

Title: Foundations of Quantum Mech. (PHYS 639) - Lecture 11

Date: Dec 14, 2009 11:00 AM

URL: <http://pirsa.org/09120088>

Abstract:

# *The deBroglie-Bohm interpretation*



Louis deBroglie  
(1892-1987)

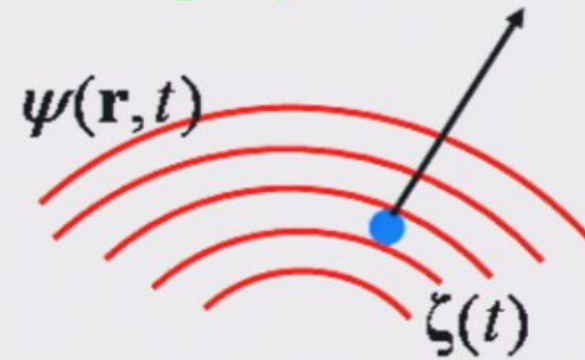


David Bohm  
(1917-1992)

## The deBroglie-Bohm interpretation for a single particle

The ontic state:  $(\psi(\mathbf{r}), \zeta)$

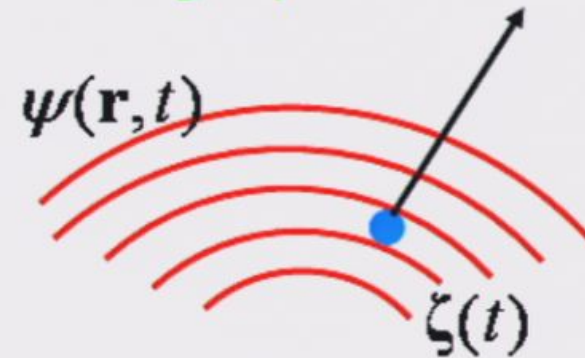
Wavefunction  $\nearrow$  Particle position  $\nwarrow$



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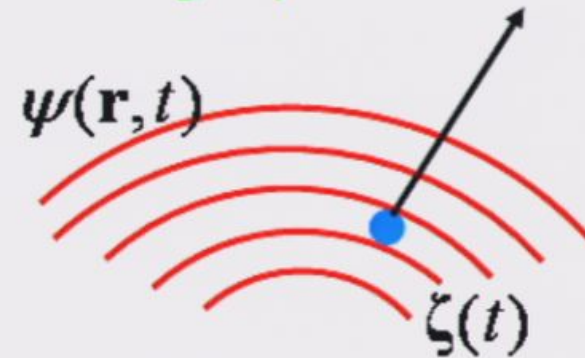
The evolution equations:

$$i\hbar \frac{\partial \psi(\mathbf{r}, t)}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi(\mathbf{r}, t) + V(\mathbf{r})\psi(\mathbf{r}, t) \quad \text{Schrödinger's eq'n}$$

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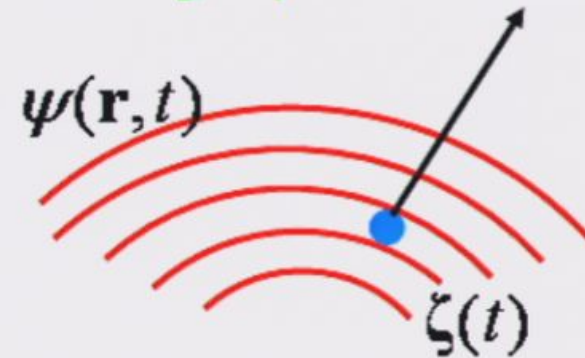
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$$\frac{\partial S}{\partial t} + \frac{(\nabla S)^2}{2m} + Q + V = 0$$

where  $Q(\mathbf{r}, t) \equiv -\frac{\hbar^2}{2m} \frac{\nabla^2 R(\mathbf{r}, t)}{R(\mathbf{r}, t)}$

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The imaginary part of the Schrodinger eq'n is:

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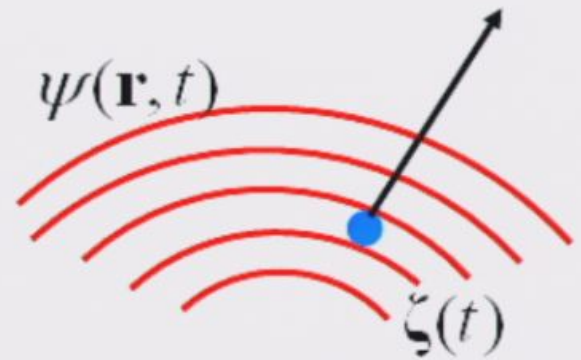
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Newtonian form of the particle dynamics:



$$m \frac{d^2 \zeta(t)}{dt^2} = -[\nabla V(\mathbf{r}) + \nabla Q(\mathbf{r}, t)]_{\mathbf{r}=\zeta(t)}$$

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(Note independence of quantum potential on amplitude)

Acting the  $\nabla$  operator on the real part of the Schrodinger eq'n gives:

$$\nabla \left[ \frac{\partial S}{\partial t} + \frac{(\nabla S)^2}{2m} + Q + V \right] = 0$$

$$\left( \frac{\partial}{\partial t} + \frac{\nabla S \cdot \nabla}{m} \right) \nabla S = -\nabla(Q + V)$$

Taking the time derivative of the guidance equation gives:

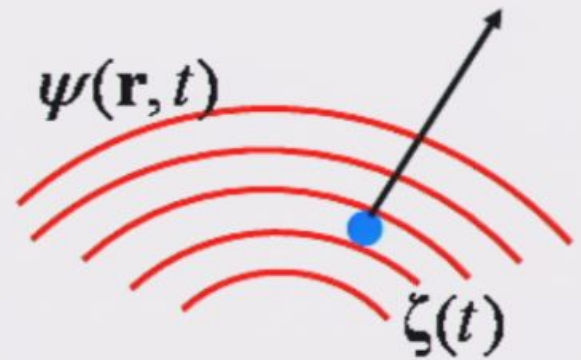
$$\frac{d\dot{\zeta}(t)}{dt} = \frac{1}{m} [\nabla S(\mathbf{r}, t)]_{\mathbf{r}=\zeta(t)}$$

$$\frac{d^2\zeta(t)}{dt^2} = \frac{1}{m} \left( \frac{\partial}{\partial t} + \frac{d\dot{\zeta}}{dt} \cdot \nabla \right) \nabla S$$

Thus

$$m \frac{d^2\zeta(t)}{dt^2} = -[\nabla V(\mathbf{r}) + \nabla Q(\mathbf{r}, t)]_{\mathbf{r}=\zeta(t)}$$

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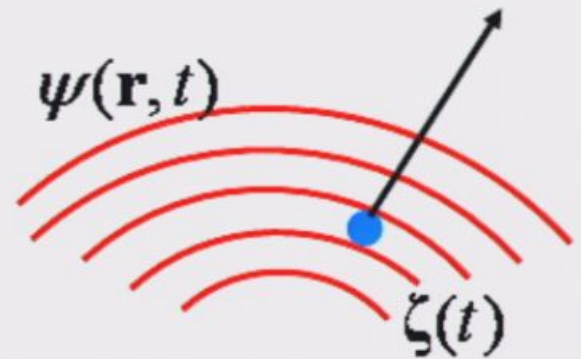


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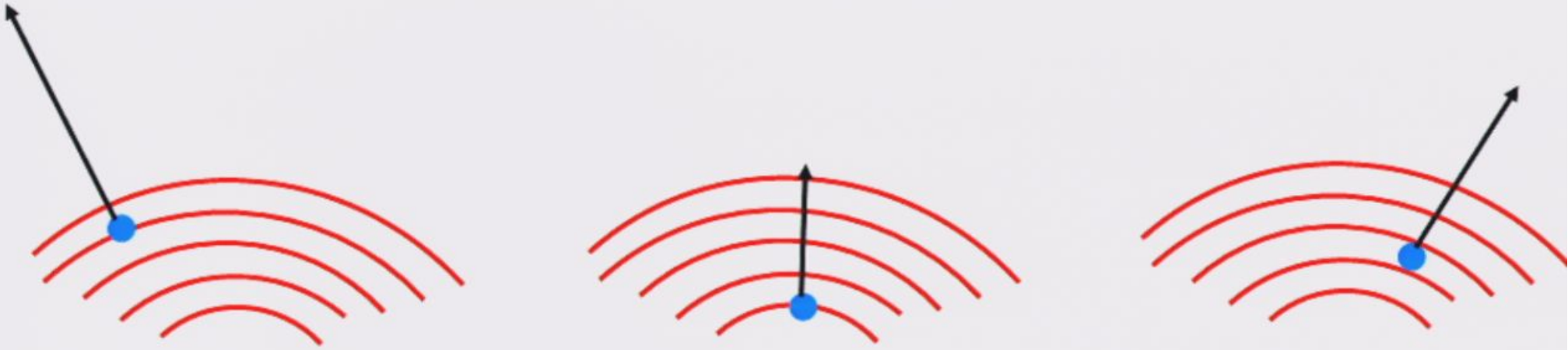
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Nonetheless the dynamics are *fundamentally first order*

$$\frac{d\zeta(t)}{dt} = \frac{1}{m} [\nabla S(\mathbf{r}, t)]_{\mathbf{r}=\zeta(t)}$$



**Epistemic state** (assuming perfect knowledge of  $\psi(\mathbf{r}, t)$ )

$\rho(\zeta)d\zeta$  = the probability the particle is within  $d\zeta$  of  $\zeta$ .

