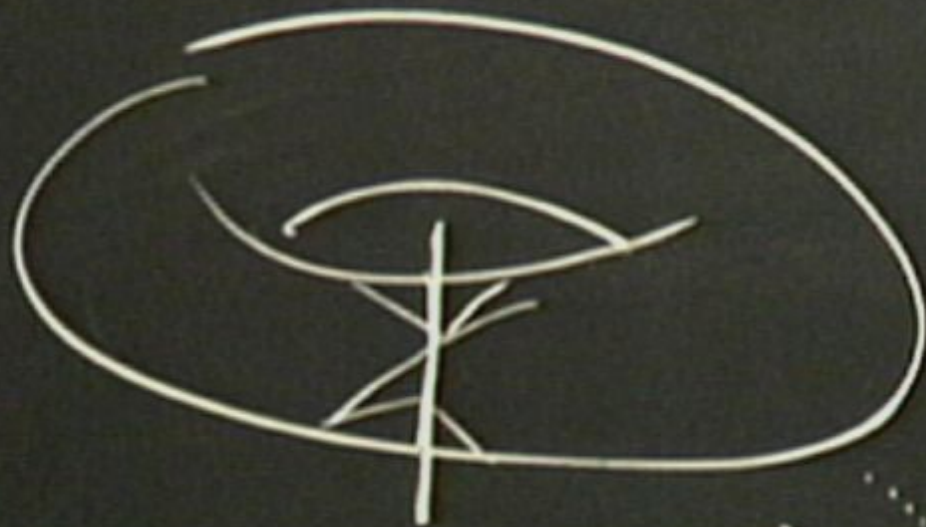
PR model was constructed to agree with boundary spin network on S^2

Works for any handlebody of genus g

Reduce Σ^3 to B^3 with a set of cuts C

At each cut, replace the projector with a group integral

$$\int_{\text{SU}(2)} dh D_{ab}^{j_1}(h) \otimes D_{cd}^{j_2}(h) \otimes D_{ef}^{j_3}(h) = \iota \iota^\dagger$$



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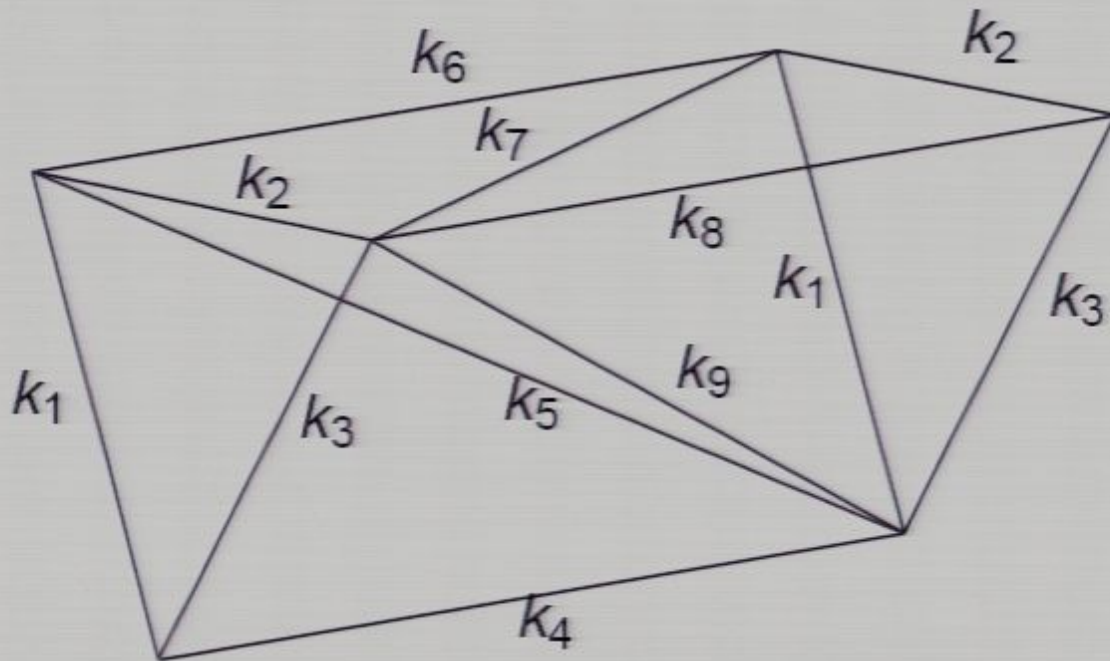
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The Ponzano-Regge model

Spin network on the boundary. Example: the solid torus



Identify k_1, k_2, k_3

$$\mathcal{Z}_{PR}(\Psi, \mathbb{T}) = \left\{ \begin{array}{ccc} k_1 & k_2 & k_3 \\ k_8 & k_9 & k_7 \end{array} \right\} \left\{ \begin{array}{ccc} k_1 & k_2 & k_3 \\ k_9 & k_4 & k_5 \end{array} \right\} \left\{ \begin{array}{ccc} k_1 & k_5 & k_6 \\ k_2 & k_7 & k_9 \end{array} \right\}.$$

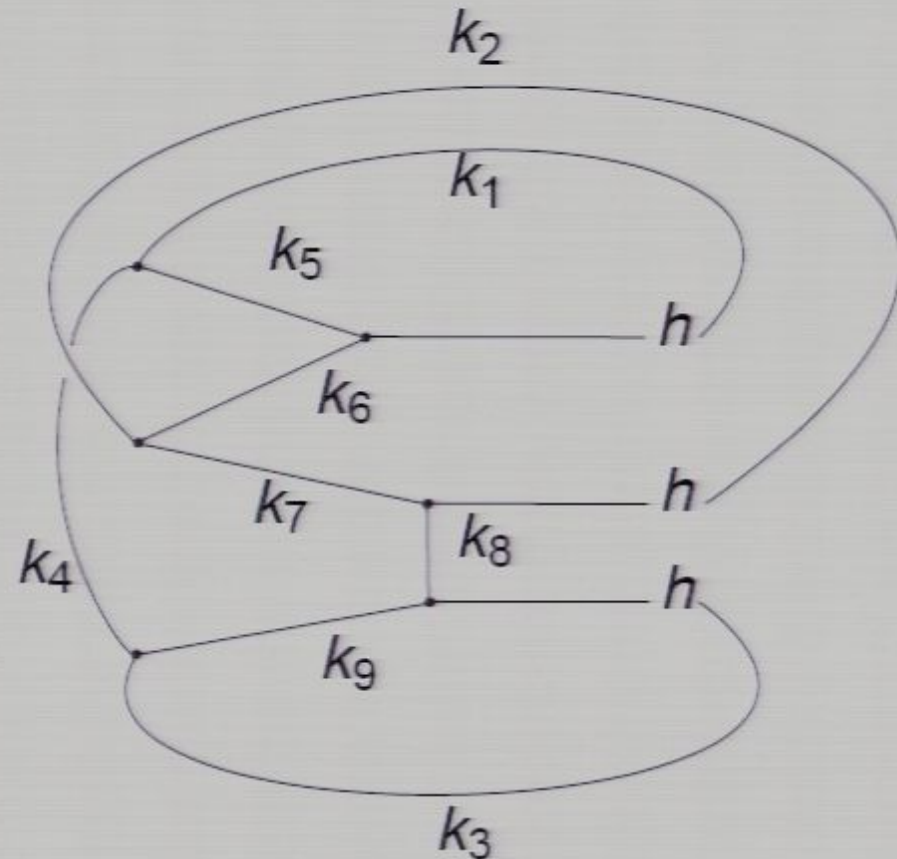
(3)

