

Title: Cosmology - Review (PHYS 621) - Lecture 8

Date: Dec 09, 2009 10:00 AM

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Abstract:

Tight coupling approximation $V_\gamma \approx V_b$

keep photon monopole, dipole, neglect higher moments

Baryon / photon ratio R v. important in governing the physics

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Keep photon monopole, dipole, neglect higher moments

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① Number continuity

δ 's are not created or destroyed!

pling approximation $v_\gamma \approx v_b$

p photon monopole, dipole, neglect higher moments

neutrino/photon ratio R v. important in governing the physics

① Number continuity n_γ

δ 's are not created or destroyed!

no expansion

$$\dot{n}_\gamma + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma) = 0$$

pling approximation $v_\gamma = v_b$

p photon monopole, dipole, neglect higher moments

ryon / photon ratio R v. important in governing the physics

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no expansion

$$\dot{n}_\gamma + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma) = 0$$

$$n_\gamma \propto a$$

pling approximation $v_\gamma \approx v_b$

p. photon monopole, dipole, neglect higher moments

photon/baryon ratio R v. important in governing the physics

① Number continuity n_γ

γ 's are not created or destroyed!

no expansion

$$\dot{n}_\gamma + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma) = 0$$

$$n_\gamma \propto a^{-3}$$

$$\dot{n}_\gamma + 3n_\gamma \frac{\dot{a}}{a} + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma) = 0$$

notation $V_\gamma \equiv v_b$

monopole, dipole, neglect higher moments

ratio R v. important in governing the physics

v continuity n_γ

are not created or destroyed!

no expansion

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$$\dot{n}_\gamma + 3n_\gamma \frac{\dot{a}}{a} + \vec{\nabla} \cdot (n_\gamma \vec{V}_\gamma) = 0$$

Linear

approximation $V_\gamma = V_b$

monopole, dipole, neglect higher moments

in ratio R v. important in governing the physics

over continuity n_γ

's are not created or destroyed!

no expansion

$$\dot{n}_\gamma + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma) = 0$$

$$n_\gamma \propto a^{-3}$$

$$\dot{n}_\gamma + 3n_\gamma \frac{\dot{a}}{a} + \vec{\nabla} \cdot (n_\gamma \vec{v}_\gamma)$$

expansion

Linearize

$$\delta n_\gamma = n_\gamma - \bar{n}_\gamma$$

$$(\delta n_\gamma)' = -3 \delta n_\gamma \frac{\dot{a}}{a} - n_\gamma \nabla \cdot \vec{v}_\gamma$$

the physics

$$n_\gamma \vec{v}_\gamma = 0$$

$$\frac{\dot{a}}{a} + \nabla \cdot (n_\gamma \vec{v}_\gamma) =$$

expansion

Linearized $\delta n_\gamma = n_\gamma - \bar{n}_\gamma$

$$(\delta \dot{n}_\gamma) = -3 \delta n_\gamma \frac{\dot{a}}{a} - n_\gamma \nabla \cdot \vec{V}_\gamma$$

$$\left(\frac{\delta \dot{n}_\gamma}{n_\gamma} \right) = -\nabla \cdot \vec{V}_\gamma$$

the physics

$$n_\gamma \vec{V}_\gamma = 0$$

$$\frac{\dot{a}}{a} + \nabla \cdot (n_\gamma \vec{V}_\gamma) = 0$$

expansion

Continuity Equation

$n\alpha$

Continuity Equation

$$n \propto T^3$$

Continuity Equation

$$n \propto T^3$$

define temperature fluctuation $\langle \delta T \rangle$

Continuity Equation

$$n_{\gamma} \propto T^3$$

define temperature fluctuation δT

$$\frac{\delta n_{\gamma}}{n_{\gamma}} = 3 \frac{\delta T}{T} \equiv 3 \delta T$$

Continuity Equation

$$n_{\gamma} \propto T^3$$

define temperature fluctuation \textcircled{H}

$$\frac{\delta n_{\gamma}}{n_{\gamma}} = 3 \frac{\delta T}{T} \equiv 3 \textcircled{H}$$