The "standard distribution"

$$\rho(\zeta, t) = |\psi(\zeta, t)|^2$$

Note: it is preserved by the dynamics:

$$\text{if } \rho(\zeta, 0) = |\psi(\zeta, 0)|^2 \text{ then } \rho(\zeta, t) = |\psi(\zeta, t)|^2$$



Proof of the preservation of the standard distribution:

The velocity field is

$$\mathbf{v}(\mathbf{r}, t) = \frac{1}{m} [\nabla S(\mathbf{r}, t)]$$

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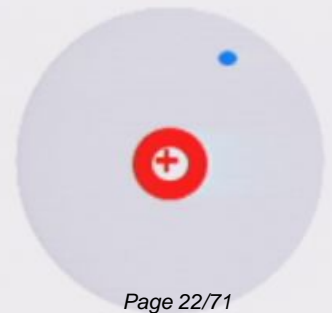
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# 1s orbital of Hydrogen atom



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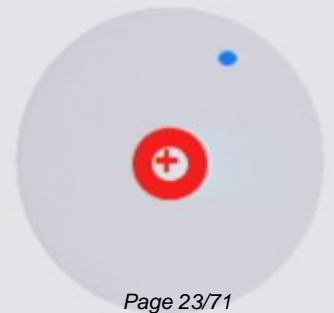
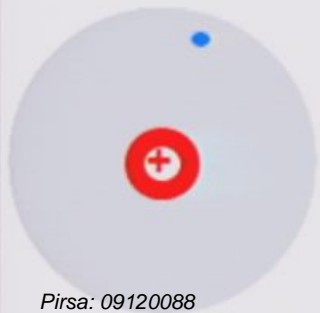


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If only the  $k$ th wave is occupied

Then the guidance equation depends only on the  $k$ th wave

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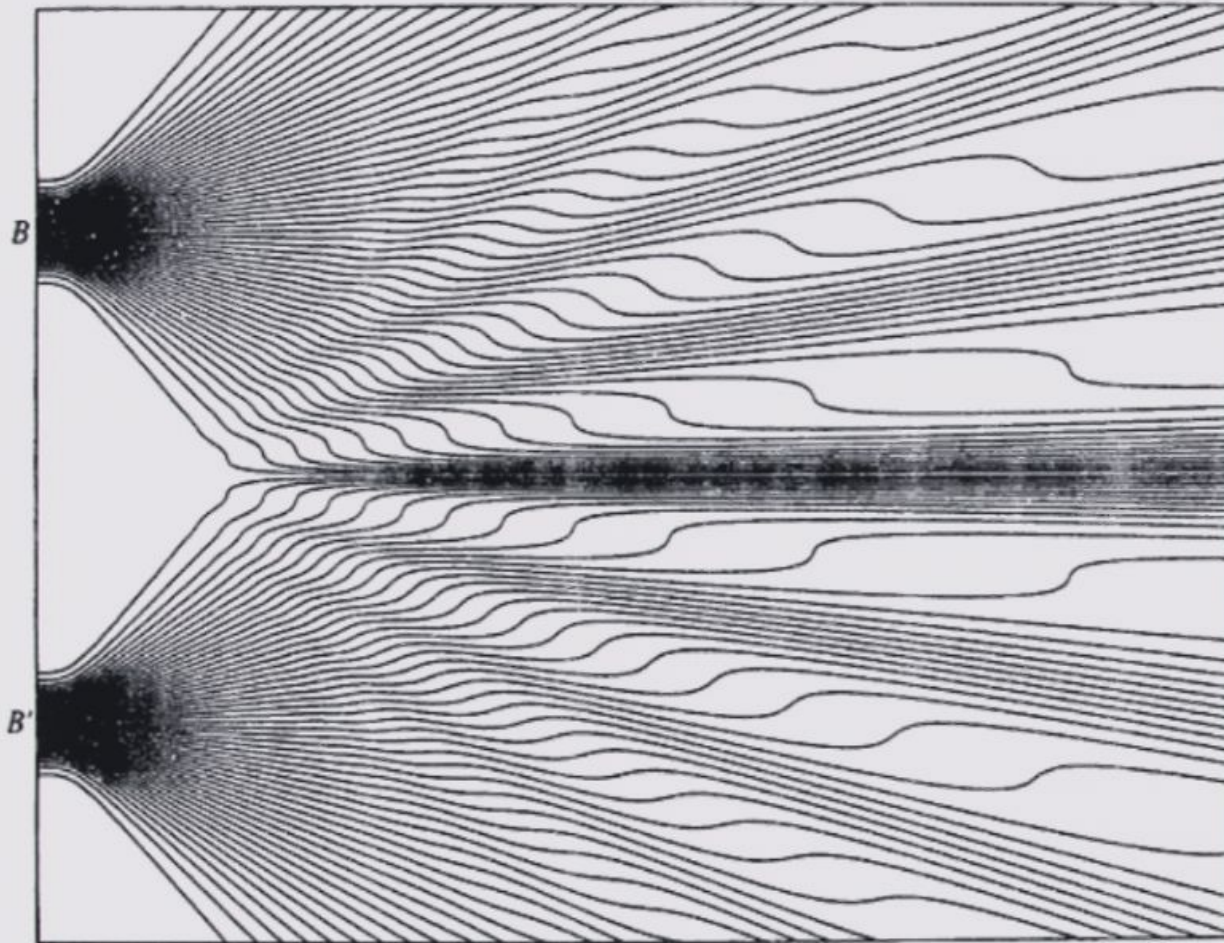
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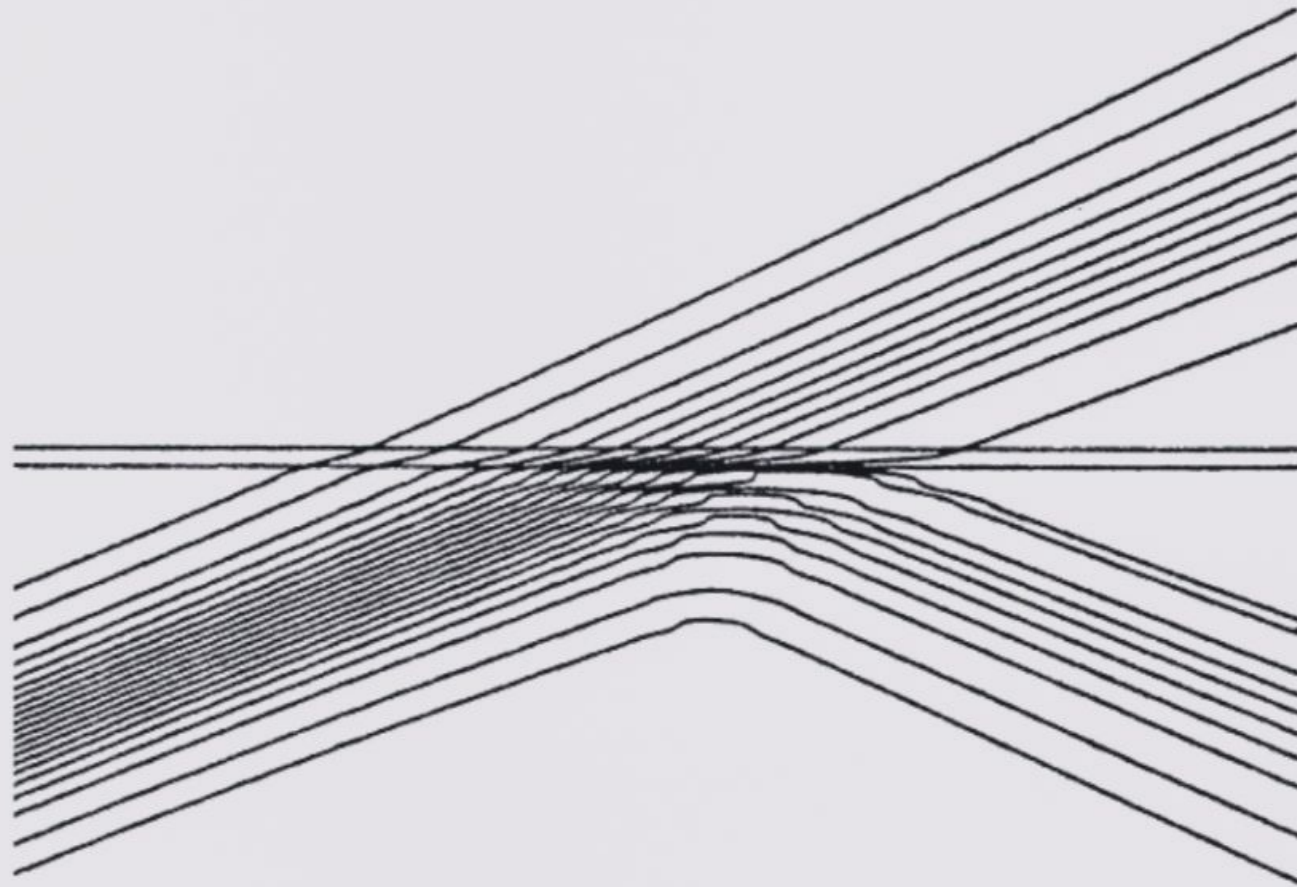
then  $R^2 = R_a^2 + R_b^2$  and  $\nabla S = \frac{R_a^2 \nabla S_a + R_b^2 \nabla S_b}{R_a^2 + R_b^2}$

$$\frac{d\zeta(t)}{dt} = \frac{1}{m} [\nabla S(\mathbf{r}, t)]_{\mathbf{r}=\zeta(t)} = \frac{\nabla S_a}{m} \quad \text{If } \zeta \in \text{Support of } \psi_a$$