The Ponzano-Regge model

Spin network on the boundary of the solid torus

PR amplitude can be expressed as spin network on the boundary

$$\mathcal{Z}_{PR}(\Psi, \mathbb{T}) = \int_{\text{SU}(2)} dh$$



Non-planar diagram

$$\int dh D^{j_1}(h) \otimes D^{j_2}(h) \otimes D^{j_3}(h)$$

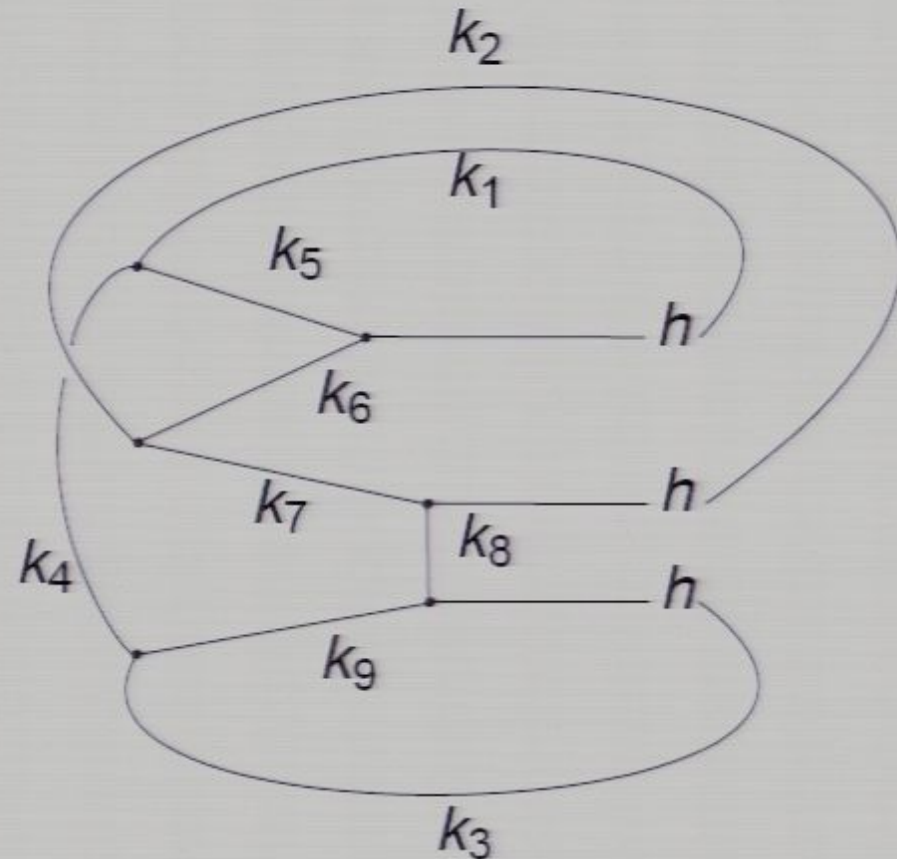
$$= \int dh \left(\begin{array}{c} j_1 \quad j_2 \quad j_3 \\ \diagdown \quad | \quad / \\ \text{---} \\ \diagup \quad | \quad \diagdown \\ j_1 \quad j_2 \quad j_3 \end{array} \right) = \int dh \left(\begin{array}{c} j_1 \quad j_2 \quad j_3 \quad j_4 \\ | \quad | \quad | \quad | \\ \text{---} \\ | \quad | \quad | \quad | \\ j_1 \quad j_2 \quad j_3 \quad j_4 \end{array} \right)$$

The Ponzano-Regge model

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Non-planar diagram

Stationary phase approximation

We want to approximate the amplitude for large boundary k_{ab}
Use stationary phase approximation

$$\int dx a(x) e^{\lambda S(x)} \simeq \sum_{x_0} a(x_0) \left(\frac{2\pi}{\lambda} \right)^{n/2} \frac{1}{\sqrt{\det(-H)}} e^{\lambda S(x_0)}$$

- S complex valued function with $\text{Re}S \leq 0$
- x_0 are critical points - $\delta S(x_0) = 0$, $\text{Re}S(x_0) = 0$
- n is dimension of integral
- H is Hessian matrix

Need to rewrite the amplitude in a different basis.

Coherent triangles

SU(2) coherent states

SU(2) spin- k representation space denoted V_k :

A coherent state $|k, k\rangle$ has

- Minimum uncertainty
- Maximum projection in $(0, 0, 1)$ direction

A state $|k, \mathbf{n}\rangle := g(\mathbf{n})|k, k\rangle$ has maximum projection in \mathbf{n} direction, ie