Real space

$$\textcircled{H} = -\frac{1}{3} \vec{\nabla} \cdot \vec{v}_{\gamma}$$

Continuity Equation

$$n_{\gamma} \propto T^3$$

define temperature fluctuation \dot{H}

$$\frac{\delta n_{\gamma}}{n_{\gamma}} = 3 \frac{\delta T}{T} \equiv 3 \dot{H}$$

Real space

$$\dot{H} = -\frac{1}{3} \vec{\nabla} \cdot \vec{v}_{\gamma}$$

Fourier space

$$\dot{H} = -\frac{1}{3} \dots$$

Equation

$$\propto T^3$$

temperature fluctuation \textcircled{H}

$$\frac{\delta n_\gamma}{n_\gamma} = 3 \frac{\delta T}{T} \equiv 3 \textcircled{H}$$

Real space

$$\textcircled{H} = -\frac{1}{3} \vec{\nabla} \cdot \vec{v}_\gamma$$

Fourier space

$$\textcircled{H} = -\frac{1}{3} i \vec{k} \cdot \vec{v}_\gamma$$

$$\textcircled{1} = -\frac{1}{3} i k \cdot \vec{v}_\sigma$$

Momentum conservation

• no expansion

$$\dot{\vec{q}} = \vec{F}$$

- wavelength stretches w/ expansion

$$\dot{\vec{q}} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

Momentum conservation

• no expansion $\dot{\vec{q}} = \vec{F}$

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$$\dot{\vec{q}} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

- m of particles

Momentum conservation

• no expansion $\dot{\vec{q}} = \vec{F}$

- wavelength stretches w/ expansion

$$\dot{\vec{q}} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

- collection of particles

momentum conservation

no expansion $\vec{q} = F$

wavelength stretches w/ expansion

$$\vec{q} + \frac{\dot{a}}{a} \vec{q} = F$$

collection of particles: $m \rightarrow m$ density (ex -)

time

$$\vec{q} = F$$

retches w/ expansion

$$\frac{\dot{a}}{a} \vec{q}$$

partic

mtm \rightarrow mtm density
for

$$(e_x + P_x) \vec{v}_x$$

time

$$\frac{d}{dt} \vec{q} = \vec{F}$$

retches w/ expansion

$$\frac{d}{dt} \vec{q} = \vec{F}$$

particles: m m → m m density $(\rho_x + P_x) \vec{v}_x$

force → pressure gradient ∇P

um conservation

expansion $\vec{q} = \vec{F}$

wavelength stretches w/ expansion

$$\vec{q} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

direction of particles: $m \vec{v} \rightarrow m \vec{v}$ density $(\rho + P)$

force \rightarrow pressure gradient ∇

$$[(\rho + P) \vec{v}]$$

um conservation

expansion $\vec{q} = \vec{F}$

wavelength stretches w/ expansion

$$\vec{q} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

collection of particles: $m \rightarrow m$ density $(\rho + P)$

force \rightarrow pressure gradient ∇

$$\nabla \cdot [(\rho + P) \vec{v}] = -4 \frac{\dot{a}}{a} (\rho + P) \vec{v} + \nabla^2 P$$

um conservation

expansion $\vec{q} = F$

wavelength stretches w/ expansion

$$\vec{q} + \frac{\dot{a}}{a} \vec{q} = F$$

collection of particles: $m \rightarrow m$ density $(\rho +$

force \rightarrow pressure gradient ∇

$$\left[(\rho + P) \vec{v} \right] = -4 \frac{\dot{a}}{a} (\rho + P) \vec{v} + \vec{\nabla} P$$
$$\frac{4}{3} \rho \vec{v} = \frac{1}{3} \vec{\nabla} P$$