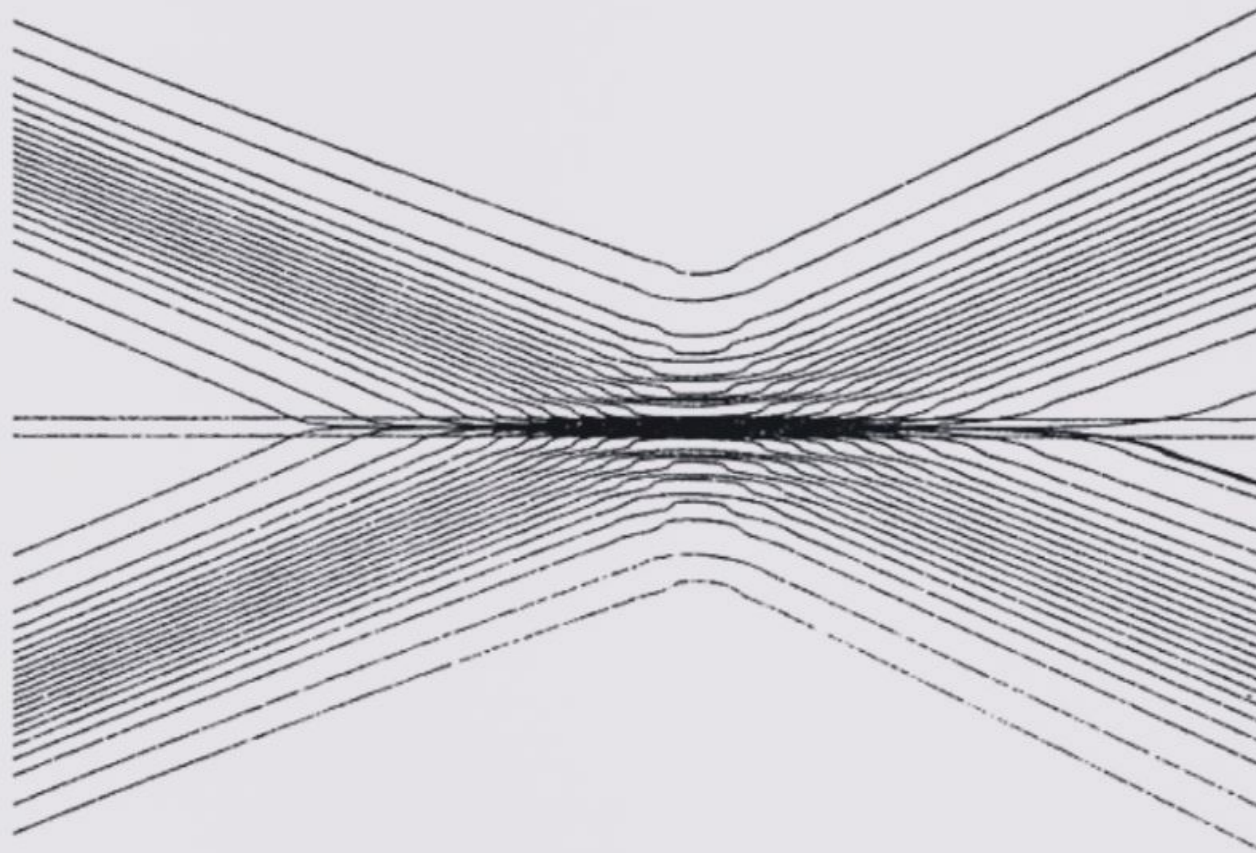




## Double slit experiment



Transmission through a barrier (probability  $\frac{1}{2}$ )



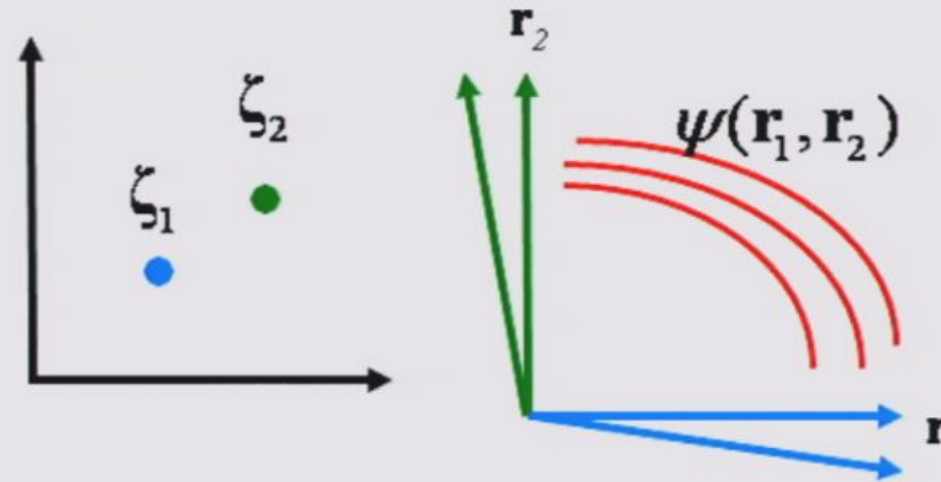
## Beam splitter experiment

## The deBroglie-Bohm interpretation for many particles

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$\nearrow$   
Wavefunction on  
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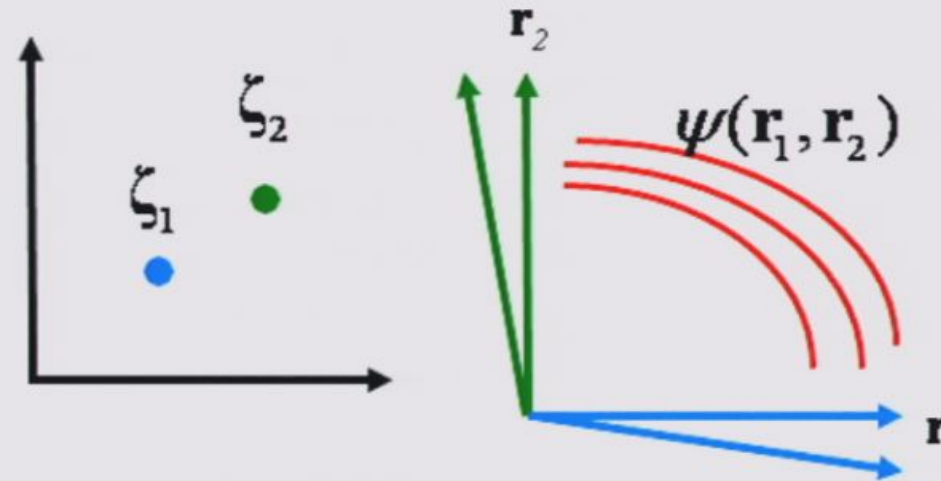


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The evolution equations:

### Schrödinger's equation

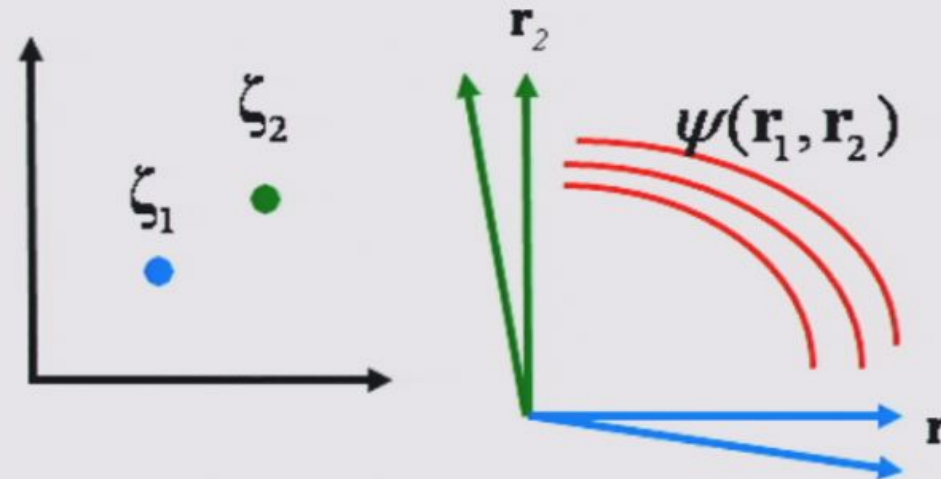
$$i\hbar \frac{\partial \psi(\mathbf{r}_1, \mathbf{r}_2, t)}{\partial t} = -\frac{\hbar^2}{2m_1} \nabla_1^2 \psi(\mathbf{r}_1, \mathbf{r}_2, t) - \frac{\hbar^2}{2m_2} \nabla_2^2 \psi(\mathbf{r}_1, \mathbf{r}_2, t) + V(\mathbf{r}_1, \mathbf{r}_2) \psi(\mathbf{r}_1, \mathbf{r}_2, t)$$

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The guidance equation

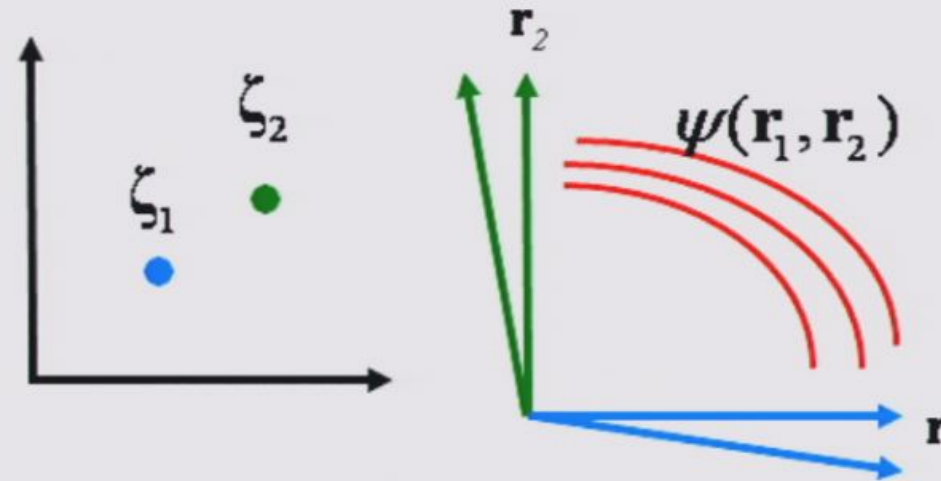


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$$\frac{d\boldsymbol{\zeta}_1(t)}{dt} = \frac{1}{m_1} [\nabla_1 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\boldsymbol{\zeta}_1(t), \mathbf{r}_2=\boldsymbol{\zeta}_2(t)} = \frac{1}{m_1} [\nabla_1 S_1(\mathbf{r}_1, t)]_{\mathbf{r}_1=\boldsymbol{\zeta}_1(t)}$$