$$L \cdot \mathbf{n} |k, \mathbf{n}\rangle = ik |k, \mathbf{n}\rangle$$

k and \mathbf{n} specify $|k, \mathbf{n}\rangle$ up to a phase

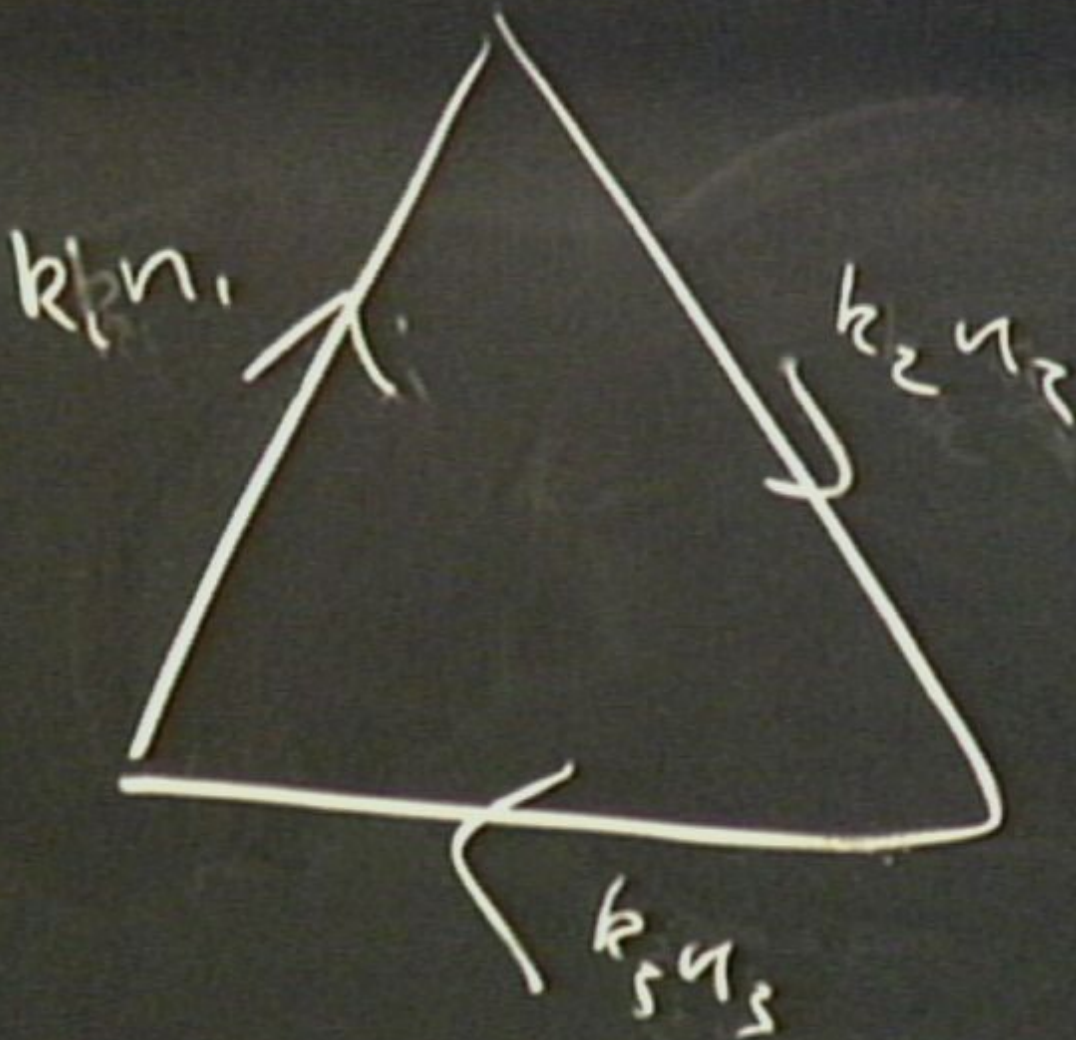
Coherent triangles

Geometry of the boundary

We specify the geometry of a triangle, a , on the boundary with

- Lengths of edges k_{ab} (spins)
- Unit edge vectors \mathbf{n}_{ab}

This gives 3 coherent states $|k_{ab}, \mathbf{n}_{ab}\rangle \in V_{k_{ab}}$



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Coherent triangles

Coherent intertwiners

Up to phase specifies coherent intertwiner

$$l_a \in \text{Inv}(V_{k_{ab}} \otimes V_{k_{ac}} \otimes V_{k_{ad}})$$

$$l_a = \int_{\text{SU}(2)} dX_a \bigotimes_{b: b \neq a} X_a |k_{ab}, \mathbf{n}_{ab}\rangle \quad (4)$$

How do we specify the phase?

Can pick the \mathbf{n}_{ab} orthogonal to some vector, say $\mathcal{N} = (0, 0, 1)$

$\exists g_{ab} \in U(1) \subset \text{SU}(2)$ such that

$$\mathbf{n}_{ba} = -g_{ab} \triangleright \mathbf{n}_{ab}$$

So define

$$|k_{ab}, \mathbf{n}_{ba}\rangle := g_{ab} |k_{ab}, -\mathbf{n}_{ab}\rangle$$

With this choice, \mathcal{Z}_{PR} does not depend on the phase choice

The Ponzano-Regge model

In the coherent state basis

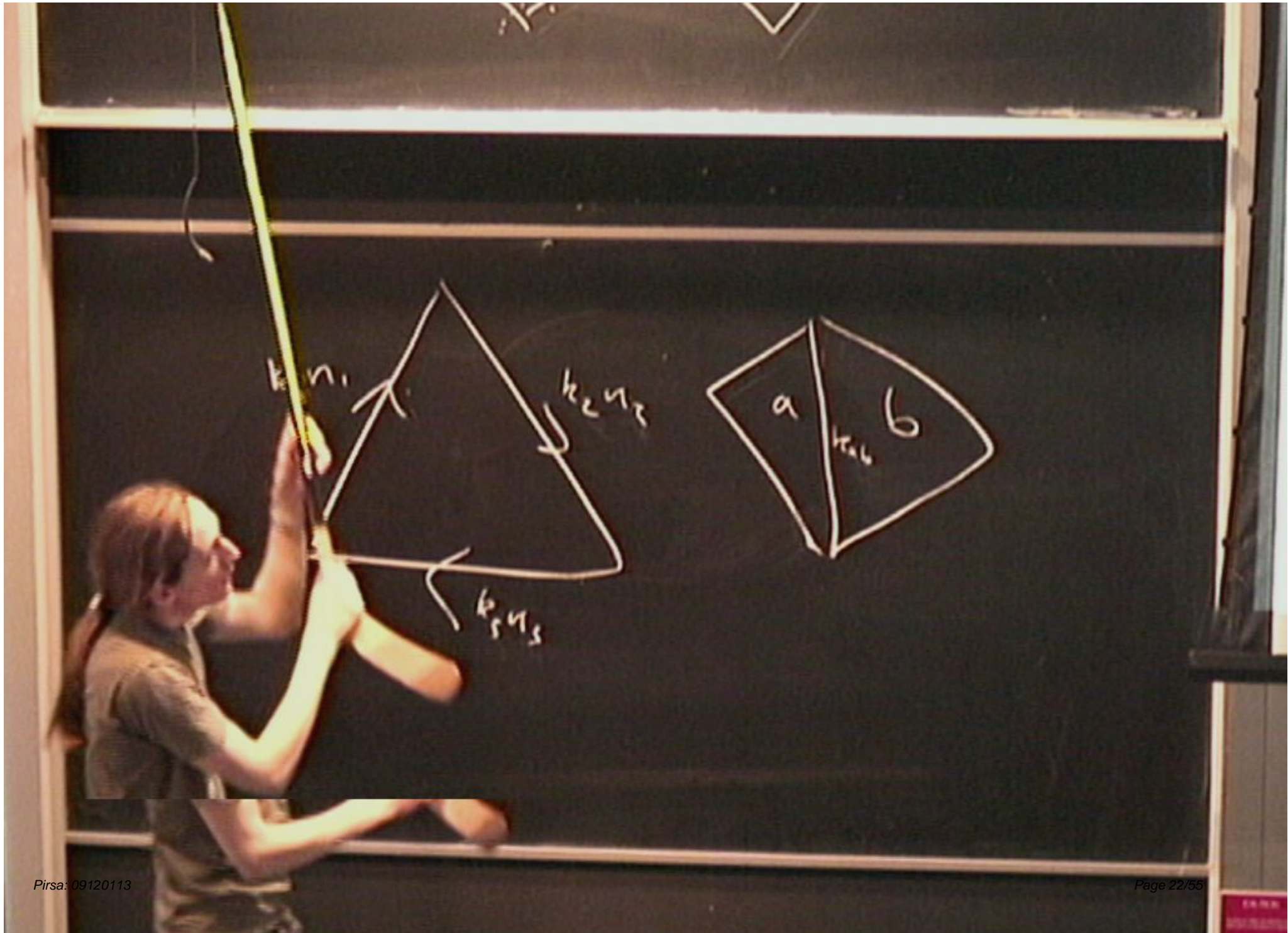
The amplitude now reads

$$\mathcal{Z}_{PR}(\Psi, \Sigma^3) = \int \prod_{\text{triangles } i} dX_i \prod_{\text{cuts } j} dh_j e^S \quad (5)$$

with the action

$$\begin{aligned} S = & \sum_{\text{edges } ab} 2k_{ab} \ln \langle -\mathbf{n}_{ab} | X_a^\dagger X_b | \mathbf{n}_{ba} \rangle \\ & + \sum_{\text{cuts } l} \sum_{\substack{\text{edges} \\ de \in l}} 2k_{de} \ln \langle -\mathbf{n}_{de} | X_d^\dagger h_l X_e | \mathbf{n}_{ed} \rangle. \end{aligned} \quad (6)$$

Can now use stationary phase analysis



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Can now use stationary phase analysis

Stationary phase

Amplitude dominated by the critical points

$$\delta S = 0, \quad \text{Re} S = 0$$

This gives the conditions:

$$\hat{X}_a \mathbf{n}_{ab} = -\hat{X}_b \mathbf{n}_{ba} \quad \text{Gluing}$$

$$\sum_{b: b \neq a} k_{ab} \mathbf{n}_{ab} = 0 \quad \text{Closure}$$

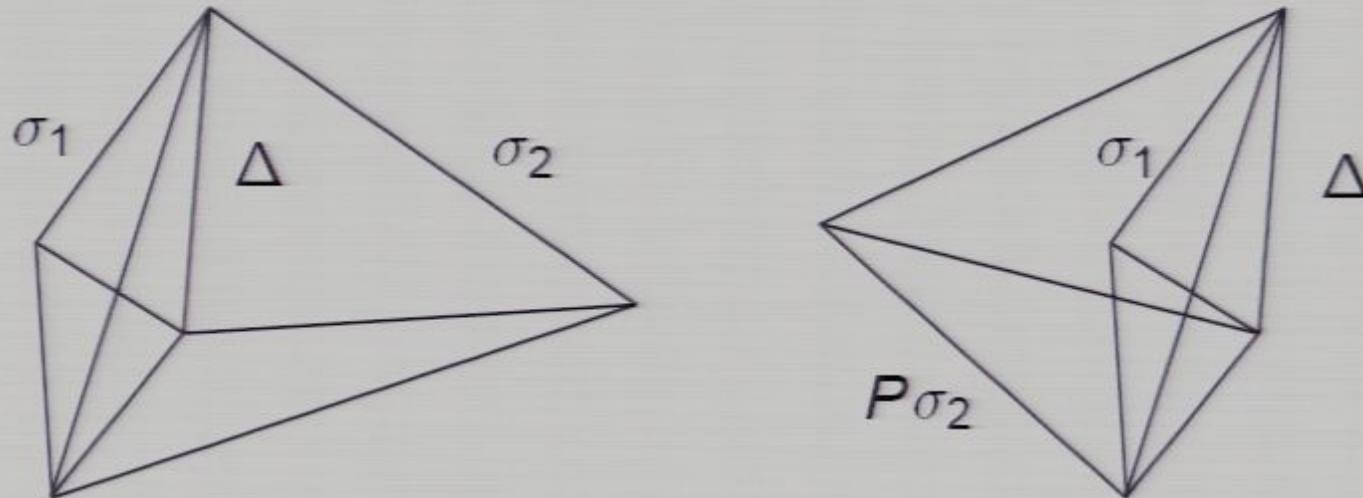
$$\sum_{ab \in C_i} k_{ab} \hat{X}_a \mathbf{n}_{ab} = 0 \quad \text{Closure on cuts}$$

Each critical point gives an immersion of boundary $k_{ab} \mathbf{n}_{ab}$ in \mathbb{R}^3

Stationary points

Example: 3-ball

Triangulate B^3 with two tetrahedra σ_1, σ_2

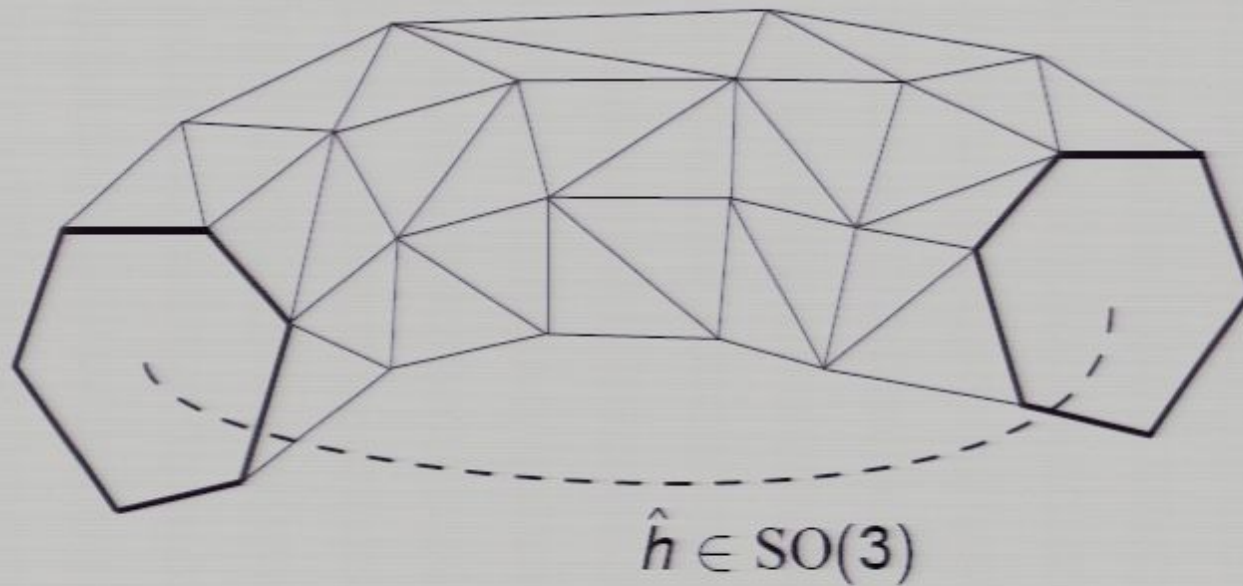


Two different possible immersions in the asymptotic formula
Parity related immersions also contribute

Stationary points

Example: $g > 0$ The solid torus

Cut immersions of the type shown below contribute



Stationary points

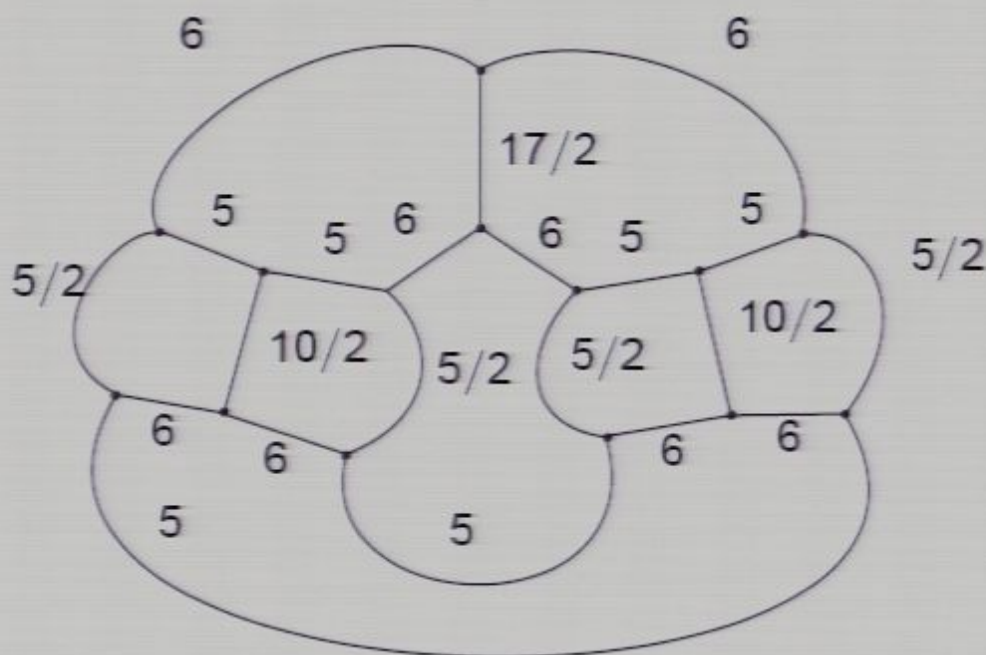
Non-isolated critical points and flexible polyhedra

Critical points may not be isolated - ie manifold of critical points

Corresponds to a flexible immersion

Example: Steffen's flexible polyhedron

<http://demonstrations.wolfram.com/SteffensFlexiblePolyhedron/>



Spin network for Steffen's polyhedron

Stationary Phase

The action at the critical points

At critical points

$$\begin{aligned} S &= \sum_{\text{edges } ab} k_{ab} \Theta_{ab}^i \\ &= S_{\text{Regge}}(i) \end{aligned}$$