momentum conservation

expansion $\vec{q} = \vec{F}$

wavelength stretches w/ expansion

$$\vec{q} + \frac{\dot{a}}{a} \vec{q} = \vec{F}$$

collection of particles: m \rightarrow m density $(\rho_\gamma + P_\gamma)$

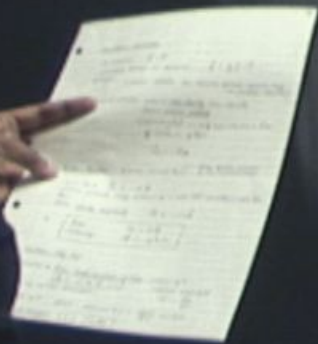
force \rightarrow pressure gradient ∇P

$$\begin{aligned} \left[(\rho_\gamma + P_\gamma) \vec{v}_\gamma \right] &= -4 \frac{\dot{a}}{a} (\rho_\gamma + P_\gamma) \vec{v}_\gamma + \vec{\nabla} P_\gamma \\ \frac{4}{3} \rho_\gamma \vec{v}_\gamma &= -\frac{1}{3} \vec{\nabla} P_\gamma \end{aligned}$$

$$\vec{v}_x = -\nabla \psi$$

continuity $(\rho_x + P_x) \vec{v}_x$
 gradient ∇P
 $\rho_x + \nabla P_x$

Euler Equation



Euler Equation

- Fourier space

$$\vec{v}_\gamma = -ik \textcircled{(-1)}$$

Euler Equation

- Fourier space

pressure gradients

$$\vec{v}_\gamma = -ik \textcircled{(-)}$$

Euler Equation

- Fourier space

$$\vec{v}_\gamma = -ik \textcircled{-1}$$

pressure gradients \rightarrow curl free flow

Euler Equation

- Fourier space

$$\vec{V}_\gamma = -i\vec{k} \textcircled{(-1)}$$

pressure gradients \rightarrow curl free flow

define velocity amplitude

$$\vec{V}_\gamma \equiv -i v_\gamma \hat{k}$$

$$\text{Euler: } \dot{V} \gamma = k \oplus H$$

$$\text{Euler: } \dot{V}_x = k \textcircled{H}$$

$$\text{Continuity: } \textcircled{H} = -\frac{1}{3} k \cdot V_x$$

$$\text{Euler : } \dot{V}_x = k \odot H$$

$$\text{Continuity : } \odot H = -\frac{1}{3} k \cdot V_x$$

Oscillator; Take One

Combine these equations

Simple harmonic oscillator

$$\ddot{H} + c_s^2 k_r^2 H = 0$$

$$c_s^2 = \frac{1}{3} \text{ (photon dominated)}$$

Oscillator: Take One

Combine these equations

Simple harmonic oscillator

$$\ddot{H} + c_s^2 k^2 H = 0$$

$$c_s^2 = \frac{1}{3} \text{ (photon dominated)}$$

adiabatic sound speed

$$c_s^2 = \frac{P_\gamma}{\rho_\gamma}$$

General Solution

$$\psi(\eta) = \psi(0) \cos(k\eta) + \frac{\dot{\psi}(0)}{kC} \sin(k\eta)$$

ψ

General Solution

$$\psi(\eta) = \psi(0) \cos(kS) + \frac{\dot{\psi}(0)}{kC_s} \sin(kS)$$

Sound horizon $S \equiv \int c_s d\eta$



General Solution

$$\Phi(\eta) = \Phi(0) \cos(kS) + \frac{\dot{\Phi}(0)}{kC_s} \sin(kS)$$

Sound horizon $S \equiv \int c_s d\eta$

Harmonic Extrema

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all modes "frozen in" at recombination (*).

Tom

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all modes "frozen in" at recombination (*).

Temperature perturbations have different amplitudes for different modes.

Harmonic Extrema

all modes "frozen in" at recombination (*).

Temperature perturbations have different amplitudes for different modes.