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \phi^{(1)}(\mathbf{r}_1, t) \chi^{(2)}(\mathbf{r}_2, t) \quad \text{Product state}$$

$$= R_1(\mathbf{r}_1, t) e^{iS_1(\mathbf{r}_1, t)/\hbar} R_2(\mathbf{r}_2, t) e^{iS_2(\mathbf{r}_2, t)/\hbar}$$

$$S(\mathbf{r}_1, \mathbf{r}_2, t) = S_1(\mathbf{r}_1, t) + S_2(\mathbf{r}_2, t)$$

$$\frac{d\zeta_1(t)}{dt} = \frac{1}{m_1} [\nabla_1 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\zeta_1(t), \mathbf{r}_2=\zeta_2(t)} = \frac{1}{m_1} [\nabla_1 S_1(\mathbf{r}_1, t)]_{\mathbf{r}_1=\zeta_1(t)}$$

$$\frac{d\zeta_2(t)}{dt} = \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\zeta_1(t), \mathbf{r}_2=\zeta_2(t)} = \frac{1}{m_2} [\nabla_2 S_2(\mathbf{r}_2, t)]_{\mathbf{r}_2=\zeta_2(t)}$$

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$$= R_1(\mathbf{r}_1, t) e^{iS_1(\mathbf{r}_1, t)/\hbar} R_2(\mathbf{r}_2, t) e^{iS_2(\mathbf{r}_2, t)/\hbar}$$

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$$\frac{d\zeta_2(t)}{dt} = \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1=\zeta_1(t), \mathbf{r}_2=\zeta_2(t)} = \frac{1}{m_2} [\nabla_2 S_2(\mathbf{r}_2, t)]_{\mathbf{r}_2=\zeta_2(t)}$$

The two particles evolve independently



$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \phi^{(1)}(\mathbf{r}_1, t) \chi^{(2)}(\mathbf{r}_2, t) \quad \text{Product state}$$

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The two particles evolve independently

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

$(\zeta_1, \zeta_2) \in$  support of  $\phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t)$   $j$ th wave is occupied

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$(\zeta_1, \zeta_2) \notin$  support of  $\phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t)$   $j$ th wave is empty

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

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$(\zeta_1, \zeta_2) \notin$  support of  $\phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t)$   $j$ th wave is empty

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

$(\zeta_1, \zeta_2) \in \text{support of } \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad j\text{th wave is occupied}$

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If only the  $k$ th wave is occupied



$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

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If only the  $k$ th wave is occupied

Then the particles evolve independently

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

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If only the  $k$ th wave is occupied

Then the particles evolve independently

But in general, they do not

$$\psi(\mathbf{r}_1, \mathbf{r}_2, t) = \sum_j c_j \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad \text{Entangled state}$$

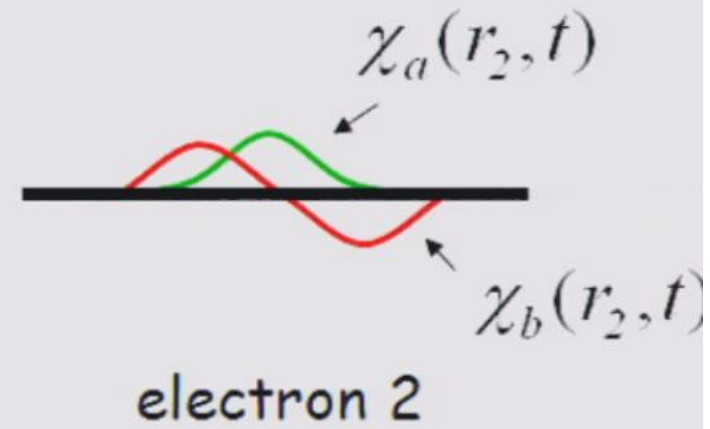
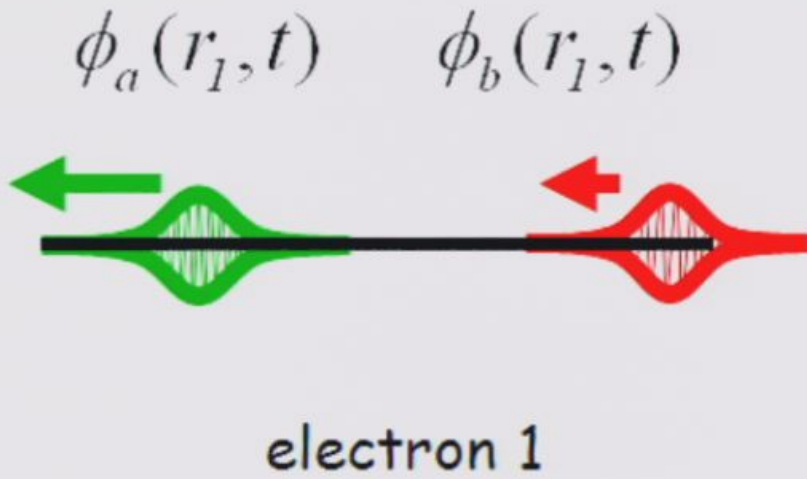
$(\zeta_1, \zeta_2) \in \text{support of } \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad j\text{th wave is occupied}$

$(\zeta_1, \zeta_2) \notin \text{support of } \phi_j^{(1)}(\mathbf{r}_1, t) \chi_j^{(2)}(\mathbf{r}_2, t) \quad j\text{th wave is empty}$

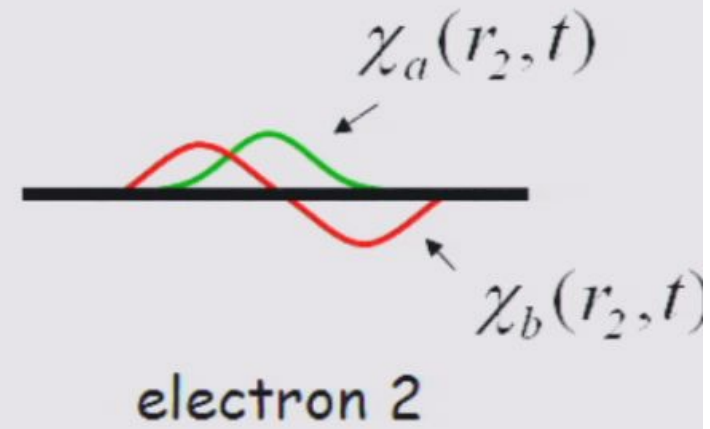
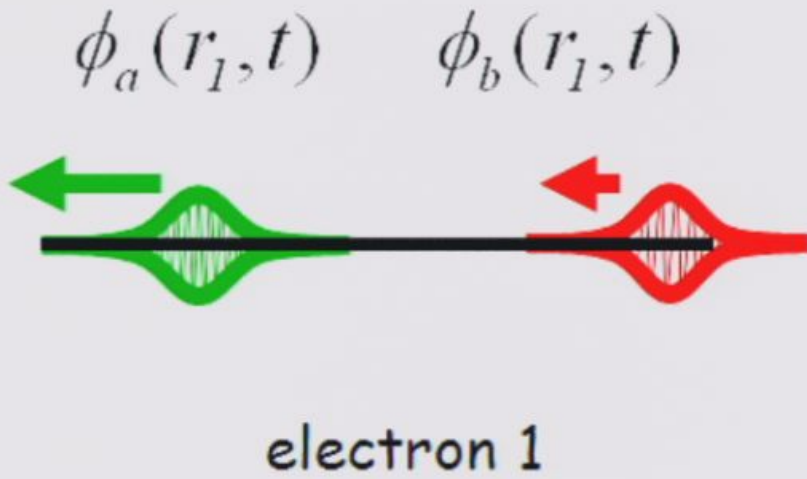
If only the  $k$ th wave is occupied

Then the particles evolve independently

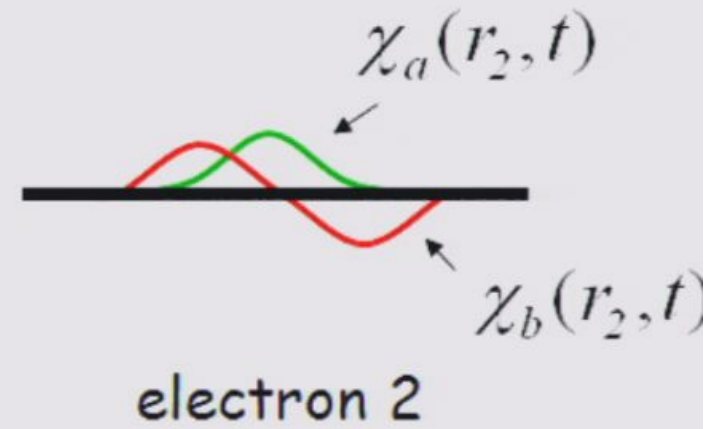
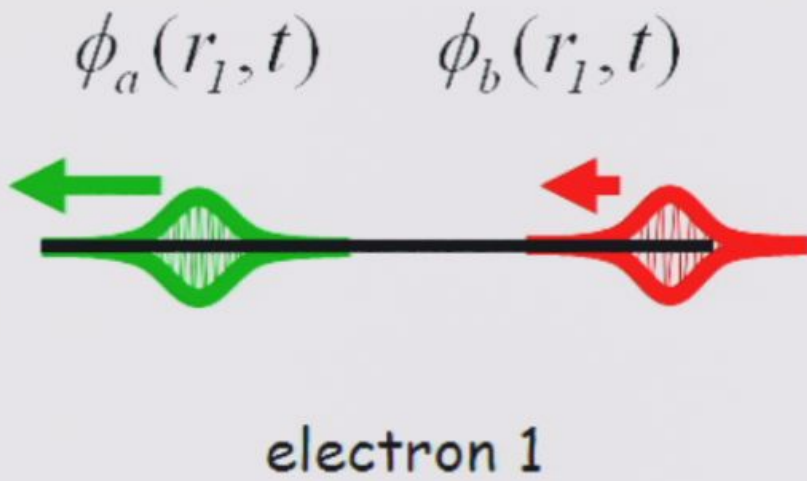
But in general, they do not