Regge action of the immersion

Use choice of spin structure to fix some minus signs

Stationary points

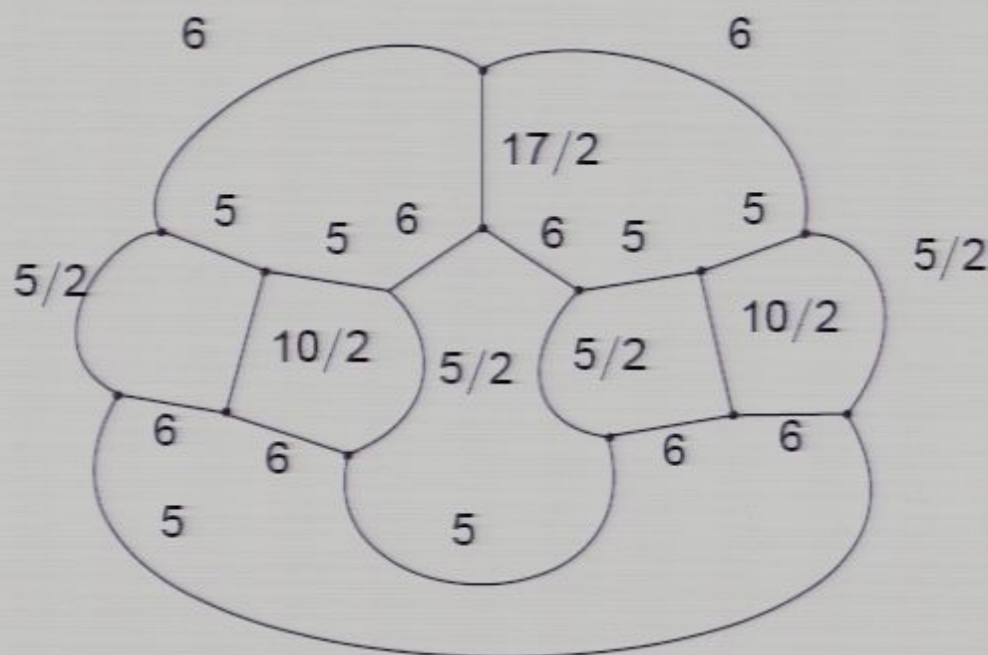
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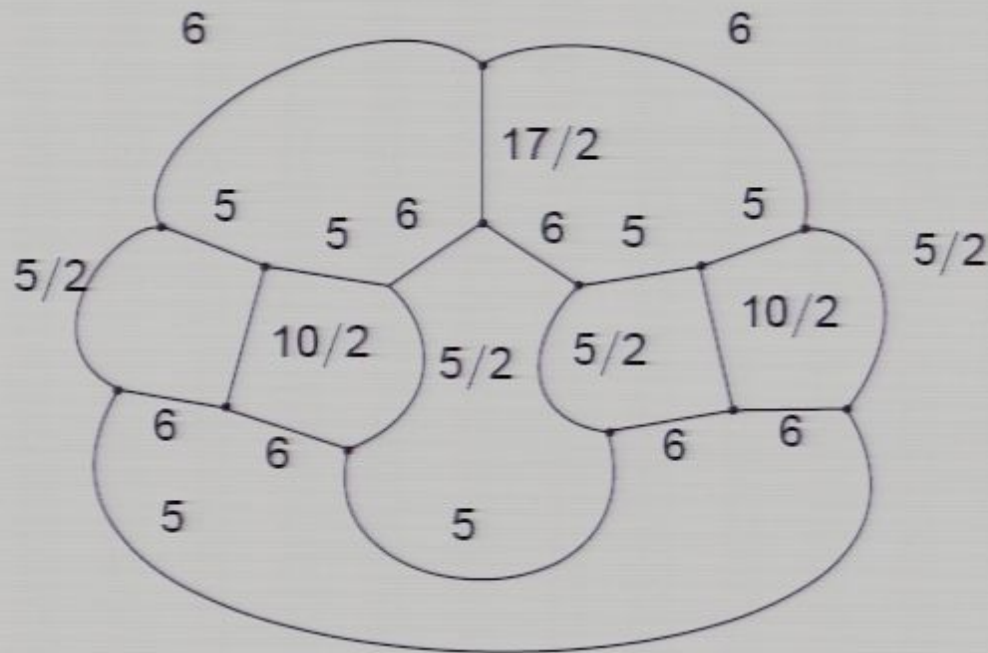
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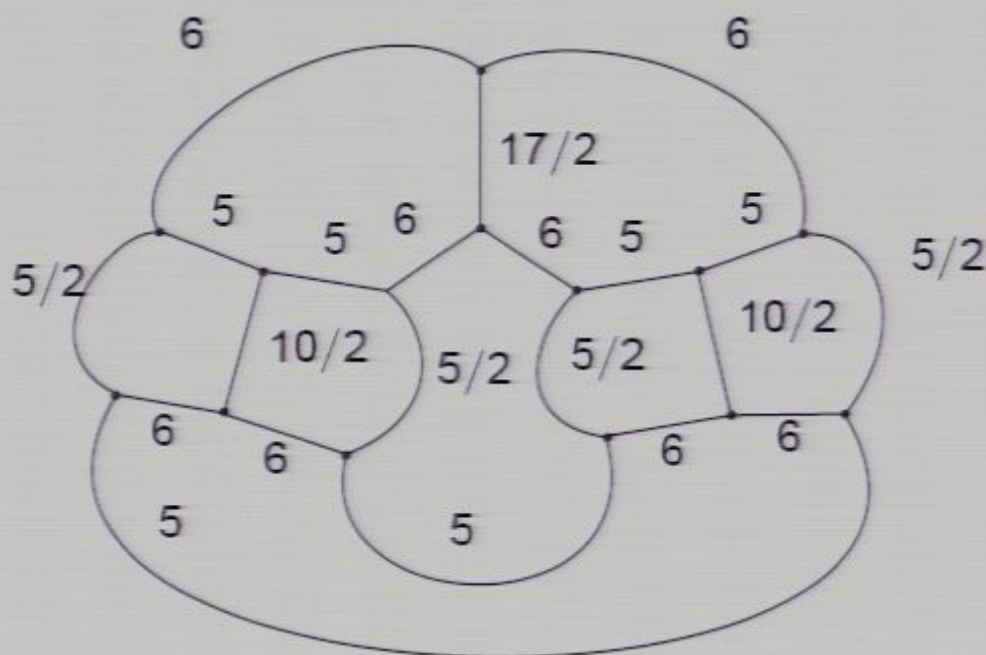
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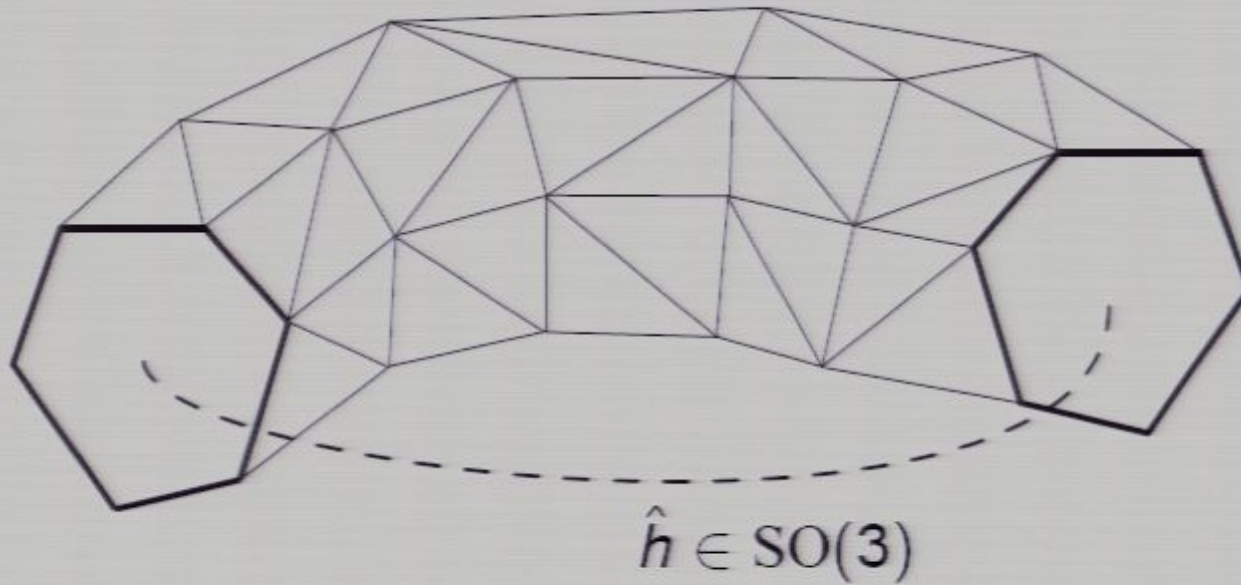
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Example: $g > 0$ The solid torus