adiabatic (curvature) mode $\delta H(0) = 0$

$$\delta H(\eta_*) = \delta H(0) \cos(kS_x)$$

modes caught at extrema \rightarrow bigger fluctuations

tide

modes caught at extrema \rightarrow bigger fluctuations

$$k_n S_x = n\pi$$

find

modes caught at extrema \rightarrow bigger fluctuations

$$k_n S_x = n\pi$$

$$\text{fundamental frequency} = \kappa_A = \frac{\pi}{S_x}$$

tide

modes caught at extrema \rightarrow bigger fluctuations

$$k_n S_x = n\pi$$

fundamental frequency =

$$k_A = \frac{\pi}{S_x}$$

harmonic relationship

to other extrema as 1 : 2 : 3

Peak location

fundamental physical scale

→ angular scale by simple projection
at angular diameter distance D_A

Peak location

fundamental physical scale

→ angular scale by simple projection
at angular diameter distance D_A

$$\theta_A = \frac{\lambda_A}{D_A}$$

Peak location

fundamental physical scale

→ angular scale by simple projection
at angular diameter distance D_A

$$\theta_A = \frac{\lambda_A}{D_A}$$

$$\ell_A = k_A D_A$$

Flat universe

$D_A = \text{conformal dist} = \eta_0 - \eta_* \approx \eta_0$

Flat universe

$$D_A = \text{conformal dist} = \int_0^z \frac{1}{H(z')} dz' \approx \frac{z}{H_0}$$

horizon distance

$$k_A = \frac{\pi}{S_D}$$

Flat universe

$$D_A = \text{conformal dist} = \eta_0 - \eta_* \approx \eta_0$$

horizon distance

$$k_A = \frac{\pi}{S_*} = \sqrt{3} \frac{\pi}{\eta_*}$$

Flat universe

$$D_A = \text{conformal dist} = \eta_0 - \eta_* \approx \eta_0$$

horizon distance

$$k_A = \frac{\pi}{S_*} = \sqrt{3} \frac{\pi}{\eta_*}$$

$$\theta_A \approx \frac{\eta_*}{\eta_0}$$

Flat universe

$$D_A = \text{conformal dist} = \eta_0 - \eta_* \approx \eta_0$$

horizon distance

$$k_A = \frac{\pi}{S_*} = \sqrt{3} \frac{\pi}{\eta_*}$$

$$\theta_A \approx \frac{\eta_*}{\eta_0}$$

MD universe

$$\eta \propto \sqrt{a}$$

$$\theta_A \approx \frac{1}{30} \approx 2^\circ$$

Flat universe

$$D_A = \text{conformal dist} = \eta_0 - \eta_* \approx \eta_0$$

horizon distance

$$k_A = \frac{\pi}{S_*} = \sqrt{3} \frac{\pi}{\eta_*}$$

$$\theta_A \approx \frac{\eta_*}{\eta_0}$$

MD universe

$$\eta \propto \sqrt{a}$$

$$\theta_A \approx \frac{1}{30} \approx 2^\circ \quad \ell_A \approx 200$$

Doppler effect

Doppler effect

Bulk motion of fluid changes observed temperature via doppler shift

Doppler effect

Bulk motion of fluid changes observed temperature via doppler shift

$$\left(\frac{\Delta T}{T}\right)_{\text{Doppler}} = \hat{n} \cdot \vec{V}_\gamma$$

Doppler effect

Bulk motion of fluid changes observed temperature via doppler shift

$$\left(\frac{\Delta T}{T}\right)_{\text{Doppler}} = \hat{n} \cdot \vec{V}_\gamma$$

avg over div^{-1} s

$$\left(\frac{\Delta T}{T}\right)_{\text{RMS}} = \frac{V_\gamma}{\sqrt{3}}$$

$$\frac{V_1}{\sqrt{3}} = \frac{\sqrt{3}}{k} \textcircled{4}$$

$$= \frac{\sqrt{3}}{k} k C_s \textcircled{4} \cos(0) \sin(kS)$$

$$\frac{V_1}{\sqrt{3}} = \frac{\sqrt{3}}{k} \textcircled{11}$$

$$= \frac{\sqrt{3}}{k} k c_s \textcircled{11}(0) \sin(ks)$$

$$= \textcircled{11}(0) \sin(ks)$$

$$\frac{V_y}{\sqrt{3}} = \frac{\sqrt{3}}{k} \textcircled{+1}$$

$$= \frac{\sqrt{3}}{k} k c_s \textcircled{+1}(0) \sin(ks)$$

$$= \textcircled{+1}(0) \sin(ks)$$

Doppler Peaks?

equal in amplitude $\frac{\pi}{2}$ out of phase

Peaks?