$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

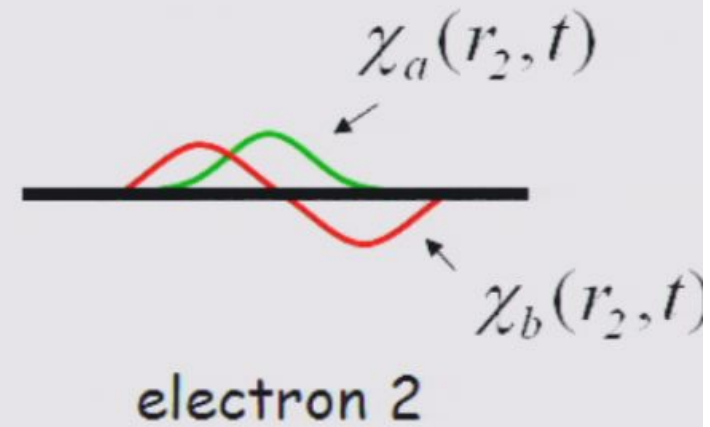
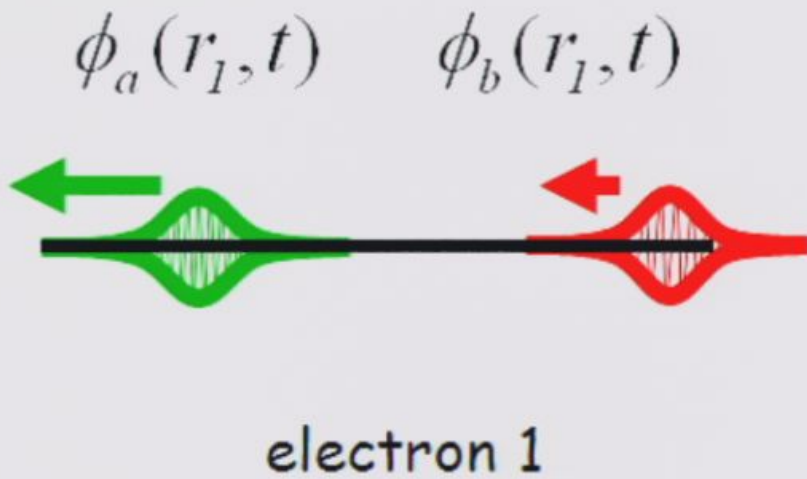


$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

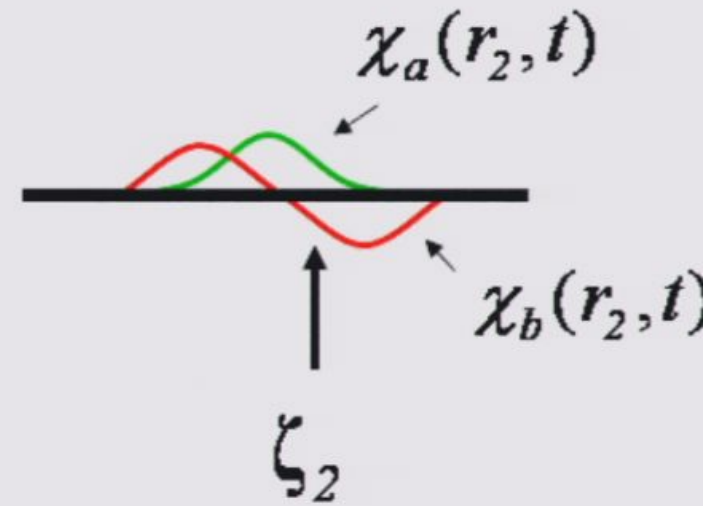
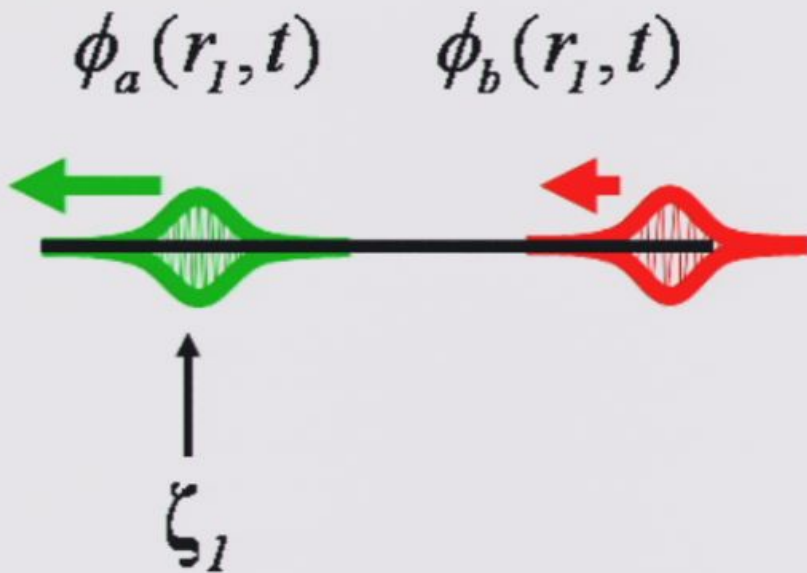


$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$



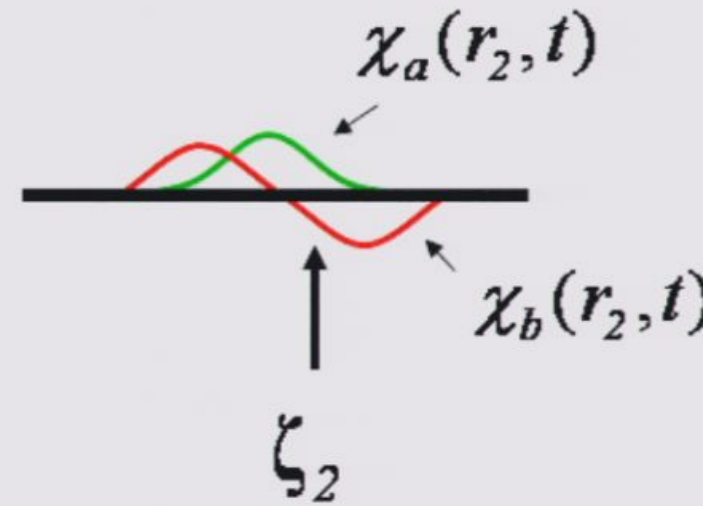
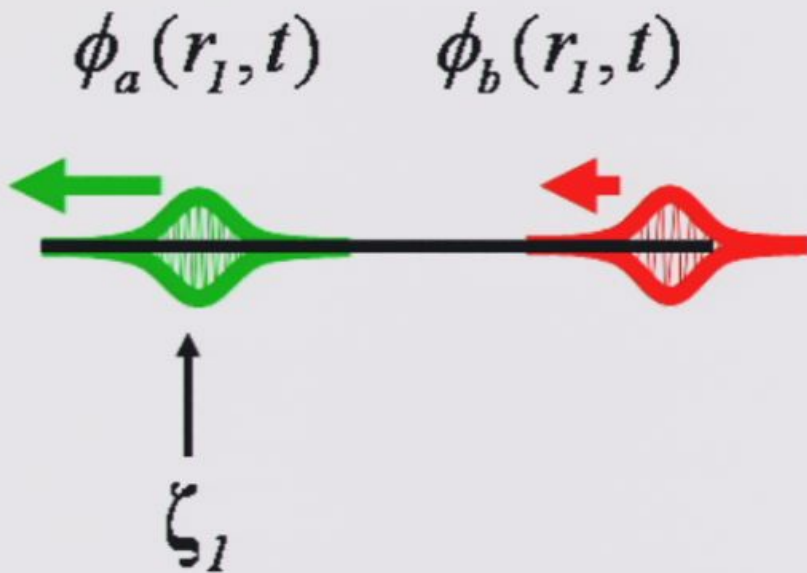