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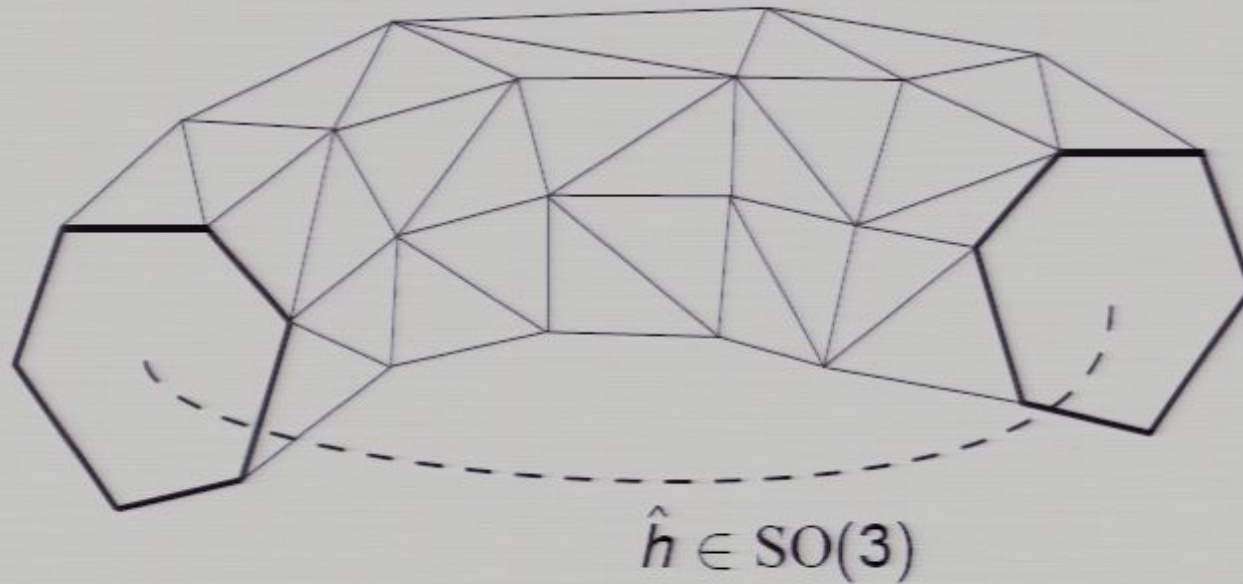
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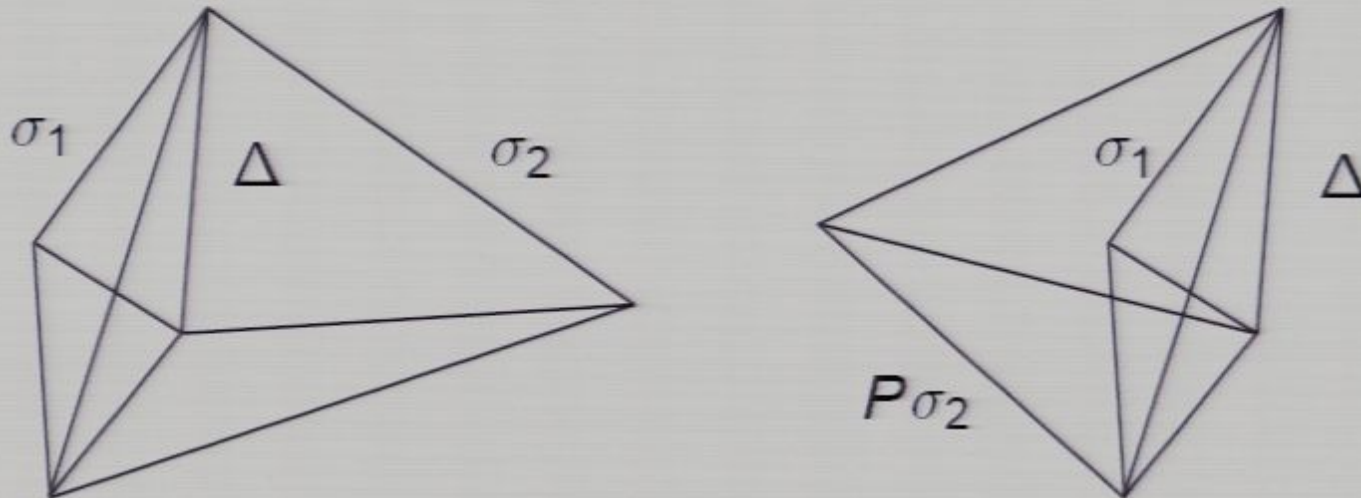
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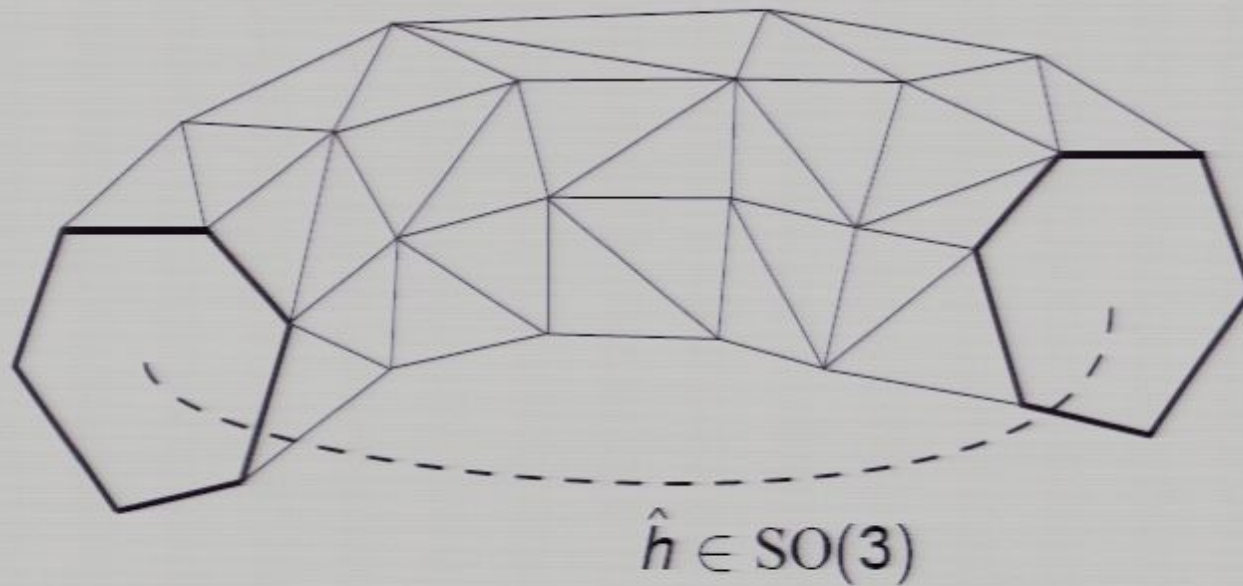


Two different possible immersions in the asymptotic formula
Parity related immersions also contribute