in amplitude $\frac{\pi}{2}$ out of phase

add in quadrature

$$\left(\frac{\Delta T}{T}\right)^2 = \textcircled{1}^2(0) \left[\cos^2(ks) + \sin^2(ks) \right] = \textcircled{1}^2(0)$$

$$\frac{V_y}{\sqrt{3}} = \frac{\sqrt{3}}{k} \textcircled{1}$$

$$= \frac{\sqrt{3}}{k} k c_s \textcircled{1}$$

$$= \textcircled{1}(0) \sin$$

Doppler Peaks?

dipole \Rightarrow doppler shifts in temp. due to motion of
the photon baryon plasma along $\hat{k} \cdot \mathbf{v}$

curl free \Rightarrow no contribution to Doppler

aks?

Doppler shifts in temp due to motion of
the photon baryon plasma along \hat{n} .

free \Rightarrow no contribution to Doppler $\hat{n} \perp \hat{k}$

Doppler Peaks?

$e \Rightarrow$ doppler shifts in temp. due to motion of the photon baryon plasma along $\hat{l} \cdot \hat{o} \cdot \hat{s}$.

curl free \Rightarrow no contribution to Doppler $\hat{n} \perp$

mathematically Doppler \rightarrow dipole $\hat{n} \cdot \hat{k} \propto Y_{1m}$

Doppler Peaks?

$e \Rightarrow$ doppler shifts in temp due to motion of the photon baryon plasma along LOS.

curl free \Rightarrow no contribution to Doppler $\hat{n} \perp$

mathematically Doppler \rightarrow dipole $\hat{n} \cdot \hat{k} \propto Y_{1m}$

Doppler Peaks?

$e \Rightarrow$ doppler shifts in temp. due to motion of the photon baryon plasma along LOS.

curl free \Rightarrow no contribution to Doppler $\hat{n} \perp$

mathematically Doppler \rightarrow dipole $\hat{n} \cdot \hat{k} \propto Y_{1m}$

couple Y_{1m} (projection) : $Y_{1m} Y_{1m}$

55?

pler shifts in temp due to motion of photon baryon plasma along LOS

$e \Rightarrow$ no contribution to Doppler $\hat{n} \perp \hat{k}$

directly Doppler \rightarrow dipole $\hat{n} \cdot \hat{k} \propto Y_{1m}$

ple Y_{lm} (projection) $Y_{1m} Y_{lm} \rightarrow Y_{l \pm 1 m}$

temp. due to motion of
plasma along LOS

Contribution to Doppler $\hat{n} \perp \hat{k}$

\rightarrow dipole $\hat{n} \cdot \hat{k} \propto Y_{1m}$

(quadrupole) $Y_{1m} Y_{2m} \rightarrow Y_{2\pm 1m} \rightarrow j_{\pm 1}$

temp. due to motion of
plasma along z axis.

tribution to Σ

$$\hat{n} \cdot \hat{k}$$

\rightarrow dipole

ction)

$$\rightarrow \int e^{i\mathbf{k}\cdot\mathbf{r}} \rightarrow j_{\ell \pm 1}(\alpha)$$

monopole couples $J_e(x)$ peaks at $\ln x$

condition of
F.O.S.

Doppler $\hat{n} \perp \hat{k}$

$\hat{k} \propto Y_{lm}$

$Y_{lm} \rightarrow Y_{l\pm 1 m} \rightarrow J_{e\pm 1}(x)$

monopole couples $j_l(x)$ peaks at $l \sim x$

dipole $\rightarrow j_{l \pm 1}(x)$

condition of
 $F=0$ s.

Doppler $\hat{n} \perp \hat{k}$

$\hat{k} \propto Y_{lm}$

$Y_{lm} \rightarrow Y_{l \pm 1 m} \rightarrow j_{l \pm 1}(x)$

monopole couples $j_l(x)$ peaks at $l \sim x$

dipole $\rightarrow j_{l \pm 1}(x)$

ion of
-0.5

pler $\hat{n} \perp \hat{k}$

$\propto Y_{lm}$

$Y_{lm} \rightarrow Y_{l \pm 1 m} \rightarrow j_{l \pm 1}(x)$