$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$



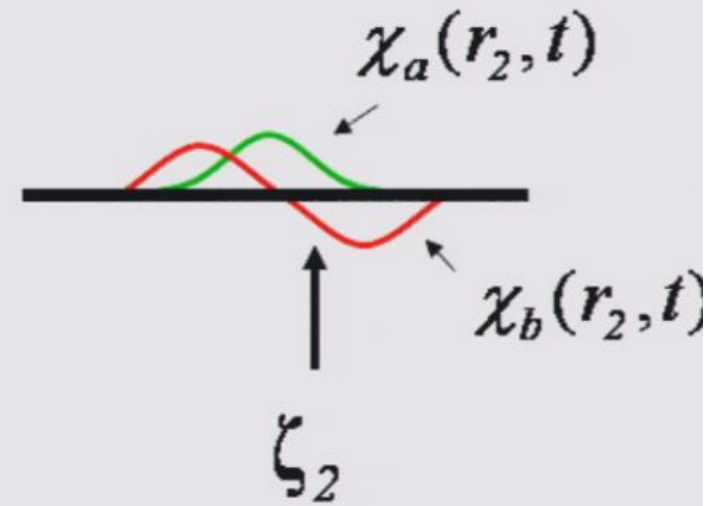
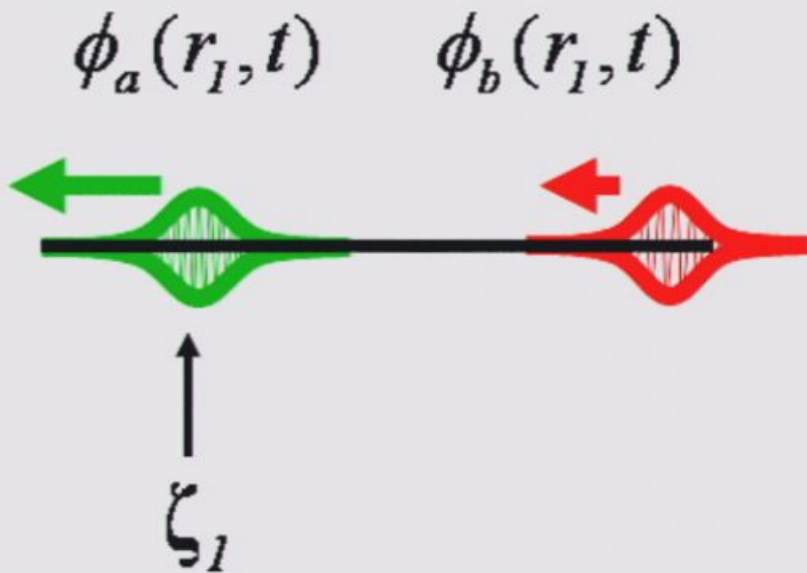
$$\psi(r_1, r_2; t) =$$

$$c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$



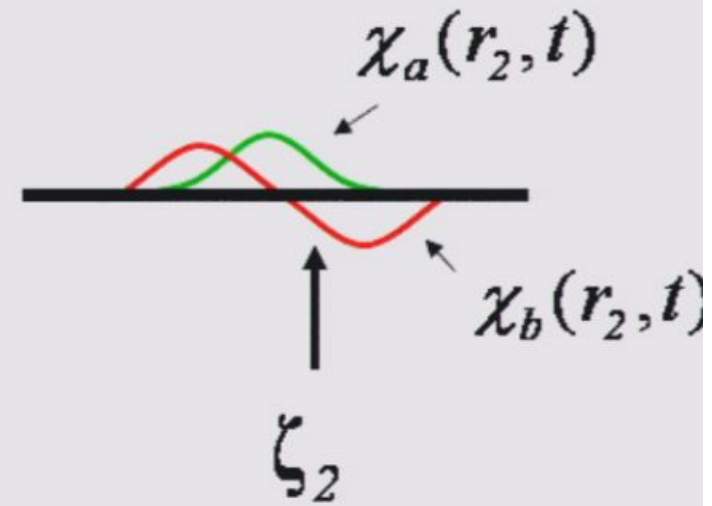
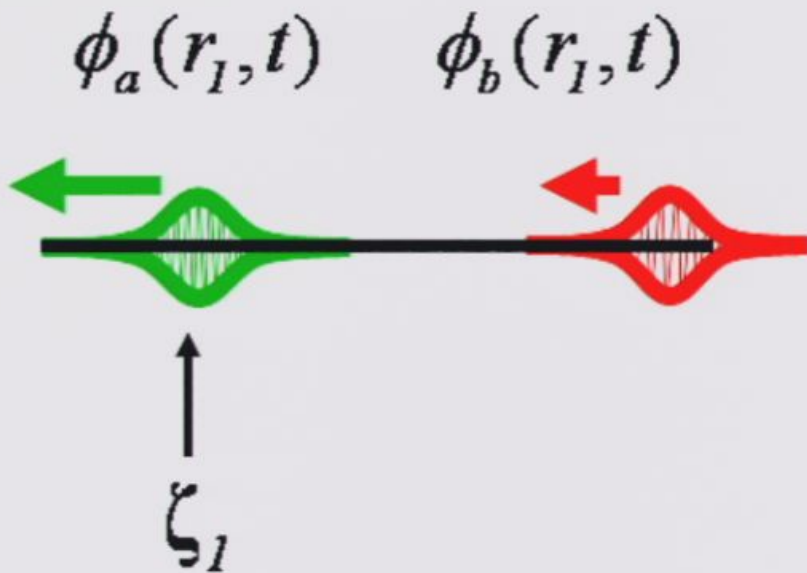
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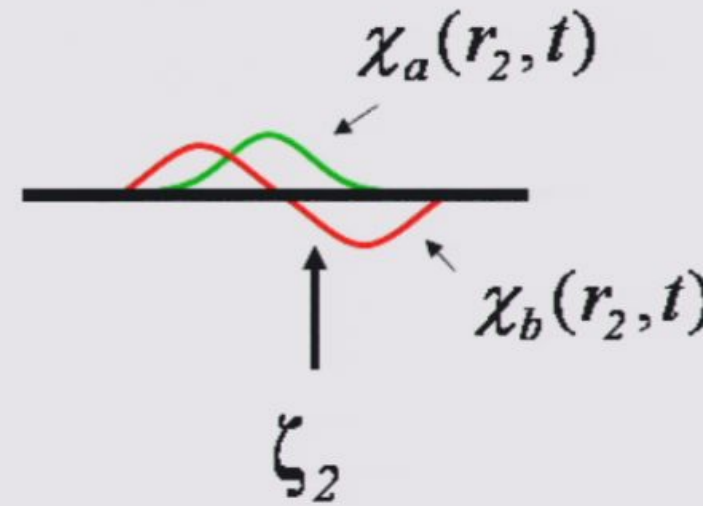
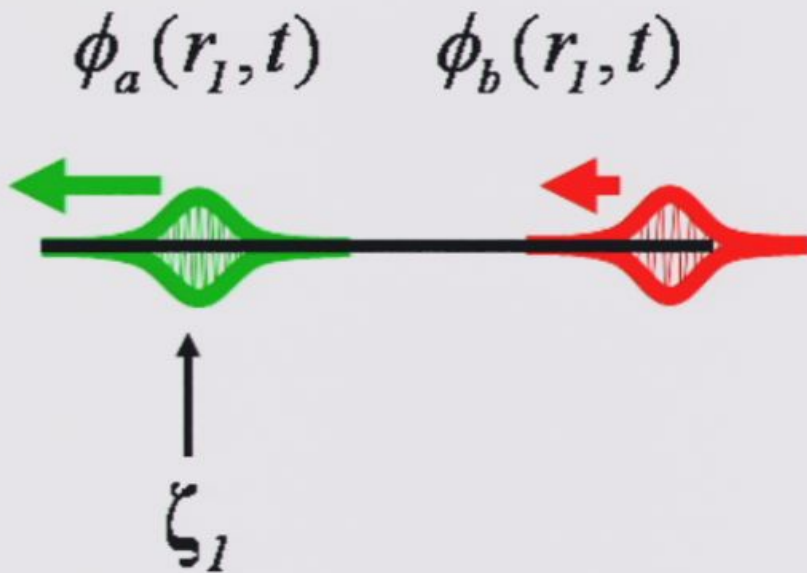


$$\psi(r_1, r_2; t) =$$

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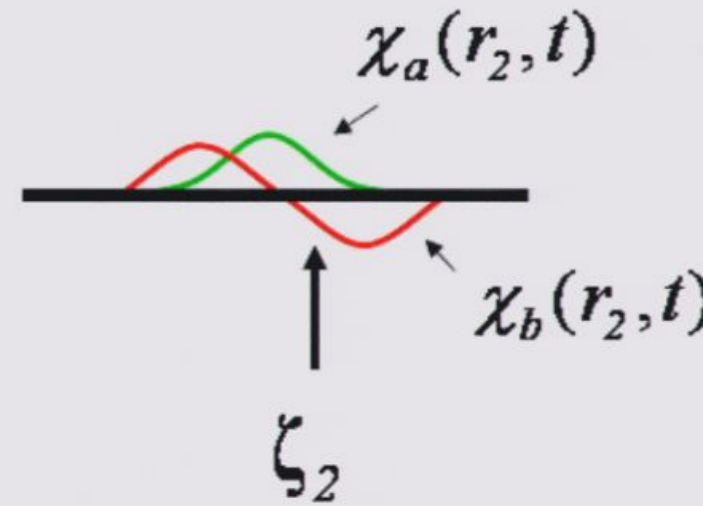
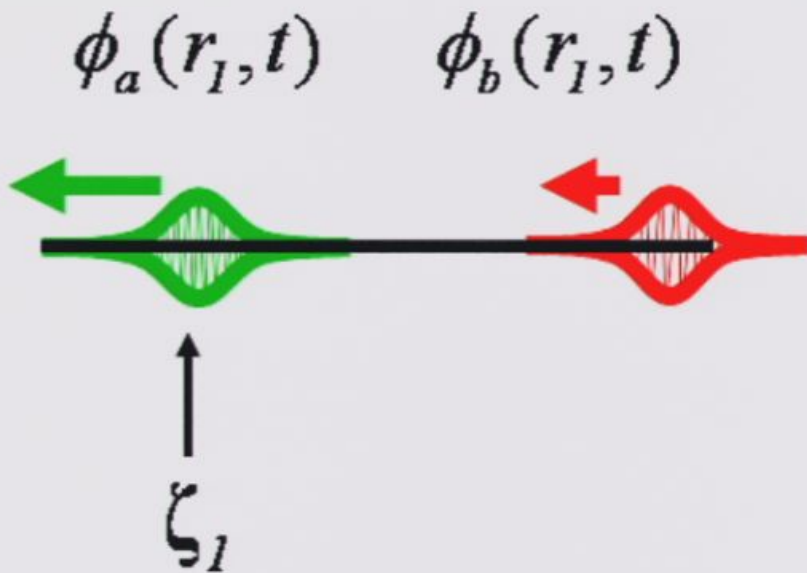
occupied wave