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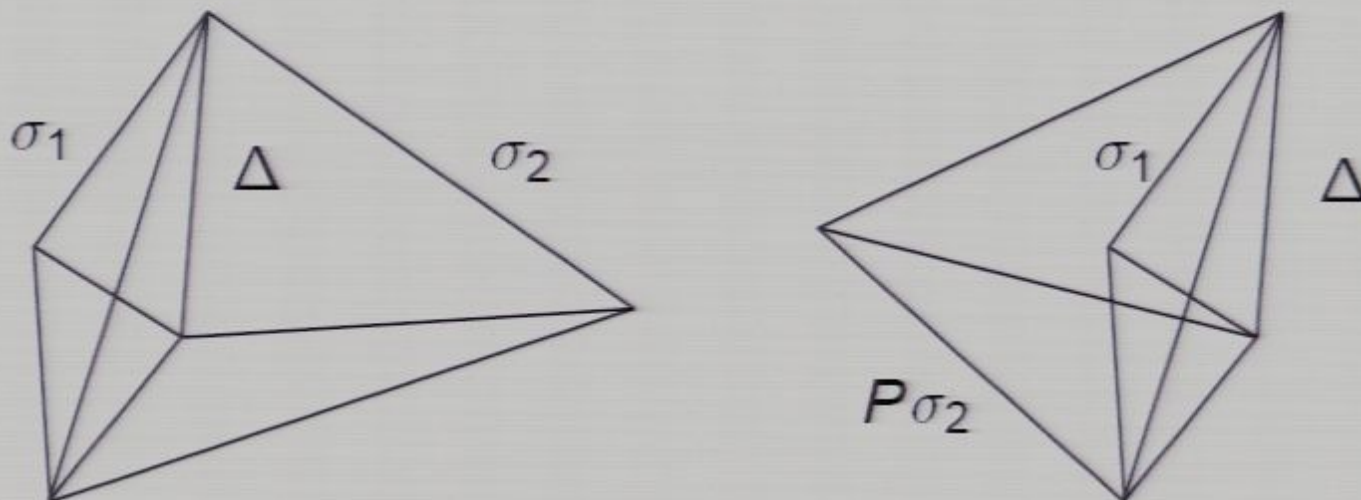
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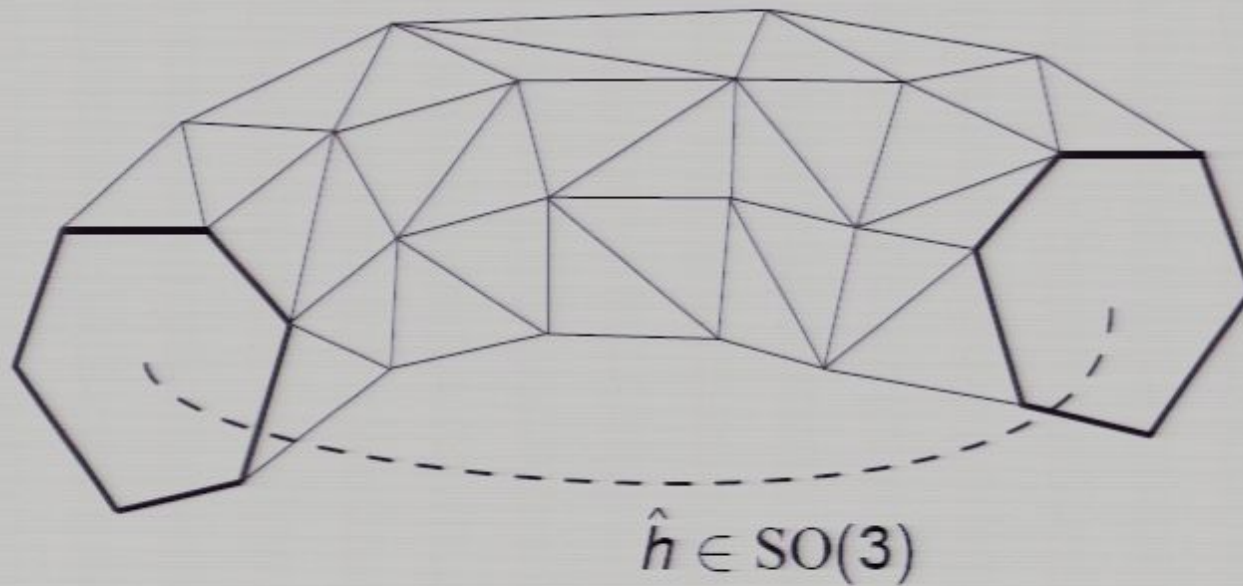