$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

}  
 occupied wave

$$\frac{d\zeta_2(t)}{dt} = \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1 = \zeta_1(t), \mathbf{r}_2 = \zeta_2(t)}$$



$$\psi(r_1, r_2; t) =$$

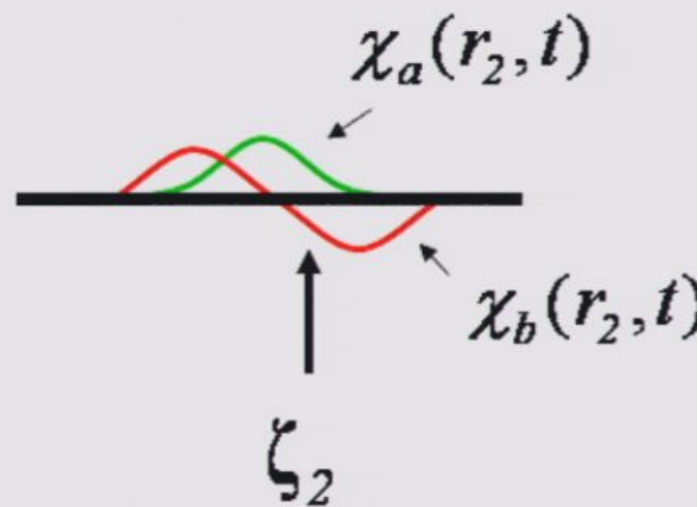
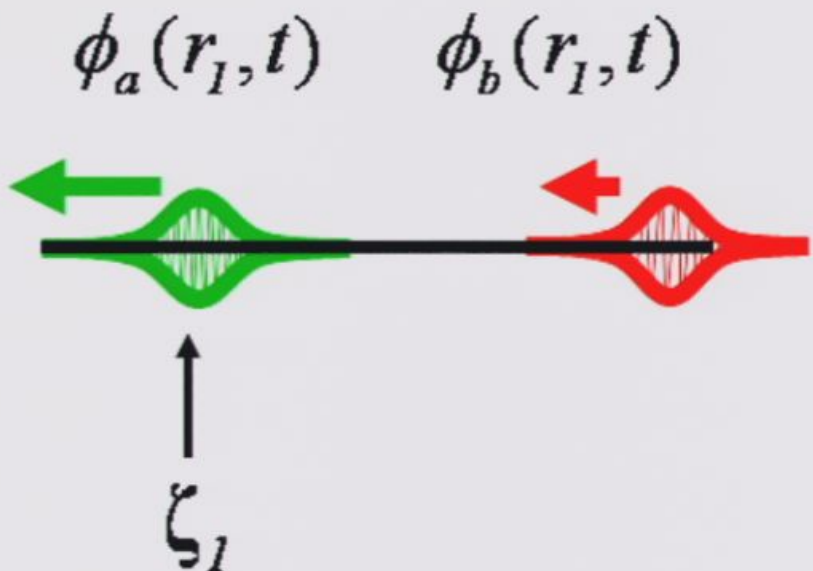
$$c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$



occupied wave

$$\frac{d\zeta_2(t)}{dt} = \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1 = \zeta_1(t), \mathbf{r}_2 = \zeta_2(t)}$$





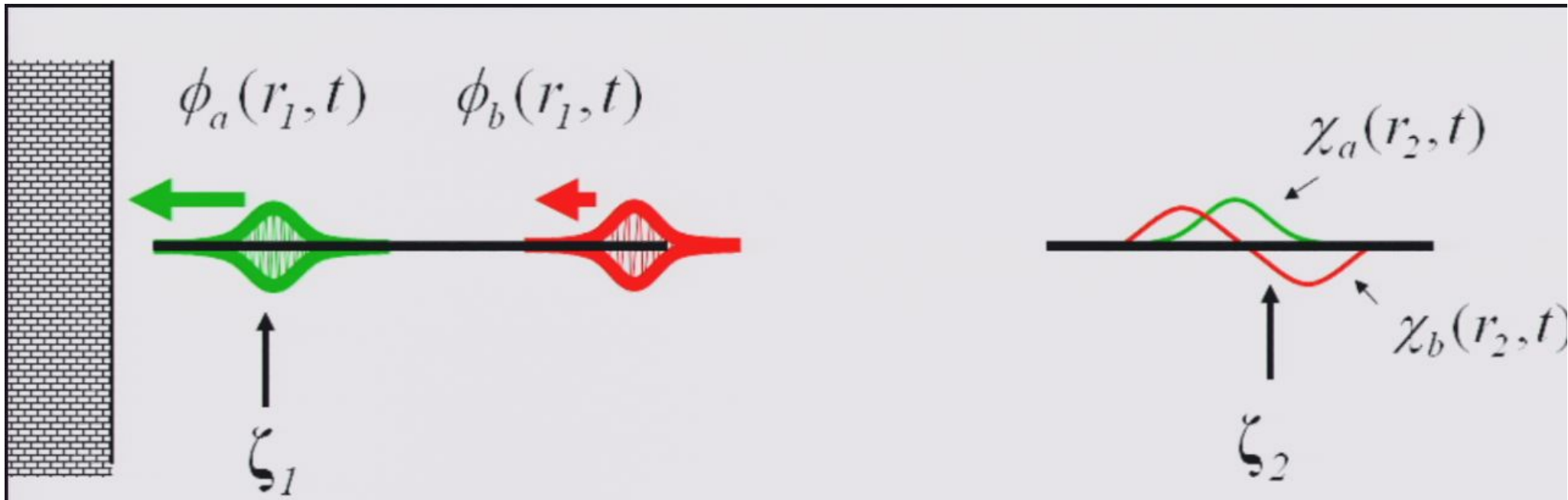
$$\psi(r_1, r_2; t) =$$

$$c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

occupied wave



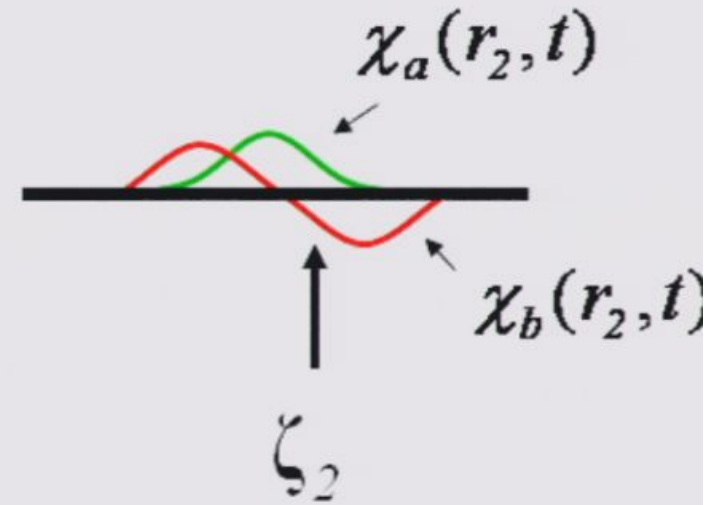
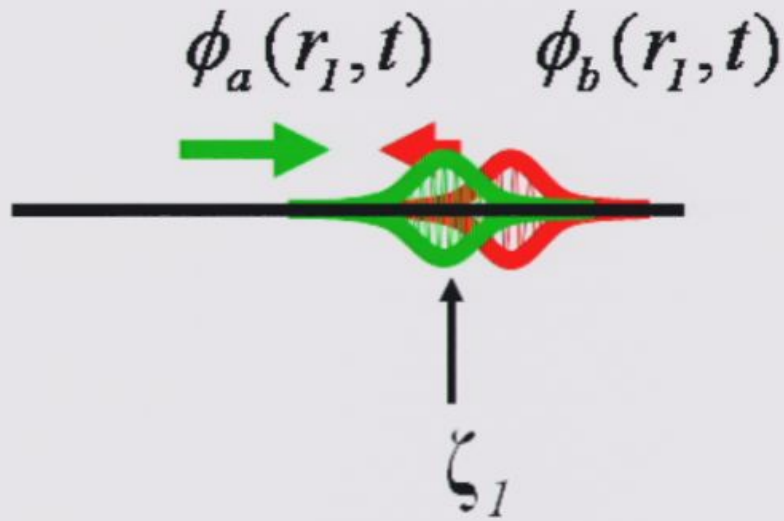
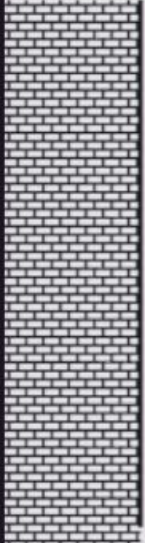
$$\frac{d\zeta_2(t)}{dt} = \frac{1}{m_2} [\nabla_2 S(\mathbf{r}_1, \mathbf{r}_2, t)]_{\mathbf{r}_1 = \zeta_1(t), \mathbf{r}_2 = \zeta_2(t)}$$



$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

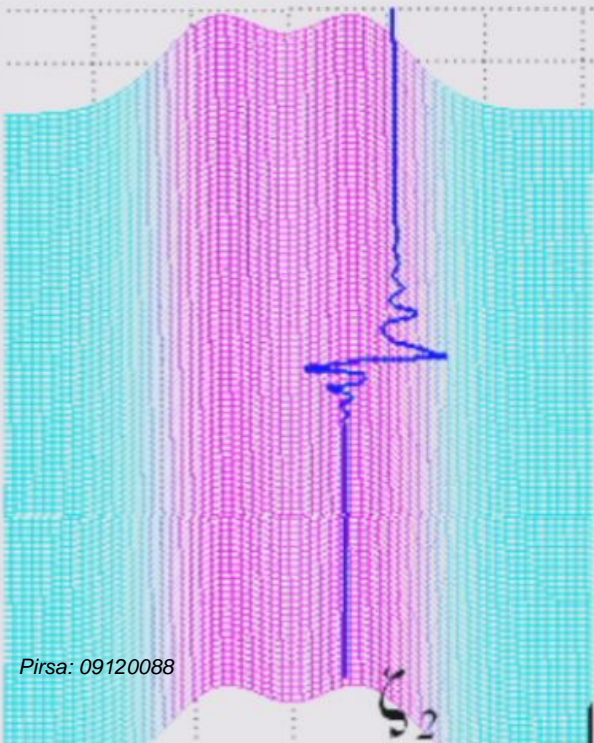
}  
 occupied wave



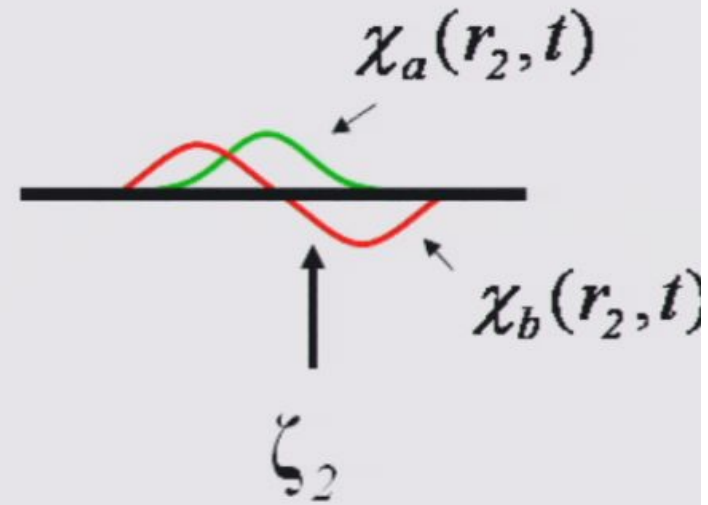
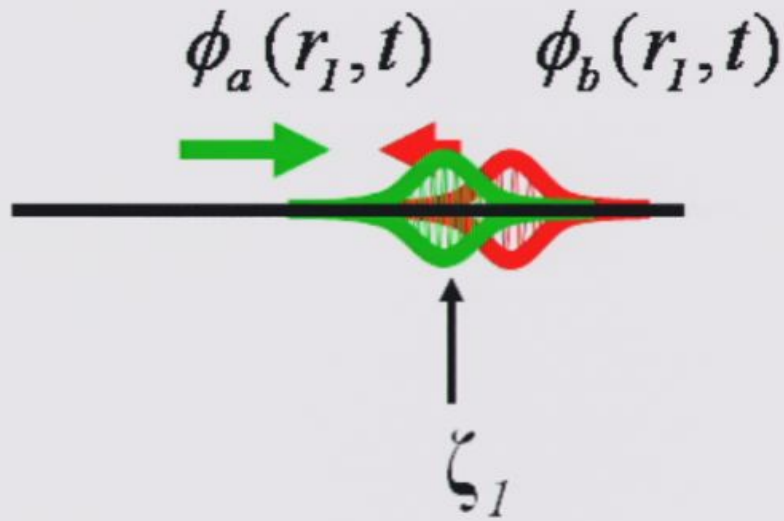
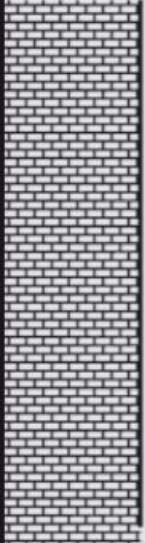


$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

both waves occupied

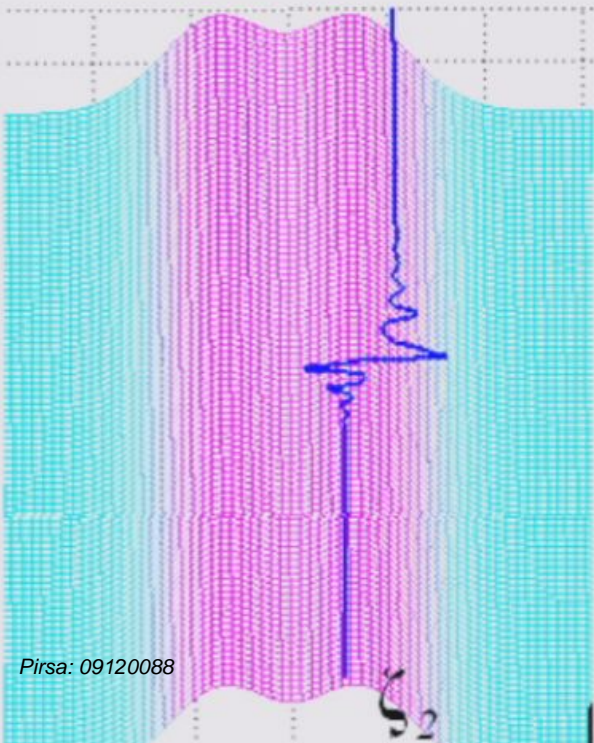


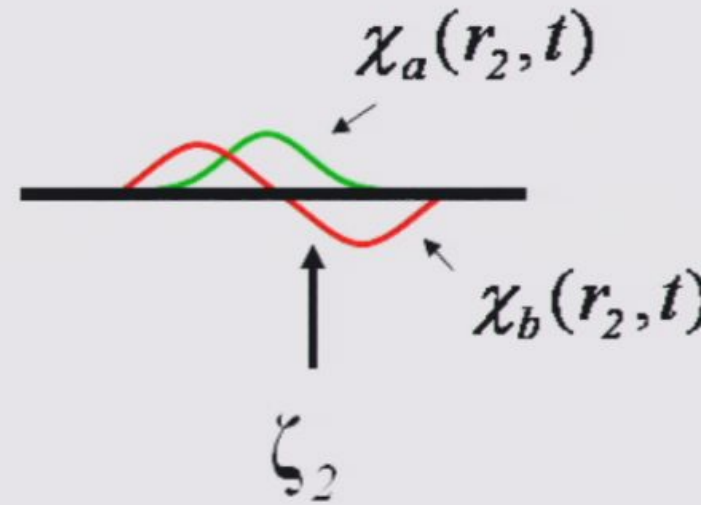
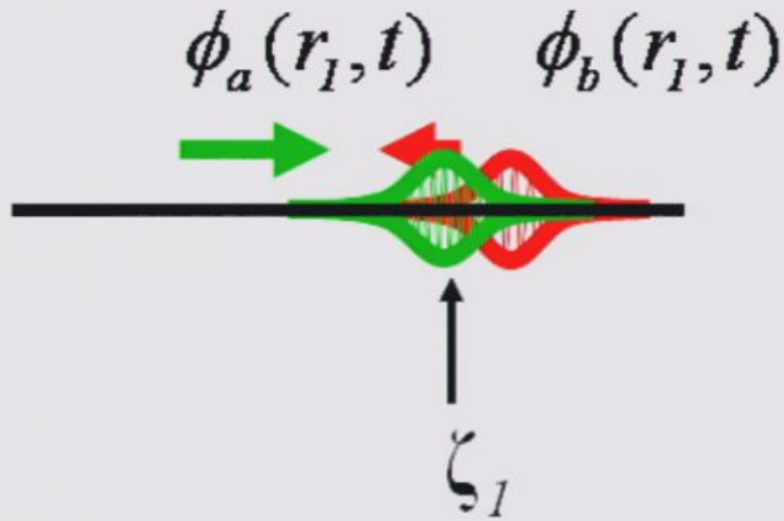
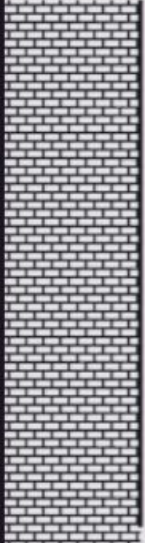




$$\psi(r_1, r_2; t) = c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$

both waves occupied



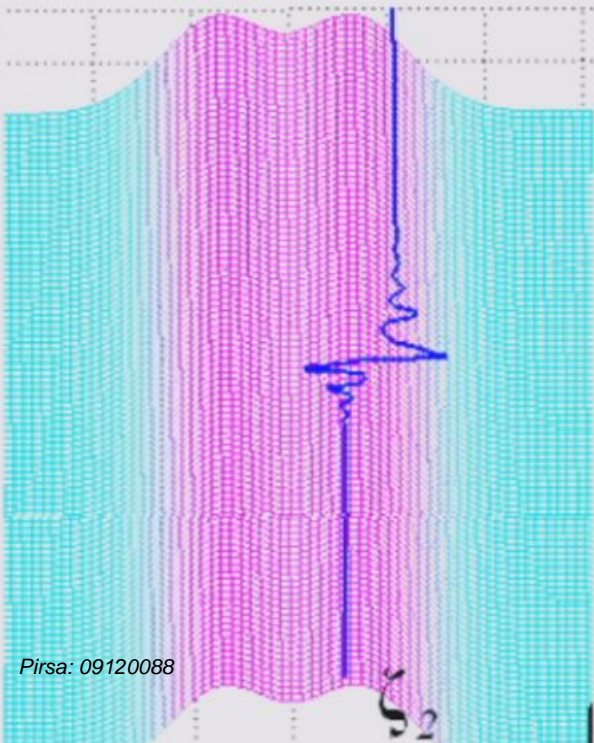


$$\psi(r_1, r_2; t) =$$

$$c_a \phi_a(r_1, t) \chi_a(r_2, t) + c_b \phi_b(r_1, t) \chi_b(r_2, t)$$



both waves occupied



Failure of local causality