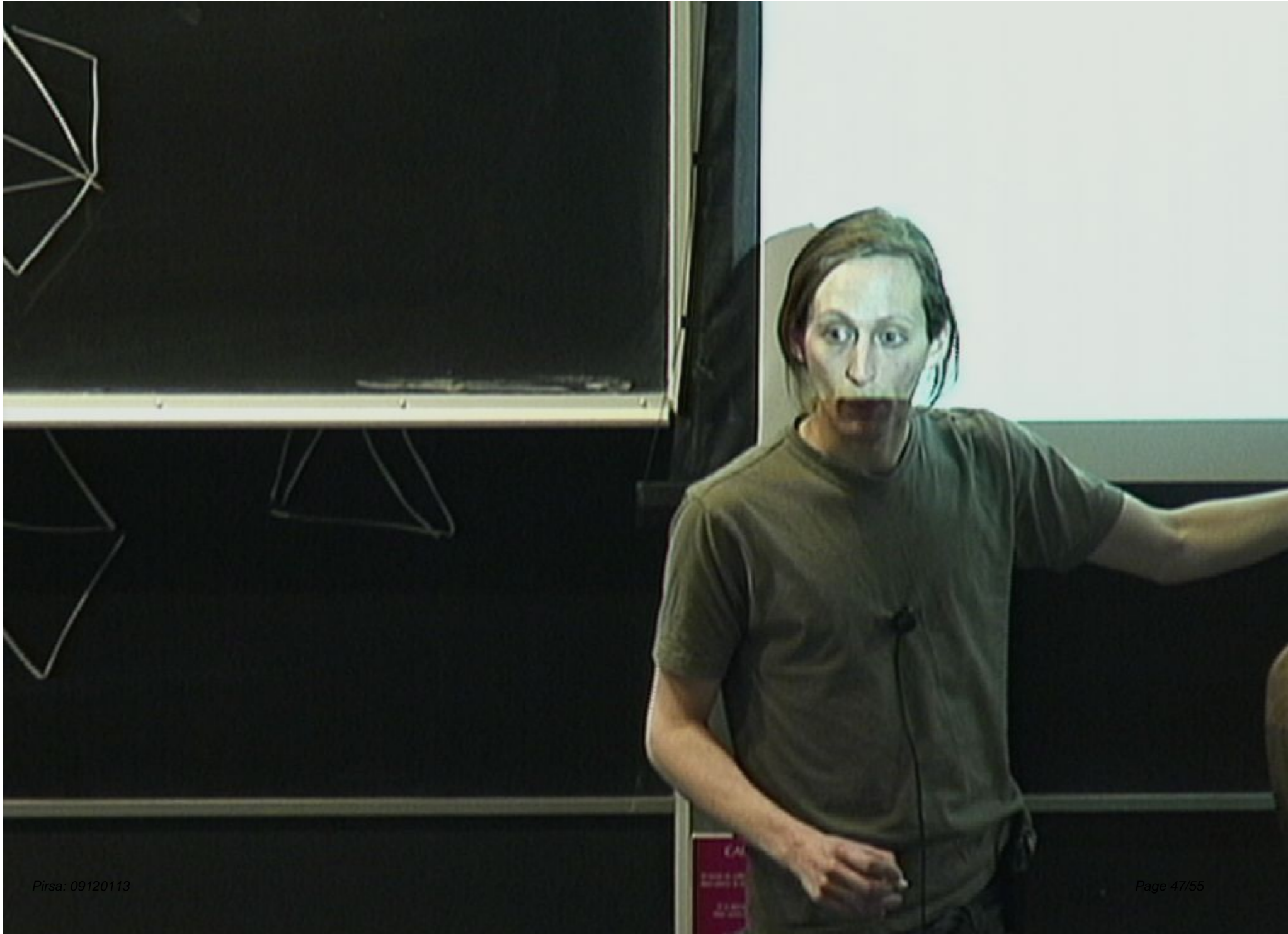
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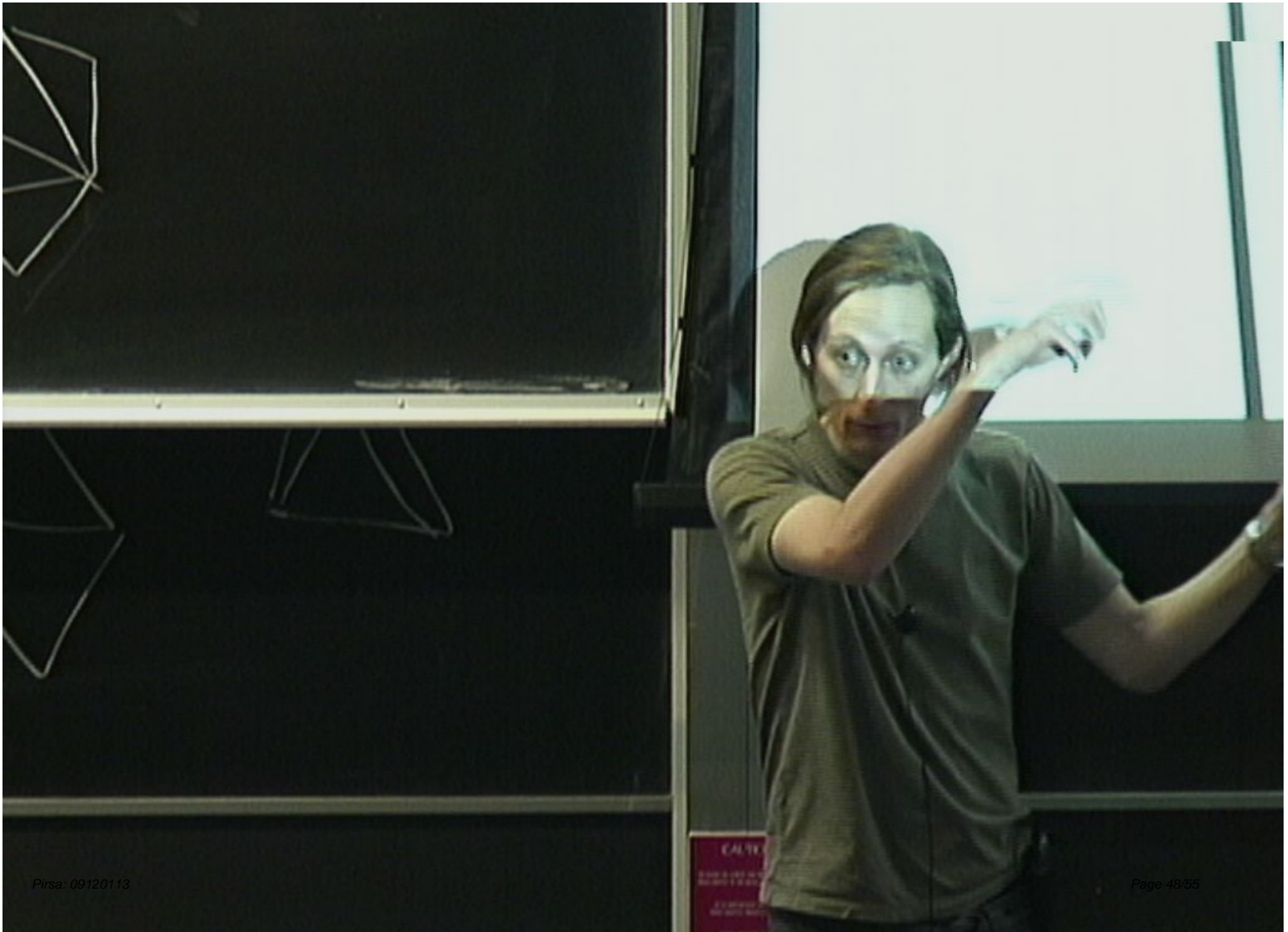
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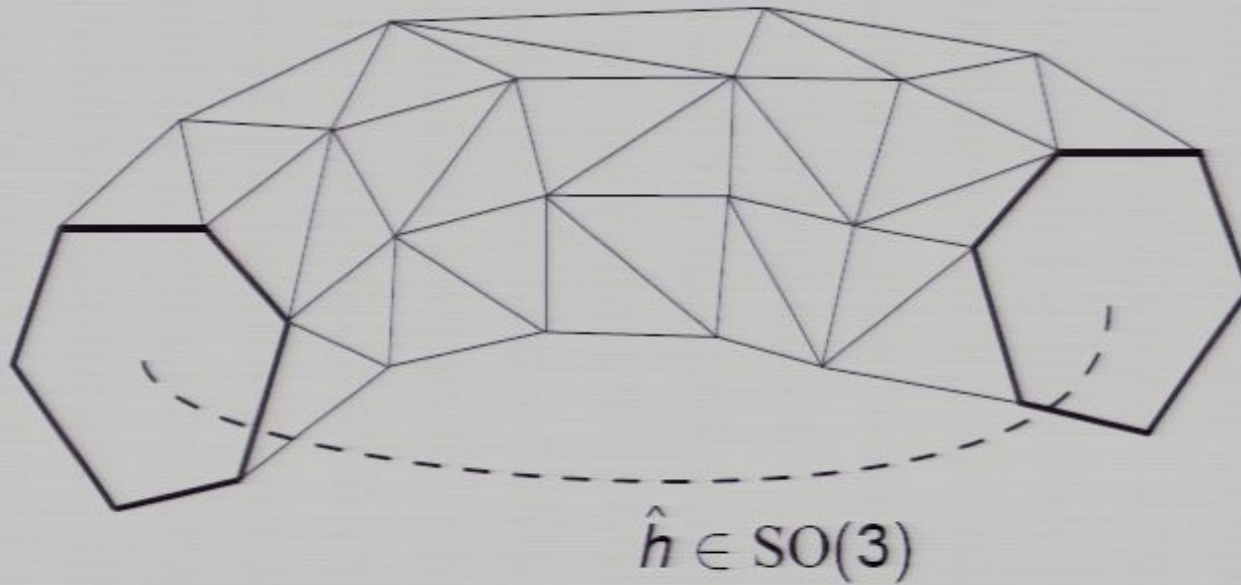




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Asymptotic formula

For isolated critical points

When $\lambda \rightarrow \infty$

$$\mathcal{Z}_{PR}(\psi_\lambda, \Sigma^3) = \left(\frac{2\pi}{\lambda}\right)^{\frac{3(|V|+|C|-1)}{2}} \sum_{\text{Immersion } i} N_i \cos \left(\lambda S_{\text{Regge}}(i) + \phi_{ab}^i \right)$$

- Θ_{ab}^i are the dihedral angles of i
- $|V|$ is number of triangles, $|C|$ is number of cuts

Stationary Phase

The action at the critical points

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Regge action of the immersion

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Asymptotic formula

Non-isolated critical points

$$\mathcal{Z}_{PR}(\psi_\lambda, \Sigma^3) = \left(\frac{2\pi}{\lambda}\right)^{\frac{3(|V|+|C|-1)-d}{2}} \sum_{\text{flexible immersions } f} N_f \cos\left(\lambda \mathcal{S}_{\text{Regge}}(f) + \phi_{ab}^f\right)$$

- d is dimension of critical manifold
- Flexible immersions dominate
- Will flexibility be an issue in 4 dimensions?

Outlook

Refinement of the triangulation

What happens when we refine the boundary triangulation?

$$S_{Regge} \rightarrow S_{Einstein}$$

- Many more terms in the sum over immersions
- Related to immersions of manifolds with smooth metrics?

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